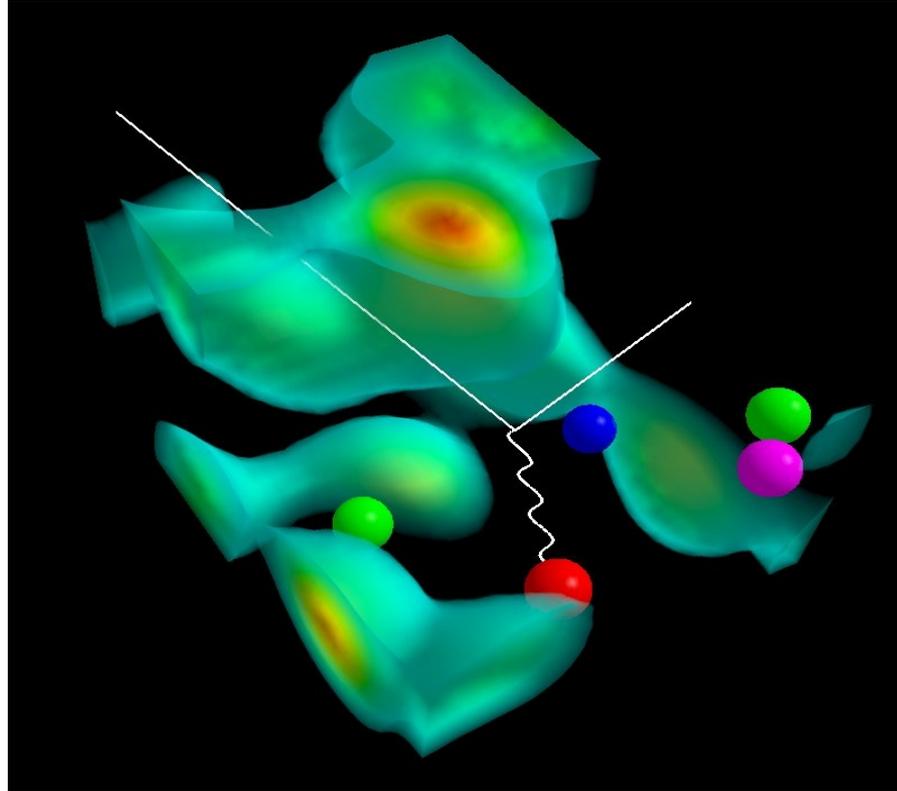


EMC effect within a quark-based theory of nuclear structure



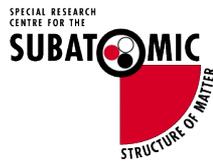
Anthony W. Thomas

HiX

Kolympari : 17th August 2019



Australian Government
Australian Research Council



Outline

I. Nuclei from Quarks

- start from a QCD-inspired model of *hadron* structure
- develop a quantitative theory of nuclear structure

II. Search for observable effects of the change in hadron structure in-medium – notably the EMC effect



I. Insights into nuclear structure

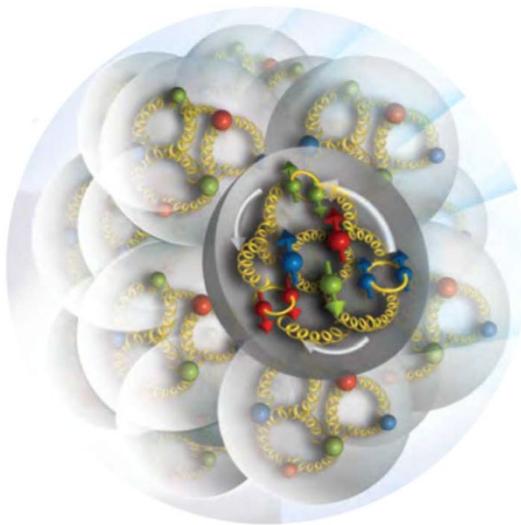
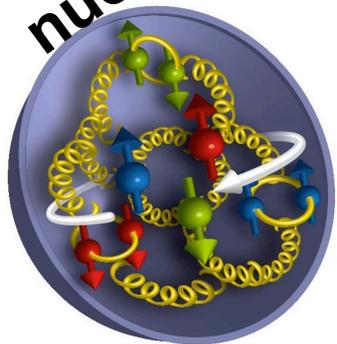
– what is the atomic nucleus?

There are two very different extremes....

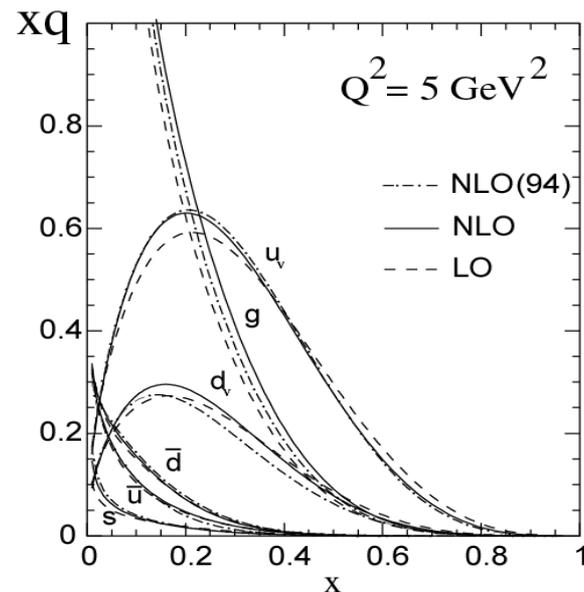
A. Nuclear Femtography

Science of mapping the position and motion of quarks and gluons in the nucleus.

Artist's Conception
of Quark and Gluons
in a proton and
nucleus



.. is just beginning



12 GeV

REQUIRES:

- High beam polarization
- High electron current
- High target polarization
- Large solid angle spectrometers

B. Extreme Chiral Effective Field Theory

- “Considering quarks is in contrast to our **modern understanding of nuclear physics...** the basic degrees of freedom of QCD (quarks and gluons) have to be considered only at higher energies. The energies relevant for nuclear physics are only a few MeV”

- anonymous referee 2017

TRUE

OR



?

- Actually not so modern.....

D. Alan Bromley (Yale) to Stan Brodsky in 1982

“Stan, you have to understand -- in nuclear physics we are only interested in how protons and neutrons make up a nucleus.

We are not interested in what is inside of a proton.”



Like this beautiful scene – very relaxing

D. Alan Bromley (Yale) to Stan Brodsky in 1982

“Stan, you have to understand -- in nuclear physics we are only interested in how protons and neutrons make up a nucleus.

We are not interested in what is inside of a proton.”



Moral: A comfortable picture is not necessarily the right one.....

What do we know?

