$B \rightarrow \gamma$ Form Factors @ NLO
LHCb UK Annual Meeting Huddersfield '20

Tadeusz Janowski, Ben Pullin, Roman Zwicky

University of Edinburgh

## Motivation

Reasons to look at $B \rightarrow \gamma$ FF's...

- $B_{s} \rightarrow \ell^{+} \ell^{-} \gamma$ as a test of Lepton Flavour Universality.
- Sensitive to the same set of operators as $b \rightarrow s \ell^{+} \ell^{-}$due to identical effective Hamiltonian.
- Shifts in the Wilson coefficient of these operators could explain anomalies in flavour data (e.g. $\left.R_{K^{(*)}}\right)$.
- Guadagnioli, Reboud, Detori '16, Hazard, Petrov, '17, Kozachuk, Melikhov, Nikitin, '18, Guadagnioli, Reboud, Zwicky, '18


## Motivation

Reasons to look at $B \rightarrow \gamma$ FF's...

- $\mathrm{B}_{\mathrm{s}} \rightarrow \ell^{+} \ell^{-} \gamma$ as a test of Lepton Flavour Universality.
- Sensitive to the same set of operators as $b \rightarrow s \ell^{+} \ell^{-}$due to identical effective Hamiltonian.
- Shifts in the Wilson coefficient of these operators could explain anomalies in flavour data (e.g. $\left.R_{K^{(*)}}\right)$.
- Guadagnioli, Reboud, Detori '16, Hazard, Petrov, '17, Kozachuk, Melikhov, Nikitin, '18, Guadagnioli, Reboud, Zwicky, '18
- Determination of the first inverse moment of the leading $B$-meson distribution amplitude, $\lambda_{B}$.
- Currently a large source of uncertainty in B-meson DA calculations.
- Lots of indirect determinations with results ranging from $\lambda_{B}=200$
- 600 MeV ...Descotes-G. Sachrajda '02, Rohrwild,Beneke '11, Braun,Khodjamirian '12, Wang '16, Braun, Beneke, Ji, Wei '18 and more.
- Should be able to give competitive results by extending the work of Ball \& Kou '03 to NLO.


## Motivation

Reasons to look at $B \rightarrow \gamma$ FF's...

- $\mathrm{B}_{\mathrm{s}} \rightarrow \ell^{+} \ell^{-} \gamma$ as a test of Lepton Flavour Universality.
- Sensitive to the same set of operators as $b \rightarrow s \ell^{+} \ell^{-}$due to identical effective Hamiltonian.
- Shifts in the Wilson coefficient of these operators could explain anomalies in flavour data (e.g. $\left.R_{K^{(*)}}\right)$.
- Guadagnioli, Reboud, Detori '16, Hazard, Petrov, '17, Kozachuk, Melikhov, Nikitin, '18, Guadagnioli, Reboud, Zwicky, '18
- Determination of the first inverse moment of the leading $B$-meson distribution amplitude, $\lambda_{B}$.
- Currently a large source of uncertainty in B-meson DA calculations.
- Lots of indirect determinations with results ranging from $\lambda_{B}=200$
- 600 MeV ...Descotes-G. Sachrajda '02, Rohrwild,Beneke '11, Braun,Khodjamirian '12, Wang '16, Braun, Beneke, Ji, Wei '18 and more.
- Should be able to give competitive results by extending the work of Ball \& Kou '03 to NLO.
- An input for flavoured axion searches. Albrecht, Stamou, Ziegler, Zwicky '19.


## Definition of the Form Factors

Form factors of interest parameterise the hadronic matrix elements of operators from the effective weak Hamiltonian.

$$
\begin{array}{ll}
\langle\gamma(k, \epsilon)| \mathcal{O}_{\mu}^{\vee}\left|\bar{B}_{q}\left(p_{B}\right)\right\rangle=P_{\mu}^{\perp} V_{\perp}\left(q^{2}\right)-P_{\mu}^{\|} V_{\|}\left(q^{2}\right) & \mathcal{O}_{\mu}^{V}=-\frac{m_{B}}{e} \bar{q} \gamma_{\mu}\left(1-\gamma_{5}\right) b, \\
\langle\gamma(k, \epsilon)| \mathcal{O}_{\mu}^{\top}\left|\bar{B}_{q}\left(p_{B}\right)\right\rangle=P_{\mu}^{\perp} T_{\perp}\left(q^{2}\right)-P_{\mu}^{\|} T_{\|}\left(q^{2}\right) & \mathcal{O}_{\mu}^{\top}=\frac{1}{e} \bar{q} i q^{\nu} \sigma_{\mu \nu}\left(1+\gamma_{5}\right) b
\end{array}
$$

