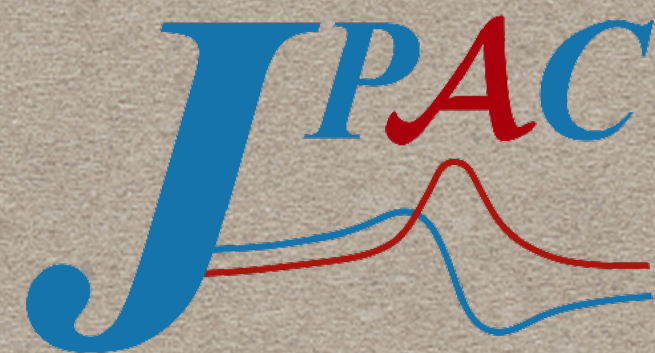


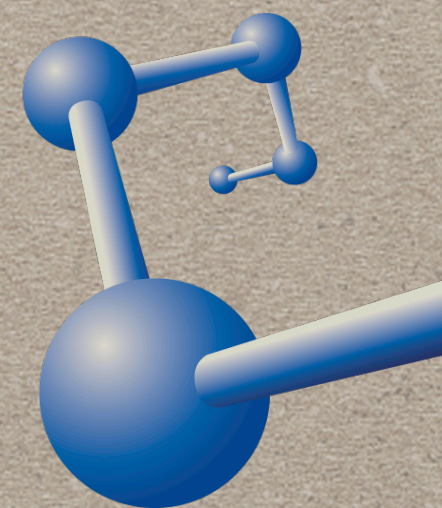
# AMPLITUDE ANALYSIS AND POLE INTERPRETATION: THE $P_c(4312)$ CASE

CÉSAR FERNÁNDEZ-RAMÍREZ

INSTITUTO DE CIENCIAS NUCLEARES – UNAM  
JOINT PHYSICS ANALYSIS CENTER (JPAC)



Instituto de  
Ciencias  
Nucleares  
UNAM



# AMPLITUDE ANALYSIS: BOTTOM-TOP APPROACH

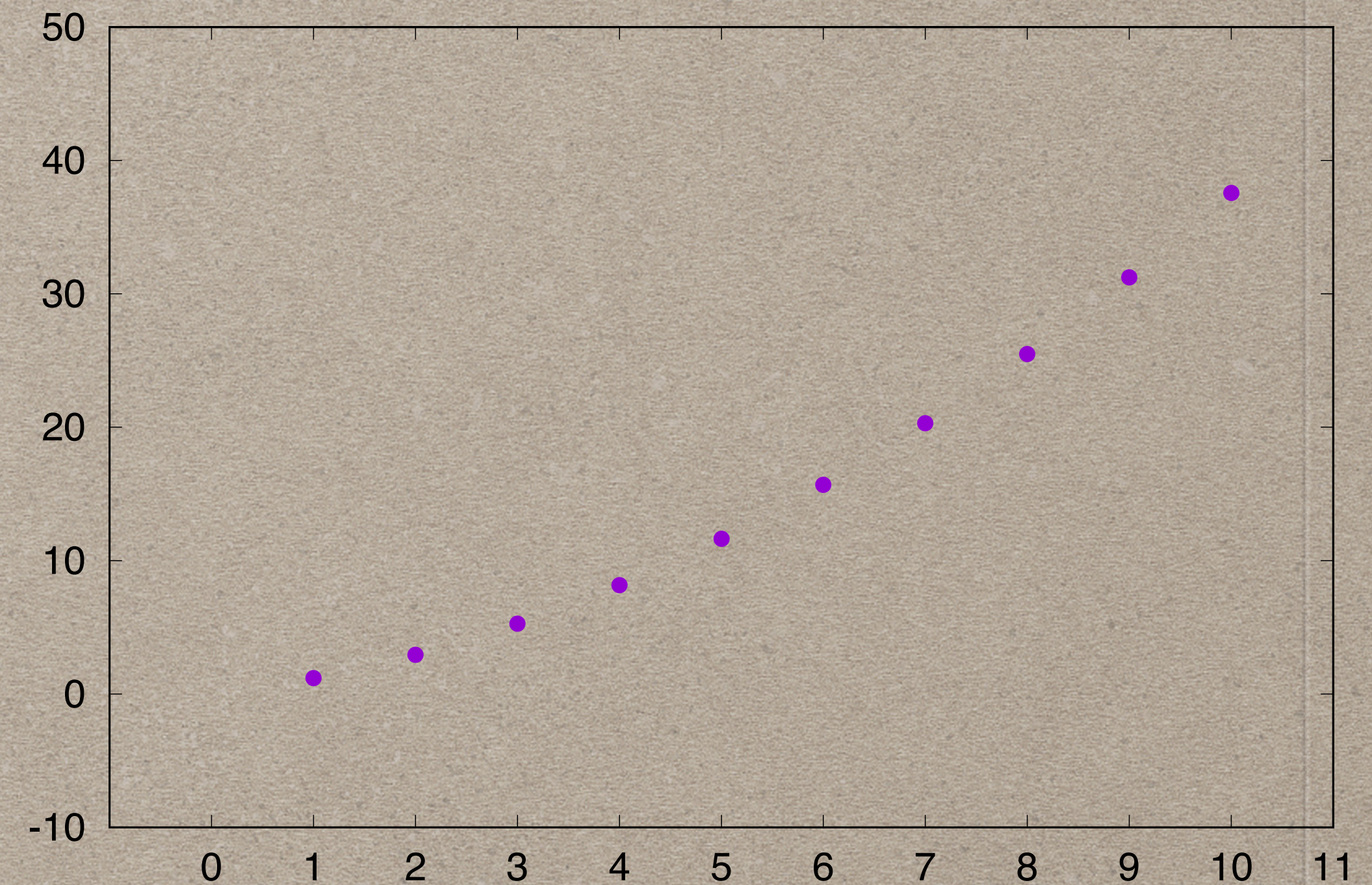
- Build the minimally-biased theory (model) with the correct physical restrictions
- Fit the experimental data and perform an error analysis
- Analytically continue the amplitude to the complex plane and the unphysical Riemann sheets
- Hunt and study poles. Two aspects:
  - Are they poles of the model only or are they also poles of the data?
  - Can we make a model-independent interpretation of the nature of the singularity?

# UNCERTAINTIES ANALYSIS: BOOTSTRAP

Alessandro will elaborate on this later

# UNCERTAINTIES ANALYSIS: BOOTSTRAP

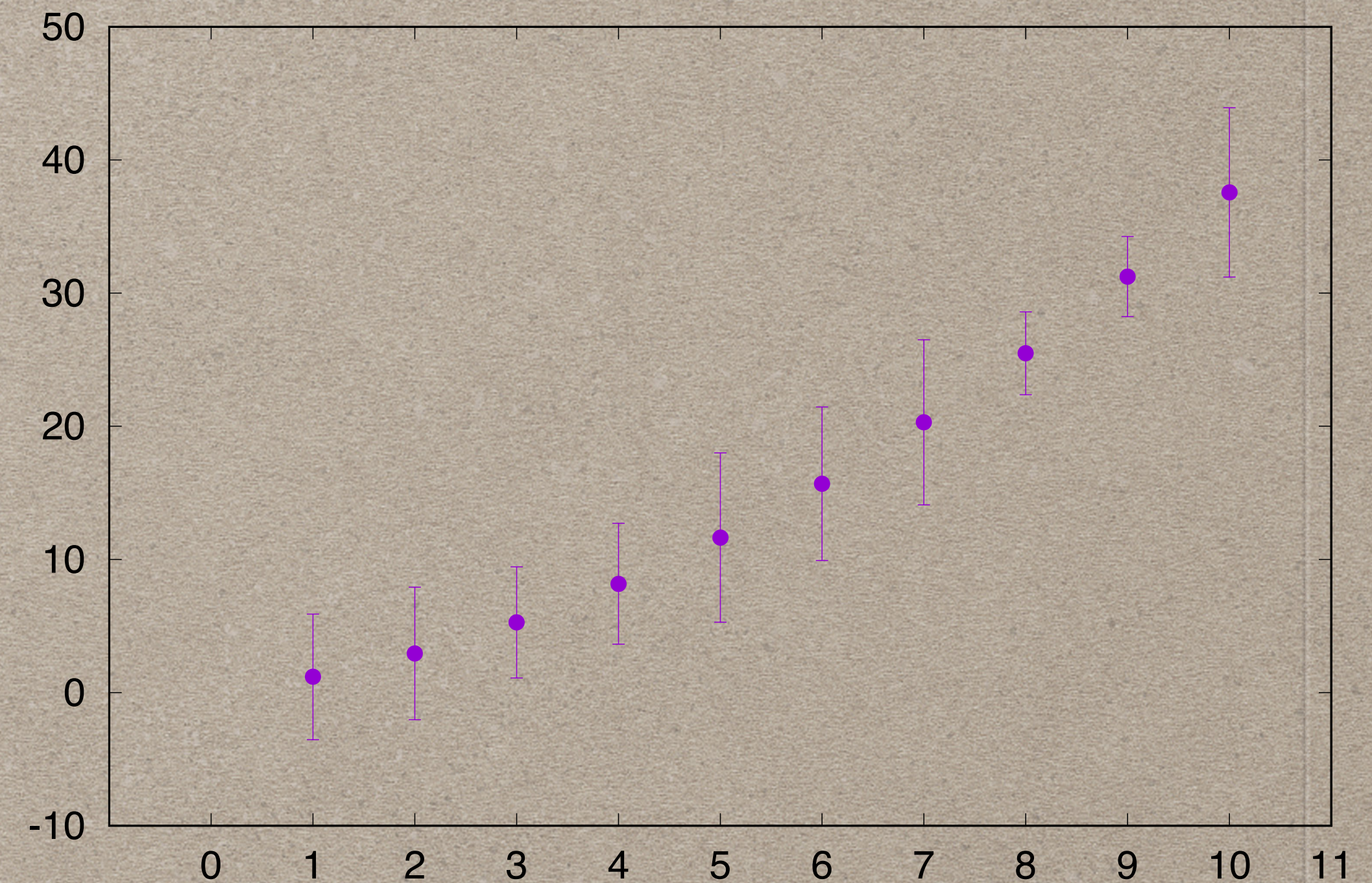
- Take the data with errors



Alessandro will elaborate on this later

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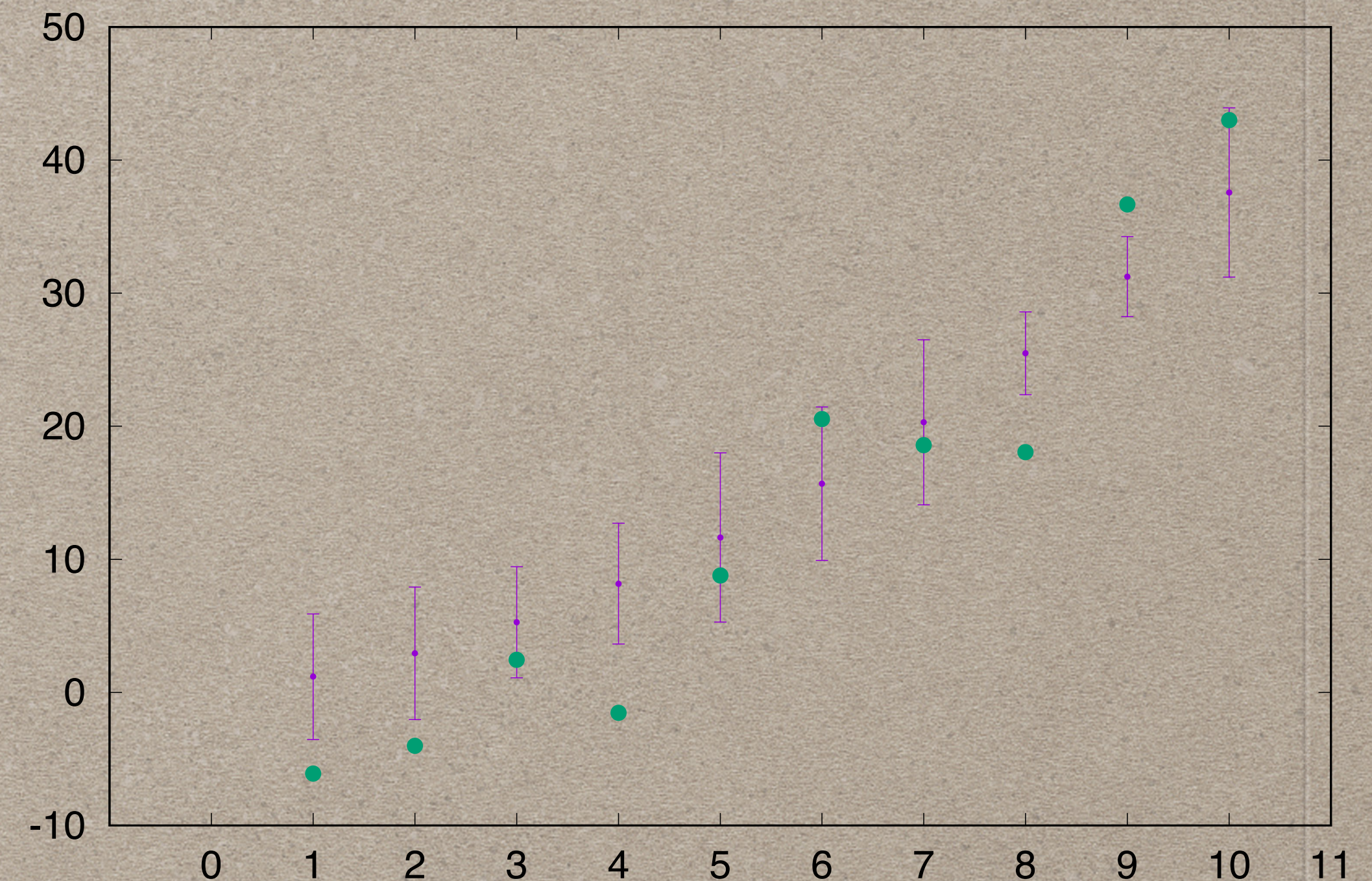
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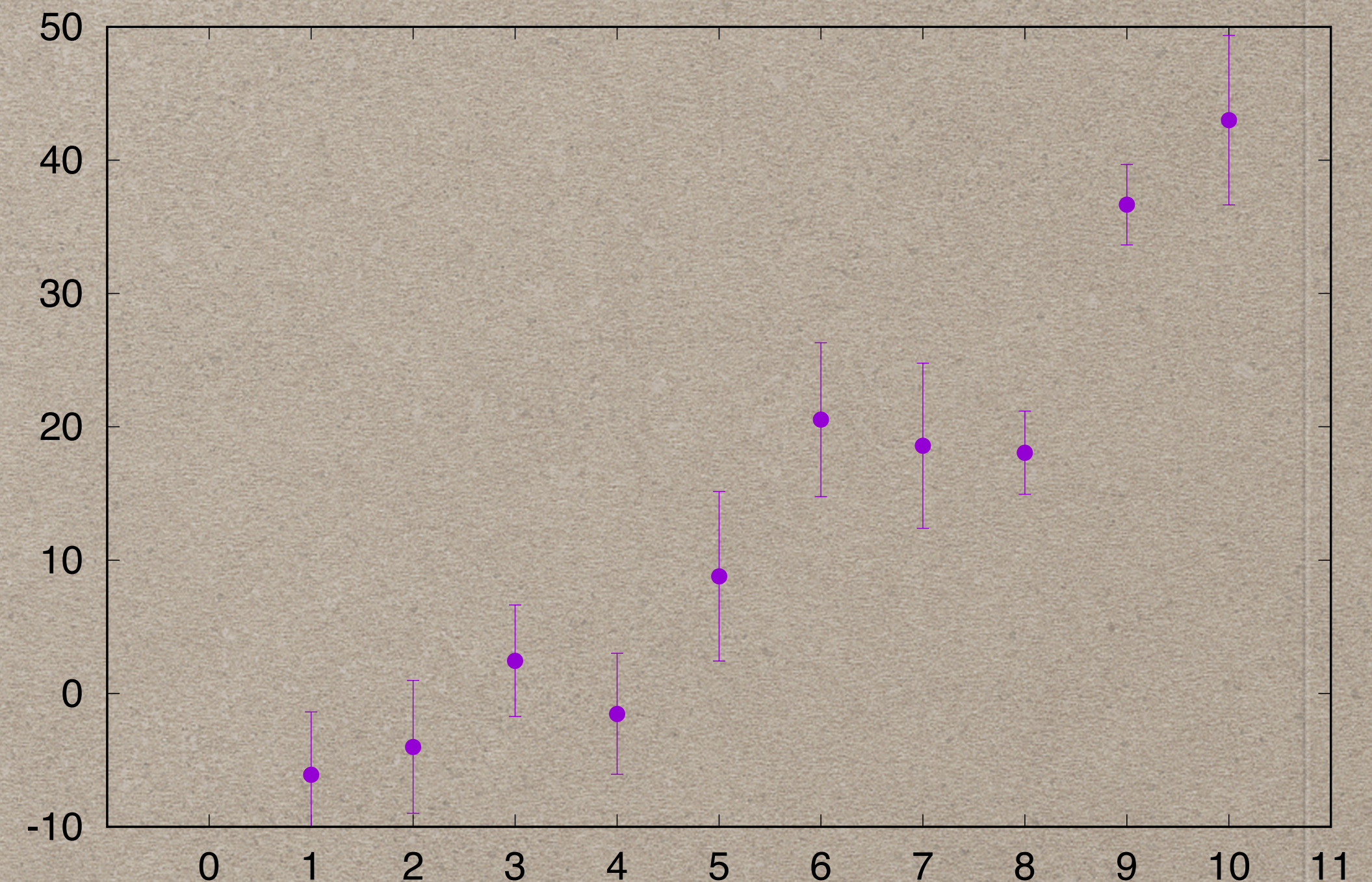
- Take the data with errors
- Randomize according to uncertainties (generate pseudodata)



Alessandro will elaborate on this later

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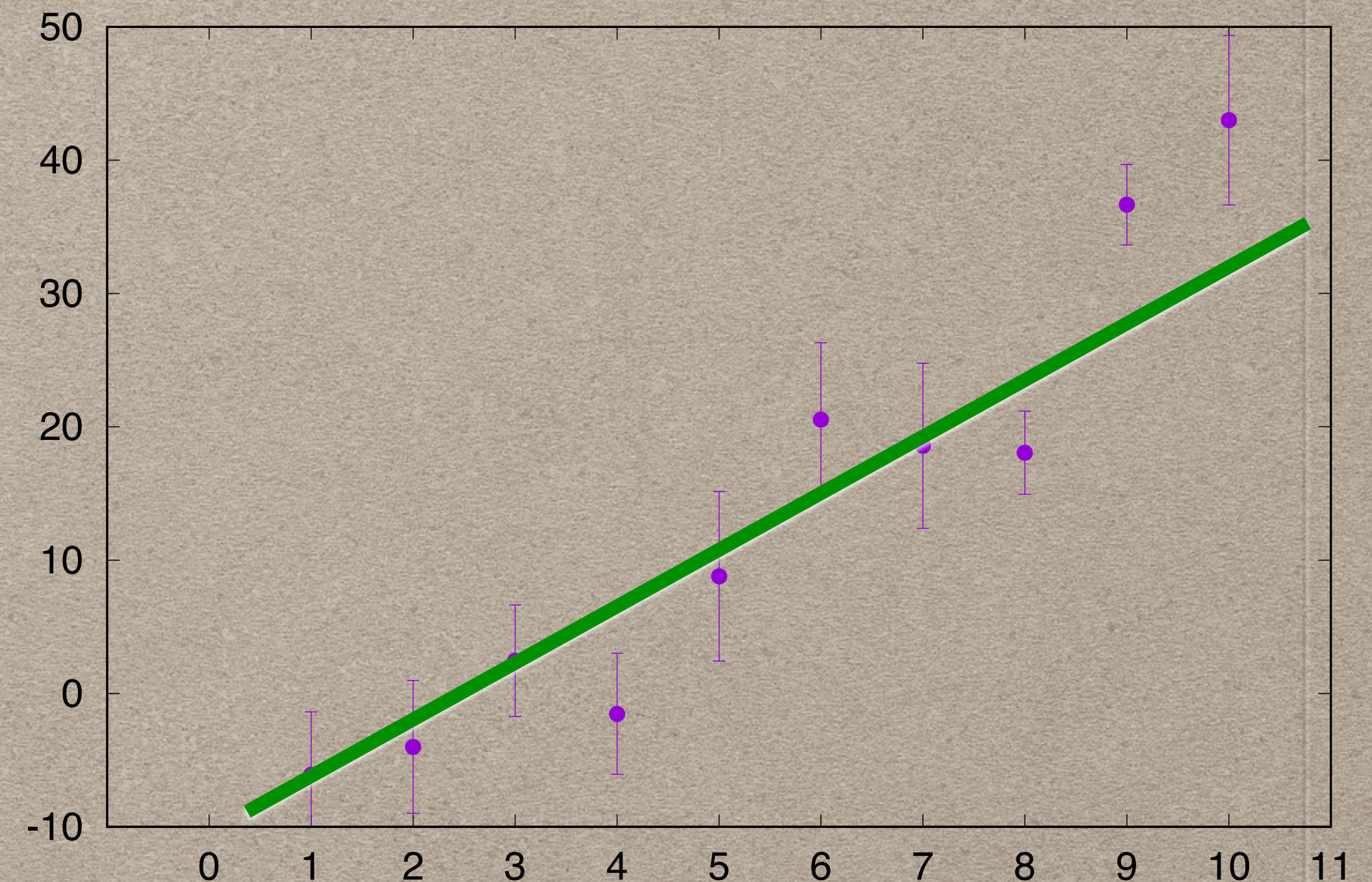
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Alessandro will elaborate on this later

# UNCERTAINTIES ANALYSIS: BOOTSTRAP

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- Fit, get parameters, compute any derivative quantity (observables, poles)

