

An aerial photograph of a cyclotron's dees and particle beam lines. The dees are large, light blue, curved structures that form the two halves of the particle accelerator. They are arranged in a circular pattern, with the particle beam lines connecting them. The background shows the complex machinery and infrastructure of the accelerator facility.

Cyclotrons

CERN Accelerator School – Introductory Course
High Tatras, September 19, 2019

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Cyclotrons/FFA - Outline

- the classical cyclotron
 - history of the cyclotron, basic concepts and scalings, focusing, stepwidth, classification of cyclotron-like accelerators
- synchro-cyclotrons
 - concept, synchronous phase, example
- isochronous cyclotrons (→ sector cyclotrons)
 - isochronous condition, focusing in Thomas-cyclotrons, spiral angle, classical extraction: pattern/stepwidth, space charge
- applications and examples of existing cyclotrons
 - TRIUMF, RIKEN SRC, PSI Ring, PSI medical cyclotron

Part II

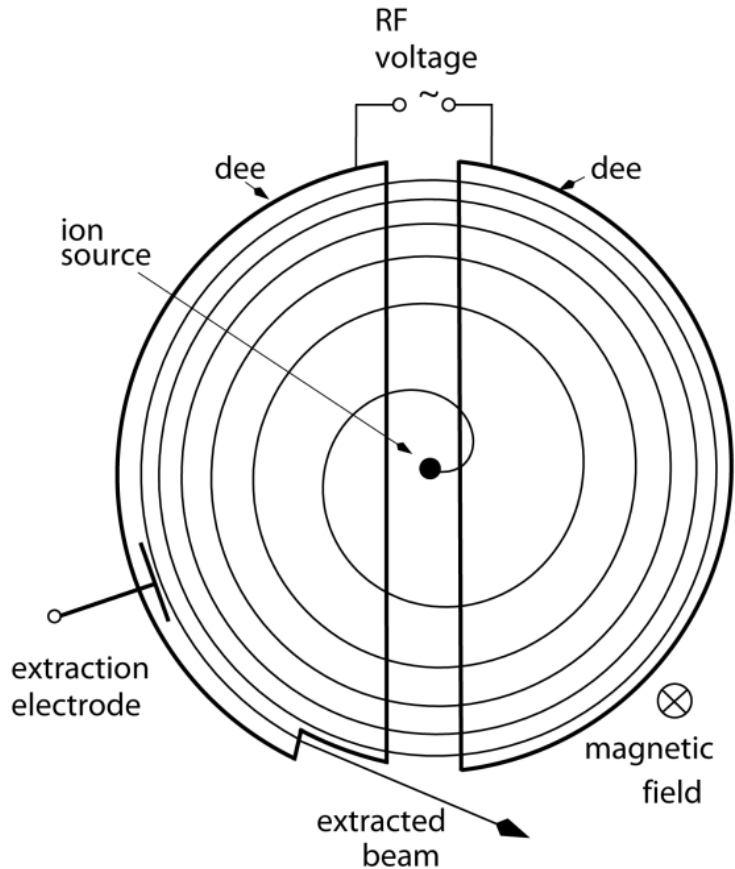
- cyclotron subsystems
 - Injection/extraction schemes, RF systems/resonators, magnets, vacuum issues, instrumentation, FFA specific magnets, FFA resonators
- FFA = Fixed Focus Alternating Gradient Accelerators
 - motivation & applications, scaling FFA's, non-scaling and linear FFA, FFA subsystems
- discussion
 - classification of circular accelerators, Pro's and Con's of cyclotrons / FFA for different applications



The Classical Cyclotron

two capacitive electrodes
„Dees“, two gaps per turn
internal ion source
homogenous B field
constant revolution time
(for low energy, $\gamma \approx 1$)

$$\omega_c = \frac{eB_z}{m}$$



powerful concept:

- **simplicity, compactness**
- **continuous injection/extraction**
- **multiple usage of accelerating voltage**

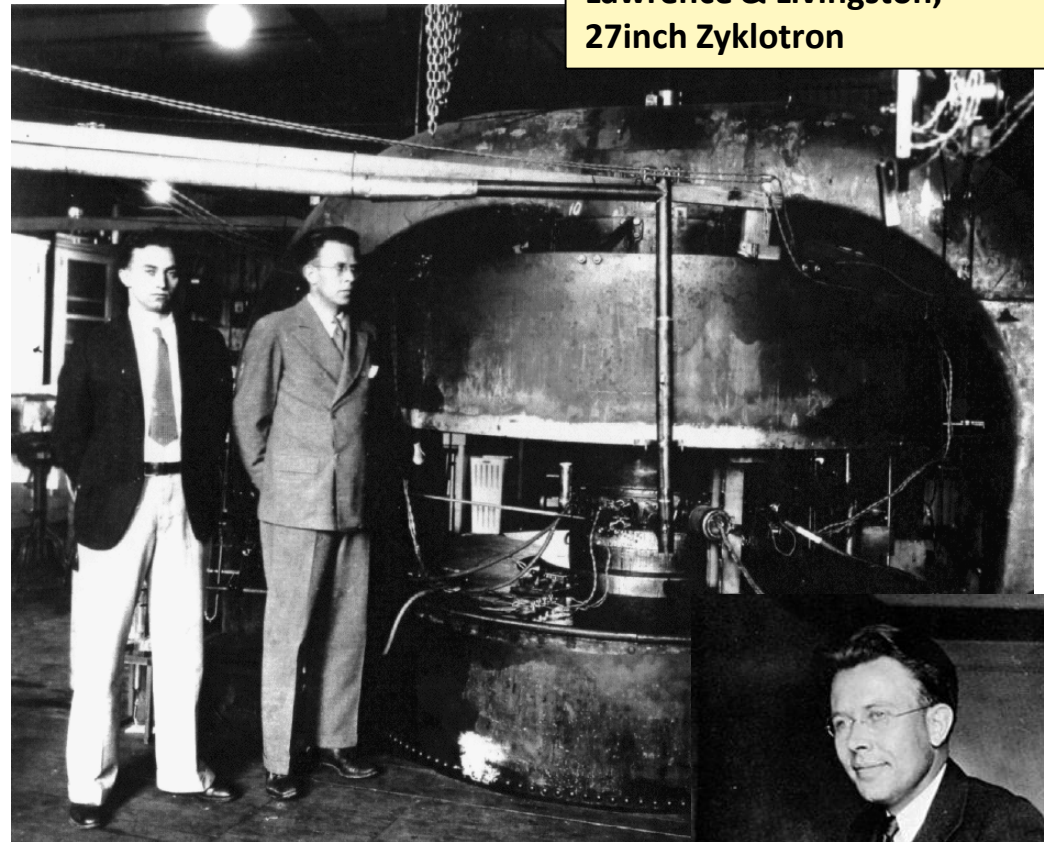


some History ...

first cyclotron: 1931, Berkeley
1kV gap-voltage 80keV Protons



Lawrence & Livingston,
27inch Zyklotron



Ernest Lawrence, Nobel Prize 1939
"for the invention and development of the cyclotron and for results obtained with it, especially with regard to artificial radioactive elements"

John Lawrence (center), 1940'ies
first medical applications: treating patients with neutrons generated in the 60inch cyclotron





PSI Ring Cyclotron & Crew



cyclotron frequency and K value

- **cyclotron frequency** (homogeneous) B-field:

$$\omega_c = \frac{eB}{\gamma m_0}$$

- **cyclotron K -value:**

→ K is the **kinetic energy reach** for protons **from bending strength** in non-relativistic approximation:

$$K = \frac{e^2}{2m_0} (B\rho)^2$$

→ K can be used to rescale the energy reach of protons to other charge-to-mass ratios:

$$\frac{E_k}{A} = K \left(\frac{Q}{A} \right)^2$$

→ K in [MeV] is often used for naming cyclotrons

examples: **K-130 cyclotron / Jyväskylä**
 cyclone C230 / IBA



cyclotron - isochronicity and scalings

continuous acceleration → revolution time should stay constant, though E_k , R vary

magnetic rigidity:

$$BR = \frac{p}{e} = \beta\gamma \frac{m_0 c}{e}$$

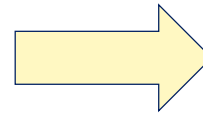
orbit radius from isochronicity:

$$R = \frac{c}{\omega_c} \beta = R_\infty \beta$$

deduced scaling of B :

$$\longrightarrow B(R) \propto \gamma(R)$$

**to be isochronous, B must be raised $\propto \gamma(R)$
→ this contradicts the focusing requirements!**



main difficulty to be overcome by cyclotron & FFA variants.

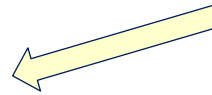


field index

the field index describes the (normalized)
radial slope of the bending field:

$$\begin{aligned}k &= \frac{R}{B} \frac{dB}{dR} \\ &= \frac{\beta}{\gamma} \frac{d\gamma}{d\beta} \\ &= \gamma^2 - 1\end{aligned}$$

from isochronous condition:
 $B \propto \gamma$, $R \propto \beta$



→ thus $k > 0$ (positive slope of field) to keep beam isochronous!



focusing in a classical cyclotron

centrifugal force mv^2/r



Lorentz force $qv \times B$



$$m\ddot{r} = mr\dot{\theta}^2 - qr\dot{\theta}B_z$$

focusing: consider small deviations x from beam orbit R ($r = R+x$):

$$\ddot{x} + \frac{q}{m}vB_z(R+x) - \frac{v^2}{R+x} = 0,$$

$$\ddot{x} + \frac{q}{m}v \left(B_z(R) + \frac{dB_z}{dR}x \right) - \frac{v^2}{R} \left(1 - \frac{x}{R} \right) = 0,$$

$$\ddot{x} + \omega_c^2(1+k)x = 0.$$

using: $\omega_c = qB_z/m = v/R$, $r\dot{\theta} \approx v$, $k = \frac{R}{B} \frac{dB}{dR}$



betatron tunes in cyclotrons

thus in radial plane:

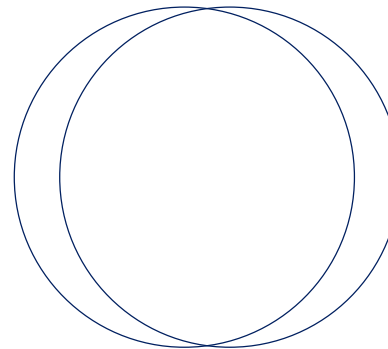
$$\omega_r = \omega_c \sqrt{1+k} = \omega_c \nu_r$$

$$\nu_r = \sqrt{1+k}$$

$$\approx \gamma$$

using isochronicity condition

note: simple case for $k = 0$: $\nu_r = 1$
 (one circular orbit oscillates w.r.t the other)

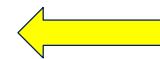


using Maxwell to relate B_z and B_R :

$$\text{rot } \vec{B} = \frac{dB_R}{dz} - \frac{dB_z}{dR} = 0$$

in vertical plane:

$$\nu_z = \sqrt{-k}$$



$k < 0$ to obtain
vertical focus.