Two independent structures

$$
P_{\mu}^{\perp}=\varepsilon_{\mu \alpha \beta \gamma} \epsilon^{* \alpha} p_{B}^{\beta} k^{\gamma}, \quad P_{\mu}^{\|}=i\left(p_{B} \cdot k \epsilon_{\mu}^{*}-p_{B} \cdot \epsilon^{*} k_{\mu}\right) .
$$

## Definition of the Form Factors

Form factors of interest parameterise the hadronic matrix elements of operators from the effective weak Hamiltonian.
$\langle\gamma(k, \epsilon)| \mathcal{O}_{\mu}^{V}\left|\bar{B}_{q}\left(p_{B}\right)\right\rangle=P_{\mu}^{\perp} V_{\perp}\left(q^{2}\right)-P_{\mu}^{\|} V_{\|}\left(q^{2}\right)$

$$
\mathcal{O}_{\mu}^{\vee}=-\frac{m_{B}}{e} \bar{q} \gamma_{\mu}\left(1-\gamma_{5}\right) b,
$$

$\langle\gamma(k, \epsilon)| \mathcal{O}_{\mu}^{\top}\left|\bar{B}_{q}\left(p_{B}\right)\right\rangle=P_{\mu}^{\perp} T_{\perp}\left(q^{2}\right)-P_{\mu}^{\|} T_{\|}\left(q^{2}\right)$

$$
\mathcal{O}_{\mu}^{T}=\frac{1}{e} \bar{q} i q^{\nu} \sigma_{\mu \nu}\left(1+\gamma_{5}\right) b
$$

Two independent structures

$$
P_{\mu}^{\perp}=\varepsilon_{\mu \alpha \beta \gamma} \epsilon^{* \alpha} p_{B}^{\beta} k^{\gamma}, \quad P_{\mu}^{\|}=i\left(p_{B} \cdot k \epsilon_{\mu}^{*}-p_{B} \cdot \epsilon^{*} k_{\mu}\right) .
$$

$B \rightarrow \gamma \ell \bar{\nu}$


## Light-Cone Sum Rules

Form factors computed within the framework of Light-Cone Sum Rules (LCSR) using photon distribution amplitudes (DA).

## Light-Cone Sum Rules

Form factors computed within the framework of Light-Cone Sum Rules (LCSR) using photon distribution amplitudes (DA).

- Sum Rule: Relate an OPE expansion to a sum over a hadronic spectral density to form a link between QCD and hadrons.


## Light-Cone Sum Rules

Form factors computed within the framework of Light-Cone Sum Rules (LCSR) using photon distribution amplitudes (DA).

- Sum Rule: Relate an OPE expansion to a sum over a hadronic spectral density to form a link between QCD and hadrons.
- Light-Cone: OPE expansion of light-like separated operators $x^{2} \sim 0$.
- Expansion in twist = dim - spin $\rightarrow$ parameterised by photon DA.
- Light-cone contributions dominant in the kinematic region $q^{2} \ll m_{b}^{2}$.


## Light-Cone Sum Rules

Form factors computed within the framework of Light-Cone Sum Rules (LCSR) using photon distribution amplitudes (DA).

- Sum Rule: Relate an OPE expansion to a sum over a hadronic spectral density to form a link between QCD and hadrons.
- Light-Cone: OPE expansion of light-like separated operators $x^{2} \sim 0$.
- Expansion in twist = dim - spin $\rightarrow$ parameterised by photon DA.
- Light-cone contributions dominant in the kinematic region $q^{2} \ll m_{b}^{2}$.
- Photon DA: Photon is not exactly point-like (hard) but also has a (soft) hadronic contribution.
- Hadronic contribution related to probability for photon to dissociate into partons.
- Reminiscent of massless vector mesons $\rightarrow$ description in terms of DA.


