

# RADIO PROPAGATION IN NON-UNIFORM MEDIA ARENA, JUNE 2022

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## WHY IS THIS IMPORTANT?

Understanding is needed on how radio propagates through natural material (polar ice, atmosphere...)  
=> typically non uniform media

How to simulate propagation?

FDTD, typically computationally expensive

Parabolic equation solver, (<https://arxiv.org/abs/2011.05997>)

Raytracing, a high frequency limit

More specific:

- 1) How to describe Cerenkov emission in non uniform media?
- 2) How to simulate horizontal propagation in ice?

## RAYTRACING WITH FERMAT'S PRINCIPLE

A good place to start

Observation of classically 'forbidden' electromagnetic wave propagation and implications for neutrino detection.

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Describes raytracing with Fermat's principle. Overdensity => Guide the raypaths

# FERMAT'S PRINCIPLE

## HOW DOES IT WORK?

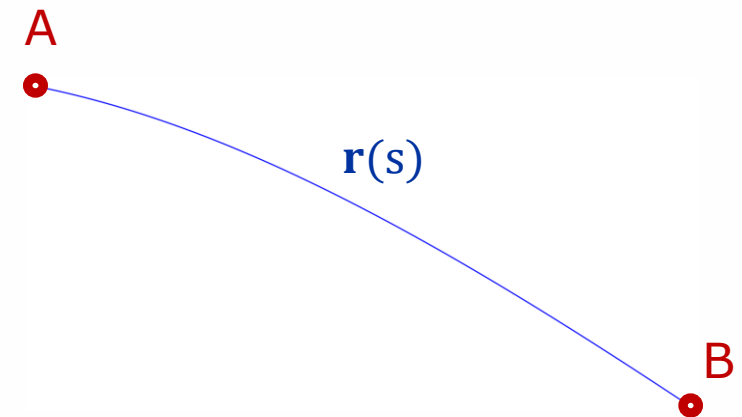
Minimize the variation in the optical path length for a path  $r$  between points A and B

$$OPL = \int_A^B n(\mathbf{r}(s)) ds$$

With  $ds^2 = \mathbf{r}(s) \cdot \mathbf{r}(s)$

Leads to:

$$\frac{\partial n}{\partial \mathbf{r}} - \frac{d}{ds} (n(\mathbf{r}(s)) \dot{\mathbf{r}}) = 0, \quad \dot{\mathbf{r}} = \frac{d\mathbf{r}}{ds}$$



# FERMAT'S PRINCIPLE

## SOLVING FOR THE RAYPATH

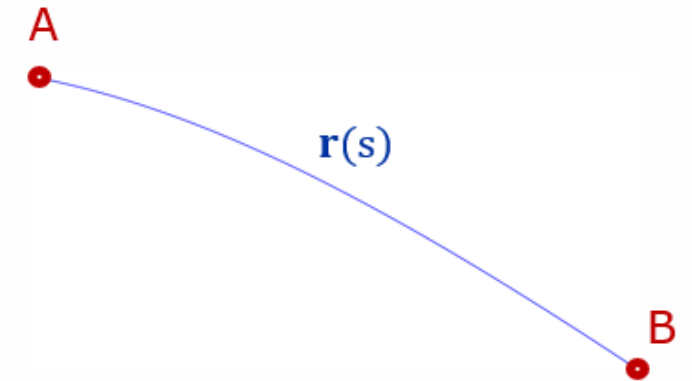
Constrict the ray to be in the  $x,z$  plane, and define the optical momenta  $p_{x_i} = n \dot{x}_i$  to get:

$$\begin{aligned} \dot{x} &= \frac{p_x}{n} & \dot{p}_x &= \frac{\partial n}{\partial x} \\ \dot{z} &= \frac{p_z}{n} & \dot{p}_z &= \frac{\partial n}{\partial z} \end{aligned}$$

For the index of refraction profile an exponential profile is used

$$n(z) = A + B e^{-Cz}$$

To include perturbations: add a differentiable function to this over a range of  $z$



# RAYTRACING WITH FERMAT'S PRINCIPLE

## WHAT INFORMATION DO WE WANT?

Assume a simple line model for the cascade, and define emission times:

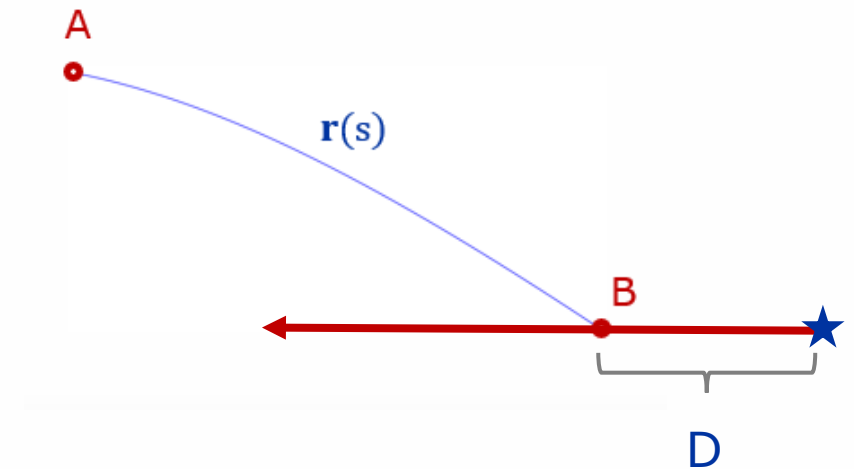
$$t' = \frac{D}{c} \longrightarrow \text{Distance from point } t' = 0$$

For each ray, use the optical path length  $L$  to get travel time:

$$\Delta t = \frac{L}{c}$$

Look at the relation between emission and arrival times:

$$t = t' + \frac{L}{c}$$



# WHAT CAN BE LEARNED FROM ARRIVAL AND EMISSION TIMES?

## THE BOOSTFACTOR

The boostfactor  $\frac{1}{\left(\frac{dt}{dt'}\right)}$  contains the information for geometric boosting.