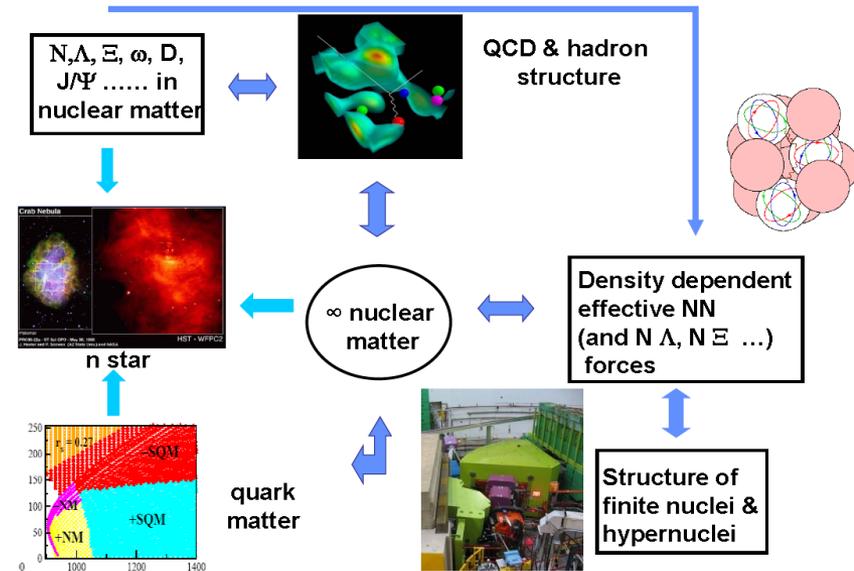
- Since 1970s, intermediate range NN attraction is strong Lorentz scalar
- In relativistic treatments (RHF, RBHF, QHD...) this leads to mean scalar field ~ 300 to 500 MeV!!
- This is not small – up to half the nucleon mass
 - death of “wrong energy scale” arguments
- Largely cancelled by large vector mean field BUT these have totally different dynamics: ω^0 just shifts energies, σ seriously modifies internal hadron dynamics

Suggests a different approach : QMC Model

(Guichon, Saito, Tsushima et al., Rodionov et al.

- see Saito et al., Prog. Part. Nucl. Phys. 58 (2007) 1 and
Prog. Part. Nucl. Phys. 100 (2018) 262-297 for reviews)

- Start with quark model (MIT bag/NJL...) for all hadrons
- Introduce a relativistic Lagrangian with σ , ω and ρ mesons coupling to non-strange quarks
- Hence **only 3 parameters** (if σ mass fixed)
 - determine by fitting to:
 - ρ_0 , E/A and symmetry energy
 - same in dense matter & finite nuclei
- Must solve self-consistently for the internal structure of baryons in-medium



Self-consistent solution of nuclear matter

$$[i\gamma^\mu\partial_\mu - (m_q - g_\sigma q\bar{\sigma}) - \gamma^0 g_\omega q\bar{\omega}] \psi = 0$$

Source of σ
changes:

$$\int_{Bag} d\vec{r} \bar{\psi}(\vec{r}) \psi(\vec{r})$$

SELF-CONSISTENCY

and hence mean scalar field changes...

and hence quark wave function changes....

**THIS PROVIDES A NATURAL SATURATION MECHANISM
(VERY EFFICIENT BECAUSE QUARKS ARE ALMOST MASSLESS)**

**source is suppressed as mean scalar field increases
(i.e. as density increases)**

Quark-Meson Coupling Model (QMC): Role of the Scalar Polarizability of the Nucleon

The response of the nucleon internal structure to the scalar field is of great interest... and importance

$$M^*(\mathbf{r}) = M - g_\sigma \sigma(\mathbf{r}) + \frac{d}{2} (g_\sigma \sigma(\mathbf{r}))^2$$

Non-linear dependence through the scalar polarizability
 $d \sim 0.22 R$ in original QMC (MIT bag)

Indeed, in nuclear matter at mean-field level (e.g. QMC), this is the **ONLY** place the response of the internal structure of the nucleon enters.

Application to nuclear structure

Derivation of Density Dependent Effective Force

Physical origin of density dependent forces of Skyrme type within the quark meson coupling model

P.A.M. Guichon ^{a,*}, H.H. Matevosyan ^{b,c}, N. Sandulescu ^{a,d,e},
A.W. Thomas ^b

Nuclear Physics A 772 (2006) 1–19

- **Start with classical theory of MIT-bag nucleons with structure modified in medium to give $M_{\text{eff}}(\sigma)$.**
- **Quantise nucleon motion (non-relativistic), expand in powers of derivatives**
- **Derive equivalent, local energy functional:**

$$\langle H(\vec{r}) \rangle = \rho M + \frac{\tau}{2M} + \mathcal{H}_0 + \mathcal{H}_3 + \mathcal{H}_{\text{eff}} + \mathcal{H}_{\text{fin}} + \mathcal{H}_{\text{so}}$$

Derivation of effective Force (cont.)

$$\mathcal{H}_0 + \mathcal{H}_3 = \rho^2 \left[\frac{-3G_\rho}{32} + \frac{G_\sigma}{8(1 + d\rho G_\sigma)^3} - \frac{G_\sigma}{2(1 + d\rho G_\sigma)} + \frac{3G_\omega}{8} \right] \\ + (\rho_n - \rho_p)^2 \left[\frac{5G_\rho}{32} + \frac{G_\sigma}{8(1 + d\rho G_\sigma)^3} - \frac{G_\omega}{8} \right],$$

$$\mathcal{H}_{\text{eff}} = \left[\left(\frac{G_\rho}{8m_\rho^2} - \frac{G_\sigma}{2m_\sigma^2} + \frac{G_\omega}{2m_\omega^2} + \frac{G_\sigma}{4M_N^2} \right) \rho_n + \left(\frac{G_\rho}{4m_\rho^2} + \frac{G_\sigma}{2M_N^2} \right) \rho_p \right] \tau_n \\ + p \leftrightarrow n,$$

$$\mathcal{H}_{\text{fin}} = \left[\left(\frac{3G_\rho}{32m_\rho^2} - \frac{3G_\sigma}{8m_\sigma^2} + \frac{3G_\omega}{8m_\omega^2} - \frac{G_\sigma}{8M_N^2} \right) \rho_n \right. \\ \left. + \left(\frac{-3G_\rho}{16m_\rho^2} - \frac{G_\sigma}{2m_\sigma^2} + \frac{G_\omega}{2m_\omega^2} - \frac{G_\sigma}{4M_N^2} \right) \rho_p \right] \nabla^2(\rho_n) + p \leftrightarrow n,$$

$$\mathcal{H}_{\text{SO}} = \nabla \cdot J_n \left[\left(\frac{-3G_\sigma}{8M_N^2} - \frac{3G_\omega(-1 + 2\mu_s)}{8M_N^2} - \frac{3G_\rho(-1 + 2\mu_v)}{32M_N^2} \right) \rho_n \right. \\ \left. + \left(\frac{-G_\sigma}{4M_N^2} + \frac{G_\omega(1 - 2\mu_s)}{4M_N^2} \right) \rho_p \right] + p \leftrightarrow n.$$

**Spin-orbit
force
predicted!**

Note the totally new, subtle density dependence

Systematic approach to finite nuclei

J.R. Stone, P.A.M. Guichon, P. G. Reinhard & A.W. Thomas:
(Phys Rev Lett, 116 (2016) 092501)

- **Constrain 3 basic quark-meson couplings ($g_\sigma^q, g_\omega^q, g_\rho^q$) so that nuclear matter properties are reproduced within errors**

$$-17 < E/A < -15 \text{ MeV}$$

$$0.14 < \rho_0 < 0.18 \text{ fm}^{-3}$$

$$28 < S_0 < 34 \text{ MeV}$$

$$L > 20 \text{ MeV}$$

$$250 < K_0 < 350 \text{ MeV}$$

- **Fix at overall best description of finite nuclei (+2 pairing pars)**
- **Benchmark comparison: SV-min 16 parameters (11+5 pairing)**

Overview of 106 Nuclei Studied – Across Periodic Table

Element	Z	N	Element	Z	N
C	6	6 - 16	Pb	82	116 - 132
O	8	4 - 20	Pu	94	134 - 154
Ca	20	16 - 32	Fm	100	148 - 156
Ni	28	24 - 50	No	102	152 - 154
Sr	38	36 - 64	Rf	104	152 - 154
Zr	40	44 - 64	Sg	106	154 - 156
Sn	50	50 - 86	Hs	108	156 - 158
Sm	62	74 - 98	Ds	110	160
Gd	64	74 - 100			

