The role of precision at the high-energy frontier Gordon Research Conference "Particle Physics" Hong Kong University of Science and Technology June 30 - July 4, 2019

Eric Laenen



## **Outline**

- ✦ Precision, accuracy, errors and uncertainty…
- ✦ …for physics at the (HL-)LHC
- ✦ …for physics at future future colliders
- ✦ Some prospects for new methods and tools towards yet further precision

Recommended: inputs to and talks at Granada Open Symposium

Apologies: credits in talk will be vastly inadequate

### Theoretical colliders

- **Hadron collider** 
	- ‣ transformed into "parton collider" via parton distribution functions



 $d\sigma_{\mathrm{H_1H_2}}(\lbrace X_n \rbrace) = \sum$  $i,j$  $\int_0^1$  $\xi_{1,\rm min}$  $d\xi_1$  $\int_0^1$  $\xi_{2,\rm min}$  $d\xi_2 f_{i/H_1}(\xi_1) f_{j/H_2}(\xi_2) d\hat{\sigma}_{ij}(\xi_1, \xi_2, \{X_n\})$ ,

- ‣ compute partonic cross section in perturbation theory
- ‣ infer (i.e. download) pdf
- Lepton collider: can do (partly) without pdf's

### Collider physics



➨ Optimize precision of inputs

### A lot of LHC data still to come



At present we only have about 190/3000  $\sqrt{6.5\%}$  of data yield after HL-LHC

### Precision, accuracy, error and uncertainty

✦ A bit of terminology: for predictions for observable O

$$
O^{[m]} = \sum_{n=1}^{m} c_n \alpha^n + \delta O^{[m]}
$$

- $\triangleright$  Precision: compute to order "m", large enough for uncertainty  $\delta O^{[m]}$  to be small enough
- $\triangleright$  But beware: it can a be small uncertainty on an incorrect result. It is then precise, but not accurate
- ▶ Errors: a measure of accuracy
	- ✓ experimental: statistical and systematical
- Uncertainty: indicates range in which true value could lie
- ← Confront prediction with measurement, all the more meaningful with small  $\delta O^{[m]}$
- ✦ This is what we *should* be doing: a highly sophisticated instance of The Scientific Method

### Example of precision vs accuracy

perturbation theory CDF Run 1 rapidity distribution of Z boson vs



Anastasiou, Dixon, Melnikov, Petriello '03

## Purpose of precision: To Measure and Explore

- Aside from exceptional moments in the development of the field, research is not about proving a theory is right or wrong, it's about finding out how things work
- We do not measure Higgs couplings precisely to **find** deviations from the SM. We measure them to **know** them!

Michelangelo Mangano at SM@LHC '19, Zurich

### Precision for SM and BSM  $\blacksquare$  **DODATA)** DODATA PRODUCED DODATA NALISE PRODUCED DODAT – [njet = 0] = 1.76 <sup>±</sup> 0.03 <sup>±</sup> 0.22 pb (data) NNLO (theory) 4.6 PRD 87, 112003 (2013) W!`⌫ = 2.77 <sup>±</sup> 0.03 <sup>±</sup> 0.36 pb (data) NNLO (theory) 4.6 PRD 87, 112003 (2013) **ATLAS** Preliminary = 44 + 3.2 4.2 pb (data) <sup>2</sup>NNLO (theory) 4.9 JHEP 01, 086 (2013)  $\mathbf{M}$

 $\mathcal{L}$  of  $\mathcal{L}$  ,  $\mathcal{L}$  of  $\mathcal{L}$  or  $\mathcal{L}$  ,  $\mathcal{L}$ 

= 1.189 <sup>±</sup> 0.009 + 0.073 0.067 pb (data) NNLO (theory) 20.3 PRD 93, 112002 (2016)

= 16.82 <sup>±</sup> 0.07 + 0.75 0.78 pb (data) <sup>2</sup>NNLO + CT10 (theory) 20.2 PRD 95 (2017) 112005

- **Falsification** 
	- ‣ Compute promising SM observables to high precision for easier falsification by data  $\frac{1}{2}$ In the 1.37 theorys of the third with the 1.31 theory. The complete the complete theory is completed to the co<br>In the 1.4 of 1.6 June 1.6 Ju<br>  $\alpha = 2.20$  and  $\alpha = 45$  for  $\alpha = 20$  are  $\alpha = 20$ . Army 20.2 army 20.2



