Lattice QCD study of Zb tetraquark channel

+ two other quarkonium(like) channels

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Zb: S.P., Bahtiyar, Petkovic, 1912.02656

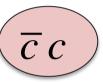
charmonium resonances: Piemonte, Collins, Padmanath, Mohler, S.P.: 1905.03506, PRD 2019

Pc: Skerbis, S.P., 1811.02285, PRD 2019

Outline

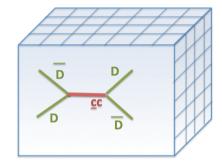
Lattice QCD study of

1) conventional charmonium resonances above <u>D</u>D threshold



$$J^{PC} = 1^{-1} \Psi(3770)$$
 know for long time

$$J^{PC}=3^{--}$$
 X(3842) discovered at LHCb 2019



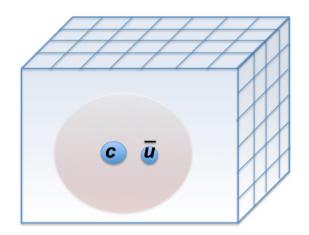
- 2) pentquark P_c channel discovered at LHCb 2015,2019
- $(\overline{c} cuud)$

3) tetraquark Z_b channel discovered at Belle 2011



Lattice QCD

$$L_{QCD} = -\frac{1}{4}G^a_{\mu\nu}G^{\mu\nu}_a + \sum_{q=u,d,s,c,b,t} \overline{q}i\gamma_\mu(\partial^\mu + ig_s\,G^\mu_a\,T^a)q - m_q\overline{q}q$$



input: g_s , m_q

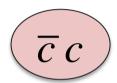
Numerical evaluation of QFT Feynman path integrals on discretized Eucledian space-time

$$\int DG \ Dq \ D\overline{q} \ e^{-S_{QCD}/\hbar}$$

$$S_{OCD} = \int d^4x \ L_{OCD}[G(x), q(x), \overline{q}(x)]$$

Extracted quantity: $E_n = energy of QCD eigenstate$ with given quantum numbers

$$E_1(p=0,J^p=0) = m_D$$



(1) Conventional charmonia

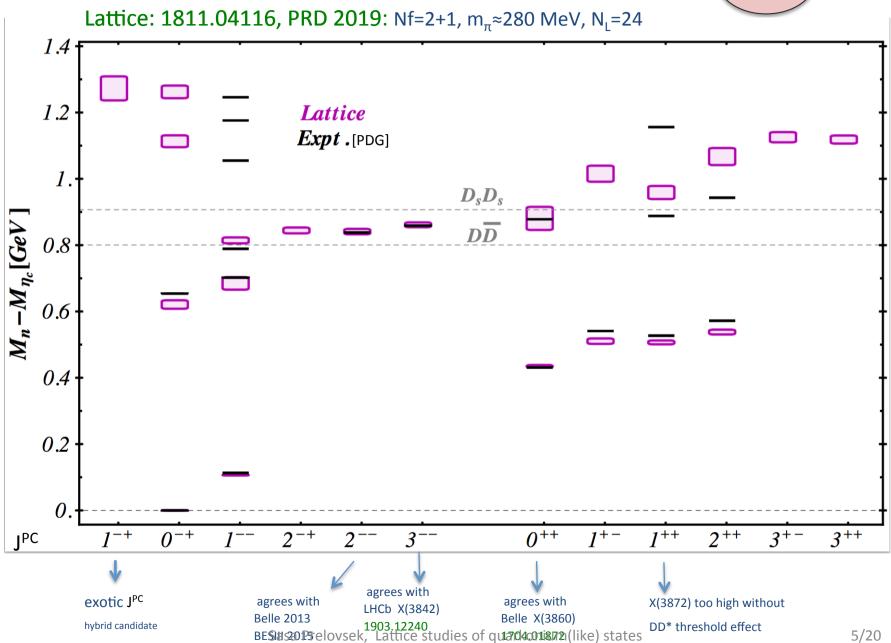
S. Piemonte, S. Collins, M. Padmanath, D. Mohler, S.P.: 1905.03506, PRD 2019

M. Padmanath, S. Collins, D. Mohler, S. Piemonte, S.P., A. Schafer, S. Weishaeupl: 1811.04116, PRD 2019

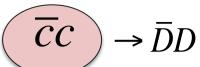
(Regensburg group)