Look at the uniform case ( $n$  constant):

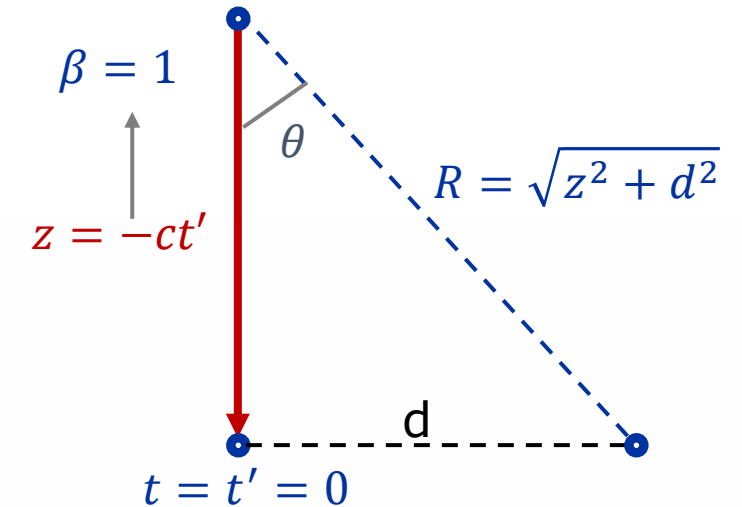
$$t = t' + n \frac{\sqrt{(-ct')^2 + d^2}}{c}$$

$$\frac{dt}{dt'} = 1 - n \cdot \cos(\theta)$$

What if the derivative is zero?

$$\frac{dt}{dt'} = 0 \Rightarrow \cos(\theta) = \frac{1}{n} \Rightarrow \theta = \theta_{\text{cerenkov}}$$

A signal emitted over an interval  $\Delta t'$  arrives in a shorter interval  $\Delta t$



## WHY DOES THE BOOSTFACTOR MATTER?

The end point formalism ([arxiv.org/abs/1112.2126](https://arxiv.org/abs/1112.2126)) :

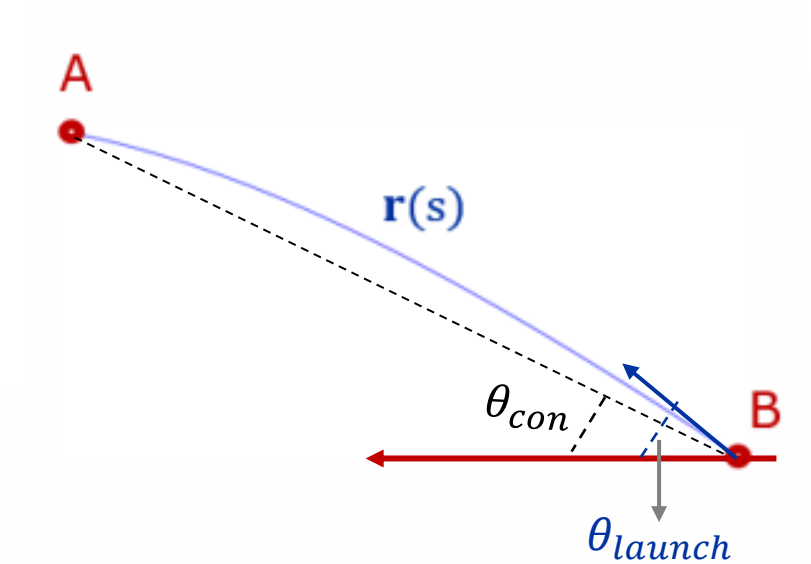
$$\vec{E}_{\pm}(\vec{x}, t) = \pm \frac{1}{\Delta t} \frac{q}{c} \left( \frac{\hat{r} \times [\hat{r} \times \vec{\beta}^*]}{\underbrace{(1 - n\vec{\beta}^* \cdot \hat{r})}_{\text{Boostfactor}^{-1}}} R \right)$$

When calculating as  $1 - n\beta \cos(\theta)$ :

What  $n$ ?

What  $\theta$ ?

Previous studies (A. Timmermans, Ba. Thesis) show that a straight line approximation might not be valid for very inclined geometries in air



