

# Elliptic multiple polylogarithms

## an algorithmic approach

The calculation of two and three loop scalar diagrams

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# Outline

- 1 Elliptic multiple polylogarithms
- 2 General solution strategy
- 3 Tree-loop banana diagram
- 4 All order sunrise diagram

# Elliptic multiple polylogarithms [Brown and Levin(2013)], [Broedel(2015)], [Weinzierl(2016)], [Duhr et al.(2018)]

- Definition eMPLs

$$E_4 \left( \begin{matrix} n_1, \dots, n_k \\ c_1, \dots, c_k \end{matrix}; x, \vec{a} \right) = \int_0^x dt \varphi_{n_1}(c_1, t, \vec{a}) E_4 \left( \begin{matrix} n_2, \dots, n_k \\ c_2, \dots, c_k \end{matrix}; t, \vec{a} \right),$$

# Integration kernels eMPLs [Duhr et al.(2018)]

$$\varphi_0(0, x) = \frac{c_4}{y}$$

$$\varphi_1(c, x) = \frac{1}{x - c} \quad \varphi_{-1}(c, x) = \frac{y c}{y(x - c)}$$

$$\varphi_{-1}(\infty, x) = \frac{x}{y} \quad \varphi_1(\infty, x) = \frac{c_4}{y} Z_4(x)$$

$$\varphi_n(\infty, x) = \frac{c_4}{y} Z_4^{(n)}(x)$$

$$\varphi_{-n}(\infty, x) = \frac{x}{y} Z_4^{(n-1)}(x) - \frac{\delta_{n2}}{c_4}$$

$$\varphi_n(c, x) = \frac{1}{x - c} Z_4^{(n-1)}(x) - \delta_{n2} \Phi_4(x)$$

$$\varphi_{-n}(c, x) = \frac{y c}{y(x - c)} Z_4^{(n-1)}(x)$$

# Integrals

- General integral

$$\int dx R(x, y) \chi(x) \quad \text{with} \quad \chi(x) = Z_4^{(m)} E_4 \left( \frac{\vec{n}}{c}; x \right)$$

- After Partial fraction decomposition

$$A_k[\chi] = \int dx x^k \chi(x) \quad B_{c,k}[\chi] = \int \frac{dx}{(x-c)^k} \chi(x)$$
$$C_k[\chi] = \int \frac{dx}{y} x^k \chi(x) \quad D_{c,k}[\chi] = \int \frac{dx}{y(x-c)^k} \chi(x)$$

# IBP reduction algorithms [Duhr et al.(2018)]

- Algorithm A

$$A_k[\chi(x)] = \frac{x^{k+1}}{k+1} \chi(x) - \frac{1}{k+1} A_{k+1}[\partial_x \chi(x)]$$

- Algorithm B

$$B_{c,k}[\chi(x)] = \frac{\chi(x)}{(k-1)(x-c)^{k-1}} - \frac{1}{k-1} B_{c,k-1}[\partial_x \chi(x)]$$

- Algorithm C

$$y x^{k-2} \chi(x) = \frac{1}{2} \sum_{l=0}^3 (-1)^l s_l(a_1, a_2, a_3, a_4) \\ [(2k-1-l) C_{k-l}[\chi(x)] + 2 C_{k+1-l}[\partial_x \chi(x)]]$$

- Algorithm D

$$\frac{y}{(x-c)^{k-1}} \chi(x) = -\frac{1}{2} \sum_{l=0}^3 [(2k-2-l) D_{c,k-l}[\chi] - 2 D_{c,k-1-l}[\partial_x \chi]]$$

$$\frac{1}{l!} \partial_c^l y_c^2$$

- Algorithm D where  $c = a_i$ , a root of  $y$

$$\frac{y}{(x-a_1)^k} \chi(x) = -\frac{1}{2} \sum_{l=0}^3 [(2k-1-l) D_{a_1,k-l}[\chi] - 2 D_{a_1,k-1-l}[\partial_x \chi]]$$

