

# Neutron-induced cross section measurements using the Liquid Lithium Target at SARAF

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NPA-X

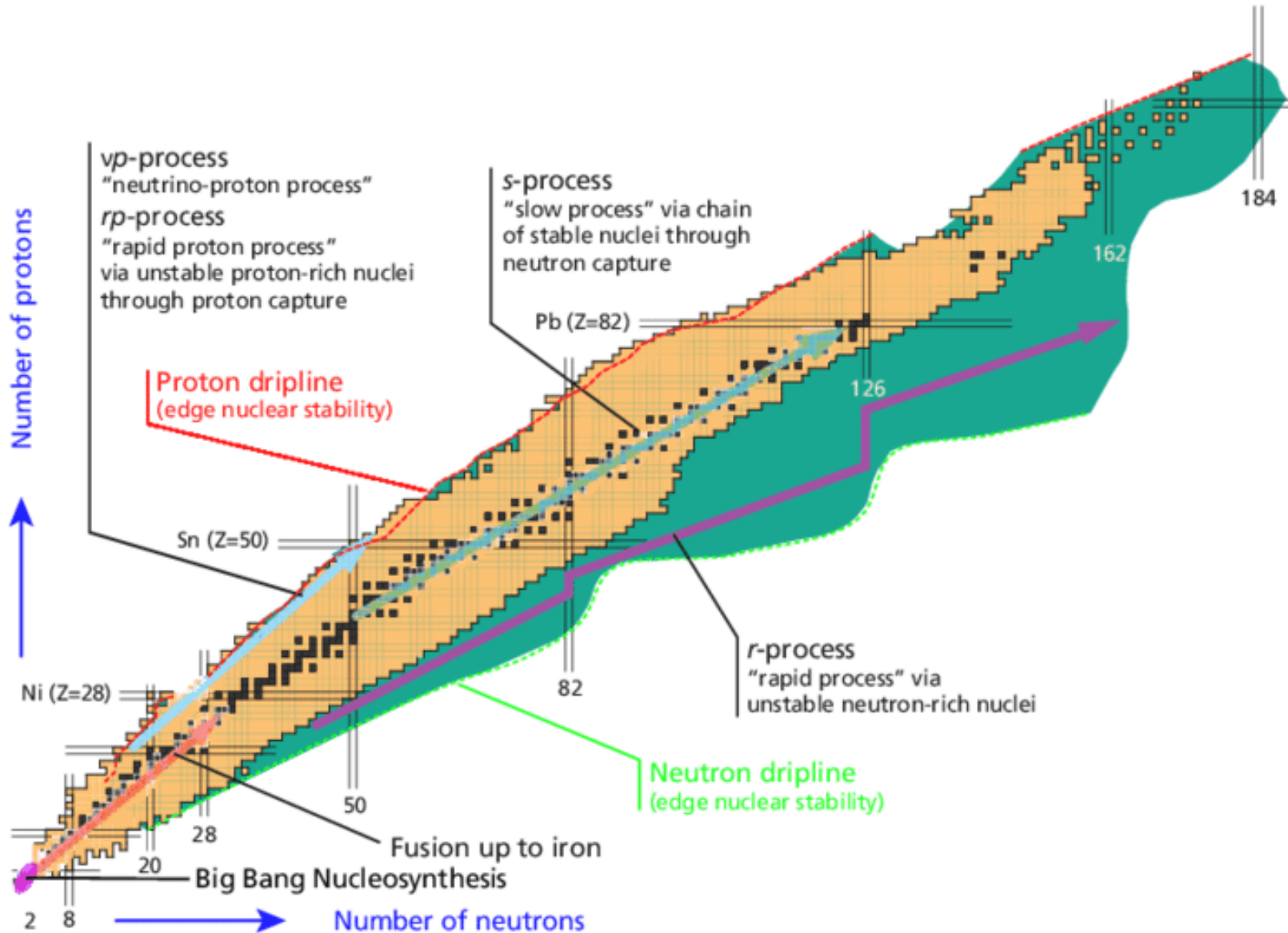
CERN 5/9/2022



האוניברסיטה העברית בירושלים  
THE HEBREW UNIVERSITY OF JERUSALEM



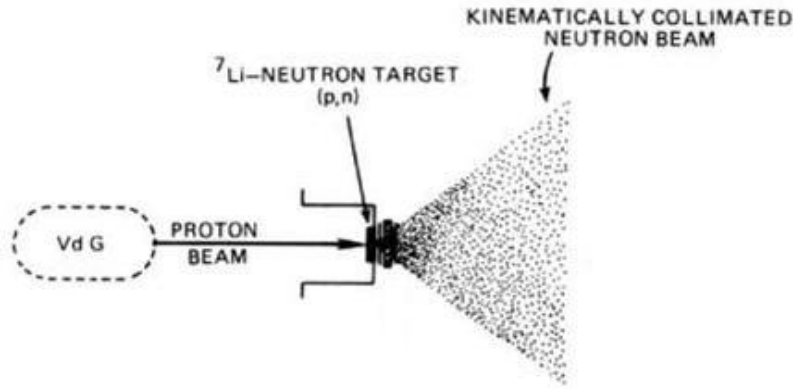
# Main stellar nucleosynthesis processes



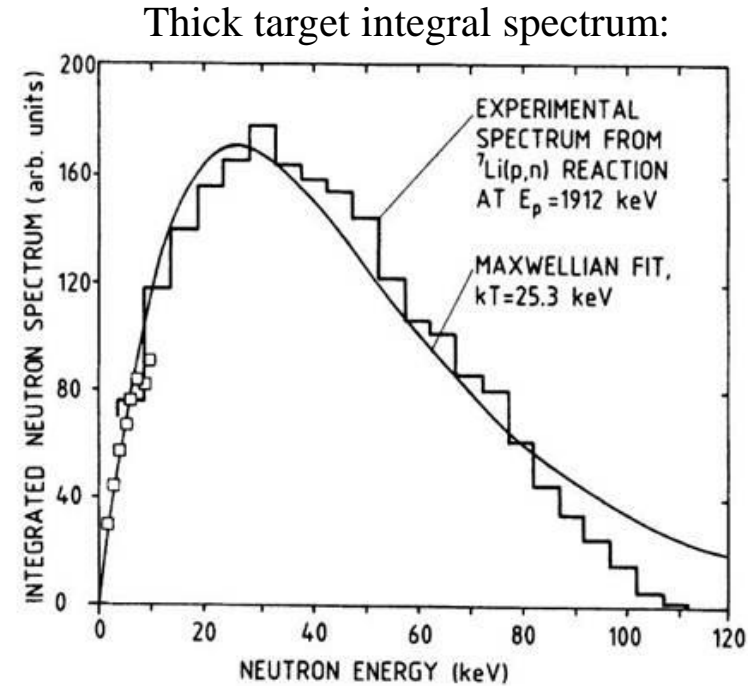
# ${}^7\text{Li}(p,n){}^7\text{Be}$ as a neutron source

$Q = -1.644 \text{ MeV}$

$E_{\text{th}}(p) = 1.881 \text{ MeV} \rightarrow E_n \approx 27 \text{ keV}$



$$\frac{dN}{dE_n} \propto E_n \exp(-E_n / kT)$$



W. Ratynski and F. Käppeler,  
PRC 37, 595 (1988)

- For  $E_p = 1912 \text{ keV} \rightarrow$  quasi-Maxwellian energy flux distribution with  $kT \approx 25 \text{ keV}$
- Neutron emission: forward cone with  $\sim \pm 60^\circ$  opening angle

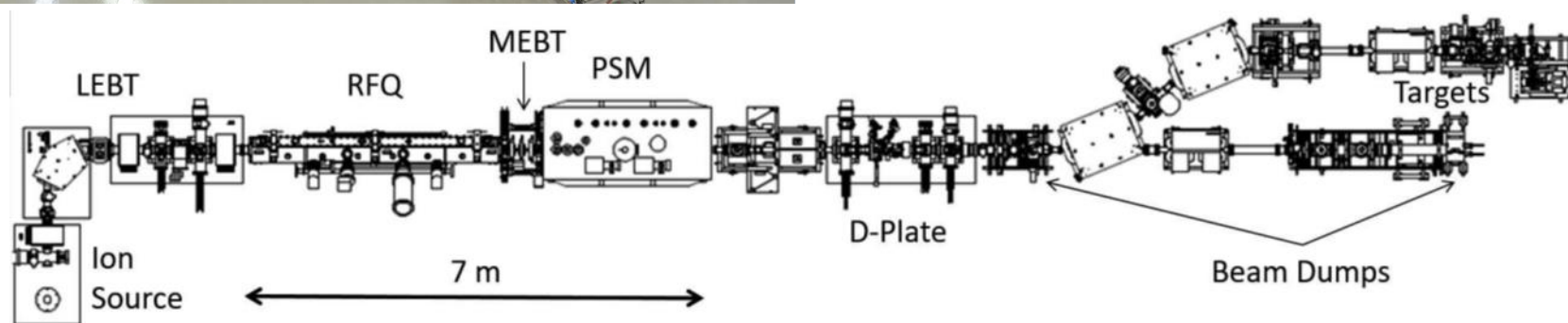
$$MACS = \frac{2}{\sqrt{\pi}} \frac{1}{(kT)^2} \int \sigma(E) \cdot E \cdot \exp(-E/kT) \cdot dE$$