- **Verification**  $\tt t$ ion $\tt t$ 
	- **Example FIGM influenced on selected observables to high precision**  $=$  374  $\pm$  7  $\pm$ mpute PCM influenced an eclected cheer vables to bigh procision NNLO QCD  $\frac{1}{4}$   $\$  $\frac{23.2}{23.2}$  and  $\frac{23.2}{20.3}$  fb (data)  $\frac{23.2}{20.3}$  provided it ected observables to high precision and sherpa (NLC) sherpa (NLC) (NLC)  $\sim$ ZZ = 6.7 <sup>±</sup> 0.7 + 0.5 0.4 pb (data) NNLO (theory) 4.6 JHEP 03, 128 (2013)
		- $\vee$  to ensure that the unique signatures are robust on HO corrections  $\sim$  onder even the employ digitate by any research in the corrections ) ensure that the unique signatures are robust on HO corrections  $\hspace{0.1cm}$  $\sim$  262.3  $\mu$  12.3  $\mu$  12.3  $\mu$  26.3  $\mu$  23.1  $\mu$  23.1  $\mu$  23.1  $\mu$  23.1  $\mu$  24.1  $\mu$ number of the distribution of the 12.7 and 1.8 fb (data) powerful theory of the control of the control of the c external and the 9.7  $\pm$  1.5  $\pm$
		- √ to extract information (measurement or exclusion) and the extract information (measurement or exclusion)  $\mathbf{r}_1$   $\mathbf{r}_2$   $\mathbf{r}_3$   $\mathbf{r}_4$   $\mathbf{r}_5$   $\mathbf{r}_6$   $\mathbf{r}_7$  (theory) 20.3 PLB 753, 552-572 (2016) 20.3 PLB 753, 552-572 (201

### Precision for (HL-)LHC

# So far excellent predictivity of SM at LHC



11

### A core process at (HL-)LHC: Drell-Yan International March 2017



 $\div$  Sub percent level experimental error for Drell-Yan  $p_T$  spectrum, and other distributions **Percent level experimental error for Drell-Yan p<sub>r</sub> spectrum and other** Toon Tovol Oxponi *•* Sensitive to SM parameters and proton PDFs!

*•* Studying lepton production in proton-proton collisions.

✦ Impact on W-mass, PDF-fits etc *•* Welliass, PUF-IIIs etc. In the candred bosons used as star as star dependence of the candred candred. Lorenzo Biandard Candres. Lorenzo Biandard Candres. Lorenzo Biandard Candres. Lorenzo Biandard Candres. Lorenzo Bian

Lorenzo Bianchini SM@LHC '19



### Drell-Yan @ (HL-)LHC Predictions at LHC8 (Z)<br>Predictions at August 2019

- $\leftarrow$  Theory challenged!
- ← NNLO + N3LL better than NNLO alone
	- ‣ NLO + resummation not sufficient **I** Scale uncertainties below the 5% level uncertainties below the 5% level uncertainties below the 5% level uncertainties of  $\mathcal{L}_\mathbf{z}$
- ← At this level many small effects must be assessed
	- ‣ N3LO + N4LL?
	- PDF uncertainties at 1%?
	- non-perturbative effects at small  $p_T$
	- QED corrections (1-2%)
	- α<sub>s</sub> uncertainty
- HL-LHC data will be much more challenging

![](_page_12_Figure_11.jpeg)

Bizon, Chen, Gehrmann-de Ridder, Gehrmann, Glover, Huss, Monni, Re, Rottoli, Torrieli '18

## Top quark pairs at (HL)-LHC

✦ We are already well into precision top quark physics era • We are well in the hadron collider precision measurement territory !!! aireauy wen muo preci

![](_page_13_Figure_2.jpeg)

- + all-QCD process, NNLO corrections for cross sections.. Czakon, **process, NNLO corrections for cross sections..** Czakon, Fiedler, Mitov '13
- ← ..and for differential distributions Czakon, Heymes, Mitov '15,'1 8 TeV 256 20 5 x 106 Czakon, Heymes, Mitov '15,'16
- good enough now to be input for PDF fits

# Precision for top quark pair production

Value of higher orders for precision, and uncertainty budget

![](_page_14_Figure_2.jpeg)

 $v \geq 11$  $ev1$ *Concurrent uncertainties:* Scales  $\sim$  3% pdf (at  $68\%$ cl) **~ 2-3%**  $\alpha_{\rm s}$  (parametric)  $\sim 1.5\%$  $m_{\text{top}}$  (parametric)  $\sim 3\%$ Soft gluon resummation makes a difference:  $5\% \rightarrow 3\%$ *MC, Fiedler, Mitov `13* <sup>6</sup>

Czakon, Fiedler, Mitov '13

- improvement due to higher orders and resummation clear
- all at the few percent level
- ✦ NLO EW also known

Czakon, Heymes, Mitov, Pagani, Tsinikos, Zaro '17

impact of photon-in-proton distribution notable

# Theory for top pairs plus more

Status of precision theory description in 3D

Markus Schulze at LHCP '18

![](_page_15_Figure_3.jpeg)

- ✦ Again, smaller effects come to the fore when precision is high 2/13
	- ‣ narrow width approximation vs. full off-shell decay, all this at higher order
	- m<sub>top</sub> definition and value a guaranteed topic for lively debates
- ✦ Also here the experimental accuracies will be challenging theory

