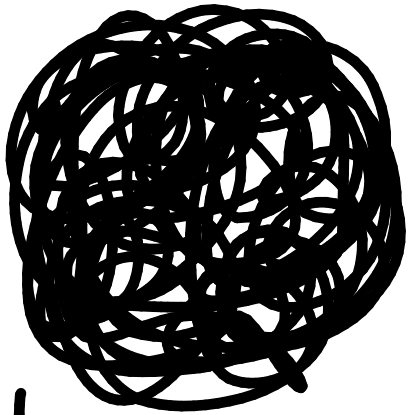


Unitarity and Entropy

Gia Dvali

LMU - MPI & NYU

Bekenstein - Bremermann Bound



$\leftarrow R \rightarrow$

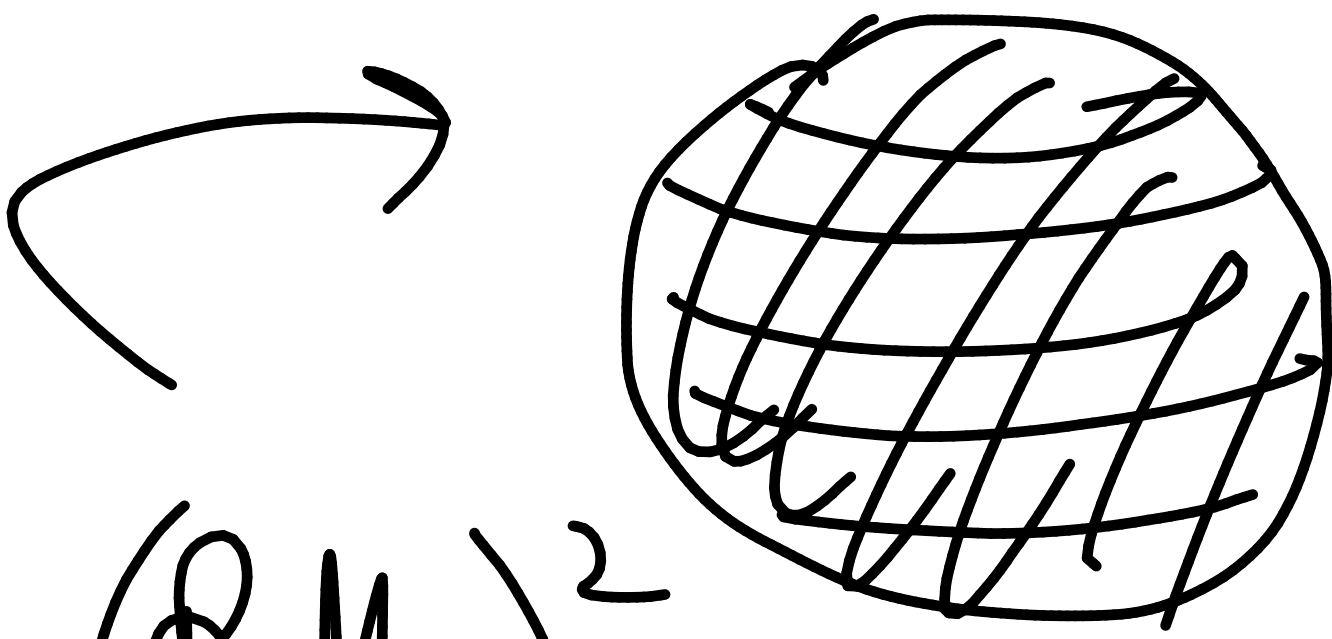
$$S \leq MR$$

In gravity is saturated by black holes.
Bekenstein entropy:

$$S_{BH} = I_{BH} R = (R M_{Pl})^2$$

⊛ What is physical meaning of bound?

⊛ What is physical meaning of Area-law?


$$S = (R_{MP})^2$$

The bound has (mostly) been discussed in gravity.

What happens beyond gravity?

In renormalizable theories?