# Soreq Applied Research Accelerator Facility (SARAF): Phase I



2 mA  $\approx 10^{16}$  protons/sec

2 mA @ 2 MeV  $\rightarrow$  4 kW beam power



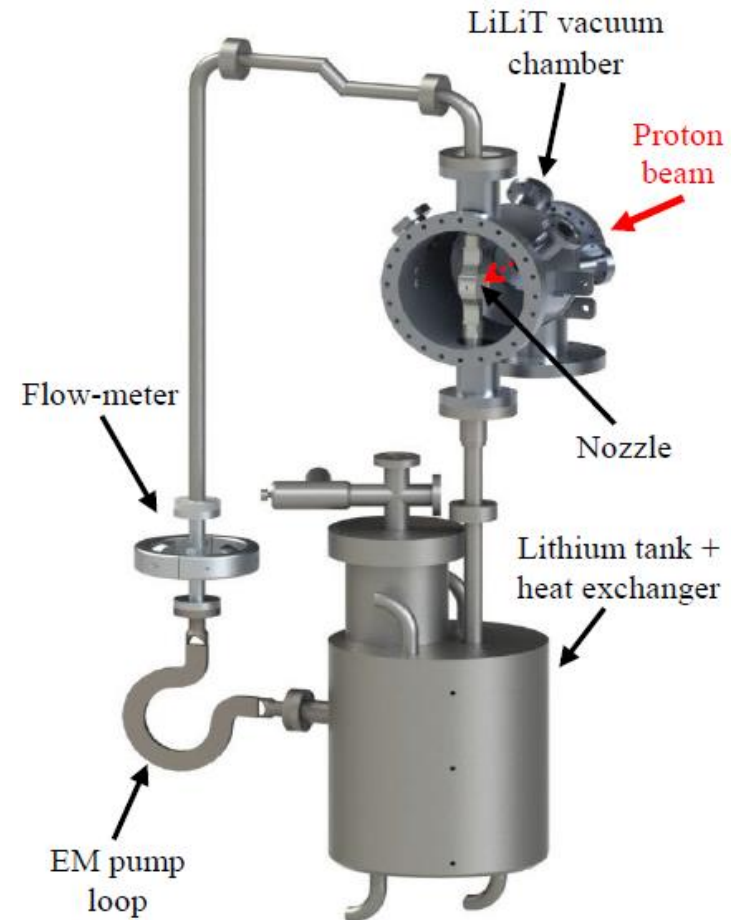
I. Mardor *et al.*, EPJ A **54**, 91 (2018)

# Liquid Lithium Target (LiLiT)



Peak power areal density: ~2.5 kW/cm<sup>2</sup>

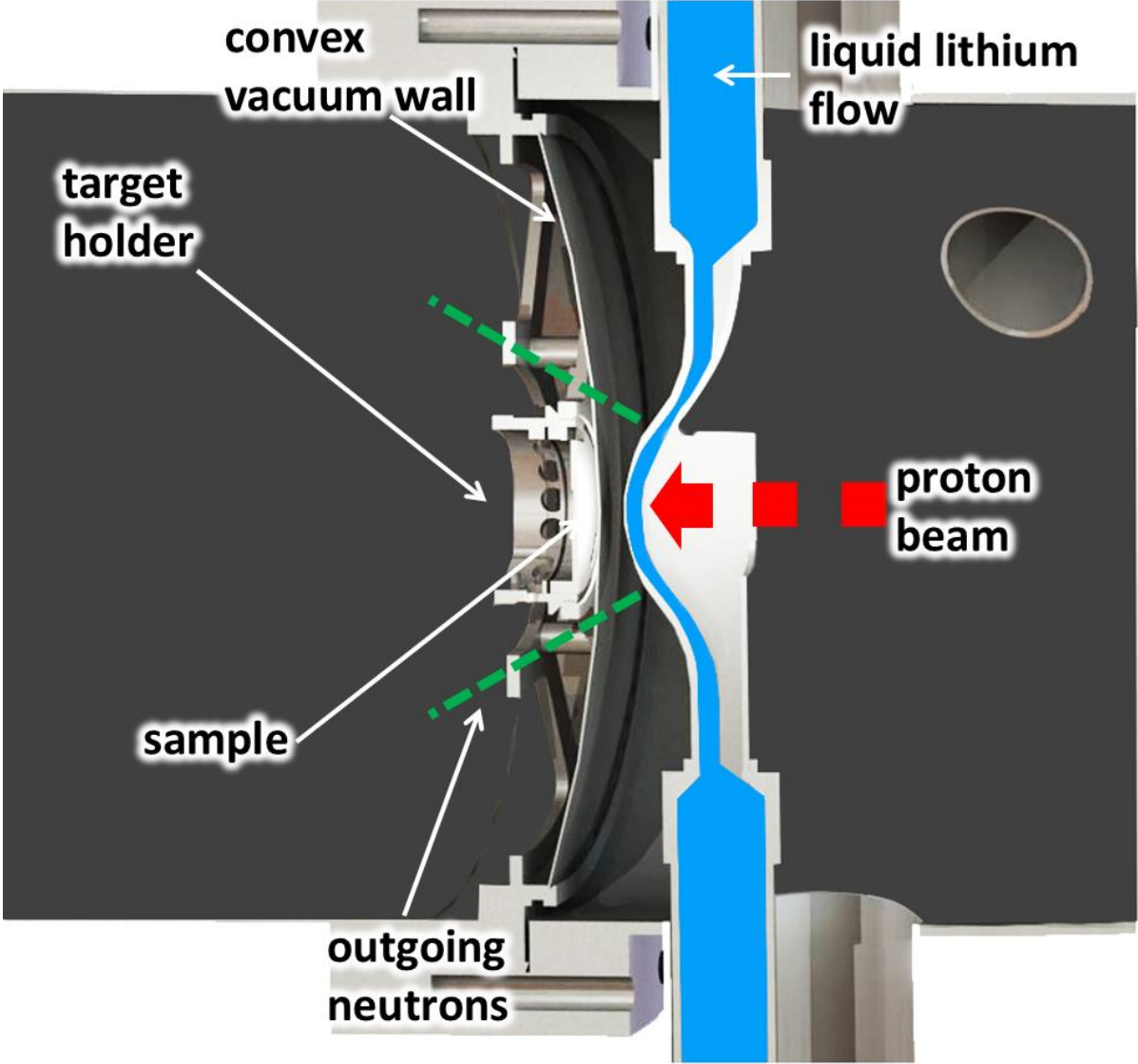
Peak power volume density: ~0.5 MW/cm<sup>3</sup>



S. Halfon *et al.*, RSI **84**, 12350 (2013)

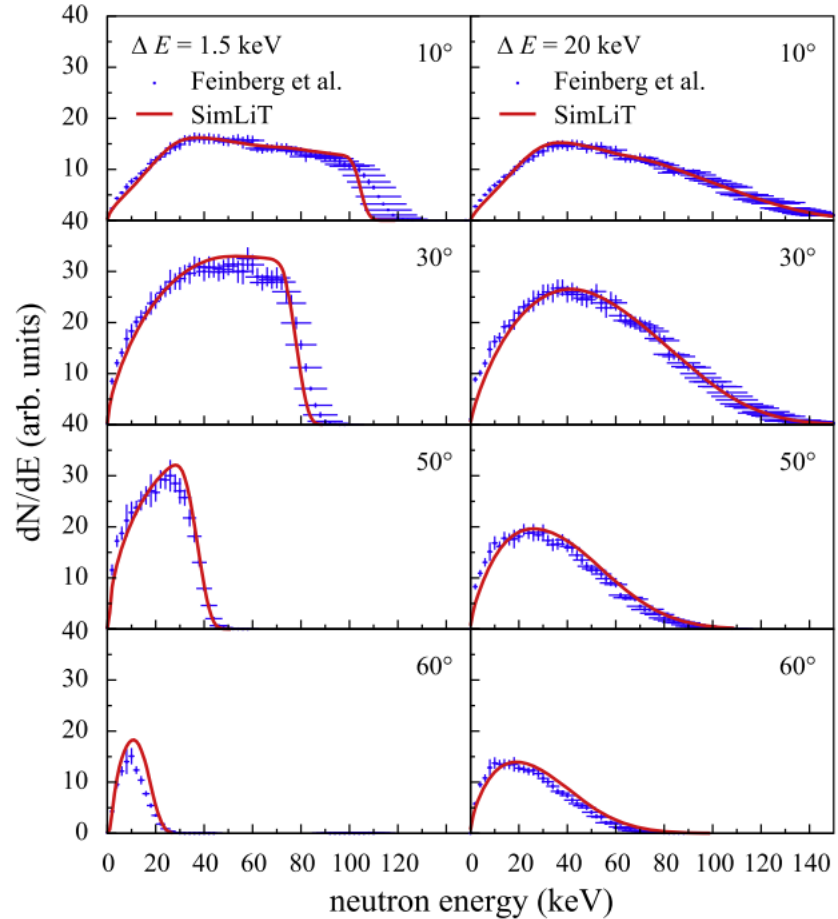
S. Halfon *et al.*, RSI **85**, 056105 (2014)

