

Z portal to a Confining Hidden Sector

Ennio Salvioni
CERN-TH



3rd FCC Physics and
Experiments Workshop

January 16, 2020

Z portal to a Confining Hidden Sector

Ennio Salvioni
CERN-TH



based on 1906.02198 [JHEP] + in progress

with H.-C. Cheng (UC Davis), C. Verhaaren (UC Irvine), L. Li (HKUST)

Introduction



A **confining hidden sector** is generic in Neutral Naturalness theories

Naturalness-inspired hidden sectors

- A **confining hidden sector** is generic in Neutral Naturalness theories
- **Neutral Naturalness** = models that solve little hierarchy problem with **color-neutral** Top Partners
- Many proposals: Twin Higgs, Folded SUSY, Tripled Top, Hyperbolic Higgs, ...
- Phenomenology is strikingly non-standard, new paradigm for searches

[Chacko, Goh, Harnik 2005]

[Burdman, Chacko, Goh, Harnik 2006]

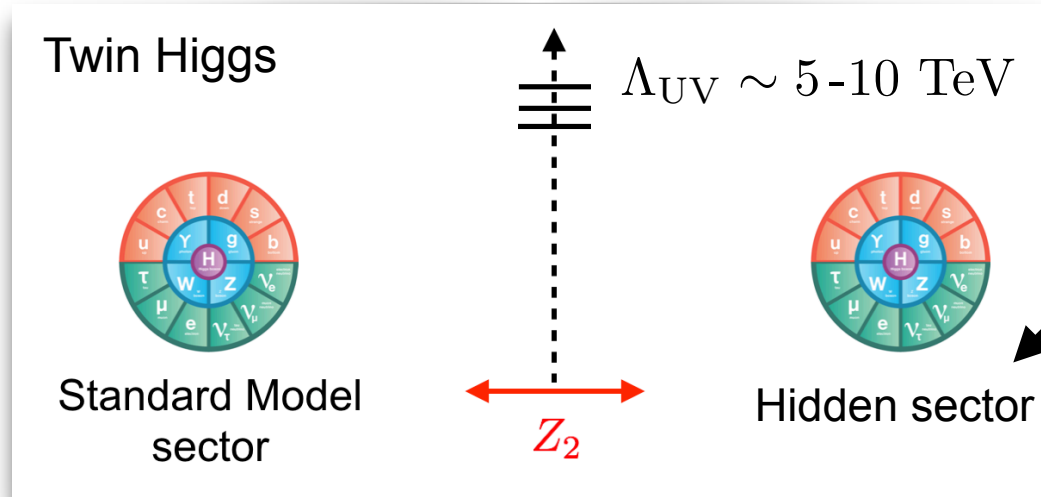
[Cheng, Li, Salvioni, Verhaaren 2018]

[Cohen, Craig, Giudice, McCullough 2018]

Neutral Naturalness

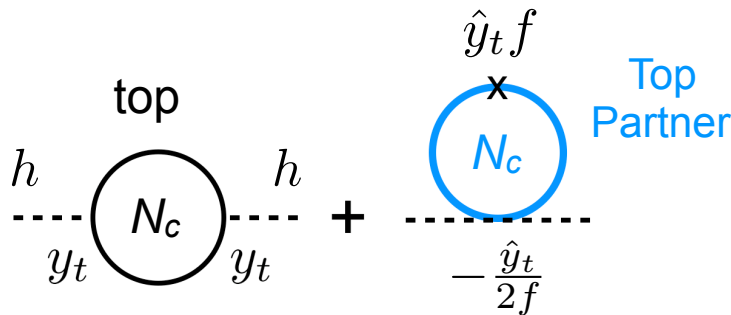
- A **discrete symmetry** relates the SM to the hidden sector

[Chacko, Goh, Harnik, 2005]



Has own $SU(3)$ color symmetry

- 1 loop:



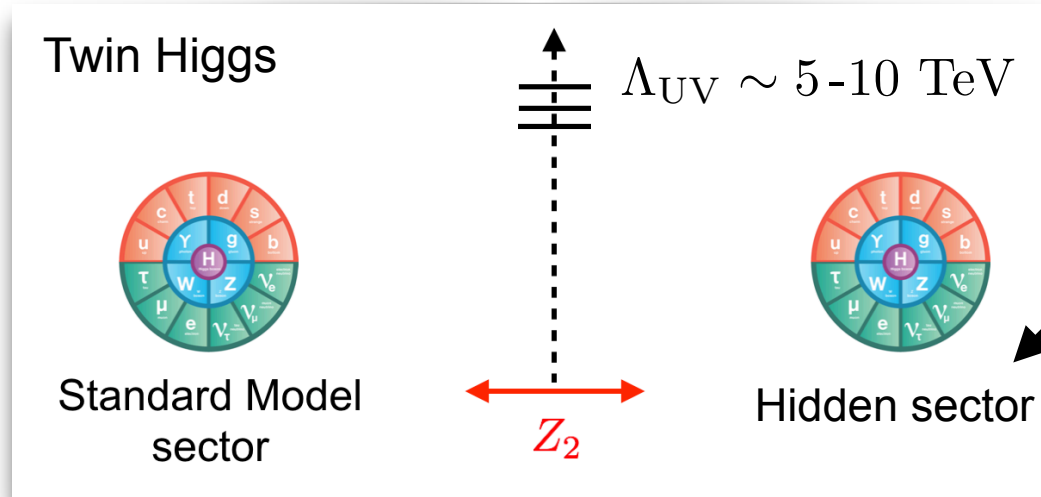
$y_t \xleftrightarrow{Z_2} \hat{y}_t$

quadratic corrections
to Higgs mass cancel out

Hidden confinement

- A **discrete symmetry** relates the SM to the hidden sector

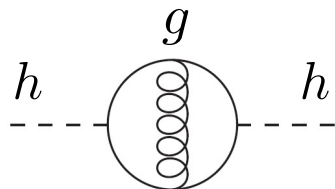
[Chacko, Goh, Harnik, 2005]



Has own $SU(3)$ color symmetry

- 2 loops:

[Craig, Katz, Strassler, Sundrum, 2015]



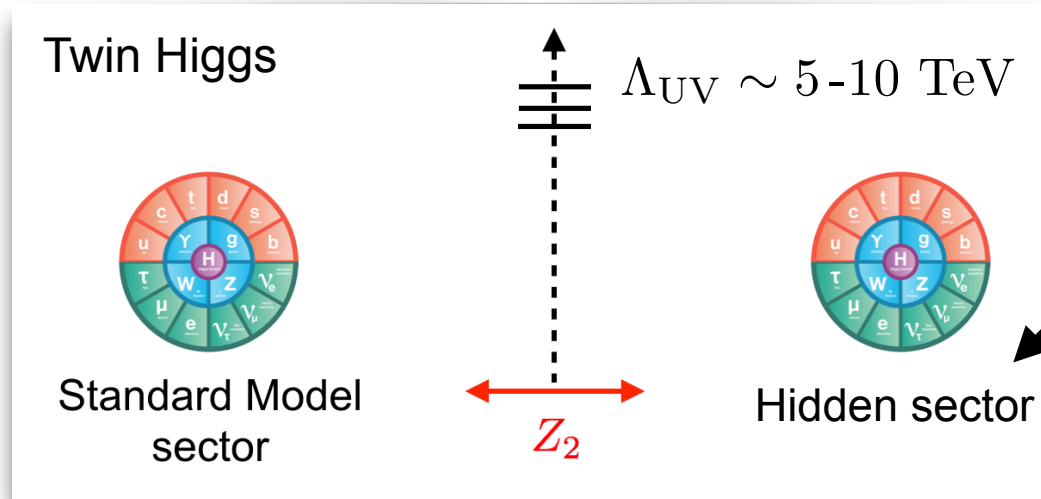
$$\rightarrow m_{h,IR}^2 \approx m_h^2(\Lambda_{UV}) + \left[\frac{3y_t^2 g_s^2}{8\pi^4} + \frac{3(y_t^2 - \hat{y}_t^2)}{4\pi^2} \right]_{\Lambda_{UV}} \Lambda_{UV}^2$$

if hidden color is only global

Hidden confinement

- A **discrete symmetry** relates the SM to the hidden sector

[Chacko, Goh, Harnik, 2005]



Has own $SU(3)$ color symmetry

- 2 loops:

[Craig, Katz, Strassler, Sundrum, 2015]

The diagram shows a Higgs boson h (dashed line) entering a loop with a gluon g (curly line) and exiting as h . A blue arrow points to the right, leading to the equation:

$$m_{h,IR}^2 \approx m_h^2(\Lambda_{UV}) + \left[\frac{3y_t^2 g_s^2}{8\pi^4} + \frac{3(y_t^2 - \hat{g}_t^2)}{4\pi^2} \right]_{\Lambda_{UV}} \Lambda_{UV}^2$$

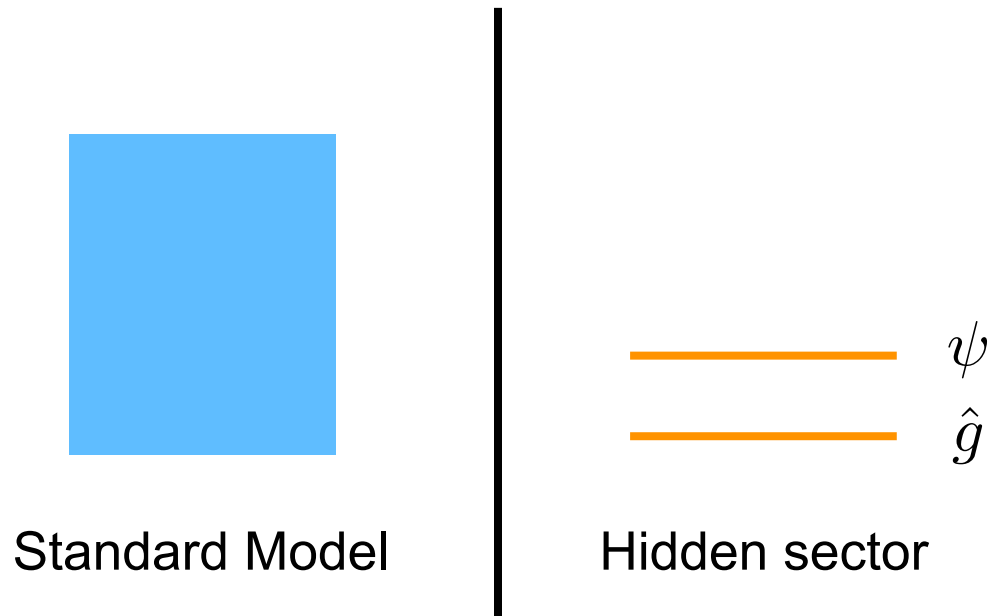
➔ gauge hidden $SU(3)$ with $g_s \simeq \hat{g}_s$ at Λ_{UV} , **confines at** $\Lambda \gtrsim \Lambda_{QCD}$

Hidden confinement

- Confining hidden sector with $0.1 \lesssim \Lambda/\text{GeV} \lesssim 10$ is generic expectation in Neutral Naturalness
- Details depend on
 - ➔ Mass spectrum of hidden matter
 - ➔ Portals to SM: Higgs is mandatory, others?
 - ➔ Additional interactions within hidden sector?
- Many models exist, aim to identify **representative cases** with distinct phenomenology, to set motivated targets for experimental searches