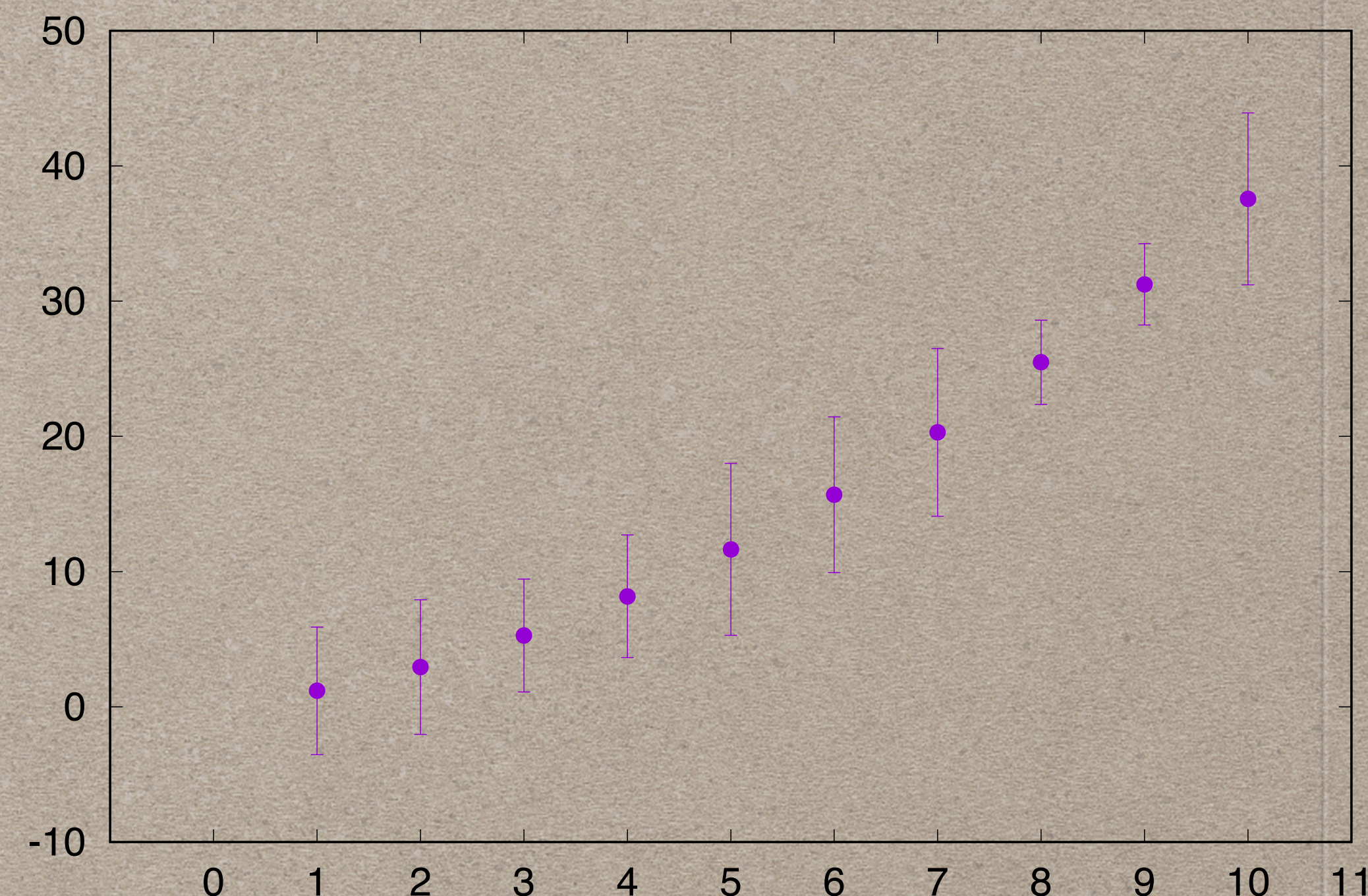


Alessandro will elaborate on this later



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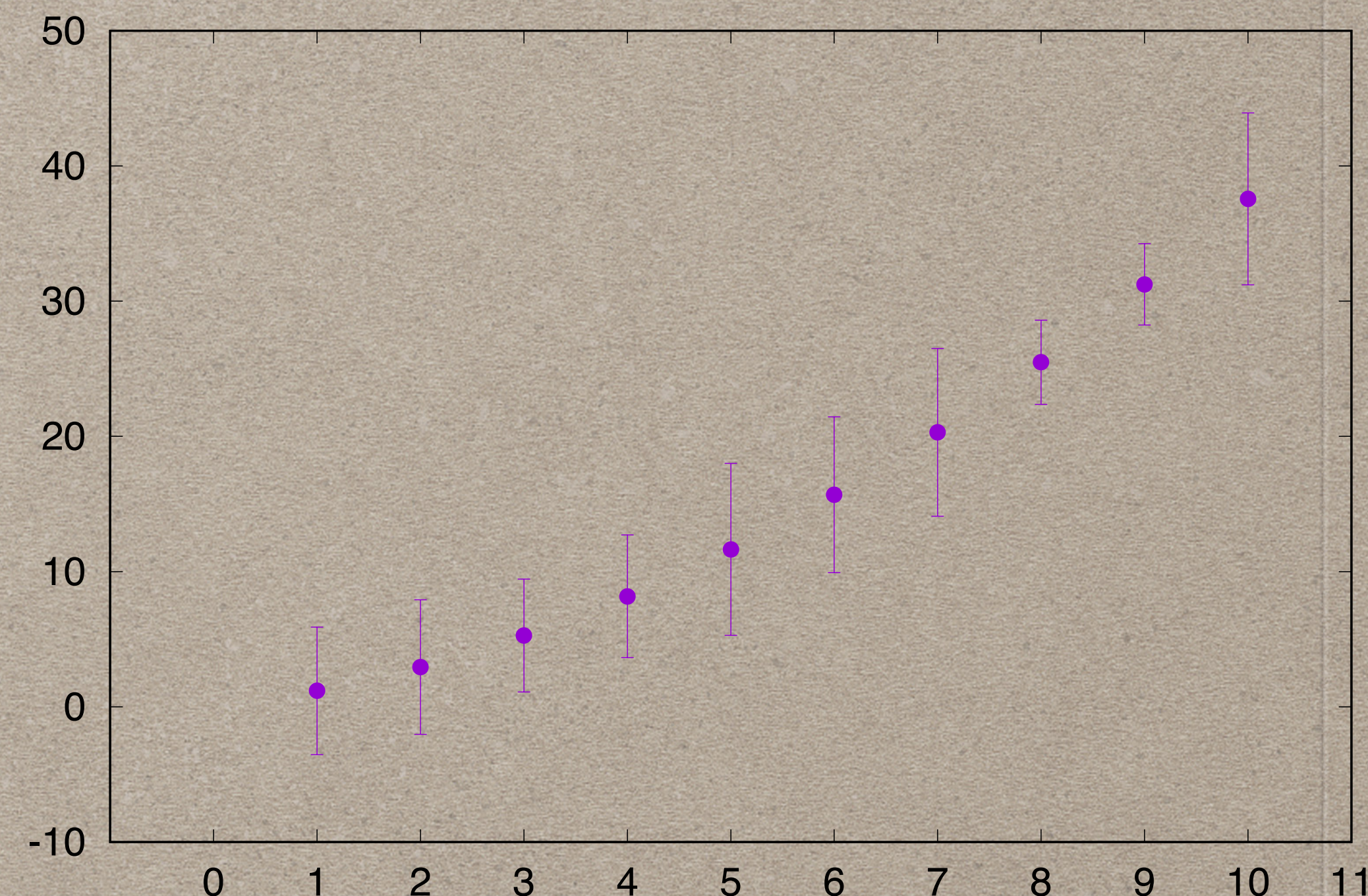
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Alessandro will elaborate on this later

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- ➔ Take the data with errors
- Randomize according to uncertainties (generate pseudodata)
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- ➔ Repeat until you have enough statistics



**You end up with N sets of parameters, and you can perform statistics on them and compute derivative quantities (poles, observables) propagating in full the errors**

Alessandro will elaborate on this later

# NATURE OF $P_c(4312)$

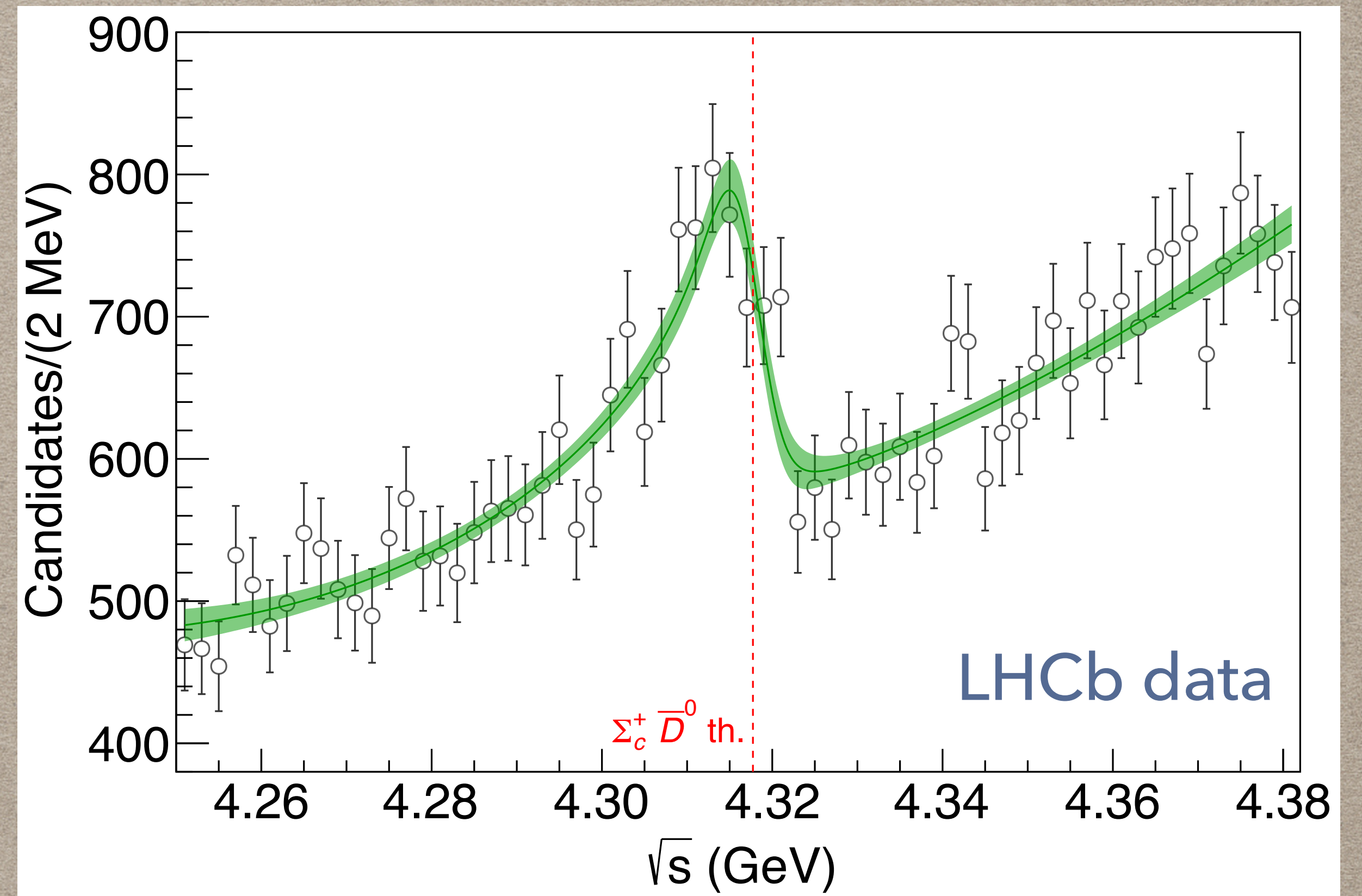
Close to a threshold

Triangle singularity


Compact pentaquark

Molecule

Virtual state

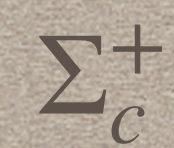
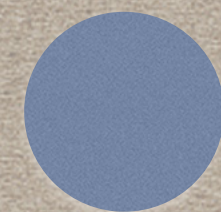


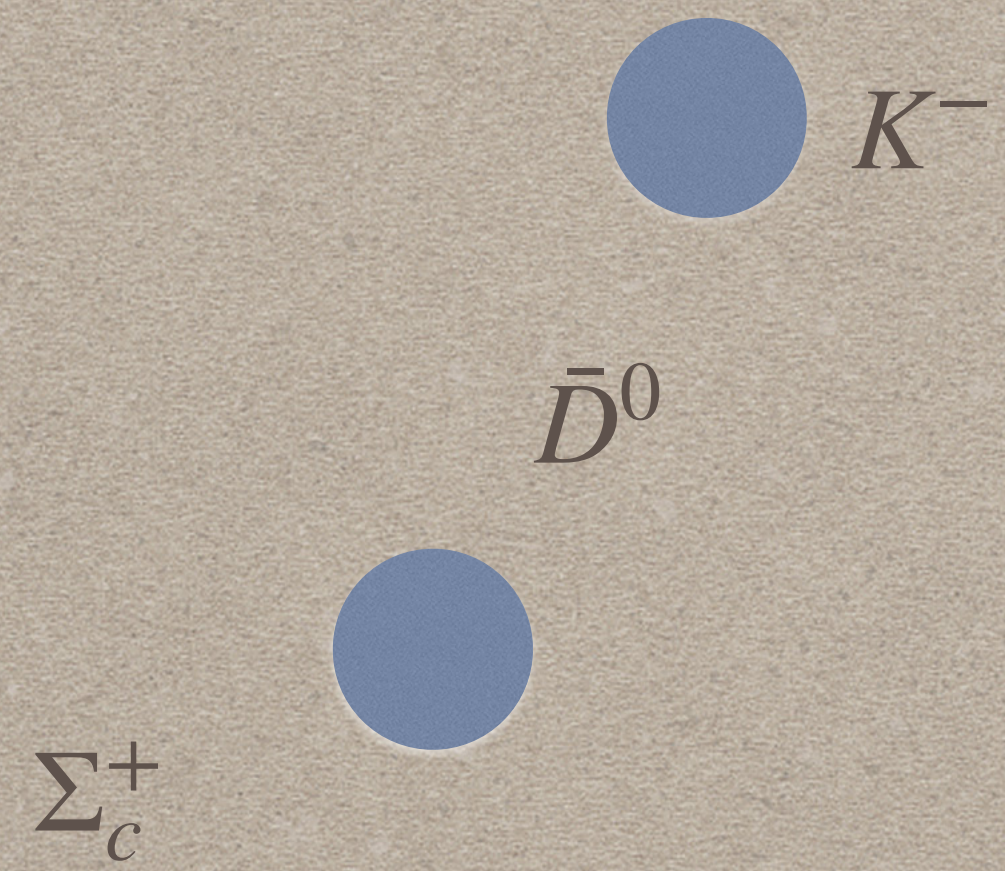
**LHCb, PRL 122 (2019) 222001**

$$\Lambda_b^0$$


$\Lambda_b^0$









$K^-$

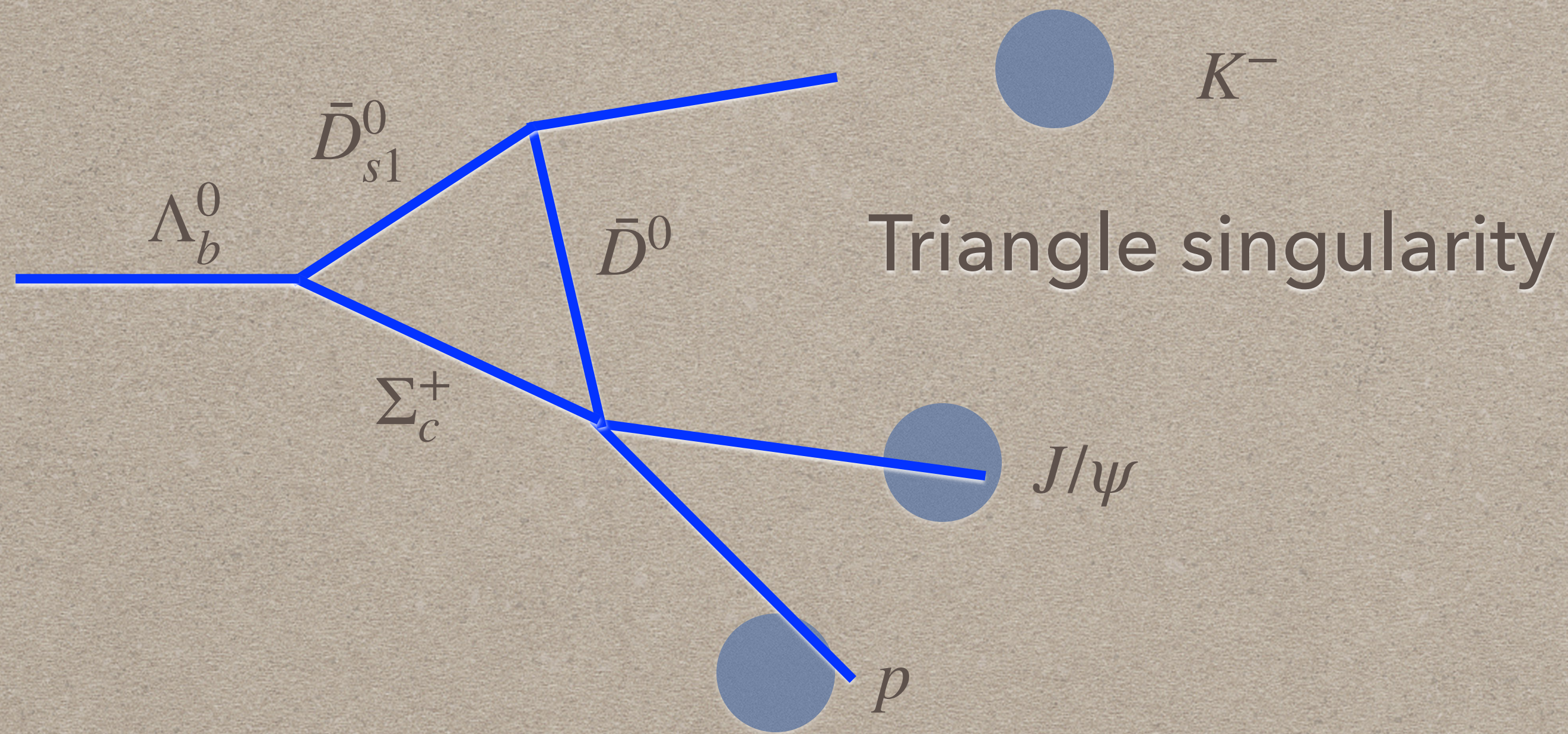


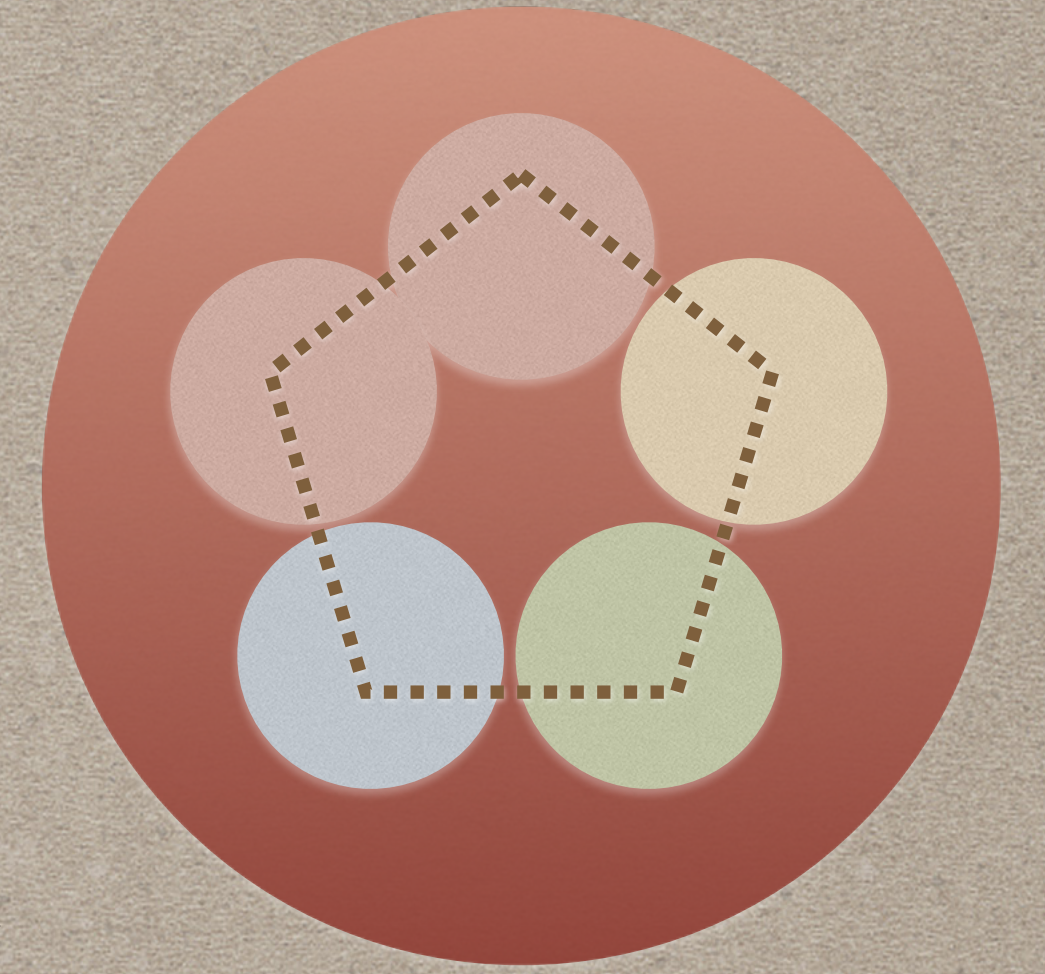
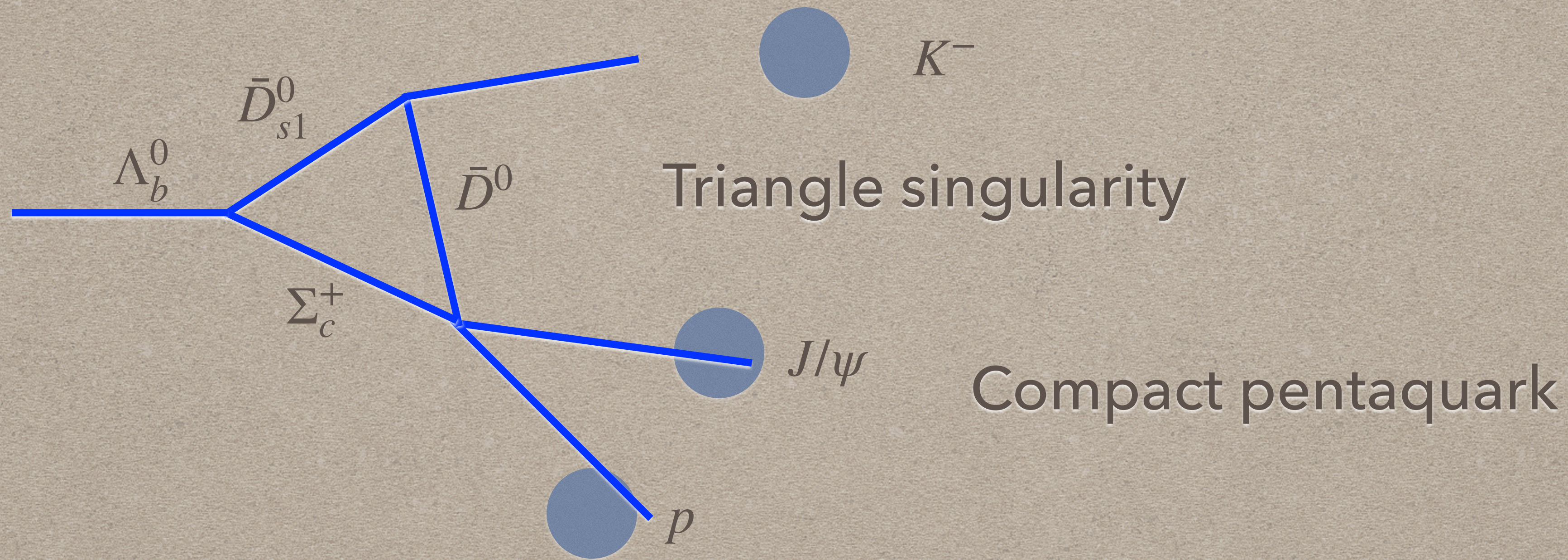
$J/\psi$

