Flavor physics: a lighthouse for navigating the hazy BSM seas

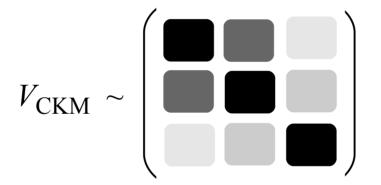
Gino Isidori

[*University of Zürich*]

- The two flavor puzzles
- ► Flavor non-universal interactions [*flavor deconstruction*]
- ► Status of the B anomalies [what we learned, what's left]
- ► Future prospects in flavor physics







There are two (long-standing) open issues in flavor physics:

- I. The observed pattern of SM Yukawa couplings does not look accidental
 - → Is there a deeper explanation for this peculiar structures?

[SM flavor puzzle]

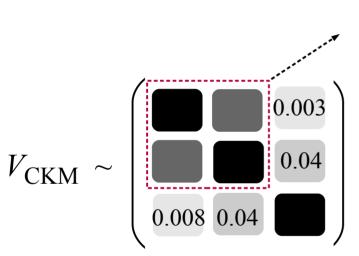
II. If the SM is only an effective theory, valid below an ultraviolet cut-off, why we do not see any deviation from the SM predictions in the (suppressed) flavor changing processes?

[NP flavor puzzle]

→ Which is the flavor structure of physics beyond the SM?

There are two (long-standing) open issues in flavor physics:

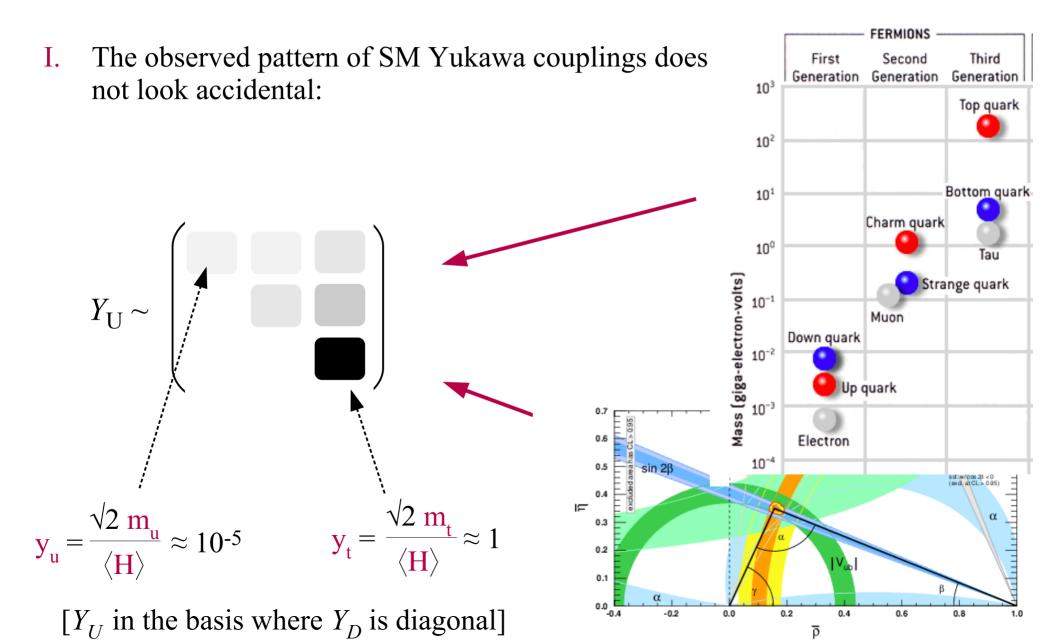
I. The observed pattern of SM Yukawa couplings does not look accidental:



unitarity violation of the 2×2 (light) block below 10⁻³!

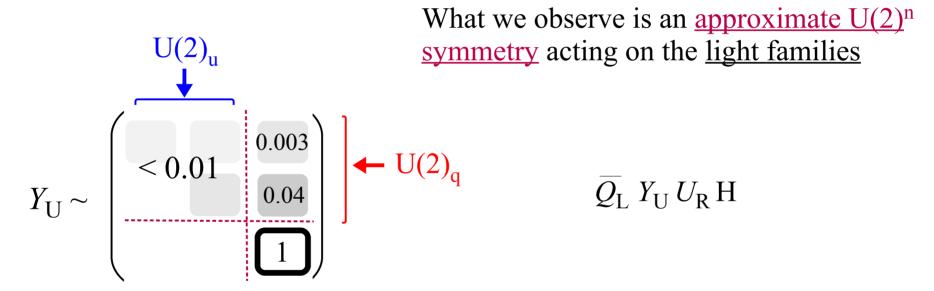
N.B.: Despite the very good knowledge we have nowadays about the CKM matrix, we are not able to detect the presence of the 3rd family by looking only at the 2×2 block (as one naively would have expected...)

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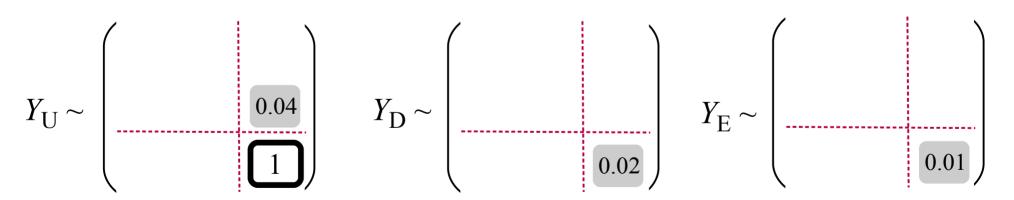


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What we observe is an <u>approximate U(2)</u>ⁿ <u>symmetry</u> acting on the <u>light families</u>

Neglecting entries < 0.01 [@ the EW scale]:



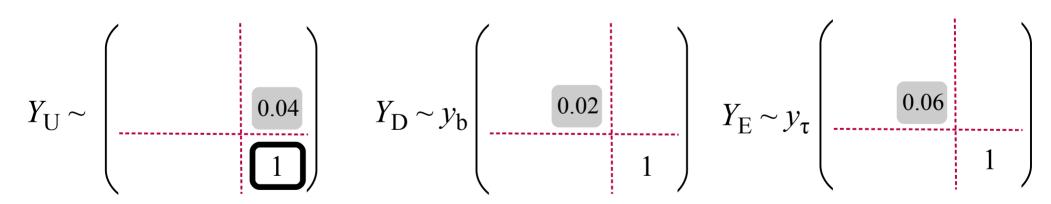
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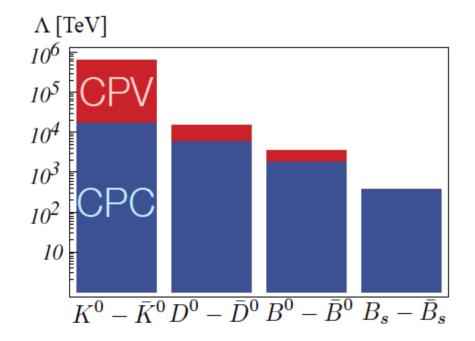
However, beside the overall normalization, in all cases there is a strong hierarchical structure: $U(3) \rightarrow U(2) \times U(1)$

There are two (long-standing) open issues in flavor physics:

II. Why we do not see any deviation from the SM predictions in the (suppressed) flavor changing processes?

$$\mathscr{L}_{\text{SM-EFT}} = \mathscr{L}_{\text{gauge}} + \mathscr{L}_{\text{Higgs}} + \left(\sum_{d,i} \frac{c_i^{[d]}}{\Lambda^{d-4}} O_i^{d \ge 5}\right)$$

Stringent bounds on the scale of possible new <u>flavor non-universal interactions</u>:

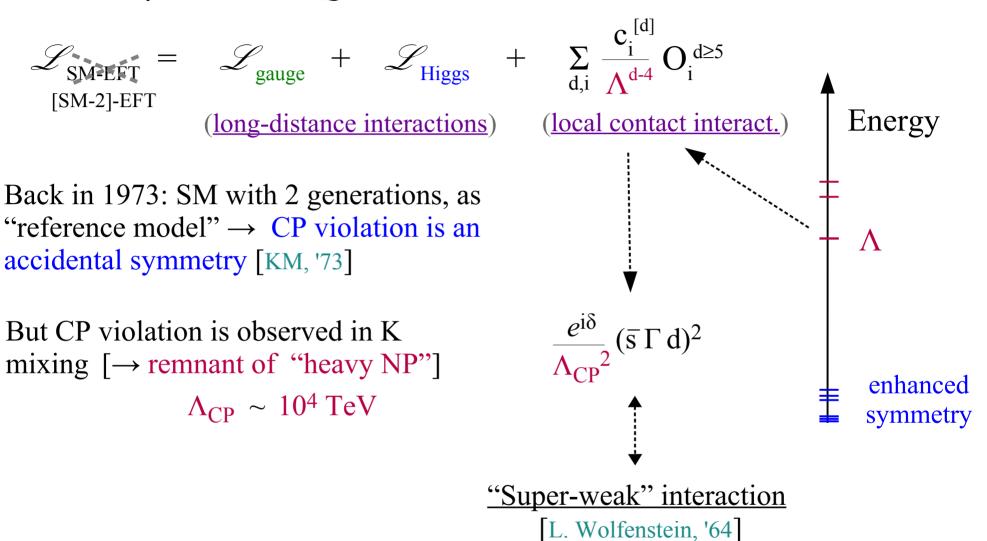


These (*over-emphasized*) high scales can be a "mirage"...

The only unambiguous message is:

No large breaking of the approximate U(2)ⁿ flavor symmetry at near-by energy scales

Accidental symmetries in QFT



Accidental symmetries in QFT

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$$(\text{long-distance interactions}) \qquad (\text{local contact interact.})$$
Back in 1973: SM with 2 generations, as "reference model" \rightarrow CP violation is an accidental symmetry [KM, '73]

But CP violation is observed in K mixing [\rightarrow remnant of "heavy NP"]
$$\Lambda_{\text{CP}} \sim 10^4 \, \text{TeV}$$
enhanced symmetry

<u>SM-3</u> [KM, '73]

 $\frac{1}{\Lambda_{\text{CP}}^{2}} \sim \frac{(G_{\text{F}} m_{\text{t}} V_{\text{ts}} V_{\text{td}})^{2}}{4\pi^{2}}$ Ellis, Gaillard,
Nanopulos, '76

Key message: beware of seemingly high scales in EFT approaches: they can be a "mirage"...

$$\mathscr{L}_{\text{SM-EFT}} = \mathscr{L}_{\text{gauge}} + \mathscr{L}_{\text{Higgs}} + \sum_{d,i} \frac{c_i^{[d]}}{\Lambda^{d-4}} O_i^{d \ge 5}$$

Flavor-degeneracy: U(3)⁵ symmetry

Yukawa couplings:

 $U(3)^5 \rightarrow V(2)^n$

peculiar breaking of the flavor symm.

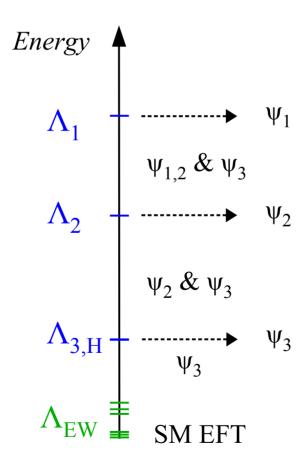
Stringent bounds on generic flavor-violating ops.

‡

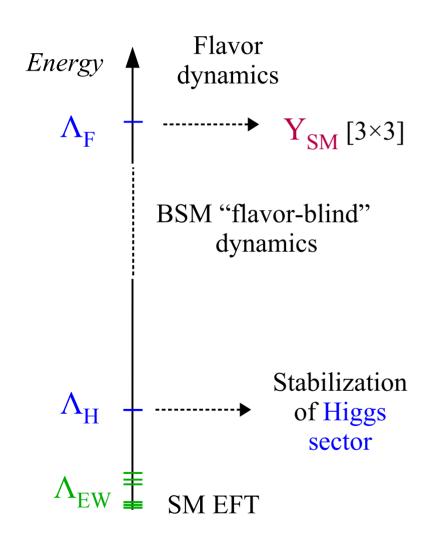
approx. U(2)ⁿ holds also beyond the SM

The big questions in flavor physics:

- Can we find an explanation for the Yukawa hierarchies?
- Are the approximate flavor symmetries accidental symmetries (*in QFT sense*)? If so, at which scale(s) do they appear, and at which scale(s) are they broken?



For a long time, the vast majority of model-building attempts to extend the SM was based on the *implicit* hypotheses of *flavor-universal* New Physics



- Concentrate on the Higgs hierarchy problem
- Postpone the flavor problem to higher scales