# Experimental setup

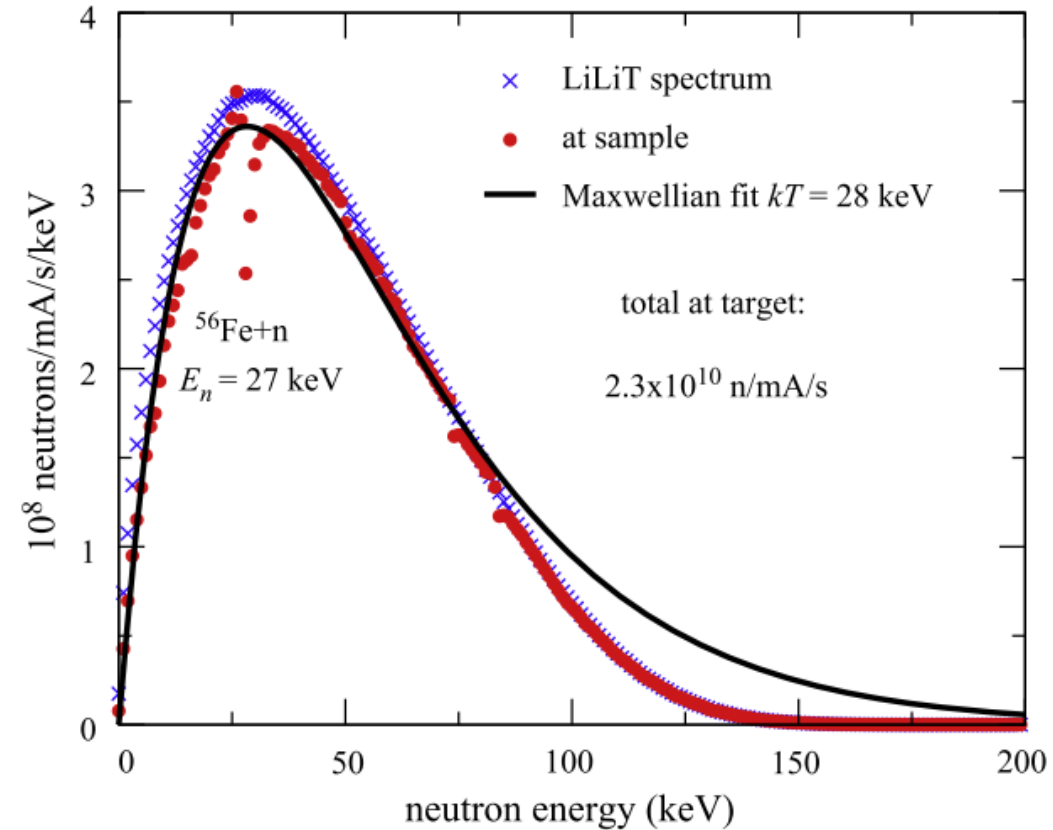


# Neutron spectrum and transport simulation

## SimLiT



## SimLiT+GEANT4

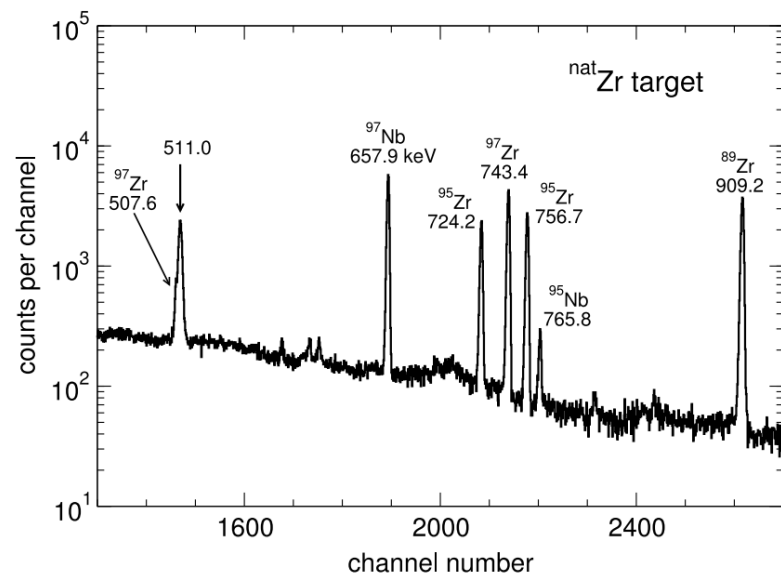


G. Feinberg *et al.*, PRC **85**, (2012)

M. Friedman *et al.*, NIM A **698**, (2013)

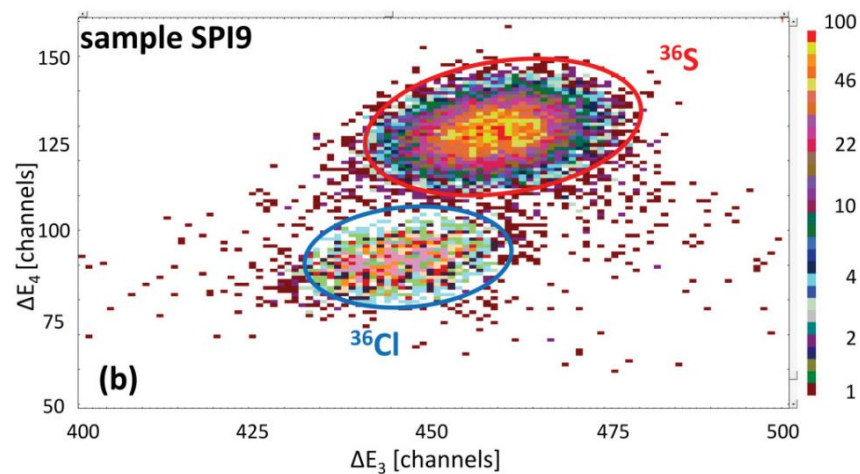
# Product measurement

## $\gamma$ -spectroscopy



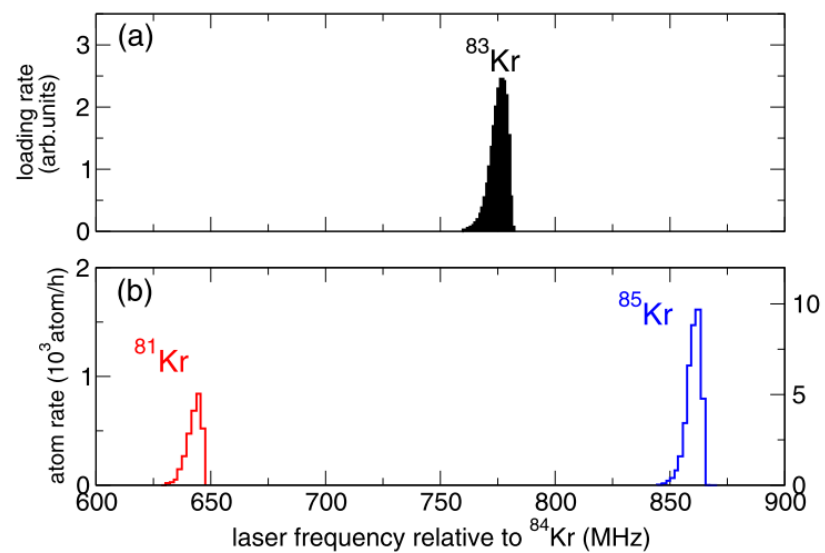
M. Tessler *et al.*, PLB **751**, (2015)

## AMS



M. Tessler *et al.*, PRL **121**, (2018)

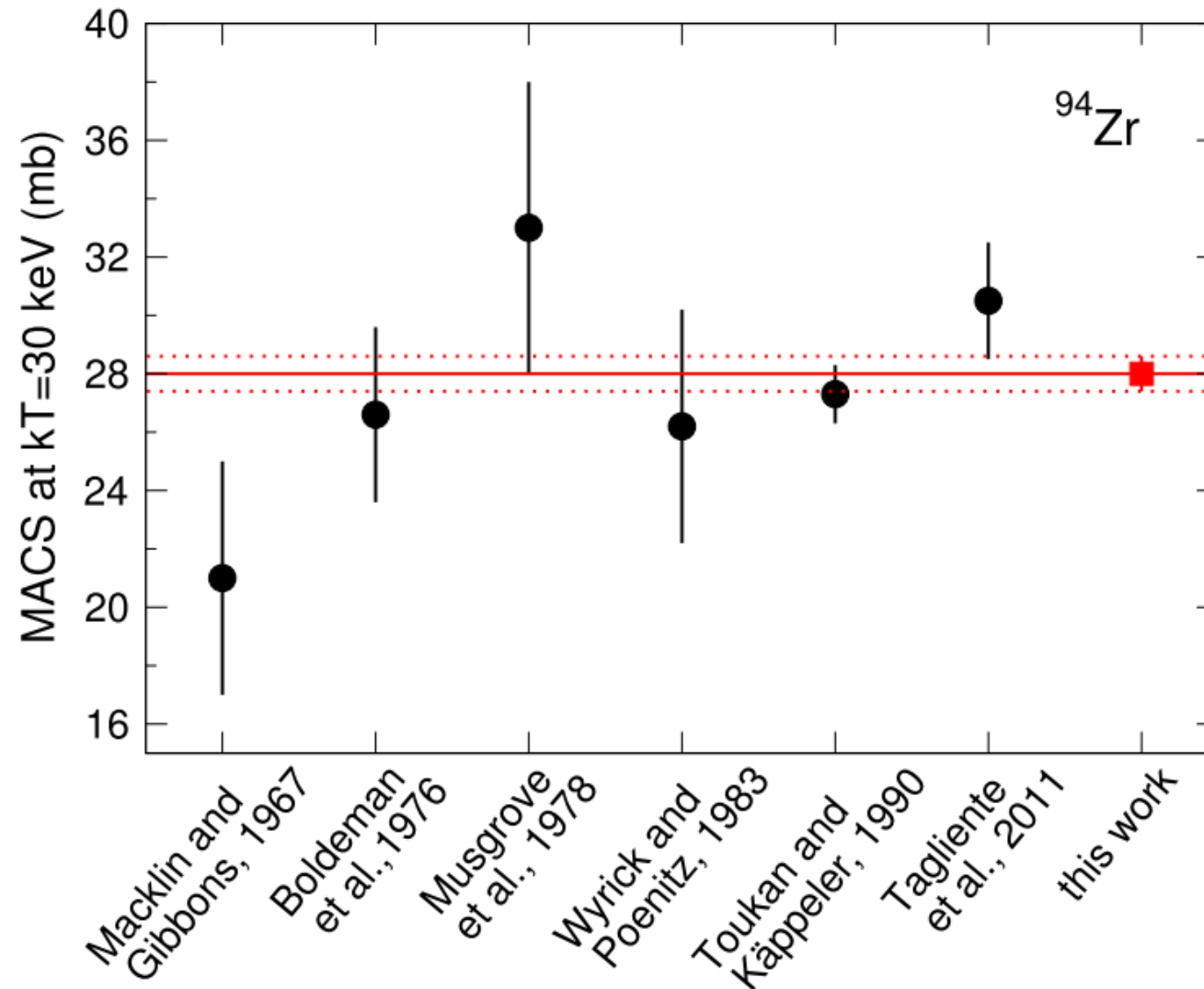
## atom trap trace analysis



M. Tessler *et al.*, PRC **104**, (2021)



# Commissioning experiment – $^{94,96}\text{Zr}(n,\gamma)$



# Phase-I experiments

Reaction	Detection tech.	Hebrew U, SARAF and collaborations below
$^{94,96}\text{Zr}(n, \gamma)$	$\gamma$ spect.	–
$^{90}\text{Zr}(\gamma, n)$	$\gamma$ spect.	–
$^7\text{Be}(n, \alpha)$	CR39	UConn, PSI, ILL WIS, CERN, TUNL
$^{23}\text{Na}, ^{35,37}\text{Cl}(n, \gamma)$	$\gamma$ spect., AMS	ANU, Goethe U, Rossendorf
$^{36,38}\text{Ar}(n, \gamma)$	AMS, LLC	ANL, Goethe U, U Bern
$^{53}\text{Mn}(n, \gamma)$	$\gamma$ spect.	PSI
$^{69,71}\text{Ga}(n, \gamma)$	$\gamma$ spect.	–
$^{74,78,80,82}\text{Se}(n, \gamma)$	$\gamma$ spect.	–
$^{78,80,84,86}\text{Kr}(n, \gamma)$	$\gamma$ spect., ATTA, LLC	ANL, Goethe U, U Bern
$^{80,82,86}\text{Kr}(\gamma, n)$	$\gamma$ spect., ATTA	ANL, Goethe U
$^{92}\text{Zr}(n, \gamma)$	AMS	ANL, ANU
$^{124,126,132,134}\text{Xe}(n, \gamma)$	$\gamma$ spect.	Goethe U
$^{136,138,140,142}\text{Ce}(n, \gamma)$	$\gamma$ spect.	–
$^{147}\text{Pm}(n, \gamma)$	$\gamma$ spect.	U Seville, ILL, PSI
$^{169,171}\text{Tm}(n, \gamma)$	$\gamma$ spect.	U Seville, nTOF, ILL, PSI
$^{208}\text{Pb}(n, \gamma)$	$\beta, \gamma$ spect.	U Seville
$^{209}\text{Bi}(n, \gamma)$	$\alpha, \beta, \gamma$ spect.	JRC, Geel