**thus: in classical cyclotron $k < 0$ required for vert. focus;
 however **this violates isochronous condition** $k = \gamma^2 - 1 > 0$**



cyclotron stepwidth classical (nonrelativistic, B const)

equation of motion for ideal centroid orbit R ,
 → relation betw. **energy and radius**

$$m\ddot{R} = m\frac{v^2}{R} - qvB_z = 0$$

$$qRB_z = \sqrt{2mE_k}$$

$$\frac{dR}{R} = \frac{1}{2} \frac{dE_k}{E_k}$$

centrifugal f.

Lorentz f.

“cyclotron language”

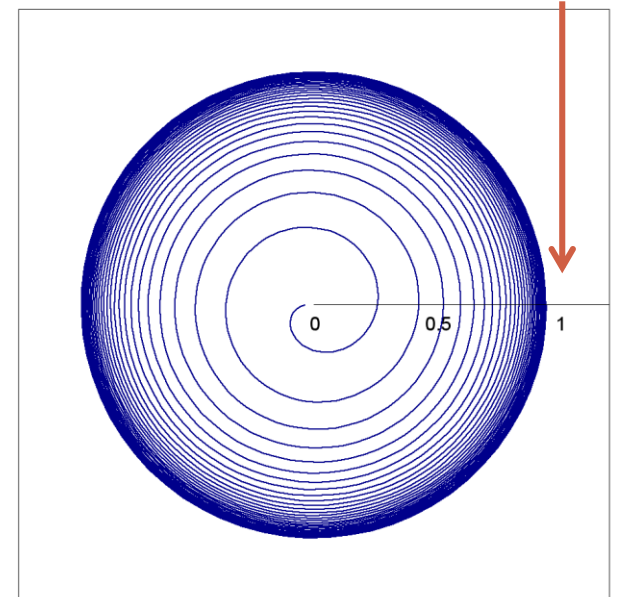
$$R_\infty = R/\beta$$

use: $\Delta E_k = \text{const}; B_z = \text{const}; E_k \propto R^2$

thus:

$$\Delta R \propto \frac{R}{E_k} \propto \frac{1}{R}$$

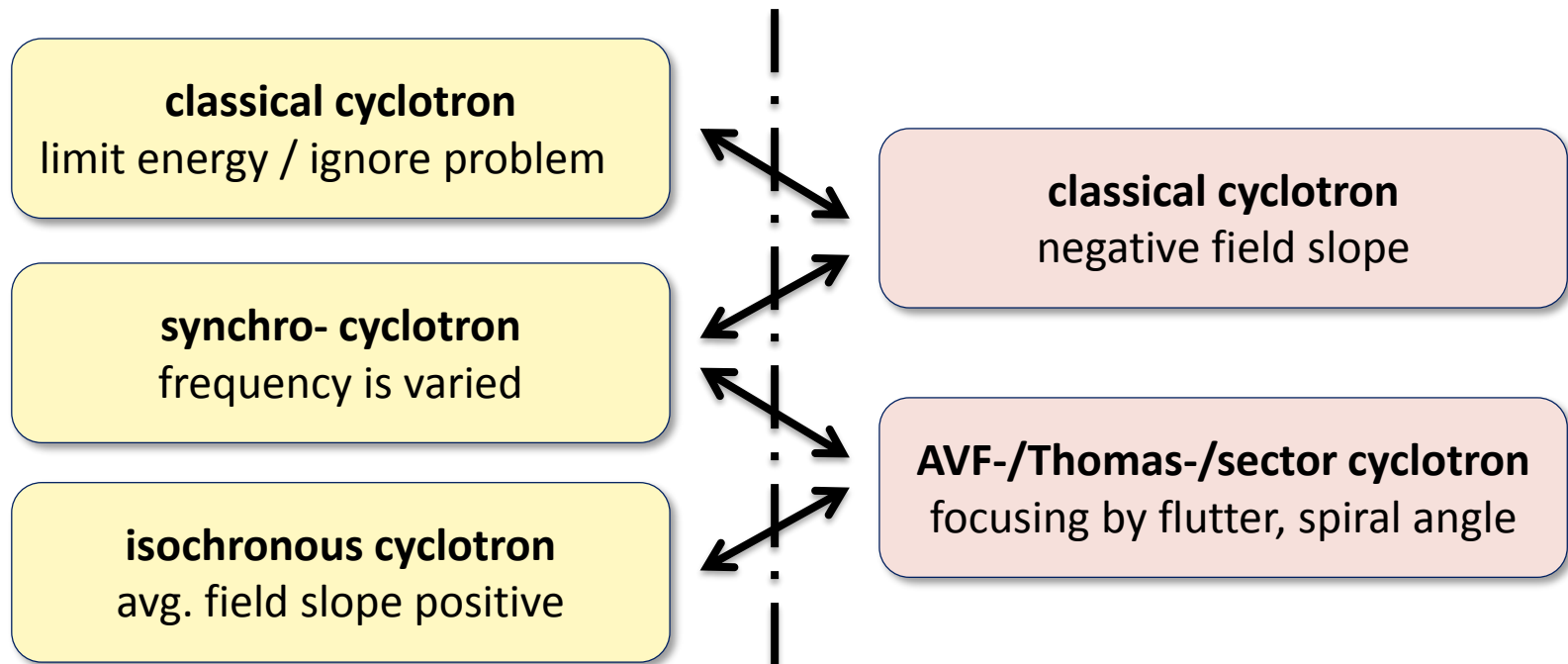
radius increment per turn decreases with increasing radius
 → **extraction becomes more and more difficult at higher energies**



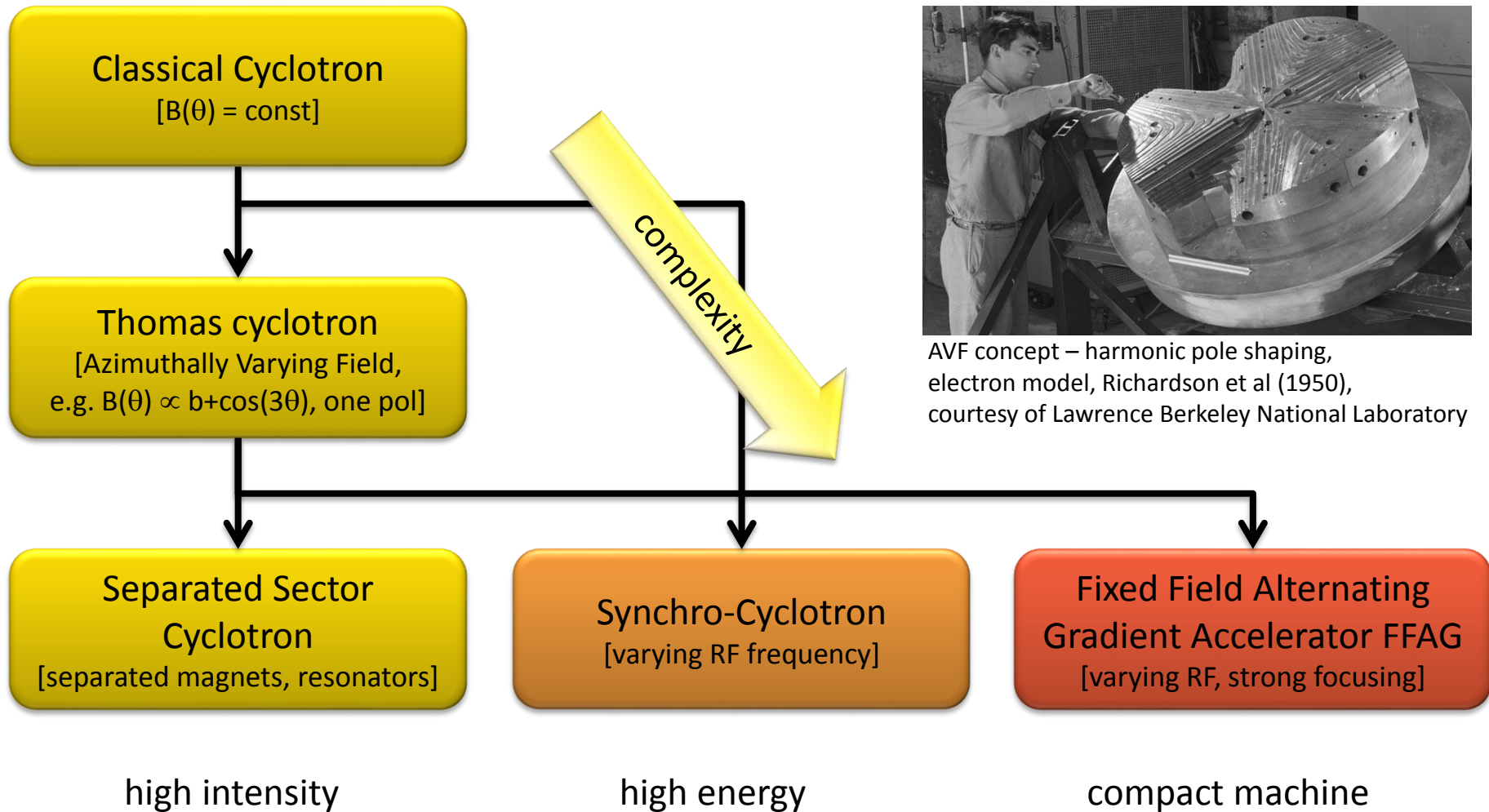
naming conventions of cyclotrons ...

