

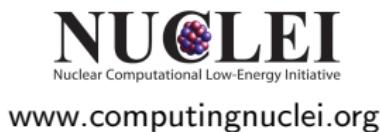
Electroweak interactions in nuclei

Stefano Gandolfi

Los Alamos National Laboratory (LANL)

NuPhys2019: Prospects in Neutrino Physics

16-18 December 2019, Cavendish Conference Centre, London, UK



www.computingnuclei.org



National Energy Research
Scientific Computing Center



At "nuclear" energies, understanding neutrino-nucleus interactions very challenging and important!

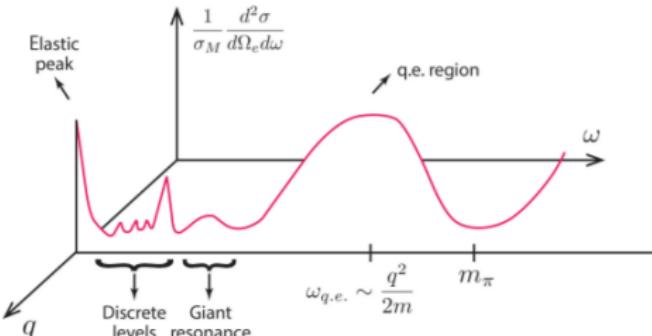
Understanding Nuclei:

- Nuclear interactions and structure
- Electroweak processes

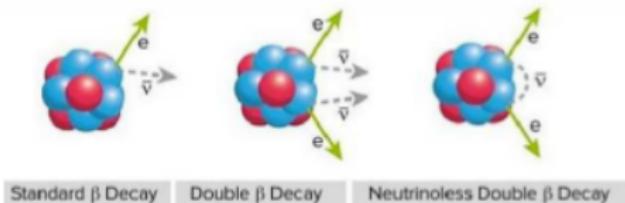
Relevance:

- Neutrino scattering in nuclei (neutrino oscillation experiments)
- Neutrinoless Double Beta Decay
- Neutrino interactions in supernovae and neutron stars, nucleosynthesis

We need a coherent picture of ν -nucleus interactions

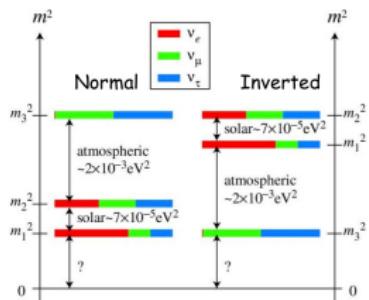


- $\omega \approx$ few MeV, $q \approx 0$: $\beta-$ and $\beta\beta-$ decays
- $\omega \approx$ few MeV, $q \approx 10^2$ MeV: Neutrinoless $\beta\beta-$ decays
- $\omega \leq$ tens MeV: Astrophysics
- $\omega \approx 10^2$ MeV: Accelerator neutrinos, ν -nucleus scattering



Motivation

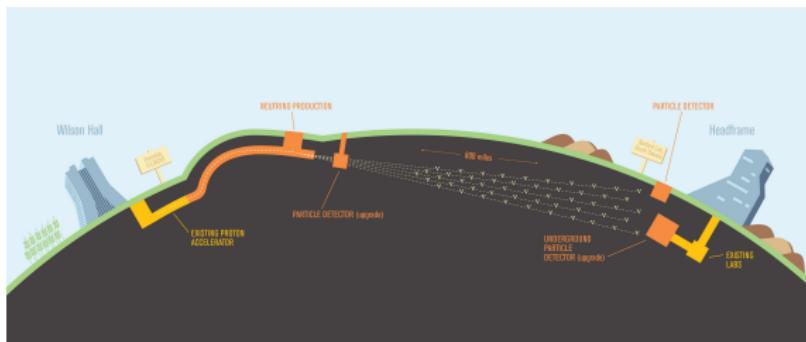
DUNE - Deep Underground Neutrino Experiment - to measure neutrino oscillations and CP violation



Simplified 2 flavors evolution (CP violation non included):

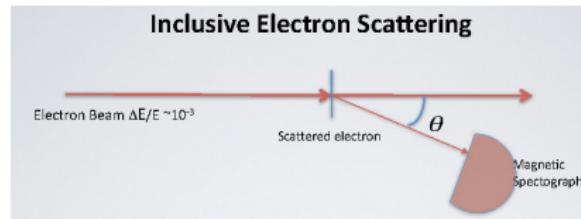
$$P_{\alpha \rightarrow \beta} = \sin^2(2\theta_{\alpha\beta}) \sin^2 \left(1.267 \frac{\Delta m_{\alpha\beta}^2 L}{E} \frac{\text{GeV}}{\text{eV}^2 \text{km}} \right)$$

Need to know E !

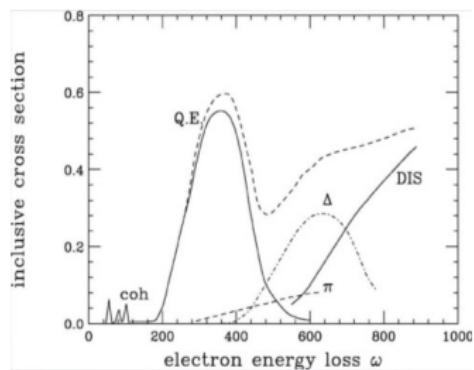


Introduction: electron energy and cross-section

Electron energy easy to know:



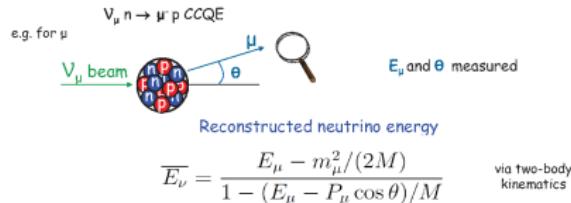
Electron scattering in nuclei:



Benhar, Day, Sick, RMP (2008)

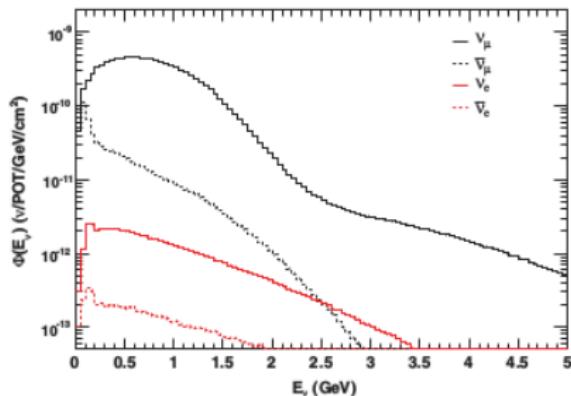
Introduction: neutrino energy and cross-section

E_ν difficult to reconstruct. Example: CCQE process



Neutral current process even more difficult.

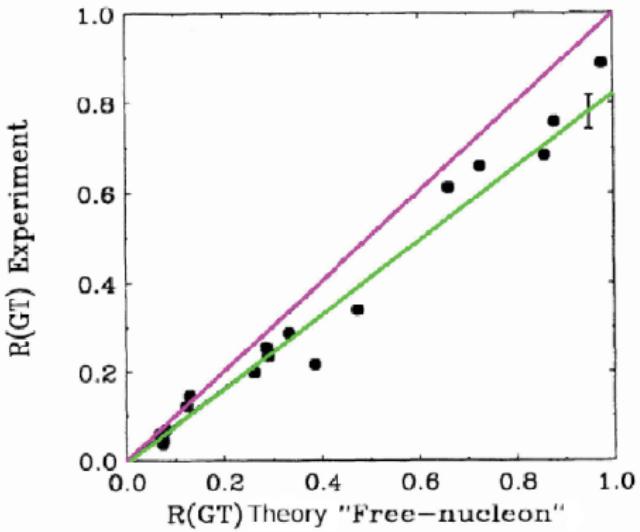
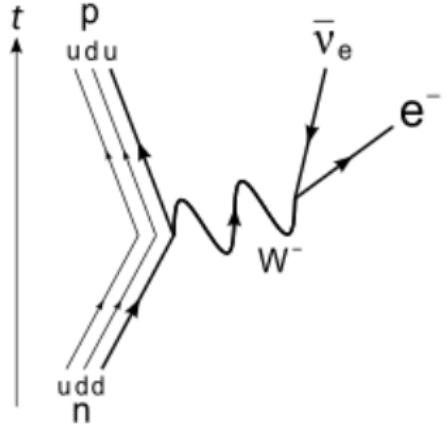
Simulation of neutrino energy distribution:



MiniBooNE Coll., PRD (2009)

Knowledge of cross-section
+ near detector
= determination of E_ν

The “quenching”- g_A problem



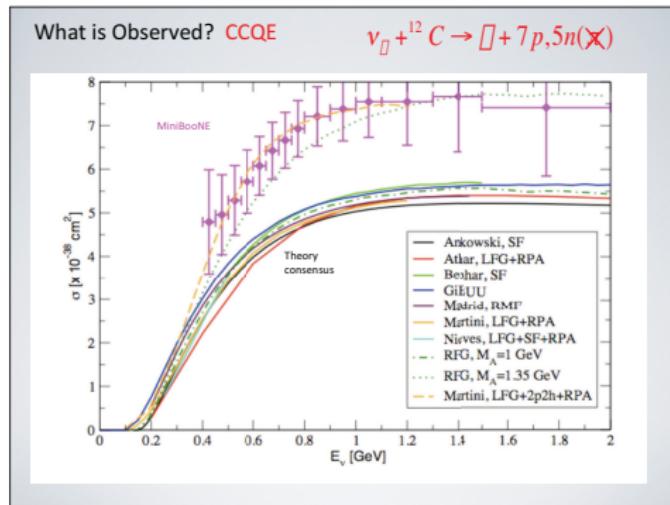
$$g_A^{\text{eff}} \simeq 0.70 g_A$$

Chou et al., PRC 47, 163 (1993)

What's the origin (or is there a need) of g_A quenching?

