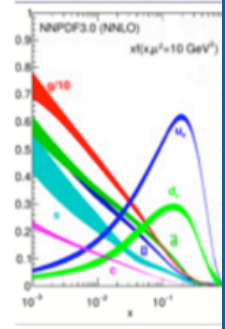


W. K. Kellogg  
Biological Station  
MICHIGAN STATE UNIVERSITY



Parton Distributions and Lattice Calculations (PDFLattice 2019)

25-27 September 2019

# GPDs – Overview

Jianwei Qiu

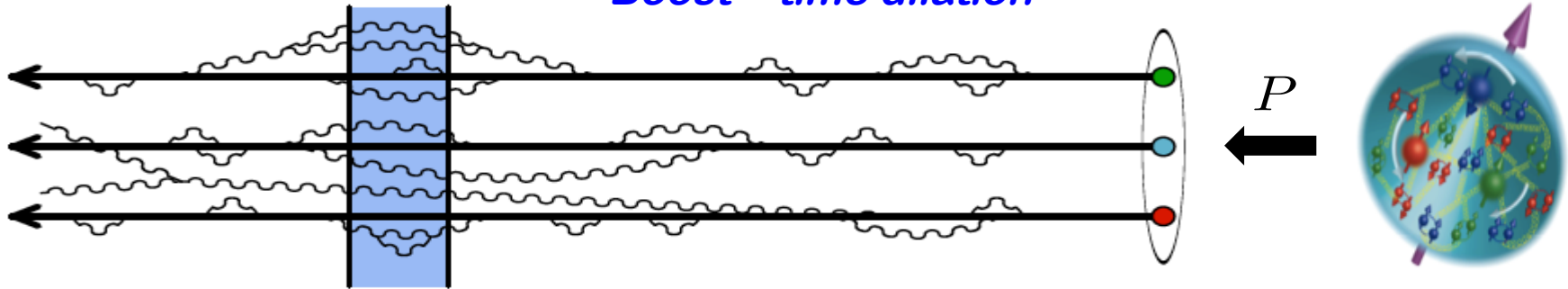
Theory Center, Jefferson Lab

September 26, 2019

# How to “see” 3D partonic structure of hadrons?

- Hard probes to “catch” the quantum fluctuation:

*Boost = time dilation*

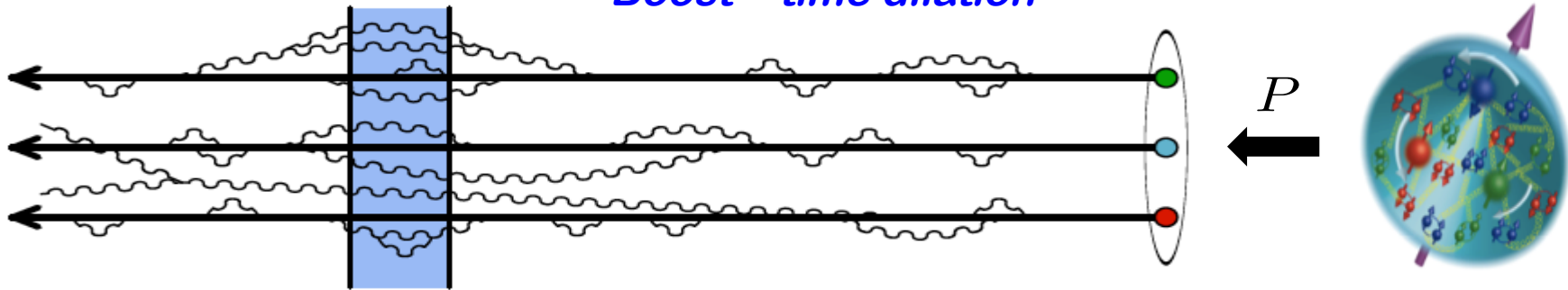


*Hard probe ( $t \sim 1/Q < fm$ )*  *Probability to “catch” the parton!*

# How to “see” 3D partonic structure of hadrons?

- Hard probes to “catch” the quantum fluctuation:

*Boost = time dilation*



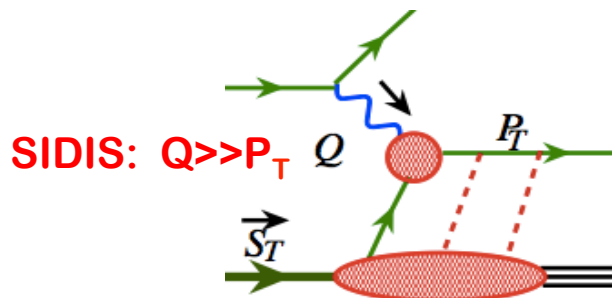
*Hard probe ( $t \sim 1/Q < fm$ ) → Probability to “catch” the parton!*

- Observables with two momentum scales:

$$Q_1 \gg Q_2 \sim 1/R \sim \Lambda_{QCD}$$

✦ Hard scale:  $Q_1$  “see” particle nature of “partons”

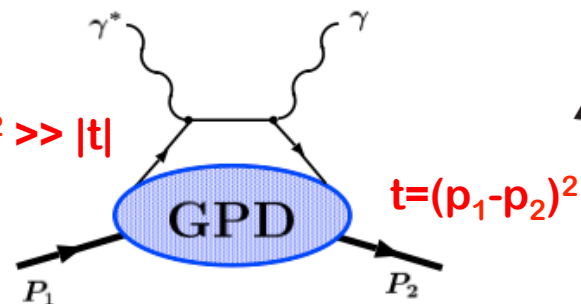
✦ Soft scale:  $Q_2$  “sensitive” to the fermi-scale structure



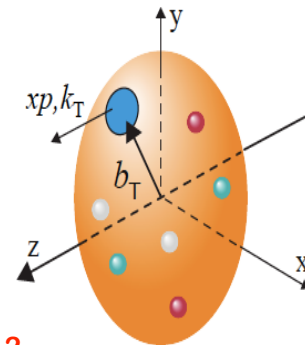
SIDIS:  $Q \gg P_T$

*TMDs – Confined motion*

DVCS:  $Q^2 \gg |t|$



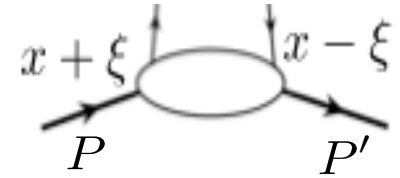
*GPDs – Spatial imaging*



# Definition of GPDs

## □ Quark “form factor”:

$$F^q = \frac{1}{2} \int \frac{dz^-}{2\pi} e^{ixP^+z^-} \langle P' | \bar{q}(-\frac{1}{2}z) \gamma^+ q(\frac{1}{2}z) | P \rangle \Big|_{z^+=0, z=0}$$

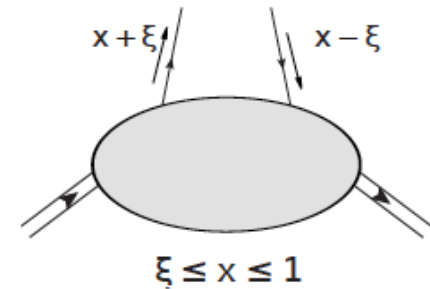
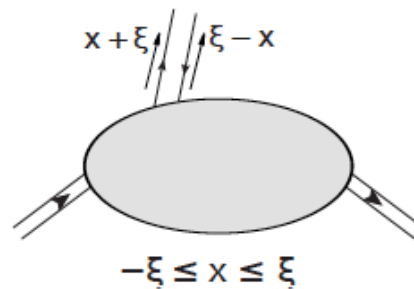
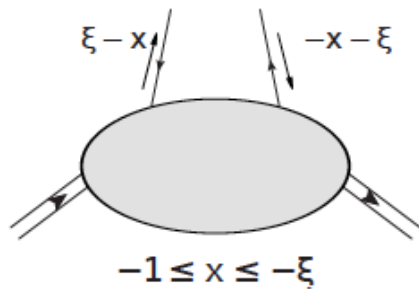


$$= \frac{1}{2P^+} \left[ H^q(x, \xi, t) \bar{u}(p') \gamma^+ u(p) + E^q(x, \xi, t) \bar{u}(p') \frac{i\sigma^{+\alpha} \Delta_\alpha}{2m} u(p) \right]$$

with  $\xi = (P' - P) \cdot n/2$  and  $t = (P' - P)^2 \Rightarrow -\Delta_\perp^2$  if  $\xi \rightarrow 0$

