

Theory benchmarks for $\sin^2\theta_W$ measurement: genuine EW and lineshape corrections

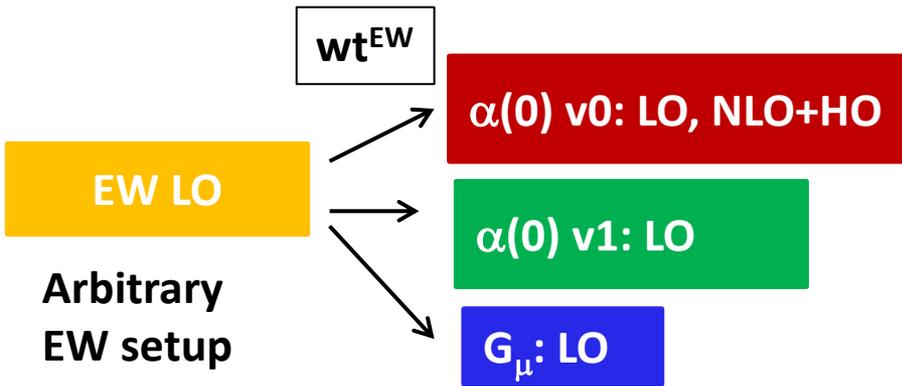
E. Richter-Was, IF UJ, Kraków

- Quick reminder
- Benchmarks status (November/December'18)
- Updates on Powheg + wt^{EW} (March'19)
- Outlook

Participants so far

PowhegZj: QCD NLO, Z+j

wt^{EW} : TauSpinner + Dizet 6.21



Powheg_ew: QCD LO, Z

$\alpha(0) v0$: LO

$\alpha(0) v1$: LO, NLO, NLO+HO

G_μ : LO, NLO, NLO+HO

MCSANC: QCD LO, Z

$\alpha(0) v1$: LO, NLO, NLO+HO

G_μ : LO, NLO, NLO+HO

DYTURBO: QCD LO, NLO, Z

$\alpha(0) v0$: LO

$\alpha(0) v1$: LO

G_μ : LO

ZGRAD2: QCD LO, EW NLO + HO

expressed interest to participate too

EW schemes: input parameters

SM fundamental relation used to calculate EW parameters at LO in different EW schemes, on-mass-shell definition.

↓ LEP legacy

↓ LHC standard

EW scheme: G_μ, α, M_Z

α, M_W, M_Z

G_μ, M_W, M_Z

$$G_\mu = \frac{\pi\alpha}{\sqrt{2}M_W^2s_W^2}$$

$$s_W^2 = 1 - m_W^2/m_Z^2$$

	$\alpha(0) \nu 0$	$\alpha(0) \nu 1$	G_μ
M_Z	91.1876 GeV	91.1876 GeV	91.1876 GeV
Γ_Z	2.4952 GeV	2.4952 GeV	2.4952 GeV
Γ_W	2.085 GeV	2.085 GeV	2.085 GeV
α	1/137.03599	1/137.03599	1/132.23323
G_μ	$1.1663787 \cdot 10^{-5} \text{ GeV}^{-2}$	$1.1254734 \cdot 10^{-5} \text{ GeV}^{-2}$	$1.1663787 \cdot 10^{-5} \text{ GeV}^{-2}$
M_W	80.93886 GeV	80.385 GeV	80.385 GeV
s_W^2	0.2121517	0.2228972	0.2228972
$\frac{G_\mu \cdot M_Z^2 \cdot \Delta^2}{\sqrt{2} \cdot 8\pi \cdot \alpha}$	1.0	1.0	1.0

Be aware: $\alpha(0) \nu 1$ comes with unphysical value of G_μ

DIZET library
exact $O(\alpha)$ + higher order terms

Powheg_ew, MCSANC
EW NLO, NLO+HO

EW schemes: input parameters

EW schemes: come with „on-shell” or „pole” definitions!

Table 44: The EW parameters used at tree-level EW, with on-mass-shell definition (LEP convention).

Parameter	$\alpha(0) \nu 0$	$\alpha(0) \nu 1$	G_μ
M_Z	91.1876 GeV	91.1876 GeV	91.1876 GeV
Γ_Z	2.4952 GeV	2.4952 GeV	2.4952 GeV
Γ_W	2.085 GeV	2.085 GeV	2.085 GeV
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Running Γ_Z in
Z-propagator

Shift:

- -30 MeV for M_Z
- change on Γ_Z
- -0.00006 for s_W^2

Scaling

- 0.99906 for α

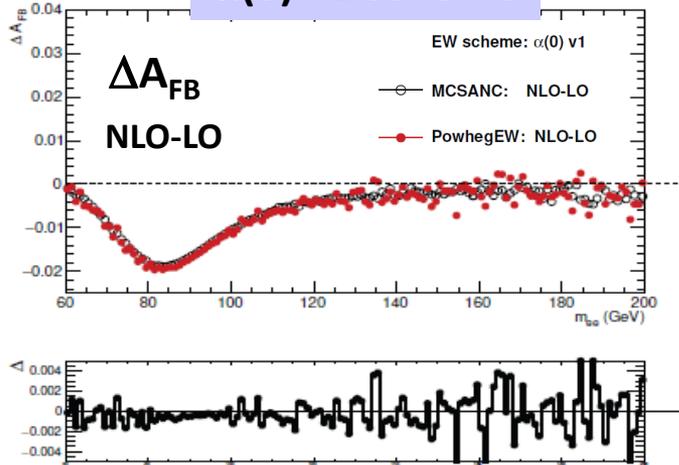
Fixed Γ_Z in
Z-propagator

Table 45: The EW parameters used at tree-level EW, with pole definition of the Z, W masses.

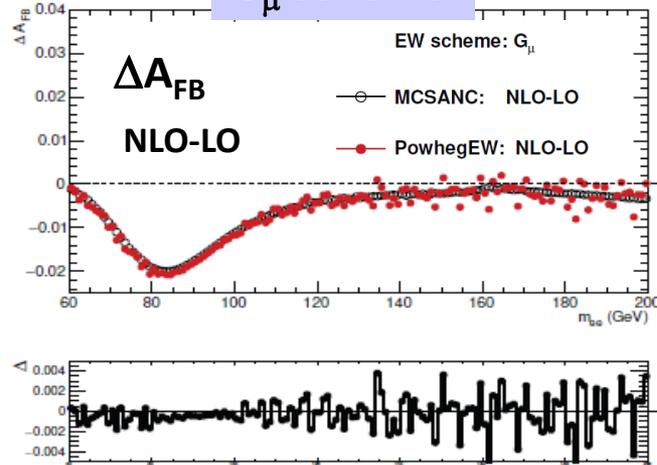
Parameter	$\alpha(0) \nu 0$	$\alpha(0) \nu 1$	G_μ
M_Z	91.15348 GeV	91.15348 GeV	91.15348 GeV
Γ_Z	2.494266 GeV	2.494266	2.494266 GeV
Γ_W	2.085 GeV	2.085 GeV	2.085 GeV
α	1/137.03599	1/137.03599	1/132.3572336357709
G_μ	$1.1663787 \cdot 10^{-5} \text{ GeV}^{-2}$	$1.126555497 \cdot 10^{-5} \text{ GeV}^{-2}$	$1.1663787 \cdot 10^{-5} \text{ GeV}^{-2}$
M_W	80.91191 GeV	80.35797 GeV	80.35797 GeV
s_W^2	0.21208680	0.22283820939	0.22283820939
$\frac{G_\mu \cdot M_Z^2 \cdot \Delta^2}{\sqrt{2} \cdot 8\pi \cdot \alpha}$	1.0	1.0	1.0

A_{FB} : EW LO, NLO, NLO+HO

$\alpha(0)$ v1 scheme

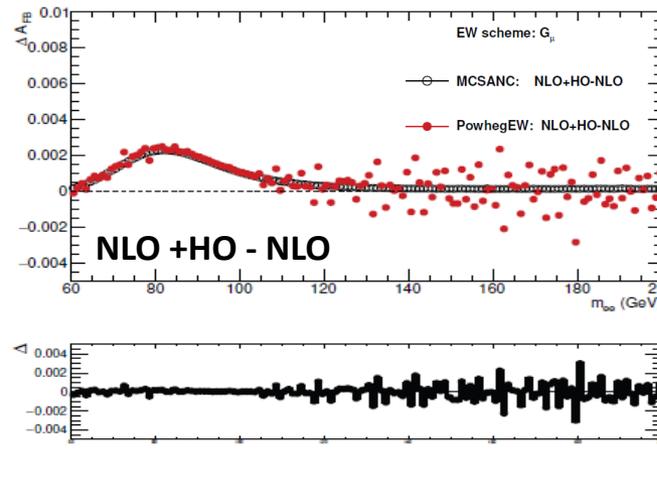
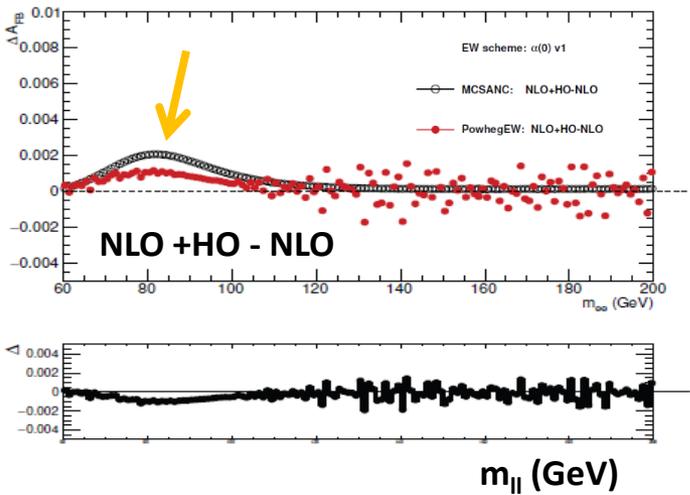


G_μ scheme



Status:
December 2018

● Powheg_ew
F. Piccinini et al.



○ MCSANC
S. Bondarenko,
L. Kalinovskaya

Pending investigating this discrepancy:

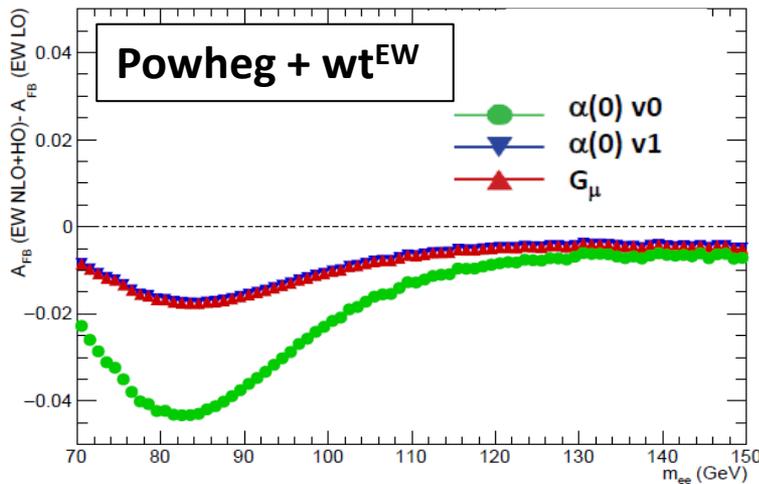
0.001 shift on ΔA_{FB} at Z-pole corresponds to shift $\sim 30 \cdot 10^{-5}$ on $\sin^2\theta_{eff}$

A_{FB} : EW LO, NLO+HO

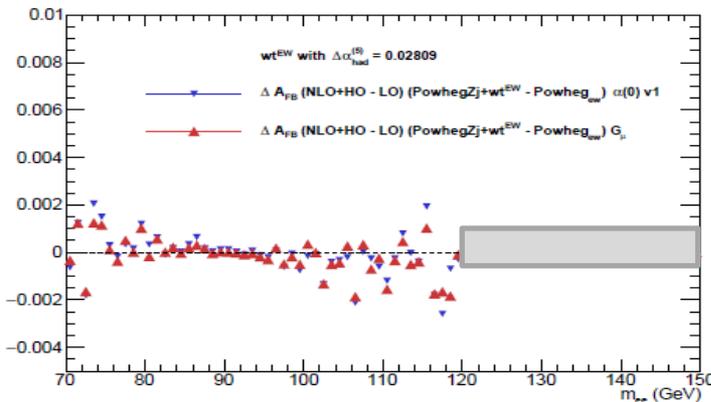
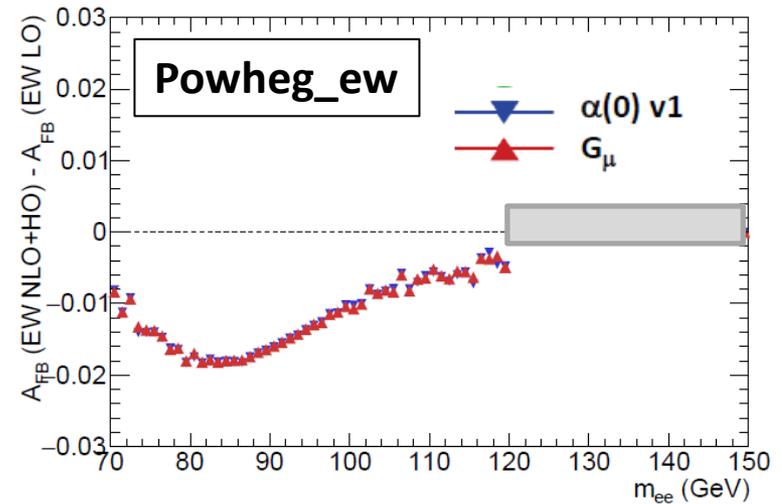
- Comparing Powheg_ew and Powheg+wt^{EW}

