

Structure of Light Λ -hypernuclei near Nucleon Drip Lines and Baryonic Interaction

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Hypernuclei

Visible baryonic matter as we know it is primarily made of neutrons and protons, made in turn from u- and d- quarks. To study baryonic interaction properties for heavier quarks (such as s-quarks), one may consider synthesizing hypernuclei, a.k.a. nuclei with protons, neutrons and hyperons.

Notation:

${}^A_{\Lambda}Z$ is a hypernucleus made of N neutrons, Z protons and 1 Λ -hyperon; mass number $A = N + Z + 1$

${}^A_{\Lambda\Lambda}Z$ is a hypernucleus made of N neutrons, Z protons and 2 Λ -hyperons; mass number $A = N + Z + 2$

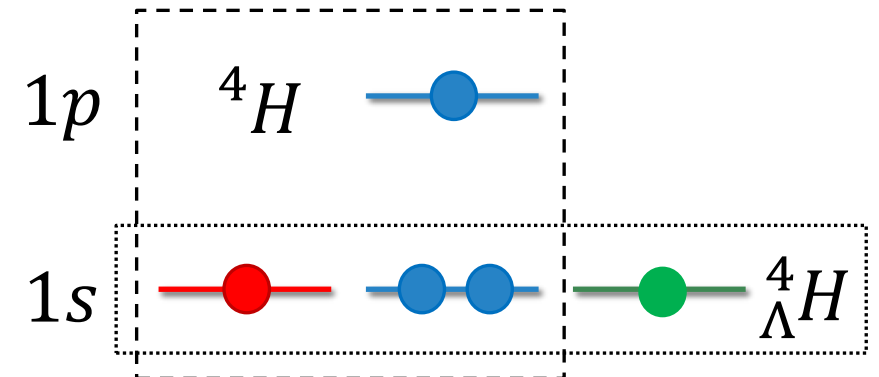
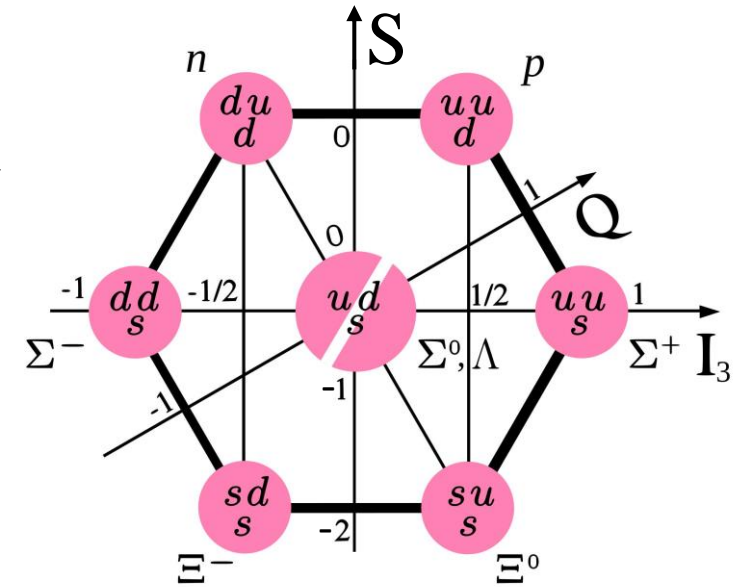
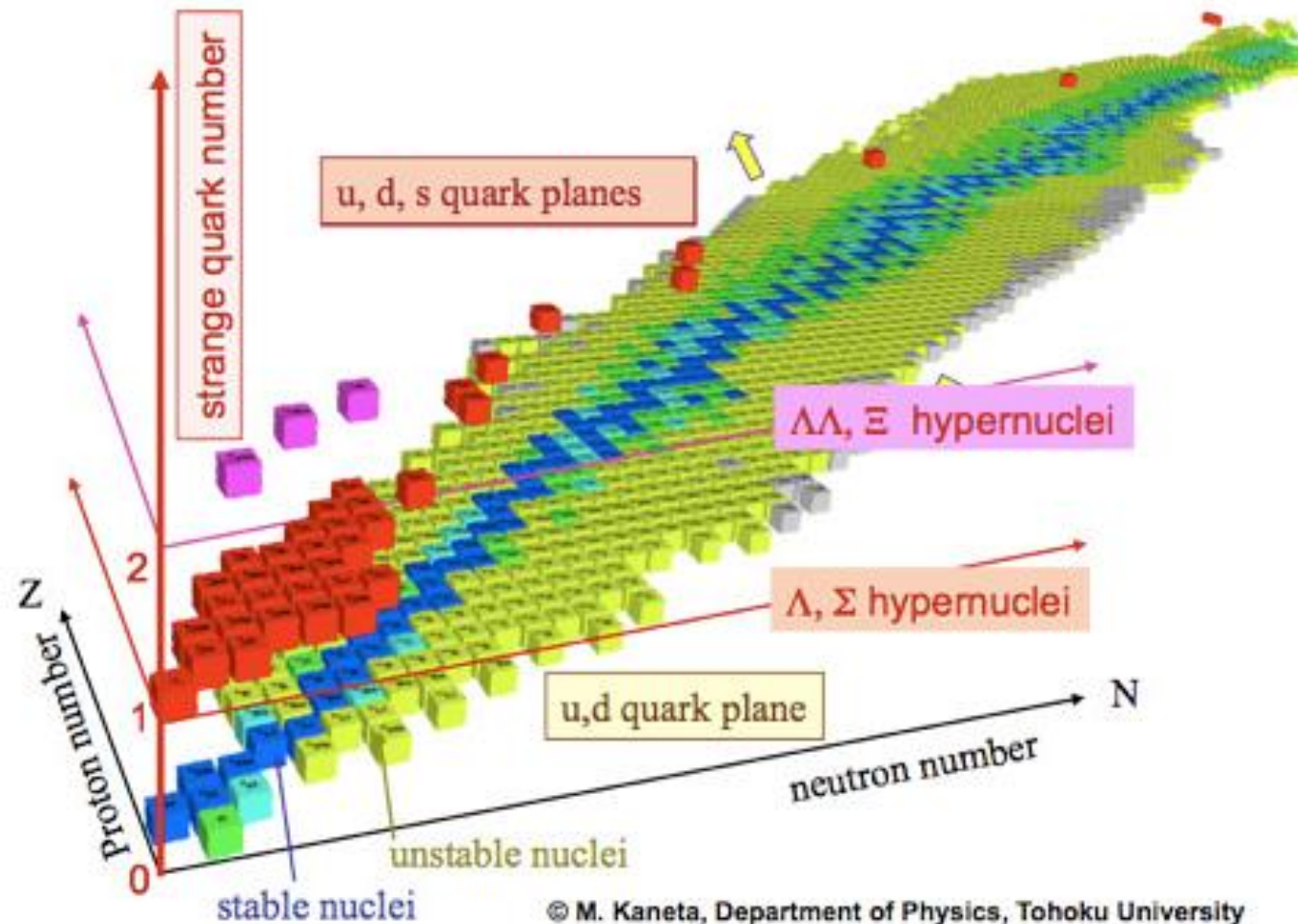


