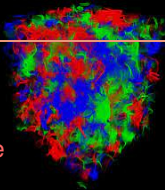




# Lattice QCD with an inhomogeneous magnetic field background

XQCD 2022, Trondheim, Norway



Dean Valois

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Gergely Endrődi Bastian Brandt Gergely Marko Francesca Cuteri

July 27, 2022

Department of Physics  
Bielefeld University

# OUTLINE

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1. Strongly magnetized physical systems
2. Magnetic field on the lattice
3. Lattice simulations
4. Summary & Conclusions

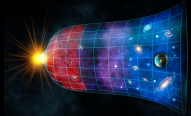
# **Strongly magnetized physical systems**

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# STRONGLY MAGNETIZED PHYSICAL SYSTEMS

Early universe

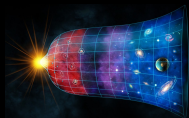
$$\sqrt{eB} \sim 1.5 \text{ GeV}$$



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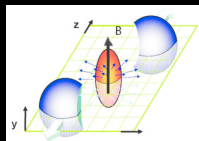
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Heavy-ion collision

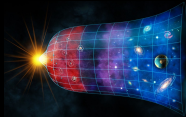
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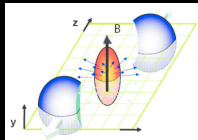
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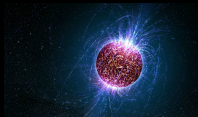
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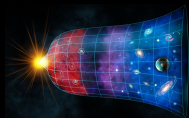
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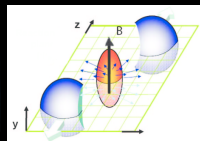
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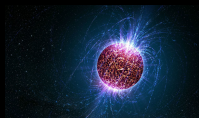
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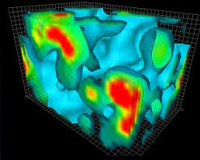


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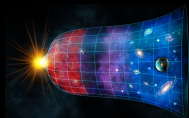


QCD vacuum

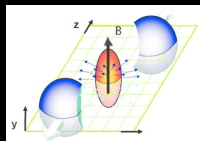


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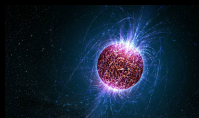
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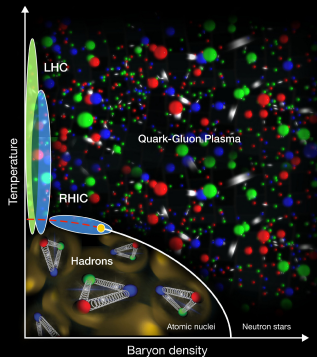
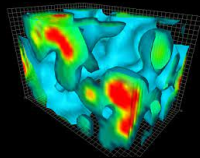
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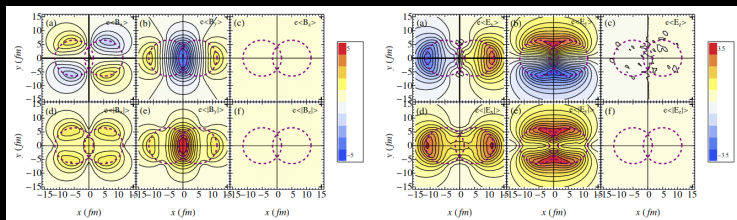





# MAGNETIC FIELDS IN HEAVY-ION COLLISIONS

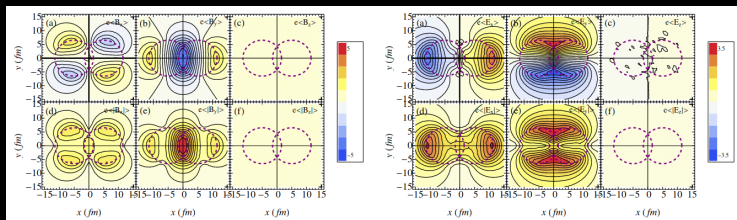
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**Figure 2:** Spatial distributions of the electric (right) and magnetic (left) fields in the transverse plane for an impact parameter  $b = 10$  fm  Deng and Huang 2012.

