Standard Model Explanations of the Short Baseline Anomalies

Anna Hayes
Los Alamos

CERN, Oct 2019

The majority of neutrino oscillation experiments converge on a consistent 3ν oscillation framework

The three lepton flavor eigenstates $\nu_\alpha=(\nu_e, \nu_\mu, \nu_\tau)$ are related to three mass eigenstates $\nu_\iota=(\nu_1, \nu_2, \nu_3)$ through a unitary transformation.

Requires that the collective set of experiments is consistent with:

[Takaaki Kajita](#)
Super-K

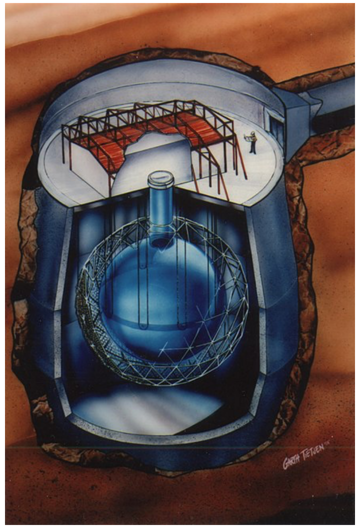
- Three mixing angles: $(\theta_{12}, \theta_{13}, \theta_{23})$;
- CP-violating phase δ
- Two mass differences: $\delta m^2=m_2^2 - m_1^2 > 0$

  $\Delta m^2 = m_{23} - (m_1^2 + m_2^2)/2$

[Art McDonald](#)
SNO

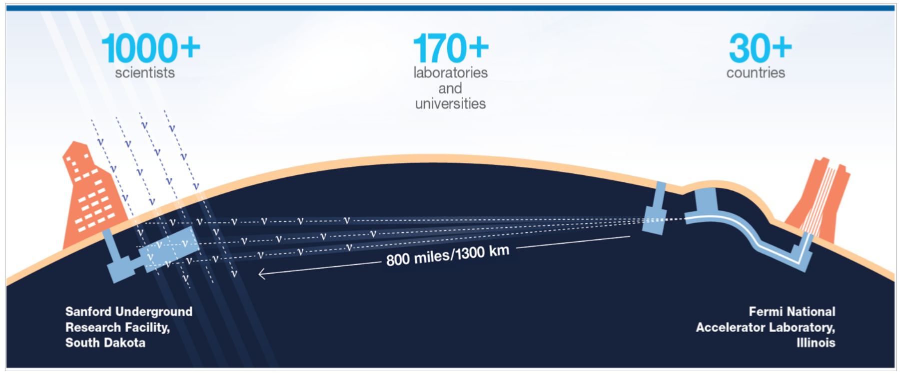
Kajita and McDonald shared the Nobel Prize in 2015

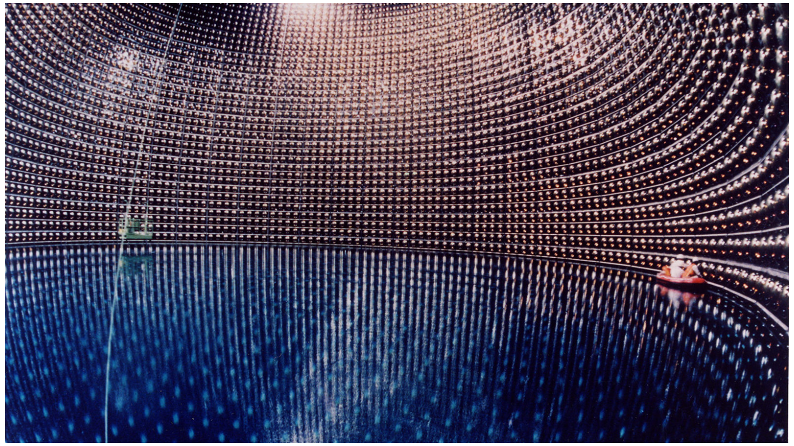
The mixing parameters are deduced from solar, atmospheric, accelerator, and reactor neutrino experiments.



$\delta m^2/10^{-5}$ eV^2	7.54	7.32 - 7.80
$\Delta m^2/10^{-3}$ eV^2	2.43 (2.38)	2.32 - 2.49
$\sin^2\theta_{12}/10^{-1}$	3.08	2.91 - 3.25
$\sin^2\theta_{13}/10^{-2}$	2.34 (2.40)	2.15 - 2.59
$\sin^2\theta_{23}/10^{-1}$	4.37 (4.55)	4.14 - 5.94
δ/π	1.39 (1.31)	0.98 - 1.77

F. Capozzi, et al. Phys. Rev. D **89**, 093018 2014





However Four Experimental Anomalies do not fit within the 3ν Mixing Picture

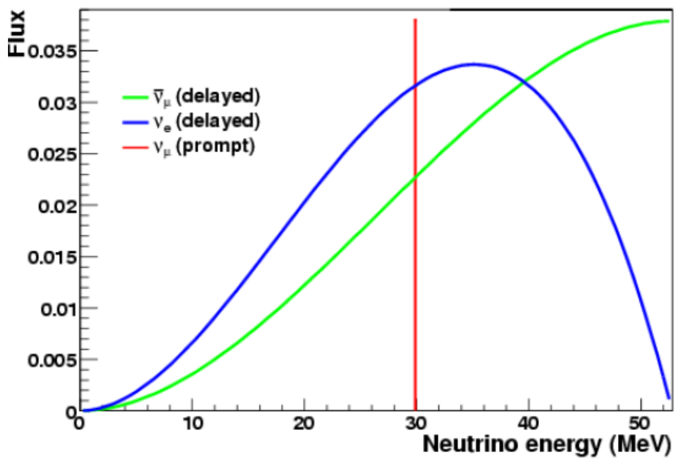
- LSND
- MiniBooNE
- The Gallium Anomaly
- The Short Base-Line Reactor Neutrino Anomaly

These anomalies possibly suggest a fourth sterile neutrino, requiring a mass on the 1 eV scale.

However, there are also complex nuclear physics issues associated with each anomaly.

LSND

LSND used neutrinos from accelerator produced stopped pions to search for neutrino oscillations with $\Delta m^2 \sim 1$ eV^2.



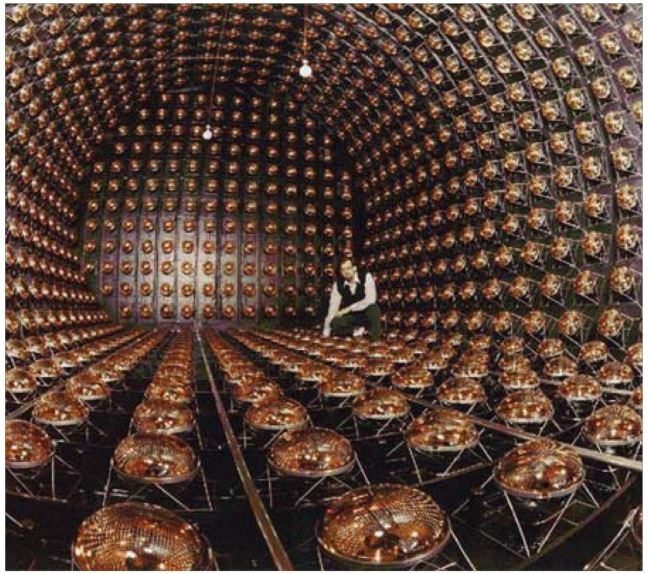
For two-state mixing:

$$P = \sin^2 2\theta \ \sin^2(1.27\Delta m^2 (L/E))$$

=> The detector was 30 m from the source and <E$_\nu$>~ 30 MeV.

800 MeV proton beam at LANSCE produces π^- (mostly get stopped) and π^+ that produce neutrinos

$$\pi^+ \to \nu_\mu \mu^+$$
$$\mu^+ \to \bar{\nu}_\mu \nu_e e^+$$



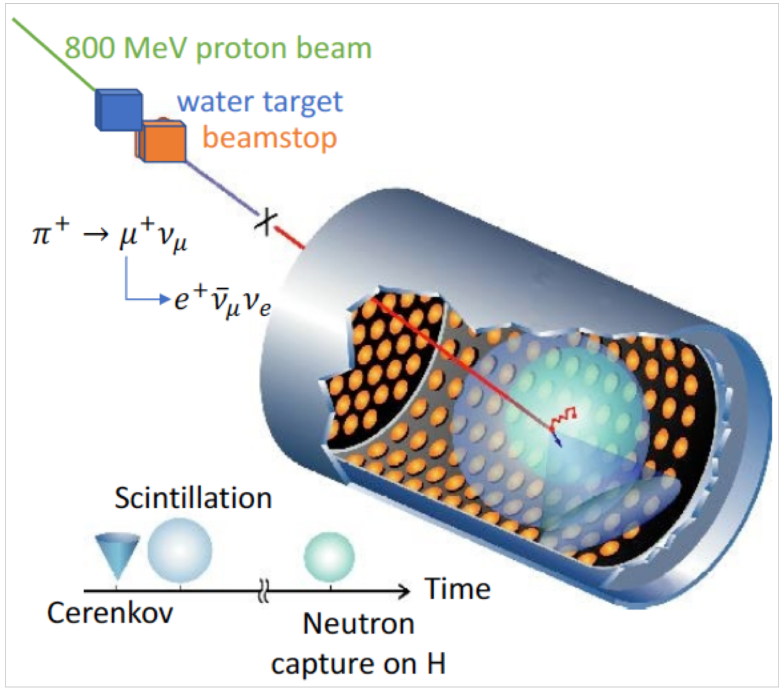
Searched for: $\bar{\nu}_\mu \to \bar{\nu}_e$

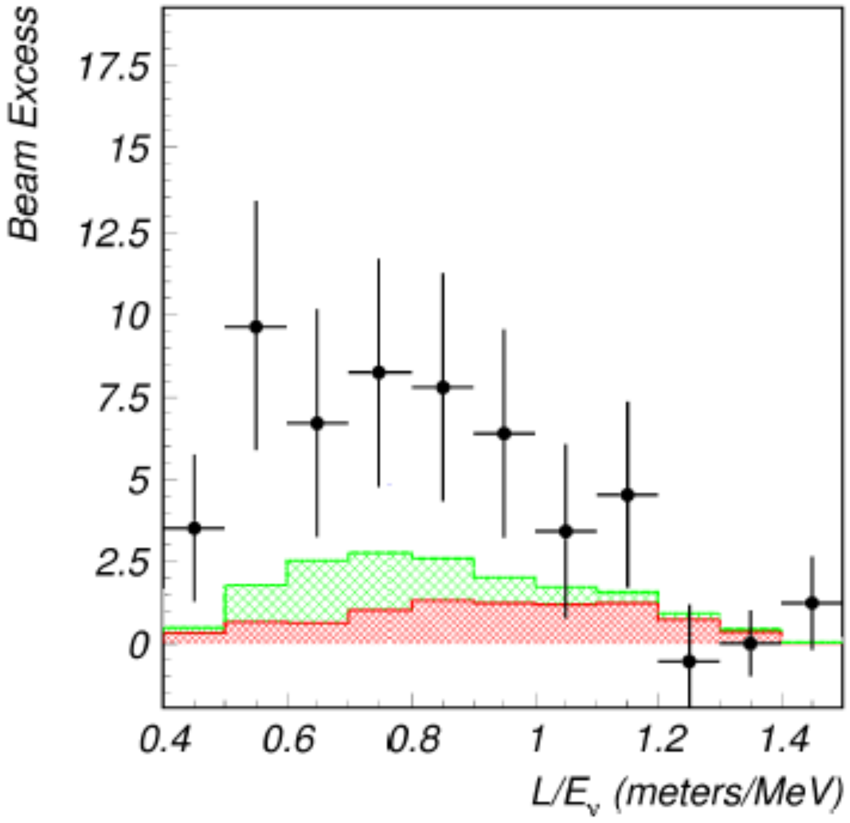
Detected via Inverse Beta Decay

$$\bar{\nu}_e + p \to n + e^+$$
$$n + p \to D + \gamma(2.2\ MeV)$$

Athanassopoulos et al., PRL. 75, 2650 (1995);PRL. 77, 3082 (1996) ;
PRL 81, 1774 (1998)

