

Dark matter in stars

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Dark Matter in the Sun

arXiv:1311.2074 / JCAP/04019 (ACV, P. Scott)

Thermal conduction by dark matter with velocity and momentum-dependent cross-sections

arXiv:1411.6626/ PRL 114.081302 (ACV, P. Scott, A. Serenelli)

Possible Indication of Momentum-Dependent Asymmetric Dark Matter in the Sun

arXiv:1504.04378/ JCAP 1508 (2015) 08, 040 (ACV, Scott, Serenelli)

Generalised form factor dark matter in the Sun

arXiv:1605.06502 /JCAP1611 (2016) 007 (ACV, Scott, Serenelli)

Updated constraints on velocity and momentum- dependent asymmetric dark matter

arXiv:1610.06737/JCAP03(2017)029 (B. Geytenbeek, S. Roa, P. Scott, A. Serenelli, ACV, M. White, A. Williams)

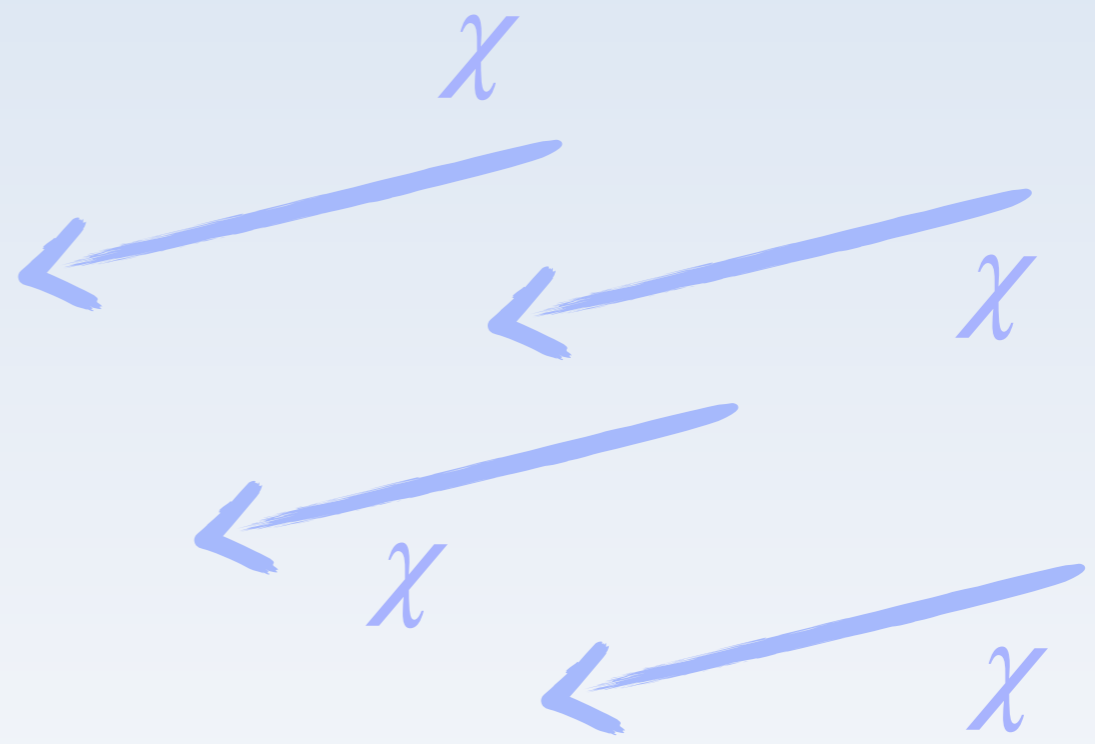
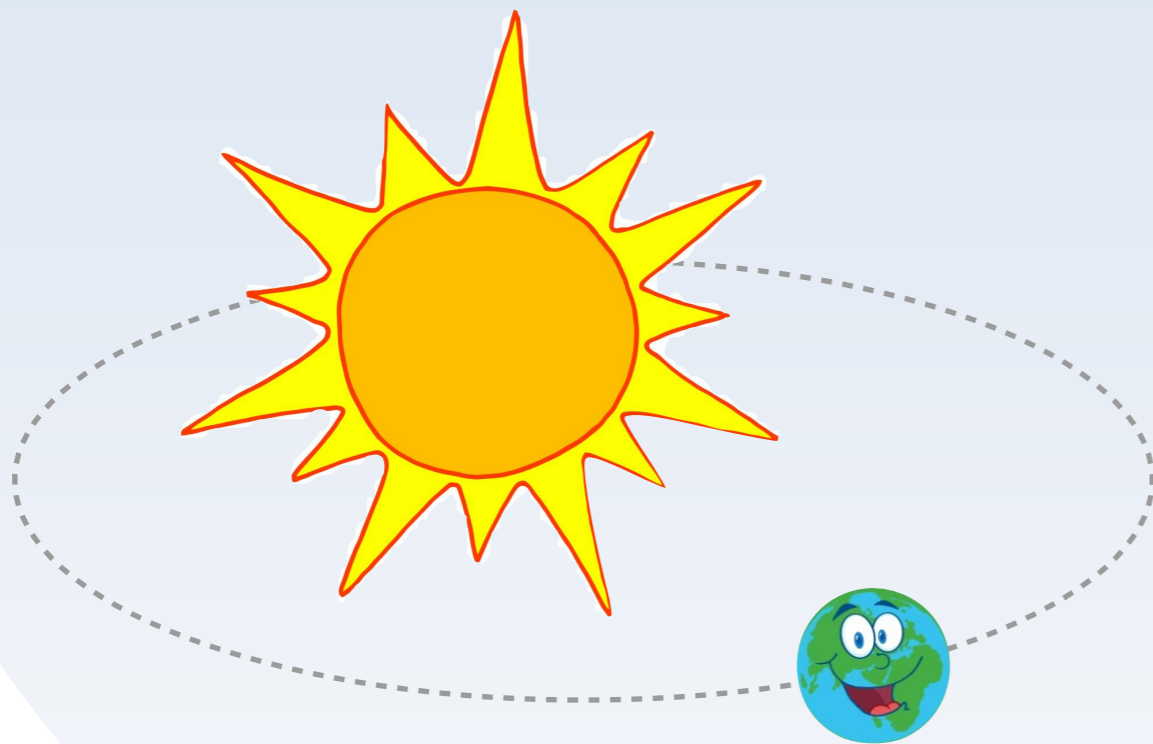
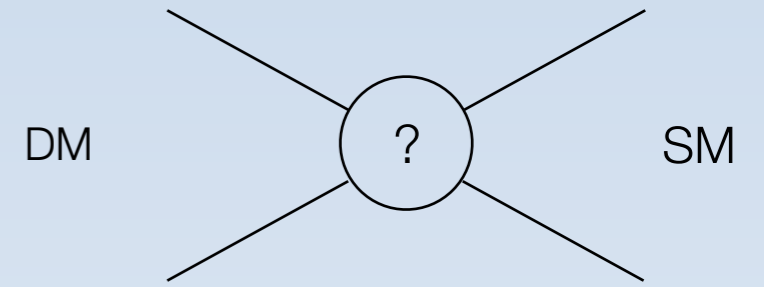
Effect of electromagnetic dipole dark matter on energy transport in the solar interior

arXiv:1703.07784/JCAP 1710 (2017)10 037: G.Busoni, A. de Simone, P. Scott, ACV

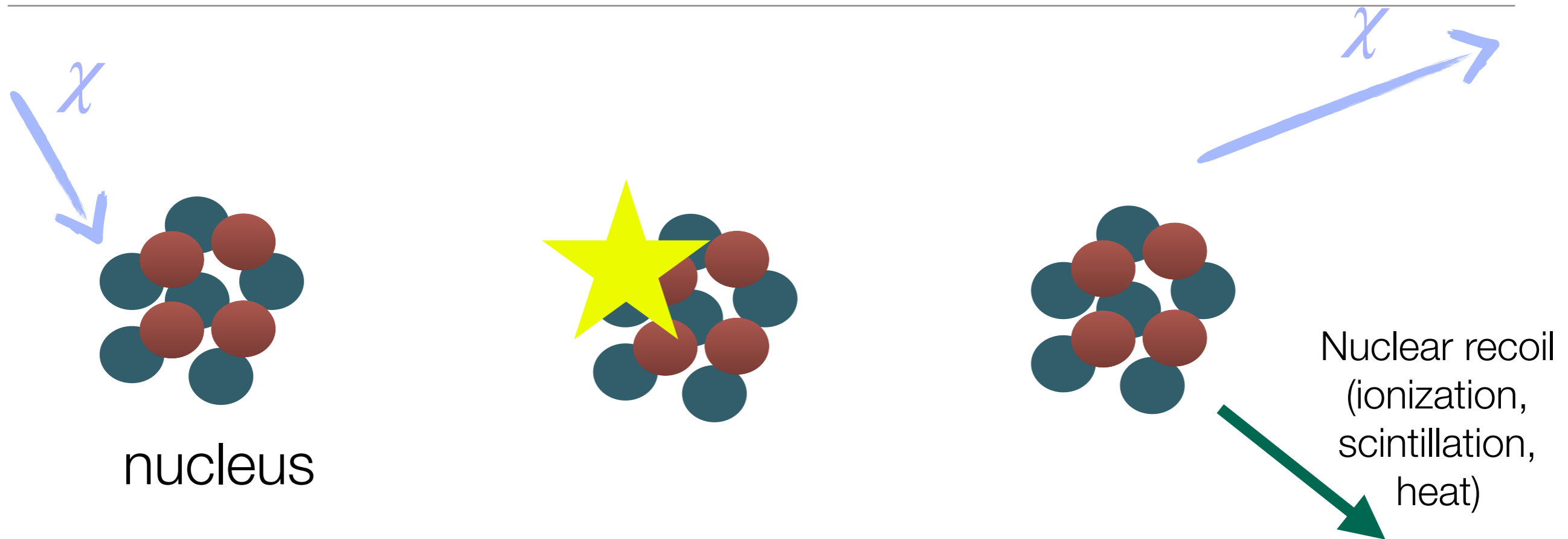
Generalised solar capture and evaporation of DM

++ work in progress w/ Hannah Banks, Siyam Ansari (Imperial), Neal Kozar (Queen's)

Direct detection



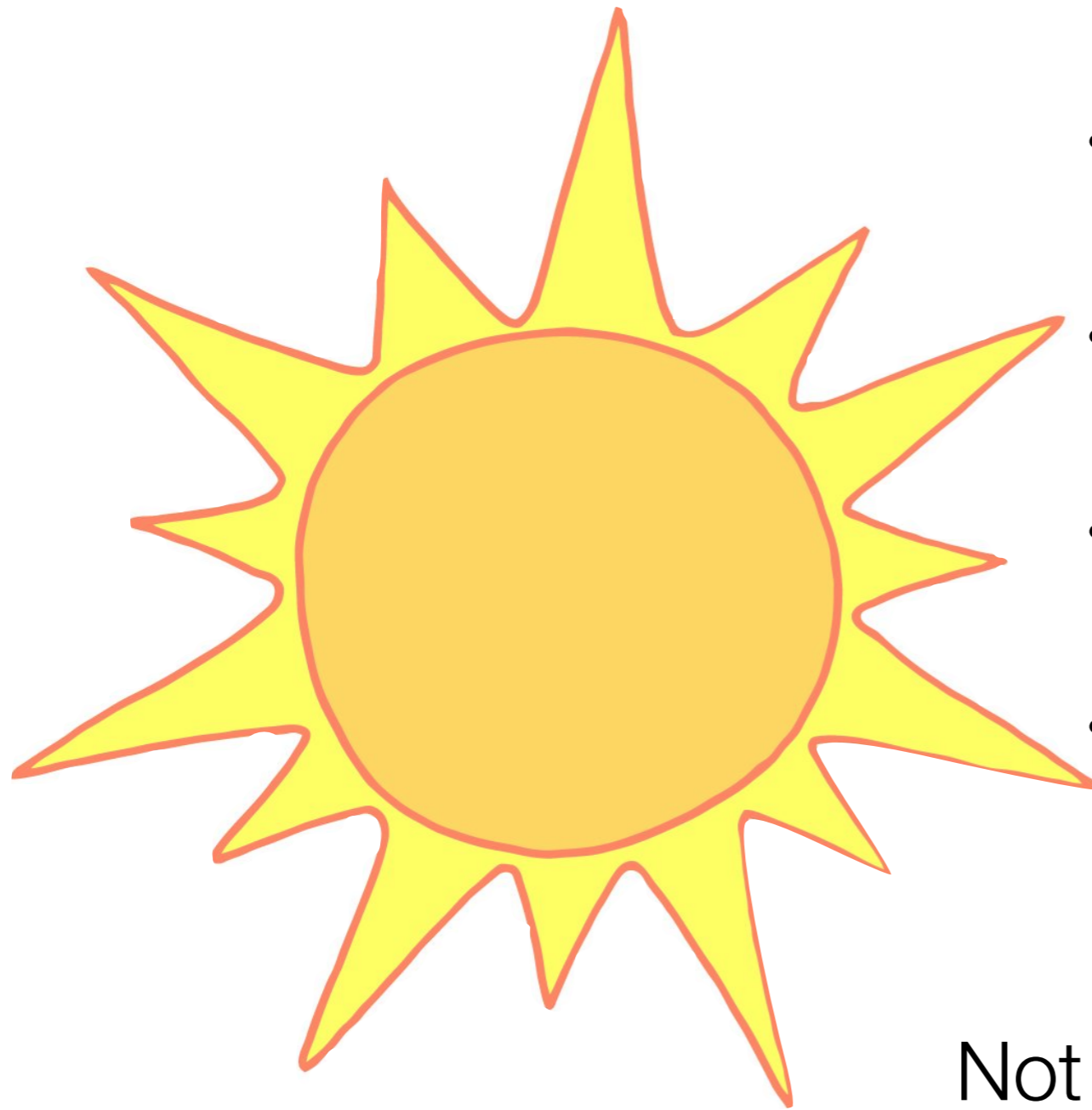
Direct detection



Most sensitive to **heavy, fast** particles \rightarrow larger recoil signal

$$R = \int_{E_T}^{\infty} dE_R \frac{\rho_0}{m_N m_\chi} \int_{v_{min}}^{\infty} v f(v) \frac{d\sigma_{WN}}{dE_R}(v, E_R) dv$$

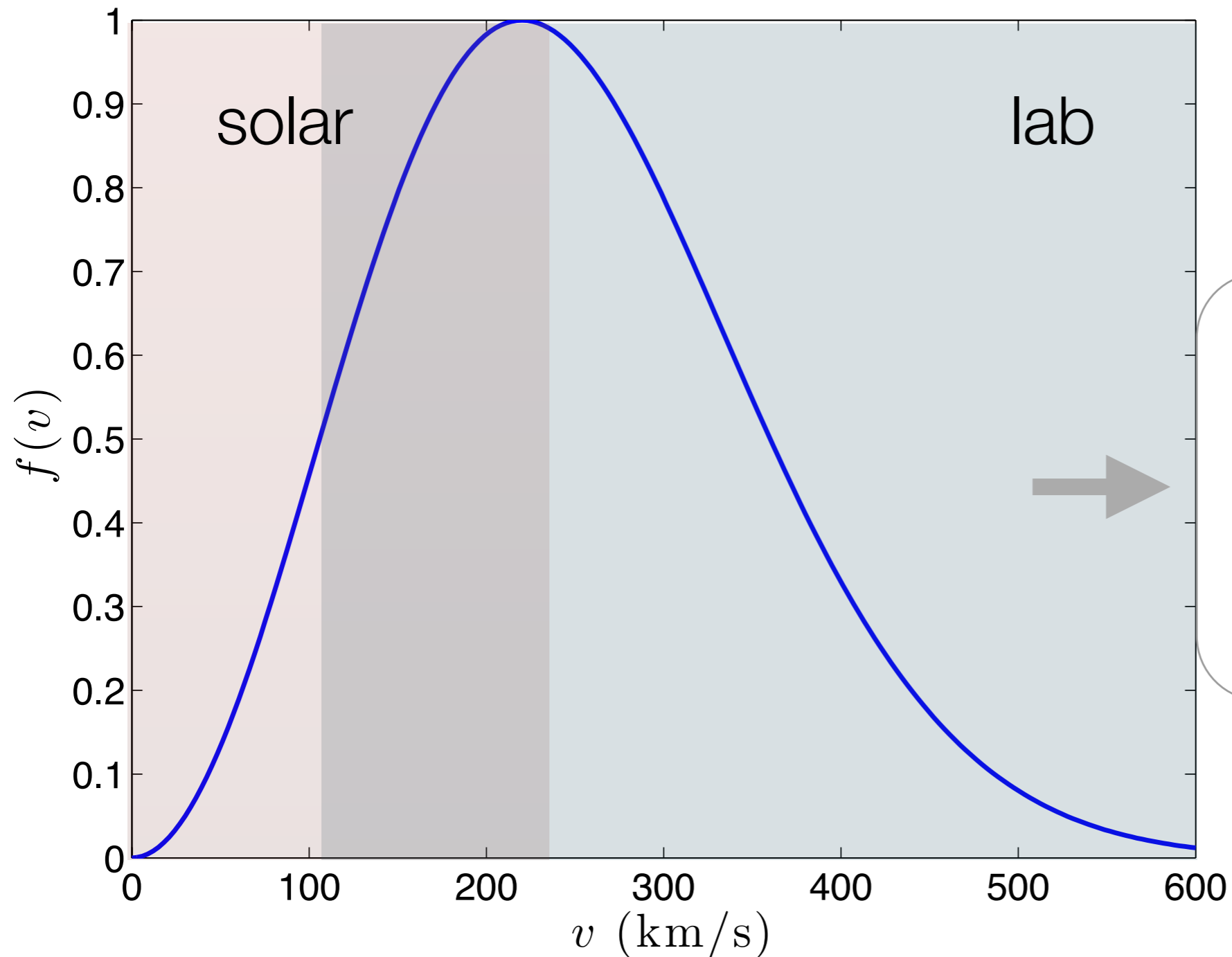
The sun is a direct detection experiment



