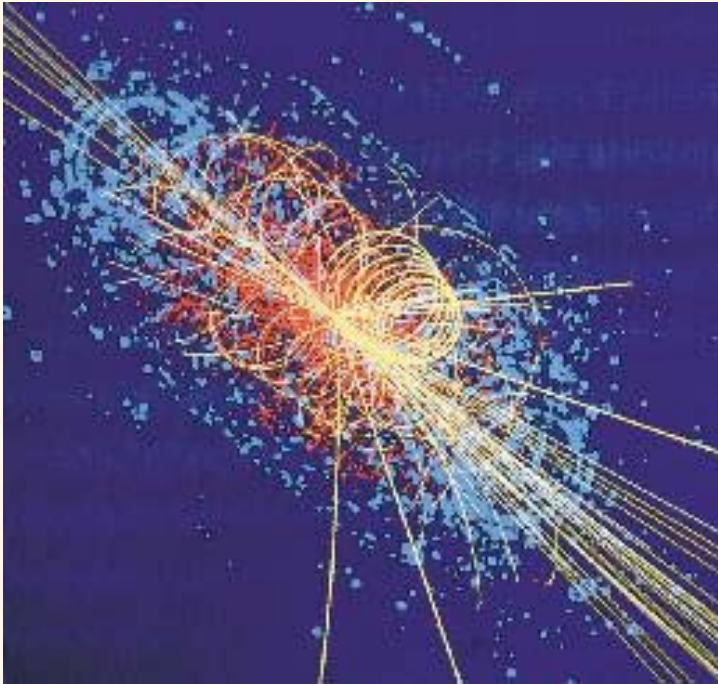


# Physics at Hadron Colliders

## Part 4



### Physics Beyond the Standard Model

- **Supersymmetry**  
(Tevatron and LHC)
- **Other Extensions of the Standard Model**
  - Extra dimensions
  - Extra gauge bosons
  - Leptoquarks ....

# Why ?

1. Gravity is not yet incorporated in the Standard Model
2. Dark Matter not accomodated
3. Many open questions in the Standard Model
  - Hierarchy problem:  $m_W$  (100 GeV)  $\rightarrow$   $m_{\text{Planck}}$  ( $10^{19}$  GeV)
  - Unification of couplings
  - Flavour / family problem
  - .....

All this calls for a **more fundamental theory** of which the Standard Model is a low energy approximation  $\rightarrow$  **New Physics**

Candidate theories: Supersymmetry  
Extra Dimensions  
Technicolor  
.....

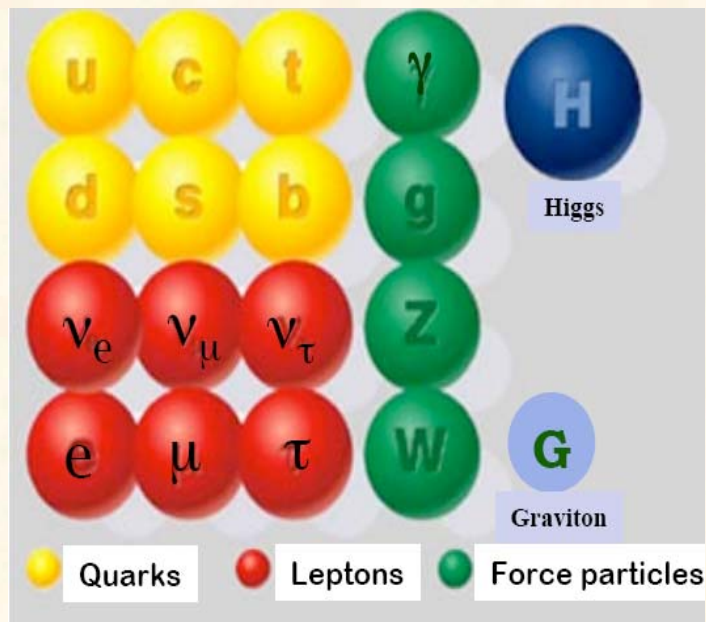
**Many extensions predict new physics at the TeV scale !!**

**Strong motivation for LHC, mass reach  $\sim 3$  TeV**

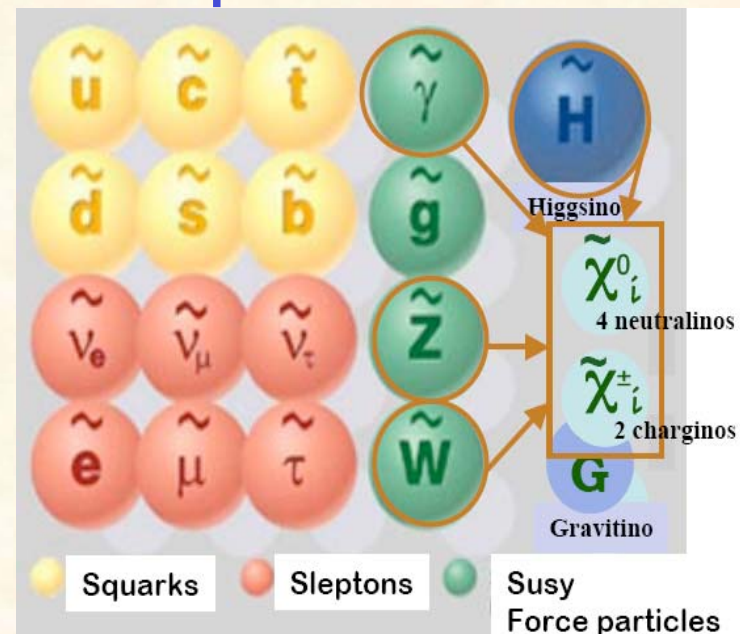
# Supersymmetry

Extends the Standard Model by predicting a new symmetry  
 Spin  $\frac{1}{2}$  matter particles (fermions)  $\Leftrightarrow$  Spin 1 force carriers (bosons)

**Standard Model particles**



**SUSY particles**



New Quantum number: R-parity:

$$R_p = (-1)^{B+L+2s} = \begin{array}{ll} +1 & \text{SM particles} \\ -1 & \text{SUSY particles} \end{array}$$

### Experimental consequences of R-parity conservation:

- SUSY particles are **produced in pairs**
- **Lightest Supersymmetric Particle (LSP)** is stable.

LSP is only **weakly interacting**:

**LSP**  $\equiv \chi^0_1$  (lightest neutralino, in many models)

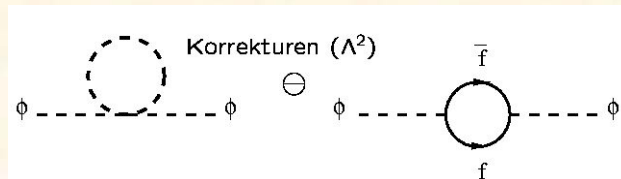
→ LSP behaves like a  $\nu$  → it escapes detection

→  **$E_T^{\text{miss}}$**  (typical SUSY signature)



# Why do we like SUSY so much?

1. Quadratically divergent quantum corrections to the Higgs boson mass are avoided

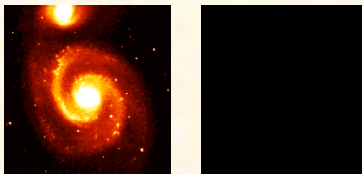


$$\Delta m_H = f(m_B^2 - m_f^2)$$

$$\rightarrow m_{\text{SUSY}} \sim 1 \text{ TeV}$$

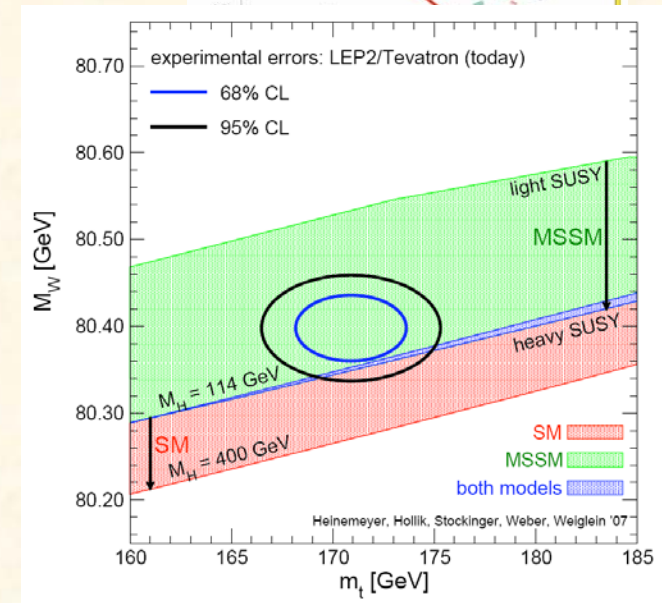
(Hierarchy or naturalness problem)

2. Unification of coupling constants of the three interactions seems possible
3. SUSY provides a candidate for dark matter



**The lightest SUSY particle (LSP)**

4. A SUSY extension is a small perturbation, consistent with the electroweak precision data



# Link to the Dark Matter in the Universe ?

Parameters of the SUSY model  $\Rightarrow$  predictions for the relic density of dark matter

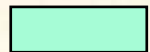
Interpretation in a simplified model

$$\rho_\chi \sim m_\chi n_\chi, \quad n_\chi \sim \frac{1}{\sigma_{ann}(\chi\chi \rightarrow \dots)}$$

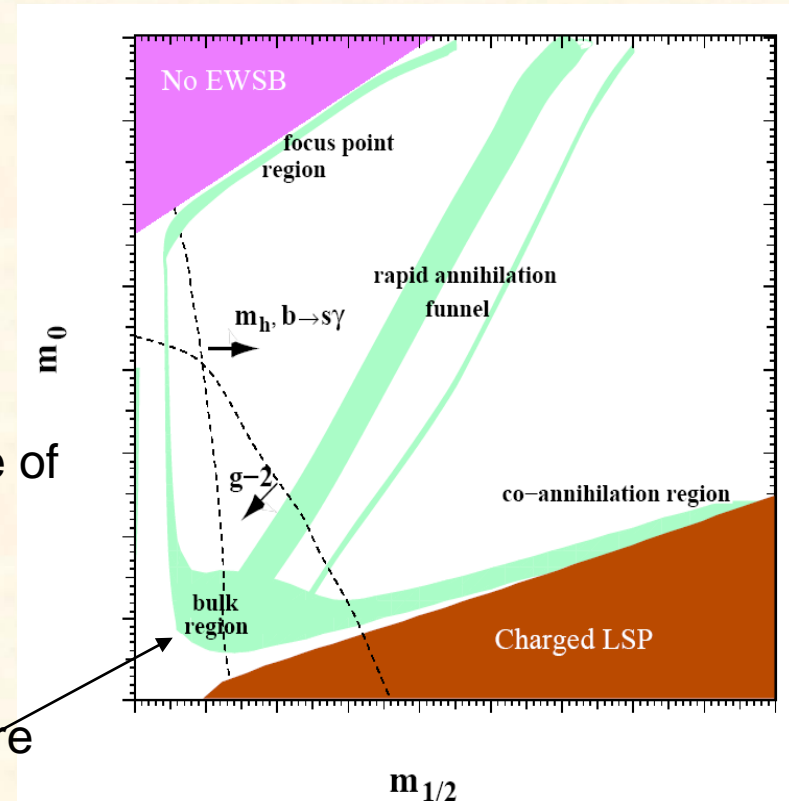
cMSSM  
(constrained Minimal Supersymmetric  
Standard Model)

Five parameters:

$m_0, m_{1/2}$  particle masses at the GUT scale  
 $A_0$  common coupling term  
 $\tan \beta$  ratio of vacuum expectation value of  
the two Higgs doublets  
 $\mu$  (sign  $\mu$ ) Higgs mass term



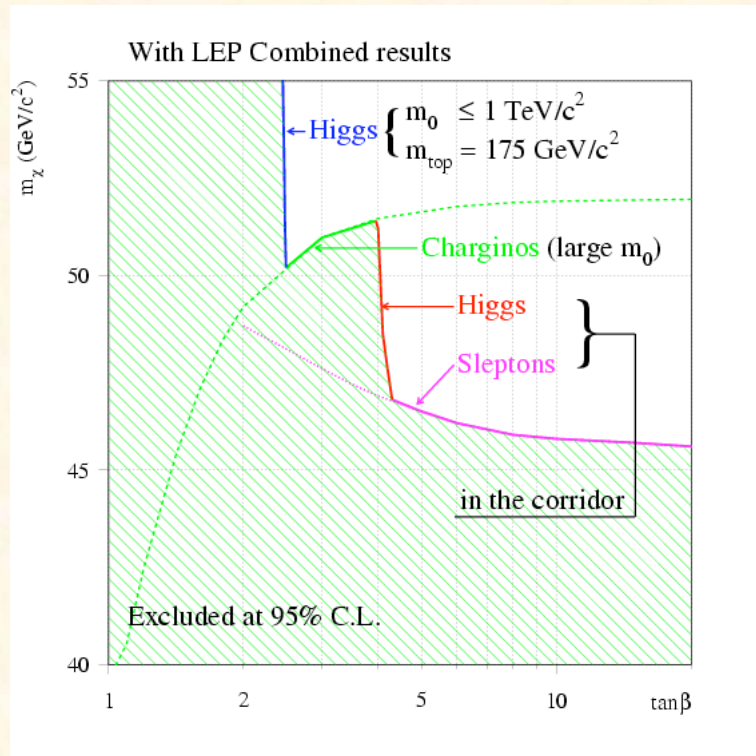
regions of parameter space which are  
consistent with the measured relic  
density of dark matter (WMAP,.....)



