Progress on Old and New Themes in cosmology

Hints for decaying dark matter from S_8 measurements

Guillermo Franco Abellán







Based on arXiv:2008.09615 with Riccardo Murgia, Vivian Poulin and Julien Lavalle

10/12/20

$$S_8 = \sigma_8 \sqrt{\frac{\Omega_m}{0.3}}$$

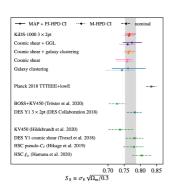
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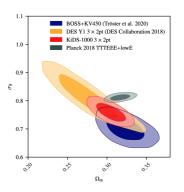
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$2-3\sigma$ tension between Weak Lensing and CMB (assuming Λ CDM)



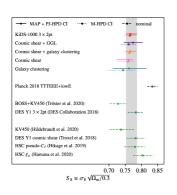


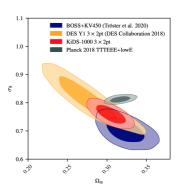
From 2008.11285

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From 2008.11285

BOSS+KIDS+2dfLenS analysis revealed tension is mainly **driven by** σ_8

Framework of the 2-body decay

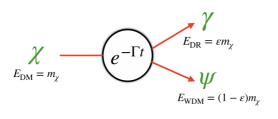
• Could the S_8 tension be related to **new physics**?

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- ullet Could it be related to the **stability of Dark Matter** (DM) on cosmological times scales?
- We explore DM decays to massless (Dark Radiation) and massive (Warm Dark Matter) particles, $\chi(\mathrm{DM}) \to \gamma(\mathrm{DR}) + \psi(\mathrm{WDM})$



Two extra parameters: Γ and ε

Evolution of perturbations and cosmological constraints in decaying dark matter models with arbitrary decay mass products

Shohei Aoyama, 1 Toyokazu Sekiguchi, 1,2 Kiyotomo Ichiki 3 and Naoshi Sugiyama 1,3,4

1402.2972

Full treatment of perts., no parameter scan

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Late universe decaying dark matter can relieve the H_0 tension

Kyriakos Vattis* and Savvas M. Koushiappas* Department of Physics, Brown University, 182 Hope St., Providence, RI 02212 and Brown Theoretical Physics Center, 182 Hope St., Providence, RI 02212

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CMB constraints on late-universe decaying dark matter as a solution to the \mathcal{H}_0 tension

Steven J. Clark, "Kyriakos Vattis," and Savvas M. Koushiappas*
Department of Physics, Brossen University, Providence, RI 02912-1843, USA and
Brown Theoretical Physics Center, Bross University, Providence, RI 02912-1843, USA

2006.03678

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Evolution of perturbations: full description

• Effects on P(k) and C_ℓ ? \to Track linear perturbations for the daughter particles: $\delta_{\rm D}$, $\theta_{\rm D}$ and $\sigma_{\rm D}$

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- Effects on P(k) and C_ℓ ? \to Track linear perturbations for the daughter particles: $\delta_{\rm D}$, $\theta_{\rm D}$ and $\sigma_{\rm D}$
- Boltzmann hierarchy of eqs. dictate the evolution of the **p.s.d.** multipoles $\Delta f_{\ell}(q, k, \tau)$
 - ♦ **DR treatment is easy**, momentum d.o.f. are integrated out
 - \diamond For WDM, one needs to follow the evolution of the full p.s.d. Computationally expensive $\to \mathcal{O}(10^8)$ ODEs to solve !

Evolution of perturbations: fluid approximation for the WDM

Based on a **fluid** description for massive neutrinos (1104.2935)

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The eqs. (valid at $k\tau\gg 1$) read

$$\dot{\delta}_{\mathrm{D}} = -3\mathcal{H}(c_{\mathrm{syn}}^2 - w)\delta_{\mathrm{D}} - (1 + w)\left(\theta_{\mathrm{D}} + \frac{\dot{h}}{2}\right) + \mathsf{a}\Gamma(1 - \varepsilon)\frac{\bar{\rho}_{\mathrm{DM}}}{\bar{\rho}_{\mathrm{D}}}(\delta_{\mathrm{DM}} - \delta_{\mathrm{D}})$$

$$\dot{ heta}_{
m D} = -\mathcal{H}(1-3c_a^2) heta_{
m D} + rac{c_{
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$$c_{a}^{2}(\tau) = w \left(5 - \frac{\mathfrak{p}_{\mathrm{D}}}{\overline{\rho}_{\mathrm{D}}} - \frac{\overline{\rho}_{\mathrm{DM}}}{\overline{\rho}_{\mathrm{D}}} \frac{a\Gamma}{3w\mathcal{H}} \frac{\varepsilon^{2}}{1 - \varepsilon} \right) \left[3(1 + w) - \frac{\overline{\rho}_{\mathrm{DM}}}{\overline{\rho}_{\mathrm{D}}} \frac{a\Gamma}{\mathcal{H}} (1 - \varepsilon) \right]^{-1}$$

and

$$c_{\rm syn}^2(k,\tau) = c_a^2(\tau) \left[1 + (1 - 2\varepsilon) T(k/k_{\rm fs}) \right]$$

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Accurate at the $\mathcal{O}(0.1\%)$ level in C_{ℓ} , and at $\mathcal{O}(1\%)$ level in P(k)

CPU time reduced from ~ 1 day to ~ 1 minute!