$$\frac{1}{l!} \partial_{a_1}^l (a_{12} a_{13} a_{14})$$

# Base integrals

- Performing partial fractioning and using the IBP relation iteratively will lead to 7 base integrals

$$A_{-1}[\chi] = \int \frac{dx}{x} \chi(x) \quad B_{c,1}[\chi] = \int \frac{dx}{x-c} \chi(x)$$

$$C_2[\chi] = \int \frac{x^2 dx}{y} \chi(x) \quad C_1[\chi] = \int \frac{x dx}{y} \chi(x)$$

$$C_0[\chi] = \int \frac{dx}{y} \chi(x) \quad C_{-1}[\chi] = \int \frac{dx}{yx} \chi(x)$$

$$D_{c,1}[\chi] = \int \frac{dx}{y(x-c)} \chi(x)$$



# Feynman integrals

- Feynman integrals with Feynman parametrization

$$I_{\nu_1, \dots, \nu_n}(\{s\}) = N \left( i\pi^{\frac{d}{2}} \right)^l \frac{\Gamma(\nu - \frac{ld}{2})}{\prod_{i=1}^n \Gamma(\nu_i)} \int_{\Delta} d^n \vec{\alpha} \left( \prod_{i=1}^n \alpha_i^{\nu_i - 1} \right) \mathcal{U}^{\nu - \frac{d}{2}(l+1)} \mathcal{F}^{-\nu + \frac{ld}{2}},$$

- Integration domain

$$\Delta = \left\{ \vec{\alpha} \mid \alpha_i \geq 0, \sum_{i \in \tilde{S}} \alpha_i = 1 \right\}$$

# Linearly reducible elliptic Feynman integrals

- Linear reducible Feynman integrals: at every integration step one can do a linear factorization  $\rightarrow$  final result in terms of MPLs. [Brown(2009), Panzer(2015)]
- Linear reducible elliptic Feynman integrals: linearly reducible if one excludes the last integration. [Hidding and Moriello(2019)]

$$\mathcal{I} = \int dx \text{MPL}(x, y) = \int dx R(x, y) G(f_1(x, y), \dots, f_n(x, y); x)$$

## Statement

Linearly reducible elliptic Feynman integrals can always be expressed as a linear combination of eMPLs.

# Solution strategy

- I Choose a Feynman diagram ( $n$  propagators)
- II Use Feynman parameter representation for the corresponding integral ( $n$  Feynman parameters)
- III Perform  $n-1$  integrations in terms of MPLs
  - ▶ Use the Cheng-Wu theorem and a convenient integration order [Cheng and Wu(1987)]
  - ▶ Use e.g. HyperInt to perform the  $n-1$  integrations [Panzer(2015)]
- IV Obtain a linearly reducible elliptic Feynman integral

# Solution strategy

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- IV Obtain a linearly reducible elliptic Feynman integral
- V Rewrite MPLs to eMPL
- VI Perform the last integration and obtain the primitive
- VII Regularize the primitive
- VIII Perform a numerical check