The **masses of the SUSY particles** are not predicted;  
 Theory has many additional new parameters (on which the masses depend)

However, charginos/neutralinos are usually lighter than squarks/sleptons/gluinos.

<u>Present mass limits</u> :	$m$ (sleptons, charginos)	$>$	90-103 GeV	LEP II
	$m$ (squarks, gluinos)	$>$	$\sim 350$ GeV	Tevatron
	$m$ (LSP, lightest neutralino)	$>$	$\sim 45$ GeV	LEP II



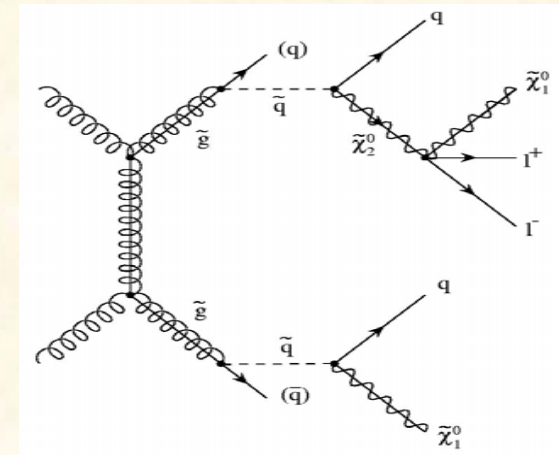
LEP-II limit on the mass of the  
 Lightest SUSY particle

assumption:  
 lightest neutralino = LSP

# Search for Supersymmetry at the LHC

- If **SUSY** exists at the electroweak scale, a discovery at the LHC should be easy
- **Squarks** and **Gluinos** are strongly produced

They decay through cascades to the lightest SUSY particle (LSP)



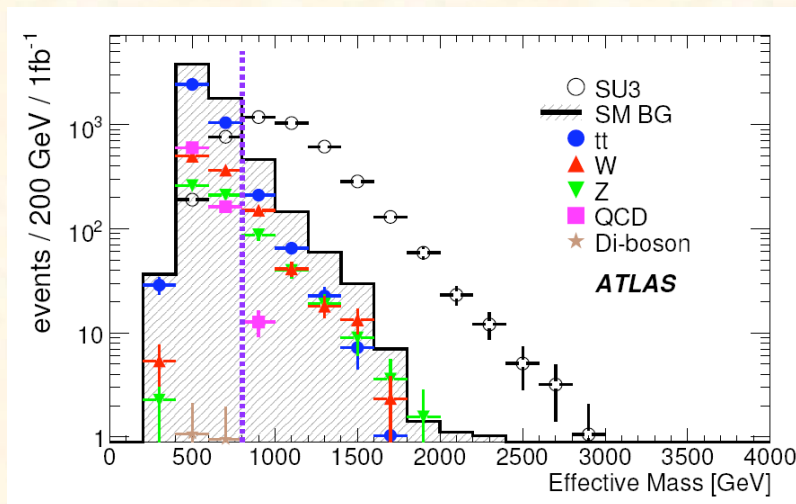
⇒ combination of  
**Jets, Leptons,  $E_T^{\text{miss}}$**

1. Step: Look for **deviations from the Standard Model**  
Example: Multijet +  **$E_T^{\text{miss}}$**  signature
2. Step: Establish the **SUSY mass scale** use inclusive variables, e.g. effective mass distribution
3. Step: Determine **model parameters** (difficult)  
Strategy: select particular decay chains and use kinematics to determine mass combinations



# Squarks and Gluinos

- If R-parity conserved, cascade decays produce distinctive events:  
multiple jets, leptons, and  $E_T^{\text{miss}}$
- Typical selection:  $N_{\text{jet}} > 4$ ,  $E_T > 100, 50, 50, 50 \text{ GeV}$ ,  $E_T^{\text{miss}} > 100 \text{ GeV}$
- Define:  $M_{\text{eff}} = E_T^{\text{miss}} + p_T^1 + p_T^2 + p_T^3 + p_T^4$  (effective mass)



LHC reach for Squark- and Gluino masses:

$0.1 \text{ fb}^{-1} \Rightarrow M \sim 750 \text{ GeV}$

$1 \text{ fb}^{-1} \Rightarrow M \sim 1350 \text{ GeV}$

$10 \text{ fb}^{-1} \Rightarrow M \sim 1800 \text{ GeV}$

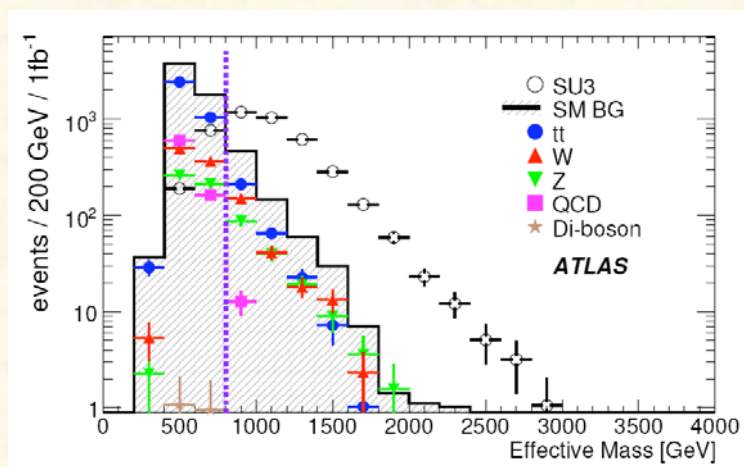
Deviations from the Standard Model  
due to SUSY at the TeV scale can be  
detected fast !

example: mSUGRA, point SU3 (bulk region)

$m_0 = 100 \text{ GeV}$ ,  $m_{1/2} = 300 \text{ GeV}$

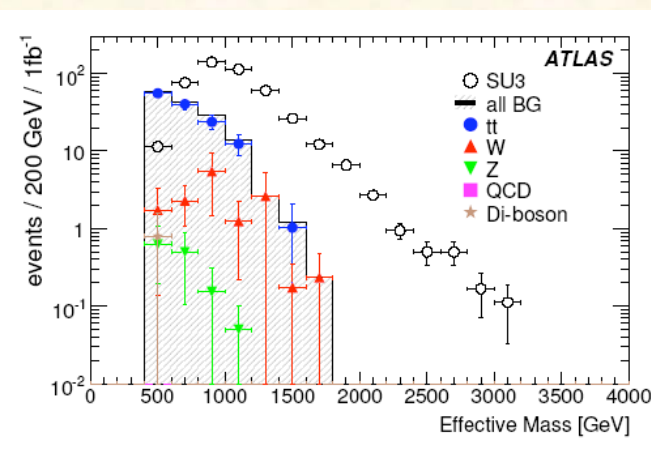
$\tan \beta = 6$ ,  $A_0 = -300 \text{ GeV}$ ,  $\mu > 0$

## ...additional potential: inclusive searches with leptons

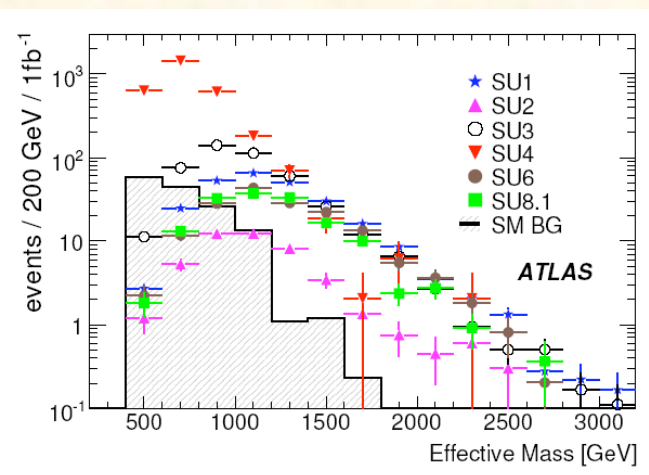


SU3, 4 jets + 0 lepton final states

- Smaller signal rates, but better S:B conditions
- Discovery potential is more robust, in particular at the beginning, when systematic uncertainties on the backgrounds are large
- Similar analyses with  $\tau$  lepton and b quark final states

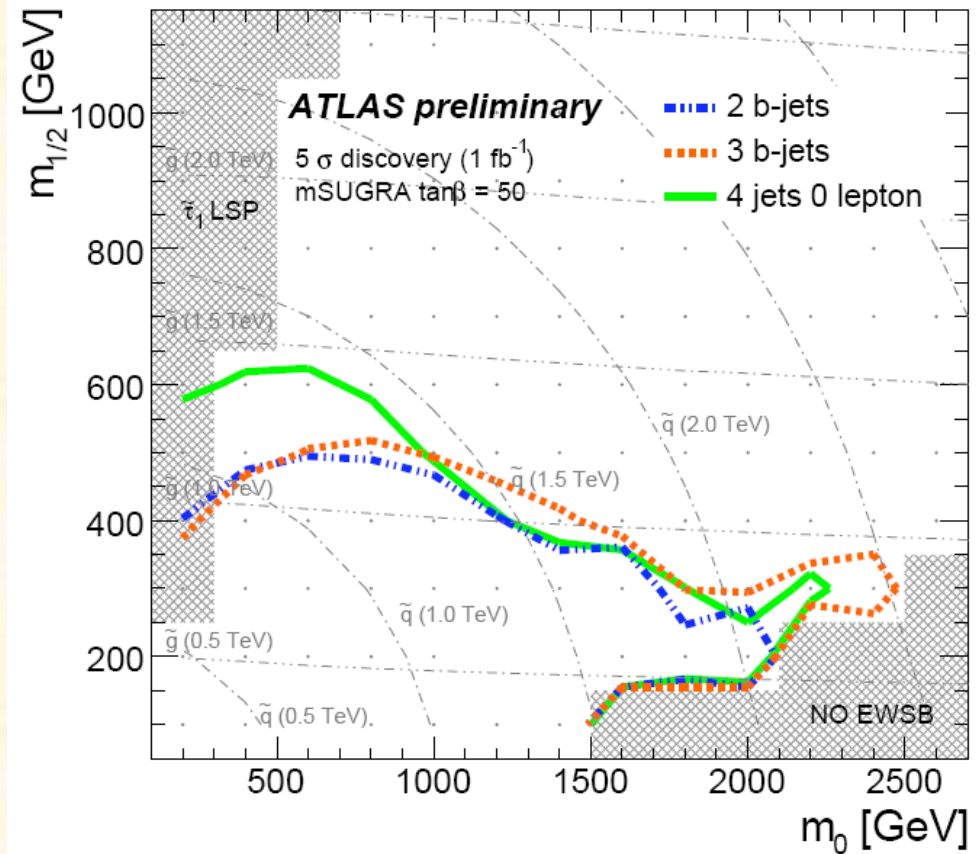
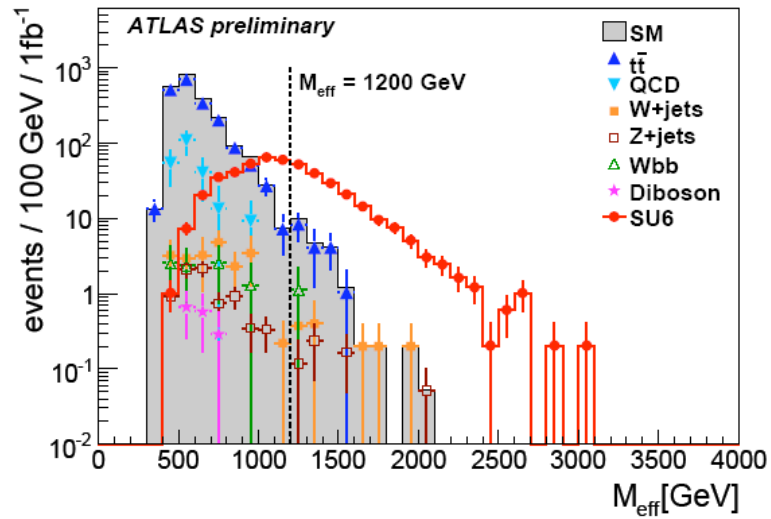