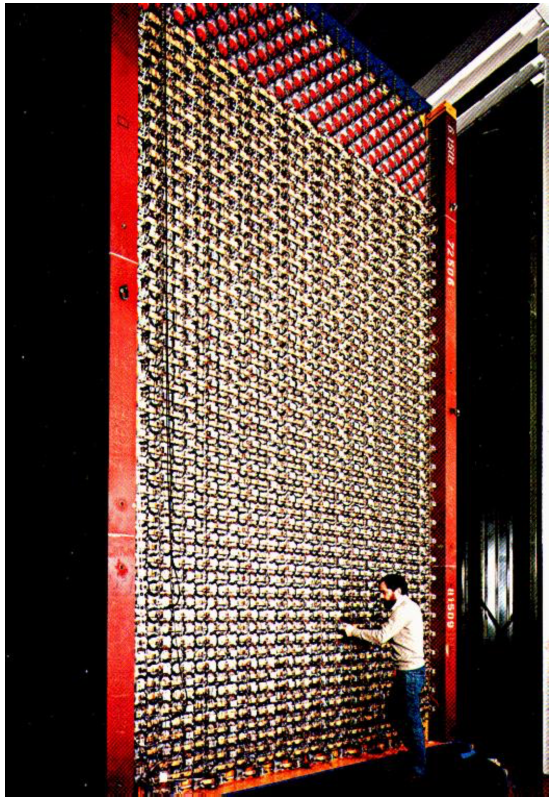
LSND Observed a 3.8σ excess

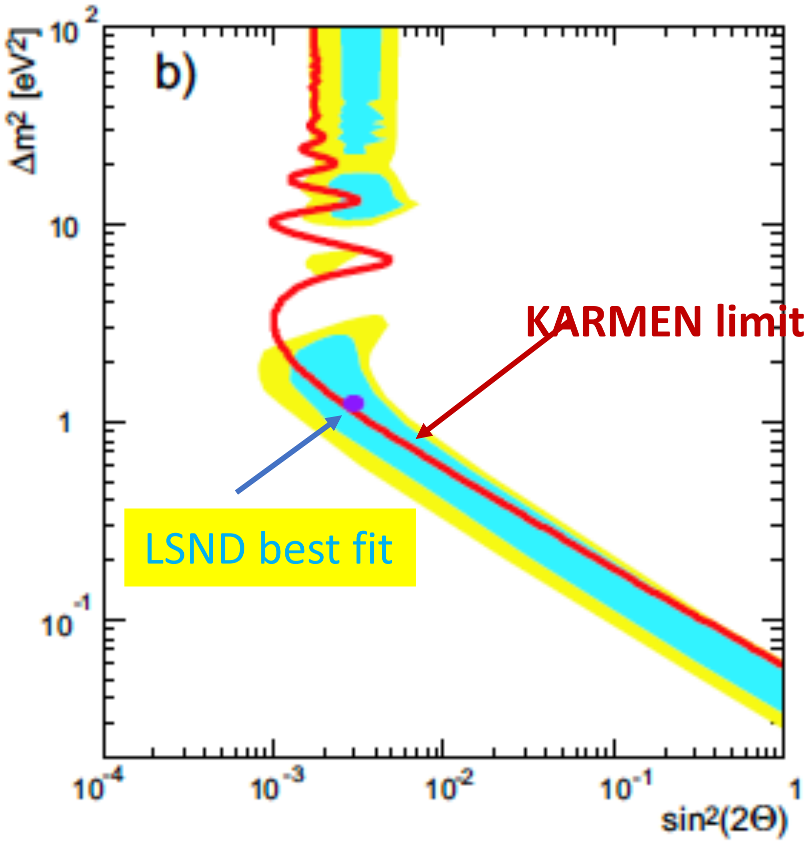




En-Chuan Huang, Neutrino 2018
Ann. Rev. Nucl. Part. Sci, 63, (1) 45

However, Karmen, at 17 m, did not see evidence for $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$





A combined analyses, with KARMEN as the near detector (17 m) and LSND as far detector (30 m), finds two possible solutions at the 64% confidence level: $\Delta m^2 \approx 7 eV^2/c^4$ or $\Delta m^2 < 1\ eV^2/c^4$

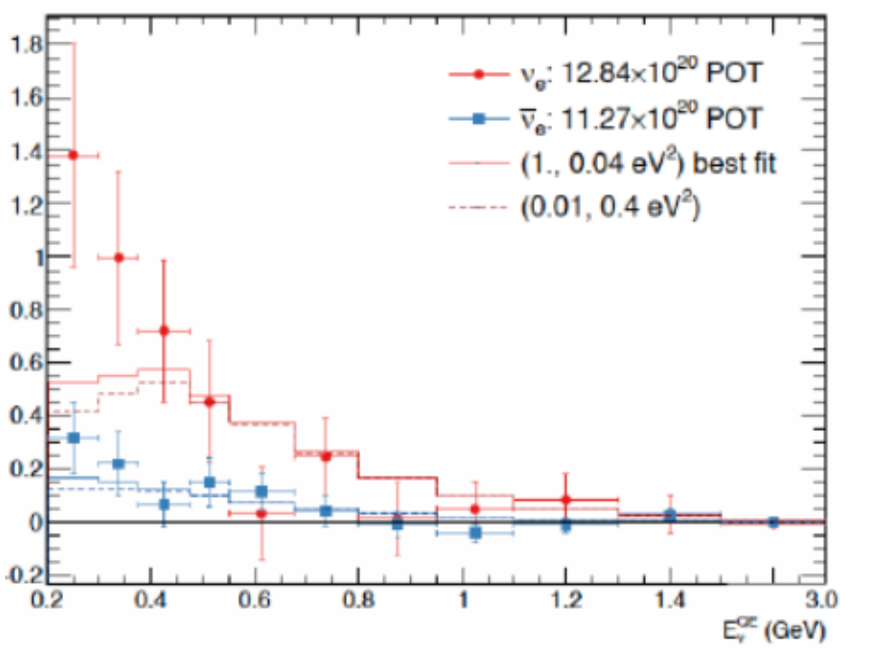
There is no known Standard Model explanation for LSND

MiniBooNE

Uses the Booster Neutrino Beam at Fermilab
Designed to test LSND , same L/E, but with <E>~ GeV, L=541 m

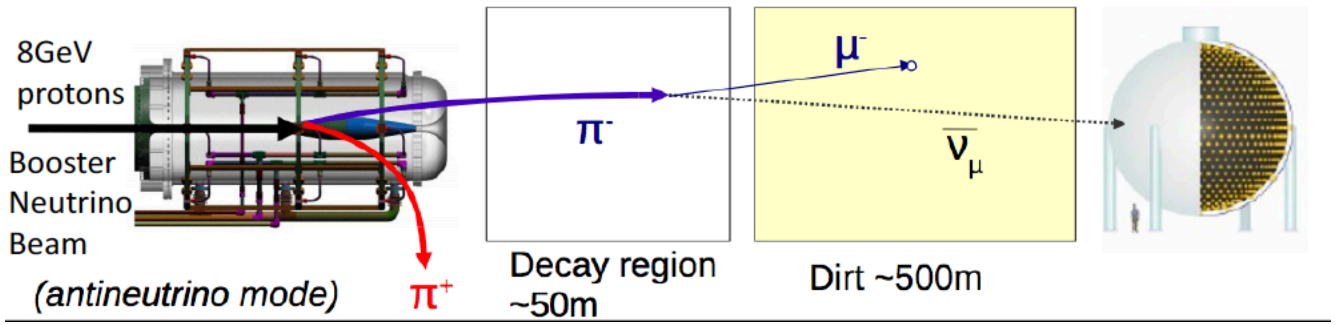
$$P = \sin^2 2\theta \ \sin^2(1.27\Delta m^2(L/E))$$

Searched for: $\nu_\mu \to \nu_e$ (or $\bar{\nu}_\mu \to \bar{\nu}_e$)



PRL, 98 (23): 231801

En-Chuan Huang, Neutrino 2018



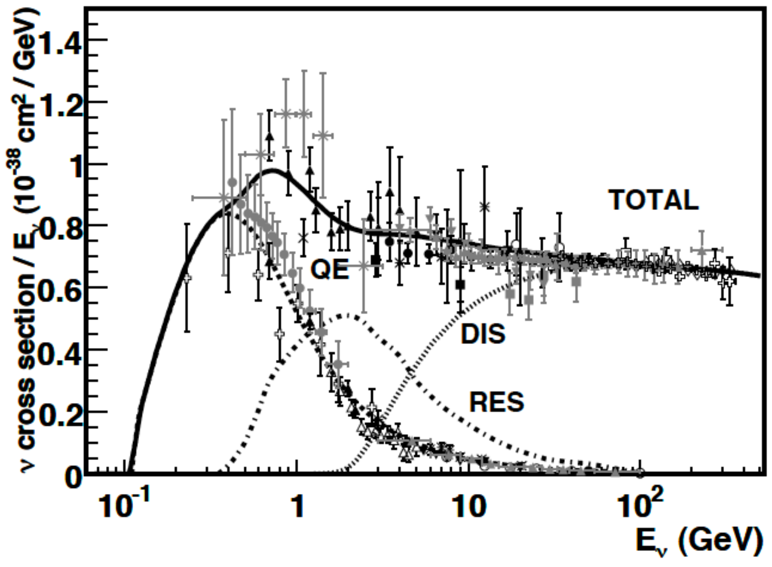
Observed an excess in both ν and $\bar{\nu}$ channels

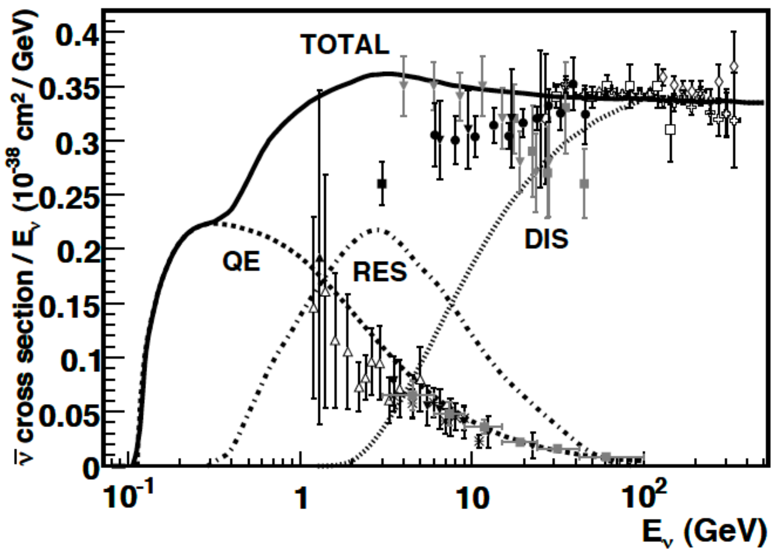
Conrad et al., Rev. Nucl. Part. Sci. 63, 45 (2013)

Kopp, JHEP05(2013)050; Gariazzo et al., JHEP, 06 (2017) 135

The Main Standard Model Issue is the imperfect knowledge of the cross sections and mis-identification of electrons versus gammas

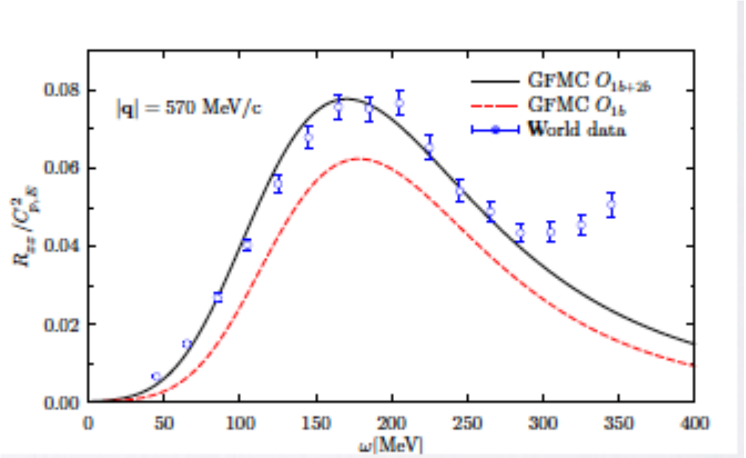
- MiniBooNE neutrino energies involve both quasi-elastic and resonance contributions.
- QE scattering involves the emission of nucleons for ^{12}C.
- RES scattering involve the production of N*, Δ and the emission of mesons and gammas.
- There is no unified first principle model to describe all of these on carbon.





Scaling arguments are used to constrain the ^{12}C cross sections from the nucleon cross sections.
Scaling works well for electron scattering, but also need to know the gamma emission for exclusive channels.

Cross section Issues: 2-body currents, Scaling with A, and misidentification of gammas as electrons



- Example, Lavato *et al.* find that 2-body currents enhance both the vector and axial contributions to the neutrino cross sections

arXiv:1509.00451 Phs Rev. Lett. 112 182502

- Unfortunately, these model currently don't cover the resonance region. They are also inclusive.