Status:
September 2018

ΔA_{FB} NLO+HO - LO



ΔA_{FB} NLO+HO - LO



wt^{EW} with $\Delta\alpha_{had}^{(5)} = 0.02809$

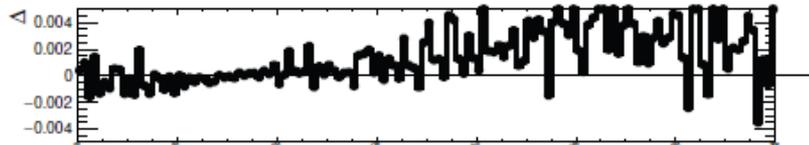
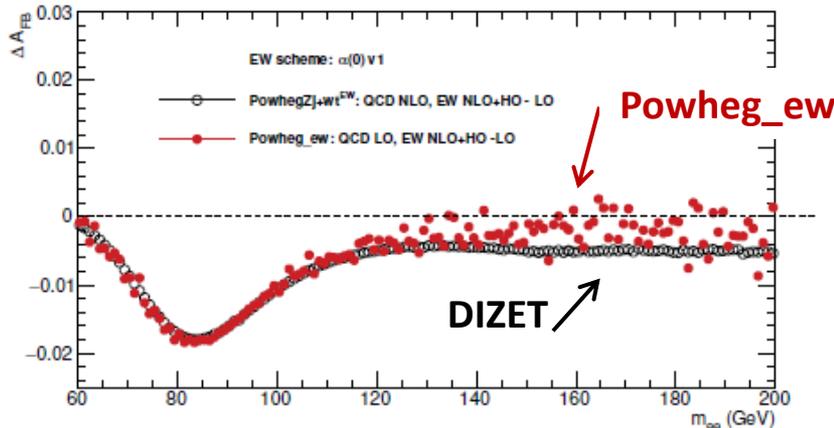
- ΔA_{FB} (NLO+HO - LO) (Powheg_{ew} - PowhegZj+wt^{EW}) $\alpha(0)$ v1
- ΔA_{FB} (NLO+HO - LO) (Powheg_{ew} - PowhegZj+wt^{EW}) G_μ

Excellent agreement on ΔA_{FB} ! Is it accidental?

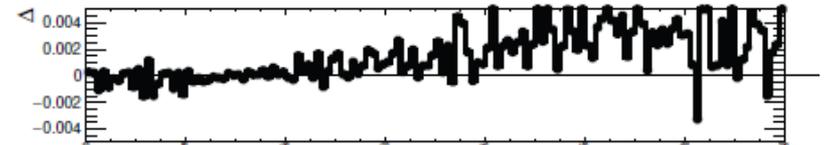
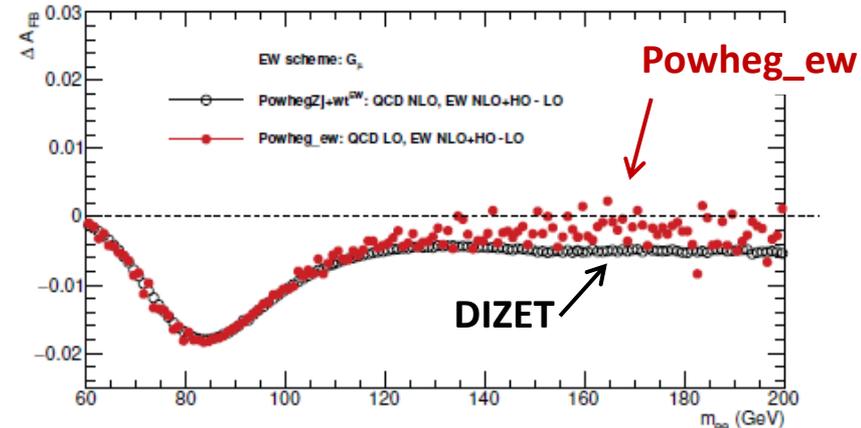
A_{FB} : EW LO, NLO+HO

Status:
December 2018

$\alpha(0)$ v1 scheme



G_μ scheme



Around Z-pole:

- > good agreement between Powheg_ew and DIZET
- > MCSANC shifted by 0.001

Pending investigating this discrepancy!

At higher masses:

- > DIZET predicts stable (NLO+HO - LO) corr. of 0.005
- > PowhegEW and MCSANC predicts (NLO+HO - LO) being close to zero.

Powheg + wt^{EW}: updates since December

- **EW weights calculated now with running Γ_Z in Z-boson propagator**
- **Updated parametrisation of $\Delta\alpha_{\text{had}}^{(5)}(s)$**
 - Jegerlehner 2017 (arXiv:1711.06089)
 - Burkhard&Pietrzyk 2001 (Phys.LettB 513 (2001) 46)
- **Updated Dizet 6.21 -> Dizet 6.42**
 - AMT4=4: Subleading two-loop corrections and re-summation recipe
 - AMT4=6: Complete two-loop corrections to M_W and fermionic two-loop corrections to $\sin^2\theta_{\text{eff}}^{\text{lep}}$

Impact of $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$

Dizet 6.21, 6.42

NEW

Parameter	$\Delta\alpha_h^{(5)}(M_Z^2) = 0.0280398$ (param. Jegerlehner 1995)	$\Delta\alpha_h^{(5)}(M_Z^2) = 0.0275762$ (param. Jegerlehner 2017)	Δ
$\alpha(M_Z^2)$	0.00775884	0.00775492	
$1/\alpha(M_Z^2)$	128.885224	128.9503292	
s_W^2	0.22351946	0.22332758	- 0.00019
$\sin^2\theta_W^{\text{eff}}(M_Z^2)$ (electron, muon)	0.23175990	0.23158294	- 0.00018
$\sin^2\theta_W^{\text{eff}}(M_Z^2)$ (up-quark)	0.23164930	0.23147645	- 0.00017
$\sin^2\theta_W^{\text{eff}}(M_Z^2)$ (down-quark)	0.23152214	0.23134945	- 0.00017
M_W	80.35281 GeV	80.36274 GeV	+10 MeV
Δr	0.03694272	0.036609	
Δr_{rem}	0.01169749	0.01170287	

$$M_W = \frac{M_Z}{\sqrt{2}} \sqrt{1 + \sqrt{1 - \frac{4A_0^2}{M_Z^2(1 - \Delta r)}}}$$

$$\Delta r = \Delta\alpha(M_Z^2) + \Delta r_{EW}$$

$$A_0 = \sqrt{\frac{\pi\alpha(0)}{\sqrt{2}G_\mu}}$$

Impact of AMT=4 → AMT=6

Updated to Dizet 6.42, which comes with more complete two-loop corrections

NEW

Parameter	AMT4= 4	AMT4 = 6	Δ
$\alpha(M_Z^2)$	0.00775492	0.00775492	
$1/\alpha(M_Z^2)$	128.9503239	128.9503239	
s_W^2	0.22332758	0.22343647	+ 0.00011
$\sin^2\theta_W^{eff}(M_Z^2)$ (electron, muon)	0.23158294	0.23153917	-0.000044
$\sin^2\theta_W^{eff}(M_Z^2)$ (up-quark)	0.23147645	0.23143261	-0.000044
$\sin^2\theta_W^{eff}(M_Z^2)$ (down-quark)	0.23134945	0.23130551	-0.000044
M_W	80.36274 GeV	80.35710 GeV	- 5.6 MeV
Δr	0.036609	0.03636609	
Δr_{rem}	0.01170287	0.01170287	

$$M_W = \frac{M_Z}{\sqrt{2}} \sqrt{1 + \sqrt{1 - \frac{4A_0^2}{M_Z^2(1 - \Delta r)}}}$$

$$\Delta r = \Delta\alpha(M_Z^2) + \Delta r_{EW}$$

$$A_0 = \sqrt{\frac{\pi\alpha(0)}{\sqrt{2}G_\mu}}$$

- **AMT4=4:** Subleading two-loop corrections and re-summation recipe (best option Dizet 6.21)
- **AMT4=6:** Complete two-loop corrections to M_W and fermionic two-loop corrections to $\sin^2\theta_{eff}^{lep}$ (best option in Dizet 6.42)

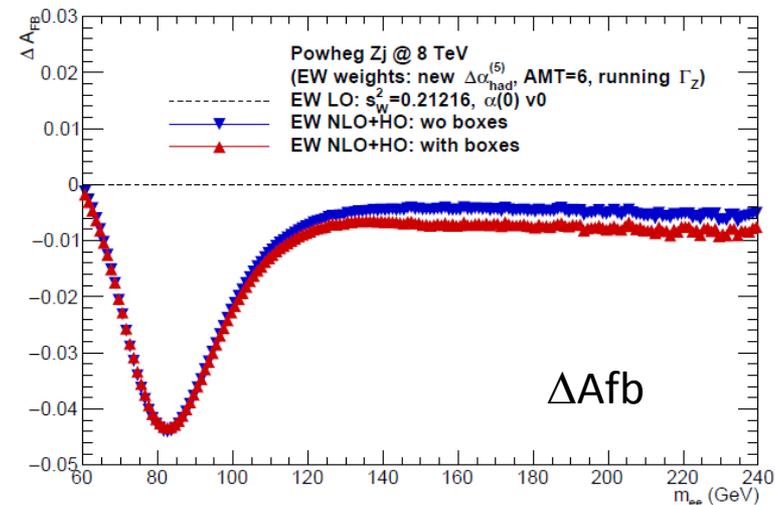
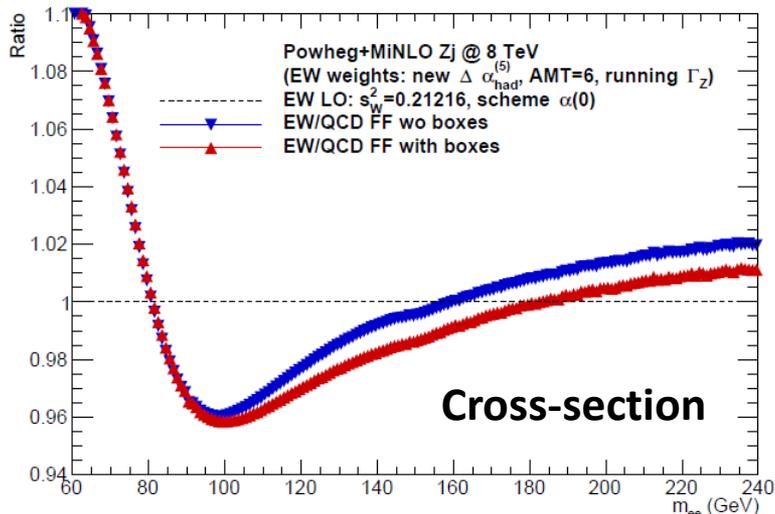
Powheg + wt^{EW}: updates since December

All updates together:

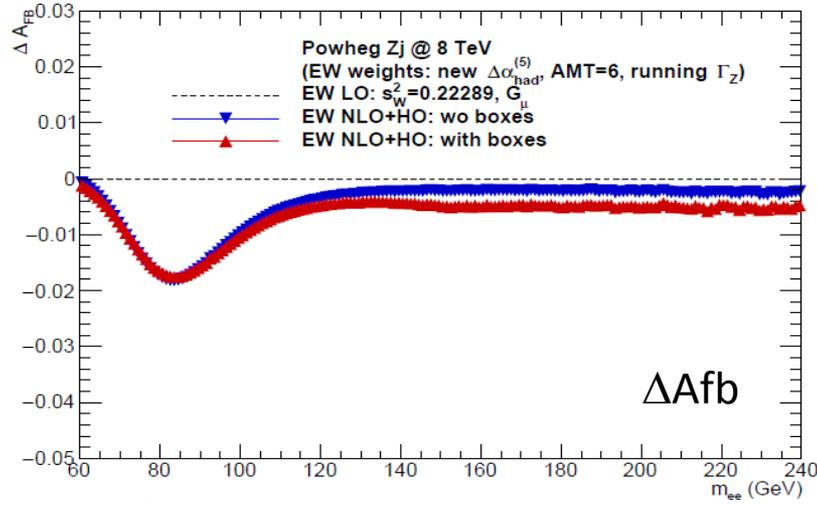
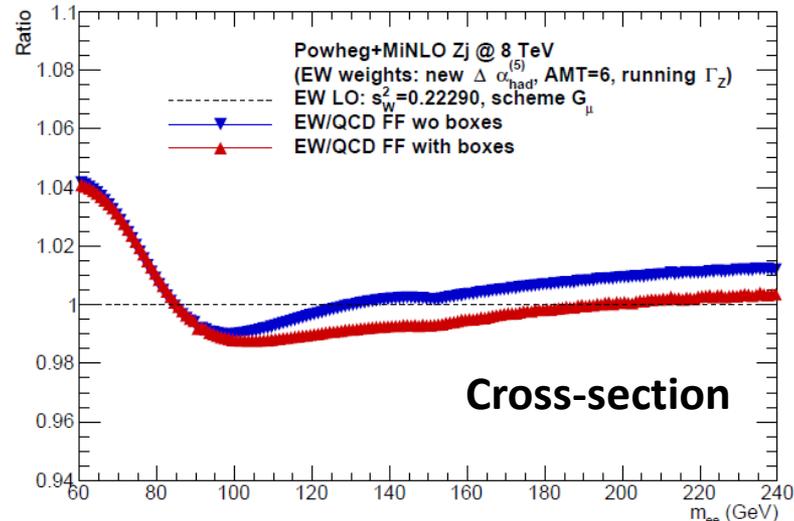
	September'18	May'19		
Parameter	$m_H = 89-93$ GeV	$m_H = 89-93$ GeV		
$\alpha(0)$ scheme				
σ NLO+HO/LO (no boxes)	0.96536	0.96503	0.03%	
σ NLO+HO/LO (with boxes)	0.96536	0.96508		
A_{fb} NLO+HO - LO (no boxes)	-0.03529	-0.03496	+0.0003	→ About + 0.00010 on \sin^2_{eff}
A_{fb} NLO+HO - LO (with boxes)	-0.03526	-0.03495		
G_μ scheme				
σ NLO+HO/LO (no boxes)	0.99243	0.99204	0.04%	
σ NLO+HO/LO (with boxes)	0.99244	0.99209		
A_{fb} NLO+HO - LO (no boxes)	-0.01553	-0.01514	+0.0004	→
A_{fb} NLO+HO - LO (with boxes)	-0.01550	-0.01512		

Powheg + wt^{EW}: EW NLO+HO corr.