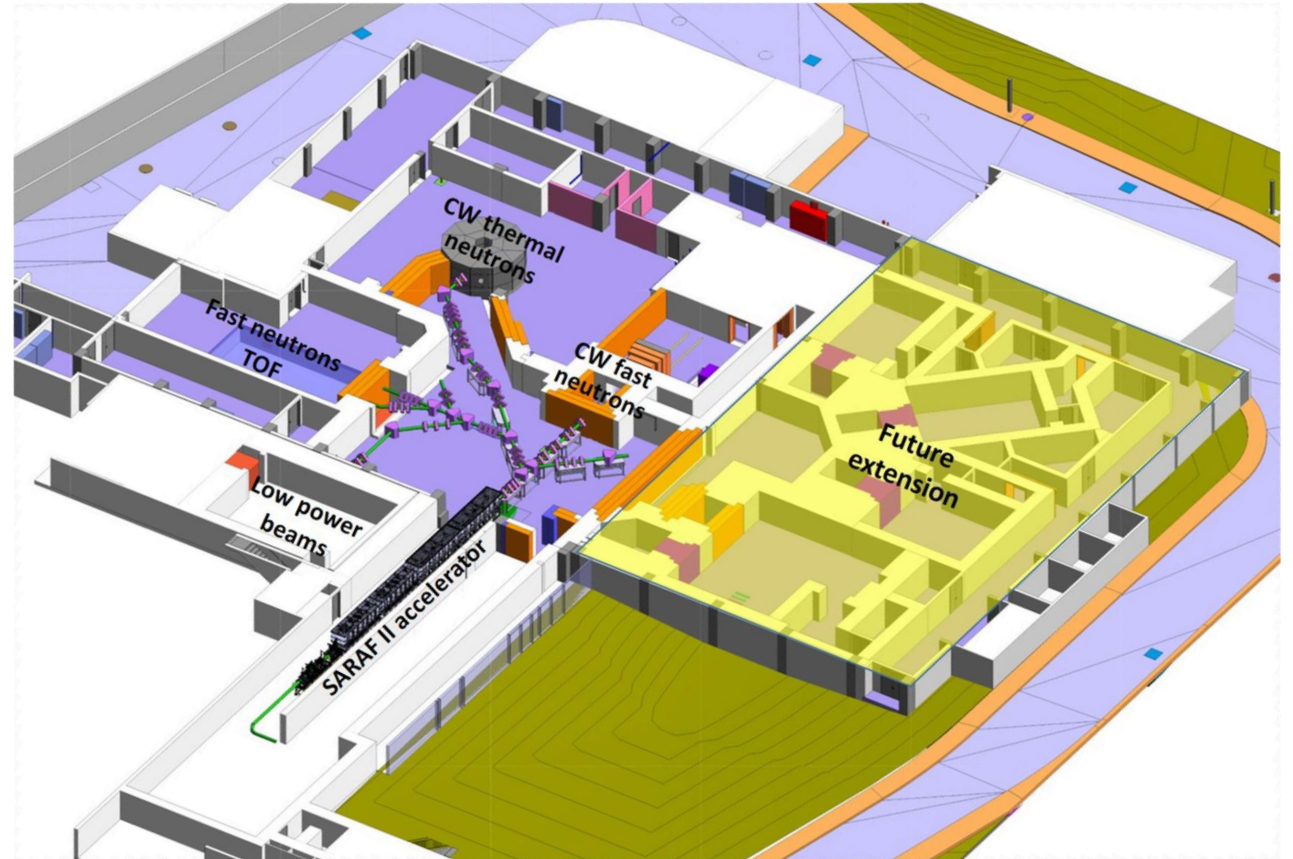
Work led by Michael Paul (HUJI)  
and Moshe Tessler (SARAF)

# SARAF Phase II

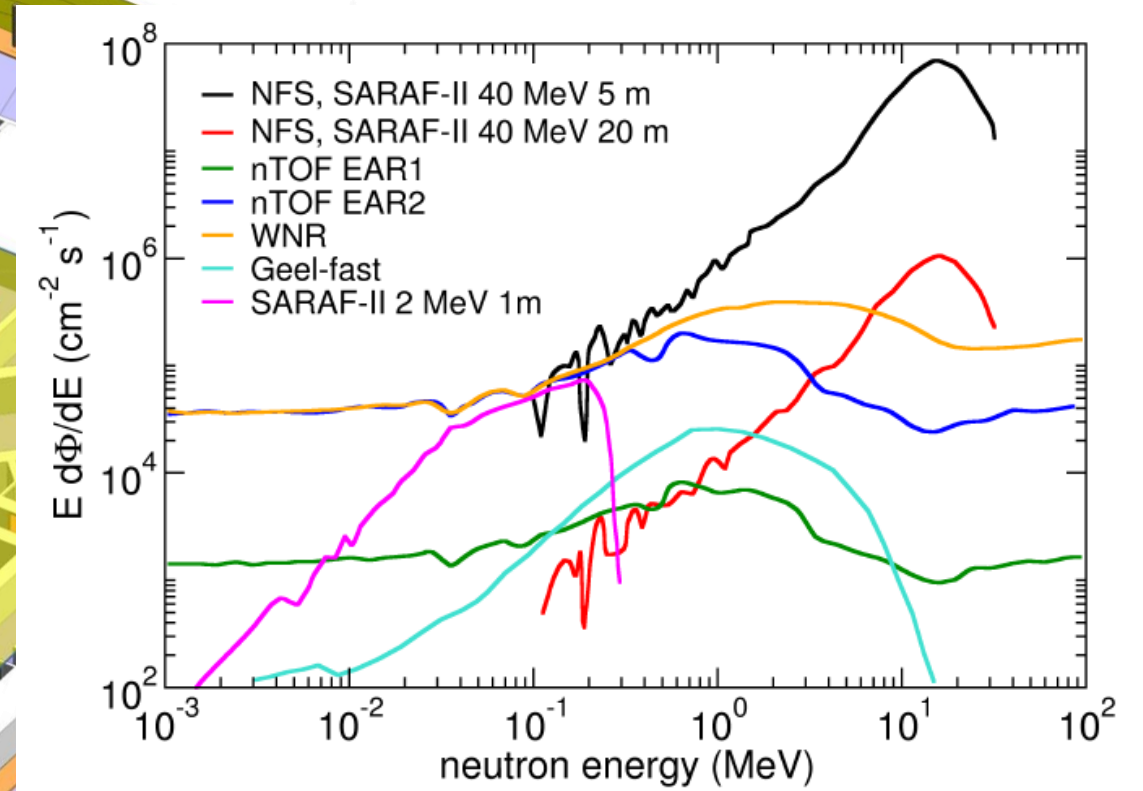
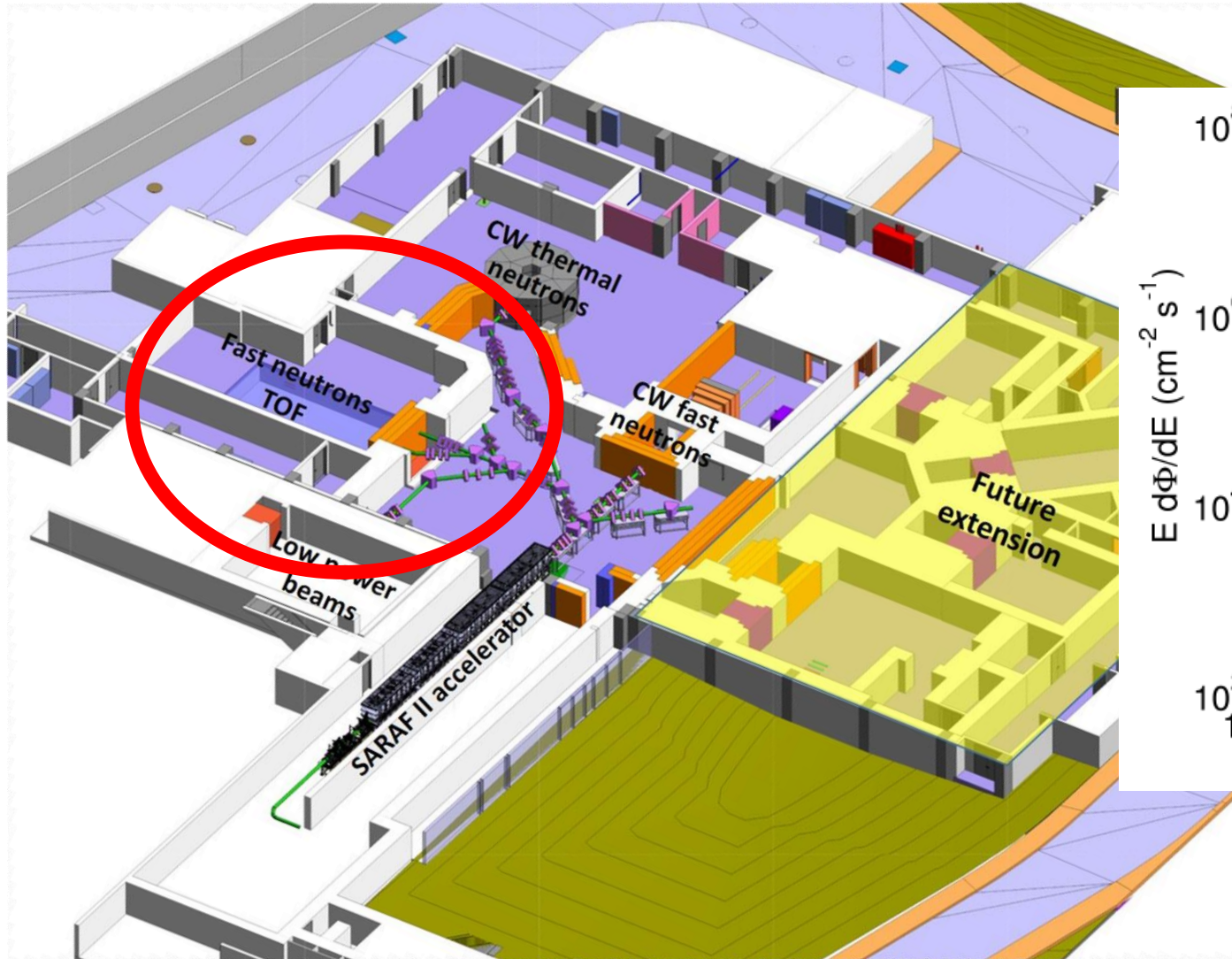
**Table 1.** SARAF-II beam top-level requirements.

