

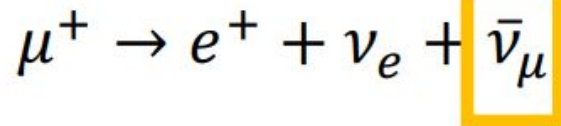
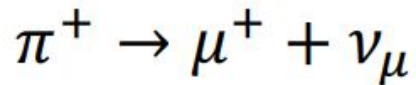
# MiniBooNE results and bounds on their non-oscillations explanations

**Carlos A. Argüelles**

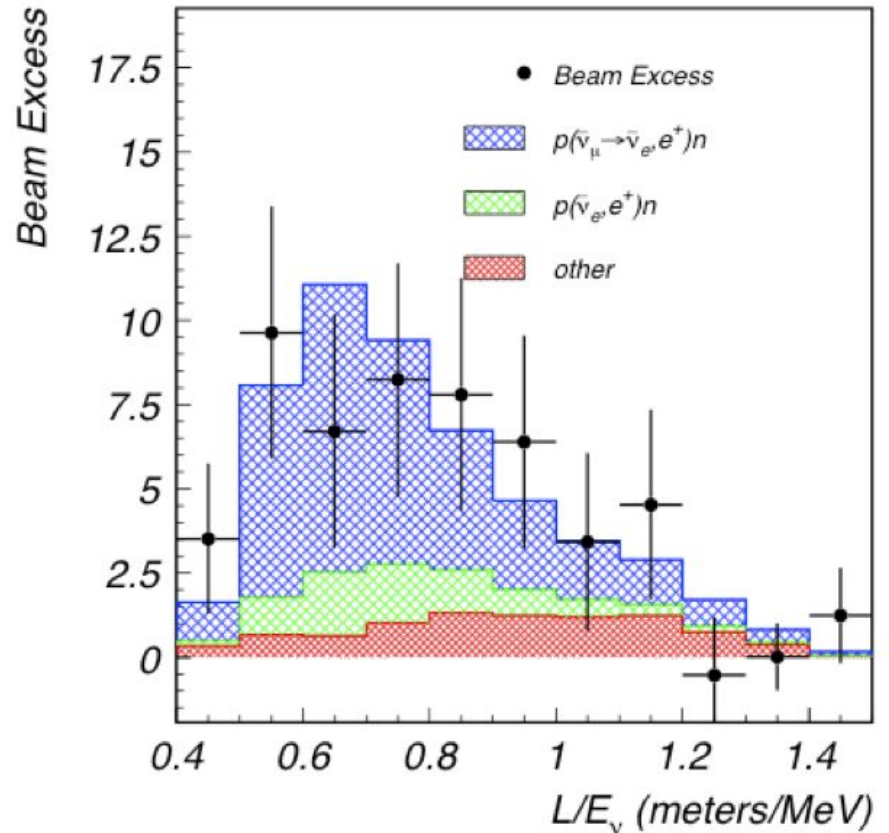
Part of this talk is based on [arXiv:1812.08768](https://arxiv.org/abs/1812.08768)

# The back story: LSND

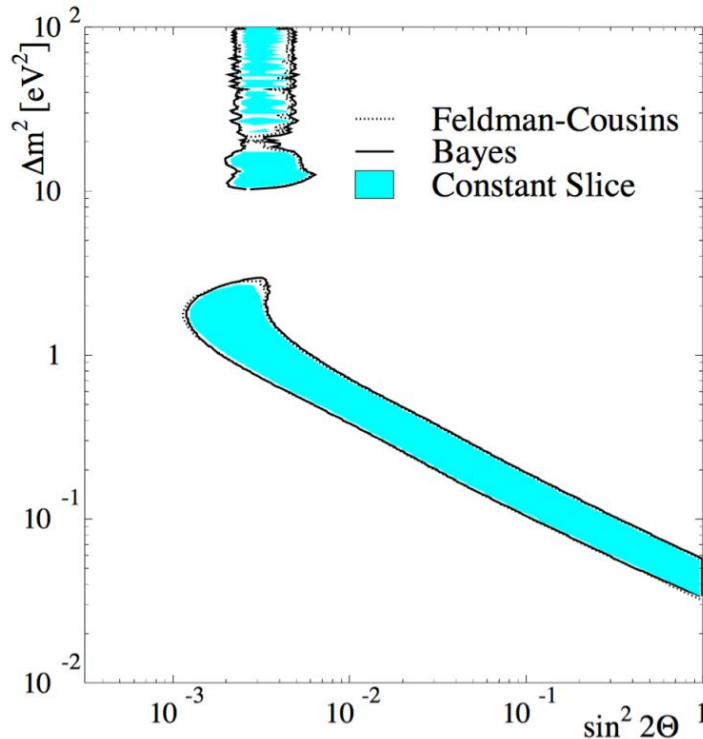
- LSND used an 800 MeV proton beam to produce pions which decay at rest.



- Search for antineutrino-electron appearance
- Observed 3.8 sigma excess of events



# If interpreted as an appearance probability ...



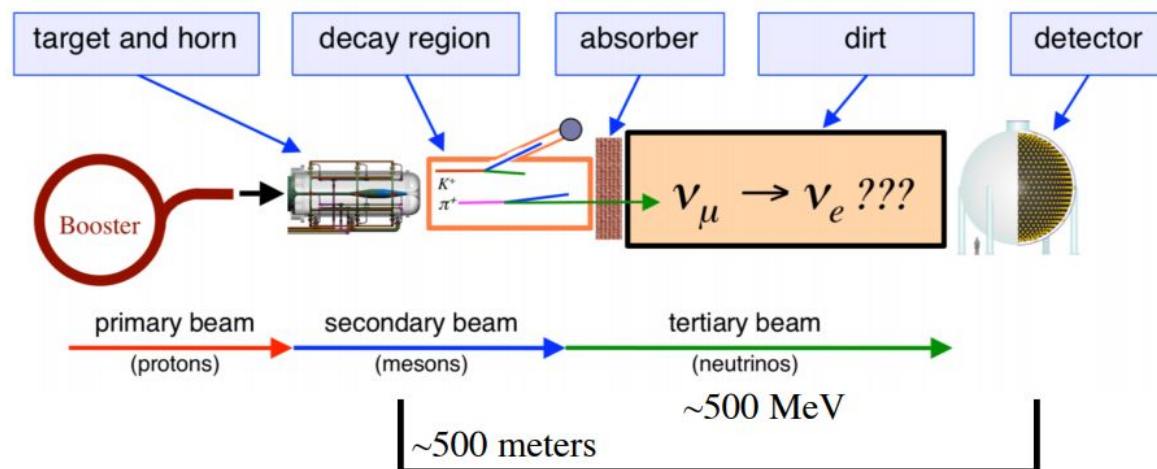
$$P(\nu_\mu \rightarrow \nu_e) = \sin^2(2\theta) \sin^2\left(1.27 \frac{\Delta m^2 L}{E_\nu}\right)$$

- The preferred parameter space of the LSND anomaly was not compatible with other known mass differences.
- If this is due to a new neutrino mass state, then we should observe a similar signal at different E and L, but same L/E!

A. Aguilar-Arevalo et al. [LSND Collaboration]  
Phys. Rev. D 64, 112007 (2001) [hep-ex/0104049].

# MiniBooNE@FNAL: proposed to test the LSND anomaly

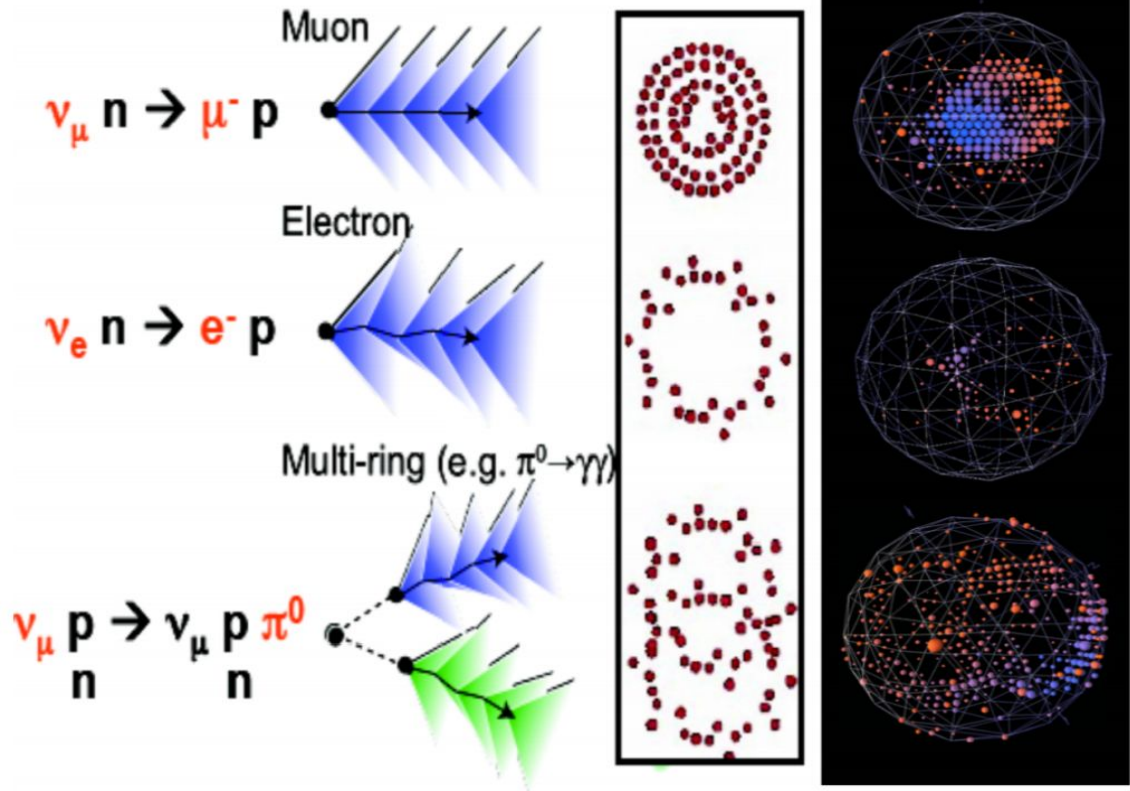
- Approximately same L/E, but  $\sim 15$  x larger energy and baseline.
- Decay-in-flight pion source.
- Higher backgrounds than LSND, but more statistics!
- Neutrino and antineutrino mode available.



# MiniBooNE experimental signatures

Three typical event signatures:

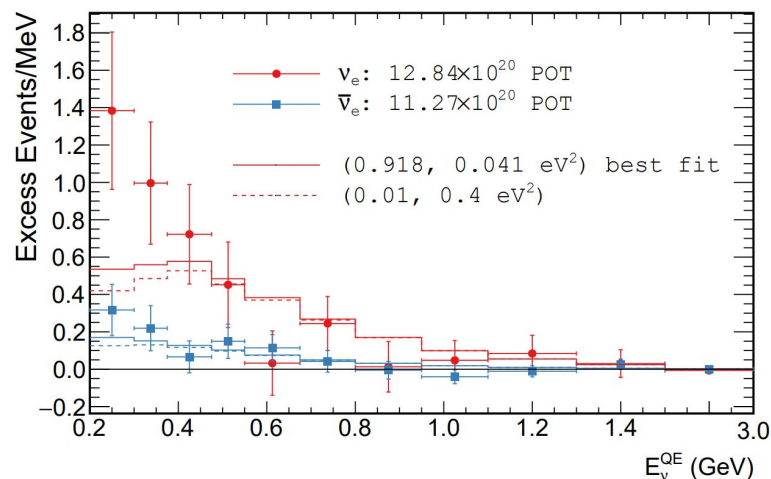
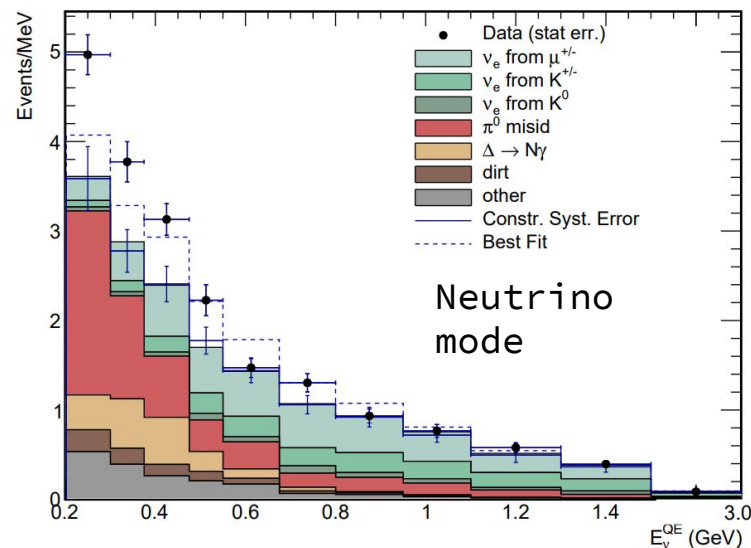
- Muon-neutrino CCQE produces sharp photon ring on PMTs,
- Electron-neutrino CCQE events produces fuzzy ring,
- Muon-neutrino NC can produce  $\pi^0$ : two gammas  $\rightarrow$  two fuzzy rings.



Cannot distinguish between electrons and photons!

# Recent MiniBooNE excess of neutrino-electron-like events

- MB has reported an excess in neutrino and antineutrino channels.
- They claim that this excess is compatible with LSND. (Previous results show tension with LSND in neutrino mode)
- Excess has remained after doubling the data: not statistical in nature.
- It has a significance of 4.7 sigma.



# This raises some more questions...

---

Do we understand all SM background/process well enough?

Are all the anomalies related? Or only some of them?

E.g., are LSND and MiniBooNE observing the same physics?

Since null results are not scrutinized as carefully as anomalous ones. Are all null results reliable?

Is there a significant signal of electron-neutrino disappearance (e.g. reactors)?

If the anomalies are confirmed as new physics, in what theories are they embedded?

# BSM-ways to approach the problem

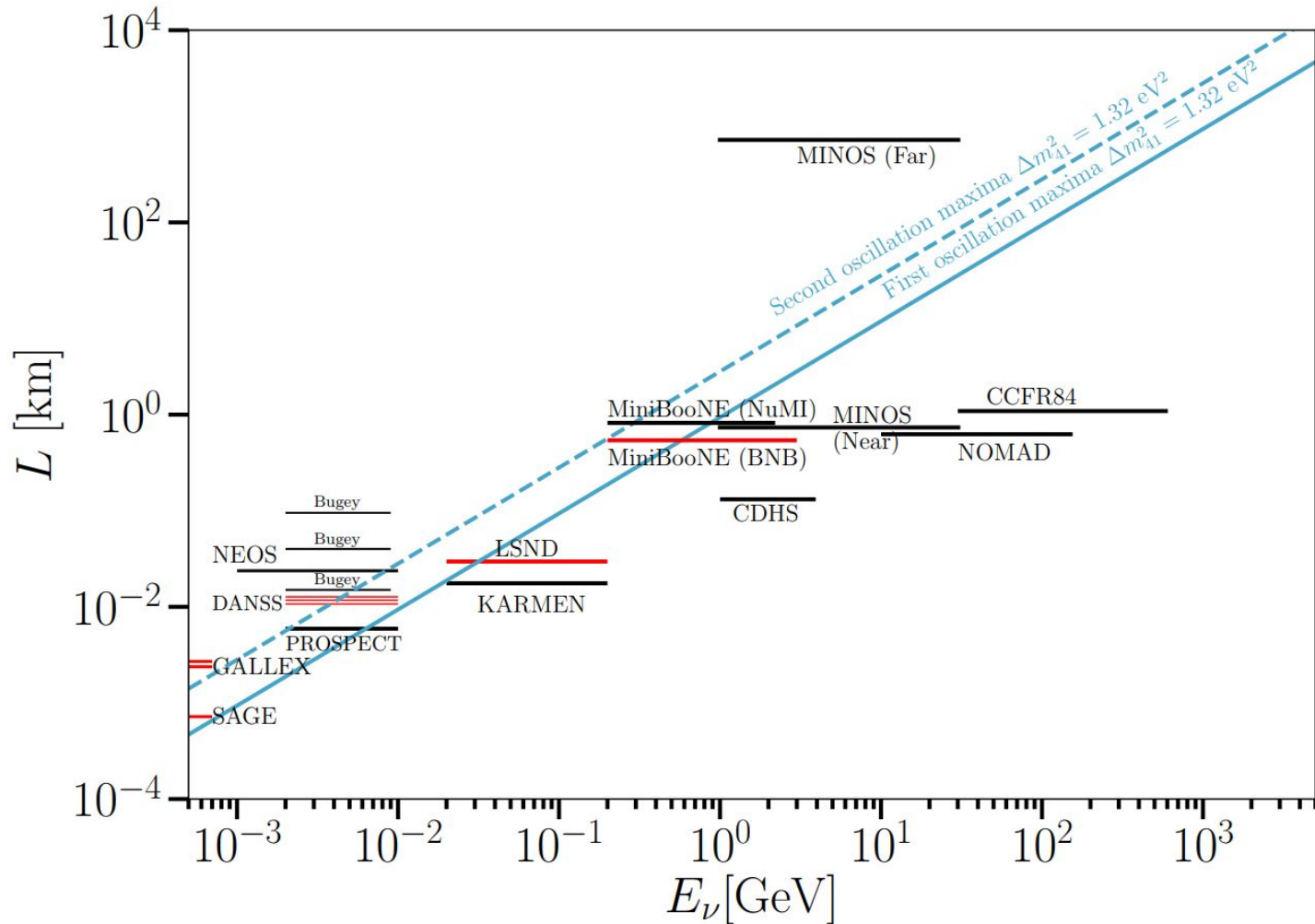
---

1. Increase the muon-neutrino to electron neutrino conversation probability (e.g. via a sterile neutrino -- Carlo's talk yesterday).

Strategy: look for signals at similar L/E.



# Sterile neutrino panorama



A. Diaz et al. 1906.00045 see also Boser et al. 1906.01739<sub>9</sub>

So ... We have discover a new particle!?

---



# If it's a “vanilla” eV-scale sterile neutrino

---

$$P_{\nu_e \rightarrow \nu_e} = 1 - 4(1 - |U_{e4}|^2)|U_{e4}|^2 \sin^2(1.27\Delta m_{41}^2 L/E)$$

$$P_{\nu_\mu \rightarrow \nu_e} = 4|U_{e4}|^2|U_{\mu4}|^2 \sin^2(1.27\Delta m_{41}^2 L/E)$$

$$P_{\nu_\mu \rightarrow \nu_\mu} = 1 - 4(1 - |U_{\mu4}|^2)|U_{\mu4}|^2 \sin^2(1.27\Delta m_{41}^2 L/E)$$

Oscillation probabilities among appearance and disappearance channels are related.

**Need to look in other channels for further confirmation!**

So ... We have discover a new particle!?

