

Swampland and TCC Constraints on Inflation and Dark Energy

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Pont Avignon Conference, Dec. 9, 2020

Outline

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- 2 The Inflationary Scenario
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Map of the Cosmic Microwave Background (CMB)

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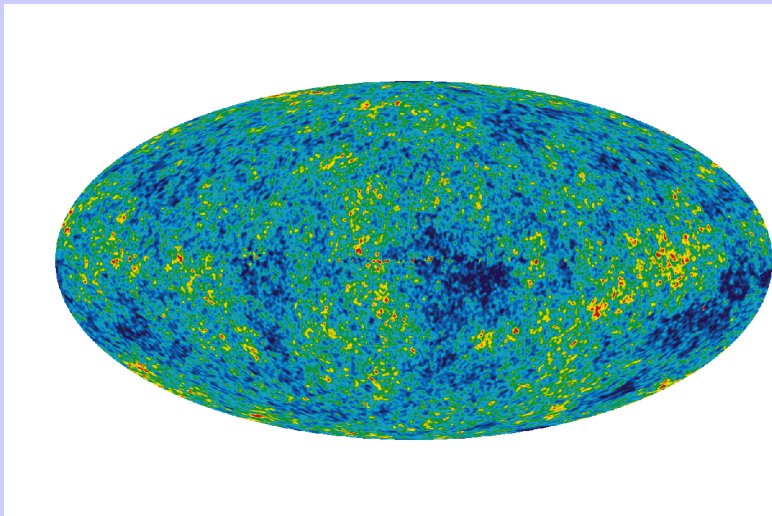
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Credit: NASA/WMAP Science Team

Angular Power Spectrum of CMB Anisotropies

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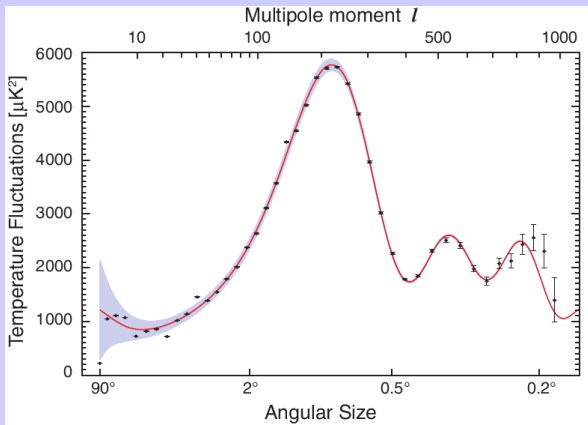
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Credit: NASA/WMAP Science Team

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1970Ap&SS...7....3S

SMALL-SCALE FLUCTUATIONS OF RELIC RADIATION

9

1970 p

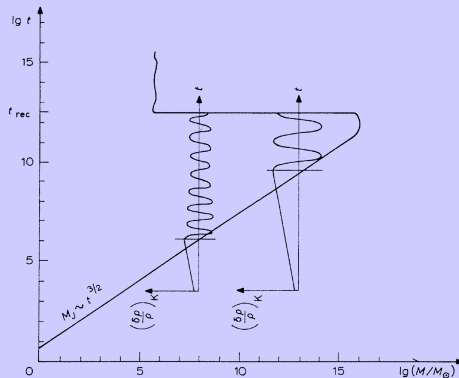


Fig. 1a. Diagram of gravitational instability in the 'big-bang' model. The region of instability is located to the right of the line $M_J(t)$; the region of stability to the left. The two additional lines of the graph demonstrate the temporal evolution of density perturbations of matter: growth until the moment when the considered mass is smaller than the Jeans mass and oscillations thereafter. It is apparent that at the moment of recombination perturbations corresponding to different masses correspond to different phases.

Predictions from 1970

R. Sunyaev and Y. Zel'dovich, *Astrophys. and Space Science* **7**, 3 (1970); P. Peebles and J. Yu, *Ap. J.* **162**, 815 (1970).

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- Given a **scale-invariant power spectrum of adiabatic fluctuations** on "super-horizon" scales before t_{eq} , i.e. standing waves.
- → "correct" power spectrum of galaxies.
- → **acoustic oscillations in CMB angular power spectrum.**
- → **baryon acoustic oscillations in matter power spectrum.**

Key Challenge

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How does one obtain such a spectrum?

- Inflationary Cosmology is the first scenario based on causal physics which yields such a spectrum.
- But it is not the only one.