Our main results:

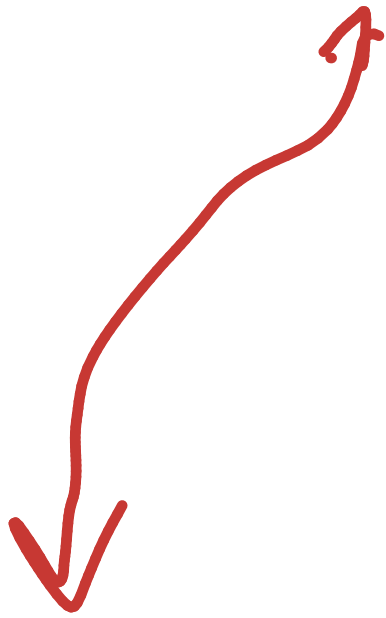
⊛ Bekenstein bound is saturated when theory saturates the bound on unitarity.

⊛ Simultaneously the entropy assumes the area law:

$$S = \ln R = (Rf)^2$$

↑ scale

Bekenstein \Rightarrow Unitarity $=$ Area



$$S_{\max} \approx \frac{1}{g^2} = \text{Area}$$



Quantum
Coupling

We shall demonstrate
for:

① 't Hooft - Polyakov
monopole:

$$S = M_m R_m = \frac{1}{g_{\text{gauge}}^2} = (R_m v)^2$$

② Baryon:

$$S = M_B R_B = \frac{1}{g_{\text{QCD}}^2} = (R_B f_\pi)^2$$

Monopole. $SO(3)$ gauge
symmetry Higgsed by

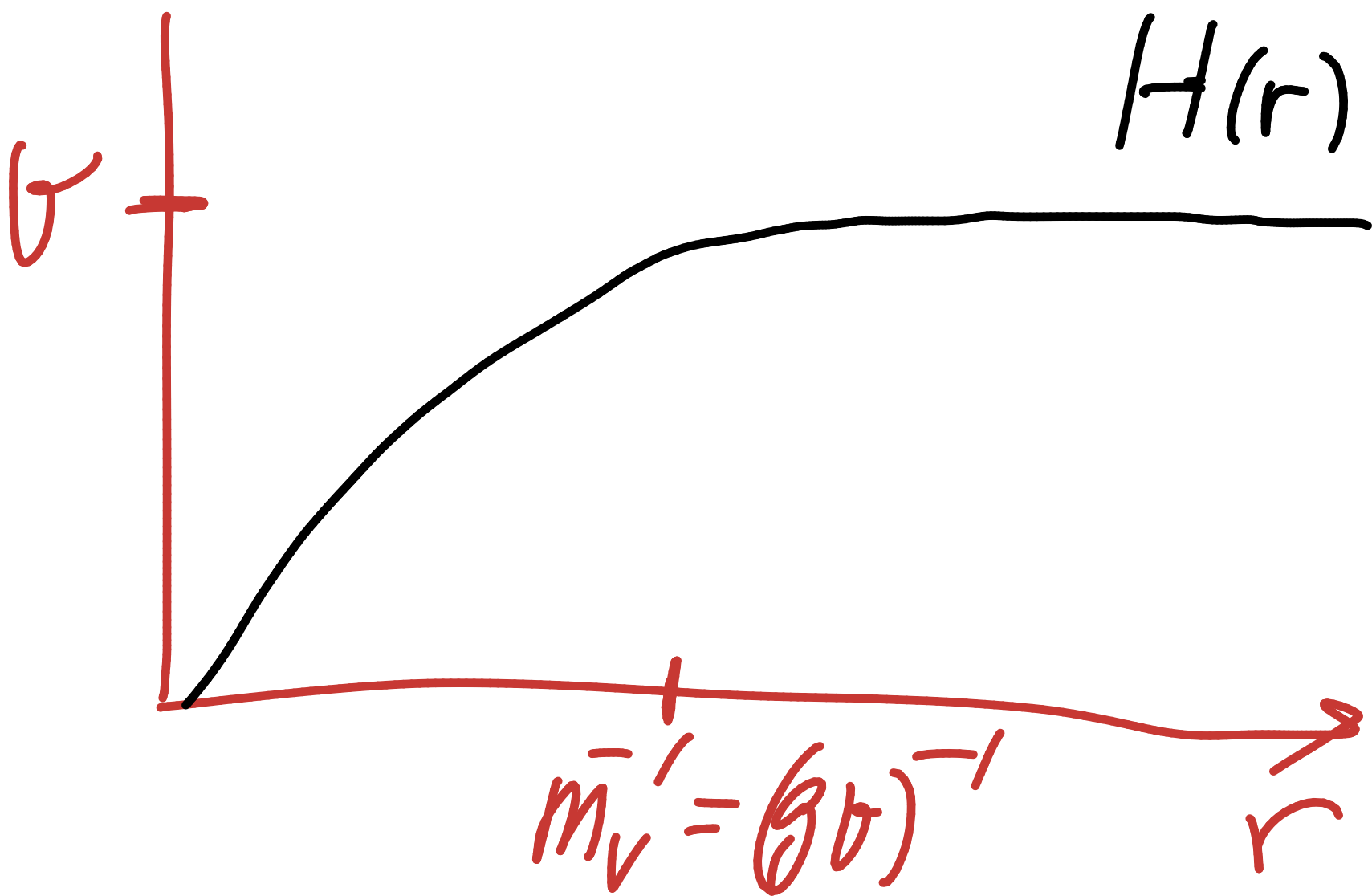
$$\phi^a \quad (a=1,2,3)$$

$$\begin{aligned} \mathcal{L} = & \partial_\mu \phi^a \partial^\mu \phi^a \\ & - \lambda^2 (\phi^a \phi^a - v^2)^2 \\ & - F_{\mu\nu}^a \tilde{F}^{a\mu\nu} \end{aligned}$$

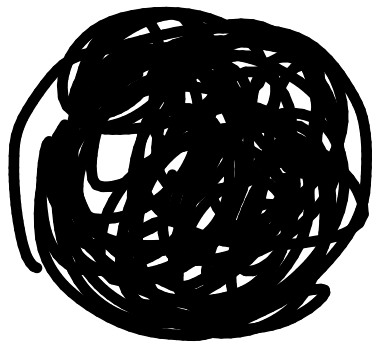
Monopole:

$$\phi^a = \frac{x^a}{r} H(r)$$

$$A_\mu^a = \epsilon_{0a\mu\nu} \frac{x^\nu}{gr^2} F(r)$$



Monopole mass and size



$$R_m = (M_V)^{-1} = (gV)^{-1}$$

← R_m → $M_m = \frac{M_V}{g^2}$

Entropy bound on monopole:

$$S \leq M_m R_m = \frac{1}{g^2}$$

Can it be saturated?