Charge-change quasi-elastic cross-section in ^{12}C

Experimental vs theory disagreement:



Alvarez-Ruso arXiv:1012.3871

Currents inconsistent with the Hamiltonian.

Nucleon-nucleon correlations and two-body processes approximately accounted for. These models do not describe electron-scattering!!!

Need of g_A “unquenching” ???

Nuclear Hamiltonian

Model: non-relativistic nucleons strongly interacting with a nucleon-nucleon (NN) and three-nucleon interaction (TNI).

$$H = -\frac{\hbar^2}{2m} \sum_{i=1}^A \nabla_i^2 + \sum_{i < j} v_{ij} + \sum_{i < j < k} V_{ijk}$$

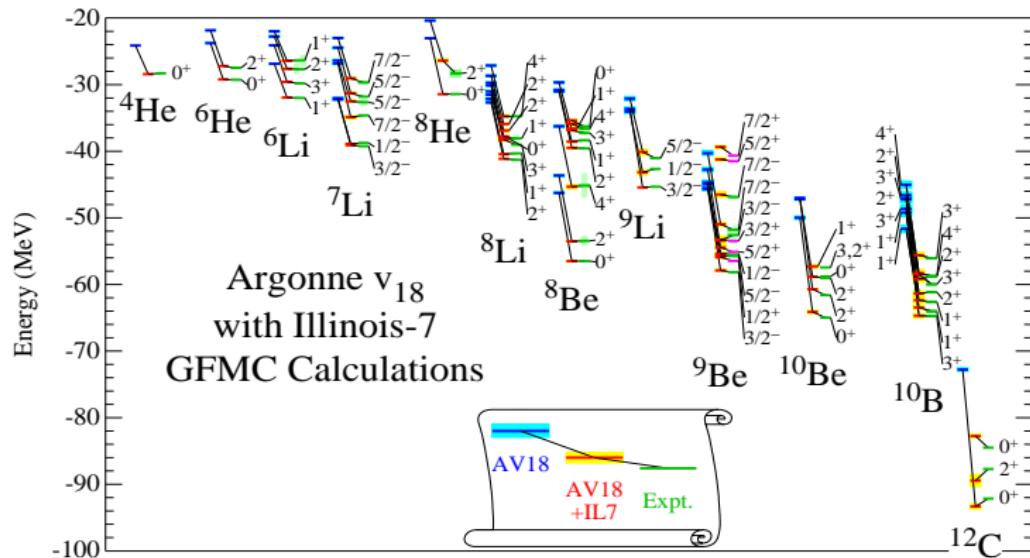
v_{ij} NN fitted on scattering data and TNI to properties of light nuclei.

Quantum Monte Carlo methods used to solve the many-body Schrödinger equation in imaginary time t :

$$H \psi(\vec{r}_1 \dots \vec{r}_N) = E \psi(\vec{r}_1 \dots \vec{r}_N) \quad \psi(t) = e^{-Ht} \psi(0)$$

Ground-state extracted in the limit of $t \rightarrow \infty$.

Light nuclei spectrum computed with GFMC



Carlson, Gandolfi, Pederiva, Pieper, Schiavilla, Schmidt, Wiringa, RMP (2015)

Also radii, densities, matrix elements, ...

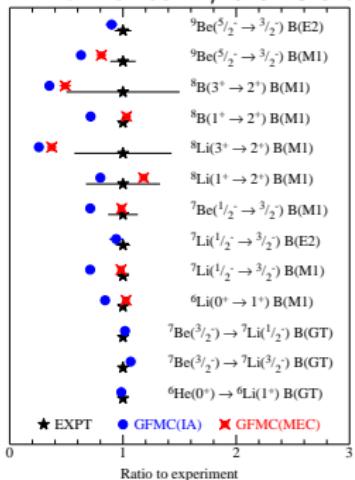
Two-body processes

$$\mathbf{j} = \mathbf{j}^{(1)} + \mathbf{j}^{(2)}(\nu) + \left[\begin{array}{c} \pi \\ \text{---} \\ \text{---} \end{array} \right] + \left[\begin{array}{c} \pi \\ \text{---} \\ \rho, \omega \end{array} \right]$$

+ $\mathbf{j}^{(3)}(\nu, 2\pi)$

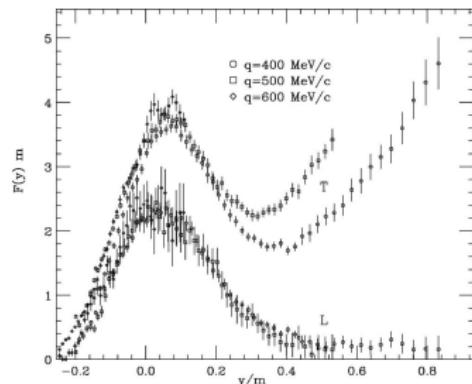
transverse

Low-momentum, transitions:



Pastore et al, PRC 2014

High-momentum, e^- scattering:
rescaled longitudinal vs transverse
electromagnetic response in ${}^{12}\text{C}$



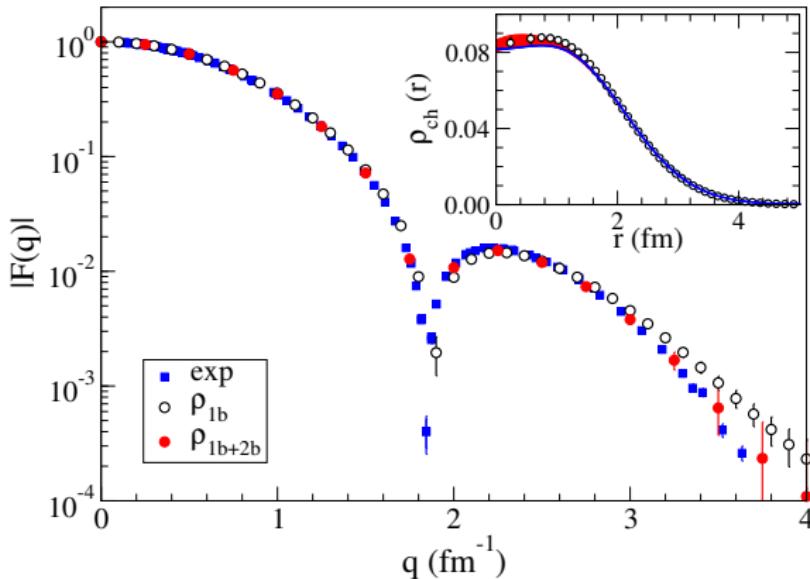
Benhar, Day, Sick, RMP (2008)

Without two-body processes, the longitudinal and transverse response is about the same

Charge form factor of ^{12}C

$$|F(q)| = \langle \psi | \rho_q | \psi \rangle$$

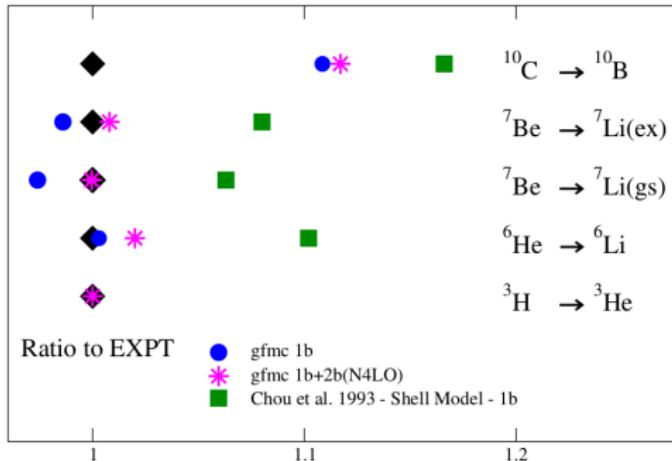
$$\rho_q = \sum_i \rho_q(i) + \sum_{i < j} \rho_q(ij)$$



Lovato, Gandolfi, Butler, Carlson, Lusk, Pieper, Schiavilla, PRL (2013)

β -decays in light nuclei

QMC calculations using a correlated wave function compared to shell-model calculations using the AV18+IL7 Hamiltonian and chiral currents.