- $M = 2 \times 10^{30}$ kg
- 73% Hydrogen
- 25% Helium
- 2% Heavier elements
(important since $\sigma_{\chi N} \propto A^2$)

Not supercooled: need to be clever about “readout”

Differences with earth-based detection

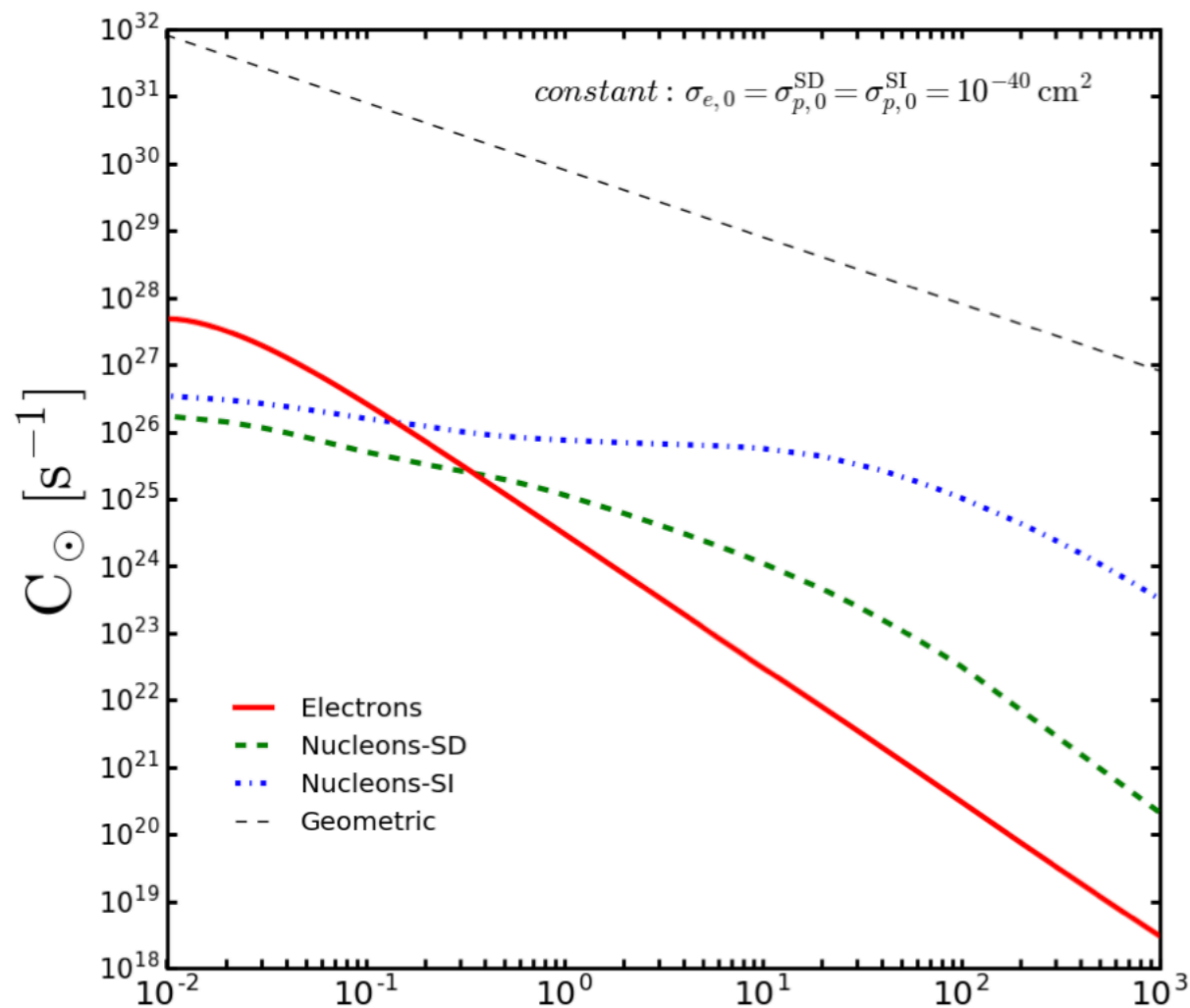


More sensitive to
lighter DM

Different particle
couplings

Capture on electrons

For certain DM masses, DM-electron scattering can efficiently capture DM



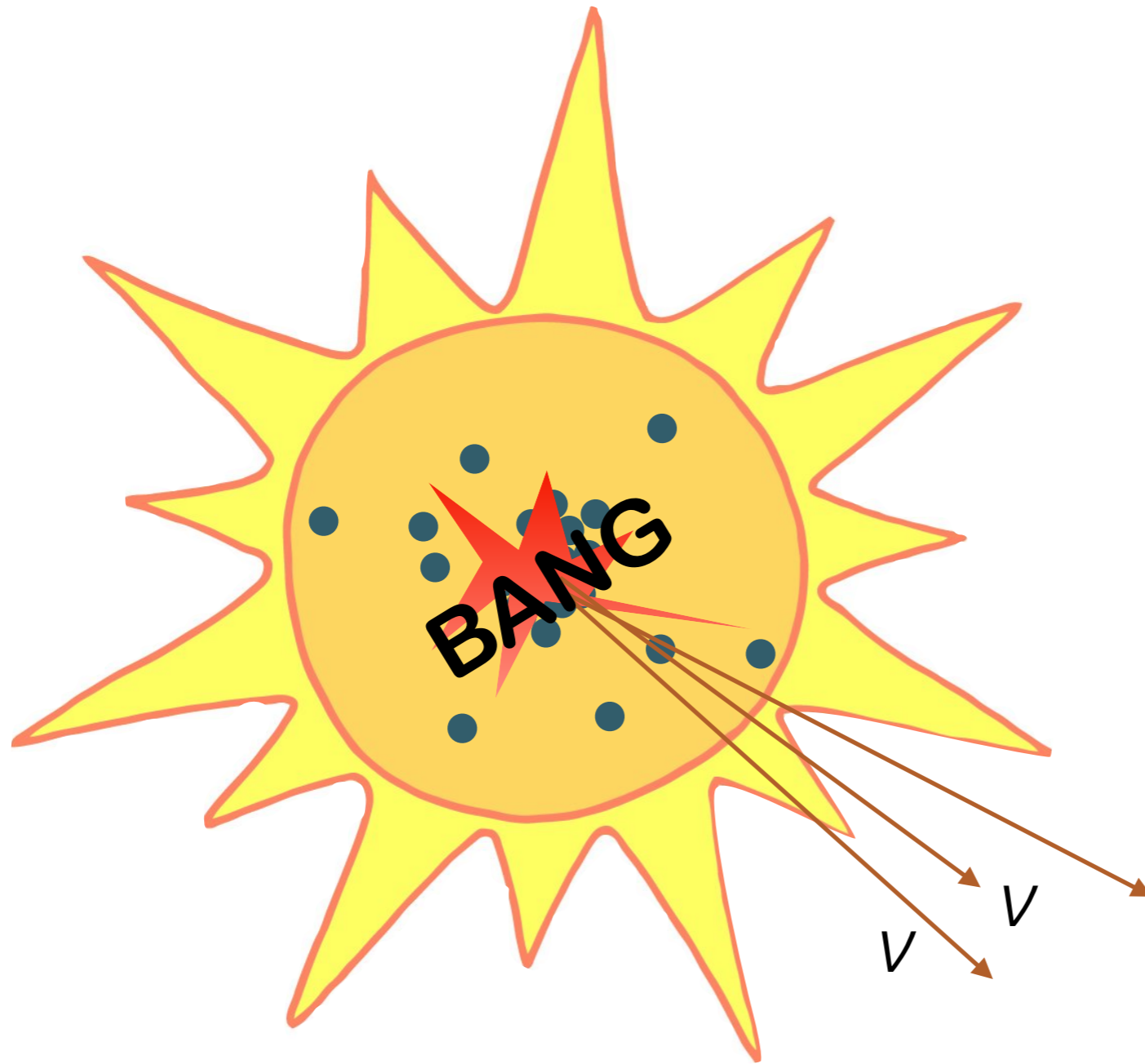
great paper

Ask Sergio



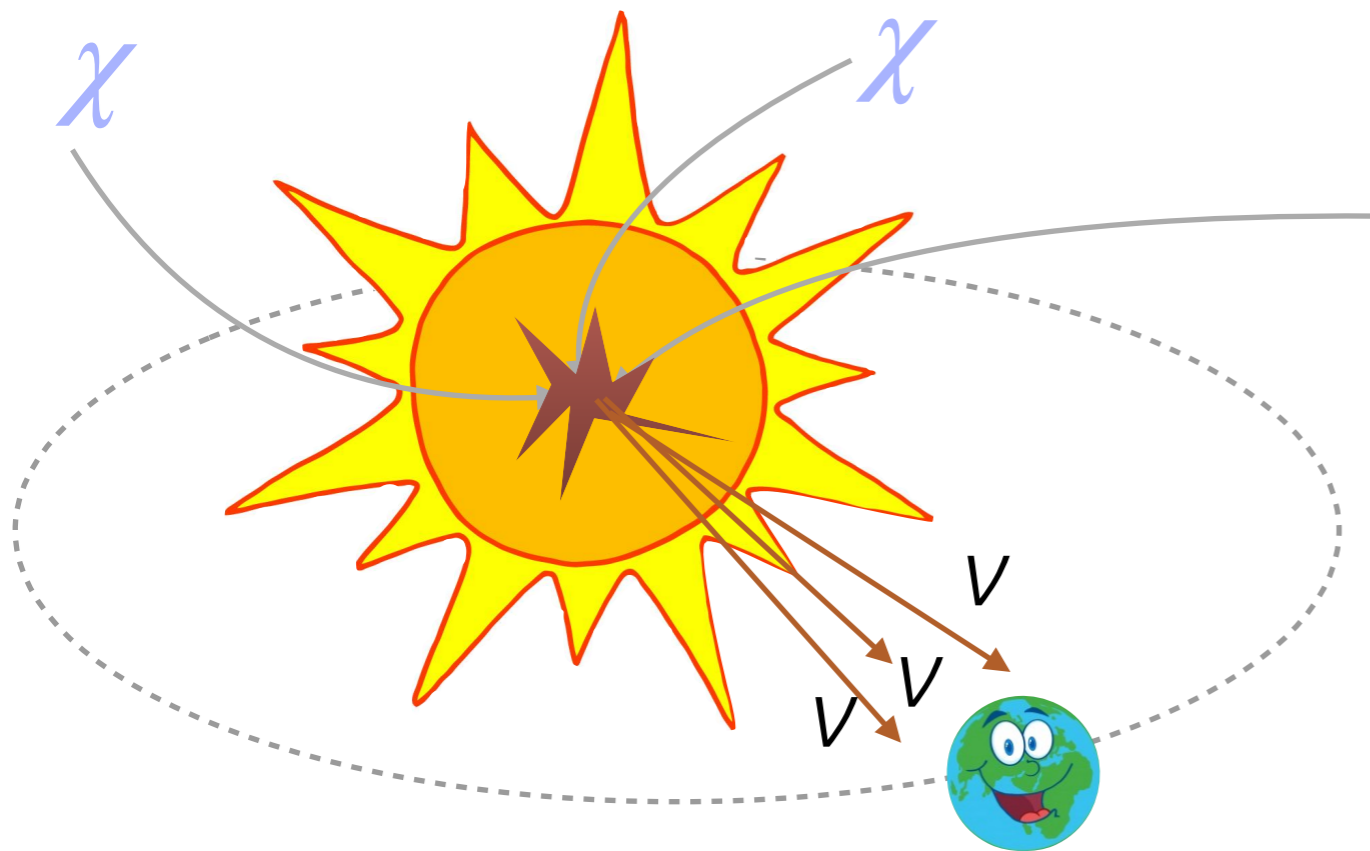
Raghuvveer Garani¹ and Sergio Palomares-Ruiz²
<https://arxiv.org/pdf/1702.02768.pdf>

If DM annihilates: look for neutrinos



Actually, you reach a steady state:
 $C(t) = A(t)$

Signals of DM in a star: annihilation

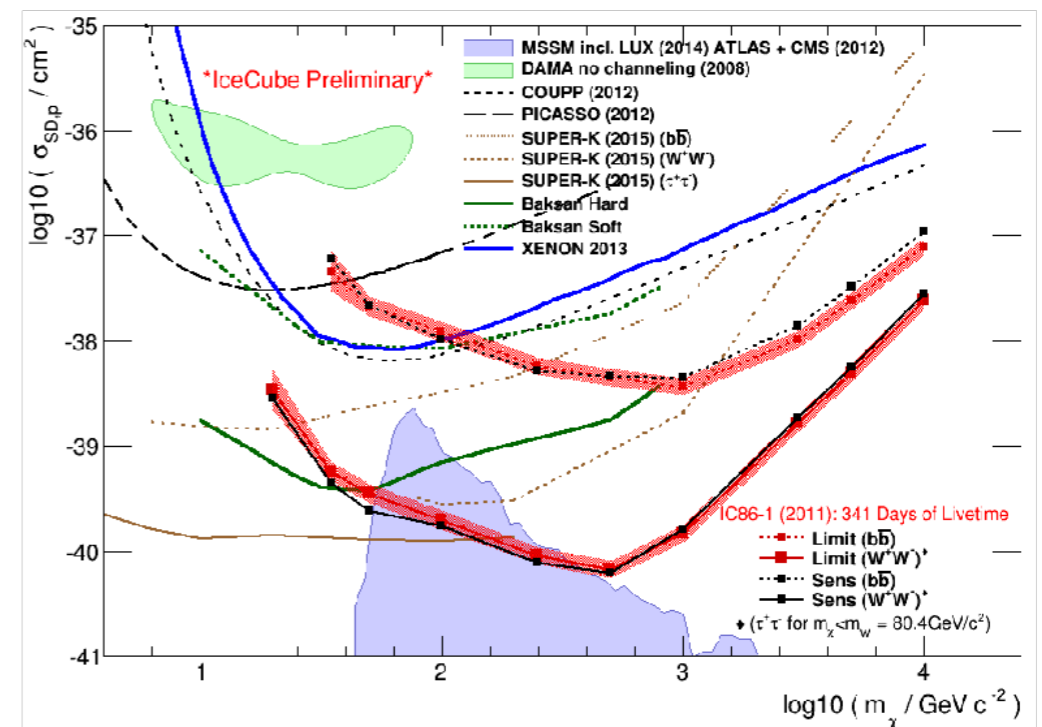


Neutrinos from annihilation can escape the sun and make their way here

Extra energy source?

