



Beam parameters and physics opportunities

- or why an upgraded EBIS and a buncher should be considered

J. Cederkall, Lund University

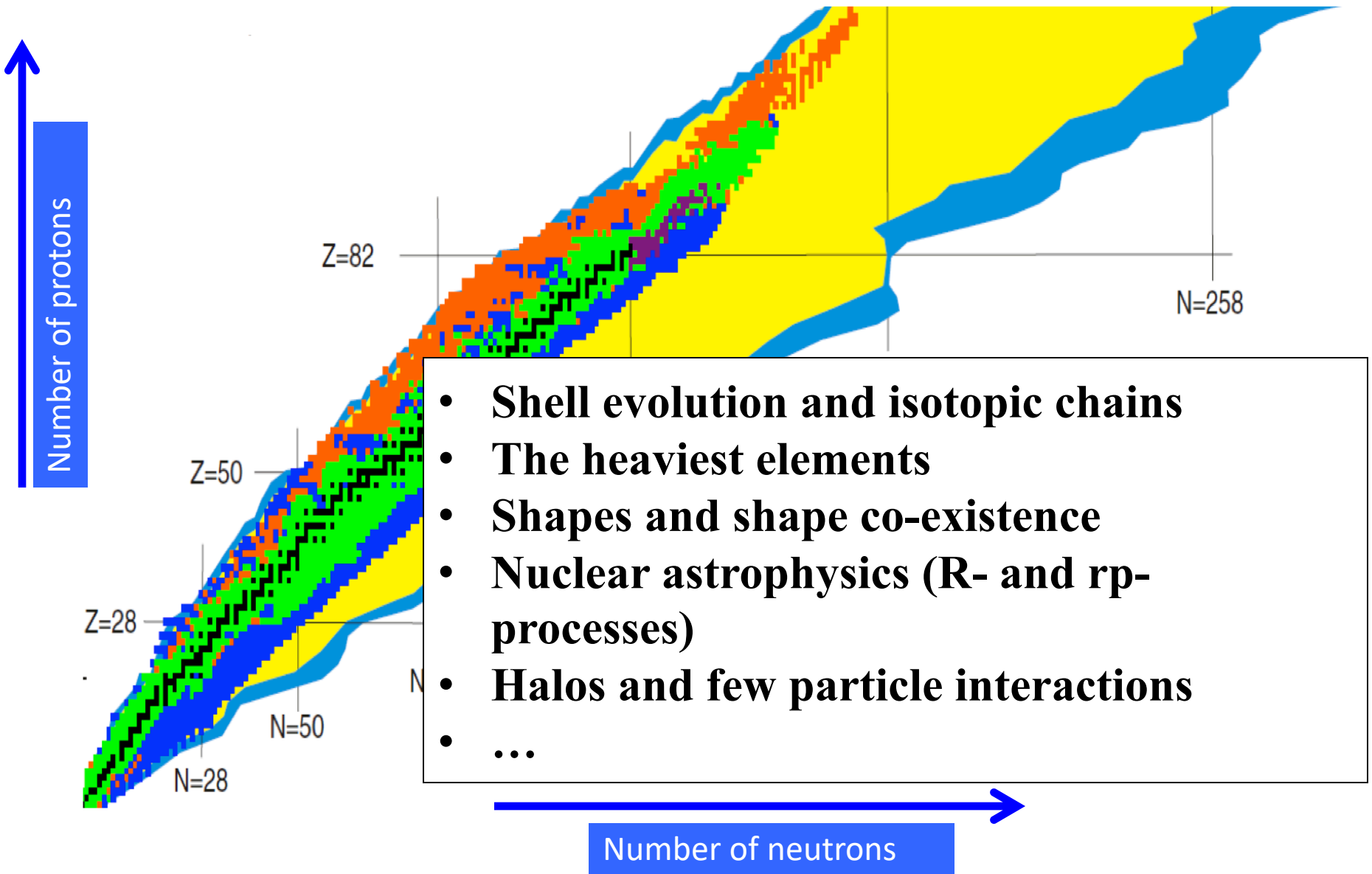


What has been achieved? What is next?

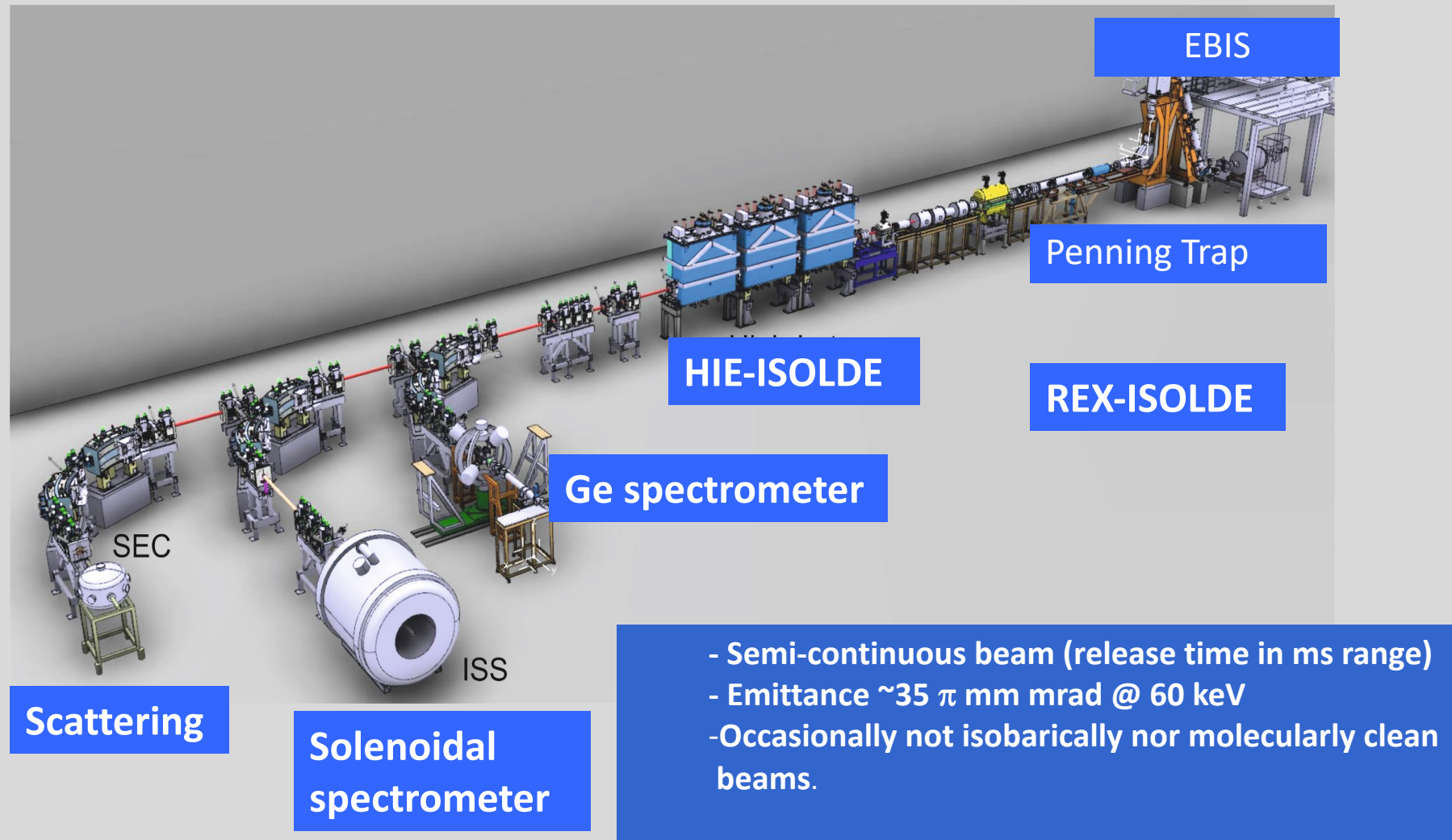
- Some goals for HIE-ISOLDE with attempt at a combined physics—and-machine view
- What is the current status?
- What can we do next?
- Which generic developments could/should we focus on?



Some physics for accelerated RIBs



REX-ISOLDE and HIE-ISOLDE (evolution...)



HIE-ISOLDE energy upgrade TDR, Y. Kadi (ed).

Table 1.1: Requested beam characteristics at HIE-ISOLDE.

Beam parameter	Description or value
Energy	continuous from < 0.7 to at least $10 \text{ MeV}/u$
Beam spot diameter	$< 1 - 3 \text{ mm FWHM}$
Beam divergence	$< 1 - 3 \text{ mrad FWHM}$
Micro-bunch structure ^a	no requirement of micro-bunching to bunch at $< 1 \text{ ns FWHM}$ with $\sim 100 \text{ ns}$ bunch spacing
Macro-bunch structure ^a	longer pulse lengths or cw operation
Energy spread	$< 0.1\%$
Absolute energy resolution ^b	no specific details given

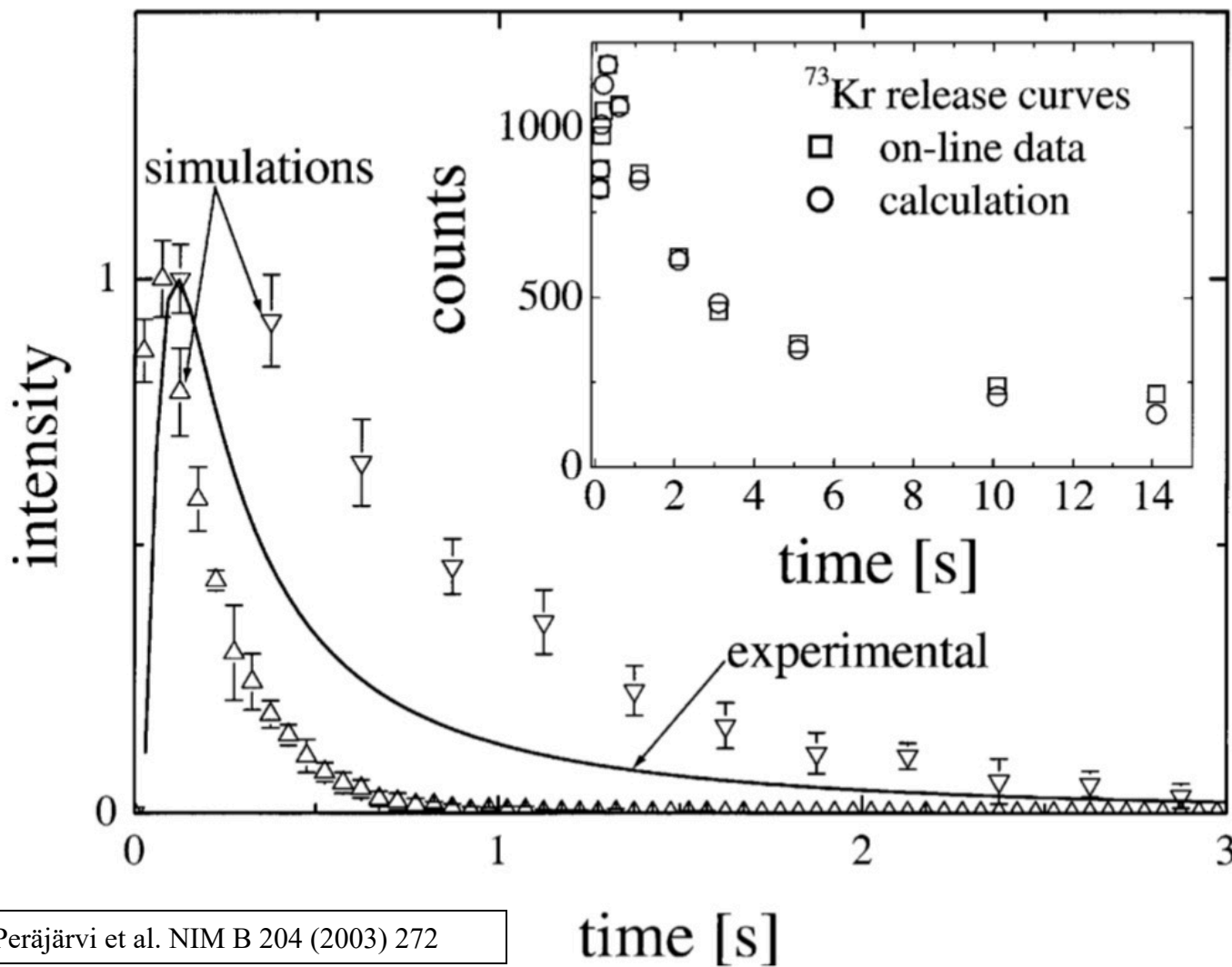
^a The time structure of the beam is determined by the charge breeder and the REX front end, see Figure 2.4.

^b No system currently exists to measure the absolute beam energy; time-of-flight systems are being discussed.

CERN-2018-002-M



Effusion, diffusion and decay



The combined time structure

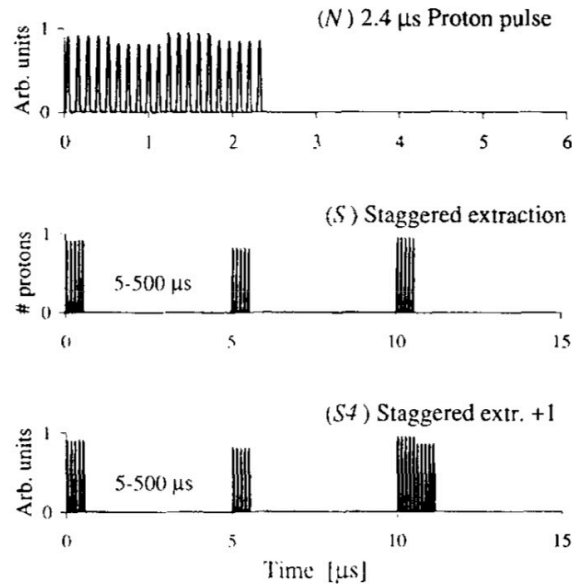


Fig. 1. The time structure of the PS-BOOSTER proton pulses is illustrated. A standard proton pulse (N) contains up to 3×10^{13} protons distributed in 20 bunches over 2.4 μ s. In the staggered mode (S) three groups of 5 bunches are extracted at time intervals ranging from 5 to 500 μ s (maximum pulse intensity: 2.2×10^{13} ppp). The maximum proton pulse intensity is obtained by immediate extraction of the last synchrotron (S4).

J. Lettry et al. NIM B 126 (1997) 170

Semi-continuous beams already from the primary target.

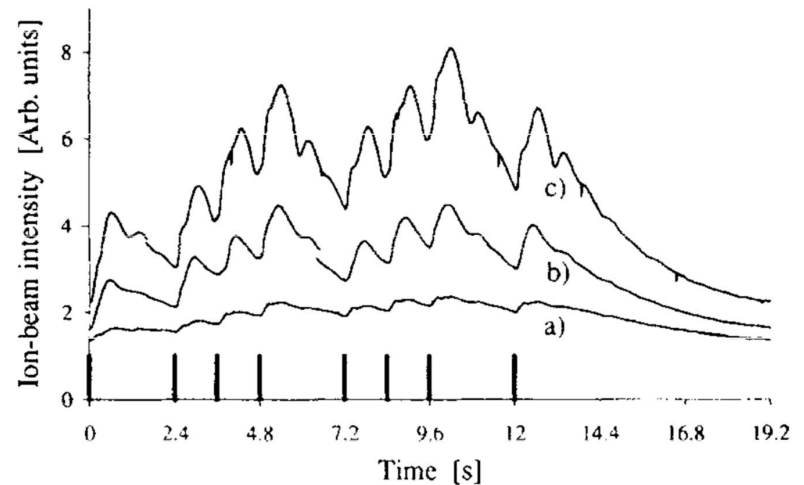


Fig. 6. The ionic current of ^{190}Hg as recorded with a 0.5 mm metallic needle is shown for different proton beam intensities in the staggered mode (a) 9×10^{12} , (b) 1.5×10^{13} and (c) 2.1×10^{13} proton per pulse. The parameters of the release function from which the release time is calculated were fitted to these data. The timing of the 8 proton pulses in the 19.2 s cycle is shown.



