

# Dynamic Jet Vetoes for Multi-lepton VBS(F) Searches <sup>1</sup>

VBS+MPI@Vienna

Richard Ruiz

Center for Cosmology, Particle Physics, and Phenomenology (CP3)  
Universite Catholique de Louvain

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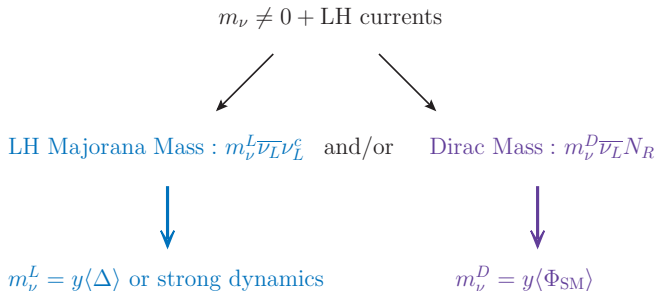
<sup>1</sup>w/ folks from Durham+Pittsburgh and NIKHEF+Paris

# Big Picture and Driving Motivation

(brief and very biased)

# Nonzero Neutrino Masses

Nonzero neutrino masses implies new degrees of freedom exist [Ma'98]:

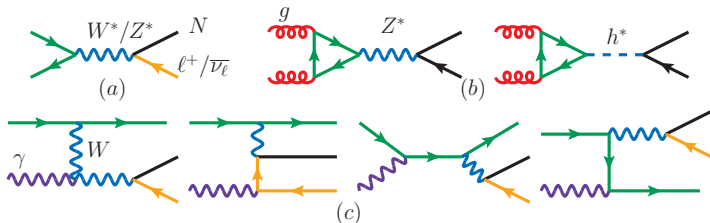


$m_\nu \neq 0 + \text{renormalizability} + \text{gauge inv.} \implies \text{new particles!}$

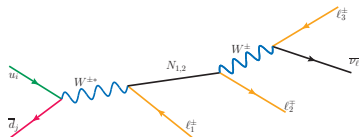
- New particles might be charged under new or old gauge symm., e.g.,  $N_R$  may have  $U(1)_{B-L}$  charge and  $\Delta_L$  is scalar  $SU(2)_L$  triplet
- Particles must couple to  $h, L$  sector, often inducing LNV/cLFV!

# Heavy Neutrinos in Hadron Collisions<sup>3</sup>

Heavy  $N$  can be produced in a variety of ways in pp collisions



$W\gamma$  fusion dominant mechanism for TeV-scale heavy neutrinos<sup>2</sup>.



● Focus on inclusive  $3\ell$  final state:

<sup>2</sup>Alva, Han, RR [1411.7305]; Degrande, Mattelaer, RR, Turner [1602.06957]

<sup>3</sup>Review on  $\nu$  mass models at colliders, Y. Cai, T. Li, T. Han, RR [1711.02180]

**Big Picture:** Collider searches for LNV (same-sign  $\ell +$  no met) / cLFV (opposite-flavor  $\ell$ ) are very sensitive since forbidden in SM.

Dominant backgrounds after typical selection criteria:

- jets mis-ID as leptons, leptons outside ID-threshold, etc.

Since 2012-13 ( $WW + 0j$  saga), new ideas for jet vetoes<sup>4</sup> in Exotica, SUSY, etc., searches have appeared, with some remarkable ( $> 10\times$ ) sensitivity improvement

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<sup>4</sup>Notably work from Bern, DESY, Durham, NIKHEF, Louvain, MIT, Paris, + others

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Reenvisioning Jet Vetoes in BSM Searches :

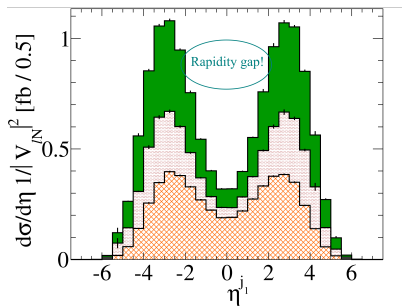
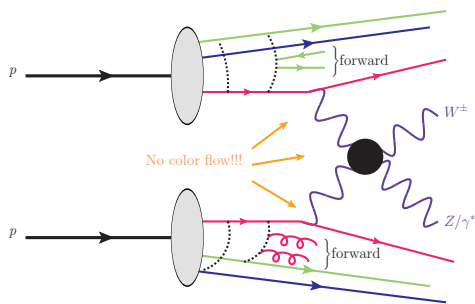
- 1 Problems with (central) jet vetoes in searches for colorless BSM
- 2 Resolution with some neat consequences
- 3 Outlook

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<sup>4</sup>Notably work from Bern, DESY, Durham, NIKHEF, Louvain, MIT, Paris, + others

## (Central) Jet Vetoes

Absence of central color flow in VBS  
 $\Rightarrow$  absence of central, high- $p_T$  jets (“rapidity gap”<sup>5</sup>)



Basis for Central Jet Veto<sup>6</sup>:

- Reject events with any jet satisfying  $p_T^j > 25 - 30$  GeV,  $|\eta^j| < 2 - 3$
- Crucial to Higgs physics but not perfect...