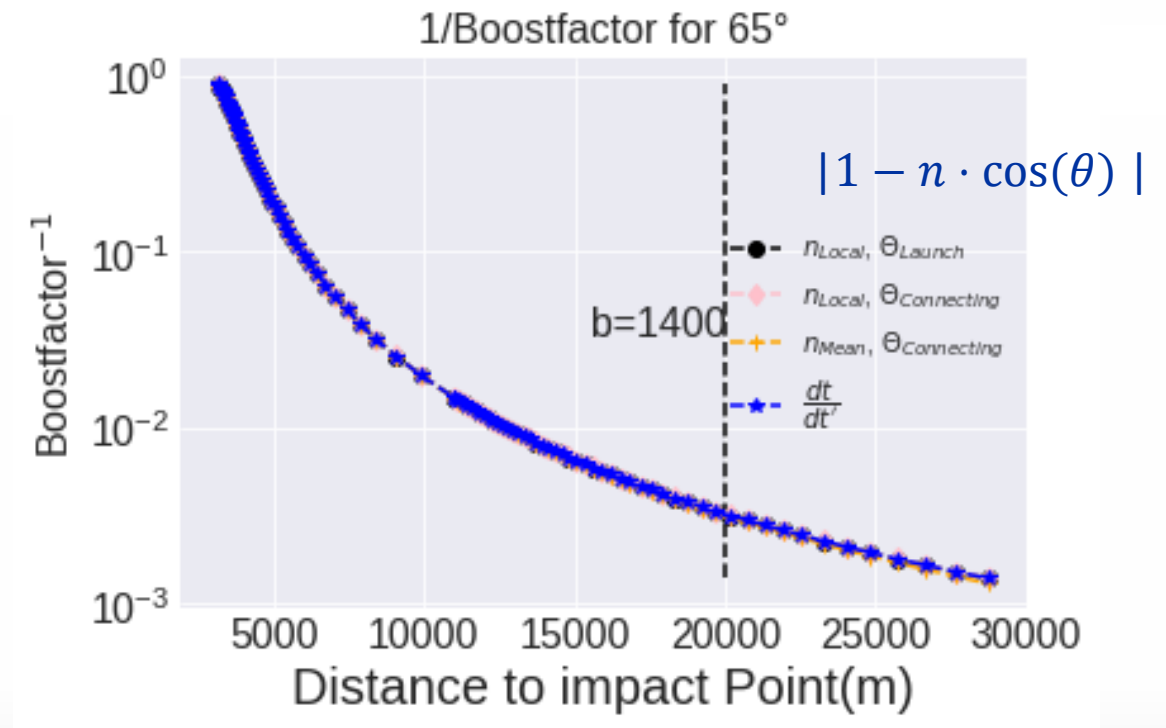
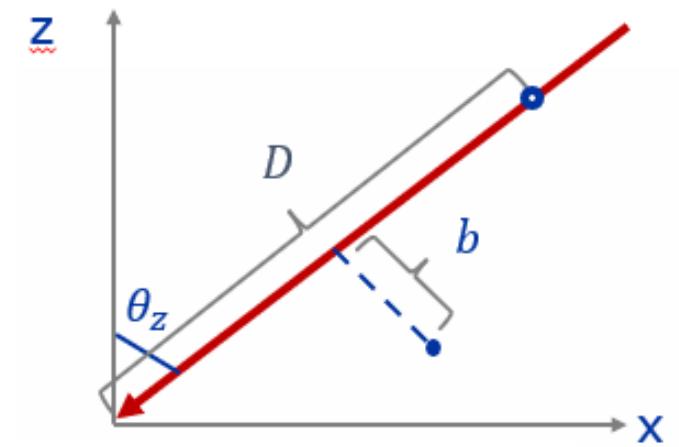
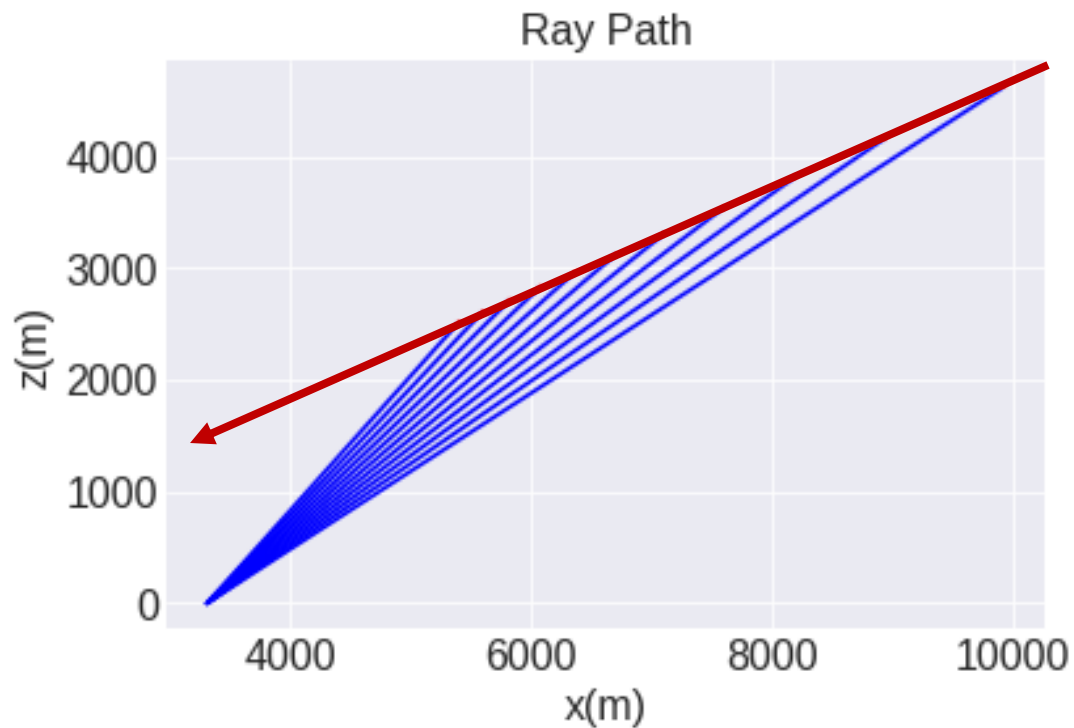