$\alpha(0)$ v0 scheme



G_μ scheme



Benchmarking: outlook

- Shall we continue with the present setup or there is different/better proposal?
- In Durham WS, Fulvio presented proposal for new EW scheme implemented in Powheg_ew. We can include it as well, need to define table for EW LO parameters below, and follow with standard plots/tables for EW NLO+HO - LO

	↓ LEP legacy		↓ LHC standard	
EW scheme:	G_μ, α, M_Z	α, M_W, M_Z	G_μ, M_W, M_Z	
	$\alpha(0) \nu 0$	$\alpha(0) \nu 1$	G_μ	$\sin 2w$
M_Z	91.1876 GeV	91.1876 GeV	91.1876 GeV	
Γ_Z	2.4952 GeV	2.4952 GeV	2.4952 GeV	
Γ_W	2.085 GeV	2.085 GeV	2.085 GeV	
α	1/137.03599	1/137.03599	1/132.23323	
G_μ	$1.1663787 \cdot 10^{-5} \text{ GeV}^{-2}$	$1.1254734 \cdot 10^{-5} \text{ GeV}^{-2}$	$1.1663787 \cdot 10^{-5} \text{ GeV}^{-2}$	
M_W	80.93886 GeV	80.385 GeV	80.385 GeV	
s_W^2	0.2121517	0.2228972	0.2228972	
$\frac{G_\mu \cdot M_Z^2 \cdot \Delta^2}{\sqrt{2} \cdot 8\pi \cdot \alpha}$	1.0	1.0	1.0	

Benchmarking: outlook

- The **wt^{EW} calculations are updated.**
- Comparison done so far still not detailed enough, agreement achieved (slide 6) seemed accidental.
- To complete, we need to break it into a few checkpoints:
 - Predicted α_{QED} at Z-pole
 - Predicted $\sin^2\theta_{\text{eff}}$ at Z-pole
 - Corrections due to Z self-energy (running width propagator)
 - What else?
- We need to complete and agree upon evaluation of parametric and theory uncertainties of each EW scheme.
 - For the $\alpha(0)$ v0 scheme (DIZET)
 - Parametric uncertainties: **$4 \cdot 10^{-5}$ for $\sin^2\theta_{\text{eff}}$**
 - Theory uncertainties: **$5 \cdot 10^{-5}$ for $\sin^2\theta_{\text{eff}}$ [LEP legacy]**
 - For the G_μ scheme
 - Parametric uncertainties: **$\pm 15 \text{ MeV}$ on M_W -> $30 \cdot 10^{-5}$ on $\sin^2\theta_W$**

Theoretical and parametric uncertainties

ALDO+ SLD + Tevatron, arXiv:1012.2367

- The remaining theoretical uncertainties were estimated to be 4 MeV on m_W and **0.000049 on $\sin^2\theta_{\text{eff}}^{\text{lep}}$** . We can use this estimate @ Z-pole.
- The parametric uncertainties were dominated by $\Delta\alpha_{\text{had}}(M_Z^2)$. The uncertainty of **± 0.00035 caused an error of ± 0.00013 on $\sin^2\theta_{\text{eff}}^{\text{lep}}$** .

EW weights from TauSpinner/DIZET, $\alpha(0)$ scheme:

- The same code as ALDO for calculating EW genuine corrections. We can use this estimate.
Theoretical uncertainties (EW, @Z-pole): $5 \cdot 10^{-5}$ on $\sin^2\theta_{\text{eff}}^{\text{lep}}$
- The parametric uncertainties should be updated for better precision on $\Delta\alpha_{\text{had}}(M_Z^2)$ and m_t .

Parametric uncertainties: $\Delta\alpha_{\text{had}}, m_t$

Parameter	$\Delta\alpha_h^{(S)}(M_Z^2) - 0.0001$	$\Delta\alpha_h^{(S)}(M_Z^2) = 0.0275762$	$\Delta\alpha_h^{(S)}(M_Z^2) + 0.0001$	$\Delta/2$
$\alpha(M_Z^2)$	0.0077540999	0.0077549240	0.0077557482	
$1/\alpha(M_Z^2)$	128.9640328306	128.9503292550	128.9366256793	
s_W^2	0.22340146	0.22343647	0.22347148	0.000035
$\sin^2\theta_W^{\text{eff}}(M_Z^2)$ (electron, muon)	0.23150412	0.23153917	0.23157421	0.000035
$\sin^2\theta_W^{\text{eff}}(M_Z^2)$ (up-quark)	0.23139759	0.23143261	0.23146763	0.000035
$\sin^2\theta_W^{\text{eff}}(M_Z^2)$ (down-quark)	0.23127052	0.23130551	0.23134049	0.000035
M_W	80.35892 GeV	80.35710 GeV	80.35529 GeV	1.8 MeV
Δr	0.03625683	0.036609	0.03647535	
Δr_{rem}	0.01170310	0.01170287	0.01170264	

Parameter	$m_t - 0.5 \text{ GeV}$	$m_t = 173.2 \text{ GeV}$	$m_t + 0.5 \text{ GeV}$	$\Delta/2$
$\alpha(M_Z^2)$	0.0077549205	0.0077549240	0.0077549274	
$1/\alpha(M_Z^2)$	128.9503873792	128.9503292550	128.9502716590	
s_W^2	0.22349450	0.22343647	0.22337836	0.000058
$\sin^2\theta_W^{\text{eff}}(M_Z^2)$ (electron, muon)	0.23155486	0.23153917	0.23152344	0.000016
$\sin^2\theta_W^{\text{eff}}(M_Z^2)$ (up-quark)	0.23144830	0.23143261	0.23141688	0.000016
$\sin^2\theta_W^{\text{eff}}(M_Z^2)$ (down-quark)	0.23132119	0.23130551	0.23128979	0.000016
M_W	80.354102 GeV	80.35710 GeV	80.360111 GeV	3 MeV
Δr	0.03654697	0.036609	0.03618477	
Δr_{rem}	0.01169343	0.01170287	0.01171229	

@ Z-pole:

Parametric uncertainties: **4 10⁻⁵** [35 10⁻⁶ (from $\Delta\alpha_{\text{had}}$) and 16 10⁻⁶ (from m_t)]

SPARE slides

EW schemes: LEP and LHC paradigms

LEP

- EW scheme: „on-mass-shell” regularisation; input: $(\alpha(0), G_F, M_Z)$, the most precisely known quantities:
 $\Delta\alpha/\alpha \sim 3.6 \cdot 10^{-9}$; $\Delta G_F/G_F \sim 8.6 \cdot 10^{-6}$; $\Delta M_Z/M_Z \sim 2.4 \cdot 10^{-5}$
- TH precision @Z-pole: 4 MeV on m_W ; $5 \cdot 10^{-5}$ on $\sin^2\theta_{\text{eff}}$
- Parametric precision on $\sin^2\theta_{\text{eff}}$: $4 \cdot 10^{-5}$ (dominated by $\Delta\alpha_{\text{had}}$)
- Used and developed still today: GFitter, FCCee preparation.

LHC

- EW scheme: „pole-mass” regularisation; input: (G_F, M_Z, M_W)
 $\Delta G_F/G_F \sim 8.6 \cdot 10^{-6}$; $\Delta M_Z/M_Z \sim 2.4 \cdot 10^{-5}$; $\Delta M_W/M_W \sim 1.9 \cdot 10^{-4}$
- **Relations to EWPO's: not established.**
- EW corrections complete at EW NLO, only some at EW NLO+HO
- **TH precision @Z-pole: not established.**
- **Parametric precision: ± 15 MeV on M_W -> $30 \cdot 10^{-5}$ on $\sin^2\theta_W$**

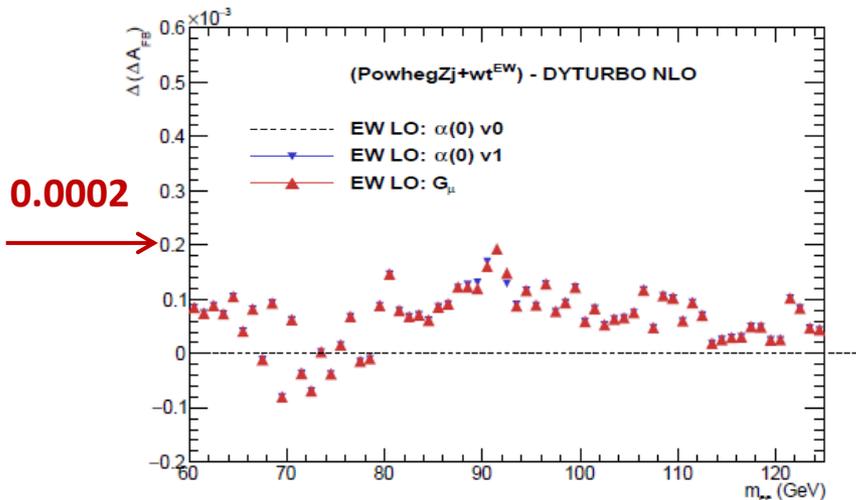
Strategy for comparisons

- **Scope:**
 - **Genuine EW and lineshape corrections** to Drell-Yan production at NLO QCD.
 - Three EW LO schemes chosen to allow for straightforward interpretation of results. We tuned EW LO parameters, otherwise out-of-the-box.
 - The highest available corrections in a given approach used.
 - QED FRS/ISR not included, so called double-deconvoluted observables.
- **Observables:**
 - Lineshape (cross-section) and forward-backward asymmetry A_{FB} in the full phase-space.
 - **Compared ratios or absolute differences** between different EW LO schemes and/or between NLO, NLO+HO predictions within each EW scheme and same MC generator. Allows to **minimize sensitivity to QCD details**.
- **Goals:**
 - Check if reweighting with **wt^{EW} (TauSpinner)** works for NLO QCD MC's. Compared distributions at EW LO (**DYTURBO, Powheg_ew**).
 - Establish how consistent are predictions between different EW schemes with EW NLO corrections (**Powheg_ew, MCSANC**).
 - Establish how consistent are EW NLO+HO corrections of **Dizet 6.21 form-factors** implemented in **wt^{EW}** and those of **Powheg_ew**.