Not fit

N	Z	N	Z
20	10 - 24	64	36 - 58
28	12 - 32	82	46 - 72
40	22 - 40	126	76 - 92
50	28 - 50		

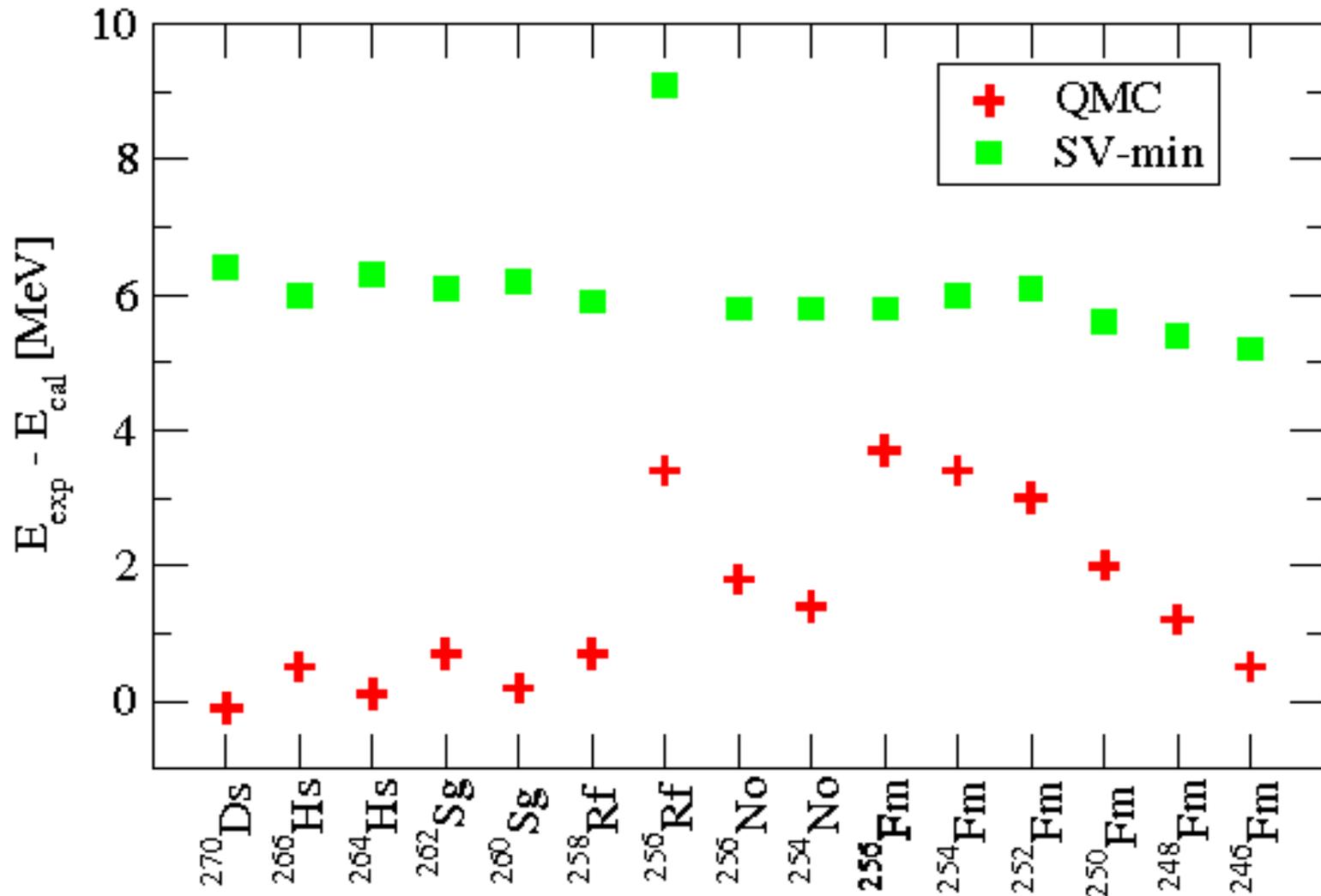
i.e. We look at most challenging cases of p- or n-rich nuclei

Overview

data	rms error %	
	QMC	SV-min
fit nuclei:		
binding energies	<u>0.36</u>	0.24
diffraction radii	1.62	0.91
surface thickness	10.9	2.9
rms radii	0.71	0.52
pairing gap (n)	57.6	17.6
pairing gap (p)	25.3	15.5
1s splitting: proton	15.8	18.5
1s splitting: neutron	20.3	16.3
superheavy nuclei:		
	<u>0.1</u>	0.3
N=Z nuclei	1.17	0.75
mirror nuclei	1.50	1.00
other	0.35	0.26

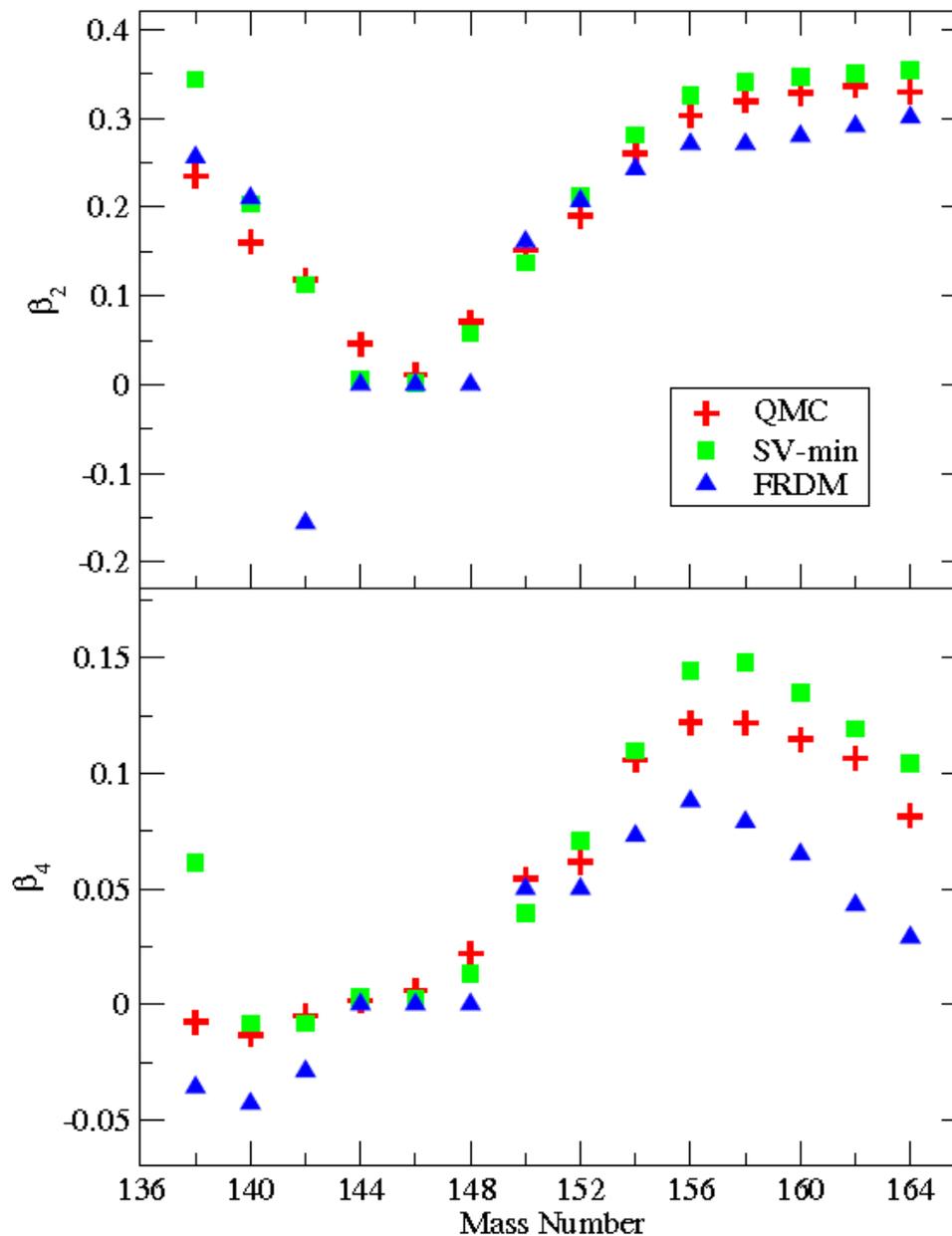
Stone et al., PRL (2016)

