

PPC 2021 Parallel talk 05/19/2021

The History and Fate of Dark Neutrons Nirmal Raj

TRIUMF

based on 2012.09865 (accepted at PRD)

David McKeen & Maxim Pospelov

YDROGEN HELIUM, CARBON

OUTER CRUST IONS, ELECTRONS

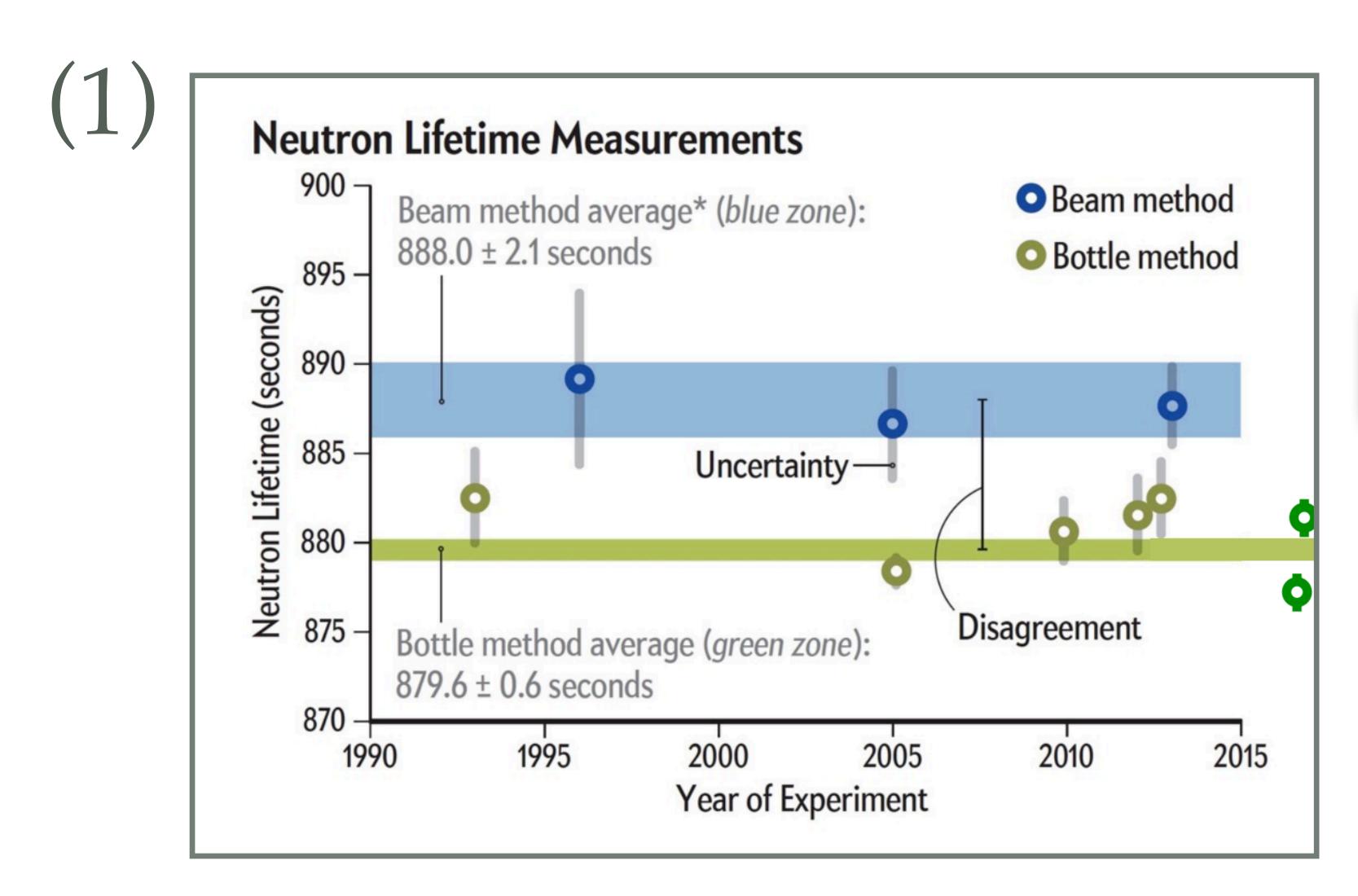
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OUTER CORE SUPERCONDUCTING PROTONS

INNER CORE



Why dark baryons? [new GeV-mass states carrying B]



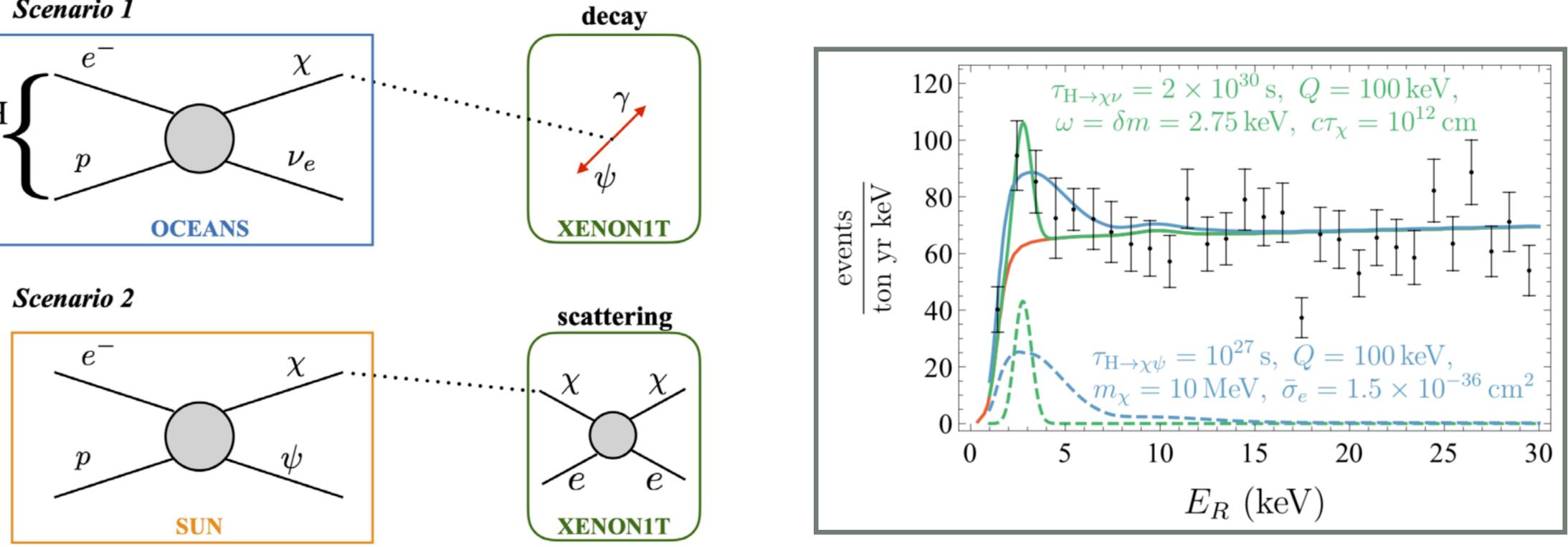
discrepancy:

$$\frac{\Delta \tau_n}{\tau_n} \approx 1\%$$

explain *n* lifetime puzzle with 1% branching to $n \rightarrow \chi$ + anything Fornal, Grinstein (2018)