# IN AIR BURSTS

## WHERE DOES THE APPROXIMATION CHANGE?

Look at estimators and compare to numerically computed derivative

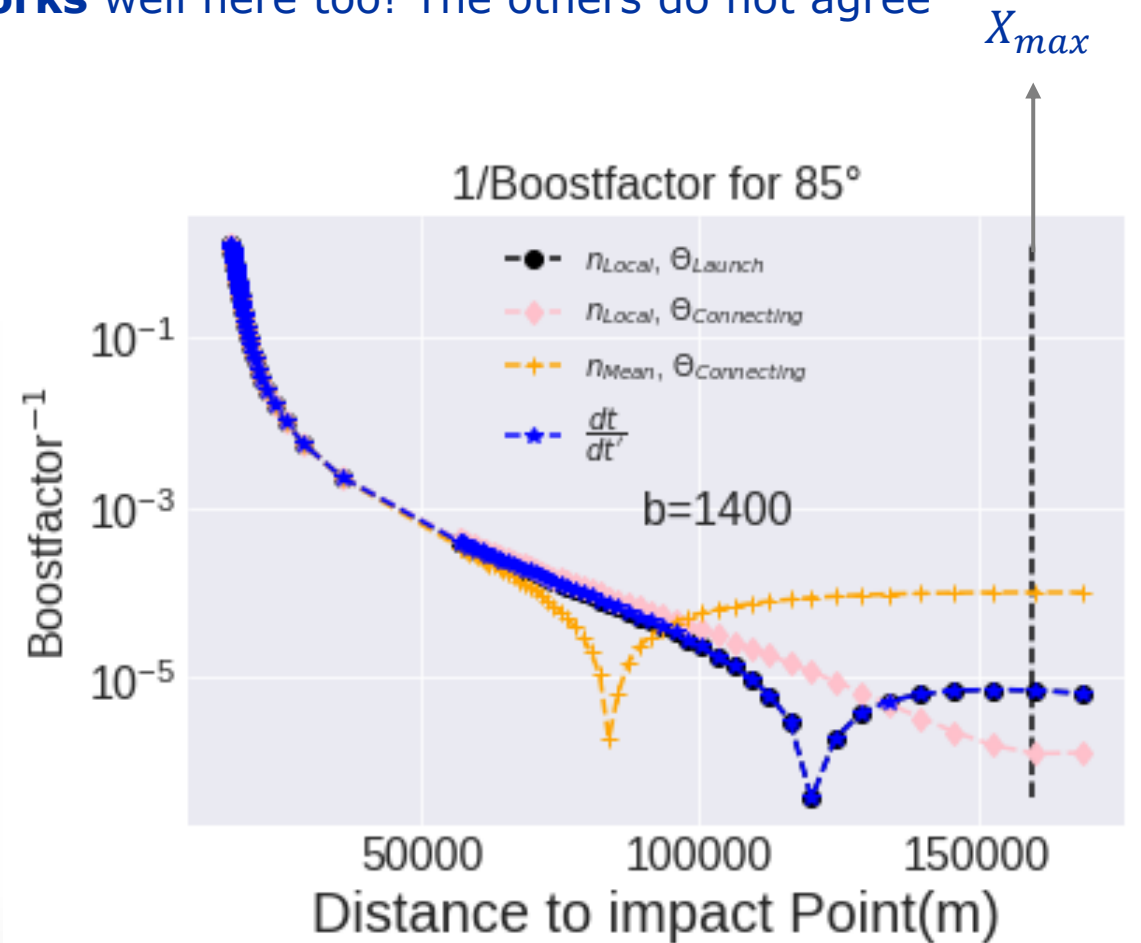
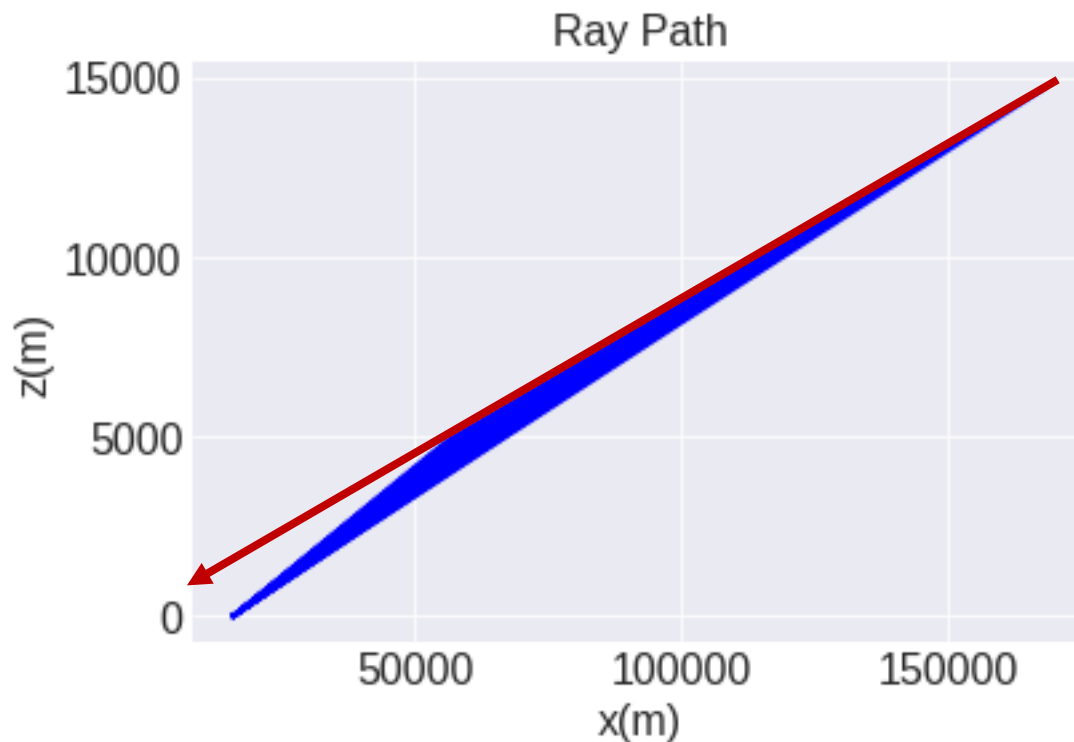
**All estimators follow the derivative**



