

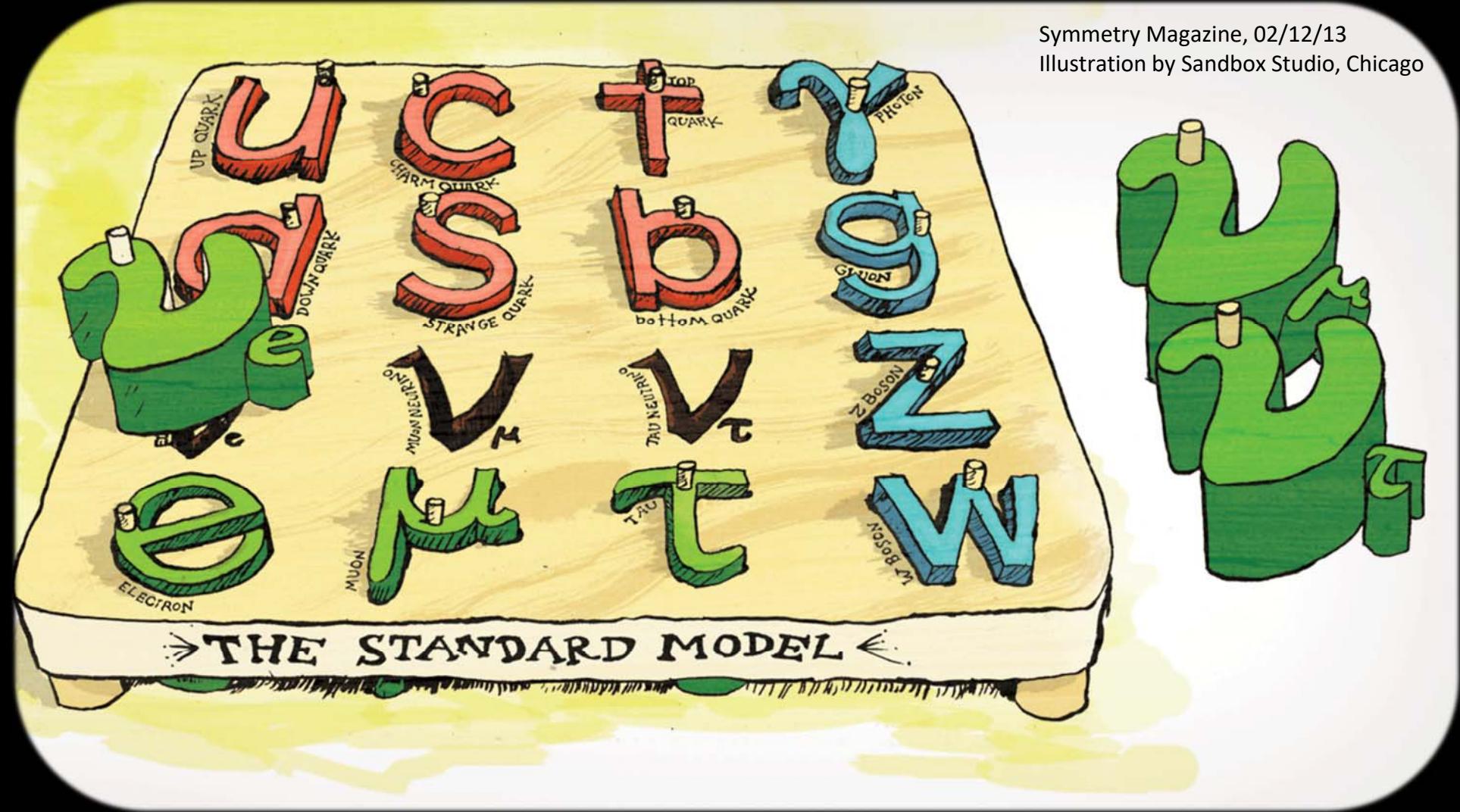
Perspectives for measuring the effective electron neutrino mass



Loredana Gastaldo

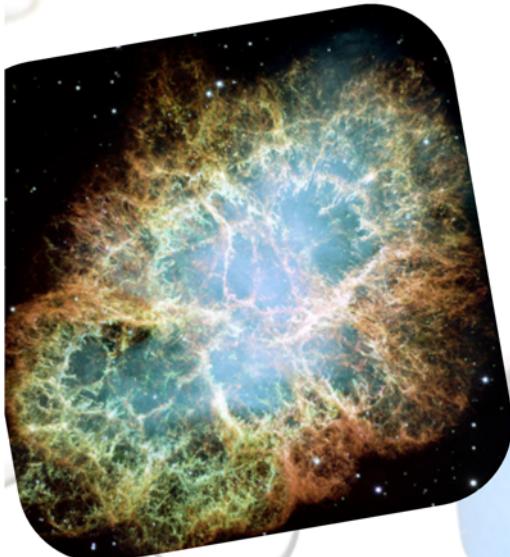
Kirchhoff Institute for Physics, Heidelberg University





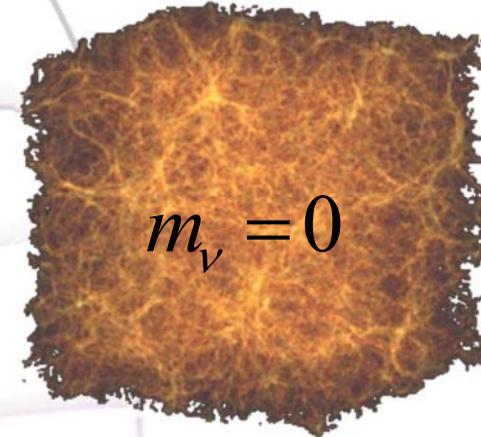
Symmetry Magazine, 02/12/13
Illustration by Sandbox Studio, Chicago

Knowing neutrino mass scale....

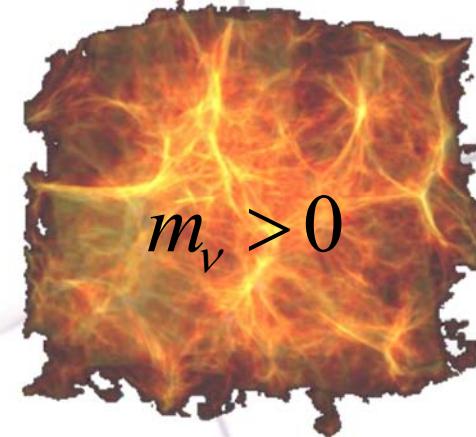


Astrophysics

Supernova neutrinos



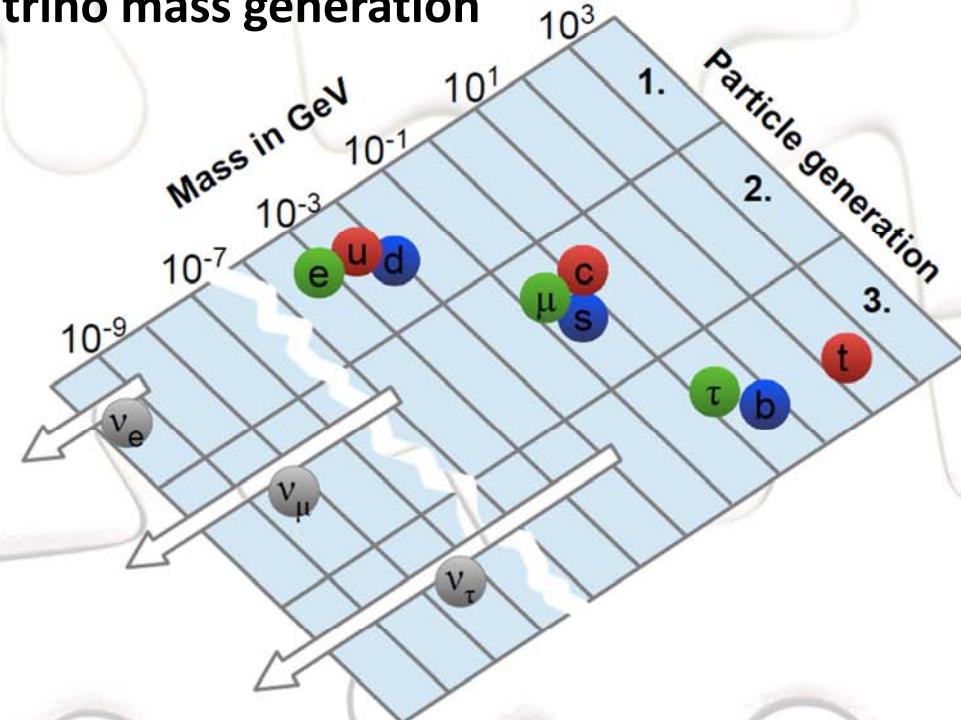
$$m_\nu = 0$$



$$m_\nu > 0$$

Particle Physics

Neutrino mass generation



Cosmology

Matter distribution
in the Universe

Outline

- Massive neutrinos
 - Observables and mass limits
 - Kinematics of EC and β decay
- Experimental approaches
 - ${}^3\text{H}$ & ${}^{163}\text{Ho}$
- Beyond the 3-neutrino scenario
- Conclusions



3-neutrino scenario

3 neutrino mass eigenstates

$$\nu_i = (\nu_1, \nu_2, \nu_3)$$

Important for Cosmology

3 neutrino flavour eigenstates

$$\nu_\alpha = (\nu_e, \nu_\mu, \nu_\tau)$$

Important for Particle Physics

Pontecorvo–Maki–Nakagawa–Sakata matrix $|\nu_\alpha\rangle = \sum_i U_{\alpha i} |\nu_i\rangle$

$$U_{\alpha i} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{bmatrix} \begin{bmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{bmatrix} \begin{bmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 \\ 0 & e^{-i\alpha/2} & 0 \\ 0 & 0 & e^{-i\beta/2} \end{bmatrix}$$

2-3 rotation

1-3 rotation
+ CP Dirac phase

1-2 rotation

extra Majorana
phases

3-neutrino scenario

Experiments studying neutrino flavor oscillations provides information on PMNS matrix parameters and difference in neutrino mass squared

Current knowledge *

$$\sin^2 \Theta_{12} \approx 0.303$$

$$\sin^2 \Theta_{23} \approx 0.545$$

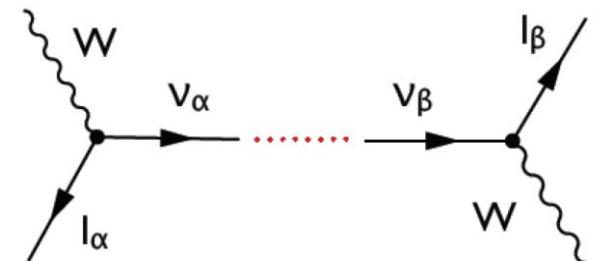
$$\sin^2 \Theta_{13} \approx 0.0225$$

$$\Delta m_{12}^2 \approx 7.34 \times 10^{-5} \text{ eV}^2$$

$$|\Delta m_{13}^2| \approx 2.48 \times 10^{-3} \text{ eV}^2$$

Open questions:

- Mass ordering ($>3\sigma$ NO)
- CP violation (1.6σ CPV)
- Majorana/Dirac
- Octant Θ_{23}
- Absolute scale of the neutrino masses



$$U_{\alpha i} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{bmatrix} \begin{bmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{bmatrix} \begin{bmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 \\ 0 & e^{-i\alpha/2} & 0 \\ 0 & 0 & e^{-i\beta/2} \end{bmatrix}$$

2-3 rotation

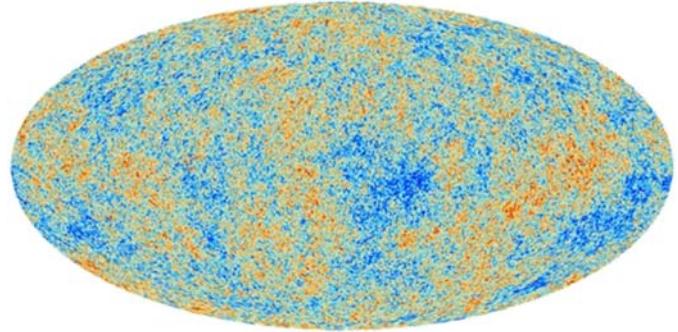
1-3 rotation
+ CP Dirac phase

1-2 rotation

extra Majorana phases

Massive neutrinos...

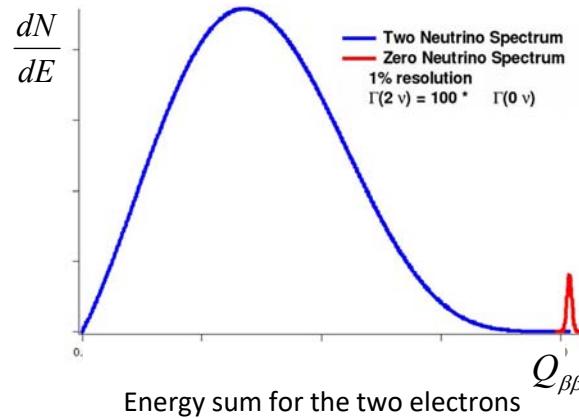
... affect the distribution of mass in the Universe



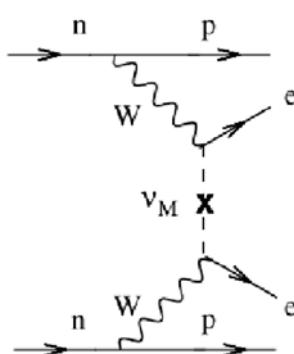
Present limit
Planck 2018

$$M = \sum_i m_i < 0.12 \text{ eV} \quad (95\% \text{ C.L.})$$

... could lead to neutrinoless double beta decay



+



$$(\tau_{1/2}^{0\nu})^{-1} = \left| \frac{m_{\beta\beta}}{m_e} \right|^2 |M_\nu^{0\nu}|^2 G^{0\nu}$$

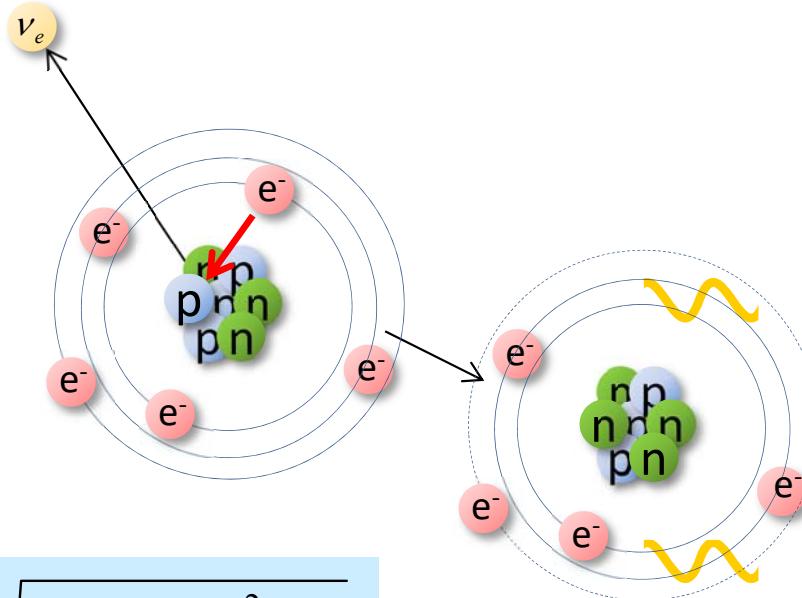
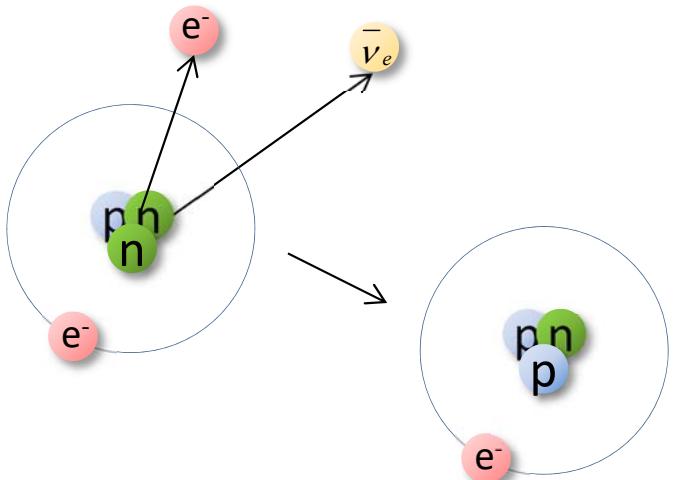
$$m_{\beta\beta}^2 = \left| \sum U_{ei}^2 m(\nu_i) \right|^2$$

Present limit $m_{\beta\beta} < 0.31 \text{ eV}$

Future $m_{\beta\beta} < 20 - 50 \text{ meV}$

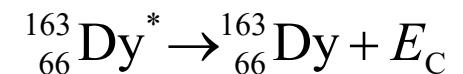
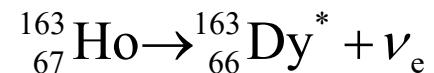
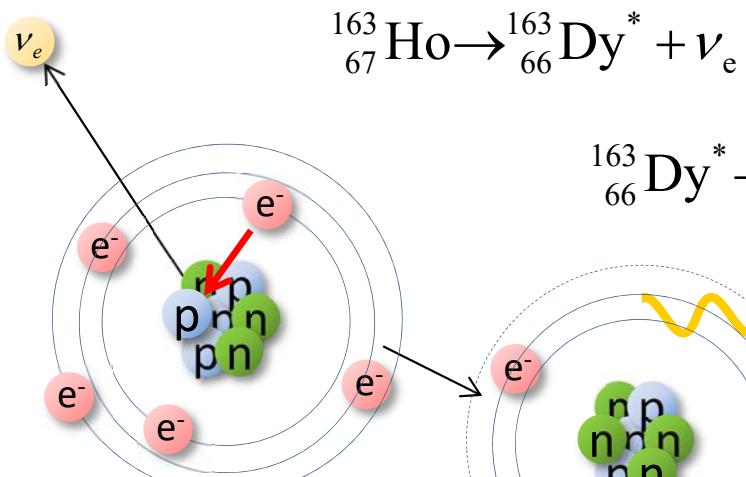
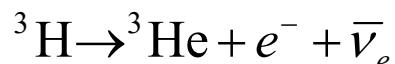
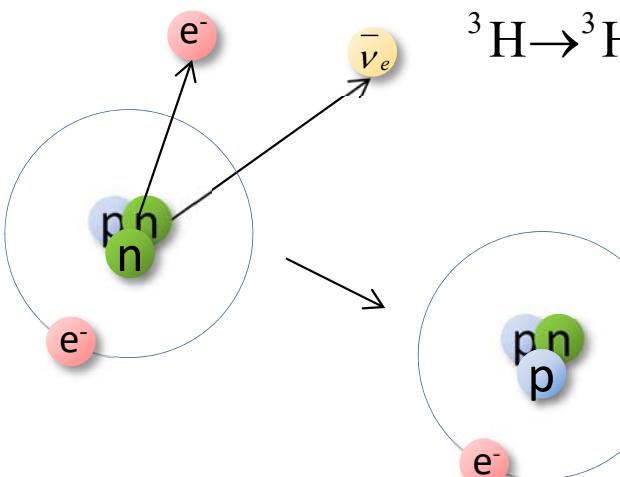
... affect the energy spectrum of low energy EC and β-decay

Beta decay and electron capture



$$\frac{dW}{dE} \propto (Q - E)^2 \sqrt{1 - \frac{m_\beta^2}{(Q - E)^2}}$$

Beta decay and electron capture



$$\frac{dW}{dE} \propto (Q - E)^2 \sqrt{1 - \frac{m_\beta^2}{(Q - E)^2}}$$

- $\tau_{1/2} \approx 12.3$ years (4*10⁸ atoms for 1 Bq)