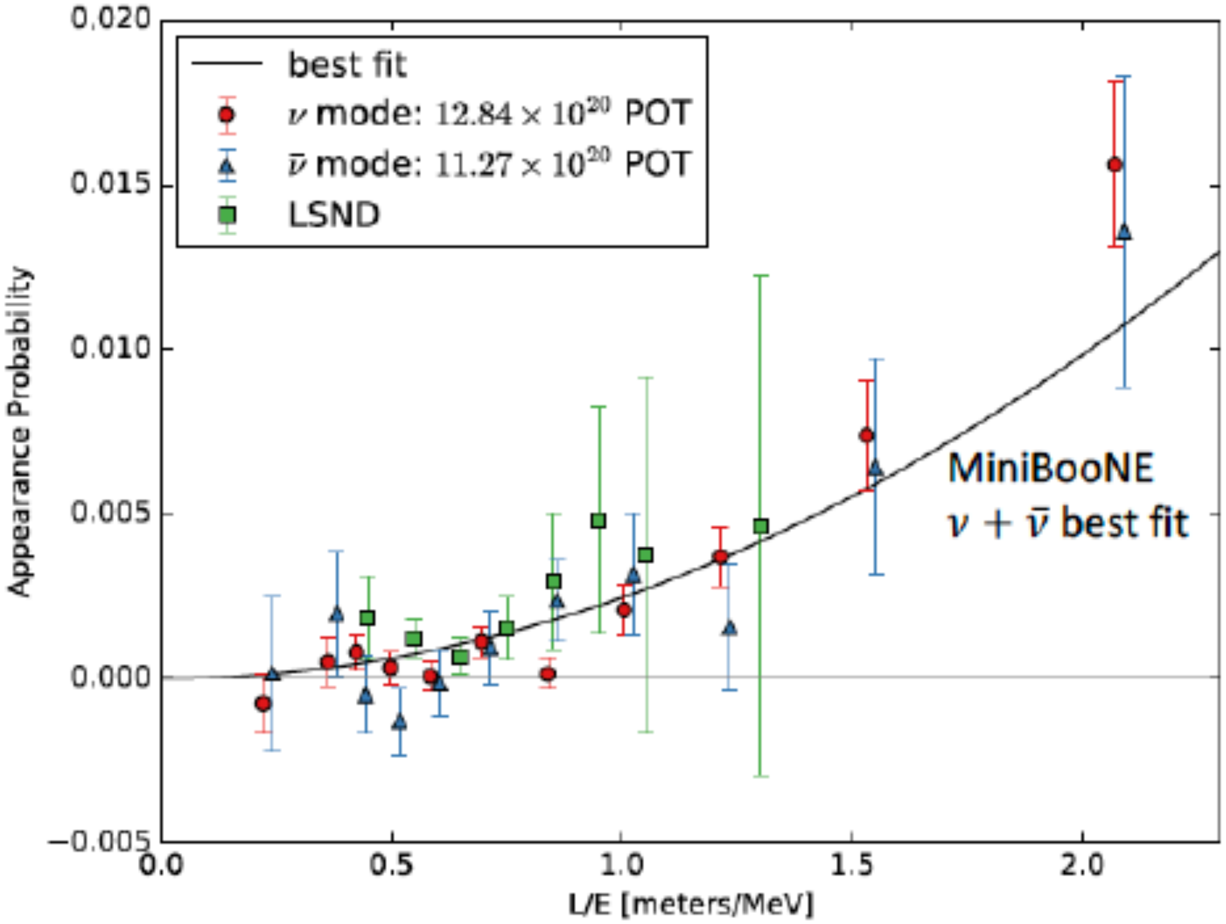
Status:

- A nuclear cross section model to explain away all of MiniBooNE's excesses has not yet been put forward.

- Misidentification of gammas from Δ and π decay ($\pi^0 \rightarrow 2\gamma$, Δ->γ+N) a suggested explanation,
- But most people agree that this is unlikely to be the explanation.

- The π^0/Δ ratio is known for the free particles, but not in the medium.

The two MiniBooNE signals ($\nu+\bar{\nu}$) and LSND are consistent with the same L/E appearance signature

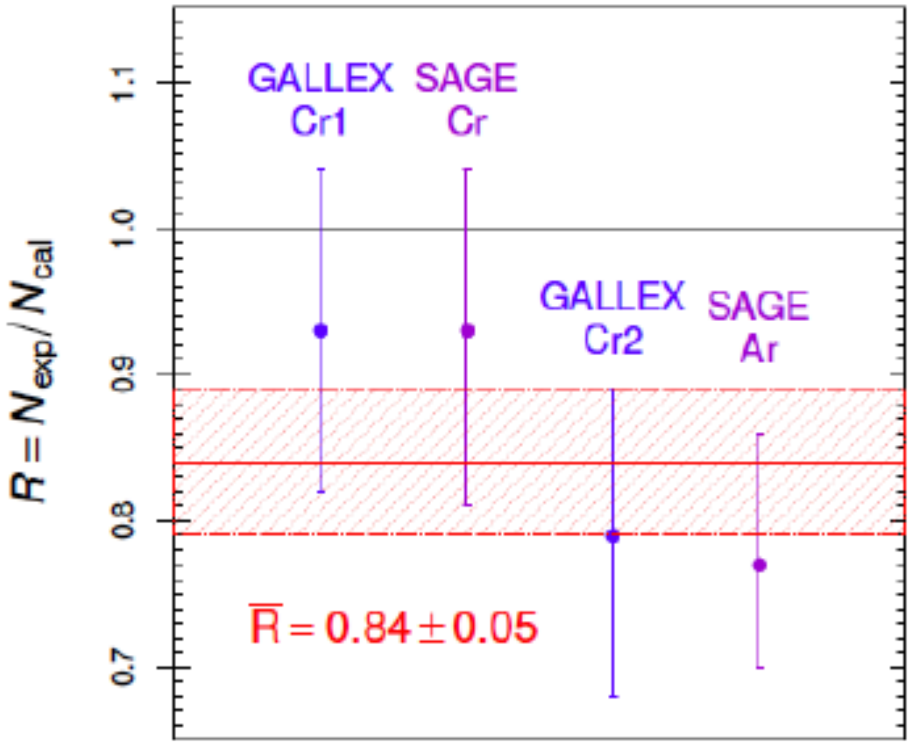


$$P = \sin^2 2\theta \ \sin^2(1.27\Delta m^2(L/E))$$

En-Chuan Huang, Neutrino 2018

The Gallium Anomaly

Monoenergetic neutrino sources used to test the SAGE and GALLEX detectors suggest too few neutrinos being detected.



$$e^- + {}^{51}Cr \rightarrow {}^{51}V + \nu_e \qquad E \simeq 0.75\,\mathrm{MeV}$$

$$e^- + {}^{37}Ar \rightarrow {}^{37}Cl + \nu_e \qquad E \simeq 0.81\,\mathrm{MeV}$$

$$\nu_e + {}^{71}Ga \rightarrow {}^{71}Ge + e^-$$

Expected comes from cross section based on $t_{1/2}$ of ^{71}Ge

Bahcall + Haxton

$\langle L \rangle_{GALLEX}$ = 1.9 m; $\langle L \rangle_{SAGE}$ = 0.6 m

=> $\Delta m^2_{\mathrm{SBL}} \gtrsim 0.5\,\mathrm{eV}^2$

Abdurashitov et al. (SAGE) 2006 PRC73 045805.; Anselmann et al. (GALLEX) 1995 PLB342; Hampel et al. (GALLEX) 1998 PLB420 114; Giunti, Phys. Rev. C **83**, 065504.

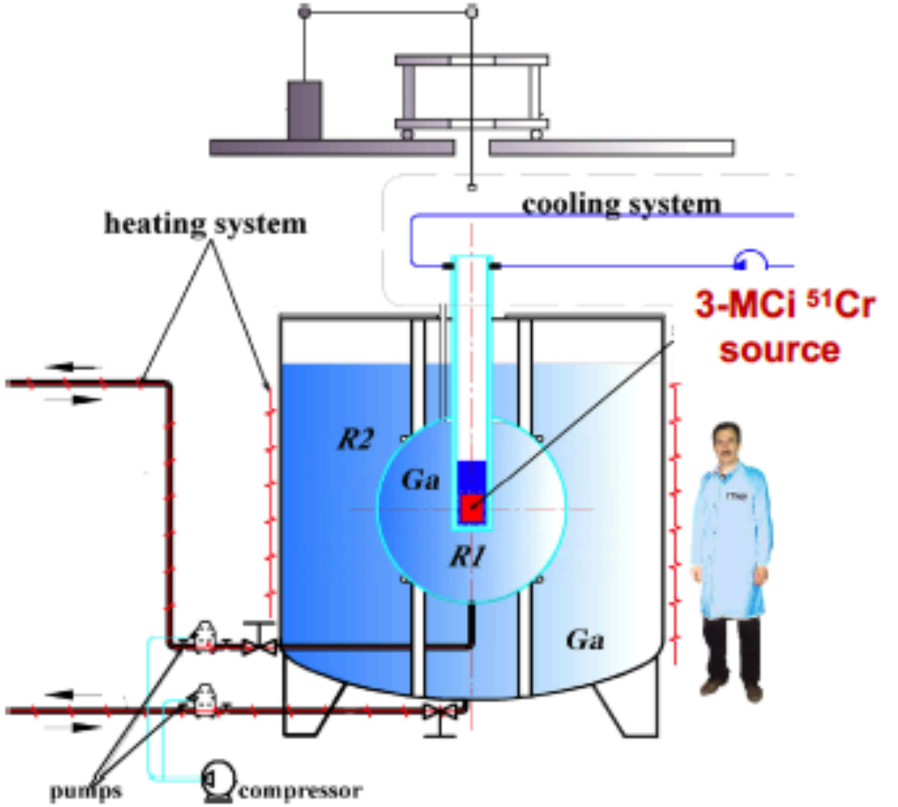
Use half-life of ^{71}Ge to determine ^{71}Ge+ν_e Zeroth-order cross section – which needs corrections at the few percent level.

$$\sigma_\nu = \frac{G_F^2 g_A^2\, p' E'\, F}{\pi(2J_i+1)} \left\{ S - \frac{2}{3} A \left[E - E' + \frac{m_e^2}{E'} \right] + \frac{2D}{3Mg_A} \left[(E+E') - \frac{m_e^2}{E'} \right] \right\}$$

Electron capture rate: $$r_{\mathrm{ec}} = \frac{2\, G_F^2\, g_A^2\, E_\nu^2\, |\phi_K(0)|^2}{2\pi(2J_i'+1)} \left[S + \frac{2}{3} A\, E_\nu + \frac{2D}{3Mg_A} E_\nu \right]$$

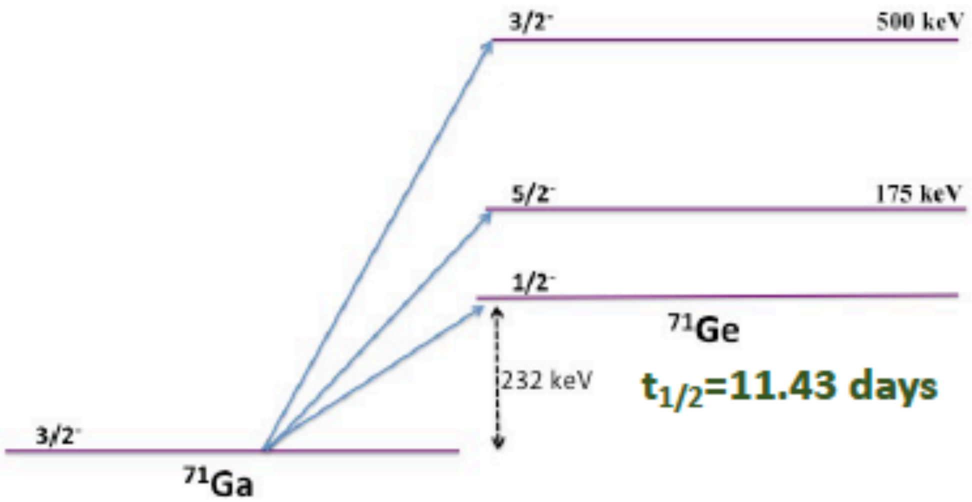
Nuclear matric elements S, A, & D	
$S = \overline{\sum_{f,i}} \lvert\langle f\lvert\Sigma\rvert i\rangle\rvert^2 = \lvert\langle J_f\lvert\lvert\Sigma_1\rvert\rvert J_i\rangle\rvert^2$	$\Sigma = \sum_{i=1}^{A} \boldsymbol{\sigma}_i \frac{\tau_z(i)}{2}$
$A = \mathrm{Re}\overline{\sum_{f,i}} \langle f\lvert\Sigma\rvert i\rangle \cdot \langle f\lvert \mathrm{i}\,\mathbf{a}\rvert i\rangle^* = \mathrm{Re}\langle J_f\lvert\lvert\Sigma_1\rvert\rvert J_i\rangle \langle J_f\lvert\lvert \mathrm{i}\, a_1\rvert\rvert J_i\rangle^*$	$\mathbf{a} = \sum_{i=1}^{A} \left\{ \frac{\boldsymbol{\sigma}_i \cdot \mathbf{p}_i}{2M}, \mathbf{x}_i \right\} \frac{\tau_z(i)}{2} + \mathbf{a}_{\mathrm{MEC}}$
$D = \mathrm{Re}\overline{\sum_{f,i}} \langle f\lvert\Sigma\rvert i\rangle \cdot \langle f\lvert\boldsymbol{\mu}_W\rvert i\rangle^* = \mathrm{Re}\langle J_f\lvert\lvert\Sigma_1\rvert\rvert J_i\rangle \langle J_f\lvert\lvert\mu_W\rvert\rvert J_i\rangle^*$	$\boldsymbol{\mu}_W = \sum_{i=1}^{A} (\boldsymbol{\sigma}_i\, \hat{\mu}_i^{\mathrm{nc}} + \mathbf{l}_i\, \hat{e}_i^{\mathrm{nc}}) + \boldsymbol{\mu}_{\mathrm{MEC}}$

Understanding the Gallium Anomaly requires new experiments and a reinvestigation of the theory



Baksan Experiment for Sterile Transitions - Two concentric zones filled with Gallium.