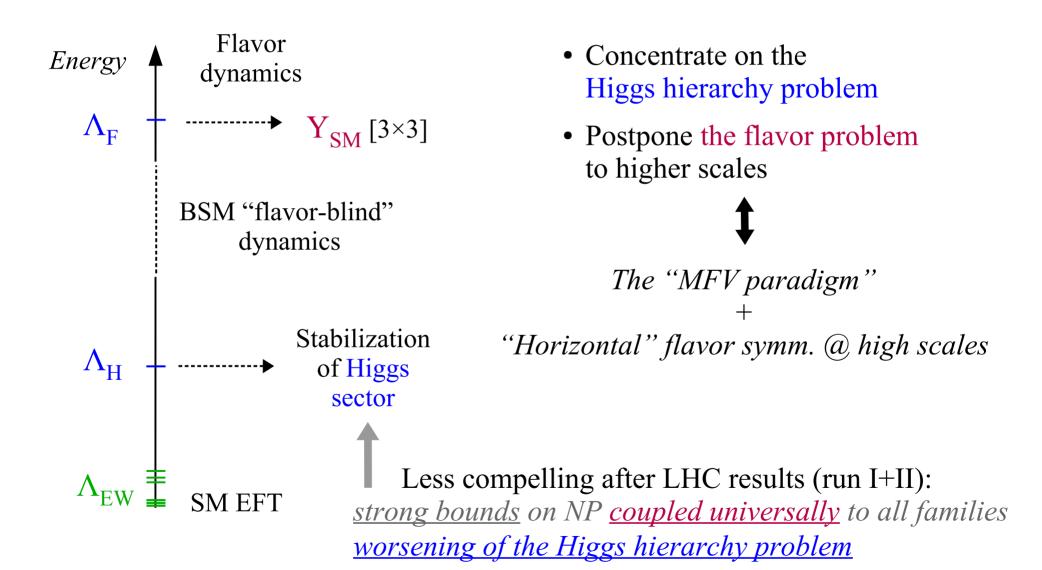


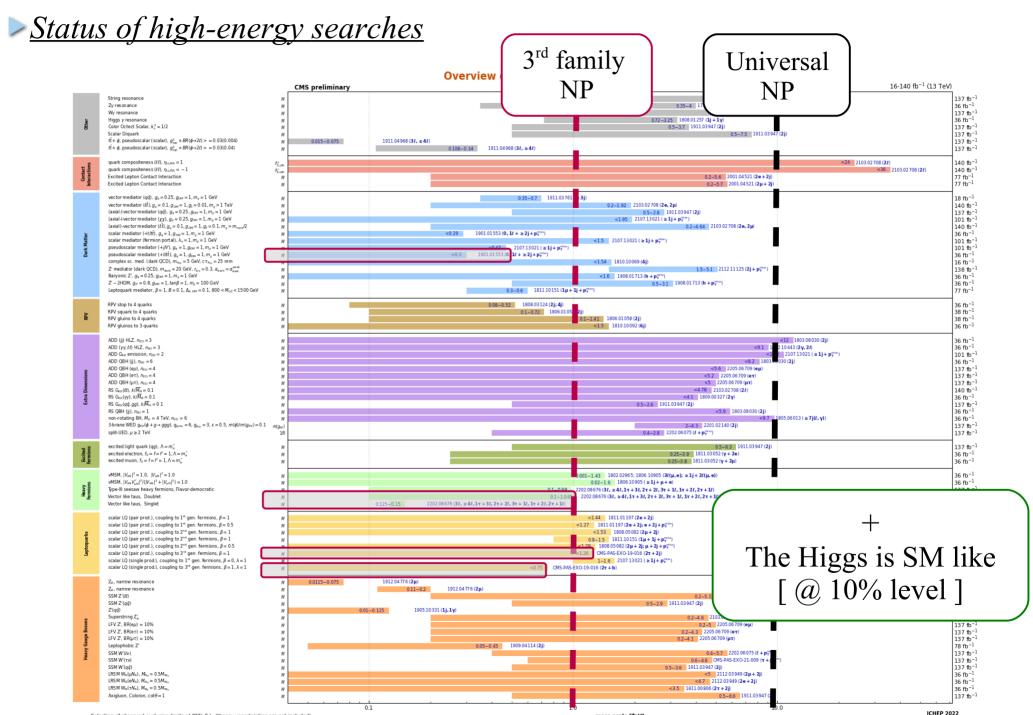
The "MFV paradigm"

"Horizontal" flavor symm. @ high scales

3 gen. = "identical copies" up to high energies

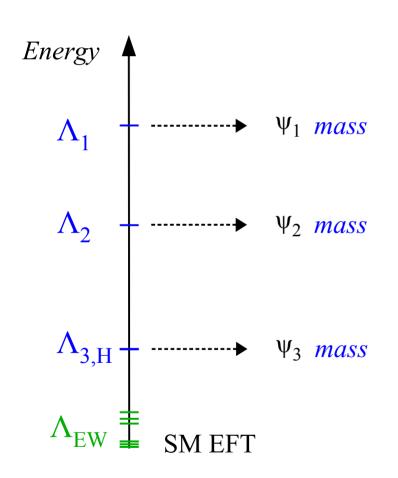
For a long time, the vast majority of model-building attempts to extend the SM was based on the *implicit* hypotheses of *flavor-universal* New Physics





mass scale [TeV]

A more efficient paradigm to address <u>both</u> flavor puzzles (I+II), & *possibly* the Higgs hierarchy, is a <u>multi-scale</u> UV with <u>flavor non-universal</u> interactions



Dvali & Shifman '00
Panico & Pomarol '16
:
Bordone *et al.* '17
Allwicher, GI, Thomsen '20
Barbieri '21
Davighi & G.I. '23

Main idea:

Flavor deconstruction of the SM gauge symm. already at the TeV scale → flavor hierarchies emerge as accidental symmetries

1st & 2nd gen. have small masses (& small coupling to NP) because they are coupled to NP at heavier scales



3 gen. = "identical copies"
up to high energies

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* *Flavor deconstruction* of the SM gauge symmetries:

Arkani-Hamed, Cohen, Georgi '01 Craig, Green, Katz, '11

$$G_{SM}^{[3]} \times G_{SM}^{[12]} \to G_{SM}^{[12]}$$

[last step @ $\sim TeV scale$]

A more intricate "weaving" vs.

- "vertical" group containing G_{SM}
- "horizontal" flavor symmetry

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Some notable points:

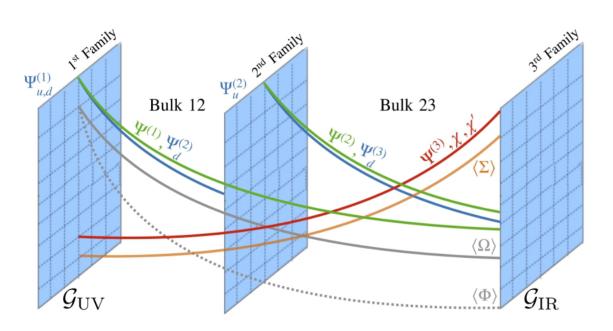
- **This symmetry breaking pattern is very general:** any scalar rep. (provided $R_{[3]}$, $R_{[12]}$ ≠ 1) breaks to the diagonal subgroup
 - → flavor universality emerges "naturally" at low energies

Craig, Garcia-Garcia, Sutherland '17

✓ Flavor hierarchies obtained without the need of peculiar choices for the "flavor charges" [e.g. U(1) charges as in Froggatt & Nielsen]

A more efficient paradigm to address <u>both</u> flavor puzzles (I+II), & *possibly* the Higgs hierarchy, is a <u>multi-scale</u> UV with <u>flavor non-universal</u> interactions

Looking more into the UV, this picture could arise from an extra-dimensional construction:



Flavor ↔ special position (topological defect) in an extra (compact) space-like dimension

Dvali & Shifman, '00

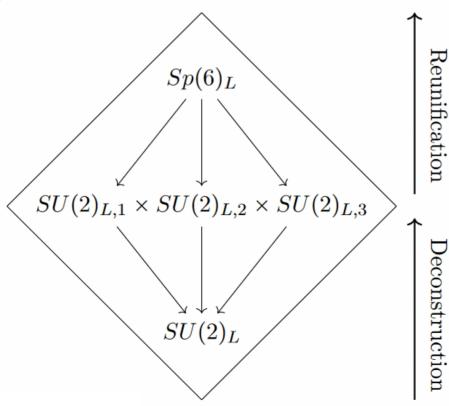
- * Flavor hierarchies related to brane distances
- * "Holographic" Higgs from appropriate choice of bulk/brane gauge symmetries
- * Anarchic neutrino masses via inverse see-saw mechanism

A more efficient paradigm to address <u>both</u> flavor puzzles (I+II), & *possibly* the Higgs hierarchy, is a <u>multi-scale</u> UV with <u>flavor non-universal</u> interactions

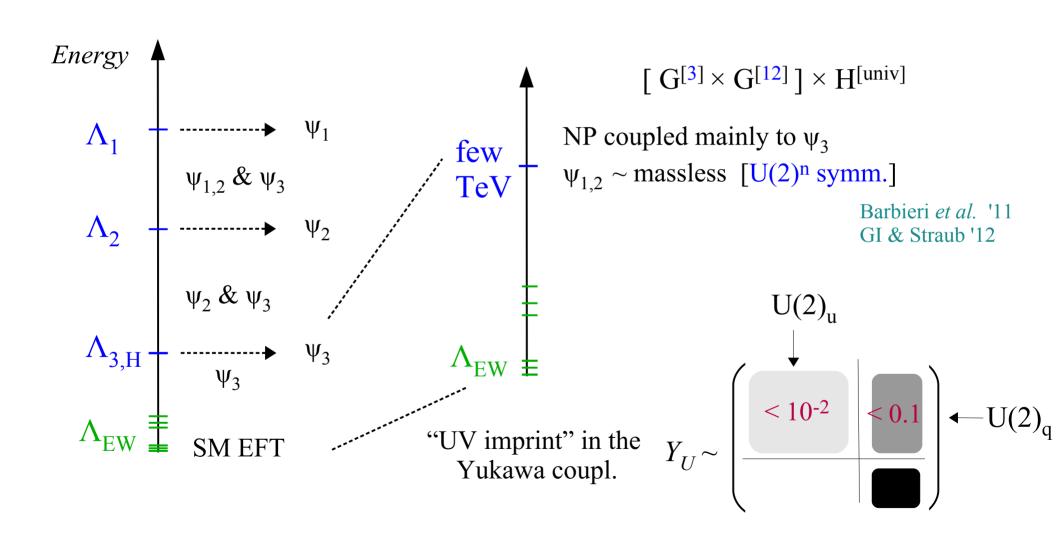
Looking more into the UV, this picture could arise from an extra-dimensional construction, or as limit of larger unified gauge groups.