- $Q_B = 18\,592.01(7)$ eV

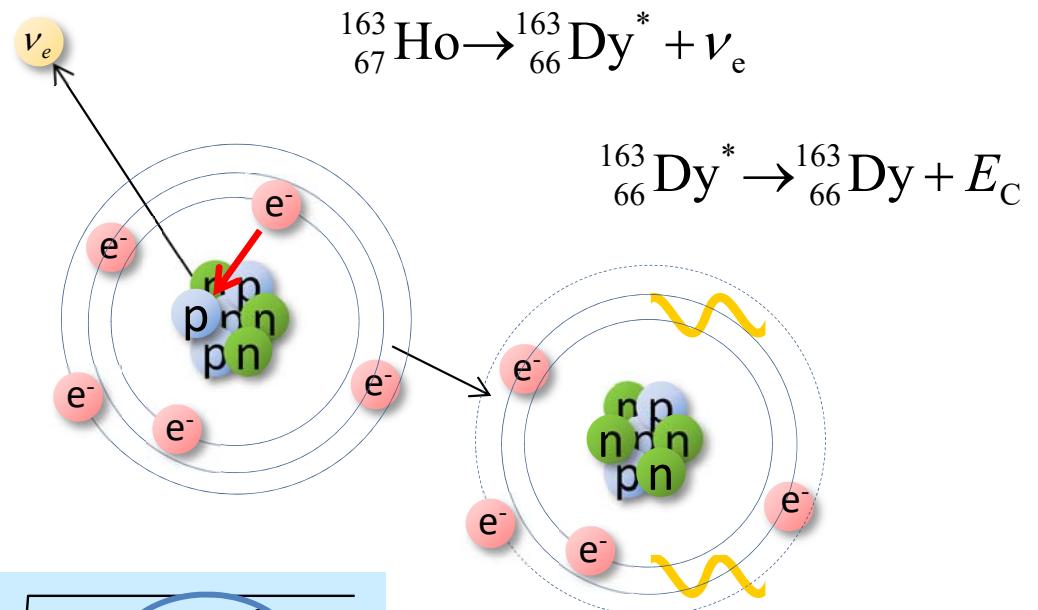
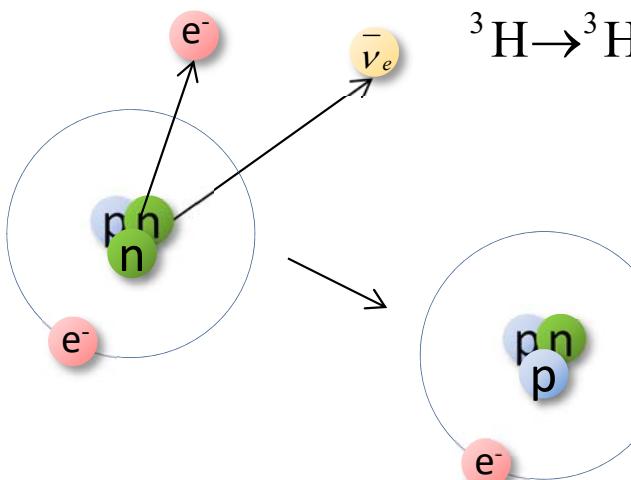
E.G. Myers et al., *Phys. Rev. Lett.* **114** (2015) 013003

- $\tau_{1/2} \approx 4570$ years (2*10¹¹ atoms for 1 Bq)

- $Q_{EC} = (2.833 \pm 0.030^{\text{stat}} \pm 0.015^{\text{syst}})$ keV

S. Eliseev et al., *Phys. Rev. Lett.* **115** (2015) 062501

Beta decay and electron capture



$$\frac{dW}{dE} \propto (Q - E)^2 \sqrt{1 - \frac{m_\beta^2}{(Q - E)^2}}$$

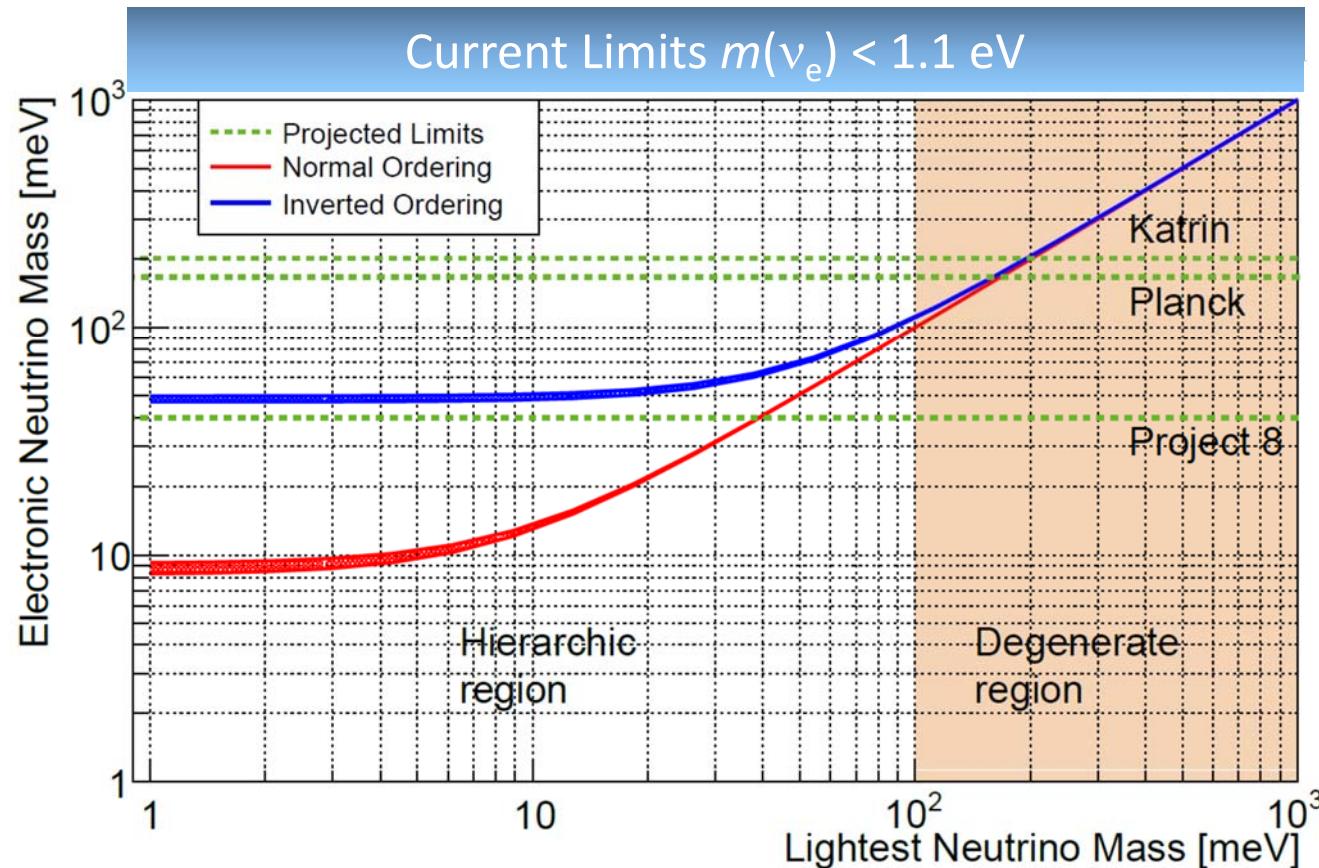
$$\frac{dW}{dE_C} \propto (Q - E)^2 \sum_i |U_{ei}|^2 \sqrt{1 - \frac{m_i^2}{(Q - E)^2}}$$

finite energy resolution $\xrightarrow{\hspace{1cm}}$

$$m_\beta^2(\nu_e) = \sum_i |U_{ei}|^2 m_i^2$$

Beta decay and electron capture

Ahtari et al., *JPhysG* **44** (2017) 5



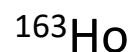
$$m(\bar{\nu}_e) < 1.1 \text{ eV}$$



Aker et al., (The KATRIN Collaboration)

Improved Upper Limit on the Neutrino Mass from a Direct Kinematic Method by KATRIN, *Phys. Rev. Lett.* **123** (2019) 221802

$$m(\nu_e) < 150 \text{ eV}$$

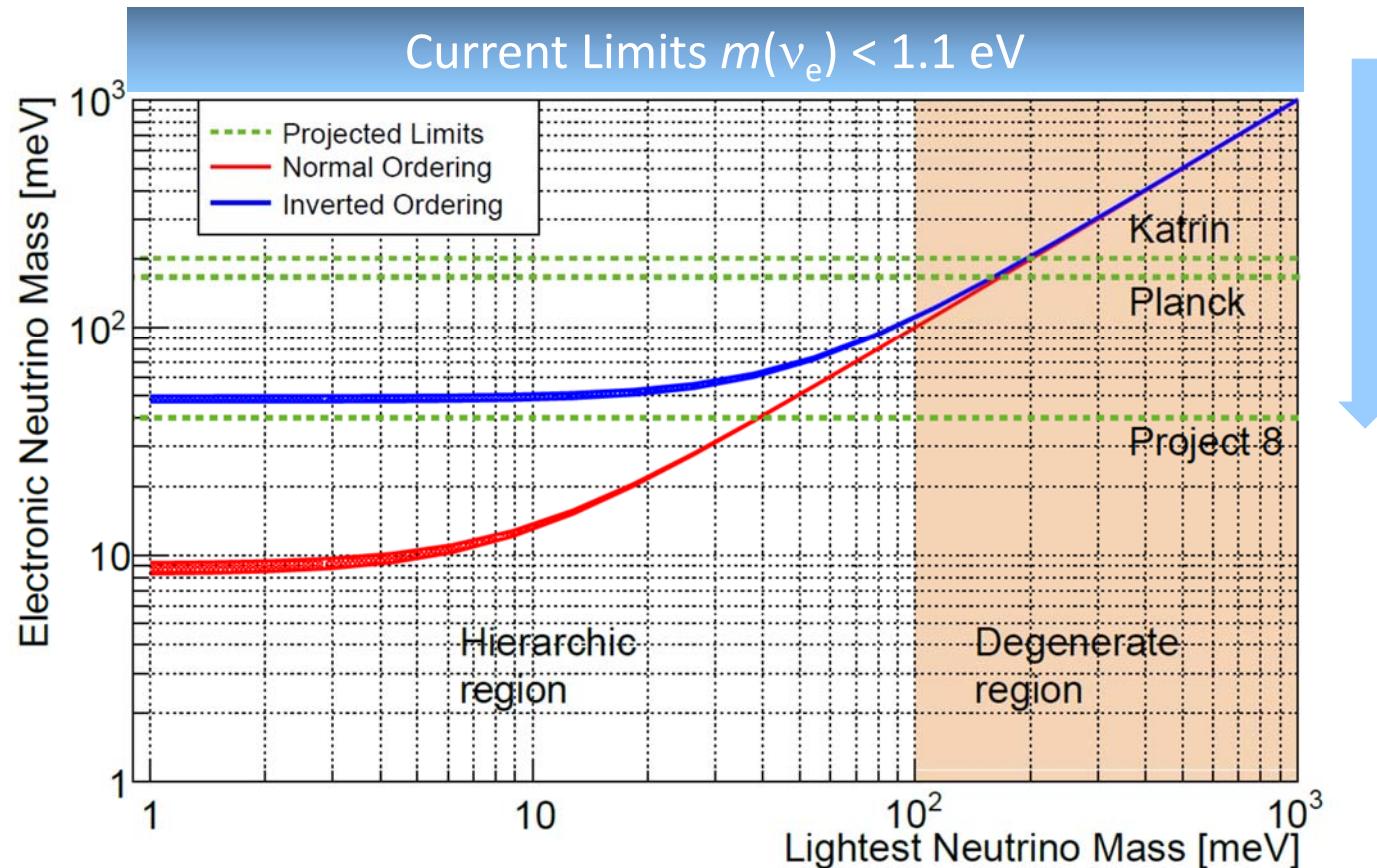


C. Velte et al., (The ECHo Collaboration)

High-resolution and low-background Ho spectrum: interpretation of the resonance tails, *Eur. Phys. J. C* **79** (2019) 1026

Beta decay and electron capture

Ahtari et al., *JPhysG* **44** (2017) 5



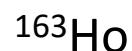
$$m(\bar{\nu}_e) < 1.1 \text{ eV}$$



Aker et al., (The KATRIN Collaboration)

Improved Upper Limit on the Neutrino Mass from a Direct Kinematic Method by KATRIN, *Phys. Rev. Lett.* **123** (2019) 221802

$$m(\nu_e) < 150 \text{ eV}$$



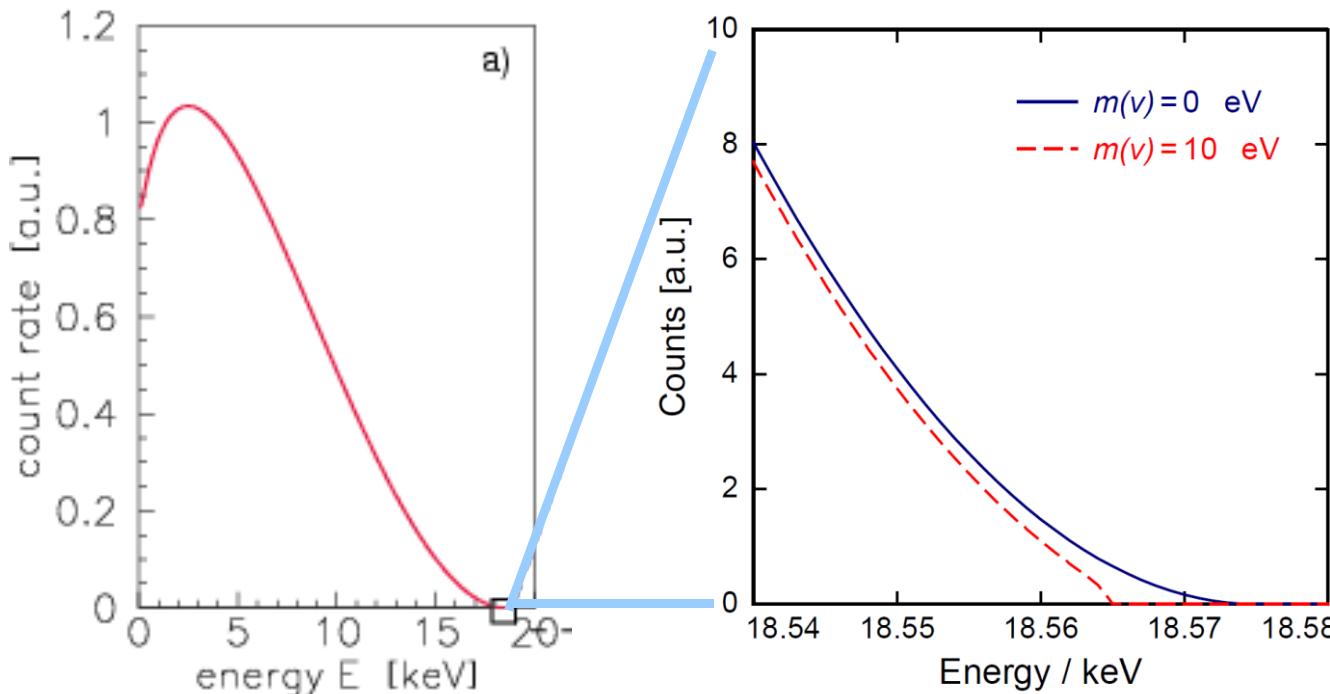
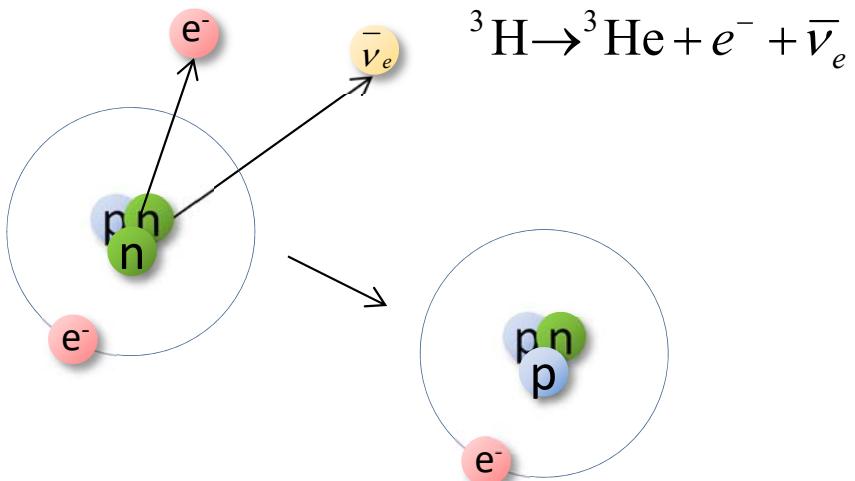
C. Velte et al., (The ECHo Collaboration)

High-resolution and low-background Ho spectrum: interpretation of the resonance tails, *Eur. Phys. J. C* **79** (2019) 1026

Experimental approaches



Beta decay of ${}^3\text{H}$



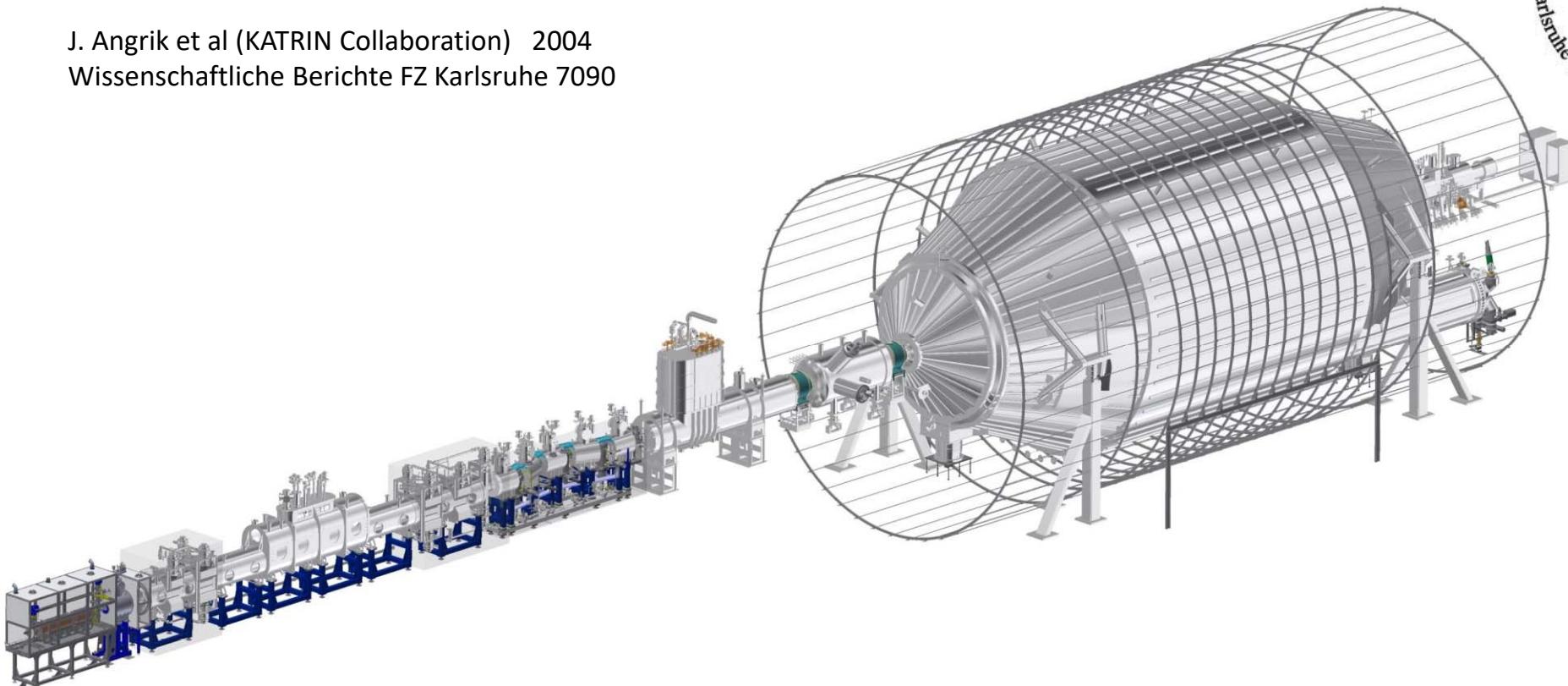
Only a small fraction of events
in the last eV below the endpoint:
 $2 * 10^{-13}$

Tritium is present as
bi-atomic molecules

^3H based experiments

The Karlsruhe Tritium Neutrino Experiment – KATRIN

J. Angrik et al (KATRIN Collaboration) 2004
Wissenschaftliche Berichte FZ Karlsruhe 7090

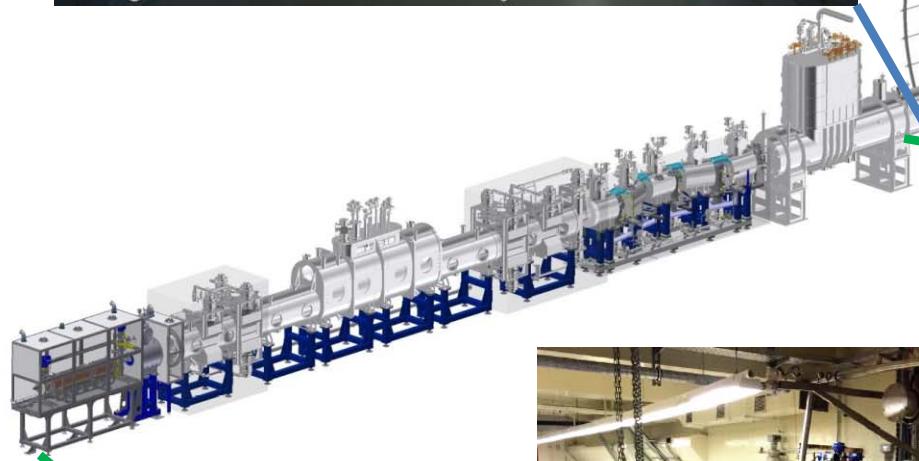
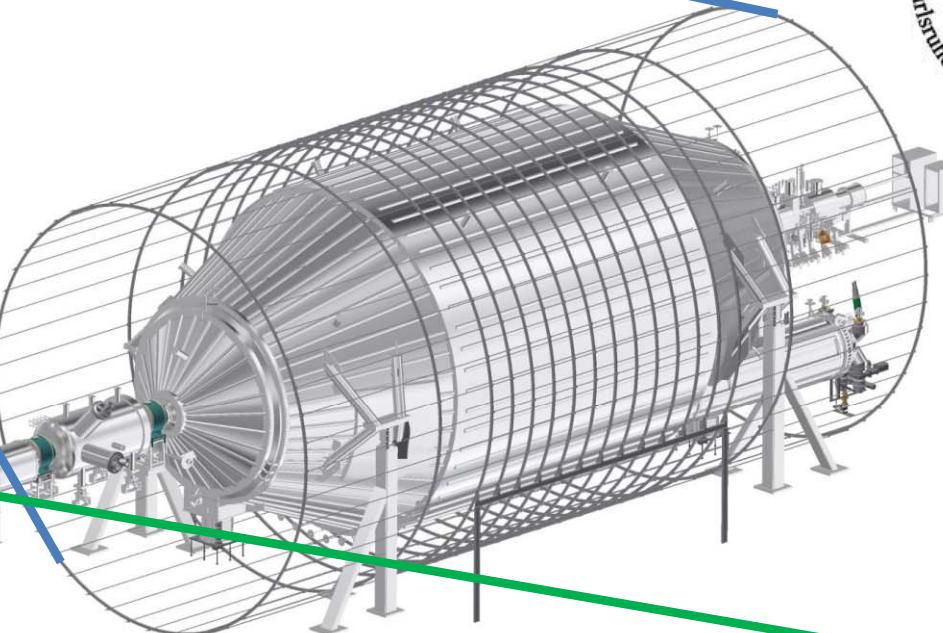