Validating reweighting with wt^{EW} : EW LO

ΔA_{FB} : driven by s^2_W value (same for $\alpha(0)$ v1 and G_μ schemes)

Benchmark for wt^{EW} reweighting



Double difference:

ΔA_{FB} (DYTURBO) - ΔA_{FB} (PowhegZj+wt^{EW})

$\alpha(0)$ v1 - $\alpha(0)$ v0

G_μ - $\alpha(0)$ v0

Agreement on $\Delta(\Delta A_{FB})$ within ± 0.0002

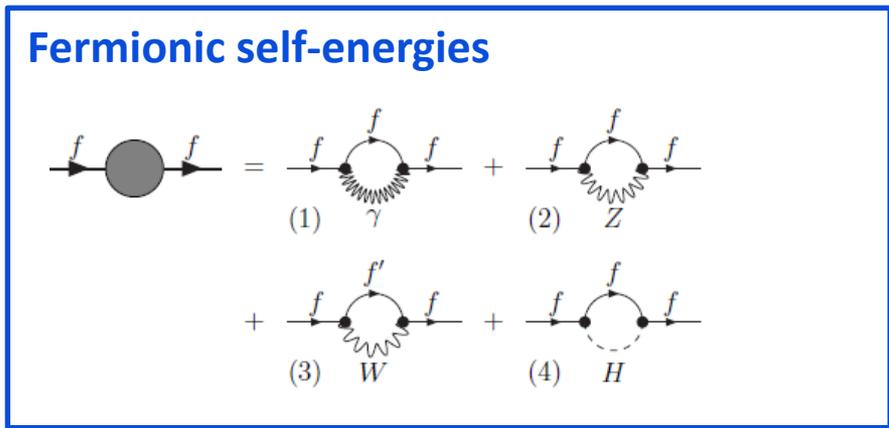
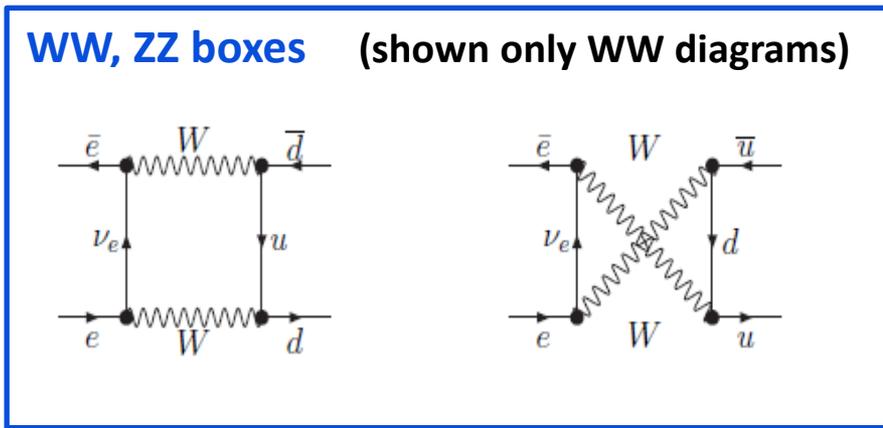
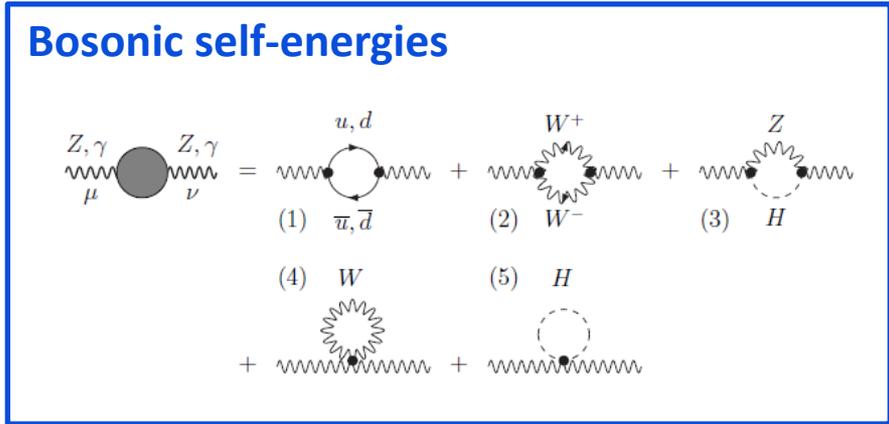
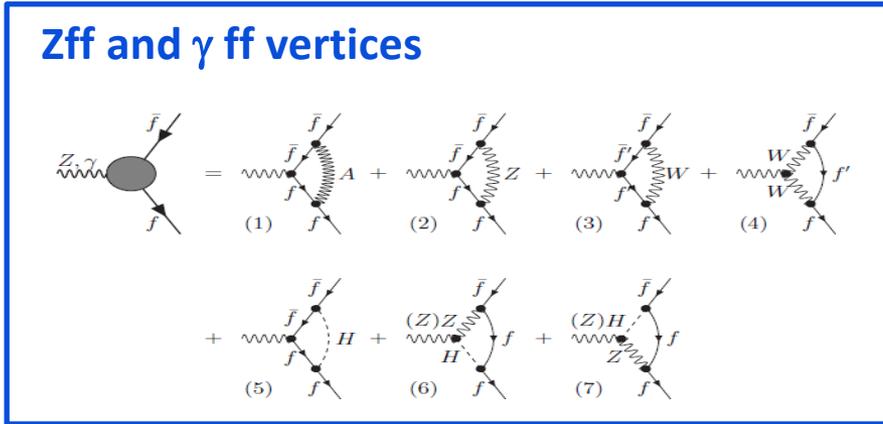
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Should redo it with much finer binning around Z-pole to better estimate precision.

Genuine EW and lineshape corrections

Gauge-invariant set of diagrams. Calculated as form-factor corrections to couplings, propagators and masses.

Eg. running $\alpha_{\text{QED}}(s)$, $\alpha_{\text{QED}}(M_Z) = 1./128.86674175$



From Zfitter/Dizet documentation

D. Bardin et al.
arXiv:9908433

Zfitter is a **semi-analytical program** for calculating total cross-sections and pseudo-observables (eg. A_{fb} , $\sin^2\theta_W^{\text{eff}}$), used by LEP1, and to a lesser degree by LEP2.

DIZET is a library for calculating form-factors and some other corrections. Provides complete EW $O(\alpha)$ weak-loop corrections supplemented with selected higher order terms (eg. vacuum polarisation, $\alpha_{\text{QED}}(Q^2)$).

For analyses at LEP1, LEP2 used always in parallel with **MC generators (KoralZ, KoralW)** eg. to evaluate systematics of simplified cuts used in analysis integration.

$$\begin{aligned}
 A_Z^{OLA}(s, t) = & i\sqrt{2}G_\mu I_e^{(3)} I_f^{(3)} M_Z^2 \chi_Z(s) \rho_{ef}(s, t) \left\{ \gamma_\mu(1 + \gamma_5) \otimes \gamma_\mu(1 + \gamma_5) \right. \\
 & - 4|Q_e|s_W^2 \kappa_e(s, t) \gamma_\mu \otimes \gamma_\mu(1 + \gamma_5) - 4|Q_f|s_W^2 \kappa_f(s, t) \gamma_\mu(1 + \gamma_5) \otimes \gamma_\mu \\
 & \left. + 16|Q_e Q_f|s_W^4 \kappa_{e,f}(s, t) \gamma_\mu \otimes \gamma_\mu \right\}. \tag{A.4.75}
 \end{aligned}$$

one loop amplitude

$$A_\gamma^{OLA} = i\chi_\gamma(s) \alpha(s) \gamma_\mu \otimes \gamma_\mu. \tag{2.2.36}$$

Dyson summation leads to the change of α into $\alpha(s)$:

$$\alpha(s) = \frac{\alpha(0)}{1 - \Delta\alpha^{\text{fer}}(s)} = \frac{\alpha(0)}{1 - \Delta\alpha^{(5)}(s) - \Delta\alpha^t(s) - \Delta\alpha^{\alpha\alpha_s}(s)}. \tag{2.2.37}$$

Vacuum polarisation corrections

LEP legacy: from Zfitter/Dizet documentation

After some trivial algebra one derives the final expressions:

$$\boxed{\rho_{ef}} = 1 + \frac{g^2}{16\pi^2} \left\{ -\Delta\rho_z^F + \mathcal{D}_z^F(s) + \frac{5}{3}B_0^F(-s; M_W, M_W) - \frac{9}{4} \frac{c_w^2}{s_w^2} \ln c_w^2 - 6 \right. \\ \left. + \frac{5}{8}c_w^2(1+c_w^2) + \frac{1}{4c_w^2}(3v_e^2+a_e^2+3v_f^2+a_f^2)\mathcal{F}_z(s) + \hat{\mathcal{F}}_w^0(s) + \hat{\mathcal{F}}_w(s) \right. \\ \left. - \frac{\tau_t}{4}[B_0^F(-s; M_W, M_W) + 1] - c_w^2(R_z - 1)s\hat{\mathcal{B}}_{ww}^d(s, t) \right\}, \quad (\text{A.4.80})$$

$$\boxed{\kappa_e} = 1 + \frac{g^2}{16\pi^2} \left\{ -\frac{c_w^2}{s_w^2}\Delta\rho^F - \Pi_{z\gamma}^F(s) - \frac{1}{6}B_0^F(-s; M_W, M_W) - \frac{1}{9} - \frac{v_e\sigma_e}{2c_w^2}\mathcal{F}_z(s) \right. \\ \left. - \hat{\mathcal{F}}_w^0(s) + (R_z - 1) \left[\frac{|Q_f|}{2}(1 - 4|Q_f|s_w^2)\mathcal{F}_z(s) + c_w^2[\hat{\mathcal{F}}_{wn}(s) \right. \right. \\ \left. \left. - |Q_f|\mathcal{F}_{wa}(s) + s\hat{\mathcal{B}}_{ww}^d(s, t) \right] \right\}, \quad (\text{A.4.81})$$

$$\boxed{\kappa_f} = 1 + \frac{g^2}{16\pi^2} \left\{ -\frac{c_w^2}{s_w^2}\Delta\rho^F - \Pi_{z\gamma}^F(s) - \frac{1}{6}B_0^F(-s; M_W, M_W) - \frac{1}{9} - \frac{v_f\sigma_f}{2c_w^2}\mathcal{F}_z(s) \right. \\ \left. - \hat{\mathcal{F}}_w(s) + (R_z - 1) \left[\frac{|Q_e|}{2}(1 - 4|Q_e|s_w^2)\mathcal{F}_z(s) + c_w^2[\hat{\mathcal{F}}_{wn}^0(s) \right. \right. \\ \left. \left. - |Q_e|\mathcal{F}_{wa}(s) + s\hat{\mathcal{B}}_{ww}^d(s, t) \right] - \frac{\tau_t}{4}[B_0^F(-s; M_W, M_W) + 1] \right\}, \quad (\text{A.4.82})$$

interference

$$\boxed{\kappa_{ef}} = 1 + \frac{g^2}{16\pi^2} \left\{ -2\frac{c_w^2}{s_w^2}\Delta\rho^F - 2\Pi_{z\gamma}^F(s) - \frac{1}{3}B_0^F(-s; M_W, M_W) - \frac{2}{9} \right. \\ \left. - \frac{1}{4c_w^2} \left[\frac{\delta_e^2 + \delta_f^2}{s_w^2}(R_w - 1) + 3v_e^2 + a_e^2 + 3v_f^2 + a_f^2 \right] \mathcal{F}_z(s) \right. \\ \left. - \hat{\mathcal{F}}_w^0(s) - \hat{\mathcal{F}}_w(s) - \frac{\tau_t}{4}[B_0^F(-s; M_W, M_W) + 1] \right. \\ \left. + c_w^2(R_z - 1) \left[\frac{2}{3} - \hat{\Pi}_{\gamma\gamma}^{\text{bos},F}(s) + s\hat{\mathcal{B}}_{ww}^d(s, t) \right] \right\}. \quad (\text{A.4.83})$$

E Fermionic loops in γ propagator

BOX

CERN, LHC EW precision, 7.05.2019

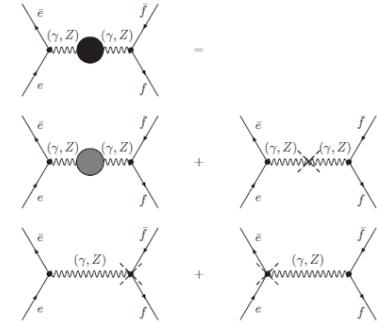


Figure A.11. Bosonic self-energies and bosonic counter-terms for $e\bar{e} \rightarrow (Z, \gamma) \rightarrow f\bar{f}$

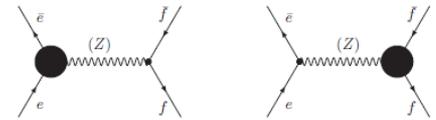


Figure A.10. Electron (a) and final fermion (b) vertices in $e\bar{e} \rightarrow (Z) \rightarrow f\bar{f}$

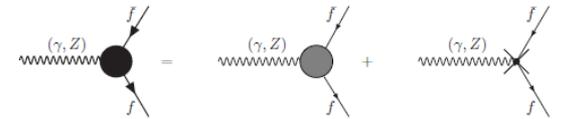


Figure A.6. Off-shell Zff and γff vertices

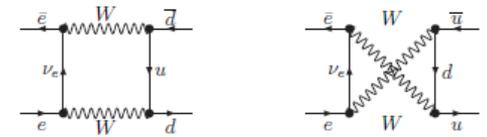


Figure A.7. The WW boxes

etc. etc.

Constructing wt^{EW} : per-event weight

ERW and Z.Was,
arXiv: 1808.08616

Define per event electroweak weight

$$wt^{EW} = \sigma_{\text{Born}}^{\text{new}} / \sigma_{\text{Born}}^{\text{old}}$$

Approach developed
in TauSpinner,
arXiv:1802.05459

$$wt^{EW} = \frac{d\sigma_{\text{Born}+EW}(x_1, x_2, \hat{s}, \cos\theta, s_W^2)}{d\sigma_{\text{Born}}(x_1, x_2, \hat{s}, \cos\theta, s_W^2)}$$

$$d\sigma_{\text{Born}}(x_1, x_2, \hat{s}, \cos\theta, s_W^2) = \sum_{q_f, \bar{q}_f} [f^{q_f}(x_1, \dots) f^{\bar{q}_f}(x_2, \dots) d\sigma_{\text{Born}}^{q_f \bar{q}_f}(\hat{s}, \cos\theta, s_W^2) + f^{q_f}(x_2, \dots) f^{\bar{q}_f}(x_1, \dots) d\sigma_{\text{Born}}^{\bar{q}_f q_f}(\hat{s}, -\cos\theta, s_W^2)]$$

$x_1, x_2, \cos\theta$ (symmetrised)
calculated using 4-momenta
of outgoing leptons;
asymmetry in sign of $\cos\theta$
from weighted average
over PDFs

Allows to reweight MC event generated between different EW LO scheme and to **Improved Born Approximation** in EW scheme used for form-factors calculation.

