SPIN OBSERVABLES IN *pN*- and *pd*- ELASTIC AND QUASI-ELASTIC SCATTERING AT ENERGIES OF SPD NICA

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PLAN:

pN- elastic scattering . Data and parametrizations of pp amplitudes <3 GeV, >3 GeV

Glauber theory of pd- elastic scattering at 0.1-3 GeV

Results of Glauber calculations of spin observables pd-elastic scattering at 3-50 GeV/c

Quaslielastic scattering pd-> {pp}(1S0)+n and complete polarization experiment

Spin-dependent pN-amplitudes in search of T-invariance violation in pd-scattering

Conclusion

NN-forces is a basis of nuclear physics .
It is important to study all their components, including spin-dependent terms at low and high energies **via the spin amplitudes of the NN elastic scattering**

PP elastic scattering

Helicity amplitudes of NN-scattering

For identical spin $\frac{1}{2}$ particles under Lorentz and P-,T- invariance:

spin non-flip $\phi_1(s,t) =< ++|M|++ >$

double spin-flip $\phi_2(s,t) =< ++|M|-- >$

spin non-flip $\phi_3(s,t) =< +-|M|+- >$

double spin-flip $\phi_4(s,t) =< +-|M|-+ >$

single spin-flip $\phi_5(s,t) =< ++|M|++- >$.

For non-identical (pn) nucleons one has 6 amplitudes,

T-reversal non-invarinace provides two additional amplitudes.

All spin-observables of NN elastic scattering are described in terms of ϕ_i

$$d\sigma/dt = N[|\phi_1|^2 + |\phi_2|^2 + |\phi_1|^3 + |\phi_4|^2 + 4|\phi_5|^2],$$

$$A_N \sim Im[(\phi_1 + \phi_2 + \phi_3 - \phi_4)\phi_5^*]$$

$$A_{NN} \sim 2|\phi_5|^2 + Re(\phi_1^*\phi_2 - \phi_3^*\phi_4)$$

$$\cdots$$

Number of linearly independent non-zero spin observables:
single-spin (asymmetries A_i, polarizations P_i) – 2
double-spin ($A_{ii}, \ldots$) – 12
triple-spin – 9
four- spin – 2

Complete polarization experiment
for pp-elastic requires 9 independent observables.
PWA GWU is performed for pp-elastic up to 3.8 GeV/c (SAID webpage:`http://gwdac.phys.gwu.edu`.
R.A. Arndt, I.I. Strakovsky, B.L. Workman PRC 56, 3005 (1997);
PWA for pn- elastic – up 1.2 GeV/c

Concerning SPD NICA, above 3 GeV/c $d\sigma/dt$ and mainly A_N (up to 50 GeV/c) and A_{NN}, C_{LL} (up to 6 GeV/c, 12 GeV/c) are measured. Data on double-spin observables D_{NN}, K_{NN} are rather poore in the region of forward angles.

Parametrizations (fit) of the pp- data:
Regge: W.P. Ford, J.W. Van Orden, Phy.Rev. **C87** (2013) 014004;
A. Sibirtsev et al. Eur.Phys.J. **A 45** (2010) 357;
Eikonal: S. Wakaizumi, M. Sawamoto,Prog. Theor. Phys. v.64 (1980) 1699

$$\phi_{ai}(s,t) = \pi\beta_{ai}(t)\frac{\xi_i(s,t)}{\Gamma(\alpha(t))}; i = \rho, \omega, a_2, f_2, P; a = 1-5;$$

$$\xi_i(t,s) = \frac{1 + S_i \exp[-i\pi\alpha(t)]}{\sin[\pi\alpha_i(t)]}\left[\frac{s}{s_0}\right]^{\alpha_i(t)},$$

$$\alpha_i(t) = \alpha_i^0 + \alpha_i^{'} t,$$

$$\beta_{1i}(t) = c_{1i}\exp(b_{1i}t),$$

$$\beta_{2i}(t) = c_{2i}\exp(b_{2i}t)\frac{-t}{4m_N^2},$$

$$\beta_{3i}(t) = c_{3i}\exp(b_{3i}t),$$

$$\beta_{4i}(t) = c_{4i}\exp(b_{4i}t)\frac{-t}{4m_N^2},$$

$$\beta_{5i}(t) = c_{5i}\exp(b_{5i}t)\left[\frac{-t}{4m_N^2}\right]^{1/2}.$$

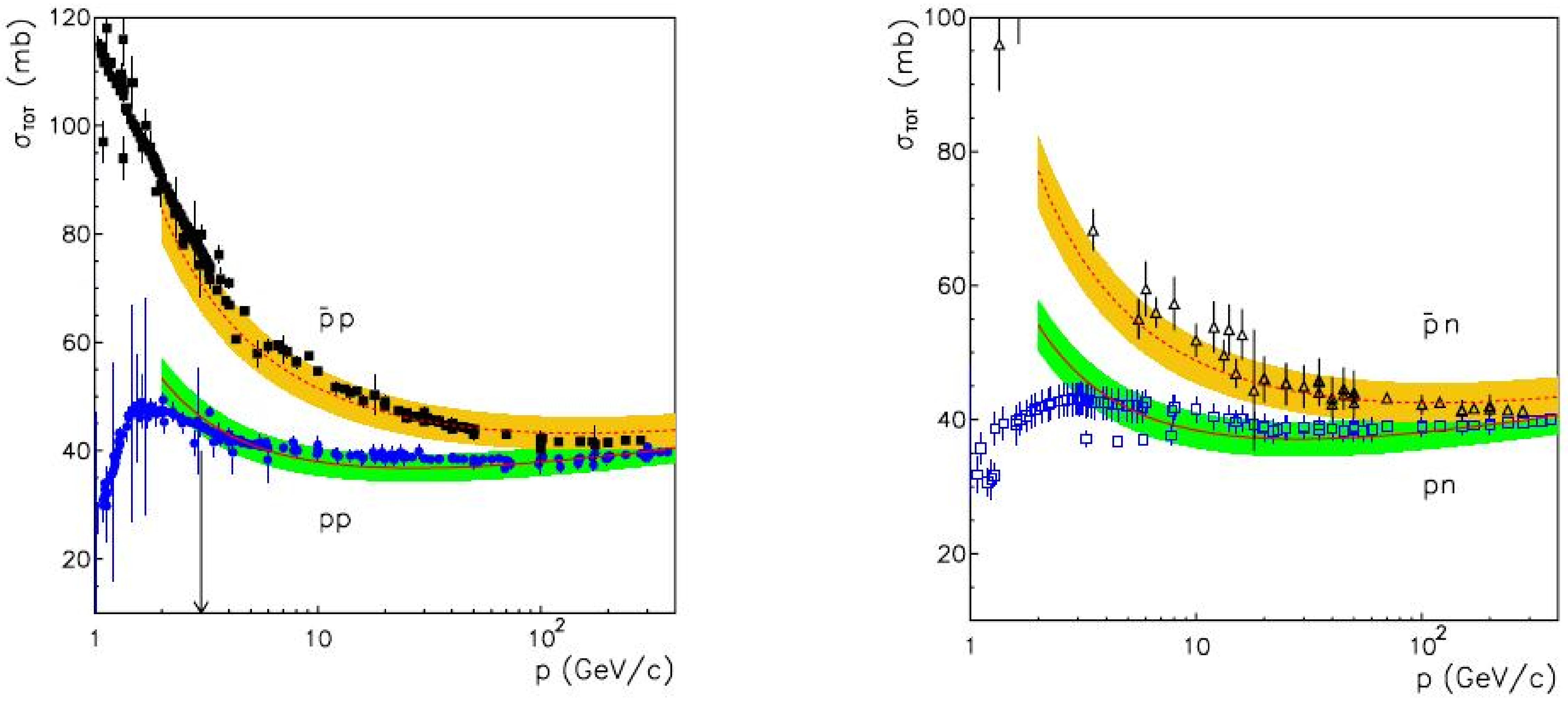
A. Sibirtsev et al, EPJ A 45 (2010) 357

Helicity amplitudes in Regge formalism

A systematic analisys of pp elastic scattering from COSY-EDDA, SATURNE, GZS ANL
A. Sibirtsev, J. Haidenbauer, H.-W. Hammer et al. EPJA 45 (2010) 357
ω, ρ, f_2, a_2 Regge exchanges and the Pomeron for P_L from 3 GeV/c up to 50 GeV/c.
Isospin structure and G-parity relations allow to obtain the $\bar{p}p$-, pn- and $\bar{p}n$ elastic amplitudes from the pp amplitudes (J.R. Pelaez, 2006):

$$\begin{aligned}
\phi(pp) &= -\phi_\omega - \phi_\rho + \phi_{f_2} + \phi_{a_2} + \phi_P \\
\phi(\bar{p}p) &= \phi_\omega + \phi_\rho + \phi_{f_2} + \phi_{a_2} + \phi_P \\
\phi(pn) &= -\phi_\omega + \phi_\rho + \phi_{f_2} - \phi_{a_2} + \phi_P \\
\phi(\bar{p}n) &= \phi_\omega - \phi_\rho + \phi_{f_2} - \phi_{a_2} + \phi_P
\end{aligned}$$

However, not all available data
($C_{NN}, C_{LL}, C_{SS}, C_{LS}, D_{NN}, D_{SS}, D_{LS}, K_{NN}, \Delta\sigma_T, H_{SNS}$... measured by ANL at 6 GeV/c, and some at 12 GeV/c) were included into the fit.



A. Sibirtsev et al, EPJA (2010): full line – pp, dashed – antip p.

PROBLEMS

A. Sibirtsev et al, EPJA (2010) 357

R- real-to-imaginary part ratio.
A problem with antip-p theory: no zero

Other problems, see O.V. Selyugin, PRD 91 (2015);
PEPA Letters 13 (2016) 116.

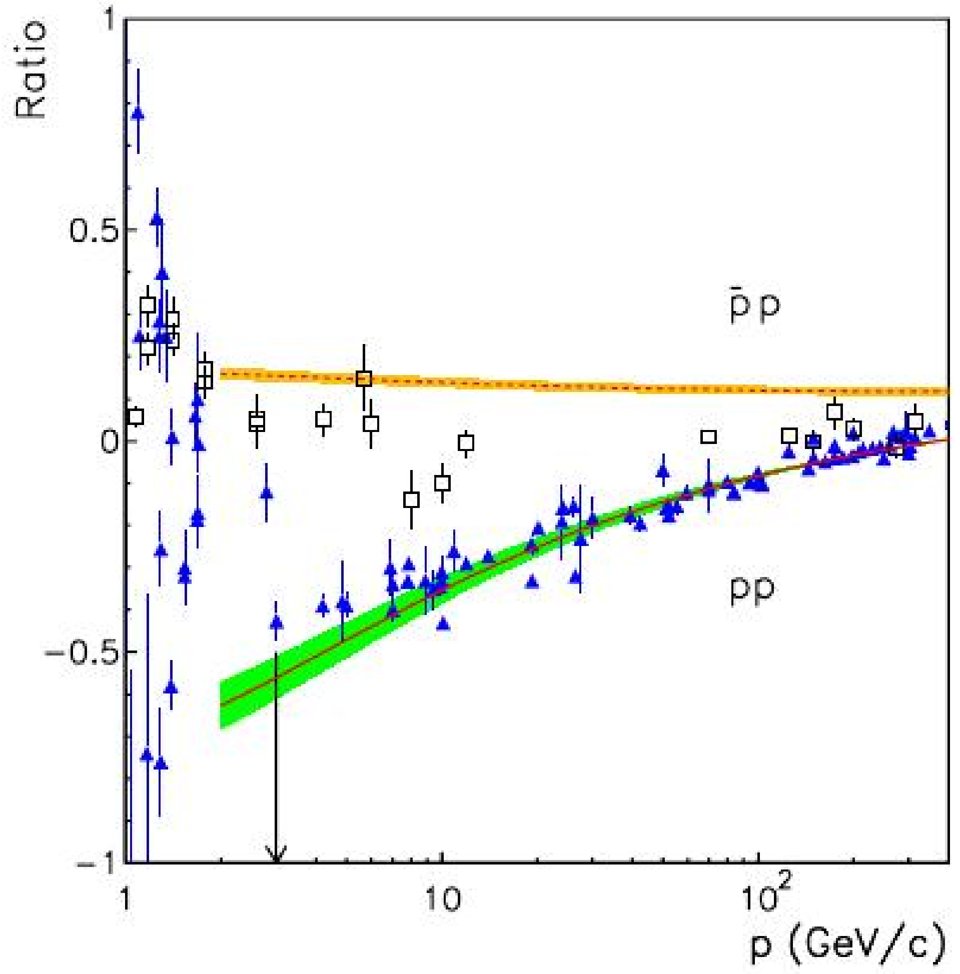
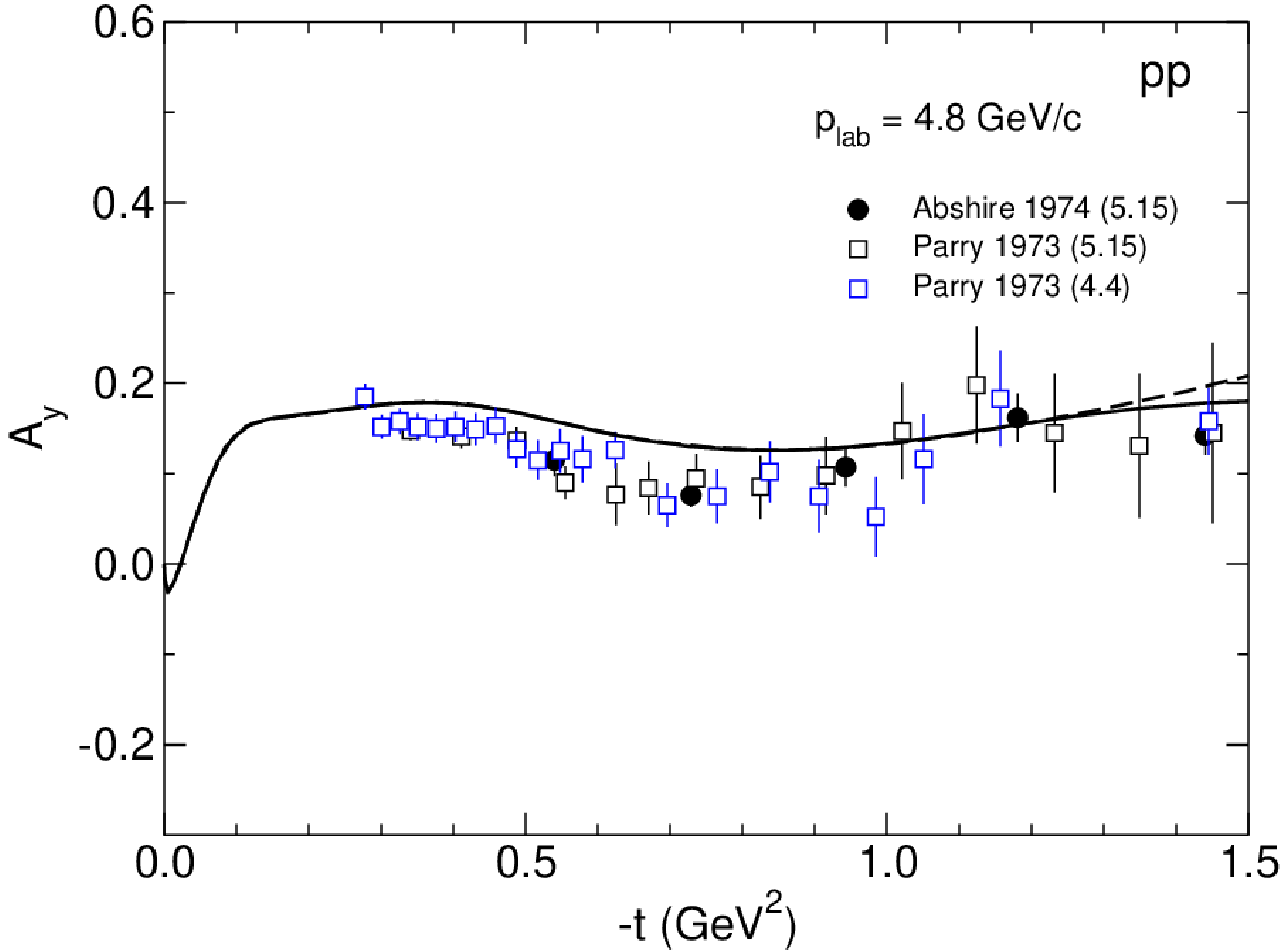
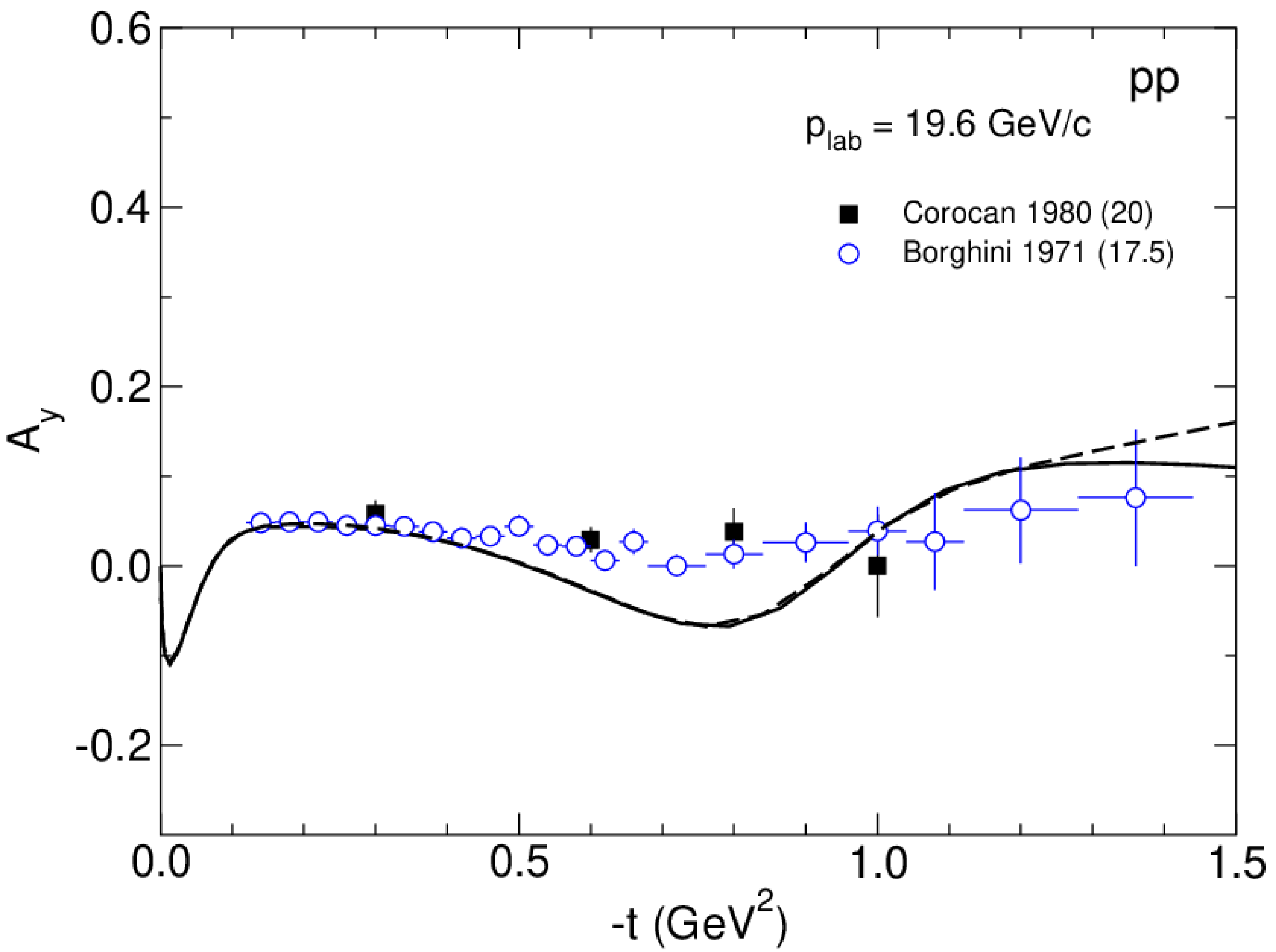


