Holographic approach to dense QCD and neutron star mergers

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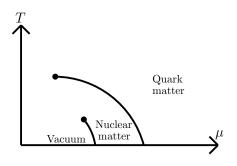
[Tuna Demircik, Christian Ecker, MJ arXiv:2112.12157] [Samuel Tootle, Christian Ecker, Konrad Topolski, Tuna Demircik, MJ, Luciano Rezzolla arXiv:2205.05691] + [earlier work]



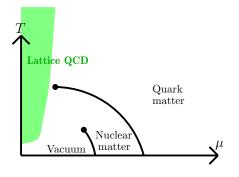
Outline

- 1. Introduction and motivation
- 2. Dense QCD matter in V-QCD▶ Both guark and nuclear matter
- 3. "Hybrid" Equations of State (EoSs)
 - Combining V-QCD with other approaches
 - Model at finite temperature and density
- 4. Application to neutron star mergers
 - Production of quark matter

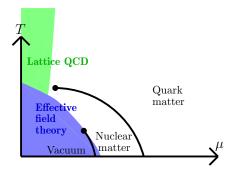
1. Introduction



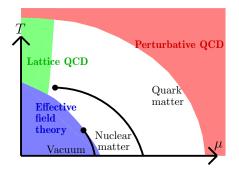
Lattice data only available at zero/small chemical potentials



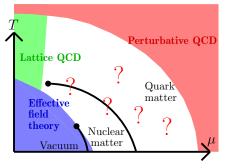
Lattice data only available at zero/small chemical potentials
 Effective field theory works at small densities



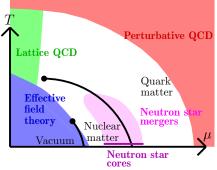
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- Perturbative QCD: only at high densities and temperatures



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- Perturbative QCD: only at high densities and temperatures
- Open questions at intermediate densities

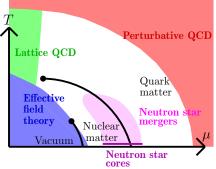


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- 1. Improving theoretical predictions important!
- 2. Incoming experimental data from neutron star measurements!

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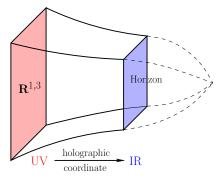


1. Improving theoretical predictions important!

2. Incoming experimental data from neutron star measurements! White region strongly coupled \Rightarrow use holography? Interpolate between known results using holography?

Gauge/gravity duality for QCD

- Motivated by the original AdS/CFT correspondence for N = 4 SYM
- Instead of conformality, confinement: non-AdS/non-CFT duality
- Field theory lives on the boundary of the 5D geometry



• Operators $O_i(x^{\mu}) \leftrightarrow$ classical bulk fields $\phi_i(x^{\mu}, r)$

$$Z_{\text{grav}}(\phi_i|_{\text{bdry}} = J_i(x^{\mu})) = \int \mathcal{D} e^{iS_{QCD} + i \int d^4 x J^i(x^{\mu})O_i(x^{\mu})}$$

 $\blacktriangleright \text{ E.g. } \bar{\psi}^{j}\psi^{i} \leftrightarrow \phi^{ij} \qquad T_{\mu\nu} \leftrightarrow g_{\mu\nu} \qquad J_{\mu} \leftrightarrow A_{\mu}$

► Thermodynamics of QCD ↔ thermodynamics of a planar bulk black hole

2. Dense holographic QCD matter

Holographic V-QCD

- A holographic model for QCD
 - Bottom-up, but trying to follow principles from string theory closely [MJ, Kiritsis 1112.1261; Review MJ 2110.08281]

The model is obtained through a fusion of two building blocks:

- 1. IHQCD: model for glue inspired by string theory [Gürsoy, Kiritsis, Nitti; Gubser, Nellore]
- Adding flavor and chiral symmetry breaking via a D-brane setup [Klebanov,Maldacena; Bigazzi,Casero,Cotrone,Iatrakis,Kiritsis,Paredes]
 Two bulk scalars: λ ↔ TrF², τ ↔ q̄q

