

New method to extract the top-quark width from Non-Resonant Production

Tomáš Ježo

University of Zürich

In collaboration with C. Herwig[†] and B. Nachman

based on [[arXiv:1903.10519](#)]

[†] my slides based on those of C. Herwig



**Universität
Zürich^{UZH}**

**LHC TOP WG meeting
May 2019**

The top-quark width

- Though the top mass has been measured to sub-% precision, current direct width measurement sensitivity is $\sim 50\%$
- The SM predicts a precise relation between the top-quark mass and width, and has been calculated at NNLO QCD

$$\Gamma_t^{\text{LO}} = \frac{G_F m_t^3}{8\sqrt{2}\pi} \left(1 - \frac{m_W^2}{m_t^2}\right)^2 \left(1 + 2\frac{m_W^2}{m_t^2}\right)$$

- ▶ Precise measurements of the top width allow stringent tests of the SM
- Modifications from BSM physics remain possible:
 - ▶ Undetected decays: $t \rightarrow H^+ b$, FCNCs, SUSY, ...
 - ▶ Radiative corrections (g_{tWb}): SUSY, 2HDMs, $W - W'$ mixing

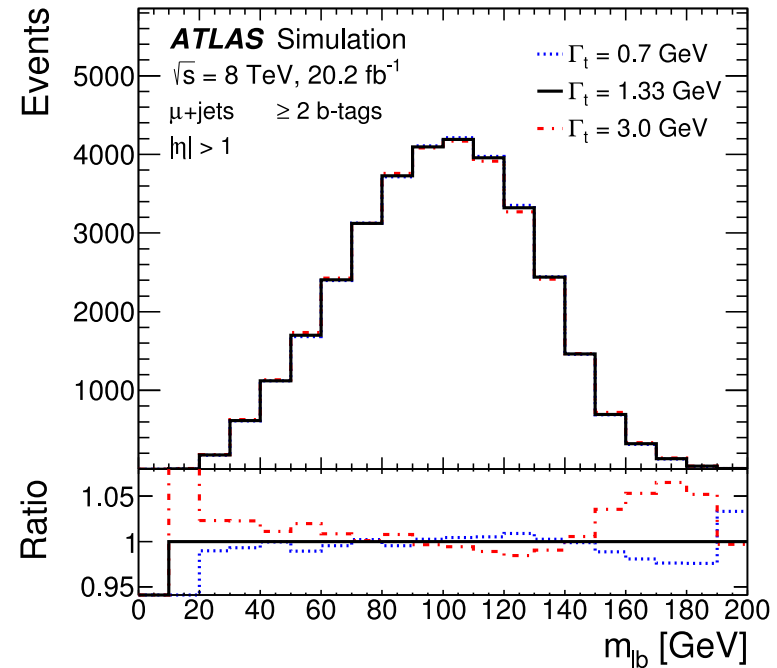
Previous measurements (I)

- Since Tevatron, there are two methods used to measure the top width
- Template fits to reconstructed mass spectra (lepton+jets) (**direct measurement**)
 - ▶ $\Gamma_t = 1.76 \pm 0.33(\text{stat.})_{-0.68}^{+0.79}(\text{syst.}) \text{ GeV}$
- Combination of t -channel single-top XS and $(t \rightarrow Wb)/(t \rightarrow Wq)$ ratio (indirect measurement)

$$\Gamma_t = \frac{\sigma_{t\text{-ch.}}}{\mathcal{B}(t \rightarrow Wb)} \cdot \frac{\Gamma(t \rightarrow Wb)^{\text{theor.}}}{\sigma_{t\text{-ch.}}^{\text{theor.}}}$$

- ▶ $\Gamma_t = 1.36 \pm 0.02(\text{stat.})_{-0.11}^{+0.14}(\text{syst.}) \text{ GeV}$
- ▶ But assumes $\mathcal{B}(t \rightarrow Wq) = 1$

[arXiv: 1709.04207]

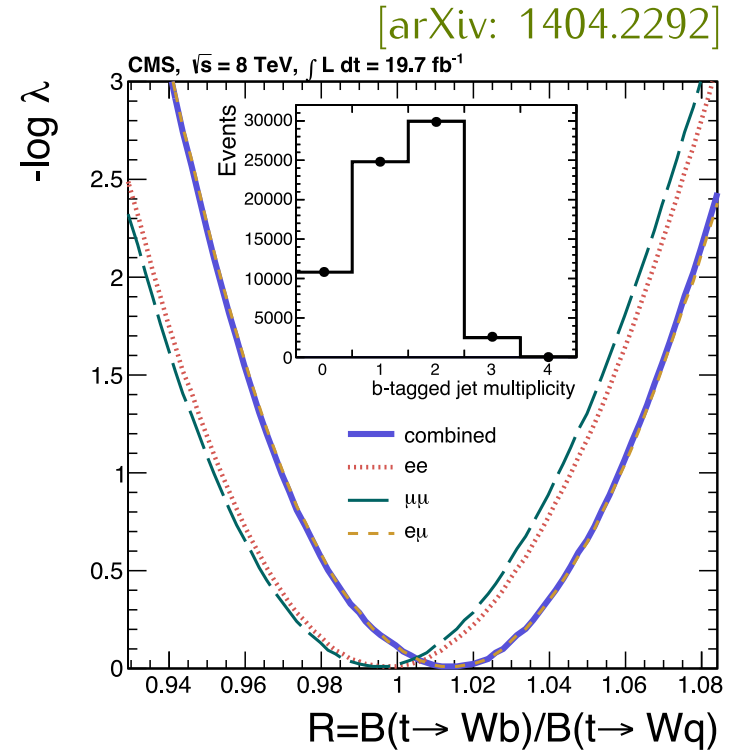


Previous measurements (II)

- Since Tevatron, there are two methods used to measure the top width
- Template fits to reconstructed mass spectra (lepton+jets) (direct measurement)
 - ▶ $\Gamma_t = 1.76 \pm 0.33(\text{stat.})_{-0.68}^{+0.79}(\text{syst.}) \text{ GeV}$
- Combination of t -channel single-top XS and $(t \rightarrow Wb)/(t \rightarrow Wq)$ ratio (indirect measurement)

