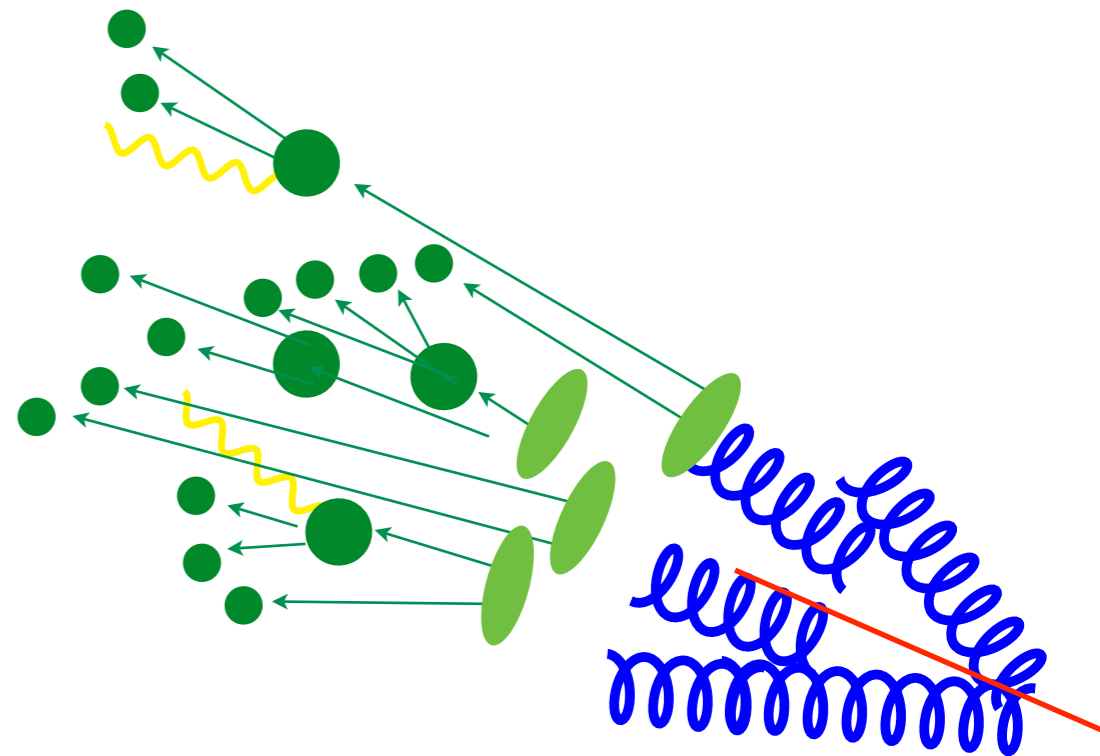


# Top quark physics at the LHC

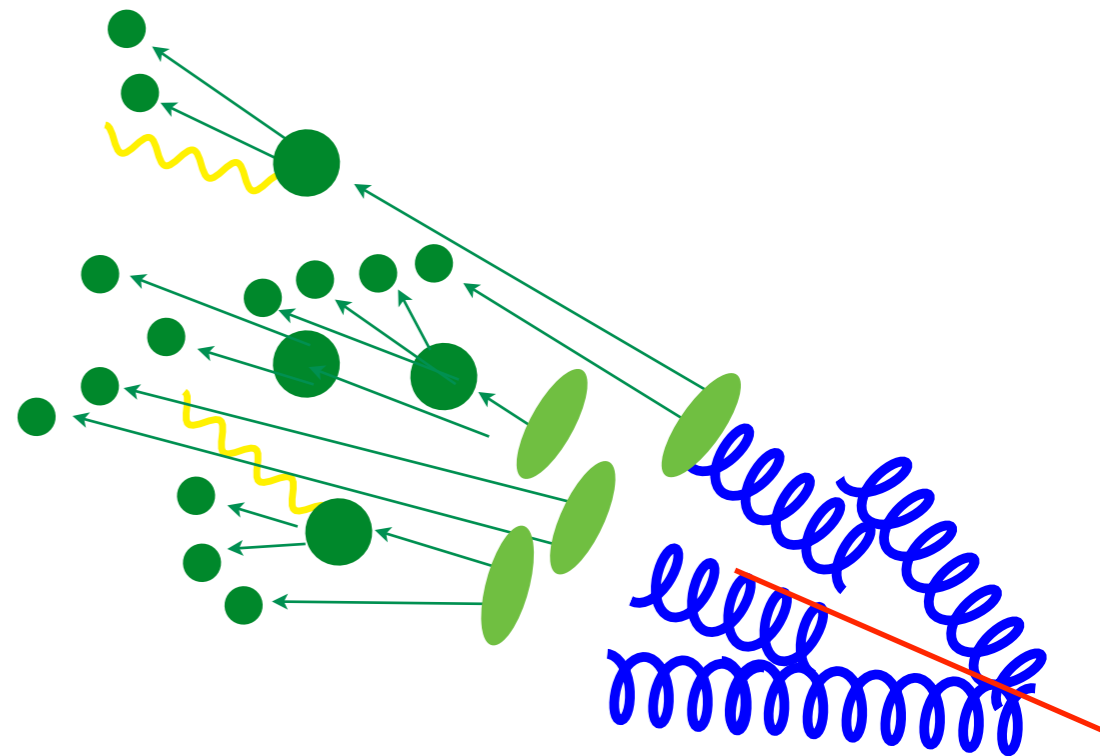
Benjamin Nachman

*Lawrence Berkeley National Laboratory*

*CTEQ school 2019*



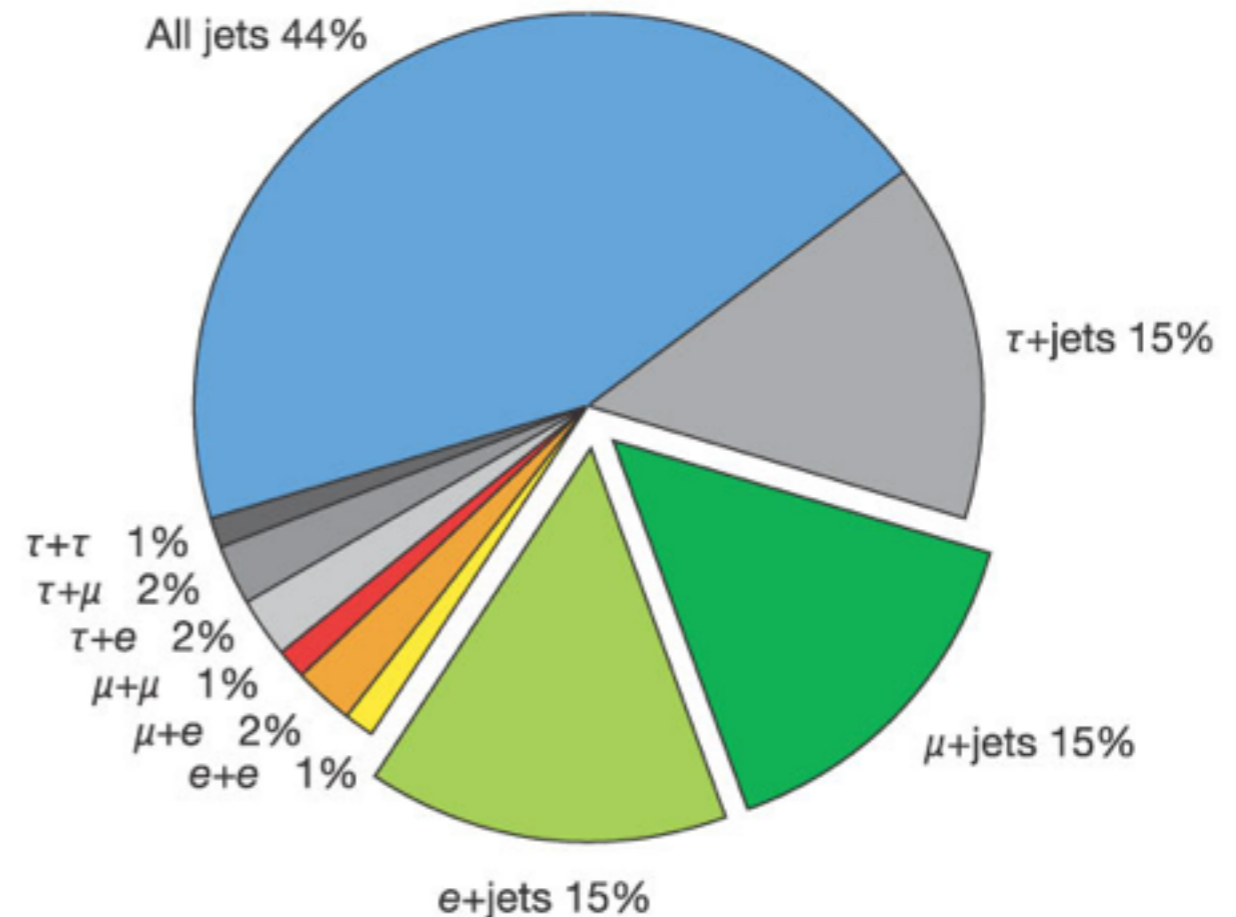
- Top quark mass
- Top quark properties
  - Jet pull
- Single top
  - Interference with di-top
- BSM with top quarks
  - stops in SUSY



# Why is the top quark special?

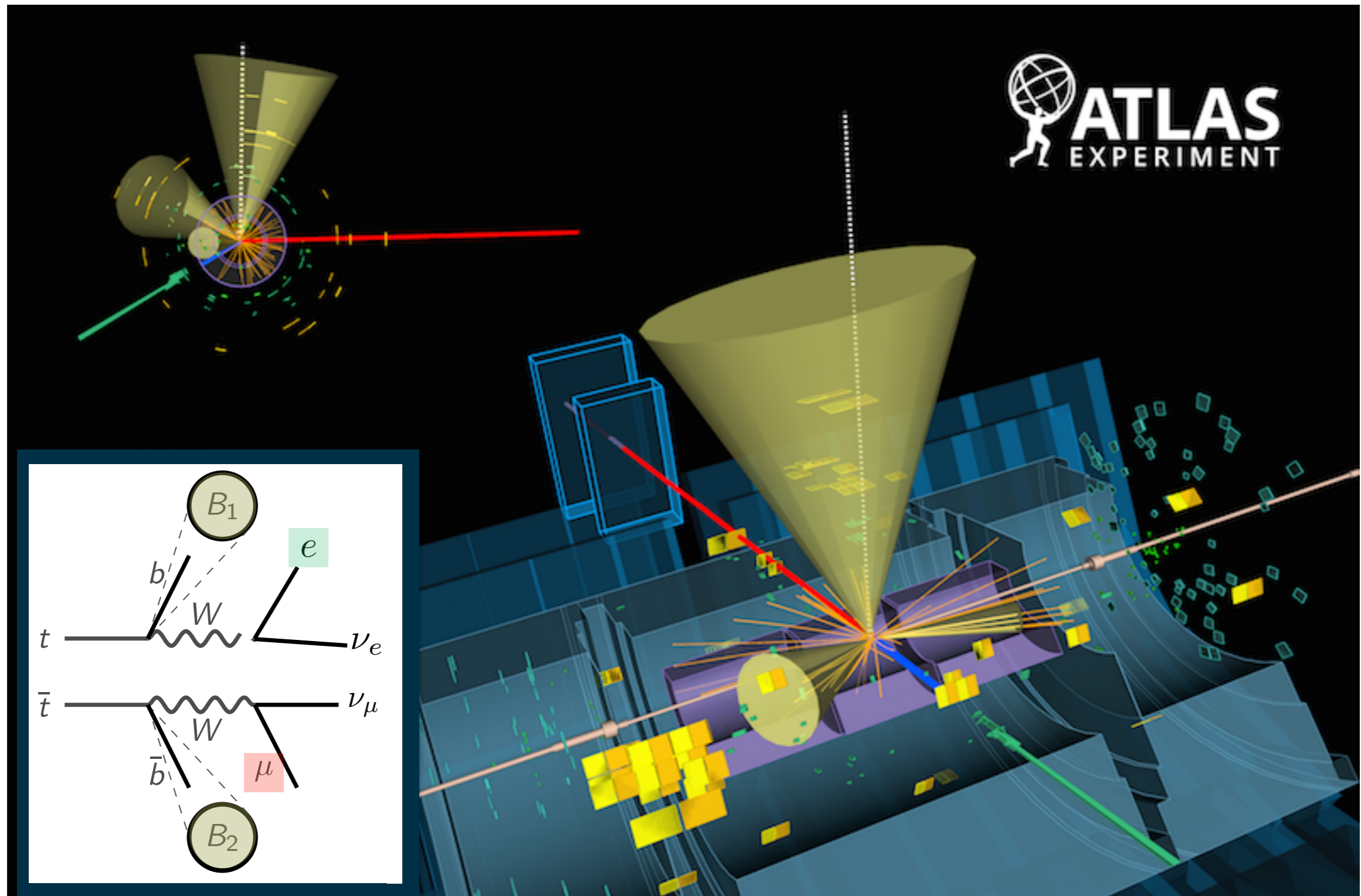


- Very heavy (mass  $\sim 175$  GeV)
  - Decays  $\sim 100\%$  of the time to  $bW$  ( $W$  on-shell)
  - Quadratic contributions to Higgs boson mass
- Short lived (lifetime  $\sim$  few  $10^{-25}$  s)
  - $\sim$  Doesn't form hadrons



# Identifying top quark final states

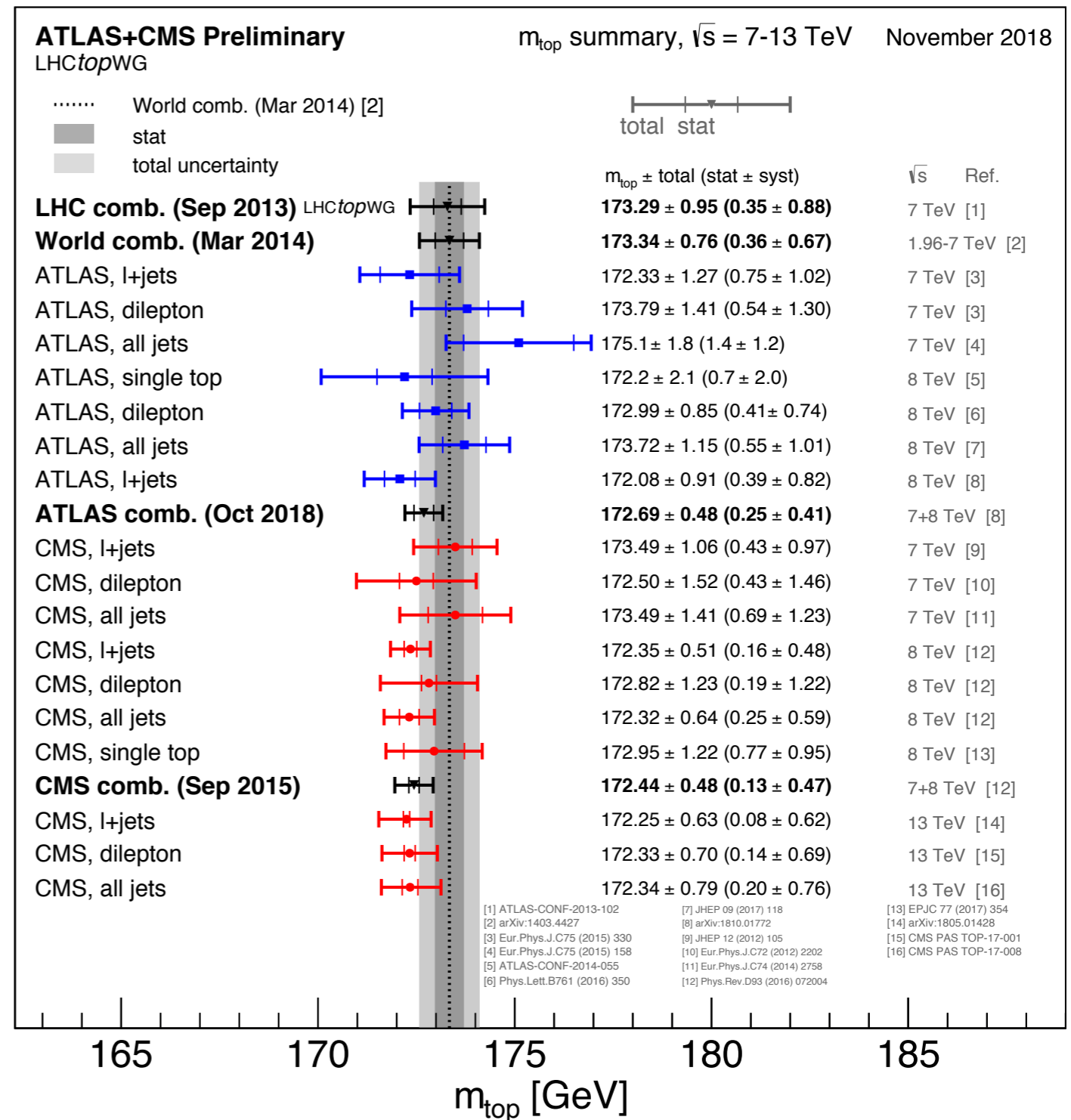
4



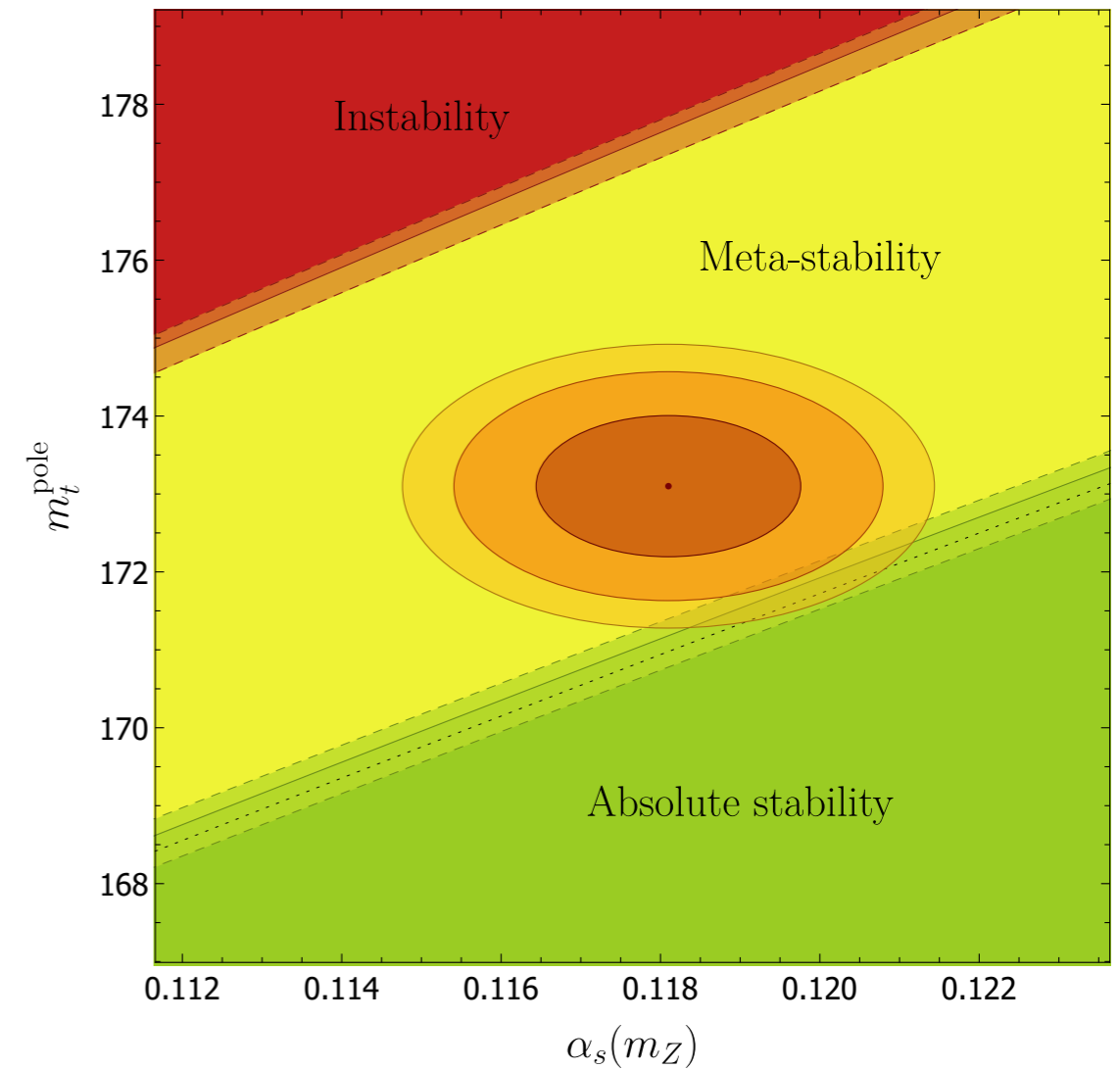
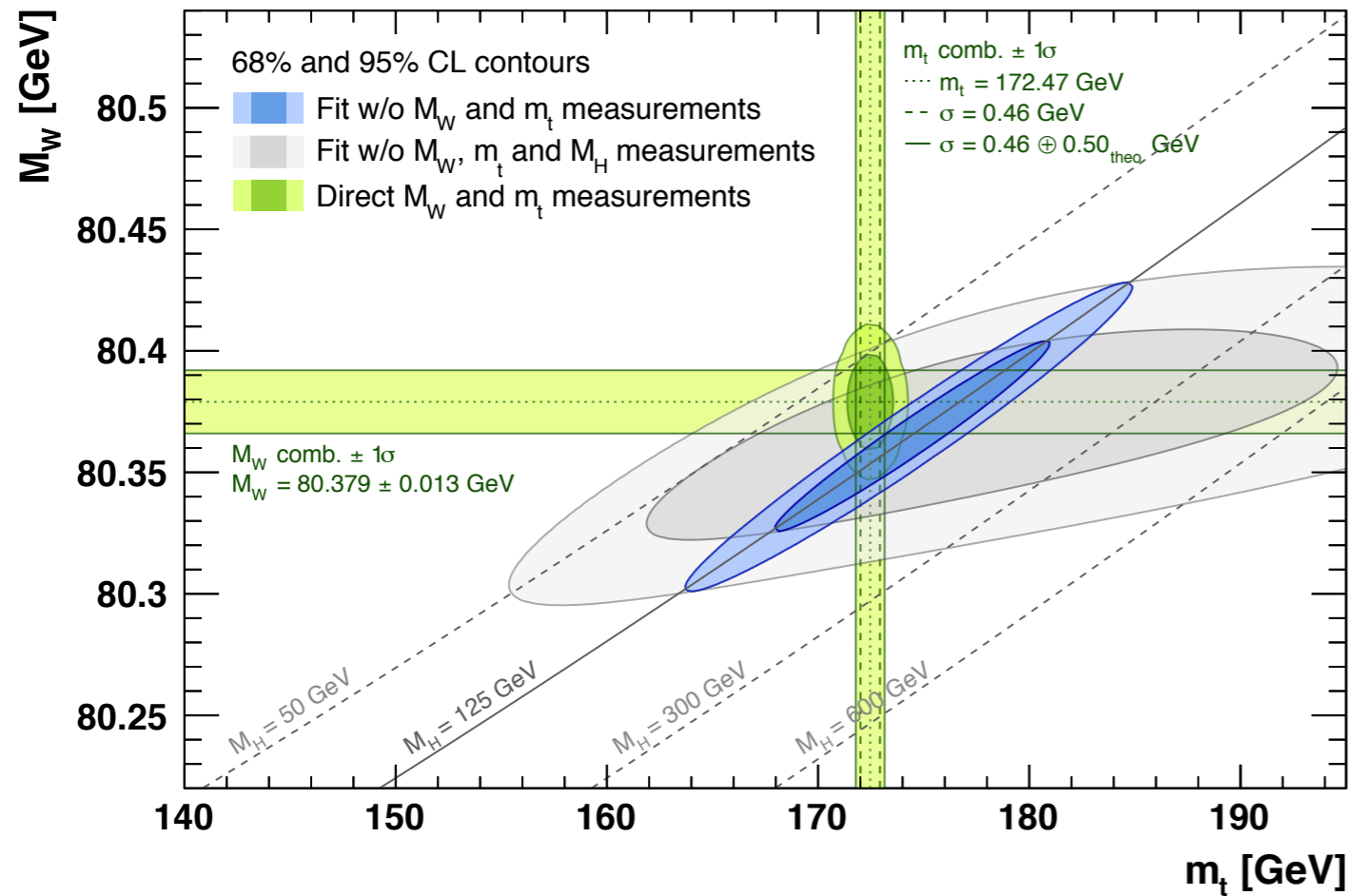
# Top quark mass



- MC mass methods
  - $m_{jjj}$  in semi-leptonic
  - $m_{bl} / m_{T2}$  in di-leptonic
  - Spectra of decay products (e.g.  $J/\psi$ )
  - ...
- Pole mass methods
  - top cross section
  - top jet mass
  - ...



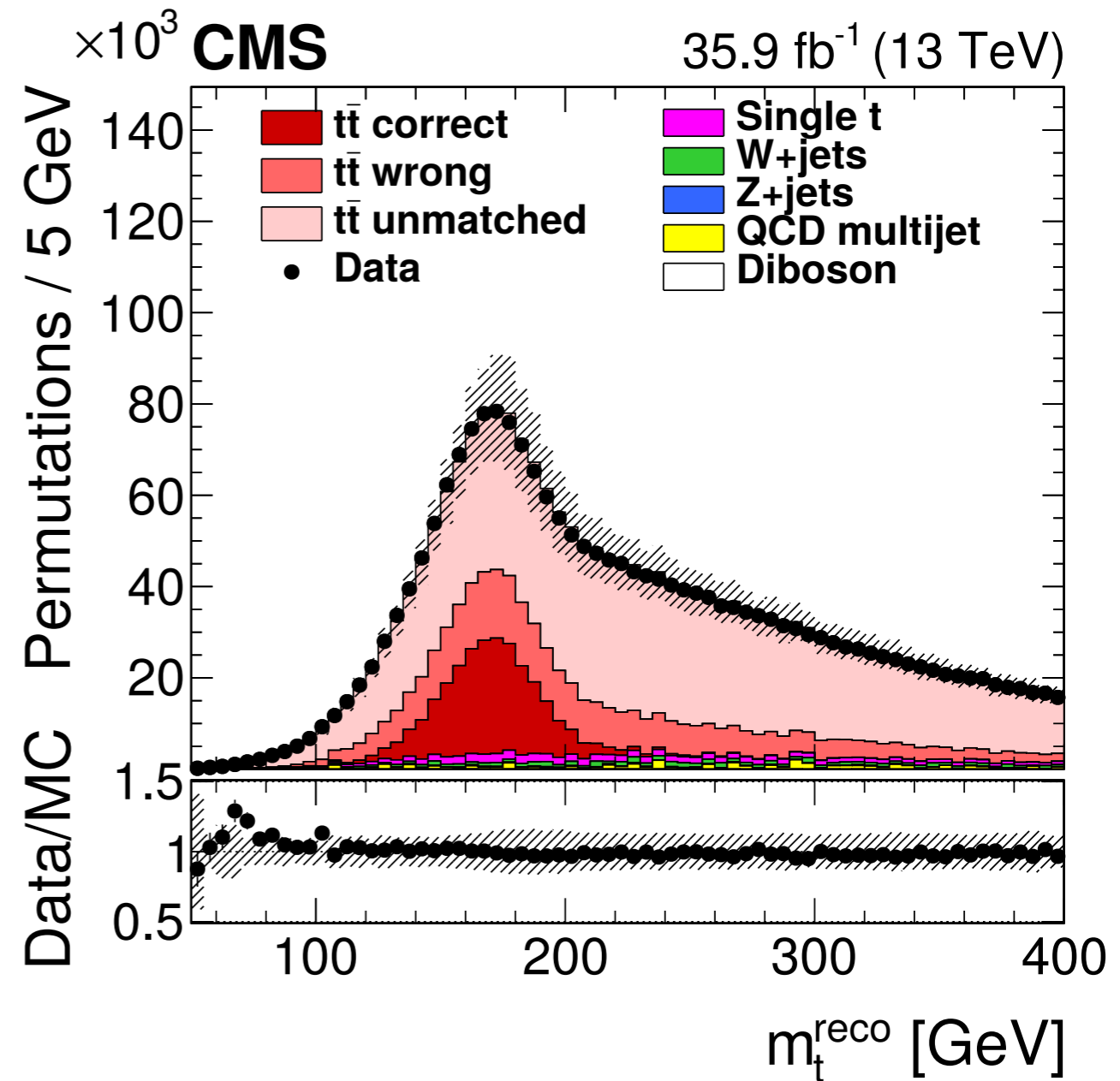
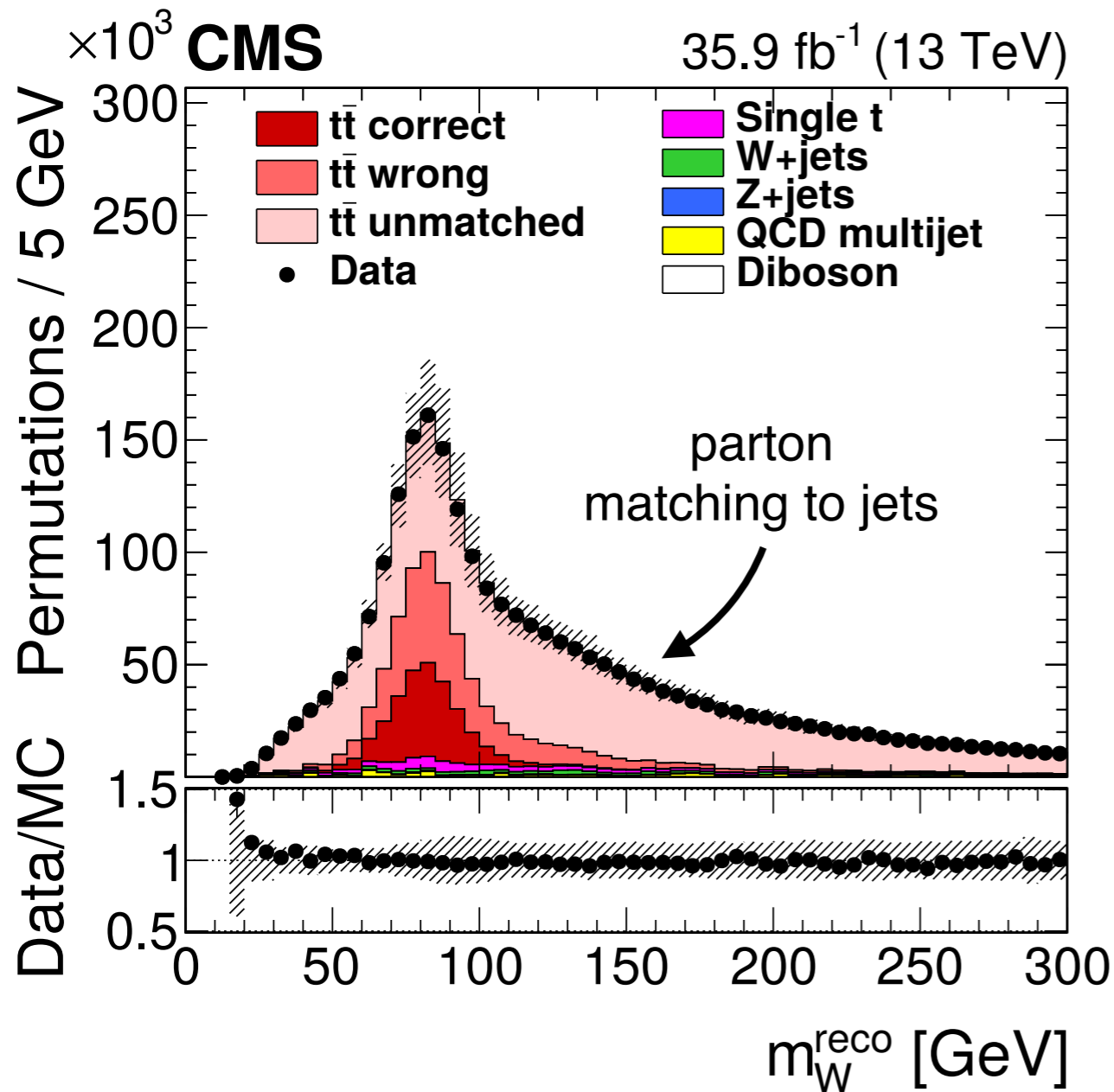
# Top quark mass



consistency of Standard Model  
(unique input from hadron collider)

