

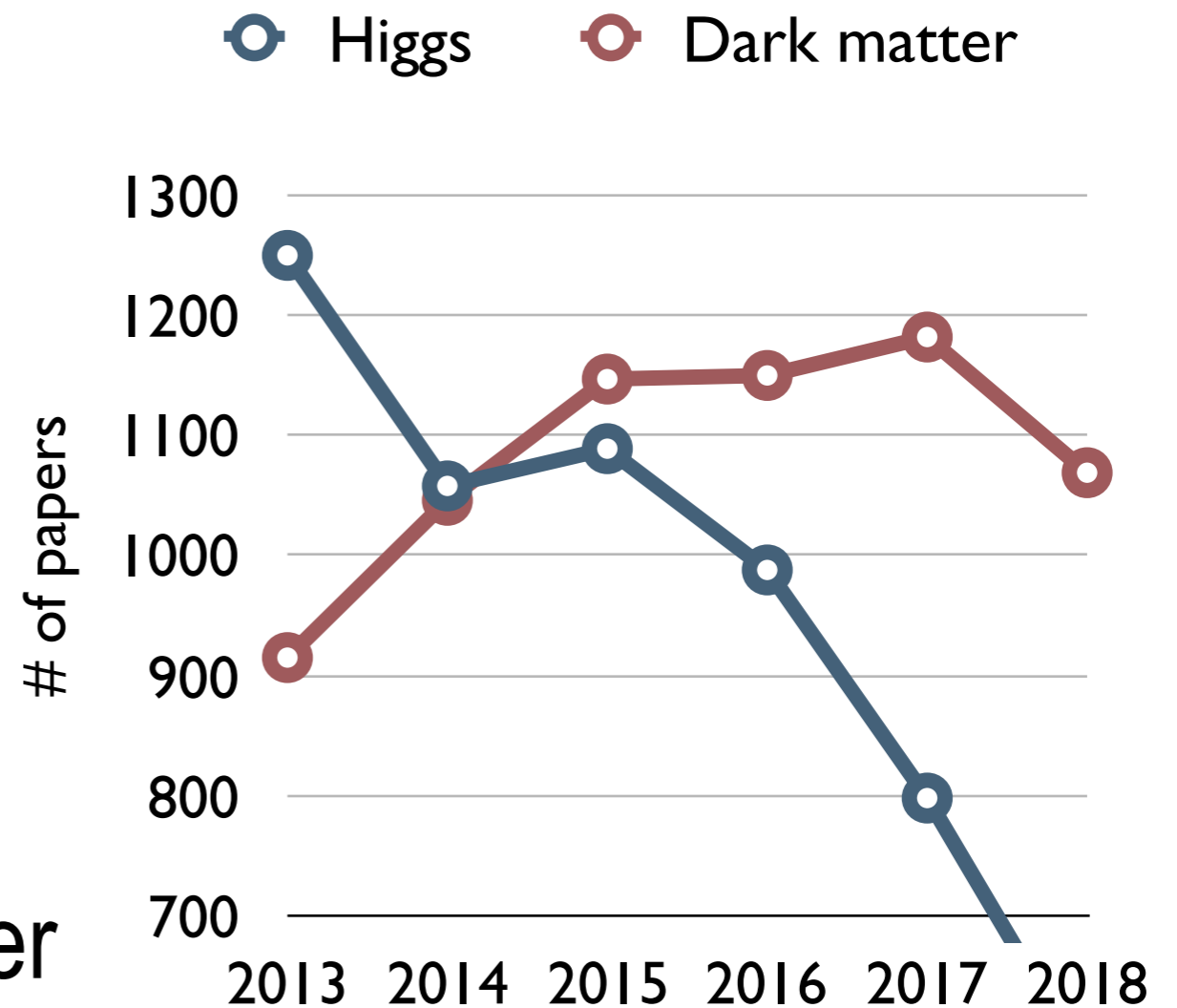
# Searching for dark matter at colliders

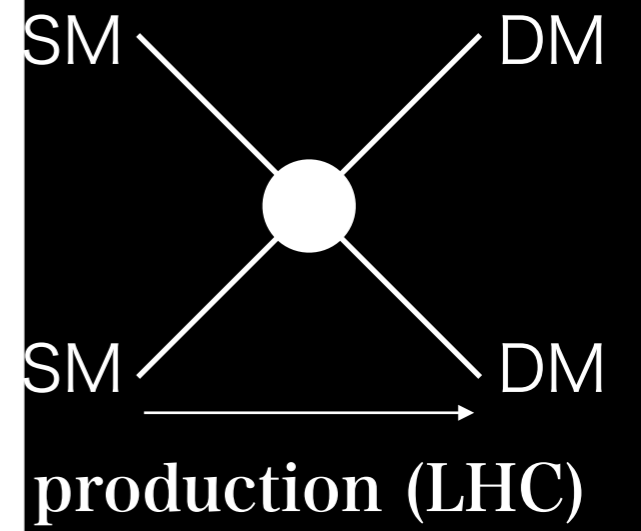
Kentarou Mawatari



# My main research field is collider phenomenology

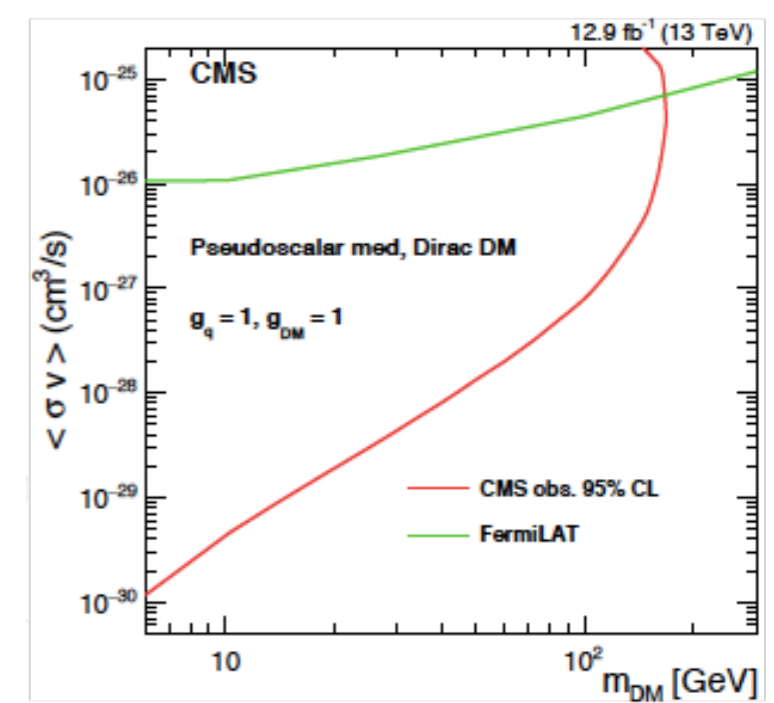
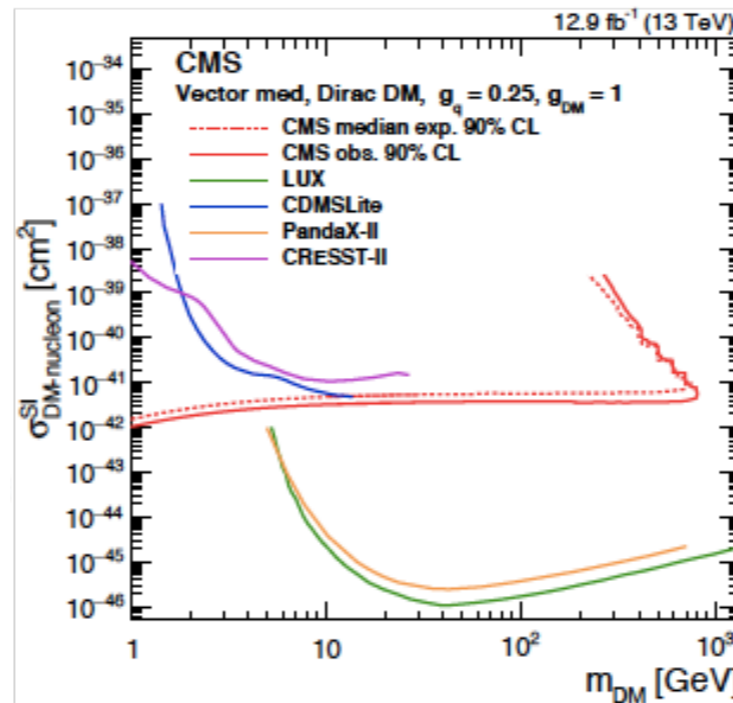
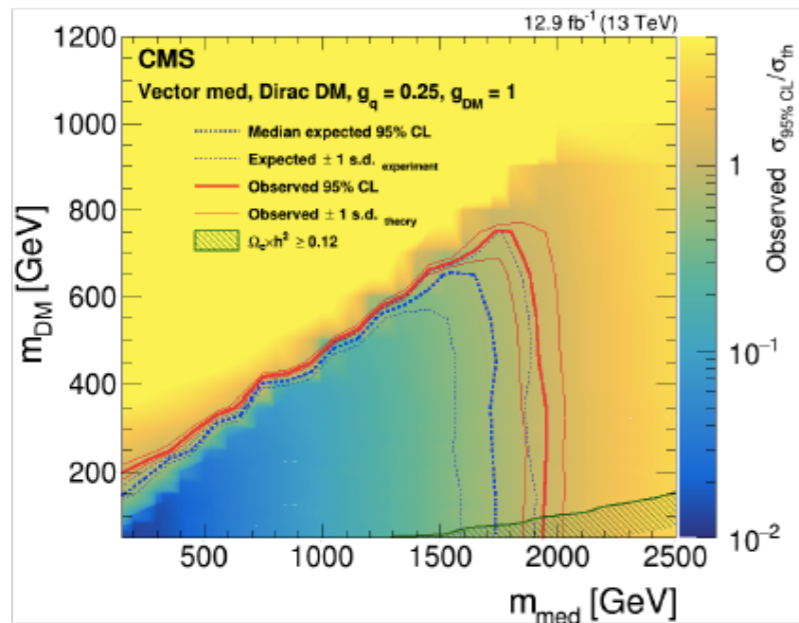
- the Higgs sector
- dark matter (DM)
- inSPIRE
  - find title higgs
  - find title dark matter





# Search for dark matter produced with an energetic jet or a hadronically decaying W or Z boson at $\sqrt{s} = 13$ TeV

The CMS Collaboration\*

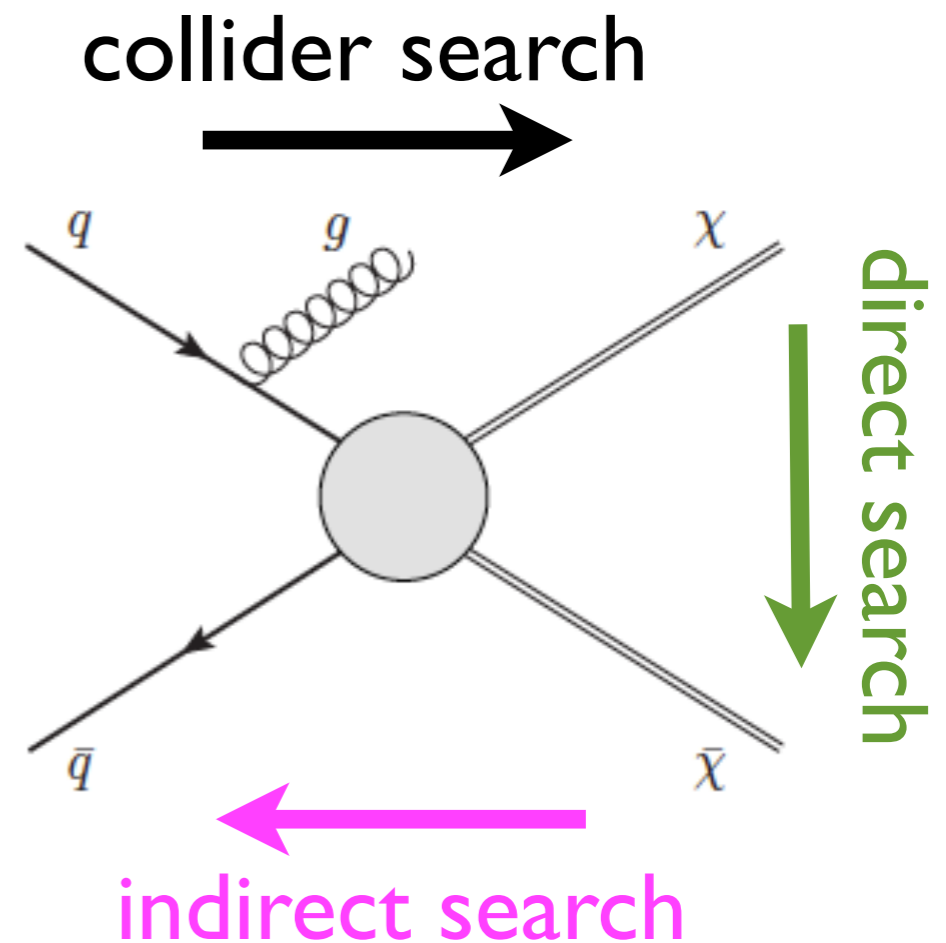


ing fraction. The results of this search provide the strongest constraints on the dark matter pair production cross section through vector and axial-vector mediators at a particle collider. When compared to the direct detection experiments, the limits obtained from this search provide stronger constraints for dark matter masses less than 5, 9, and 550 GeV, assuming vector, scalar, and axial-vector mediators, respectively. The search yields stronger constraints for dark matter masses less than 200 GeV, assuming a pseudoscalar mediator, when compared to the indirect detection results from Fermi-LAT.

arXiv:1703.01651v2 [hep-ex]

# Looking for dark matter

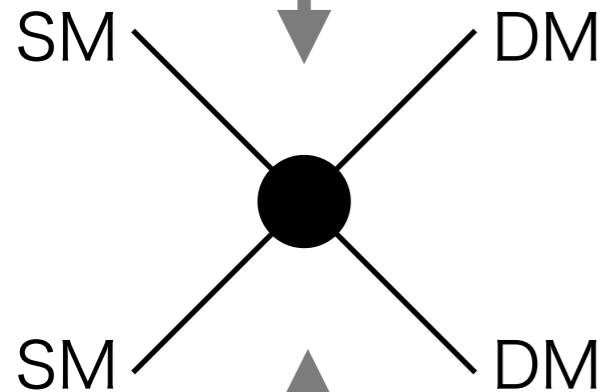
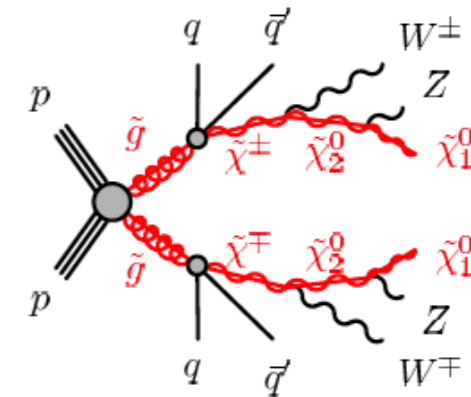
- If dark matter has non-gravitational interactions with ordinal matter, we can observe it directly and/or indirectly.
- Indeed, many types of dark matter search experiments are currently on-going all over the world.