## The Computation

Information on the FF is contained in the correlation function:

$$
\begin{aligned}
\Pi_{\mu}^{\Gamma}\left(p_{B}^{2}, q^{2}\right) & =i \int_{x} e^{-i p_{B} \cdot x}\langle\gamma(k, \epsilon)| T\left\{J_{B}(x) \mathcal{O}_{\mu}^{\Gamma}(0)\right\}|0\rangle, \\
& =P_{\mu}^{\perp} \Pi_{\perp}^{\ulcorner }\left(p_{B}^{2}, q^{2}\right)-P_{\mu}^{\|} \Pi_{\|}^{\Gamma}\left(p_{B}^{2}, q^{2}\right)
\end{aligned}
$$

Interpolating operator for the $B$-meson $J_{B}=m_{b} \bar{b} \gamma_{5} q, \Gamma=T, V$.
For the perturbative photon contribution:
$\langle\gamma(k, \epsilon)| \rightarrow-i e \int_{y} e^{i k \cdot y}\langle 0| \sum_{\psi=b, q} Q_{\psi} \bar{\psi} \gamma_{\rho} \psi(y)$

## The Computation

Information on the FF is contained in the correlation function:

$$
\begin{aligned}
\Pi_{\mu}^{\Gamma}\left(p_{B}^{2}, q^{2}\right) & =i \int_{x} e^{-i p_{B} \cdot x}\langle\gamma(k, \epsilon)| T\left\{J_{B}(x) \mathcal{O}_{\mu}^{\Gamma}(0)\right\}|0\rangle, \\
& =P_{\mu}^{\perp} \Pi_{\perp}^{\Gamma}\left(p_{B}^{2}, q^{2}\right)-P_{\mu}^{\|} \Pi_{\|}^{\Gamma}\left(p_{B}^{2}, q^{2}\right)
\end{aligned}
$$

Interpolating operator for the $B$-meson $J_{B}=m_{b} \bar{b} \gamma_{5} q, \Gamma=T, V$.
For the perturbative photon contribution:
$\langle\gamma(k, \epsilon)| \rightarrow-i e \int_{y} e^{i k \cdot y}\langle 0| \sum_{\psi=b, q} Q_{\psi} \bar{\psi} \gamma_{\rho} \psi(y)$
Step-by-step:

- Compute the LCOPE expansion $\rightarrow$ Up to twist-4 3-particle and $O\left(\alpha_{s}\right)$ for twist-1\&2.


## The Computation

Information on the FF is contained in the correlation function:

$$
\begin{aligned}
\Pi_{\mu}^{\Gamma}\left(p_{B}^{2}, q^{2}\right) & =i \int_{x} e^{-i p_{B} \cdot x}\langle\gamma(k, \epsilon)| T\left\{J_{B}(x) \mathcal{O}_{\mu}^{\Gamma}(0)\right\}|0\rangle, \\
& =P_{\mu}^{\perp} \Pi_{\perp}^{\ulcorner }\left(p_{B}^{2}, q^{2}\right)-P_{\mu}^{\|} \Pi_{\|}^{\Gamma}\left(p_{B}^{2}, q^{2}\right)
\end{aligned}
$$

Interpolating operator for the $B$-meson $J_{B}=m_{b} \bar{b} \gamma_{5} q, \Gamma=T, V$.
For the perturbative photon contribution:
$\langle\gamma(k, \epsilon)| \rightarrow-i e \int_{y} e^{i k \cdot y}\langle 0| \sum_{\psi=b, q} Q_{\psi} \bar{\psi} \gamma_{\rho} \psi(y)$
Step-by-step:

- Compute the LCOPE expansion $\rightarrow$ Up to twist-4 3-particle and O $\left(\alpha_{s}\right)$ for twist-1\&2.
- Write LCOPE expansion as a dispersion integral in $p_{B}$ \& match to the physical hadronic dispersion relation.


## The Computation

Information on the FF is contained in the correlation function:

$$
\begin{aligned}
\Pi_{\mu}^{\Gamma}\left(p_{B}^{2}, q^{2}\right) & =i \int_{x} e^{-i p_{B} \cdot x}\langle\gamma(k, \epsilon)| T\left\{J_{B}(x) \mathcal{O}_{\mu}^{\Gamma}(0)\right\}|0\rangle, \\
& =P_{\mu}^{\perp} \Pi_{\perp}^{\ulcorner }\left(p_{B}^{2}, q^{2}\right)-P_{\mu}^{\|} \Pi_{\|}^{\Gamma}\left(p_{B}^{2}, q^{2}\right)
\end{aligned}
$$

Interpolating operator for the $B$-meson $J_{B}=m_{b} \bar{b} \gamma_{5} q, \Gamma=T, V$.
For the perturbative photon contribution:
$\langle\gamma(k, \epsilon)| \rightarrow-i e \int_{y} e^{i k \cdot y}\langle 0| \sum_{\psi=b, q} Q_{\psi} \bar{\psi} \gamma_{\rho} \psi(y)$
Step-by-step:

- Compute the LCOPE expansion $\rightarrow$ Up to twist-4 3-particle and O $\left(\alpha_{\mathrm{s}}\right)$ for twist-1\&2.
- Write LCOPE expansion as a dispersion integral in $p_{B}$ \& match to the physical hadronic dispersion relation.
- Employ quark-hardron duality.