Constructing wt^{EW} : EW Improved Born (IBA)

ERW and Z.Was,
arXiv: 1808.08616

$$\mathcal{A}^{Born+EW} = \frac{\alpha}{s} \left\{ \begin{aligned} & [\bar{u}\gamma^\mu v g_{\mu\nu} \bar{v}\gamma^\nu u] \cdot (q_\ell \cdot q_f) \Gamma_{V\Pi} \chi_\gamma(s) \\ & + [\bar{u}\gamma^\mu v g_{\mu\nu} \bar{v}\gamma^\nu u \cdot (v_\ell \cdot v_f \cdot vv_{\ell f}) + \bar{u}\gamma^\mu v g_{\mu\nu} \bar{v}\gamma^\nu \gamma^5 u \cdot (v_\ell \cdot a_f) \\ & + \bar{u}\gamma^\mu \gamma^5 v g_{\mu\nu} \bar{v}\gamma^\nu u \cdot (a_\ell \cdot v_f) + \bar{u}\gamma^\mu \gamma^5 v g_{\mu\nu} \bar{v}\gamma^\nu \gamma^5 u \cdot (a_\ell \cdot a_f)] \cdot Z_{V\Pi} \chi_Z(s) \end{aligned} \right\}$$

$$\chi_\gamma(s) = 1$$

$$\chi_Z(s) = \frac{G_\mu \cdot M_Z^2 \cdot \Delta^2}{\sqrt{2} \cdot 8\pi \cdot \alpha} \cdot \frac{s}{s - M_Z^2 + i \cdot \Gamma_Z \cdot M_Z}$$

$$v_\ell = (2 \cdot T_3^\ell - 4 \cdot q_\ell \cdot s_W^2) \cdot K_\ell(s, t) / \Delta$$

$$v_f = (2 \cdot T_3^f - 4 \cdot q_f \cdot s_W^2) \cdot K_f(s, t) / \Delta$$

$$a_\ell = (2 \cdot T_3^\ell) / \Delta$$

$$a_f = (2 \cdot T_3^f) / \Delta$$

$$\Delta = \sqrt{16 \cdot s_W^2 \cdot (1 - s_W^2)}$$

$$Z_{V\Pi} = \rho_{e,f}(s, t)$$

**EW form-factors, functions of $(s, t) = (m_{\Pi}, \cos\theta)$
Calculated with Dizeit 6.21 library.**

$$\Gamma_{V\Pi} = \frac{1}{2 - (1 + \Pi_{\gamma\gamma}(s))}$$

Vacuum polarisation corrections, used low-energy experiment input.

Warning: problem for analytic continuation.

$$vv_{\ell f} = \frac{1}{v_\ell \cdot v_f} \left[(2 \cdot T_3^\ell)(2 \cdot T_3^f) - 4 \cdot q_\ell \cdot s_W^2 \cdot K_f(s, t) (2 \cdot T_3^\ell) - 4 \cdot q_f \cdot s_W^2 \cdot K_\ell(s, t) (2 \cdot T_3^f) + (4 \cdot q_\ell \cdot s_W^2)(4 \cdot q_f \cdot s_W^2) K_{\ell f}(s, t) \right] \frac{1}{\Delta^2}$$

EW schemes

- **LEP legacy: ($\alpha(0)$, G_μ , M_Z)**

D. Bardin et al.
arXiv:9908433

- Inputs are very precisely measured physics quantities
- M_Z , M_W are on-shell masses
- Genuine EW and lineshape corrections in form of (multiplicative) form-factors to LO couplings

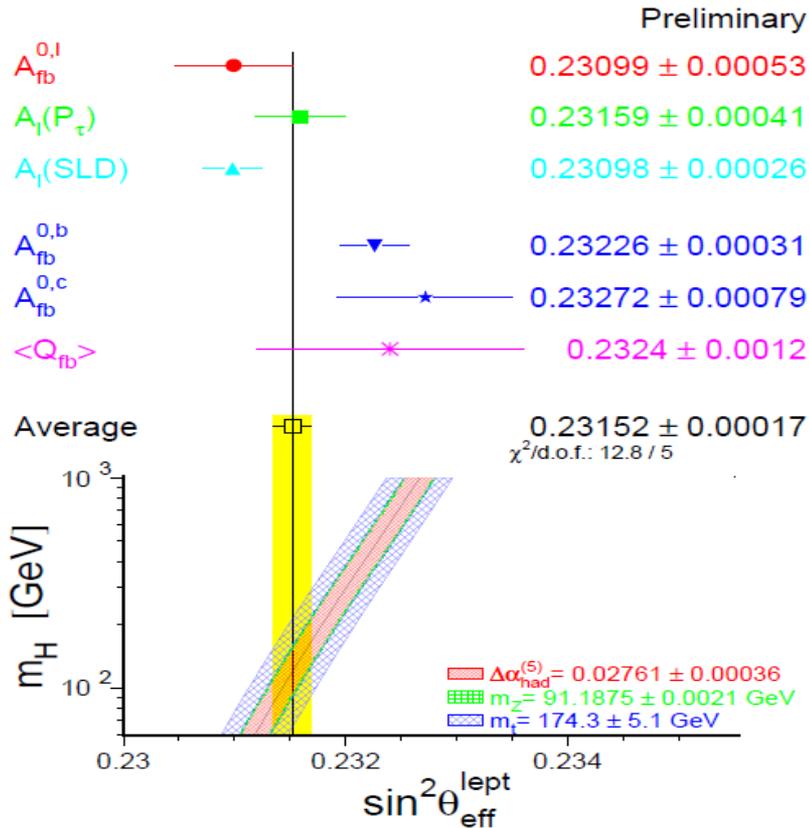
- **LHC paradigm: (G_μ , M_Z , M_W).**

S. Dittmaier, M. Huber
arXiv:0911.2329

- M_Z , M_W are pole-masses or complex masses.
- Absorbs most of universal corrections into lowest-order couplings
- Higher-order corrections redefine couplings in non-multiplicative manner

Re-discovered impact of $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$

arXiv:hep-ex/0112021



Recent measurements:

arXiv: 1706.09436

$$\Delta\alpha_{\text{had}}^{(5)}(M_Z) = 0.02753 \pm 0.00009$$

arXiv: 1711.06089

$$\Delta\alpha_{\text{had}}^{(5)}(M_Z) = 0.02774 \pm 0.00016$$

$$= 0.02752 \pm 0.00012$$

arXiv: 1802.02995

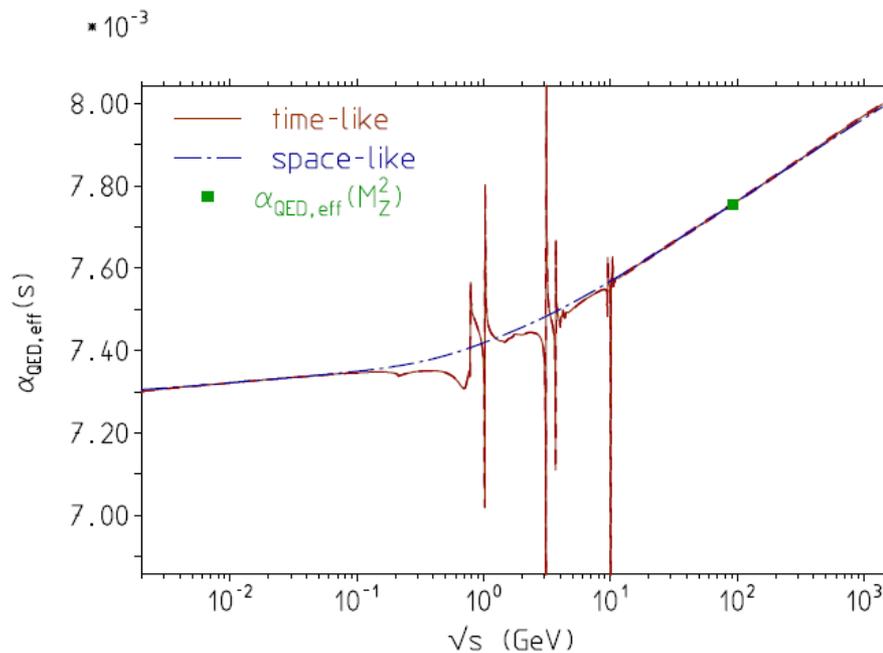
$$\Delta\alpha_{\text{had}}^{(5)}(M_Z) = 0.02761 \pm 0.00011$$

± 0.00036 on $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$ \rightarrow ± 0.00013 on $\sin^2\theta_{\text{eff}}^{\text{lept}}$

$\alpha_{\text{QED}}(s)$ and $\sin^2\theta_W(s)$

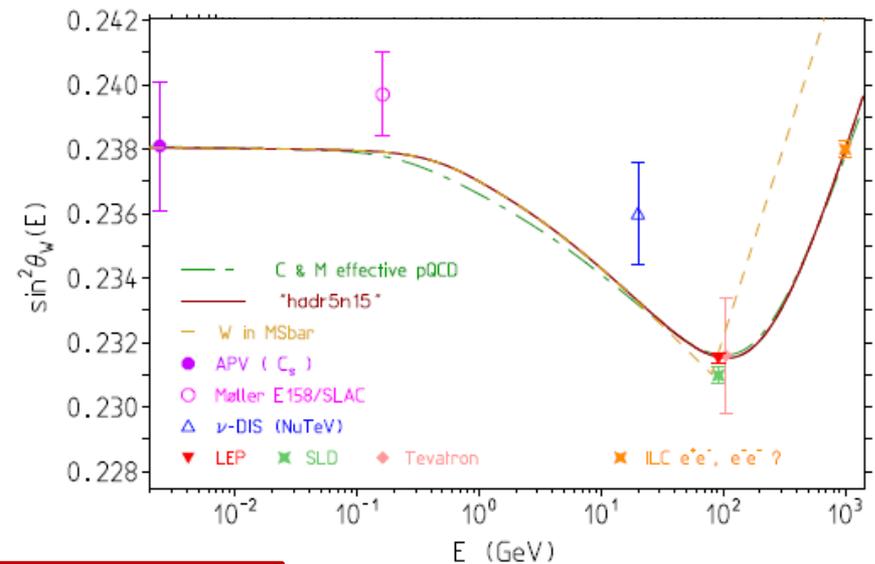
F. Jegerlehner
FCCee WS 2019

$$\alpha(s) = \frac{\alpha}{1-\Delta\alpha(s)}; \quad \Delta\alpha(s) = \Delta\alpha_{\text{lep}}(s) + \Delta\alpha_{\text{had}}^{(5)}(s) + \Delta\alpha_{\text{top}}(s)$$



$$\sin^2\Theta_i \cos^2\Theta_i = \frac{\pi\alpha}{\sqrt{2}G_\mu M_Z^2} \frac{1}{1-\Delta r_i}$$

$$\Delta r_i = \Delta r_i(\alpha, G_\mu, M_Z, m_H, m_{f \neq t}, m_t)$$



$$\Delta\alpha_{\text{hadrons}}^{(5)}(M_Z^2) = 0.027738 \pm 0.000158 [0.027523 \pm 0.000119]$$

