

Photoproduction of ρ mesons on nuclei in ultraperipheral nuclear collisions at the LHC



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Outline:

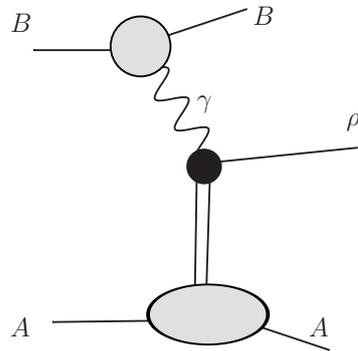
- Heavy-ion ultraperipheral collisions (UPCs) at the LHC
- Nuclear shadowing in coherent photoproduction of ρ mesons on nuclei and comparison to ALICE Pb-Pb data
- Incoherent photoproduction of ρ mesons in Pb-Pb UPCs

Based on [Frankfurt, Guzey, Strikman, Zhalov, Phys. Lett. B 752 \(2016\) 51](#);
[Guzey, Kryshen, Zhalov, Phys. Rev C 102 \(2020\) 1, 015208](#)

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Ultrapерipheral collisions (UPCs)

- In ion-ion collisions, nuclei can interact at large impact parameters $b \gg R_A + R_B$ → **ultrapерipheral collisions** (UPCs) → strong interaction suppressed → in equivalent photon approximation, interaction via quasi-real photons, Budnev, Ginzburg, Meledin, Serbo, Phys. Rept. 15 (1975) 181 → photon-proton and **photon-nucleus** scattering at unprecedentedly high energies, Baltz et al, Phys. Rept. 480 (2008) 1



- UPCs correspond to empty detector with only two pion (lepton) tracks
- Nuclear coherence by veto on neutron production by Zero Degree Calorimeters and selection of small p_T

- Coherent and incoherent (target nucleus breaks up) photoproduction of vector mesons in UPCs:

$$\frac{d\sigma_{AA \rightarrow \rho AA'}}{dy} = N_{\gamma/A}(y) \sigma_{\gamma A \rightarrow \rho A'}(y) + N_{\gamma/A}(-y) \sigma_{\gamma A \rightarrow \rho A'}(-y)$$

Photon flux from QED+ suppression of strong interactions at small b:

- high intensity $\sim Z^2$
- high photon energy $\sim \gamma_L$

Photoproduction cross section

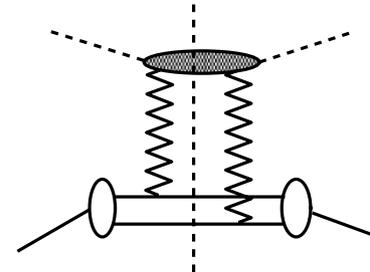
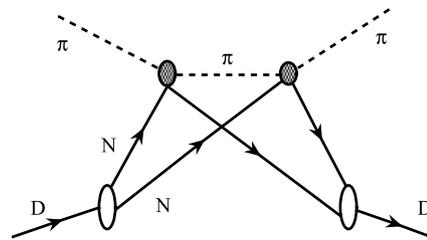
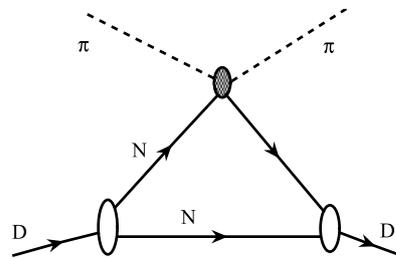
$y =$ rapidity of ρ

Nuclear shadowing

- **Nuclear shadowing (NS)** = suppression of cross section on a nucleus compared to sum of cross sections on individual nucleons: $\sigma_A < A \sigma_N$.
- Observed for various beams (p , π , γ , γ^* , ν) of large energies (> 1 GeV).
- Explained by simultaneous interaction of projectile with target nucleons \rightarrow destructive interference among amplitudes for interaction with 1, 2, ... nucleons \rightarrow nucleons in rear of the nucleus “see” smaller (shadowed) flux: $\sigma_A \sim A^{2/3}$.

elastic intermediate state, Glauber (1955)

inelastic intermediate state, Gribov (1969)