The pulse structure

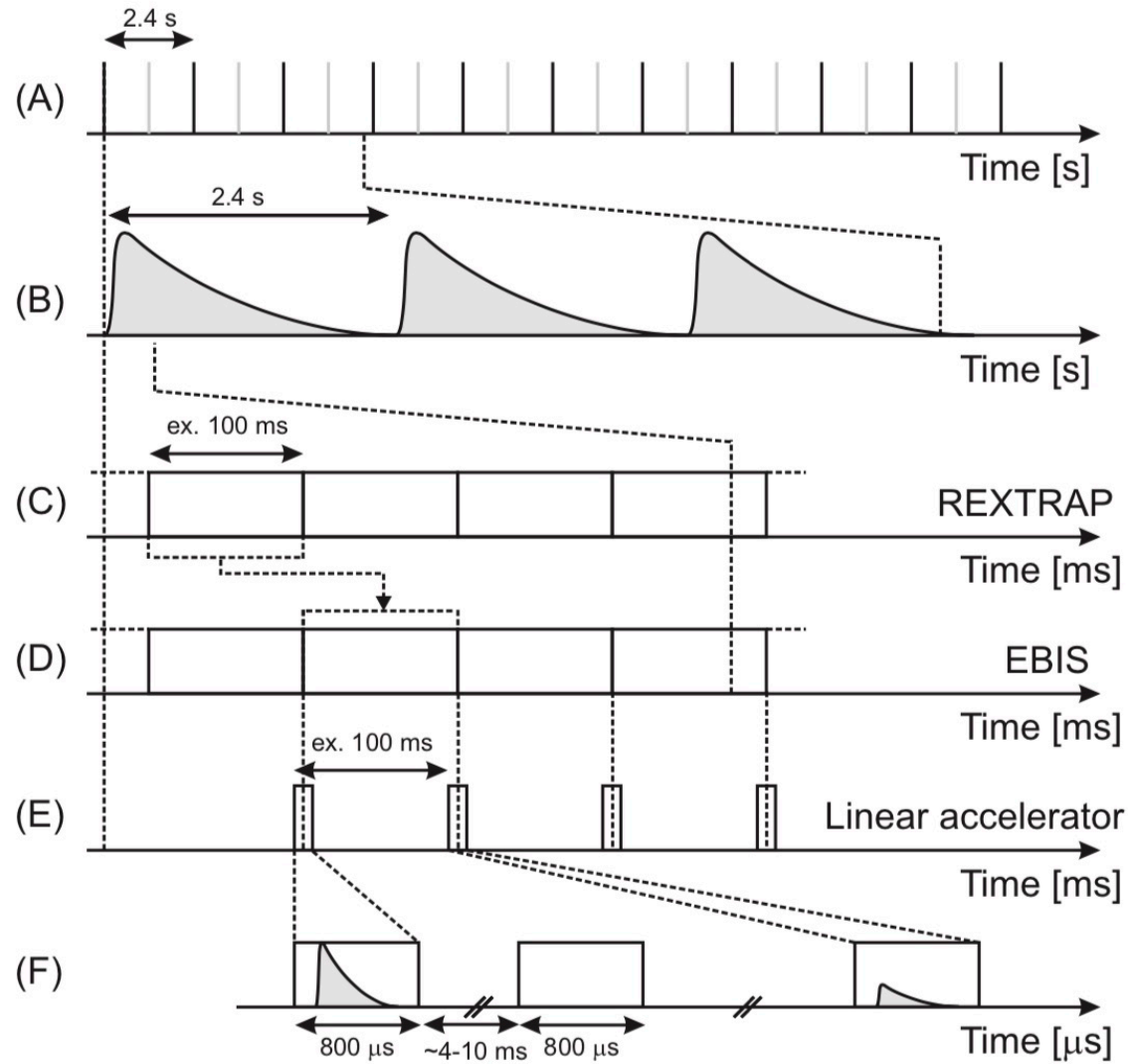


Fig. 2.4: REX beam time structure. Figure courtesy of J. van de Walle.



The pulse structure

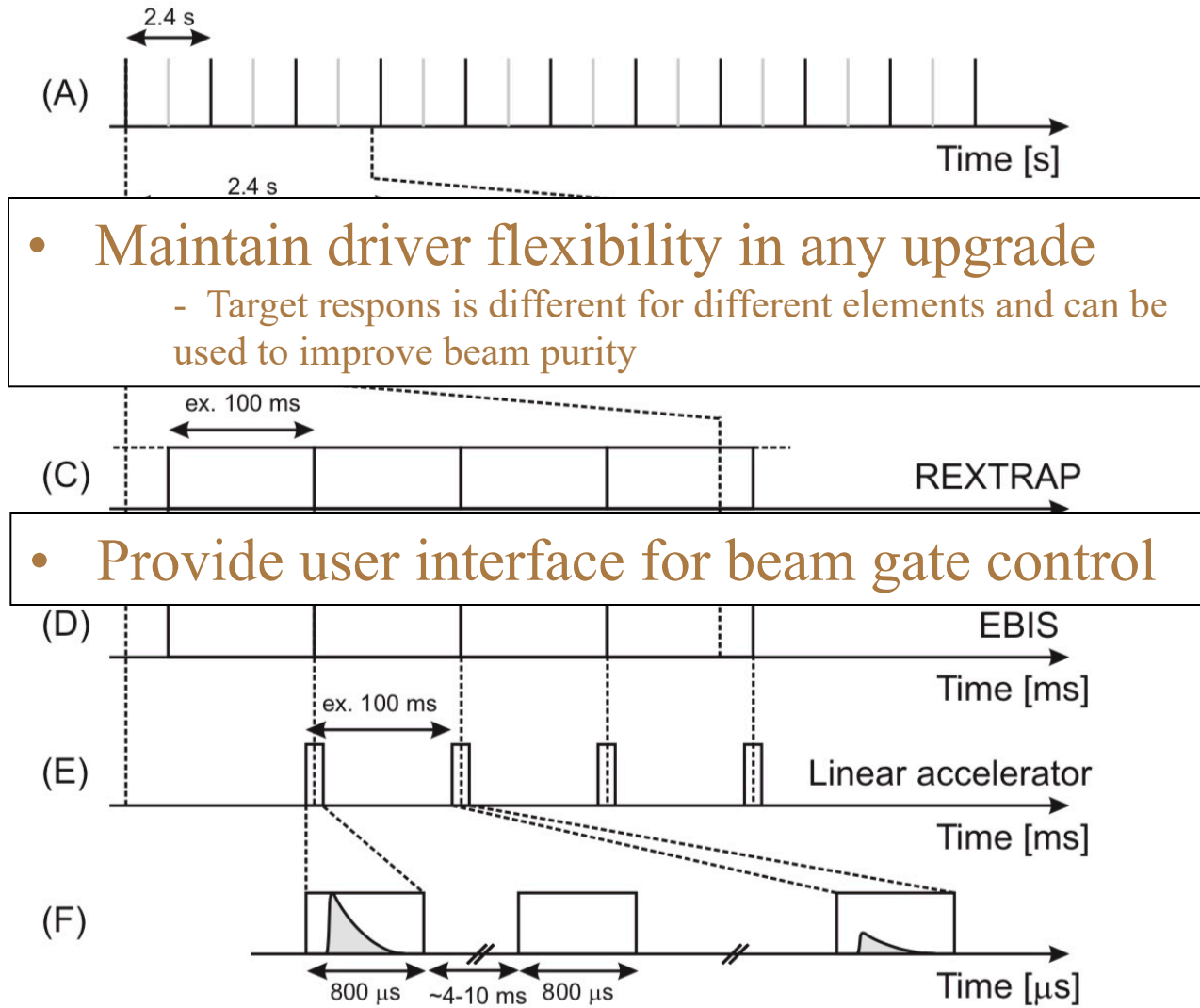
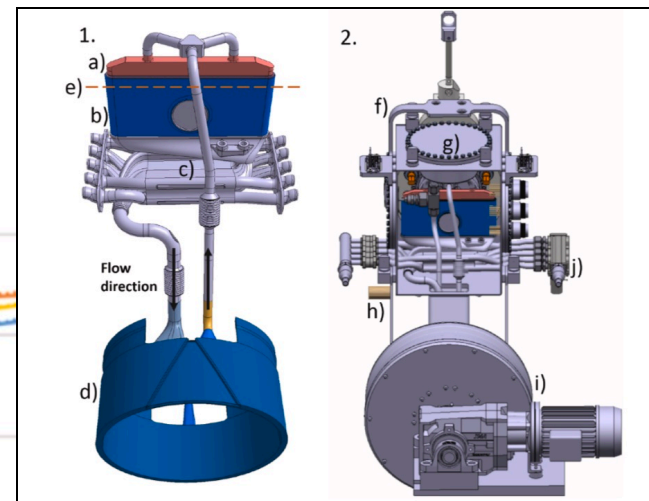
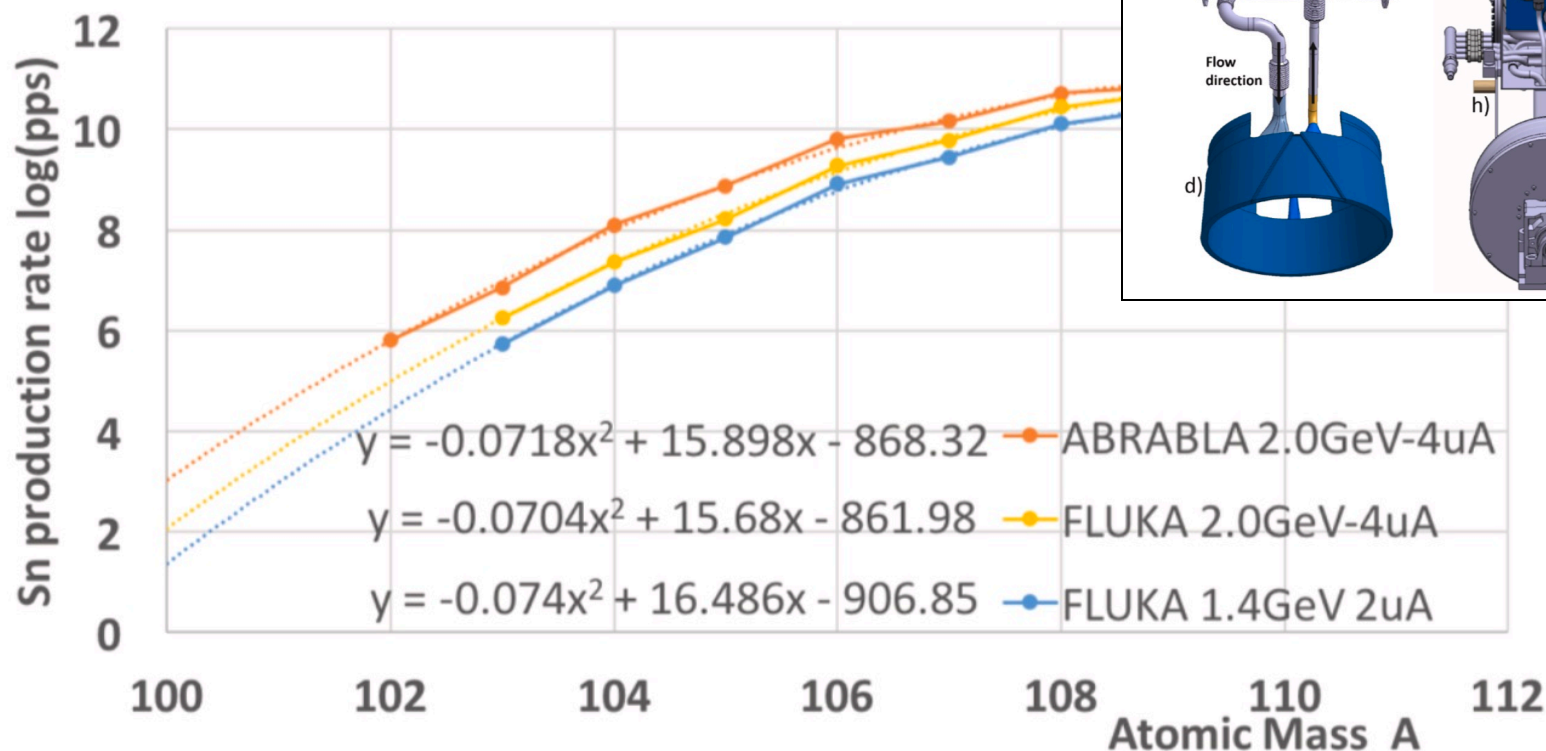


Fig. 2.4: REX beam time structure. Figure courtesy of J. van de Walle.



Target development for future improved yields?

LIEBE target

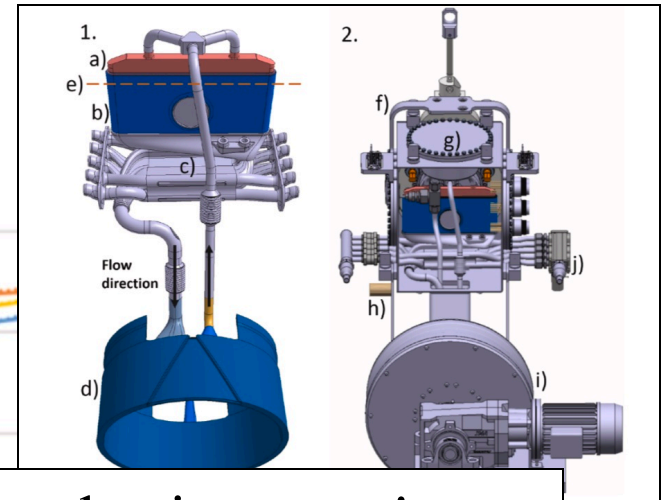
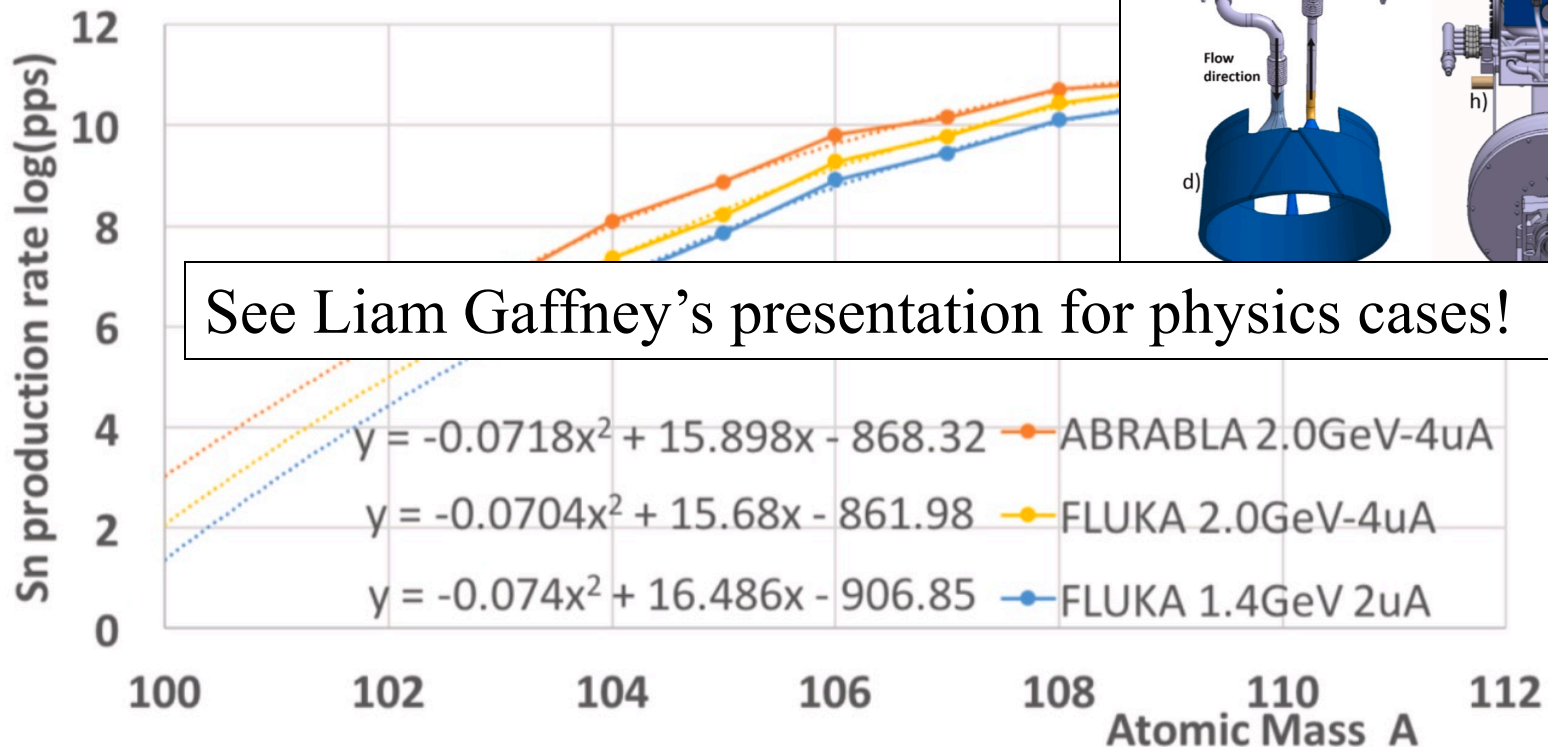


F.Boix Pamies, T.Stora, E.Barbero et al. NIM B <https://doi.org/10.1016/j.nimb.2019.06.043>

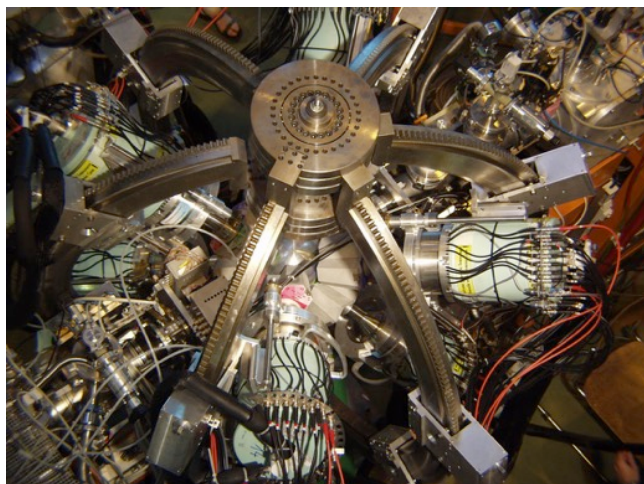
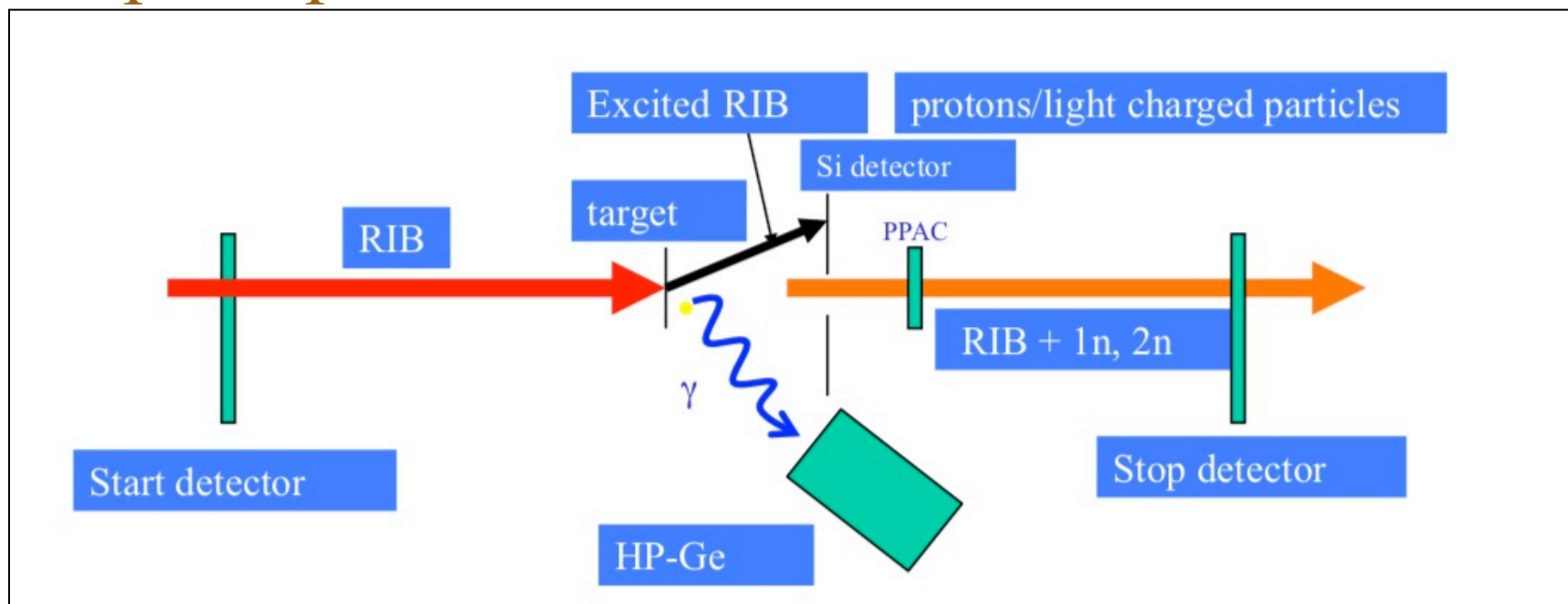


Future Improved Yields?