Fig. 16. Ratio of the real-to-imaginary parts of the forward amplitudes for *pp* (triangles, solid line) and $\bar{p}p$ (squares, dashed line), respectively. The data are taken from the PDG [32].



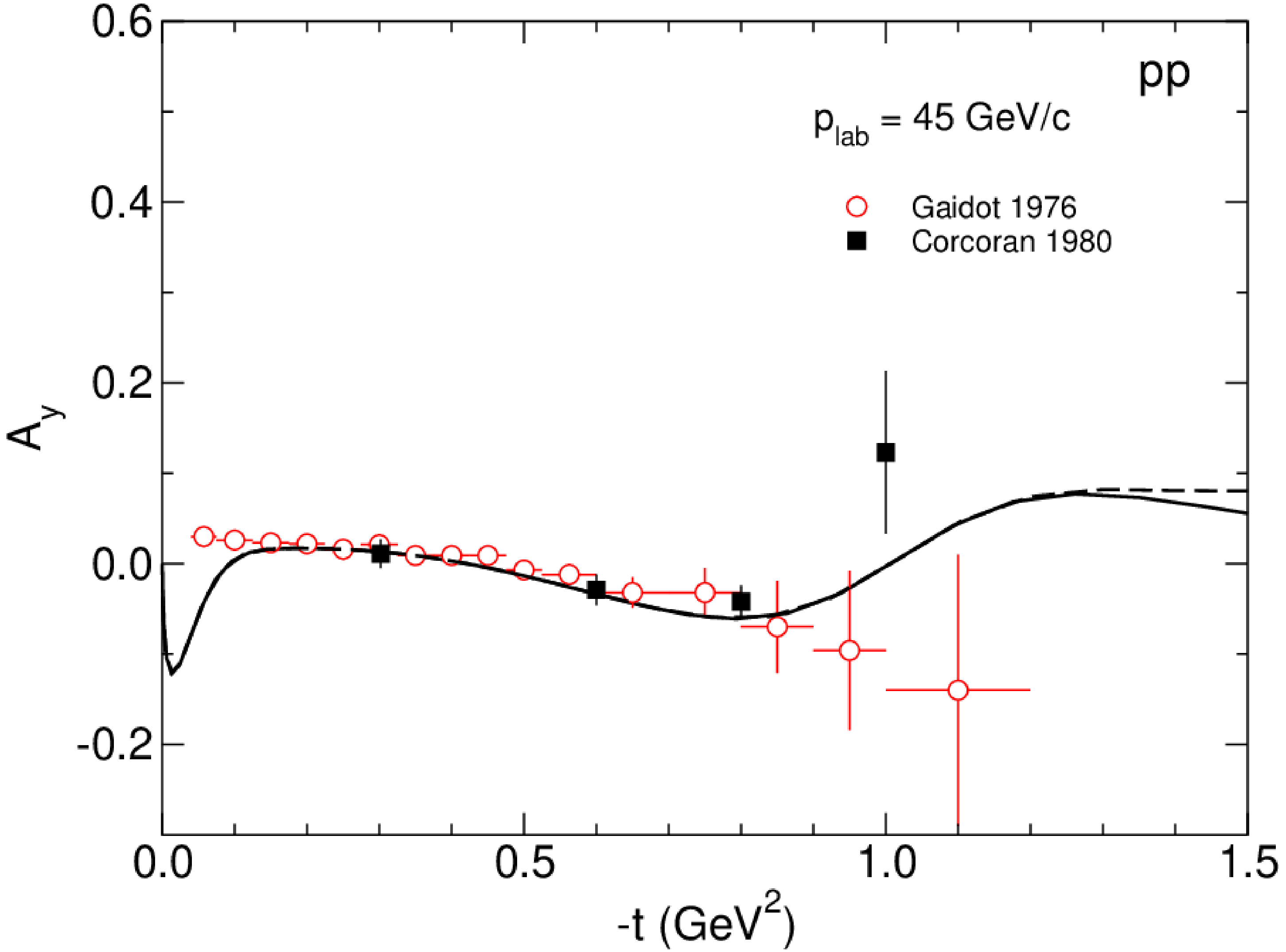
pp- helicity amplitudes
by A.Sibirtsev:
Regge parametrization
Gaussian (dashed)

pp- helicity amplitudes
A.Sibirtsev:
Regge parametrization;
Gaussian (dashed)



https://www.hepdata.net/record/ins113043

A.Gaidat et al. Phys.Lett. B61 (1976) 103



pp- helicity amplitudes
A.Sibirtsev: Regge parametrization (full); Gaussians (dashed)

pd elastic scattering

Phenomenology of the $pd \to pd$ transition

$$\frac{1}{2} + 1 \to \frac{1}{2} + 1$$

$(2+1)^2(2\frac{1}{2}+1)^2 = 36$ transition amplitudes

P-parity $\Longrightarrow$ 18 independent amplitudes

T-invariance $\Longrightarrow$ 12 independent amplitudes

At $\theta_{cm} = 0 \Longrightarrow$ 4 (for T-inv. P-inv.) + 1 (T- viol. P-inv.)

Elastic $pd \to pd$ transitions

$$\hat{M}(\mathbf{q},\mathbf{s}) =$$
$$\exp\left(\frac{1}{2}i\mathbf{q}\cdot\mathbf{s}\right)M_{pp}(\mathbf{q}) + \exp\left(-\frac{1}{2}i\mathbf{q}\cdot\mathbf{s}\right)M_{pn}(\mathbf{q}) +$$
$$+\frac{i}{2\pi^{3/2}}\int \exp\left(i\mathbf{q}'\cdot\mathbf{s}\right)\Big[M_{pp}(\mathbf{q}_1)M_{pn}(\mathbf{q}_2) + p \leftrightarrow n\Big]d^2\mathbf{q}'.$$

On-shell elastic pN scattering amplitude (**T-even, P-even**)

$$\begin{aligned} M_{pN} &= A_N + (C_N\boldsymbol{\sigma}_1 + C'_N\boldsymbol{\sigma}_2)\cdot\hat{\mathbf{n}} + B_N(\boldsymbol{\sigma}_1\cdot\hat{\mathbf{k}})(\boldsymbol{\sigma}_2\cdot\hat{\mathbf{k}}) + \\ &+ (G_N - H_N)(\boldsymbol{\sigma}_1\cdot\hat{\mathbf{n}})(\boldsymbol{\sigma}_2\cdot\hat{\mathbf{n}}) + (G_N + H_N)(\boldsymbol{\sigma}_1\cdot\hat{\mathbf{q}})(\boldsymbol{\sigma}_2\cdot\hat{\mathbf{q}}) \end{aligned}$$

M. Platonova, V. Kukulin, PRC **81** (2010) 014004:

$\hat{\mathbf{q}} = (\mathbf{p} - \mathbf{p}')$, $\hat{\mathbf{k}} = (\mathbf{p} + \mathbf{p}')/$, $\hat{\mathbf{n}} = [\mathbf{k} \times \mathbf{q}]$ – unit vect. ($Z \uparrow\uparrow \hat{\mathbf{k}}$, $X \uparrow\uparrow \hat{\mathbf{q}}$ $Y \uparrow\uparrow \hat{\mathbf{n}}$)

$$
\begin{aligned}
M = {} & (A_1 + A_2 \boldsymbol{\sigma}\hat{\mathbf{n}}) + (A_3 + A_4 \boldsymbol{\sigma}\hat{\mathbf{n}})(\mathbf{S}\hat{\mathbf{q}})^2 + (A_5 + A_6 \boldsymbol{\sigma}\hat{\mathbf{n}})(\mathbf{S}\hat{\mathbf{n}})^2 + A_7(\boldsymbol{\sigma}\hat{\mathbf{k}})(\mathbf{S}\hat{\mathbf{k}}) + \\
& A_8(\boldsymbol{\sigma}\hat{\mathbf{q}})\left[(\mathbf{S}\hat{\mathbf{q}})(\mathbf{S}\hat{\mathbf{n}}) + (\mathbf{S}\hat{\mathbf{n}})(\mathbf{S}\hat{\mathbf{q}})\right] + (A_9 + A_{10}\boldsymbol{\sigma}\hat{\mathbf{n}})(\mathbf{S}\hat{\mathbf{n}}) + A_{11}(\boldsymbol{\sigma}\hat{\mathbf{q}})(\mathbf{S}\hat{\mathbf{q}}) + \\
& A_{12}(\boldsymbol{\sigma}\hat{\mathbf{k}})\left[(\mathbf{S}\hat{\mathbf{k}})(\mathbf{S}\hat{\mathbf{n}}) + (\mathbf{S}\hat{\mathbf{n}})(\mathbf{S}\hat{\mathbf{k}})\right]
\end{aligned}
$$

$$- -$$

$$
\begin{aligned}
& +(T_{13} + T_{14}\boldsymbol{\sigma}\hat{\mathbf{n}})\left[(\mathbf{S}\hat{\mathbf{k}})(\mathbf{S}\hat{\mathbf{q}}) + (\mathbf{S}\hat{\mathbf{q}})(\mathbf{S}\hat{\mathbf{k}})\right] + T_{15}(\boldsymbol{\sigma}\hat{\mathbf{q}})(\mathbf{S}\hat{\mathbf{k}}) + T_{16}(\boldsymbol{\sigma}\hat{\mathbf{k}})(\mathbf{S}\hat{\mathbf{q}}) + \\
& T_{17}(\boldsymbol{\sigma}\hat{\mathbf{k}})\left[(\mathbf{S}\hat{\mathbf{q}})(\mathbf{S}\hat{\mathbf{n}}) + (\mathbf{S}\hat{\mathbf{n}})(\mathbf{S}\hat{\mathbf{q}})\right] + T_{18}(\boldsymbol{\sigma}\hat{\mathbf{q}})\left[(\mathbf{S}\hat{\mathbf{k}})(\mathbf{S}\hat{\mathbf{n}}) + (\mathbf{S}\hat{\mathbf{n}})(\mathbf{S}\hat{\mathbf{k}})\right]
\end{aligned}
$$

T-even P-even: $A_1 \div A_{12}$
(see M. Platonova, V.I. Kukulin,PRC **81** (2010) 014004)

