

# Measurement of R ratios at BESIII

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# Outline

- **Introduction**
- **BESIII data sets**
- **Progress in  $R$  ratio measurement**
  - ▶ Signal and backgrounds
  - ▶ Simulation of inclusive hadronic events
  - ▶ Systematics due to simulation model
- **Exclusive measurements**
  - ▶  $e^+e^- \rightarrow \phi\eta/\eta'$
  - ▶  $e^+e^- \rightarrow \omega\eta/\pi^0$
  - ▶  $e^+e^- \rightarrow K^+K^-\pi^0\pi^0$
- **Summary**

# The definition of $R$ ratio

$R$  is defined as the ratio of the production rate of hadron and muon pairs:

$$R \equiv \frac{\sigma^0(e^+e^- \rightarrow \text{hadrons})}{\sigma^0(e^+e^- \rightarrow \mu^+\mu^-)} \equiv \frac{\sigma_{\text{had}}^0}{\sigma_{\mu\mu}^0}$$

That is:

$$R \equiv \frac{\sigma(e^- e^+ \rightarrow \text{hadrons})}{\sigma(e^- e^+ \rightarrow \mu^+ \mu^-)}$$

According to the QED theory:

$$\sigma_{\mu\mu}^0(s) = \frac{4\pi\alpha^2}{3s} \frac{\beta_\mu(3 - \beta_\mu^2)}{2}$$

# Why $R$ ratios?

R ratios can contribute to:

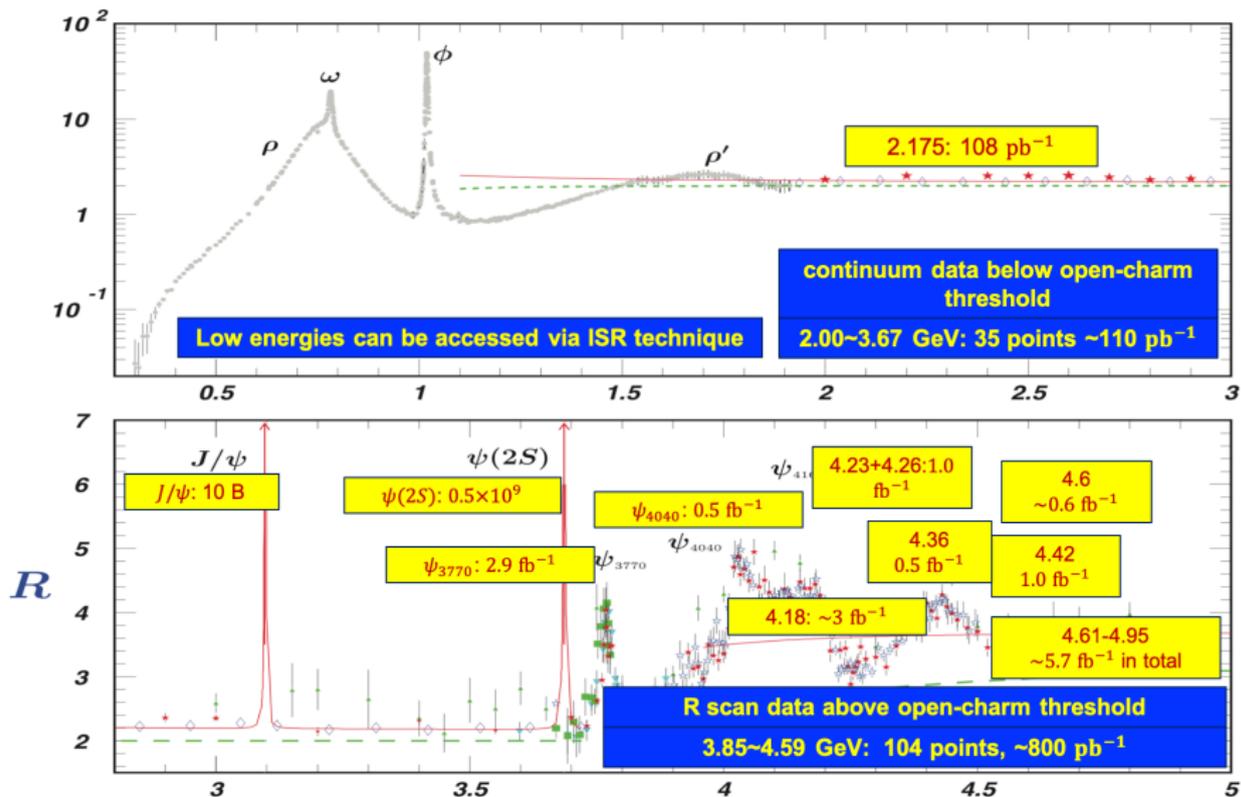
- Calculation of anomalous magnetic moment of Muon:  $a_\mu = (g_\mu - 2)/2$ .