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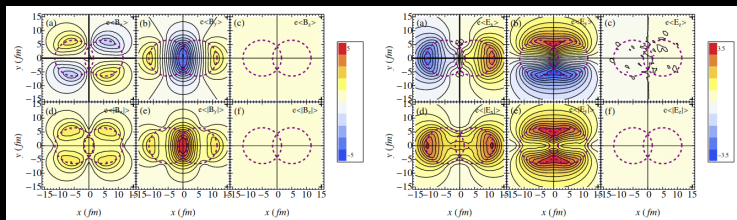



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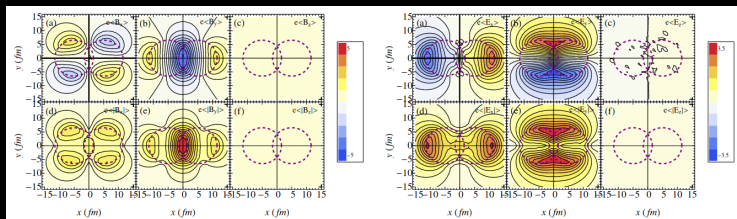



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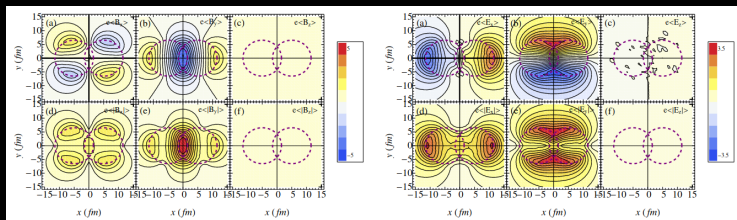



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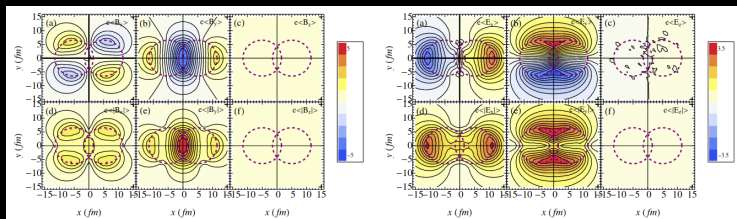
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$B(x)$  as  
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# Magnetic field on the lattice

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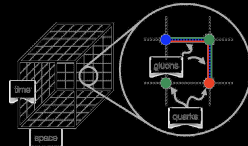


# UNIFORM MAGNETIC FIELD ON THE LATTICE

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gluon fields  $\longrightarrow U_\mu = e^{iagA_\mu^b T_b} \in \text{SU}(3)$