# Banana like diagrams

# Tree-loop Banana A and Banana B

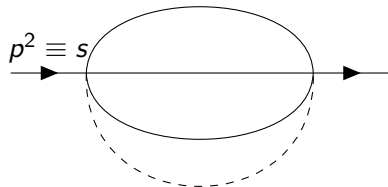


Figure: Banana A

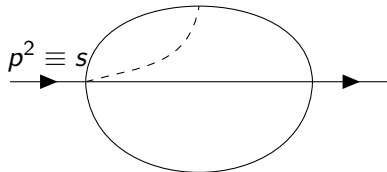
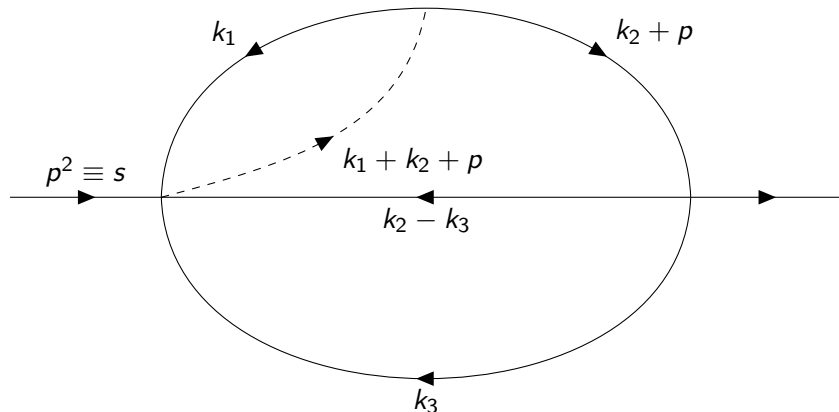


Figure: Banana B

# Calculation banana B

I Choose Feynman diagram, 5 propagators



# Calculation banana B

II Integral in Feynman parameter representation, 5 parameters. Finite integral with  $d = 4 - 2\epsilon$  and  $S_{21221}$

$$S_{21221}(S, m^2) \Big|_{\epsilon^0} = \int_{\Delta} d^5 \vec{\alpha} \alpha_1 \alpha_3 \alpha_4 \mathcal{U}^0 \mathcal{F}^{-2}$$

• Symanzik polynomials

$$\mathcal{U} = \alpha_2 (\alpha_4 \alpha_5 + \alpha_3 (\alpha_4 + \alpha_5)) + \alpha_1 (\alpha_2 (\alpha_3 + \alpha_4) + \alpha_4 \alpha_5 + \alpha_3 (\alpha_4 + \alpha_5))$$

$$\begin{aligned} \mathcal{F} = & (\alpha_2 (\alpha_3 + \alpha_4) + \alpha_4 \alpha_5 + \alpha_3 (\alpha_4 + \alpha_5)) \alpha_1^2 m^2 + \alpha_1 (\alpha_3^2 (\alpha_4 + \alpha_5) m^2 \\ & + \alpha_4 \alpha_5 (\alpha_4 + \alpha_5) m^2 + \alpha_3 (\alpha_4^2 m^2 + \alpha_5^2 m^2 + \alpha_5 \alpha_4 (3m^2 + S)) \\ & + \alpha_2 (\alpha_3^2 m^2 + \alpha_4 (\alpha_4 + 2\alpha_5) m^2 + \alpha_3 (2\alpha_5 m^2 + \alpha_4 (3m^2 + S))) \\ & + \alpha_2 (\alpha_3^2 (\alpha_4 + \alpha_5) m^2 + \alpha_4 \alpha_5 (\alpha_4 + \alpha_5) m^2 \\ & + \alpha_3 (\alpha_4^2 m^2 + \alpha_5^2 m^2 + \alpha_5 \alpha_4 (3m^2 + S))) \end{aligned}$$



# Calculation banana B

III Perform  $n-1$  integrations in terms of MPLs by using Cheng-Wu

# Calculation banana B

IV Obtain a linearly reducible elliptic Feynman integral

$$\begin{aligned} S_{21221}(S, m^2) \Big|_{\epsilon^0} = & \int_0^1 dx \left[ P_1(x) G(0, \gamma(x), 1) \right. \\ & - \frac{1}{2} P_1(x) \log^2 \left( \frac{m^2 x}{m^2 - S(x-1)x} \right) + P_2(x, y) G(0, \phi(x, y), 1) \\ & + P_2(x, -y) G(0, \phi(x, -y), 1) - P_2(x, y) G(\phi(x, -y), \phi(x, y), 1) \\ & - P_2(x, -y) G(\phi(x, y), \phi(x, -y), 1) - P_1(x) G(\phi(x, y), \gamma(x), 1) \\ & - P_1(x) G(\phi(x, -y), \gamma(x), 1) - P_2(x, y) \log^2(\rho(x, y)) \\ & \left. - P_2(x, -y) \log^2(\rho(x, -y)) \right] \end{aligned}$$

