

CKM metrology from Unitarity Triangle fits



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4th Conference on the Interplay between
Particle and Astroparticle physics (IPA2022)
6th September 2022
Technische Universität
Vienna, Austria





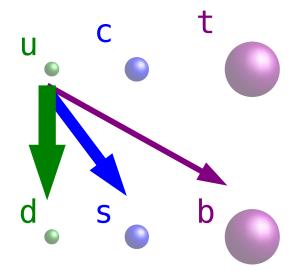


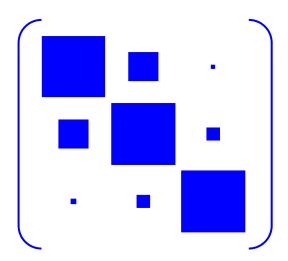


Flavour mixing and CP violation in the Standard Model

- The CP symmetry is violated in any field theory having in the Lagrangian at least one phase that cannot be re-absorbed
- The **mass eigenstates** are not eigenstates of the weak interaction. This feature of the Standard Model Hamiltonian produces the (unitary) **mixing matrix** V_{CKM} .

$$egin{pmatrix} egin{pmatrix} egi$$





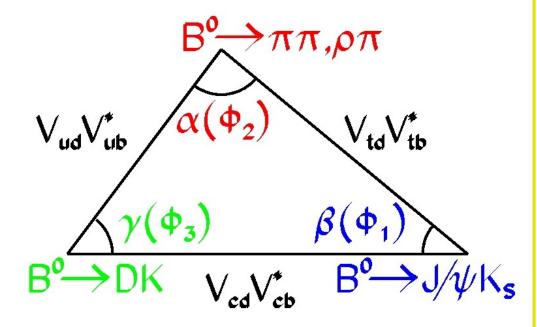
The CKM matrix and the Unitarity Triangle

$$egin{pmatrix} egin{pmatrix} egi$$

With three families of quarks, the unitary CKM matrix has four independent parameters: three rotation angles and one phase. This phase allows CP violation in the SM. All the flavour mixing processes are related (through the unitarity of the V_{CKM}) to this phase.

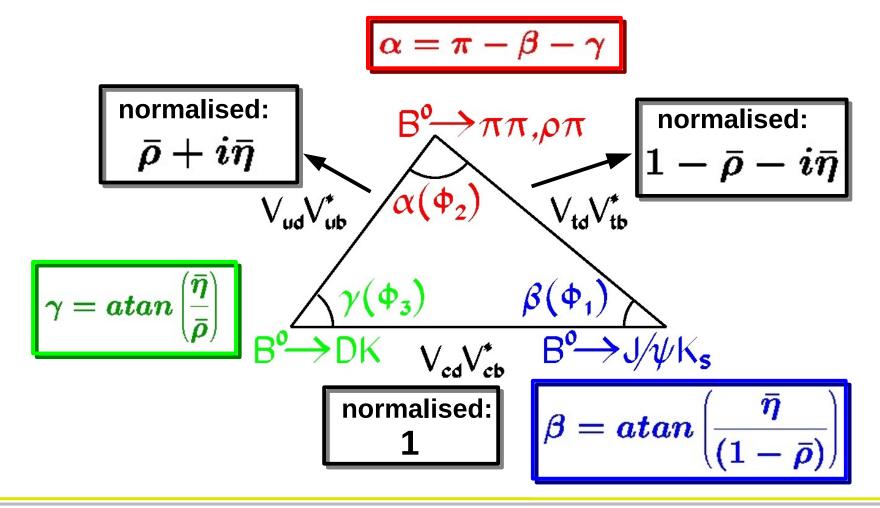
Unitarity Triangle
$$V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + V_{td}V_{tb}^* = 0$$

All the angles are related to the CP asymmetries of specific B decays



The CKM matrix and the Unitarity Triangle

- Description The Wolfenstein-Buras parameterisation of the CKM matrix allows to obtain a simplified (and approximate) form of the matrix, maintaining its unitarity. The four independent parameters being λ, A, ρ and η.
- The unitarity triangle can be drawn in the ρ-η plane and its sides and angles can be linked to various processes involving B mesons



Unitarity Triangle analysis in the Standard Model

- Standard Model (SM) Unitarity Triangle analysis:
 - All updated with Summer 2022 inputs
 - provide the best determination of CKM parameters
 - test the consistency of the SM ("direct" vs "indirect" determinations)
 - provide predictions (from data..) for SM observables

.. and beyond

- New Physics (NP) Unitarity Triangle analysis:
 - Also all updated with Summer 2022 inputs
 - model-independent analysis
 - provides limit on the allowed deviations from the SM
 - obtain the NP scale







M.Bona, M. Ciuchini, D. Derkach, F. Ferrari, E. Franco, V. Lubicz, G. Martinelli, M. Pierini, L. Silvestrini, S. Simula, A. Stocchi, C. Tarantino, V. Vagnoni, M. Valli and L. Vittorio

Plots and numbers in this talk are obtained with inputs updated this summer hence they are labelled "summer22".

Some changes have been included in July 22 for ICHEP22 and for this talk with respect to the results presented in May 2022 at

LHCP22 and FPCP22.

Statistical method and inputs:

$$f(ar
ho,ar\eta,X|c_1,...,c_m) \sim \prod_{j=1,m} f_j(\mathcal{C}|ar
ho,ar\eta,X) *$$
Bayes Theorem

Bayes Theorem

$$\prod_{i=1,N}^{J-1,m}\!\!f_i(x_i)f_0(ar
ho,ar\eta)$$

$$X\equiv x_1,...,x_n=m_t,B_K,F_B,...$$

$$\mathcal{C} \equiv c_1,...,c_m = \epsilon, \Delta m_d/\Delta m_s, A_{CP}(J/\psi K_S),...$$

 $(1-\bar{\rho})^2+\bar{\eta}^2$ Δm_d

 $(1-\bar{\rho})^2+\bar{\eta}^2$ $\Delta m_d/\Delta m_s$

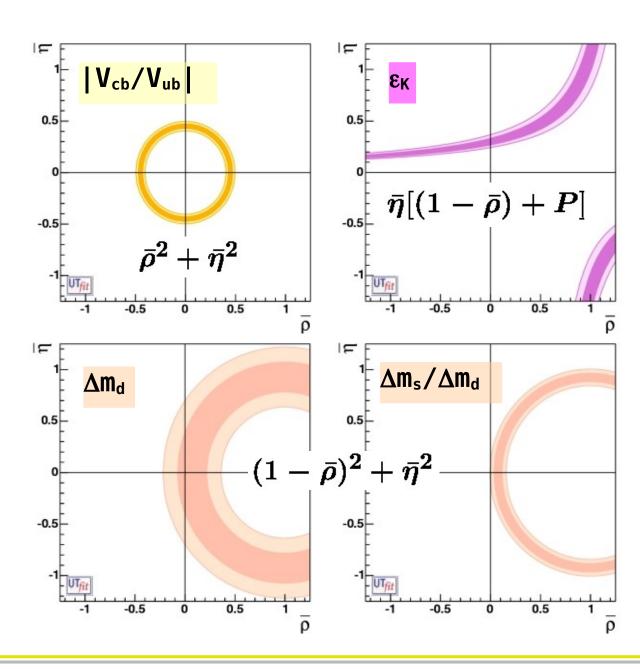
 $A_{CP}(J/\psi K_S)$ $\sin 2\beta$ $f_B^2 B_B$

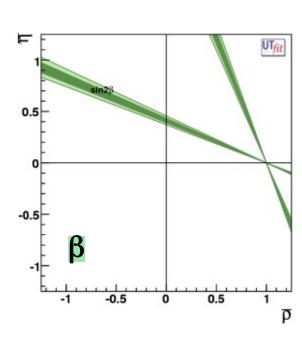
Standard Model + OPE/HQET/ Lattice QCD to go from quarks to hadrons

M. Bona et al. (UTfit Collaboration) JHEP 0507:028,2005 hep-ph/0501199 M. Bona et al. (UTfit Collaboration)

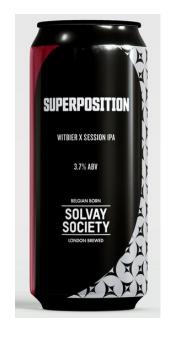
JHEP 0603:080,2006 hep-ph/0509219

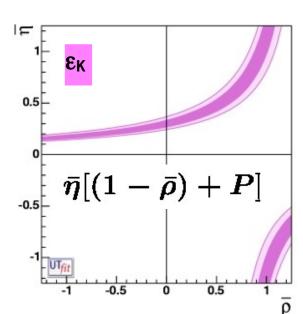
Inputs mapped on the ρ - η plane:

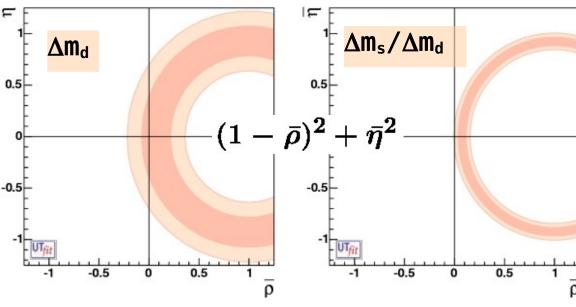




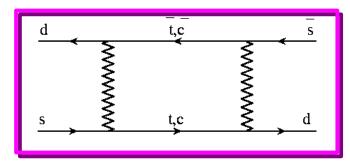
Meson mixing: K and B_{d/s}







ε_{K} from \overline{K} -K mixing

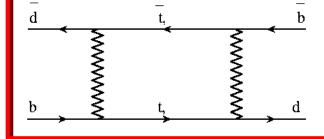


$$\varepsilon_{\rm K} = (2.228 \pm 0.011) \cdot 10^{-3}$$

PDG

∆m_q from B_q-B_q mixing

q=d.s



 $\Delta m_d = 0.5065 \pm 0.0019 \text{ ps}^{-1}$

 $\Delta m_S = 17.765 \pm 0.006 \text{ ps}^{-1}$

HFLAV

Lattice QCD inputs:

$$|\epsilon_{K}| = C_{\epsilon} \frac{B_{K} A^{2} \lambda^{6} \overline{\eta} \{-\eta (S_{0}(x_{c})(1-\lambda^{2}/2) + \eta_{3} S_{0}(x_{c},x_{t}) + \eta_{4} (S_{0}(x_{t})) A^{2} \lambda^{4} (1-\overline{\rho}) \}$$

S₀ = Inami-Lim functions for c-c, c-t, e t-t contributions (from perturbative calculations)

$$\Delta m_{d} = \frac{G_{F}^{2}}{6\pi^{2}} m_{W}^{2} \eta \left(S(x_{t}) \right) m_{B_{d}} f_{B_{d}}^{2} \mathring{B}_{B_{d}} |V_{tb}|^{2} |V_{td}|^{2} =$$

$$= \frac{G_{F}^{2}}{6\pi^{2}} m_{W}^{2} \eta_{b} S(x_{t}) m_{B_{d}} f_{B_{d}}^{2} \mathring{B}_{B_{d}} |V_{cb}|^{2} \lambda^{2} ((1-\overline{\rho})^{2} + \overline{\eta}^{2})$$

$$\Delta m_d \approx [(1-\rho)^2 + \eta^2] \frac{f_{B_s}^2 B_{B_s}}{\xi^2}$$

$$\Delta m_s \approx f_{B_s}^2 B_{B_s}$$

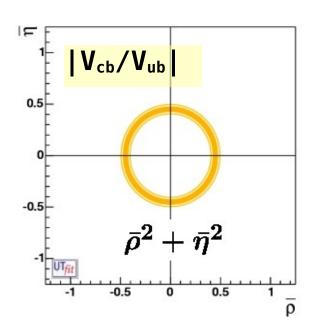
lattice inputs updated for this summer

Observables	Measurement
B _K	0.756 ± 0.016
f _{Bs}	0.2301 ± 0.0012
f _{Bs} /f _{Bd}	1.208 ± 0.005
B _{Bs} /B _{Bd}	1.015 ± 0.021
B _{Bs}	1.284 ± 0.059

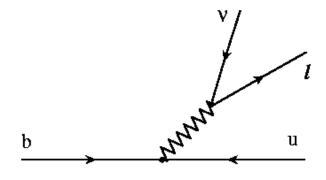
We quote the weighted average of the $N_f=2+1+1$ and $N_f=2+1$ results with the error rescaled when chi2/dof > 1, as done by FLAG for the $N_f=2+1+1$ and $N_f=2+1$ averages separately

[new HPQCD (2+1+1) result 1907.01025]

V_{cb} and V_{ub}



$$\left|\frac{V_{ub}}{V_{cb}}\right| = \frac{\lambda}{1 - \frac{\lambda^2}{2}} \sqrt{\bar{\rho}^2 + \bar{\eta}^2}$$



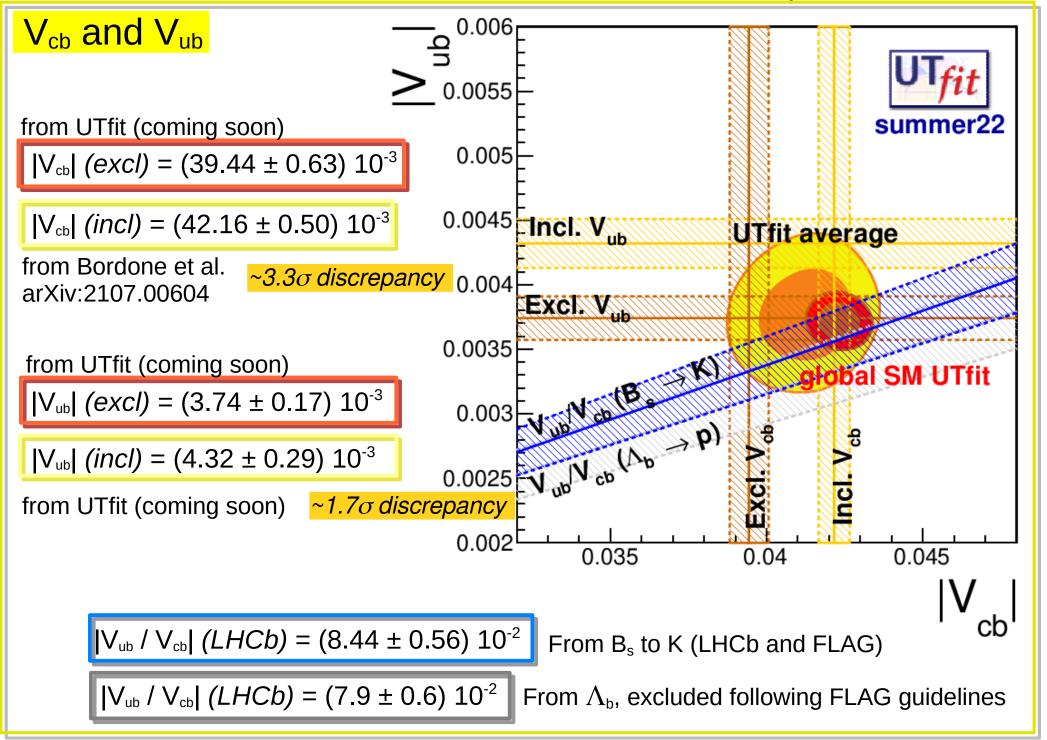
tree diagrams

 $b \rightarrow c$ and $b \rightarrow u$ transition

- negligible new physics contributions
- inclusive and exclusive semileptonic
 B decay branching ratios

QCD corrections to be included

- inclusive measurements: OPE
- exclusive measurements: form factors from lattice QCD





A-la-D'Agostini two-dimensional average procedure:

$$|V_{cb}| = (41.25 \pm 0.95) \, 10^{-3}$$

uncertainty ~ 2.3%

$$|V_{ub}| = (3.77 \pm 0.24) \ 10^{-3}$$

uncertainty ~ 6.4%

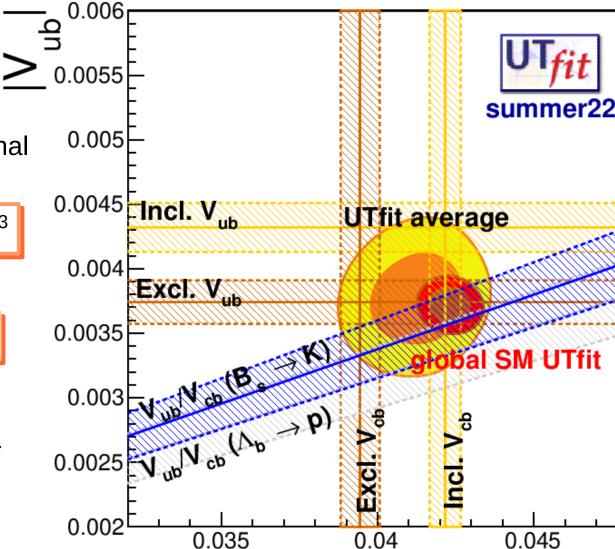
Correlation $\rho = 0.11$

Updated averages including correlation

From global SM fit

$$|V_{cb}| = (42.3 \pm 0.4) \cdot 10^{-3}$$

$$|V_{ub}| = (3.72 \pm 0.09) \, 10^{-3}$$



UTfit prediction:

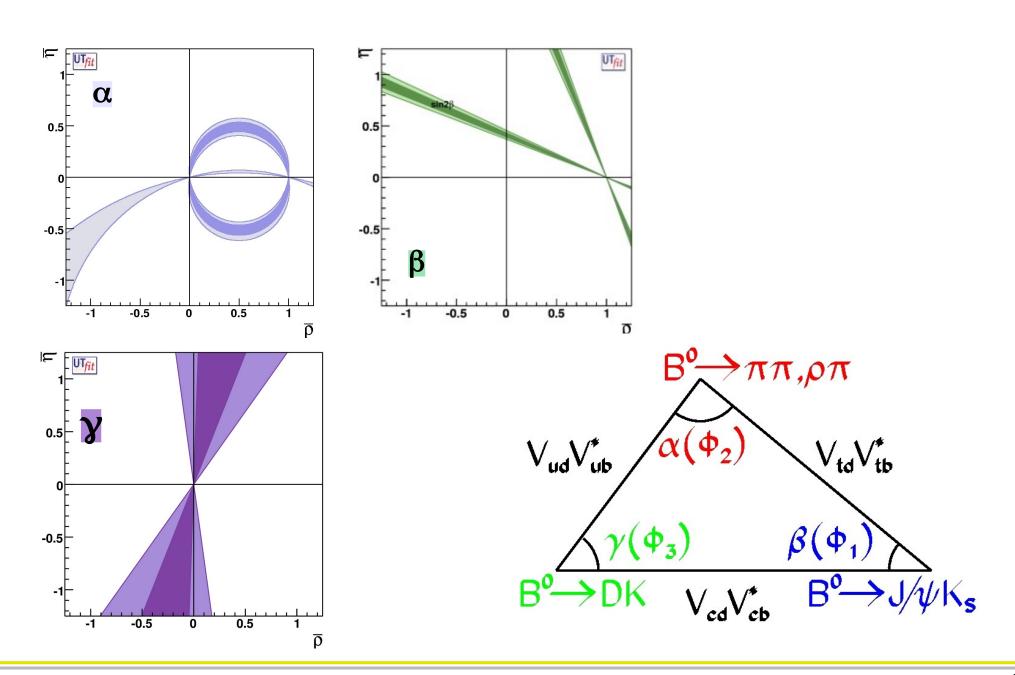
$$|V_{cb}| = (42.6 \pm 0.5) \, 10^{-3}$$

0.04

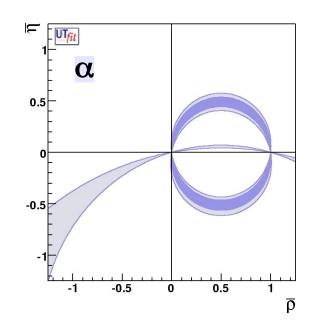
$$|V_{ub}| = (3.70 \pm 0.10) \cdot 10^{-3}$$

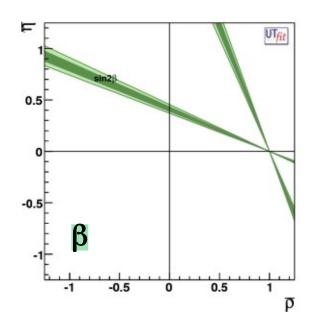
0.045

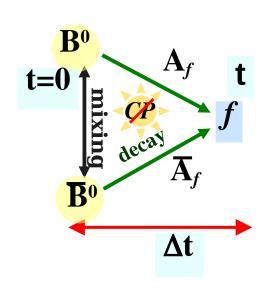
Angle constraints in the $\overline{\rho}$ - $\overline{\eta}$ plane:



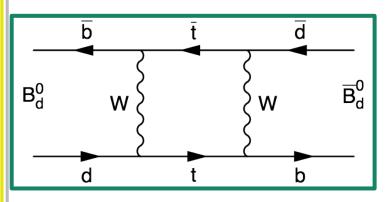
Time-dependent CP asymmetry:

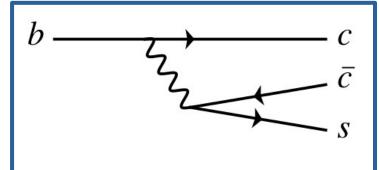


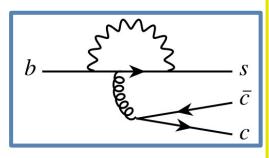




CP Violation in interference between mixing and decay:

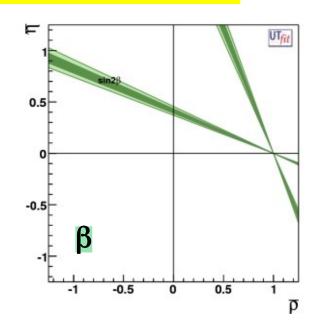


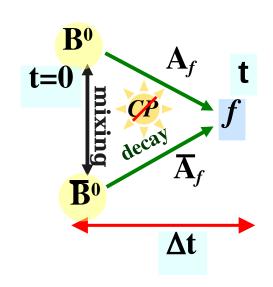




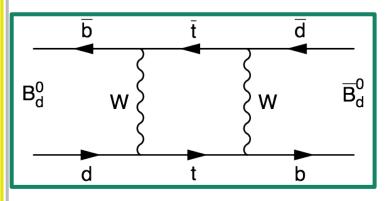
Time-dependent CP asymmetry:

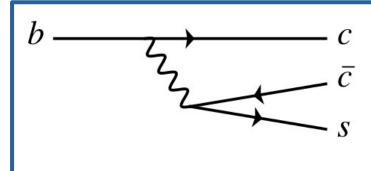
sin2β from time-dependent A_{CP} in $B \rightarrow J/\psi K^0$

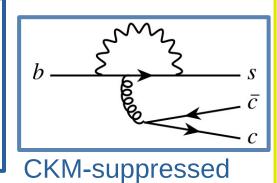




CP Violation in interference between mixing and decay:







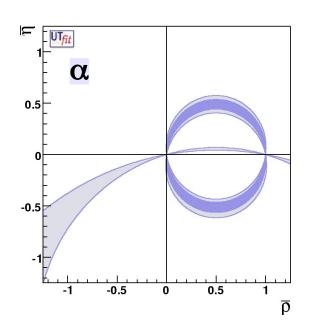
pollution by penguins

 $\Delta S = -0.01 \pm 0.01$

$$a_{f_{CP}}(t) = -\eta_{CP} \sin \Delta m_d \Delta t \sin 2\beta$$

Time-dependent CP asymmetry:

α : CP violation in B⁰ $\rightarrow \pi^+\pi^-$



considering the tree (T) only:

$$C_{\pi\pi} = 0$$

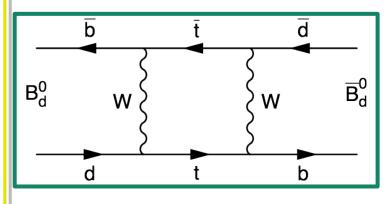
 $S_{\pi\pi} = \sin(2\alpha)$

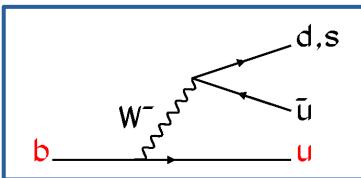
adding the penguins (P):

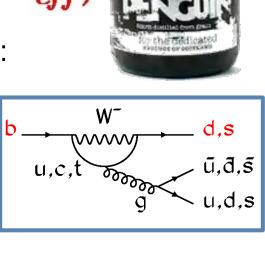
$$C_{\pi\pi} \propto \sin(\delta)$$

$$S_{\pi\pi} = \sqrt{1-C_{\pi\pi}^2\sin(2lpha_{e\!f\!f})}$$





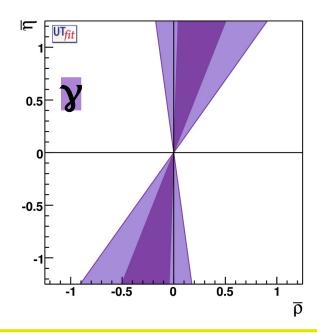


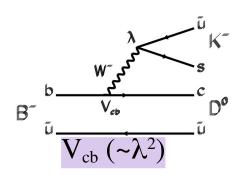


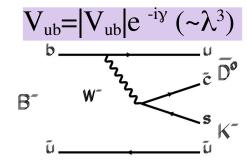
Direct CP asymmetry:

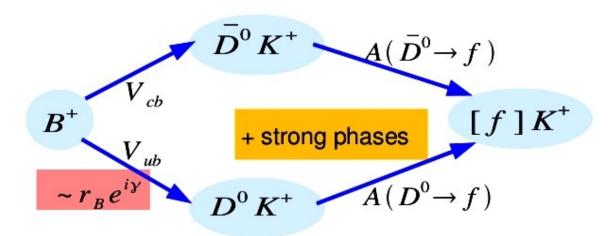
γ and DK trees

- D^(*)K^(*) decays: from BRs and BR ratios, no time-dependent analysis, just rates
- the phase γ is measured exploiting interferences: two amplitudes leading to the same final states
- some rates can be really small: ~ 10⁻⁷





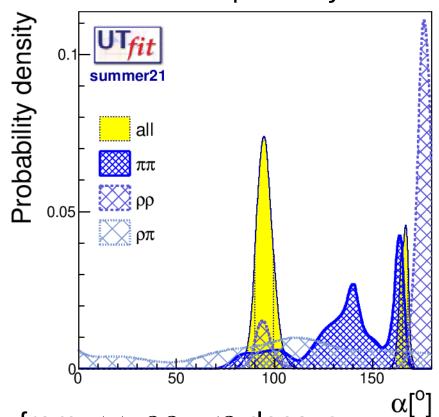




 $B \Rightarrow D^{(*)0}(D^{(*)0})K^{(*)}$ decays can proceed both through V_{cb} and V_{ub} amplitudes

