

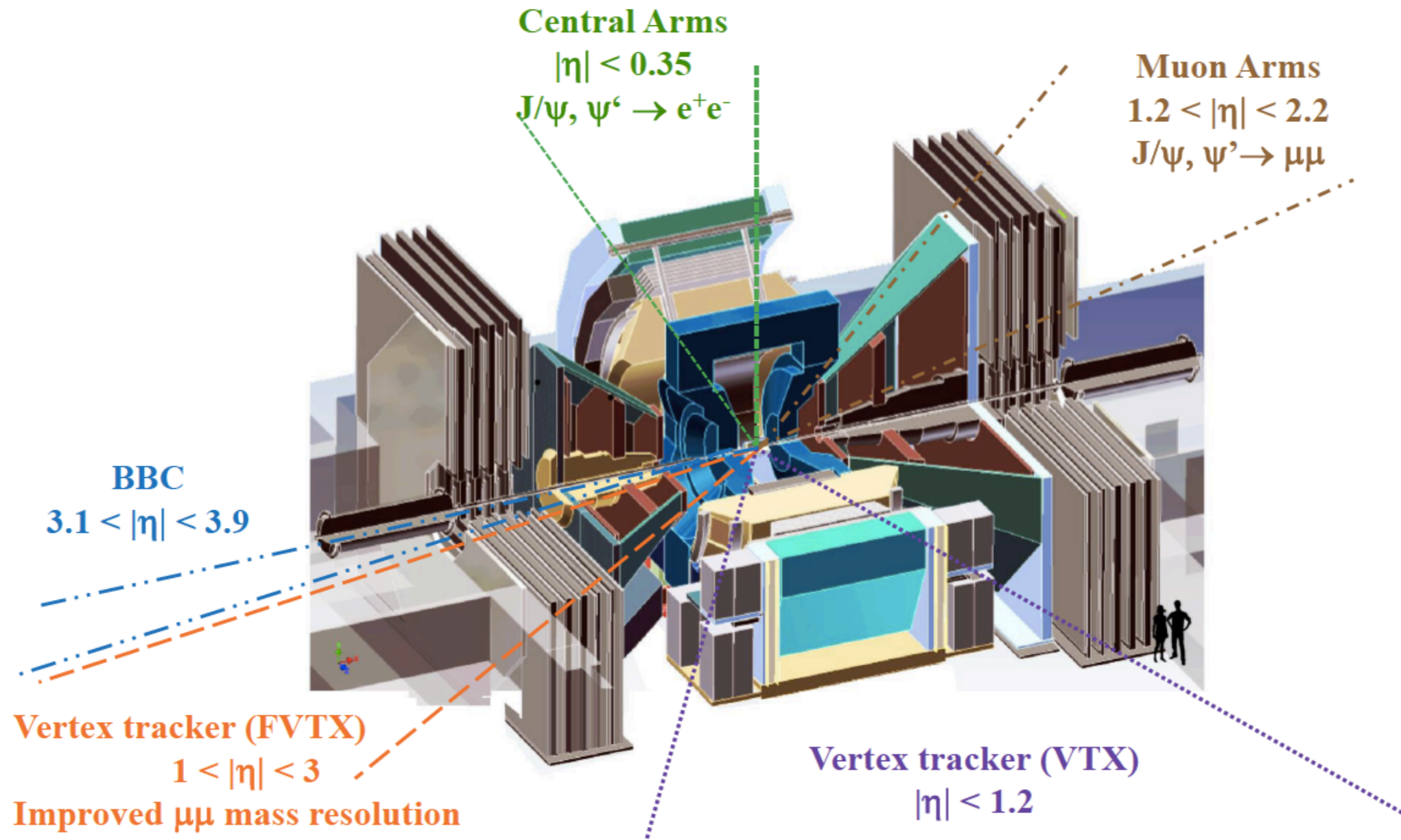
PHENIX Results on J/ψ Production in $p+Al$, $p+Au$ and ^3He+Au Collisions

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Quarkonia - PHENIX



Di-leptons measured in $y = (-2.2, -1.2) (-0.35, 0.35) (1.2, 2.2)$

- No triggering required in Au+Au
- Efficient triggers in p+p, (p,d, ^3He)+A

Introduction

Recent quarkonium analyses in PHENIX have focused on small systems results for J/ψ and $\psi(2S)$.

$p+Al$, $p+Au$ and ^3He+Au data from the 2014 and 2015 RHIC runs have been added to our $d+Au$ data from 2008.

This allows us to address two questions about J/ψ production:

- Do we see evidence of differences between $p+Au$ and ^3He+Au
 - i.e. evidence of final state effects?
- How well do we understand J/ψ modification in light systems?

See [PHENIX, arXiv:1910.14487](#) for a description of the measurement.

Introduction

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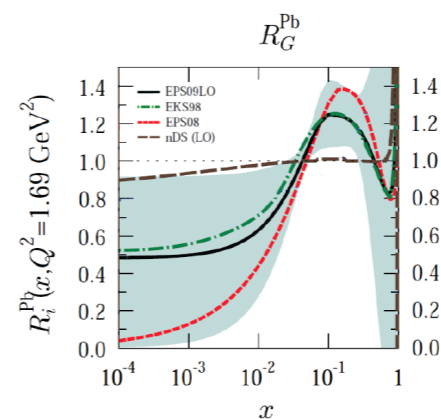
See [PHENIX, arXiv:1910.14487](#) for a description of the measurement.

We need to start by briefly reviewing effects that modify quarkonia production in a nuclear target.

Quarkonia production in a nucleus

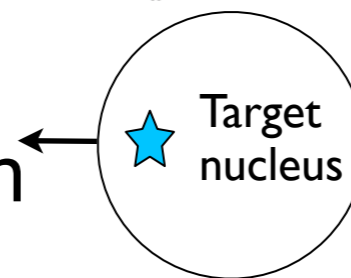
Processes that modify the quarkonia yield in a nuclear target - called **cold nuclear matter** (CNM) processes.

Gluon shadowing - parton distributions are modified in a nucleus



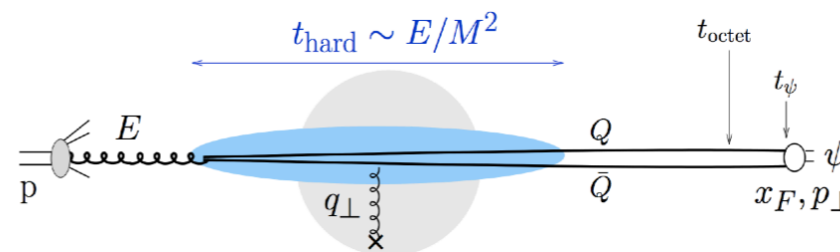
Affects underlying heavy quark yield

Absorption - breakup of the precursor quarkonium by collision with a target nucleon



Breaks up forming mesons

Initial state energy loss of a parton in cold nuclear matter



Changes rapidity distribution

Cronin effect - multiple elastic scattering of partons

Modifies the p_T distribution

There is also a possibility that quarkonium states may be broken up in the final state by interactions with particles produced in the collision.

Breaks up bound mesons

Shadowing

Recent shadowing parameterizations

- [EPPS16](#) (Eskola et. al., Eur. Phys. J. C 77, 163 (2017))
- [nCTEQ15](#) (Kovarik et. al., Phys. Rev. D 93, 085037 (2016))

Bayesian re-weighting of EPPS16 and nCTEQ15 gluon nPDF's

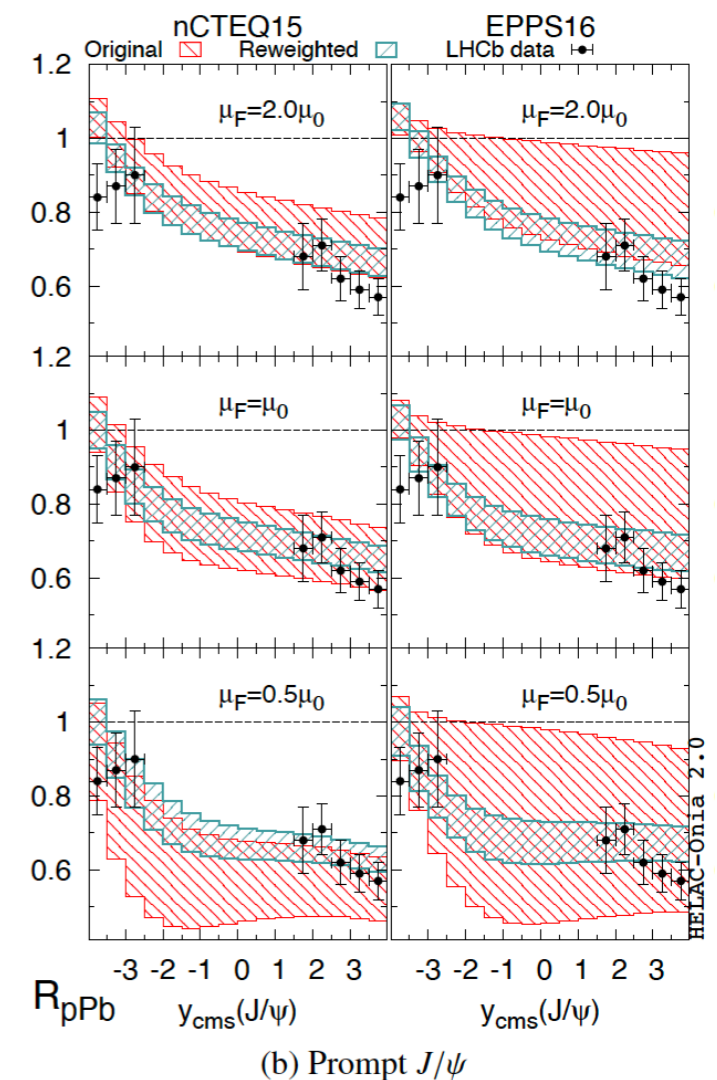
- (Kusina et. al., Phys. Rev. Lett. 121, 052004 (2018))
- Adds **LHC pPb data** - gluon dominated processes
 - D_0 , J/ψ , $B \rightarrow J/\psi$, and $Y(1S)$ mesons

See also Eskola et. al. arXiv:1906.03943

- Considerably narrows uncertainty band
- Reduces R_g at forward rapidity
- “Absorbs” initial state energy loss into nPDF?

Fitted to centrality integrated data only

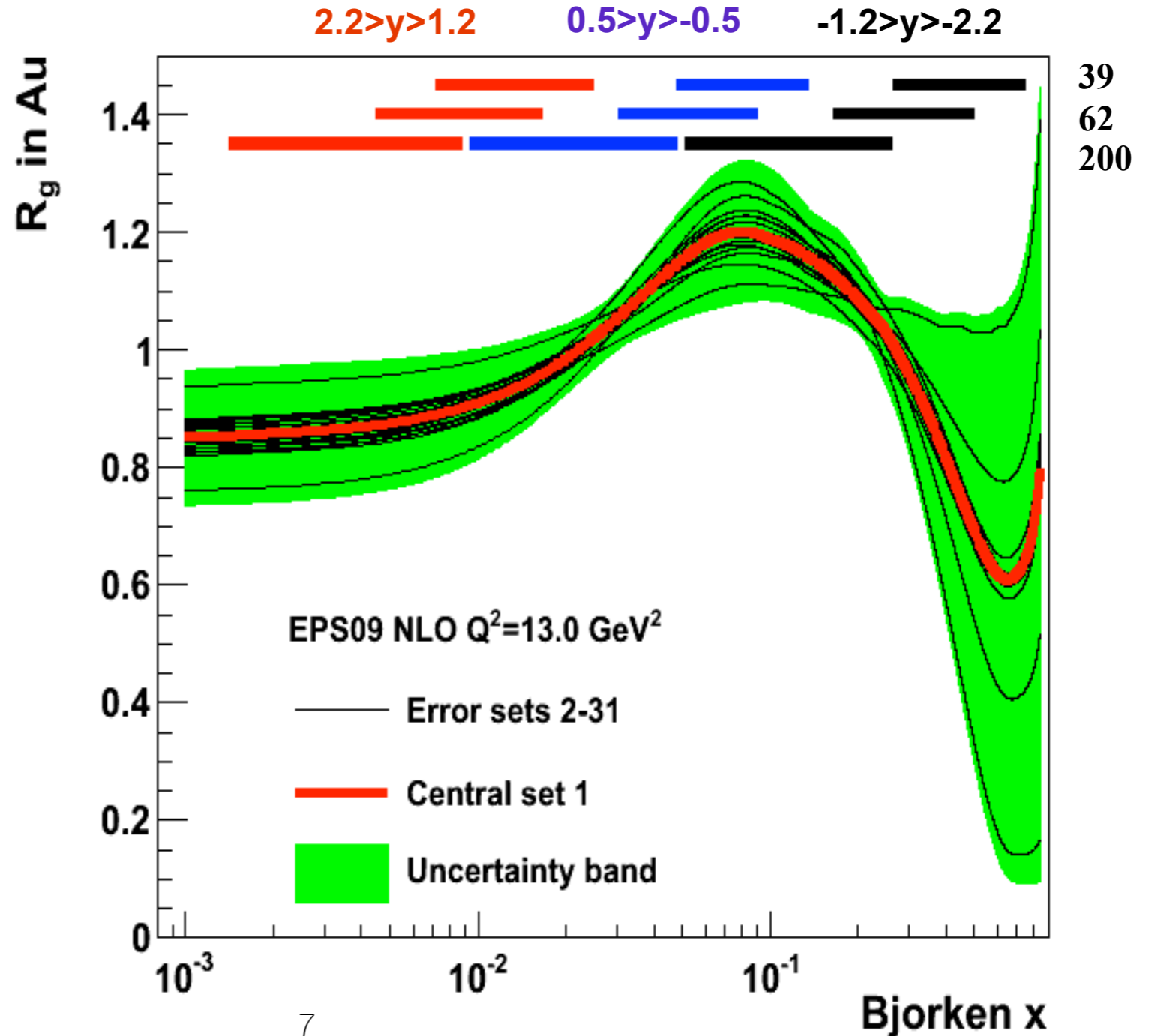
- Has no information about centrality dependence
- **Centrality dependence has to be invented**



x versus rapidity in PHENIX for J/ψ

$$x_2 = \frac{\sqrt{M^2 + p_T^2}}{\sqrt{s_{NN}}} e^{-y}, \quad Q^2 = M^2 + p_T^2$$

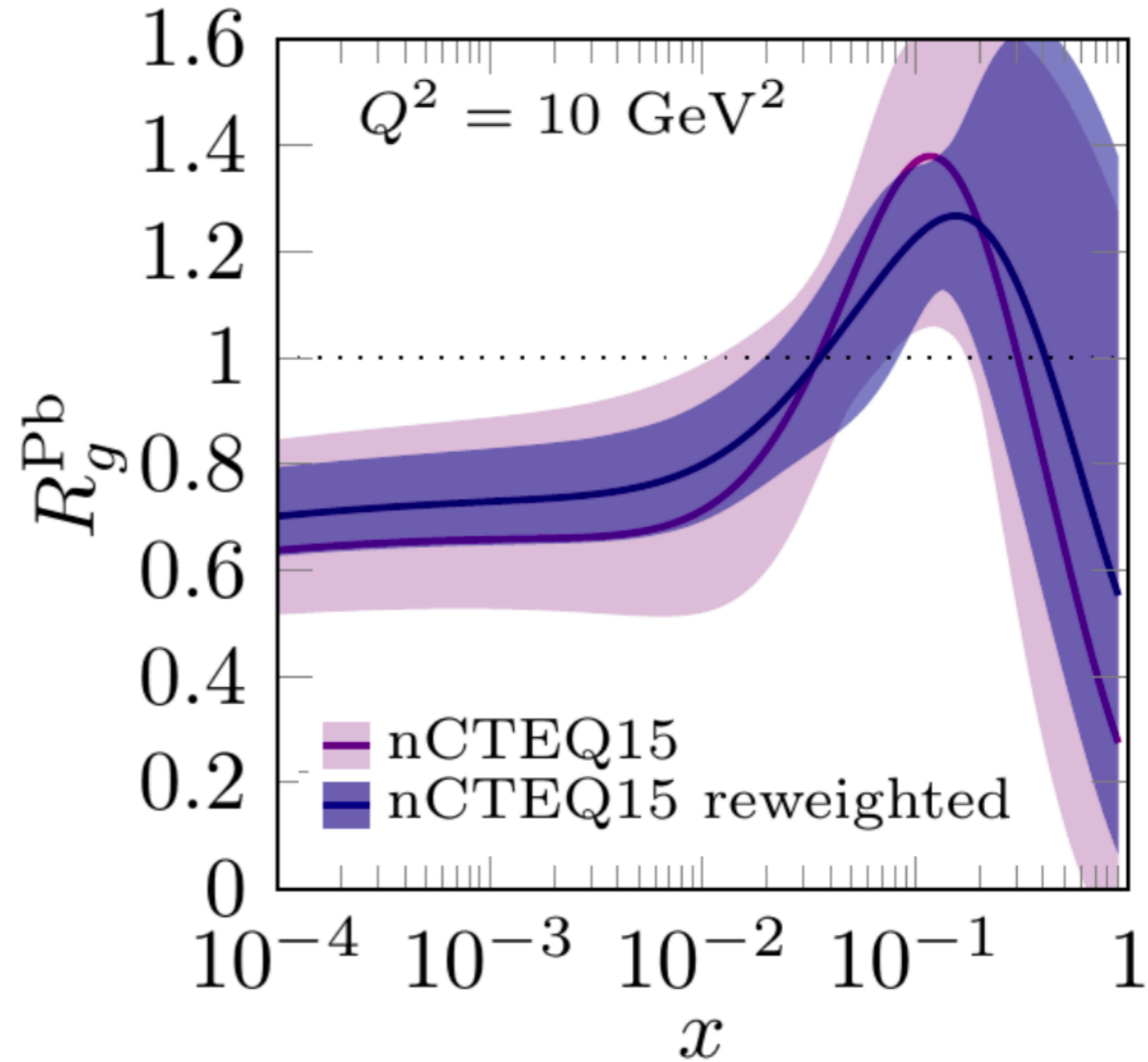
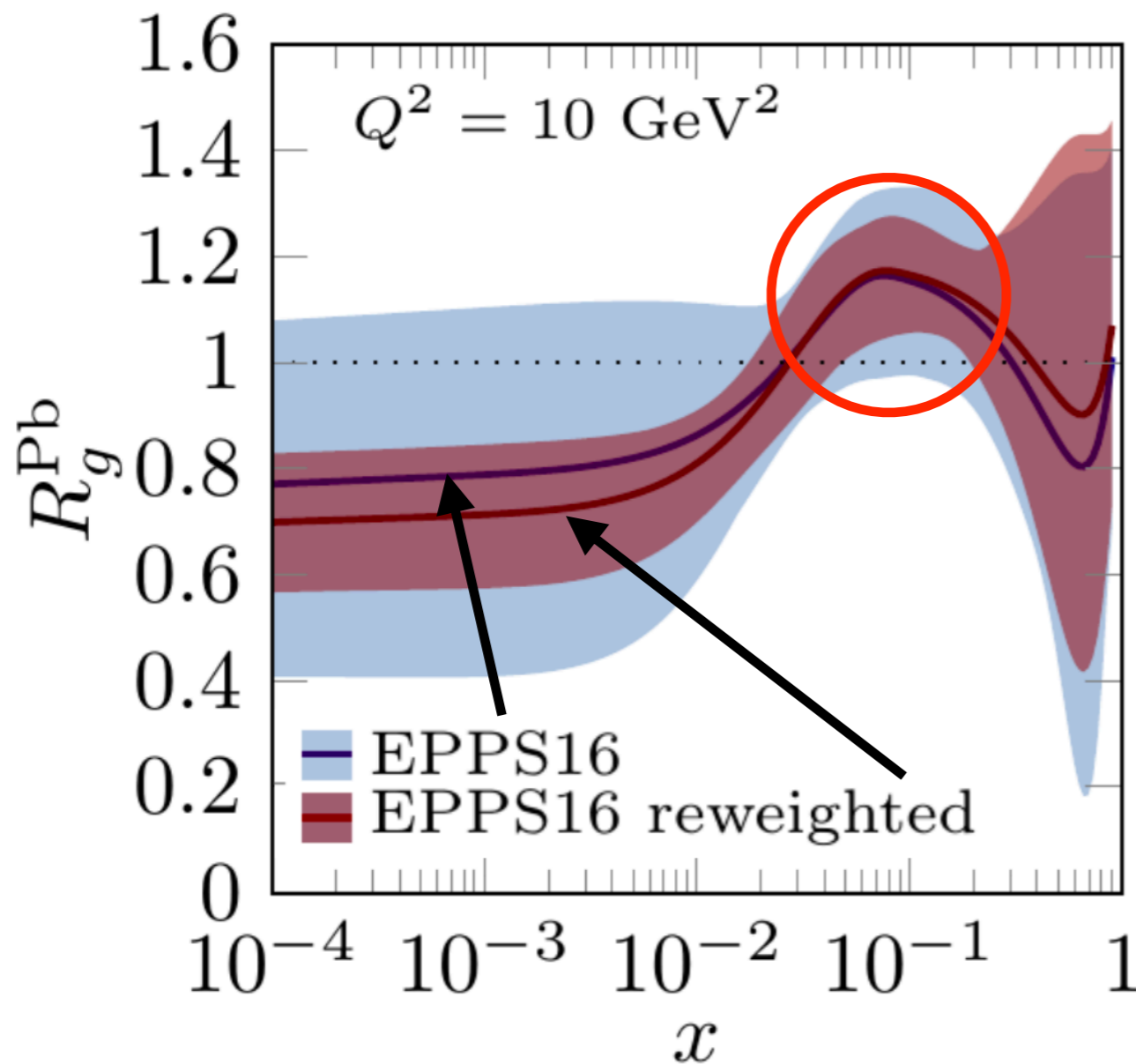
Ranges of x values sampled by J/ψ production in PHENIX.