V Rewrite MPLs to eMPLs by taking the derivative and integrate back

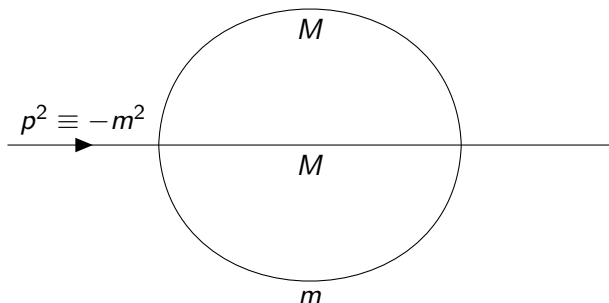
VI Perform the last integration and obtain primitive

# Calculation banana B

$$\begin{aligned}
 S_{212221}(S, m^2) \Big|_{\epsilon^0} = & -\frac{1}{2c_4 m^2 S} \left[ 2 \left( E_4 \left( \begin{smallmatrix} 0 & 1 & 1 \\ 0 & a_3 & 0 \end{smallmatrix}; 1 \right) + E_4 \left( \begin{smallmatrix} 0 & 1 & 1 \\ 0 & a_4 & 0 \end{smallmatrix}; 1 \right) \right) + 2E_4 \left( \begin{smallmatrix} 0 & 1 & -1 \\ 0 & a_3 & \infty \end{smallmatrix}; 1 \right) \right. \\
 & + 2E_4 \left( \begin{smallmatrix} 0 & 1 & -1 \\ 0 & a_3 & 0 \end{smallmatrix}; 1 \right) + 2E_4 \left( \begin{smallmatrix} 0 & 1 & -1 \\ 0 & a_4 & \infty \end{smallmatrix}; 1 \right) + 2E_4 \left( \begin{smallmatrix} 0 & 1 & -1 \\ 0 & a_4 & 0 \end{smallmatrix}; 1 \right) + E_4 \left( \begin{smallmatrix} 0 & 1 & -1 \\ 0 & c_1 & \infty \end{smallmatrix}; 1 \right) \\
 & + E_4 \left( \begin{smallmatrix} 0 & 1 & -1 \\ 0 & c_1 & 0 \end{smallmatrix}; 1 \right) + E_4 \left( \begin{smallmatrix} 0 & 1 & 1 \\ 0 & c_1 & 0 \end{smallmatrix}; 1 \right) - 3E_4 \left( \begin{smallmatrix} 0 & 1 & -1 \\ 0 & 0 & \infty \end{smallmatrix}; 1 \right) + E_4 \left( \begin{smallmatrix} 0 & 1 & -1 \\ 0 & 1 & \infty \end{smallmatrix}; 1 \right) \\
 & - 3E_4 \left( \begin{smallmatrix} 0 & 1 & -1 \\ 0 & 0 & 0 \end{smallmatrix}; 1 \right) - 3E_4 \left( \begin{smallmatrix} 0 & 1 & 1 \\ 0 & 0 & 0 \end{smallmatrix}; 1 \right) + E_4 \left( \begin{smallmatrix} 0 & 1 & -1 \\ 0 & 1 & 0 \end{smallmatrix}; 1 \right) + E_4 \left( \begin{smallmatrix} 0 & 1 & 1 \\ 0 & 1 & 0 \end{smallmatrix}; 1 \right) \\
 & - c_4 \left( 2E_4 \left( \begin{smallmatrix} -1 & 1 & -1 \\ \infty & a_3 & 0 \end{smallmatrix}; 1 \right) + 2E_4 \left( \begin{smallmatrix} -1 & 1 & -1 \\ \infty & a_3 & \infty \end{smallmatrix}; 1 \right) + 2E_4 \left( \begin{smallmatrix} -1 & 1 & 1 \\ \infty & a_3 & 0 \end{smallmatrix}; 1 \right) \right. \\