- **NS in photoproduction of light vector mesons ρ , ω , ϕ :**
 - dynamics of soft γp and γA interaction at high energies
 - test of VMD model and role of inelastic (Gribov) shadowing
- **NS in photoproduction of heavy vector mesons J/ψ , $\psi(2S)$, Y :**
 - mechanism of nuclear shadowing: leading twist vs. HT vs. saturation
 - new constraints on nuclear gluon distribution $g_A(x, \mu^2)$ at small x

Coherent photoproduction of ρ on nuclei

- Measured with fixed targets (SLAC, $W < 6$ GeV), in Au-Au UPCs@RHIC ($W < 12$ GeV), and Pb-Pb UPCs@LHC at 2.76 TeV ($W=46$ GeV) and 5.02 TeV ($W=62$ GeV)
- The method based on the vector meson dominance (VMD) model for $\gamma \rightarrow \rho$ transition and Glauber model for shadowing in ρA scattering fails to describe the data.

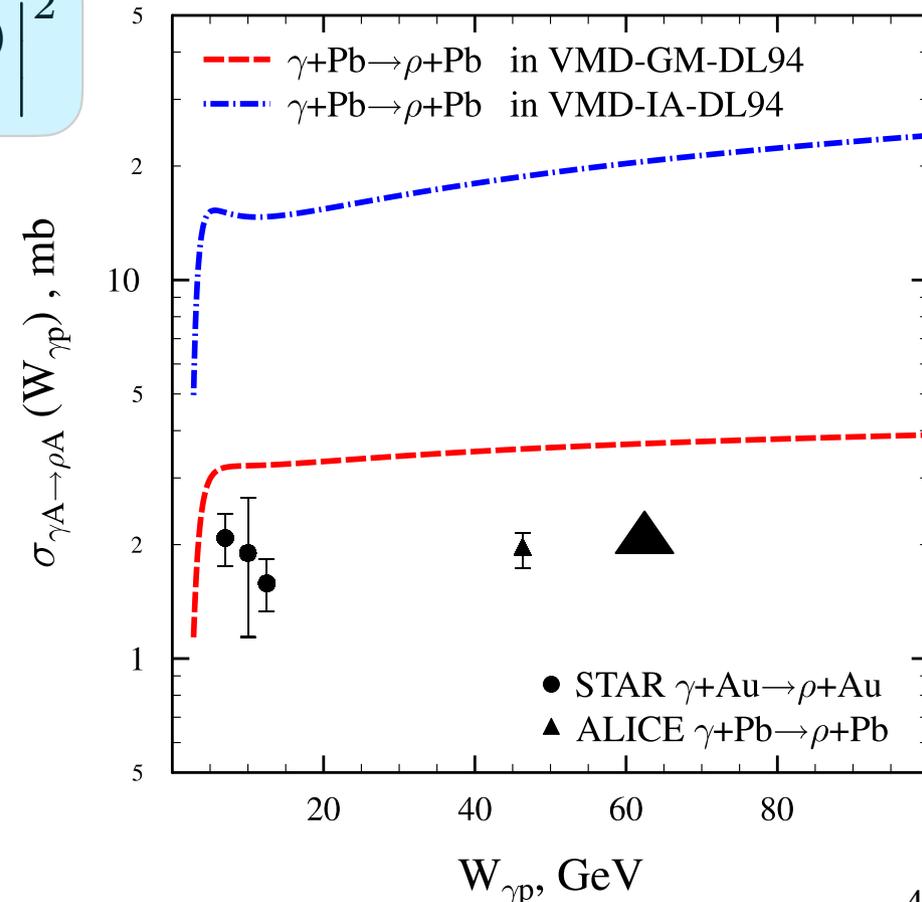
$$\sigma_{\gamma A \rightarrow \rho A}^{\text{VMD}} = \left(\frac{e}{f_\rho} \right)^2 \int d^2b \left| 1 - e^{-\frac{1}{2} \sigma_{\rho N} T_A(b)} \right|^2$$

$\sigma_{\rho N}$ = ρ -nucleon total cross section from VMD and constituent quark counting

nuclear optical density:

$$T_A(b) = \int dz \rho_A(b, z)$$

- The disagreement is by factor ~ 1.5 for large- W RHIC (STAR), Adler, et al, Phys. Rev. Lett. 89 (2002) 272302; Abelev et al., Phys. Rev. C 77 (2008) 034910; Agakishiev, et al., Phys. Rev. C 85 (2012) 014910 and ALICE data, Adam et al (ALICE), JHEP 1509 (2015) 095; Acharya et al, JHEP 06 (2020) 035.
- \rightarrow modified model is required.