(anti-)periodic BC



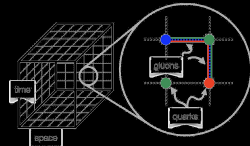
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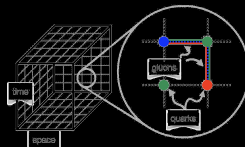
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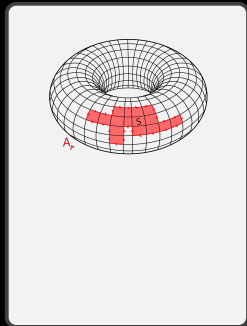
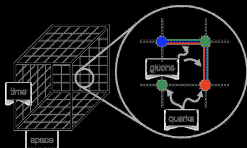
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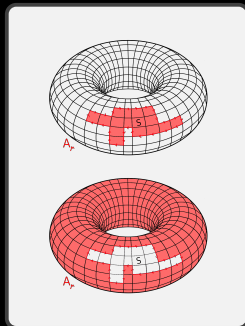
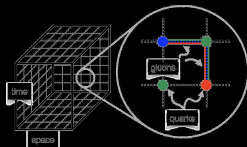
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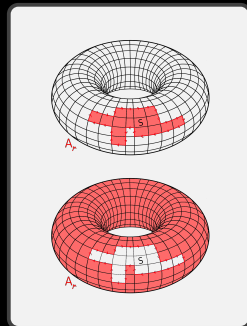
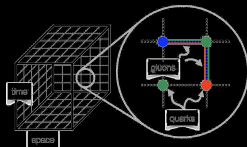
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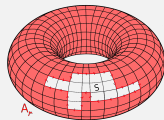
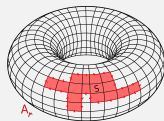
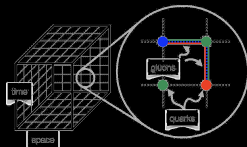
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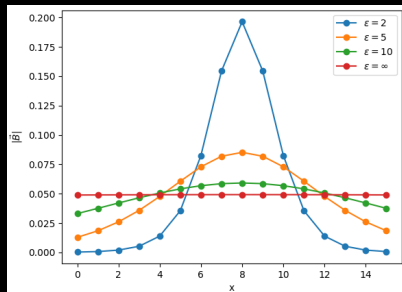
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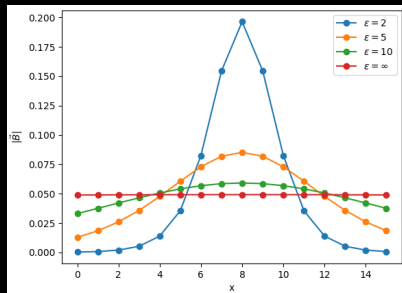


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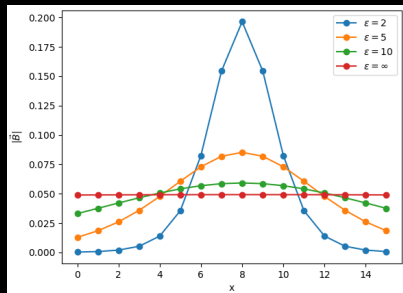


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# Lattice simulations

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- Temperature range  $68 \text{ MeV} \leq T \leq 300 \text{ MeV}$  (crossover transition at  $T_c \sim 155 \text{ MeV}$ ).



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- Local Polyakov loop

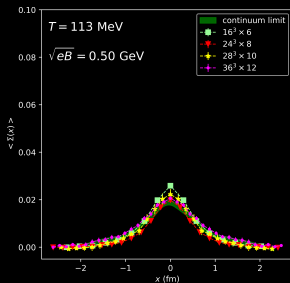
$$P = \frac{1}{L_x L_y} \sum_{y,z} \text{Re Tr} \prod_n U_t(x, y, z, n) \xrightarrow{\text{renormalization}} \frac{P(x, T, B)}{P(x, T, 0)}$$

- Local electric currents (**u**, **d** and **s** quarks!)

$$\langle J_i(x) \rangle = e \left\langle \frac{2}{3} \bar{u} \gamma^i u - \frac{1}{3} \bar{d} \gamma^i d - \frac{1}{3} \bar{s} \gamma^i s \right\rangle$$