^{71}Ge in each Ga zone analyzed separately



$$\sigma = \sigma_{gs}\left(1 + \xi_{175}\frac{\mathrm{BGT}_{175}}{\mathrm{BGT}_{gs}} + \xi_{500}\frac{\mathrm{BGT}_{500}}{\mathrm{BGT}_{gs}}\right)$$

Bachall PRC55 3391 (1997); Haxton,PLB B353, 422 (1995) and PLB 431, 110 (1998).

Subdominant corrections to the cross section need to be recalculated:

- **Weak magnetism, Finite size**
- **1st forbidden currents**

There are three standard model theory issues

- The energy available to the neutrino emitted in electron capture Is the nuclear Q-value minus the atomic binding energies differences between Ga and Ge.
  - ➢ We do not agree with Bahcall on this number.

- The form of the energy weighting of the axial current term in the cross section is:

$$\frac{2}{3}A\left[E - E' + \frac{m_e^2}{E'}\right]$$

But Bahcall writes: $$\frac{2}{3}A\left[E - E' + \frac{m_e^2}{E'}\right]$$

(A typo? ~1% issue)

- The excited state cross sections come from a model calculation
- Earlier experimental efforts to determine theses via (p,n) scattering questioned.

Haxton et al.

$$\langle J_f \| O^{J=1}_{(\mathrm{p,n})} \| J_i \rangle = \langle J_f \| O^{J=1}_{\mathrm{GT}} \| J_i \rangle + \delta \langle J_f \| O^{J=1}_{L=2} \| J_i \rangle_{\mathrm{SM}}$$

Gallium Status

- Several few percent Standard Model corrections need reevaluation:

  E_ν for both Ga, and sources Cr and Ar; corrections to cross sections.

  (Friar +Hayes , In prep)

But to date no calculations presented to explain the anomaly away

Reactor Neutrino Anomaly

Requires an understanding of reactor anti-neutrino spectra.

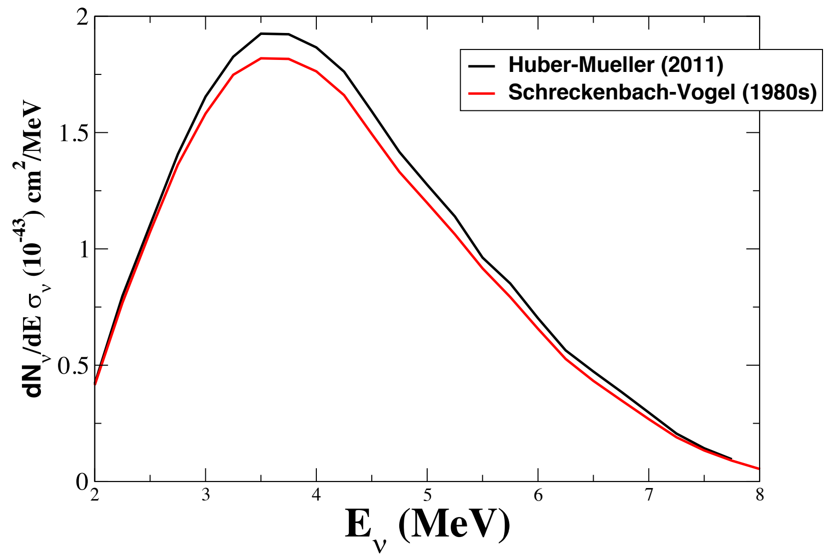
The predicted number of detectable reactor antineutrinos has evolved upward over time

In the 1980s two predictions became the standards for the field:

- Schreckenbach *et al*. converted their measured fission b-spectra for ^{235}U, ^{239}Pu and ^{241}Pu into antineutrino spectra

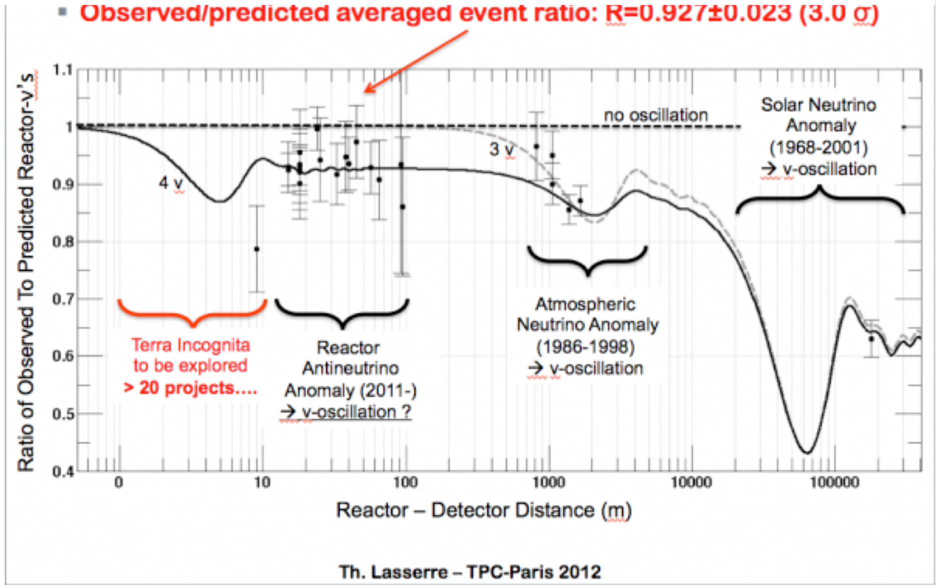
- Vogel *et al*. used the nuclear databases to predict the spectrum for ^{238}U

In 2011 both Mueller *et al.* and Huber predicted that improvements in the description of the spectra increase the expected number of antineutrinos by 5-6%.



This led to a 5-6% shortfall in the antineutrino flux in all short baseline reactor experiments - Reactor Neutrino Anomaly

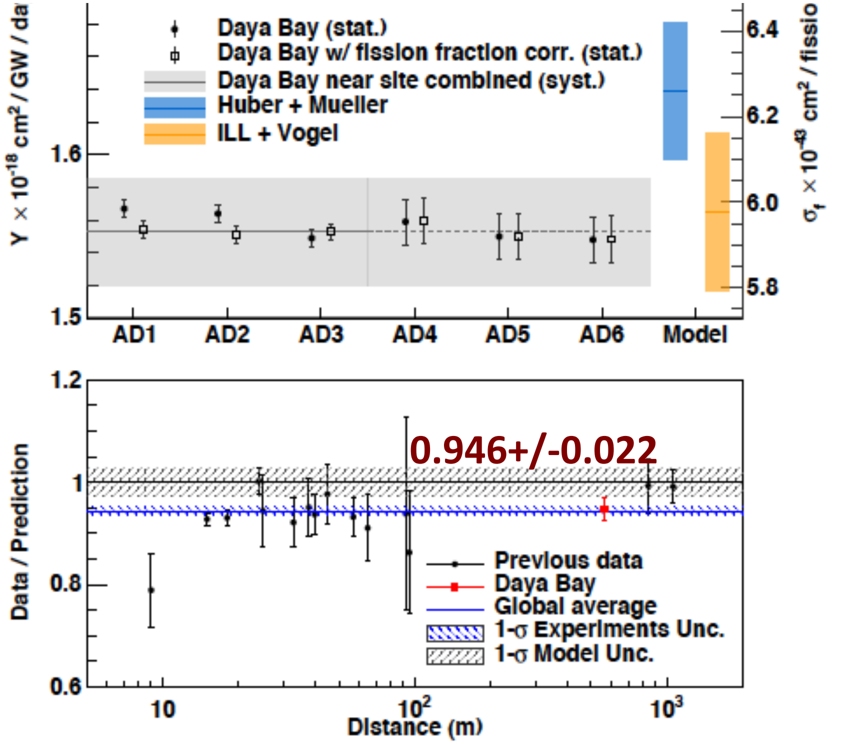
From Th. Lasserre, 2012



If this is an oscillation phenomenon, it requires a 1 eV sterile neutrino.

Results from Daya Bay, 2016

PRL,116 (2016) 061801



Accurate measurements of the total flux at Daya Bay, RENO and Double Chooz confirms the shortfall.

The issue then becomes ones of:

- **Confirming/re-examining the expectations and their uncertainties**
- **Confirming/denying the existence of 1 eV sterile neutrinos**

Two ways to determine the antineutrino spectra

- The summation method – Sum up all the beta decays, weighted by their fission yields.

$$\frac{dN_\nu}{dE_\nu} = \sum Y(Z,A) \sum B(E_0) P(E, E_0, Z, A)$$

Cumulative Fission yields

Beta branching ratios

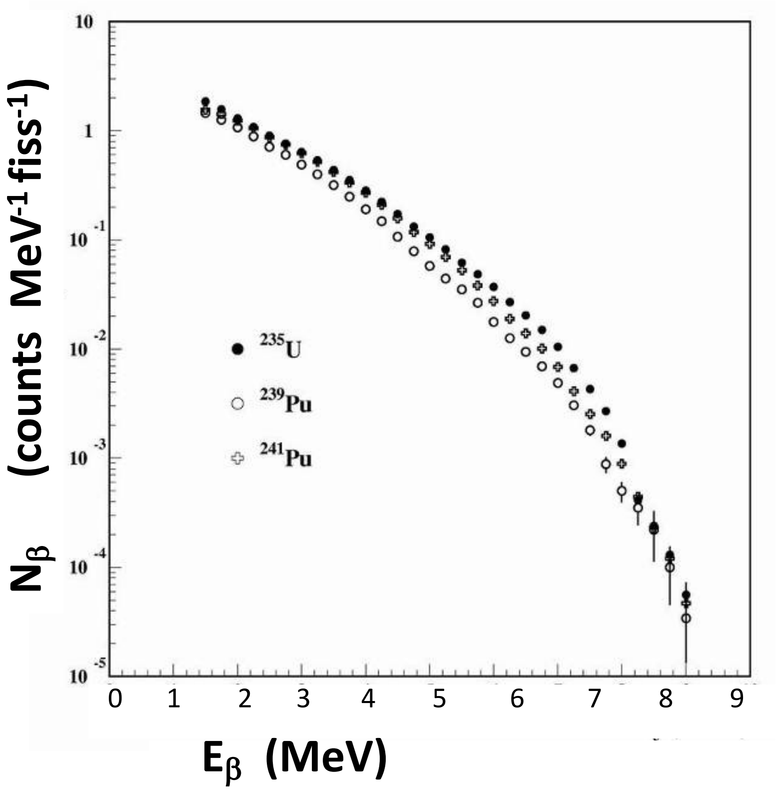
Individual spectra

- The conversion method - Measure the beta electron spectrum and convert into an anti-neutrino spectrum

$$\frac{dN_e}{dE_e} \rightarrow \frac{dN_\nu}{dE_\nu}$$

Both methods introduce uncertainties

The Original Expected Fluxes were determined via the Conversion Method using β-Spectra (electrons) made at the ILL Reactor in the 1980s



K. Schreckenbach et al. PLB118, 162 (1985)
A.A. Hahn et al. PLB160, 325 (1989)

- The thermal fission beta spectra for ^{235}U, ^{239}Pu, ^{241}Pu were measured at ILL.
- These β-spectra were converted to antineutrino spectra by fitting to 30 end-point energies
- Vogel *et al*. used the ENDF-5 nuclear database to estimate ^{238}U, which requires fast neutron fission

Vogel, et al., Phys. Rev. C24, 1543 (1981).