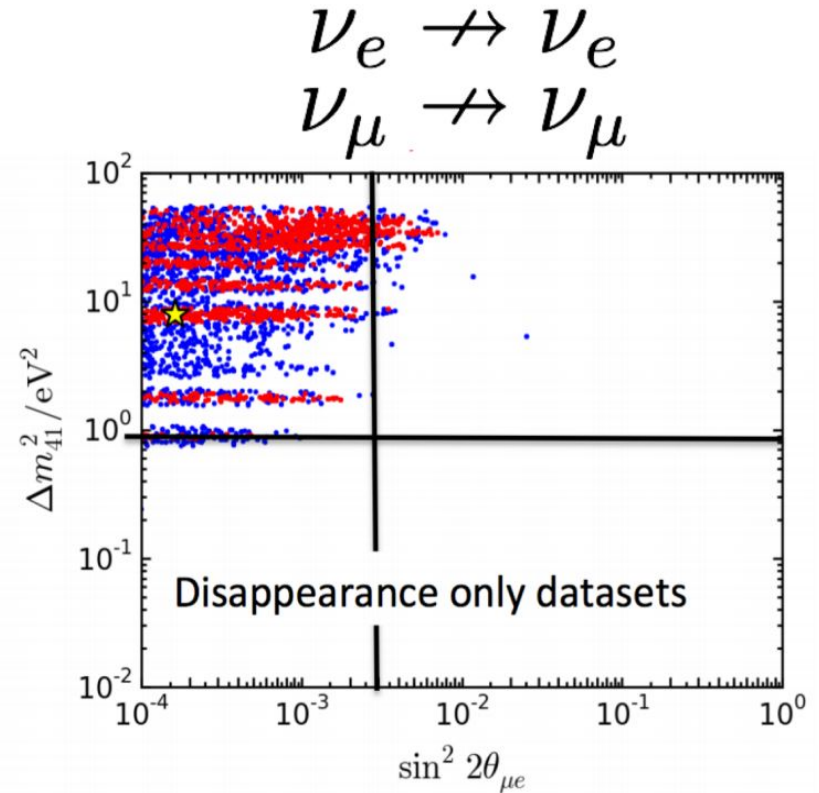
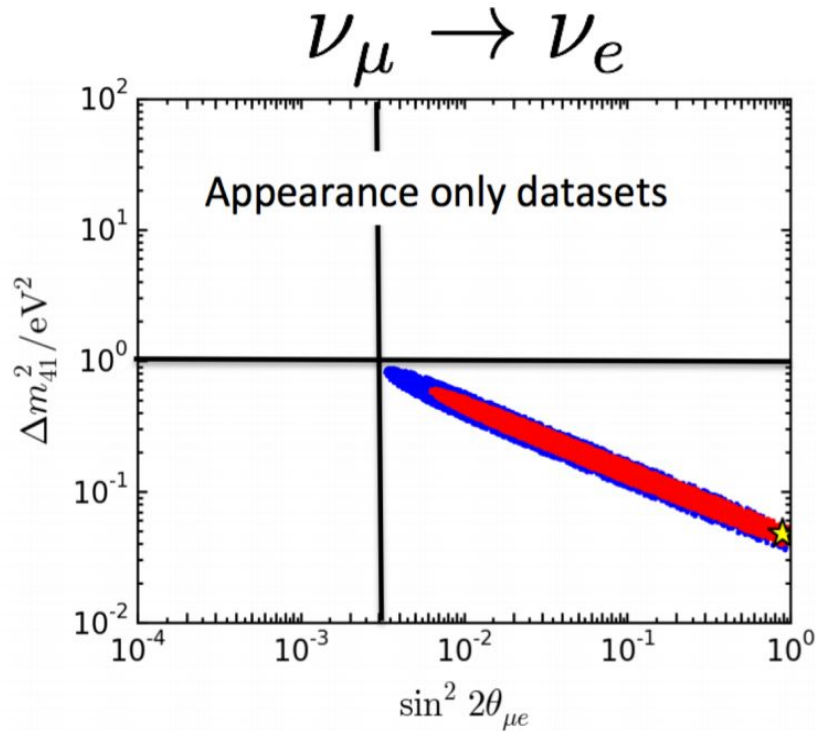
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\* R.I.P. grumpy cat.

# App and Dis preference regions don't match!

---



From Collin et al 1602.00671, similar conclusions from other groups see Gariazzo et al. 1703.00860, and Dentler et al JHEP 1808 (2018)

# Tension in the global data!!!

PG: parameter goodness-of-fit.  
Larger is better. Small is bad.  
Very small is very bad.

Analysis	$\chi^2_{\min,\text{global}}$	$\chi^2_{\min,\text{app}}$	$\Delta\chi^2_{\text{app}}$	$\chi^2_{\min,\text{disapp}}$	$\Delta\chi^2_{\text{disapp}}$	$\chi^2_{\text{PG}}/\text{dof}$	PG
Global	1120.9	79.1	11.9	1012.2	17.7	29.6/2	$3.71 \times 10^{-7}$
<b>Removing anomalous data sets</b>							
w/o LSND	1099.2	86.8	12.8	1012.2	0.1	12.9/2	$1.6 \times 10^{-3}$
w/o MiniBooNE	1012.2	40.7	8.3	947.2	16.1	24.4/2	$5.2 \times 10^{-6}$
w/o reactors	925.1	79.1	12.2	833.8	8.1	20.3/2	$3.8 \times 10^{-5}$
w/o gallium	1116.0	79.1	13.8	1003.1	20.1	33.9/2	$4.4 \times 10^{-8}$
<b>Removing constraints</b>							
w/o IceCube	920.8	79.1	11.9	812.4	17.5	29.4/2	$4.2 \times 10^{-7}$
w/o MINOS(+)	1052.1	79.1	15.6	948.6	8.94	24.5/2	$4.7 \times 10^{-6}$
w/o MB disapp	1054.9	79.1	14.7	947.2	13.9	28.7/2	$6.0 \times 10^{-7}$
w/o CDHS	1104.8	79.1	11.9	997.5	16.3	28.2/2	$7.5 \times 10^{-7}$
<b>Removing classes of data</b>							
$\bar{\nu}_e$ dis vs app	628.6	79.1	0.8	542.9	5.8	6.6/2	$3.6 \times 10^{-2}$
$\bar{\nu}_\mu$ dis vs app	564.7	79.1	12.0	468.9	4.7	16.7/2	$2.3 \times 10^{-4}$
$\bar{\nu}_\mu$ dis + solar vs app	884.4	79.1	13.9	781.7	9.7	23.6/2	$7.4 \times 10^{-6}$

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**A vanilla 3+1 sterile neutrino fails to explain all the data!**

# BSM-ways to approach the problem

---

1. Increase the muon-neutrino to electron neutrino conversation probability (e.g. via a sterile neutrino -- Carlo's talk yesterday).

Strategy: look for signals at similar L/E.

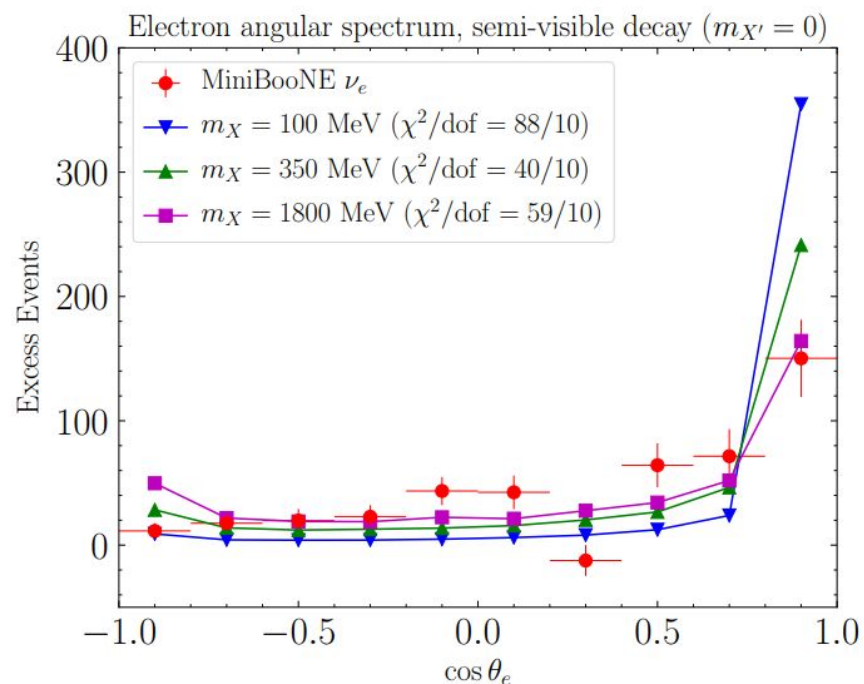
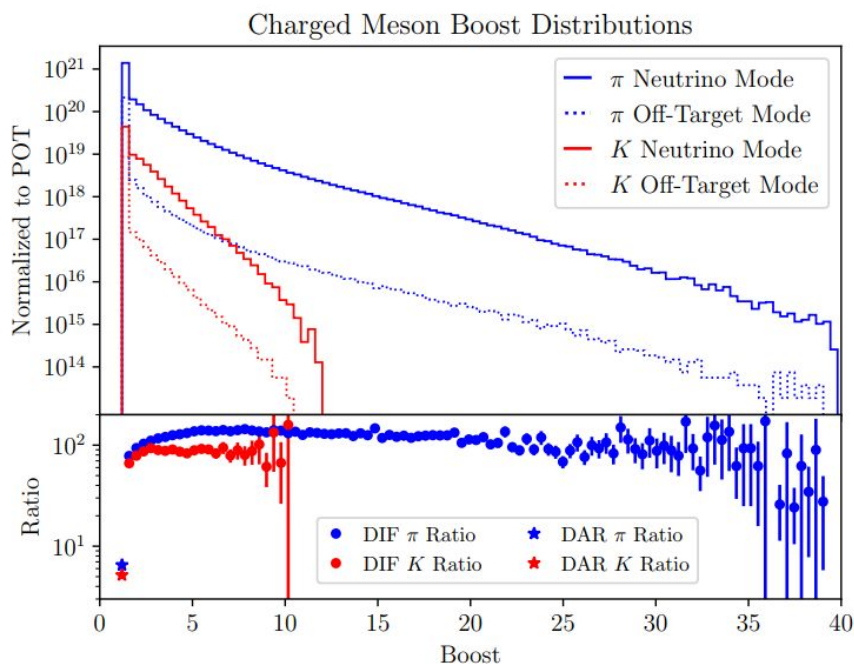
2. Production of new particles in the target that produce signatures in the detector directly or by their decay products.

Strategy: change the beam configuration, see if signal remains.



# Consider null results from MB beam-dump+kinematics

Jordan et al. arXiv:1810.07185



No observation of an excess in MB beam dump configuration and broad angular distribution of the excess disfavors scenarios in which the particle produced in the target:

1. either decays (visibly or semi-visibly) or
2. scatters elastically

# BSM-ways to approach the problem

---

1. Increase the muon-neutrino to electron neutrino conversation probability (e.g. via a sterile neutrino -- Carlo's talk yesterday).

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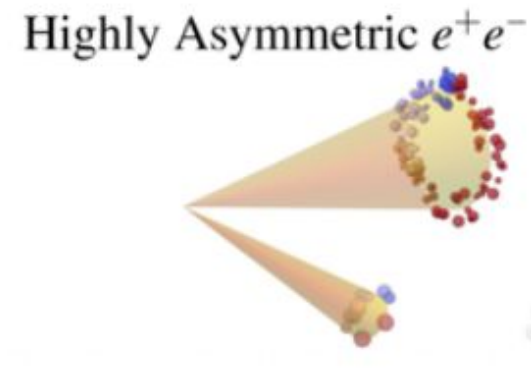
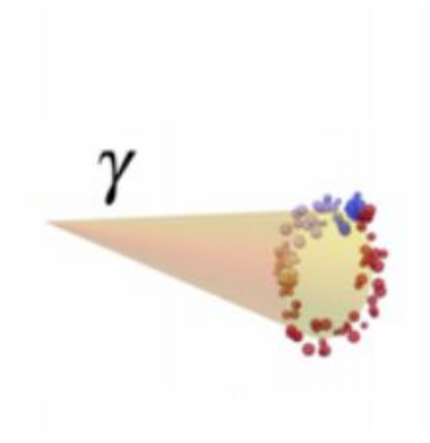
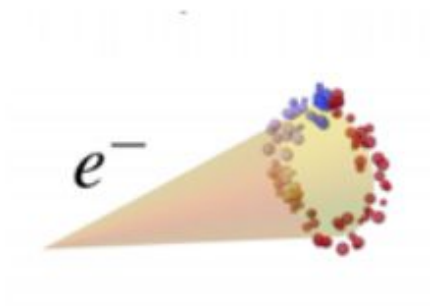
3. New neutrino interactions that mimic the signature process

Strategy: this talk!

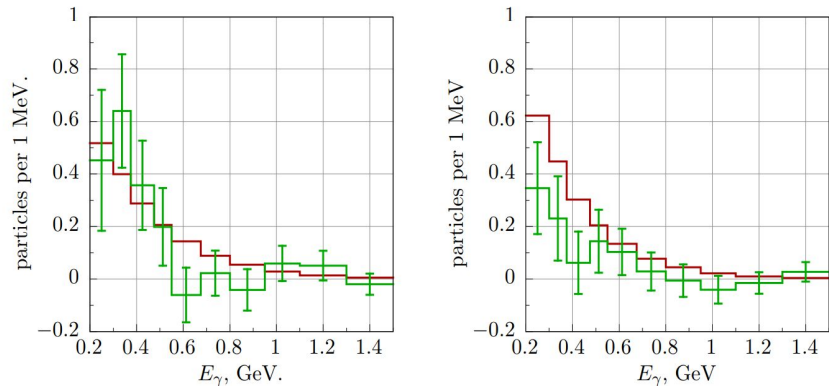
# Making the MB signature

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If not statistical fluctuation, then (**probably**) one of the options below:



# Old transition magnetic moment explanation



see, e.g.,

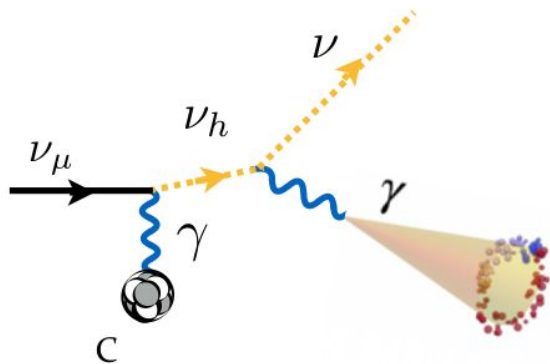
*S.N. Gninenko, 0902.3802*

*S.N. Gninenko, 1009.5536*

*M. Masip et al, 1210.1519*

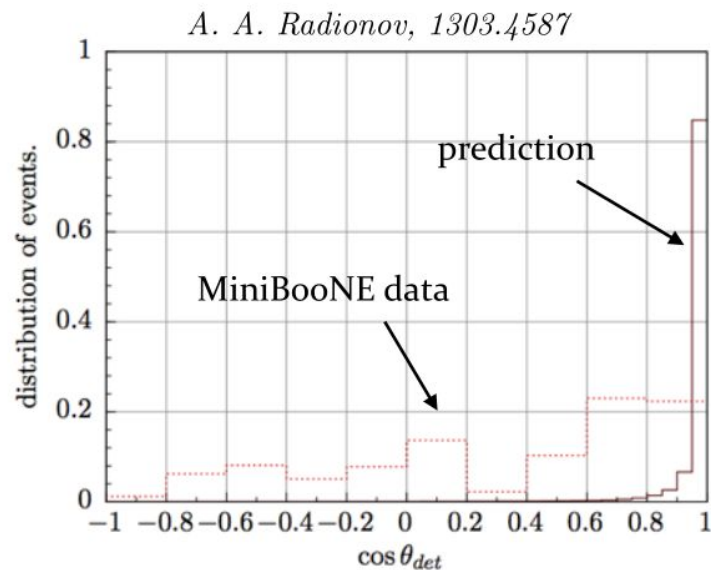
**transition magnetic moment**

models are very similar to light dark photon case.



$$\mathcal{L}_{int} = \frac{1}{2} \sum_{\substack{i,j=1 \\ i \leq j}}^{N_\nu} \bar{\nu}_j (\mu_{ij} + \gamma_5 d_{ij}) \sigma_{\mu\nu} \mathcal{F}^{\mu\nu} \nu_i + \text{h.c.}$$

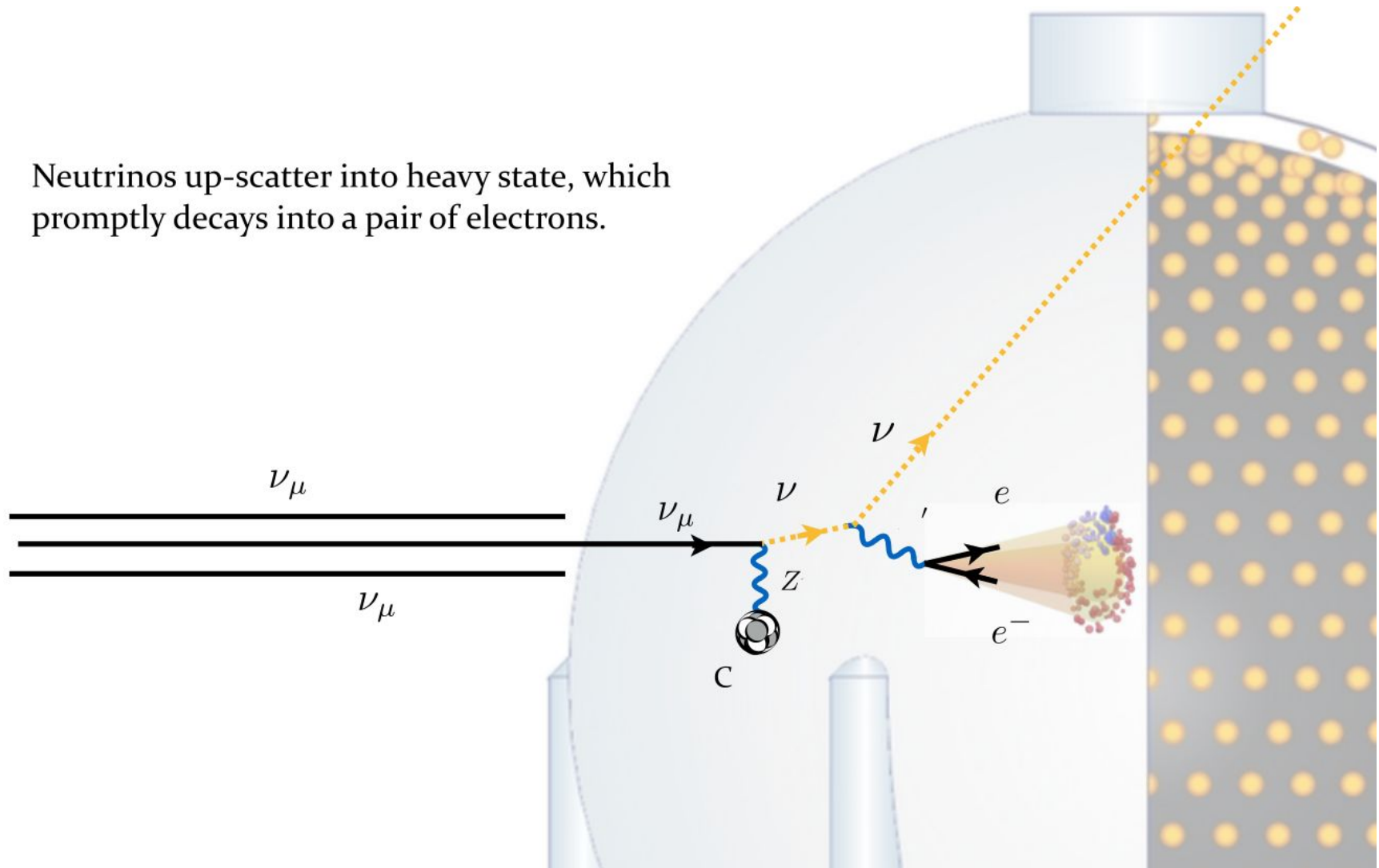
Signal is too forward!