E.g.: "electroweak-flavor" unification

Davighi & Tooby-Smith '22 :



A more efficient paradigm to address <u>both</u> flavor puzzles (I+II), & *possibly* the Higgs hierarchy, is a <u>multi-scale</u> UV with <u>flavor non-universal</u> interactions



Effective organizing principle for the flavor structure of the SMEFT

To understand which are the viable options for TeV-scale dynamics, we recently analysed all the extensions of the SM gauge group compatible with the following three general assumptions:

Davighi & G.I. '23

- I. Obtain the U(2)ⁿ flavor symmetry as accidental symmetry of the (non- universal) gauge sector
- II. Elementary Higgs up to (at least) the TeV scale → New states should preserve Higgs-mass stability → NP coupled to 3rd generation should occur at the TeV scale
- III. Explain charge-quantization \rightarrow Semi-simple embedding in the UV [i.e. no U(1) groups in the UV]

 $\overline{\psi}_L Y \psi_R H$

- Flavor hierarchies from gauge non-universality [a brief detour]
 - I. $U(2)^n$ flavor symmetry as accidental symmetry of the gauge sector.
 - Classify the allowed Yukawa structures under a flavor-deconstruction of three basic factors characterizing the SM fermions and the EW gauge group: $SU(2)_L \times U(1)_R \times U(1)_{B-L}$



• Deconstructing <u>any pair of the three</u> (or all of them) leads to the desired U(2)ⁿ flavor symmetry

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 - Deconstructing <u>any pair of the three</u> (or all of them) leads to the desired U(2)ⁿ flavor symmetry:
 - ✓ Part of the EW group necessarily need to be deconstructed
 - Minimal choice represented by SM hypercharge [$Y=T^3_R+(B-L)/2$]. However, $U(1)^{[3]}_Y \times U(1)^{[2]}_Y \times U(1)^{[1]}_Y$ has two drawbacks:

 Navarro & King '23
 Davighi & Stefanek '23
 - * No immediate semi-simple embedding
 - Same deconstruction mechanism protecting both mixing & light masses → not vey efficient in suppressing flavor bounds

II.+III. Explain charge-quantization → Semi-simple embedding in the UV

But new states should preserve Higgs-mass stability \rightarrow NP coupled to 3rd generation should occur at the TeV scale



Semi-simple embeddings of the SM have been classified and there are very few possibilities, all featuring one of the possible 3 basic options:

Allanach, Gripaios, Tooby-Smith '23

- $SU(4)\times SU(2)\times SU(2)$ [Pati & Salam '74]
- SU(5) [Georgi & Glashow, '74]
- SO(10) [Georgi '75, Fritzsch & Minkowski '75]



Proton stability → only the Pati-Salam option is possible at low scales

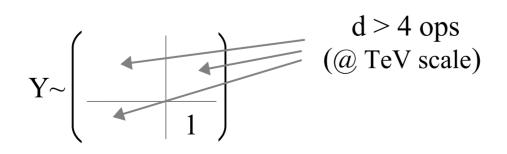
$$SU(3)_{c} \times U(1)_{B-L} \hookrightarrow SU(4) \sim \begin{bmatrix} SU(3)_{c} & 0 \\ \hline 0 & 0 \end{bmatrix} \begin{bmatrix} 0 & LQ \\ \hline LQ & 0 \end{bmatrix} \begin{bmatrix} 1/3 & 0 \\ \hline 0 & -1 \end{bmatrix}$$

I. + II. + III. : four basic options:

	TeV-scale gauge group: $G_U \times G_3 \times H_{12}$				
	G_U	G_3	H_{12}		
1	$SU(2)_L$	$SU(4)^{[3]} \times SU(2)_R^{[3]}$	$SU(3)^{[12]} \times U(1)_{B-L}^{[12]} \times U(1)_{R}^{[12]}$		
2	$SU(2)_R$	$SU(4)^{[3]} \times SU(2)_L^{[3]}$	$SU(3)^{[12]} \times SU(2)_L^{[12]} \times U(1)_{B-L}^{[12]}$		
3	SU(4)	$SU(2)_L^{[3]} \times SU(2)_R^{[3]}$	$SU(2)_L^{[12]} \times U(1)_R^{[12]}$		
4	Ø	$SU(4)^{[3]} \times SU(2)_L^{[3]} \times SU(2)_R^{[3]}$	$\mathrm{SU}(3)^{[12]} \times \mathrm{SU}(2)^{[12]}_L \times \mathrm{U}(1)^{[12]}_{B-L} \times \mathrm{U}(1)^{[12]}_R$		

Higgs & 3rd gen. fields charged only under these groups

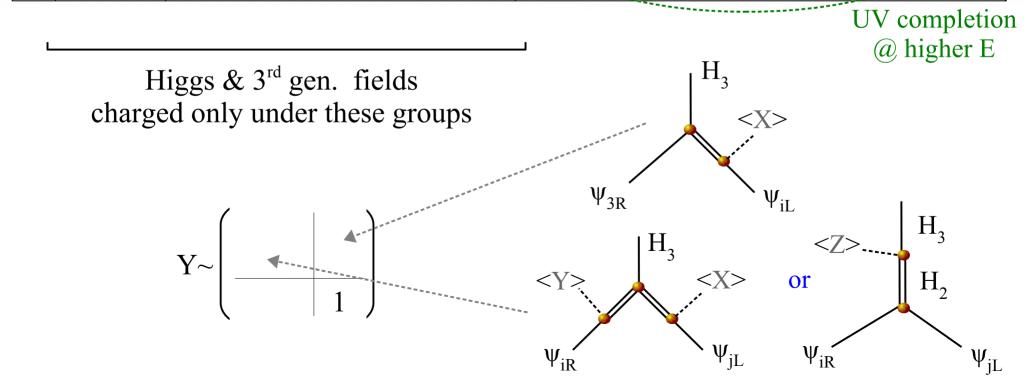
 $\begin{array}{c} & \text{UV completion} \\ & \text{@ higher E} \\ \\ \text{small impact on } \delta m_h \end{array}$



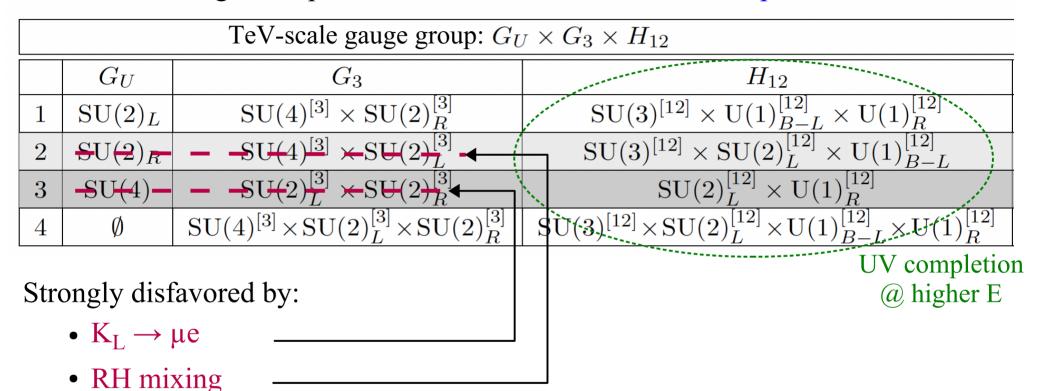
Allwicher, GI, Thomsen '20 Davighi, G.I., Pesut '22 Davighi & G.I. '23