<sup>5</sup>Dokshitzer, Khoze, Troyan ('86)

<sup>6</sup>Barger, et al, PRD44, 2701 ('91) + PLB346, 106 ('95); Bjorken ('94)



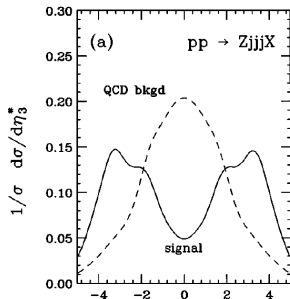
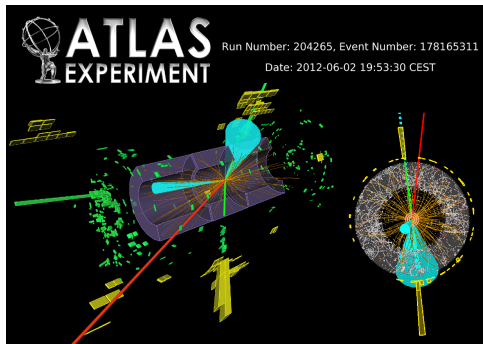
## Three Imperfections

# 1. Central Vetoes vs Weak Boson Scattering

Despite rapidity gap, VBS **not immune** to central jet veto

- Central region for ( $VV$ )-system  $\neq$  central region for lab-frame
- Motivation for alternative gap definitions<sup>7</sup>, e.g.,  $\eta^* = \eta^{j3} - \left(\frac{\eta^{j1} + \eta^{j2}}{2}\right)$

**Example:** 2012 VBF  $\rightarrow h \rightarrow \tau_{had} \tau_{\mu}$  candidate with 1 very **central**  $j^{VBS}$ !



D. Rainwater, et al.,  
PRD 54, 6680 (1996)  $\eta_3^* = \eta_3 - \bar{\eta}$

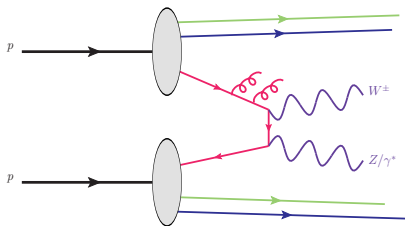
<sup>7</sup>Rainwater, Szalapski, Zeppenfeld, PRD54, 6680 (1996)

## 2. Static Vetoes vs $pp \rightarrow WW + 0j$

Difficulty with  $pp \rightarrow WW + 0j$  is more technical, but more widespread:

- Analysis-quality jets ( $j_k$ ) are "hard" and "central(ish)"
- $\implies pp \rightarrow WW + 0j$  is inclusive with respect to  $p_T^{j_k} < p_T^{\text{Veto}}$

$$\sigma(WW + 0j) = \sigma(pp \rightarrow WW + X) - \sigma(pp \rightarrow WW + 1j + X)$$



Sensitivity to higher-order  $\mathcal{O}(\alpha_s^k)$  terms grows if  $p_T^{\text{Veto}}$  too restrictive:

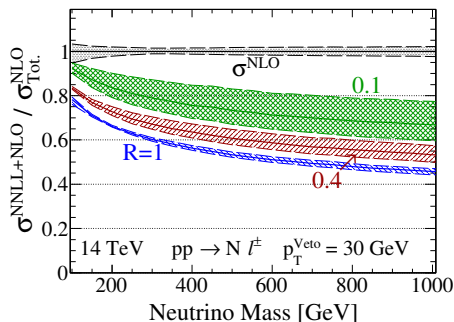
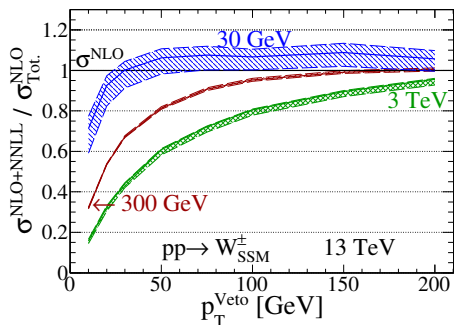
$$\begin{aligned} \sigma(WW + 1j + X) &\sim \int \frac{dPS_1 \alpha_s}{(p_f + p_g)^2 (p_{f'} + p_{g'})^2} \sim \int \frac{\alpha_s d \cos \theta dE_g^2}{E_f E_g (1 - \cos \theta') E_{f'} E_{g'} (1 - \cos \theta)} \\ &\sim \int \frac{\alpha_s d\theta^2 dE_g^2}{\theta^2 E_g^2} \sim \alpha_s \log^2(M_{WW}/p_T^{\text{Veto}})^2 \rightarrow \infty, \text{ if } (M_{WW}/p_T^{\text{Veto}}) \rightarrow 0 \end{aligned}$$