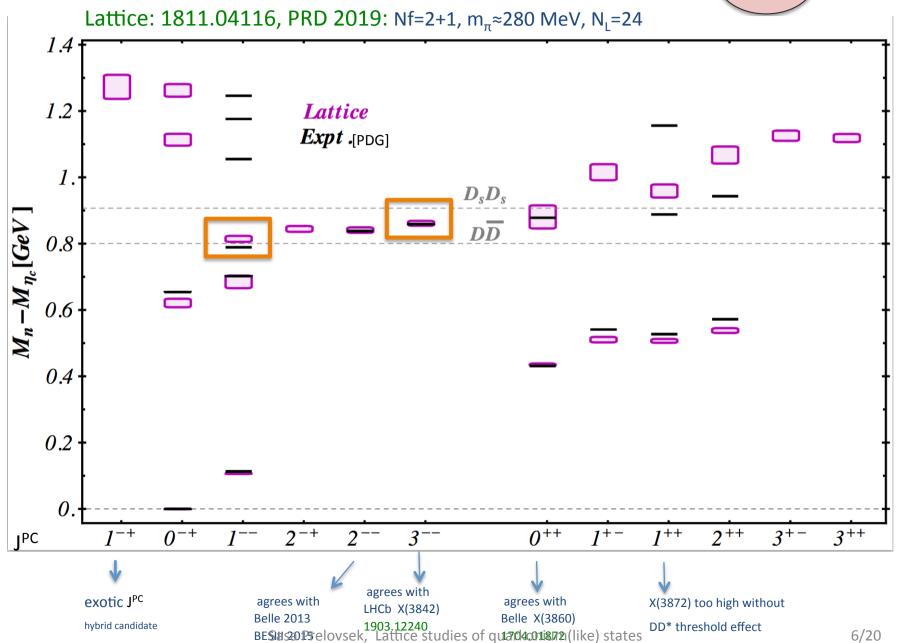




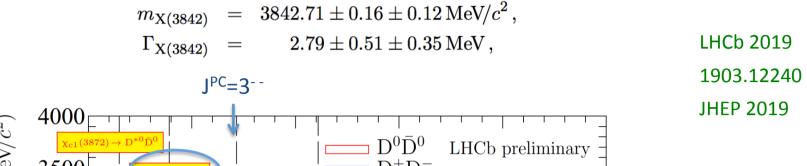


Next: strong decays of resonances





First exp. discovery of a charmonium with spin J=3



JPC not experimentally measured

LHCb paper:

"The narrow natural width and the mass of the X(3842) state suggest the interpretation as charmonium state with JPC = 3 -- "

Quark model quantum numbers:

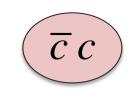
$$n^{2s+1}l_{J} = 1^{3}D_{3}$$

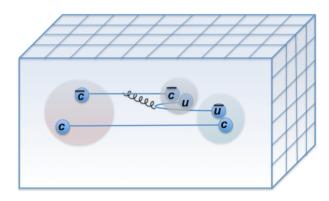
$$\bar{c}c \rightarrow \bar{D}D$$

Candidates/ $(1 \text{ MeV}/c^2)$ D^+D^- 3500 ⊡ $X(3842) \rightarrow D\bar{D}$ 3000 $\chi_{c2}(3930) \rightarrow D\bar{D}$ 2500 2000 1500 1000 500 3.8 3.7 3.9 4.1 4.2 $\left[\text{GeV}/c^2 \right]$ $m_{
m Dar{D}}$

Charmonia with JPC=1- and 3-

taking into account strong transitions to **DD** Strategy





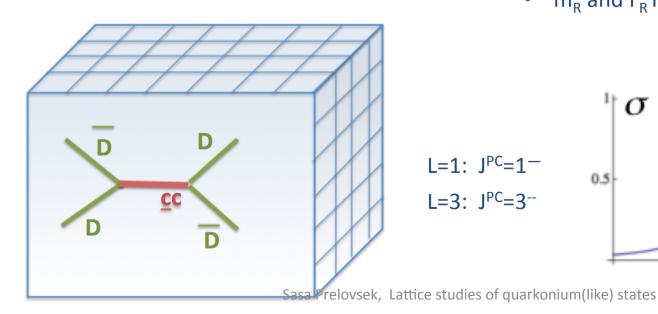
only 1 previous lattice study extracted width of charmonium resonances (0⁺⁺ and 1⁻⁻) Lang, Leskovec, Mohler, S.P., 1503.05363, JHEP 2015