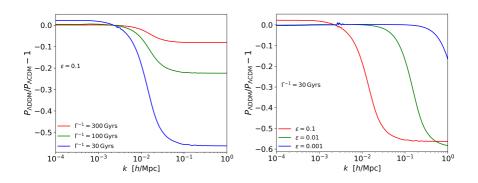
Physical effects of the WDM daughter

The WDM daughter leads to a **power suppression** in P(k) at small scales, $k > k_{\rm fs}$, where $k_{\rm fs} \sim \mathcal{H}/c_a$

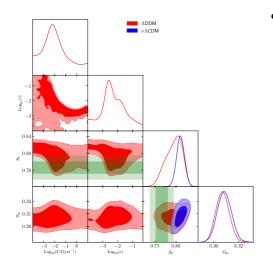
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- Γ controls the **depth** of the power suppression
- ullet controls the **cut-off** scale (k_{fs})



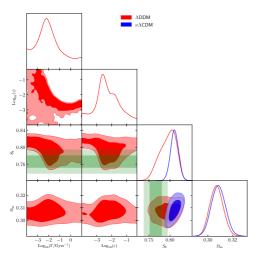
Resolving the S_8 tension with the 2-body decay



• Modified version of CLASS Run MCMC against **Planck**, **BAO SNIa**, $f\sigma_8 \& S_8^{-1}$

 $^{^{1}}S_{8} = 0.766^{+0.020}_{-0.014}$ from KIDS+BOSS+2dfLenS

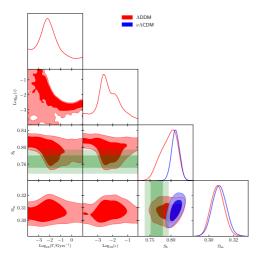
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	ν Λ CDM	ΛDDM
$\chi^2_{\rm CMB}$	1015.9	1015.2
$\chi^2_{S_8}$	5.64	0.002

$$\longrightarrow \Delta\chi^2_{\rm min} \simeq -5.5$$

 $arepsilon \simeq 0.7$ % and $\Gamma^{-1} \simeq 55$ Gyrs

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 \bullet In the best-fit cosmology, $\Omega_{\rm wdm}/\Omega_{\rm m}\sim 18\%$

 $^{^2} And$ a smaller $\omega_{\rm cdm},$ to keep $\Omega_{\rm m}$ fixed.

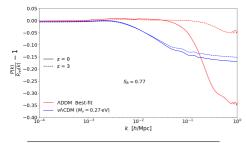
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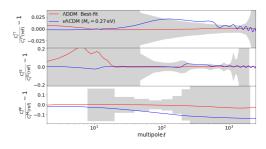
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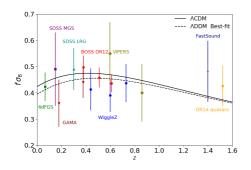
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Accurate measurements of $f\sigma_8$ at $0 \lesssim z \lesssim 1$ will further test the 2-body decay

Conclusions

- First thorough cosmological analysis of the 2-body decay scenario
- It fully restores cosmological concordance for S_8 (but not for H_0)
- Many interesting implications (DM model building, small-scale crisis, Xenon-1T excess)
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THANKS FOR YOUR ATTENTION

$$f(q, k, \mu, \tau) = \bar{f}(q, \tau) + \Delta f(q, k, \mu, \tau)$$

Expand Δf in multipoles. The Boltzmann eq. leads to the following **hierarchy** (in synchronous gauge comoving with the mother)

$$\begin{split} &\frac{\partial}{\partial \tau} \left(\Delta f_{0} \right) = -\frac{\mathbf{q}k}{\mathcal{E}} \Delta f_{1} + q \frac{\partial \bar{f}}{\partial q} \frac{\dot{h}}{6} + a \frac{\Gamma \bar{N}_{\mathrm{M}}(\tau)}{4\pi q^{3} \mathcal{H}} \delta(\tau - \tau_{q}) \delta_{\mathrm{M}}, \\ &\frac{\partial}{\partial \tau} \left(\Delta f_{1} \right) = \frac{\mathbf{q}k}{3\mathcal{E}} \left[\Delta f_{0} - 2\Delta f_{2} \right], \\ &\frac{\partial}{\partial \tau} \left(\Delta f_{2} \right) = \frac{\mathbf{q}k}{5\mathcal{E}} \left[2\Delta f_{1} - 3\Delta f_{3} \right] - q \frac{\partial \bar{f}}{\partial q} \frac{(\dot{h} + 6\dot{\eta})}{15}, \\ &\frac{\partial}{\partial \tau} \left(\Delta f_{l} \right) = \frac{\mathbf{q}k}{(2l+1)\mathcal{E}} \left[l\Delta f_{l-1} - (l+1)\Delta f_{l+1} \right] \qquad \text{(for } l \geq 3\text{)}. \end{split}$$

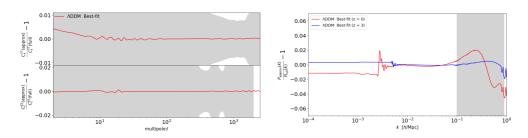
where $q=\mathit{a}(au_{q})p_{\mathrm{max}}.$ In the relat. limit $\mathbf{q}/\mathcal{E}=1$, so we can take

 $F_I \equiv {4\pi\over
ho_c} \int dq \ q^3 \Delta f_I$ and integrate out the dependency on q

Bonus II: Checking the accuracy of the fluid approximation

We compare two configurations (at the **best-fit** values)

- Full: Solve Boltzmann hierarchy with $N_q = 10^4$
- **Approx**: Solve Boltzmann hierarchy with $N_q=300$ and switch-on fluid eqs. at $k\tau>25$



The residual error on S_8 is $\sim 0.65\%$, smaller than the $\sim 1.8\%$ error of the measurement from BOSS+KIDS+2dfLenS