Why dark baryons? [new GeV-mass states carrying B] could explain recent XENON1T excess

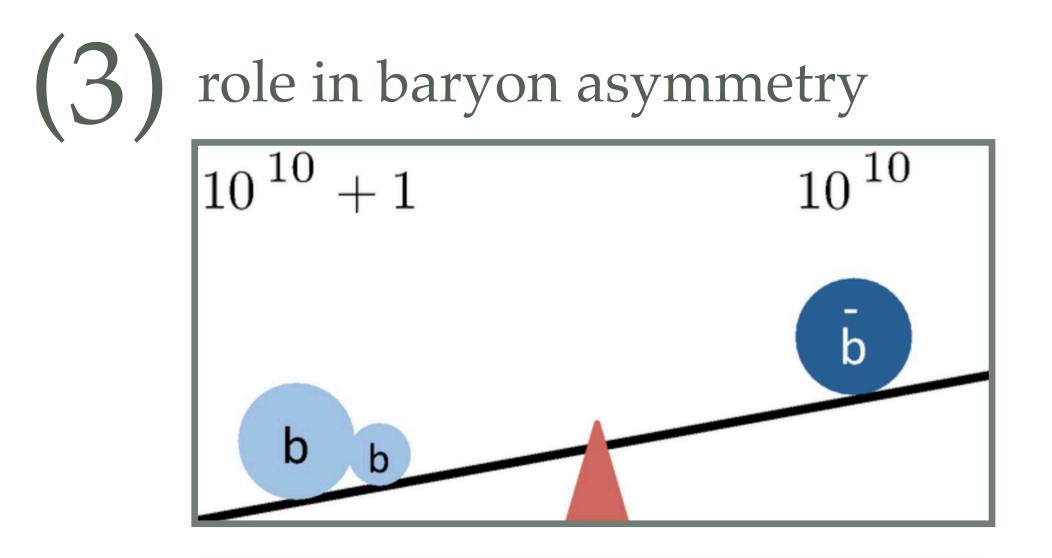
Scenario 1 e^{-} •••••• Η u_e p**OCEANS**



2006.15140 [PRL 125, 231803 (2020)] : McKeen, Pospelov, Raj



Why dark baryons? [new GeV-mass states carrying B]



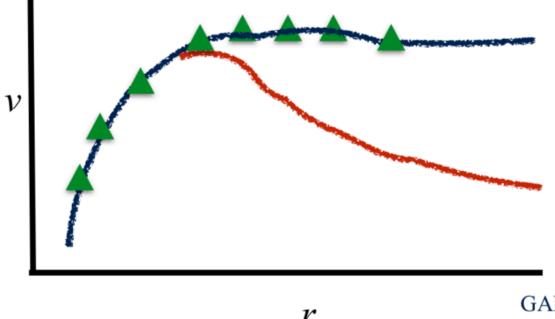
D. McKeen and A. E. Nelson, Phys. Rev. D **94**, 076002 (2016), arXiv:1512.05359 [hep-ph].

K. Aitken, D. McKeen, T. Neder, and A. E. Nelson, Phys. Rev. D **96**, 075009 (2017), arXiv:1708.01259 [hepph].

K. Babu, P. Bhupal Dev, E. C. Fortes, and R. Mohapatra, Phys. Rev. D 87, 115019 (2013), arXiv:1303.6918 [hep-ph]; R. Allahverdi, P. S. B. Dev, and B. Dutta, Phys. Lett. B 779, 262 (2018), arXiv:1712.02713 [hep-ph]; G. Elor, M. Escudero, and A. Nelson, Phys. Rev. D 99, 035031 (2019), arXiv:1810.00880 [hep-ph]; A. E. Nelson and H. Xiao, Phys. Rev. D 100, 075002 (2019), arXiv:1901.08141 [hep-ph]; G. Alonso-Álvarez, G. Elor, A. E. Nelson, and H. Xiao, JHEP 03, 046 (2020), arXiv:1907.10612 [hep-ph].

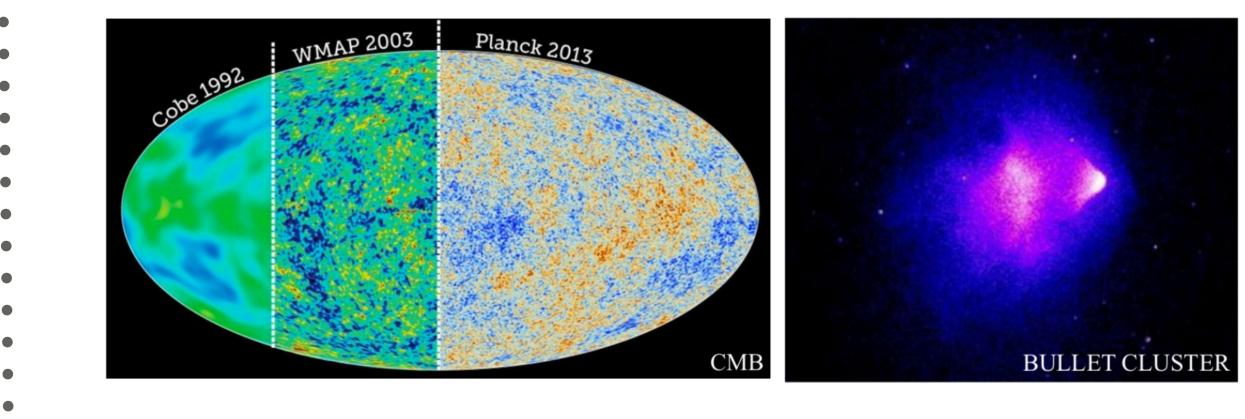
T. Bringmann, J. M. Cline, and J. M. Cornell, Phys. Rev. D **99**, 035024 (2019), arXiv:1810.08215 [hep-ph].