—> need a **lot** of DM

but “Dark stars” could exist in very high-DM density environments



Some of the strongest limits on spin-dependent wimps are from this process (SuperKamiokande, IceCube)

Population: $\frac{dN_\chi}{dt} = C(t) - 2A(t) - E(t)$

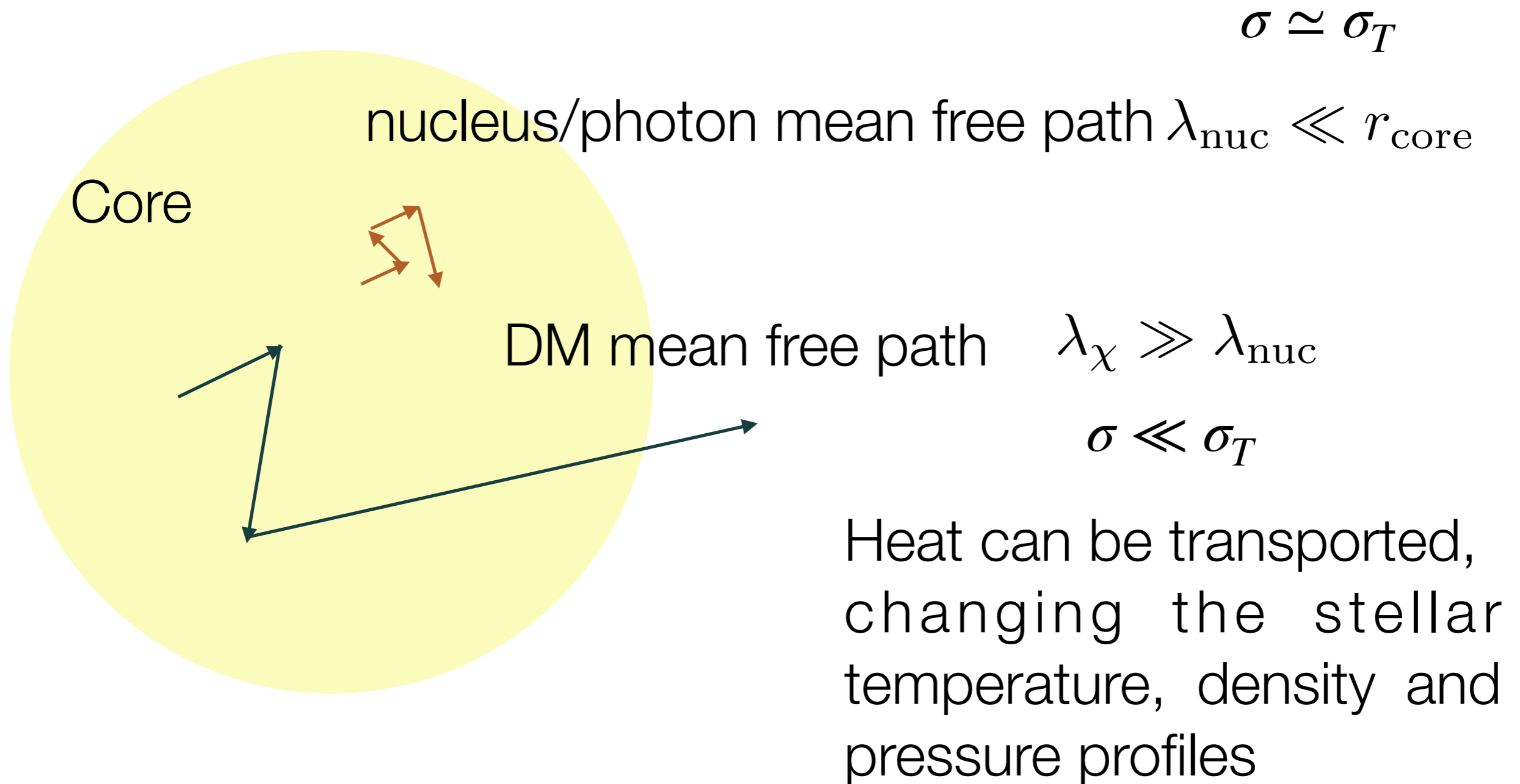
$C(t)$ Capture rate

~~$A(t)$ Annihilation rate~~

~~$E(t)$ Evaporation rate (low m)~~

No anti-DM leftover in
the universe:
asymmetric dark
matter

Asymmetric DM in stars



$$\sigma \simeq \sigma_T$$

$$\sigma \ll \sigma_T$$

Heat can be transported, changing the stellar temperature, density and pressure profiles

Can this be observed?

Probes of Solar structure

Obvious

Mass, age, radius, luminosity are extremely well-measured and are the first thing any solar model must satisfy.

Neutrinos*

pp constrained by overall luminosity, but other byproducts of pp chain extremely sensitive

to T . e.g.

$$\phi_{\nu, ^8\text{B}} \propto T_c^{25}$$

*actually this mechanism was first studied as part of the solar neutrino problem: lower than expected solar core $T \rightarrow$ fewer neutrinos

Helioseismology?

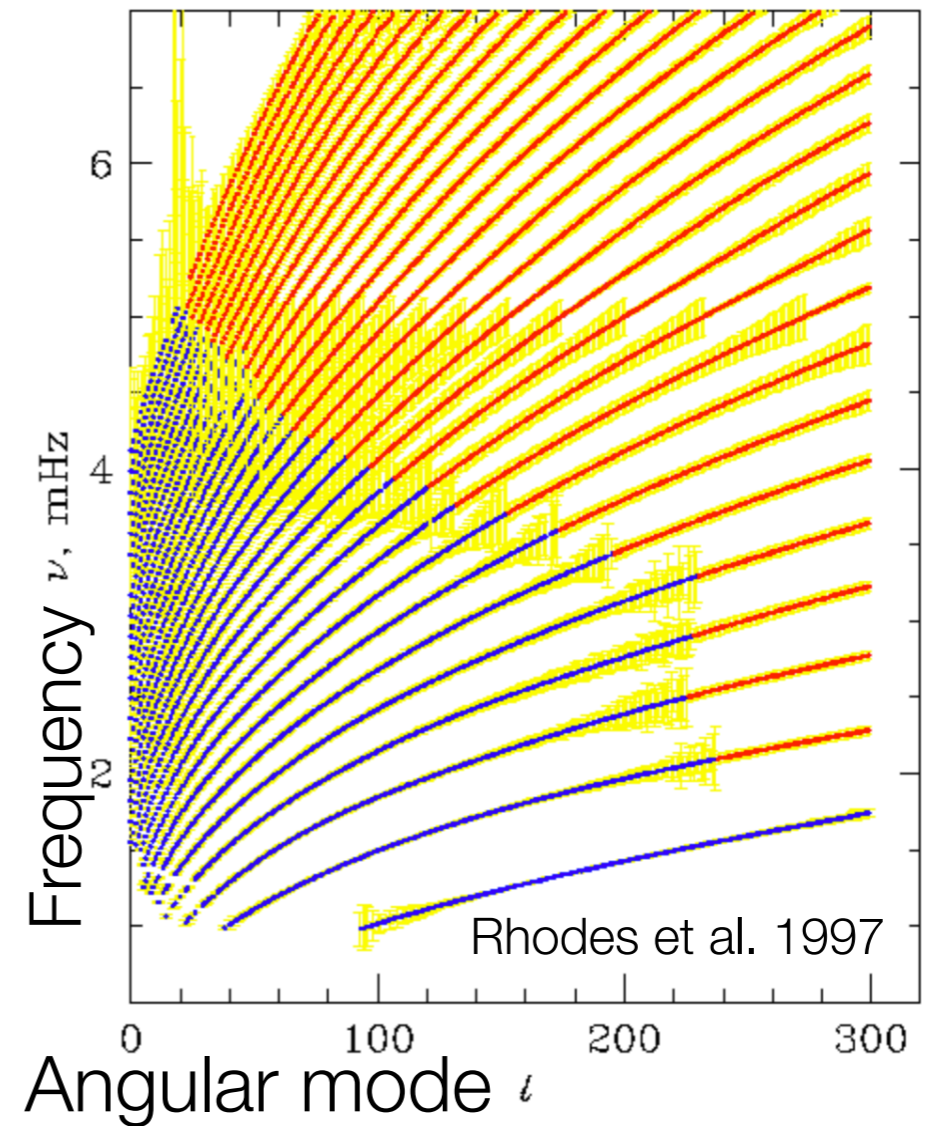
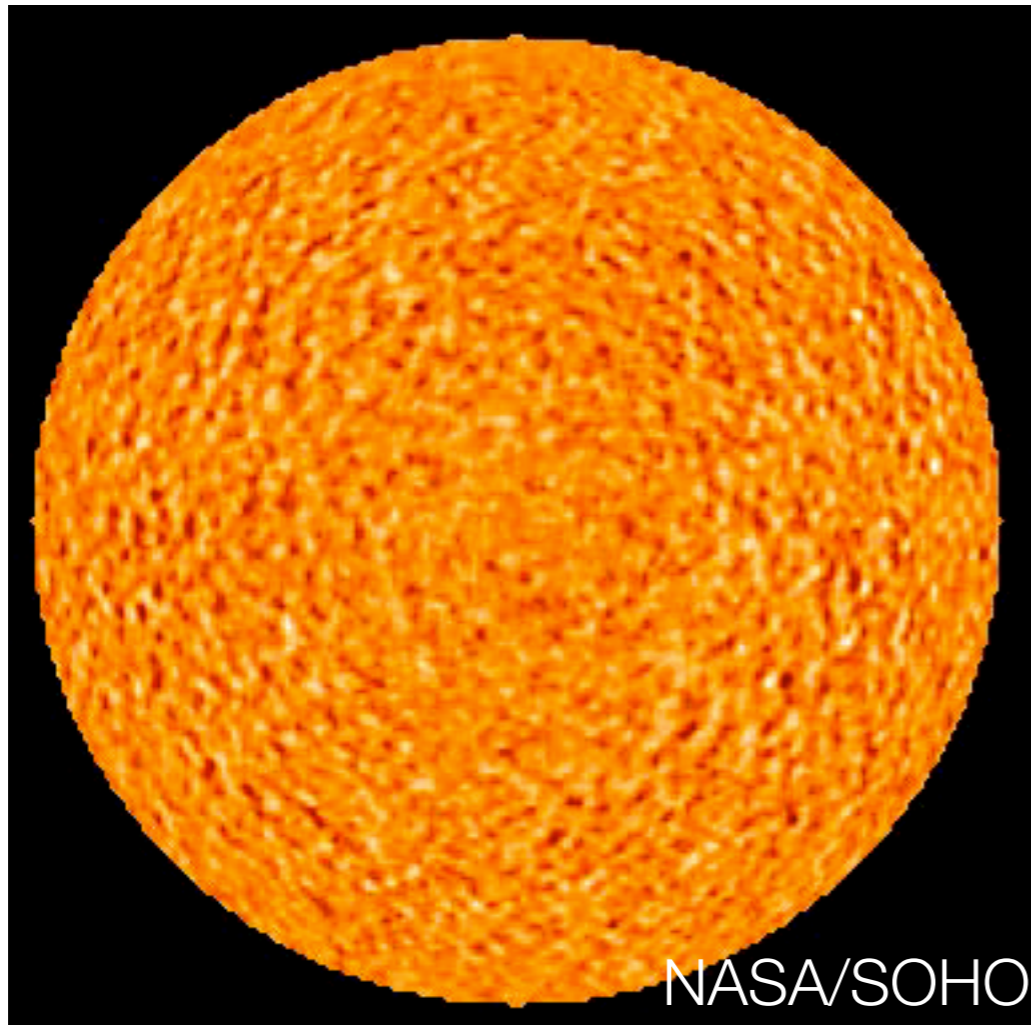


No, it's because neutrinos change flavour as they propagate to Earth

Art McDonald (Queen's)

Nobel Prize

Helioseismology



The **frequencies** of these **eigenmodes** should be predictable from:

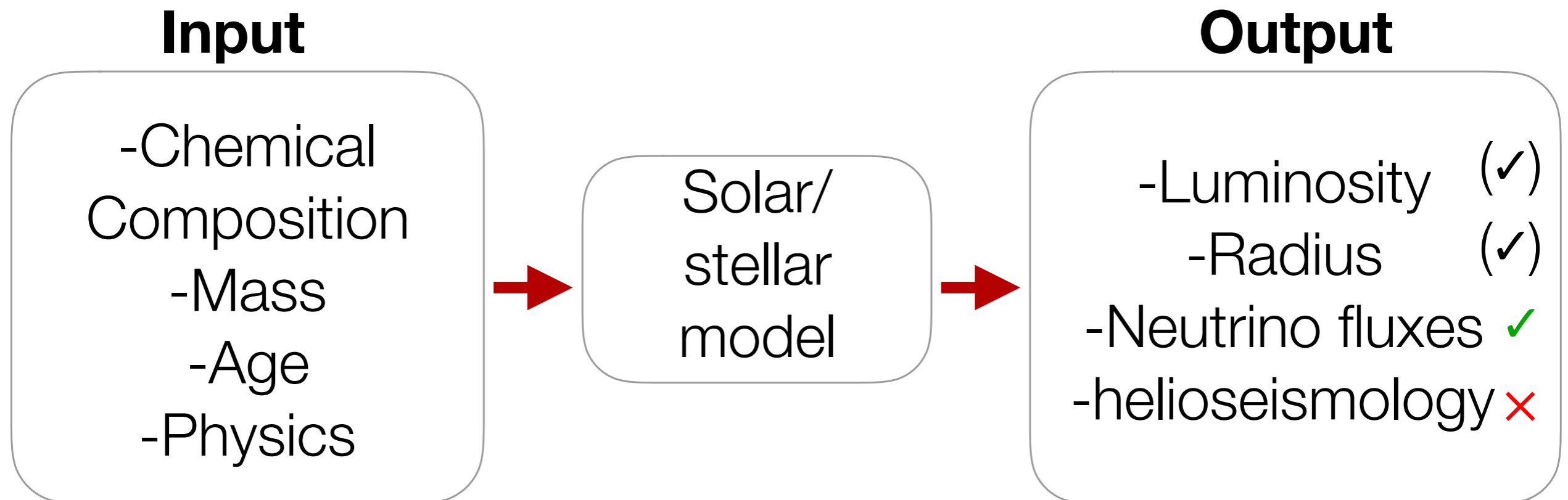
-density

-sound speed

-molecular weight (i.e. **elemental composition**)

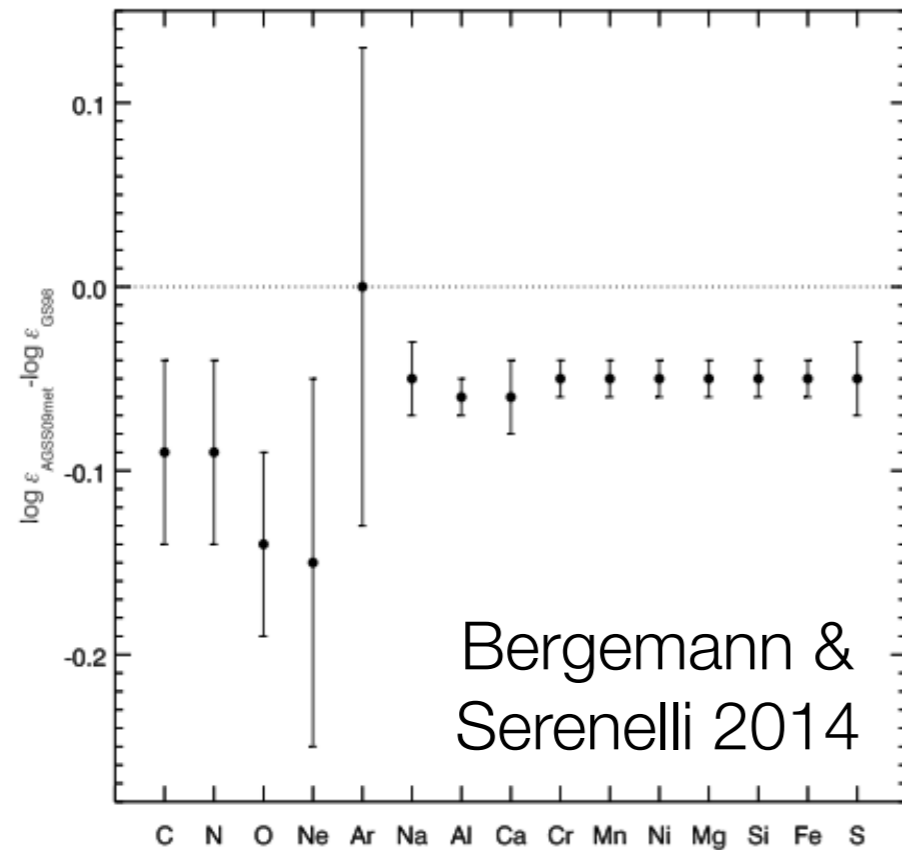
-convective zone radius

Standard Solar Model (since 1970s)

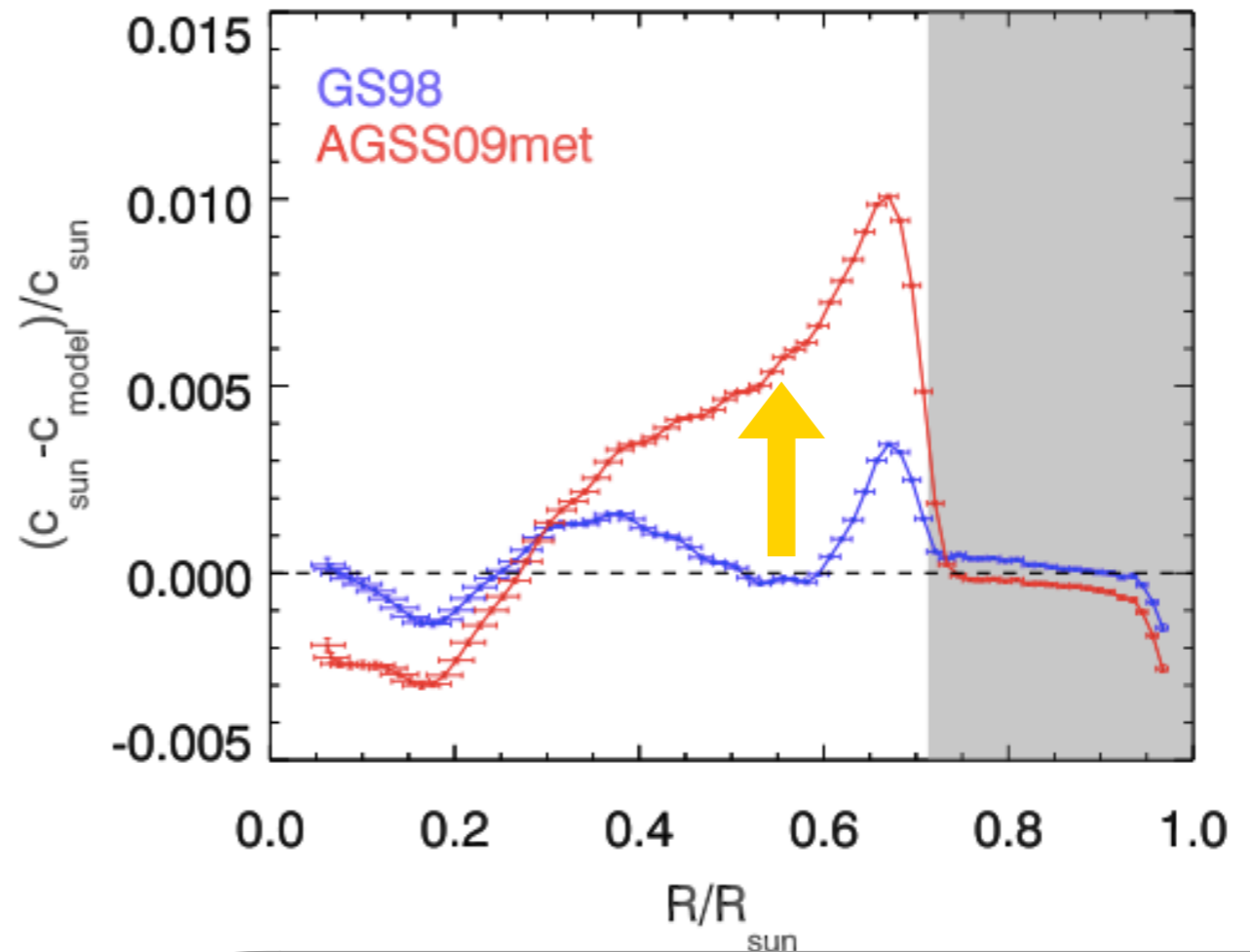


Known problem since 2004: still no solution!