Main ideas:

- high activity source $10^{11} \text{ e}^-/\text{s}$
- high resolution MAC-E* filter to select electrons close to the end point
- count electrons as function of retarding potential
→ integral spectrum

*MAC-E: Magnetic Adiabatic Collimation with Electrostatic Filter

The KATRIN experiment: present status



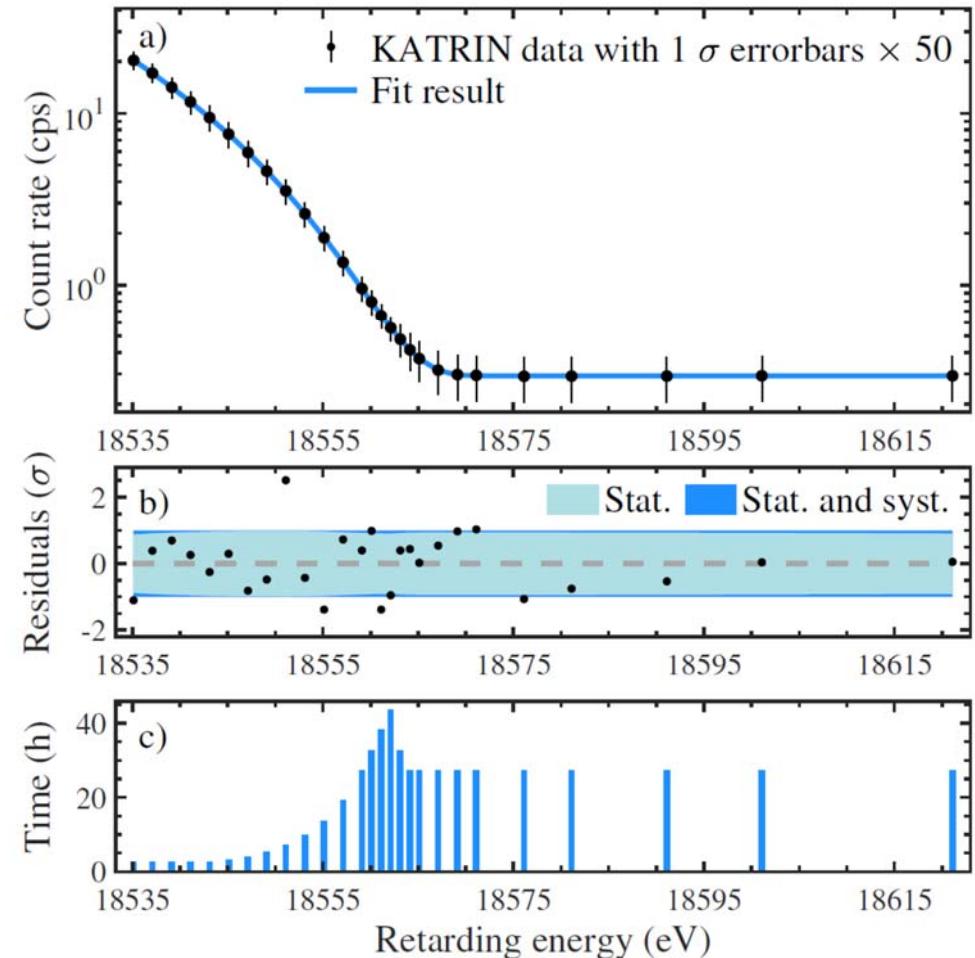
The KATRIN experiment: recent results

Neutrino Mass campaign #1

- April 10 – May 13 2019 → 780 h
- 22% nominal activity → $2.45 \cdot 10^{10}$ Bq
- Integral spectrum in the energy range [Q - 40 eV; Q + 50 eV] → $2 \cdot 10^6$ events in 90-eV-wide interval

Fit of the integral spectrum

- excellent goodness-of-fit $\chi^2 = 21.4$ for 23 d.o.f. (p-value = 0.56)



The KATRIN experiment: recent results

Two independent analysis methods

- covariance matrix
- MC propagation

→ both methods agree to a few percent

Best fit results - ν -mass and E_0

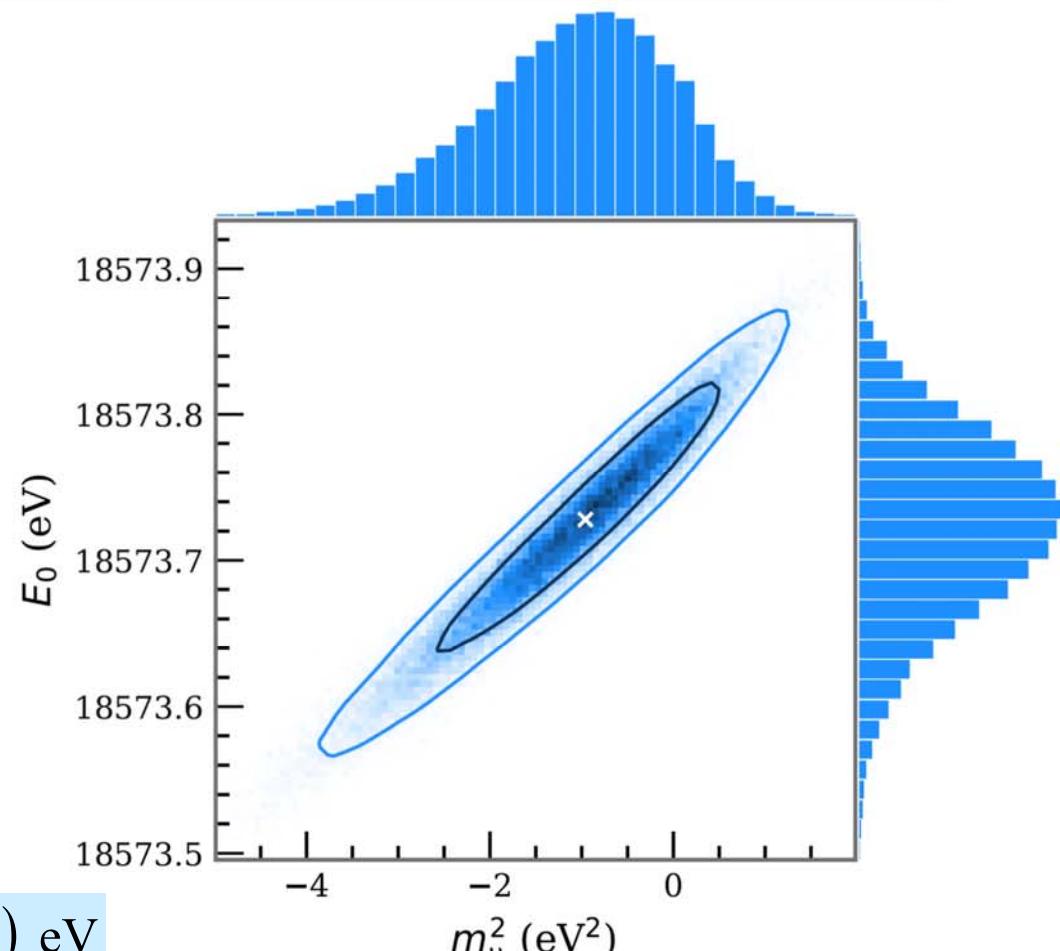
$$m^2(\nu_e) = (-1.0^{+0.9}_{-1.1}) \text{ eV}^2 \quad (90\% \text{ CL})$$

$$E_0 = (18573.7 \pm 0.1) \text{ eV} \Rightarrow Q = (18575.2 \pm 0.5) \text{ eV}$$

To be compared to $Q[\Delta M(^3H, ^3He)] = (18575.72 \pm 0.07) \text{ eV}$

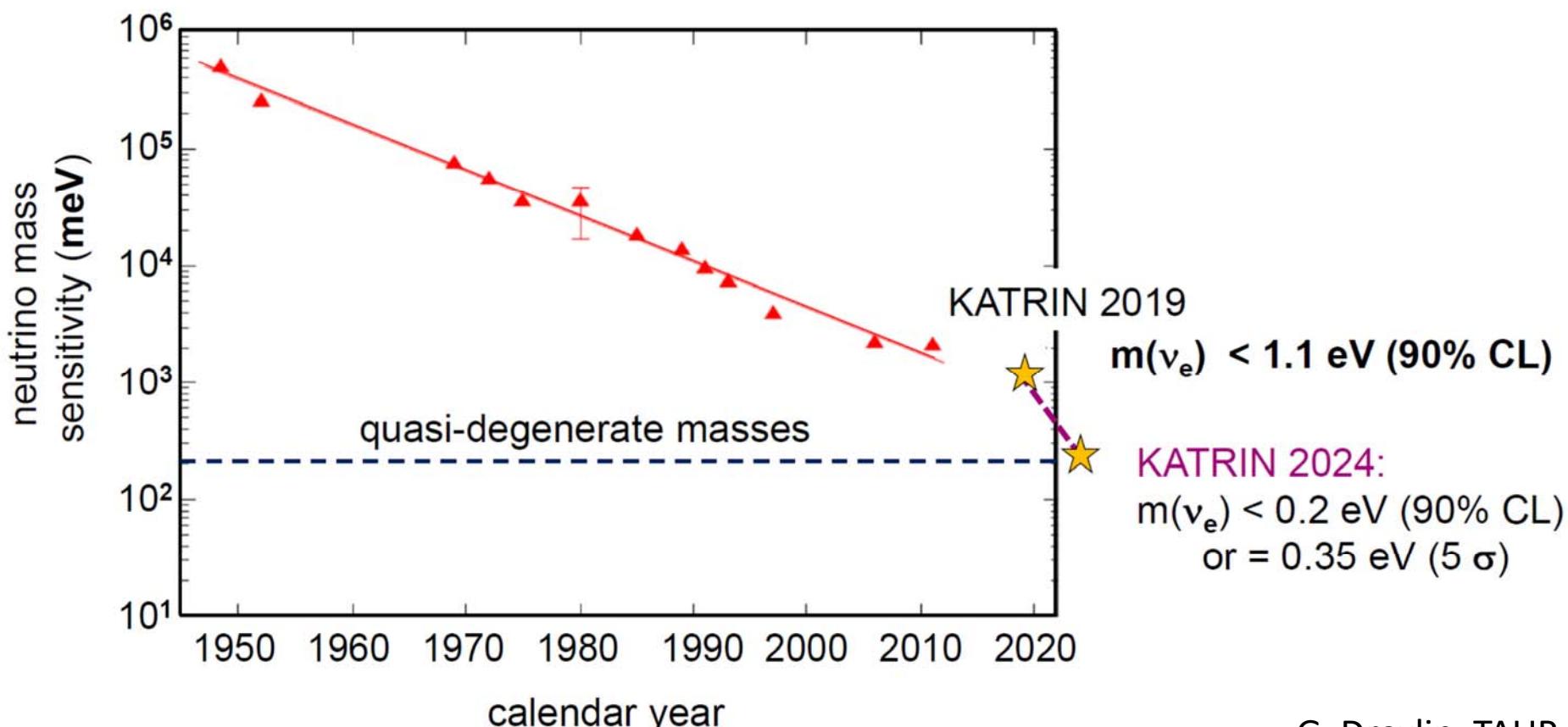
Upper limit on the effective electron neutrino mass:

$$m(\nu_e) < 1.1 \text{ eV} \quad (90\% \text{ CL})$$



The KATRIN experiment: future plans

- Reduction of background and systematics
- 1000 days of measurements at nominal ρd ($5 \cdot 10^{17}$ molecules cm^{-2}) corresponding to 3 tritium campaigns (65 days each) per calendar year over the next 5 years



³H based experiments

❖ KATRIN - Karlsruhe Tritium Neutrino Experiment

Main ideas: - high activity source: 10^{11} e⁻/s

- high resolution MAC-E filter to select electrons close to the end point

- count electrons as function of retarding potential

→ integral spectrum

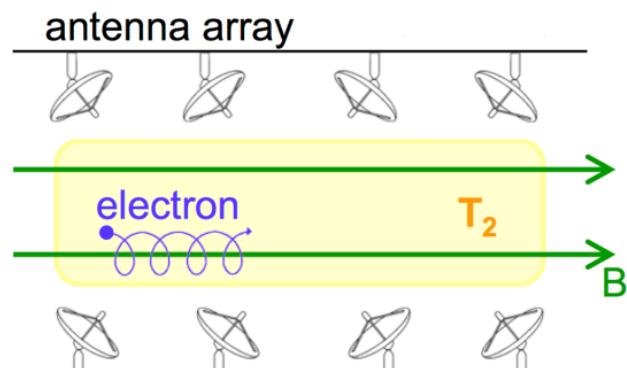


❖ Project8

Main ideas: - Source = detector: $10^{11} - 10^{13}$ ³T₂ molecules /cm³

- Use cyclotron frequency to extract electron energy

- Differential spectrum



$$\omega_\gamma = \frac{\omega_0}{\gamma} = \frac{eB}{K + m_e}$$

@ 1 Tesla $\omega(18 \text{ keV}) \sim 26 \text{ GHz}$
 $P(18 \text{ keV}) = 1.2 \text{ fW}$

Project 8

Phase I (2010 - 2016)

Demonstration of the CRES method
Conversion electron lines from ^{83m}Kr source

Phase II (2015 - 2019)

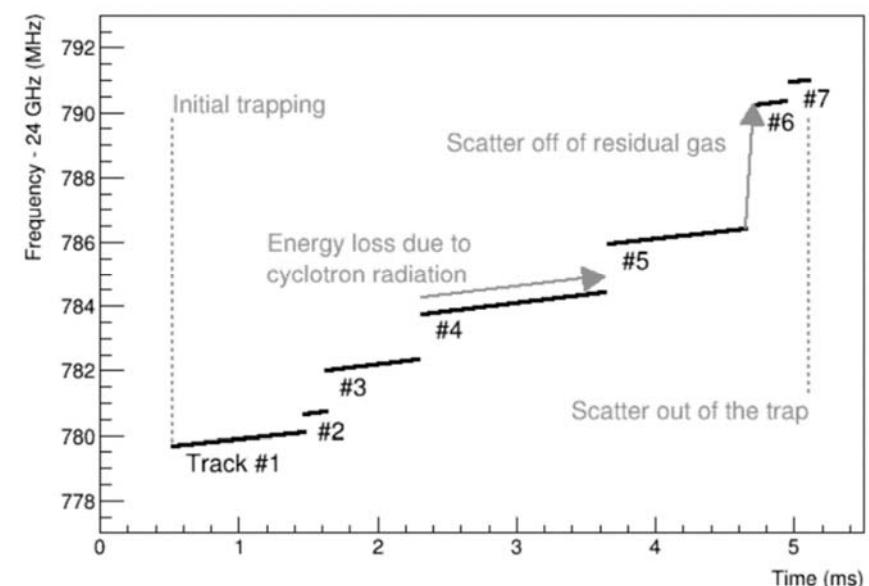
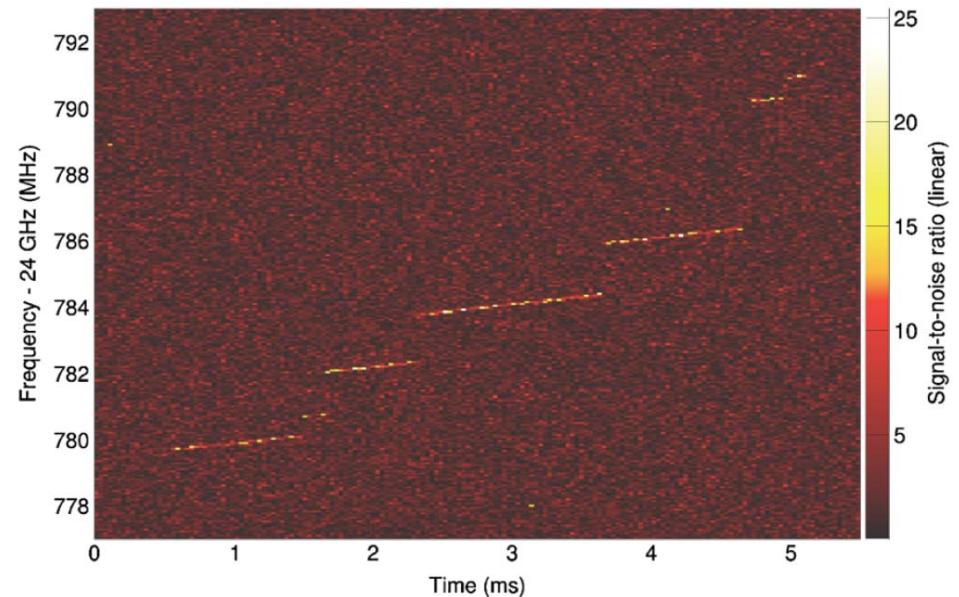
Spectroscopy of continuous T_2 spectrum
Study of systematics
improvement of the energy resolution

Phase III (2016 - 2023)

200 cm³ effective source volume (1 year)
Phased array antenna
Sensitivity goal: $m(v_e) < 2 \text{ eV } 90\% \text{ C.L.}$

Phase IV (2022 +)

Large scale exp. with atomic ^3H source for
Sensitivity goal: $m(v_e) < 40 \text{ meV } 90\% \text{ C.L.}$



Project 8

Phase I (2010 - 2016)

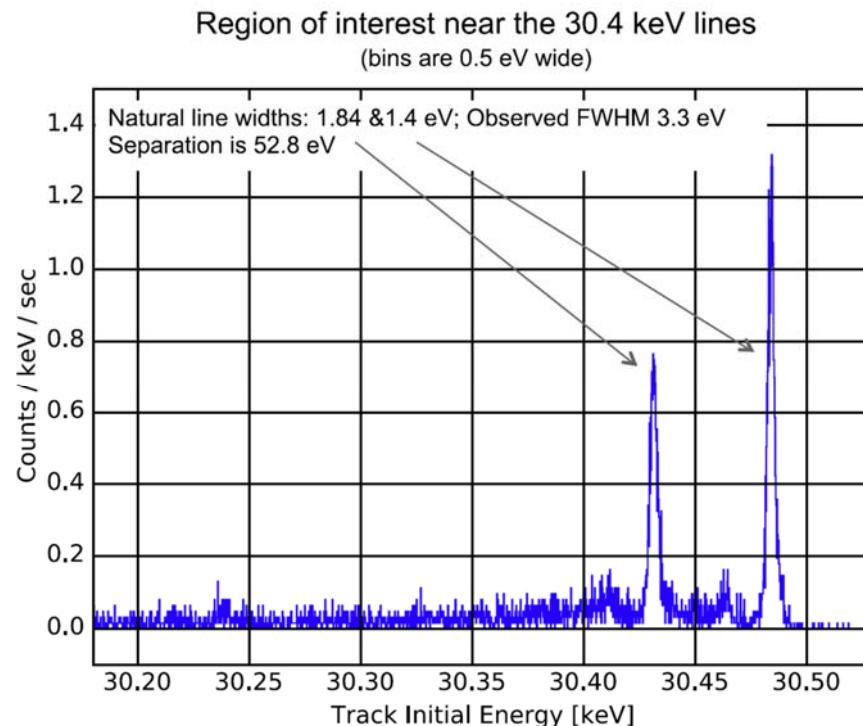
Demonstration of the CRES method
Conversion electron lines from ^{83m}Kr source

Phase II (2015 - 2019)

Spectroscopy of continuous T_2 spectrum
Study of systematics
improvement of the energy resolution

Phase III (2016 - 2023)

200 cm³ effective source volume (1 year)
Phased array antenna
Sensitivity goal: $m(v_e) < 2 \text{ eV} 90\% \text{ C.L.}$



Phase IV (2022 +)

Large scale exp. with atomic ^3H source for
Sensitivity goal: $m(v_e) < 40 \text{ meV} 90\% \text{ C.L.}$

Project 8

Phase I (2010 - 2016)