In this talk

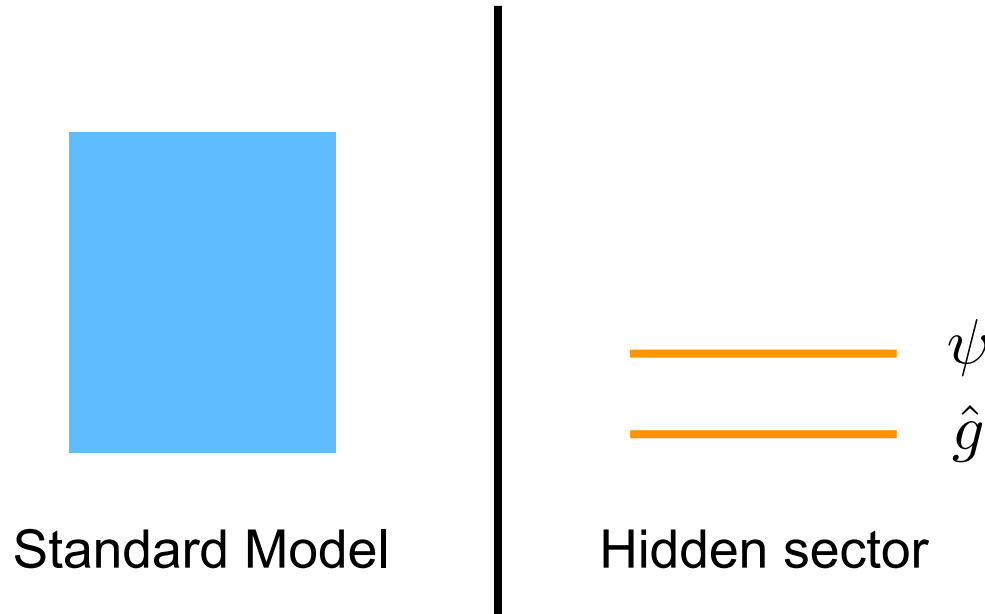
- Hidden QCD with 1 light fermion flavor



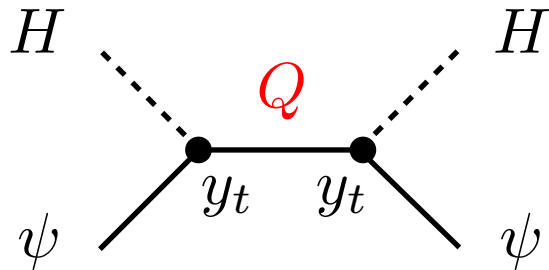
In this talk

- Hidden QCD with 1 light fermion flavor
- Connected to SM by Higgs and Z portals

$$\mathcal{L}_6 \sim \frac{y_t^2}{M^2} \left(iH^\dagger \overleftrightarrow{D}_\mu H \bar{\psi}_R \gamma^\mu \psi_R + |H|^2 m_\psi \bar{\psi} \psi + \frac{c_g \alpha_d}{12\pi} |H|^2 \hat{G}_{\mu\nu} \hat{G}^{\mu\nu} \right)$$



Ultraviolet origin?



ψ is only light flavor in hidden sector

Q are TeV-scale fermions charged under SM electroweak + hidden color

- Fairly common ingredients in Neutral Naturalness:
 - ✓ Typically only 3rd generation is copied, hidden sector has fewer flavors
 - ✓ TeV-scale particles play roles in Higgs mass stabilization
- Examples: **models realizing accidental SUSY**, composite Twin Higgs

[Cheng, Li, Salvioni, Verhaaren, 2019]

Description of Hidden Sector

Phenomenology

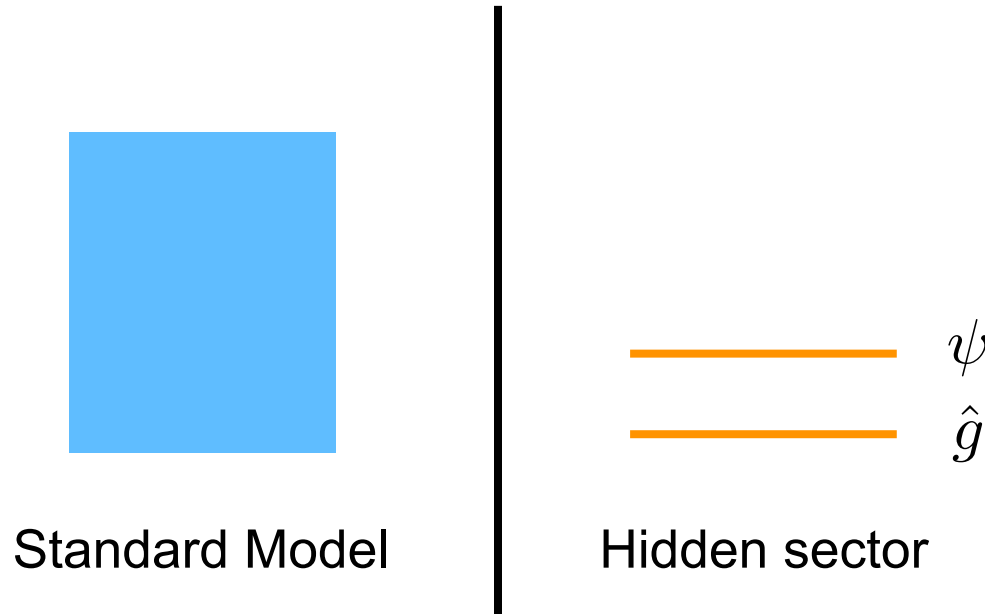
Description of Hidden Sector

Phenomenology

In this talk

- Hidden QCD with 1 light fermion flavor
- Connected to SM by Higgs and Z portals

$$\mathcal{L}_6 \sim \frac{y_t^2}{M^2} \left(iH^\dagger \overleftrightarrow{D}_\mu H \bar{\psi}_R \gamma^\mu \psi_R + |H|^2 m_\psi \bar{\psi} \psi + \frac{c_g \alpha_d}{12\pi} |H|^2 \hat{G}_{\mu\nu} \hat{G}^{\mu\nu} \right)$$



In this talk

- Hidden QCD with 1 light fermion flavor
- Connected to SM by Higgs and Z portals

$$\mathcal{L}_6 \sim \frac{y_t^2}{M^2} \left(iH^\dagger \overleftrightarrow{D}_\mu H \bar{\psi}_R \gamma^\mu \psi_R + |H|^2 m_\psi \bar{\psi} \psi + \frac{c_g \alpha_d}{12\pi} |H|^2 \hat{G}_{\mu\nu} \hat{G}^{\mu\nu} \right)$$

In this talk

- Hidden QCD with 1 light fermion flavor
- Connected to SM by Higgs and Z portals

$$\mathcal{L}_6 \sim \frac{y_t^2}{M^2} \left(iH^\dagger \overleftrightarrow{D}_\mu H \bar{\psi}_R \gamma^\mu \psi_R + |H|^2 \cancel{m_\psi \bar{\psi} \psi} + \frac{c_g \alpha_d}{12\pi} |H|^2 \hat{G}_{\mu\nu} \hat{G}^{\mu\nu} \right)$$

Higgs portal

$$\text{BR}(h \rightarrow \hat{g}\hat{g}) \approx 1.3 \times 10^{-5} \left(\frac{2 \text{ TeV}}{M} \right)^4 \left(\frac{\alpha_d}{0.18} \right)^2 \left(\frac{c_g}{1} \right)^2$$

In this talk

- Hidden QCD with 1 light fermion flavor
- Connected to SM by Higgs and Z portals

$$\mathcal{L}_6 \sim \frac{y_t^2}{M^2} \left(iH^\dagger \overleftrightarrow{D}_\mu H \bar{\psi}_R \gamma^\mu \psi_R + |H|^2 \cancel{m_\psi} \bar{\psi} \psi + \frac{c_g \alpha_d}{12\pi} |H|^2 \hat{G}_{\mu\nu} \hat{G}^{\mu\nu} \right)$$

Z portal

Higgs portal

$$\text{BR}(h \rightarrow \hat{g}\hat{g}) \approx 1.3 \times 10^{-5} \left(\frac{2 \text{ TeV}}{M} \right)^4 \left(\frac{\alpha_d}{0.18} \right)^2 \left(\frac{c_g}{1} \right)^2$$

$$\text{BR}(Z \rightarrow \bar{\psi}\psi) \approx 2.2 \times 10^{-5} \left(\frac{2 \text{ TeV}}{M} \right)^4$$

In this talk

- At the LHC, $N_Z / N_h \sim 10^3$
- At FCC-ee, up to 10^{12} Z's versus $< 10^7$ Higgses

 **Z decays to hidden sector dominate**

Z portal

Higgs portal

$$\text{BR}(h \rightarrow \hat{g}\hat{g}) \approx 1.3 \times 10^{-5} \left(\frac{2 \text{ TeV}}{M}\right)^4 \left(\frac{\alpha_d}{0.18}\right)^2 \left(\frac{c_g}{1}\right)^2$$

$$\text{BR}(Z \rightarrow \bar{\psi}\psi) \approx 2.2 \times 10^{-5} \left(\frac{2 \text{ TeV}}{M}\right)^4$$

The hidden hadrons

- $SU(3)$ QCD with 1 flavor: no chiral symmetry  no Goldstone bosons

- Partial results available from lattice:

[Farchioni, Montvay, Münster,
Scholz, Sudmann, Wuilloud, 2007]

- Baryons significantly heavier than mesons
- Ratio of scalar/pseudoscalar meson masses: $m_{\hat{S}}/m_{\hat{P}} \approx 1.5$
- No info on vector meson

The hidden hadrons

- $SU(3)$ QCD with 1 flavor: no chiral symmetry \rightarrow no Goldstone bosons
- Partial results available from lattice: [Farchioni, Montvay, Münster, Scholz, Sudmann, Wuilloud, 2007]
 - Baryons significantly heavier than mesons
 - Ratio of scalar/pseudoscalar meson masses: $m_{\hat{S}}/m_{\hat{P}} \approx 1.5$
 - No info on vector meson

- **Assume**

$$m_{\hat{P}} \lesssim m_{\hat{V}} < m_{\hat{S}}$$

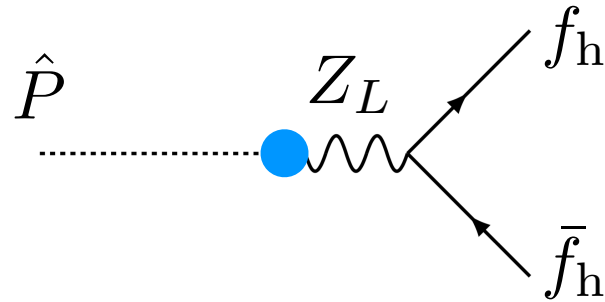
0^{-+} 1^{--} 0^{++} (J^{PC})