$\sin 2\alpha \ (\phi_2) \ \text{and} \ \gamma \ (\phi_3)$

 α with $\pi\pi/\rho\rho$ BR and C/S results and $\rho\pi$ analysis

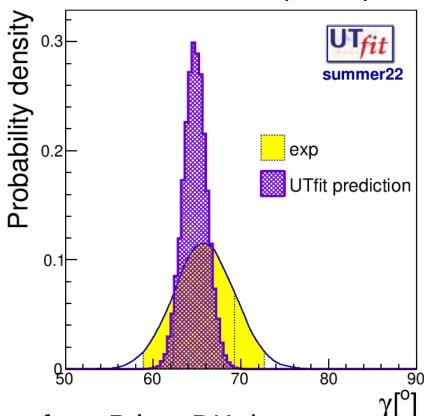


α from $\pi\pi$, $\rho\rho$, $\pi\rho$ decays: combined SM: (95.0 ± 4.7)°

UTfit prediction: (92.3 ± 1.5)°

 α from HFLAV: 85.5 ± 4.6

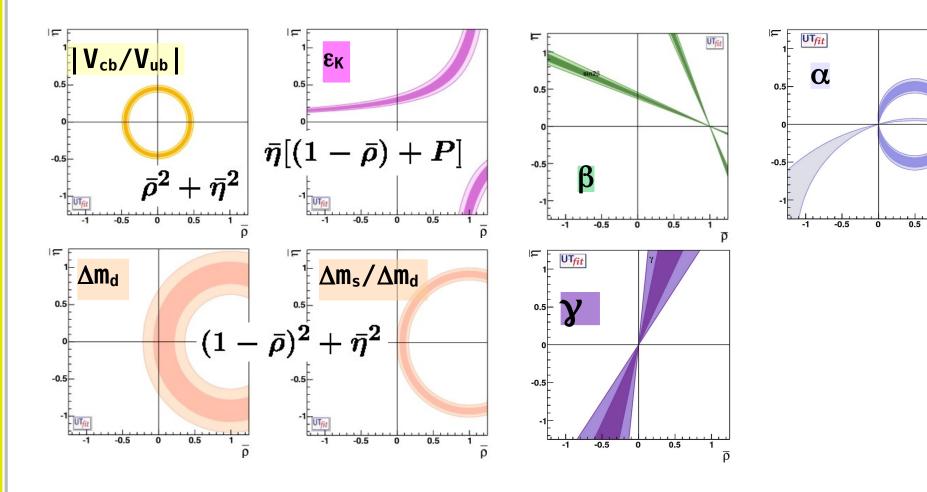
y updated with all the latest results (LHCb)

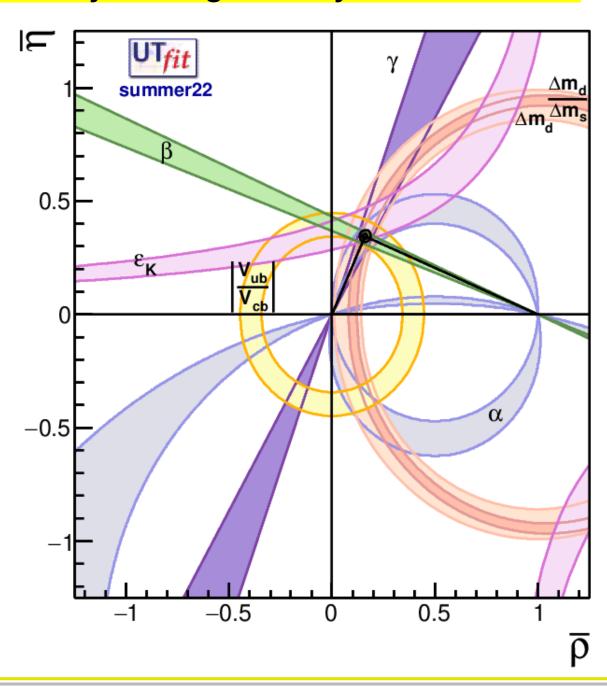


 γ from B into DK decays:

HFLAV: $(65.8 \pm 3.4)^{\circ}$

UTfit prediction: $(64.9 \pm 1.3)^{\circ}$





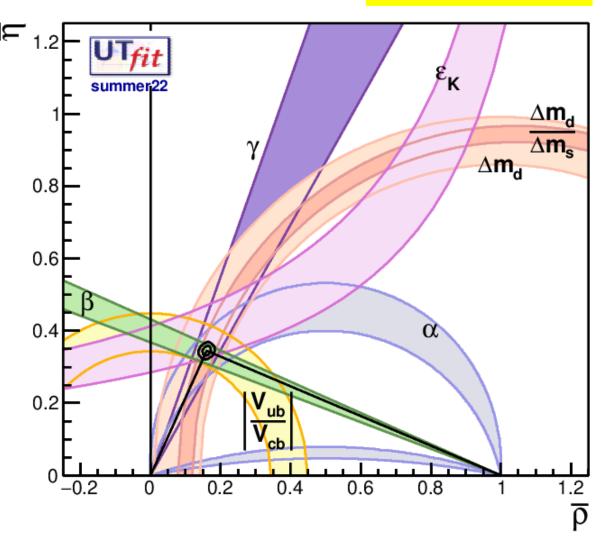
levels @ 95% Prob

~6%

$$\overline{\rho}$$
 = 0.160 ± 0.009 $\overline{\eta}$ = 0.345 ± 0.009

~3%



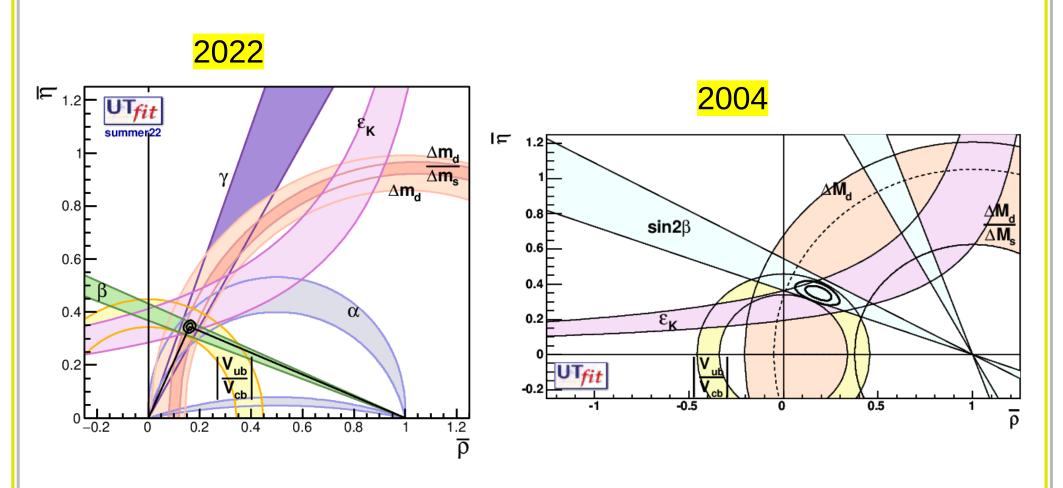


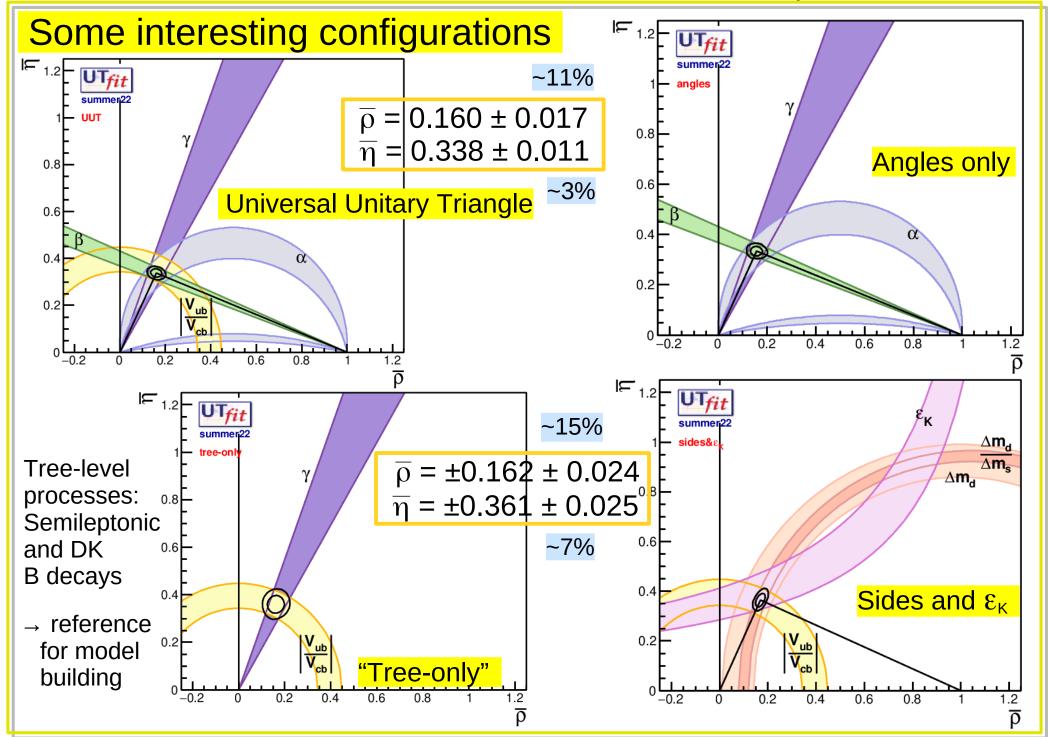
levels @ 95% Prob

~6%

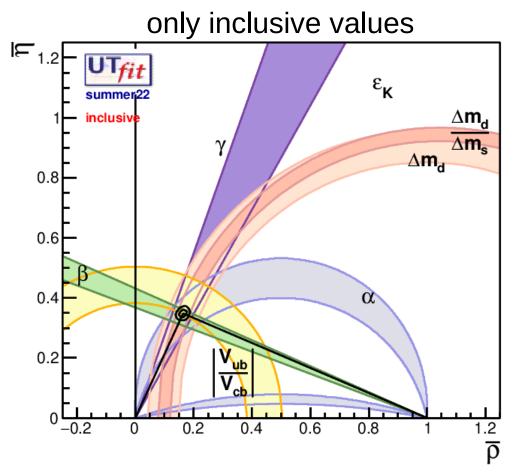
$$\overline{\rho}$$
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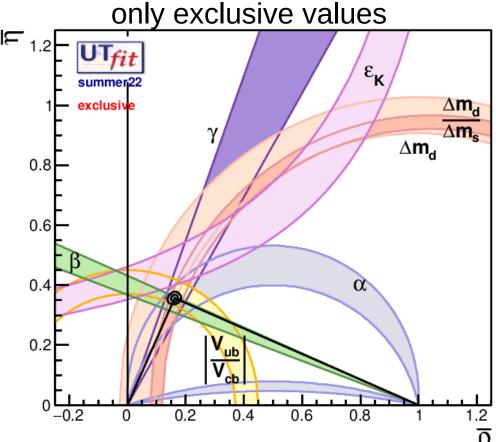
~3%





Inclusive vs Exclusive





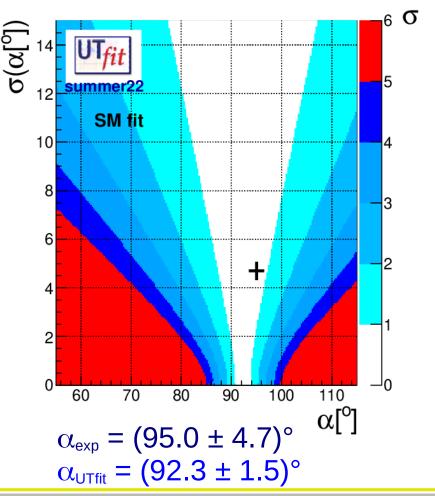
$$\overline{\rho}$$
 = 0.164 ± 0.009
 $\overline{\eta}$ = 0.348 ± 0.009
 $\sin 2\beta$ = 0.753 ± 0.028

$$\overline{
ho}$$
 = 0.162 ± 0.009
 $\overline{\eta}$ = 0.356 ± 0.009
 $\sin 2\beta$ = 0.755 ± 0.020

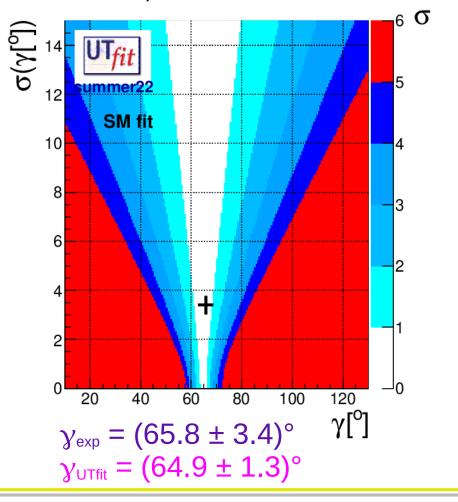
compatibility plots

A way to "measure" the agreement of a single measurement with the indirect determination from the fit using all the other inputs: test for the SM description of the flavour physics

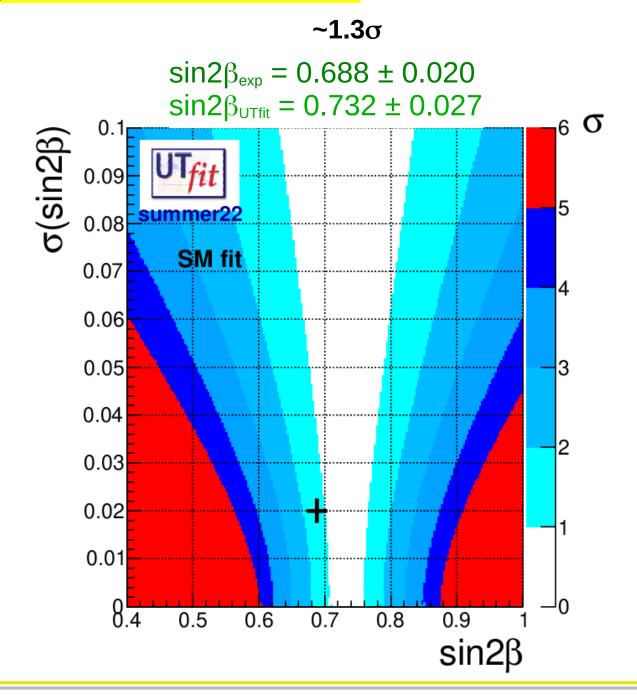
Color code: agreement between the predicted values and the measurements at better than 1, 2, ... $n\sigma$

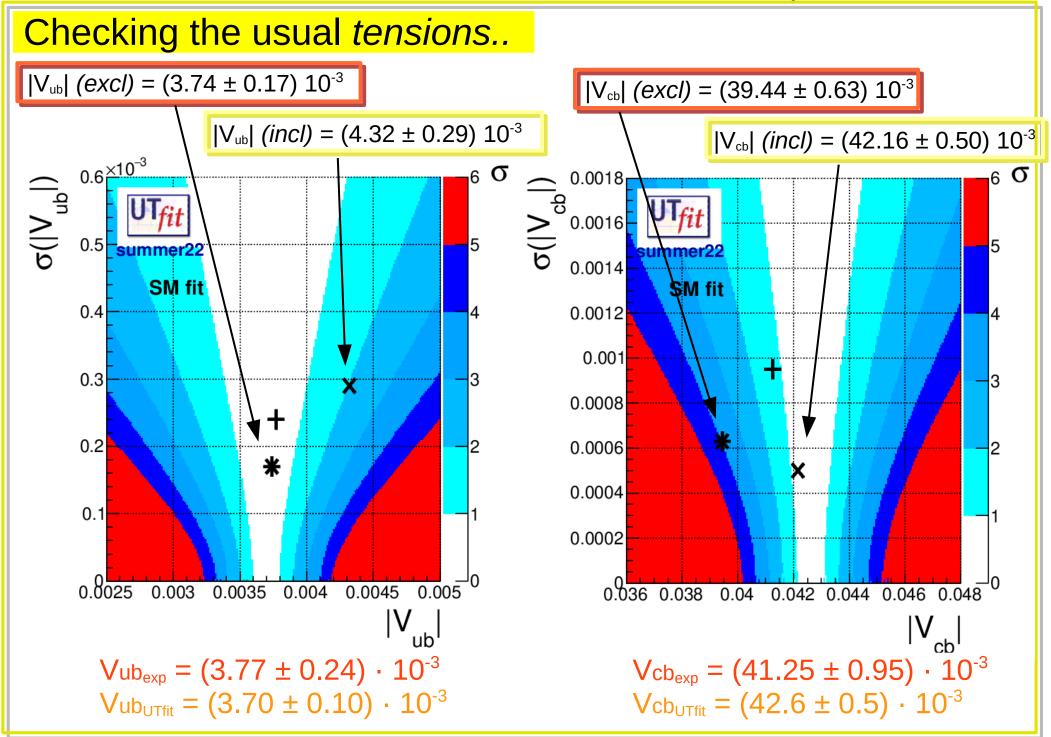


The cross has the coordinates (x,y)=(central value, error) of the direct measurement



Checking the usual tensions...