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Hubble Radius vs. Horizon

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- **Horizon**: Forward light cone of a point on the initial Cauchy surface.
- Horizon: region of causal contact.
- **Hubble radius**: $l_H(t) = H^{-1}(t)$ inverse expansion rate.
- Hubble radius: local concept, relevant for dynamics of cosmological fluctuations.
- In Standard Big Bang Cosmology: Hubble radius = horizon.
- In any theory which can provide a mechanism for the origin of structure: Hubble radius \neq horizon.

Criteria for a Successful Early Universe Scenario

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Conclusions

- **Horizon \gg Hubble radius** in order for the scenario to solve the “horizon problem” of Standard Big Bang Cosmology.
- Scales of cosmological interest today **originate inside the Hubble radius at early times** in order for a causal generation mechanism of fluctuations to be possible.
- **Squeezing** of fluctuations on super-Hubble scales in order to obtain the acoustic oscillations in the CMB angular power spectrum.
- Mechanism for producing a **scale-invariant spectrum of curvature fluctuations** on super-Hubble scales.

Inflation as a Solution

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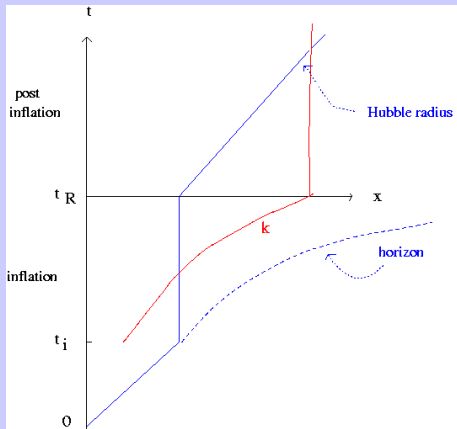
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Bounce as a Solution

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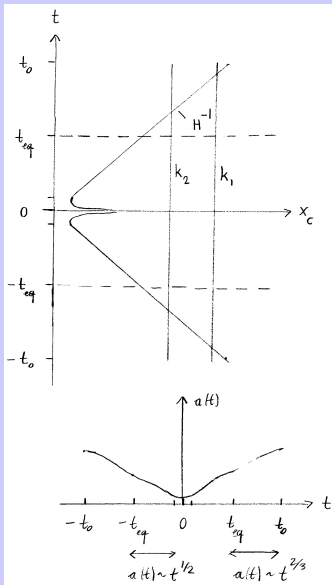
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Emergent Universe

R.B. and C. Vafa, *Nucl. Phys. B*316:391 (1989)

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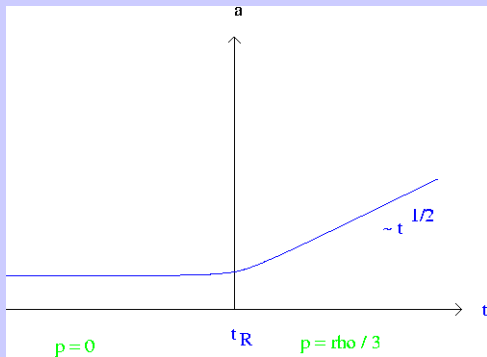
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Emergent Universe as a Solution

A. Nayeri, R.B. and C. Vafa, *Phys. Rev. Lett.* 97:021302 (2006)

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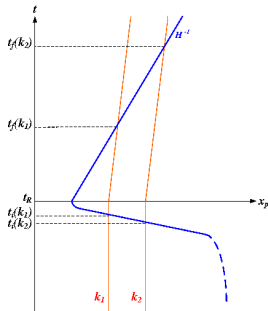
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Cosmological Inflation: The Current Paradigm of Early Universe Cosmology

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- **Cosmological Inflation** has become the **paradigm of early universe cosmology**.
- Idea: phase of accelerated (almost exponential) expansion of space in the early universe, for $t_i < t < t_R$.
- Solves some problems of Standard Big Bang Cosmology (explains the isotropy and spatial flatness of the universe).
- Provides a **causal mechanism for structure formation**.
- Inflationary cosmology has an **initial singularity** [Borde and Vilenkin].

Cosmological Inflation: The Current Paradigm of Early Universe Cosmology

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- Space-time described by the Einstein-Hilbert action.
- Matter dominated by a **slowly rolling scalar field φ** .
- in order to obtain equation of state $p \simeq -\rho$.
- **Quantum vacuum fluctuations** during the inflationary phase evolve into an **almost scale-invariant spectrum of cosmological perturbations** (V. Mukhanov and G. Chibisov, 1982).