(4) could constitute the <u>dark matter</u> of the universe

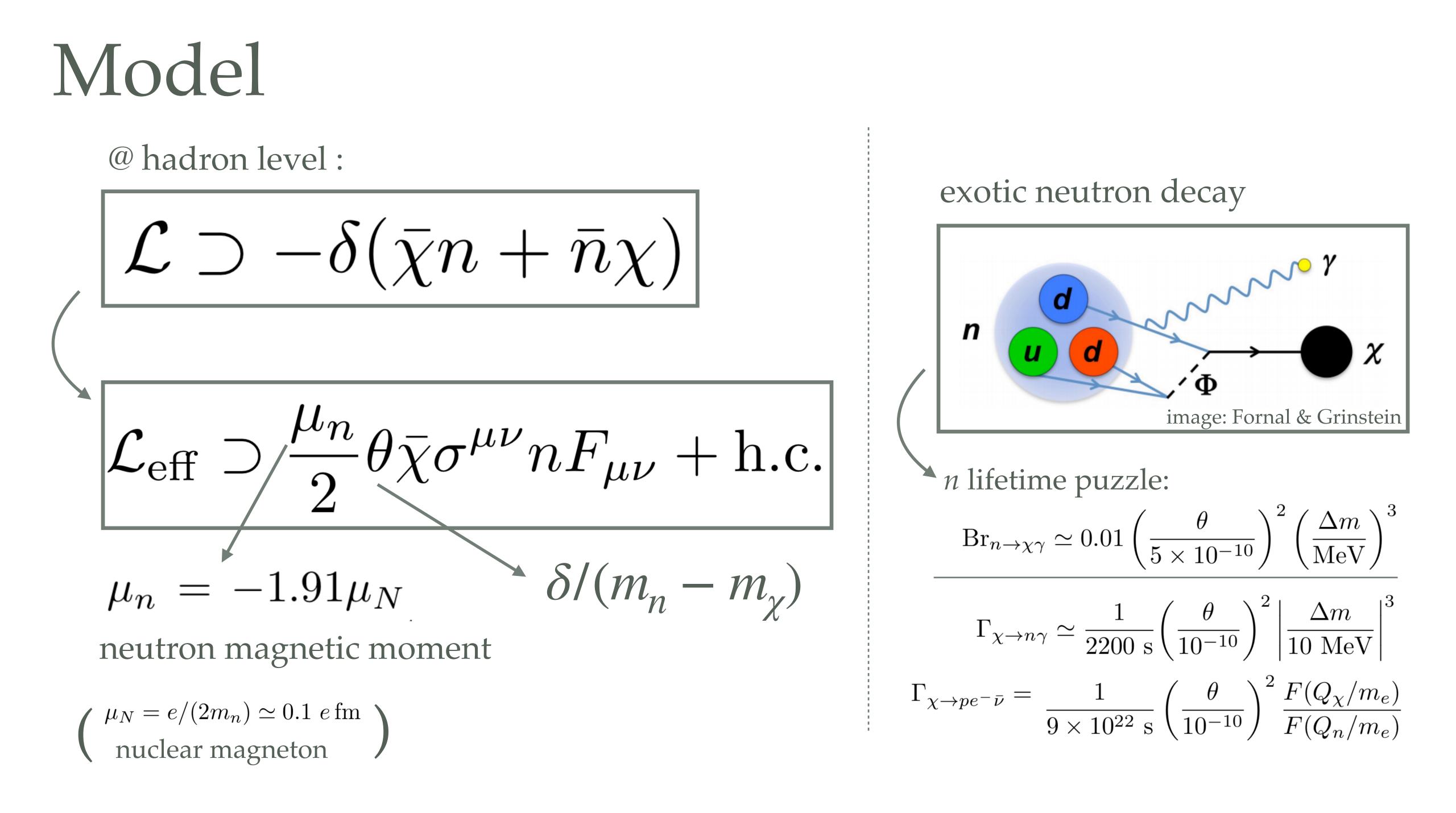




GALACTIC ROTATION

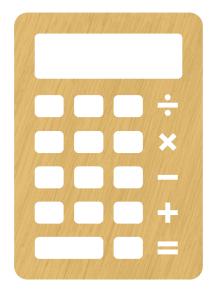


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Prehistoric census

(i)
$$n_{\chi}^{0} = 5.4(n_{p}^{0} + n_{n}^{0})$$
 (χ is the dark matter if $\tau_{\chi} > t_{U}$)
(ii) $n_{\chi}^{0} = 0.01(n_{p}^{0} + n_{n}^{0})$ (perhaps never chem eqbm)