SU3, 4 jets + 1 lepton final states



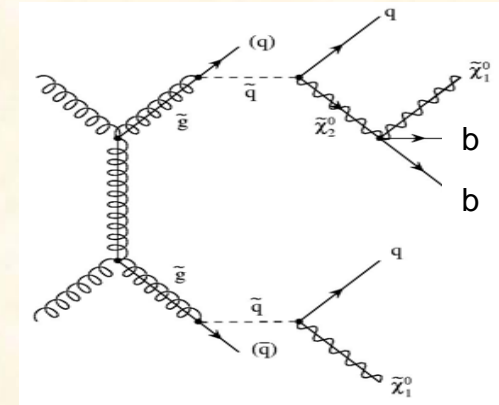
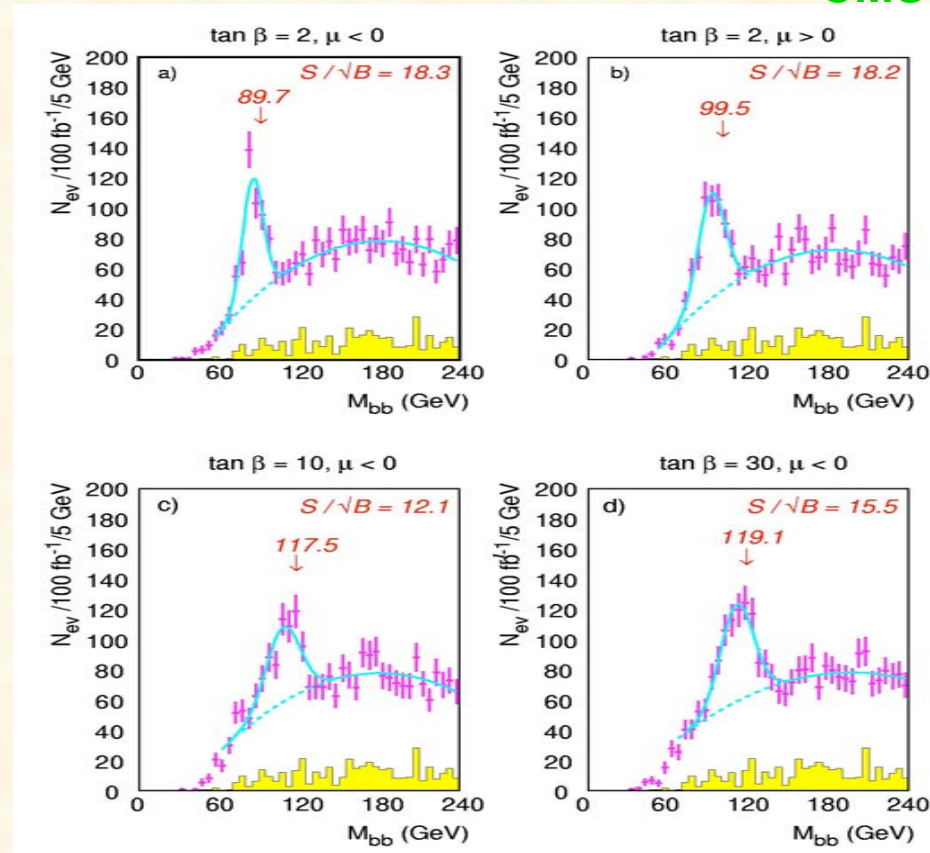
4 jets + 1 lepton final states for other benchmark points

## Example: search for final states with $E_T^{\text{miss}}$ and two b-jets



$h \rightarrow bb$ :

CMS



important if  $\chi^0_2 \rightarrow \chi^0_1 h$  is open;  
bb peak can be reconstructed in  
many cases

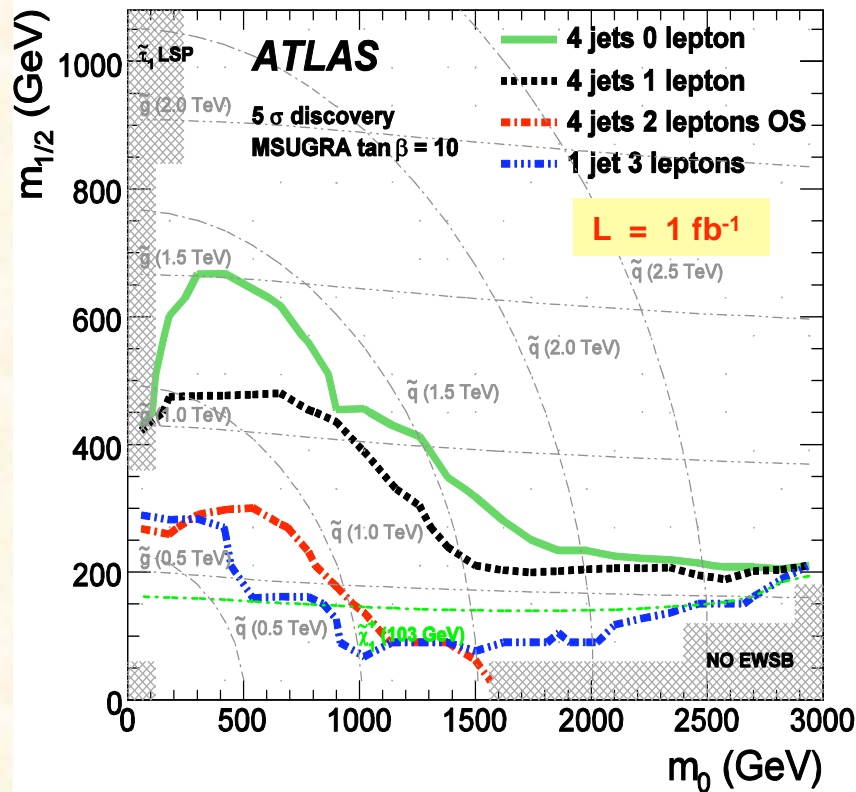
**Could be a Higgs discovery mode !**

**SM background can be reduced  
by applying a cut on  $E_T^{\text{miss}}$**

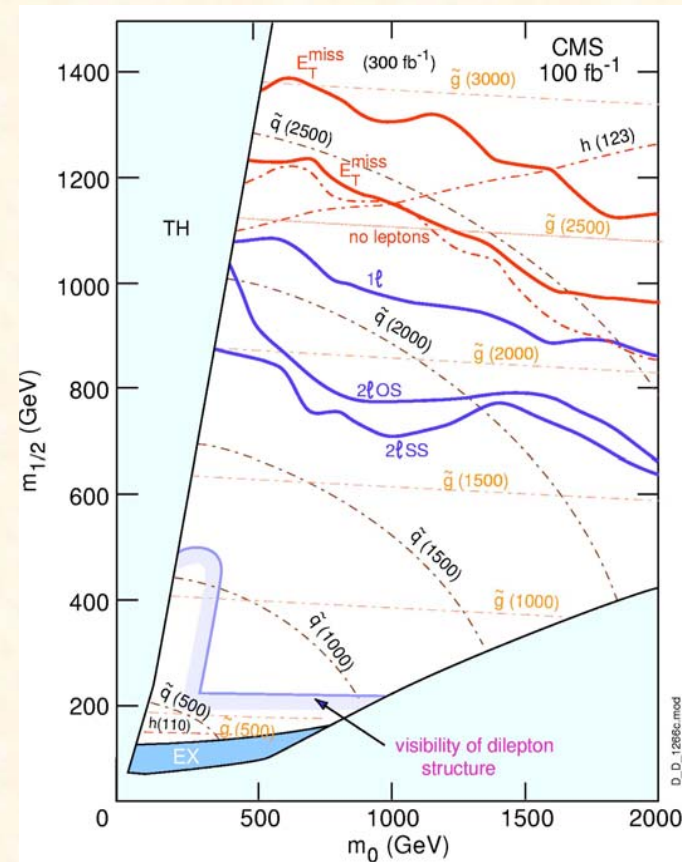


## LHC reach in the $m_0 - m_{1/2}$ mSUGRA plane:

Multijet +  $E_T^{\text{miss}}$  signature



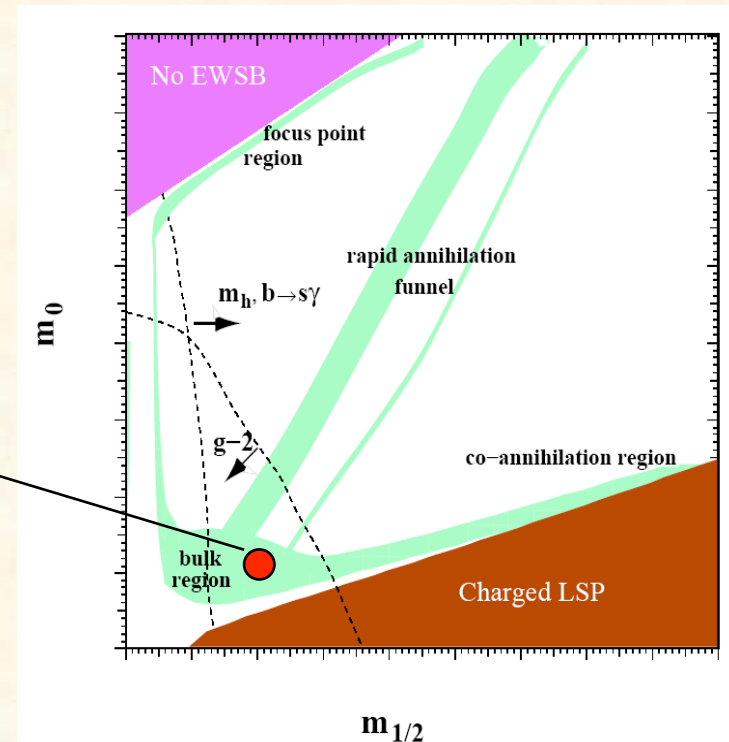
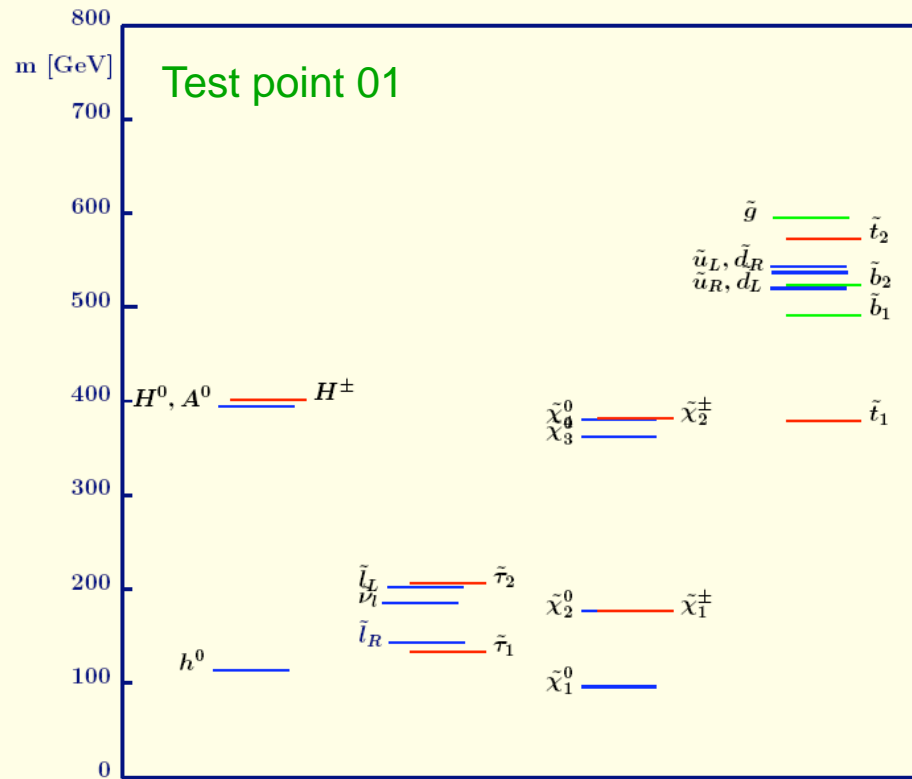
SUSY cascade decays give also rise to many other inclusive signatures: **leptons**, **b-jets**,  **$\tau$ 's**