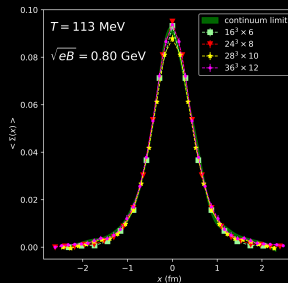
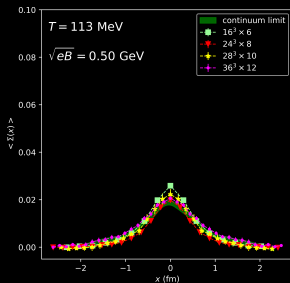
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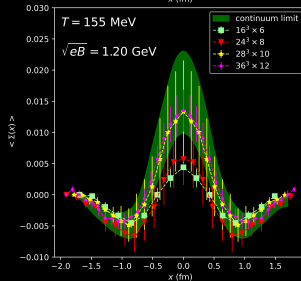
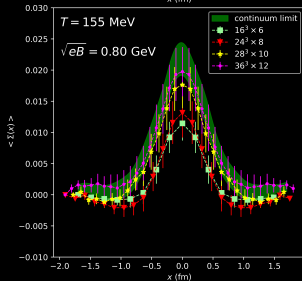
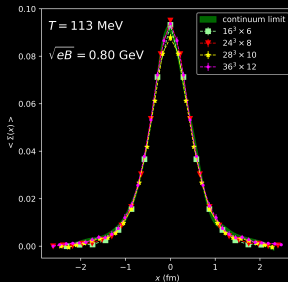
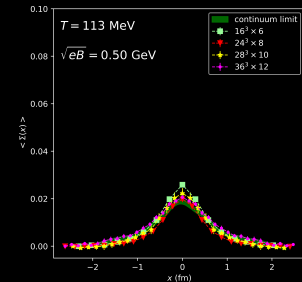




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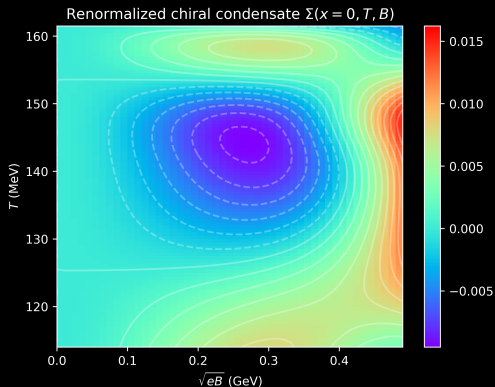


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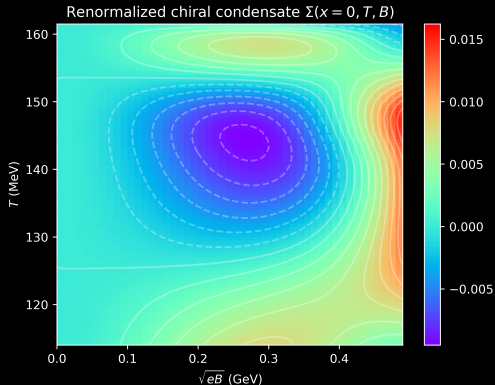
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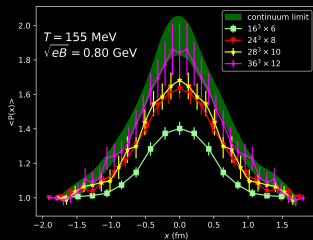
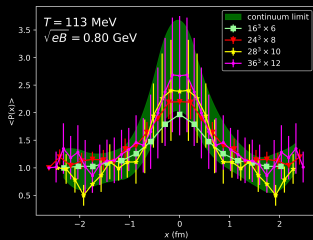
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- Magnetic catalysis  $T$  away from  $T_c$
- Inverse catalysis for  $T$  around  $T_c$   Endródi et al. 2019