Modified vector meson dominance (mVMD) model

- At large beam energies E_γ , the photon can be viewed as superposition of long-lived ($l_c \sim E_\gamma$) fluctuations interacting with hadrons with different cross sections, [Good, Walker, Phys. Rev. 120 \(1960\) 1857](#)

- It can be realized by introducing the probability distribution $P(\sigma)$, [Blattel et al, Phys. Rev. D 47 \(1993\) 2761](#)

$$\int d\sigma P(\sigma) = 1,$$

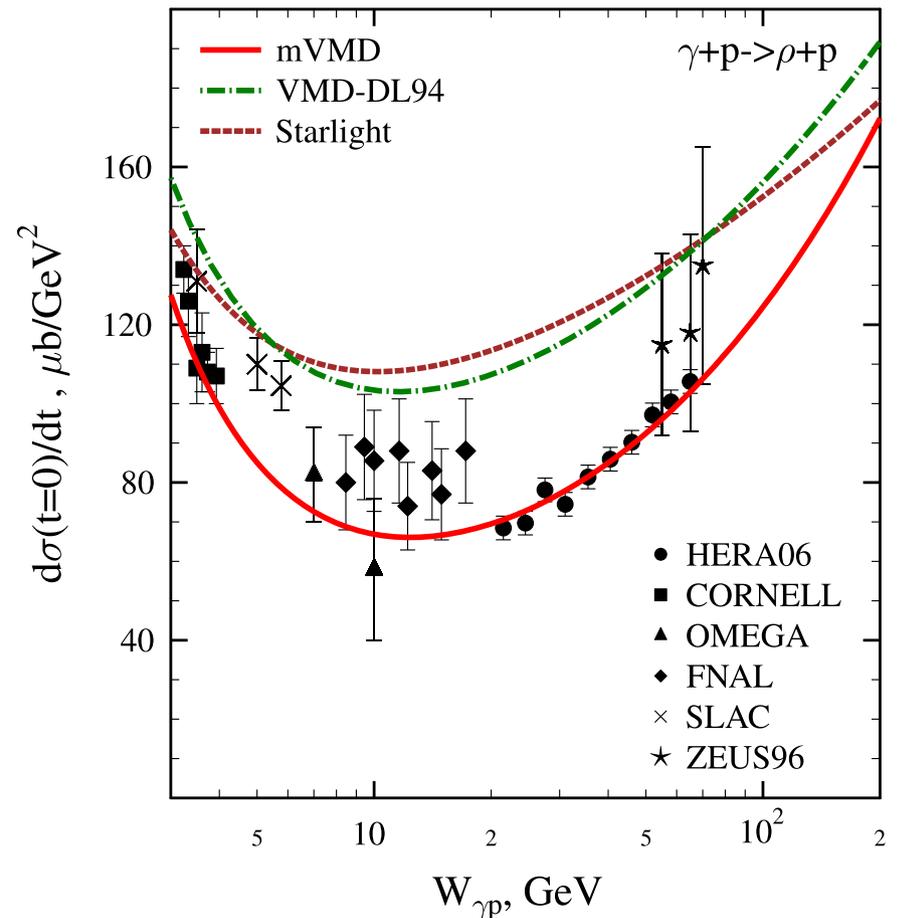
$$\int d\sigma P(\sigma)\sigma = \langle\sigma\rangle, \rightarrow \text{from } d\sigma(\gamma p \rightarrow \rho p)/dt$$

$$\int d\sigma P(\sigma)\sigma^2 = \langle\sigma\rangle^2(1 + \omega_\sigma) \rightarrow \text{from measured } \gamma \text{ diffract. dissociation into large masses, } \text{Chapin 1985}$$

- Shape like for pion + small- σ enhancement to take into account smaller size of ρ in $\gamma p \rightarrow \rho p$ than in $\sigma_{\pi N} \rightarrow$

$$P(\sigma) = C \frac{1}{1 + (\sigma/\sigma_0)^2} e^{-(\sigma/\sigma_0 - 1)^2/\Omega^2}$$