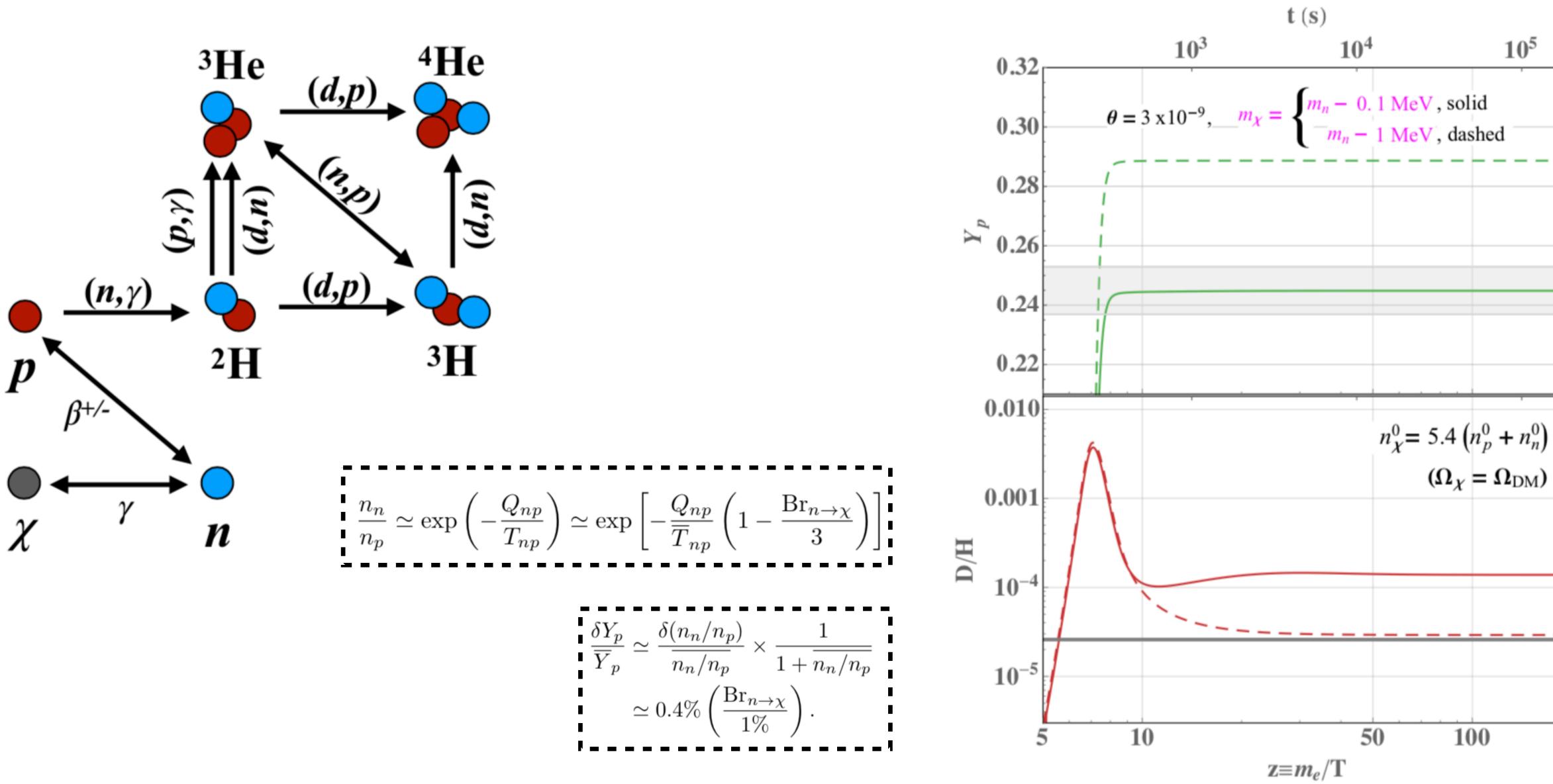


$$\frac{\mu_n}{2} \theta \bar{\chi} \sigma^{\mu\nu} n F_{\mu\nu} \longrightarrow \text{number-changing rate} \\ \Gamma_{\Delta\chi} \sim \theta^2 \mu_n^2 T^3 \gtrsim H \text{ for } T \gtrsim 100 \text{ MeV} \left(\frac{10^{-9}}{\theta}\right)^2 \\ \text{above QCD transition} => \text{quark level description required} \\ -\delta(\bar{\chi}n + \bar{n}\chi) \longleftarrow \bar{\chi} q q q / \Lambda^2 => \Gamma_{\Delta\chi} \sim T^5 / \Lambda^4 \\ \text{chemical equilbrium keepable down to } T \sim \text{GeV}-\text{PeV} \end{cases}$$

for $\theta \sim 10^{-20} - 10^{-10}$ and $\Delta m \sim 1 - 100$ MeV ----- $n_{\chi} \sim n_p = n_n$ reasonable since universe was probably that hot



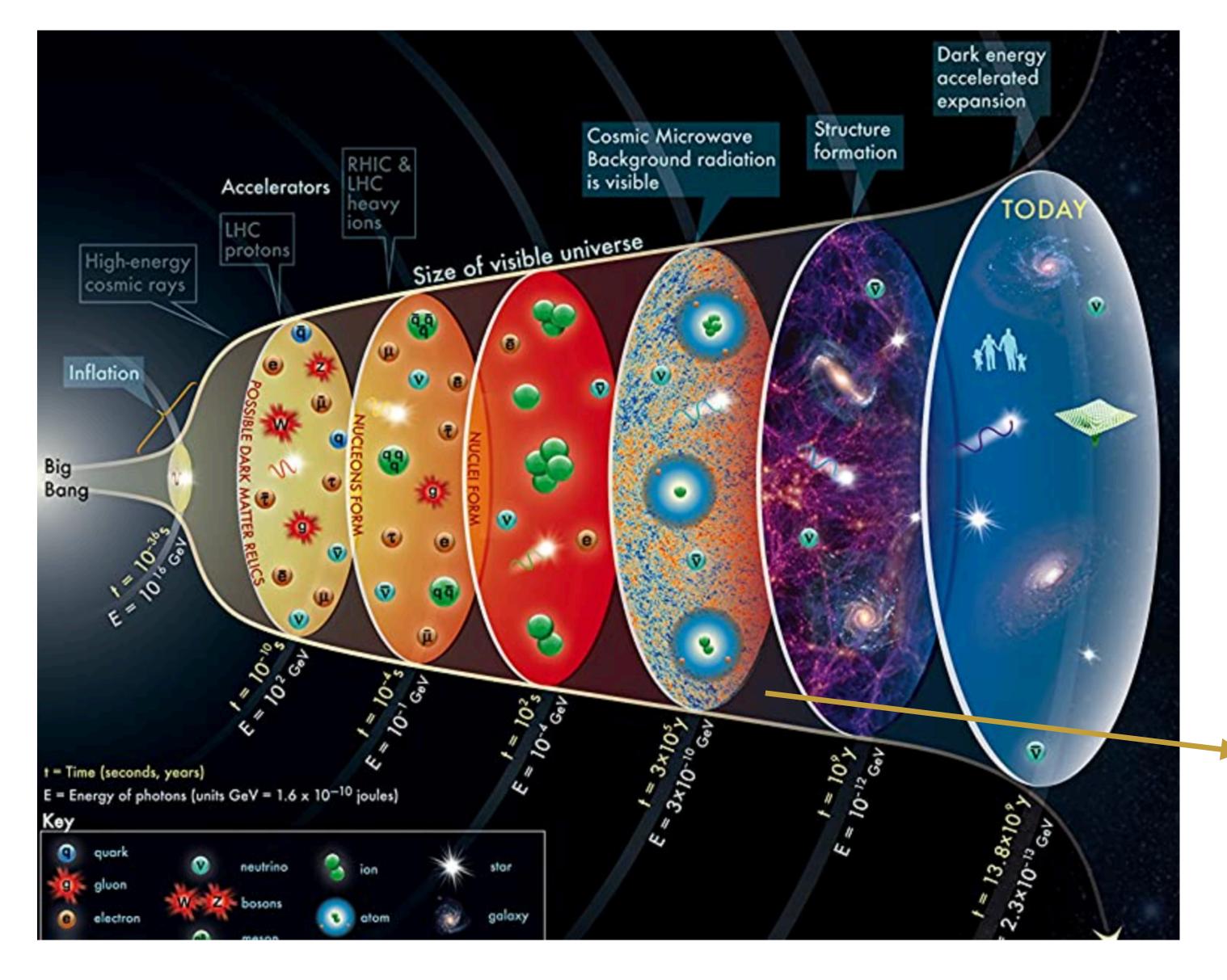
Probes: [1] primordial nucleosynthesis



$$\frac{\delta Y_p}{\overline{Y}_p} \simeq \frac{\delta(n_n/n_p)}{\overline{n_n/n_p}}$$
$$\simeq 0.4\% \left(\frac{1}{\overline{Y}_p}\right)$$



Probes: [2] relic radiation

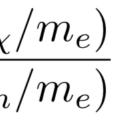


When kinematically open:

$$\Gamma_{\chi \to pe^-\bar{\nu}} = \frac{1}{9 \times 10^{22} \text{ s}} \left(\frac{\theta}{10^{-10}}\right)^2 \frac{F(Q_{\chi})}{F(Q_{\pi})}$$

$$\Gamma_{\chi \to n\gamma} \simeq \frac{1}{2200 \text{ s}} \left(\frac{\theta}{10^{-10}}\right)^2 \left|\frac{\Delta m}{10 \text{ MeV}}\right|$$

e or γ could "rewrite" reionization history by dumping EM energy in Dark Ages (i.e. modify optical depth)