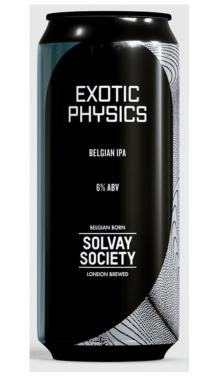


UT analysis including new physics

fit simultaneously for the CKM and the NP parameters (generalized UT fit)

- add most general loop NP to all sectors
- use all available experimental info
- find out NP contributions to $\Delta F=2$ transitions

B_d and B_s mixing amplitudes (2+2 real parameters):



$$A_{q} = C_{B_{q}} e^{2i\phi_{B_{q}}} A_{q}^{SM} e^{2i\phi_{q}^{SM}} = \left(1 + \frac{A_{q}^{NP}}{A_{q}^{SM}} e^{2i(\phi_{q}^{NP} - \phi_{q}^{SM})}\right) A_{q}^{SM} e^{2i\phi_{q}^{SM}}$$

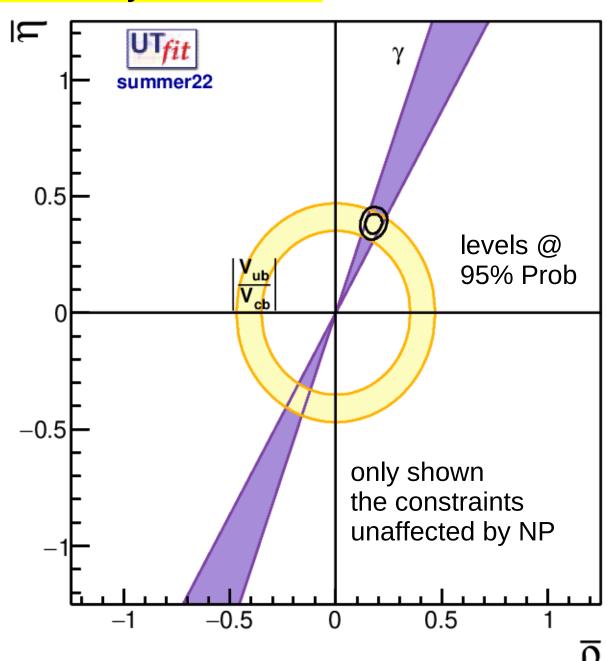
$$\begin{array}{l} \Delta \, m_{q/K} = C_{B_q/\Delta \, m_K} (\Delta \, m_{q/K})^{SM} \\ A_{CP}^{B_d \to J/\psi K_s} = \sin 2(\beta + \varphi_{B_d}) \\ A_{SL}^q = \operatorname{Im} \left(\Gamma_{12}^q / A_q \right) \end{array} \quad \begin{array}{l} \epsilon_K = C_\epsilon \epsilon_K^{SM} \\ A_{CP}^{B_s \to J/\psi \, \varphi} \sim \sin 2(-\beta_s + \varphi_{B_s}) \\ \Delta \, \Gamma^q / \Delta \, m_q = \operatorname{Re} \left(\Gamma_{12}^q / A_q \right) \end{array}$$

$$\epsilon_{K} = C_{\epsilon} \epsilon_{K}^{SM}$$

$$A_{CP}^{B_{s} \to J/\psi \phi} \sim \sin 2(-\beta_{s} + \phi_{B_{s}})$$

$$\Delta \Gamma^{q} / \Delta m_{q} = \text{Re} \left(\Gamma_{12}^{q} / A_{q} \right)$$

NP analysis results



$$\overline{\rho}$$
 = 0.169 ± 0.025
 $\overline{\eta}$ = 0.365 ± 0.026

SM is

$$\overline{\rho}$$
 = 0.160 ± 0.009 $\overline{\eta}$ = 0.345 ± 0.009

NP parameter results

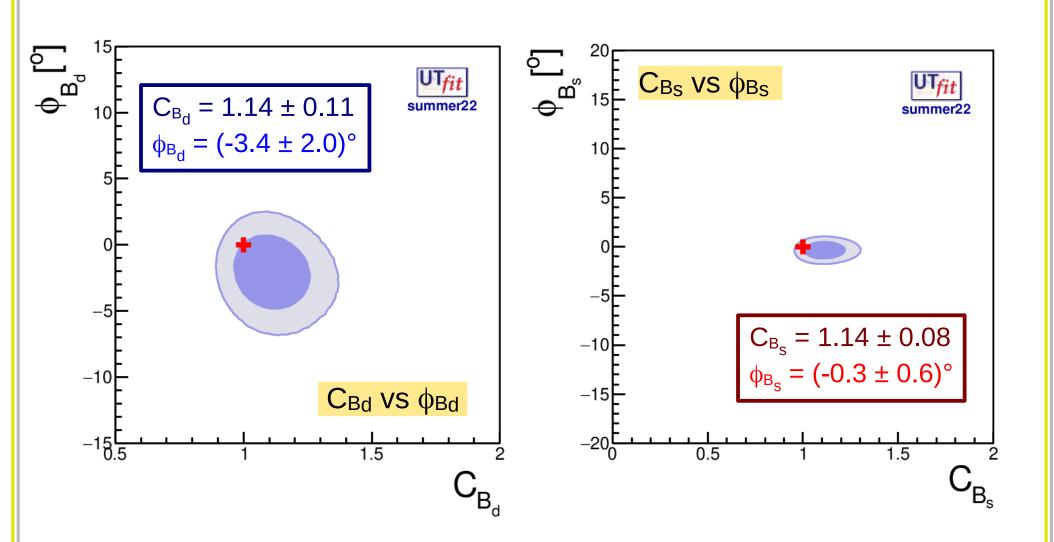
$$A_q = C_{B_q} e^{2i\phi_{B_q}} A_q^{SM} e^{2i\phi_q^{SM}}$$

dark: 68% light: 95%

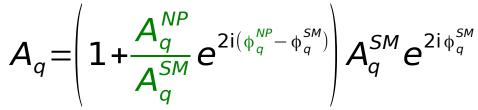
SM: red cross

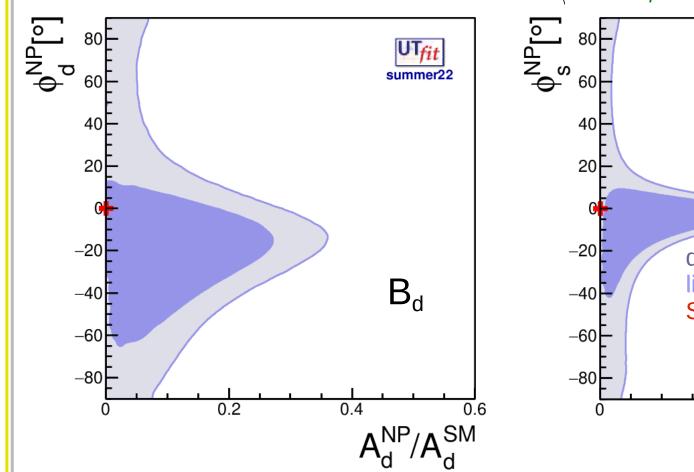
K system

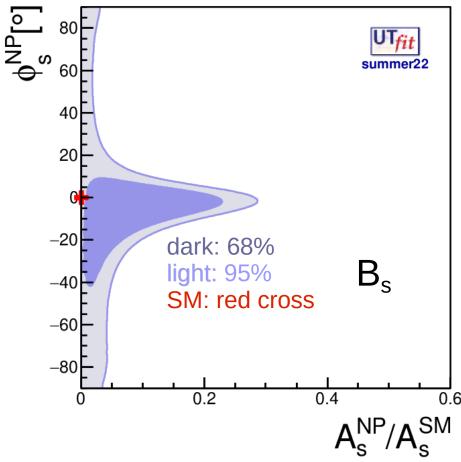
 $C_{e_{K}} = 1.12 \pm 0.12$



NP parameter results







The ratio of NP/SM amplitudes is:

- < 25% @68% prob. (35% @95%) in B_d mixing
- < 25% @68% prob. (30% @95%) in B_s mixing

testing the new-physics scale

M. Bona *et al.* (UTfit) JHEP 0803:049,2008 arXiv:0707.0636

R G E

At the high scale

new physics enters according to its specific features

At the low scale

use OPE to write the most general effective Hamiltonian. the operators have different chiralities than the SM NP effects are in the Wilson Coefficients C

$$C_i(\Lambda) = F_i \frac{L_i}{\Lambda^2}$$

$$\mathcal{H}_{\text{eff}}^{\Delta B=2} = \sum_{i=1}^{5} C_i Q_i^{bq} + \sum_{i=1}^{3} \tilde{C}_i \tilde{Q}_i^{bq}$$

$$Q_1^{q_i q_j} = \bar{q}_{jL}^{\alpha} \gamma_{\mu} q_{iL}^{\alpha} \bar{q}_{jL}^{\beta} \gamma^{\mu} q_{iL}^{\beta} ,$$

$$Q_2^{q_i q_j} = \bar{q}_{jR}^{\alpha} q_{iL}^{\alpha} \bar{q}_{jR}^{\beta} q_{iL}^{\beta} ,$$

$$Q_3^{q_i q_j} = \bar{q}_{jR}^{\alpha} q_{iL}^{\beta} \bar{q}_{jR}^{\beta} q_{iL}^{\alpha} ,$$

$$Q_4^{q_i q_j} = \bar{q}_{jR}^{\alpha} q_{iL}^{\alpha} \bar{q}_{jL}^{\beta} q_{iR}^{\beta} ,$$

$$Q_5^{q_i q_j} = \bar{q}_{jR}^{\alpha} q_{iL}^{\beta} \bar{q}_{jL}^{\beta} q_{iR}^{\alpha} .$$

