## Towards the gauge beta function at $\mathcal{O}\left(1 / N_{f}^{2}\right)$ and $\mathcal{O}\left(1 / N_{f}^{3}\right)$

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Nicola Dondi, Gerald Dunne, MR, Francesco Saninno: arXiv:1903.02568
Nicola Dondi, MR, Francesco Saninno: in preparation


Cosmology \& Particle Physics

## Which matter systems are asymptotically safe in $d=4 ?$

- Gauge-Yukawa theories at large $N_{f} \& N_{c}$ (perturbatively) [Lutim, Sannino '14]
- How far does this extend to small $N_{c}$ ?
- Test gauge theories at large $N_{f}$ non-perturbatively

Standard QCD picture:

- Small $N_{f}$ : asymptotic freedom \& confinement in the IR
- Medium $N_{f}$ : asymptotic freedom \& IR Banks-Zaks fixed point
- Large $N_{f}$ : asymptotic freedom lost

[Antipin, Sannino '17]


## Beta functions of (S)QED and (S)QCD

$$
\beta(K)=\beta^{(0)}(K)+\frac{\beta^{(1)}(K)}{N_{f}}+\ldots
$$



UV fixed point for QED \& QCD
Landau pole for SQED \& SQCD

## How physical are these fixed points?

- The fermion mass anomalous dimension goes to zero in QCD and to infinity in QED
- Hints for FP in QCD at medium $N_{f}$ from resummations with Meijer G-functions
[Antipin, Maiezza, Vasquez '18]
- Lattice studies inconclusive so far
[Leino, Rindlisbacher, Rummukainen, Sannino, Tuominen '19]
- Poles might be resummable within the $1 / N_{f}$ series


## How to go beyond $1 / N_{f}$

- The next orders in the $1 / N_{f}$ expansion would test the physical nature of the FP
- No known resummation formula for two bubble-chains, needed for $1 / N_{f}^{2}$ and higher orders
- Can we extract the location of the pole, the residuum, etc., with a finite amount of coefficients?


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Two methods:

- Large-order behaviour of expansion coefficients
- Padé approximants


## Large-order behaviour: Darboux's Theorem

The nearby singularity determines the large-order growth of the expansion coefficients $a_{n}$. E.g. for expansion around $z=0$

- pole of order $p$ at $z_{0}\left(f(z) \sim \phi(z)\left(1-z / z_{0}\right)^{p}+\right.$ finite $)$

$$
a_{n} \sim \frac{1}{z_{0}^{n}}\binom{n+p-1}{n} \phi\left(z_{0}\right)+\ldots
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- logarithmic branch cut at $z_{0}\left(f(z) \sim \phi(z) \ln \left(1-z / z_{0}\right)+\right.$ finite $)$

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Expectation for QED $F_{Q E D}=\sum_{n} f_{n} x^{n}$

$$
f_{n} \sim\left[R_{0}\left(\frac{2}{15}\right)^{n}+R_{1}\left(\frac{2}{21}\right)^{n}+R_{2}\left(\frac{2}{27}\right)^{n}+\ldots\right]
$$

## Large-order behaviour of $F_{\text {QED }}$



Ratio test $\frac{f_{n+1}}{f_{n}}$ reveals location of the first pole

## Large-order behaviour of $F_{\text {QED }}$



## Large-order behaviour of $F_{\text {QED }}$



With the knowledge of the pole the residuum is computable

## Large-order behaviour of $F_{\text {QED }}$



Subtracting the first pole reveals the second pole

$$
\tilde{f}_{n}=f_{n}+\frac{28}{45 \pi^{2}}\left(\frac{15}{2}\right)^{-n-1}
$$

## Large-order behaviour of $F_{\text {QED }}$



After $\sim 30$ terms the large-order behaviour sets in (for subleading behaviour later)

## How many coefficients are needed?


"Closer" to the origin $\rightarrow$ less coefficients are needed

## Padé methods

Analytic continuation of truncated Taylor series by ration of two polynomials

$$
F_{\mathrm{QED}}(x) \approx \sum_{n=0}^{M} f_{n} x^{n} \quad \longrightarrow \quad \mathcal{P}^{[R, S]}(x)=\frac{P_{R}(x)}{Q_{S}(x)}
$$

with $R+S=M$.

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Rewriting of resummed $F_{Q E D}(x)$

$$
F_{\mathrm{QED}}(x) \sim \frac{\Gamma\left(1+\frac{x}{3}\right)}{\Gamma\left(\frac{1}{2}+\frac{x}{3}\right)} \frac{\sin ^{2}\left(\frac{\pi x}{3}\right)}{\cos \left(\frac{\pi x}{3}\right)}
$$

Padé approximant with $2 R \approx S$ should lead to best results.