Solar composition problem



revised - old abundances



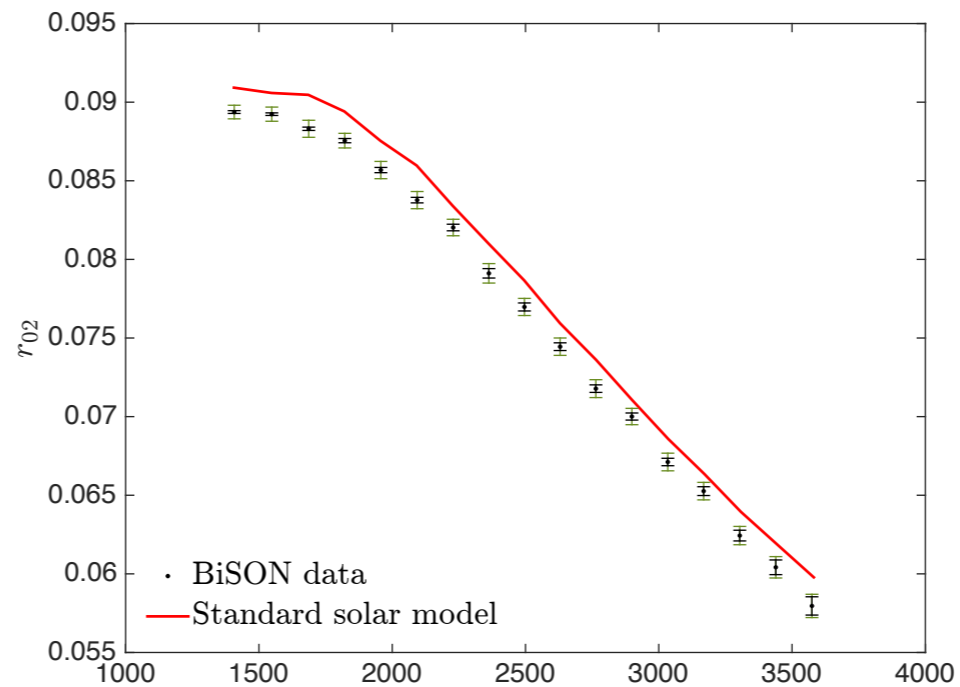
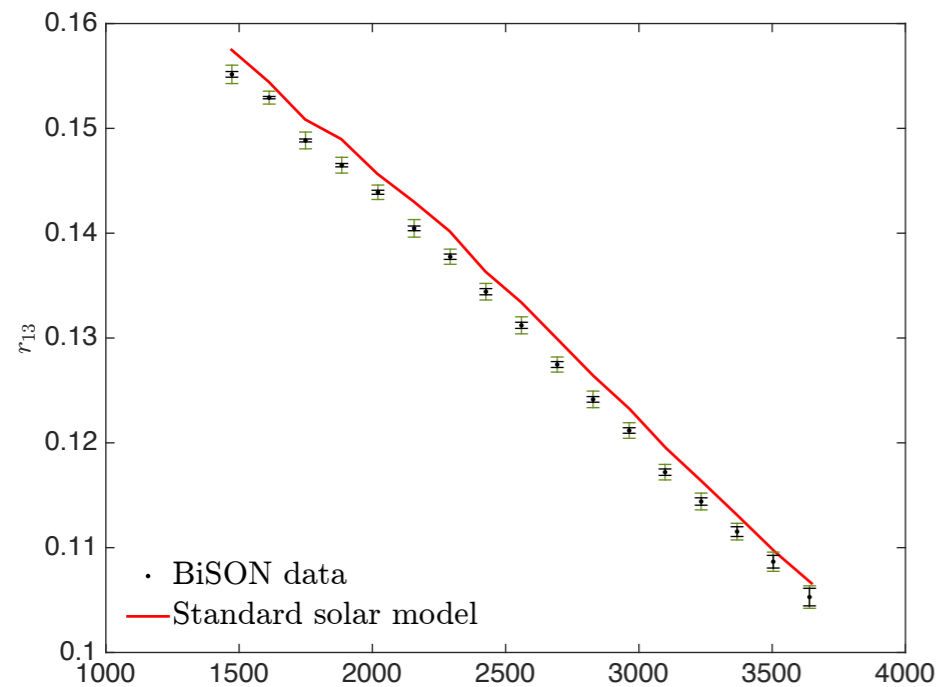
$$R_{CZ,\odot} = 0.713 \pm 0.001 R_{\odot}$$

$$R_{CZ,SSM} = 0.722 \pm 0.004 R_{\odot}$$

Mainly: smaller mean molecular weight, which shifts temperature, pressure, density gradients

Solar composition problem

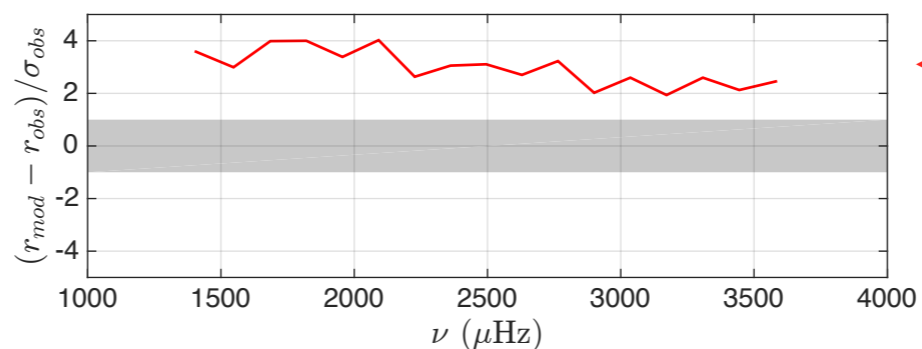
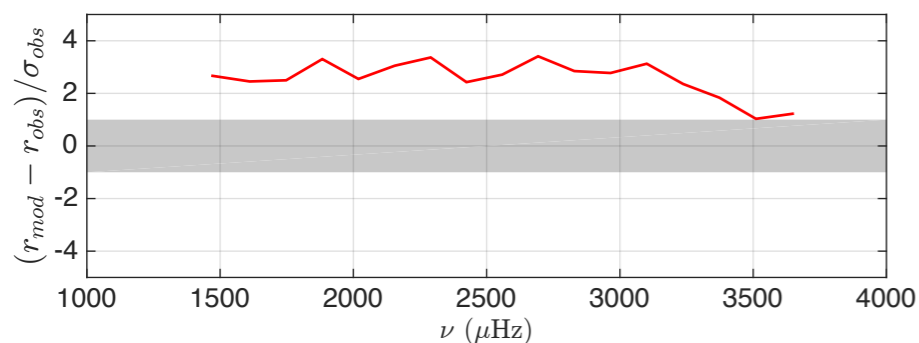
Small frequency separations: a probe of the core



$$\Delta_l(n) \equiv \nu_{n,l} - \nu_{n-1,l}$$

$$d_{l,l+2}(n) \equiv \nu_{n,l} - \nu_{n-1,l+2}$$

$$\simeq -(4l + 6) \frac{\Delta_l(n)}{4\pi^2 \nu_{n,l}} \int_0^{R_\odot} \frac{dc_s}{dr} \frac{dr}{r}$$



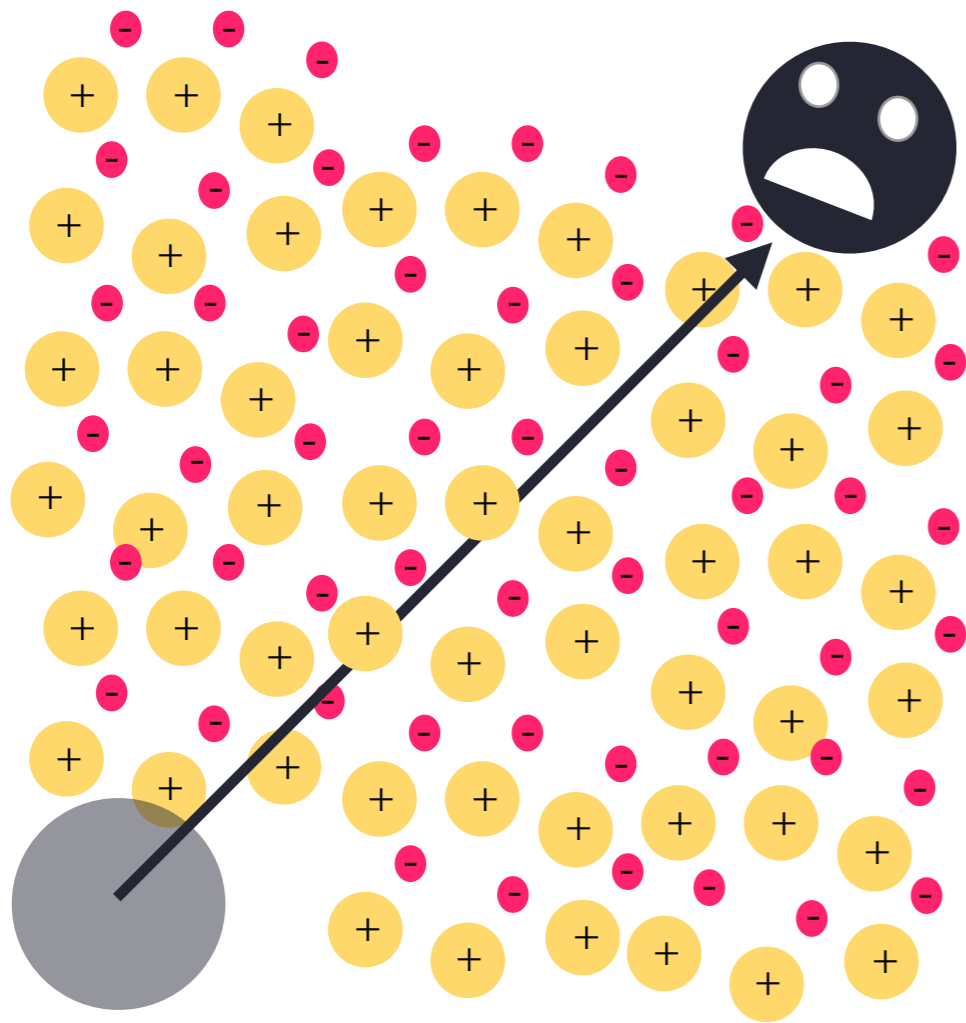
← SSM describes the core **very badly**

$$r_{02}(n) = \frac{d_{02}(n)}{\Delta_1(n)}$$

$$r_{13}(n) = \frac{d_{13}(n)}{\Delta_0(n+1)}$$

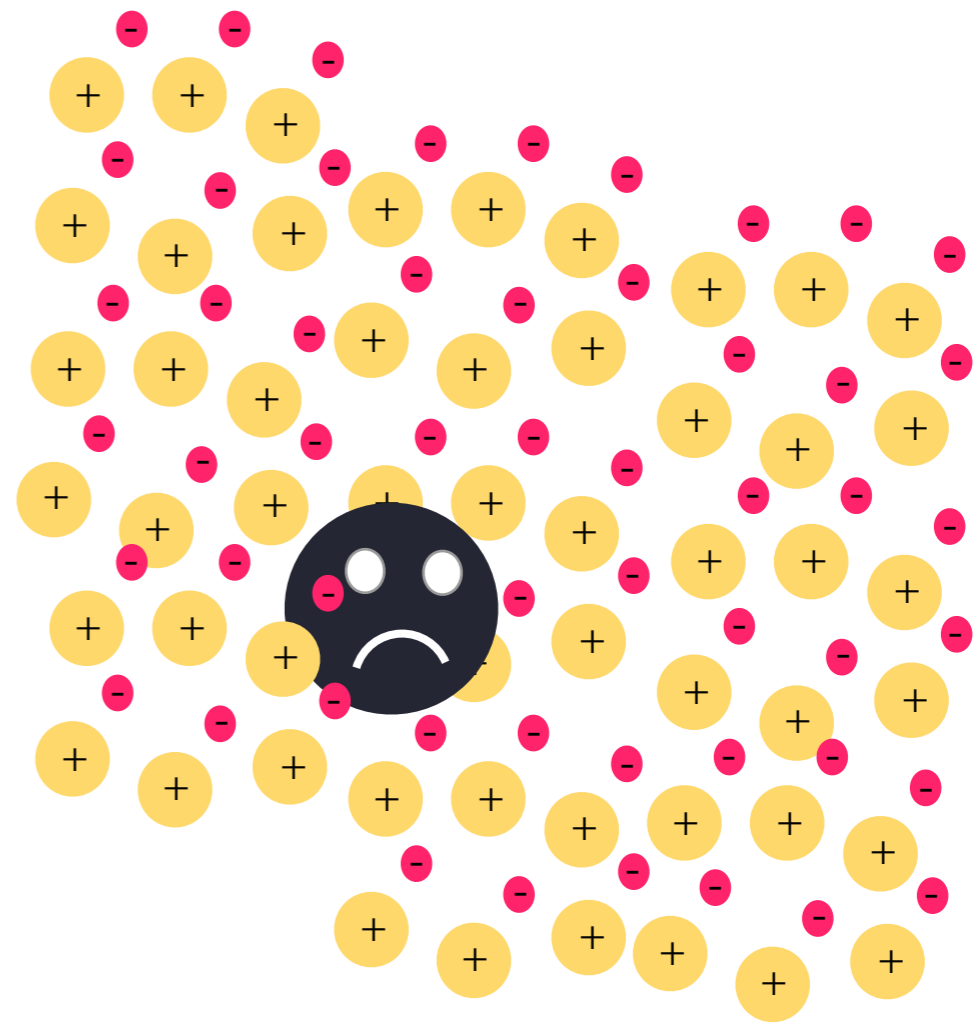
Heat transport: two regimes

Interactions too weak

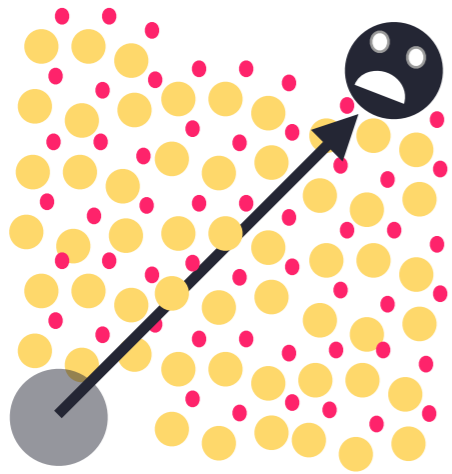


DM goes far but cannot efficiently transfer momentum

Interactions too strong



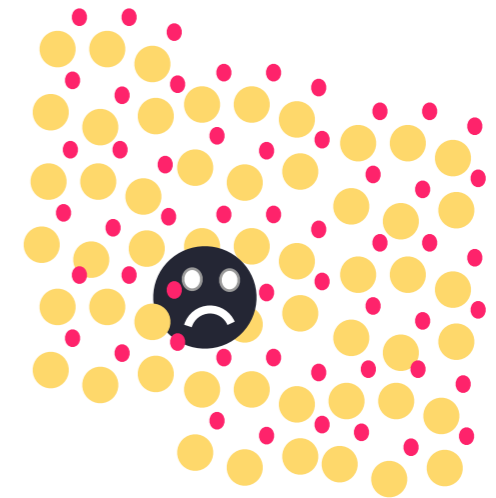
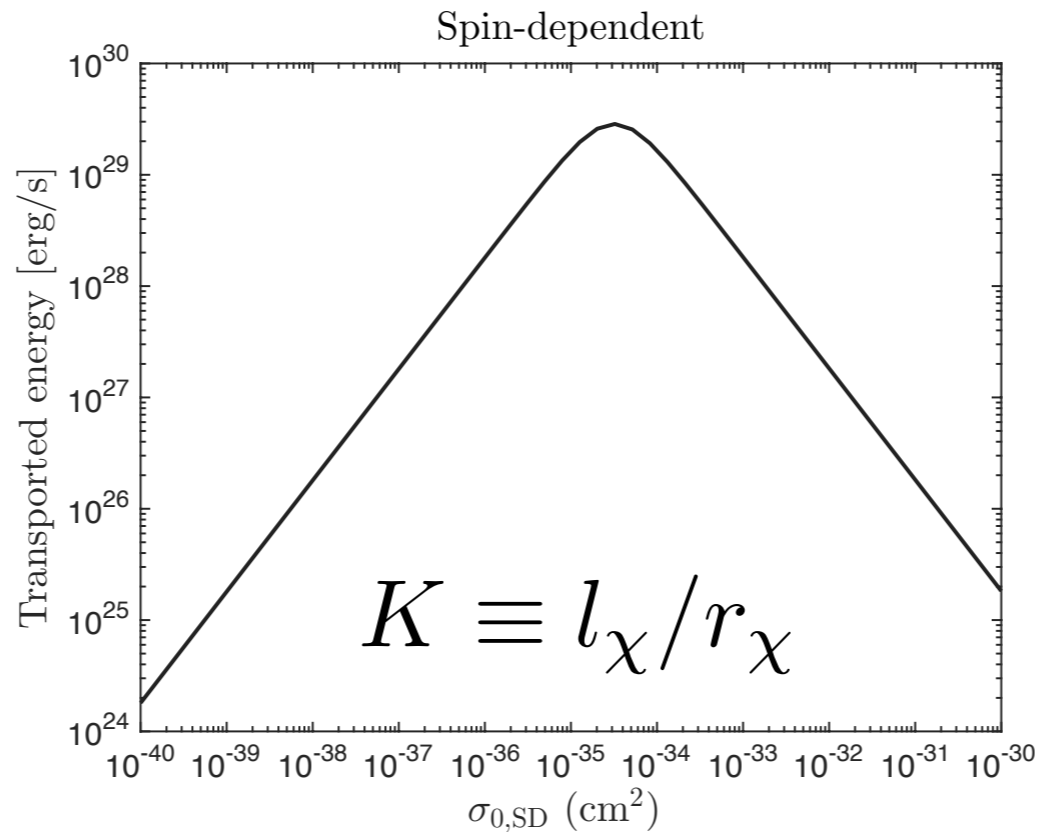
Efficient momentum transfer but DM is "stuck"