Demonstration of the CRES method
Conversion electron lines from ^{83m}Kr source

Phase II (2015 - 2019)

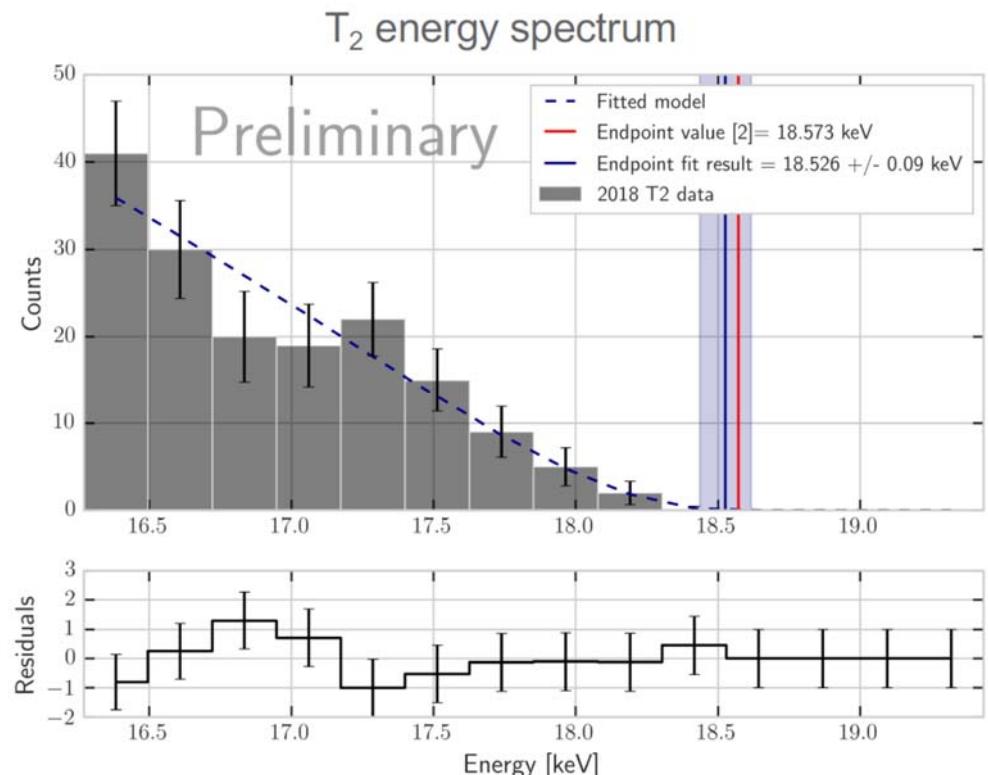
Spectroscopy of continuous T_2 spectrum
Study of systematics
improvement of the energy resolution

Phase III (2016 - 2023)

200 cm³ effective source volume (1 year)
Phased array antenna
Sensitivity goal: $m(v_e) < 2 \text{ eV} 90\% \text{ C.L.}$

Phase IV (2022 +)

Large scale exp. with atomic ^3H source for
Sensitivity goal: $m(v_e) < 40 \text{ meV} 90\% \text{ C.L.}$



Project 8

Phase I (2010 - 2016)

Demonstration of the CRES method
Conversion electron lines from ^{83m}Kr source

Phase II (2015 - 2019)

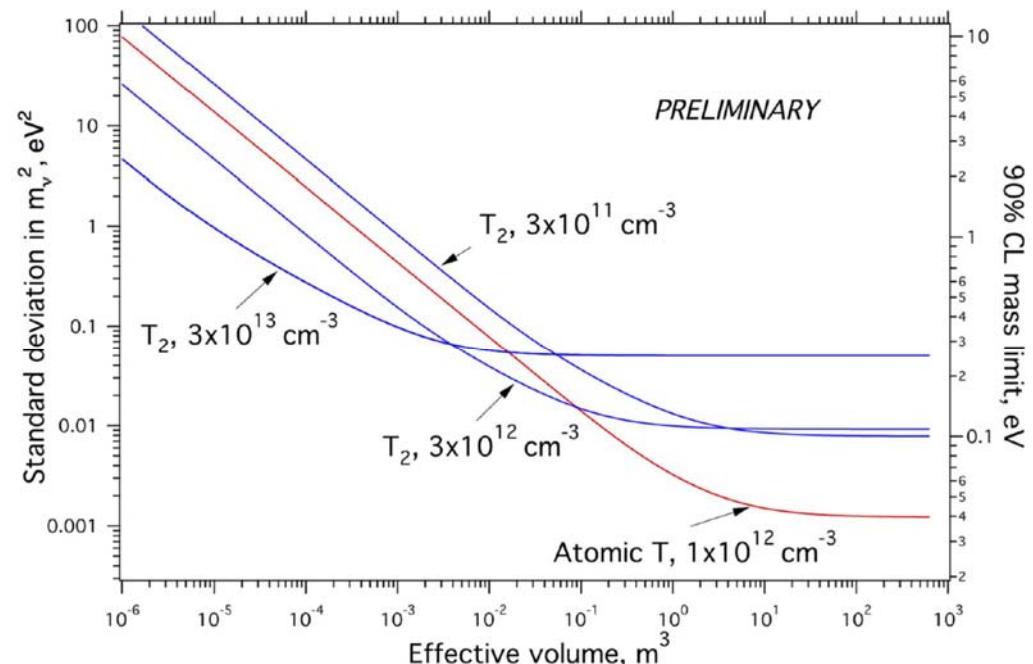
Spectroscopy of continuous T_2 spectrum
Study of systematics
improvement of the energy resolution

Phase III (2016 - 2023)

200 cm³ effective source volume (1 year)
Phased array antenna
Sensitivity goal: $m(v_e) < 2 \text{ eV } 90\% \text{ C.L.}$

Phase IV (2022 +)

Large scale exp. with atomic ^3H source for
Sensitivity goal: $m(v_e) < 40 \text{ meV } 90\% \text{ C.L.}$



^3H based experiments

❖ KATRIN - Karlsruhe Tritium Neutrino Experiment

Main ideas: - high activity source: $10^{11} \text{ e}^-/\text{s}$

- high resolution MAC-E filter to select electrons close to the end point
- count electrons as function of retarding potential
→ integral spectrum



❖ Project8

Main ideas: - Source = detector: $10^{11} - 10^{13} \text{ }^3\text{H}_2 \text{ molecules/cm}^3$

- Use cyclotron frequency to extract electron energy
- Differential spectrum



❖ PTOLEMY - Princeton Tritium Observatory for Light, Early-Universe, Massive-Neutrino Yield

Main ideas: large area tritium source: 100 g atomic ^3H

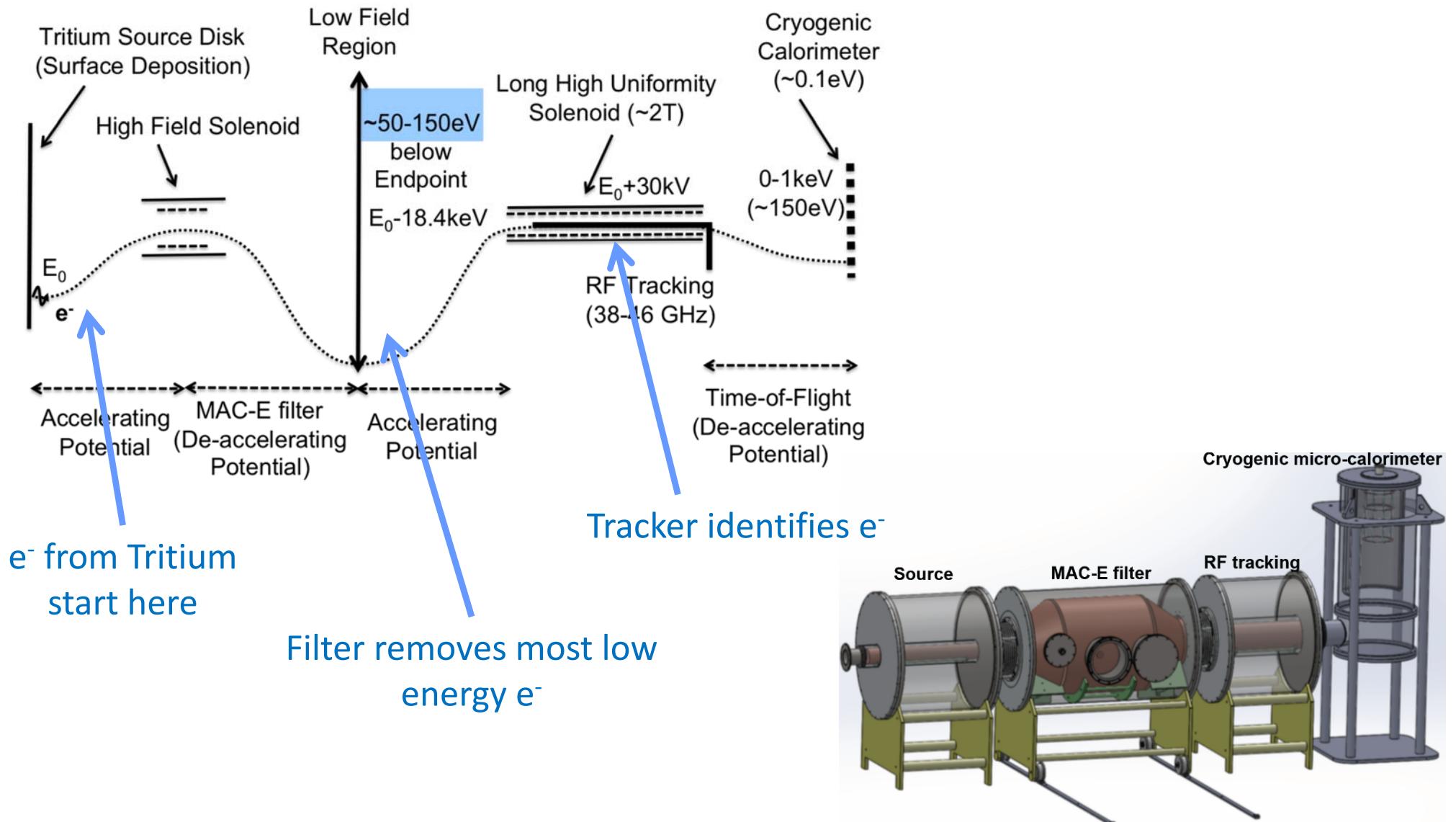
MAC-E Iter methods

RF tracking and time-of-flight systems

cryogenic calorimetry → differential spectrum



PTOLEMY



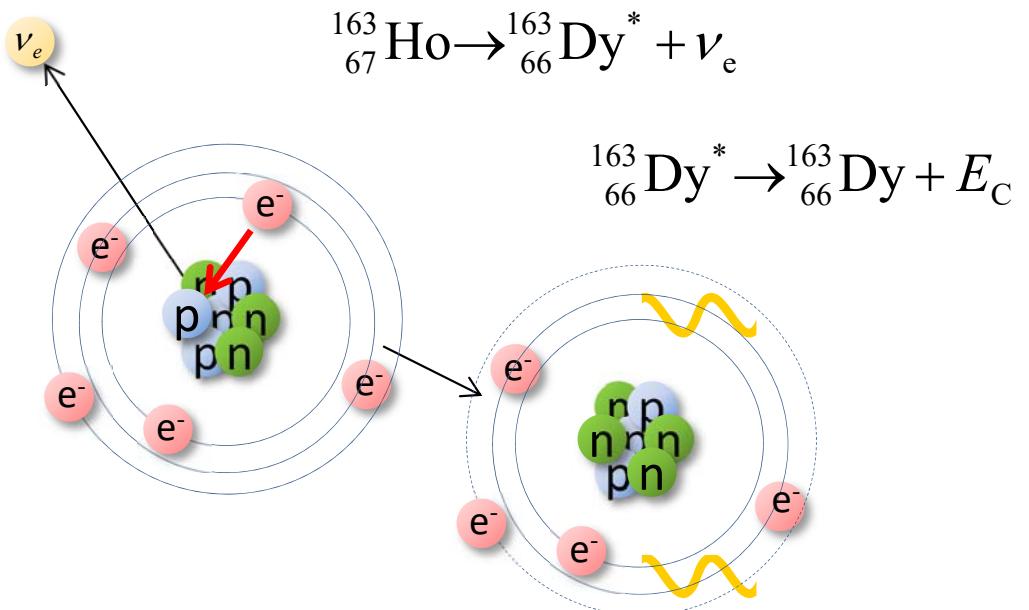
S. Betts et al., [astro-ph.IM] arXiv:1307.4738v2

Ptolemy Collaboration JCAP 07 (2019) 047

Electron capture in ^{163}Ho

Atomic de-excitation:

- X-ray emission
- Auger electrons
- Coster-Kronig transitions



- $\tau_{1/2} \cong 4570 \text{ years}$ (2×10^{11} atoms for 1 Bq)

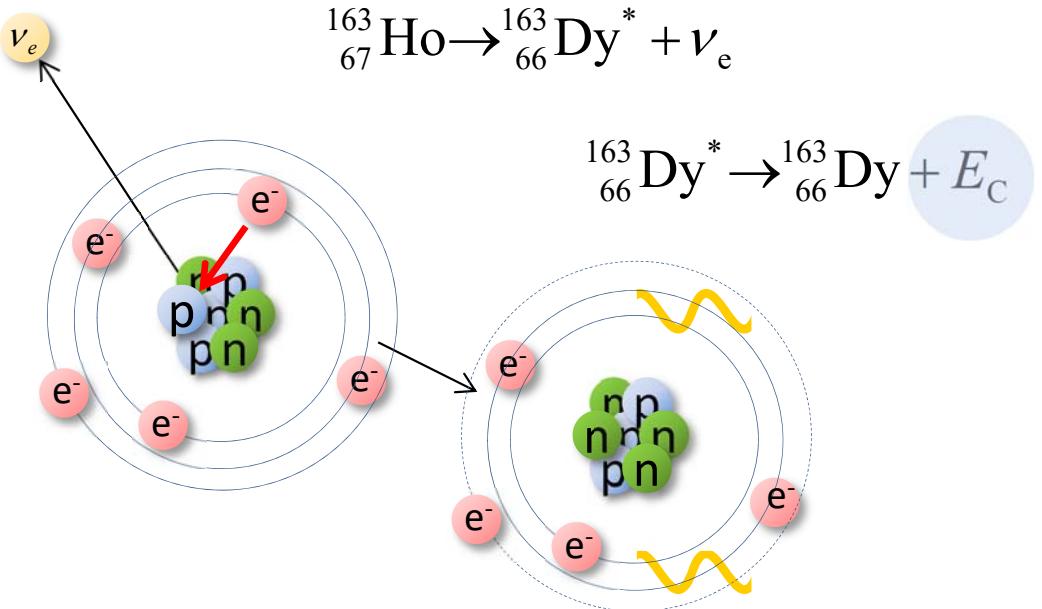
- $Q_{EC} = (2.833 \pm 0.030^{\text{stat}} \pm 0.015^{\text{syst}}) \text{ keV}$

S. Eliseev et al., *Phys. Rev. Lett.* **115** (2015) 062501

Electron capture in ^{163}Ho

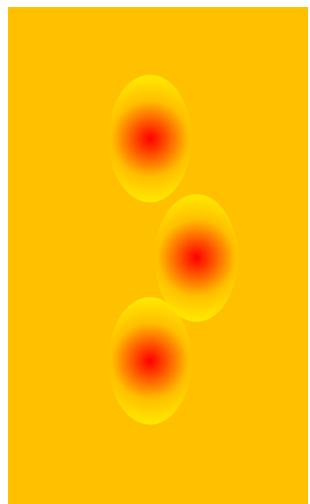
Atomic de-excitation:

- X-ray emission
- Auger electrons
- Coster-Kronig transitions



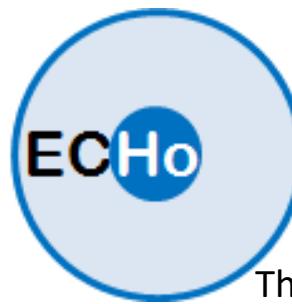
Calorimetric measurement

A. De Rujula and M. Lusignoli, *Phys. Lett.* **118B** (1982)



ν_e

ν_e



The ECHO Collaboration EPJ-ST 226 8 (2017) 1623



Source = Detector

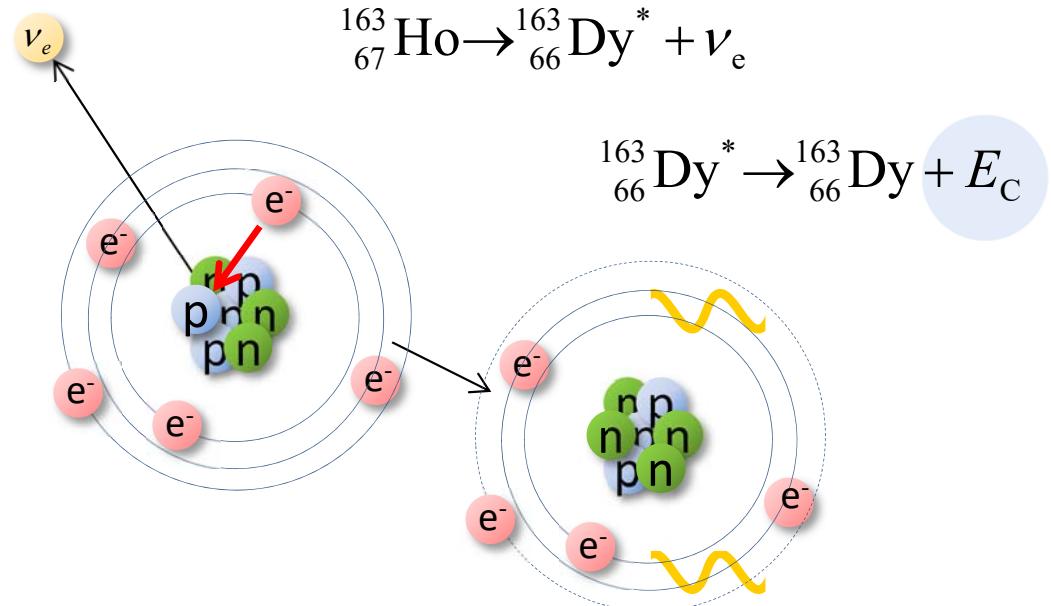
ν_e

B. Alpert et al, Eur. Phys. J. C 75 (2015) 112

Electron capture in ^{163}Ho : spectrum

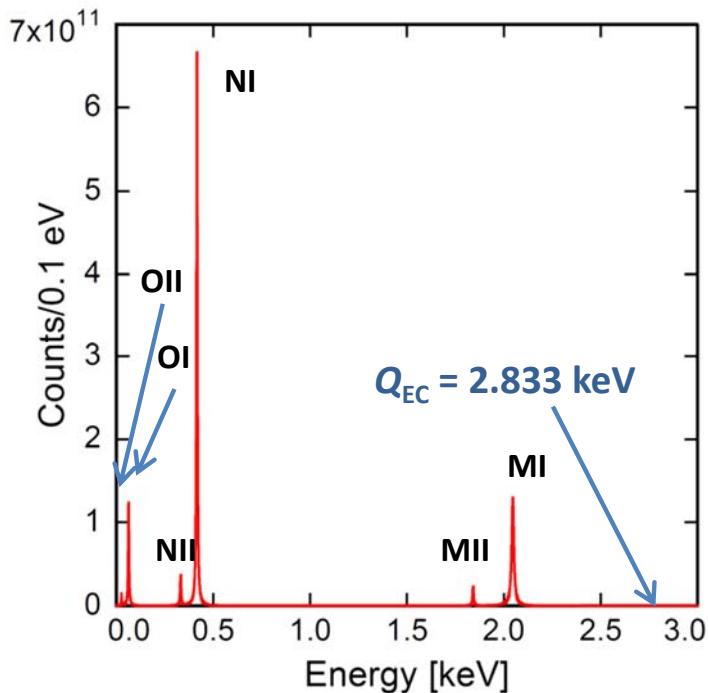
Atomic de-excitation:

- X-ray emission
- Auger electrons
- Coster-Kronig transitions



Calorimetric measurement

A. De Rujula and M. Lusignoli, *Phys. Lett.* **118B** (1982)

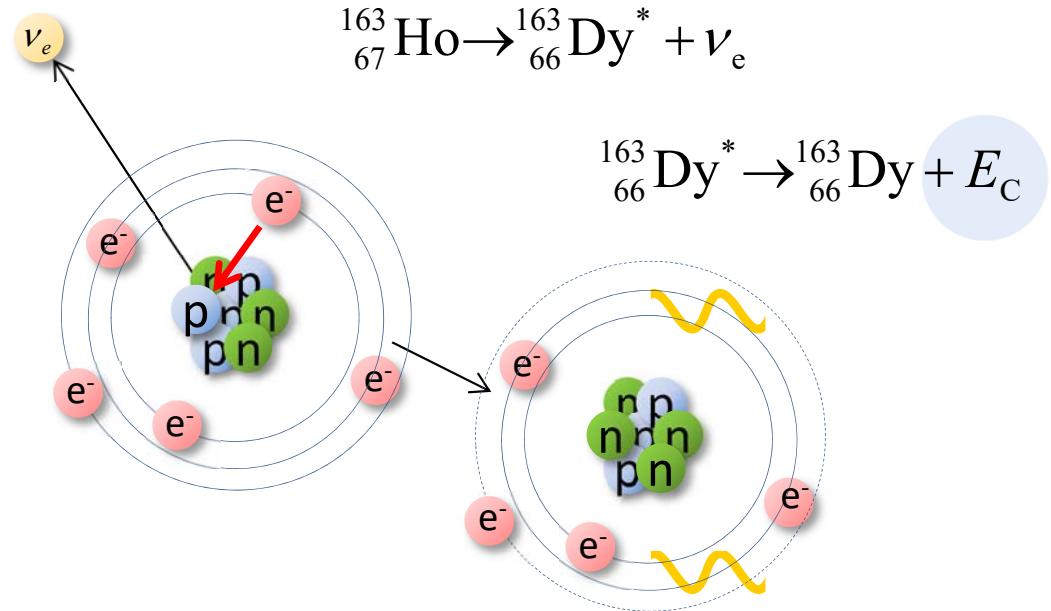


$$\frac{dW}{dE_C} = A(Q_{EC} - E_C)^2 \sqrt{1 - \frac{m_\nu^2}{(Q_{EC} - E_C)^2}} \sum_H B_H \varphi_H^2(0) \frac{\frac{\Gamma_H}{2\pi}}{(E_C - E_H)^2 + \frac{\Gamma_H^2}{4}}$$

Electron capture in ^{163}Ho : spectrum

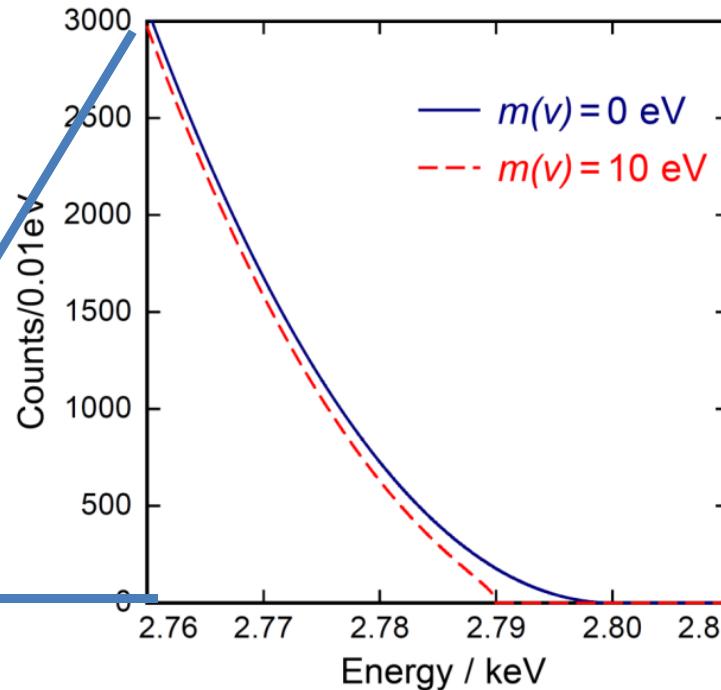
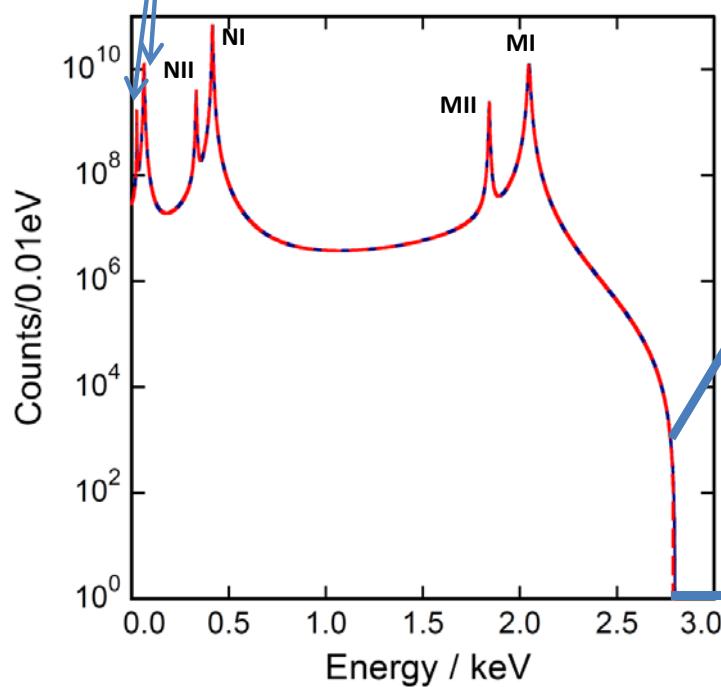
Atomic de-excitation:

- X-ray emission
- Auger electrons
- Coster-Kronig transitions



Calorimetric measurement

A. De Rujula and M. Lusignoli, *Phys. Lett.* **118B** (1982)



Requirements for sub-eV sensitivity in ECHO

Statistics in the end point region

- $N_{\text{ev}} > 10^{14} \rightarrow A \approx 1 \text{ MBq}$

Unresolved pile-up ($f_{\text{pu}} \sim a \cdot \tau_r$)

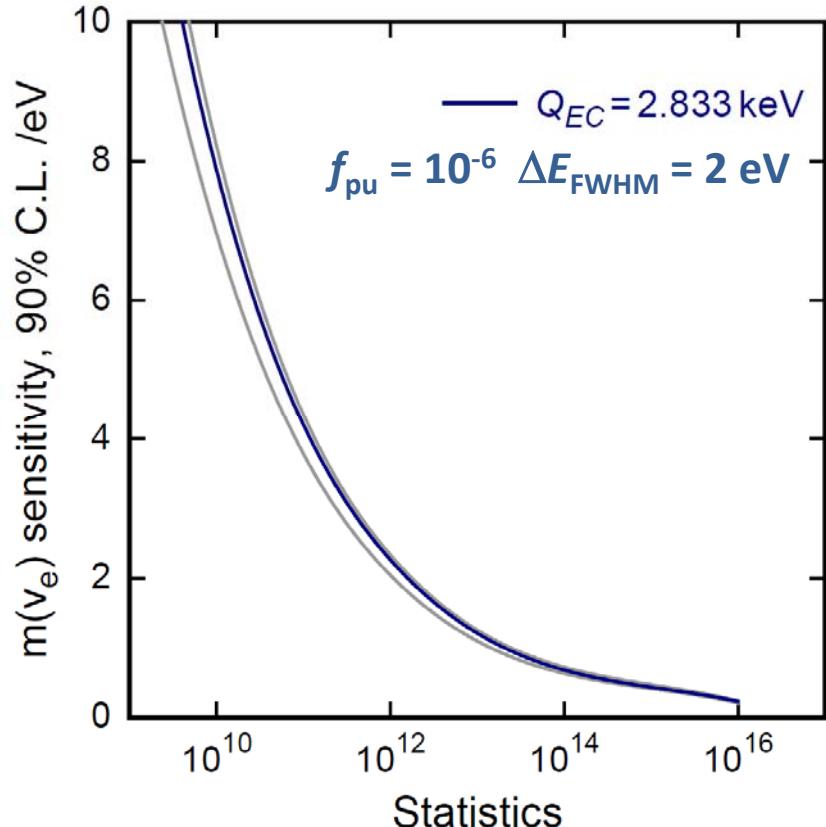
- $f_{\text{pu}} < 10^{-5}$
- $\tau_r < 1 \mu\text{s} \rightarrow a \sim 10 \text{ Bq}$
- 10^5 pixels

Precision characterization of the endpoint region

- $\Delta E_{\text{FWHM}} < 3 \text{ eV}$

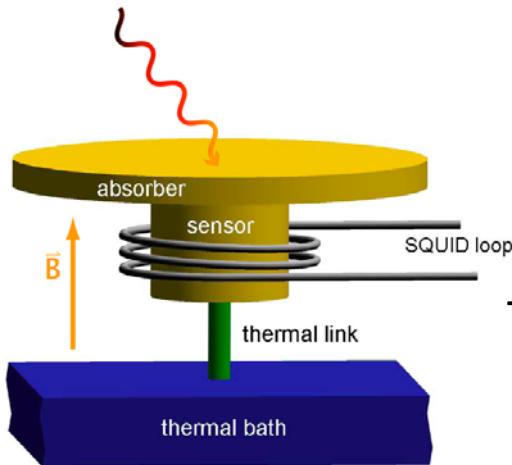
Background level

- $< 10^{-6} \text{ events/eV/det/day}$

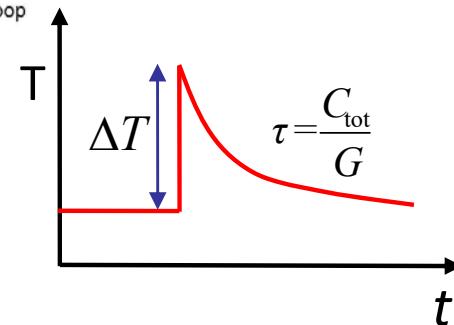


ECHo detectors

ECHo uses large arrays of low T **metallic magnetic calorimeters** with enclosed ^{163}Ho



$$\Delta T \approx \frac{E}{C_{\text{tot}}}$$



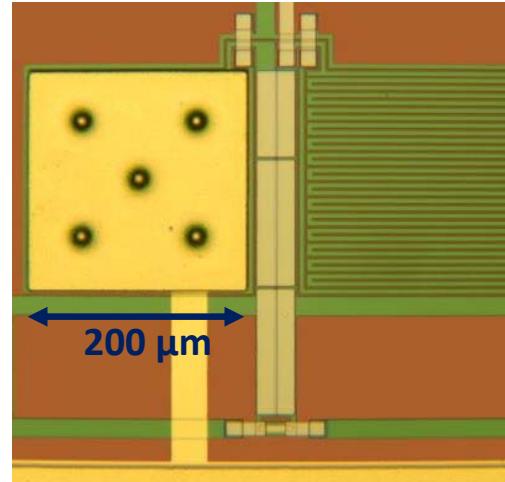
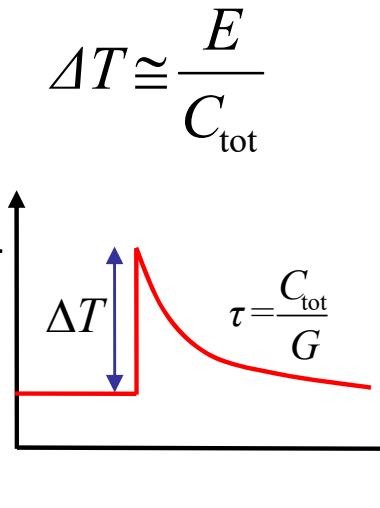
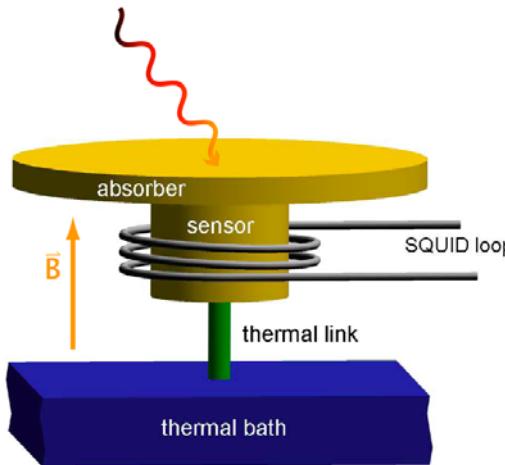
A.Fleischmann, C. Enss and G. M. Seidel,
Topics in Applied Physics **99** (2005) 63

A.Fleischmann et al.,
AIP Conf. Proc. **1185** (2009) 571

L. Gastaldo et al.,
Nucl. Inst. Meth. A, **711** (2013) 1

ECHo detectors

ECHo uses large arrays of low T **metallic magnetic calorimeters** with enclosed ^{163}Ho



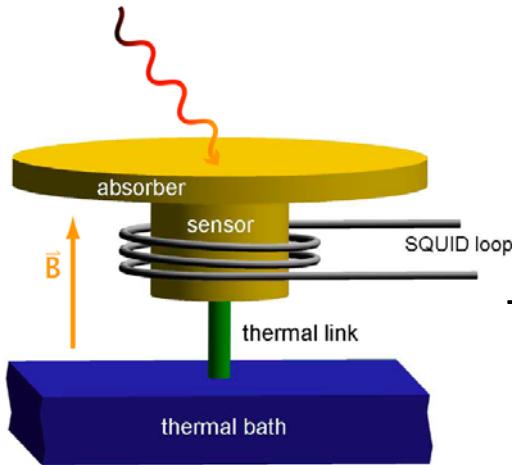
A. Fleischmann, C. Enss and G. M. Seidel,
Topics in Applied Physics **99** (2005) 63

A. Fleischmann et al.,
AIP Conf. Proc. **1185** (2009) 571

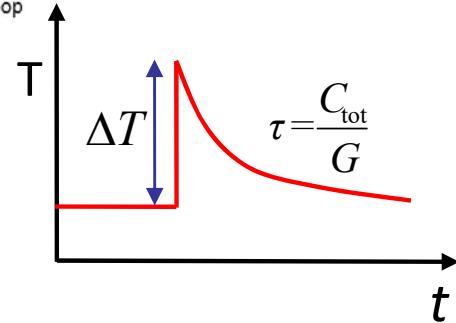
L. Gastaldo et al.,
Nucl. Inst. Meth. A, **711** (2013) 1

ECHO detectors

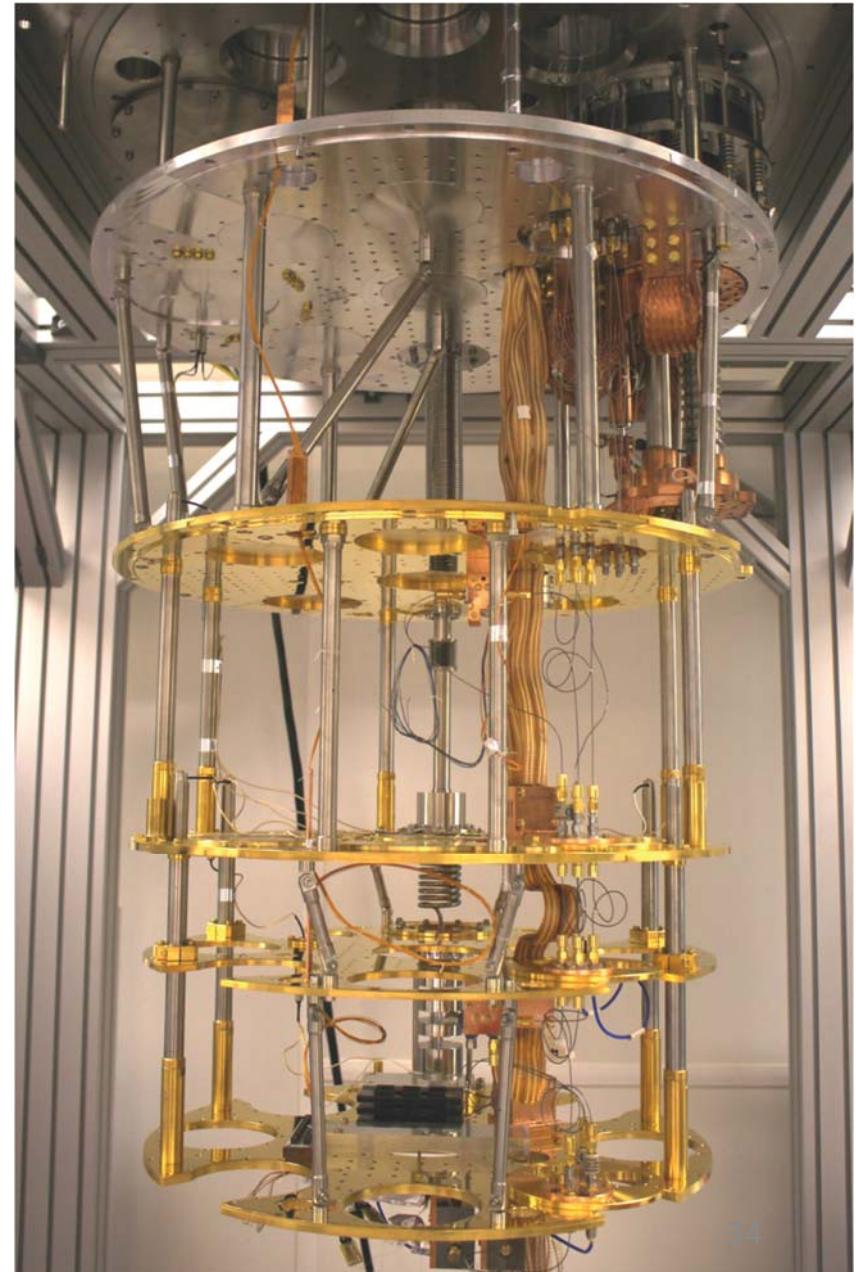
ECHO uses large arrays of low T **metallic magnetic calorimeters** with enclosed ^{163}Ho



$$\Delta T \approx \frac{E}{C_{\text{tot}}}$$

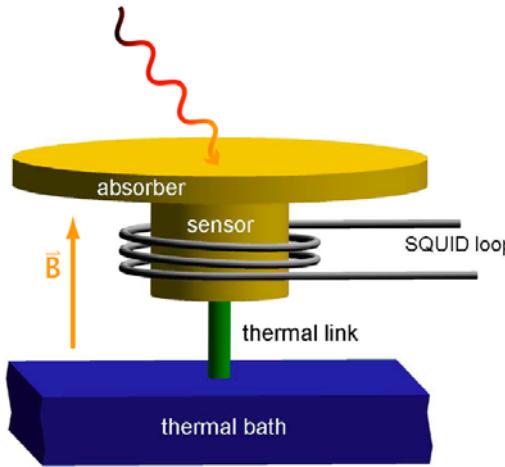


Operated at $T \approx 20 \text{ mK}$

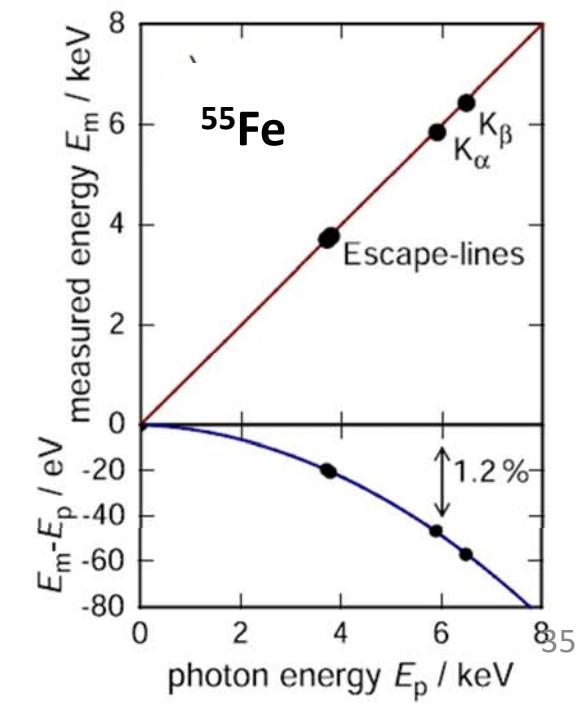
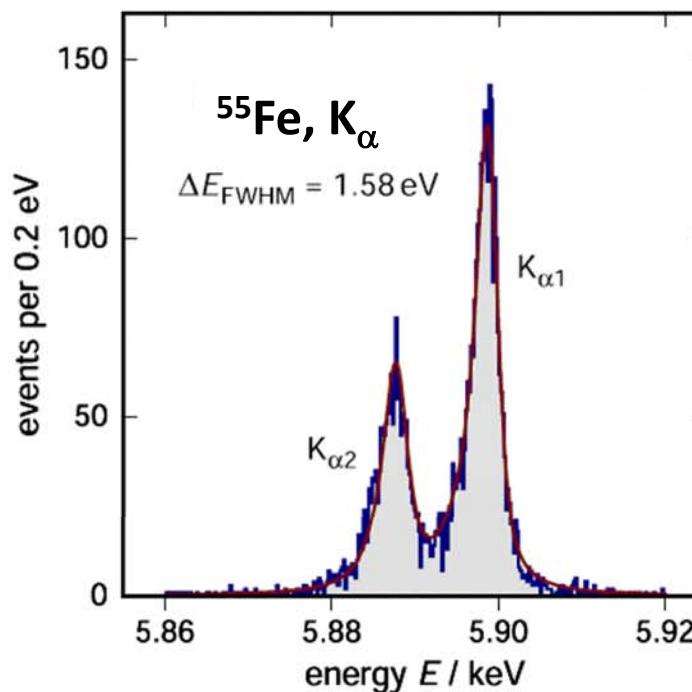
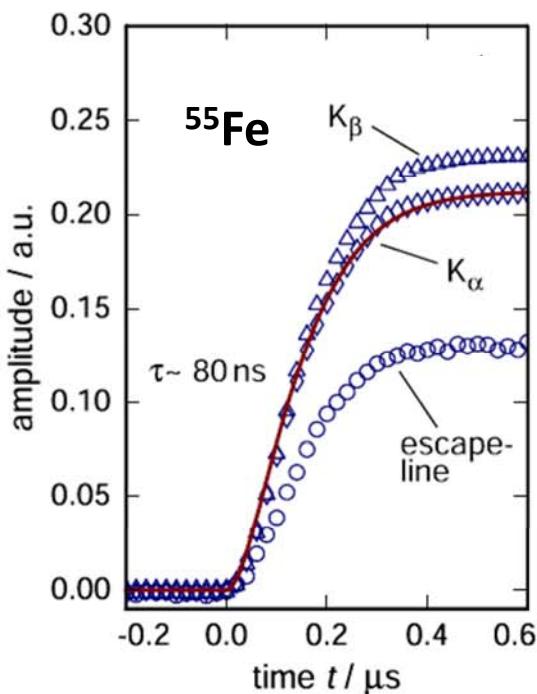
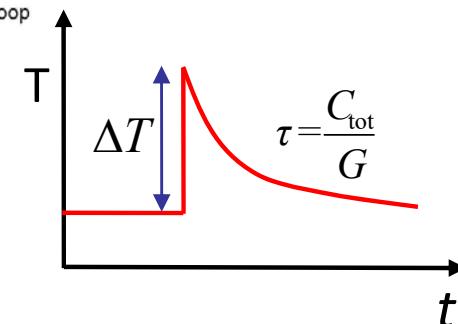