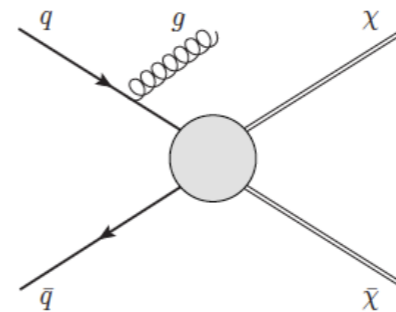
# DM (or MET) searches at LHC Run-I

Top-down approach

**UV model** : SUSY, ExtraDim, ...  
 $\{\mathcal{L}; m_{\text{DM}}, m_1, m_2, \dots, g_1, g_2, \dots\}$



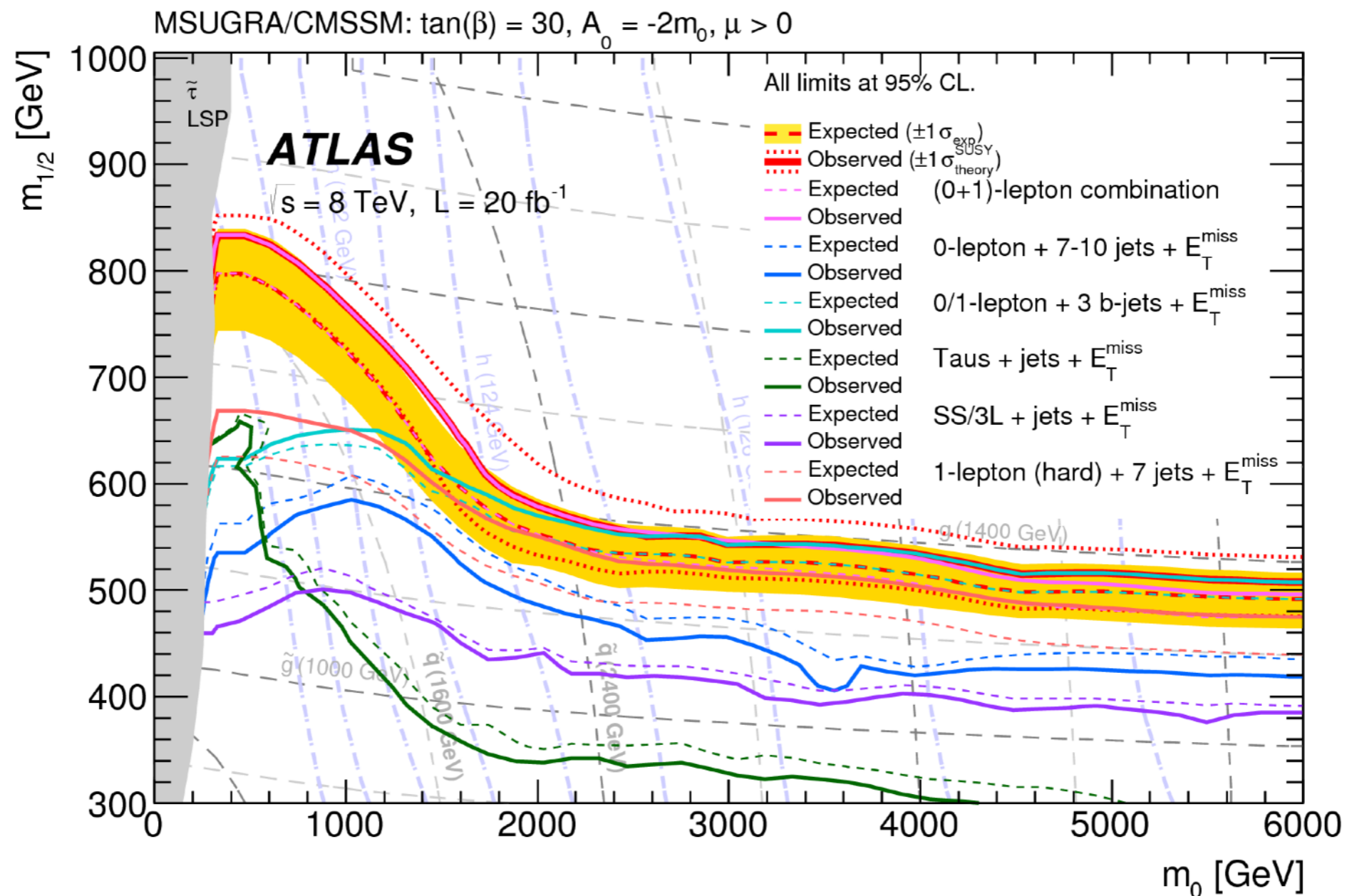
**EFT** : SM + DM particles  
 $\{\mathcal{L}; m_{\text{DM}}, M_*\}$   $\mathcal{L} = \frac{1}{M_*^2} \bar{\chi} \Gamma^\mu \chi \bar{q} \Gamma_\mu q$



Bottom-up approach

$\tilde{\chi}_1^0$ 

# DM (or MET) searches at LHC Run-I: top-down (e.g. Supersymmetry)

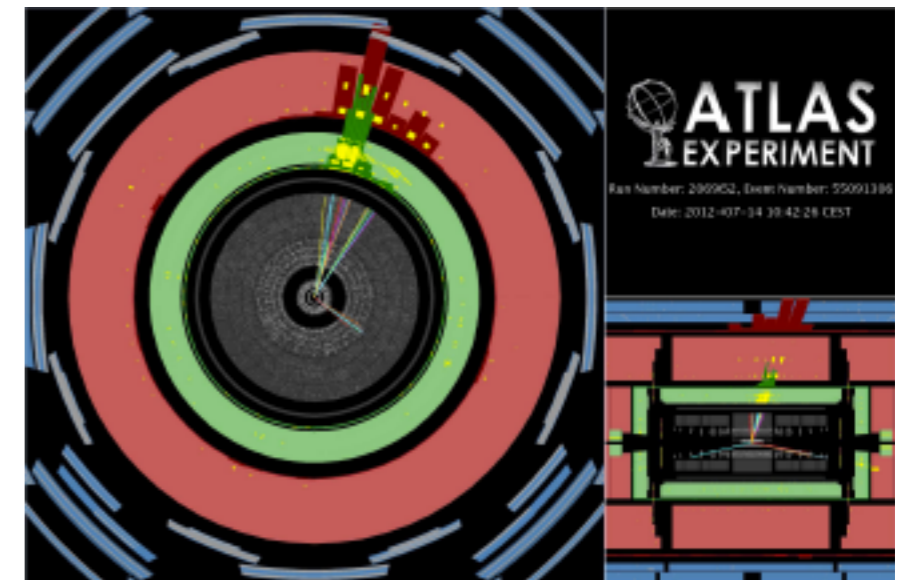
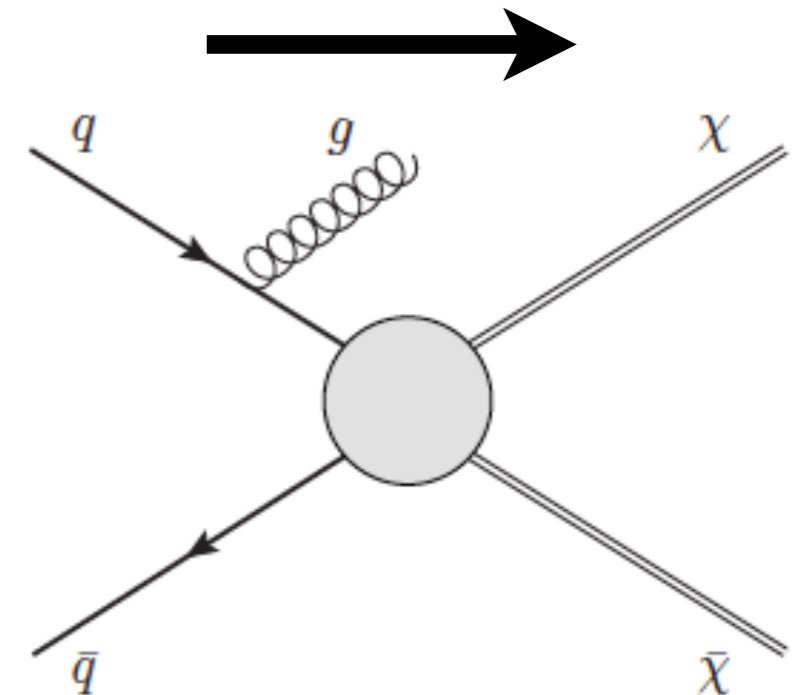


predictive, but many parameters and many final states...

# DM (or MET) searches at LHC Run-I: bottom-up (Effective field theory)

- employed contact interaction operators in EFTs (effective field theories).
- vector  $\frac{1}{M_*^2} \bar{\chi} \gamma^\mu \chi \bar{q} \gamma_\mu q$
- axial-vector  $\frac{1}{M_*^2} \bar{\chi} \gamma^\mu \gamma^5 \chi \bar{q} \gamma_\mu \gamma^5 q$
- The signal is determined by the Lorentz structure, the DM mass, and the overall coupling (or the cutoff scale).
- Simple final states
- easy interpretation to non-collider DM searches
- 

mono- $X$  search





# DM (or MET) searches at LHC Run-I: bottom-up (Effective field theory)

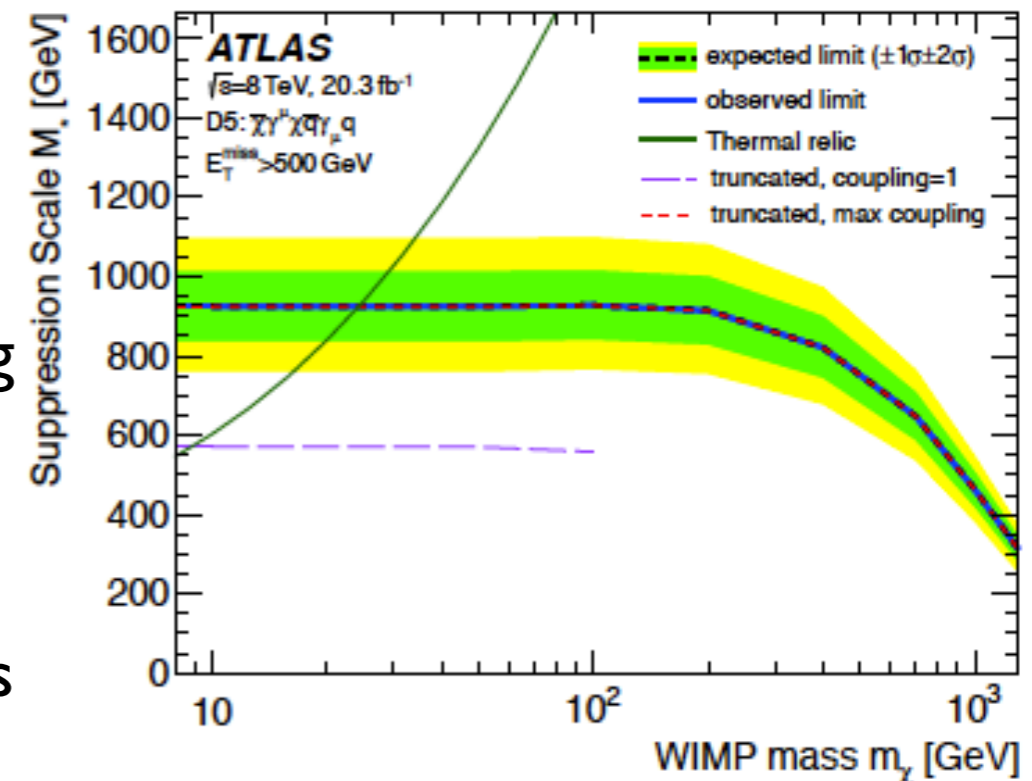
- employed contact interaction operators in EFTs (effective field theories).

- vector  $\frac{1}{M_\star^2} \bar{\chi} \gamma^\mu \chi \bar{q} \gamma_\mu q$

- axial-vector  $\frac{1}{M_\star^2} \bar{\chi} \gamma^\mu \gamma^5 \chi \bar{q} \gamma_\mu \gamma^5 q$

- The signal is determined by the Lorentz structure, the DM mass, and the overall coupling (or the cutoff scale).
- Simple final states
- easy interpretation to non-collider DM searches
- **EFT validation;  $M_\star \leq$  (LHC accessible energy)**

[1502.01518]



CERN, Dec 10-11 2015

The WG activity builds on the experience of the previous ATLAS-CMS Dark Matter Forum, whose findings are documented in <http://arxiv.org/abs/1507.00966>

## The WG

- brings together theorists and experimentalists to define guidelines and recommendations for the benchmark models, interpretation, and characterisation necessary for broad and systematic searches for dark matter at the LHC.
- develops and promotes well-defined signal models, specifying the assumptions behind them and describing the conditions under which they should be used.
- works to improve the set of tools available to the experiments, such as higher-precision calculations of the backgrounds.
- assists theorists with understanding and making use of LHC results.
- develops and maintains close connections with theorists and other experimental particle DM searches (e.g. Direct and Indirect Detection experiments) in order to help verify and constrain particle physics models of astrophysical excesses, to understand how collider searches and non-collider experiments complement one another, and to help build a comprehensive understanding of viable dark matter models.

# LHC DMWG documents

- [arXiv:1507.00966](#)  
“DM benchmark models for early LHC Run-2”  
(Report of the ATLAS/CMS Dark Matter Forum)
- [arXiv:1603.04156](#)  
“Recommendations on presenting LHC searches for MET signals using simplified s-channel models of DM”
- [arXiv:1703.05703](#)  
“Comparing LHC searches for heavy mediators of DM production in visible and invisible decay channels”
- [arXiv:1705.04664](#)  
“Precise predictions for V+jets DM backgrounds”

# Summary of the LHC DMWG activities

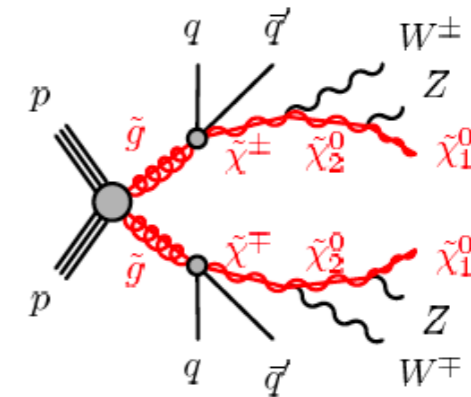
- develops well-defined benchmark models
- provides (accurate) simulation tools for signal and background
- defines guidelines for presentation and interpretation of the LHC results

# DM (or MET) searches at LHC Run-II

Top-down approach

**UV model** : SUSY, ExtraDim, ...

$$\{\mathcal{L}; m_{\text{DM}}, m_1, m_2, \dots, g_1, g_2, \dots\}$$

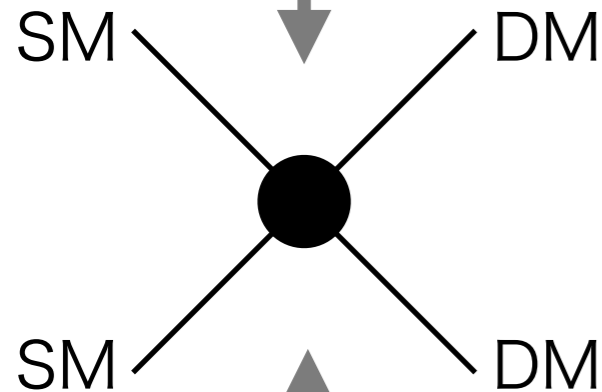
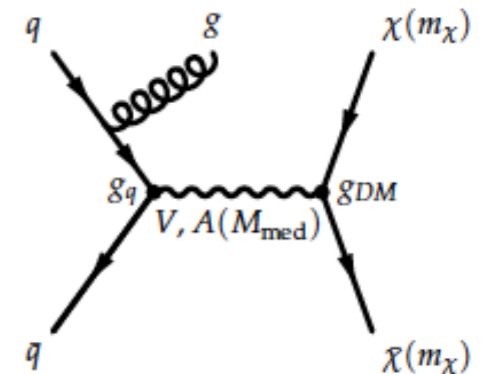
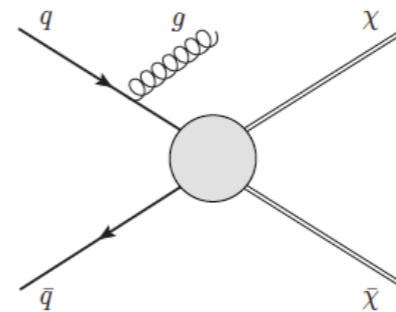


**Simplified model** : SM + DM + Mediator particles

$$\{\mathcal{L}; m_{\text{DM}}, m_{\text{med}}, g_{\text{DM}}, g_q\} \quad \mathcal{L} = g_{\text{DM}} Z'_\mu \bar{\chi} \gamma^\mu \chi + g_q Z'_\mu \bar{q} \gamma^\mu q$$

**EFT** : SM + DM particles

$$\{\mathcal{L}; m_{\text{DM}}, M_*\} \quad \mathcal{L} = \frac{1}{M_*^2} \bar{\chi} \Gamma^\mu \chi \bar{q} \Gamma_\mu q$$



Bottom-up approach

# Simplified DM models

- Run-II is employing simplified DM models.

- $$\mathcal{L}_{\text{vector}} = g_q \sum_{q=u,d,s,c,b,t} Z'_\mu \bar{q} \gamma^\mu q + g_\chi Z'_\mu \bar{\chi} \gamma^\mu \chi$$

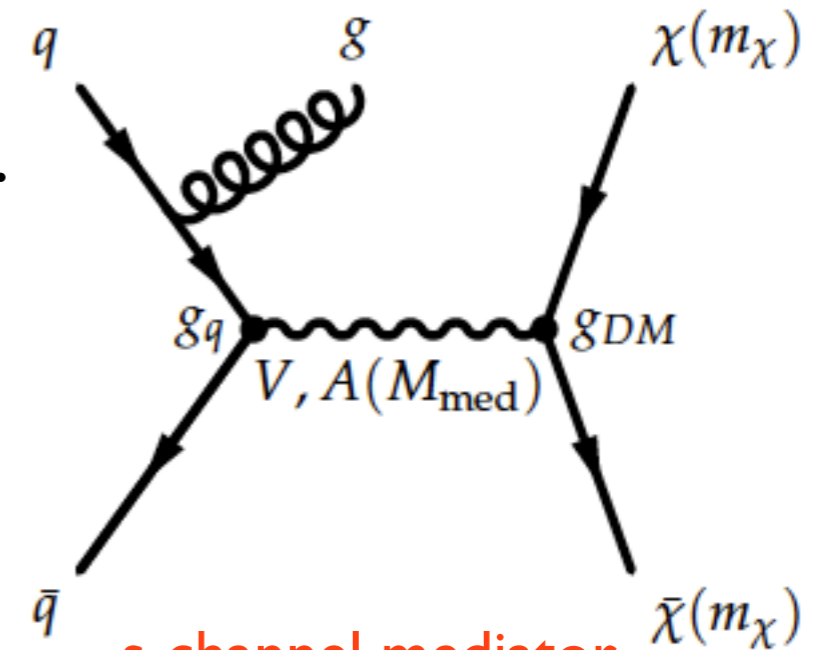
- $$\mathcal{L}_{\text{axial-vector}} = g_q \sum_{q=u,d,s,c,b,t} Z'_\mu \bar{q} \gamma^\mu \gamma^5 q + g_\chi Z'_\mu \bar{\chi} \gamma^\mu \gamma^5 \chi$$

- The signal is determined by the mediator type, the DM and mediator masses, and one or two coupling(s).