# IN AIR BURSTS

## WHAT ABOUT INCLINED SHOWERS?

The estimator with **local n and launch angle works** well here too! The others do not agree  
Similar results found by A.Timmermans



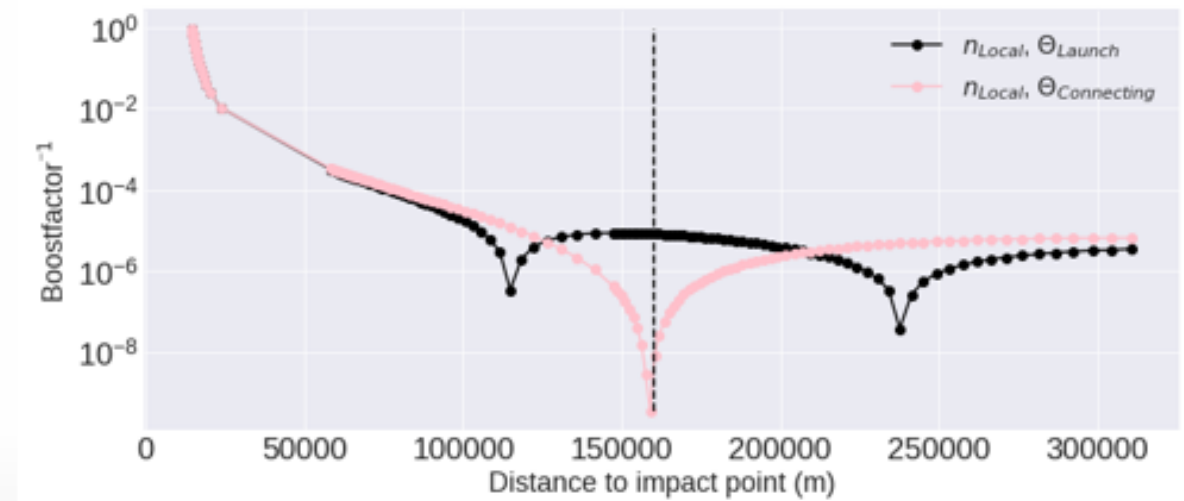
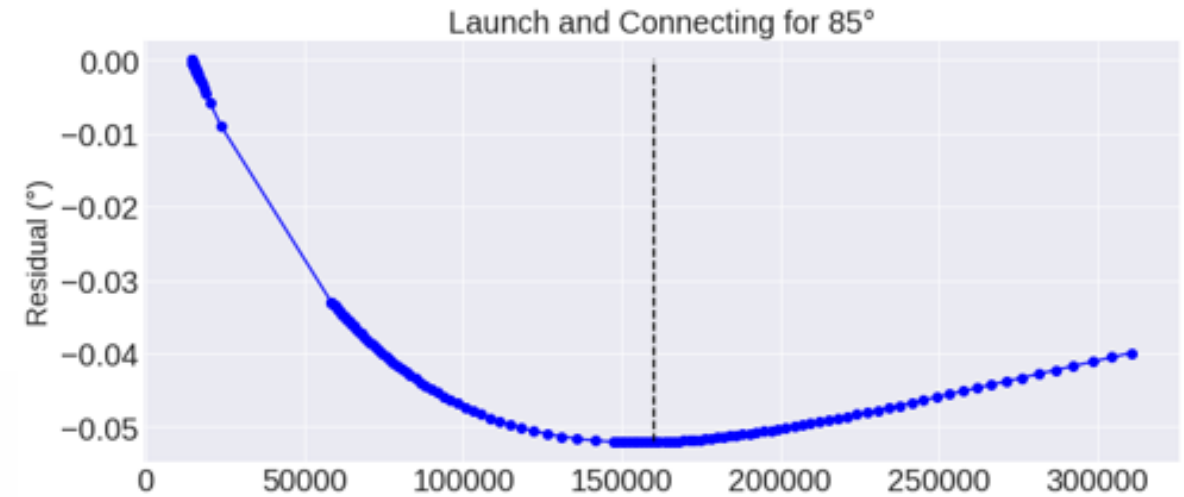
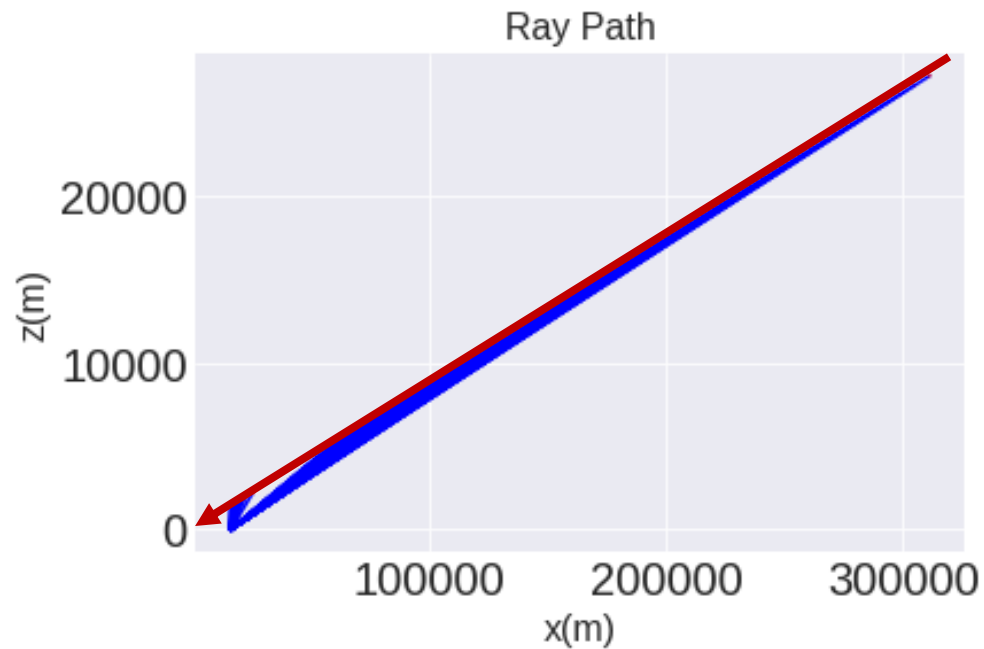
# IN AIR BURSTS