Re-weighting of gluon nPDF's

Eskola et. al., arXiv:1906.03943.

Hessian re-weighting of nPDF's using LHCb D^0 data.



d+Au J/ ψ and open HF results vs rapidity

PHENIX, PRL 112 252301 (2014)

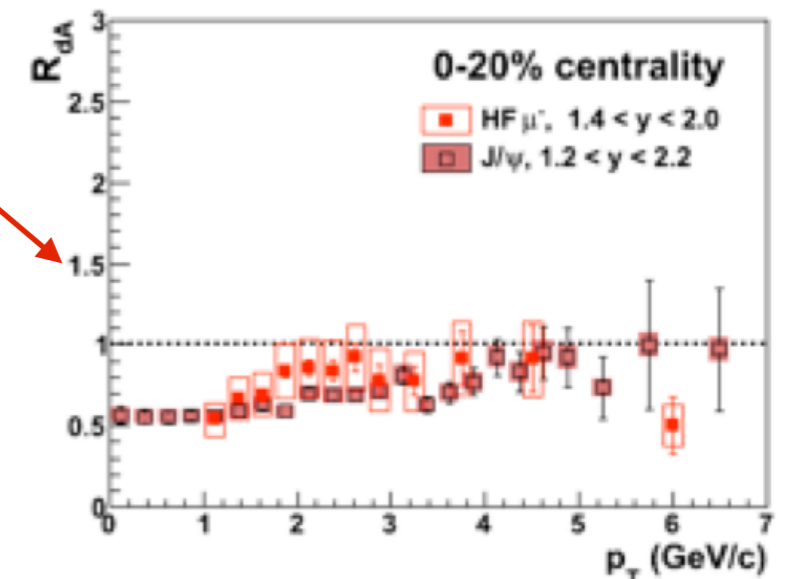
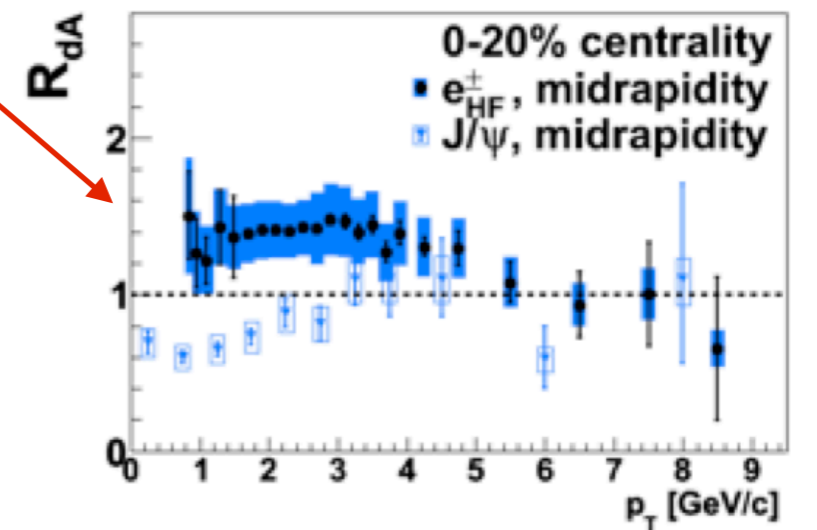
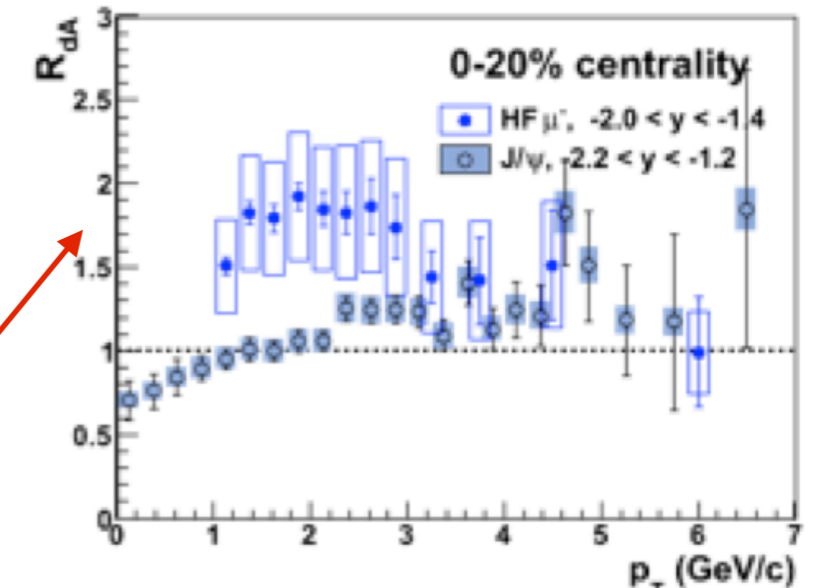
Caveat: Different kinematics!

The J/ ψ suppression at **backward & mid rapidity** is much stronger than for HF.

- Implies J/ ψ is suppressed **beyond** the underlying HF production.

At **forward rapidity** they are similar.

- Implies J/ ψ suppressed at forward rapidity **because** the underlying HF is suppressed.



J/ψ absorption

Backward rapidity J/ψ in PHENIX experience a significant “absorption” cross section from breakup due to collisions with target nucleons

- in addition to substantial anti-shadowing.

Parameterized using model of expansion of color neutral charmonium precursor as it crosses the target (Arleo et. al., PRC 61, 054906 (2000)).

- Fitted to world’s σ_{abs} data for nuclear crossing time $\tau > 0.05 \text{ fm/c}$
- All data **corrected for shadowing** with EKS98 or EPS09

Provides good description of data from $\sqrt{s_{\text{NN}}} = 17 \text{ to } 200 \text{ GeV}$

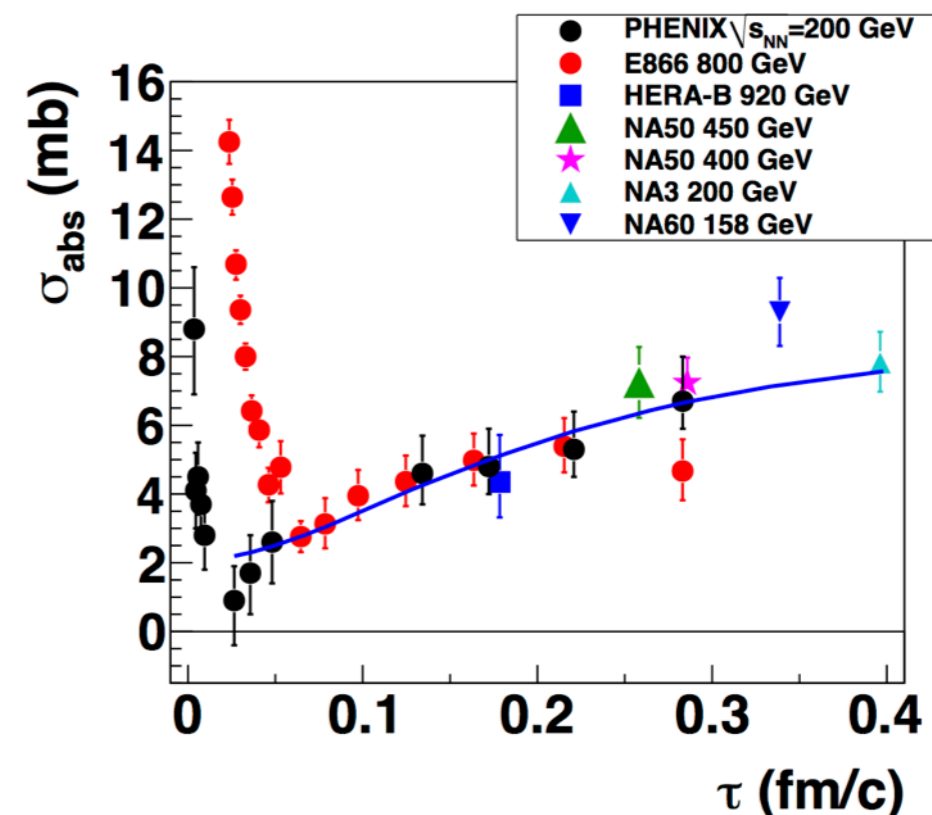
Anti-shadowing parameterizations have remained stable over several generations.

- Well constrained by DIS data.

Not expected to be important at LHC

- Nuclear crossing times very short at all y

Phys.Rev. C87 (2013) 5, 054910



Initial state energy loss

Incoming parton energy loses energy due to gluon radiation associated with p_T broadening.

Examples:

- Arleo et. al. JHEP 05 (2013) 155.
- Sharma and Vitev, PRC 87 (2013) 044905.
- Kopeliovich et al., Phys.Rev. C95 (2017) 065203.

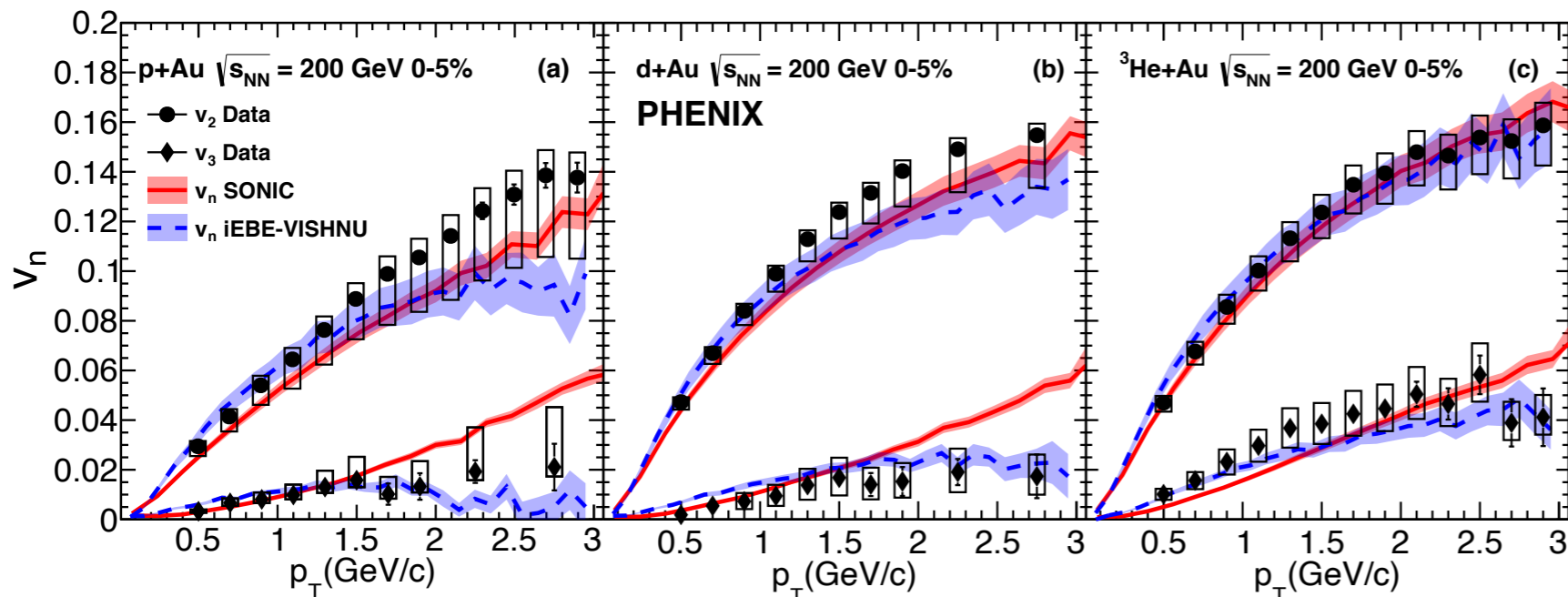
The Bayesian re-weighted shadowing seems to explain p+A data reasonably well without additional effects from initial state energy loss.

- Absorbs initial state energy loss effects into the shadowing parameterization?

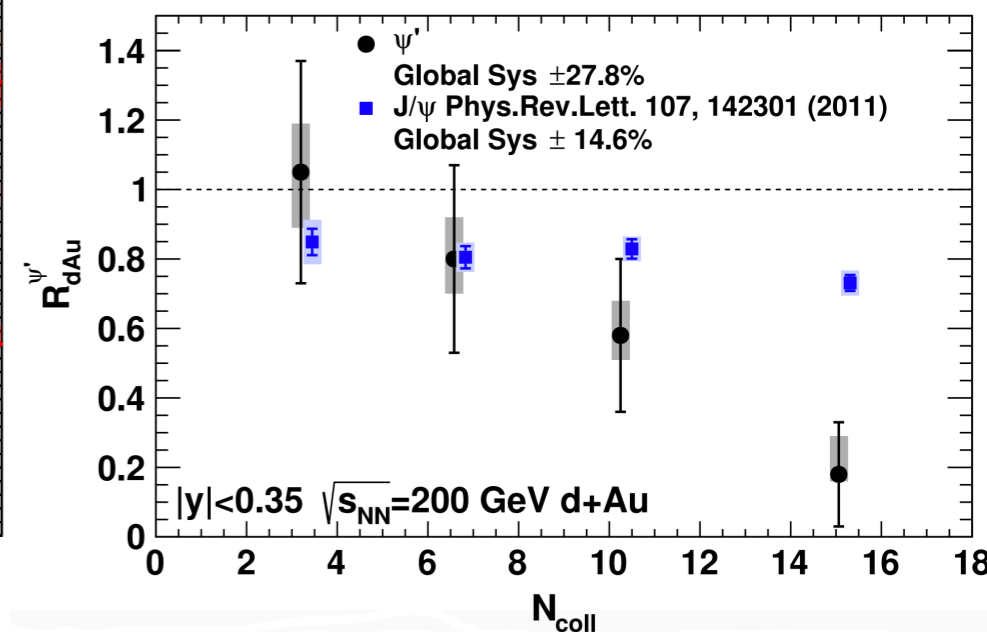
Motivation for studying multiple light systems

- Interest in p+A partly motivated by the A+A program
 - p+A is sensitive to initial state effects that are not theoretically well understood.
 - Assumed for a long time **only** initial state effects important in p+A.
- **But:**
 - Unexpectedly strong suppression of the $\psi(2S)$ observed in d+Au collisions at RHIC, and then in p+Pb collisions at LHC.
 - Evidence of flow in small systems observed, first at LHC and then at RHIC.
 - Final state effects on quarkonia production in p+A?

