

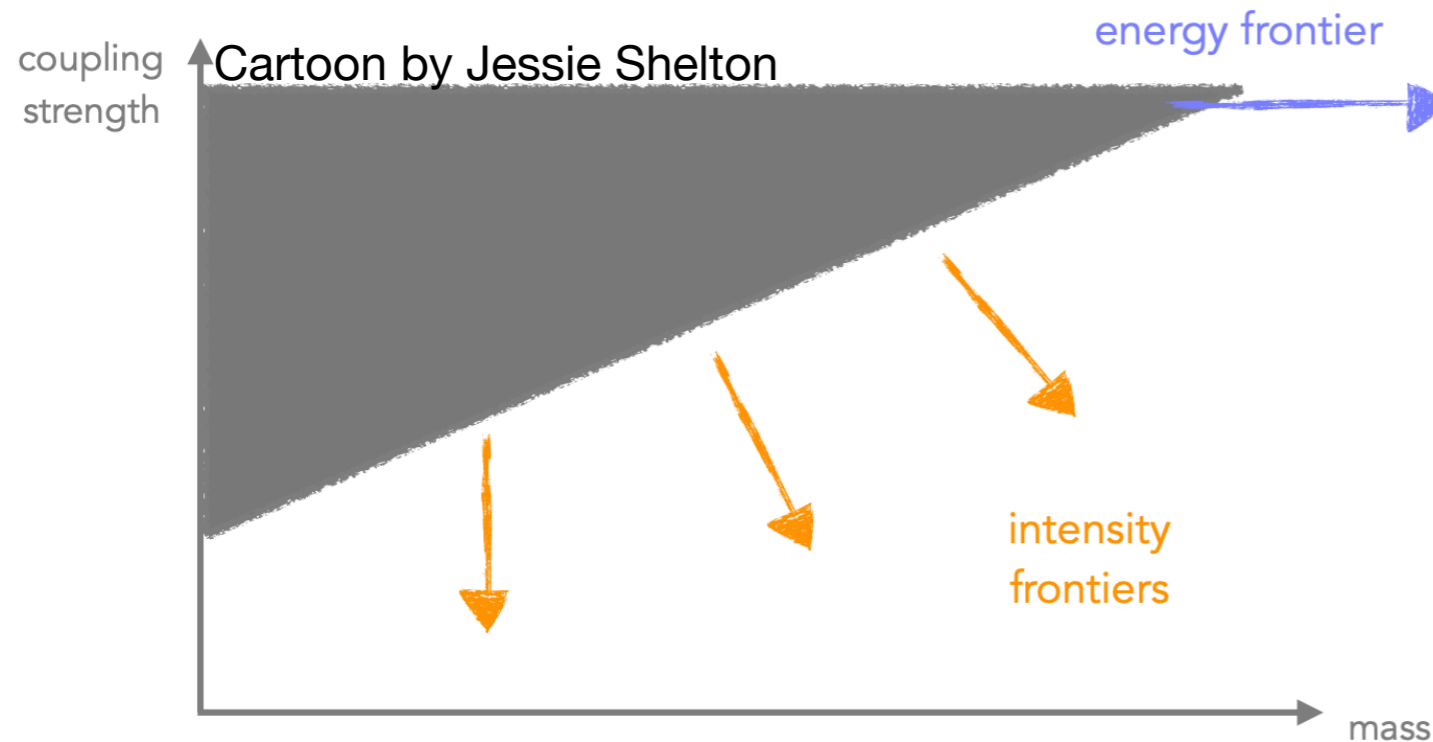
Exploring the lifetime frontier at the LHC and beyond

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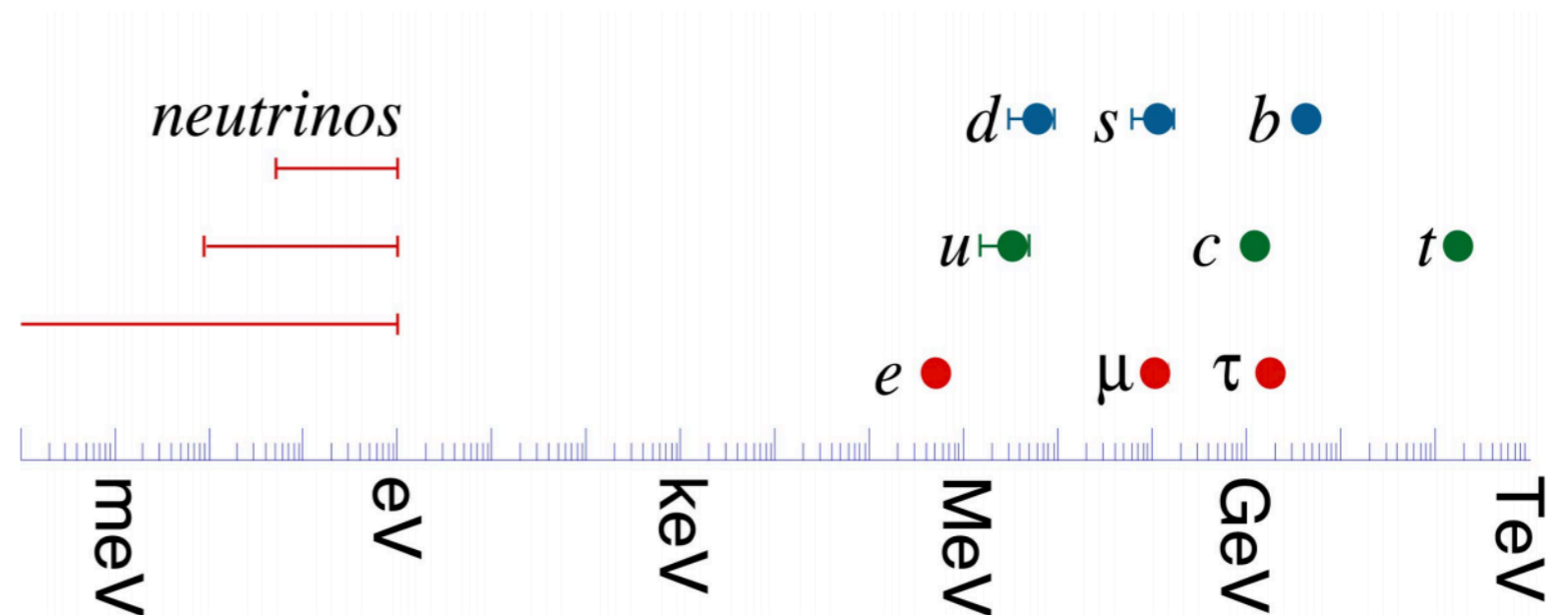
 @suchi_kulkarni

Based on: arXiv:1905.11889 and arXiv:1908.11741 with F. Deppisch, W. Liu



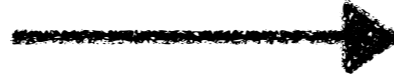
- Third dimension: lifetime!
- Not well explored at the LHC
- Long lived particle (LLP) searches: very little to no background

- LLP occur when either mass splitting between two particles is small or the coupling is suppressed
- Neutrino masses: Evidence of BSM physics
- Intimately related to new physics at lifetime frontier



Tiny Yukawa

- SM + right handed fields



- Small Dirac neutrino masses

Lepton number violation

- RH fields mass term



- Baryogenesis through leptogenesis

$$\mathcal{M}_\nu = \begin{pmatrix} \mathbf{0} & M_D \\ M_D^\top & M_N \end{pmatrix}$$

Dirac mass, usually small Y_ν

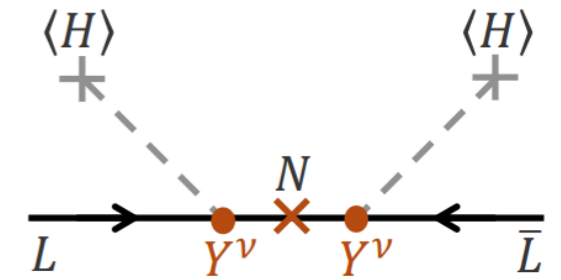
Majorana mass, can be heavy, can have $Y_N \sim 1$, introduces Lepton number violation

$$\mathcal{L} \supset Y_{ij}^\nu \bar{N}_i L_j \cdot H - \frac{1}{2} M_{ij} \bar{N}_i N_j^c$$

See e.g. Deppisch, *New J. Phys.* 17 (2015) 075019

- Let us go to one extreme $M_N \approx 10^{14}$ GeV (GUT scale)

$$\mathcal{L} \supset Y_{ij}^\nu \bar{N}_i L_j \cdot H - \frac{1}{2} M_{ij} \bar{N}_i N_j^c \xrightarrow{\langle H \rangle \ll M_N} \frac{1}{2} (Y_{ki}^\nu M_{kl}^{-1} Y_{lj}^\nu) (\bar{L}_i^c \cdot H) (H^T \cdot L_j)$$



- Light neutrino mass

$$m_\nu \approx 0.1 \text{ eV} \left(\frac{Y_\nu \langle H \rangle}{100 \text{ GeV}} \right)^2 \left(\frac{10^{14} \text{ GeV}}{M} \right)$$

- Heavy neutrino lifetime

$$L_N \approx 0.025 \text{ m} \cdot \left(\frac{10^{-6}}{V_{\mu N}} \right)^2 \cdot \left(\frac{100 \text{ GeV}}{m_N} \right)^5$$

- Sterile neutrino mass scale M_N unknown

- $\approx 10^{14}$ GeV Naive seesaw, GUTs
- $\approx 10^9$ GeV Thermal leptogenesis
- $\approx 10^3$ GeV Production at the LHC
- ≈ 1 keV Dark matter candidate
- ≈ 1 eV Oscillations, cosmology, $0\nu\beta\beta$

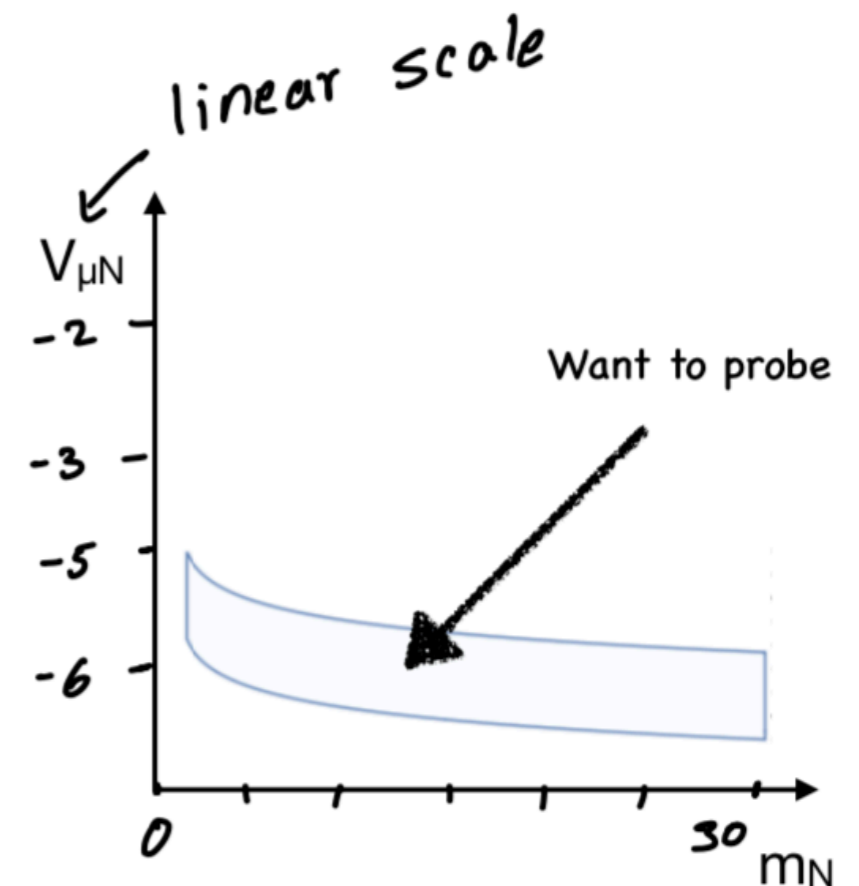
Region of interest for this talk

- Gauge group: $SU(3)_C \times SU(2)_L \times U(1)_Y \times U(1)_{B-L}$ Mohapatra, Marshak (PRL 44 (1980) 1316(319))
- Characteristics
 - **Particle content:** B-L gauge boson (Z'), Higgs boson (χ_{B-L}), 3 heavy neutrinos (N)
 - **Couplings:** g'_{B-L} (B-L coupling), $\sin\alpha$ (χ_{B-L} , Higgs mixing), V_{lN} (neutrino mixing)
 - **Free parameters:** 5 masses, 5 couplings (diagonal V_{lN})
 - Assume only light muon neutrino \rightarrow 3 masses, 3 couplings
 - **Charges:** χ : +2; N: -1; q: 1/3; l: -1

$$\mathcal{M} = \begin{pmatrix} 0 & m_D \\ m_D & M_R \end{pmatrix} \quad m_\nu \approx -\frac{M_D^2}{M_R} = -V_{lN}^2 M_R$$

- Heavy neutrino lifetime

$$L_N \approx 0.025 \text{ m} \cdot \left(\frac{10^{-6}}{V_{\mu N}} \right)^2 \cdot \left(\frac{100 \text{ GeV}}{m_N} \right)^5$$



Higgs sector

$$\mathcal{L} \supset (D^\mu H)^\dagger (D_\mu H) + (D^\mu \chi)^\dagger D_\mu \chi - \mathcal{V}(H, \chi),$$

$$\mathcal{V}(H, \chi) = m^2 H^\dagger H + \mu^2 |\chi|^2 + \lambda_1 (H^\dagger H)^2 + \lambda_2 |\chi|^4 + \lambda_3 H^\dagger H |\chi|^2$$

$$D_\mu = \partial_\mu + ig_s \mathcal{T}_\alpha G_\mu^\alpha + ig T_a W_\mu^a + ig_1 Y B_\mu + i(\tilde{g} Y + g'_1 Y_{B-L}) B'_\mu$$

Kinetic term

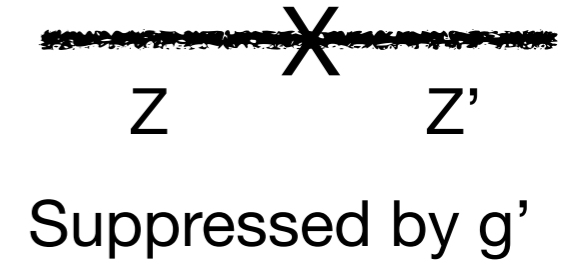
$$\mathcal{L} \supset -\frac{1}{4} F'^{\mu\nu} F'_{\mu\nu} \quad \text{Abelian hyper-charge and B-L mixing terms set to zero}$$