Knudsen
(non-local)

$$K \gtrsim 1$$

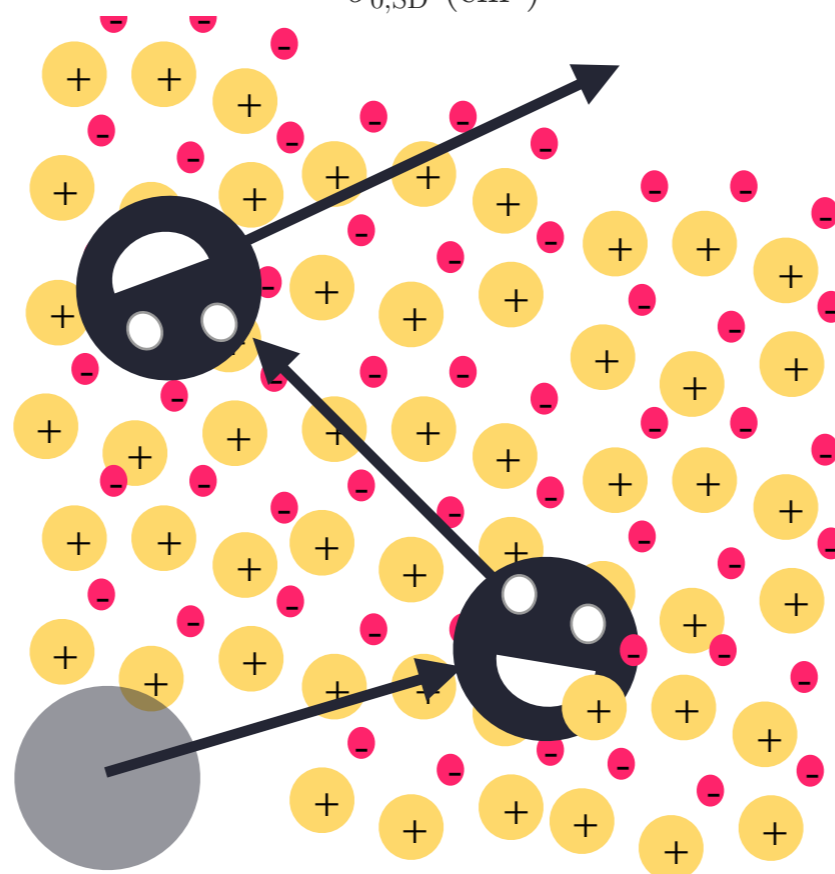
Calculable
(but wrong)



LTE

$$K < 1$$

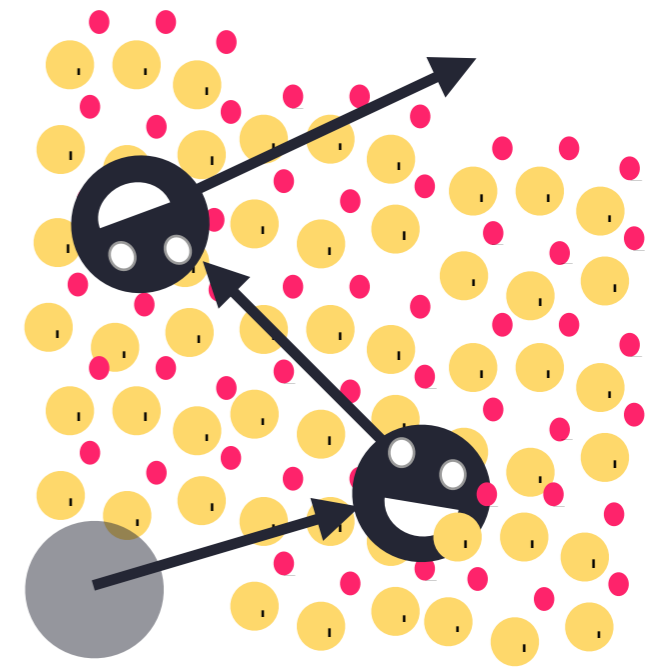
Somewhat
calculable
but unstable



optimal heat transport:
“Knudsen Peak”

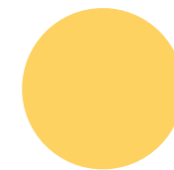
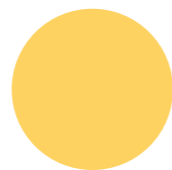
not calculable

Knudsen peak depends sensitively on the microphysical interaction (i.e. the structure of the cross section)



E.g. “billiard ball”

1/r force



1 is the loneliest operator

\mathcal{O}_1	$1_\chi 1_N$
\mathcal{O}_2	$(\vec{v}^\perp)^2$
\mathcal{O}_3	$i\vec{S}_N \cdot (\frac{\vec{q}}{m_N} \times \vec{v}^\perp)$
\mathcal{O}_4	$\vec{S}_\chi \cdot \vec{S}_N$
\mathcal{O}_5	$i\vec{S}_\chi \cdot (\frac{\vec{q}}{m_N} \times \vec{v}^\perp)$
\mathcal{O}_6	$(\frac{\vec{q}}{m_N} \cdot \vec{S}_N)(\frac{\vec{q}}{m_N} \cdot \vec{S}_\chi)$
\mathcal{O}_7	$\vec{S}_N \cdot \vec{v}^\perp$
\mathcal{O}_8	$\vec{S}_\chi \cdot \vec{v}^\perp$
\mathcal{O}_9	$i\vec{S}_\chi \cdot (\vec{S}_N \times \frac{\vec{q}}{m_N})$
\mathcal{O}_{10}	$i\frac{\vec{q}}{m_N} \cdot \vec{S}_N$
\mathcal{O}_{11}	$i\frac{\vec{q}}{m_N} \cdot \vec{S}_\chi$
\mathcal{O}_{12}	$\vec{S}_\chi \cdot (\vec{S}_N \times \vec{v}^\perp)$
\mathcal{O}_{13}	$i(\vec{S}_\chi \cdot \vec{v}^\perp)(\frac{\vec{q}}{m_N} \cdot \vec{S}_N)$
\mathcal{O}_{14}	$i(\vec{S}_N \cdot \vec{v}^\perp)(\frac{\vec{q}}{m_N} \cdot \vec{S}_\chi)$
\mathcal{O}_{15}	$-(\vec{S}_\chi \cdot \frac{\vec{q}}{m_N}) \left((\vec{S}_N \times \vec{v}^\perp) \cdot \frac{\vec{q}}{m_N} \right)$

Theories of particle interactions can give scattering cross sections that depend on the kinematical quantities.

In the non-relativistic limit

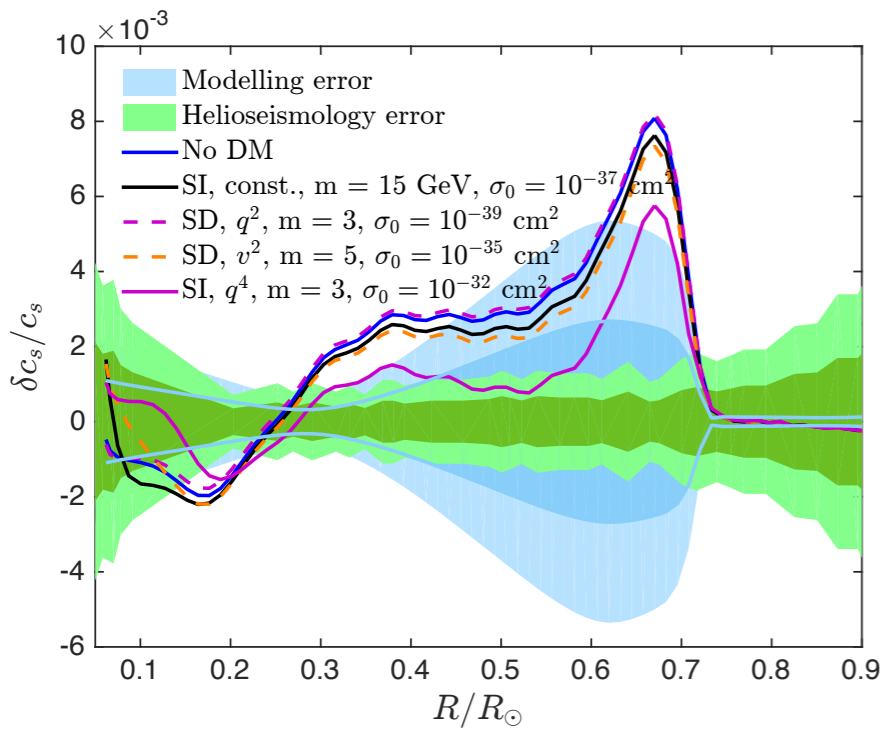
v_{rel} Relative velocity

q Exchanged momentum (scattering angle)

S_χ Dark matter spin

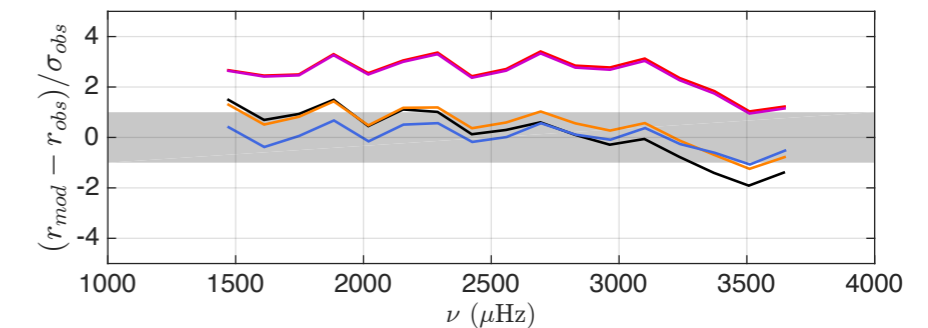
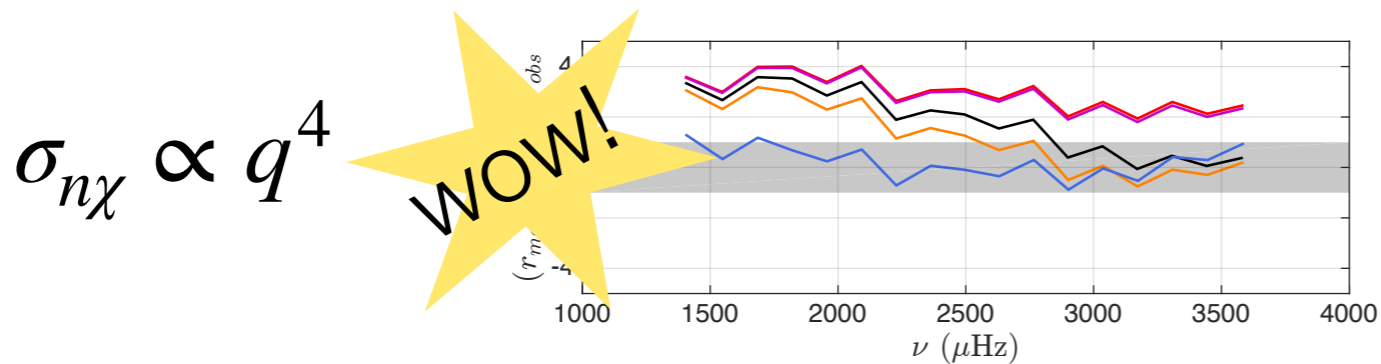
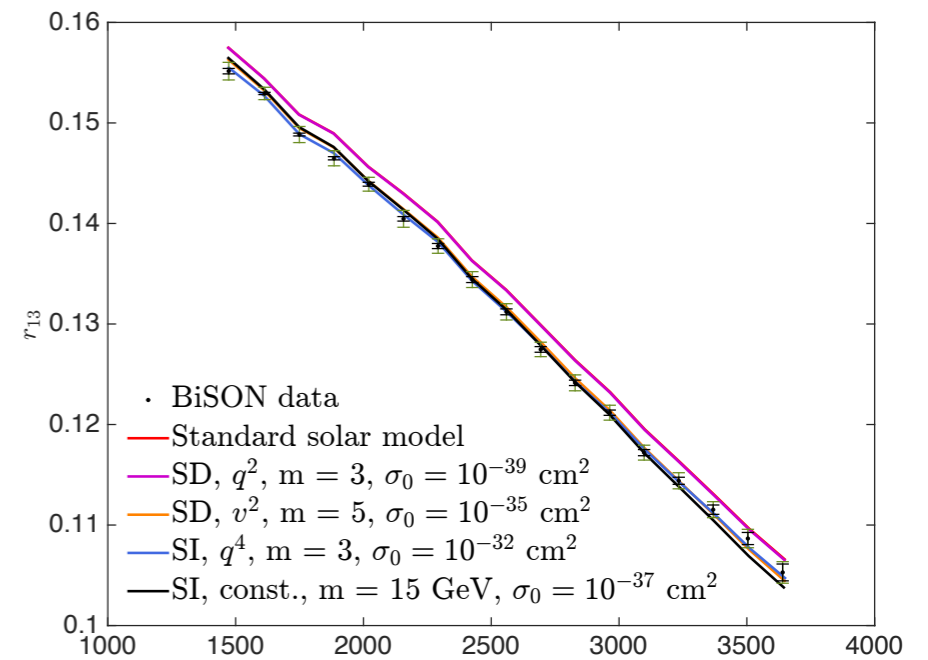
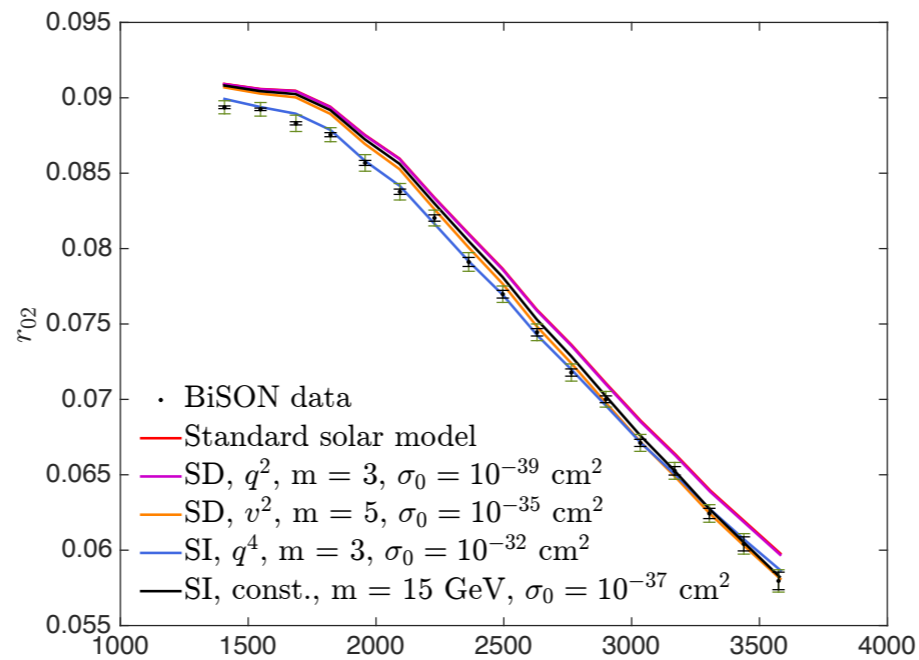
S_N Nuclear spin

In the Sun: expect very different sensitivity vs direct detection experiments.