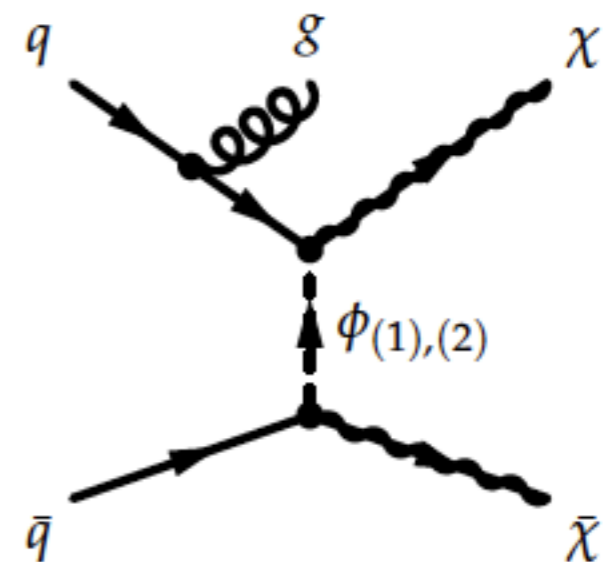
- Richer phenomenology

- Interpretations to non-collider DM searches are more complicated.

DM Forum [1507.00966]



s-channel mediator



t-channel mediator

# s-channel simplified models: spin-0, 1

- Simplified DM models (s-channel):

- spin-1 mediator

$$\mathcal{L}_{\text{vector}} = -g_{\text{DM}} Z'_\mu \bar{\chi} \gamma^\mu \chi - g_q \sum_{q=u,d,s,c,b,t} Z'_\mu \bar{q} \gamma^\mu q,$$

$$\mathcal{L}_{\text{axial-vector}} = -g_{\text{DM}} Z'_\mu \bar{\chi} \gamma^\mu \gamma_5 \chi - g_q \sum_{q=u,d,s,c,b,t} Z'_\mu \bar{q} \gamma^\mu \gamma_5 q.$$

- spin-0 mediator

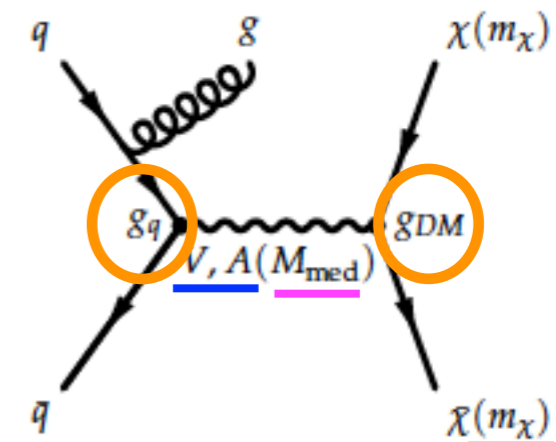
$$\mathcal{L}_{\text{scalar}} = -g_{\text{DM}} \phi \bar{\chi} \chi - g_q \frac{\phi}{\sqrt{2}} \sum_{q=u,d,s,c,b,t} y_q \bar{q} q,$$

$$\mathcal{L}_{\text{pseudo-scalar}} = -ig_{\text{DM}} \phi \bar{\chi} \gamma_5 \chi - ig_q \frac{\phi}{\sqrt{2}} \sum_{q=u,d,s,c,b,t} y_q \bar{q} \gamma_5 q,$$

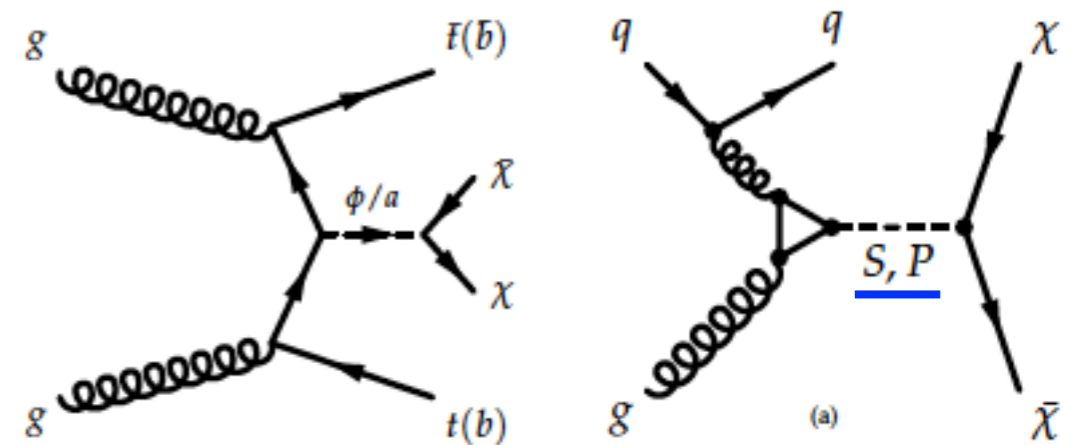
- The signal is determined by

- the mediator type (V, A, S, P)
- the DM and mediator masses
- the two couplings

DM Forum [1507.00966]



spin-1 mediator



spin-0 mediator

👉 Top-philic DM models

# s-channel simplified models: spin-2

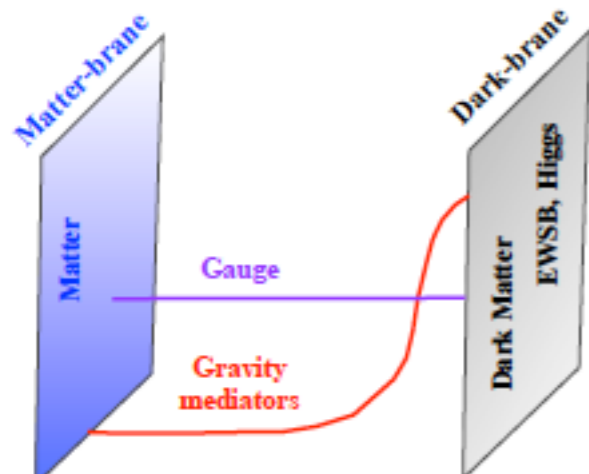
Gravity-mediated DM:

Lee, Park, Sanz [1306.4107, 1401.5301]

Rueter, Rizzo, Hewett [1706.07540]

Das, Degrande, Hirschi, Maltoni, Shao [1605.09359]

Kraml, Laa, KM, Yamashita [1701.07008, EPJC]



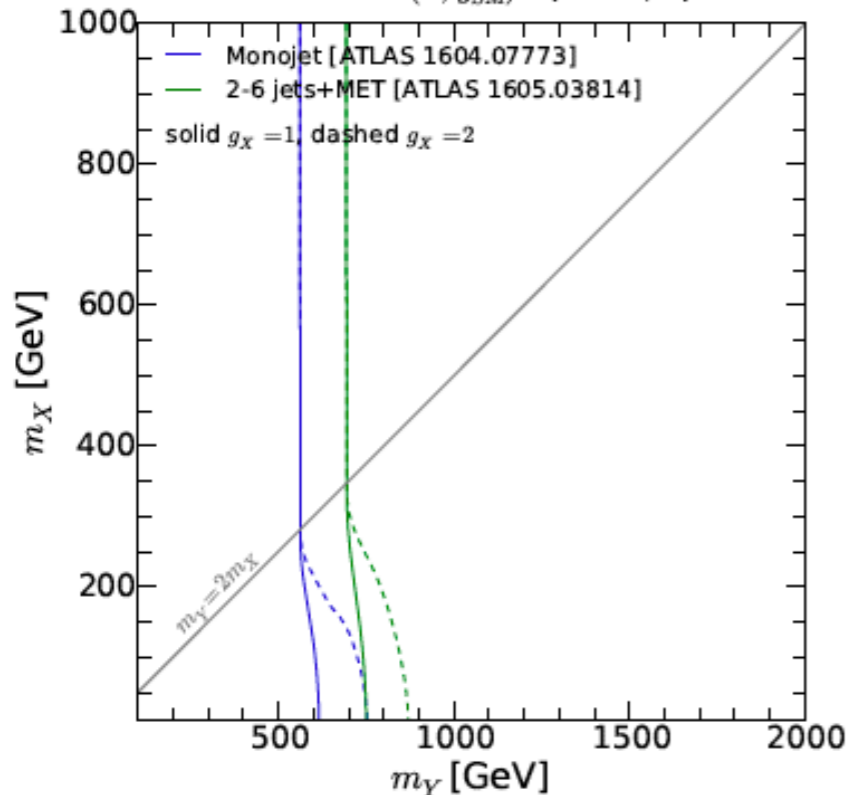
$$\mathcal{L}_X^{Y_2} = -\frac{1}{\Lambda} g_X^T T_{\mu\nu}^X Y_2^{\mu\nu}$$

$$\mathcal{L}_{SM}^{Y_2} = -\frac{1}{\Lambda} \sum_i g_i^T T_{\mu\nu}^i Y_2^{\mu\nu}$$

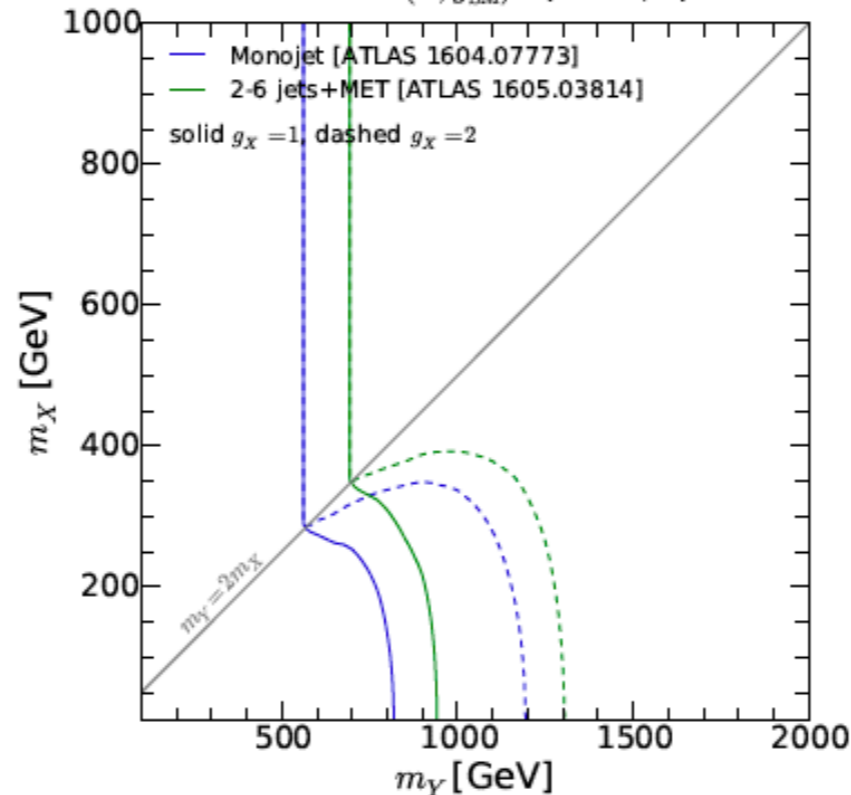
$$g_{SM} \equiv g_H^T = g_q^T = g_l^T = g_g^T = g_W^T = g_B^T$$

$$\{m_X, m_Y, g_X/\Lambda, g_{SM}/\Lambda\}$$

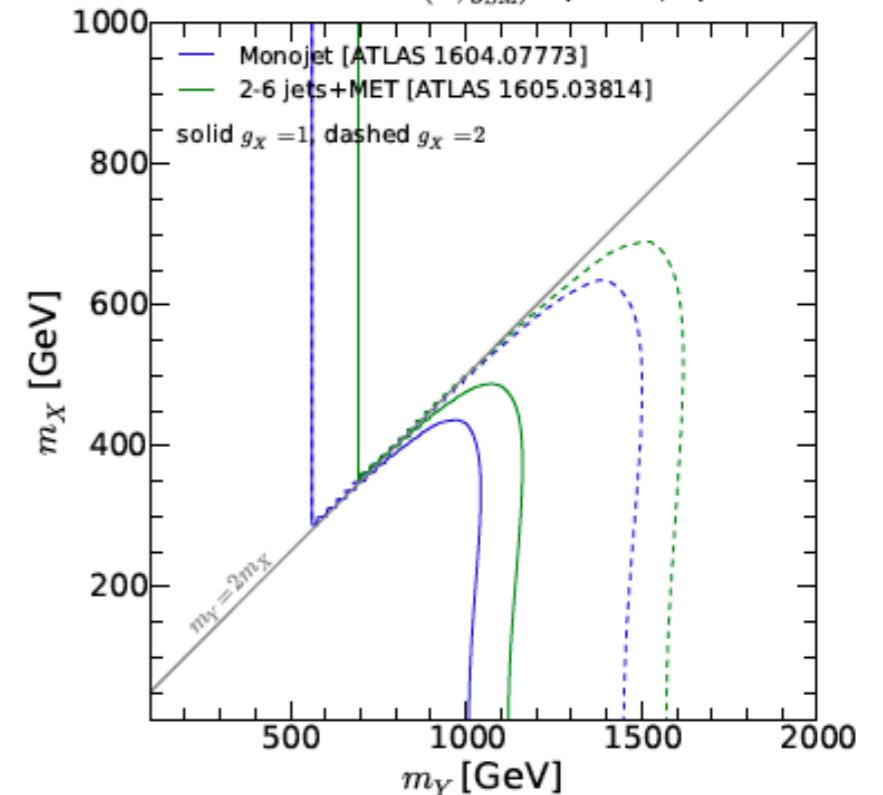
Scalar DM ( $\Lambda, g_{SM}$ ) = (3 TeV, 1)



Dirac DM ( $\Lambda, g_{SM}$ ) = (3 TeV, 1)



Vector DM ( $\Lambda, g_{SM}$ ) = (3 TeV, 1)



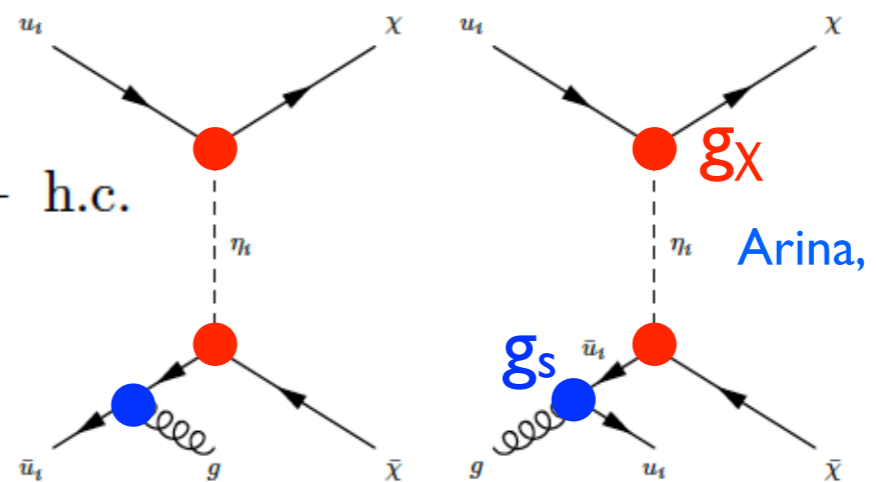


# t-channel simplified models: colored scalar

See, e.g., De Simone, Jacques [1603.08002, EPJC]

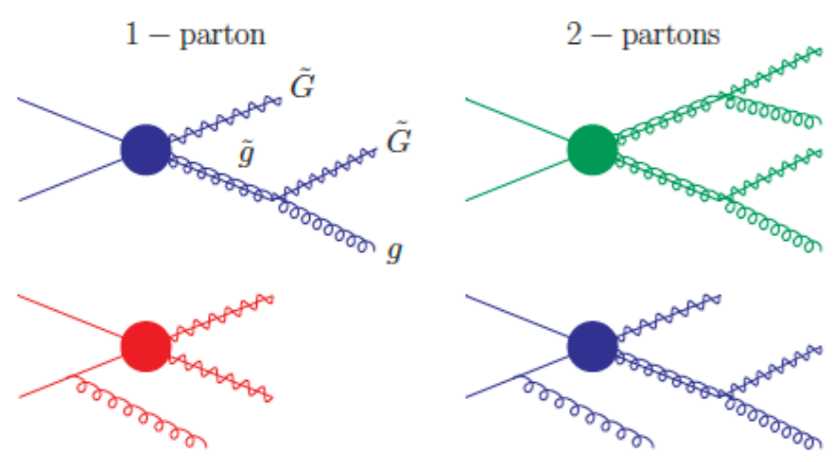
$$\mathcal{L} = \sum_{i=1,2,3} g_i \chi \left( \bar{Q}_L^{(i)} \eta_L^{(i)} + \bar{u}_R^{(i)} \eta_{u,R}^{(i)} + \bar{d}_R^{(i)} \eta_{d,R}^{(i)} \right) + \text{h.c.}$$