- simulate DD scattering on the lattice
- determine scattering amplitude (via Luscher's method from En)

$$S_l(E) = \exp[2i\delta_l(E)], \quad l = 1,3$$

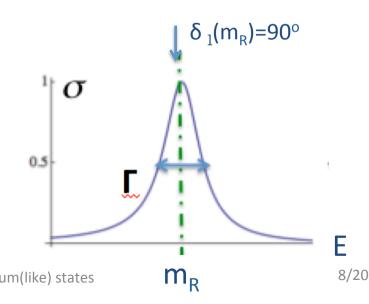
$$\sigma(E) \propto |S(E) - 1|^2 \propto |t(E)|^2$$

 m_R and Γ_R from Breit-Wigner type fits

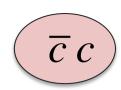


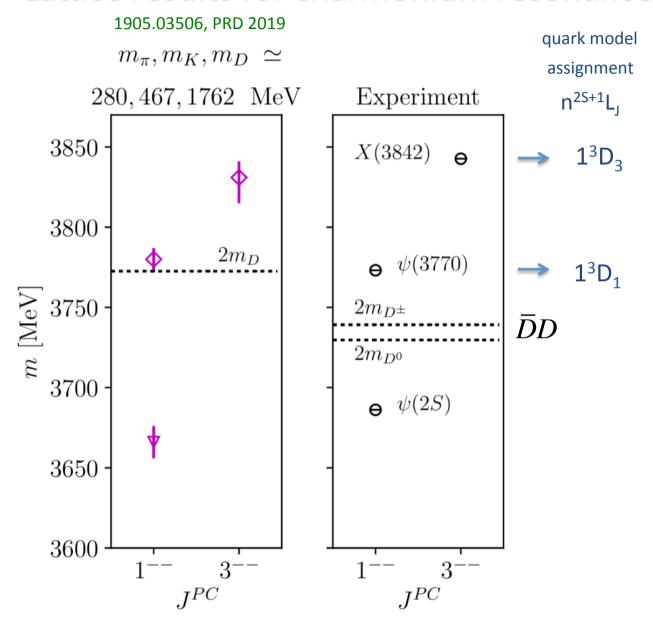
L=1: J^{PC}=1⁻

L=3: J^{PC}=3⁻⁻



Lattice results for charmonium resonances





widths of resonances:

ψ(3770)

$$\Gamma = \frac{g^2 \ p^3}{6\pi \ s}$$

	g
lat	$16.0^{+2.1}_{-0.2}$
exp	18.7 ± 0.9

X(3842)
 to narrow to
 resolve in this lat. sim.

Note: all exotic hadrons are strongly decaying resonances!



(2) P_c pentaquark channel

U. Skerbis, S.P., 1811.02285, PRD 2019

Lattice study of P_c pentaquark channel



$$P_c = \text{uud}\overline{c}c \rightarrow (\text{uud})(\overline{c}c)$$
light-baryon charmonium

$$\rightarrow$$
 (uuc) ($\overline{c}d$) charmed-baryon charmed-meson

Question we address:

Do Pc resonances appear in one-channel p J/ ψ scattering on the lattice (in approximation where this channel is decoupled from other channels)

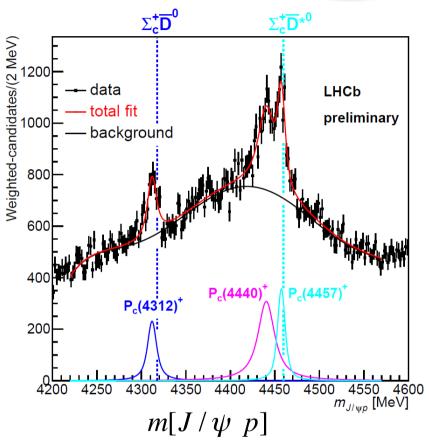
$$p J/\psi \rightarrow P_c \rightarrow p J/\psi$$

We simulate this scattering and cover also the energy region of P_c for the first time.

U. Skerbis, S.P., 1811.02285, PRD 2019

The answer from our lattice simulation: No.