Right handed neutrino term

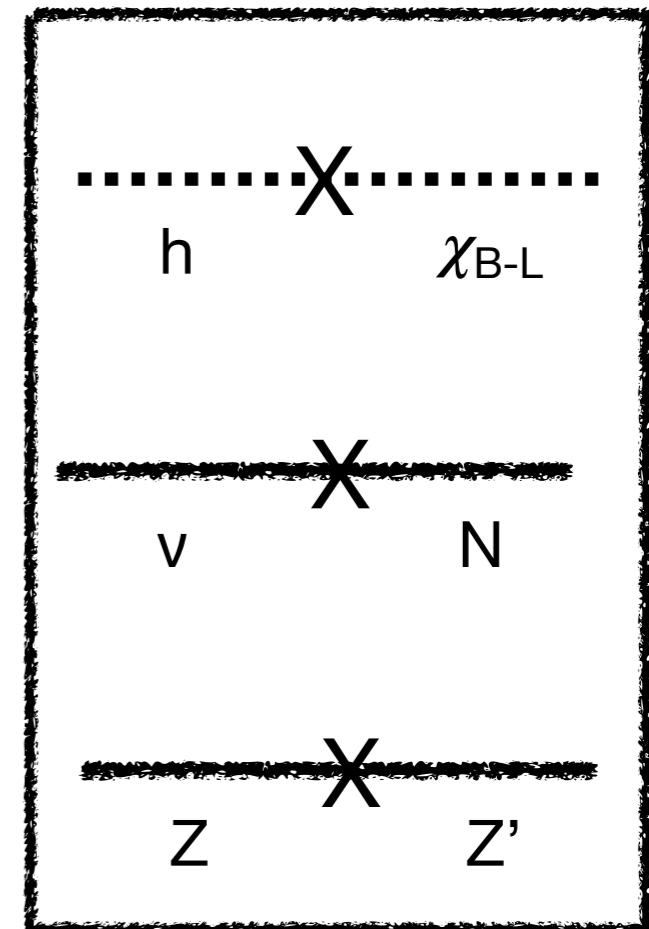
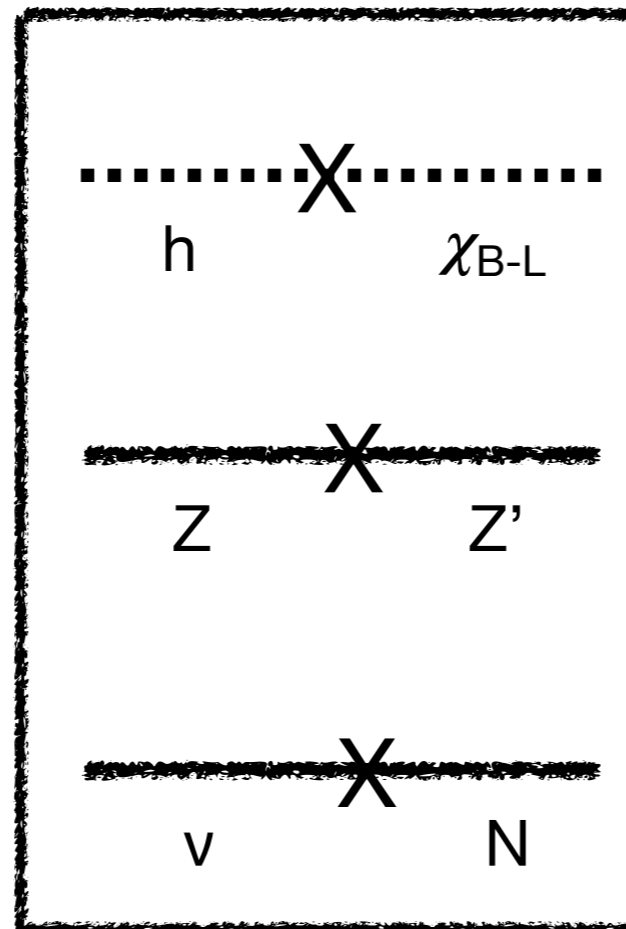
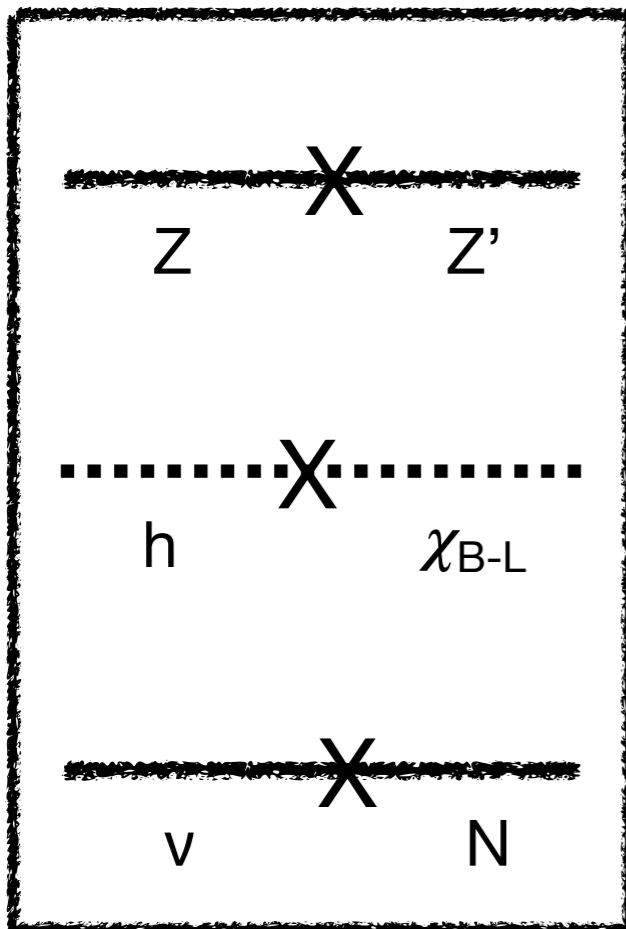
$$\mathcal{L} \supset i \bar{\nu}_{Ri} \gamma_\mu D^\mu \nu_{Ri}$$

Additional Yukawa terms

$$\mathcal{L} \supset -y_{ij}^\nu \bar{L}_i \nu_{Rj} \tilde{H} - y_{ij}^M \bar{\nu}_{Ri}^c \nu_{Rj} \chi + \text{h.c.}$$



Different B-L mass hierarchies lead to different phenomenology





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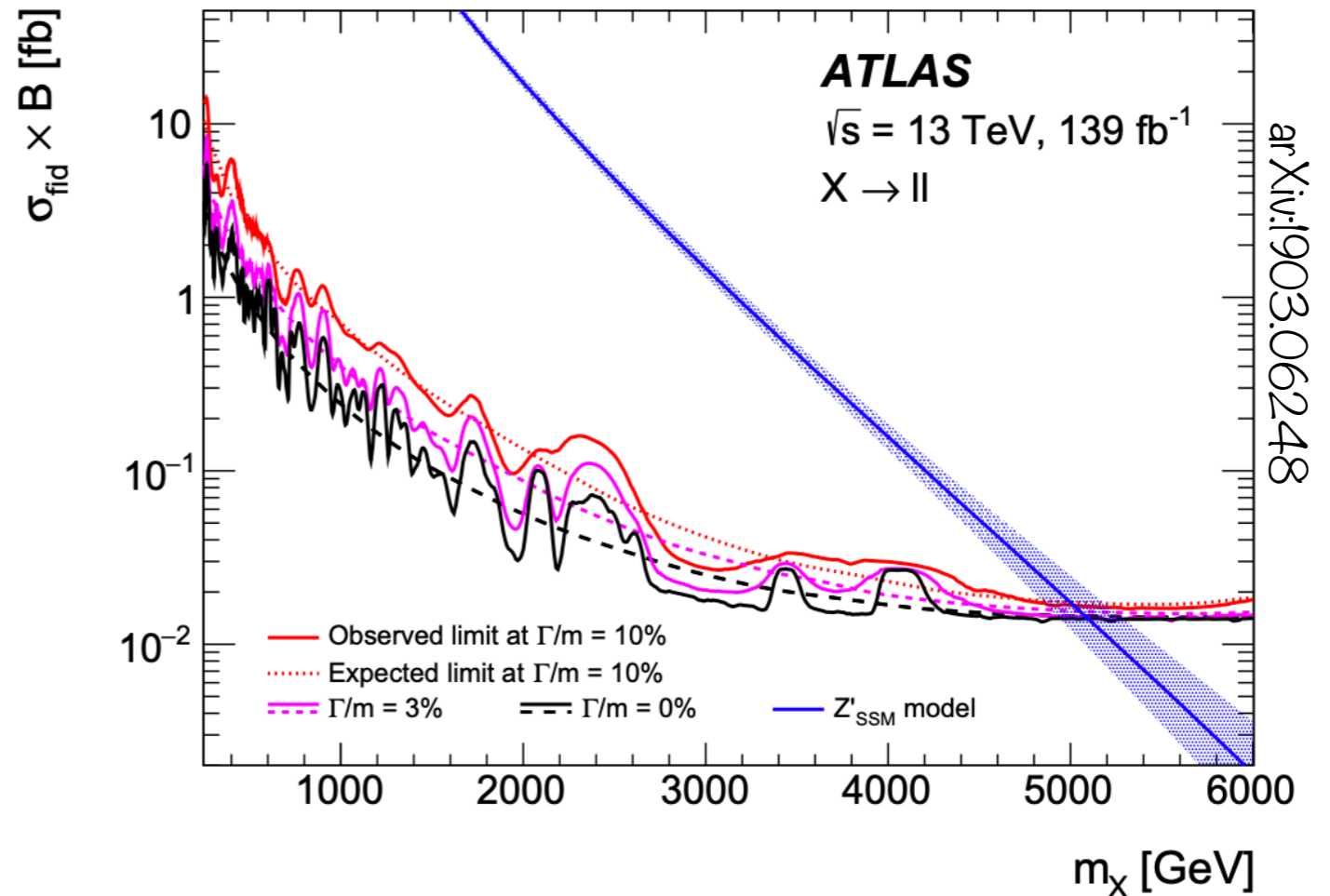
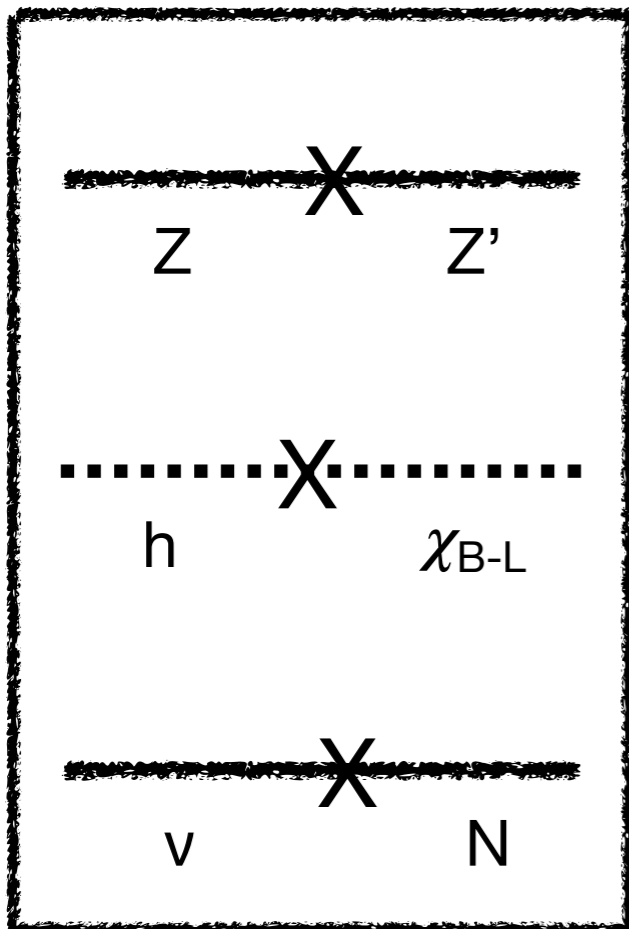


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$M_{Z'}$ must at least be 125 GeV \rightarrow strong constraints from dilepton searches





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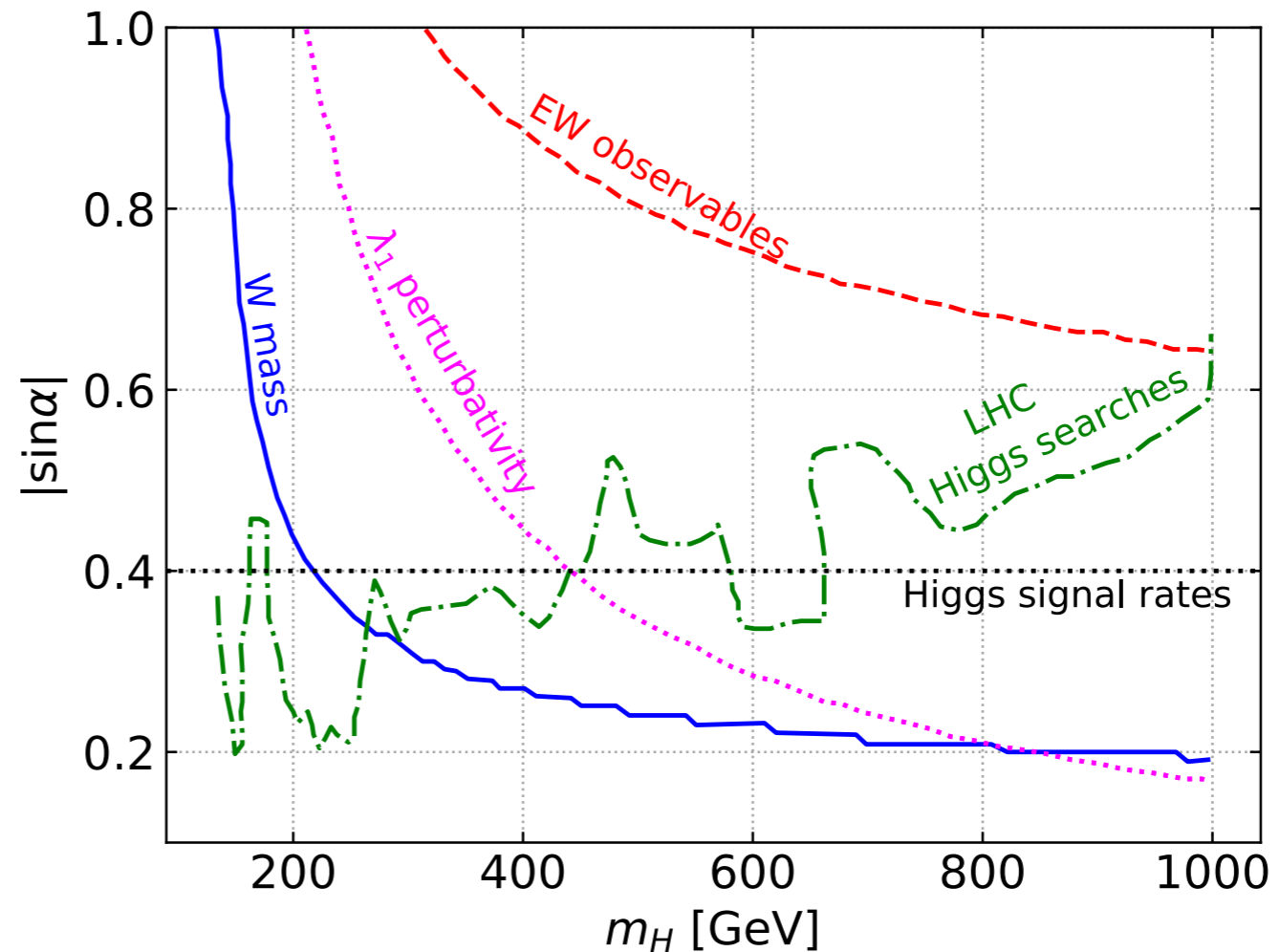
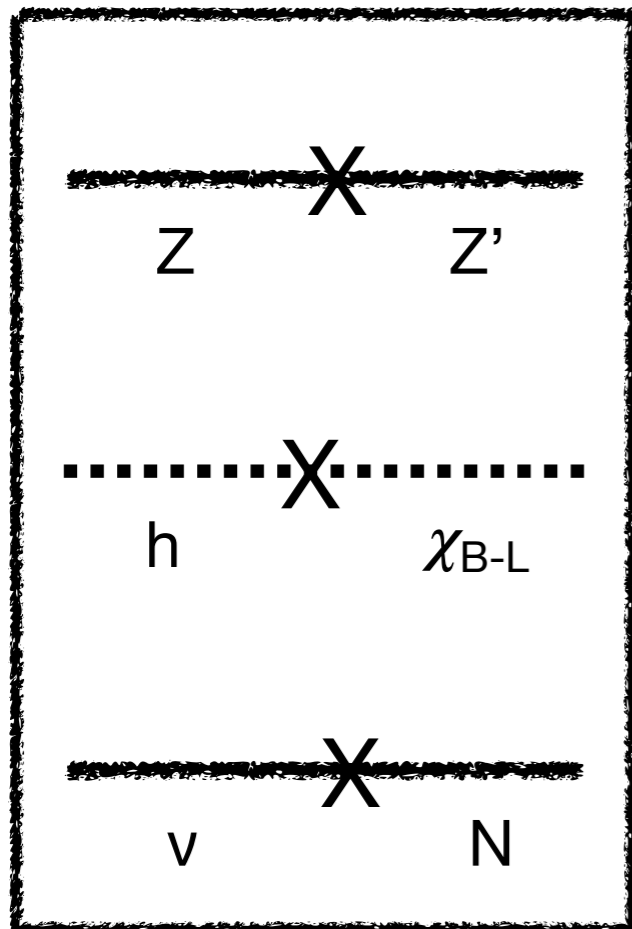


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Constraints from heavy Higgs searches, EW observables and theory considerations