Parameter	Value	Comment
Ion species	protons/deuterons	$M/q \leq 2$
Energy range	5–40 MeV deuterons 5–35 MeV protons	variable energy
Current range	0.04–5 mA	CW (and pulsed)
Operation	6000 hours/year	
Maintenance	hands-on	low beam loss

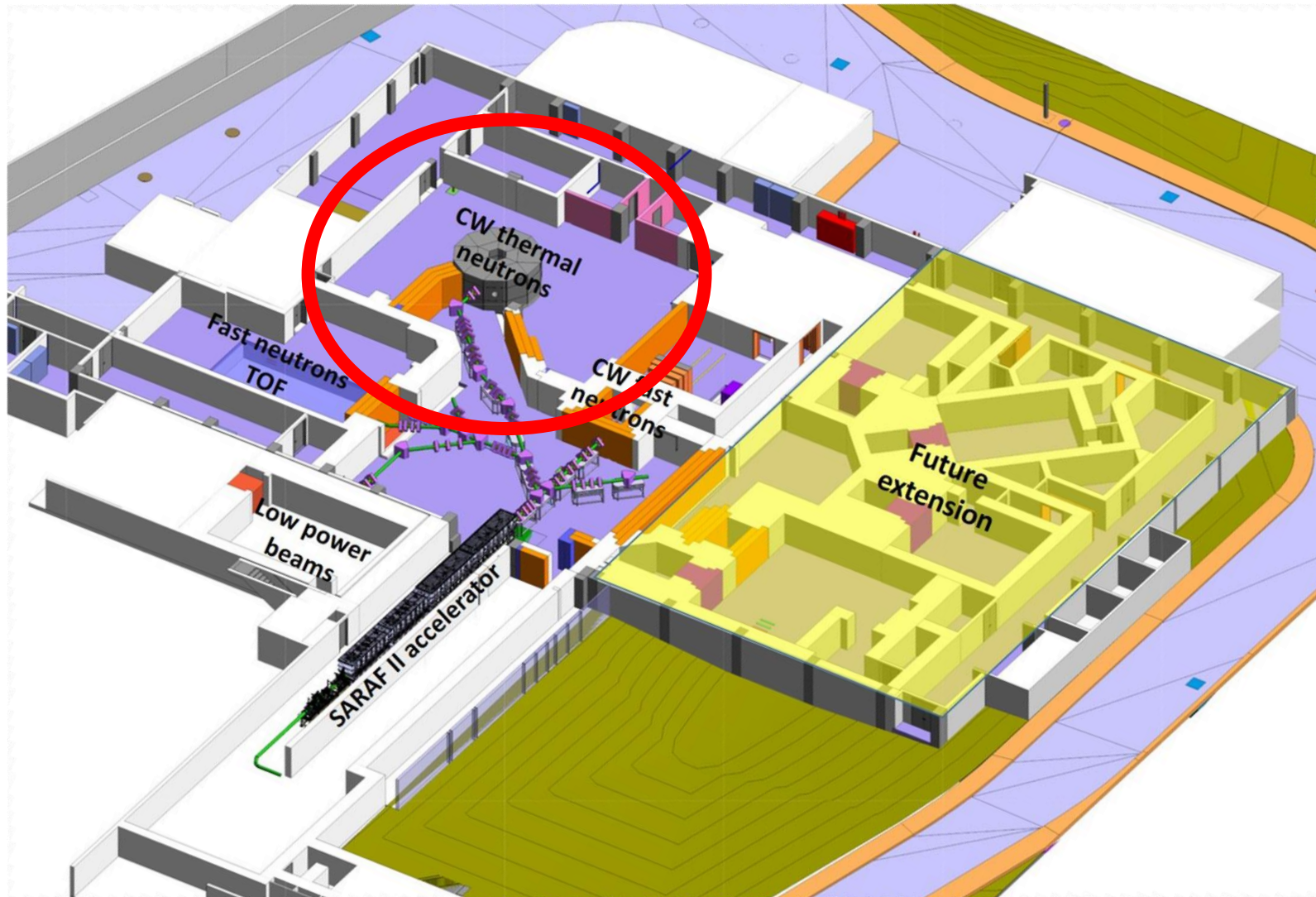
I. Mardor *et al.*, EPJA **54** (2018)