## The Computation

Information on the FF is contained in the correlation function:

$$
\begin{aligned}
\Pi_{\mu}^{\Gamma}\left(p_{B}^{2}, q^{2}\right) & =i \int_{x} e^{-i p_{B} \cdot x}\langle\gamma(k, \epsilon)| T\left\{J_{B}(x) \mathcal{O}_{\mu}^{\Gamma}(0)\right\}|0\rangle, \\
& =P_{\mu}^{\perp} \Pi_{\perp}^{\ulcorner }\left(p_{B}^{2}, q^{2}\right)-P_{\mu}^{\|} \Pi_{\|}^{\Gamma}\left(p_{B}^{2}, q^{2}\right)
\end{aligned}
$$

Interpolating operator for the $B$-meson $J_{B}=m_{b} \bar{b} \gamma_{5} q, \Gamma=T, V$.
For the perturbative photon contribution:
$\langle\gamma(k, \epsilon)| \rightarrow-i e \int_{y} e^{i k \cdot y}\langle 0| \sum_{\psi=b, q} Q_{\psi} \bar{\psi} \gamma_{\rho} \psi(y)$
Step-by-step:

- Compute the LCOPE expansion $\rightarrow$ Up to twist-4 3-particle and O $\left(\alpha_{\mathrm{s}}\right)$ for twist-1\&2.
- Write LCOPE expansion as a dispersion integral in $p_{B}$ \& match to the physical hadronic dispersion relation.
- Employ quark-hardron duality.
- Apply Borel transform - reduces contributions from higher resonances.


## The Computation

Information on the FF is contained in the correlation function:

$$
\begin{aligned}
\Pi_{\mu}^{\Gamma}\left(p_{B}^{2}, q^{2}\right) & =i \int_{x} e^{-i p_{B} \cdot x}\langle\gamma(k, \epsilon)| T\left\{J_{B}(x) \mathcal{O}_{\mu}^{\Gamma}(0)\right\}|0\rangle, \\
& =P_{\mu}^{\perp} \Pi_{\perp}^{\Gamma}\left(p_{B}^{2}, q^{2}\right)-P_{\mu}^{\|} \Pi_{\|}^{\Gamma}\left(p_{B}^{2}, q^{2}\right)
\end{aligned}
$$

Interpolating operator for the $B$-meson $J_{B}=m_{b} \bar{b} \gamma_{5} q, \Gamma=T, V$.
For the perturbative photon contribution:
$\langle\gamma(k, \epsilon)| \rightarrow-i e \int_{y} e^{i k \cdot y}\langle 0| \sum_{\psi=b, q} Q_{\psi} \bar{\psi} \gamma_{\rho} \psi(y)$

Sum rule:

$$
\Gamma_{\perp, \|}\left(q^{2}\right)=\frac{1}{m_{B} f_{B}} \int_{m_{b}^{2}}^{s_{0}} \mathrm{ds} \frac{1}{\pi} \operatorname{Im}\left[\Pi_{\perp, \|}^{\ulcorner }\left(s, q^{2}\right)\right] e^{\left(m_{B}^{2}-s\right) / M^{2}}
$$

## OPE Expansion @ LO

Leading contributions to the correlation function:

- Pertubative

b)


## OPE Expansion @ LO

Leading contributions to the correlation function:

- Pertubative

- Soft*

*soft contribution proportional to $Q_{b}$ is treated using local OPE


## OPE Expansion @ NLO

$O\left(\alpha_{s}\right)$ corrections to twist-2 have been computed:

a)

d)

b)

e)

c)

f)

## OPE Expansion @ NLO



- Main focus has been on calculating the $O\left(\alpha_{s}\right)$ corrections to the perturbative contribution.
- Achieved by using and extending the basis of master integrals for $\mathrm{H} \rightarrow$ WW Di Vita, Mastrolia, Primo, Schubert '17