Source	Contribution( $\times 10^{11}$ )
$a_\mu^{\text{QED}}$	$116584718.92 \pm 0.03$
$a_\mu^{\text{Weak}}$	$153.6 \pm 1.0$
$a_\mu^{\text{had}} [\text{LO}]$	$6939 \pm 40$
$a_\mu^{\text{had}} [\text{NLO}]$	$-98.7 \pm 0.9$
$a_\mu^{\text{had}} [\text{NNLO}]$	$12.4 \pm 0.1$
$a_\mu^{\text{had, l-l}}$	$105 \pm 26$
$a_\mu^{\text{SM}}$	$116591830 \pm 48$
$a_\mu^{\text{exp}}$	$116592091 \pm 63$
$\Delta a_\mu$	$261 \pm 79$

- Determination of running coupling constant of QED theory at  $M_Z$ :  $\Delta\alpha(M_Z^2)$ .

Source	Contribution( $\times 10^4$ )
$\Delta\alpha_{\text{lepton}}(M_Z^2)$	$314.979 \pm 0.002$
$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$	$276.0 \pm 1.0$
$\Delta\alpha_{\text{top}}(M_Z^2)$	$-0.7180 \pm 0.0054$

Prog. Theor. Exp. Phys. 2020, 083C01 (2020); PRD 98, 030001 (2018); PRD 97, 114025 (2018)



# Determination of $R$ ratio in experiment

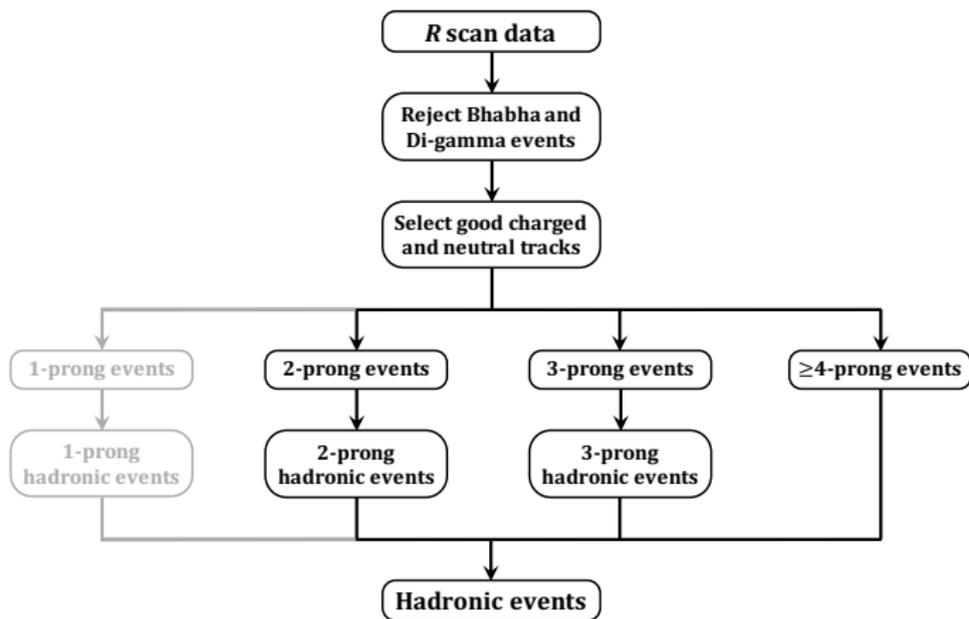
Experimentally,  $R$  ratio is determined by

$$R = \frac{N_{\text{had}}^{\text{obs}} - N_{\text{bkg}}}{\mathcal{L}_{\text{int}} \epsilon_{\text{had}} \epsilon_{\text{trig}} (1 + \delta) \sigma_{\mu\mu}^0}$$