# Novel explanations of the MiniBooNE anomaly

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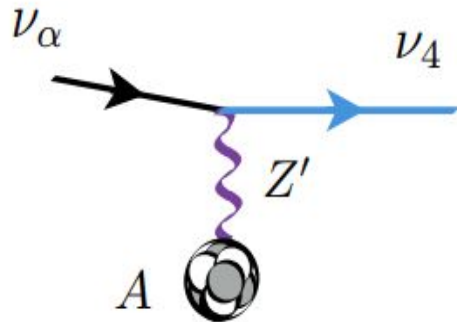
Neutrinos up-scatter into heavy state, which promptly decays into a pair of electrons.



# Producing the MiniBooNE signature: two implementations

---

$$\alpha_D \alpha_{\text{QED}} \varepsilon^2 |U_{\alpha 4}|^2$$



Small  $Mz'$ :

- Larger cross section
- Small  $Q^2$ : more forward  $\nu_4$ , larger coherent to diffractive contributions

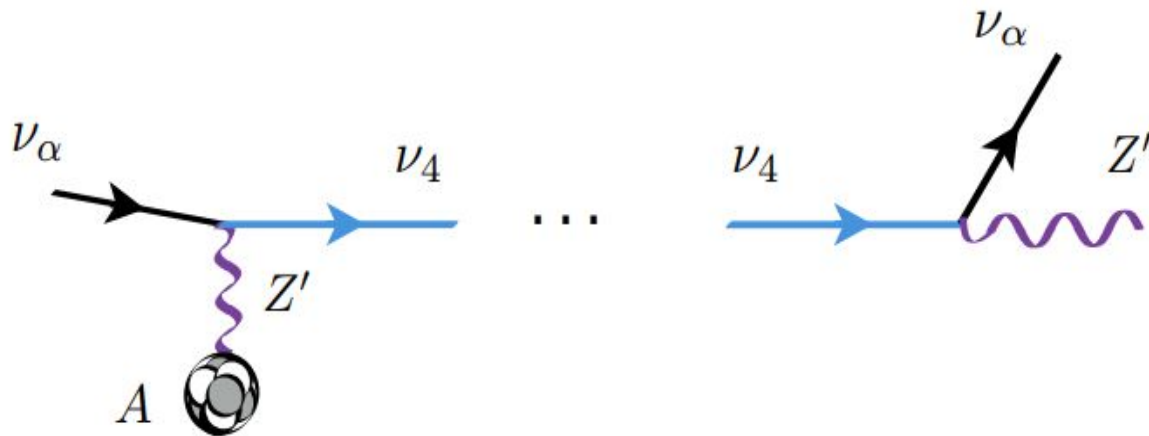
Large  $Mz'$ :

- Smaller cross section
- Larger  $Q^2$ : more isotropic  $\nu_4$  production, more diffractive contribution

# Producing the MiniBooNE signature

---

**Model from Bertuzzo et al.  
Phys. Rev. Lett. 121, 241801 (2018)**

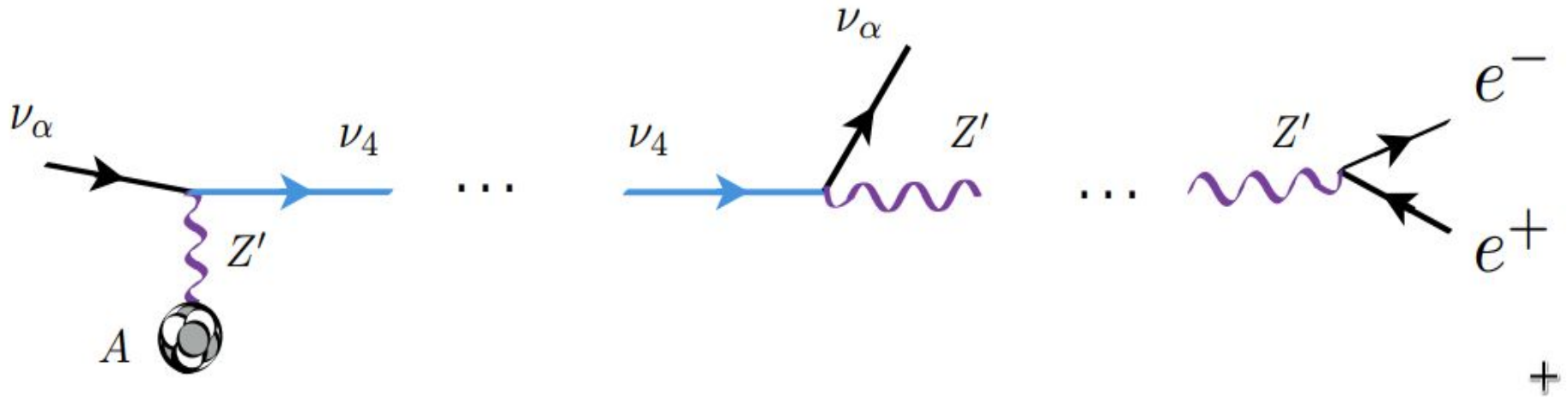


If  $M_4 > M_{Z'}$ , two body decay is the dominant decay channel.  
 $M_4 \gtrsim 100$  MeV so the decay products are not so boosted in order to reproduce angular distribution.

# Producing the MiniBooNE signature

---

**Model from Bertuzzo et al.  
Phys. Rev. Lett. 121, 241801 (2018)**

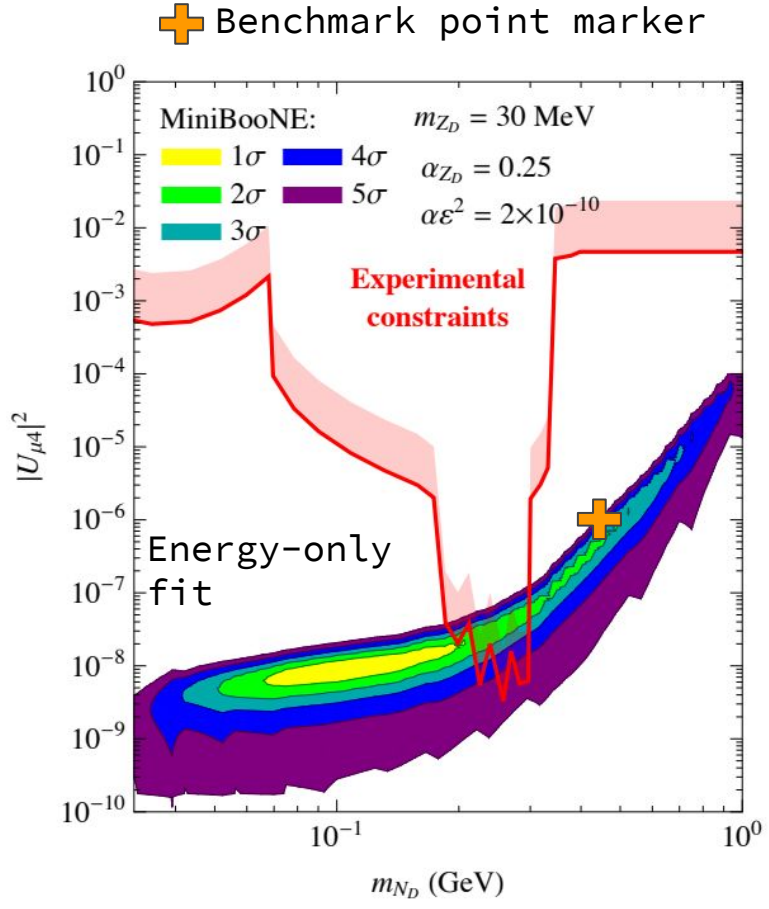
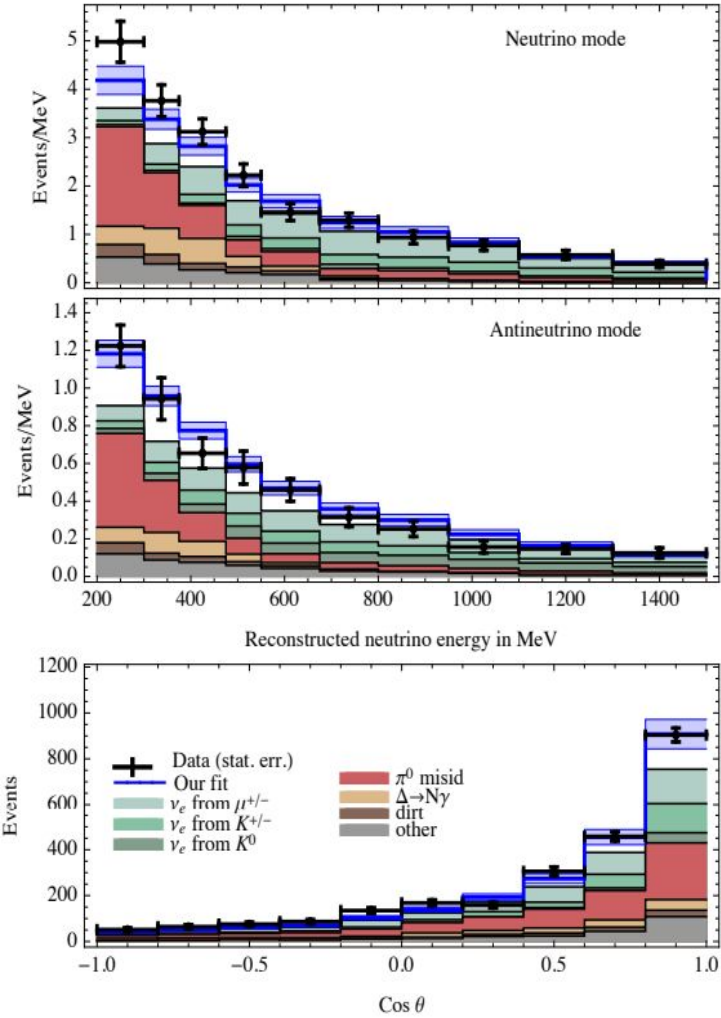


$M_{Z'}$  should be light ( $< 60$  MeV) so that the electron pair is collimated and can “fake” an electron-neutrino ring



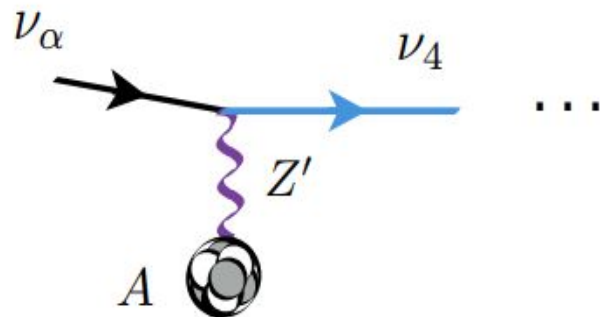
# Model by Bertuzzo et al. parameter space

Distributions at benchmark point

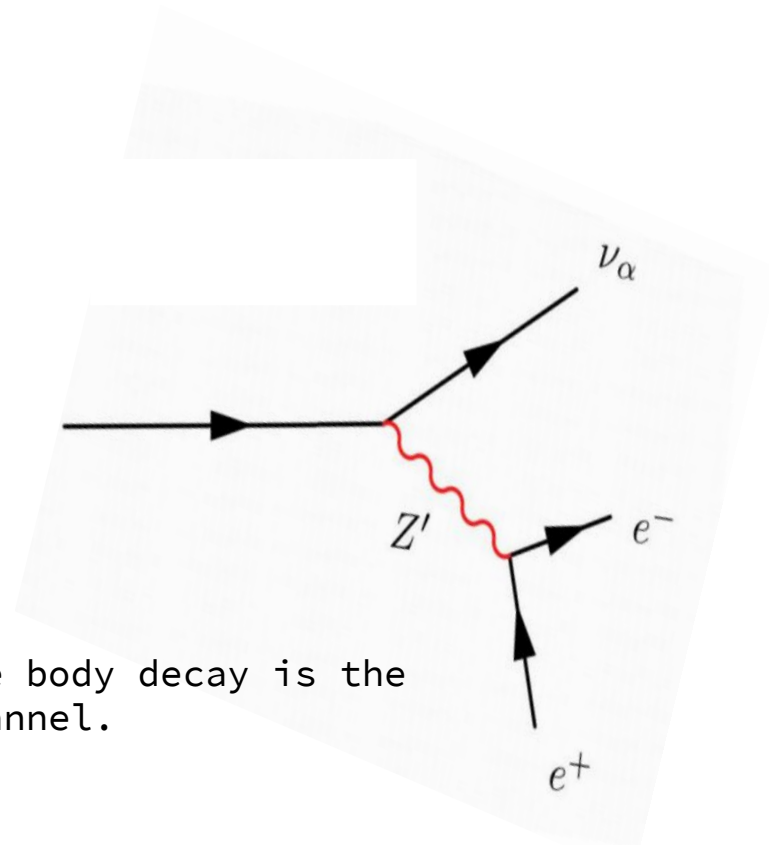


# Producing the MiniBooNE signature

Model from Ballett et al.  
arXiv:1808.02915; see also 1903.07589.



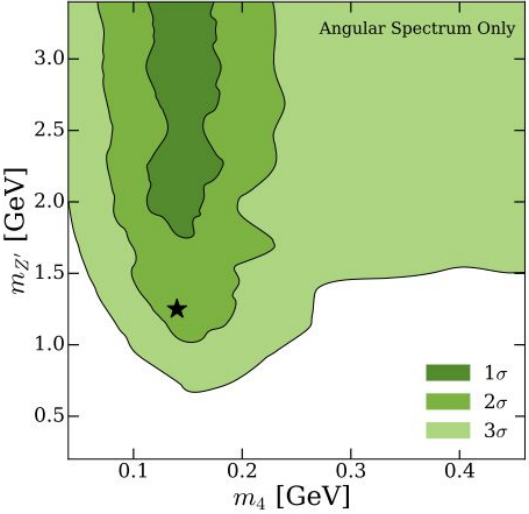
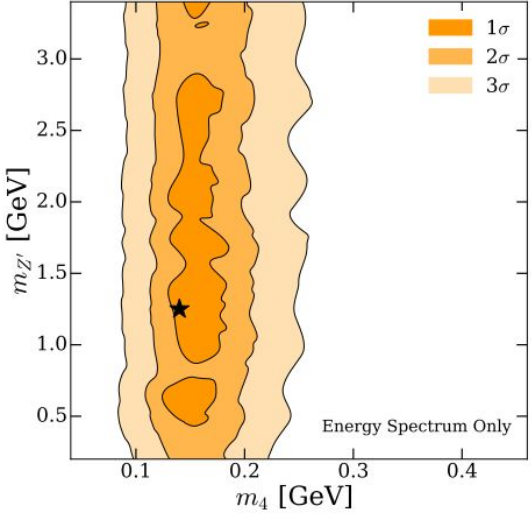
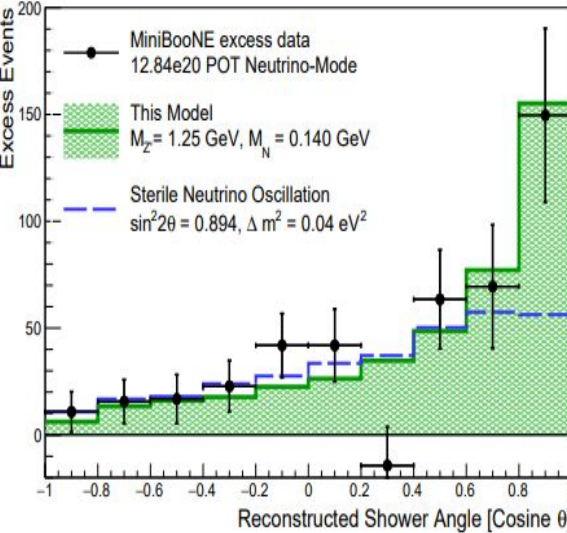
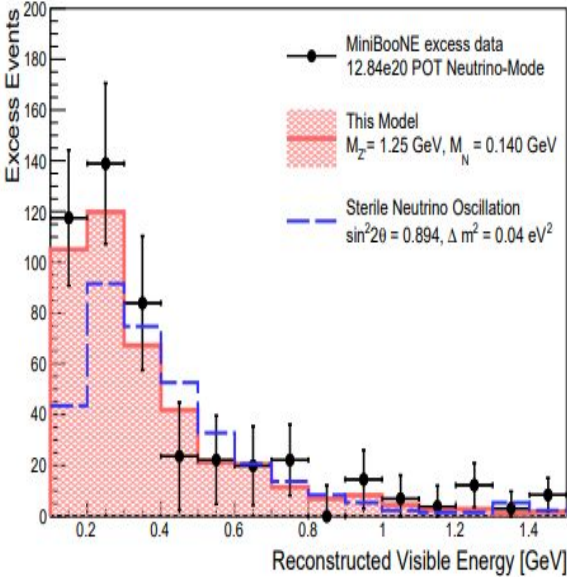
If  $M_4 < M_{Z'}$ , three body decay is the dominant decay channel.



$M_4 \sim > 300$  MeV: large too many high-energy events  
 $M_4 \sim < 50$  MeV most events in lowest energy bin.

$M_{Z'} \sim < 1$  GeV: spectrum is too forward.

