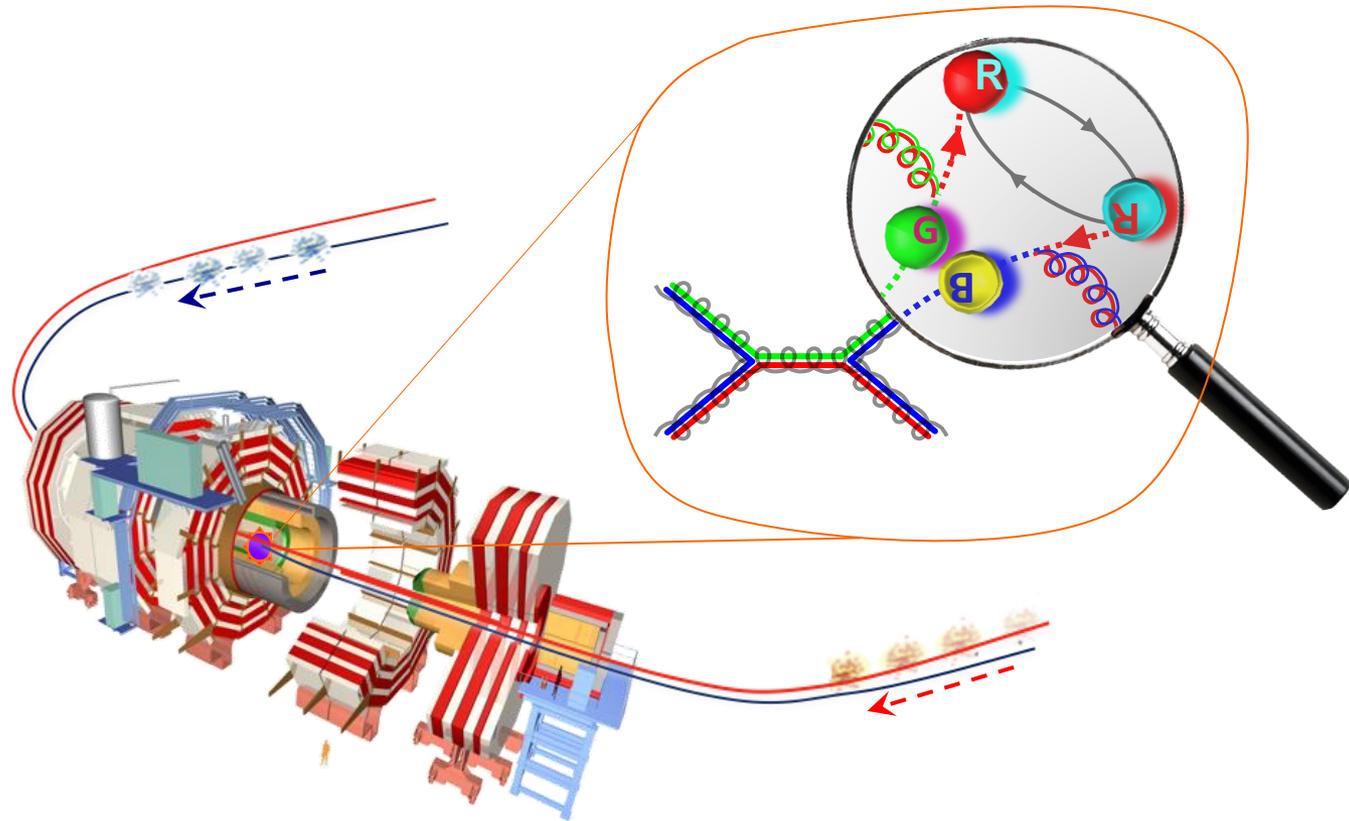


# Advancing the understanding of quarkonium production through global-fit analyses of LHC data



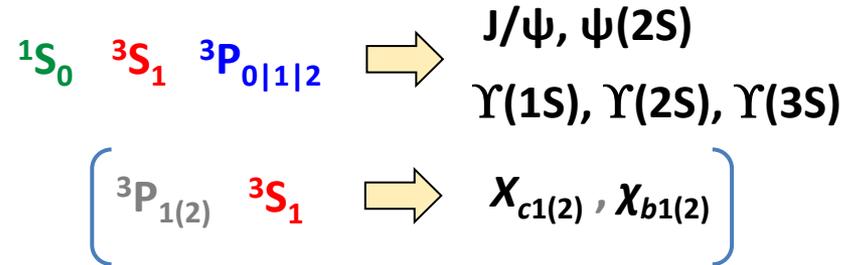
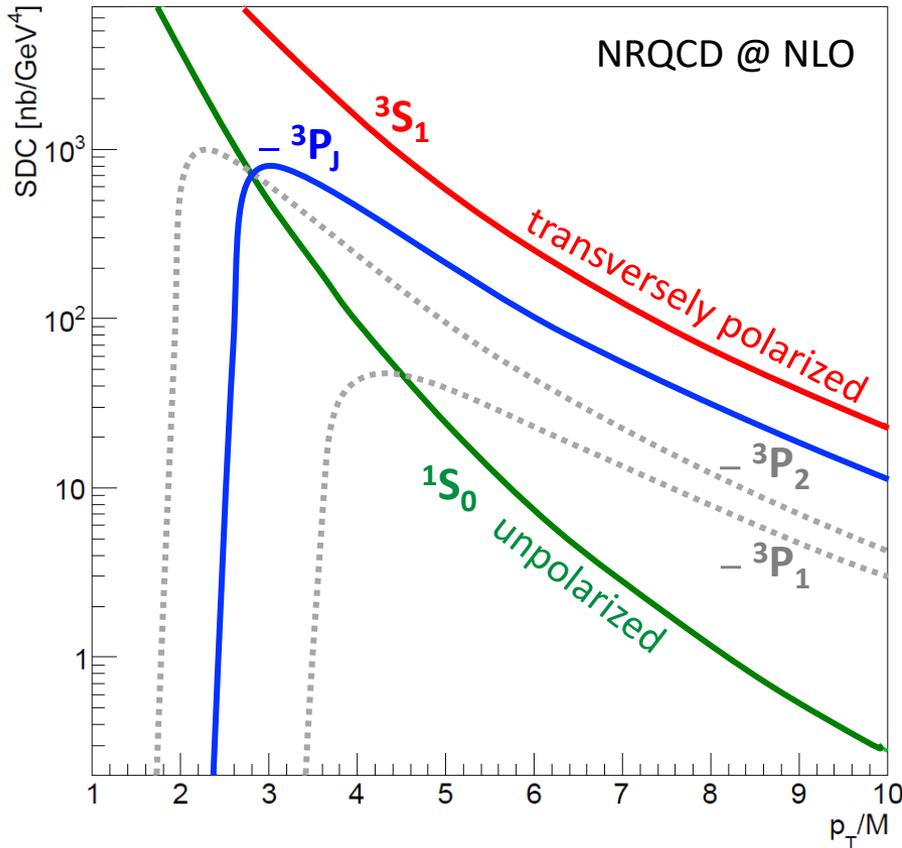
**FCT** Fundação  
para a Ciência  
e a Tecnologia

14<sup>th</sup> International Workshop on Heavy Quarkonium  
March 2021

**Mariana Araújo,**  
Pietro Faccioli, Carlos Lourenço  
and Thomas Madlener

# NRQCD: octet terms dominate quarkonium production

In NRQCD, S-wave quarkonia are mainly produced through a superposition of colour octet terms, with computed SDCs (functions of  $p_T$ ) and unknown LDMEs (normalizations)

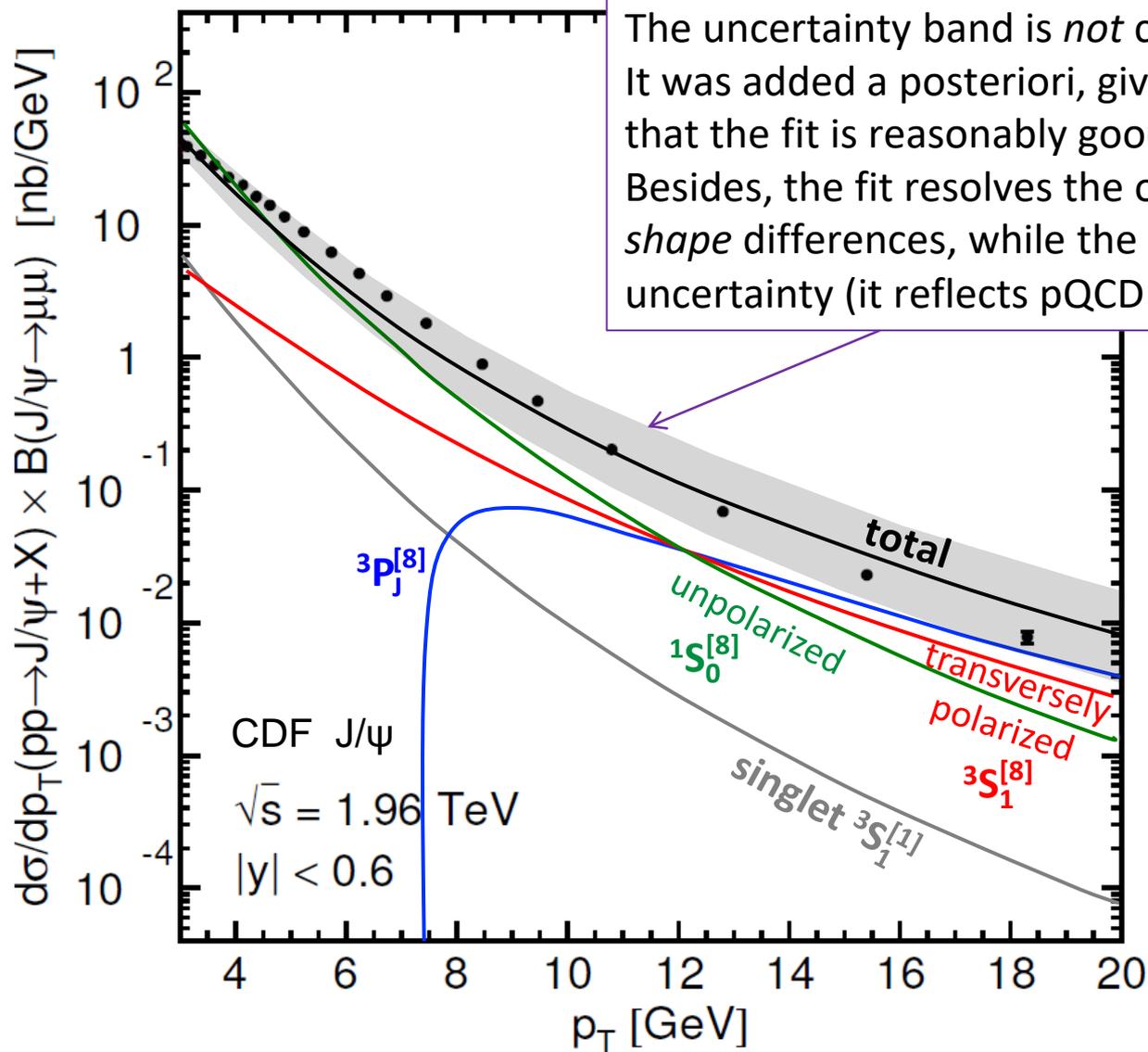


Fits to measured  $p_T$  distributions provide the LDMEs, which are then used to predict the polarizations

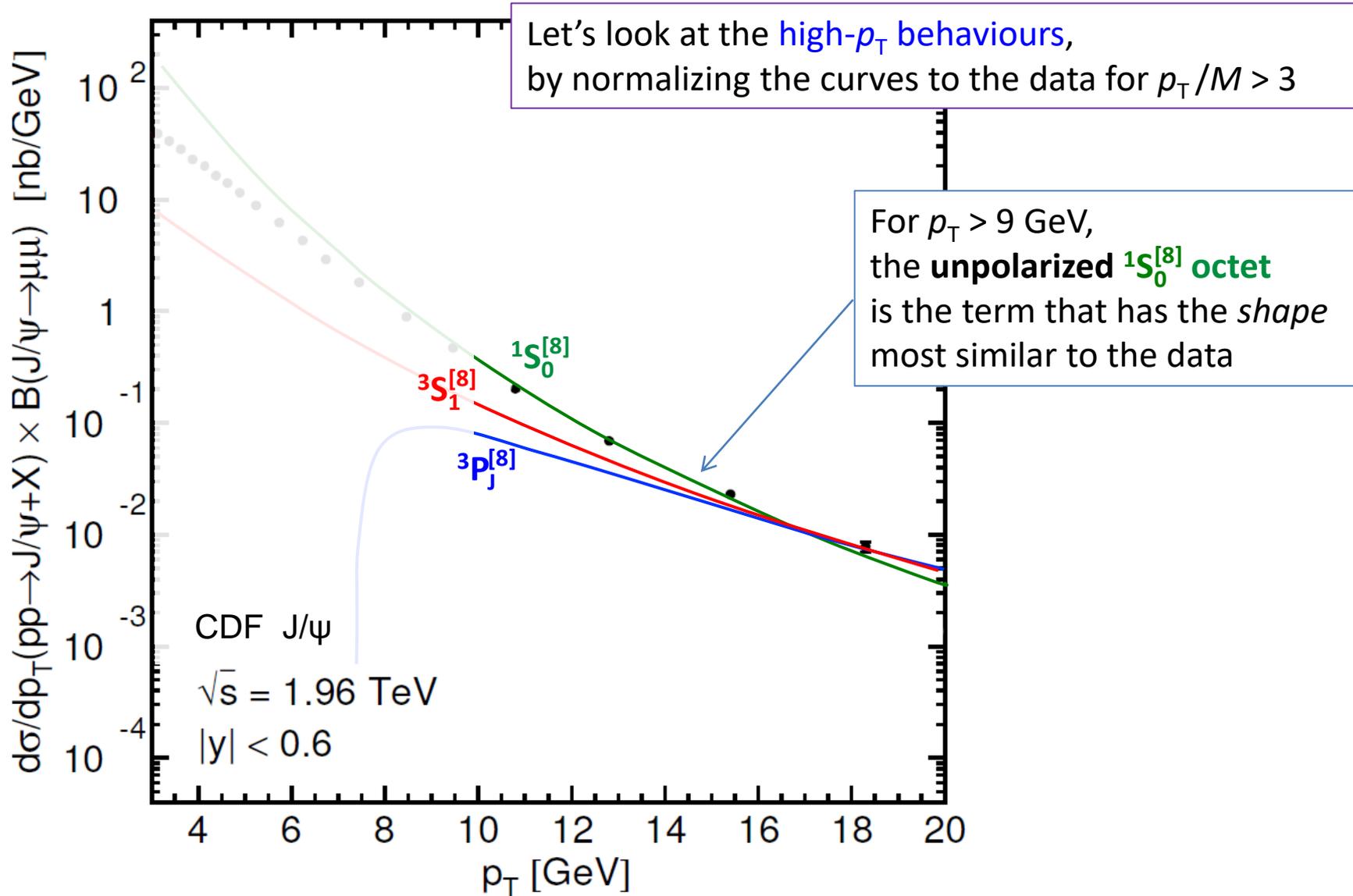
If the fits of the S-wave quarkonium cross sections indicate  $^3S_1$  and  $^3P_J$  dominance, then one predicts transverse polarization ( $\lambda_g \approx +1$ ) at high  $p_T$

The famous “strong transverse polarization” prediction was not a “first principles” result of NRQCD but the outcome of fits of cross section measurements... *and those fits were biased*

# Example of a biased fit



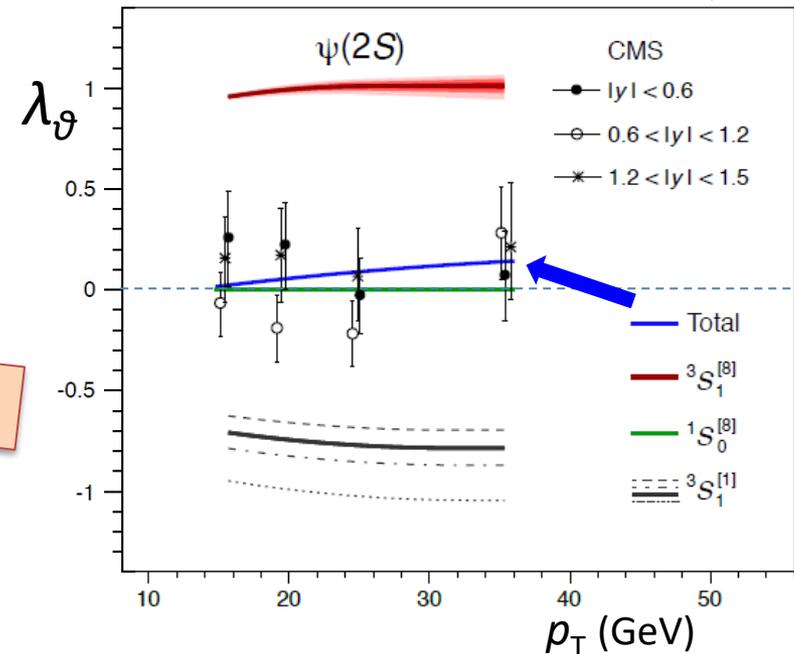
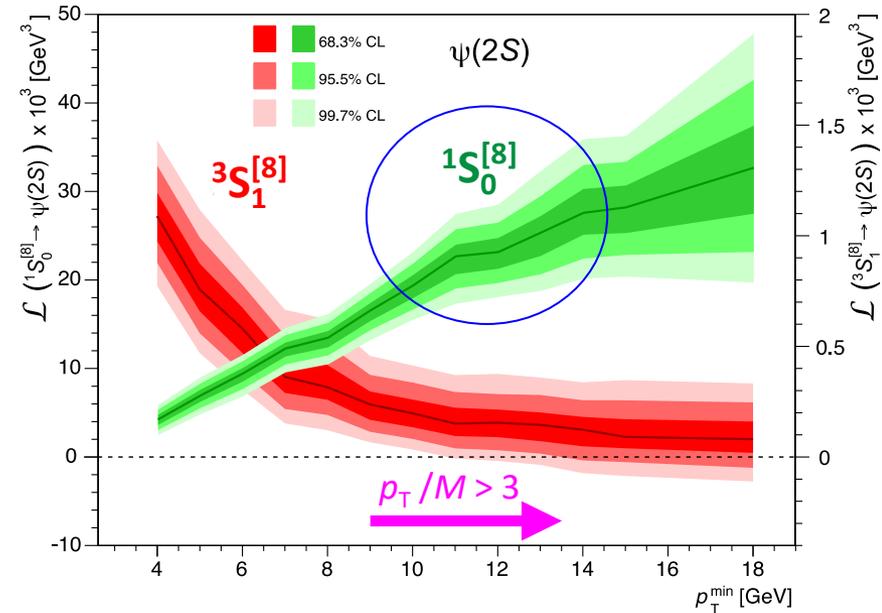
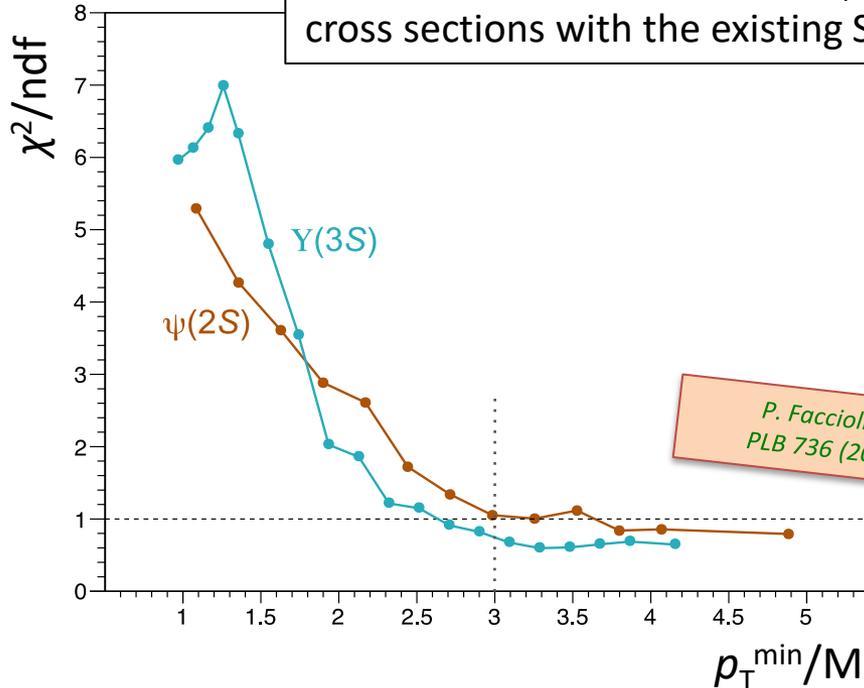
# What happens if we exclude the low $p_T$ region?



# There is no quarkonium polarization puzzle...

For  $p_T/M > 3$ ,  
 the cross sections and polarizations *can* be  
 simultaneously and consistently fitted;  
 the fit results are stable and show that  
 the measurements imply  $1S_0^{[8]}$  dominance

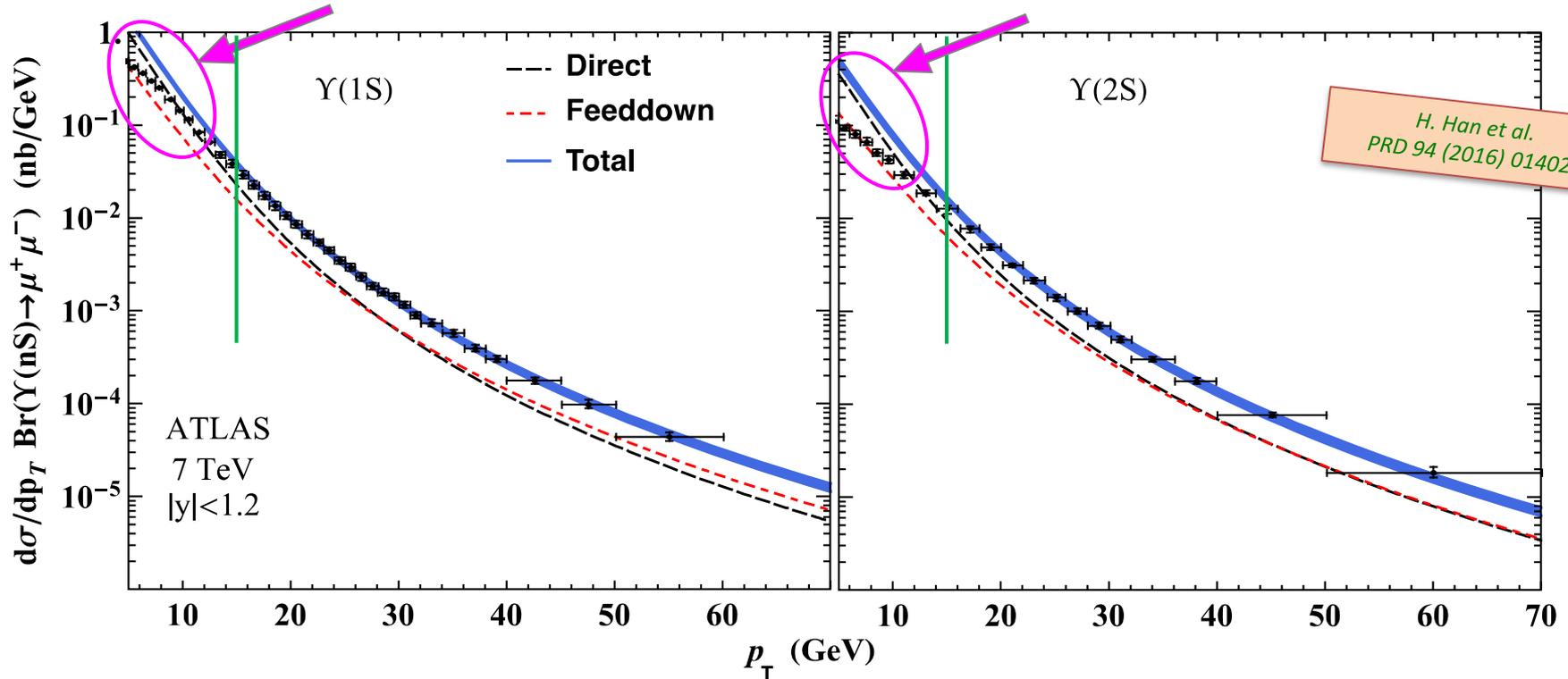
The fit quality improves dramatically  
 if we do not try to fit the low  $p_T/M$   
 cross sections with the existing SDCs



# The low- $p_T$ SDC problem

Current NLO SDCs are not valid in the low  $p_T$  region.

Fitting data with a wrong fit model necessarily leads to wrong fit results.



# A model-independent global charmonium fit (1)

We made a simultaneous fit of the  $J/\psi$ ,  $\psi(2S)$  and  $\chi_{c1,2}$  cross sections and polarizations measured by ATLAS and CMS

Accounting for the **momentum and polarization transfer** from mother to daughter in the **feed-down** decays:

$$\begin{aligned}\psi(2S) &\rightarrow \chi_{c1,2} \gamma \\ \psi(2S) &\rightarrow J/\psi X \\ \chi_{c1,2} &\rightarrow J/\psi \gamma\end{aligned}$$

**Momentum propagation:**  $p_{\top}/m$  (daughter) =  $P_{\top}/M$  (mother)

**Polarization propagation:** calculated in the electric dipole approximation; precisely accounts for the observable dilepton distribution, without higher-order terms

Theory calculations of the production kinematics are not used anywhere; the fit is *exclusively based on empirical parametrizations*

# A model-independent global charmonium fit (2)

1)  $J/\psi$  and  $\psi(2S)$  directly produced cross sections: a superposition of two terms

$$\sigma_{\text{dir}} \propto [(1 - f_p) g_u + f_p g_p] \quad \text{with} \quad g_{u,p} = (p_T/M) \left( 1 + \frac{1}{\beta_{u,p} - 2} \frac{(p_T/M)^2}{\gamma} \right)^{-\beta_{u,p}}$$

- The unpolarized ( $g_u$ ;  $\lambda_\theta = 0$ ) and transversely polarized ( $g_p$ ;  $\lambda_\theta = +1$ ) terms have a common  $\gamma$ .
- The  $g_u$  and  $g_p$  shapes and their relative contributions  $f_p$  are constrained by the polarization data.

2)  $\chi_{c1}$  and  $\chi_{c2}$  cross sections (and their  $J/\psi$  feed-down contributions):

- independent from the (direct)  $\psi$  terms;
- no separation of polarized and unpolarized contributions (study made before any  $\chi_c$  polarization measurements)

The global fit has, hence, four terms and 6 free parameters:

- the unpolarized and polarized  $\psi$  terms, defined by  $\beta_u(\psi)$ ,  $\beta_p(\psi)$ ,  $f_p$ ,  $\gamma$
- plus the  $\chi_{c1}$  and  $\chi_{c2}$  cross sections, defined by  $\beta(\chi_1)$ ,  $\beta(\chi_2)$ ,  $\gamma$

# Cross sections and polarizations are correlated (1)

The detection acceptances depend on the polarization scenario assumed in their evaluation

For each set of parameter values considered while running the fit, the polarizations and cross sections are calculated, for all states, as functions of  $p_T$

The resulting  $\lambda_\theta$  values can be immediately compared to the measured ones

For the cross sections, we first scale the measured values by correction factors calculated for the  $\lambda_\theta$  value under consideration

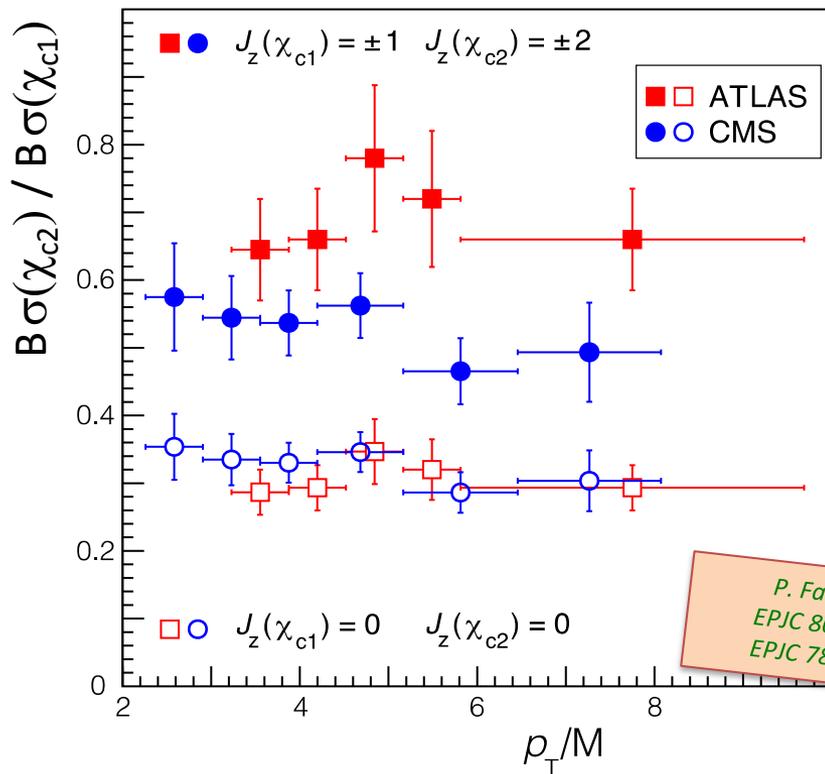
These acceptance corrections are computed using tables published by the experiments for the cross sections of particles produced with fully transverse or fully longitudinal polarization

# Cross sections and polarizations are correlated (2)

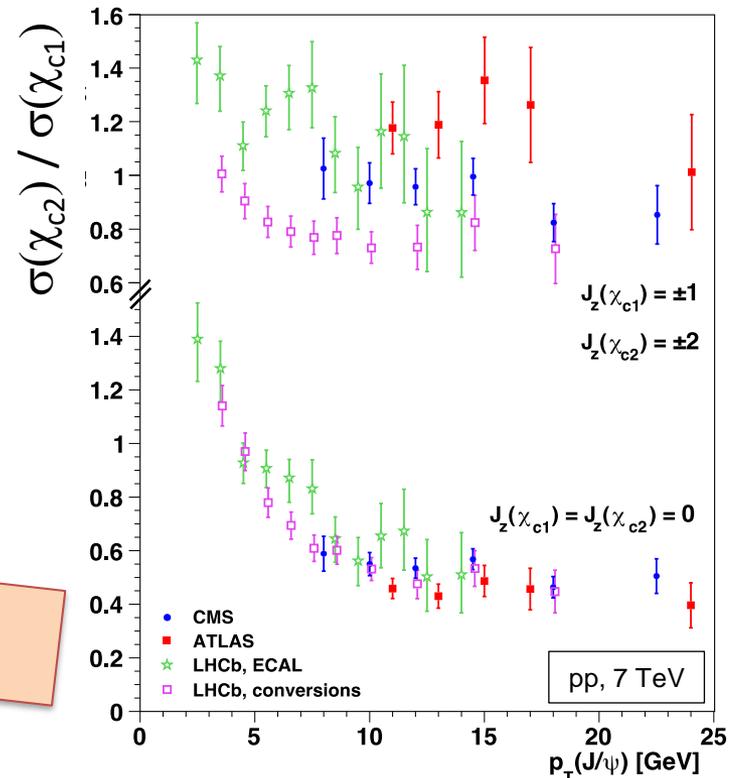
The  $\chi_{c2}$  over  $\chi_{c1}$  cross section ratio provides a good example of the crucial importance of the polarization scenario assumed in the evaluation of the acceptance corrections

Very different patterns and consistencies among data sets are seen for two hypotheses: spin alignments  $J_z(\chi_{c1}) = \pm 1, J_z(\chi_{c2}) = \pm 2$  and  $J_z(\chi_{c1}) = J_z(\chi_{c2}) = 0$

The “default” unpolarized hypothesis leads to intermediate values



*P. Faccioli et al.  
EPJC 80 (2020) 623  
EPJC 78 (2018) 268*



# Fit results

The fit has 100 data points and 10 parameters:

- $\beta_u(\psi)$ ,  $\beta_p(\psi)$ ,  $f_p$ ,  $\beta(\chi_1)$ ,  $\beta(\chi_2)$ ,  $\gamma$
- 4 normalizations

plus 10 nuisance parameters (\*)

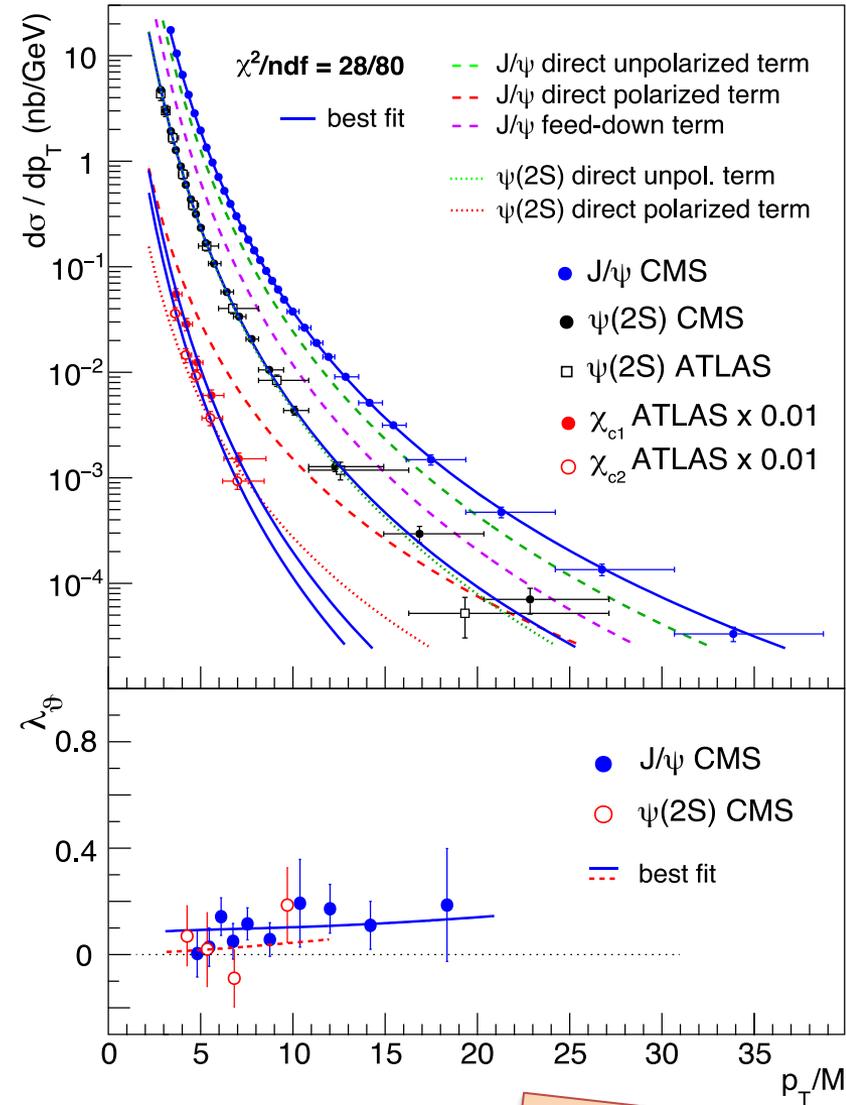
The  $\chi_{c1}$  and  $\chi_{c2}$   $p_T/M$  distributions are very similar to the unpolarized term dominating  $\psi$  production

$$\begin{aligned}\beta_u(\psi) &= 3.42 \pm 0.05 \\ \beta(\chi_1) &= 3.46 \pm 0.08 \\ \beta(\chi_2) &= 3.49 \pm 0.10\end{aligned}$$

The high precision  $\psi$  data constrain the  $\chi_c$  results

The polarized term has a weak contribution: the charmonium states are nearly unpolarized

(\*) The nuisance parameters are: ATLAS and CMS integrated-luminosity uncertainties; uncertainties on the branching ratios (B) used to derive the cross sections

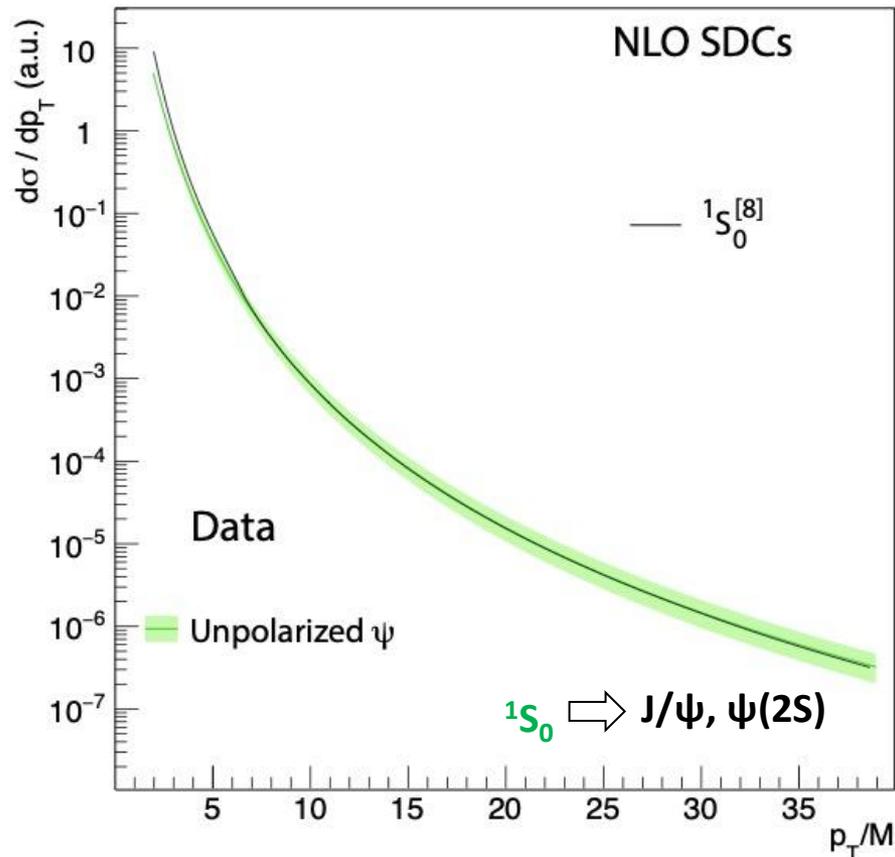


*P. Faccioli et al.  
EPJC 78 (2018) 268*

# Data fit vs. NRQCD (1)

The shape functions from the global fit agree **very well** with their NRQCD counterparts, over 8 orders of magnitude !

The data bands and the NLO SDCs were obtained in **completely independent ways**

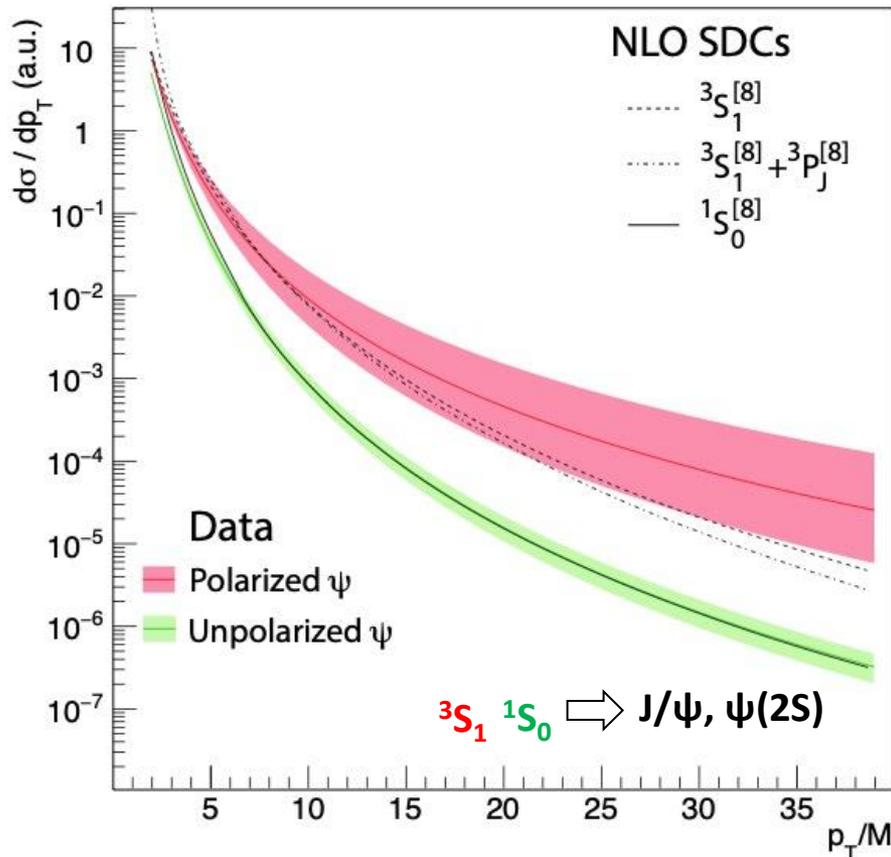


The width of the data bands only reflects *shape* uncertainties

# Data fit vs. NRQCD (2)

The shape functions from the global fit agree **very well** with their NRQCD counterparts, over 8 orders of magnitude !

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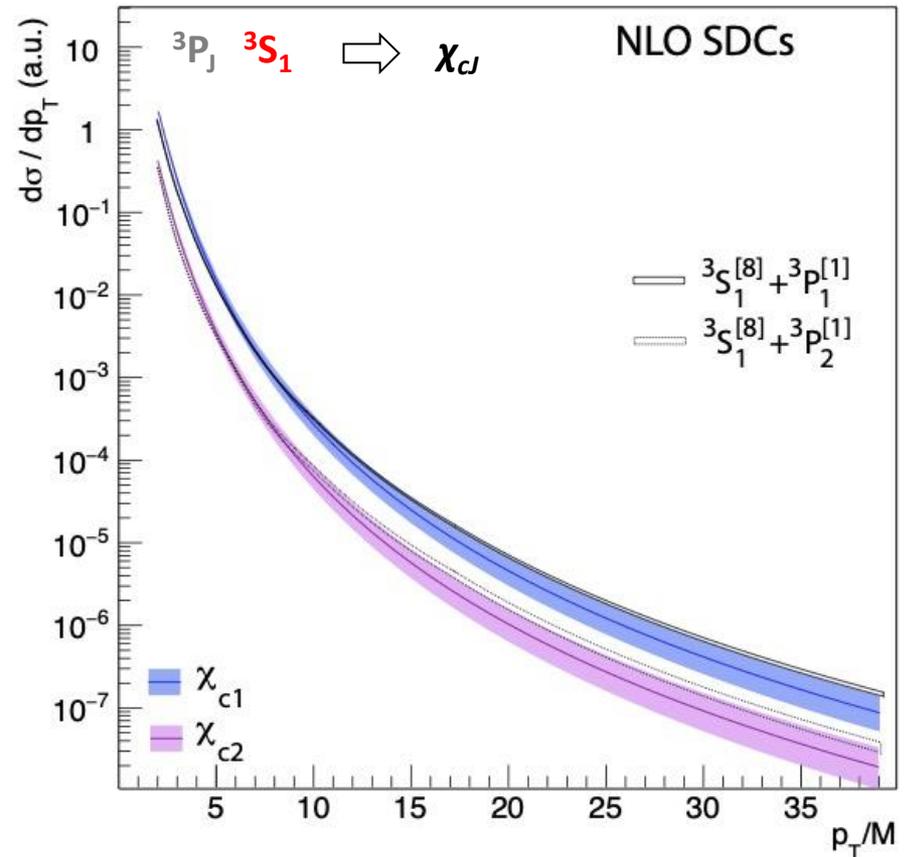
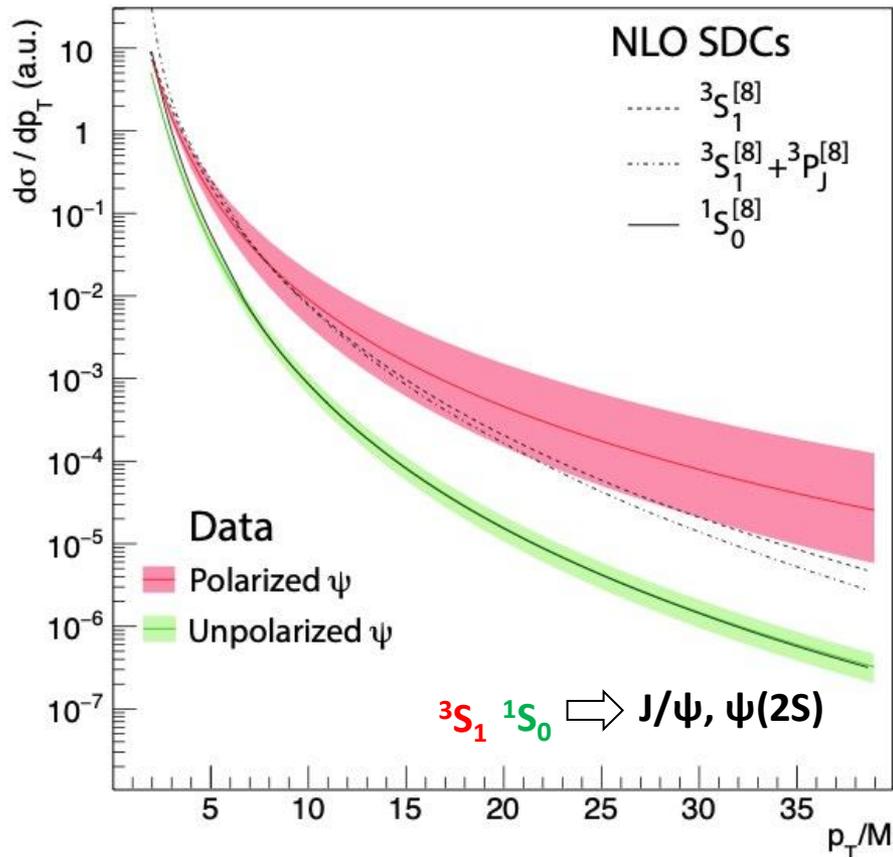


The width of the data bands only reflects *shape* uncertainties

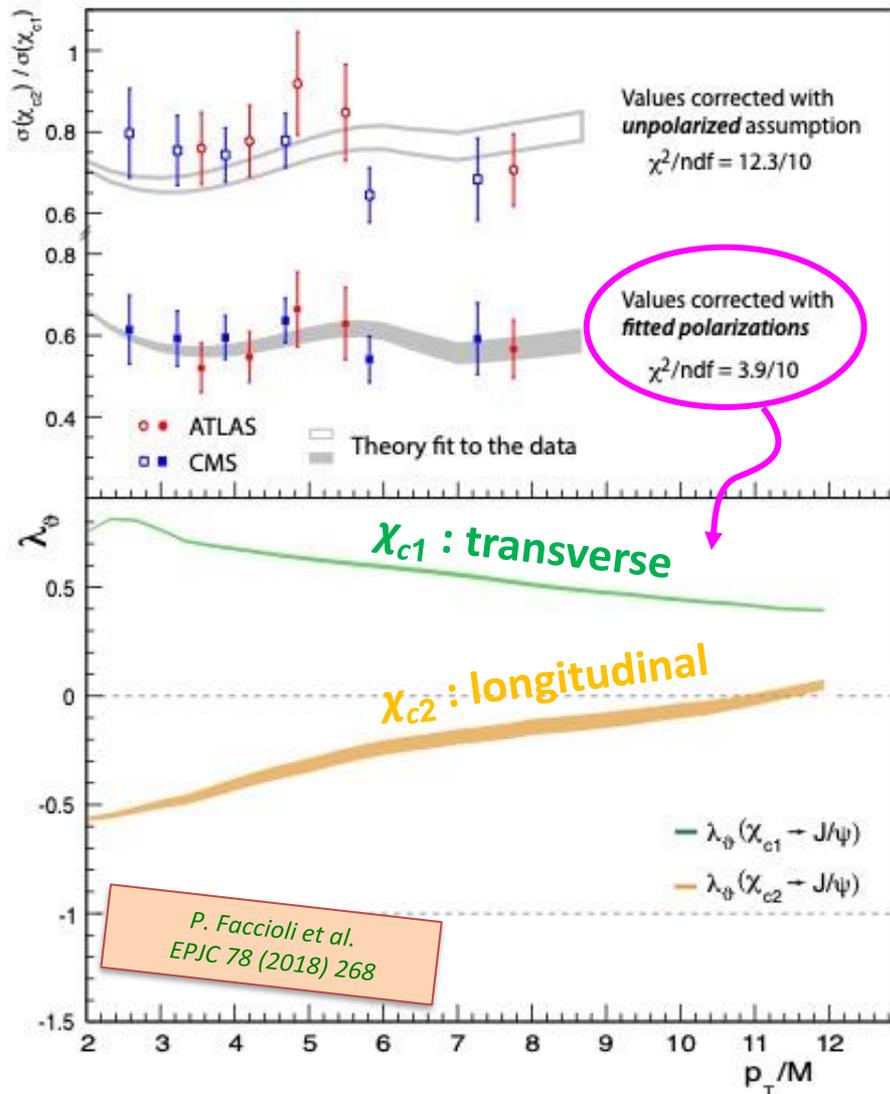
# Data fit vs. NRQCD (3)

The shape functions from the global fit agree **very well** with their NRQCD counterparts, over 8 orders of magnitude !

The data bands and the NLO SDCs were obtained in **completely independent ways** !



# A precise NRQCD prediction



In NRQCD,  $\chi_{c1,2}$  production has two terms: the  $^3S_1$  octet and the  $^3P_{1,2}$  singlet.

One single parameter  $r$  determines

- 1) the  $\chi_{c2} / \chi_{c1}$  yield ratio
- 2)  $\lambda_\theta(\chi_{c1})$
- 3)  $\lambda_\theta(\chi_{c2})$

$$r \equiv m_c^2 \left\langle \mathcal{O}^{\chi_{c0}}(^3S_1^{[8]}) \right\rangle / \left\langle \mathcal{O}^{\chi_{c0}}(^3P_0^{[1]}) \right\rangle$$

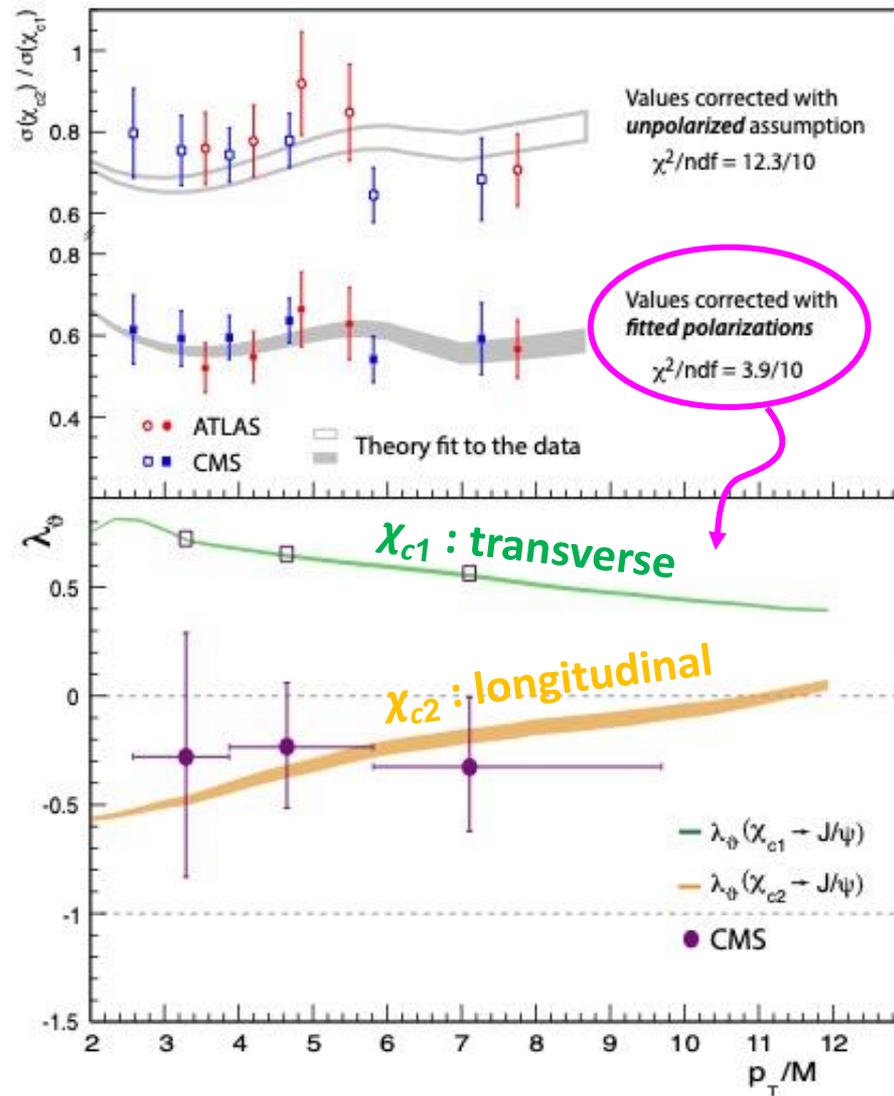
ATLAS and CMS  $\chi_{c2} / \chi_{c1}$  ratios imply

$$r = \mathbf{0.217 \pm 0.003}$$

(reminder: acceptance correction depends on  $\lambda_\theta$ )

Conclusion: within NRQCD, the  $\chi_{c1}$  and  $\chi_{c2}$  polarizations are predicted to be **very different** from one another; a strongly constrained, precise and unambiguous prediction

# A precise NRQCD prediction vs. CMS data



CMS measured the ratio between the  $\cos\vartheta$  distributions of the dimuons associated with the  $\chi_{c2}$  and  $\chi_{c1}$  radiative ( $J/\psi \gamma$ ) decays (\*)

This provides a constraint on the *difference* between the  $\chi_{c2}$  and  $\chi_{c1}$  polarizations

We fix  $\lambda_\theta(\chi_{c1})$  to the prediction, and derive  $\lambda_\theta(\chi_{c2})$

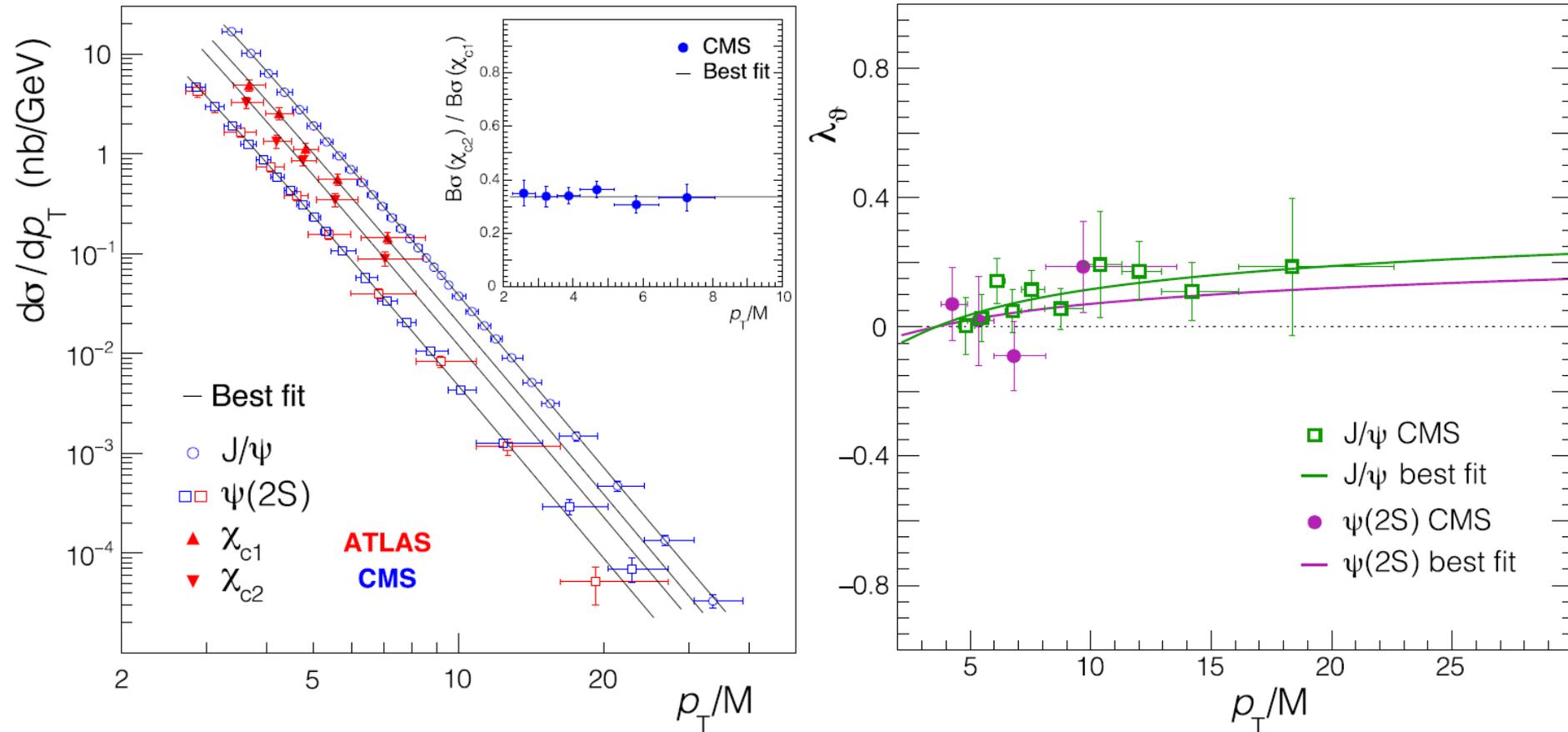
The data is in good agreement with the (quite extreme) **predicted polarizations**

(\*) See previous talk for all the details

# A constraint on the sum of the $\chi_{c1}$ and $\chi_{c2}$ polarizations

The  $J/\psi$ ,  $\psi(2S)$ ,  $\chi_{c1}$  and  $\chi_{c2}$  cross sections measured by ATLAS and CMS, together with the  $J/\psi$  and  $\psi(2S)$  polarizations, constrain *the sum* of the  $\chi_{c1}$  and  $\chi_{c2}$  polarizations

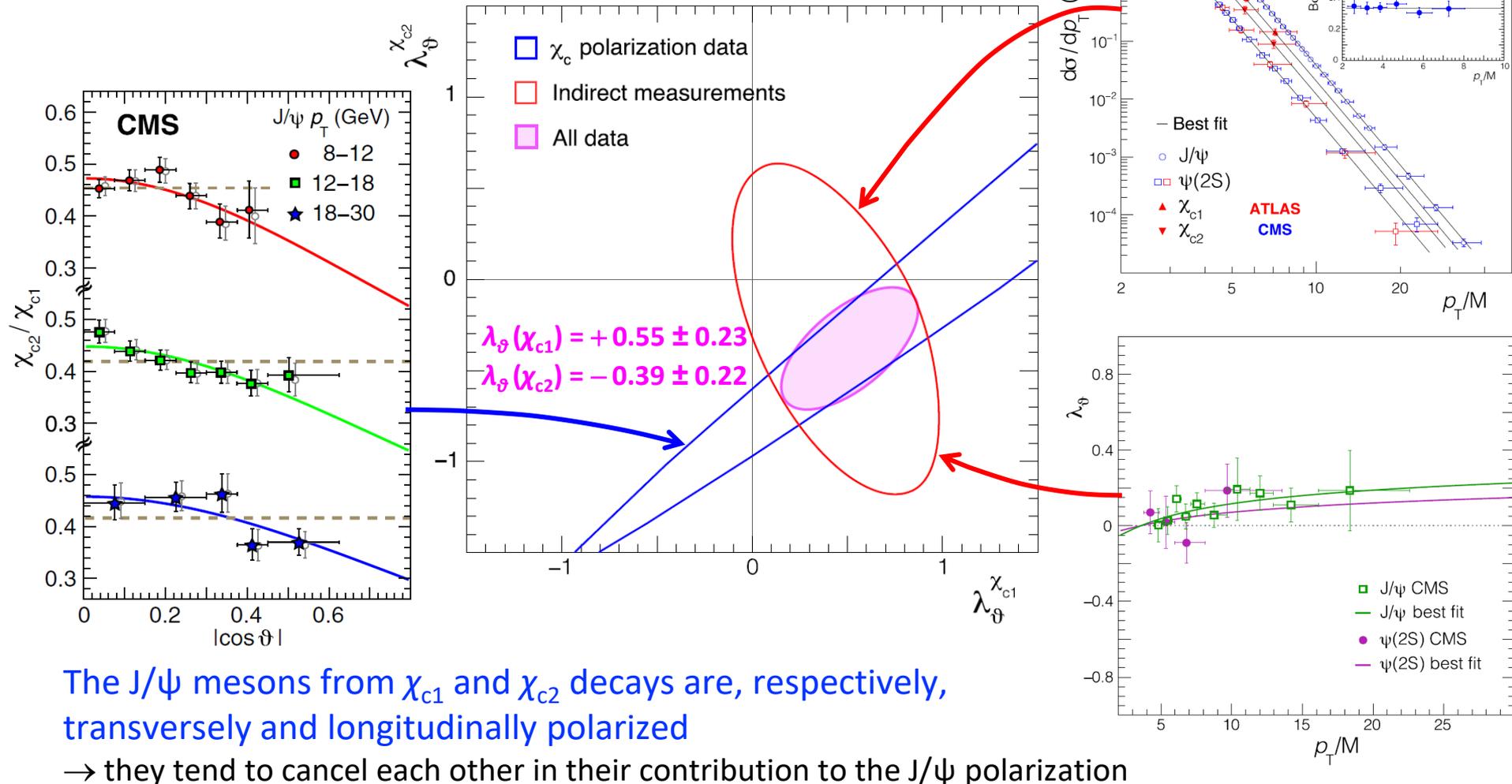
Only assumption: the *directly produced*  $J/\psi$  and  $\psi(2S)$  states have identical polarizations, vs.  $p_T/M$



# The $\chi_{c1}$ and $\chi_{c2}$ states are strongly polarized !

*P. Faccioli et al.  
EPJC 80 (2020) 623*

The combination of these “orthogonal” experimental constraints determines the two individual  $\chi_{c1}$  and  $\chi_{c2}$  polarizations; this is a purely experimental result: no theory involved

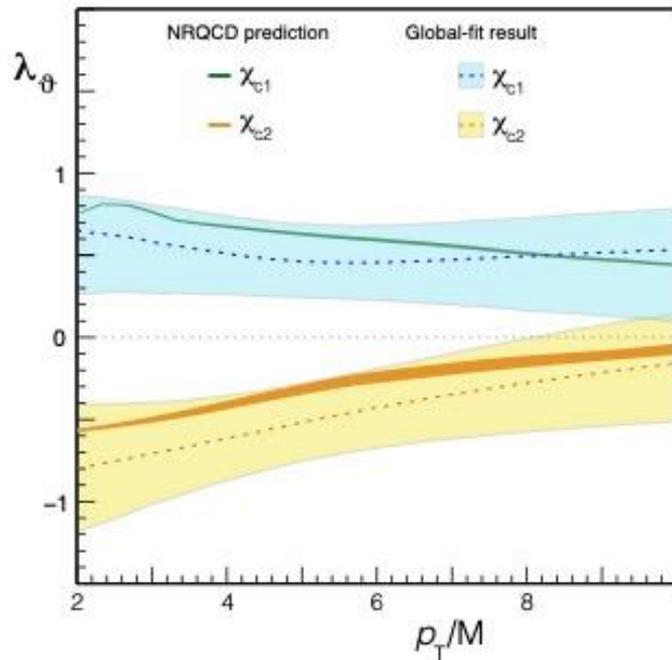


The  $J/\psi$  mesons from  $\chi_{c1}$  and  $\chi_{c2}$  decays are, respectively, transversely and longitudinally polarized

→ they tend to cancel each other in their contribution to the  $J/\psi$  polarization

# The $\chi_{c1}$ and $\chi_{c2}$ polarizations vs. $p_T$

The global fit of all charmonium data also provides results (as a function of  $p_T$ ) for the individual  $\chi_{c1}$  and  $\chi_{c2}$  polarizations



Narrow bands: NRQCD prediction obtained from the ATLAS and CMS  $\chi_{c2} / \chi_{c1}$  ratios;  $r = \mathbf{0.217 \pm 0.003}$

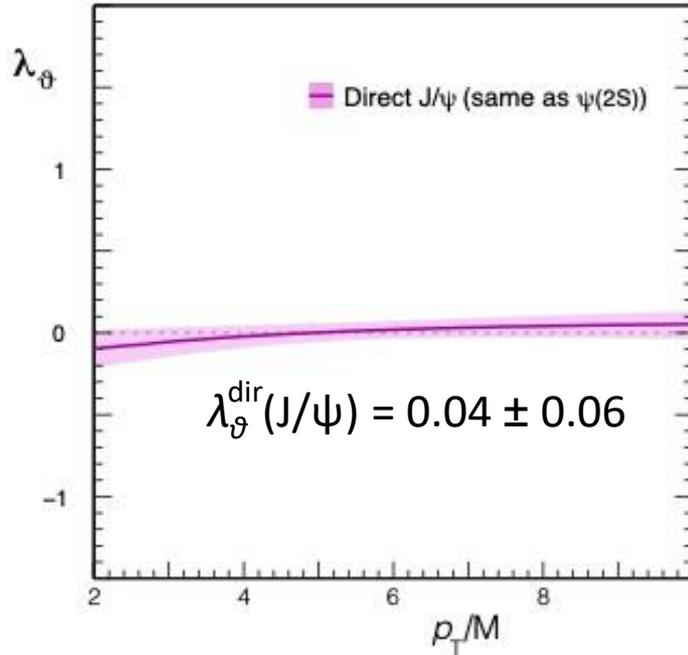
Wide bands: result of the global fit of all charmonium cross sections and polarizations, with no theory inputs

The NRQCD prediction agrees perfectly with the data

An out-of-the-box success of NRQCD !

# The (direct) $J/\psi$ polarizations vs. $p_T$

The global fit of all charmonium data also provides results (as a function of  $p_T$ )  
for the *directly produced*  $J/\psi$  and  $\psi(2S)$  polarizations



Very strong evidence of  
unpolarized  $J/\psi$  and  $\psi(2S)$  production !

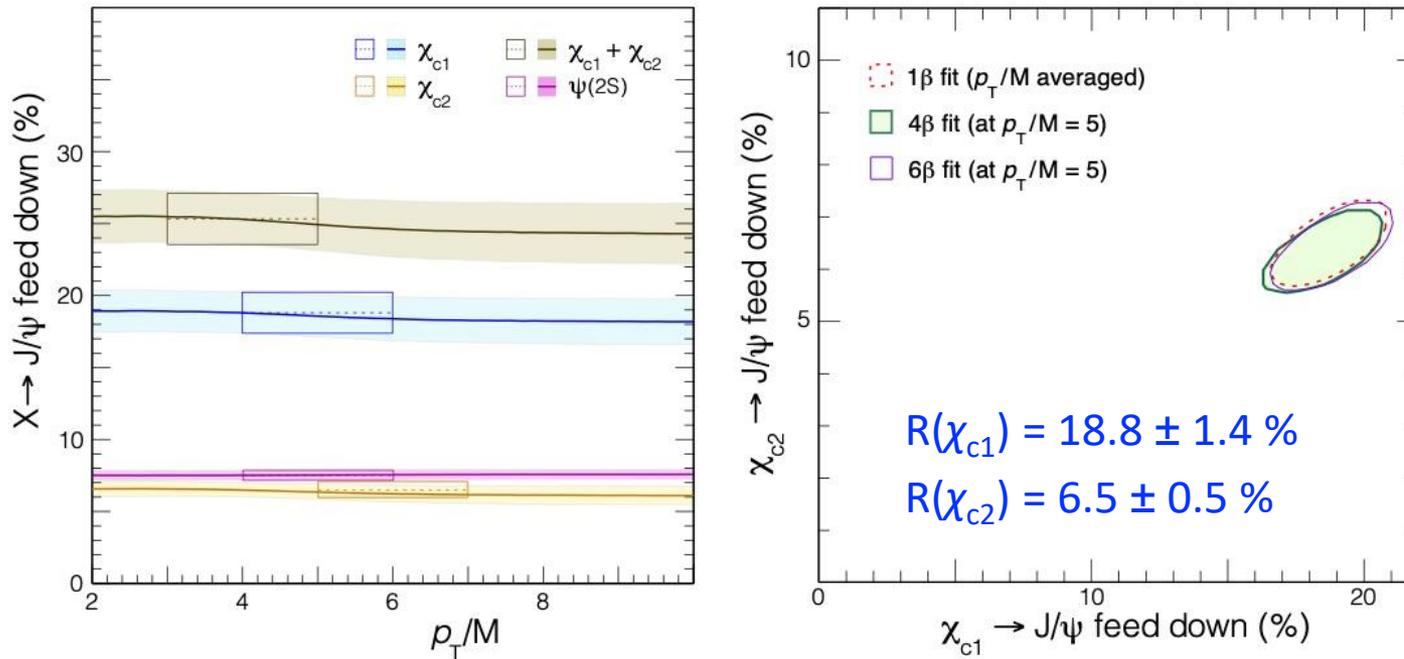
**Zero** and **constant** polarization: a big challenge  
to production models of vector quarkonia

It is very unlikely that we are seeing a fine-tuned  
cancellation of a mixture of subprocesses

→ a clear sign of the **unique nature and production mechanism** of heavy quarkonia

# The $\chi_{c1}$ and $\chi_{c2}$ feed-down fractions

Finally, the global-fit analysis also determines the fractions of  $J/\psi$  mesons produced from the feed-down decays of the  $\chi_{c1}$  and  $\chi_{c2}$  mesons



The fraction of the prompt  $J/\psi$  yield due to directly-produced mesons is  $67.2 \pm 1.9 \%$ , a remarkably precise value

# Summary

1) The methodology of global-fit analyses has a strong impact on the obtained results

The handling of correlations and uncertainties, both experimental and theoretical (including limitations in the validity domain of the perturbative calculations) is not trivial and can easily bias the results

The puzzling coexistence of several very different “NRQCD predictions” for the polarization was entirely due to the inconsistent treatments of data and theory

2) With a correct treatment of experimental and theoretical correlations and limitations, NRQCD describes remarkably well the  $J/\psi$ ,  $\psi(2S)$ ,  $\chi_{c1}$  and  $\chi_{c2}$  yields and polarizations measured at mid rapidity

3) The new CMS measurement of  $\chi_{c1}$  and  $\chi_{c2}$  polarizations allows us to extract, for the first time, the polarization of the *directly* produced  $J/\psi$  mesons

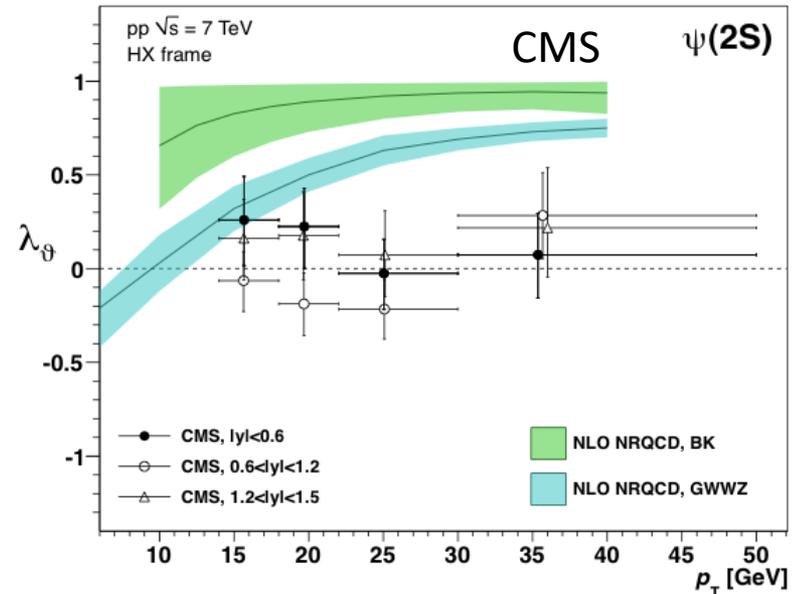
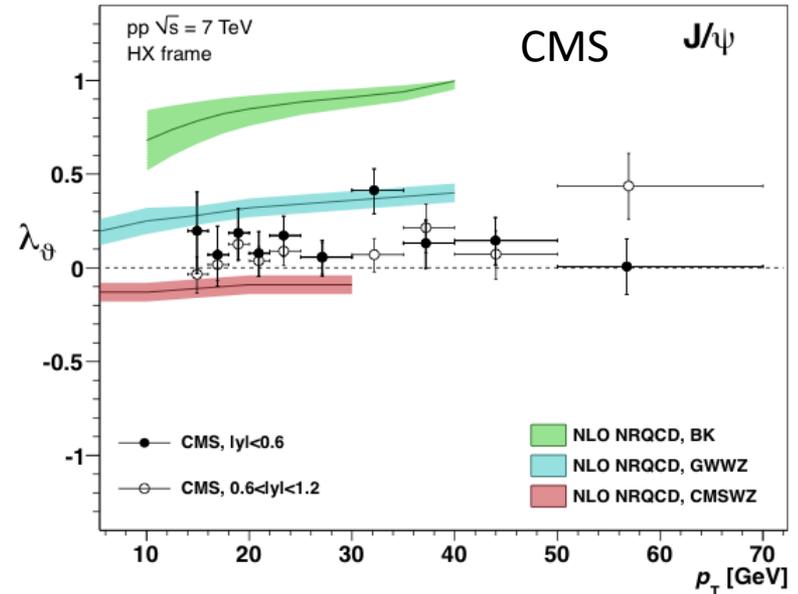
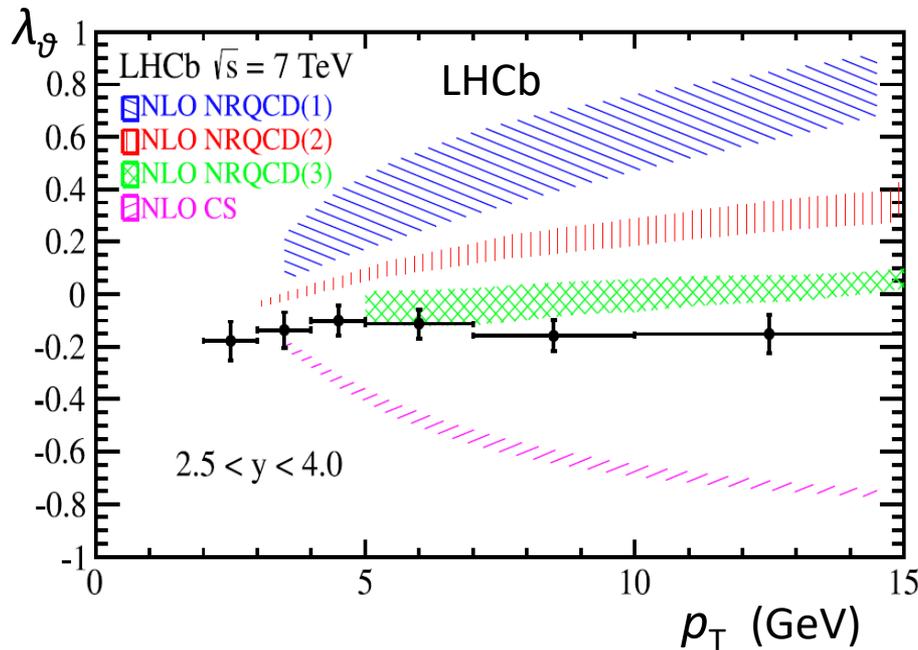
The result,  $\lambda_9^{\text{dir}}(J/\psi) = 0.04 \pm 0.06$ , with no  $p_T$  dependence, is a *conceptual challenge* for NRQCD and any other production model

Backup

# Inaccurate fit methodologies give unreliable results

The same theory inputs (NLO SDCs) result in very different quarkonium polarization “NRQCD predictions” simply because of inconsistent treatments of data and theory

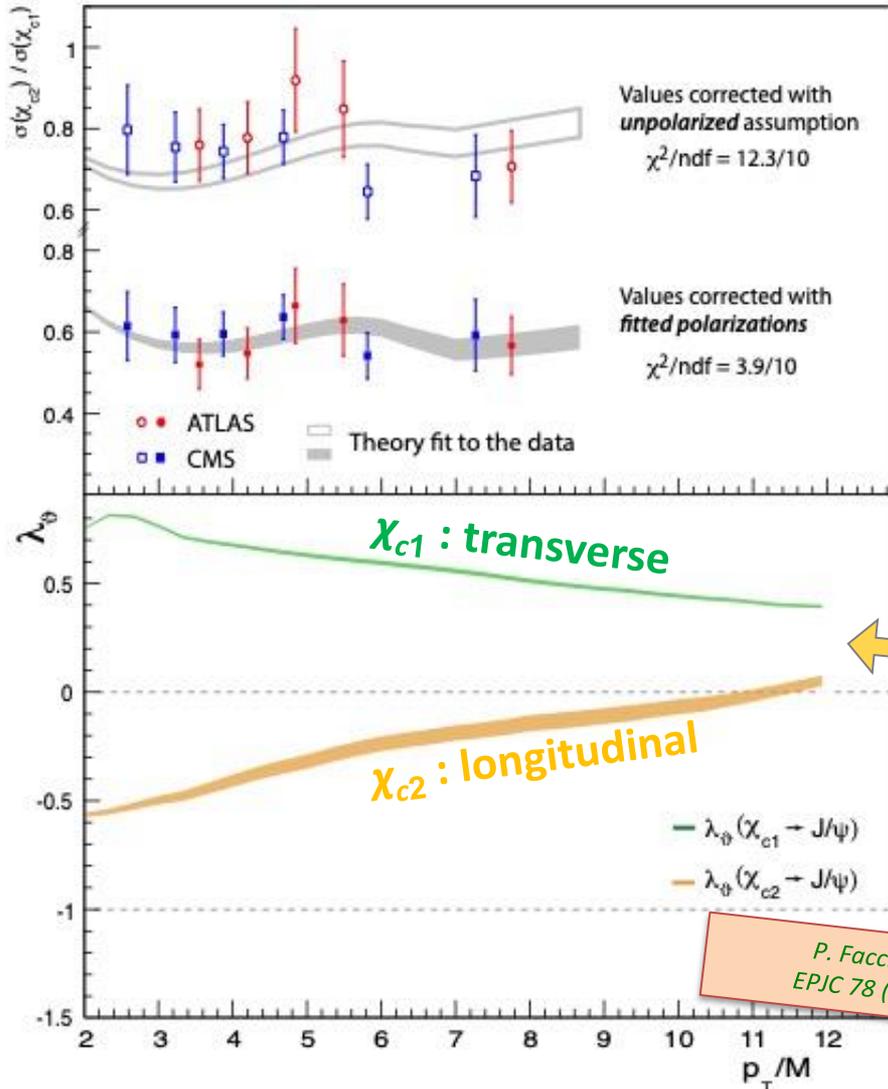
With robust analysis techniques and high quality data, quarkonium production measurements can become a **high precision probe of QCD** !



# Comparison between two predictions (1)

In NRQCD, one single parameter determines *both* the  $\chi_{c2} / \chi_{c1}$  ratio and the two polarizations

$$r \equiv m_c^2 \left\langle \mathcal{O}^{\chi_{c0}}(^3S_1^{[8]}) \right\rangle / \left\langle \mathcal{O}^{\chi_{c0}}(^3P_0^{[1]}) \right\rangle$$



Faccioli et al. derive  $r = 0.217 \pm 0.003$  from CMS + ATLAS data (averaged) with acceptance corrections corresponding to the *final* polarization prediction (*iterative* procedure) and, thus, no added “polarization uncertainty”

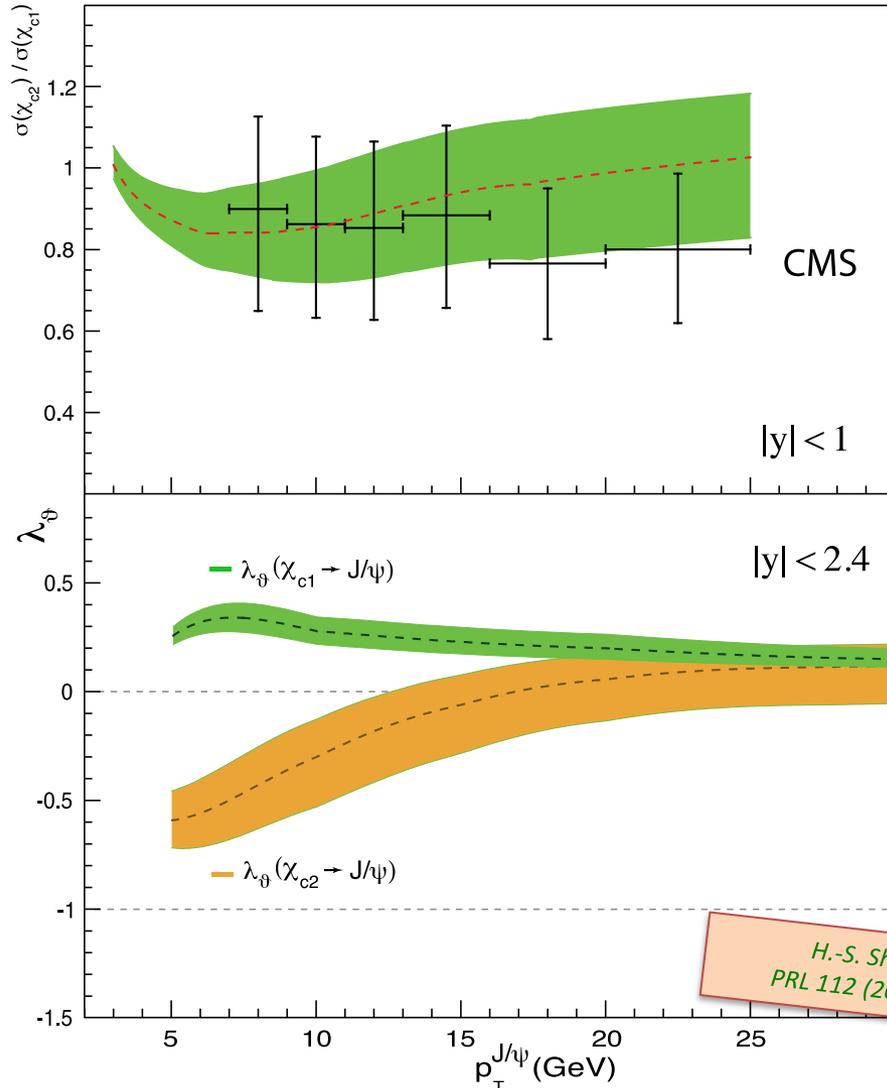
← A very precise prediction

P. Faccioli et al.  
 EPJC 78 (2018) 268

# Comparison between two predictions (2)

In NRQCD, one single parameter determines *both* the  $\chi_{c2} / \chi_{c1}$  ratio and the two polarizations

$$r \equiv m_c^2 \left\langle \mathcal{O}^{\chi_{c0}}({}^3S_1^{[8]}) \right\rangle / \left\langle \mathcal{O}^{\chi_{c0}}({}^3P_0^{[1]}) \right\rangle$$



H.-S. Shao et al.  
PRL 112 (2014) 182003

Shao et al. derive  $r = 0.27 \pm 0.06$  from CDF or CMS data with the following procedure:

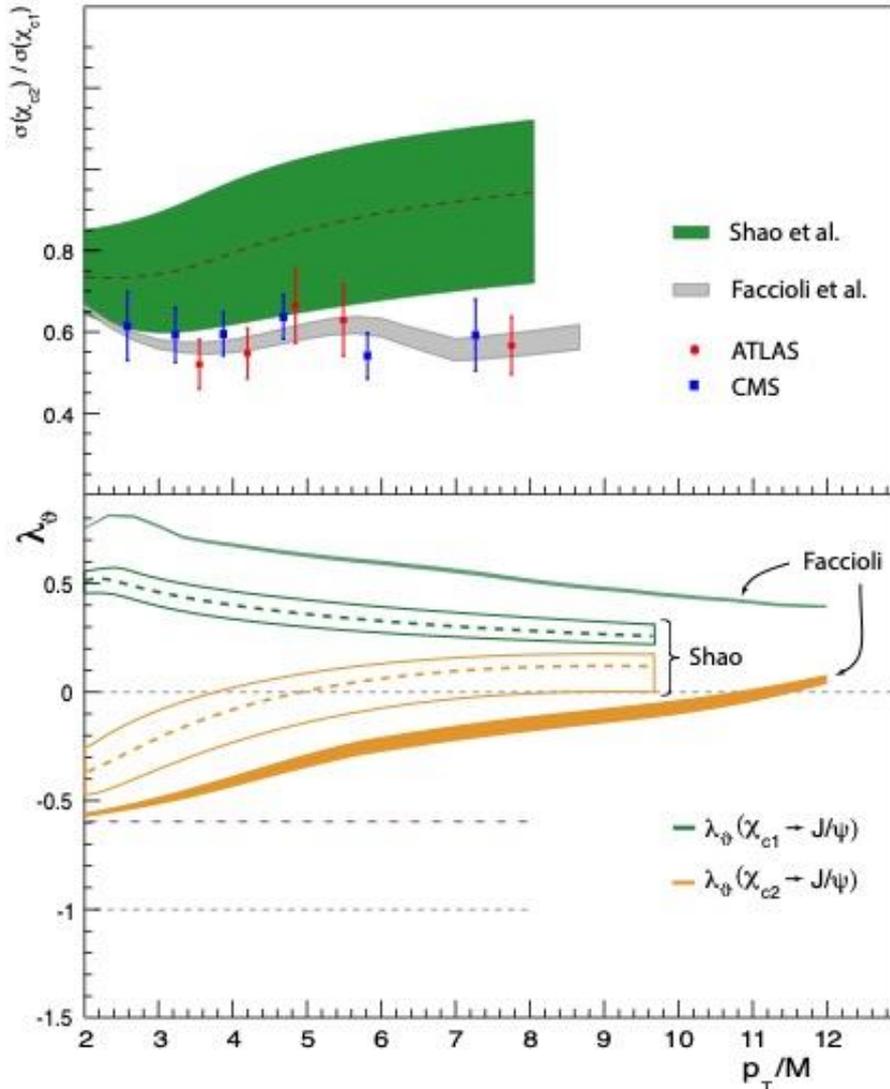
- CDF:
- = central values using  $\lambda_\theta = 0.13 \pm 0.15$  for  $\chi_{c1}$  and  $\chi_{c2}$
  - = no correlated variations considered
  - = uncertainty added in quadrature with all others

- CMS:
- = central values using  $\lambda_\theta = 0$  for  $\chi_{c1}$  and  $\chi_{c2}$
  - = polarization uncertainty from *maximum* range of correlated variations of  $\lambda_\theta(\chi_{c1})$  and  $\lambda_\theta(\chi_{c2})$

# Comparison between two predictions (3)

In NRQCD, one single parameter determines *both* the  $\chi_{c2} / \chi_{c1}$  ratio and the two polarizations

$$r \equiv m_c^2 \left\langle \mathcal{O}^{\chi_{c0}}(^3S_1^{[8]}) \right\rangle / \left\langle \mathcal{O}^{\chi_{c0}}(^3P_0^{[1]}) \right\rangle$$



Same theory inputs but *different analyses of the experimental data* lead to very different determinations of  $r$

Shao et al., PRL 112 (2014) 182003  
 $r = 0.27 \pm 0.06$

Faccioli et al., EPJC 78 (2018) 268  
 $r = 0.217 \pm 0.003$

This shows how crucial it is to rigorously treat the correlations between the cross sections and the polarizations and to properly account for the uncertainties

There is no “*polarization uncertainty*”; instead, there is a *correlation* between the polarization and the cross sections