$\mathbf{T_{13} \div T_{18}}$: TVPC

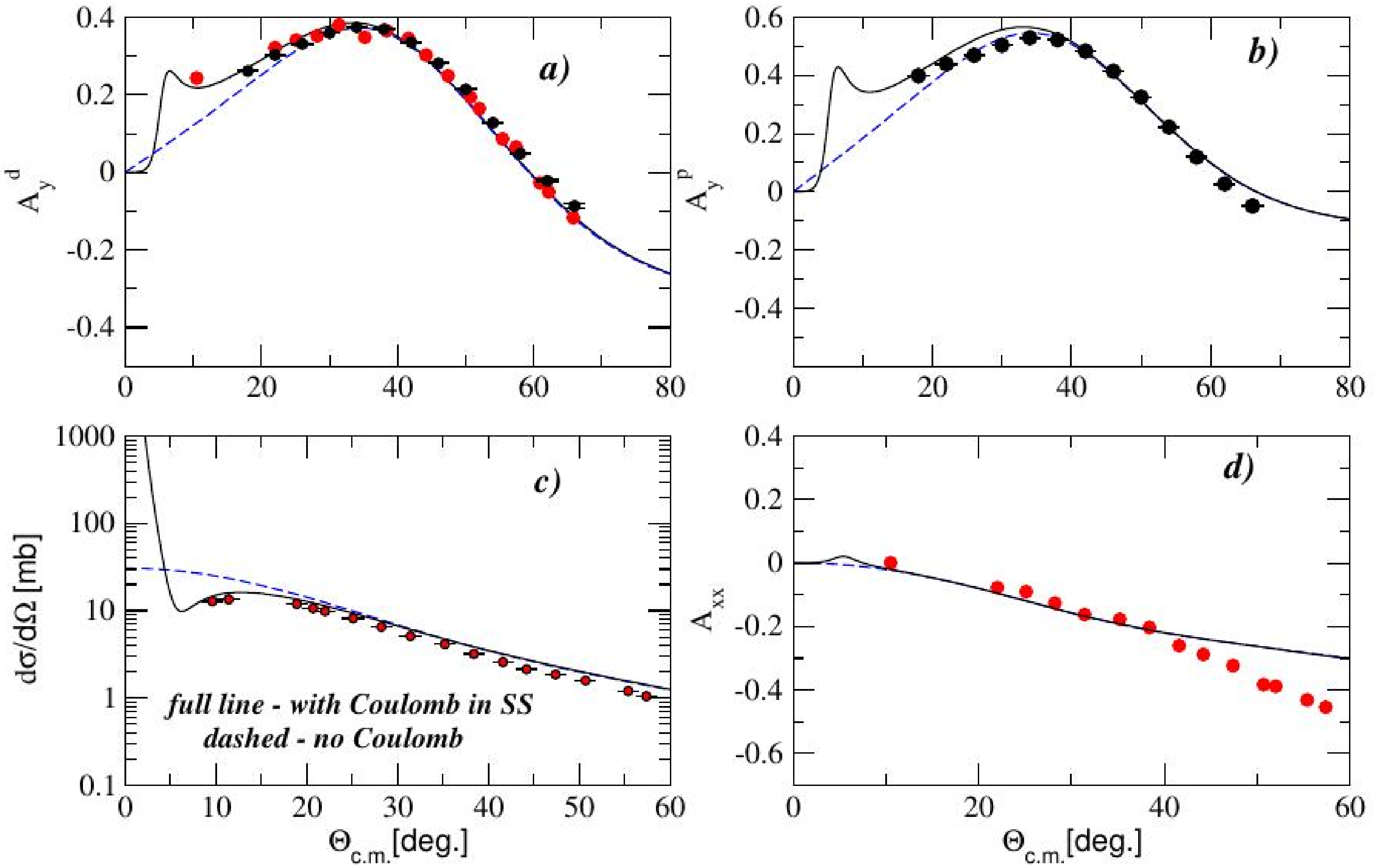
The polarized elastic differential pd cross section

$$\left(\frac{d\sigma}{d\Omega}\right)_{pol} = \left(\frac{d\sigma}{d\Omega}\right)_0 \left[1 + \frac{3}{2} p_j^p p_i^d C_{j,i} + \frac{1}{3} P_{ij}^d A_{ij} + \dots\right]. \tag{2}$$

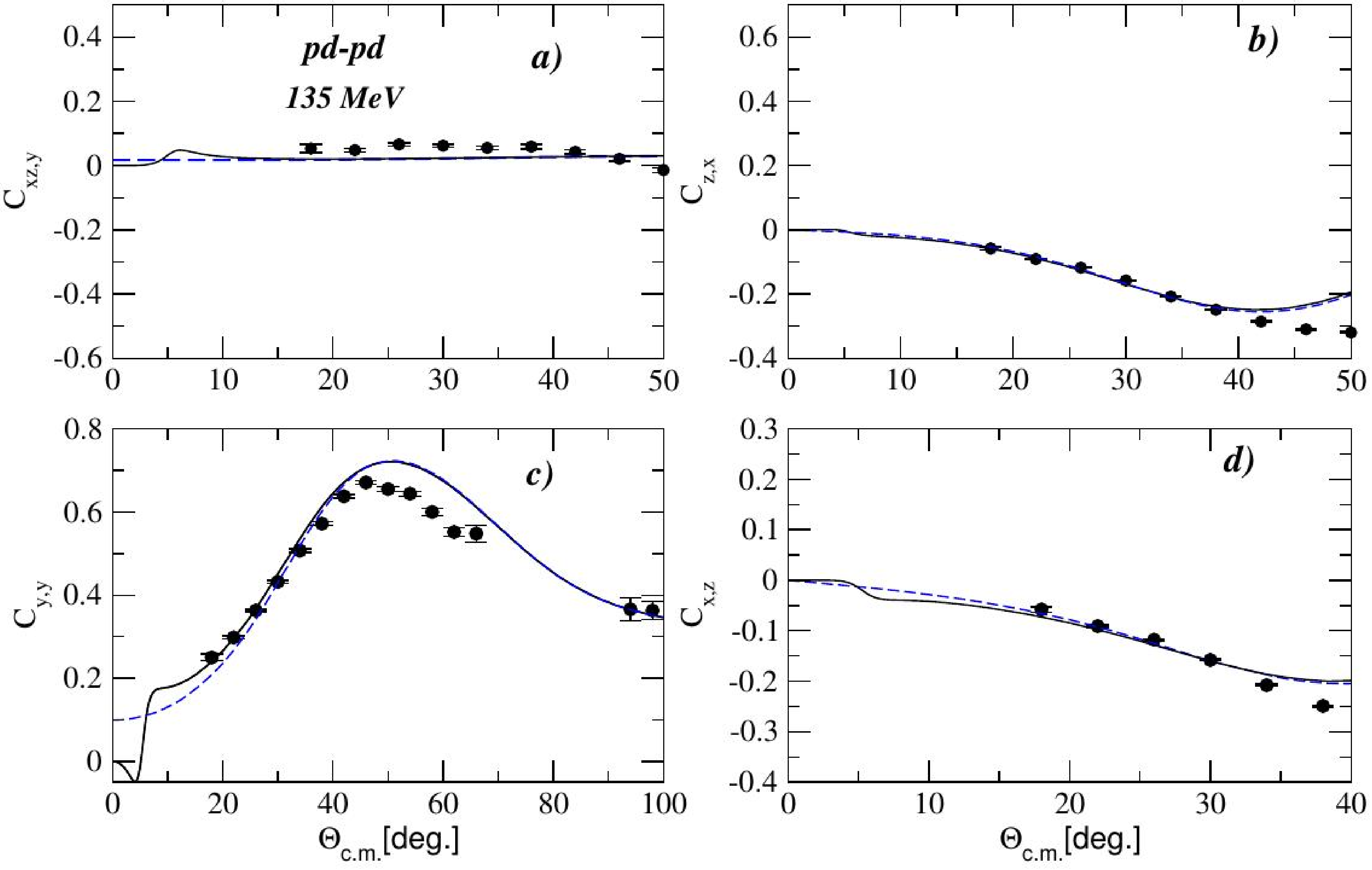
$$C_{y,y} = Tr M S_y \sigma_y M^+ / Tr M M^+, \quad \dots \tag{3}$$

Test calculations: pd elastic scattering at 135 MeV

A.A. Temerbayev, Yu.N.Uzikov, Yad. Fiz. **78** (2015) 38

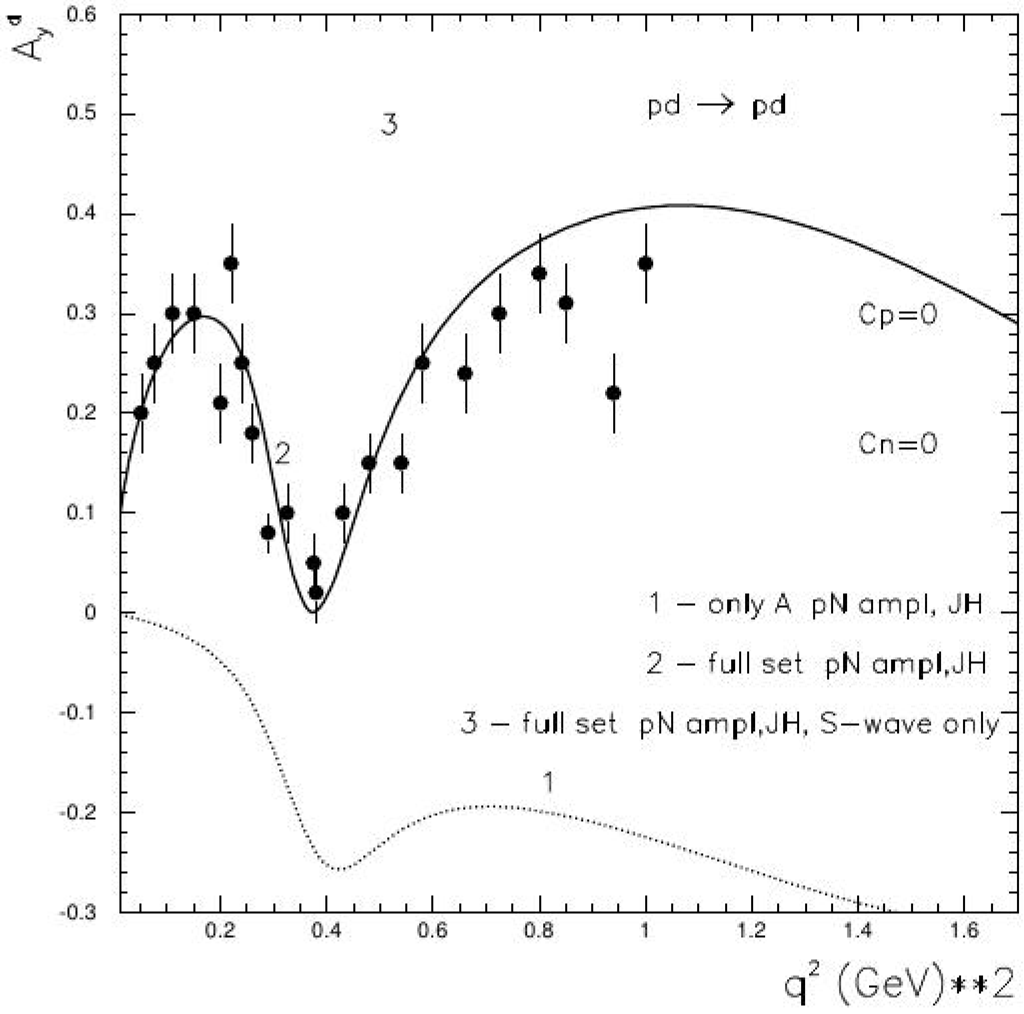


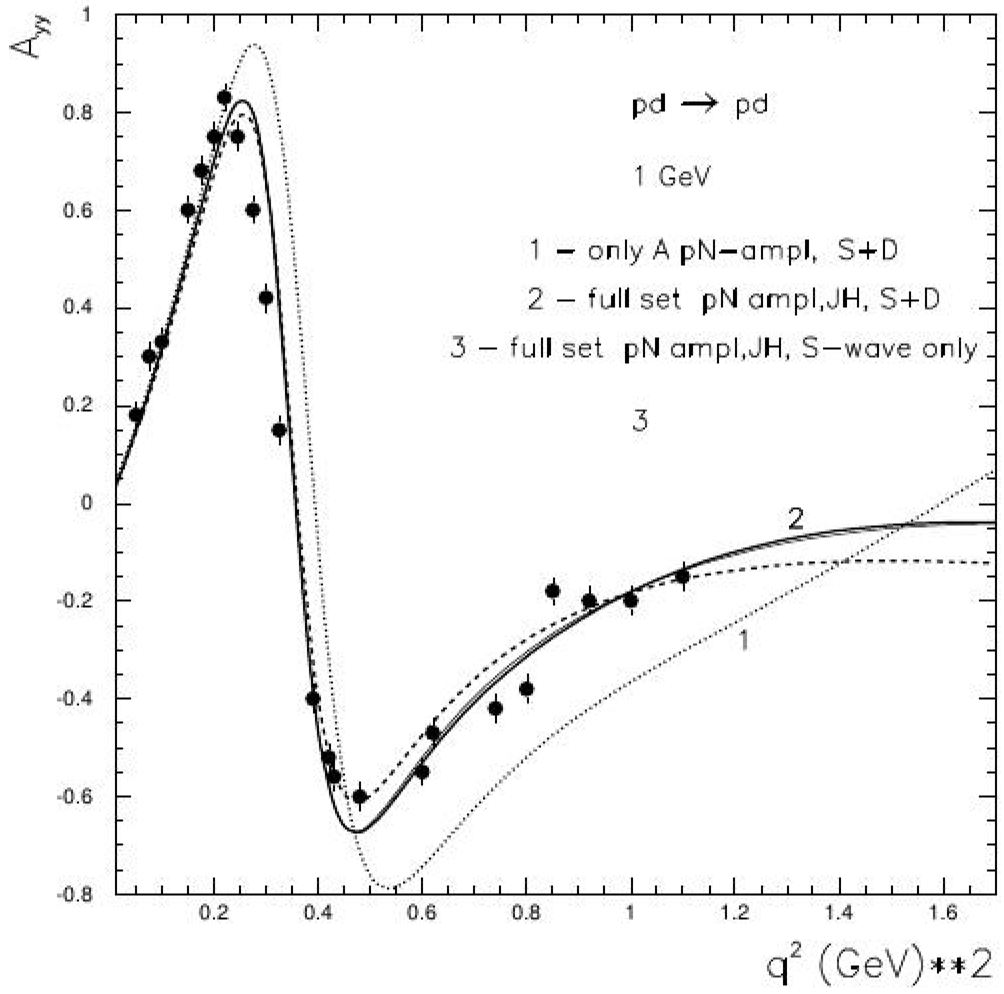
Data: K. Sekiguchi et al. PRC (2002); B. von Przewoski et al. PRC (2006)
See also Faddeev calculations: A.Deltuva, A.C. Fonseca, P.U. Sauer, PRC 71 (2005) 054005.