 $S_{V-QCD} = N_c^2 M^3 \int d^5 x \sqrt{g} \Big[R - \frac{4}{3} \frac{(\partial \lambda)^2}{\lambda^2} + V_g(\lambda) \Big] \\ - N_f N_c M^3 \int d^5 x V_{f0}(\lambda) e^{-\tau^2} \sqrt{-\det(g_{ab} + \kappa(\lambda)\partial_a \tau \partial_b \tau + w(\lambda)F_{ab})}$

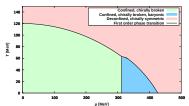
Effective model, many potentials V_g , V_{f0} , w, κ – essential to fix them by fitting QCD data \rightarrow predictions for other observables

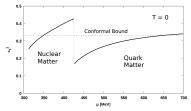
Quark matter:

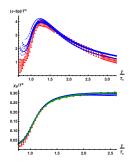
- Require agreement with confinement, asymptotic freedom . . .
- (Good) fit to lattice data at $\mu \approx 0$
- Extrapolate to high μ and down to $T \approx 0$
- Extrapolated results consistent with pQCD and EFT! [MJ, Jokela, Remes, 1809.07770]

Nuclear matter:

- In holographic setups, baryons ↔ solitons of gauge fields in the gravity dual [Talk and poster by Edwan Préau]
- ► However, solitons technically involved ... use a simple homogeneous approximation ≈ smeared solitons [Ishii, MJ, Nijs 1903.06169]
- Stiff EoS: "violation" of the conformal bound $c_s^2 = 1/3!$







3. Hybrid EoSs

Combining with other approaches

The V-QCD EoS as such is however not fully satisfactory:

- 1. Our (homogeneous) approach for nuclear matter only works at high densities
- 2. Temperature dependence is trivial in the confined phases, and therefore also for holographic nuclear matter

This is a large N_c issue, T dependence would arise from loops Solutions:

1. At low densities for nuclear matter, use "traditional" nuclear theory results

 \Rightarrow choose the Hempel-Schaffner-Bielich model with DD2 interactions (HS(DD2))

[Typel et al. 0908.2344; Hempel, Schaffner-Bielich 0911.4073]

2. Since no reliable results available, borrow T dependence from basically the simplest reasonable model

⇒ use van der Waals (vdW) gas (protons, neutrons, electrons) [Ecker, MJ, Nijs, van der Schee 1908.03213] [Jokela, MJ, Nijs, Remes 2006:01141] [Demircik, Ecker, MJ 2112.12157]

Overview of the hybrid model

Meson gas

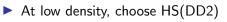
HS(DD2)

vdW NM

APR V-QCD NN

'n,

- V-QCD for quark matter and cold dense nuclear matter
- Van der Waals model extrapolates dense V-QCD nuclear matter to finite T



- At medium density, use APR cold EoS (using only HS(DD2) would lead to tension with neutron star observations)
- Add QCD mesons to HS(DD2), important to describe the critical point

[Demircik, Ecker, MJ 2112.12157]

V-OCD OM

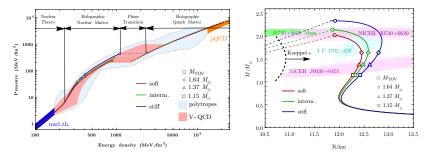
Mixed

phase

ní

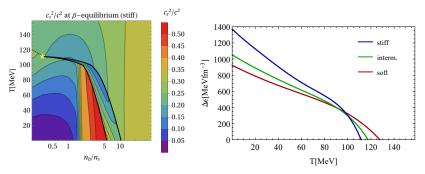
Cold EoS and known constraints

- ► Three choices of EoSs: soft, intermediate, and stiff ↔ the degrees of freedom of V-QCD left free by fit to lattice data
- Compared to bands of all feasible cold matter EoS: Without and with holography