- Mesons decay back to visible sector through Z and h portals

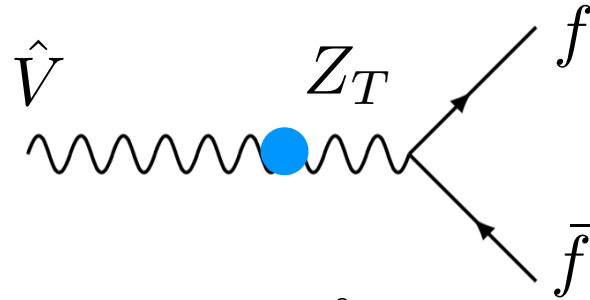
Hidden meson decays

- It turns out, **all** lightest mesons decay dominantly via **Z portal**

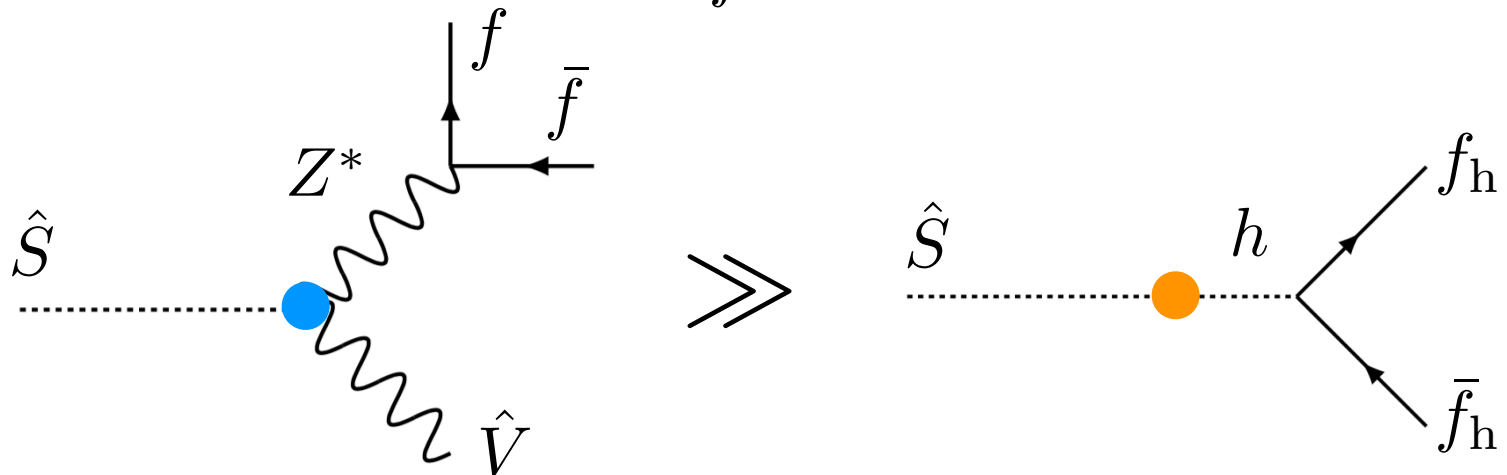
- Pseudoscalar



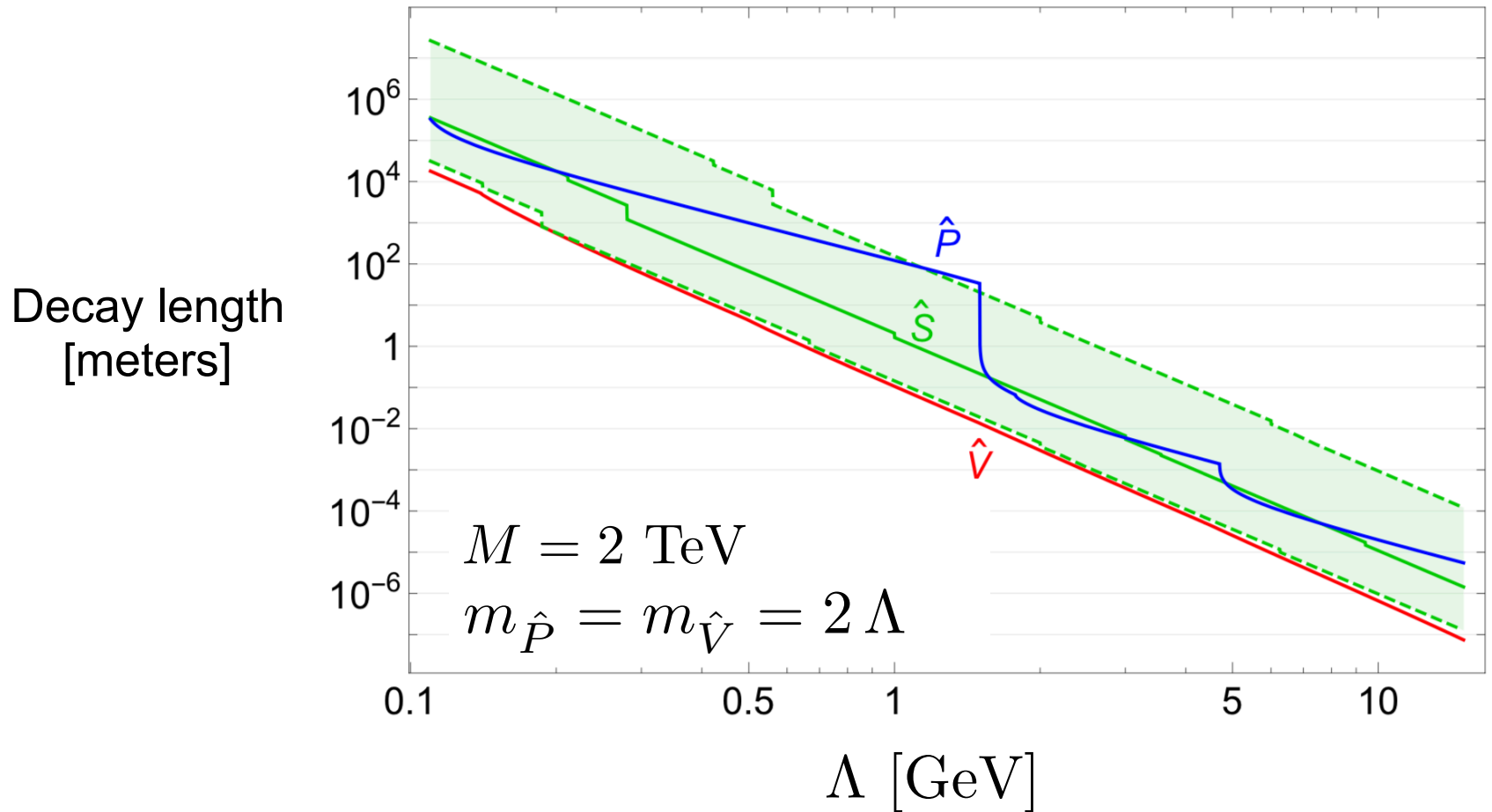
- Vector



- Scalar



Hidden meson decays



Two main parameters determine pheno:

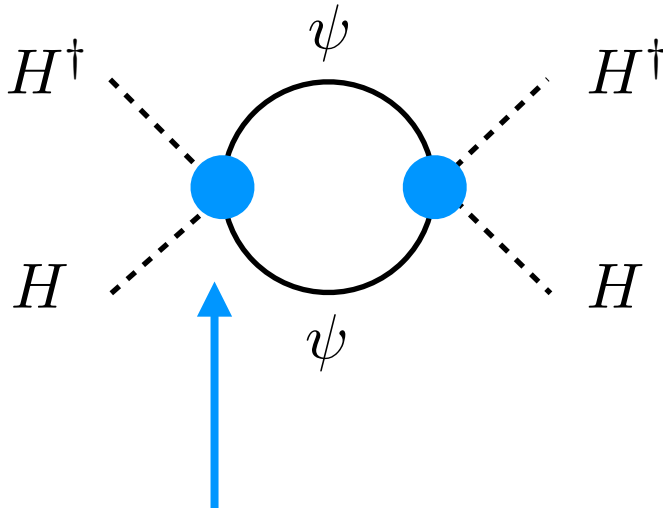
- Hidden confinement scale Λ
- Mediation scale M

Description of Hidden Sector

Phenomenology

Indirect vs direct probes

- EW precision tests:



$$\Delta\hat{T} = \frac{4}{3} \frac{N_d y_t^2}{16\pi^2} \frac{m_t^2}{M^2}$$

$$\Delta\hat{T} \lesssim 10^{-3} \quad \rightarrow \quad M \gtrsim 0.9 \text{ TeV}$$

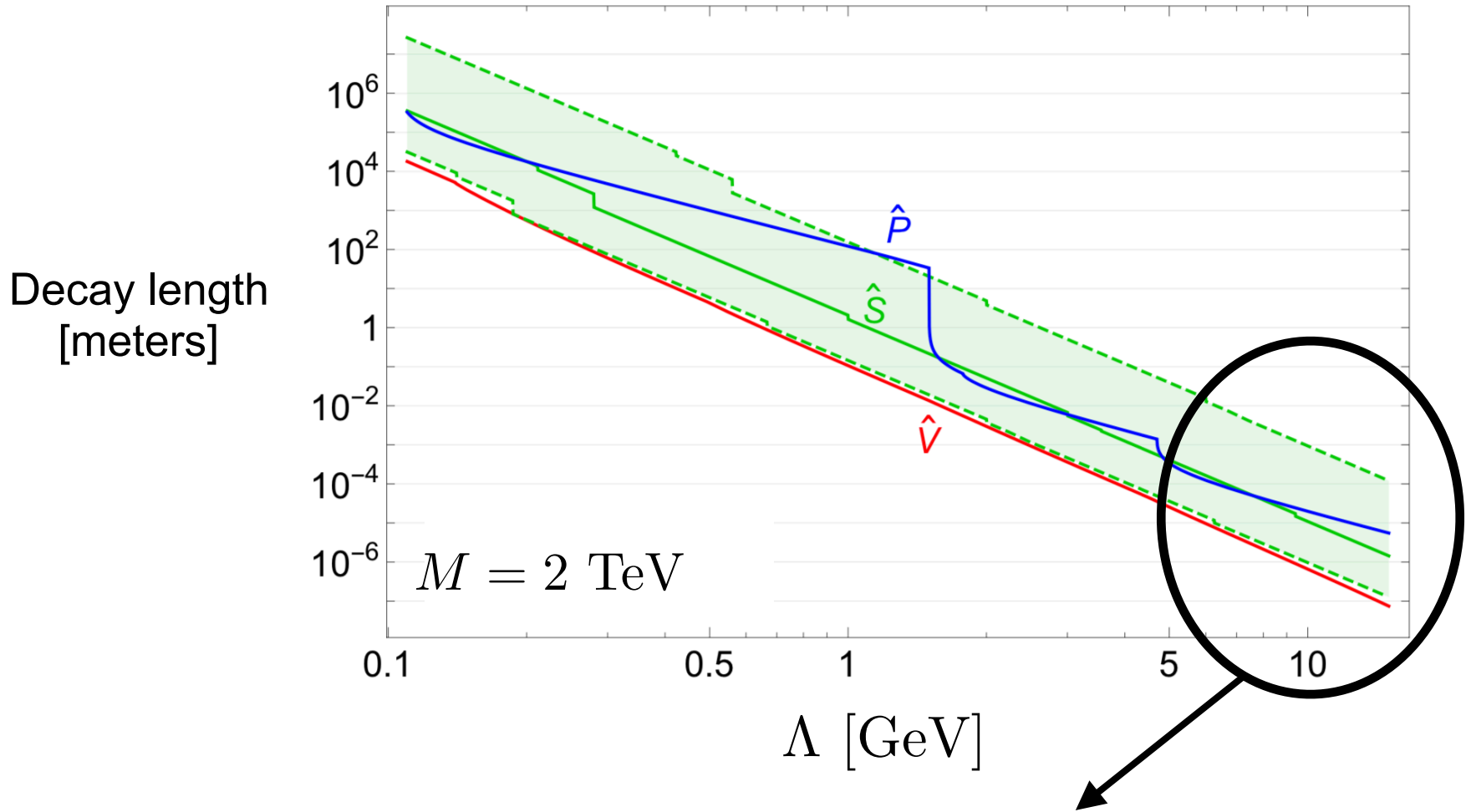
$$\frac{y_t^2}{M^2} (iH^\dagger \overleftrightarrow{D}_\mu H \bar{\psi}_R \gamma^\mu \psi_R)$$

FCC-ee:

$$\Delta\hat{T} \lesssim 10^{-4} \quad \rightarrow \quad M \gtrsim 2.7 \text{ TeV}$$

- Can direct probes do better? Z decays to hidden hadrons

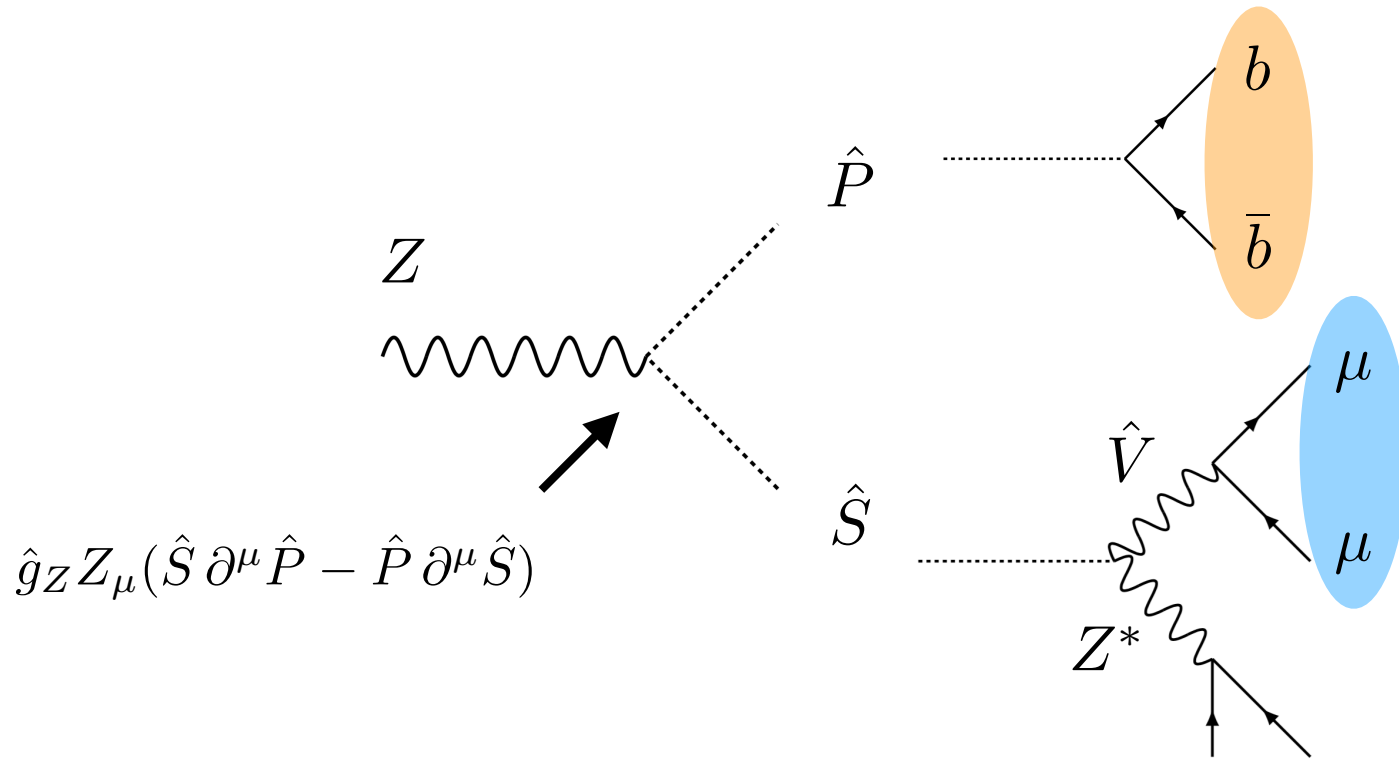
Phenomenology/1



$\Lambda \gtrsim 10 \text{ GeV}, M \sim \text{few TeV} : Z \rightarrow 2 \text{ mesons, prompt decays}$