### Higgs production at (HL-)LHC *igara Higgs* pro  $\overline{1}$ ✓↵*s*(*µr*) ⇡ ◆`  $\int_{0}^{1}$   $\int$  $\mathbf{r}$ after. We note, though, that we observe a very small, but systematic, increase of the expansion in the range

+ Status of Higgs production mechanisms vs theory: theory (just) ahead for now ahead for now

54

![](_page_16_Figure_2.jpeg)

expansion at  $\frac{1}{2}$  and  $\frac{1}{2}$  and  $\frac{1}{2}$  are systematic increase original expansion of  $\frac{1}{2}$  and  $\frac{1}{2}$  are systematic increase original expansion of  $\frac{1}{2}$  and  $\frac{1}{2}$  are systematic increase origina from values of the participants of the part of the p<br>The part of the part of t Ig OILZJN COLES

Anastasiou, Duhr, Dulat, Herzog, Mistlberger '15

![](_page_16_Figure_5.jpeg)

bilises starting from *N* = 4, leaving a negligible trun-

← Calculation was done in "1-z" = soft expansion

$$
\hat{\sigma}_{ij}^{(3,N)} = \delta_{ig} \,\delta_{jg} \,\hat{\sigma}_{SV}^{(3)} + \sum_{n=0}^{N} c_{ij}^{(n)} \,(1-z)^n
$$

 $\rightarrow$  to N=37.. The including substantial pileup. As an example, measurements of  $N=37$ . reached half-percent precision (aside from an overall luminosity uncertainty) for a range of kine-

"

**EXECUTE: That is a can expect that substantial full analytic result also available in the substantial functions in the needed in the needed in the substantial functions in the needed in the substantial functions in the su** the limit and you to be a wide a wallance range  $\mu$  Mistlberger '18

### Higgs boson studies at future colliders J. de Blas et al, arXiv: 1905.03764

← In general, Higgs uncertainties at HL-LHC = 1/2 those of LHC

to the current precision  $\mathbf{r}_{\text{max}}$  for the future are taken from ref. [117]. For the future uncertainties, the fu

 $\leftarrow$  FCC: well below 1% uncertainties arise either from intrinsic limitations in the theoretical calculation (ThIntr) and parametric uncertainties (ThPar).

F. Caola @ Granada

![](_page_17_Picture_539.jpeg)

ggF cross section: many sources of small uncertainties in the cross section. many sources of small uncertainties in the case of small uncertainty of the cross section. be beaten down

![](_page_17_Figure_6.jpeg)

# Status of QCD corrections for LHC

- The theory community has responded to the precision challenge very impressively in the last 20 years
- ✦ Started in 2005 with highly ambitious (at the time) Les Houches **wishlist** for NLO calculations. All done by 2011
- ← This lead to NLO revolution
- ← New frontier: NNLO and even N3LO
	- ‣ many new methods have been developed: healthy marketplace
- ✦ Resummation, parton showers
	- ‣ much progress here in each (higher precision), and their combination
- ✦ Also much improved:
	- ‣ PDFs, heavy quark masses, computing methods

### The NLO "revolution"

- Clever new methods have led to a breakthrough in NLO calculations. Particular the calculation of the oneloop diagrams has been "solved" in full generality, and has been automatized.
	- ▶ Results now in codes such as aMC@NLO, Powheg Box, MCFM,... *d*⇤*pp<sup>X</sup> d*<sup>3</sup>*p*<sup>1</sup> *. . . d*<sup>3</sup>*p<sup>n</sup>* IC@NL ⌅ *Fowheg Box, MCFM,...*
- ✦ Basic notion: all one-loop amplitudes can written as a sum of boxes, triangle, bubbles and tadpoles s, triang

$$
\sum_{i=1}^{n} \sum_{j=1}^{n} a_j \sqrt{1 + \sum_i b_i} + \sum_i c_i \sqrt{1 + \sum_i c_i}
$$

✓ In essence because we live in 4 dimensions: every vector can be decomposed in at maximum four independent vectors Vermaseren, van Neerven; Bern, Dixon, Kosower,...

$$
\mathcal{M} = \sum_{i} a_i(D) \text{Boxes}_i + \sum_{i} b_i(D) \text{Triangles}_i + \sum_{i} c_i(D) \text{Bubbles}_i + \sum_{i} d_i(D) \text{Tadpoles}_i
$$

- ✦ Job: find coefficients
	- **Example 2** can use (generalized) unitarity: determined them from cuts and poles

![](_page_19_Figure_9.jpeg)

### Status of higher order calculations in QCD

![](_page_20_Figure_1.jpeg)