1.) resonant acceleration

2.) transverse focusing



classification of cyclotron like accelerators



AVF concept – harmonic pole shaping, electron model, Richardson et al (1950), courtesy of Lawrence Berkeley National Laboratory



next: **synchro-cyclotrons**

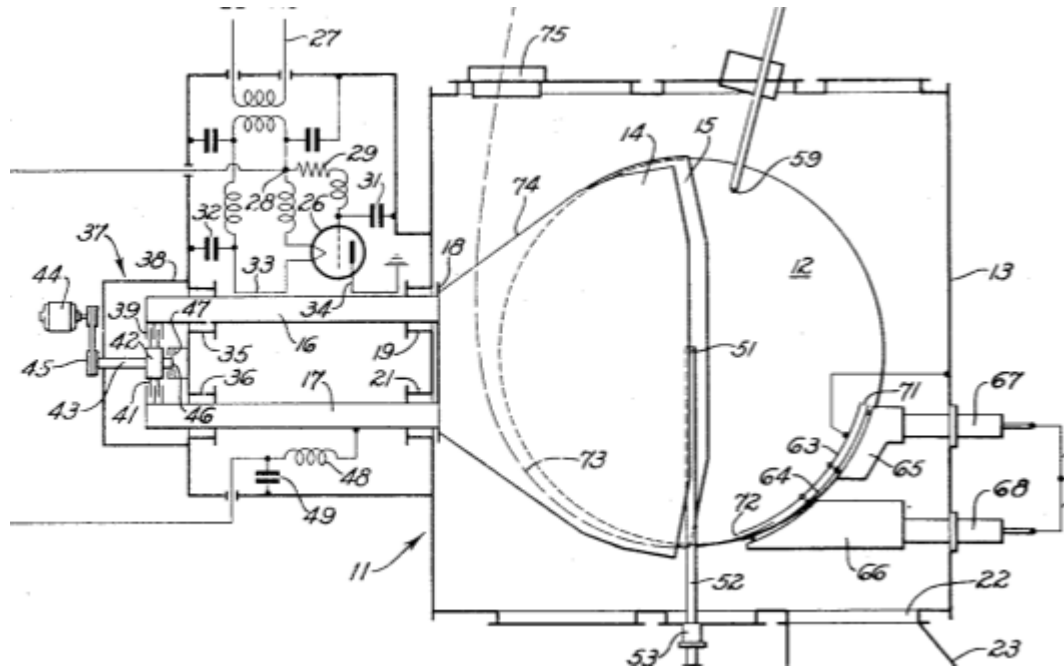
- concept and properties
- frequency variation and synchronous phase
- an example for a modern synchrocyclotron

*exciting
coil*

*pole
piece*



Synchrocyclotron -concept



first proposal by
Mc.Millan, Berkeley

- accelerating frequency is variable, is reduced during acceleration
- negative field index (= negative slope) ensures sufficient focusing
- operation is pulsed, thus avg. intensity is low
- bending field constant in time, thus rep. rate high, e.g. 1kHz

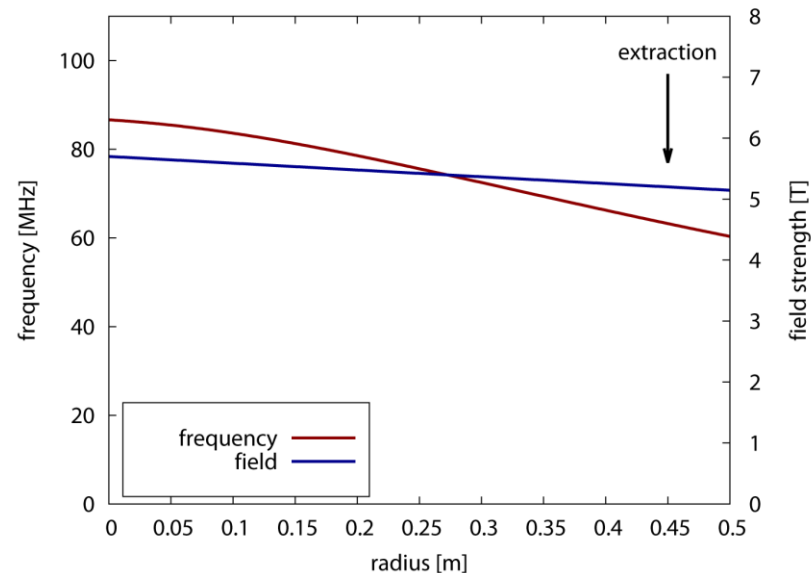


Synchrocyclotron continued

advantages	disadvantages
<ul style="list-style-type: none">- high energies possible ($\geq 1\text{Gev}$)- focusing by field gradient, no complicated flutter required \rightarrow thus compact magnet- only RF is cycled, fast repetition as compared to synchrotron	<ul style="list-style-type: none">- low intensity, at least factor 100 less than CW cyclotron- complicated RF control required- weak focusing, large beam

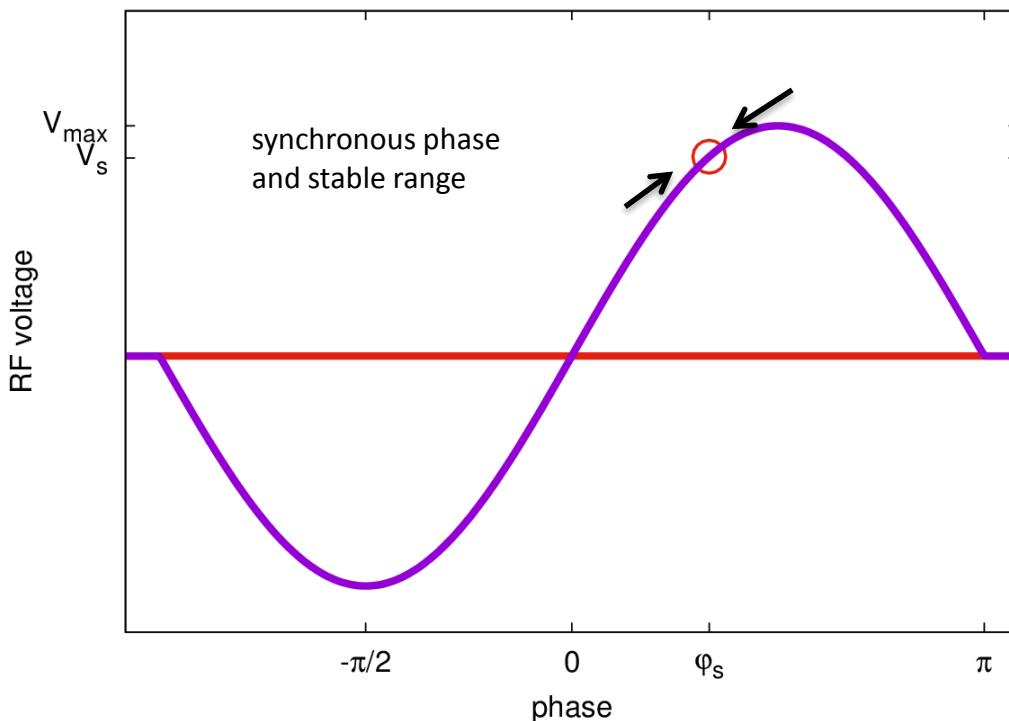
**numerical example
field and frequency vs.
radius:**

- 230MeV p, strong field
- RF curve must be programmed in some way



Synchrocyclotron and synchronous phase

- internal source generates continuous beam; only a fraction is captured by RF wave in a phase range around a synchronous particle
- in comparison to a synchrotron the “storage time” is short, thus in practice no synchrotron oscillations



relation of
energy gain per turn and
rate of frequency change

$$\frac{qU_0 N \cos \varphi_s}{E_k + E_0} = -\frac{2\pi}{\omega^2} \frac{d\omega}{dt}$$



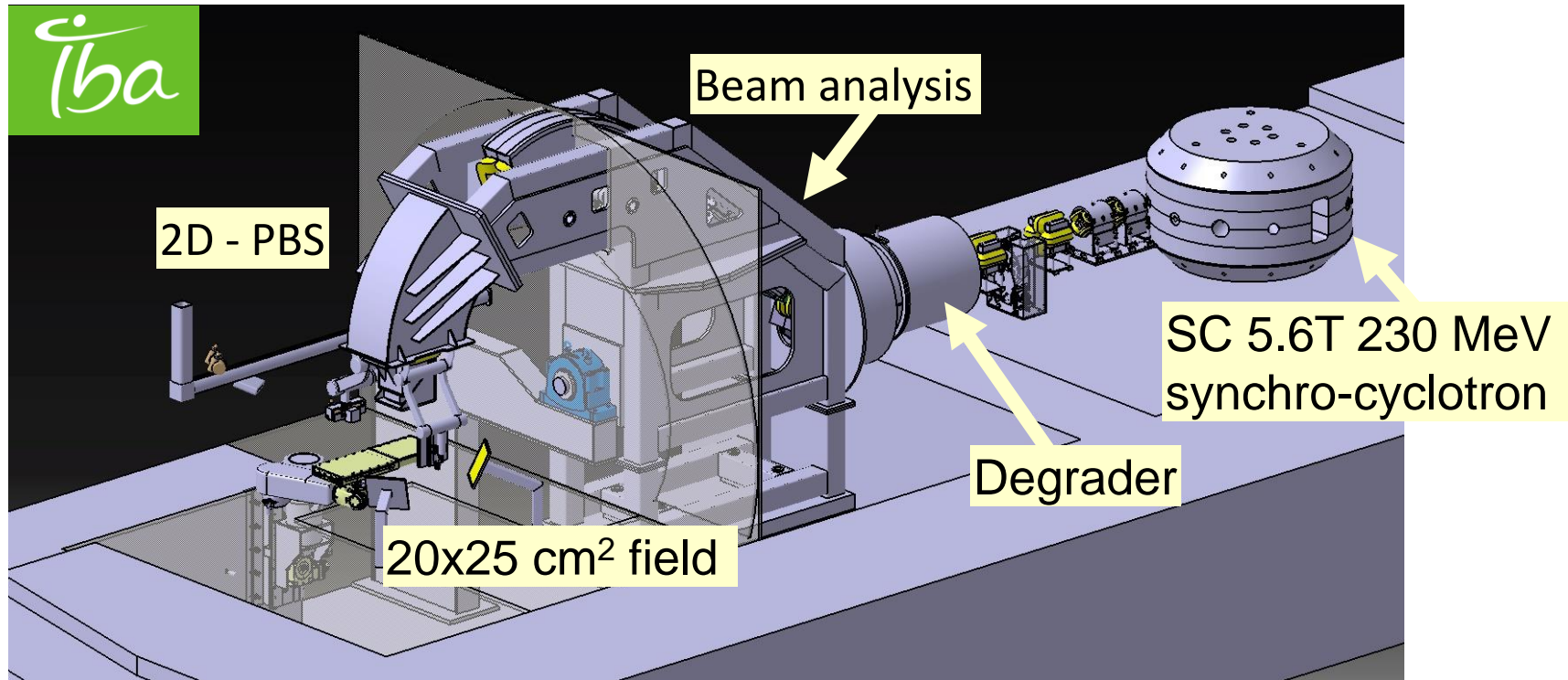
A modern synchrocyclotron for medical application – IBA S2C2

→ at the same energy synchrocyclotrons can be build more compact and with lower cost than sector cyclotrons; however, the achievable current is significantly lower

energy	230 MeV
current	130 nA
dimensions	Ø2.5 m x 2 m
weight	< 50 t
extraction radius	0.45 m
s.c. coil strength	5.6 Tesla
RF frequency	90...60 MHz
repetition rate	1 kHz