Gauge link:  $W[a, b] = P \exp\left( ig \int_b^a dx^- A^+(x^- n_-) \right)$

## □ Kinematics:



Two more for quarks:  $\tilde{H}_q(x, \xi, t, Q)$ ,  $\tilde{E}_q(x, \xi, t, Q)$

with  $\gamma \cdot n \rightarrow \gamma \cdot n \gamma_5$

# Definition of GPDs

## □ Gluon “form factor”:

Mueller et al., 94;  
Ji, 96;  
Radyushkin, 96

$$F^g = \frac{1}{P^+} \int \frac{dz^-}{2\pi} e^{ixP^+z^-} \langle p' | G^{+\mu}(-\frac{1}{2}z) G_{\mu}^+(\frac{1}{2}z) | p \rangle \Big|_{z^+=0, \mathbf{z}=0}$$

$$= \frac{1}{2P^+} \left[ H^g(x, \xi, t) \bar{u}(p') \gamma^+ u(p) + E^g(x, \xi, t) \bar{u}(p') \frac{i\sigma^{+\alpha} \Delta_\alpha}{2m} u(p) \right]$$

Two more for gluons:  $\tilde{H}^g(x, \xi, t)$   $\tilde{E}^g(x, \xi, t)$

with the two gluon field strength contracted anti-symmetrically

## □ Forward limit – connection to collinear PDFs:

$$H^q(x, 0, 0) = q(x), \quad \tilde{H}^q(x, 0, 0) = \Delta q(x) \quad \text{for } x > 0$$

$$H^q(x, 0, 0) = -\bar{q}(-x), \quad \tilde{H}^q(x, 0, 0) = \Delta \bar{q}(-x) \quad \text{for } x < 0$$

$$H^g(x, 0, 0) = xg(x), \quad \tilde{H}^g(x, 0, 0) = x\Delta g(x) \quad \text{for } x > 0$$

The factorization scale dependence is suppressed

# Properties of GPDs

## □ Connection to Dirac and Pauli form factors:

$$\int_{-1}^1 dx H^q(x, \xi, t) = F_1^q(t), \quad \int_{-1}^1 dx E^q(x, \xi, t) = F_2^q(t)$$

Where  $\langle p' | \bar{q}(0) \gamma^\mu q(0) | p \rangle = \bar{u}(p') \left[ F_1^q(t) \gamma^\mu + F_2^q(t) \frac{i\sigma^{\mu\alpha} \Delta_\alpha}{2m} \right] u(p)$

And the axial and pseudoscalar version:

$$\int_{-1}^1 dx \tilde{H}^q(x, \xi, t) = g_A^q(t), \quad \int_{-1}^1 dx \tilde{E}^q(x, \xi, t) = g_P^q(t)$$

Where  $\langle p' | \bar{q}(0) \gamma^\mu \gamma_5 q(0) | p \rangle = \bar{u}(p') \left[ g_A^q(t) \gamma^\mu \gamma_5 + g_P^q(t) \frac{\gamma_5 \Delta^\mu}{2m} \right] u(p)$

## □ Some symmetry properties:

$$H^g(x, \xi, t) = H^g(-x, \xi, t) \quad E^g(x, \xi, t) = E^g(-x, \xi, t)$$

$$\tilde{H}^g(x, \xi, t) = -\tilde{H}^g(-x, \xi, t) \quad \tilde{E}^g(x, \xi, t) = -\tilde{E}^g(-x, \xi, t)$$

$$H^{q,g}(x, \xi, t) = H^{q,g}(x, -\xi, t), \dots$$

$$H^{q,g}(x, \xi, t)^* = H^{q,g}(x, -\xi, t), \dots$$

**GPDs are real value functions**

# Properties of GPDs

## □ QCD energy-momentum tensor:

$$T^{\mu\nu} = \sum_{i=q,g} T_i^{\mu\nu} \quad \text{with} \quad T_q^{\mu\nu} = \bar{q} \gamma^{(\mu} \overleftrightarrow{D}^{\nu)} q$$

$$T_g^{\mu\nu} = G^{\mu\alpha} G_{\alpha}{}^{\nu} + \frac{1}{4} g^{\mu\nu} G^{\alpha\beta} G_{\alpha\beta}$$

## □ Form factors:

$$\langle p' | T_{q,g}^{\mu\nu} | p \rangle = A_{q,g}(t) \bar{u} P^{(\mu} \gamma^{\nu)} u + B_{q,g}(t) \bar{u} \frac{P^{(\mu} i \sigma^{\nu)\alpha} \Delta_{\alpha}}{2m} u$$

$$+ C_{q,g}(t) \frac{\Delta^{\mu} \Delta^{\nu} - g^{\mu\nu} \Delta^2}{m} \bar{u} u + \bar{C}_{q,g}(t) m g^{\mu\nu} \bar{u} u$$

## □ Light-cone helicity operator:

$$J^3 = \int dx^- d^2\mathbf{x} M^{+12}(x) \quad \text{with} \quad M^{\alpha\mu\nu} = T^{\alpha\nu} x^{\mu} - T^{\alpha\mu} x^{\nu}$$

## □ Connection to the proton spin:

$$\langle J_q^3 \rangle = \frac{1}{2} [A_q(0) + B_q(0)], \quad \langle J_g^3 \rangle = \frac{1}{2} [A_g(0) + B_g(0)]$$

Ji, PRL78, 1997

$$A_q(t) + B_q(t) = \int_{-1}^1 dx x [H_q(x, \xi, t) + E_q(x, \xi, t)]$$

$$A_g(t) + B_g(t) = \int_0^1 dx [H_g(x, \xi, t) + E_g(x, \xi, t)]$$

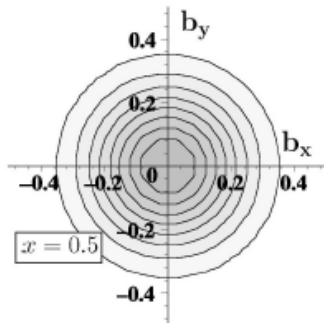
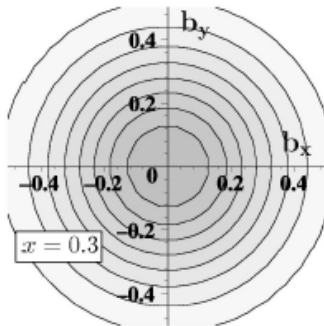
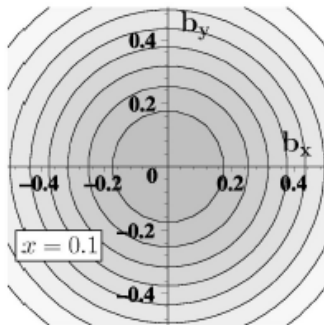
# Spatial imaging from GPDs

M. Burkardt, PRD 2000

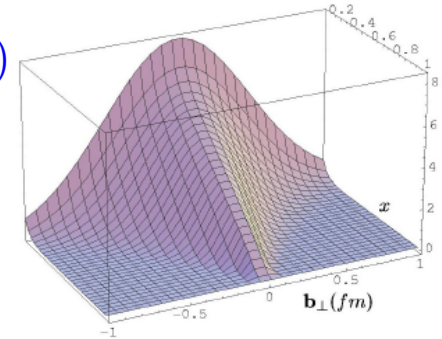
## □ Impact parameter dependent quark distribution:

$$q(x, b_{\perp}, Q) = \int d^2\Delta_{\perp} e^{-i\Delta_{\perp} \cdot b_{\perp}} H_q(x, \xi = 0, t = -\Delta_{\perp}^2, Q)$$