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Lagrangian for a Scalar Field

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Lagrangian:

$$\mathcal{L} = \frac{1}{2} \partial_\mu \varphi \partial^\mu \varphi - V(\varphi)$$

Energy density and pressure of a scalar field:

$$\rho = \frac{1}{2} \dot{\varphi}^2 + \frac{1}{2} (\nabla \varphi)^2 + V(\varphi)$$

$$p = \frac{1}{2} \dot{\varphi}^2 - \frac{1}{6} (\nabla \varphi)^2 - V(\varphi)$$

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- Equation of motion

$$\ddot{\varphi} + 3H\dot{\varphi} = -V'(\varphi)$$

- Slow roll condition 1:

$$\frac{V'}{V} m_{pl} \ll 1.$$

- Slow roll condition 2:

$$\frac{V''}{V} m_{pl}^2 \ll 1.$$

- Large field range $\delta\varphi > m_{pl}$ required if inflation is to be a local attractor in initial condition space (R.B., arXiv:1601.01918 for a review).

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Question

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- **Self-consistent** as an **effective field theory** (EFT).
- Does inflation emerge from a UV complete theory?

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Obstacles to Obtaining Inflation

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Conclusions

- **Swampland conjectures:** Inflation hard to realize in superstring theory.
- **Trans-Planckian Censorship Conjecture:** causality and unitarity problem for inflation.

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Conclusions

- There is a vast **landscape** of **effective field theories**.
- Any space-time dimension, and number of fields, any shape of the potential, any field range.
- **Superstring theory** is very **constraining**.
- Only a **small subset** of all EFTs is consistent with string theory.
- The rest lie in the **swampland**.

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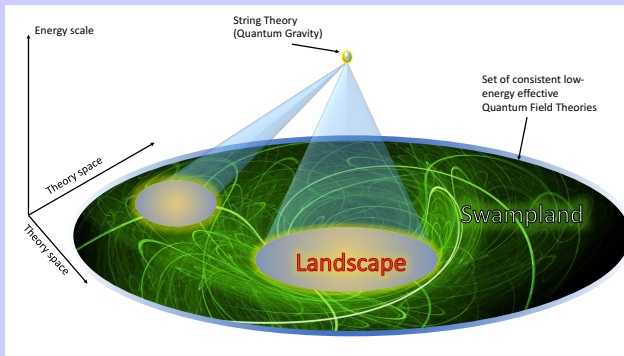
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Swampland Conjectures

H. Ooguri and C. Vafa, hep-th/0605264; G. Obied, H. Ooguri, L. Spodyneiko and C. Vafa, arXiv:1806.08362; S. Garg and C. Krishnan, arXiv:1807.05193; H. Ooguri, E. Palti, G. Shiu and C. Vafa, arXiv:1810.05506.

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Conclusions

What are conditions for habitable islands sticking out from the swamp?

- The effective field theory is only valid for $\Delta\varphi < dm_{pl}$ (field range condition).
- The potential of φ obeys (de Sitter conjecture)

$$\left| \frac{V'}{V} \right| m_{pl} \geq c_1 \quad \text{or} \quad \frac{V''}{V} m_{pl}^2 \leq -c_2$$

Note: d, c_1, c_2 constants of order 1.

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Conclusions

- φ is a string theory modulus field.
- **Field range condition:** move $\delta\varphi \sim m_{pl} \rightarrow$ tower of string states becomes massless and must be included in low energy EFT.
- **De Sitter conjecture:** moduli stabilization \rightarrow steep potential.
- Example: S. Laliberte and R.B., arXiv:1911.00199:

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Inflation in the Swamp?

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- **Slow roll inflation in the swampland.**
- False vacuum inflation in the swampland.
- Note: Warm inflation can escape from the swamp.

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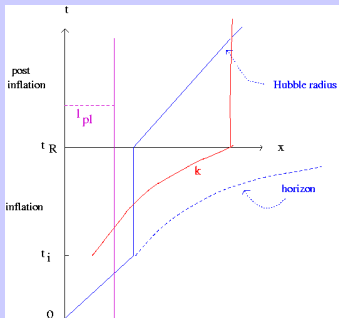
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Trans-Planckian Problem

J. Martin and R.B., *Phys. Rev. D* 63, 123501 (2002)



- **Success of inflation:** At early times scales are inside the Hubble radius \rightarrow causal generation mechanism is possible.
- **Problem:** If time period of inflation is more than $70H^{-1}$, then $\lambda_p(t) < l_{pl}$ at the beginning of inflation.
- \rightarrow new physics **MUST** enter into the calculation of the fluctuations.