Nature Physics 15, 214-220 (2019)



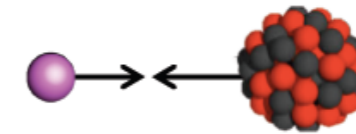
PHENIX, PRL 111 202301 (2013)



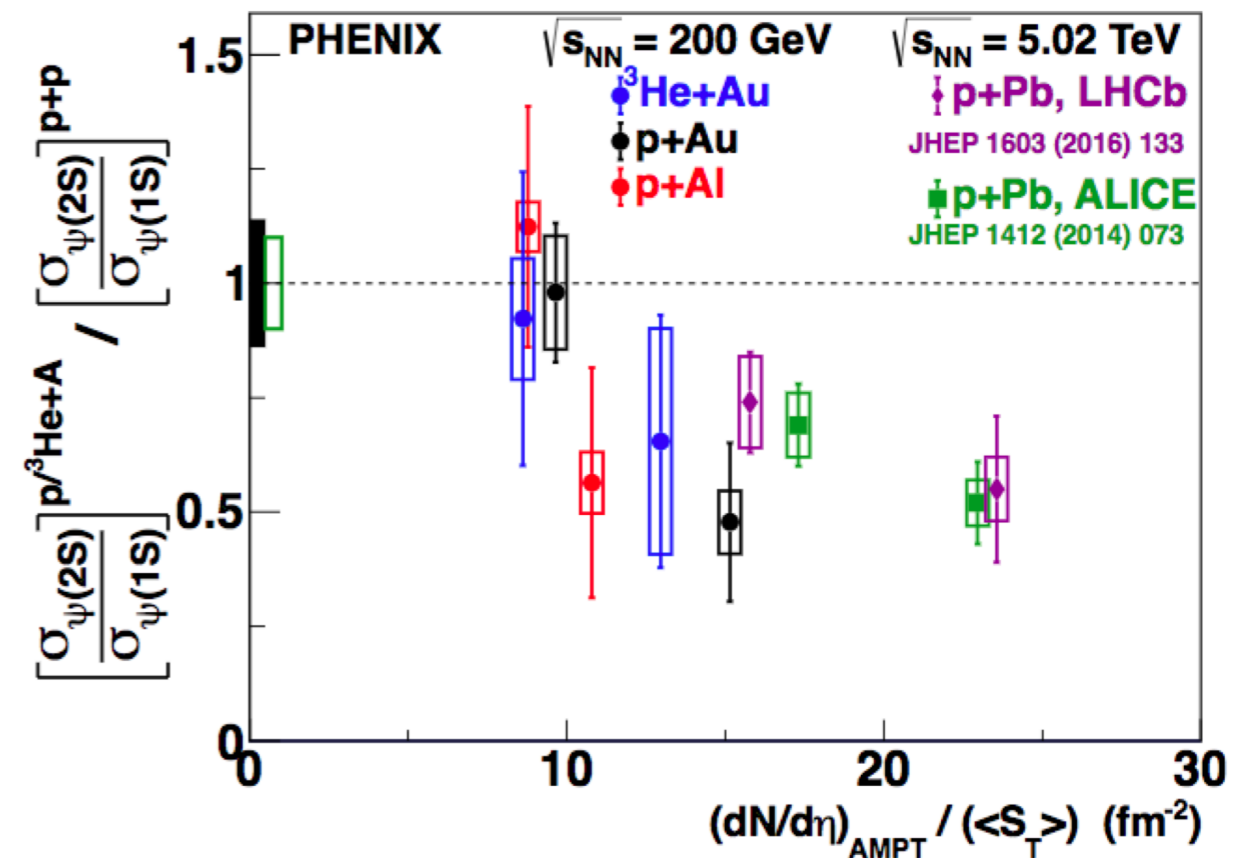
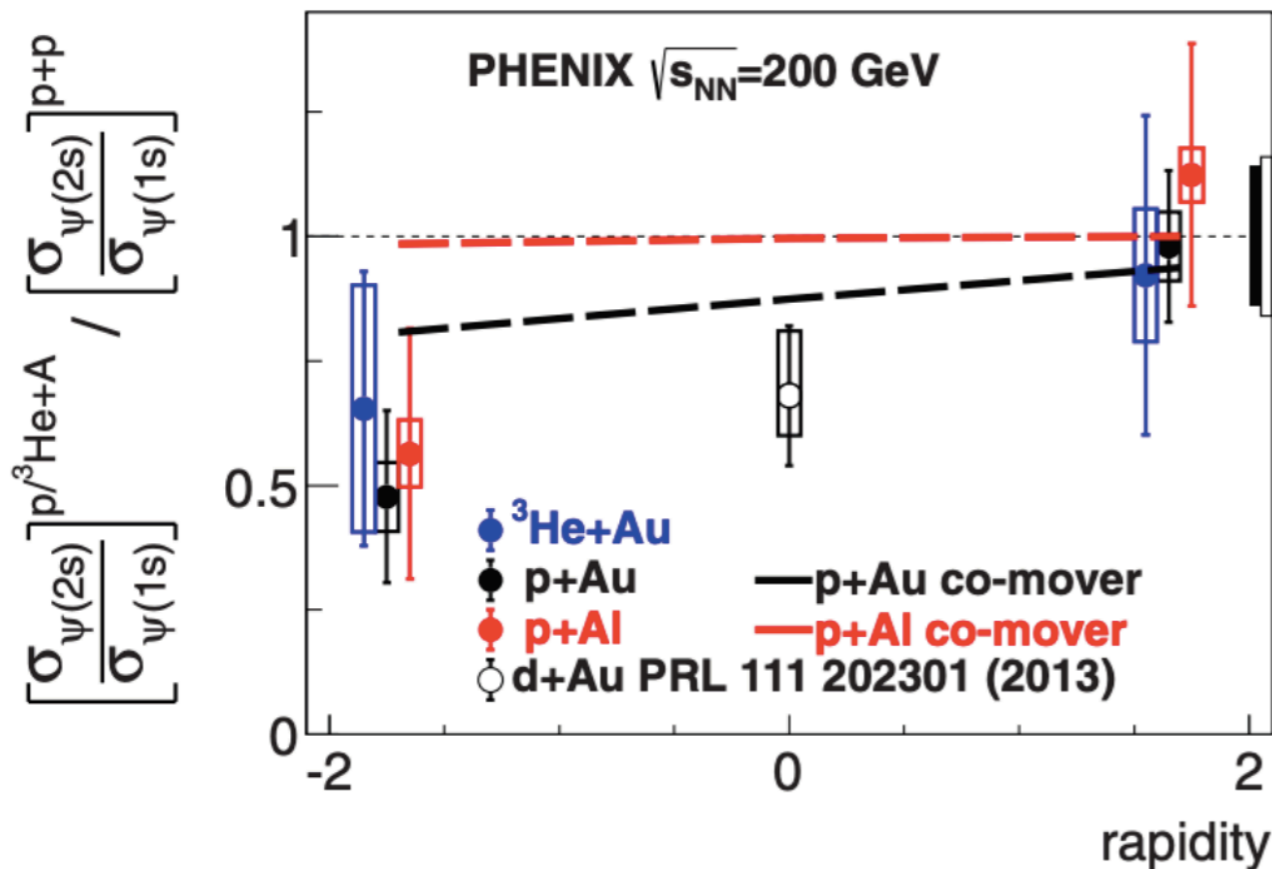
$\psi(2S)$ results

Strong suppression at mid & backward rapidity relative to J/ψ

- Not explained by CNM effects (CNM very similar for both states)



PHENIX, PRC95 034904 (2017)



The differential suppression is correlated with particle multiplicity

- Suggestive of a final state effect

Transport model

Du and Rapp ([JHEP 1903 \(2019\) 015](#)) have adapted their transport model, used to describe heavy ion collisions, for use in small systems.

They try to describe all available charmonium J/ψ and $\psi(2S)$ data from RHIC and LHC, including the J/ψ v_2 .

The transport model uses

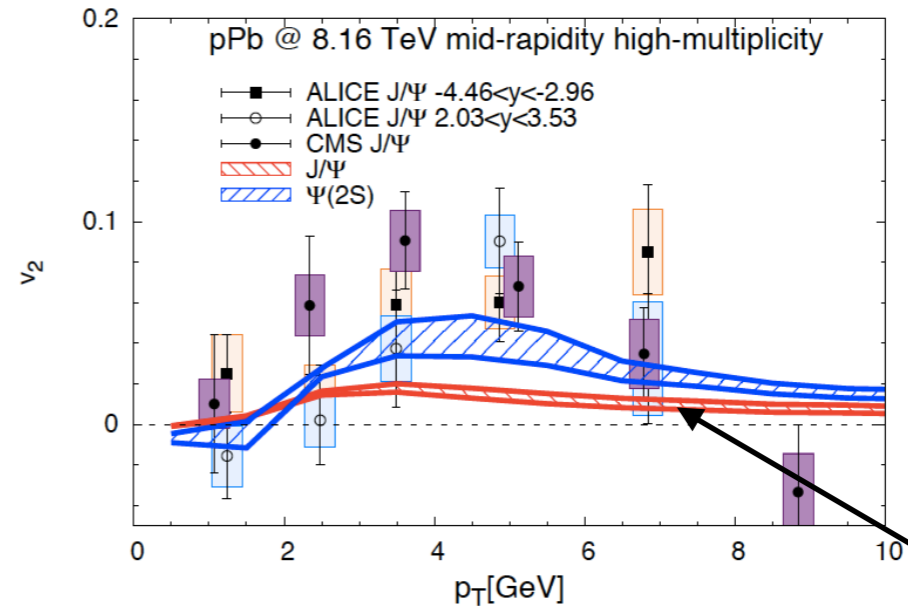
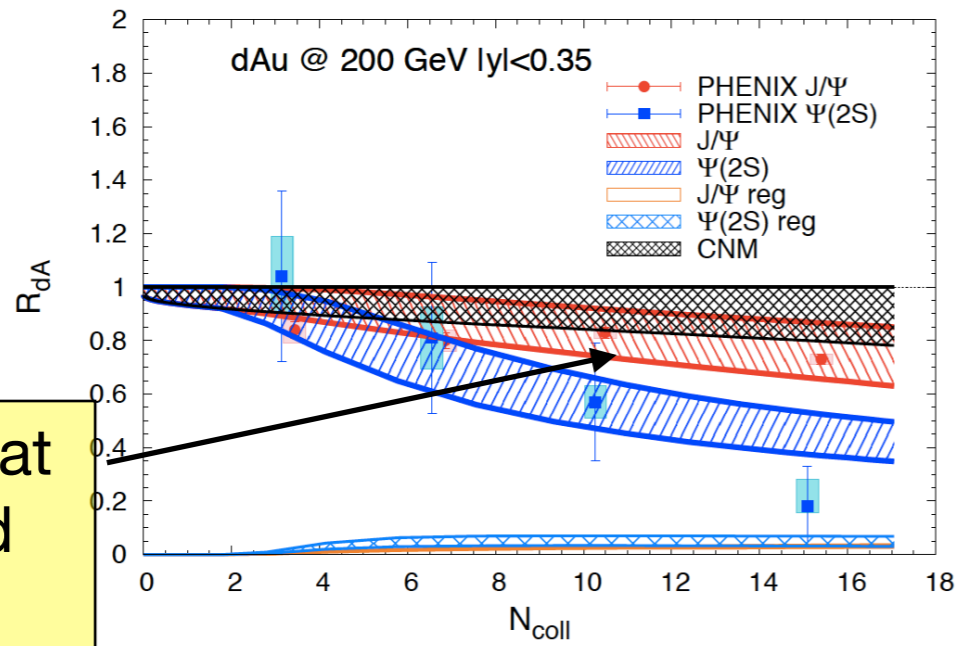
- A rate equation approach within a fireball model
- Initial geometry of the fireball from a Monte-Carlo event generator
- Initial anisotropies are caused by fluctuations
- Includes corrections for CNM effects
 - EPS09 shadowing with assumed linear centrality dependence
 - Constant nuclear absorption at backward rapidity

Some results for J/ψ , $\psi(2S)$ centrality dependence and J/ψ v_2 are shown on the next slide.

- The calculations also provide a good description of the p_T dependence — not shown here.

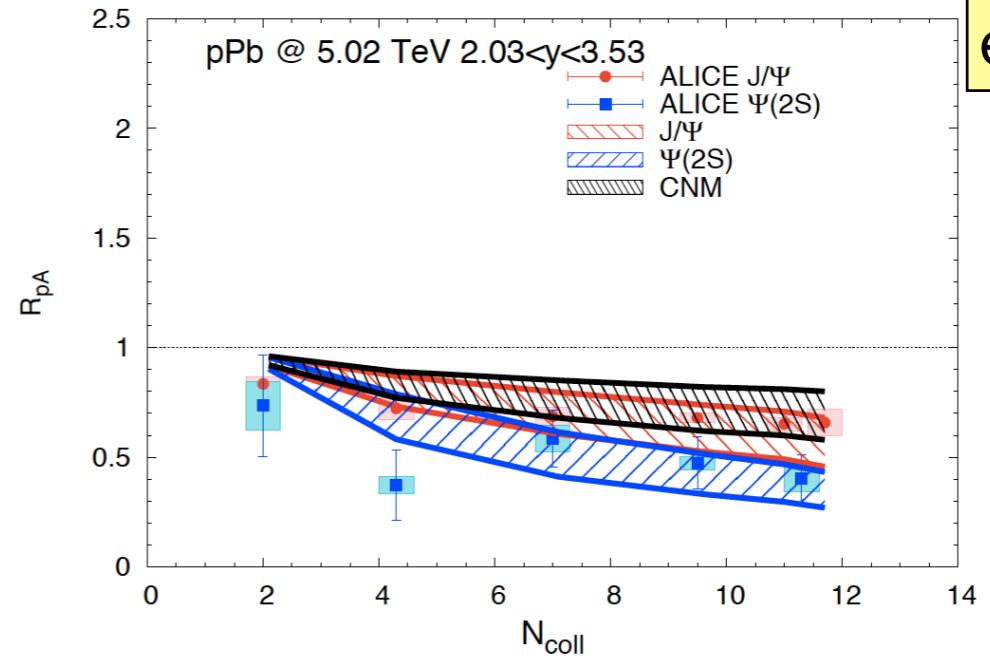
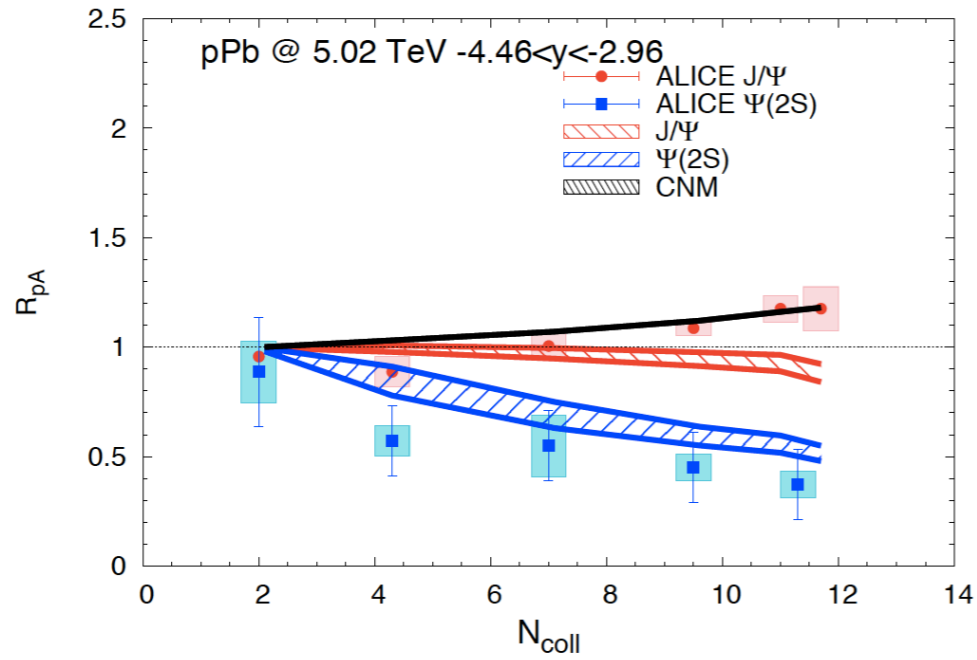
200 GeV

~20% effect at RHIC beyond CNM on J/ψ

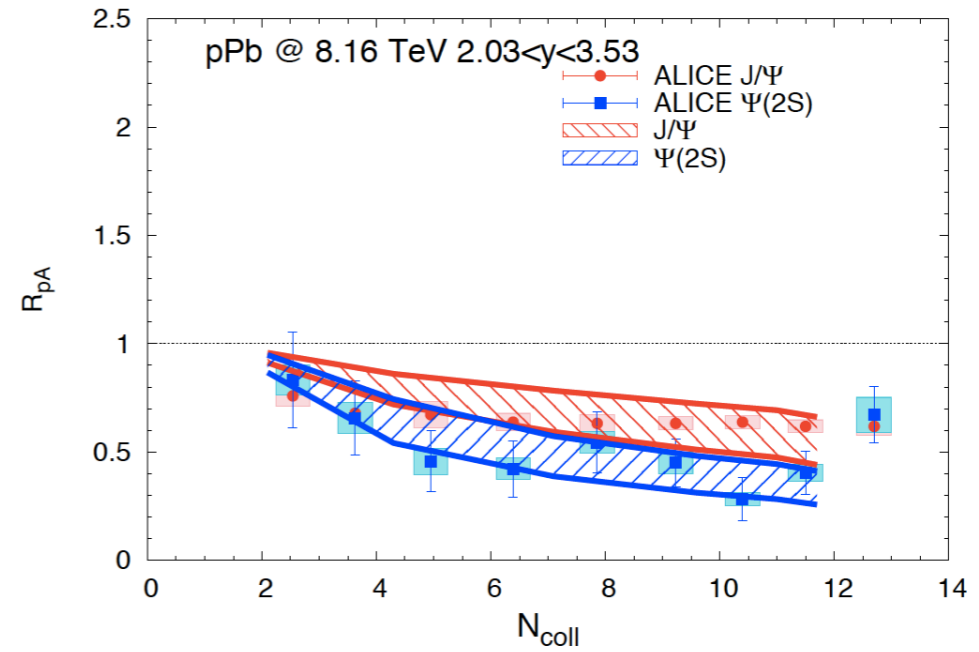
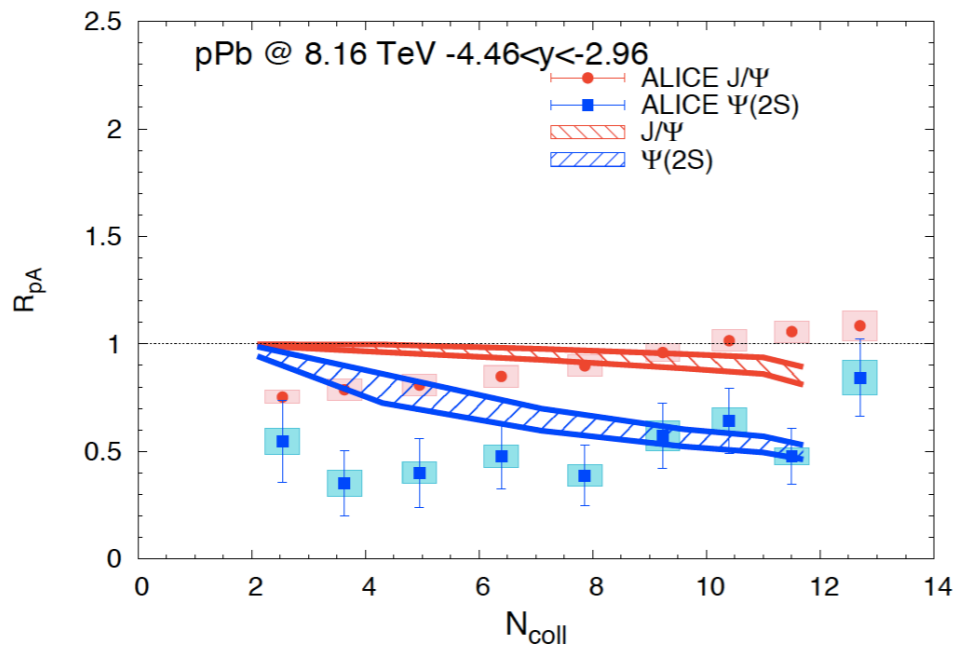