Phenomenology/1

- If mesons have masses of 10-30 GeV, the Z decays to 2-body final states
- Example:



➔ $Z \rightarrow (b\bar{b})(\mu\mu) + X$

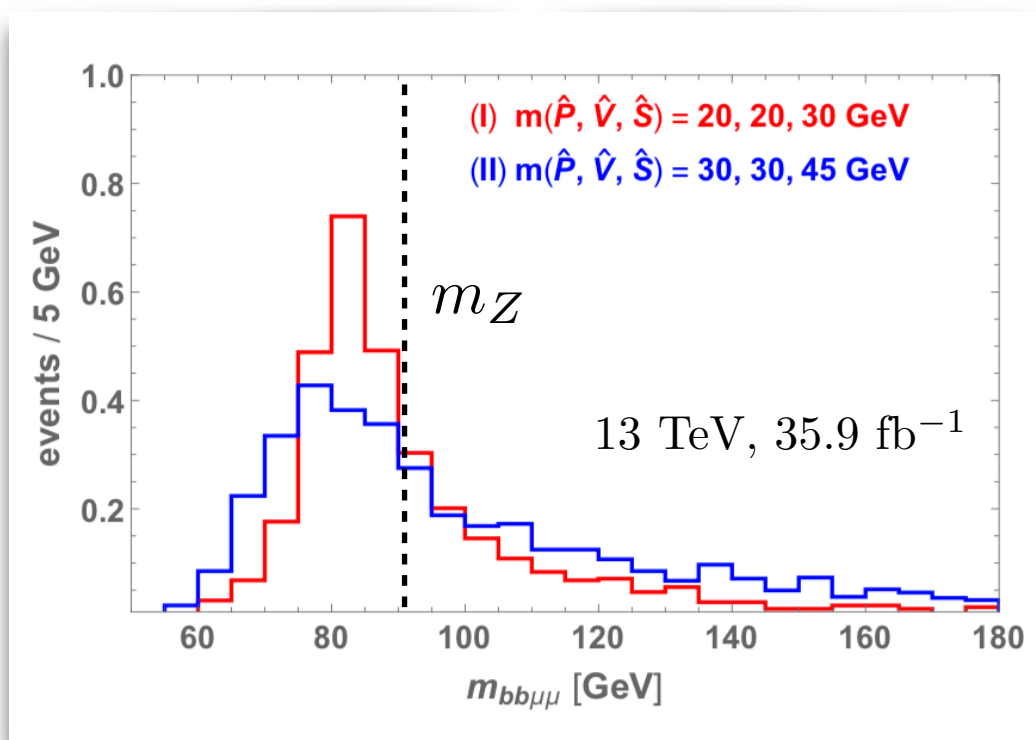
Phenomenology/1

- No dedicated LHC search yet, but can learn from $h \rightarrow aa \rightarrow (b\bar{b})(\mu\mu)$
- CMS analysis:

[CMS, 1812.06359]

$$p_T^{\mu 1,2} > 20, 9 \text{ GeV}, \quad p_T^{b 1,2} > 20, 15 \text{ GeV}$$

- For Z signal, selection efficiency ~ 10 times smaller than for Higgs (!)



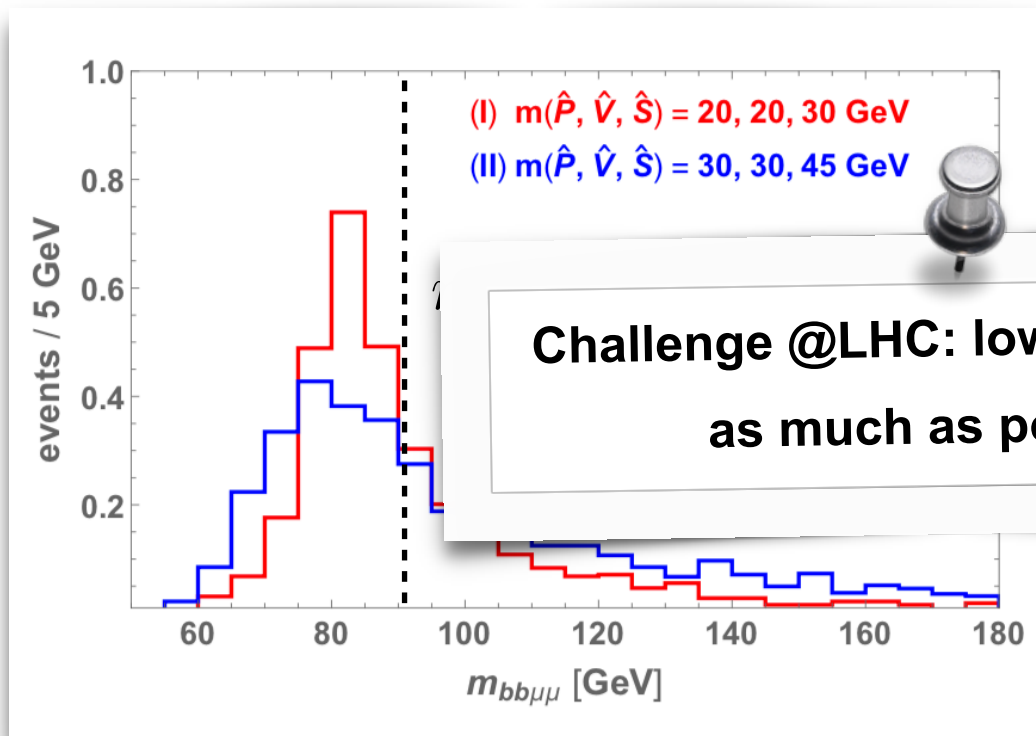
Phenomenology/1

- No dedicated LHC search yet, but can learn from $h \rightarrow aa \rightarrow (b\bar{b})(\mu\mu)$
- CMS analysis:

[CMS, 1812.06359]

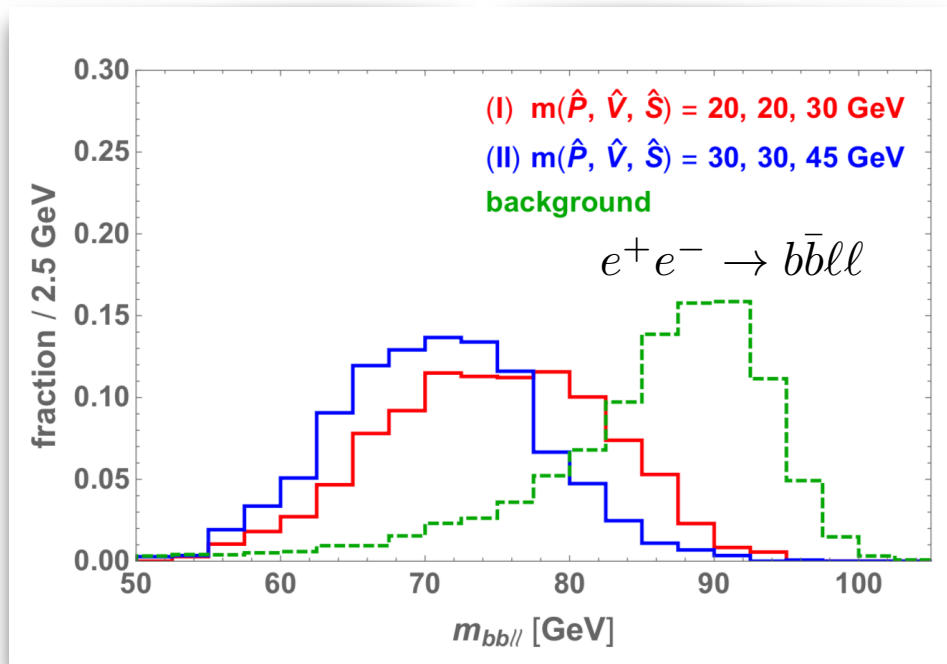
$$p_T^{\mu^{1,2}} > 20, 9 \text{ GeV}, \quad p_T^{b^{1,2}} > 20, 15 \text{ GeV}$$

- For Z signal, selection efficiency ~ 10 times smaller than for Higgs (!)



Phenomenology/1

- At GigaZ factory, background is negligible after cuts



$$|m_{\ell\ell} - m_{\hat{V}}| < 0.5 \text{ GeV}$$

$$m_{bb} \in [m_{\hat{P}} - 10 \text{ GeV}, m_{\hat{P}} + 5 \text{ GeV}]$$

$$m_{bbll} < 85 \text{ GeV}$$

$$N_B < 0.1 \text{ events}$$

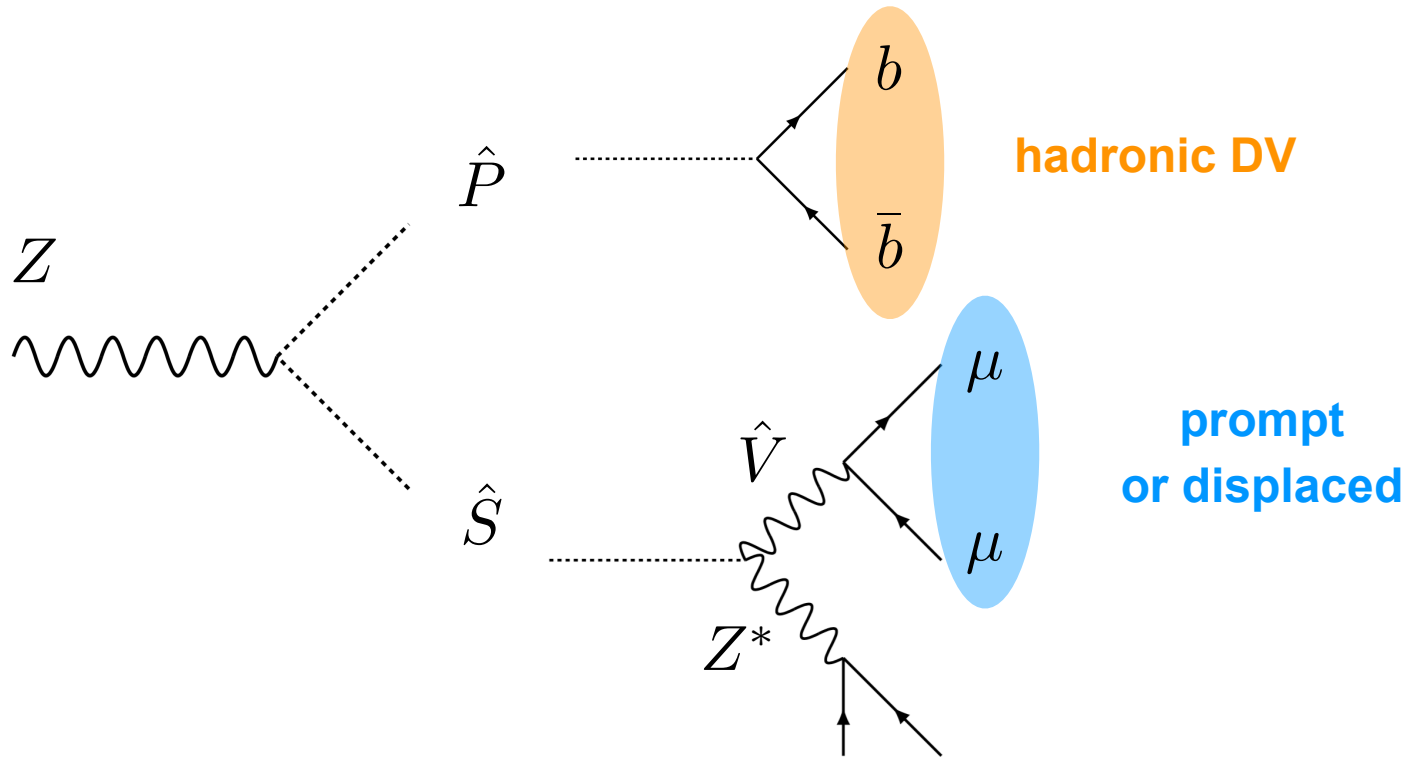
- Sensitivity on mediation scale:

$$(I) \quad M \gtrsim 5.4 \text{ TeV} \left(\frac{\mathcal{B}_{\hat{P}\hat{S}}}{1} \right)^{1/4}$$

$$(II) \quad M \gtrsim 5.2 \text{ TeV} \left(\frac{\mathcal{B}_{\hat{P}\hat{S}}}{1} \right)^{1/4} \quad (\text{GigaZ})$$