## 2.5 Static Vetoes vs Mass<sup>8</sup>

**Plotted:** veto efficiency  $\varepsilon(p_T^j < p_T^{\text{Veto}}) = \sigma^{\text{NLO+NNLL}} / \sigma^{\text{NLO}}$  for

(L) SSM  $pp \rightarrow W'$  ( $M_{W'} = 30 - 3000$  GeV)

(R)  $pp \rightarrow N\ell$



- Significant dependence on hard scale / mass scale
- Utility of jet veto unclear when  $p_T^{\text{Veto}} \gtrsim 60 - 70$  GeV due to top quarks

<sup>8</sup>w/ Fuks [[1701.05263](#)] and w/ Pascoli, et al [[1805.09335](#), [1812.08750](#)]

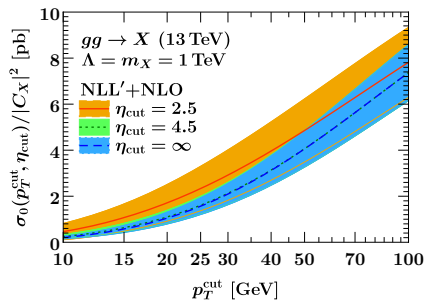
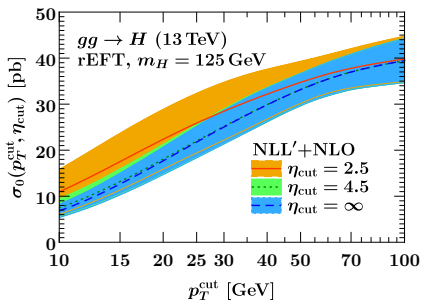
### 3. Central Vetoes vs Higgs+0j

Often, jet vetoes are restricted to  $|\eta^j| < \eta^{\text{cut}}$ , where  $\eta^{\text{cut}} \approx 2.4$

- smaller  $\eta^{\text{cut}} \implies$  less sensitivity pile up
- $\implies$  larger sensitivity to higher-order  $\mathcal{O}(\alpha_s^k)$  terms

Until recently<sup>9</sup>, analytic calculations implicitly assumed  $\eta^{\text{cut}} \rightarrow \infty$

- For  $pp \rightarrow H + 0j$ , slight change in  $\sigma$ , but big change to uncertainty



<sup>9</sup>very nice study by Michel, Pietrulewicz, and Tackmann [[1810.12911](#)]

# Dynamic Jet Vetoes<sup>10</sup>:

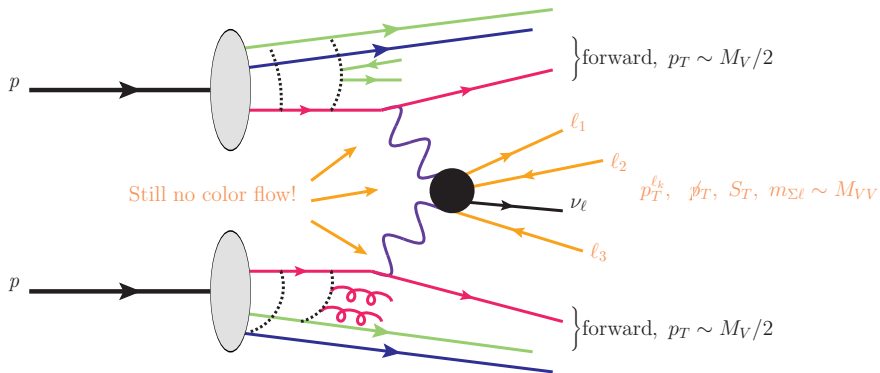
## Reenvisioning the Jet Veto

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<sup>10</sup>w/ Pascoli, et al [[1805.09335](#)]