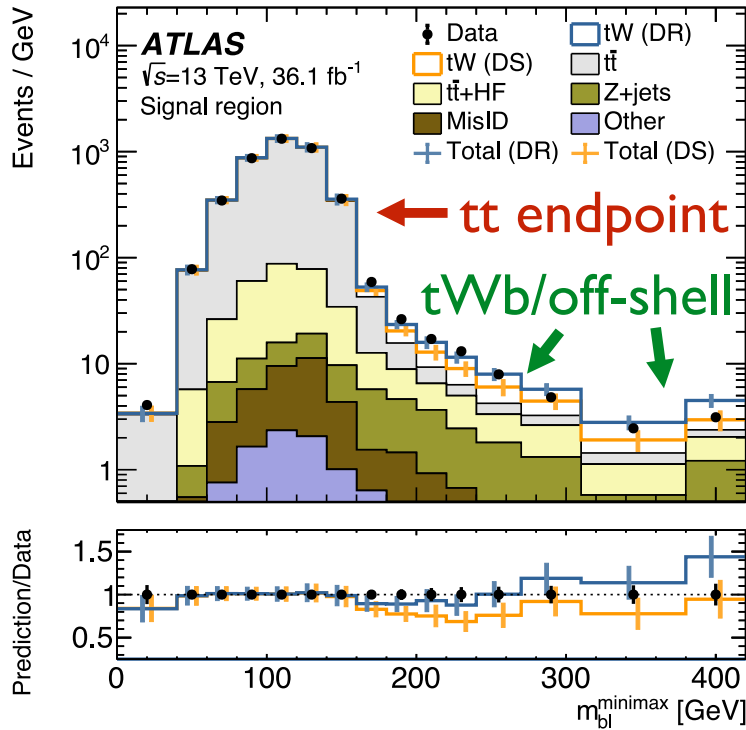
$$\Gamma_t = \frac{\sigma_{t\text{-ch.}}}{\mathcal{B}(t \rightarrow Wb)} \cdot \frac{\Gamma(t \rightarrow Wb)^{\text{theor.}}}{\sigma_{t\text{-ch.}}^{\text{theor.}}}$$

- ▶ $\Gamma_t = 1.36 \pm 0.02(\text{stat.})_{-0.11}^{+0.14}(\text{syst.}) \text{ GeV}$
- ▶ But assumes $\mathcal{B}(t \rightarrow Wq) = 1$

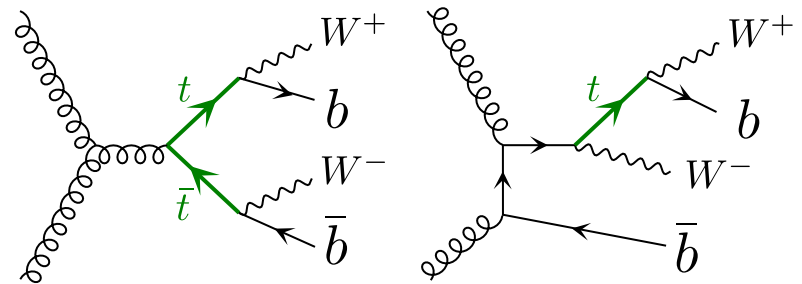


Our proposal

- Use the measurement of m_{bl}^{minimax} in the dilepton channel [TOPQ-2017-05] to extract Γ_t using off-shell events (direct measurement)



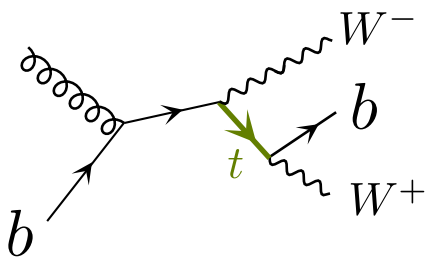
- $m_{bl}^{\text{minimax}} = \min\{\max(m_{b_1, \ell_1}, m_{b_2, \ell_2}), \max(m_{b_1, \ell_2}, m_{b_2, \ell_1})\}$
- Kinematic endpoint at $\sqrt{m_t^2 - m_W^2}$ at LO
- Large m_{bl}^{minimax} dominated by off-shell $t\bar{t}$ and tW production



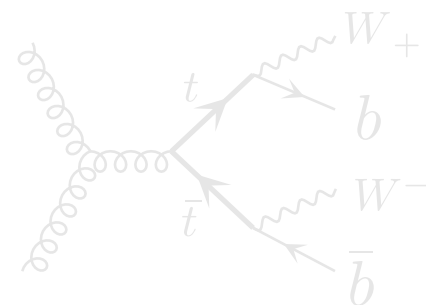
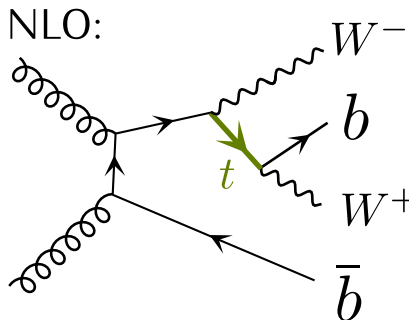
$t\bar{t}$ and tW interplay

- tW 5FNS (b in proton, $m_b = 0$)

LO:



NLO:



- Requires a procedure to remove $t\bar{t}$ contribution

$$\mathcal{M} = \mathcal{M}^{tW} + \mathcal{M}^{t\bar{t}}$$

Diagram Removal: $\mathcal{R}^{\text{DR}} = \frac{|\mathcal{M}^{tW}|^2}{2s}$

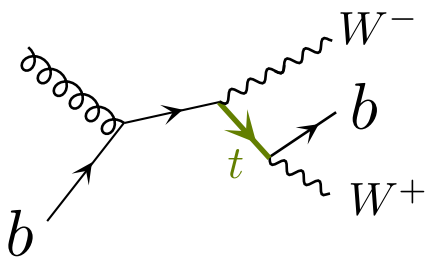
Diagram Subtraction: $\mathcal{R}^{\text{DS}} = \frac{|\mathcal{M}^{tW} + \mathcal{M}^{t\bar{t}}|^2 - \mathcal{C}}{2s}$

Or others, see [arXiv: 1607.05862]

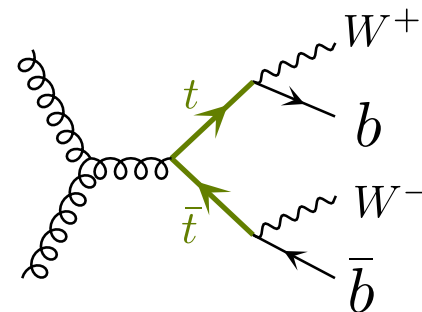
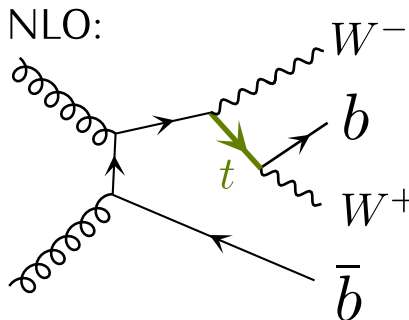
$t\bar{t}$ and tW interplay

