

O). Periods.

$$I = \int_{\gamma} \omega$$

X alg. variety / \mathbb{Q}
 ω closed algebraic
 n -form
 $\gamma \subseteq X(\mathbb{C})$ closed
 n -chain.

View: $[\omega] \in H_{dk}^n(X; \mathbb{Q})$

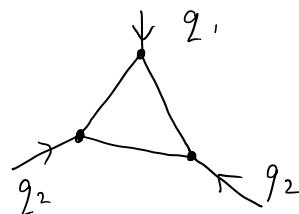
$[\gamma] \in H_n(X(\mathbb{C}); \mathbb{Q})$

I has a cohomological interpretation.

Variants: • ω, γ depend on parameters, $X \rightarrow S$

• γ has boundary $\partial\gamma \subseteq D$ X replaced with (X, D)
relative cohomology

Example: G any Feynman graph



$$I_G(q_i, m_i) = \text{pre factors} \cdot \int_0^\infty \int_0^\infty \frac{S}{(T+U)^p}$$

explicit polynomials depending on
parameters m_i , particle masses
 q_i particle momenta.

Theorem (Bloch-Esnault-Kreimer '07 ne masses, momenta,
B — 'IS , generic q_i, m_i)

I_G is a period of a canonical cohomology spaces :

$$\left(H_{\text{dk}}^n(X, \mathbb{D}), H^n(X(\mathbb{C}), \mathbb{D}(\mathbb{C})) \right)$$

↗ local resolution of singularities of graph
hypersurface complement, relative
to boundaries.

What do we gain ?

- Canonical ODE's (Picard-Fuchs equations)
- Weights
- Action of "Motivic" Galois group
- ⋮
- lots more

Coachian in brief.

a "mechanic period" is an equivalence class

$$I^m = [H, [\omega], [\gamma]]^m$$

Coachian is a rule to break it up into simpler pieces "atoms"

$$\Delta I^m = \sum_{i=1}^n a_i^m \otimes b_i^{\partial R}$$

RHS knows
"everything"!

respects

- all algebraic relations
- all ODE's
- Monodromy

Today: Integrals of the form

$$I(s_1, \dots, s_r) = \int_{\mathcal{X}} f_1^{s_1} \dots f_r^{s_r} \omega$$

viewed as a function
of s_i .

where $f_1, \dots, f_r : X \rightarrow \mathbb{G}_m$.

joint work with C. Dupont
in progress with C. Dupont, J. Fresán, M. Tapuskin

1967.06603, 1910.01107, 1810.07682

- Examples:
- Hypergeometric functions
 - String perturbation
 - Dim Reg.

Two different rebus of view

"GLOBAL" $s_1, \dots, s_r \in \mathbb{C}$ generic, fixed.

"LOCAL" Laurent expand in s_i at
non-generic point
 $s_1 = \dots = s_r = 0$

GLOBAL

$$I(s_1, \dots, s_r) = \int_{\gamma} w f_1^{s_1} \dots f_r^{s_r}$$

\$s_i\$ fixed \$\in \mathbb{C}\$
 generic

$I(s_1, \dots, s_r)$ is a period of cohomology with coefficients

$$H_{dk}^r(X, (\mathcal{V}, \nabla)) \quad (\mathcal{V}, \nabla) \quad \text{alg. v. bundle on } X$$

$$H_r(X(\mathbb{C}), \mathcal{L}_s^{\vee}) \quad \mathcal{L}_s^{\vee} \quad \text{local syst. of solns.}$$

FINITE RANK

FINITELY MANY PERIODS

LOCAL

$s_1, \dots, s_r \in \mathbb{C}$ formal parameters
at non-generic point

Taylor (Laurent) expansion

$$I(s_1, \dots, s_r) = \sum_{k_i \geq 0} \frac{s_1^{k_1} \dots s_r^{k_r}}{k_1! \dots k_r!} \int_{\gamma} \log^{k_1} f_1 \dots \log^{k_r} f_r w$$

classical periods of auxilliary varieties $X_{k_1, \dots, k_r} \subseteq X \times \mathbb{G}_m^{k_1 + \dots + k_r}$

INFINITELY MANY PERIODS

Toy example : Euler Beta function

$$\frac{s_0 s_1}{s_0 + s_1} \beta(s_0, s_1) = s_1 \int_0^1 x^{s_0} (1-x)^{s_1} \frac{dx}{1-x} = \frac{\Gamma(s_0) \Gamma(s_1)}{\Gamma(s_0 + s_1)}$$