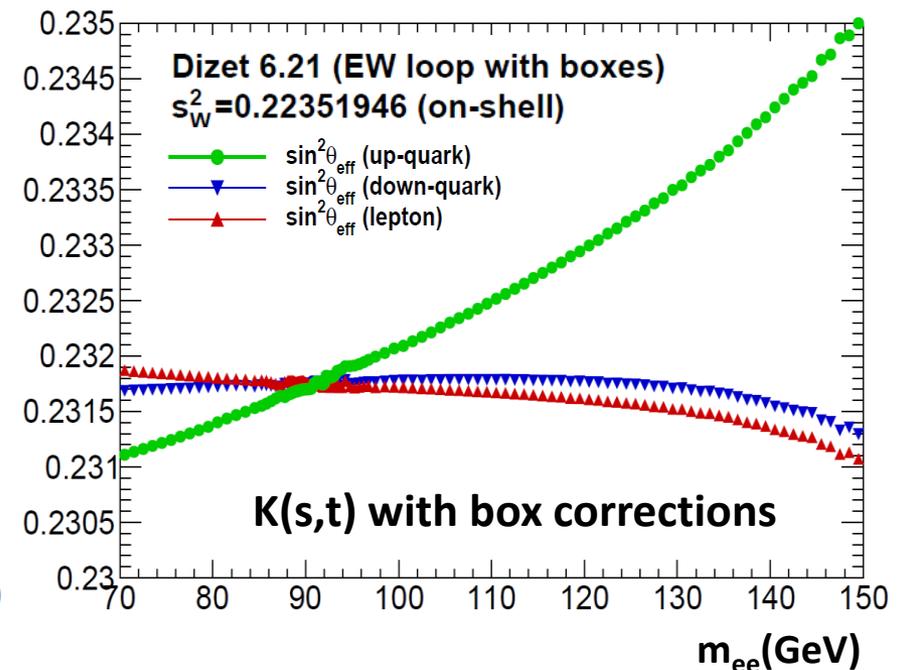
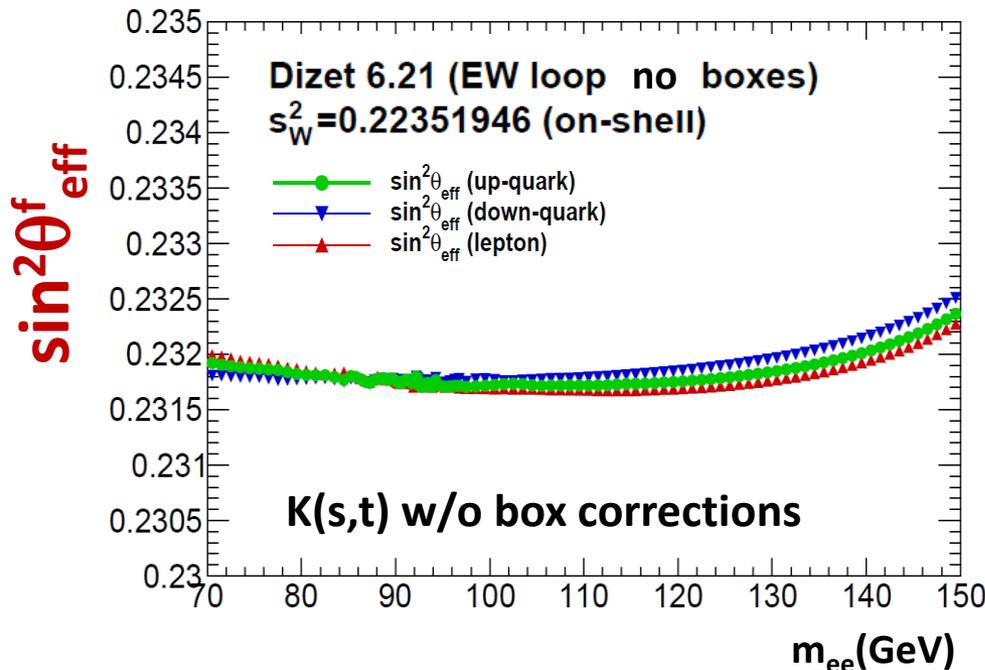
Effective weak mixing angle @ LHC

Here convoluted with line-shape and $\cos\theta^*$ distribution of MC events.

$$\sin^2 \theta_W^{eff} = \underset{\uparrow(s,t)}{\text{Re}(K_Z^f) s_W^2} + I_f^2 \quad I_f^2 = \alpha^2(s) \frac{35}{18} \left[1 - \frac{8}{3} \text{Re}(K_Z^f) s_W^2 \right] = \sim 10^{-4}$$

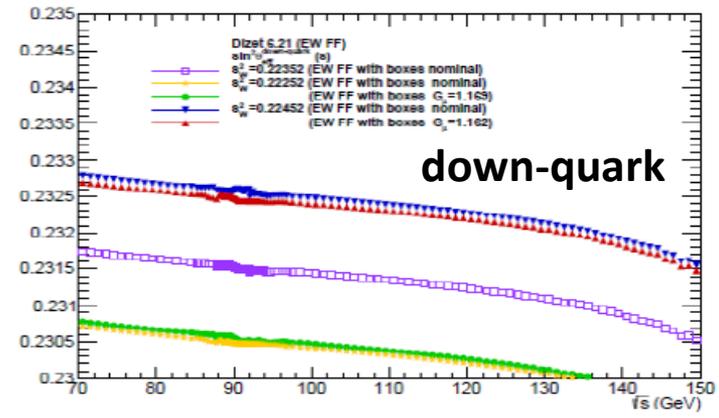
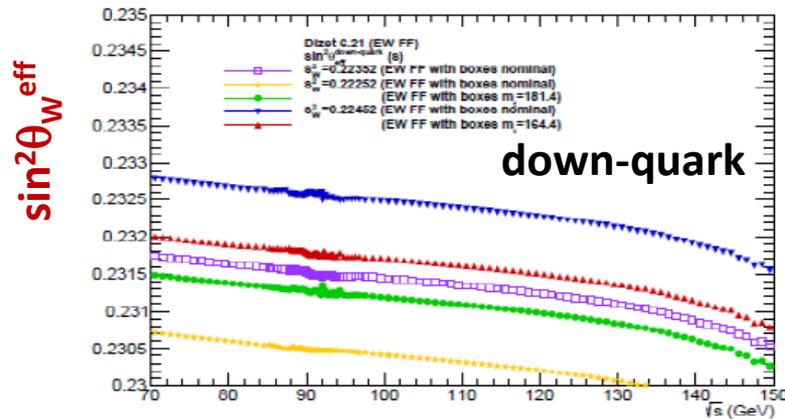
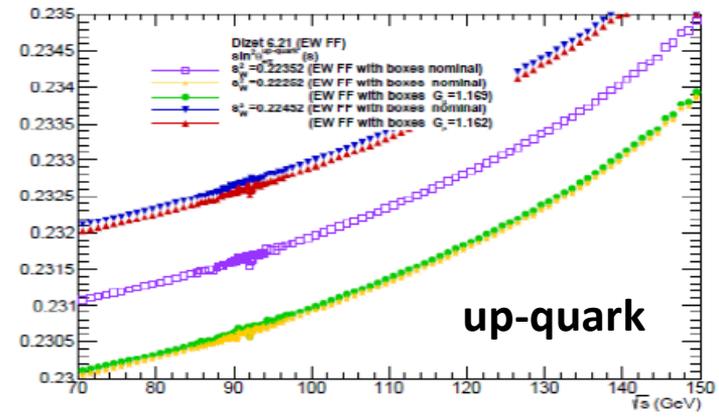
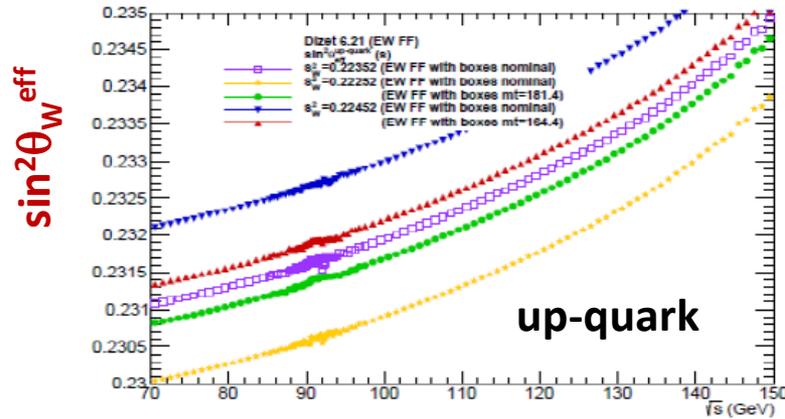
$$s_W^2 = 1 - M_W^2/M_Z^2$$

$\sin^2 \theta_W^{eff}$ defines ratio of couplings $g_Z^f = \frac{v_f}{a_f} = 1 - 4|q_f|(K_Z^f s_W^2 + I_f^2)$



Effect on $\sin^2\theta_W^{\text{eff}}$ from recalculating form-factors

Variation in m_t or G_μ corresponding to $\pm 100 \cdot 10^{-5}$ variation in s_W^2



- | | | | |
|--|-----------------------------------|--|--|
| | $s_W^2 = 0.22352$ (EW FF nominal) | | $s_W^2 = 0.22352$ (EW FF nominal) |
| | $s_W^2 = 0.22252$ (EW FF nominal) | | $s_W^2 = 0.22252$ (EW FF nominal) |
| | (EW FF $m_t = 181.4$) | | (EW FF $G_\mu = 1.169 \cdot 10^{-5}$) |
| | $s_W^2 = 0.22452$ (EW FF nominal) | | $s_W^2 = 0.22452$ (EW FF nominal) |
| | (EW FF $m_t = 164.4$) | | (EW FF $G_\mu = 1.162 \cdot 10^{-5}$) |

How to scan s_W^2 around its nominal value

We want to stay consistent within SM QED/EW field theory

optME=1 : change $s_W^2 + \delta_{s_{2W}}$ as if due to new V-coupling X-particle
no SM parameters modifications

$$v_\ell = (2 \cdot T_3^\ell - 4 \cdot q_\ell \cdot (s_W^2 + \delta_{s_{2W}})) \cdot K_\ell(s, t) / \Delta$$

$$v_f = (2 \cdot T_3^f - 4 \cdot q_f \cdot (s_W^2 + \delta_{s_{2W}})) \cdot K_f(s, t) / \Delta$$

$$vv_{\ell f} = \frac{1}{v_\ell \cdot v_f} [(2 \cdot T_3^\ell)(2 \cdot T_3^f) - 4 \cdot q_\ell \cdot (s_W^2 + \delta_{s_{2W}}) \cdot K_f(s, t)(2 \cdot T_3^\ell) - 4 \cdot q_f \cdot (s_W^2 + \delta_{s_{2W}}) \cdot K_\ell(s, t)(2 \cdot T_3^f) + (4 \cdot q_\ell \cdot s_W^2)(4 \cdot q_f \cdot s_W^2)K_{\ell f}(s, t) + 2 \cdot (4 \cdot q_\ell)(4 \cdot q_f) \cdot s_W^2 \cdot \delta_{s_{2W}} K_{\ell f}(s, t)] \frac{1}{\Delta^2}$$

optME=2 : change s_W^2 as if due to changes in the SM input parameters, like m_t or G_μ .
Followed (option) by form-factors recalculation.

optME=3 : add δ_V directly in the vertex, almost equivalent to optME=1 for $s_W^2 + \delta_V/K_\ell$ and $s_W^2 + \delta_V/K_f$

$$v_\ell = (2 \cdot T_3^\ell - 4 \cdot q_\ell \cdot (s_W^2 \cdot K_\ell(s, t) + \delta_V)) / \Delta$$

$$v_f = (2 \cdot T_3^f - 4 \cdot q_f \cdot (s_W^2 \cdot K_f(s, t) + \delta_V)) / \Delta$$

$$vv_{\ell f} = \frac{1}{v_\ell \cdot v_f} [(2 \cdot T_3^\ell)(2 \cdot T_3^f) - 4 \cdot q_\ell \cdot (s_W^2 \cdot K_f(s, t) + \delta_V)(2 \cdot T_3^\ell) - 4 \cdot q_f \cdot (s_W^2 \cdot K_\ell(s, t) + \delta_V)(2 \cdot T_3^f) + (4 \cdot q_\ell \cdot s_W^2)(4 \cdot q_f \cdot s_W^2)K_{\ell f}(s, t) + 2 \cdot (4 \cdot q_\ell)(4 \cdot q_f) \cdot s_W^2 \cdot K_{\ell f}(s, t) \cdot \delta_V] \frac{1}{\Delta^2}$$

How to scan s_W^2

April 2018

Around Z-pole slope of $A_4(s_W^2)$ or $A_{fb}(s_W^2)$ depends on how variation is done in practice.

Sensitivity:

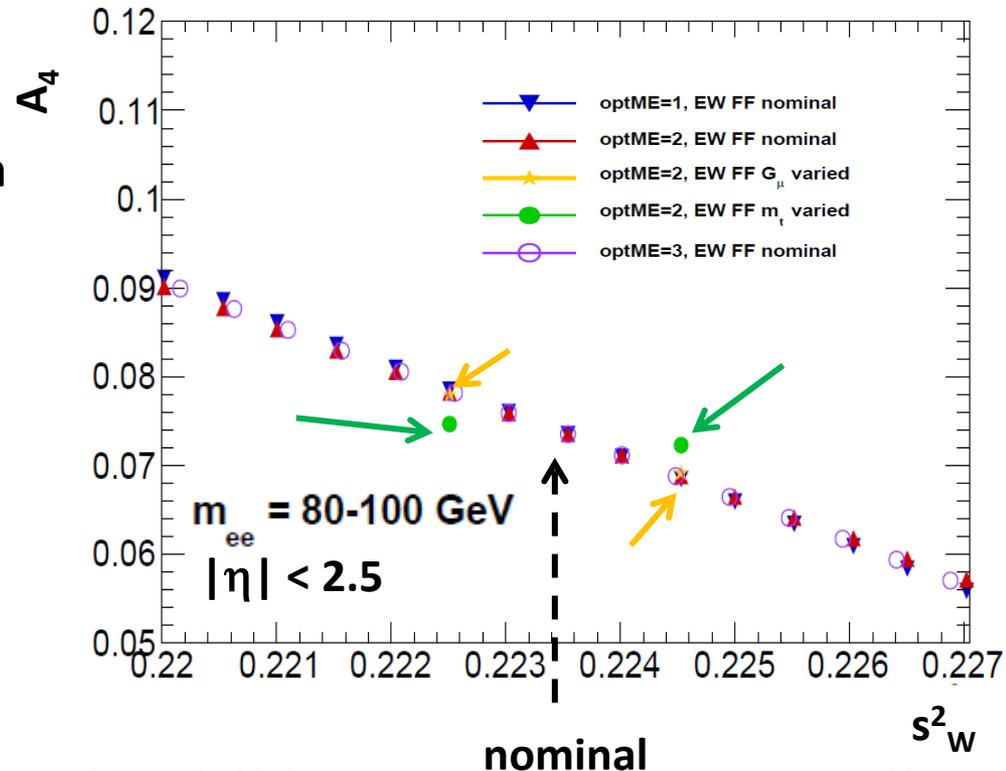
1% relative uncertainty on A_{fb} or A_4 corresponds to **20 * 10⁻⁵** uncertainty on s_W^2 (except case of m_t variation).