[Frankfurt, Guzey, Strikman, Zhilov, PLB 752 \(2016\) 51](#)

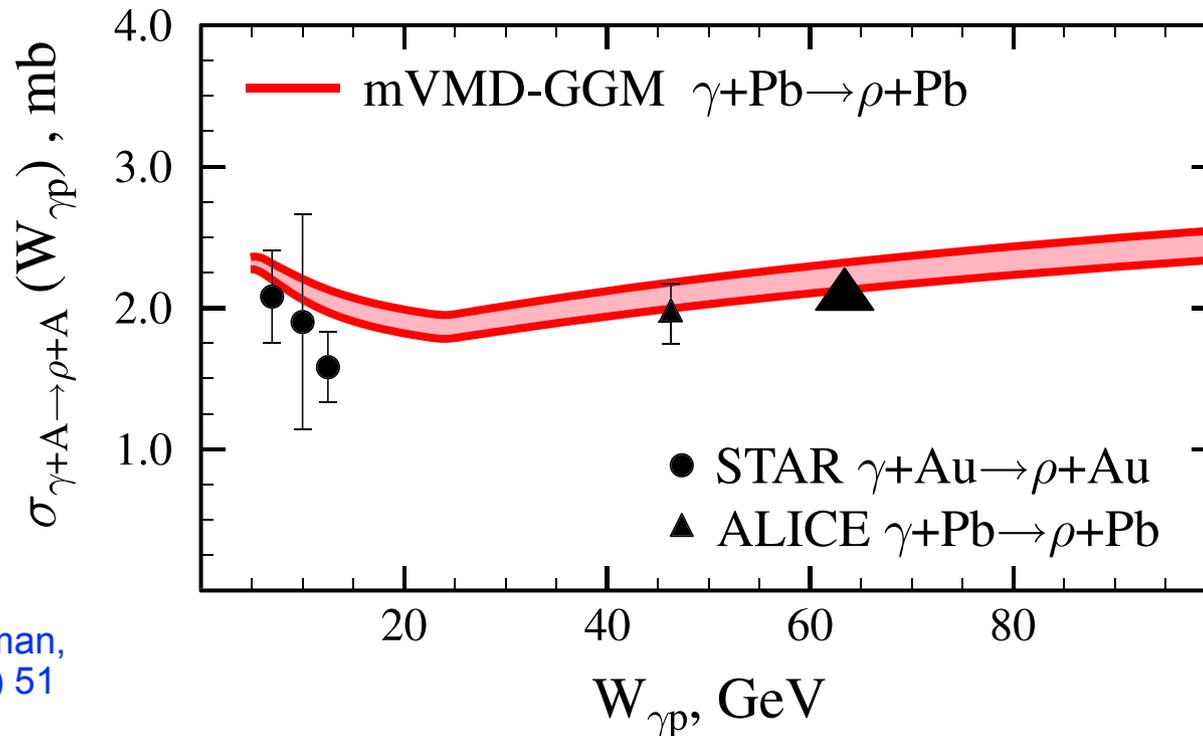


Photoproduction of ρ on Pb in modified VMD with Gribov-Glauber model of shadowing

- With cross section fluctuations:

$$\sigma_{\gamma A \rightarrow \rho A}^{\text{mVMD-GGM}} = \left(\frac{e}{f_\rho} \right)^2 \int d^2\vec{b} \left| \int d\sigma P(\sigma) \left(1 - e^{-\frac{\sigma}{2} T_A(b)} \right) \right|^2$$