Superheavy Binding : 0.1% accuracy

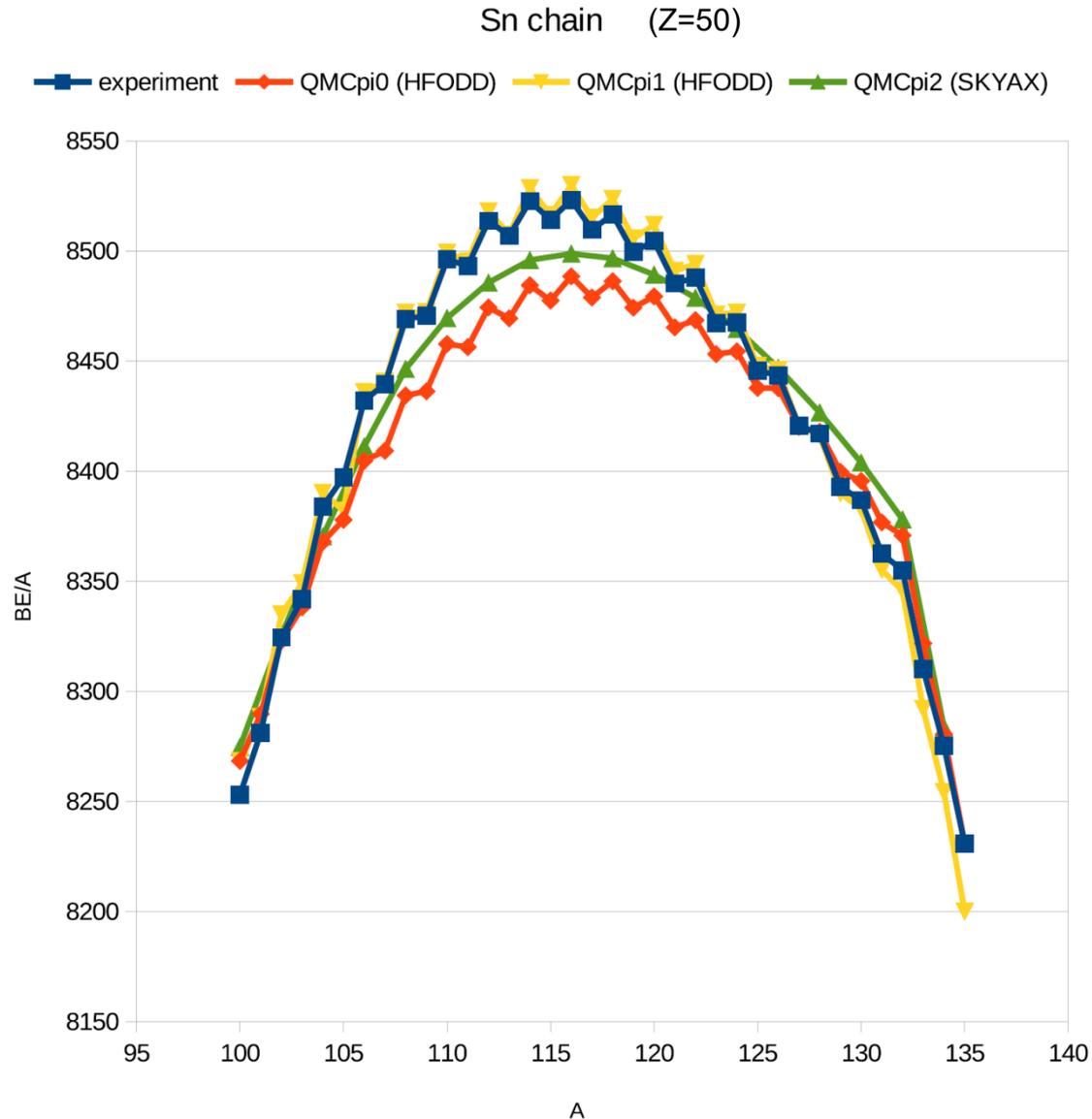


Stone et al., PRL 116 (2016) 092501

Deformation in Gd (Z=64) Isotopes



Martinez, Konieczka, Bąszyk *et al.* – HFODD Implementation

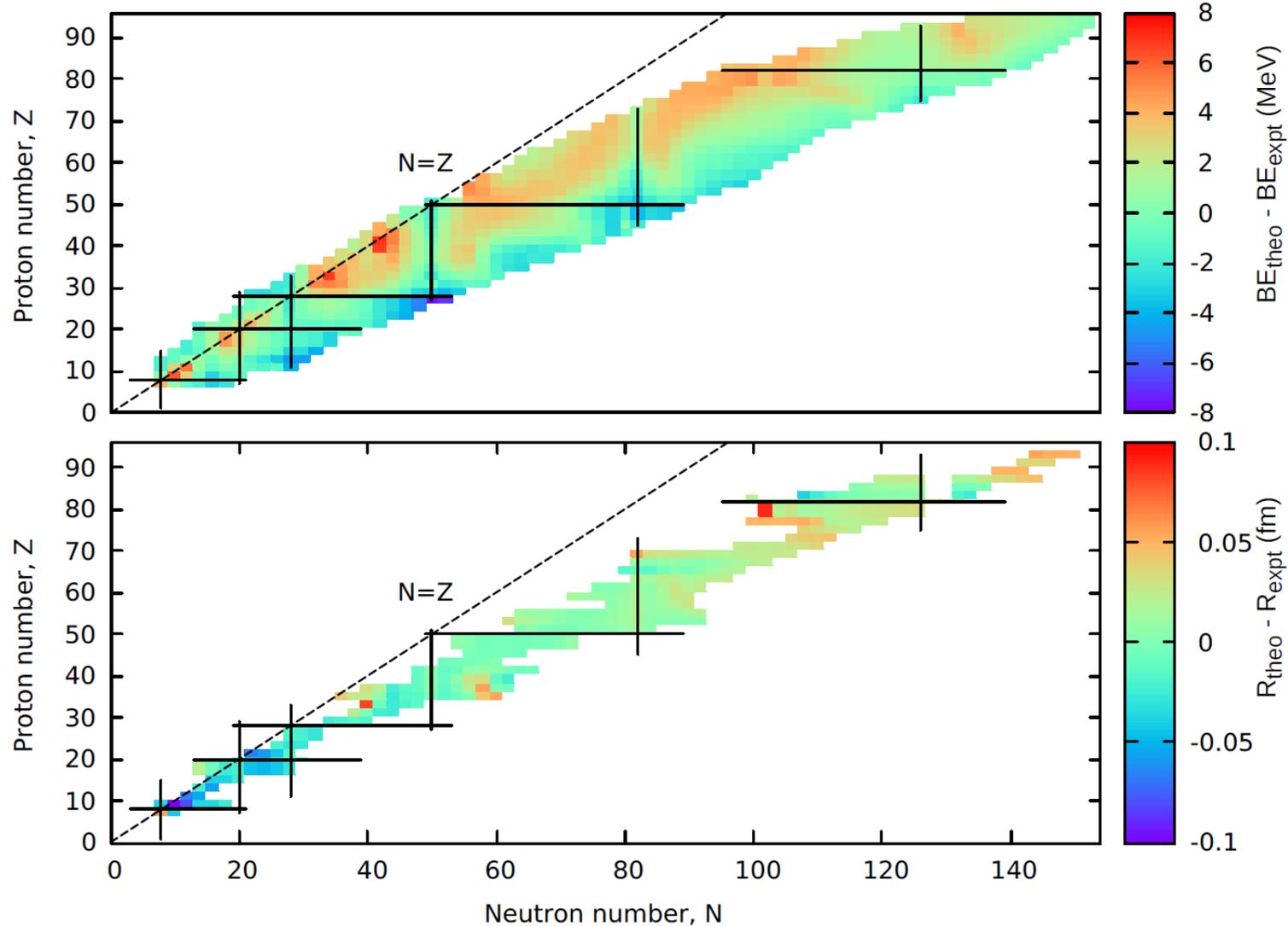


Publication in preparation....