# Model by Ballett et al. parameter space



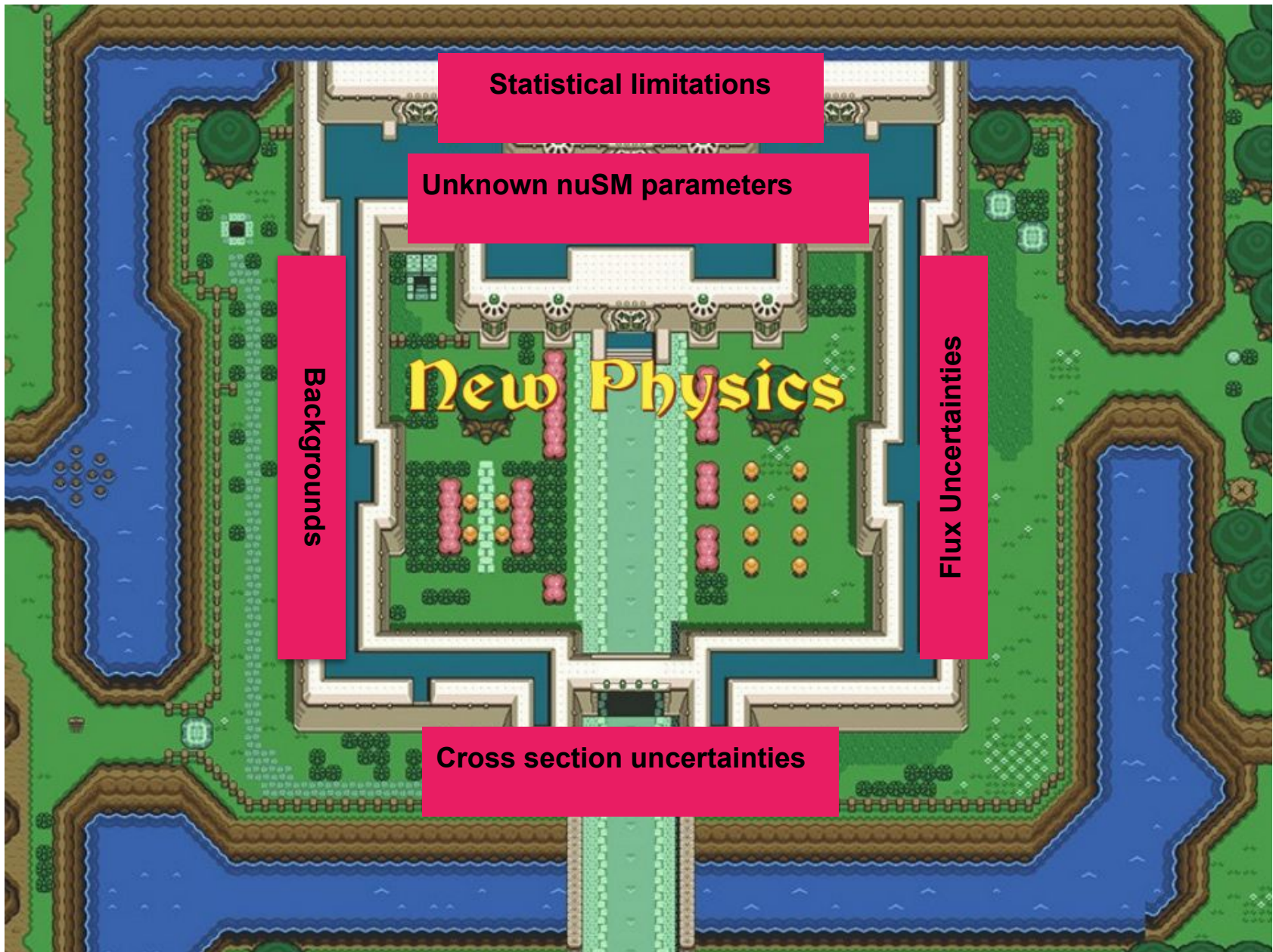


# Siege

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*noun*

a military operation in which enemy forces surround a town or building, cutting off essential supplies, with the aim of compelling the surrender of those inside.



**Statistical limitations**

**Unknown nuSM parameters**

**Backgrounds**

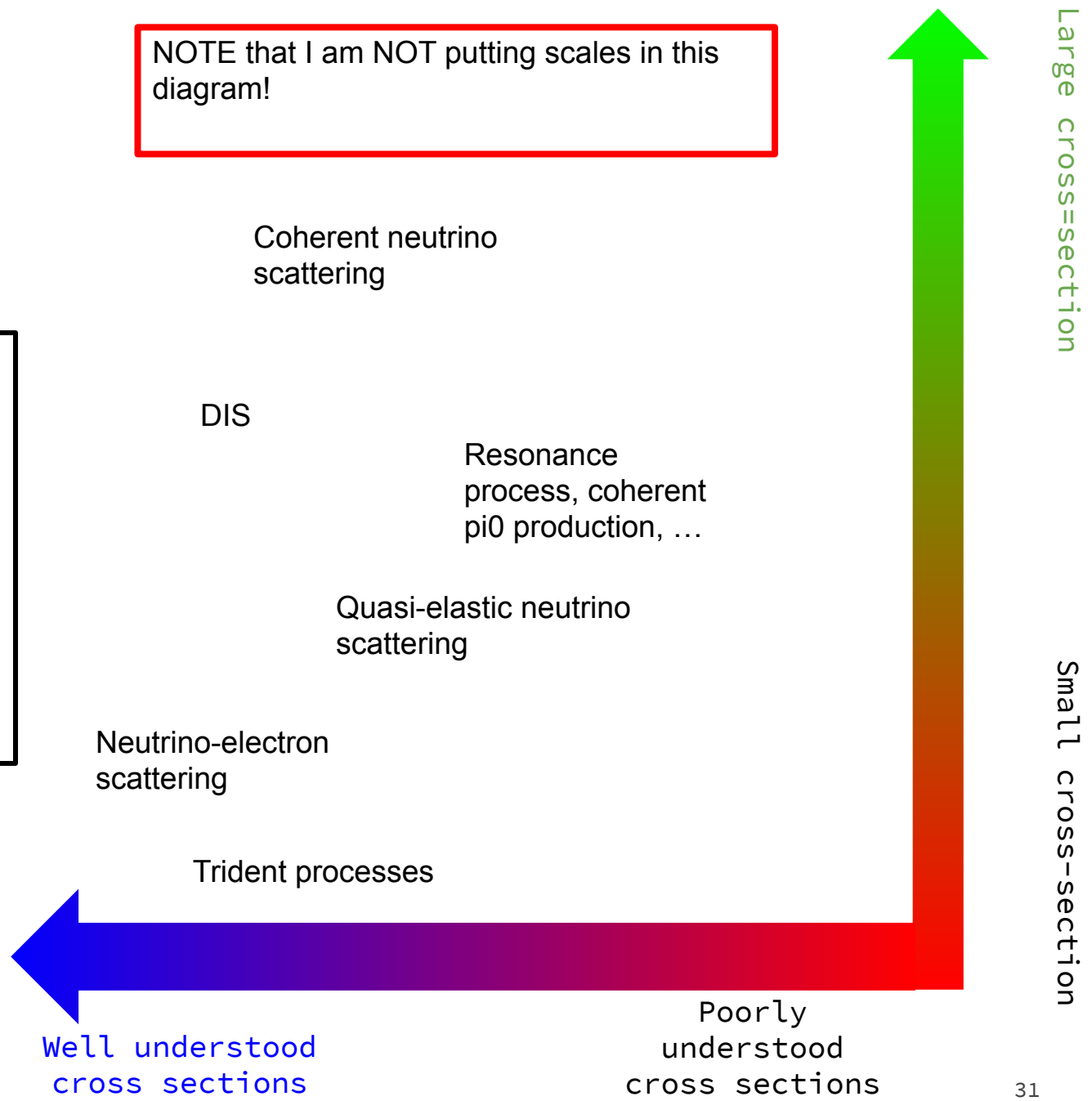
**Flux Uncertainties**

**Cross section uncertainties**

# Strategy

NOTE that I am NOT putting scales in this diagram!

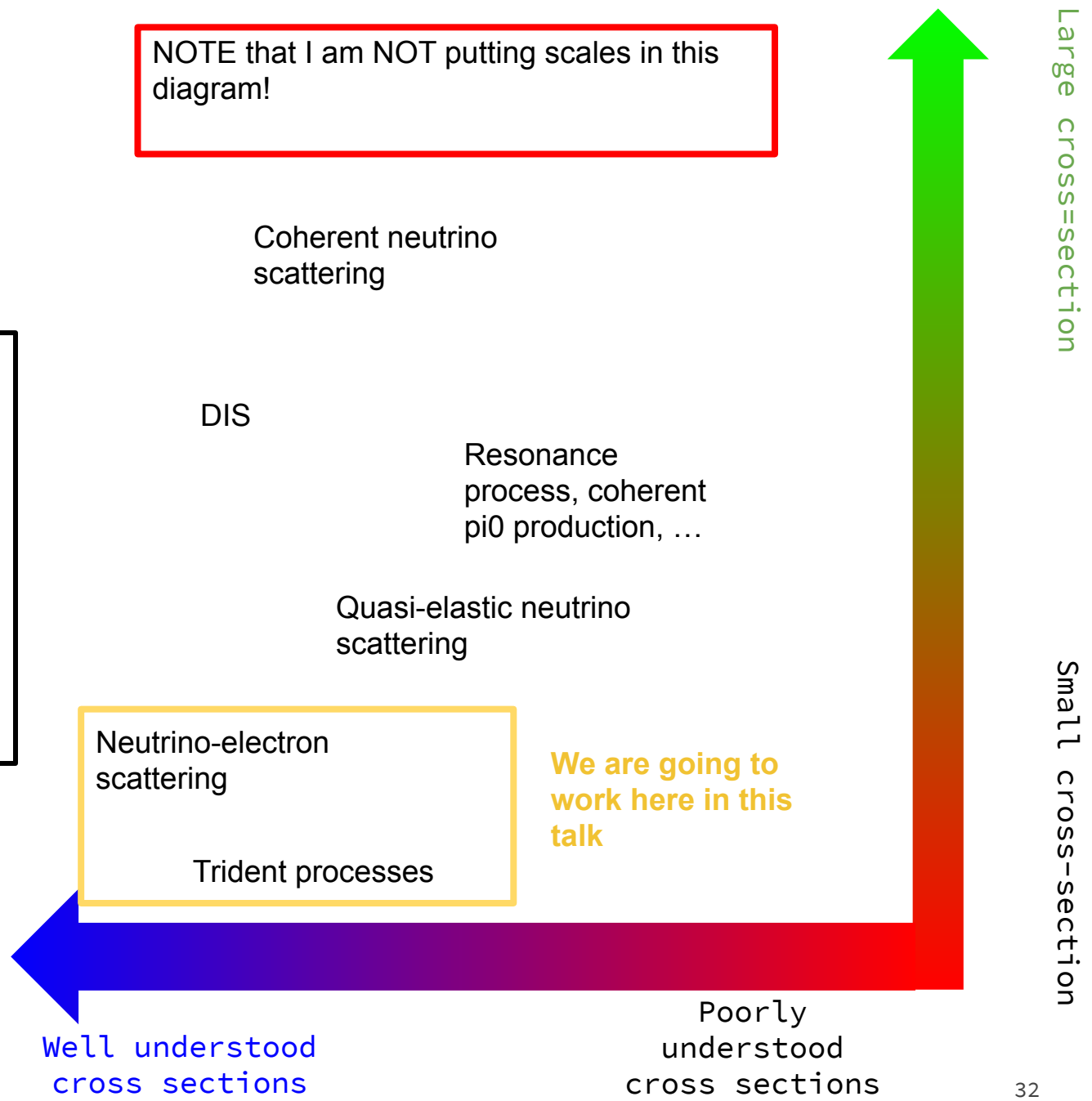
Looking for novel neutrino interactions requires understanding of our SM neutrino interactions. This is tough!



# Strategy

NOTE that I am NOT putting scales in this diagram!

Looking for novel neutrino interactions requires understanding of our SM neutrino interactions. This is tough!



Well understood cross sections

Poorly understood cross sections

Large cross=section

Small cross=section

We are going to work here in this talk



# What kind of signature are we hunting for?

---

- Production cross section is dominated by coherent processes; i.e. little/no hadronic activity at vertex.
- Products angular distribution is broad at Booster beam energies, but less so at higher energy beams.
- Electron pair produced by  $Z'$  decay is very collimated.

**How can we confirm or rule out a model like this?**

# What kind of signature are we hunting for?

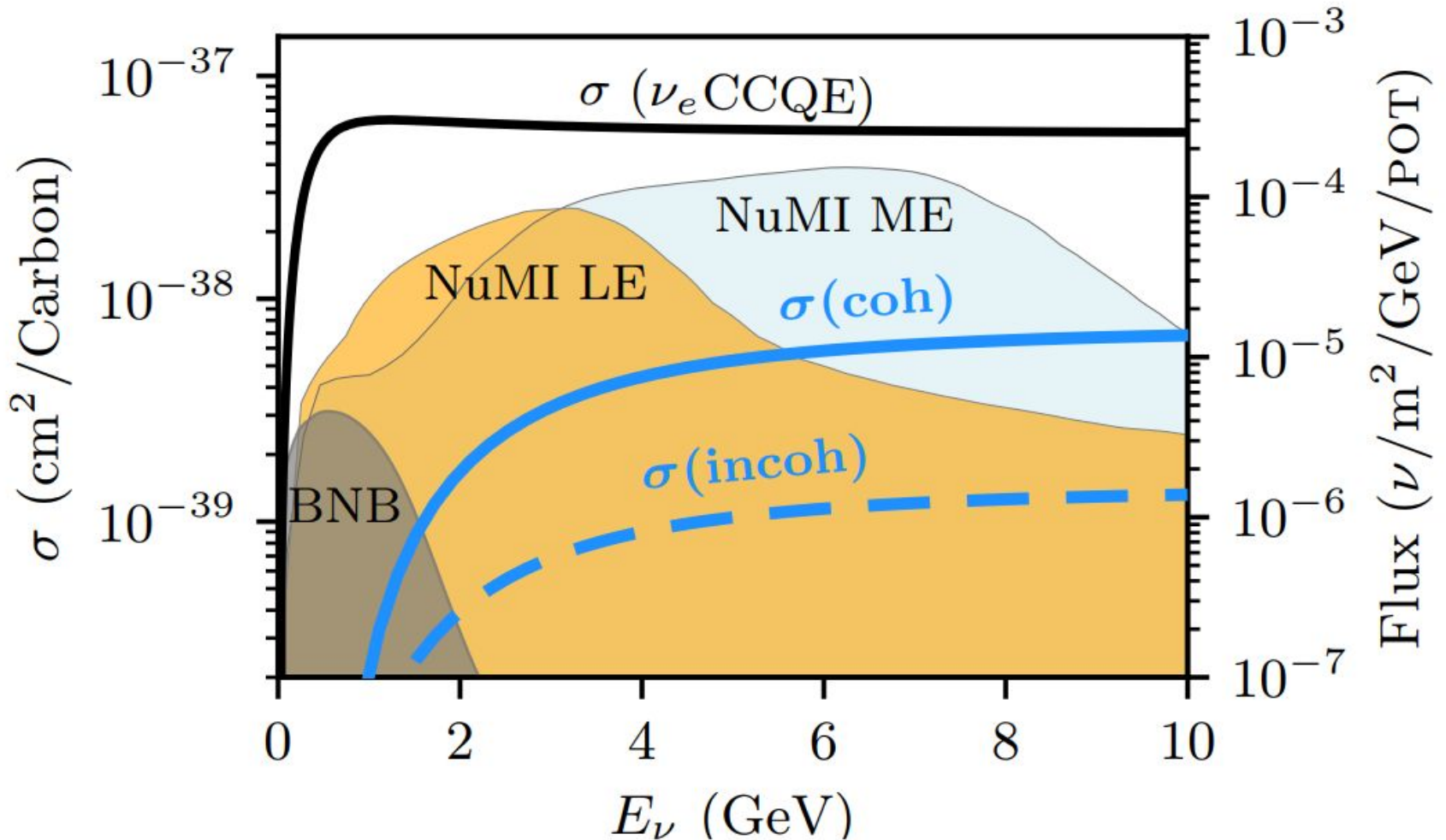
---

- Production cross section is dominated by coherent processes; i.e. little/no hadronic activity at vertex.
- Products angular distribution is broad at Booster beam energies, but less so at higher energy beams.
- Electron pair produced by  $Z'$  decay is very collimated.

## How can we confirm or rule out a model like this?

We are going to use neutrino scattering data to look for evidence of this process

# How big of a cross section are we talking about?



Here we use benchmark (BP) point parameters reported by Bertuzzo et al. 1807.09877.

# Neutrino-electron scattering measurements

---

We have measured neutrino-electron scattering @:

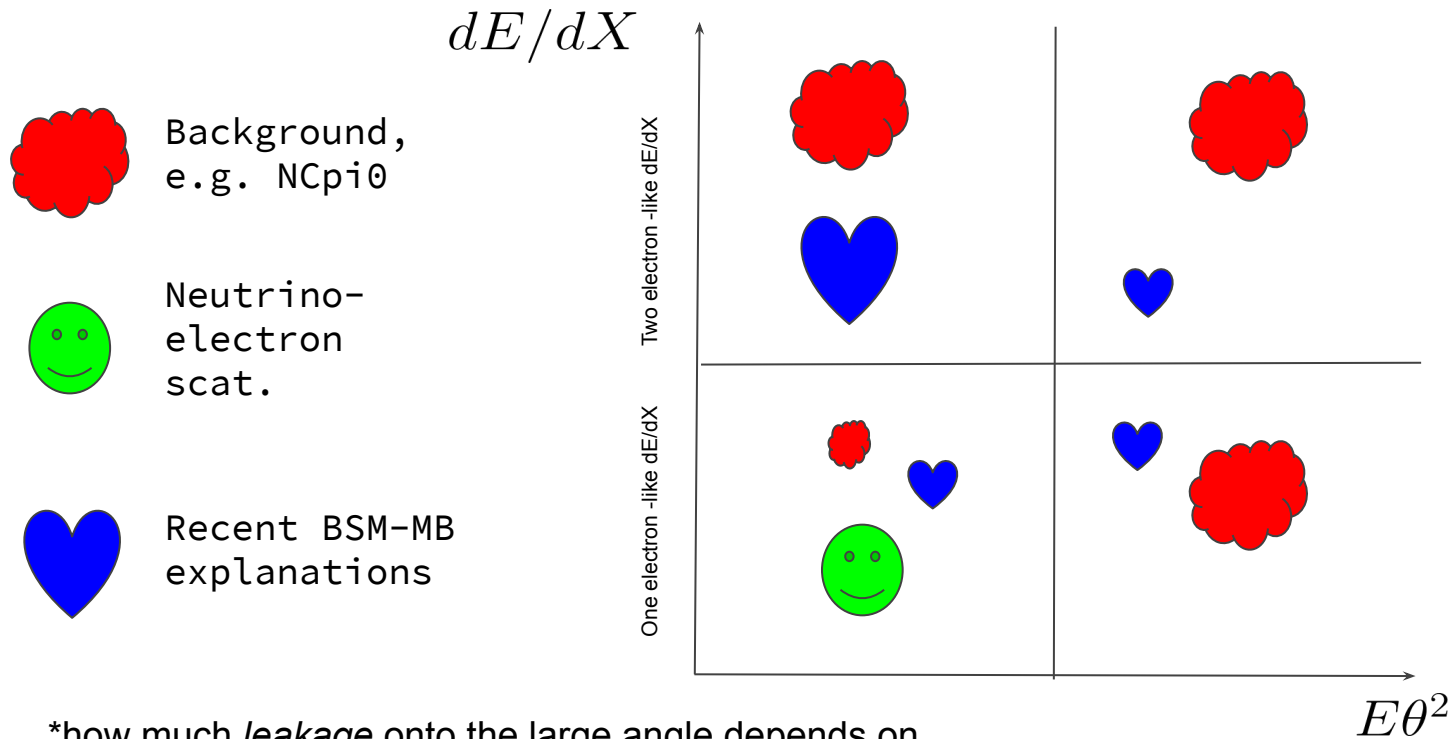
- **LSND**
  - **TEXONO**
  - **Borexino**
  - **SuperK**
  - **MINERvA Low-Energy**
  - **CHARM-II**
- Too low energies for BSM case of interest
- We will focus on these experiments.