- $N_{\text{had}}^{\text{obs}}$ : Numbers of observed hadronic events.
- $N_{\text{bkg}}$ : Number of the residual background events.
- $\mathcal{L}_{\text{int}}$ : Integrated luminosity.
- $\epsilon_{\text{trig}}$ : Trigger efficiency.
- $\epsilon_{\text{had}}$ : Detection efficiency of the hadronic events.
- $(1 + \delta)$ : ISR correction factor.
- $\sigma_{\mu\mu}^0(s) = 86.85 \text{ nb/s}$ : Leading order QED cross section for  $e^+e^- \rightarrow \mu^+\mu^-$ .

# Signal and background estimation

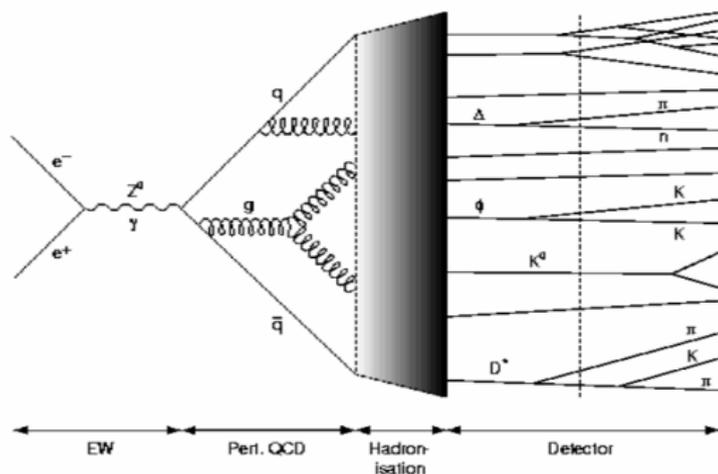
Selection of inclusive hadronic events from data:



- ▶ Residual QED-related backgrounds are estimated by using dedicated MC generators.
- ▶ Beam-associated backgrounds are studied by separate-beam data.

# Signal simulation: LUARLW model

Hadronization procedure in LUARLW model:

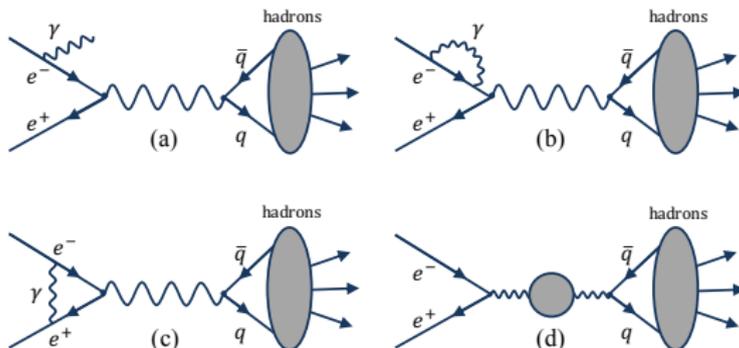


Features of LUARLW model:

- ▶ Development of **JETSET** for low energy experiments.
- ▶ Simulate production and decay of both **continuum** and **resonance** states.
- ▶ Kinematics of initial hadrons are determined by **Lund Area Law** [[arXiv:hep-ph/9910285](https://arxiv.org/abs/hep-ph/9910285)].
- ▶ **Phenomenological parameters** should be tuned according to experiment data.
- ▶ Integrated the **Initial-state radiation (ISR)** and **Vacuum Polarization (VP)** corrections.

# ISR in LUARLW: Feynman Diagram

In LUARLW, the **Feynman Diagram (FD)** scheme is used to simulate ISR correction and calculate  $(1 + \delta)$ . Feynman diagrams related to ISR procedure are:



The total hadronic cross section measured by experiment is the total effect of all these diagrams:

$$\sigma_{\text{had}}^{\text{tot}}(s) = \beta \int_0^{x_m} dx \frac{x^\beta}{x} \left(1 - x + \frac{x^2}{2}\right) \frac{\sigma_{\text{had}}^0(s')}{|1 - \Pi(s')|^2} + \delta_{\text{vert}} \frac{\sigma_{\text{had}}^0(s)}{|1 - \Pi(s)|^2},$$

and

$$1 + \delta \equiv \sigma_{\text{had}}^{\text{tot}} / \sigma_{\text{had}}^0$$

According to definition of  $R$ , its uncertainty is expressed as

$$\left(\frac{\Delta R}{R}\right)_{\text{sys}}^2 = \left(\frac{\Delta \tilde{N}}{\tilde{N}}\right)^2 + \left(\frac{\Delta \mathcal{L}_{\text{int.}}}{\mathcal{L}_{\text{int.}}}\right)^2 + \left(\frac{\Delta \varepsilon_{\text{had}}}{\varepsilon_{\text{had}}}\right)^2 + \left[\frac{\Delta(1+\delta)}{(1+\delta)}\right]^2 + \left(\frac{\Delta \varepsilon_{\text{trig}}}{\varepsilon_{\text{trig}}}\right)^2,$$

where

$$\tilde{N} = N_{\text{had}}^{\text{obs}} - N_{\text{bkg}}$$

In practice, its uncertainties are addressed in different aspects:

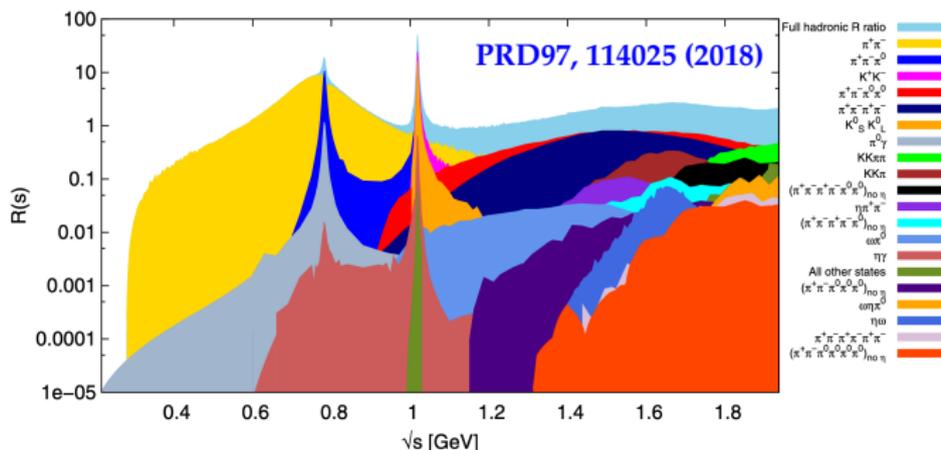
- **Event selection:** all implemented selection criteria are slightly changed.
- **Background estimation:** different methods and background simulation models are used.
- **Integrated luminosity:** uncertainty is evaluated in [CPC 41, 063001 \(2017\)](#).
- **Signal simulation:** Hybrid model is development and used as a cross check.
- **ISR correction factor:** considered in calculation precision, different  $\sigma_{\text{had}}^0$ .
- **Trigger efficiency:**  $\varepsilon_{\text{trig}}$  approaches 100% with an uncertainty less than 0.1%.

# Hybrid model

The Hybrid model is developed as an alternative simulation model of hadronic events.

Hybrid model is consisted of **THREE** components:

- **ConExc**: simulates 51 exclusive processes with known cross section line-shapes.
- **Phokhara**: describes 10 well known processes such as  $e^+e^- \rightarrow 2\pi, 3\pi, 4\pi$  etc..
- **LUARLW**: responsible for remain unknown processes with tuned parameters.

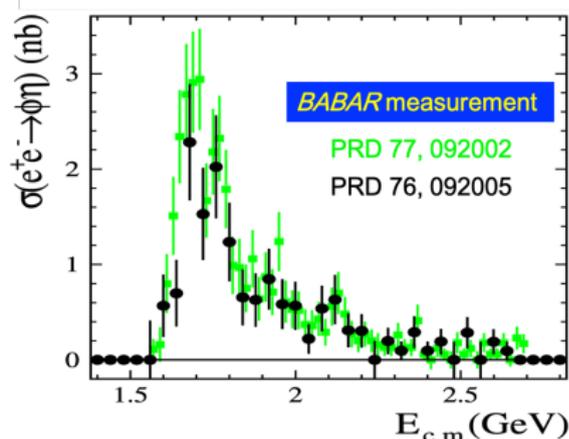
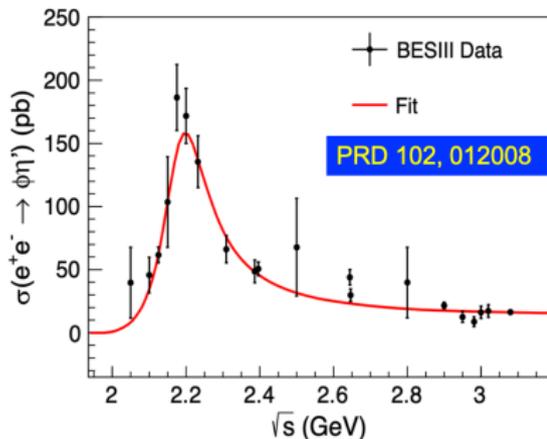


- ▶ As much as known experimental knowledge is implemented in ConExc.
- ▶ Residual double-countings among the three components are negligible.

# Latest exclusive measurements at BESIII

# Exclusive measurement: $e^+e^- \rightarrow \phi\eta/\eta'$

BESIII measured the  $e^+e^- \rightarrow \phi\eta/\eta'$  processes at  $\sqrt{s} = 2.00 \sim 3.08$  GeV:



If the resonance structure observed in these two processes are  $\phi(2170)$ :

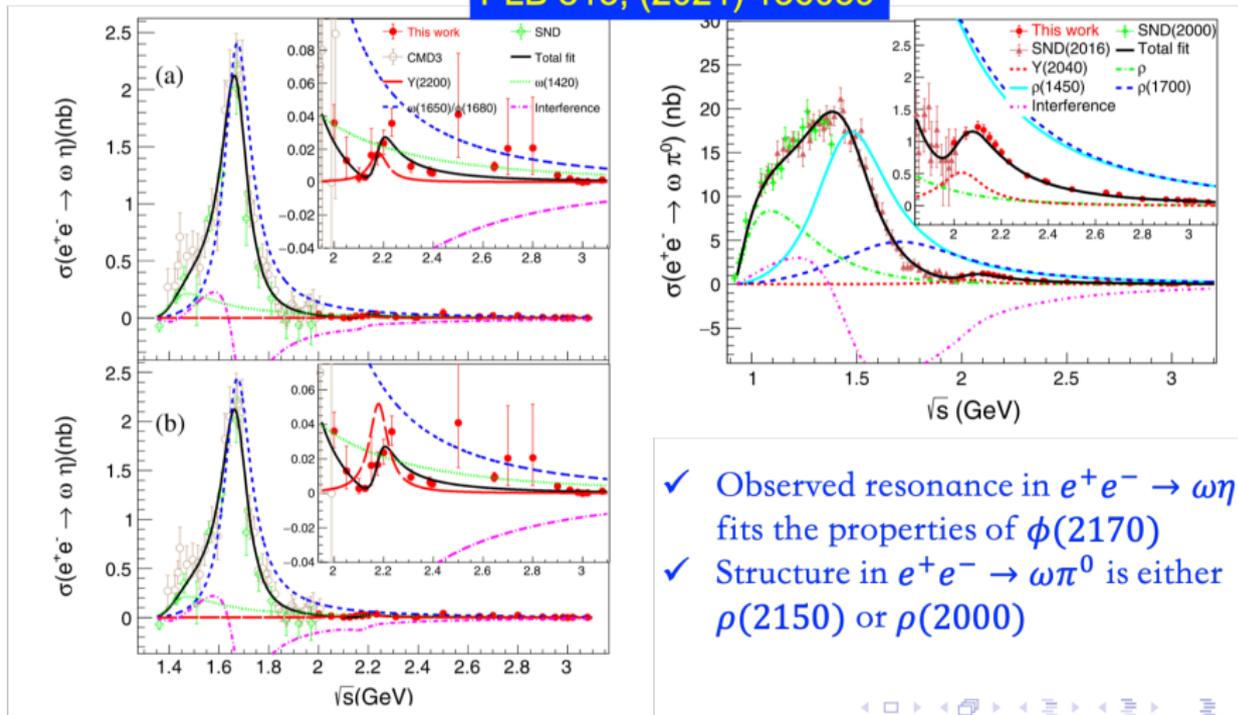
$$\frac{Br[\phi(2170) \rightarrow \phi\eta] \cdot \Gamma_{ee}^{\phi(2170)}}{Br[\phi(2170) \rightarrow \phi\eta'] \cdot \Gamma_{ee}^{\phi(2170)}} = 0.23 \pm 0.10 \pm 0.18$$

which disagrees the predictions of several models.

# Exclusive measurement: $e^+e^- \rightarrow \omega\eta/\pi^0$

BESIII measured the  $e^+e^- \rightarrow \omega\eta/\pi^0$  processes at  $\sqrt{s} = 2.00 \sim 3.08$  GeV:

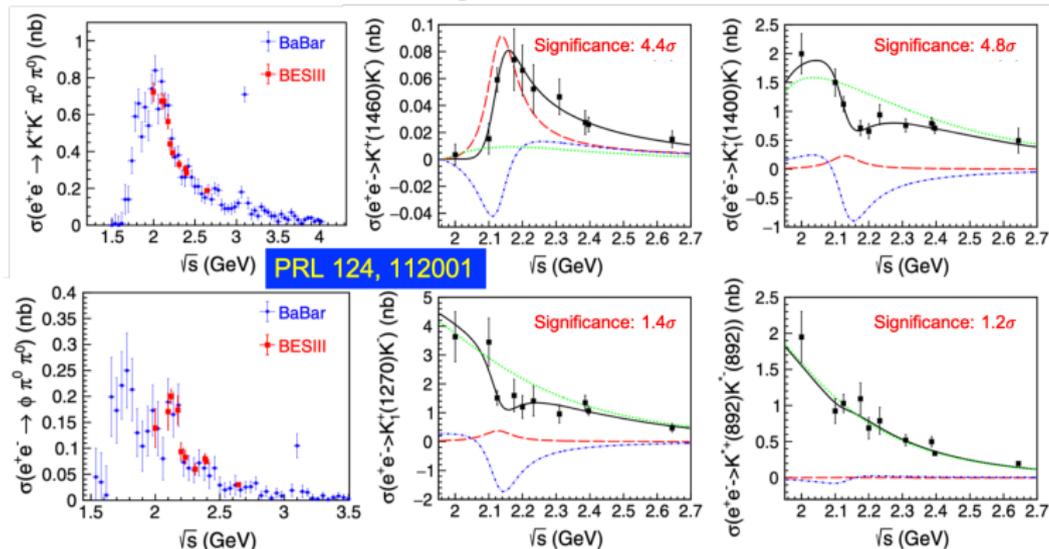
PLB 813, (2021) 136059



- ✓ Observed resonance in  $e^+e^- \rightarrow \omega\eta$  fits the properties of  $\phi(2170)$
- ✓ Structure in  $e^+e^- \rightarrow \omega\pi^0$  is either  $\rho(2150)$  or  $\rho(2000)$

# Exclusive measurement: $e^+e^- \rightarrow K^+K^-\pi^0\pi^0$

BESIII measured the  $e^+e^- \rightarrow K^+K^-\pi^0\pi^0$  processes at  $\sqrt{s} = 2.000 \sim 2.644$  GeV:



- ✓ The PWA study for the  $K^+K^-\pi^0\pi^0$  final states, with the cross sections of subprocesses  $e^+e^- \rightarrow K^+(1460)K^-, K_1^+(1400)K^-, K_1^+(1270)K^-$  and  $K^*(892)K^*(892)$  extracted.
- ✓ Assuming the structure observed in these subprocesses is  $\phi(2170)$ , above estimated significances in these processes can help to distinguish corresponding theoretic models.

# Summary

- After numerous tuning and checking on signal simulation models, significant progress of  $R$  measurement is made at BESIII.
- The current nominal and alternative signal simulation models can both describe data well and result in consistent  $R$  ratios.
- Latest exclusive measurements provide valuable input to the vector meson spectroscopy.

**Thanks for your attention!**

# Backups

$E_{\text{beam}}: 1.000\sim 2.475 \text{ GeV}$   
 $\sigma_E: 5.16 \times 10^{-4}$   
 $L: 1.0 \times 10^{33} \text{ cm}^{-2}\text{s}^{-1} @3770$

**Linac**

**BES**

**Storage ring**

<b>Leptons</b>	$u$ <small>up</small>	$c$ <small>charm</small>	$t$ <small>top</small>
	$d$ <small>down</small>	$s$ <small>strange</small>	$b$ <small>bottom</small>
	$\nu_e$ <small>e- neutrino</small>	$\nu_\mu$ <small><math>\mu</math>- neutrino</small>	$\nu_\tau$ <small><math>\tau</math>- neutrino</small>
	$e$ <small>electron</small>	$\mu$ <small>muon</small>	$\tau$ <small>tau</small>
	Three Generations of Matter		

**Main Drift Chamber**

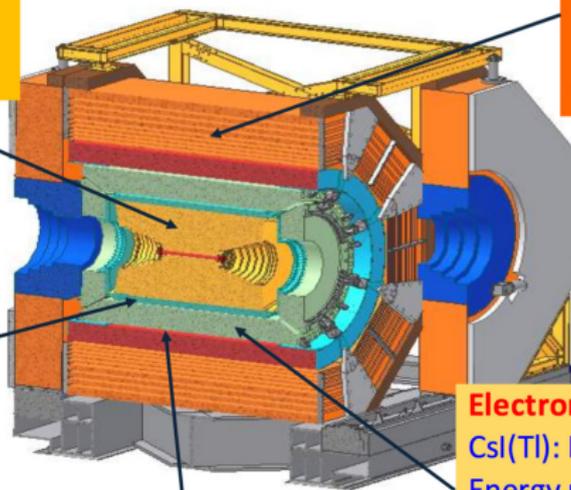
Small cell, 43 layer

 $\sigma_{xy} = 130 \mu\text{m}$ ,  $dE/dx \sim 6\%$  $\sigma_p/p = 0.5\%$  at 1 GeV**Muon Counter**

Resistive plate chamber

Barrel: 9 layers

Endcaps: 8 layers

 $\sigma_{\text{spatial}} = 1.48 \text{ cm}$ **Time Of Flight**

Plastic scintillator

 $\sigma_T(\text{barrel}) = 80 \text{ ps}$  $\sigma_T(\text{endcap}) = 110 \text{ ps}$ 

65 ps with MRPC ETOF

**Electromagnetic Calorimeter**CsI(Tl):  $L = 28 \text{ cm}$  ( $15X_0$ )

Energy range: 0.02-2 GeV

Barrel  $\sigma_E = 2.5\%$ ,  $\sigma_I = 6 \text{ mm}$ Endcap  $\sigma_E = 5.0\%$ ,  $\sigma_I = 9 \text{ mm}$ **SC Magnet 1.0T**

# ISR in Hybrid: Structure Function

In Hybrid generator, the **Structure Function (SF)** scheme is implemented:

$$\sigma_{\text{had}}^{\text{tot}}(s) = \int_0^{x_m} dx F_{\text{SF}}(x, s) \frac{\sigma_{\text{had}}^0(s')}{|1 - \Pi(s')|^2},$$

where

$$F_{\text{SF}}(x, s) = \beta x^{\beta-1} \Delta - \beta \left(1 - \frac{1}{2}x\right) - \frac{1}{8} \beta^2 \left[ 4(2-x) \ln x + \frac{1+3(1-x)^2}{x} \ln(1-x) + 6-x \right],$$

$$\Delta = 1 + \frac{\alpha}{\pi} \left( \frac{3}{2}L + \frac{\pi^2}{3} - 2 \right) + \left( \frac{\alpha}{\pi} \right)^2 \left\{ \left[ \frac{9}{8} - 2\zeta(2) \right] L^2 + \left[ -\frac{45}{16} + \frac{11}{2} \zeta(2) + 3\zeta(3) \right] L - \frac{6}{5} [\zeta(2)]^2 - \frac{9}{2} \zeta(3) - 6\zeta(2) \ln 2 + \frac{3}{8} \zeta(2) + \frac{57}{12} \right\}.$$

- ▶ Exactly same input hadronic cross section  $\sigma_{\text{had}}^0$  as that of LUARLW generator.
- ▶ Fred Jegerlehner group calculated  $\Pi(s)$  is used, which is different from LUARLW.
- ▶ This ISR correction scheme is integrated in Hybrid model.