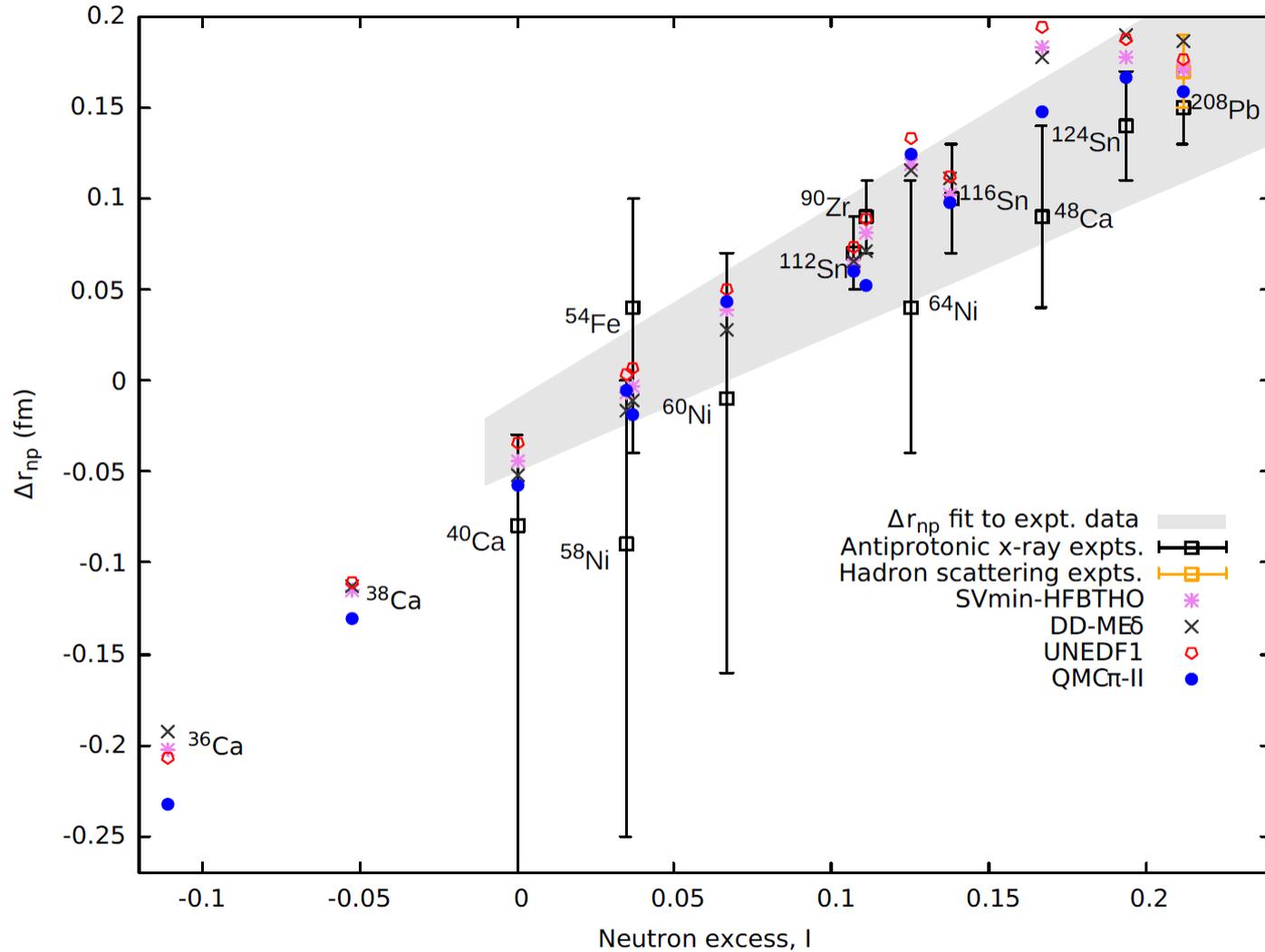
Most recent development

- Includes pion Fock terms
- Includes σ^3 term – now EDF has 5 free parameters

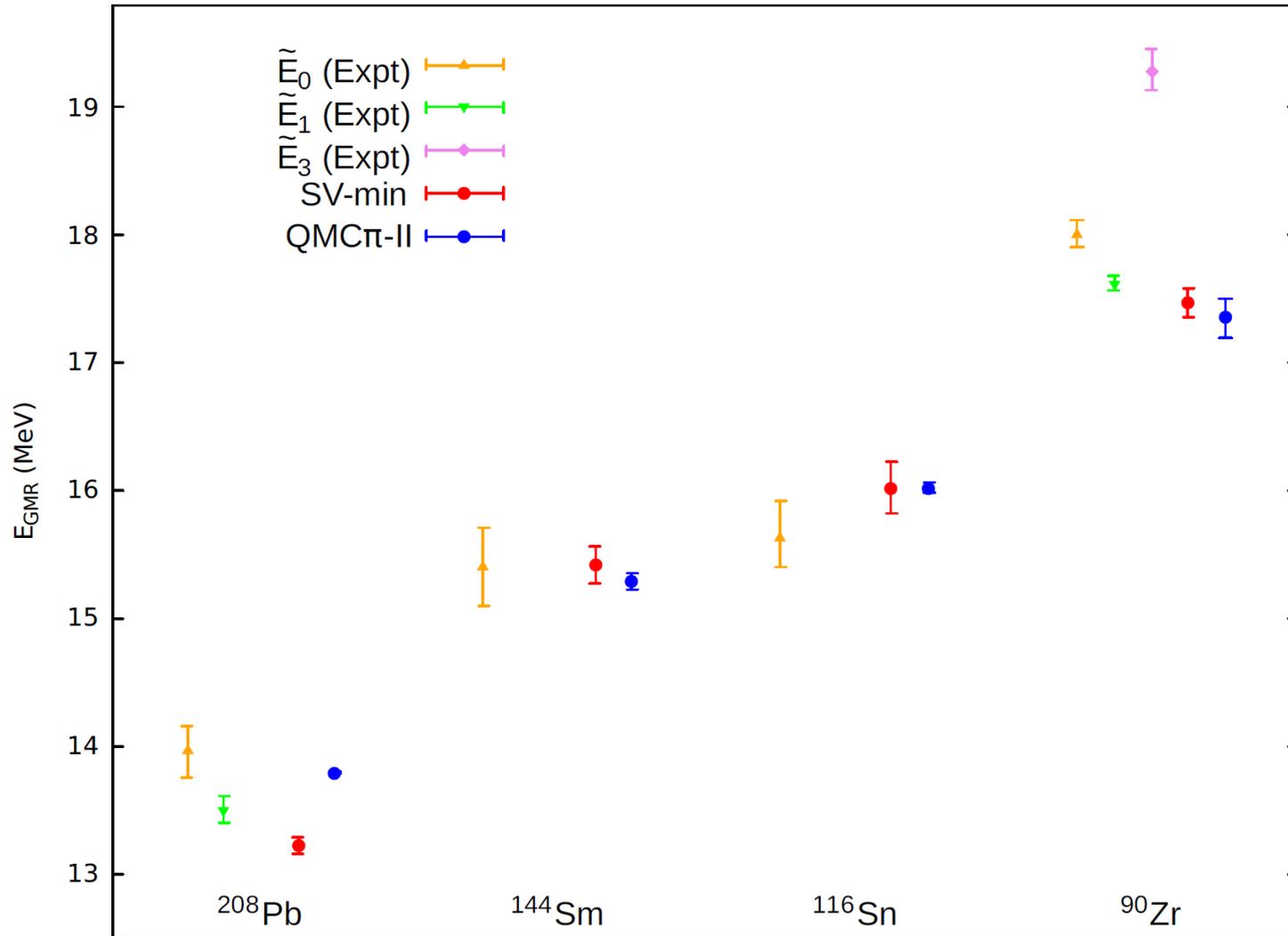
Overview of 739 even-even nuclei



Neutron-proton radius difference



Giant Monopole Resonance



Martinez et al., arXiv:1811.06628

Summary: Finite Nuclei

- The effective force was *derived* at the quark level *based upon changing structure of bound nucleon*
- Has many less parameters but reproduces nuclear properties at a level comparable with the best phenomenological Skyrme forces
- Looks like standard nuclear force
- BUT underlying theory also predicts modified internal structure and hence modified
 - DIS structure functions
 - elastic form factors.....

Nuclear DIS Structure Functions : The EMC Effect

To address questions like this one **MUST** start with a theory that quantitatively describes nuclear structure and allows calculation of structure functions