- F_i : function of the NP flavour couplings
- L_i: loop factor (in NP models with no tree-level FCNC)
- Λ : NP scale (typical mass of new particles mediating $\Delta F=2$ processes)

testing the TeV scale

The dependence of C on Λ changes depending on the flavour structure. We can consider different flavour scenarios:

• Generic: $C(\Lambda) = \alpha/\Lambda^2$

 $F_i\sim 1$, arbitrary phase

• NMFV: $C(\Lambda) = \alpha \times |F_{SM}|/\Lambda^2$ $F_i \sim |F_{SM}|$, arbitrary phase • MFV: $C(\Lambda) = \alpha \times |F_{SM}|/\Lambda^2$ $F_1 \sim |F_{SM}|$, $F_{i\neq 1} \sim 0$, SM phase

 α (L_i) is the coupling among NP and SM

- \odot α ~ 1 for strongly coupled NP
- $\odot \alpha \sim \alpha_w$ (α_s) in case of loop coupling through weak (strong) interactions

If no NP effect is seen lower bound on NP scale Λ

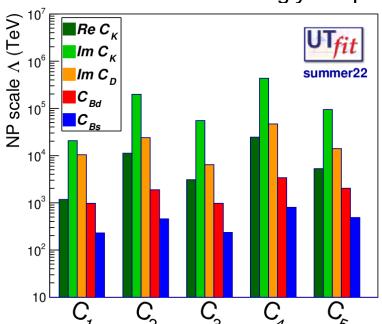
F is the flavour coupling and so

F_{SM} is the combination of CKM factors for the considered process

results from the Wilson coefficients

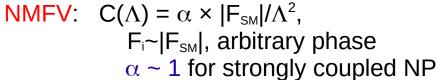
Generic: $C(\Lambda) = \alpha/\Lambda^2$, $F_i \sim 1$, arbitrary phase

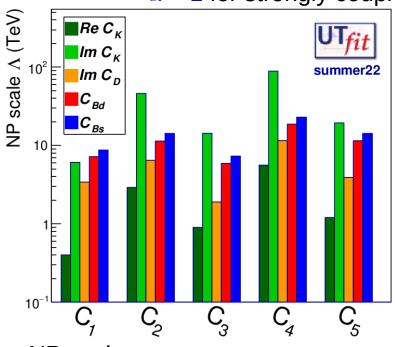
 $\alpha \sim 1$ for strongly coupled NP



 $\Lambda > 4.4 \ 10^5 \ TeV$

Lower bounds on NP scale (at 95% prob.)





 $\Lambda > 95 \text{ TeV}$

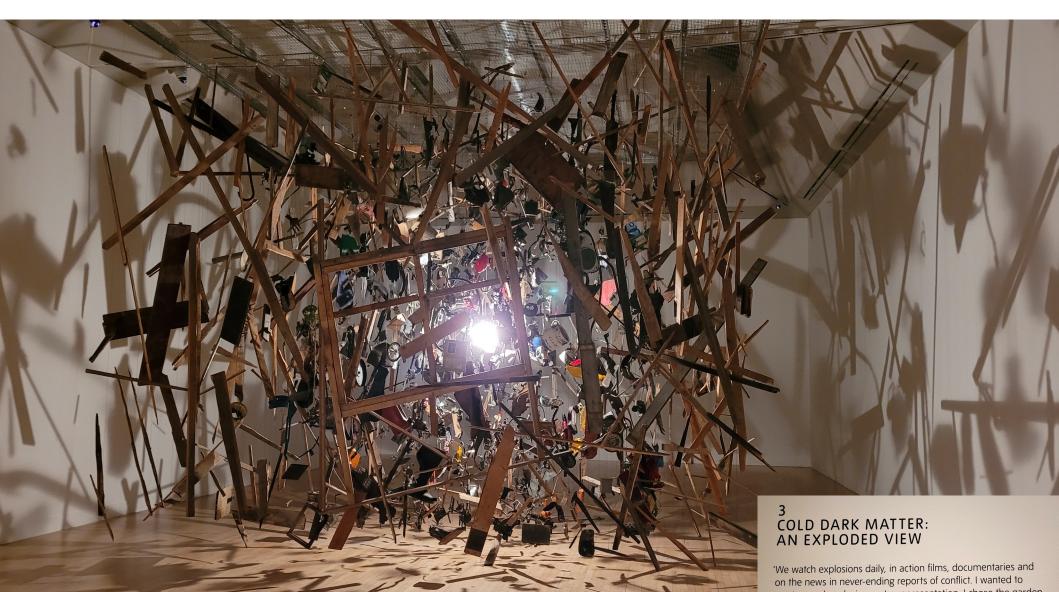
 $\alpha \sim \alpha_w$ in case of loop coupling through weak interactions

 $\Lambda > 1.3 \ 10^4 \ TeV$

 $\alpha \sim \alpha_w$ in case of loop coupling through weak interactions

 $\Lambda > 2.9 \text{ TeV}$

for lower bound for loop-mediated contributions, simply multiply by α_s (~ 0.1) or by α_w (~ 0.03).



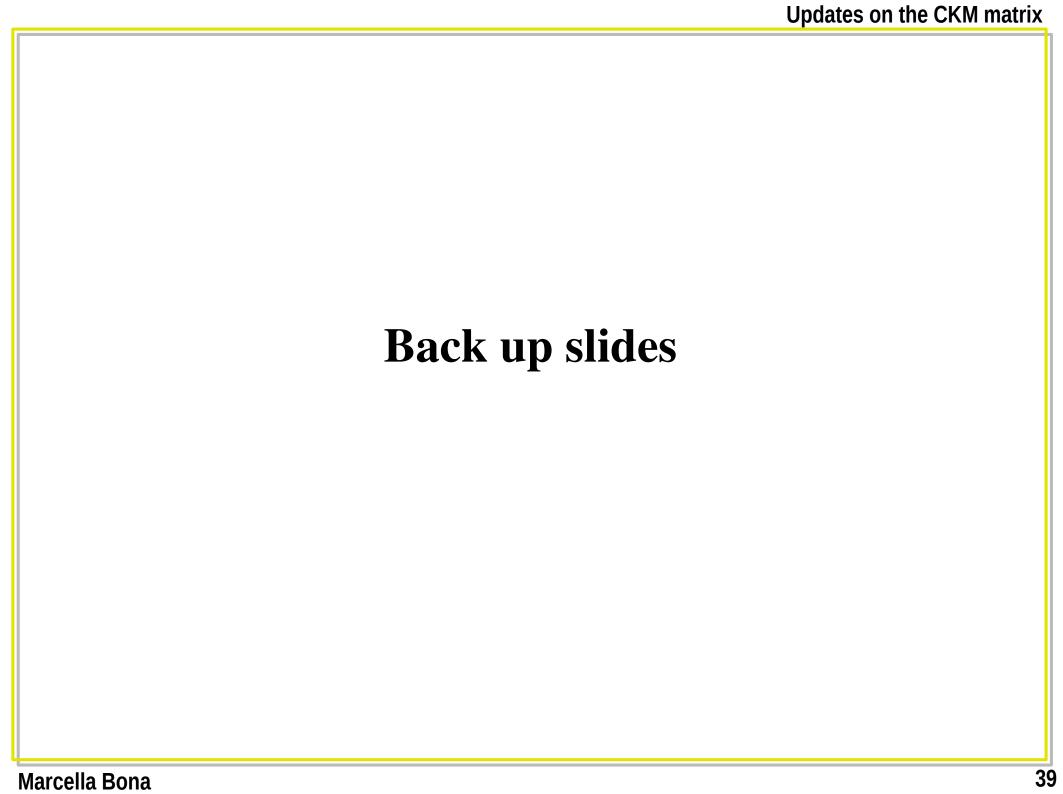
Cornelia Parker

create a real explosion, not a representation. I chose the garden shed because it's the place where you store things you can't quite throw away.

The shed was blown up at the Army School of Ammunition. We used Semtex, a plastic explosive popular with terrorists. I pressed the plunger that blew the shed skywards. The soldiers helped me comb the field afterwards, picking up the blackened, mangled objects.

conclusions

- SM analysis displays very good (improved) overall consistency
- Still open discussion on semileptonic inclusive vs exclusive: exclusive fit shows tension, V_{cb} now showing the biggest discrepancy..
- UTA provides determination of NP contributions to $\Delta F=2$ amplitudes. It currently leaves space for NP at the level of 25-35%
- So the scale analysis points to high scales for the generic scenario and at the limit of LHC reach for weak coupling. Indirect searches are not only complementary to direct searches, but they might be the main way to glimpse at new physics.