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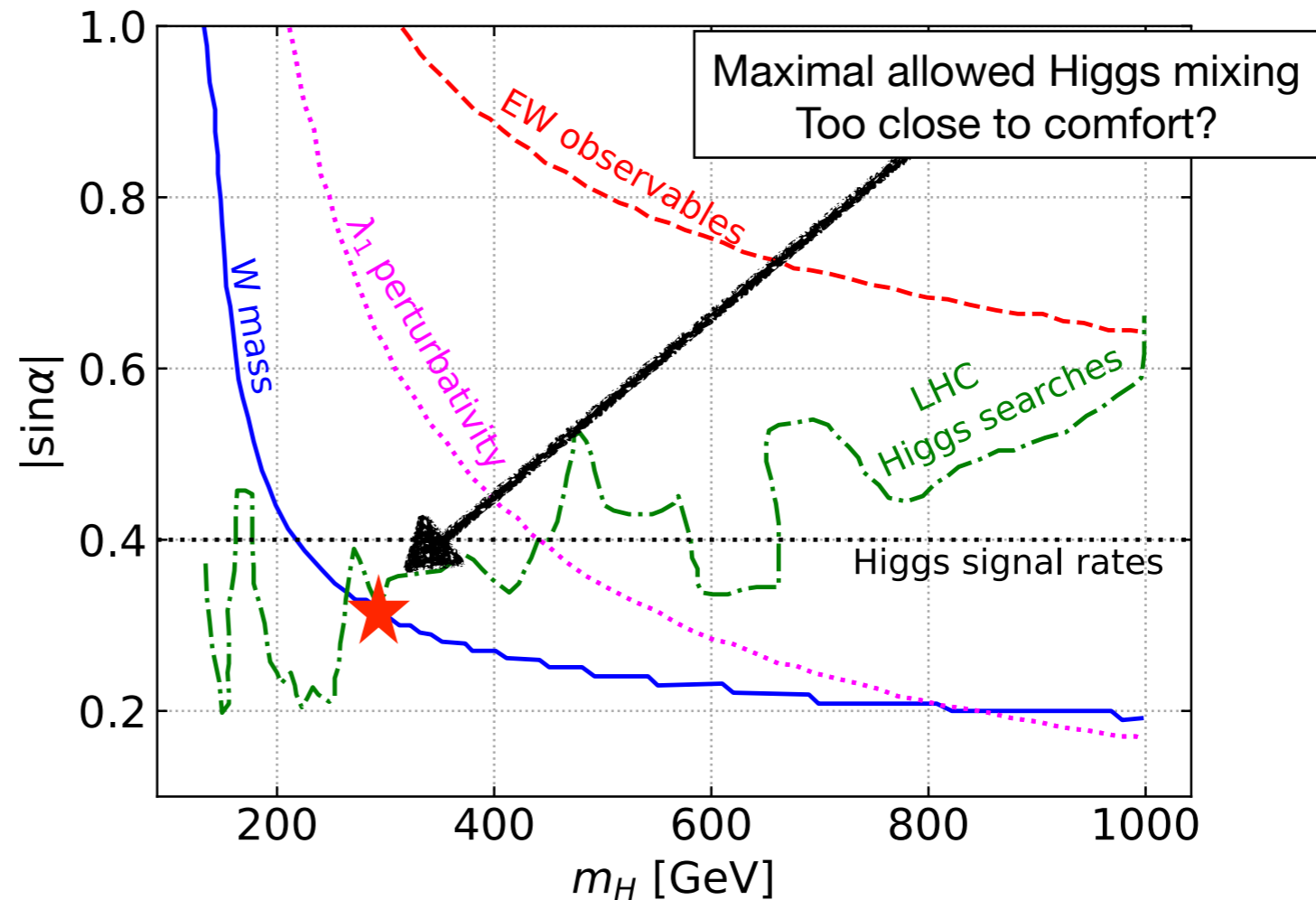
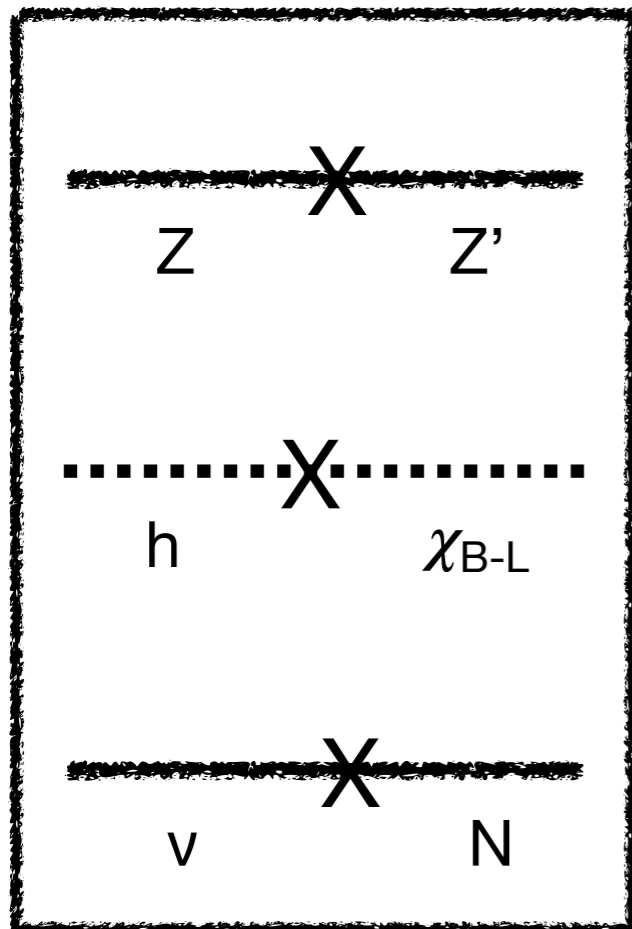


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Constraints from heavy Higgs searches, EW observables and theory considerations





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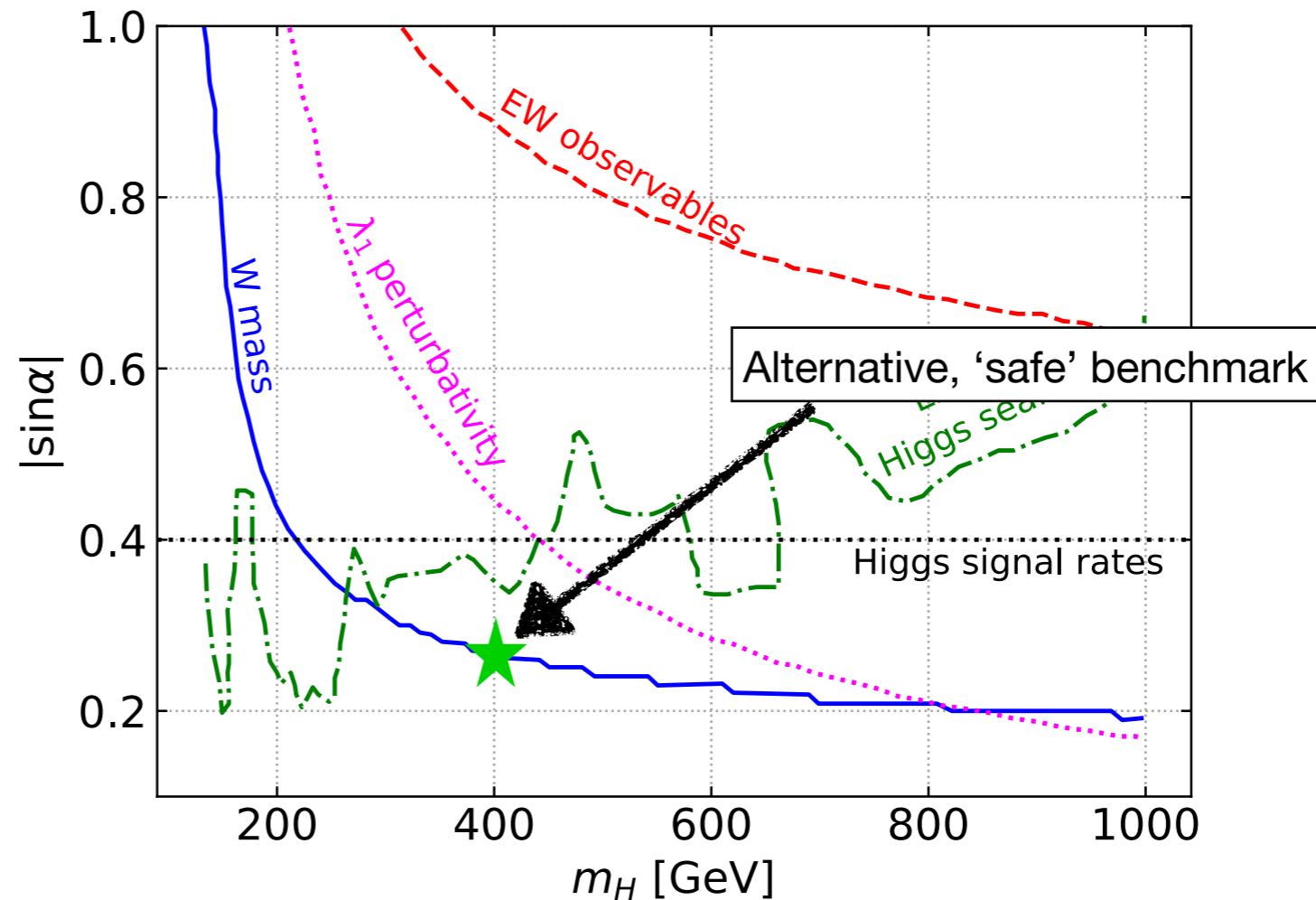
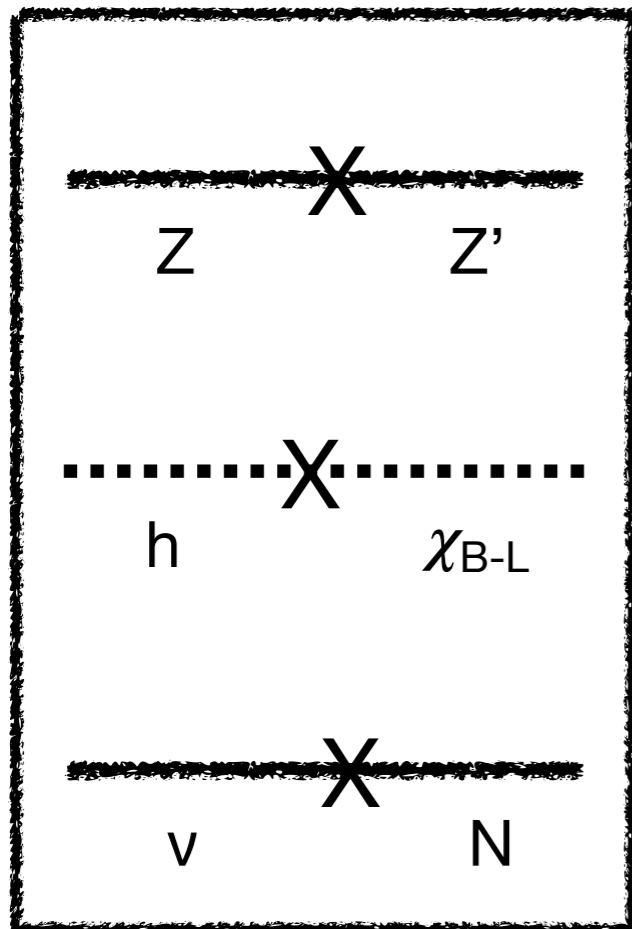


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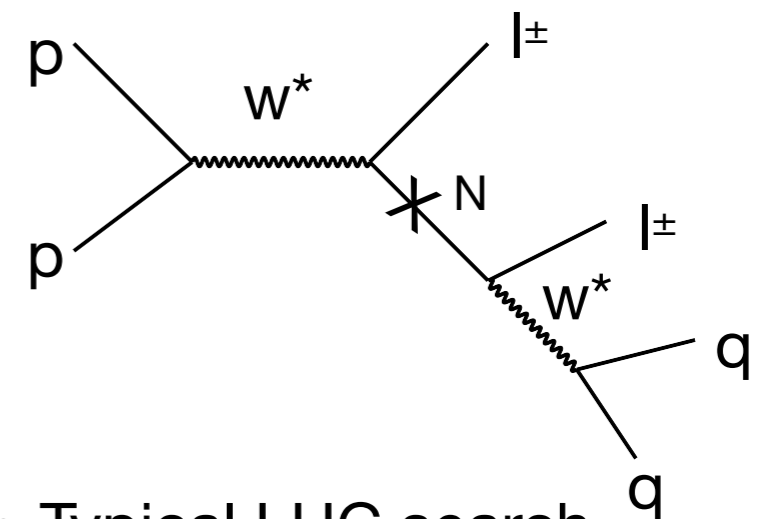
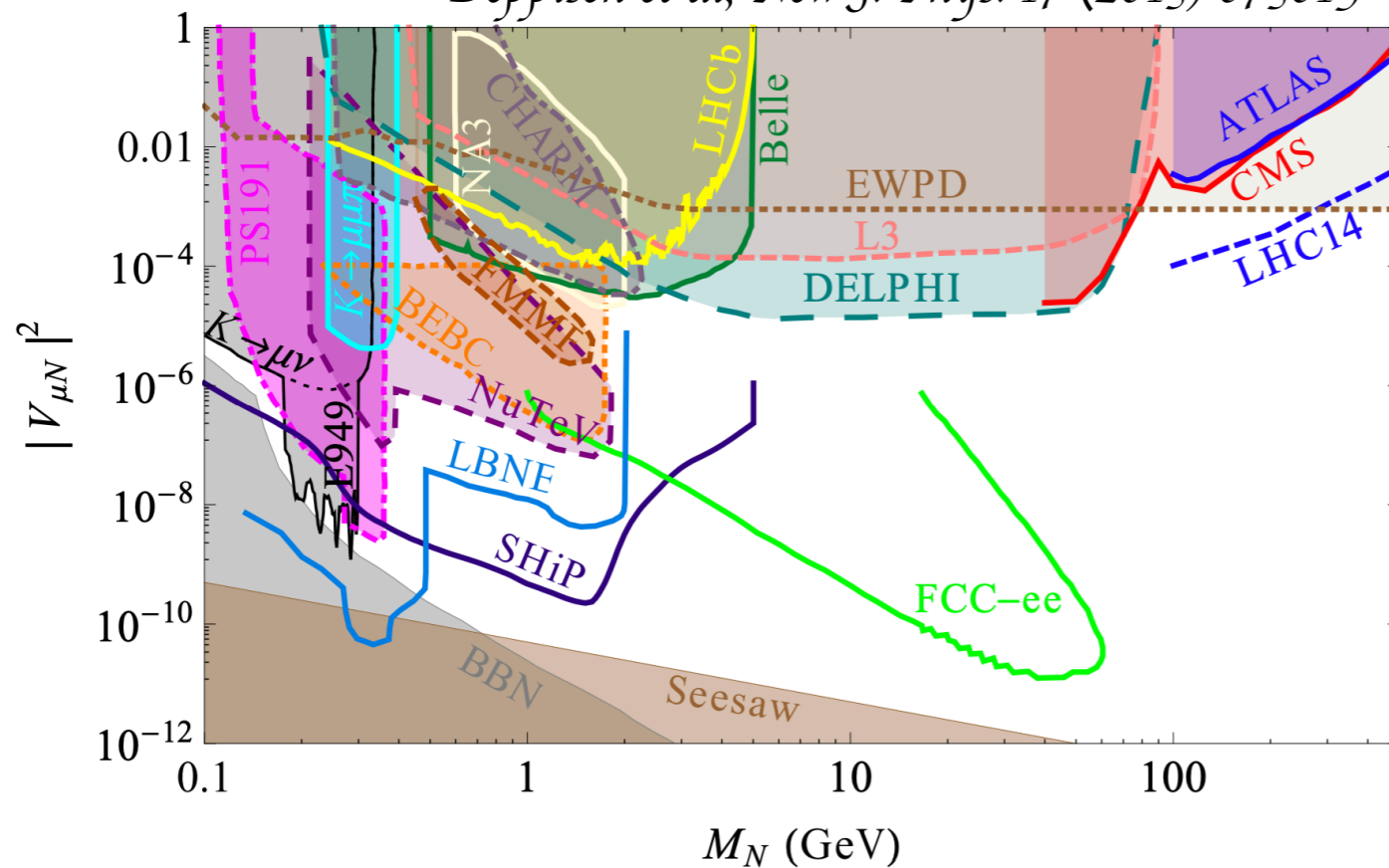


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Constraints from heavy Higgs searches, EW observables and theory considerations