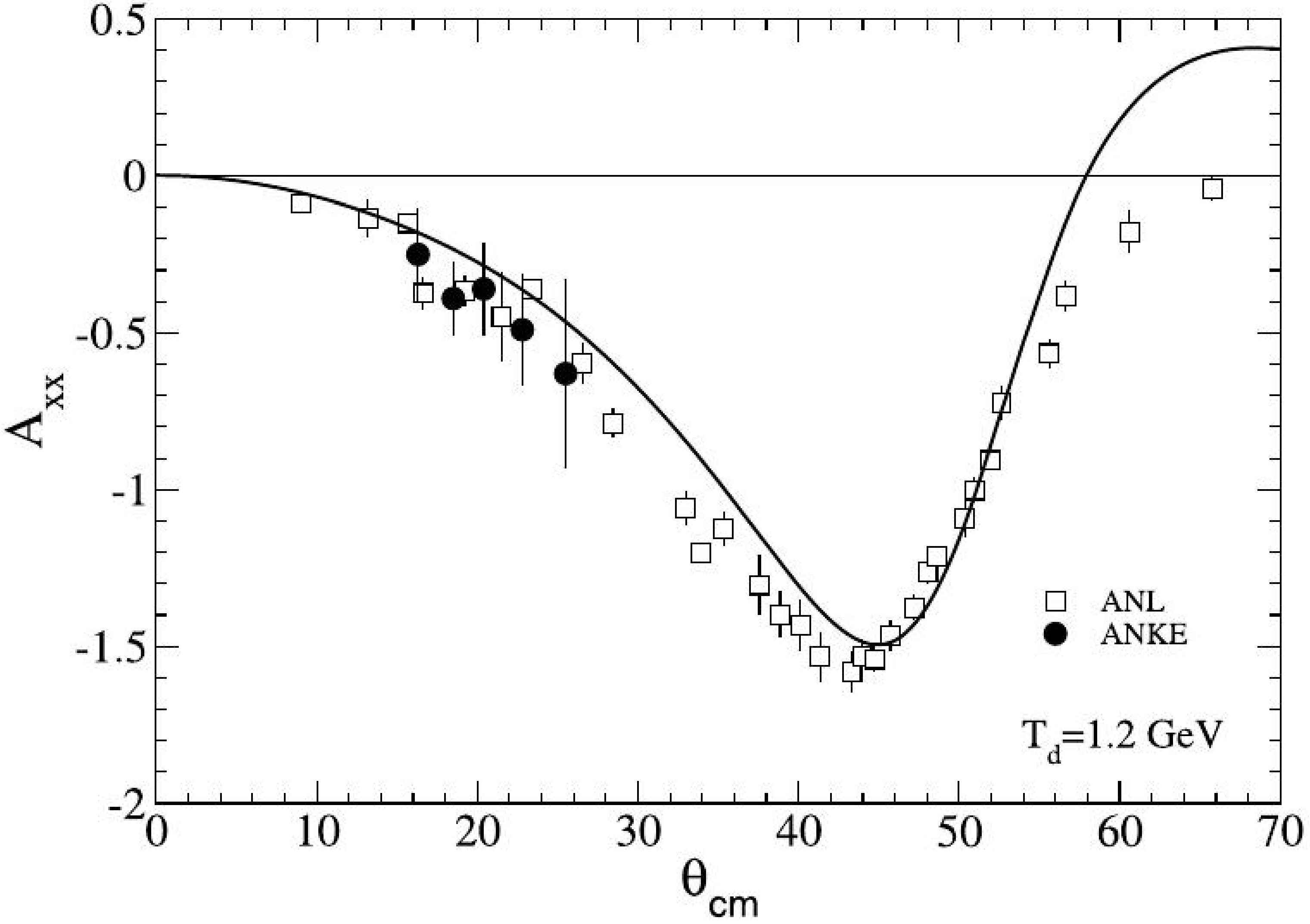


Curves: the modified Glauber model; A.A. Temerbayev, Yu.N.Uzikov, Yad. Fiz. **78** (2015) 38
Data: von B.Przewoski et al. PRC 74 (2006) 064003

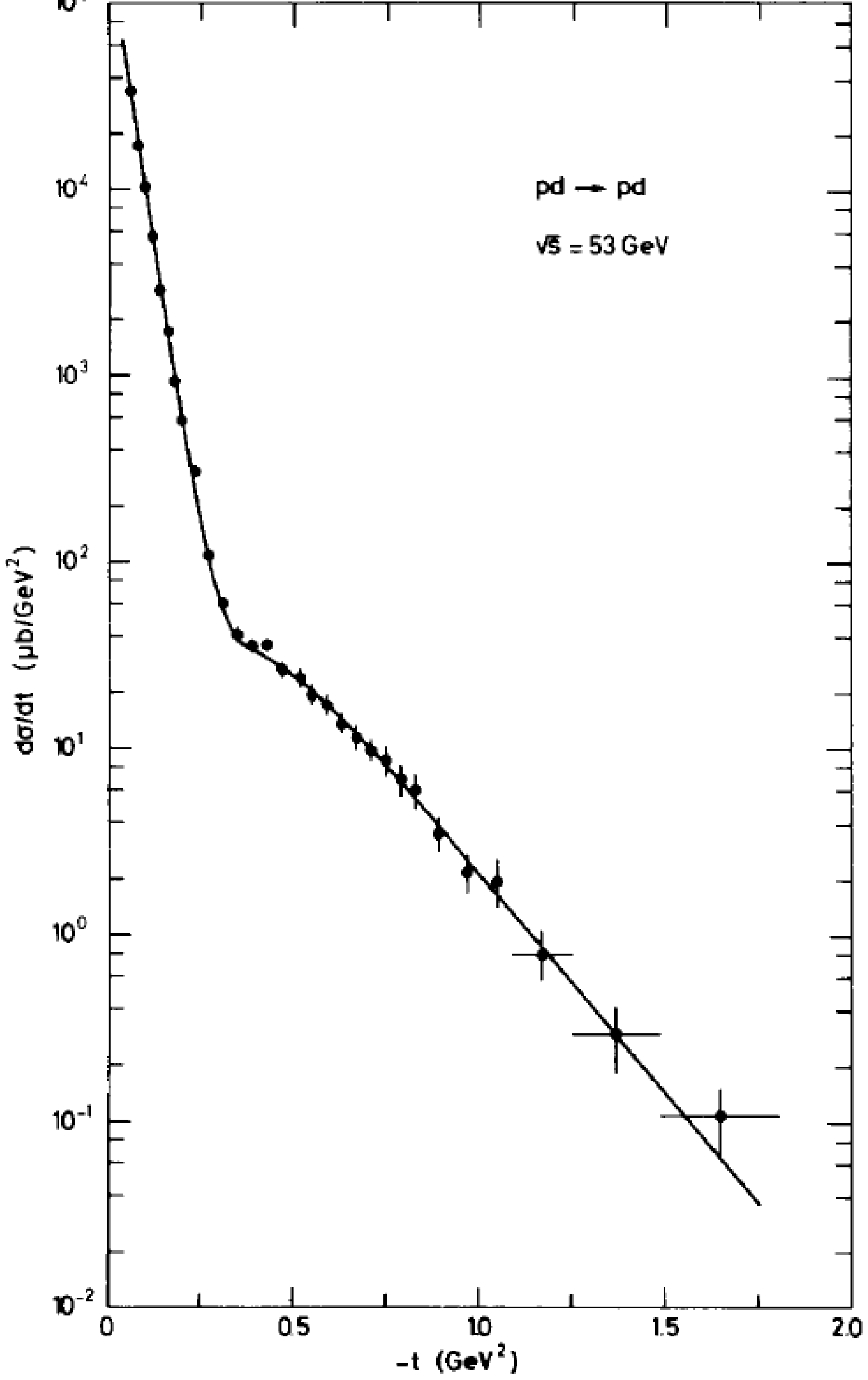
Test calculations: pd elastic scattering at 1 GeV



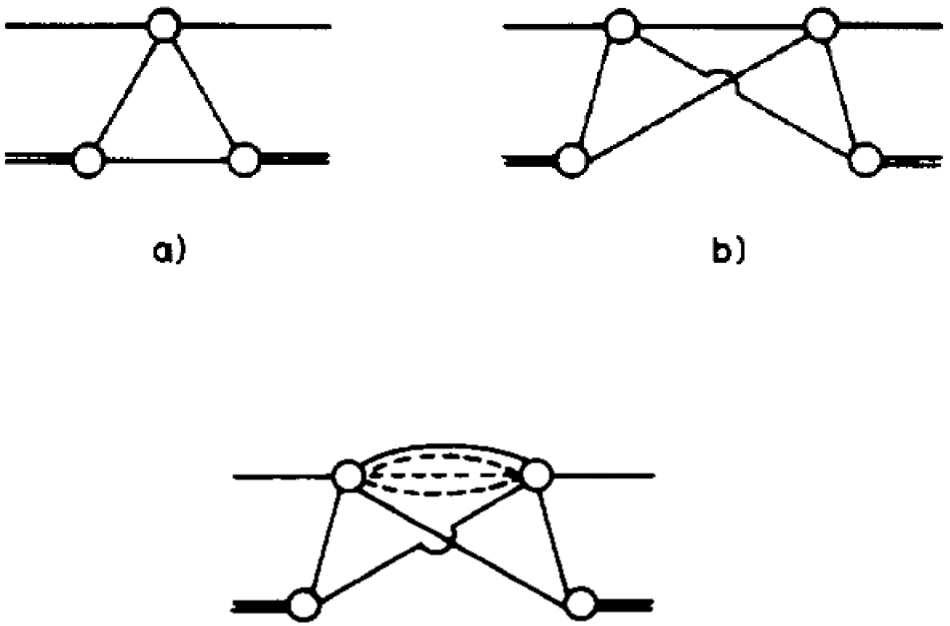




D. Mchedlishvili et al. (ANKE@COSY) Nucl.Phys. A977 (2018)

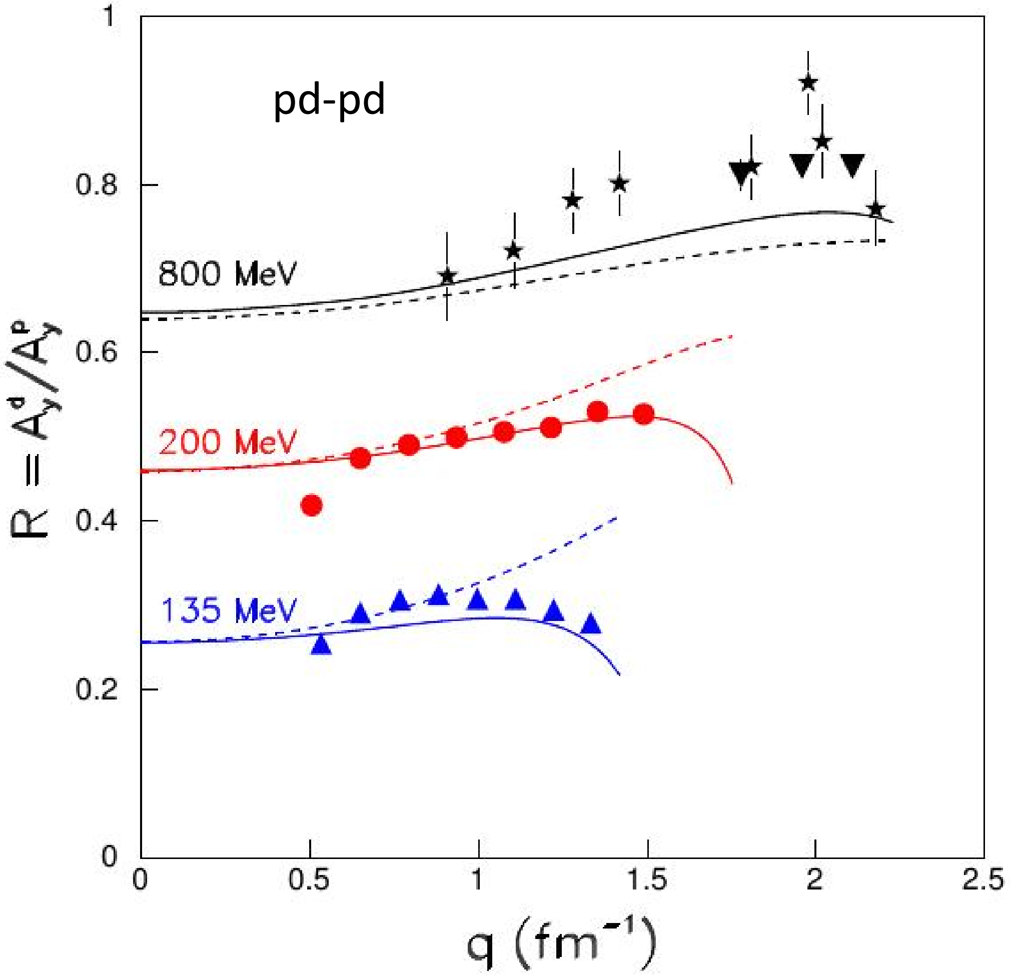


Glauber calculations of the d\sigma/dt for pd-elastic
at ISR energies sqrt{s}=53 and 63 GeV
G.Goggi et al. Nucl. Phys. B149 (1979) 381



$R = A\_y^d / A\_y^p$ Yu.N.U, C. Wilkin, Phys. Lett. B793 (2019) 224,

$R \neq 2/3$ is sensitive to spin-spin NN terms



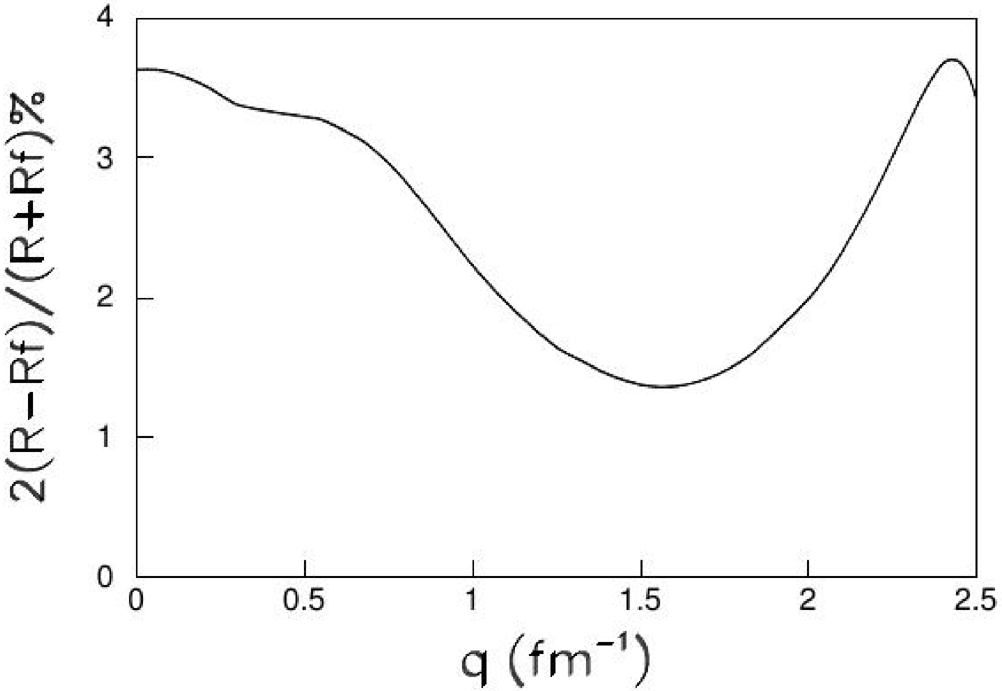


Figure 2: Difference between the predictions of the refined Glauber model [10] without (R) and with (Rf) the NN spin-spin contribution at 800 MeV expressed as a percentage of their average.

