



ALICE

Inclusive p_T spectrum of b-jets in p-Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV

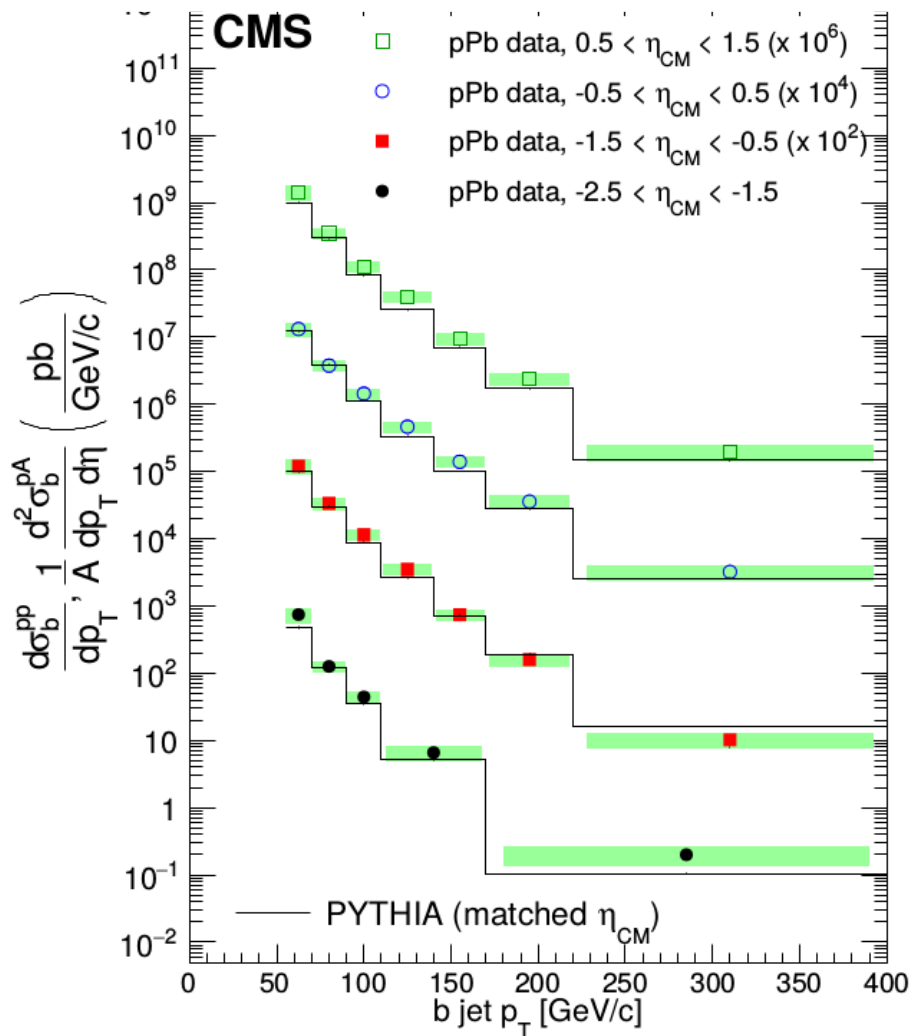
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CMS Results (pPb, 5.02 TeV, full jets 2018)



Properties of b-quark:

- **large mass** ($4.62 \text{ GeV}/c^2$) → it can be created only in initial hard scatterings. Its production rate can be calculated from pQCD
- **long lifetime** → it survives through the whole evolution of QGP

ALICE wants to study b-jets at **lower momenta**

[The CMS Collaboration - "Transverse momentum spectra of inclusive b jets in pPb collisions at $\sqrt{s_{NN}} = 5.02 \text{ TeV}$ ", CERN-PH-EP/2013-037, 2018/10/09]

Event selection

- Minimum bias trigger (V0 scintillator arrays)
- $|z_{\text{vtx}}| < 10$ cm
- Pileup rejection

After event selection we have $6 \cdot 10^8$ minimum bias events

Track selection

- Hybrid tracks
- $|\eta_{\text{track}}| < 0.9$

Jet selection:

- Charged anti- k_T , $R=0.4$
- $p_{T, \text{constituent}} > 0.15$ GeV/c
- p_T recombination scheme
- $|\eta_{\text{jet}}| < 0.9 - R < 0.5$

Background density correction:

Two leading k_T jets are excluded

$$p_{T, \text{charged jet}}^{\text{corrected}} = p_{T, \text{charged jet}}^{\text{RAW}} - \rho \cdot A_{\text{jet}}$$
$$\rho = \frac{A_{\text{physical jets}}}{A_{\text{all jets}}} \times \text{median}_{k_T \text{ physical jet}} \left\{ \frac{p_{T, \text{jet}}^{\text{ch, raw}}}{A_{\text{jet}}} \right\}$$

b-jet candidate selection:

- 3 prong SV is made out of jet constituents
- In each event we consider the most displaced SV

Discrimination variables:

- 1) Significance of the distance between primary and secondary vertices

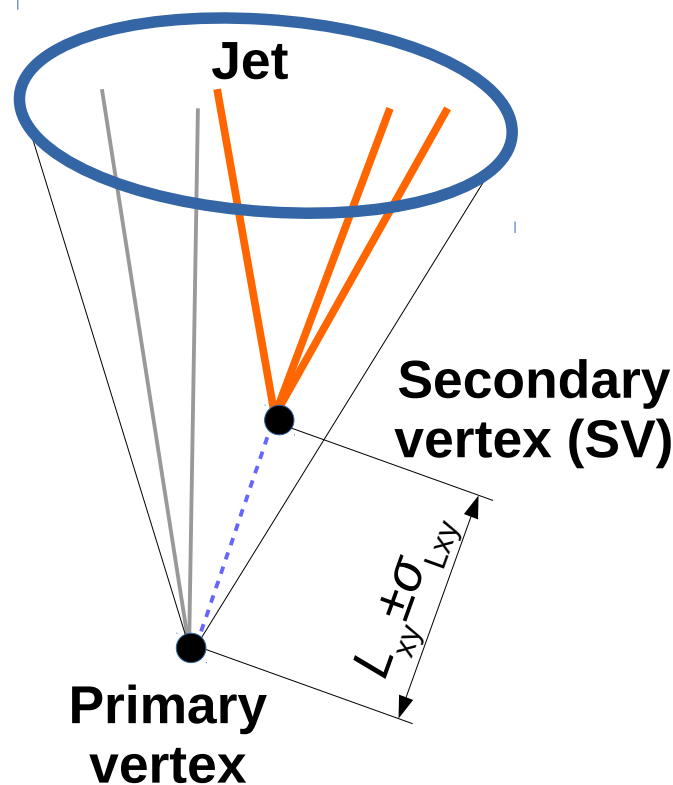
$$SL_{xy} = L_{xy} / \sigma_{L_{xy}} > 5 \text{ to } 9$$

- 2) Dispersion of the SV $\sigma_{SV} < 0.02 \text{ to } 0.05 \text{ cm}$

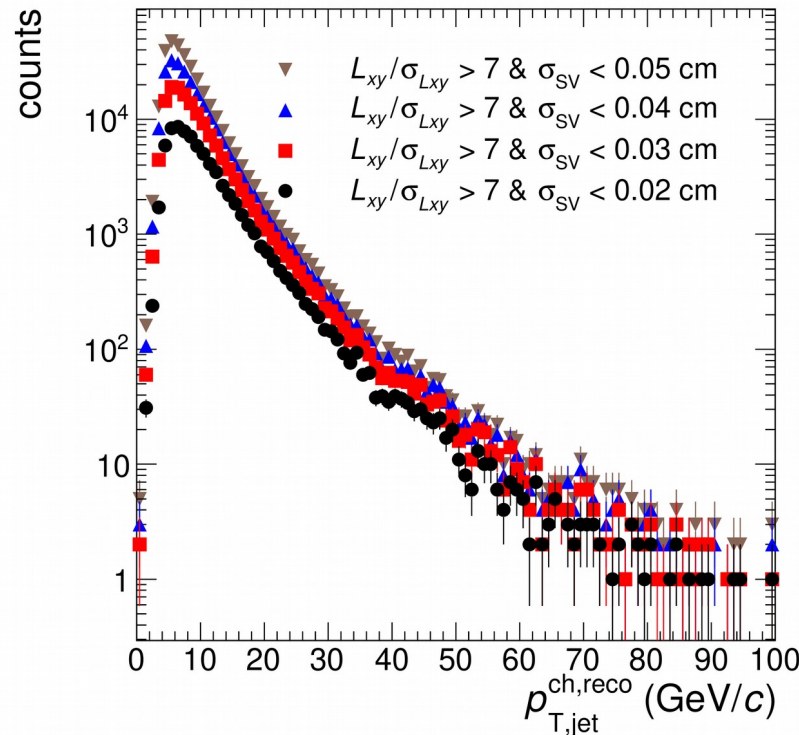
$$\sigma_{SV} = \sqrt{\sum_{i=1}^3 d_i^2}$$

d_i – distance of the closest approach (DCA) of i-th prong to SV

- 3) Invariant mass in SV (*reserved for purity estimation*)