$\{m_\chi, M, g\}$

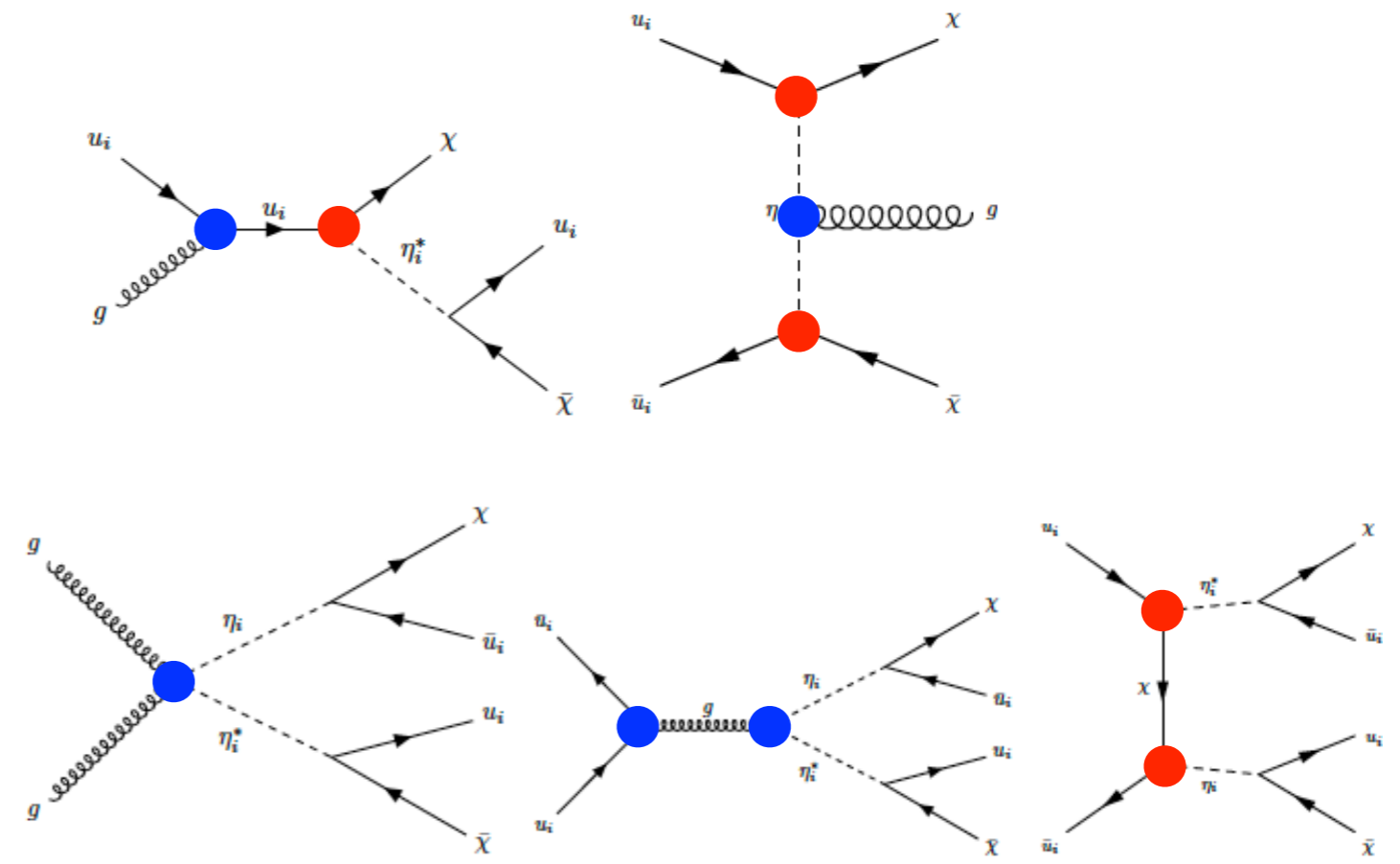


Arina, Fuks, Mantani [1911.?????]

Gravitino production at the LHC  
Maltoni, Martini, KM, Oexl [1502.01637, JHEP]



ME+PS merging,  
proper treatment of resonances,  
...



# Summary of the LHC DMWG activities

- develops well-defined benchmark models
- provides (accurate) simulation tools for signal and background
- defines guidelines for presentation and interpretation of the LHC results

Search for new physics in final states with an energetic jet or a hadronically decaying  $W$  or  $Z$  boson and transverse momentum imbalance at  $\sqrt{s} = 13$  TeV

The CMS Collaboration\*

### 3 Simulated samples

To model the SM backgrounds, simulated Monte Carlo (MC) samples are produced for the  $Z$ +jets,  $W$ +jets,  $\gamma$ +jets, and quantum chromodynamics (QCD) multijet processes at leading order (LO) using the MADGRAPH5\_aMC@NLO 2.2.2 [56] generator and are generated with up to four additional partons in the matrix element calculations. The samples for the  $t\bar{t}$  and single top quark background processes are produced at next-to-leading order (NLO) using POWHEG 2.0 and 1.0, respectively [57, 58], and the set of diboson ( $WW$ ,  $WZ$ ,  $ZZ$ ) samples is produced at LO with PYTHIA 8.205 [59].

Vector and axial-vector monojet and mono- $V$  dark matter signals are simulated at NLO using the DMSIMP models [60, 61] with the MADGRAPH5\_aMC@NLO generator. Both scalar and pseudoscalar monojet production contain gluon-initiated loop processes. These samples are generated at LO with one or two additional partons in the matrix element calculations, taking into account finite top mass effects using the MADGRAPH5\_aMC@NLO generator in conjunction with the DMSIMP models.

# Constraints on mediator-based dark matter and scalar dark energy models using $\sqrt{s} = 13$ TeV $pp$ collision data collected by the ATLAS detector

The ATLAS Collaboration

Table 4: Details of the generation setup and Universal FeynRules Output (UFO) model used for the spin-1 mediator simplified models, for each signature considered in this paper.

Model and Final State	UFO	Generator and Parton Shower	Cross-section	Additional details
$Z'(\chi\bar{\chi}) + j$	DMV [26, 215]	POWHEG-BOX v2 [216] + PYTHIA 8.205 [217]	NLO	Particle-level rescaling of leptophobic $Z'_A$ scenario of Ref. [26] (see Appendix A.1)
$Z'(\chi\bar{\chi}) + \gamma$	<u>DMSimp</u> [116, 218]	MG5_AMC@NLO 2.4.3 (NLO) [213] + PYTHIA 8.212	NLO	Leptophobic $Z'_A$ scenario simulated, other scenarios obtained by cross-section rescaling (see Appendix A.1)
$Z'(\chi\bar{\chi}) + V$	<u>DMSimp</u>	MG5_AMC@NLO 2.5.3 (NLO) + PYTHIA 8.212	NLO	Particle-level rescaling of LO samples of Ref. [20] to each of the four NLO scenarios (see Appendix A.1)
$Z'(qq)$ or $Z'(qq)+ISR$	<u>DMSimp</u>	MG5_AMC@NLO 2.2.3 (NLO) + PYTHIA 8.210	NLO	Leptophobic $Z'_A$ scenario simulated, other scenario obtained by Gaussian resonance limits and cross-section rescaling [214]
$Z'(b\bar{b})$	<u>DMSimp</u>	MG5_AMC@NLO 2.2.3 (NLO) + PYTHIA 8.210	NLO	Leptophobic $Z'_A$ scenario simulated, other scenario obtained by Gaussian resonance limits and cross-section rescaling [214]

## Simplified dark matter models

### Authors

- s-channel (spin-0 and spin-1)
  - Antony Martini (Université catholique de Louvain) & Kentarou Mawatari (LPSC Grenoble)
    - Emails: kentarou.mawatari @ lpsc.in2p3.fr
- s-channel (spin-0 and spin-1 electroweak)
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    - Emails: cenzhang @ bnl.gov
- s-channel (spin-2)
  - Goutam Das (Saha Inst.), Celine Degrande (CERN) & Kentarou Mawatari (LPSC Grenoble)
    - Emails: goutam.das @ saha.ac.in, celine.degrande @ cern.ch, kentarou.mawatari @ lpsc.in2p3.fr

### Description of the model

This is simplified dark matter models for NLO. Our lagrangian consists of different types of DM:

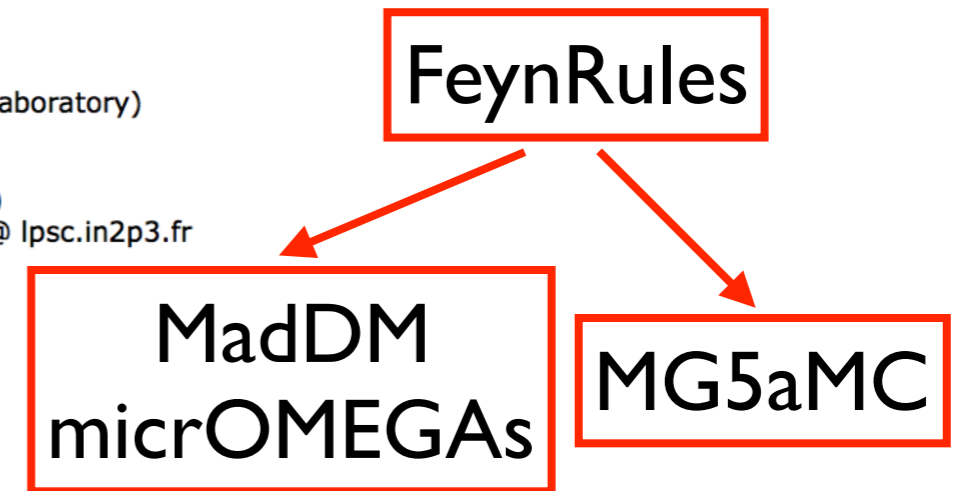
- $X_r$  (real scalar DM)
- $X_c$  (complex scalar DM)
- $X_d$  (Dirac spinor DM)
- $X_m$  (Majorana spinor DM) [to be done.]
- $X_v$  (vector DM)
- ...

and different types of mediators:

- s-channel
  - $Y_0$  (spin-0)
  - $Y_1$  (spin-1)
  - $Y_2$  (spin-2)
  - ...
- t-channel [to be done.]

See more details in

- [1508.00564](#) : O. Mattelaer, E. Vryonidou, "Dark matter production through loop-induced processes at the LHC: the s-channel mediator case" (EPJC75(2015)436).
- [1508.05327](#) : M. Backovic, M. Kramer, F. Maltoni, A. Martini, K. Mawatari, M. Pellen, "Higher-order QCD predictions for dark matter production at the LHC in simplified models with s-channel mediators" (EPJC75(2015)482).
- [1509.05785](#) : M. Neubert, J. Wang, C. Zhang, "Higher-order QCD predictions for dark matter production in mono-Z searches at the LHC" (JHEP1602(2016)082).
- [1605.09359](#) : G. Das, C. Degrande, V. Hirschi, F. Maltoni, H. Shao, "NLO predictions for the production of a spin-two particle at the LHC"
- [1701.07008](#) : S. Kraml, U. Laa, K. Mawatari, K. Yamashita, "Simplified dark matter models with a spin-2 mediator at the LHC".



# I-min MadGraph5\_aMC@NLO tutorial

```
./bin/mg5_aMC
>import model DMSimp_s_spin1
>generate p p > xd xd~ j [QCD]
>output
>launch
```

- ➔ Start the MG5\_aMC shell
- ➔ Import the model
- ➔ Generate the process
- ➔ Write the code (including html)
- ➔ Generate the LO/NLO events

param\_card.dat

run\_card.dat

```
#####
## INFORMATION FOR DMINPUTS
#####
Block dminputs
 1 0.000000e+00 # gVXc
 2 1.000000e+00 # gVXd
 3 0.000000e+00 # gAXd
 4 2.500000e-01 # gVd11
 5 2.500000e-01 # gVu11
 6 2.500000e-01 # gVd22
 7 2.500000e-01 # gVu22
 8 2.500000e-01 # gVd33
 9 2.500000e-01 # gVu33
10 0.000000e+00 # gAd11
11 0.000000e+00 # gAu11
12 0.000000e+00 # gAd22
13 0.000000e+00 # gAu22
14 0.000000e+00 # gAd33
15 0.000000e+00 # gAu33
```

$$\mathcal{L}_{X_D}^{Y_1} = \bar{X}_D \gamma_\mu (g_{X_D}^V + g_{X_D}^A \gamma_5) X_D Y_1^\mu$$

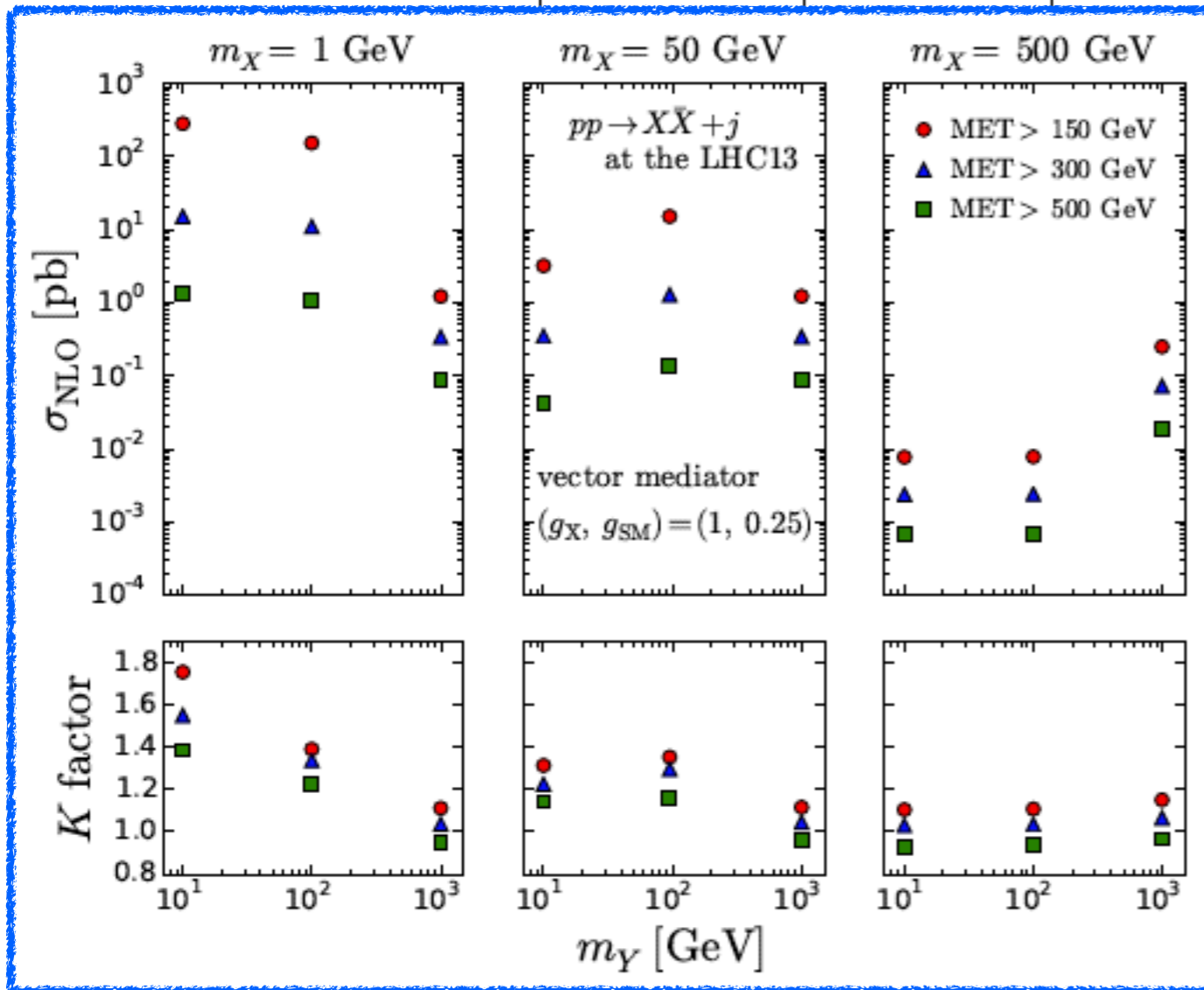
$$\mathcal{L}_{SM}^{Y_1} = \sum_{i,j} [\bar{q}_i \gamma_\mu (g_{qij}^V + g_{qij}^A \gamma_5) q_j] Y_1^\mu$$

```
#####
## INFORMATION FOR MASS
#####
Block mass
 6 1.720000e+02 # MT
15 1.777000e+00 # MTA
23 9.118760e+01 # MZ
25 1.250000e+02 # MH
51 1.000000e+01 # MXc
52 1.000000e+01 # MXd
55 1.000000e+03 # MY1
5000001 1.000000e+01 # MXr
```

```
#####
# Collider type and energy
#####
 1 = lpp1 ! beam 1 type (0 = no PDF)
 1 = lpp2 ! beam 2 type (0 = no PDF)
6500 = ebeam1 ! beam 1 energy in GeV
6500 = ebeam2 ! beam 2 energy in GeV
#####
# PDF choice: this automatically fixes also alpha_s(MZ) and its evol.
#####
nn23nlo = pdlabel ! PDF set
230000 = lhaid ! if pdlabel=lhapdf, this is the lhapdf number
#####
# Include the NLO Monte Carlo subtr. terms for the following parton
# shower (HERWIG6 | HERWIGPP | PYTHIA6Q | PYTHIA6PT | PYTHIA8)
# WARNING: PYTHIA6PT works only for processes without FSR!!!!
#####
HERWIG6 = parton_shower
```