Will measure it very soon @ **MINERvA Medium-Energy, NOvA,**  
and later @ **DUNE**.

# Strategy

---

## Electron-neutrino-like scattering search



\*how much *leakage* onto the large angle depends on model parameter and neutrino energy.

# Working around limited information!

---

By design at final cut level CHARM-II and Minerva measurements have small backgrounds: also means small amount of BSM-signal leaking in. We cannot use the final event samples to constrain the new models :(!

Would be great if we had access to the reconstructed electron energy and angular distributions at different cut levels.



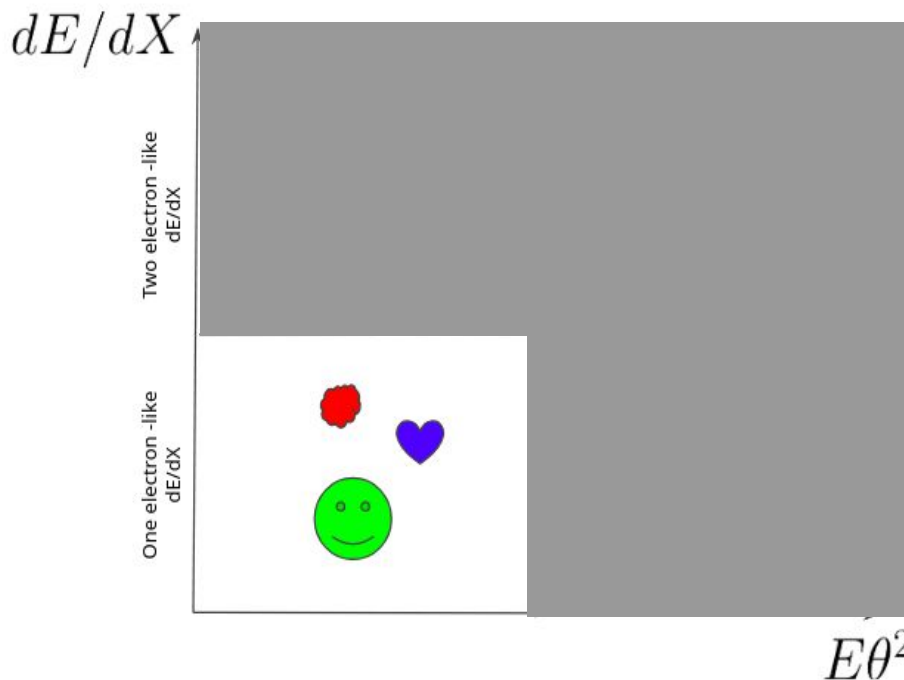
Background,  
e.g. NC,  $\pi^0$



Neutrino-  
electron  
scat.



Recent BSM-MB  
explanations

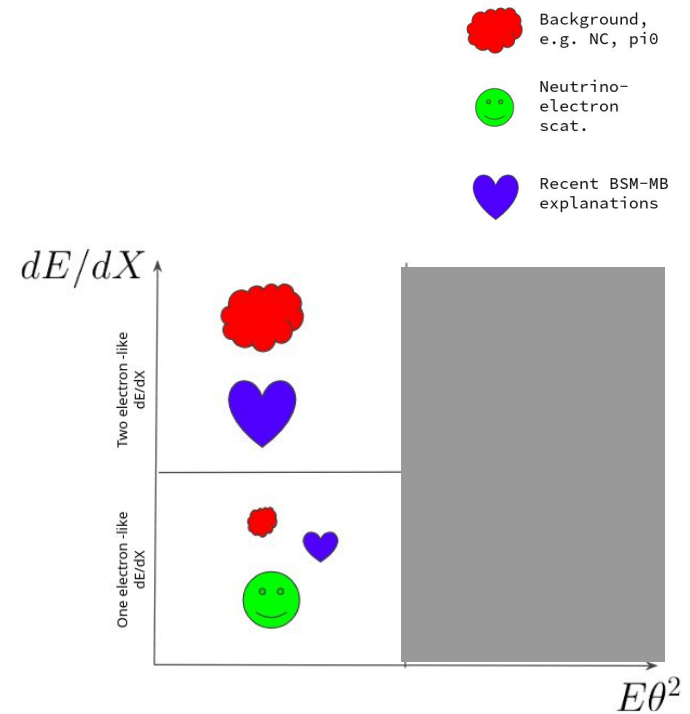
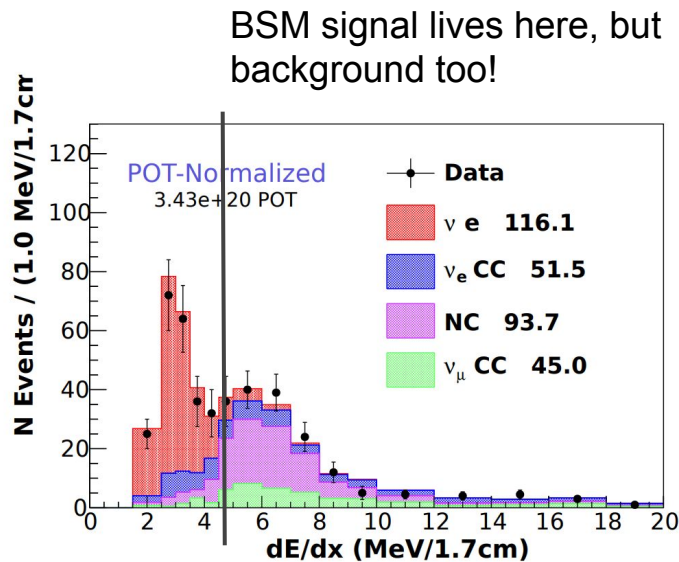


# MINERvA analysis strategy

For MINERvA we are going to use the  $dE/dX$  distribution of candidate electron-neutrino scattering events.

All MINERvA cuts applied, except for the final  $dE/dX$  cut!

Note that backgrounds have been tuned here!



# MINERvA analysis strategy

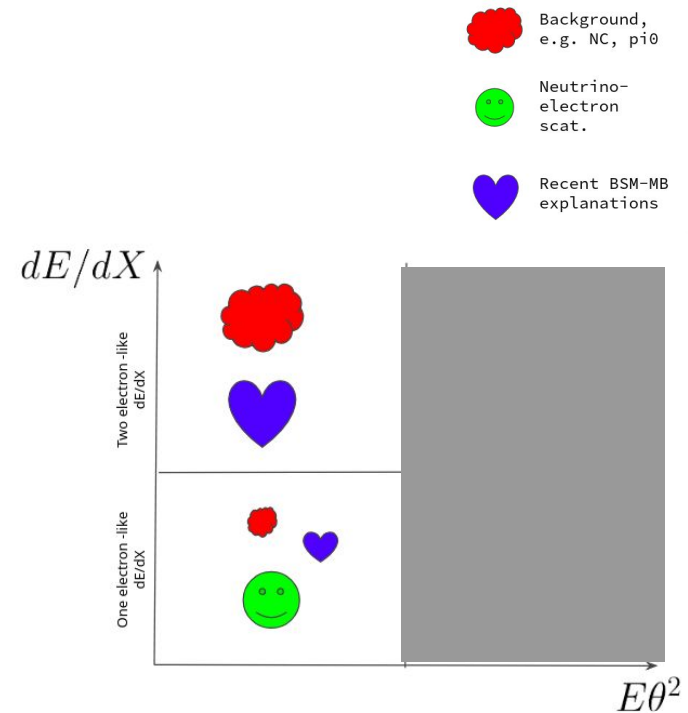
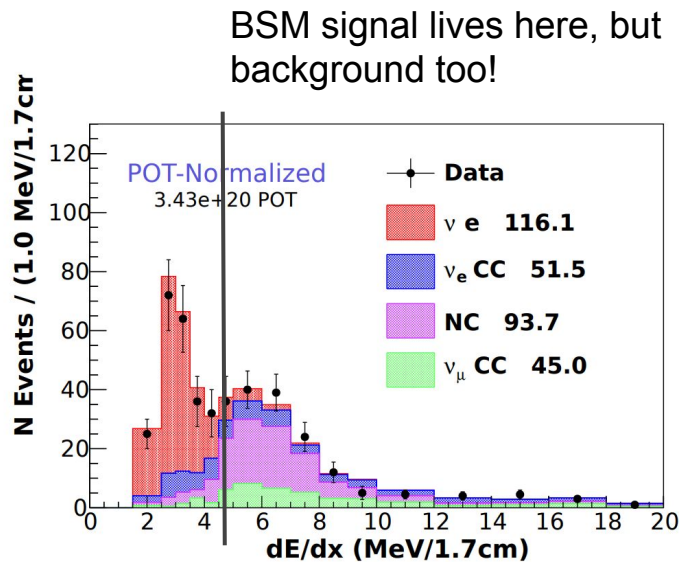
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Parameter	Tuned value
$\nu_e$	$0.76 \pm 0.03$
$\nu_\mu$ NC	$0.64 \pm 0.03$
$\nu_\mu$ CC	$1.00 \pm 0.02$

For MINERvA we are going to use the  $dE/dX$  distribution of candidate electron-neutrino scattering events.

All MINERvA cuts applied, except for the final  $dE/dX$  cut!

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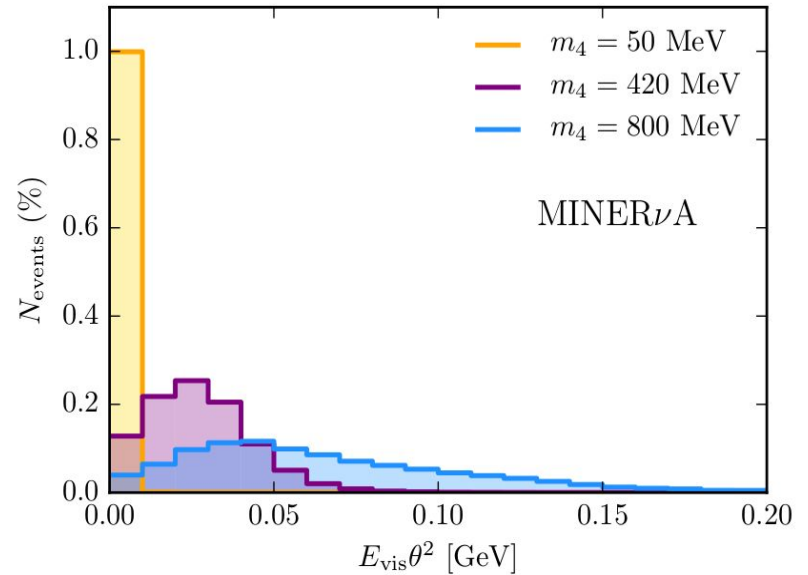
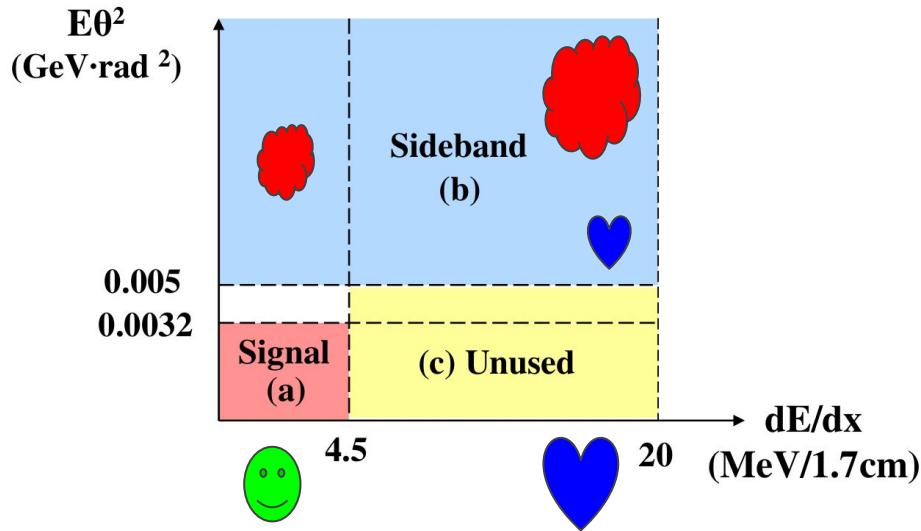




# Sidebands used for tuning background on MINERvA

Tuning parameters diagram from J. Park thesis:

<http://lss.fnal.gov/archive/thesis/2000/fermilab-thesis-2013-36.pdf>



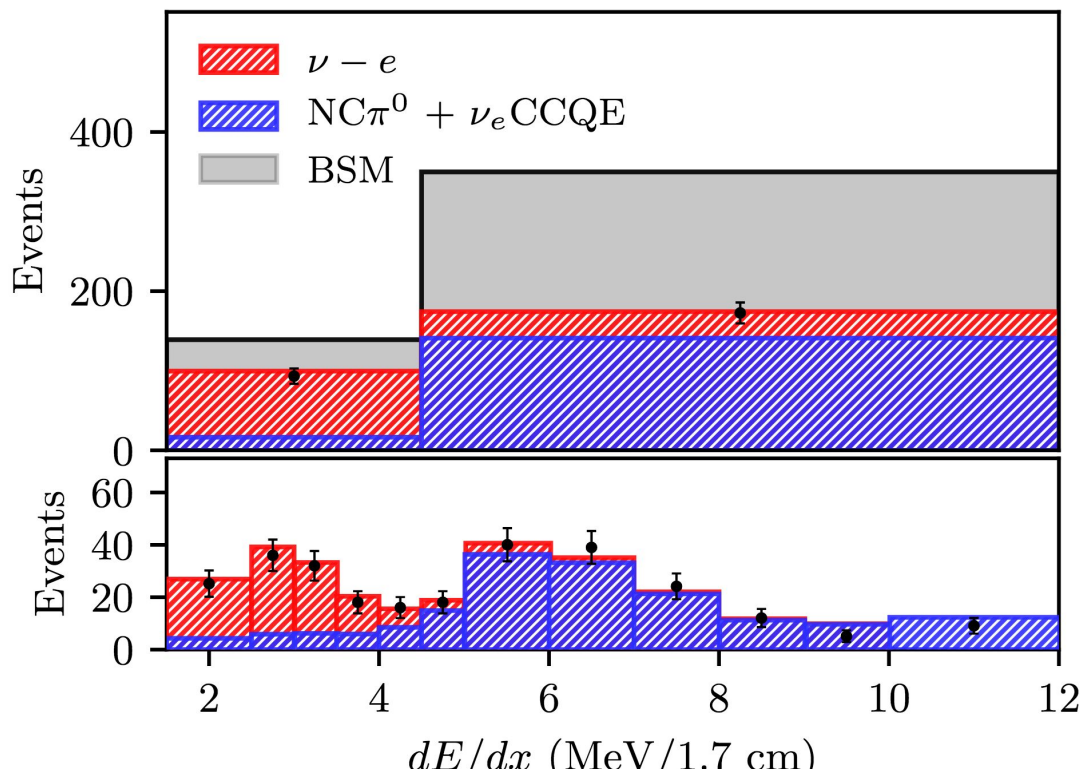
For large heavy neutrino masses the BSM contribution leaks the sideband used to constrain the background on the neutrino electron scattering region.

# MINERvA: Our Analysis setup

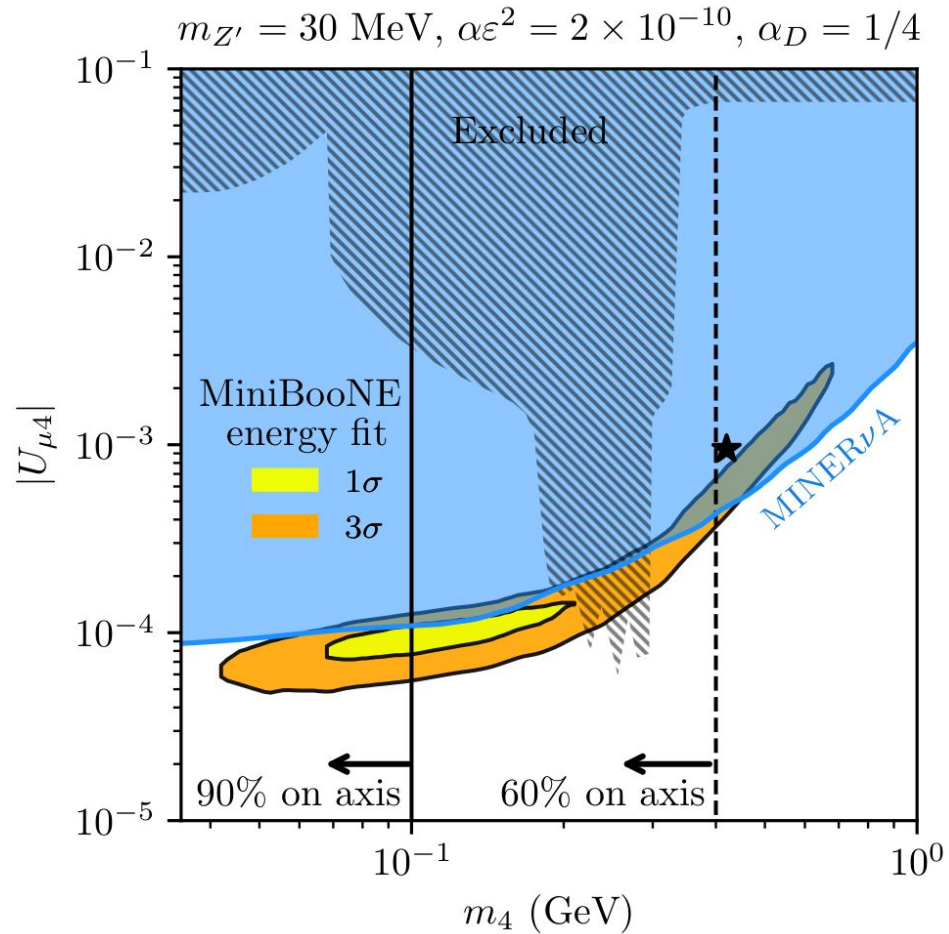
- We do a rate-only analysis on the single bin with

$dE/dx > 4.5 \text{ MeV}/(1.7 \text{ cm})$

- We use  $3.43 \times 10^{20}$  POTs, Assume fiducial mass of 6.10 tons.