Entropy from Goldstones

$$\sigma_\alpha \quad \alpha = 1, 2, \dots, N$$

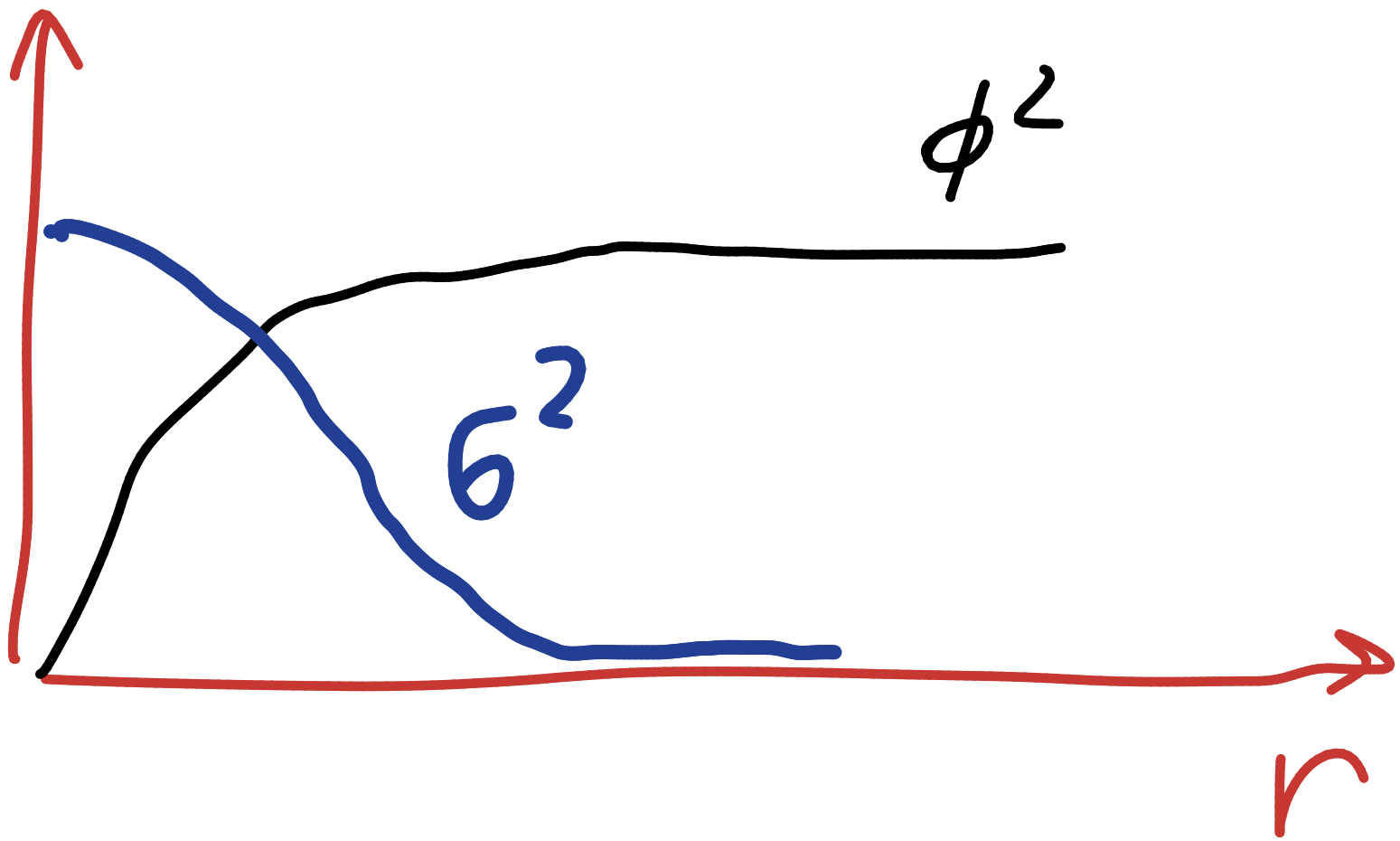
\leftarrow scalar

$SO(N)$ - global symmetry
spontaneous breaking
in monopole:

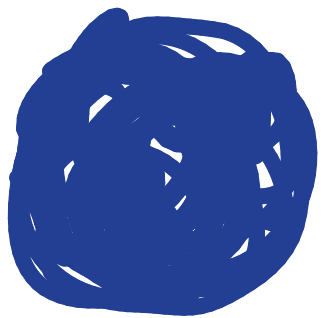
$$L = \partial_\mu \sigma_\alpha \partial^\mu \sigma_\alpha$$

$$- (g^2 \phi^2 - m^2) \sigma_\alpha \sigma_\alpha$$

$$- \frac{g^2}{6} (\sigma_\alpha \sigma_\alpha)^2$$



$\langle \sigma \rangle \neq 0$



$SO(N) \rightarrow SO(N-1)$



$\sim N$ localized Goldstones!