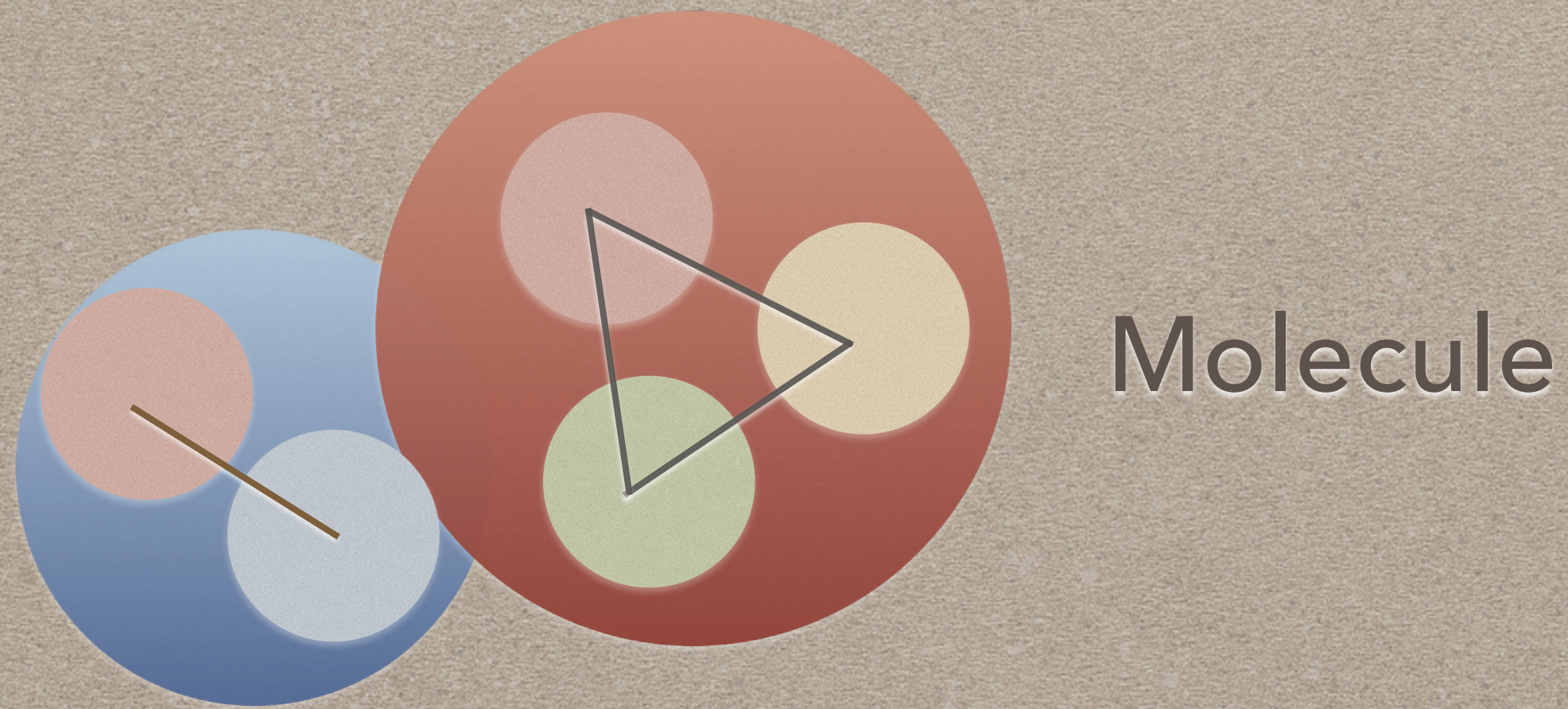
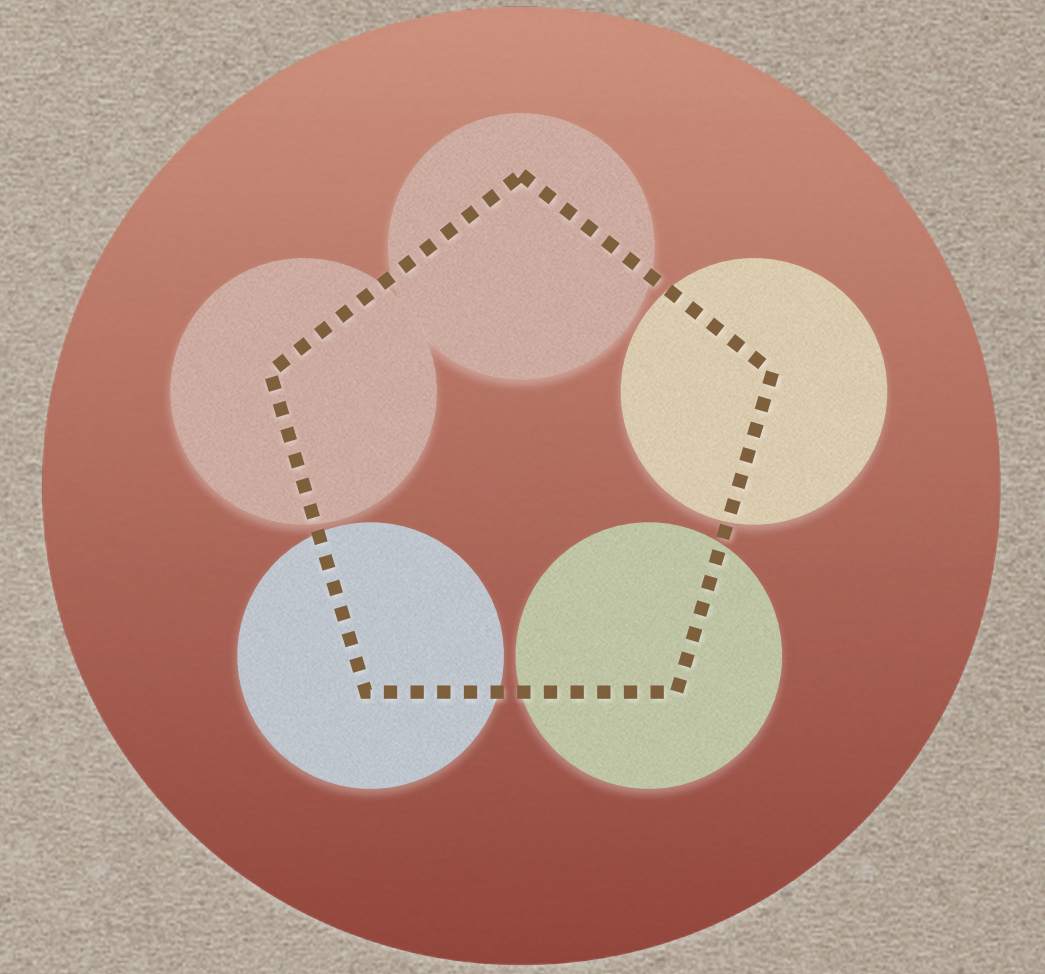
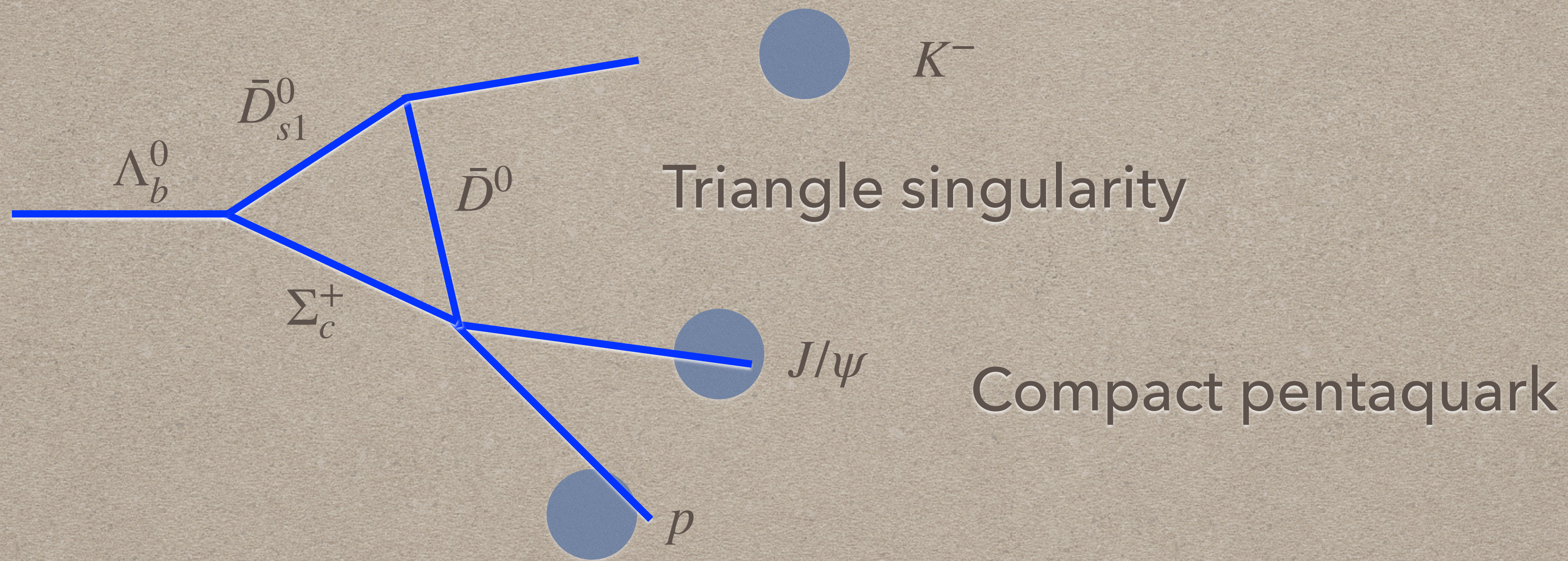


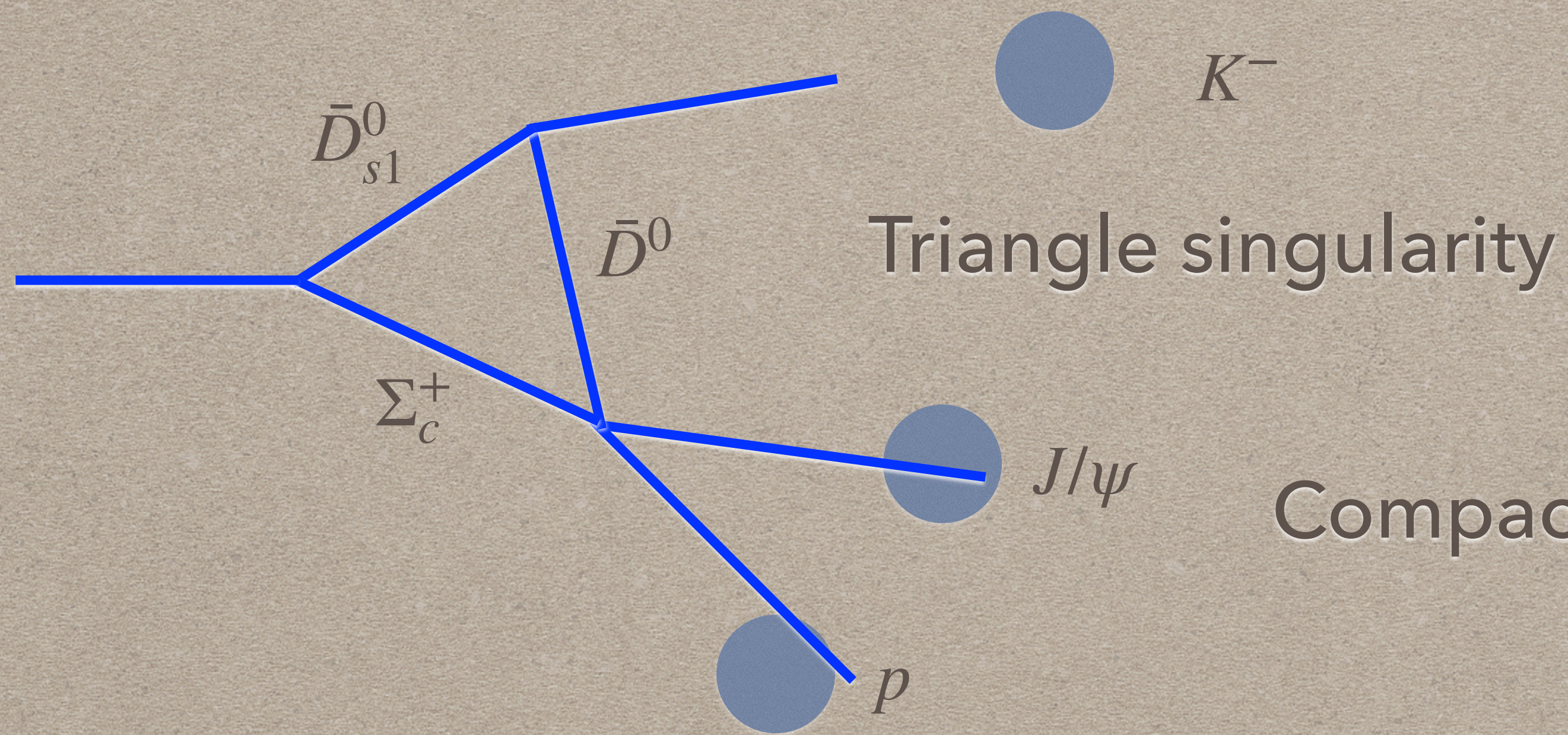
$p$



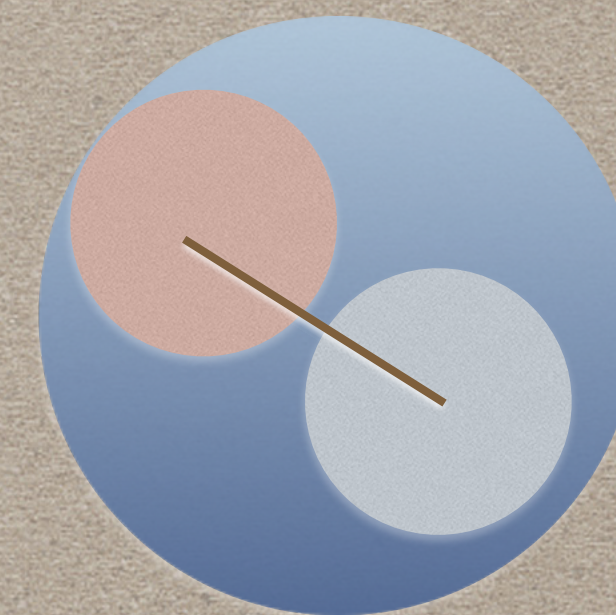
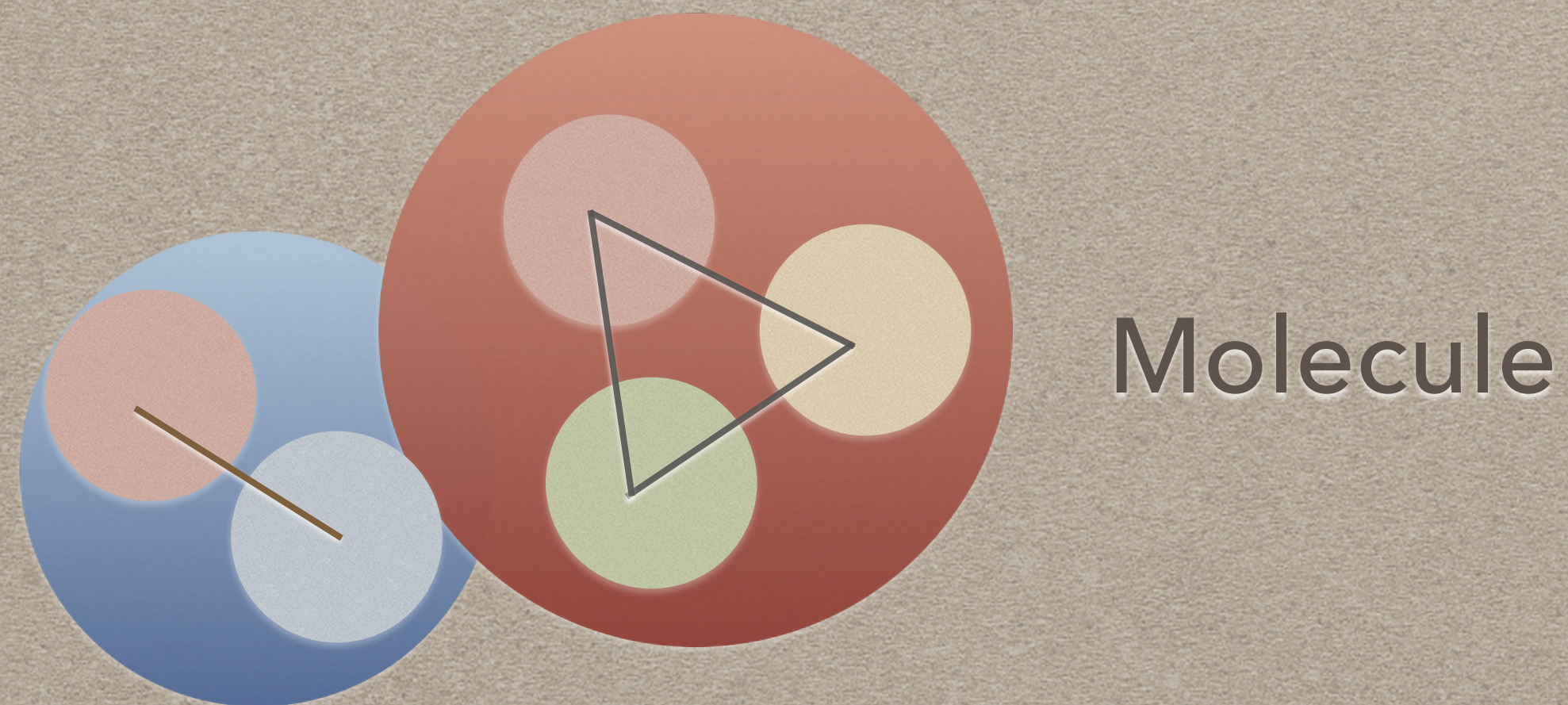
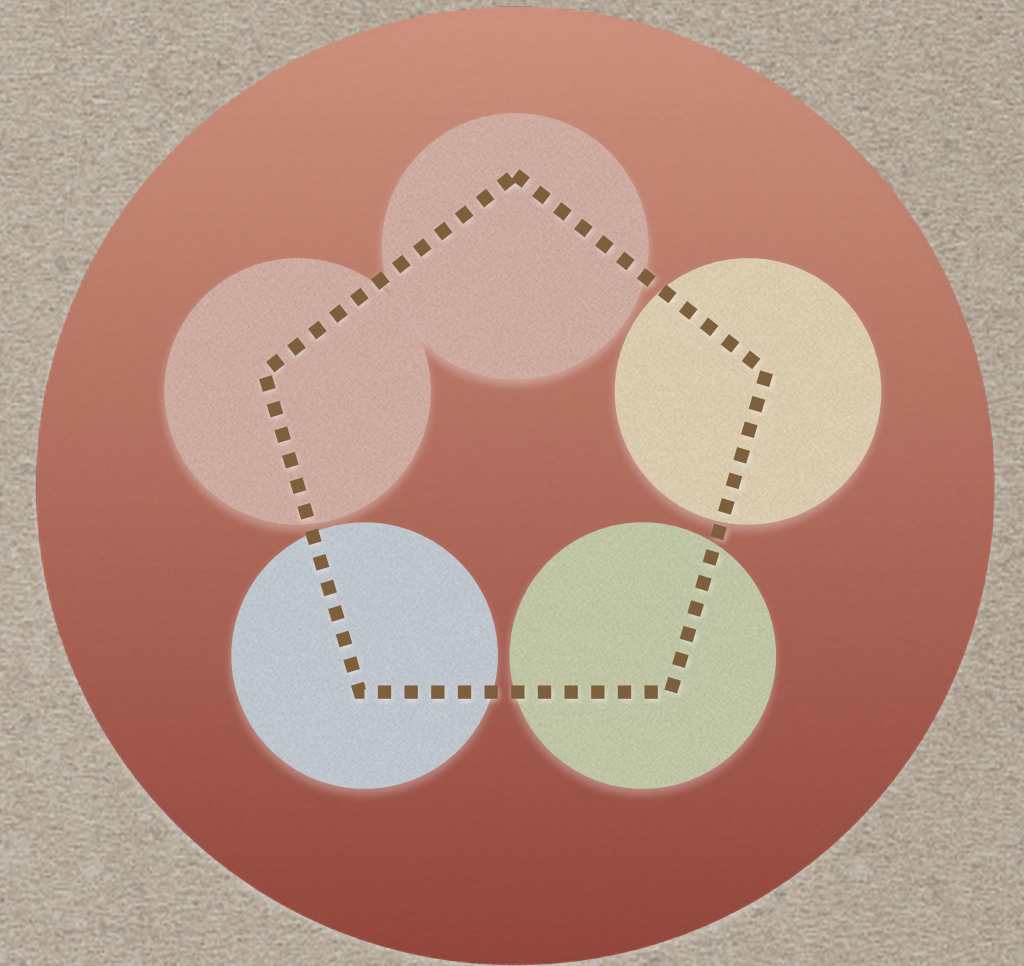




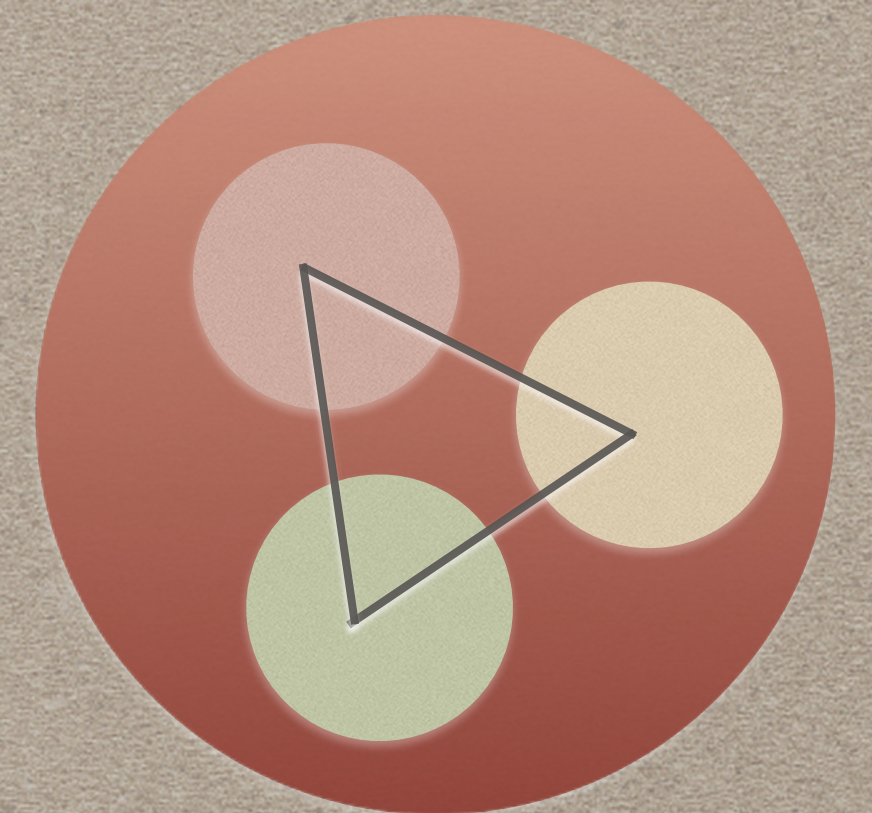




Compact pentaquark



Virtual state



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The virtual state is actually a very well-established physics case

Consider Nucleon-Nucleon interaction:

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$p + n$  (singlet)  $\rightarrow d$  (bound state, deuteron)

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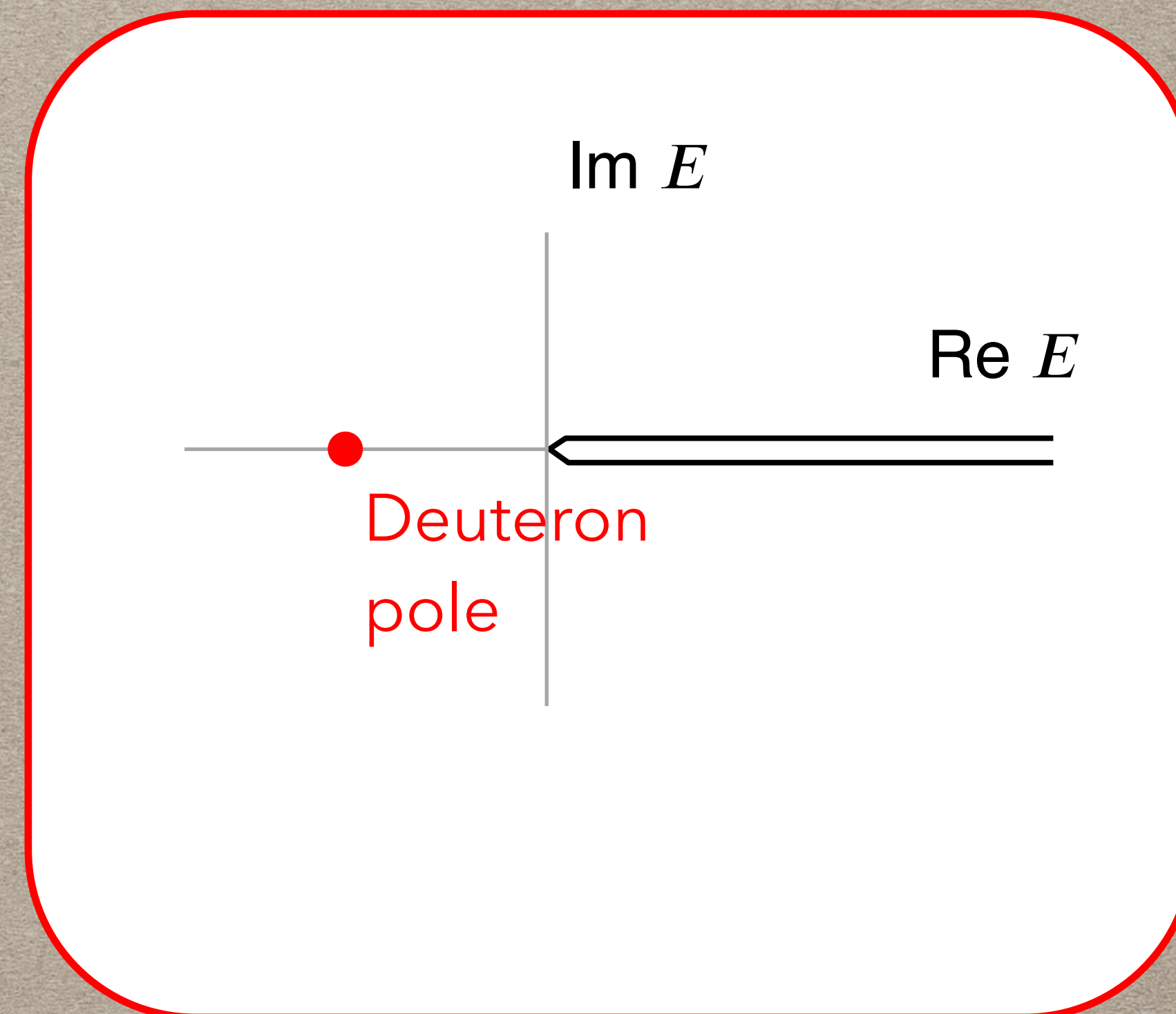
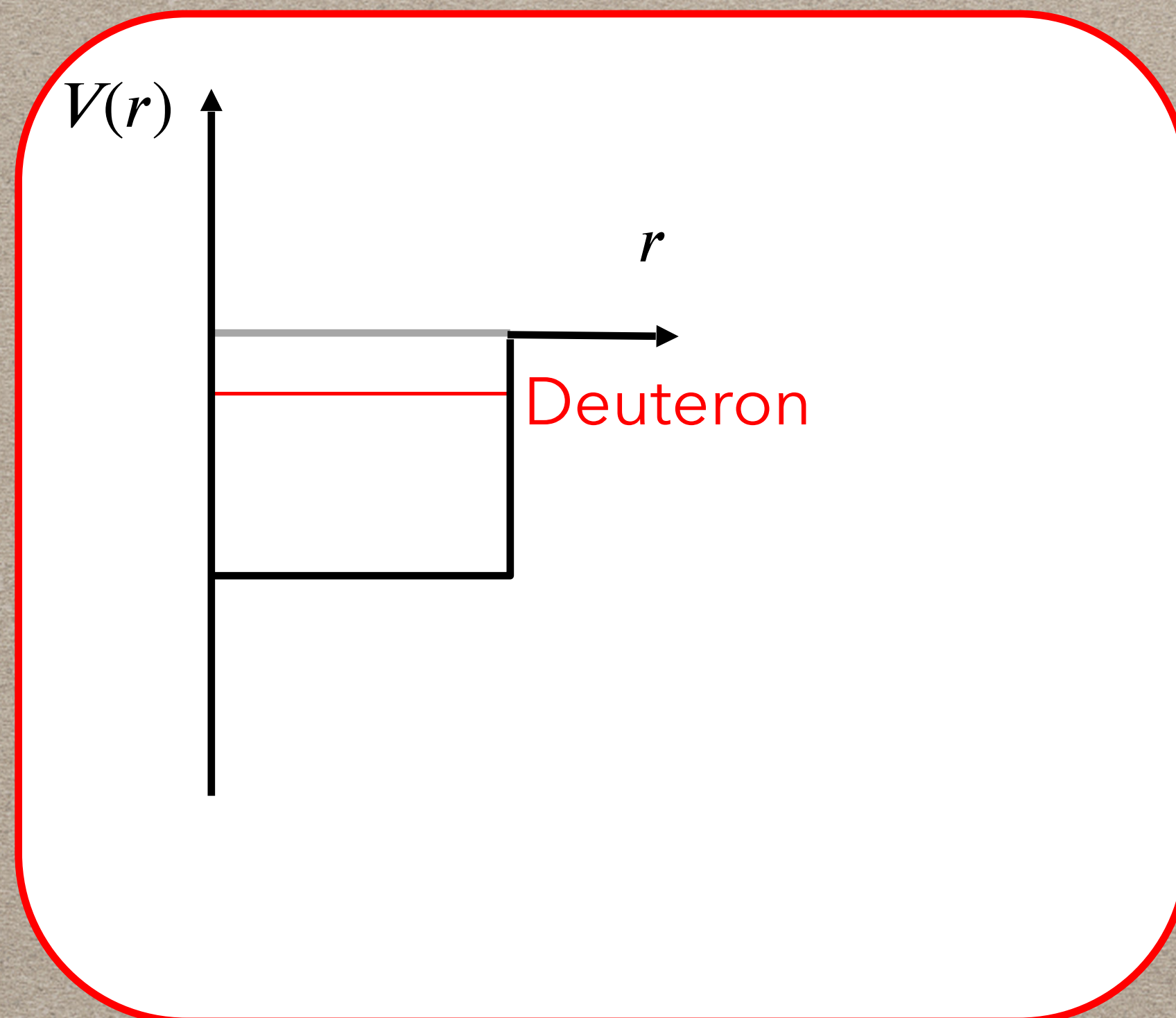
$nn$  scattering never generates a bound state, but generates a signal that can be seen in the scattering lengths