– very, very few examples.....

Theoretical Understanding

- Still numerous proposals but few consistent theories
- Initial studies used MIT bag¹ to estimate effect of self-consistent change of structure in-medium – but better to use a covariant theory
- For that Bentz and Thomas² re-derived change of nucleon structure in-medium in the NJL model
- This set the framework for sophisticated studies by Bentz, Cloët and collaborators over the last decade

¹ Thomas, Michels, Schreiber and Guichon, Phys. Lett. B233 (1989) 43

² Bentz and Thomas, Nucl. Phys. A696 (2001) 138

EMC Effect for Finite Nuclei

(There is also a spin dependent EMC effect - as large as unpolarized)

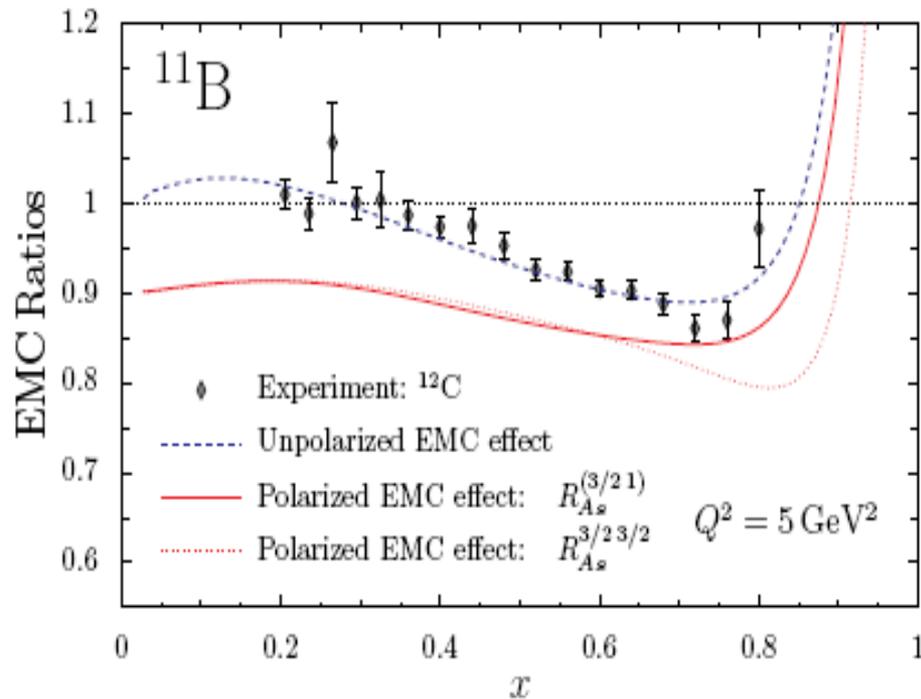


FIG. 7: The EMC and polarized EMC effect in ^{11}B . The empirical data is from Ref. [31].

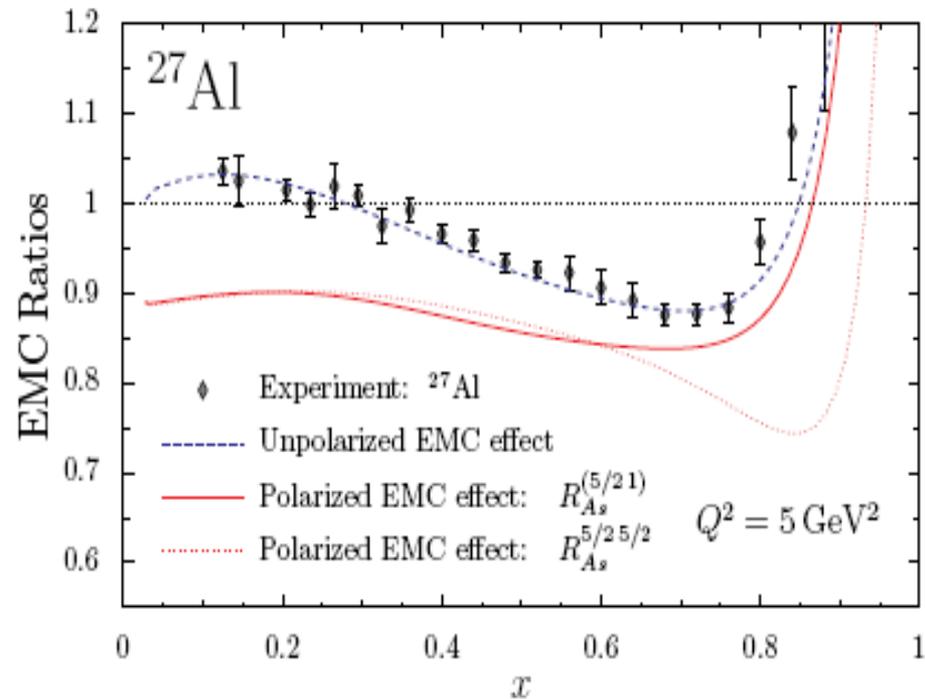
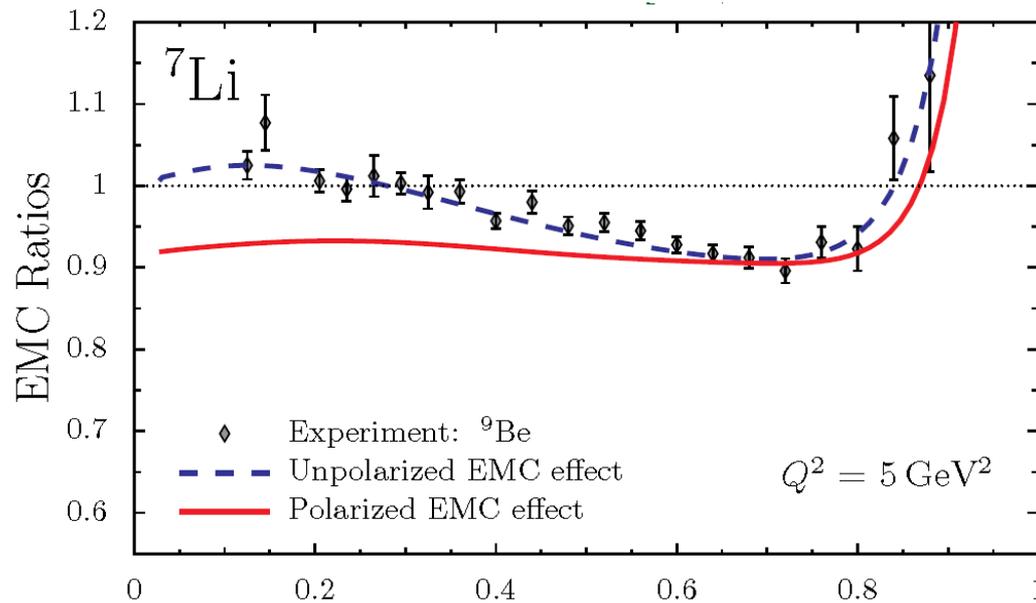