Plug EoSs in TOV: neutron star M(R) curves (left plot)
 Compares well with mass/radius observations

Results: phase transition and critical point



- Low T: strong 1st order nuclear to quark matter transition and mixed phase
- High T: weak first order transition \approx crossover
- Critical point with 110 MeV $\lesssim T_c \lesssim 130$ MeV 480 MeV $\lesssim \mu_{bc} \lesssim 580$ MeV

 Close to results in other (simpler) holographic models [DeWolfe et al. 1012.1864; Knaute et al. 1702.06731; Critelli et al. 1706.00455]_{13/19}

4. (Holographic) Neutron Star Mergers

Simulating Binary Neutron Star Mergers

Have to solve the 3+1D General Relativistic hydrodynamics equations:

$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} = 8\pi G_N T_{\mu\nu} , \quad \nabla_\mu T^{\mu\nu} = 0 , \quad \nabla_\mu J^\mu = 0$$

with initial spacetime and fluid distribution modelling a NS binary system

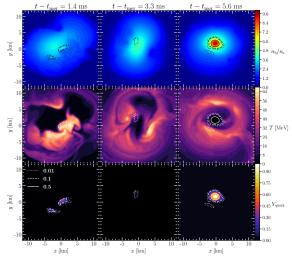
- Equation of State $p = p(n_b, T, Y_e)$ as input use V-QCD hybrid EoS
- Spectral code Frankfurt University/Kadath (FUKA) for initial data [Papenfort, Tootle, Grandclement, Most, Rezzolla 2103.09911]
- Frankfurt/Illinois (FIL) code for binary evolution with tabulated EoS [Most, Papenfort, Rezzolla 1907.10328]
- Implemented in the Einstein Toolkit

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[http://einsteintoolkit.org]
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Need supercomputing: Project BNSMIC with 100 million core-hours on HAWK at the High-Performance Computing Center Stuttgart

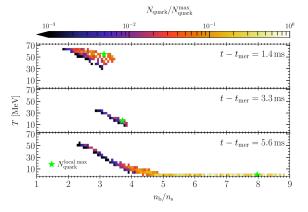
Hot, Warm and Cold Quarks

Simulations with parameters chosen to match with GW170817



Hypermassive neutron star after merger (soft EoS, $M_1/M_2 = 0.7$) [Tootle, Ecker, Topolski, Demircik, MJ, Rezzolla 2205.05691] 16/19 Hot quarks: in the hottest region at early times

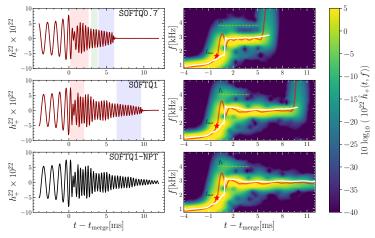
- Warm quarks: at intermediate times due to complicated post-merger dynamics
- Cold quarks: in the densest core at late times



[Tootle, Ecker, Topolski, Demircik, MJ, Rezzolla 2205.05691]

details

Imprint on Gravitational Waves



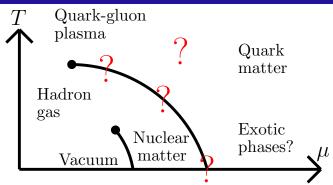
- Most significant signature of the phase transition: short lifetime of remnant
- ► Early collapse in tension with electromagnetic signal from GW170817 ⇒ constrains the EoS – soft model disfavored

Summary

- (Effective) holography, combined with other approaches, is useful to study dense QCD
- Using V-QCD with simple approximations, many details work really well:
 - $\checkmark\,$ Precise fit of lattice thermodynamics at $\mu\approx$ 0
 - ✓ Extrapolated EoS for cold quark matter reasonable
 - ✓ Simultaneous model for nuclear and quark matter
 - ✓ Stiff EoS for nuclear matter
- We constructed an EoS at finite temperature and density using V-QCD (+other models)
 - Predictions for the critical point
 - Input for merger simulations
- State-of-the-art binary neutron star merger simulations with our EoS
 - Production of hot, warm and cold quark matter

Thank you!