But: J/ψ v₂ not explained

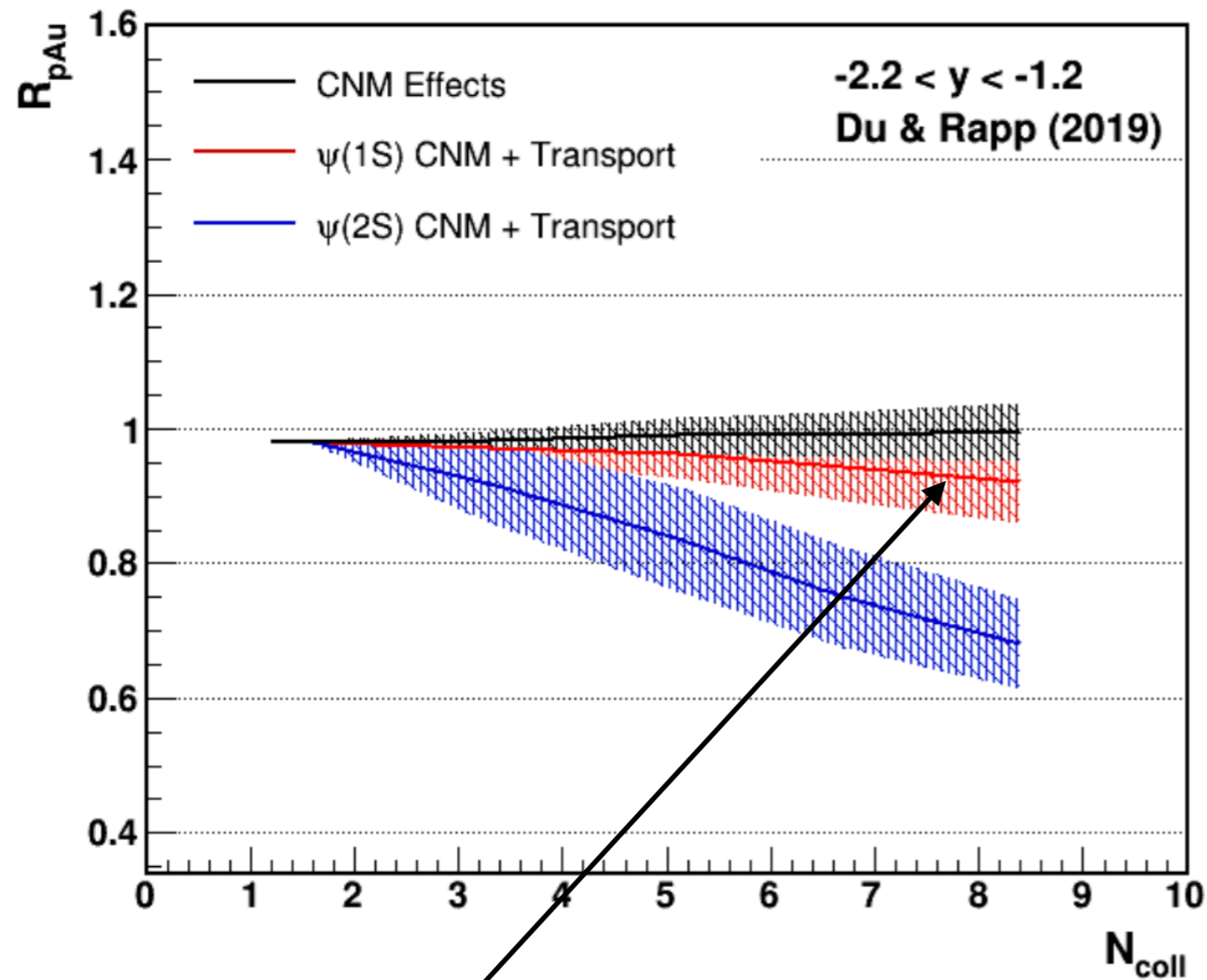
5.02 TeV



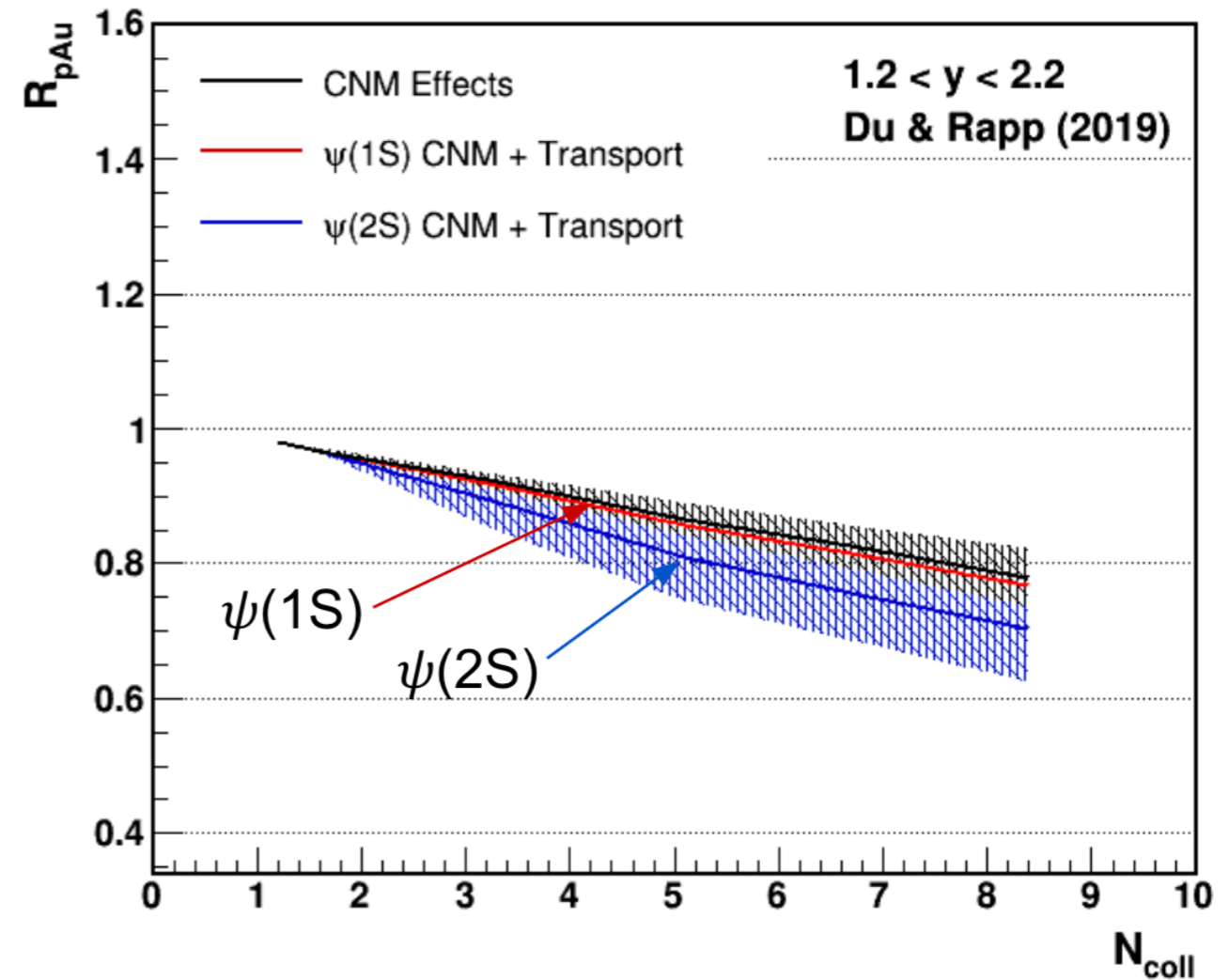
8.16 TeV



Perspective: transport model prediction for p+Au



~10% effect in p+Au beyond CNM on J/ψ



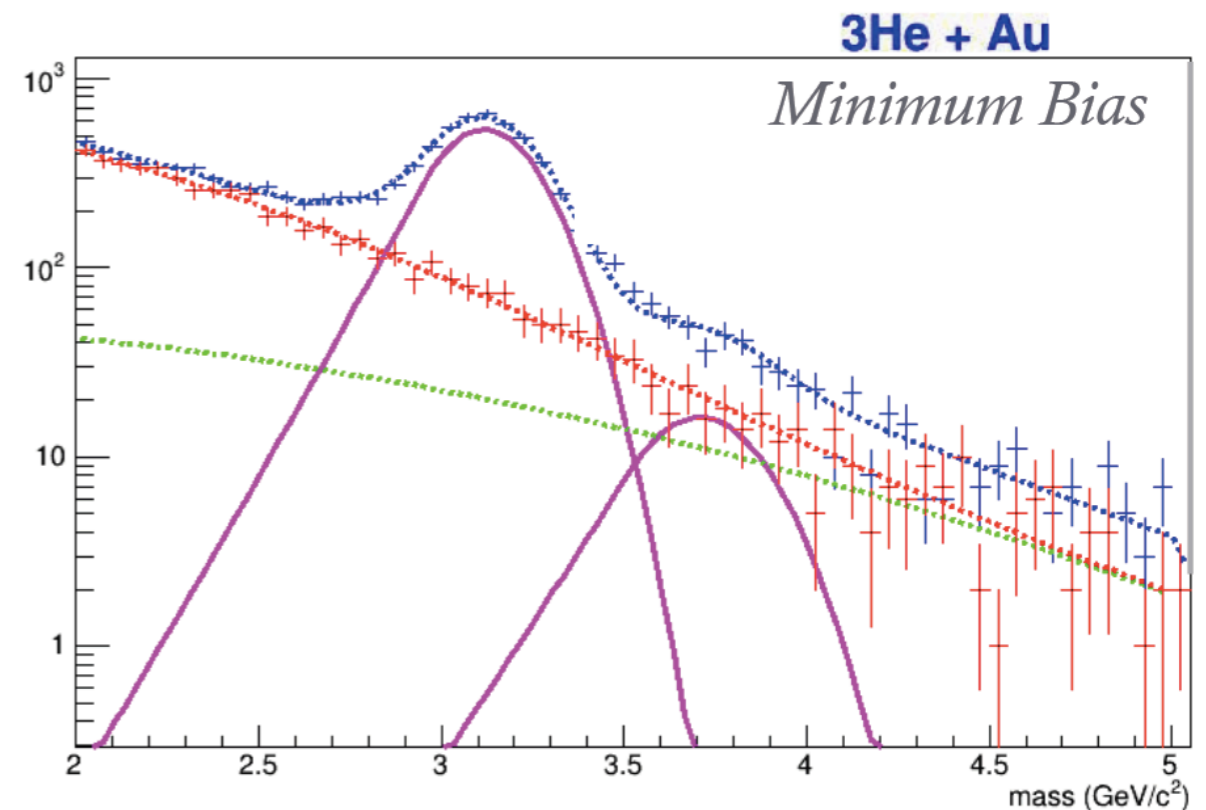
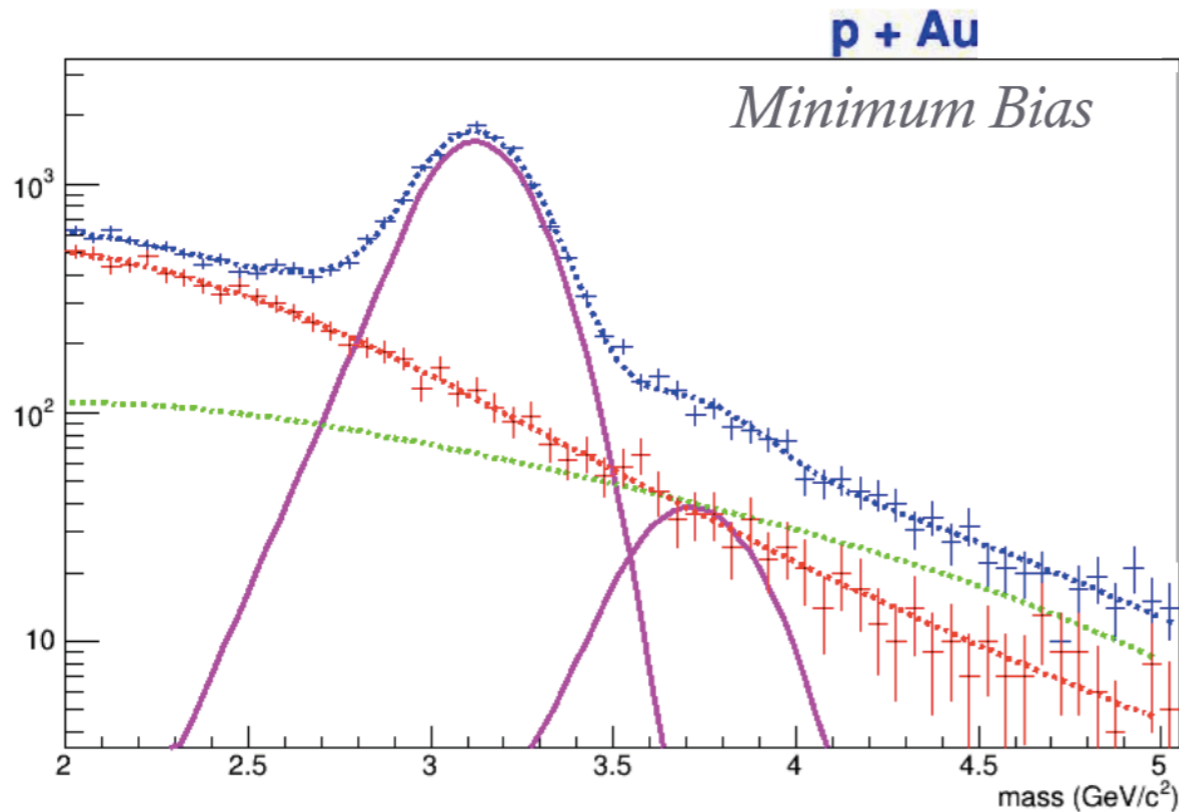
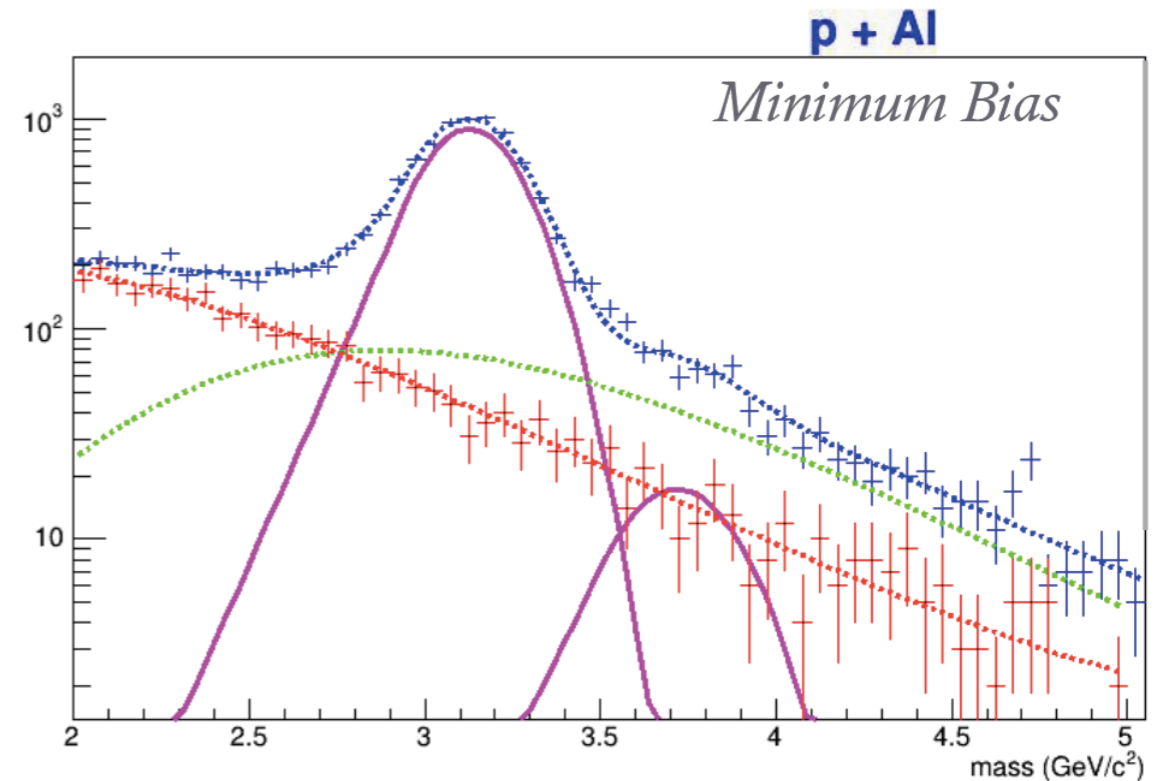
Little effect in p+Au beyond CNM on J/ψ

J/ ψ Measurement

Dimuon invariant mass spectra for p+Al, p+Au and ^3He +Au

With fitted J/ ψ peak and various background sources shown

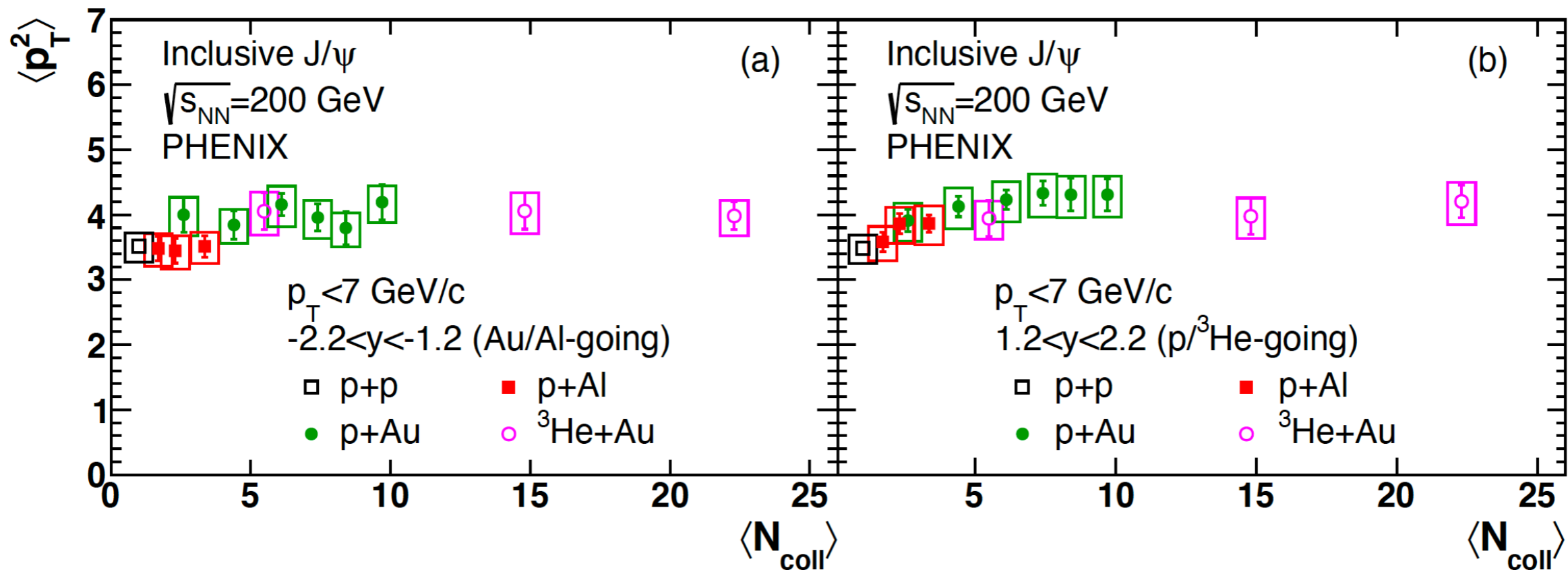
- **Red** : combinatoric from like-sign.
- **Green**: non-combinatoric (physics) background estimate.