LIEBE target

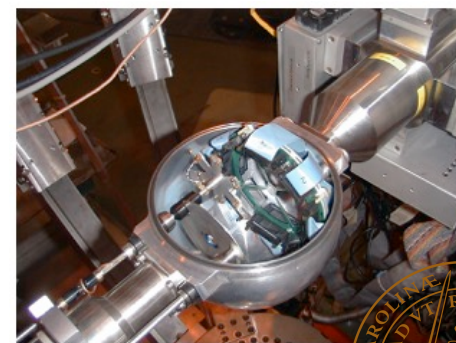


The principle

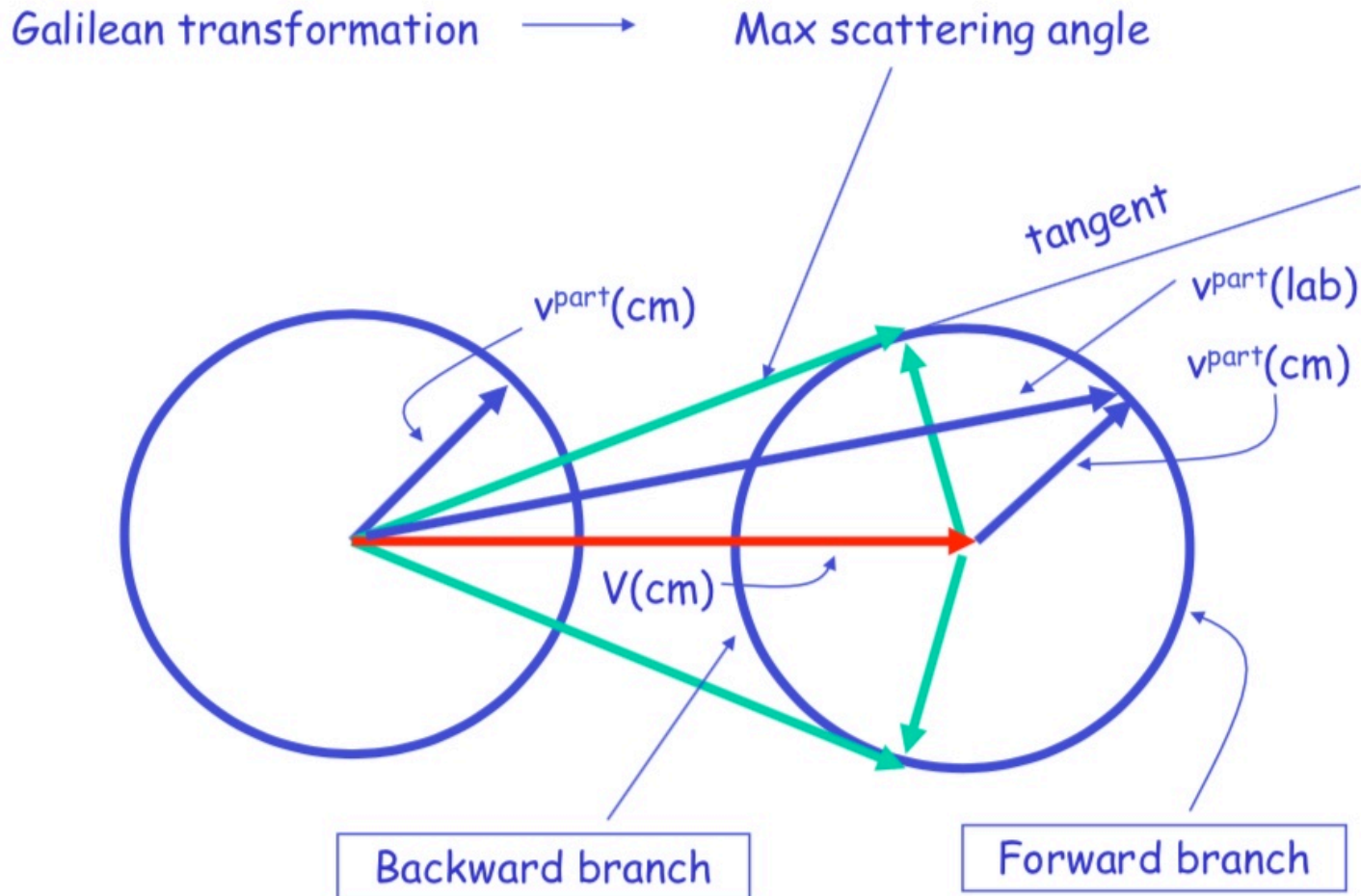


Worst case:
assume ca 10 mm

A geometric diagram showing a beam diverging from a source (green vertical bar) to a target (grey vertical bar). The distance between the source and target is labeled 'ca 30 mm'. The height of the target is labeled 'ca 15 mm'. The angle of divergence is labeled $\Delta\phi = (14 - 5) \text{ deg}$.

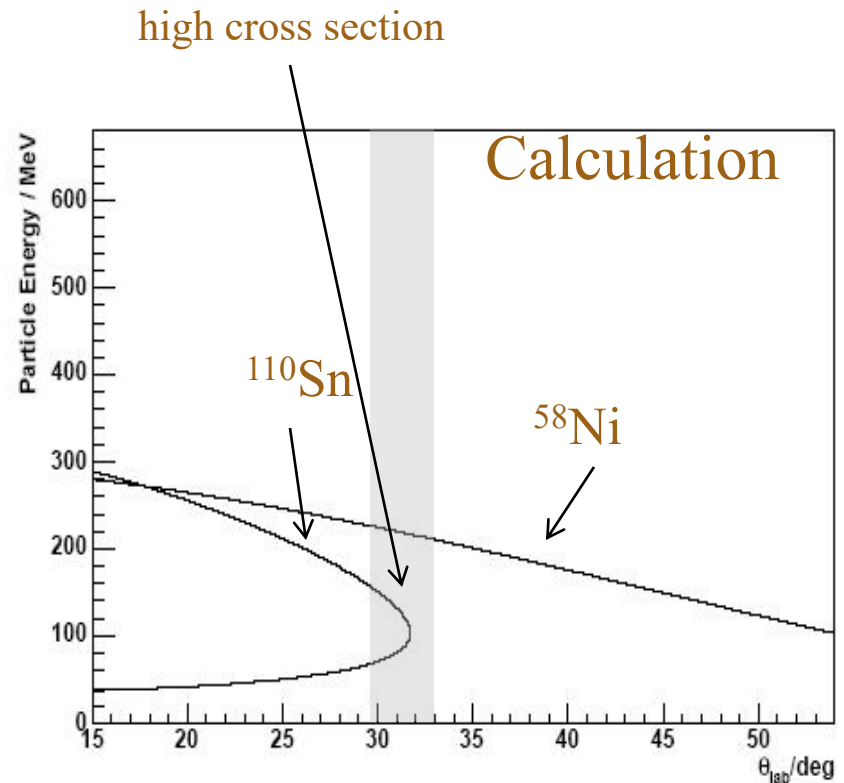
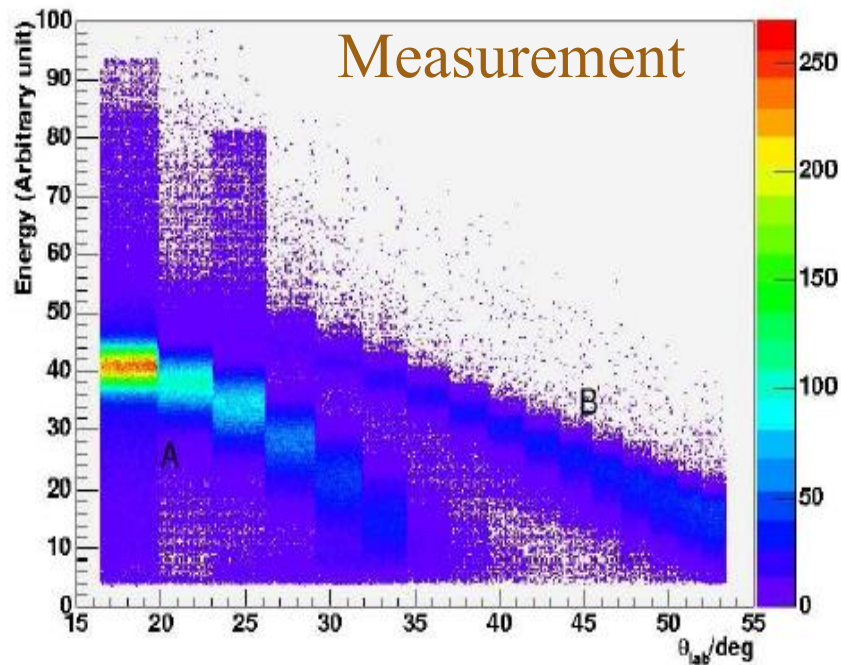


Kinematics



Kinematics I

$^{110,108,106}\text{Sn}$ on ^{58}Ni at 2.8 MeV/u

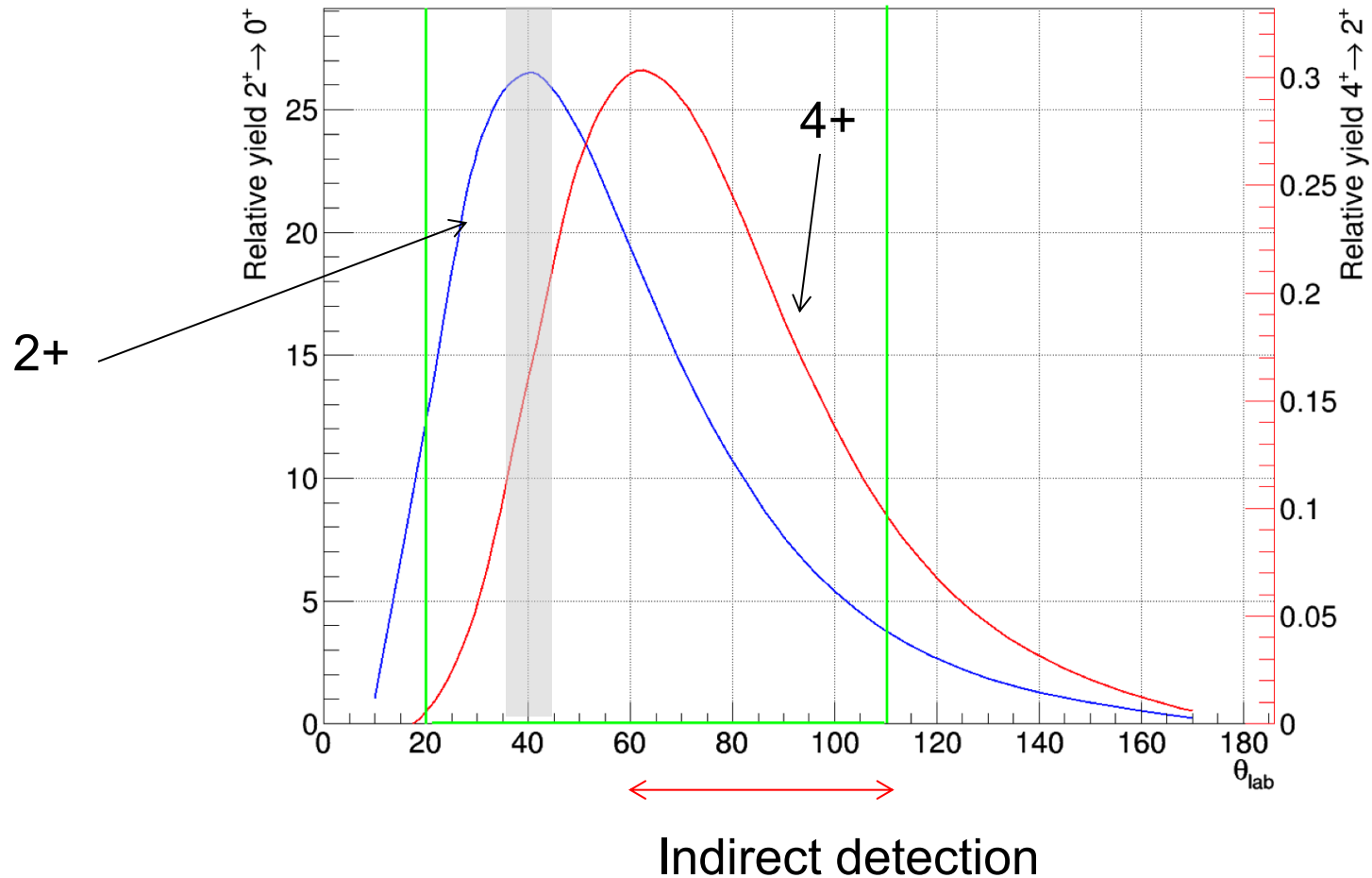


Energy in lab vs scattering angle



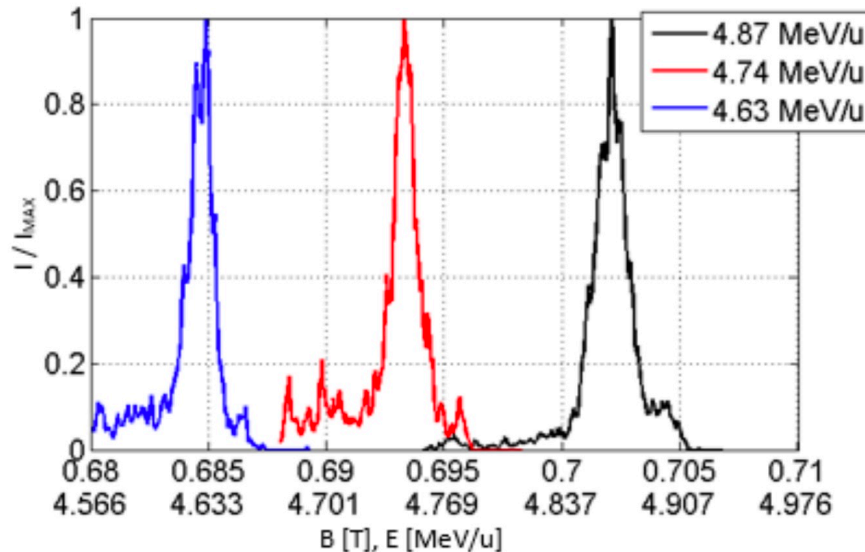
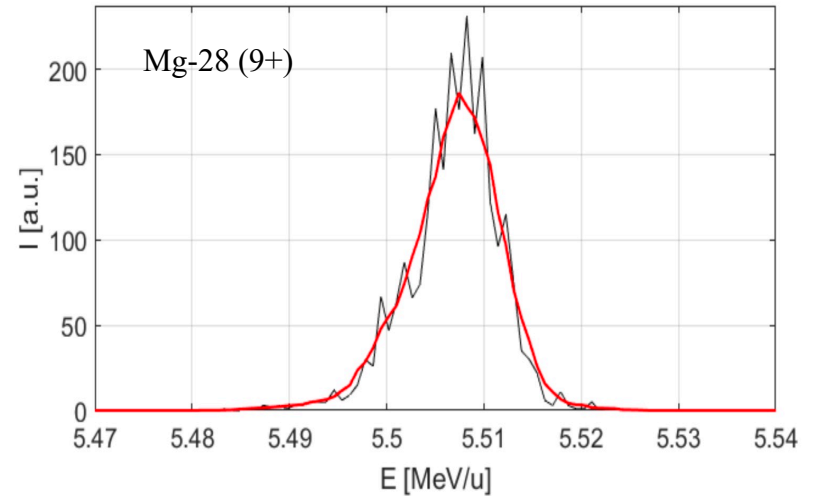
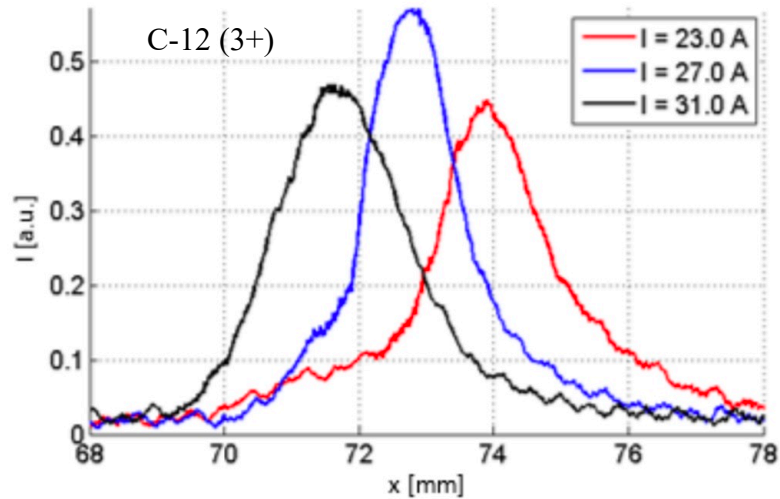
Cross section for a 2^+ and 4^+ excitation

Sn-110 and Pb-206



Energy spread and beam spot size

HIE-ISOLDE TDR CERN-2018-002-M Y. Kadi, M. A. Fraser A. Papageorgiou-Koufidou (eds.)