Phenomenology/1

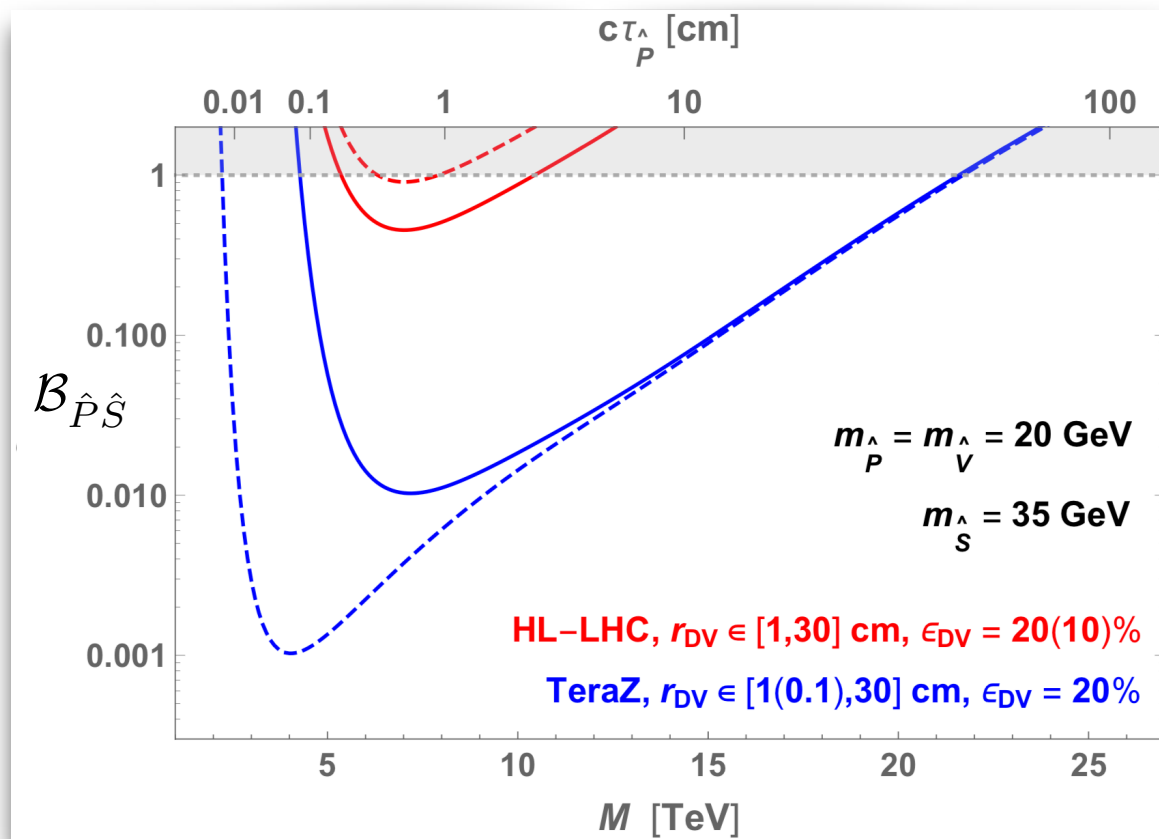
- For larger $M \gtrsim 5 \text{ TeV}$, at least some of the mesons become **long-lived**



$$\text{BR}(Z \rightarrow \bar{\psi}\psi) \sim M^{-4}, \quad c\tau_{\text{meson}} \sim M^4$$

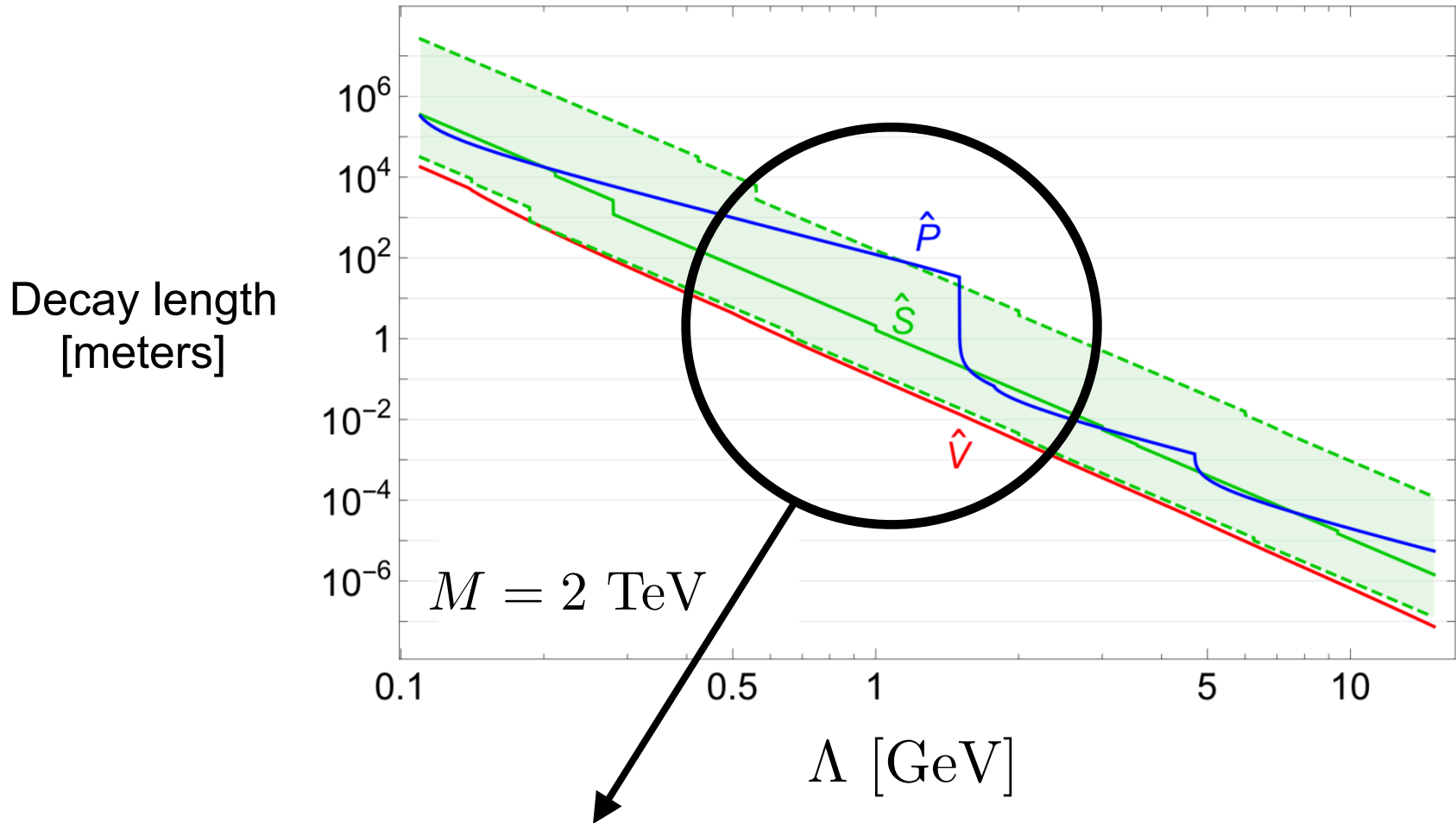
Phenomenology/1

- For larger $M \gtrsim 5$ TeV, at least some of the mesons become **long-lived**



At TeraZ, reach up to $M \sim 20$ TeV

Phenomenology/2



$\Lambda \lesssim 2\text{-}3 \text{ GeV} : Z \rightarrow \text{hidden jets, long-lived mesons}$

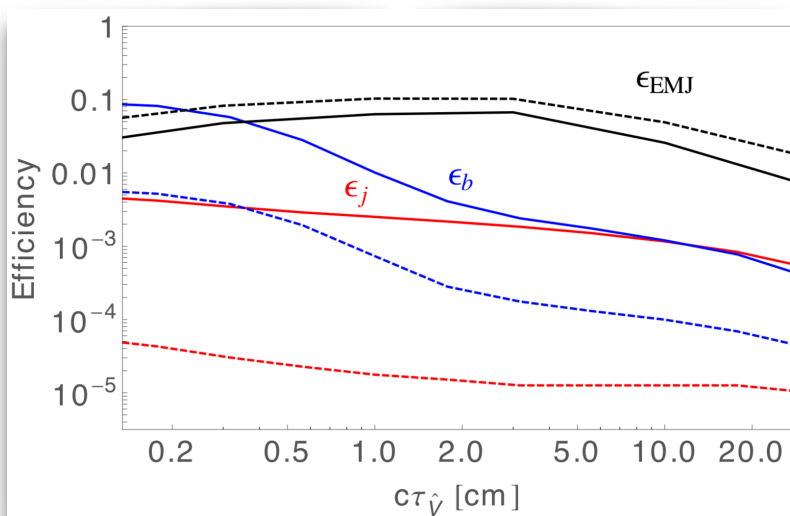
Phenomenology/2

- Hidden jets made of long-lived mesons
- Similarities with emerging/semi-visible jets, **but** relatively soft production mode and democratic $\hat{V} \rightarrow f_{\text{SM}} \bar{f}_{\text{SM}}$
[Schwaller, Stolarski, Weiler, 1502.05409]
[Cohen, Lisanti, Lou, 1503.00009]

- **LHC:** ATLAS and CMS searches not sensitive, **best reach @ LHCb**
Resolve single $\hat{V} \rightarrow \mu\mu$ decay inside the VELO

- **Z factory:** use LHC tracker resolution for conservative estimate
(checked it against CMS search for emerging jets) [CMS 1810.10069]