- tW 5FNS (b in proton, $m_b = 0$)

LO:



NLO:



► Requires a procedure to remove $t\bar{t}$ contribution

$$\mathcal{M} = \mathcal{M}^{tW} + \mathcal{M}^{t\bar{t}}$$

Diagram Removal: $\mathcal{R}^{\text{DR}} = \frac{|\mathcal{M}^{tW}|^2}{2s}$

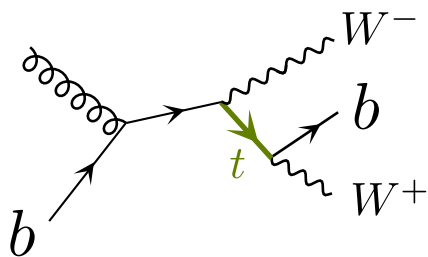
Diagram Subtraction: $\mathcal{R}^{\text{DS}} = \frac{|\mathcal{M}^{tW} + \mathcal{M}^{t\bar{t}}|^2 - \mathcal{C}}{2s}$

Or others, see [arXiv: 1607.05862]

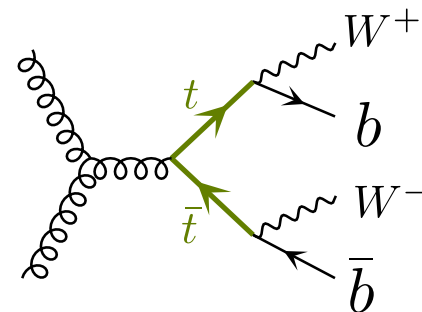
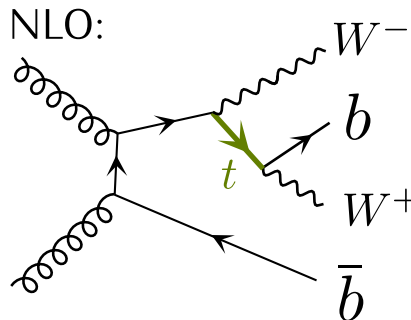
$t\bar{t}$ and tW interplay

- tW 5FNS (b in proton, $m_b = 0$)

LO:



NLO:



- Requires a procedure to remove $t\bar{t}$ contribution

$$\mathcal{M} = \mathcal{M}^{tW} + \mathcal{M}^{t\bar{t}}$$

Diagram Removal: $\mathcal{R}^{\text{DR}} = \frac{|\mathcal{M}^{tW}|^2}{2s}$

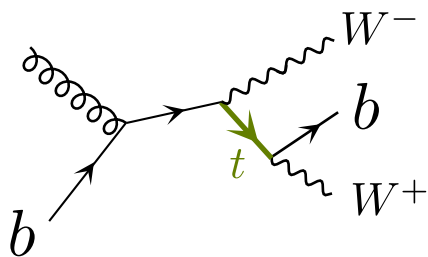
Diagram Subtraction: $\mathcal{R}^{\text{DS}} = \frac{|\mathcal{M}^{tW} + \mathcal{M}^{t\bar{t}}|^2 - \mathcal{C}}{2s}$

Or others, see [\[arXiv: 1607.05862\]](https://arxiv.org/abs/1607.05862)

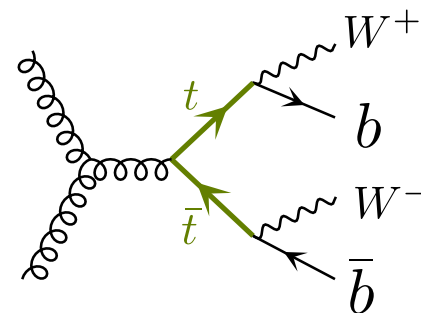
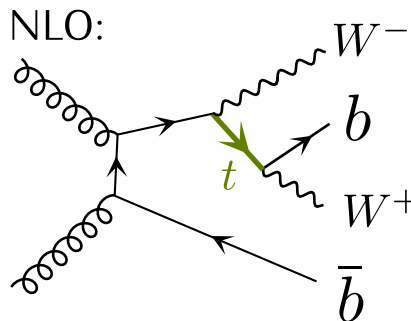
$t\bar{t}$ and tW interplay

- tW 5FNS (b in proton, $m_b = 0$)

LO:



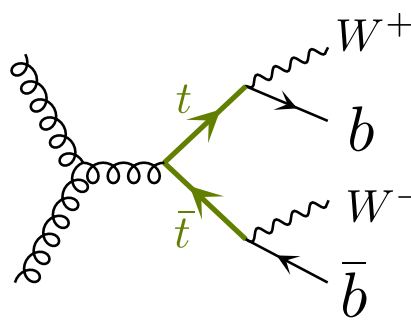
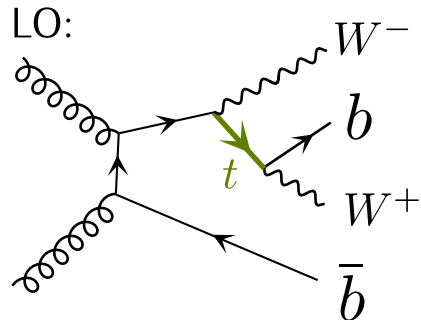
NLO:



- Requires a procedure to remove $t\bar{t}$ contribution

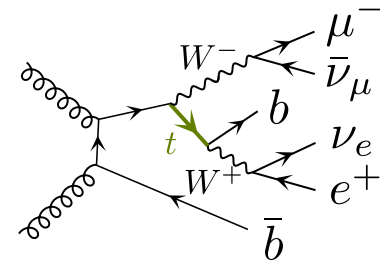
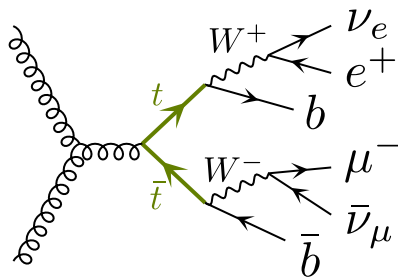
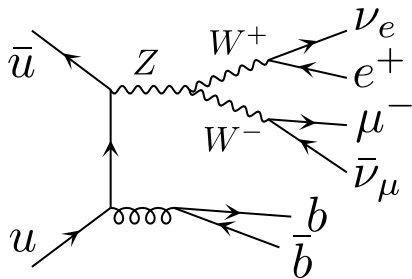
- tW 4FNS (no b in proton, $m_b > 0$)

LO:



- Unified treatment of $t\bar{t}$ and tW

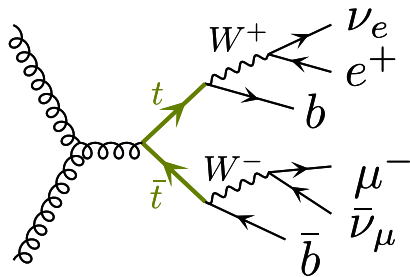
$t\bar{t}$ & tW in POWHEG BOX RES: b_bbar_4l



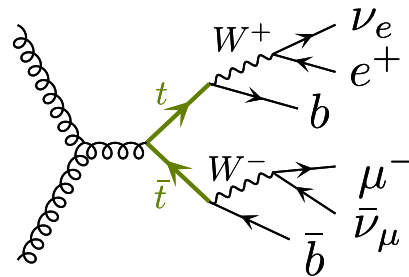
- $pp \rightarrow \ell^+ \nu_\ell l^- \bar{\nu}_l b \bar{b}$ production at NLO
- Resonance-aware matching to parton showers [TJ, Nason, 2015]
- Exact spin correlations* and exact off-shell effects
- W hadronic decays work in progress
- Generator: b_bbar_4l [TJ, Lindert, Nason, Oleari, Pozzorini, 2016]

$t\bar{t}$ & tW in POWHEG BOX RES: b_bbar_4l

Traditional NLOPS



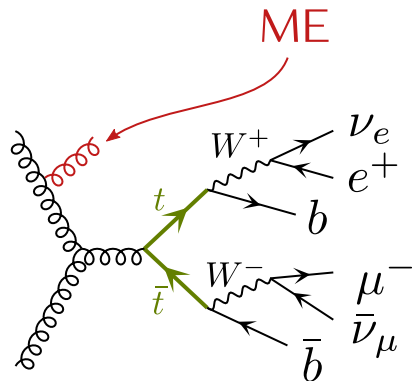
Multiple-radiation-improved
NLOPS (allrad)



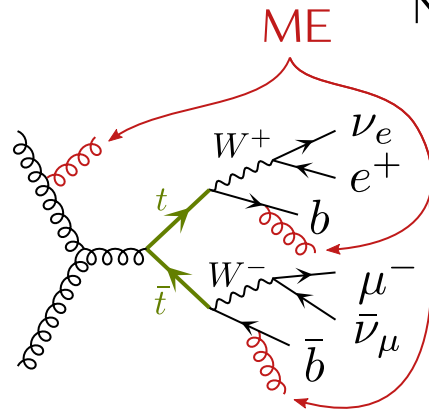
- $pp \rightarrow \ell^+ \nu_\ell l^- \bar{\nu}_l b \bar{b}$ production at NLO
- Resonance-aware matching to parton showers [TJ, Nason, 2015]
- Exact spin correlations* and exact off-shell effects
- W hadronic decays work in progress
- Generator: b_bbar_4l [TJ, Lindert, Nason, Oleari, Pozzorini, 2016]
- Multiple emissions described using the matrix element via the allrad feature

$t\bar{t}$ & tW in POWHEG BOX RES: b_bbar_4l

Traditional NLOPS



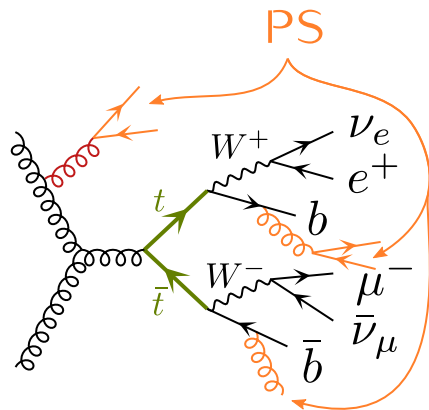
Multiple-radiation-improved
NLOPS (allrad)



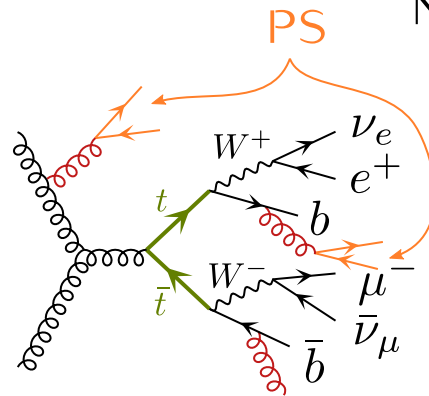
- $pp \rightarrow \ell^+ \nu_\ell l^- \bar{\nu}_l b \bar{b}$ production at NLO
- Resonance-aware matching to parton showers [TJ, Nason, 2015]
- Exact spin correlations* and exact off-shell effects
- W hadronic decays work in progress
- Generator: b_bbar_4l [TJ, Lindert, Nason, Oleari, Pozzorini, 2016]
- Multiple emissions described using the matrix element via the allrad feature

$t\bar{t}$ & tW in POWHEG BOX RES: b_bbar_4l

Traditional NLOPS



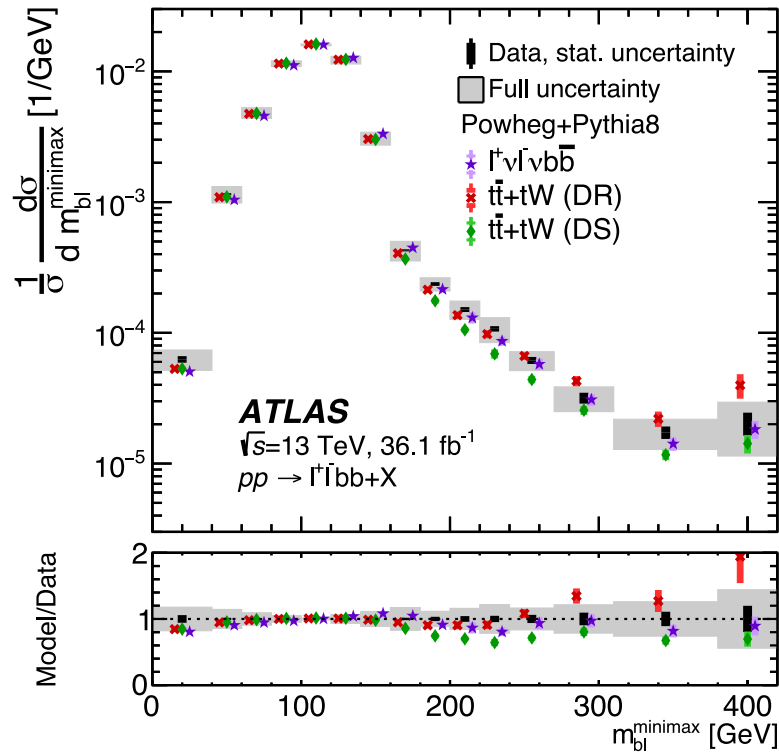
Multiple-radiation-improved
NLOPS (allrad)