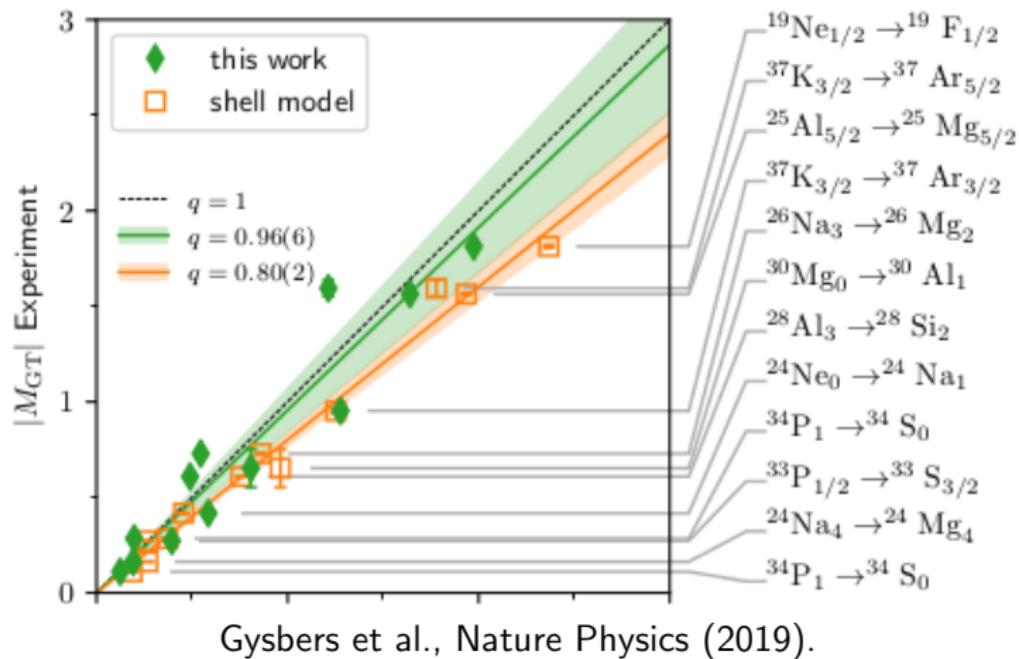


Pastore, et al., PRC 97, 022501 (2018).

The effect of correlations in the nuclear wave function is critical!

β -decays in *sd*-shell nuclei

VS-IMSRG calculations using NN-N⁴LO+3N_{lnl}



Gysbers et al., Nature Physics (2019).

Inclusive scattering

Electron scattering:

$$\left(\frac{d^2\sigma}{d\epsilon' d\Omega} \right)_{\nu/\bar{\nu}} = \left(\frac{d\sigma}{d\Omega} \right)_M \left[\frac{Q^4}{q^4} R_L(q, \omega) + \left(\frac{Q^2}{2q^2} + \tan^2 \frac{\theta}{2} \right) R_T(q, \omega) \right]$$

R_T and R_L transverse and longitudinal response functions.

Neutrino scattering:

$$\begin{aligned} \left(\frac{d^2\sigma}{d\epsilon' d\Omega} \right)_{\nu/\bar{\nu}} &= \frac{G^2}{2\pi^2} k' \epsilon' \cos^2 \frac{\theta}{2} \left[R_{00}(q, \omega) + \frac{\omega^2}{q^2} R_{zz}(q, \omega) - \frac{\omega}{q} R_{0z}(q, \omega) + \right. \\ &\quad \left. \left(\tan^2 \frac{\theta}{2} + \frac{Q^2}{2q^2} \right) R_{xx+yy}(q, \omega) \mp \tan \frac{\theta}{2} \sqrt{\tan^2 \frac{\theta}{2} + \frac{Q^2}{q^2}} R_{xy}(q, \omega) \right] \end{aligned}$$

R_{00} , R_{zz} , R_{0z} , R_{xx+yy} , and R_{xy} neutrino response functions.

R_{xy} is important for ν vs $\bar{\nu}$ processes.

Response functions

$$\begin{aligned} R(q, \omega) &= \sum_n \langle \Psi | j^\dagger(q) | n \rangle \langle n | j(q) | \Psi \rangle \delta(\omega - E_n + E_0) \\ &= \int dt \langle \Psi | j^\dagger(q) \exp[i(H - \omega)t] j(q) | \Psi \rangle \\ &= \int dt E(q, \tau) \end{aligned}$$

Using QMC we can calculate exactly $E(q, \tau)$ and then reconstruct $R(q, \omega)$.

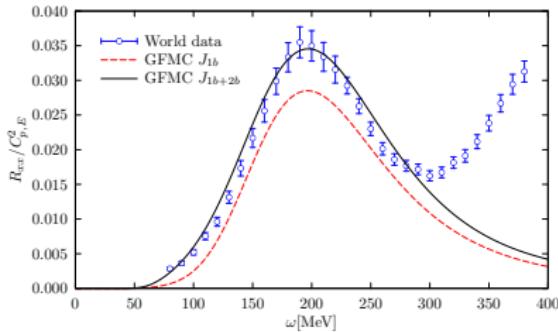
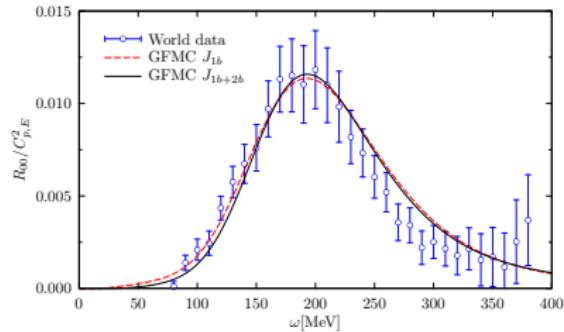
Ingredients:

- Hamiltonian H
- Ground-state Ψ (H)
- Currents described by the electroweak operators $\mathbf{j}(q)$, constructed consistently with H .

Response functions

Using the maximum entropy method, we can reconstruct the response functions.

Longitudinal and transverse response functions of ${}^4\text{He}$ ($q=600$ MeV)

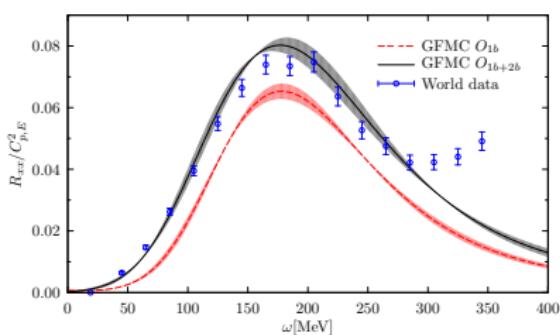
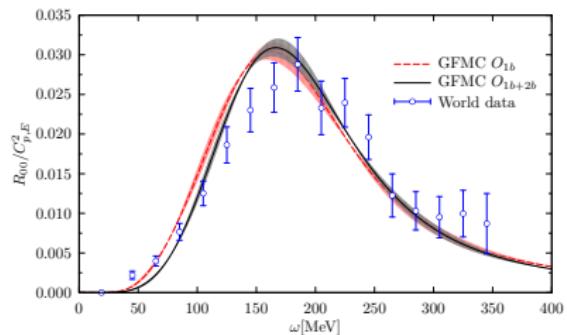


Lovato, Gandolfi, Carlson, Pieper, Schiavilla, PRC (2015)

Similar agreement also with other kinematics, $q=400, 500$, and 700 MeV.

Electromagnetic response functions of ^{12}C

Electromagnetic longitudinal and transverse response functions of ^{12}C
($q=570$ MeV)

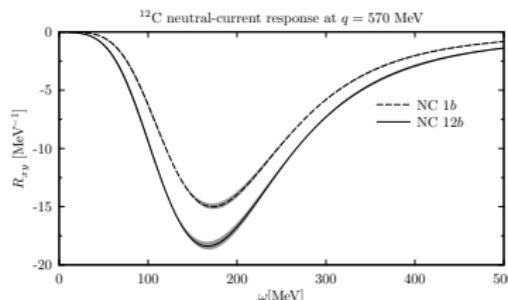
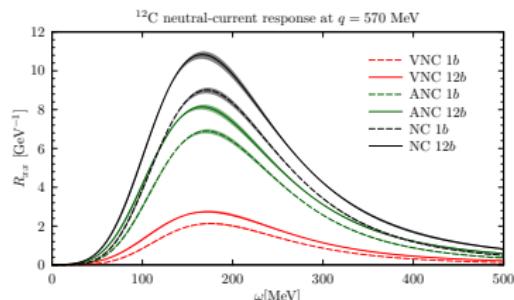
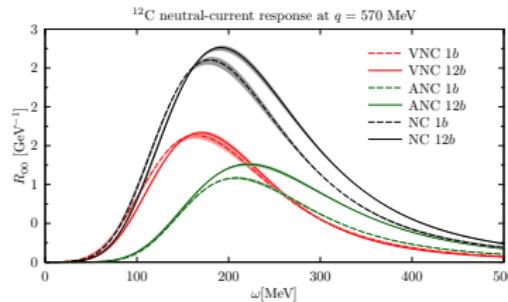
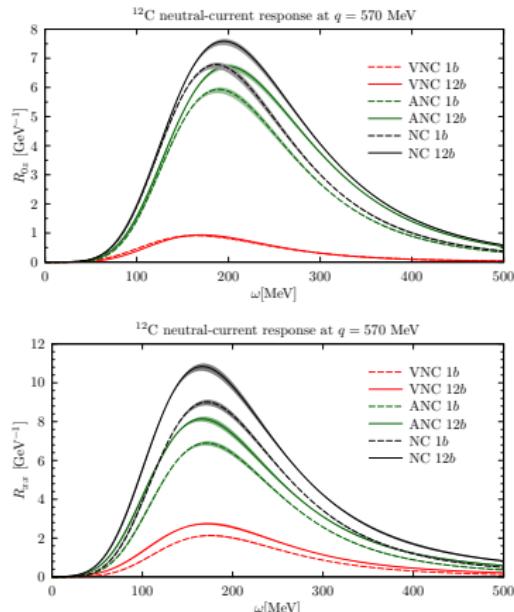


Lovato, Gandolfi, et al., PRL (2016).

