Coupling to matter in degenerate scalar-tensor theories

Sebastian Garcia-Saenz

Imperial College London

Based on 2004.11619 with C. Deffayet

Outline and results

- Degenerate higher-order scalar-tensor theories (DHOST)
 - Review

see also N. Frusciante's talk

- Matter coupling problem
- Pathological matter fields
 - Loss of degeneracy
 - Non-commutation of constraints
- Spinor fields in DHOST
 - Not pathological (so far)

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have drawbacks

- tension with fifth-force experiments
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- This motivates the study of generalized scalar-tensor theories

$$\mathcal{L} = F(\phi,
abla \phi,
abla^2 \phi)$$

Horndeski (1974)

These models are potentially pathological

$$\mathcal{L} = F(\phi, \nabla \phi, \nabla^2 \phi) \qquad \stackrel{?}{\longrightarrow} \qquad \frac{\delta \mathcal{L}}{\delta \phi} \supset \partial^3 \phi \,, \, \partial^4 \phi$$

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Nicolis et al. (2008) Deffayet et al. (2009)

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- Many known cases evade this issue
 - Second-order eqs. of motion

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Higher-order eqs. of motion + degeneracy

Beyond Horndeski, DHOST Gleyzes, Langlois, Piazza, Vernizzi (2015) Crisostomi, Koyama, Tasinato (2016) Ben Achour, Langlois, Noui (2016)

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16(diff symmetry) + 8 = 24

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Result

 $\mathbf{2}+\mathbf{1}$ dynamical degrees of freedom

Key observation is that matter fields may spoil constraints

de Rham & Matas (2016)

DHOST 3 DoF matter N DoF DHOST + matter $\geq 3 + N$ DoF

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We will say that the coupling to a given matter sector is inconsistent if the total degree of freedom count is

> 3 + N

If the matter sector is consistent within GR, but inconsistent within DHOST, then we will conclude that the ghost has reappeared

Remarks

Already in GR not every matter field is consistent

Isenberg & Nester (1977)

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Our framework

- Minimal coupling to the metric (but in a general frame)
- Quadratic DHOST $\mathcal{L} \supset \nabla^2 \phi$, $(\nabla^2 \phi)^2$

Two ways for matter to spoil constraints

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(I) Hessian matrix

$$\mathcal{H}_{ij} = rac{\partial^2 \mathcal{L}}{\partial \dot{\Psi}^i \partial \dot{\Psi}^j} , \qquad \Psi^i \equiv \{ g_{\mu\nu} , \ \phi \ , \
abla_\mu \phi \ , \ ext{matter} \}$$

fails to have the required rank

 \rightarrow direct loss of degeneracy (primary) constraint

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Example: non-Maxwell vector field

$$\mathcal{L}_m = \sqrt{-g} \, \nabla^\mu B^\nu \nabla_\mu B_\nu$$

 $\mathsf{rank}\; \mathcal{H}^{(\mathrm{DHOST})} + \mathsf{rank}\; \mathcal{H}^{(\mathrm{matter})} \; < \; \mathsf{rank}\; \mathcal{H}$

because ${\mathcal H}$ is not block-diagonal

Two ways for matter to spoil constraints

(II) Some constraint in the matter sector fails to Poisson-commute with the DHOST constraint

$$\left\{ \mathcal{C}^{(\mathrm{DHOST})} \;,\; \mathcal{C}^{(\mathrm{matter})} \right\} \neq 0$$

 $\rightarrow~$ loss of secondary constraints

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Example: cubic Galileon/KGB

$$\mathcal{L}_m = \sqrt{-g} \left[-\frac{1}{2} (\nabla \pi)^2 + \kappa (\nabla \pi)^2 \Box \pi \right]$$

Hessian constraint \checkmark Poisson bracket constraint X

Spinor fields

Spinor fields in DHOST are *potentially* problematic

- they couple to the spin connection
 - \rightarrow Hessian matrix not block-diagonal
- they have constraints
 - $\rightarrow\,$ must Poisson-commute with DHOST constraint

$$\mathcal{L}_{m}=-\frac{1}{2}\sqrt{-g}\,\bar{\lambda}\gamma^{\mu}\left(\partial_{\mu}+\Omega_{\mu}\right)\lambda$$

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We have shown that a large class of spinor models is actually consistent

- any quadratic action with arbitrarily many Majorana and/or Dirac spin-1/2 fields
- most general self-interacting Majorana spin-1/2 field, linear in $abla \lambda$

- Beyond minimal coupling
 - curvature couplings
 - scalar-matter couplings $\phi T^{\mu}_{\ \mu}$, $\partial_{\mu} \phi \partial_{\nu} \phi T^{\mu\nu}$

- disformal couplings $\mathcal{L}_m[Ag_{\mu\nu} + B\partial_\mu\phi\partial_\nu\phi, \psi_m]$

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Cubic DHOST

Ben Achour et al. (2016)

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- Other spinor fields e.g. spin-3/2
- Other physically relevant matter fields e.g. fluids with higher derivative corrections

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Thank you