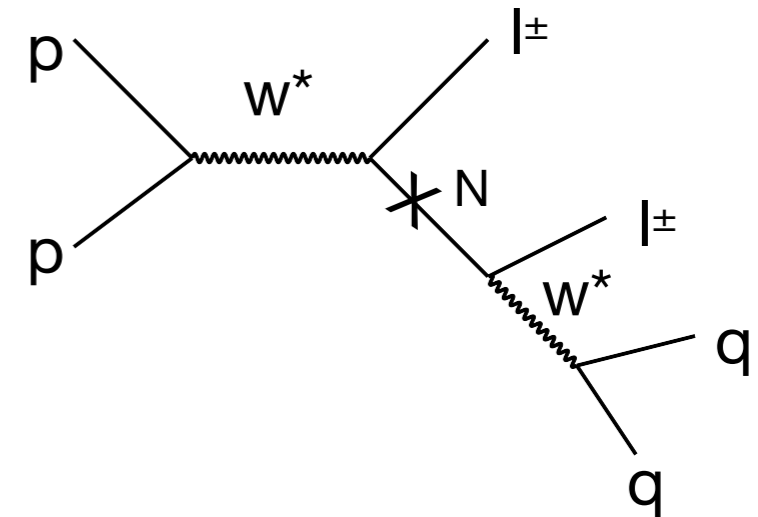
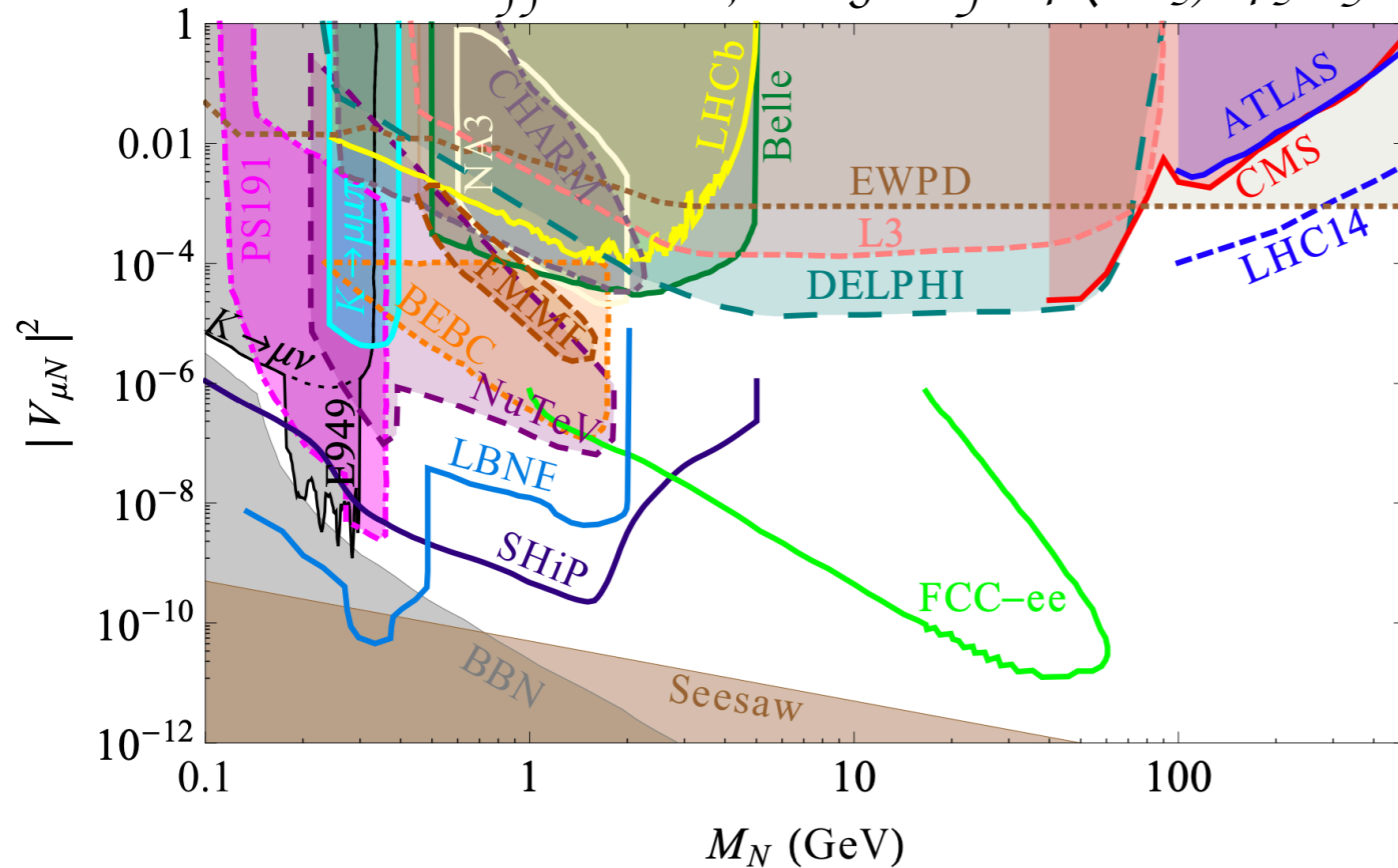
Deppisch et al, *New J. Phys.* 17 (2015) 075019



- Typical LHC search
- Minimal process
- Prompt lepton requirement

- Several probes for heavy neutrinos (aka heavy neutral leptons; HNL)
- Intensity frontier typically covers low masses; small mixing angles
- Limit plot corresponds to HNL production via SM mediators, B-L charges not taken into account
- At the LHC, same sign leptons from production (decays) of heavy neutrinos via SM W boson
- Current LHC limits weak, rapidly changing situation

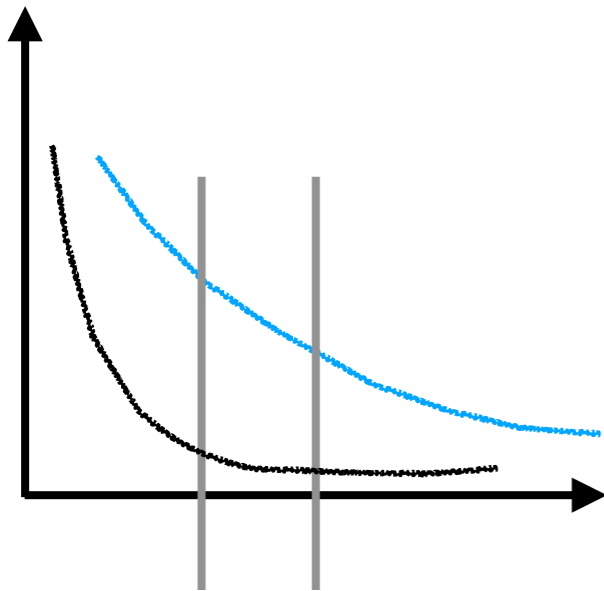
Deppisch et al, *New J. Phys.* 17 (2015) 075019



- Heavy neutrino decay length

$$L_N \approx 0.025 \text{ m} \cdot \left(\frac{10^{-6}}{V_{\mu N}} \right)^2 \cdot \left(\frac{100 \text{ GeV}}{m_N} \right)^5$$

- Build FCC!
- Heavy neutrino phenomenology at LHC necessitates exploration at lifetime frontier
- For neutrino masses ≈ 100 's of GeV, $V_{\mu N} \approx 10^{-6} \rightarrow L_N \approx 25\text{mm}$
- Problem: $V_{\mu N}$ suppression \rightarrow have higher luminosity
- Neutrino mass of order 10 GeV, $V_{\mu N} \approx 10^{-6}$, $L_N \approx 100\text{m}$, decays outside of the LHC \rightarrow build bigger detectors; several proposals exist



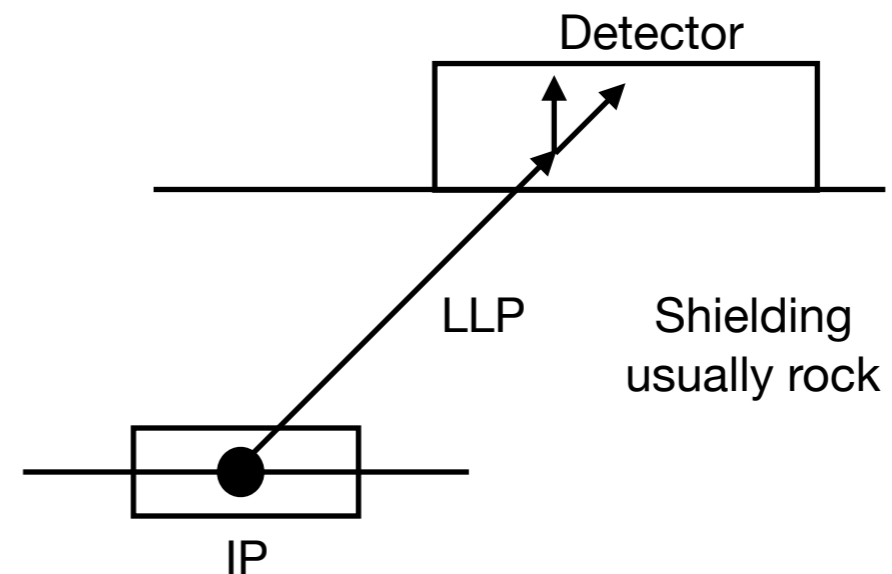
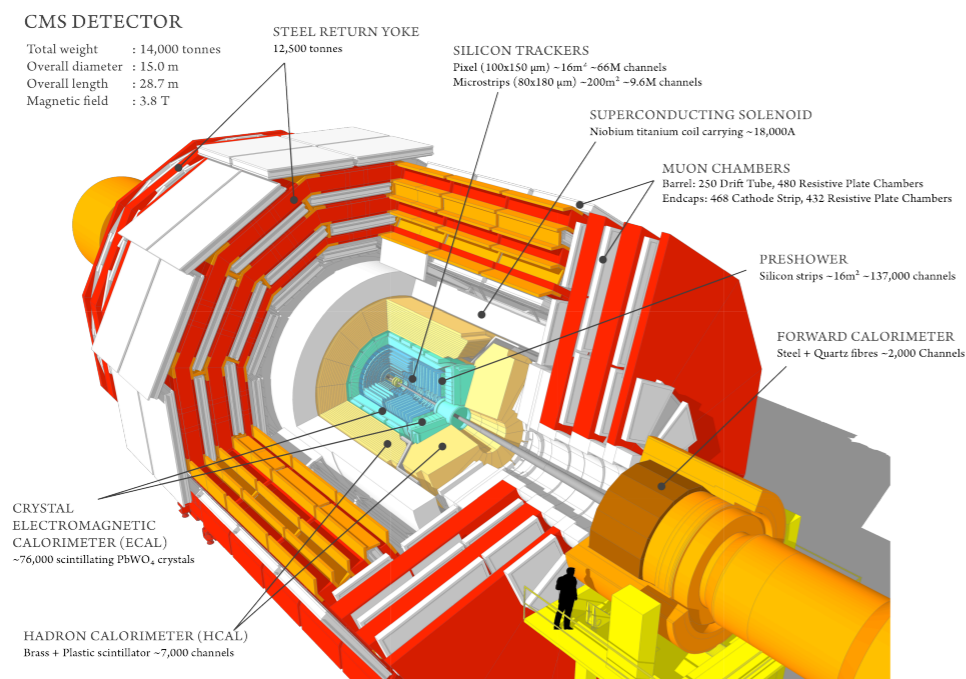
$$P_{\text{decay}}(bc\tau, L_1, L_2) = e^{-\frac{L_1}{bc\tau}} - e^{-\frac{L_2}{bc\tau}}$$

$$\approx \frac{L_2 - L_1}{bc\tau} \quad \text{for } (L_2 - L_1) \ll bc\tau$$

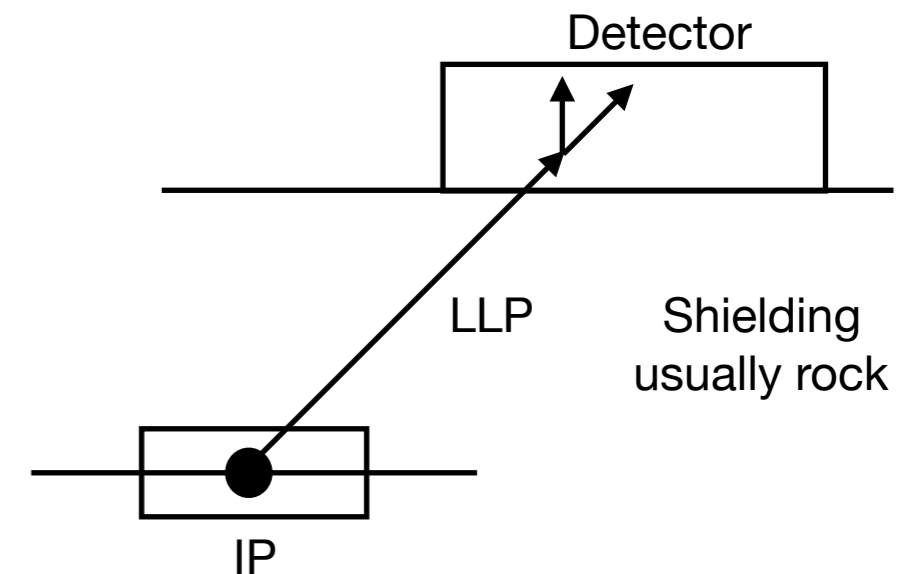
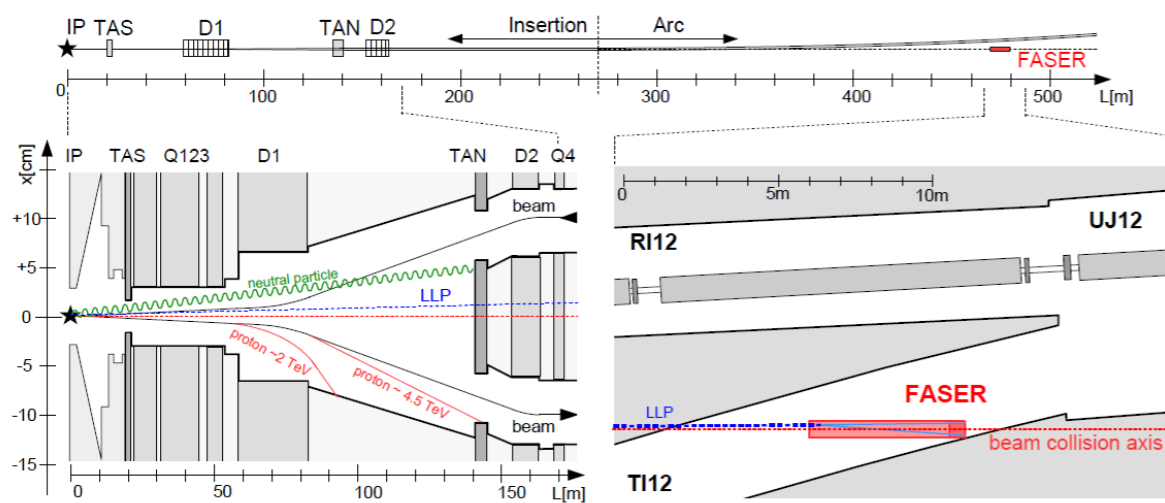
- Boost depends on production mechanism and mass hierarchy between progenitor and decay product

$$N_{\text{obs}} \approx (\sigma_{\text{sig}}^{LHC} \mathcal{L}) \epsilon_{LLP}^{\text{detector}} n_{LLP} \epsilon_{\text{geometric}} P_{\text{decay}}(\bar{bc}\tau, L_1, L_2)$$

- Geometric acceptance depends on the distance and geometry of the detector



Detector	Location	Distance from IP (m)	Dimensions (m)	Luminosity (fb ⁻¹)
FASER-2	ATLAS	480	Cylinder 5 X 1	3000
CODEX-b	LHC cavity	3	10 X 10 X 10	300
MAPP	LHCb/ MoEDAL	50	7 - 10 tunnel 5 - 25 degrees angle	300
MATHUSLA	CMS	100	200 X 200 X 20	3000



Approved experiment



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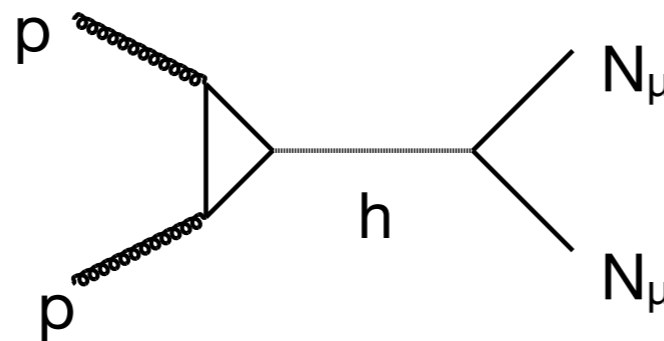
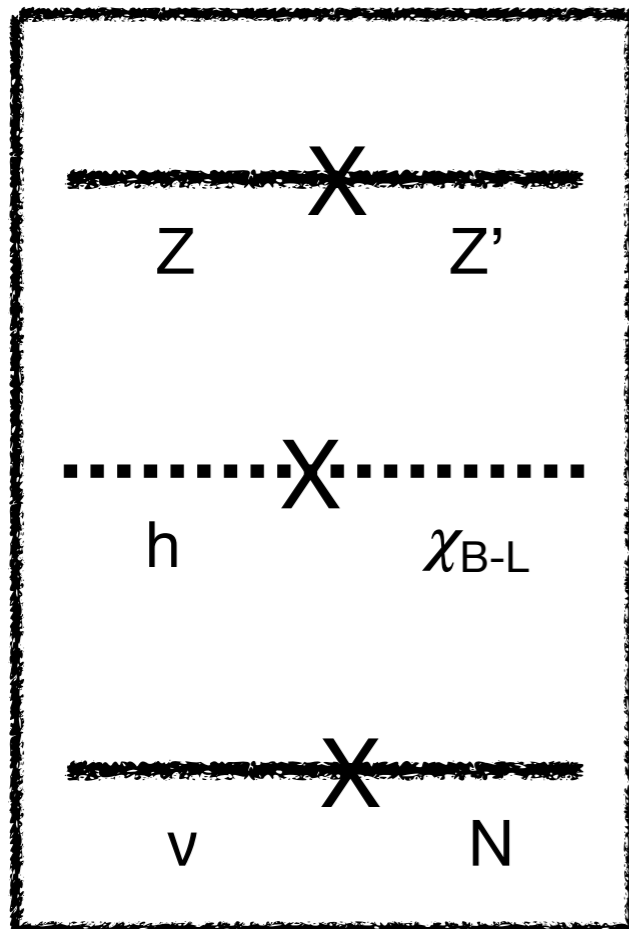
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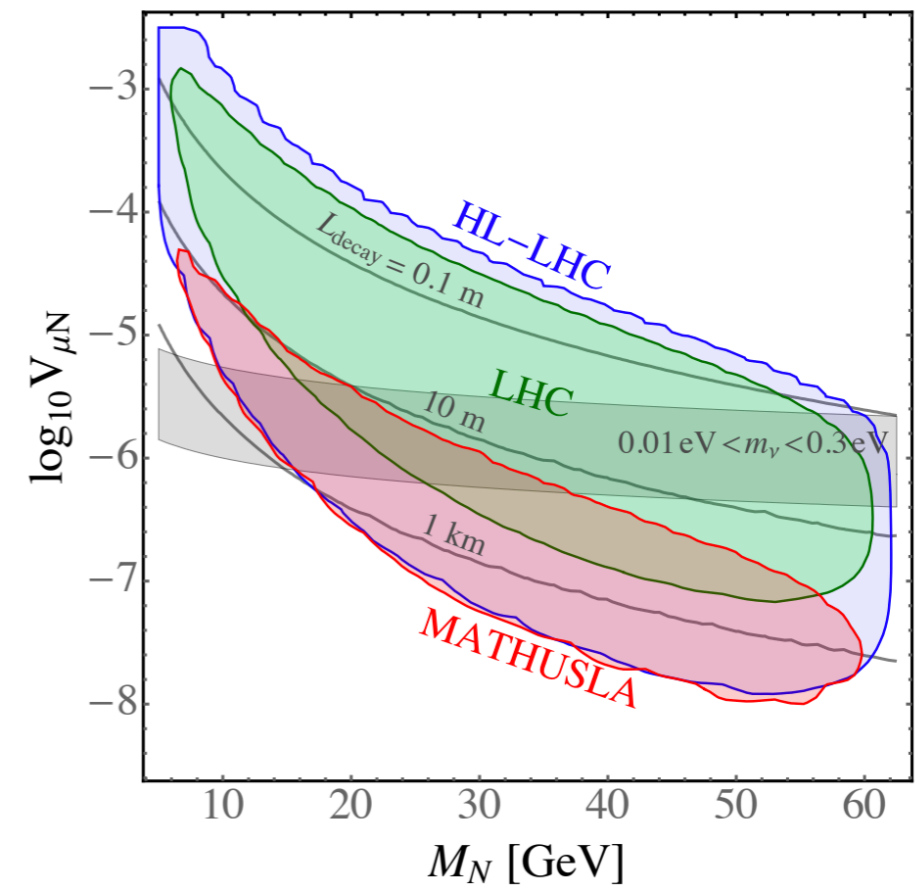
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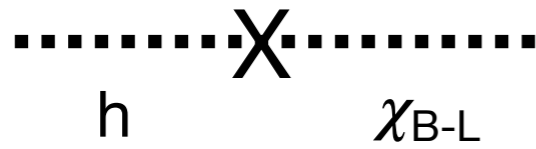
$h \rightarrow N N$ still possible if $\sin\alpha$ is large