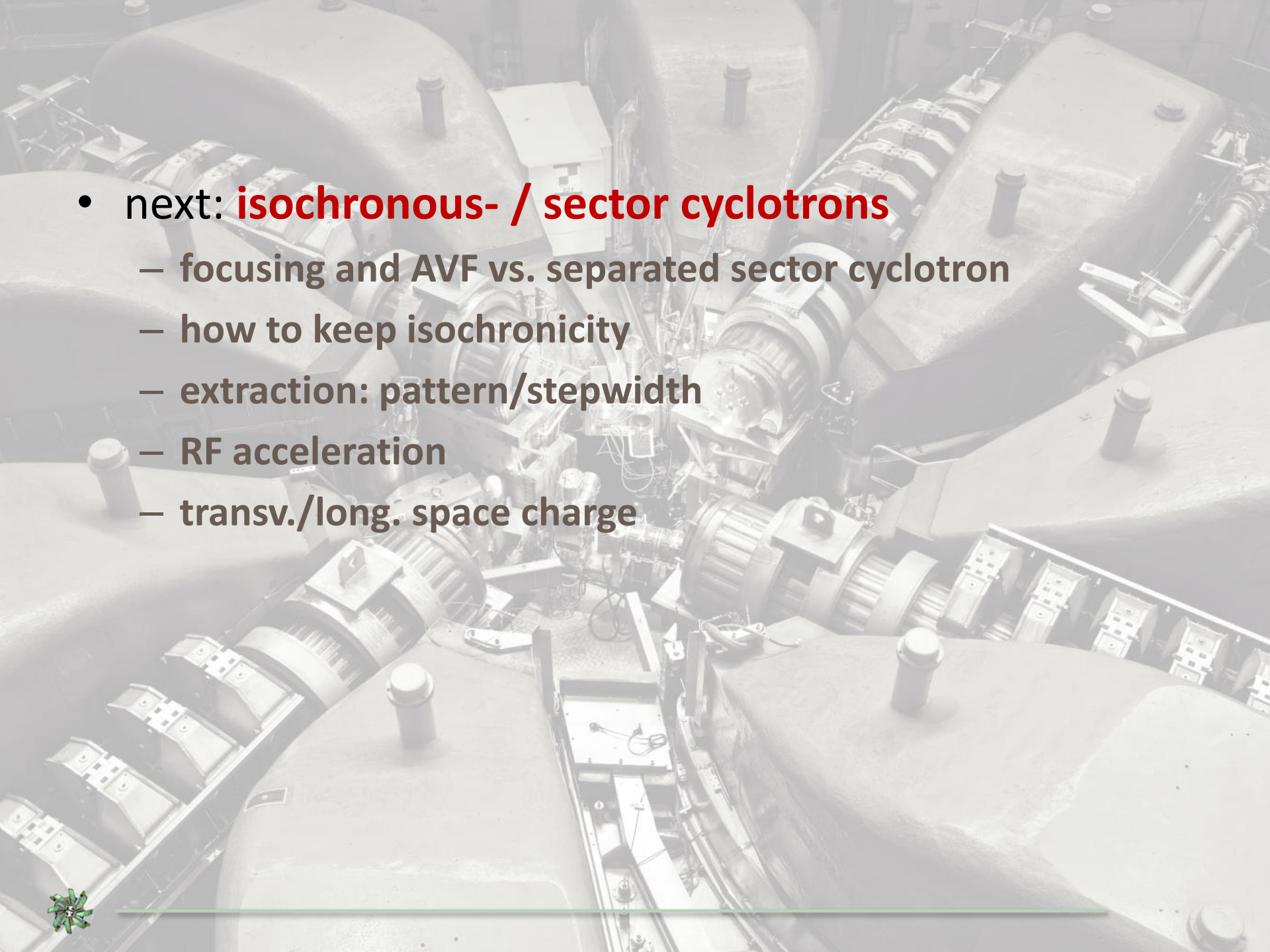


compact treatment facility using the high field synchro-cyclotron



- required area: 24x13.5m² (is small)
- 2-dim pencil beam scanning



- 
- next: **isochronous- / sector cyclotrons**
 - focusing and AVF vs. separated sector cyclotron
 - how to keep isochronicity
 - extraction: pattern/stepwidth
 - RF acceleration
 - transv./long. space charge

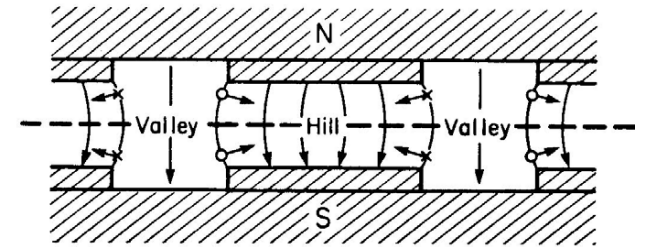


focusing in sector cyclotrons

hill / valley variation of magnetic field (Thomas focusing) makes it possible to design cyclotrons for higher energies

Flutter factor:

$$F^2 = \frac{\overline{B_z^2} - \overline{B_z}^2}{\overline{B_z}^2}$$

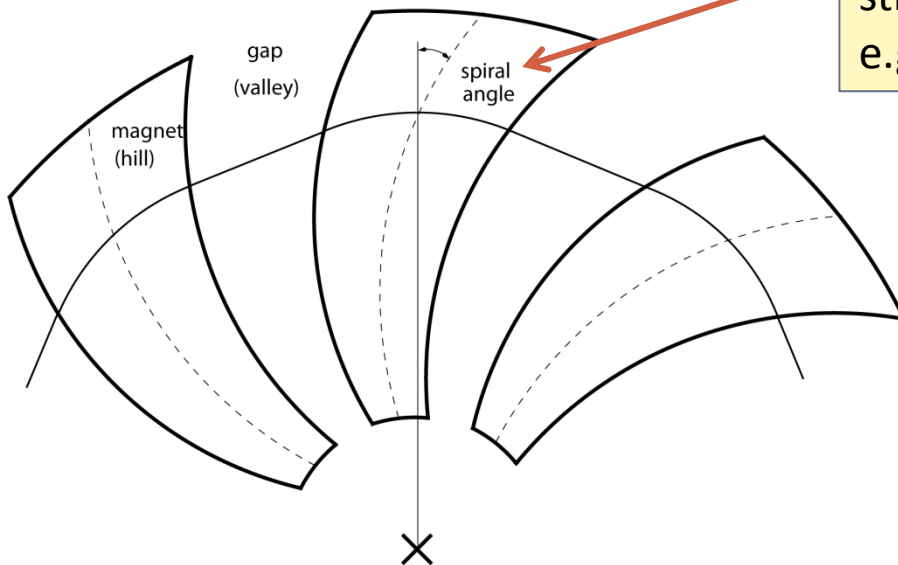


[illustration of focusing at edges]

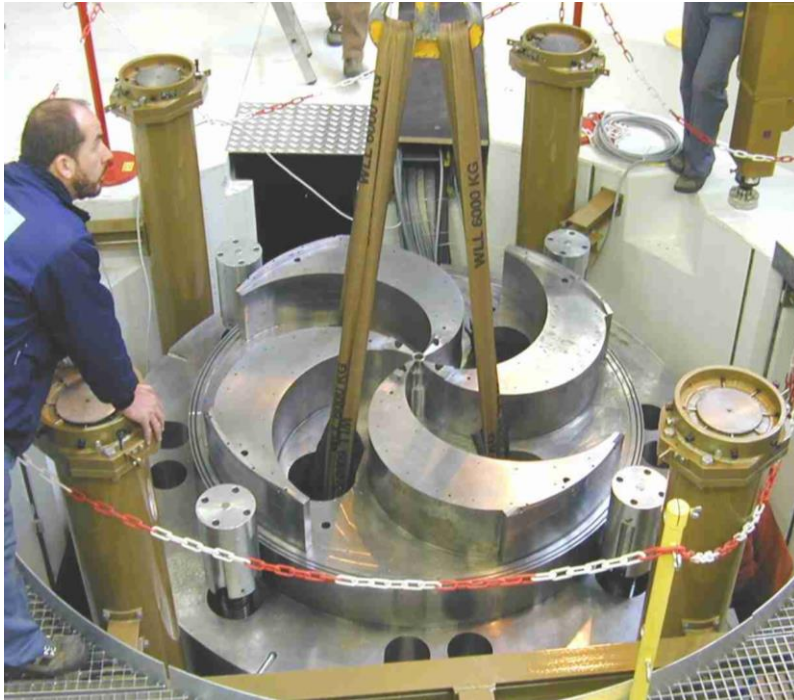
with flutter and additional spiral angle of bending field:

$$\nu_z^2 = -\frac{R}{B_z} \frac{dB_z}{dR} + F^2(1 + 2 \tan^2 \delta)$$

strong term
e.g.: $\delta=27^\circ$: $2 \tan^2 \delta = 1.0$

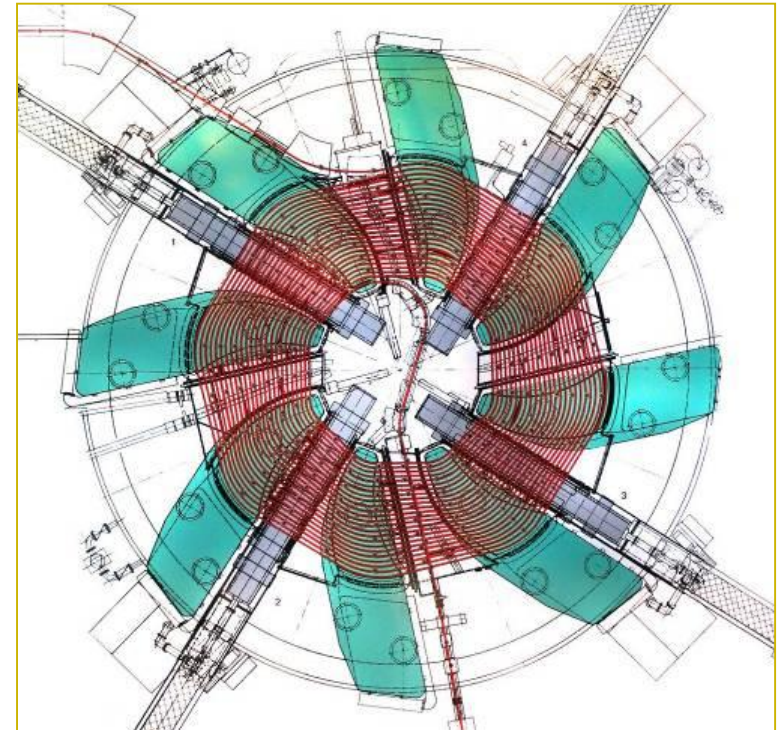


Azimuthally Varying Field vs. Separated Sector Cyclotrons



PSI/Varian comet: 250MeV sc. medical cyclotron

- **AVF = single pole with shaping**
- often **spiral poles** used
- **internal source** possible
- **D-type RF electrodes**, rel. low energy gain
- **compact**, cost effective
- depicted Varian cyclotron: 80% extraction efficiency; **not suited for high power**



