

SEP. 23 - 27, 2019

Korea Institute for Advanced Study (KIAS)

1F Auditorium

**Search for muon-philic
new light gauge boson at
Belle II**

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based on Yongsoo Jho, Yongjoon
Kwon, SCP, Po-Yan Tseng,
1904.13053

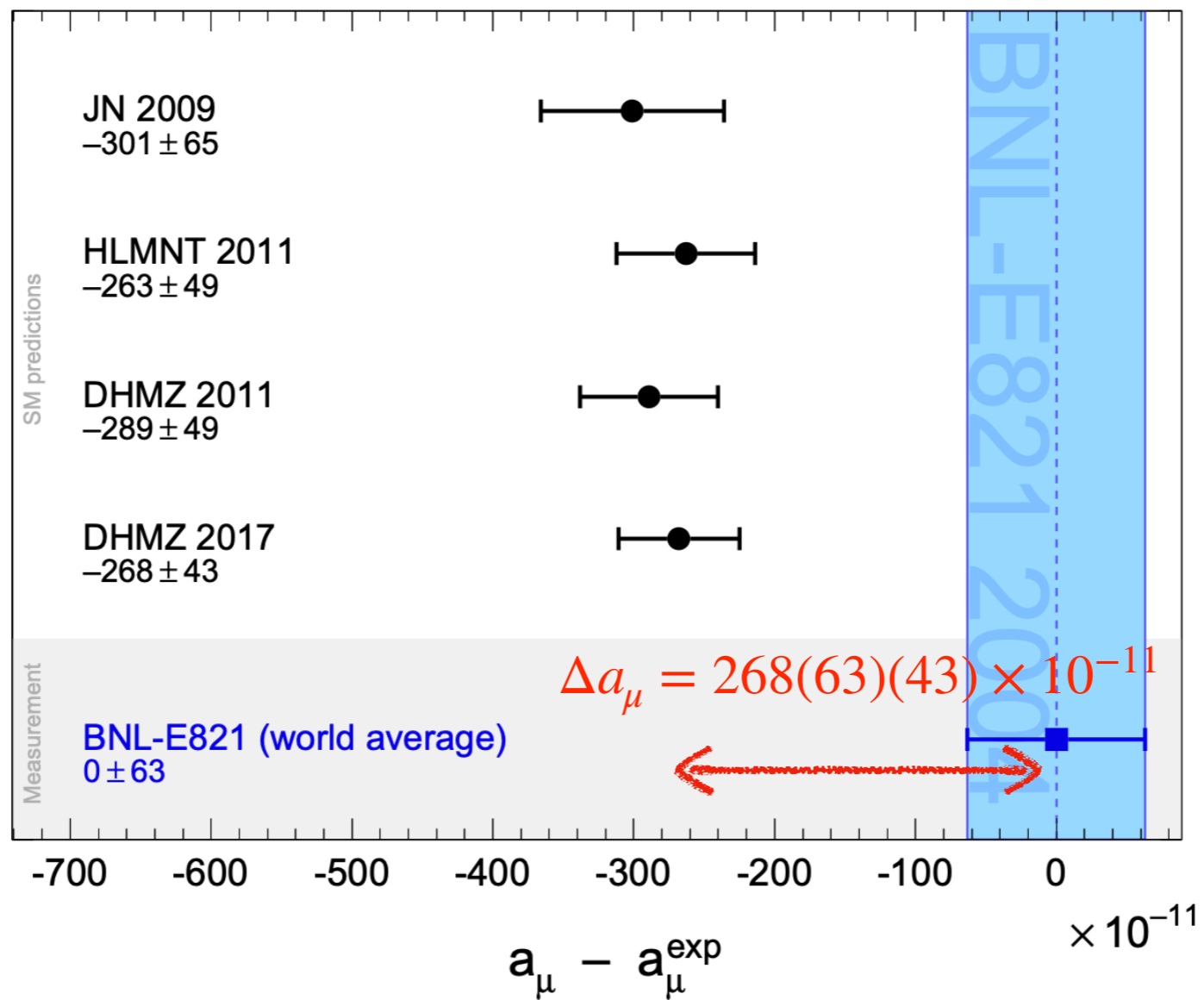
**INTERNATIONAL WORKSHOP ON
NEW PHYSICS AT THE LOW ENERGY SCALES**

NEPLES-2019

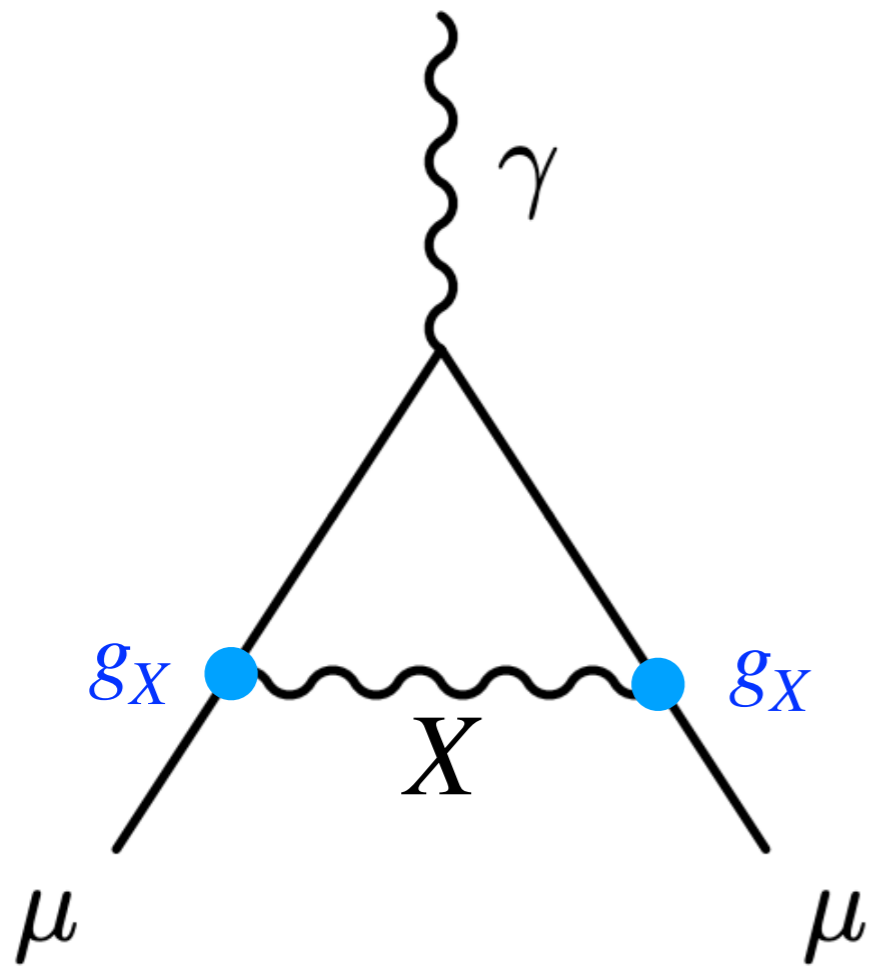
outline

- Why μ -philic light boson?
- Constraints
- Search strategy at Belle-II
- Conclusion

Motivation



X contribution to $(g - 2)_\mu$



$$\Delta a_\mu^X = \frac{g_X^2}{8\pi^2} \int_0^1 dz \frac{2z(1-z)^2}{(1-z)^2 + (m_X/m_\mu)^2 z}$$

$$\approx \begin{cases} 1 & \text{if } m_X \ll m_\mu \\ \frac{2}{3} \frac{m_\mu^2}{m_X^2} & \text{if } m_X \gg m_\mu \end{cases}$$

similar for a scalar

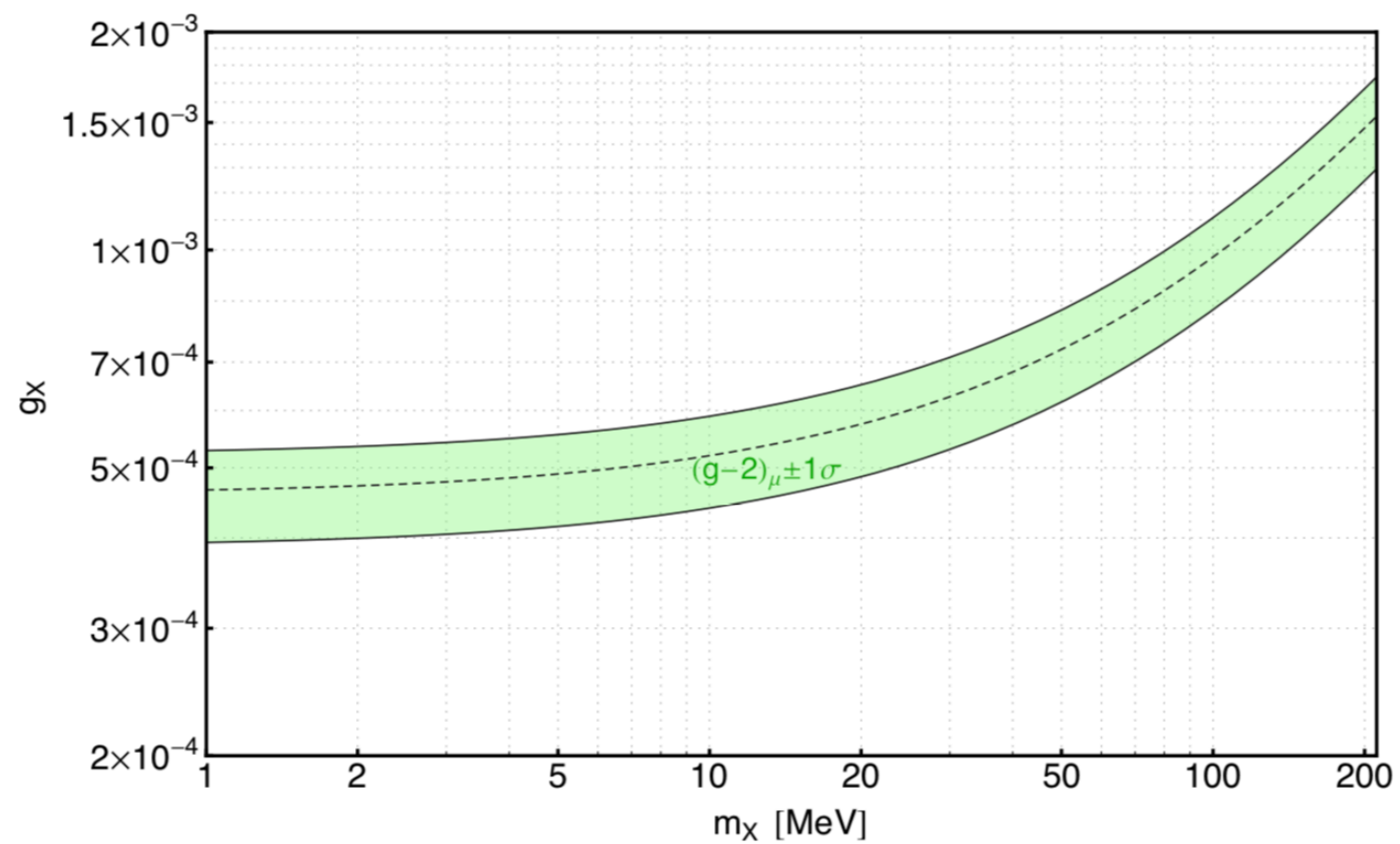
‘Heavy’ : $m_X \gg m_\mu$

$$\Delta a_\mu^X = \frac{g_X^2}{8\pi^2} \int_0^1 dz \frac{2z(1-z)^2}{(1-z)^2 + (m_X/m_\mu)^2 z} \approx \frac{g_X^2}{12\pi^2} \frac{m_\mu^2}{m_X^2}$$