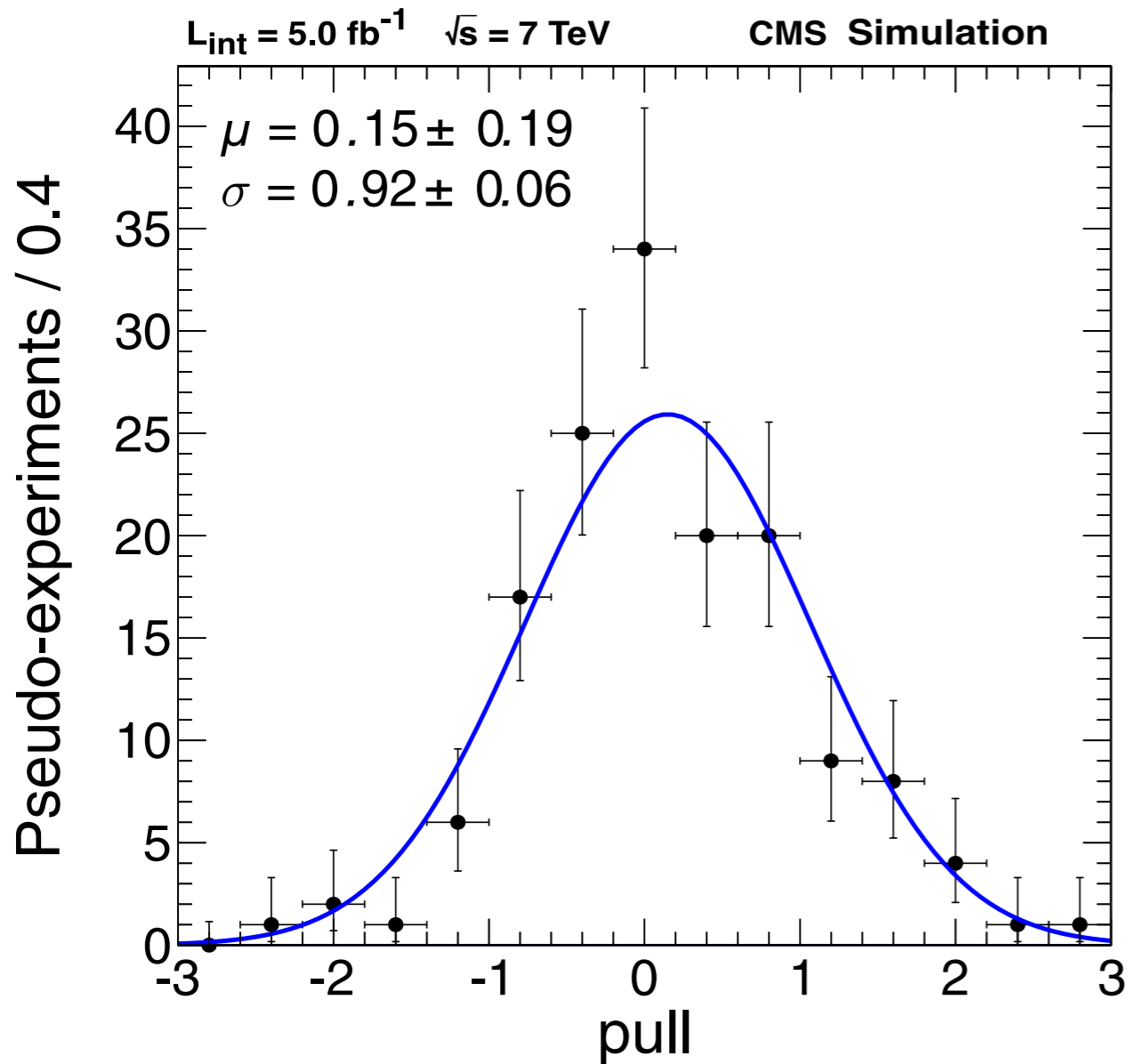
lifetime of the universe

# Top quark mass: lepton + jets

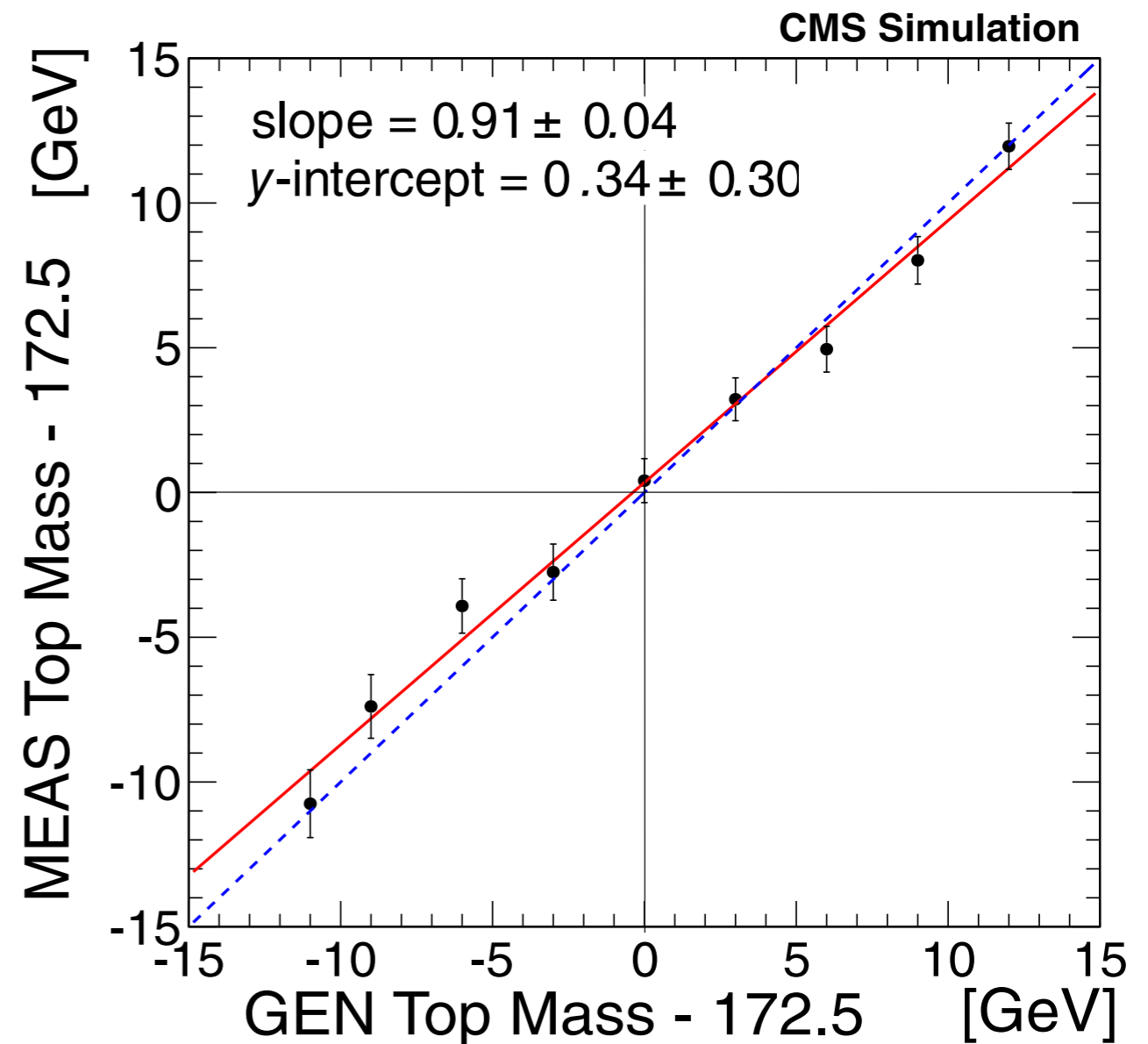


Use templates to fit  $m_{jjb}$ , constrain jet energy with fit to  $m_{jj}$

# Top quark mass: lepton + jets



(measured - true)/uncertainty



Remember:  $\langle c \times x \rangle = c \times \langle x \rangle$

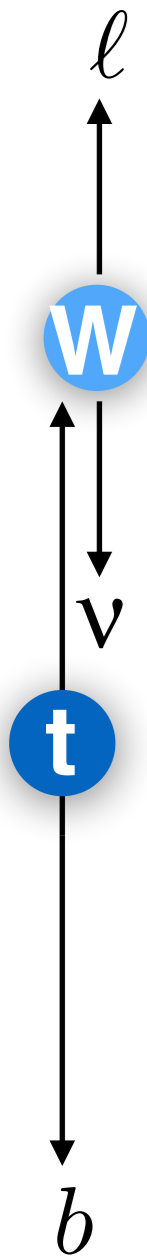
$$\sigma(c \times x) = c \times \sigma(x)$$



# Top quark mass: dilepton

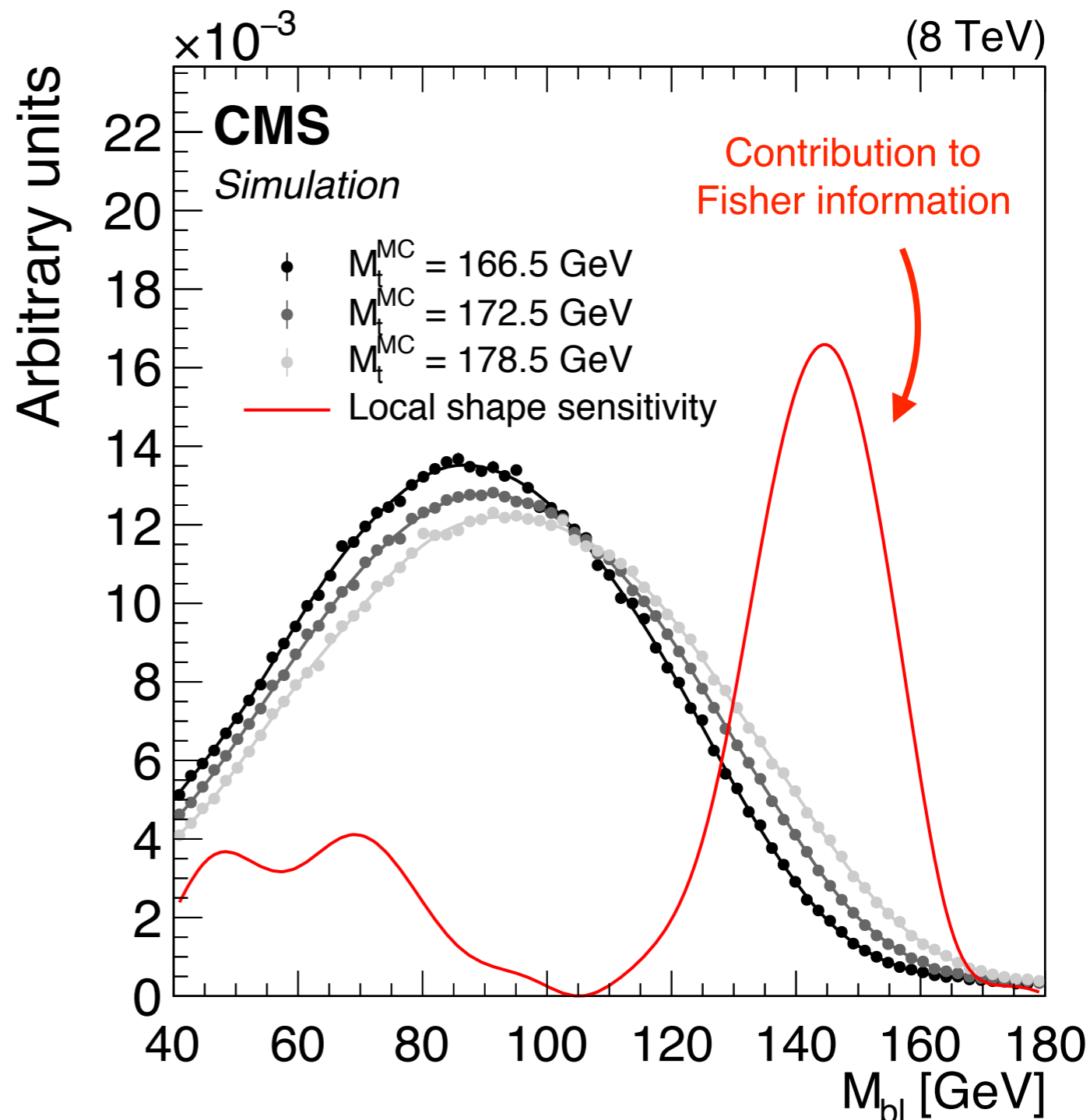


Neutrinos reduce the sensitivity (+ lower stats to begin with), but less background.



$$m_{bl}^{\max} = \sqrt{m_t^2 - m_W^2}$$

(can you derive this? easiest to use the W rest frame)



# MC-template-based top quark mass



Warning: the “kinematic mass”, aka MC mass is not the pole mass (or any well-defined mass in QFT)

There is a  $O(1)$  GeV ambiguity between the quantity we measure (which is MC-dependent) and the top quark mass in the Standard Model.

|                                   | 2D approach       |                   |                         |                        |
|-----------------------------------|-------------------|-------------------|-------------------------|------------------------|
|                                   | $\delta m_t^{2D}$ | $\delta JSF^{2D}$ |                         |                        |
|                                   | [GeV]             | [%]               |                         |                        |
| <i>Experimental uncertainties</i> |                   |                   |                         |                        |
| Method calibration                | 0.05              | <0.1              | Underlying event        | $-0.10 \pm 0.08$ +0.1  |
| JEC (quad. sum)                   | 0.13              | 0.2               | Early resonance decays  | $-0.22 \pm 0.09$ +0.8  |
| – InterCalibration                | (-0.02)           | (<0.1)            | Color reconnection      | $+0.34 \pm 0.09$ -0.1  |
| – MPFInSitu                       | (-0.01)           | (<0.1)            | <b>Total systematic</b> | <b>0.72</b> <b>1.0</b> |
| – Uncorrelated                    | (-0.13)           | (+0.2)            | Statistical (expected)  | 0.09 0.1               |
| Jet energy resolution             | -0.08             | +0.1              | <b>Total (expected)</b> | <b>0.72</b> <b>1.0</b> |
| b tagging                         | +0.03             | <0.1              |                         |                        |

# Alternative approaches

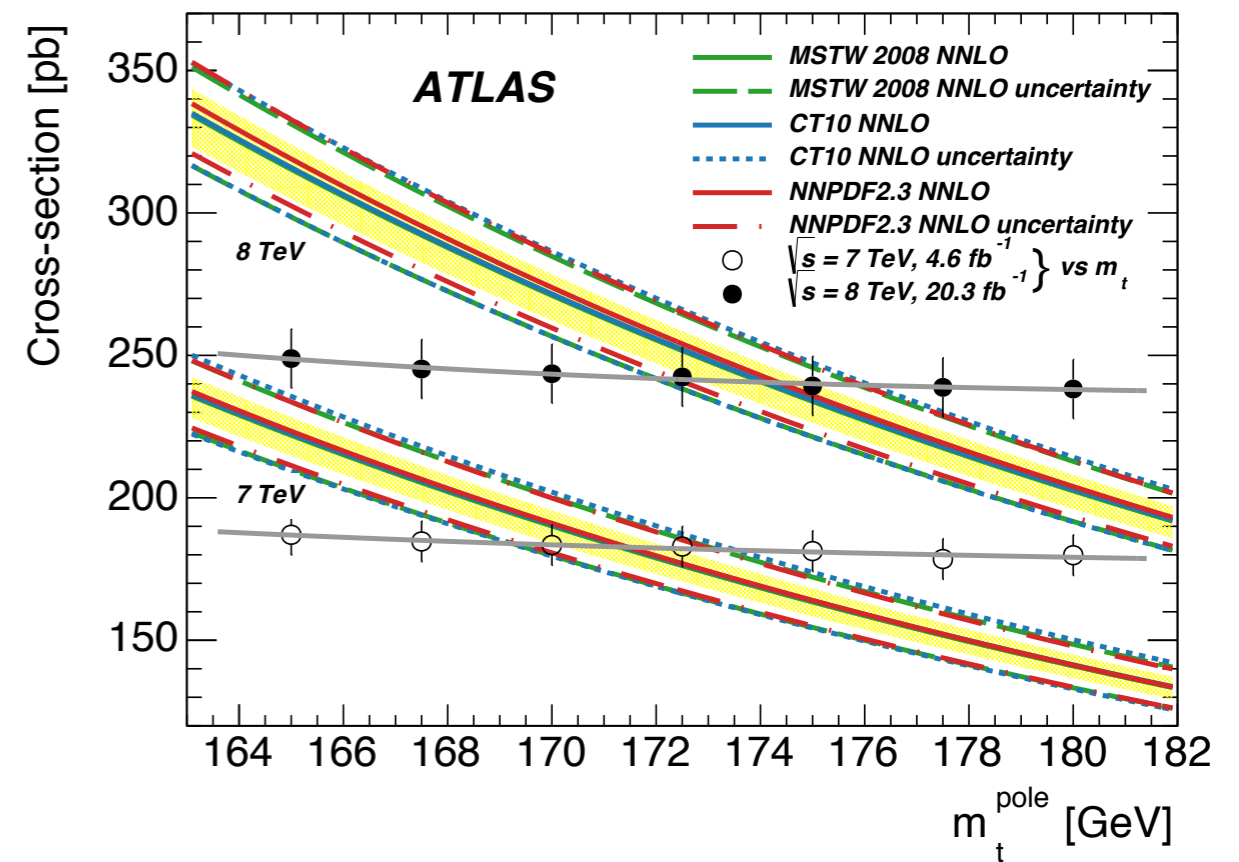
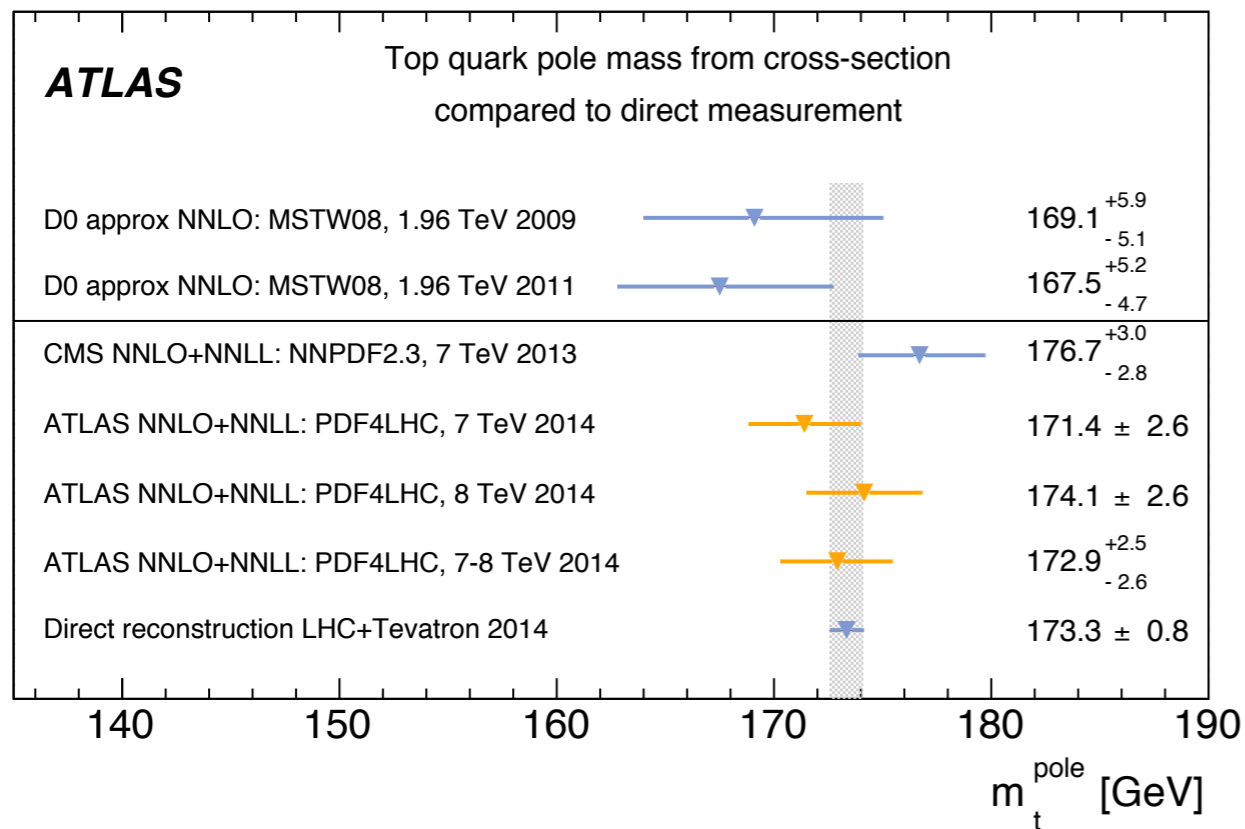
11

How would you measure the mass?

# Alternative approaches

12

How would you measure the mass?

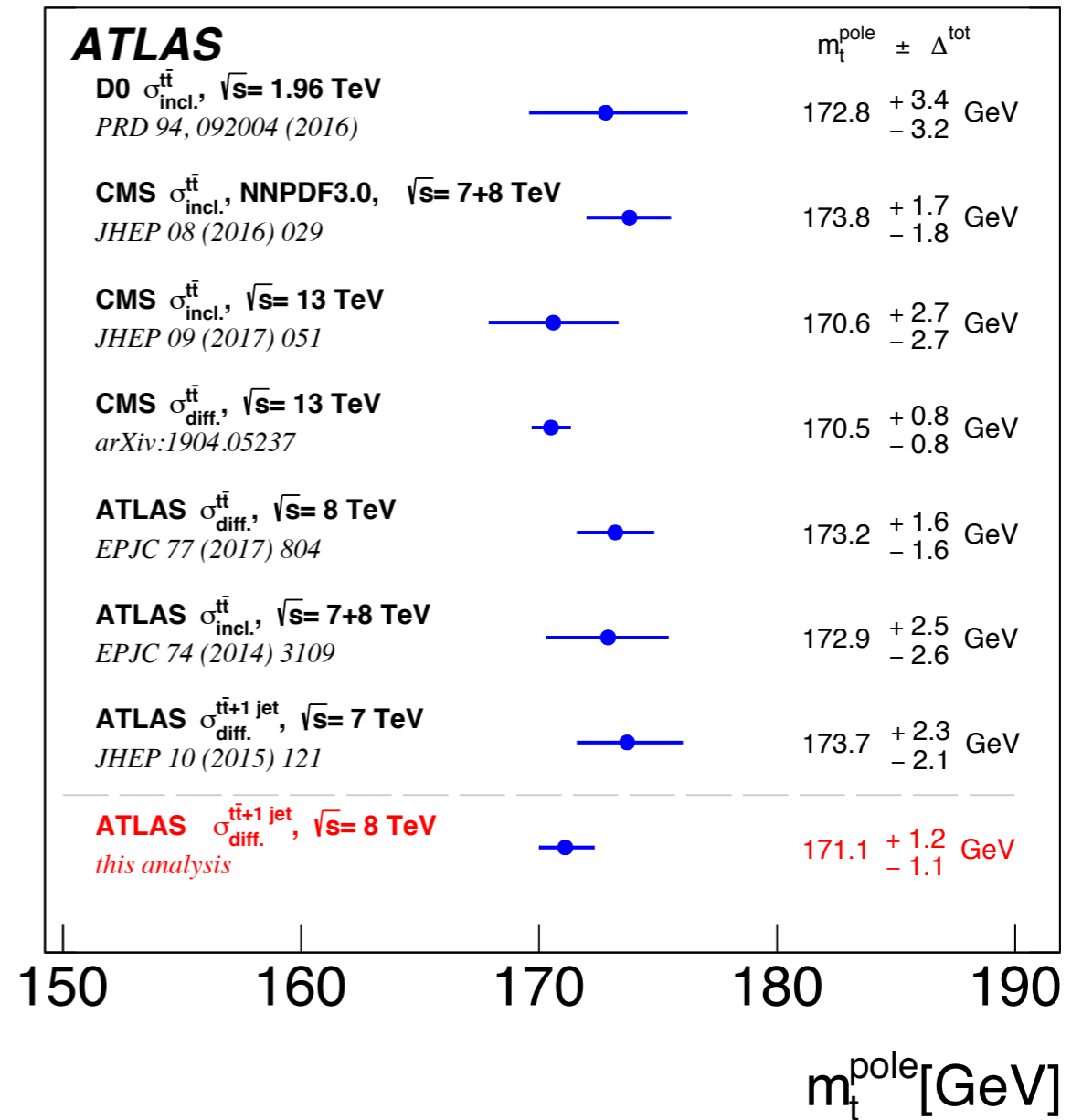
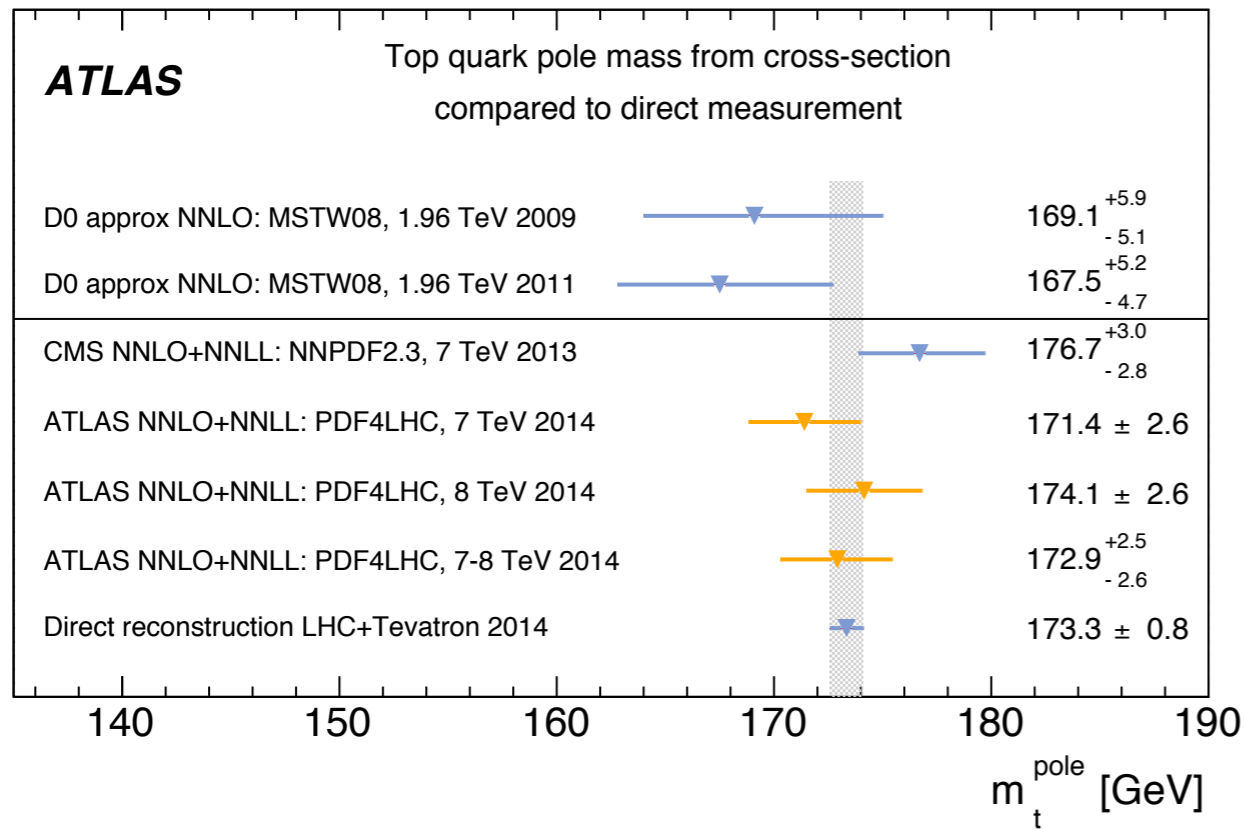


One possibility is to use the inclusive cross-section.  
This is the safest approach, but least sensitive.  
Would be nice to have an energy scan in  $e^+e^-$ !