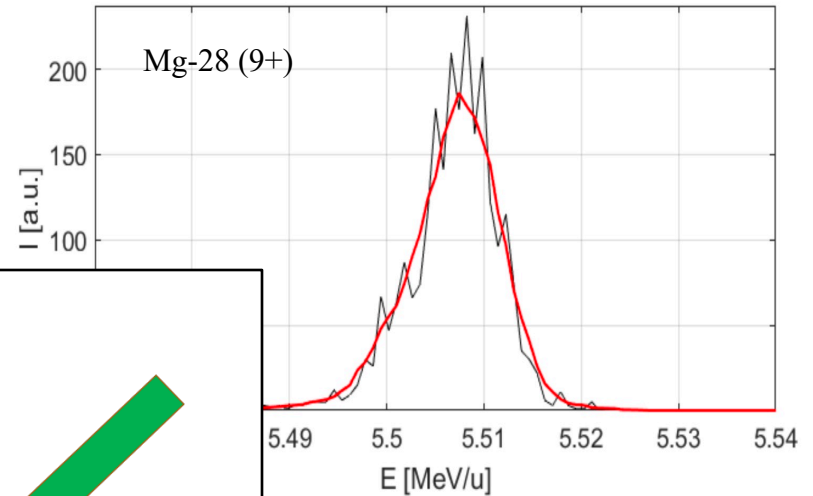
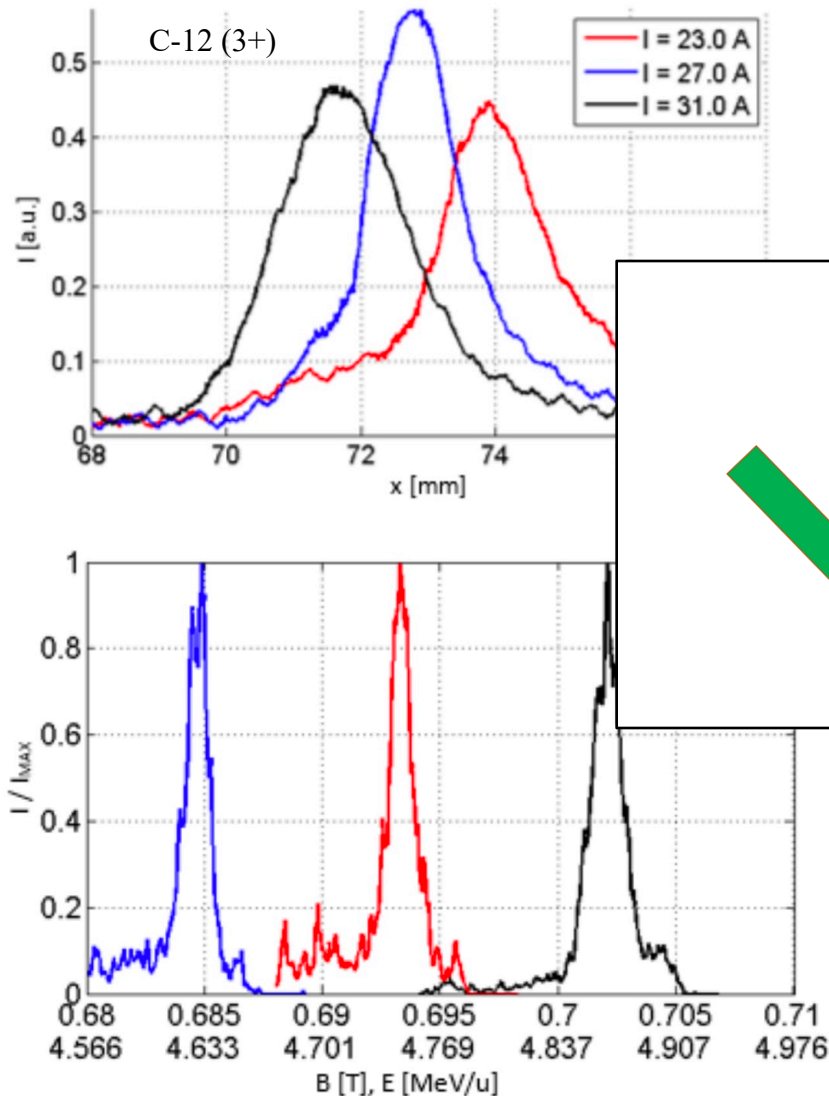
	Nominal Energy [MeV/u]	Measured Energy [MeV/u]	Energy Spread [%]
66Ni16+	4.5	4.47	0.2
9Li3+	6.9	6.72	0.5
132Sn31+	5.5	5.49	0.4
78Zn20+	4.3	4.27	0.3

J. Rodríguez et al. priv comm.



Energy spread and beam spot size

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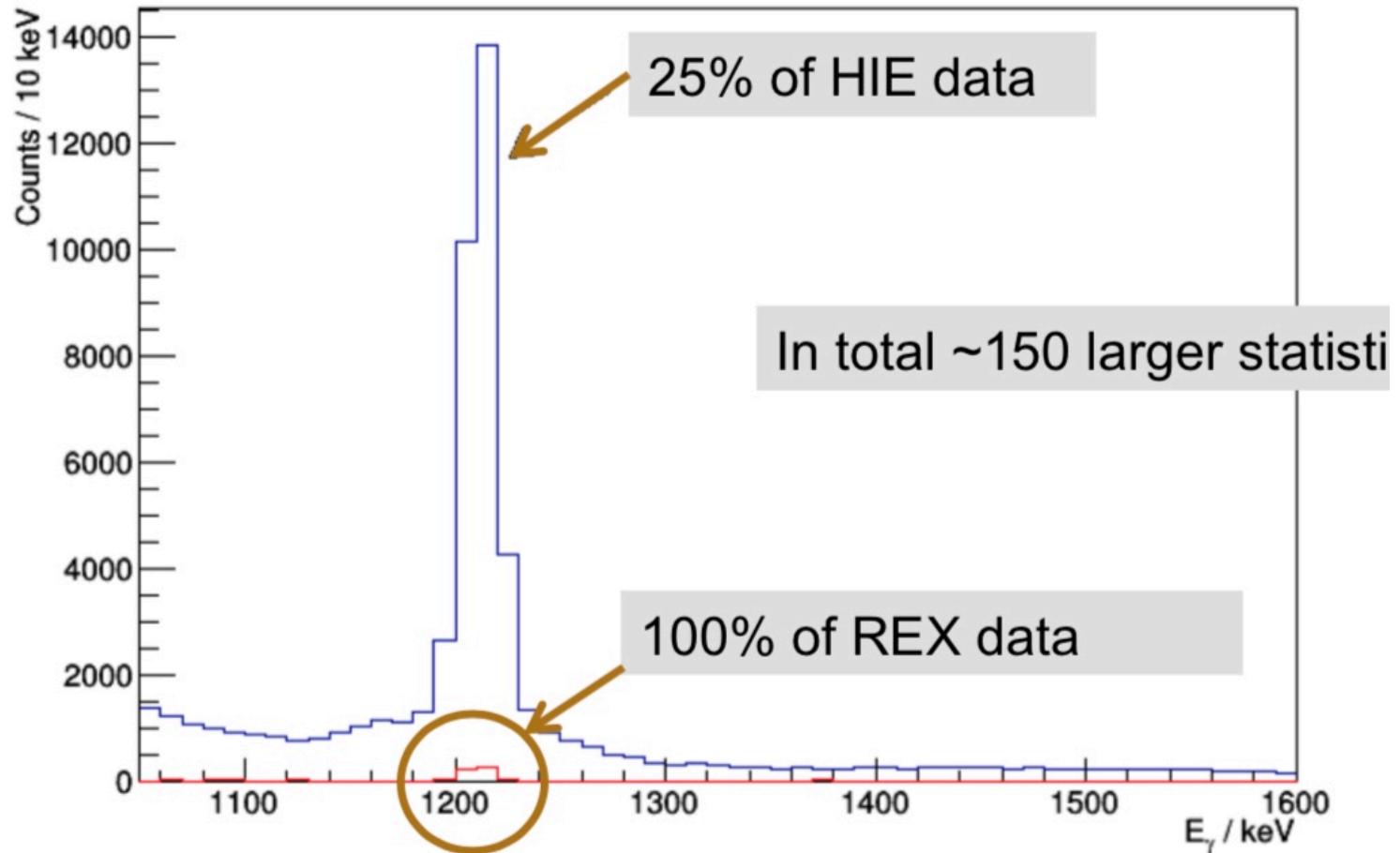


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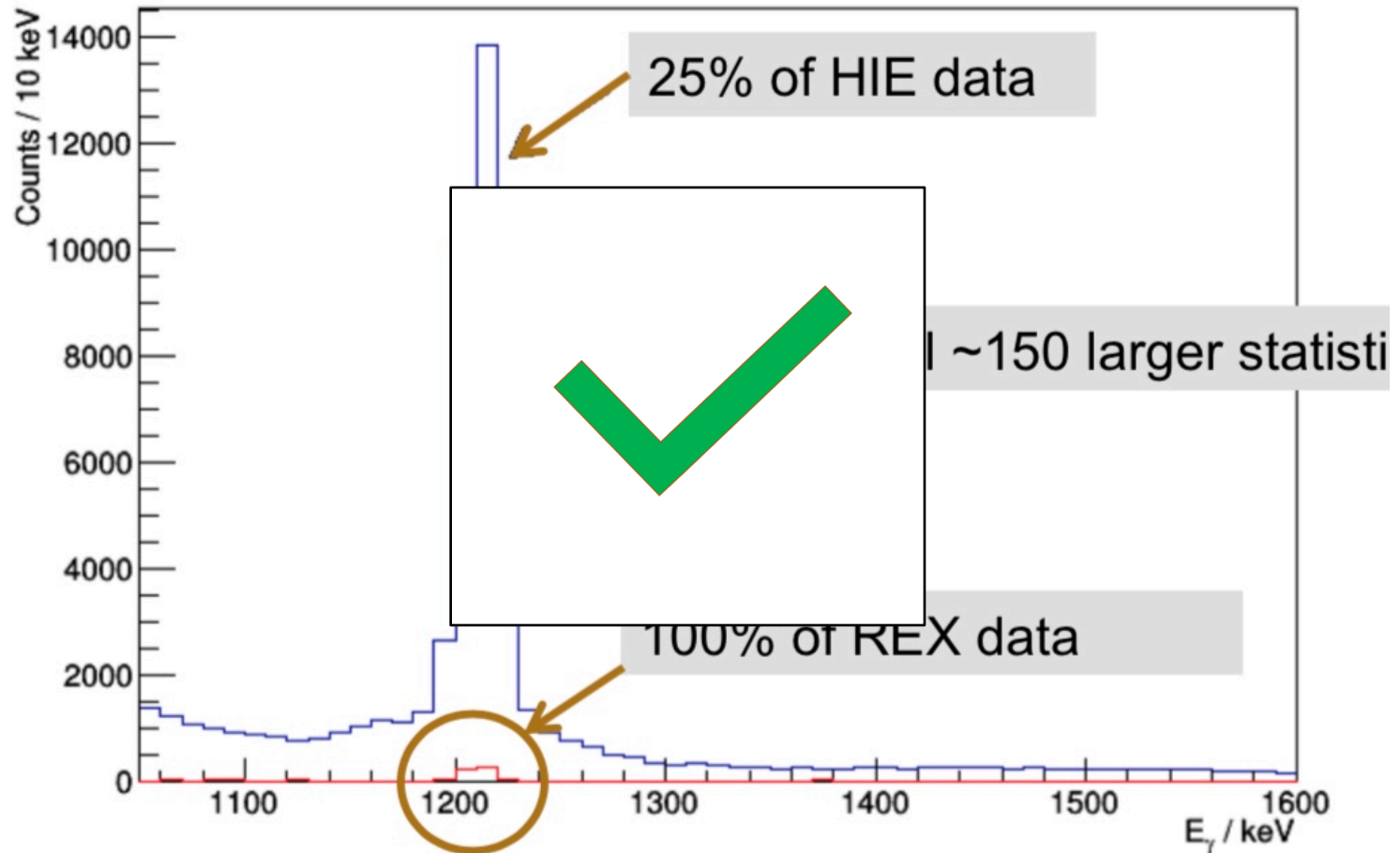
J. A. Rodriguez Rodriguez et al. priv comm.