Phenomenology/2



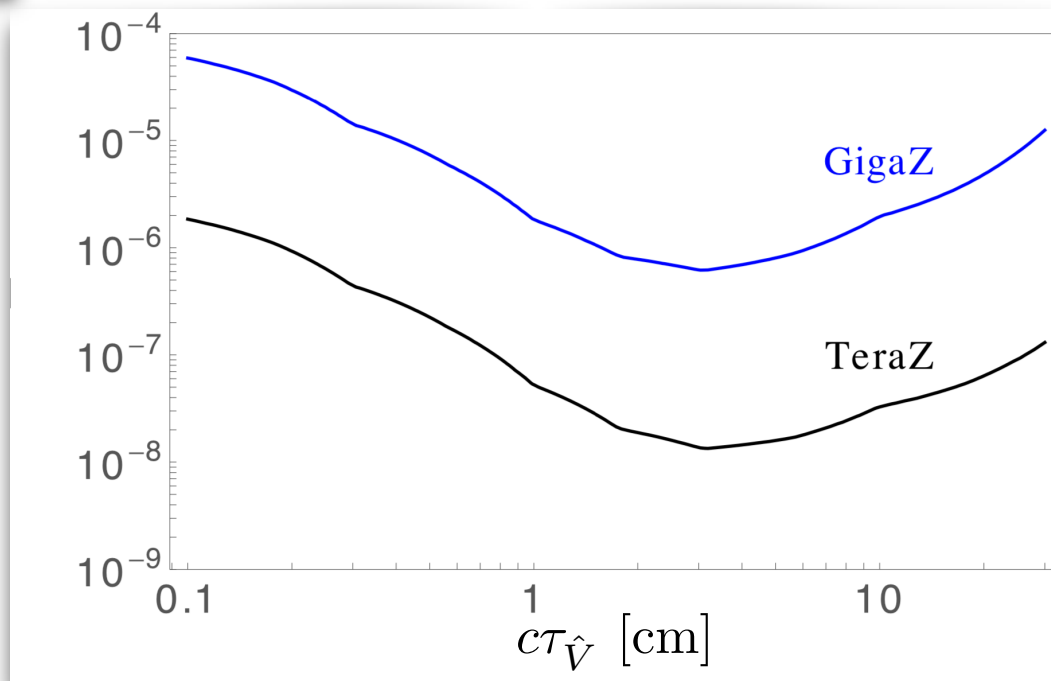
Solid: $N_{\text{tracks}} \geq 5, N_\ell = 0, 1$

Dashed: $N_{\text{tracks}} \geq 3, N_\ell \geq 2$

(leverage $\hat{V} \rightarrow \ell\ell$ decays)

$$\text{BR}(Z \rightarrow \bar{\psi}\psi)$$

[Cheng, Li, Salvioni, Verhaaren 2019]



Message

- Neutral Naturalness motivates a **confining hidden sector** at $0.1 \lesssim \Lambda/\text{GeV} \lesssim 10$
- Here: 1-flavor hidden QCD, coupled to SM via Higgs **and Z portal**
- Indirect sensitivity via EWPO, but **much stronger direct potential @FCC-ee**
A trillion Z bosons give extraordinary reach on exotic decays: $M \sim 20 \text{ TeV}$
- More to learn: multi-flavor case in progress
[Cheng, Li, Salvioni, Verhaaren]
- For Z decays to hidden sectors that contain Dark Matter, extensive analysis done in [Liu, Wang, Wang, Xue, 1712.07237]

Backup

Phenomenology/2

- For $m_{\hat{V}} \lesssim \text{GeV}$, other production mechanisms become important: Bremsstrahlung, meson decays, Drell-Yan, ...
- Hidden fermions have no electric charge, vector meson couples to neutral current

$$\mathcal{L} = -A_D^\mu \left(\varepsilon e J_\mu^{\text{EM}} + \varepsilon_Z \frac{g_Z}{2} J_\mu^{\text{NC}} \right)$$

- No global analysis available yet.

To gain first impression, approximate \hat{V} as “kinetically mixed dark photon” with

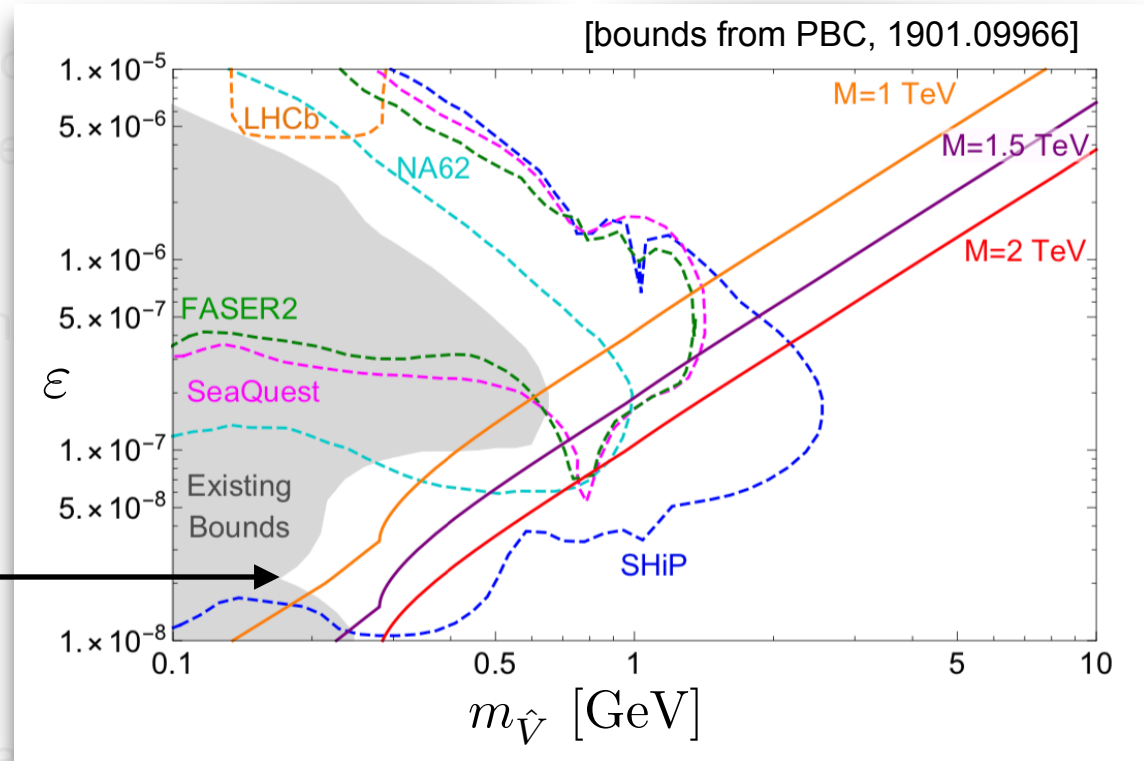
$$\varepsilon \approx 10^{-7} \left(\frac{m_{\hat{V}}}{\text{GeV}} \right)^2 \left(\frac{2 \text{ TeV}}{M} \right)^2$$

Phenomenology/2

- For $m_{\hat{V}} \lesssim \text{GeV}$, ϵ is constrained by Bremsstrahlung, meson decays

- Constituents have no direct current

- No global analysis available



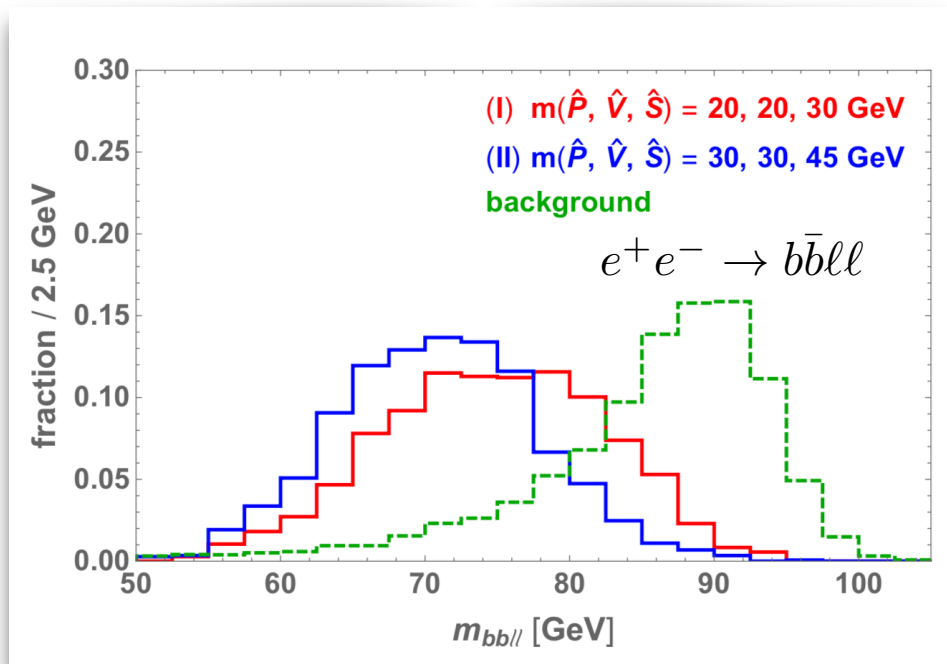
effect of
 $\text{BR}(\hat{V} \rightarrow \nu\nu)$

To gain first impression, approximate \hat{V} as “kinetically mixed dark photon” with

$$\epsilon \approx 10^{-7} \left(\frac{m_{\hat{V}}}{\text{GeV}} \right)^2 \left(\frac{2 \text{ TeV}}{M} \right)^2$$

Phenomenology/1

- At GigaZ factory, background is negligible after cuts



$$|m_{\ell\ell} - m_{\hat{V}}| < 0.5 \text{ GeV}$$

$$m_{bb} \in [m_{\hat{P}} - 10 \text{ GeV}, m_{\hat{P}} + 5 \text{ GeV}]$$

$$m_{bbll} < 85 \text{ GeV}$$

$$N_B < 0.1 \text{ events}$$

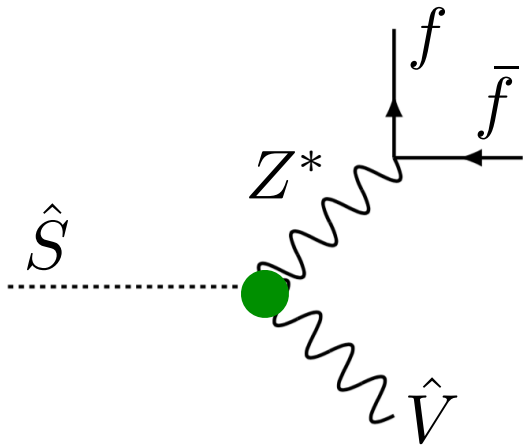
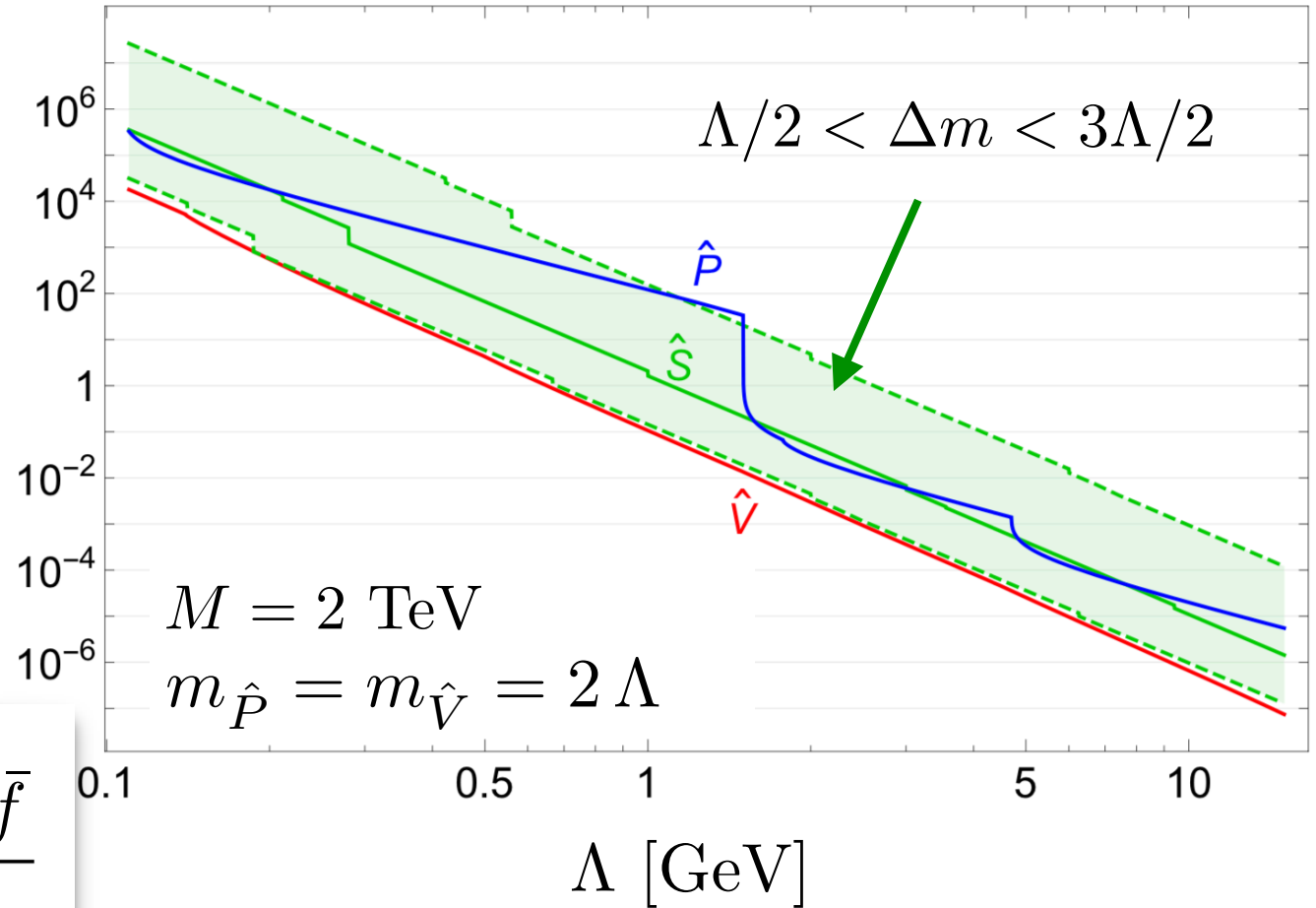
- **At LEP1 ?** $N_Z^{\text{total}} \approx 2.2 \times 10^7$