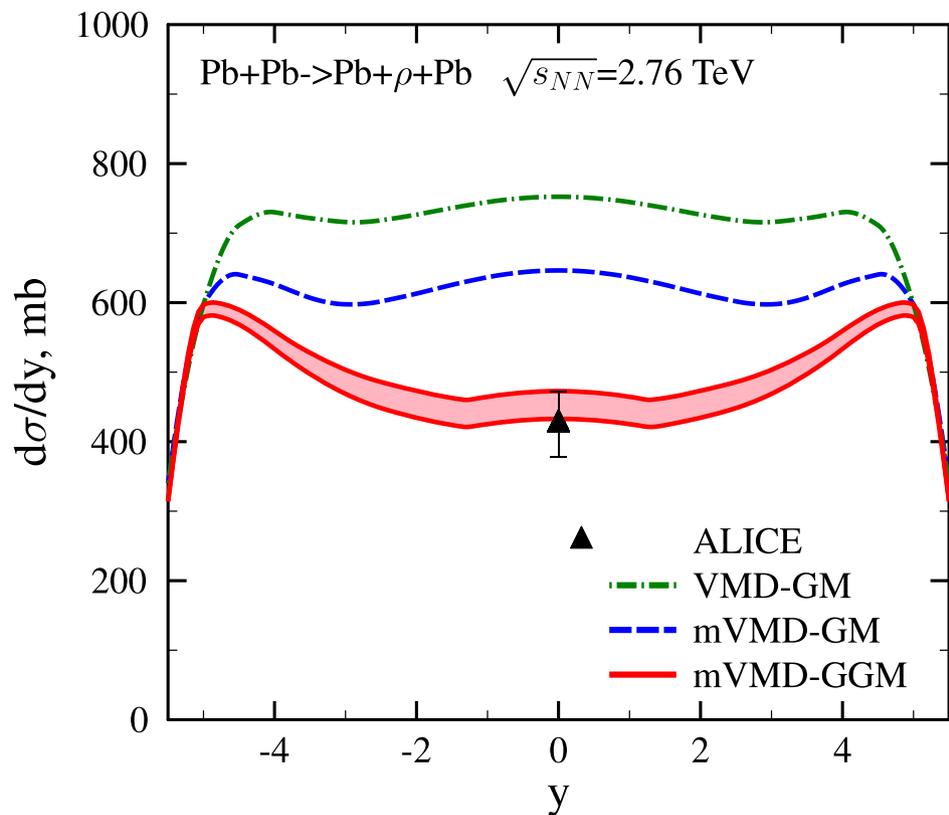
- Two additional effects compared to standard methods: correct description of the $\gamma p \rightarrow \rho p$ cross section and inclusion of inelastic Gribov shadowing in $\sigma_{\gamma A \rightarrow \rho A}$
- \rightarrow good description of normalization and W -dependence of $\sigma_{\gamma A \rightarrow \rho A}$



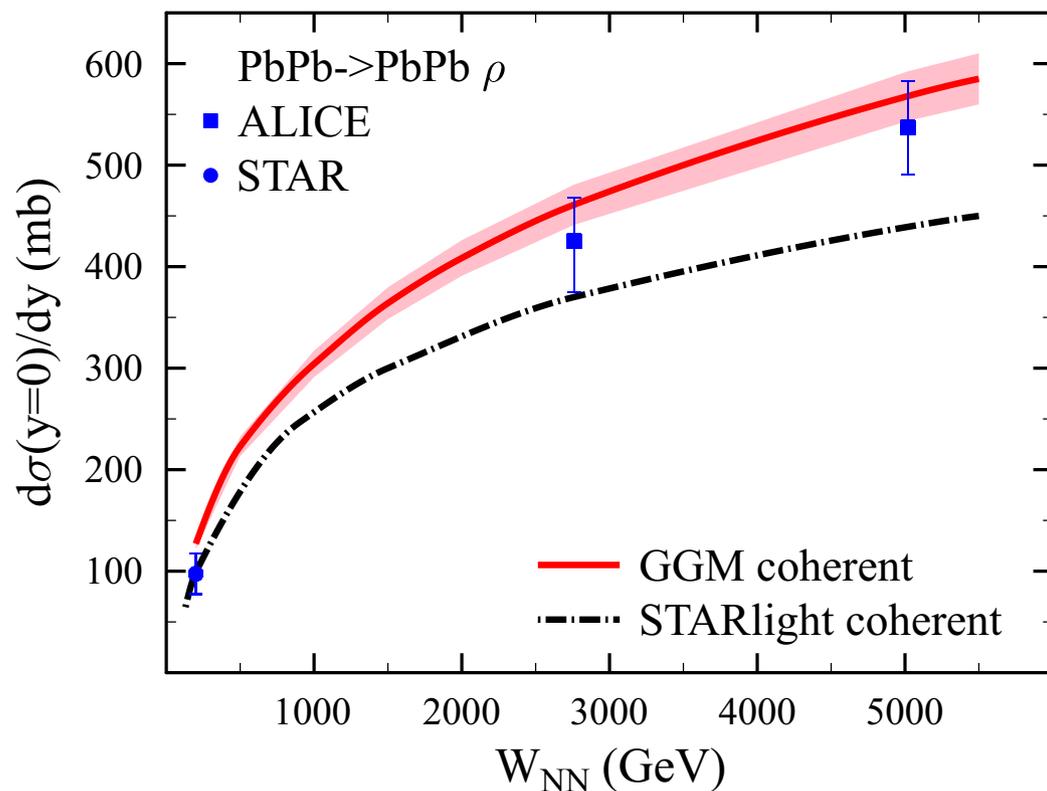
Photoproduction of ρ in Pb-Pb UPCs at LHC