FIT

Parameterized

$$S_\beta(E) = \sum_{i=1,30} a_i \; S^i(E, E_o^i)$$

$$S^i(E, E_0^i) = E_\beta p_\beta (E_0^i - E_\beta)^2 F(E, Z_{eff})(1 + \delta_{corrections})$$

Two inputs are needed to convert β-spectra to antineutrino spectra:
(1) Z of the fission fragments for the Fermi function, (2) sub-dominant corrections

$$S^i(E,E_0^i) = E_\beta p_\beta (E_0^i - E_\beta)^2 F(E,\ Z)(1+\ \delta_{corrections})$$

The Fermi function:

Z_{eff} used for Fermi function

The corrections:

$$\delta_{correction}(E_e, Z, A) = \delta_{FS} + \delta_{WM} + \delta_R + \delta_{rad}$$

δ_{FS} = Finite size correction to Fermi function

δ_{WM} = Weak magnetism

δ_R = Recoil correction

δ_{rad} = Radiative correction

A change to the approximations used for these effects led to the anomaly

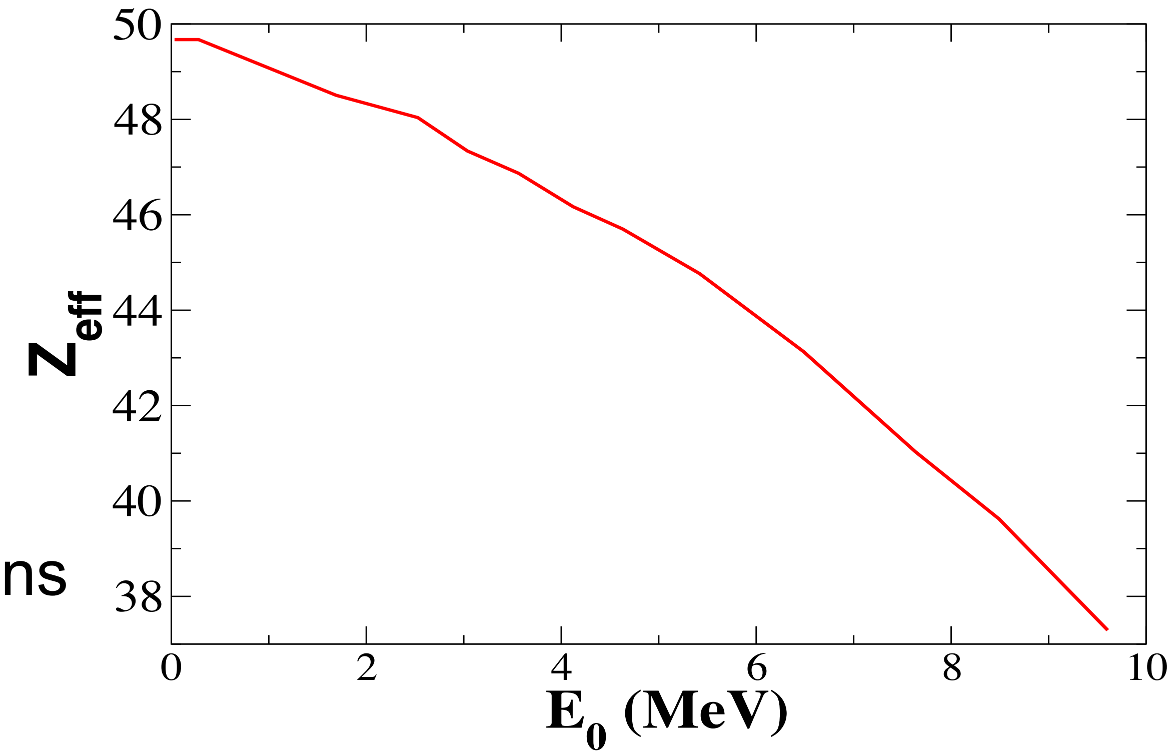
An energy-dependent Zeff representing the fission fragments is needed to determine the Fermi function

- On average, higher end-point energy means lower Z.
- Comes from nuclear binding energy differences

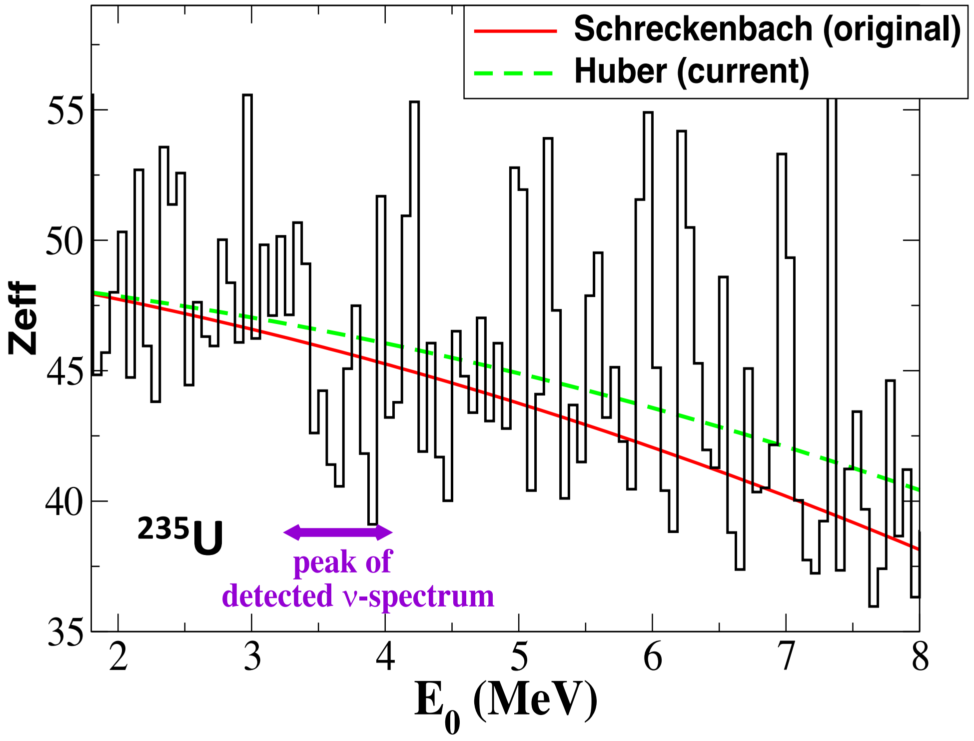
Parameterized used by both Schreckenbach and Huber involved a quadratic function:

$$Z_{eff} \sim a + b\ E_0 + c\ E_0^2$$

But the difference in their parameterizations is a large part of the anomaly.



The newer fit to Zeff used in the Fermi function for the conversion of the aggregate β-spectrum, led to a higher ν-spectrum



- **Huber's new parameterization of Zeff with end-point energy E_0 changes the Fermi function, accounting for 50% of the current anomaly.**
- **But the data do not follow a simple quadratic form.**

The corrections to b-decay are half the anomaly

$$S^i(E,E_0^i) = E_\beta p_\beta (E_0^i - E_\beta)^2 F(E,\ Z)(1+\ \underline{\delta_{corrections}})$$

$$\delta_{correction}(E_e,Z,A) = \delta_{FS} + \delta_{WM} + \delta_R + \delta_{rad}$$

$$\delta_{FS} = \text{Finite size correction to Fermi function}$$

$$\delta_{\mathrm{WM}} = \text{ Weak magnetism}$$

$$\delta_{\mathrm{R}} = \text{Recoil correction}$$

$$\delta_{\mathrm{rad}} = \text{ Radiative correction}$$

- Recoil and radiative corrections are well-known and <u>nucleus independent.</u>

- The finite size and weak magnetism corrections are nucleus dependent and should be applied to each β-decay transition, which is a problem for the conversion method.

Corrections for GT Transitions

1. Finite size of the nucleus

Vogel, (Mueller)

$$\delta_{FS} = A_c = -\frac{10Z\alpha R}{9\hbar c}E_\beta \; ; \; R = 1.2A^{1/3}$$

Friar, Holstein

$$\delta_{FS} = A_c = -\frac{3Z\alpha R}{2\hbar c}\left(E_\beta - \frac{E_\nu}{27} + \frac{m_e^2}{E_\beta}\right); \; R = \frac{36}{35}(1.2A^{1/3})$$

2. Weak magnetism

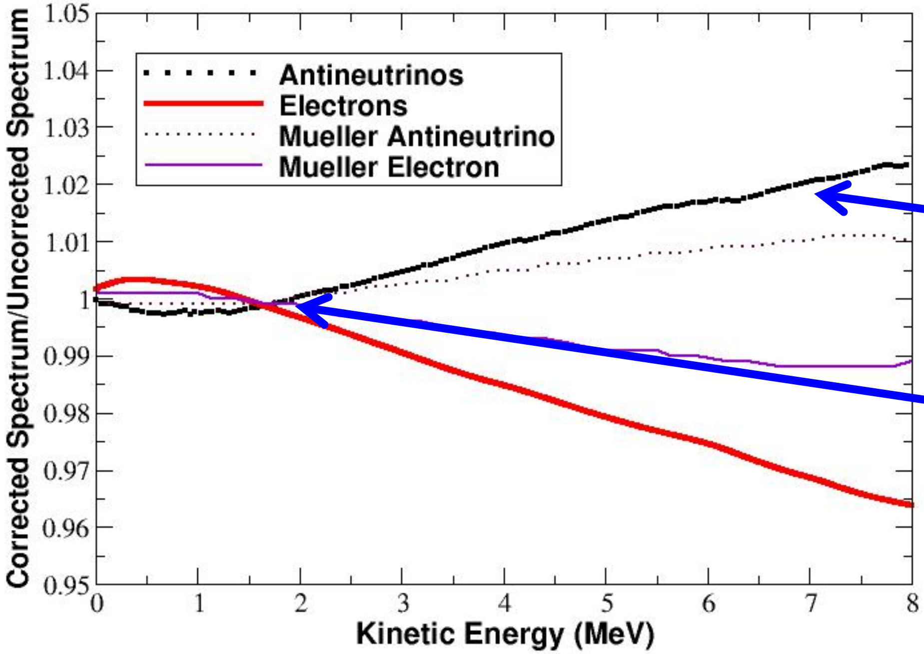
Vogel, (Mueller)

$$\delta_{WM} = A_w = \frac{4(\mu_V - 1/2)}{3M_n}2E_\beta$$

Friar, Holstein

$$\delta_{WM} = A_w = \frac{4(\mu_V - 1/2)}{6M_n}(E_\beta \beta^2 - E_\nu)$$

Effect of FS and MW Corrections to Spectrum using ENDF/B-VII, and assuming that all Transitions are allowed



$$\cong \frac{1}{2}\left[\frac{4(\mu_\nu - \frac{1}{2})}{3} - \frac{3Z\alpha R}{2\hbar c}\right]$$

$$\overline{E}_0 / 2$$

Originally approximated by a parameterization: $\delta_{FS} + \delta_{WM} = 0.0065(E_\nu - 4MeV))$

30% of the beta-decay transitions involved are so-called forbidden

Allowed transitions $\Delta L=0$; Forbidden transitions $\Delta L\neq 0$

Forbidden transitions introduce a shape factor C(E):

$$S(E_e,Z,A)=\frac{G_F^2}{2\pi^3}p_eE_e(E_0-E_e)^2\underline{C(E)}F(E_e,Z,A)(1+\delta_{corr}(E_e,Z,A))$$

The corrections for forbidden transitions are also different and sometimes unknown :

Classification	ΔJ^π	Operator	Shape Factor $C(E)$	Fractional Weak Magnetism Correction $\delta_{WM}(E)$
Allowed GT	1^+	$\Sigma\equiv\sigma\tau$	1	$\frac{2}{3}\left[\frac{\mu_v-1/2}{M_N g_A}\right](E_e\beta^2-E_\nu)$
Non-unique 1^{st} Forbidden GT	0^-	$[\Sigma,r]^{0-}$	$p_e^2+E_\nu^2+2\beta^2E_\nu E_e$	0
Non-unique 1^{st} Forbidden ρ_A	0^-	$[\Sigma,r]^{0-}$	λE_0^2	0
Non-unique 1^{st} Forbidden GT	1^-	$[\Sigma,r]^{1-}$	$p_e^2+E_\nu^2-\frac{4}{3}\beta^2E_\nu E_e$	$\left[\frac{\mu_v-1/2}{M_N g_A}\right]\left[\frac{(p_e^2+E_\nu^2)(\beta^2E_e-E_\nu)+2\beta^2E_eE_\nu(E_\nu-E_e)/3}{(p_e^2+E_\nu^2-4\beta^2E_\nu E_e/3)}\right]$
Unique 1^{st} Forbidden GT	2^-	$[\Sigma,r]^{2-}$	$p_e^2+E_\nu^2$	$\frac{3}{5}\left[\frac{\mu_v-1/2}{M_N g_A}\right]\left[\frac{(p_e^2+E_\nu^2)(\beta^2E_e-E_\nu)+2\beta^2E_eE_\nu(E_\nu-E_e)/3}{(p_e^2+E_\nu^2)}\right]$
Allowed F	0^+	τ	1	0
Non-unique 1^{st} Forbidden F	1^-	$r\tau$	$p_e^2+E_\nu^2+\frac{2}{3}\beta^2E_\nu E_e$	0
Non-unique 1^{st} Forbidden $\vec{J}_V$	1^-	$r\tau$	E_0^2	-

The forbidden transitions increase the uncertainty in the expected spectrum.