Chart of nuclei and hypernuclei



Motivation

- Study of the properties of baryonic interactions beyond nucleon-nucleon interactions
- Predictions for bound hypernuclei with an unbound nucleon core
- Research in neutron star structure

Hartree-Fock approach for hypernuclei

- Energy density functional:

$$E = \langle g.s. | T + V_{12} | g.s. \rangle = \int H(\rho, \tau, J) dr, \quad |g.s.\rangle = \frac{1}{\sqrt{A!}} \det|\phi_i(\mathbf{r}_j)|$$

- Hartree-Fock single particle states:

$$\frac{\delta}{\delta\phi_i} (E - \sum_i e_i \int |\phi_i(\mathbf{r})|^2 dr) = 0$$

- Hartree-Fock equations:

$$\frac{\hbar^2}{2m_{q,\Lambda}^*(r)} \left[-R_\alpha''(r) + \frac{l_\alpha(l_\alpha + 1)}{r^2} R_\alpha(r) \right] - \left(\frac{\hbar^2}{2m_{q,\Lambda}^*(r)} \right)' R_\alpha'(r) + U_{q,\Lambda}(r) R_\alpha(r) = e_\alpha R_\alpha(r)$$

Skyrme interaction for hypernuclei

- Nucleon-nucleon Skyrme potential:

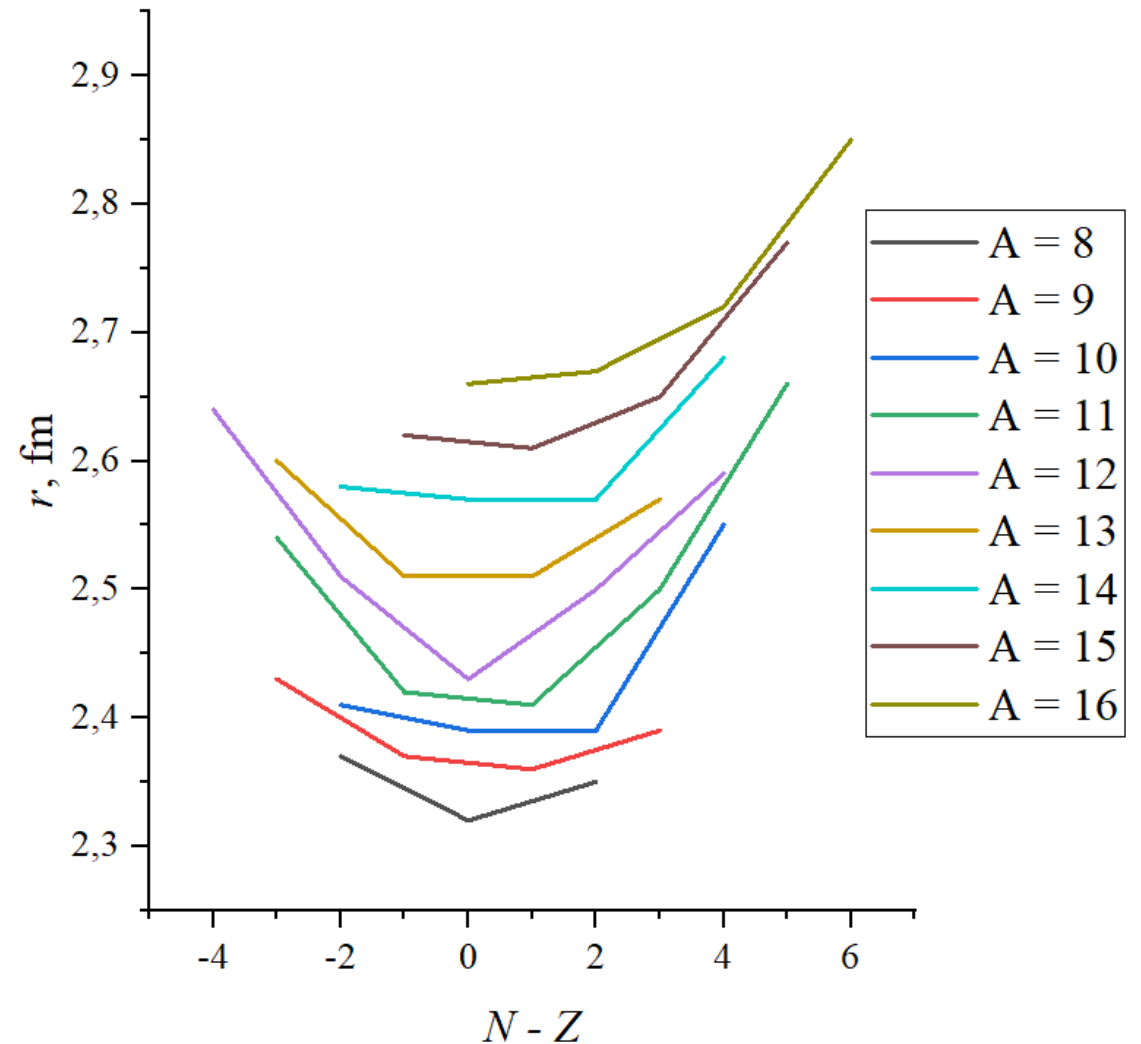
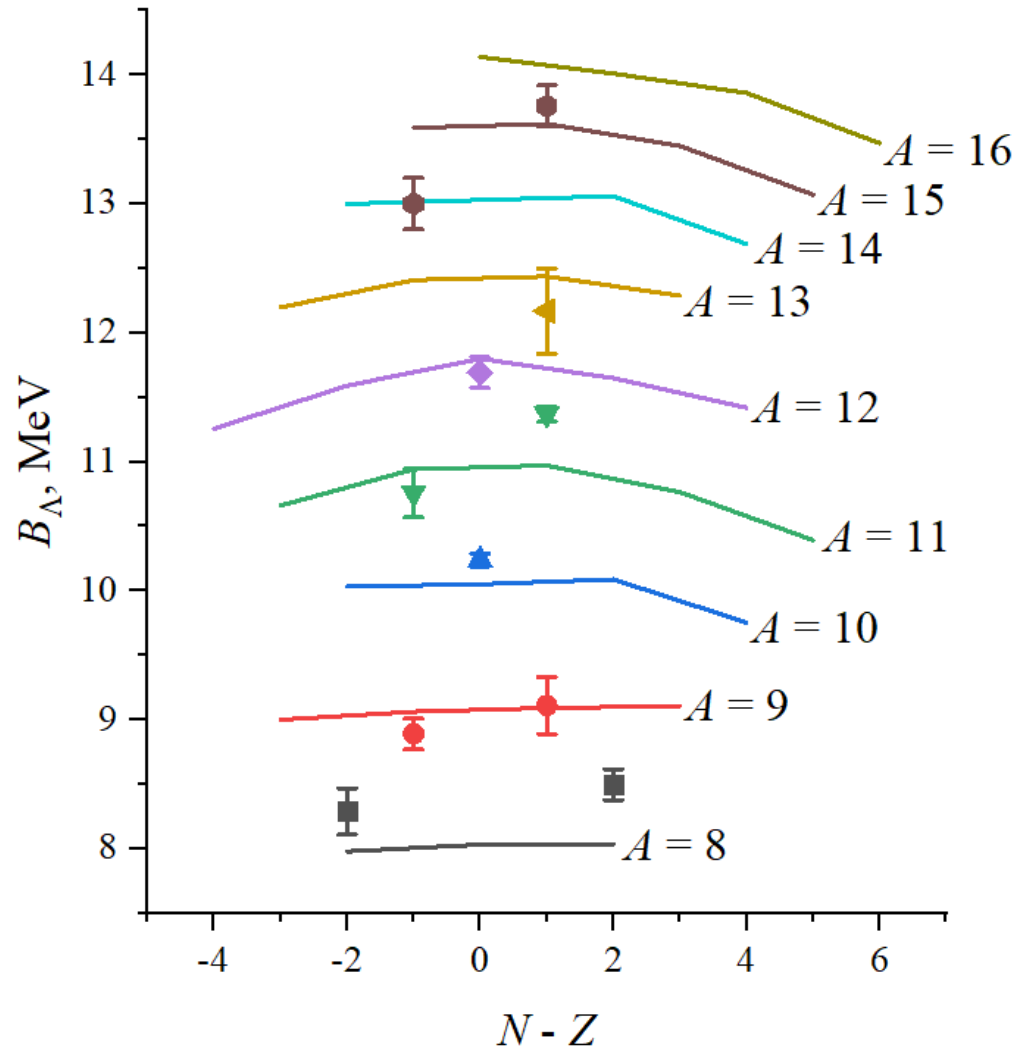
$$V_{NN}(\mathbf{r}_1, \mathbf{r}_2) = t_0(1 + x_0 P_\sigma)\delta(\mathbf{r}_{12}) + \frac{1}{2}t_1(1 + x_1 P_\sigma)(\mathbf{k}^2\delta(\mathbf{r}_{12}) + \delta(\mathbf{r}_{12})\mathbf{k}'^2) \\ + t_2(1 + x_2 P_\sigma)\mathbf{k}'\delta(\mathbf{r}_{12})\mathbf{k} + \frac{1}{6}t_3\rho^\alpha(\mathbf{R})(1 + x_3 P_\sigma)\delta(\mathbf{r}_{12}) + iW(\sigma_1 + \sigma_2)[\mathbf{k}' \times \delta(\mathbf{r})\mathbf{k}]$$

- Hyperon-nucleon Skyrme potential:

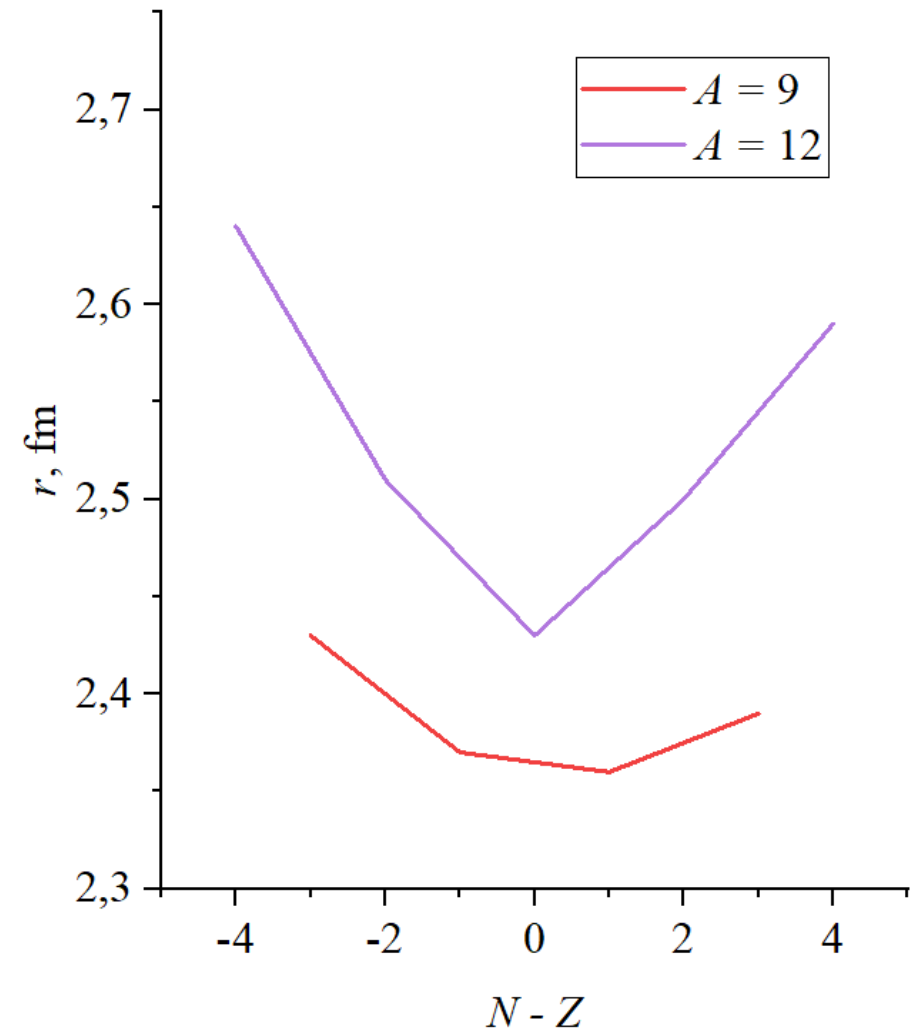
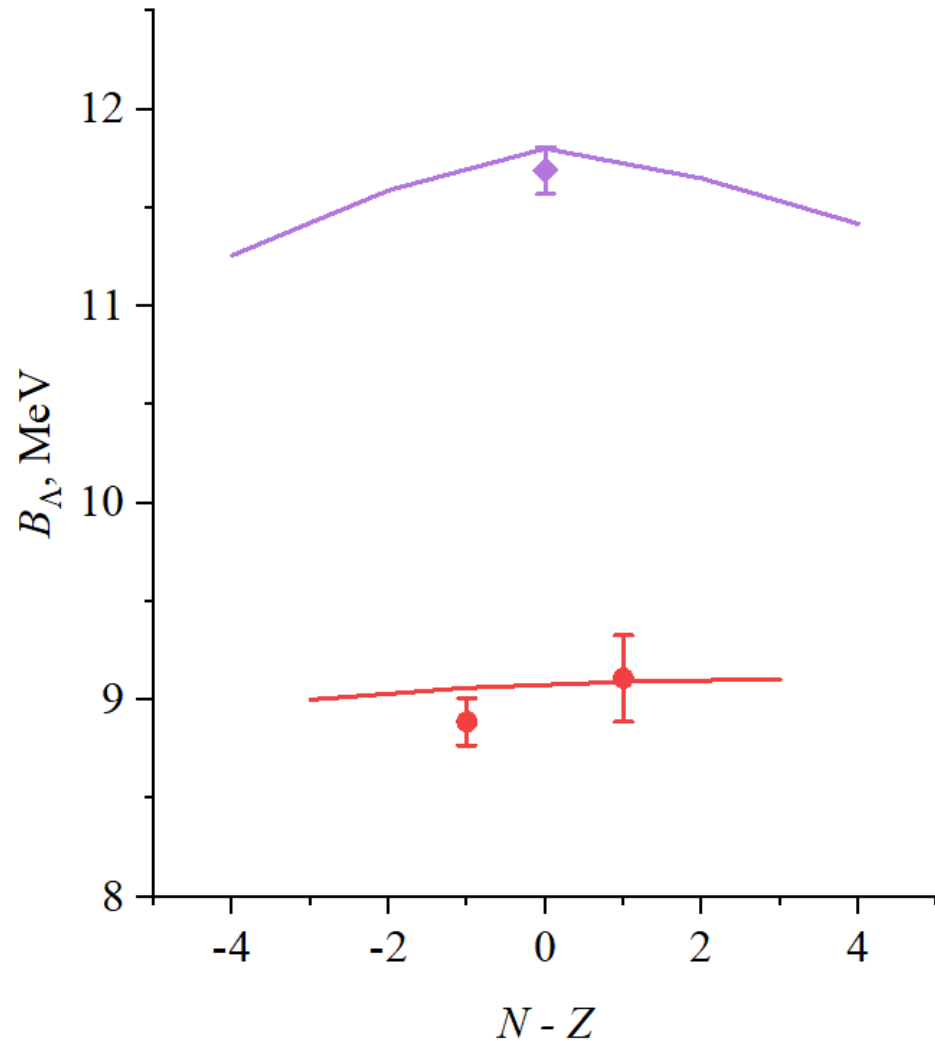
$$V_{\Lambda N}(\mathbf{r}_\Lambda, \mathbf{r}_q) = t_0^\Lambda(1 + x_0^\Lambda P_\sigma)\delta(\mathbf{r}_{\Lambda q}) + \frac{1}{2}t_1^\Lambda(\mathbf{k}^2\delta(\mathbf{r}_{\Lambda q}) + \delta(\mathbf{r}_{\Lambda q})\mathbf{k}'^2) \\ + t_2^\Lambda\mathbf{k}'\delta(\mathbf{r}_{\Lambda q})\mathbf{k} + \frac{1}{6}t_3^\Lambda\rho^\alpha(\mathbf{R})\delta(\mathbf{r}_{\Lambda q})$$