Role of two-nucleon currents very important (as expected).

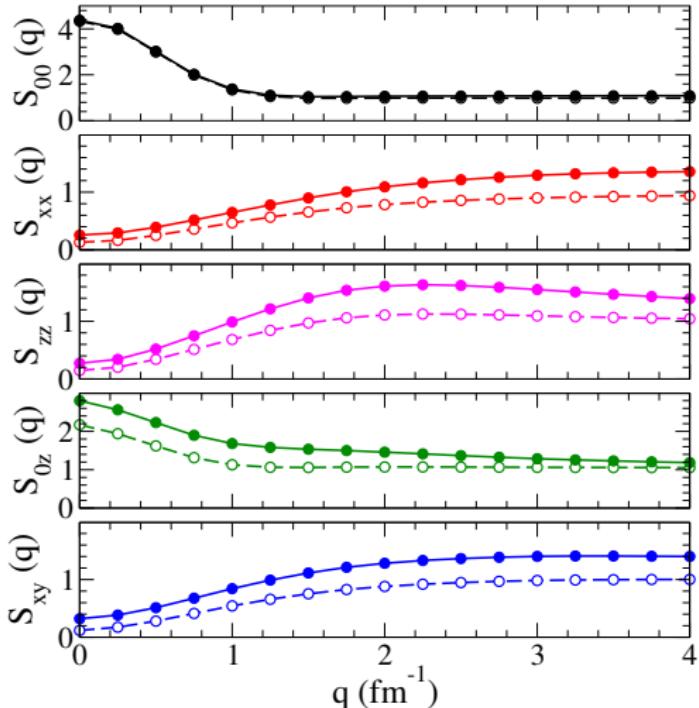
Neutral Electroweak response functions of ^{12}C

Transverse vector, axial, and neutral current of ^{12}C ($q=570$ MeV)



Lovato, Gandolfi, et al., PRC 97, 022502 (2018)

Neutral Electroweak sum-rules in ^{12}C

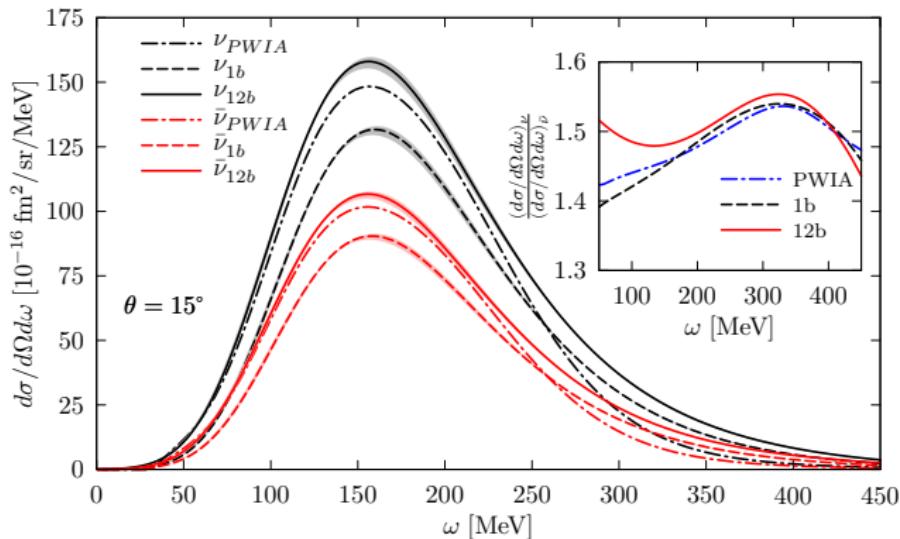


Lovato, Gandolfi, Carlson, Pieper, Schiavilla, PRL (2014).

Two-body operators enhance sum-rules up to 50%.

Neutral Electroweak cross-section of ^{12}C

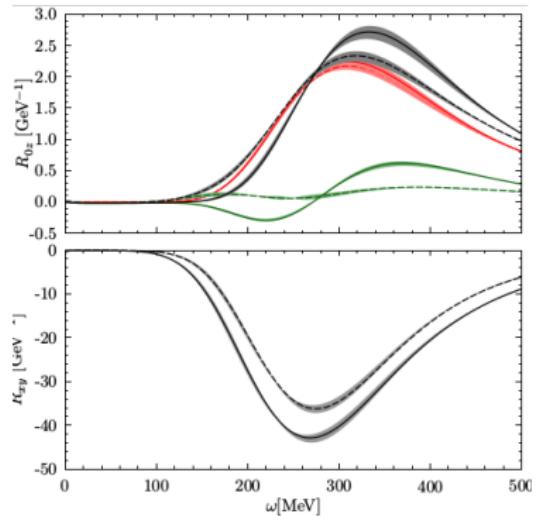
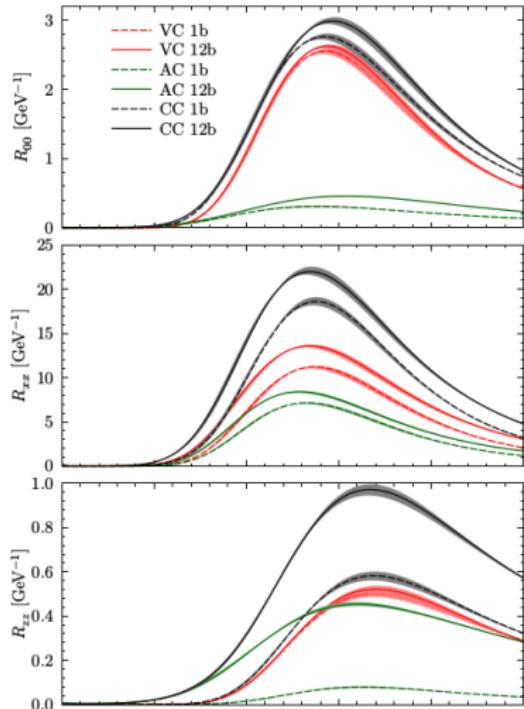
From the response functions, we can reconstruct the cross-section:



Lovato, Gandolfi, et al., PRC 97, 022502 (2018)

PRELIMINARY!

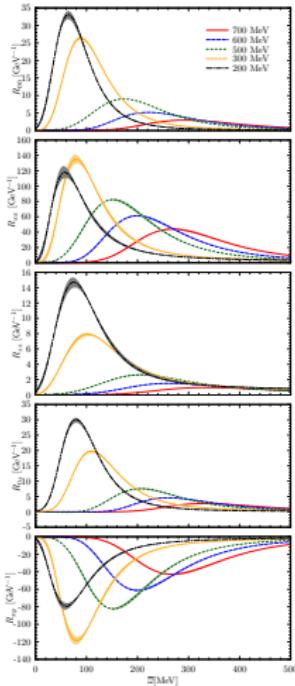
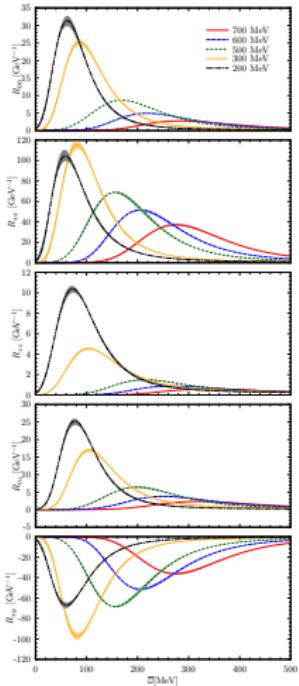
Vector, axial, and charge changing current of ^{12}C , $q=700$ MeV



Lovato, Rocco, Carlson, Gandolfi, Schiavilla, in preparation.

PRELIMINARY!

Vector, axial, and charge changing current of ^{12}C



Lovato, Rocco,
Carlson, Gandolfi,
Schiavilla, in
preparation.