$$g_X \sim 1, \quad m_X \sim 330 \text{ GeV}$$

ruled out by LHC

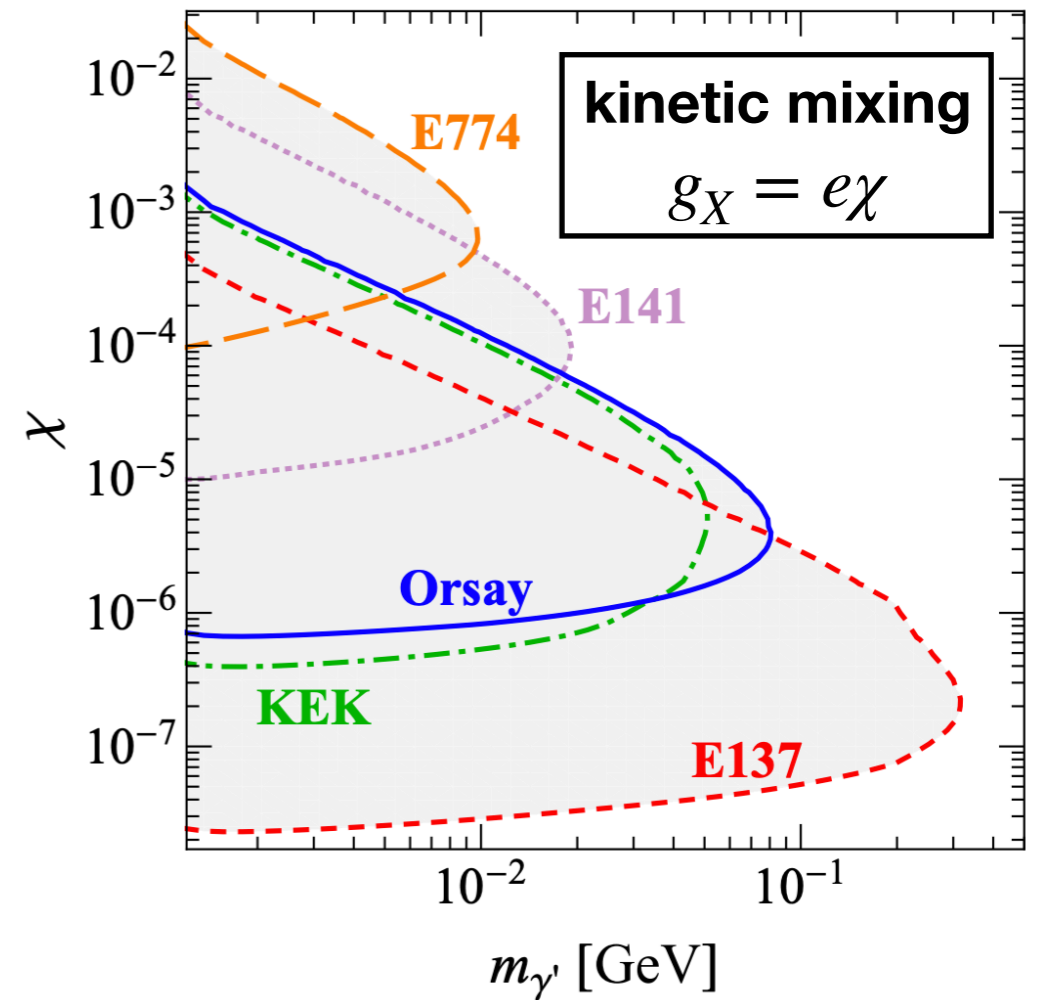
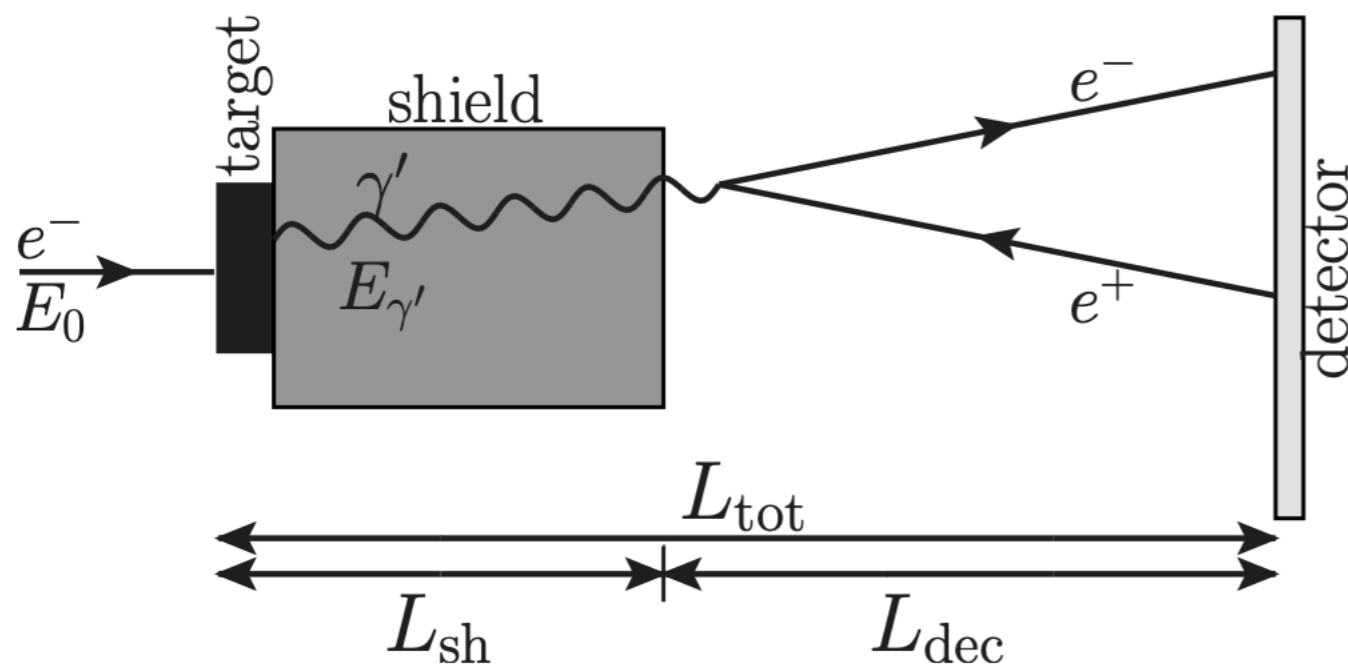
‘Light’: $m_X \leq 2m_\mu$



Ball park: $g_X \sim 5 \times 10^{-4}$, $m_X \sim 10$ MeV

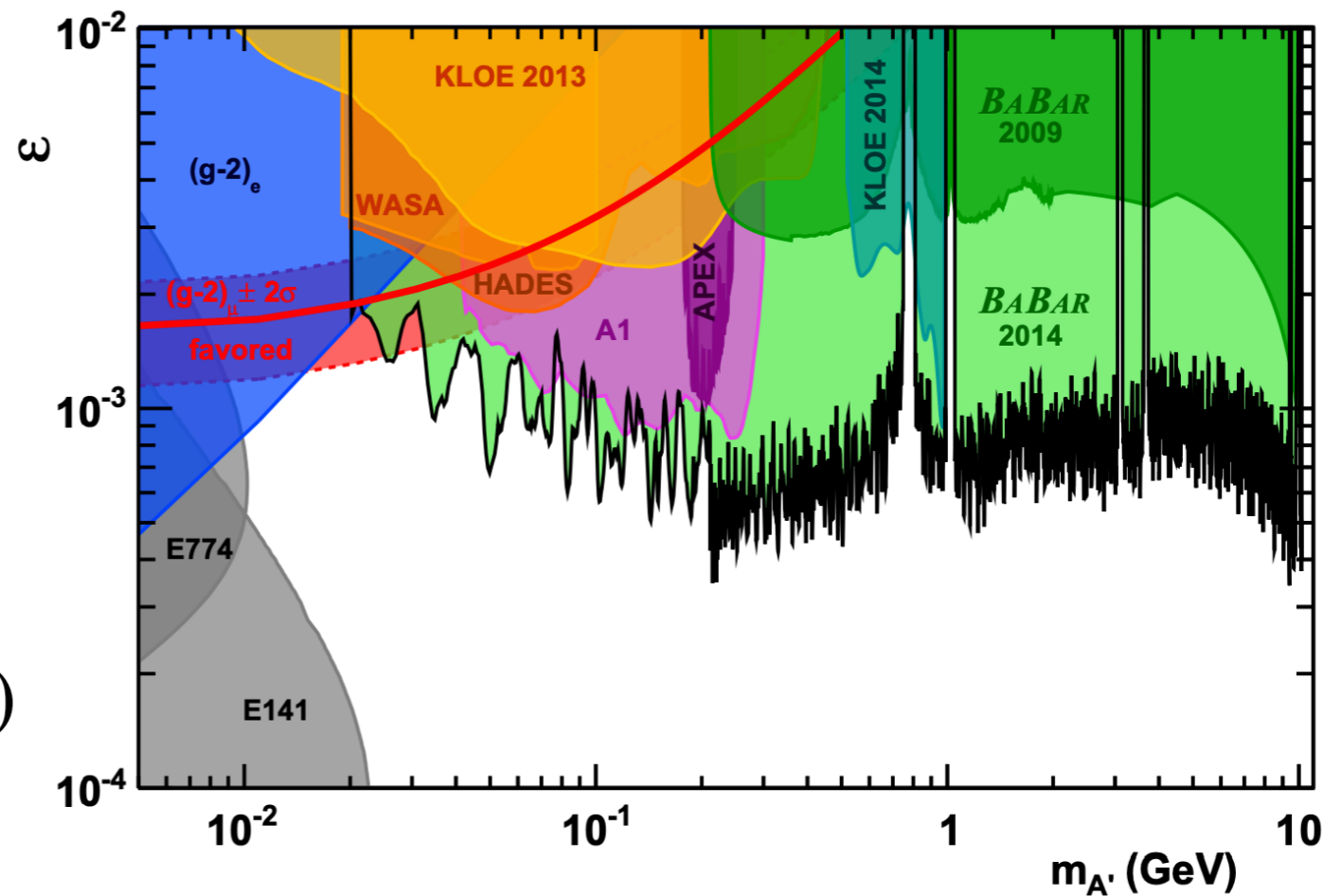
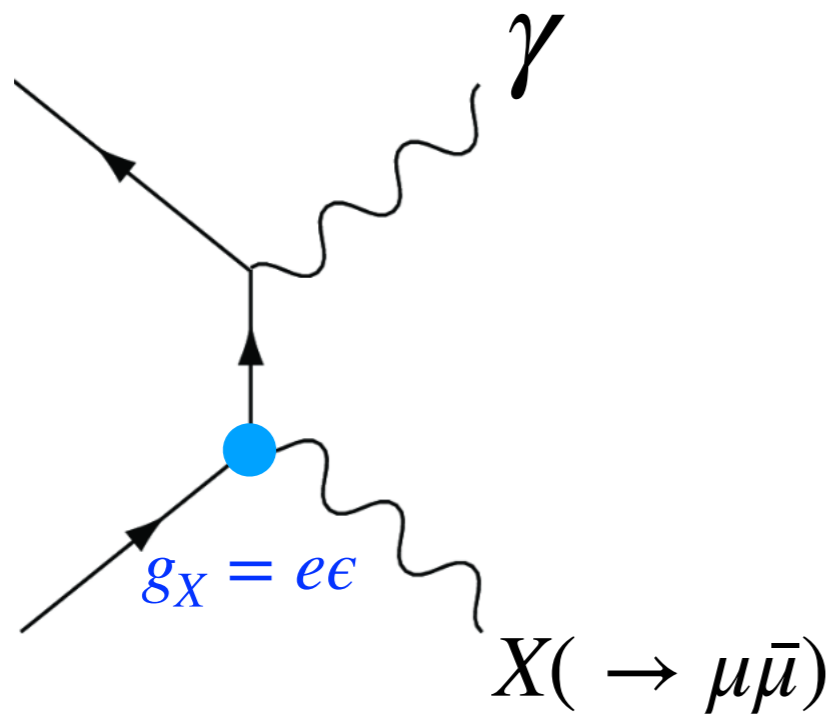
Light, flavor universal X:

Beam dump $e^+e^- \rightarrow \gamma X(e^+e^-)$



Light, flavor universal case:

$$\text{BarBar } e^+e^- \rightarrow \gamma X (X \rightarrow \mu^+\mu^-)$$



**No interesting parameter for $(g - 2)_\mu$
anomaly exists for light, flavor universal X**

~~No~~ interesting parameter for $(g - 2)_\mu$
anomaly exists for ~~flavor universal X~~
 μ - philic

e.g. $U(1)_{L_\mu}$, $U(1)_{L_\mu - B_i}$, $U(1)_{L_\mu - L_\tau}$

L_μ

- Lepton numbers (L_e, L_μ, L_τ) are accidental symmetries of the SM.

- Anomalous

- WZW contribution:

$$g_X \epsilon^{\mu\nu\rho\sigma} X_\mu (g'^2 B_\nu \partial_\rho B_\sigma + g^2 W_\nu^a \partial_\rho W_\sigma^a + \frac{g^2}{3} \epsilon^{abc} W_\nu^a W_\rho^b W_\sigma^c)$$

induce $Z \rightarrow X\gamma, B \rightarrow KX, B \rightarrow \pi X \dots$

- highly constrained

$$L_\mu - B_{i=1,2,3}$$

- Anomaly free
- Has additional hadronic interactions
- highly constrained

$$L_\mu - L_\tau$$

- The differences $(L_i - L_j)$, $i \neq j$ are anomaly free in the SM.
- Economical extension of the SM:
 $G_{SM} \times U(1)_{L_i - L_j}$ does not need any exotic fermion.
- $L_\mu - L_\tau$ explains $(g - 2)_\mu$: our X-boson!

Model Lagrangian

Only two free parameters

$$\mathcal{L}_{\text{minimal}} \ni -\frac{1}{4} X_{\mu\nu} X^{\mu\nu} - \frac{1}{2} m_X^2 X^\mu X_\mu - g_X X_\mu J_X^\mu$$

$$J_X^\mu = \bar{\mu} \gamma^\mu \mu - \bar{\tau} \gamma^\mu \tau + \bar{\nu}_\mu \gamma^\mu \nu_{\mu L} + \bar{\nu}_\tau \gamma^\mu \nu_{\tau L}$$

(possible)
Extension with dark sector

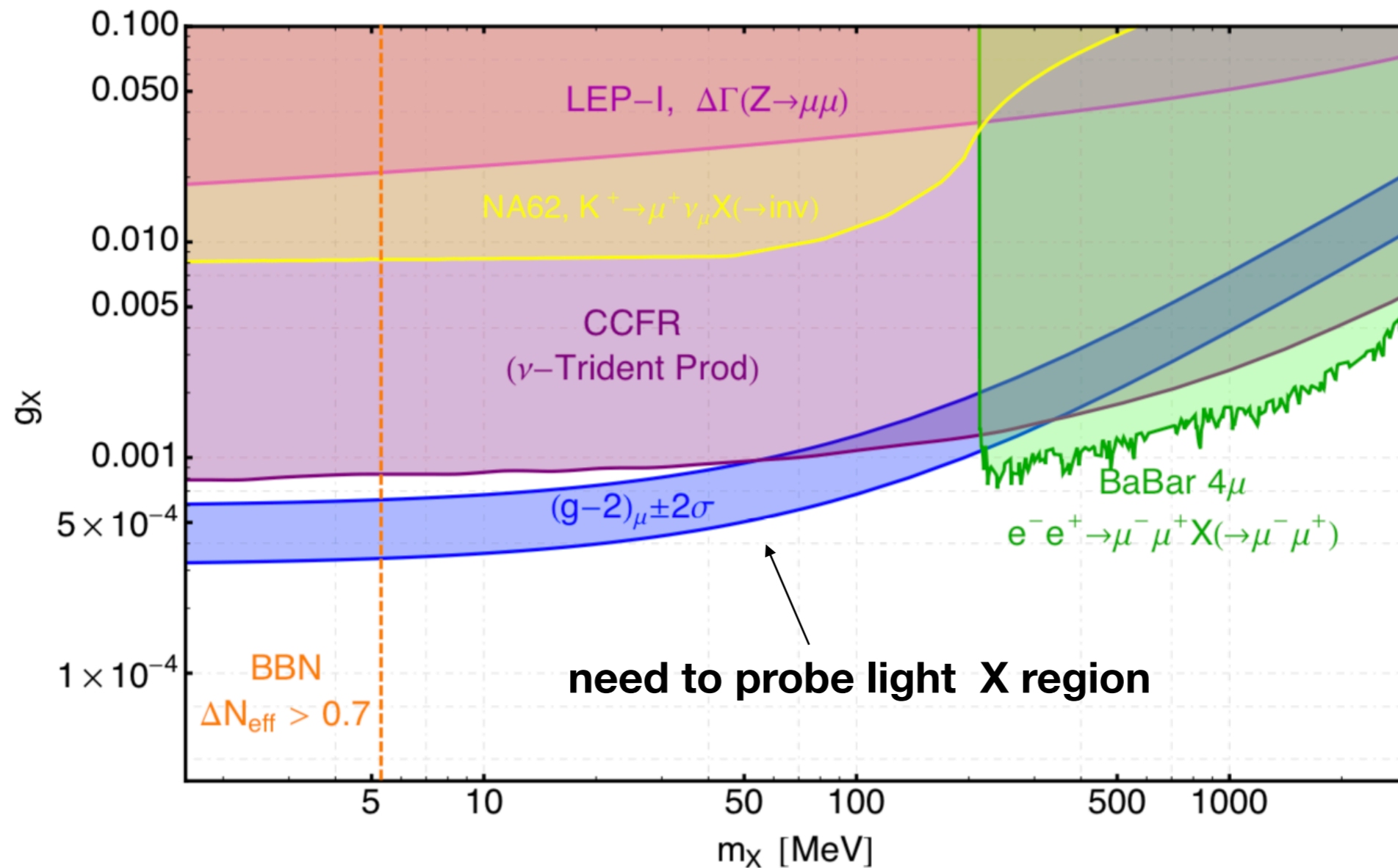
$$\mathcal{L} = \mathcal{L}_{\text{minimal}} + \sum_{i=1}^{N_\chi} \left[\bar{\chi}_i (i\gamma^\mu \partial_\mu - m_{\chi_i}) \chi_i - g_D (\bar{\chi}_i \gamma^\mu \chi_i) X_\mu \right]$$

χ_i : **dark sector fermions**

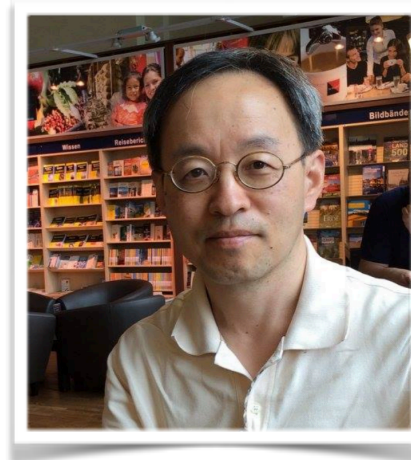
$$\Gamma_{X,\text{total}} = (1 + \delta_{\text{NM}}) \cdot \Gamma_{\text{Minimal}} \quad \text{:enhanced decay width}$$

$$\delta_{\text{NM}} = \sum_{i=1}^{N_\chi} \frac{g_D^2}{g_X^2} \left(1 + \frac{m_{\chi_i}^2}{m_X^2}\right) \sqrt{1 - 4 \frac{m_{\chi_i}^2}{m_X^2}}$$