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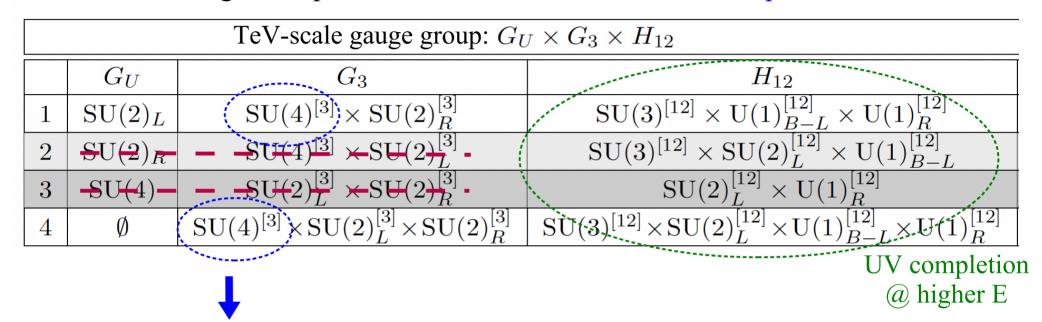


I. + II. + III. + general pheno bounds: two viable TeV-scale options:



$$Y \sim \left(\begin{array}{c|c} & & & \\ \hline & & 1 \end{array}\right)$$

I. + II. + general pheno bounds: two viable TeV-scale options:

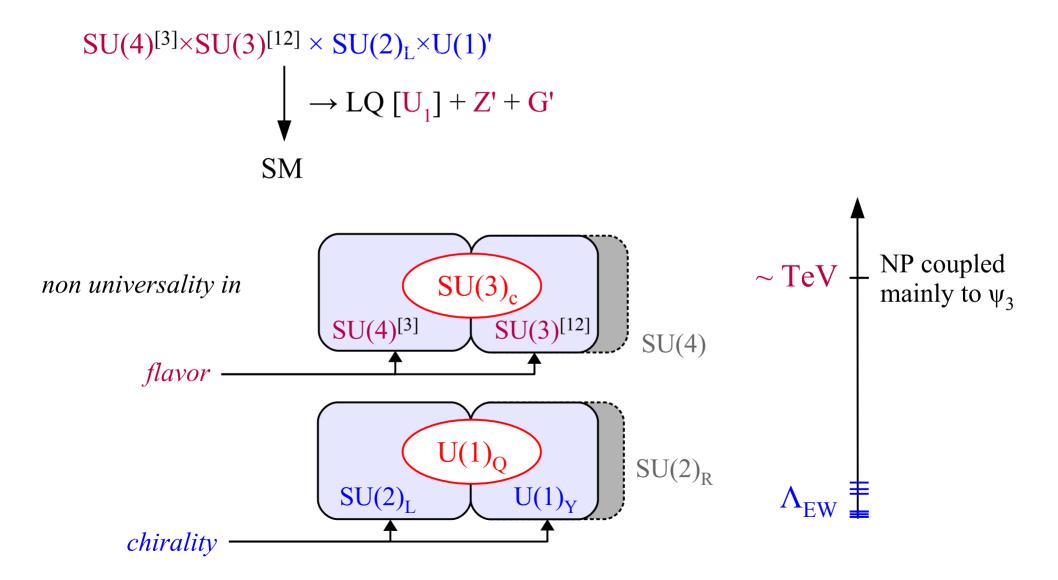


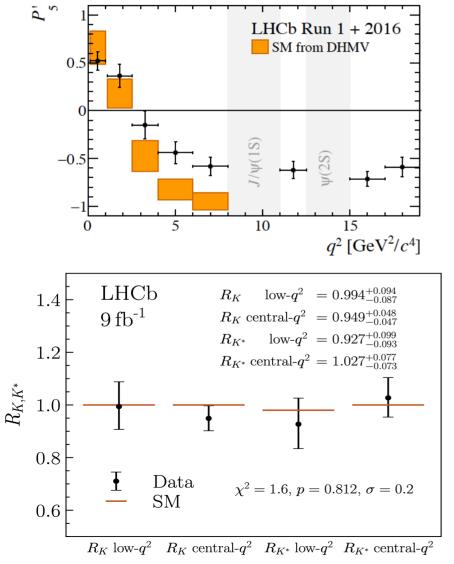
General feature:

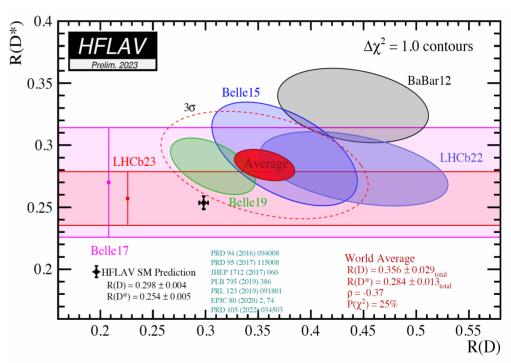
SU(4) group acting on the 3rd family, with low-energy breaking scale to avoid fine-tuning on the Higgs mass:

$$\delta m_h^2/m_h^2 < 1 \rightarrow \Lambda_U = M_U/g_U \lesssim 5 \text{ TeV}$$

This connects with the class of consistent TeV-scale models proposed to address the B-physics anomalies...







From 2013 results in (various) semi-leptonic B decays started to exhibit tensions with the SM predictions. Several exclusive channels are involved, but they are all sensitive only to the following two classes of partonic transitions:

$$b \rightarrow c lv$$
 (Charged Currents) $b \rightarrow s l^+l^-$ (Neutral Currents)

Most of the anomalies are connected to a possible breaking of Lepton Flavor Universality = <u>accidental symmetry</u> of the SM Lagrangian in the limit where we neglect the lepton Yukawa couplings

Even if the significance went down recently (*not completely...*), worth to discuss as example of consistent TeV-scale (new) physics that could be revealed by precision flavor experiments

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The anomalies can be grouped into 3 categories:

- I. LFU anomaly in CC [τ vs. (μ , e)] $b \rightarrow c lv$
- II. ΔC_9 (*lepton-universal*) anomaly in NC modes
- (III.) LFU anomaly in NC [μ vs. e] & BR(B_s $\rightarrow \mu\mu$)

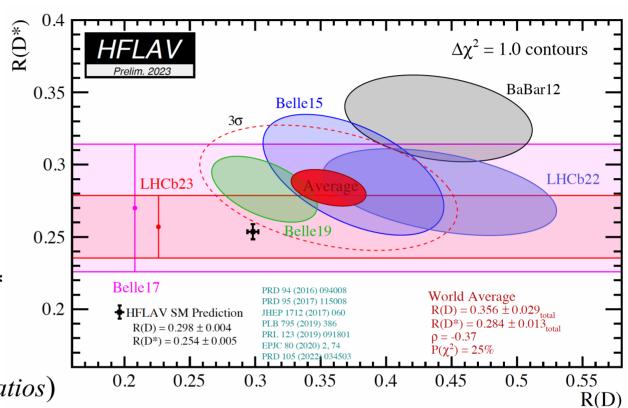
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I. LFU anomaly in CC[τ vs. (μ, e)]

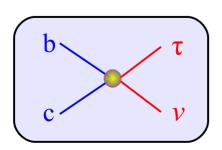
$$R(X) = \frac{\Gamma(B \to X \tau v)}{\Gamma(B \to X l v)}$$

$$X = D \text{ or } D^*$$

Clean SM predictions
 (uncertainties cancel in the ratios)



- 3.0σ excess over SM
- Compete with SM @ tree-level $\rightarrow low scale of NP$