- This translates into good description of RHIC and LHC (Run 1 and 2) data on cross section of coherent photoproduction of ρ mesons $d\sigma(AA \rightarrow \rho AA)/dy$ at $y=0$
- Left panel: Rapidity y dependence within Glauber (GM) and Gribov-Glauber (GGM) approaches to shadowing
- Right panel: Invariant energy $W_{NN} = \sqrt{s_{NN}}$ dependence

Frankfurt, Guzey, Strikman, Zhalov, PLB 752 (2016) 51



Guzey, Kryshen, Zhalov, PRC 102 (2020) 1, 015208



Incoherent photoproduction of ρ on nuclei

- The discussed framework can be generalized to **incoherent** ρ meson photoproduction on nuclei using **completeness of intermediate nuclear states**:

$$\begin{aligned}\sigma_{\gamma A \rightarrow \rho A'}^{\text{GG}} &= \left(\frac{e}{f_\rho}\right)^2 \int d^2\vec{b} \int d\sigma P(\sigma) \int d\sigma' P(\sigma') \frac{\sigma\sigma'}{16\pi B} T_A(b) \exp\left[-\frac{\sigma+\sigma'}{2}T_A(b) + \frac{\sigma\sigma'}{16\pi B}T_A(b)\right] \\ &= \left(\frac{e}{f_\rho}\right)^2 \int d^2\vec{b} T_A(b) \left(\int d\sigma P(\sigma) \frac{\sigma}{\sqrt{16\pi B}} \exp\left[-\frac{\sigma^{\text{in}}}{2}T_A(b)\right]\right)^2.\end{aligned}$$