$q(x, b_{\perp})$  for unpol. p



## Unpolarized proton



- $F_1(-\Delta_{\perp}^2) = \int dx H(x, 0, -\Delta_{\perp}^2)$
  - $x$  = momentum fraction of the quark
  - $b_{\perp}$  relative to  $\perp$  center of momentum
  - small  $x$ : large 'meson cloud'
  - larger  $x$ : compact 'valence core'
  - $x \rightarrow 1$ : active quark becomes center of momentum
- $\vec{b}_{\perp} \rightarrow 0$  (narrow distribution) for  $x \rightarrow 1$

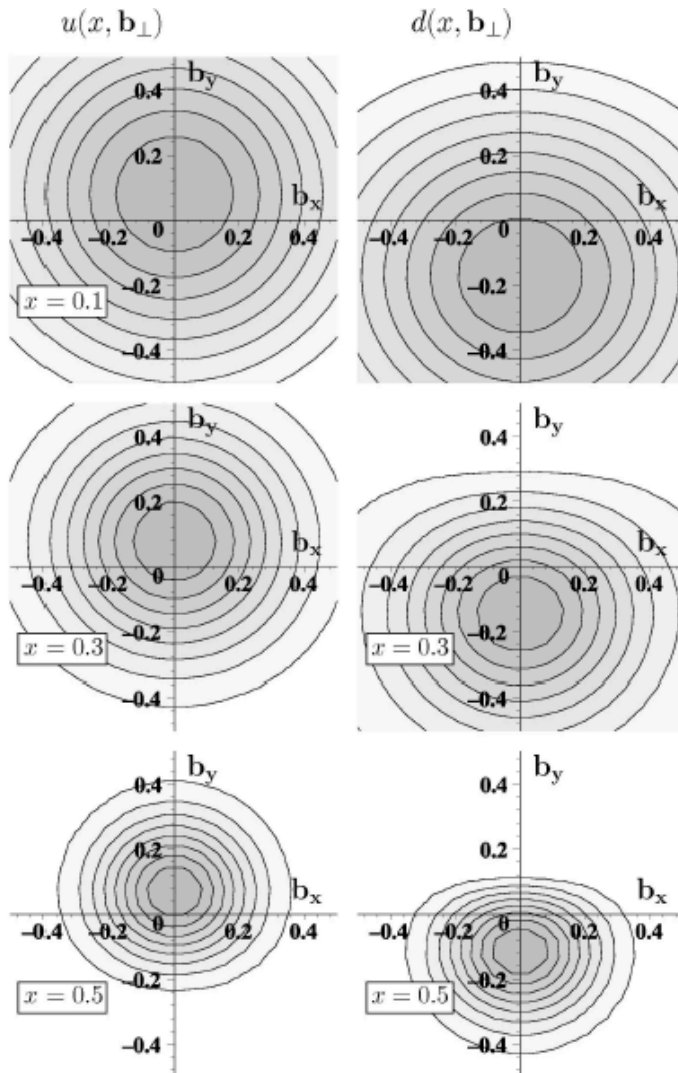


# Spatial imaging from GPDs

M. Burkardt, PRD 2000

## □ Impact parameter dependent quark distribution:

Proton polarized in +x direction



$$q(x, \mathbf{b}_\perp) = \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H_q(x, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp} - \frac{1}{2M} \frac{\partial}{\partial b_y} \int \frac{d^2 \Delta_\perp}{(2\pi)^2} E_q(x, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp}$$

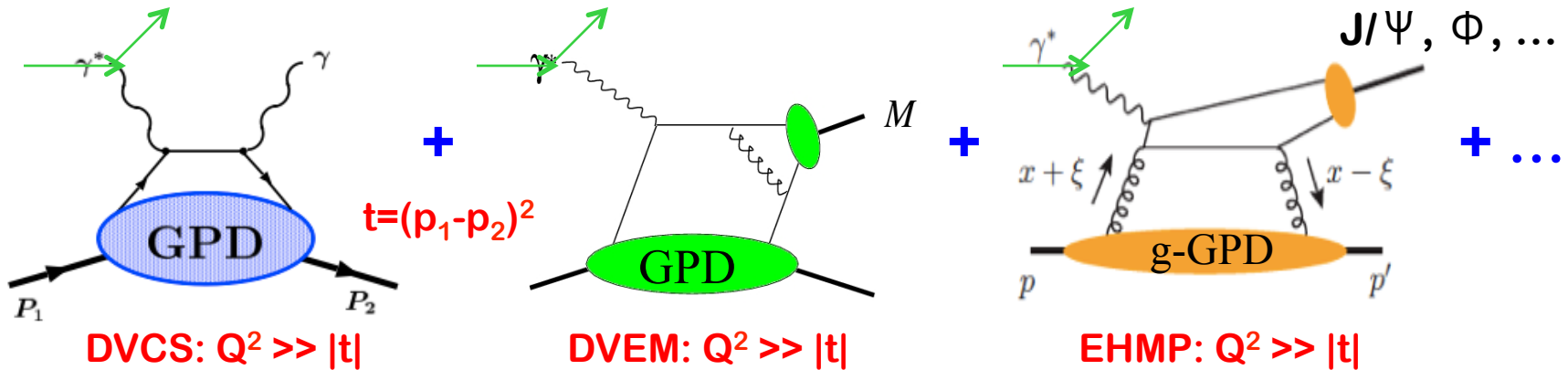
Physics: relevant density in DIS is  $j^+ \equiv j^0 + j^3$  and left-right asymmetry from  $j^3$

Sign and magnitude of the averaged shift related to the hadron's magnetic moment:

$$\begin{aligned} \langle b_y^q \rangle &\equiv \int dx \int d^2 b_\perp q(x, \mathbf{b}_\perp) b_y \\ &= \frac{1}{2M} \int dx E_q(x, 0) = \frac{\kappa_q}{2M} \end{aligned}$$

# Hunting for GPDs – Exclusive DIS

## Experimental access to GPDs:



## Much more complicated – $(x, \xi, t)$ variables:

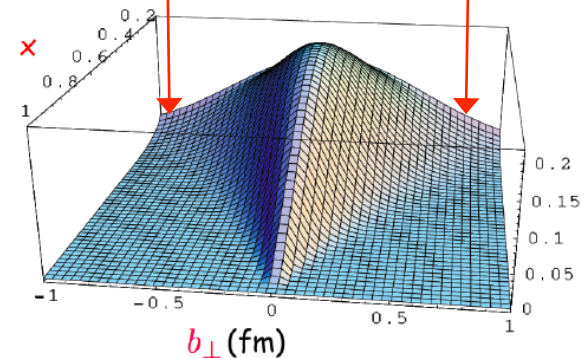
✧ Challenge to derive GPDs from data

## GPDs could tell us:

- ✧ Orbital contribution to proton's spin
- ✧ Proton radius of quark & gluon density
- ✧ Hints for color confining radius/mechanism
- ✧ Origin of nuclear force, ...
- ✧ ...

How far does glue density spread?

How fast does glue density fall?



