

# Structure functions generated by zero sound excitations

V.A. Sadovnikova

*National Research Center "Kurchatov Institute"*

*B. P. Konstantinov Petersburg Nuclear Physics Institute*

*Gatchina, St. Petersburg 188300, Russia*

We study the form of the structure functions connected to the zero sound excitations in the symmetric and asymmetric nuclear matter (ANM). The density response  $\Pi(\omega, k)$  (the retarded polarization operator) of ANM on the weak external field  $V_0(\omega, k) = \tau_z e^{i\vec{q}\vec{r} - i(\omega + i\eta)t}$  is considered. The structure functions  $S(\omega, k)$  is defined as  $S(\omega, k) = -\frac{1}{\pi} \text{Im} \Pi(\omega, k)$  [1].

In [2] the three complex branches of the zero sound excitations in ANM were obtained:  $\omega_{si}(k)$ ,  $i = n, p, np$ . We calculate these branches as solutions of the dispersion equation  $E(\omega, k) = 0$ . Calculations were made in the framework of RPA with the Landau-Migdal quasiparticle-quasihole isovector interaction  $F'(\vec{\tau}, \vec{\tau}')$ .

It was shown that in the external field  $V_0(\omega, k)$  the total polarization operator is the sum [4]:  $\Pi = \Pi^{pp} + \Pi^{nn} - \Pi^{pn} - \Pi^{np}$ . Expressions for  $\Pi^{\tau, \tau'}$  are obtained from the system of the type similar to the system for the effective fields in [3, 4]:  $\Pi^{pp} = \Pi_0^p (1 - \Pi_0^n F^{nn}) / \det M(\omega, k) \equiv D^{pp} / \det M(\omega, k)$ ,  $\Pi^{np} = \Pi_0^p \Pi_0^n F^{np} / \det M(\omega, k) \equiv D^{np} / \det M(\omega, k)$ . Changing  $p \leftrightarrow n$  we obtain  $\Pi^{nn}$ ,  $\Pi^{pn}$ . Dispersion equation for the frequencies of zero sound excitations is  $E(\omega, k) \equiv \det M(\omega, k) = 0$ . So, the branches  $\omega_{si}(k)$  are the zeros of  $\det M(\omega, k)$  and the poles of  $\Pi^{\tau, \tau'}$  by construction.

In our approach  $S(\omega, k)$  must be considered as a sum over three independent processes: the widths of the different  $\omega_{si}(k)$  correspond to the different decays of excitations. The imaginary part of  $\omega_{sn}(k)$  describes in nuclei the semidirect decay due to emission of a neutron, reaction  $(\gamma, n)$ . Decay of  $\omega_{sp}(k)$  accompanied by emission of proton. About of  $\omega_{snp}(k)$  we can say that one nucleon is emitted and its isospin is not fixed [2]. We rewrite  $S(\omega, k) = \sum_i S(\omega, k)_i$ .

Near the pole at  $\omega \approx \text{Re}(\omega_{si})$  we approximate  $\det M(\omega, k)^{-1} = R^i(\omega_{si}, k) / (\omega - \omega_{si}) + \text{Reg}(\omega, k)$ . Here  $\text{Reg}(\omega, k)$  is a smooth function near the pole. This permits us to write  $S(\omega, k)_i = -\frac{1}{\pi} \text{Im} \left[ \sum_{\tau, \tau'} (D^{\tau\tau'}(\omega, k)) R^i(\omega_{si}, k) / (\omega - \omega_{si}) + \text{Reg} \right]$ . Then, let define the envelope curve of the pole terms  $S^e(\omega, k) = -\frac{1}{\pi} \sum_{\tau, \tau'} \text{Im} \left[ D^{\tau, \tau'}(\omega, k) \sum_i R^i(\omega_{si}, k) / (\omega - \omega_{si}) \right]$ .

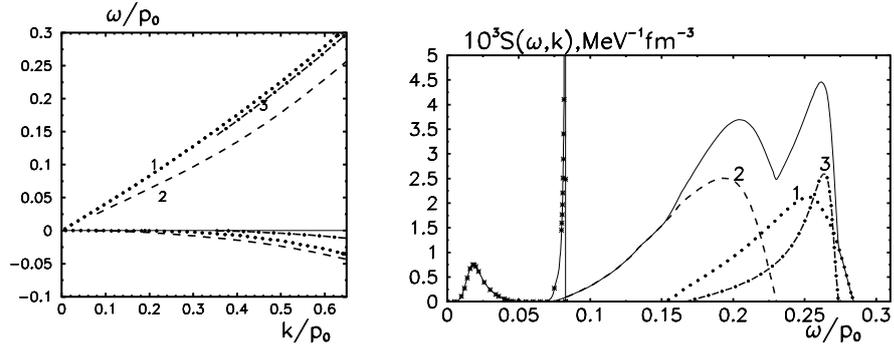


Figure 1: *left* :  $\omega_{sn}(k)$  (dotted curve, 1),  $\omega_{sp}(k)$  (dashed curve, 2),  $\omega_{snp}(k)$  (dash-dotted, 3). At  $\omega > 0$  ( $\omega < 0$ ) the real (imaginary) parts of  $\omega_{si}(k)$  are demonstrated; *right* : the pole terms of  $S(\omega, k)_i$ ,  $i = n, p, np$  (numbers 1, 2, 3) are shown. The envelope curve  $S^e(\omega, k)$  is marked by the solid curve for  $k/p_0 = 0.6$  and by the solid with stars for  $k/p_0 = 0.2$ .

We demonstrate results for ANM with asymmetry parameter  $\beta = 0.2$ . In the left figure the branches  $\omega_{sn}(k)$ ,  $i = n, p, np$  are shown [2]. In the right figure  $S^e(\omega, k)$  are presented for  $k/p_0 = 0.6$  and  $k/p_0 = 0.2$  ( $p_0 = 0.268 \text{ GeV}$ ). For  $k/p_0 = 0.6$  the structure functions for the different processes  $S(\omega, k)_i$ ,  $i = n, p, np$  are presented (the numbers 1, 2, 3, correspondingly). As it was expected the form of the structure function is decomposed over the contributions of the definite processes, corresponding to  $\omega_{si}(k)$ . The widths of maxima (*right*) are determined by the imaginary parts of  $\omega_{si}$  (*left*).

## References

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