# Leptons in Vector Boson Scattering

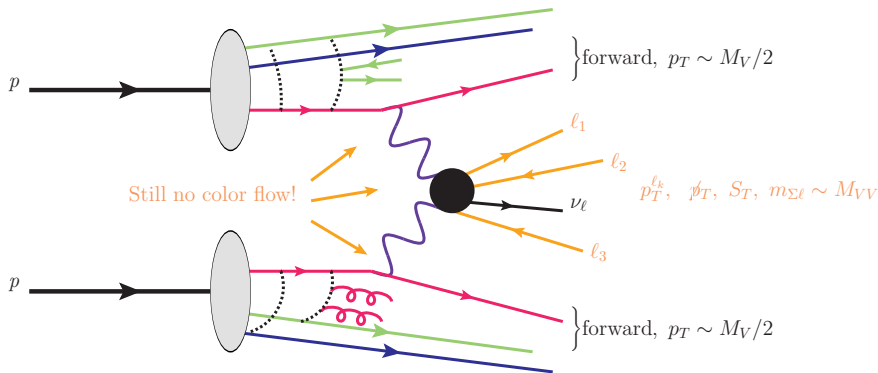
For production of leptons in VBS:  $p_T^{\ell k}, S_T \sim M_{VV} \gg p_T^{jVBS} \sim M_V/2$



<sup>11</sup>Inspired by CMS using the ratio  $r_j^\ell = (p_T^\ell/p_T^j)$  for lepton isolation [[1701.06940](#)]

# Leptons in Vector Boson Scattering

For production of leptons in VBS:  $p_T^{\ell k}, S_T \sim M_{VV} \gg p_T^{jVBS} \sim M_V/2$



**An idea<sup>11</sup>:** on event-by-event basis, set  $p_T^{\text{Veto}} = p_T^{\ell_1}$

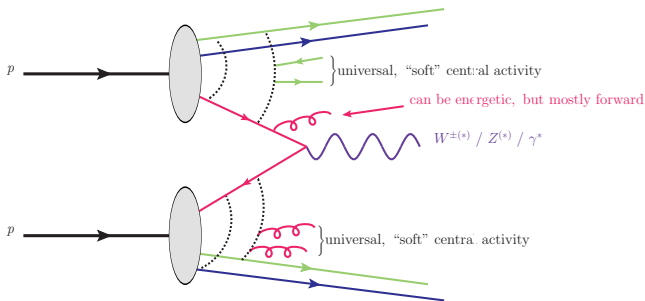
- VBS events pass by construction.
- What about other color-singlet processes, e.g., high-mass Drell-Yan?

<sup>11</sup>Inspired by CMS using the ratio  $r_j^\ell = (p_T^\ell/p_T^j)$  for lepton isolation [[1701.06940](#)]



# Leptons in High-Mass Drell-Yan and Gluon Fusion

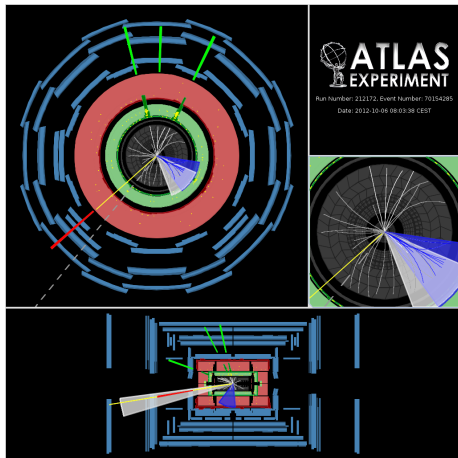
For production of leptons in DY:  $p_T^{\ell k}, S_T \sim Q \gg p_T^j \sim$  Sudakov shoulder



On event-by-event basis, set  $p_T^{\text{Veto}} = p_T^{\ell_1}$

- High-mass, DY- and GF-type processes pass by construction
- Since  $(Q/p_T^{\text{Veto}}) \sim 1$ , higher-order jet veto logarithms negligible
- What about other background processes, e.g., top quarks?

# Top Quark Background vs Dynamical Jet Vetoes



$pp \rightarrow t\bar{t}Z \rightarrow 1\mu + 3e + 2j_b + \cancel{E}_T$   
candidate event [[1509.05276](#)]

Classic kinematics:

- $m_{ee} = 93$  GeV
- $\cancel{E}_T = 57$  GeV

Typically,

- $p_T^{e1} \sim \frac{m_t}{4} \left(1 + \frac{M_W^2}{m_t^2}\right) \sim 50$  GeV
  - $p_T^{e3} \sim \frac{M_Z}{2} \sim 45$  GeV
  - $p_T^{b1} \sim \frac{m_t}{2} \left(1 - \frac{M_W^2}{m_t^2}\right) \sim 60$  GeV
- $p_T^{b1} > p_T^{\ell k} \implies$  event vetoed!