## OPE Expansion @ NLO



- Main focus has been on calculating the $O\left(\alpha_{s}\right)$ corrections to the perturbative contribution.
- Achieved by using and extending the basis of master integrals for $\mathrm{H} \rightarrow$ WW Di Vita, Mastrolia, Primo, Schubert '17

Work-flow:

- Compute traces (FeynCalc) $\rightarrow$ Scalar integrals


## OPE Expansion @ NLO



- Main focus has been on calculating the $O\left(\alpha_{s}\right)$ corrections to the perturbative contribution.
- Achieved by using and extending the basis of master integrals for $\mathrm{H} \rightarrow$ WW Di Vita, Mastrolia, Primo, Schubert '17

Work-flow:

- Compute traces (FeynCalc) $\rightarrow$ Scalar integrals
- Reduction to master integrals (Kira, LiteRed)


## OPE Expansion @ NLO



- Main focus has been on calculating the $O\left(\alpha_{s}\right)$ corrections to the perturbative contribution.
- Achieved by using and extending the basis of master integrals for $H \rightarrow$ WW Di Vita, Mastrolia, Primo, Schubert '17

Work-flow:

- Compute traces (FeynCalc) $\rightarrow$ Scalar integrals
- Reduction to master integrals (Kira, LiteRed)
- Master integrals computed using differential equation method $\rightarrow$ Expression in terms of Goncharov functions.


## OPE Expansion @ NLO



- Main focus has been on calculating the $O\left(\alpha_{s}\right)$ corrections to the perturbative contribution.
- Achieved by using and extending the basis of master integrals for $H \rightarrow$ WW Di Vita, Mastrolia, Primo, Schubert '17

Work-flow:

- Compute traces (FeynCalc) $\rightarrow$ Scalar integrals
- Reduction to master integrals (Kira, LiteRed)
- Master integrals computed using differential equation method $\rightarrow$ Expression in terms of Goncharov functions .
- Transform to dispersive representation $\rightarrow$ take Im


## OPE Expansion @ NLO



- Main focus has been on calculating the $O\left(\alpha_{s}\right)$ corrections to the perturbative contribution.
- Achieved by using and extending the basis of master integrals for $H \rightarrow$ WW Di Vita, Mastrolia, Primo, Schubert '17

Work-flow:

- Compute traces (FeynCalc) $\rightarrow$ Scalar integrals
- Reduction to master integrals (Kira, LiteRed)
- Master integrals computed using differential equation method $\rightarrow$ Expression in terms of Goncharov functions .
- Transform to dispersive representation $\rightarrow$ take Im
- Goncharov functions $\rightarrow$ Classical polylogs (PolyLogTools)


## OPE Expansion @ NLO



- Main focus has been on calculating the $O\left(\alpha_{s}\right)$ corrections to the perturbative contribution.
- Achieved by using and extending the basis of master integrals for $H \rightarrow$ WW
Di Vita, Mastrolia, Primo, Schubert '17
Work-flow:
- Compute traces (FeynCalc) $\rightarrow$ Scalar integrals
- Reduction to master integrals (Kira, LiteRed)
- Master integrals computed using differential equation method $\rightarrow$ Expression in terms of Goncharov functions .
- Transform to dispersive representation $\rightarrow$ take Im
- Goncharov functions $\rightarrow$ Classical polylogs (PolyLogTools)
- Numerical tests (SecDec, PolyLogTools)


## Sanity Checks

1) Any FF determination has to obey the EOM:

$$
\begin{gathered}
\bar{q} i\left(D^{\nu} i \sigma_{\mu \nu}+\overleftrightarrow{B_{\mu}}\right)\left[\gamma_{5}\right] b+\left(m_{q} \mp m_{b}\right) \bar{q} \gamma_{\mu}\left[\gamma_{5}\right] b=0 \\
\text { Tensor "New" }
\end{gathered}
$$

$\rightarrow$ Provides non-trival check of the photon DA classification of Ball, Braun, Kivel '02.