[hep-ex/0509008]

$$N_S \approx 28 \mathcal{B}_{\hat{P}\hat{S}} \left(\frac{2 \text{ TeV}}{M} \right)^4 (\mathcal{A}\epsilon)_{\text{tot}} \quad \text{but } \epsilon_b \sim 0.3, \text{ sensitivity marginal at best}$$

Hidden meson decays

Decay length
[meters]



$$\Gamma_{\hat{S}} \propto (\Delta m)^7$$

$$\Delta m \equiv m_{\hat{S}} - m_{\hat{V}}$$

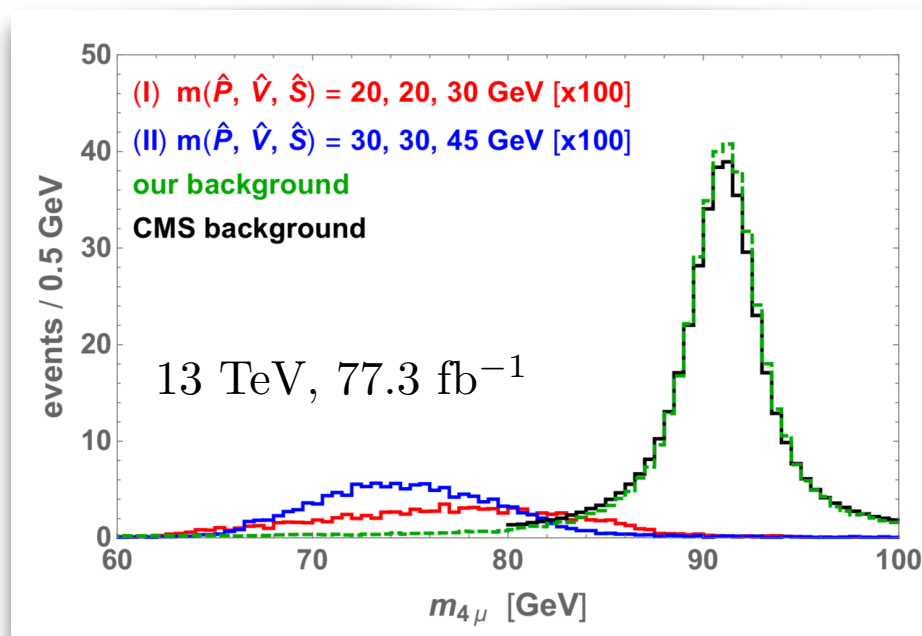
Phenomenology/1

- At LHC, very clean final state from $Z \rightarrow \hat{V}\hat{S} \rightarrow (\mu\mu)(\mu\mu) + X$
- Start from CMS search for light Z' , but implement our own analysis

[CMS, 1808.03684]

$$|m(\mu_1^+ \mu_1^-) - m(\mu_1^+ \mu_1^-)| \text{ or}$$

$$|m(\mu_1^+ \mu_2^-) - m(\mu_2^+ \mu_1^-)| < 1 \text{ GeV}$$



now

$$(I) \quad M \gtrsim 1.5, 2.0, 3.3 \text{ TeV} \left(\frac{\mathcal{B}_{\hat{V}\hat{S}}}{0.1} \right)^{1/4}$$

HL-LHC

Meson decays

$$\Gamma(\hat{V} \rightarrow f\bar{f}) = N_d N_c^f \frac{\pi \alpha_Z^2}{12} \frac{m_t^4}{M^4} \frac{m_{\hat{V}}^2 |\psi(0)|^2}{m_Z^4} \frac{\left(1 - \frac{4m_f^2}{m_{\hat{V}}^2}\right)^{1/2}}{\left(1 - \frac{m_{\hat{V}}^2}{m_Z^2}\right)^2} \left[v_f^2 \left(1 + \frac{2m_f^2}{m_{\hat{V}}^2}\right) + a_f^2 \left(1 - \frac{4m_f^2}{m_{\hat{V}}^2}\right) \right],$$

$$\Gamma(\hat{P} \rightarrow f\bar{f}) = N_d N_c(f) 2\pi \alpha_Z^2 \frac{m_t^4}{M^4} a_f^2 \frac{\mu_\psi^2 m_f^2}{m_Z^4} \frac{|\psi(0)|^2}{m_{\hat{P}}^2} \left(1 - \frac{4m_f^2}{m_{\hat{P}}^2}\right)^{1/2},$$

$$\Gamma(\hat{S} \rightarrow f\bar{f}) = \frac{18 N_d N_c^f}{\pi} (\lambda_{h\psi\psi} \lambda_{hff})^2 \frac{|\psi'(0)|^2}{m_h^4} \frac{\left(1 - \frac{4m_f^2}{m_{\hat{S}}^2}\right)^{3/2}}{\left(1 - \frac{m_{\hat{S}}^2}{m_h^2}\right)^2} \quad \left(|\psi(0)|^2 = \frac{\Lambda^3}{4\pi}\right)$$

$$\lambda_{h\psi\psi} = 2c_g \mu_\psi m_t^2 / (3bvM^2) \text{ where } b = 11 - 2N_l/3$$

$$\Gamma(\hat{S} \rightarrow \hat{V} f\bar{f}) \sim \frac{\alpha_Z^2 N_f}{16\pi} \frac{m_t^4}{M^4} \frac{k^7}{m_Z^4} |\varepsilon_{if}|^2 \quad k = \frac{m_{\hat{S}}^2 - m_{\hat{V}}^2}{2m_{\hat{S}}} = \Delta m \left(1 - \frac{\Delta m}{2m_{\hat{S}}}\right)$$

$$\frac{\Gamma(\hat{S} \rightarrow \bar{b}b)}{\Gamma(\hat{S} \rightarrow \hat{V} f\bar{f})} \sim \frac{c_g^2}{b^2} \frac{8N_d N_c}{\pi N_f} \frac{y_t^2 y_b^2}{\alpha_Z^2} \frac{m_t^4}{m_h^4} \frac{\mu_\psi^2}{m_t^2} \frac{\Lambda^7}{k^7} \approx 10^{-5} \left(\frac{\Lambda}{5 \text{ GeV}}\right)^2 \left(\frac{\mu_\psi}{\Lambda}\right)^2 \left(\frac{\Lambda}{k}\right)^7 \left(\frac{c_g}{4}\right)^2,$$

Long-lived heavy mesons

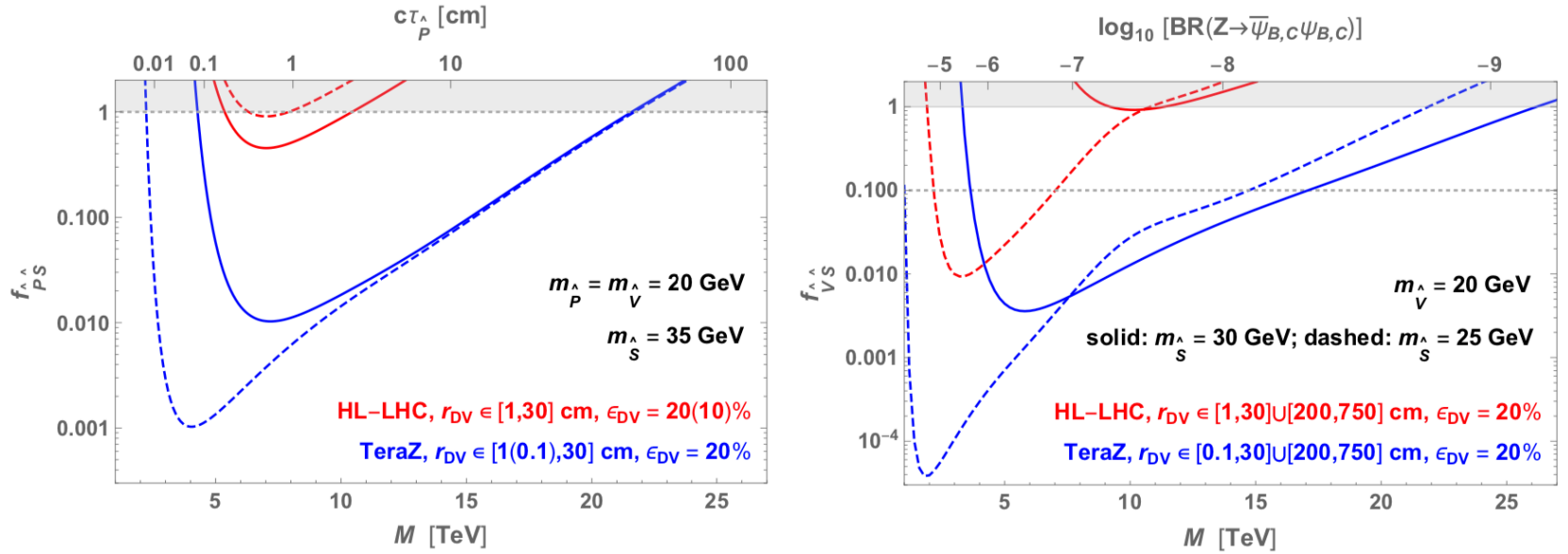


FIG. 5: *Left:* projected bounds on $f_{\hat{P}\hat{S}}$ from future searches for $Z \rightarrow \hat{P}\hat{S}$, where \hat{P} is long-lived while $\hat{S} \rightarrow (\hat{V} \rightarrow \ell\ell)f\bar{f}$. *Right:* projected bounds on $f_{\hat{V}\hat{S}}$ from future searches for $Z \rightarrow \hat{V}\hat{S}$, where \hat{S} is long lived while $\hat{V} \rightarrow \ell\ell$. All bounds are at 95% CL, assuming negligible SM background. We take $\ell = \mu$ (e or μ) at the LHC (TeraZ), and the dilepton pair must be prompt at the LHC, but can be displaced at TeraZ. The dotted gray lines correspond to the educated guesses $f_{\hat{P}\hat{S}, \hat{V}\hat{S}} \sim 1, 0.1$ discussed in Subsec. [IV A](#).