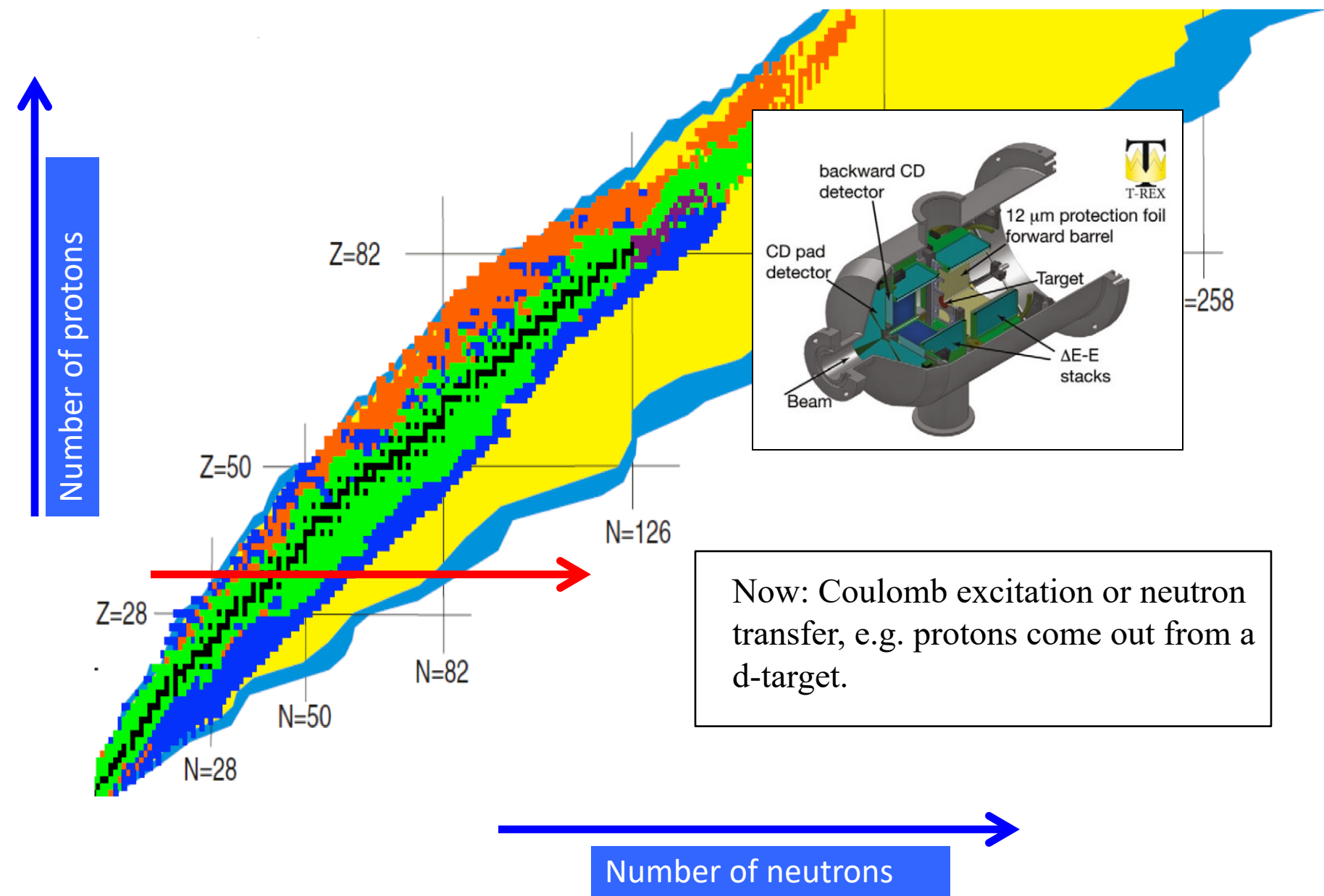
Statistics now and then



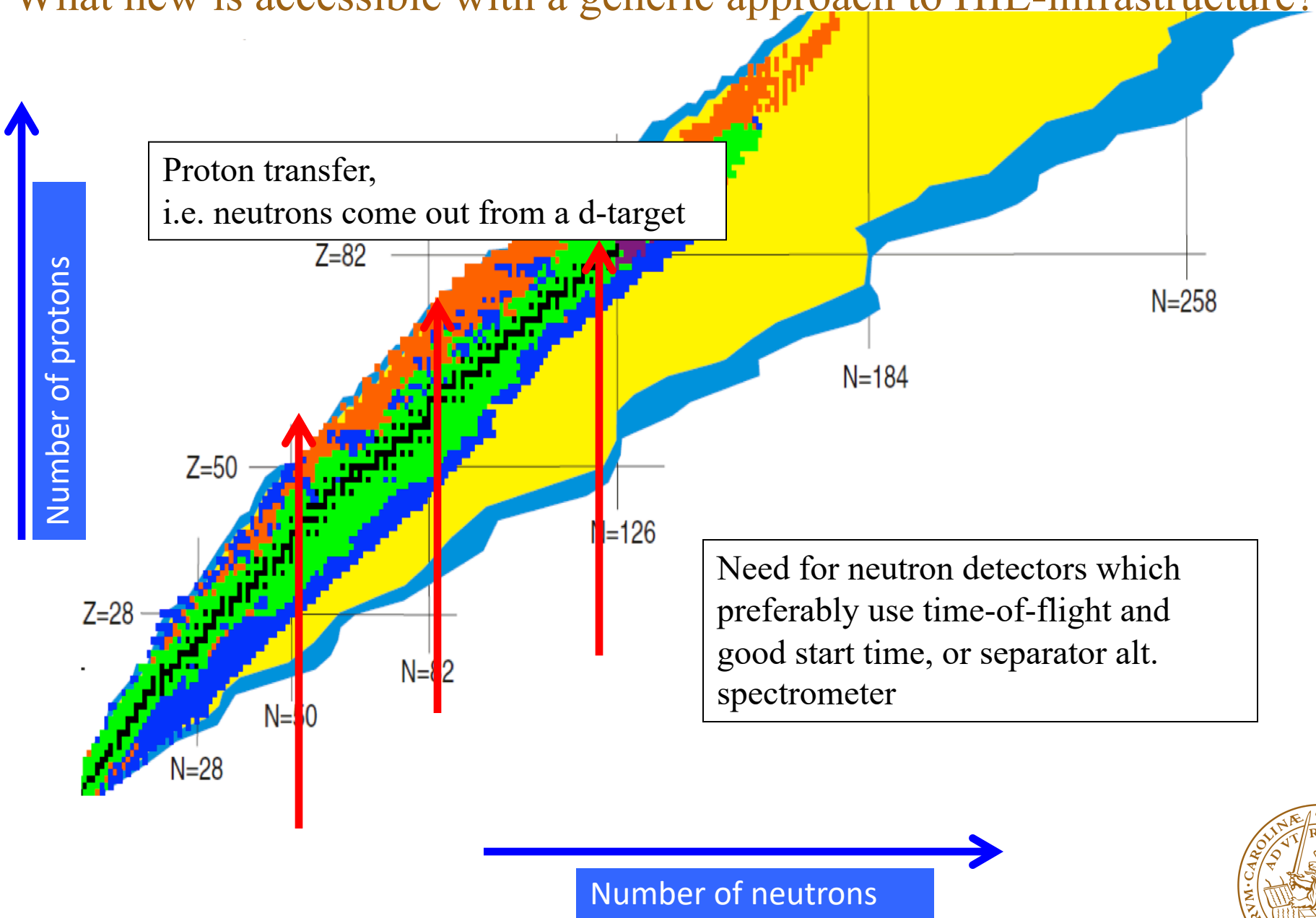
Statistics now and then



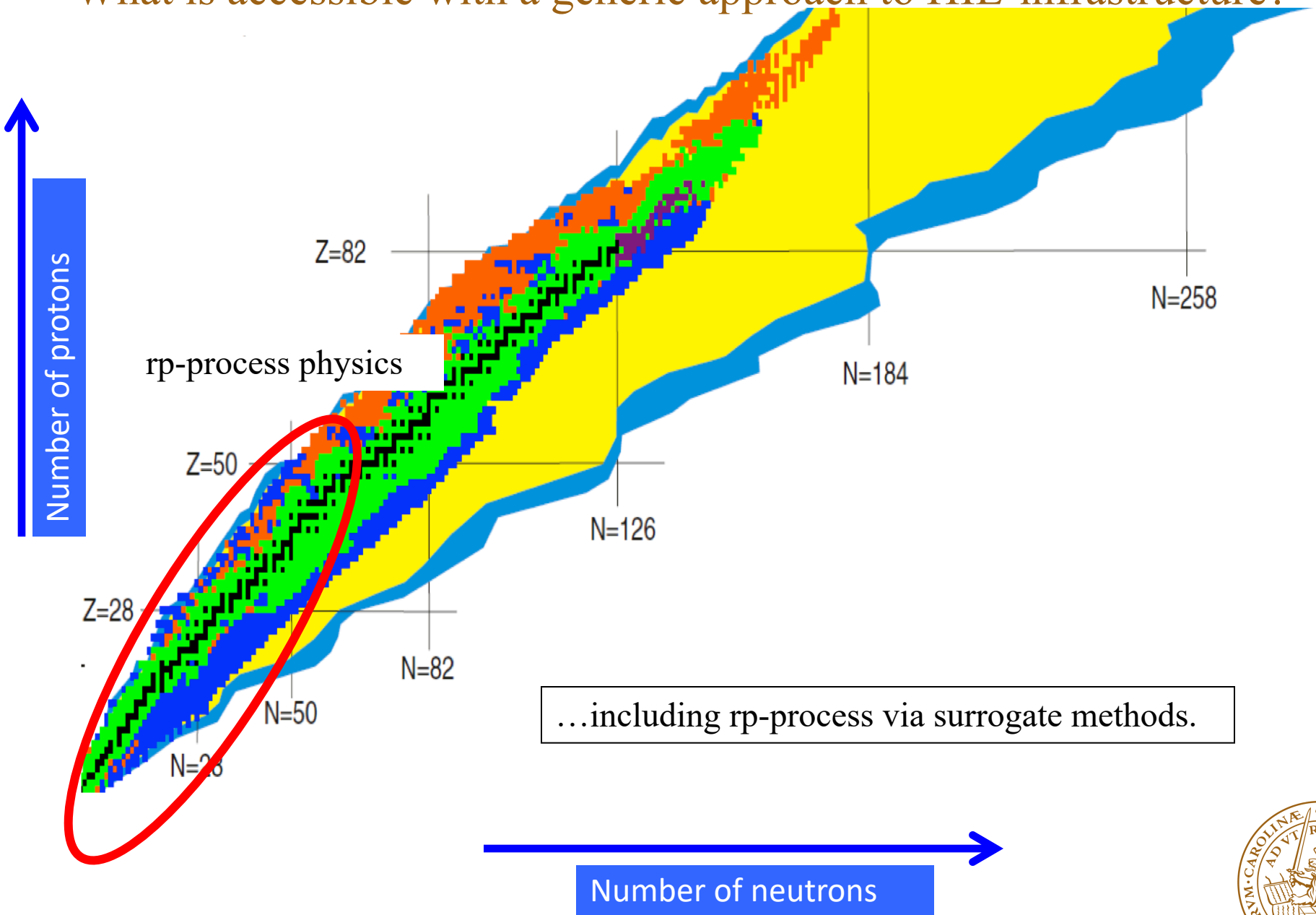
What is accessible with a generic approach to HIE-infrastructure?



What new is accessible with a generic approach to HIE-infrastructure?



What is accessible with a generic approach to HIE-infrastructure?



Single particle dominated states

For two particles outside a core:

$$\begin{aligned}
 H &= \underbrace{\sum_{k=3}^A [T_k + U(\vec{r}_k)] + \sum_{k=3}^A \sum_{l=k+1}^A W(\vec{r}_k, \vec{r}_l) - \sum_{k=3}^A U(\vec{r}_k)}_{H_{core}} \\
 &+ \underbrace{\sum_{k=1}^2 [T_k + U(\vec{r}_k)]}_{H_0^1 + H_0^2 \text{ independent motion}} + \underbrace{\sum_{k=1}^2 \sum_{l=k+1}^A W(\vec{r}_k, \vec{r}_l) - \sum_{k=1}^2 U(\vec{r}_k)}_{H_{12} \text{ interaction}} \\
 H_{12} &= \underbrace{\sum_{l=3}^A W(\vec{r}_1, \vec{r}_l) - U(r_1)}_{\approx 0} + \underbrace{\sum_{l=3}^A W(\vec{r}_2, \vec{r}_l) - U(r_2)}_{\approx 0} + W(\vec{r}_1, \vec{r}_2) \\
 &= V(\vec{r}_1, \vec{r}_2)
 \end{aligned}$$



Single particle dominated states

$$\begin{aligned}
 E = \langle \Phi_{J,T}^0 | H | \Phi_{J,T}^0 \rangle &= \underbrace{\langle \Phi_{J,T}^{core} | H_{core} | \Phi_{J,T}^{core} \rangle}_{\text{Binding energy of core}} \\
 &+ \underbrace{\langle \Phi_{J,T}^{\alpha_1, \alpha_2} | H_1 + H_2 | \Phi_{J,T}^{\alpha_1, \alpha_2} \rangle}_{\text{E1+E2 single-particle energies}} + \underbrace{\langle \Phi_{J,T}^{\alpha_1, \alpha_2} | V(\vec{r}_1, \vec{r}_2) | \Phi_{J,T}^{\alpha_1, \alpha_2} \rangle}_{\text{Interaction energy, two-body matrix element}}
 \end{aligned}$$

Transfer reactions

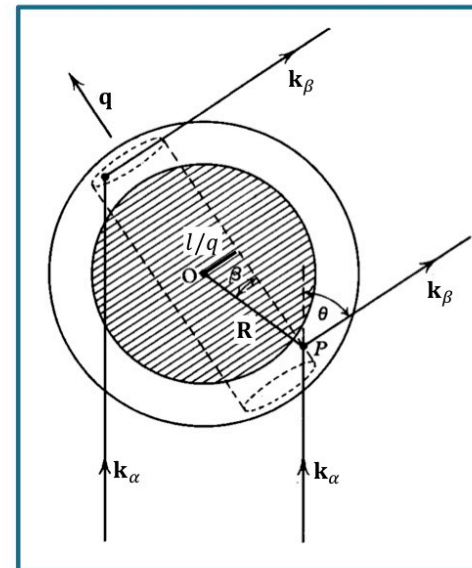
Angular momentum transfer: $\mathbf{l} = \mathbf{q} \times \mathbf{R}$,
with q the transferred momentum.

Cross section ~ 1 mb

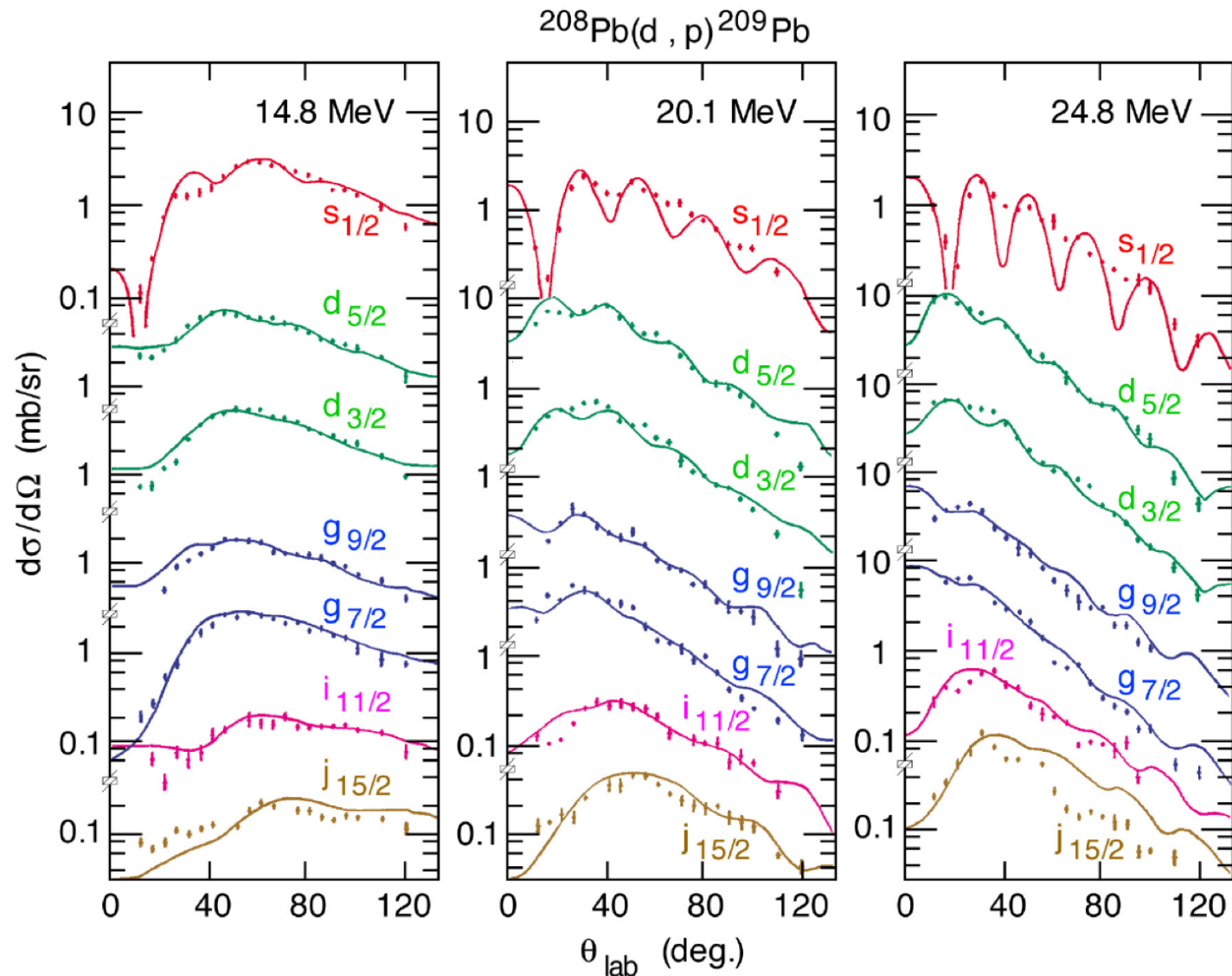
Target thickness 100 ug/cm^2 to mg/cm^2

Intensity $\sim 10^4$ pps

One can show that transferred angular momentum maps onto the scattering angle



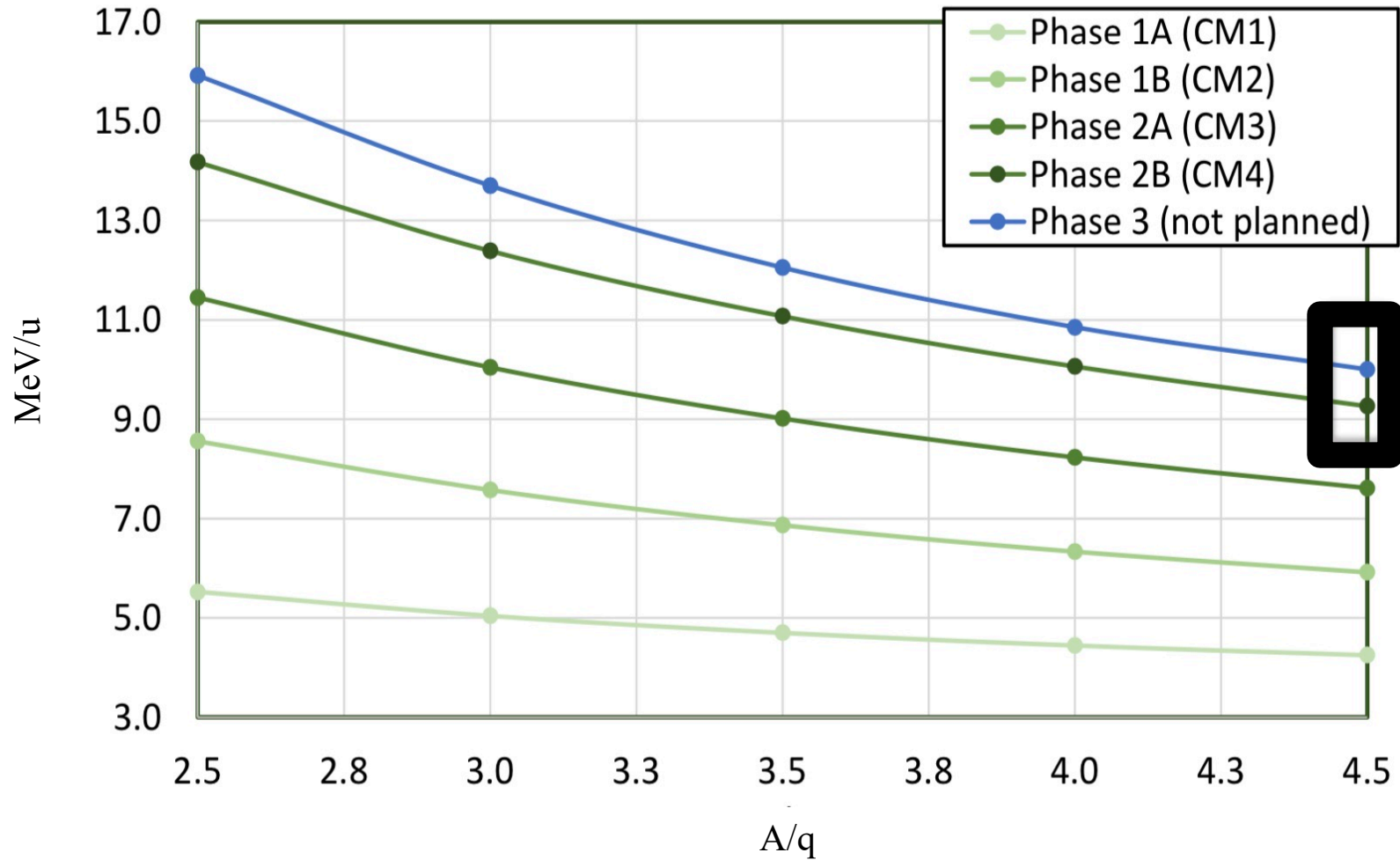
Angular distributions from (d,p)