The QCD phase diagram



Focus in this talk: phases at high density

- \blacktriangleright Nuclear matter: dense liquid of protons and neutrons density \gtrsim density of large nuclei
- Quark matter: densely packed phase of free quarks and gluons

Laboratory experiments challenging ($T_{QCD} \sim 10^{12}~{\rm K})$, in particular at high density – lots of effort

Recent and future progress: LHC, RHIC, FAIR, NICA, ...

Neutron stars

Neutron stars: extremely dense cold QCD matter

- Tolman-Oppenheimer-Volkoff (TOV) equations map equation of state (EoS) to mass-radius relation $^{2\rho}$
- EoS can be constrained by measuring masses and radii

Mass measurements: dozens of results using various methods

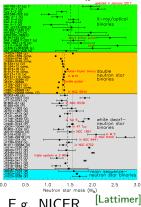
Highest masses from Shapiro delay measurement of NS – white dwarf binaries J0348+0432 and J0740+6620: Antoniadis et al 1304.6875 $M_{\rm max} \gtrsim 2 M_{\odot}$

Cromartie et al 1904.06759]

Radius measurements: more challenging, high uncertainties

► Cooling after X-ray bursts ⇒ radii around 10-15 km

More and better results expected in near future! E.g. NICER



Crust e Z n

Outer Core n-p Fermi liquid

Inner

Core

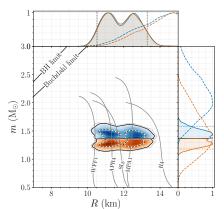
 $\sim 0.5a$

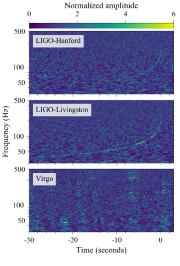
//~1-2km

-10km

LIGO/Virgo constraints from GW170817

- The tidal deformability A measures how strongly neutron stars deform in gravitational field
- Inspiral phase GW signal gives an upper bound Λ ≤ 580
- Implies a rough upper bound for neutron star radius: $R \lesssim 13.5$ km

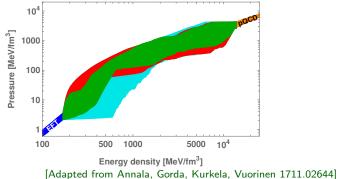




Constraints on equation of state (EoS)

State of the art for QCD EoS at T = 0: interpolations between nuclear EoS and pQCD, constrained by

- 1. Mass bound $M_{
 m max} > 2 M_{\odot}$ (excludes cyan area)
- 2. LIGO constraint from GW170817: (excludes red area)



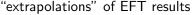
Source of uncertainties: physics at strong coupling \Rightarrow Can holographic methods be used to reduce uncertainties further?

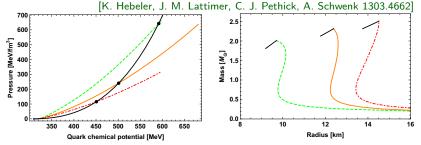
Recent progress on dense holographic QCD

For quark matter, use D3-D7 top down model: $\epsilon = 3p + \frac{\sqrt{3}m^2}{2\pi}\sqrt{p}$ [Karch, O'Bannon, 0709.0570]

▶ N = 4 SYM + $N_f = 3$ probe hypermultiplets in the fundamental representation

For nuclear matter use with stiff, intermediate, and soft





Strong first order nuclear to quark matter transitions

 Neutron stars with "holographic" quark matter core (black curves) are unstable

[Hoyos, Rodriguez, Jokela, Vuorinen 1603.02943]25/19

Varying the quark mass *m* one can get quark stars and hybrid stars [Annala, Ecker, Hoyos, Jokela, Rodriguez-Fernandez, Vuorinen 1711.06244]