- NN: SLy4
- Λ N: SLL4', YBZ1, YBZ5, LY1, YMR, SkSH1

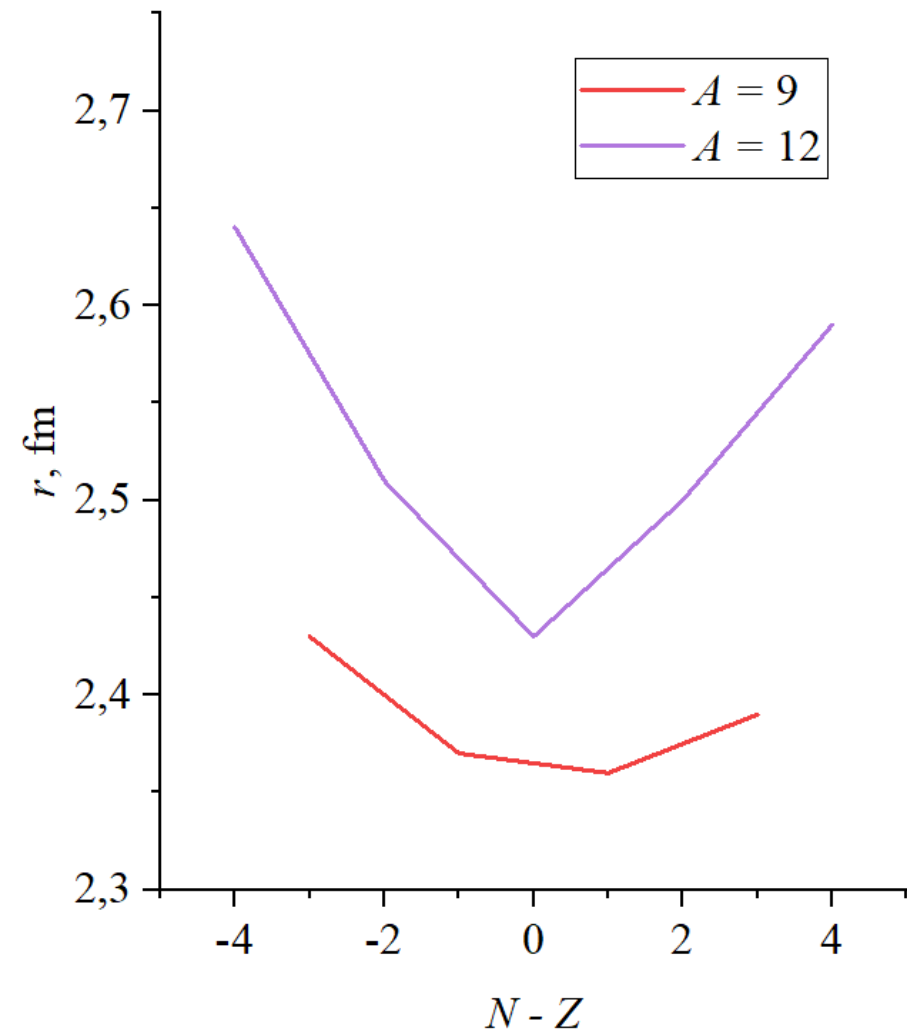
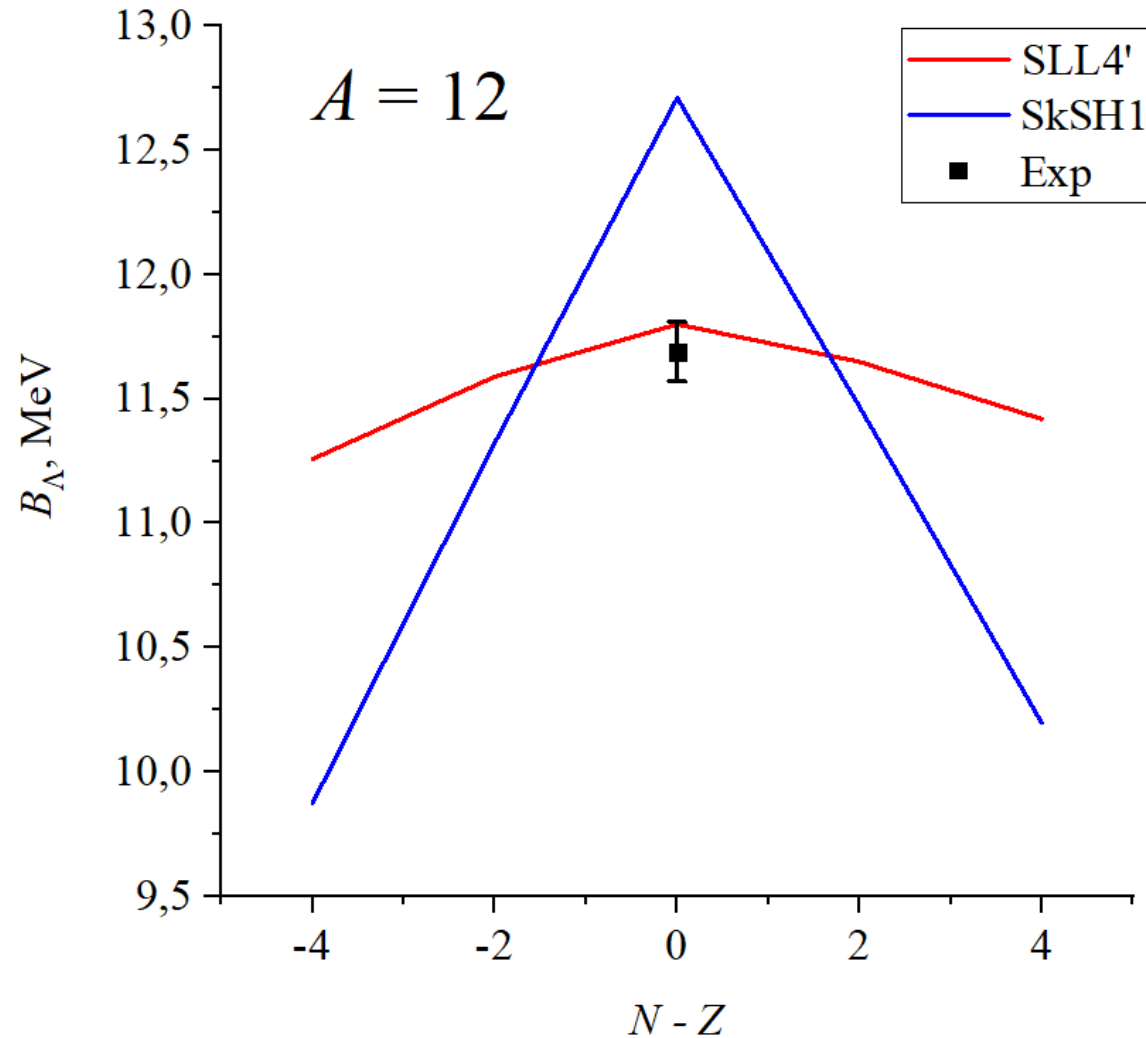
Hyperon binding energies and radii of nuclear cores in $A+1_{\Lambda}Z$