## HOW BAD IS IT REALLY?

This boostfactor is for particles with momentum along the cascade axis

Momentum in a real cascade is distributed

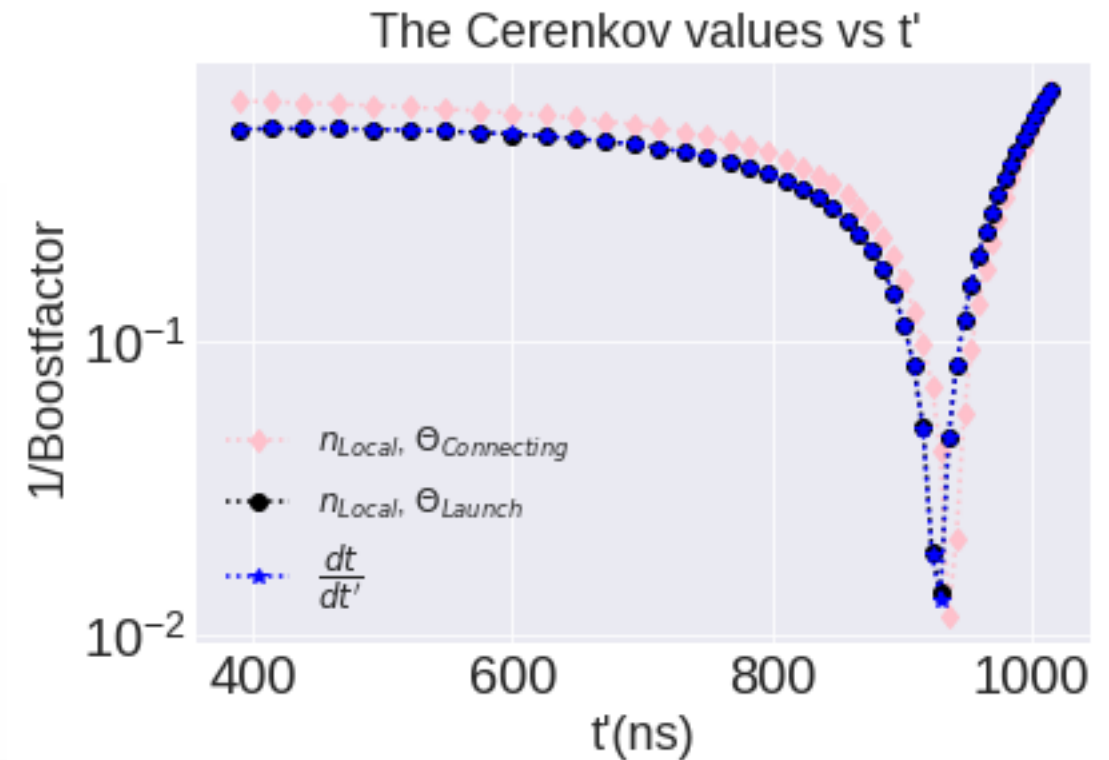
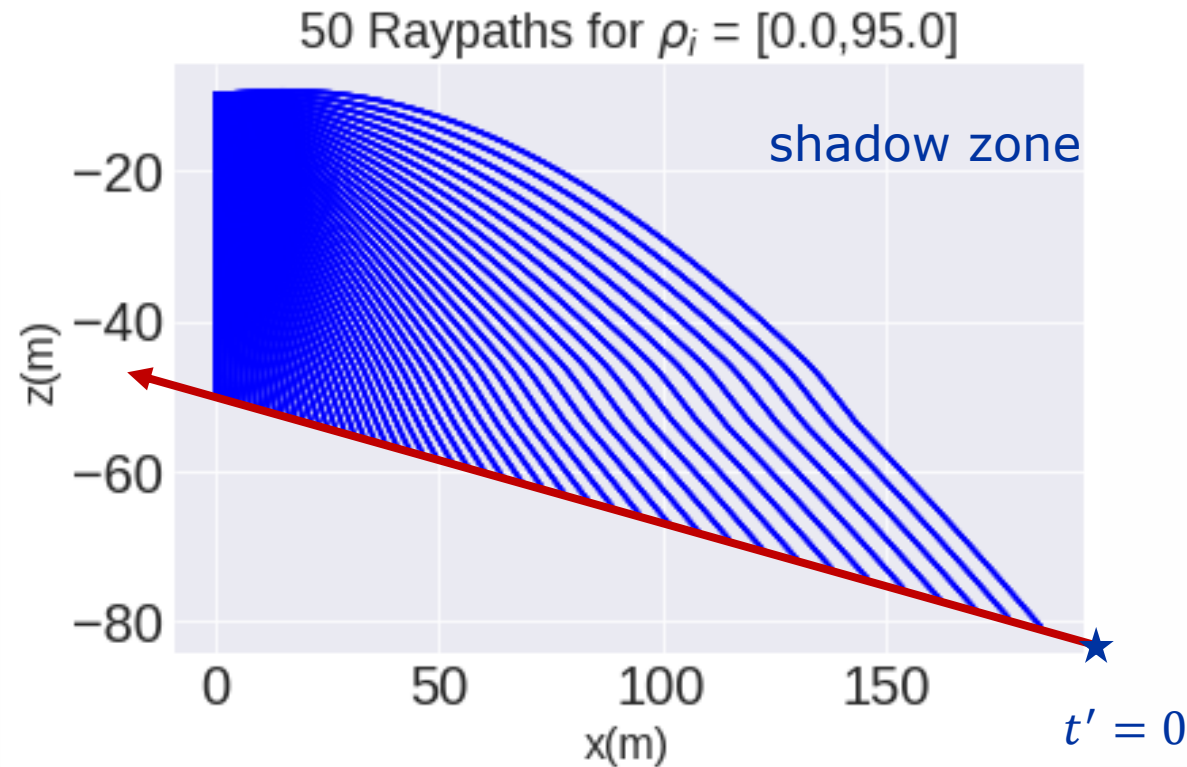
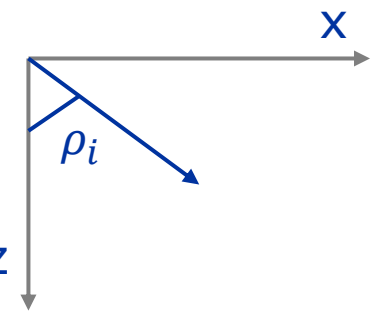
⇒ Non trivial result for realistic radio simulation