 & + 2E_4 \left( \begin{smallmatrix} -1 & 1 & -1 \\ \infty & a_4 & 0 \end{smallmatrix}; 1 \right) + 2E_4 \left( \begin{smallmatrix} -1 & 1 & -1 \\ \infty & a_4 & \infty \end{smallmatrix}; 1 \right) + 2E_4 \left( \begin{smallmatrix} -1 & 1 & 1 \\ \infty & a_4 & 0 \end{smallmatrix}; 1 \right) \\
 & + E_4 \left( \begin{smallmatrix} -1 & 1 & -1 \\ \infty & c_1 & 0 \end{smallmatrix}; 1 \right) + E_4 \left( \begin{smallmatrix} -1 & 1 & -1 \\ \infty & c_1 & \infty \end{smallmatrix}; 1 \right) + E_4 \left( \begin{smallmatrix} -1 & 1 & 1 \\ \infty & c_1 & 0 \end{smallmatrix}; 1 \right) - 3 \left( E_4 \left( \begin{smallmatrix} 1 & -1 & -1 \\ c_1 & 0 & \infty \end{smallmatrix}; 1 \right) \right. \\
 & + E_4 \left( \begin{smallmatrix} 1 & -1 & -1 \\ c_1 & \infty & 0 \end{smallmatrix}; 1 \right) + E_4 \left( \begin{smallmatrix} 1 & -1 & -1 \\ c_1 & \infty & \infty \end{smallmatrix}; 1 \right) + E_4 \left( \begin{smallmatrix} 1 & -1 & 1 \\ c_1 & \infty & 0 \end{smallmatrix}; 1 \right) + E_4 \left( \begin{smallmatrix} 1 & 1 & -1 \\ c_1 & 0 & \infty \end{smallmatrix}; 1 \right) \\
 & + E_4 \left( \begin{smallmatrix} 1 & -1 & -1 \\ c_1 & 0 & 0 \end{smallmatrix}; 1 \right) + E_4 \left( \begin{smallmatrix} 1 & -1 & 1 \\ c_1 & 0 & 0 \end{smallmatrix}; 1 \right) + E_4 \left( \begin{smallmatrix} 1 & 1 & -1 \\ c_1 & 0 & 0 \end{smallmatrix}; 1 \right) + E_4 \left( \begin{smallmatrix} 1 & 1 & 1 \\ c_1 & 0 & 0 \end{smallmatrix}; 1 \right) \Big) \\
 & - 3E_4 \left( \begin{smallmatrix} -1 & 1 & -1 \\ \infty & 0 & 0 \end{smallmatrix}; 1 \right) - 3E_4 \left( \begin{smallmatrix} -1 & 1 & -1 \\ \infty & 0 & \infty \end{smallmatrix}; 1 \right) - 3E_4 \left( \begin{smallmatrix} -1 & 1 & 1 \\ \infty & 0 & 0 \end{smallmatrix}; 1 \right) + E_4 \left( \begin{smallmatrix} -1 & 1 & -1 \\ \infty & 1 & 0 \end{smallmatrix}; 1 \right) \\
 & \left. + E_4 \left( \begin{smallmatrix} -1 & 1 & -1 \\ \infty & 1 & \infty \end{smallmatrix}; 1 \right) + E_4 \left( \begin{smallmatrix} -1 & 1 & 1 \\ \infty & 1 & 0 \end{smallmatrix}; 1 \right) \right] .
 \end{aligned}$$