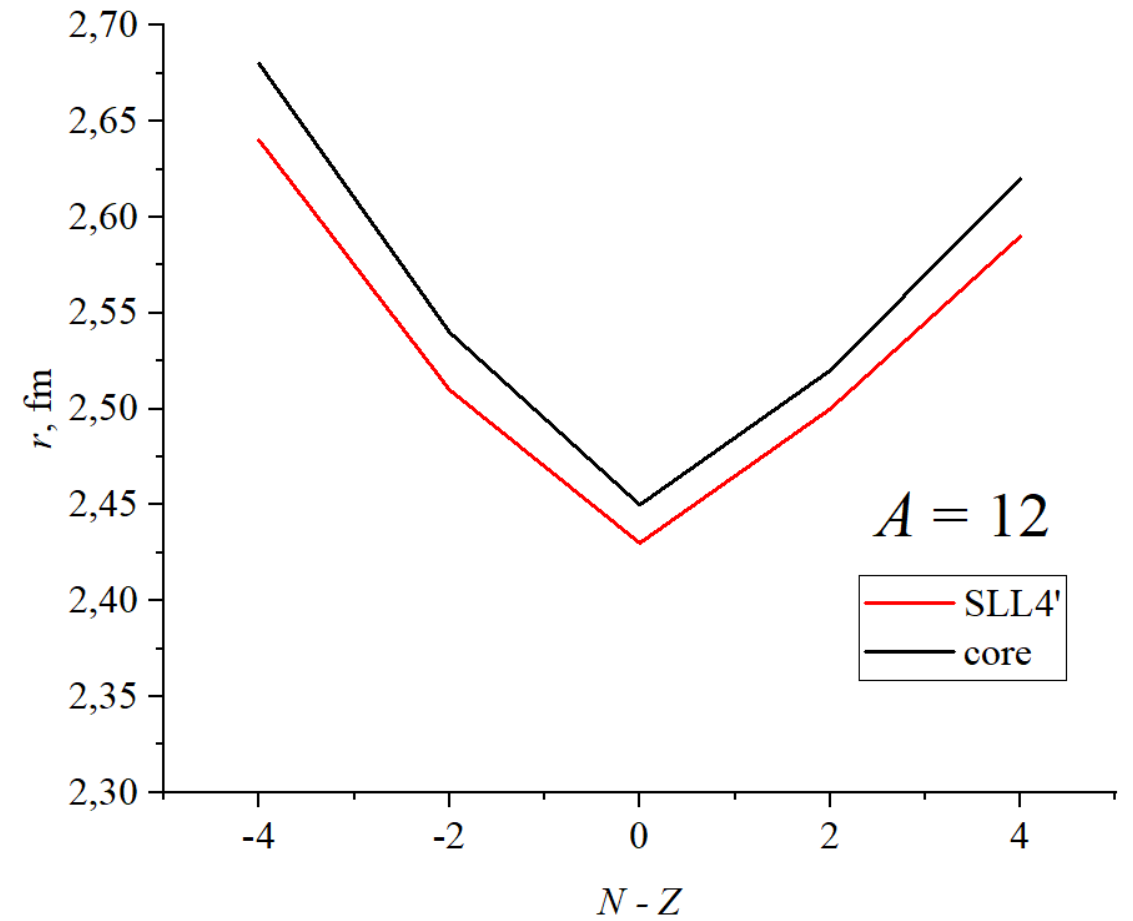
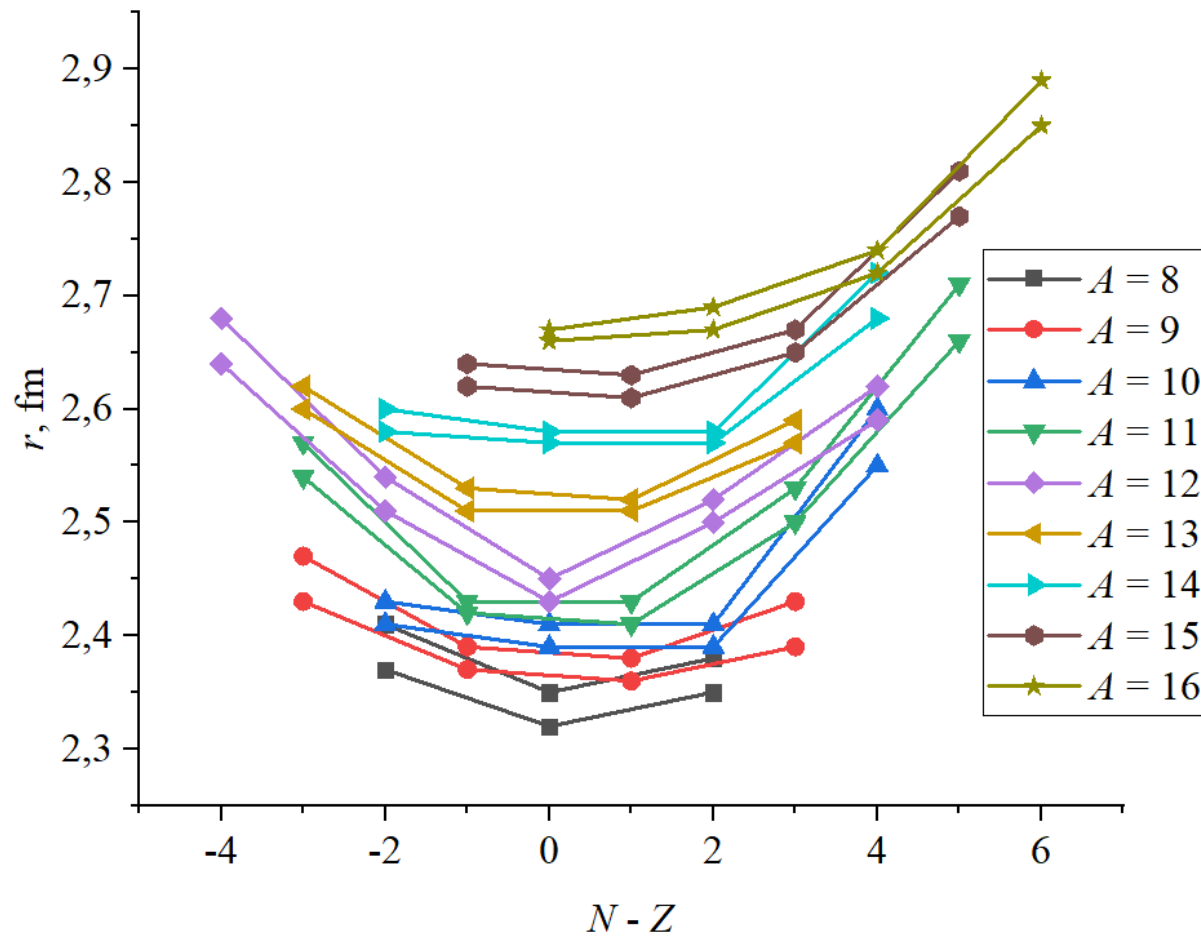
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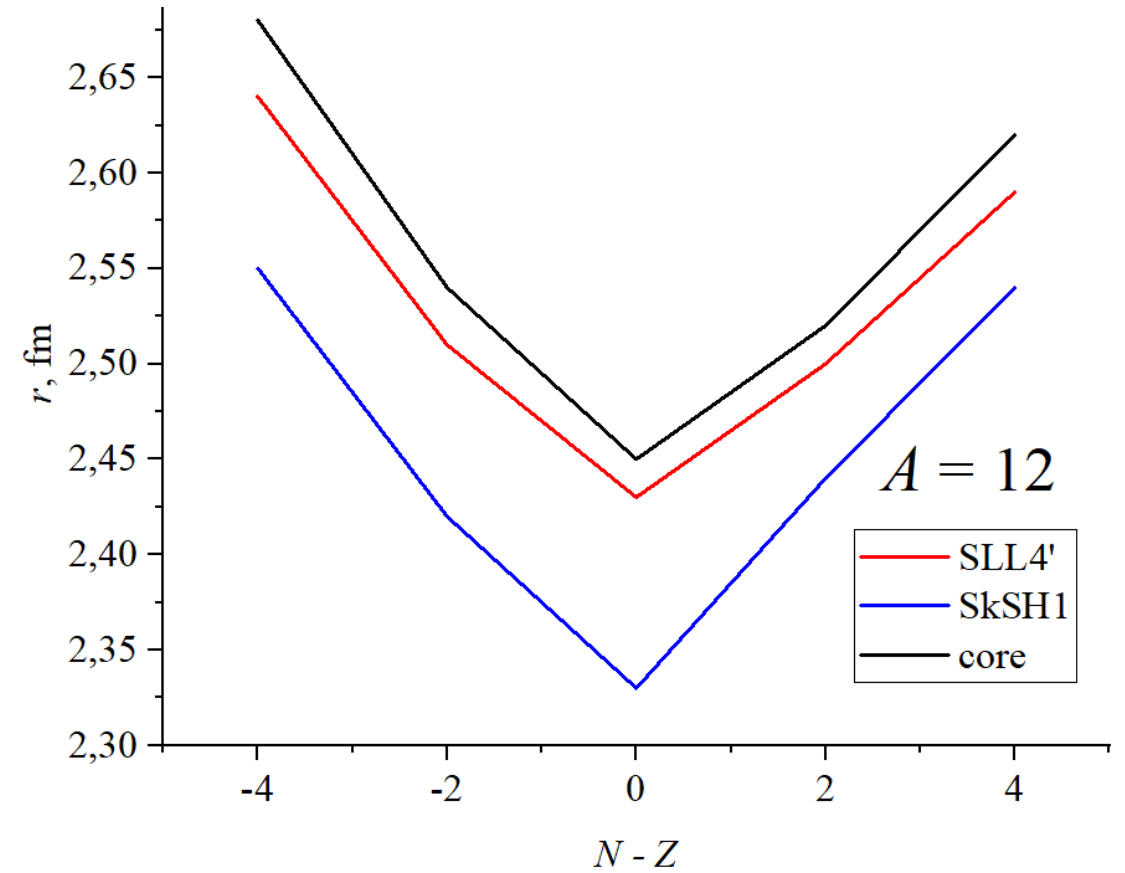