Known constraints

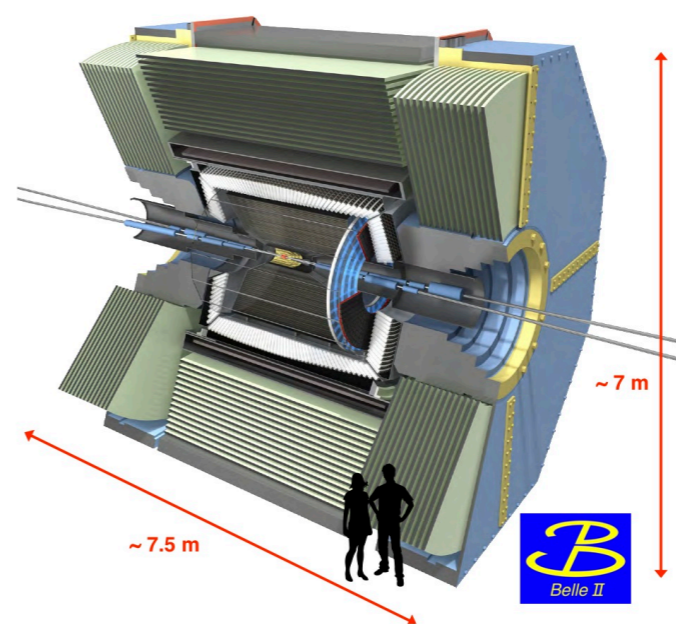
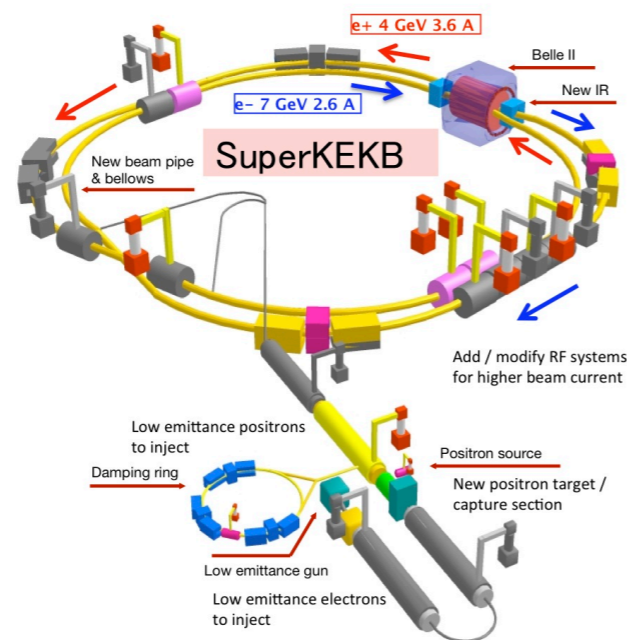


Ball park: $g_X \sim 5 \times 10^{-4}, m_X \sim 10 \text{ MeV}$



Belle-II(50ab^{-1}) = 50 Belle

:an ideal place to look for a light-X



i) $\epsilon(p_{\mu^\pm} > 0.6\text{GeV}/c) = 0.9$ at KLM(K_L, μ) detector

Excellent muon detector

$$\sigma_{p_{\mu^\pm}} / p_{\mu^\pm} = 0.0011 p_{\mu^\pm} [\text{GeV}] \oplus 0.0025 / \beta$$

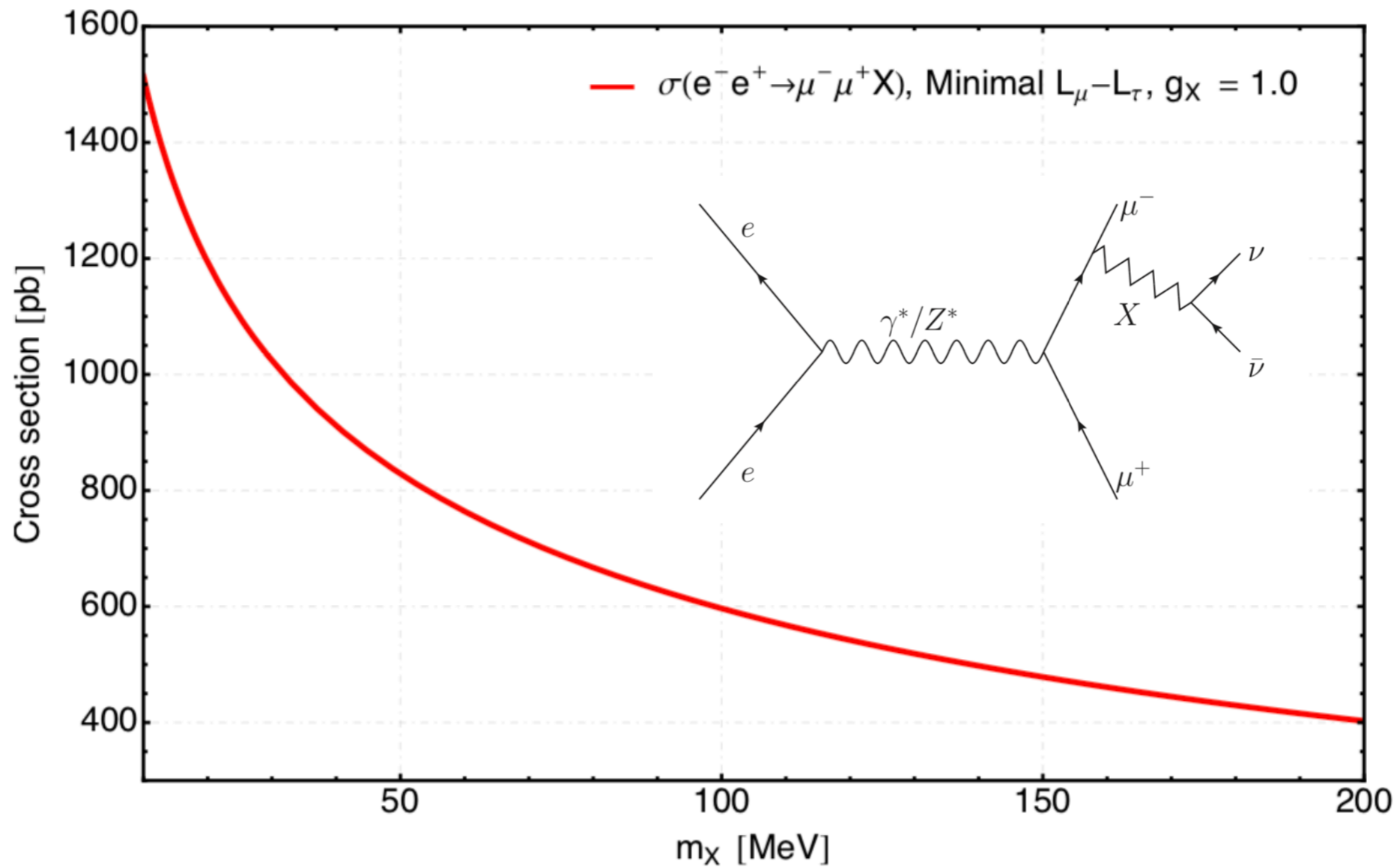
Excellent muon tracker (Central Drift Chamber)

ii) $1 - \epsilon_\gamma = 10^{-6}$ Low intrinsic probability of missing the photon detection inside the ECL crystals +KLM detectability (end cap region)

