

#### Stockholm University



centre

# 21 CM COSMOLOGY AND Spin-dependent dark matter

#### **AXEL WIDMARK**

# JCAP 06(2019)014 arXiv:1902.09552

21 cm cosmology and spin temperature reduction via spin-dependent dark matter interactions

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# BACKGROUND

# **HYPERFINE STRUCTURE**



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AA (triplet) t↓ (singlet)

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44 (triplet) 44 (singlet)

 $E_{*} = 5.9 \mu eV$  $T_{*} = 0.068 K$  $\lambda_{\star} = 21 \text{ cm}$ 

### **SPIN TEMPERATURE**

nat = 3 e Ts

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# INTRODUCING A DARK SECTOR

# DARK SECTOR COOLING

- Can a cold dark matter component explain the EDGES signal?
  - Reduce the kinetic temperature of the hydrogen gas, ~2 % DM subcomponent of milli-charged dark matter (Berlin et al. 2018, Muñoz et al. 2018, Barkana et al. 2018)
- Alternate idea: affect the spin temperature directly, via spin dependent dark matter interactions

# SPIN DEPENDENT DARK MATTER

- Light dark matter fermion: χ (~MeV)
- Very light pseudo-vector force mediator: V(~eV)
- Hydrogen and dark fermion masses are mismatched – ineffective momentum transfer
- Cross section is enhanced like  $v^{-4}$ , down to  $v/c \simeq m_V/m_{\chi}$
- Further enhancement from collision inelasticity

# **DEEXCITATION (1)**

DM fermion



Hydrogen atom



### **DEEXCITATION (2)**



# **DEEXCITATION (3)**



(singlet)

## **DEEXCITATION (4)**



excess energy is carried away by DM fermion



(singlet)

# **INELASTIC FORM FACTOR**



# **EFFECTIVE TEMPERATURE**

• The hydrogen and DM gas temperatures are different  $(T_K \neq T_{\chi})$ , but the relative velocity between them can be parametrised by an effective temperature

$$\tilde{T}(m_H^{-1} + m_{\chi}^{-1}) = T_H m_H^{-1} + T_{\chi} m_{\chi}^{-1}$$

In the ideal case, where  $m_{\chi} \ll m_H$  and  $T_{\chi} = 0$ , we get

$$\tilde{T} = \frac{m_{\chi}}{m_H} T_H$$

# DARK SECTOR LIMITS

## **COUPLING CONSTANTS**

$$\mathcal{L} \supset g_{\chi} V_{\mu} \bar{\chi} \gamma^{\mu} \gamma^{5} \chi + g_{N} V_{\mu} \bar{N} \gamma^{\mu} \gamma^{5} N$$

DM-nucleon interaction:

$$\sigma_{\chi N} = \frac{g_{\chi}^2 g_N^2 m_{\chi}^2}{4\pi [(m_{\chi} v/c)^2 + m_V^2]^2}$$

DM self-interaction:

$$\sigma_{\chi\chi} = \frac{g_{\chi}^4 m_{\chi}^2}{8\pi [(m_{\chi}v/c)^2 + m_V^2]^2}$$

# LIMITS FROM DM SELF-INTERACTION

Bullet Cluster: 
$$g_{\chi}^4 \lesssim 1.8 \times 10^{-13} \left(\frac{m_{\chi}}{\text{MeV}}\right)^3$$

If we consider a DM subcomponent (<30 %), self-interaction can be arbitrarily strong

$$g_{\chi} \simeq 1$$

Fan et al., Phys. Dark Univ. 2 (2013) 139-156.

#### **SCALAR MEDIATOR** $10^{-3}$ $B \to K \phi$ $10^{-6}$ n-Xe $K \to \pi \; \phi$ 5th force $g_N$ 10<sup>-9</sup> SN1987a

HB stars

RG stars

keV

 $m_V$ 

 $10^{-12}$ 

 $10^{-15}$ 

eV

 $\phi \overline{t} t$ 

MeV

GeV

# **PSEUDO-VECTOR MEDIATOR**

- Similar limits apply in the spin-dependent case, although there are some additional complications for pseudo-vector mediators
- Energy/ $m_V^2$  enhanced coupling to the mediator's longitudinal mode in stellar cooling bounds, giving

$$g_N < 10^{-12} \times \frac{m_V}{10 \text{ keV}}$$

- Such bounds are model dependent, a more complete particle model is highly constrained
- The dominant process of stellar cooling is the Compton process with He, which is suppressed for spin-dependent interactions

RESULTS



# **DOES IT WORK?**

- Significant spin temperature cooling by this mechanism requires  $g_N > 10^{-11}$ 
  - Excluded by a few orders of magnitude (stellar cooling bounds give  $g_N < 10^{-12}$ , or even stronger with longitudinal mode enhancement) similar scenario for coupling to the electron
  - Maybe some limits can be relaxed (helium Compton process suppression, more complete dark sector)
- Perhaps the same mechanism can cool the spin gas in an alternative model, subject to a different set of constraints