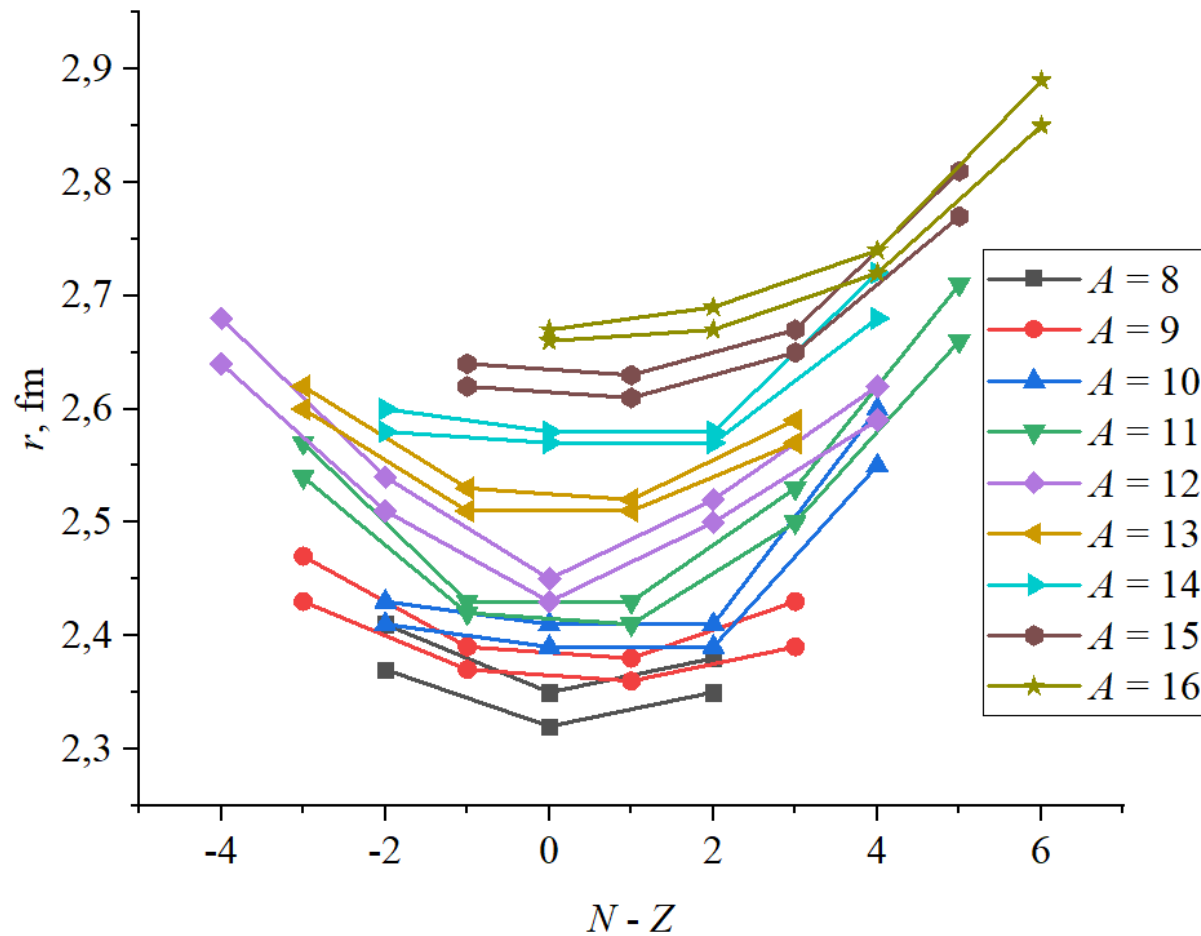
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Nuclear core distortion (polarization) by hyperon