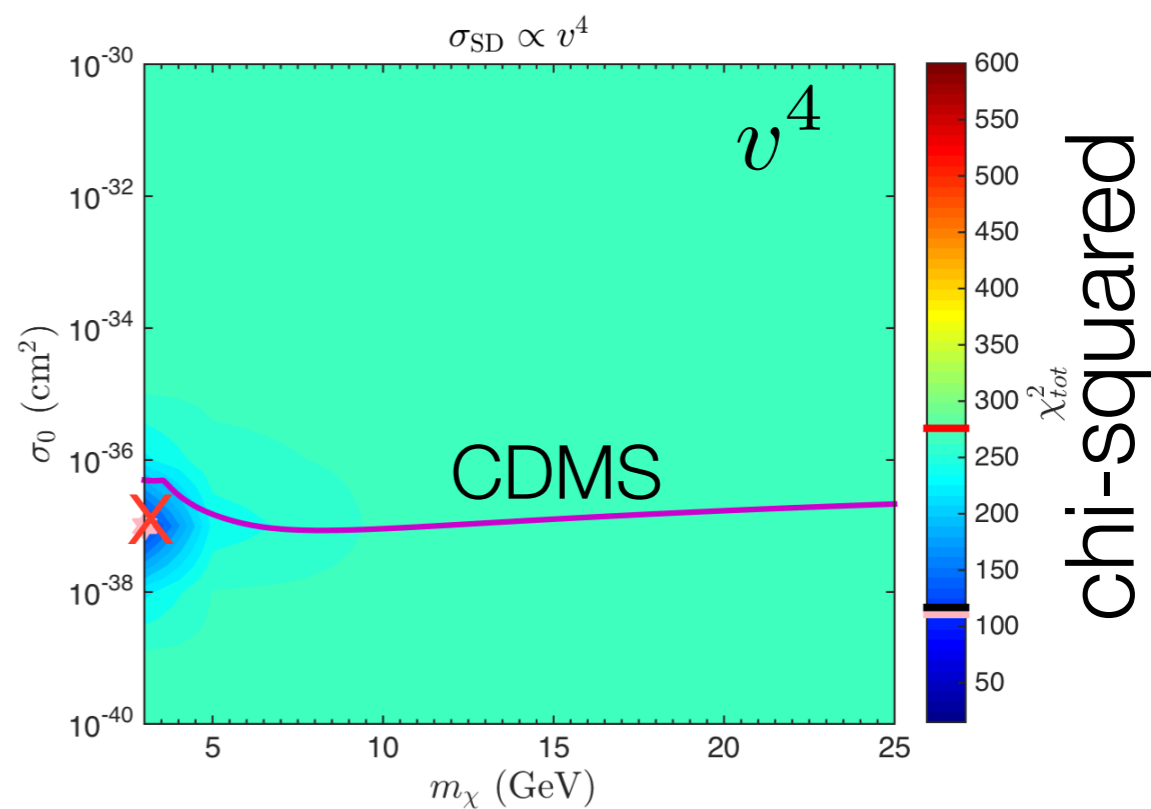
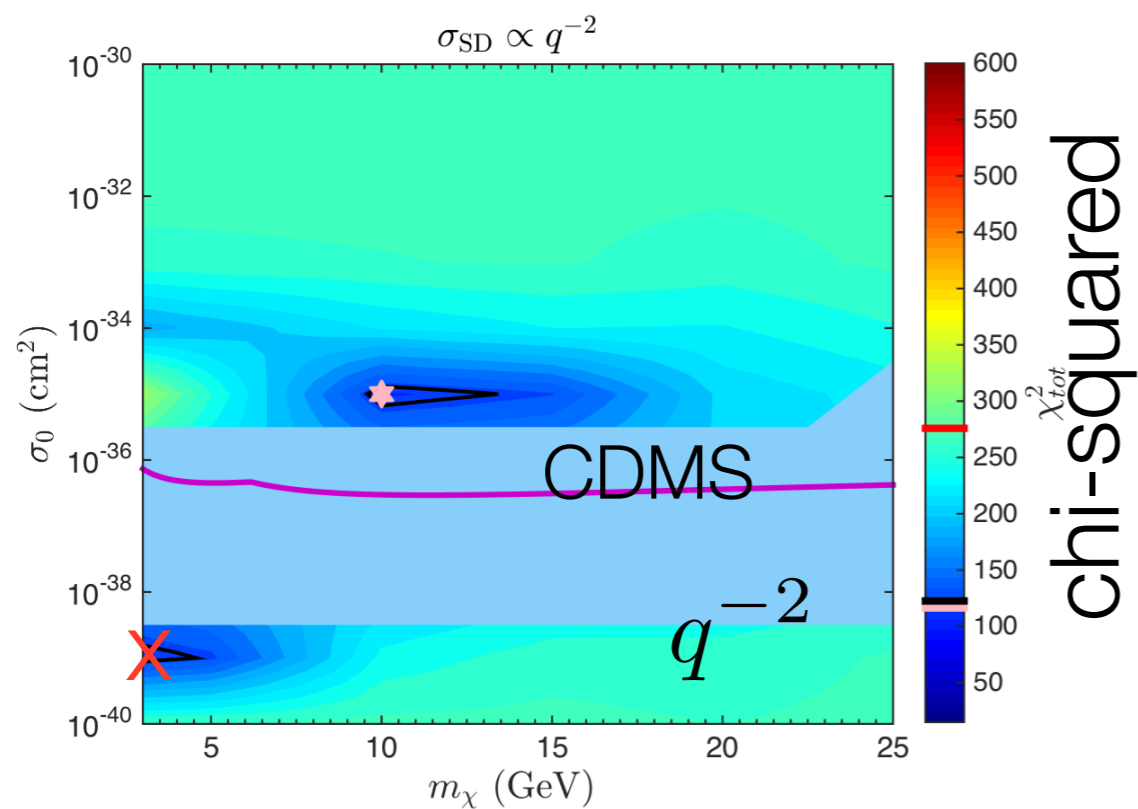
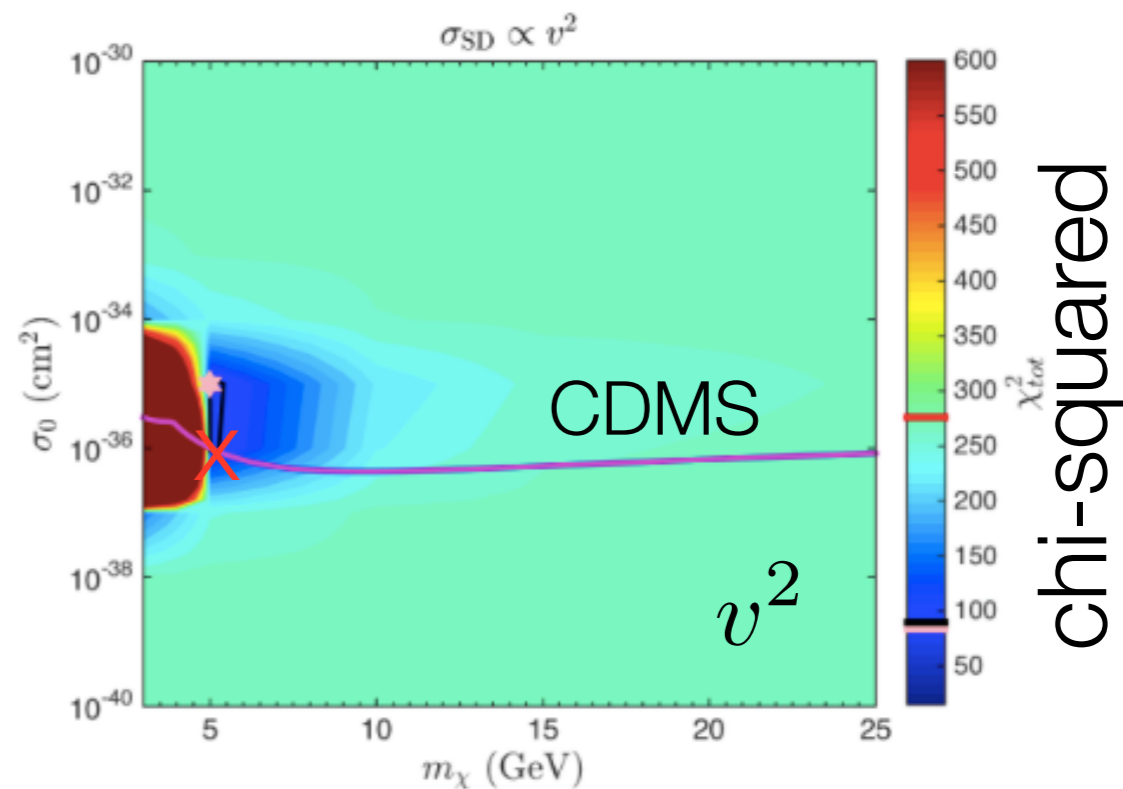
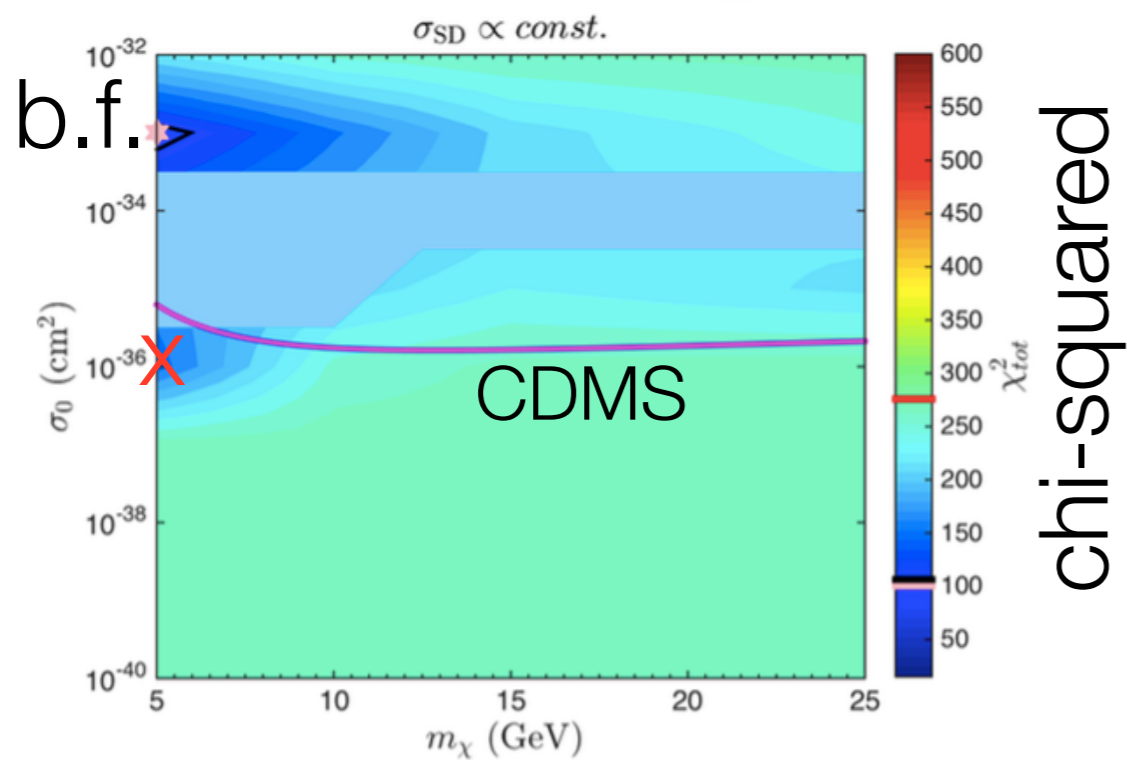
Sound speed

Core



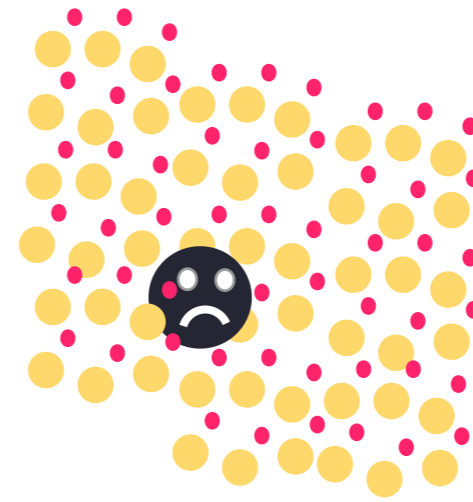
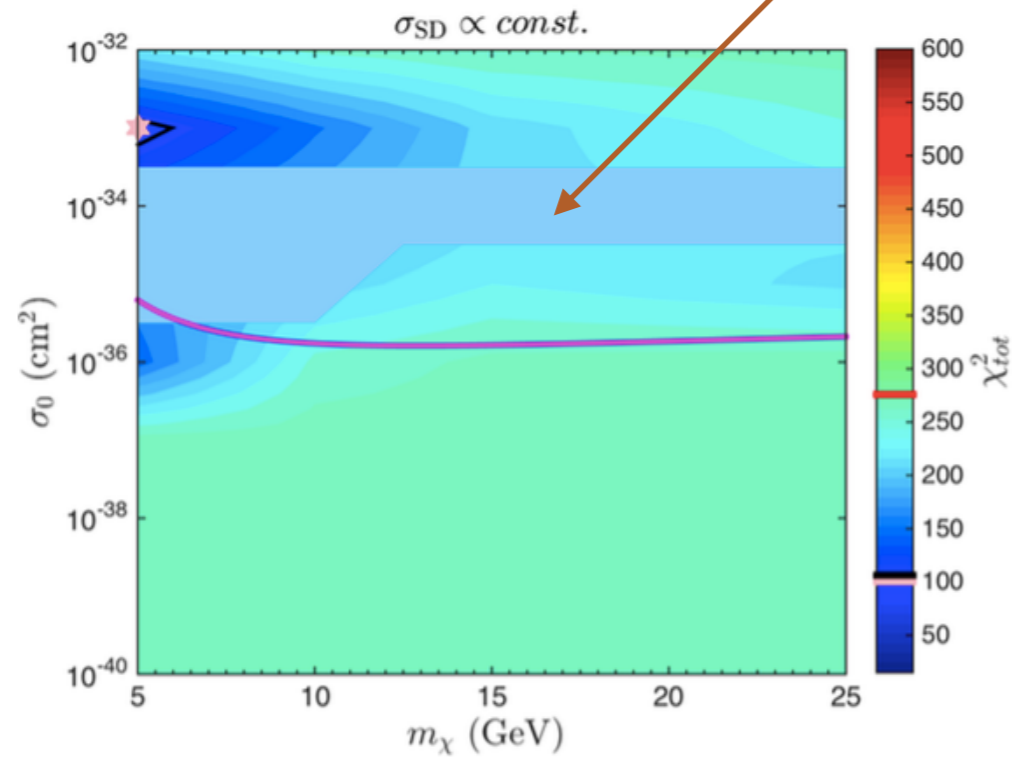
SI scattering ruled out by direct detection

Spin-dependent scenarios



— SSM — Within 3 sigma of BF \times Allowed

What's going on here?



