# Towards the gauge beta function at $\mathcal{O}(1/N_f^2)$ and $\mathcal{O}(1/N_f^3)$

#### Manuel Reichert

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CP3-Origins, SDU Odense, Denmark

Nicola Dondi, Gerald Dunne, MR, Francesco Saninno: arXiv:1903.02568 Nicola Dondi, MR, Francesco Saninno: in preparation



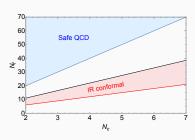


#### Which matter systems are asymptotically safe in d = 4?

- ullet Gauge-Yukawa theories at large  $N_f$  &  $N_c$  (perturbatively) [Litim, Sannino '14]
- How far does this extend to small  $N_c$ ?
- $\bullet$  Test gauge theories at large  $N_f$  non-perturbatively

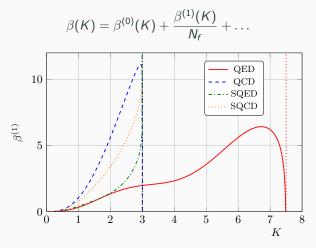
#### Standard QCD picture:

- Small N<sub>f</sub>: asymptotic freedom & confinement in the IR
- Medium N<sub>f</sub>: asymptotic freedom
   & IR Banks-Zaks fixed point



[Antipin, Sannino '17]

#### Beta functions of (S)QED and (S)QCD



UV fixed point for QED & QCD

Landau pole for SQED & SQCD

#### How physical are these fixed points?

- The fermion mass anomalous dimension goes to zero in QCD and to infinity in QED

  [Antipin, Sannino '17]
- Hints for FP in QCD at medium  $N_f$  from resummations with Meijer G-functions [Antioin, Majezza, Vasquez '18]
- Lattice studies inconclusive so far

[Leino, Rindlisbacher, Rummukainen, Sannino, Tuominen '19]

• Poles might be resummable within the  $1/N_f$  series

[Alanne, Blasi, Dondi '19]

## How to go beyond $1/N_f$

- ullet The next orders in the  $1/N_f$  expansion would test the physical nature of the FP
- No known resummation formula for two bubble-chains, needed for  $1/N_f^2$  and higher orders
- Can we extract the location of the pole, the residuum, etc., with a finite amount of coefficients?

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#### Two methods:

- Large-order behaviour of expansion coefficients
- Padé approximants

#### Large-order behaviour: Darboux's Theorem

The nearby singularity determines the large-order growth of the expansion coefficients  $a_n$ . E.g. for expansion around z = 0

• pole of order p at  $z_0$   $(f(z) \sim \phi(z)(1-z/z_0)^p + \text{finite})$ 

$$a_n \sim \frac{1}{z_0^n} \binom{n+p-1}{n} \phi(z_0) + \dots$$

• logarithmic branch cut at  $z_0$  ( $f(z) \sim \phi(z) \ln(1-z/z_0) + \text{finite}$ )

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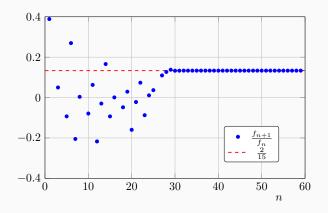
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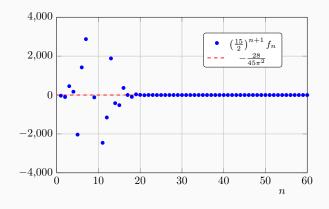
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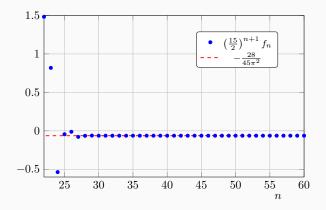
Expectation for QED  $F_{QED} = \sum_{n} f_n x^n$ 

$$f_n \sim \left[ R_0 \left( \frac{2}{15} \right)^n + R_1 \left( \frac{2}{21} \right)^n + R_2 \left( \frac{2}{27} \right)^n + \dots \right]$$

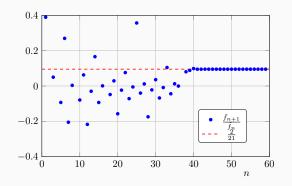


Ratio test  $\frac{f_{n+1}}{f_n}$  reveals location of the first pole



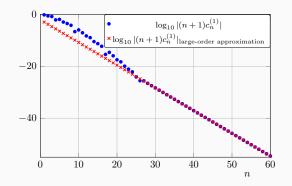


With the knowledge of the pole the residuum is computable



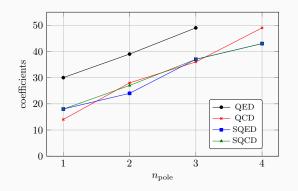
Subtracting the first pole reveals the second pole

$$\tilde{f}_n = f_n + \frac{28}{45\pi^2} \left(\frac{15}{2}\right)^{-n-1}$$



After  $\sim$  30 terms the large-order behaviour sets in (for subleading behaviour later)