This indicates that the coupling of p J/ ψ channel with other two-hadron channels is likely responsible for Pc resonances in experiment.



This is in line with LHCb results, where Pc's are found near other thresholds. This by itself indicates that other channels are important.



(3) Z_b⁺ tetraquark channel

S.P., H. Bahtiyar, J. Petkovic, 1912.02656

Zb in experiment

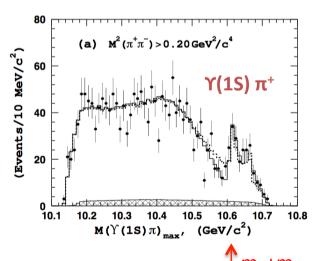
discovered by Belle in 2011 [PRL 108 (2012) 122001]

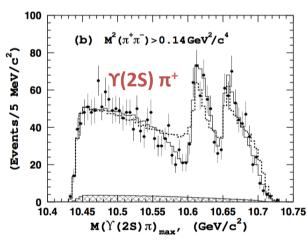
$$Z_{b}^{+}(10610)$$
, $Z_{b}^{+}(10650)$

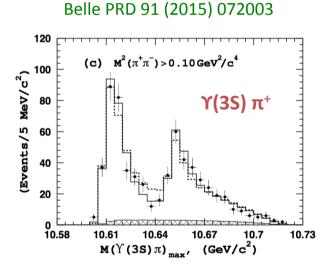
$$Z_b^+ \to \Upsilon \pi^+$$

$$\overline{bb} \overline{d}u$$

 Z_b observed in strong decays $~Y(1S)\,\pi$, $Y(2S)\,\pi$, $Y(3S)\,\pi$, $h_b(1S)\,\pi$, $h_b(2S)\,\pi$, $B~B^*$, $B^*~B^*$







Z_b on the lattice with static b and <u>b</u>

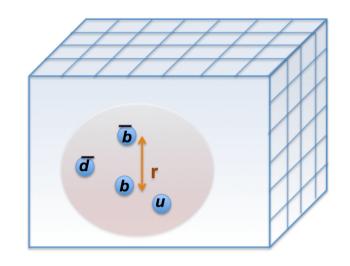
Only previous lat study

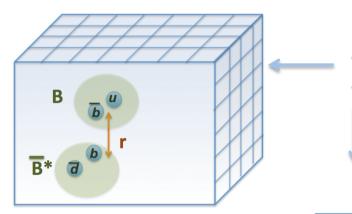
Bicudo, Cichy, Peters, Wagner [proceedings Lat16: 1602.07621

proceedings Lat17: 1709.03306]

Born-Oppenheimer approach

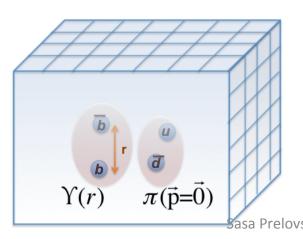
Fock components incorporated for $S_{\bar{b}b} = 1$

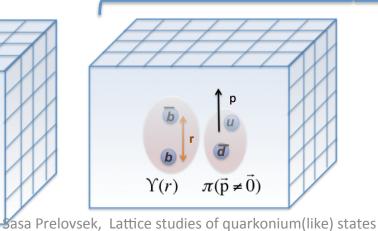


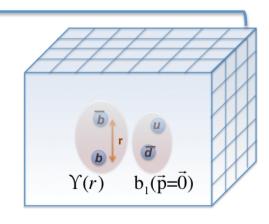


- main aim: extract static potential V(r) between B and \underline{B}^*
- momentum of light degrees of freedom not conserved in presence of static quarks