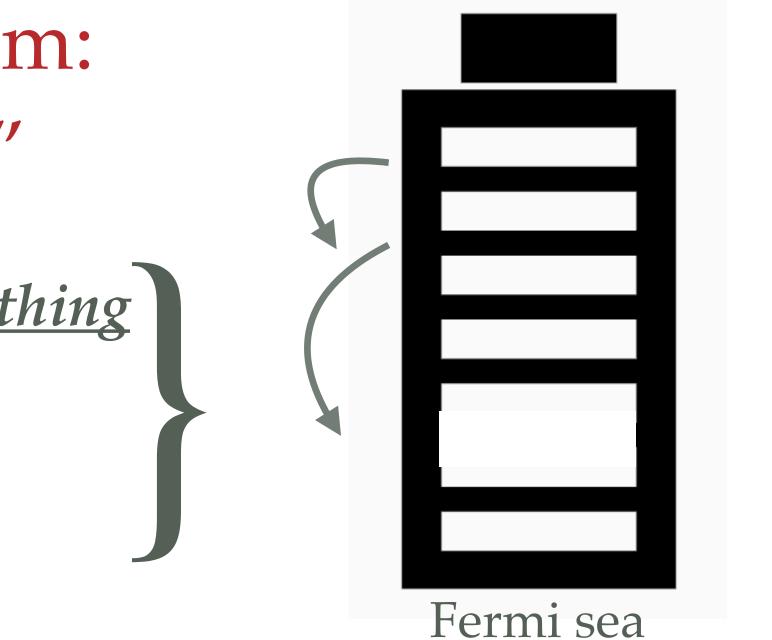


Probes: [3] neutron star temperatures

new heating mechanism: nucleon "Auger effect"

> $n \rightarrow \chi + anything$ $n n \rightarrow n \chi$ $p n \rightarrow p \chi$

10⁵⁷ neutrons 10⁵⁶ protons



neutron Fermi energy ~ 100 MeV

=> explosive liberation of energy!



Probes: [3] neutron star temperatures

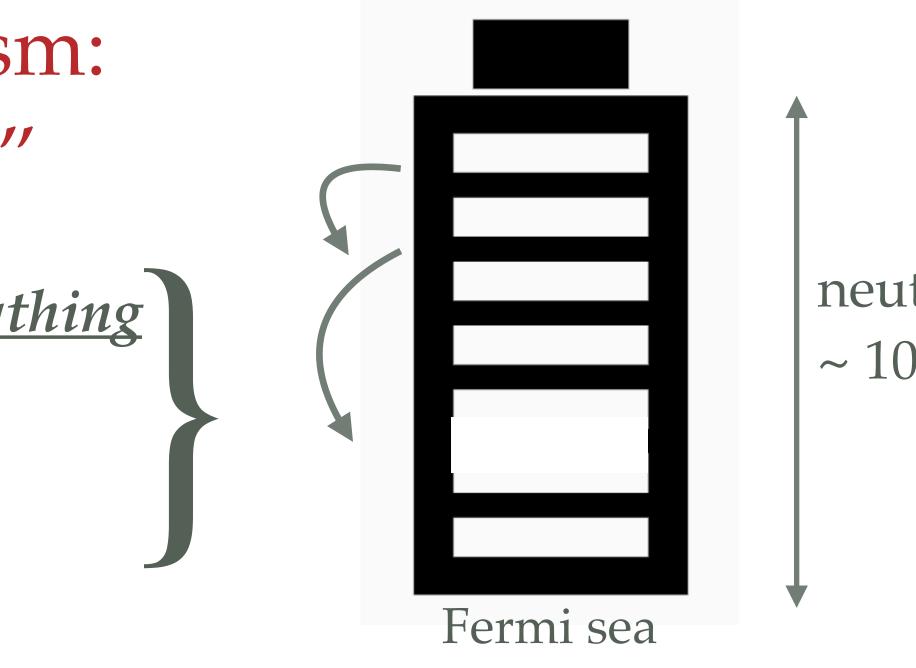
new heating mechanism: nucleon "Auger effect"

> $n \rightarrow \chi + \underline{anything}$ $n n \rightarrow n \chi$ $p n \rightarrow p \chi$

Sebastien Guillot^{1,2,3,8}, George G. Pavlov⁴, Cristobal Reyes³, Andreas Reisenegger³, Luis E. Rodriguez⁵, Blagoy Rangelov⁶, and Oleg Kargaltsev⁷

We report nondetections of the $\sim 3 \times 10^8$ yr old slow, isolated, rotation-powered pulsar PSR J2144–3933 in observations with the Hubble Space Telescope in one optical band (F475X) and two far-ultraviolet bands (F125LP and F140LP), yielding upper bounds $F_{F475X} < 22.7$ nJy, $F_{F125LP} < 5.9$ nJy, and $F_{F140LP} < 19.5$ nJy, at the pivot wavelengths 4940 Å, 1438 Å and 1528 Å, respectively. Assuming a blackbody spectrum, we deduce a conservative upper bound on the surface (unredshifted) temperature of the pulsar of T < 42,000 K. This makes

Suitable lab:



neutron Fermi energy ~ 100 MeV

Hubble Space Telescope Nondetection of PSR J2144–3933: The Coldest Known **Neutron Star***



Probes: [3] neutron star temperatures

new heating mechanism: nucleon "Auger effect"

> $n \rightarrow \chi + \underline{anything}$ $n n \rightarrow n \chi$ $p n \rightarrow p \chi$

Masha Baryakhtar,¹ Joseph Bramante,¹ Shirley Weishi Li,² Tim Linden,² and Nirmal Raj³ ¹Perimeter Institute for Theoretical Physics, Waterloo, Ontario N2L 2Y5, Canada ²CCAPP and Department of Physics, The Ohio State University, Columbus, Ohio 43210, USA ³Department of Physics, University of Notre Dame, Notre Dame, Indiana 46556, USA (Received 10 April 2017; revised manuscript received 20 July 2017; published 26 September 2017)

We identify a largely model-independent signature of dark matter (DM) interactions with nucleons and electrons. DM in the local galactic halo, gravitationally accelerated to over half the speed of light, scatters against and deposits kinetic energy into neutron stars, heating them to infrared blackbody temperatures. The resulting radiation could potentially be detected by the James Webb Space Telescope, the Thirty Meter Telescope, or the European Extremely Large Telescope. This mechanism also produces optical emission

Future lab:

Fermi sea Dark Kinetic Heating of Neutron Stars and an Infrared Window on WIMPs, SIMPs, and Pure Higgsinos