Setting  $p_T^{\text{Veto}} = p_T^{\ell_1}$  can eliminate top background **without**  $b$ -jet tagging!

# Jet Vetoes and SM Backgrounds

Associated Top Quark Production:  $pp \rightarrow t\bar{t}l\ell, t\bar{t}l\nu, tqll$

- Typical  $p_T$  of lepton from top:  $p_T^\ell \sim \frac{m_t}{4} \left(1 + \frac{M_W^2}{m_t^2}\right) \sim 50$  GeV
- Typical  $p_T$  of  $b$  from top:  $p_T^b \sim \frac{m_t}{2} \left(1 - \frac{M_W^2}{m_t^2}\right) \sim 65$  GeV
- $p_T^\ell < p_T^b \implies$  top events vetoed without need of  $b$ -tagging

Fake Leptons:

- Low- $p_T$  jet (in  $t\bar{t}$  events) identified as  $e^\pm$  or  $\tau^\pm$
- Low- $p_T$  charged  $\ell^\pm$  from weak decays of hadrons (in  $t\bar{t}$  events)
- Color conservation  $\implies$  second jet with comparable  $p_T$  likely exist

Electroweak Production<sup>12</sup>:  $pp \rightarrow 4\ell, 3\ell\nu, WWW, WW\ell\ell$

- Jet veto  $\implies$  EW bosons at rest since no recoil
- Typical  $S_T \equiv \sum_\ell |\vec{p}_T^\ell|$  for  $3W$  or  $WZ$ :  $S_T \sim \frac{3M_W}{2} \sim 120 - 130$  GeV

<sup>12</sup>Note: by itself, a dynamic jet veto is **not** a “silver bullet”

**Proof of Concept I:**

**LHC Searches for Heavy Sterile Neutrinos<sup>13</sup>**

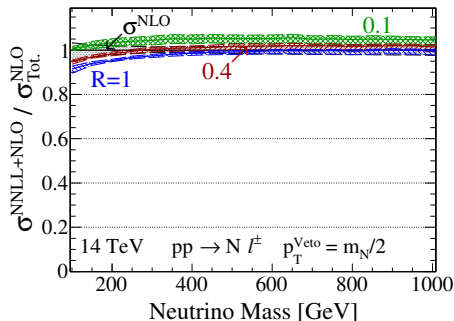
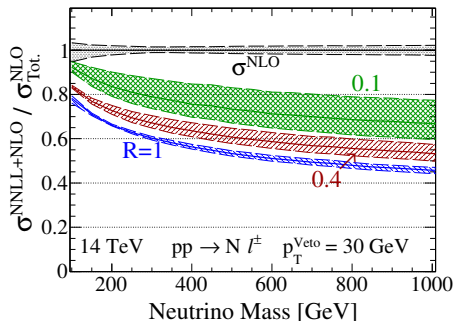
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<sup>13</sup>w/ Pascoli, etc, 8pg letter [[1805.09335](#)], 99pg analysis [[1812.08750](#)]

# Static Vetoes vs Dynamic Vetoes vs Mass

**Plotted:** veto efficiency  $\varepsilon(p_T^j < p_T^{\text{Veto}}) = \sigma^{\text{NLO+NNLL}} / \sigma^{\text{NLO}}$

for  $pp \rightarrow N \ell$  with (L)  $p_T^{\text{Veto}} = 30 \text{ GeV}$  (R)  $p_T^{\text{Veto}} = \frac{m_N}{2} \sim p_T^{\ell 1}$



- Sensitivity to higher order  $\alpha_s^k$  reduced since veto logarithms scale as  $\alpha_s \log^2(m_N/p_T^{\text{Veto}})^2 \sim \alpha_s \log^2(4) \sim \frac{4}{10} \ll 1$
- veto dependence on mass scale jet radius largely vanish

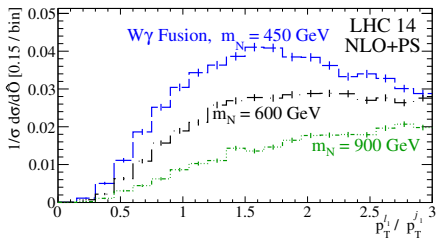
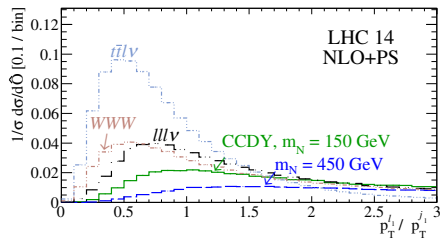
# Leptons vs Hadrons I

The proposed jet veto scheme discriminates leading  $j$  activity against leading  $\ell$  activity. Is this even sensible?