- ‣ LO well-understood, now more efficient than ever
- ▶ NLO: automatized, a flood of results
- ‣ NNLO: top quark production (single and pair), dijet production, 1-jet inclusive, …. *p i* or *i esults* the *ion (cinels and pair)* direction
- $\triangleright$  NNNLO: Higgs production,  $F_2(x,Q)$ ,  $\overline{\mathcal{A}}$  and  $\overline{\mathcal{A}}$  and  $\overline{\mathcal{A}}$  are production *nf A*[1*/*2]

## Automatic higher order calculations

- **QCD NLO is automatized** 
	- ‣ no limitations in principle, but high multiplicity means longer running
	- including matching to parton showers: aMC@NLO in MadGraph5 framework

Allwal, Frederix, Frixione, Hirschi, Maltoni, Mattelaer, Shao, Torrieli, Zaro '14 **approaching 4000 citations..**

Frederix, Frixione, Hirschi, Maltoni, Pagani, Shao, Zaro '18

- NLO EW now also included
- ✦ POWHEG Box for NLO + PS

Nason, Oleari; + Frxione, Aioli, Re '04 ff

- general framework for NLO + PS
- **HELAC-NLO**

Bevilacqua, Czakon, Garzelli, van Hameren, Kardos, Malamos Papadopoulos, Pitta, Worek, Shao '14

Codes increasingly incorportated into exp'tl frameworks

**High Energy Physics** Na<br>Register Tools Enthrance Chairs<br>Backs

### Generate processes online using MadGraph5\_aMC@NLO

It improve our web services we sequest that you register. Registert on is quick and free. You may register for a password by allowing here.<br>Please note the connect coherence for MadKingdof, aMC:0NI.0, ar XIrc1405.0201 Bag-

![](_page_21_Picture_102.jpeg)

![](_page_21_Picture_103.jpeg)

### **NNLO QCD and NLO EW Les Houches Wishlist**

Wishlist part 1 - Higgs (V=W,Z)

![](_page_22_Picture_11.jpeg)

Wishlist part 2 - jets and heavy quarks

![](_page_22_Picture_12.jpeg)

# Precision for BSM

- Top down BSM
	- assume new physics model, or a simplified version, compute signals including (QCD) corrections
	- ‣ much work here has been done for MSSM, compositie models etc
	- Example: NLO + NNLL resummed for squark-gluino production  $\frac{1}{10^2}$ 
		- ✓ includes threshold and Coulomb corrections
	- ‣ improves limit-setting for gluino masses etc.

Beenakker, Borchesnky, Kraemer, Kulesza, EL '16

Beneke, Piclum, Schwinn, Wever, '16

- Bottom-up BSM
	- be agnostic about new physics, parametrize it as effective theory

![](_page_23_Figure_11.jpeg)

![](_page_24_Figure_0.jpeg)

### Precision for other future colliders

### Precision of α<sup>s</sup> Fig. 9.2. We recall that the set are extractions which are extractions which are based from extractions which are based fr on  $P$  recision of  $\alpha_{\rm c}$

experimental as well as theoretical issues. The six pre-averages are summarized in

![](_page_26_Figure_1.jpeg)

![](_page_27_Picture_0.jpeg)

![](_page_27_Figure_1.jpeg)

Top mass exp't error: 25 MeV expected Theory one larger at present.

![](_page_27_Figure_3.jpeg)

### Electroweak corrections

- QED corrections are also becoming quite important at the LHC
	- ‣ NNLO in QCD = NLO in EW
- Manohar, Nason, Salam, Zanderighi '17
- ‣ Photons in protons: LUXqed formalism
	- $\checkmark$  Extract from precise ep data  $\to$  large increase in precision! (% level)

 $L^{\mu\nu}W_{\mu\nu}$  (DIS) =  $\hat{\sigma}_{e\gamma} \otimes f_{\gamma/p}$  ( $\gamma$ PDF)

- EW loop corrections: many scales
	- ‣ particular relevant for EW precision observables
	- at large  $p_T$  in hadron colliders

![](_page_28_Figure_11.jpeg)

 $10^{-4}$ 

0.2 0.4 0.6 0.8 1 1.2 1.4 1.6 1.8 2

![](_page_28_Figure_12.jpeg)

### High precision at lepton collider: Higgs couplings

### Improvements w.r.t. HL-LHC

![](_page_29_Figure_2.jpeg)

### Some prospects for more precision

### News from PDFs: theory errors  $\text{Var}$  is the manufacture  $\text{Var}$ process, *µ<sup>f</sup>* is fully correlated among processes. This assumption is based on the observation that *µ<sup>f</sup>* variations

are universal (process-independent). In the appendix we have a set of th

✦ PDF errors dominant for many LHC predictions. Thus far they only reflect the uncertainties of measurements and their correlations spectral to the central theory prediction for the *i-t*h crossdominant for many LHC predictions. Thus far they only reflect the provide expressions for *Sij* in the 3- and 9-point cases.