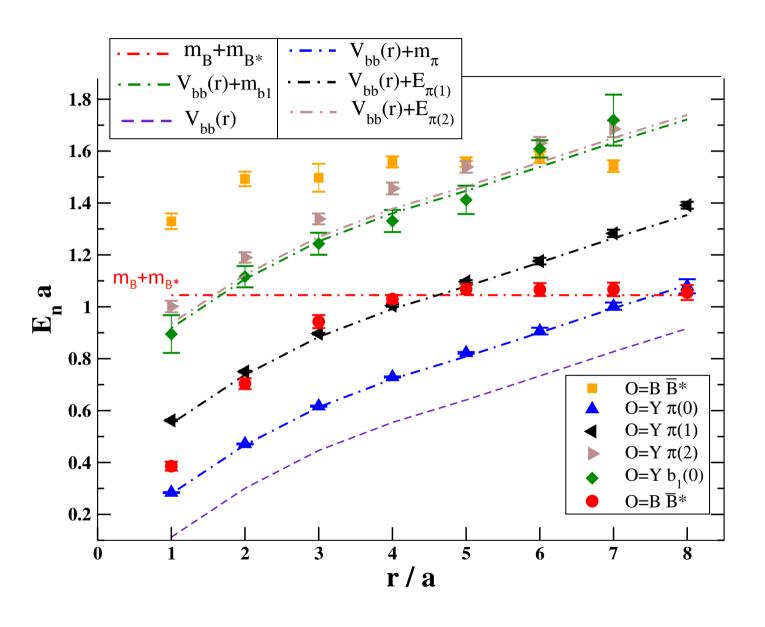


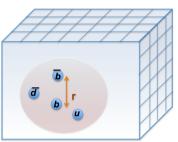






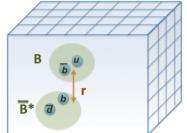
Eigen-energies of Z_b system

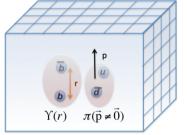


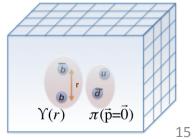


dot-dashed-lines:

 $E_n^{\text{non-int}}$

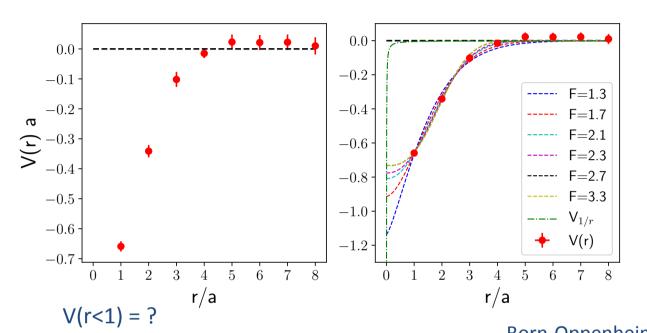


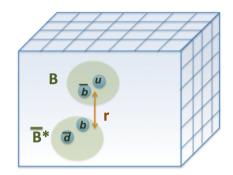




Static potential V(r) for interaction between B and B*

We assume that B \underline{B}^* eigenstate is decoupled from other channels (overlaps support that).





Born-Oppenheimer approach: B and <u>B</u>* move in

$$\left[-\frac{\hbar^2}{2\mu} \frac{d^2}{dr^2} + \frac{\hbar^2 L(L+1)}{2\mu r^2} + V(r) \right] u(r) = Eu(r)$$

$$\mu = \frac{1}{2} m_B^{\text{exp}}, \quad \psi \propto \frac{u}{r} Y_{LM}$$

We focus on most relevant: s-wave (L=0)

$$V(r) = E_n(r) - m_B - m_{B*} (m_{B*} = m_B)$$

parametrizing V:
$$V(r) = -A \exp[-(\frac{r}{d})^F]$$

Need for theory input!

- · analytic form of potential
- behavior at very small r

Results on Zb based on extracted V(r)

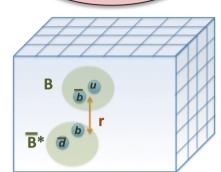
$$\left(\overline{b}b\ \overline{d}u\right)$$

$$V(r) = -A \exp[-(\frac{r}{d})^F]$$
 for F=1.3

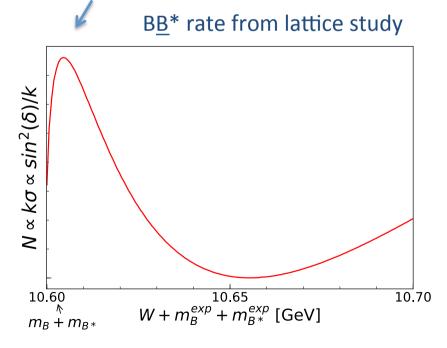
Zb found to be virtual bound state (pole of S-matrix for k=-i|k|)