# HOW ABOUT IN ICE?

## DOES THE ESTIMATION HOLD HERE?

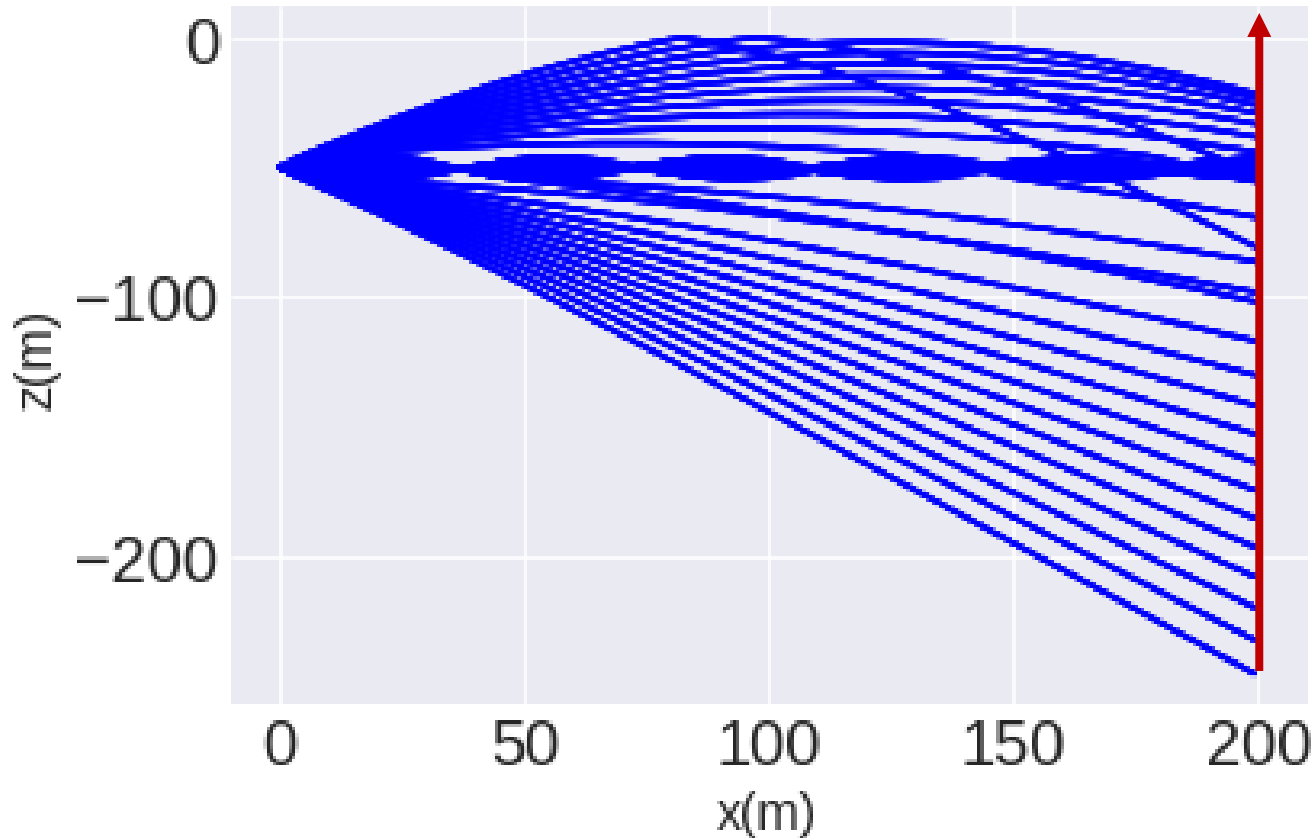
$1 - n_{local} \cdot (\cos(\theta_{launch})) = \frac{dt}{dt'}$  seems to work here too, even with perturbation



# HOW ABOUT IN ICE?

## CAN AN OVERDENSITY GUIDE THE RAYPATHS?

50 Raypaths for  $\rho_i = [45.0, 135.0]$



Add a cosine window at -50 m depth (peak strength 0.15)

=> Raypaths get trapped in the perturbation

What does this mean for the shadow zone?