Excellent photon detector for background rejection

X-boson at Belle II

$$e^+e^- \rightarrow \mu^+\mu^-X(\rightarrow \nu\bar{\nu})$$

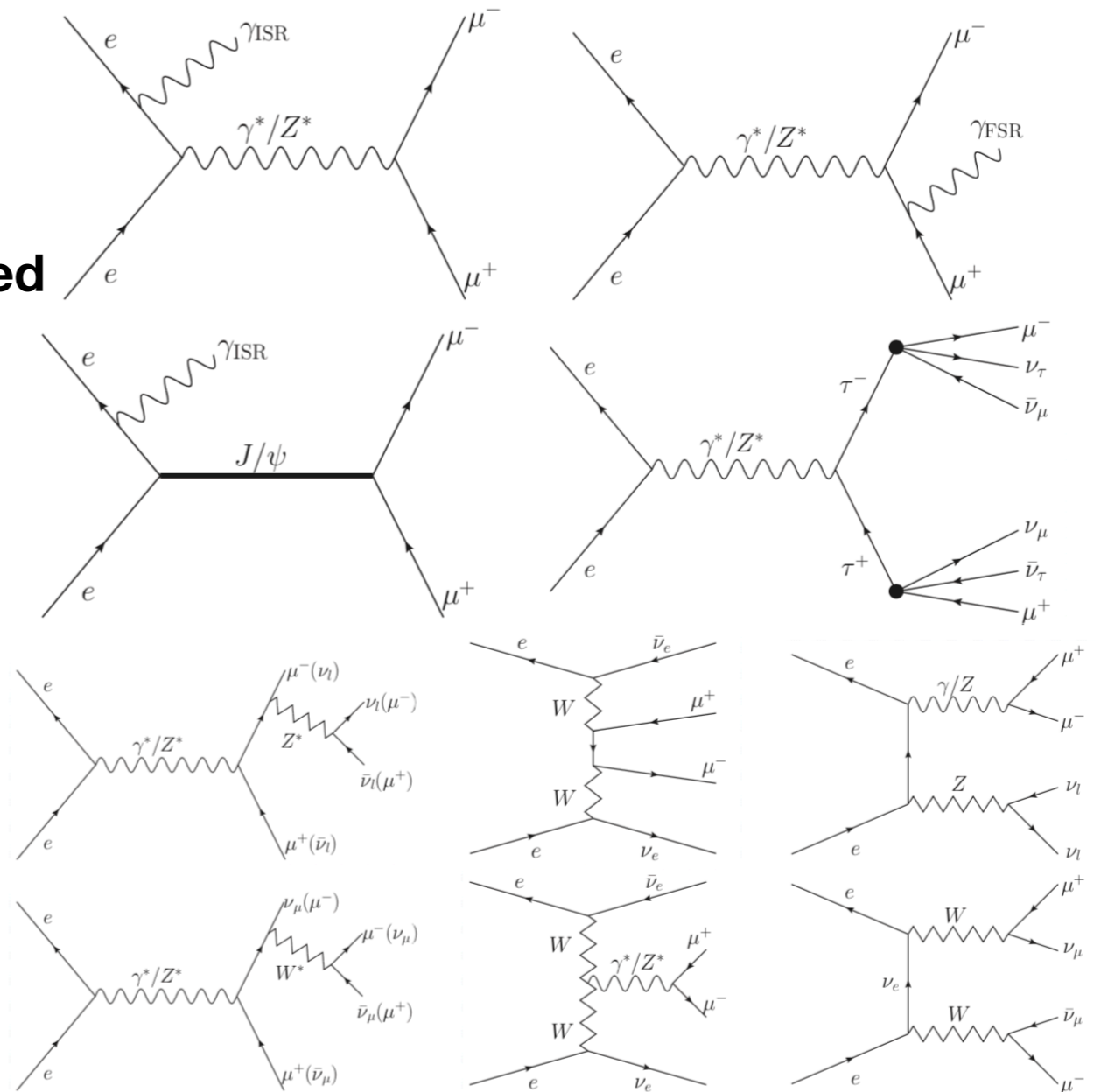


SM background: $\mu^+ \mu^- + MET$

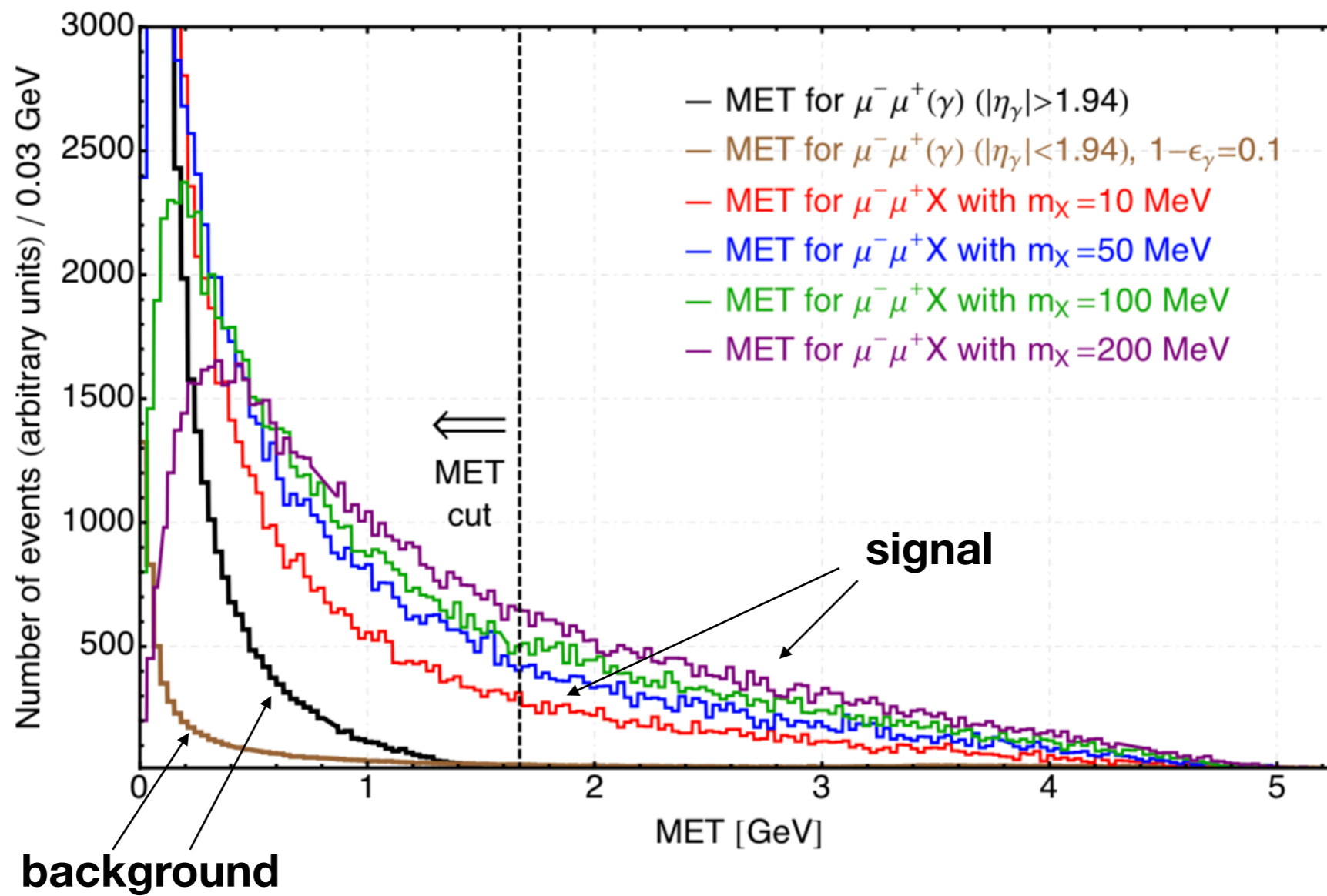
- $e^- e^+ \rightarrow \mu^- \mu^+ (\gamma_{ISR,FSR})$: **soft γ not detected**

- $e^- e^+ \rightarrow \tau^- \tau^+ \rightarrow \mu^- \mu^+ \nu_\mu \bar{\nu}_\mu \nu_\tau \bar{\nu}_\tau$

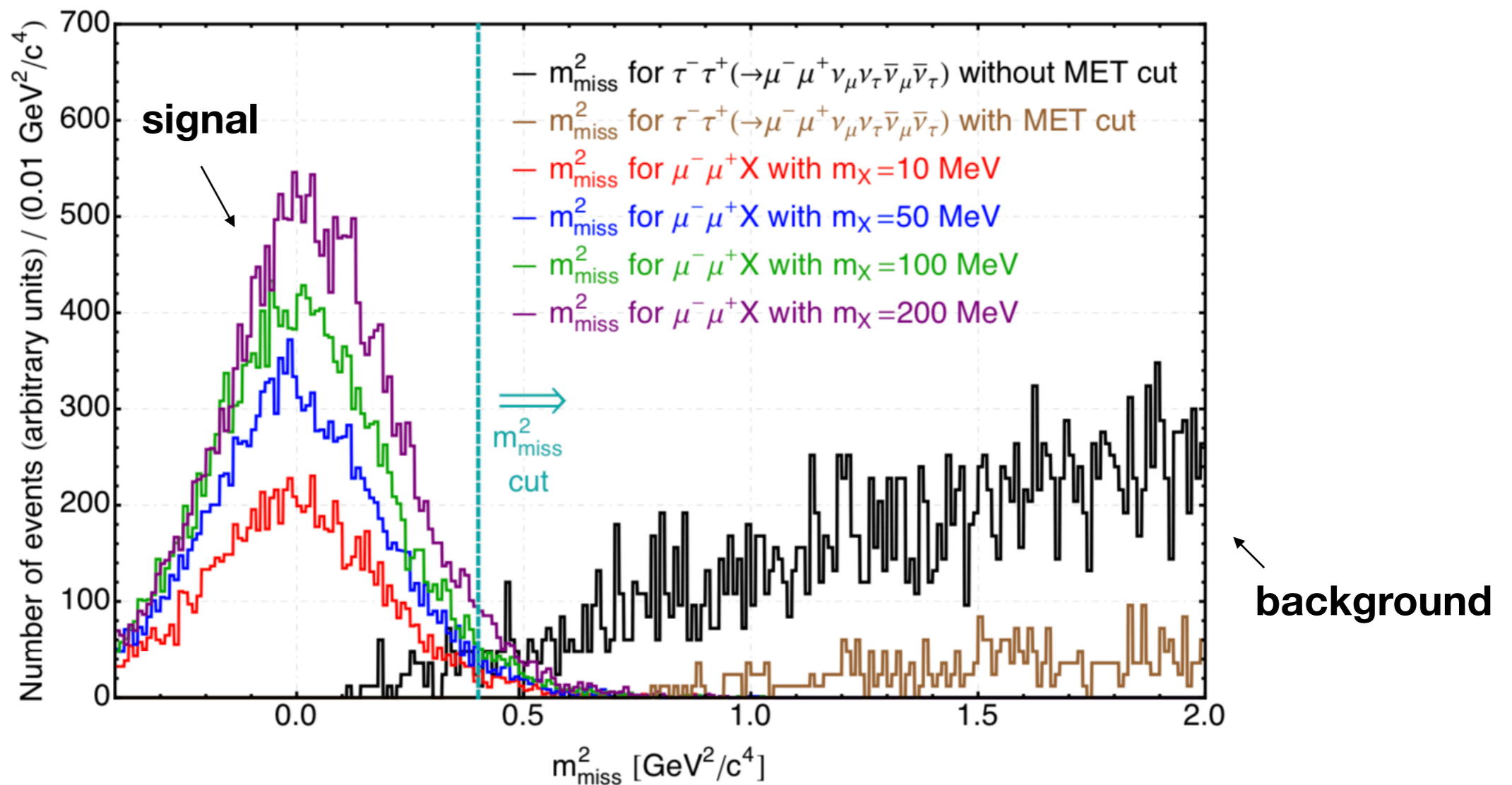
- $e^- e^+ \rightarrow \mu^- \mu^+ \nu_l \bar{\nu}_l$ by off-shell W and Z