# All order sunrise diagram

# All order sunrise diagram



- Computed in  $d=4-2\epsilon$

# An all order sunrise solution [Kalmykov and Kniehl(2009)]

$$J_{122}(m^2, M^2) = - (M^2)^{-2-\epsilon} (m^2)^{1-\epsilon} \gamma^2 (1+\epsilon) \frac{(1+\epsilon)}{\epsilon(1-\epsilon)}$$
$$\left[ \frac{1}{6} {}_4F_3 \left( \begin{matrix} 1, \frac{3}{2}, 1+\frac{\epsilon}{2}, \frac{3}{2}+\frac{\epsilon}{2} \\ 2-\epsilon, \frac{5}{4}, \frac{7}{4} \end{matrix} \middle| -t^2 \right) \right.$$
$$- \left( \frac{M^2}{m^2} \right)^{1-\epsilon} \frac{\epsilon}{(1+\epsilon)(1+2\epsilon)} {}_4F_3 \left( \begin{matrix} 1, \frac{1}{2}+\epsilon, 1+\frac{\epsilon}{2}, 1+\epsilon \\ \frac{3}{2}-\frac{\epsilon}{2}, \frac{3}{4}+\frac{\epsilon}{2}, \frac{5}{4}+\frac{\epsilon}{2} \end{matrix} \middle| -t^2 \right)$$
$$\left. - \left( \frac{M^2}{m^2} \right)^{-\epsilon} \frac{(1-\epsilon)}{(2-\epsilon)(3+2\epsilon)} {}_4F_3 \left( \begin{matrix} 1, \frac{3}{2}+\frac{\epsilon}{2}, \frac{3}{2}+\epsilon, 1+\epsilon \\ 2-\frac{\epsilon}{2}, \frac{5}{4}+\frac{\epsilon}{2}, \frac{7}{4}+\frac{\epsilon}{2} \end{matrix} \middle| -t^2 \right) \right]$$

$$t = \frac{m^2}{2M^2}$$

- For an all order sunrise with equal masses see [Adams and Weinzierl(2018)]

# An all order sunrise solution [Kalmykov and Kniehl(2009)]

$$J_{122}(m^2, M^2) = - (M^2)^{-2-\epsilon} (m^2)^{1-\epsilon} \gamma^2 (1 + \epsilon) \frac{(1 + \epsilon)}{\epsilon(1 - \epsilon)}$$
$$\left[ \frac{1}{6} {}_4F_3 \left( \begin{matrix} 1, \frac{3}{2}, 1 + \frac{\epsilon}{2}, \frac{3}{2} + \frac{\epsilon}{2} \\ 2 - \epsilon, \frac{5}{4}, \frac{7}{4} \end{matrix} \middle| -t^2 \right) \right.$$
$$- \left( \frac{M^2}{m^2} \right)^{1-\epsilon} \frac{\epsilon}{(1 + \epsilon)(1 + 2\epsilon)} {}_4F_3 \left( \begin{matrix} 1, \frac{1}{2} + \epsilon, 1 + \frac{\epsilon}{2}, 1 + \epsilon \\ \frac{3}{2} - \frac{\epsilon}{2}, \frac{3}{4} + \frac{\epsilon}{2}, \frac{5}{4} + \frac{\epsilon}{2} \end{matrix} \middle| -t^2 \right)$$
$$\left. - \left( \frac{M^2}{m^2} \right)^{-\epsilon} \frac{(1 - \epsilon)}{(2 - \epsilon)(3 + 2\epsilon)} {}_4F_3 \left( \begin{matrix} 1, \frac{3}{2} + \frac{\epsilon}{2}, \frac{3}{2} + \epsilon, 1 + \epsilon \\ 2 - \frac{\epsilon}{2}, \frac{5}{4} + \frac{\epsilon}{2}, \frac{7}{4} + \frac{\epsilon}{2} \end{matrix} \middle| -t^2 \right) \right]$$

$$t = \frac{m^2}{2M^2} \quad {}_4F_3^{(1)}(t)$$

- For an all order sunrise with equal masses see [Adams and Weinzierl(2018)]

# Hypergeometric function [Kotikov, Moriello, Campert]

- Hypergeometric function expressed as a one-fold integral

$$\frac{{}_4F_3^{(1)}(t)}{K(\epsilon, t)} = \int_0^1 dx \left[ -\frac{2}{(1-x^2)} \left(\frac{1-x^2}{x^2}\right)^\epsilon + \frac{2}{(1-x^2)ty} \left(\frac{(1-x^2)(1+ty)^2}{4x^2(ty)^2}\right)^\epsilon - \frac{2}{(1-x^2)ty} \left(\frac{(1-x^2)(1+ty)^2}{4x^2(ty)^2}\right)^\epsilon \mathcal{S}_\epsilon \right]$$