Factorization: Short-Time Approximation

$$R_{\alpha}(q, \omega) = \sum_f \delta(\omega + E_0 - E_f) \langle 0 | O_{\alpha}^{\dagger}(\mathbf{q}) | f \rangle \langle f | O_{\alpha}(\mathbf{q}) | 0 \rangle$$

$$R_{\alpha}(q, \omega) = \int dt \langle 0 | O_{\alpha}^{\dagger}(\mathbf{q}) e^{i(H-\omega)t} O_{\alpha}(\mathbf{q}) | 0 \rangle$$

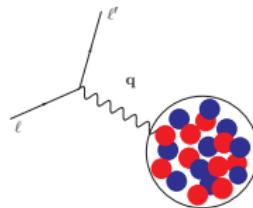
At short time, expand $P(t) = e^{i(H-\omega)t}$ and keep up to 2b-terms

$$H \sim \sum_i t_i + \sum_{i < j} v_{ij}$$

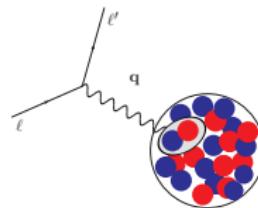
and

$$O_i^{\dagger} P(t) O_i + O_i^{\dagger} P(t) O_j + O_i^{\dagger} P(t) O_{ij} + O_{ij}^{\dagger} P(t) O_{ij}$$

1b



2b



PWIA vs Short-time-approximation

PWIA: Response functions given by incoherent scattering off **single nucleons that propagate freely in the final state** (plane waves)

STA: Response functions are given by the scattering off **pairs of fully interacting nucleons that propagate into a correlated pair of nucleons**

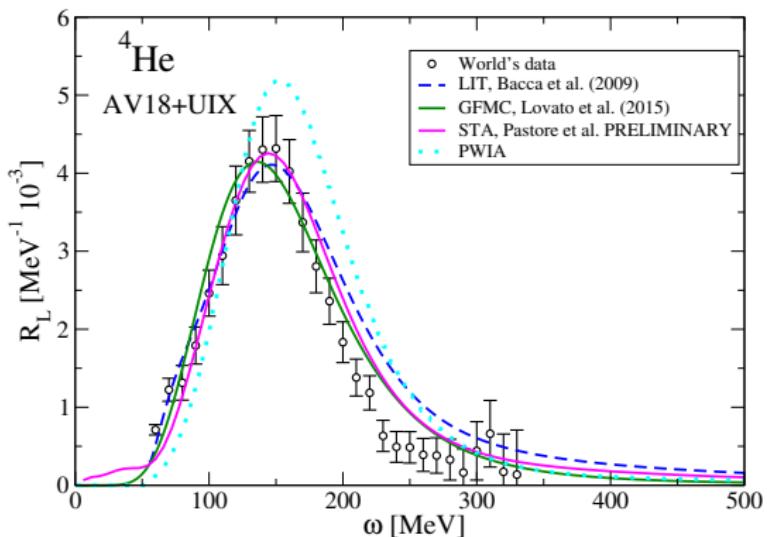
$$R_\alpha(q, \omega) = \sum_f \delta(\omega + E_0 - E_f) \langle 0 | O_\alpha^\dagger(\mathbf{q}) | f \rangle \langle f | O_\alpha(\mathbf{q}) | 0 \rangle$$

$$O_\alpha(\mathbf{q}) = O_\alpha^{(1)}(\mathbf{q}) + O_\alpha^{(2)}(\mathbf{q}) = 1b + 2b$$

$$|f\rangle \sim |\psi_{p,P,J,M,L,S,T,M_T}(r, R)\rangle = \text{correlated two-nucleon w.f.}$$

- * We retain **two-body physics** consistently **in the nuclear interactions** and **electroweak currents**
- * $R_\alpha(q, \omega)$ requires only direct calculation of g.s. $|0\rangle$ w.f.'s *
- * STA can be implemented to accommodate for more two-body physics, e.g., pion-production induced by e and ν

The Short-Time Approximation

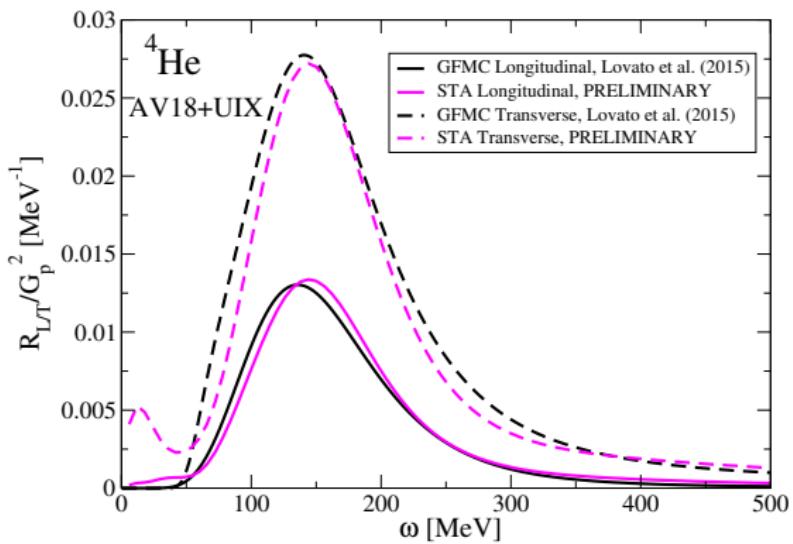


Longitudinal Response function at $q = 500$ MeV

Excellent agreement with full GFMC and EXPT at $q > 500$ MeV

Pastore, Carlson, et al., arXiv:1909.06400.

The Short-Time Approximation



Longitudinal vs Transverse Response Function at $q = 500$ MeV
Pastore, Carlson, et al., arXiv:1909.06400.

Summary and future work

Conclusions:

- “Quenching” of g_A maybe understood. Two-body currents and nuclear correlations very important.
- Electron scattering in ^{12}C calculated using GFMC. Good agreement with experiments. One- and two-body vector currents tested.
- Two-body axial currents show a similar enhancement in response functions and sum rules.
- STA approximation beyond PWIA, very powerful, promising results.

In progress/future work:

- Calculation of charge changing weak currents almost complete. Cross-section next.
- Extension to larger nuclei with STA.
- Extension to exclusive processes.

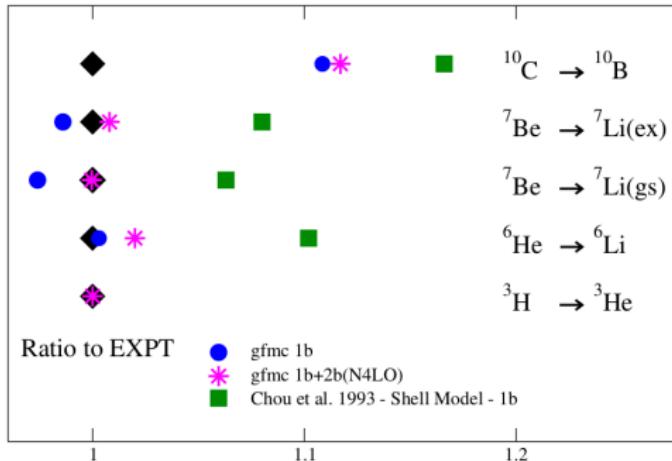
Acknowledgments:

- J. Carlson (LANL)
- S. Pastore (WUSTL)
- A. Lovato, N. Rocco, S. Pieper (ANL)
- R. Schiavilla (Jlab/ODU)

Extra slides

β -decays in light nuclei

QMC calculations using a correlated wave function compared to shell-model calculations using the AV18+IL7 Hamiltonian and chiral currents.

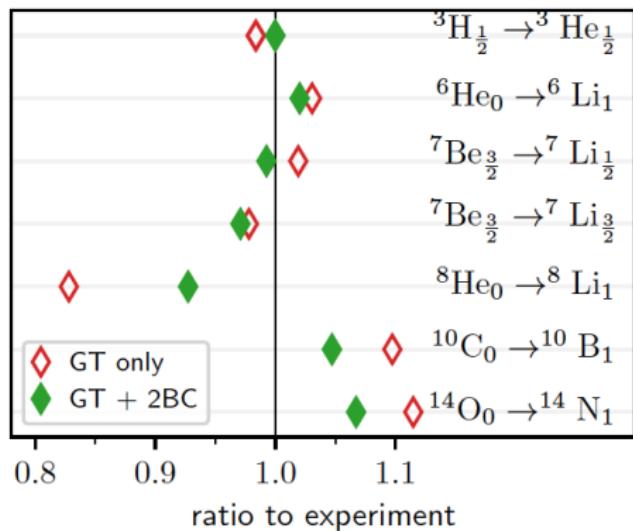


Pastore, et al., PRC 97, 022501 (2018).

The effect of correlations in the nuclear wave function is critical!

β -decays in light nuclei

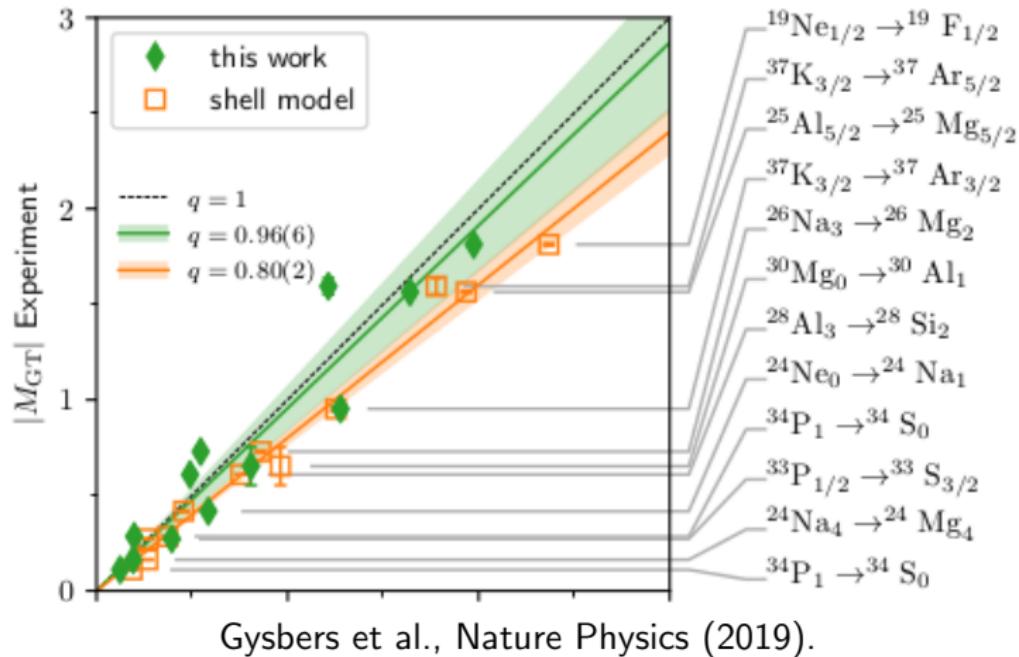
NCSM calculations using NN-N⁴LO+3N_{lnl}



Gysbers et al., Nature Physics (2019).