FIG. 9: The EMC and polarized EMC effect in ^{27}Al . The empirical data is from Ref. [31].

Cloët, Bentz & Thomas, Phys. Lett. B642 (2006) 210
(nucl-th/0605061)

Approved JLab Experiment

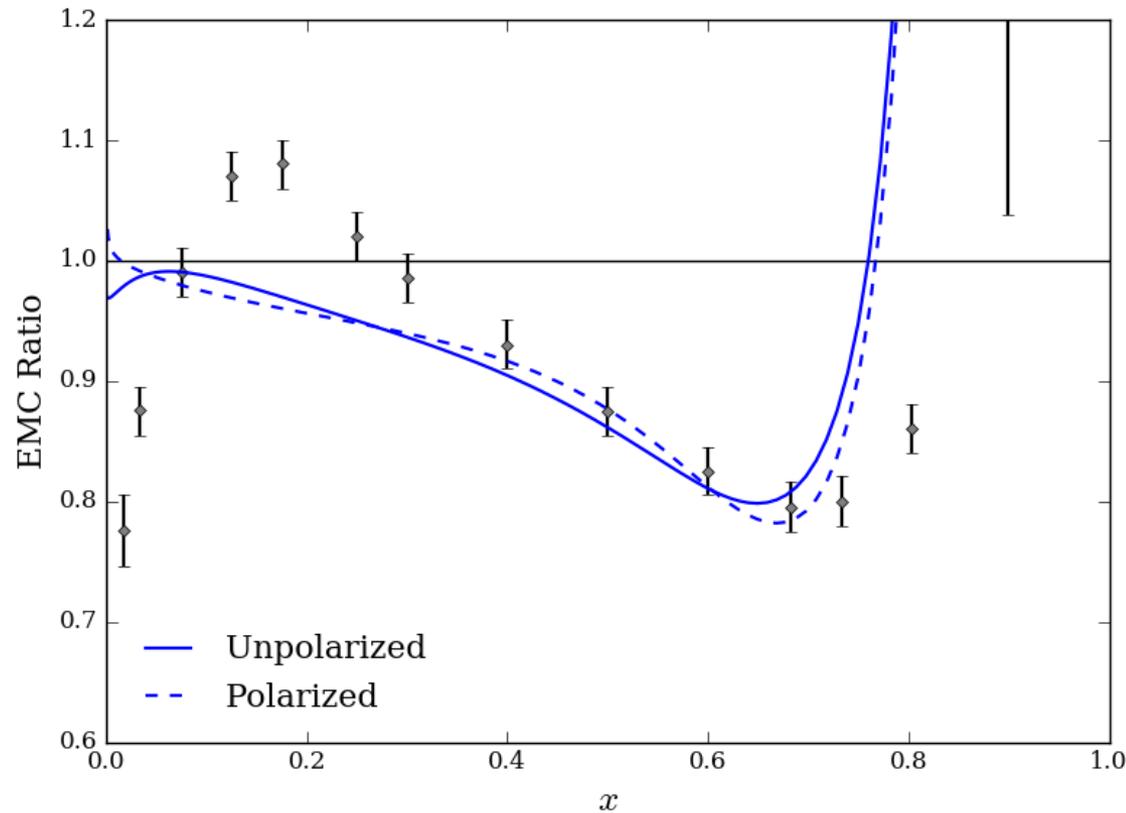
- Effect in ${}^7\text{Li}$ is slightly suppressed because it is a light nucleus and proton does not carry all the spin (simple WF: $P_p = 13/15$ & $P_n = 2/15$)
- Experiment now approved at JLab [E12-14-001] to measure spin structure functions of ${}^7\text{Li}$ (GFMC: $P_p = 0.86$ & $P_n = 0.04$)
- *Everyone with their favourite explanation for the EMC effect should make a prediction for the polarized EMC effect in ${}^7\text{Li}$*



Other tests (e.g. Isovector EMC effect)

Model dependence of spin-EMC effect

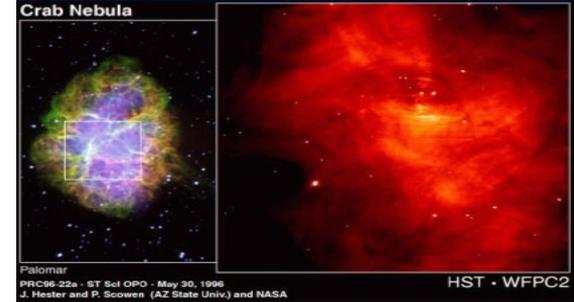
Went back to QMC, with defects of bag model (especially too small at large- x). Simply examine, without details of nuclear structure, at ρ_0 , how the polarized EMC effect compares with the unpolarized effect.



Spin-EMC Effect is a crucial test

- **Tensor correlations leading to high momentum components in nuclear wave function have been proposed as an alternate explanation of the EMC effect**
- **The tensor force scatters 3S_1 pairs almost entirely into 3D_1 at high momentum ($\sim 84\%$ at $p > 400$ MeV/c)**
- **Nucleons in SRC are depolarized – simple Clebsch-Gordan coefficients - and cannot contribute to spin-EMC effect**
- **That is SRC predicts essentially NO spin-EMC effect**

I. Summary



- Intermediate range NN attraction is **STRONG Lorentz scalar**
- This modifies the intrinsic structure of the bound nucleon
 - profound change in shell model :
what occupies shell model states are **NOT** free nucleons
- Scalar polarizability is a natural source of three-body forces (NNN, HNN, HHN...)
 - clear physical interpretation
- Naturally generates effective HN and HNN forces with no new parameters and predicts heavy neutron stars

II. Summary

- **Initial systematic study of finite nuclei very promising**
 - **Binding energies typically within 0.3% across periodic table**
 - **Super-heavies ($Z > 100$) especially good**
- **Need empirical confirmation:**
 - **Response Functions & Coulomb sum rule (soon?)**
 - **Spin EMC – definitive test of SRC proposal**

Special Mentions.....