- Tevatron reach can be extended with early data
  - Expect multiple signatures for TeV-scale SUSY
- Long term mass reach (300  $\text{fb}^{-1}$ ): 2.5 – 3 TeV

# How can the underlying theoretical model be identified ?

Measurement of the SUSY spectrum  $\rightarrow$  Parameter of the theory

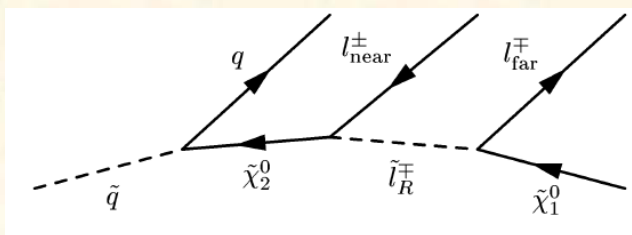


LHC: strongly interacting squarks and gluinos

ILC / CLIC: precise investigation of electroweak SUSY partners

# LHC Strategy: End point spectra of cascade decays

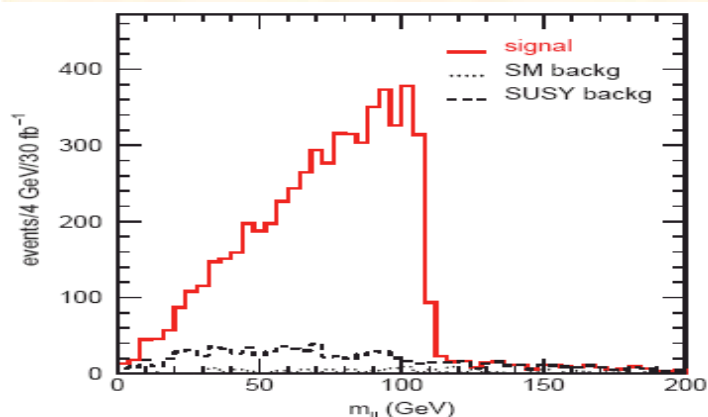
Example:  $\tilde{q} \rightarrow q\tilde{\chi}_2^0 \rightarrow q\tilde{\ell}^\pm \ell^\mp \rightarrow q\ell^\pm \ell^\mp \tilde{\chi}_1^0$



$$M_{\ell^+\ell^-}^{\max} = \frac{\sqrt{(m_{\tilde{\chi}_2^0}^2 - m_{\tilde{\ell}}^2)(m_{\tilde{\ell}}^2 - m_{\tilde{\chi}_1^0}^2)}}{m_{\tilde{\ell}}}$$

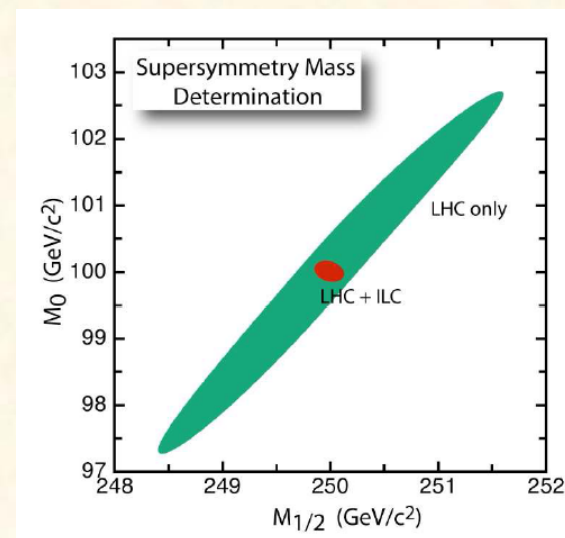
$$M_{\ell_1 q}^{\max} = \frac{\sqrt{(m_{\tilde{\chi}_2^0}^2 - m_{\tilde{\ell}}^2)(m_{\tilde{q}}^2 - m_{\tilde{\chi}_2^0}^2)}}{m_{\tilde{\chi}_2^0}}$$

Results for point 01:

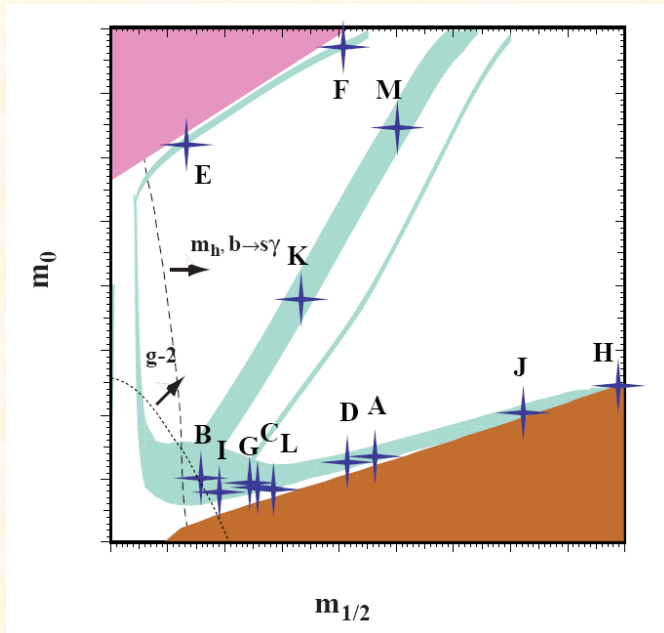


	LHC	LHC+ILC
$\Delta m_{\tilde{\chi}_1^0}$	4.8	0.05 (input)
$\Delta m_{\tilde{l}_R}$	4.8	0.05 (input)
$\Delta m_{\tilde{\chi}_2^0}$	4.7	0.08
$\Delta m_{\tilde{q}_L}$	8.7	4.9
$\Delta m_{\tilde{q}_R}$	11.8	10.9
$\Delta m_{\tilde{g}}$	8.0	6.4
$\Delta m_{\tilde{b}_1}$	7.5	5.7
$\Delta m_{\tilde{b}_2}$	7.9	6.2
$\Delta m_{\tilde{l}_L}$	5.0	0.2 (input)
$\Delta m_{\tilde{\chi}_4^0}$	5.1	2.23

$L = 300 \text{ fb}^{-1}$

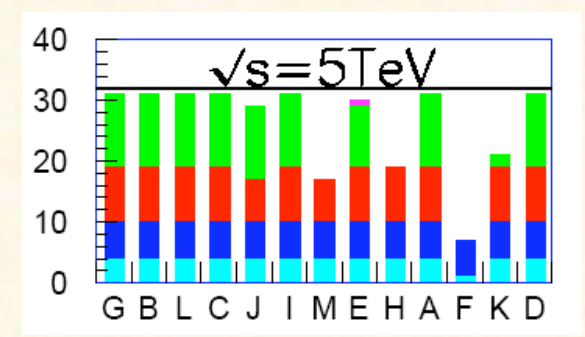
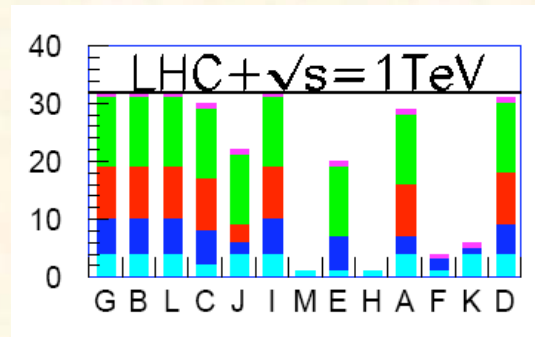
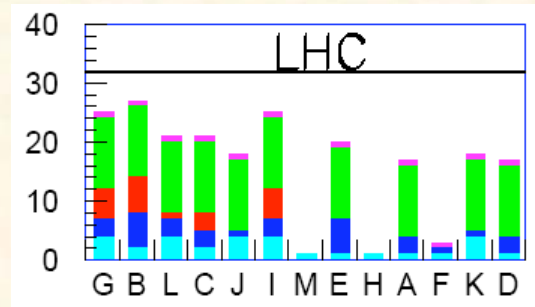


# The LHC and the ILC (International Linear Collider, in study/planning phase) are complementary in SUSY searches



■ gluino 
 ■ squarks 
 ■ sleptons 
 ■  $\chi^{0,\pm}$ 
■ H

Number of observable SUSY particles:



)\* Study by J. Ellis et al., hep-ph/0202110

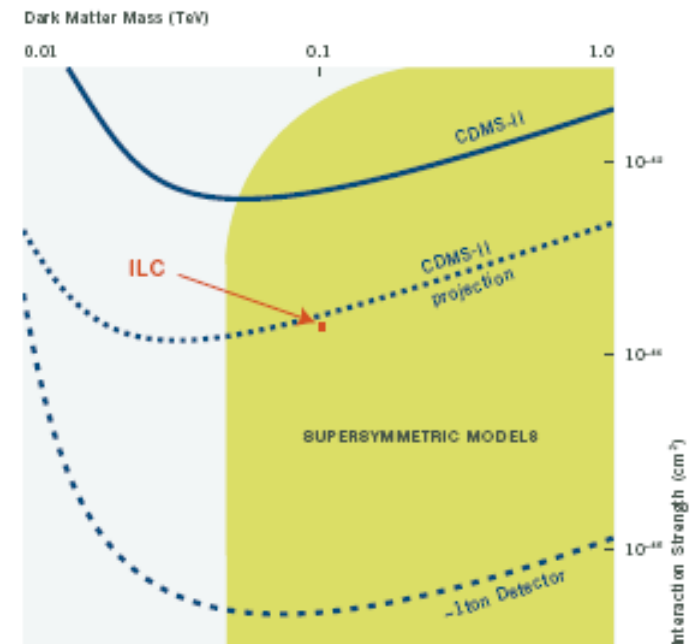
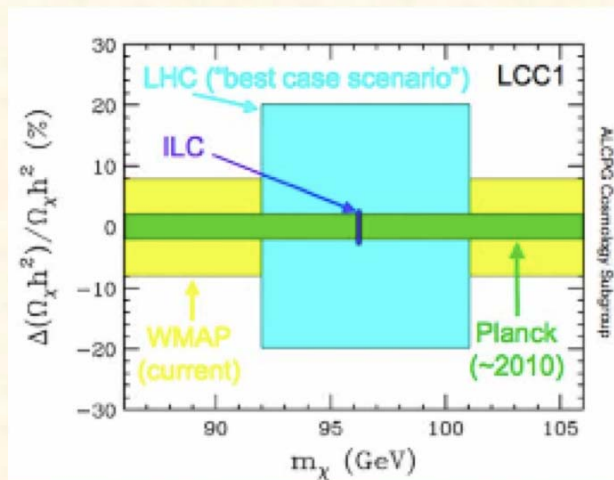


## Strategy in SUSY Searches at the LHC:

- Search for multijet +  $E_T^{\text{miss}}$  excess
- If found, select SUSY sample (simple cuts)
- Look for special features ( $\gamma$ 's , long lived sleptons)
- Look for  $\ell^\pm$ ,  $\ell^+ \ell^-$ ,  $\ell^\pm \ell^\pm$ , b-jets,  $\tau$ 's
- End point analyses, global fit  $\rightarrow$  SUSY model parameters