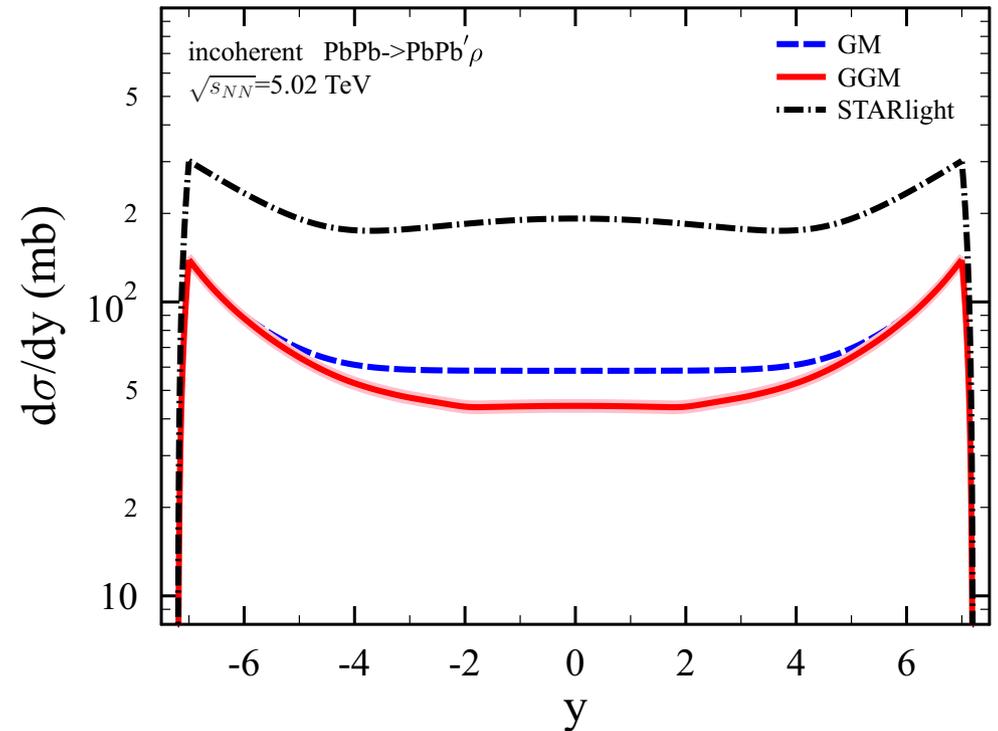
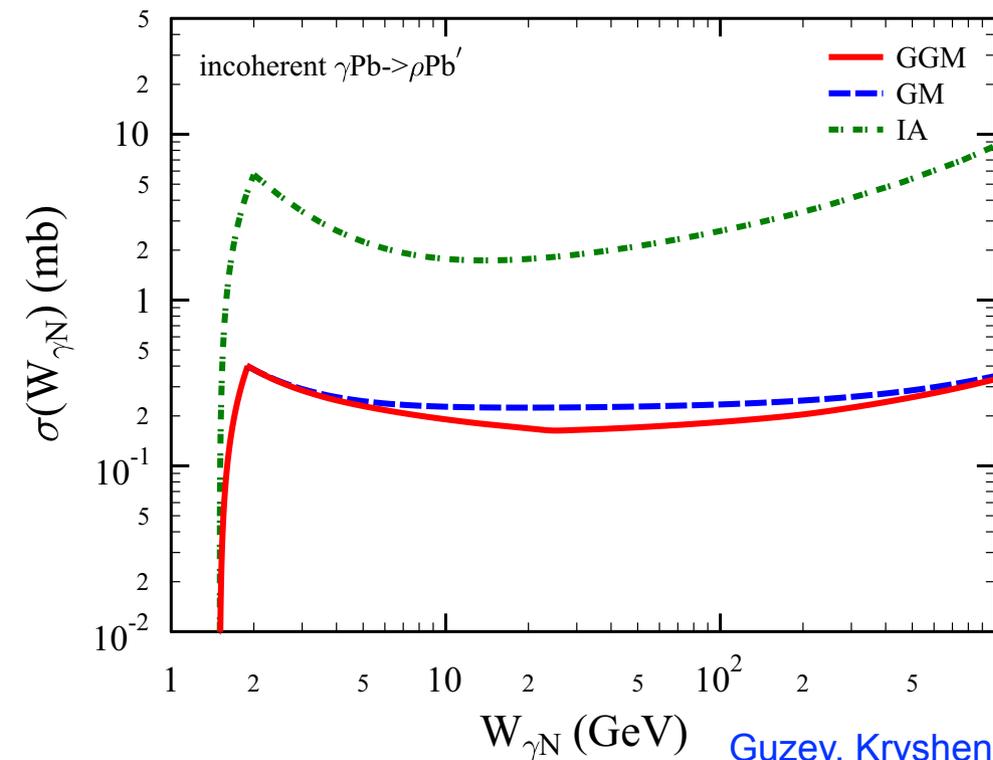
- In the absence of cross section fluctuations $P(\sigma)=\delta(\sigma-\sigma_{\rho N}) \rightarrow$ reproduced standard Glauber model formula:

$$\sigma_{\gamma A \rightarrow \rho A'}^{\text{Glauber}} = \left(\frac{e}{f_\rho}\right)^2 \frac{\sigma_{\rho N}^2}{16\pi B} \int d^2\vec{b} T_A(b) e^{-\sigma_{\rho N}^{\text{in}} T_A(b)} = \sigma_{\gamma p \rightarrow \rho p} \int d^2\vec{b} T_A(b) e^{-\sigma_{\rho N}^{\text{in}} T_A(b)}$$

- Like in the coherent case, **cross section fluctuations increase nuclear shadowing** \rightarrow decrease nuclear cross section.