Trans-Planckian Censorship Conjecture (TCC)

A. Bedroya and C. Vafa., arXiv:1909.11063

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Conclusions

No trans-Planckian modes exit the Hubble horizon.

$$ds^2 = dt^2 - a(t)^2 d\mathbf{x}^2$$

$$H(t) \equiv \frac{\dot{a}}{a}(t)$$

$$\frac{a(t_R)}{a(t_i)}|_{pl} < H(t_R)^{-1}$$

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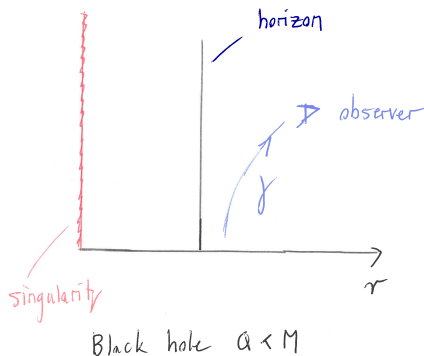
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Justification

R.B. arXiv:1911.06056

Analogy with Penrose's Cosmic Censorship Hypothesis:



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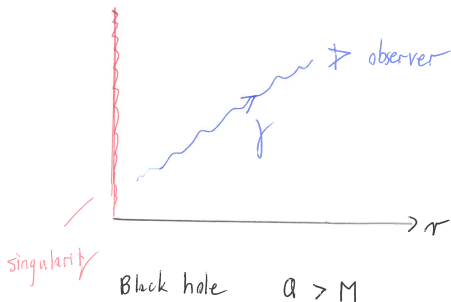
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- General Relativity allows for solutions with **timelike singularities**: super-extremal black holes.
- → Cauchy problem not well defined for observer external to black holes.
- Evolution **non-unitary** for external observer.
- Conjecture: ultraviolet physics → **external observer** shielded from the **singularity** and **non-unitarity** by **horizon**.

Justification

R.B. arXiv:1911.06056

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Cosmological Version of the Censorship Conjecture

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Translation

- Singularity \rightarrow trans-Planckian modes
- Black Hole horizon \rightarrow Hubble horizon.
- Position space \rightarrow momentum space.

Observer outside of Hubble horizon must be shielded from the trans-Planckian modes.

Cosmological Version of the Censorship Conjecture

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Why Hubble Horizon?

R.B. arXiv:1911.06056; A. Bedroya and C. Vafa., arXiv:1909.11063

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- Recall: Fluctuations only oscillate on sub-Hubble scales.
- Recall: Fluctuations freeze out, become **squeezed states** and **classicalize** on super-Hubble scales.
- **Demand:** classical region be insensitive to trans-Planckian region.
- → no trans-Planckian modes ever exit Hubble horizon.

Why Hubble Horizon?

R.B. arXiv:1911.06056; A. Bedroya and C. Vafa., arXiv:1909.11063

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Justification

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- Recall: **non-unitarity** of quantum field theory in an expanding universe (N. Weiss, Phys. Rev. D32, 3228 (1985)).
- \mathcal{H} is the product Hilbert space of a harmonic oscillator Hilbert space for all comoving wave numbers k
- Fixed k_{min} , time dependent k_{max} : $k_{max}(t)a(t)^{-1} = m_{pl}$
- **Demand: classical region be insensitive to non-unitarity.**
- \rightarrow no trans-Planckian modes ever exit Hubble horizon.

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Application to Inflation

A. Bedroya, R.B., M. Loverde and C. Vafa., arXiv:1909.11106

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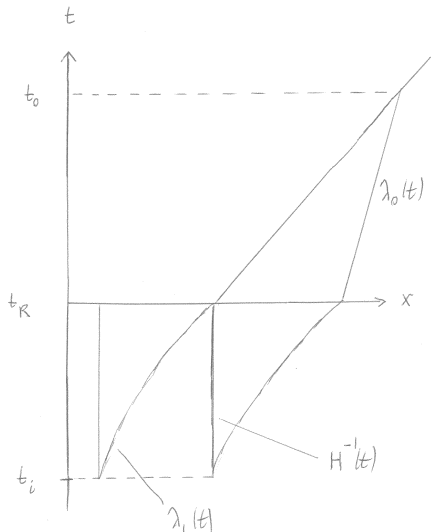
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TCC implies:

$$\frac{a(t_R)}{a(t_*)} l_{pl} < H(t_R)^{-1}$$

Demanding that inflation yields a causal mechanism for generating CMB anisotropies implies:

$$H_0^{-1} \frac{a(t_0)}{a(t_R)} \frac{a(t_R)}{a(t_*)} < H^{-1}(t_*)$$

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Upper bound on the **energy scale of inflation**:

$$V^{1/4} < 3 \times 10^9 \text{GeV}$$

→ **upper bound** on the **primordial tensor to scalar ratio** r :

$$r < 10^{-30}$$

Note: Secondary tensors will be larger than the primary ones.

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For **slow roll inflation**:

$$\epsilon < 10^{-31}$$

Up to logarithmic factors:

$$\Delta\varphi \sim \epsilon^{1/2} m_{pl}$$

→ extreme fine tuning of the initial velocity is required

$$\frac{\dot{\varphi}_{SR}}{\dot{\varphi}_N} \sim \epsilon^{1/2}$$

Bouncing Cosmologies Consistent with TCC

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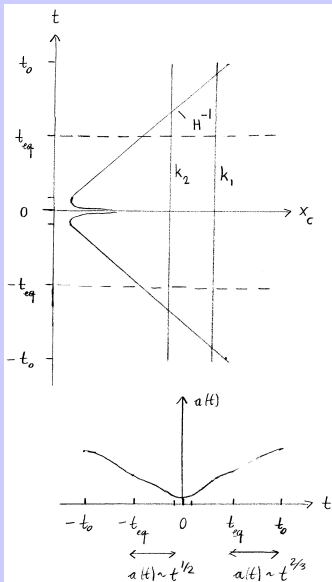
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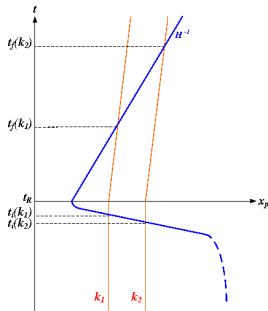
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- Inflation does not **naturally** emerge from superstring theory.
- Inflation is **highly fine tuned** assuming validity of the TCC.
- Alternatives to inflation are consistent with the TCC.
- **String Gas Cosmology** is a specific realization of the emergent scenario which directly comes from fundamental principles of string theory.
- The **Ekpyrotic bouncing scenario** can be easily embedded in string theory.

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- 3 Swampland Criteria
- 4 Trans-Planckian Censorship Conjecture
- 5 Implications for Dark Energy**
- 6 Alternatives
 - Ekpyrotic S-Brane Bounce
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Constraints from Swampland Considerations

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Conclusions

Dark Energy cannot be a Cosmological Constant.

Note: This also follows from the TCC!

Quintessence is **constrained by observations.**

Consider a quintessence model with potential

$$V(\phi) = V_0 e^{-\lambda\phi/m_{pl}}$$

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Current Constraints

L. Heisenberg, M. Bartelmann, RB and A. Refregier, arXiv:1808.02877

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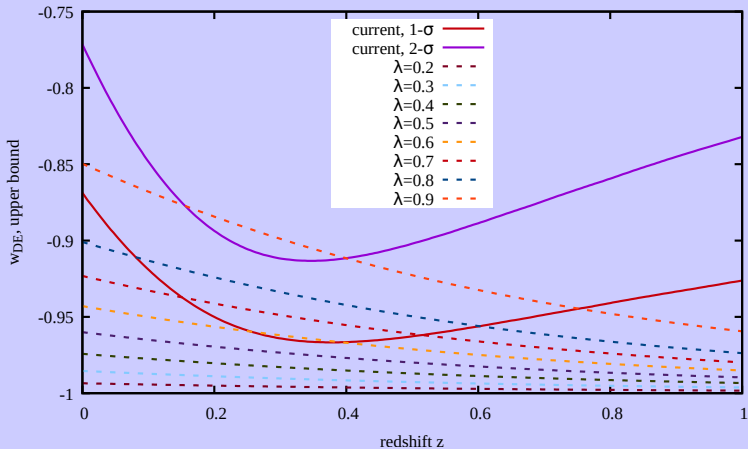
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Projected Euclid Constraints