Consider the collider signature  $pp \rightarrow 3e/\mu + X$  at NLO+PS(LL)

- Generator cuts + anti- $k_T$  ( $R = 1$ ) +  $p_T^\ell > 15$  GeV +  $|\eta^\ell| < 2.4$

**Plotted:**  $r_{j_1}^{\ell_1} = p_T^{\ell_1} / p_T^{j_1}$  for (L) DY signal + SM Bkg and (R)  $W\gamma$  fusion



Separation between SM and BSM is clear

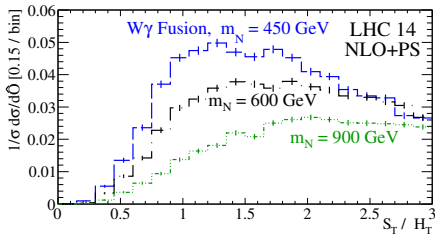
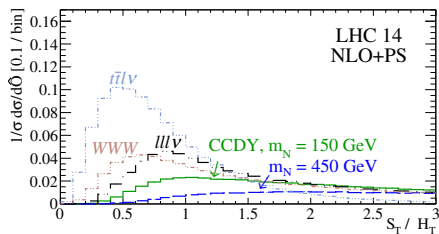
- Fine, but what about other measures of lep. and had. activity?

# Leptons vs Hadrons II

Leading  $p_T$  measures local activity. Global measures include

$$\begin{aligned} \text{Leptons:} \quad S_T^{\text{Excl.}} &= \sum_{k=1}^3 |\vec{p}_T^{\ell k}|, \quad \text{where } p_T^{\ell k} > p_T^{\ell_{k+1}} \\ \text{Hadrons:} \quad H_T^{\text{Incl.}} &\approx \sum_{k \in \{\text{clusters}\}} |\vec{p}_T^{jk}|, \quad \text{for } |\eta^{jk}| < 4.5 \end{aligned}$$

**Plotted:**  $r_{H_T}^{S_T} = S_T^{\text{Excl.}} / H_T^{\text{Incl.}}$  for (L) DY signal + SM Bkg and (R)  $W\gamma$



Similar behavior to  $r_{j_1}^{\ell_1}$ , but we will come back to this (or see [[1901.09937](#)])

# Leptons vs Hadrons III

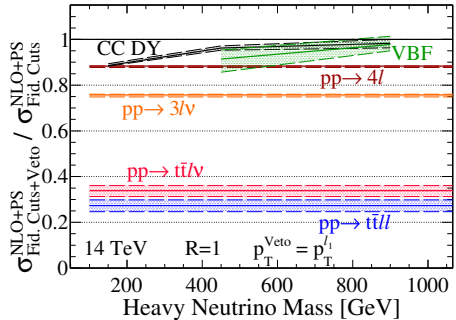
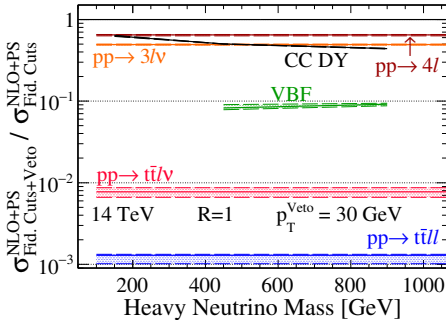
Static jet vetoes  $\implies$  poor signal efficiency for high-mass BSM

- Setting  $p_T^{\text{Veto}} = p_T^{\ell_1}$  can alleviate this

**Plotted:** veto efficiency  $\varepsilon(p_T^j < p_T^{\text{Veto}}) = \sigma_{\text{Fid. Cuts+Veto}}^{\text{NLO+PS}} / \sigma_{\text{Fid. Cuts}}^{\text{NLO+PS}}$

for (L) Veto if  $p_T^j > 30$  GeV

(R) Veto if  $p_T^j > p_T^{\ell_1}$



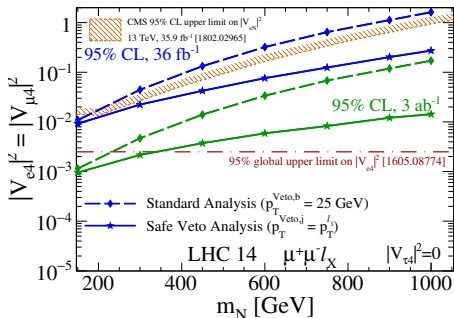
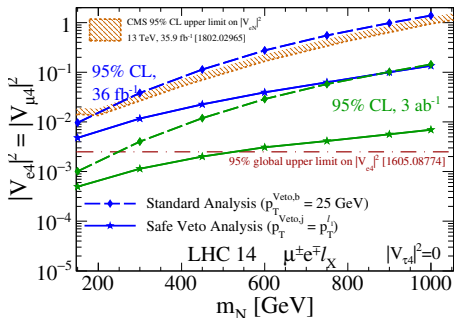
Background grows, too, but with additional lepton cuts, dynamic  $\gg$  static