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- where we define

$$\mathcal{S}_\epsilon = \sum_{l=0}^{\infty} (-2)^{l+1} \sum_{n=0}^{\infty} \epsilon^{n+l+2} \frac{(-1)^{n+l+1}}{n!(l+1)!} S_{n+1, l+1}$$

$$S_{n+1, l+1} = \int_0^\chi \frac{dz}{z} (\log(z) - \log(\chi))^n (\log(1+z))^{l+1}$$

# Hypergeometric function [Kotikov, Moriello, Campert]

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- where we define

$$\mathcal{S}_\epsilon = \sum_{l=0}^{\infty} (-2)^{l+1} \sum_{n=0}^{\infty} \epsilon^{n+l+2} \frac{(-1)^{n+l+1}}{n!(l+1)!} S_{n+1, l+1}$$

$$S_{n+1, l+1} = \int_0^\chi \frac{dz}{z} (\log(z) - \log(\chi))^n (\log(1+z))^{l+1}$$

## An all order result

$$\begin{aligned} \frac{{}_4F_3^{(1)}(t)}{K(\epsilon, t)} &= -\frac{\Gamma\left(\frac{1}{2} - \epsilon\right)\Gamma(\epsilon)}{\sqrt{\pi}} + P_1 | \sum_{i=0}^{\infty} \epsilon^i \left(\frac{L_A^i}{i!}\right) \\ &\quad - P_1 | \sum_{i=2}^{\infty} \epsilon^i \sum_{k=0}^{i-2} \frac{L_A^k}{k!} \sum_{n=0}^{i-k-2} (-2)^{i-k-n-1} S_{n+1, i-k-n-1}. \end{aligned}$$

- Where the Nielson polylogs in terms of eMPLs

$$S_{n+1, l+1} = \sum_{k=0}^n \binom{n}{k} (-L_{S_A})^{n-k} \left[ P_0 | (L_{S_A}^k L_{S_B}^{l+1}) + I_{k, l}^{(S)}(0) \right]$$

- The concatenation  $|$  is defined as

$$E_4\left(\frac{\vec{n}}{\vec{c}}; x\right) | E_4\left(\frac{\vec{m}}{\vec{d}}; x\right) = E_4\left(\frac{\vec{n} \vec{m}}{\vec{c} \vec{d}}; x\right)$$

# Results in terms of eMPLs

$$P_1 | \sum_{i=0}^{\infty} \epsilon^i \left( \frac{L_A^i}{i!} \right)$$

$$P_1 = E_4 \left( \begin{matrix} -1 \\ -1 \end{matrix}; x \right) - E_4 \left( \begin{matrix} -1 \\ 1 \end{matrix}; x \right),$$

$$\begin{aligned} L_A = & -E_4 \left( \begin{matrix} 1 \\ a_1 \end{matrix}; x \right) - E_4 \left( \begin{matrix} 1 \\ a_2 \end{matrix}; x \right) - E_4 \left( \begin{matrix} 1 \\ a_3 \end{matrix}; x \right) - E_4 \left( \begin{matrix} 1 \\ a_4 \end{matrix}; x \right) - 2E_4 \left( \begin{matrix} -1 \\ -1 \end{matrix}; x \right) \\ & - 2E_4 \left( \begin{matrix} -1 \\ 1 \end{matrix}; x \right) + 3E_4 \left( \begin{matrix} 1 \\ -1 \end{matrix}; x \right) - 2E_4 \left( \begin{matrix} 1 \\ 0 \end{matrix}; x \right) + 3E_4 \left( \begin{matrix} 1 \\ 1 \end{matrix}; x \right) \\ & - \log(t^2 + 1) + 2 \log \left( \sqrt{t^2 + 1} + 1 \right) - \log(4), \end{aligned}$$