# Alternative approaches



How would you measure the mass?

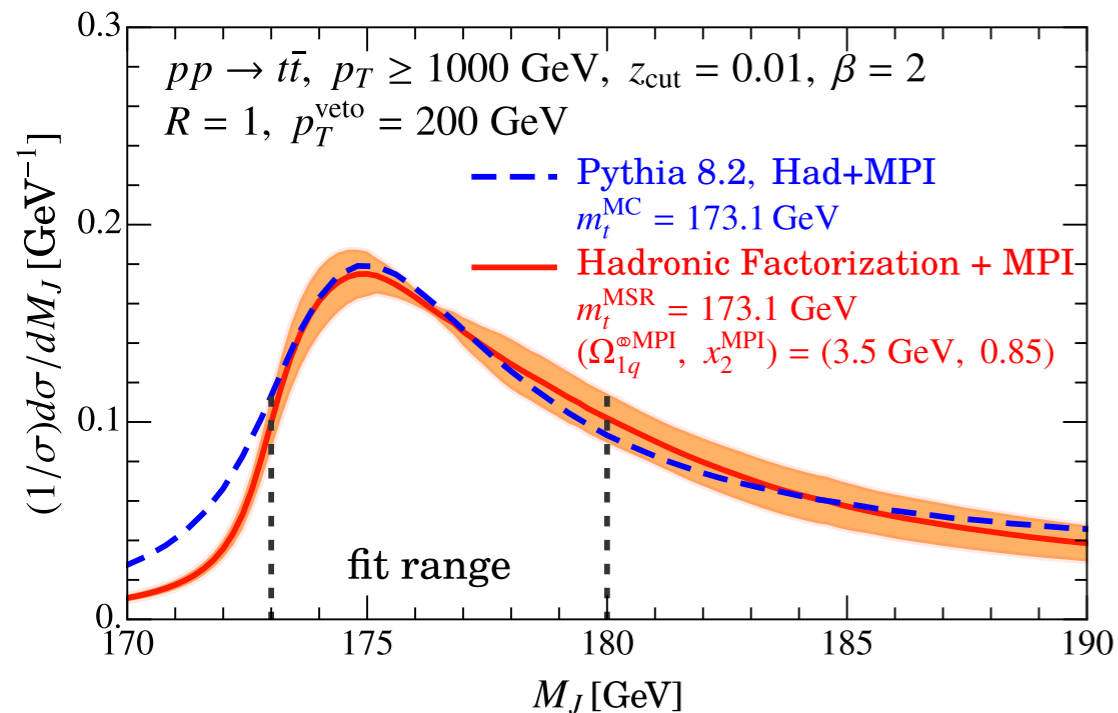
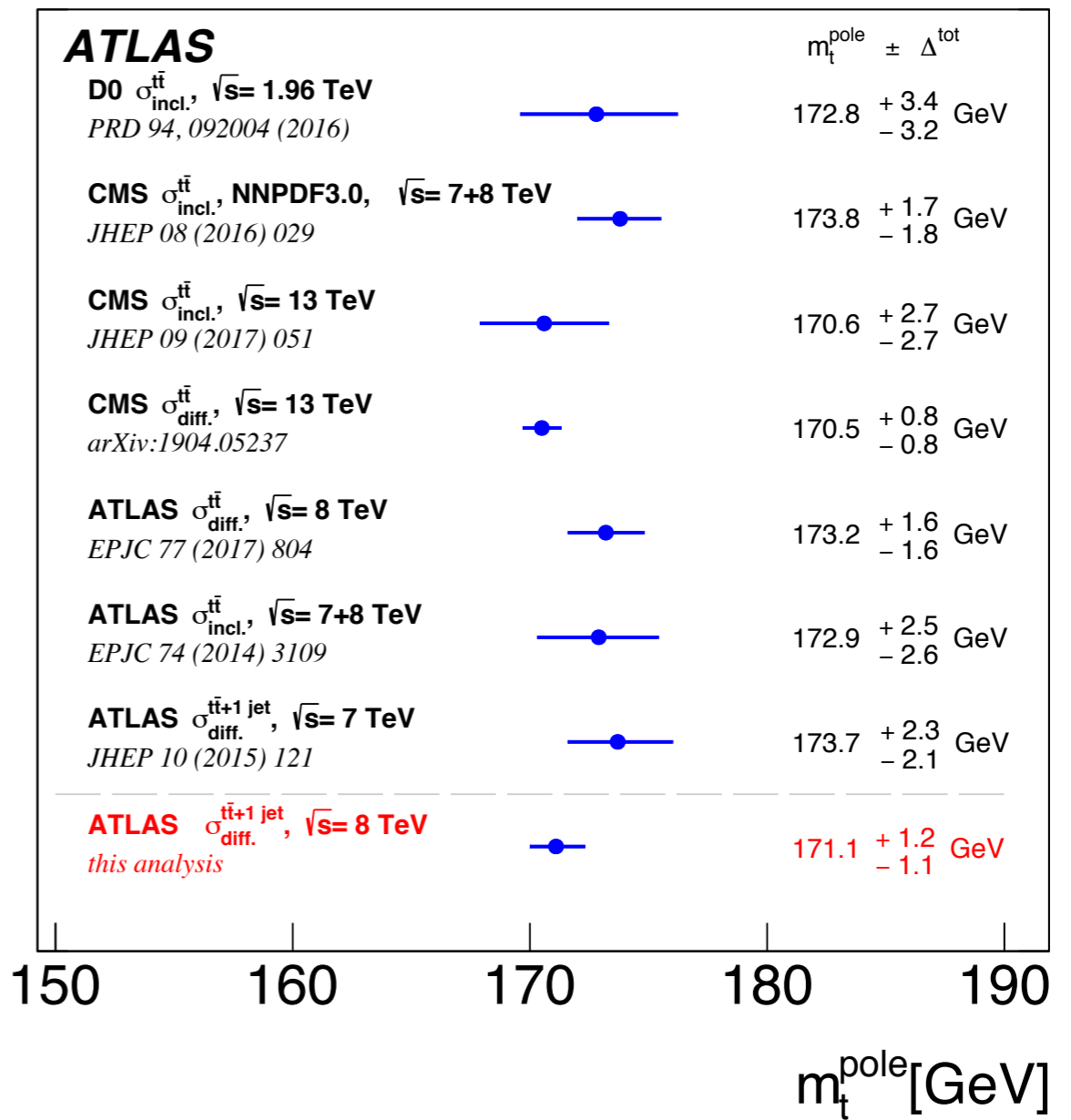
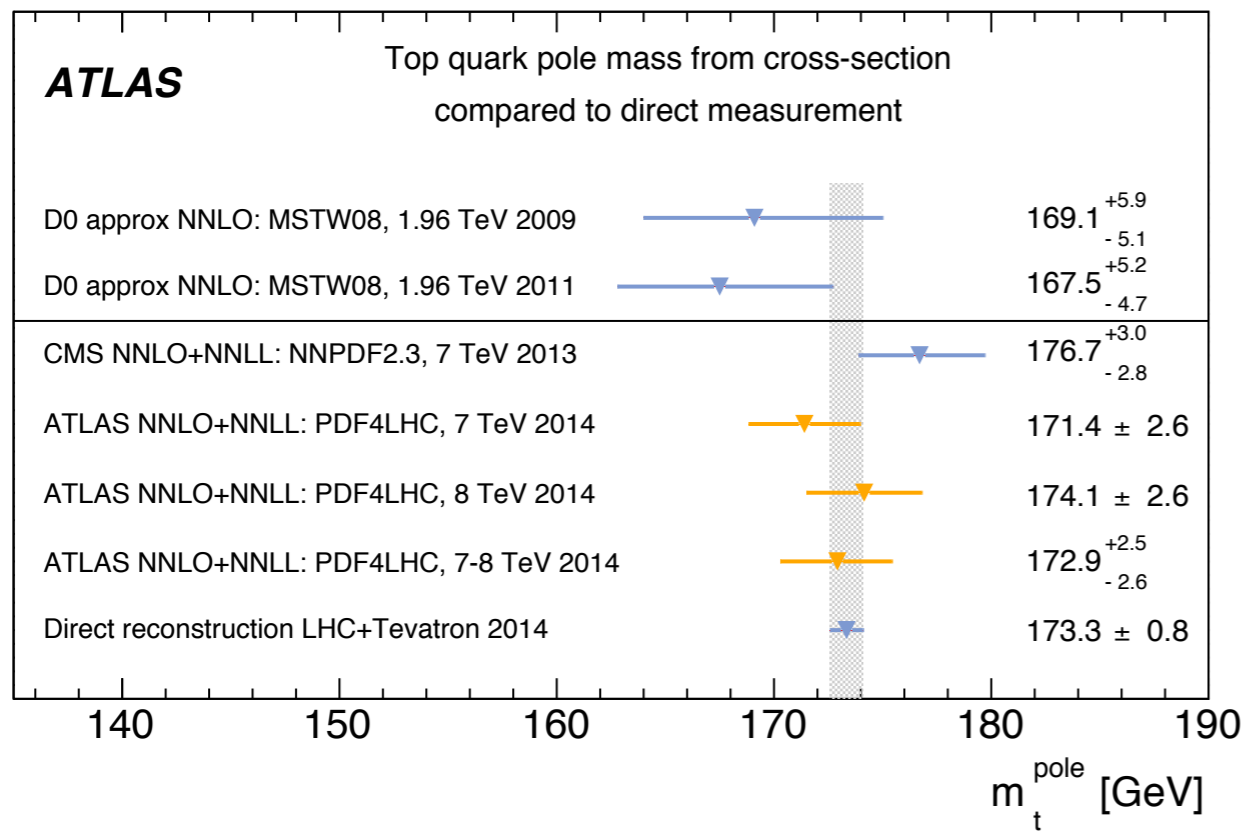


Another possibility is to use **differential cross-sections**. These require **unfolding to parton level** in order to compare with calculations\*.

\*also true for the inclusive cross section, but less sensitive.

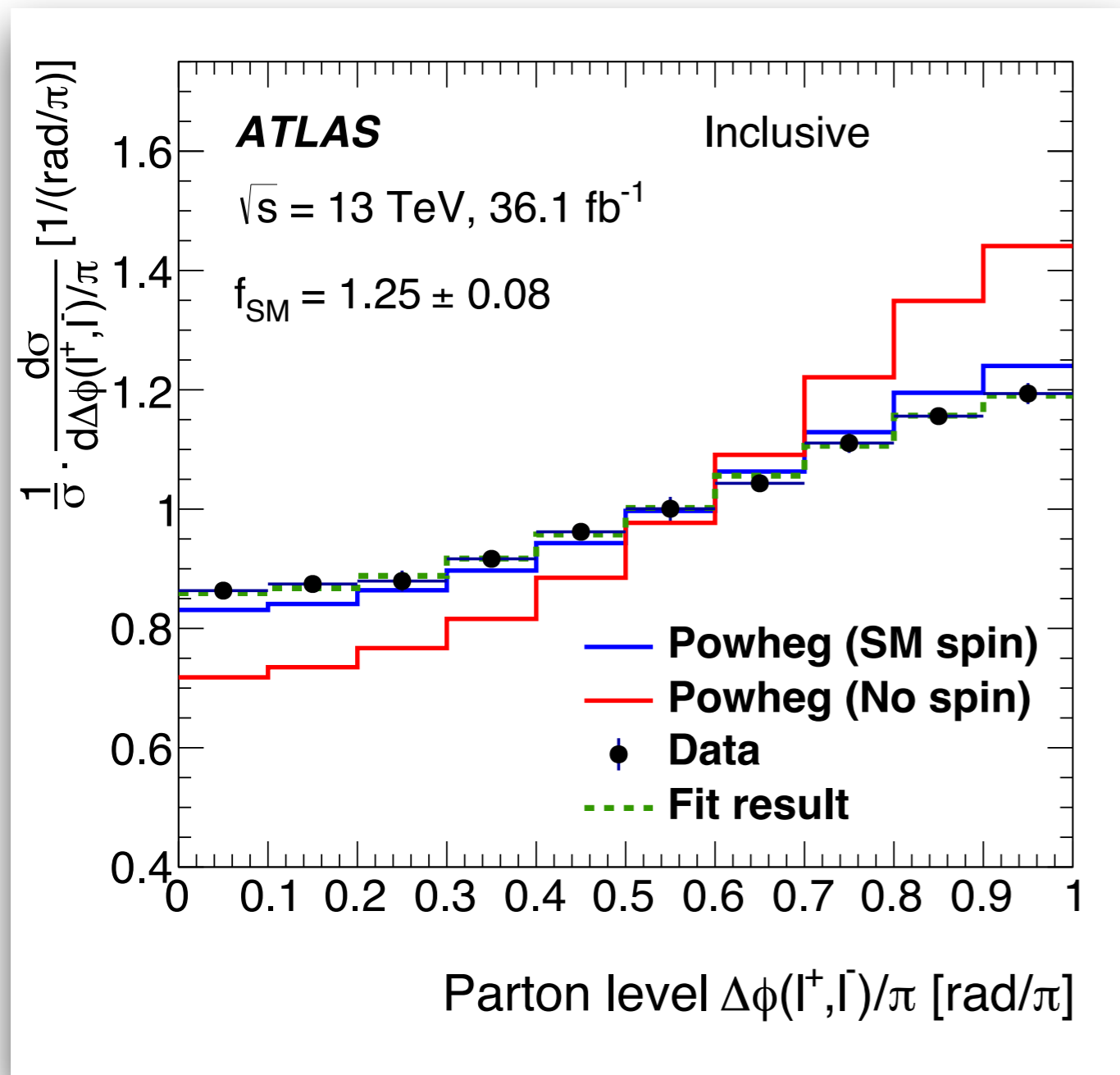
# Alternative approaches

How would you measure the mass?



What about combining theoretical soundness + kinematic shape + jet substructure ??

The top quark has many interesting properties that can be measured at the LHC.

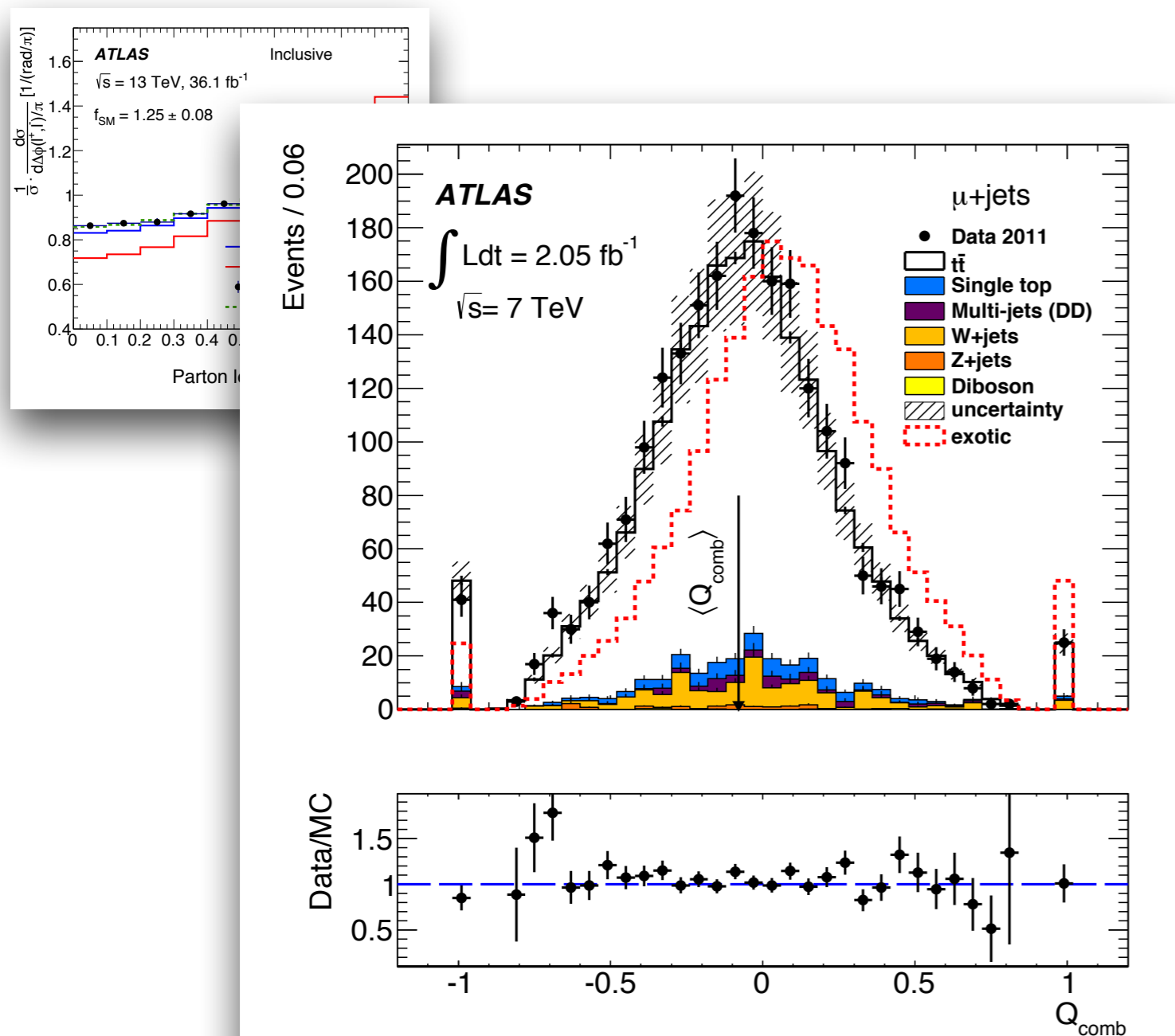


Spin correlations via angular correlations in dilepton events

(BSM may add an uncorrelated component - more on that later)

Unfolded to parton level

The top quark has many interesting properties that can be measured at the LHC.



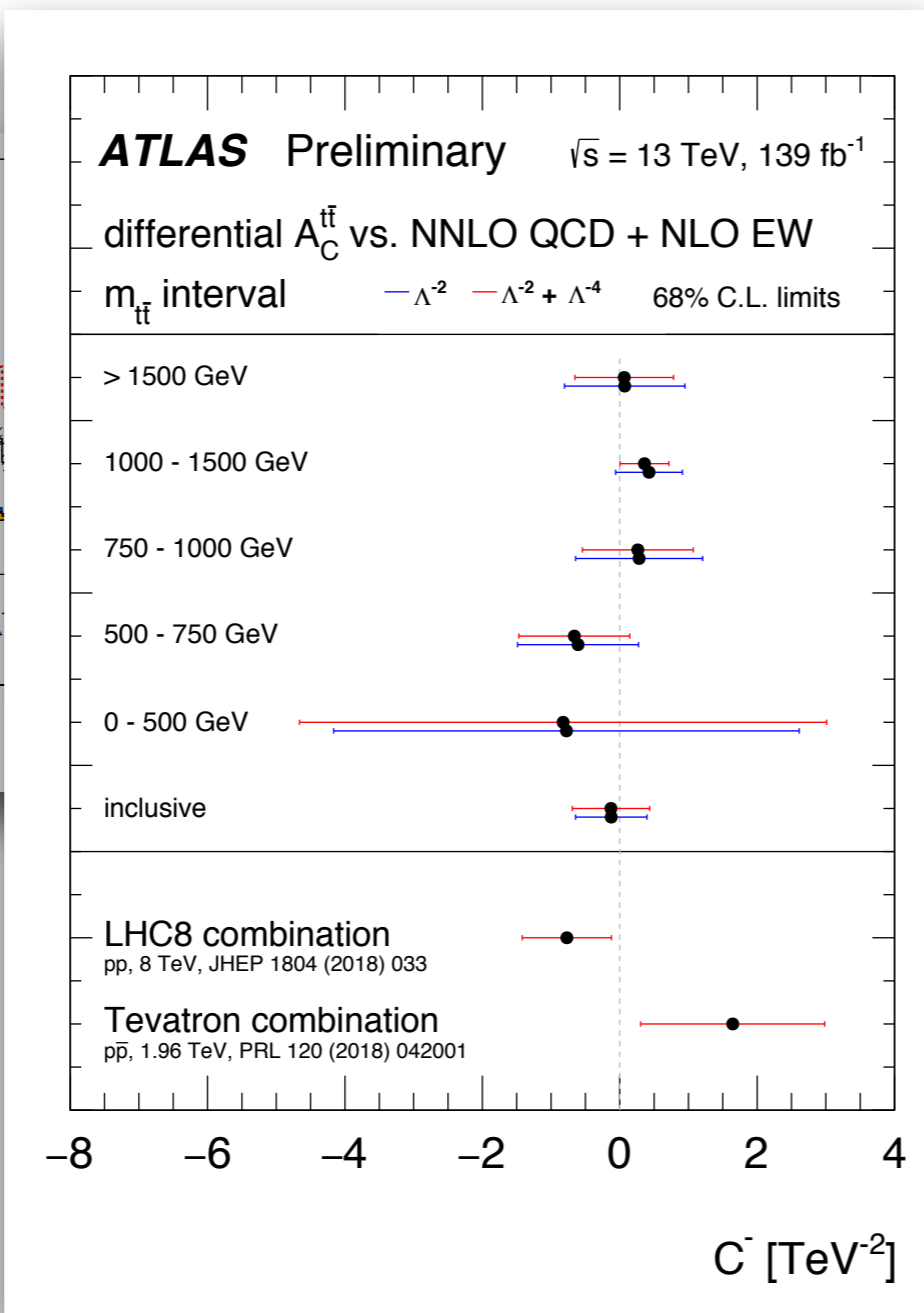
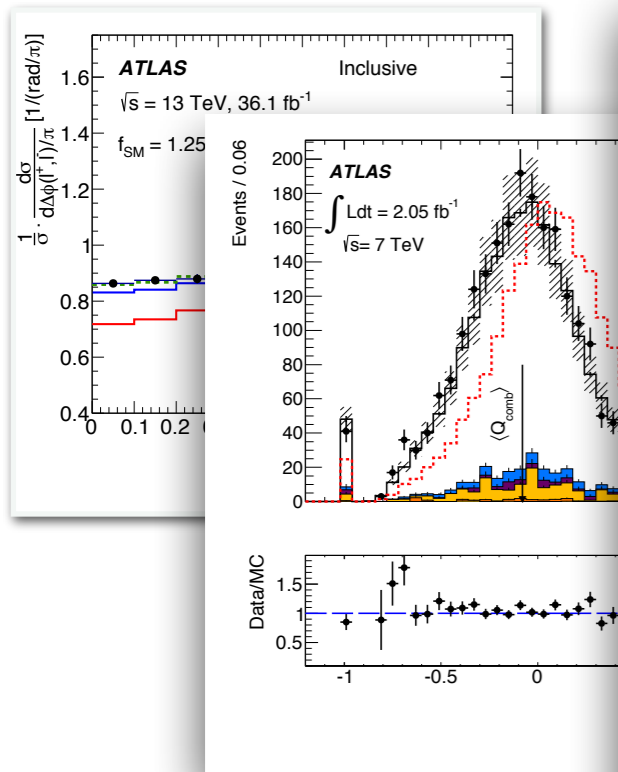
We are quite sure of the charge now, but early in the LHC, it was fashionable to rule out exotic electric charges.

(requires methods for estimating the “jet charge”)



# Beyond the mass

The top quark has many interesting properties that can be measured at the LHC.



The Tevatron observed an excess in the forward-backward asymmetry (now ~gone)

This spurred a big program to probe the “charge asymmetry”

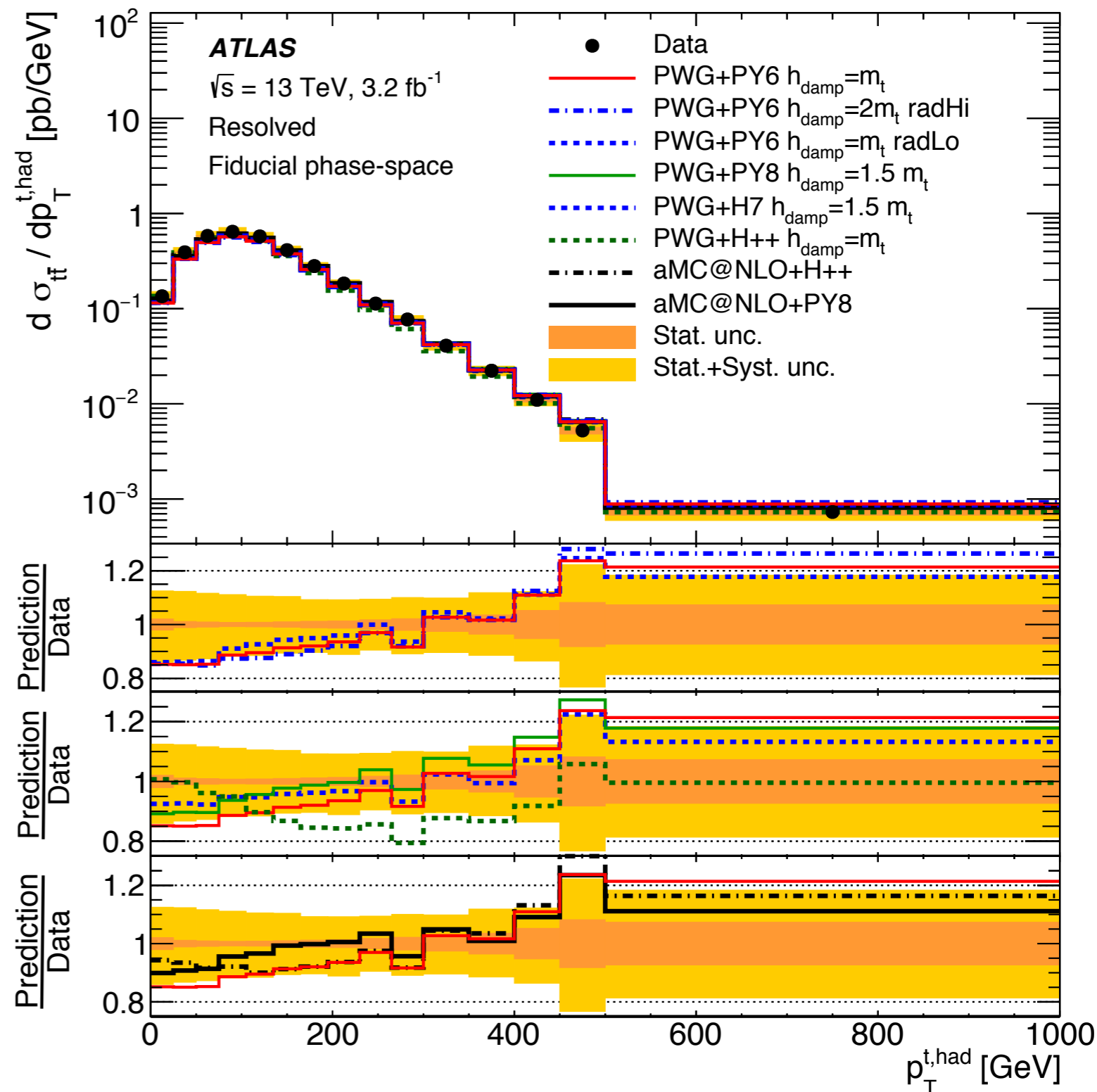
$$A_C = \frac{N(\Delta|y| > 0) - N(\Delta|y| < 0)}{N(\Delta|y| > 0) + N(\Delta|y| < 0)}$$

Which is much less sensitive than the Tevatron because we are gg dominated.

As top-quark pair production is a background for many searches, it is important to study the kinematic tails.