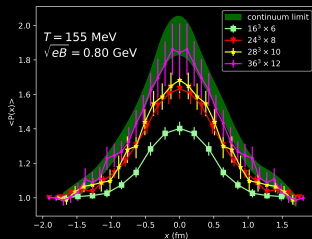
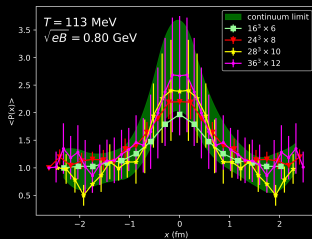
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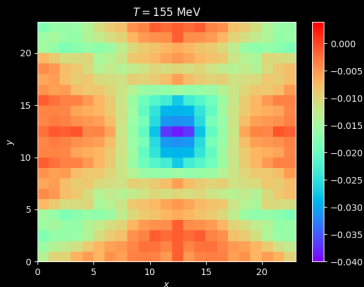
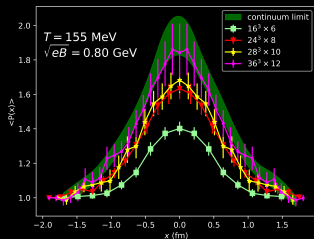
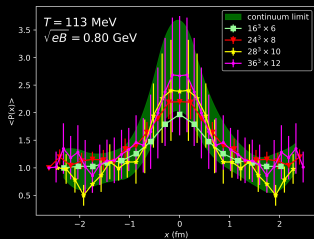




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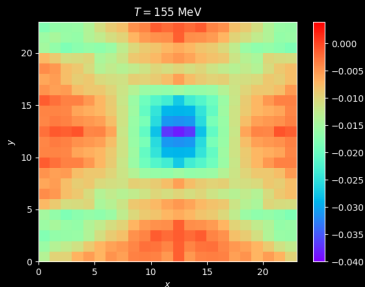
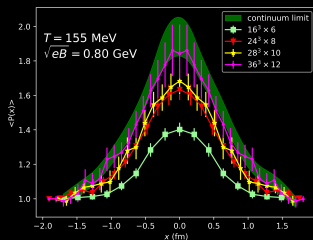
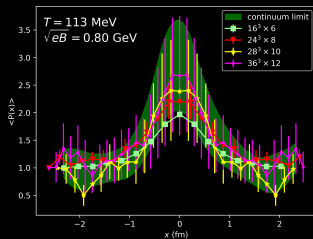
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The interaction of the condensate with  $P$  causes the dips!

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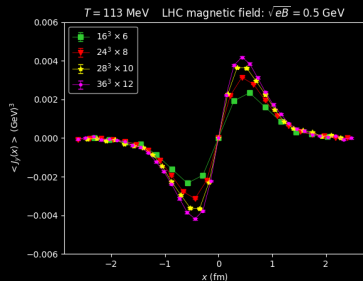
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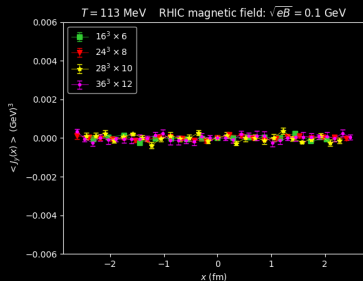
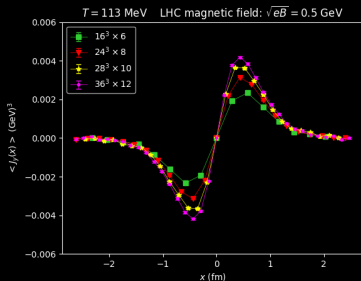
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**Figure 6:** Lattice electric currents for LHC-like ( $\sqrt{eB} = 0.5 \text{ GeV}$ ) and RHIC-like ( $\sqrt{eB} = 0.1 \text{ GeV}$ ) magnetic fields, respectively.