Raw p_T spectrum of b-jets



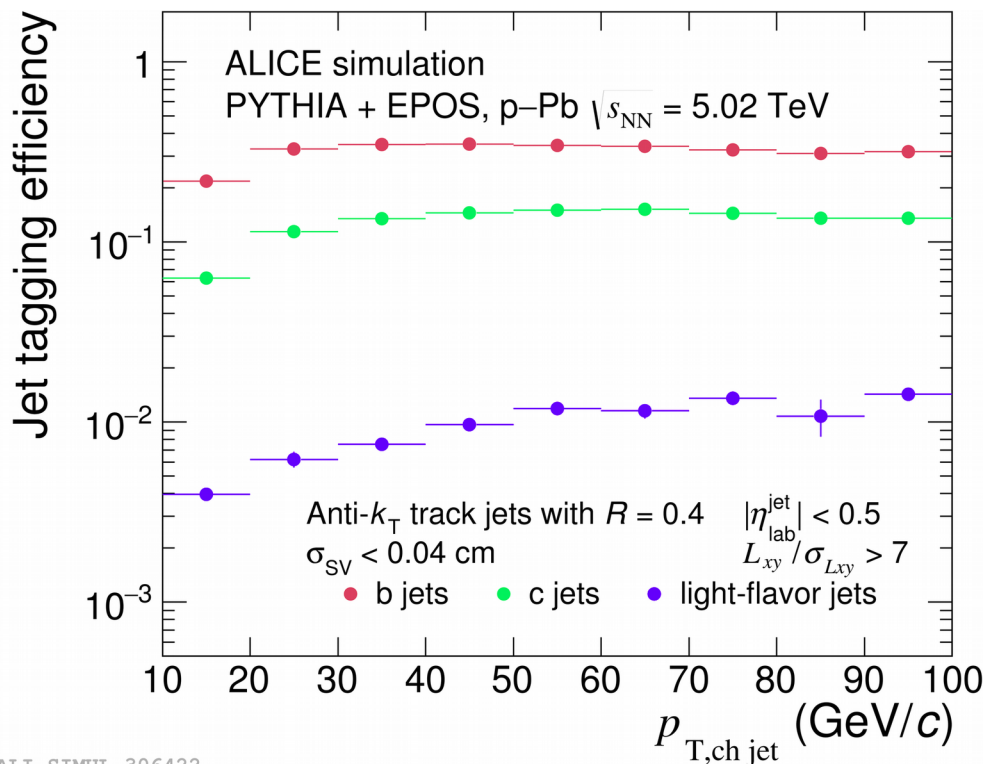
- Selected sample of b-jet candidates contains b, c and LF-jets.
- To get RAW transverse momentum spectra of b-jets, the spectrum of b-jets candidates needs to be corrected:

$$\frac{dN_{b\text{-jet}}^{\text{primary}}}{dp_{T, \text{jet ch}}} = \frac{dN_{b\text{-jet candidates}}^{\text{raw}}}{dp_{T, \text{jet ch}}} \times \frac{P_b}{\varepsilon_b}$$

P_b – purity of the b-jet candidates

ε_b – efficiency of the b-jet selection after applying cuts

Efficiency of SV tagging



ALI-SIMUL-306422

$$\epsilon_b \approx 35 \%, \epsilon_c \approx 11 \%, \epsilon_{LF} \approx 1 \%$$

b-jet tagging efficiency is estimated from EPOS+PYTHIA detector-level simulation:

$$\epsilon_b = \frac{N_{b-jets}^{\text{selected}}}{N_{b-jets}^{\text{all}}}$$

N_{b-jets}^{all} – the number of b-jets without any constraint on presence and parameters of SV

$N_{b-jets}^{\text{selected}}$ – the number of b-jets that were reconstructed when applying cuts on b-jets candidates

b-jet purity from data driven template fit method



- The data driven method is based on representation of the measured distribution of invariant mass of SV as a linear combination of MC templates:

$$\begin{cases} n_{SV} = P_b \cdot T_b + P_c \cdot T_c + P_{LF} \cdot T_{LF} \\ 1 = P_b + P_c + P_{LF} \end{cases}$$

n_{SV} – measured SV invariant mass distribution in given jet- p_T bin

T_b, T_c, T_{LF} – MC template spectra for each jet flavor

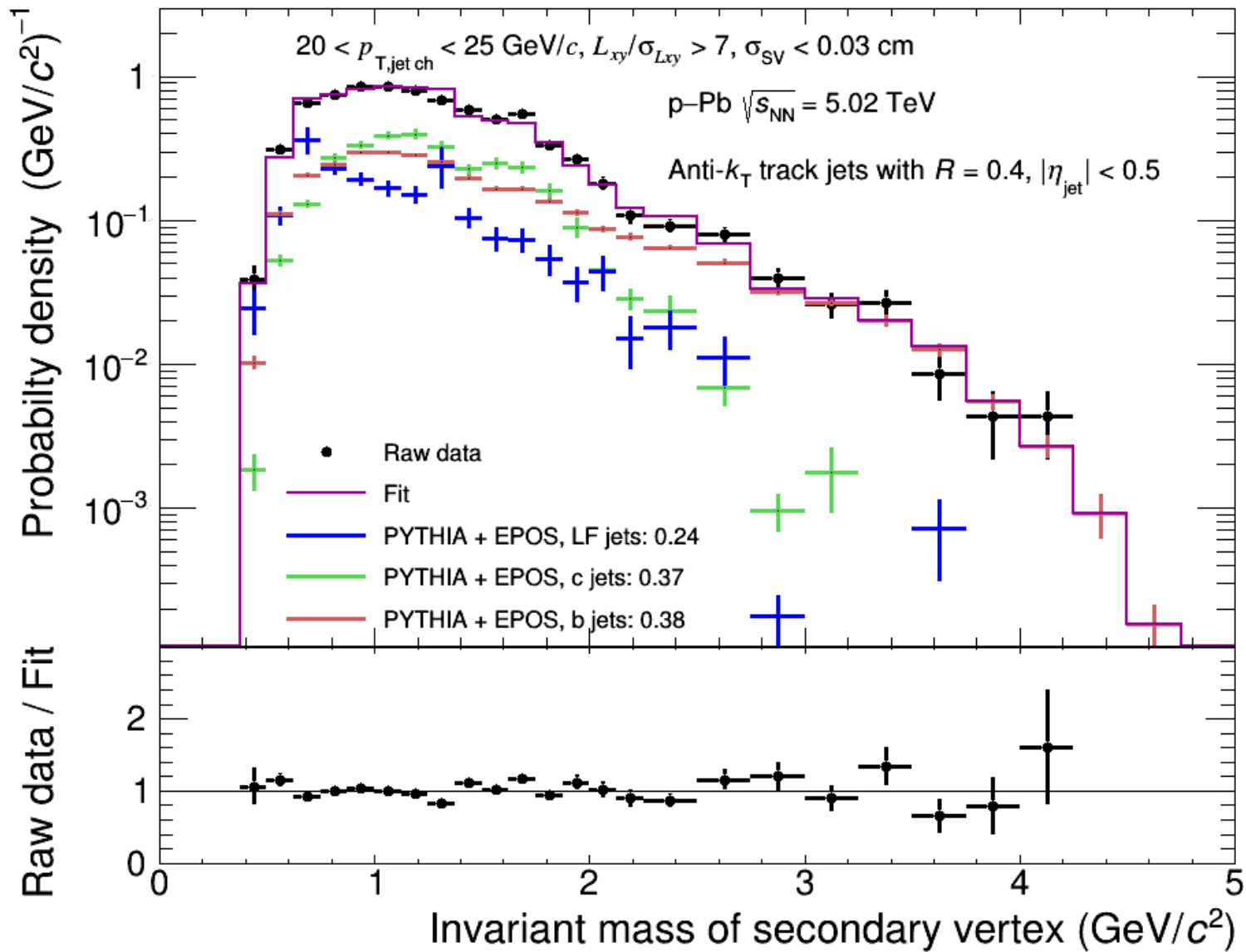
P_b, P_c, P_{LF} – purity for each jet flavor

- Purity is evaluated in 5 GeV/c wide $p_{T, \text{jet ch}}$ bins
- **TMinuit library** was used to fit MC templates to measured distribution