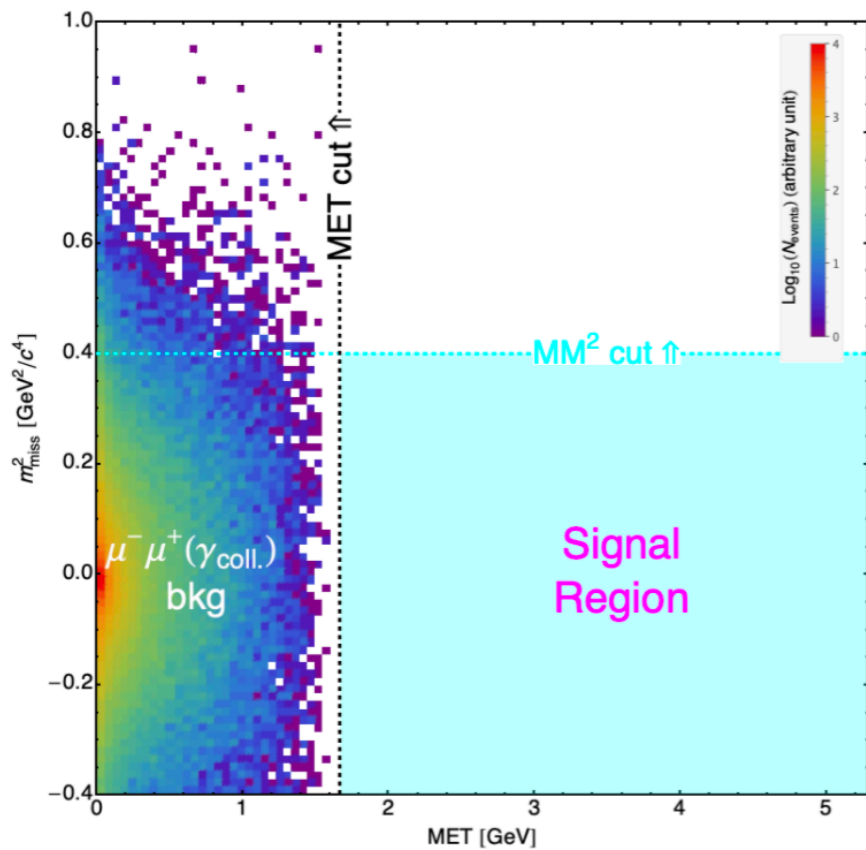
MET cut



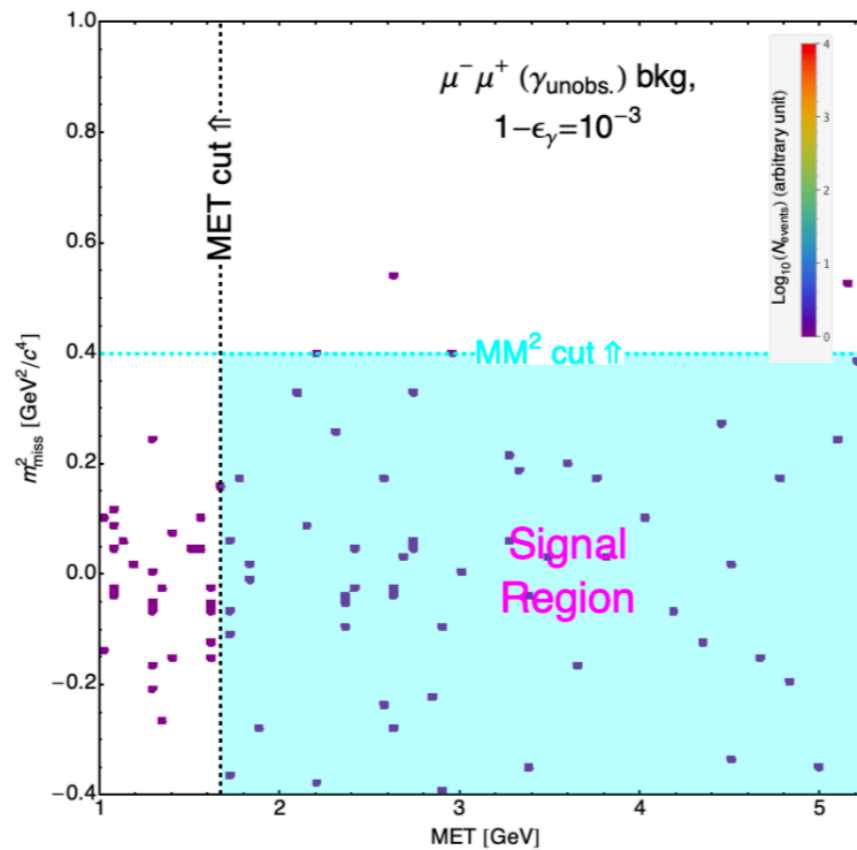
Missing mass cut



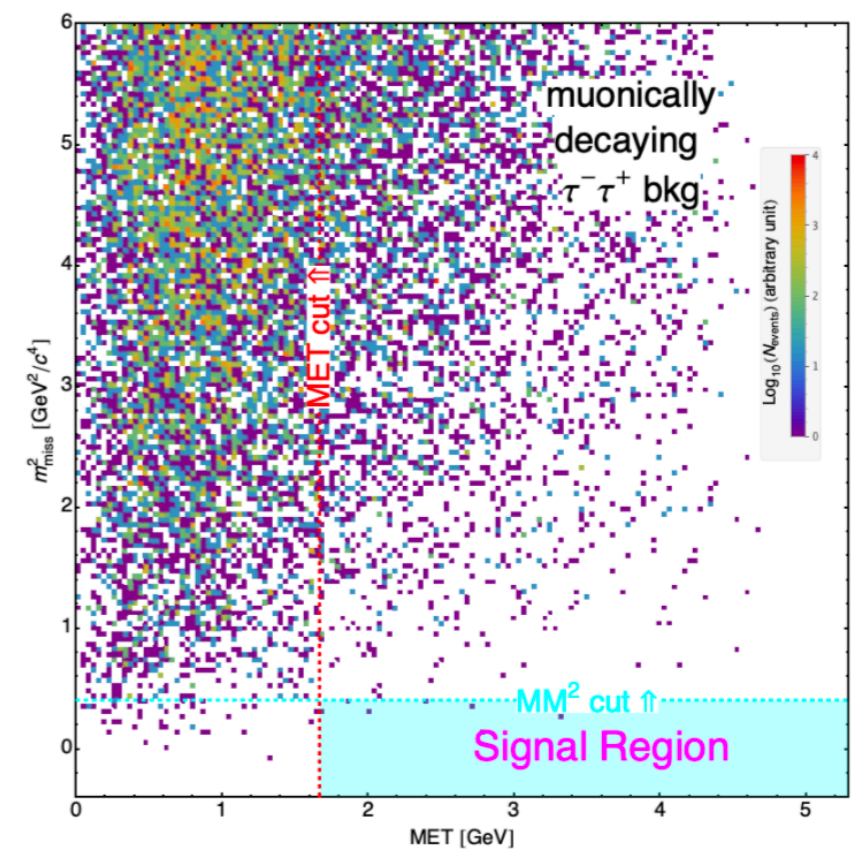
Cutting background



(a) $\mu^- \mu^+ \gamma$ with $|\eta_\gamma^*| > 1.94$
photon to beam pipe



(b) $\mu^- \mu^+ \gamma$ with $|\eta_\gamma^*| < 1.94$

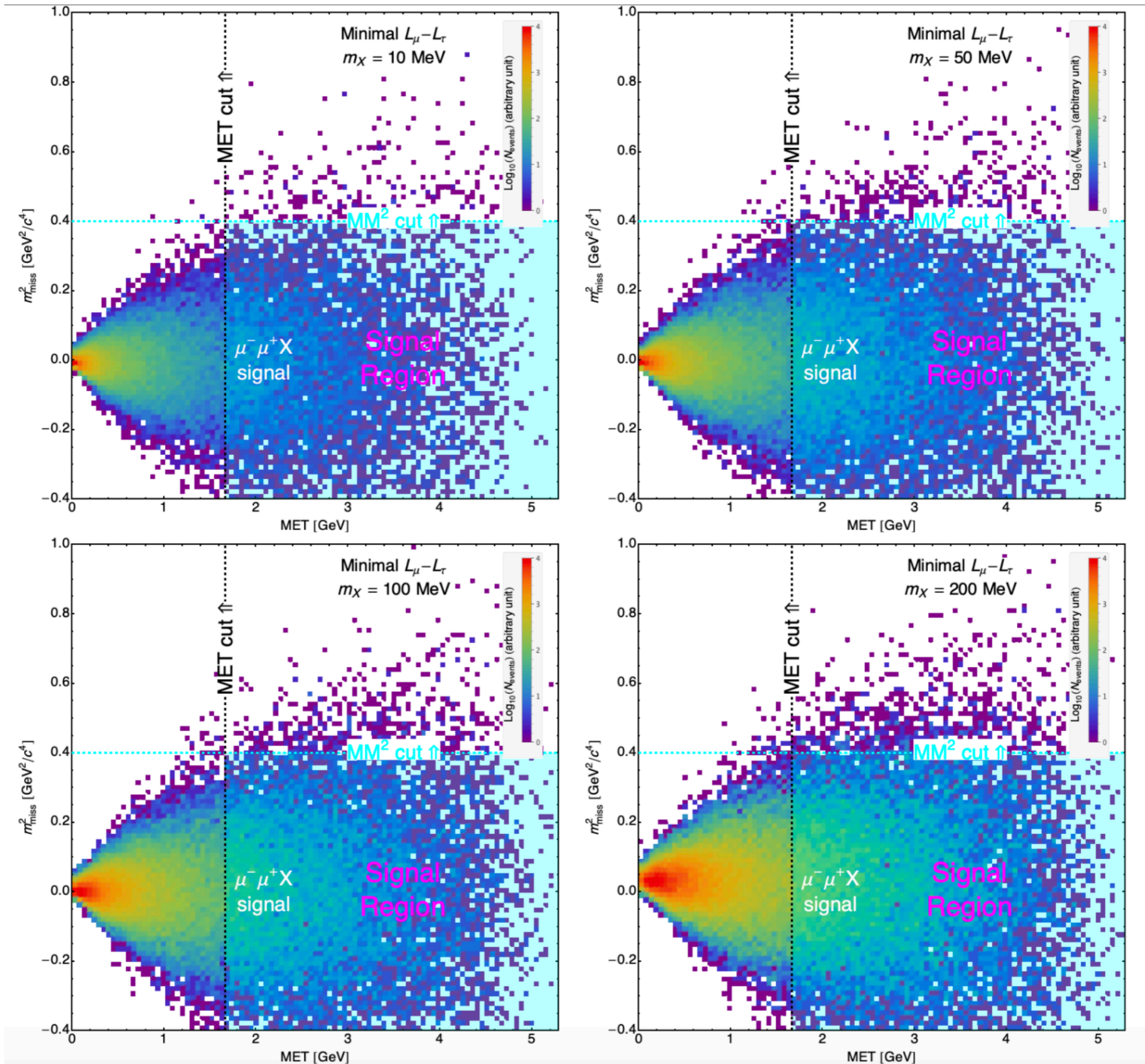


(c) $\tau^- \tau^+ (\rightarrow \mu^- \mu^+ \nu \bar{\nu} \bar{\nu} \bar{\nu})$

i) $\cancel{E}_T > 1.67$ GeV,

ii) $m_{\text{miss}}^2 < 0.4$ GeV²/c⁴

Signal



Signal/background

$$\mathcal{S} = \left[\int_{\cancel{E}_T^{\text{cut}}}^{\sqrt{s}/2} d\cancel{E}_T \int_{(m_{\text{miss}}^2)_{\text{cut}}^{\text{min}}}^{(m_{\text{miss}}^2)_{\text{cut}}^{\text{max}}} dm_{\text{miss}}^2 [\epsilon(p_{\mu^\pm})]^2 \frac{d^2\sigma_{\mu\mu X}^{\text{signal}}(s)}{d\cancel{E}_T dm_{\text{miss}}^2} \right] \cdot \int \mathcal{L} dt,$$
$$\mathcal{B} = \left[\int_{\cancel{E}_T^{\text{cut}}}^{\sqrt{s}/2} d\cancel{E}_T \int_{(m_{\text{miss}}^2)_{\text{cut}}^{\text{min}}}^{(m_{\text{miss}}^2)_{\text{cut}}^{\text{max}}} dm_{\text{miss}}^2 [\epsilon(p_{\mu^\pm})]^2 \frac{d^2\sigma_{\mu\mu\gamma,\tau\tau,W^*/Z^*}^{\text{background}}(s)}{d\cancel{E}_T dm_{\text{miss}}^2} \right] \cdot \int \mathcal{L} dt,$$