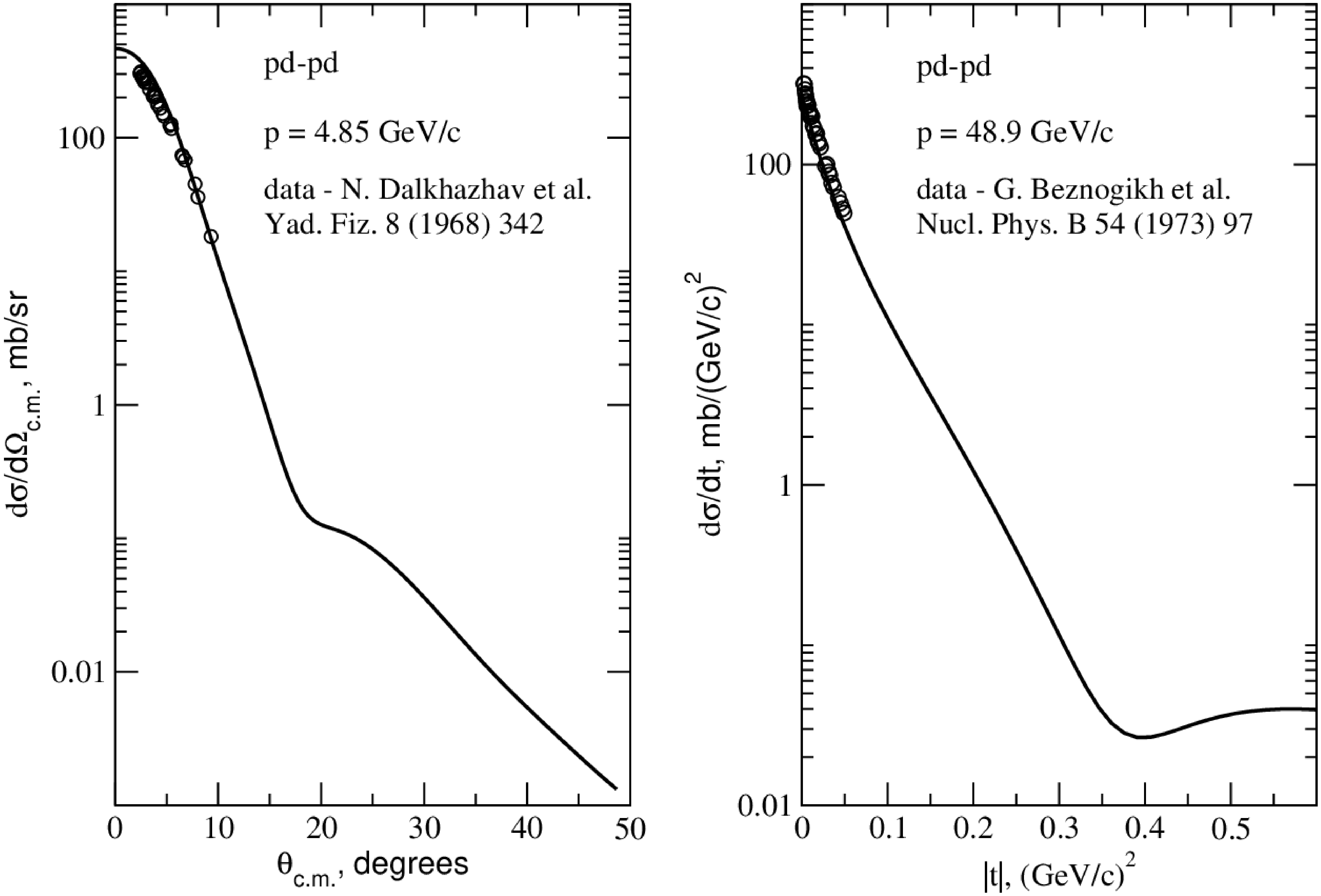
RESULTS OF pd-pd CALCULATIONS within THE GLAUBER THEORY FOR SPD NICA

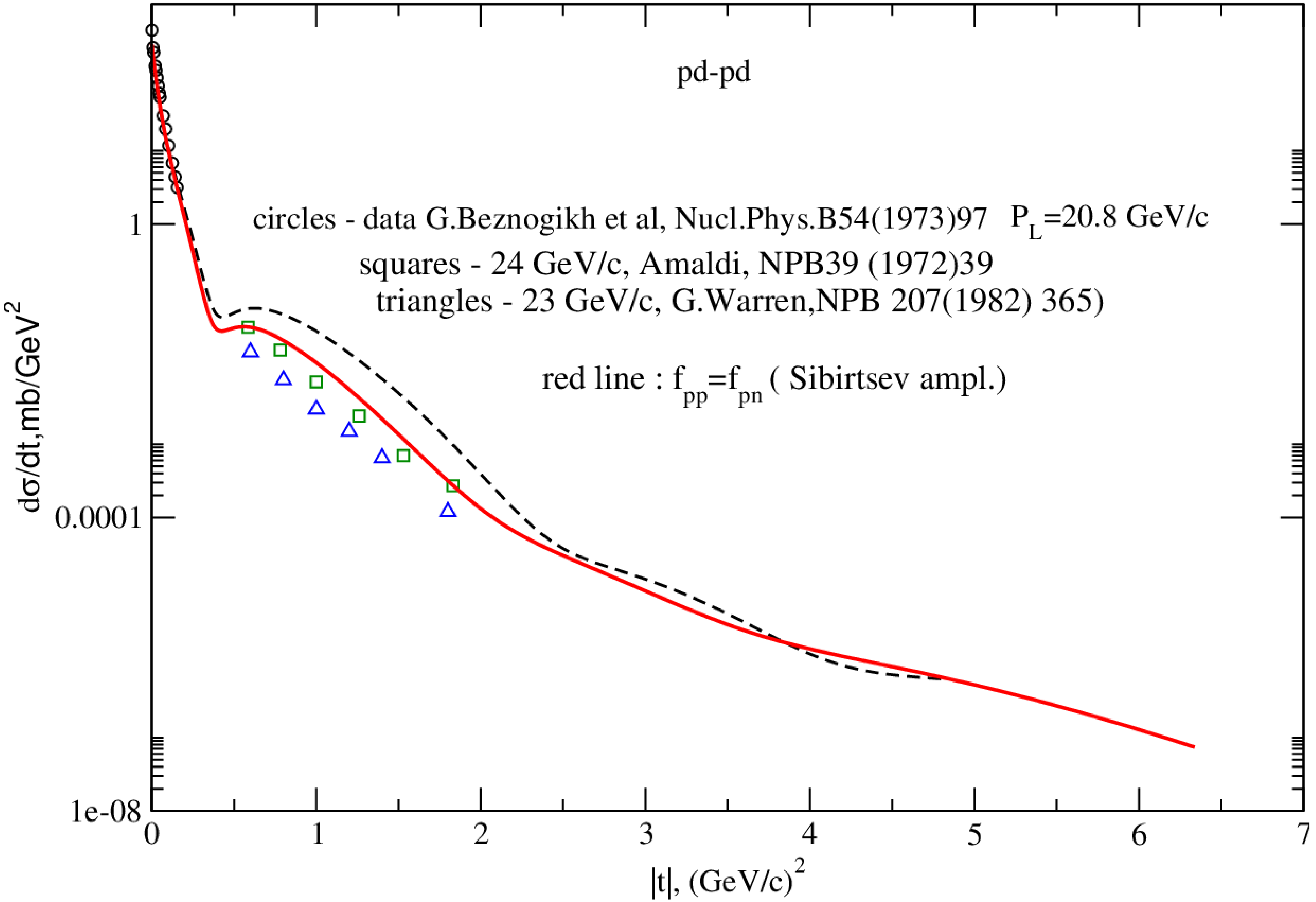
pp $\sqrt{s_{NN}} = 3.4 - 27 GeV$

dd $\sqrt{s_{dd}} = 6.7 - 28 GeV$

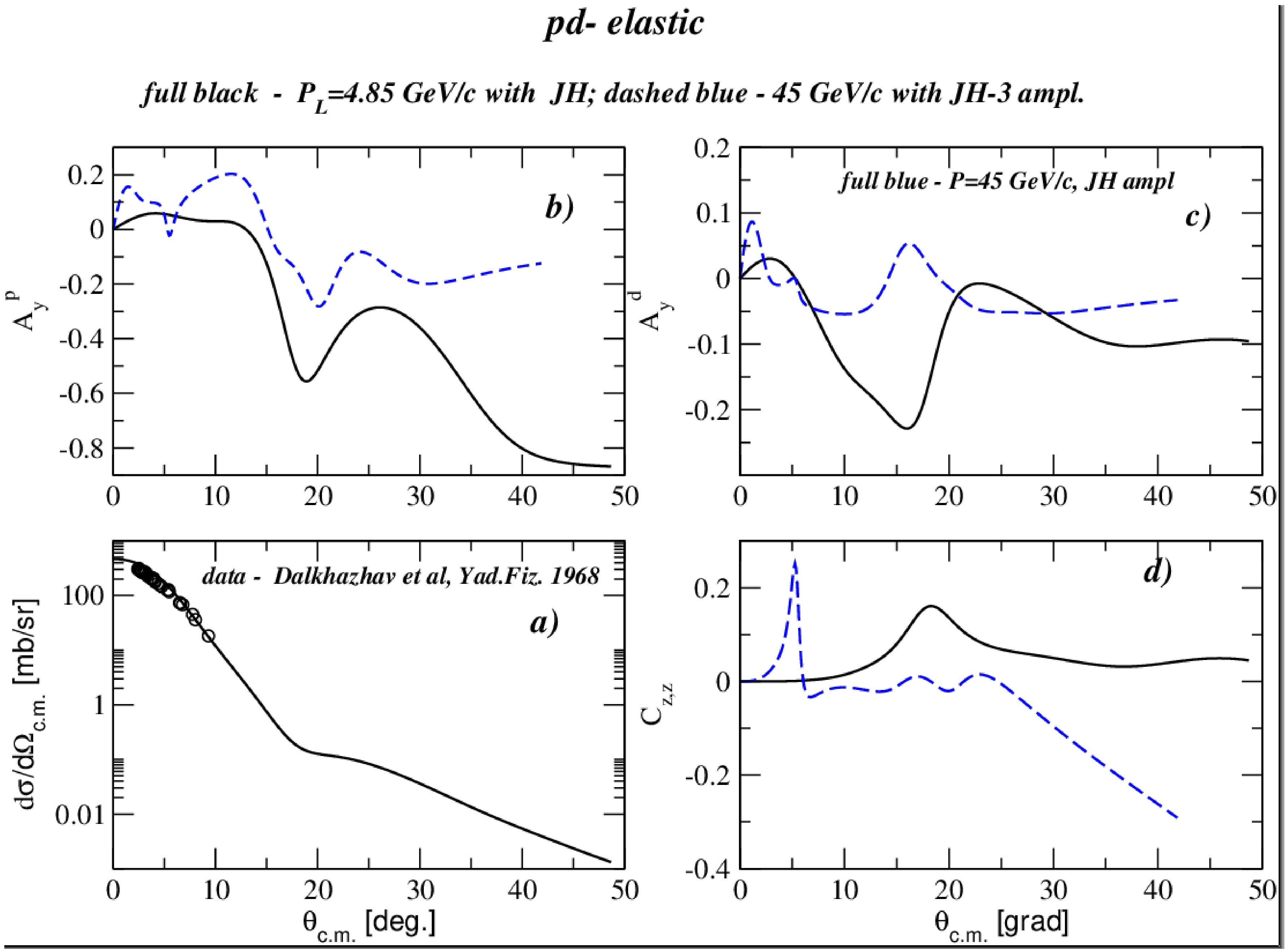
pd $\sqrt{s_{pd}} = 4.76 - 15.3 GeV$

$$L \sim 10^{27} - 10^{32} cm^{-2} s^{-1}$$


pd-pd
p = 4.85 GeV/c
data - N. Dalkhazhav et al.
Yad. Fiz. 8 (1968) 342
dσ/dΩc.m., mb/sr
100
1
0.01
0
10
20
30
40
50
θc.m., degrees
pd-pd
p = 48.9 GeV/c
data - G. Beznogikh et al.
Nucl. Phys. B 54 (1973) 97
dσ/dt, mb/(GeV/c)2
100
1
0.01
0
0.1
0.2
0.3
0.4
0.5
|t|, (GeV/c)2



pd-pd
circles - data G.Beznogikh et al, Nucl.Phys.B54(1973)97 P_L=20.8 GeV/c
squares - 24 GeV/c, Amaldi, NPB39 (1972)39
triangles - 23 GeV/c, G.Warren,NPB 207(1982) 365)
red line : f_pp=f_pn (Sibirtsev ampl.)
dσ/dt,mb/GeV^2
1
0.0001
1e-08
0
1
2
3
4
5
6
7
|t|, (GeV/c)^2

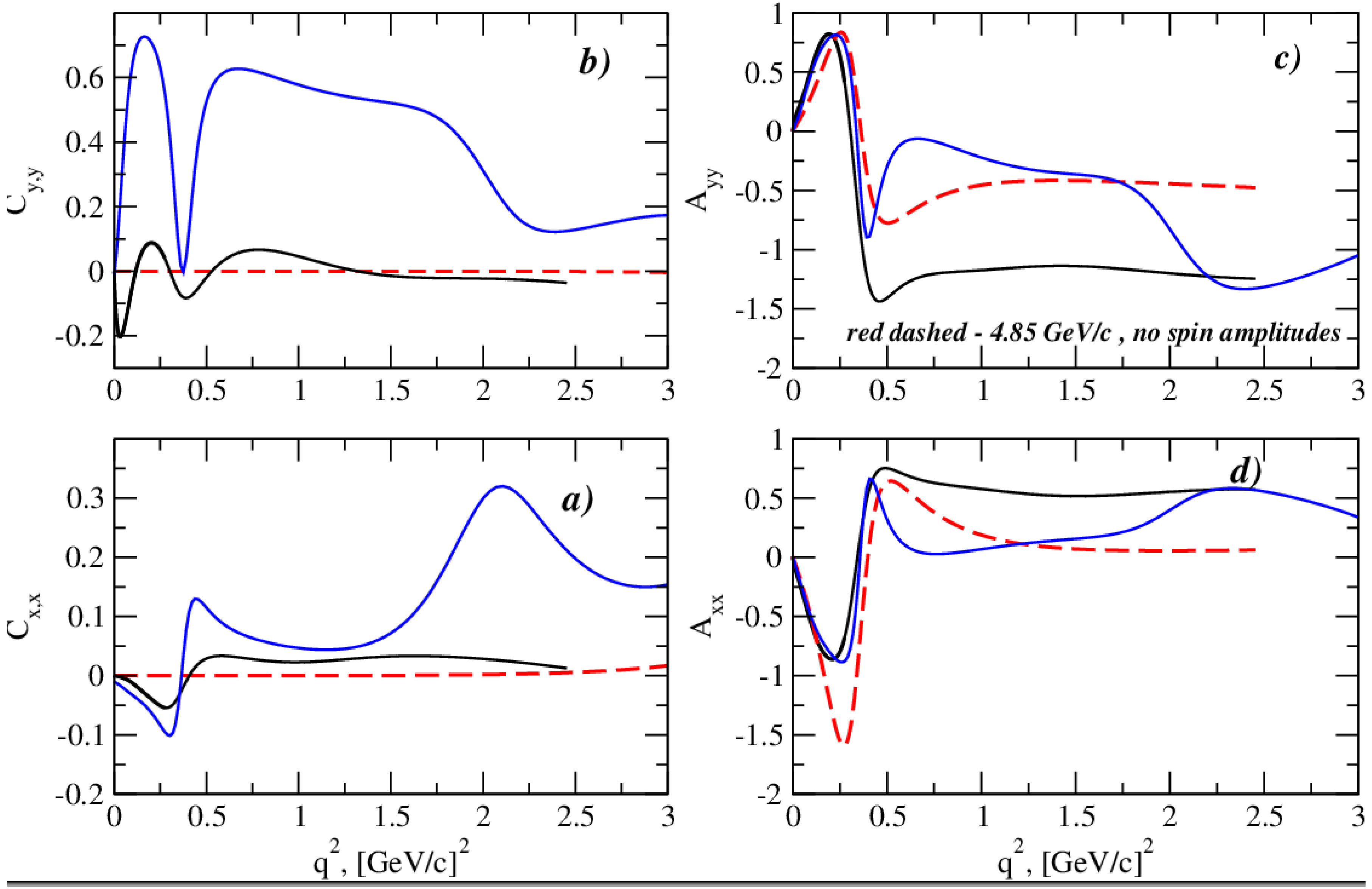


pd- elastic
full black - P_L=4.85 GeV/c with JH; dashed blue - 45 GeV/c with JH-3 ampl.
b)
A_y^p
0.2
0
-0.2
-0.4
-0.6
-0.8
0
10
20
30
40
50
full blue - P=45 GeV/c, JH ampl
c)
A_y^d
0.2
0.1
0
-0.1
-0.2
data - Dalkhazhav et al, Yad.Fiz. 1968
a)
$d\sigma/d\Omega_{c.m.}$ [mb/sr]
100
1
0.01
$\theta_{c.m.}$ [deg.]
d)
$C_{z,z}$
0.2
0
-0.2
-0.4
$\theta_{c.m.}$ [grad]