 $T$ op Atlas, c $T$ erential c $\sim$ di $T$ erential c $\sim$ 

← First step in including uncertainties from Missing Higher Orders, by fitting PDF's for a set of  $\mu_R$ ,  $\mu_F$  values, and from there sum the exp.("C") and th. ("S") covariance matrix  $u = 1/2$  $s, \mu_F$  values, and from there sum the exp.("C") and th. ("S") covariance *N* dat the lower kinematic cut has been increased from *Q*<sup>2</sup>

32

$$
\chi^2 = \sum_{i,j=1}^{N_{\text{dat}}} \left( D_i - T_i^{(0)} \right) \left( S + C \right)_{ij}^{-1} \left( D_j - T_j^{(0)} \right)
$$

← Introduces much more correlations between experimental inputs The theory covariance matrix *Sij* has been constructed QCD-induced cross-talk

![](_page_31_Figure_6.jpeg)

where (*k*)

*<sup>i</sup>* <sup>=</sup> *<sup>T</sup>*(*k*)

*<sup>i</sup> <sup>T</sup>*(0)

![](_page_31_Figure_7.jpeg)

## News from PDFs: theory errors

- Validation: see at NLO if the NNLO central value is in the MHO uncertainty band
- Looks ok:

![](_page_32_Figure_3.jpeg)

### QCD precision at REALLY high order I more convenient for extending analyses of the strong coupling constant into the non-perturbative *R*∗ *and five-loop calculations* Ben Ruijl

the values  $\overline{\phantom{a}}$  ,  $\overline{\phantom{a}}$  ,

- ✦ **five**-loop QCD beta function  $r_{\rm c}$  regime, using the Force calculations of four-loop vertex functions in ref. [22]. For the gauge-loop vertex functions in ref. [22]. For the gauge-loop vertex functions in ref. [22]. For the gauge-loop vertex funct  $\epsilon$ -roop general ref. (B.4) of ref.  $\ell$ 
	- $\triangleright$  first done for SU(3), then for SU(N) Baikov, Chetyrkin, Külm 16
		- ✓ using R\* method to extract divergences include  $\Gamma$  and  $\Gamma$  and  $\Gamma$  is the same was observed for the moment of the moment of  $\Gamma$  $\sqrt{2}$   $\frac{1}{2}$  and  $\frac{1}{2}$  is the induction value of the  $\frac{1}{2}$  and  $\frac{1}{2}$
		- $\checkmark$  6 days on 32 core machine  $\checkmark$
	- ive loop expansion in MSbar scheme very benign ∕

 $\beta(\alpha_s, n_f = 4) = 1 + 0.490197 \alpha_s + 0.308790 \alpha_s^2 + 0.485901 \alpha_s^3 + 0.280601 \alpha_s^4 + \dots$  $\overline{a}$ 

V less than 1% change due to 5-loop term, even at α<sub>s</sub>=0.47 rm, even at  $\alpha_{\rm s}$ =0.47

 $\leftarrow$  implications for running coupling:  $\leftarrow$  <sup>1.06</sup>  $\leftarrow$  11000 contributed (N<sub>1</sub>

![](_page_33_Figure_9.jpeg)

Baikov, Chetyrkin, Kühn '16

Herzog, Ruijl, Ueda, Vermaseren, Vogt '17,

gante v, επειγρατή, παρτιστό<br>Herzog, Ruiil, Ueda, Vermaseren, Vogt '17,

 $\frac{1}{\sqrt{2}}$ 

 $\overbrace{\phantom{aaaaa}}^{x}$ 

<sup>s</sup> +0.080921α<sup>4</sup>

### QCD precision at REALLY high order II *Oep* = *fi* ⊗ *c*<sup>o</sup> *<sup>i</sup>* , *Opp* = *fi* ⊗ *fk* ⊗ *c*<sup>o</sup> *ik* (1.1) in terms of the respective partonic cross sections (coefficient functions) *c*<sup>o</sup> and the universal parton *LLY* high *dy*  $\bf{M}$ \$ *x* large contributions of *d* (4) *AA* , which enter at N4LO for the first time, the series would look much e entire benefits the best with constant and non-the new process of 1.4 to 1.6 km security ratio and non-the N coefficients. This sizeably *d* (4) *AA* contribution (∼ *n*<sup>2</sup> *<sup>c</sup>* + 36 ) also implies that the leading large-*nc*

*FA* /*nc* <sup>=</sup> <sup>5</sup>/2 and *<sup>d</sup>* (4)