# QCD factorization

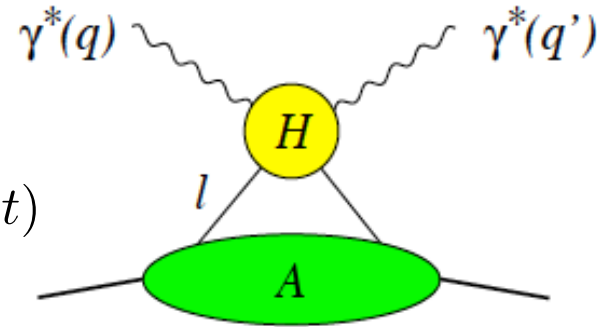
## □ Deep Virtual Compton Scattering (DVCS):

$$\gamma^*(q) + p(p) \rightarrow \gamma^*(q') + p(p')$$

## □ Factorization:

$$\mathcal{A}(\gamma^* p \rightarrow \gamma p) = \sum_i \int_{-1}^1 dx T^i(x, \xi, \rho, Q^2) F^i(x, \xi, t)$$

$$\rho = -(q + q')^2 / 2(p + p') \cdot (q + q')$$



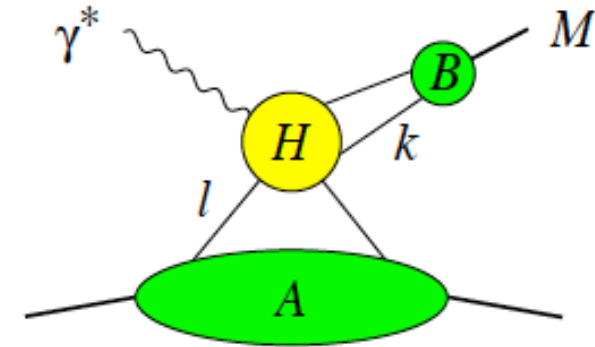
## □ Deep Virtual Meson Production (DVMP):

$$\gamma^*(q) + p(p) \rightarrow M(q') + p(p')$$

## □ Factorization:

$$\mathcal{A}(\gamma_L^* p \rightarrow M_L p) = \frac{1}{Q} \sum_{ij} \int_{-1}^1 dx$$

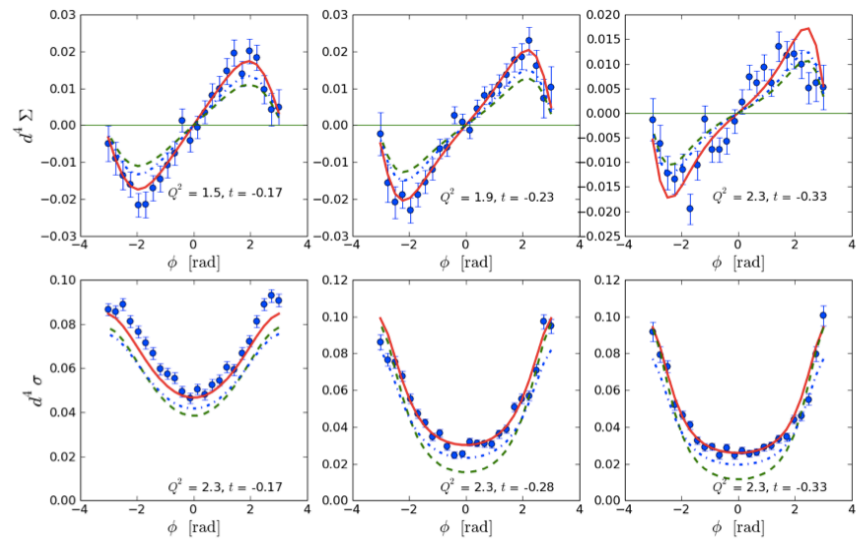
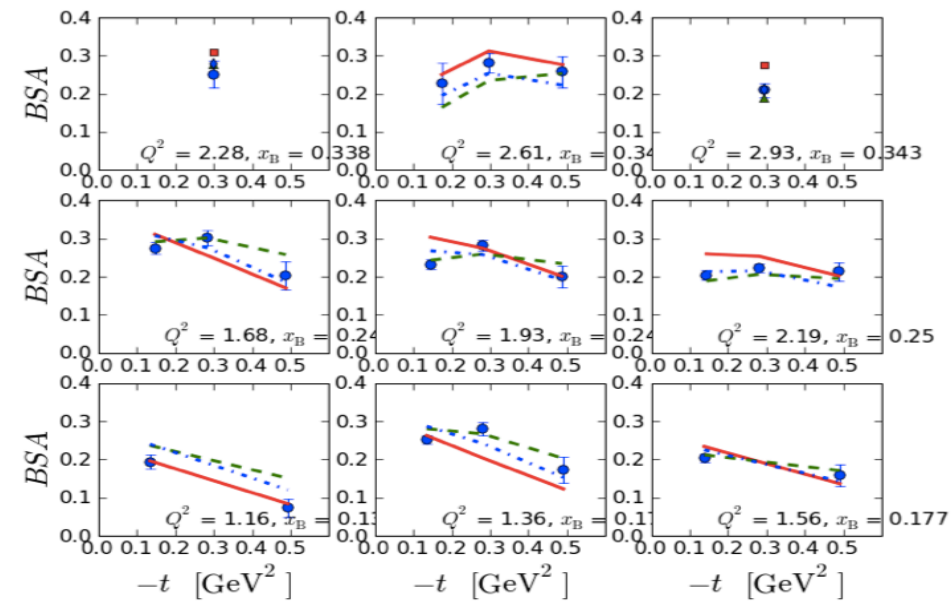
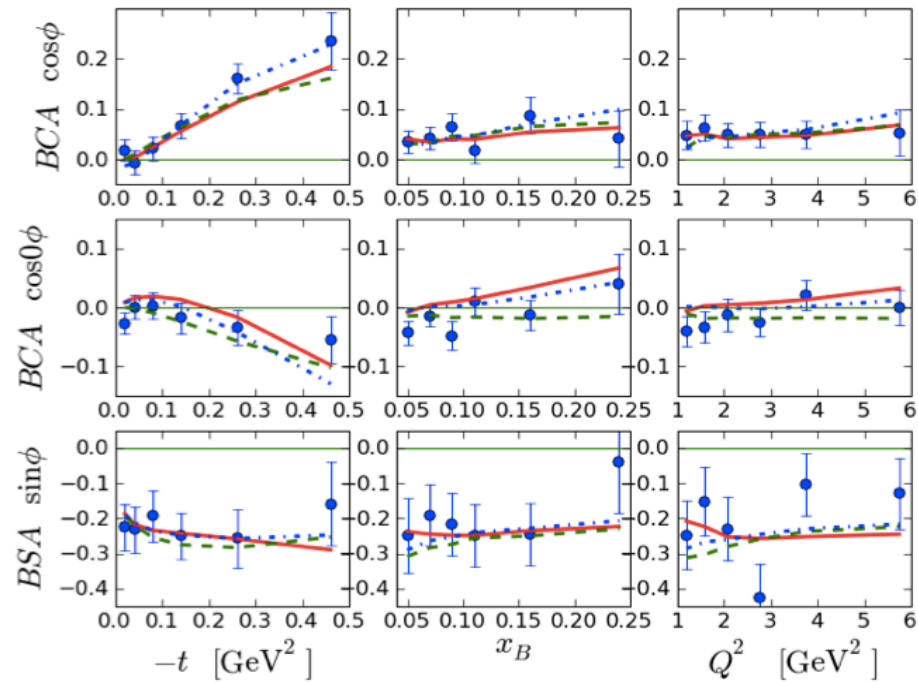
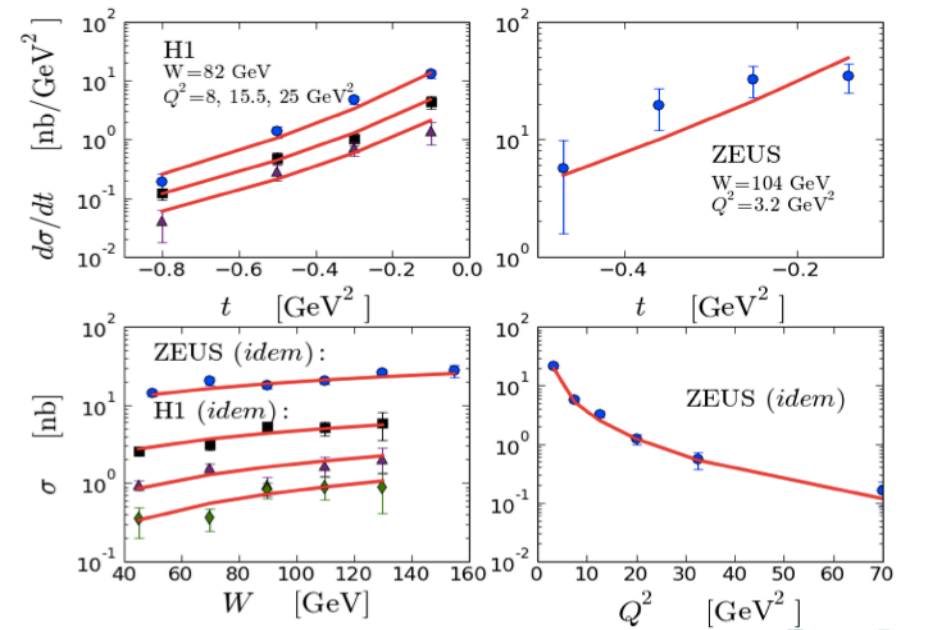
$$\times \int_0^1 dz T^{ij}(x, \xi, z, Q^2) F^i(x, \xi, t) \Phi^j(z)$$