Some updated inputs

lattice inputs updated for this summer

Observables	Measurement
B _K	0.756 ± 0.016
f _{Bs}	0.2301 ± 0.0012
f _{Bs} /f _{Bd}	1.208 ± 0.005
B _{Bs} /B _{Bd}	1.015 ± 0.021
B _{Bs}	1.284 ± 0.059

We quote, instead, the weighted average of the $N_f=2+1+1$ and $N_f=2+1$ results with the error rescaled when chi2/dof > 1, as done by FLAG for the $N_f=2+1+1$ and $N_f=2+1$ averages separately [new HPQCD (2+1+1) result 1907.01025]

V_{ud} and V_{us} updated for this summer

Observables	Measurement
V_{ud}	0.97433 ± 0.00019
V_{us}	0.2249 (± 0.0004)

 V_{ud} is taken from the PDG average of V_{ud} FLAG numbers (for 2+1+1 and 2+1) and superallowed beta decays value. PDG scale factor S=2.0

V_{us} is not used in the fit

obtained excluding the given constraint from the fit

Observables	Measurement	Prediction A	Pull (#σ)
sin2β	0.688 ± 0.020	0.732 ± 0.027	~ 1.3
γ	65.8 ± 3.4	64.9 ± 1.3	< 1
α	95.0 ± 4.7	92.3 ± 1.5	< 1
ε _κ · 10³	2.228 ± 0.001	2.04 ± 0.14	< 1
V _{cb} · 10 ³	41.25 ± 0.95	42.6 ± 0.5	< 1
$ V_{cb} \cdot 10^3$ (incl)	42.16 0.50		< 1
V _{cb} • 10 ³ (excl)	39.44 0.63		~ 4.0
V _{ub} · 10 ³	3.77 ± 0.24	3.70 ± 0.10	< 1
V _{ub} • 10 ³ (incl)	4.32 ± 0.29	-	~ 2.0
V _{ub} • 10 ³ (excl)	3.74 ± 0.17	-	< 1
BR(B $\rightarrow \tau \nu$)[10 ⁻⁴]	1.09 ± 0.24	0.88 ± 0.05	< 1
A _{SL} ^d · 10 ³	-2.1 ± 1.7	-0.33 ± 0.02	< 1
A _{SL} ^s · 10 ³	-0.6 ± 2.8	0.014 ± 0.001	< 1

We obtain the
predictions for the lattice
parameters in different
configurations in the fit:

- only lattice parameters ratios
 - (F_{Bs}/F_B, B_{Bs}/B_{Bd} used)
- only B parameters
 - (B_{Bs}¹, B_{Bs}/B_{Bd} used)
- nlv decav constants f

UI	ily u	ecay	Constants
•	$(f_{Bs},$	$f_{\text{Bs}}/f_{\text{B}}$	included)

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Observables	Measurement	Prediction
B _K	0.756 ± 0.016	0.832 ± 0.054
No B lattice		
$f_B\sqrt{B_{Bd}}$	(0.2142 ± 0.0056)	0.212 ± 0.010
$f_{Bs}\sqrt{B_{Bs}}$	(0.2607 ± 0.0061)	0.259 ± 0.010
ξ	(1.217 ± 0.014)	1.225 ± 0.033
Ratios only		
f _{Bs}	0.2301 ± 0.0012	0.227 ± 0.009
B _{Bs}	1.284 ± 0.059	1.30 ± 0.10
B pars only		
f _{Bs} /f _{Bd}	1.208 ± 0.005	1.215 ± 0.028
f _{Bs}	0.2301 ± 0.0012	0.228 ± 0.008
f pars only		
B _{Bs} /B _{Bd}	1.015 ± 0.021	1.017 ± 0.028
B_Bs	1.284 ± 0.059	1.290 ± 0.065

new-physics-specific constraints

$$A_{\rm SL}^s \equiv \frac{\Gamma(\bar{B}_s \to \ell^+ X) - \Gamma(B_s \to \ell^- X)}{\Gamma(\bar{B}_s \to \ell^+ X) + \Gamma(B_s \to \ell^- X)} = \operatorname{Im}\left(\frac{\Gamma_{12}^s}{A_s^{\rm full}}\right)$$

semileptonic asymmetries in B⁰ and B_s: sensitive to NP effects in both size and phase. Taken from the latest HFLAV.

Cleo, BaBar, Belle,

same-side dilepton charge asymmetry:

admixture of B_s and B_d so sensitive to NP effects in both.

$$A_{\rm SL}^{\mu\mu} \times 10^3 = -7.9 \pm 2.0$$

lifetime τ^{FS} in flavour-specific final states:

average lifetime is a function to the width and the width difference

$$\tau^{FS}(B_s) = 1.527 \pm 0.011 \text{ ps}$$

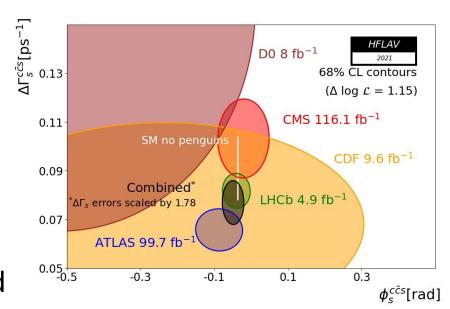
 ϕ_s =2 β_s vs $\Delta\Gamma_s$ from $B_s \rightarrow J/\psi \phi$ angular analysis as a function of proper time and b-tagging

$$\phi_s = -0.049 \pm 0.019 \text{ rad}$$

D0 arXiv:1106.6308

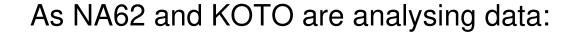
$$A_{\rm SL}^{\mu\mu} = \frac{f_d \chi_{d0} (A_{\rm SL}^d) + f_s \chi_{s0} (A_{\rm SL}^s)}{f_d \chi_{d0} + f_s \chi_{s0}}$$

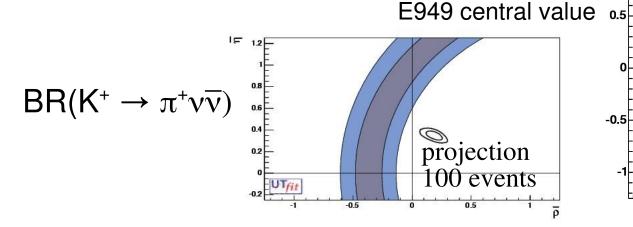
D0 and LHCb

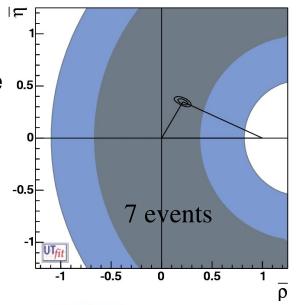


some old plots coming back to fashion:

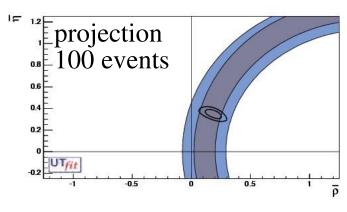
2007 global fit area

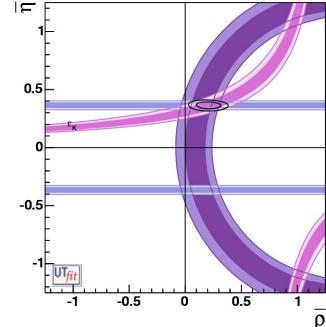






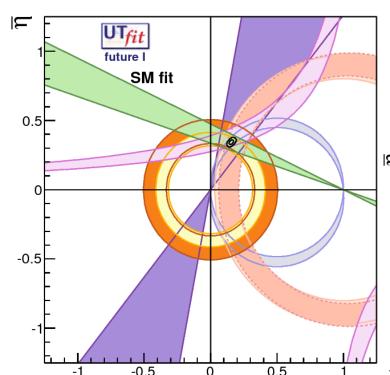
SM central value





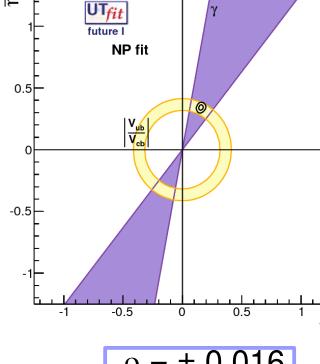
including BR($K^0 \rightarrow \pi^0 \nu \overline{\nu}$) SM central value

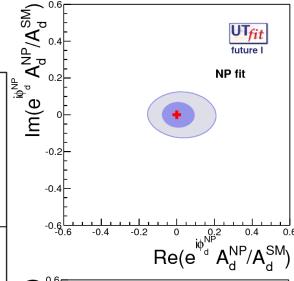


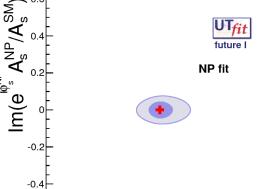


 $\rho = \pm 0.015$ $\eta = \pm 0.015$ **future I** scenario: errors from

Belle II at 5/ab + LHCb at 10/fb







 $\frac{\overline{\rho}}{\eta} = 0.154 \pm 0.015$ $\frac{\overline{\rho}}{\eta} = 0.344 \pm 0.013$

current sensitivity

$$\frac{\overline{\rho}}{\eta}$$
 = 0.150 ± 0.027
 $\frac{\overline{\rho}}{\eta}$ = 0.363 ± 0.025

 $\rho = \pm 0.016$

 $\eta = \pm 0.019$