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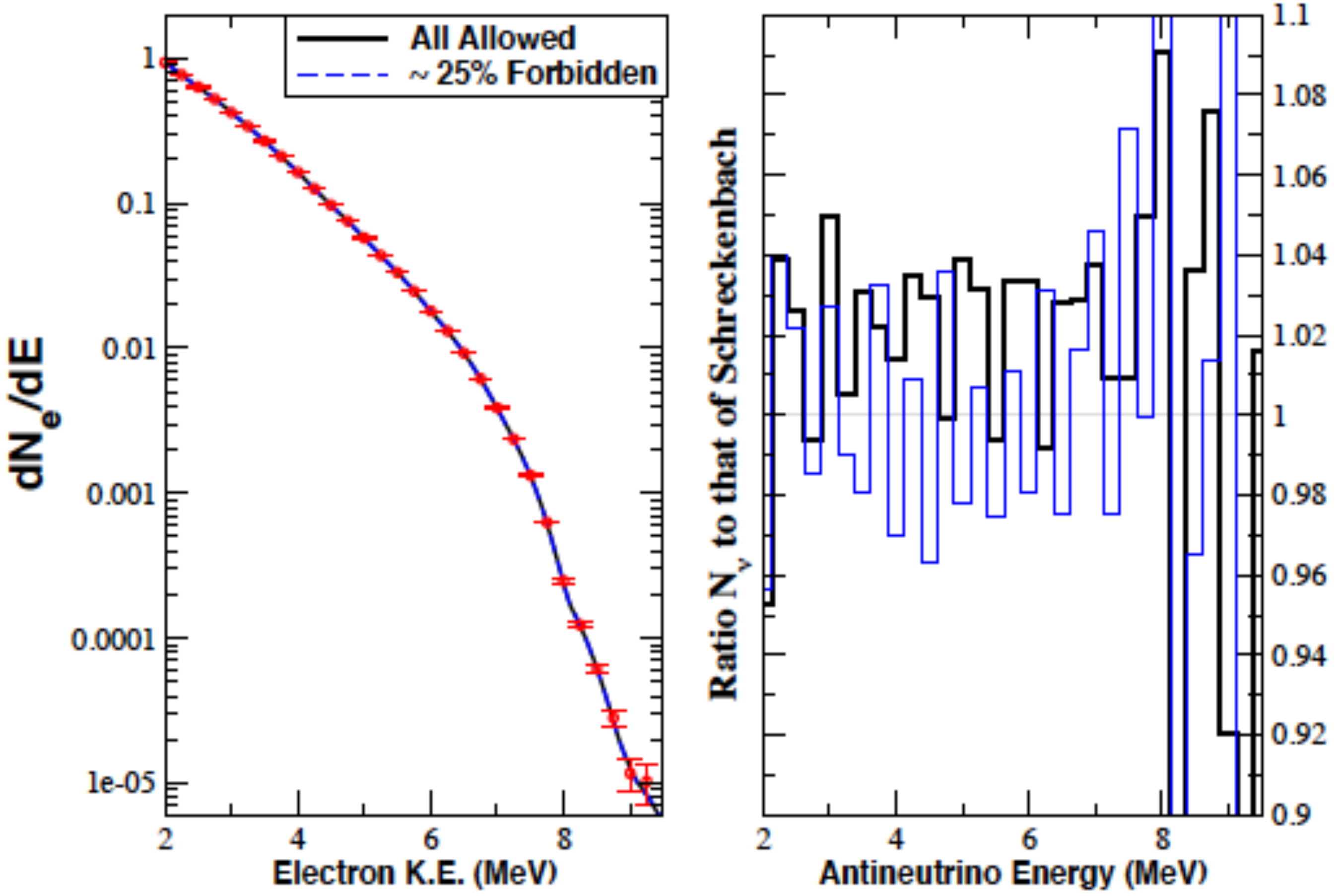
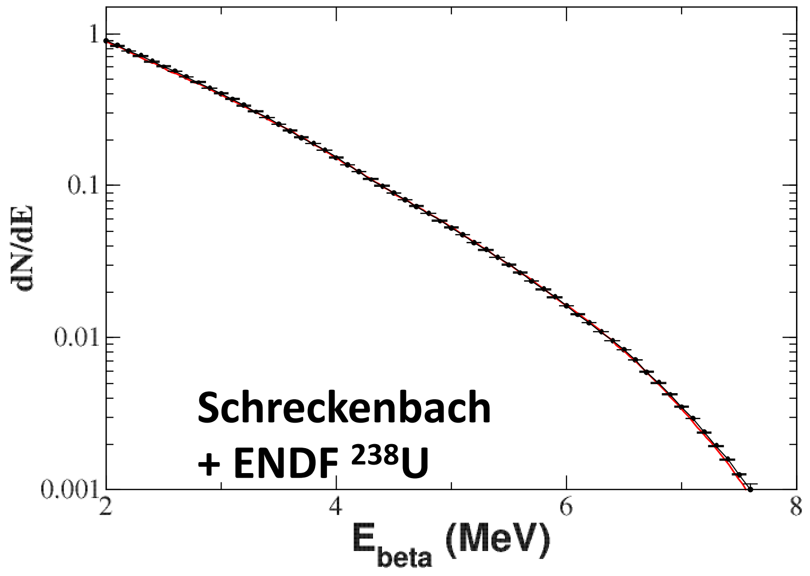
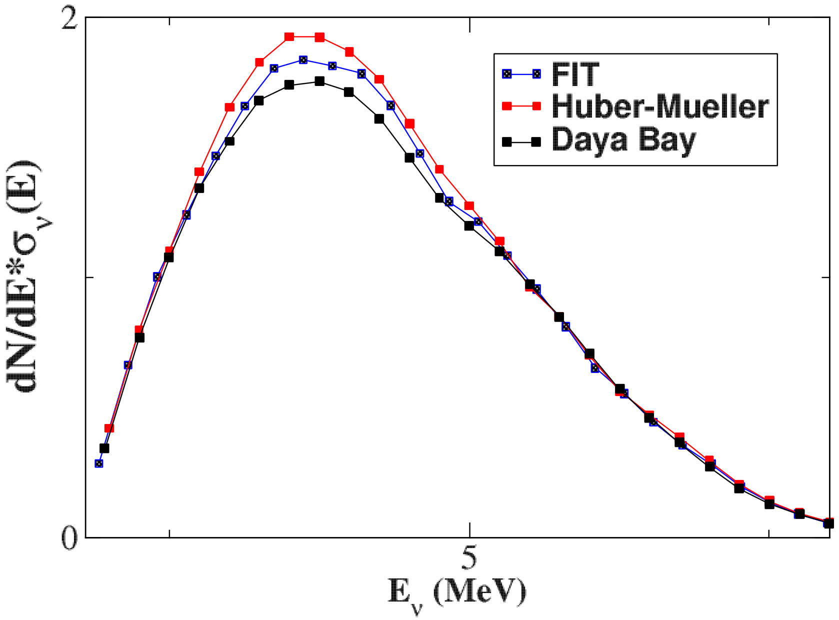


Two equally good fits to the Schreckenbach β-spectra, lead to ν-spectra that differ by 4%.

Simultaneous fit of the Daya Bay antineutrino spectrum and the equivalent aggregate β-spectrum with (1) point-wise Z_{eff} and (2) improved descriptions of forbidden transitions reduces the anomaly from 5% to 2.5%



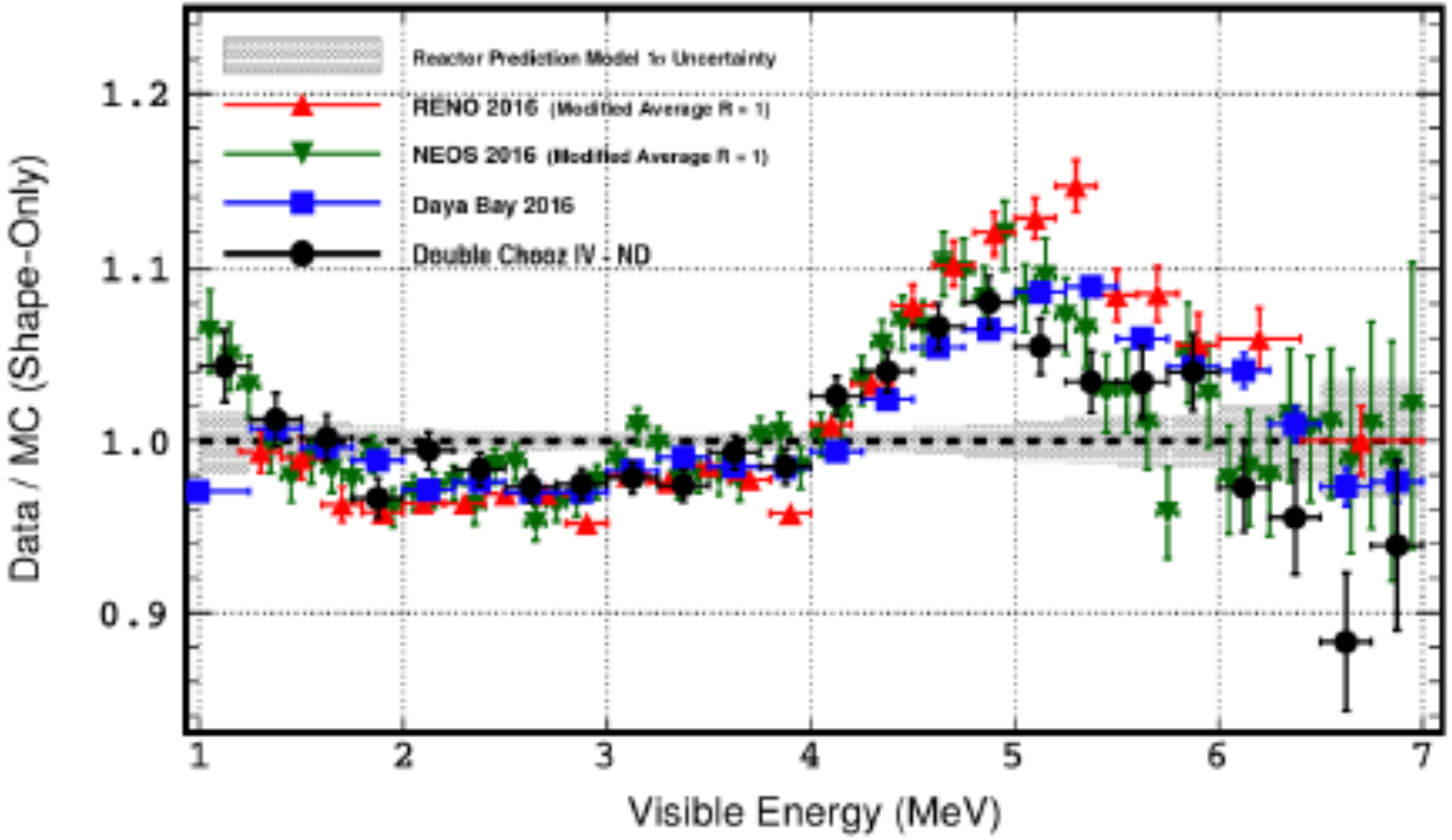


The magnitude of the IBD cross sections change, depending on assumptions, but not the ratio of one isotope to another

	all allowed $Z_{\mathrm{eff}}^{\mathrm{Huber}}$	all allowed Z_{eff}	allow.+forbid. Z_{eff}	allow.+forbid. $(Z_{\mathrm{eff}}^2)^{1/2}$
^{235}U	6.69	6.58	6.47	6.48
^{239}Pu	4.36	4.3	4.22	4.23
ratio	1.534	1.530	1.533	1.532

A 'BUMP' is seen in the Measured Spectra compared to predictions

The Reactor Neutrino ‘BUMP’



Expectations:
P. Huber, Phys. Rev. C 84, 024617 (2011);
Th. A. Mueller et al., Phys. Rev. C 83, 054615 (2011);
N. Haag, Phys. Rev. Lett. 112, 122501 (2014).