+ QCD splitting functions at four and five loops distribution functions (PDFs) *fi*(*x*,µ2) of the proton at a scale <sup>µ</sup> of the order of a physical scale.  $\frac{1}{2}$  and five loope operation-product expansion (OPE), and the component functions in the coefficient functions in the coeffic contribution a less good approximation at the provides and five leader. The numerical impact of the higher-order contributions to the splitting functions *P*<sup>±</sup>

with *d* (4)

their scale dependence is given by the renormalization-group evolution equations

‣

*P* = *a*<sup>s</sup> *P*(0) + *a*<sup>2</sup>

$$
\frac{\partial}{\partial \ln \mu^2} f_i(x, \mu^2) = \int_x^1 \frac{dy}{y} P_{ik}(y, \alpha_s(\mu^2)) f_k\left(\frac{x}{y}, \mu^2\right) \qquad P = a_s P^{(0)} + a_s^2 P^{(1)} + a_s^3 P^{(2)} + a_s^4 P^{(3)} + \dots
$$

at three loops (NNLO) known analytically; at 4 loops in part numerically, at 5 loops some moments  $\sim$  come moments in the light-cone operator-product expansion (OPE), and the coefficient functions in th expanded in powers of the strong coupling and strong coupling  $\overline{a}$ fown analytically, at 4 loops in part humencally, at 5 loops  $JNI$  () known analytically: at 4 loops in part numerically, at 5 loops  $\mathbf{r}$  full width of along about  $\mathbf{r}$ , the  $\mathbf{r}$  rections are about twice as large at a lower at a lower twice at a lower twice and  $\mathbf{r}$ *<sup>f</sup>* ) = 0.25 and *nf* = 3.

see refs. [1, 2] for the corresponding helicity-averaged and helicity-dependent splitting functions. Herzog, Moch, Ruijl, Ueda, Vermaseren, Vogt '18,

*AA* /*na* = 135/8 in QCD, see, e.g., app. C of ref. [32]: Without the rather

*<sup>r</sup>*=<sup>1</sup> (*qr* −*q*¯*r*) . (1.5)

eq. (1.1) can be expanded in powers of the strong coupling *a*<sup>s</sup> ≡ <sup>α</sup>s(µ2)/(4π),

![](_page_34_Figure_5.jpeg)

# Prospects for further QCD accuracy

- There is a "vibrant" community addressing NNLO for  $2 \rightarrow 3$ , N3LO and beyond
- ✦ Involves progress in
	- loop diagrams: analytical and numerical approaches
	- ▶ R divergence management, new subtraction mechanisms, phase space slicing making a comeback
	- $\triangleright$  Shuffle/Hopf algebra of polylogs to 3rd order  $\rightarrow$  elliptic integrals
	- ‣ threshold expansions
	- ‣ automation, computing methods

### Soft logarithms at next-to-leading power East considered in form to a result in form to Eq. (90) for the same space of the space of the space of the same solid in form to Eq. (90) for the same space of the space

✦ General soft expansion for 2→1 processes NLP (*p*1*, p*2*, k*) *M*(*ggg*) ć  $\int$ <sup>2</sup> $\int$ *p*<sup>1</sup> *· k p*<sup>2</sup> *· k*  $for 2 \rightarrow 1$  *processes* 

*s*ˆ

$$
\frac{d\sigma}{dz} = \sum_{n=0}^{\infty} \left(\frac{\alpha_s}{\pi}\right)^n \sum_{m=0}^{2n-1} \left\{ c_{nm}^{(-1)} \frac{\log^m(1-z)}{1-z} \bigg|_+ + c_{nm}^{(0)} \log^m(1-z) + \dots \right\}
$$

- ← NLP logarithm organization
- ‣ exhibit all-order patterns. Leading logarithmic resummation now achieved for a number of reacions and side is defined in the right-hand side is defined in the soft side is defined in the adjoint of the adjoint o  $\frac{1}{2}$ representation. One may then follow similar arguments to the extension of the Eq. (106), yielding to Eq. (106),  $\frac{1}{2}$

$$
\hat{\sigma}_{\text{LO}}^{(gg)}(Q^2) \exp\left\{\frac{2\alpha_s C_A}{\pi}\log^2(N)\right\} \left(1+\frac{2\alpha_s C_F}{\pi}\frac{\log N}{N}\right)
$$

Beneke, Broggio, Garny, Jaskiewicz, Vernazza, Szafron, Wang '18 Bahjat-Abbas, Bonocore, EL, Magnea, Sinninghe-**Damsté, Vernazza, White '19** Moult, Stewart, Vita, Xhu '18

- ‣ technology also used to extend phase space slicing methods for NNLO calculations consists of an elective consists of an election of an election of an election of an election of gluons, as shown a pair of  $\mathbb{R}^n$ A check of these results is that it reproduces known LP, and conjectured NLP results for Het large boson production, in the large top mass limit. As is well-known, the LO process limit in the LO process in Figure 7. Higher-order contributions near threshold have been discussed for example in
- ■ 33, and the formuch better numerical behavior

