

$B \to \gamma$ Form Factors @ NLO

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Motivation

Reasons to look at $B \to \gamma$ FF's...

- $B_s \to \ell^+ \ell^- \gamma$ as a test of Lepton Flavour Universality.
 - Sensitive to the same set of operators as $b \to s \ \ell^+ \ell^-$ due to identical effective Hamiltonian.
 - Shifts in the Wilson coefficient of these operators could explain anomalies in flavour data (e.g. $R_{K(*)}$).
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- Determination of the first inverse moment of the leading B-meson distribution amplitude, λ_B .
 - · Currently a large source of uncertainty in B-meson DA calculations.
 - · Lots of indirect determinations with results ranging from $\lambda_B=200$
 - 600 MeV ...Descotes-G. Sachrajda '02, Rohrwild, Beneke '11, Braun, Khodjamirian '12, Wang '16, Braun, Beneke, Ji, Wei '18 and more.
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 - Should be able to give competitive results by extending the work of Ball & Kou '03 to NLO.
- · An input for flavoured axion searches. Albrecht, Stamou, Ziegler, Zwicky '19.

Definition of the Form Factors

Form factors of interest parameterise the hadronic matrix elements of operators from the effective weak Hamiltonian.

$$\langle \gamma(k,\epsilon) | \mathcal{O}_{\mu}^{V} | \bar{B}_{q}(p_{B}) \rangle = P_{\mu}^{\perp} V_{\perp}(q^{2}) - P_{\mu}^{\parallel} V_{\parallel}(q^{2}) \qquad \mathcal{O}_{\mu}^{V} = -\frac{m_{B}}{e} \bar{q} \gamma_{\mu} (1 - \gamma_{5}) b,$$

$$\langle \gamma(k,\epsilon) | \mathcal{O}_{\mu}^{\mathsf{T}} | \bar{B}_{q}(p_{B}) \rangle = P_{\mu}^{\perp} T_{\perp}(q^{2}) - P_{\mu}^{\parallel} T_{\parallel}(q^{2}) \qquad \mathcal{O}_{\mu}^{\mathsf{T}} = \frac{1}{e} \bar{q} i q^{\nu} \sigma_{\mu\nu} (1 + \gamma_{5}) b$$

Two independent structures

$$P_{\mu}^{\perp} = \varepsilon_{\mu\alpha\beta\gamma} \, \epsilon^{*\alpha} p_{\mathsf{B}}^{\beta} k^{\gamma}, \qquad \qquad P_{\mu}^{\parallel} = \mathsf{i}(p_{\mathsf{B}} \cdot k \, \epsilon_{\mu}^{*} - p_{\mathsf{B}} \cdot \epsilon^{*} \, k_{\mu}).$$

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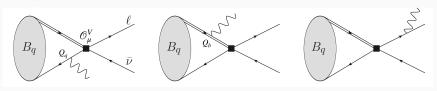
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$$P_{\mu}^{\perp} = \varepsilon_{\mu\alpha\beta\gamma} \, \epsilon^{*\alpha} p_{\rm B}^{\beta} k^{\gamma}, \qquad \qquad P_{\mu}^{\parallel} = i(p_{\rm B} \cdot k \, \epsilon_{\mu}^* - p_{\rm B} \cdot \epsilon^* \, k_{\mu}).$$





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- Light-Cone: OPE expansion of light-like separated operators $x^2 \sim 0$.
 - Expansion in twist = \dim spin \rightarrow parameterised by photon DA.
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 - Expansion in twist = \dim spin \rightarrow parameterised by photon DA.
 - Light-cone contributions dominant in the kinematic region $q^2 \ll m_b^2$.
- Photon DA: Photon is not exactly point-like (hard) but also has a (soft) hadronic contribution.
 - Hadronic contribution related to probability for photon to dissociate into partons.
 - Reminiscent of massless vector mesons → description in terms of DA.

Information on the FF is contained in the correlation function:

$$\begin{split} \Pi^{\Gamma}_{\mu}(p_B^2, q^2) &= i \int_{x} e^{-ip_B \cdot x} \left\langle \gamma(k, \epsilon) \right| T\{J_B(x) \mathcal{O}^{\Gamma}_{\mu}(0)\} \left| 0 \right\rangle, \\ &= P^{\perp}_{\mu} \Pi^{\Gamma}_{\perp}(p_B^2, q^2) - P^{\parallel}_{\mu} \Pi^{\Gamma}_{\parallel}(p_B^2, q^2) \end{split}$$

Interpolating operator for the *B*-meson $J_B=m_b\bar{b}\gamma_5q$, $\Gamma=T,V$.

For the perturbative photon contribution:

$$\langle \gamma(k,\epsilon)|
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• Compute the LCOPE expansion \rightarrow Up to twist-4 3-particle and $O(\alpha_s)$ for twist-1&2.

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- Employ quark-hardron duality.
- Apply Borel transform reduces contributions from higher resonances.