The interaction is strong enough to generate pole but not to bind the system

# BOUND AND VIRTUAL STATES

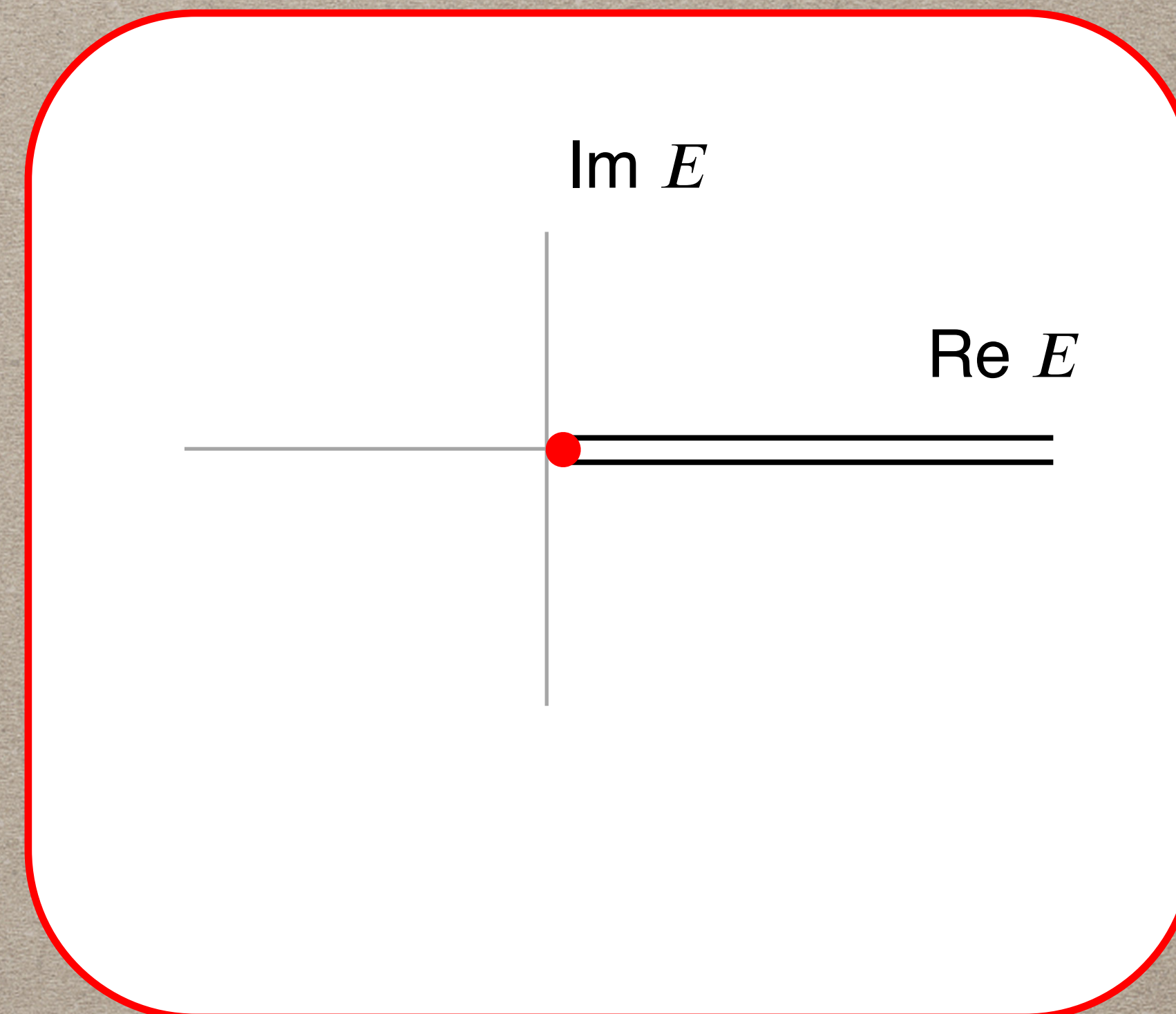
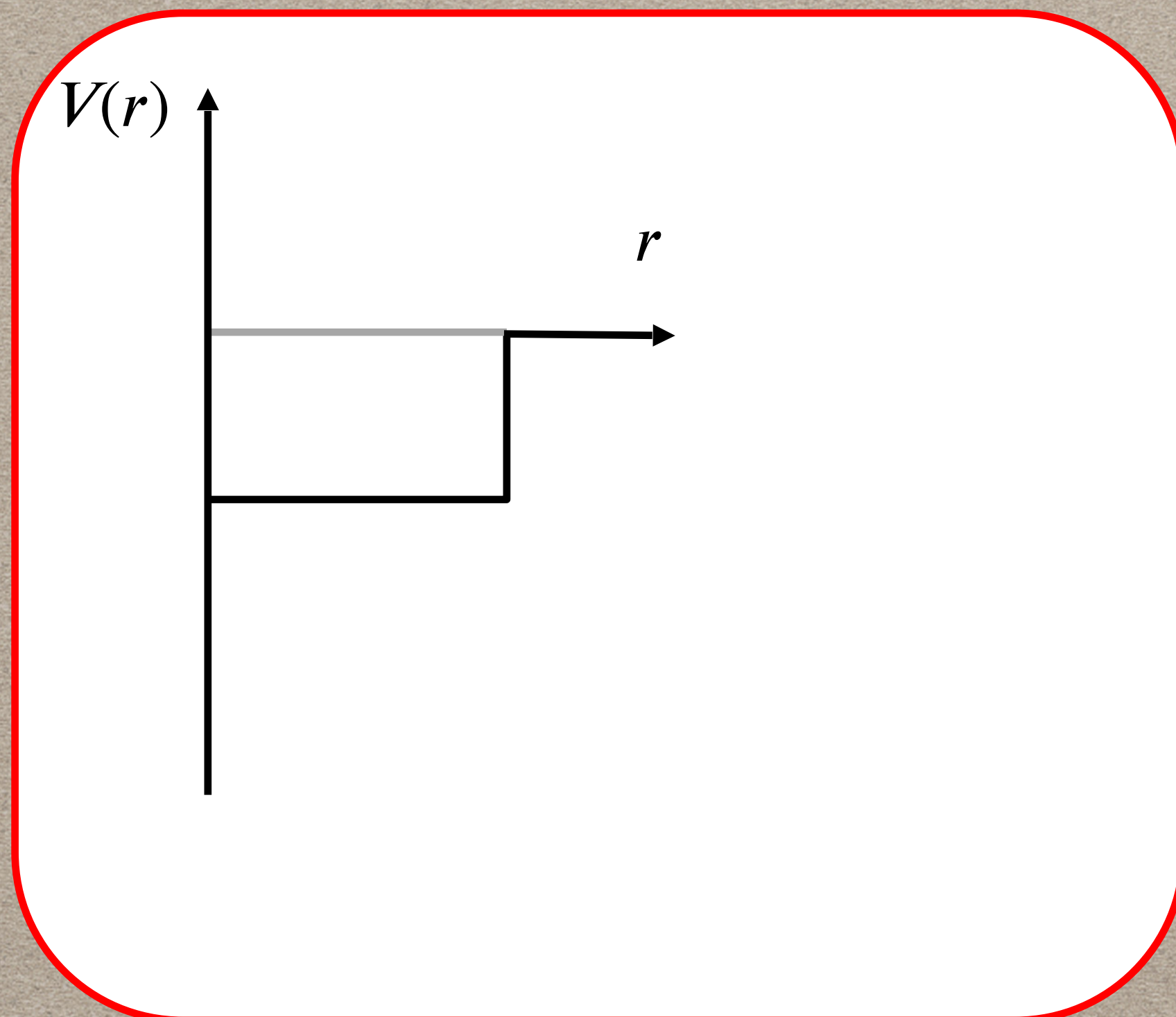
Example from  $pn$  scattering

Bound state on the real axis I sheet (deuteron)



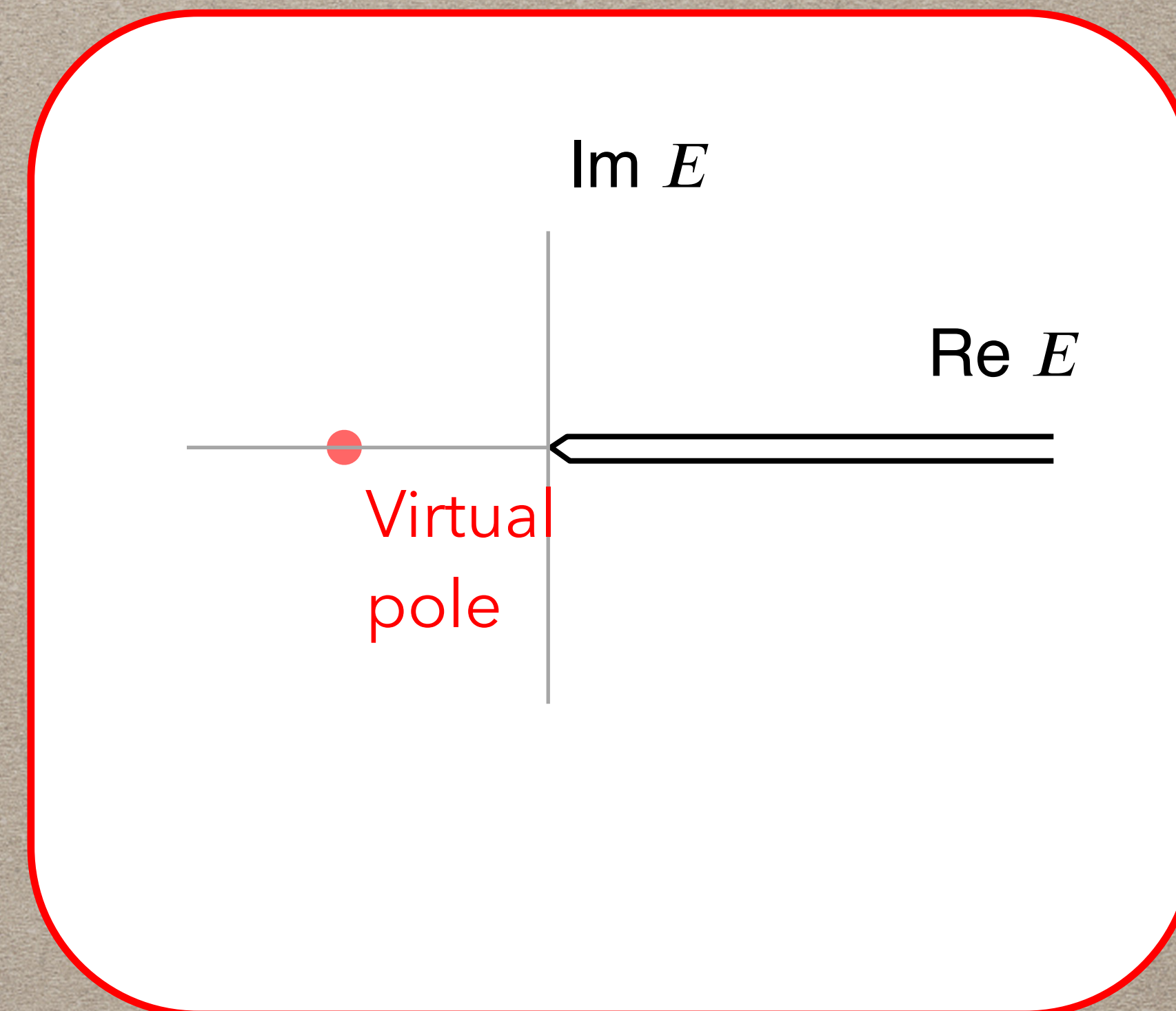
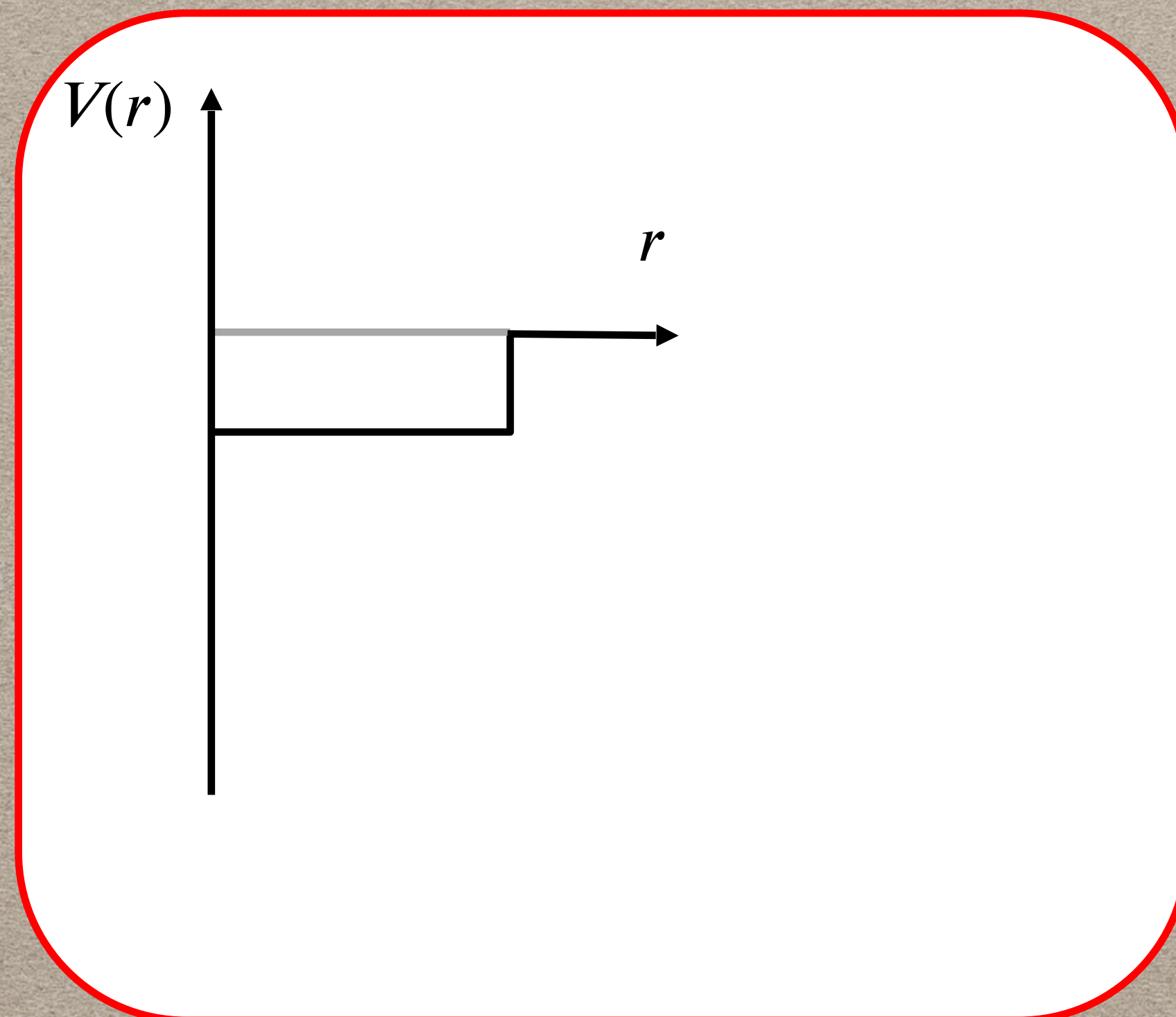
# BOUND AND VIRTUAL STATES

Decreasing the potential strength,  
the pole reaches threshold

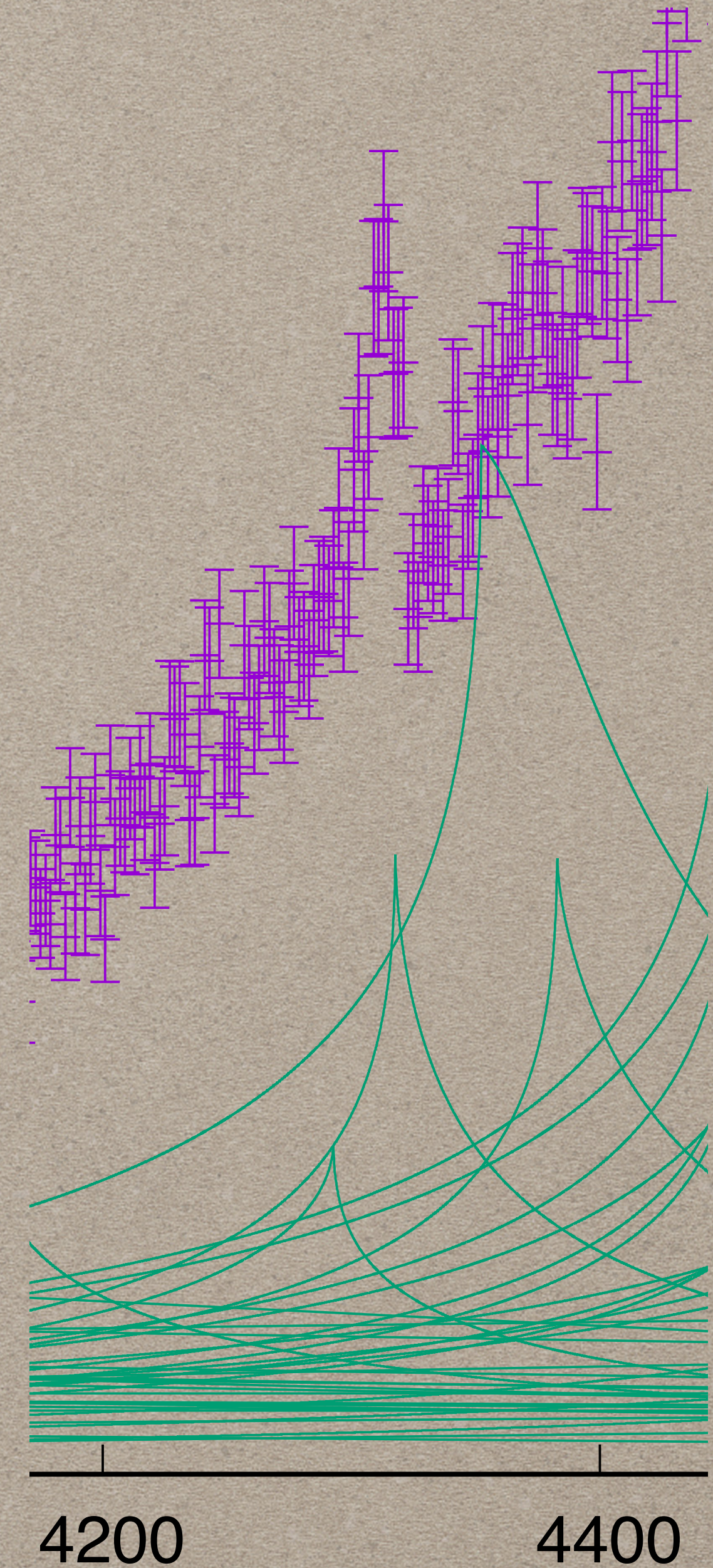
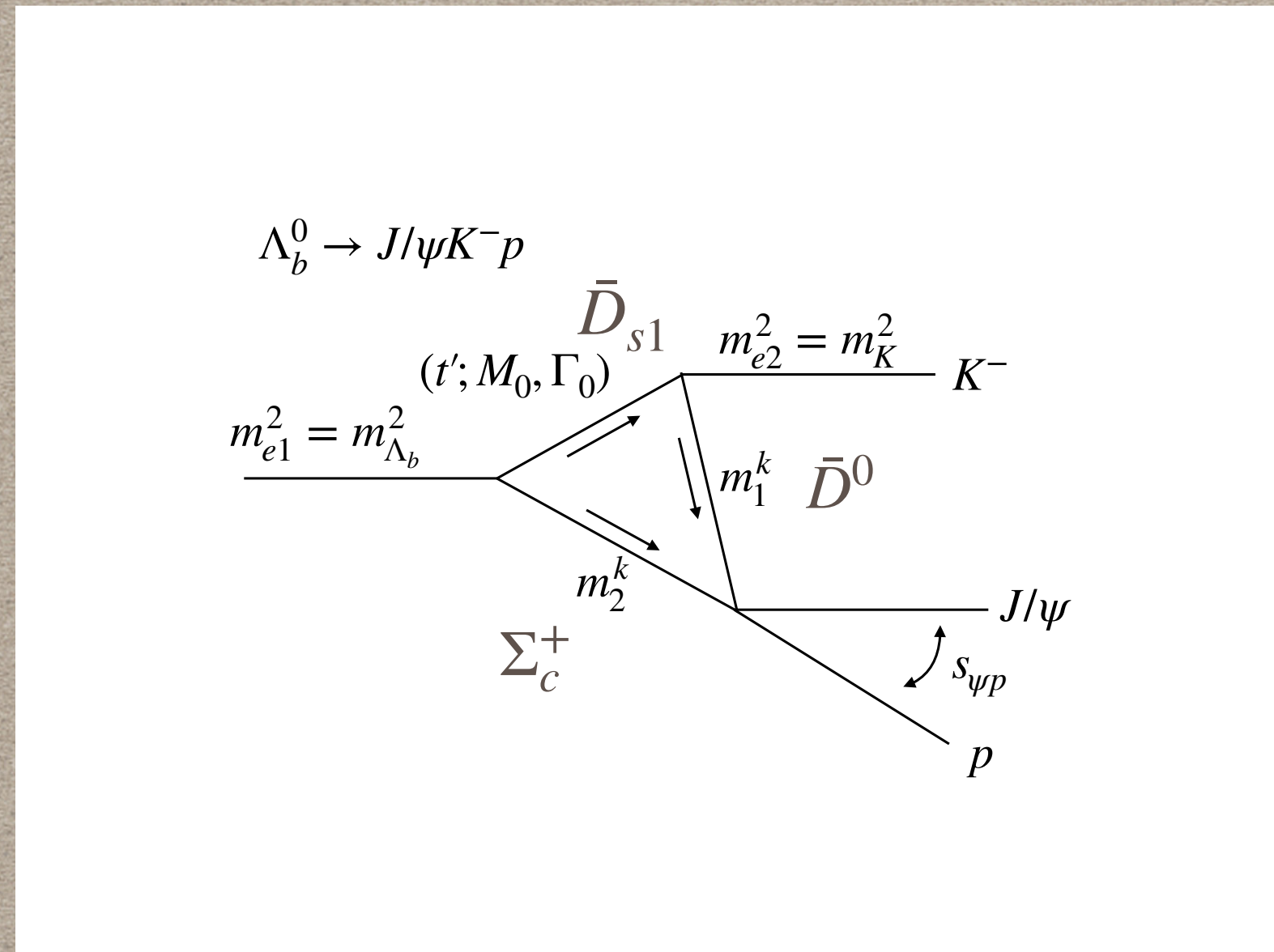