Sizeable deviations from universal I-Love-Q relations
 [Yagi, Yunes, 1303.1528]

Including running of the quark mass + color superconductivity [Bitaghsir Fadafan, Cruz Rojas, Evans, 1911.12705; 2009.14079]

- ▶ Possibility of an intermediate χ SB deconfined phase
- Stiffer holographic equations of state (high speed of sound)
- Quark matter cores

Using Einstein-Maxwell-dilaton for quark matter [Mamani, Flores, Zanchin, 2006.09401]

(Largish) quark stars also studied in Witten-Sakai-Sugimoto and in D4-D6 models [Burikham, Hirunsirisawat, Pinkanjanarod, 1003.5470 Kim, Shin, Lee, Wan, 1108.6139, 1404.3474]

This talk: towards more realistic model of quark matter?

Ansatz for potentials, (x = 1)

$$\begin{split} V_{g}(\lambda) &= 12 \left[1 + V_{1}\lambda + \frac{V_{2}\lambda^{2}}{1 + \lambda/\lambda_{0}} + V_{\mathrm{IR}}e^{-\lambda_{0}/\lambda}(\lambda/\lambda_{0})^{4/3}\sqrt{\log(1 + \lambda/\lambda_{0})} \right] \\ V_{f0}(\lambda) &= W_{0} + W_{1}\lambda + \frac{W_{2}\lambda^{2}}{1 + \lambda/\lambda_{0}} + W_{\mathrm{IR}}e^{-\lambda_{0}/\lambda}(\lambda/\lambda_{0})^{2} \\ \frac{1}{w(\lambda)} &= w_{0} \left[1 + \frac{w_{1}\lambda/\lambda_{0}}{1 + \lambda/\lambda_{0}} + \bar{w}_{0}e^{-\lambda_{0}/\lambda_{W_{s}}}\frac{(w_{s}\lambda/\lambda_{0})^{4/3}}{\log(1 + w_{s}\lambda/\lambda_{0})} \right] \\ V_{1} &= \frac{11}{27\pi^{2}} , \quad V_{2} &= \frac{4619}{46656\pi^{4}} \\ W_{1} &= \frac{8 + 3W_{0}}{9\pi^{2}} ; \quad W_{2} &= \frac{6488 + 999W_{0}}{15552\pi^{4}} \end{split}$$

Fixed UV/IR asymptotics \Rightarrow fit parameters only affect details in the middle

Constraining the potentials

In the UV ($\lambda \rightarrow 0$):

► UV expansions of potentials matched with perturbative QCD beta functions ⇒ asymptotic freedom and logarithmic flow of the coupling and quark mass, as in QCD

[Gürsoy, Kiritsis 0707.1324; MJ, Kiritsis 1112.1261]

In the IR $(\lambda \to \infty)$: various qualitative constraints

- Linear confinement, discrete glueball & meson spectrum, linear radial trajectories
- Existence of a "good" IR singularity
- Correct behavior at large quark masses
- Working potentials often string-inspired power-laws, multiplied by logarithmic corrections (i.e, first guesses usually work!)

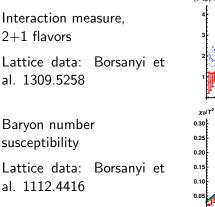
[Gürsoy, Kiritsis, Nitti 0707.1349; MJ, Kiritsis 1112.1261; Arean, Iatrakis, MJ, Kiritsis 1309.2286, 1609.08922; MJ 1501.07272]

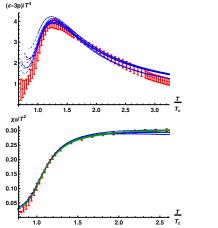
Final task: determine the potentials in the middle, $\lambda = \mathcal{O}(1)$

Qualitative comparison to lattice/experimental data

Constraining the model at $\mu\approx 0$

- Many parameters already fixed by requiring qualitative agreement with QCD
- Good description of lattice data nontrivial result!