# Cross sections for mono-j at LHC13

$(m_Y, m_X)$ [GeV]			MET > 150 GeV	MET > 300 GeV
(100, 1)	$m_Y > 2m_X$	$\sigma_{\text{LO}}$ [pb]	$1.100 \times 10^2$ $^{+10.6}_{-9.3} \pm 1.5\%$	$0.822 \times 10^1$ $^{+14.4}_{-12.0} \pm 1.1\%$
		$\sigma_{\text{NLO}}$ [pb]	$1.530 \times 10^2$ $^{+6.5}_{-5.7} \pm 0.5\%$	$1.100 \times 10^1$ $^{+7.4}_{-7.2} \pm 0.6\%$
		$K$ factor	1.39	1.34



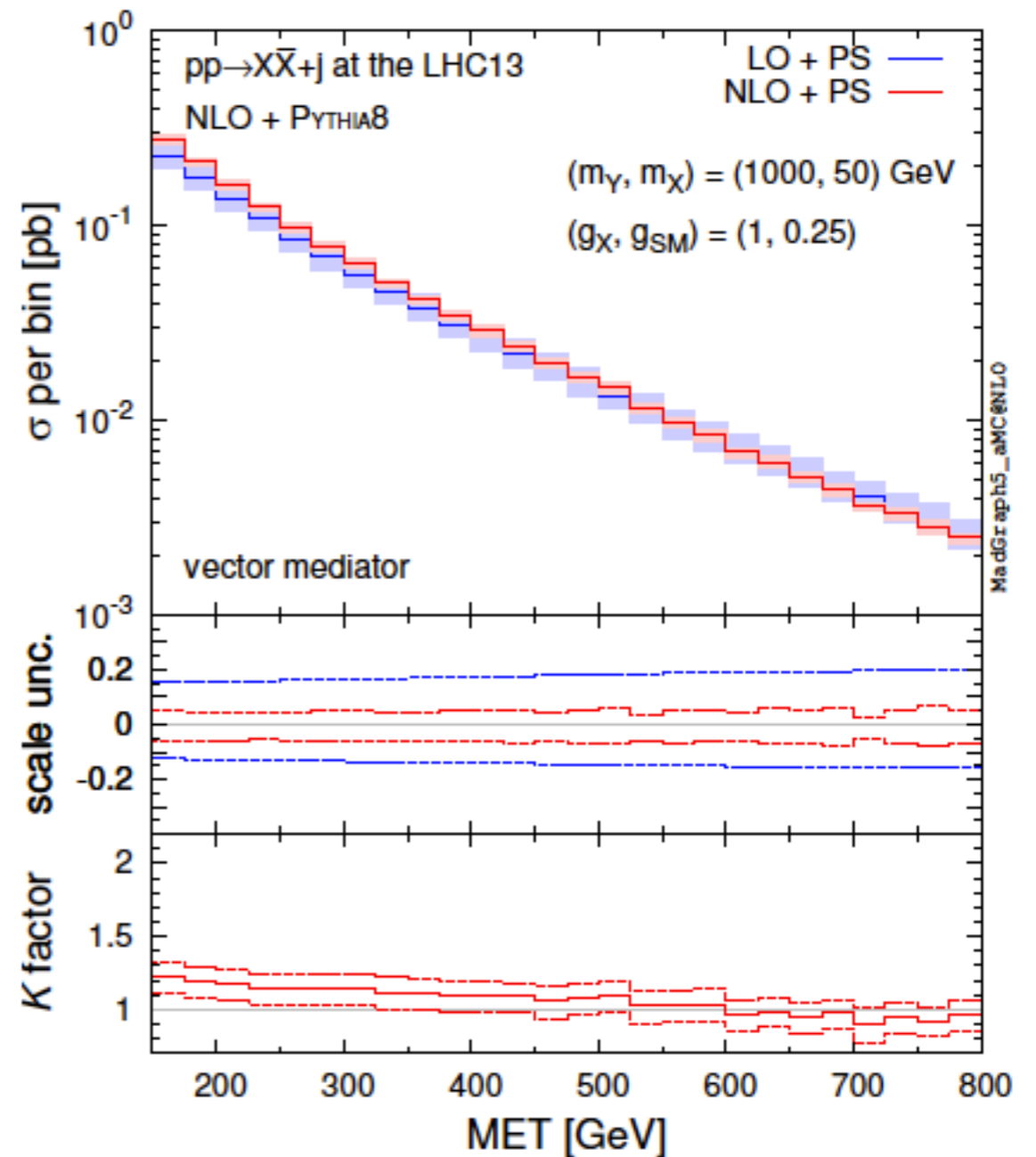
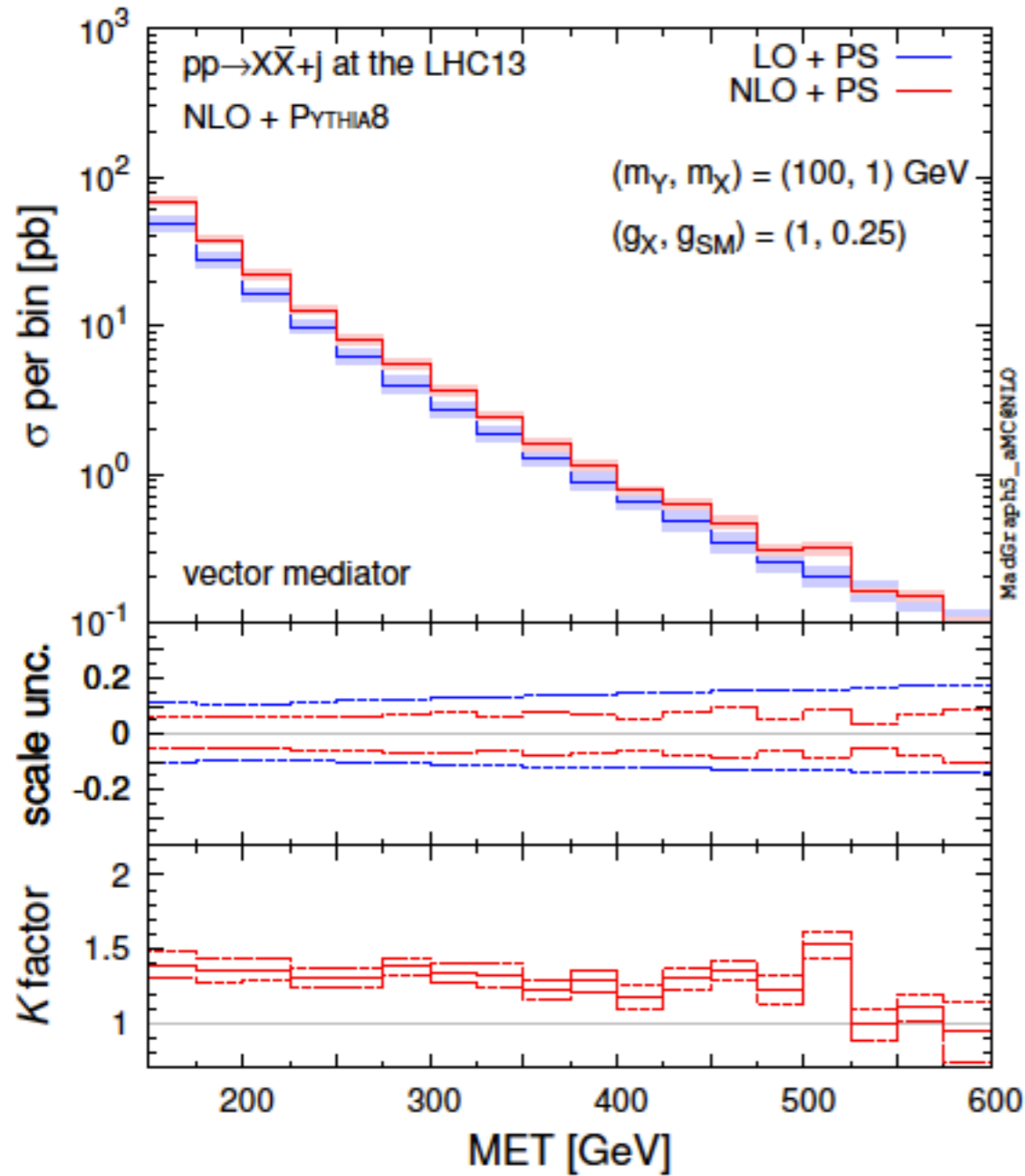
$10^1$ $^{+11.0}_{-9.6} \pm 1.5\%$	$0.988 \times 10^0$ $^{+14.7}_{-12.2} \pm 1.1\%$
$10^1$ $^{+6.0}_{-5.5} \pm 0.5\%$	$1.281 \times 10^0$ $^{+6.8}_{-6.8} \pm 0.6\%$
$10^{-3}$ $^{+17.4}_{-14.0} \pm 4.3\%$	$2.329 \times 10^{-3}$ $^{+18.9}_{-15.0} \pm 4.6\%$
$10^{-3}$ $^{+5.3}_{-6.4} \pm 2.2\%$	$2.411 \times 10^{-3}$ $^{+5.5}_{-6.8} \pm 2.3\%$
	1.04

PDF uncertainty  
scale (renormalization and factorization) uncertainty

## NLO corrections

- strongly depends on the mass spectrum and the kinematical cuts.
- sizeably reduces of the scale and PDF uncertainties.

# DM production at NLO+PS

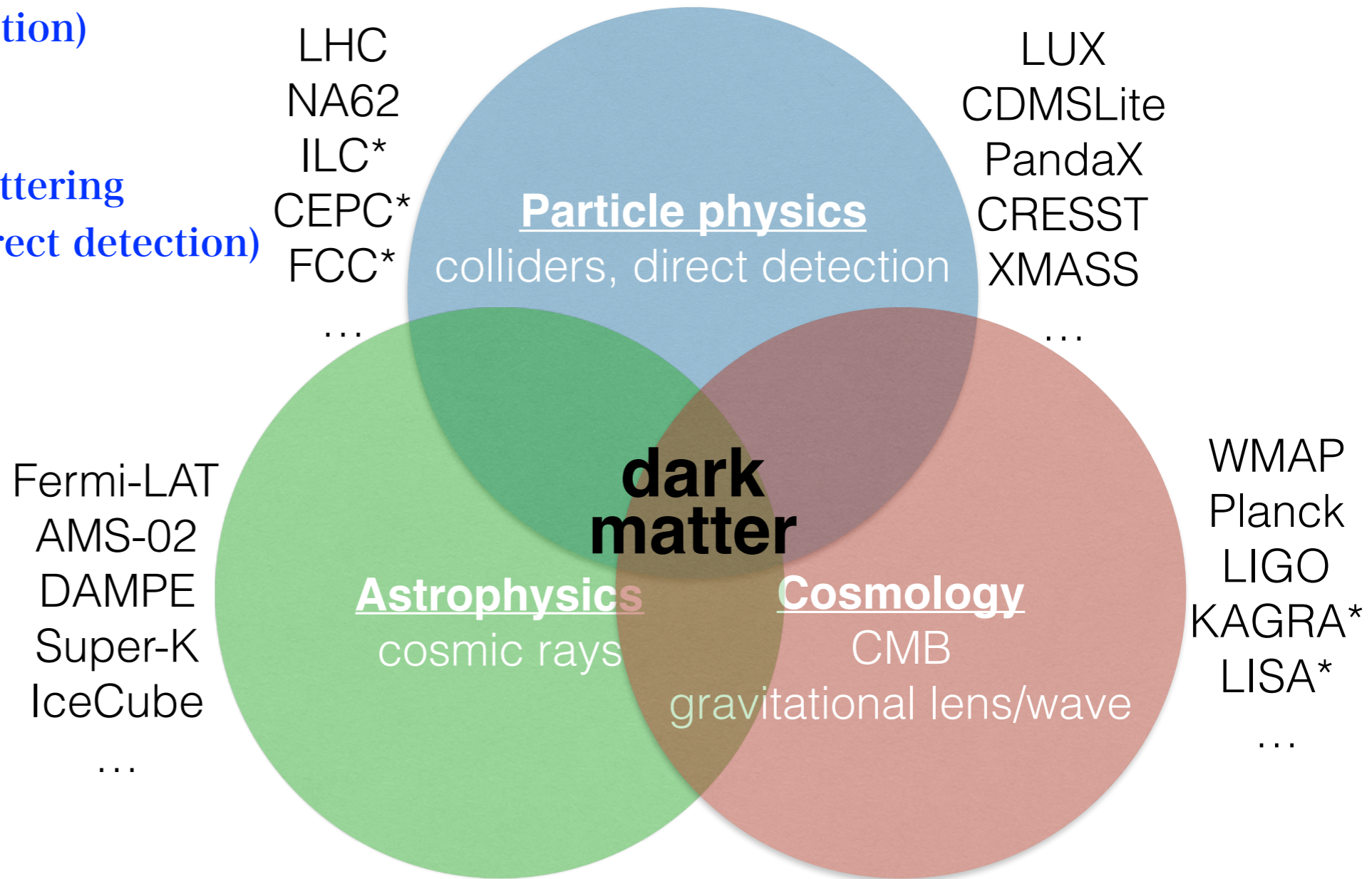
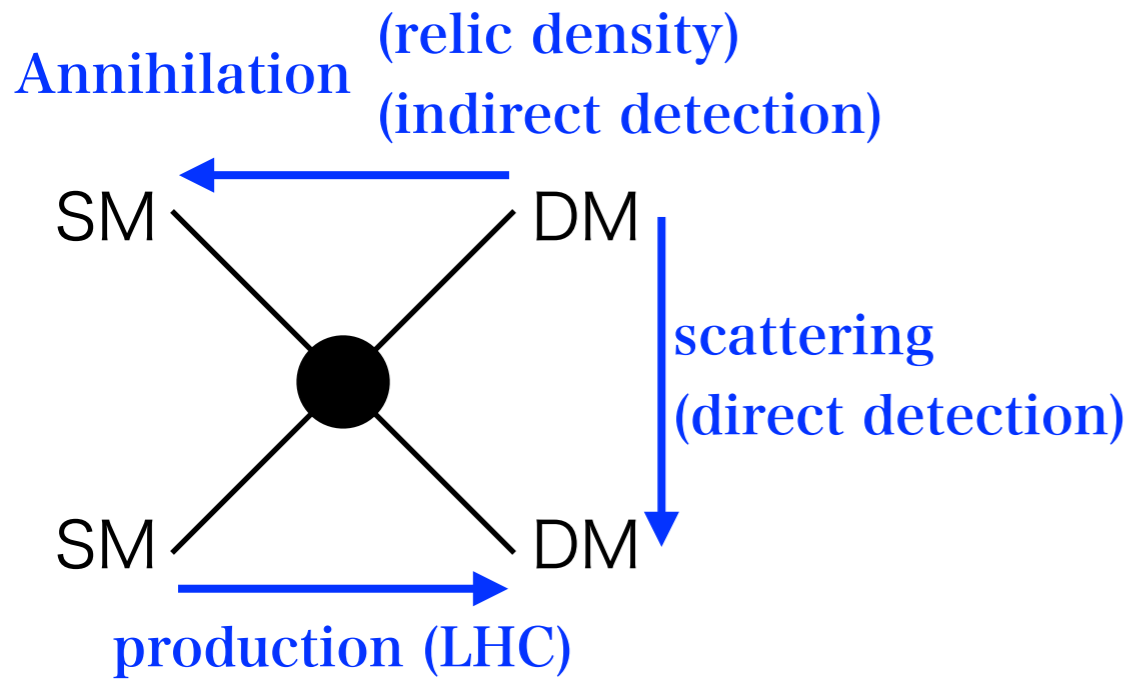




# Summary of the LHC DMWG activities

- develops well-defined benchmark models
- provides (accurate) simulation tools for signal and background
- defines guidelines for presentation and interpretation of the LHC results