GLOBAL point of view

$$s_0, s_1 \in \mathbb{C}$$

$$s_0, s_1, s_0 + s_1 \notin \mathbb{Z},$$

$$X = \mathbb{P}^1 \setminus \{0, 1, \infty\}$$

$$V = \mathcal{O}_X, \quad \nabla_s = d + s_0 \frac{dx}{x} + s_1 \frac{dx}{1-x}$$

$$\mathcal{L}_s^\vee \text{ rank 1 loc. sys over } \mathbb{Q}(e^{2\pi i s_j}), \quad x^{s_0} (1-x)^{s_1} \mathbb{Q}(e^{2\pi i s_j})$$

$$H_{\text{dR}}^1(X, (V, \nabla_s)) \cong \left[\frac{dx}{1-x} \right] \mathbb{Q}(s_0, s_1)$$

$$H_{\text{et}}^1(X(\mathbb{C}), \mathcal{L}_s^\vee) \cong (0, 1) \otimes x^{s_0} (1-x)^{s_1}$$

} rank 1

period matrix :

1x1 matrix

$$\left(\frac{s_0 s_1}{s_0 + s_1} \beta(s_0, s_1) \right)$$

LOCAL POINT OF VIEW : EXPAND AROUND BAD POINT

$$s_0 = s_1 = 0$$

$$s_1 \int_0^1 x (1-x)^{s_1} \frac{dx}{1-x} = \sum \frac{s_0^{k_0} s_1^{k_1+1}}{k_0! k_1!} \int_{\text{dch}} \log(x_0) \log(1-x_1)^{k_1} \frac{dx}{1-x}$$

↑
regularized integral .

How to interpret the expansion coefficients ? Write

$$\log x = \int_1^x \frac{dv}{v}$$

Coefficients of $s_0^{k_0}, s_1^{k_1}$ of the form:

$$\int \frac{du_1}{1-u_1} \dots \frac{du_{k_1}}{1-u_{k_1}} \frac{dx}{1-x} \frac{dv_1}{v_1} \dots \frac{dv_{k_0}}{v_{k_0}}$$

$\{0 < u_1 < \dots < u_{k_1} < x < v_1 < \dots < v_{k_0} < 1\}$

Classical periods on $M_{0,n}$

(iterated integrals on $X = \mathbb{P}^1 \setminus \{0, 1, \infty\}$)
Multiple zeta values $\zeta(l_1, \dots, l_m)$

Not obvious, but classical:
 $\text{period (GLOBAL)} = \text{period (LOCAL)}$

$$\frac{s_0 s_1}{s_0 + s_1} \beta(s_0, s_1) = \exp \left(\sum_{n \geq 2} \frac{(-1)^{n-1} J(n)}{n} \left((s_0 + s_1)^n - s_0^n - s_1^n \right) \right)$$

\uparrow

$\mu_{0,4}$ single period or
with coeffs.

∞ sequence of periods of $M_{0,n+3}$

Goal is to check that natural "relativistic" constructions
on local & global sides agree.

New structures

(i) Single-valued periods

(ii) Coactions

1) "single-valued periods"

Complex beta function :

$$\frac{s_0 s_1}{s_0 + s_1} \beta_C(s_0, s_1) = \frac{-s_1}{2\pi i} \int_{\mathbb{C}} |z|^{2s_0} |1-z|^{2s_1} \left(\frac{d\bar{z}}{\bar{z}-1} - \frac{d\bar{z}}{\bar{z}} \right) \wedge \frac{dz}{1-z}$$

Three ways to compute $\beta_C(s_0, s_1)$ from $\beta(s_0, s_1)$

①

"Double copy":

$$\frac{s_0 s_1}{s_0 + s_1} \beta_C(s_0, s_1) = - \frac{\beta(s_0, s_1)}{\beta(-s_0, -s_1)}$$

Applying twisted period relations:

②

(KLT formula) :

$$\beta_C(s_0, s_1) = -\frac{1}{2\pi i} \left(\sum_i \frac{2 \sin(\pi s_0) \sin(\pi s_1)}{\sin \pi(s_0 + s_1)} \right) \beta(s_0, s_1)^2$$

③

(Local-global)

$$\frac{s_0 s_1}{s_0 + s_1} \beta_C(s_0, s_1) = \exp \left(\sum_{\substack{n \geq 2 \\ n \text{ odd}}} \frac{2 \Im(n)}{n} \left((s_0 + s_1)^n - s_0^n - s_1^n \right) \right)$$

"throw away all even zeta values
& double odd zeta values"

2).