$\sin\alpha = 0.3, g' \sim 0$

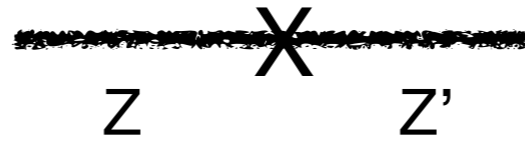


Deppisch et al, JHEP 1808 (2018) 181

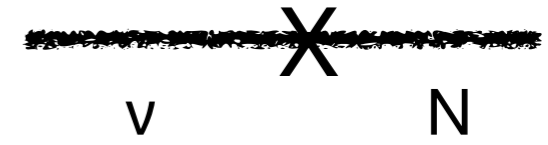




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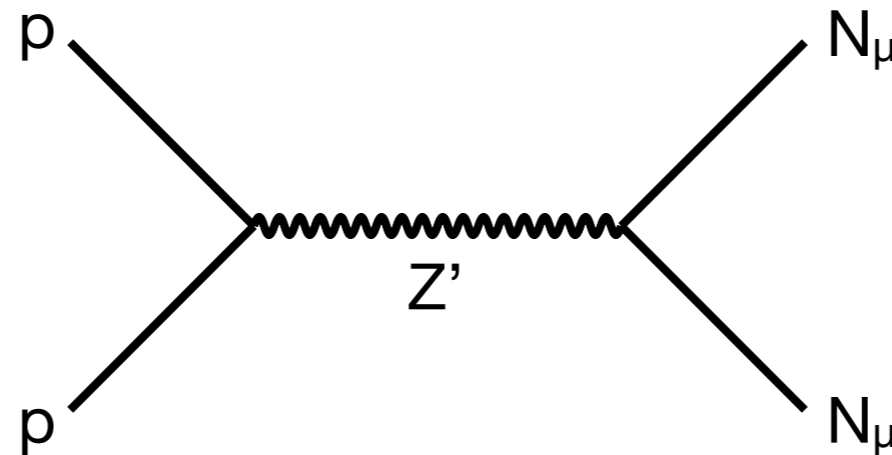
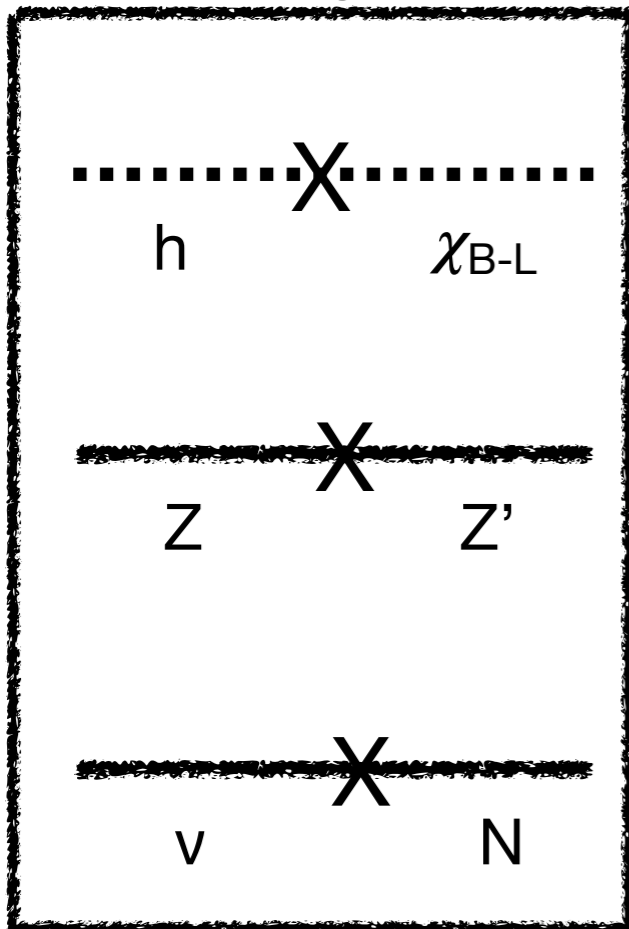


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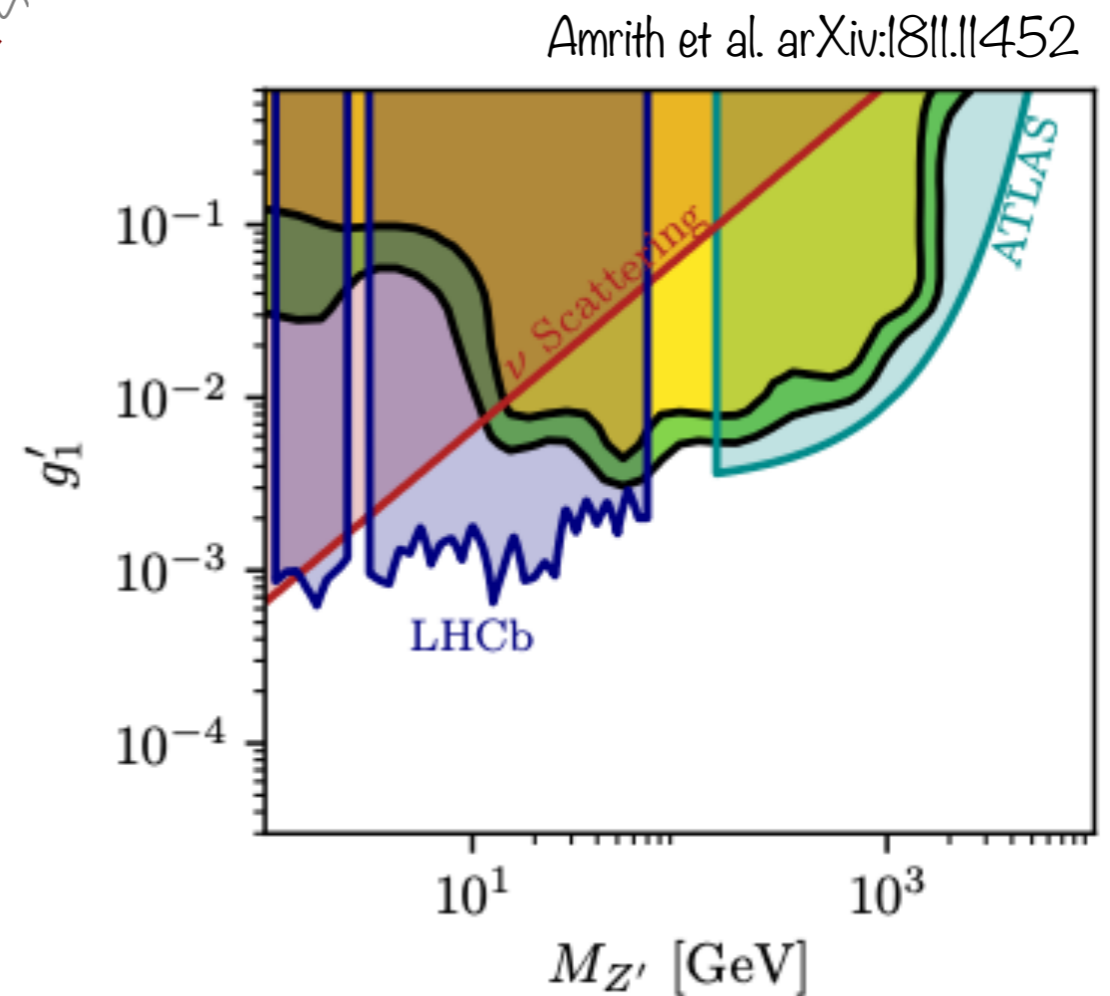
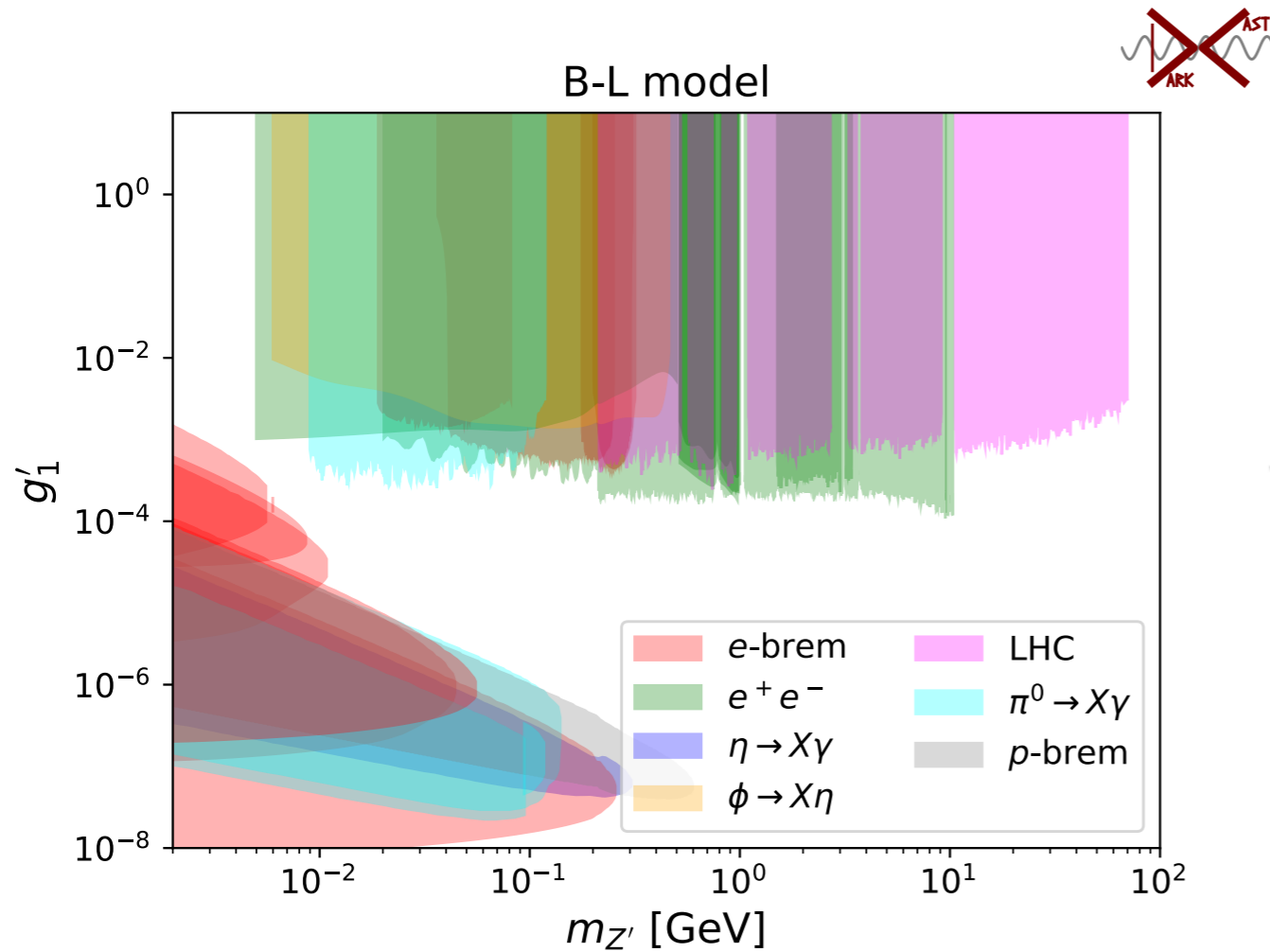
$\sin\alpha = 0, g' = 10^{-3}$

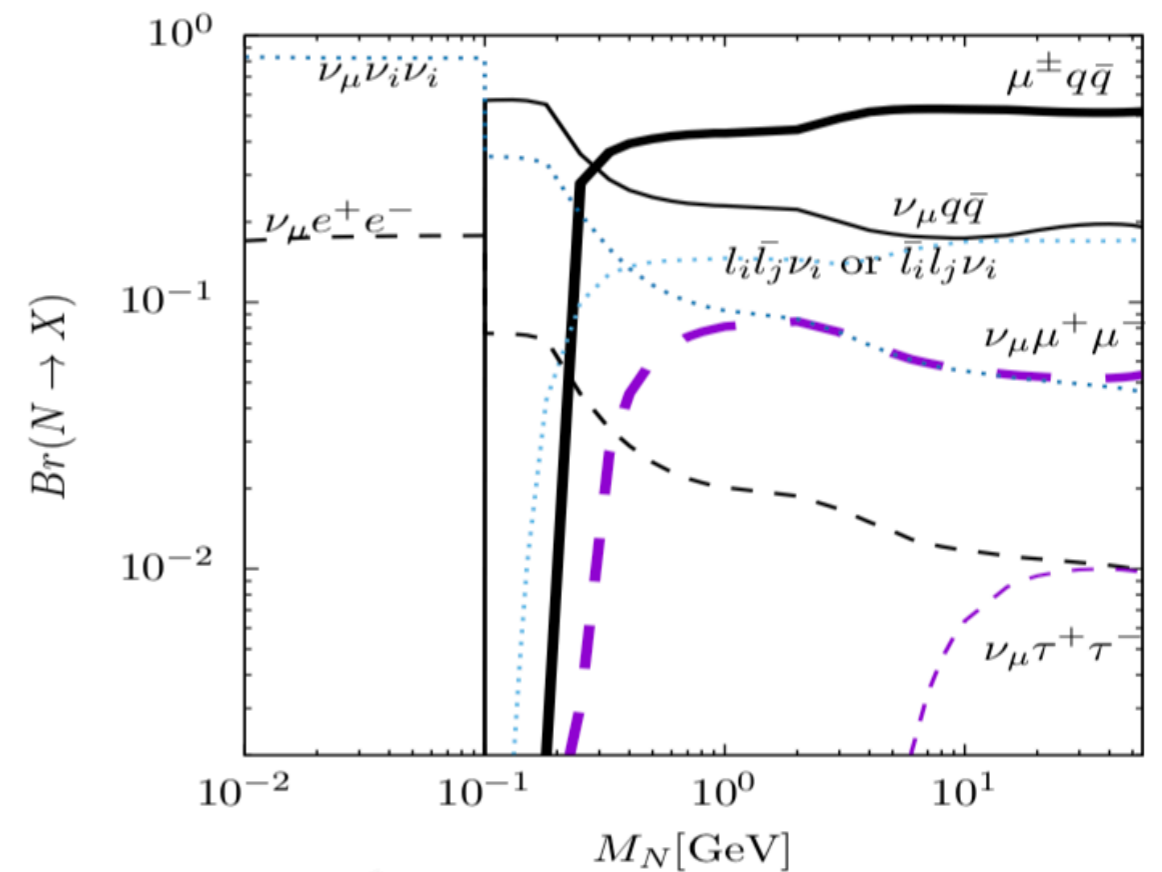
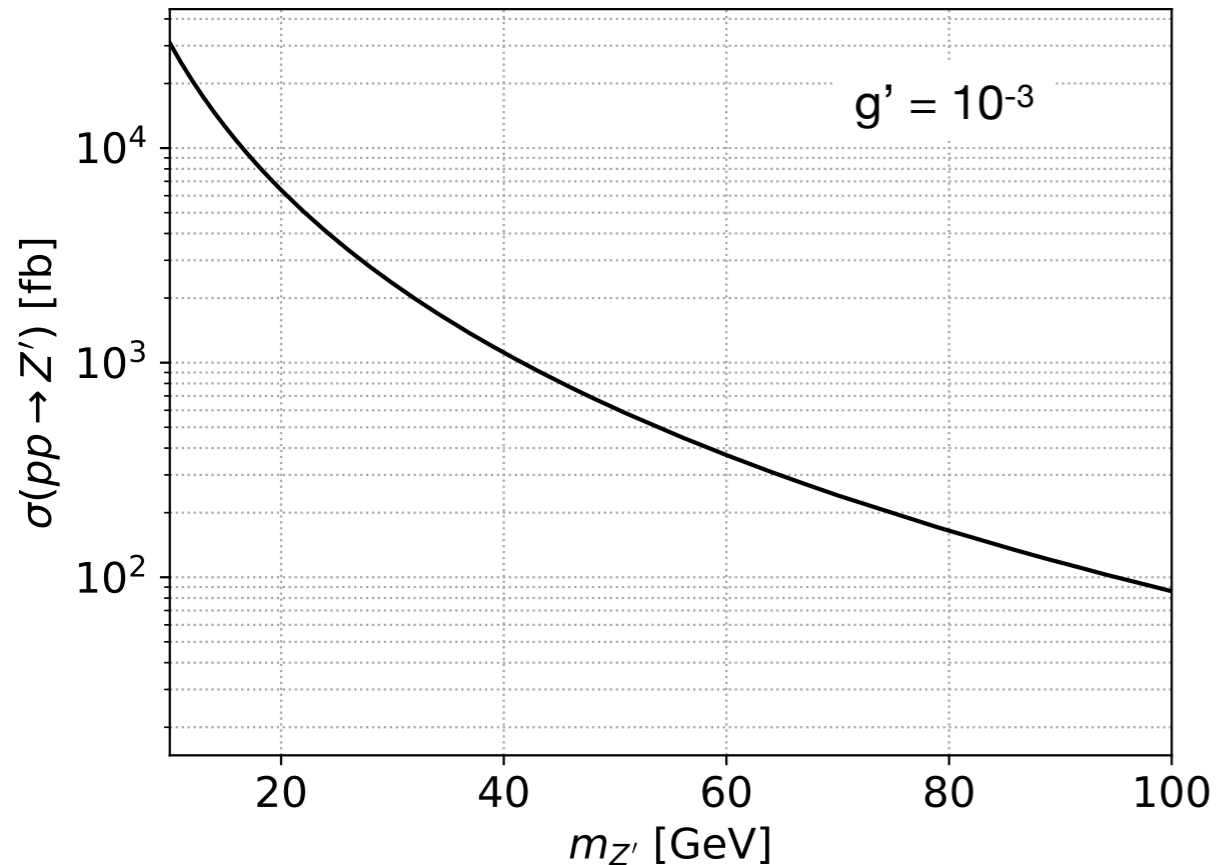