pd- elastic

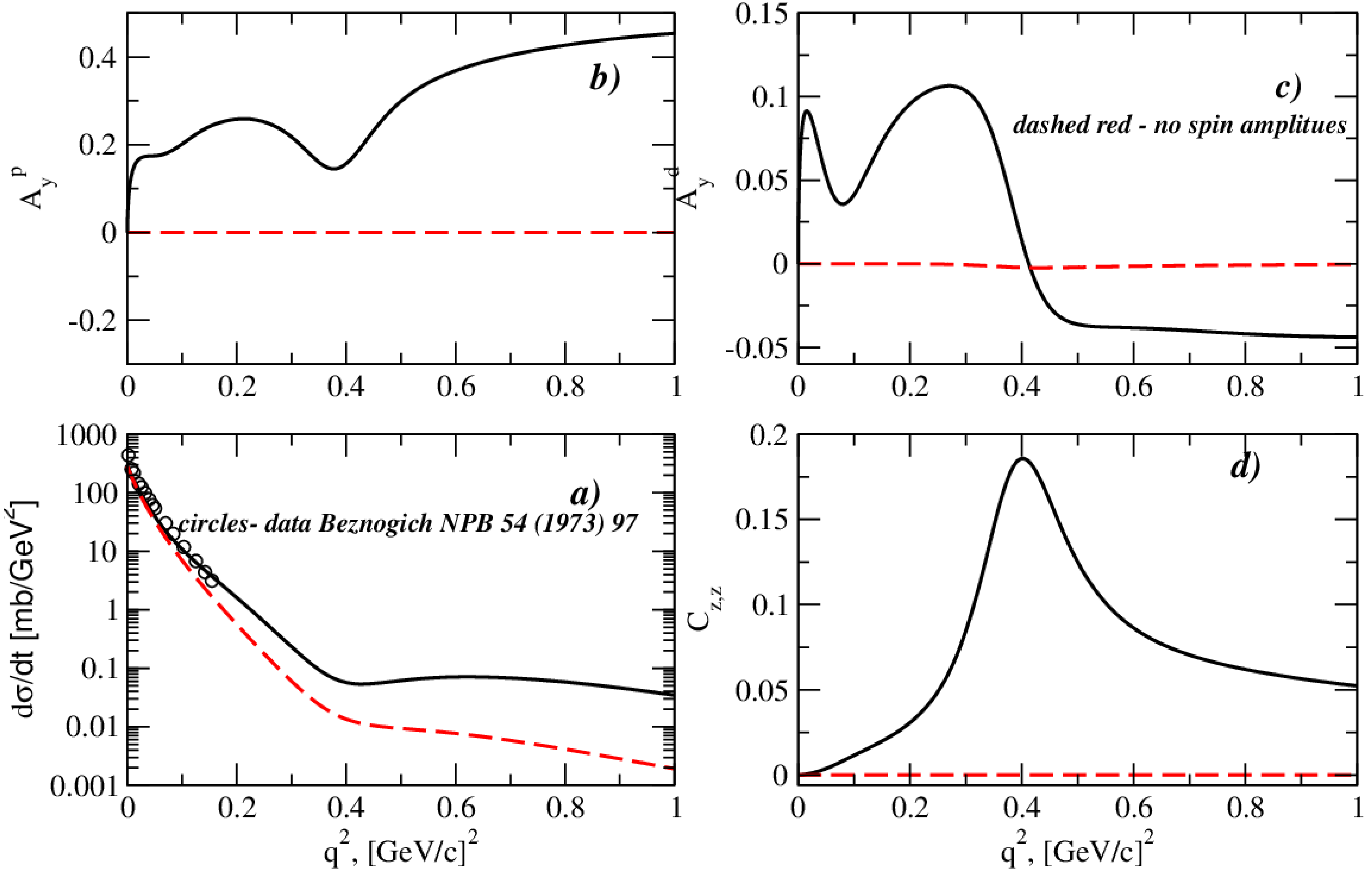
full black - P_L=4.85 GeV/c with Sibirtsev amplitudes; full blue - 45 GeV/c (JH)

dashed red - 4.85 GeV/c without spin dependent pN amplitudes



pd- elastic

full black - P_L=20.4 GeV/c with Sibirtsev amplitudes



Quasielastic scattering

$$p + d \to \{pp\}(^1S_0) + n$$

$$p + d \rightarrow \{pp\}_s + n$$

$$\mathcal{F} = \mathcal{A}(\mathbf{e} \cdot \mathbf{k})(\boldsymbol{\sigma} \cdot \mathbf{k}) + \mathcal{B}\mathbf{e} \cdot \boldsymbol{\sigma},$$

$$p + n \rightarrow n + p$$

$$f_{12}^{collin} = \alpha + \beta(\boldsymbol{\sigma}_1 \cdot \boldsymbol{\sigma}_2) + (\varepsilon - \beta)(\boldsymbol{\sigma}_1 \cdot \mathbf{k})(\boldsymbol{\sigma}_2 \cdot \mathbf{k}).$$

$$\mathcal{A} = (\epsilon - \beta)S(Q/2), \quad \mathcal{B} = \beta S(Q/2)$$

$$d\sigma_0 = \frac{1}{3}\mathcal{K}\left\{|\varepsilon|^2 + 2|\beta|^2\right\},$$

$$T_{20} = \frac{1}{\sqrt{2}}A_{zz} = \sqrt{2}\frac{|\beta|^2 - |\varepsilon|^2}{|\varepsilon|^2 + 2|\beta|^2},$$

$$C_{x,x} = C_{y,y} = -2\frac{Re\varepsilon\beta^*}{|\varepsilon|^2 + 2|\beta|^2},\ C_{xz,y} = -C_{yz,x} = 3\frac{Im\beta\varepsilon^*}{|\varepsilon|^2 + 2|\beta|^2}.$$

Complete polarization experiment gives: $|\varepsilon|, |\beta|, \mathrm{Re}\,\varepsilon\beta^*, \mathrm{Im}\,\varepsilon\beta^*$

Search for T-invariance violation in double polarized pd - scattering

Yu.N. U. , J. Haidenbauer, PRC 94 (2016) 035501

Yu. N. U., A. Temerbayev , PRC 92 (2015) 014002

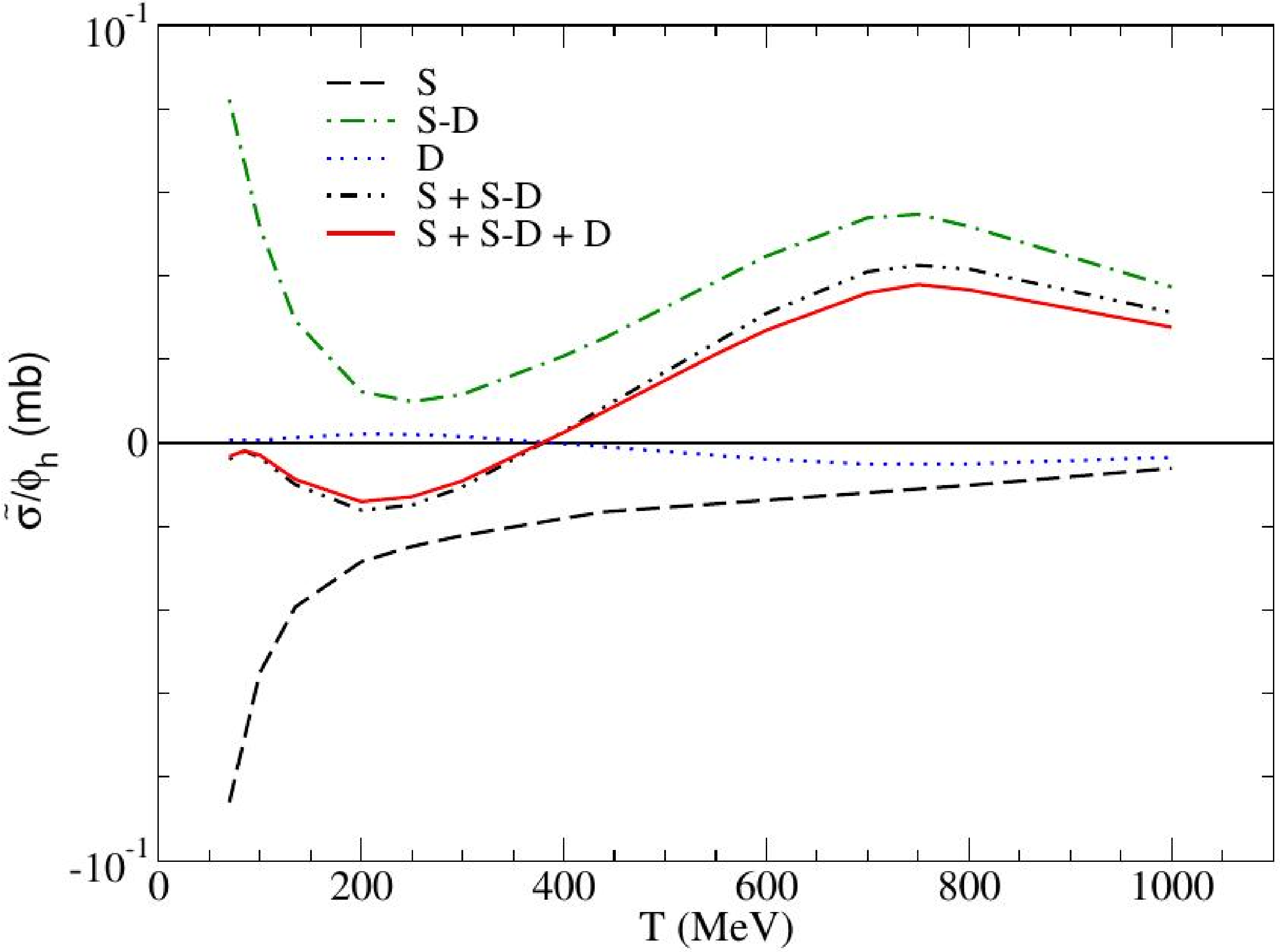
N. Nikolaev, F.Rathman, A. Silenko, Yu.N. Uzikov, 2004.09943 [nucl-th]

$$C' \approx i\phi_5 + iq/2m(\phi_1 + \phi_3)/2$$

Yu.N.U., A.A. Temerbayev, PRC 92 (2015) 014002;
Yu.N.U., J. Haidenabuer, PRC 94 (2016) 035501.

$$\sigma_{tot} = \underbrace{\sigma_0 + \sigma_1 \mathbf{p}^p \cdot \mathbf{P}^d + \sigma_2 (\mathbf{p}^p \cdot \hat{\mathbf{k}})(\mathbf{P}^d \cdot \hat{\mathbf{k}}) + \sigma_3 P_{zz}}_{T-even, P-even} + \underbrace{\widetilde{\sigma}_{tvpc} p_y^p P_{xz}^d}_{T-odd, P-even}$$

TVPC. The S- and D- wave contributions-I



Yu.N. Uzikov, J.Haidenbauer, PRC 94 (2016) 035501

CONCLUSION

- Spin-dependent Glauber theory of pd- elastic scattering works well in describing of pd spin observables at moderate energies 0.1-3 GeV. Hopefully, this theory should work also at higher energies, 3-50 GeV .
- Existing data on pp- and especially pn- elastic spin amplitudes above 3 GeV are rather poor, although different phenomenological parametrizations are available.
- Test of T-invariance in pd double polarized scattering at NICA energies requires knowledge of pN-spin amplitudes.
- Measurement of single and *double spin observables of pd- elastic and quasielastic scattering at SPD NICA energies and comparison with our results of the Glauber calculations* will give a clean test for pp- and pn- amplitudes.

Thank you for attention!

APPENDIX

T-even P-even

$$M_N(\mathbf{p},\mathbf{q};\boldsymbol{\sigma},\boldsymbol{\sigma}_N)$$
$$= A_N + C_N\boldsymbol{\sigma}\hat{\mathbf{n}} + C'_N\boldsymbol{\sigma}_N\hat{\mathbf{n}} + B_N(\boldsymbol{\sigma}\hat{\mathbf{k}})(\boldsymbol{\sigma}_N\hat{\mathbf{k}})$$
$$+ (G_N + H_N)(\boldsymbol{\sigma}\hat{\mathbf{q}})(\boldsymbol{\sigma}_N\hat{\mathbf{q}}) + (G_N - H_N)(\boldsymbol{\sigma}\hat{\mathbf{n}})(\boldsymbol{\sigma}_N\hat{\mathbf{n}})$$

T-odd P-even , M.Simonius,PRL 1997

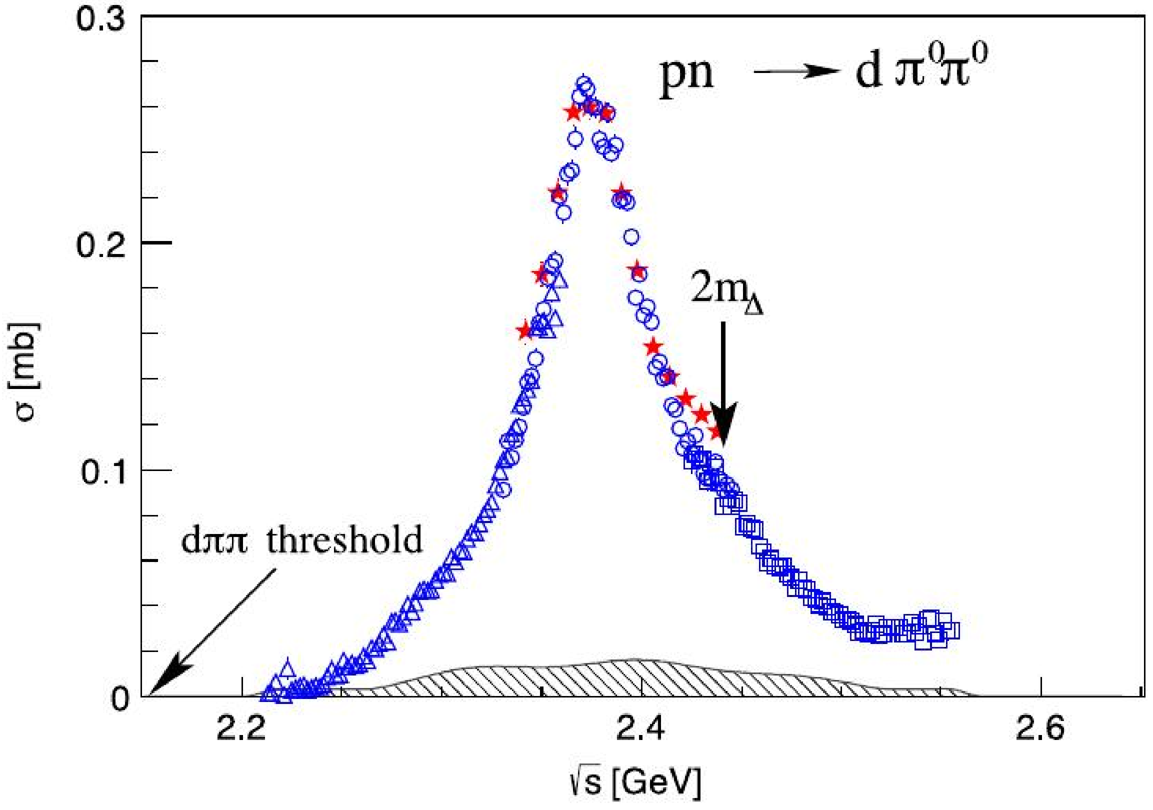
$$t_{pN} = h_N[(\boldsymbol{\sigma}\cdot\mathbf{k})(\boldsymbol{\sigma}_N\cdot\mathbf{q}) + (\boldsymbol{\sigma}_N\cdot\mathbf{k})(\boldsymbol{\sigma}\cdot\mathbf{q})$$
$$- \tfrac{2}{3}(\boldsymbol{\sigma}_N\cdot\boldsymbol{\sigma})(\mathbf{k}\cdot\mathbf{q})]/m_p^2$$
$$+ g_N[\boldsymbol{\sigma}\times\boldsymbol{\sigma}_N]\cdot[\mathbf{q}\times\mathbf{k}][\boldsymbol{\tau}-\boldsymbol{\tau}_N]_z/m_p^2$$
$$+ g'_N(\boldsymbol{\sigma}-\boldsymbol{\sigma}_N)\cdot i\,[\mathbf{q}\times\mathbf{k}][\boldsymbol{\tau}\times\boldsymbol{\tau}_N]_z/m_p^2.$$