# BOUND AND VIRTUAL STATES

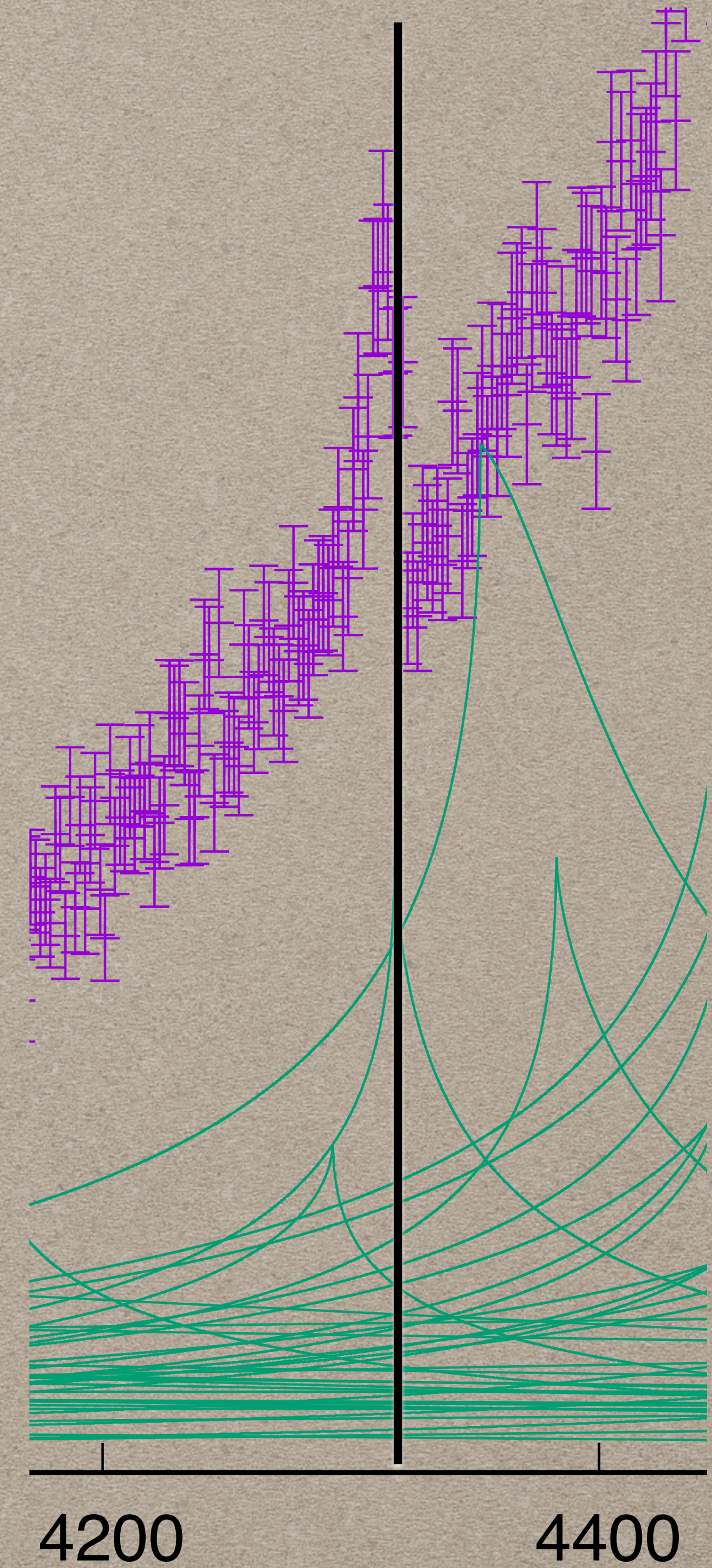
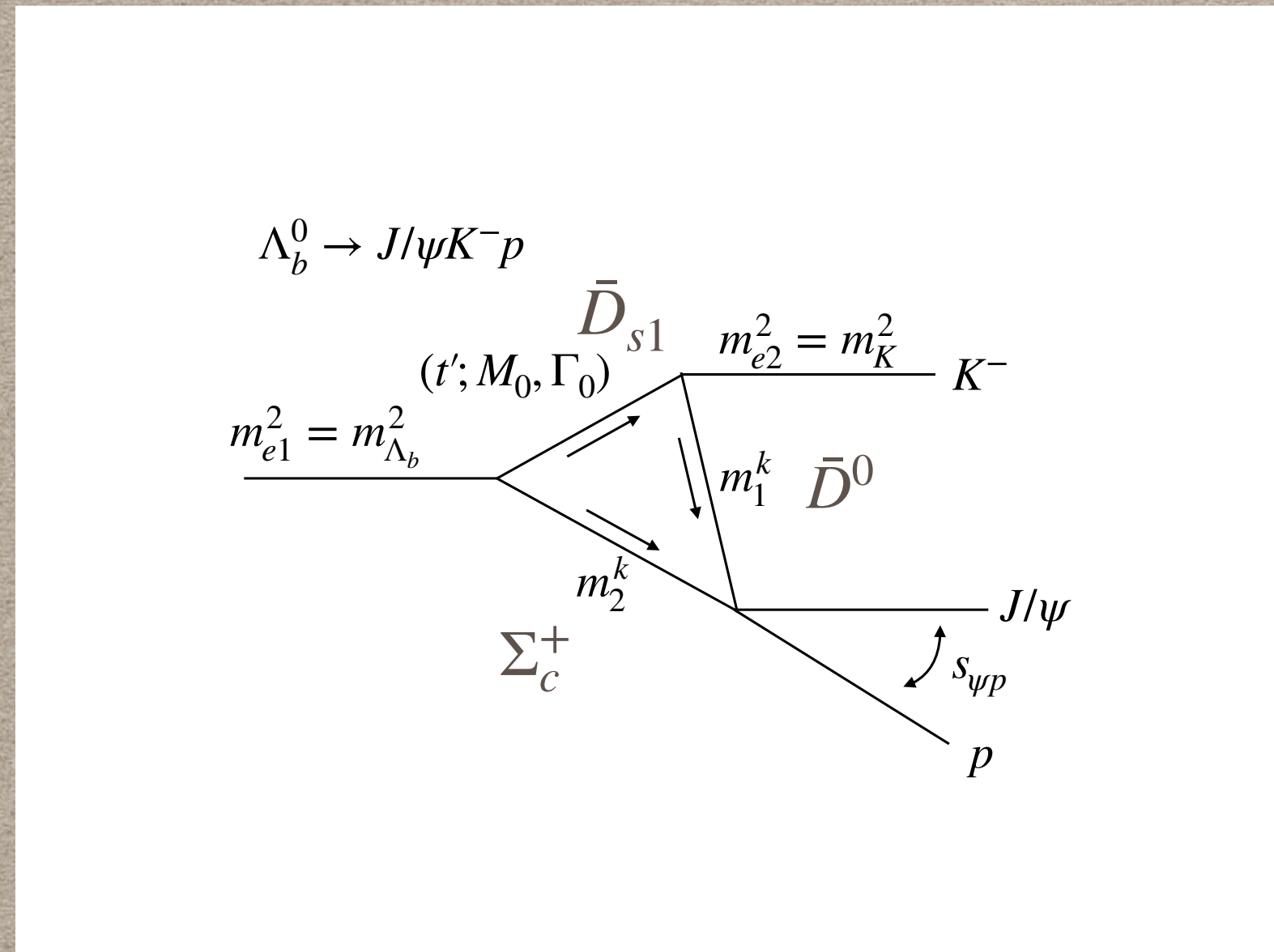
The pole jumps on the II sheet,  
it becomes a virtual state



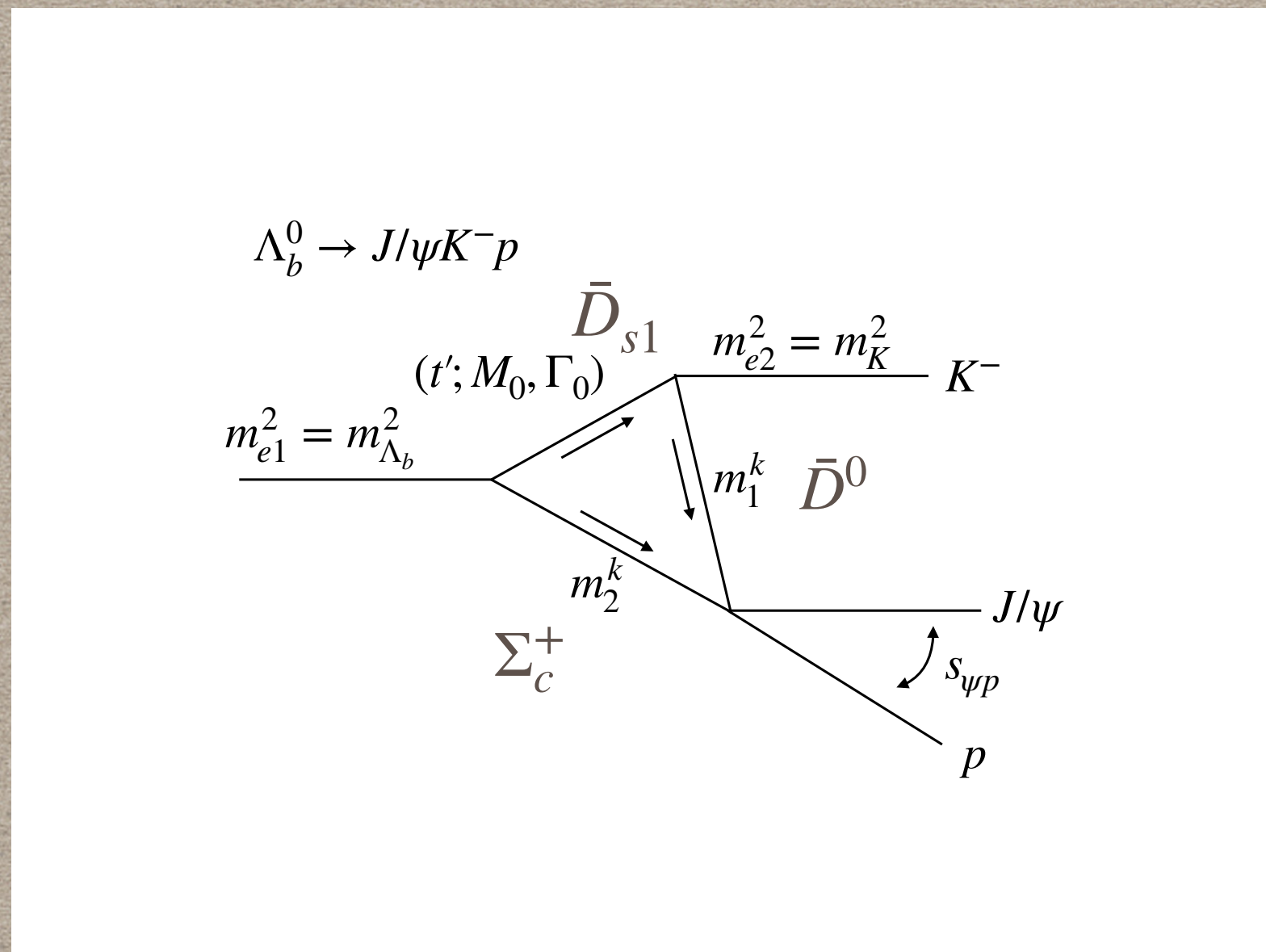
# TRIANGLE SINGULARITY



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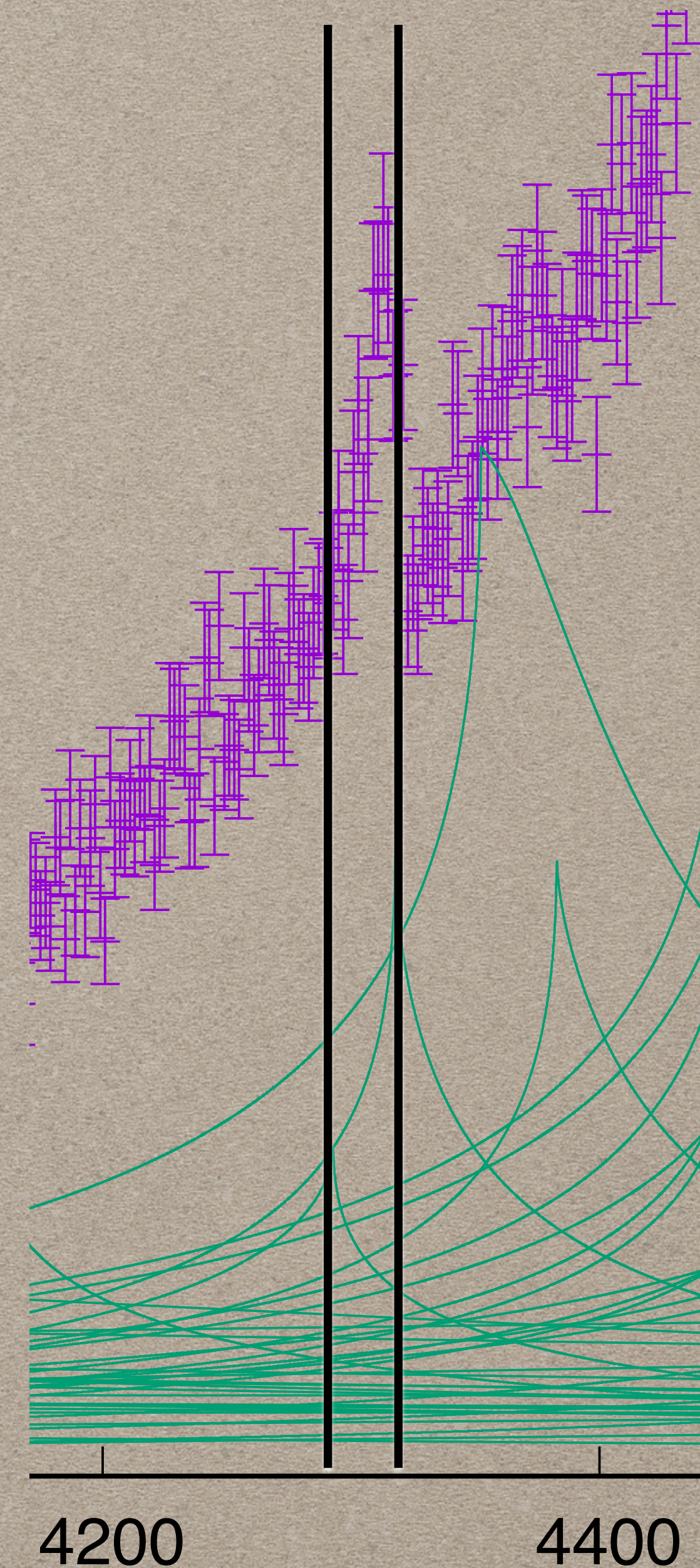


# TRIANGLE SINGULARITY



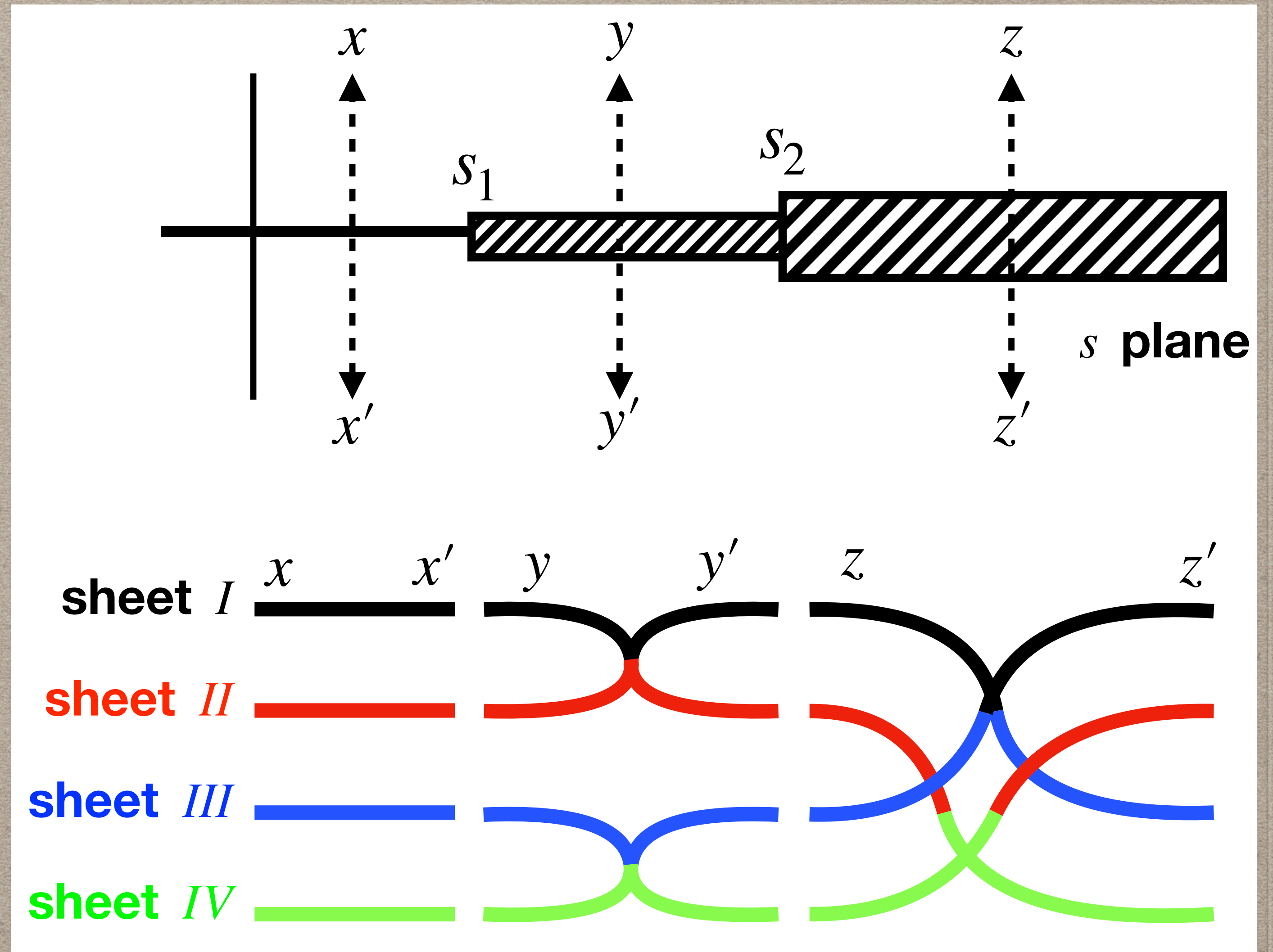
Triangle singularities do not generate poles but the phase motion is the same as for a pole (i.e. the Argand plot is going to be the same)

[Remember Bernhard's talk]



# RIEMANN SHEETS STRUCTURE

First threshold is  $J/\psi$   $p$  channel  
and the second is  $\Sigma_c^+ \bar{D}^0$



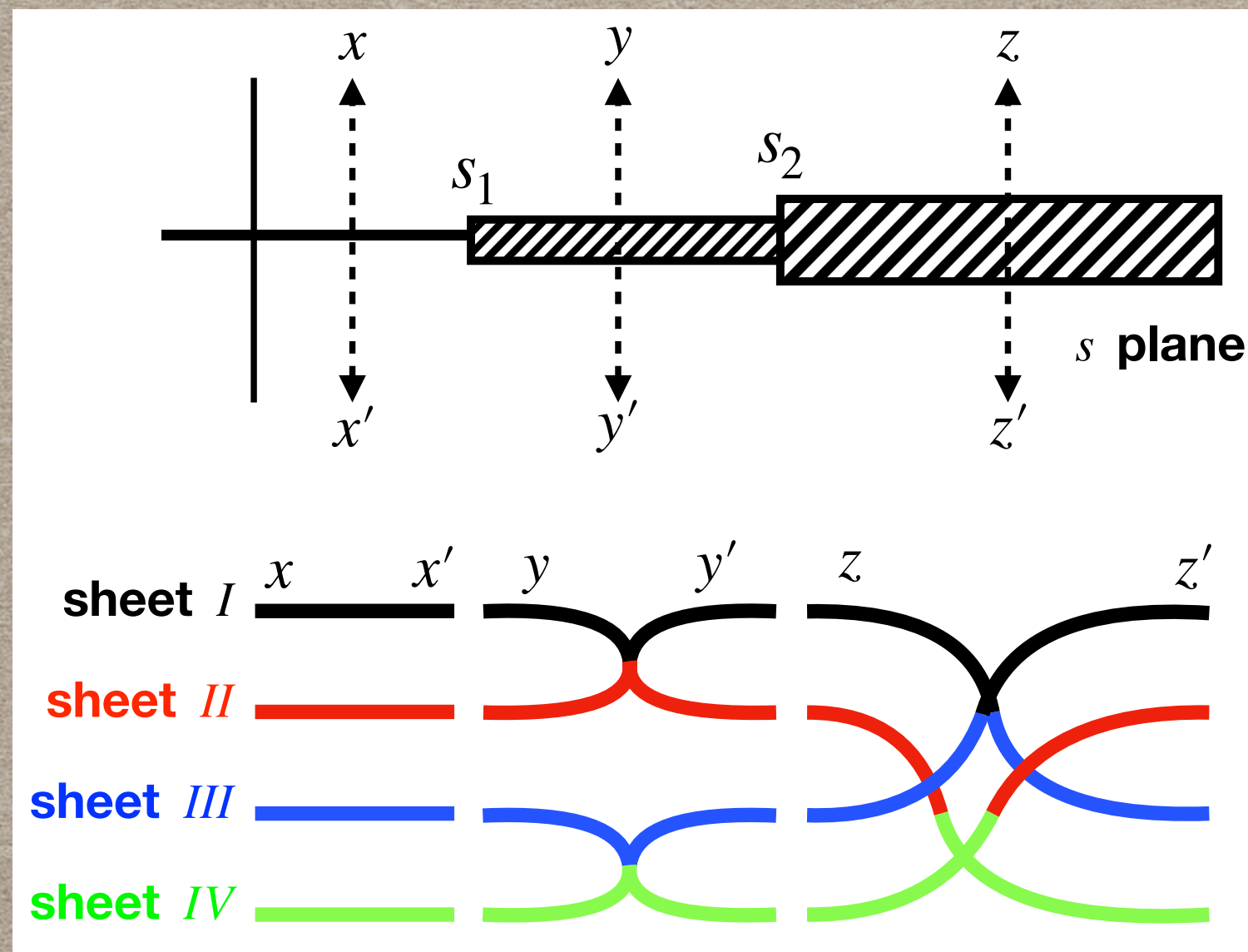


# COMPACT PENTAQUARK

$$\Sigma_c^+ \bar{D}^0$$



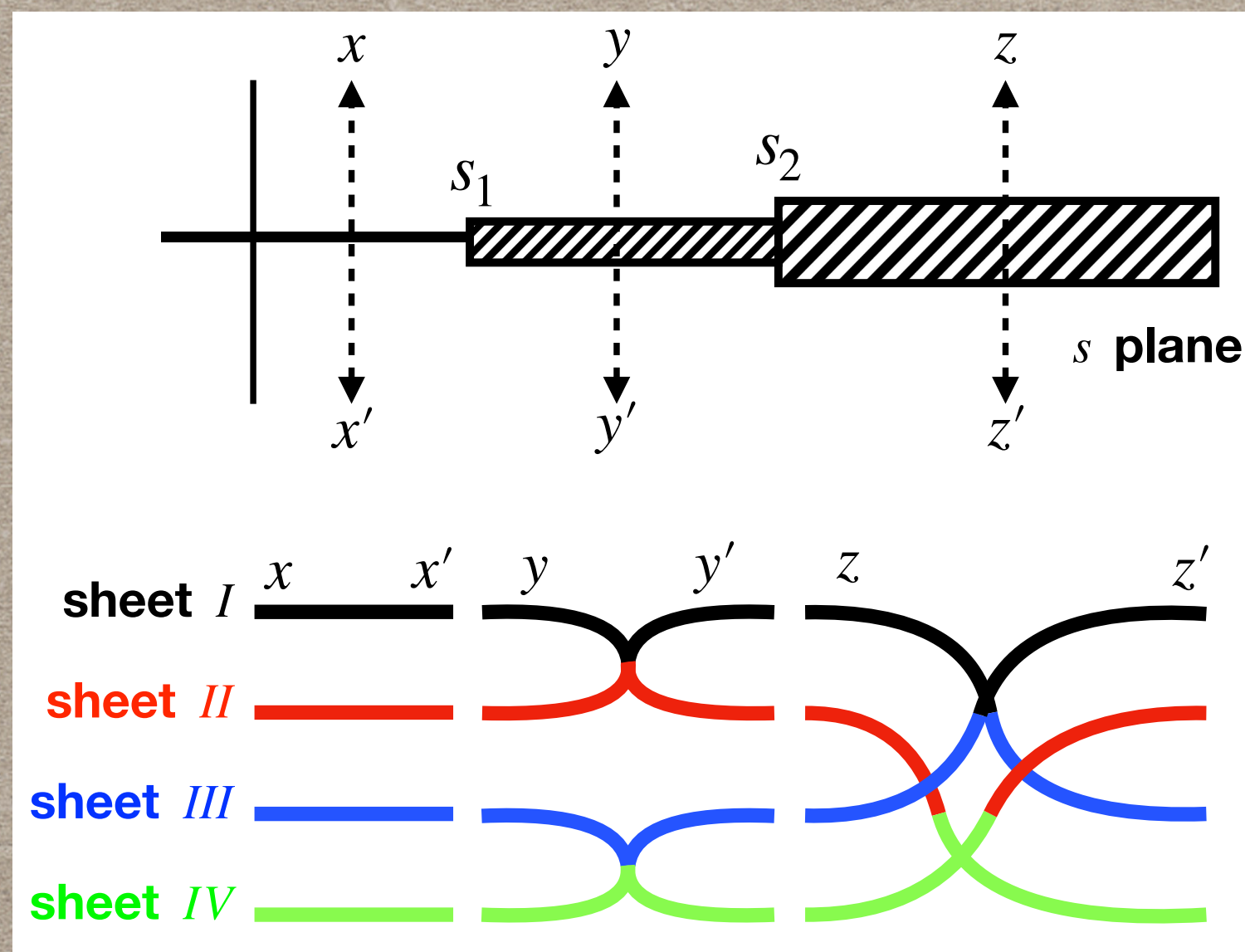
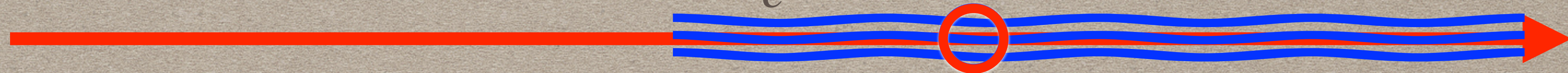
*J/ψ p*