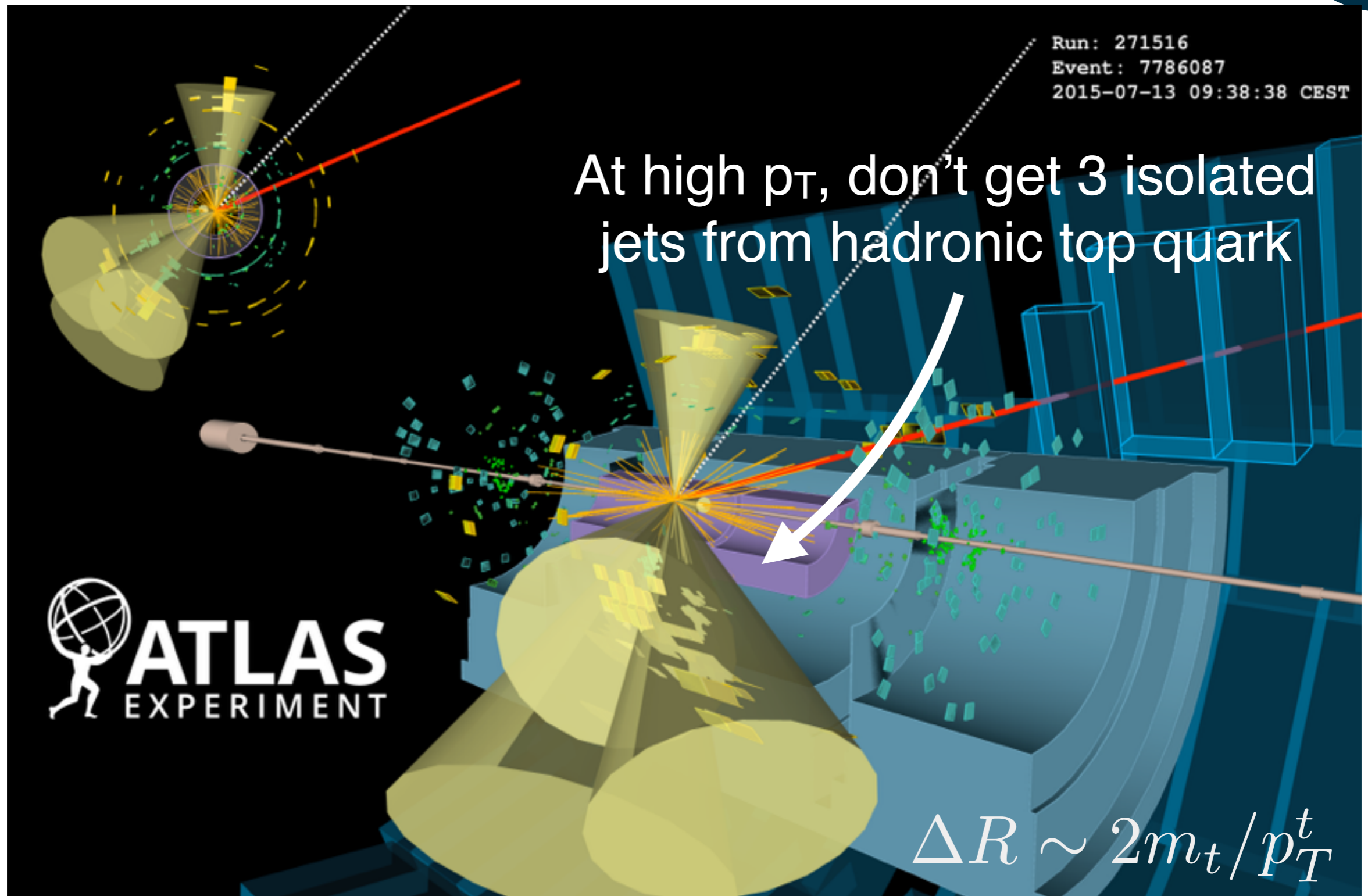
Long-standing observation that the spectrum is harder in simulation than data.

This seems to be partially accounted for at NNLO, but full story still under investigation.



# Boosted top quarks

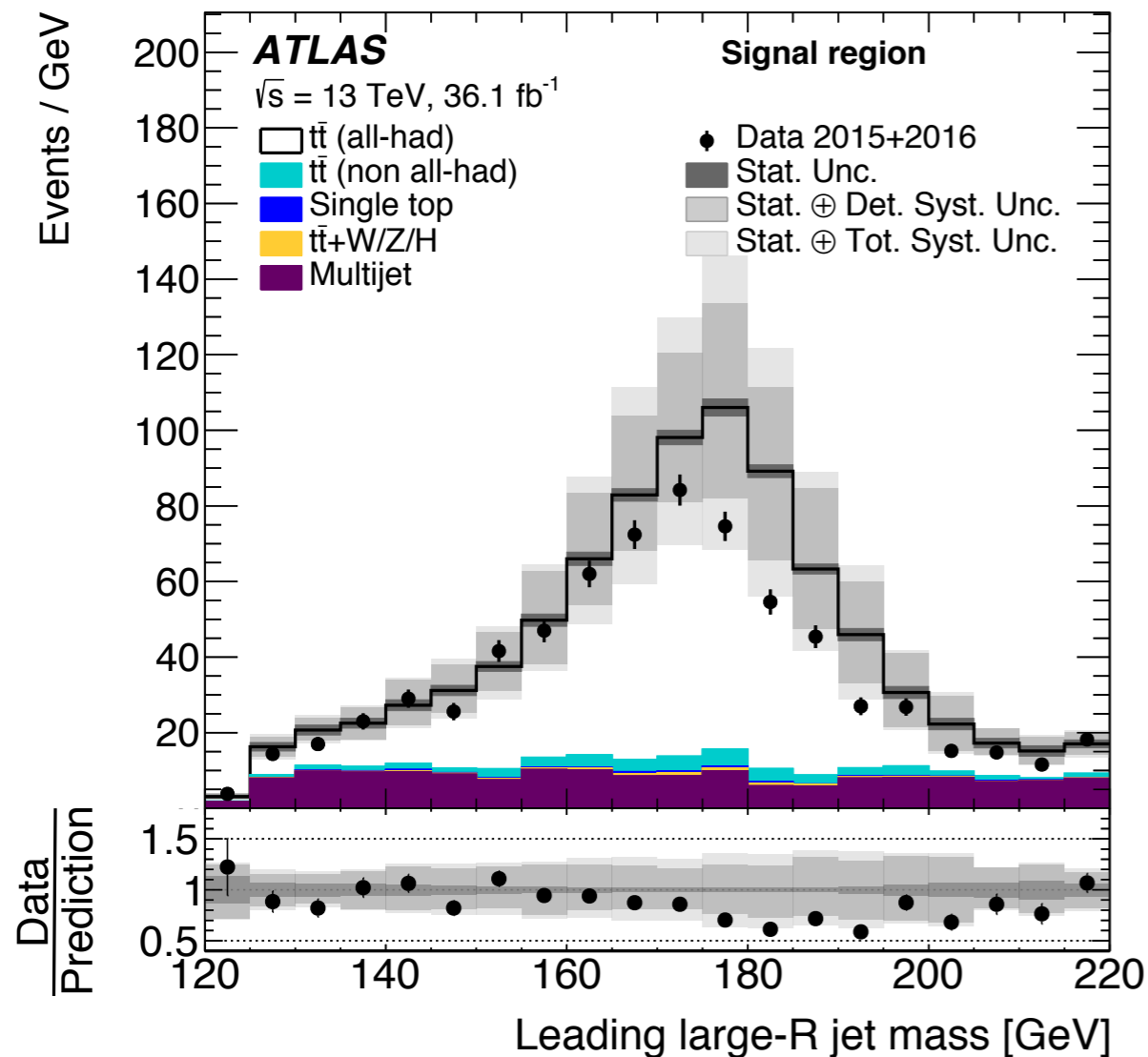
19



(where does this come from?)

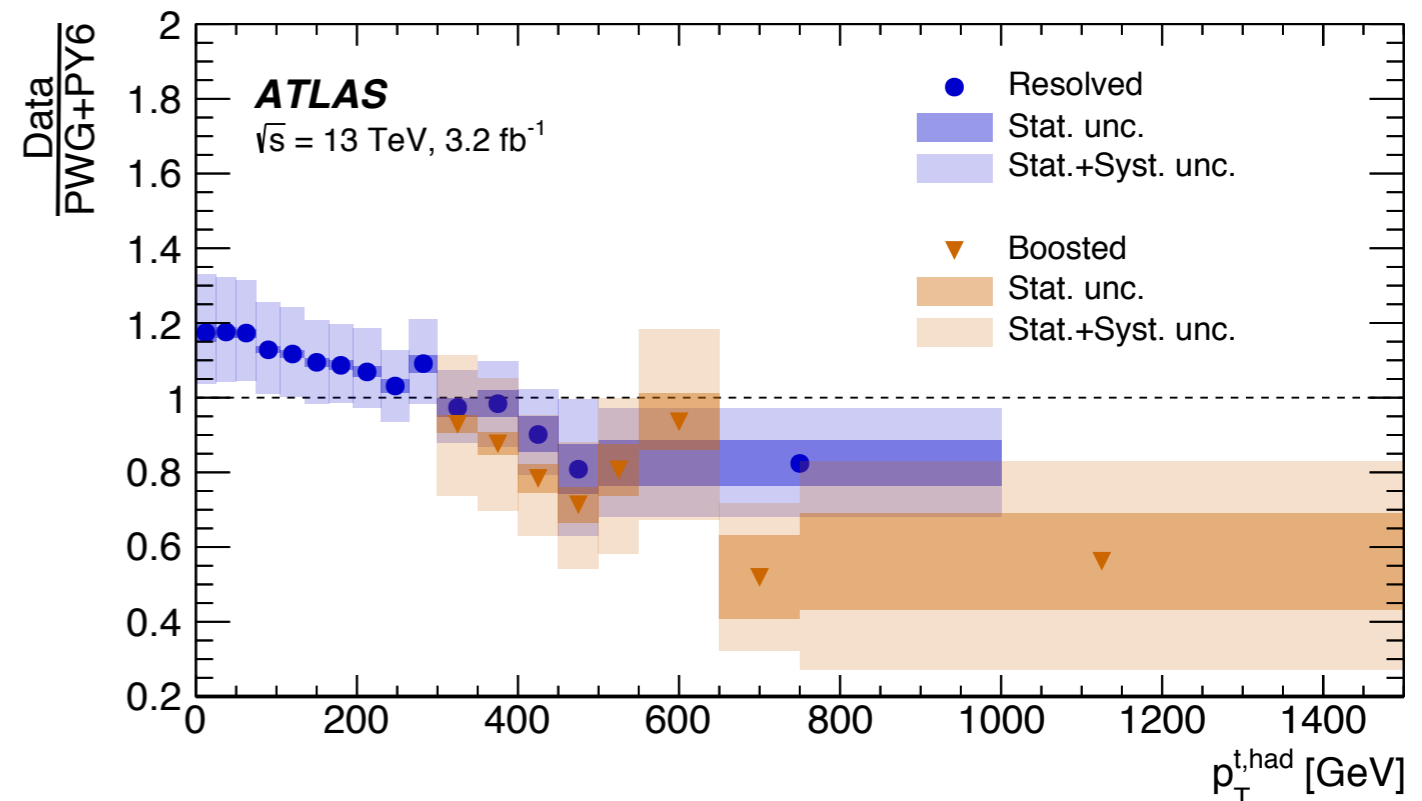
# Boosted top quarks

20



In the “boosted” regime, need to use large( $r$ )-radius jets

These jets have a 3-prong substructure and mass  $\sim m_t$



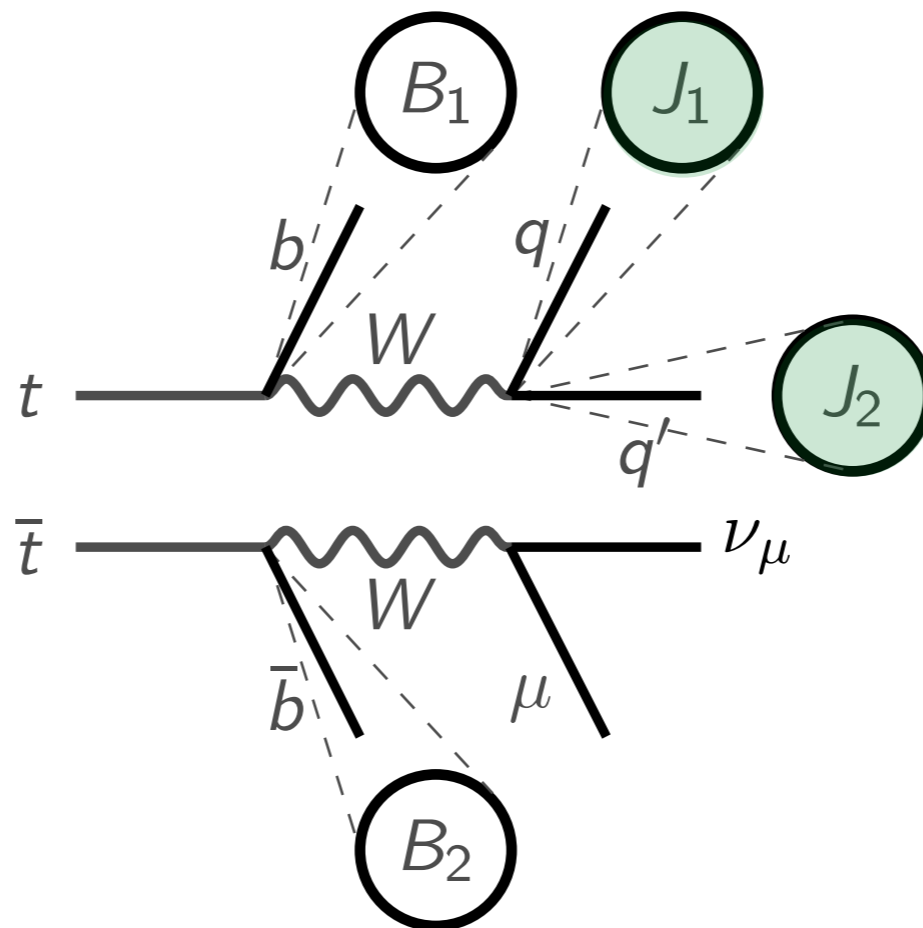
Jet substructure tagging is therefore important part for SM measurements, not just BSM searches!

# Tops as a source of color singlets

21

As the  $W$  bosons produced from top quarks are on-shell, top quark pairs offer unique LEP-like laboratory for QCD studies.

One can identify ('tag') events using the leptonic  $W$  decay and the study ('probe') the structure of the hadronic  $W$  decay.

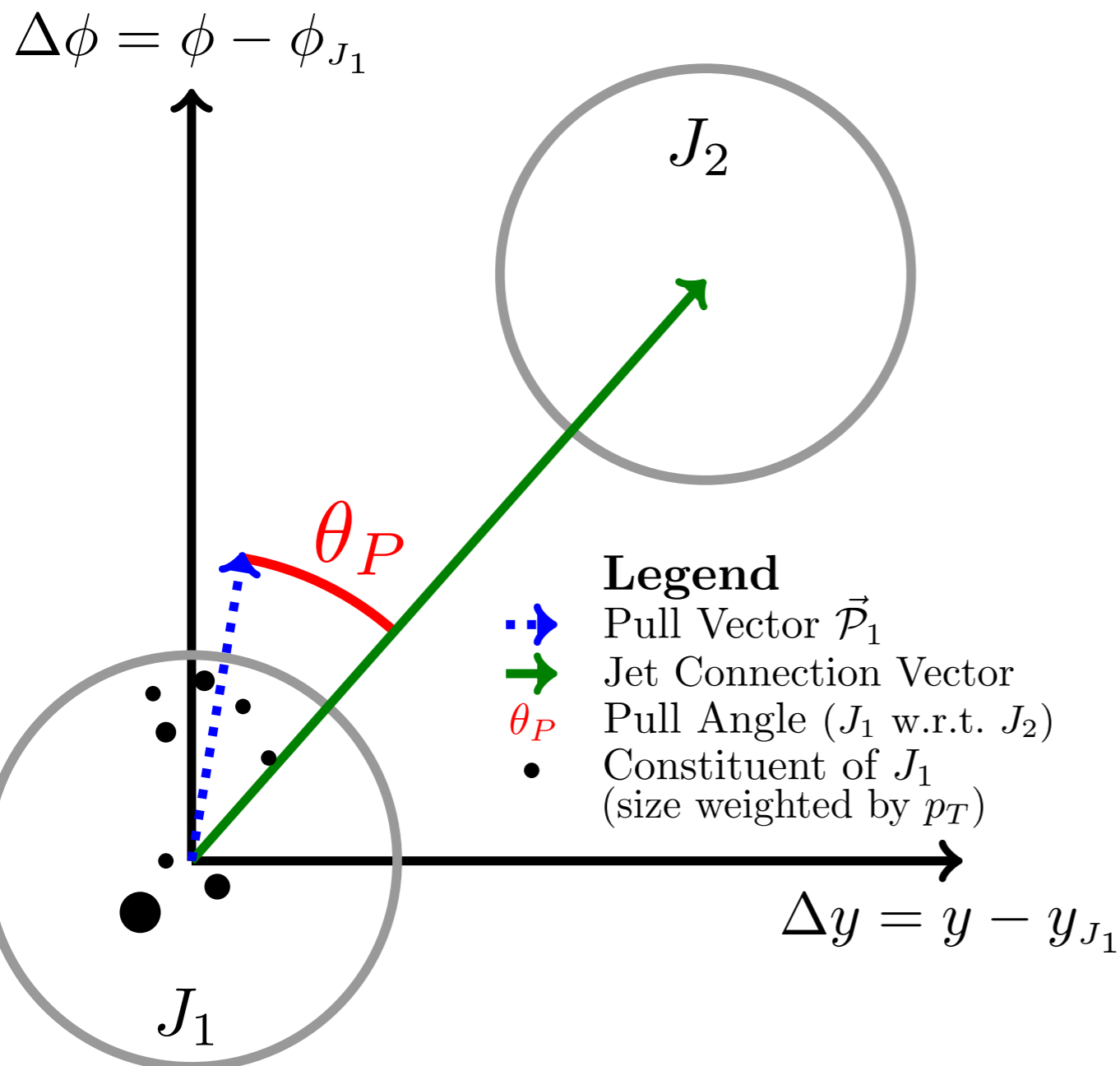


# Tops as a source of “singlet jets”

The only way to get a clean sample of jet originating from a singlet decay (since LEP)

Step 1: How singlet-like is it?

Tool: Jet superstructure - the interplay between jet substructure and global radiation / kinematics.



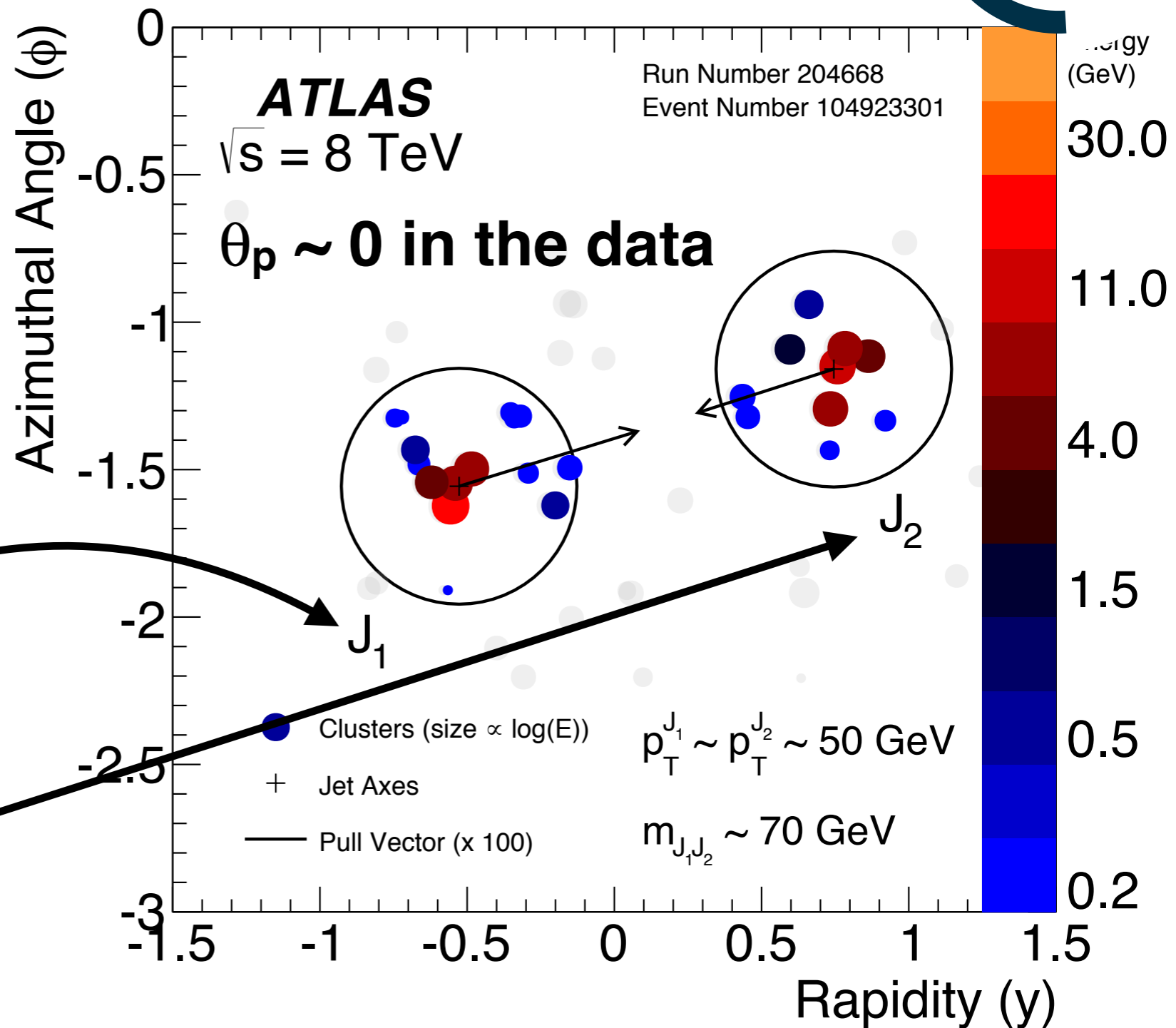
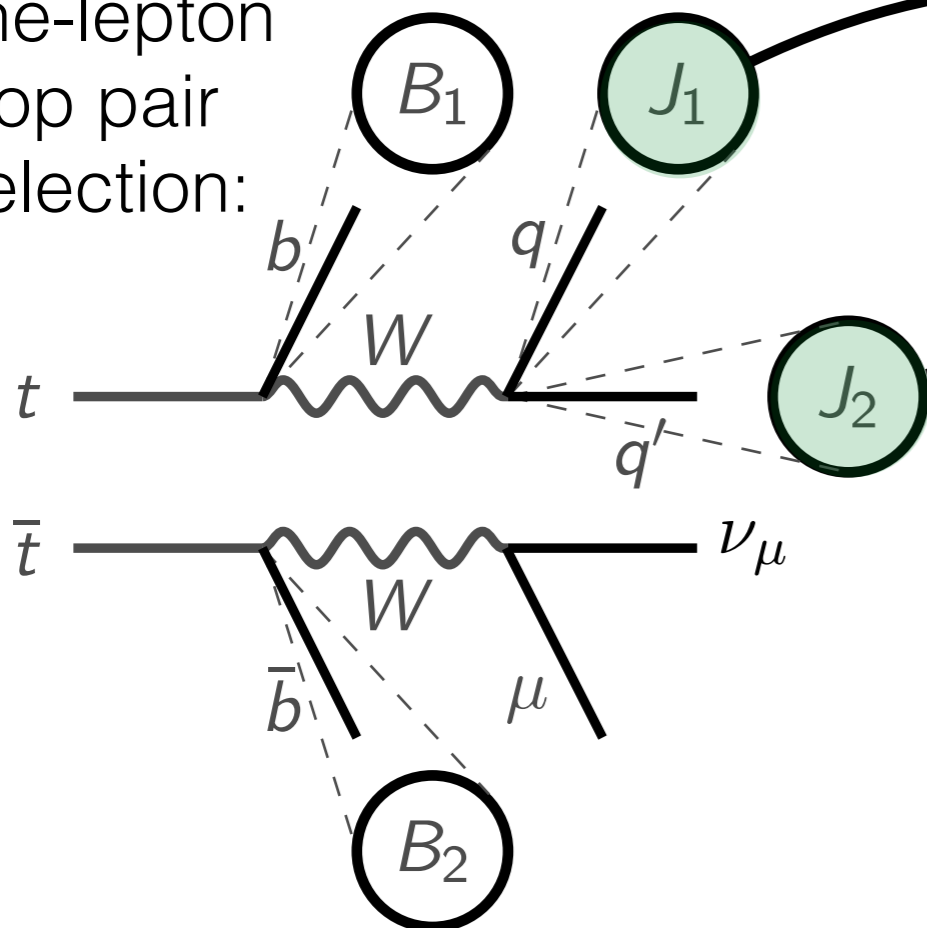
$$\vec{\Phi}(J) = \sum_{i \in J} \frac{|\vec{\Delta r}_i| \cdot p_T^i}{p_T^J} \vec{\Delta r}_i$$

Question: how much does the radiation from one jet lean towards the other?

# Hadronic Correlations: Jet Superstructure

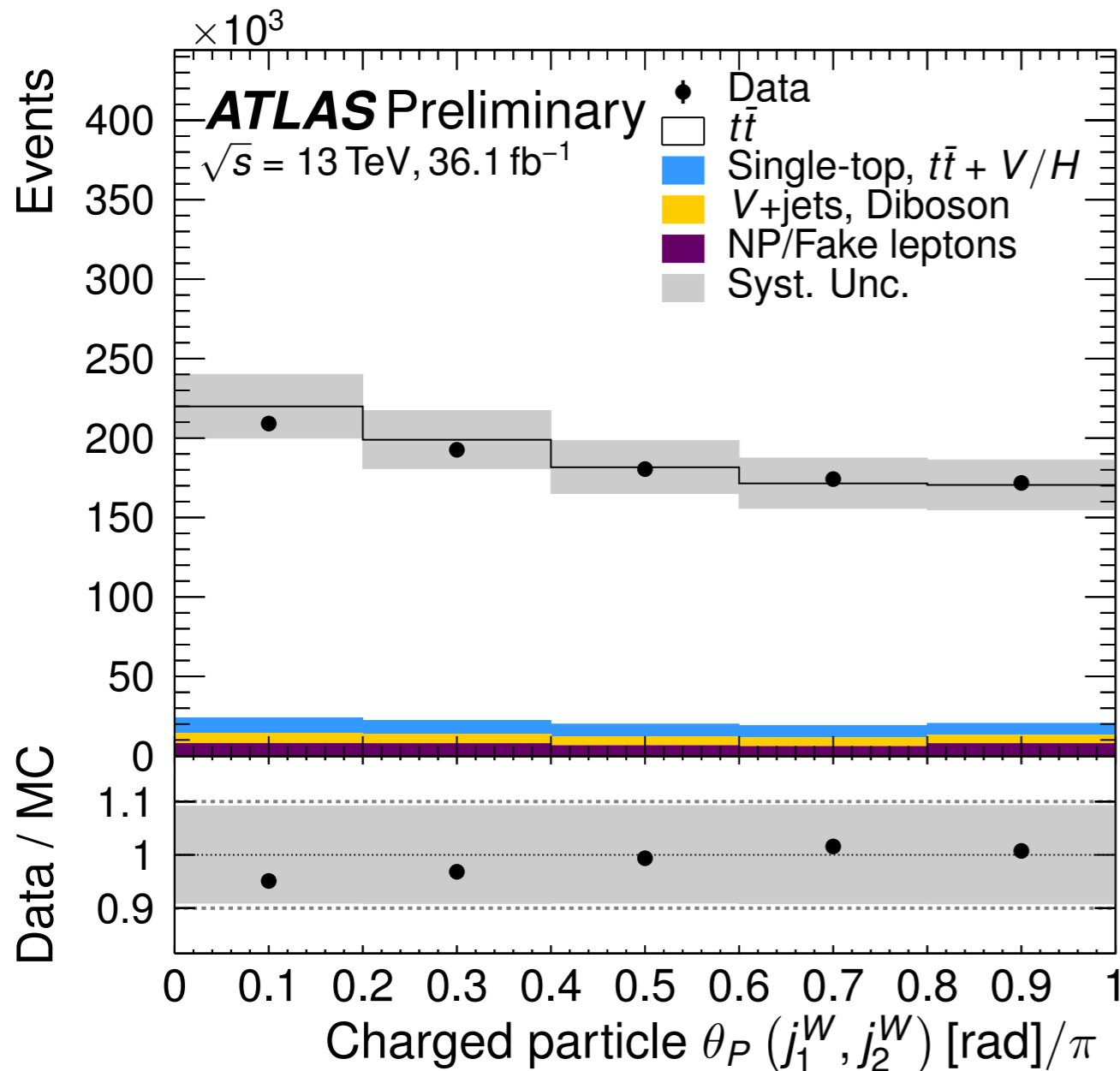
23

One-lepton  
top pair  
selection:

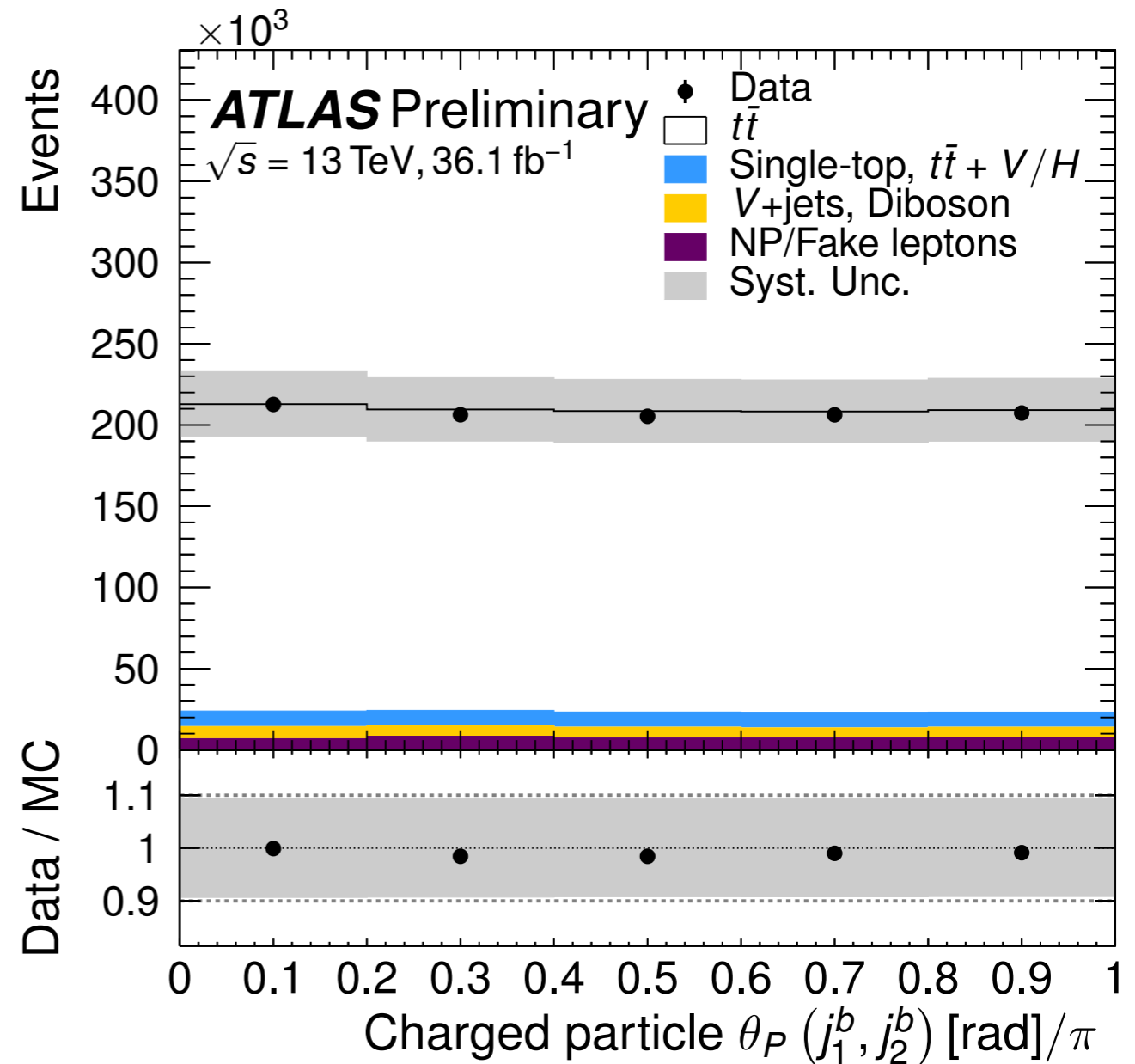


# Hadronic Correlations: Jet Superstructure

24



peak at 0 - W daughters  
are “connected”

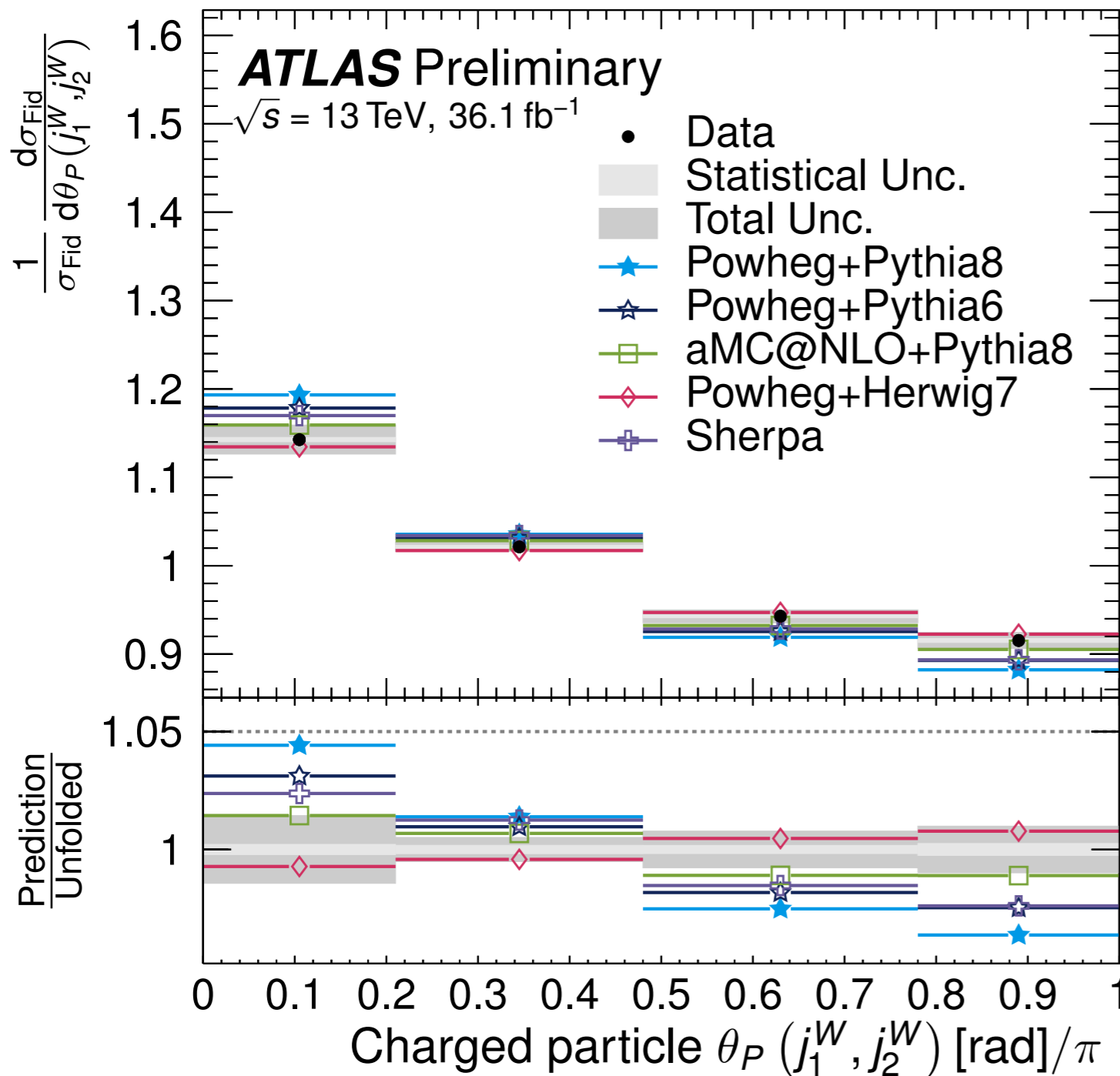


no peak - b's are  
not connected



# Jet Pull for MC Tuning

25

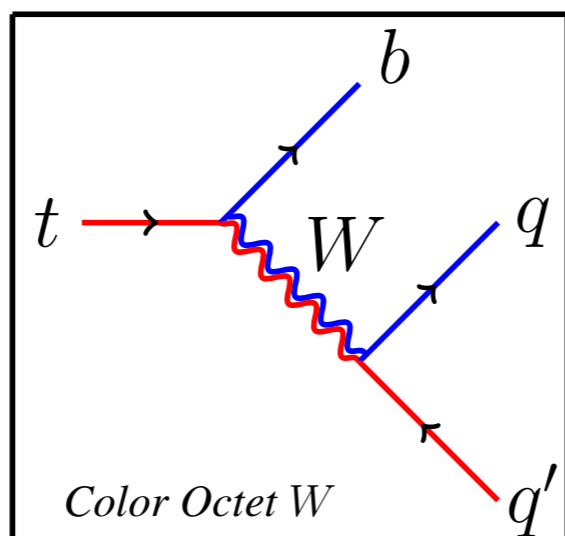
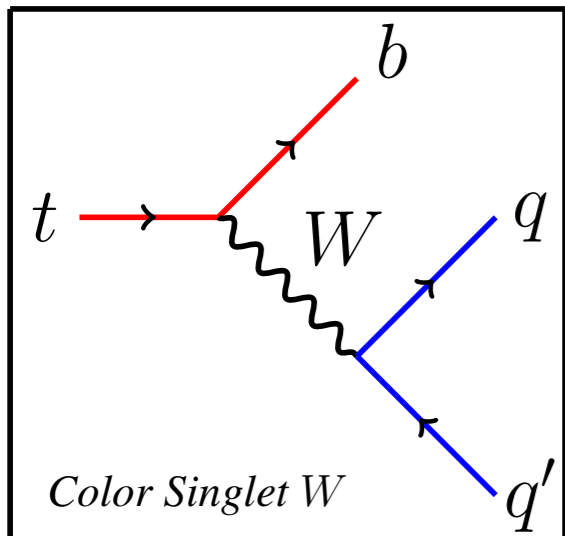


Challenging measurement, but we can reach to the 1%-level for precision.

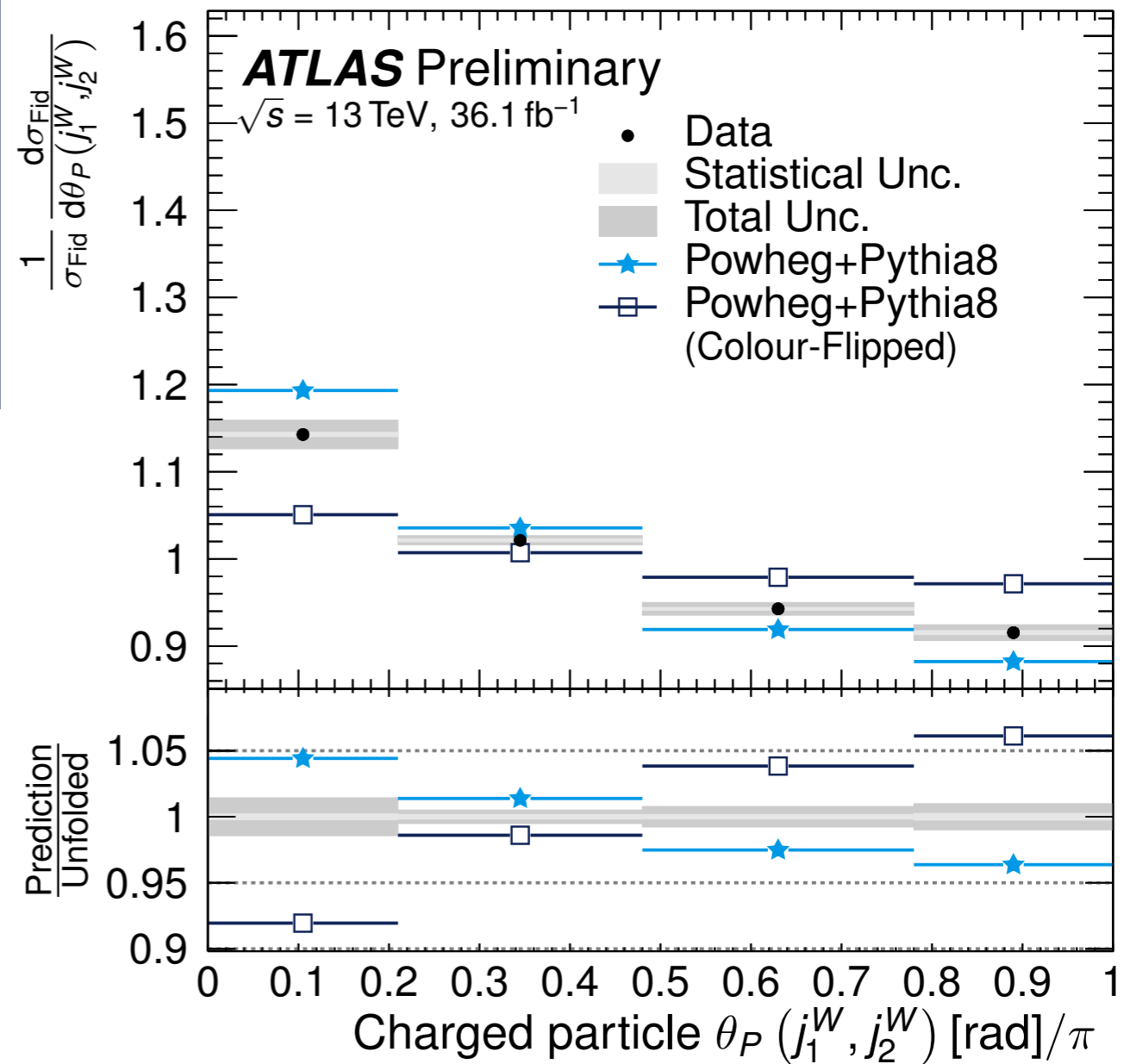
Interestingly, there is a rather large spread in the MC predictions.

(useful for tuning!)

# Jet Pull: Much more to learn!

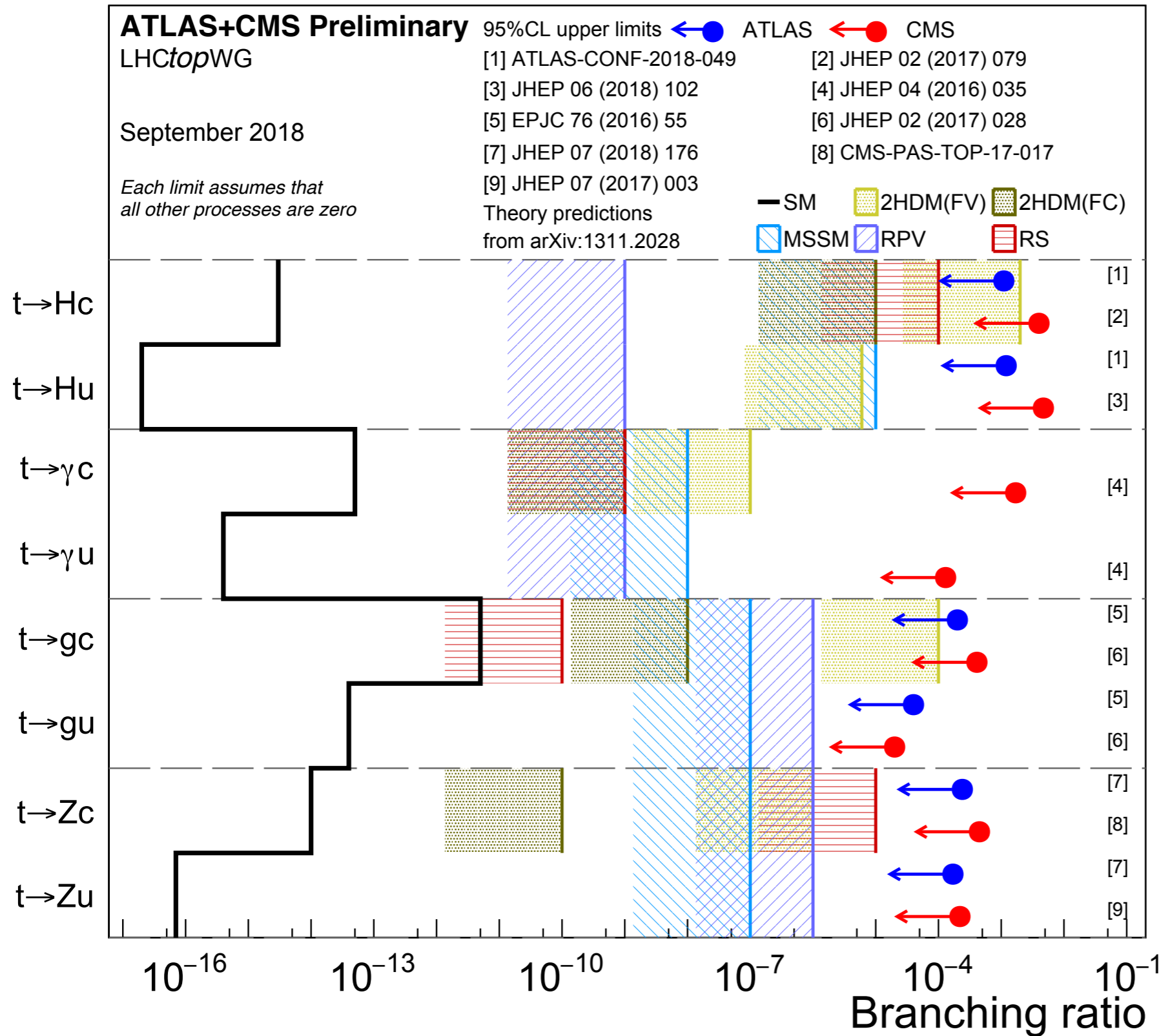


An important message: if we artificially make W's in Pythia octet-like, then we only marginally prefer the singlet.  
clearly there is something we can learn for here!

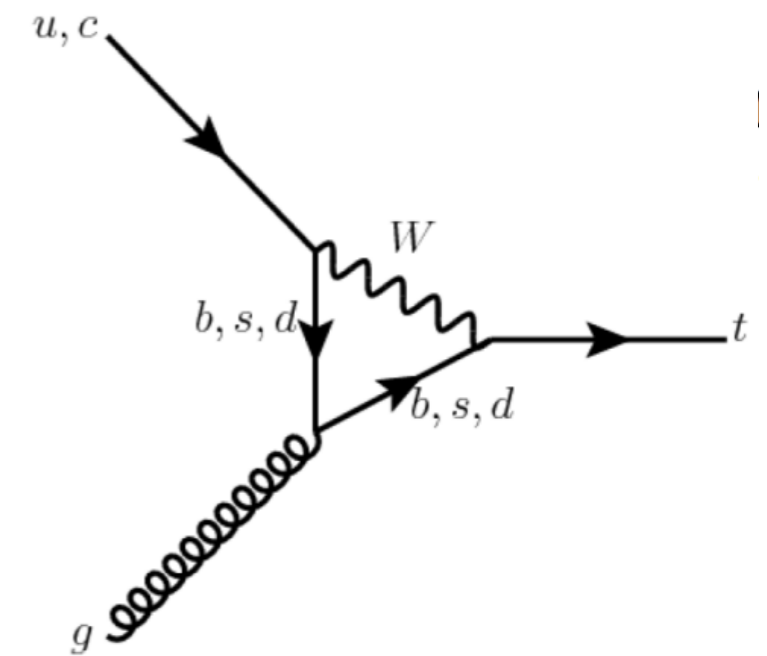


*How do we know the W is not an octet?*

# Back to top properties: decays



top decays to anything other than  $Wb$  are (highly) suppressed in the Standard Model.

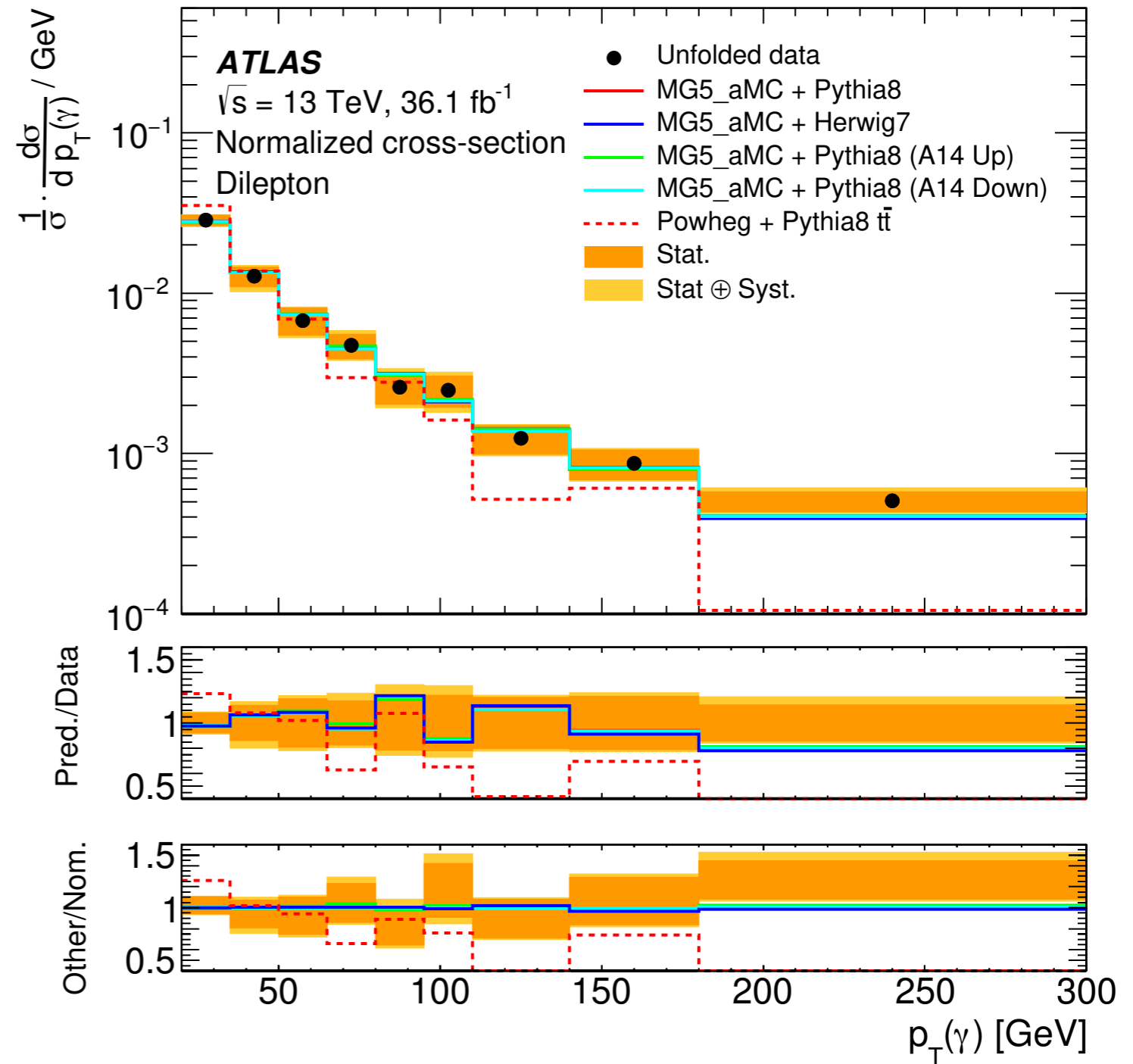
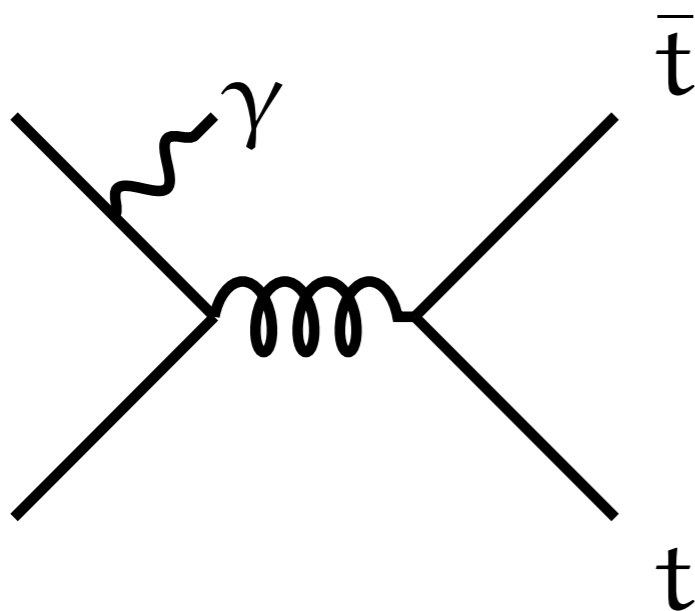


# Associated production: photons

28

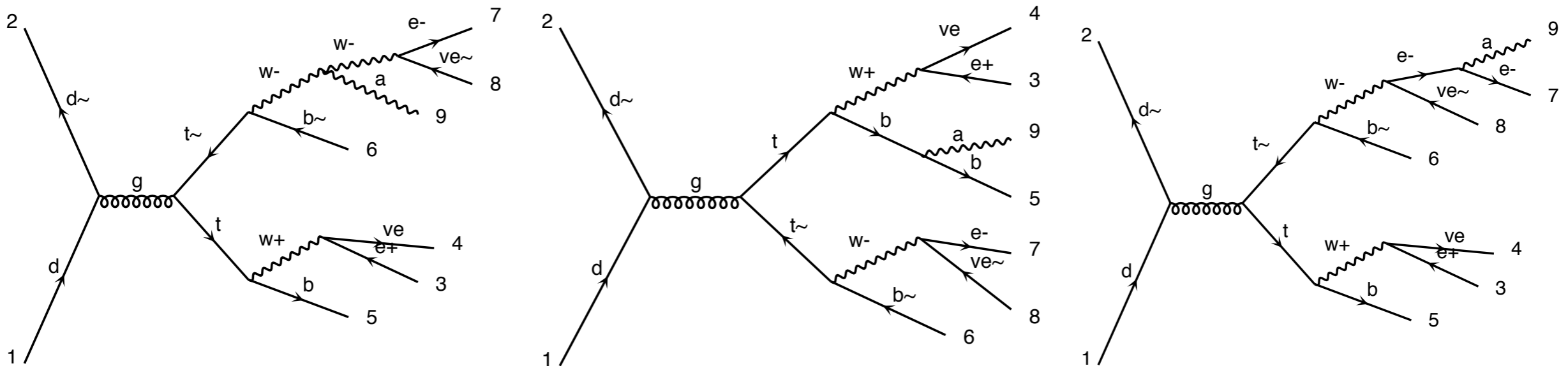
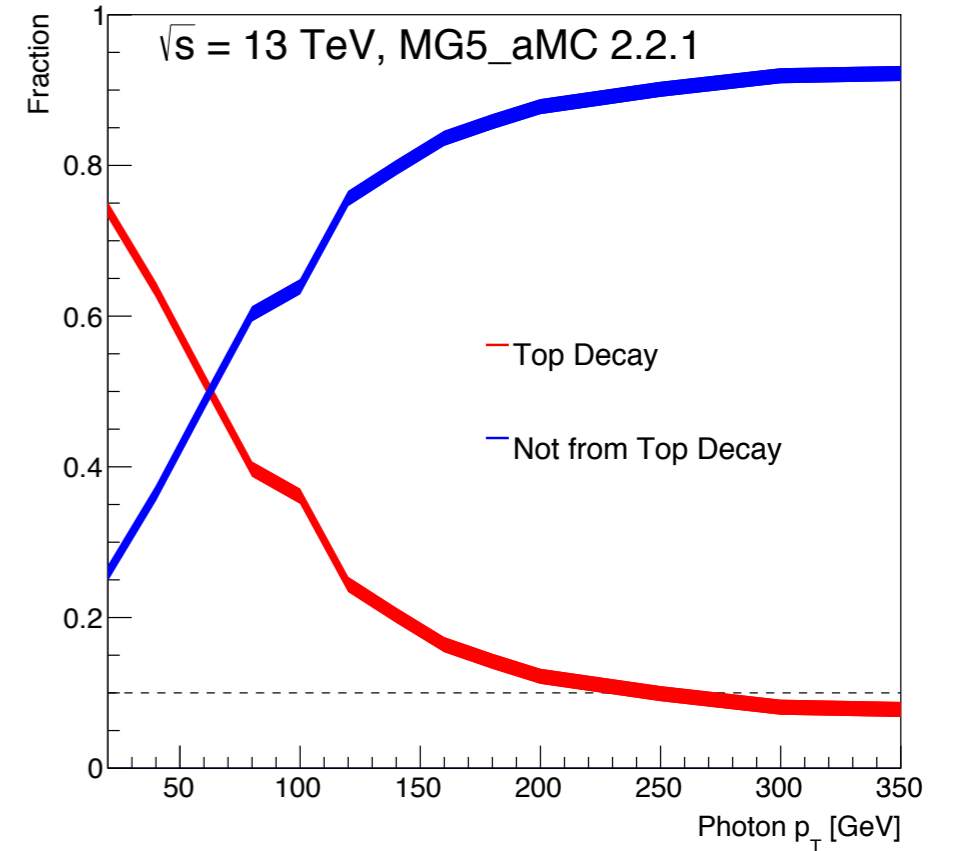
Dominated by  
qq production

*Powerful tool for  
EFT, etc.*



# Associated production: photons

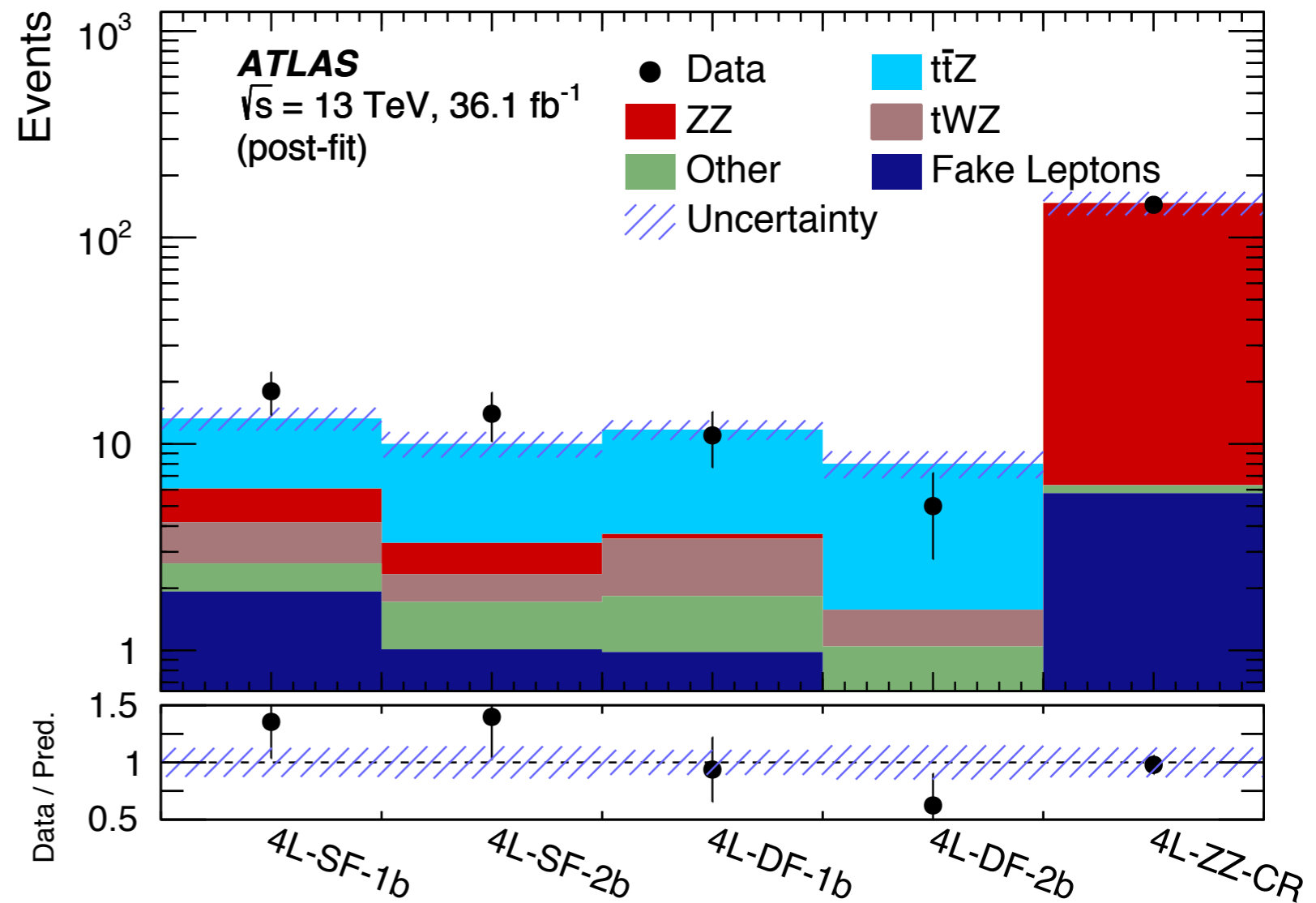
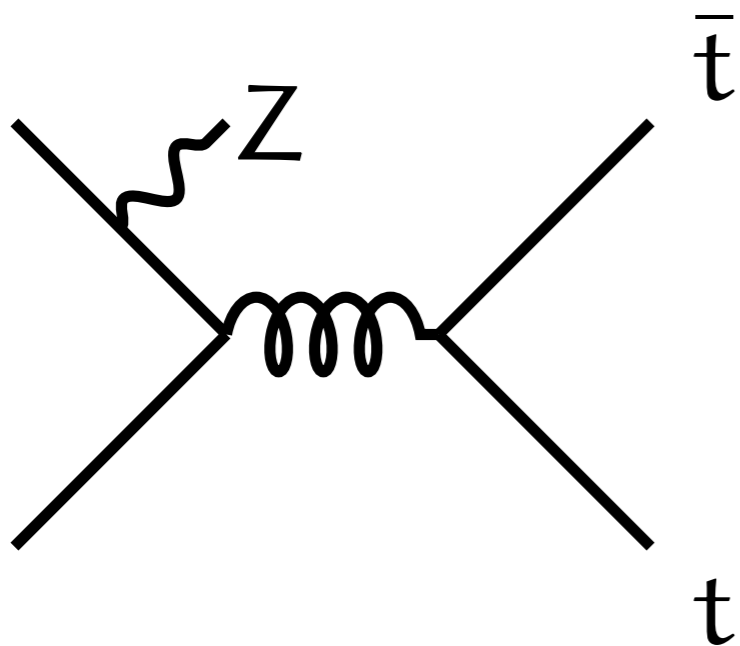
Warning: any electrically charged object can radiate a photon!