L. Heisenberg, M. Bartelmann, RB and A. Refregier, arXiv:1808.02877

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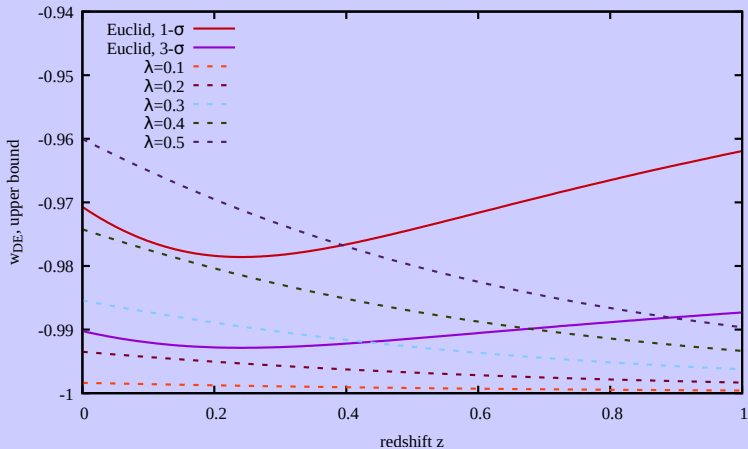
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Ekpyrotic Bounce

J. Khoury, B. Ovrut, P. Steinhardt and N. Turok, Phys. Rev. D, 2001

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Conclusions

- Space-time described by Einstein-Hilbert action.
- Idea: **Slow contraction** given by matter with equation of state $w \gg 1$.
- Obtained by assuming that matter is dominated by a **scalar field φ with a negative exponential potential**.
- Anisotropies diluted, creates spatial flatness
- Local attractor in initial condition space.
- Note: **Negative exponential potentials are ubiquitous in string theory.**

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Obtaining an Ekpyrotic Universe

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Conclusions

$$V(\varphi) = -V_0 \exp(-\sqrt{2/p} \varphi / m_{pl}) \quad p \ll 1$$

$$a(t) \sim (-t)^p$$

$$w \simeq \frac{4}{3p}$$

$$\varphi(t) = \sqrt{2p} m_{pl} \log\left(-\sqrt{\frac{V_0}{m_{pl}^2 p(1-3p)}} t\right)$$

Ekpyrosis: Small Field and Large Slope

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Conclusions

Consider $\delta t = H^{-1}$:

$$\varphi \sim p^{1/2} \log(p^{-1/2}) m_{pl} \ll m_{pl}.$$

Relative slope of the potential:

$$\left| \frac{V'}{V} \right| m_{pl} \sim p^{-1/2} \gg 1.$$

Relative curvature of the potential:

$$\frac{V''}{V} m_{pl}^2 = \frac{2}{p} \gg 1.$$

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Safe Ekpyrosis

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Conclusions

- Ekpyrotic scenario consistent with the swampland criteria.
- Ekpyrotic scenario consistent with the TCC.

S-Brane and Ekpyrosis

RB and Z. Wang, arXiv:2001.00638, arXiv:2004.06437

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Challenges for Ekpyrotic Cosmology:

- How do we get the bounce?
- How do we obtain a scale-invariant spectrum of curvature fluctuations?
- Can we obtain a spectrum of gravitational waves relevant to current observations?

Adding an S-Brane to the EFT action can solve all three problems, and leads to two consistency relations for cosmological observables.

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RB and Z. Wang, arXiv:2001.00638, arXiv:2004.06437

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S-Brane and Ekpyrosis II

RB and Z. Wang, arXiv:2004.06437; RB, K. Dasgupta and ZW, arXiv:2007.01203

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Consistency relations:

$$\begin{aligned}n_t &= 1 - n_s \\ r &\sim 36(k_{CTB})^{4q}(1 - n_s)^2\end{aligned}$$

Reheating after Ekpyrosis

Coupling of the S-brane to Standard Model gauge fields \rightarrow gauge field production during S-brane decay.

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Idea: At the string scale a new tower of string states becomes comparable in mass to the usual low energy degrees of freedom;

→ they must be included in the low energy effective action.

Included as an **S-Brane**.

$$S = \int d^4x \sqrt{-g} \left[R + \frac{1}{2} \partial_\mu \varphi \partial^\mu \varphi - V(\varphi) \right] \\ - \int d^4x \kappa \delta(\tau - \tau_B) \sqrt{\gamma},$$

$$\kappa \equiv N \eta_S^3,$$

Note: The S-brane has $\rho = 0$ and $p < 0 \rightarrow$ can mediate the transition between contraction and expansion.