β -decays in *sd*-shell nuclei

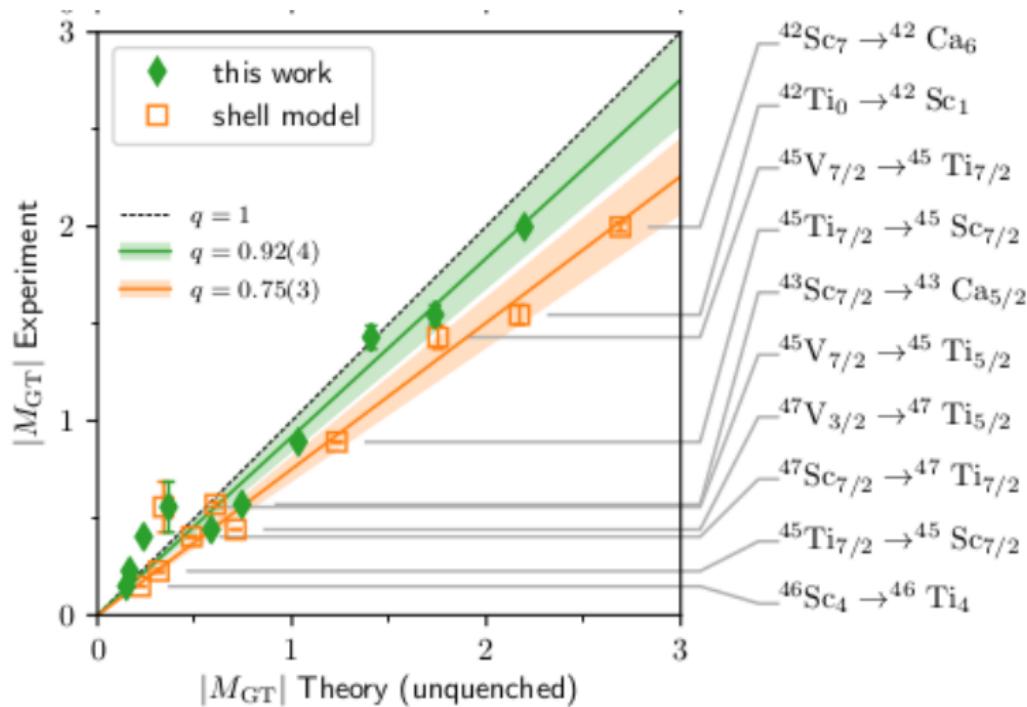
VS-IMSRG calculations using NN-N⁴LO+3N_{lnl}



Gysbers et al., Nature Physics (2019).

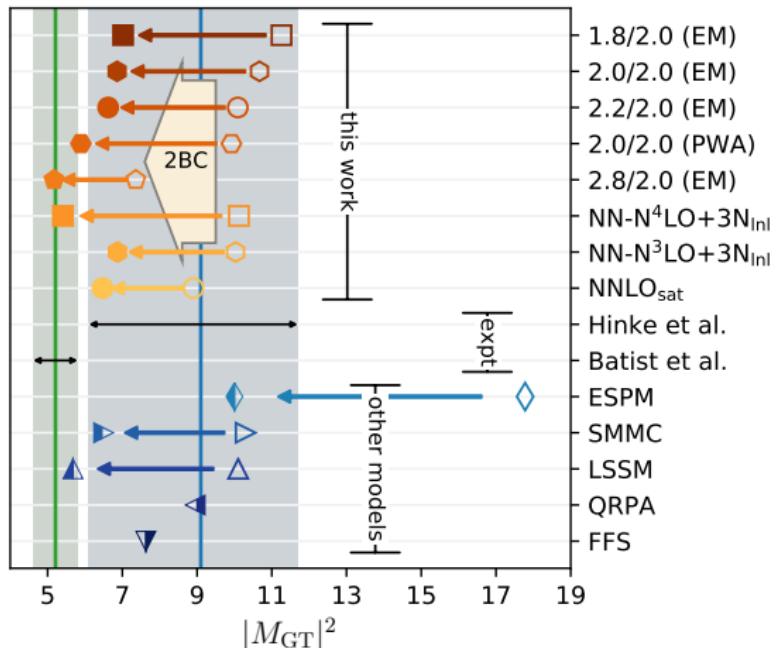
β -decays in pf -shell nuclei

VS-IMSRG calculations using NN-N⁴LO+3N_{lnl}



Gysbers et al., Nature Physics (2019).

β -decay in ^{100}Sn



ESPM: Extreme Single Particle Model

SMMC: Shell Model MC

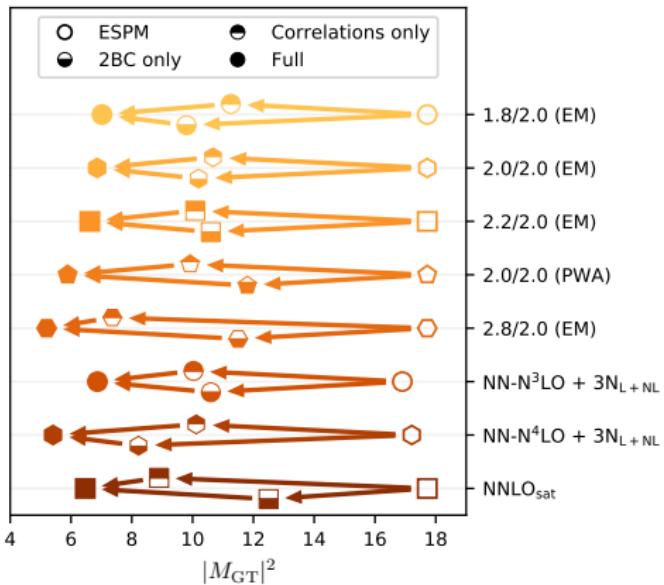
LSSM: Large Space Shell Model

QRPA: quasiparticle random phase approximation

FFS: finite Fermi systems

Gysbers et al., Nature Physics (2019).

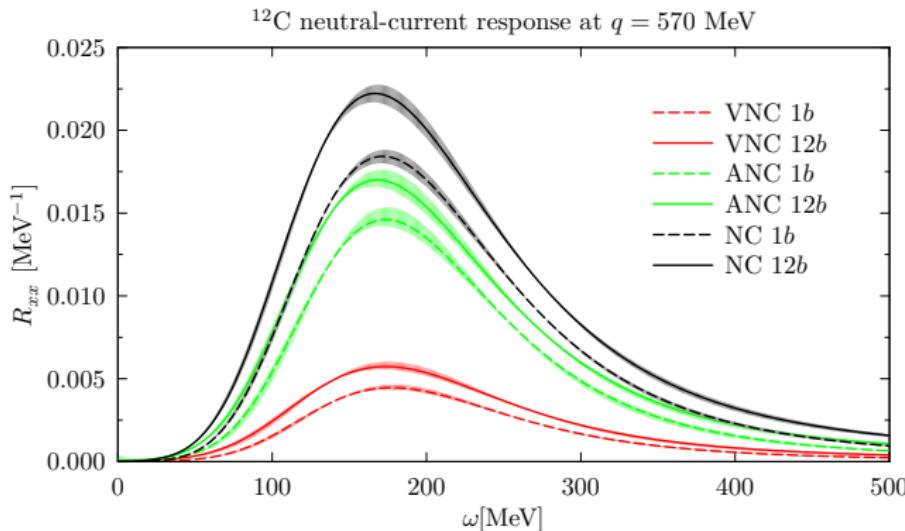
Role of correlations vs 2BC



Gysbers et al., Nature Physics (2019).