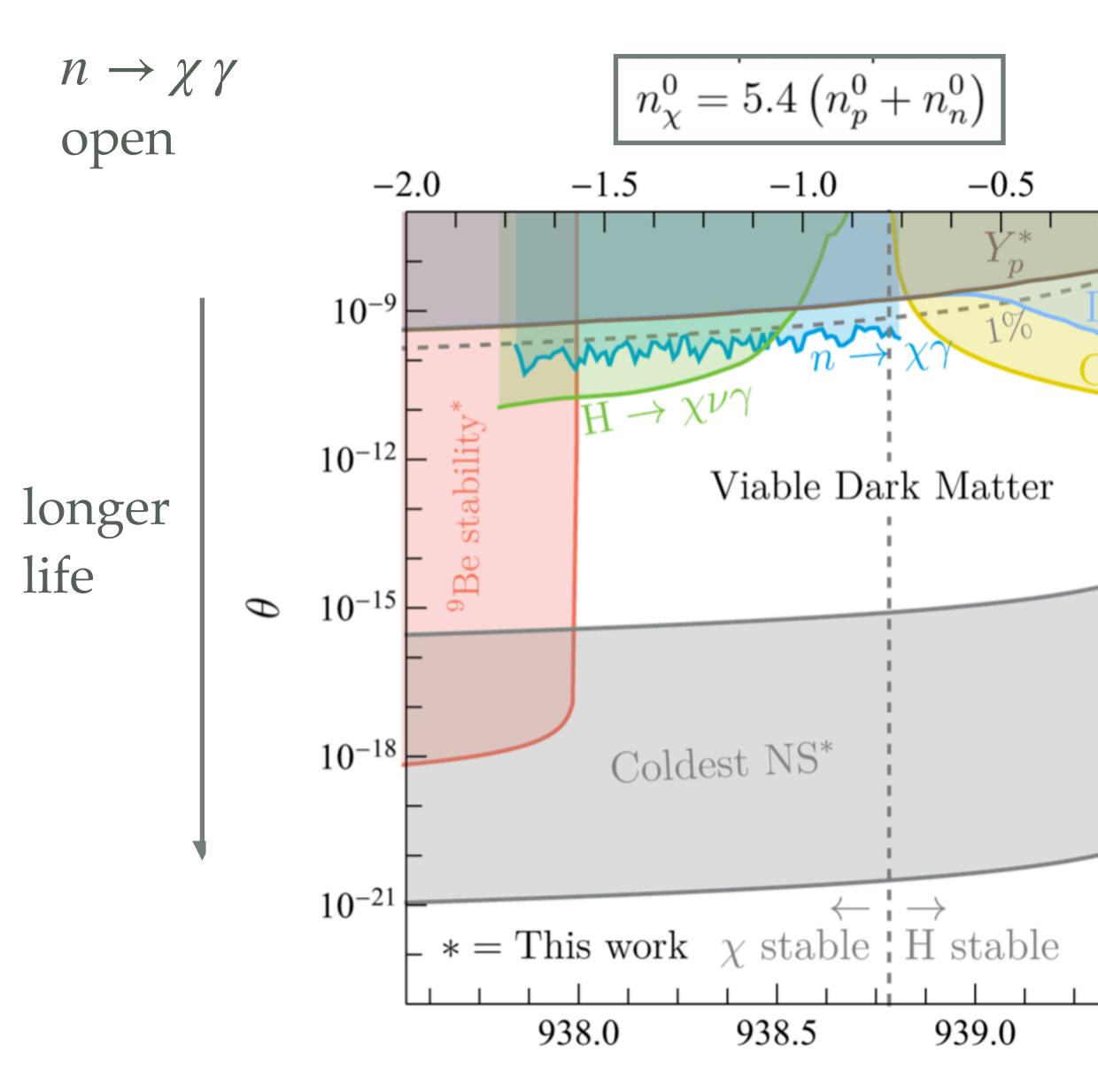
nucleon Fermi sea

neutron Fermi energy ~ 100 MeV

optimized for ~2000 K



Constraints



• BBN data: $Y_p = 0.245 \pm 0.004$, $D/H = (2.55 \pm 0.03) \times 10^{-5}$, ${}^{3}\text{He/H} = (1.0 \pm 0.5) \times 10^{-5}$,

• CMB limit:
$$f_{\chi}/\tau_{\chi} \lesssim 10^{-25} \text{ s}^{-1}$$

T. R. Slatyer, Physical Review D 87 (2013), 10.1103/physrevd.87.123513. J. M. Cline and P. Scott, JCAP 03, 044 (2013), [Erratum: JCAP 05, E01 (2013)], arXiv:1301.5908 [astro-ph.CO].

• $n \rightarrow \chi \gamma$ direct search: 1802.01595 [nucl-ex]

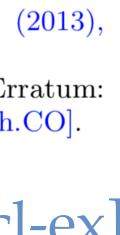
• $H \rightarrow \chi \nu \gamma$: Borexino recast by McKeen, Pospelov (2003.02270)

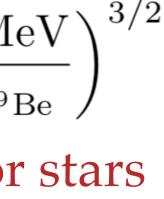
• ${}^{9}\text{Be} \rightarrow 2 \,{}^{4}\text{He} + \chi$:

()

Limited by: $\tau_{^{9}\text{Be}} \sim 4 \times 10^{10} \text{ yr} \left(\frac{10^{-19}}{\theta}\right)^2 \left(\frac{1 \text{ MeV}}{\Omega_{^{9}\text{Pe}}}\right)^{3/2}$ < 3 x 10⁹ yr in metal-poor stars

940 • NS: J2144-3933





Highlights

- neutron lifetime puzzle. □ small 100 keV-ish window left for UCN experiments to target!

Thank you! Questions?