PSI Ring cyclotron

- **modular layout**, larger cyclotrons possible, sector magnets, box resonators, stronger focusing, injection/extraction in straight sections
- **external injection** required, i.e. pre-accelerator
- **box-resonators** (high voltage gain)
- high **extraction efficiency** possible:
e.g. PSI: 99.98% = $(1 - 2 \times 10^{-4})$

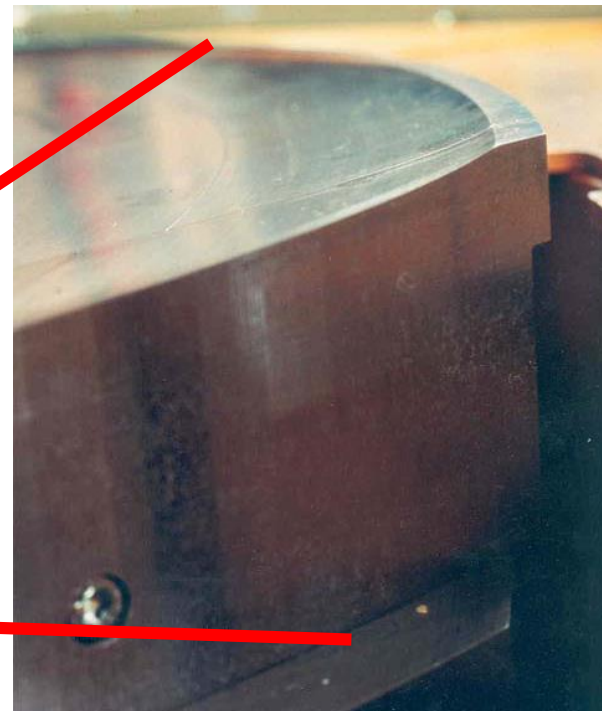


three methods to raise the average magnetic field with γ

remember:

$$\begin{aligned} \text{rev.time} &: R \propto \beta \\ \text{momentum} &: BR \propto \beta\gamma \\ \text{thus} &: B \propto \gamma \end{aligned}$$

- 1.) broader hills (poles) with radius
- 2.) **decrease pole gap with radius**
- 3.) s.c. coil arrangement to enhance field at large radius (in addition to iron dominated field)

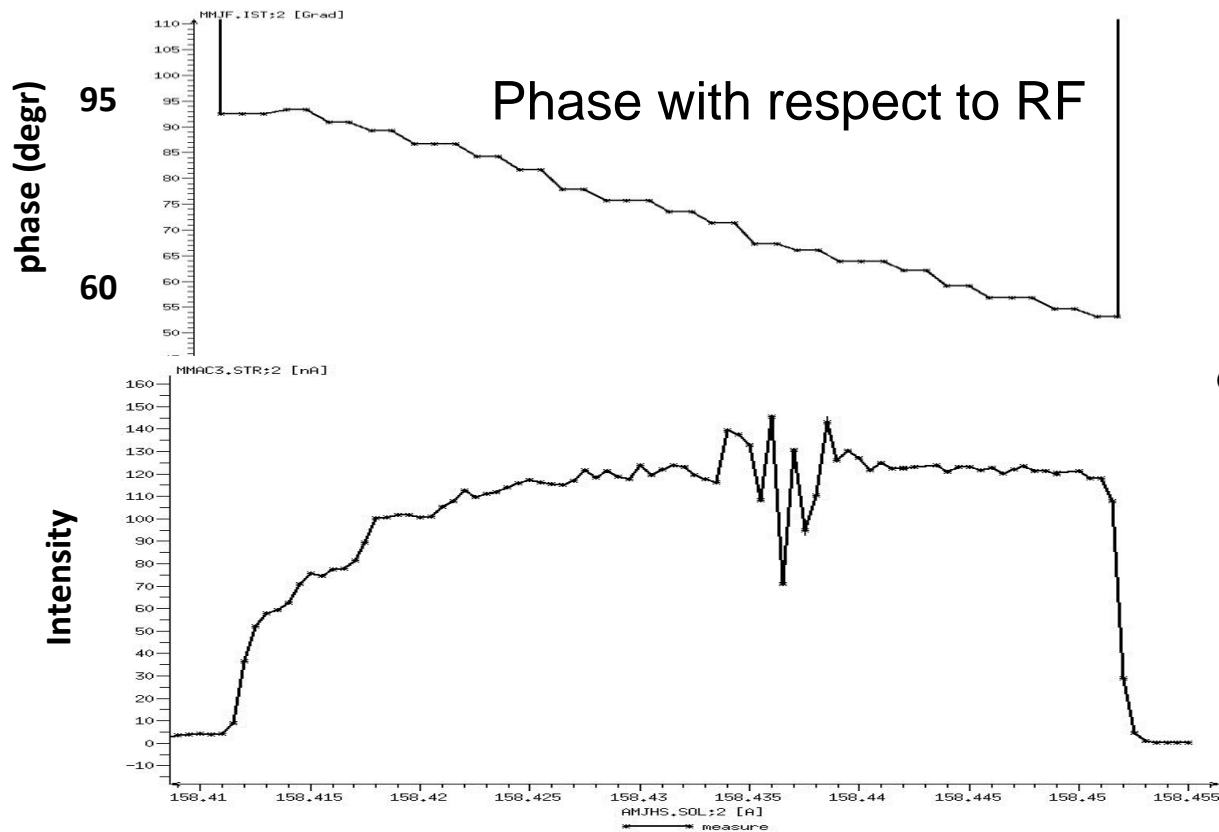


(photo: S. Zaremba, IBA)



field stability is critical for isochronicity

example: medical Comet cyclotron (PSI)



$$\Delta\phi_{RF} \propto n_{\text{turn}} \frac{\Delta B}{B}$$

e.g. : $n_{\text{turn}} = 600$

158.41

158.43

158.45

Current in main coil (A)



derivation of (relativistic) turn separation in a cyclotron

starting point: bending strength

→ compute total log.differential

→ use field index $k = R/B \cdot dB/dR$

$$BR = \sqrt{\gamma^2 - 1} \frac{m_0 c}{e}$$

$$\frac{dB}{B} + \frac{dR}{R} = \frac{\gamma d\gamma}{\gamma^2 - 1}$$

$$\frac{dR}{d\gamma} = \frac{\gamma R}{\gamma^2 - 1} \frac{1}{1 + k}$$

radius change per turn

$$\frac{dR}{dn_t} = \frac{dR}{d\gamma} \frac{d\gamma}{dn_t} \quad [U_t = \text{energy gain per turn}]$$

$$= \frac{U_t}{m_0 c^2} \frac{\gamma R}{(\gamma^2 - 1)(1 + k)} \quad \left. \vphantom{\frac{U_t}{m_0 c^2}} \right\} \text{isochronicity not conserved (last turns)}$$

$$= \frac{U_t}{m_0 c^2} \frac{R}{(\gamma^2 - 1)\gamma} \quad \left. \vphantom{\frac{U_t}{m_0 c^2}} \right\} \text{isochronicity conserved (general scaling)}$$



turn separation - discussion

for clean extraction a large stepwidth (turn separation) is of utmost importance; in the PSI Ring most efforts were directed towards maximizing the turn separation

general scaling at extraction:

$$\Delta R(R_{\text{extr}}) = \frac{U_t}{m_0 c^2} \frac{R_{\text{extr}}}{(\gamma^2 - 1)\gamma}$$

desirable:

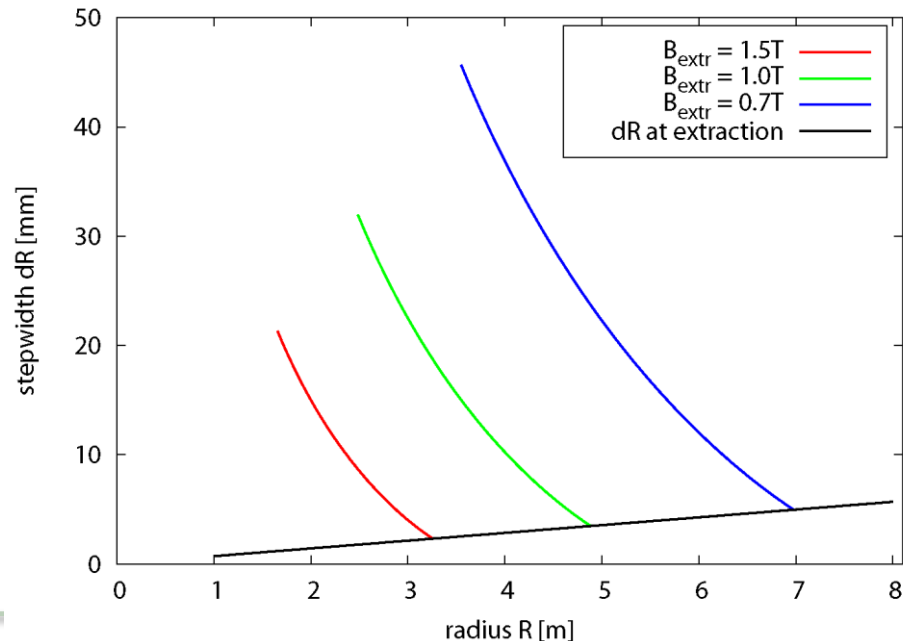
- limited energy (< 1GeV)
- large radius R_{extr}
- high energy gain U_t

scaling during acceleration:

$$\frac{dR}{dn_t} \approx \frac{U_t}{m_0 c^2} \frac{R}{\beta^2} \rightarrow \Delta R(R) \propto \frac{1}{R}$$

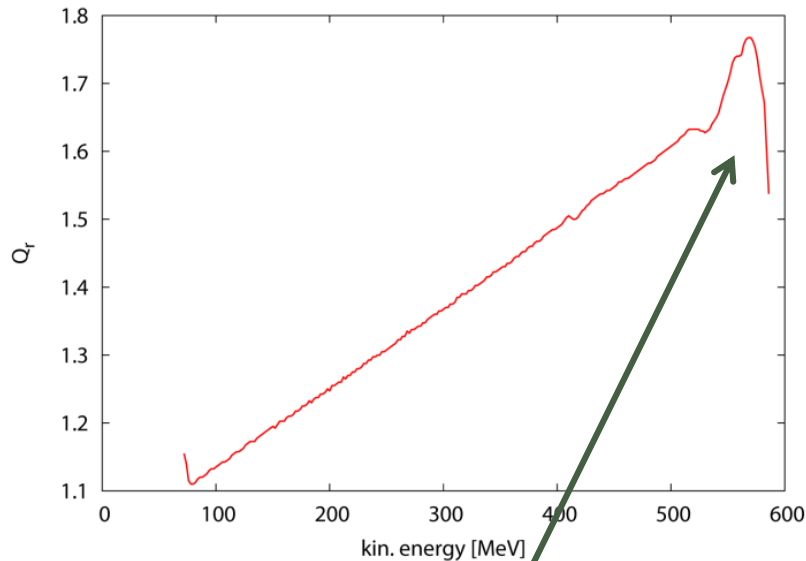
illustration:

stepwidth vs. radius in cyclotrons of different sizes but same energy;
100MeV inj \rightarrow 800MeV extr



extraction with off-center orbits

betatron oscillations around the “closed orbit” can be used to increase the radial stepwidth by a factor 3 !