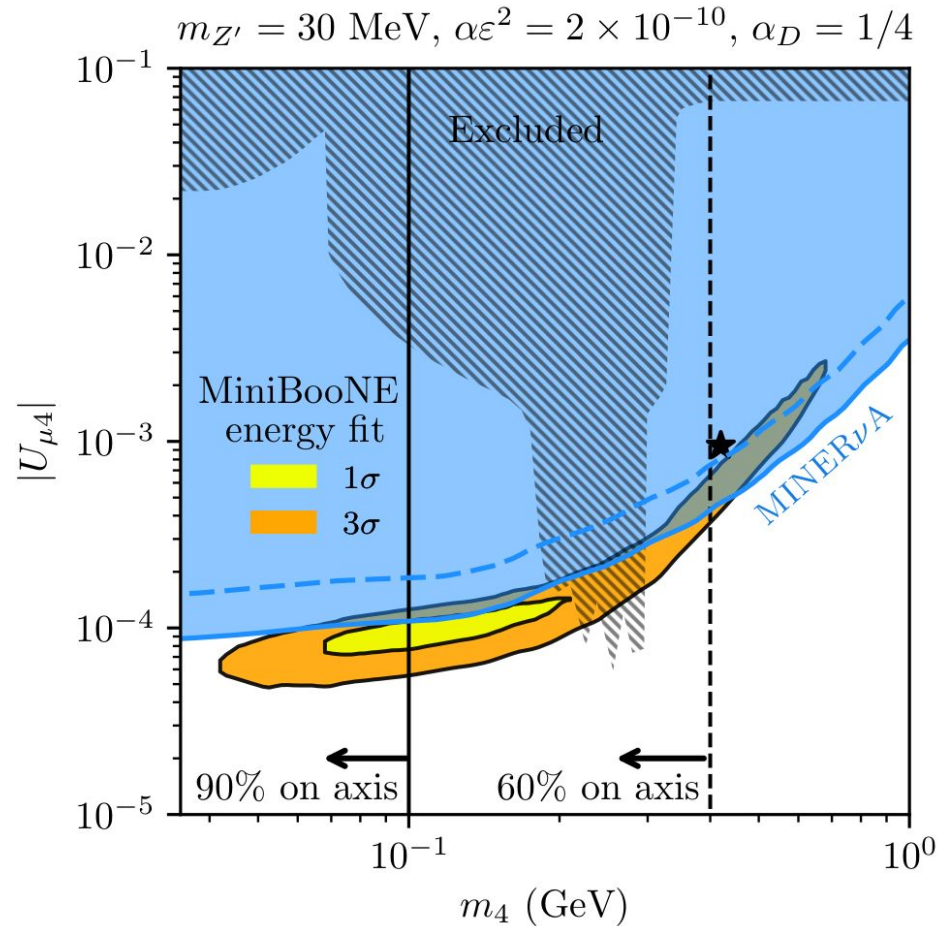
# MINERvA result



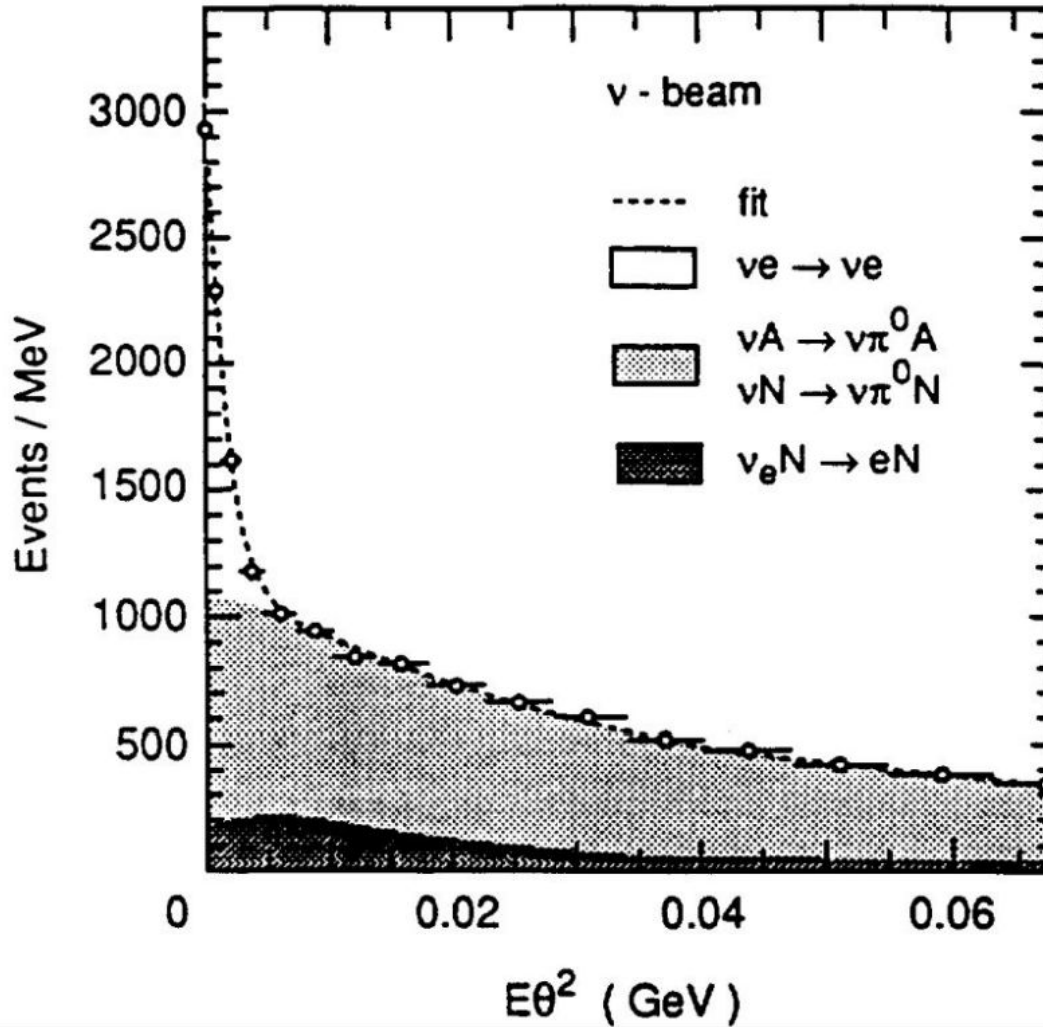
# MINERvA result

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We checked that changing the background uncertainty from 30% to 100% changes the result by no more than a factor of two. The constraint power is coming from the BSM signal overshooting the data.

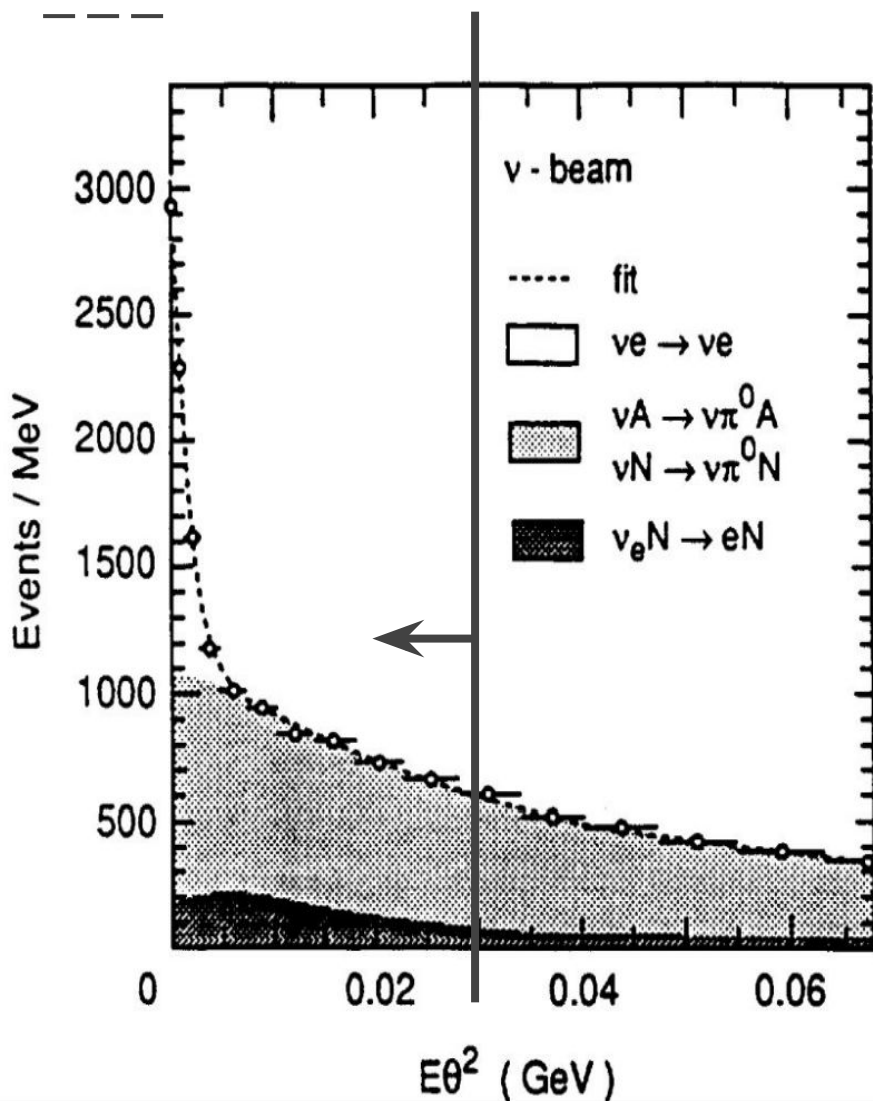


# CHARM-II: complementary measurement

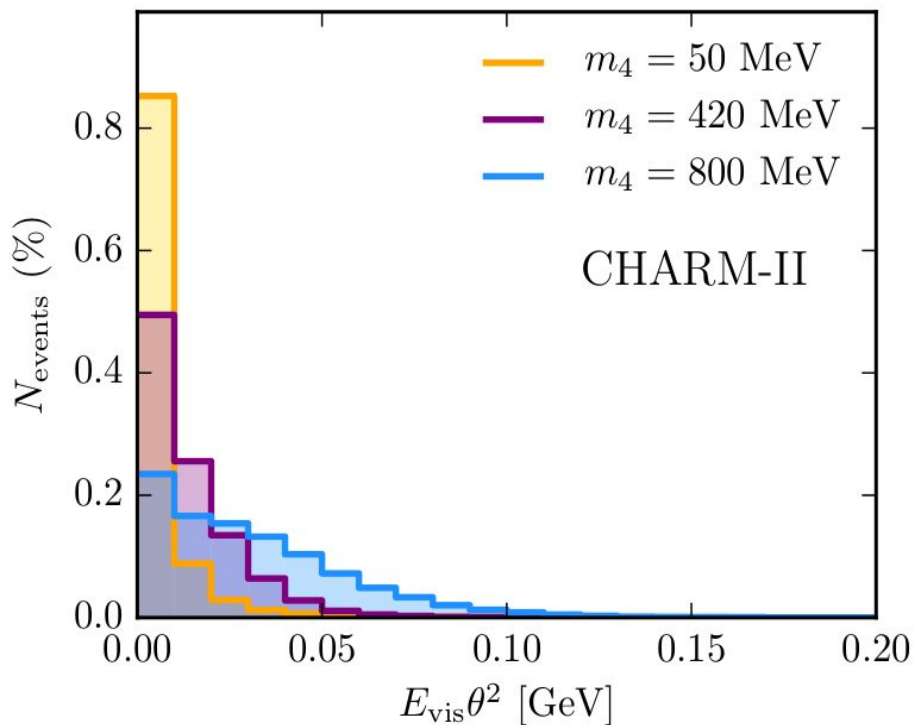


For CHARM-II we are going to use the  $E\theta^2$  distribution before the final  $dE/dX$  cut is applied.

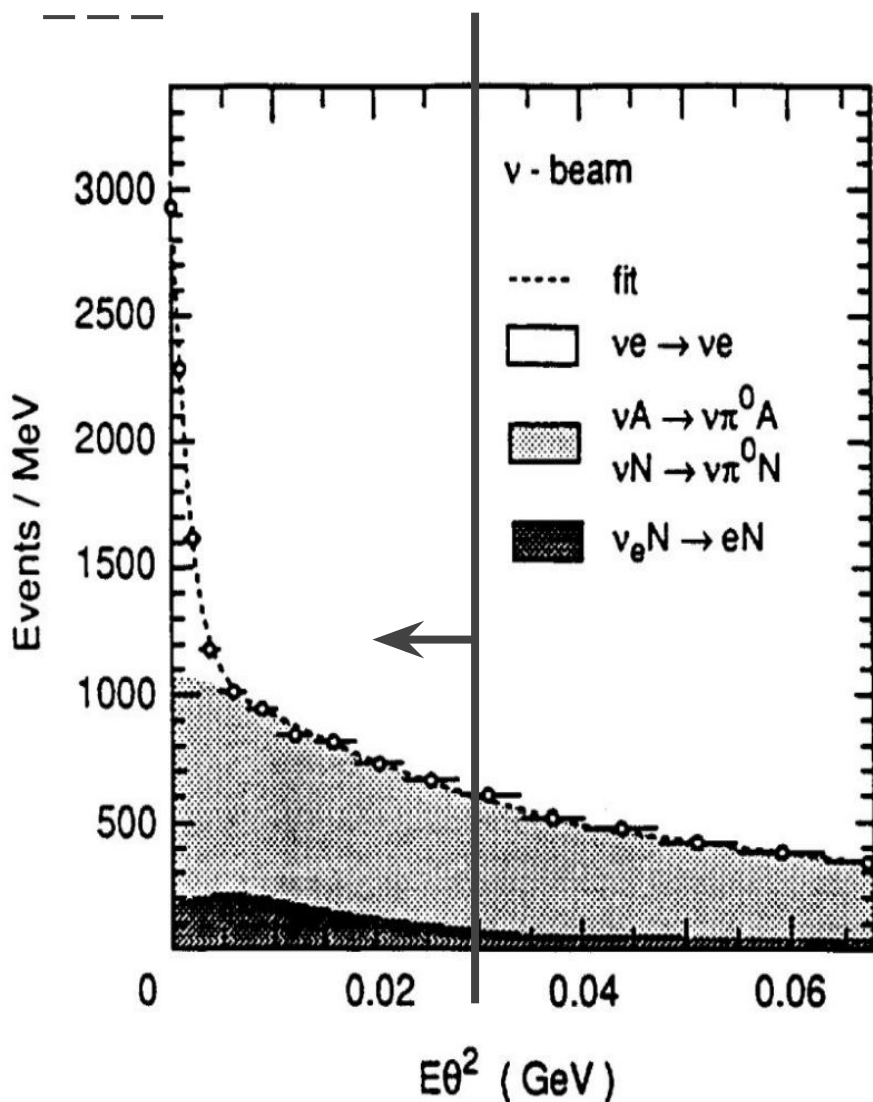
# Finding “BSM-safe” sideband to measure background



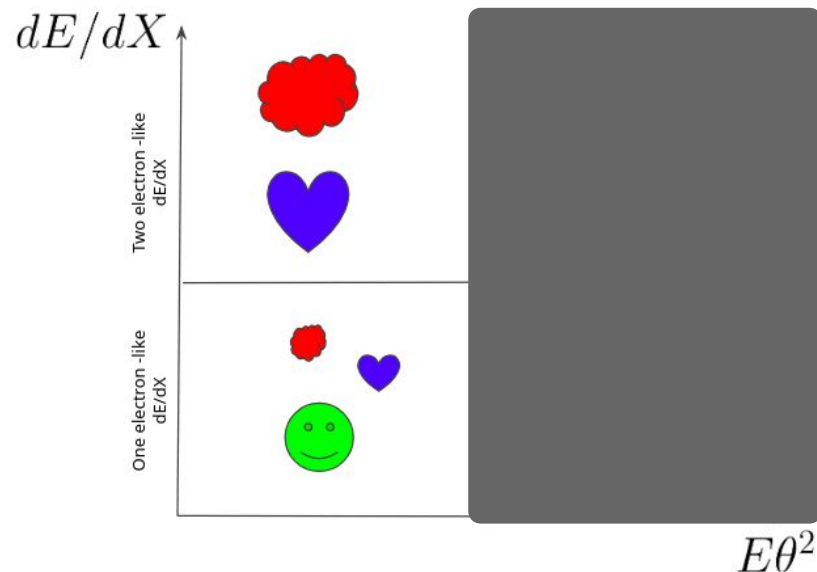
Angular distribution of BSM-signal



# CHARM-II: complementary measurement



For CHARM-II we use the distribution before the angular cut and  $dE/dX$  were applied



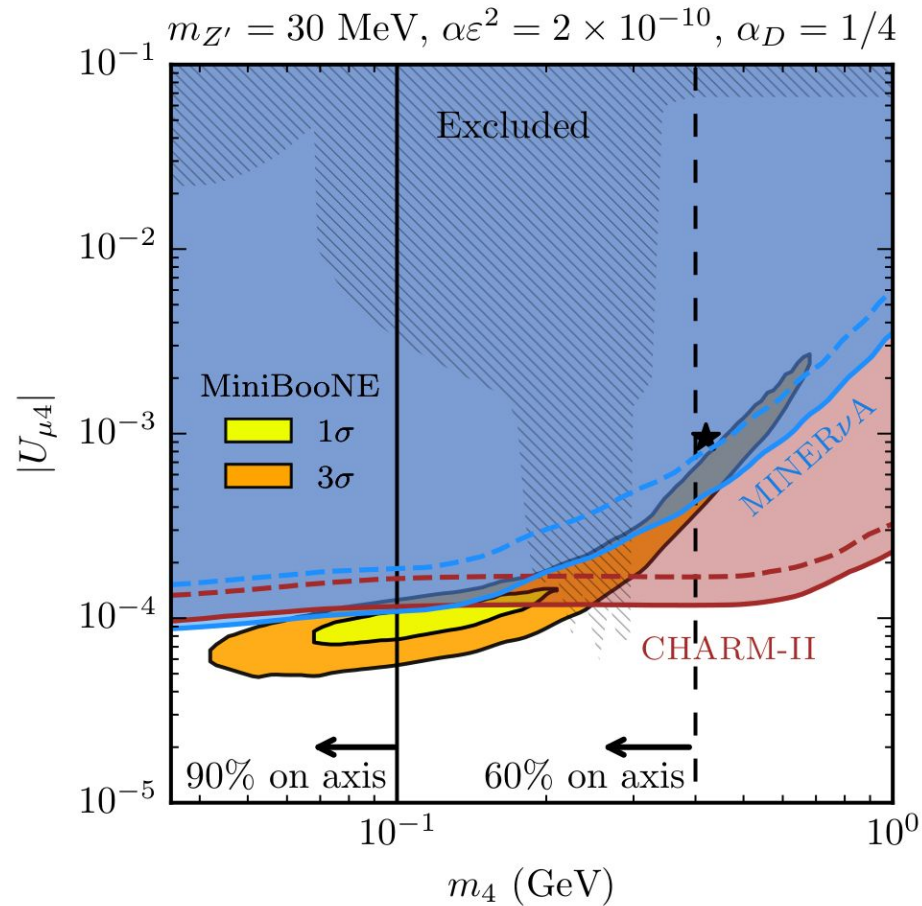
Use the region with  $E\theta^2 > 0.03$  to obtain the background uncertainty.