Cosmology (BBN + CMB) stringently limits dark neutron explanation of

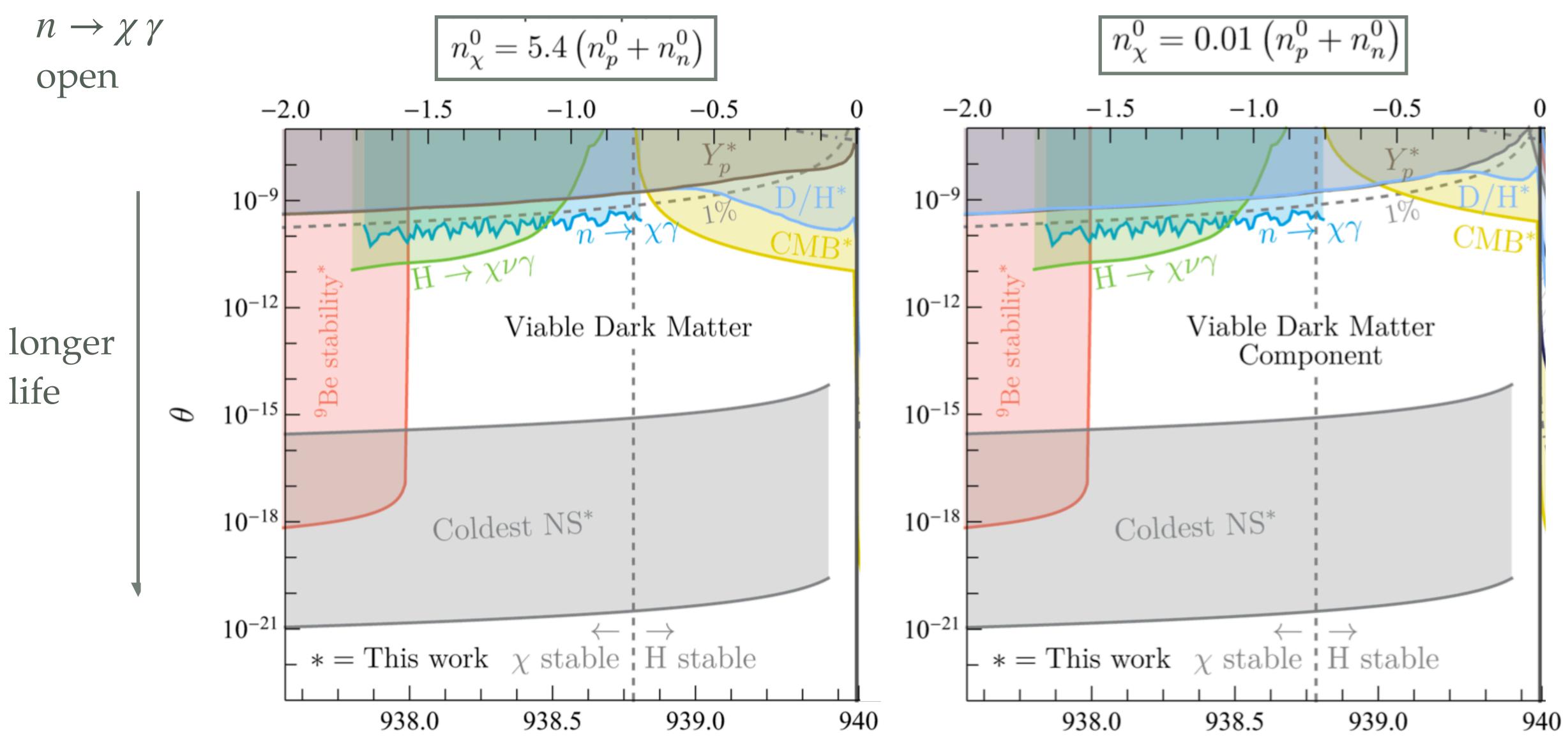
very slow exotic neutron decays => explosive heating of neutron stars.

Heavier-than-neutron dark neutrons (see back-up slides): cosmology sole probe.