- $pp \rightarrow \ell^+ \nu_\ell l^- \bar{\nu}_l b \bar{b}$ production at NLO
- Resonance-aware matching to parton showers [TJ, Nason, 2015]
- Exact spin correlations* and exact off-shell effects
- W hadronic decays work in progress
- Generator: b_bbar_4l [TJ, Lindert, Nason, Oleari, Pozzorini, 2016]
- Multiple emissions described using the matrix element via the allrad feature

$t\bar{t}$ and tW interplay

- Probing the quantum interference between singly and doubly resonant top-quark production in pp collisions at $\sqrt{s} = 13$ TeV with the ATLAS detector [TOPQ-2017-05]:



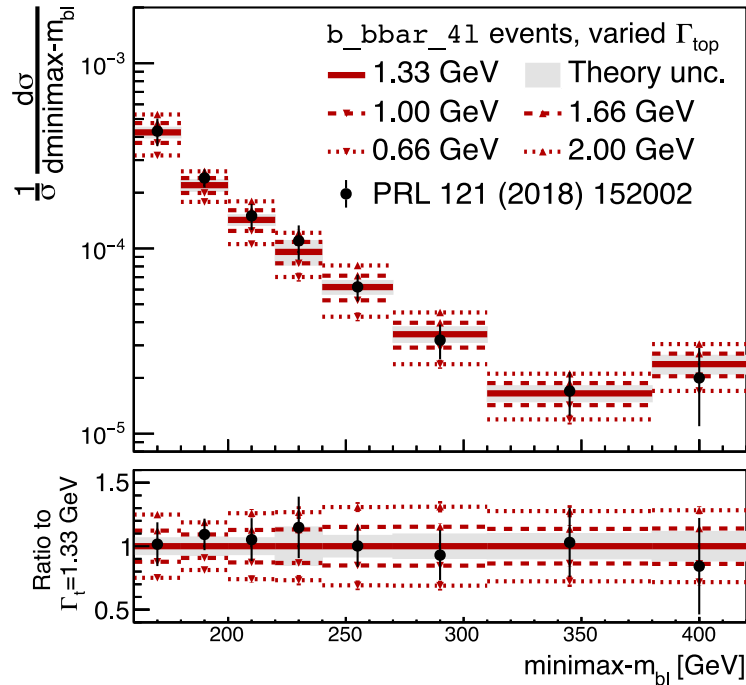
- $b_b\bar{b}_41$ prediction describes the data in the tail much better than predictions obtained using DR and DS prescriptions

Simulated samples

- Study width-dependence using the `b_bbar_4l` process implemented in POWHEG BOX RES
- This includes:
 - Consistent NLO+PS treatment of top resonances
 - Corrections to top propagators, and off-shell top-decay chains
 - Full NLO accuracy in production, decay, and their interference
 - Exact spin correlation at NLO
- Nominal sample uses $m_t = 172.5$ GeV and $\Gamma_t = 1.33$ GeV
 - Alternatives with $\Gamma_t = 0.66, 1.0, 1.66, 2.0$ GeV
 - Systematics: scale, PDF, α_S , and top mass variations
- Also generated a LO MG5_aMC@NLO sample for comparison

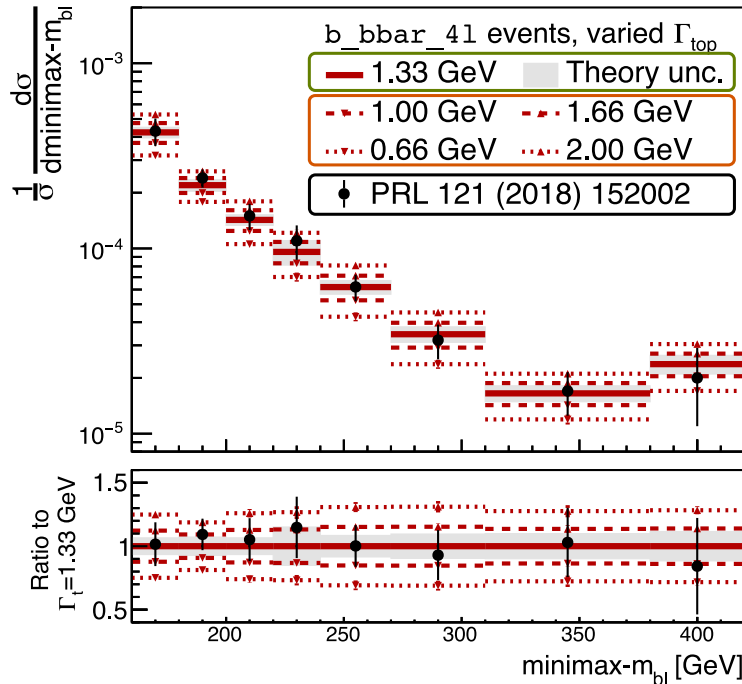
Extraction procedure (I)

- Calculate minimax- $m_{b\ell}$ using the dedicated [TOPQ-2017-05](#) Rivet routine:
- Interpolate yields in each bin as a function of the width



Extraction procedure (I)

- Calculate $\text{minimax-}m_{b\ell}$ using the dedicated **TOPQ-2017-05** Rivet routine:
- Interpolate yields in each bin as a function of the width



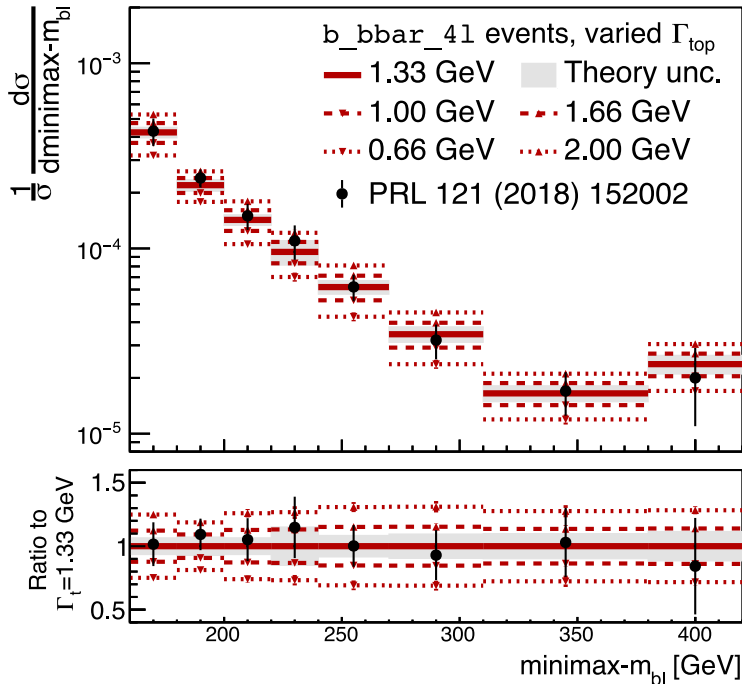
b_bbar_4l – nominal
(scale, PDF, α_S , m_t uncertainties)

b_bbar_4l – varied Γ_t

ATLAS data
TOPQ-2017-05

Extraction procedure (I)