## Importance for the interplay between direct and indirect Dark Matter searches

- Following a discovery of New Physics at the LHC (deviation from the Standard Model) the LHC will aim to test the Dark Matter hypothesis
- Estimation of relic density in a simple model-dependent scenario will be the first goal
- Less model-dependent scenarios will follow, detailed studies probably require the ILC
- Conclusive result is only possible in conjunction with astroparticle physics experiments
- Ultimate goal: observation of LSP at the LHC, confirmed by a signal in a direct dark matter experiment with predicted mass and cross-section



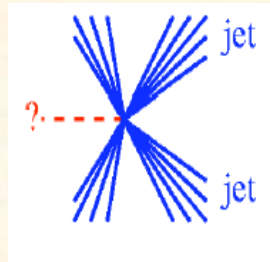
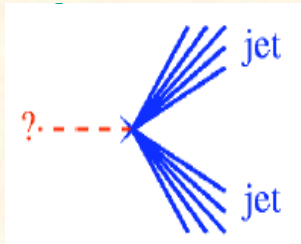
The Search for



**SUSY at the Tevatron**

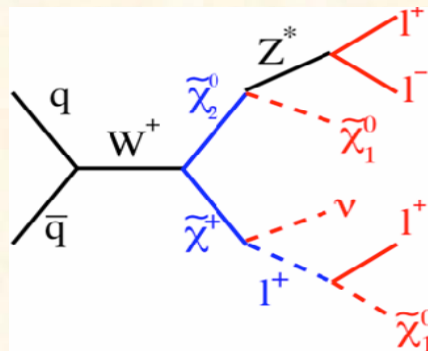
## The two classical signatures

1. Search for Squarks and Gluinos: **Jet +  $E_T^{\text{miss}}$**  signature  
produced via QCD processes



2. Search for Charginos and Neutralinos: **Multilepton +  $E_T^{\text{miss}}$**  signature  
produced via electroweak processes (associated production)

$$\tilde{\chi}_2^0 \tilde{\chi}_1^\pm \rightarrow l^\pm l^\mp l^\pm \tilde{\chi}_1^0 \tilde{\chi}_1^0 X$$







## Search for Squarks and Gluinos



- Three different analyses, depending on squark / gluinos mass relations:

(i) dijet analysis

small  $m_0$ ,  $m(\text{squark}) < m(\text{gluino})$

$$\tilde{q} \bar{\tilde{q}} \rightarrow q \tilde{\chi}_1^0 \bar{q} \tilde{\chi}_1^0$$

(ii) 3-jet analysis

intermediate  $m_0$   $m(\text{squark}) \approx m(\text{gluino})$

$$\tilde{q} \tilde{g} \rightarrow q \tilde{\chi}_1^0 \bar{q} \tilde{q} \tilde{\chi}_1^0$$

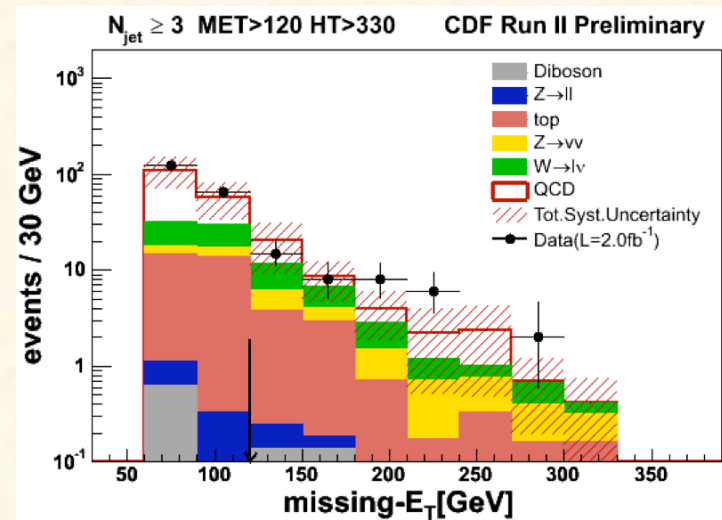
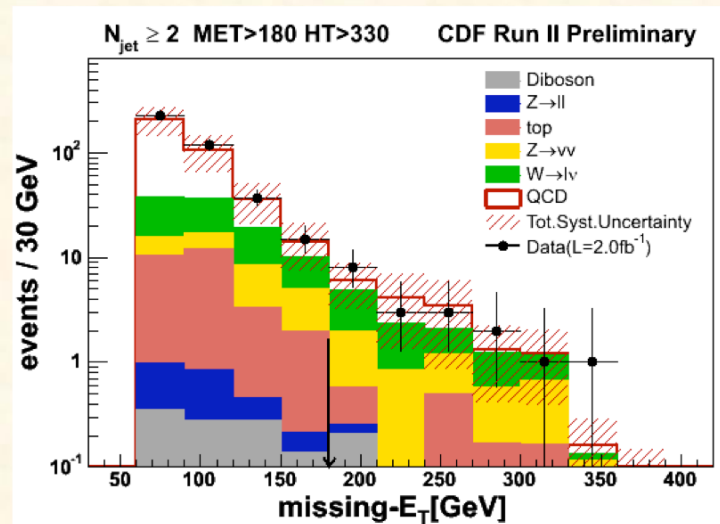
(iii) Gluino analysis

large  $m_0$ ,  $m(\text{squark}) > m(\text{gluino})$

$$\tilde{g} \tilde{g} \rightarrow q \bar{q} \tilde{\chi}_1^0 q \bar{q} \tilde{\chi}_1^0$$

- **Main backgrounds:**  $Z \rightarrow \nu\nu + \text{jets}$ ,  $t\bar{t}$ ,  $W + \text{jet production}$
- **Event selection:**
  - \* require at least 2, 3 or 4 jets with  $P_T > 60 / 40 / 30 / 20$  GeV
  - \* veto on isolated electrons and muons
  - \* isolation of  $E_T^{\text{miss}}$  and all jets
  - \* optimization of the final cuts  $\rightarrow$  discriminating variables

# Search for Squarks and Gluinos (cont.)



Expected background:

samples	2-jets	3-jets	4-jets
QCD	$4.37 \pm 2.01$	$13.34 \pm 4.67$	$15.26 \pm 7.60$
top	$1.35 \pm 1.22$	$7.56 \pm 3.85$	$22.14 \pm 7.29$
$Z \rightarrow \nu\nu + \text{jets}$	$3.95 \pm 1.09$	$5.39 \pm 1.74$	$2.74 \pm 0.95$
$Z \rightarrow ll + \text{jets}$	$0.09 \pm 0.04$	$0.16 \pm 0.11$	$0.14 \pm 0.08$
$W \rightarrow l\nu + \text{jets}$	$6.08 \pm 2.15$	$10.69 \pm 3.84$	$7.68 \pm 2.85$
WW/WZ/ZZ	$0.21 \pm 0.19$	$0.35 \pm 0.17$	$0.49 \pm 0.34$
tot SM	$16 \pm 5$	$37 \pm 12$	$48 \pm 17$

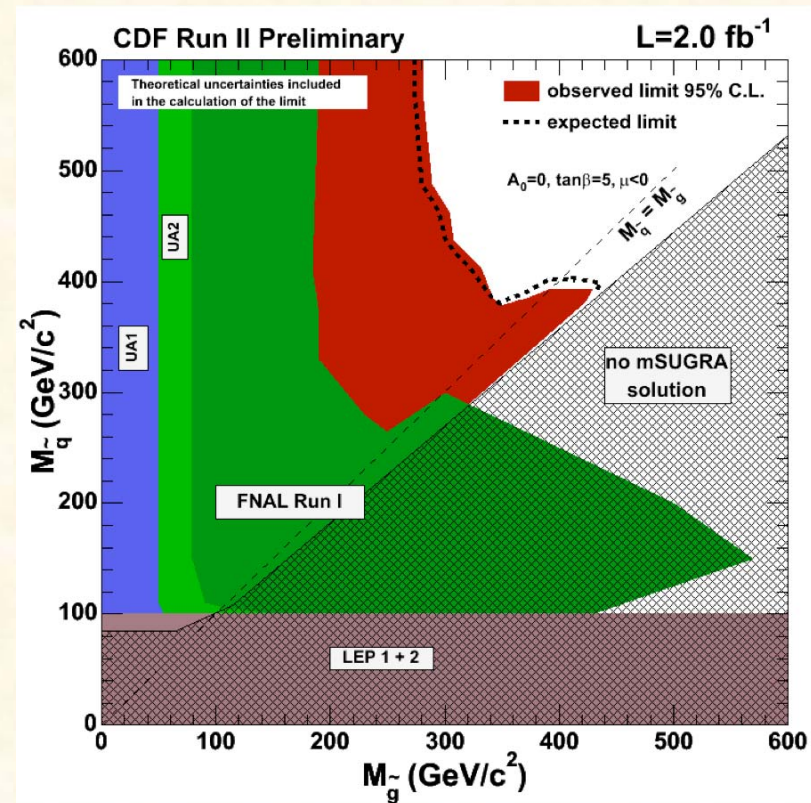
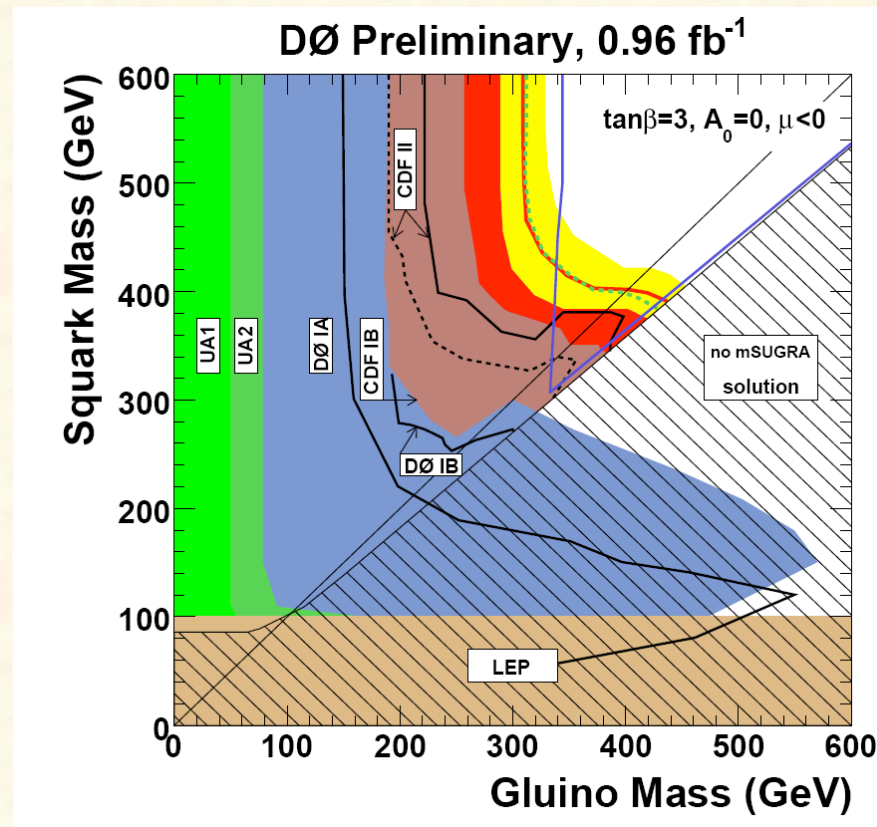
Observed events in data:

Region	Observed data
4-jets	45
3-jets	38
2-jets	18

No excess above background from Standard Model processes

→ No evidence for SUSY (yet) → Set limits on masses of SUSY particles

# Excluded regions in the $m(\text{squark})$ vs. $m(\text{gluino})$ plane



## Exclusion limits

(incl. systematic uncertainties)\*:

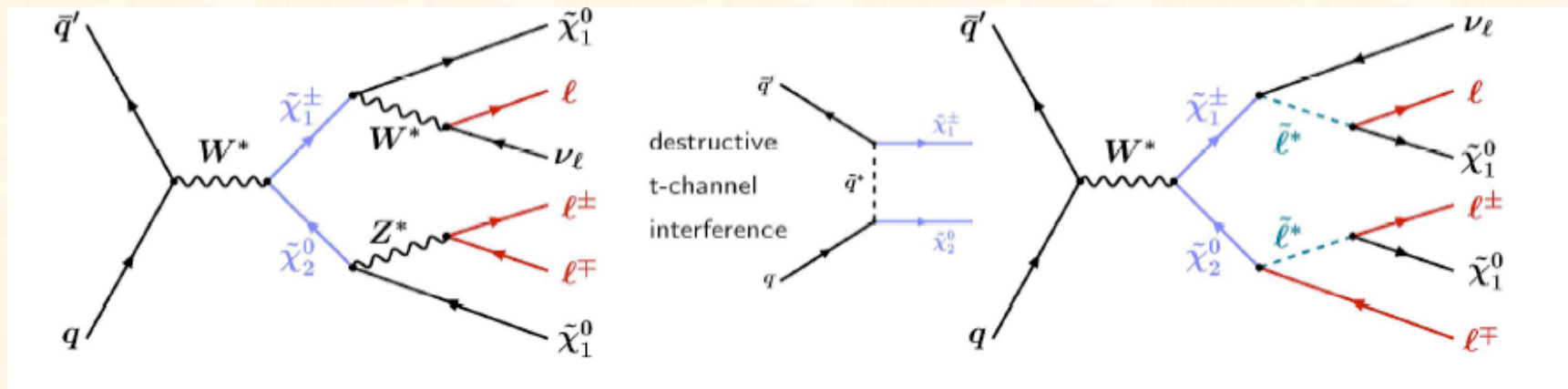
$$m(\text{gluino}) > 290 \text{ GeV}/c^2$$

$$m(\text{squark}) > 375 \text{ GeV}/c^2$$

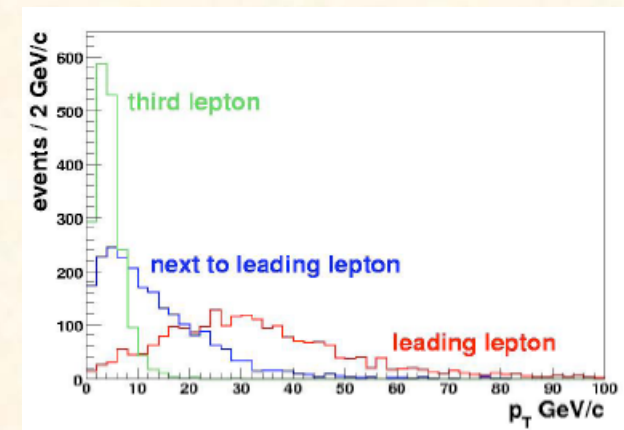
)\* uncertainties from structure functions, change of renormalization and factorization scale  $\mu$  by a factor of 2, NLO calculation, default choice:  $\mu = m(\text{gluino})$ ,  $m(\text{squark})$  or  $\frac{1}{2}(m(\text{gluino})+m(\text{squark}))$  for  $gg$ ,  $qq$ ,  $qg$  production

# Search for Charginos and Neutralinos - the tri-lepton channel-

- Gaugino pair production via electroweak processes  
(small cross sections,  $\sim 0.1 - 0.5$  pb, however, small expected background)



- For small gaugino masses ( $\sim 100$  GeV/ $c^2$ ) one needs to be sensitive to low  $P_T$  leptons

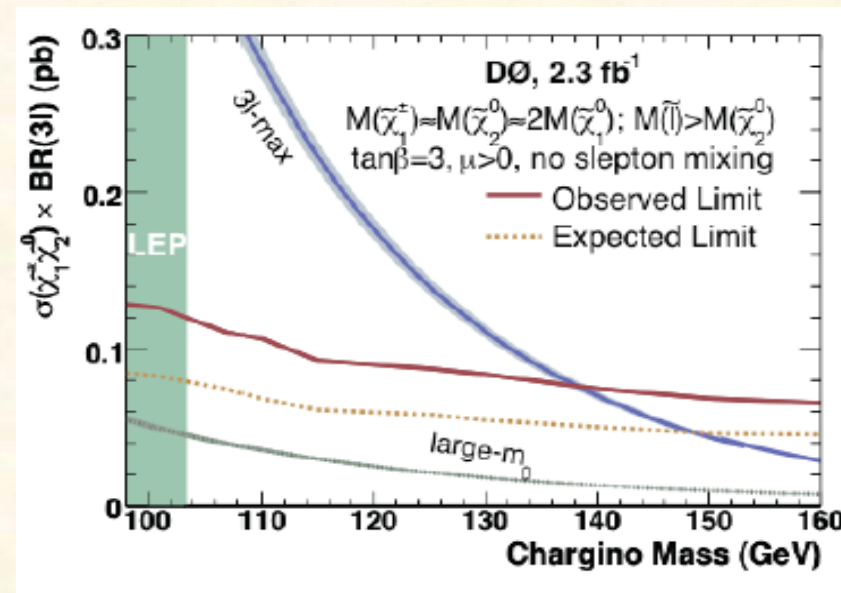
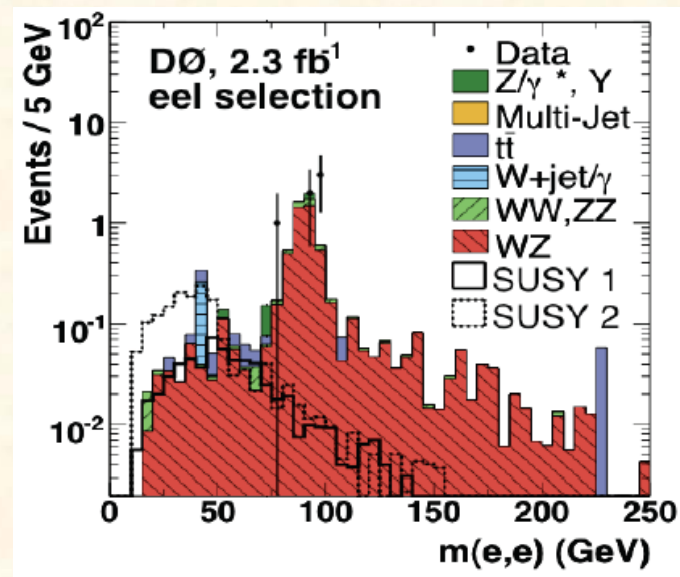




## Analysis:

- Search for different  $(\ell\ell\ell)$  + like-sign  $\mu\mu$  final states with missing transverse momentum
- In order to gain efficiency, no lepton identification is required for the 3<sup>rd</sup> lepton, select: two identified leptons + a track with  $p_T > 4$  GeV/c

mSUGRA interpretation



For specific scenarios: sensitivity / limits above LEP limits;  
e.g.,  $M(\chi^\pm) > 140$  GeV/c<sup>2</sup> for the 3l-max scenario

## Can LHC probe extra dimensions ?

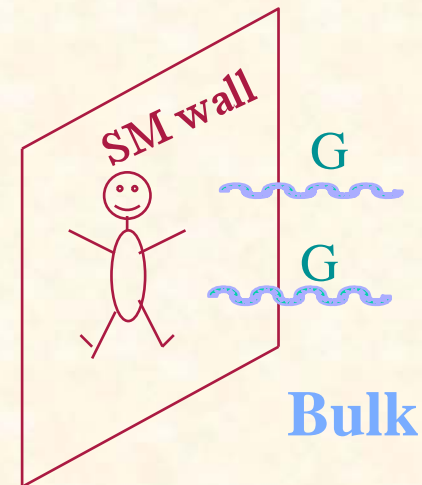
- Much recent theoretical interest in models with extra dimensions  
(Explain the weakness of gravity (or hierarchy problem) by extra dimensions)
- New physics can appear at the TeV-mass scale,  
i.e. accessible at the LHC

Example: Search for direct Graviton production

$$gg \rightarrow gG, qg \rightarrow qG, q\bar{q} \rightarrow Gg$$

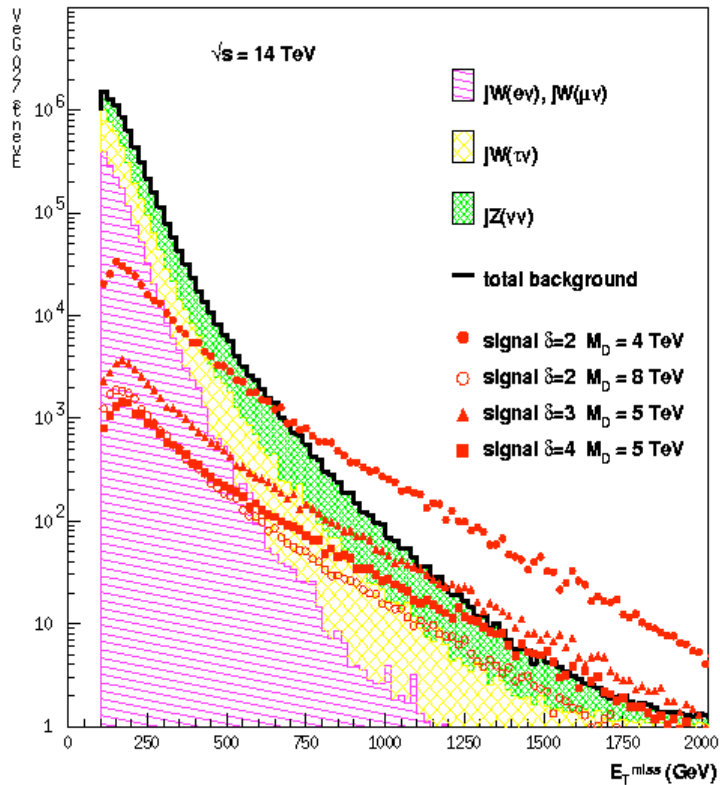
$$q\bar{q} \rightarrow G\gamma$$

$\Rightarrow$  Jets or Photons with  $E_T^{\text{miss}}$



# Search for escaping gravitons:

Jet +  $E_T^{\text{miss}}$  search:



Main backgrounds:

jet+Z( $\rightarrow \nu\nu$ ), jet+W $\rightarrow$ jet+(e,  $\mu$ ,  $\tau$ ) $\nu$

$$G_N^{-1} = 8\pi R^\delta M_D^{2+\delta}$$

$\delta$  : # extra dimensions

$M_D$  = scale of gravitation

$R$  = radius (extension)

$M_D^{\text{max}}$	=	9.1,	7.0,	6.0 TeV
	for			
$\delta$	=	2,	3,	4

LHC experiments are sensitive, but conclusions on the underlying theory are difficult and require a detailed measurement program

## More ideas?

### 1. New resonances decaying into lepton pairs

examples:  $W'$  and  $Z'$  or Graviton resonances (extra dimensions)

use again leptonic decay mode to search for them:  $W' \rightarrow \ell \nu$   
 $Z' \rightarrow \ell \ell$

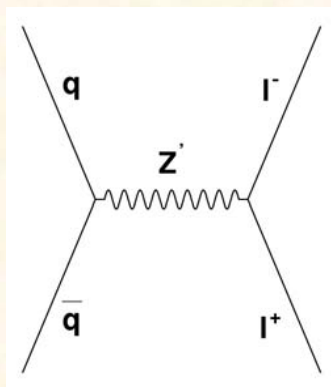
### 2. Leptoquarks ?