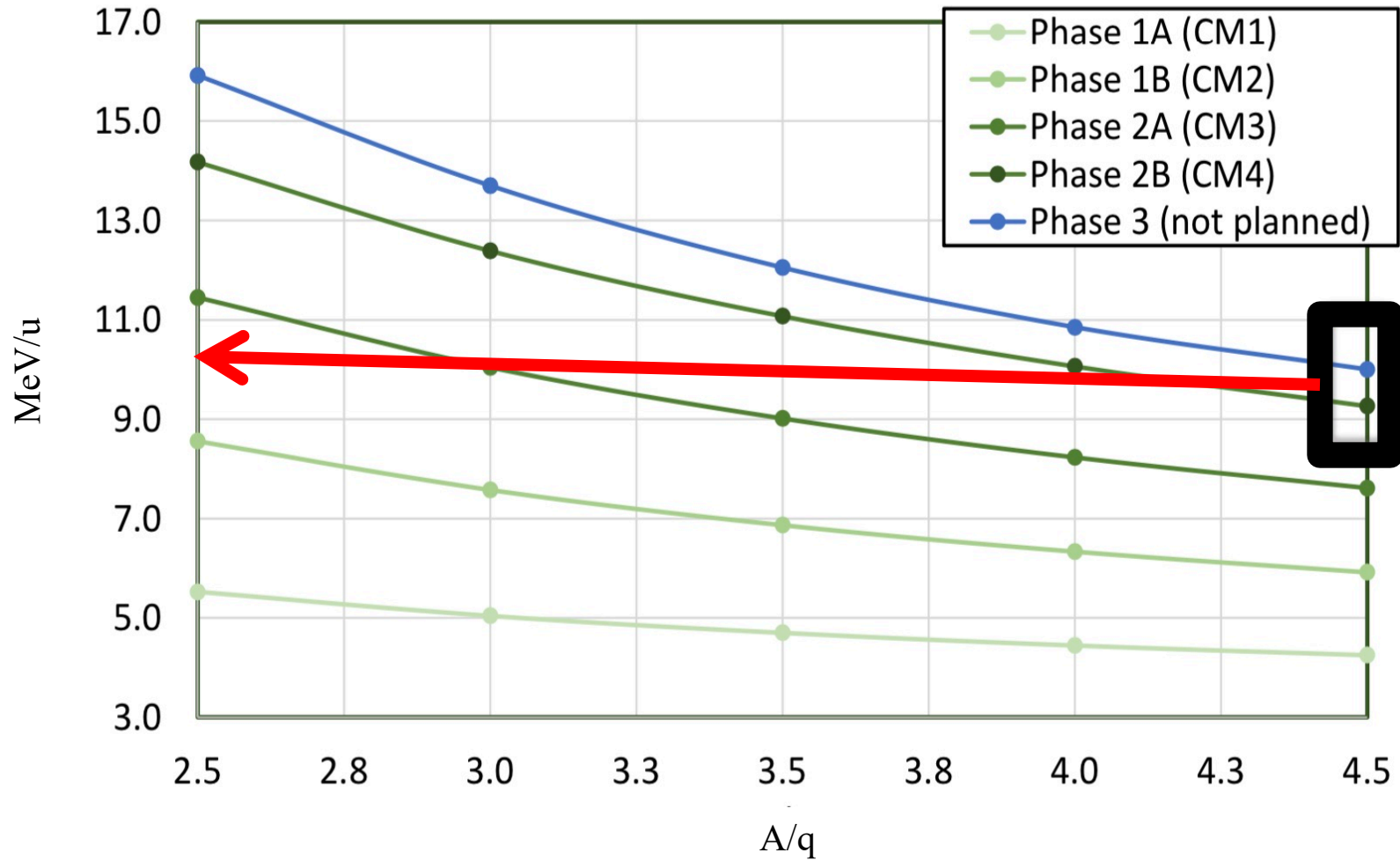
Muehlener et al.
Phys. Rev. 159, 1043 (1967)



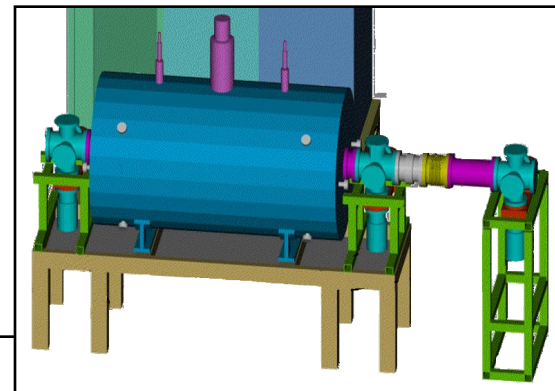
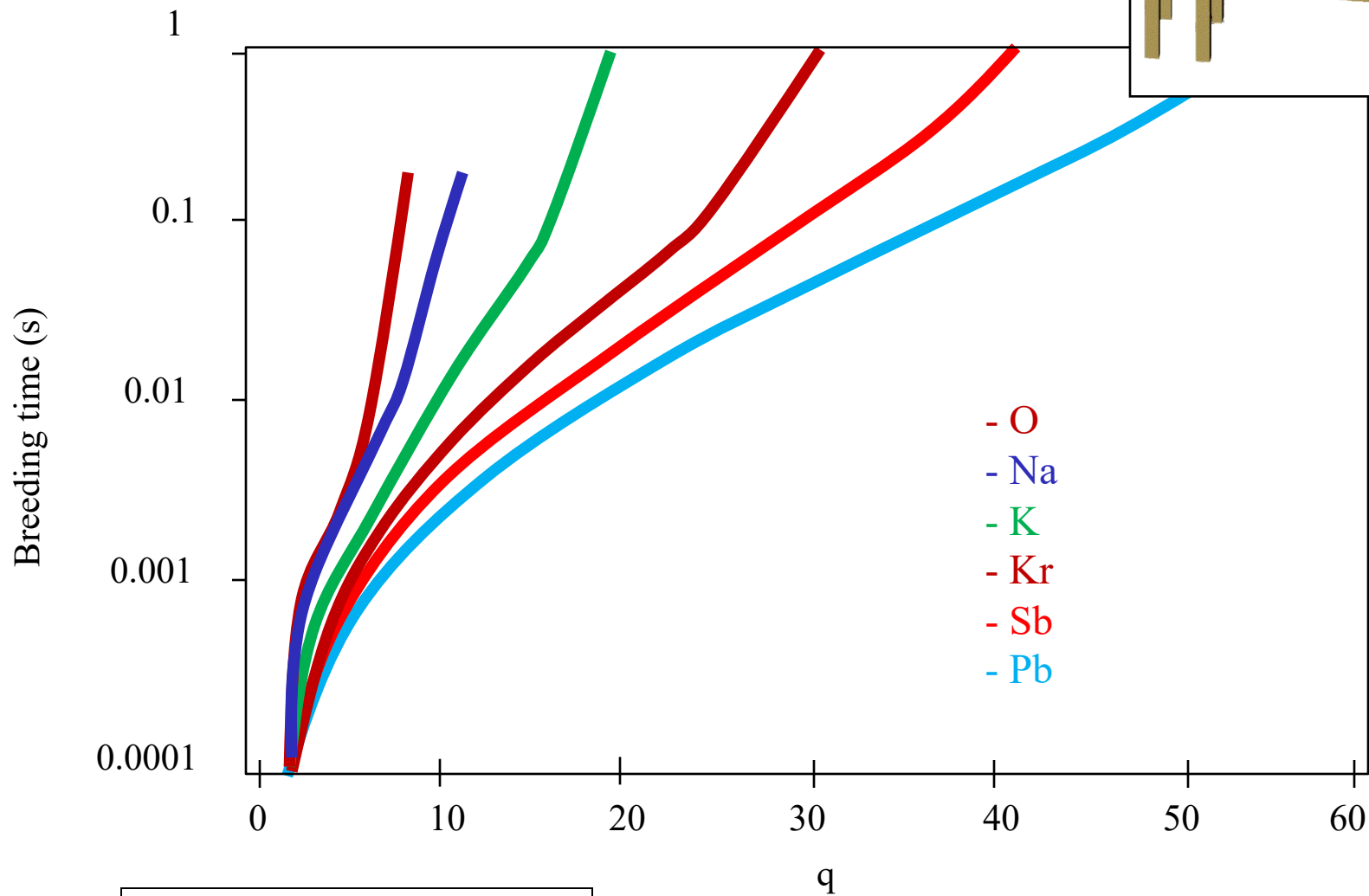
Maximum energy



Maximum energy



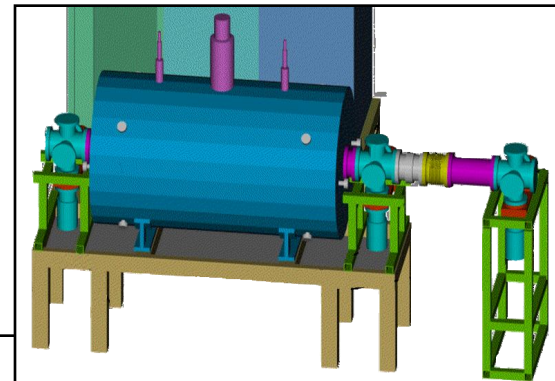
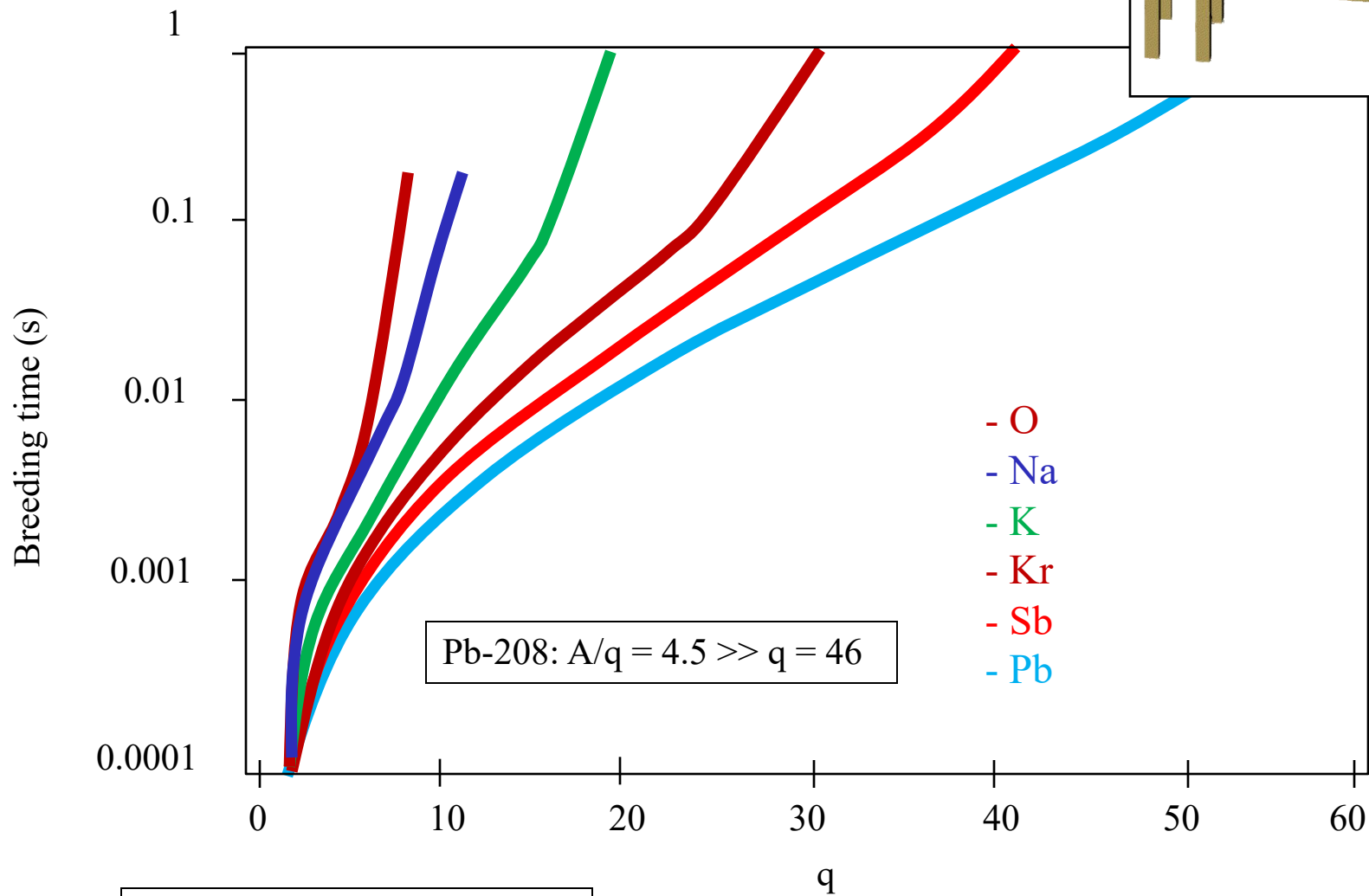
EBIS current situation



Courtesy F. Wenander



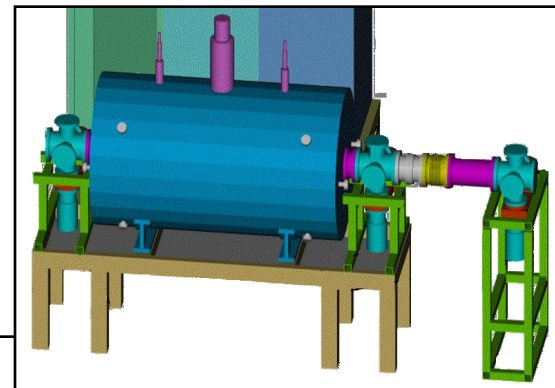
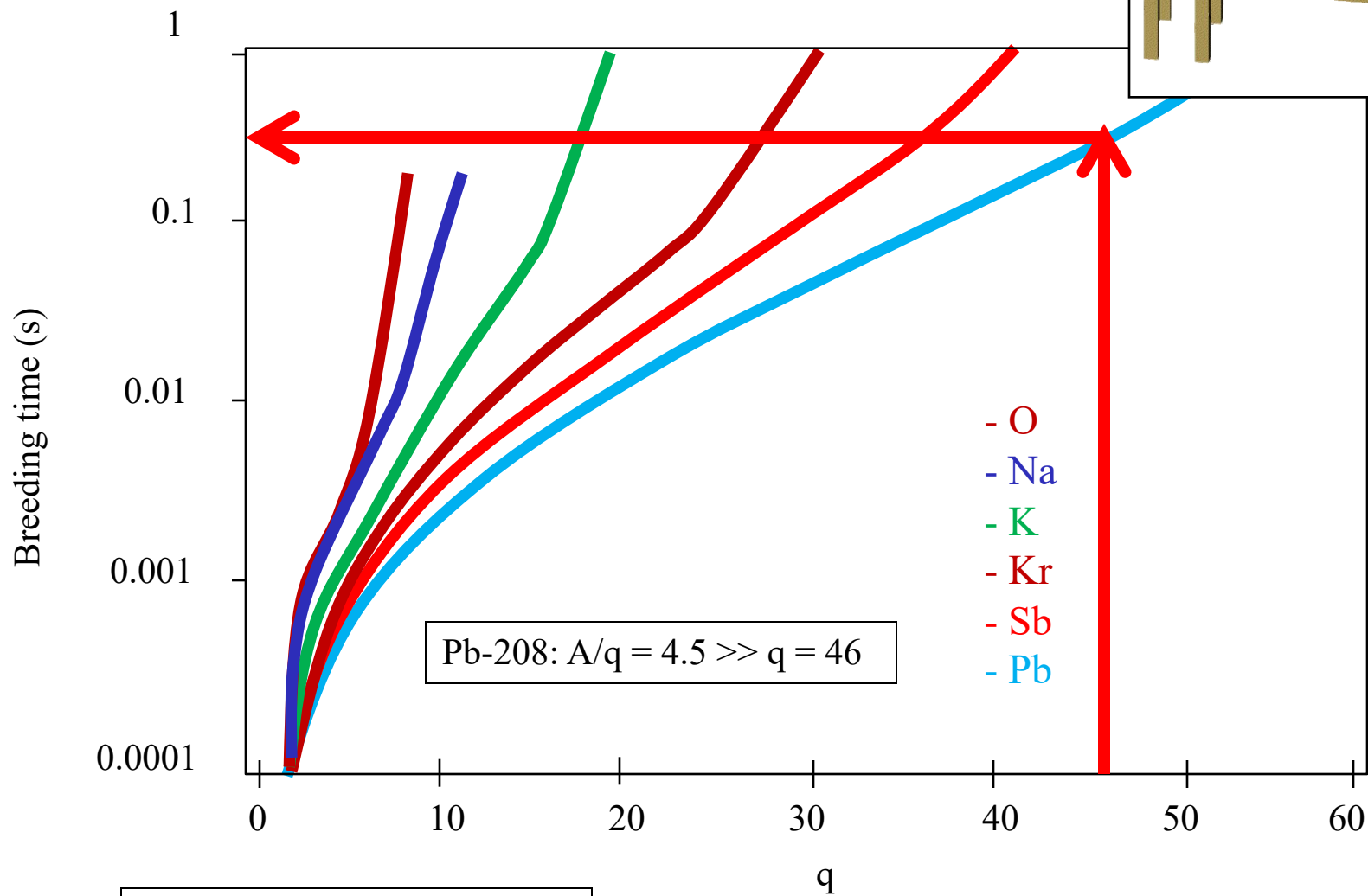
EBIS current situation