### Elliptic progress: from math for loops tegrals over kernels which are rational functions with at most simple poles. The most  $\Gamma$ lintians over  $\alpha$  functions  $\Gamma$ well-known examples are the classic polylogarithms Li*n*(*x*), of which the logarithm is a  $\Gamma$ lling is over a most simple poles.  $\Gamma$ well-known examples are the classic polylogic space and the logarithms Lincoln to Lincoln the logarithms Lincoln special case, constraint *y*2 = *P*4(*x*)=(*x a*1)(*x a*2)(*x a*3)(*x a*4) *.* (2.6)

 $\mathcal{P}_\mathcal{P}$  , we have polynogarithms are multi-valued functions defined records as iterated in  $\mathcal{P}_\mathcal{P}$ 

✦ Polylogarithms appear after doing loop integrals \* Polylogarithms appear after doing loop integrals  $\overline{\phantom{0}}$ *dx*0 *<sup>x</sup>*<sup>0</sup> Li*n*1(*<sup>x</sup>* elliptic curve. The periods are periods and periods and periods are chosen are chosen and periods are chose

$$
\text{Equivalences:} \quad \text{Equivalences:} \quad
$$

well-known control are the constructed are the classic polynomials are the logarithms and classic polynomials a<br>An article of which the logarithm is a set which the logarithm is a set which the logarithm is a set which the

and more generallly multiple polylogarithms (MPL's) and more generally multiple polylogarithms (MPL's)  $\Box$  and more generally multiple polylogarithms (MPI's) generally multiple polyic *<u>v</u> multi i*c polylogant<br>*\** 

$$
G(a_1,...,a_n;x) = \int_0^x \frac{dt}{t-a_1} G(a_2,...,a_n;t)
$$

+ They obey a "shuffle algebra" 2 *a*2

discussions.

special case,

Thus, any loop of a similar argument:

\n
$$
G(a_1, \ldots, a_k; x) G(a_{k+1}, \ldots, a_{k+l}; x) = \sum_{\sigma \in \Sigma(k,l)} G(a_{\sigma(1)}, \ldots, a_{\sigma(k+l)}; x)
$$
\nwhich can exactly simplify result: the math, invariance, then compute long in-

- **Example 20** which can greatly simplify results. Its math properties help compute loop integrals 2⌃(*k,l*) *G*(*a*(1)*,...,a*(*k*+*l*); *x*)*,* (2.3) where  $\alpha$ *i*,  $\beta$ *i* cair*f* on *ping* is conto: it indirepresence not permutations of *a*  $\mathbf{w}$  = *a*23 **a**<sub>24</sub> *a*<sub>2</sub> *, c*<sup>4</sup> =
	- At two loop and certainly beyond "elliptic" functions start appearing, including a<sup>1</sup> *alon and certainly heyond* "al It is possible to assign notions of *length* and *weight* to MPLs. The length of an iterated and to up that contently to uy once the first and second up potential, motorem,

$$
K(\lambda) = \int_0^1 \frac{dt}{\sqrt{(1 - t^2)(1 - \lambda t^2)}}, \quad E(\lambda) = \int_0^1 dt \sqrt{\frac{1 - \lambda t^2}{1 - t^2}}
$$

**Extensions of MPL technology to elliptic case are appearing and analyzing to MPLS. The length of an iterated of an i** length of an MPL *G*(*a*1*,...,an*; *x*) is *n*. The notion of weight, however, is more subtle. Extensions of MPL technology to elliptic case are appearing its elliptic version, this is not the general case. One can also assign a notion of weight for

 $\mathbf{p}$  integral (polynomia  $\mathbf{p}$  is always defined as the number of integrations, thus the number of integrations, the number of integrations, thus the number of integrations, thus the number of integrations, thus th length of an MPL *G*(*a*1*,...,an*; *x*) is *n*. The notion of weight, however, is more subtle. elliptic version, the general case of weight for a non-tegral case of weight for a non-tegral case of weight for weight for a sign of weight  $\alpha$  of  $\alpha$  and  $\beta$  and  $\beta$  are no integrals left to perform; see  $\alpha$  of  $\beta$ the weight remembers that of the iterated integral it originated from. For example, Broedel, Duhr, Dulat, Penante, Tancredi <mark>'19</mark>

# Computing progress: analytical, numerical

- Computer algebra!
	- ‣ Many dedicated mathematica packages for categorizing loop diagrams (FIRE, REDUZE, …)
	- ‣ Most powerful language, especially for high loops: FORM J. Vermaseren
		- ✓ still under active development. Recent: FORCER, code that writes other code
- Monte Carlo technology
	- ‣ improved parton showers, matching to fixed order

IMPORTANT:

As for other areas, for future progress we cannot take the **new talent** entering precision calculations for granted.

Support, including **recognition**, for creating/maintaining tools and technical innovations will be needed!

# SMEFT and top physics

Brown, Buckley,Englert,Ferrando,Galler,Miller, More,Russell,White,Warrack

### **TopFitter**

- ‣ Confront LHC and Tevatron top data with theory, including operators  $(14)$ 
	- pair production (+ vector boson) and single top
- Experimental uncertainties as given
- ‣ Theoretical ones: vary scales, and use PDF uncertainties
- ‣ NLO effects via SM K-factors
- Top flavour-changing interactions, global analysis
	- Top pair and single top contributions
	- Include NLO for SM included
	- Include also running and mixing for operators
- ✦ Recent note on common standards in EFT approach by all involved We note that neglecting *|M*HDO*|* order, if the interference of the operator with the operator with the SM is small or  $\sim$ conservation) then this neglected term could provide the leading contribution of new physics. Fur-

![](_page_39_Figure_13.jpeg)

**Aguilar-Saavedra et al it may provide the leading contribution in some regions of phase space. However, one regions of phase space. The leading contribution in some regions of phase space. However, one of phase space. How** could avoid avoid avoid these issues by a series by also calculating  $\alpha$  *arXiv*:1802.07237

### Dimension-6 operators for single top in t-channel

![](_page_40_Picture_1.jpeg)

de Beurs, EL, Vreeswijk,Vryonidou '18

For single top production in the t-channel, only 3 operators matter!

$$
O_{\varphi Q}^{(3)} = i \frac{1}{2} y_t^2 \left( \varphi^{\dagger} \overleftrightarrow{D}_{\mu} \varphi \right) (\overline{Q} \gamma^{\mu} \tau^I Q)
$$
  

$$
O_{tW} = y_t g_w (\overline{Q} \sigma^{\mu \nu} \tau^I t) \tilde{\varphi} W_{\mu \nu}^I
$$
  

$$
O_{qQ,rs}^{(3)} = (\overline{q}_r \gamma^{\mu} \tau^I q_s) (\overline{Q} \gamma_{\mu} \tau^I Q)
$$

- Easier to get info on each operator separately
- ✦ Effects will be small, hence we should also include QCD corrections to SM
- They will affect the cross section, and differential distributions, in different ways
	- used MadGraph5\_aMC@NLO

### Effect on amplitudes and cross sections

- How do these operators modify single top scattering amplitudes and the cross section?
	- ‣ Amplitude

$$
\mathcal{M} = \mathcal{M}_{\rm SM} + \sum_i \frac{1\text{TeV}^2}{\Lambda^2} C_i\, \mathcal{M}_i
$$

‣ Squared amplitude

$$
|\mathcal{M}|^2 = |\mathcal{M}|^2_{\mathrm{SM}} + \sum_i \frac{1\mathrm{TeV}^2}{\Lambda^2} C_i 2\mathrm{Re}\left(\mathcal{M}_{\mathrm{SM}}^*\mathcal{M}_i\right) + \sum_{i\leq j} \frac{1\mathrm{TeV}^4}{\Lambda^4} C_i C_j \left|\mathcal{M}\right|^2_{i,j}
$$

▶ Cross sections

$$
\sigma = \sigma_{\rm SM} + \sum_{i} \frac{1 \text{TeV}^2}{\Lambda^2} C_i \,\sigma_i + \sum_{i \le j} \frac{1 \text{TeV}^4}{\Lambda^4} C_i C_j \,\sigma_{i,j}
$$

 $\triangleleft$  NLO QCD needed in order to make C<sub>i</sub> corrections stand out from  $\sigma_{\text{SM}}(\mu, \alpha_s(\mu))$ 

# SMEFT single top distributions at NLO

- Wbj production then W-decay via MadSpin, parton showering with Pythia8
- Top quark  $p_T$  and  $n$  normalized distributions
	- ‣ Four-fermion operator different shape. QCD corrections at large pT and central rapidity notable
	- Allow for new physics operators in production and decay (EFT=2)

![](_page_42_Figure_5.jpeg)

### Next steps in automation in MG5\_aMC@NLO

- (Not yet NNLO...)
- SMEFT@NLO is fully included
	- ‣ but renormalization group running and mixing of all the operators still to be done
- MSSM@NLO is under construction. Two example plots [thanks to M. Zaro]
	- susy pair production at NLO; # jets in gluino pair production after decay and ps

![](_page_43_Figure_6.jpeg)

## Imprecise and somewhat uncertain outlook

- With only the first few percent of LHC data acquired, experiment will demand high theoretical precision
	- ‣ not just NLO or NNLO, small other effects come into play
- → not just NLO or NNLO, small other effects come into play<br>← Theory community is meeting the challenge, with quite spectacular progress in the last 15 years

S. Dittmaier @ Granada

### Can theory provide the necessary precision?

→ Optimists: "Yes. No show-stoppers seen, great progress can be anticipated."

Sceptics: "Enormous challenge! Conceptual progress difficult to extrapolate."

- New ideas, methods and talent give reason for optimism
	- ‣ with sufficient support, recognition and resources