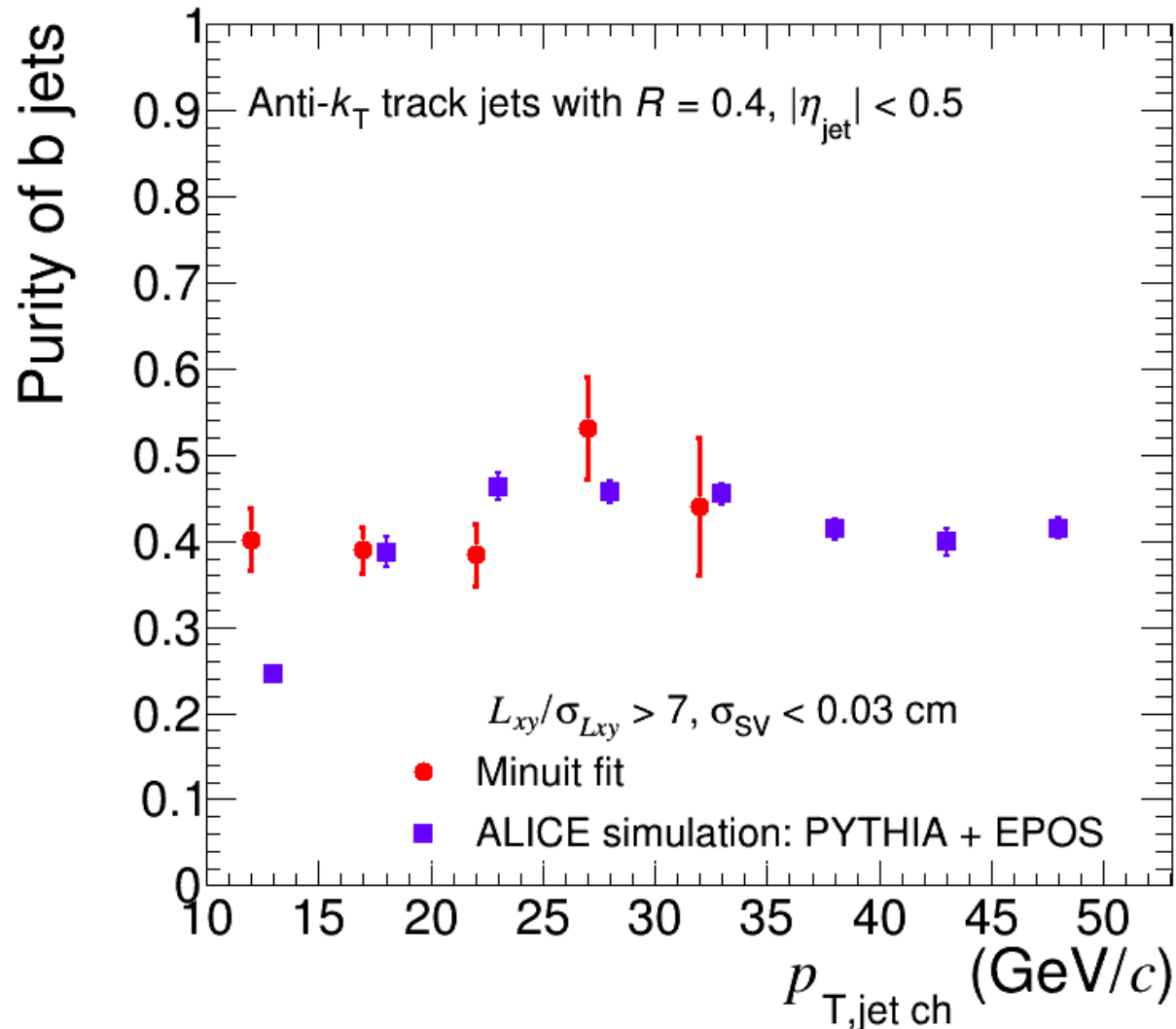
$$\chi^2 = \sum_{i=1}^{nbis} \frac{(n_{SV,i} - P_b \cdot T_{b,i} - P_c \cdot T_{c,i} - P_{LF} \cdot T_{LF,i})^2}{\sigma_{n_{SV,i}}^2 + (\sigma_{T_{b,i}} \cdot P_b)^2 + (\sigma_{T_{c,i}} \cdot P_c)^2 + (\sigma_{T_{LF,i}} \cdot P_{LF})^2}$$

$\sigma_{n_{SV}}, \sigma_{T_b}, \sigma_{T_c}, \sigma_{T_{LF}}$ – statistical error for each jet flavor

Results of TMinuit fitting



Data-driven method results



- Bad convergence for $p_{T,\text{jet ch}} > 35$ GeV/c → larger statistics for MC and Real Data is required

POWbc method - based on b and c -jets spectra calculated by Next-to-Leading Order (NLO) POWHEG generator:

- 1) Fold generated b and c jet spectra (**particle level**) with jet response matrix which accounts for momentum smearing due to local background fluctuations and instrumental effects
- 2) Purity estimate:

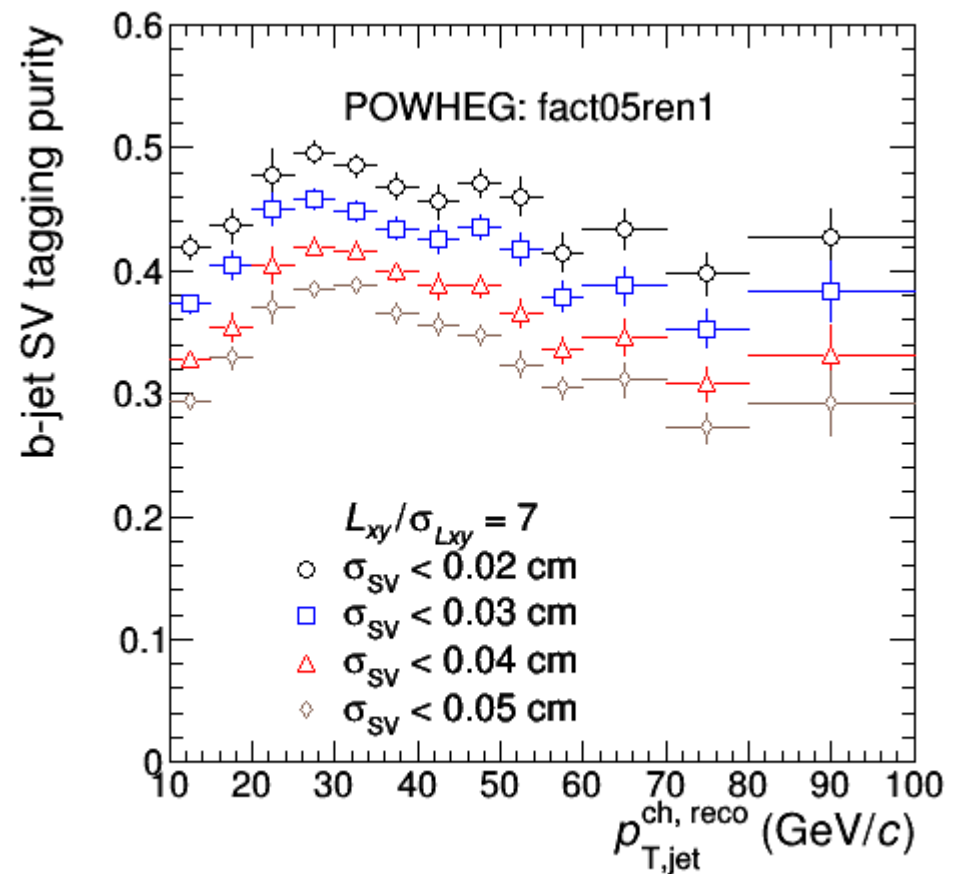
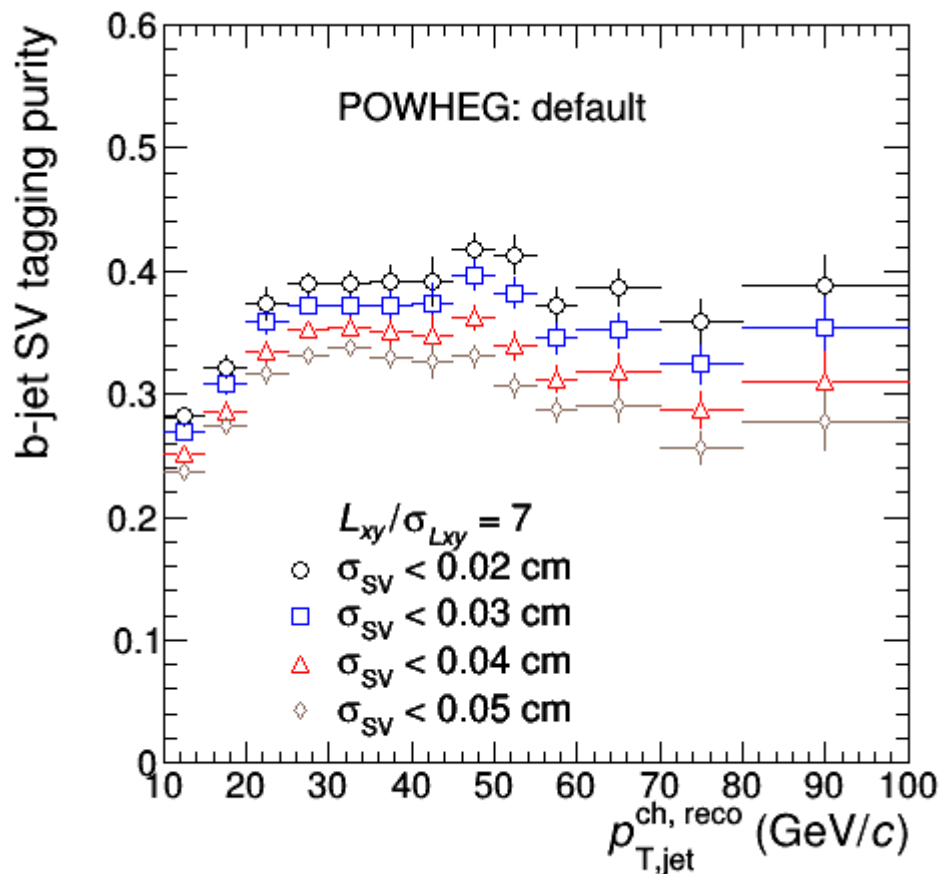
$$P_b = \frac{N_b \varepsilon_b}{N_b \varepsilon_b + N_c \varepsilon_c + N_{LF} \varepsilon_{LF}}$$