$M_{Z'} < M_h \rightarrow M_{Z'} < 100 \text{ GeV};$

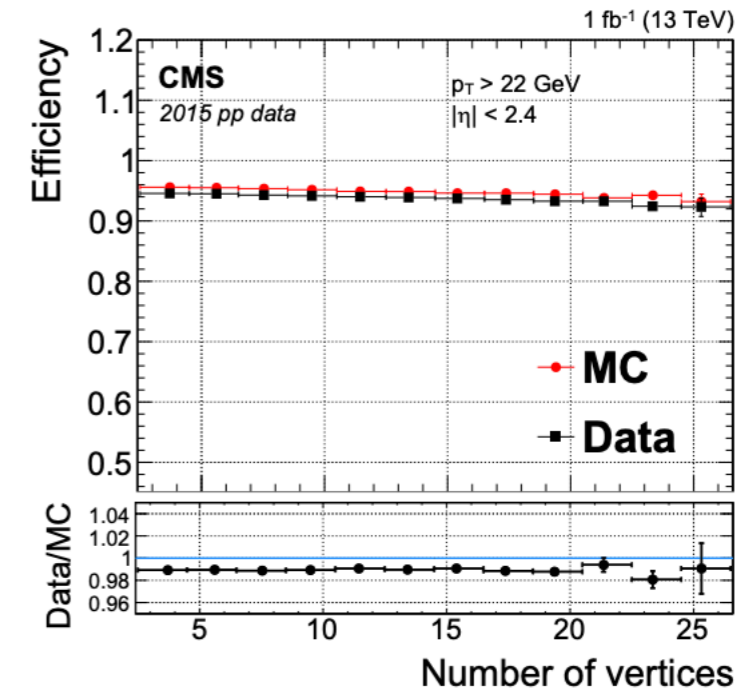
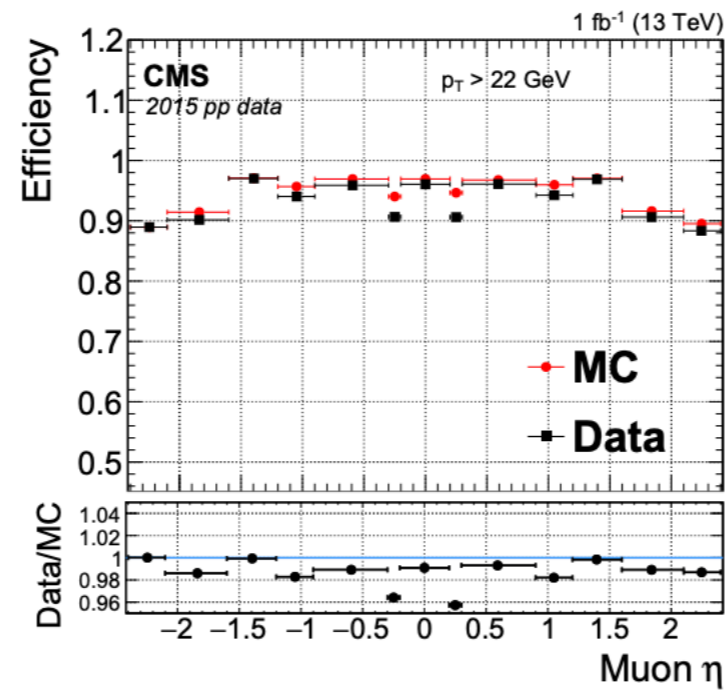
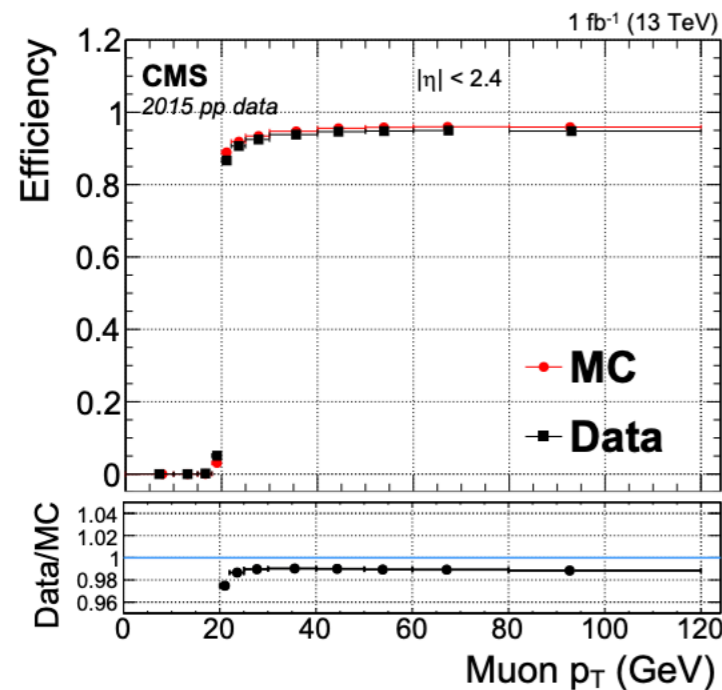
What do the experiments have to say about this mass range?

- Current constraints on the Z' masses come from LHC as well as fixed target experiments
- Reinterpretation of LHC SM searches done via CONTUR





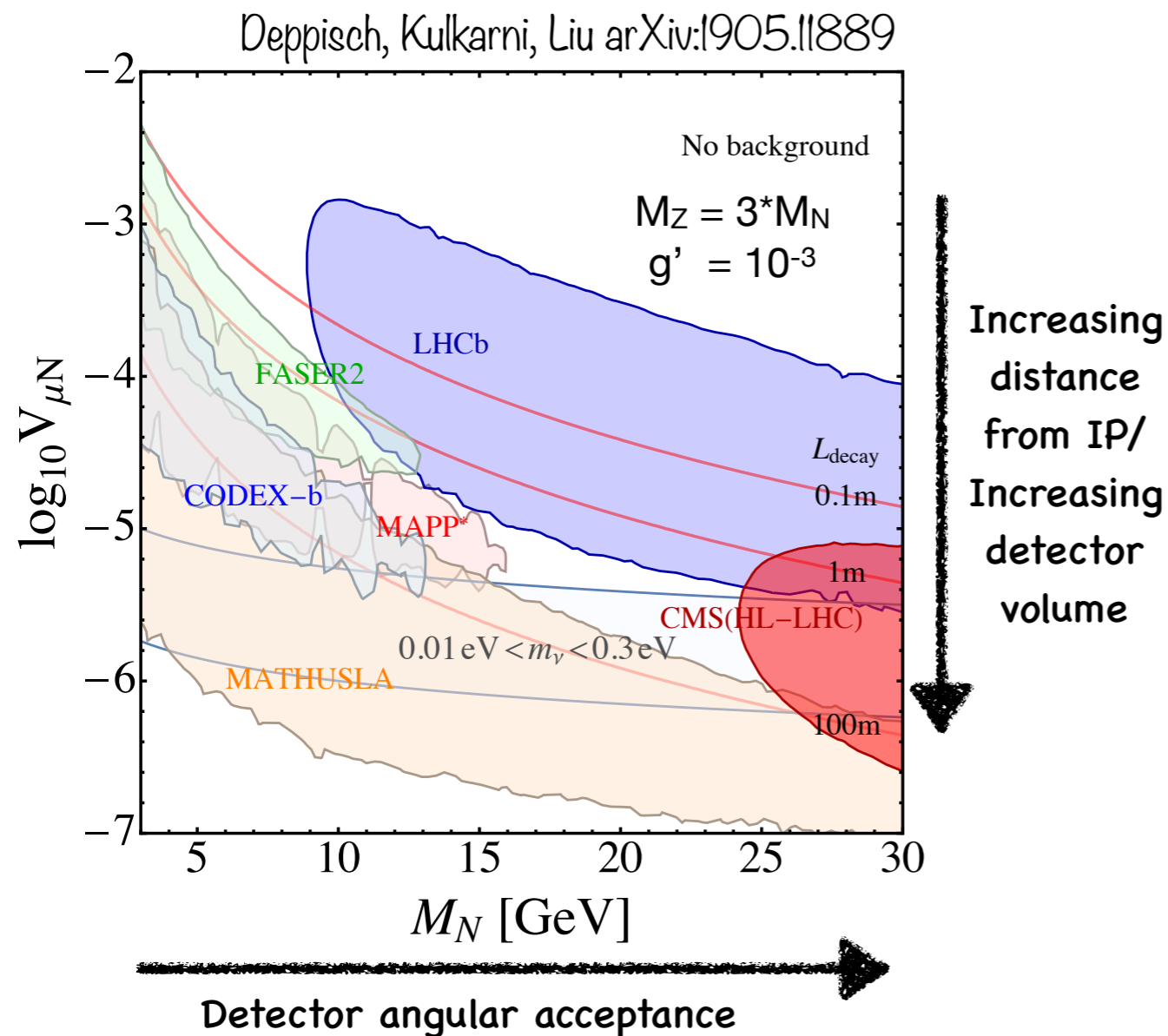
- Significant Z' production cross section
- $BR(Z' \rightarrow N N) \sim 8\%$ (per specie)
- Branching ratio N to at least one muon final state between 10 to 30%
- Typical HNL decay lengths ($V_{\mu N} \approx 10^{-6}$): O(100) m
- Potential for good reach in neutrino mixing angle at future facilities
- Concentrate on final state containing muons



arXiv:1804.04528.

- Electrons: absorbed in calorimeter (Electrons in muon system can appear as noise)
- Taus: decay, penalty due to tau branching fraction, challenging at LHC
- Jets: Not good resolution, larger trigger requirements compared to leptons
- Muons:
 - Excellent efficiencies for a large part
 - Lower background compared to hadrons
 - NB: situation is different for highly displaced muons

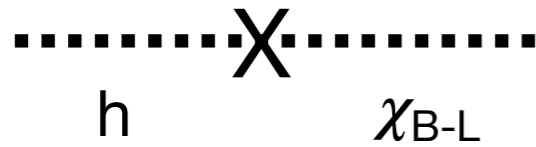
- Model implementation in MadGraph
- Decays and hadronization using Pythia8
- No detector simulation
- Final states under considerations:
 - LHCb: $p p \rightarrow N N \rightarrow N \mu j j$
 - General purpose LHC: $p p \rightarrow N N \rightarrow N \mu \mu \nu$
 - MATHUSLA, FASER, CODEX-b: $p p \rightarrow N N$
- Cuts:
 - $L_z = 480\text{m}$, $L_d = 1.5\text{m}, 5\text{m}$, $R = 1\text{m}, 5\text{m}$ [FASER]
 - $L_z \sim 30 - 60\text{m}$, $L_x \sim 4 - 15\text{m}$, $L_y \sim -10 - 10\text{m}$ [MAPP*]
 - $L_x = -100 \sim 100\text{m}$ $L_y = 100 \sim 120\text{m}$, $L_z = 100 \sim 300\text{m}$ [MATHUSLA]



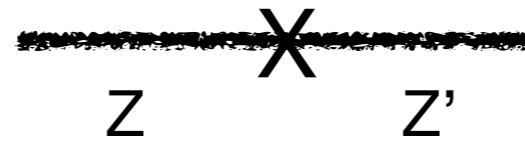
- For LHCb, use μjj final state; CMS $\mu \mu \nu$
- For other detector any final state allowed
- Look at the decay of only one heavy neutrino
- Apply some minimal cuts on the p_T and $|\eta|$ of final state particles
- Assume all neutrino decays within the detector volume are detected
- Nice interplay of boost and lifetime
- Detector of maximal interest: MATHUSLA
- ATLAS/CMS trigger requirements too high

B-L portals

Deppisch, Kulkarni, Liu arXiv:1908.11741



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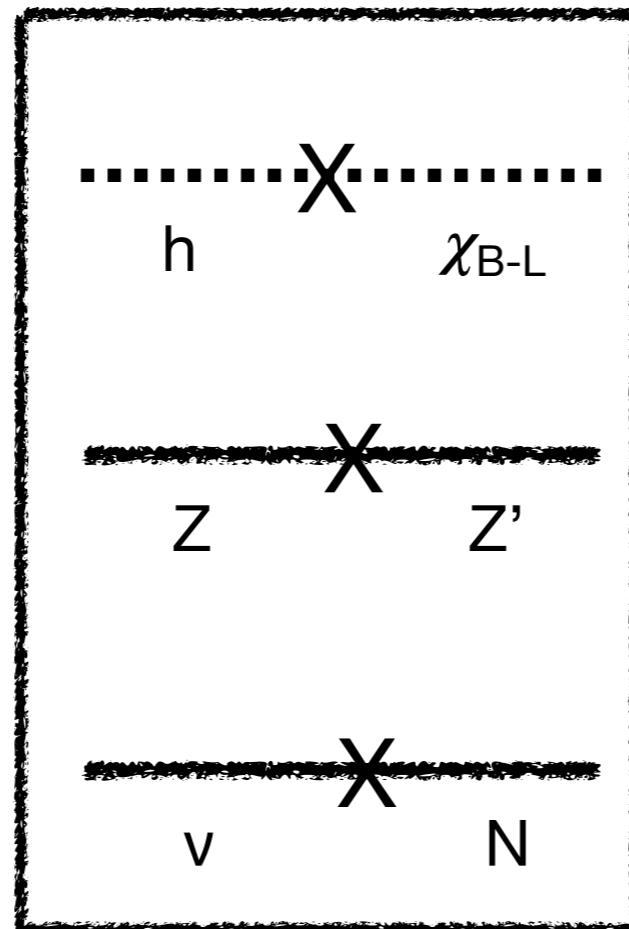