# COMPACT PENTAQUARK

$J/\psi p$

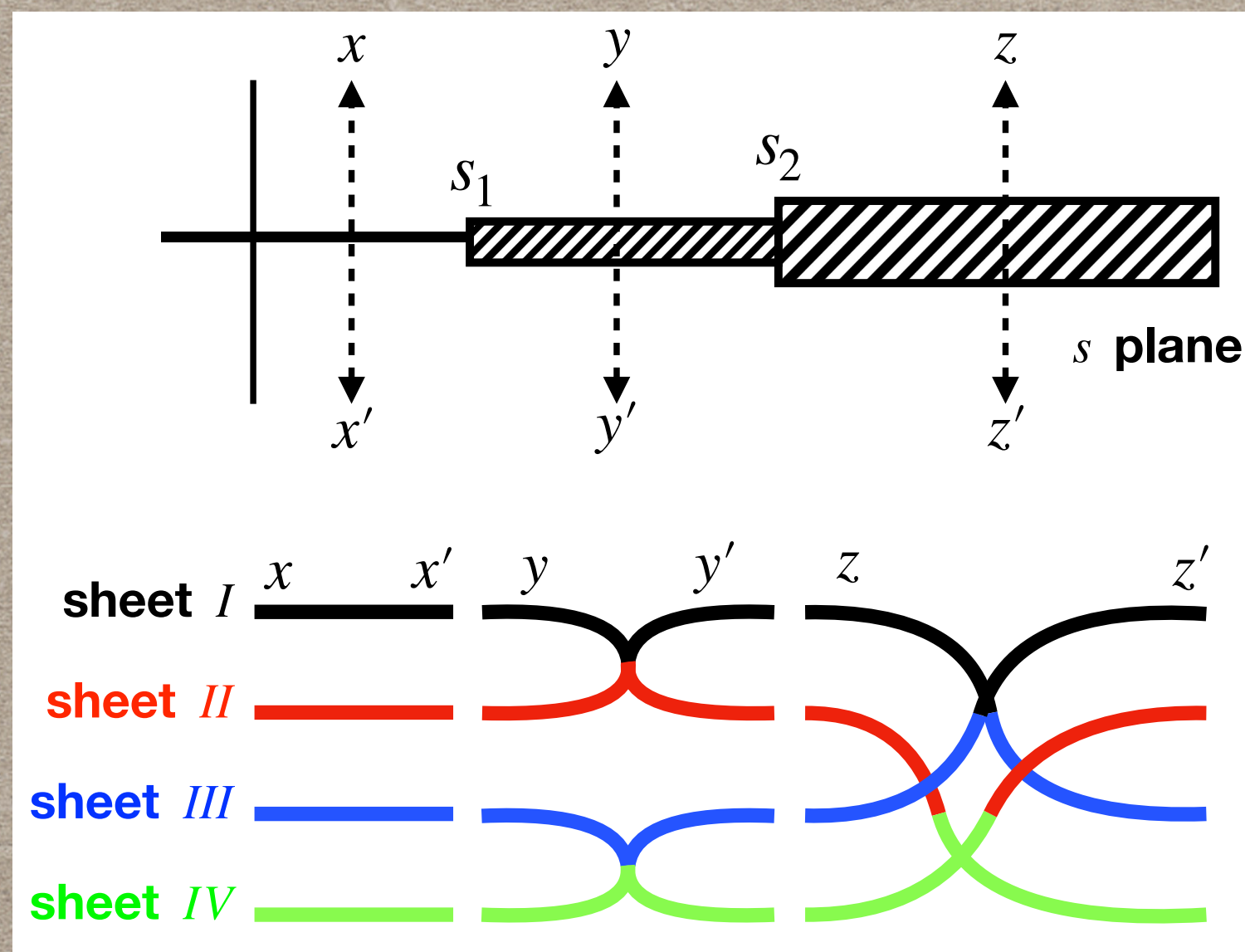
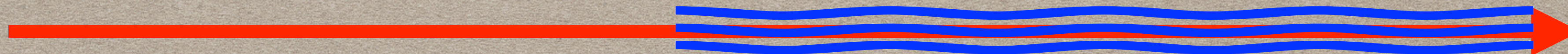
$\Sigma_c^+ \bar{D}^0$





# COMPACT PENTAQUARK

$J/\psi p$

$\Sigma_c^+ \bar{D}^0$

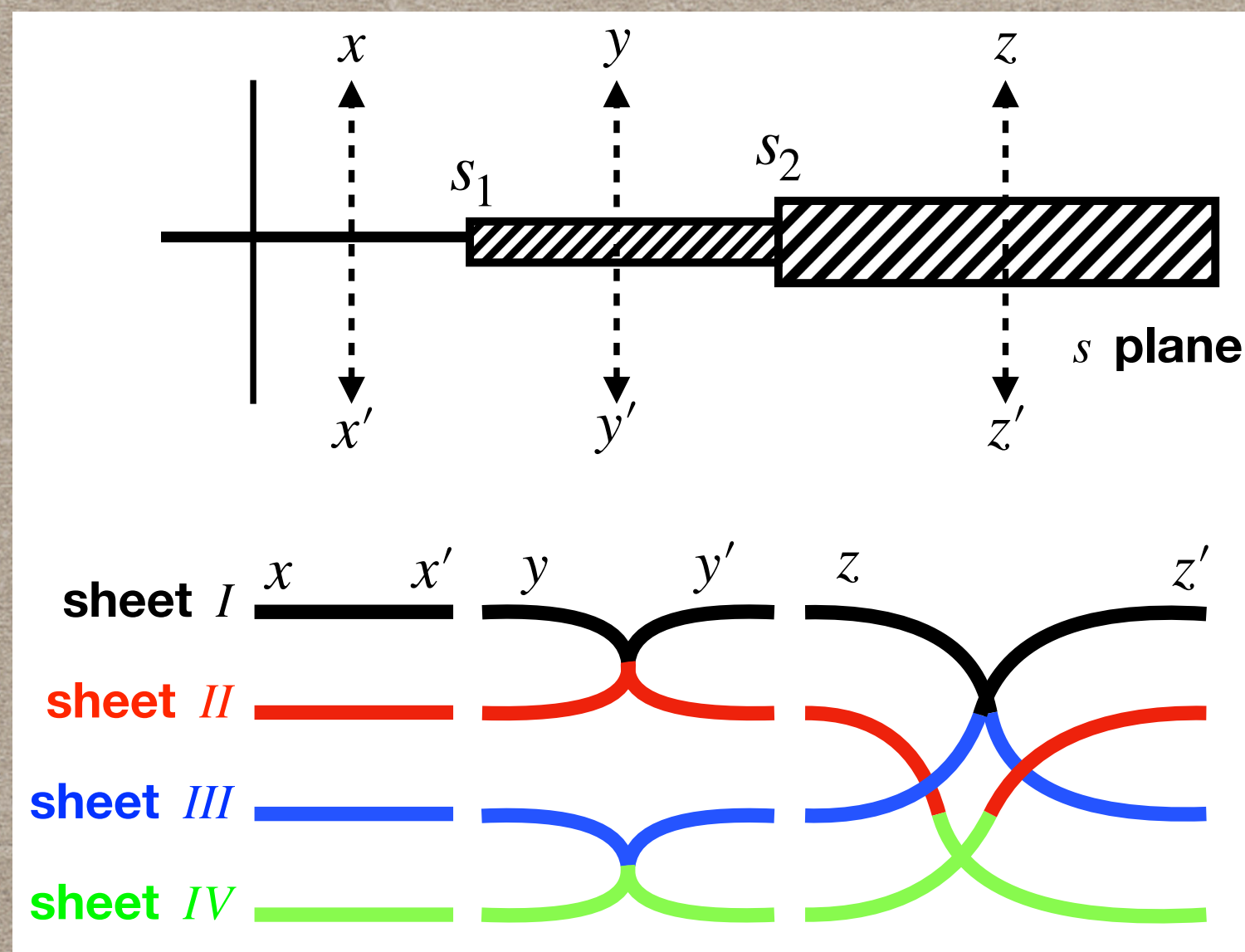
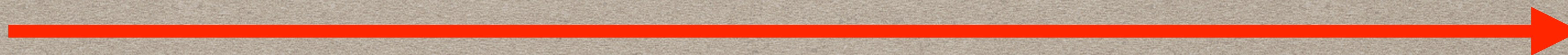


II sheet  III sheet 

# MOLECULE



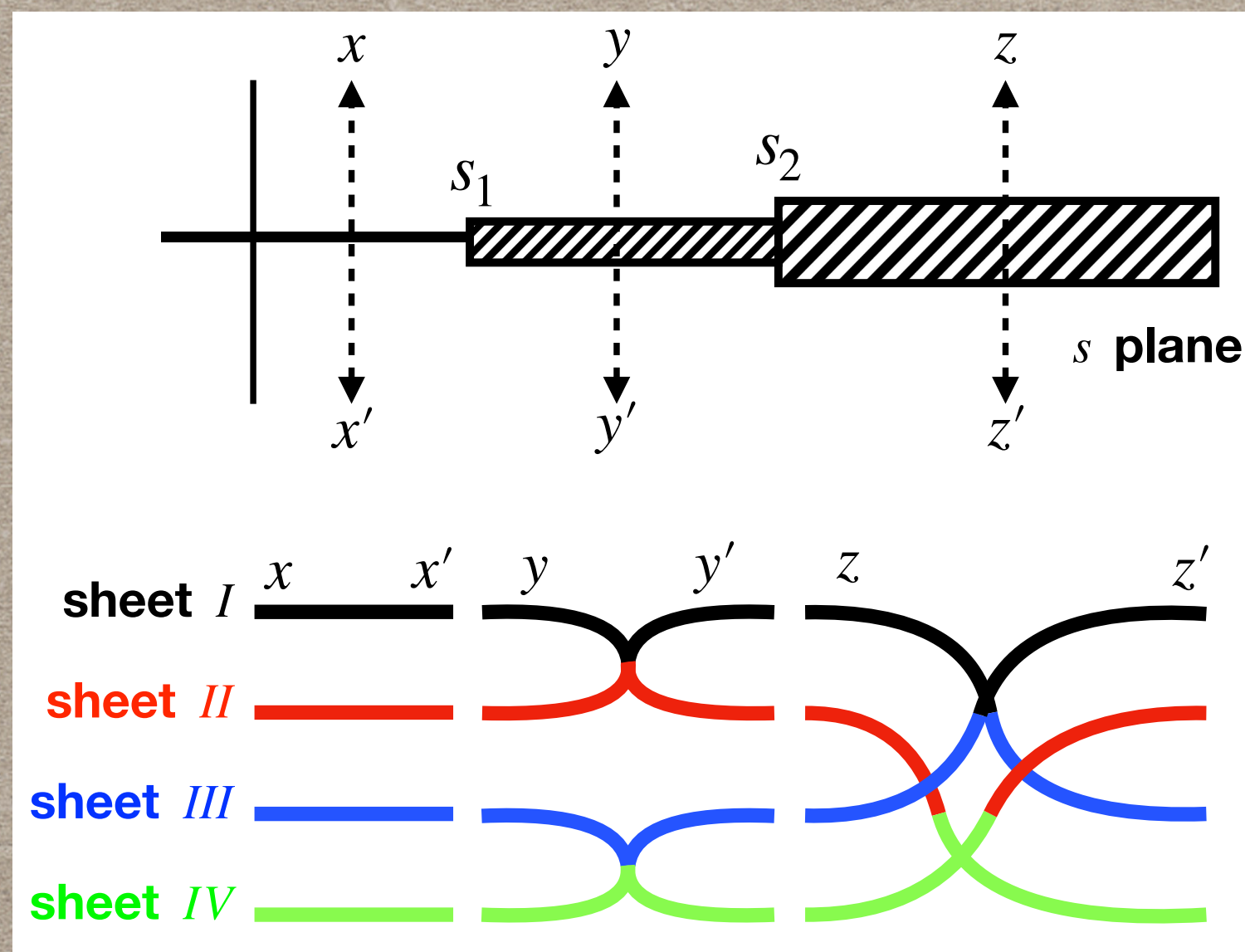
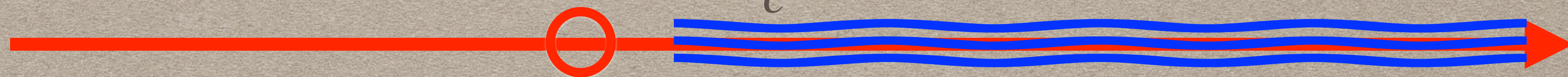
$J/\psi p$



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$J/\psi p$

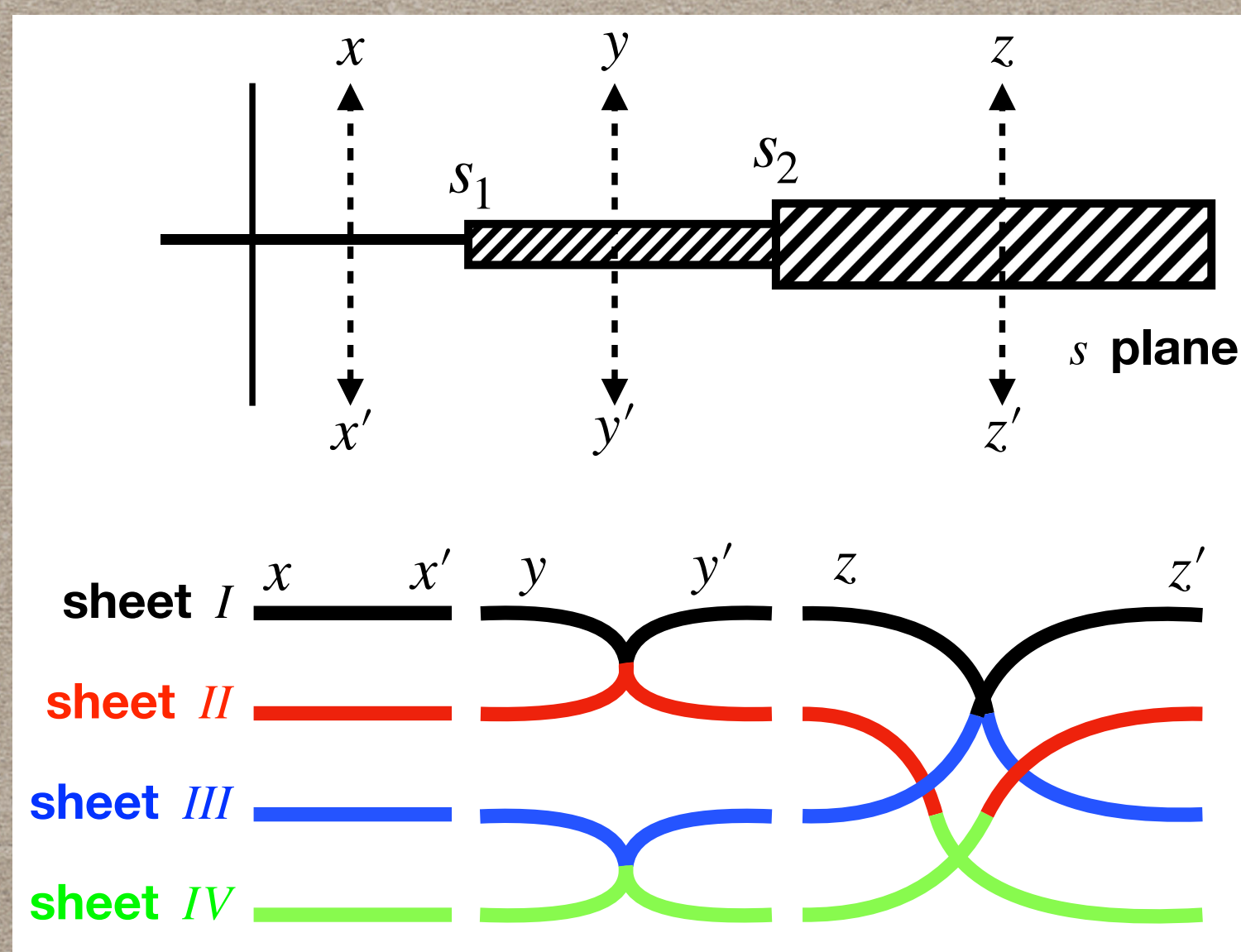
$\Sigma_c^+ \bar{D}^0$



# MOLECULE

$J/\psi p$

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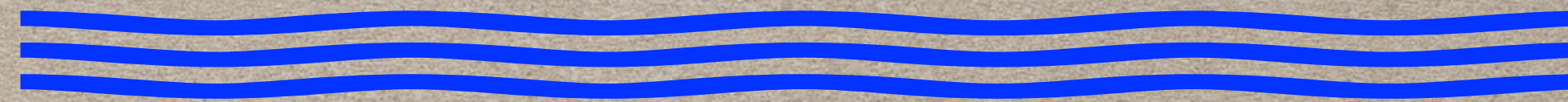
II sheet

Either nothing on the III sheet  
or shadow pole

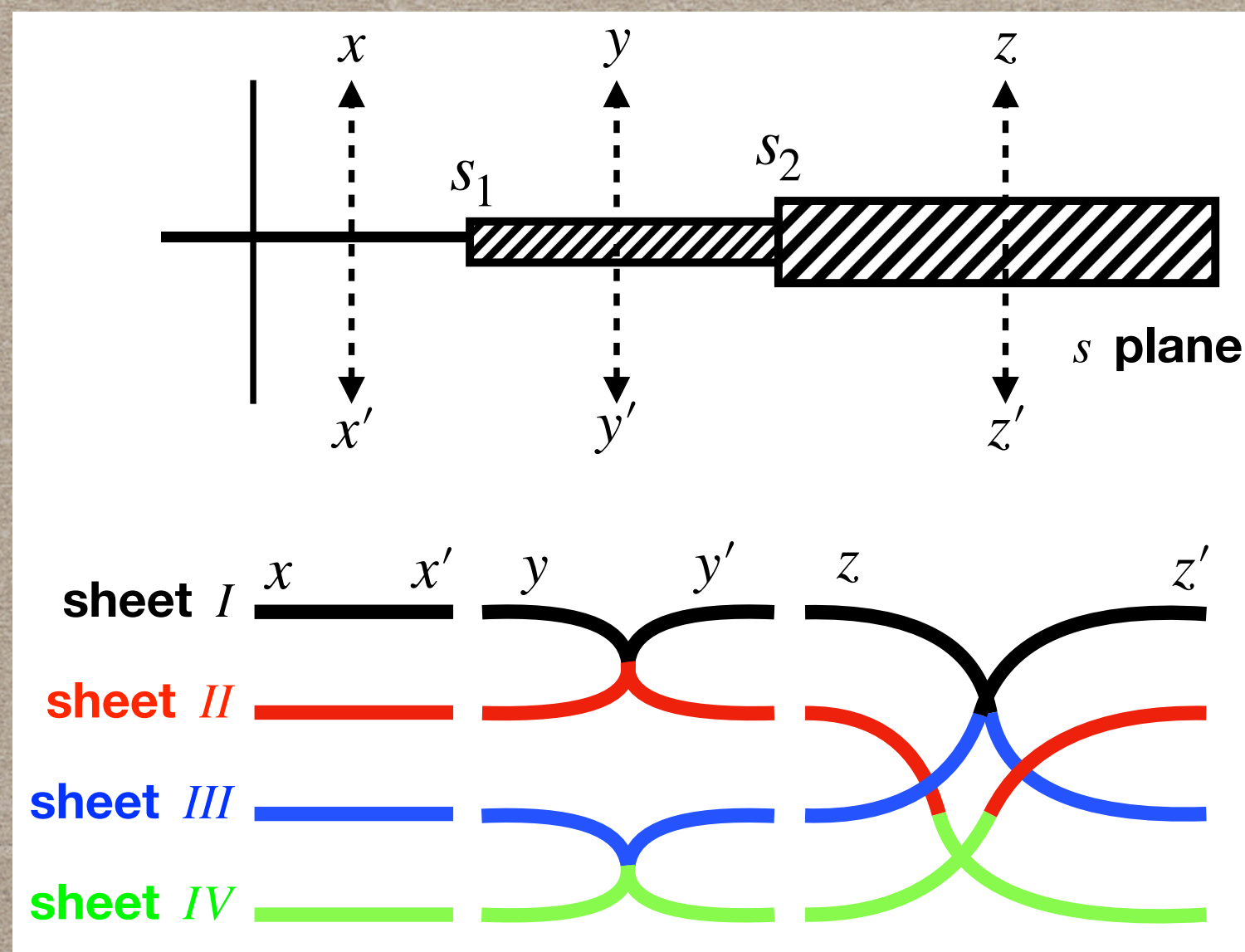
# VIRTUAL STATE

II sheet

$$\textcircled{\text{O}} \Sigma_c^+ \bar{D}^0$$



$J/\psi p$

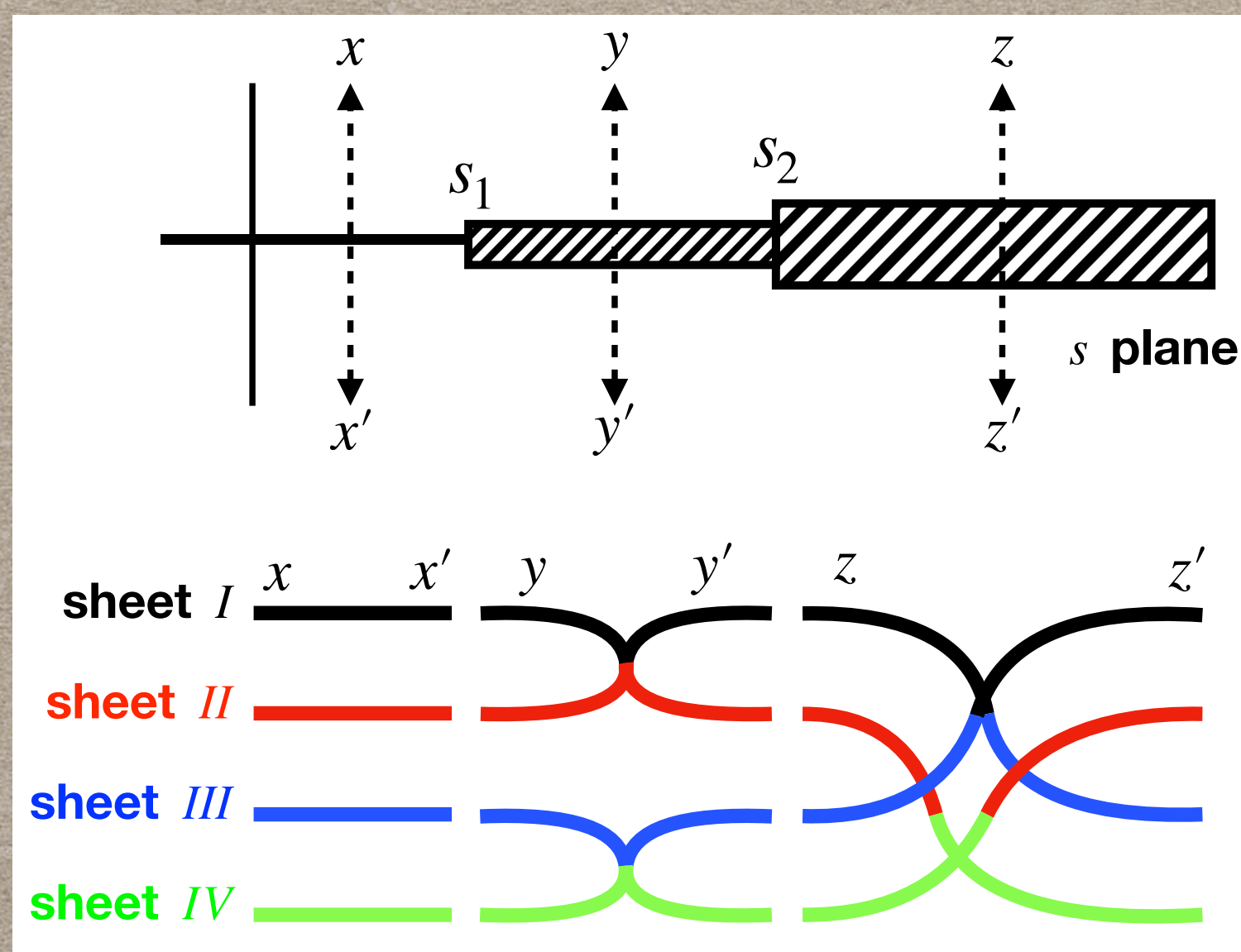


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$J/\psi p$

IV sheet

$\Sigma_c^+ \bar{D}^0$





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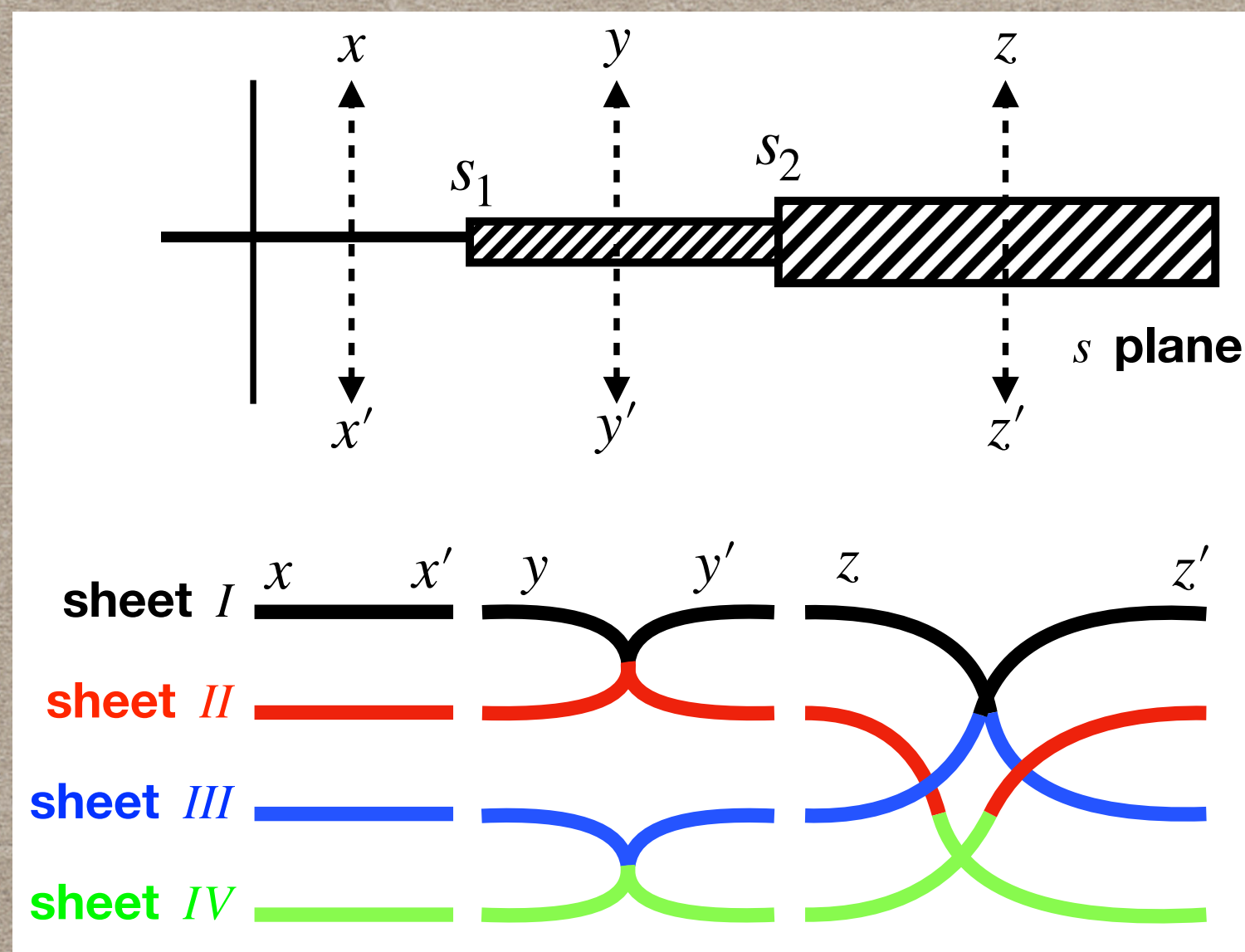
IV sheet