Allow for rate/slope to change; with this we estimate its rate to be constrain to be  $\sim 3\%$ .

# Putting it all together: the money plot

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Here for CHARM-II we also consider 3 times larger background uncertainty (dashed)

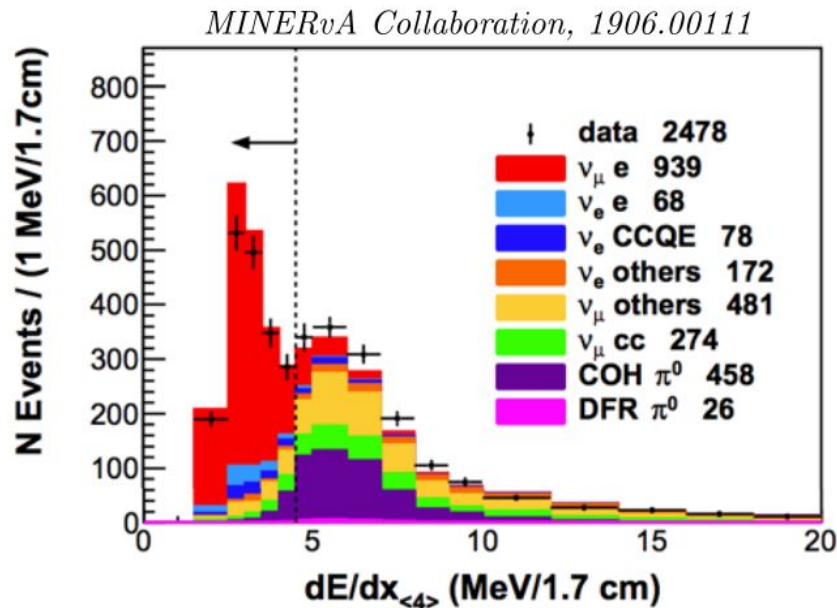




# But wait ... there is more! New Minerva ME set this summer!

BUT! MINERvA actually sees an excess in the region of interest!

Backgrounds are large — tuning takes place.



After tuning.

Excess attributed to coherent +0 events

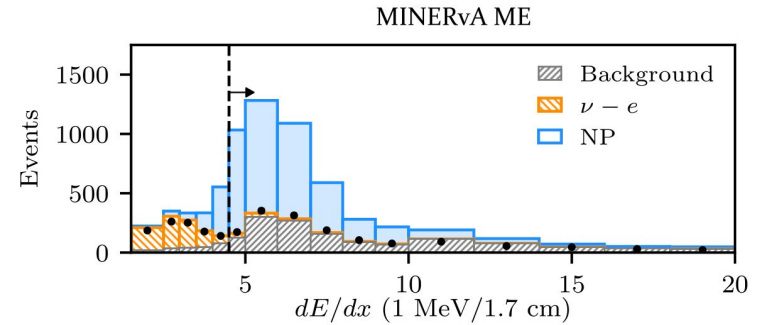
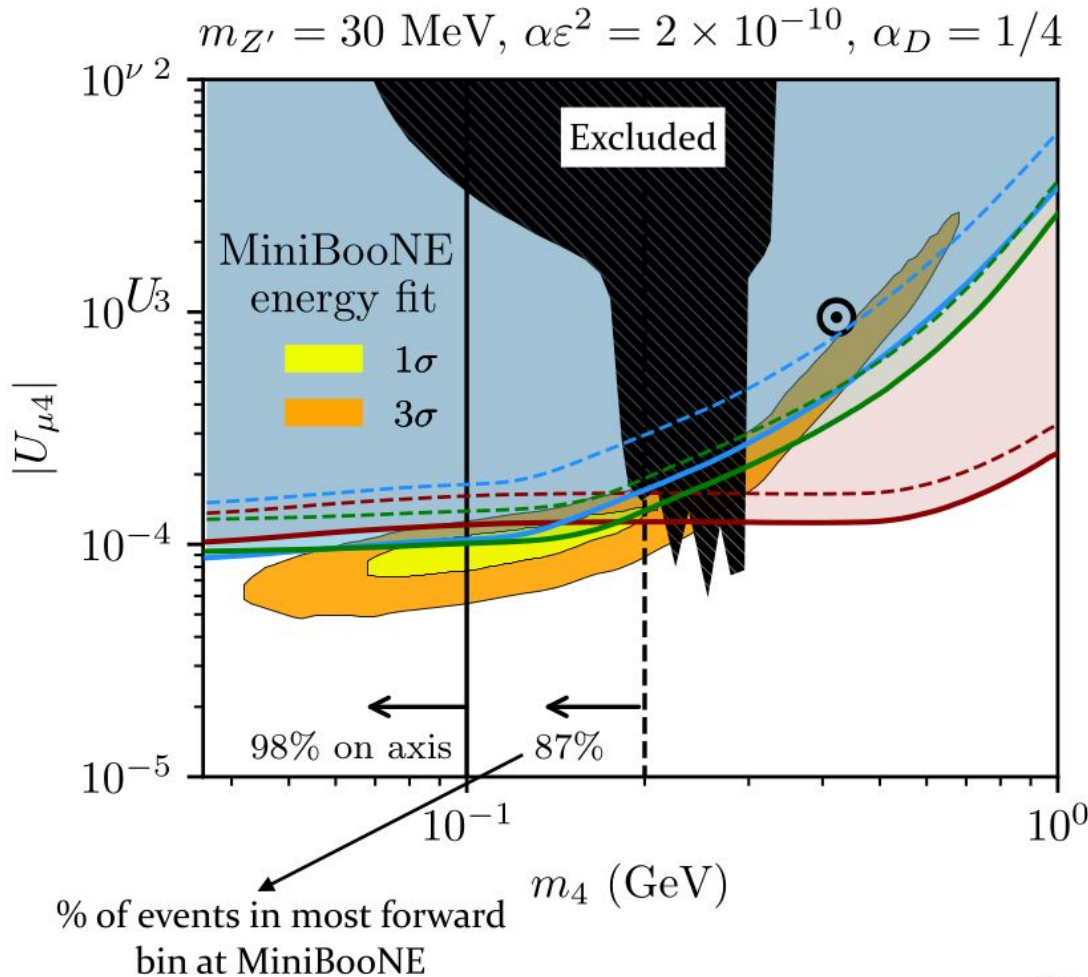
— disagreement with GENIE prediction claimed in NC and CC channels —

MINERvA Collaboration, 1906.00111

Process	Normalization
ν <sub>e</sub>	0.87 ± 0.03
ν <sub>μ</sub> CC	1.08 ± 0.04
ν <sub>μ</sub> NC	0.86 ± 0.04
NC COH 0.8 < E <sub>e</sub> < 2.0 GeV	0.9 ± 0.2
NC COH 2.0 < E <sub>e</sub> < 3.0 GeV	1.0 ± 0.3
NC COH 3.0 < E <sub>e</sub> < 5.0 GeV	1.3 ± 0.2
NC COH 5.0 < E <sub>e</sub> < 7.0 GeV	1.5 ± 0.3
NC COH 7.0 < E <sub>e</sub> < 9.0 GeV	1.7 ± 0.8
NC COH 9.0 < E <sub>e</sub>	3.0 ± 0.9

TABLE I. Background normalization scale factors extracted from the fits to kinematic sidebands, with statistical uncertainties.

# Updated constraints with Minerva ME



Vary uncertainty on backgrounds:

30% — 100% for MINERvA LE

40% — 100% for MINERvA ME\*

3 — 10% for CHARM-II

\*To appear in v2 of *C. Argüelles, 1812.08768*

# Take home message: lessons learned

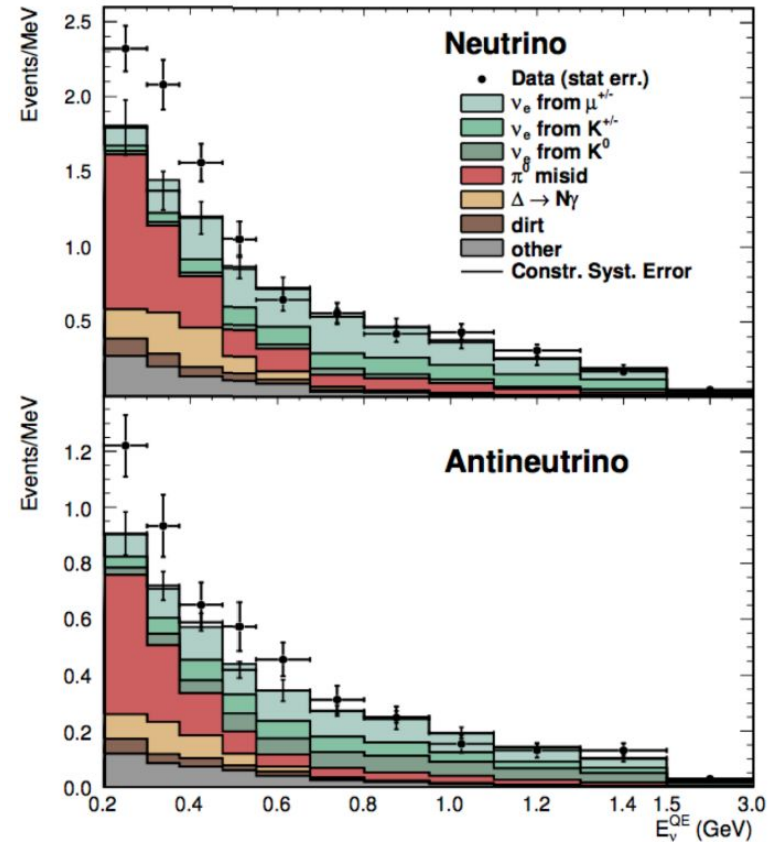
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- We are excited to see upcoming neutrino-electron scattering analyses by NOvA!
- We have used three different experiments to constrain recent MiniBooNE explanations. Tensions constrain large part of parameter space, but variations of the models are alive.
- Other variations can produce signatures in other experiments. E.g. 3+1 with large tau mixing leads to many double bangs at IceCube (P. Coloma, 1906.02106).
- Neutrino electron scattering is a powerful tool to constraint new physics interactions. But we need to understand the backgrounds better as we move into the sidebands to get all the power.

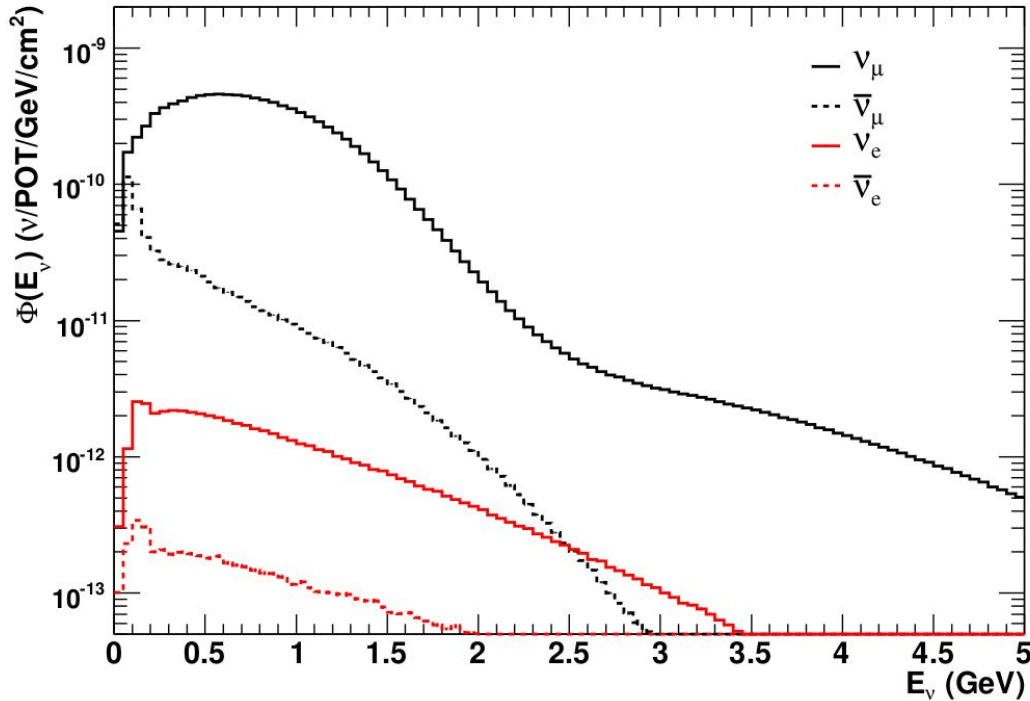
Bonus  
slides!

# MiniBooNE previous results

- Neutrino mode excess above the null primarily under 475 MeV. This is not expected given the LSND signal.
- Antineutrino mode saw excess above and below 475 MeV.
- Antineutrino mode was consistent with LSND, neutrino mode had tension.



# Booster beam

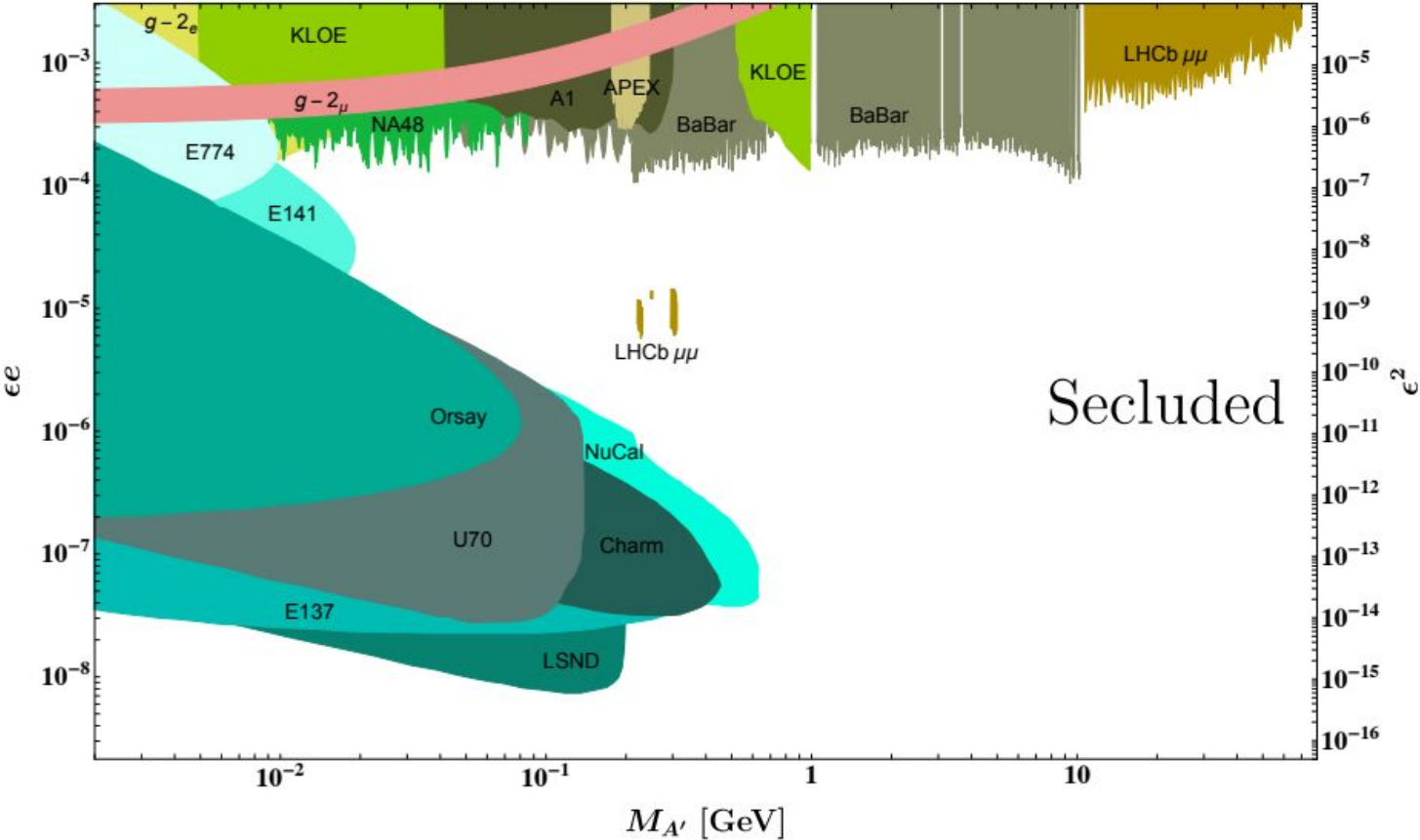


- Maxima  $< 1$  GeV neutrino energy
- Production of heavy states via neutrino interaction: hard.
- Heavier BSM physics look like “effective”-interactions; then angular distribution of excess wont work.