$p+Al, p+Au, {}^3He+Au < p_T^2 > vs N_{coll}$

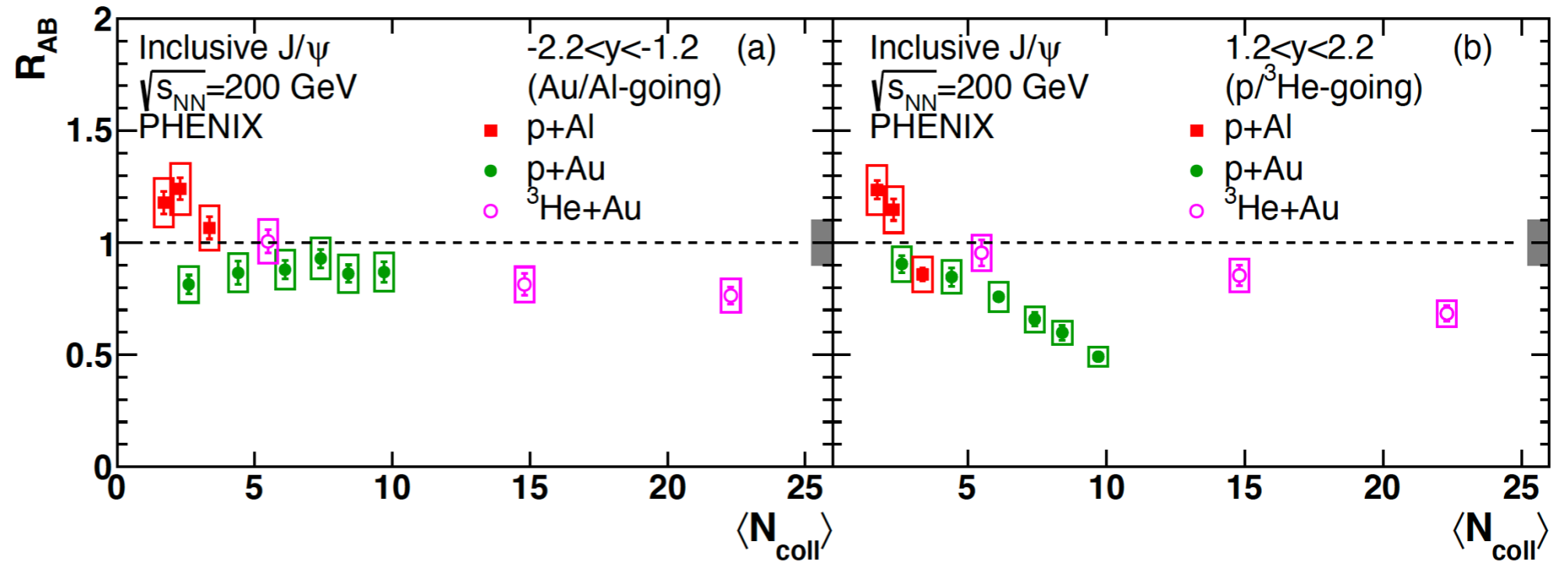
Compare $\langle p_T^2 \rangle$ for the three collision systems.

- Limited to $p_T < 7 \text{ GeV}/c$ because of statistics.
- Slight enhancement for larger N_{coll}
- No evidence of any effects due to projectile size.

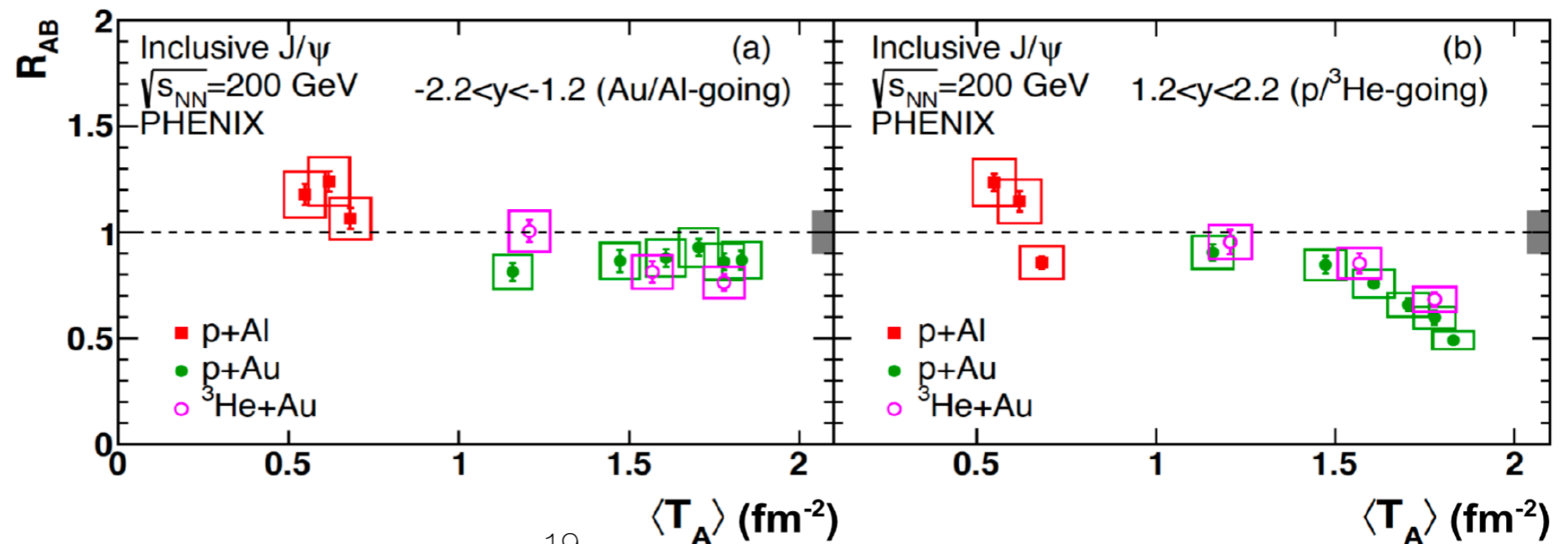


N_{coll} and T_A dependence

No N_{coll} scaling expected between p+Au and ^3He +Au.



Plot modification versus nuclear thickness of target (T_A) instead.



Final state effects on J/ψ production?

To look for evidence of final state effects on J/ψ in light systems we can compare Run 14 $^3\text{He}+\text{Au}$ and Run 15 p+Au data.

Systematic uncertainties reduced because:

- The runs are close in time.
- There were no detector upgrades in between.
- The same simulations model was used for analysis.

The particle multiplicity in $^3\text{He}+\text{Au}$ is \sim twice that in p+Au (Phys. Rev. Lett. 121, 222301 (2018)).

Make the ratio of $^3\text{He}+\text{Au}$ to p+Au modifications.

- Should be < 1 if greater energy produced in $^3\text{He}+\text{Au}$ is important.
- Expect the greatest effect at backward rapidity

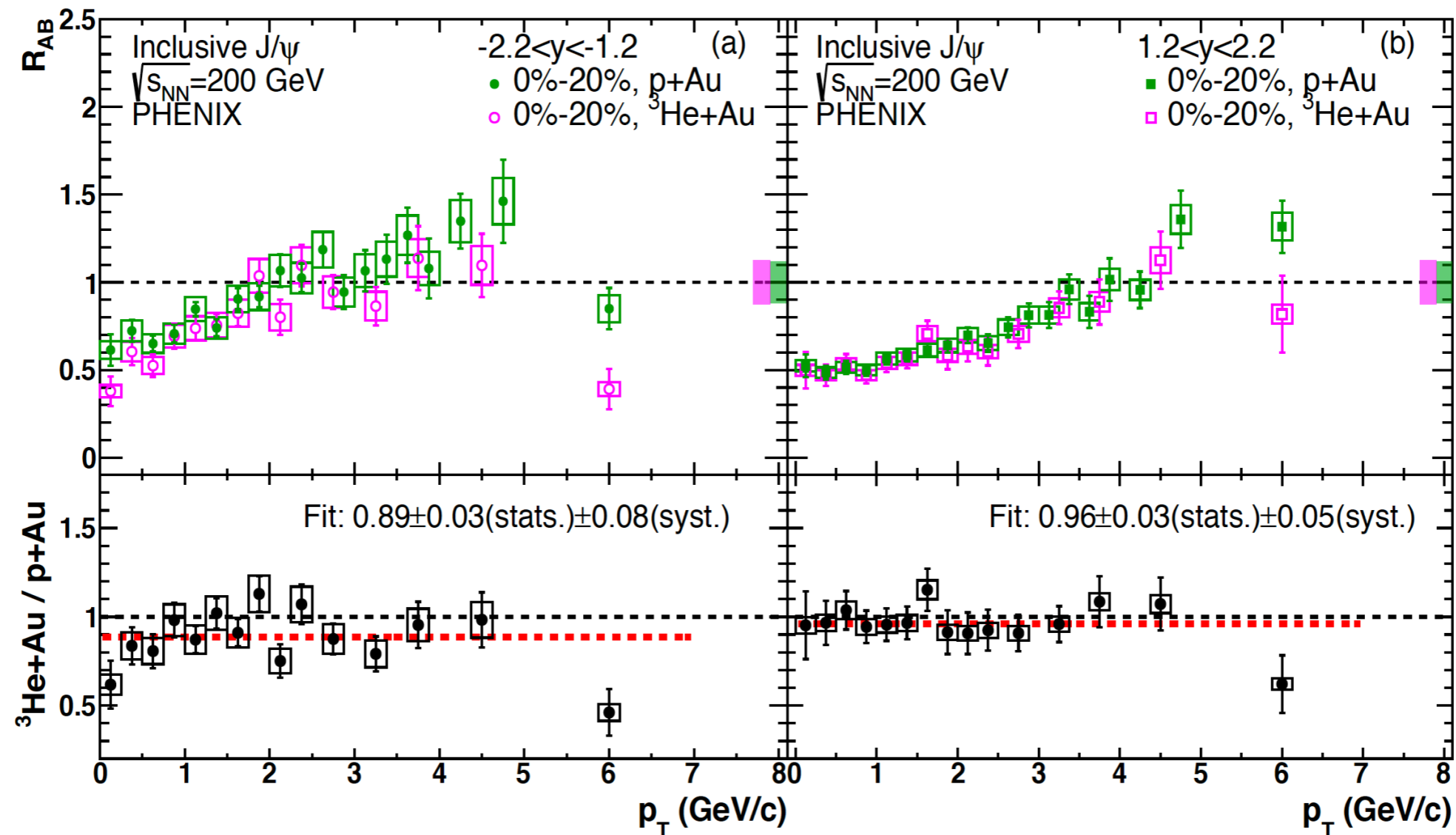
$^3\text{He}+\text{Au}$ to $p+\text{Au}$ ratio (0-20% centrality)

Backward rapidity ratio $0.89 \pm 0.03 \pm 0.08$

- Consistent with some additional suppression (90% probability).
- But not far beyond the systematic uncertainty.

Forward rapidity ratio $9.96 \pm 0.03 \pm 0.05$

- Consistent with 1



$^3\text{He}+\text{Au}$ to $p+\text{Au}$ ratio (0-20% centrality)

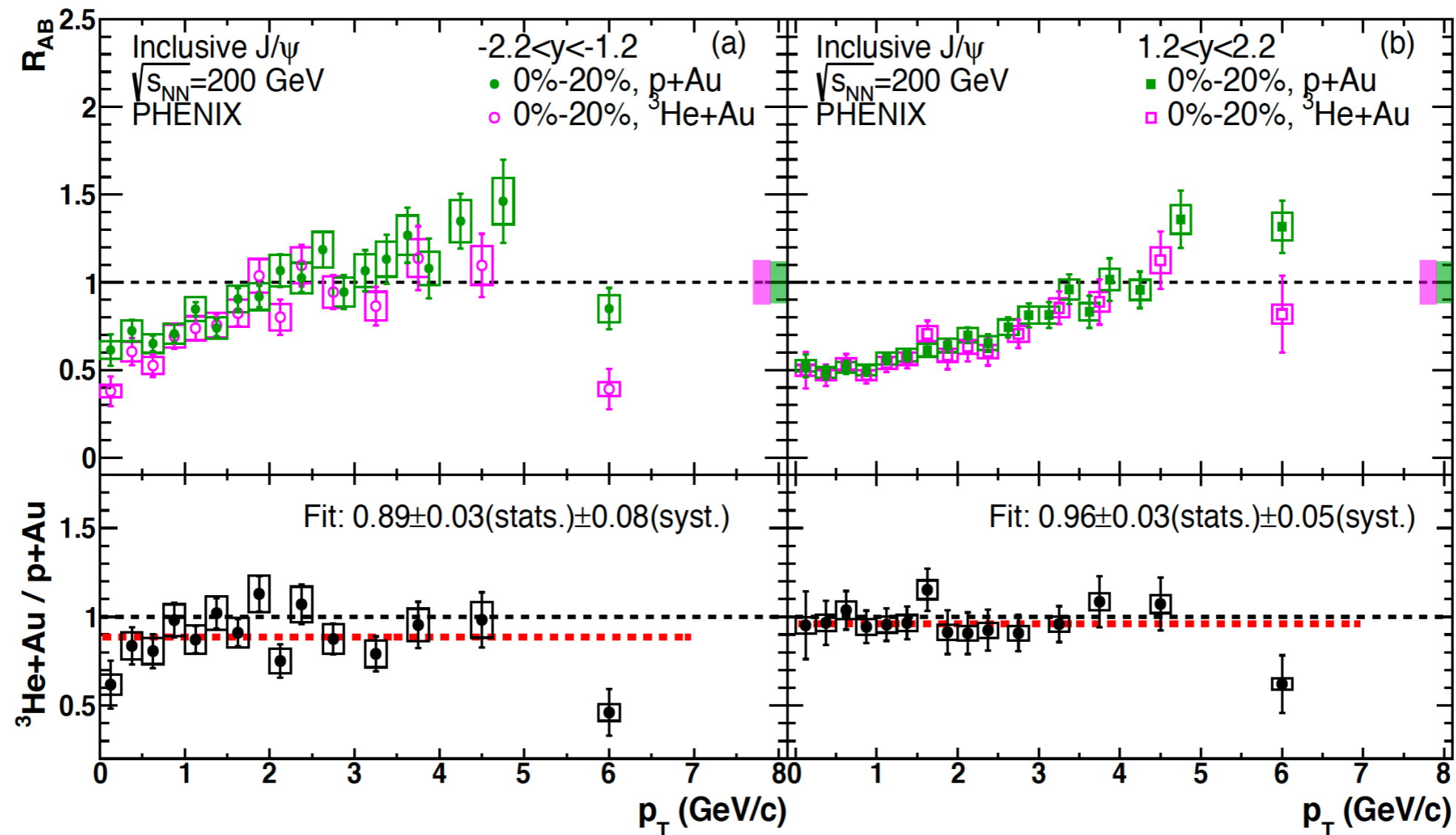
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Forward rapidity ratio $0.96 \pm 0.03 \pm 0.05$

- Consistent with 1

Both are consistent with the transport model



How well do we understand J/ψ modification?

We have a new p+Au J/ψ data set with high integrated luminosity.

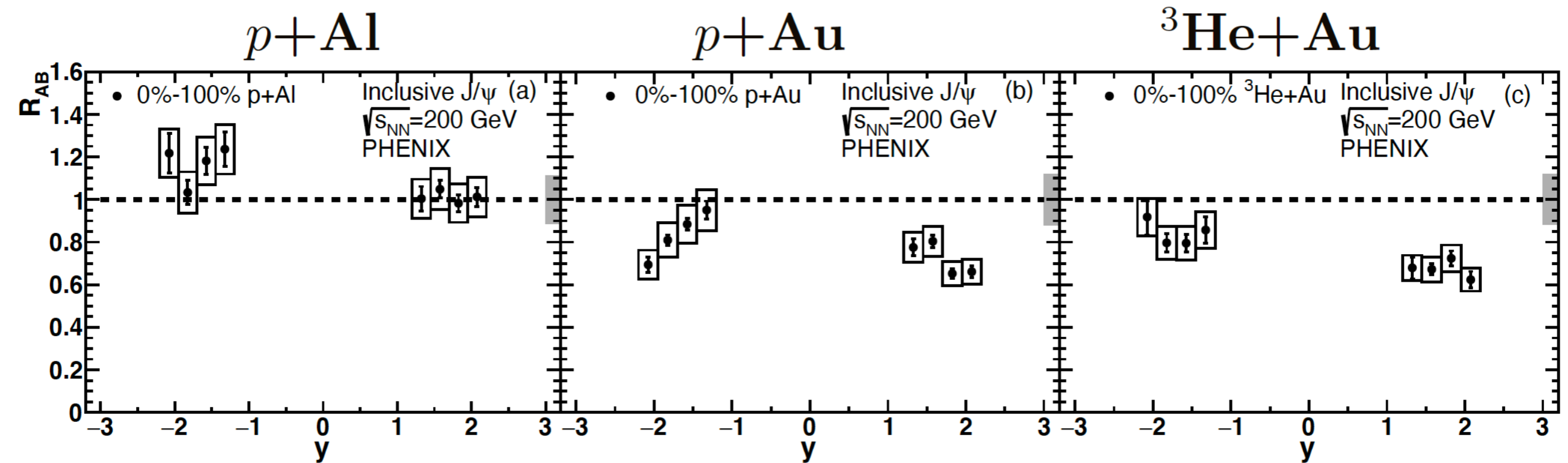
We can use it to test our understanding of the effects that contribute to J/ψ modification.