29/19

Extrapolated EoSs of cold quark matter

The V-QCD cold quark matter result compares nicely to known constraints:

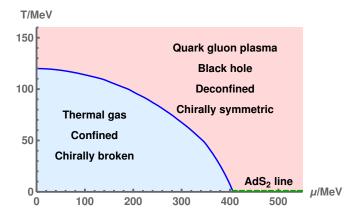


Band of allowed 10⁴ equations of state Pressure [MeV/fm³. 1000 (EoSs) (gray, polytropic interpolations) 100 Stiff. intermediate. and 10 soft nuclear EoSs [Hebeler, Lattimer, Pethick, 10⁴ 100 500 1000 5000 Schwenk 1303.4662] Energy density [MeV/fm³]

Approach similar in spirit to studies of the QCD critical point

[DeWolfe,Gubser,Rosen 1012.1864; Knaute,Yaresko,Kämpfer 1702.06731; Critelli, Noronha, Noronha-Hostler, Portillo, Ratti, Rougemont, 1706.00455]

Phase diagram with quark matter



- With quark matter only, expected phase diagram
- Cold QM equation of state (EoS) and location of the T = 0 phase transition agree with contraints

Nuclear matter in holographic models

Each baryon maps to a solitonic "instanton" configuration of gauge fields in the bulk [Witten: Gross, Ooguri; ...]

- Such instantons have been studied in many models, including V-QCD (see the poster/talk by Edwan Préau)
- However, dense nuclear matter requires studying many-instanton solutions
- Extremely challenging!
- This talk: set N_f = 2 and try first a simple approximation scheme (homogeneous), reasonable at high densities?
 [Rozali, Shieh, Van Raamsdonk, Wu 0708.1322]
 Aⁱ = h(r)σⁱ

[Li,Schmitt,Wang 1505.04886; Elliot-Ripley,Sutcliffe,Zamaklar 1607.04832]

[Kovensky, Poole, Schmitt, 2111.03374]

Homogeneous nuclear matter in V-QCD

Nuclear matter in the probe limit: consider full brane action $S = S_{\text{DBI}} + S_{\text{CS}}$ where

[Bigazzi, Casero, Cotrone, Kiritsis, Paredes; Casero, Kiritsis, Paredes]

$$S_{\text{DBI}} = -\frac{1}{2}M^3 N_c \,\mathbb{T}r \int d^5 x \, V_{f0}(\lambda) e^{-\tau^2} \left(\sqrt{-\det A^{(L)}} + \sqrt{-\det A^{(R)}} \right) \\ A_{MN}^{(L/R)} = g_{MN} + \delta_M^r \delta_N^r \kappa(\lambda) \tau'(r)^2 + \delta_{MN}^{rt} w(\lambda) \Phi'(r) + w(\lambda) F_{MN}^{(L/R)} \\ \text{gives the dynamics of the solitons (will be expanded in } F^{(L/R)}) \text{ and}$$

$$S_{\rm CS} = \frac{N_c}{8\pi^2} \int \Phi(r) e^{-b\tau^2} dt \wedge \left(F^{(L)} \wedge F^{(L)} - F^{(R)} \wedge F^{(R)} + \cdots\right)$$

sources the baryon number for the solitons

Extra parameter, b > 1, to ensure regularity of solutions Set $N_f = 2$ and consider the homogeneous SU(2) Ansatz [Rozali, Shieh, Van Raamsdonk, Wu 0708.1322]

$$A_L^i = -A_R^i = h(r)\sigma^i$$

[Ishii, MJ, Nijs, 1903.06169]

Discontinuity and smeared instantons

With the homogeneous Ansatz $A_i^a(r) = h(r)\delta_i^a$ baryon number vanishes for any smooth h(r):

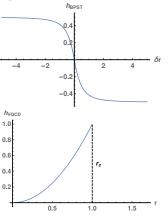
$$N_b \propto \int dr \frac{d}{dr} \left[\text{CS} - \text{term} \right] = 0$$

How can this issue be avoided?