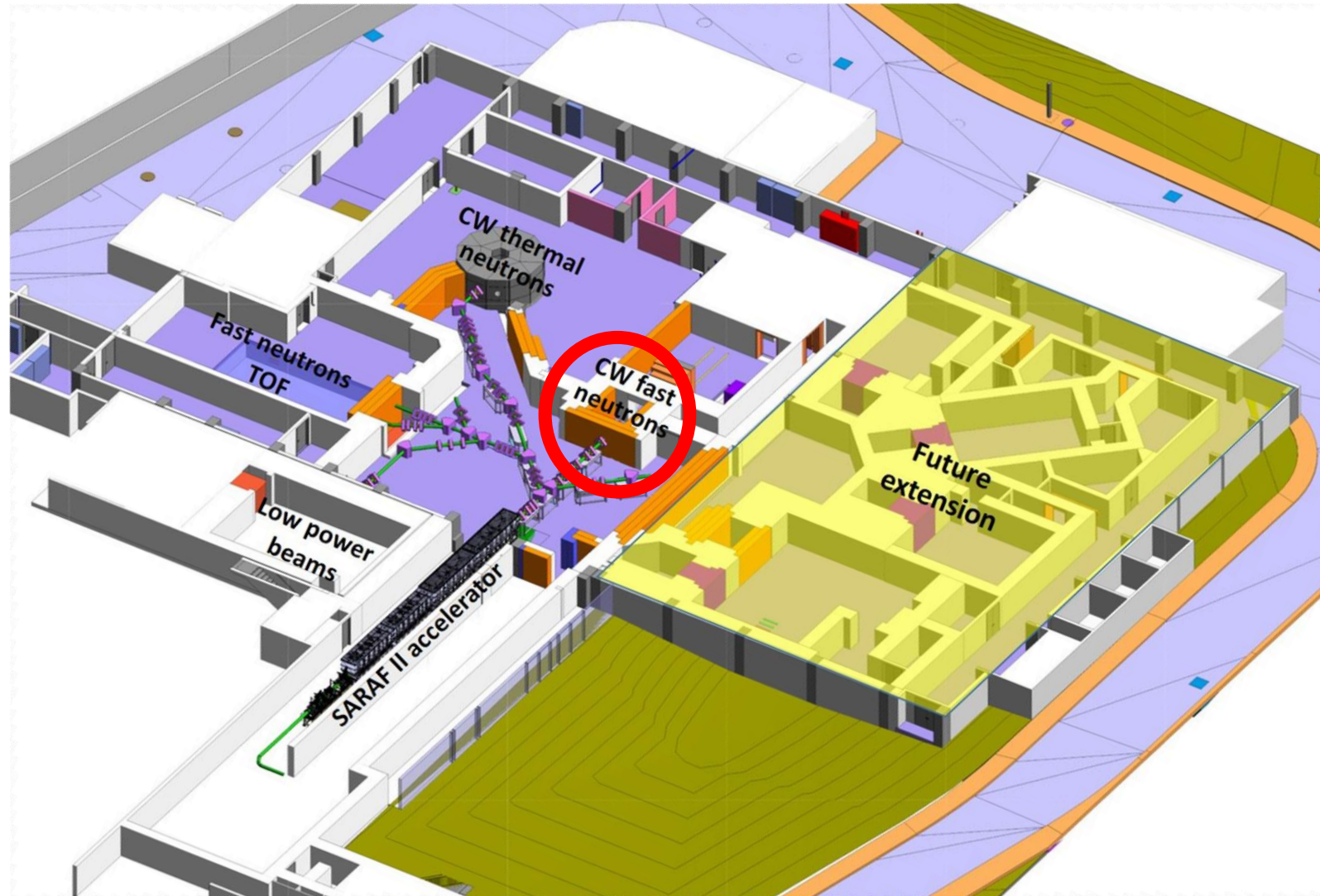
# SARAF Phase II – neutron TOF area



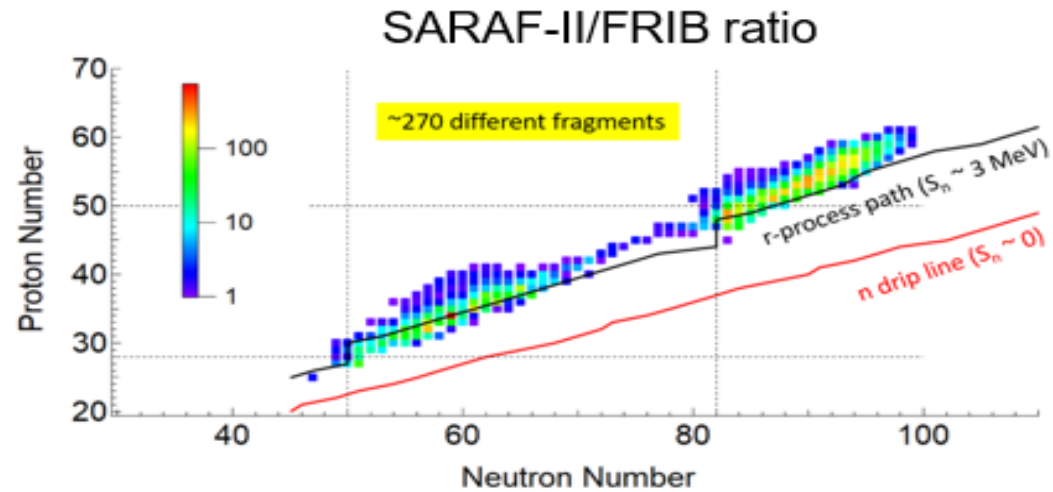
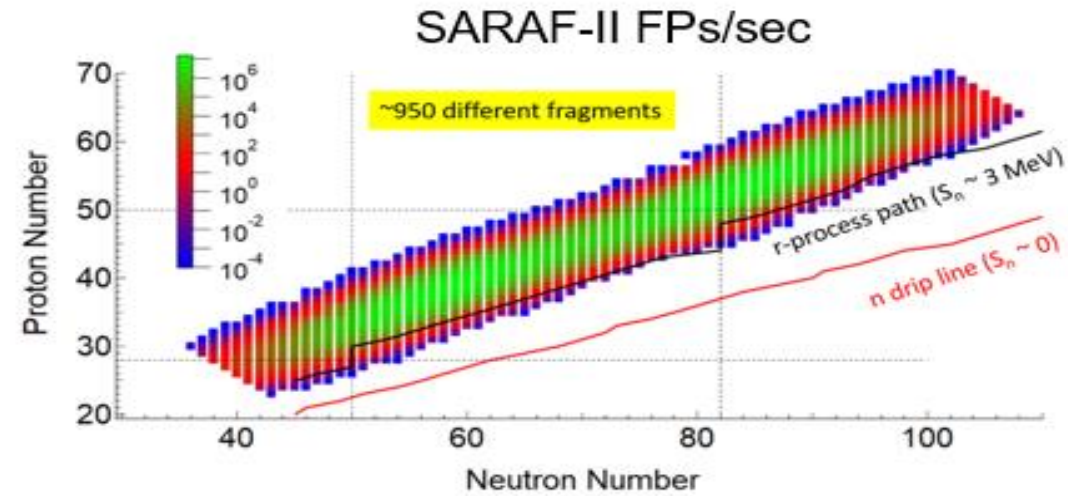
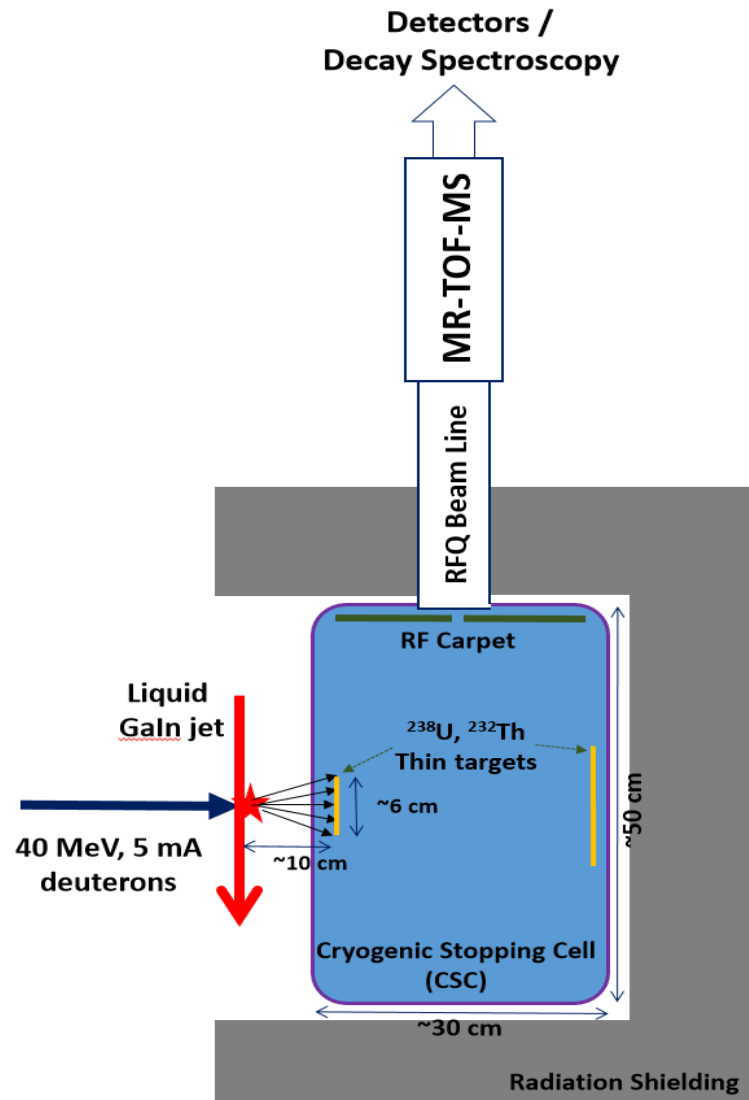
# SARAF Phase II – neutron camera



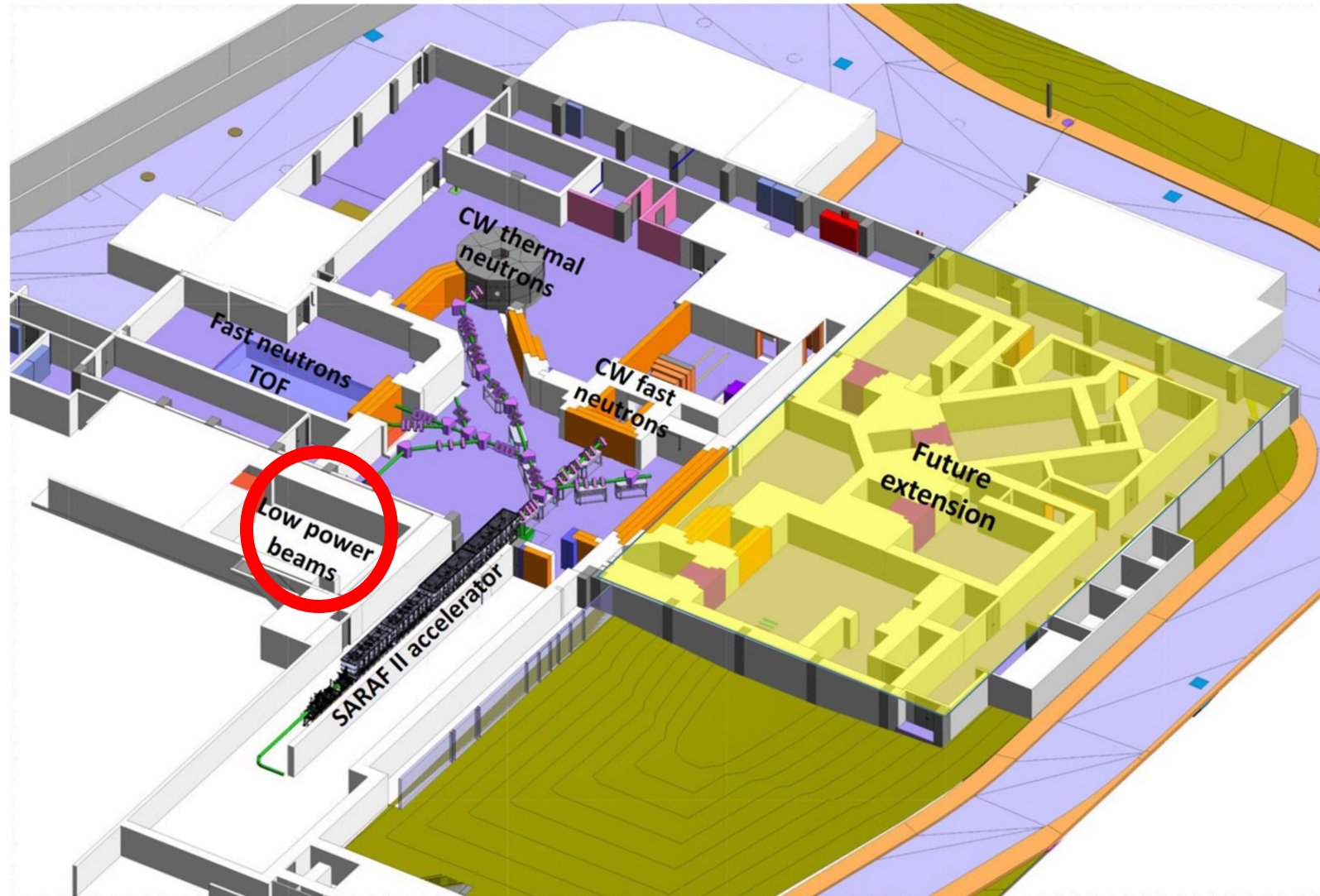
# SARAF Phase II – exotic nuclide facility



# SARONA - The SARaf exotic Nuclide fAcility



# SARAF Phase II – neutron induced cross-sections





# Why $(n,p)$ and $(n,\alpha)$ ?

- Supernovae and other explosive scenarios are fast.
- Nucleosynthesis follows sequences of capture-decay processes.
- Waiting points: Long-lived ( $>$ minutes) nuclei tend to create bottlenecks. High impact on nucleosynthesis. High abundancies.
- For proton-rich nuclei,  $(n,p)$  and  $(n,\alpha)$  enhance destruction of waiting-point nuclei.