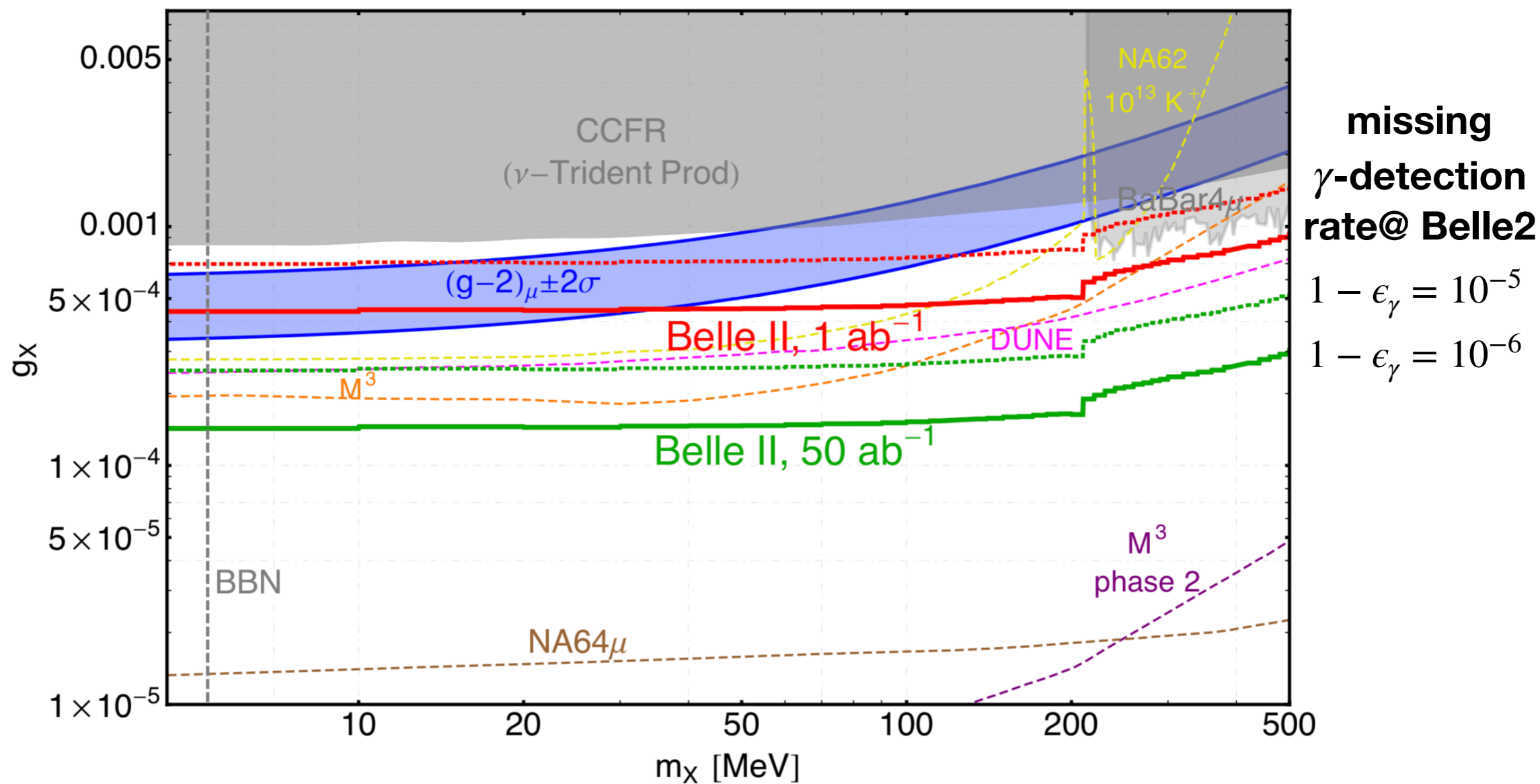
i) $\cancel{E}_T > 1.67 \text{ GeV},$

ii) $m_{\text{miss}}^2 < 0.4 \text{ GeV}^2/c^4$

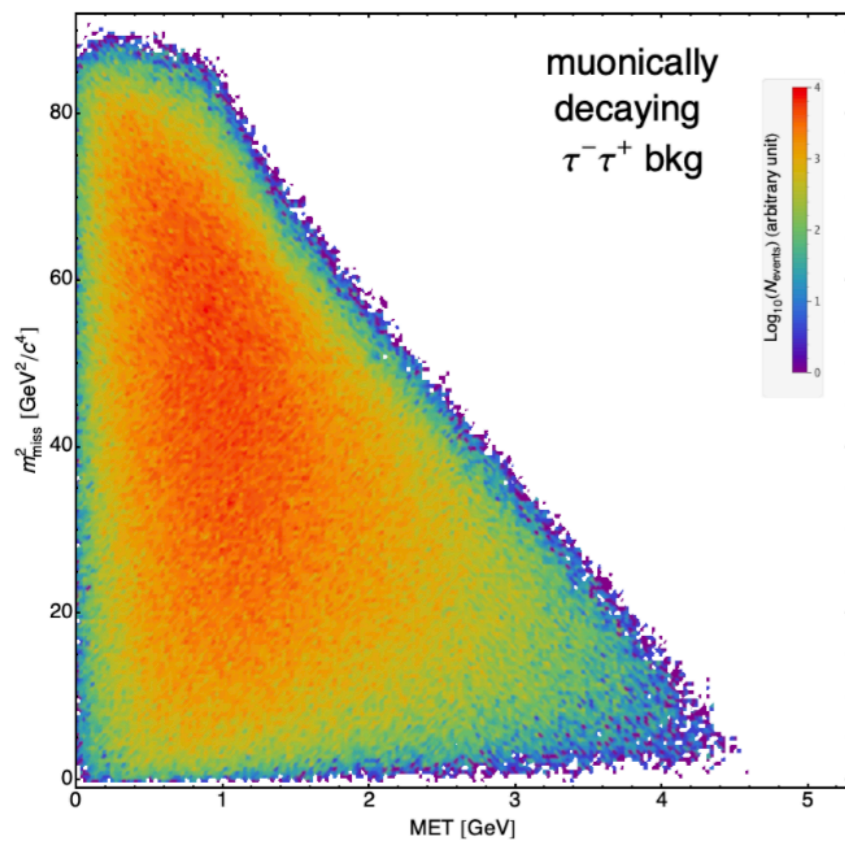
iii) $\epsilon(p_{\mu^\pm} > 0.6\text{GeV}/c) = 0.9$ at KLM(K_L, μ) detector

Sensitivity limit:

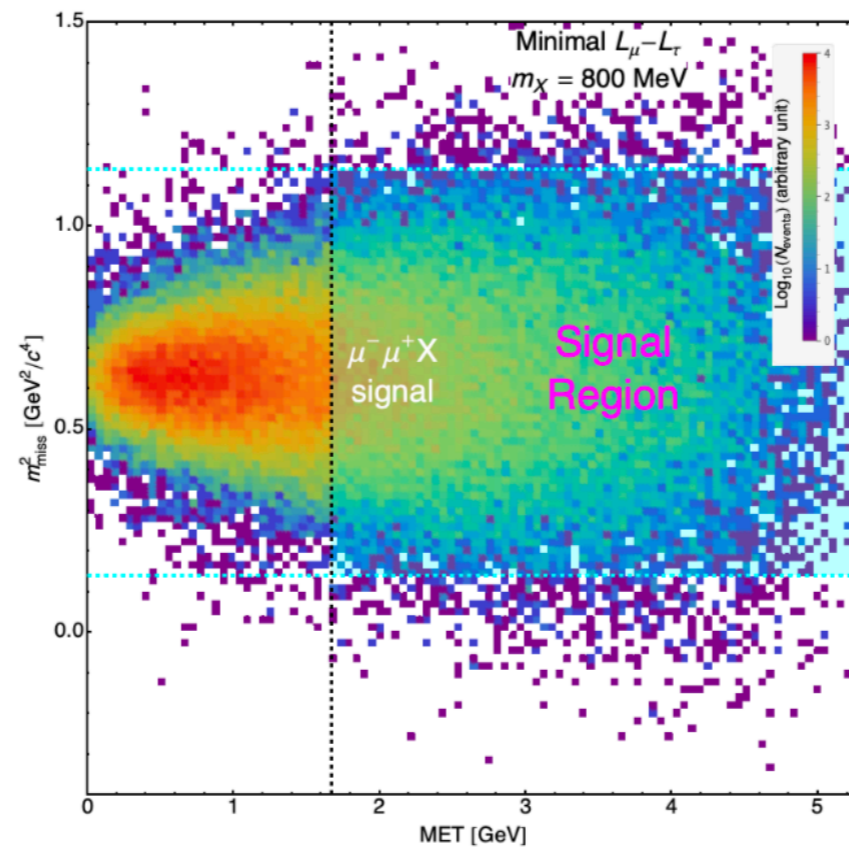
$$S/\sqrt{S+B} = 3$$



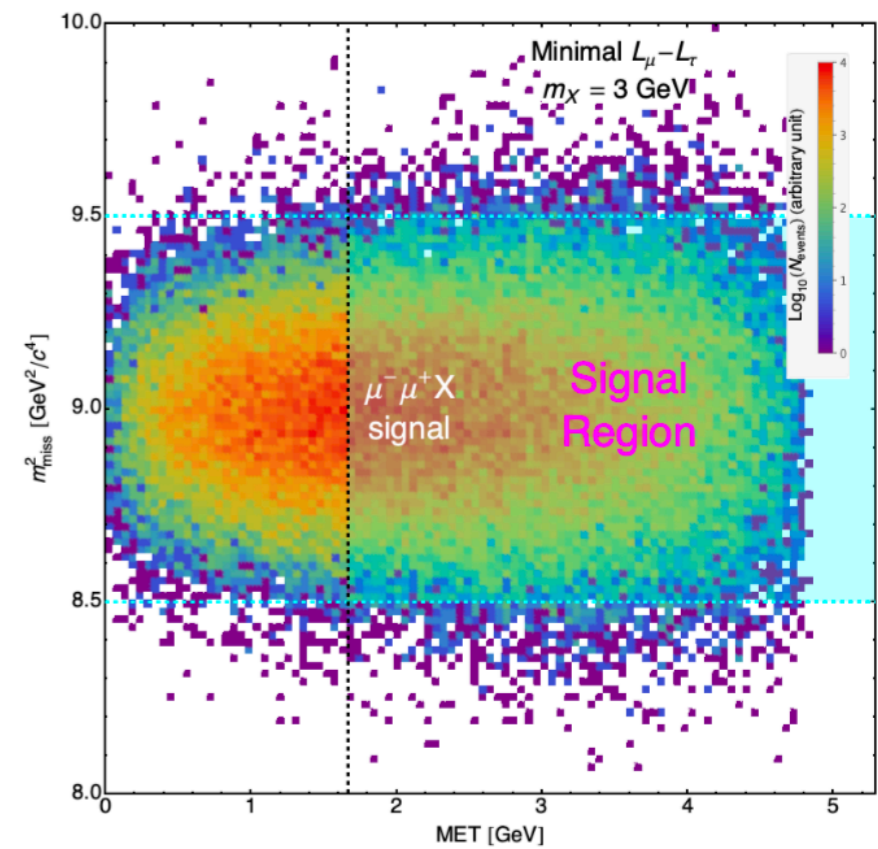
For heavier X ($m_X > 2m_\mu$)