- Calculate minimax- $m_{b\ell}$ using the dedicated **TOPQ-2017-05** Rivet routine:
- Interpolate yields in each bin as a function of the width



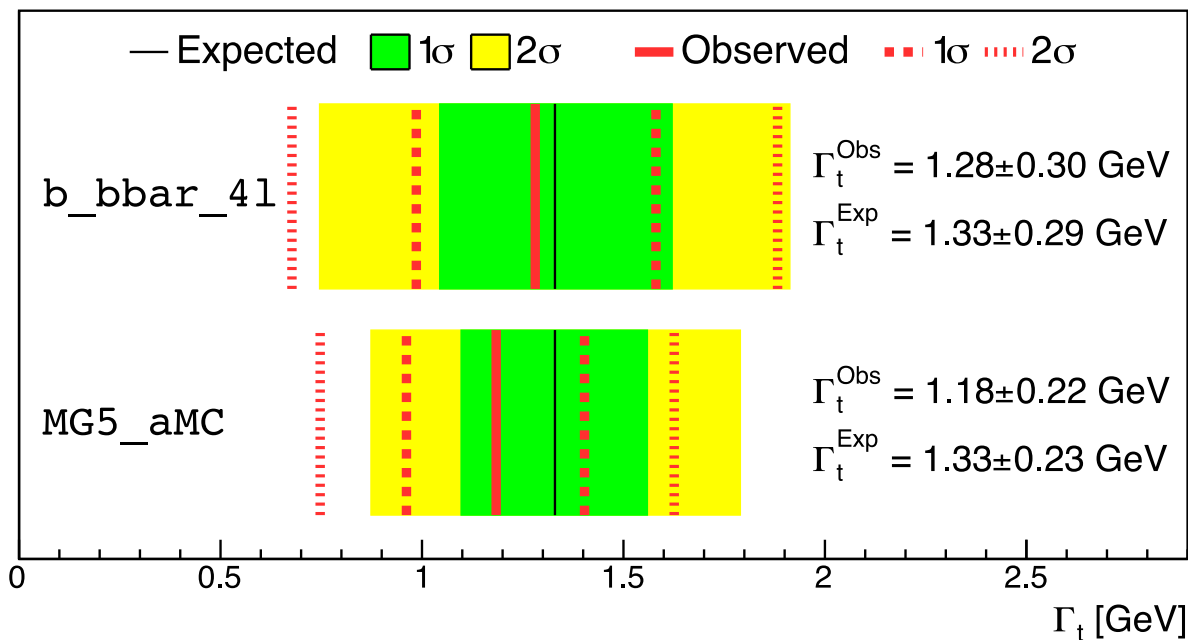
- Predictions in excellent agreement with the data
- Calculate a χ^2 using data + covariance matrix

$$\chi^2 = \sum_{i,j} (d_i - m_i) \cdot V_{ij}^{-1} \cdot (d_j - m_j)$$

- The best-fit width minimizes the χ^2
- Vary data (d_i) and model (m_i) to assess uncertainties

Results

- We find that $\Gamma_t = 1.28 \pm 0.30$ GeV using b_bbar_4l model
- Measurement uncertainty (0.27 GeV) dominates theory (0.14 GeV)
- The LO MG5_aMC@NLO model extracts $\Gamma_t = 1.18 \pm 0.22$ GeV
 - Predicts larger sensitivity to the width than the NLO model!



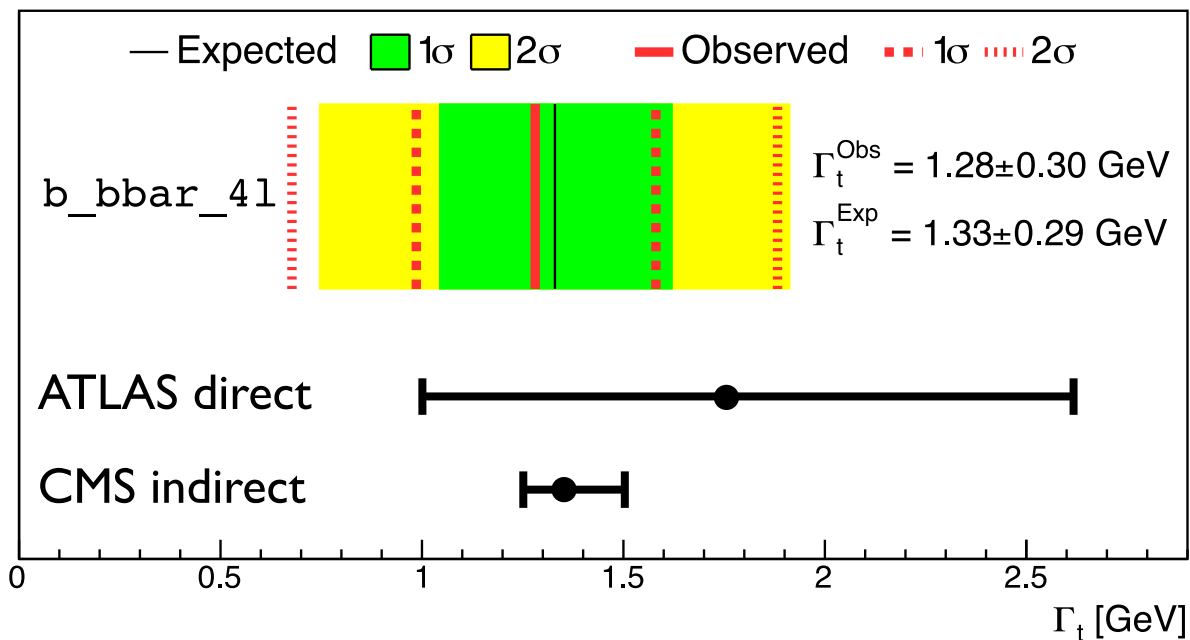
Compare to:

best direct result
 $1.76 + 0.86 / - 0.76$ GeV
(ATLAS: [1709.04207])

best indirect result
 $1.36 + 0.14 / - 0.11$ GeV
(CMS: [1404.2292])

Results

- We find that $\Gamma_t = 1.28 \pm 0.30$ GeV using b_bbar_41 model
- Measurement uncertainty (0.27 GeV) dominates theory (0.14 GeV)
- The LO MG5_aMC@NLO model extracts $\Gamma_t = 1.18 \pm 0.22$ GeV
 - ▶ Predicts larger sensitivity to the width than the NLO model!