Predictions for incoherent ρ photoproduction in Pb-Pb UPCs at LHC at 5.02 TeV

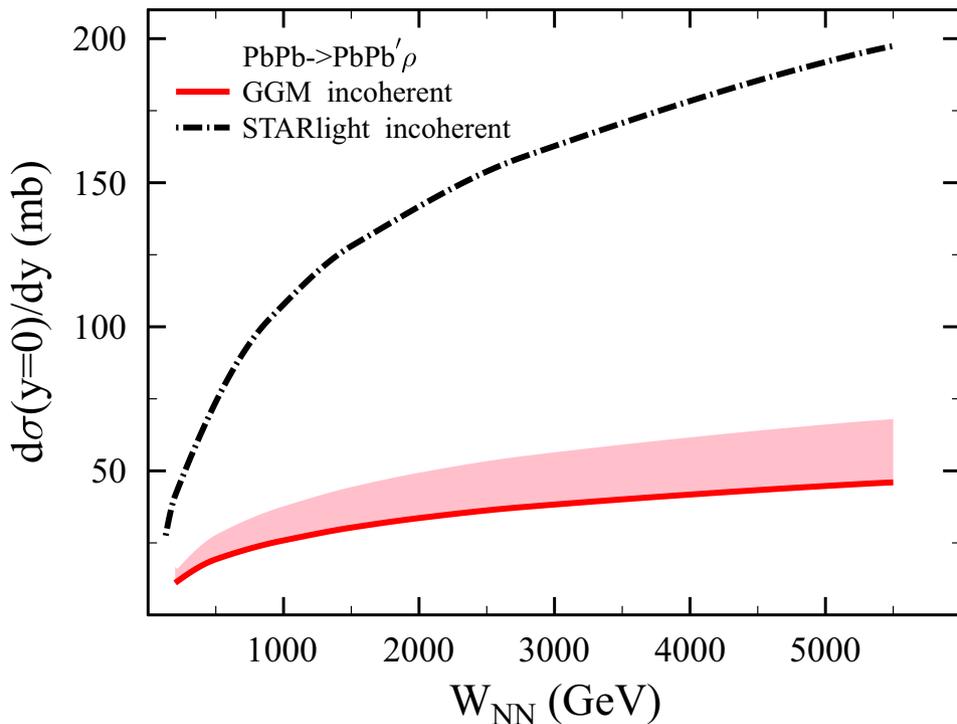
- Left: W dependence of $\sigma(\gamma A \rightarrow \rho A')$ using 3 different approaches to nuclear shadowing: no shadowing (IA), Glauber (GM) and Gribov-Glauber models (GGM)
- Right panel: Rapidity y dependence of UPC cross section $d\sigma(AA \rightarrow \rho AA')/dy$
- If no special data selection is made, these predictions can be increased by 50% by the contribution of nucleon dissociation $\gamma N \rightarrow \rho Y$ (next slide)



Comparison to STARlight MC

- Our predictions significantly differ from those of STARlight Monte Carlo both in coherent and incoherent cases.
- STARlight is the main MC generator used in data analysis of UPCs at RHIC and LHC, [Klein et al, Comput. Phys. Commun. 212 \(2017\) 258](#)
- STARlight incoherent cross section:

$$\sigma_{\gamma A \rightarrow \rho A'}^{\text{STARlight}} = \frac{\sigma_{\gamma p \rightarrow \rho p}}{\sigma_{\rho N}^{\text{in}}} \sigma_{\rho A}^{\text{in}} = \left(\frac{e}{f_\rho} \right)^2 \frac{\sigma_{\rho N}^{\text{el}}}{\sigma_{\rho N}^{\text{in}}} \int d^2 \vec{b} (1 - e^{-\sigma_{\rho N} T_A(b)})$$



	GM	GGM	STARlight
Incoherent, mb	58	44	192
Coherent, mb	840	570	440

Summary

- Coherent and incoherent photoproduction of ρ mesons in Pb-Pb UPCs at the LHC allows one to study nuclear shadowing and hadronic fluctuations of the photon at unprecedentedly high energies.
- Inelastic (Gribov) nuclear shadowing is essential for describing the normalization and energy dependence of the Run 1 and 2 ALICE data.
- Presented framework allowed us to make predictions for incoherent photoproduction of ρ mesons in Pb-Pb UPCs at the LHC, which are very different from frequently used STARlight MC.