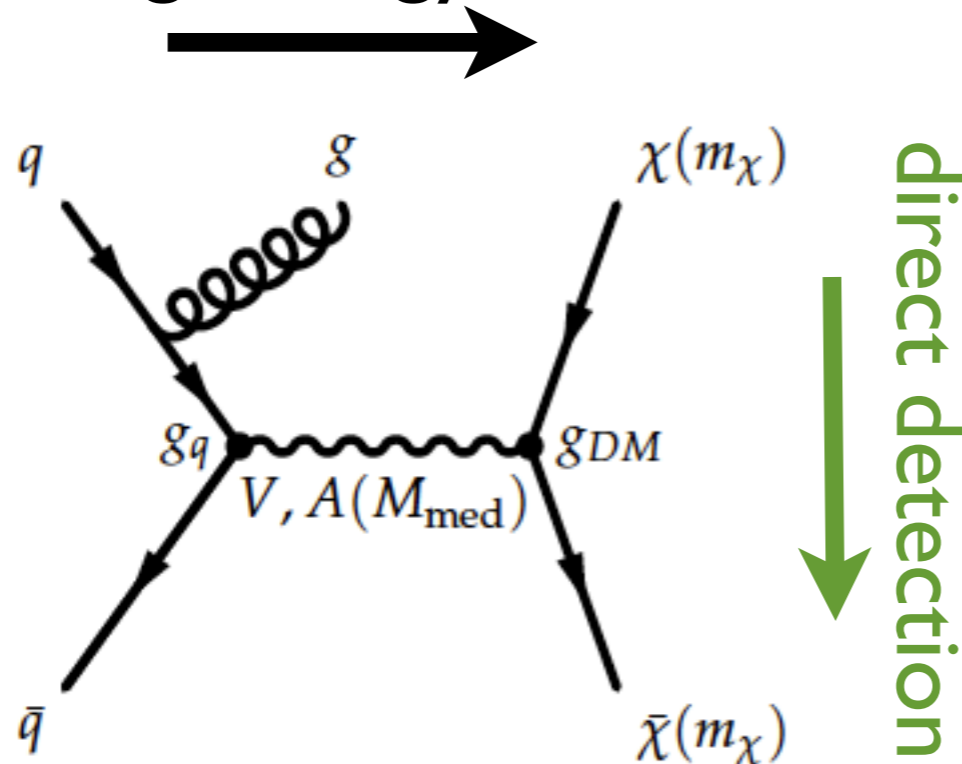
# Interpretation of the LHC MET search



# Signatures of simplified DM models

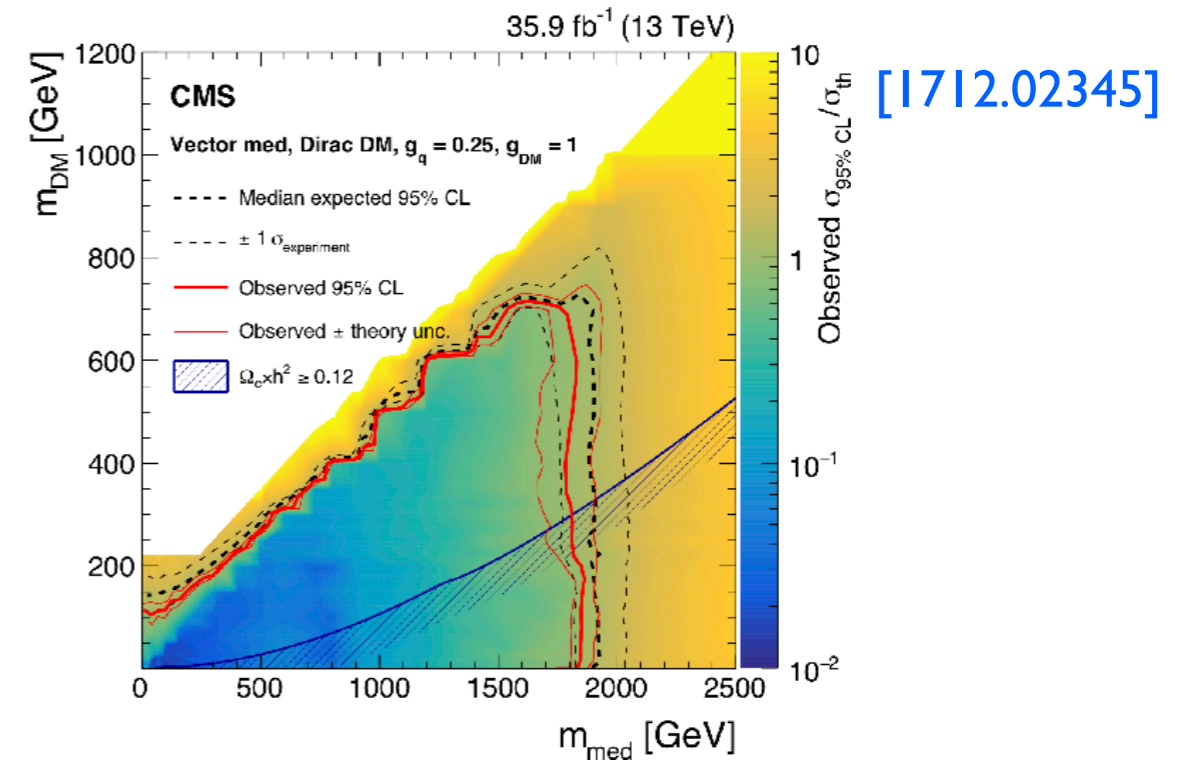
LHC DMWG [1603.04156, 1703.05703]

missing-energy search

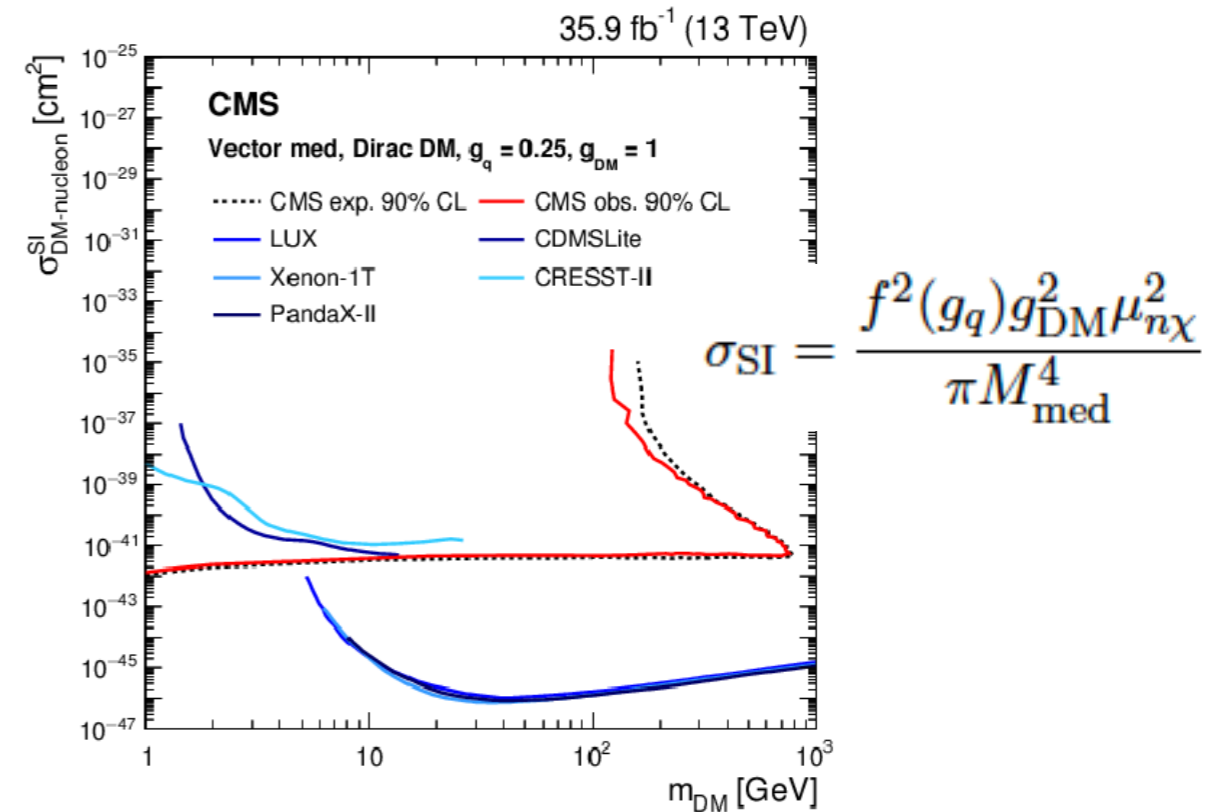


relic density  
indirect detection

direct detection



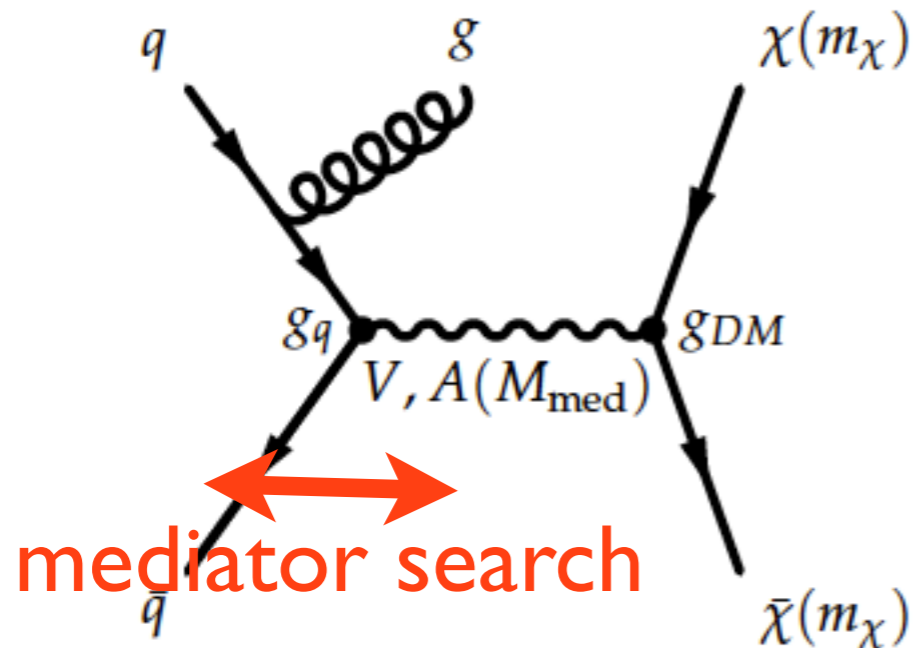
[1712.02345]



# Signatures of simplified DM models

LHC DMWG [1603.04156, 1703.05703]