Compare to:

best direct result
 $1.76 + 0.86 / - 0.76$ GeV
(ATLAS: [1709.04207])

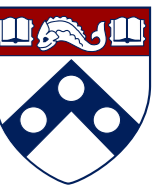
best indirect result
 $1.36 + 0.14 / - 0.11$ GeV
(CMS: [1404.2292])

Summary

- Direct measurements of the top-quark width are currently at the level of accuracy of 50%
- We present a new method that considers only off-shell events
- $W^+W^-b\bar{b}$ analysis strategy of TOPQ-2017-05 provides a clean extraction of the cross section in the off-shell regime
- The `b_bbar_4l` generator is the first even generator capable of describing the $t\bar{t}$ and tWb processes required for this analysis at NLO+PS accuracy
- The combination of the two yields $\Gamma_t = 1.28 \pm 0.30$ GeV
- This strategy should be pursued by LHC experiments to design new, more powerful analyses!

Direct determination of Γ_t with bb41

Backup slides

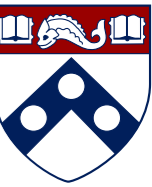


b_bbar_4l generator setup

- Nominal sample produced with $m_t=172.5$ GeV, $\Gamma_H=1.33$ GeV,
- Alternative widths simulated: 0.66, 1.0, 1.66, 2.0 GeV
- NNPDF3.0 NLO $\alpha_s=0.118$ set, with A14 tune
- Scale: geometric mean of (anti)top transverse mass, $h_{\text{damp}}=m_t$
- All different-family lepton flavor combs ("channel 7")
- Three hardest emissions kept ("allrad 1")
- Uncertainties
 - Top mass variations $m_t=171.5, 173.5$
 - α_s variations: 0.115 (0.121) in PDF + Var3C in shower
 - Weights: 7-point scale variations, PDF eigenvectors
 - For statistically independent events, fit dependence per bin:

$$m_i(\alpha_s, m_t) = \hat{m}_i(\alpha_s^{\text{SM}}, m_t^{\text{SM}}) + \hat{a}_i(\alpha_s - \alpha_s^{\text{SM}}) + \hat{b}_i(m_t - m_t^{\text{SM}})$$