Courtesy F. Wenander



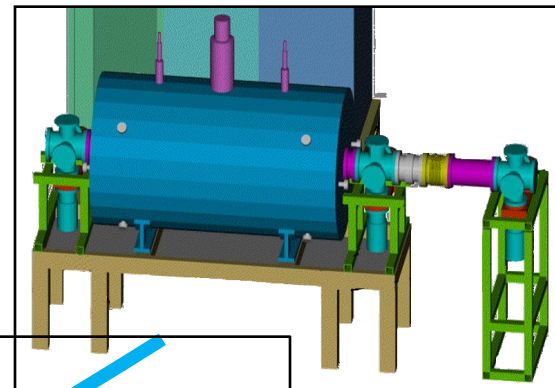
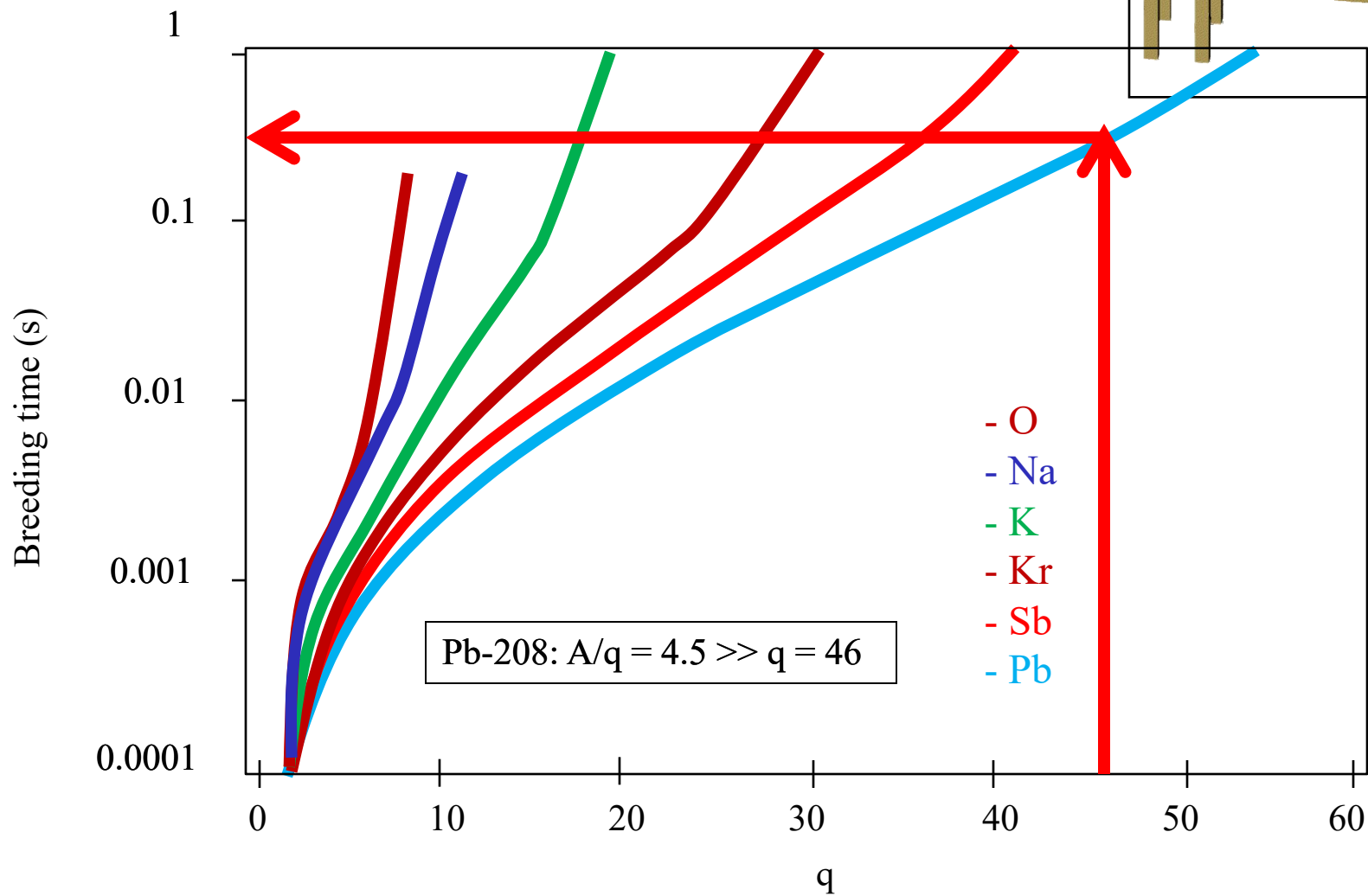
EBIS current situation



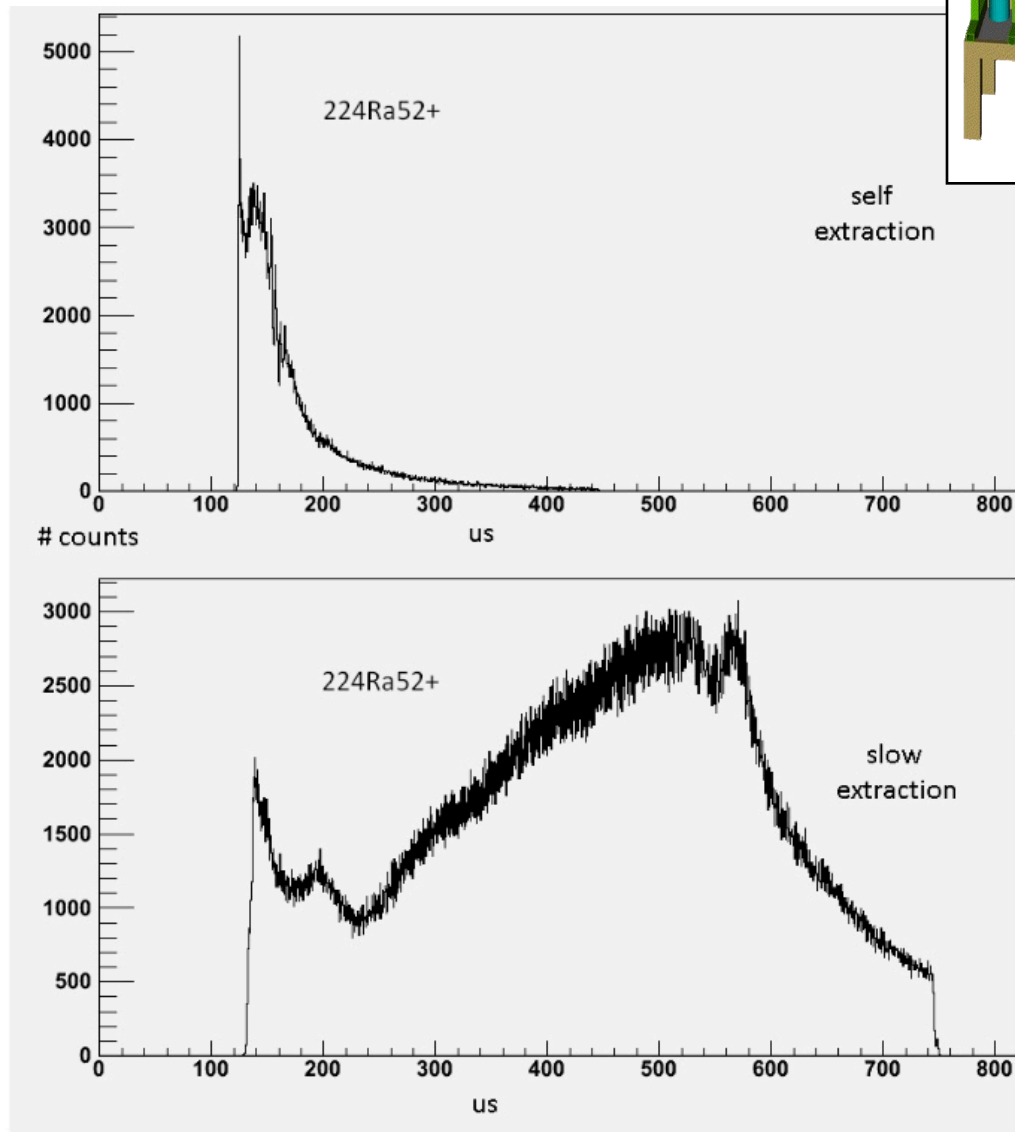
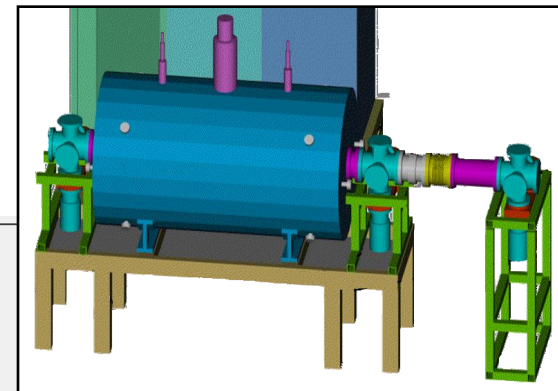
Courtesy F. Wenander



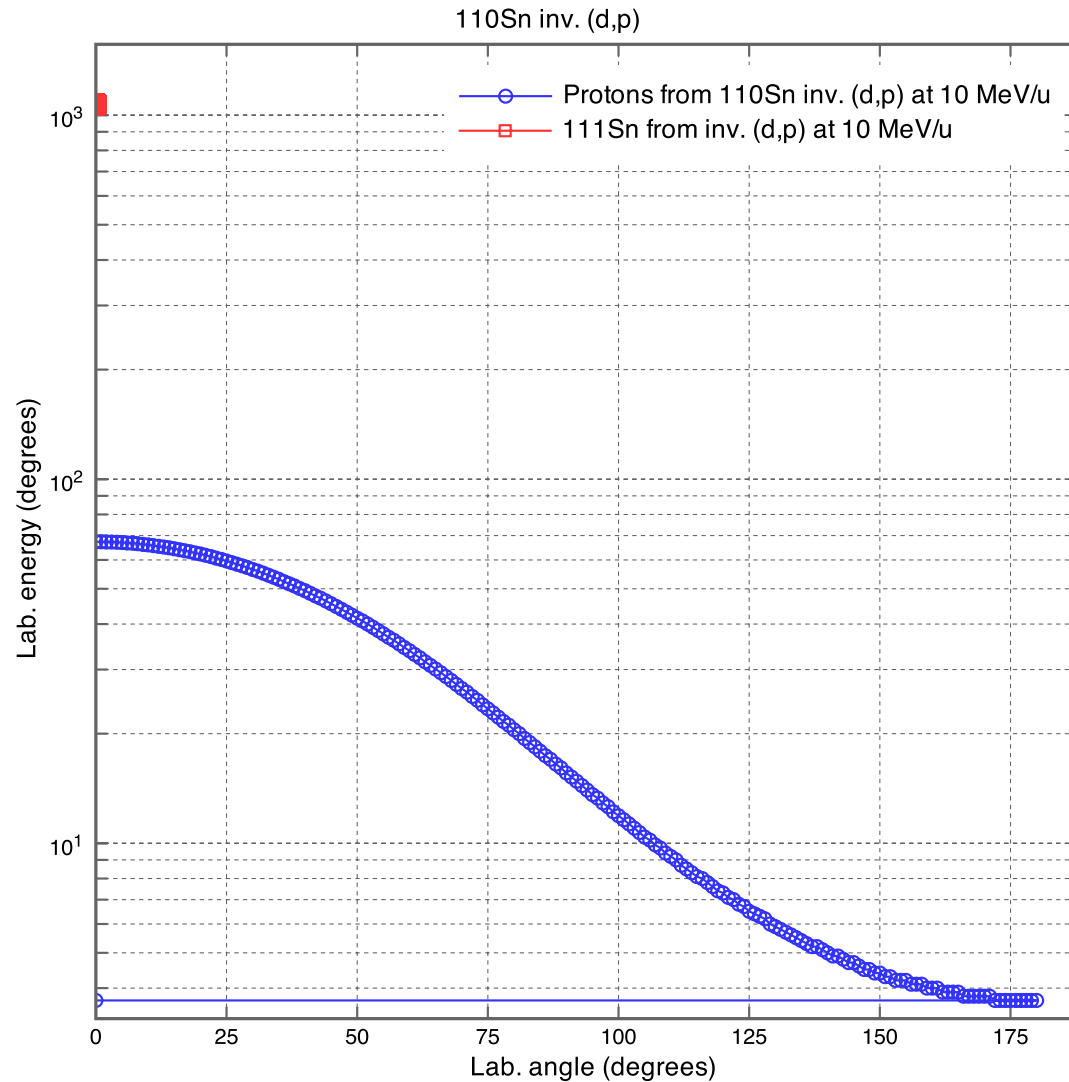
EBIS current situation



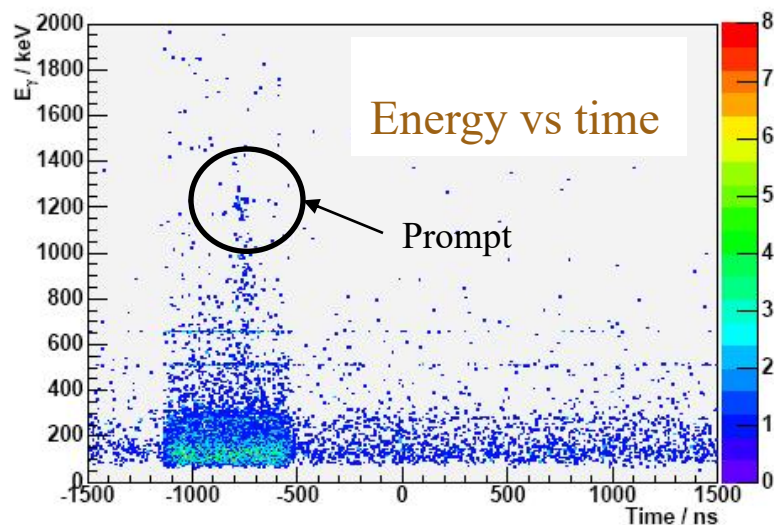
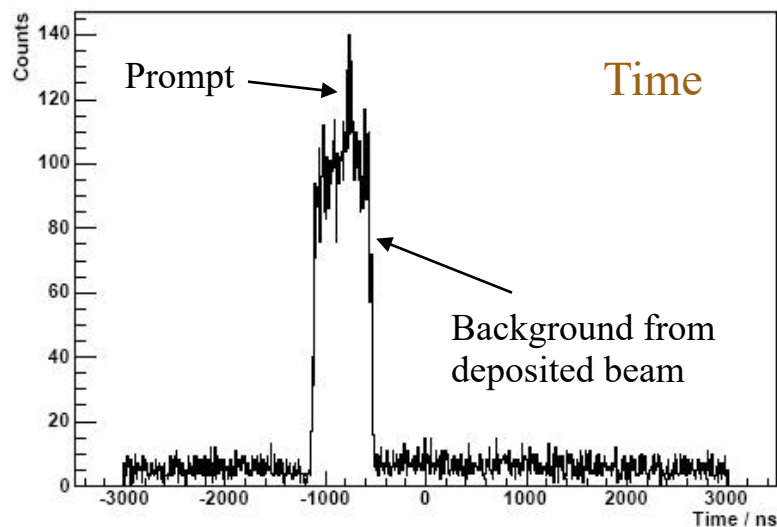
EBIS current situation