Number of degenerate
micro-states

$$N_{st} \sim \binom{2N}{N} \sim 2^{2N}$$

Monopole entropy:

$$S_{\text{mon}} = \ln(N_{st}) \sim N$$

Unitarity bound:

$$g^2 N \leq 1$$

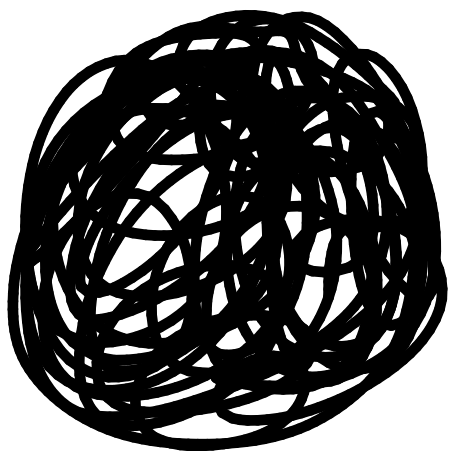
For $N = \frac{1}{g^2}$ entropy bound is saturated by area-law!

$$S_{\text{max}} = N = \frac{1}{g^2} = M_m R_m = (R_m b)^2$$

The same is achieved
by coupling to fermion
flavors:

$$\phi^a \psi_\alpha^b \lambda_\alpha^c \epsilon_{abc}$$

↙ $\alpha = 1, 2 \dots N$
 $SO(N)$



↙ $\sim N$ localized
fermion zero modes!

Again:

$$n_{st} \sim 2^N$$

Entropy: $S_{\text{mon}} \sim N$

Unitarity bound

$$g^2 N \leq 1$$

$S_{\text{mon}} = N = \frac{1}{g^2} = M_m R_m = (R_m \alpha')^2$

Baryons in $SU(N_c)$

QCD with

N -flavors of quarks.

't Hooft limit

$$N_c \rightarrow \infty$$

$$g^2 N_c = \text{fixed}$$

$$\Lambda = \text{fixed}$$

Spontaneous chiral
symmetry breaking

$$U(N)_L \otimes U(N)_R \rightarrow U(N)_F$$

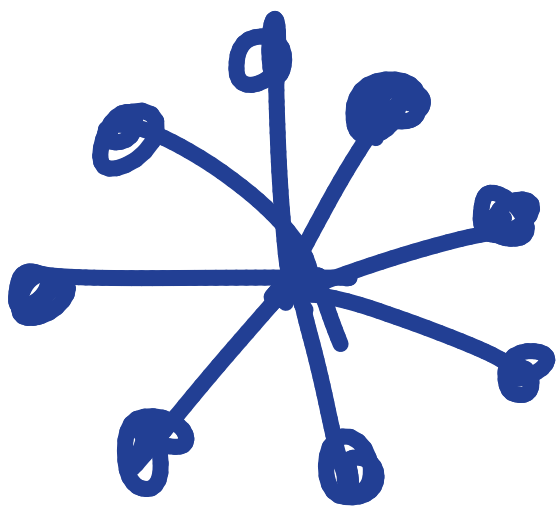
$N^2 - 1$ Goldstones (pions)

+ η' -meson

Pion decay constant:

$$f_\pi = \sqrt{N_c} \Lambda$$

Baryons (Witten).



} N_c quarks

$$\left\langle R_B \sim \Lambda^{-1} \right\rangle$$

Mass: $M_B \sim N_c \Lambda$

Entropy bound:

$$S_{MAX} = M_B R_B = N_c$$

I_s is saturated at
the unitarity bound:

$$N_c \sim N \sim \frac{1}{g^2}$$

Baryon entropy:

$$S_B = \ln \binom{N_c + N}{N_c} \sim N$$

Thus, at the unitarity bound we have:

$$S_B = M_B R_B = \frac{1}{g^2} = N =$$
$$= (R_B \sqrt{\pi})^2$$

Area low!

Conclusions:

The universal
phenomenon:

B.B. = unitarity = Area



Bound = $\frac{1}{\text{coupling}}$ = Area

Can confinement in QCD (at least for large N_c) be understood as preventive mechanism against violation of Bekenstein bound by free colored states?

Entropy of a free quark

$$S_q = \ln(2N_c N_f)$$

$$S_{\max} = M_{\text{QR}}$$