ECHO detectors

ECHO uses large arrays of low T **metallic magnetic calorimeters** with enclosed ^{163}Ho



$$\Delta T \approx \frac{E}{C_{\text{tot}}}$$



Fast risetime

→ Reduction un-resolved pile-up

Extremely good energy resolution

→ Reduced smearing in the end point region

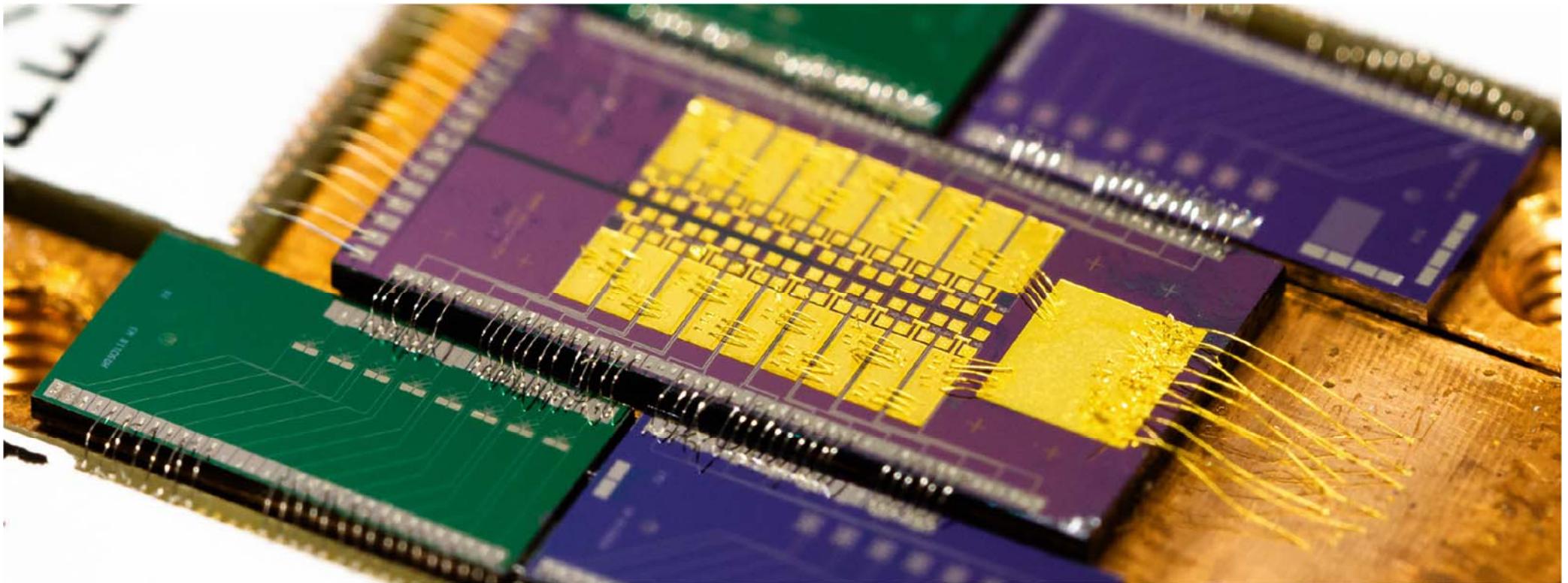
Excellent linearity

→ precise definition of the energy scale

ECHo detectors

ECHo-1k chip-Au

- High purity ^{163}Ho source * → activity per pixel $a \approx 1 \text{ Bq}$
- 4 Front-end chips each with 8 dc-SQUIDs for parallel readout

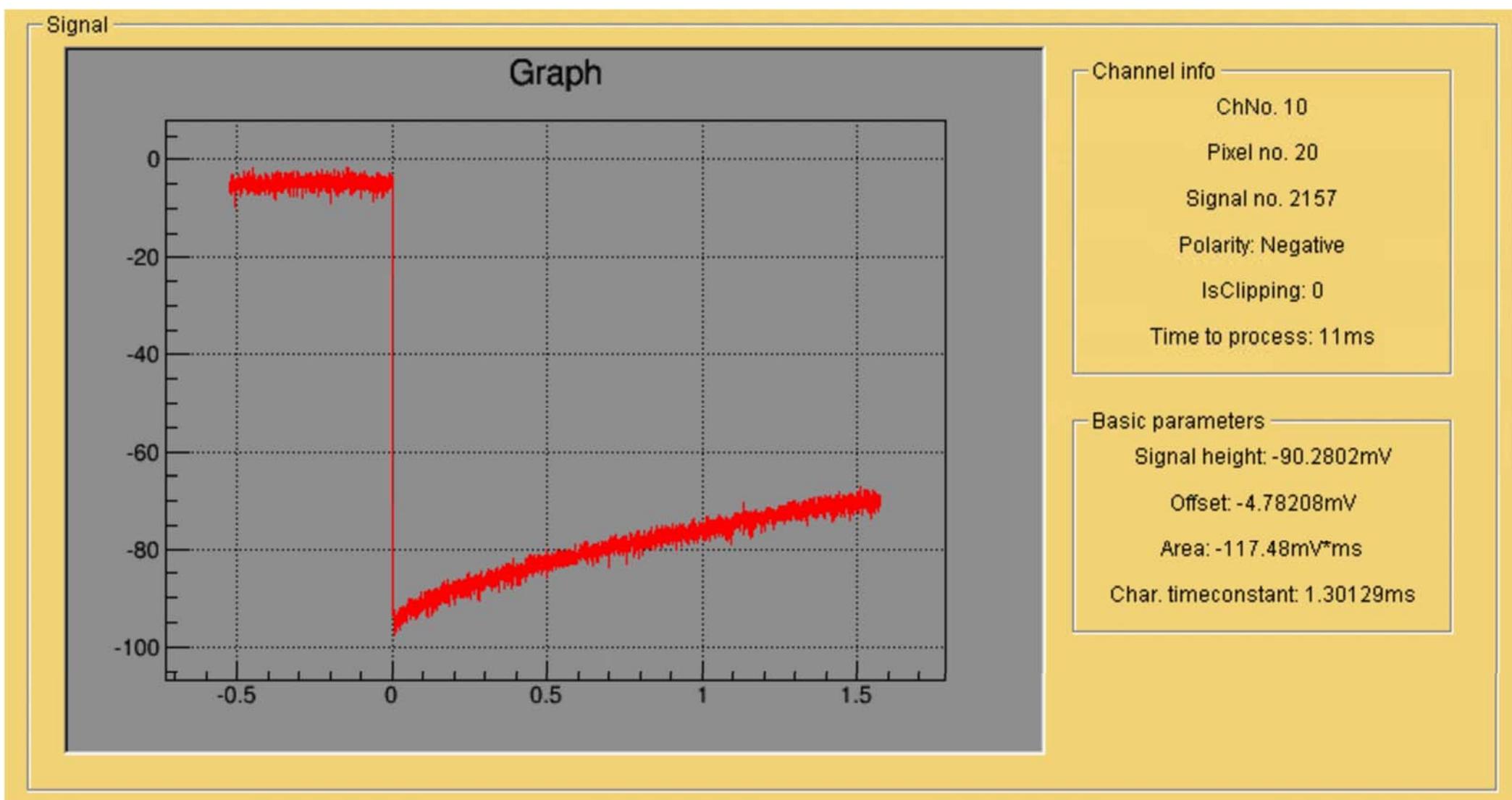


* H. Dorrer et al, Radiochim. Acta (2018)

F. Schneider et al., *NIM B* **376** (2016) 388

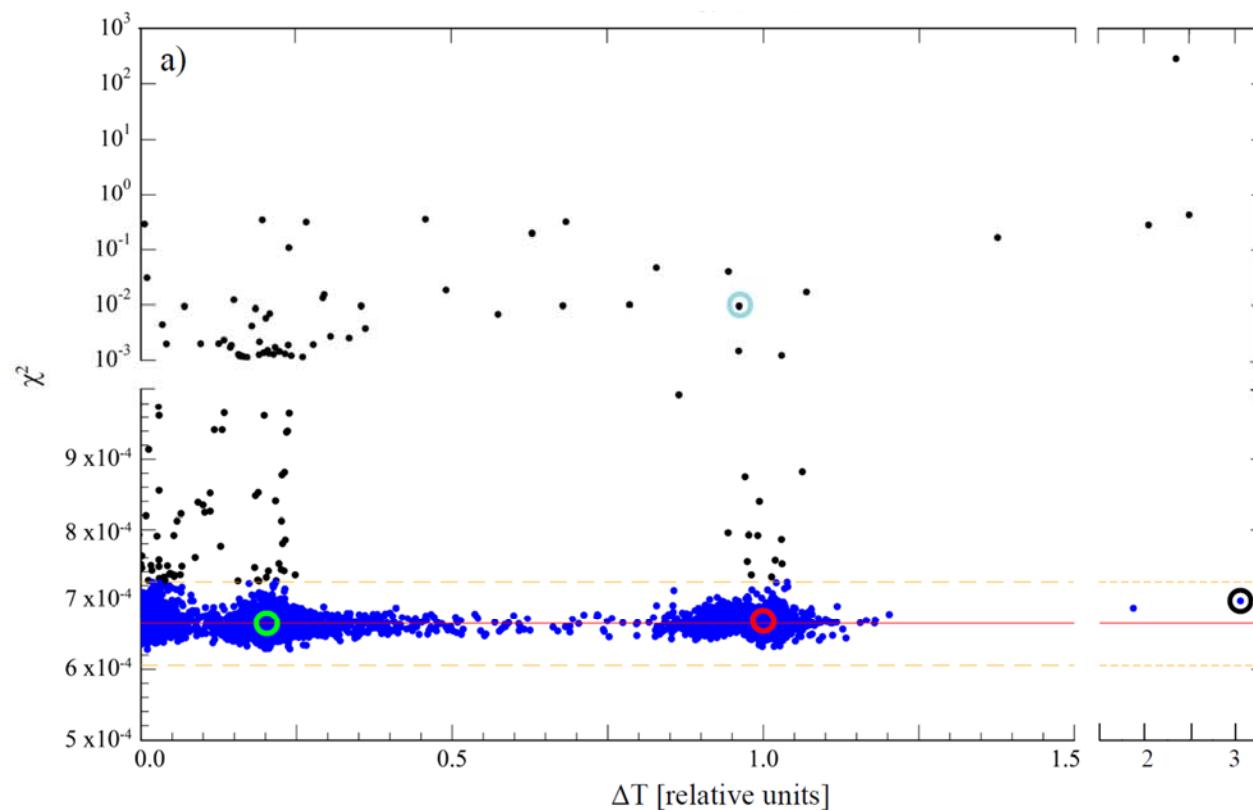
T. Kieck et al., *Rev. Sci. Instrum.* **90** (2019) 053304

ECHO signals

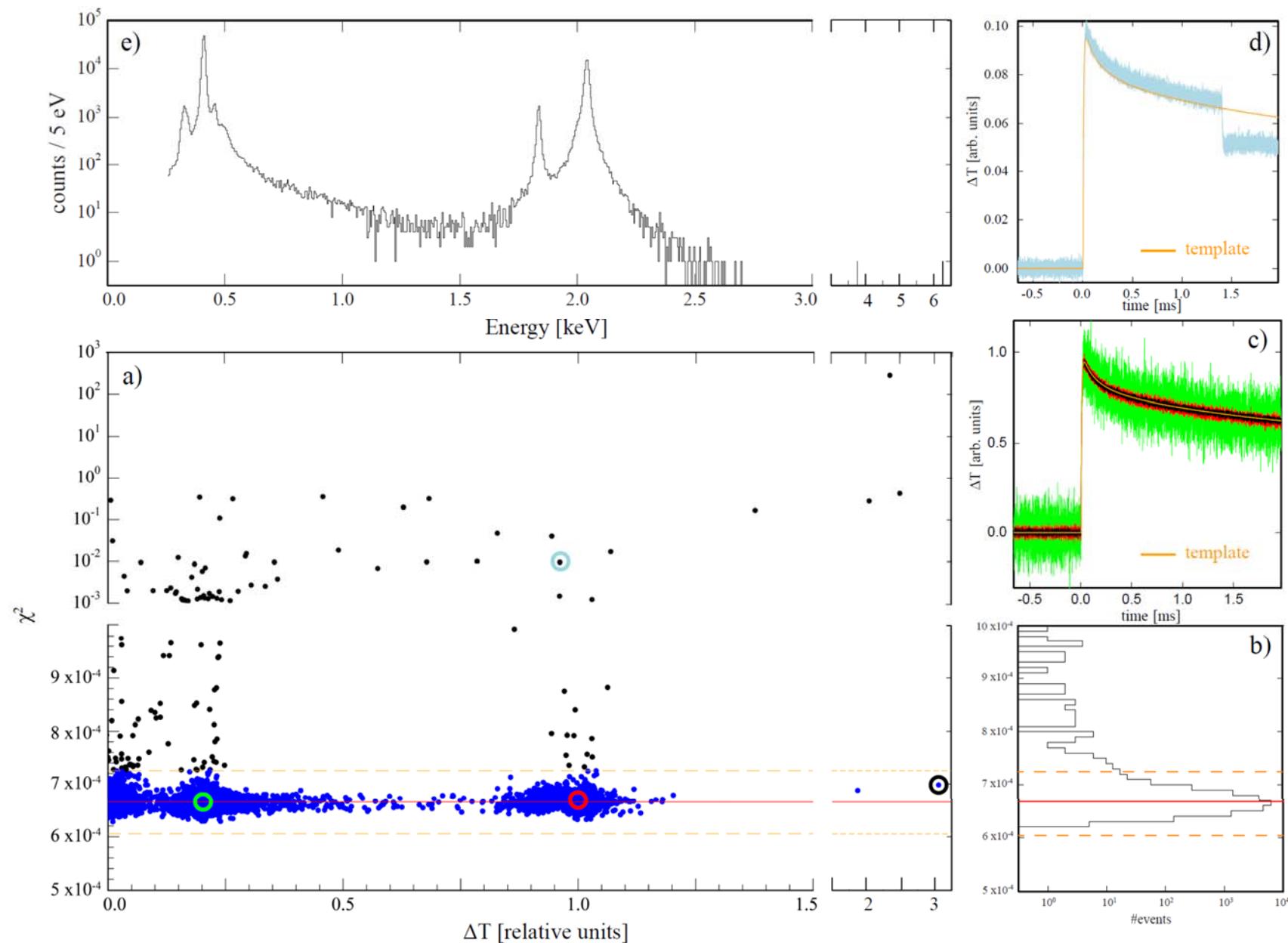


From signals to spectrum

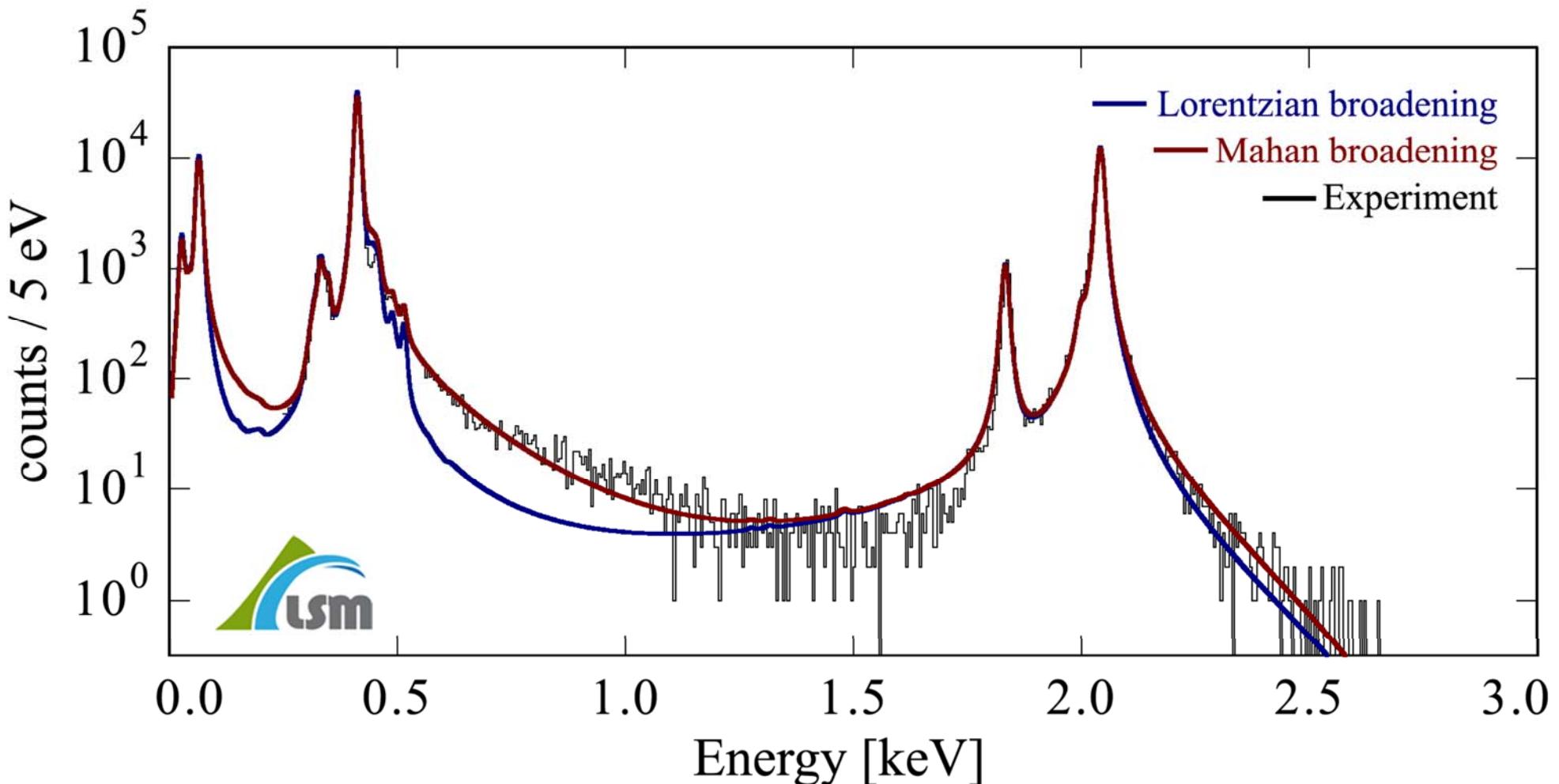
Fit all pulses with time template
key parameters are extracted to perform cuts