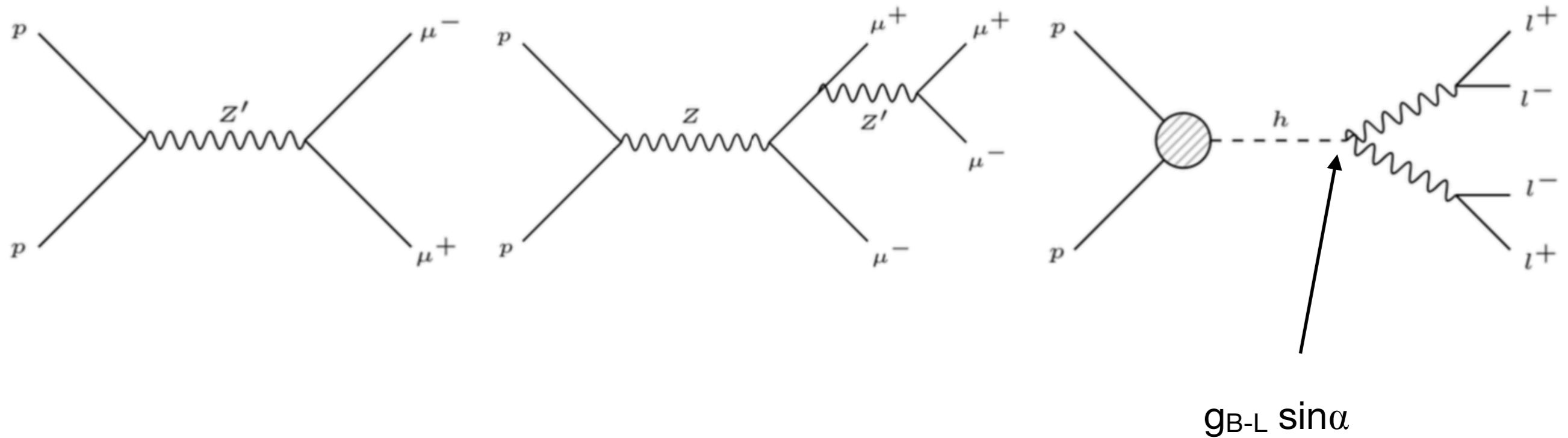
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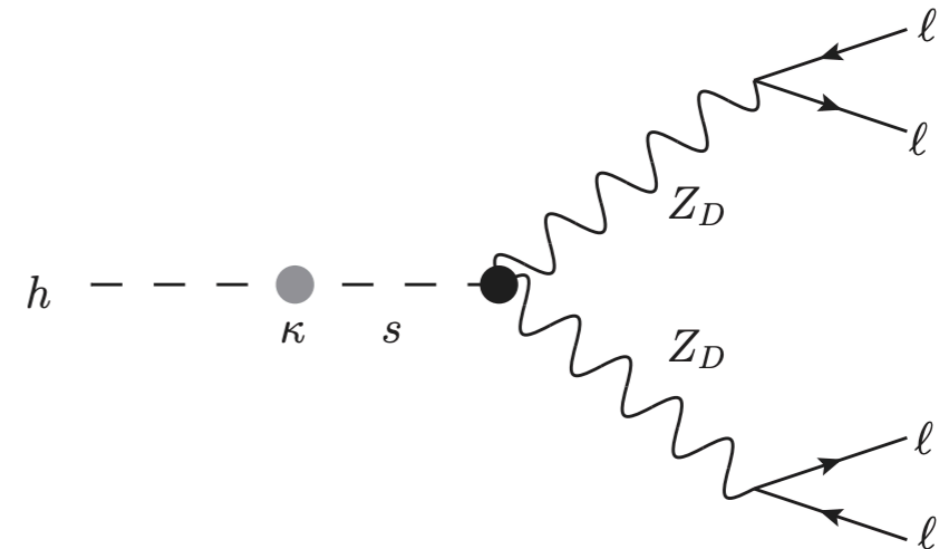
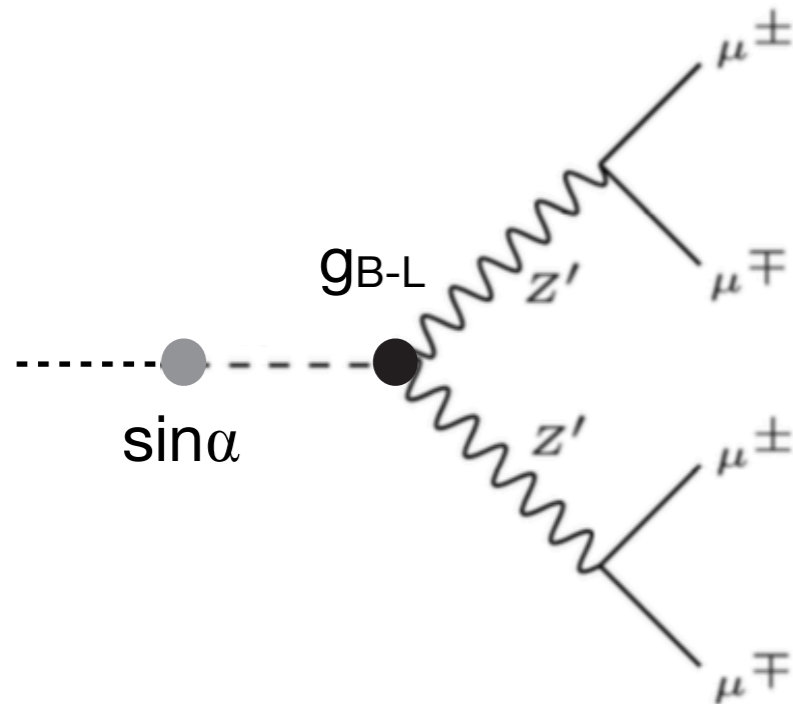
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$$\sin\alpha = 0.3, g' = 10^{-3}$$

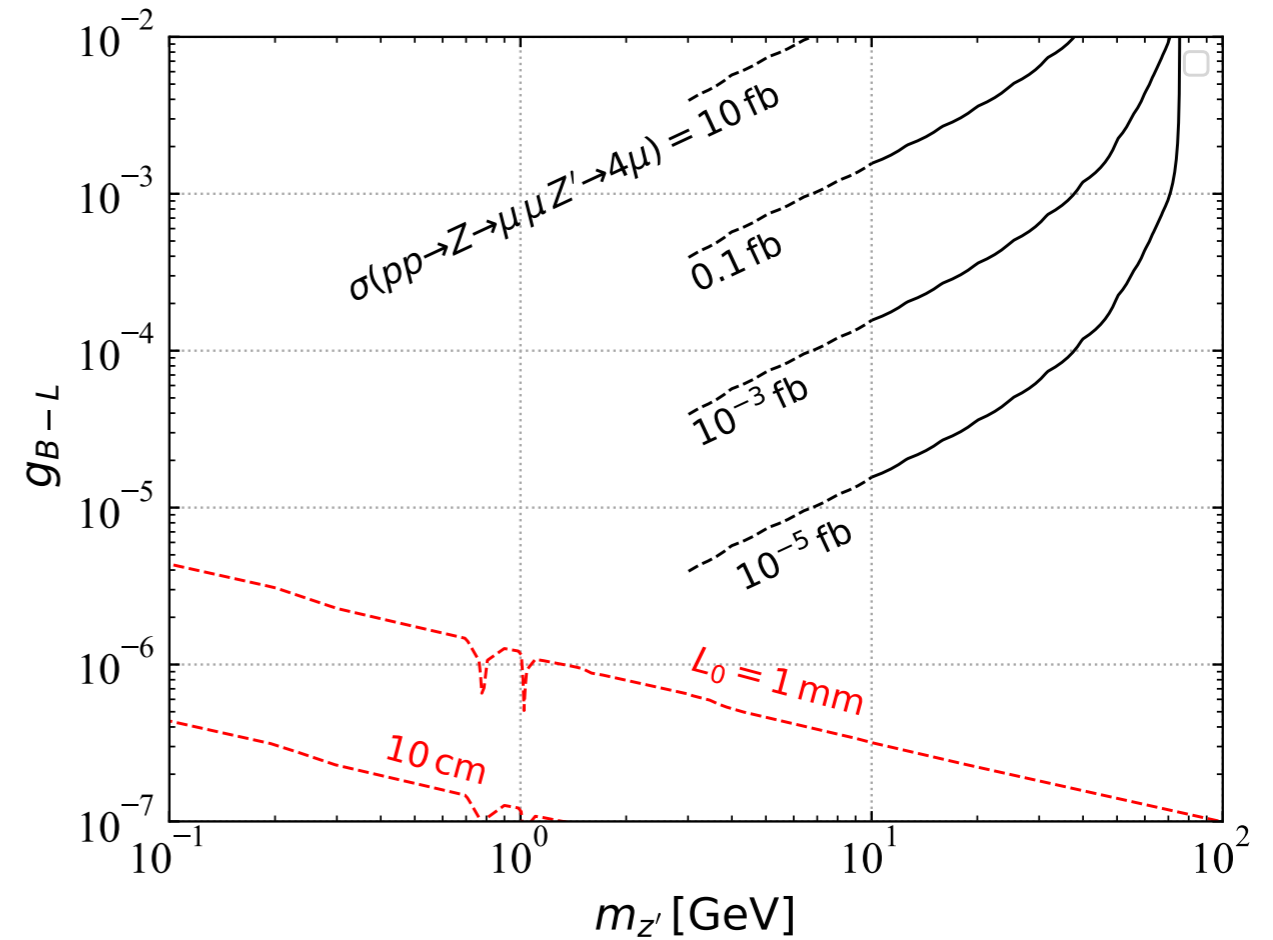
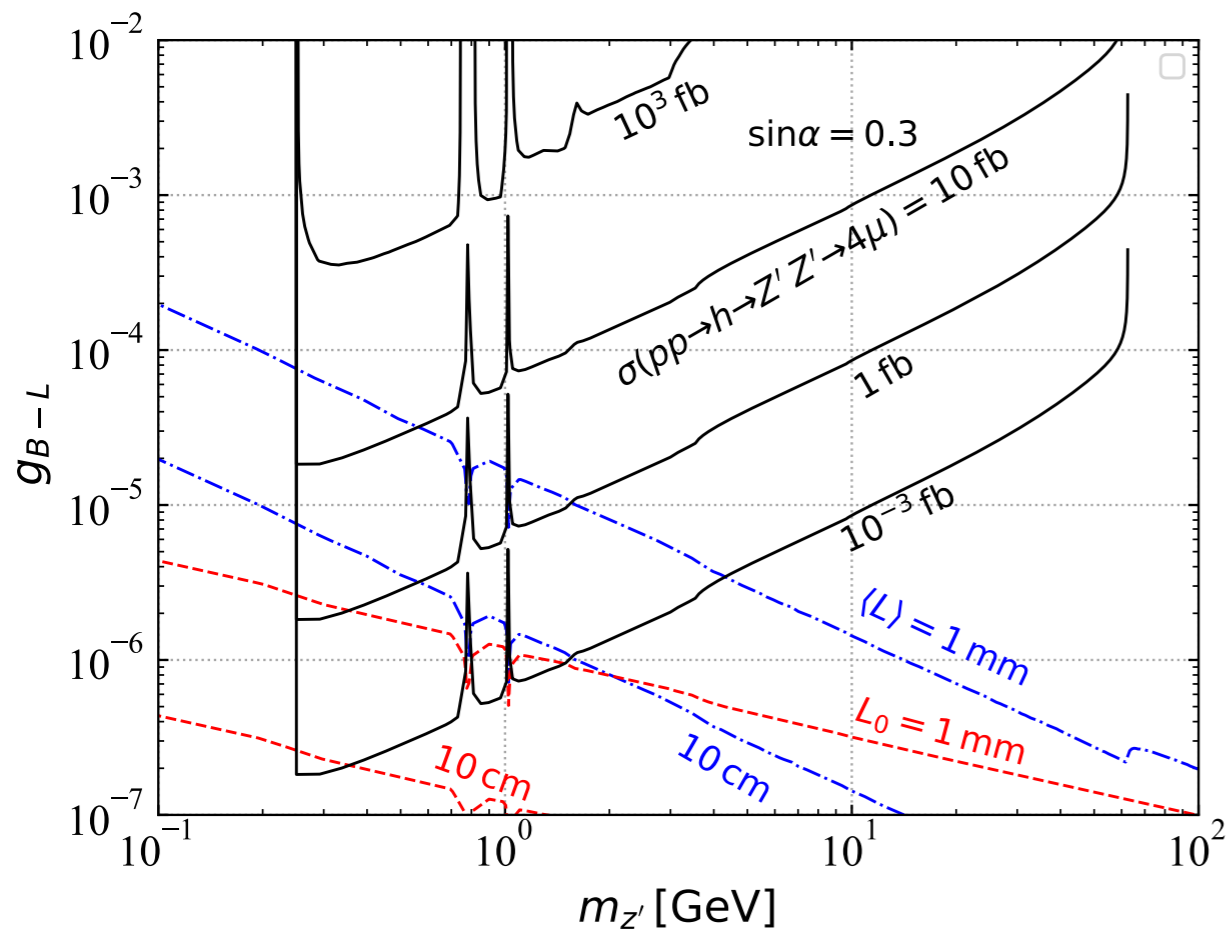
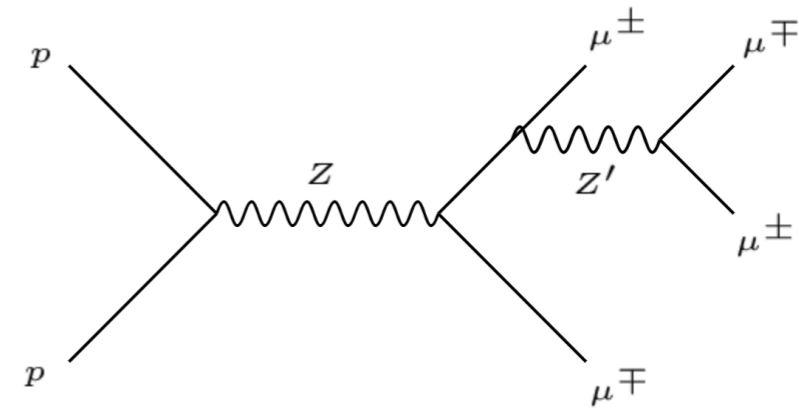
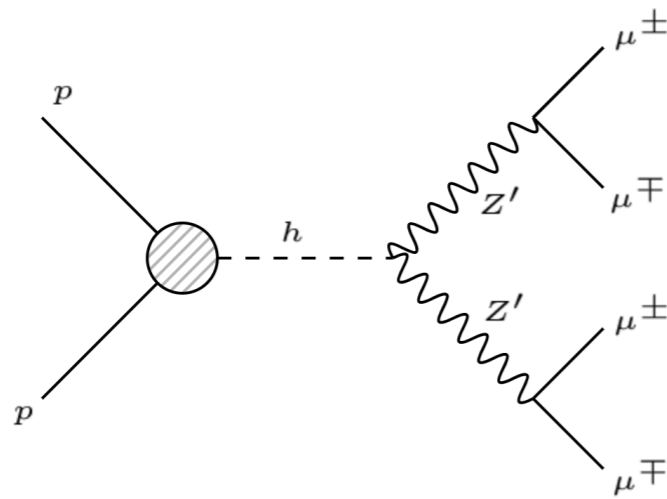




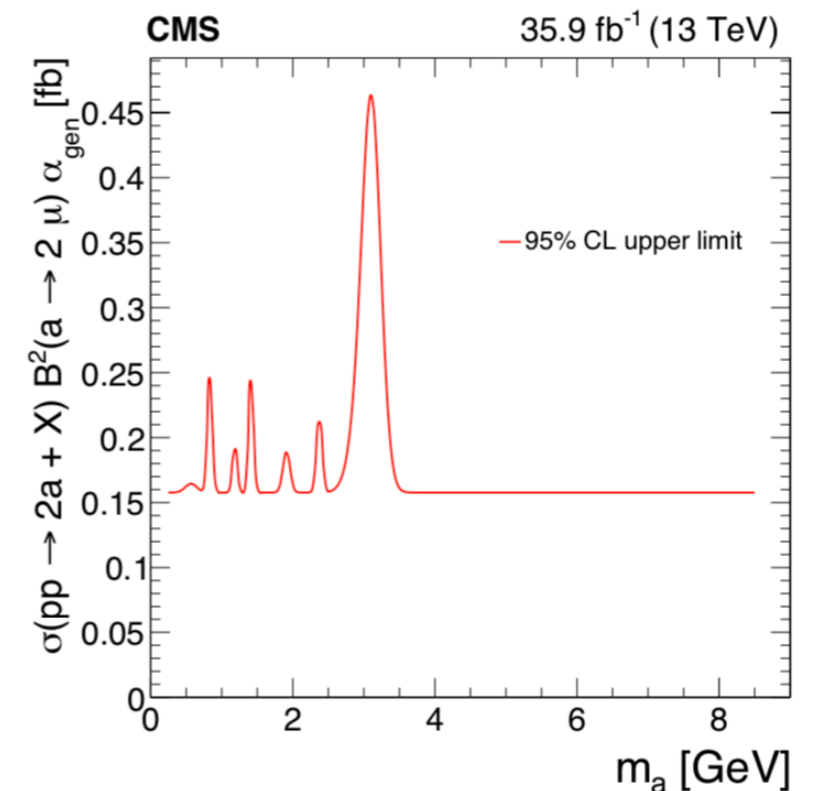
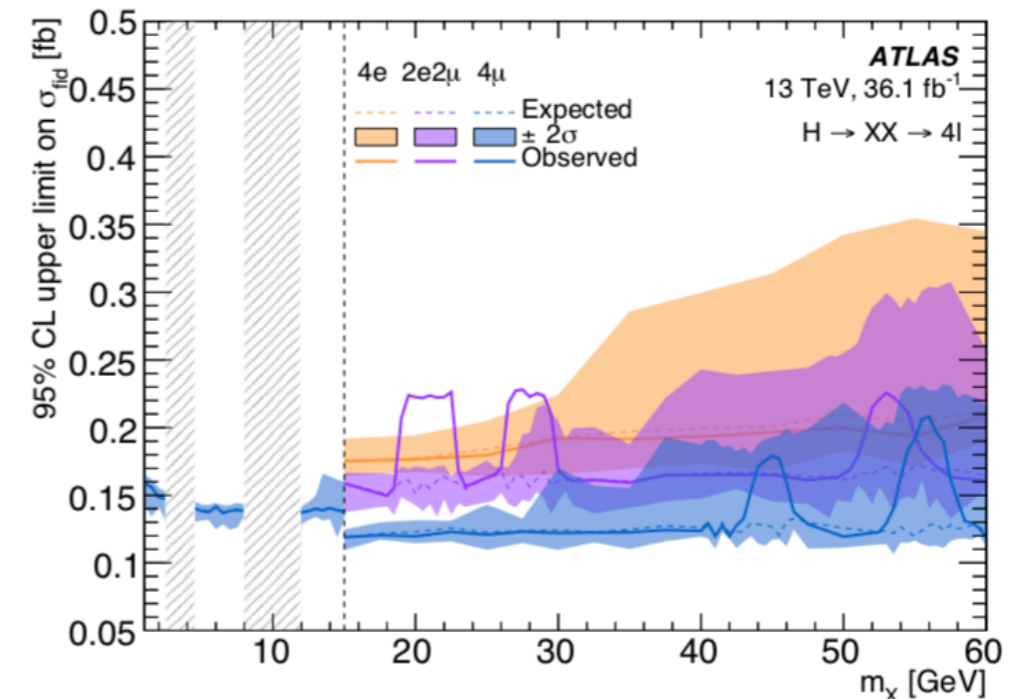
- New channels to probe
- FSR process usually most attractive for L_μ - L_τ models, included here for completion
- Both CMS and ATLAS search for 4 muon final state, interpretations usually in NMSSM or dark photon scenarios
- Higgs and Z' mediated heavy neutrino production still a possibility