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Either nothing on the  
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# ANALYSIS OF THE $P_c(4312)$ SIGNAL

- Build a theory in the near threshold region
- Analyze the three datasets provided by LHCb  $P_c(4312)$
- 66 experimental data
- Experimental resolution incorporated
- Error analysis through bootstrap

**JPAC, PRL 123 (2019) 092001**

# NEAR-THRESHOLD THEORY

Hypotheses:

- 📌 Only one partial wave contributes to the signal
- 📌 The threshold drives the physics (testable)
- 📌 Other effects are absorbed in the parameters (testable)

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- 📌 Only one partial wave contributes to the signal
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Caveat:

- 📌 We fit the  $J/\psi$  p projection (no info on quantum numbers)

# NEAR-THRESHOLD THEORY

$$\frac{dN}{d\sqrt{s}} = \rho(s) \left[ |F(s)|^2 + B(s) \right]$$

$$(T^{-1})_{ij} = M_{ij} - ik_i \delta_{ij}$$

$$M_{ij}(s) = m_{ij} - c_{ij}s$$

The matrix elements  $M_{ij}$  are singularity free near threshold and can be expanded in a Taylor series

$$F(s) = P_1(s)T_{11}(s)$$

$$B(s) = b_0 + b_1s$$

$$P_1(s) = p_0 + p_1s$$

**Frazer & Hendry, PR 134 (1964) B1307**

# AMPLITUDE IN THE NEAR THRESHOLD REGION

$$\frac{dN}{d\sqrt{s}} = \rho(s) \left[ |F(s)|^2 + B(s) \right]$$

$$B(s) = b_0 + b_1 s$$

$$F(s) = (p_0 + p_1 s) \frac{[m_{22} - c_{22}s - ik_2]}{[m_{22} - c_{22}s - ik_2][m_{11} - c_{11}s - ik_1] - m_{12}^2}$$

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$$\frac{dN}{d\sqrt{s}} = \rho(s) \left[ |F(s)|^2 + B(s) \right]$$

Production, hyperons and effects due to other (far) singularities

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Channel coupling

Effective range approximation if  $c_{ij}=0$

Under the effective range approximation only poles in the II and IV sheet can happen

When  $c_{ij} \neq 0$ , poles can appear in any sheet (no threshold domination hypothesis)

# FIT RESULTS

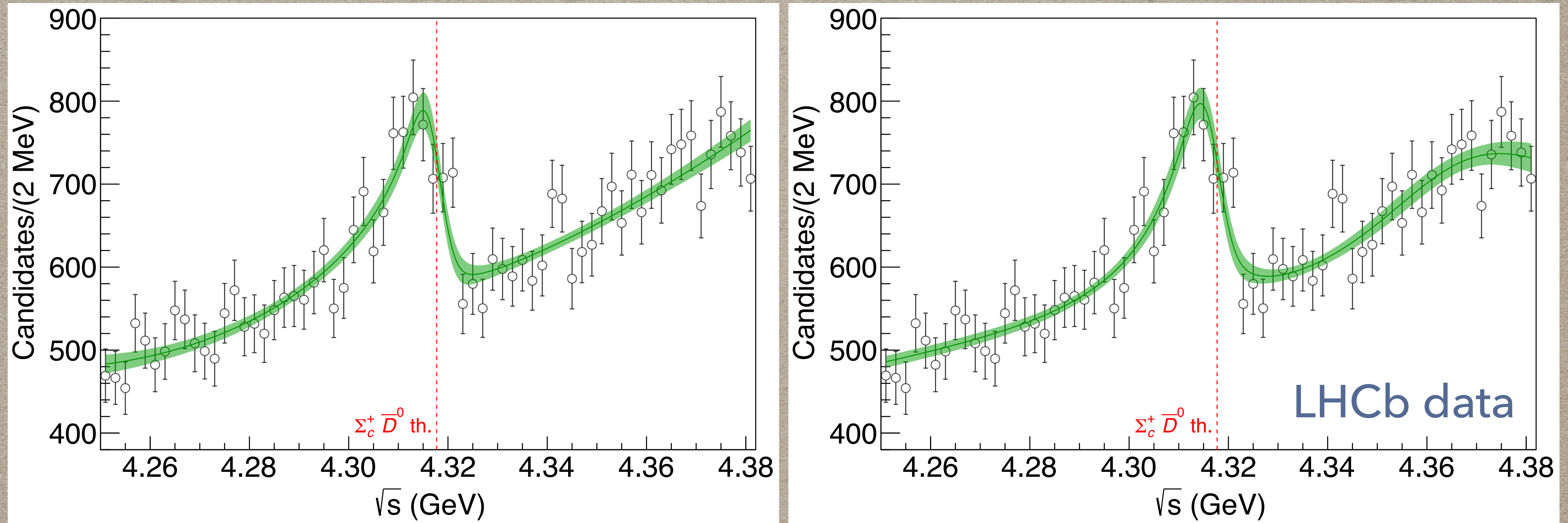
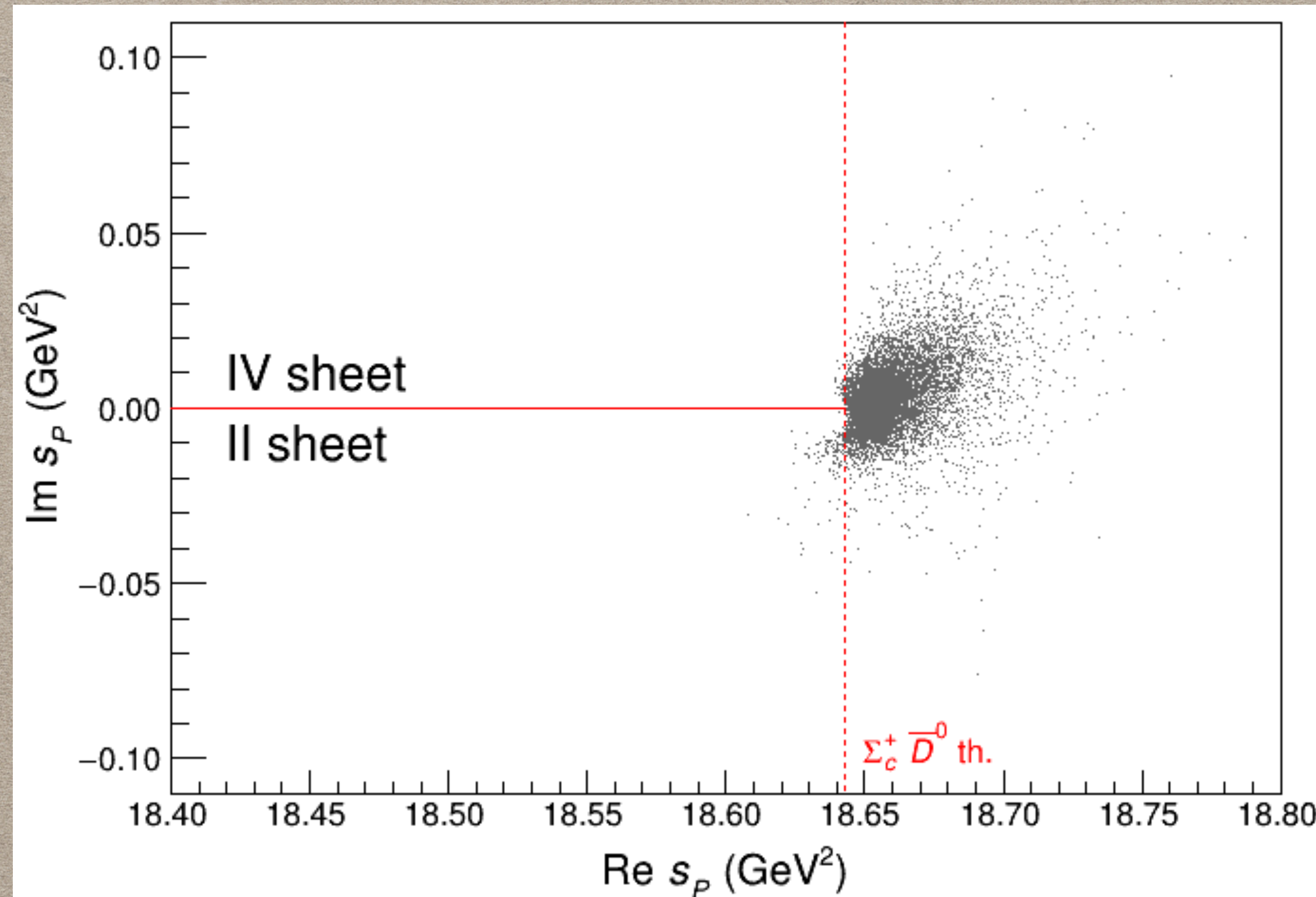
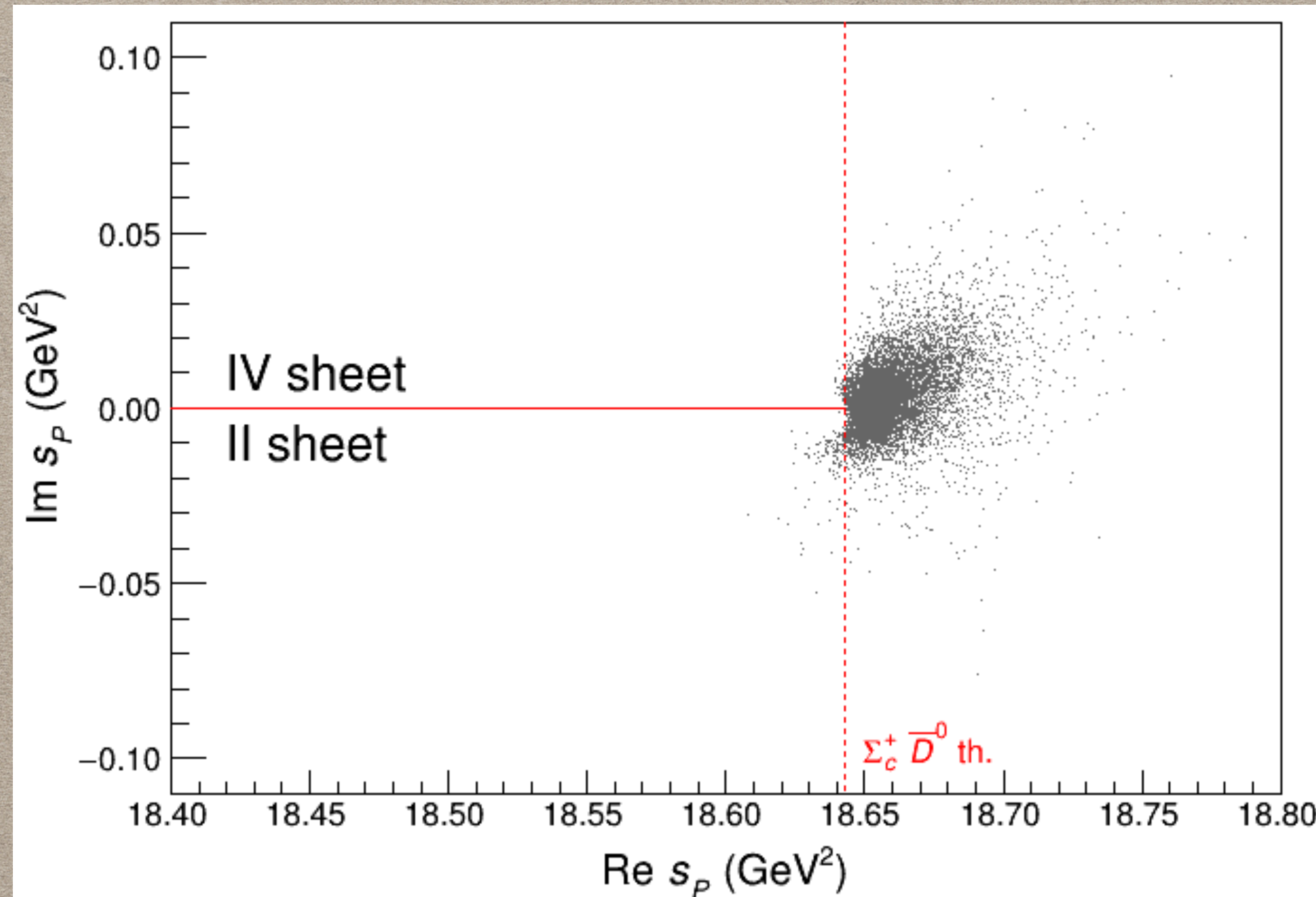


FIG. 1. Fits to the  $\cos \theta_{P_c}$ -weighted  $J/\psi p$  mass distribution from LHCb [9] according to cases A (left) and B (right). The amplitude of case A is expressed in the scattering length approximation, *i.e.*  $c_{ij} = 0$  in Eq. (3), and is able to describe either bound (molecular) or virtual states. The amplitude of case B is given in the effective range approximation, *i.e.* finite  $c_{ii}$ , and extends the description to genuine pentaquark states. The solid line and green band show the result of the fit and the  $1\sigma$  confidence level provided by the bootstrap analysis, respectively.

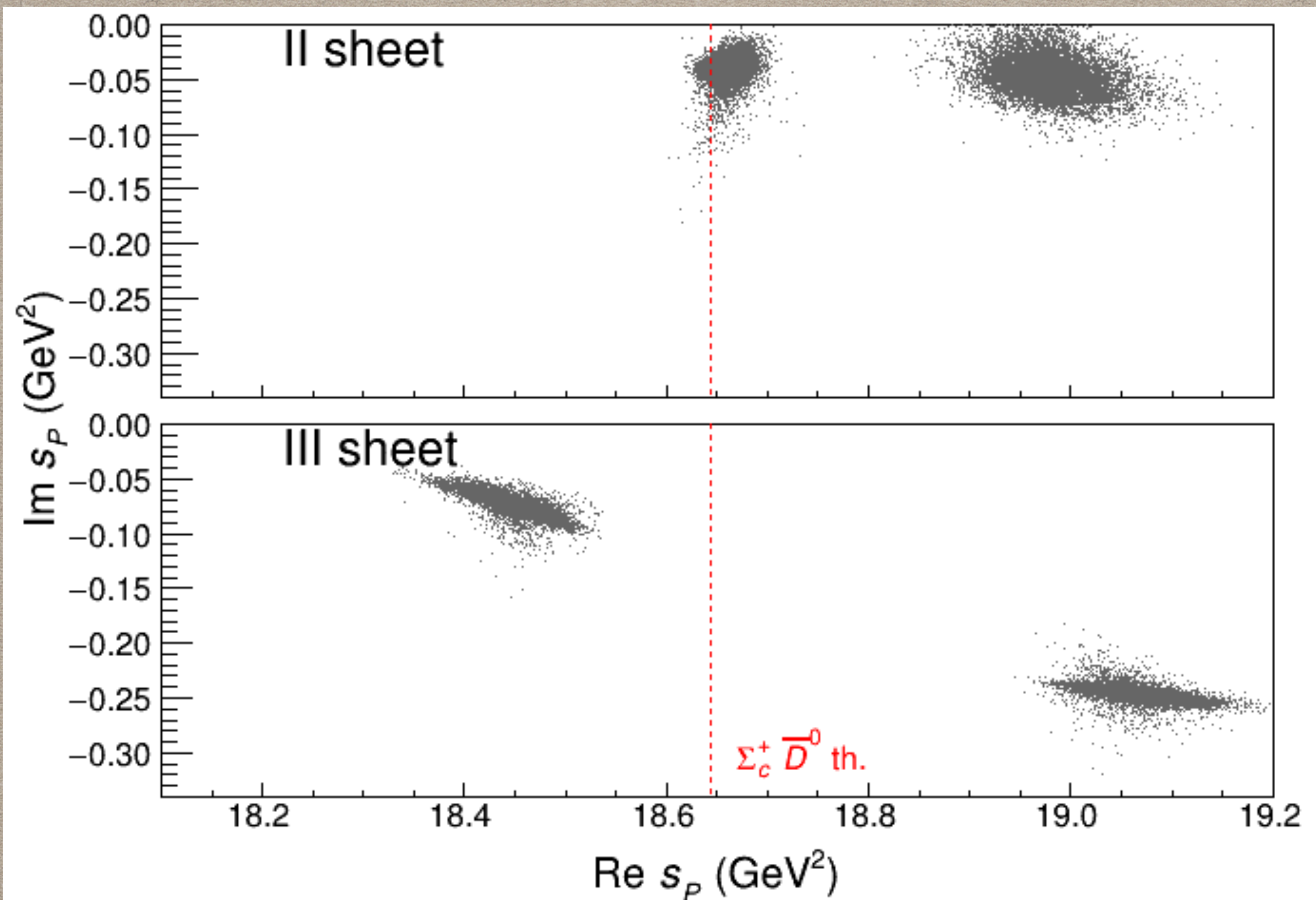
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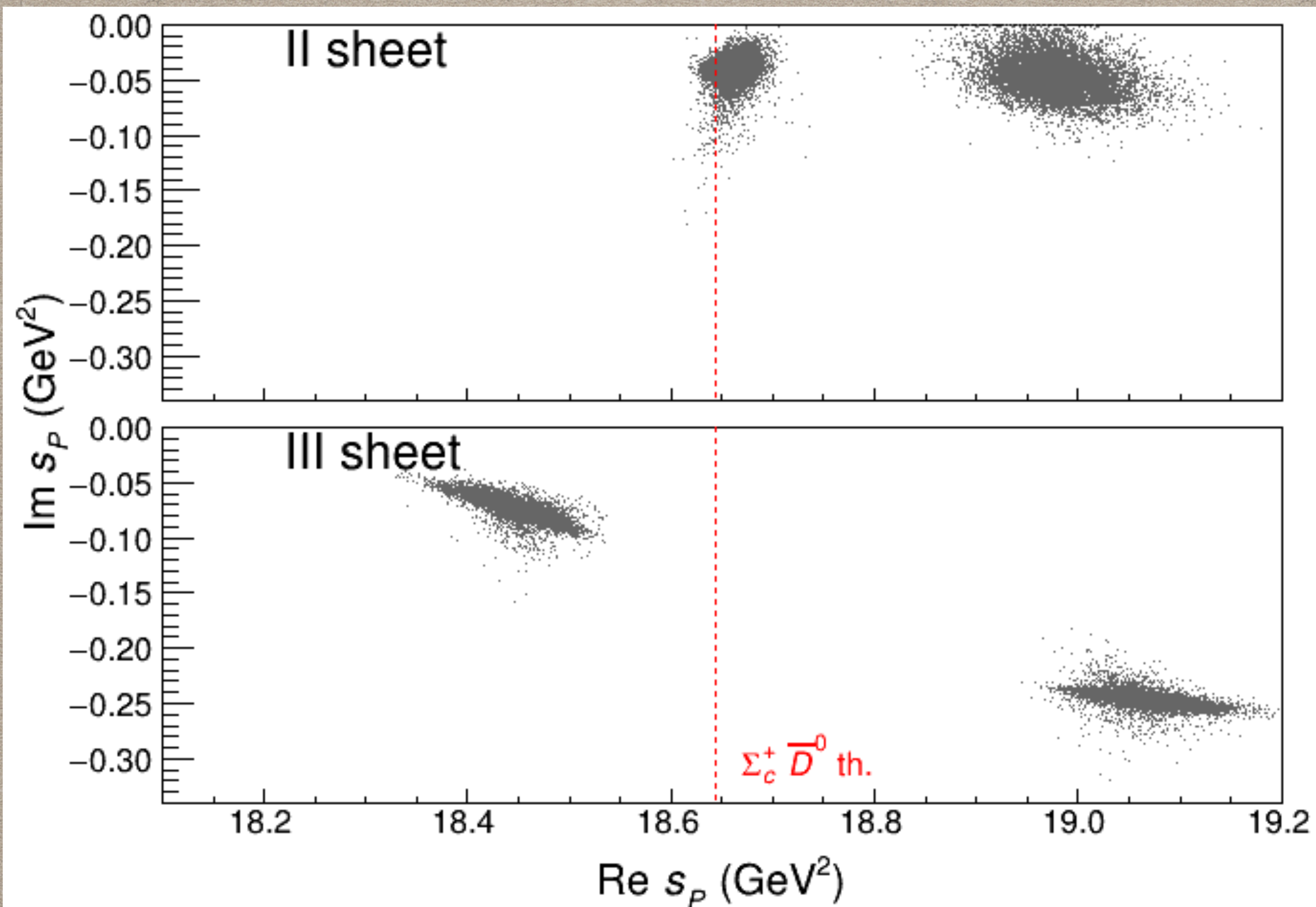
# POLE MOVEMENT: $C \neq 0$



When  $m_{12}=0$  both channels decouple and

$$\frac{[m_{22} - c_{22}s - ik_2]}{[m_{22} - c_{22}s - ik_2][m_{11} - c_{11}s - ik_1] - m_{12}^2}$$

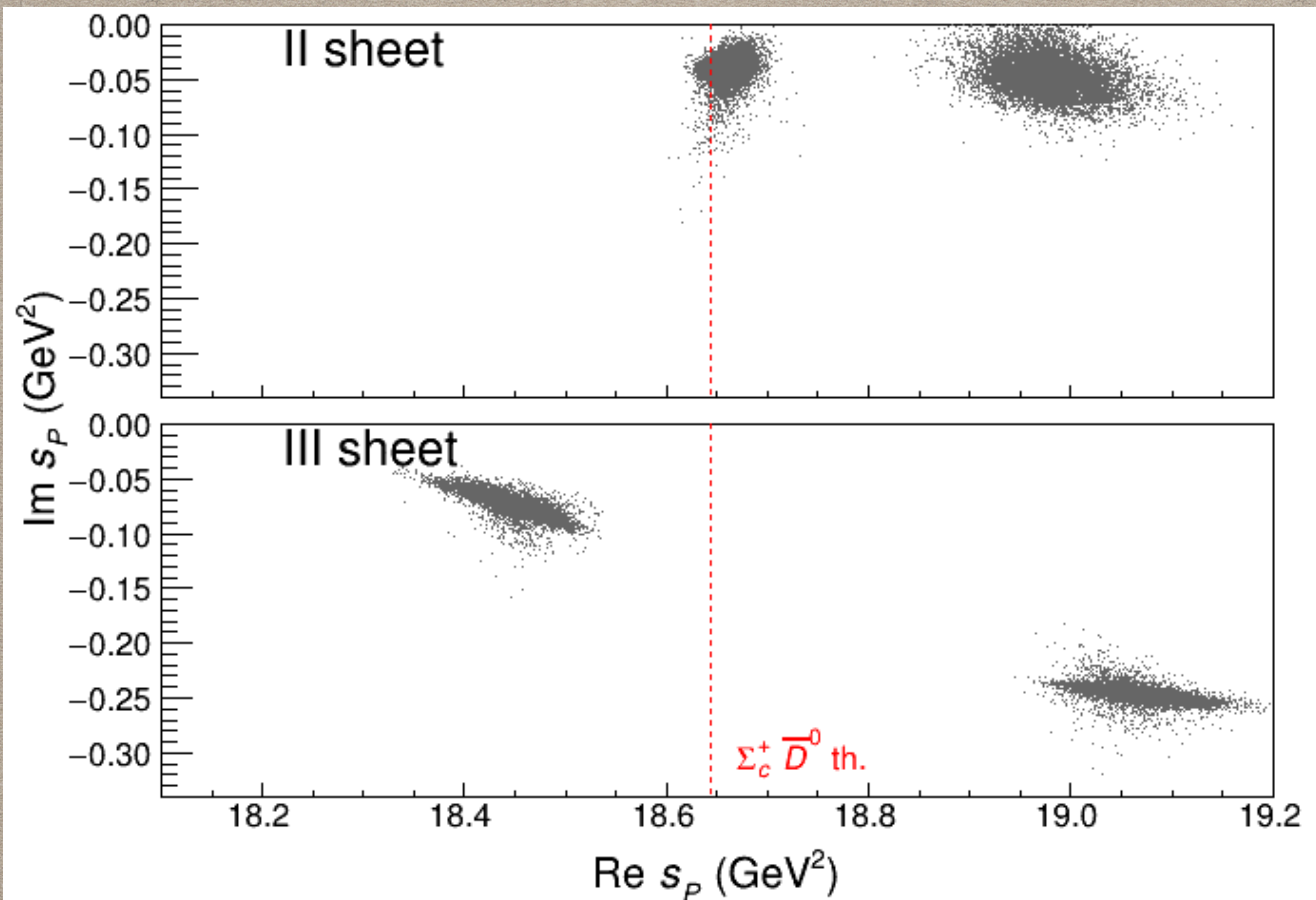
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$$\frac{[m_{22} - c_{22}s - ik_2]}{[m_{22} - c_{22}s - ik_2][m_{11} - c_{11}s - ik_1] - m_{12}^2}$$