# Associated production: W/Z

30

Powerful  
probe for EFTs

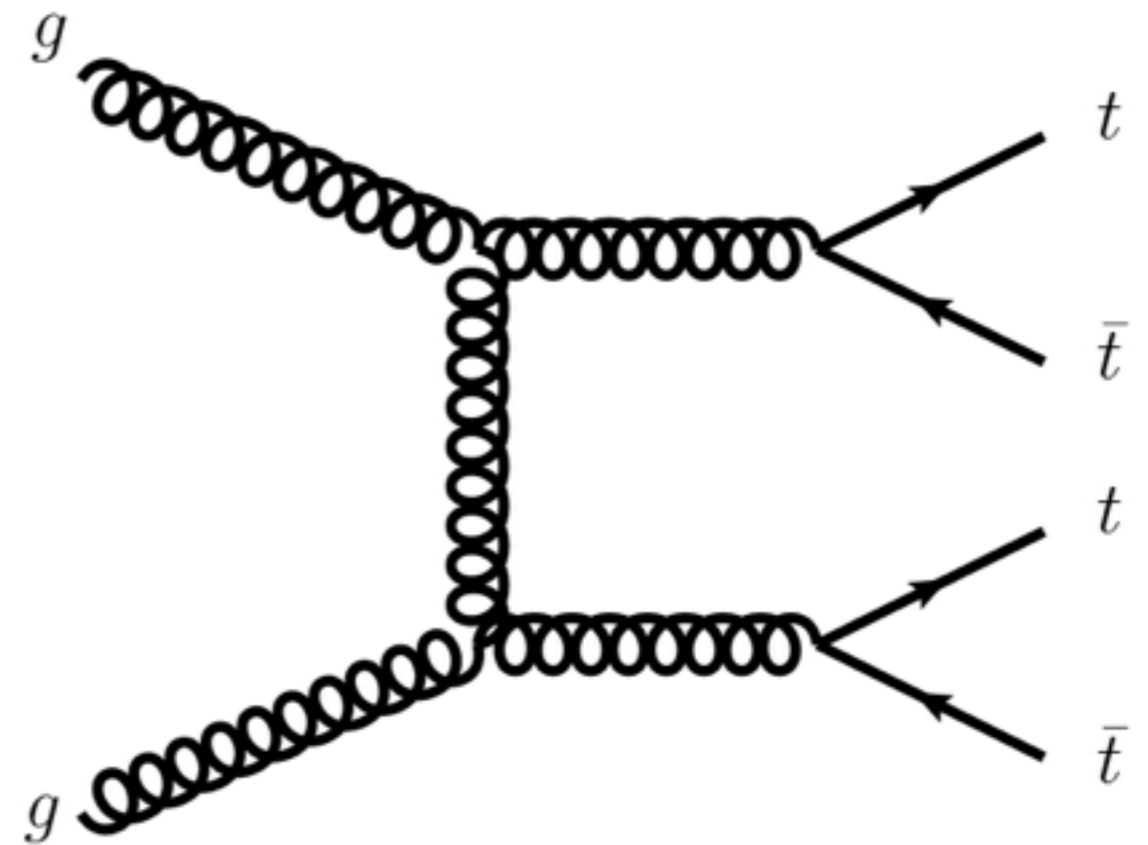
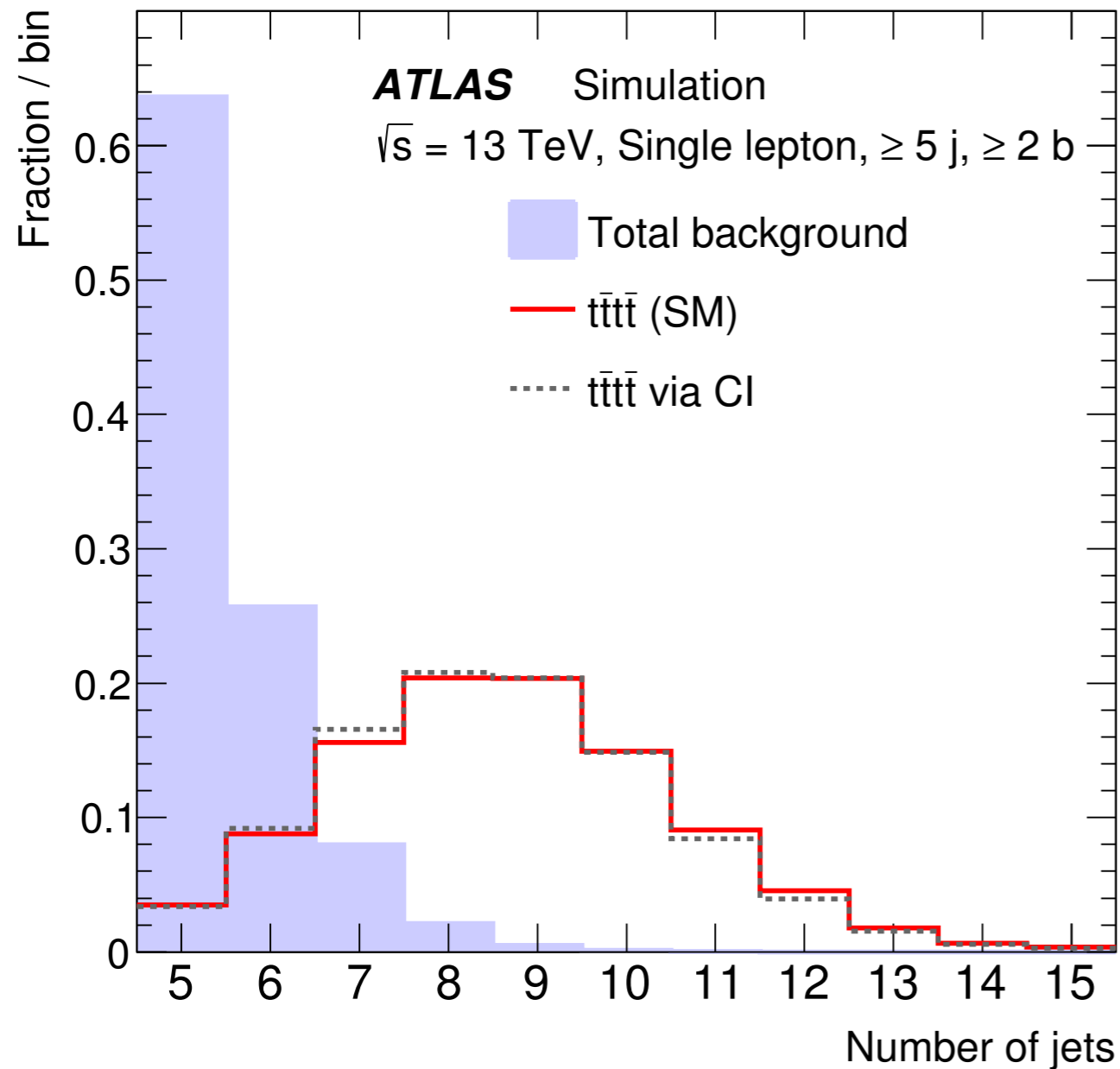


Also a significant background for BSM (both for same-sign lepton searches and for top + MET ( $Z \rightarrow \nu\nu$ ))

# Associated production: many tops

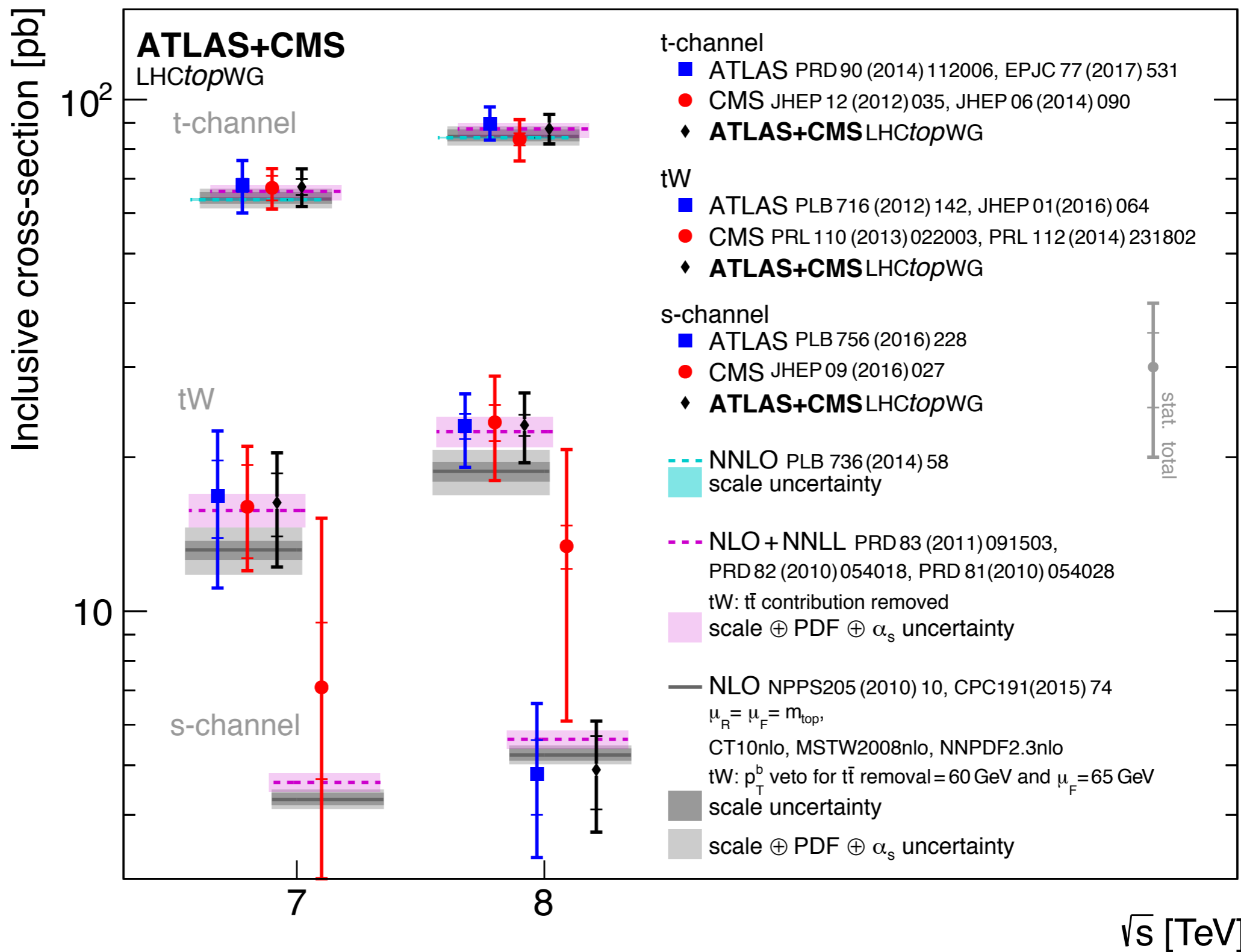
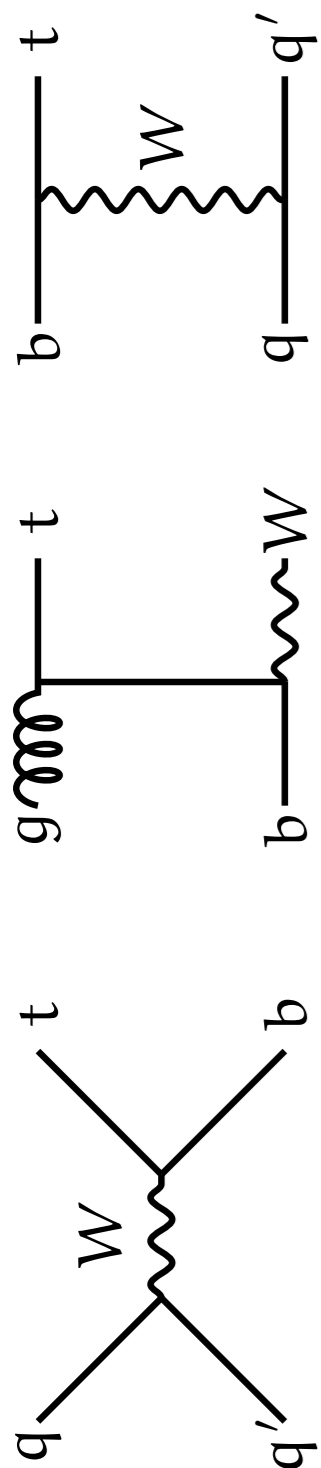
31

Another interesting probe of EFTs with a diversity of final states.



(current sensitivity is  
~few x Standard Model  
prediction of ~10 fb)

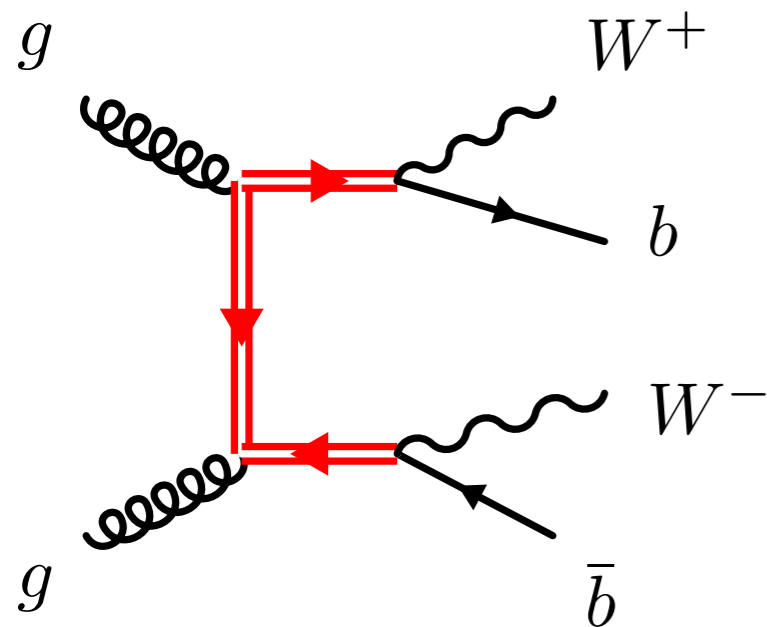
# Single top





# Interference with pair production

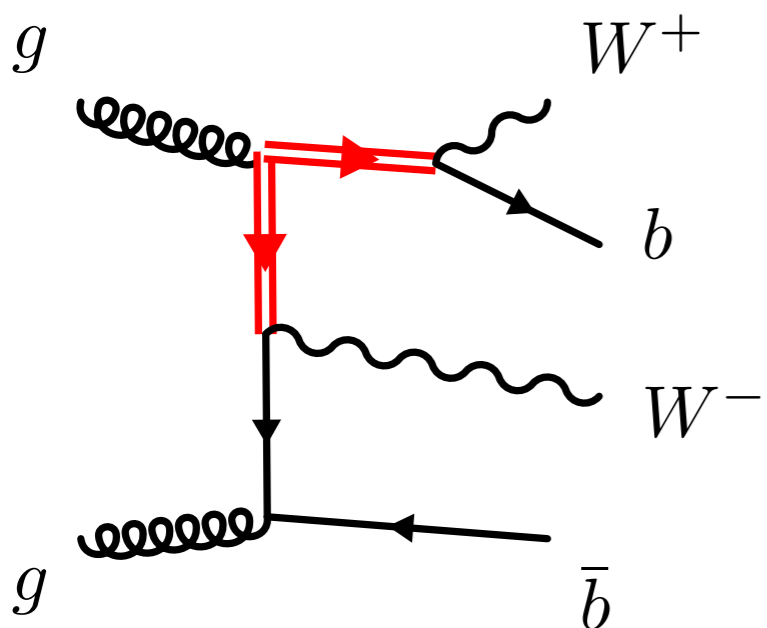
33



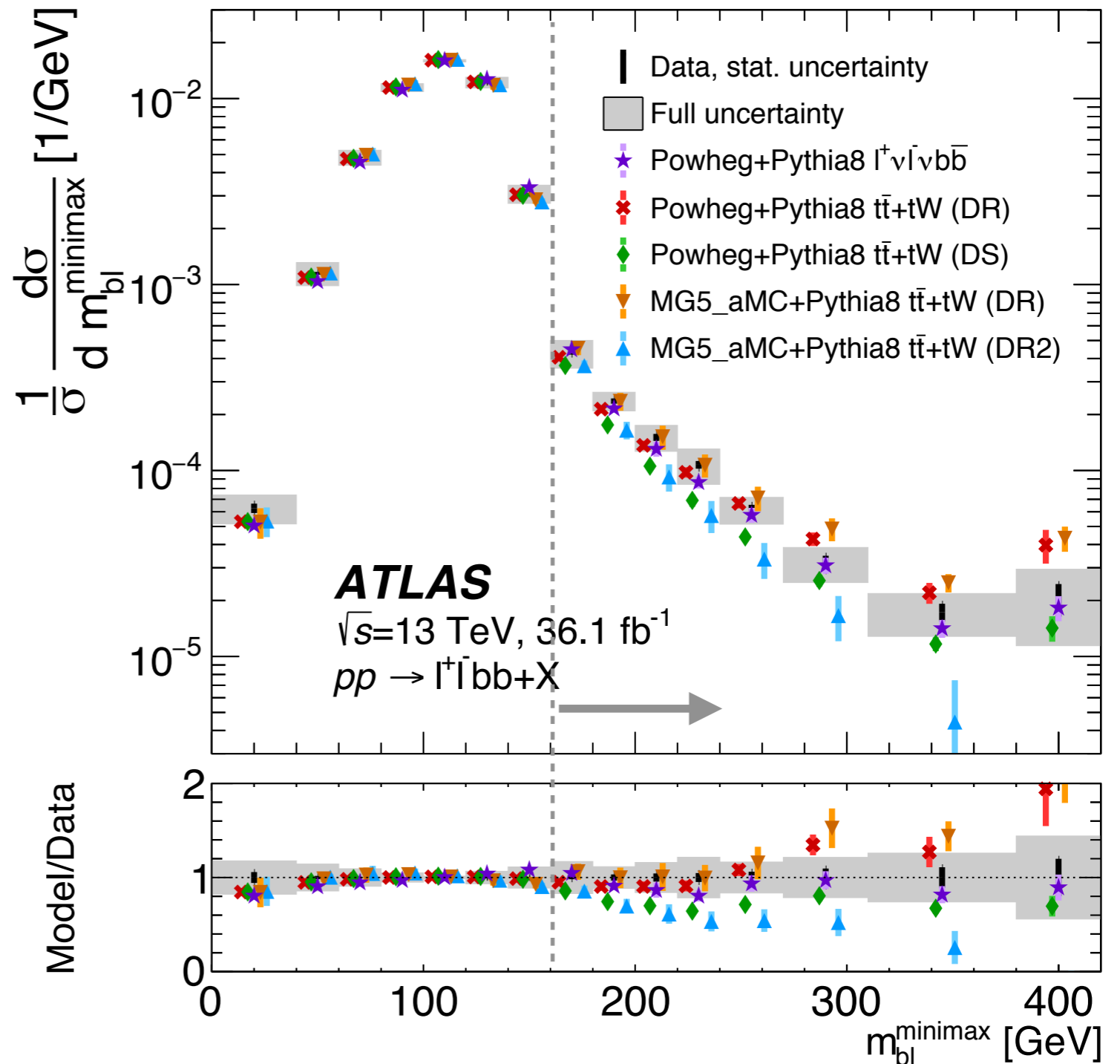
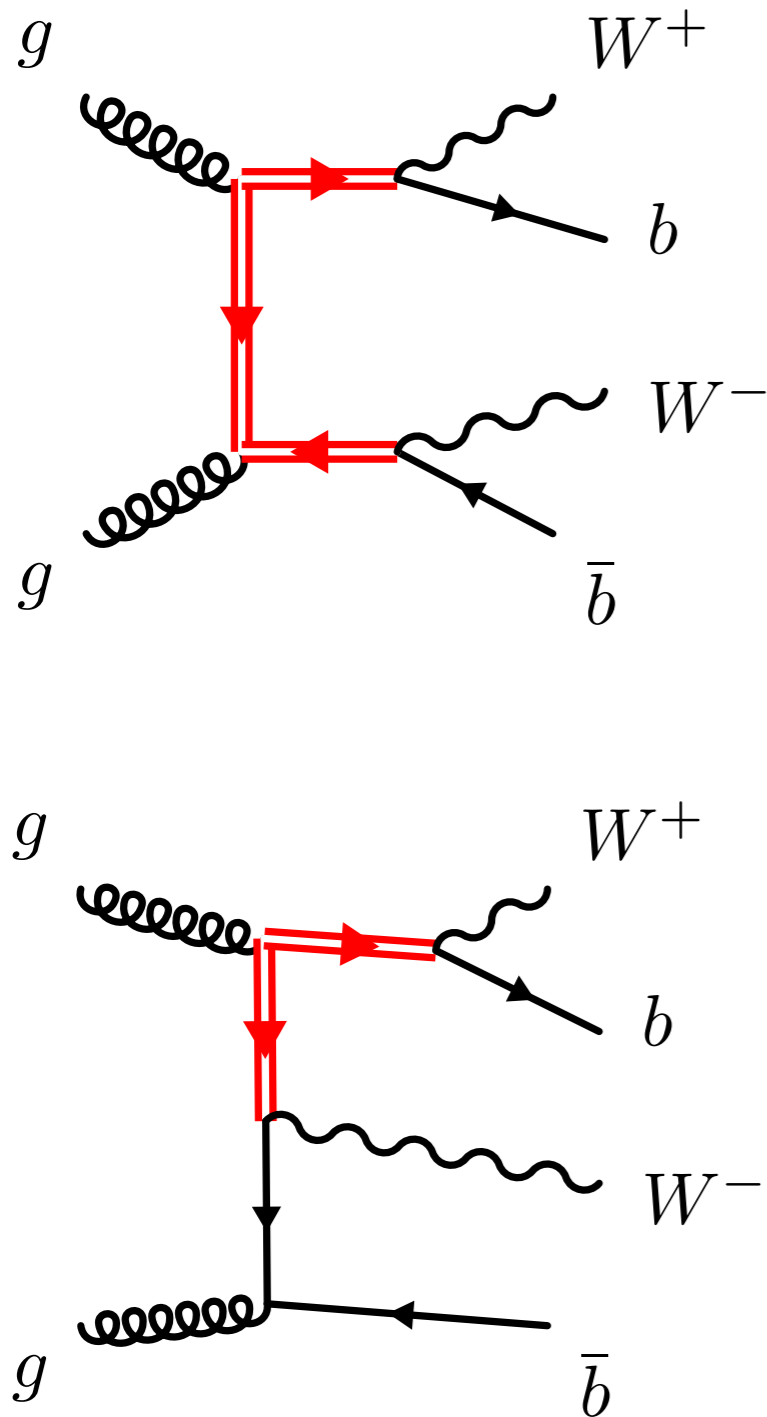
Many searches where top-pairs are a background are finding that  $Wtb$  is a significant residual background.

However, beyond leading order, single top and top pair production interfere!

This makes the modeling of this process complicated!