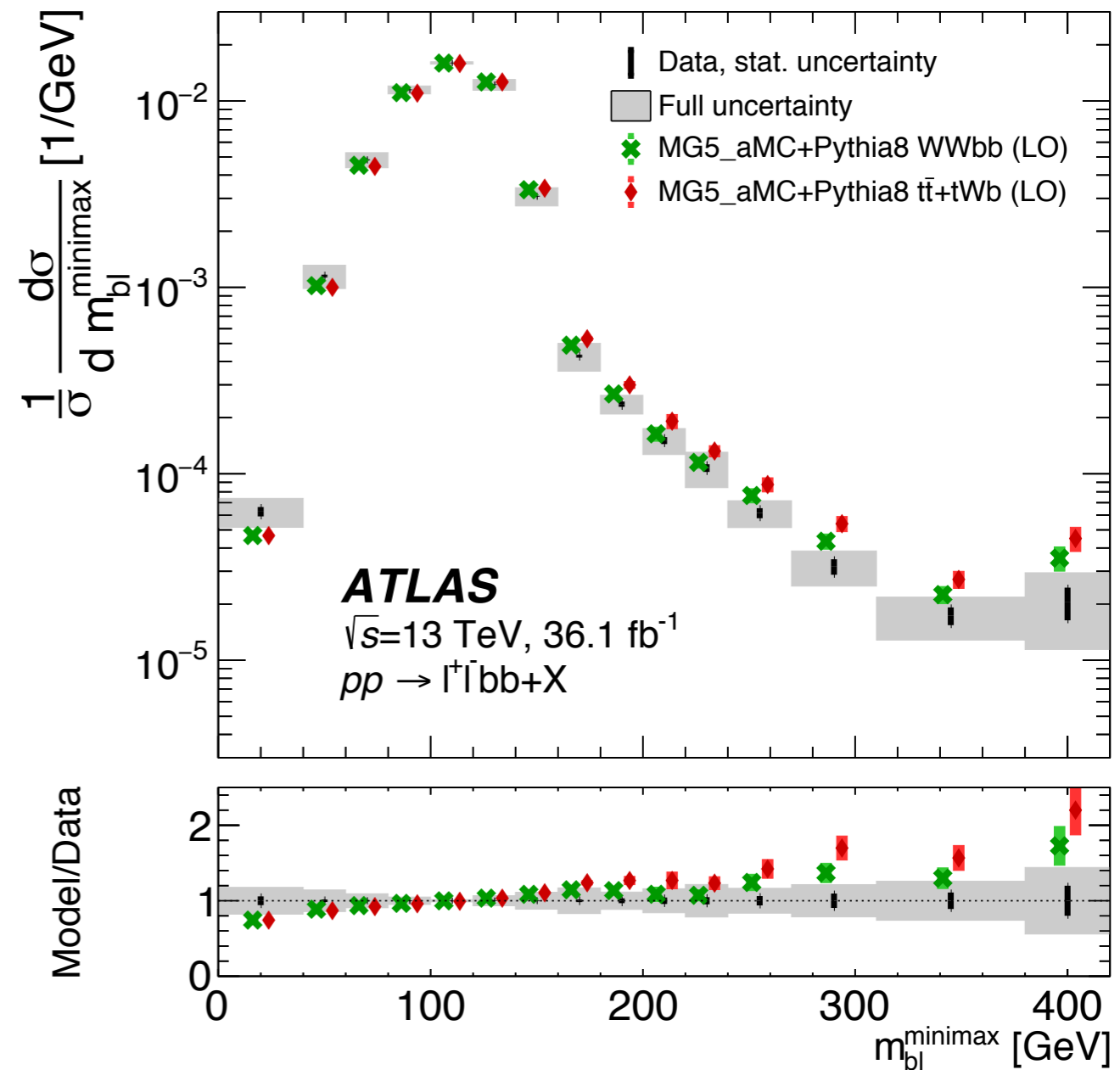
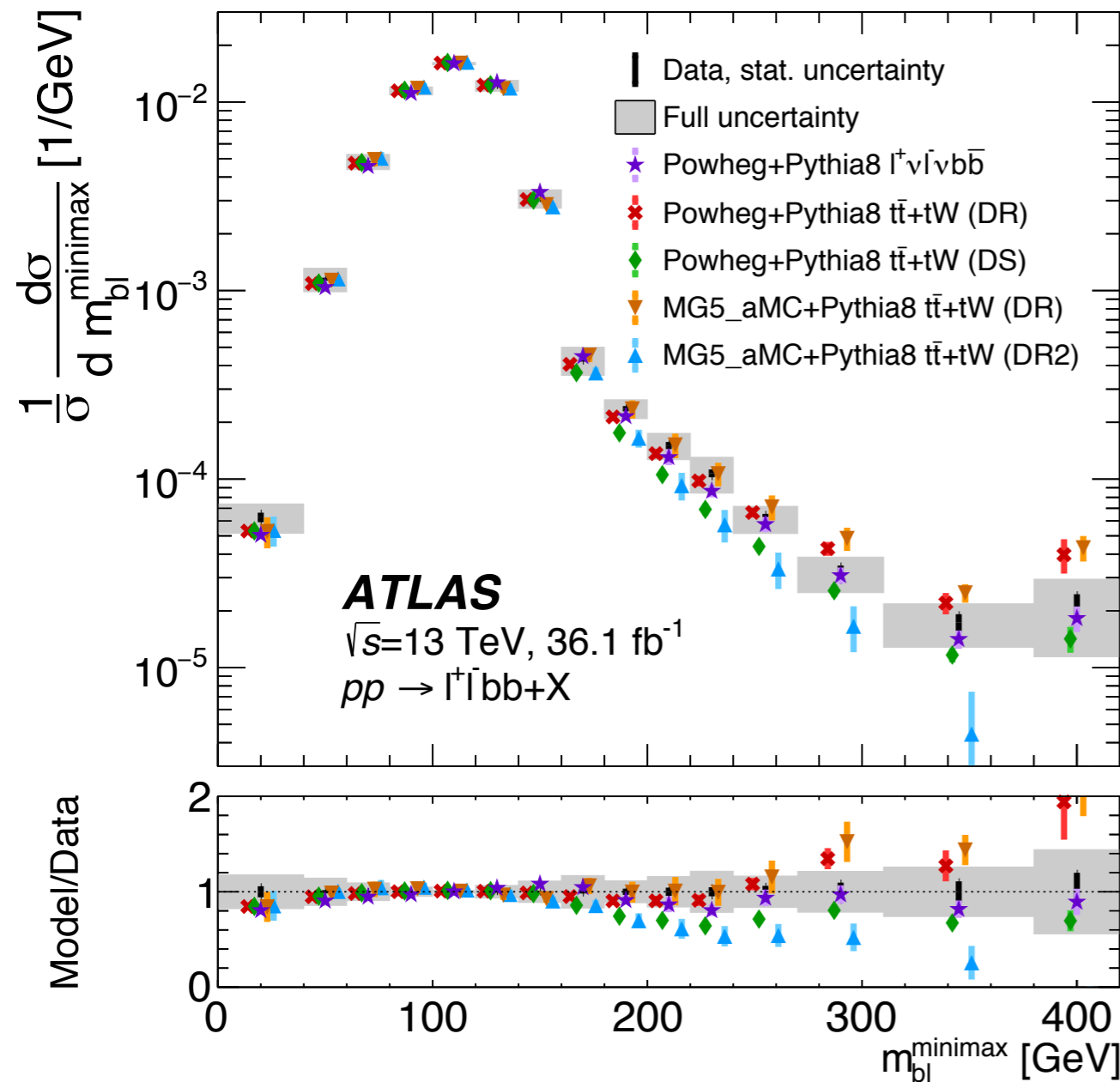
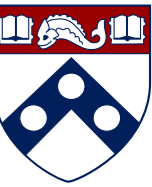
Uncertainty breakdown



Uncertainty [GeV]		b_bbar_4l	MG5_aMC@NLO
Experimental		+0.27/-0.26	± 0.20
Theory	PDF	± 0.06	± 0.04
	Scale	± 0.10	± 0.06
	m_t	± 0.03	± 0.03
	α_s	± 0.06	± 0.04
	Combined	± 0.14	± 0.10
Simulation Stats.		± 0.04	± 0.04
Total		± 0.30	± 0.22

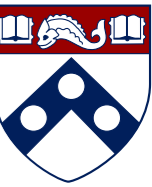
Uncertainties on the (observed) extracted widths, in GeV

Various generator comparisons



(1806.04667 / TOPQ-2017-05)

Goodness of fit



Model	Full Distribution		$m_{bl}^{\text{minimax}} > 160 \text{ GeV}$	
	χ^2 / nDOF	$p\text{-value}$	χ^2 / nDOF	$p\text{-value}$
Powheg+Pythia8 $t\bar{t} + tW$ (DR)	10 / 14	0.71	8.5 / 8	0.40
Powheg+Pythia8 $t\bar{t} + tW$ (DS)	10 / 14	0.77	6.6 / 8	0.56
Powheg+Pythia8 $\ell^+ \nu \ell^- \nu bb$	5.9 / 14	0.92	2.0 / 8	0.95
MG5_aMC+Pythia8 $t\bar{t} + tW$ (DR1)	26 / 14	0.14	13 / 8	0.17
MG5_aMC+Pythia8 $t\bar{t} + tW$ (DR2)	36 / 14	0.02	20 / 8	0.08
Powheg+Herwig++ $t\bar{t} + tW$ (DR)	26 / 14	0.07	7.3 / 8	0.48
MG5_aMC+Herwig++ $t\bar{t} + tW$ (DR)	30 / 14	0.04	11 / 8	0.23
Powheg+Pythia6 $t\bar{t} + tW$ (DR)	14 / 14	0.49	11 / 8	0.23
Powheg+Pythia6 $t\bar{t} + tW$ (DS)	14 / 14	0.49	10 / 8	0.32
MG5_aMC+Pythia8 (LO) $WWbb$	12 / 14	0.68	8.2 / 8	0.42
MG5_aMC+Pythia8 (LO) $WWbb$, no int.	28 / 14	0.05	22 / 8	0.005

(1806.04667 / TOPQ-2017-05)