missing-energy search



mediator search

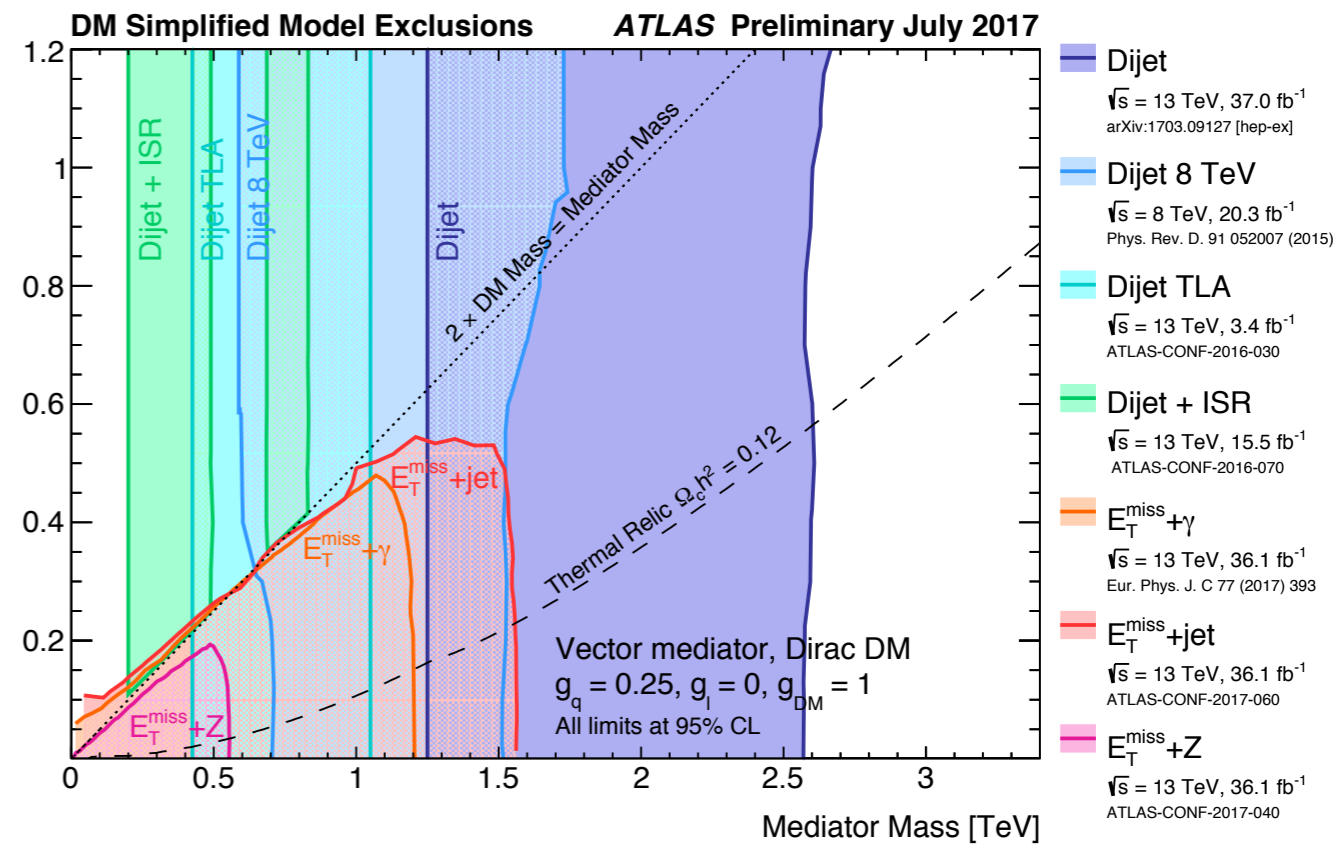


relic density

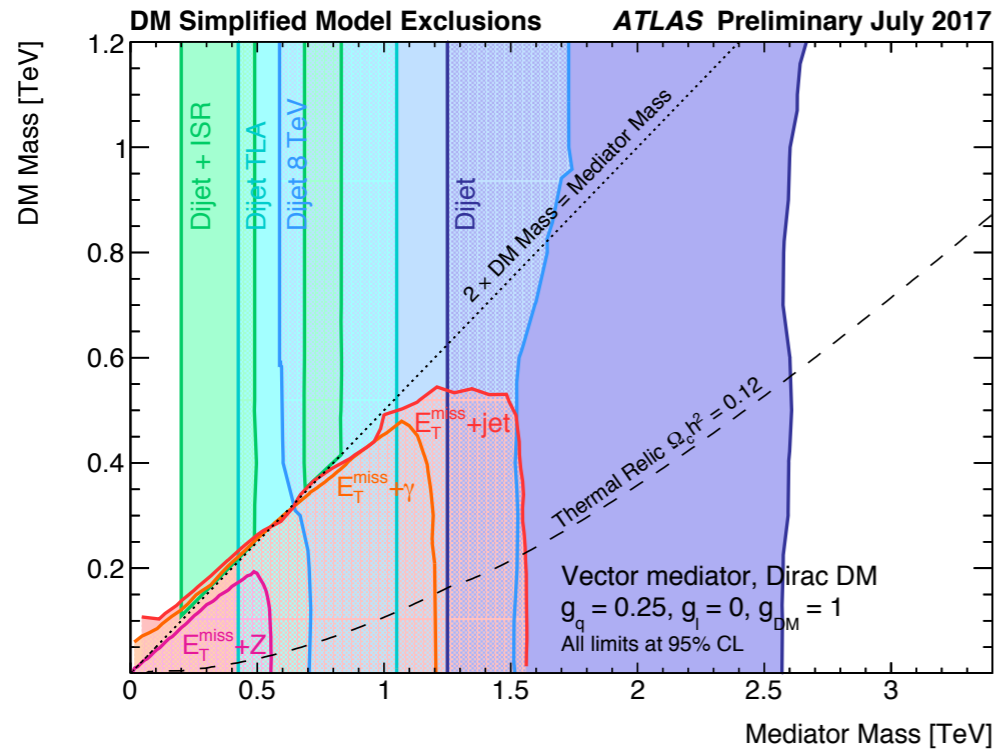
indirect detection



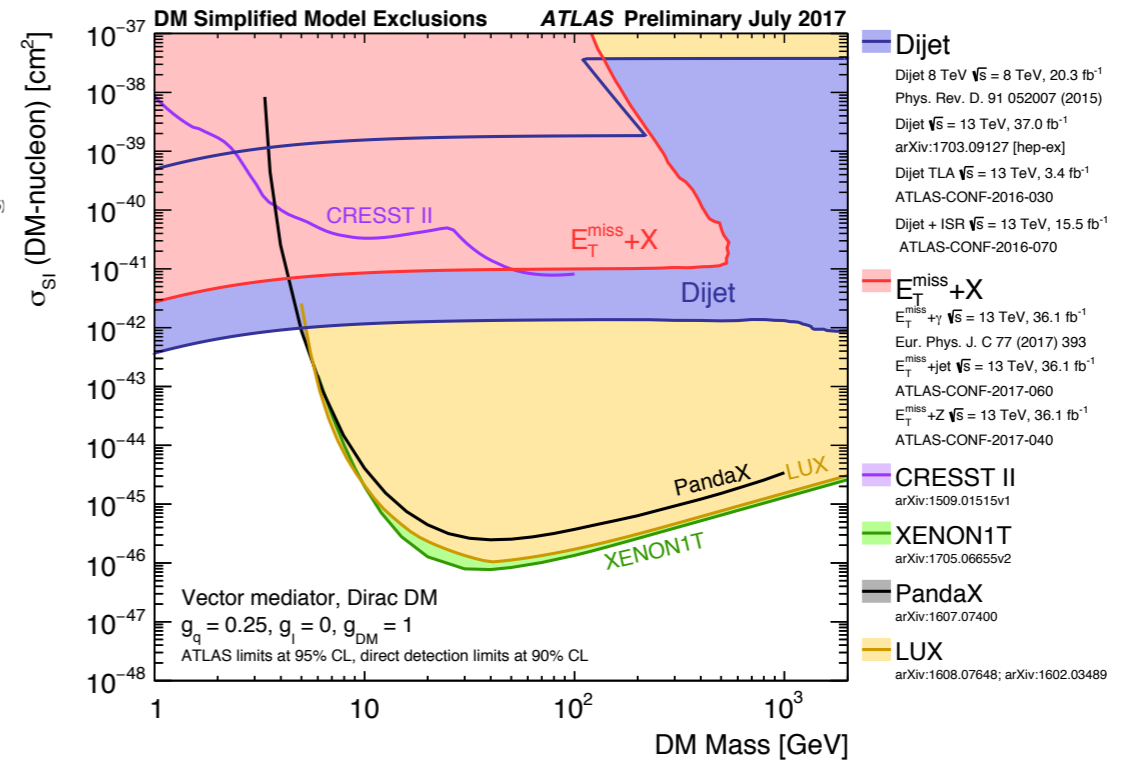
direct detection



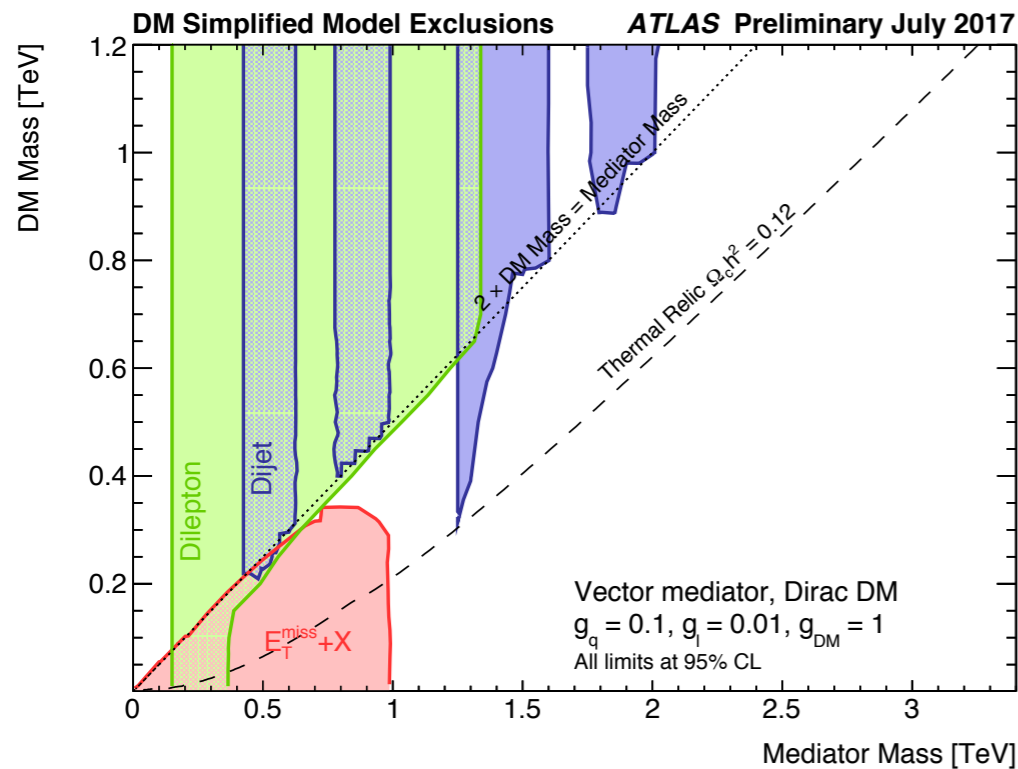
# Interpretation of the LHC searches



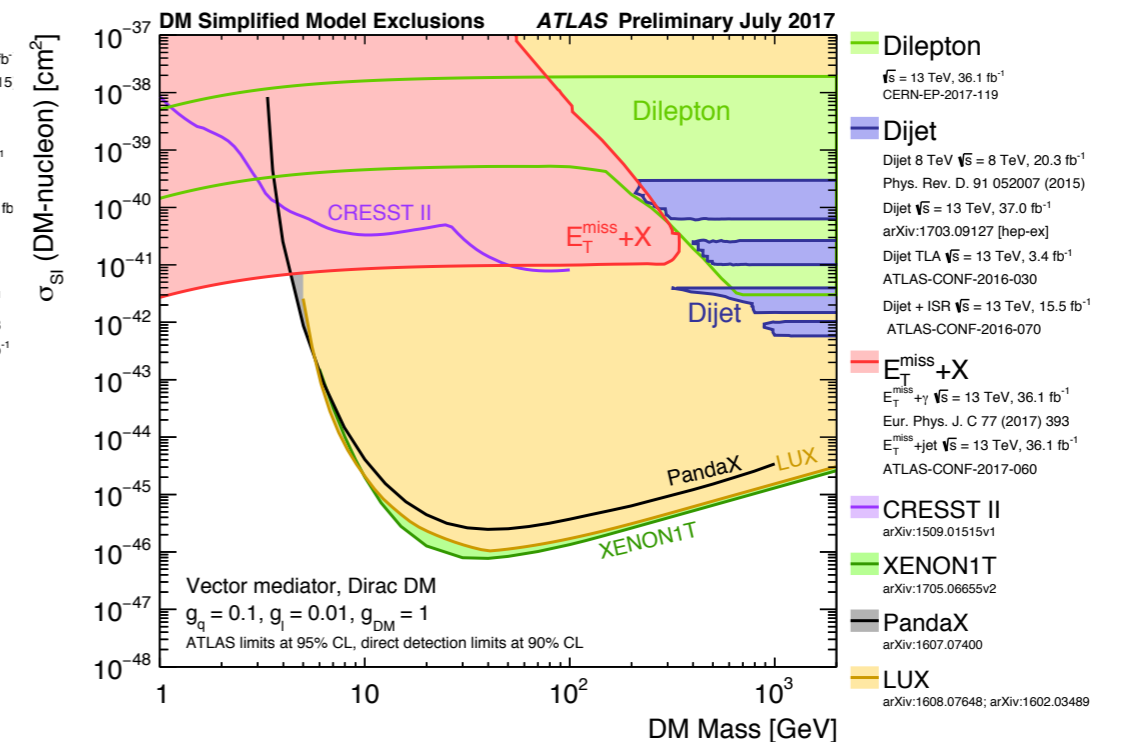
- Dijet**  
 $\sqrt{s} = 13 \text{ TeV}, 37.0 \text{ fb}^{-1}$   
 arXiv:1703.09127 [hep-ex]
- Dijet 8 TeV**  
 $\sqrt{s} = 8 \text{ TeV}, 20.3 \text{ fb}^{-1}$   
 Phys. Rev. D. 91 052007 (2015)
- Dijet TLA**  
 $\sqrt{s} = 13 \text{ TeV}, 3.4 \text{ fb}^{-1}$   
 ATLAS-CONF-2016-030
- Dijet + ISR**  
 $\sqrt{s} = 13 \text{ TeV}, 15.5 \text{ fb}^{-1}$   
 ATLAS-CONF-2016-070
- $E_T^{\text{miss}} + \gamma$**   
 $\sqrt{s} = 13 \text{ TeV}, 36.1 \text{ fb}^{-1}$   
 Eur. Phys. J. C 77 (2017) 393
- $E_T^{\text{miss}} + \text{jet}$**   
 $\sqrt{s} = 13 \text{ TeV}, 36.1 \text{ fb}^{-1}$   
 ATLAS-CONF-2017-060
- $E_T^{\text{miss}} + Z$**   
 $\sqrt{s} = 13 \text{ TeV}, 36.1 \text{ fb}^{-1}$   
 ATLAS-CONF-2017-040



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 Dijet 8 TeV  $\sqrt{s} = 8 \text{ TeV}, 20.3 \text{ fb}^{-1}$   
 Phys. Rev. D. 91 052007 (2015)  
 Dijet  $\sqrt{s} = 13 \text{ TeV}, 37.0 \text{ fb}^{-1}$   
 arXiv:1703.09127 [hep-ex]  
 Dijet TLA  $\sqrt{s} = 13 \text{ TeV}, 3.4 \text{ fb}^{-1}$   
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 ATLAS-CONF-2017-040
- CRESST II**  
 arXiv:1509.01515v1
- XENON1T**  
 arXiv:1705.06655v2
- PandaX**  
 arXiv:1607.07400
- LUX**  
 arXiv:1608.07648; arXiv:1602.03489

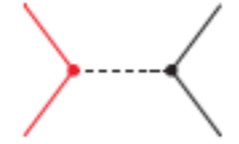
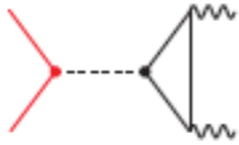
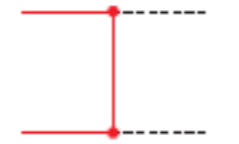
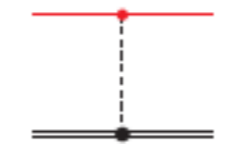
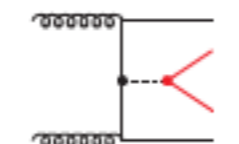
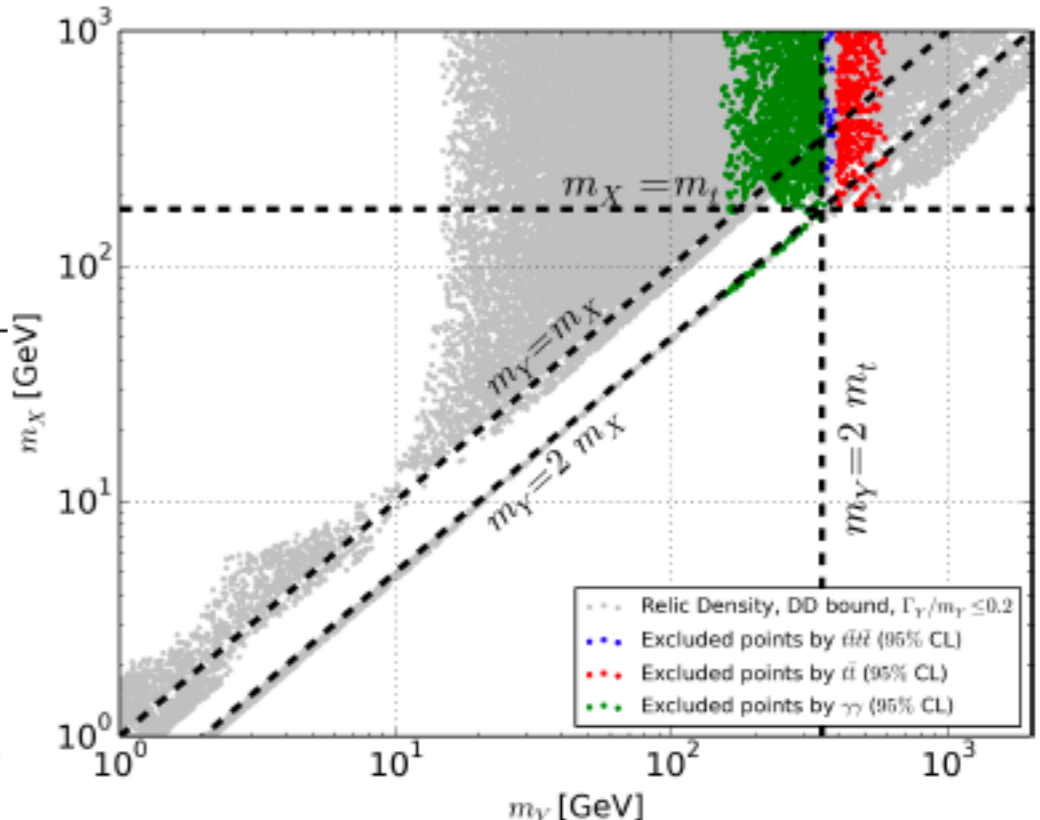
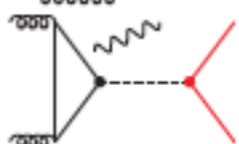
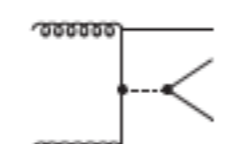
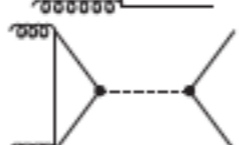
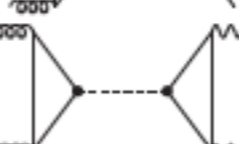


- Dijet**  
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 ATLAS-CONF-2017-060
- Dilepton**  
 $\sqrt{s} = 13 \text{ TeV}, 36.1 \text{ fb}^{-1}$   
 CERN-EP-2017-119



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 CERN-EP-2017-119
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 Dijet 8 TeV  $\sqrt{s} = 8 \text{ TeV}, 20.3 \text{ fb}^{-1}$   
 Phys. Rev. D. 91 052007 (2015)  
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- PandaX**  
 arXiv:1607.07400
- LUX**  
 arXiv:1608.07648; arXiv:1602.03489

# A comprehensive simplified DM study: s-channel spin-0

Cosmology	relic		$m_X > m_t$	Arina, Backovic, Conte, Fuks, Guo, Heisig, Hespel, Kraemer, Maltoni, Martini, KM, Pellen, Vryonidou [1605.09242, JHEP] Planck, FermiLAT
	indirect		$m_X < m_t$	
Astrophysics			$m_X > m_Y$	
	direct		$m_X > 1 \text{ GeV}$	LUX, CDMSLite
Colliders	$\cancel{E}_T$		$m_Y > 2m_X$	
	no $\cancel{E}_T$		$m_Y > 2m_X$	
			$m_Y > 2m_t$	
			$m_Y > 2m_t$	
			$m_Y < 2m_X,$	

$$\mathcal{L}_{t,X}^{Y_0} = -\left(g_t \frac{y_t}{\sqrt{2}} \bar{t}t + g_X \bar{X}X\right)Y_0$$