N_b, N_c – folded POWHEG p_T spectrum of b and c -jets

$N_{LF} = \text{RAW } p_T \text{ spectrum of inclusive jets} - N_b \varepsilon_b - N_c \varepsilon_c$

$\varepsilon_b, \varepsilon_c, \varepsilon_{LF}$ – efficiency of SV tagging for b , c and LF -jets for corresponding SL_{xy} and σ_{SV}

POWbc method results

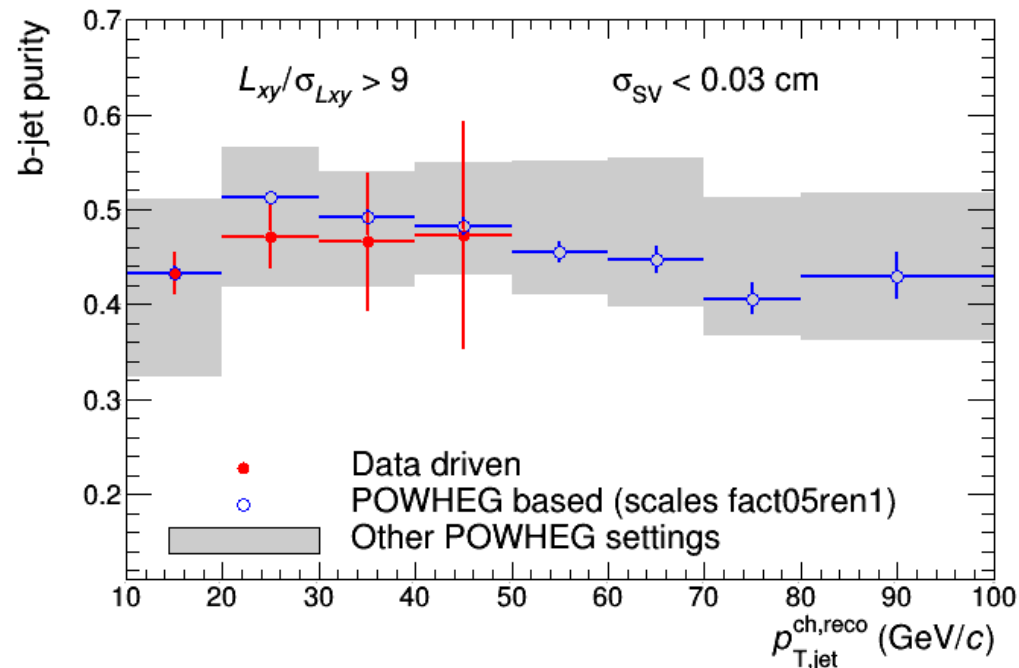
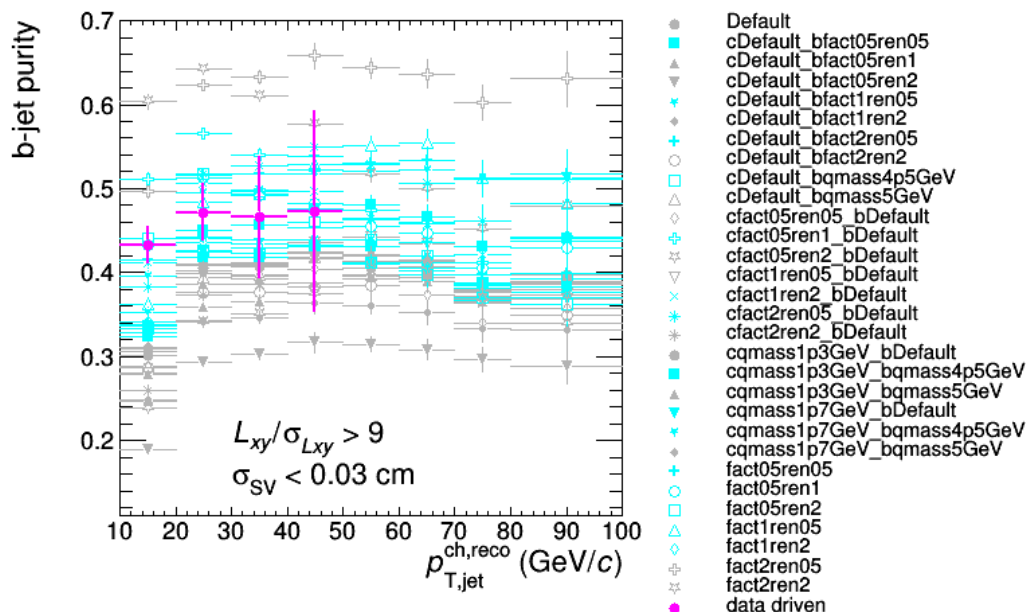


- Problem with a significant scale uncertainty from POWHEG
- Results strongly depends on POWHEG settings → need to choose the optimal one

Hybrid method for purity estimation



Hybrid method: different POWHEG settings were tested against the template fit results to find plausible POWHEG settings (regularization and renormalization scale)

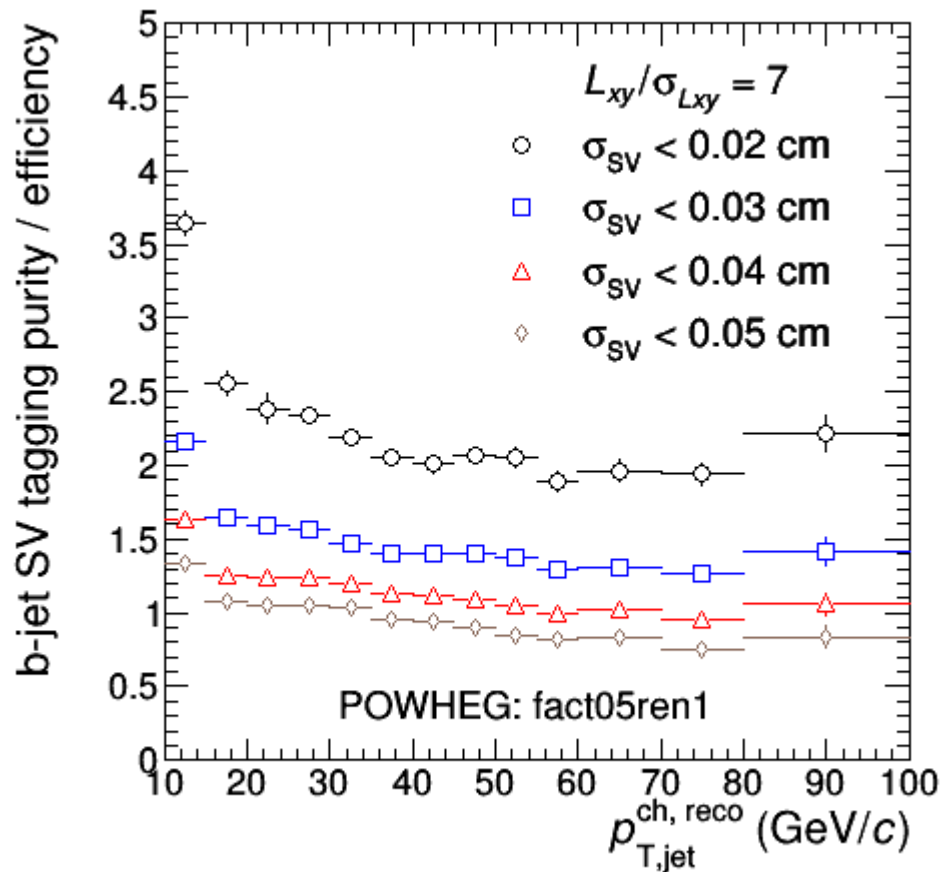


χ^2 compatibility of POWHEG with data was

$$\chi^2 = \sum_{bins} \frac{(f_b^{POWbc} - f_b^{data-driven})^2}{\sigma_{POWbc}^2 + \sigma_{data-driven}^2}$$

- χ^2 was computed for all tagging settings (SL_{xy} and σ_{SV})
- Only settings for which $\chi^2 / \text{n.d.f.} < 10$ were used