LTE

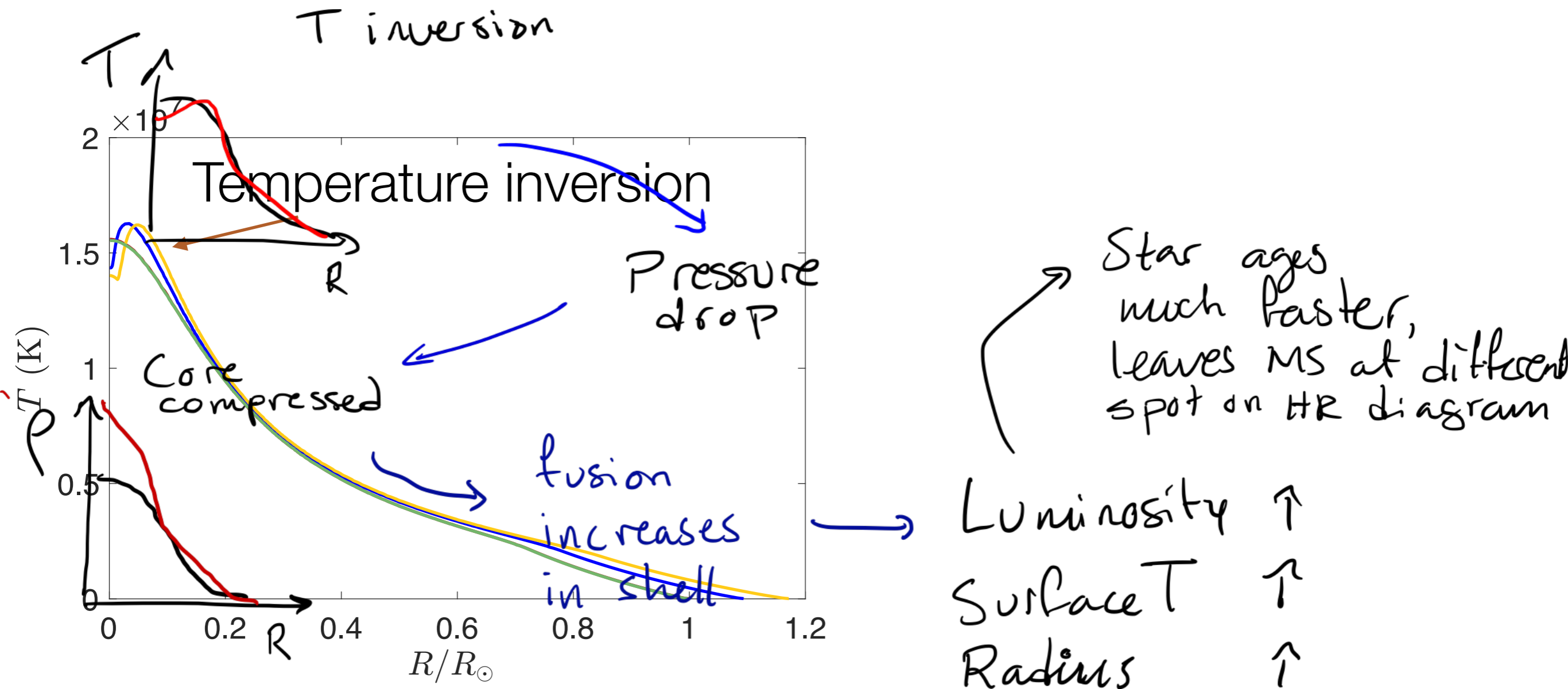
Somewhat calculable but unstable

- Very fast changes with respect to typical time scales
- Changes in structure are so strong that viable solar models (correct radius, temperature, luminosity) can't be obtained.
- **Similar effect occurs for smaller cross sections but to higher DM densities, e.g. in the galactic centre**

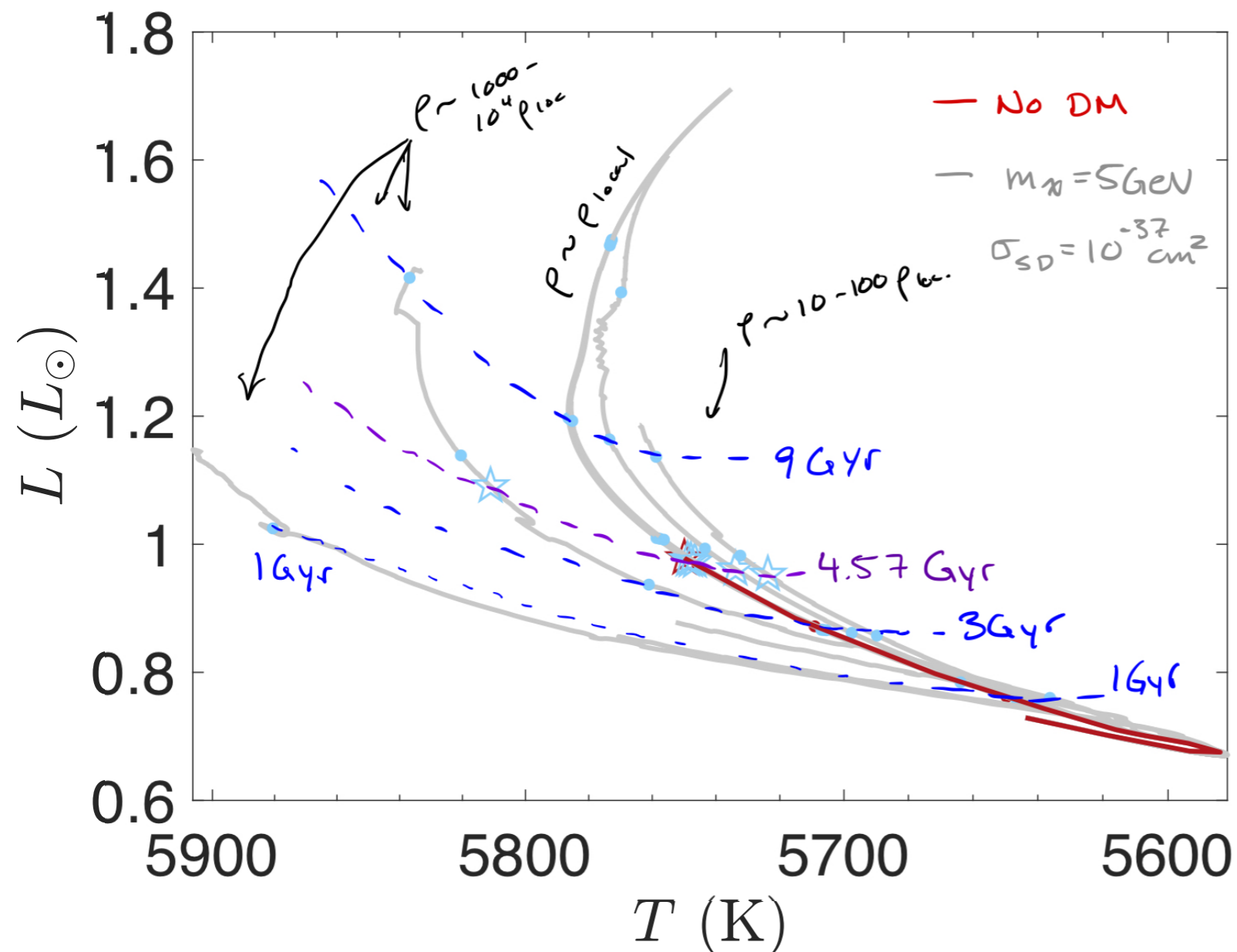


A bit more dark matter: shallower T gradient, pp fusion is “spread” a little more: star evolves a little slower with increasing DM concentrations.

A lot more DM: crazy things start to happen



Main Sequence evolution



A little more DM (10-100 times the local density): evolution slows down, star remains slightly cooler than Sun

A lot more DM (>500 times the local density): Temperature inversion causes very fast main sequence evolution: a little like the red giant phase

Weak interactions in DM-rich regions

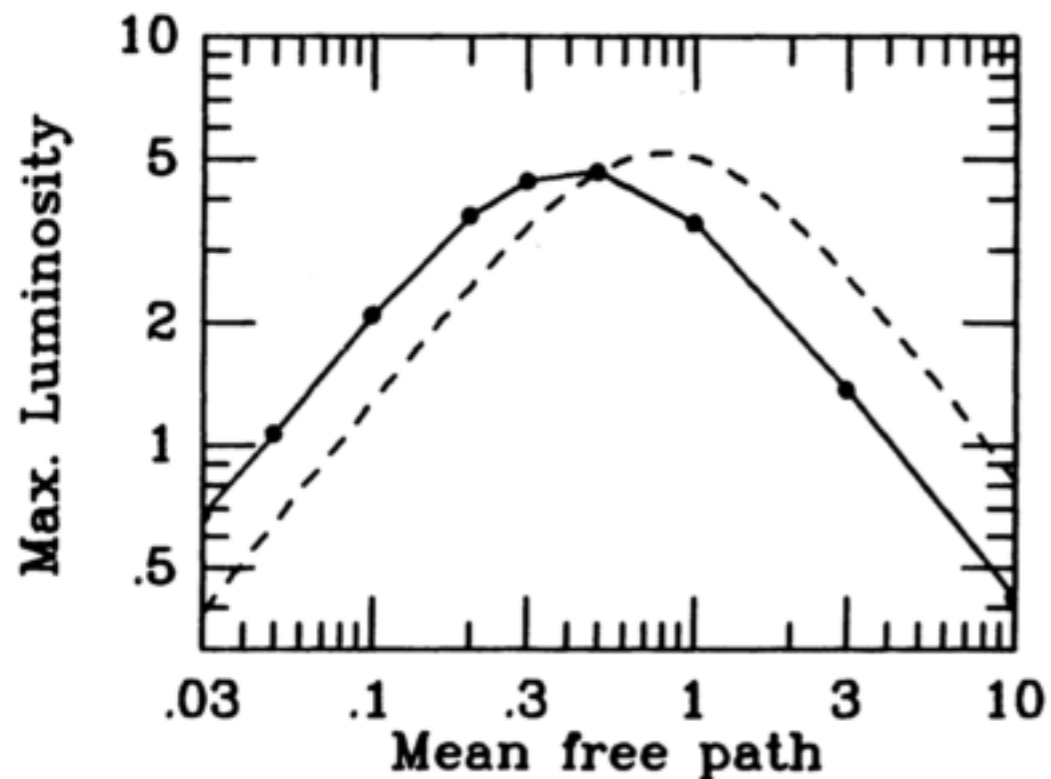
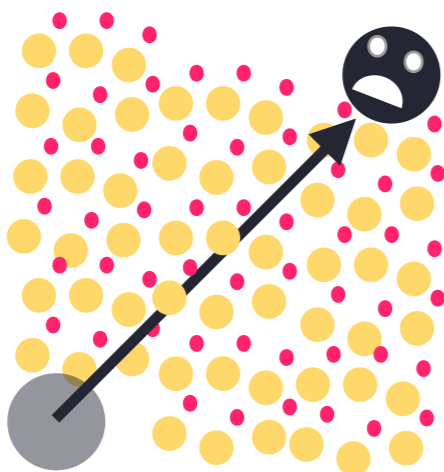
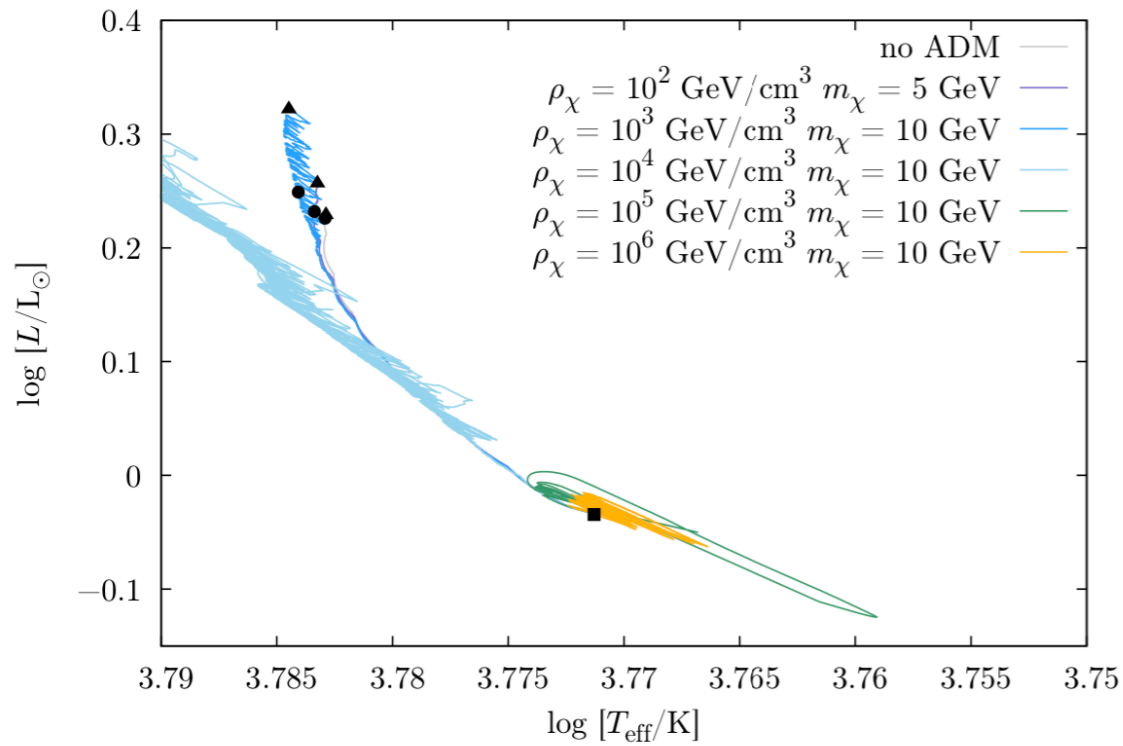


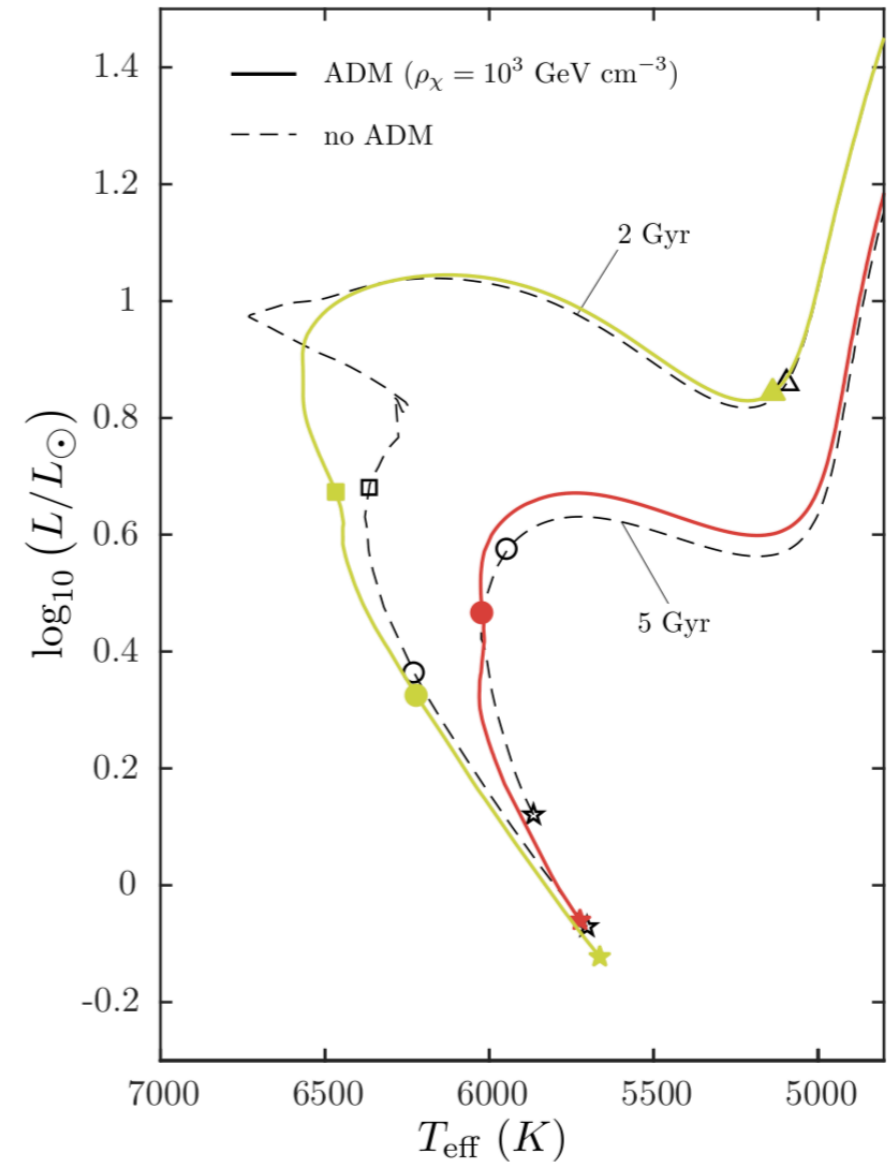
FIG. 11.—Maximum luminosity carried by WIMPs in our “generic” stellar model described in Fig. 8. The maximum refers to the luminosity curve 8c and the corresponding curves for other mean free paths. Each of the black dots corresponds to a Monte Carlo run for that value of the mfp. The dashed line is the formula used by Gilliland *et al.* (1986) in their investigation of WIMP heat transport in the Sun.