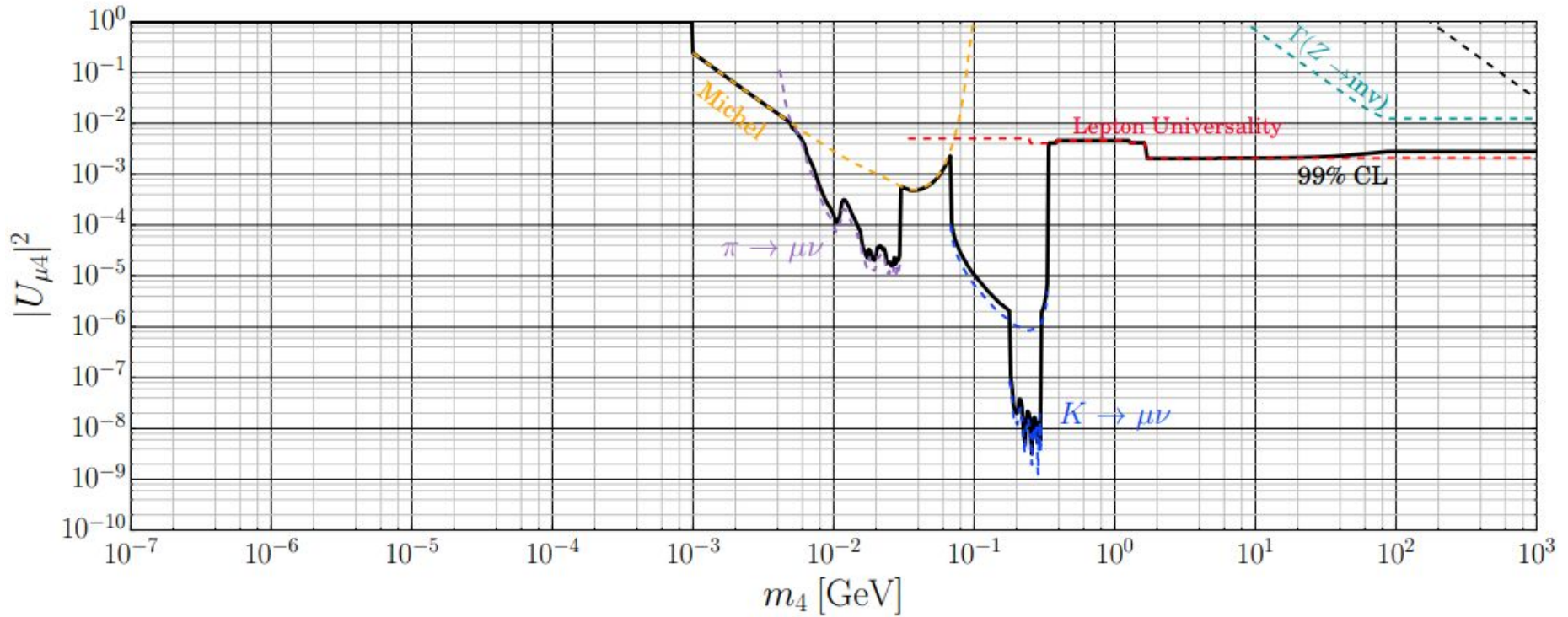
Recent proposals:  
Light new physics  $\sim < \text{GeV}$ .

# Dark Photon Searches

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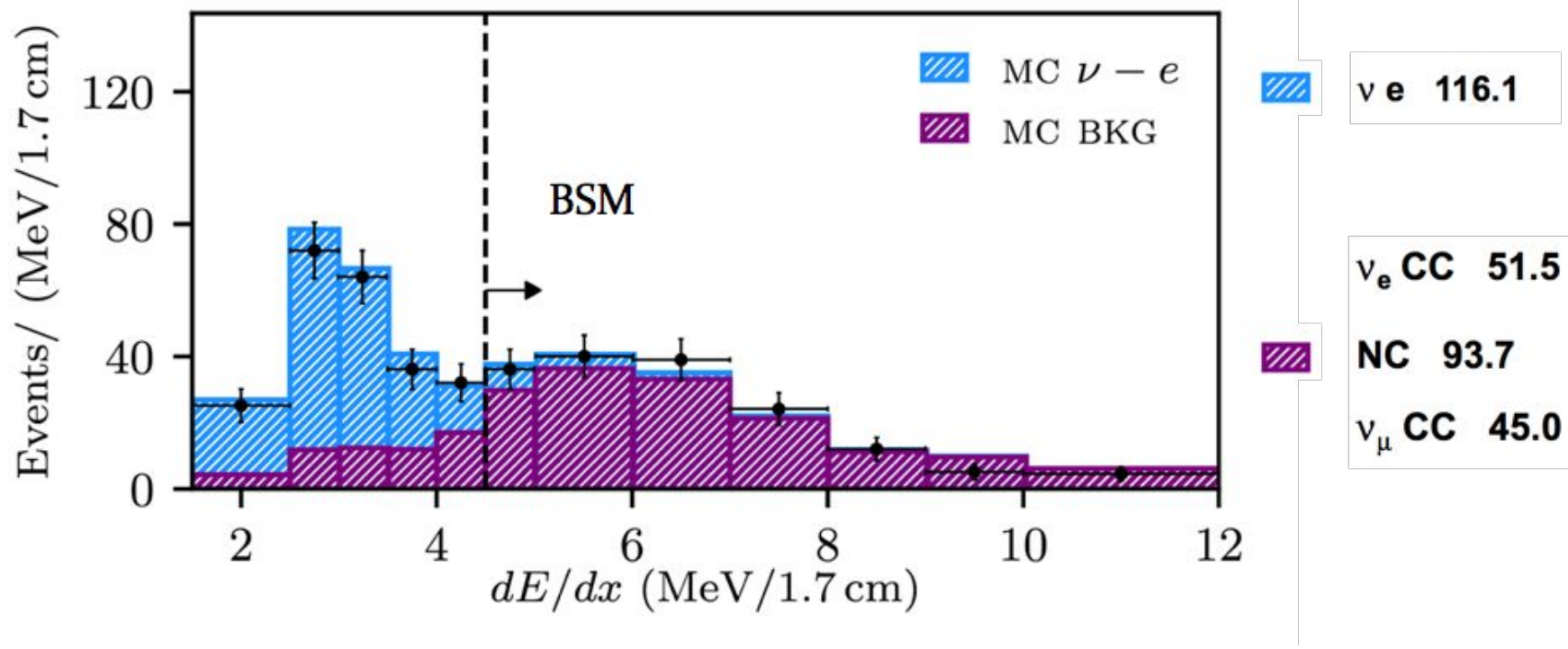
# Heavy Neutrino Limits



See arXiv:1511.00683



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# MINERvA: Our Analysis setup

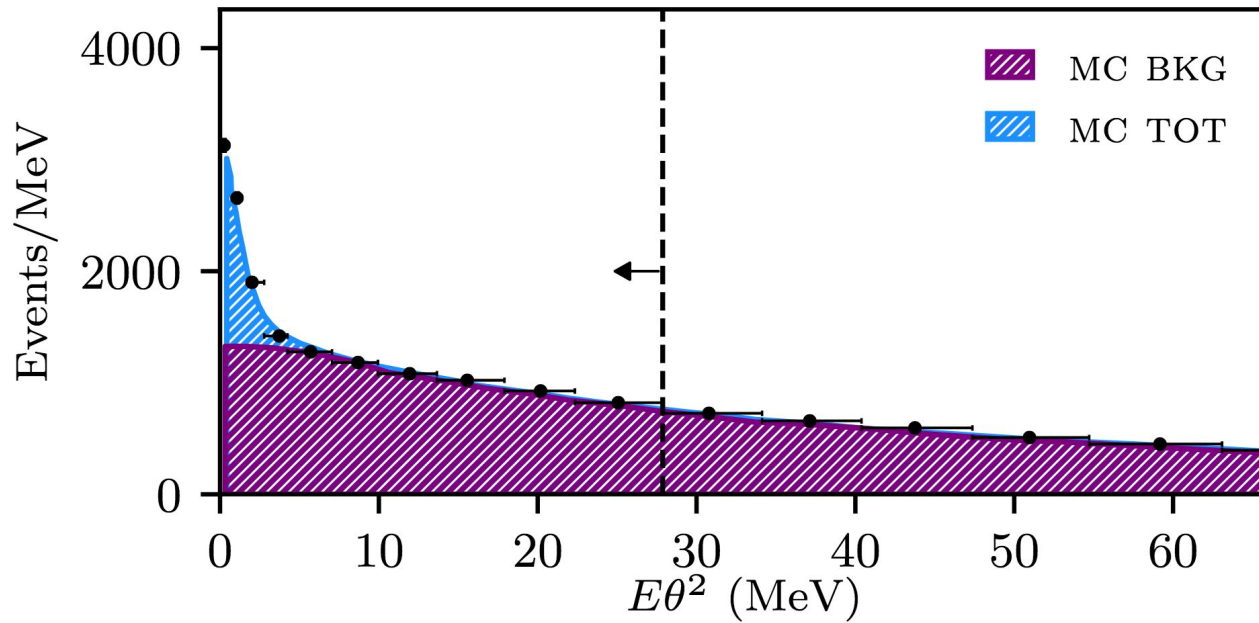
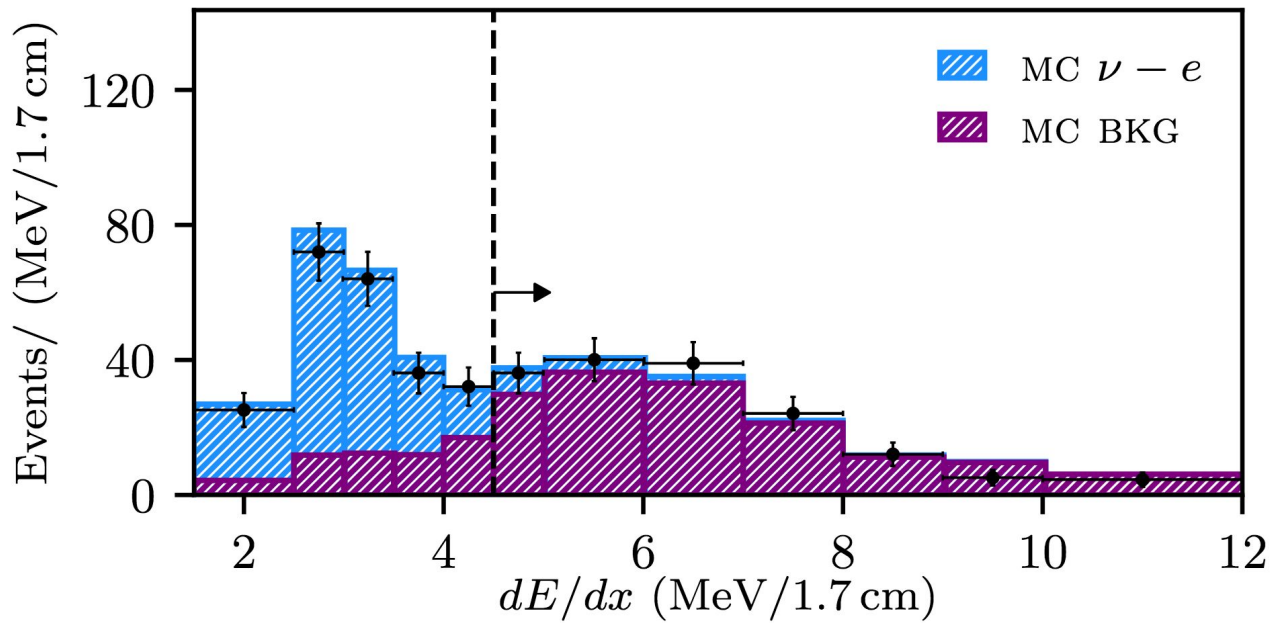
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We use the following  $\chi^2$  definition:

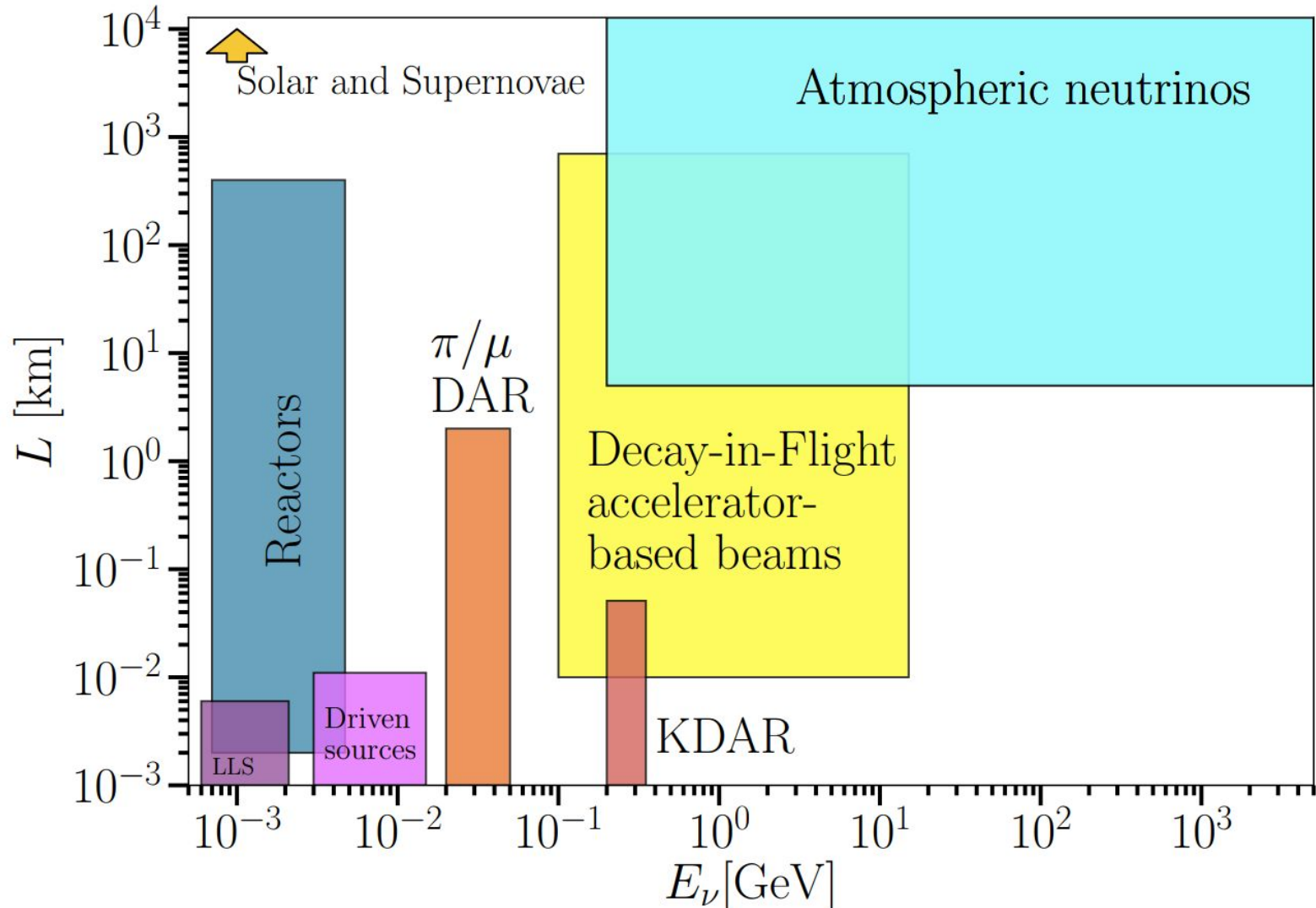
$$\chi_{\alpha\beta}^2 = \frac{(N_{\text{data}} - (1 + \alpha + \beta)\mu_{\text{MC}}^{\text{BKG}} - (1 + \alpha)\mu_{\text{MC}}^{nu-e} - (1 + \alpha)\mu_{\text{BSM}})^2}{N_{\text{data}}} + \left(\frac{\alpha}{\sigma_{\alpha}}\right)^2 + \left(\frac{\beta}{\sigma_{\beta}}\right)^2$$

- We set  $\sigma_{\alpha} = 10\%$  account for beam uncertainties.
- We set  $\sigma_{\beta} = 30\%$  motivated by the amount of tuning; conservative with respect to tune normalization uncertainty.
- We include only coherent contribution to the BSM signal to avoid hadronic activity cuts.

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# Neutrino flux panorama



A. Diaz, CA, G.Collin, JM. Conrad, M. Shaevitz to appear 60  
 \*very\* soon.

# Novel alternative explanations of the MiniBooNE anomaly

Assume that the SM gauge group is extended by a new U(1)', which mixes kinetically with the SM hypercharge

$$\mathcal{L}_{\text{kin}} \supset \frac{1}{4} \hat{Z}'_{\mu\nu} \hat{Z}'^{\mu\nu} + \frac{\sin \chi}{2} \hat{Z}'_{\mu\nu} \hat{B}^{\mu\nu} + \frac{m_{\hat{Z}'}^2}{2} \hat{Z}'^\mu \hat{Z}'_\mu$$

Also introduce a new SM-gauge singlet, charged under the new U(1)', which is allowed to mix with SM active neutrinos.

Bertuzzo et al. Phys. Rev. Lett. 121, 241801 (2018)

Ballett et al. arXiv:1808.02915

# Our CHARM-II analysis setup details

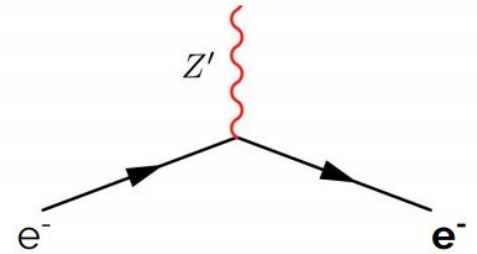
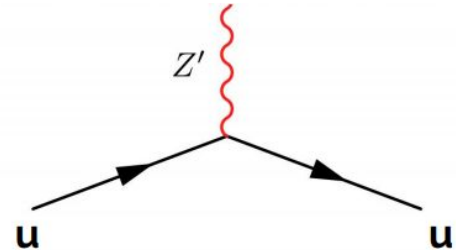
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- Rate-only analysis on a single bin with  $E\theta^2 < 0.03$  GeV.
- Same  $\chi^2$  definition as in MINERvA, but updated uncertainties.
- Background norm. from sideband  $\sim 3\%$ ; flux uncertainty  $\sim 4\%$ .
- We assume a fiducial mass of 547tons,  $\langle A \rangle \sim 20.7$ , and  $2.5e19$  POT.

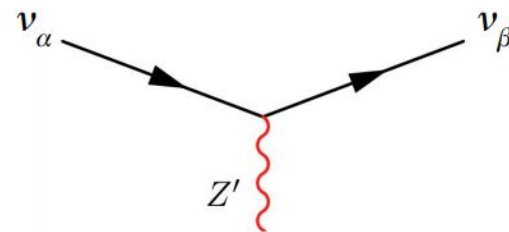
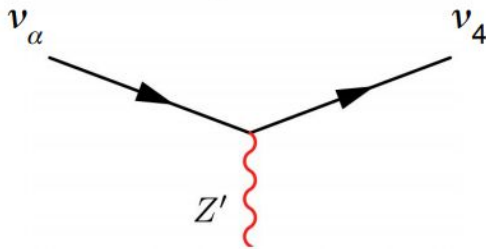
# Interaction Lagrangian

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$$\mathcal{L} \supset -eq_f c_W \chi \bar{f} \gamma^\mu f Z'_\mu$$



$$U_{\alpha 4}^* g' \bar{\nu}_\alpha \gamma^\mu P_L \nu_4 Z'_\mu + U_{\alpha 4}^* U_{\beta 4} g' \bar{\nu}_\alpha \gamma^\mu P_L \nu_\beta Z'_\mu$$



Much more detail discussion of the model building aspect of these models in next speakers talk!