#### How many coefficients are needed?



"Closer" to the origin  $\rightarrow$  less coefficients are needed

#### Padé methods

Analytic continuation of truncated Taylor series by ration of two polynomials

$$F_{\text{QED}}(x) \approx \sum_{n=0}^{M} f_n x^n \longrightarrow \mathcal{P}^{[R,S]}(x) = \frac{P_R(x)}{Q_S(x)}$$

with R + S = M.

#### Padé methods

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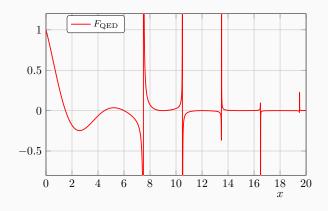
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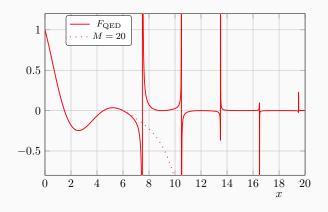
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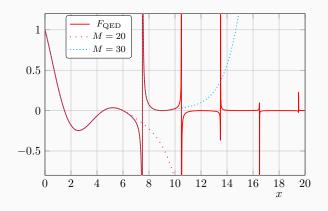
Rewriting of resummed  $F_{QED}(x)$ 

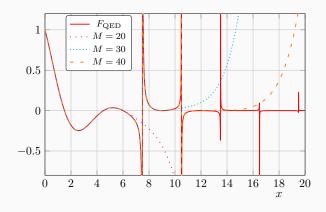
$$F_{\text{QED}}(x) \sim \frac{\Gamma(1+\frac{x}{3})}{\Gamma(\frac{1}{2}+\frac{x}{3})} \frac{\sin^2(\frac{\pi x}{3})}{\cos(\frac{\pi x}{3})}$$

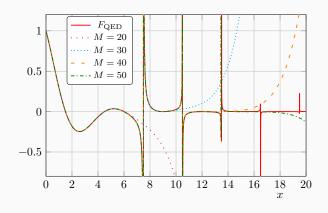
Padé approximant with  $2R \approx S$  should lead to best results.



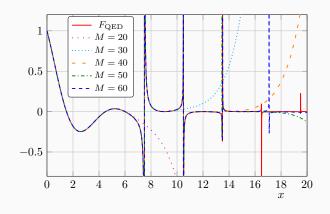


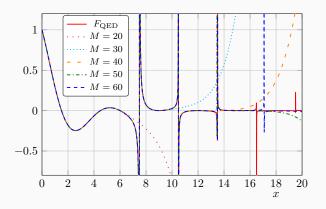






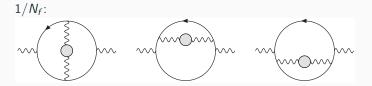
## Padé approximants of $\overline{F}_{QED}$



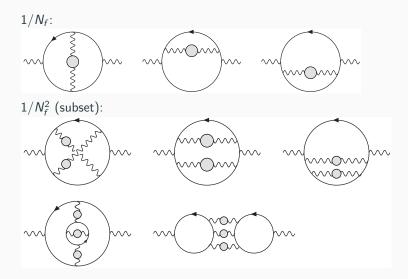


- Need  $\sim$  30 coefficients to resolve first singularity (similar to large-order growth analysis)
- Can resolve function beyond the first singularity

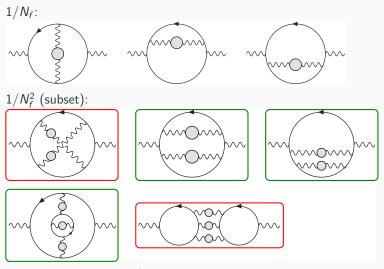
## **QED** beta function



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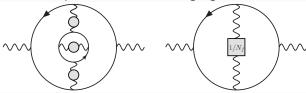
## **QED** beta function



Master integral known / not know

## Beyond $1/N_f$ : nested diagrams

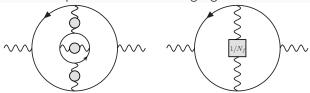
Nested sub-part of beta function: gauge & RG scale independent