Purity to efficiency correction



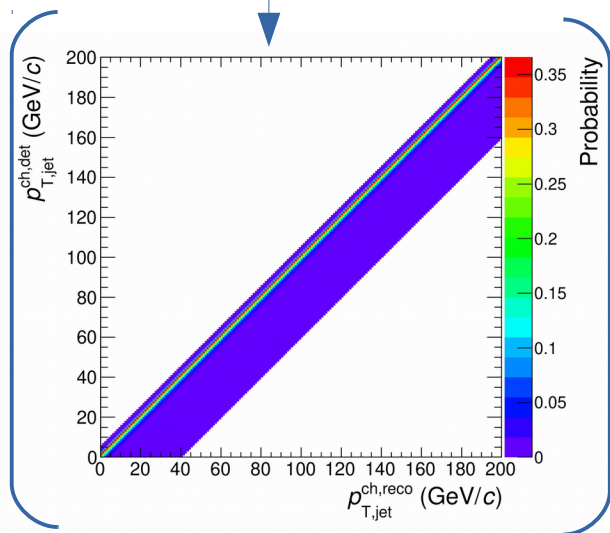
Unfolding procedure

δp_T matrix represents smearing of p_T spectrum due to local fluctuations in background density

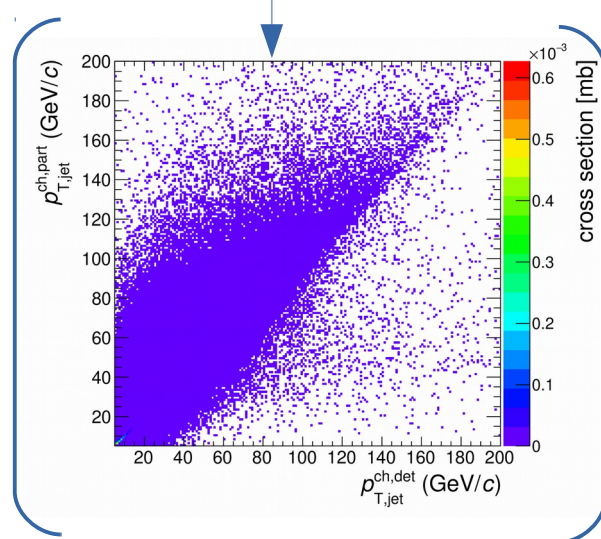
Instrumental matrix represents detector effects

b-jets p_T spectrum on a particle level this is our goal

Measured b-jet p_T spectrum corrected for SV efficiency and purity



\times



$$\times \left[N_{T, jet}^{ch, particle level} \right] = \left[N_{T, jet}^{ch, detector level} \right]$$

The problem of searching of b-jet p_T spectrum on a particle level can be solved with unfolding technics

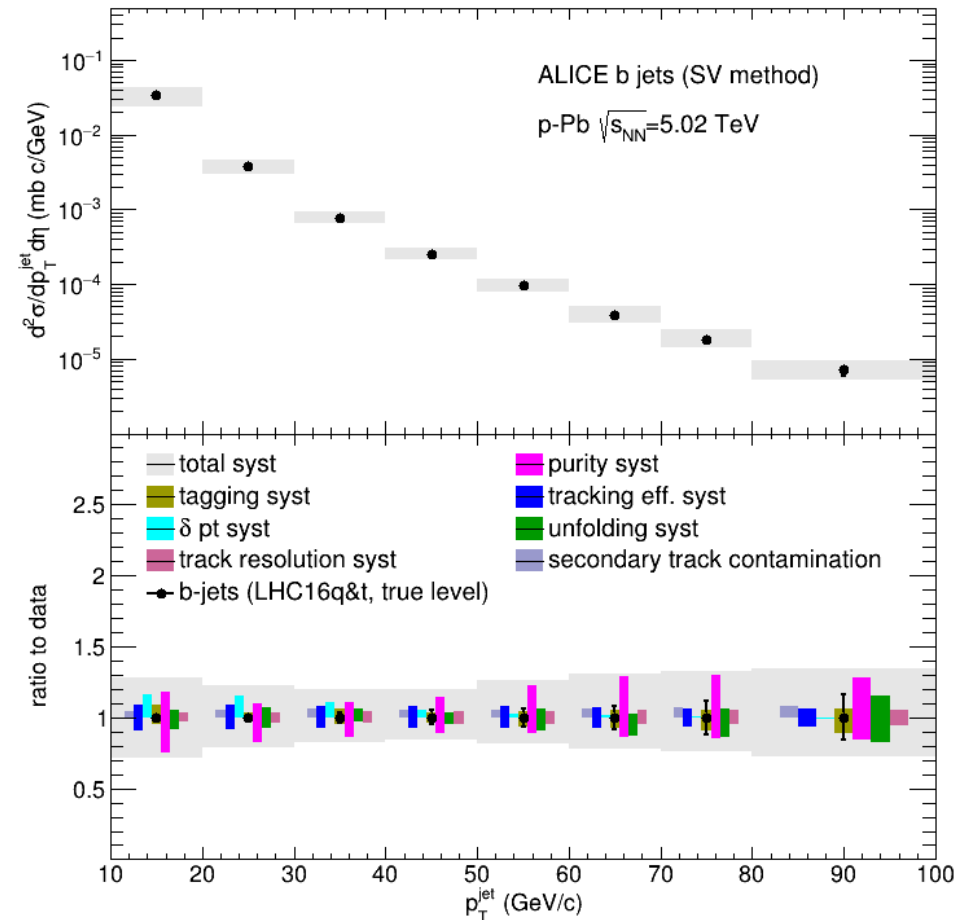
Systematic uncertainties



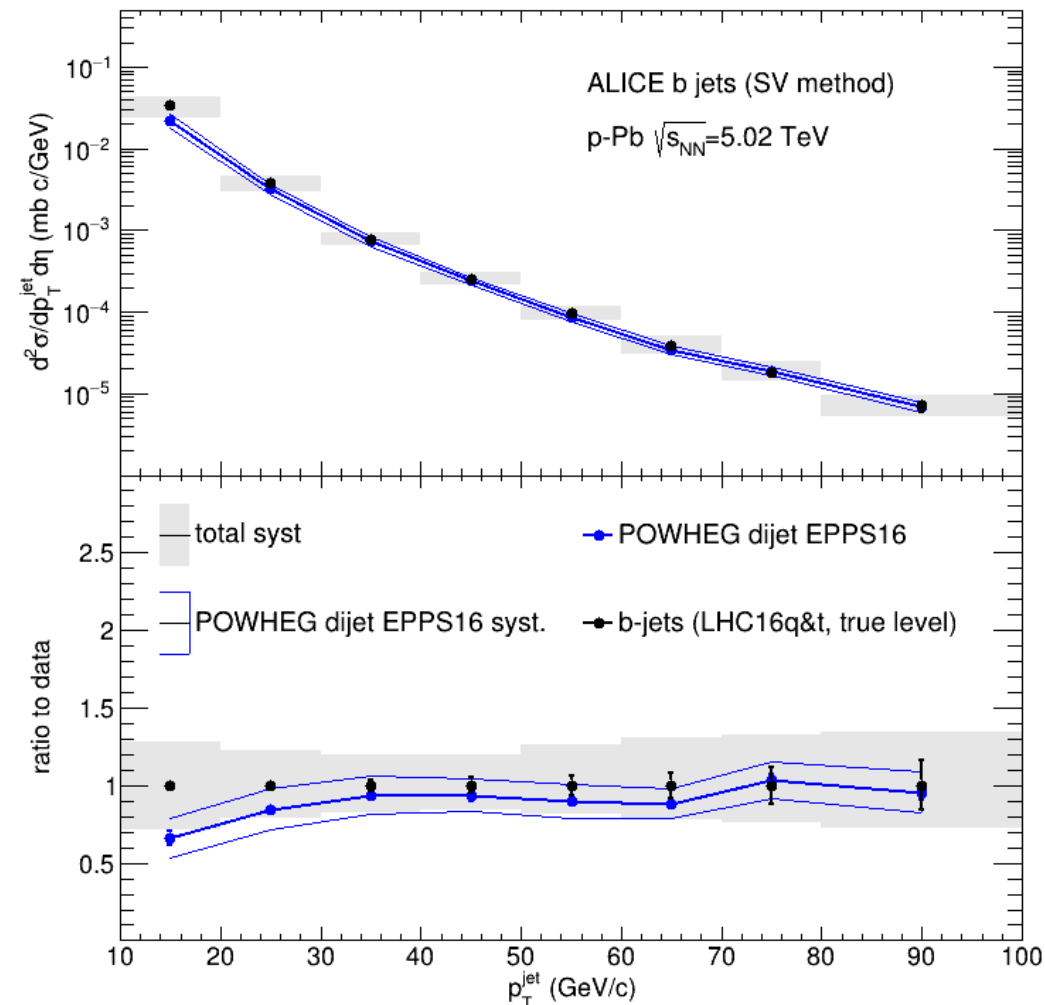
Total systematics calculated with the formula:

$$\Delta_{\text{sys}}^{\text{tot}} = \sqrt{\sum_i (\Delta_{\text{sys}}^i)^2}$$

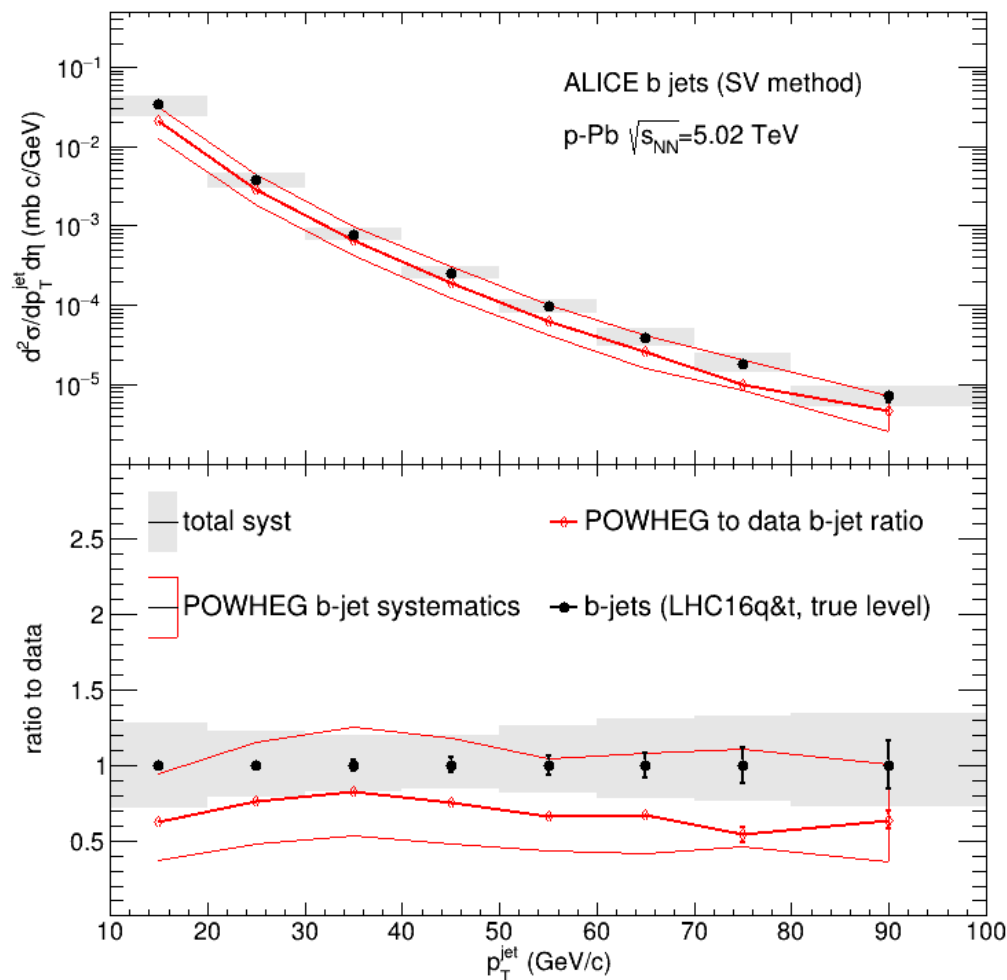
Source of syst. uncert.	$10 < p_{T,\text{ch jet}} < 20 \text{ GeV}/c$	$40 < p_{T,\text{ch jet}} < 50 \text{ GeV}/c$	$80 < p_{T,\text{ch jet}} < 100 \text{ GeV}/c$
Purity (%)	-24.7 / +18.2	-11.0 / +14.8	-15.7 / +28.5
SV Tagging (%)	-4.7 / +9.1	-4.3 / +3.1	-11.2 / +5.9
Unfolding (%)	-8.5 / +5.6	-4.4 / +3.6	-17.8 / +14.9
Tracking eff. (%)	-9.2 / +9.2	-7.9 / +7.9	-6.5 / +6.5
$p_{T,\text{track}}$ smearing (%)	-3.3 / +3.3	-4.5 / +4.5	-5.3 / +5.3
Secondaries (%)	-0 / +4.1	-0 / 5.4	-0 / +7.8
δp_T	-0 / +16.0	-0 / +5.4	-0 / +0.2
Total (%)	-28.2 / +28.5	-15.5 / +19.6	-27.5 / +34.7



Comparison SV results to POWHEG simulations

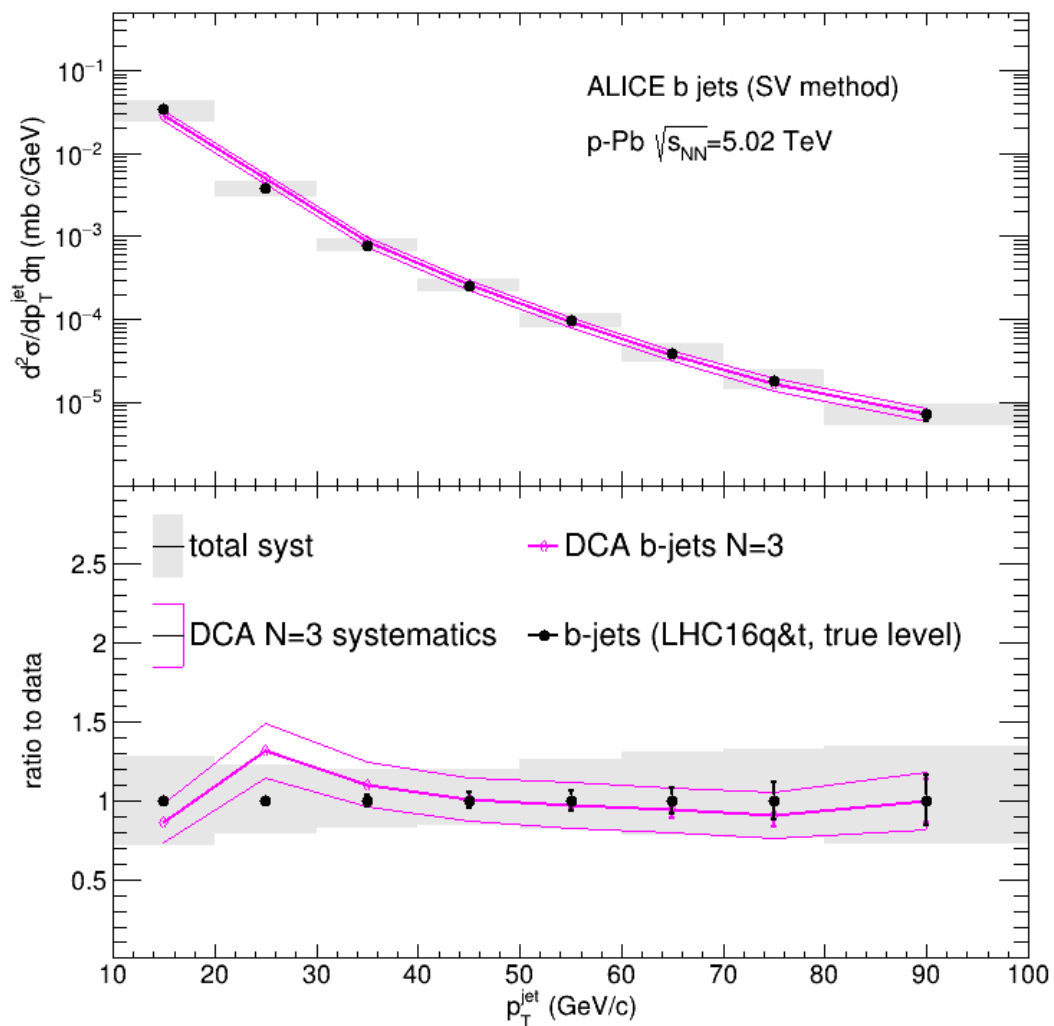
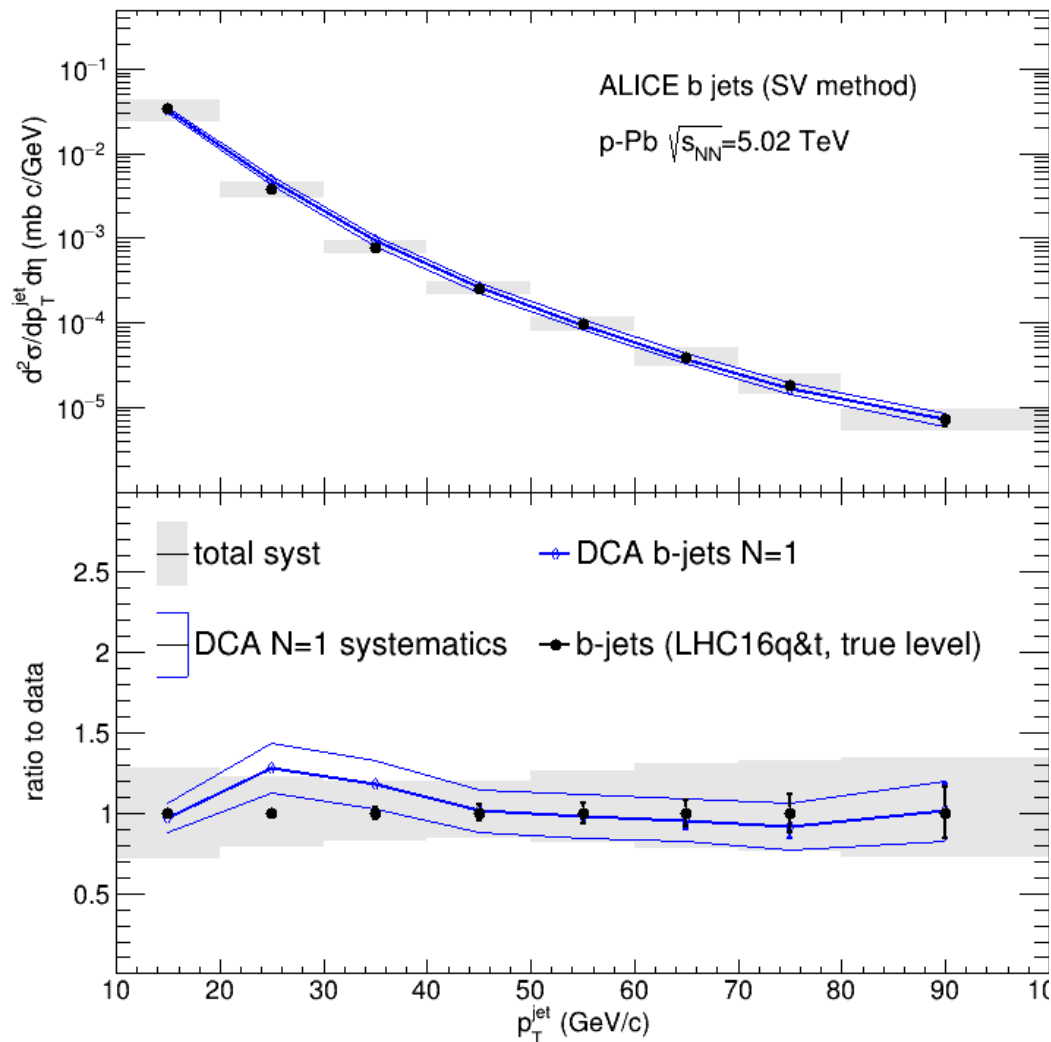