Null-test signal:

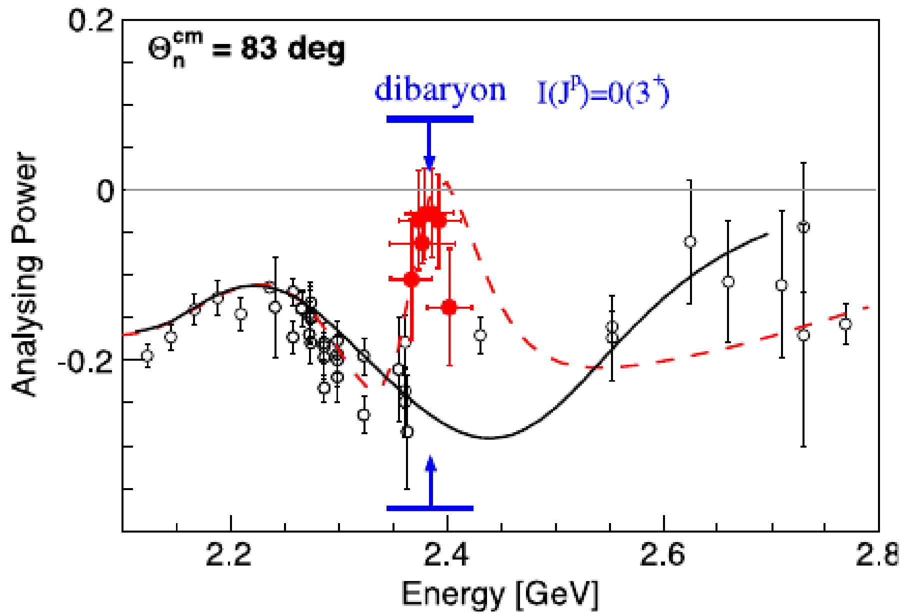
$$\tilde{g} = \frac{i}{4\pi m_p}\int_0^\infty dq q^2\Big[S_0^{(0)}(q) - \sqrt{8}S_2^{(1)}(q) - 4S_0^{(2)}(q)$$
$$+\sqrt{2}\frac{4}{3}S_2^{(2)}(q)+9S_1^{(2)}(q)\Big][-C'_n(q)\,h_p+C'_p(q)(g_n - h_n)]$$

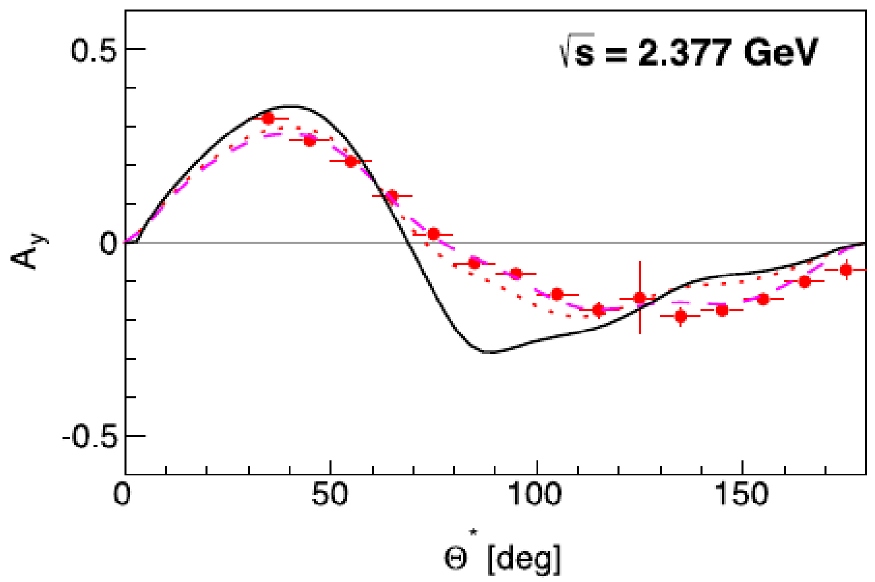
d*(2380) dybarion $J^P = 3^+$

H. Clement / Progress in Particle and Nuclear Physics 93 (2017) 195–242



H. Clement / Progress in Particle and Nuclear Physics 93 (2017) 195–242





M.N. Platonova, V.I. Kukulin, EPJ A (2020) 56:132

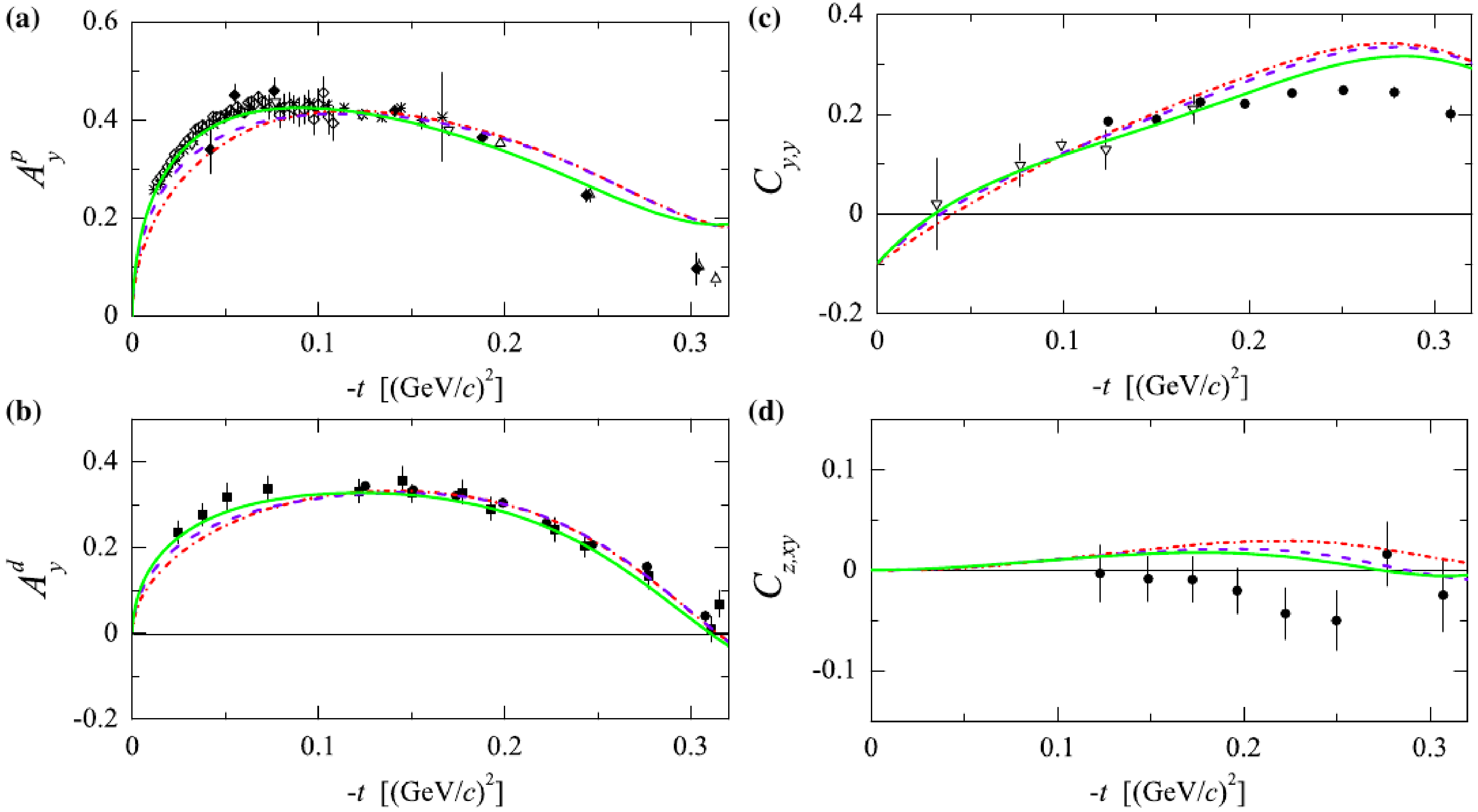


Fig. 8 Proton analyzing power A_y^p (**a**), deuteron analyzing power A_y^d (**b**), vector spin-correlation parameter $C_{y,y}$ (**c**) and tensor spin-correlation parameter $C_{z,xy}$ (**d**) in pd (dp) elastic scattering at the incident (equivalent) proton energy T_p = 800 MeV. Dash-dotted (red), dashed (violet) and solid (green) lines show the refined Glauber model calculations with NN amplitudes corresponding to the SAID PWA solution SM16 [35,41], with the modified amplitude C_p and with modified amplitudes C_p and C_n (see Fig. 6), respectively. Experimental data at T_p = 796 and 800 MeV are the same as in Figs. 2, 4 and 5

– *Planned experiments to search for CP violation beyond the SM*

• Detecting a non-zero **EDM** of elementary fermion (neutron, atoms, charged particles). The current experimental limit

$$|d_n| \leq 2.9 \times 10^{-26} e\,cm$$

is much less as compared the SM estimation (B.H.J. McKellar et al. PLB 197 (1987)

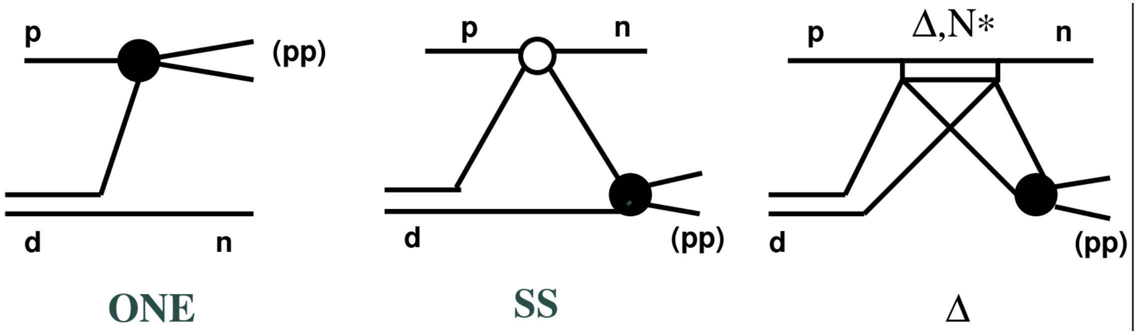
$$1.4 \times 10^{-33} e\,cm \leq |d_n| \leq 1.6 \times 10^{-31} e\,cm$$

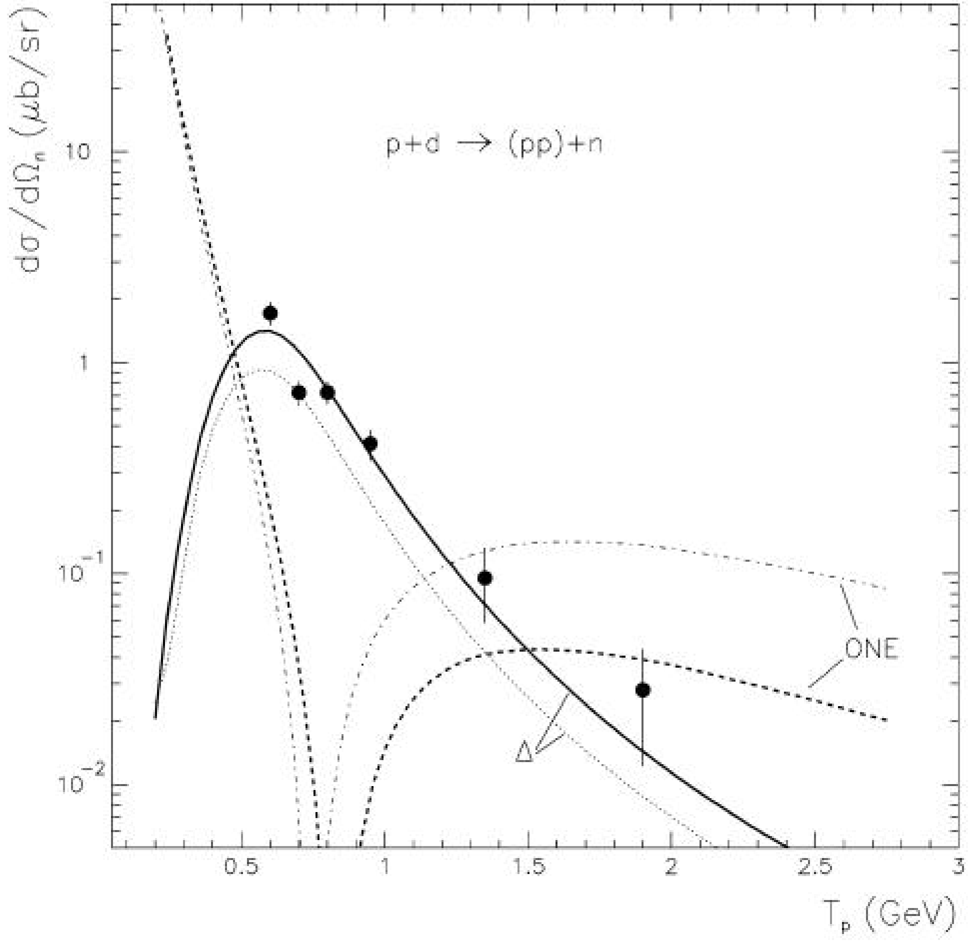
• Search for CP violation in the **neutrino sector** ($\theta_{13} \neq 0$, then generation of lepton asymmetry and via $B - L$ conservation to get the BAU).

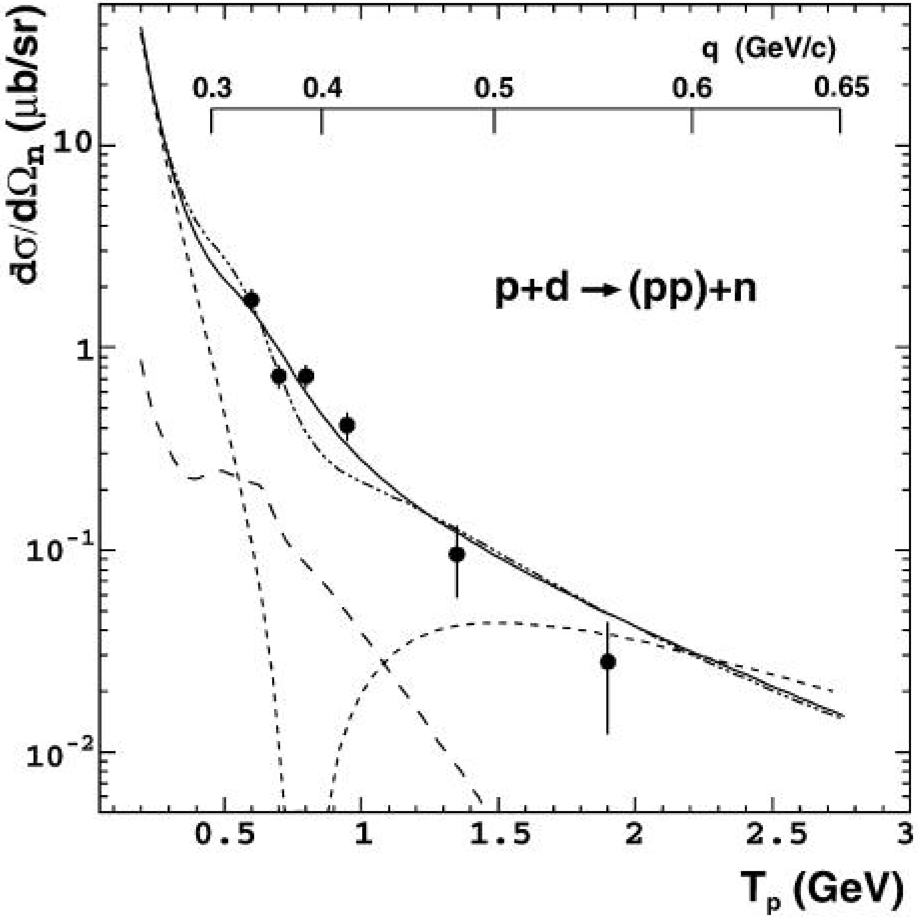
Thouse are T-violating and Parity violating (TVPV) effects.