## □ Evolution:

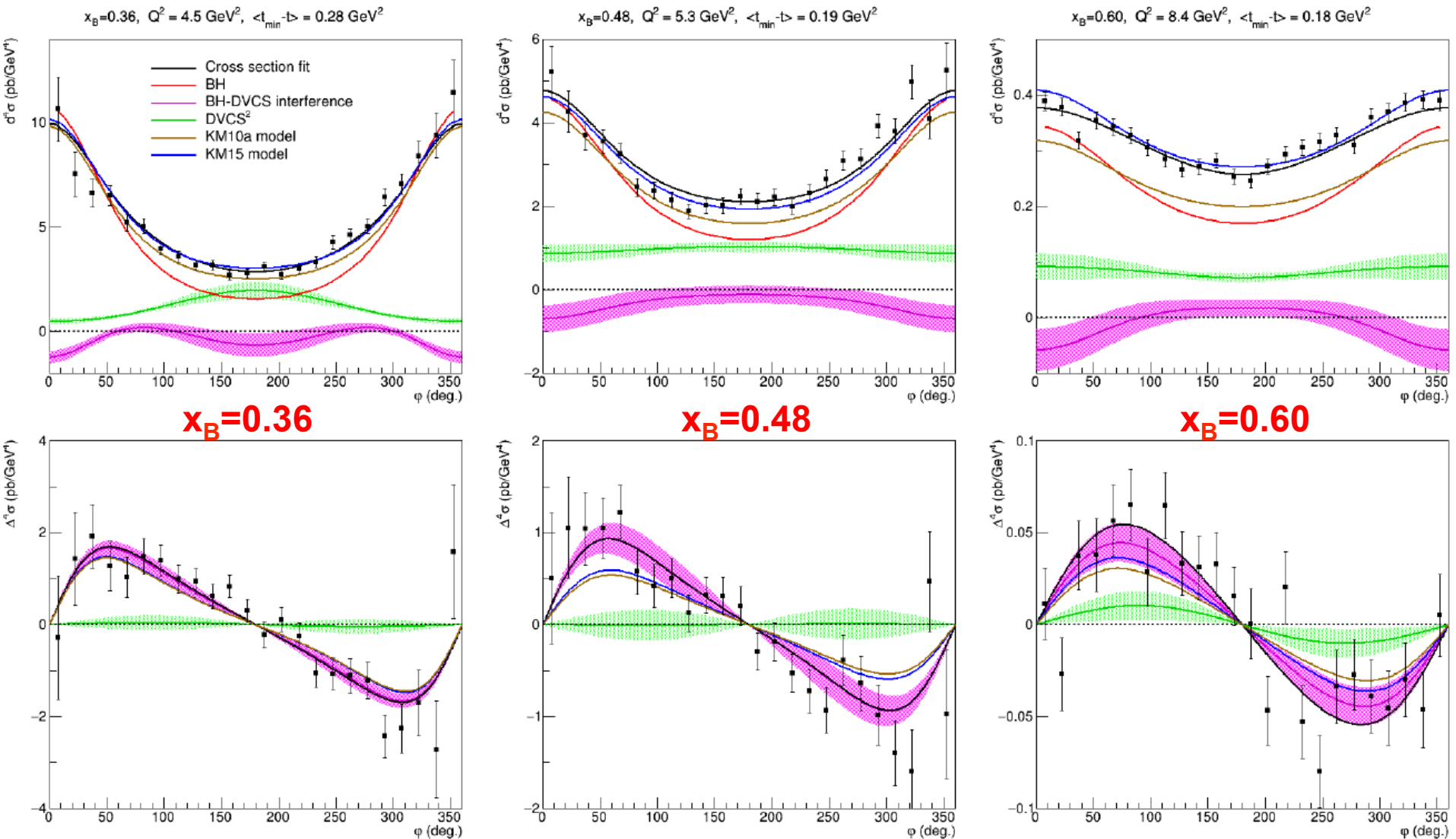
Factorization naturally lead to evolution equations for GPDs

# Data: just the beginning



# JLab E12-06-114 DVCS/Hall A Experiment at 11 GeV

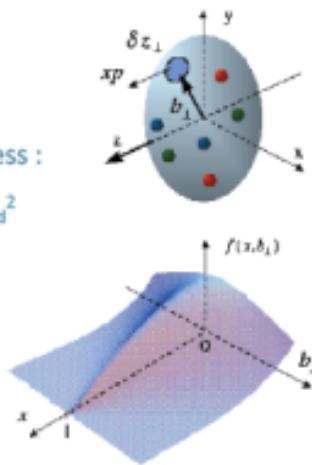
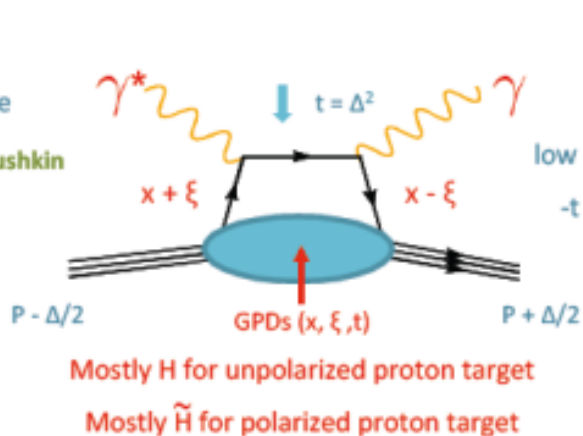
## Sample of cross-section results:



# JLab DVCS/Hall B Experiment at 11 GeV

## CLAS12:

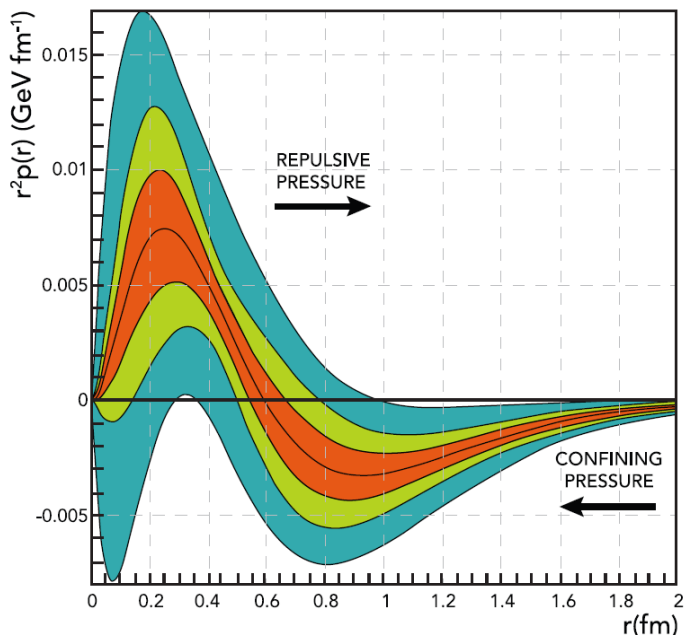
$Q_{\text{hard}}^2$  large  
Mueller, Ji, Radyushkin  
(94-96)