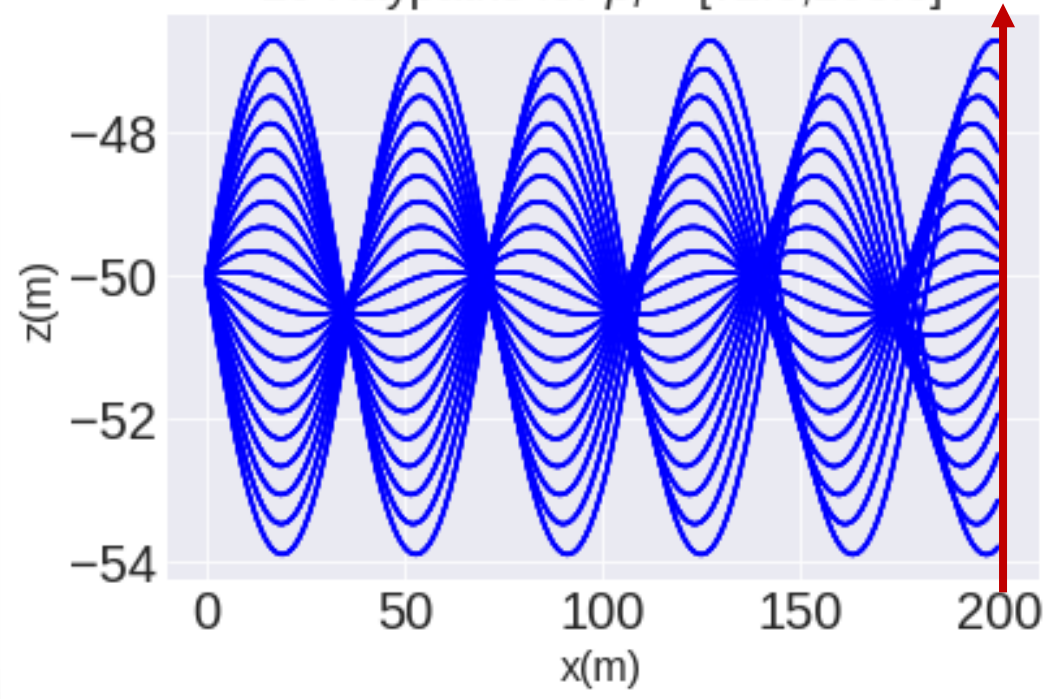
# HOW ABOUT IN ICE?

## IS THE BOOSTFACTOR STILL RELEVANT?

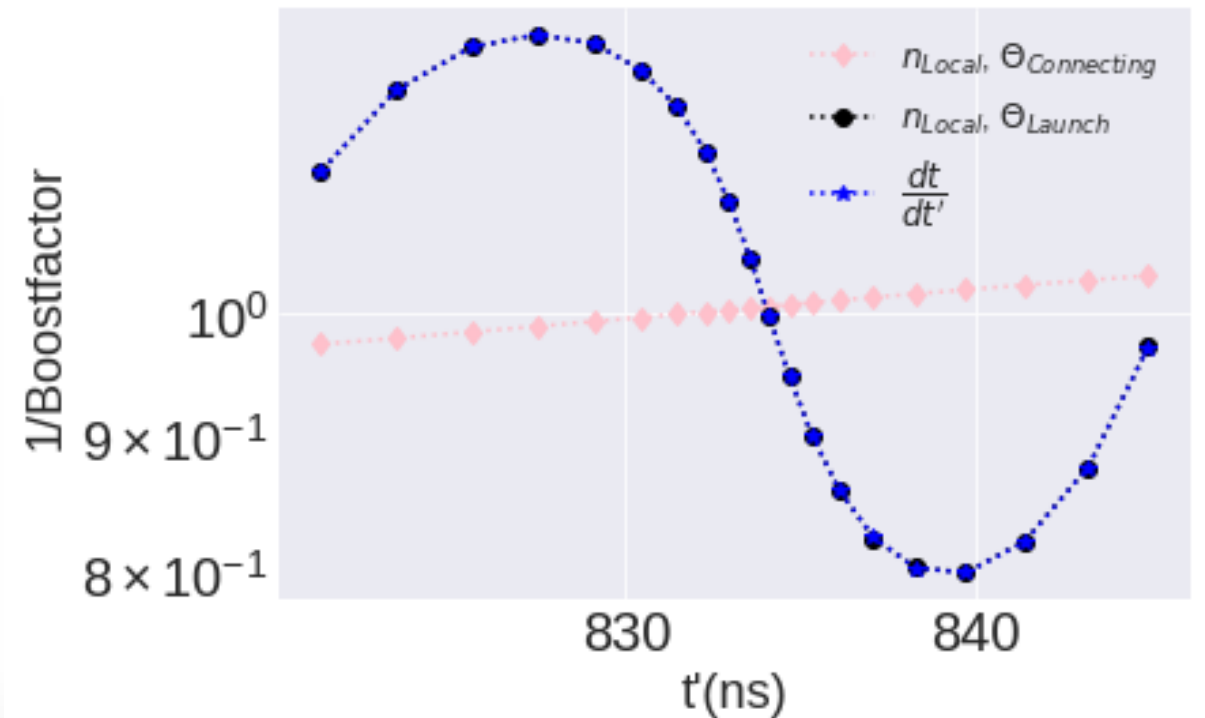
Look at the rays that remain trapped in the overdensity, can we use the same estimator here?

$$1 - n_{local} \cdot (\cos(\theta_{launch})) = \frac{dt}{dt'} \text{ holds}$$

20 Raypaths for  $\rho_i = [72.0, 108.0]$



The Cerenkov values vs  $t'$



## WHAT ARE THE MAIN RESULTS? HOW TO PROCEED?

When describing the boostfactor, using a Cerenkov **estimator with local n and launch angle works (for now)**

The raytracer finds **waveguide like behavior for overdensities**

Waveguide propagation is typically frequency dependent => needs to be taken into account

What effect does launch vs connecting have? => Simulate worst case scenarios