From signals to spectrum



^{163}Ho spectral shape analysis

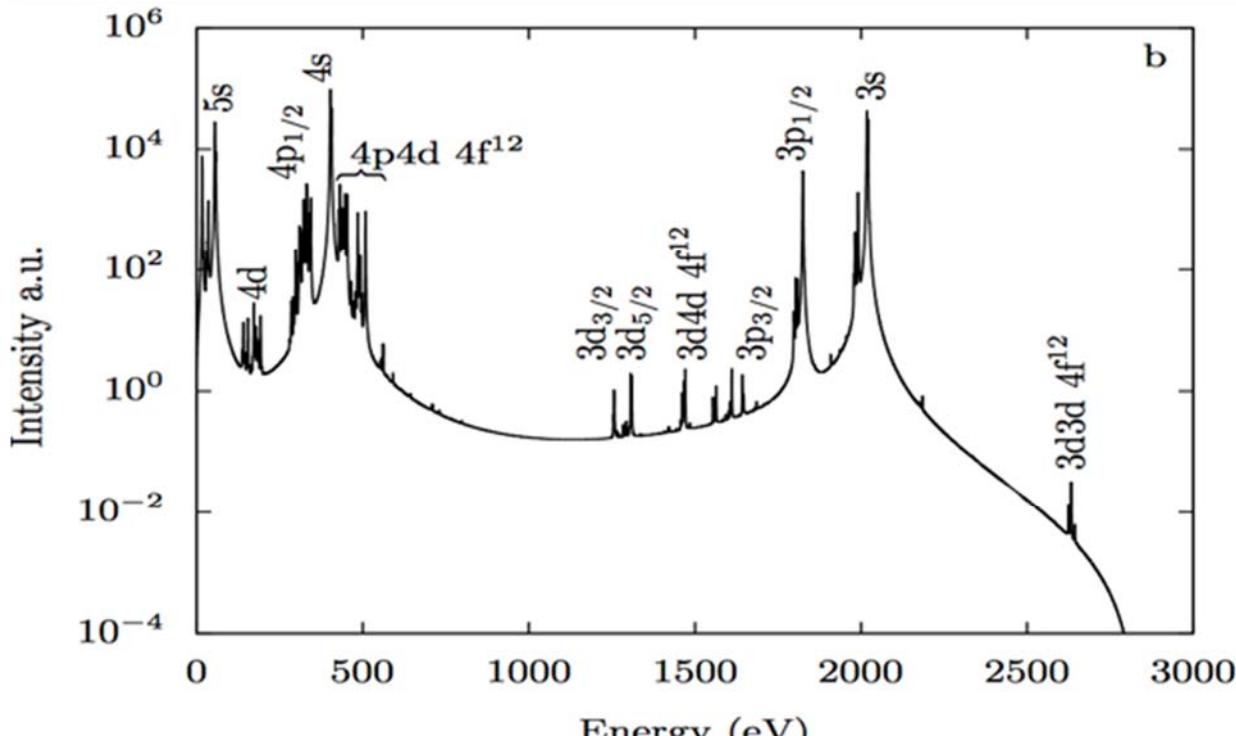


test of analysis routines:

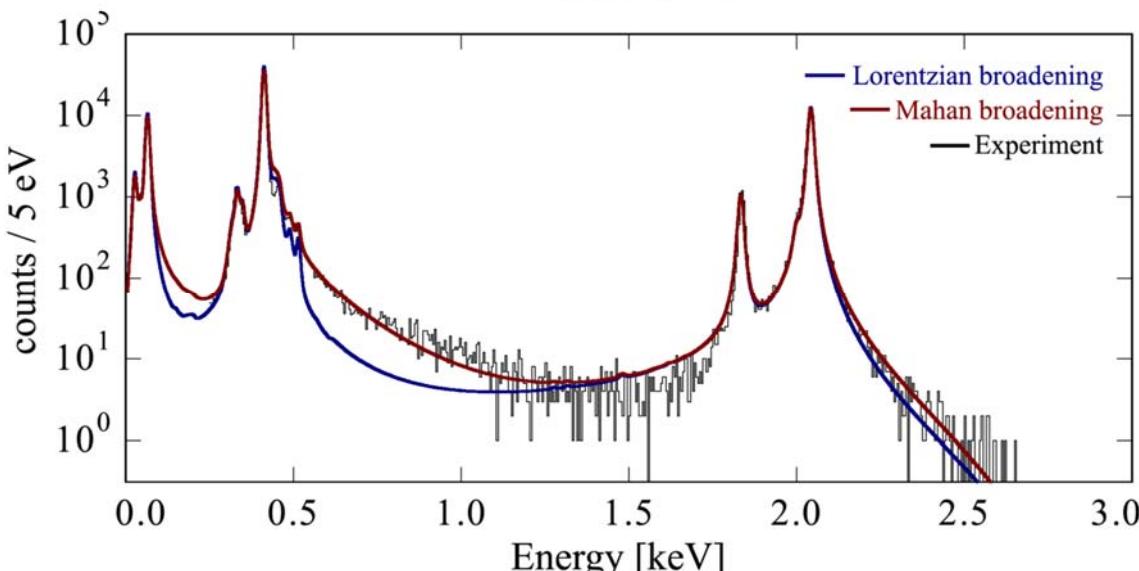
$$Q_{\text{EC}} = (2838 \pm 14) \text{ eV}$$

$m(v_e) < 150 \text{ eV}$ (95% C.L.) profile log-likelihood ratio hypothesis test

^{163}Ho spectral shape analysis

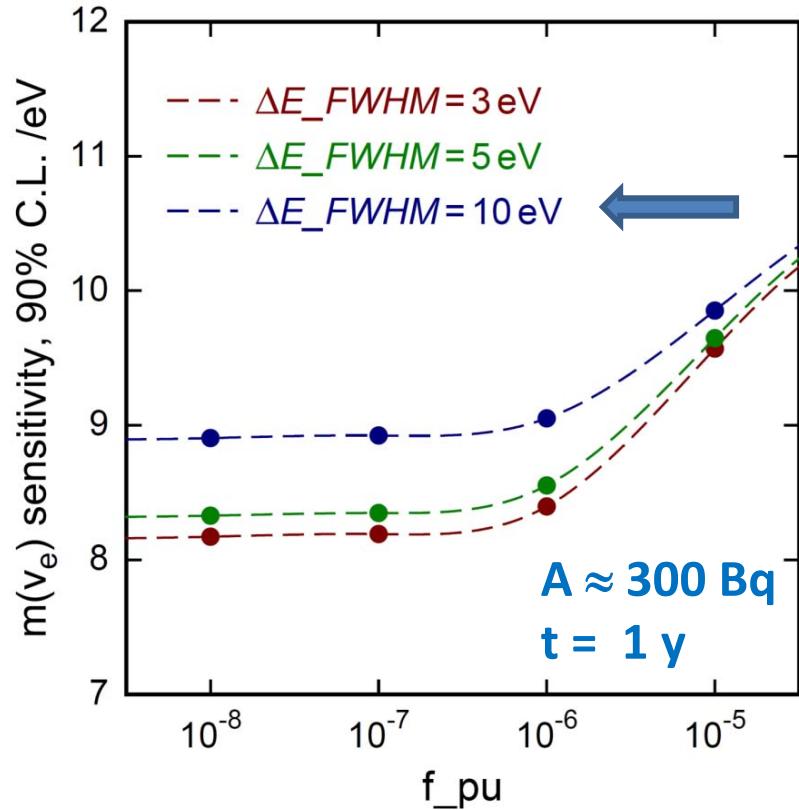


- A. Faessler et al.
J. Phys. G **42** (2015) 015108
 - R. G. H. Robertson
Phys. Rev. C **91**, 035504 (2015)
 - A. Faessler et al.
Phys. Rev. C **91**, 064302 (2015)
 - A. Faessler et al.
Phys. Rev. C **95**, (2017) 045502
 - A. Faessler and F. Simkovic
Phys. Rev. C **91**, 045505 (2015)
 - A. De Rujula and M. Lusignoli
JHEP 05 (2016) 015, arXiv:1601.04990v1
- ❖ Brass et al., *Phys. Rev. C* **97** (2018) 054620
Ab-initio calculation



ECHo timeplan

ECHo-1k – revised (2015 – 2018+)



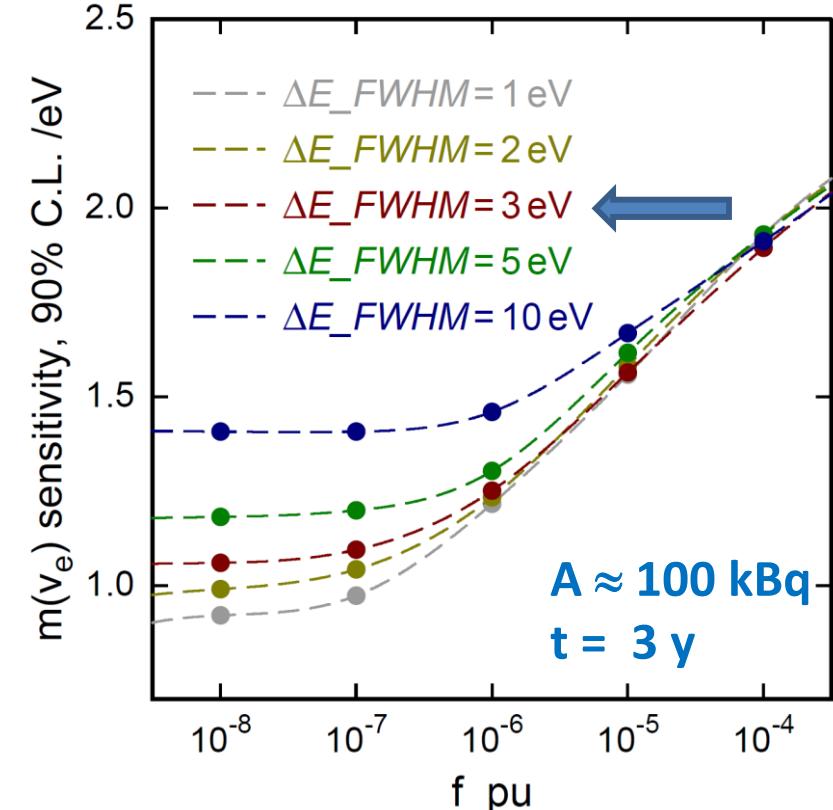
$m(\nu_e) < 10$ eV 90% C.L.

Activity per pixel: 1 - 5 Bq

Number of detectors: 60 - 100

Readout: parallel two stage SQUID

ECHo-100k (2018 – 2021+)



$m(\nu_e) < 1.5$ eV 90% C.L.

Activity per pixel: 10 Bq

Number of detectors: 12000

Readout: microwave SQUID multiplexing

Towards ECHo-100k

ECHo-100k chip in fabrication

- single pixel optimization: ^{163}Ho activity per pixel $a \approx 10 \text{ Bq}$
reduced absorber thickness → increase signal to noise ratio
 - suitable for parallel and multiplexed readout
- ^{163}Ho implantation on several chips foreseen before the end of the year

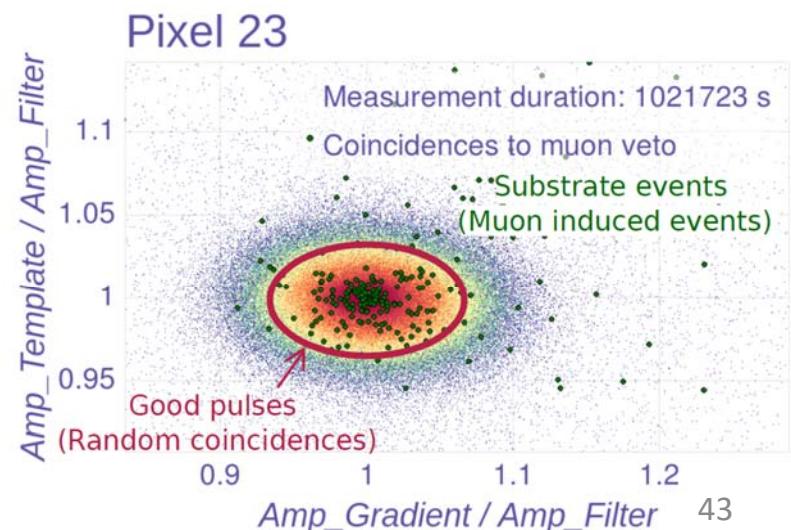
Microwave multiplexed readout of MMC demonstrated

- adapt room temperature electronics for larger number of detectors/channel (goal 400 det/ch)
 - install 13 new microwave channels in ECHo cryostat
- ^{163}Ho spectrum acquired at GHz frequency foreseen early 2020

Preparation of background model for ECHo

- Experiments with muon veto demonstrate that muon related events **discriminated via pulse shape**
- Effect of **low energy secondary** radiation is being investigated via Monte Carlo simulations

A. Göggelmann et al. Muon induced background in ECHo,
in preparation



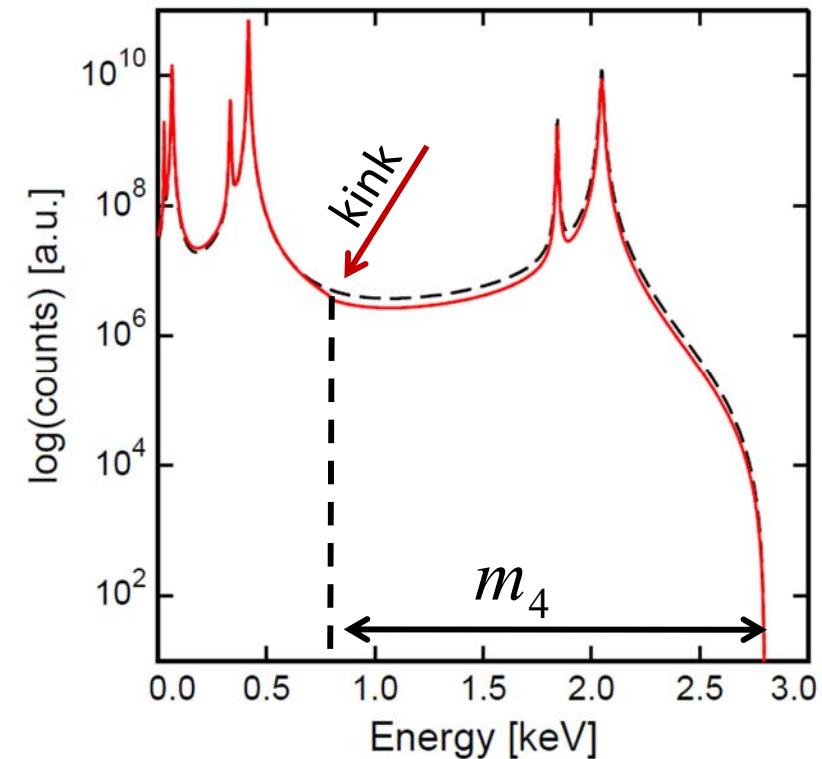
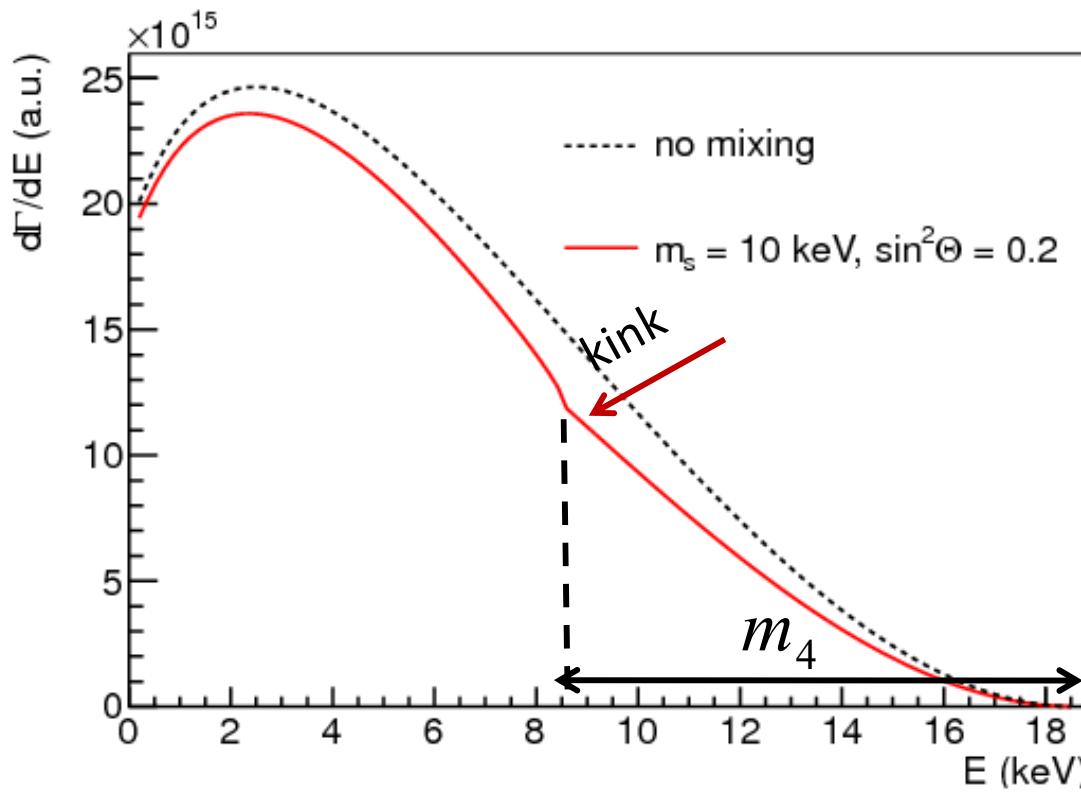
Beyond the standard 3-neutrino scenario



Sterile neutrinos in β - and EC spectra

The existence of a sterile neutrino implies the existence of a fourth mass eigenstate

$$\frac{dN}{dE} = \cos^2 \vartheta_s \frac{dN}{dE}(m_{\text{active}}) + \sin^2 \vartheta_s \frac{dN}{dE}(m_{\text{sterile}})$$

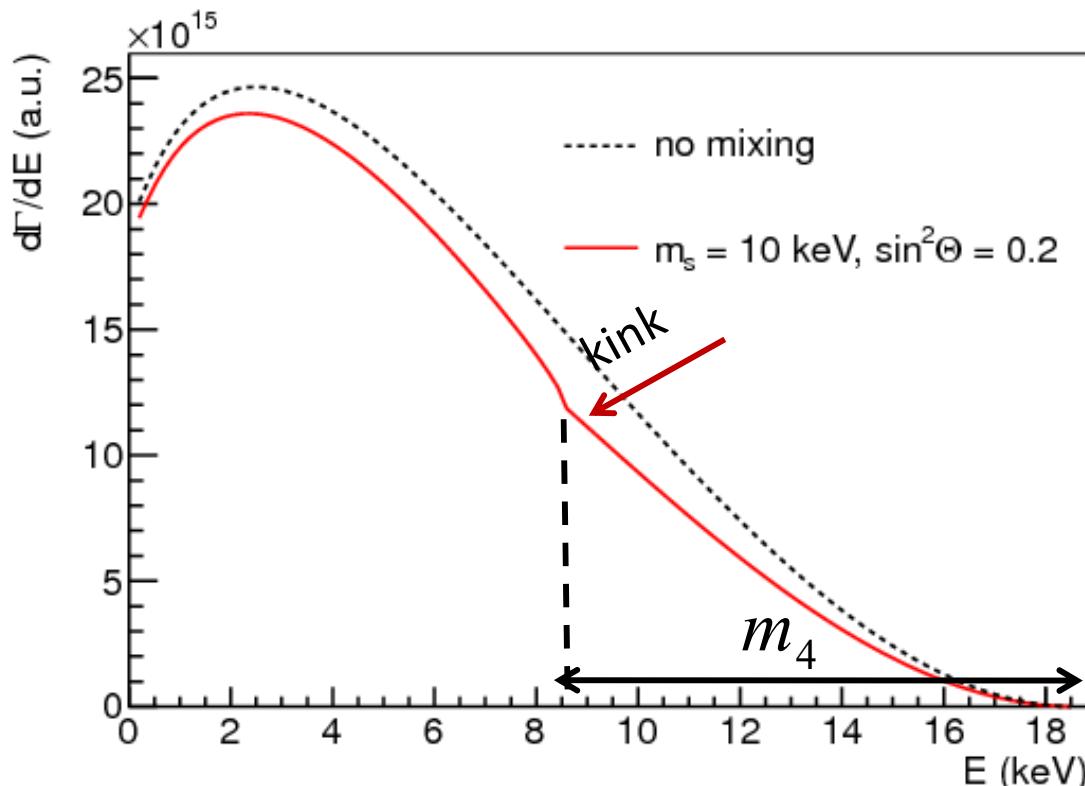


Search for sterile neutrino signature both at the eV scale as well as keV scale

Sterile neutrinos in β - and EC spectra

The existence of a sterile neutrino implies the existence of a fourth mass eigenstate

$$\frac{dN}{dE} = \cos^2 \vartheta_s \frac{dN}{dE}(m_{active}) + \sin^2 \vartheta_s \frac{dN}{dE}(m_{sterile})$$



eV-scale

Esmaili & Peres, *Phys. Rev. D* **85** (2012) 117301

Giunti, Li, Zhang, arXiv:1912.12956

keV-scale

TRISTAN experiment

Mertens et al., *J. Phys. G: Nucl. Part. Phys.* **46** (2019) 065203

Benso et al., *Phys. Rev. D* **100** (2019) 115035

Sterile neutrinos in β - and EC spectra

The existence of a sterile neutrino implies the existence of a fourth mass eigenstate