You can either use the nice, analytically tractable, wrong results from **Press and Spergel** (1987) or the correct, semiempirical, unstable solution from **Gould and Raffelt** (1990)

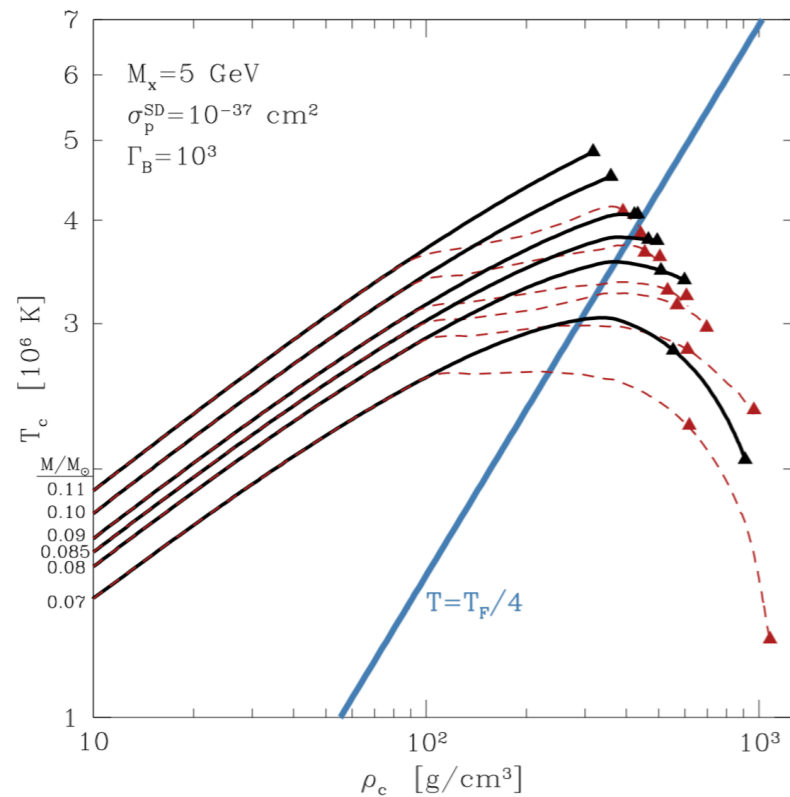




Lopés and Lopés (2019)
use Spergel & Press transport



locco et al. 1201.5387
Dangerous instabilities

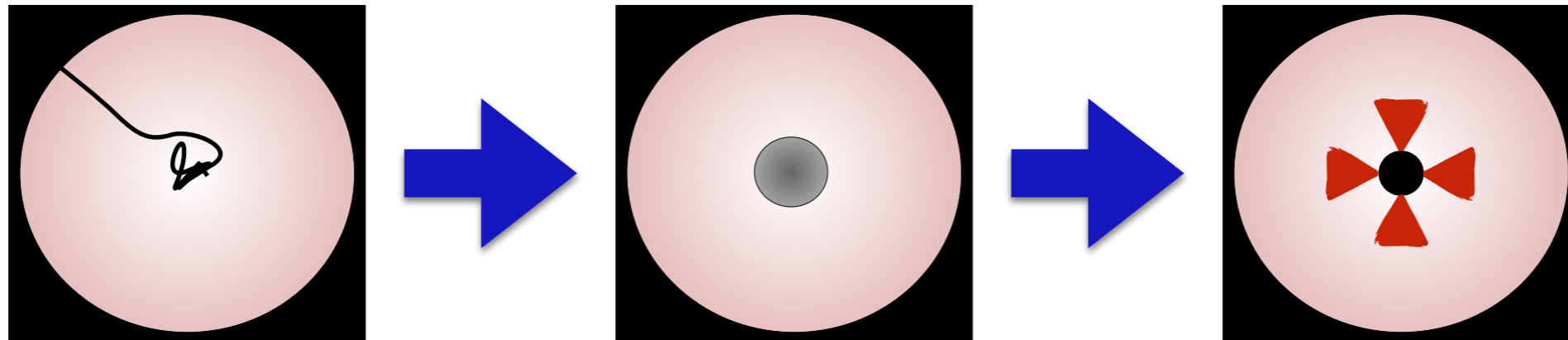


Hurst & Zentner
2011: Spergel &
Press transport,
some iffy
approximations

Something is going on
but nobody has managed to
compute it correctly yet

Weird stars

Dark matter ignition of Type Ia SNe



Heavy asymmetric DM
captured by White Dwarf

DM thermalizes at
small radius

DM sphere
collapses

DM sphere collapse

OR

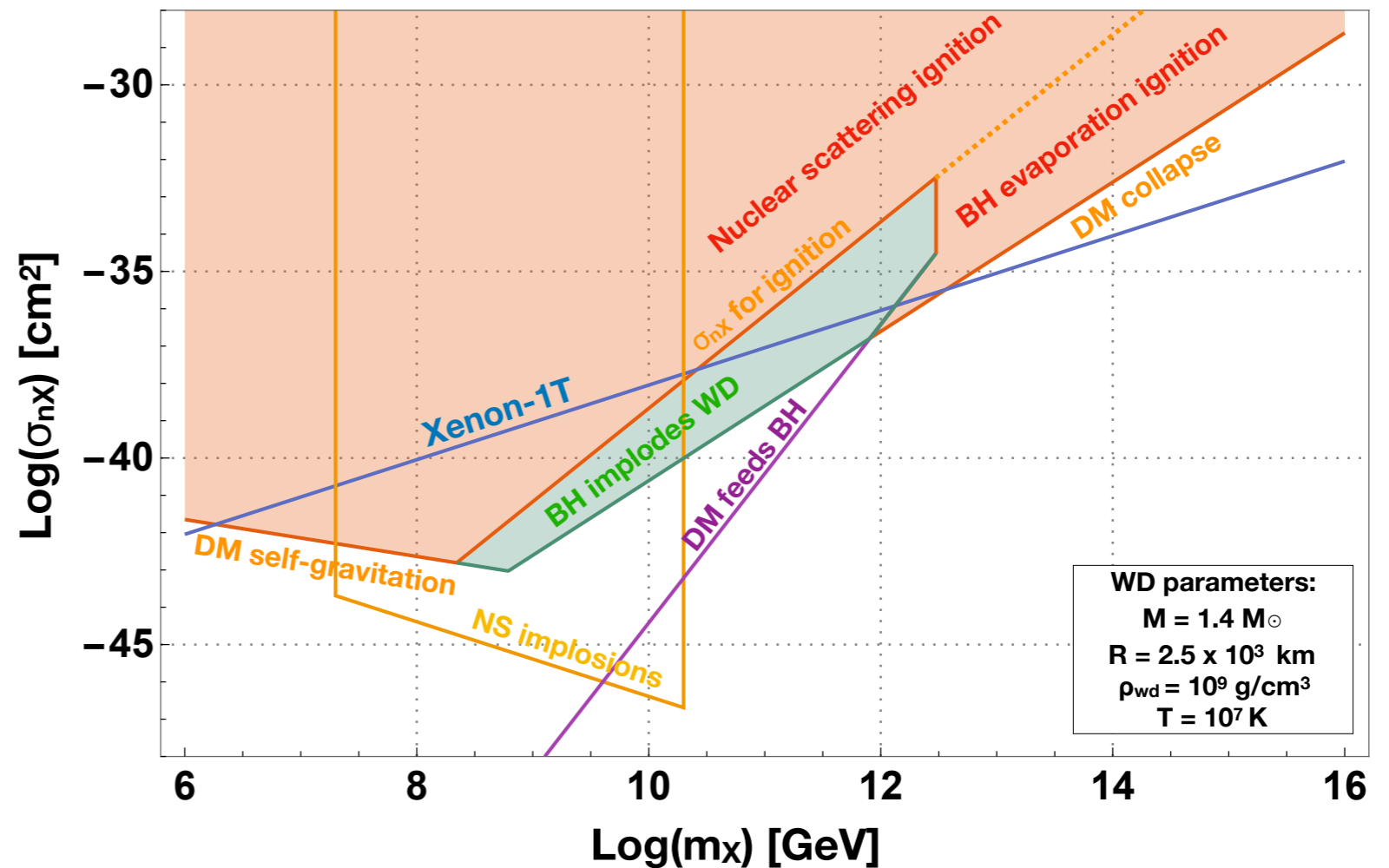
BH evaporation

heats core to $\sim 10^{10}$ K

**Type Ia SN
is ignited in sub-
Chandra
mass WDs**

Bounds from observation of a not-exploded WD

WD SDSS
J160420.40+05
5542.3
Age **~3.3 Gyrs**

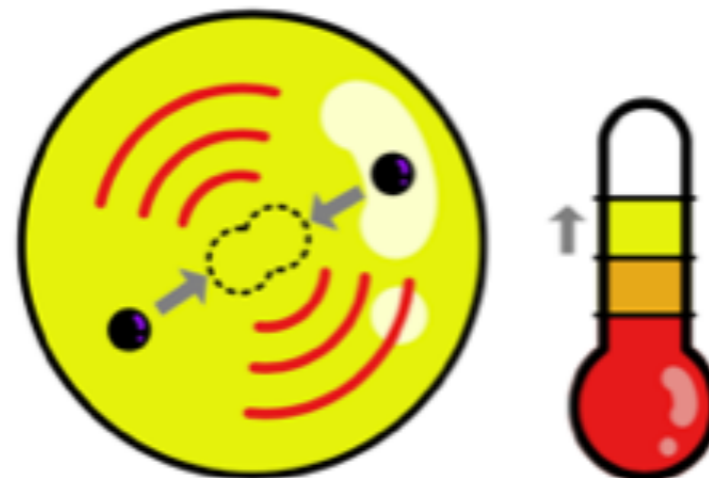


Neutron stars (Nirmal Raj)

Minimum signature

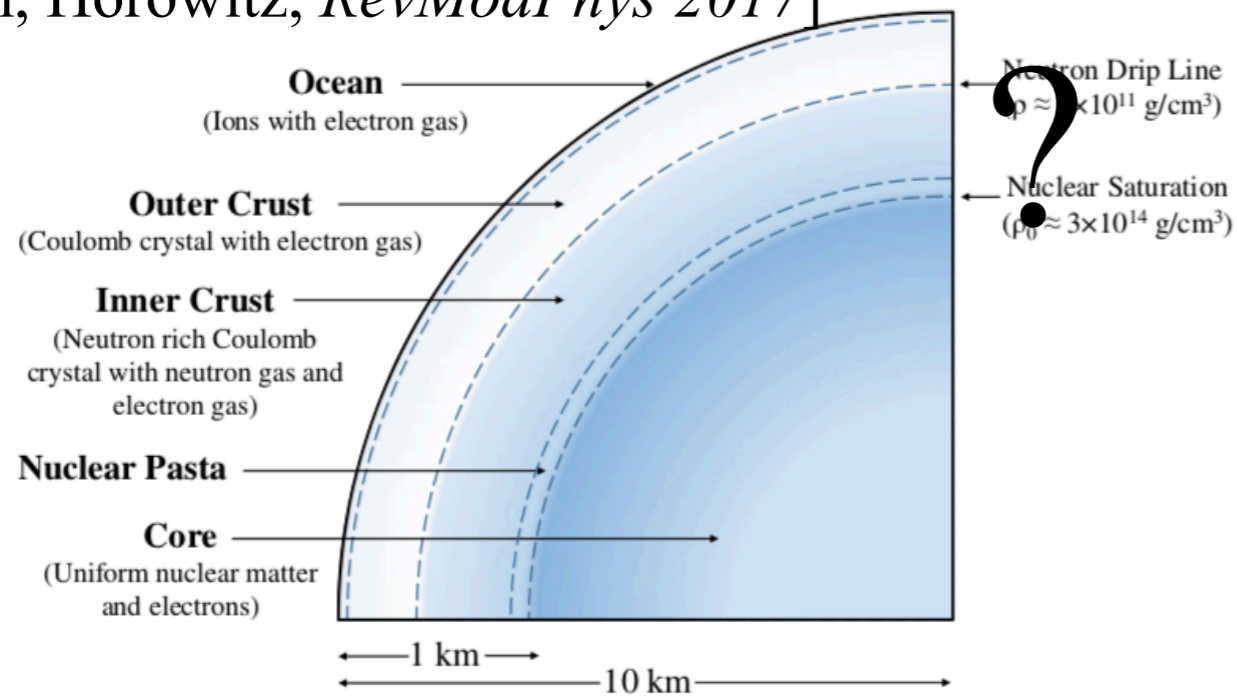


Possible bonus

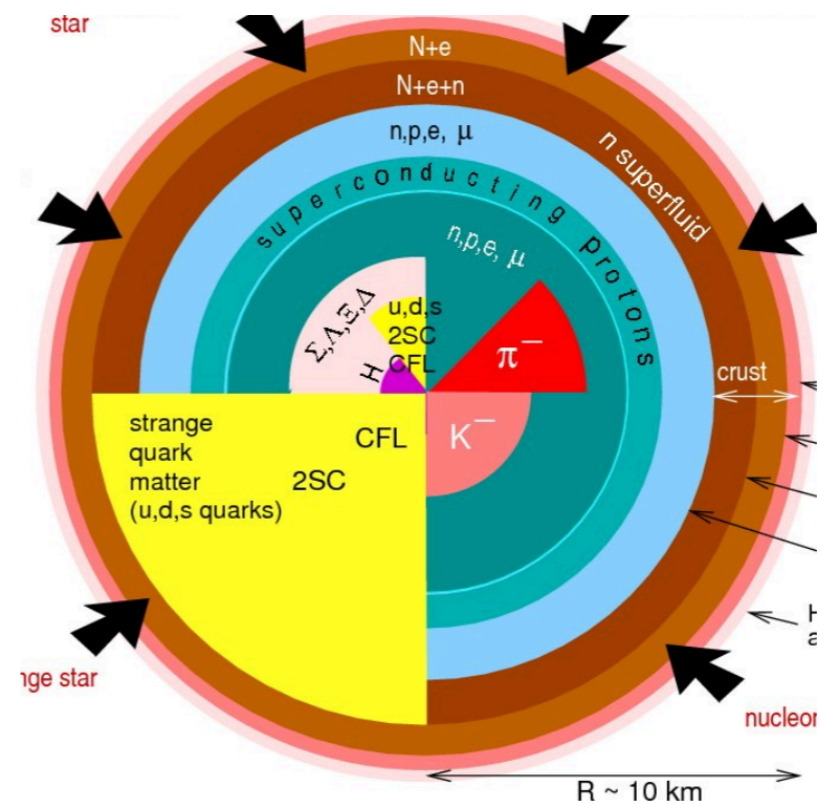
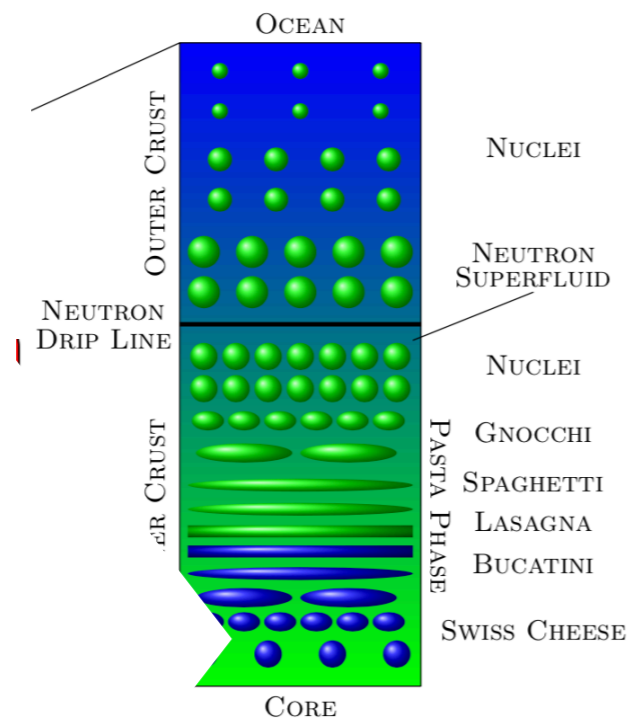


Heating nuclear pasta with DM (Nirmal Raj + friends)

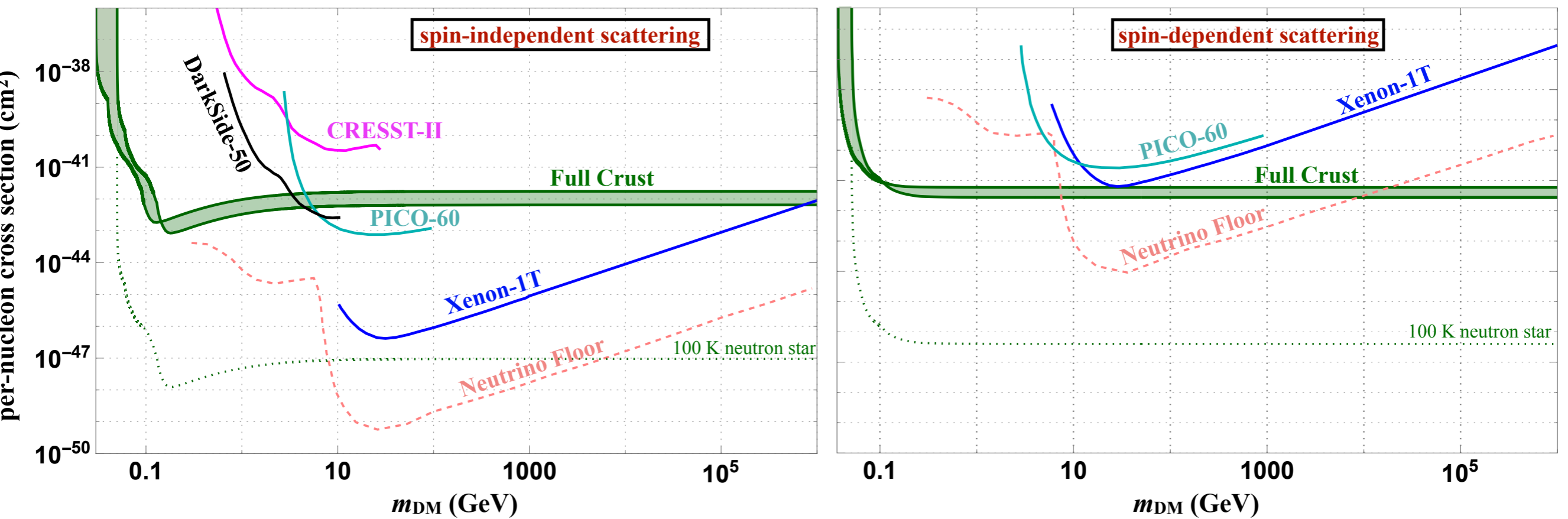
[Caplan, Horowitz, *RevModPhys* 2017]



deeper =>
knowledge of structure
more uncertain



If you see a cool neutron star: limit on DM heating



Conclusions

There may be a connection between the solar composition and asymmetric dark matter

Can we see DM in distant stars? What happens when they stop burning hydrogen? More work needed!

Asymmetric dark matter: currently looking at **all** data with GAMBIT

DM can prevent cooling in WDs and neutron stars...
and can ignite WDs and act as a mechanism for type Ia supernova detonation