# Discovery Prospects

**Signal I:**  $pp \rightarrow \mu^\pm e^\mp \ell_X + \text{MET}$  and **Signal II:**  $pp \rightarrow \mu^+ \mu^- \ell_X + \text{MET}$

**Plot:** LHC 14 sensitivity to  $(\text{coupling})^2$  vs heavy neutrino mass



Improved sensitivity up to  $10 - 11 \times$  with  $\mathcal{L} = 3 \text{ ab}^{-1}$ .

- Dash = standard search<sup>14</sup> with  $b$ -jet veto (13 TeV CMS for  $e/\mu$ )
- Solid = “improved” analysis with special type of jet veto

<sup>14</sup>More aggressive cuts on individual leptons: e.g.,  $p_T^{\ell_1} > 55 \text{ GeV}$ ,  $m_{3\ell} > 80 \text{ GeV}$

# Generalizing Dynamic Jet Vetoes:

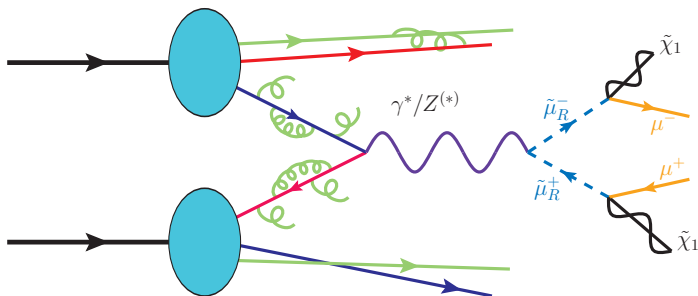
## Sleptons vs Hadrons<sup>15</sup>

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<sup>15</sup>w/ Fuks, Nordstrom, & Williamson [[1901.09937](#)]

# Sleptons without Hadrons

Clear from heavy neutrino study that a dynamic jet veto can improve slepton searches



Decided to have fun and explore alternative measures  
of hadronic vs leptonic activities  
- Generalizing definition of a dynamic jet veto

# Leptons vs Hadrons IV

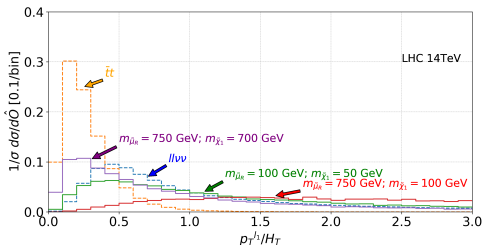
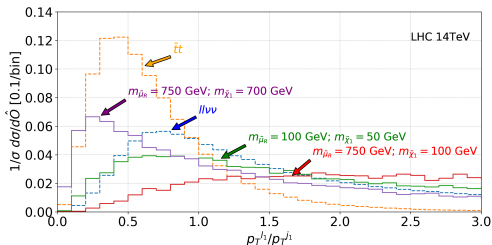
Consider the signature  $pp \rightarrow \mu^+ \mu^- + X$  at NLO+PS(LL)

- Generator cuts +  $p_T^\ell > 15$  GeV  
 +  $|\eta^\ell| < 2.4$  + anti- $k_T$  ( $R = 1$ )

- (top)  $r_{j_1}^{\ell_1} = p_T^{\ell_1} / p_T^{j_1}$

- (bottom)  $r_{H_T}^{\ell_1} = p_T^{\ell_1} / H_T$

Turns out  $p_T^{\text{Veto}} = p_T^{\ell_1}$  is not best  
 discriminate :)

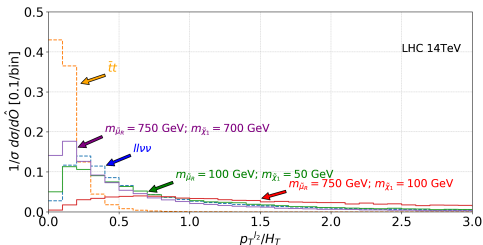
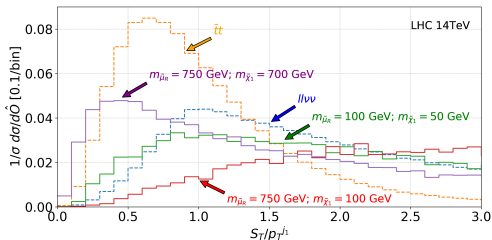