# Data status

THE ASTROPHYSICAL JOURNAL, 653:474–489, 2006 December 10  
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## SENSITIVITY OF $p$ -PROCESS NUCLEOSYNTHESIS TO NUCLEAR REACTION RATES IN A $25 M_{\odot}$ SUPERNOVA MODEL

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 Received 2005 August 5; accepted 2006 August 14

TABLE 2  
 SELECTED  $(\gamma, p)$  OR  $(n, p)$  REACTIONS

Reactions		
$^{126}\text{Ba}(\gamma, p)^{125}\text{Cs}^*$	$^{92}\text{Mo}(\gamma, p)^{91}\text{Nb}^*$	$^{75}\text{Se}(n, p)^{75}\text{As}^*$
$^{110}\text{Sn}(\gamma, p)^{109}\text{In}^*$	$^{86}\text{Rb}(n, p)^{86}\text{Kr}^*$	$^{74}\text{Se}(\gamma, p)^{73}\text{As}^*$
$^{106}\text{Cd}(\gamma, p)^{105}\text{Ag}$	$^{85}\text{Sr}(n, p)^{85}\text{Rb}^*$	$^{76}\text{As}(n, p)^{76}\text{Ge}^*$
$^{104}\text{Cd}(\gamma, p)^{103}\text{Ag}$	$^{84}\text{Sr}(\gamma, p)^{83}\text{Rb}^*$	$^{75}\text{As}(\gamma, p)^{74}\text{Ge}^*$
$^{100}\text{Pd}(\gamma, p)^{99}\text{Rh}$	$^{78}\text{Kr}(\gamma, p)^{77}\text{Br}^*$	$^{73}\text{As}(\gamma, p)^{72}\text{Ge}$
$^{96}\text{Ru}(\gamma, p)^{95}\text{Tc}^*$	$^{77}\text{Se}(n, p)^{77}\text{As}$	$^{71}\text{Ge}(n, p)^{71}\text{Ga}$

NOTES.—These are reactions that, together with their respective inverse reactions, were found to exhibit the strongest influence on the final  $p$ -abundances. Their impact is illustrated in Fig. 15a. Particularly important rates are marked with an asterisk (\*).

**-NO DATA FOUND-**

# Data status

THE ASTROPHYSICAL JOURNAL, 729:46 (18pp), 2011 March 1

doi:[10.1088/0004-637X/729/1/46](https://doi.org/10.1088/0004-637X/729/1/46)

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## UNCERTAINTIES IN THE $\nu p$ -PROCESS: SUPERNOVA DYNAMICS VERSUS NUCLEAR PHYSICS

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*Received 2010 April 14; accepted 2010 December 27; published 2011 February 8*

### ABSTRACT

We examine how the uncertainties involved in supernova dynamics, as well as in nuclear data inputs, affect the  $\nu p$ -process in the neutrino-driven winds. For the supernova dynamics, we find that the wind termination by the preceding dense ejecta shell, as well as the electron fraction ( $Y_{e,3}$ ; at  $3 \times 10^9$  K), plays a crucial role. A wind termination within the temperature range of  $(1.5\text{--}3) \times 10^9$  K greatly enhances the efficiency of the  $\nu p$ -process. This implies that the early wind phase, when the innermost layer of the preceding supernova ejecta is still  $\sim 200\text{--}1000$  km from the center, is most relevant to the  $\nu p$ -process. The outflows with  $Y_{e,3} = 0.52\text{--}0.60$  result in the production of the  $p$ -nuclei up to  $A = 108$  with interesting amounts. Furthermore, the  $p$ -nuclei up to  $A = 152$  can be produced if  $Y_{e,3} = 0.65$  is achieved. For the nuclear data inputs, we test the sensitivity to the rates relevant to the breakout from the  $p$ - $p$  chain region ( $A < 12$ ), to the  $(n, p)$  rates on heavy nuclei, and to the nuclear masses along the  $\nu p$ -process pathway. **We find that a small variation of the rates of triple- $\alpha$  and of the  $(n, p)$  reaction on  $^{56}\text{Ni}$  leads to a substantial change in the  $p$ -nuclei production.** We also find that  $^{96}\text{Pd}$  ( $N = 50$ ) on the  $\nu p$ -process path plays a role as a second seed nucleus for the production of heavier  $p$ -nuclei. The uncertainty in the nuclear mass of  $^{82}\text{Zr}$  can lead to a factor of two reduction in the abundance of the  $p$ -isotope  $^{84}\text{Sr}$ .

**-NO DATA FOUND-**

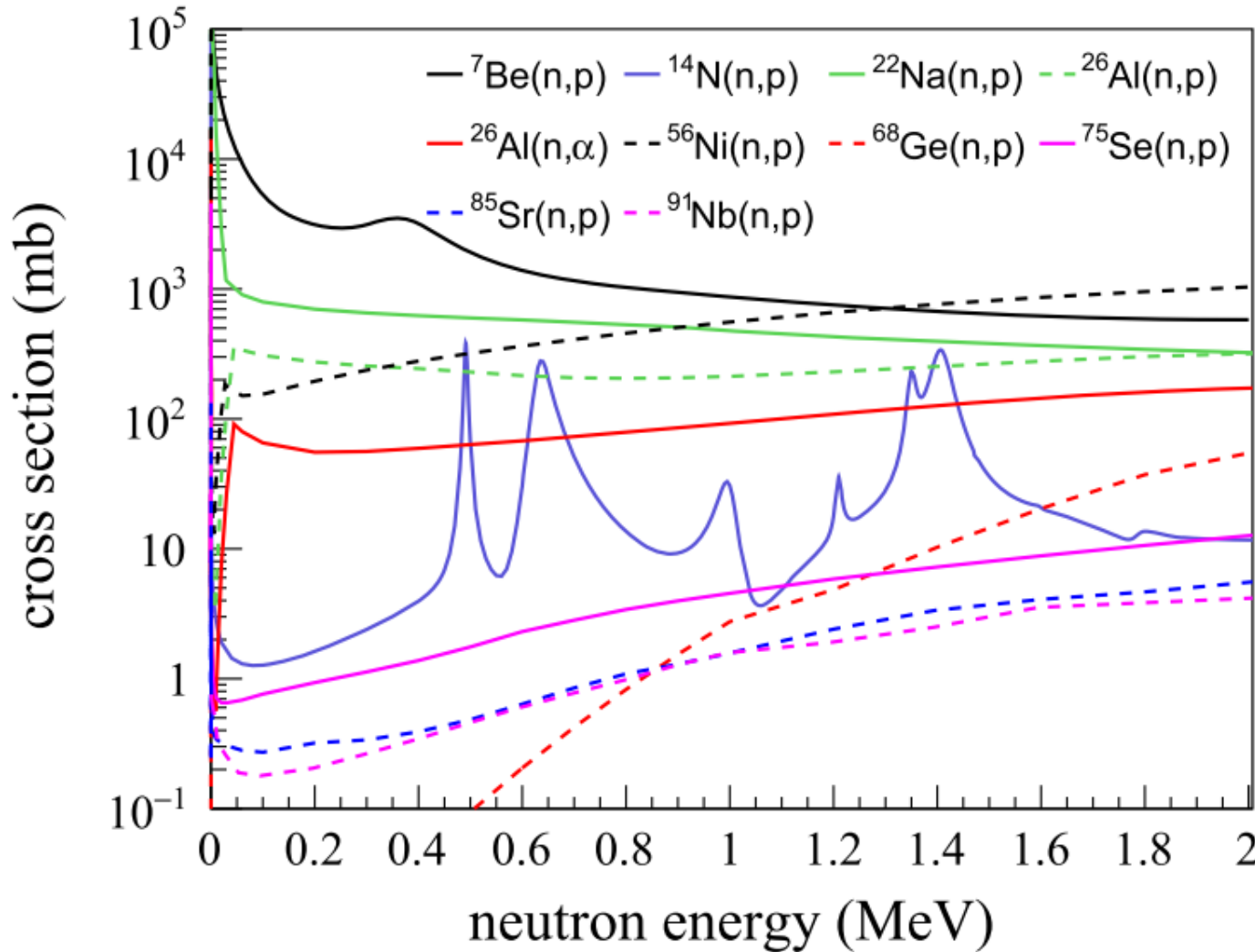
# Data status

THE ASTROPHYSICAL JOURNAL SUPPLEMENT SERIES, 19  
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THE EFFECTS OF THERMONUCLEAR REACTIONS IN MASSIVE STARS

CHRISTIAN ILIADIS<sup>1,2</sup>, ANDREW L. COHEN<sup>1</sup>, ANDREW W. CLAYTON<sup>1</sup>

proximate order of  
 Tables 2, 4, and 6  
 five reactions,  $^{26}\text{Al}(n,\alpha)^{23}\text{Na}$ ,  $^{26}\text{Al}(n,p)^{26}\text{Mg}$ , and  
 future measurements  
 which the rate needs  
 Ne/C burning,  $\approx 1.4$



7 C 104, L022803 (2021)

massive stars: Study of the key  $^{26}\text{Al}(n, p)$  reaction

W C 104, L032803 (2021)

massive stars: Study of the key  $^{26}\text{Al}(n, \alpha)$  reaction

and  $^{26}\text{Al}(n, \alpha)$   
 to  $\sim 150$  keV

21). PRC 104 L022803

U. LEUTNER ET AL. (2021). PRC 104 L032803

**Motivation:**  $(n,p)$  and  $(n,\alpha)$  data scarce, especially for unstable proton-rich isotopes. ( $^{26}\text{Al}(n,p)$ ,  $^{56}\text{Ni}(n,p)$  and more).

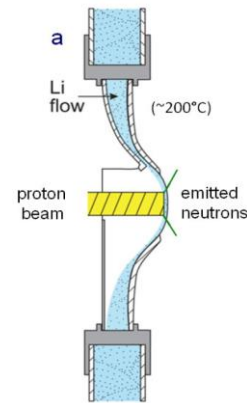
**Goal:** Establish an apparatus for  $(n,p)$  and  $(n,\alpha)$  measurements on stable and unstable isotopes at explosive stellar temperatures ( $\sim 1.5\text{-}3.5$  GK,  $\sim 10\text{-}2000$  keV).