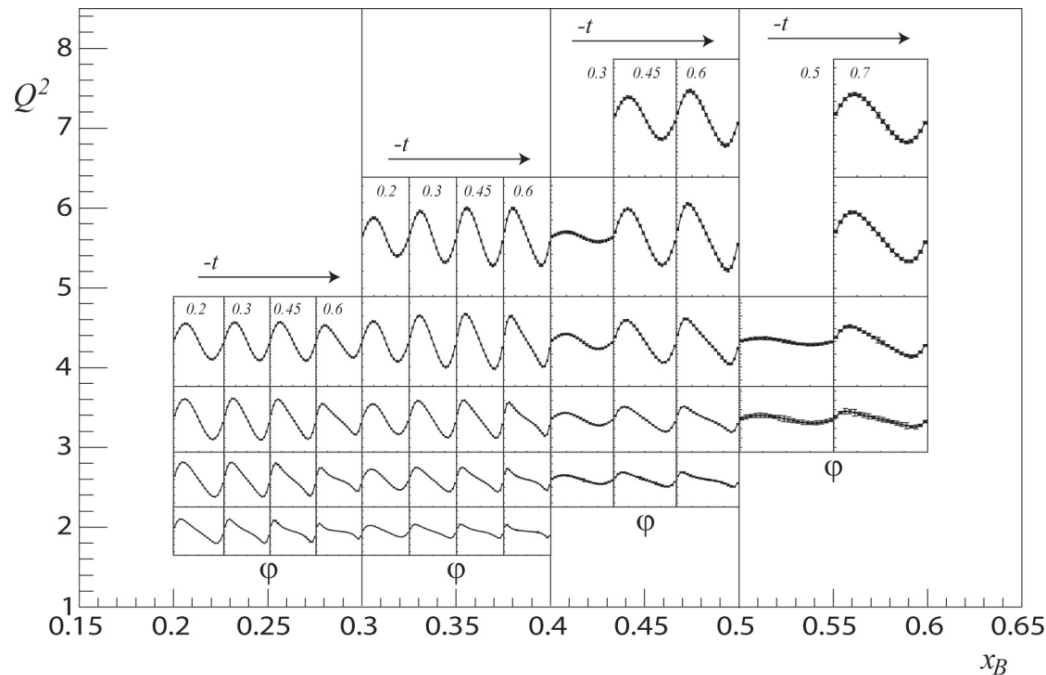
Hall A DVCS  
scaling check  
completed

Hall B DVCS on  
H 50% complete

CLAS12 (projected)

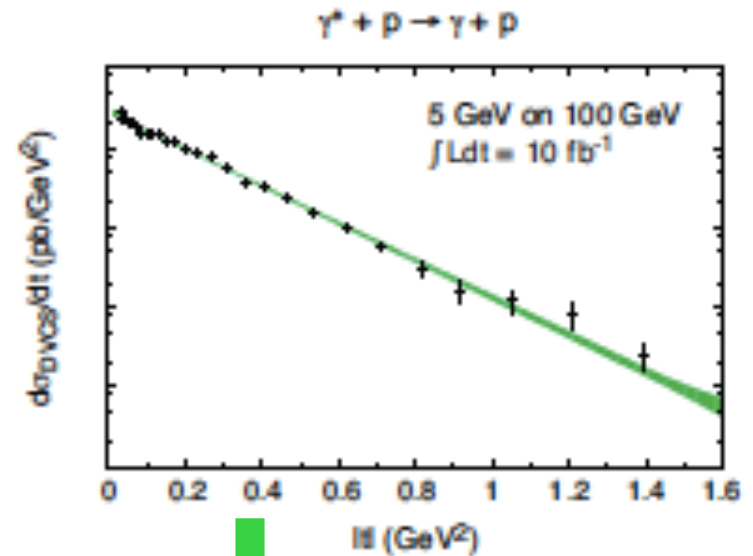
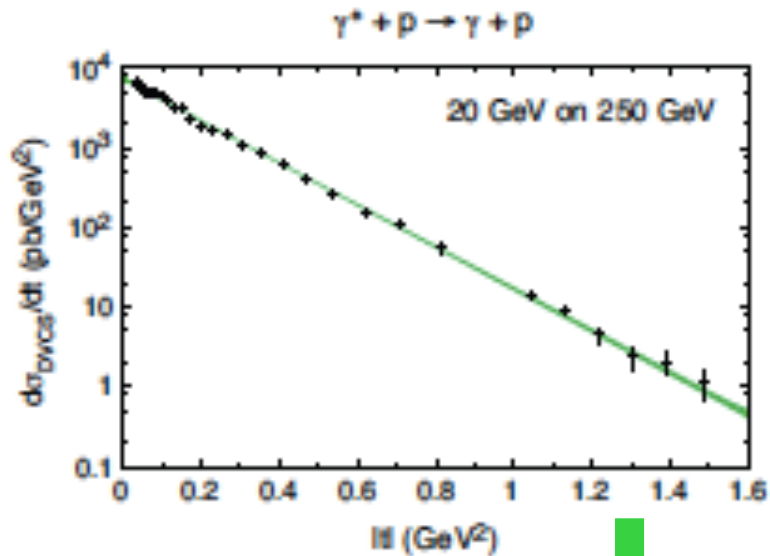


Nature **557**, 396-399 (2018)

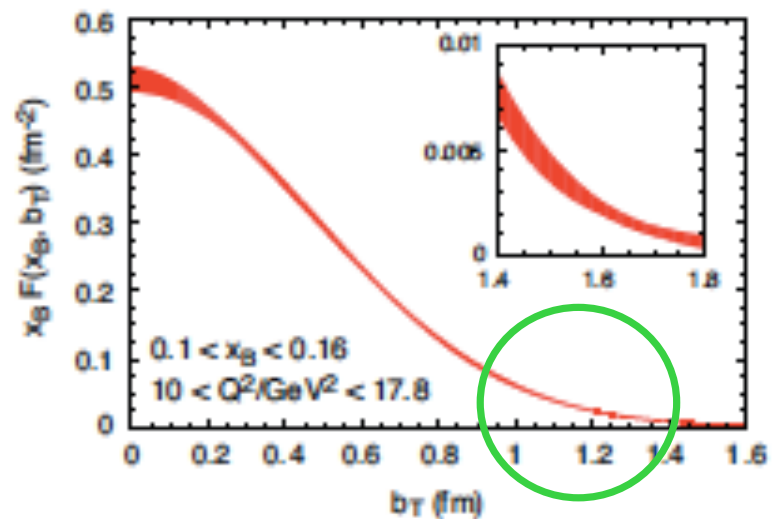
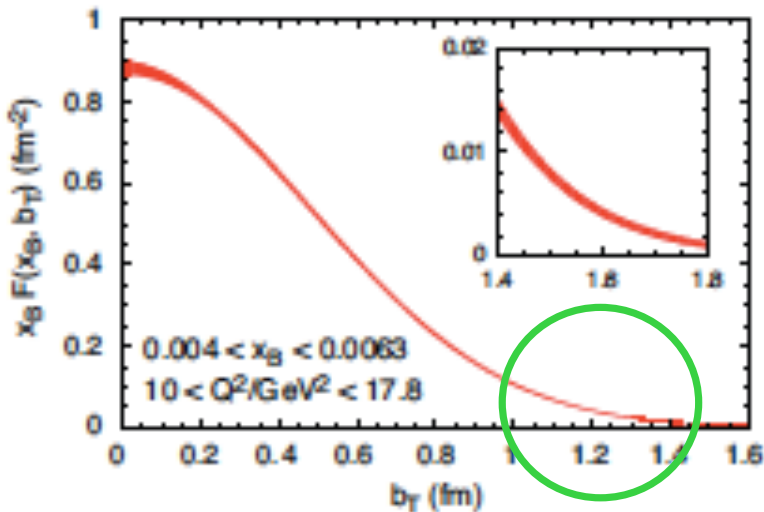