## Sanity Checks

1) Any FF determination has to obey the EOM:

$$
\begin{gathered}
\bar{q} i\left(D^{\nu} i \sigma_{\mu \nu}+\overleftrightarrow{D_{\mu}}\right)\left[\gamma_{5}\right] b+\left(m_{q} \mp m_{b}\right) \bar{q} \gamma_{\mu}\left[\gamma_{5}\right] b=0 \\
\text { Tensor "New" }
\end{gathered}
$$

$\rightarrow$ Provides non-trival check of the photon DA classification of Ball, Braun, Kivel '02.
2) Algebraic Identity:

$$
T_{\perp}(0)=T_{\|}(0)
$$

$\rightarrow$ Checks consistency of $\gamma_{5}$ prescription

## Sanity Checks

1) Any FF determination has to obey the EOM:

$$
\begin{gathered}
\bar{q} i\left(D^{\nu} i \sigma_{\mu \nu}+\overleftrightarrow{D_{\mu}}\right)\left[\gamma_{5}\right] b+\left(m_{q} \mp m_{b}\right) \bar{q} \gamma_{\mu}\left[\gamma_{5}\right] b=0 \\
\text { Tensor "New" }
\end{gathered}
$$

$\rightarrow$ Provides non-trival check of the photon DA classification of Ball, Braun, Kivel '02.
2) Algebraic Identity:

$$
T_{\perp}(0)=T_{\|}(0)
$$

$\rightarrow$ Checks consistency of $\gamma_{5}$ prescription
3) Renormalisation:
$\rightarrow$ Non-trivial mixing of operators
$\rightarrow$ Cancellation of thousands of Goncharov functions
$\rightarrow$ Cancellation of explicit scale dependence at NLO

## Scale Dependence

## Cancellation of explicit scale dependence at NLO




Similar results for the other FF

## Breakdown and Borel stability

## Breakdown

|  | $V_{u}^{\perp}(0)$ | $V_{d}^{\perp}(0)$ |
| :--- | :---: | :---: |
| $T w=1, L O$ | $82 \%$ | $85 \%$ |
| $T w=1$, NLO | $-11 \%$ | $-10 \%$ |
| $T w=2, L O$ | $49 \%$ | $41 \%$ |
| $T w=2$, NLO | $-11 \%$ | $-9 \%$ |
| $T w=3, L O$ | $-1 \%$ | $<1 \%$ |
| $T w=4, L O$ | $-8 \%$ | $-6 \%$ |

- Both charged and neutral FF has similar breakdown with reasonable convergence
- Appears that radiative corrections act to reduce FF but the relative sign depends on the scale/mass scheme.


## Breakdown and Borel stability

| Breakdown |  |  |
| :--- | :---: | :---: |
|  | $V_{\frac{\perp}{u}}^{\perp}(0)$ | $V_{d}^{\perp}(0)$ |
| $T w=1, L O$ | $82 \%$ | $85 \%$ |
| $T w=1$, NLO | $-11 \%$ | $-10 \%$ |
| $T w=2$, LO | $49 \%$ | $41 \%$ |
| $T w=2$, NLO | $-11 \%$ | $-9 \%$ |
| $T w=3, L O$ | $-1 \%$ | $<1 \%$ |
| $T w=4, L O$ | $-8 \%$ | $-6 \%$ |

- Both charged and neutral FF has similar breakdown with reasonable convergence
- Appears that radiative corrections act to reduce FF but the relative sign depends on the scale/mass scheme.

- Stability under variation of Borel parameter signals "good" sum rule.
- Variation increases with $q^{2}$ but stays under control (<10\%).


## Preliminary Plots - Vector



- $Q_{q}$ Dominates $\rightarrow$ difference between charged and neutral


## Preliminary Plots - Tensor



## Summary

- Buds $\rightarrow \gamma$ FF have been computed @ NLO for low $q^{2} \ll m_{b}^{2}$. $\rightarrow$ Full numerics with z-expansion fit and error correlation are on the way.
- $D_{\text {uds }} \rightarrow \gamma$ FF can also be extracted but only for a very small kinematic range.
- NLO corrections significantly reduce the scale dependence.
- Should be able to provide a competitive $1 / \lambda_{B}$.
- Experimental data on $B_{s} \rightarrow \ell \ell \gamma$ and $B \rightarrow \ell^{+} \nu \gamma$ anticipated.


## Summary

- Buds $\rightarrow \gamma$ FF have been computed @ NLO for low $q^{2} \ll m_{b}^{2}$. $\rightarrow$ Full numerics with z-expansion fit and error correlation are on the way.
- $D_{\text {uds }} \rightarrow \gamma$ FF can also be extracted but only for a very small kinematic range.
- NLO corrections significantly reduce the scale dependence.
- Should be able to provide a competitive $1 / \lambda_{B}$.
- Experimental data on $B_{s} \rightarrow \ell \ell \gamma$ and $B \rightarrow \ell^{+} \nu \gamma$ anticipated.

Thanks for you attention!