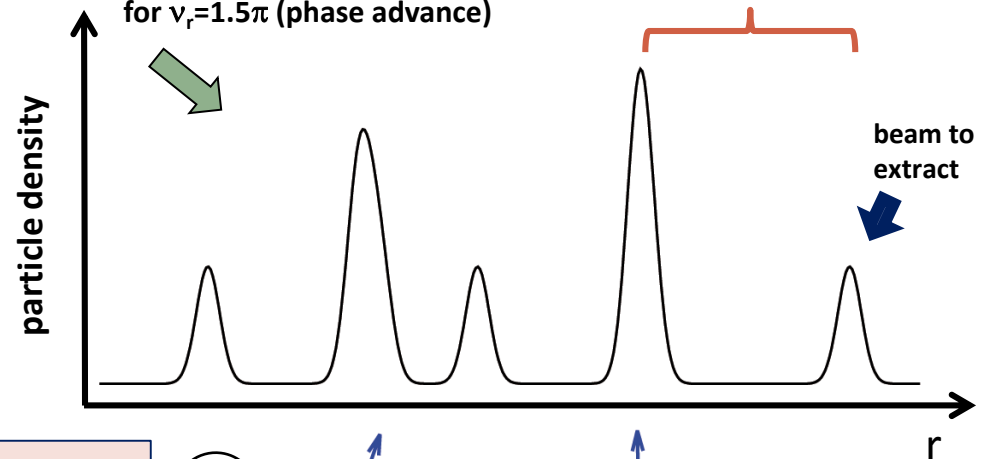


radial tune vs. energy (PSI Ring)
typically $\nu_r \approx \gamma$ during acceleration;
but decrease in outer fringe field

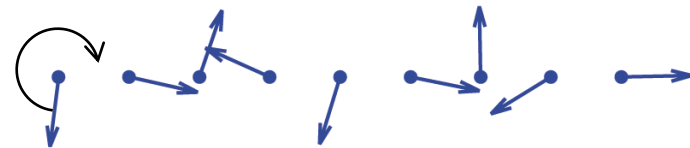
without orbit oscillations: stepwidth from E_k -gain (PSI: 6mm)



with orbit oscillations: extraction gap; up to 3 x stepwidth possible for $\nu_r = 1.5\pi$ (phase advance)

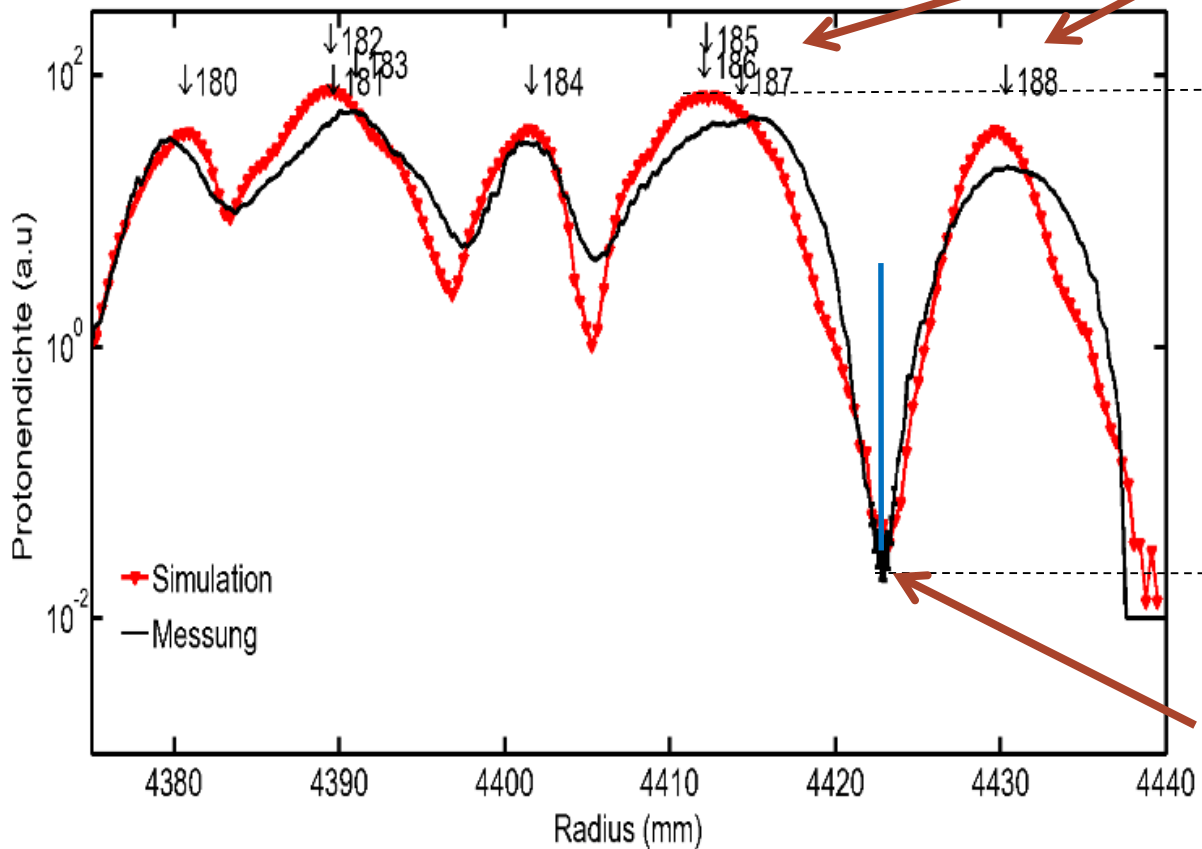


phase vector of orbit oscillations (r, r')



extraction profile measured at PSI Ring Cyclotron

red: tracking simulation [OPAL]
black: measurement



turn numbers
from simulation

dynamic range:
factor 2.000 in
particle density

position of extraction septum
 $d=50\mu\text{m}$

[Y.Bi et al]



longitudinal space charge

sector model (W.Joho, 1981):

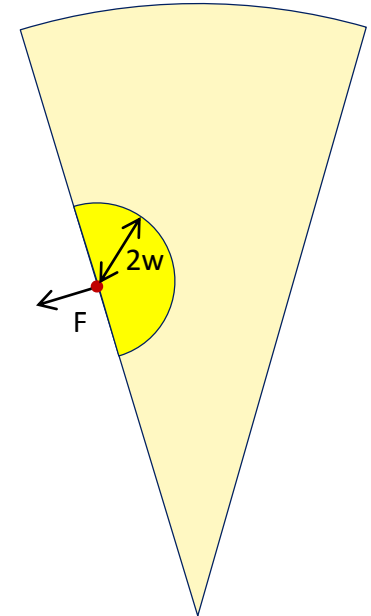
- accumulated energy spread transforms into transverse tails
- consider rotating uniform sectors of charge (overlapping turns)
- test particle “sees” only fraction of sector due to shielding of vacuum chamber with gap height $2w$

two factors are proportional to the number of turns:

- 1) the charge density in the sector
- 2) the time span the force acts

$$\Delta U_{sc} = \frac{8}{3} e I_p Z_0 \ln \left(4 \frac{w}{a} \right) \cdot \frac{n_{\max}^2}{\beta_{\max}} \approx 2.800 \Omega \cdot e I_p \cdot \frac{n_{\max}^2}{\beta_{\max}}$$

derivation see: [High Intensity Aspects of Cyclotrons, ECPM-2012, PSI](#)



in addition:

- 3) the inverse of turn separation at extraction: $\frac{1}{\Delta R_{\text{extr}}} \propto n_{\max}$

► thus the attainable current at constant losses scales as n_{\max}^{-3}

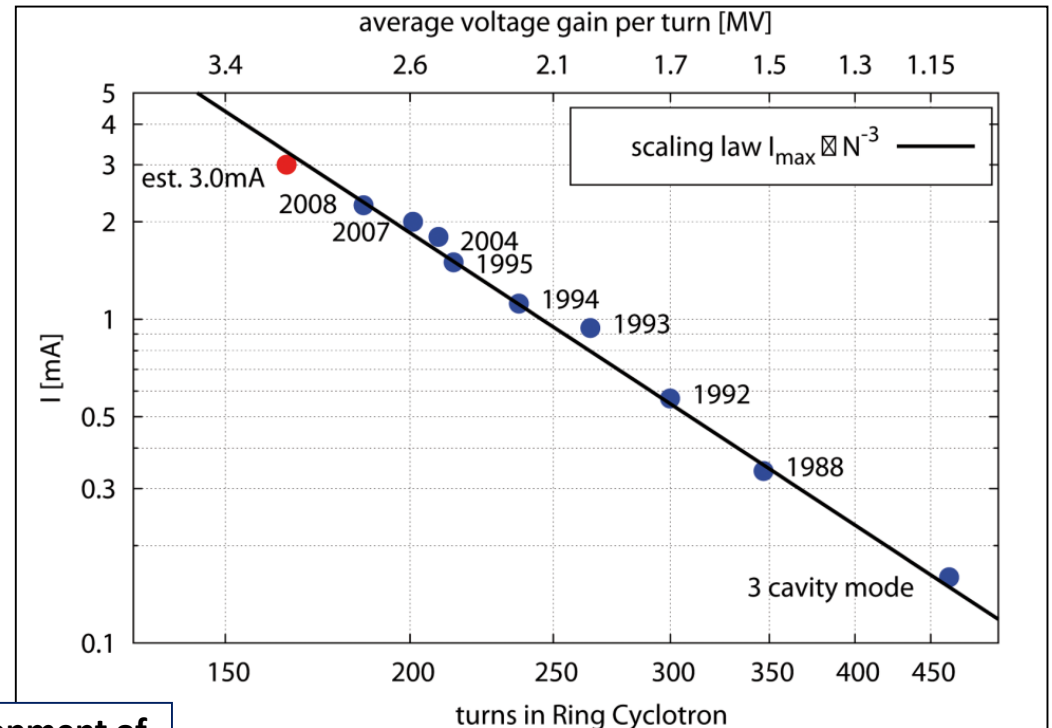


longitudinal space charge; evidence for third power law

- at PSI the maximum attainable current indeed scales with the third power of the turn number
- maximum energy gain per turn is of utmost importance in this type of high intensity cyclotron