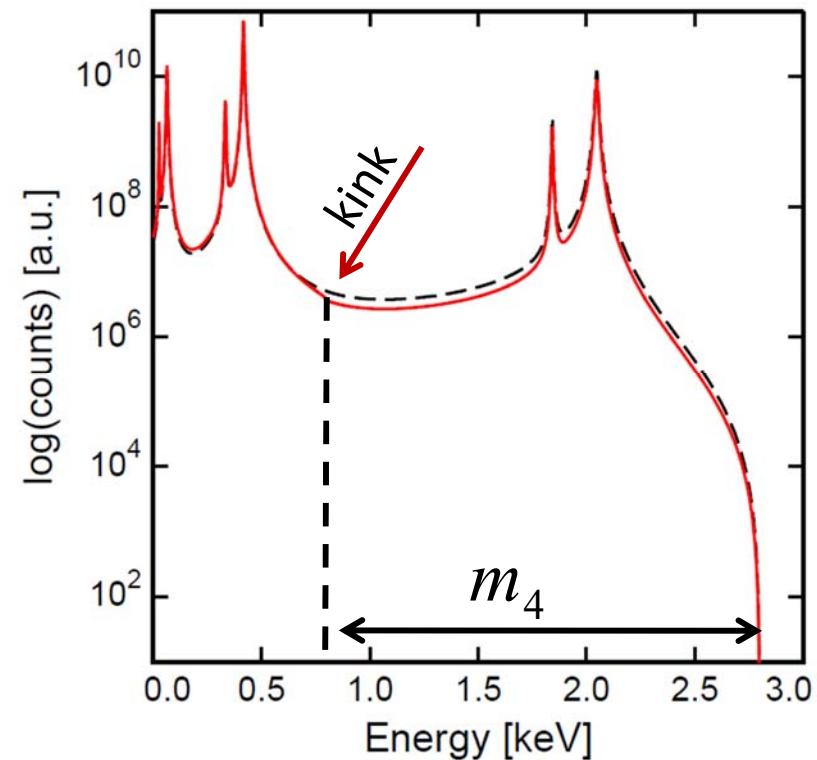
$$\frac{dN}{dE} = \cos^2 \vartheta_s \frac{dN}{dE}(m_{active}) + \sin^2 \vartheta_s \frac{dN}{dE}(m_{sterile})$$

eV-scale

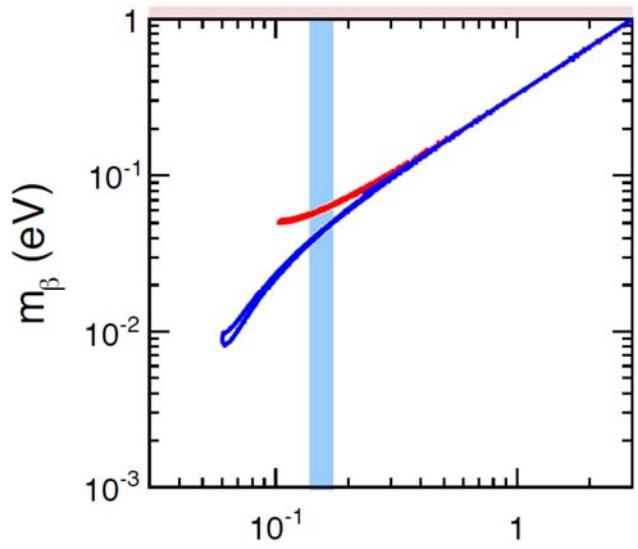
Gastaldo, Giunti, Zavanin, J. High Energy Physics
2016 (2016) 61

keV-scale

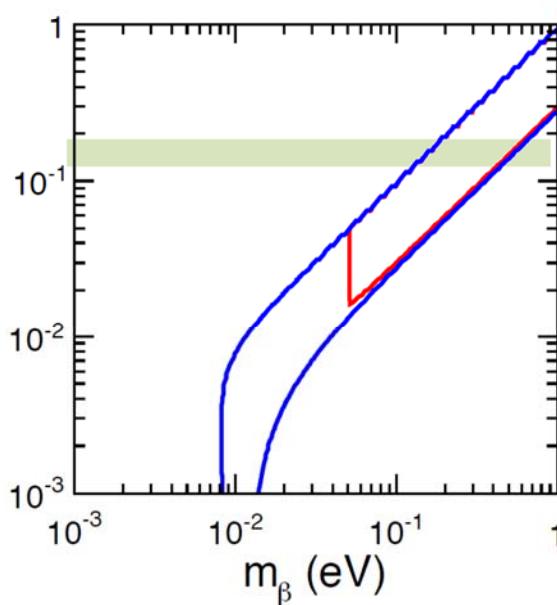
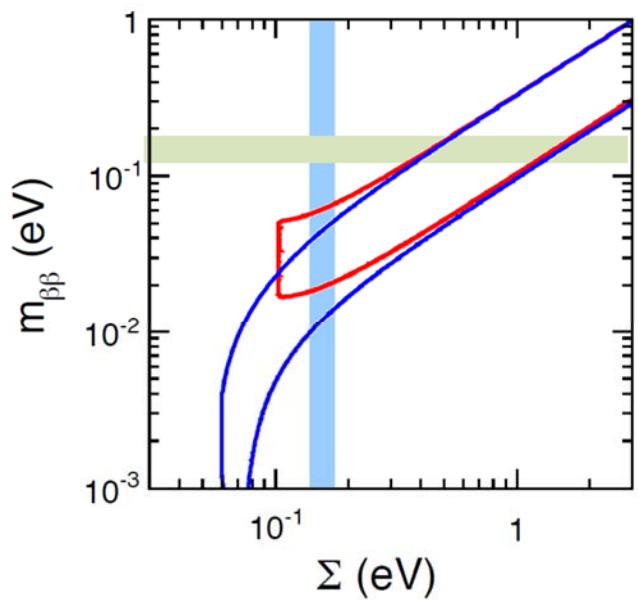
A White Paper on keV Sterile Neutrino Dark Matter
JCAP 01 (2017) 025



Conclusions



(m_β , $m_{\beta\beta}$, Σ)

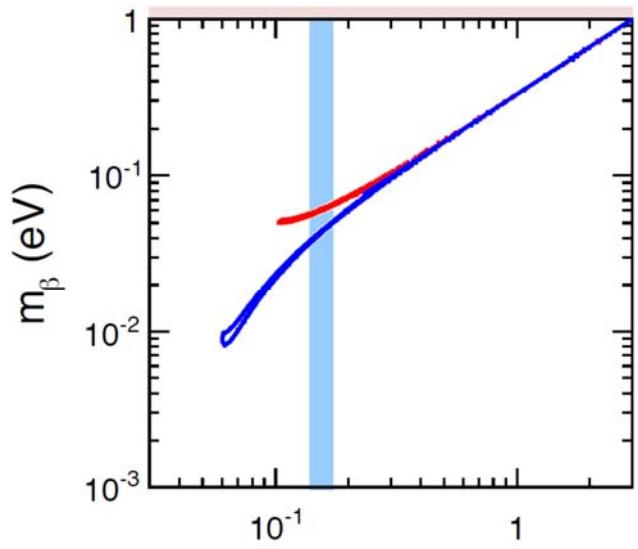


$$\Omega_\nu h^2 \simeq \sum \frac{m_\nu}{93 \text{ eV}}$$

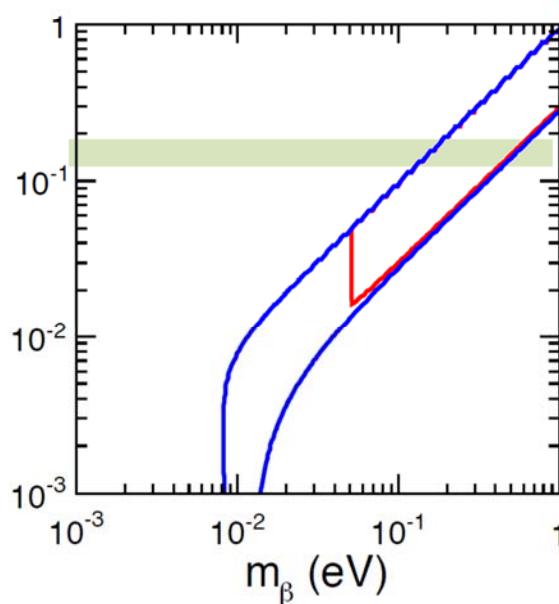
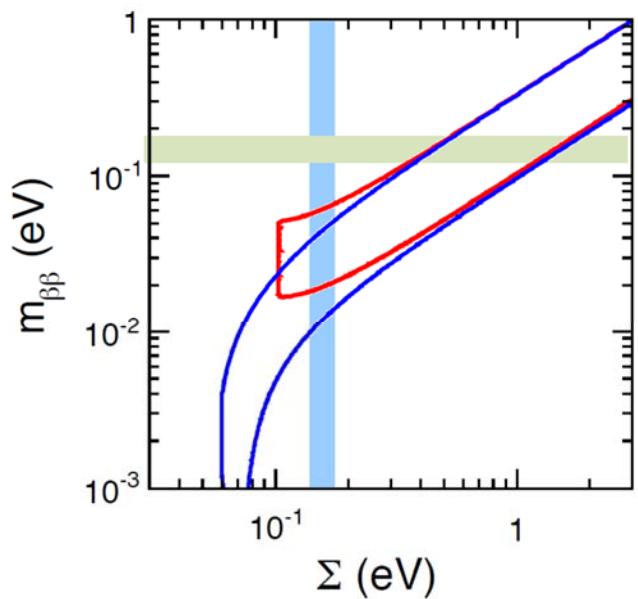
$$m_{\beta\beta} = \left| \sum_i U_{ei}^2 m_i \right|$$

$$m_\beta^2(\nu_e) = \sum_i |U_{ei}|^2 m_i^2$$

Conclusions



(m_β , $m_{\beta\beta}$, Σ)



$$\Omega_\nu h^2 \simeq \sum \frac{m_\nu}{93 \text{ eV}}$$

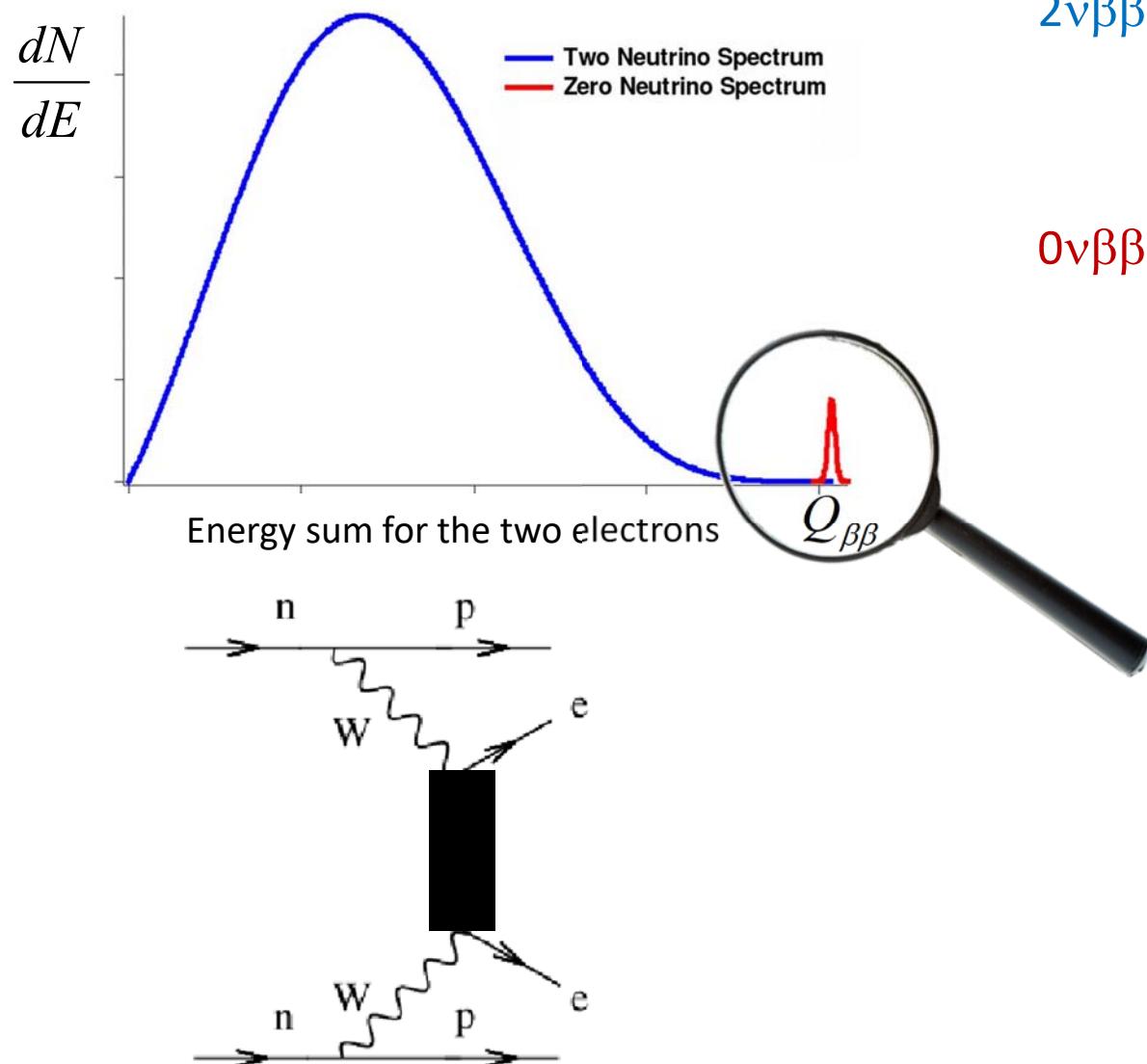
$$m_{\beta\beta} = \left| \sum_i U_{ei}^2 m_i \right|$$

$$m_\beta^2(\nu_e) = \sum_i |U_{ei}|^2 m_i^2$$

Thank you!

Neutrinoless double beta decay - ν mass

The halflife for $0\nu 2\beta$ decay depends on the neutrino mass



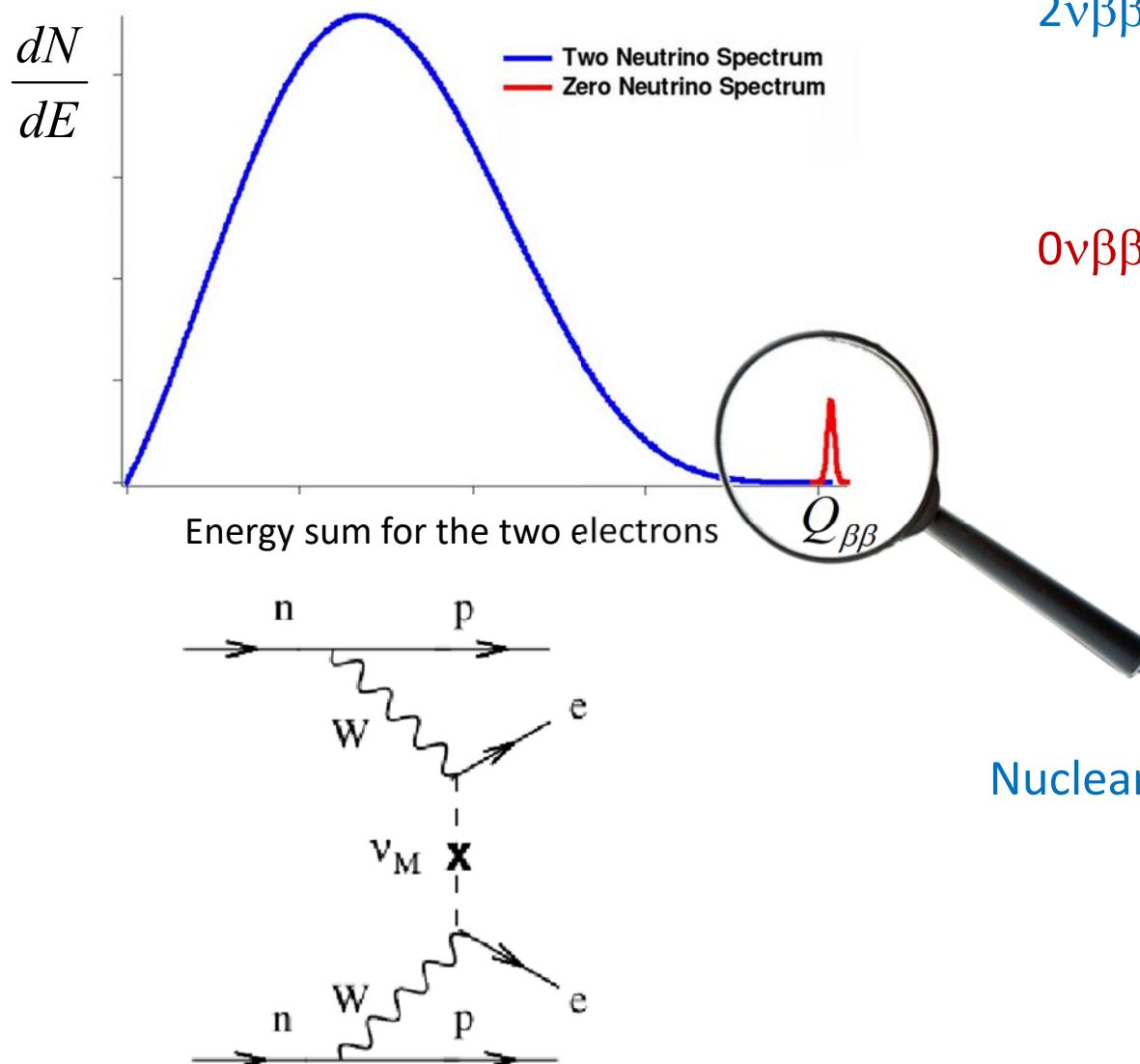
$$\tau_{1/2} \approx 10^{20} \text{ years}$$



$$\tau_{1/2} > 10^{25} \text{ years}$$

Neutrinoless double beta decay - ν mass

The halflife for $0\nu 2\beta$ decay depends on the neutrino mass



$$\tau_{1/2} \approx 10^{20} \text{ years}$$



$$\tau_{1/2} > 10^{25} \text{ years}$$

$$m_{\beta\beta}^2 = \left| \sum U_{ei}^2 m(\nu_i) \right|^2$$

Nuclear matrix element

Phase space term

$$(\tau_{1/2}^{0\nu})^{-1} = \left| \frac{m_{\beta\beta}}{m_e} \right|^2 |M_{\nu}^{0\nu}|^2 G^{0\nu}$$

Neutrinoless double beta decay - sensitivity

Typically an excess of events is not found...

A limit on the halflife for 0ν2e decay can be defined as function of:

Mass of the isotope	M	[kg]	}	Exposure $M \times T$	[kg × year]
Measuring time	T	[year]			
Energy resolution	ΔE	[keV]			
Background index	b	[keV ⁻¹ ton ⁻¹ year ⁻¹]			

Two limits defined by the background index

> 1 background events in ROI

$$\left(\tau_{1/2}^{\text{exp}}\right)^{-1} = (\ln 2) N_a \frac{a}{A} \epsilon \sqrt{\frac{MT}{b\Delta E}}$$

< 1 background events in ROI

$$\left(\tau_{1/2}^{\text{exp}}\right)^{-1} = (\ln 2) N_a \frac{a}{A} \epsilon \frac{MT}{n_{CL}}$$

Fight against background

Direct reduction of background activity

- Select and use ultra-pure materials
- Minimize all passive (non “source”) materials
- Avoid material re-contamination (machining, manipulation, storage)
- Fabricate ultra-clean materials (underground fab if needed)
- underground labs — reduced muon flux & related induced activations

Discrimination techniques

- Energy resolution
- Active veto detector
- Tracking (topology)
- Particle ID, angular, spatial, time correlations
- Fiducial Fits
- Granularity (arrays)
- Pulse shape discrimination (PSD)
- Ion Identification

Methods	
TPCs (liquid, gas)	^{136}Xe
Doped Liquid Scintillators	$^{136}\text{Xe}, ^{130}\text{Te}$
Solid state detectors	$^{76}\text{Ge}, ^{116}\text{Cd}$
Bolometers (+ enhancements)	$^{130}\text{Te}, ^{82}\text{Se}, ^{100}\text{Mo}, ^{116}\text{Cd}$
Foils with tracking chambers	$^{82}\text{Se}, ^{150}\text{Nd}, ^{100}\text{Mo}$

Both approaches are needed

Many different experiments

Experiment	Isotope	Technique	Mass $\beta\beta(0\nu)$ isotope
CUORICINO	130Te	TeO ₂ Bolometer	10 kg
NEMO3	100Mo/82Se	Foils with tracking	6.9/0.9 kg
GERDA I	76Ge	Ge diodes in LAr	15 kg
EXO200	136Xe	Xe liquid TPC	160 kg
KamLAND-ZEN	136Xe	2.7% in liquid scint.	380 kg
CUORE-0	130Te	TeO ₂ Bolometer	11 kg
GERDA II	76Ge	Point contact Ge in LAr	30+35 kg
Majorana D	76Ge	Point contact Ge	30 kg
CUORE	130Te	TeO ₂ Bolometer	206 kg
SNO+	130Te	0.3% natTe suspended in Scint	55 kg
NEXT-100	136Xe	High pressure Xe TPC	80 kg
SuperNEMO D	82Se	Foils with tracking	7 kg
CANDLES	48Ca	305 kg of CaF ₂ crystals - liq. scint	0.3 kg
LUCIFER	82Se	ZnSe scint. bolometer	18 kg
1TGe (GERDA+MJ)	76Ge	Best technology from GERDA and MAJORANA	~ tonne
CUPID	-	Hybrid Bolometers	~ tonne
nEXO	136Xe	Xe liquid TPC	~ tonne
SuperNEMO	82Se	Foils with tracking	100 kg
AMoRE	100Mo	CaMoO ₄ scint. bolometer	50 kg
MOON	100Mo	Mo sheets	200 kg
COBRA	116Cd	CdZnTe detectors	10 kg/183 kg
CARVEL	48Ca	48CaWO ₄ crystal scint.	~ tonne
DCBA	150Nd	Nd foils & tracking chambers	20 kg



Best results