We look at the

- Rapidity dependence
- p_T dependence
- Centrality dependence

Rapidity dependence, p_T and centrality integrated

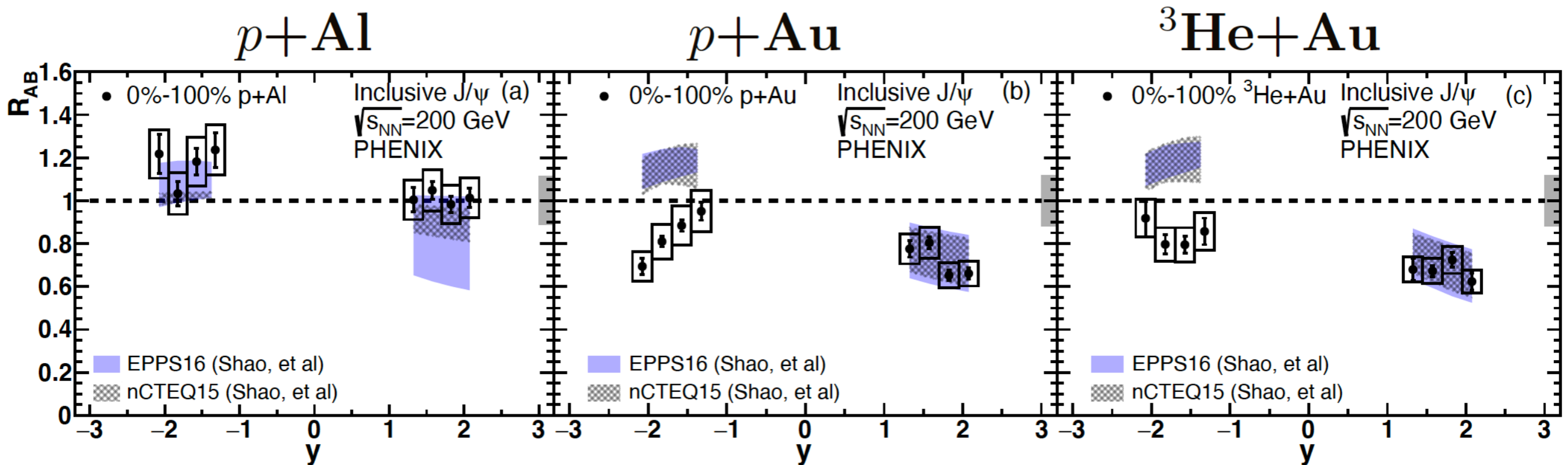
0-100% centrality, p_T integrated.



Rapidity dependence, p_T and centrality integrated

0-100% centrality, p_T integrated.

- Add EPPS16 and nCTEQ15 shadowing, with Bayesian re-weighting.



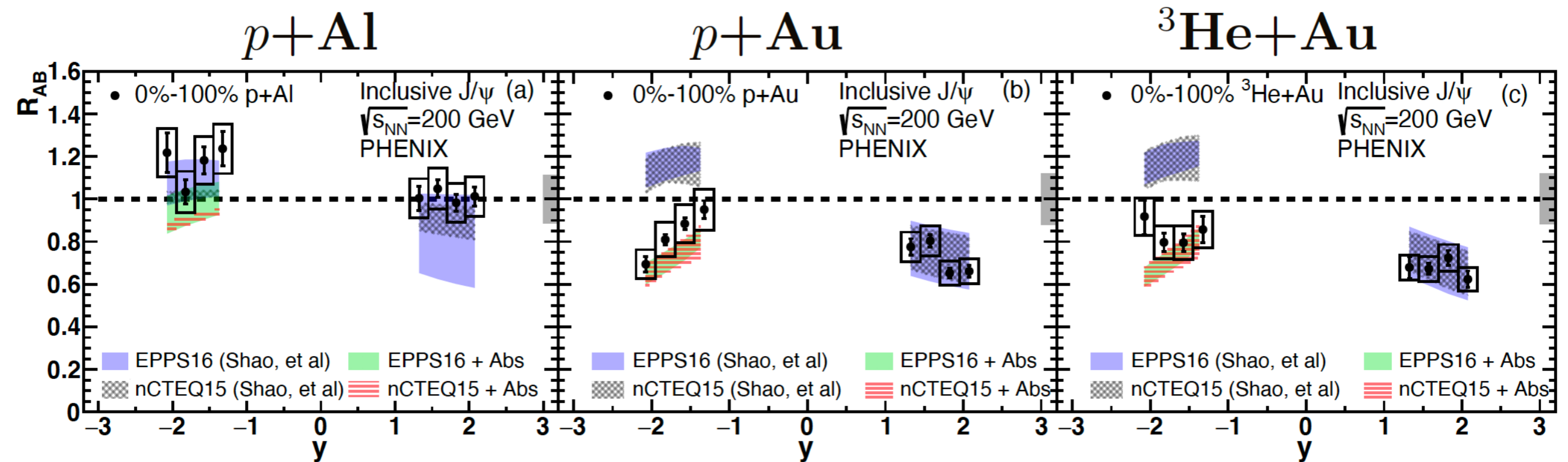
Nice agreement at forward rapidity.

Anti-shadowing alone at backward rapidity does not describe the data.

Rapidity dependence, p_T and centrality integrated

0-100% centrality, p_T integrated.

- Add EPPS16 and nCTEQ15 shadowing, with Bayesian re-weighting.
- Fold absorption prediction with shadowing at backward rapidity.



Not so bad!

How about the p_T dependence?

p_T dependence, 0-100% centrality

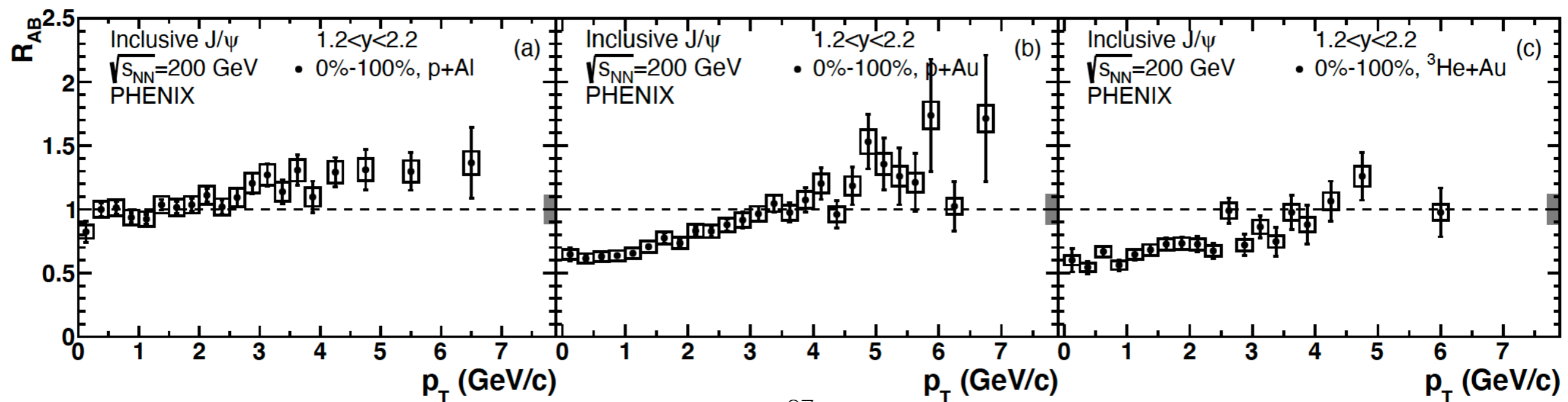
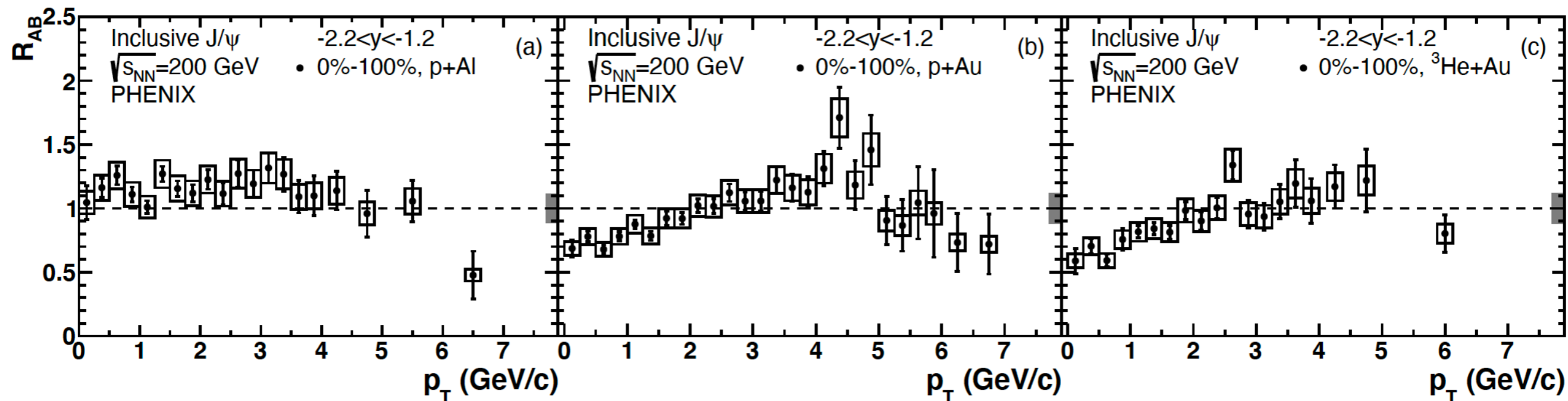
Top row: Backward rapidity

Bottom row: Forward rapidity

$p+Al$

$p+Au$

^3He+Au



p_T dependence, 0-100% centrality

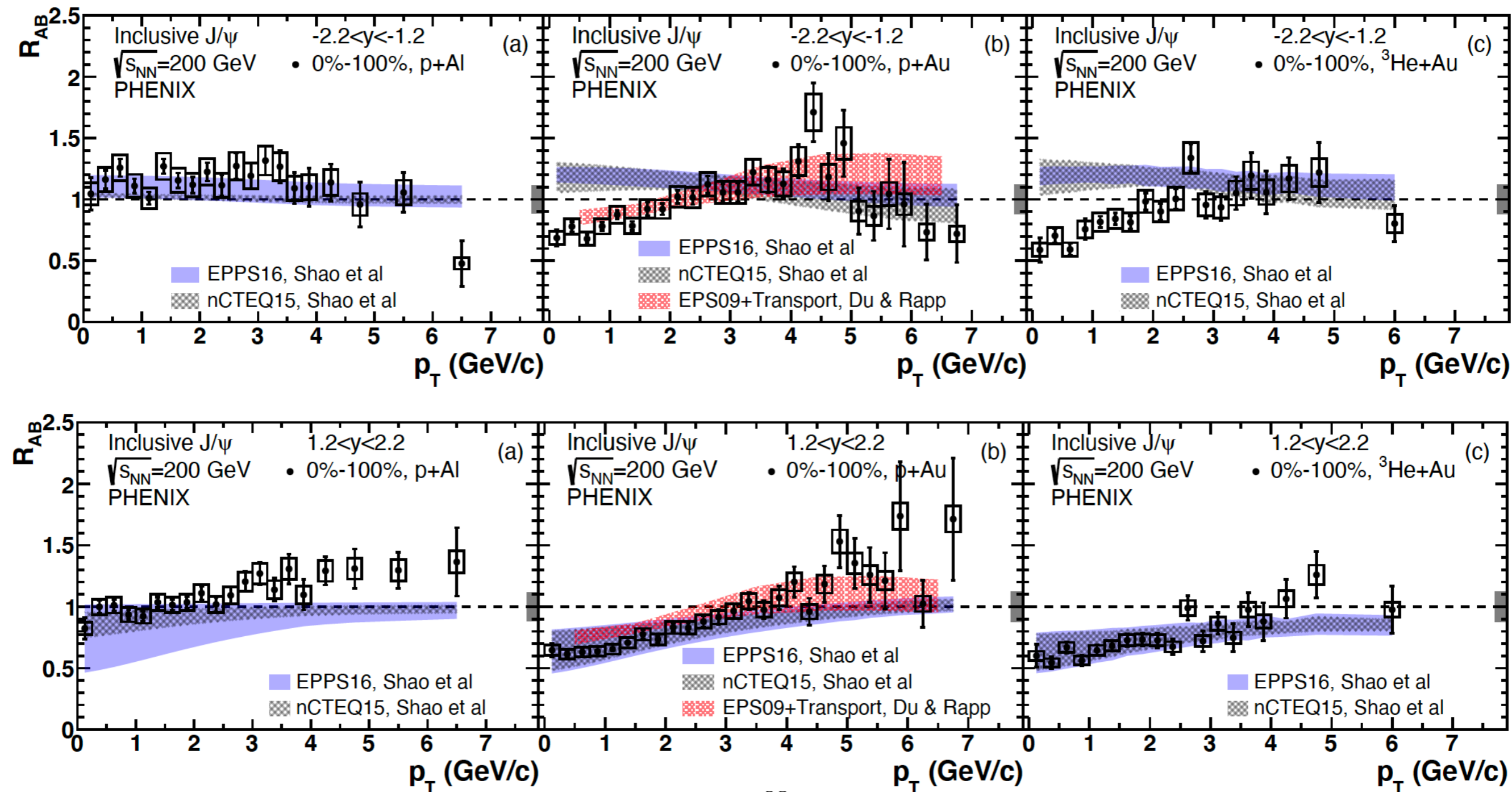
Blue: Bayesian re-weighted shadowing only

Red: Transport + EPS09 + absorption (-y) + p_T broadening

$p+Al$

$p+Au$

^3He+Au



p_T dependence, 0-100% centrality

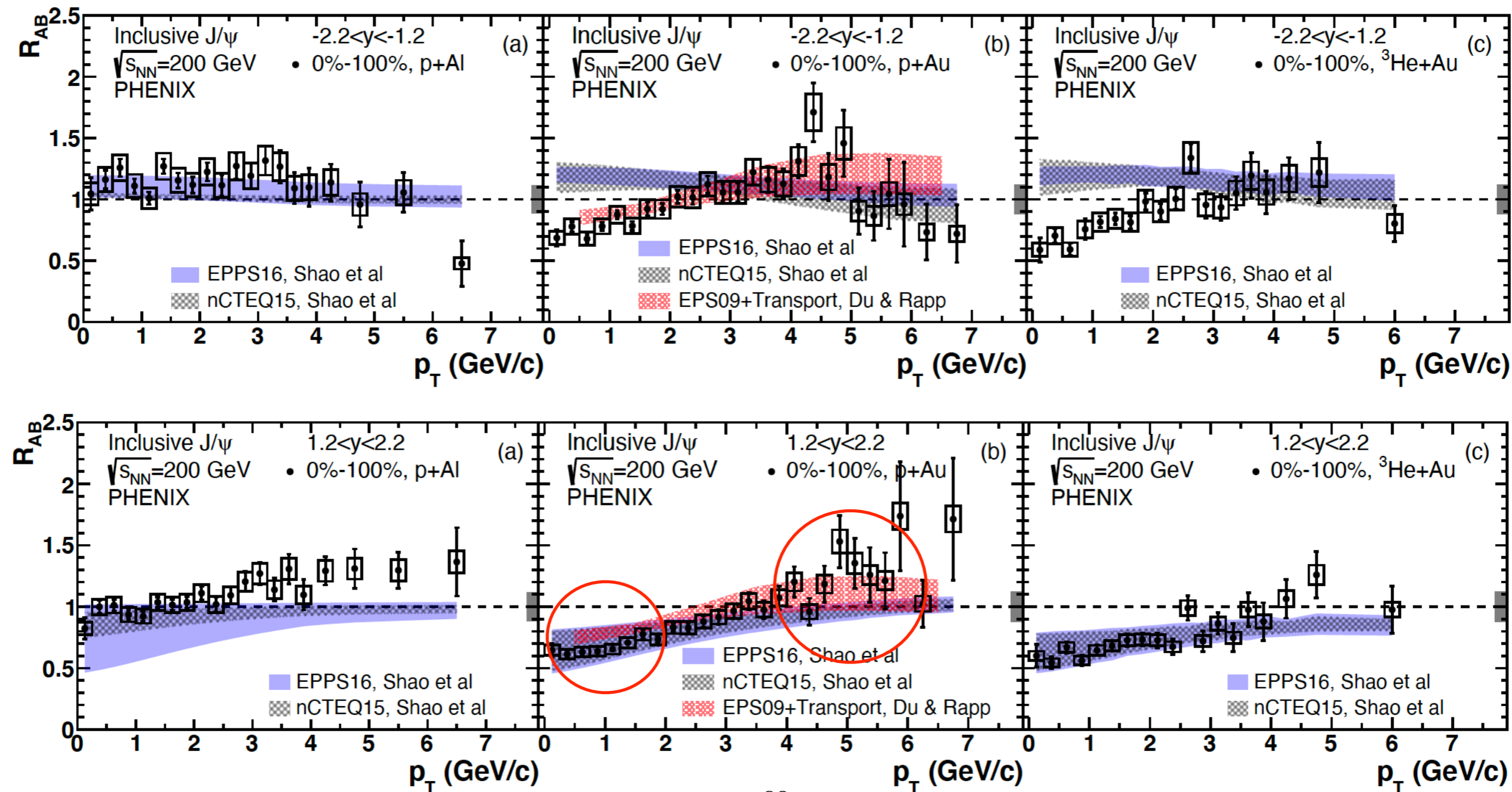
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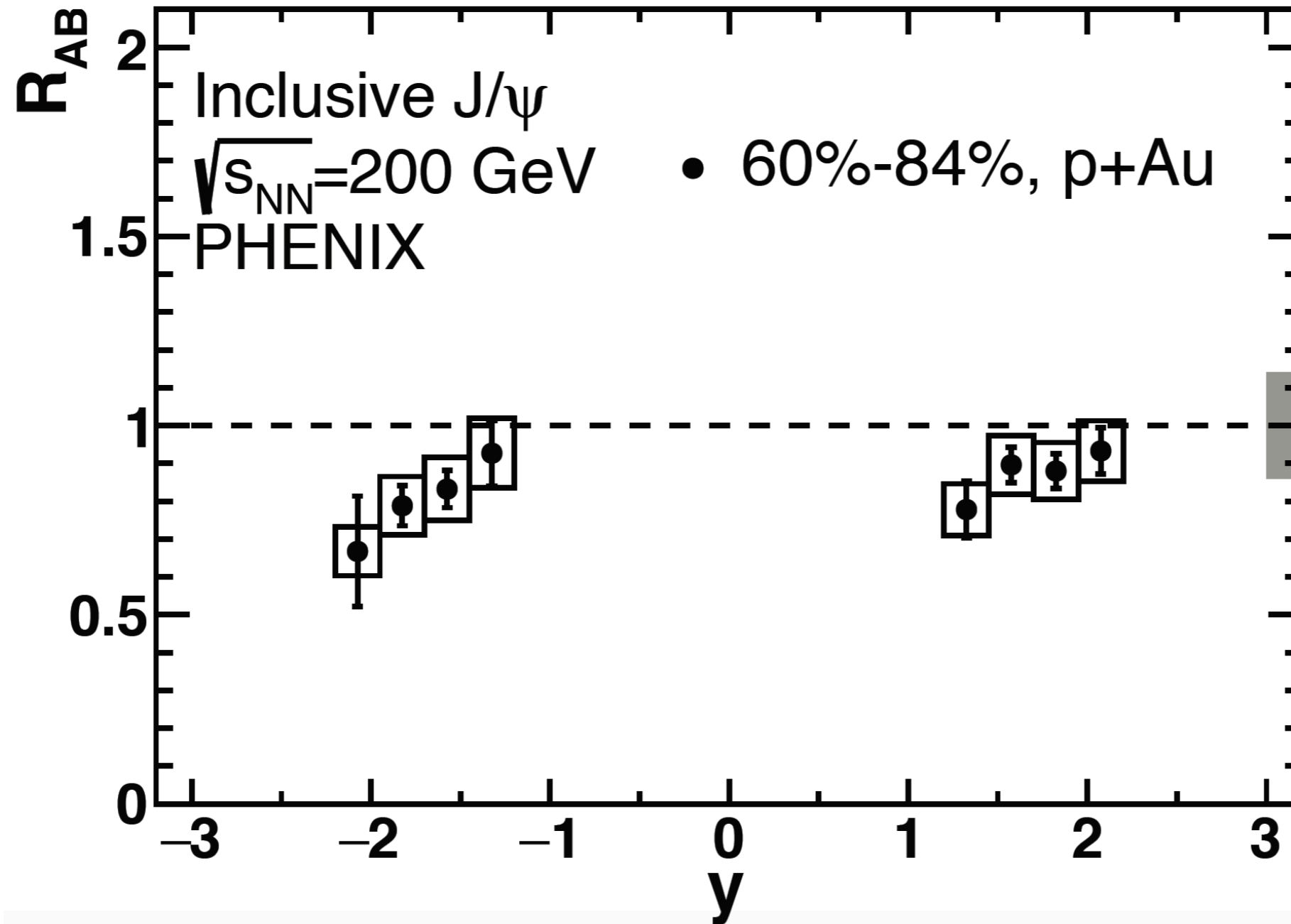
$p+Al$