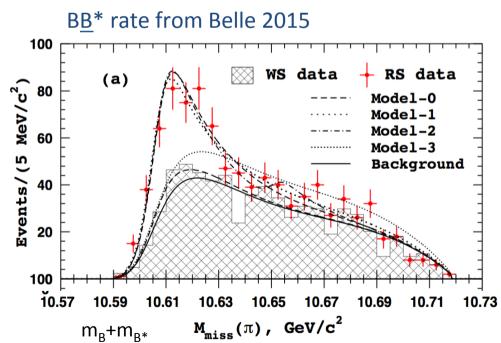
$$M_{Zb} = m_B + m_{B^*} - 13 \pm 10 \text{ MeV}$$

consistent with re-analysis of exp data: Hanhart, et al 1805.07453, PRD 2018



Zb peak is a consequence of virtual bound state





$$M \approx m_B + m_{B^*} - 400 \text{ MeV}$$

 $M_{Zb} \approx m_B + m_{B^*} - 13 \text{ MeV}$

lattice

experiment
Belle PRD 91 (2015) 072003

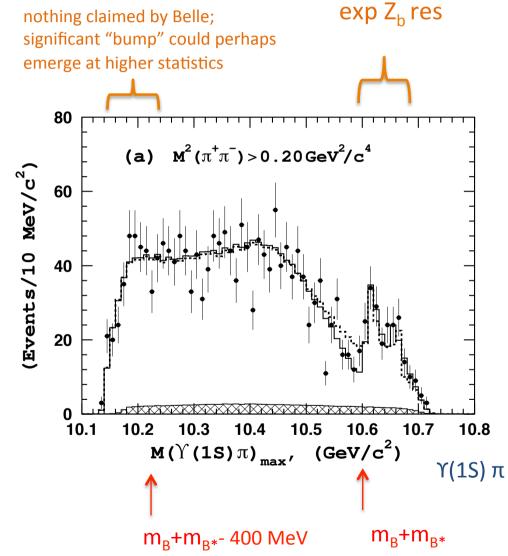
Relating lattice results to Belle experiment

LHCb: try to look for Zb in BB* final state (exclusive or inclusive)

A possible deep bound state !?

If it exist: it could be perhaps visible only in $\Upsilon(1S)$ π , since it is located below all other thresholds.