# DVCS at future EIC (White Paper)

## □ Cross Sections:

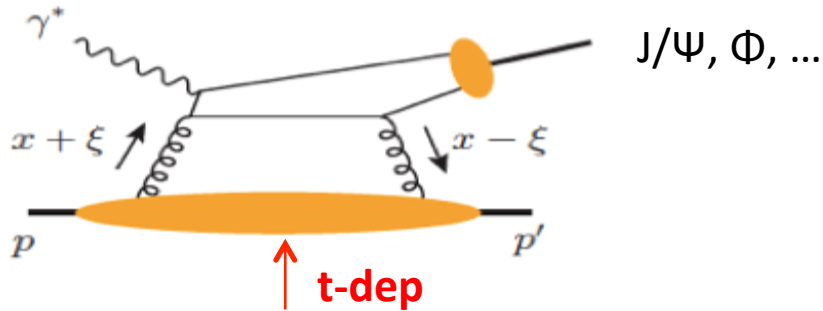


## □ Spatial distributions:



# Imaging the gluon (White Paper)

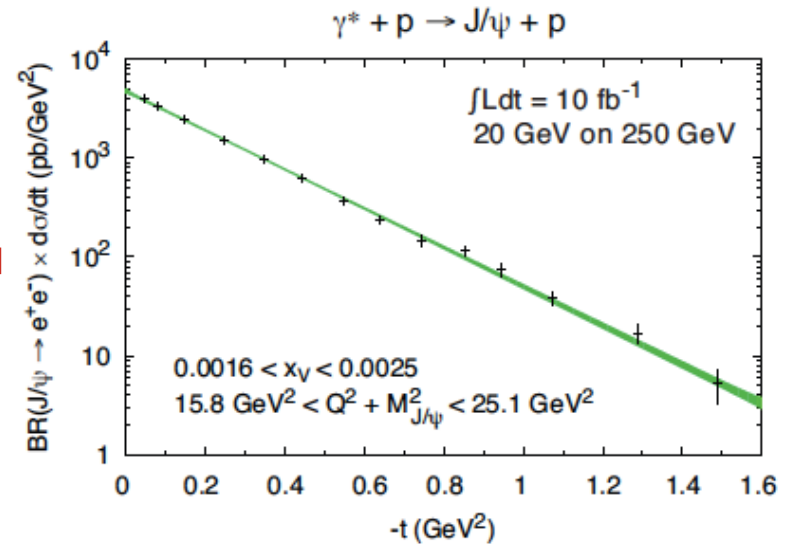
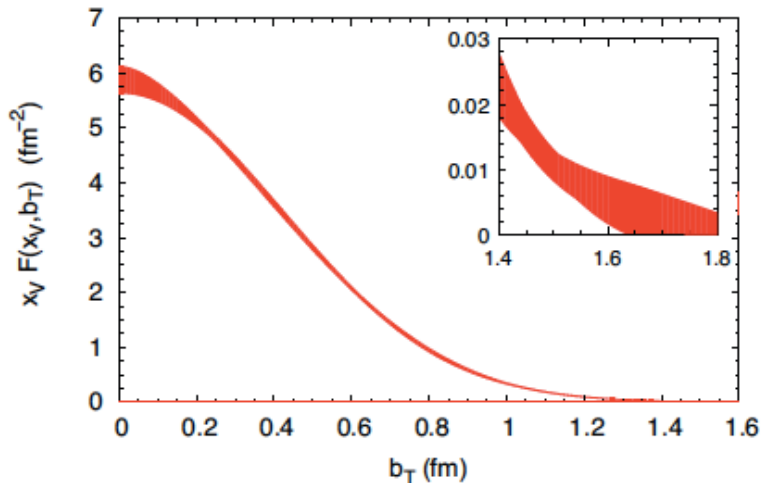
## Exclusive vector meson production:



$$\frac{d\sigma}{dx_B dQ^2 dt}$$

- ✧ Fourier transform of the t-dep
- ➡ Spatial imaging of glue density
- ✧ Resolution  $\sim 1/Q$  or  $1/M_Q$

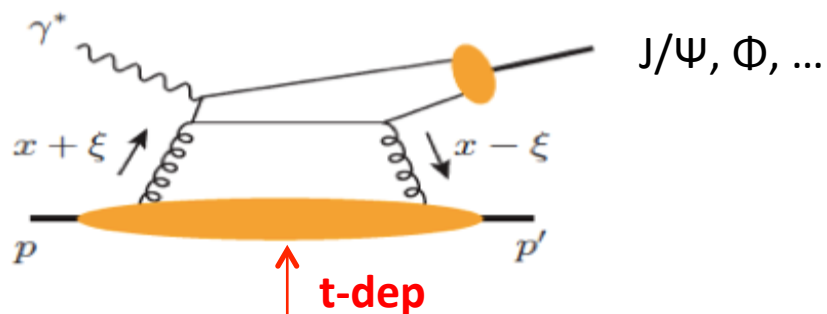
## Gluon imaging from simulation:





# Imaging the gluon (White Paper)

## Exclusive vector meson production:



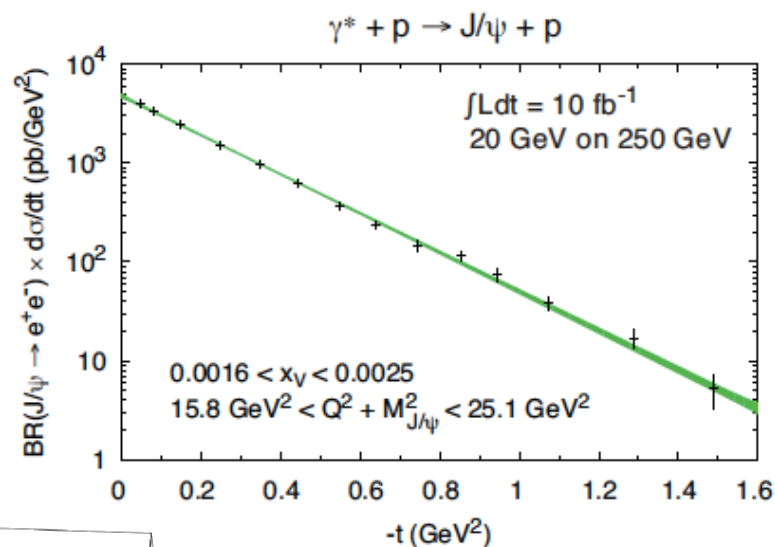
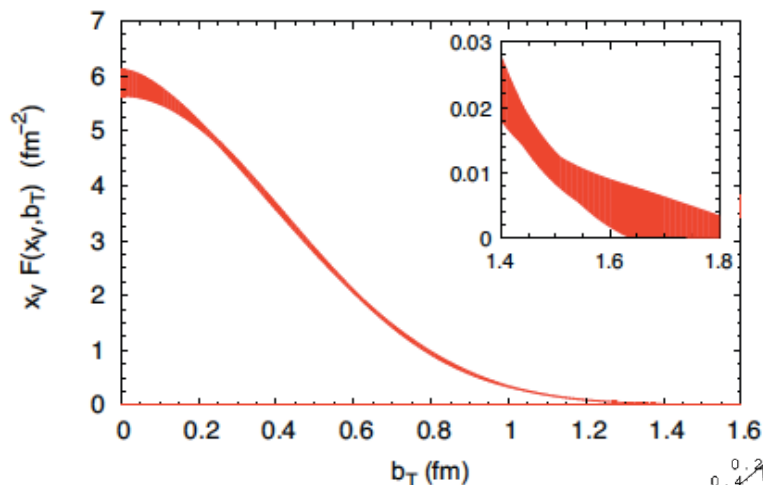
$$\frac{d\sigma}{dx_B dQ^2 dt}$$