$p+Au$

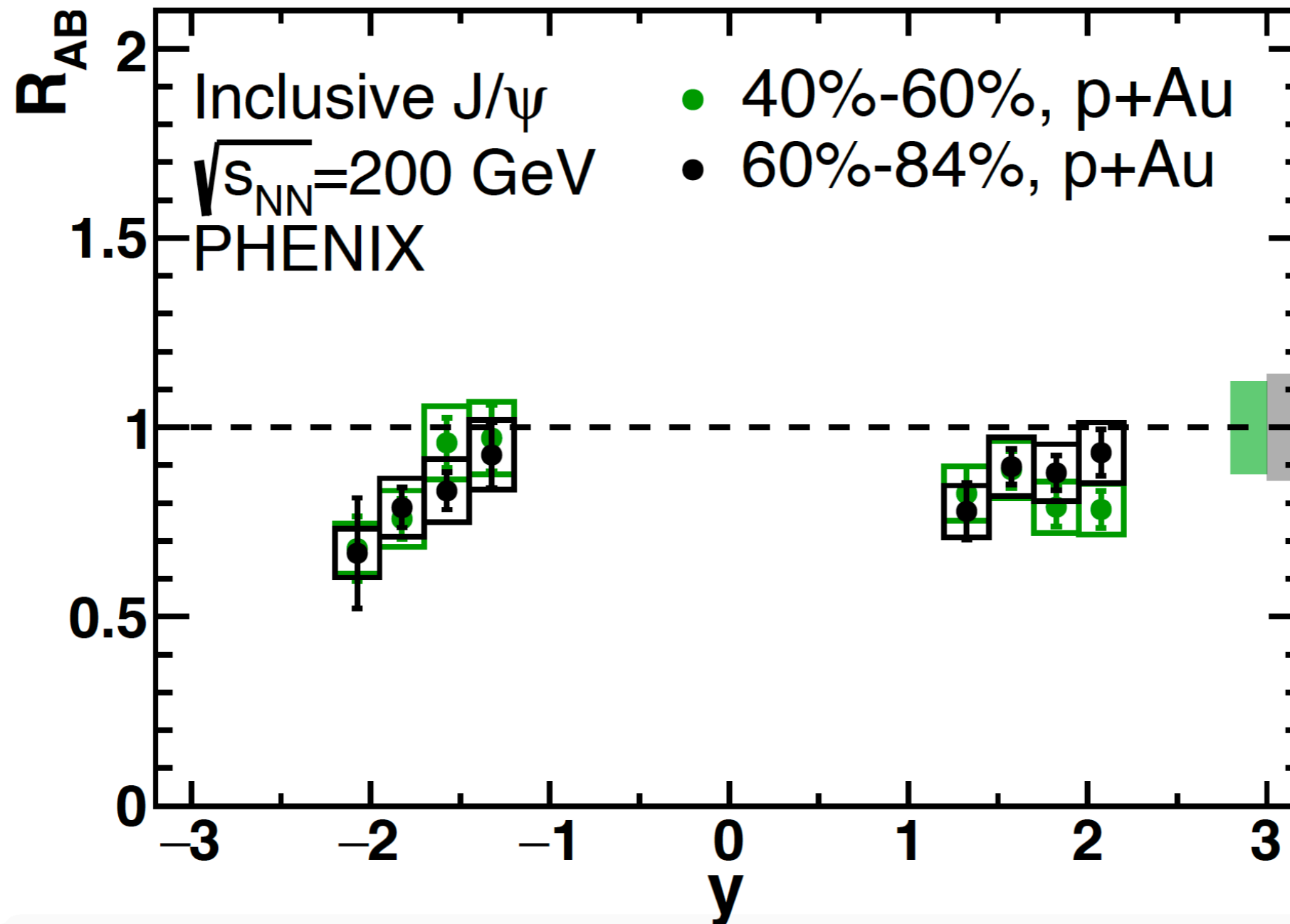
^3He+Au



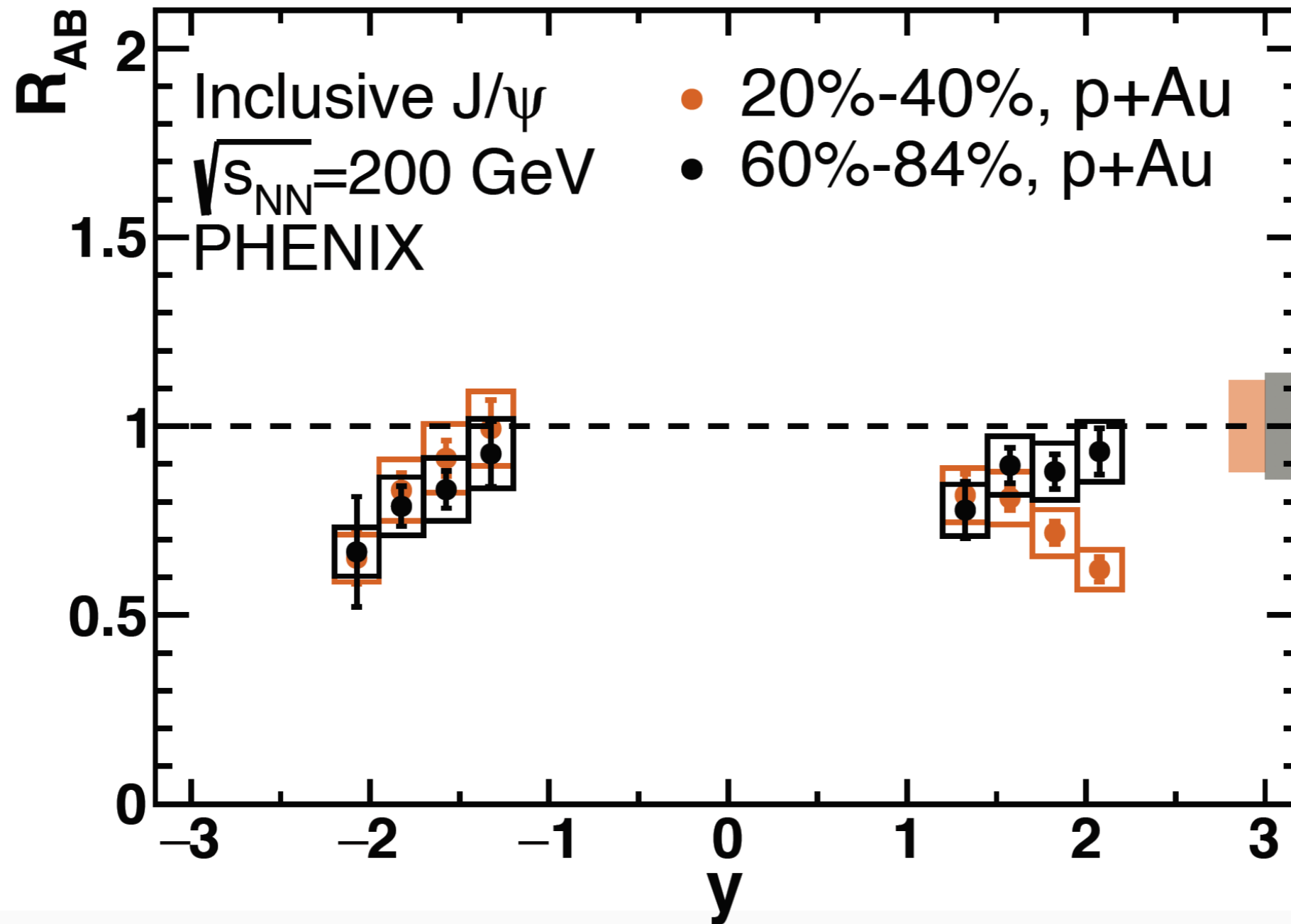
p+Au centrality dependence



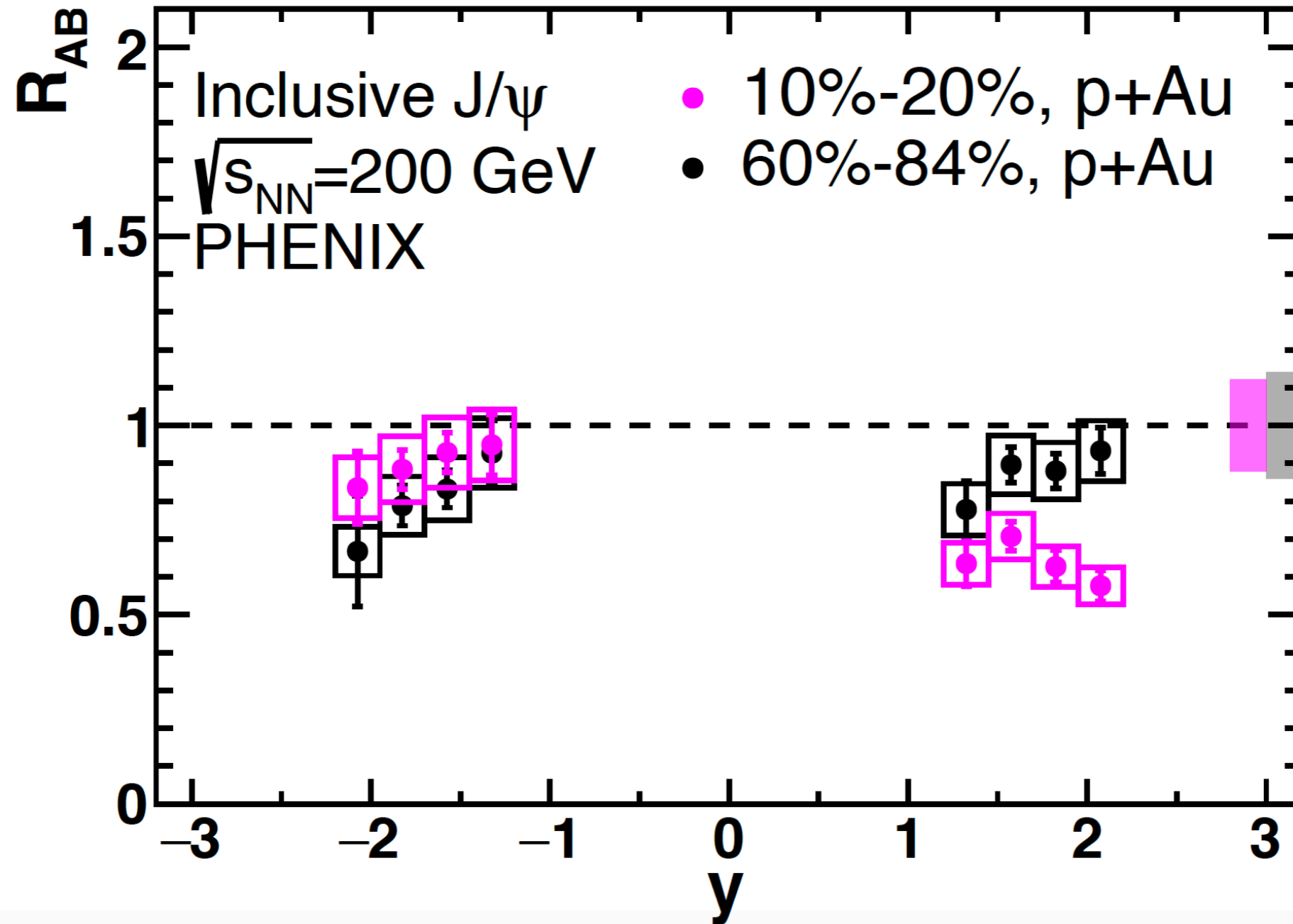
p+Au centrality dependence



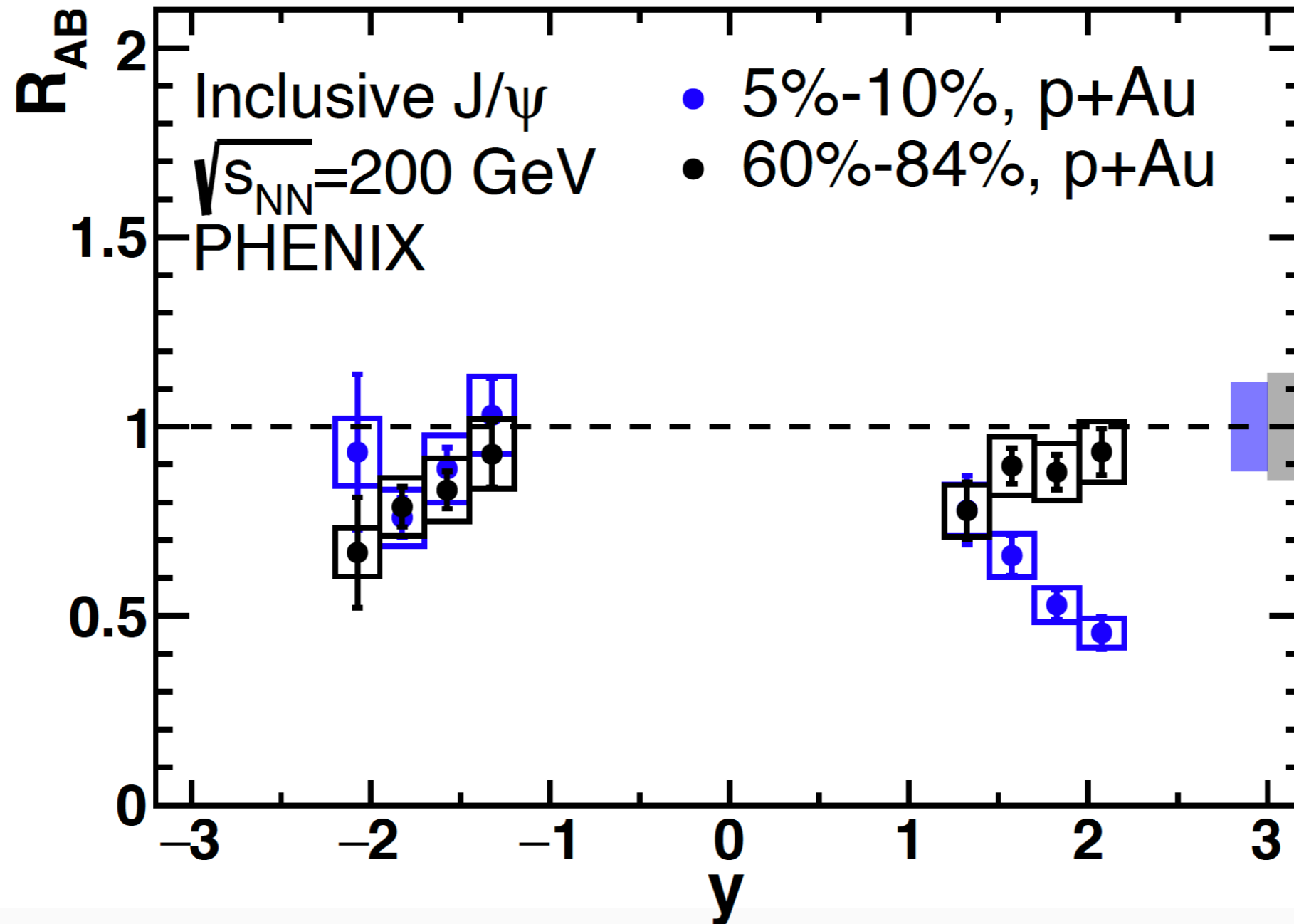
p+Au centrality dependence



p+Au centrality dependence



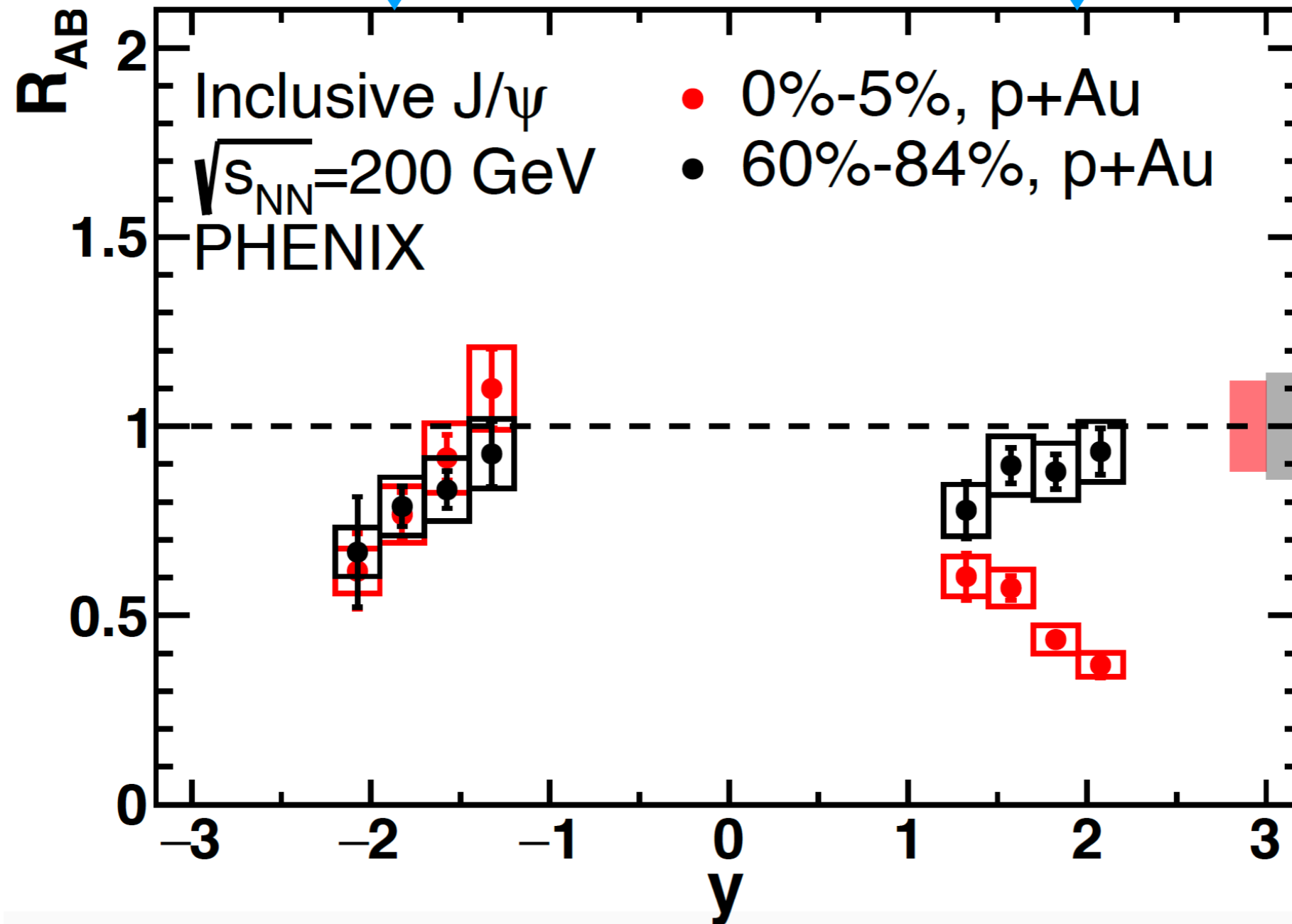
p+Au centrality dependence



p+Au centrality dependence

Trade-off between anti-shadowing and absorption.

Very strong centrality dependence of suppression.

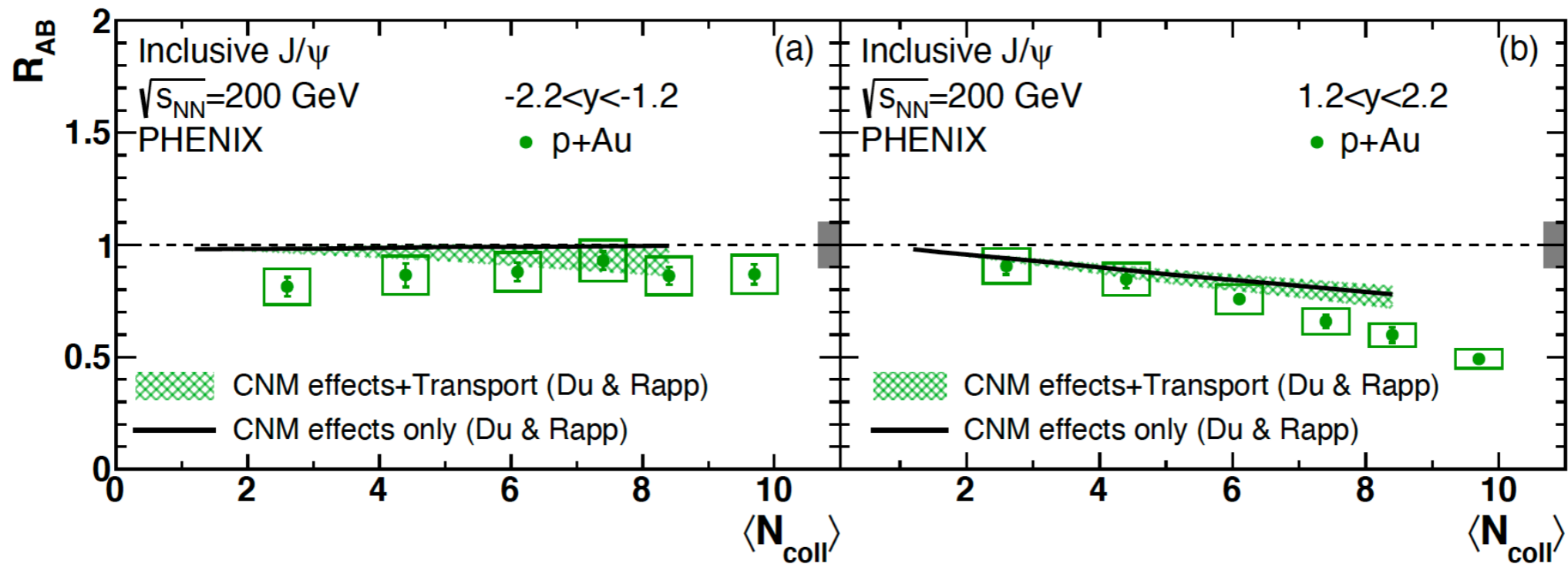


p+Au N_{coll} dependence

Compare transport calculation with N_{coll} dependence of p_T integrated data.

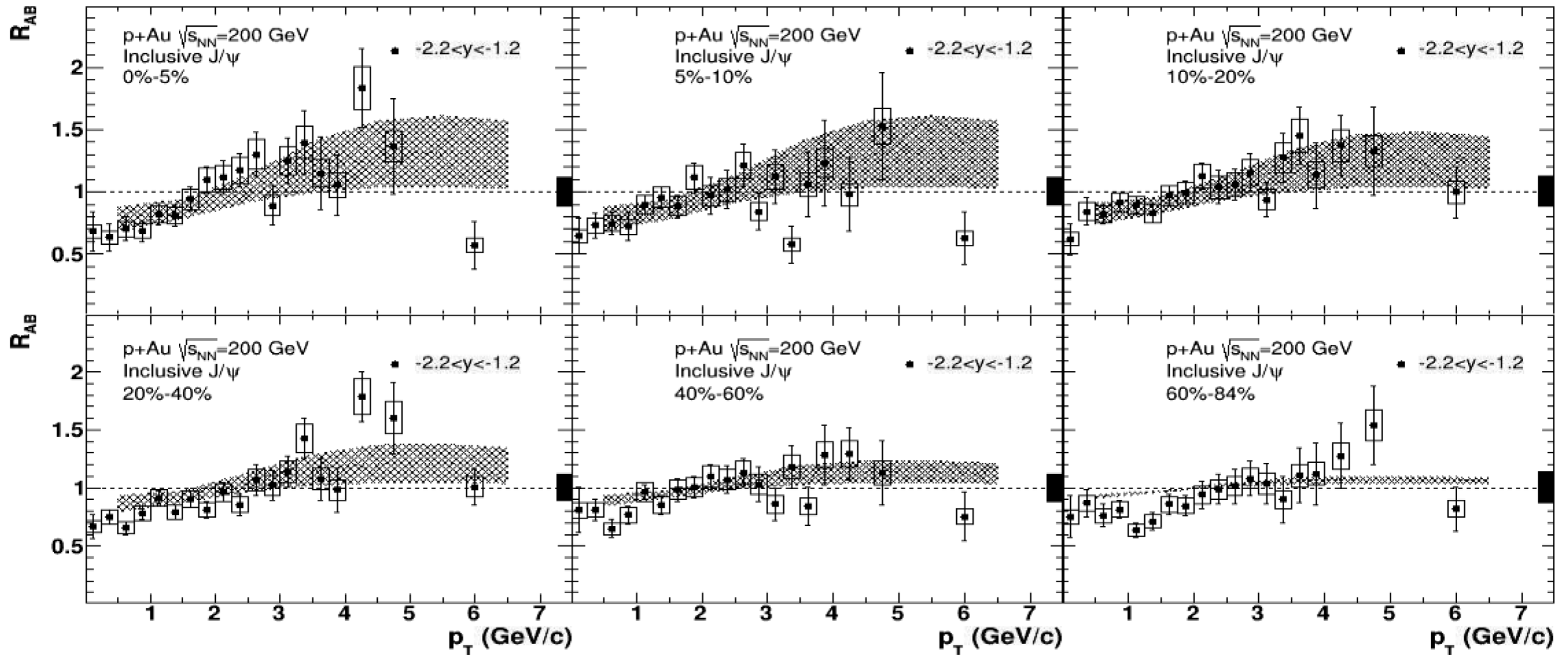
Forward rapidity suppression is dominated by EPS09 shadowing.

- Centrality dependence is assumed to be linear.
- Underpredicts suppression considerably!



$p+Au$ p_T dependence, binned in centrality

Backward rapidity $-2.2 < y < -1.2$
Compared with transport model

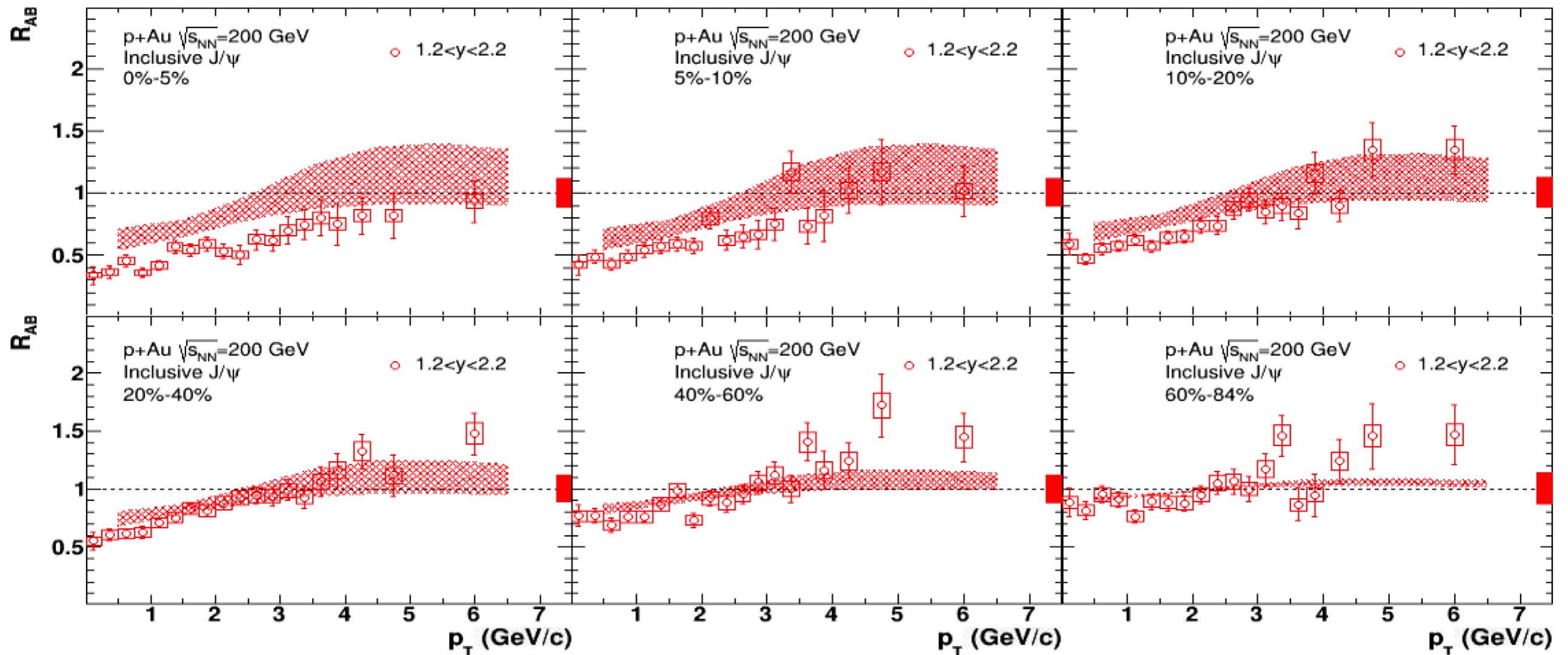


Transport model does a good job at backward rapidity.

$p+Au$ p_T dependence, binned in centrality

Forward rapidity $1.2 < y < 2.2$

Compared with transport model



- Remember:** Transport model has very small effect from medium at forward y .
So: EPS09 shadowing does not have a strong enough centrality dependence.
Bayesian re-weighted shadowing describes 0-100% data well
- But has no centrality prediction.

Comments

The **centrality and p_T integrated** data are described well by the Bayesian re-weighted shadowing parameterizations.

- Requires the addition of nuclear absorption at backward rapidity.

The p_T dependence of the **centrality integrated** data at forward rapidity is well described.

- Shadowing alone needs some p_T broadening added, but does pretty well.

The **centrality dependence** at forward rapidity is not understood.

- The shadowing parameterizations are fitted to centrality integrated data.
- An attempt to determine centrality dependence from the target mass dependence of data (EPS09s: Helenius et. al., JHEP 1207 (2012) 073) does not describe data.