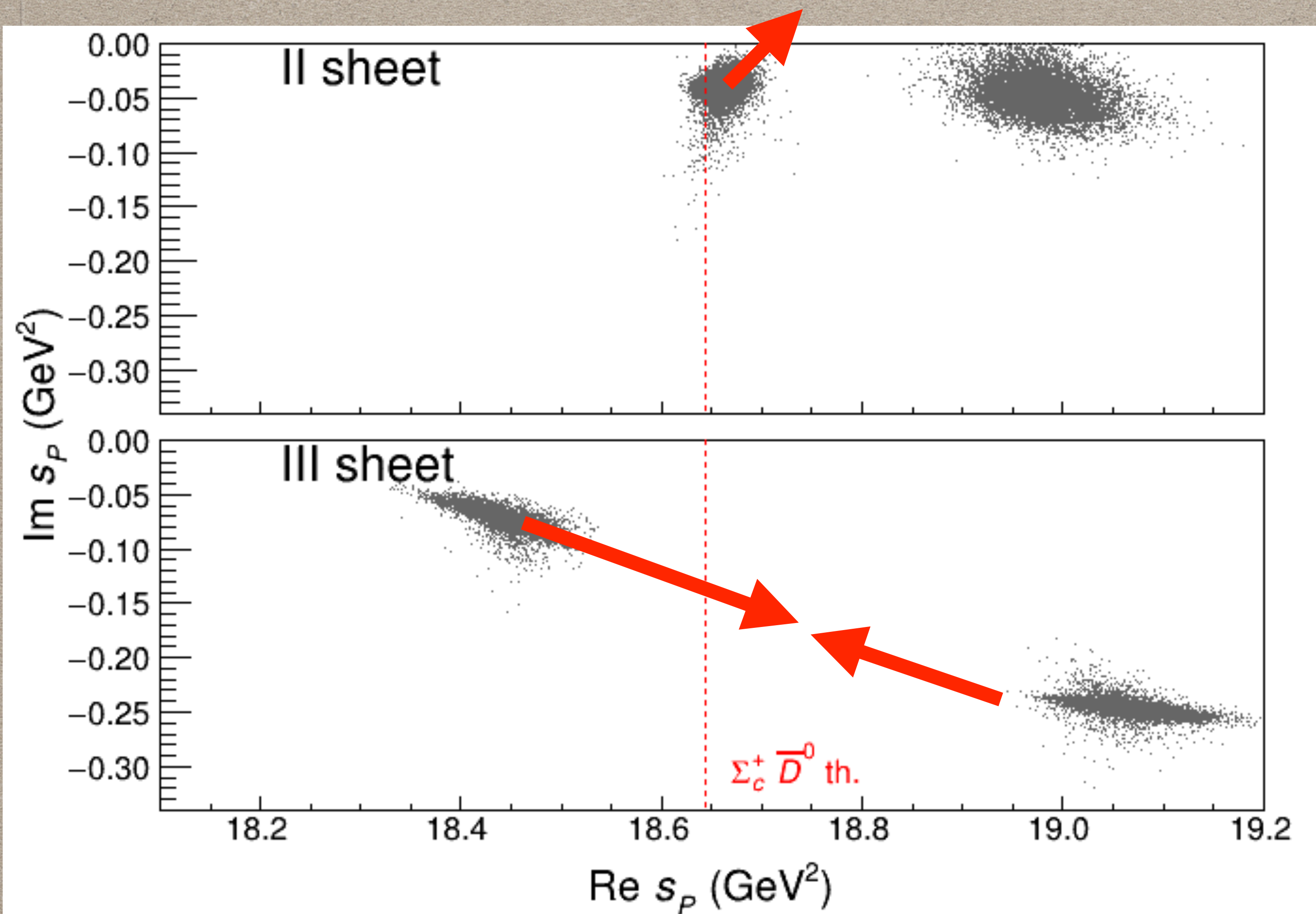
# POLE MOVEMENT: $C \neq 0$



When  $m_{12}=0$  both channels decouple and

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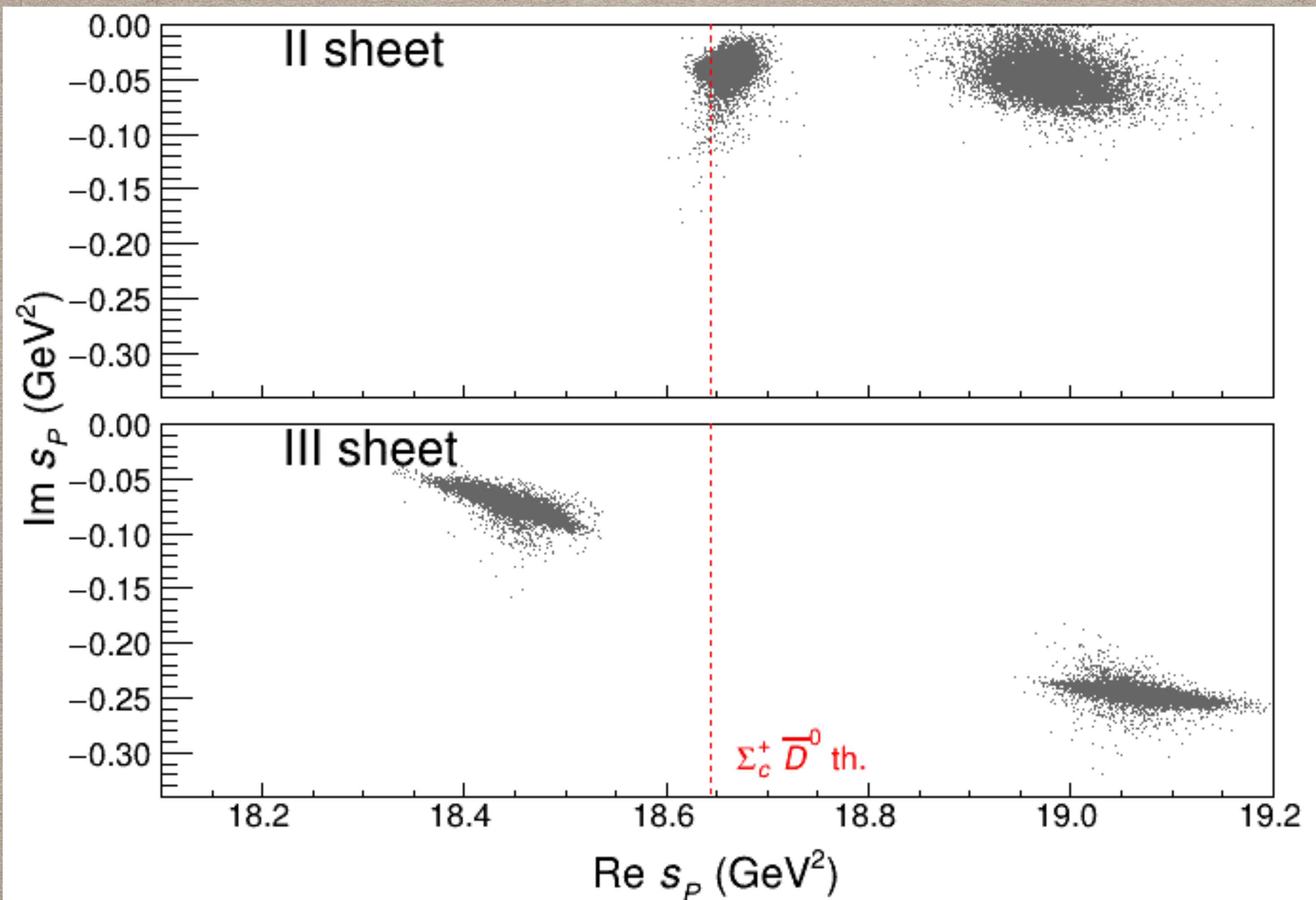


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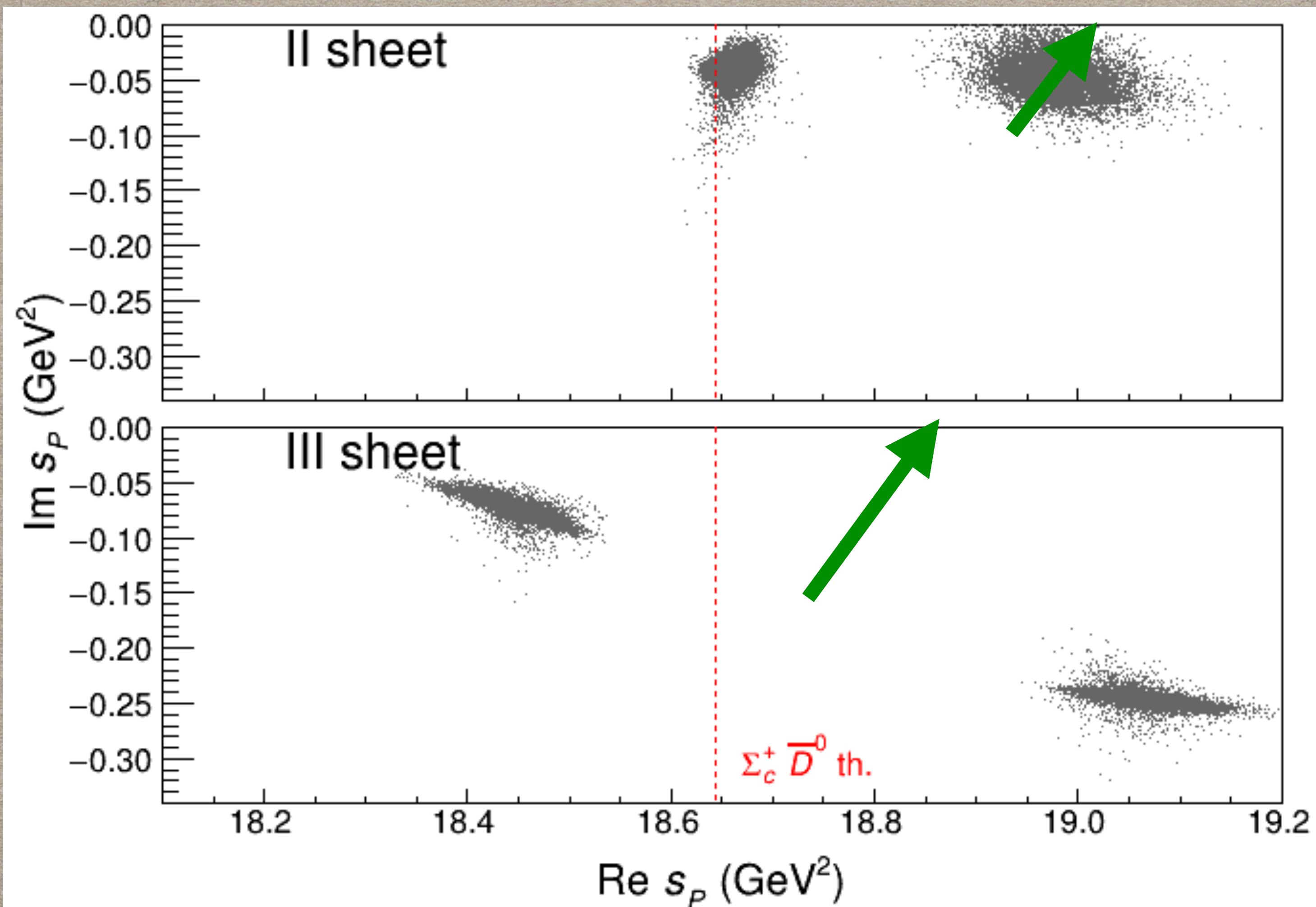
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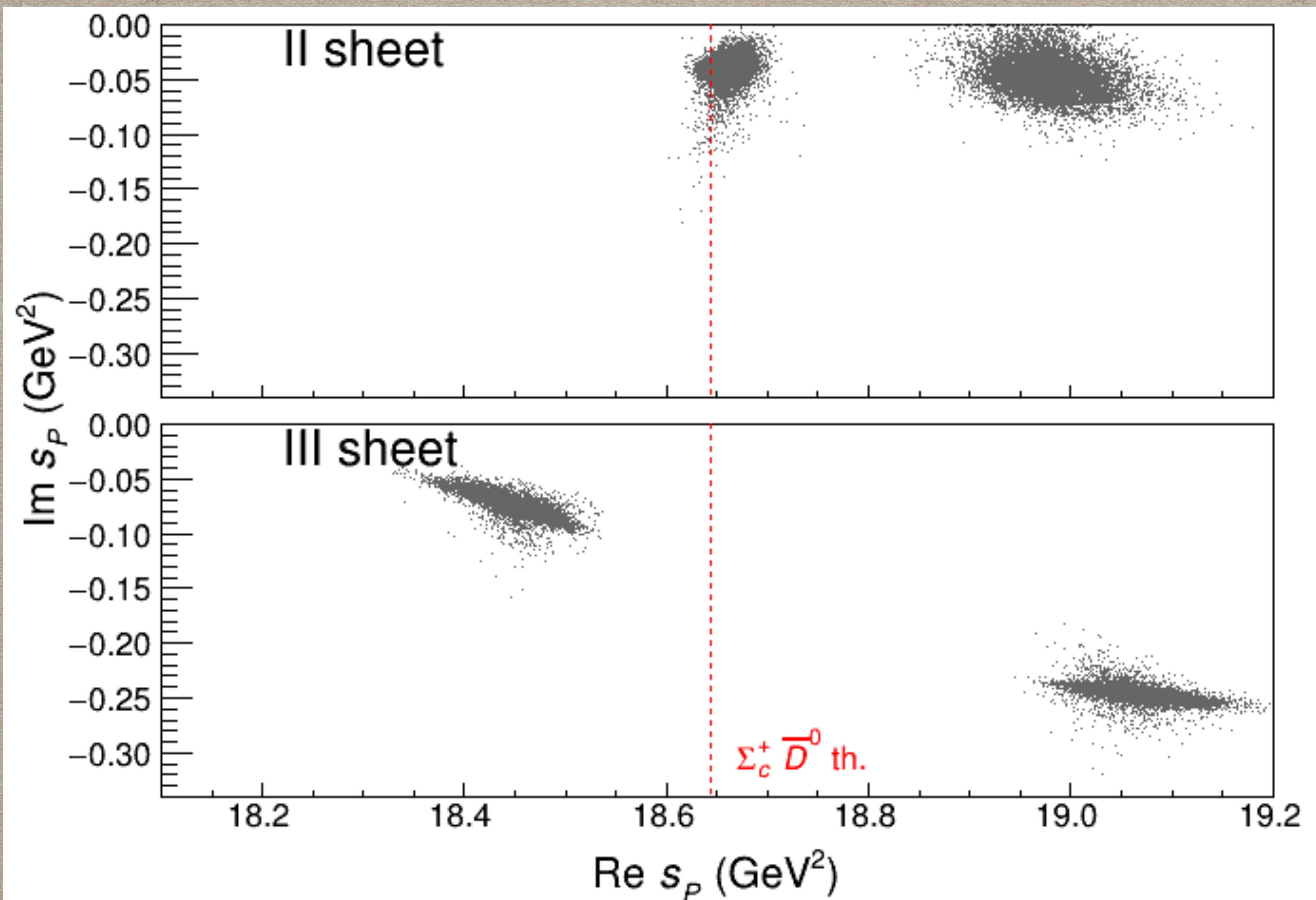
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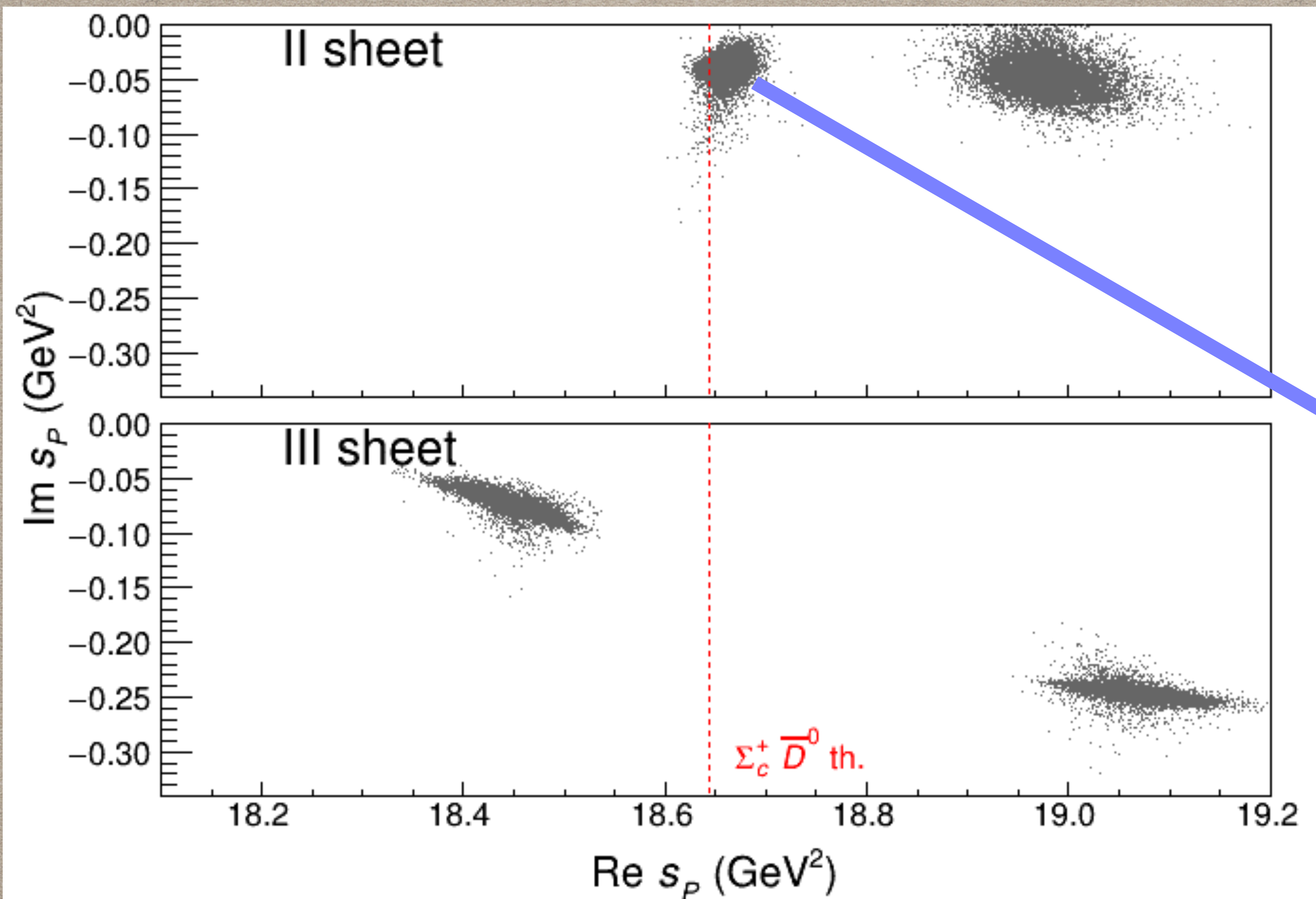
# POLE MOVEMENT: $C \neq 0$



When  $m_{12}=0$  both channels decouple and

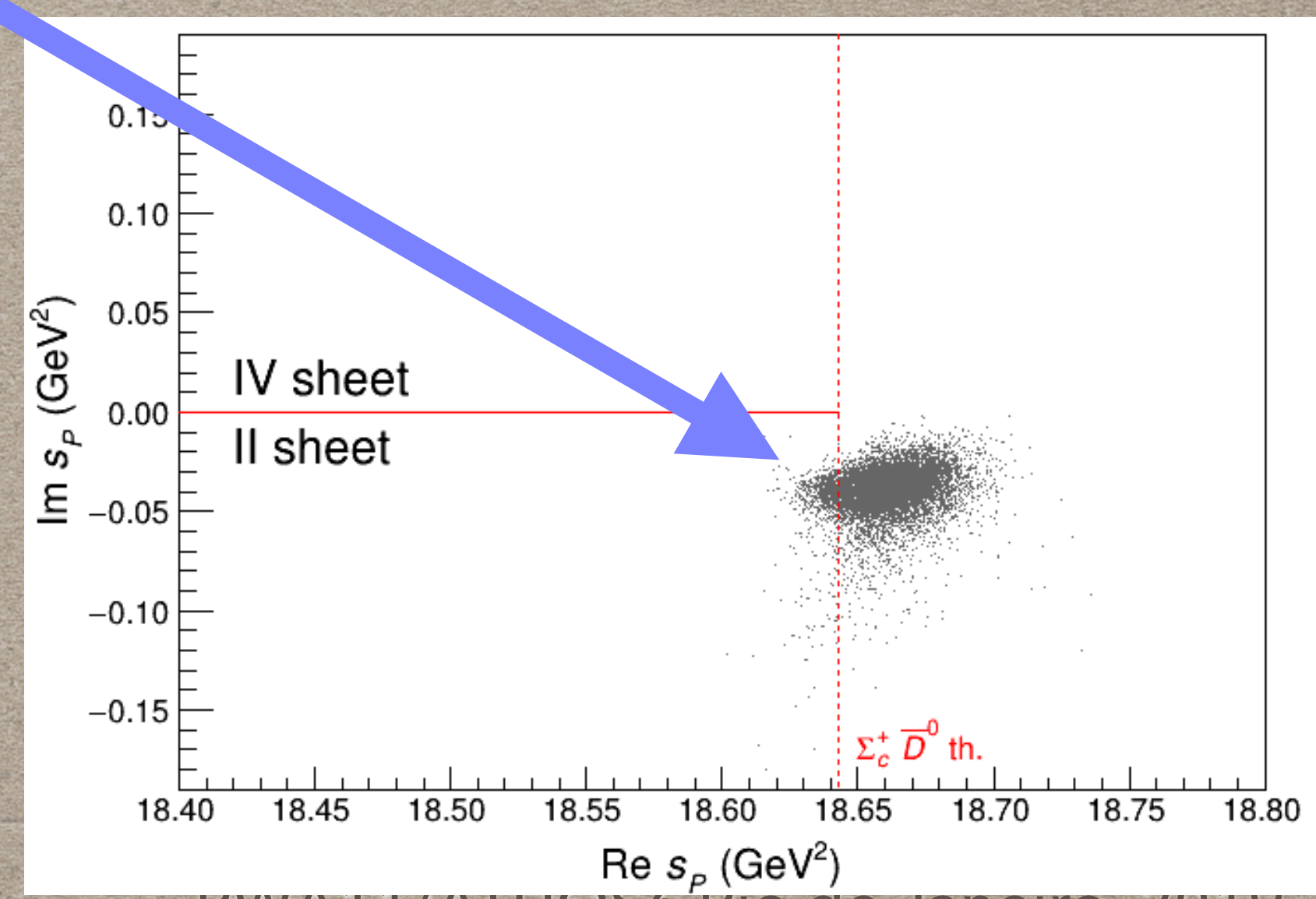
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# CONCLUSIONS

- In a near-threshold region we can build a minimally biased approach
- We can study pole stability against changes in the parameters compatible with the experimental uncertainties through bootstrap (more in Alessandro's talk)
- We can study pole motion, getting insight in the nature of the signal **without any prior model if the situation is simple enough**
- The  $P_c(4312)$  is a very suitable test case. Pole obtained :
  - $[M= 4319.7(1.6) \text{ MeV}; \Gamma=-0.8(2.4) \text{ MeV}; M= 4319.8(1.5) \text{ MeV}; \Gamma= 9.2(2.9) \text{ MeV}]$
- The favored interpretation based on pole motion is **Virtual state**

# ADVERTISEMENT



**CHARM 2020**  
in Mexico City,  
May 18-22  
at the Main Campus  
of UNAM  
(UNESCO World Heritage Site)

<https://indico.nucleares.unam.mx/e/charm20>

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