→ with constant losses at the extraction electrode the maximum attainable current indeed scales as:

$$I_{\max} \propto n_t^{-3}$$



historical development of current and turn numbers in PSI Ring Cyclotron





next: **cyclotron examples**

- compact cyclotrons
- TRIUMF, RIKEN SRC, PSI-Comet, PSI-HIPA

compact cyclotrons for Isotope production



CYCLONE 30 (IBA) : H⁻ 15 à 30 MeV

Vertical setup



some cyclotrons

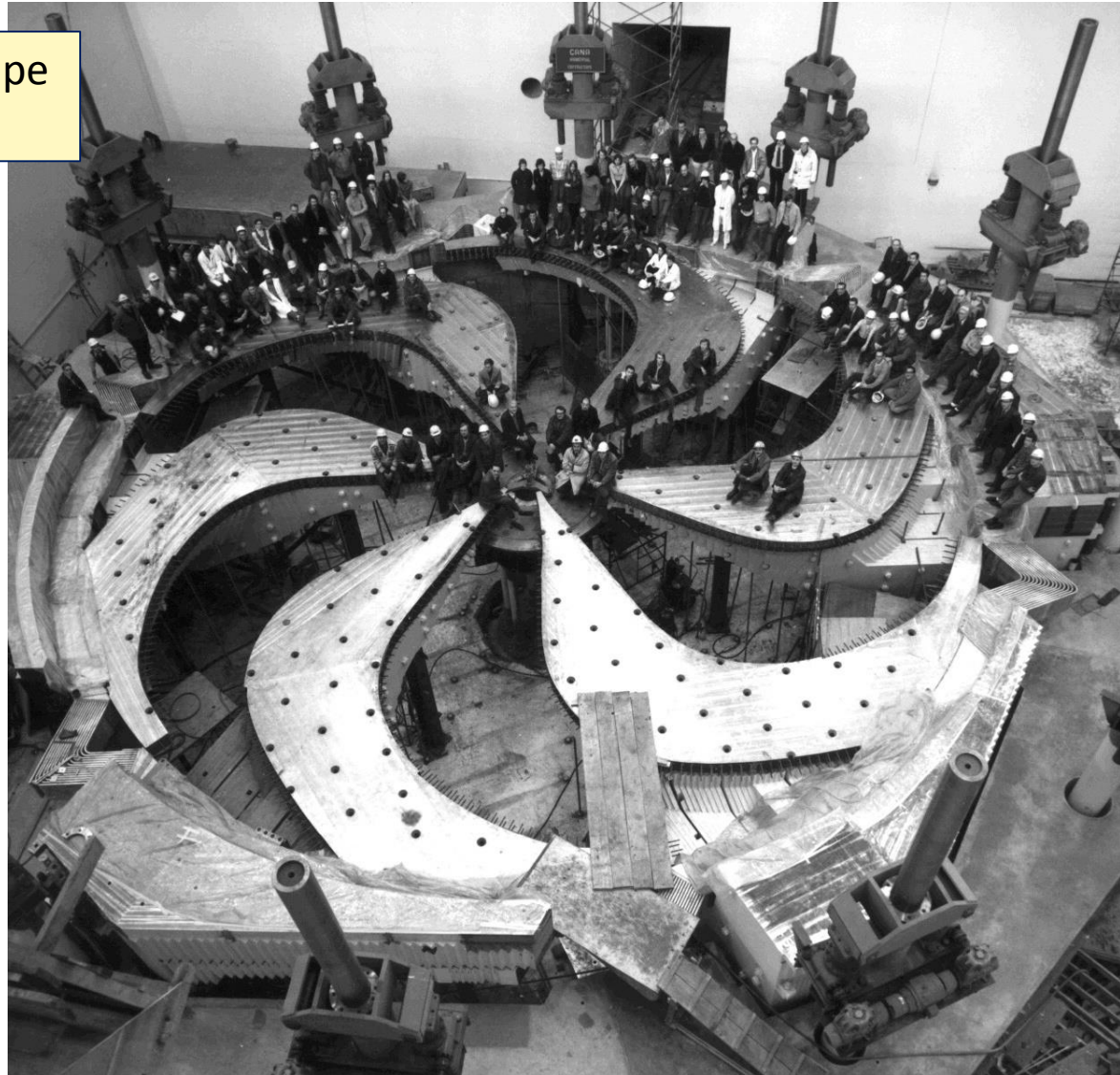
	TRIUMF	RIKEN SRC (supercond.)	PSI Ring	PSI medical (supercond.)
particles	H- → p	ions	p	p
K [MeV]	520	2600	592	250
magnets (poles)	(6)	6	8	(4)
peak field strength [T]	0.6	3.8	2.1	3.8
R_{inj}/R_{extr} [m]	0.25/3.8...7.9	3.6/5.4	2.4/4.5	-/0.8
P_{max} [kW]	110	1 (86Kr)	1300	0.25
extraction efficiency (tot. transmission)	0.9995 (0.70)	(0.63)	0.9998	0.80
extraction method	stripping foil	electrostatic deflector	electrostatic deflector	electrostatic deflector
comment	variable energy	ions, flexible	high intensity	compact



cyclotron examples: TRIUMF / Vancouver

photo: iron poles with spiral shape
($\delta_{\max}=70\text{deg}$)

- p, 520MeV, up to 110kW beam power
- diameter: 18m (largest n.c. cyclotron worldwide)
- extraction by stripping H^-
→ variable energy;
multiple extraction points possible

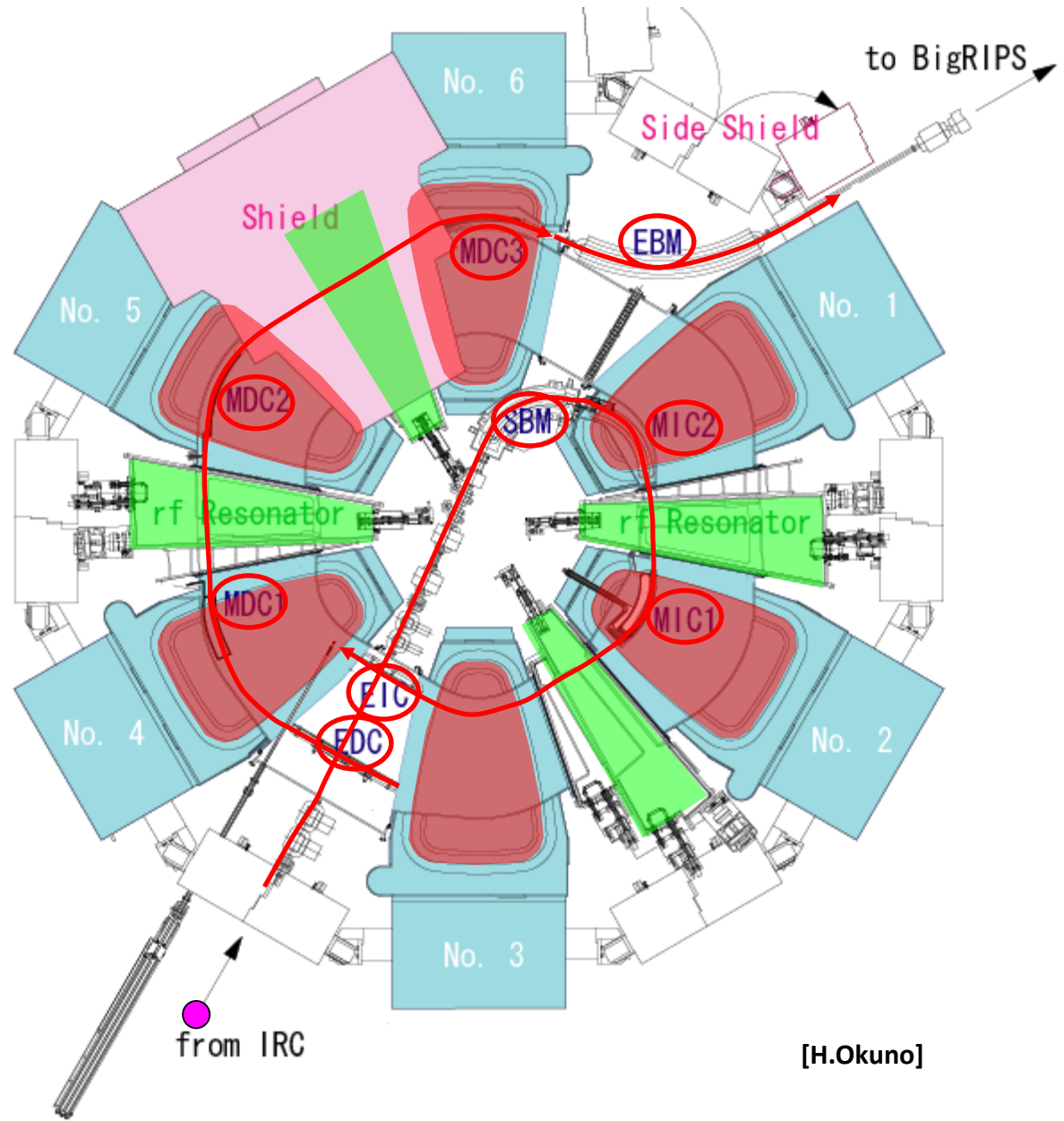


example: RIKEN (Jp) superconducting cyclotron

K = 2,600 MeV
Max. Field: 3.8T (235 MJ)
RF frequency: 18-38 MHz
Weight: 8,300 tons
Diameter: 19m
Height: 8m

superconducting
Sector Magnets :6
RF Resonator :4
Injection elements.
Extraction elements.