## Padé approximants of $F_{\mathrm{Q} E \mathrm{D}}$



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## Padé approximants of $F_{\mathrm{Q} E \mathrm{D}}$



## Padé approximants of $F_{Q E D}$



- Need $\sim 30$ coefficients to resolve first singularity (similar to large-order growth analysis)
- Can resolve function beyond the first singularity


## QED beta function

$1 / N_{f}$ :


## QED beta function

$1 / N_{f}$ :

$1 / N_{f}^{2}$ (subset):


## QED beta function

$1 / N_{f}$ :

$1 / N_{f}^{2}$ (subset):


Master integral known / not know

## Beyond $1 / N_{f}$ : nested diagrams

Nested sub-part of beta function: gauge \& RG scale independent


Computation up to $K^{44}$

## Beyond $1 / N_{f}$ : nested diagrams

Nested sub-part of beta function: gauge \& RG scale independent


Computation up to $K^{44}$
At $\mathcal{O}\left(1 / N_{f}^{3}\right)$


Computation up to $K^{32}$

## Ratio test at $\mathcal{O}\left(1 / N_{f}^{2}\right)$

$$
\beta_{\text {nested }}^{(2)}=\sum_{n} b_{n} K^{n}
$$



New finite radius of convergence

## Ratio test at $\mathcal{O}\left(1 / N_{f}^{2}\right)$

$$
\beta_{\text {nested }}^{(2)}=\sum_{n} b_{n} K^{n}
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New finite radius of convergence but extreme slow convergence

## Richardson extrapolation

Enhance the convergence of the series

$$
a_{n}=s+\frac{A}{n}+\frac{B}{n^{2}}+\frac{C}{n^{3}}+\ldots
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\mathrm{R}^{(1)} a_{n} \equiv s=(n+1) a_{n+1}-n a_{n}
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For oscillating series: Shanks transformation

## Ratio test at $\mathcal{O}\left(1 / N_{f}^{2}\right)$



Bare series: $K^{*}=3.14$

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Bare series: $K^{*}=3.14$
First Richardson: $K^{*}=3.003$
Second Richardson: $K^{*}=3.00008$

## Residue at $\mathcal{O}\left(1 / N_{f}^{2}\right)$



Bare series: $3^{n} n^{2} b_{n}=-0.512$
Second Richardson: $3^{n} n^{2} b_{n}=-0.500007$

## Subleading behaviour

$$
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Bare series: $K^{*}=3.215$
Second Richardson: $K^{*}=3.0003$

## Subleading behaviour

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Bare series: $3^{n} n^{3} \tilde{b}_{n}=-0.512$
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## Large-order behaviour

Large-order behaviour

$$
\begin{aligned}
b_{n} & \sim-\frac{1}{2} \frac{1}{3^{n}}\left(\frac{1}{n^{2}}+\frac{1}{n^{3}}+\ldots\right)+\mathcal{O}\left(\frac{1}{(x>3)^{n}}\right) \\
& =-\frac{1}{2} \frac{1}{3^{n}} \frac{1}{n(n-1)}
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Resummation

$$
\sum_{n=4}^{\infty} b_{n} K^{n} \sim \frac{1}{6}(K-3) \ln \left(1-\frac{K}{3}\right)+\text { finite }
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Logarithmic branch cut but no pole!

## Nested beta function at $1 / N_{f}^{2}$


"Exact" nested beta function up to $K=3$
Beta function ambiguous beyond $K=3$ or magic cancellation needed

## Nested beta function at $1 / N_{f}^{2}$ beyond the first branch cut



No singularity before $K=15 / 2$
Positive pole at $K=15 / 2$ ?

## Nested beta function at $1 / N_{f}^{3}$



No singularity before $K=3$
Branch cut at $K=3$ ?

## Outlook 1: Conformal Padé

Can we do more with the coefficients that we have?

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Can we do more with the coefficients that we have?
Conformal map:

$$
K=\frac{6 z}{1+z^{2}} \quad \longleftrightarrow \quad z=\frac{K / 3}{1+\sqrt{1-K^{2} / 9}}
$$




## Outlook 1: Conformal Padé



Improvement over standard Padé
Requires knowledge on the location of the branch cut

## Outlook 2: Renormalons



Two factorially divergent contributions but the sum goes to zero
Are we picking up renormalon contributions?

## Outlook 2: Renormalons

Borel transform of the finite part of $1 / N_{f}$


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Borel transform of the finite part of $1 / N_{f}$


Do we pick up the renormalon at $t=3$ ?
Why not the renormalon at $t=6$ ?

## Outlook 2: Renormalons



Large-order behaviour

$$
a_{n} \sim \frac{n!}{n^{3} 3^{n}}\left(-3-9 \frac{1}{\ln (n)^{3}}\right)+\ldots
$$

## Summary and outlook

- Large-order behaviour \& Padé methods constitute powerful tools
- First partial result beyond $\mathcal{O}\left(1 / N_{f}\right)$ for QED:

New logarithmic branch cut at $K^{*}=3$ without pole

- Ideas: Conformal Padé \& tracking renormalons
- Future: Remaining diagrams (Master integrals?) \& QCD


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- Large-order behaviour \& Padé methods constitute powerful tools
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## Thank you for your attention