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String Gas Cosmology: A Realization of the Emergent Scenario

R.B. and C. Vafa, *Nucl. Phys. B*316:391 (1989)

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Based on new degrees of freedom and new symmetries of string theory.

New Degrees of Freedom

Assumption: All spatial dimensions toroidal, radius R .

String states:

- momentum modes: $E_n = n/R$
- winding modes: $E_m = mR$
- oscillatory modes: E independent of R

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New Symmetries: T-Duality

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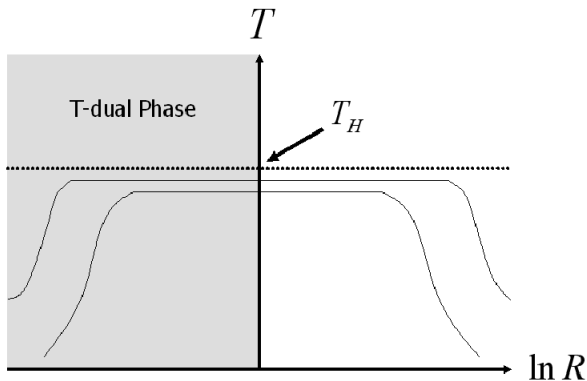
T-Duality

- Momentum modes: $E_n = n/R$
- Winding modes: $E_m = mR$
- Duality: $R \rightarrow 1/R$ $(n, m) \rightarrow (m, n)$
- Mass spectrum of string states unchanged
- Symmetry of vertex operators
- Symmetry at non-perturbative level \rightarrow existence of D-branes

Absence of a Temperature Singularity in String Cosmology

R.B. and C. Vafa, *Nucl. Phys. B*316:391 (1989)

Temperature-size relation in string gas cosmology



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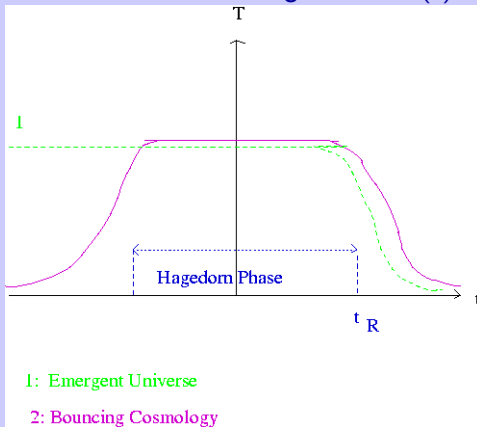
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Assume some action gives us $R(t)$



Position Operators

R.B. and C. Vafa, *Nucl. Phys. B*316:391 (1989)

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Conclusions

Position operators (dual to momenta)

$$|x\rangle = \sum_p \exp(ix \cdot p) |p\rangle$$

Dual position operators (dual to windings)

$$|\tilde{x}\rangle = \sum_w \exp(i\tilde{x} \cdot w) |w\rangle$$

Position Operators

R.B. and C. Vafa, *Nucl. Phys. B*316:391 (1989)

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Heavy vs. Light Modes

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Conclusions

- $R \gg 1$: momentum modes light.
- $R \ll 1$: winding modes light.
- $R \gg 1$: length measured in terms of $|x >$.
- $R \ll 1$: length measured in terms of $|\tilde{x} >$
- $R \sim 1$: both $|x >$ and $|\tilde{x} >$ important.

Conclusion: At string scale densities usual effective field theory (EFT) based on supergravity will break down.

Conclusion: If an effective field theory description is valid, it must be an EFT in 18 spatial dimensions.

Heavy vs. Light Modes

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Physical length operator

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$$l_p(R) = R \quad R \gg 1$$

$$l_p(R) = \frac{1}{R} \quad R \ll 1$$

Structure formation in string gas cosmology

A. Nayeri, R.B. and C. Vafa, *Phys. Rev. Lett.* 97:021302 (2006)

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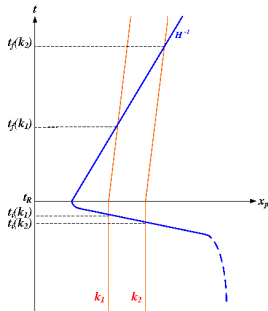
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N.B. Perturbations originate as thermal string gas fluctuations.