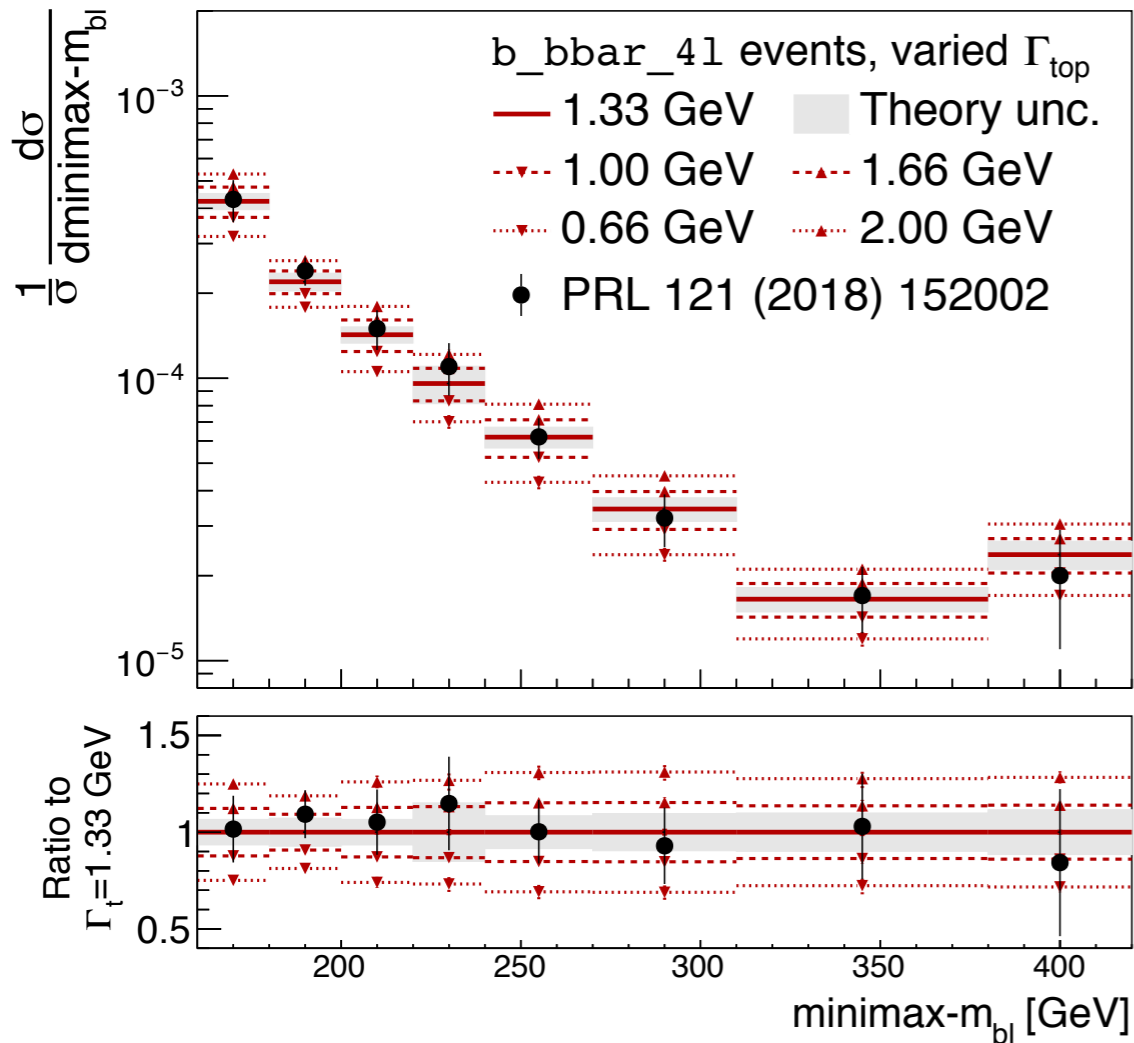
# Interference with pair production



Interference-sensitive region

# Top quark width

35

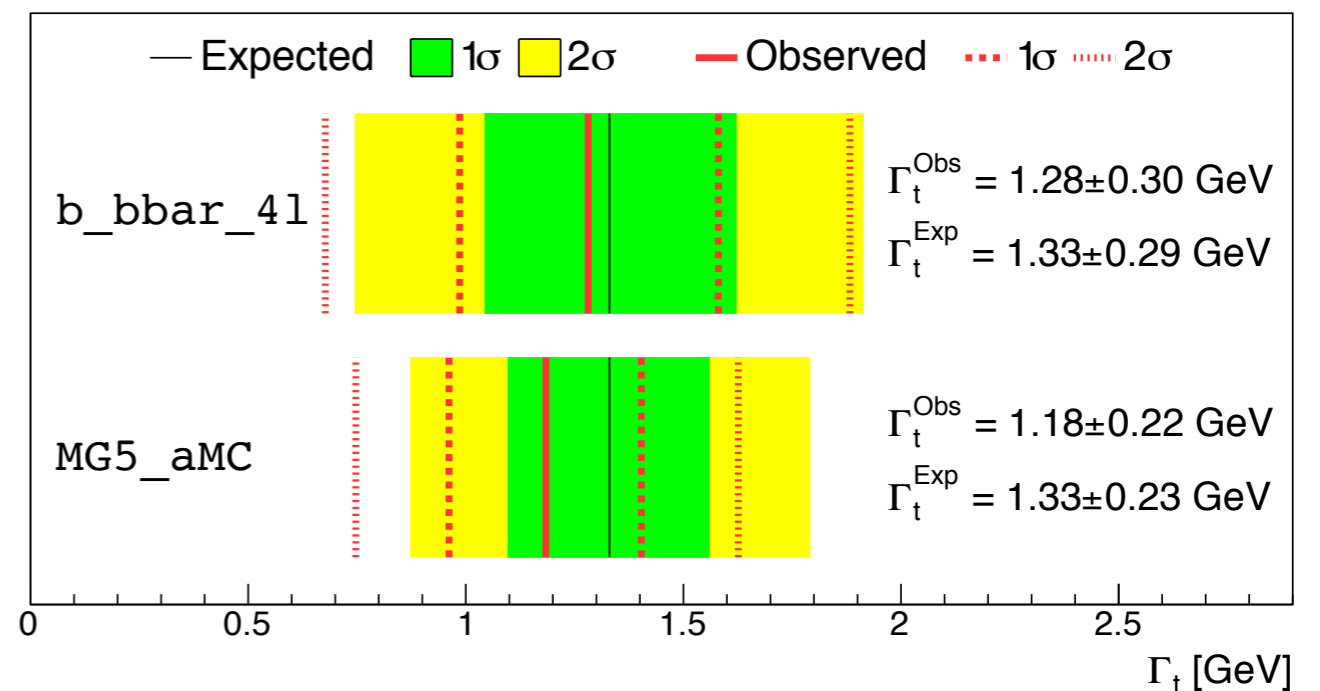


The top quark width is precisely known, but experimental precision is only  $\sim 20\%$

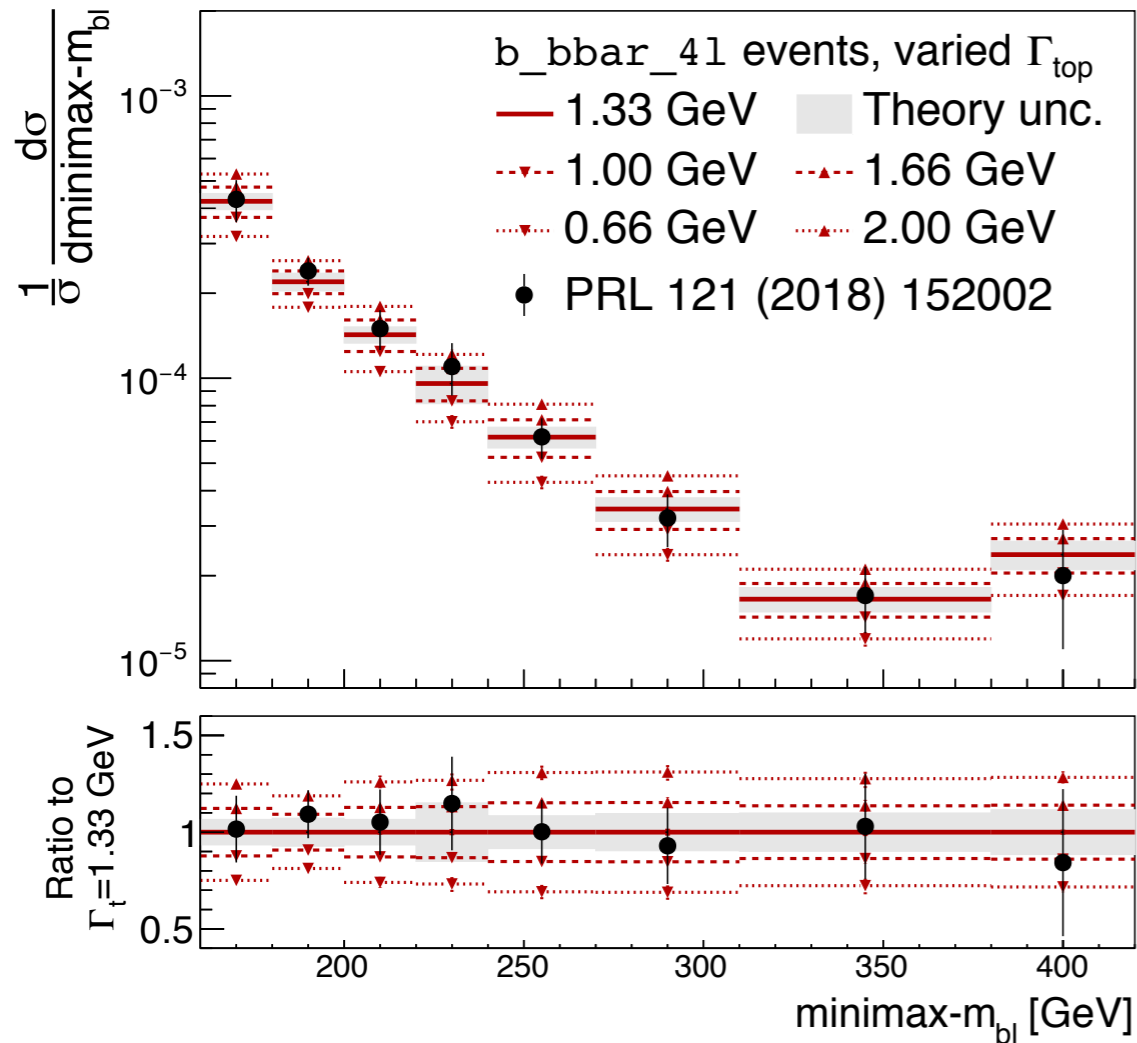
In fact, this region is rather sensitive to the top quark width!

Stronger sensitivity than the peak region (traditional method)

(why?)



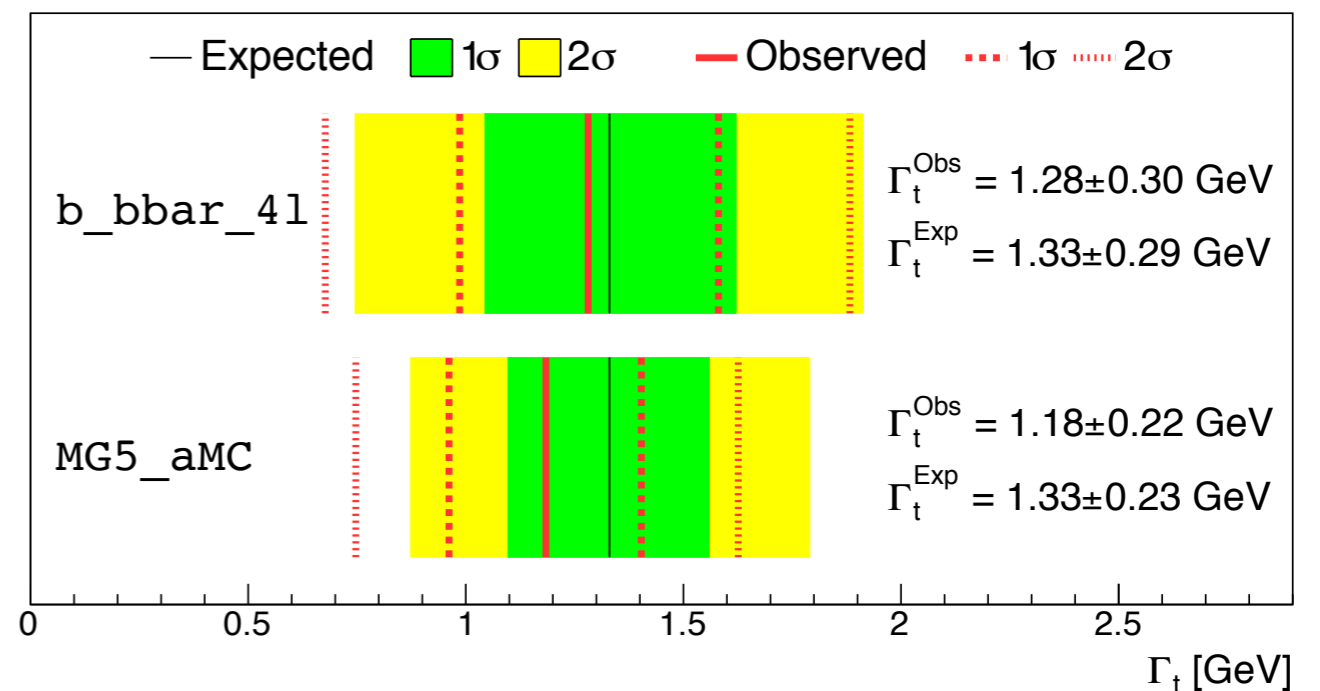
# Top quark width



$$BR(t \rightarrow \text{BSM}) < \frac{\Gamma_{\text{ext}}^{+95\%} - \Gamma_{t \rightarrow bW}^{\text{SM}}}{\Gamma_{\text{ext}}^{+95\%}}$$

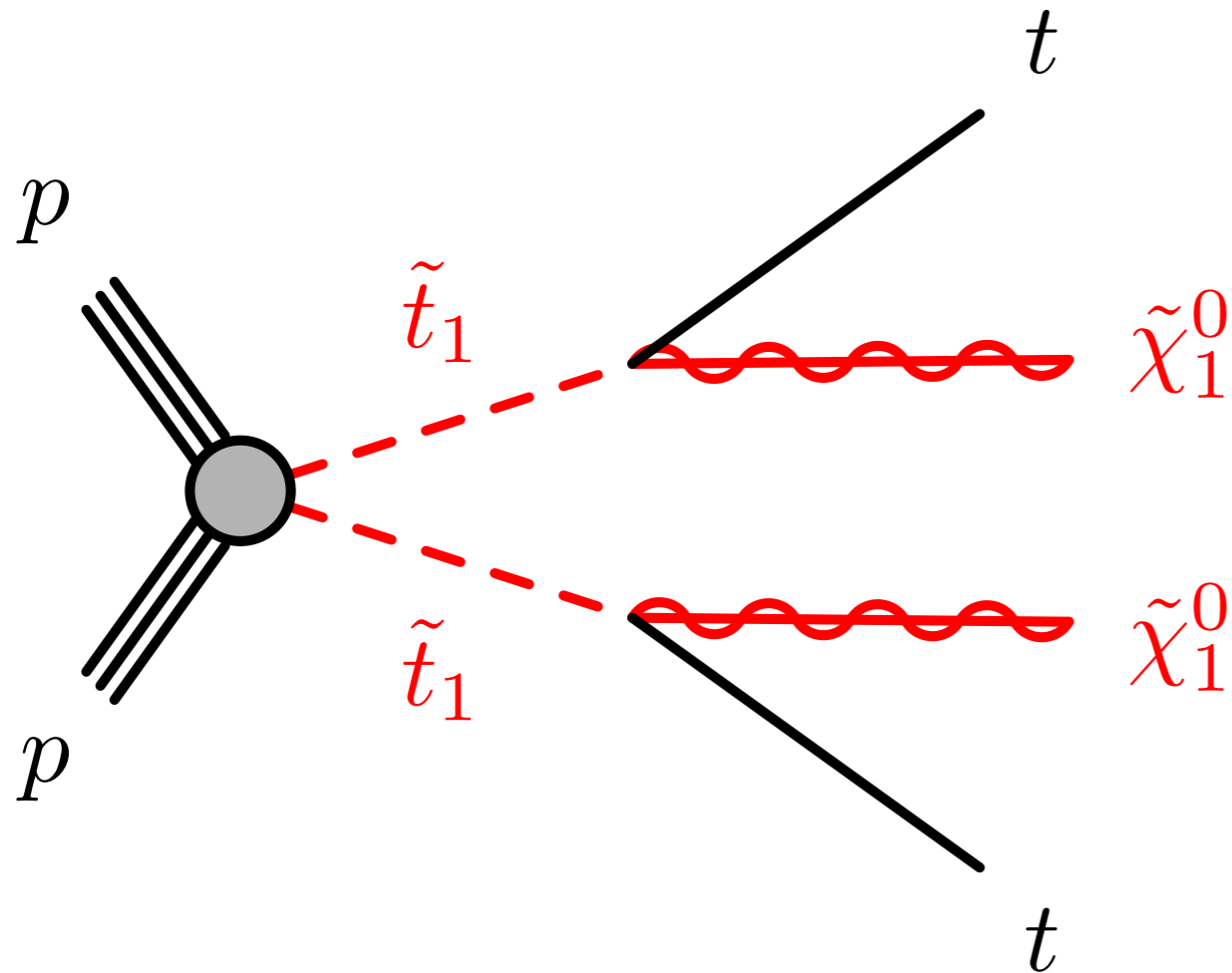
One can use this to set model-independent limits on the BSM branching ratio.

These are quite weak compared to the FCNC results shown earlier!



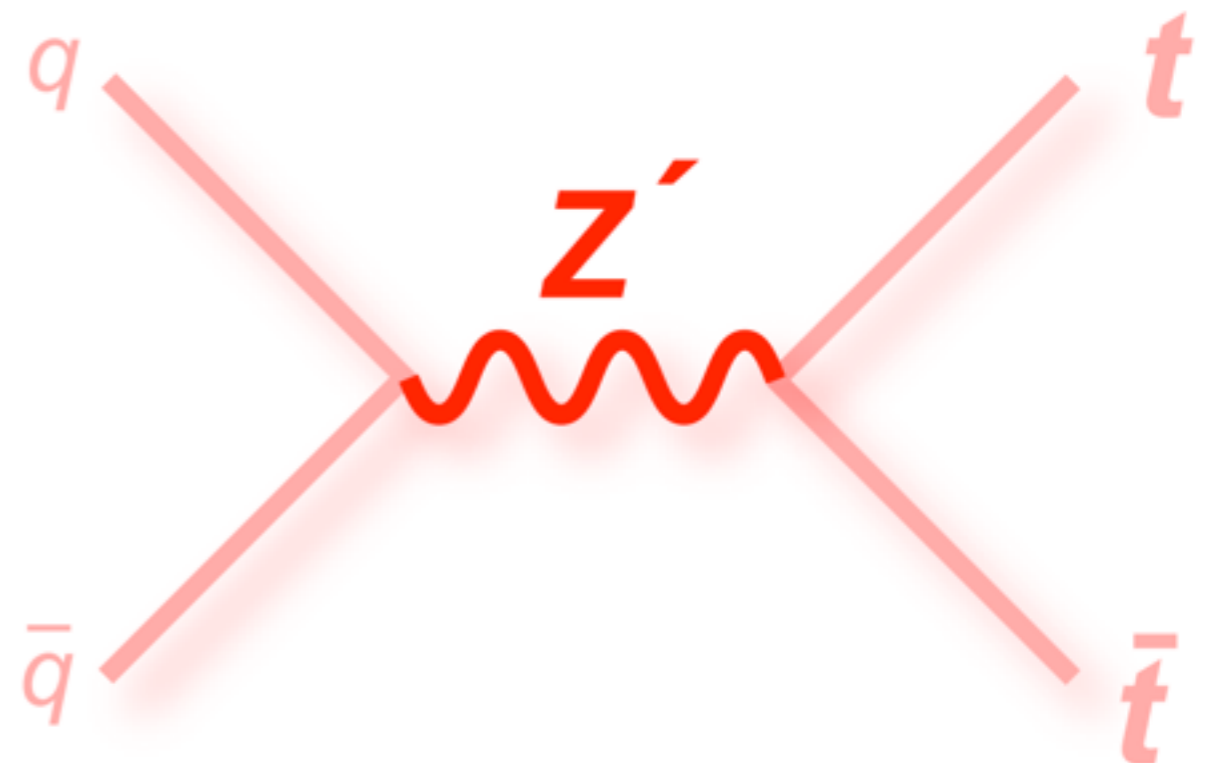
# Last topic: BSM with top quarks

37

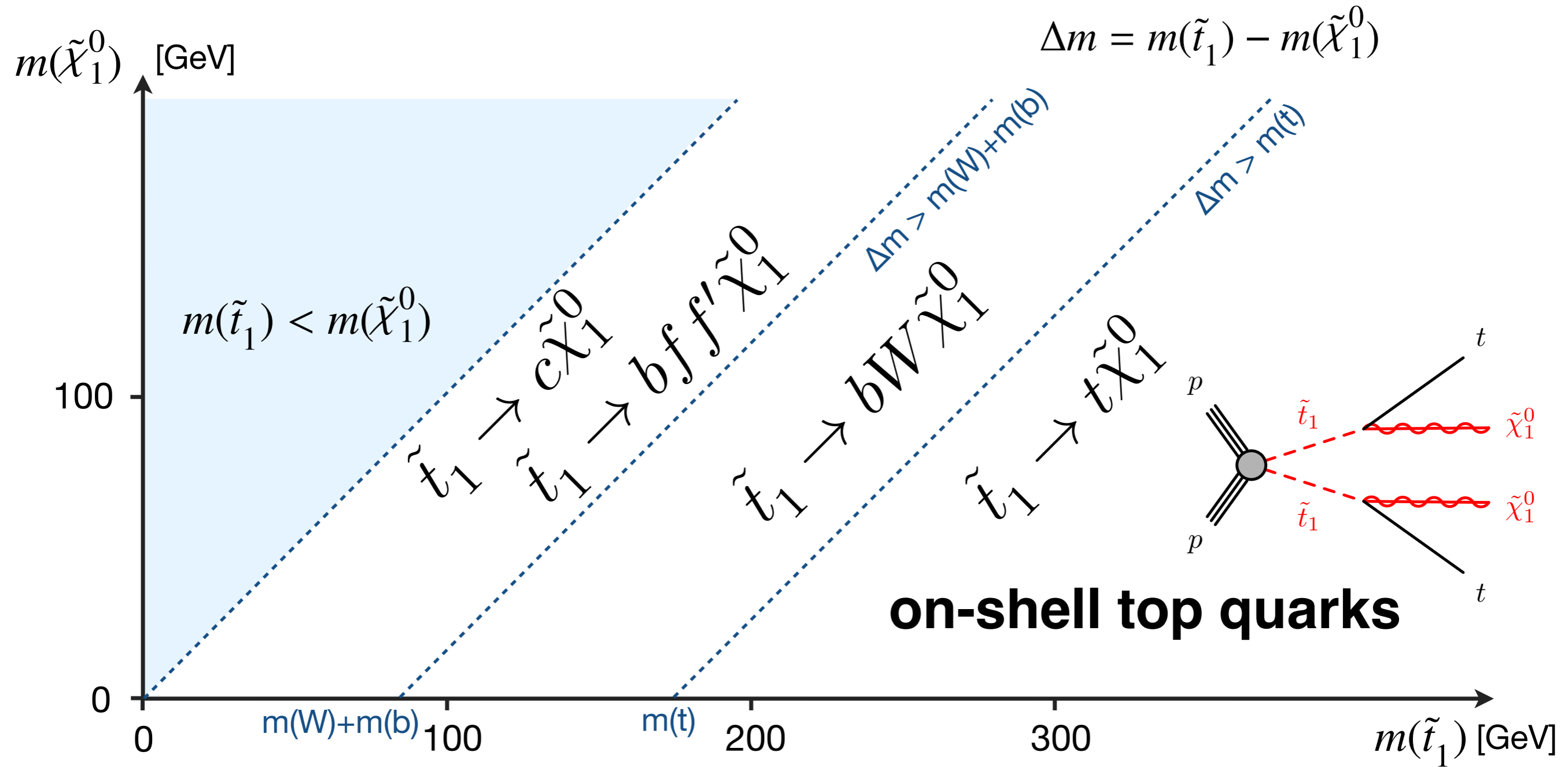


Top quarks feature prominently in many searches, either as the target or as the main background.

I'll briefly discuss the interplay between the stop search and SM top measurements.