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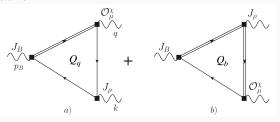
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Sum rule:

$$\Gamma_{\perp,\parallel}(q^2) = \frac{1}{m_B f_B} \int_{m_b^2}^{s_0} \mathrm{d}s \, \frac{1}{\pi} \mathrm{Im}[\Pi_{\perp,\parallel}^{\Gamma}(s,q^2)] e^{(m_B^2 - s)/M^2}.$$

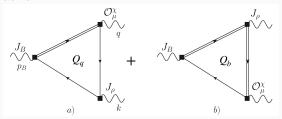
Leading contributions to the correlation function:

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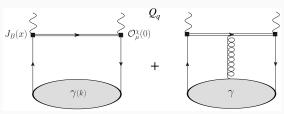


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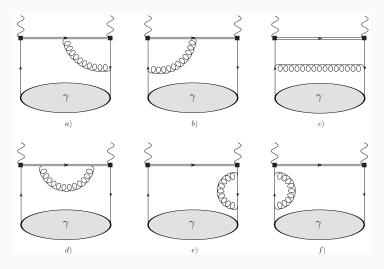


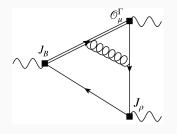
· Soft*



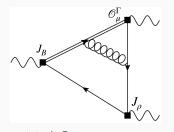
^{*}soft contribution proportional to Q_b is treated using local OPE

 $O(\alpha_s)$ corrections to twist-2 have been computed:





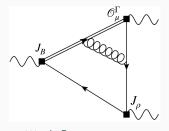
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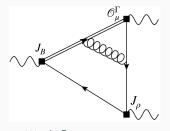
Work-flow:

• Compute traces (FeynCalc) \rightarrow Scalar integrals



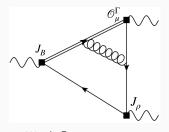
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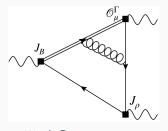
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- Master integrals computed using differential equation method
 - → Expression in terms of Goncharov functions .



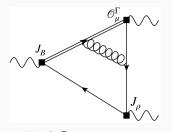
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- Numerical tests (SecDec, PolyLogTools)

Sanity Checks

1) Any FF determination has to obey the EOM:

$$ar{q}i(D^{
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Tensor "New" Vector

→ Provides non-trival check of the photon DA classification of Ball, Braun, Kivel '02.

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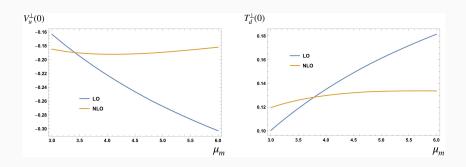
3) Renormalisation:

- → Non-trivial mixing of operators
- → Cancellation of thousands of Goncharov functions
- ightarrow Cancellation of explicit scale dependence at NLO

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Scale Dependence

Cancellation of explicit scale dependence at NLO



Similar results for the other FF

Breakdown and Borel stability

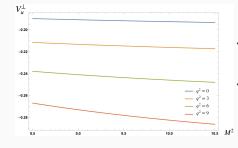
Breakdown			
	$V_u^{\perp}(0)$	$V_d^{\perp}(0)$	
Tw=1, LO	82%	85%	
Tw=1, NLO	-11%	-10 %	
Tw=2, LO	49%	41%	
Tw=2, NLO	-11%	-9%	
Tw=3, LO	-1%	<1%	
Tw=4, LO	-8%	-6%	

- Both charged and neutral FF has similar breakdown with reasonable convergence
- Appears that radiative corrections act to reduce FF but the relative sign depends on the scale/mass scheme.

Breakdown and Borel stability

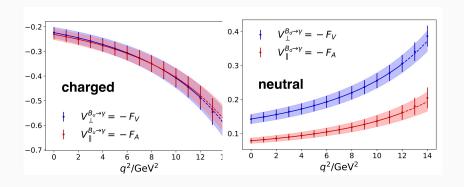
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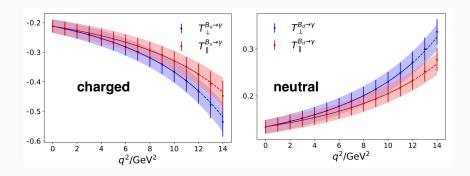
- Stability under variation of Borel parameter signals "good" sum rule.
- Variation increases with q^2 but stays under control (< 10%).

Preliminary Plots - Vector



 \cdot Q_q Dominates o difference between charged and neutral

Preliminary Plots - Tensor



• Recall
$$T_{\perp}(0) = T_{\parallel}(0)$$

Summary

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- $D_{uds} \rightarrow \gamma$ FF can also be extracted but only for a very small kinematic range.
- · NLO corrections significantly reduce the scale dependence.
- · Should be able to provide a competitive $1/\lambda_B$.
- Experimental data on $B_{\rm s} \to \ell\ell\gamma$ and $B \to \ell^+\nu\gamma$ anticipated.

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