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Conclusions

- Calculate matter correlation functions in the Hagedorn phase (neglecting the metric fluctuations)
- For fixed k , convert the matter fluctuations to metric fluctuations at Hubble radius crossing $t = t_i(k)$
- Evolve the metric fluctuations for $t > t_i(k)$ using the usual theory of cosmological perturbations

Extracting the Metric Fluctuations

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Conclusions

Ansatz for the metric including cosmological perturbations and gravitational waves:

$$ds^2 = a^2(\eta) \left((1 + 2\phi) d\eta^2 - [(1 - 2\phi)\delta_{ij} + h_{ij}] dx^i dx^j \right).$$

Inserting into the perturbed Einstein equations yields

$$\langle |\phi(k)|^2 \rangle = 16\pi^2 G^2 k^{-4} \langle \delta T^0_0(k) \delta T^0_0(k) \rangle,$$

$$\langle |h(k)|^2 \rangle = 16\pi^2 G^2 k^{-4} \langle \delta T^i_j(k) \delta T^i_j(k) \rangle.$$

Power Spectrum of Cosmological Perturbations

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Key ingredient: For **thermal fluctuations**:

$$\langle \delta \rho^2 \rangle = \frac{T^2}{R^6} C_V.$$

Key ingredient: For **string thermodynamics** in a compact space

$$C_V \approx 2 \frac{R^2 / \ell_s^3}{T (1 - T/T_H)}.$$

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Power spectrum of cosmological fluctuations

$$\begin{aligned}P_{\Phi}(k) &= 8G^2 k^{-1} \langle |\delta\rho(k)|^2 \rangle \\&= 8G^2 k^2 \langle (\delta M)^2 \rangle_R \\&= 8G^2 k^{-4} \langle (\delta\rho)^2 \rangle_R \\&= 8G^2 \frac{T}{\ell_s^3} \frac{1}{1 - T/T_H}\end{aligned}$$

Key features:

- **scale-invariant** like for inflation
- **slight red tilt** like for inflation

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Prediction: Spectrum of Gravitational Waves

R.B., A. Nayeri, S. Patil and C. Vafa, *Phys. Rev. Lett.* (2007)

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Key ingredient for **string thermodynamics**

$$\langle |T_{ij}(R)|^2 \rangle \sim \frac{T}{\ell_s^3 R^4} (1 - T/T_H)$$

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Requirements

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Conclusions

- Static Hagedorn phase (including static dilaton) \rightarrow new physics required.
- $C_V(R) \sim R^2$ obtained from a thermal gas of strings provided there are winding modes which dominate.
- Cosmological fluctuations in the IR are described by Einstein gravity.

Conclusions concerning String Gas Cosmology

R.B. and C. Vafa, *Nucl. Phys. B*316:391 (1989)

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Conclusions

- Consistent with swampland criteria.
- Emerges from string theory considerations.
- Consistent with the TCC

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Conclusions

- Inflation does not **naturally** emerge from superstring theory.
- Inflation is **highly fine tuned** assuming validity of the TCC.
- Ekpyrosis and String Gas Cosmology are consistent with the swampland criteria and the TCC.
- **Alternatives to Inflation** appear more promising in light of fundamental physics.

Conclusions Concerning Inflation

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- Dark Energy **cannot be a Cosmological Constant.**
- **Quintessence as Dark Energy** is constrained.
- It is likely that **radical new ideas** are required in order to explain Dark Energy.

Conclusions Concerning Dark Energy

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Size Moduli [S. Watson, 2004; S. Patil and R.B., 2004, 2005]

- winding modes prevent expansion
- momentum modes prevent contraction
- $\rightarrow V_{\text{eff}}(R)$ has a minimum at a finite value of R , $\rightarrow R_{\text{min}}$
- in heterotic string theory there are **enhanced symmetry states** containing both momentum and winding which are massless at R_{min}
- $\rightarrow V_{\text{eff}}(R_{\text{min}}) = 0$
- \rightarrow **size moduli stabilized** in Einstein gravity background

Shape Moduli [Y-K. Cheung, S. Watson and R.B., 2005]

- enhanced symmetry states
- \rightarrow harmonic oscillator potential for θ
- \rightarrow **shape moduli stabilized**

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Dilaton stabilization in SGC

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Conclusions

- The only remaining modulus is the dilaton.
- Make use of **gaugino condensation** to give the dilaton a potential with a unique minimum.
- \rightarrow dilaton is stabilized.
- Dilaton stabilization is consistent with size stabilization [R. Danos, A. Frey and R.B., 2008].
- Gaugino condensation induces (high scale) **supersymmetry breaking** [S. Mishra, W. Xue, R.B. and U. Yajnik, 2012].