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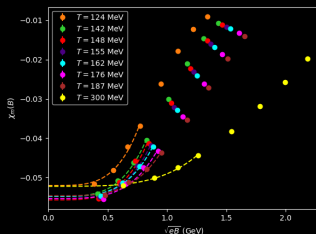
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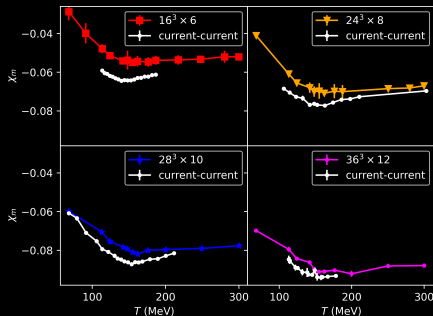
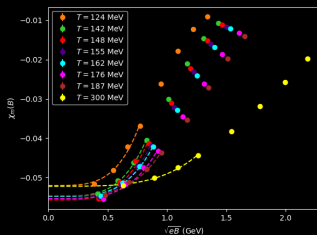
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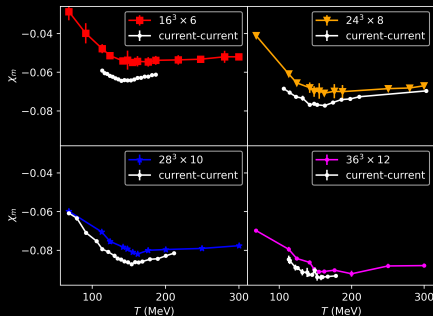
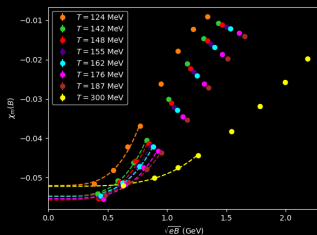
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The susceptibility contains an additive divergence  $\chi_m \sim \log(a)$

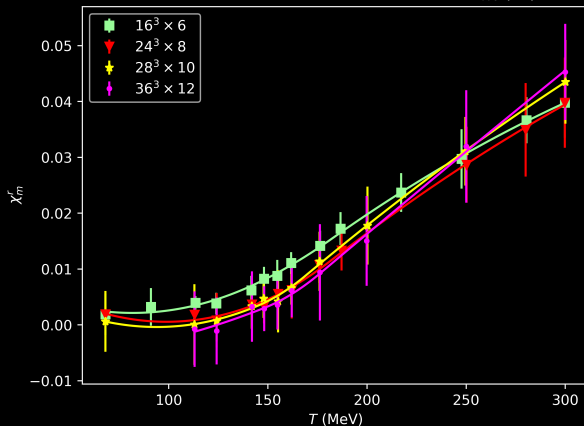
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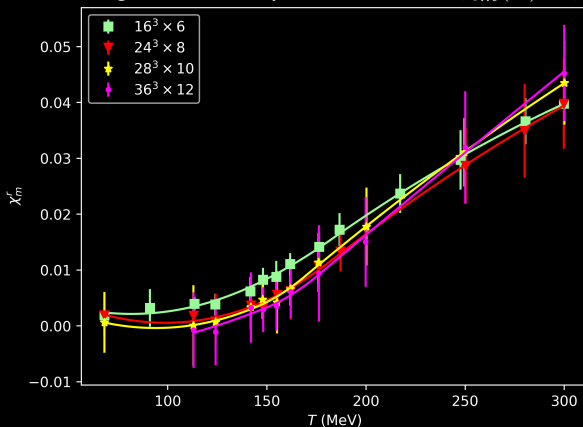
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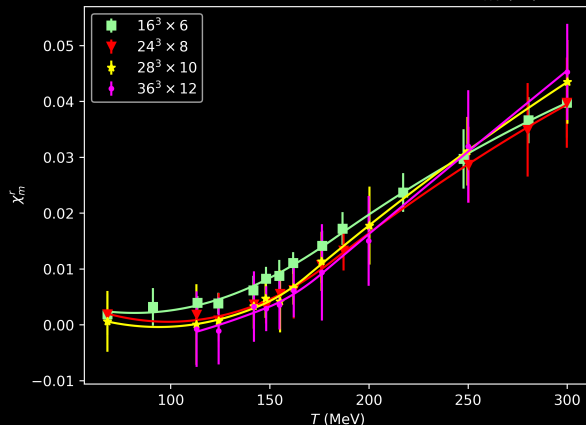
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Great agreement with the current-current method!  Bali, Gergely Endrődi,

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More on electromagnetic

- fields in lattice QCD:  
posters by



J. J. H. Hernandez




E. Garnacho Velasco

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