Euclidean electroweak response functions of ^{12}C

Transverse vector, axial, and neutral current of ^{12}C ($q=570$ MeV)

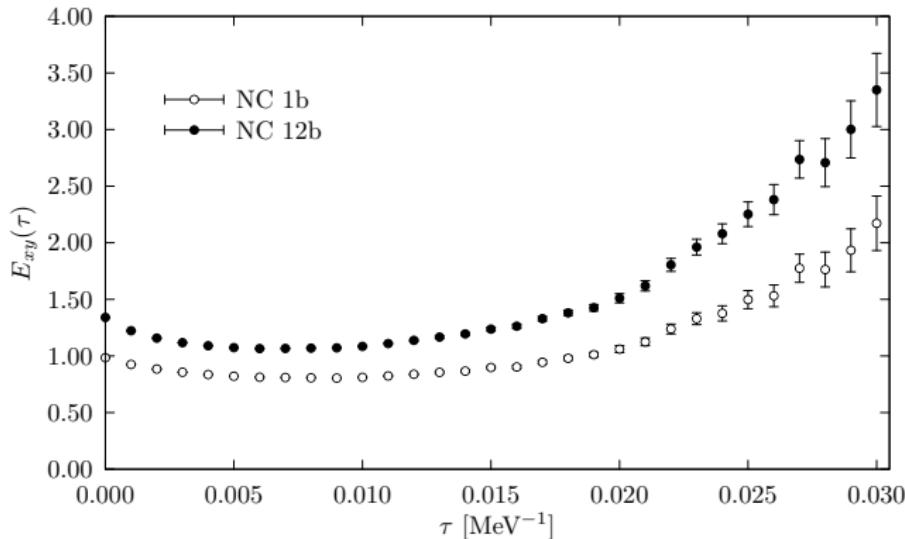


Axial currents give the largest contribution.

Role of axial form factor?

Euclidean electroweak response functions of ^{12}C

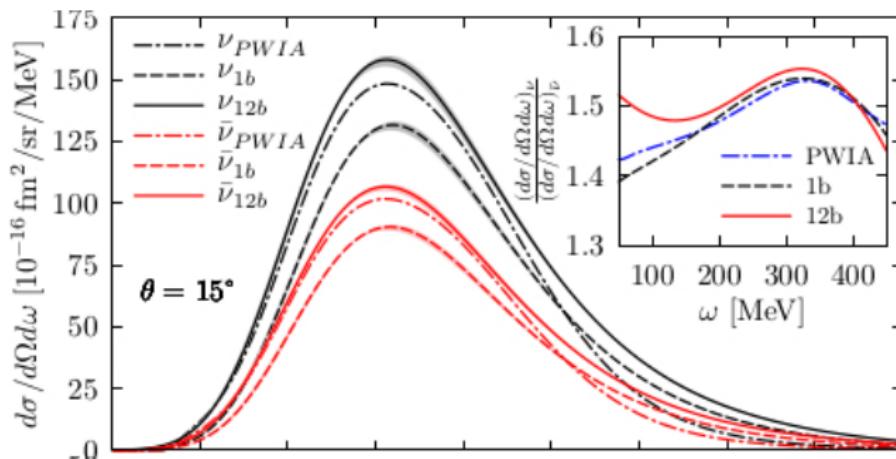
R_{xy} term responsible for ν vs $\bar{\nu}$ response. ^{12}C , $q=570$ MeV



Lovato, Gandolfi, Carlson, Pieper, Schiavilla, PRC (2015)

Electroweak cross-section of ^{12}C

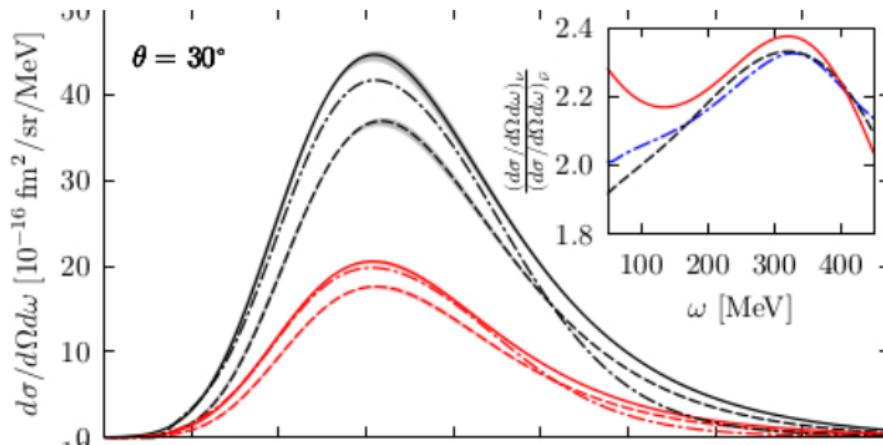
From the response functions, we can calculate the cross-section:



Lovato, et al., PRC 97, 022502 (2018)

Electroweak cross-section of ^{12}C

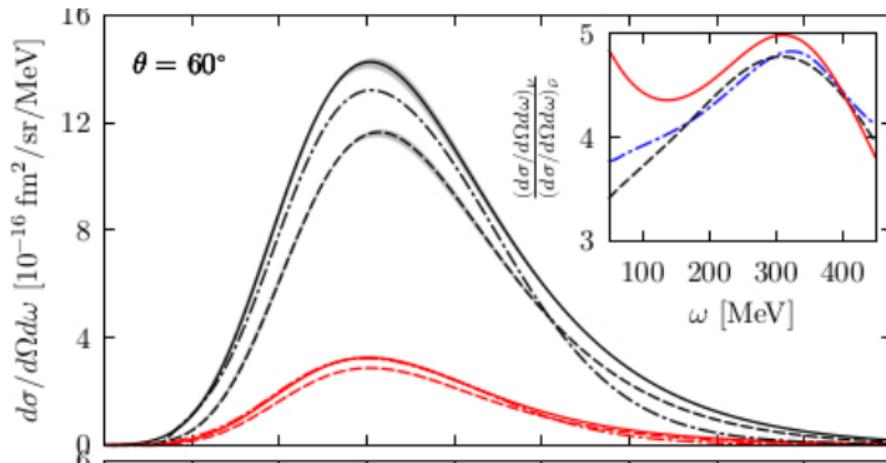
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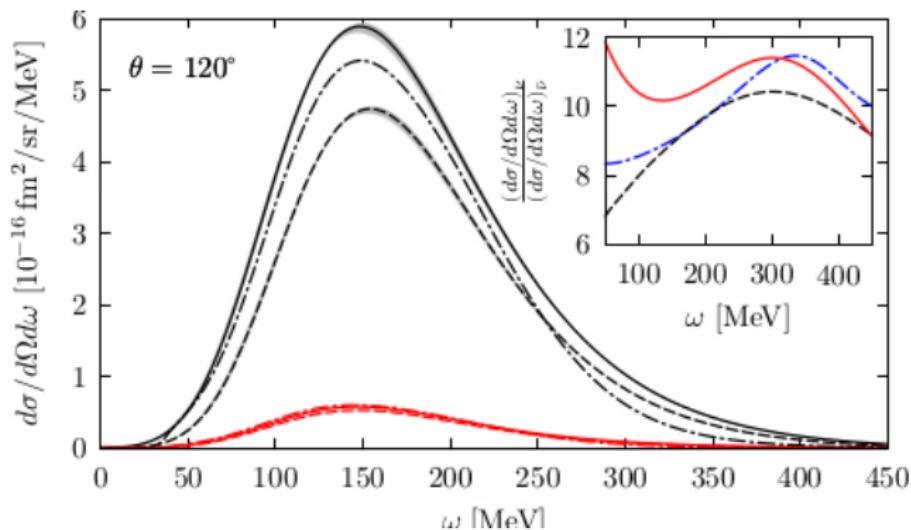
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Electroweak cross-section of ^{12}C

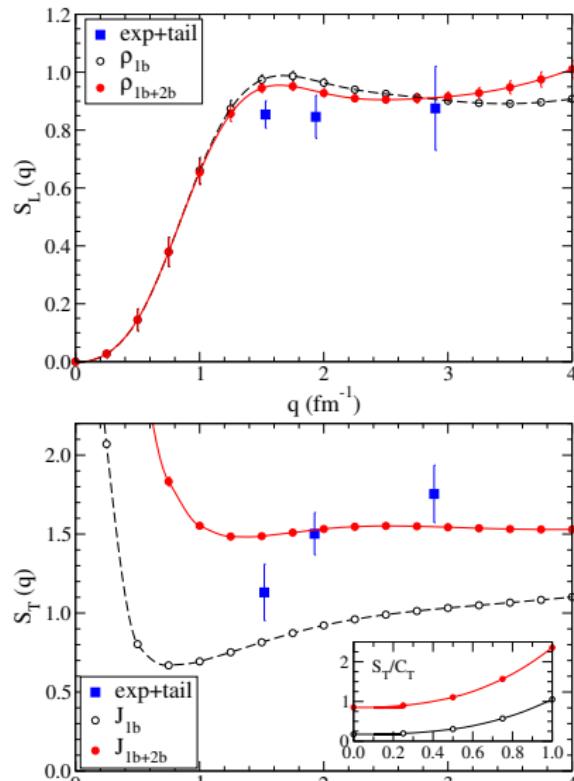
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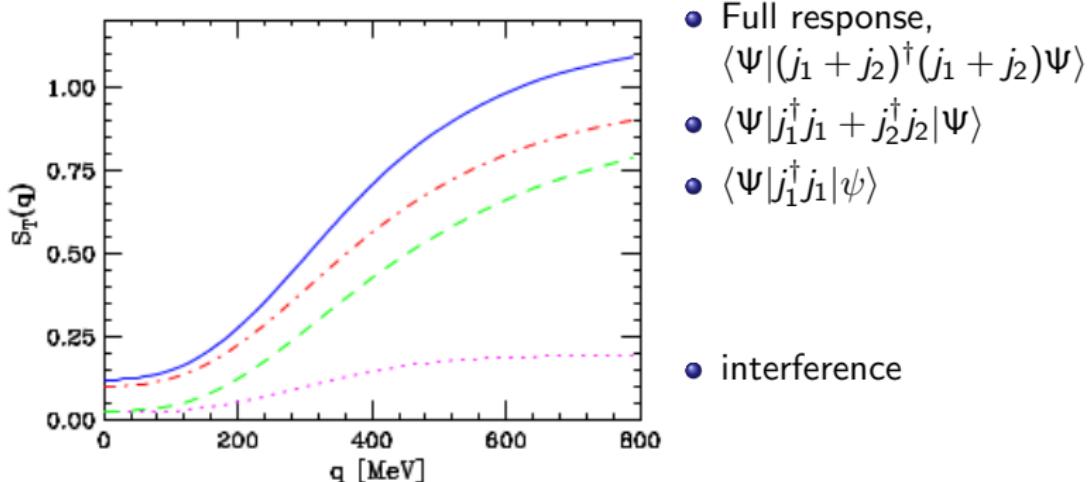
Lovato, et al., PRC 97, 022502 (2018)