Sensitivity in case of m_t variation much weaker due to compensating corrections from recalculated FF

Recalculating FF makes procedure formally gauge-invariant, but requires m_t or G_μ unreasonably far from measured values.

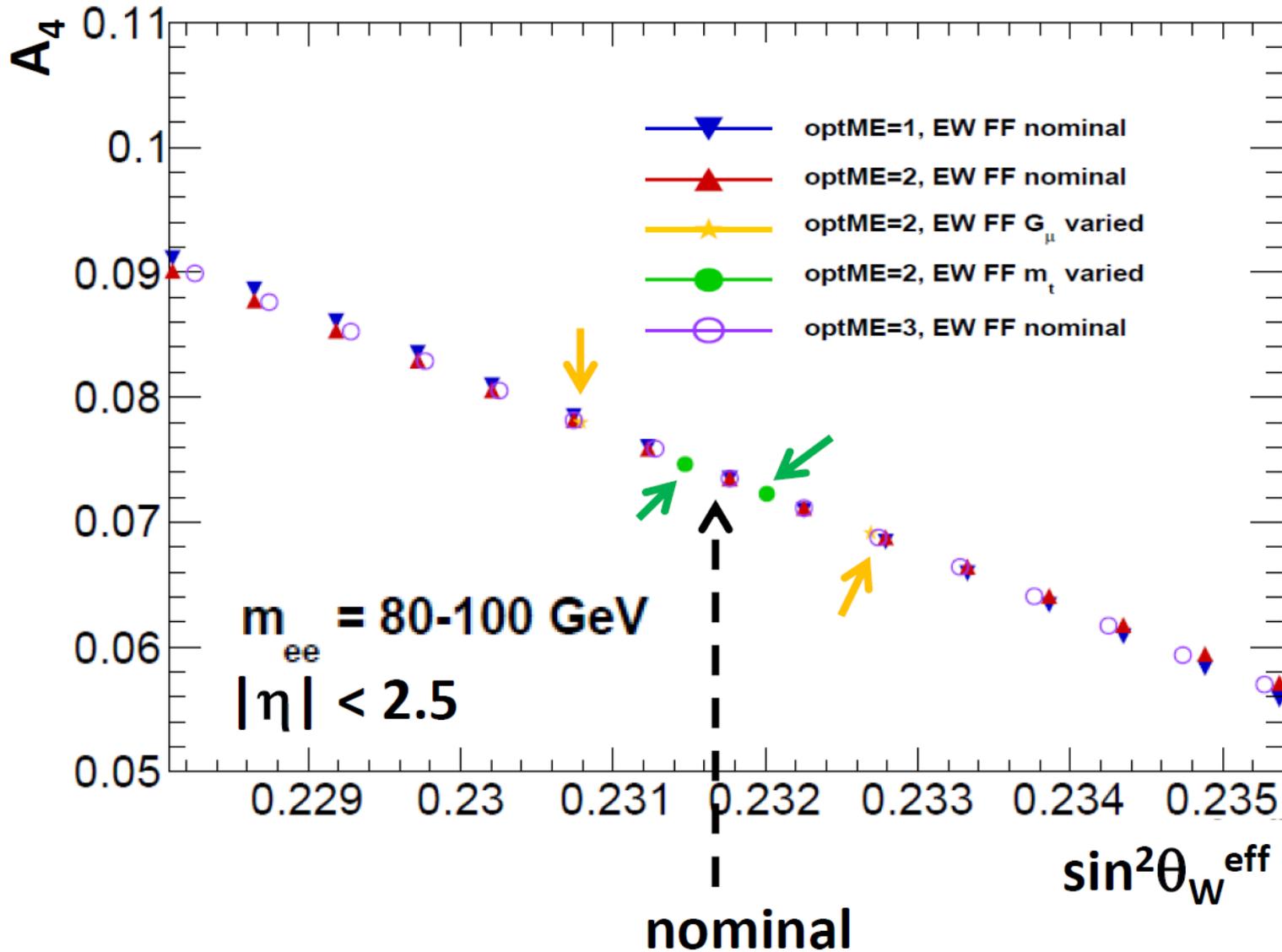
Scan was performed using per event weight w_t^{EW} introduced earlier.

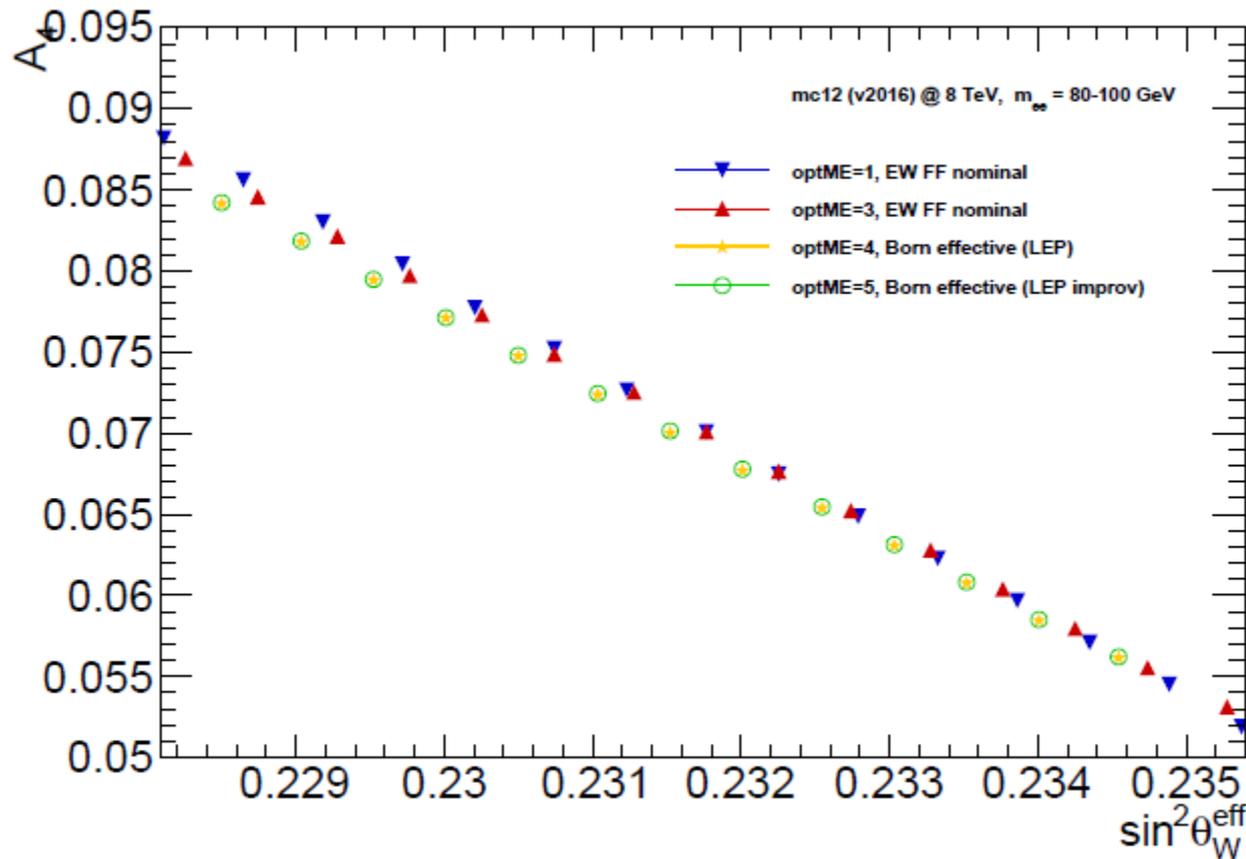
@ 8 TeV	A_4	$\Delta A_4/A_4 = 0.01$ Δs_W^2
<i>optME=2</i>	0.0735	0.00017
<i>optME=2, G_μ varied</i>		0.00019
<i>optME=2, m_t varied</i>		0.00070
<i>optME=3</i>		0.00017
<i>optME=1</i>		0.00016



How to scan $\sin^2\theta_{\text{eff}}$

April 2018





Scan with EW LO:

optME = 4,5 use Born EW LO amplitude, parametrised with $\sin^2\theta_{\text{eff}}$ and $\alpha(M_Z)$
Shift by $20 \cdot 10^{-5}$ in predicted $\sin^2\theta_{\text{eff}}$ corresponding to the same A_{fb} or A_4 .

EW schemes: LEP and LHC paradigms

LEP

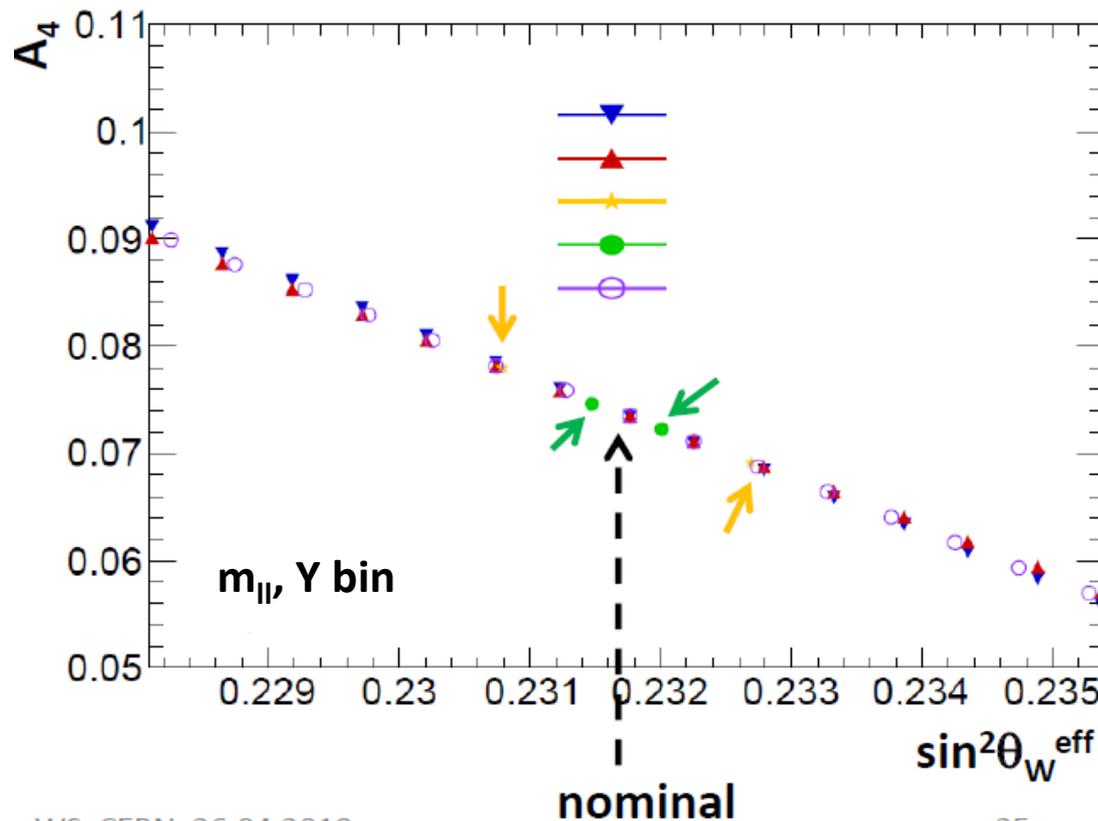
- EW scheme: „on-mass-schell” regularisation; input: $(\alpha(0), G_F, M_Z)$, the most precisely known quantities:
 $\Delta\alpha/\alpha \sim 3.6 \cdot 10^{-9}$; $\Delta G_F/G_F \sim 8.6 \cdot 10^{-6}$; $\Delta M_Z/M_Z \sim 2.4 \cdot 10^{-5}$
- TH precision @Z-pole: 4 MeV on m_W ; $5 \cdot 10^{-5}$ on $\sin^2\theta_{\text{eff}}$
- Parametric precision on $\sin^2\theta_{\text{eff}}$: $4 \cdot 10^{-5}$ (dominated by $\Delta\alpha_{\text{had}}$)
- Used and developed still today: GFitter, FCCee preparation.

LHC

- EW scheme: „pole-mass” regularisation; input: (G_F, M_Z, M_W)
 $\Delta G_F/G_F \sim 8.6 \cdot 10^{-6}$; $\Delta M_Z/M_Z \sim 2.4 \cdot 10^{-5}$; $\Delta M_W/M_W \sim 1.9 \cdot 10^{-4}$
- **Relations to EWPO's: not established.**
- EW corrections complete at EW NLO, only some at EW NLO+HO
- **TH precision @Z-pole: not established.**
- **Parametric precision: ± 15 MeV on M_W -> $30 \cdot 10^{-5}$ on $\sin^2\theta_W$**

How to scan s^2w : outlook

- Similarly here: compare different options, on vertical scale either A_{FB} or A_4 . Expose, how other SM parameters need to vary in you scheme.
- We need it with granularity $5 \cdot 10^{-5}$, in a range $\pm 100 \cdot 10^{-5}$ around central value, in different mass and Y bins.