(a) $\tau^-\tau^+(\rightarrow\mu^-\mu^+\nu\nu\bar{\nu}\bar{\nu})$

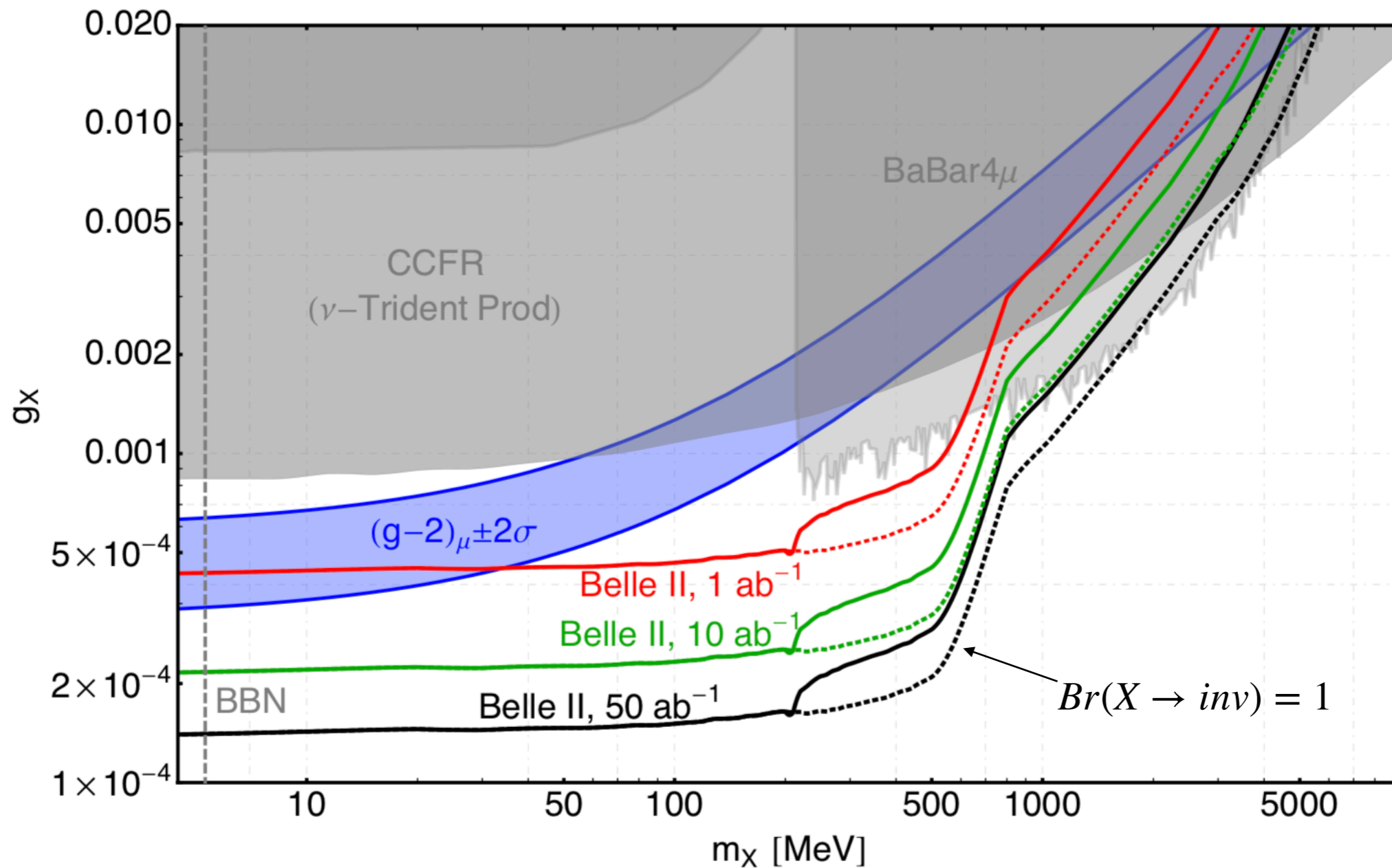


(b) $\mu^-\mu^+X$, $m_X = 800$ MeV



(c) $\mu^-\mu^+X$, $m_X = 3$ GeV

Sensitivity bound: $S/\sqrt{S+B} = 3$ with & without dark matter

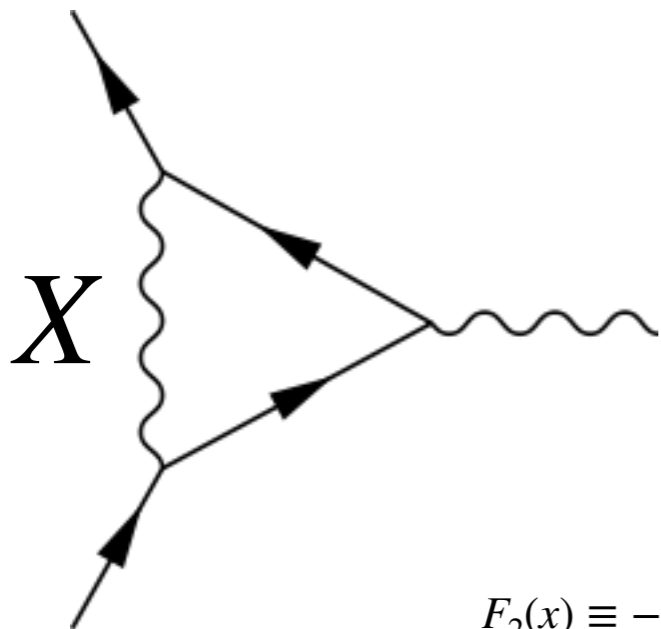


Summary

- The gauged $L_\mu - L_\tau$ can explain $(g - 2)_\mu$ anomaly with $g_X \sim 10^{-4}$ and $m_X \sim 10$ MeV.
- Belle-2 with 50/ab will completely cover the preferred range in $\mu^+ \mu^- + MET$ channel
- Possible extension with dark matter can be also interesting
- Stay tuned!

Back up

$Z - \mu\bar{\mu}$



$$\frac{\Delta\Gamma(Z \rightarrow \mu^- \mu^+)}{\Gamma(Z \rightarrow \mu^- \mu^+)} = \frac{g_X^2}{16\pi^2} F_2 \left(\frac{m_X^2}{m_Z^2} \right)$$

$$F_2(x) \equiv -2 \left\{ \frac{7}{4} + x + \left(x + \frac{3}{2} \right) \ln x + (1+x)^2 \left[\text{Li}_2 \left(\frac{x}{1+x} \right) + \frac{1}{2} \ln^2 \left(\frac{x}{1+x} \right) - \frac{\pi^2}{6} \right] \right\}$$

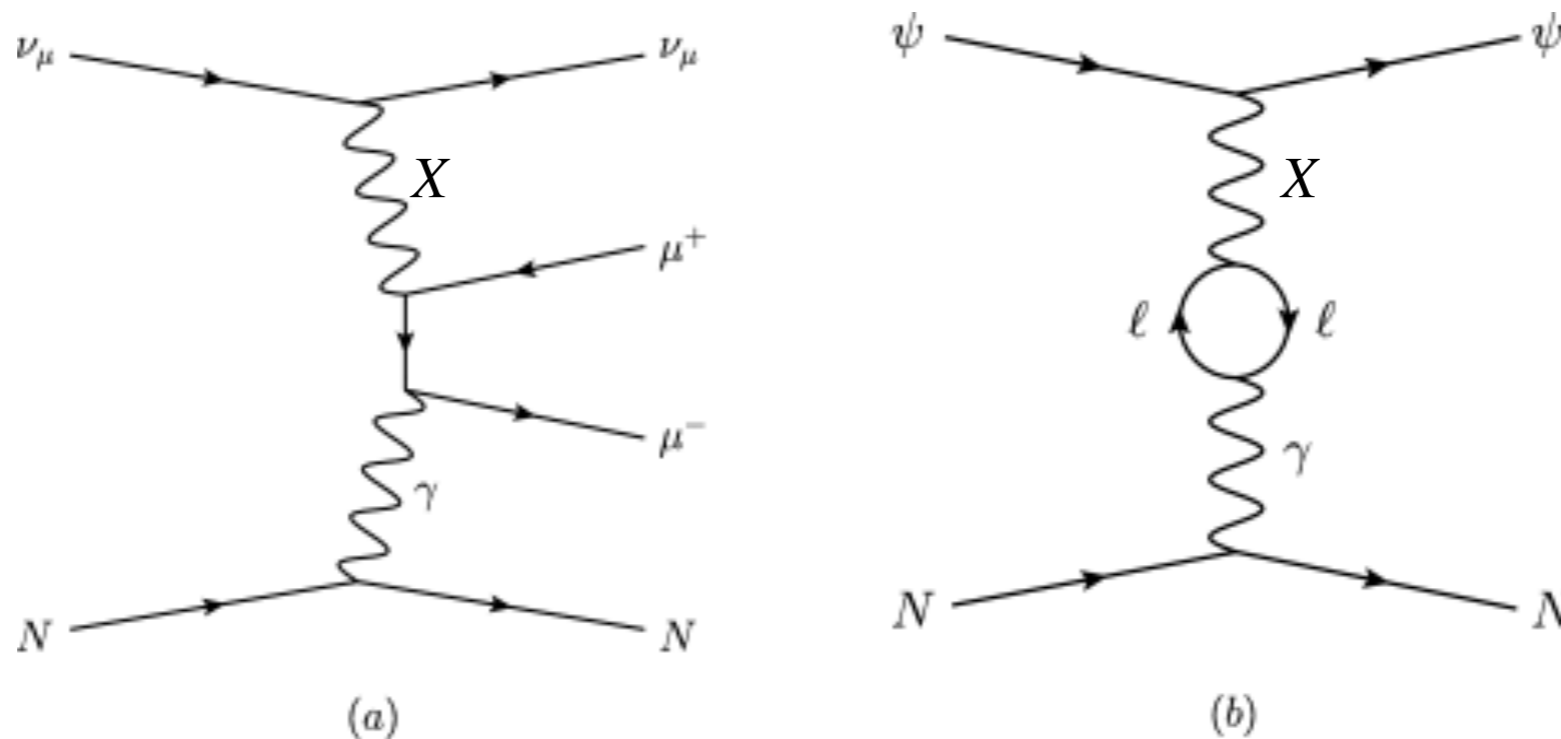
$$\left| \frac{g_X^2}{16\pi^2} F_2 \left(\frac{m_X^2}{m_Z^2} \right) \right| < \left| \frac{\Gamma(Z \rightarrow \mu^- \mu^+)}{\Gamma(Z \rightarrow e^- e^+)} - 1 \right|$$

$$\mathbf{Br}(Z \rightarrow e^- e^+) = 3.3632 \pm 0.0042 \%,$$

$$\mathbf{Br}(Z \rightarrow \mu^- \mu^+) = 3.3662 \pm 0.0066 \%$$

Neutrino trident

$(\nu N \rightarrow \nu N \mu^+ \mu^-)$



$$\sigma_{\text{CCFR}} / \sigma_{\text{SM}} = 0.82 \pm 0.28 \quad E_{\mu} = 160 \text{ GeV}$$

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