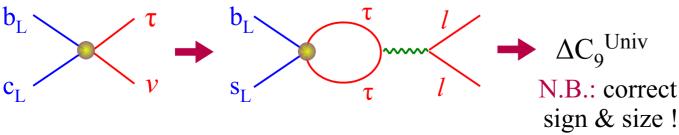


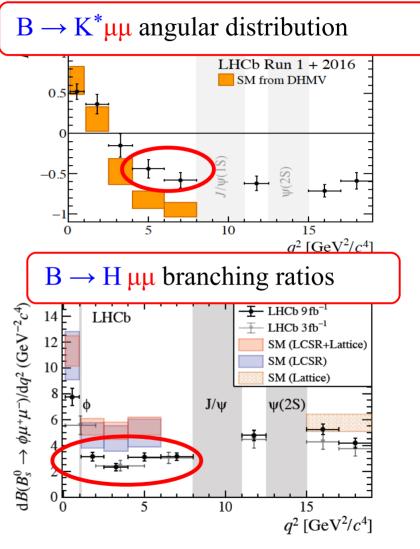
II) ΔC_9 (*lepton-universal*) anomaly in NC modes

$$\mathcal{O}_9^{\ell} = (\bar{s}_L \gamma_\mu b_L)(\bar{\ell} \gamma^\mu \ell)$$

- Possible contamination from SM longdistance (charming penguins)
- All attempts to <u>compute</u> the effect agree on $\sim 3\sigma$ deviation from SM
- Compete with SM @ loop-level

Possible explanation connected to CC (hence 3rd family LFU violation):

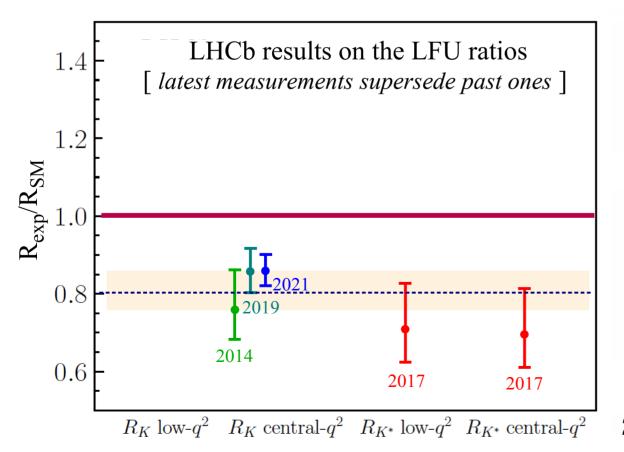


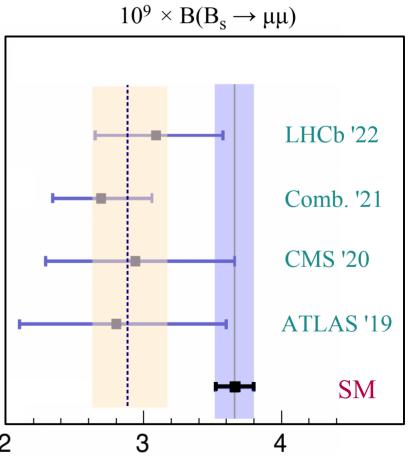


Bobeth & Haisch '11 Crivellin *et al.* '18

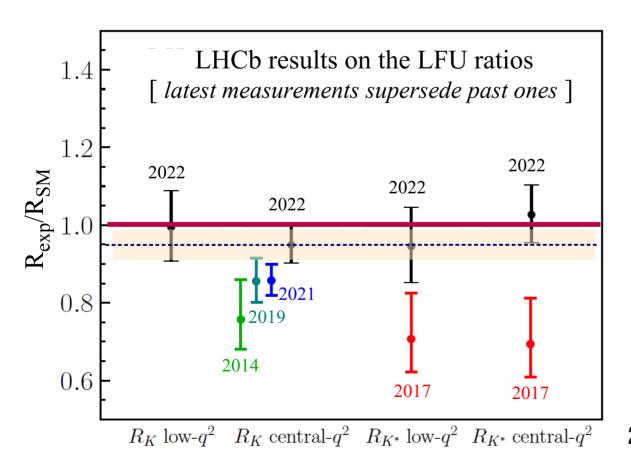
Alguero et al. '18

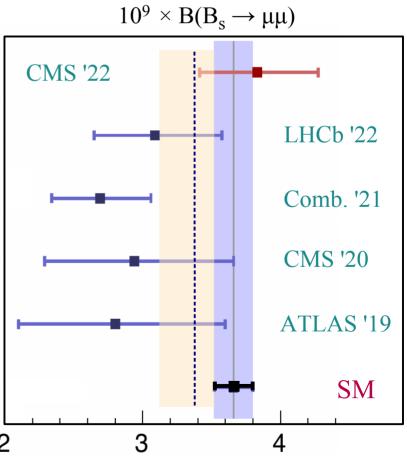
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 - Clean SM predictions $(LFU\ ratios + no\ long-distance\ in\ B_s \rightarrow \mu\mu)$
 - Highest significance till summer 2022





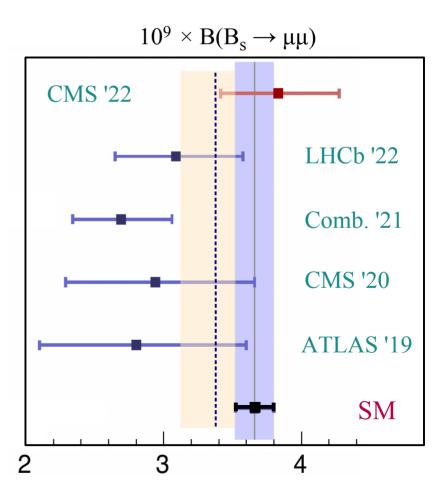
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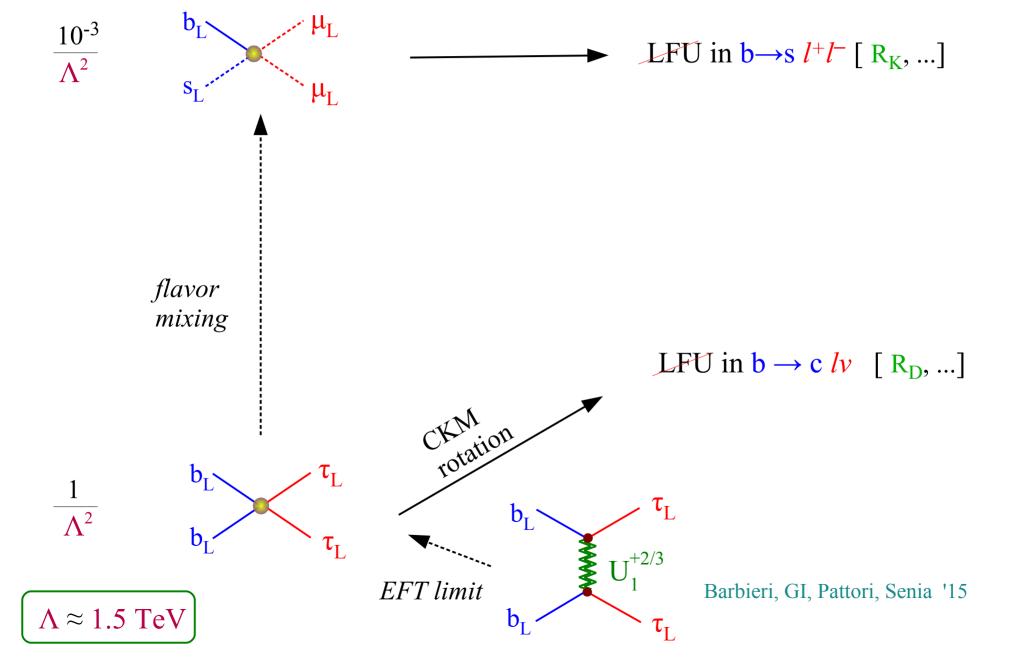