Electromagnetic sum-rules in ^{12}C

Sum rules: $S_{L,T}(q) = C_{L,T} \int R_{L,T}(\omega, q) d\omega$



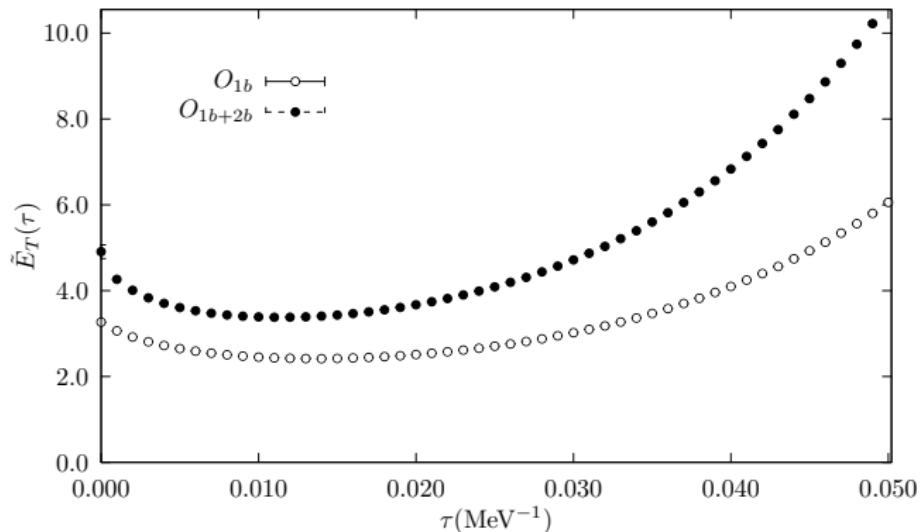
Transverse sum rule of ^{12}C



Benhar, Lovato, Rocco, PRC (2015)

Euclidean response

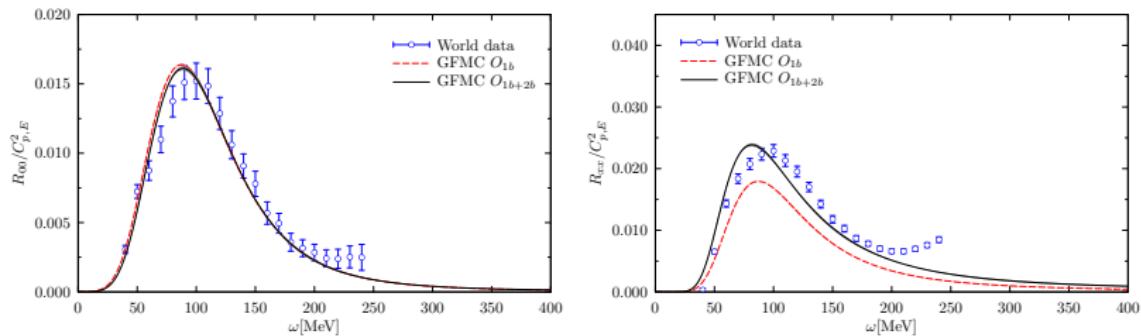
Transverse electromagnetic (euclidean) response functions of ${}^4\text{He}$ ($q=500$ MeV)



Note: results multiplied by $\exp(\tau q^2/2m)$

Longitudinal and transverse response response

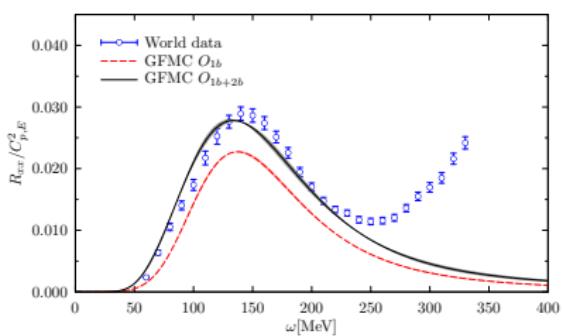
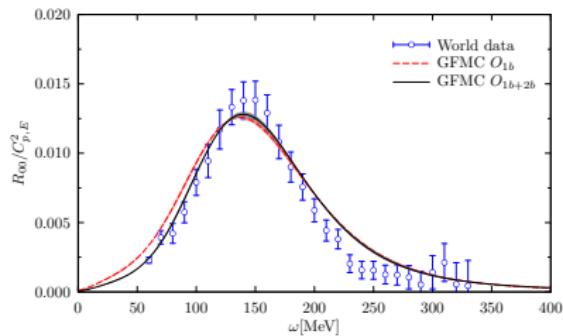
Longitudinal and transverse electromagnetic response functions of ${}^4\text{He}$
($q=400$ MeV)



Note: results multiplied by $\exp(\tau q^2/2m)$

Longitudinal and transverse response response

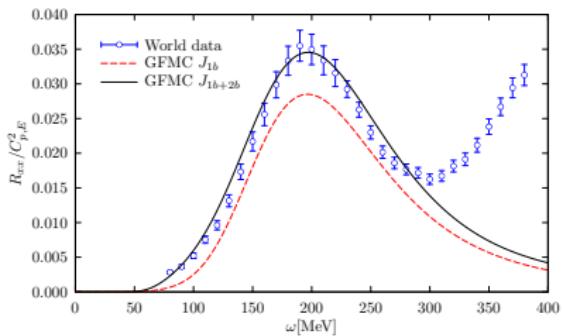
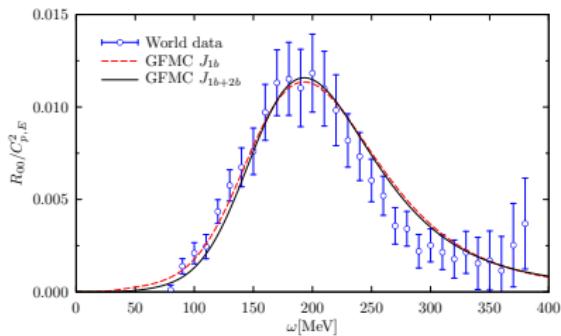
Longitudinal and transverse electromagnetic response functions of ${}^4\text{He}$
($q=500$ MeV)



Note: results multiplied by $\exp(\tau q^2/2m)$

Longitudinal and transverse response response

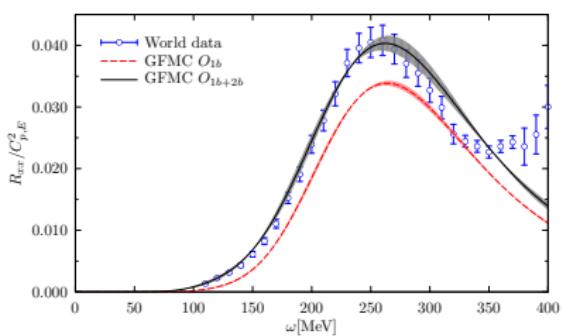
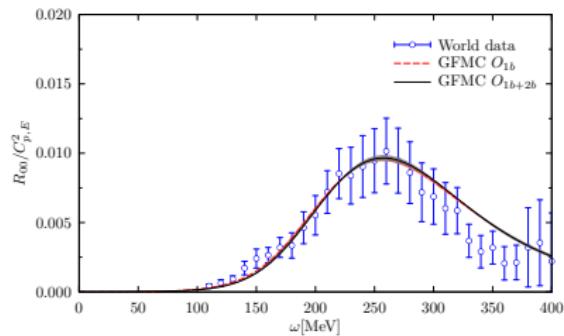
Longitudinal and transverse electromagnetic response functions of ${}^4\text{He}$ ($q=600$ MeV)



Note: results multiplied by $\exp(\tau q^2/2m)$

Longitudinal and transverse response response

Longitudinal and transverse electromagnetic response functions of ${}^4\text{He}$
($q=700$ MeV)



Note: results multiplied by $\exp(\tau q^2/2m)$

Transverse response response

Transverse electromagnetic response functions of ${}^4\text{He}$ ($q=500$ MeV).
Role of the interference:

