

Determining the initial conditions and transport properties of quark-gluon plasma by flow measurements at the LHC

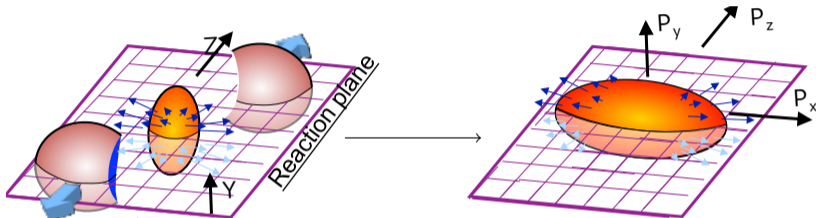
Jasper Parkkila, Dong Jo Kim

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Anisotropic Flow

Initial geometry fluctuations → **Transport $\delta_\mu T^{\mu\nu} = 0$** → final-state particles



$$\varepsilon_n e^{in\Phi_n} \equiv -\frac{\langle r^n e^{in(\phi - \Phi_n)} \rangle}{\langle r^n \rangle}, n \geq 2. \quad (1)$$

(theory only - Initial State models)

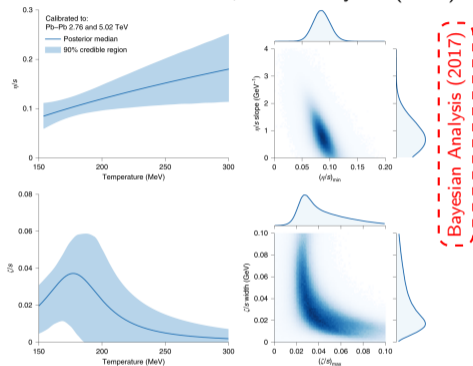
$$\frac{dN}{d\phi} \propto \frac{1}{2\pi} \sum_{n=-\infty}^{\infty} \underbrace{\langle e^{in\phi} \rangle}_{V_n} e^{-in\phi}, \quad (2)$$

where $V_n \equiv \langle e^{in\phi} \rangle = v_n e^{in\psi_n}$. (experiments, theory - hydro+hadronization models with $\eta/s(T)$, $\zeta/s(T)$)

- Collectivity as a probe to the properties of the medium – transport properties such as $\eta/s(T)$, $\zeta/s(T)$, IS

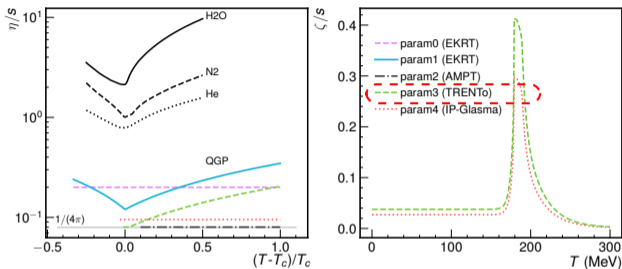
Current understanding of the medium properties

Steffen A. Bass et. al, Nature Physics (2019)



- ALICE data on multiplicity, spectra and flow are key inputs to estimate the properties of the QGP (including p–Pb data), i.e Global Bayesian Analysis and other theory groups.

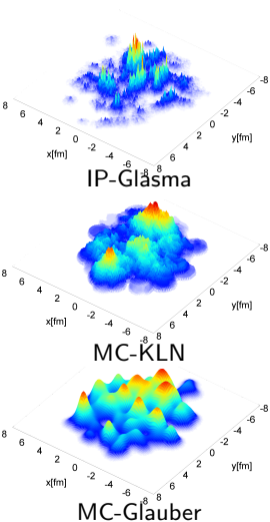
Various parameterizations used in latest hydrodynamic calculations



- Best fit seems to indicate $\eta/s \approx 0.12$ around $T_c \approx 150$ MeV, very close to $1/4\pi$ (≈ 0.08) from string theory¹(AdS/CFT correspondence).
- $\eta/s(T)$ and $\zeta/s(T)$ should be constrained further (larger uncertainties) by separating the effects from the initial conditions.

¹D. T. Son et. al. Phys. Rev. Lett. 94 (2005) 111601

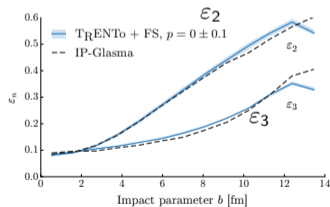
Model summary and parameterizations



B. Schenke, P. Tribedy and R. Venugopalan, Phys.Rev.Lett. 108 252301 (2012)

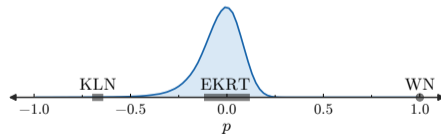
Model	Hydrodynamic code	Initial conditions	η/s	ζ/s
EKRT+param0	EbyE	EKRT	0.20	0
EKRT+param1	EbyE	EKRT	$\eta/s(T)$	0
AMPT ¹ +param2	iEBE-VISHNU	AMPT	0.08	0
T _R ENTo+param3	iEBE-VISHNU	T _R ENTo($p = 0$)	$\eta/s(T)$	$\zeta/s(T)$
IP-Glasma+param4	MUSIC	IP-Glasma	0.095	$\zeta/s(T)$

- State of the art models compared to latest measurements
- Many models calibrated to reproduce lower order harmonic observables



$$\epsilon \sim f(T_A, T_B)$$

$$\epsilon_n e^{in\Phi_n} = -\{r^n e^{in\varphi}\} / \{r^n\}$$

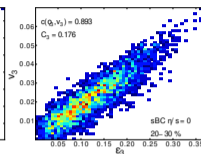
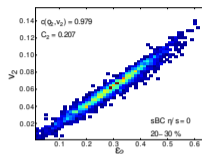


Phys.Rev. C94, 024907 (2016)

¹HIJING with AMPT for initial thermalization

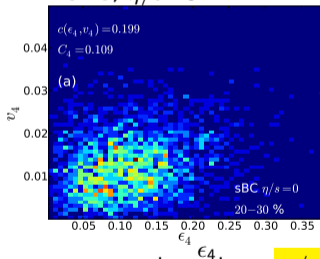
Non-linearity of the higher order flow, $\varepsilon_n \propto v_n$ holds only for $n = 2, 3$

$$n = 2, 3 (v_n \propto \varepsilon_n)$$

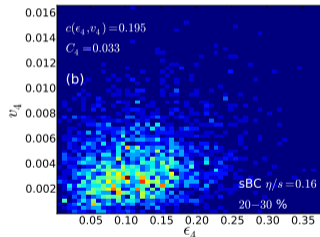


sBC, $\eta/s=0$.

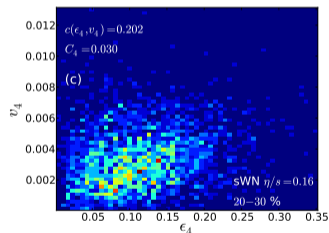
$n = 4$



sBC, $\eta/s=0.16$



sWN, $\eta/s=0.16$



20 – 30%

- v_4 response to ε_4 depends on η/s as well as the initial conditions.¹
- For a rather minimal value of $\eta/s = 1/4\pi$, larger contributions from non-linear corrections.²

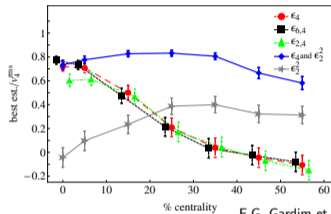
¹H. Niemi et al., Phys. Rev. C 87, 054901 (2013)

²D. Teaney, L. Yan Phys. Rev. C 86 (2012) 044908

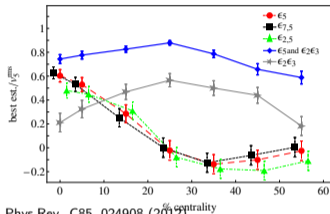
Non-linearity of the higher order flow and cross-harmonic decomposition

- Decomposition into **linear** and **non-linear** contributions

$$v_4 e^{im\psi_4} = k \epsilon_4 e^{4i\Phi_4} + k' \epsilon_2^2 e^{4i\Phi_2}$$



$$v_5 e^{im\psi_5} = k \epsilon_5 e^{5i\Phi_5} + k' \epsilon_2 e^{2i\Phi_2} \epsilon_3 e^{3i\Phi_3}$$



$$\begin{aligned}
 V_4 &= V_{4L} + \chi_{4,22} V_2^2 \rightarrow v_{4,22} = \chi_{4,22} \langle |V_2|^4 \rangle^{\frac{1}{2}} \\
 V_5 &= V_{5L} + \chi_{5,32} V_2 V_3 \rightarrow \dots \\
 V_6 &= V_{6L} + \chi_{6,222} V_2^3 + \chi_{6,33} V_3^2 + \chi_{6,24} V_2 V_{4L} \\
 &\dots
 \end{aligned}$$

(3)

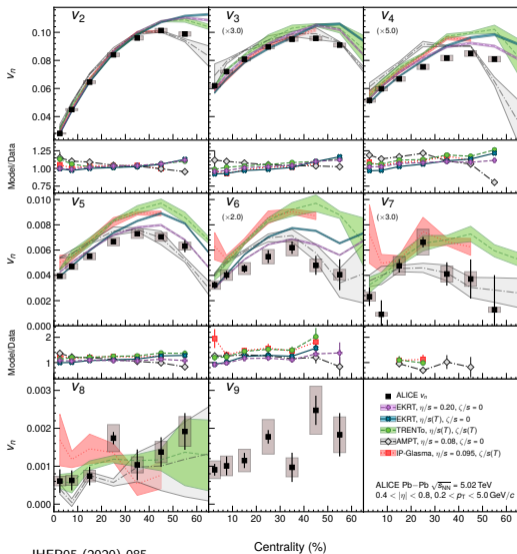
The magnitude of the **non-linear** contribution and non-linear flow mode coefficients:

$$\begin{aligned}
 v_{4,22} &= \frac{\Re \langle V_4 (V_2^*)^2 \rangle}{\sqrt{\langle |V_2|^4 \rangle}} \\
 &\approx \langle v_4 \cos(4\psi_4 - 4\psi_2) \rangle, \quad (4) \\
 \chi_{4,22} &= \frac{v_{4,22}}{\sqrt{\langle v_2^4 \rangle}}.
 \end{aligned}$$

Linear part is extracted from the total and **non-linear** contributions:

$$\underbrace{\langle |V_{4L}|^2 \rangle^{\frac{1}{2}}}_{v_{4L}} = \left(\underbrace{\langle |V_4|^2 \rangle}_{v_4^2} - \underbrace{\chi_{4,22}^2 \langle |V_2|^4 \rangle}_{v_{4,NL}^2} \right)^{\frac{1}{2}}. \quad (5)$$

Anisotropic flow in heavy-ion collisions

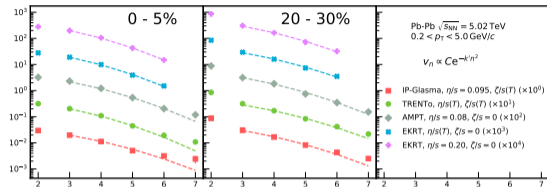
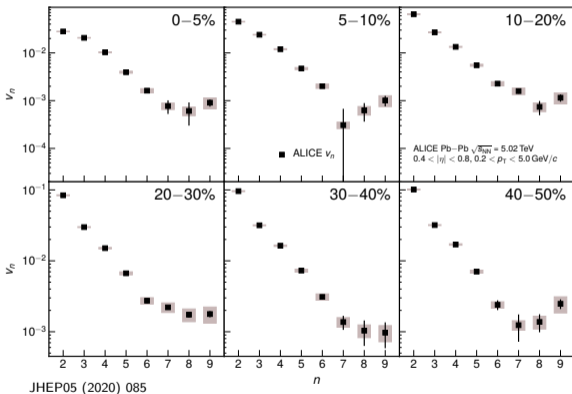


JHEP05 (2020) 085

Centrality (%)

- Hydrodynamic calculations agree well with the data up to v_3 .
- AMPT ($\eta/s = 0.08$) describes the data best in mid-peripheral collisions, but fails to describe the central collisions for $n = 4, 5$
- IP-Glasma, TRENTo: good description for $n = 2, 3$, overestimations at $n \geq 4$
- EKRT ($\eta/s = 0.20$): best agreement among all model configurations at $n \geq 4$
- EKRT ($\eta/s(T)$): comparable up to mid-central collisions, but overestimate the data for peripheral collisions for $n = 4, 5$
- As n increases, the sensitivity to model parameterizations gets larger

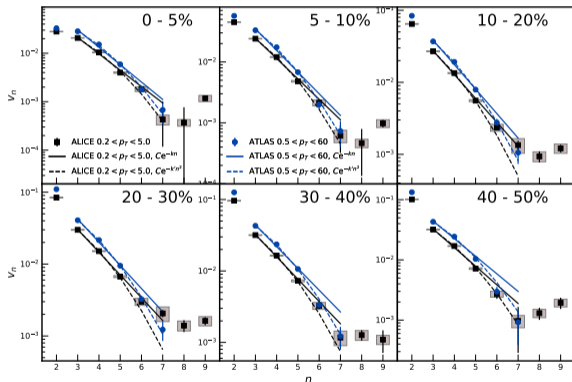
Flow power spectrum



- Power spectra is measured up to v_9 in various centrality bins.
- Clear decrease in magnitude $v_n \propto e^{-k'n^2}$ (viscous damping¹) is observed as n increases
 - Clear damping also observed in hydrodynamical calculations
 - Slope of the calculations is dependent on the model parameterizations
- Interesting feature predicted in acoustic model² for $n \geq 7$ can be further investigated with 2018 data

¹E. Shuryak, PRC84,044912 (2011), R. Lacey et. al. arXiv:1301.0165

²E. Shuryak, PRC84,044912 (2011), Universe 3 (2017) no.4, 75

Damping depends on p_T 

$$v_n \propto e^{-kn}$$

$$v_n \propto e^{-k'n^2}$$

ATLAS, Eur. Phys. J. C 78 (2018) 997

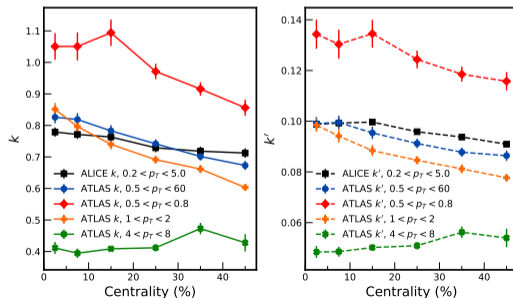
- Both functions work well for $3 \leq n \leq 5$. Note that p_T ranges are different.

- $0.2 < p_T < 5.0 \text{ GeV}/c$

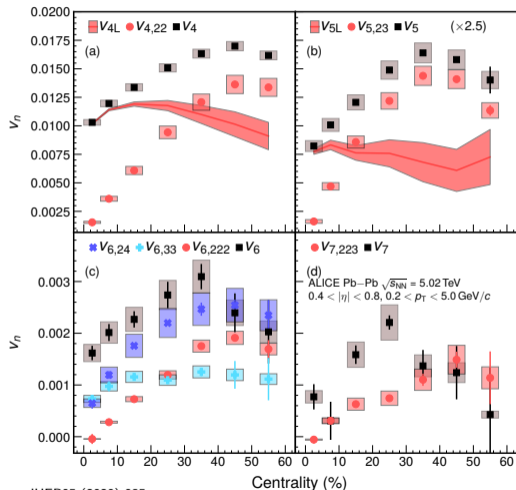
- $0.5 < p_T < 60.0 \text{ GeV}/c$

- $e^{-k'n^2}$ works better for $n > 5$

- For $n > 7$, clear deviation from both fits.



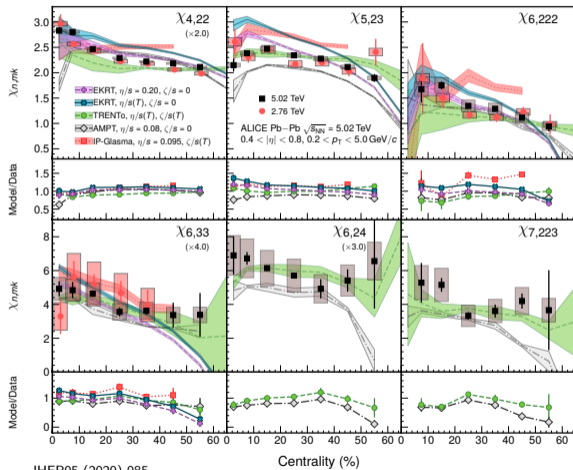
Linear and non-linear decomposition up to 7th order



- **Linear** component dominant in central collisions
- **Non-linear** component increasingly dominant in mid-central to peripheral collisions
- Strength of the non-linear flow mode depends on the harmonic order.

Measured quantities v_n and $v_{n,m}$, from which the **linear** component is derived.

Non-linear flow mode coefficients, hydrodynamic predictions

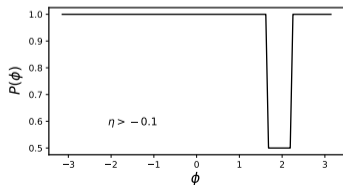
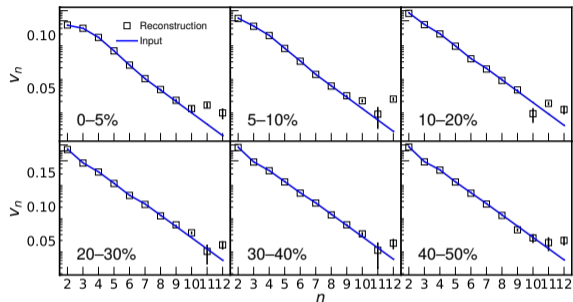


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- Clear centrality dependence for all harmonics
 - Non-linear response at high harmonics larger
 - Large disagreements between the data and the models for $\chi_{5,23}$.
- Model parameterizations need further tuning to capture the magnitude and the centrality dependence.

Toward measurement of ultra-high harmonics

ToyMC:



- 1 Damped input according to $v_n \propto e^{-k'n^2}$, $n \geq 7$
- 2 ϕ -gap modulation
- 3 Measured multiplicity

-
- Plausibility of ultra-high harmonic results. Reconstruction to be verified with ToyMC simulations.
 - 35M event ToyMC closure test, converged up to v_{11}
 - $n > 11$ needs further investigation

Summary

General:

- The flow coefficients, flow modes and non-linear flow mode coefficients of the charged hadrons are measured up to the 9th and 7th harmonic, respectively.
- Cross harmonic decomposition: higher order harmonic non-linear flow mode is more sensitive η/s , ζ/s parameterizations.
- Better constraints on initial conditions and $\eta/s(T)$, $\zeta/s(T)$ with improved precision and extended harmonic orders.

Hydrodynamic models:

- Good agreement for low order harmonic v_n . However, higher orders are not well reproduced.
- Set for the future Bayesian analysis

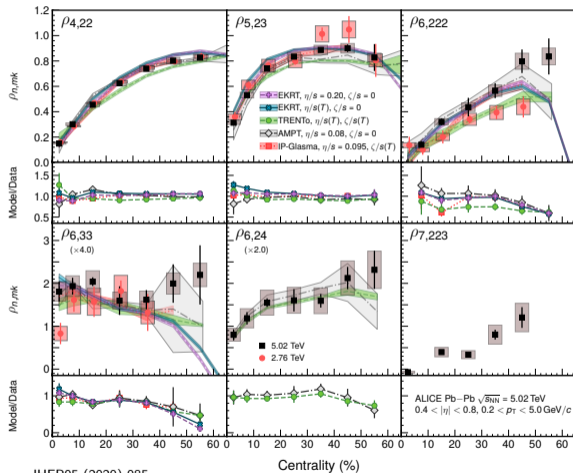
Our data can help to further constrain η/s and ζ/s in model calculations.

Related:

- E-b-E two or three harmonic flow magnitude correlations [A. Bilandzic, today]
- What can we learn from small systems? [Y. Zhou, today]

Thank you!

Symmetry-plane correlations, hydrodynamic predictions



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- Clear centrality dependence for all harmonics
- Good agreement between the data and the models