# Leptons vs Hadrons V

Some choices were rubbish (top), others were fantastic (bottom)

- (top)  $r_{j_1}^{S_T} = S_T/p_T^{j_1}$
- (bottom)  $r_{H_T}^{\ell_2} = p_T^{\ell_2}/H_T$

Separation between  $t\bar{t}$  from  $V^{(*)}V^{(*)}$  is clear :)



# Improved Discovery Prospects

As a benchmark analysis, we considered CMS's Run II search, which uses  $p_T^{\text{Veto}} = 25$  GeV for  $|\eta^j| < 2.4$  [1806.05264]

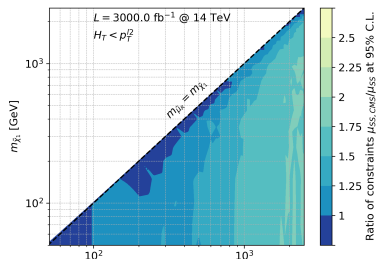
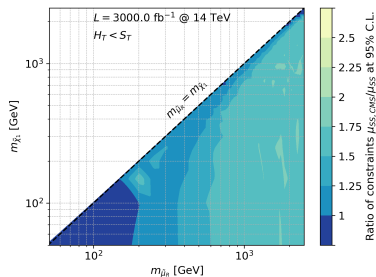
No attempt to optimize cuts:

- Relaxed  $p_T^{\ell_k}$  cuts
- swapped static veto for dynamic

Improvement was clear (lighter  $\implies$  improvement)

- (top)  $H_T^{\text{Veto}} = S_T$
- (bottom)  $H_T^{\text{Veto}} = p_T^{\ell_2}$

Using  $H_T$  and  $p_T^{\ell_k}$  give better improvement



# Summary and Next Steps

*A class of jet vetoes has been developed for high-mass, multi- $\ell$  searches, one based on comparing lepton vs hadronic activities, e.g.,  $H_T^{\text{Veto}} = p_T^{\ell_1}$*

- **New scheme reveals**  $> 90 - 95\%$  **signal acceptance with little-to-no dependence on mass scales** (contrary to previous schemes)
- Substantial reduction in QCD theory uncertainty  
 $\implies$  **less need for high-precision resummation**
- Redesigned search analysis gives better reduction of background

Application to new physics searches appears very promising:

- **Sterile neutrinos:** improved sensitivity  $\sim 10 - 11\times$
- **Sleptons:** setting  $p_T^{\text{Veto}}$ ,  $H_T^{\text{Veto}} = p_T^{\ell_k}$ ,  $S_T$  reveal generic improvement
- **To Do:** Check impact of MPI and add'l uncertainties
- **To Do:** Rapidity dependencies?

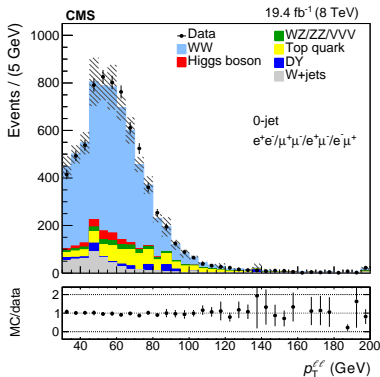
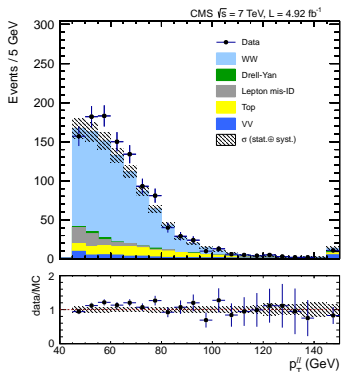


**Thank you.**



## Backup I: Jet Vetoes in LHC Collisions

Plotted is  $p_T$  of the  $(\ell_1^+ \ell_2^-)$ -system in  $pp \rightarrow W^+W^- + 0j \rightarrow \ell_1^+ \ell_2^- \nu \bar{\nu} + 0j$



- During Run I, ATLAS and CMS [1306.1126; 1507.03268] consistently disagreed with SM predictions for  $pp \rightarrow WW + 0j$
- Did not "go away" with more data, was not statistical
- Resolved by noting little suppression of ultra-low  $p_T$  QCD radiation
  - ▶ Add appropriate terms  $\implies$  excellent theory control over jet vetoes!