- **So you have to make up your own centrality dependence of shadowing.**

Shadowing centrality dependence

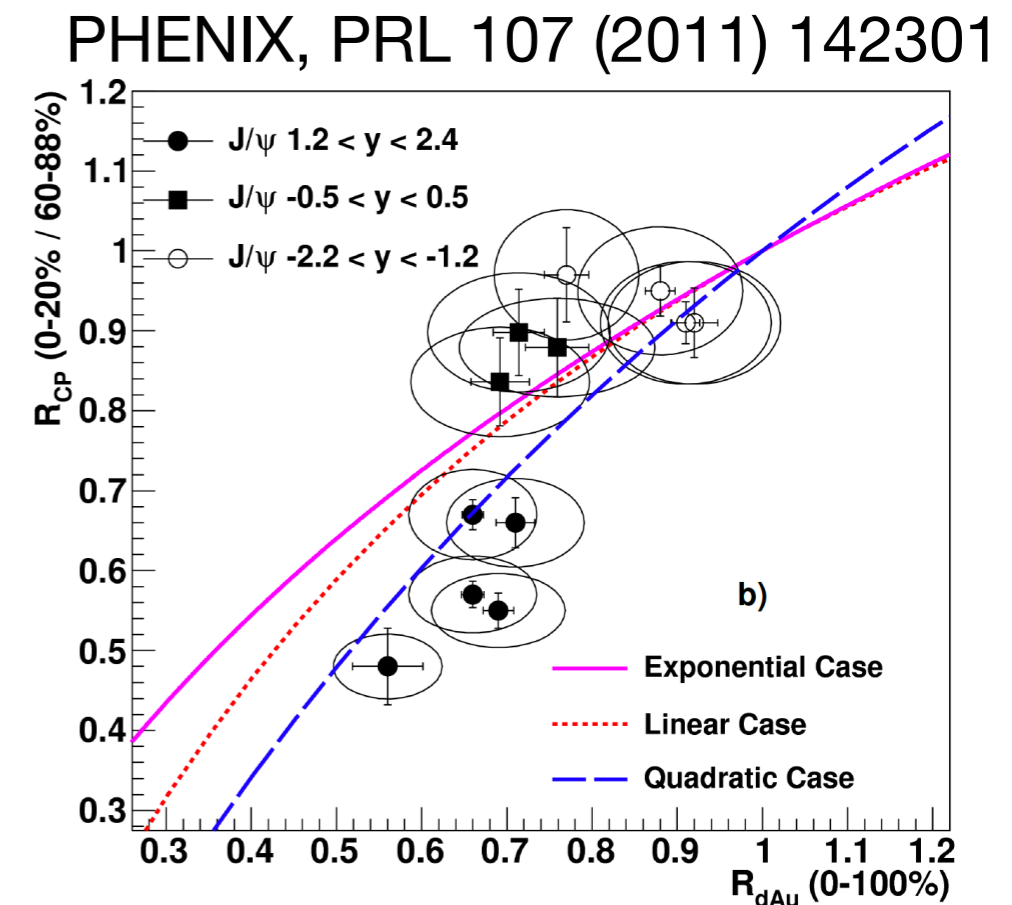
The centrality dependence of shadowing is an unsolved problem.

This plot from PHENIX d+Au J/ ψ data indicates that the forward rapidity shadowing needs a **stronger than quadratic** dependence on nuclear thickness to match the rapidity dependence of the data.

A later study (McGlinchey et.al. PRC 87 (2013) 5, 054910) concluded that shadowing was **heavily concentrated at small impact parameter** for d+Au.

Our new p+Au data set has advantages for studying this.

- The connection of impact parameter to centrality is more direct in p+Au.
- High statistical precision allows fine binning for central collisions.



Summary

Measured J/ψ modification in p+Al, p+Au and $^3\text{He}+\text{Au}$.

Increased modification in $^3\text{He}+\text{Au}$ over p+Au is small

- Consistent with Du and Rapp transport model
- Future work: try to extract $\psi(2S)$ in p+Au (and maybe $^3\text{He}+\text{Au}$).

0-100% centrality, p_T integrated data well described:

- $1.2 < y < 2.2$: by Bayesian re-weighted shadowing alone
- $-2.2 < y < -1.2$: needs addition of absorption prediction from model

p_T dependence of 0-100% centrality data

- $1.2 < y < 2.2$: Bayesian re-weighted shadowing, needs p_T broadening
- $-2.2 < y < -1.2$: Bayesian re-weighted shadowing needs absorption
- Transport model does well at both rapidities

Centrality dependence of p+Au data compared with transport model

$1.2 < y < 2.2$: **very** strong increase with centrality, not described by EPS09

$-2.2 < y < -1.2$: Does well

Backup

J/ ψ in p+Pb at 5.02 TeV

ALICE, JHEP 1511 (2015) 127

Left: Negative rapidity

Right: Positive rapidity

Compared with calculations using
EPS09 shadowing (blue)

Energy loss in cold matter (red)

No absorption cross section at
LHC energy (crossing time).

- EPS09 shadowing works OK at backward rapidity.
- But underpredicts suppression at forward rapidity.