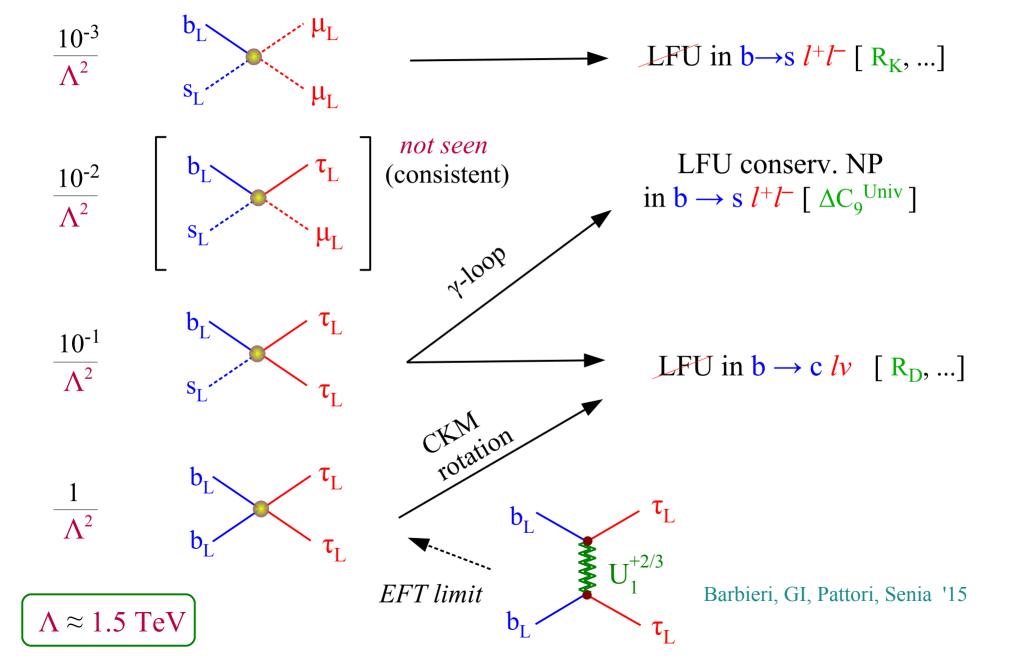


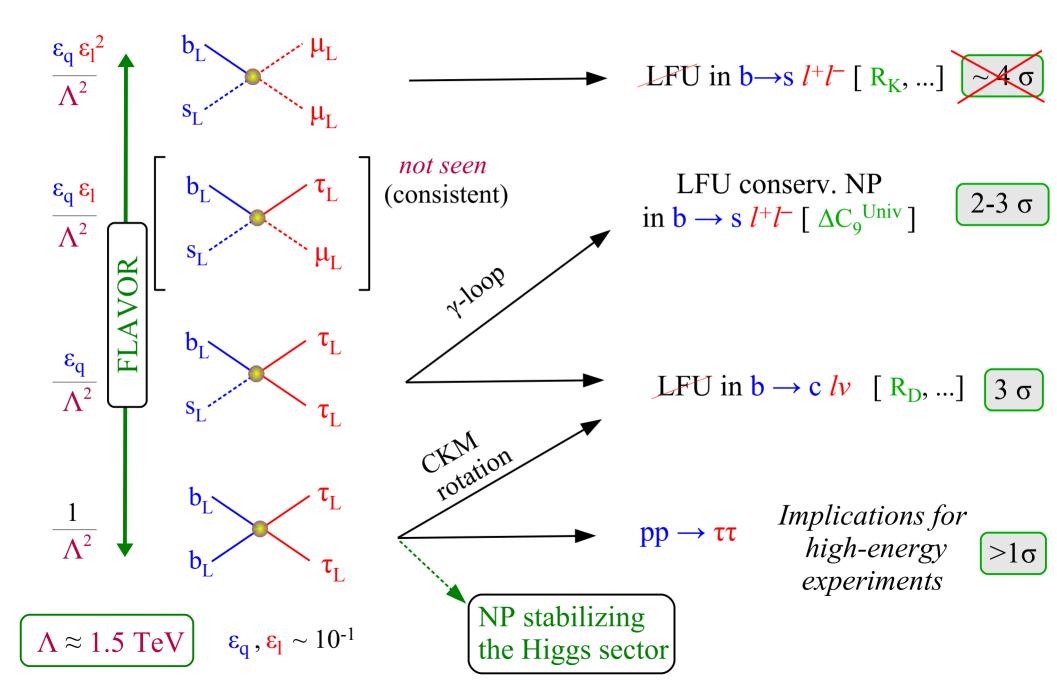
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LHCb results on the LFU ratios 1.4 [latest measurements supersede past ones] 2022 2022 $R_{\rm exp}/R_{\rm SM}$ 2022 2022 0.8 2019 2014 modes 2017 2017 R_K low- q^2 R_K central- q^2 R_{K^*} low- q^2 R_{K^*} central- q^2 N.B.: While the overall loss of significance is high, the overall implications for the class of NP models I advocate, are modest

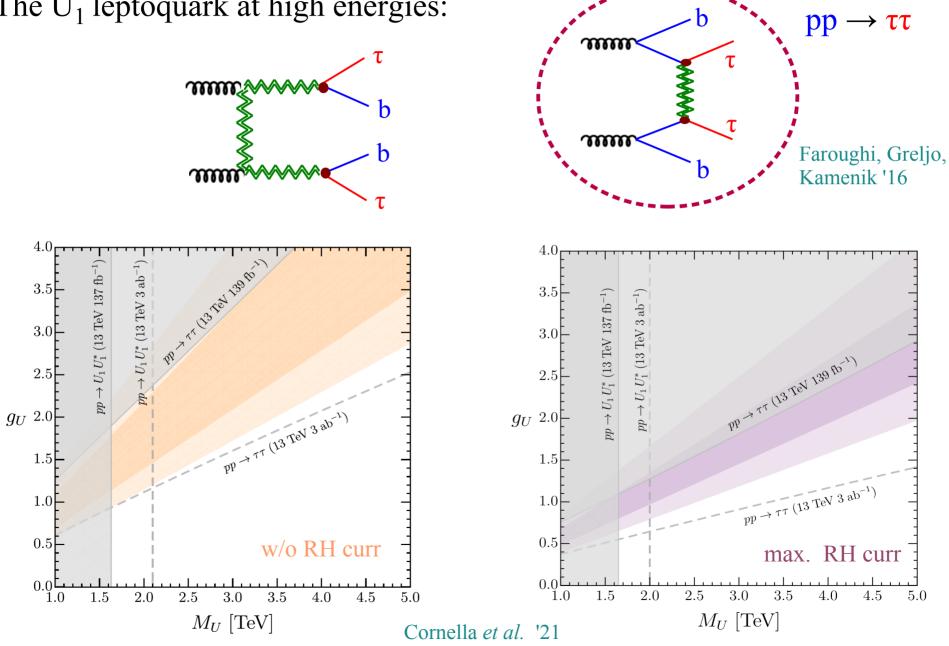








The U₁ leptoquark at high energies:

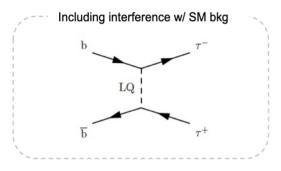




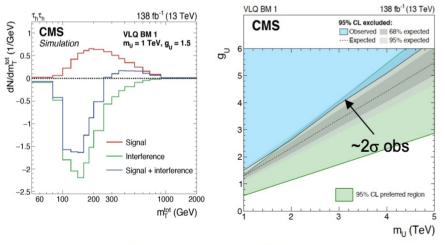


LQ-b-τ: Comparison of recent results

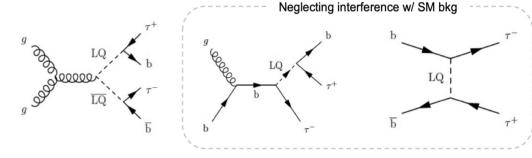




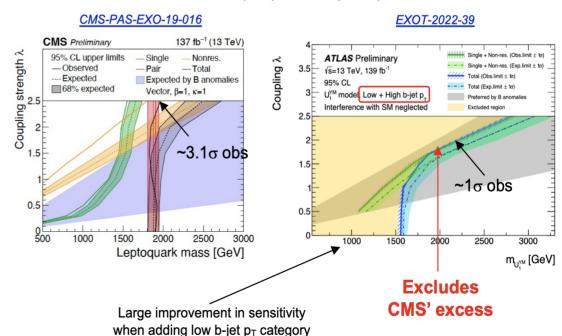
CMS-HIG-21-001



Shown at Moriond EW 2022

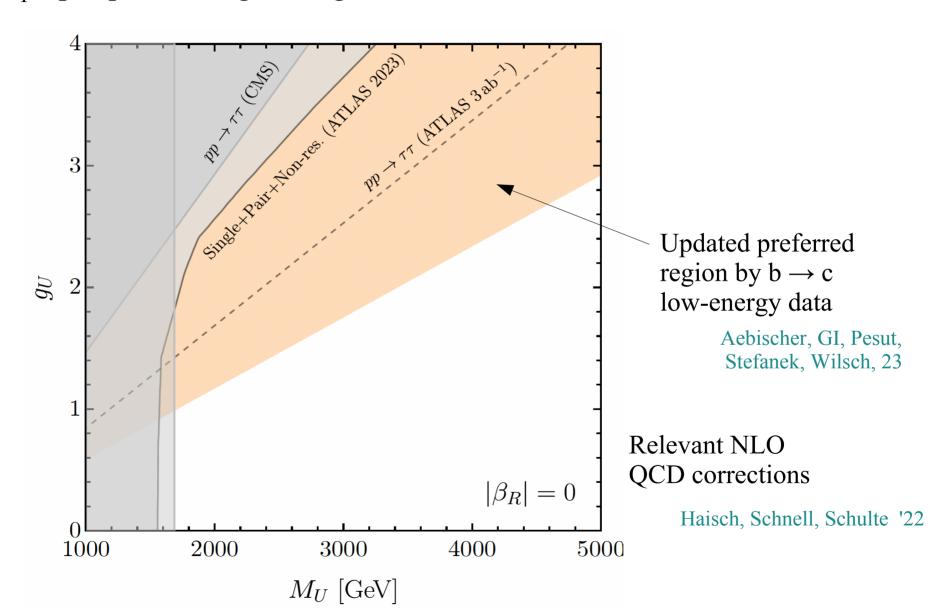


Caveat: BR=1 (CMS) vs BR=0.5 (ATLAS)



Need to clarify interference issue for future interpretations

The U₁ leptoquark at high energies:



In the next ~ 10 years we can expect significant progress in flavor physics from more than one experiment:

- LHCb-II [b & c physics @ hadron collider]: ~ 10 × present statistics
- Belle-II [b & t physics @ e+e-]: $\sim 10 \times \text{statistics of Belle-I}$
- NA62 / HIKE [$K \rightarrow \pi \nu \nu$]: $\sim 3 / 30 \times \text{present statistics}$
- $\mu \rightarrow e \left[\mu \rightarrow 3e \, (PSI) + \mu N \rightarrow e N(FNAL) \right]$: ~ 1000 × present sensitivity

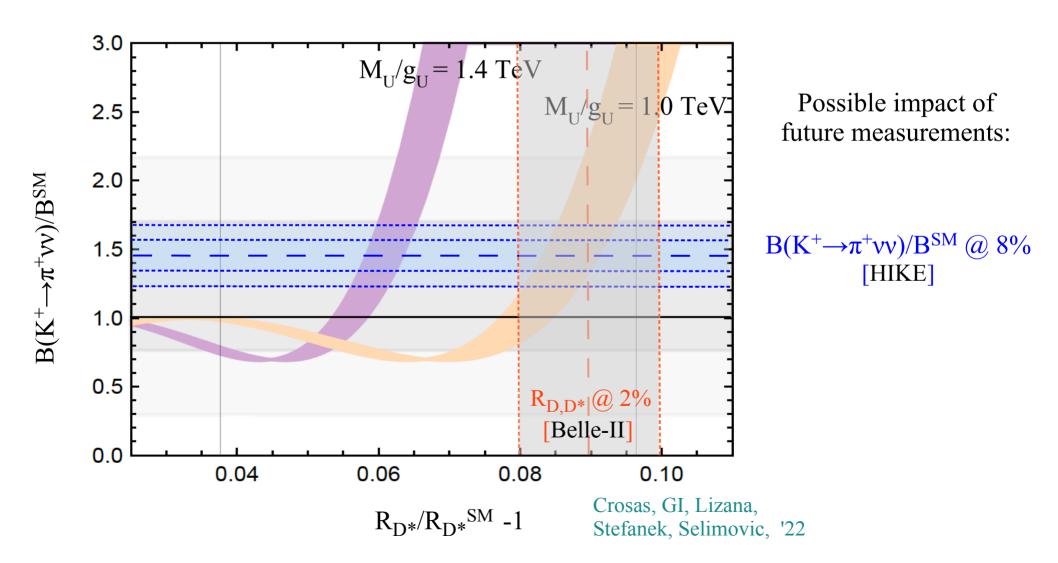
→ Increase in NP sensitivity significantly higher vs. direct searches

$$A(\psi_i \to \psi_j + X) = A_0 \left[1 + \frac{c_{NP} m_W^2}{c_{SM} \Lambda^2} \right]$$

What can we expect

- I. Modest improvements [up to $\sim 1/2$ on the amplitude]:
 - $B_{s,d}$ mixing & $B(b \rightarrow s\gamma) \rightarrow$ dominated by th. errors
- II. Significant improvements:
 - LFU tests in neutral currents $[R_K,...]$: ~ (5-10)% \rightarrow ~ (2-3) %
 - LFU tests in charged currents $[R_D,...]$: ~ (5-10)% \rightarrow ~ (2-3) %
 - B(K $\rightarrow \pi \nu \nu$) & B(B $\rightarrow K \nu \nu$) @ 10% level

E.g.: Correlation between $B(K^+ \to \pi^+ \nu \nu)$ & R_{D,D^*} in models addressing current charged-current anomalies:



What can we expect

- I. Modest improvements [up to $\sim 1/2$ on the amplitude]:
 - $B_{s,d}$ mixing & $B(b \to s\gamma) \to dominated by th. errors$

II. Significant improvements:

- LFU tests in neutral currents $[R_{\kappa},...]$: ~ (5-10)% \rightarrow ~ (2-3) %
- LFU tests in charged currents $[R_D,...]$: $\sim (5-10)\% \rightarrow \sim (2-3)\%$
- B(K $\rightarrow \pi \nu \nu$) & B(B $\rightarrow K \nu \nu$) @ 10% level

III. Expected largest improvements:

- LFV in B decays, especially $b \rightarrow s + \tau \mu$
- LFV in τ decays
- FCNC of type $b \rightarrow s + \tau \tau$

Conclusions

- Flavor physics represents one the most intriguing aspects of the SM and, at the same time, a great opportunity to investigate the nature of physics beyond the SM.
- The idea of a *multi-scale construction & flavor deconstruction* has several appealing aspects: model-building direction largely unexplored so far, that could help shedding light not only on the origin of the flavor hierarchies, but also on the EW hierarchy problem.
- The model-building efforts along this direction, triggered by the B-physics anomalies are still very motivated (*independently of the anomalies fate*).
- If the these ideas corrects, <u>new non-standard effects could emerge soon</u> both at low and at high energies and (some of the) anomalies could just be the first hint...