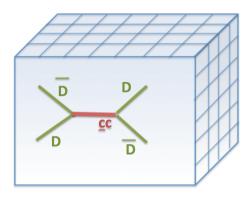
virtual bound state

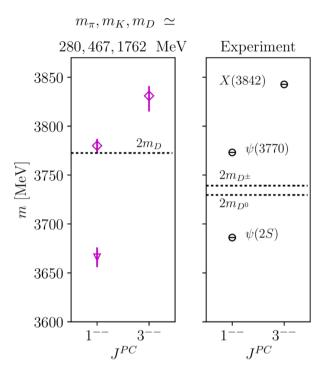


Conclusions

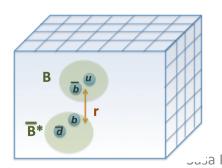
Lattice QCD study of

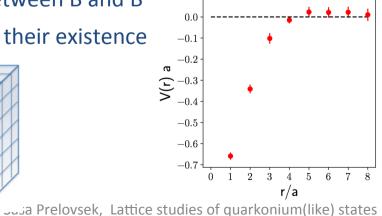
charmonium resonances

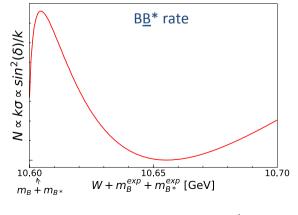




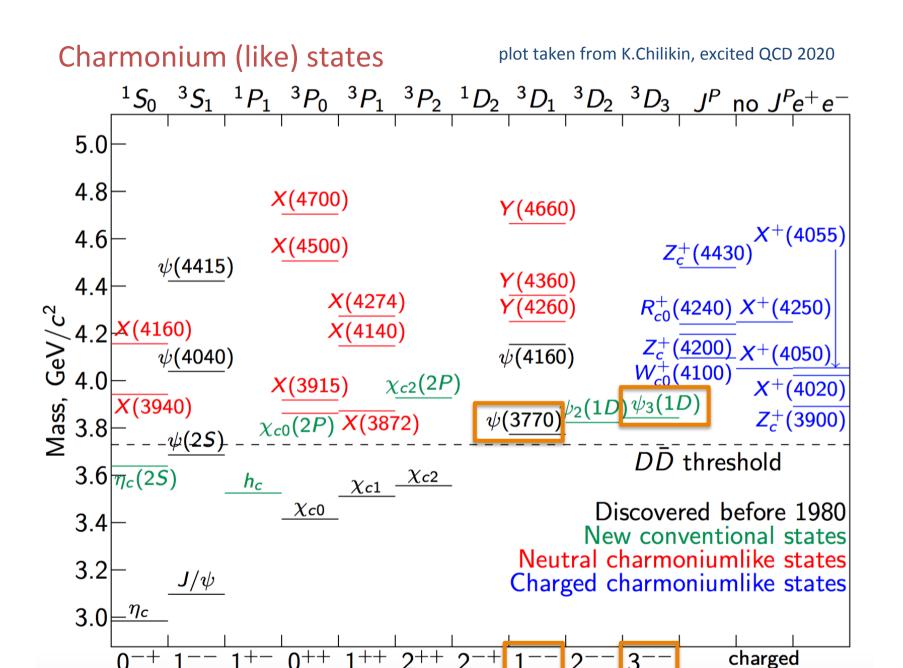
- Pc does not appear in one-channel scattering $p \ J/\psi \to P_c \to p \ J/\psi$ This indicates that the coupling to other channels (for example $\Sigma_c^+ \underline{D}^{(*)}$) is important
- Zb: strong attraction between B and B* is likely responsible for their existence







Backup



Charmonium resonances in DD scattering fits of phase shifts for l=1,3

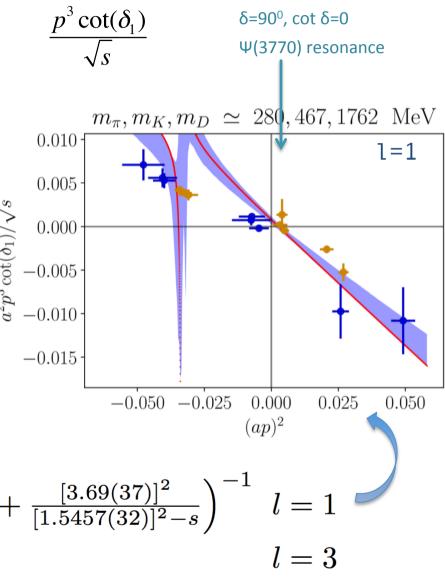
$$E_{cm} = \sqrt{s} = 2\sqrt{m_D^2 + p^2}$$

p = relative momenta of D-mesons in CMF

$$\frac{p^{2l+1}\cot(\delta_l)}{\sqrt{s}} = \frac{m^2 - s}{G^2}$$
 Breit-Wigner
$$\delta_1(m_R) = 90^\circ$$

Result:

$$\frac{p^{2l+1}\cot(\delta_3)}{\sqrt{s}} = \frac{m_3^2 - s}{g_3^2}$$
esult:
$$\frac{p^{2l+1}\cot(\delta)}{\sqrt{s}} = \begin{cases} \left(\frac{[0.63(33)]^2}{[1.4966(30)]^2 - s} + \frac{[3.69(37)]^2}{[1.5457(32)]^2 - s}\right)^{-1} & l = 1\\ \frac{[1.568(11)]^2 - s}{[0.07(3)]^2} & l = 3 \end{cases}$$



Scattering amplitude t(E) in complex energy plane

$$S_l(E) = \exp[2i\delta_l(E)] = 1 + 2i\rho t_l(E)$$

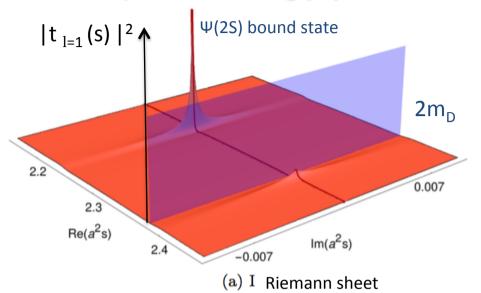
$$t_l(s) = \frac{1}{\rho \cot(\delta_l) - i \rho}$$

$$\rho = \frac{2p}{\sqrt{s}} = \sqrt{1 - 4\frac{m_D^2}{s}} \, .$$

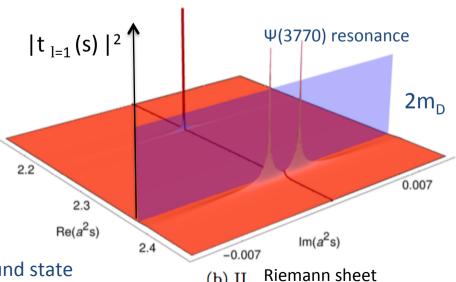
poles in t(s) related to resonance and bound states