**X=DM**  
**Y=mediator**

# Spin-2 model: MET vs. resonance searches

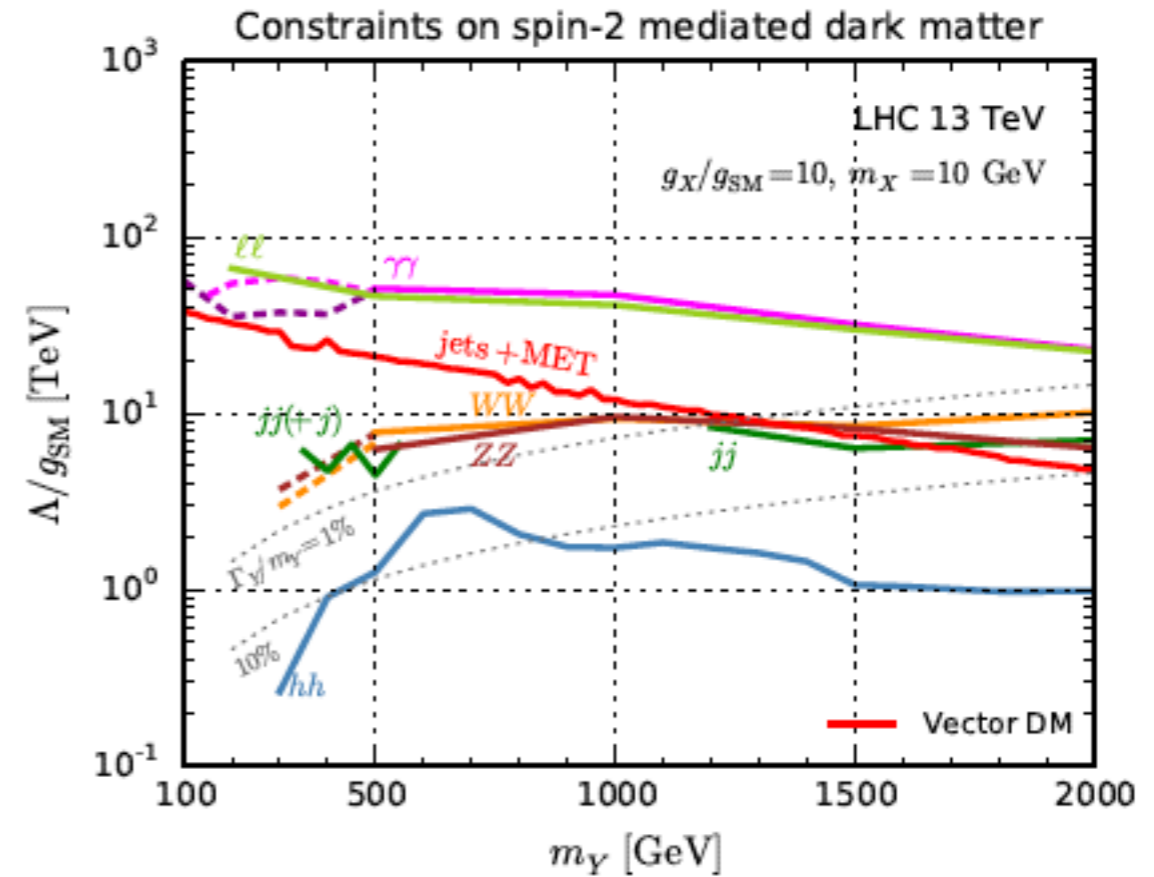
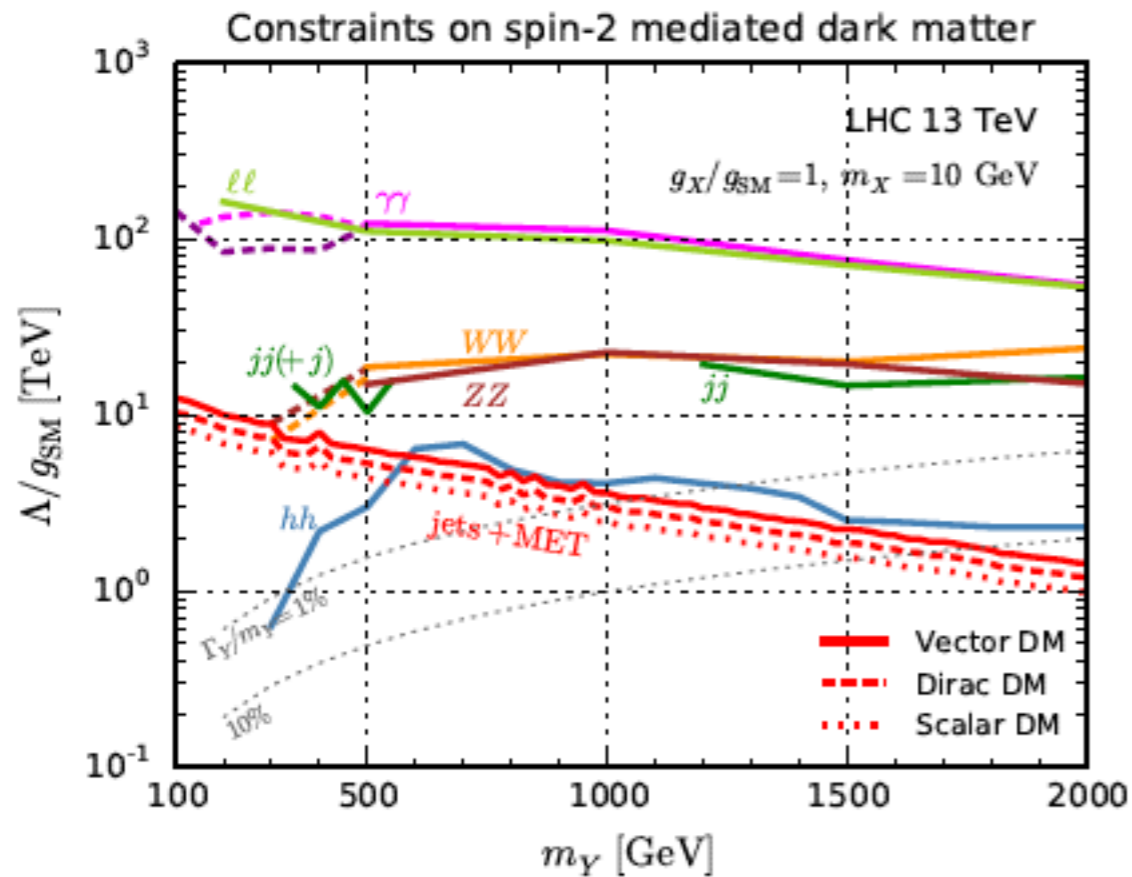
Kraml, Laa, KM, Yamashita [1701.07008, EPJC]

$$\mathcal{L}_X^{Y_2} = -\frac{1}{\Lambda} g_X^T T_{\mu\nu}^X Y_2^{\mu\nu}$$

$$\mathcal{L}_{\text{SM}}^{Y_2} = -\frac{1}{\Lambda} \sum_i g_i^T T_{\mu\nu}^i Y_2^{\mu\nu}$$

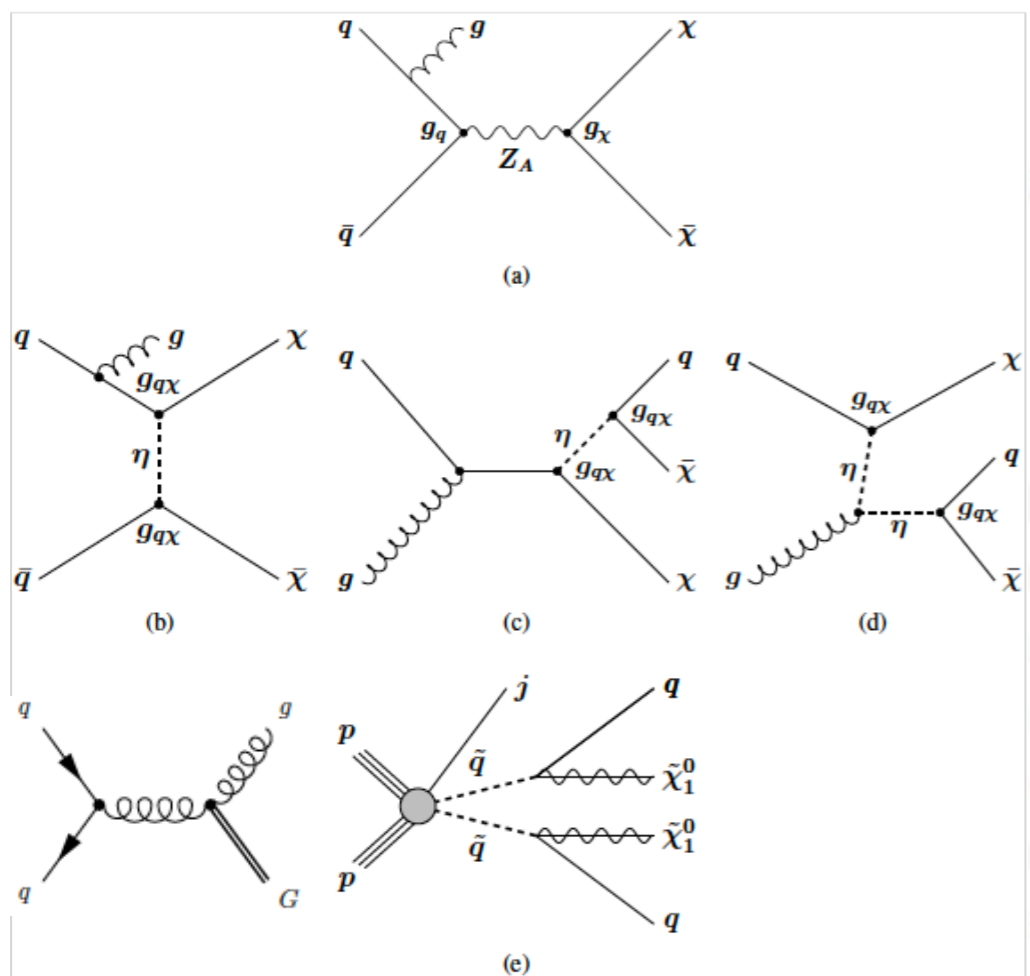
$$g_{\text{SM}} \equiv g_H^T = g_q^T = g_\ell^T = g_g^T = g_W^T = g_B^T$$

$$\{m_X, m_Y, g_X/\Lambda, g_{\text{SM}}/\Lambda\}$$



# Search for dark matter and other new phenomena in events with an energetic jet and large missing transverse momentum using the ATLAS detector

ATLAS Collaboration

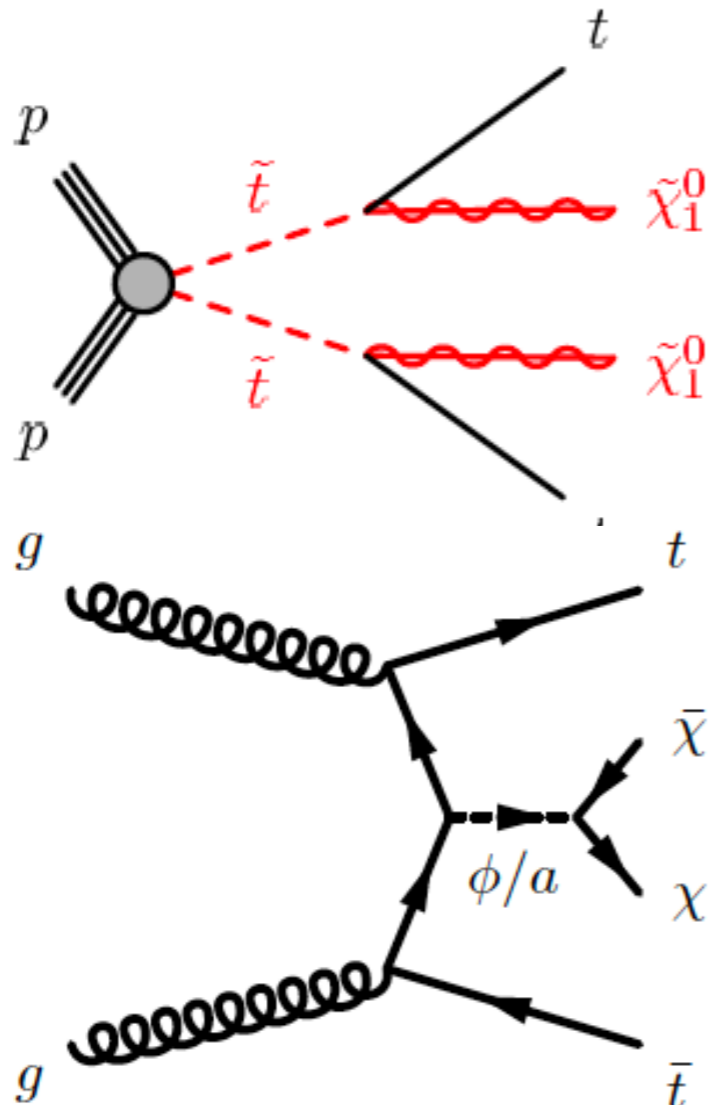


phenomena in final states with an energetic jet and large missing transverse momentum have been reported. The search uses proton–proton collision data with an integrated luminosity of  $36.1 \text{ fb}^{-1}$  at a centre-of-mass energy of  $13 \text{ TeV}$  using the ATLAS detector at the Large Hadron Collider. Events with a leading jet with a transverse momentum above  $250 \text{ GeV}$  and no other jets with a transverse momentum above  $250 \text{ GeV}$  are considered with increasing requirements on the missing transverse momentum. Good agreement is observed between the data and the Standard Model predictions. The results are translated into exclusion limits in models with pair-produced weakly interacting dark-matter candidates, large extra spatial dimensions, and supersymmetric particles in several compressed scenarios.

exclusion limits in models with pair-produced weakly interacting dark-matter candidates, large extra spatial dimensions, and supersymmetric particles in several compressed scenarios.

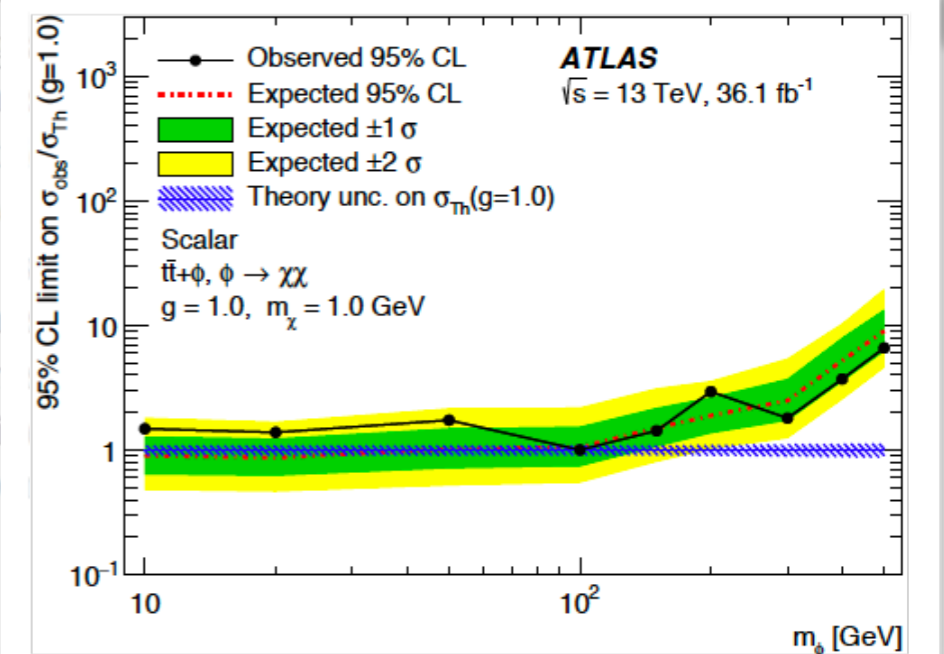
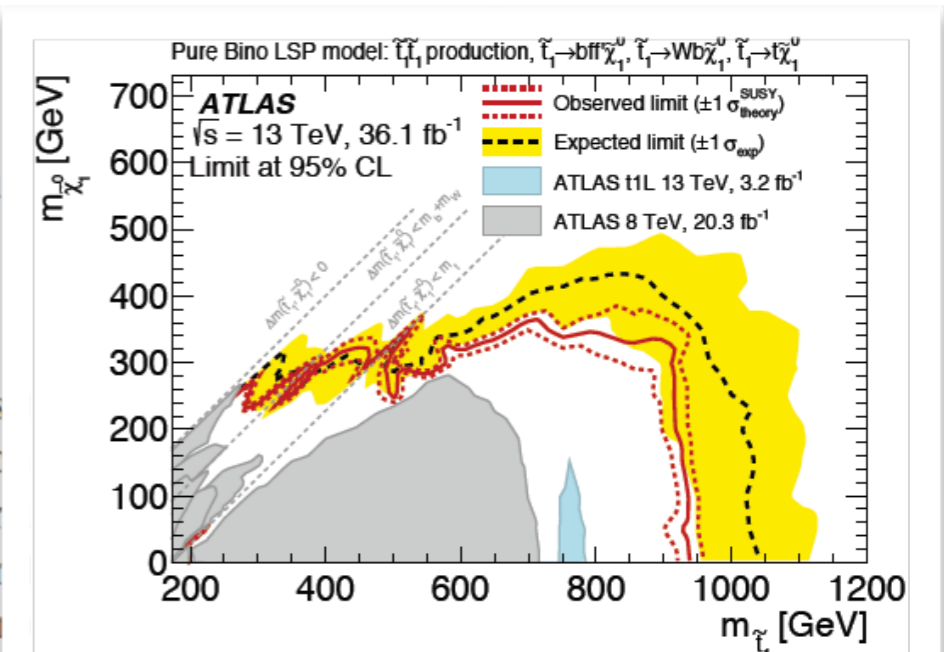


# Search for top-squark pair production in final states with one lepton, jets, and missing transverse momentum using $36 \text{ fb}^{-1}$ of $\sqrt{s} = 13 \text{ TeV } pp$ collision data with the ATLAS detector



The ATLAS Collaboration

Direct pair production of top squarks in final states with one isolated electron or muon and missing transverse momentum are reported. The analysis also takes into account top-squark pair production that decays into a pair of dark-matter particles and a top quark. The search uses data from proton-proton collisions in 2015 and 2016 at a centre-of-mass energy of  $\sqrt{s} = 13 \text{ TeV}$  recorded with the ATLAS detector, corresponding to an integrated luminosity of  $36.1 \text{ fb}^{-1}$ . Scenarios with different mass-splitting parameters are considered, including possible intermediate supersymmetric particles such as  $V$  bosons or the top quarks produce additional jets. Exclusion limits at 95% confidence level in several top-squark decay scenarios. For the spin-0 mediator models, upper limits are also derived for the top-squark production cross-section.



# Summary of the LHC DMWG activities

## [Discussions]

- develops well-defined benchmark models

Are the simplified models too simple/complicated?

- provides (accurate) simulation tools for signal and background

Are the current available tools satisfactory?

- defines guidelines for presentation and interpretation of the LHC results

Are the current experimental informations enough for reinterpretation/recasting?

# Summary and outlook

- Dark matter is one of the biggest mysteries in our world.
- The systematic simulation framework for dark matter searches have been developed not only for LHC but also for non-collider experiments.
- Higher-order predictions not only provide reliable rate but also reduce the theoretical uncertainty.
- A single model can be constrained by many different LHC searches and also by non-collider searches.
- Recasting business becomes more and more important.