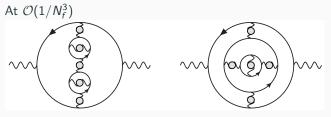
Computation up to  $K^{44}$ 

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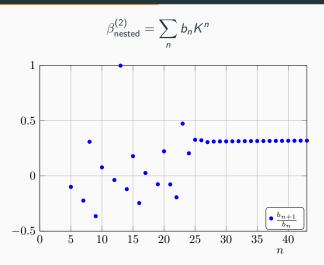


Computation up to  $K^{44}$ 



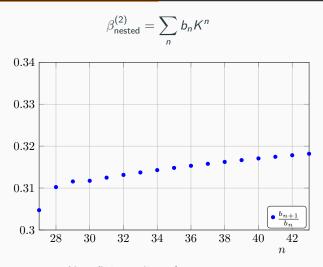
Computation up to  $K^{32}$ 

## Ratio test at $\mathcal{O}(1/N_f^2)$



New finite radius of convergence

## Ratio test at $\mathcal{O}(1/N_f^2)$



New finite radius of convergence but extreme slow convergence

Enhance the convergence of the series

$$a_n = s + \frac{A}{n} + \frac{B}{n^2} + \frac{C}{n^3} + \dots$$

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First Richardson (
$$B = C = ... = 0$$
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$$R^{(1)}a_n \equiv s = (n+1)a_{n+1} - na_n$$

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Second Richardson ( $C = \ldots = 0$ )

$$R^{(2)}a_n \equiv s = \frac{1}{2} \left( (n+2)^2 a_{n+2} - 2(n+1)^2 a_{n+1} + n^2 a_n \right)$$

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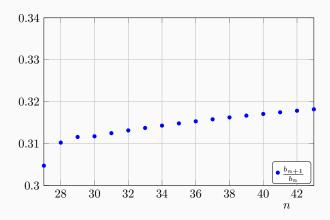
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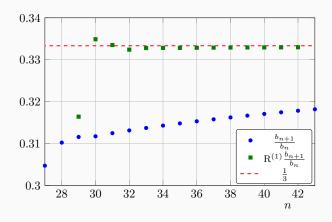
For oscillating series: Shanks transformation

## Ratio test at $\mathcal{O}(1/N_f^2)$



Bare series:  $K^* = 3.14$ 

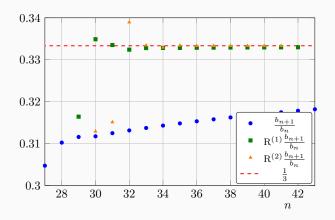
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Bare series:  $K^* = 3.14$ 

First Richardson:  $K^* = 3.003$ 

# Ratio test at $\mathcal{O}(1/N_f^2)$

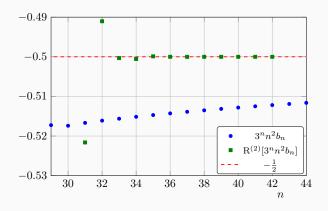


Bare series:  $K^* = 3.14$ 

First Richardson:  $K^* = 3.003$ 

Second Richardson:  $K^* = 3.00008$ 

# Residue at $\mathcal{O}(1/N_f^2)$



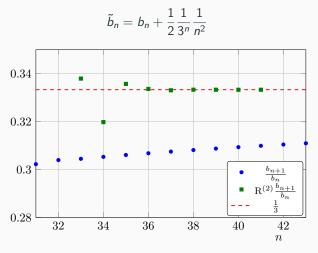
Bare series:  $3^n n^2 b_n = -0.512$ 

Second Richardson:  $3^n n^2 b_n = -0.500007$ 

## **Subleading behaviour**

$$\tilde{b}_n = b_n + \frac{1}{2} \frac{1}{3^n} \frac{1}{n^2}$$

# Subleading behaviour

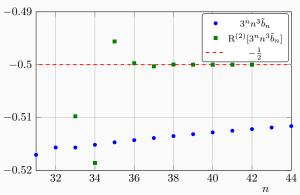


Bare series:  $K^* = 3.215$ 

Second Richardson:  $K^* = 3.0003$ 

## Subleading behaviour

$$\tilde{b}_n = b_n + \frac{1}{2} \frac{1}{3^n} \frac{1}{n^2}$$



Bare series:  $3^n n^3 \tilde{b}_n = -0.512$ 

Second Richardson:  $3^n n^3 \tilde{b}_n = -0.500007$ 

### Large-order behaviour

Large-order behaviour

$$b_n \sim -\frac{1}{2} \frac{1}{3^n} \left( \frac{1}{n^2} + \frac{1}{n^3} + \dots \right) + \mathcal{O}\left( \frac{1}{(x > 3)^n} \right)$$
$$= -\frac{1}{2} \frac{1}{3^n} \frac{1}{n(n-1)}$$