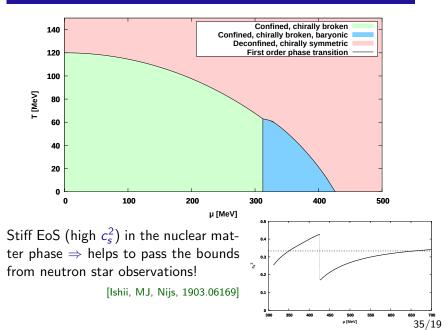
Smearing the BPST soliton in singular Landau gauge:

$$\langle A_i^a \rangle \sim \int \frac{d^3 \times \eta_{i4}^a \, \delta r}{(\delta r^2 + x^2 + \rho^2)(\delta r^2 + x^2)} \\ \sim -\frac{\delta_i^a \, \delta r}{\sqrt{\delta r^2 + \rho^2} + |\delta r|}$$

- This suggests a solution: introduce a discontinuity in h(r) at r = r_c
- The discontinuity sources nonzero baryon charge!



Phase diagram at zero quark mass



Van der Waals model

Ideal gas of protons, neutrons and electrons with

Excluded volume correction for nucleons

$$p_{\text{ex}}(T, \{\mu_i\}) = p_{\text{id}}(T, \{\tilde{\mu}_i\})$$

$$\tilde{\mu}_i = \mu_i - v_0 p_{\text{ex}}(T, \{\mu_i\}) \quad (i = p, n)$$

 $v_0 \sim$ volume of one nucleon

(Mostly) attractive potential term to match with (APR and)
 V-QCD at T = 0

$$p_{\mathrm{vdW}}(T, \{\mu_i\}) = p_{\mathrm{ex}}(T, \{\mu_i\}) + \Delta p(\{\mu_i\})$$

schematically:

$$\Delta p(\{\mu_i\}) = p_{\mathrm{V-QCD}}(T = 0, \{\mu_i\}) - p_{\mathrm{ex}}(T = 0, \{\mu_i\})$$

[Rischke, Gorenstein, Stoecker, Greiner, Z Phys. C 51, 485 (1991)]

[Vovchenko, Gorenstein, Stoecker, 1609.03975]

[Vovchenko, Motornenko, Alba, Gorenstein, Satarov, Stoecker, 1707.09215]

Hempel-Schaffner-Bielich DD2 model

A widely used general purpose model for the EoS

> Parameters: temperature, density, charge fraction Y_q

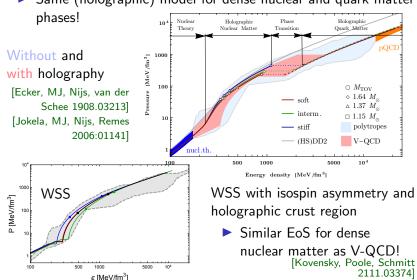
Combines two approaches (in thermodynamically consistent way):

- For n < n_s, statistical method with excluded volume corrections and interactions, including light and heavy nuclei [Hempel, Schaffner-Bielich, 0911.4073]
- For n > n_s, relativistic mean field theory of nucleons interacting with σ, ρ, and ω mesons (DD2)

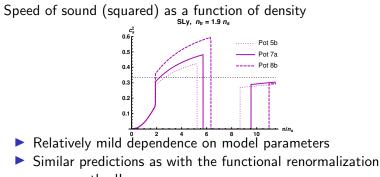
[Typel, Ropke, Klahn, Blaschke, Wolter, 0908.2344]

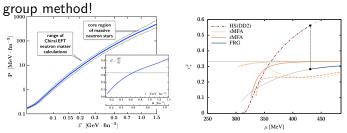
Results: Cold Hybrid Equations of State

Variations in model parameters give rise to the band
 Same (holographic) model for dense nuclear and quark matter