Guichon



Tsushima



Saito



Stone



Krein



Bentz



Matevosyan



Cloët



Whittenbury



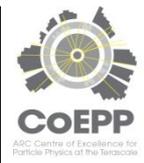
Simenel



Martinez



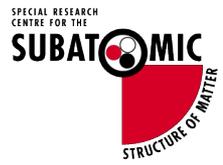
Motta



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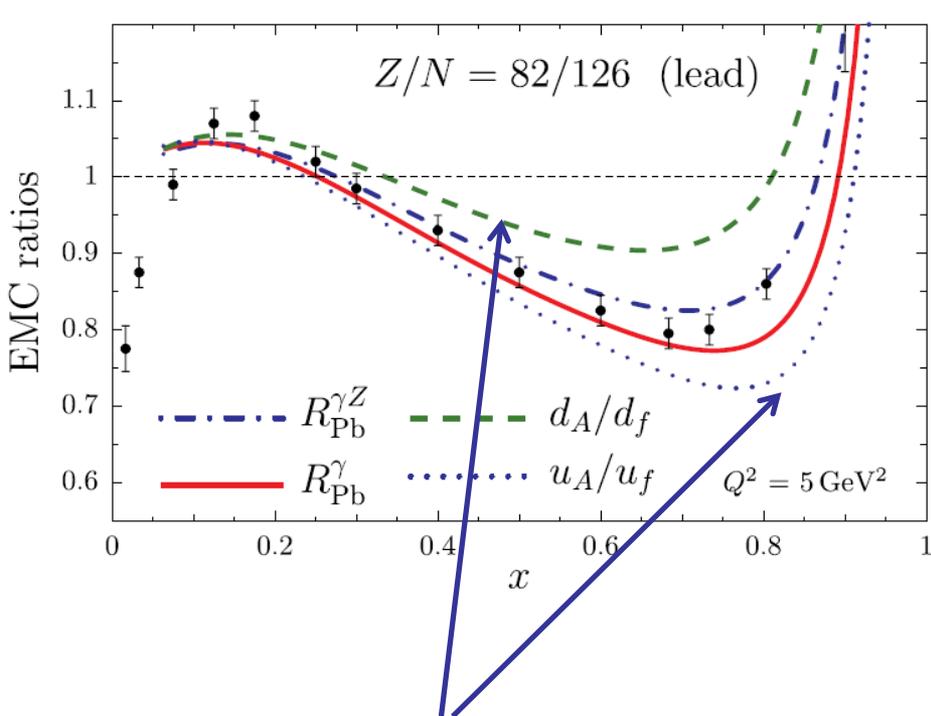
Isovector EMC Effect

- New realization concerning EMC effect in this approach:
 - isovector force in nucleus (like Fe) with $N \neq Z$ effects ALL u and d quarks in the nucleus
 - subtracting structure functions of extra neutrons is not enough
 - *there is a shift of momentum from all u to all d quarks*
- Sign and magnitude of this effect exhibits little model dependence

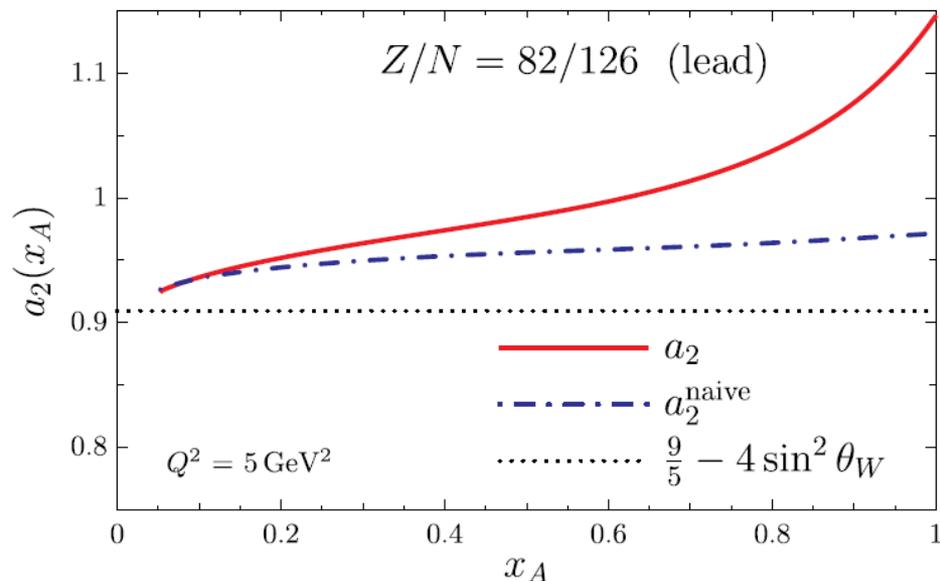
Cloet *et al.*, Phys.Rev.Lett.102:252301,2009
Londergan et al., Phys Rev D67 (2003) 111901

Parity-Violating Deep Inelastic Scattering and the Flavor Dependence of the EMC Effect

I. C. Cloët,¹ W. Bentz,² and A. W. Thomas¹



$$A_{\text{PV}} = \frac{G_F Q^2}{4\sqrt{2}\pi\alpha_{\text{em}}} \left[a_2(x_A) + \frac{1 - (1 - y)^2}{1 + (1 - y)^2} a_3(x_A) \right]$$



Ideally tested at EIC with CC reactions

Parity violating EMC will test this at JLab 12 GeV

Effect of scalar field on quark spinor

- MIT bag model: quark spinor modified in bound nucleon

$$\psi = \frac{\mathcal{N}}{4\pi} \begin{pmatrix} j_0(xu'/R_B) \\ i\beta_q \vec{\sigma} \cdot \hat{u}' j_1(xu'/R_B) \end{pmatrix} \chi_m$$

- Lower component enhanced by attractive scalar field

$$\beta_q = \sqrt{\frac{\Omega_0 - m_q^* R_B}{\Omega_0 + m_q^* R_B}}$$

- This leads to a *very small* ($\sim 1\%$ at ρ_0) *increase in bag radius*
- It also *suppresses the scalar coupling to the nucleon as the scalar field increases*

$$\frac{\Omega_0/2 + m_q^* R_B (\Omega_0 - 1)}{\Omega_0 (\Omega_0 - 1) + m_q^* R_B / 2} = \int \bar{\psi} \psi \, dV$$

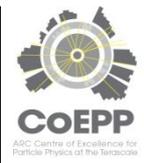
- This is the “scalar polarizability”: a new saturation mechanism for nuclear matter

Key papers on QMC

- **Two major, recent papers:**
 1. Guichon, Matevosyan, Sandulescu, Thomas, Nucl. Phys. A772 (2006) 1.
 2. Guichon and Thomas, Phys. Rev. Lett. 93 (2004) 132502
- **Built on earlier work on QMC: e.g.**
 3. Guichon, Phys. Lett. B200 (1988) 235
 4. Guichon, Saito, Rodionov, Thomas, Nucl. Phys. A601 (1996) 349
- **Major review of applications of QMC to many nuclear systems:**
 5. Saito, Tsushima, Thomas, Prog. Part. Nucl. Phys. 58 (2007) 1-167 (hep-ph/0506314)

References to: Covariant Version of QMC

- **Basic Model: (Covariant, chiral, confining version of NJL)**
- **Bentz & Thomas, Nucl. Phys. A696 (2001) 138**
- **Bentz, Horikawa, Ishii, Thomas, Nucl. Phys. A720 (2003) 95**
- **Applications to DIS:**
- **Cloet, Bentz, Thomas, Phys. Rev. Lett. 95 (2005) 052302**
- **Cloet, Bentz, Thomas, Phys. Lett. B642 (2006) 210**
- **Applications to neutron stars – including SQM:**
- **Lawley, Bentz, Thomas, Phys. Lett. B632 (2006) 495**
- **Lawley, Bentz, Thomas, J. Phys. G32 (2006) 667**



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