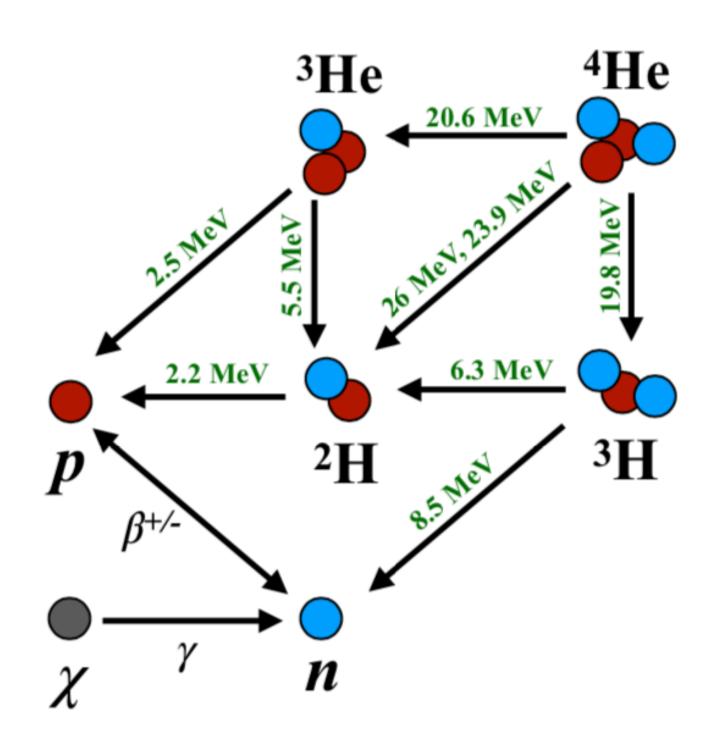


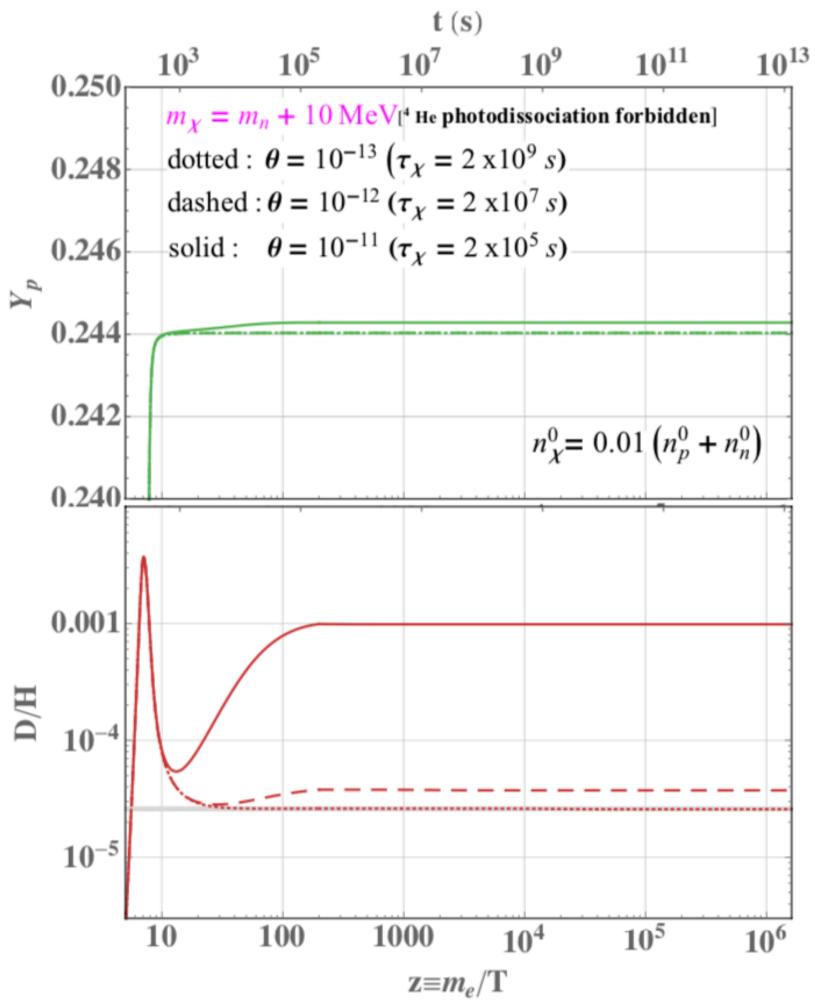
Back-up slides

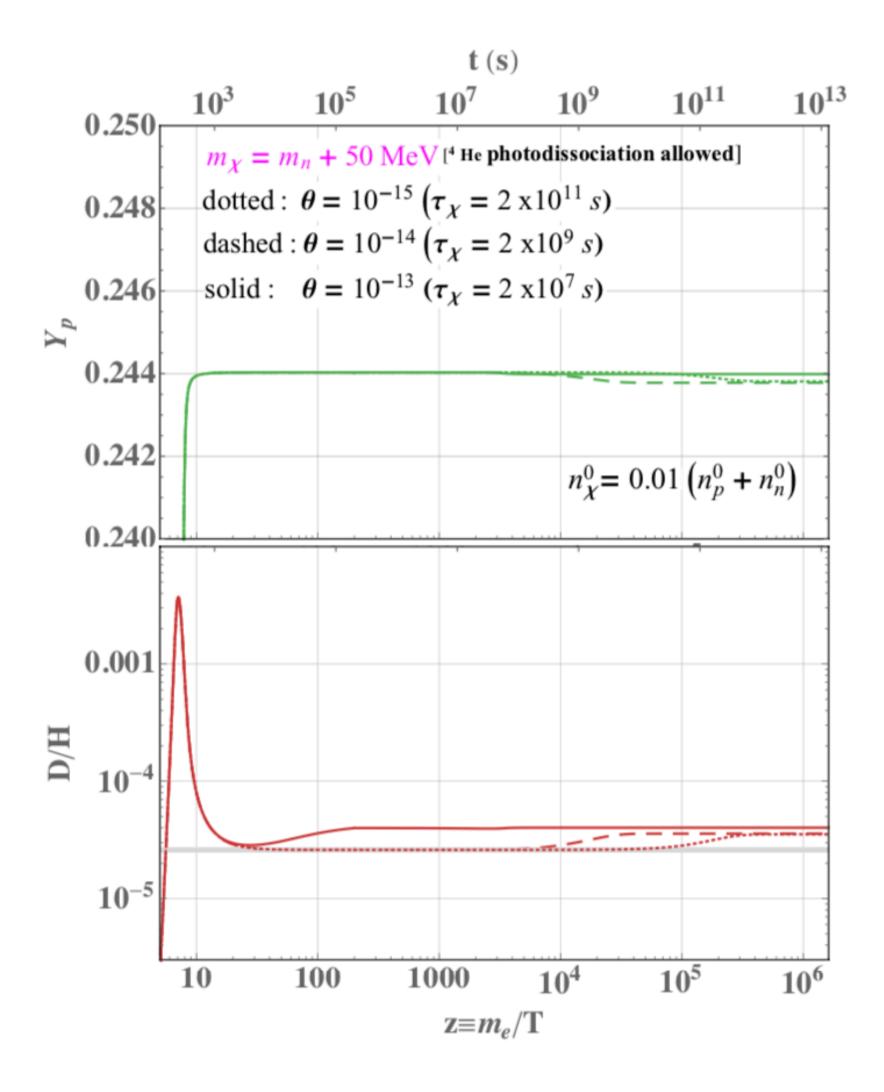
Constraints



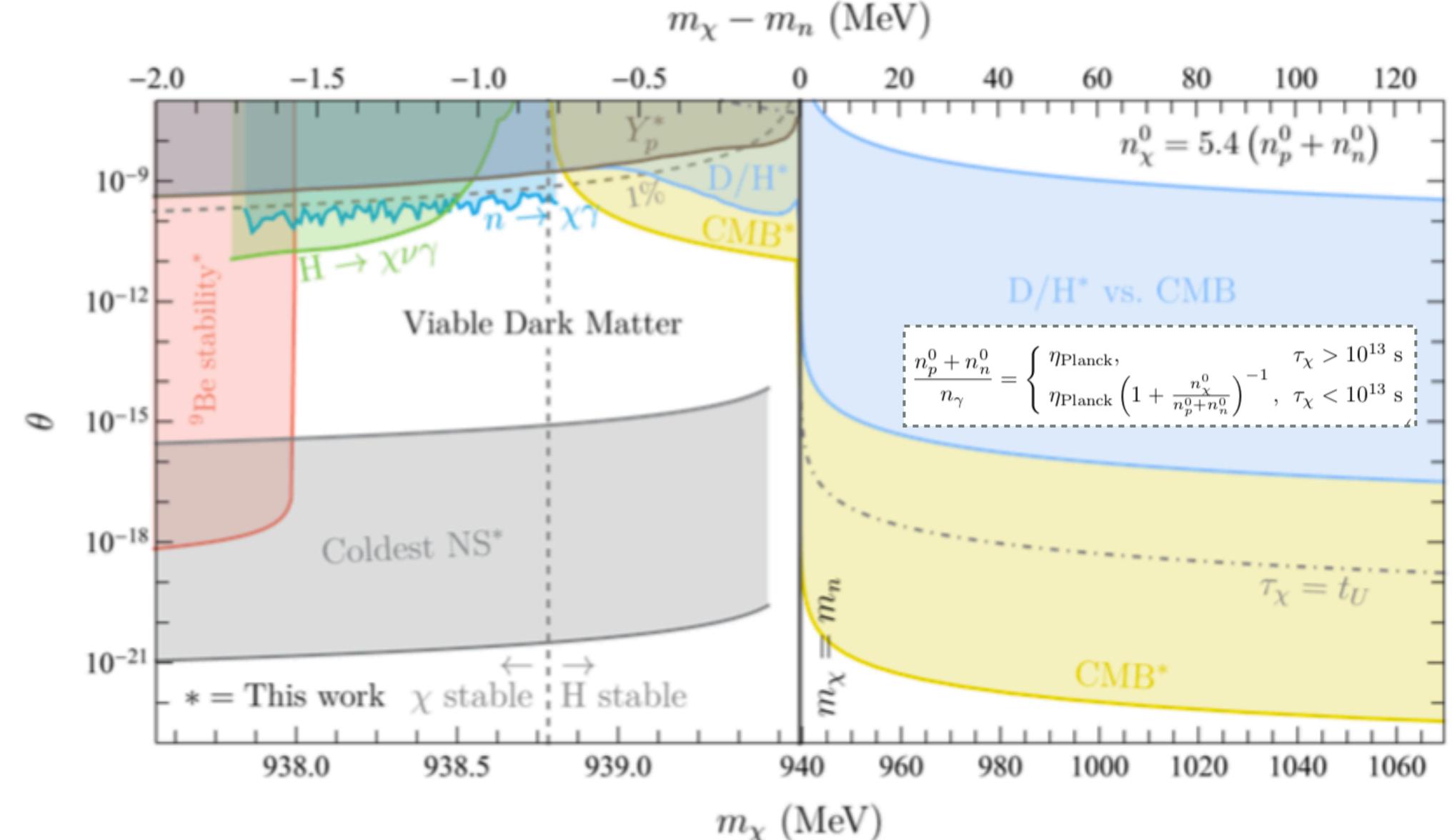
Signals: photodissociation post-nucleosynthesis







Constraints: χ all the dark matter



Constraints: χ percent-level dark matter