HVQ EPPS16 PYTHIA8



HVQ EPPS09NLO PYTHIA6

Comparison SV results to DCA analysis



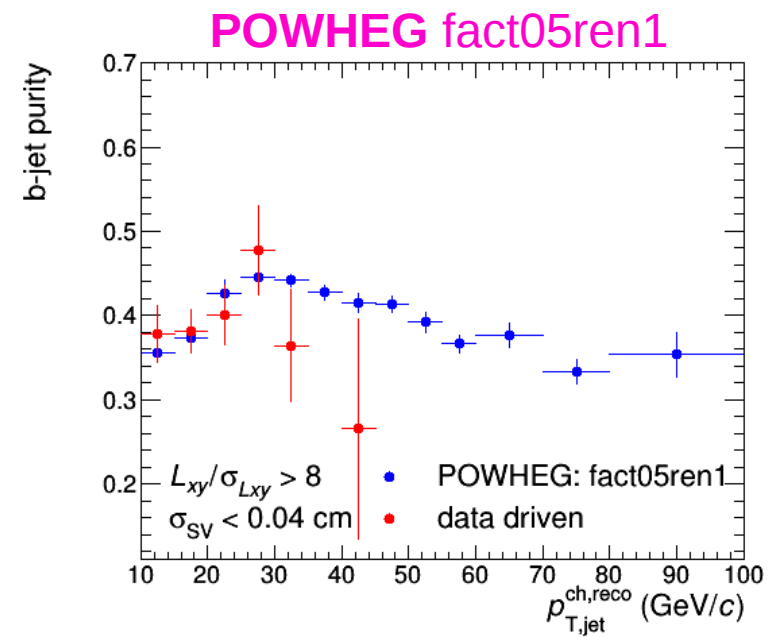
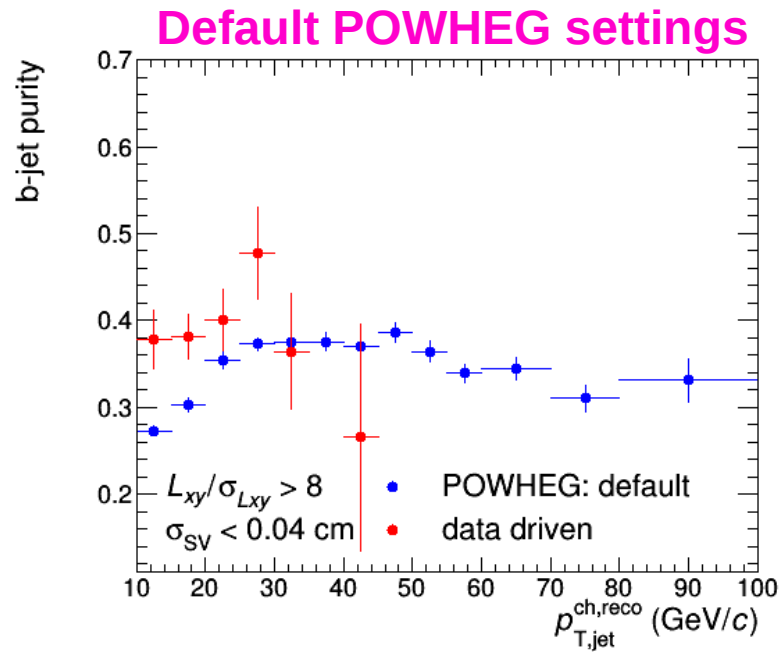
DCA analysis was done by Hadi Hassan in parallel [?]

- SV method was applied to measure b-jet spectra in p-Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV
- Correction on efficiency / purity of SV tagging was done with hybrid method (data-driven + POWbc)
- Results successfully are compatible with different POWHEG simulations and an independent analysis based on the DCA

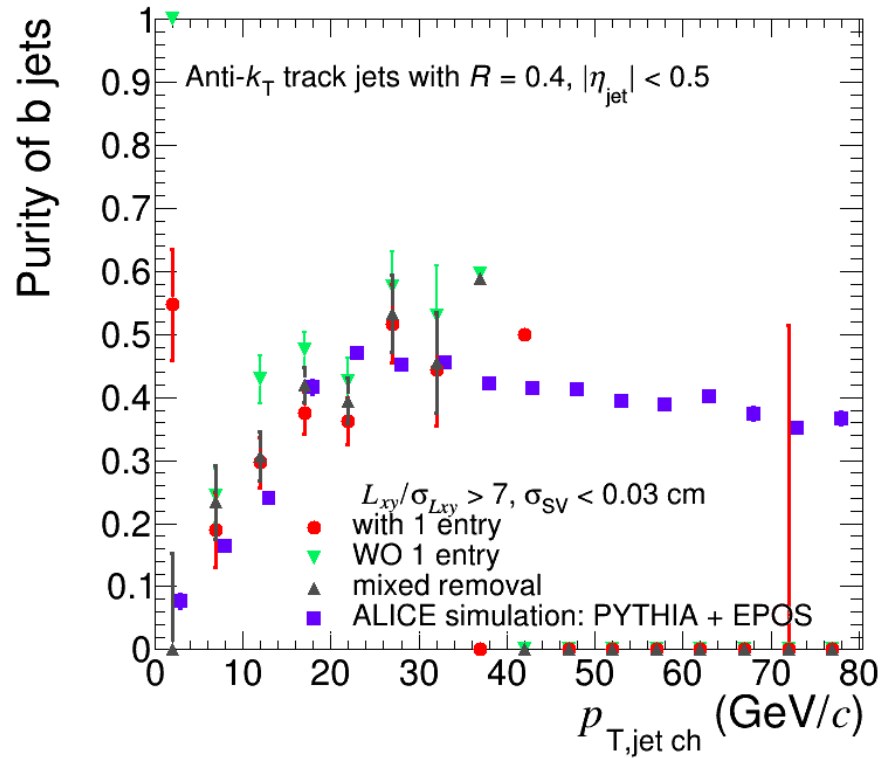
Further steps:

- Calculate R_{pB} for b-jets
- Apply MVA methods to estimate SV purity to efficiency correction and decrease systematic uncertainties