Particles that decay into leptons and quarks  
(violate lepton and baryon number; appear in Grand Unified theories)

here: search for low mass Leptoquarks (TeV scale)

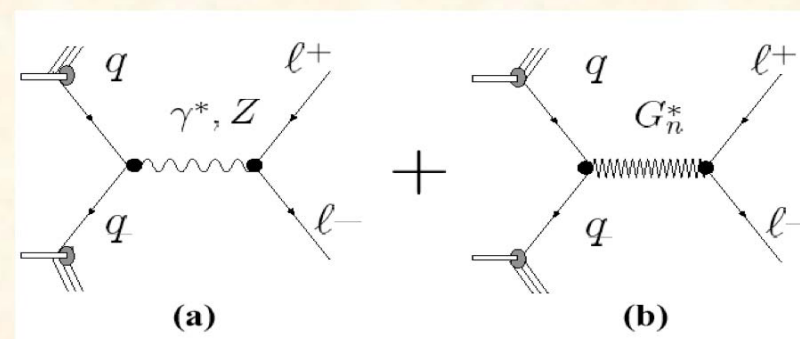
## Fermilab Search for New Resonances in High Mass Di-leptons

- **Neutral Gauge Boson  $Z'$**   
assume SM-like couplings



- **Randall-Sundrum narrow Graviton resonances decaying to di-lepton**

appear in Extra Dim. Scenarios



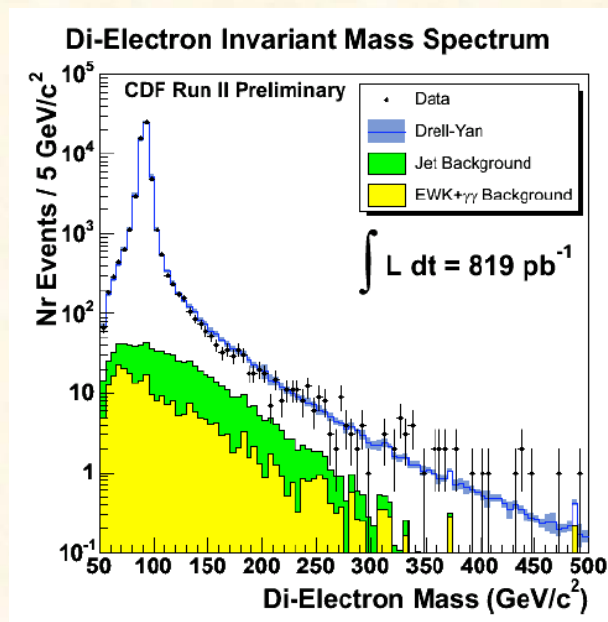
Main background from Drell-Yan pairs



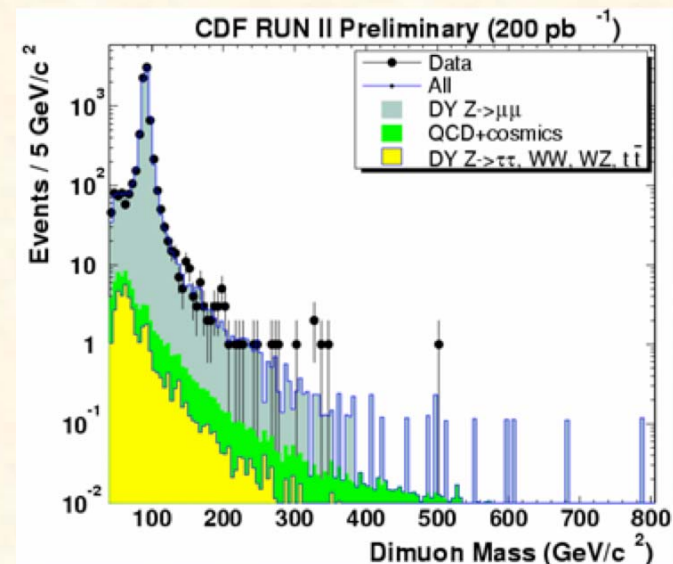
# Search for New Resonances in High Mass Di-leptons



## Di-electron Invariant Mass



## Di-muon Invariant Mass



Data are consistent with background from SM processes. No excess observed.

**Z' mass limits (SM couplings)**

**95% C.L.**

**CDF /D0:**

**ee**

**965**

**$\mu\mu$**

**835**

**$\tau\tau$**

**394**

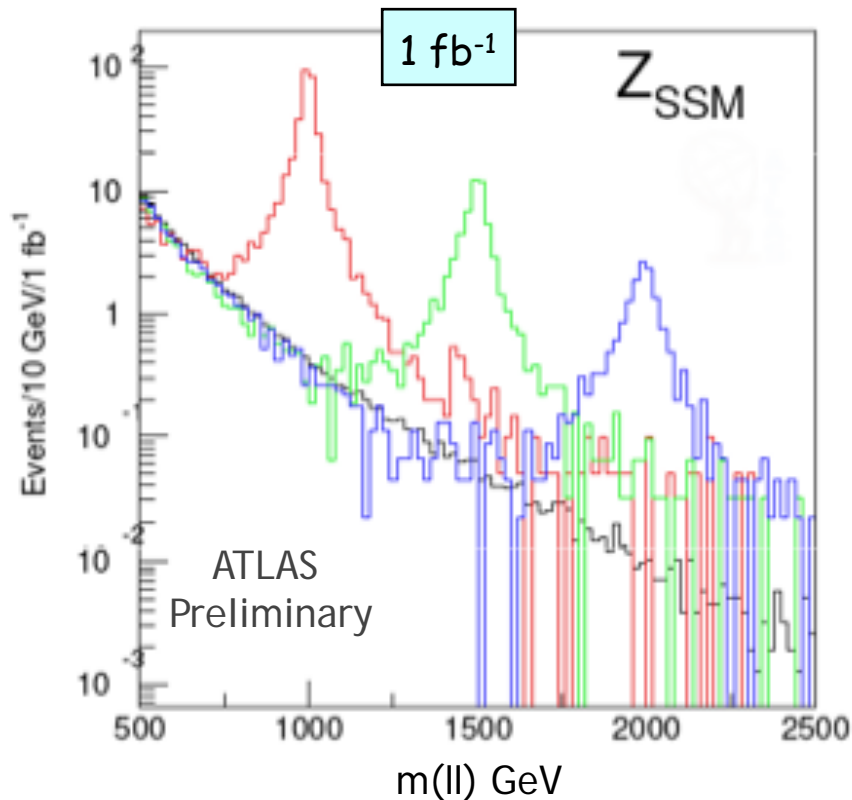
**GeV/c<sup>2</sup>**

# Early Surprises at the LHC??

- as already mentioned, the experiments must be open for surprises /  
unknowns /  
unexpected discoveries
- requires unbiased measurements of
  - inclusive lepton spectra
  - dileptons spectra.....
  - $E_T^{\text{miss}}$  spectrum.....
  - .....

## One example of many....

$Z' \rightarrow e^+e^-$  with SM-like couplings ( $Z_{\text{SSM}}$ )



$\sqrt{s} = 14 \text{ TeV}$

Mass (TeV)	Events / fb <sup>-1</sup> (after cuts)	Luminosity needed for a 5 $\sigma$ discovery + (10 obs. events)
1	~160	~70 pb <sup>-1</sup>
1.5	~30	~300 pb <sup>-1</sup>
2	~7	~1.5 fb <sup>-1</sup>

Discovery window above Tevatron limits  
 $m \sim 1 \text{ TeV}$ , perhaps even in 2011... (?)

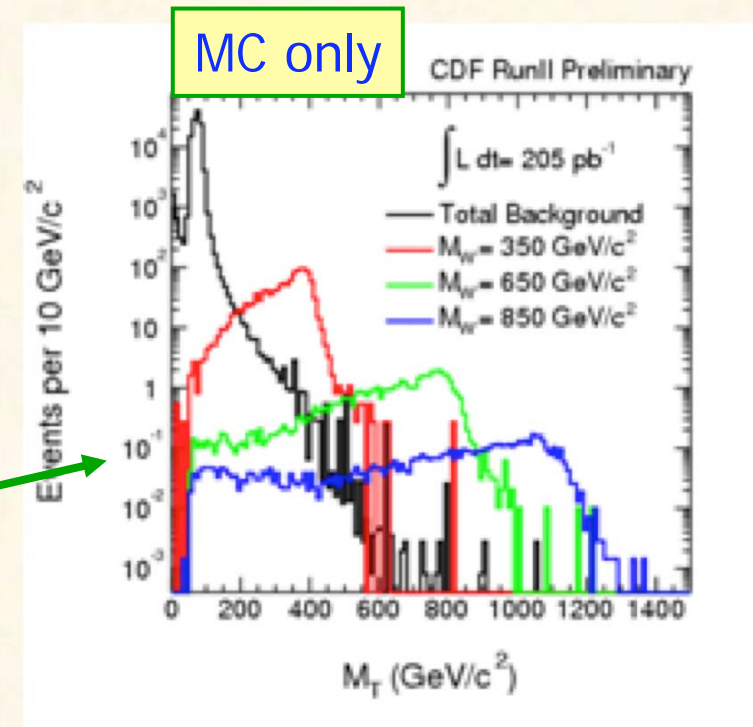


## Search for $W' \rightarrow e\nu$

- $W'$ : additional charged heavy vector boson
- appears in theories based on the extension of the gauge group
- e.g. Left-right symmetric models:  
 $SU(2)_R \quad W_R$
- assume: the neutrino from  $W'$  decay is light and stable.

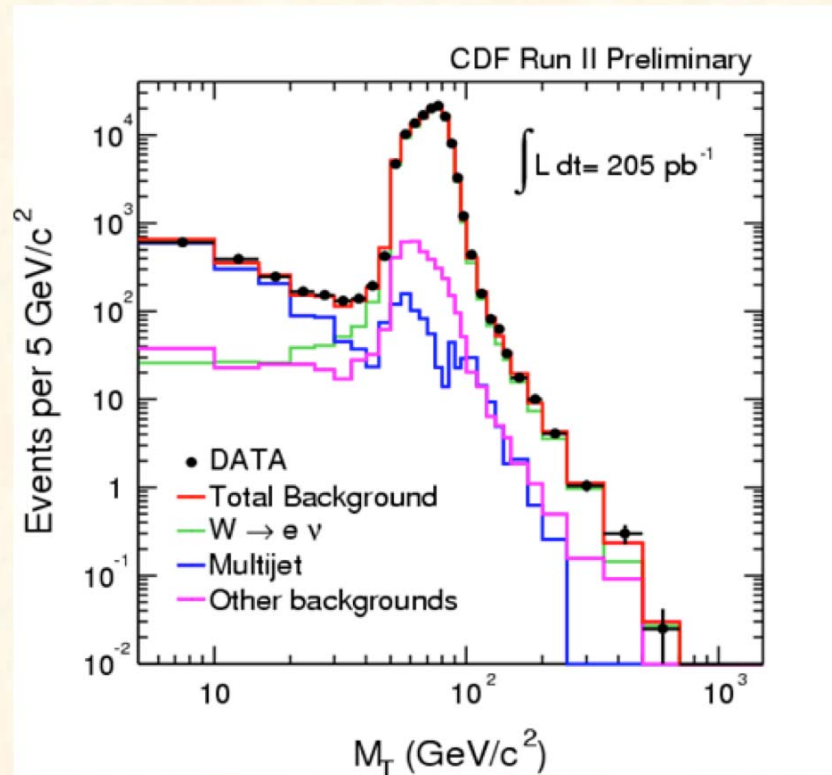
Signature: high  $p_T$  electron + high  $E_T^{\text{miss}}$

→ peak in transverse mass distribution





## Search for $W' \rightarrow e\nu$



Data:

consistent with one well known  $W$   
+ background



Limit:  $M(W') > 842 \text{ GeV}/c^2$

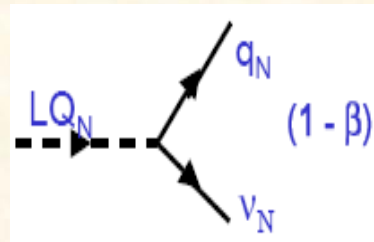
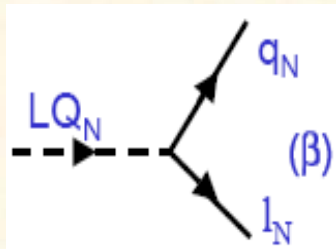
(assuming Standard Model couplings)



# Search for Scalar Leptoquarks (LQ)

- Production:  
pair production via QCD processes  
(qq and gg fusion)