Fourier transform of the t-dep

Spatial imaging of glue density

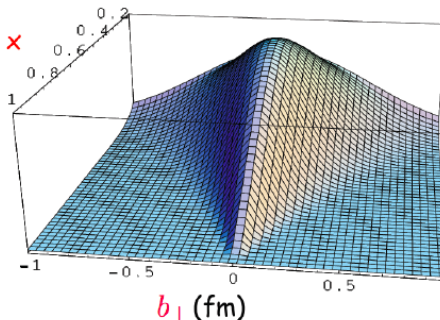
Resolution  $\sim 1/Q$  or  $1/M_Q$

## Gluon imaging from simulation:



Only possible at the EIC  
Gluon radius?

Gluon radius (x)!



How spread  
at small-x?  
Color confinement

Y.-S. Liu et al.  
arXiv:1902.00307

## □ Definition:

$$\tilde{F}(\Gamma, x, \tilde{\xi}, t, P^z, \tilde{\mu}) = \frac{1}{N} \int \frac{dz}{4\pi} e^{ixzP^z} \langle P'', S'' | \bar{\psi} \left( \frac{z}{2} \right) \Gamma \lambda^a W_z \left( \frac{z}{2}, -\frac{z}{2} \right) \psi \left( -\frac{z}{2} \right) | P', S' \rangle$$

GPDs could be extracted from  
LQCD calculation of j-j correlation

## □ Decomposition:

$$\tilde{F}(\Gamma, x, \xi, t, P^z, \tilde{\mu}) = \frac{1}{2P^t} \bar{u}(P'', S'') \left\{ \tilde{H}(\Gamma, x, \xi, t, P^z, \tilde{\mu}) \Gamma + \tilde{E}(\Gamma, x, \xi, t, P^z, \tilde{\mu}) \frac{[\Delta, \Gamma]}{4M} + \tilde{H}'(\bar{\Gamma}, x, \xi, t, \mu) \frac{P^{[z} \Delta^{\perp]}}{M^2} + \tilde{E}'(\bar{\Gamma}, x, \xi, t, \mu) \frac{\gamma^{[z} P^{\perp]}}{M} \right\} u(P', S')$$

## □ Factorization:

$$\tilde{F}(\Gamma, x, \xi, t, P^z, \mu_R, p_R^z) = \int_{-1}^1 \frac{dy}{|y|} C_\Gamma \left( \frac{x}{y}, \frac{\xi}{y}, r, \frac{yP^z}{\mu}, \frac{yP^z}{p_R^z} \right) F(\bar{\Gamma}, y, \xi, t, \mu) + \mathcal{O} \left( \frac{M^2}{P_z^2}, \frac{t}{P_z^2}, \frac{\Lambda_{\text{QCD}}^2}{x^2 P_z^2} \right)$$

## □ Matching coefficient – NLO – RI/MOM renormalization:

$$C_\Gamma \left( x, \xi, r, \frac{p^z}{\mu}, \frac{p^z}{p_R^z} \right) = \delta(1-x) + \left[ f_1 \left( \Gamma, x, \xi, \frac{p^z}{\mu} \right) - \left| \frac{p^z}{p_R^z} \right| f_2 \left( \Gamma, \frac{p^z}{p_R^z} (x-1) + 1, r \right) \right]_+ + \delta_{\Gamma, i\sigma^z \pm} \delta(1-x) \frac{\alpha_s C_F}{4\pi} \ln \left( \frac{\mu^2}{\mu_R^2} \right) + \mathcal{O}(\alpha_s^2)$$

where the functions  $f_1$ ,  $f_2$  and  $\delta_\Gamma$  are given in the paper (1902.00307)

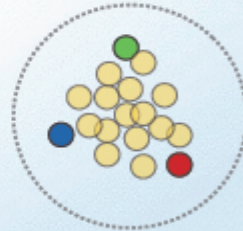
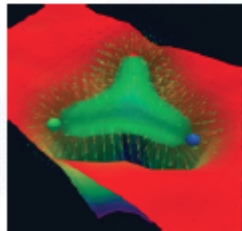
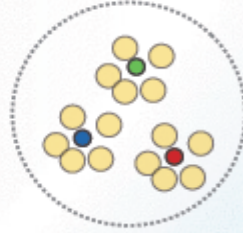
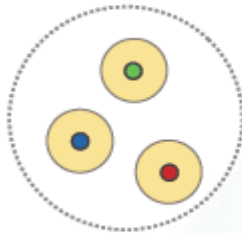
# Beyond the 3D picture

## □ Spatial distributions of quarks and gluons:

Static



Boosted



**Bag Model:**

Gluon field distribution is wider than the fast moving quarks.

**Gluon radius > Charge Radius**

**Constituent Quark Model:**

Gluons and sea quarks hide inside massive quarks.

**Gluon radius ~ Charge Radius**

**Lattice Gauge theory (with slow moving quarks):**

Gluons more concentrated inside the quarks

**Gluon radius < Charge Radius**

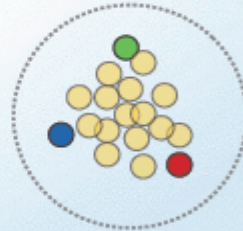
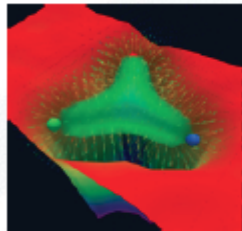
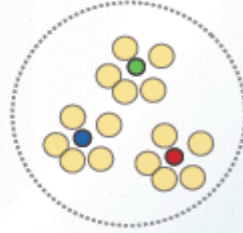
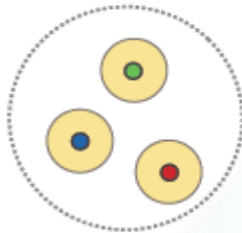
# Beyond the 3D picture

## □ Spatial distributions of quarks and gluons:

Static



Boosted



**Bag Model:**

Gluon field distribution is wider than the fast moving quarks.

**Gluon radius > Charge Radius**

**Constituent Quark Model:**

Gluons and sea quarks hide inside massive quarks.

**Gluon radius ~ Charge Radius**

**Lattice Gauge theory (with slow moving quarks):**

Gluons more concentrated inside the quarks

**Gluon radius < Charge Radius**

***3D confined motion (TMDs) + spatial distribution (GPDs)***

***Hints on the color confining mechanism***

**Relation between charge radius, quark radius (x), and gluon radius (x)?** Jefferson Lab

# Summary

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## ❑ GPDs are fundamental quantum probability distributions

Carry important information on spatial imaging of hadron's partonic structure

## ❑ Need exclusive processes with a unbroken hadron under hard collisions

✧ Need lepton-hadron facilities

✧ Need well-controlled exclusive processes in lepton-hadron collisions

✧ DVCS, DDVCS, DVMP, Diffractive heavy vector boson production, ...

✧ JLab12, COMPASS, and future EIC will produce a lot of data on GPDs

## ❑ GPDs could be extracted from LQCD calculations

Work just got started – more efforts are needed!

**Thank you!**