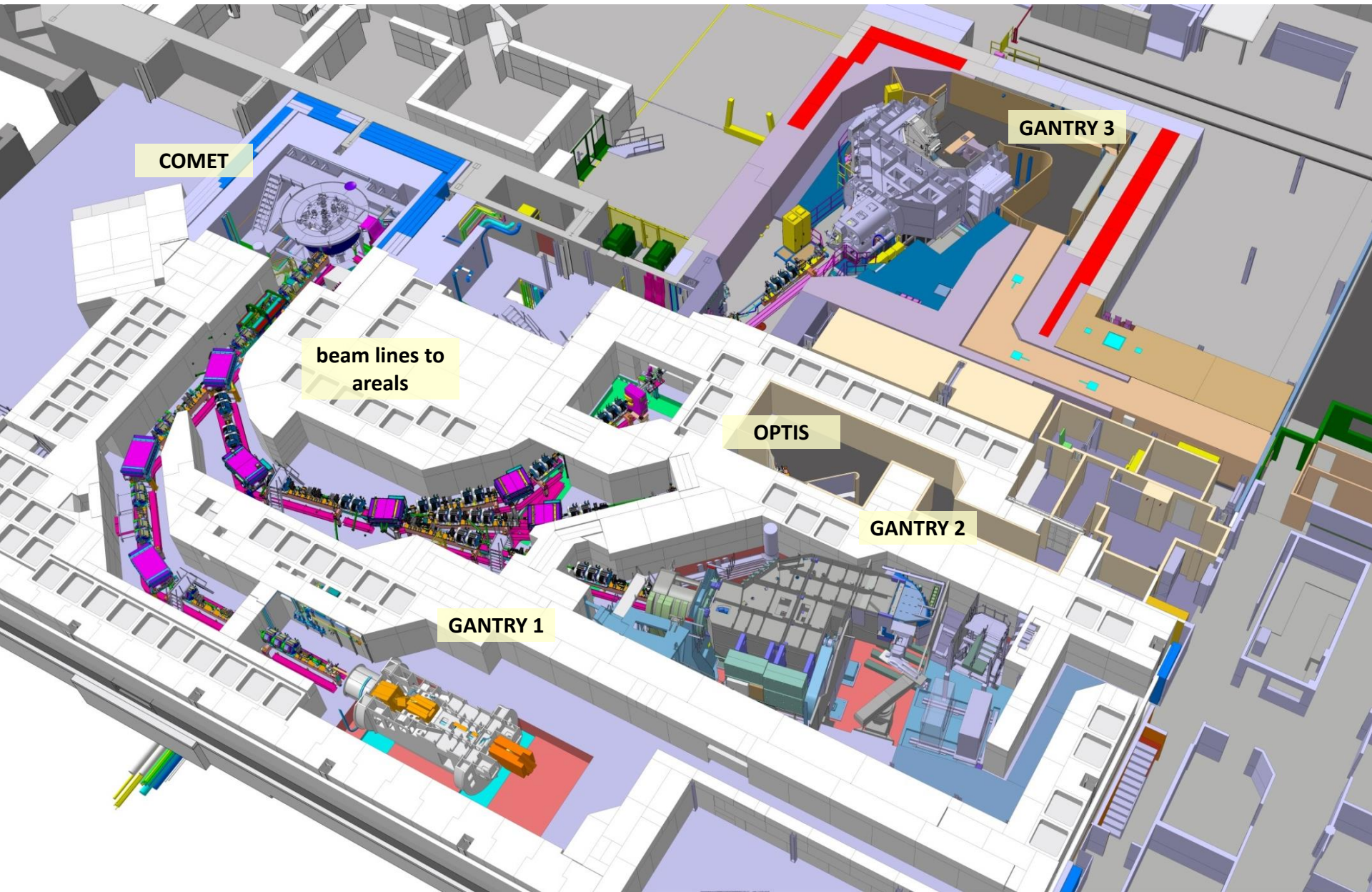
utilization:
**broad spectrum of
ions up to Uranium**



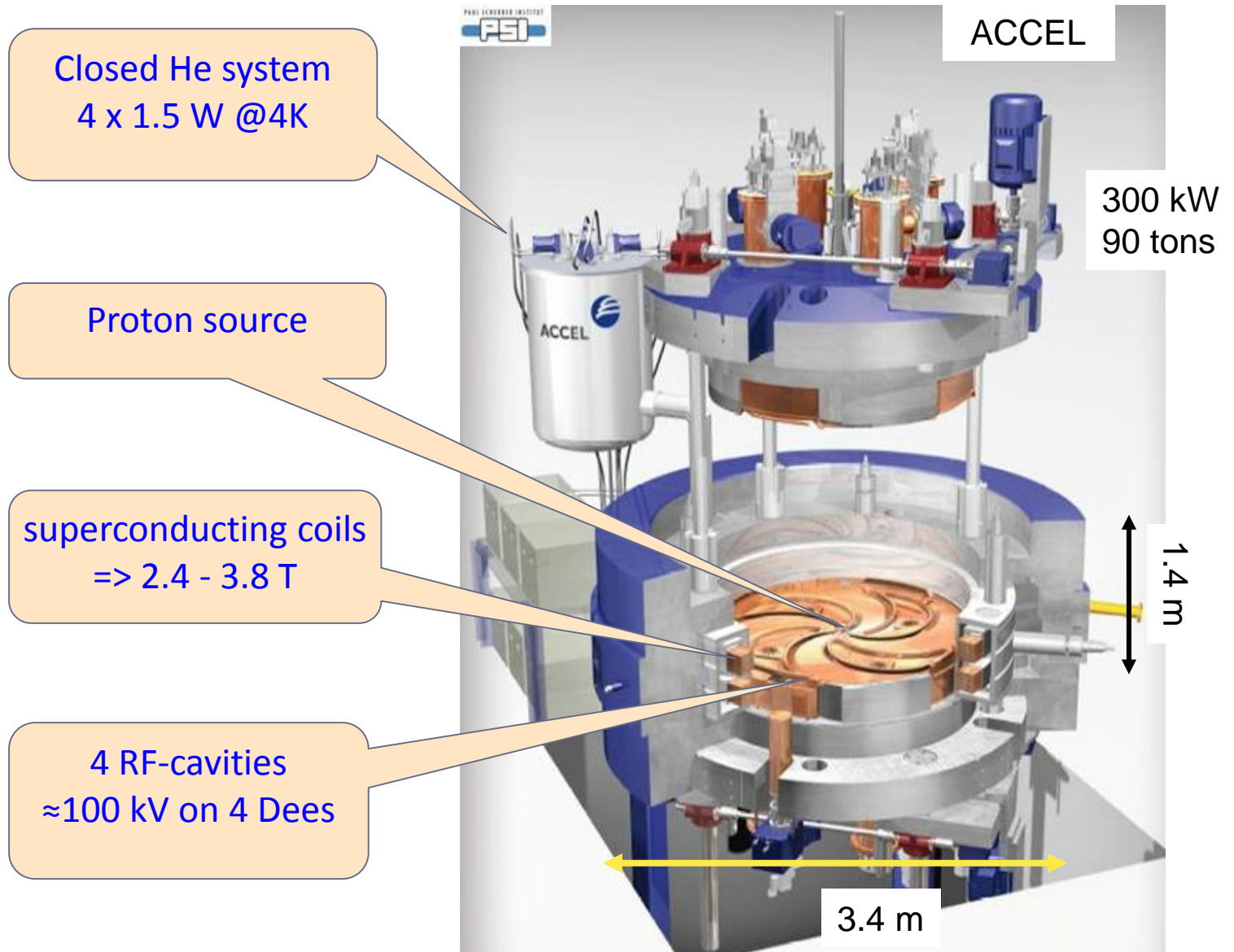
RIKEN SRC in the vault



PSI Proton Therapy Facility

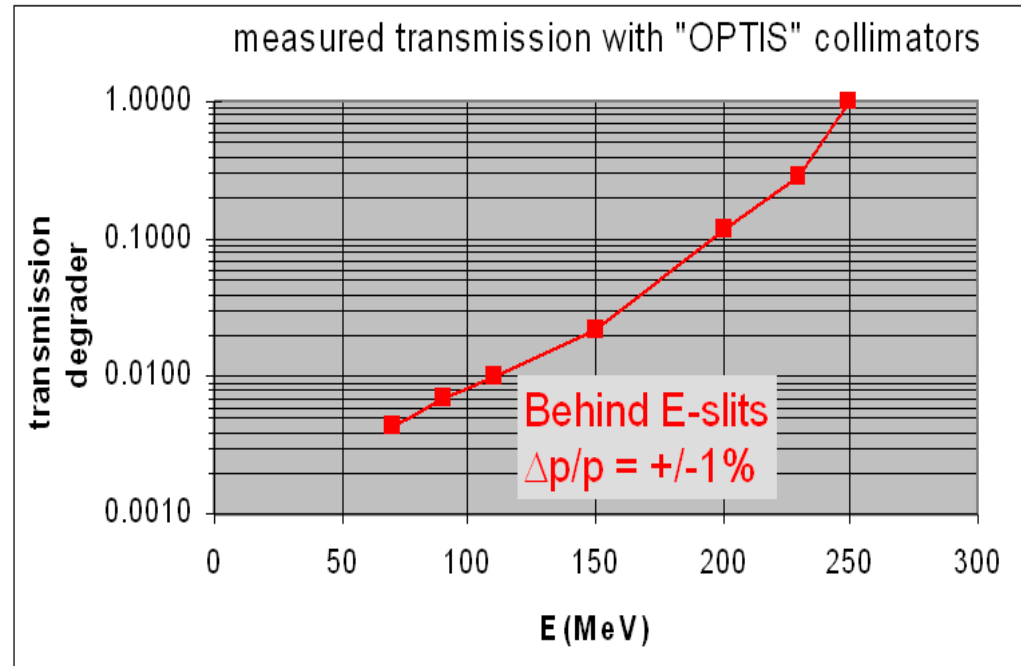


250 MeV proton cyclotron (ACCEL/Varian)

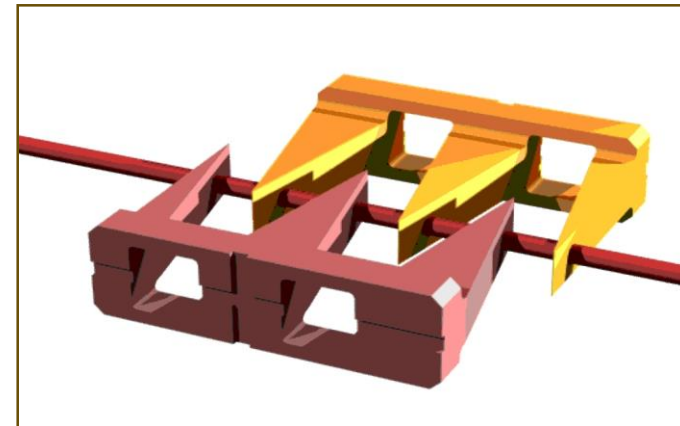


Cyclotron needs degrader :

- cyclotron has fixed energy; need **degrader** for energies down to 70MeV
- collimation after degrader to keep emittance → lose intensity with degrader



degrader: (carbon wedges in vacuum)
and laminated beam line magnets for
fast energy changes < 80 ms / step



examples: PSI High Intensity Proton Accelerator

Ring Cyclotron 590 MeV
2.4mA / 1.4MW
diameter: 15m

meson production targets

SINQ
spallation source

proton therapie center
[250MeV sc. cyclotron]

dimensions:
120 x 220m²



Outlook: Cyclotrons II & FFA

- **cyclotron subsystems**
extraction schemes, RF systems/resonators, magnets, vacuum issues, instrumentation
- **FFA = Fixed Focus Alternating Gradient Accelerators**
motivation & applications, scaling FFA's, non-scaling and linear FFA, FFA subsystems
- **discussion**
classification of circular accelerators, cyclotron vs. FFAG, Pro's and Con's of cyclotrons for different applications