Nuclear core distortion (polarization) by hyperon



Conclusions

Hypernuclear Hartree-Fock approach was utilized to study the properties of light neutron-rich Λ -hypernuclei.

- Hyperon binding energy depends on excess of neutrons or protons in the nucleus (isospin)
- Various choices of Λ N-interaction can lead to significantly different predictions of B_Λ
- Hyperon binding energy vs isospin dependence is more pronounced for the cases with more dramatic density changes, and the corresponding nuclei should be chosen as primary sources of information on Λ N-interaction

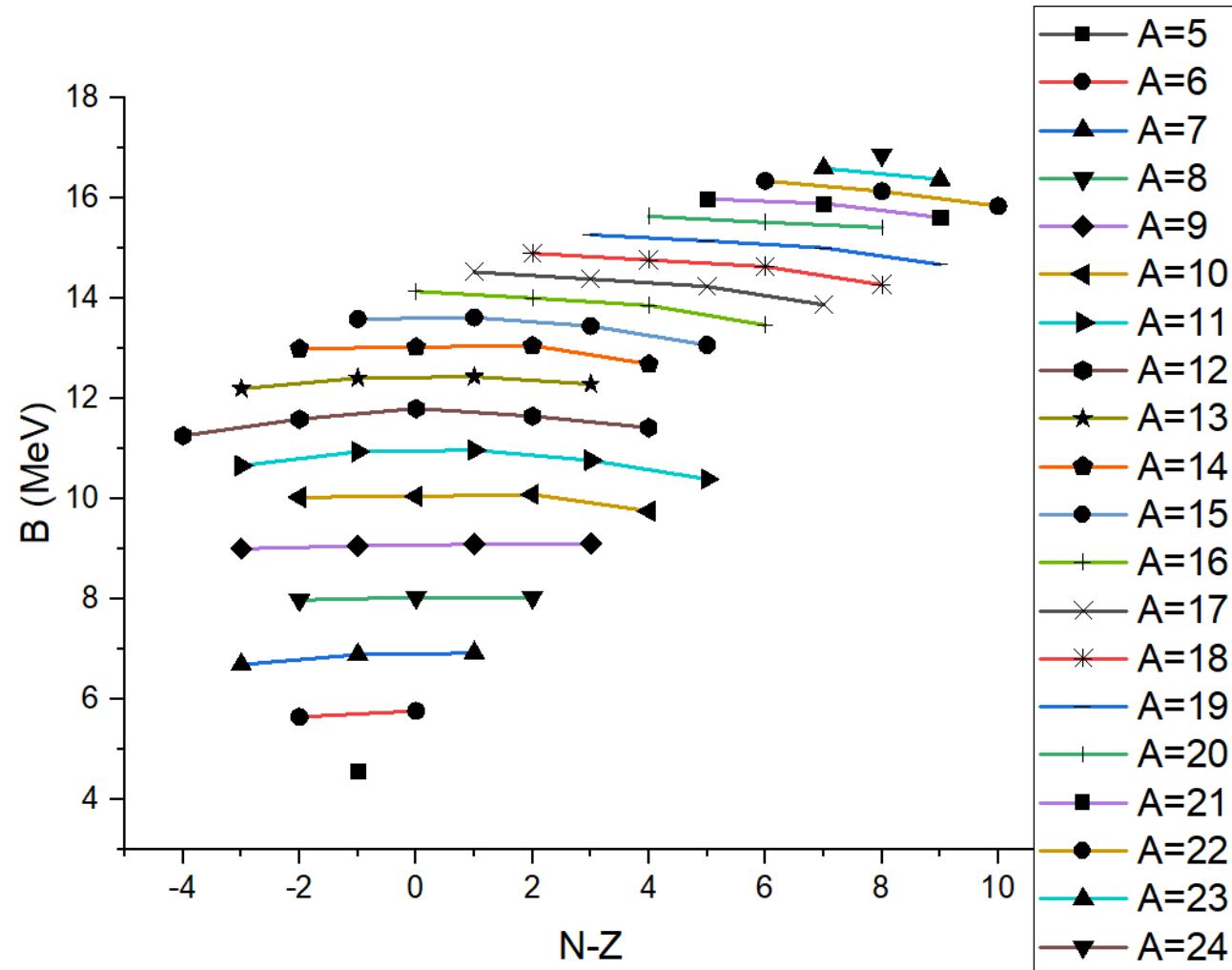
It is planned to improve the predictive power of the approach by taking charge symmetry breaking into consideration.

Thank you for your attention

Back-up slides



Hyperon binding energy



$$B_{\Lambda}({}_{\Lambda}^AZ) = B_{tot}({}_{\Lambda}^AZ) - B_{tot}({}^{A-1}Z)$$

- The difference in neighboring isobar chains is around 1 MeV for lighter hypernuclei, smaller as A increases
- Symmetric character of B_{Λ} with respect to isospin $N - Z$
- B_{Λ} is almost constant for nuclei in the same isobar chain

Hypernuclear radii