Coactions

General formalism forces:

$$\Delta \left(\frac{s_0 s_1}{s_0 + s_1} \beta^m(s_0, s_1) \right) = \frac{s_0 s_1}{s_0 + s_1} \beta^m(s_0, s_1) \otimes \frac{s_0 s_1}{s_0 + s_1} \beta^{dk}(s_0, s_1)$$

 Does this mean anything? Compare with local version.

$$\frac{s_0 s_1}{s_0 + s_1} \beta^m(s_0, s_1) = \exp \left(\sum_{n \geq 2} (-1)^{n-1} \frac{\zeta(n)}{n} \left((s_0 + s_1)^n - s_0^n - s_1^n \right) \right) \quad \bullet = m dk$$

 $\Delta \zeta^m(n) = \zeta^m(n) \otimes 1 + 1 \otimes \zeta^{dk}(n) \quad (\text{B. 2011})$

↑
non-trivial.

Check:  \Leftrightarrow 

(Proof: Recall

$$\Delta \alpha = \alpha \otimes 1 + 1 \otimes \alpha$$

$$\Delta \alpha^n = (\alpha \otimes 1 + 1 \otimes \alpha)^n$$

$$\Leftrightarrow \Delta e^\alpha = e^\alpha \otimes e^\alpha$$

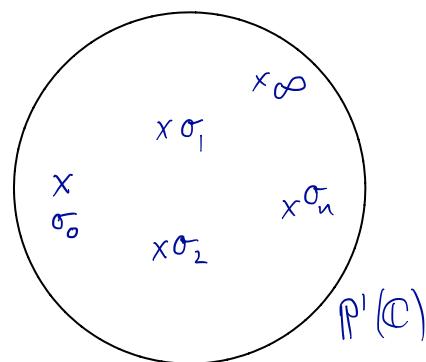
.)

Theorem (B-Dupont '19)

All the above results hold for Lamecella hypergeometric functions, in particular ${}_2F_1$.

$$\sum = \{\sigma_0, \dots, \sigma_n\}$$

$$\sigma_0 = 0$$



$$(\mathcal{L}_\sum)_{ij} = -s_j \int_0^{\sigma_i} x^{s_0} \prod_{k=1}^n \left(1-x^{-\sigma_k}\right)^{s_k} \frac{dx}{x-\sigma_j}$$

$$1 \leq i, j \leq n$$

Matrix of Lamecella functions.

Multivalued functions of $\sigma_0, \dots, \sigma_n$.

$$\text{Example : } \sum = \{0, 1\}, \quad \mathcal{L}_{\{0, 1\}} = \frac{s_0 s_1}{s_0 + s_1} \beta(s_0, s_1)$$

$$\sum = \{0, 1, \gamma y\}$$

$$\beta(b, c-b) {}_2F_1(a, b, c; y) = \int_0^1 x^{b-1} (1-x)^{c-b-1} (1-yx)^{-a} dx$$

(i) Single-valued Lamicella fns

$$\left(L_{\Sigma}^{\$}\right)_{ij} = \frac{s_j}{2\pi i} \int_{\mathbb{C}} |z|^{2s_0} \prod_{k=1}^n \left|1 - z \bar{\sigma}_k^{-1}\right|^{2s_k} \left(\frac{d\bar{z}}{\bar{z} - \bar{\sigma}_i} - \frac{d\bar{z}}{\bar{z} - \bar{\sigma}_j} \right)$$

$$s_i \in \mathbb{R}, \quad s_i > 0 \quad 2(s_0 + \dots + s_n) < 1$$

Double copy!

$$L_{\Sigma}^{\$} = \left(L_{\bar{\Sigma}}(-s_0, \dots, -s_n) \right)^{-1} L_{\Sigma}(s_0, \dots, s_n)$$

\uparrow
complex conjugate $\bar{\Sigma} = \{\bar{\sigma}_i\}$

(ii) Coactions

$$\Delta L_{\Sigma}^m = L_{\Sigma}^m \otimes L_{\Sigma}^{\partial R}$$

(cf formula conjectured by Abreu, Bittó, Duhr, Gardi)

Idea of proof:

- Tangential base point regularization & "renormalisation" of poles in twisted cohomology
- Prove that L_{Σ} is meta-abelian quotient of generalised Drinfeld associators.