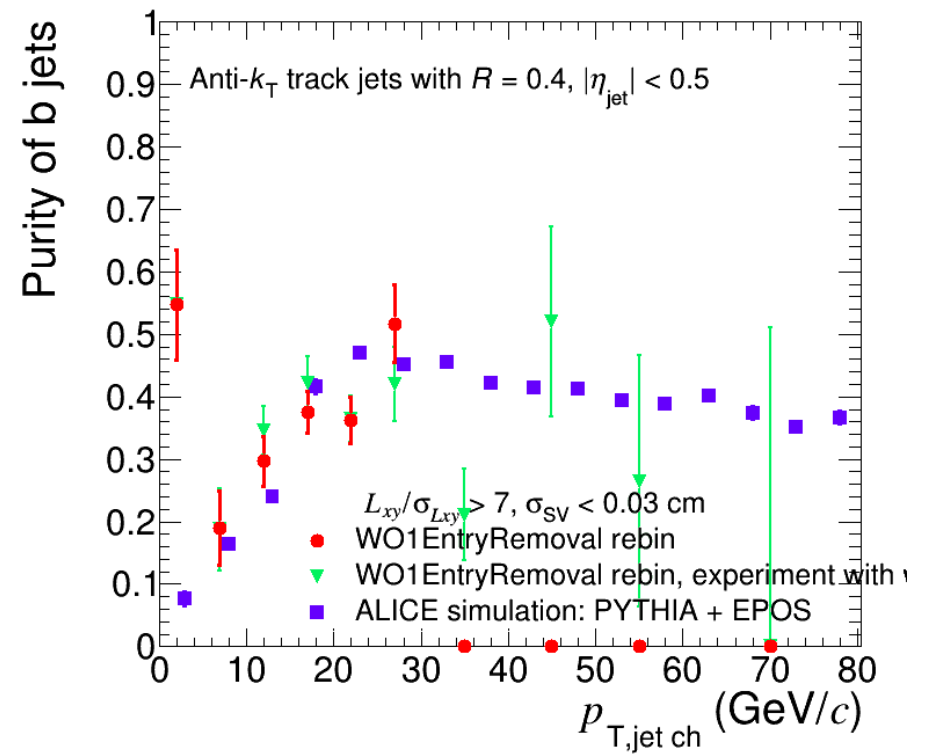
Backup



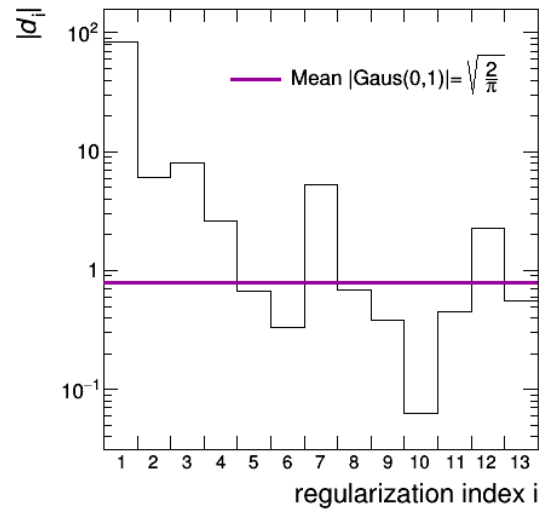
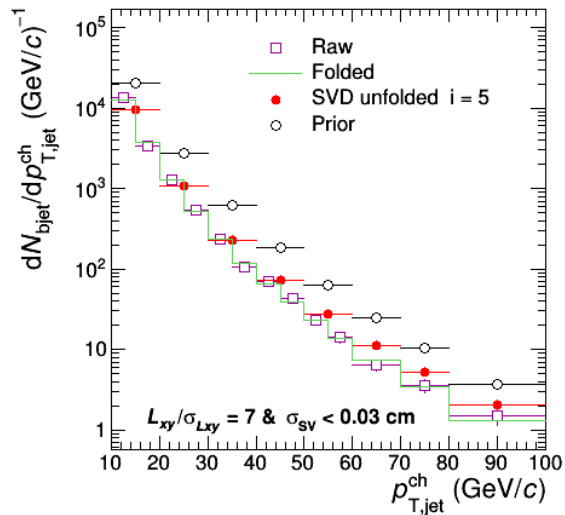
Effect of removal of 1-entry bin



Different binning

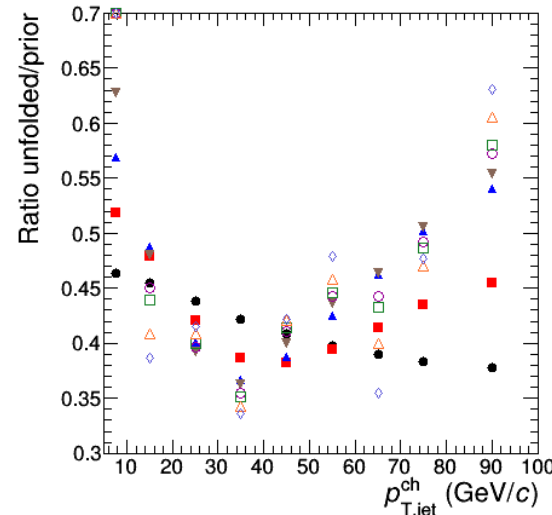
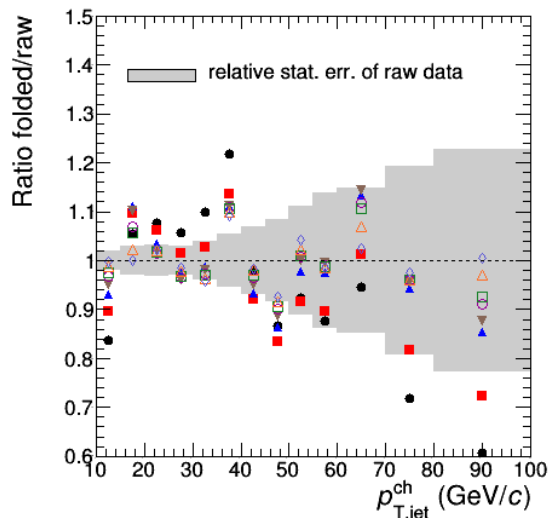


Unfolding of raw b-jet spectrum



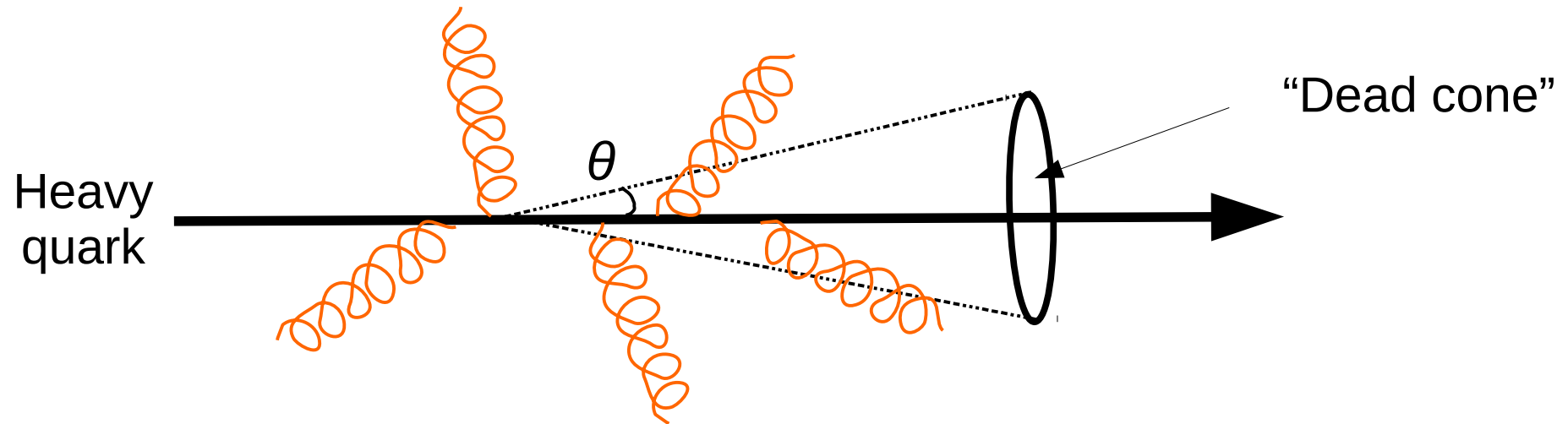
Analysis settings:

- $SL_{xy} > 7$, $\sigma_{SV} < 0.03$ cm
- SVD unfolding with $k=5$
- Pur/Eff correction from hybrid method
- prior POWHEG+PYTHIA6 b-jet spectrum
- δp_T matrix from real data events with SV
- instrumental matrix from PYTHIA



Dead cone effect

“Gluonsstrahlung” - process of gluon radiation by quarks (or gluons)



“Dead cone” effect – gluon radiation from massive quarks is suppressed at angles $\theta < m/E \rightarrow$ **Less E loss** inside the medium for heavy quarks expected

Gluonsstrahlung probability

$$\sim \frac{\theta^2}{[\theta^2 + (m/E)^2]^2}$$

[Yu.L. Dokshitzer, D.E. Kharzeev - “Heavy Quark Colorimetry of QCD Matter”, arXiv:hep-ph/0106202]

Probability of gluon emission

For light quarks:

$$dP_0 \simeq \frac{\alpha_s C_F}{\pi} \frac{d\omega}{\omega} \frac{dk_T^2}{k_T^2} = \frac{\alpha_s C_F}{\pi} \frac{d\omega}{\omega} \frac{d\theta^2}{\theta^2}$$

For heavy quarks:

$$dP_{HQ} = \frac{\alpha_s C_F}{\pi} \frac{d\omega}{\omega} \frac{k_T^2 dk_T^2}{(k_T^2 + \omega^2 \theta_0^2)^2} = \frac{\alpha_s C_F}{\pi} \frac{d\omega}{\omega} \frac{\theta^2 d\theta^2}{(\theta^2 + \theta_0^2)^2}$$

$$\theta_0 = \frac{M}{E}$$

Where

ω - Energy, C_F - “color charge”, k_T - transverse momenta

dP_0 - Probability to radiate gluon

Probability of gluon emission

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