EW LO schemes: running/fixed width

Running and fixed Z-boson width in the propagator: taking into account photon-loop corrections to Γ_Z

- Fixed width $\chi_Z(s) = \frac{1}{s - M_Z^2 + i \cdot \Gamma_Z \cdot M_Z}$.

- Running width (LEP legacy)

$$\chi'_Z(s) = \frac{1}{s - M_Z^2 + i \cdot \Gamma_Z \cdot s/M_Z}$$



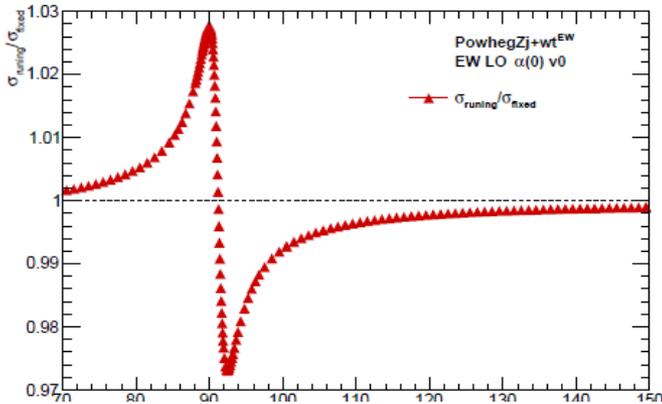
$$\begin{aligned} \chi'_Z(s) &= \frac{1}{s(1 + i \cdot \Gamma_Z/M_Z) - M_Z^2} \\ &= \frac{(1 - i \cdot \Gamma_Z/M_Z)}{s(1 + \Gamma_Z^2/M_Z^2) - M_Z^2(1 - i \cdot \Gamma_Z/M_Z)} \\ &= \frac{1}{(1 + \Gamma_Z^2/M_Z^2) \left(s - \frac{M_Z^2}{1 + \Gamma_Z^2/M_Z^2} + i \cdot \frac{\Gamma_Z M_Z}{1 + \Gamma_Z^2/M_Z^2} \right)} \\ &= N_Z \frac{1}{s - M_Z'^2 + i \Gamma_Z' M_Z'} \\ M_Z' &= \frac{M_Z}{\sqrt{1 + \Gamma_Z^2/M_Z^2}} \\ \Gamma_Z' &= \frac{\Gamma_Z}{\sqrt{1 + \Gamma_Z^2/M_Z^2}} \\ N_Z &= \frac{(1 - i \cdot \Gamma_Z/M_Z)}{(1 + \Gamma_Z^2/M_Z^2)} = \frac{(1 - i \cdot \Gamma_Z'/M_Z')}{(1 + \Gamma_Z'^2/M_Z'^2)} \end{aligned}$$

Both equivalent if redefined parameters m_Z , Γ_Z , N_Z (normalization). Change in the normalisation can (?) be absorbed into G_μ redefinition. In case of „pole” convention (last slide) it was absorbed into α .

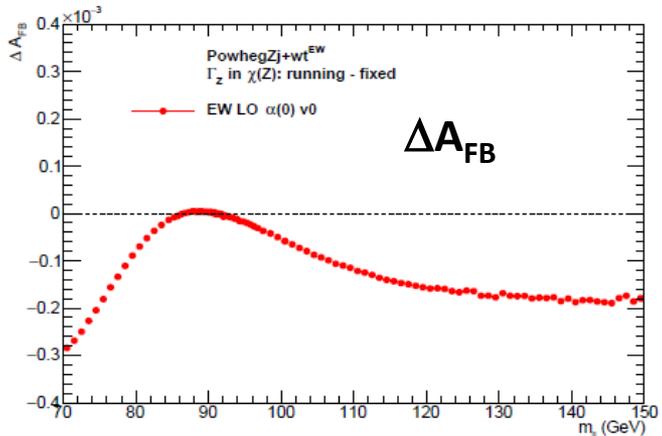
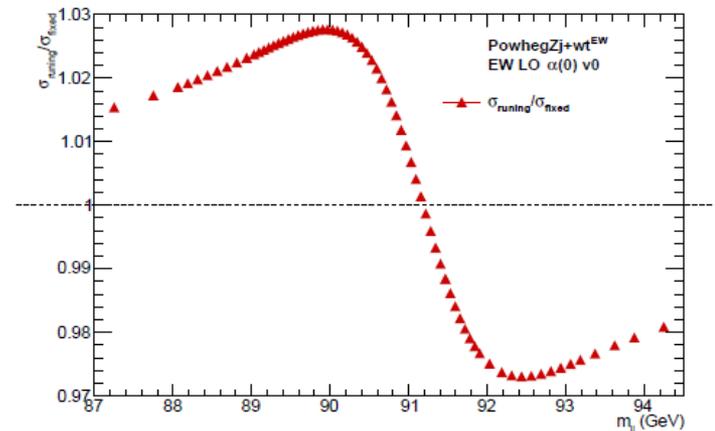
EW LO schemes: running/fixed width

Nowdays MC's are using fixed-width propagators with on-shell $M_Z = 91.1876$.
 How does it affect predictions, if running-> fixed without reparametrizing?
 $\Delta A_{FB} (m_{II}=80-100 \text{ GeV}) = 0.0005$ (thanks to D. Walker for pointing it out)

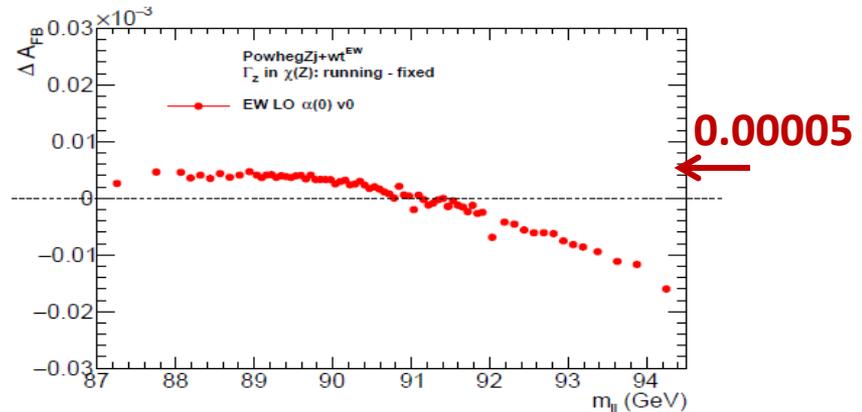
Lineshape: ratio



zoom



zoom



Dictionary

EW LO Born (LO = lowest order):

tree-level vertex and propagator of the Z/γ^* bosons, setting of SM EW parameters defines the EW scheme.

EW effective Born:

tree-level vertex and propagator of the Z/γ^* bosons, EW couplings: $\alpha(m_Z)$, $\sin^2\theta_W(m_Z)$, m_Z , set of best measured values.

EW Improved Born Approximation (IBA):

tree-level vertex and propagator of the Z/γ^* bosons, EW couplings and propagators multiplied by form-factors dependent on the scattering angle of the lepton (choice of frame) and virtuality of Z/γ^* .

QED/EW corrections: D. Bardin et al. arXiv:9908433

separate set of „QED corrections” and „genuine EW + lineshape corrections”

Tree level

Standard Model relations:

$$G_F = \frac{\pi\alpha}{\sqrt{2}m_W^2 \sin^2 \theta_W^{\text{tree}}},$$

$$\rho_0 = \frac{m_W^2}{m_Z^2 \cos^2 \theta_W^{\text{tree}}} = 1$$

Powheg_ew: EW LO, NLO, NLO+HO

Status:

September 2018

F. Piccinini et al.

	EW order	$m_{ee} = 89 - 93$ GeV	$m_{ee} = 80 - 100$ GeV	$m_{ee} = 70 - 120$ GeV
$A_{FB} \alpha(0) v0$	LO	0.06691361	0.06392369	0.06253754
$A_{FB} \alpha(0) v1$	LO	0.04653886	0.04343789	0.04212883
$A_{FB} G_\mu$	LO	0.04653886	0.04343789	0.04212883
$A_{FB} \alpha(0) v1$	NLO	0.03004289	0.02690785	0.02569858
$A_{FB} G_\mu$	NLO	0.02905841	0.02592168	0.02471918
$A_{FB} \alpha(0) v1$	NLO+HO	0.03083234	0.02770533	0.02649700
$A_{FB} G_\mu$	NLO+HO	0.03090286	0.02777783	0.02656851

A_{FB}

$\Delta A_{FB} \alpha(0) v1$	NLO-LO	-0.0164959	-0.0165300	-0.0164302
$\Delta A_{FB} G_\mu$	NLO-LO	-0.0174805	-0.0175162	-0.0174096

ΔA_{FB} (NLO - LO)

$\Delta A_{FB} \alpha(0) v1$	NLO+HO-LO	-0.0157065	-0.0157326	-0.0156318
$\Delta A_{FB} G_\mu$	NLO+HO-LO	-0.0156360	-0.0156596	-0.0155603

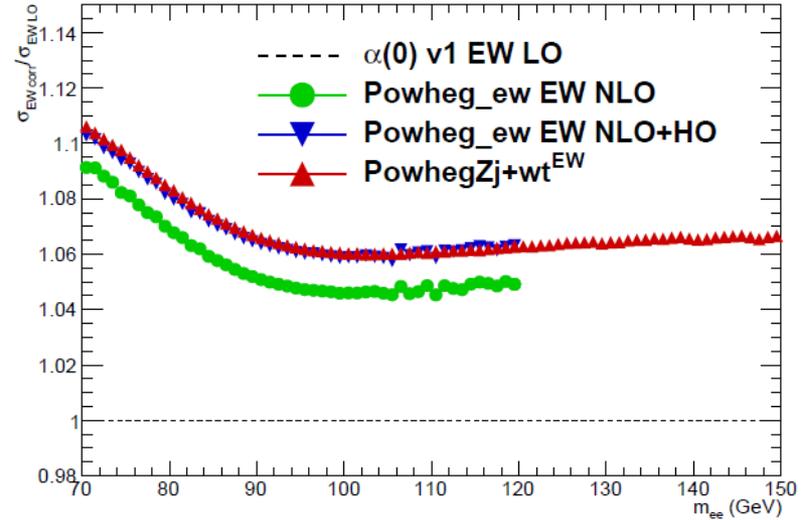
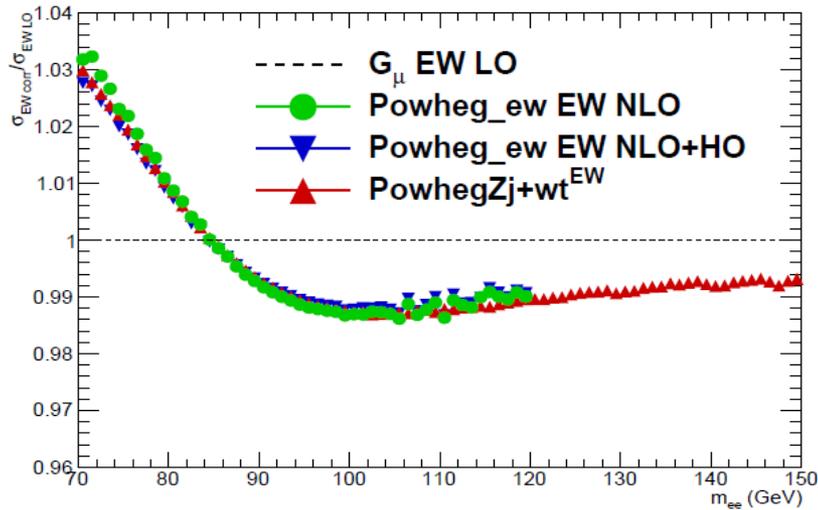
ΔA_{FB} (NLO+HO - LO)

ΔA_{FB}	EW order	$m_{ee} = 89 - 93$ GeV	$m_{ee} = 80 - 100$ GeV	$m_{ee} = 70 - 120$ GeV
$\alpha(0) v1 - \alpha(0) v0$	LO	-0.020375	-0.020486	-0.020487
$G_\mu - \alpha(0) v0$	LO	-0.020375	-0.020486	-0.0204871
$G_\mu - \alpha(0) v1$	LO	0.0	0.0	0.0
$G_\mu - \alpha(0) v1$	NLO	-0.00098	-0.00098	-0.00098
$G_\mu - \alpha(0) v1$	NLO + HO	-0.00007	-0.00007	-0.00007

ΔA_{FB} between
EW schemes at
LO, NLO, NLO+HO

Better than 0.0001 agreement on A_{FB} at NLO+HO between two EW schemes !

Cross-section: EW LO, NLO, NLO+HO



Powheg_ew

	EW order	$m_{ee} = 89 - 93 \text{ GeV}$
$\alpha(0) \text{ v1}$	NLO/LO	1.050350
G_μ	NLO/LO	0.991230
$\alpha(0) \text{ v1}$	NLO+HO/LO	1.063247
G_μ	NLO+HO/LO	0.991038

	EW order	$m_{ee} = 89 - 93 \text{ GeV}$
Powheg_ew	NLO+HO/LO	
$\alpha(0) \text{ v1}$		1.06325
G_μ		0.99104
PowhegZj+wt ^{EW}	NLO+HO/LO	
$\alpha(0) \text{ v0}$		0.96452
$\alpha(0) \text{ v1}$		1.06506
G_μ		0.99167

In G_μ scheme, NLO corrections < 1%

HO corrections < 0.02%