Sn-110, (d,p) kinematics

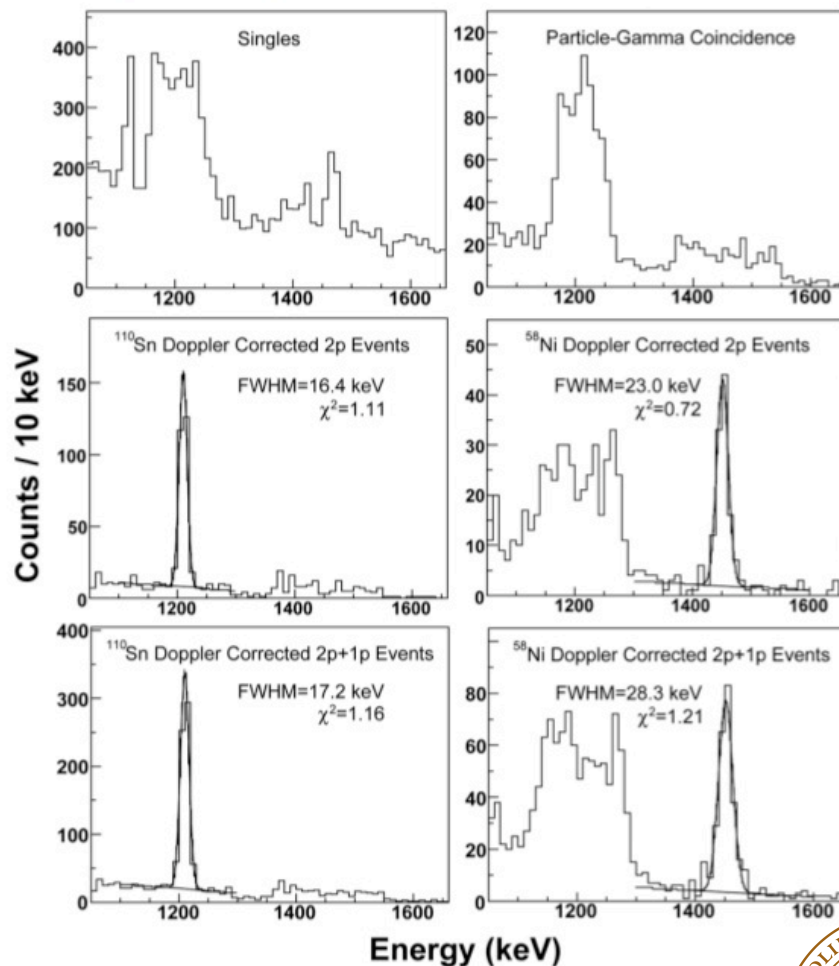


Particle – γ coincidence and Doppler shift

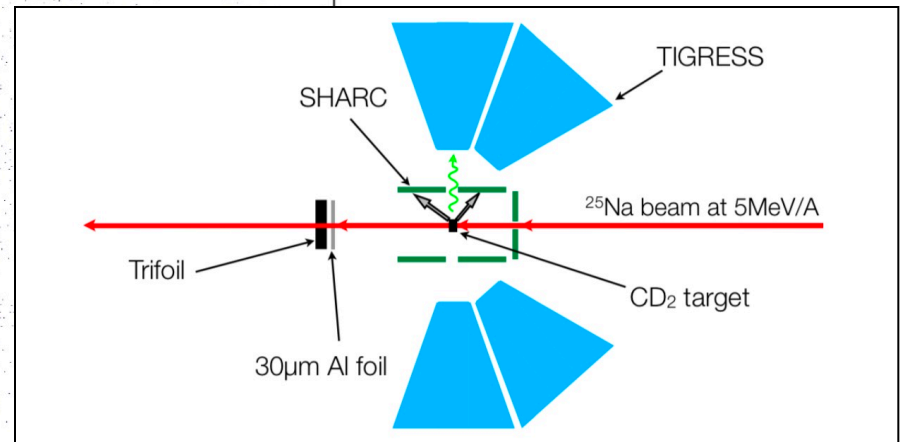
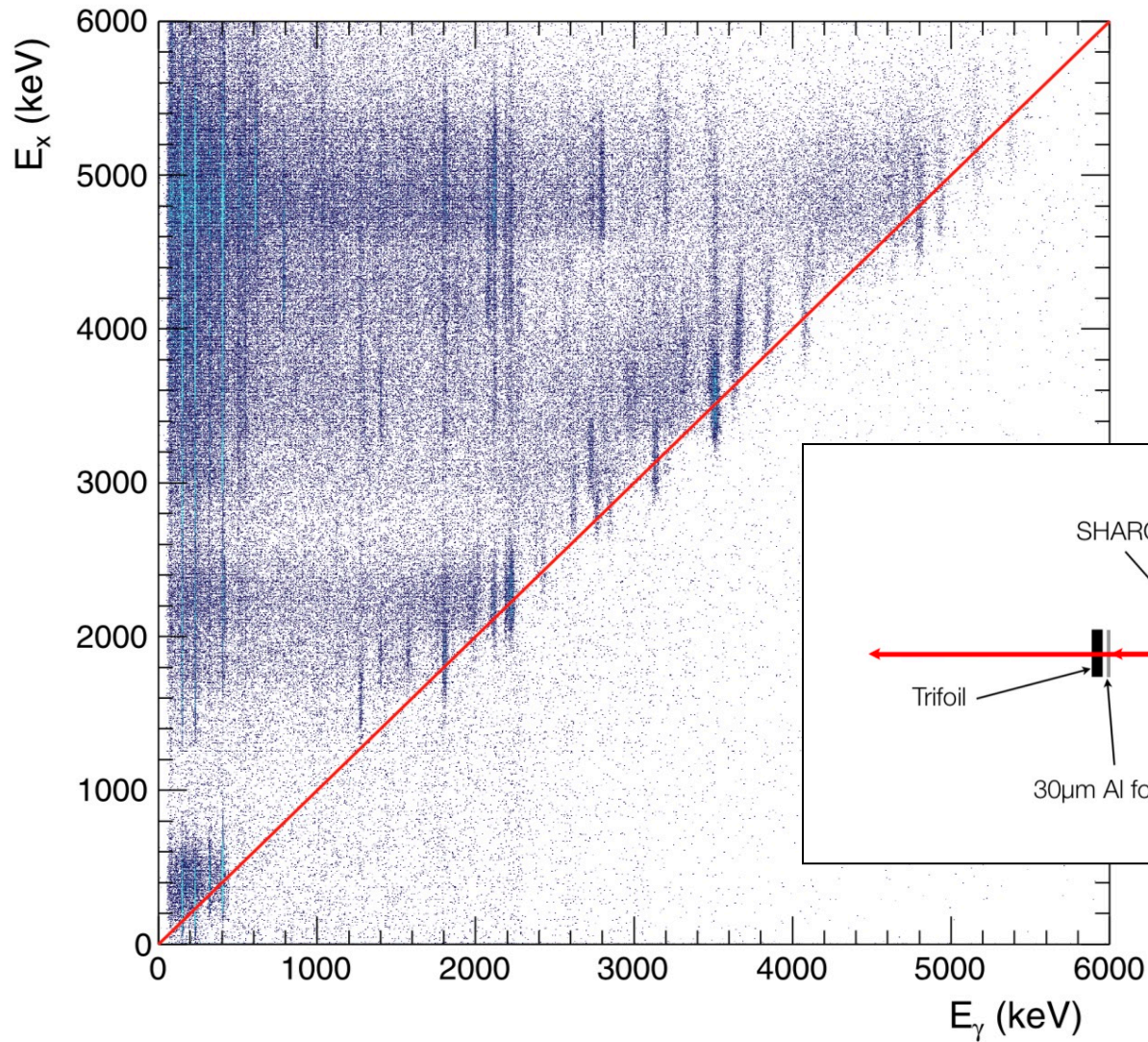


PRL **98**, 172501 (2007)

PHYSICAL REV



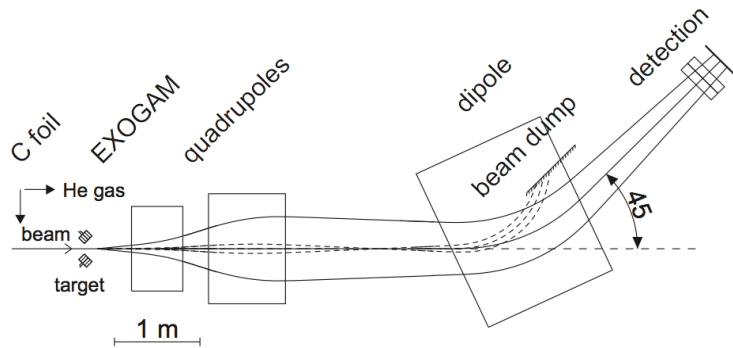
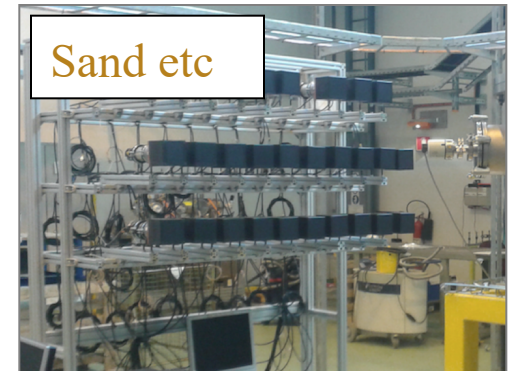
Recoil detection



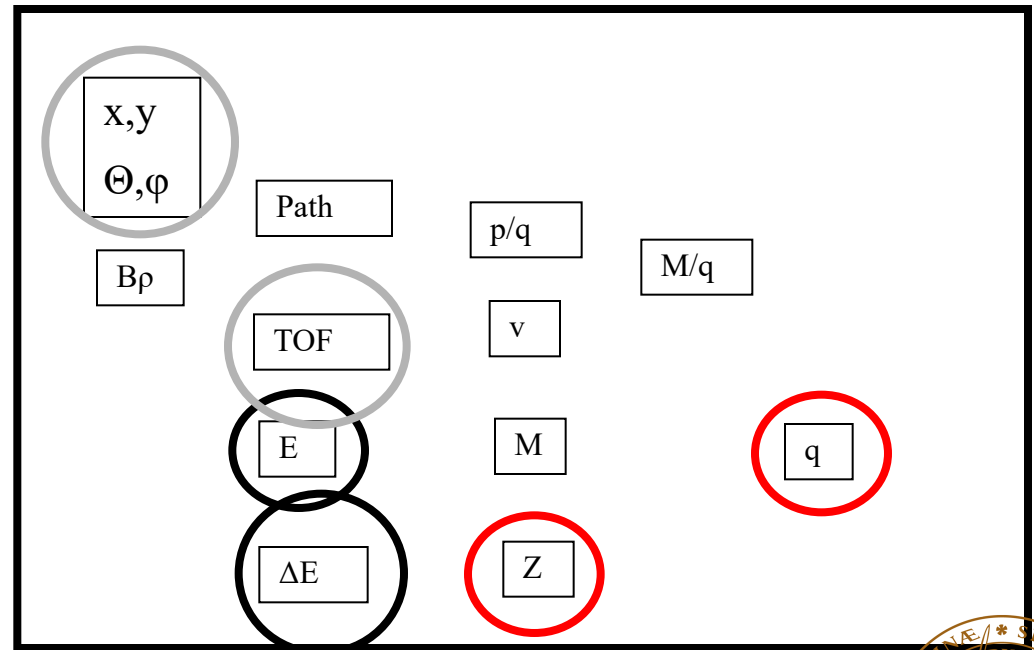
G. Wilson et al.
Physics Letters B 759 (2016)
417–423



TOF spectrometers etc...



M. Rejmund et al. NIMA 621 558 (2010)

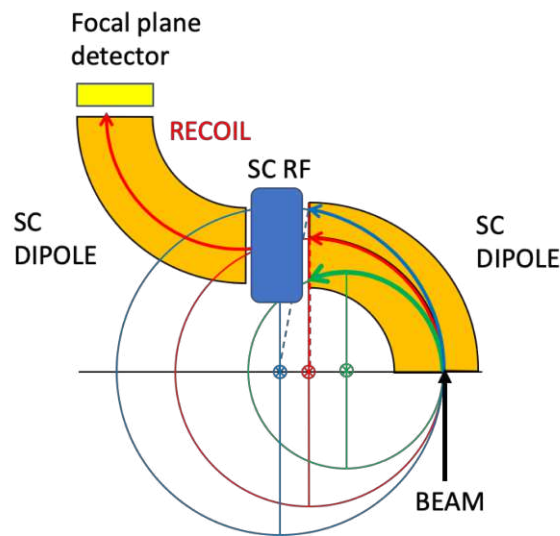


New ideas...

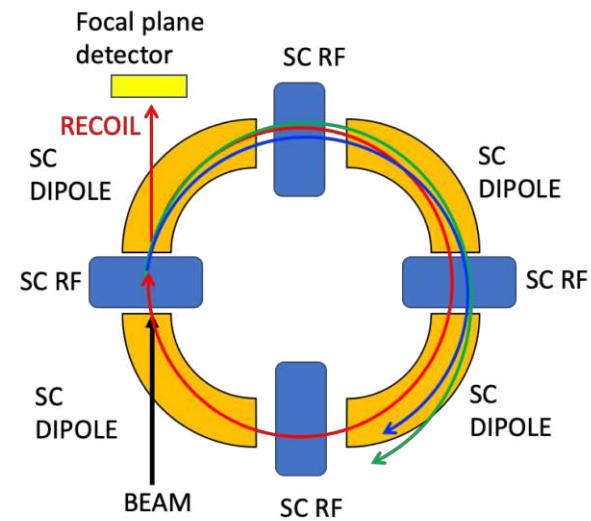
Ismael Bravo, Olof Tengblad

Proposal for a design study using SC elements

Explore new design concept using SC coils and RF cavities. Produce a compact, efficient and high-selectivity recoil separator. Study size, weight, efficiency, selectivity, cost and running cost.



Classical design concept



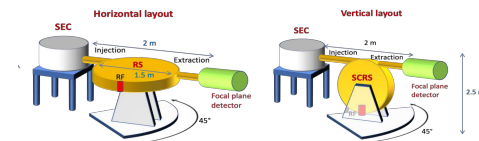
Ring design concept

SC solenoids

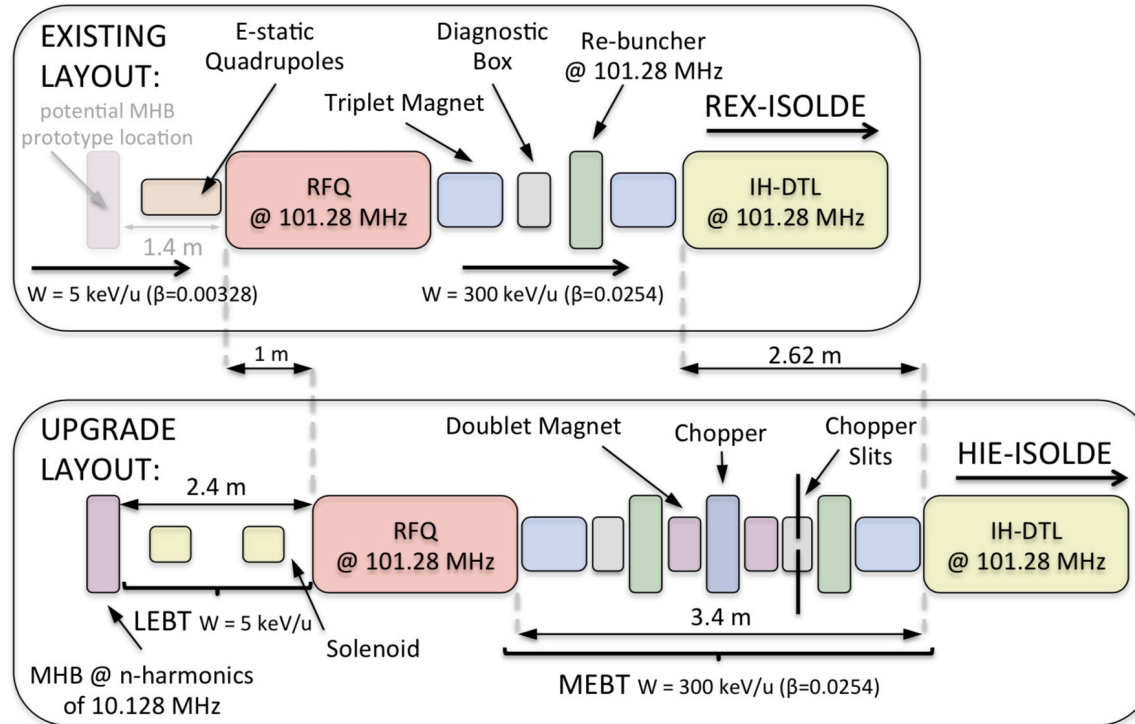
- Combined function magnets for bending and focussing
- High fields ~ 8 T

SC RF cavities

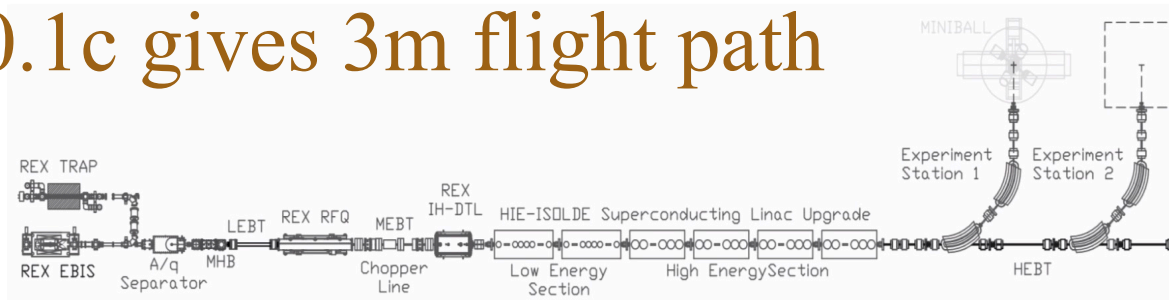
- High gradients ~ 10 MV/m
- HTS materials (> 77 K)
- Rebuncher ~ 10 MHz



100 ns Buncher



@0.1c gives 3m flight path



Resolution in spectrometer and Si-setup

Table 2

Major contributions in keV to the resolution of the excitation energy spectra of single neutron stripping and pickup reactions in inverse kinematics, where the heavy ion is detected in a spectrometer. The detection angle corresponds to 10°_{cm} . The last column is an approximate estimate as a sum in quadrature of the net effect of five non-Gaussian contributions. Other symbols are explained in the text

Reaction	E_i/A (MeV)	θ_{lab}	Origin of contribution					Σ_{quad}
			$\Delta\theta$	Δp	E_{stragg}	$\theta_{1/2}$	dE/dx	
$p(^{12}\text{Be}, ^{11}\text{Be})d$	30	1.07°	172	147	101	74	23	259
$p(^{12}\text{Be}, ^{11}\text{Be})d$	15	1.06°	84	71	99	74	37	169
$p(^{77}\text{Kr}, ^{76}\text{Kr})d$	30	0.16°	1404	811	808	723	56	1952
$p(^{77}\text{Kr}, ^{76}\text{Kr})d$	10	0.10°	334	143	502	570	268	883
$d(^{76}\text{Kr}, ^{77}\text{Kr})p$	10	0.21°	1140	614	2177	1859	1321	3408

Table 3

Major contributions in keV to the resolution of the excitation energy spectra of single neutron pickup and stripping reactions in inverse kinematics, where the light particle is detected in a silicon detector. Symbols as described in text and Table 2

Reaction	E_i/A (MeV)	θ_{lab}	Origin of contribution					Σ_{quad}
			$\Delta\theta$	ΔE_f	ΔE_i	$\theta_{1/2}$	dE/dx	
$p(^{12}\text{Be}, d)^{11}\text{Be}$	30	19.0°	136	74	114	96	649	685
$p(^{12}\text{Be}, d)^{11}\text{Be}$	15	17.8°	66	72	55	89	984	995
$p(^{77}\text{Kr}, d)^{76}\text{Kr}$	30	15.0°	124	55	64	63	186	249
$p(^{77}\text{Kr}, d)^{76}\text{Kr}$	10	6.0°	26	24	23	19	775	777
$d(^{76}\text{Kr}, p)^{77}\text{Kr}$	10	155.3°	52	93	37	60	1309	1310



What is next?

- **Bunch beam for any form of TOF measurement or timed injection**
 - time-of-flight dependent detectors (eg. for neutrons)
 - spectrometer/recoil separator
 - ring
- **Upgrade EBIS**
 - reach higher charge states for heavy nuclei quicker
 - higher charge states can reduce pressure on accelerator to reach highest voltages for all cavities.
- **Pre-studies already exists but need to be a priority for next step for the high-energy program together with higher intensity**