Laura Bernard, STEREO (from Moriond 2019)

Possible Origins of the 'Bump'

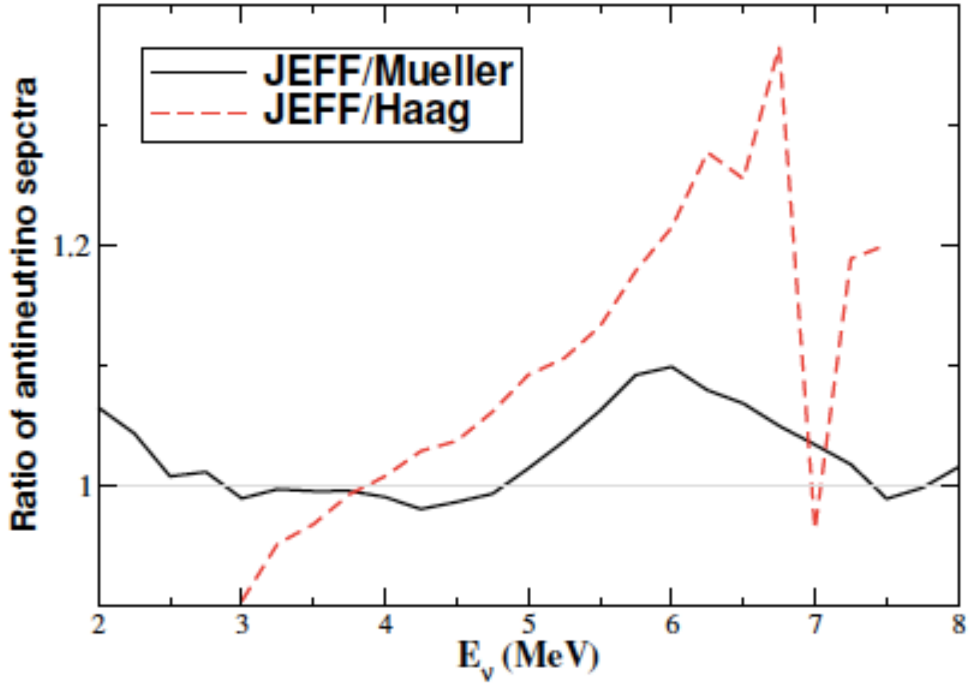
- ^{238}U as a source of the shoulder
  – Possible because ^{238}U has a hard spectrum and contributes significantly in the Bump energy region. It is also the most uncertain actinide.

- A possible error in the ILL b-decay measurements
  - Possible but not predicted by current updated nuclear databases.

- The harder PWR Neutron Spectrum
  - Possible but not predicted by standard fission theory.
  - no convincing experimental data either way.

All of these are nuclear physics explanations pointing to the problem lying with the 'expected spectra'.

For example, if the BUMP does not change with the fuel evolution, ^{238}U is a likely source



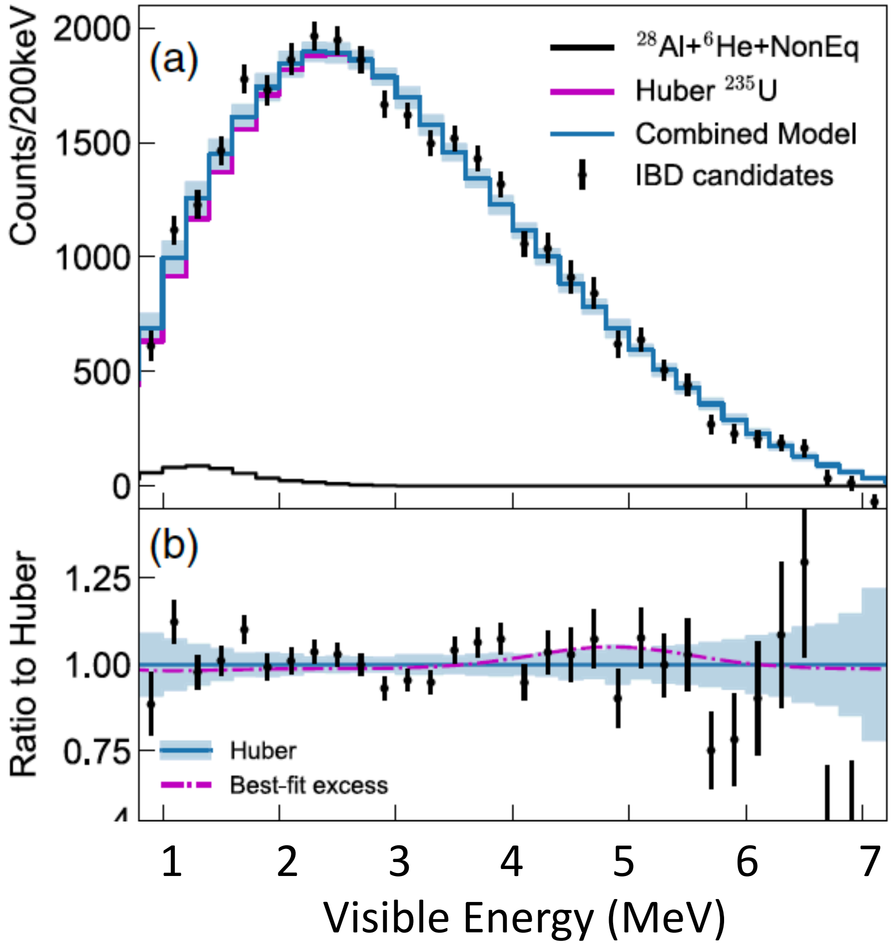
Relative to the JEFF database, both Mueller and Haag show a BUMP.

The harder spectrum of ^{238}U increases it's relative importance.

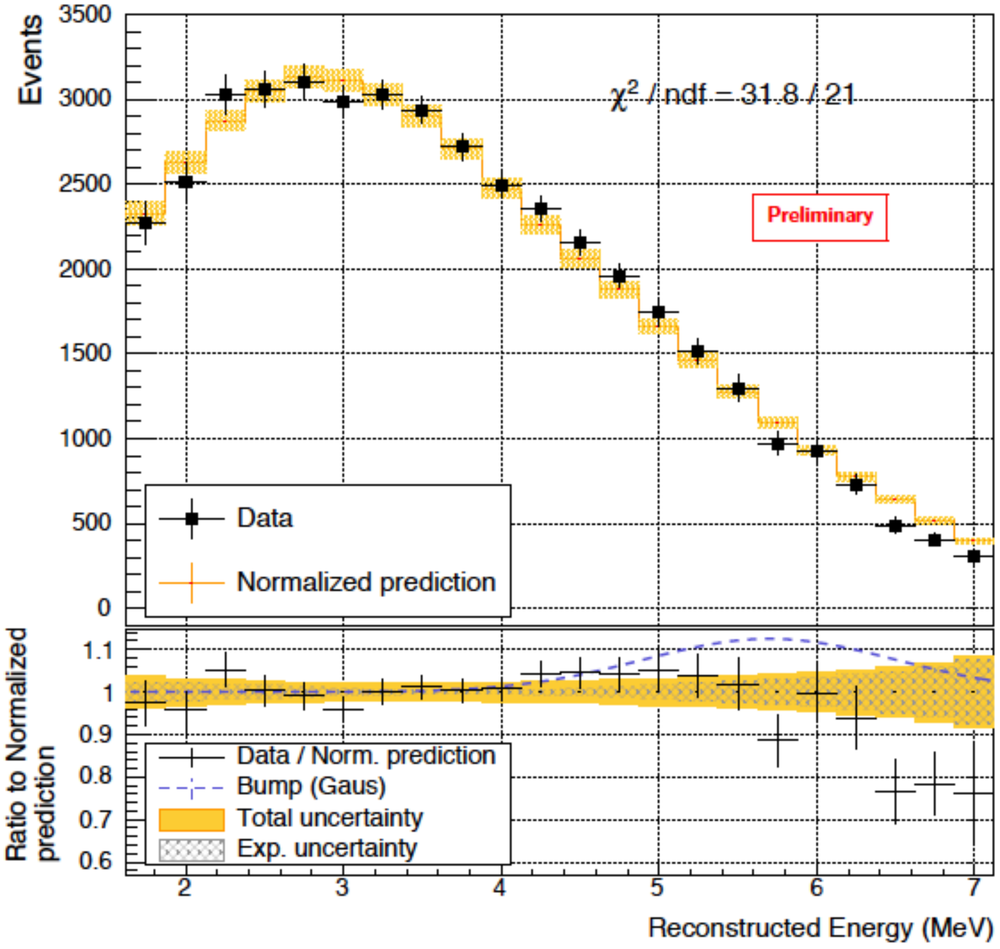
- If this is the correct explanation, the current VSBL experiments with highly enriched ^{235}U reactor will not see a BUMP.
- If, on the other hand, the ILL data are responsible all VSBL expts will see the Bump.

Hayes + Vogel, Ann. Rev. of Nucl & Part. Sci, 66 219 (2016)
Mohanty, arXiv: 1711.1.02801

VSBL Measurements by STEREO (ILL) and PROSPECT (HIFR), both 100% ^{235}U



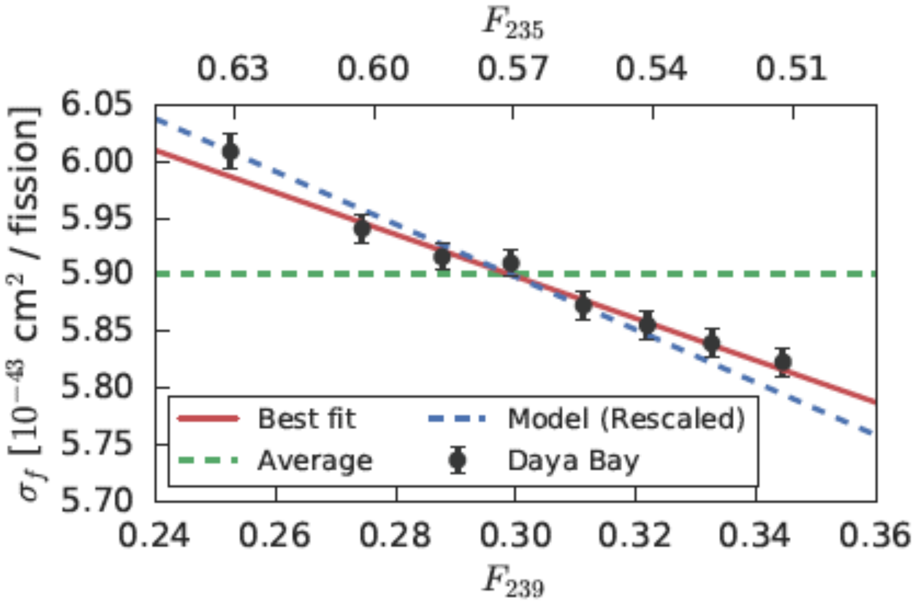
PROSPECT

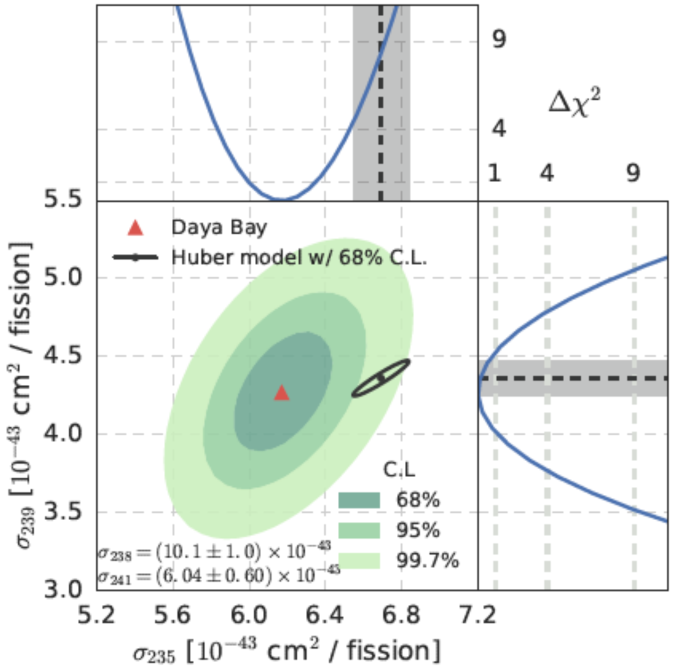


STEREO

Changes in the Antineutrino Spectra with the Reactor Fuel Burnup

The Total Number of Antineutrinos Decreases with Burnup, but the Huber-Mueller Model does not agree with the measured slope



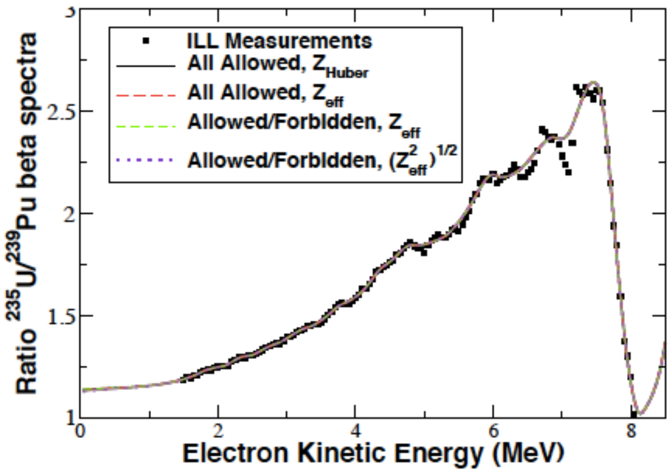


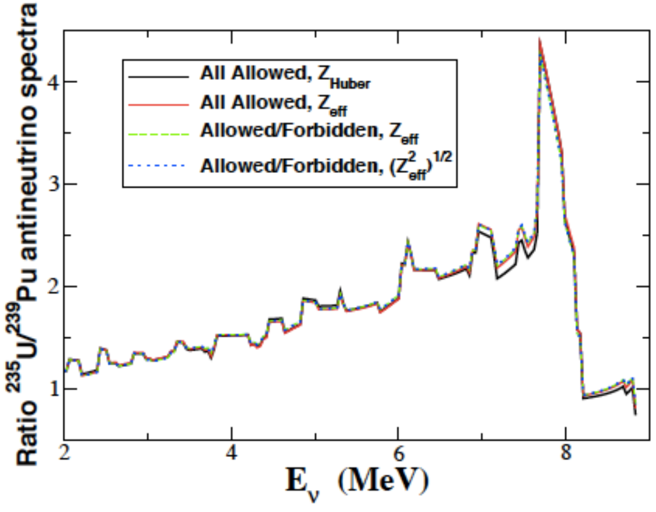
$$\sigma_f(F_{239}) = \bar{\sigma}_f + \frac{d\sigma_f}{dF_{239}}(F_{239} - \overline{F}_{239})$$

$$d\sigma_f/dF_{239} = (-1.86 \pm 0.18) \times 10^{-43} \text{ cm}^2/\text{fission}$$ **Experiment**

$$(-2.46 \pm 0.06) \times 10^{-43} \text{cm}^2/\text{fission}$$ **Expected**

The discrepancy between current Huber-Mueller model predictions and the Daya Bay results can be traced to the original Schreckenbach measured $^{235}U/^{239}Pu$ ratio





Using different will change the IBD cross sections for ^{235}U and ^{239}Pu.