Meson FCNC decays

By matching to Eq. (27) we identify

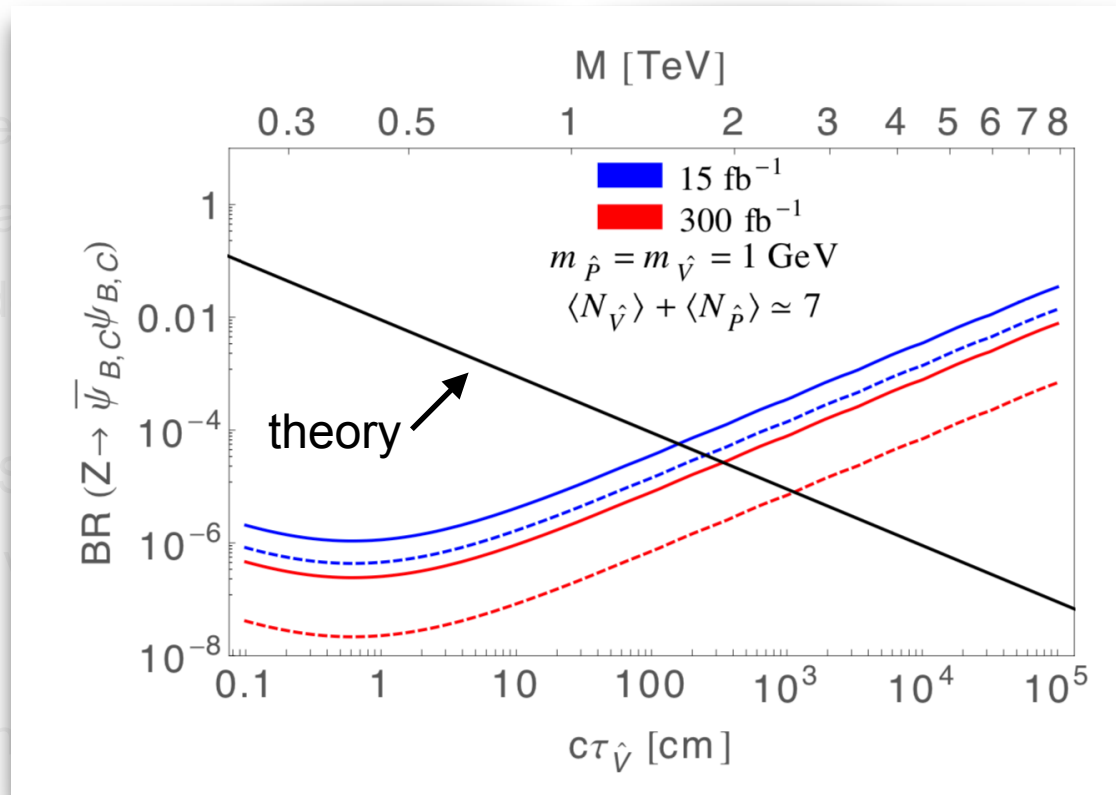
$$\varepsilon_Z \simeq g_Z \sqrt{\frac{N_d}{2}} \frac{m_t^2}{M^2} \frac{|\psi(0)| m_{\hat{V}}^{1/2}}{m_Z^2} \approx 3.2 \times 10^{-7} \left(\frac{\Lambda}{1 \text{ GeV}} \right)^{3/2} \left(\frac{m_{\hat{V}}}{2 \text{ GeV}} \right)^{1/2} \left(\frac{2 \text{ TeV}}{M} \right)^2, \quad (53)$$

where we have assumed $m_{\hat{V}} \ll m_Z$ and included an extra factor of $\sqrt{2}$ to account for the two hidden sectors B, C .¹⁸ Very recently, Ref. [67] derived strong FCNC decay bounds on ε_Z for a generic light vector X_μ . In our setup, though, the NP contribution to a given SM final state should be smaller than the corresponding perturbative rate, $\text{BR}(B \rightarrow K f \bar{f})_{\text{NP}} \lesssim \text{BR}(B \rightarrow K \psi \bar{\psi})$, where f is a SM fermion and we have taken B decays as example. This allows us to conservatively estimate

$$\frac{\text{BR}(B \rightarrow K f \bar{f})_{\text{NP}}}{\text{BR}(B \rightarrow K f \bar{f})_{\text{SM}}} \lesssim \left(\frac{m_t}{M} \right)^4 \frac{\text{BR}(\hat{V} \rightarrow \nu \bar{\nu})}{\text{BR}(\hat{V} \rightarrow f \bar{f})}, \quad (54)$$

where we assumed dominance of the Z penguin over the box amplitudes and used Eq. (2). An analogous expression applies to K decays. For any f , the RHS of Eq. (54) is below a percent if $M > 1 \text{ TeV}$, showing that the precision needed to probe the NP is well beyond the current one [48].¹⁹ We do not expect non-perturbative corrections to change this conclusion.

Phenomenology/2



soft production mode

$M > 1.6 (2.0) \text{ TeV}$
 for $L = 15 (300) \text{ fb}^{-1}$

[see also: future displaced dimuon trigger @ CMS, Gershtein and Knapen, 1907.00007]

Theory motivation

- Add **two** copies of MSSM top sector,

[Cheng, Li, **ES**, Verhaaren, 2019]

$$SU(3)_A \times SU(3)_B \times SU(3)_C \times SU(2) \times U(1)$$

Superpotential:

$$W = y_t(Q_A H u_A^c + Q_B H u_B^c + Q_C H u_C^c) + M(Q_B Q_B'^c + Q_C Q_C'^c)$$

Z_3 \nearrow Z_2

few TeV

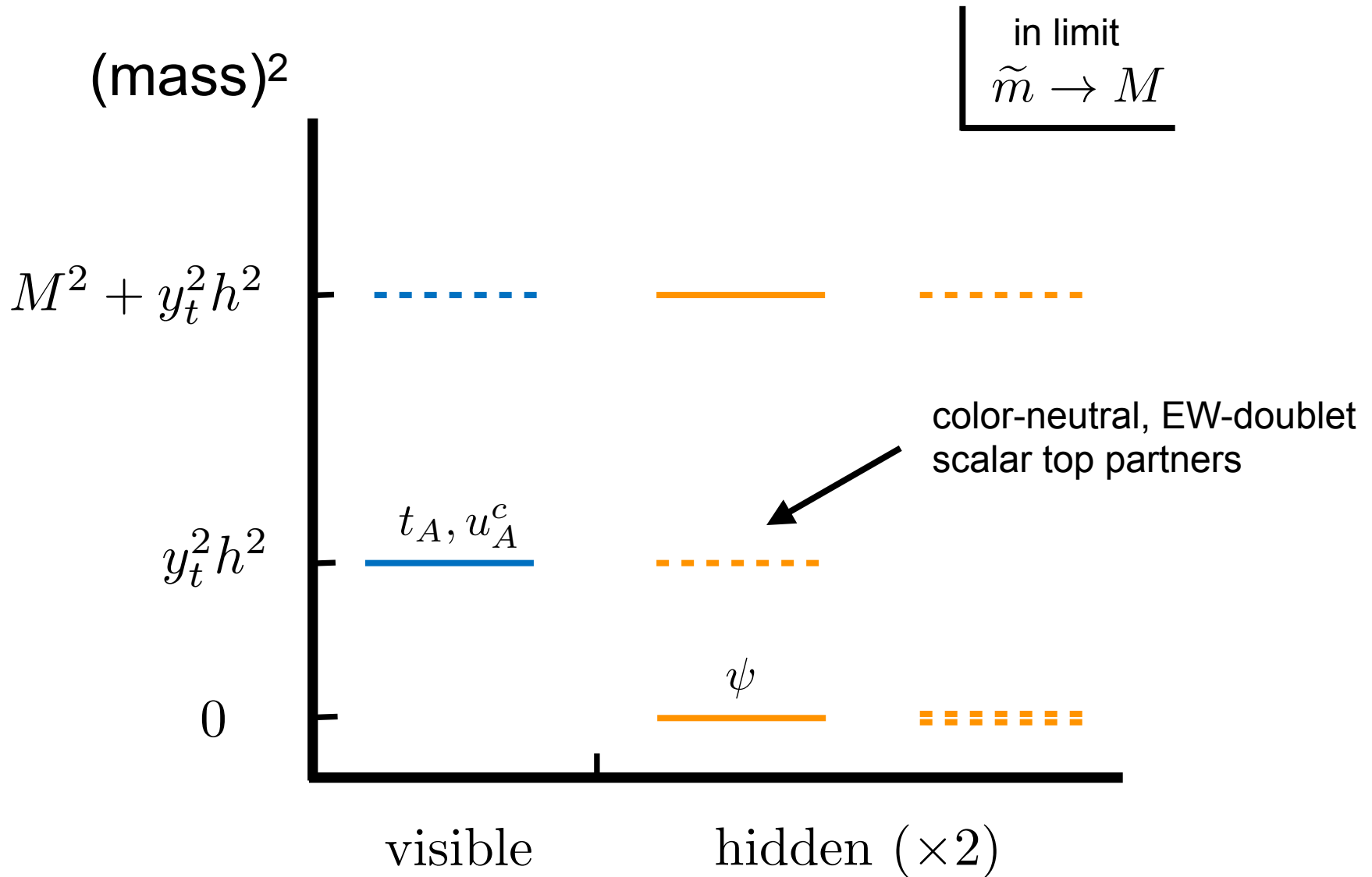
Soft masses:

$$V_s = +\tilde{m}^2(|\tilde{Q}_A|^2 + |\tilde{u}_A^c|^2) - \tilde{m}^2(|\tilde{Q}_B|^2 + |\tilde{Q}_C|^2)$$

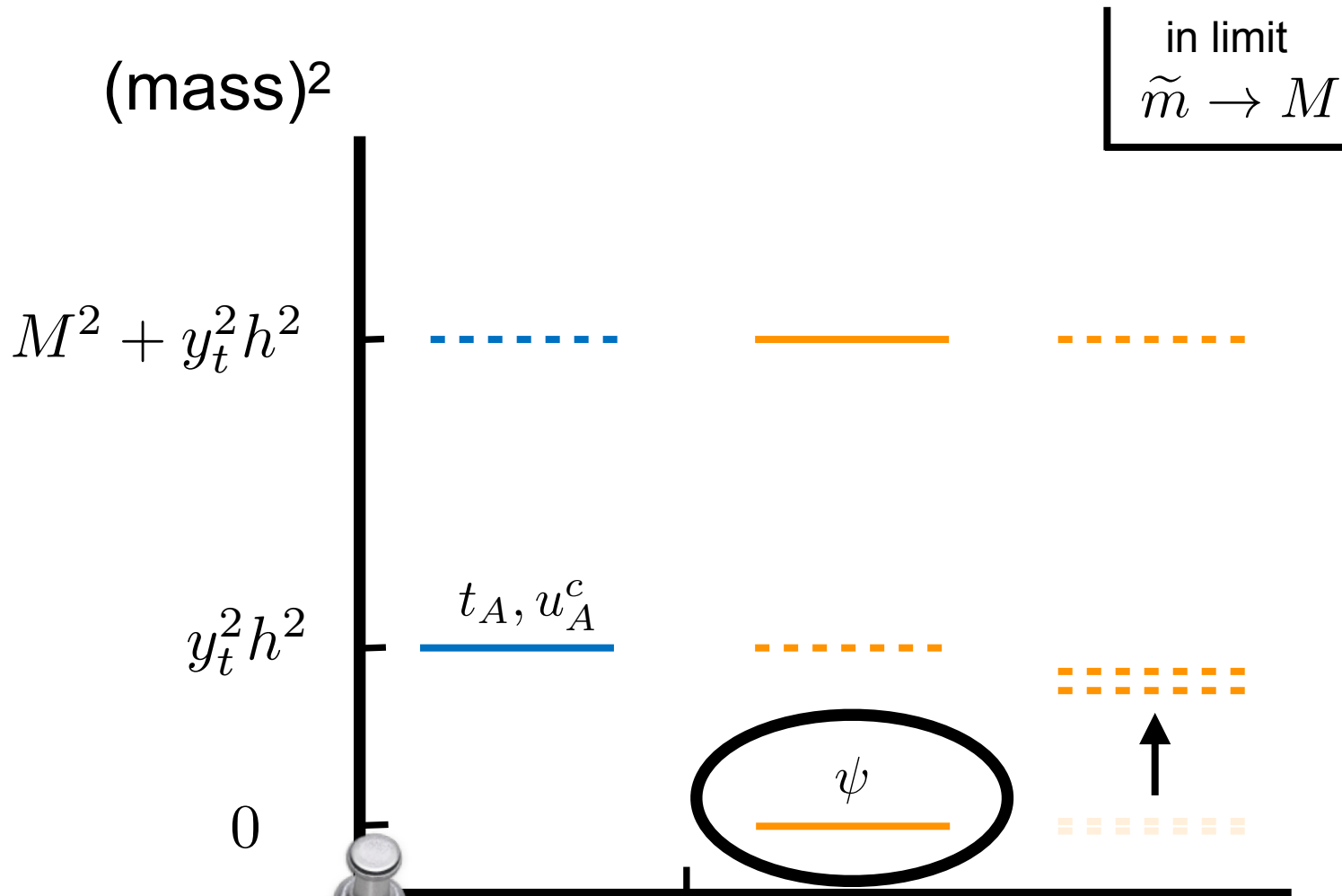
raise colored stops

lower color-neutral “stops”

Accidental supersymmetry



Accidental supersymmetry



A very light Dirac fermion
at bottom of spectrum

hidden ($\times 2$)