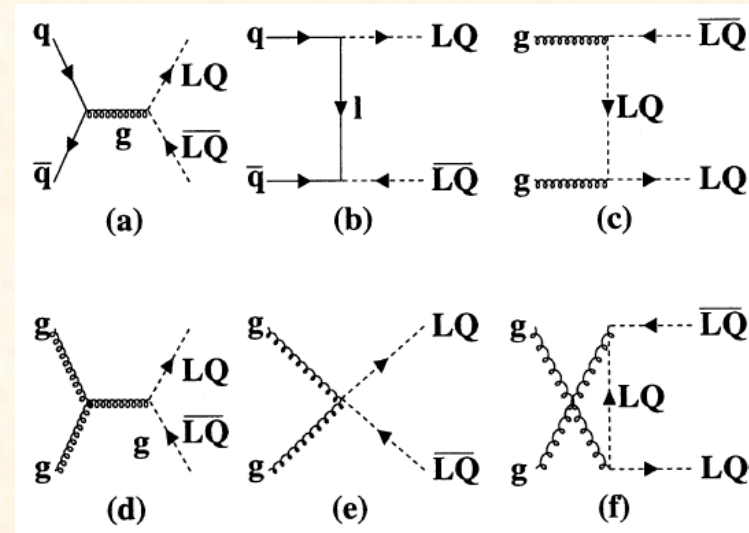
- Decay: into a lepton and a quark



$\beta$  = LQ branching fraction to charged lepton and quark

N = generation index

Leptoquarks of 1., 2., and 3. generation



## Experimental Signatures:

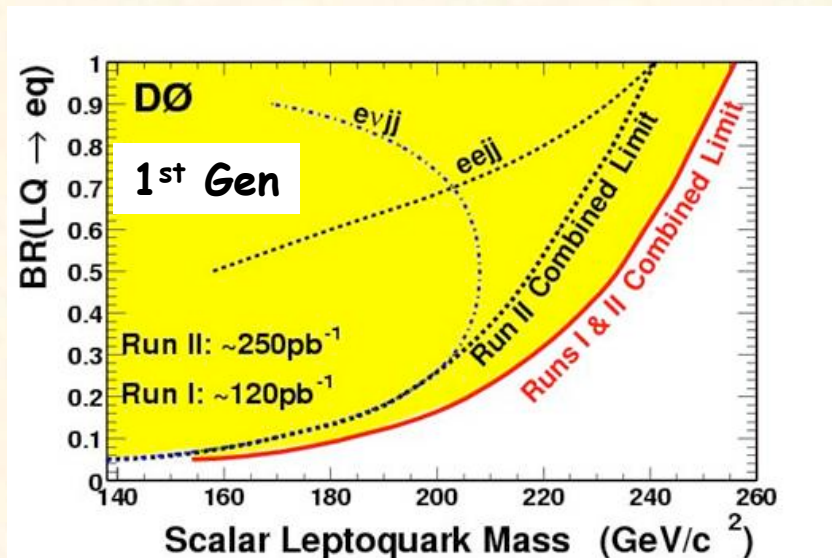
- two high  $p_T$  isolated leptons + jets .OR.
- one isolated lepton +  $P_{T}^{miss}$  + jets .OR.
- $P_{T}^{miss}$  + jets



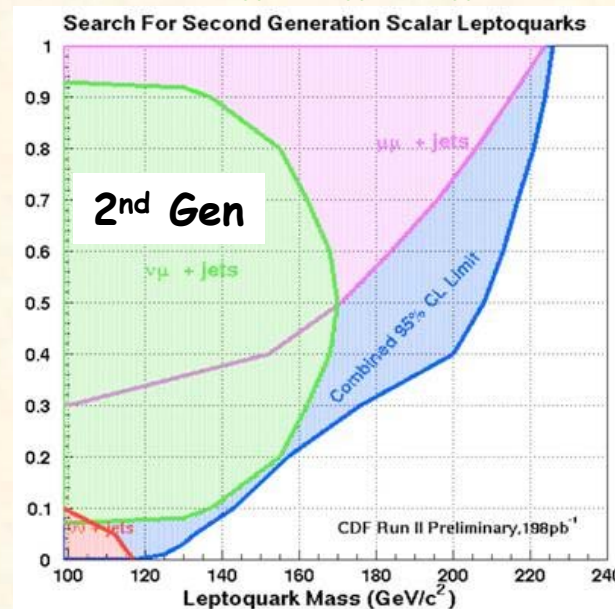
# 1<sup>st</sup>, 2<sup>nd</sup> and 3<sup>rd</sup> generation Leptoquarks



channels:  $ee\ jj$ ,  $ev\ jj$



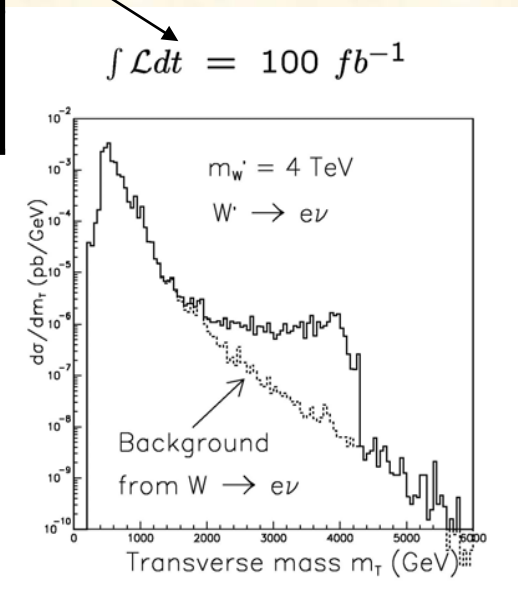
channels:  $\mu\mu\ jj$ ,  $\epsilon\nu\ jj$ ,  $\nu\nu\ jj$



95% C.L. Mass Limits	1. Generation LQ	2. Generation LQ	3. Generation LQ
CDF (Run II)	235 GeV/c <sup>2</sup>	224 GeV/c <sup>2</sup>	129 GeV/c <sup>2</sup>
DØ (Run I + II)	256 GeV/c <sup>2</sup>	200 GeV/c <sup>2</sup> (Run I)	

# LHC reach for other BSM Physics (a few examples for 30 and 100 fb<sup>-1</sup>)

	30 fb <sup>-1</sup>	100 fb <sup>-1</sup>
Excited Quarks $Q^* \rightarrow q \gamma$	$M(q^*) \sim 3.5 \text{ TeV}$	$M(q^*) \sim 6 \text{ TeV}$
Leptoquarks	$M(\text{LQ}) \sim 1 \text{ TeV}$	$M(\text{LQ}) \sim 1.5 \text{ TeV}$
$Z' \rightarrow \ell\ell, jj$ $W' \rightarrow \ell \nu$	$M(Z') \sim 3 \text{ TeV}$ $M(W') \sim 4 \text{ TeV}$	$M(Z') \sim 5 \text{ TeV}$ $M(W') \sim 6 \text{ TeV}$
Compositeness (from Di-jet)	$\Lambda \sim 25 \text{ TeV}$	$\Lambda \sim 40 \text{ TeV}$



# Sensitivity to New Physics with jets in Early LHC data

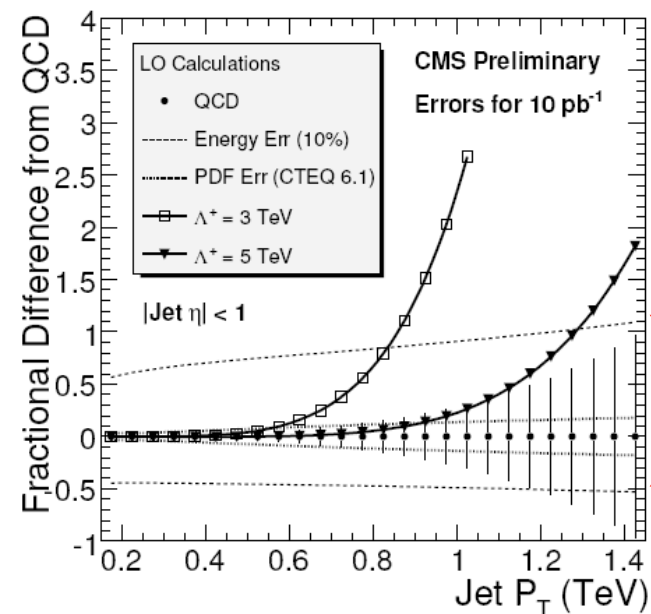
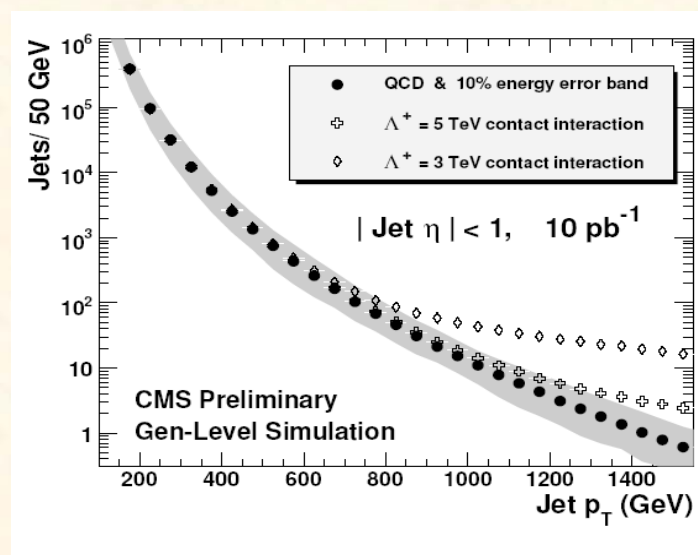
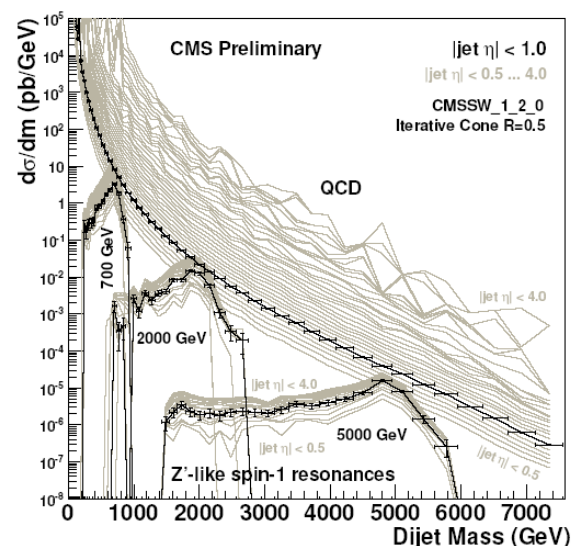
- Even with JES uncertainties expected with early data and an int. luminosity of only  $10 \text{ pb}^{-1}$  compositeness scales of  $\sim 3 \text{ TeV}$  can be reached

(close to the present Tevatron reach of  $\Lambda > 2.7 \text{ TeV}$ )

- Resonances decaying into two jets:

Discovery sensitivity around  $2 \text{ TeV}$   
(Spin-1  $Z'$  like resonance) for  $\sim 200 \text{ pb}^{-1}$

Present Tevatron limits:  $320 < m < 740 \text{ GeV}$



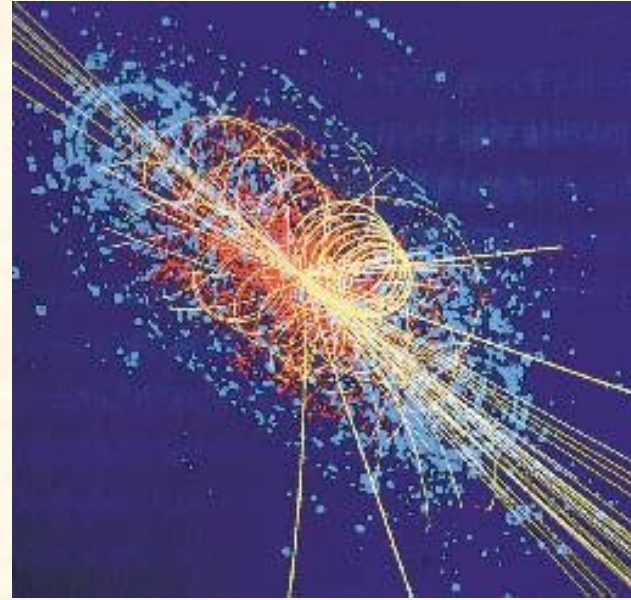
10% JES

# Conclusions

1. Experiments at Hadron Colliders have a huge discovery potential
  - **SM Higgs**: full mass range, already at low luminosity;  
Vector boson fusion channels improve the sensitivity significantly
  - **MSSM Higgs**: parameter space covered
  - **SUSY**: discovery of TeV-scale SUSY should be easy,  
determination of model parameters is more difficult
  - **Exotics**: experiments seem robust enough to cope with new scenarios
2. Experiments have also a great potential for precision measurements
  - $m_W$  to  $\sim 10 - 15$  MeV
  - $m_t$  to  $\sim 1$  GeV
  - $\Delta m_H / m_H$  to 0.1% (100 - 600 GeV)
  - + gauge couplings and measurements in the top sector .....



# End of lectures



- In case you have any questions:  
please do not hesitate to contact me: [karl.jakobs@uni-freiburg.de](mailto:karl.jakobs@uni-freiburg.de)
- Transparencies will be made available as .pdf files on the web  
(school pages)