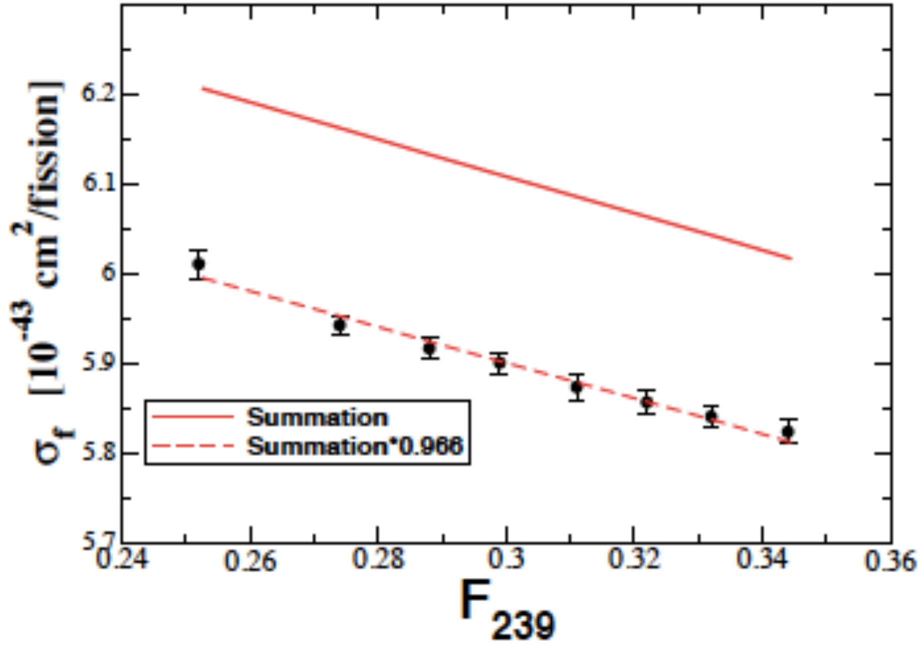
	all allowed $Z_{\rm eff}^{\rm Huber}$	all allowed $Z_{\rm eff}$	allow.+forbid. $Z_{\rm eff}$	allow.+forbid. $(Z_{\rm eff}^2)^{1/2}$
^{235}U	6.69	6.58	6.47	6.48
^{239}Pu	4.36	4.3	4.22	4.23
ratio	1.534	1.530	1.533	1.532

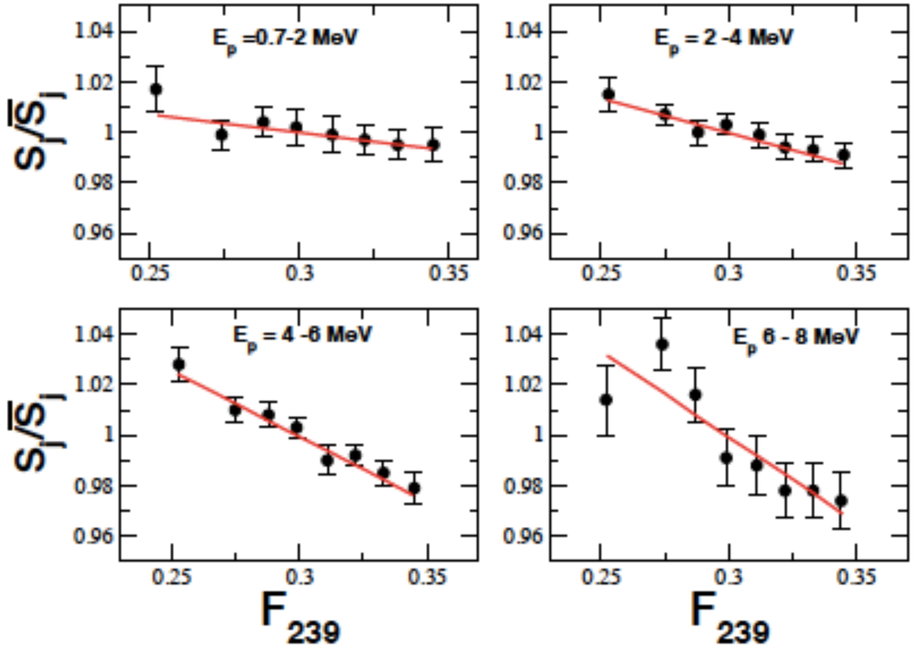
But the ratio of $^{235}U/^{239}Pu$ is fixed.

$\sigma5/\sigma9$ = 1.53 +/- 0.05 (Schreckenbach)

$\sigma5/\sigma9$ = 1.445 +/- 0.097 (Daya Bay)

The Nuclear database explains all of the Daya Bay fuel evolution data, but still allows for a (smaller) anomaly



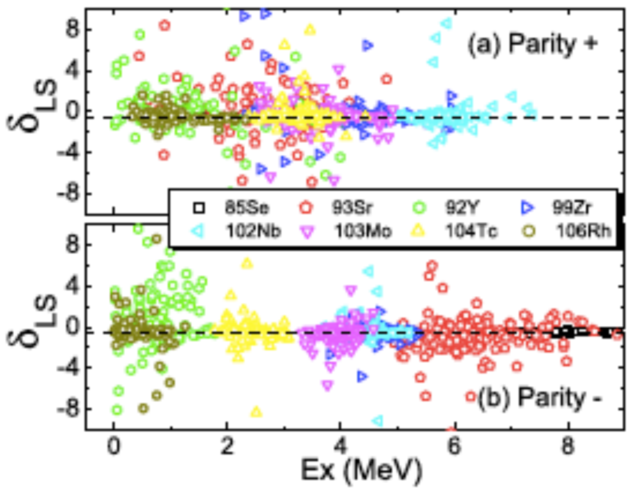


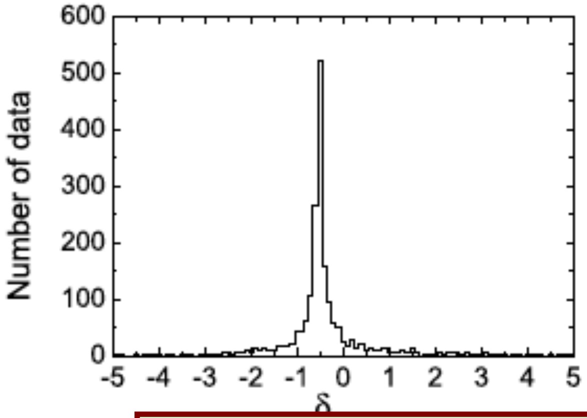
- The IBD yield is predicted to change with the correct slope.
- But the absolute predicted value is high by 3.5%.
- This anomaly is not statistically significant but it means that Daya Bay evolution data do not rule out sterile neutrinos.

Summary

- Four classes of experiments do not fall into the 3-ν –standard picture: LSND, MiniBooNE, the Reactor Anomaly and the Gallium Anomaly .
- Several of the Anomalies involve detailed nuclear and atomic physics issues.
- Improve treatments of the standard model physics tends to reduce the size of the anomalies.
- But more experimental and theoretical investigation is needed.

Weak Magnetism has an uncertainty arising from the approximation used for the orbital contribution and from omitted 2-body currents.
But, dominant 0+→0- transitions have zero δ_{WM}, with no uncertainty





$$\delta_{WM}^{GT} = \frac{4(\mu_V - 1/2)}{6M_N g_A}(E_e\beta^2 - E_\nu)$$

$$\delta_{LS}^{j_f j_i} \equiv \frac{\langle J_f \,||\, \vec{\Lambda} \,||\, J_i\rangle}{\langle J_f \,||\, \vec{\Sigma} \,||\, J_i\rangle} \simeq -\frac{1}{2}$$

- **Checked for a subset of fission fragments.**
- **A check for all fission fragments, including 2-body terms, requires a large super-computing effort.**

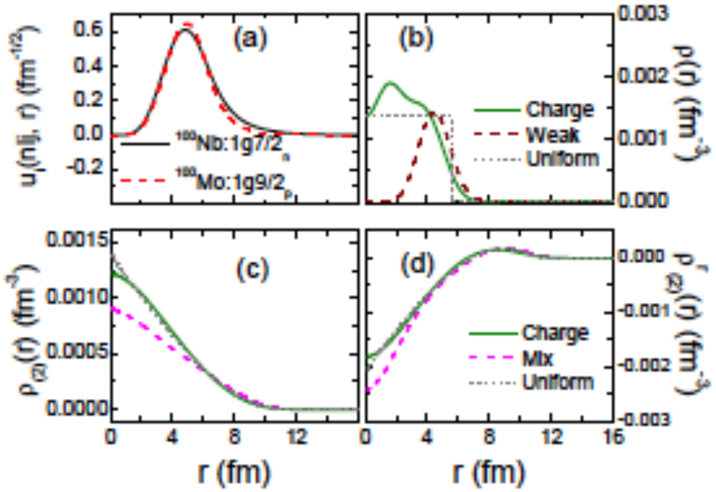
Estimated uncertainty ~ 30% for this 4% correction to the spectra

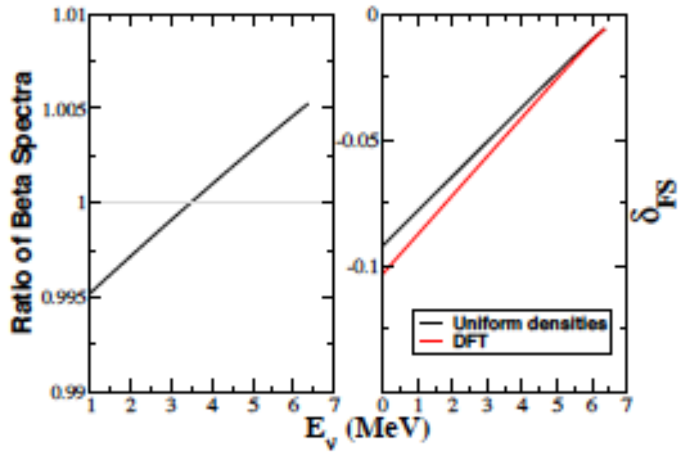
Wang and Hayes, Phys. Rev. C 95, 064313 (2017) .

The Finite Size Correction can be expressed in terms of Zemach moments

$$\delta_{FS} = \Delta F_{\mathrm{REL}}/F_{\mathrm{REL}} = -\frac{Z\alpha}{3\hbar c}\left(4E\,\langle r\rangle_{(2)} + E\,\langle r\rangle^{r}_{(2)} - \frac{E_\nu \langle r\rangle^{r}_{(2)}}{3} + \frac{m^2c^4}{E}\left(2\,\langle r\rangle_{(2)} - \langle r\rangle^{r}_{(2)}\right)\right)$$

Approximated as : $$\delta_{FS} = -\frac{3Z\alpha}{2\hbar c} < r >_{(2)} (E_e - \frac{E_\nu}{27} + \frac{m^2c^4}{3E_e})$$





- **Found to be a good approximation for allowed transitions.**
- **Not checked for forbidden transitions.**

Estimated uncertainty ~ 20% for this 5% correction to the spectra

Wang, *et al.* PRC, 94, 034314 (2016)