Much less attention was paid to T-violating P-conserving (TVPC) flavor conserving effects.

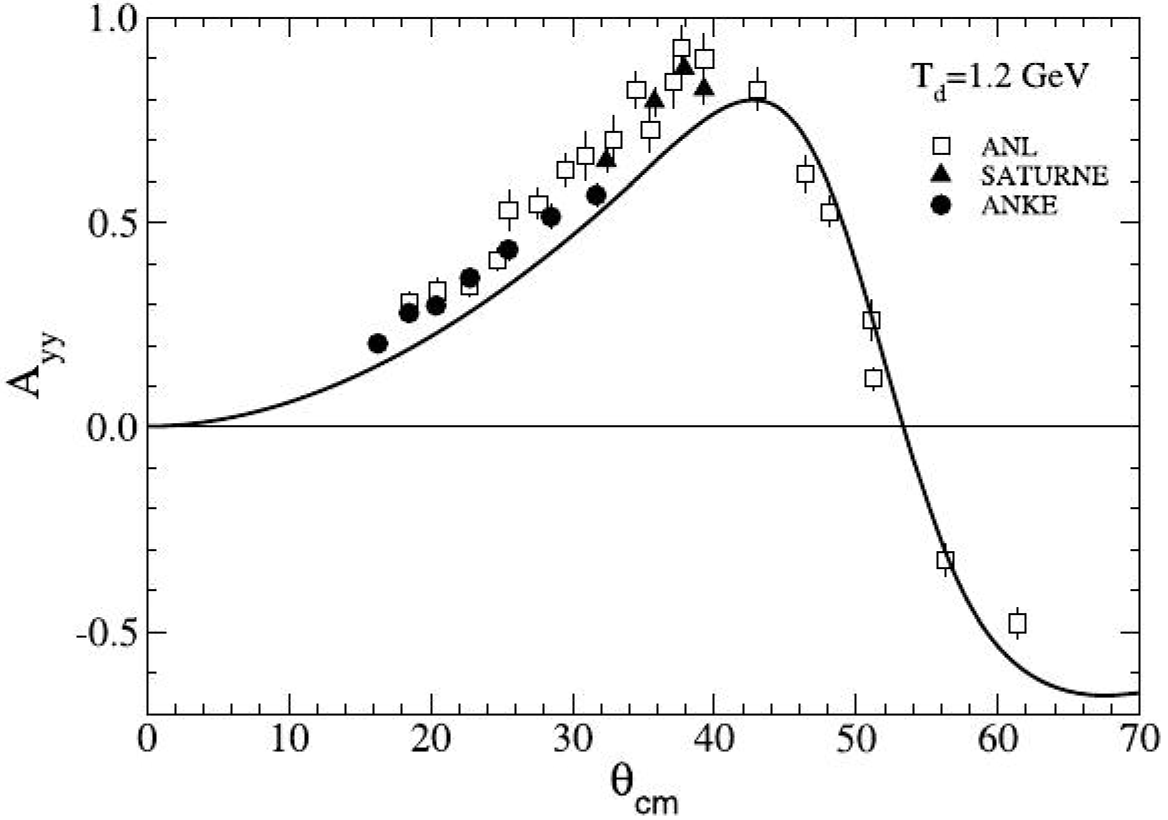
J. Haidenbauer, Yu. Uzikov, PLB (2003)

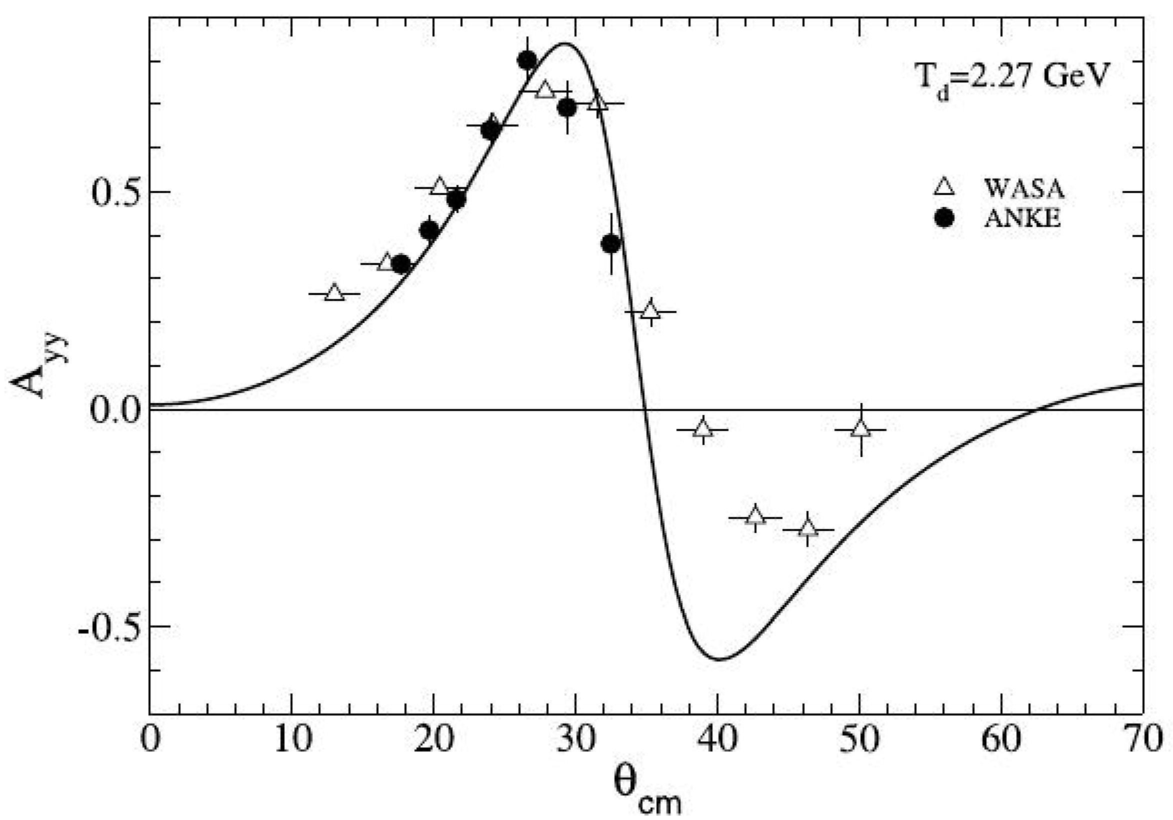


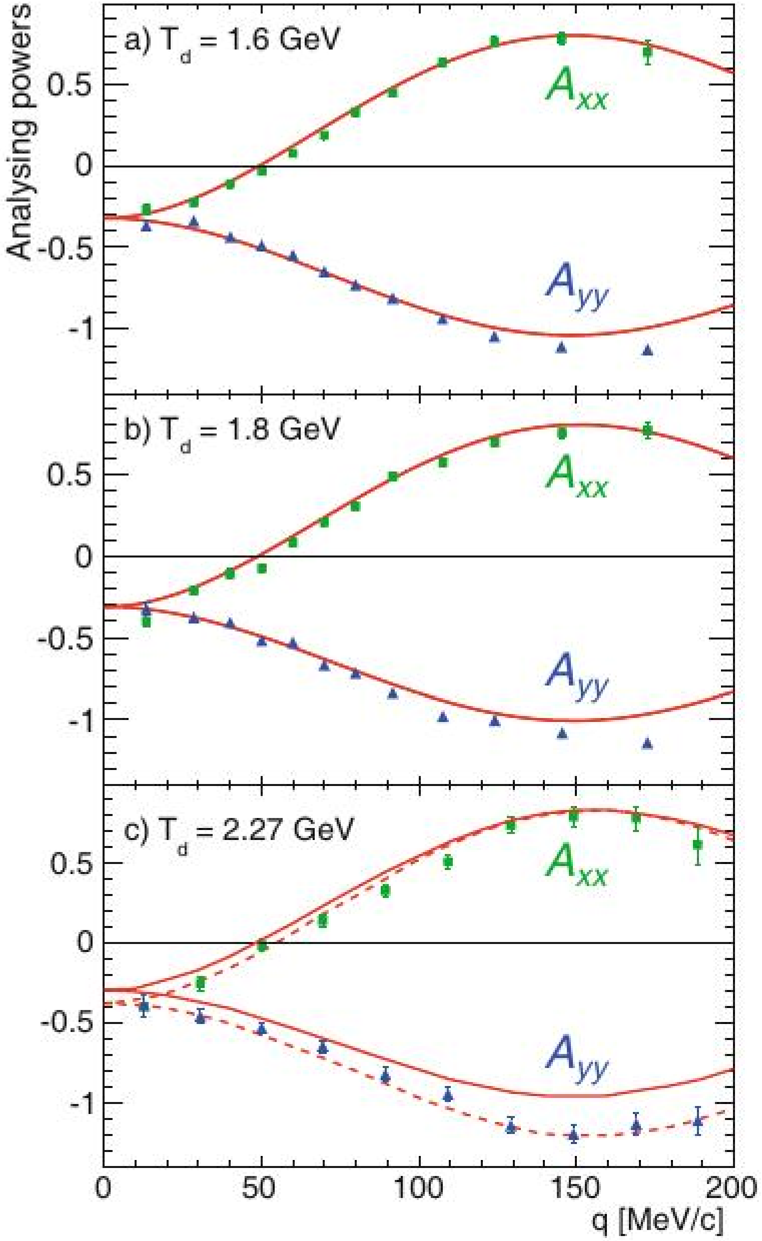




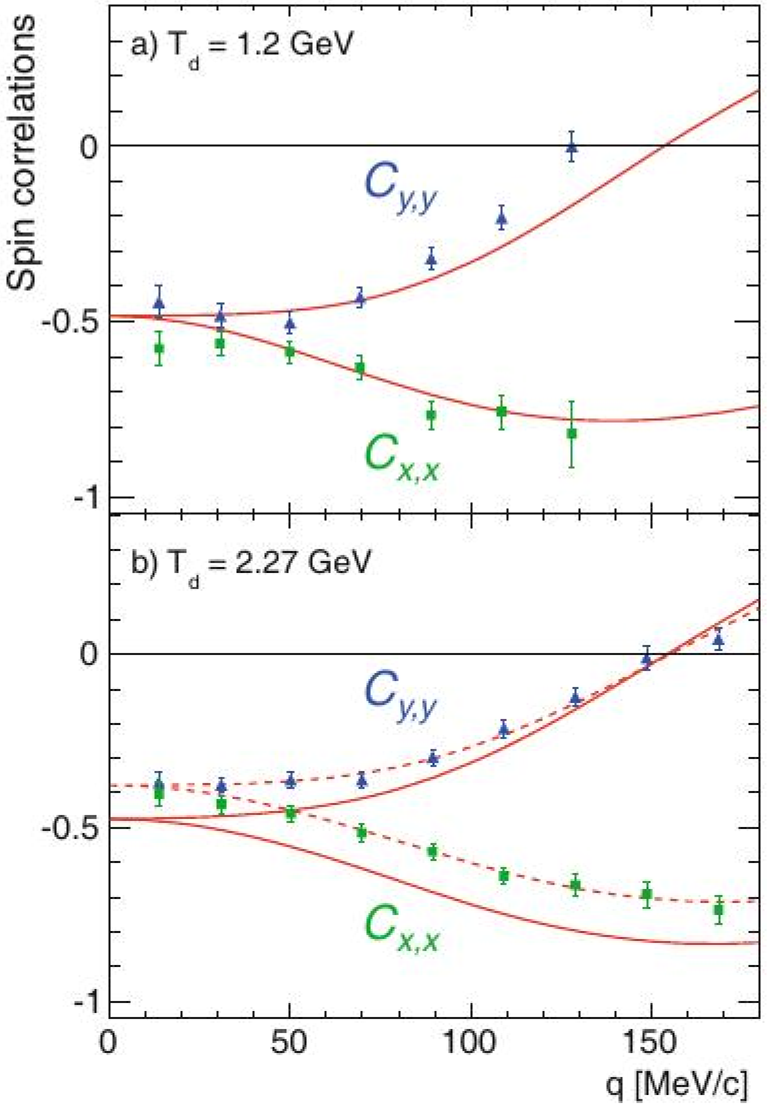
D. Mchedlishvili et al. Nucl.Phys. A 977 (2018) 14 , pd-elastic







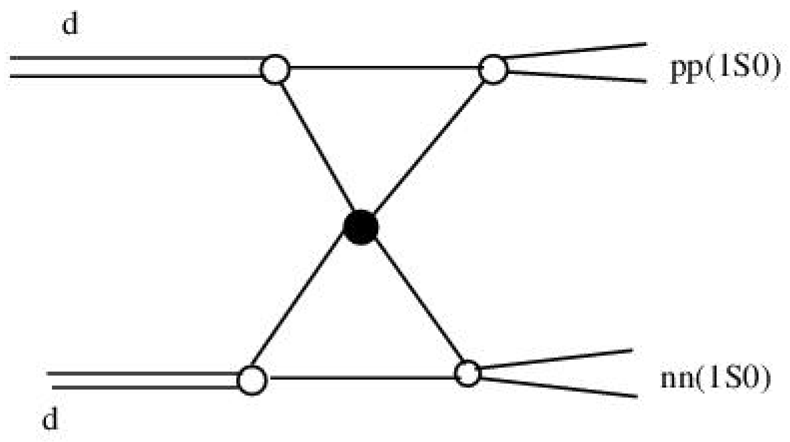
'ig. 8. Tensor analysing powers A_{xx} (squares) and A_{yy} (tri-
ngles) of the $\vec{d}p \to \{pp\}_s n$ reaction at three beam energies for
ɔw diproton excitation energy, $E_{pp} < 3\,\mathrm{MeV}$, compared to im-

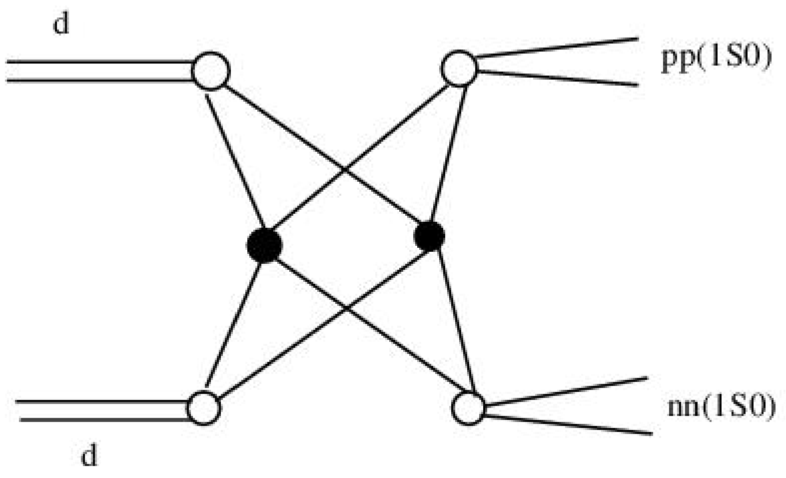


pd-> (pp)+n, E_pp< 3MeV, 1S0 ANKE

D. Mchedlishvili, et al. EPJA 49 (2013)

dd- elastic and quasi-elastic scattering





$$Del = \frac{d\sigma / dt_{data.} - d\sigma / dt_{theor-\exp.}}{d\sigma / dt_{theor-\exp.}}$$

P. Gauron, B. Nicolescu, O.V. Selyugin,PLB 397 (1997)