## Results in terms of eMPLs

$$S_{n+1,l+1} = \sum_{k=0}^n \binom{n}{k} (-L_{S_A})^{n-k} \left[ P_0 | (L_{S_A}^k L_{S_B}^{l+1}) + I_{k,l}^{(S)}(0) \right]$$

$$P_0 = 2E_4 \left( \begin{matrix} -1 \\ -1 \end{matrix}; x \right) + 2E_4 \left( \begin{matrix} -1 \\ 1 \end{matrix}; x \right)$$

$$L_{S_A} = 2E_4 \left( \begin{matrix} -1 \\ -1 \end{matrix}; x \right) + 2E_4 \left( \begin{matrix} -1 \\ 1 \end{matrix}; x \right) + \log \left( \sqrt{t^2 + 1} - 1 \right) \\ - \log \left( \sqrt{t^2 + 1} + 1 \right)$$






$$L_{S_B} = E_4 \left( \begin{matrix} -1 \\ 1 \end{matrix}; x \right) + E_4 \left( \begin{matrix} -1 \\ -1 \end{matrix}; x \right) - E_4 \left( \begin{matrix} 1 \\ -1 \end{matrix}; x \right) + \frac{1}{2} E_4 \left( \begin{matrix} 1 \\ a_1 \end{matrix}; x \right) \\ + \frac{1}{2} E_4 \left( \begin{matrix} 1 \\ a_2 \end{matrix}; x \right) + \frac{1}{2} E_4 \left( \begin{matrix} 1 \\ a_3 \end{matrix}; x \right) + \frac{1}{2} E_4 \left( \begin{matrix} 1 \\ a_4 \end{matrix}; x \right) - E_4 \left( \begin{matrix} 1 \\ 1 \end{matrix}; x \right) \\ + \frac{1}{2} \log(t^2 + 1) - \log \left( \sqrt{t^2 + 1} + 1 \right) + \log(2)$$

# Summary

- General solution strategy for linearly reducible elliptic Feynman integrals.
- Solutions of Banana like diagrams.
- An all order epsilon result for a sunrise diagram.
- Future perspective: This strategy can be used to calculate diagrams of phenomenological interest.






Thank you for your attention :)

# Bibliography I

-  F. Brown and A. Levin, *Multiple elliptic polylogarithms*, (2013), arXiv:hep-ph/11106917 [hep-ph] .
-  J. Broedel, *Elliptic multiple zeta values and one-loop superstring amplitudes*, JHEP **07**, 112 (2015), arXiv:1412.5535 [hep-th] .
-  S. Weinzierl, *The iterated structure of the all-order result for the two-loop sunrise integral*, J. Math. Phys. **57**, 032304 (2016), arXiv:1512.05630 [hep-ph] .
-  C. Duhr et al., *Elliptic polylogarithms and iterated integrals on elliptic curves. Part I: general formalism*, JHEP **05**, 093 (2018), arXiv:1712.07089 [hep-th] .
-  F. C. S. Brown, *On the periods of some Feynman integrals*, arXiv e-prints , arXiv:0910.0114 (2009), arXiv:0910.0114 [math.AG] .



# Bibliography II

-  E. Panzer, *Feynman integrals and hyperlogarithms*, Ph.D. thesis, Humboldt U., Berlin, Inst. Math. (2015), arXiv:1506.07243 [math-ph] .
-  M. Hidding and F. Moriello, *All orders structure and efficient computation of linearly reducible elliptic Feynman integrals*, JHEP **01**, 169 (2019), arXiv:1712.04441 [hep-ph] .
-  H. Cheng and T. T. Wu, *Expanding protons: Scattering at high-energies*, Cambridge, USA: MIT-PR. (1987) 285p (1987).
-  L. Adams and S. Weinzierl, Commun. Num. Theor. Phys. **12**, 193 (2018), arXiv:1704.08895 [hep-ph] .
-  M. Yu. Kalmykov and B. A. Kniehl, *Towards all-order Laurent expansion of generalized hypergeometric functions around rational values of parameters*, Nucl. Phys. **B809**, 365 (2009), arXiv:0807.0567 [hep-th] .