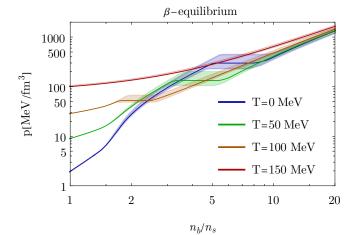
Speed of sound and comparison to FRG





[Drews, Weise 1610.07568; Otto, Oertel, Schaefer 1910.11929]39/19

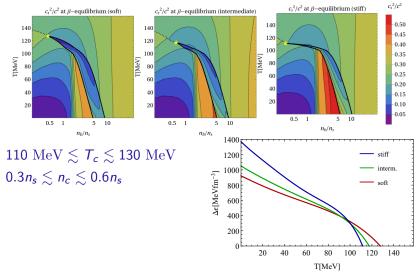
Results: EoS at Finite T



- Bands: variation of the V-QCD model (soft/intermediate/stiff)
- With increasing T, weaker transition at lower pressure [Demircik, Ecker, MJ 2112.12157]

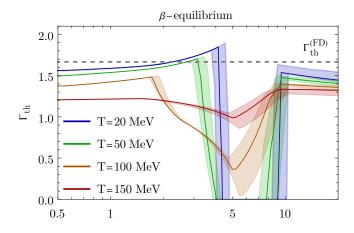
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Results: critical point



Critical point is determined by fitting the latent heat in the region of strong phase transition and extrapolating

Results: thermal index



 $\Gamma_{\rm th}(n_b, T) = 1 + rac{p(n_b, T) - p(n_b, 0)}{e(n_b, T) - e(n_b, 0)}$

- Values in expected range
- Low values in the mixed phase

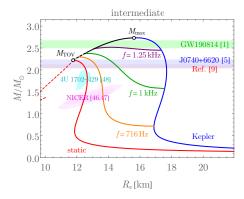
Rapidly spinning holographic neutron stars

GW190814: LIGO/Virgo observed a merger of a $23M_{\odot}$ black hole with a $2.6M_{\odot}$ compact object

[2006.12611]

▶ $2.6M_{\odot}$ falls in the "gap": a black hole or a neutron star?

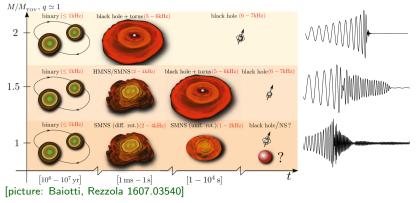
- Holographic EoSs easily compatible with the neutron star interpretation
- ► However requires fast rotation, f ≥ 1 kHz



[Demircik, Ecker, MJ, 2009.10731]

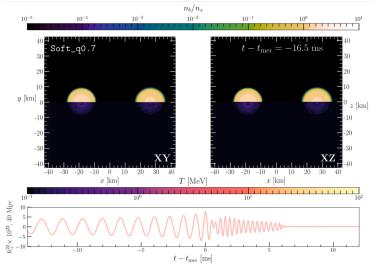
Neutron star mergers

- Significant sources of gravitational radiation
- Microscopic properties of dense matter encoded in GW and EM signal



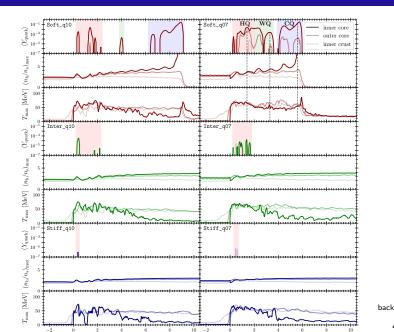
One good event (GW170817) and a few other events already observed! [LIGO/Virgo, 1710.05832]

Simulations with parameters chosen to match with GW170817



Soft EoS with $M_1/M_2 = 0.7$ [Tootle, Ecker, Topolski, Demircik, MJ, Rezzolla 2205.05691] [Movie by K. Topolski]

Details on quark formation



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Mechanical Toy Model