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Resummation

$$\sum_{n=4}^{\infty} b_n K^n \sim \frac{1}{6} (K-3) \ln \left(1 - \frac{K}{3}\right) + \text{finite}$$

### Large-order behaviour

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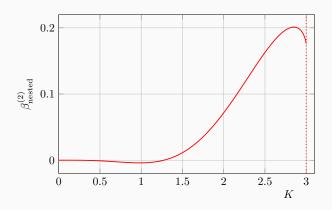
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Logarithmic branch cut but no pole!

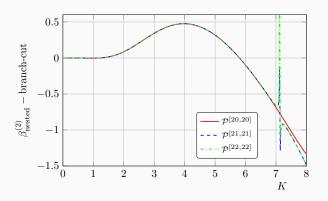
### Nested beta function at $1/N_f^2$



"Exact" nested beta function up to K=3

Beta function ambiguous beyond  $\mathcal{K}=3$  or magic cancellation needed

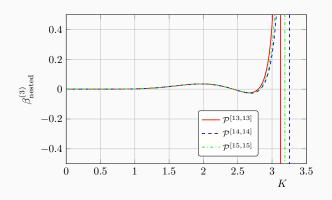
# Nested beta function at $1/N_f^2$ beyond the first branch cut



No singularity before K = 15/2

Positive pole at K = 15/2?

# Nested beta function at $1/N_f^3$



No singularity before K = 3

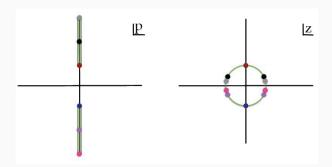
Branch cut at K = 3?

Can we do more with the coefficients that we have?

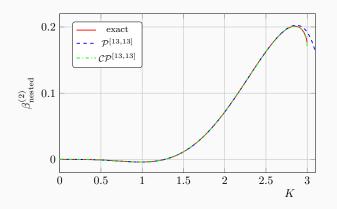
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#### Conformal map:

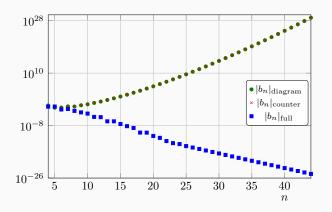
$$K = rac{6z}{1+z^2} \qquad \longleftrightarrow \qquad z = rac{K/3}{1+\sqrt{1-K^2/9}}$$



#### **Outlook 1: Conformal Padé**

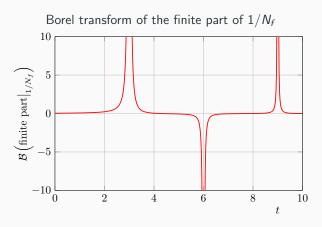


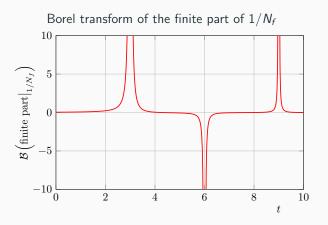
Improvement over standard Padé
Requires knowledge on the location of the branch cut



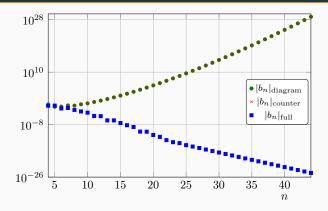
Two factorially divergent contributions but the sum goes to zero

Are we picking up renormalon contributions?





Do we pick up the renormalon at t=3? Why not the renormalon at t=6?



Large-order behaviour

$$a_n \sim \frac{n!}{n^3 \, 3^n} \left( -3 - 9 \frac{1}{\ln(n)^3} \right) + \dots$$

### **Summary and outlook**

- Large-order behaviour & Padé methods constitute powerful tools
- First partial result beyond \$\mathcal{O}(1/N\_f)\$ for QED:
   New logarithmic branch cut at \$K^\* = 3\$ without pole
- Ideas: Conformal Padé & tracking renormalons
- Future: Remaining diagrams (Master integrals?) & QCD

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# Thank you for your attention