### Protons - SARAF



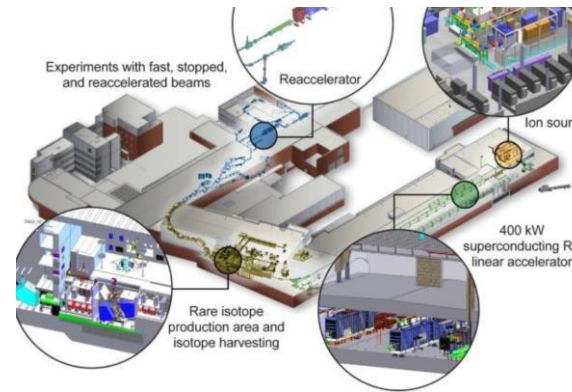
2-5 MeV protons  
Beam current > 1 mA.

### Neutrons - LiLiT



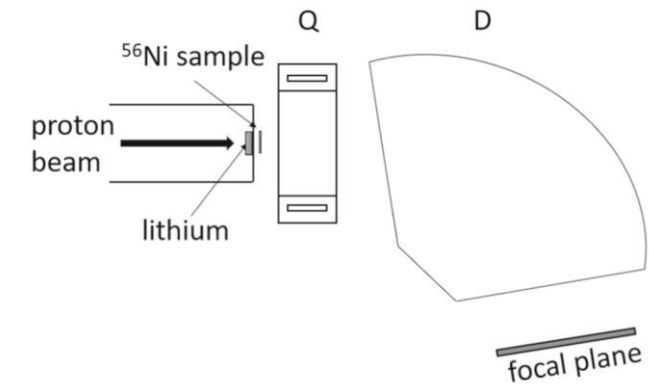
keV neutrons via  $^7\text{Li}(p,n)$   
Operational

### Targets - FRIB



Isotope Harvesting project  
E. Abel *et al.* J. Phys. G 46 (2019)

### Detection

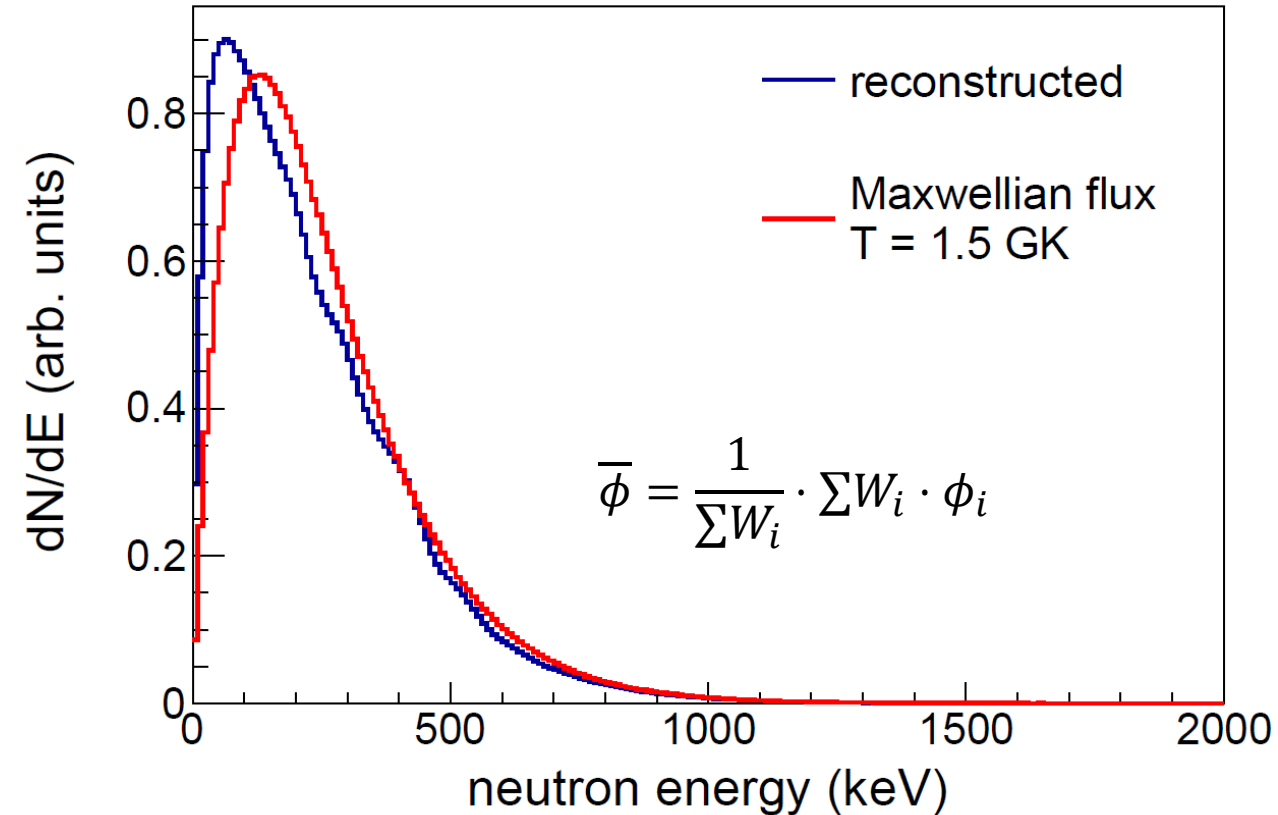
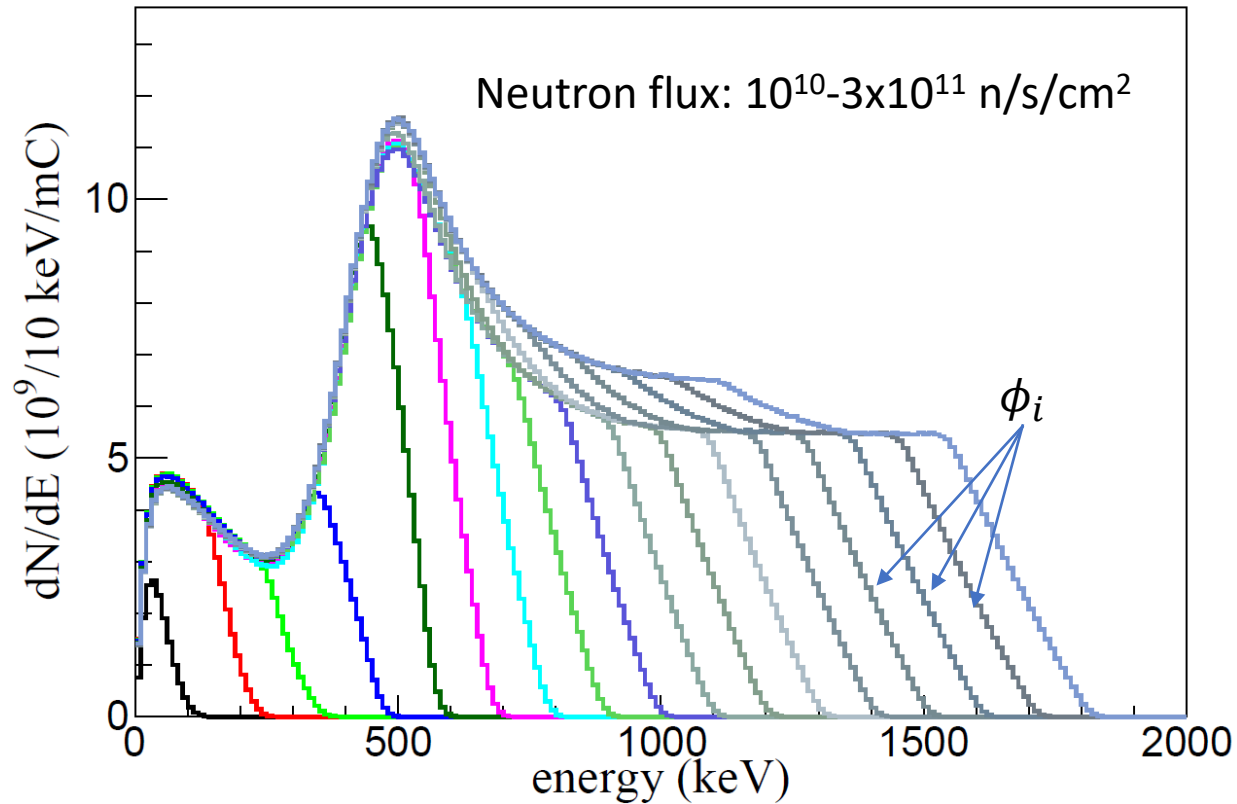


Design stages  
efficiency  $\sim 0.5\%$

->See also a talk by Pelagia Tsintari about  $(n,p)$  studies with SECAR!

# Neutrons for studies of explosive nucleosynthesis

neutron energy spectra for different proton energies  
 $E_p = 1900\text{-}3600$  keV



$^{56}\text{Ni}(n,p)$  ( $t_{1/2} = 6$  d)

Sample size: 1 mCi ( $3 \times 10^{13}$  atoms)

Detection rate: 1-150 counts/hour (TENDL2019)

# Effect of differences in the neutron spectrum

reaction	calculated cross section (mb)					
	T = 1.5 GK			T = 3.5 GK		
	Maxwell.	reconst.	corr.	Maxwell.	reconst.	corr.
$^{14}\text{N}(n, p)$	9.10	9.08	1.00	27.1	27.9	0.97
$^{26}\text{Al}(n, p)$	266	261	1.02	242	241	1.00
$^{26}\text{Al}(n, \alpha)$	61.2	60.6	1.01	73.3	72.9	1.01
$^{40}\text{K}(n, p)$	8.21	9.17	0.90	11.47	11.89	0.96
$^{40}\text{K}(n, \alpha)$	22.63	26.49	0.85	25.34	27.59	0.92



# Summary

- LiLiT @ SARAF is a high-intensity neutron source for s-process measurements, with a neutron flux on the order of  $10^{10}$  n/s on the sample.
- During the operation of SRARAF Phase-I, LiLiT produced many MACS values of relevance for the s-process, and is expected to continue conducting similar measurement in the future.
- SARAF Phase-II will provide new opportunities for nuclear nucleosynthesis studies.
- LiLiT will also be used to produce higher energy neutrons, which will allow direct (n,p) and (n, $\alpha$ ) cross section at explosive stellar temperatures.

Thanks for listening!