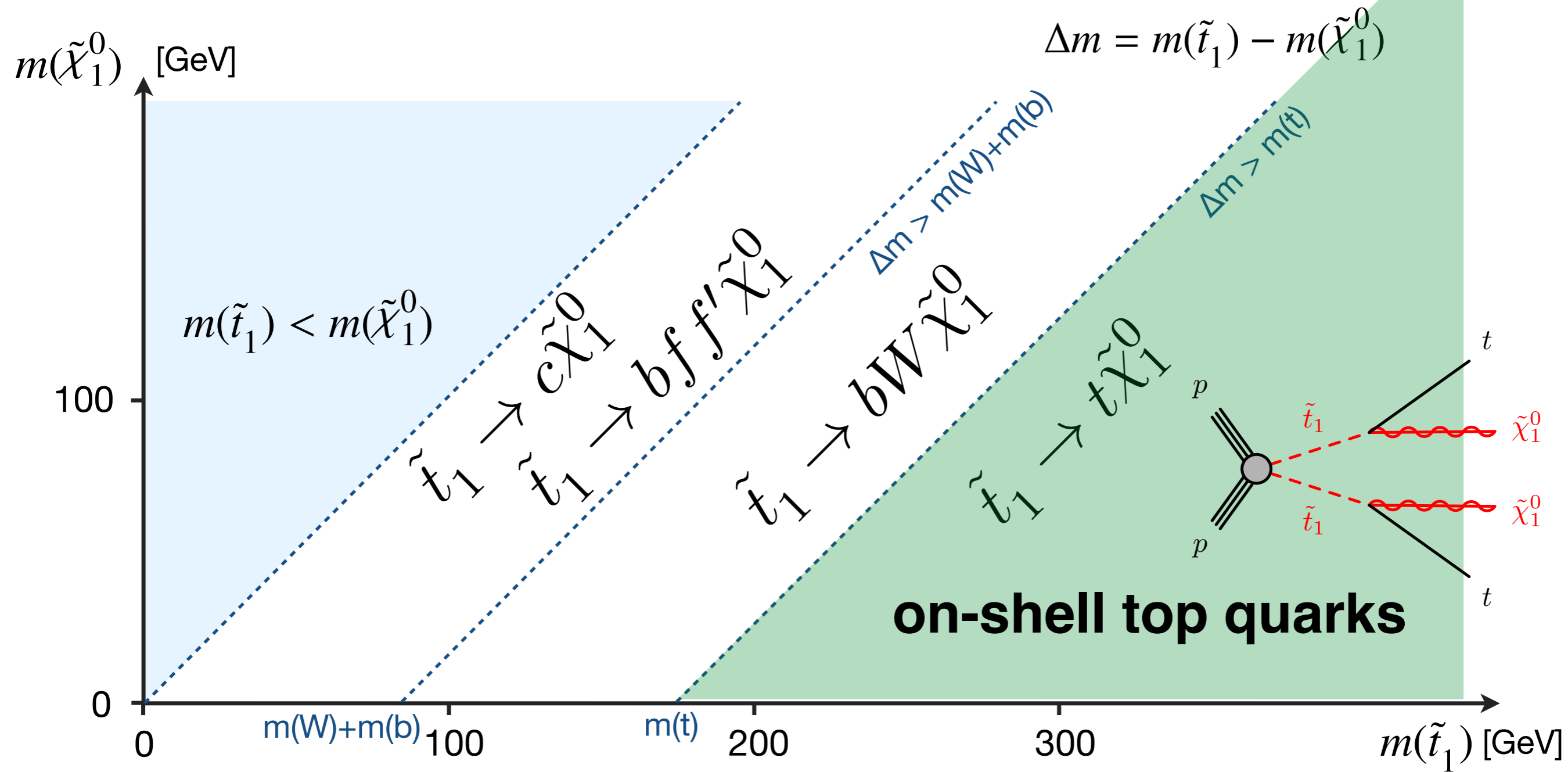
# BSM with top quarks



I won't motivate why SUSY/stops - see **Beyond SM** tomorrow

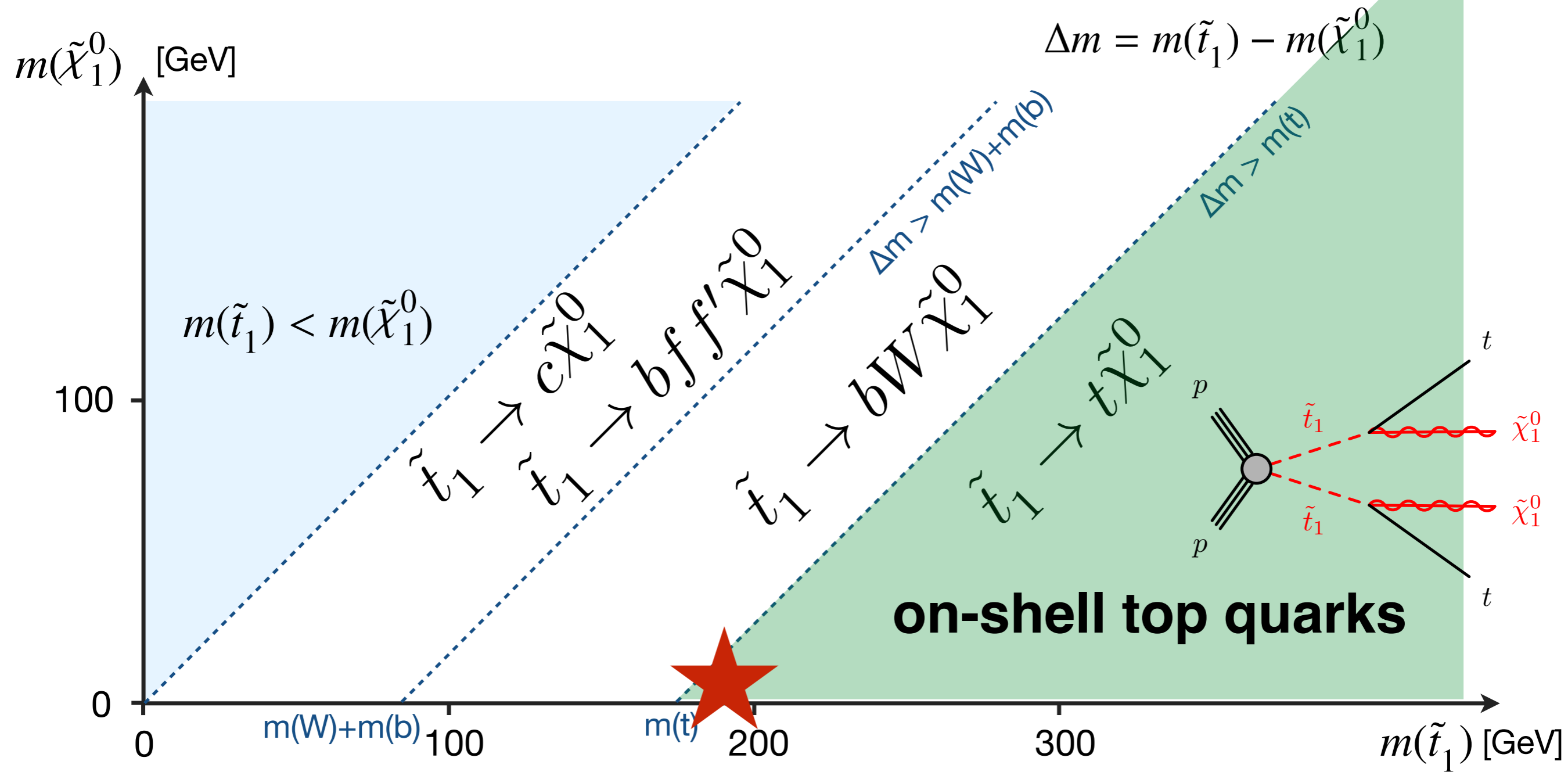
# BSM with top quarks

39



# BSM with top quarks

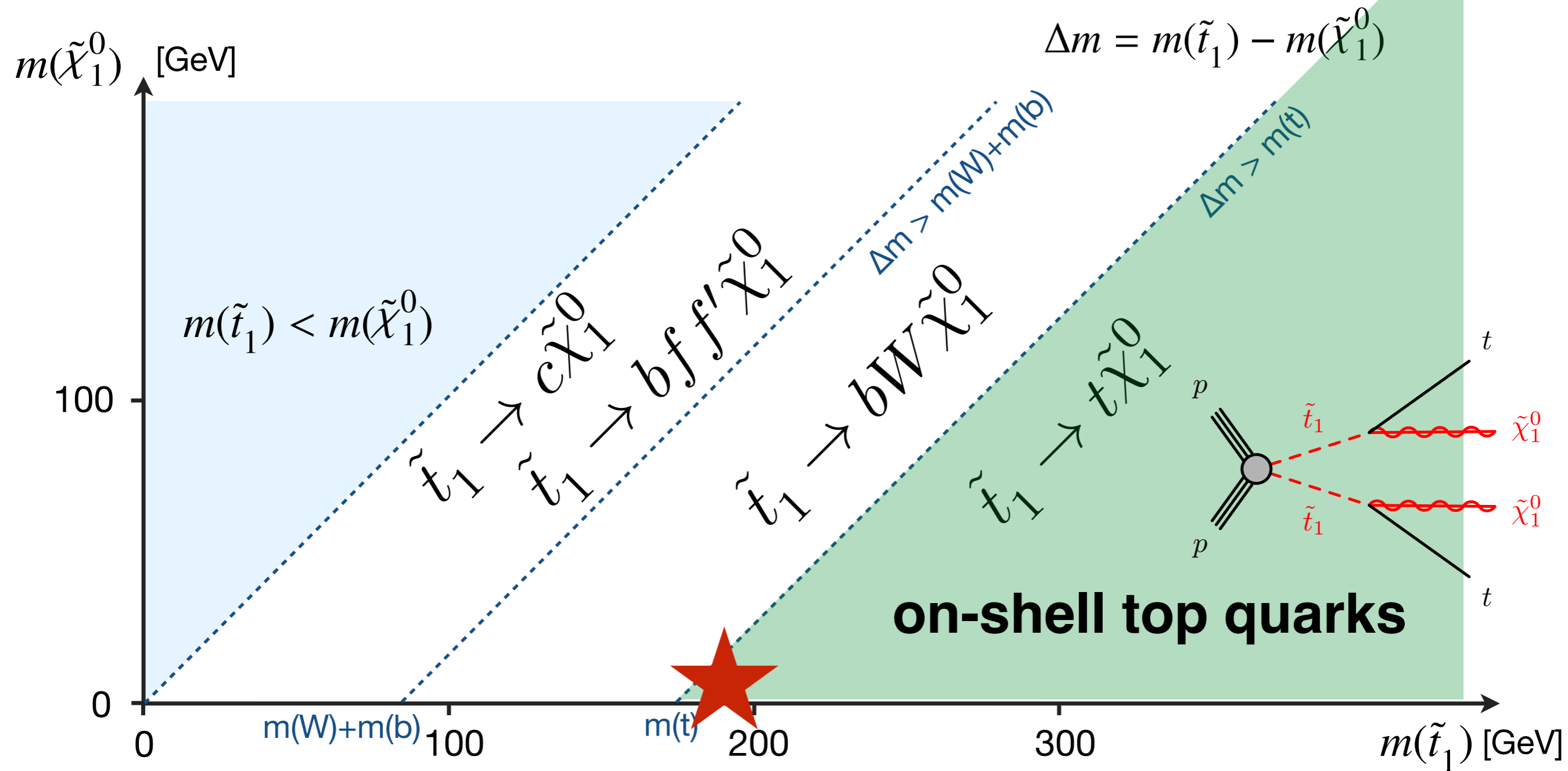
40





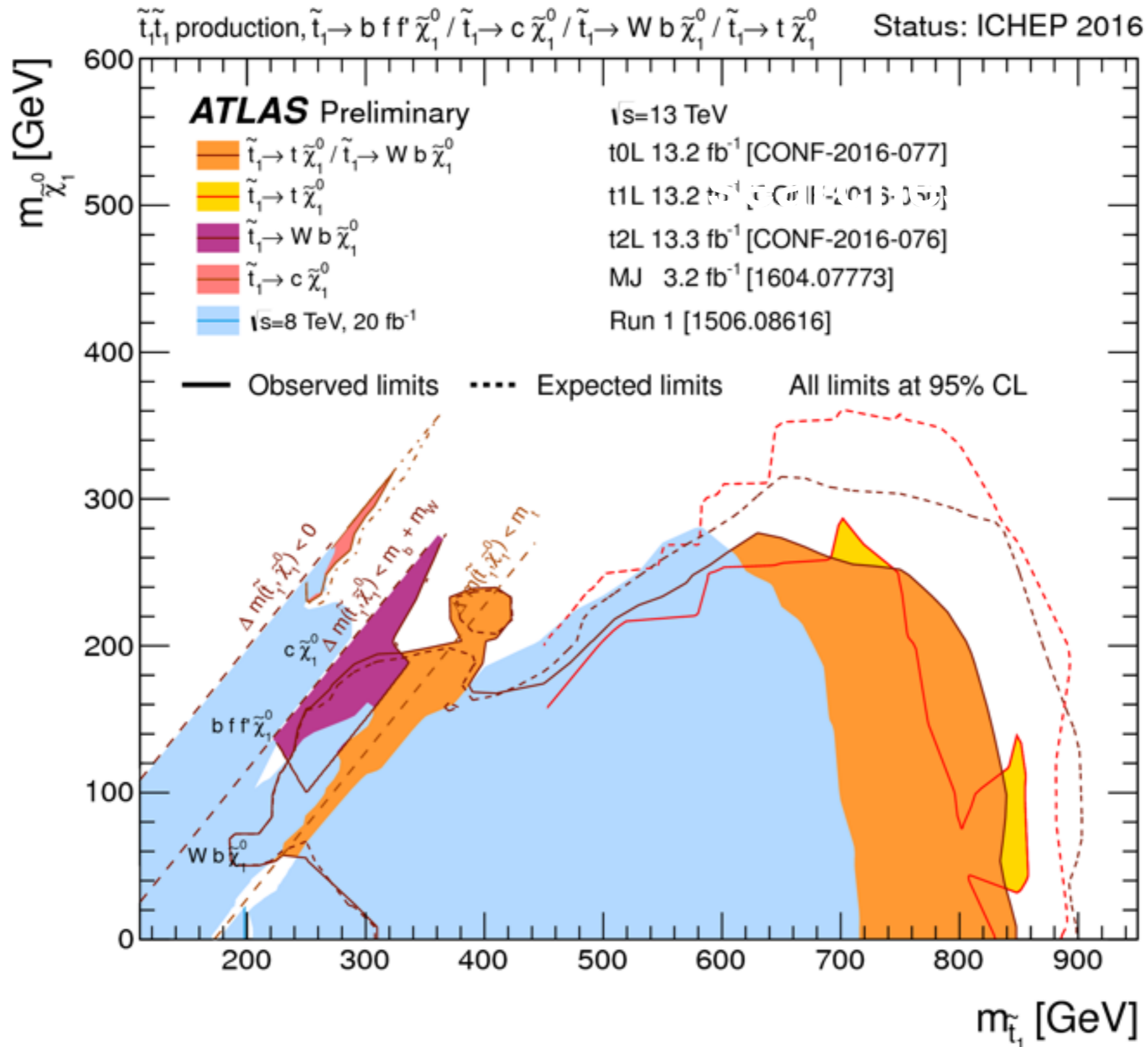
# BSM with top quarks

41

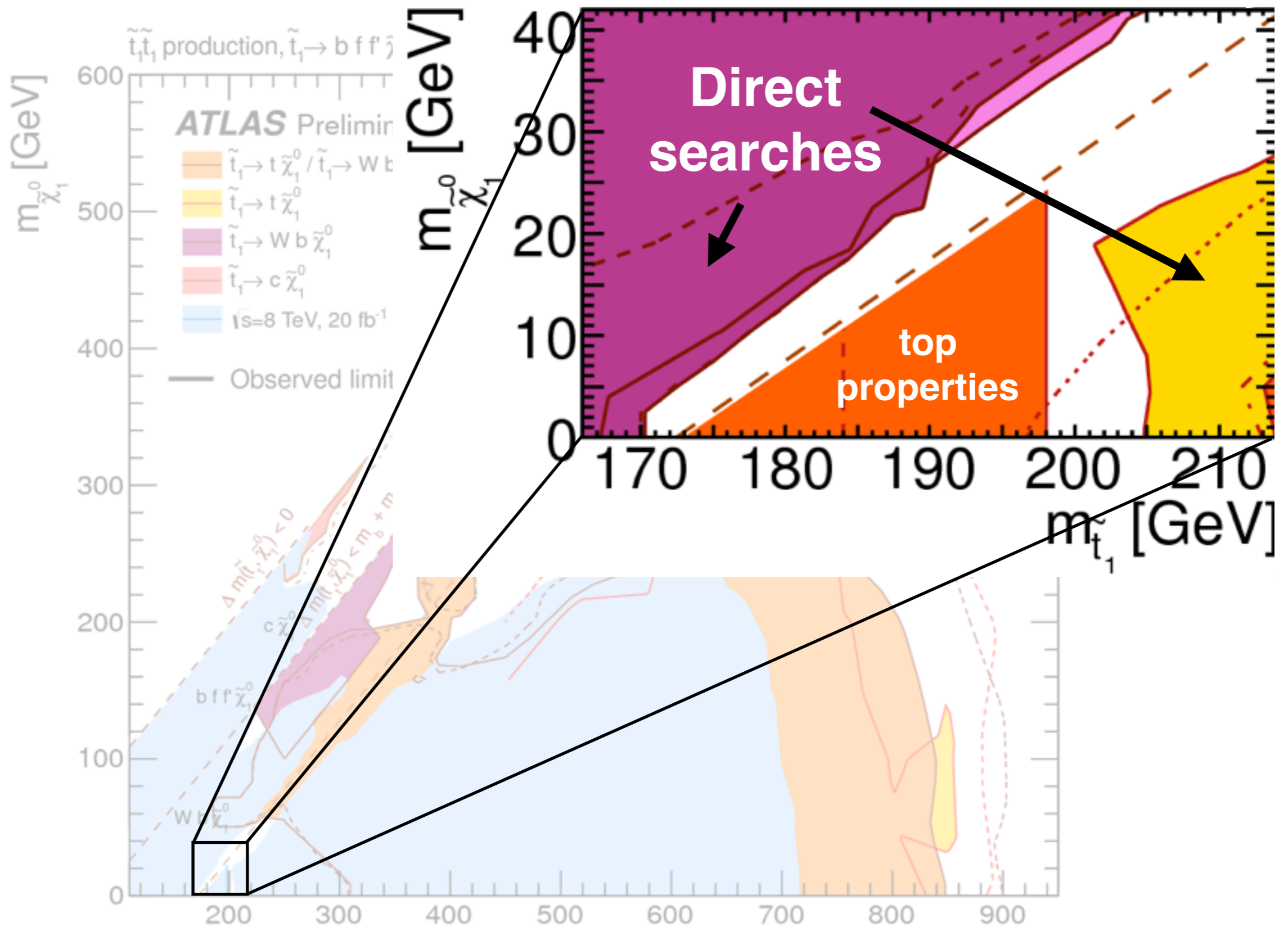


In this “stealth” region, the neutralinos have little momentum.

# BSM with top quarks



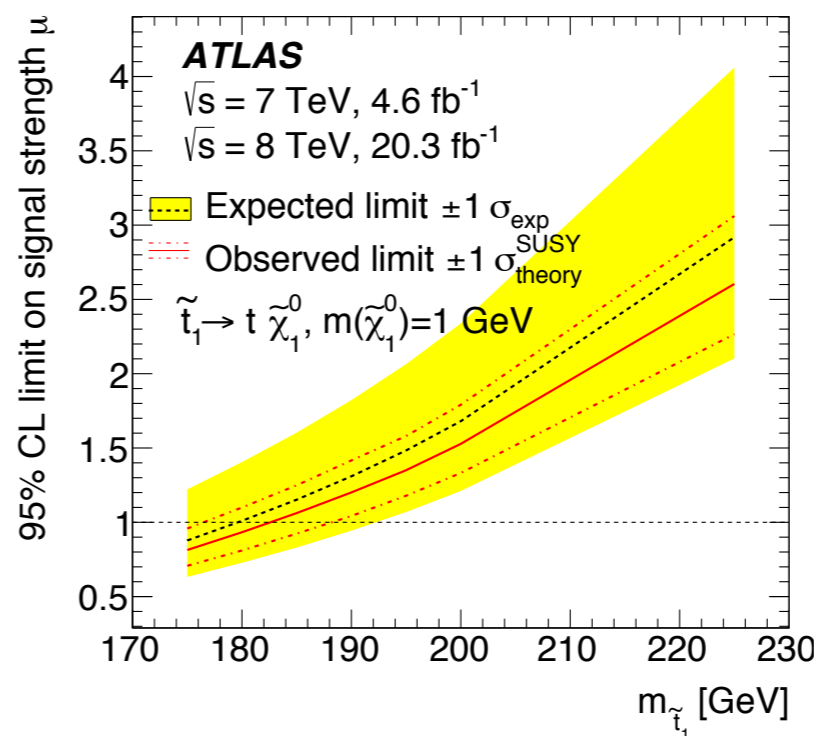
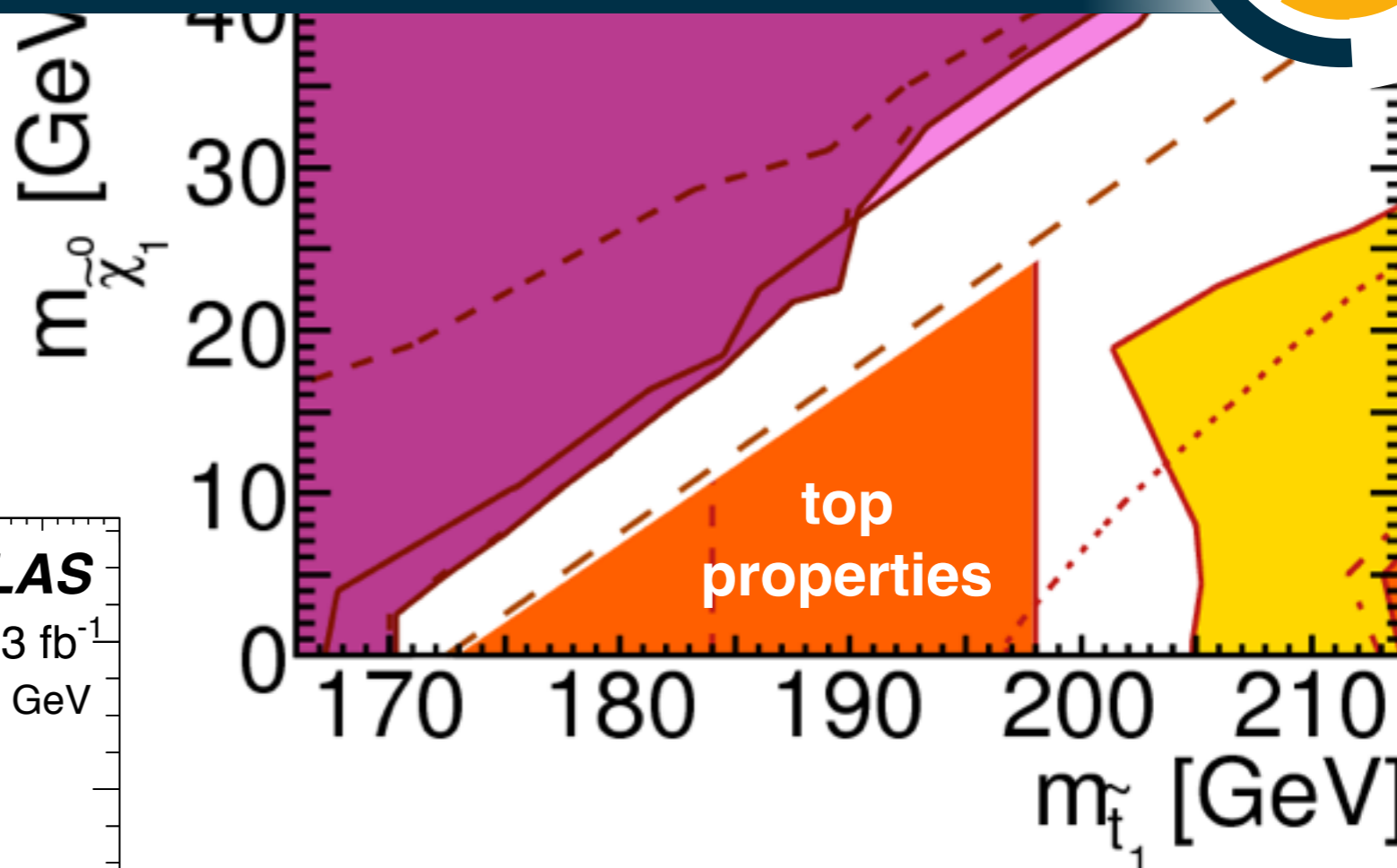
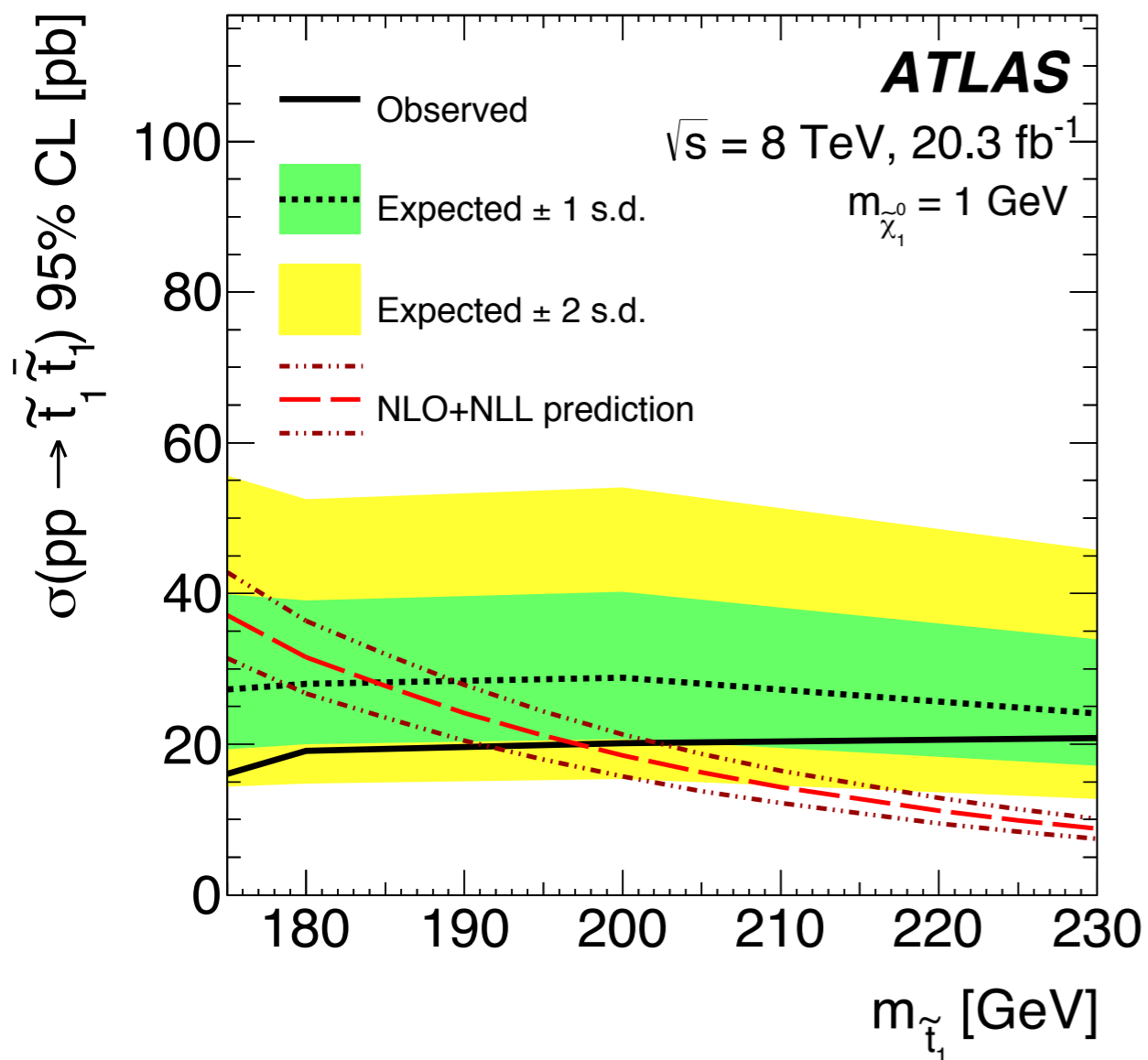
# Stop $\sim$ top



# Stop $\sim$ top

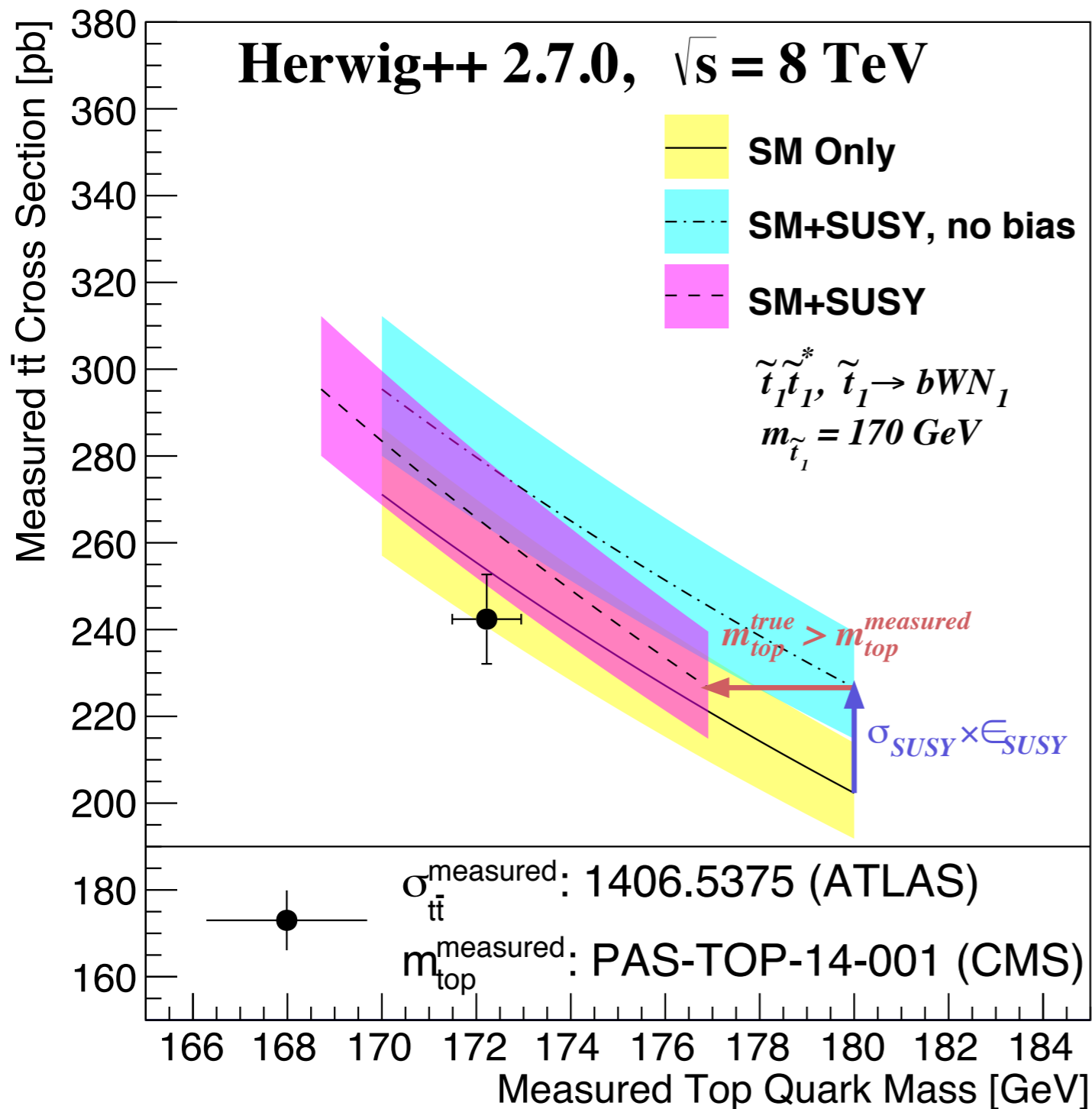
44

spin correlations  
measurement



Top cross-section  
measurement

# Warning: sneaky light stop



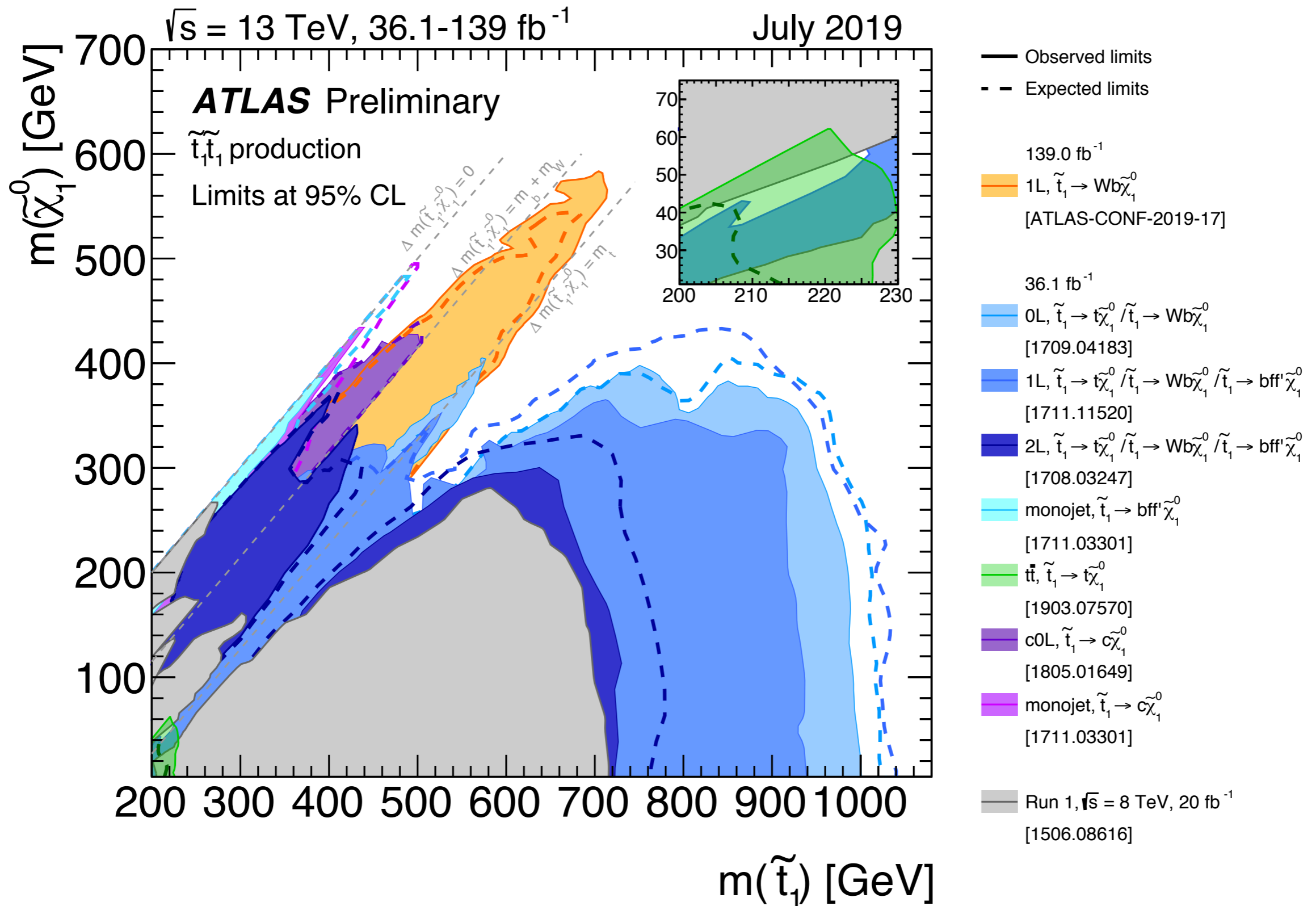
A light stop could introduce a bias in the top mass measurement.

If **we measure the mass too low\***, we will predict too high a  $t\bar{t}$  cross-section.

The **added cross-section from the stop** can then be hidden!

\*it is always low because  $m(b\ell\nu)$  or  $m(bjj)$  are always less for the stop than the top.

# Status today





Questions?