Fig for l=1 and $m_D \approx 1762$ MeV:

one resonance, one bound state, one virtual bound state

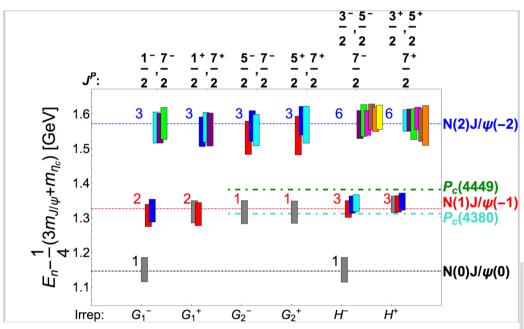


virtual bound state

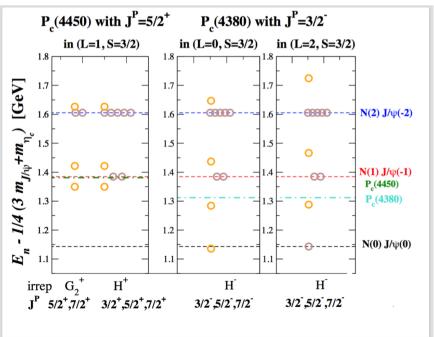


Sasa Prelovsek, Lattice studies of quarkonium(like) states

proton J/Ψ scattering in lattice QCD in Pc channels



U. Skerbis, S. Prelovsek, 1811.02285



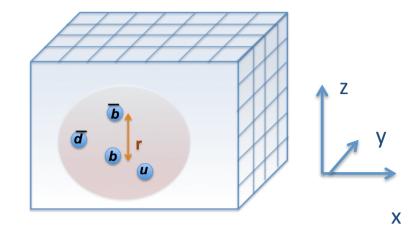
Sasa Prelovsek, Lattice studies of quarkonium(like) states

Zb with static b b

Good symmetries and quantum numbers:

$$I=1$$
 $I_3=0$ (consider neutral Z_b)

$$S_{\text{heavy}} = 1 \quad (Sz)_{\text{heavy}} = 0 \qquad \overline{b}(\uparrow)b(\downarrow) - \overline{b}(\downarrow)b(\uparrow)$$



heavy quark can not flip spin via gluon exchange

note: transition is not possible to final states with $S_{heavy}=0$ (η_b , h_b)

$$(Jz)_{light} = 0$$
 [Jx and Jy not conserved]

 $C \cdot P = -1$ (P= inversion over midpoint between b and <u>b</u>)

$$R_{light}$$
 = reflection over yz plane = P_{light} * R_{light} (y, π) : ϵ =-1

momentum of light degrees of freedom: not conserved

0.0 V(r) -0.2Masses of bound states in Zb channel ----- F=1.3 -0.4F=1.7 F=2.1 -0.6F = 2.3F = 2.7-0.8F=3.3 -1.0 $V_{1/r}$ V(r)-1.2(a) $V_{reg} + V_{1/r}$, r/a = [1, 4](b) $V_{rea} + V_{1/r}$, r/a = [2, 4]r/a $-m_B - m_{B^*}$ [MeV] - m_{B*} [MeV] -200 -100 -400 -200 Bound state Bound state -600 Virtual bound state Virtual bound state -300-800 m_B -1000-400Σ ₹ -1200 -5002.5 3.0 3.5 2.5 3.0 1.5 2.0 1.5 2.0 3.5 Parameter F of potential V Parameter F of potential V (c) V_{reg} , r/a = [1, 4](d) V_{reg} , r/a = [2, 4] $-m_B - m_{B^*}$ [MeV] - m_{B*} [MeV] -200 -100 -400-200Bound state Bound state -600Virtual bound state Virtual bound state -300 -800 m_B -1000-400Σ ₹ -1200 -5003.0 3.5 2.5 3.0 2.0 2.5 1.5 2.0 3.5 1.5 Parameter F of potential V Parameter F of potential V