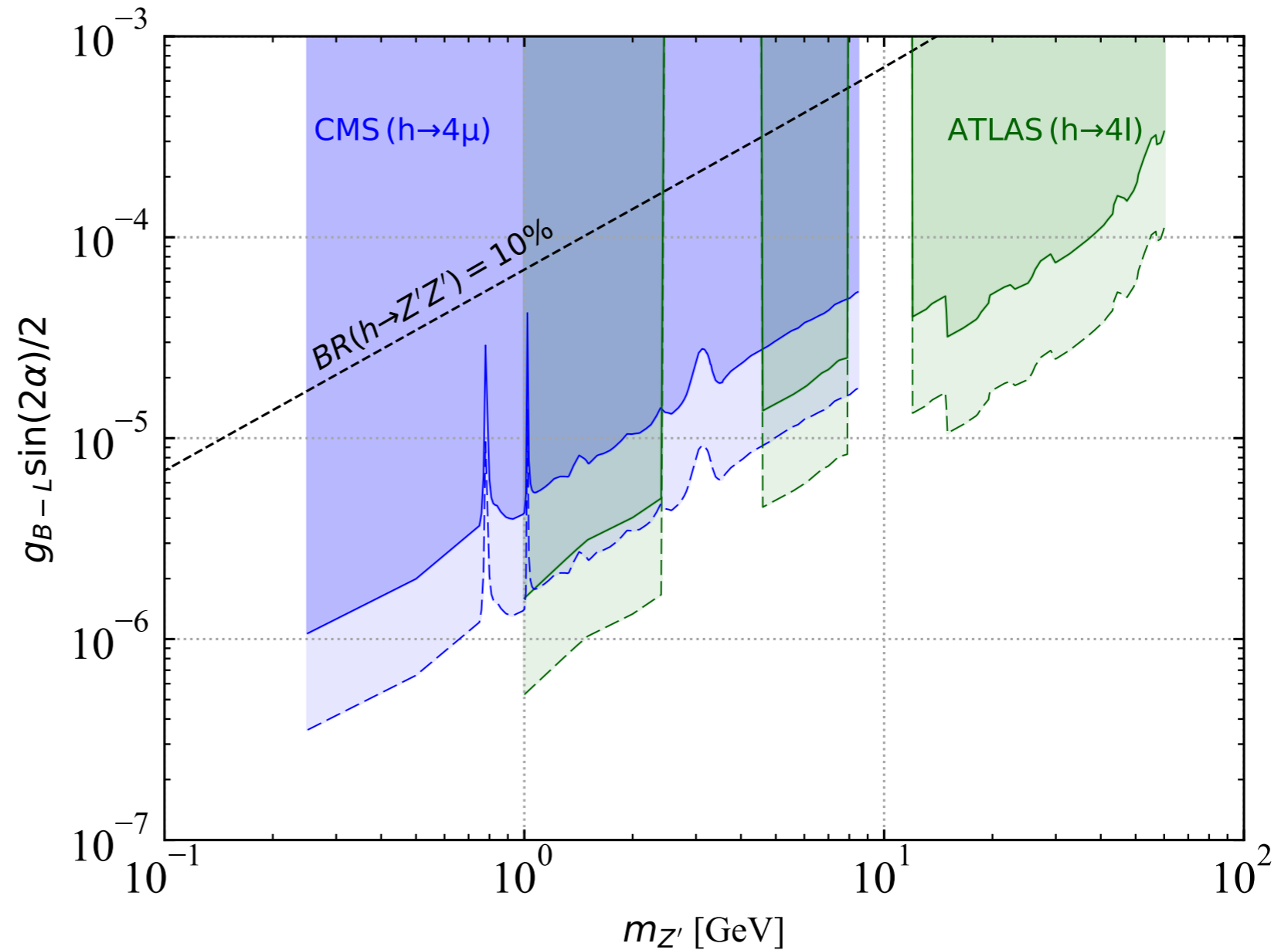
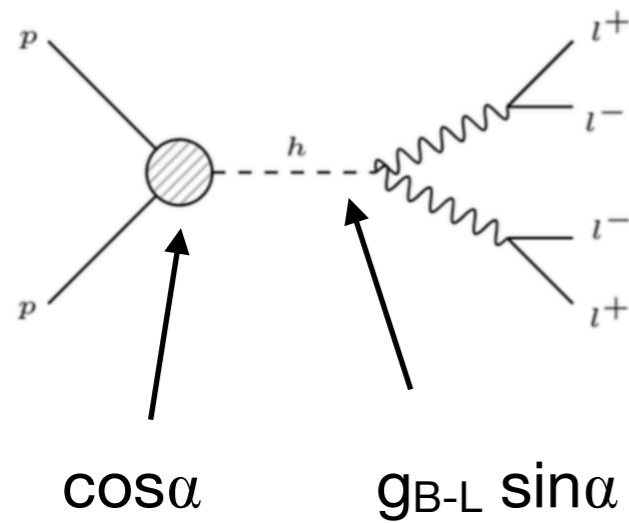


- In B-L model the term responsible for generating Z' mass is also responsible for mediating SM - B-L interactions
- Z' lifetime also controlled by g_{B-L} .
- This is not the case for dark photon models, mass generated by dark Higgs also mediated interactions
- Lifetime of Z_D controlled by mixing parameter ϵ



- **ATLAS-EXOT-2016-22 (13 TeV, 36.1 fb⁻¹):**
 - Analysis for Higgs decays to pair of dark photons
 - Searches for Z_D decays to pair of electrons or muons
 - Presents fiducial cross sections and limits on Higgs Z_D branching fraction
 - Z_D mass between 1 to 60 GeV
- **CMS-HIG-18-003 (13 TeV, 35.9 fb⁻¹):**
 - Searches for pair production of light bosons
 - Considers NMSSM models
 - Also sensitive to moderate displacements $L_{xy} < 10$ cm
 - Presents fiducial cross sections and model specific limits
 - Z_D mass between 0.25 to 8.5 GeV

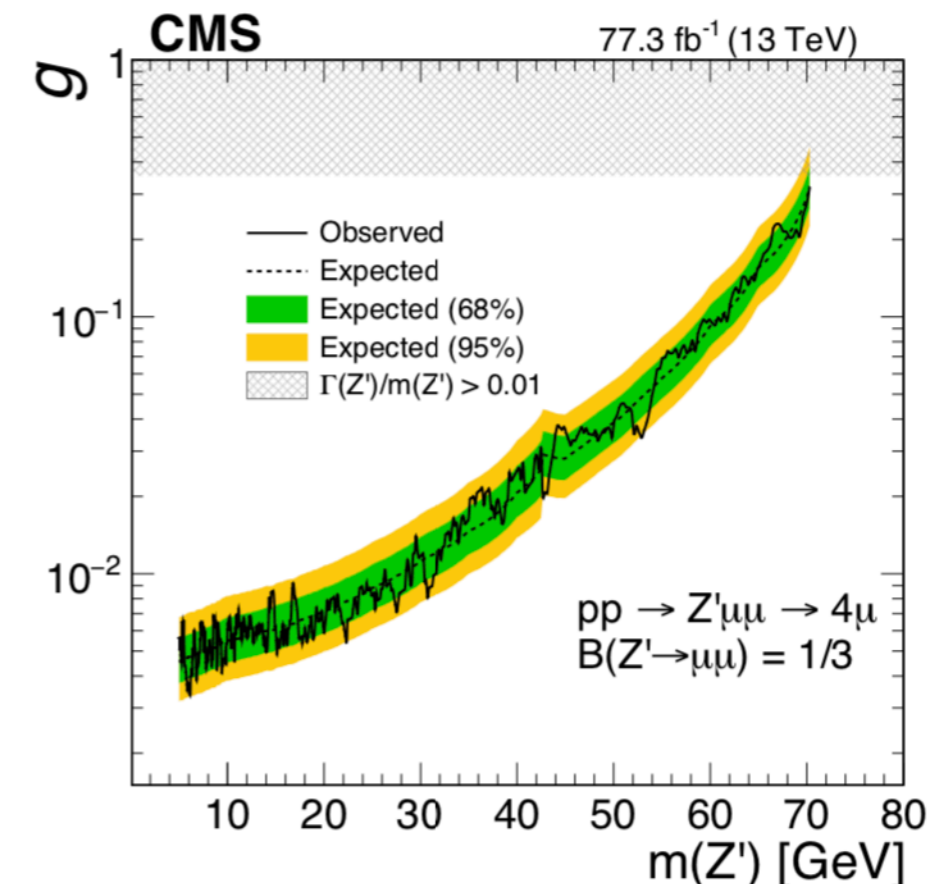




- For Higgs mediated Z' production, constraints on the product of coupling
- Easy way to rescale limits given g_{B-L} or $\sin \alpha$ value
- Important: rescaling to be taken with caution, different treatment for displaced regime

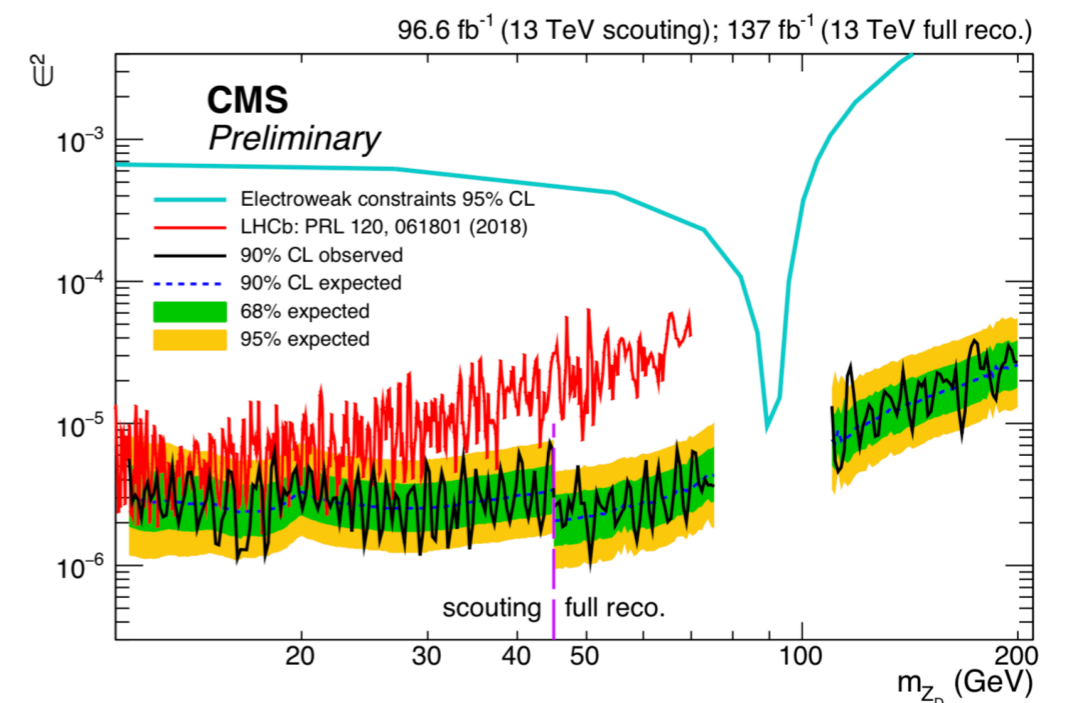
- **CMS-EXO-18-008 (13TeV, 77.3 fb⁻¹):**

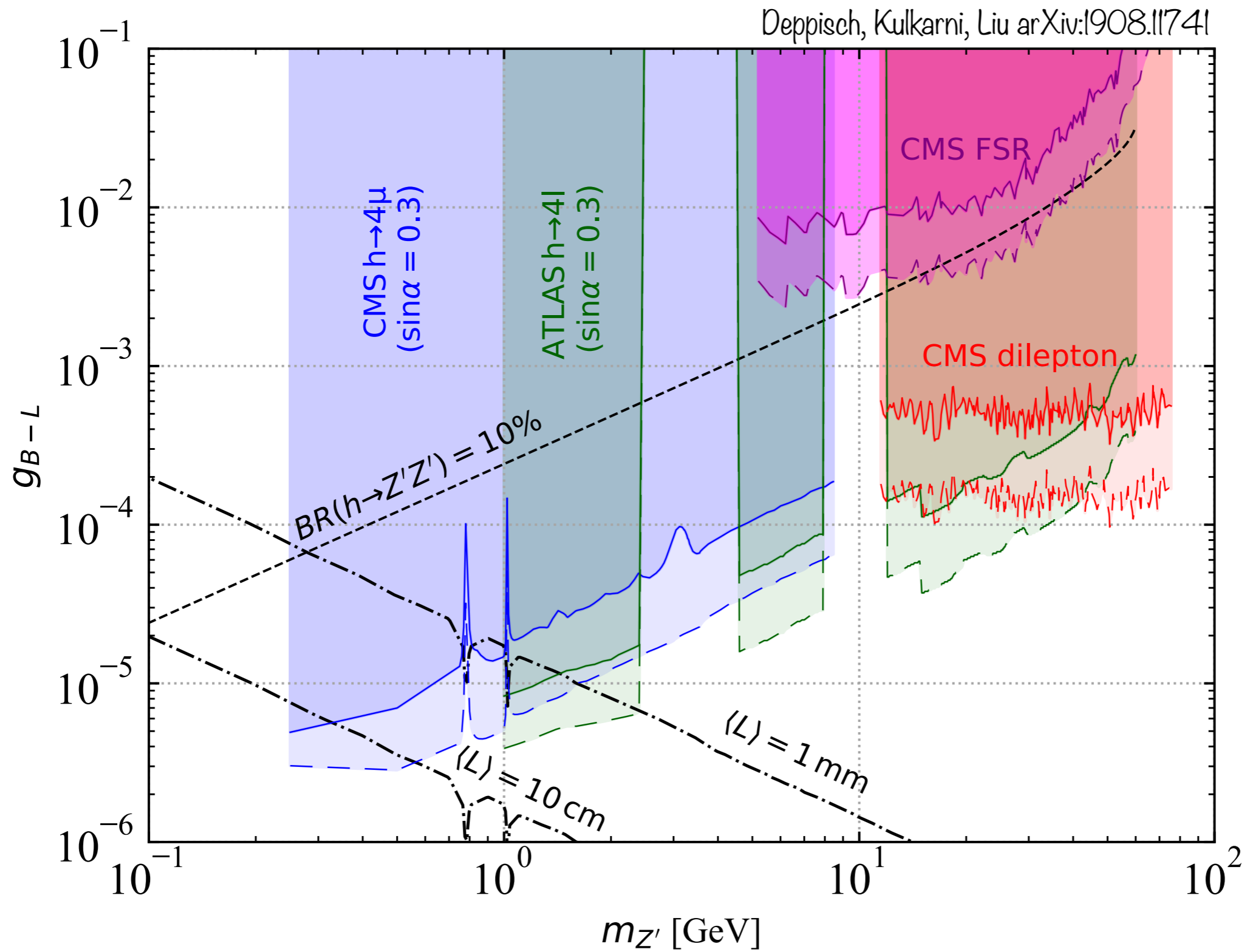
- Final state radiation of Z' in DY Z production
- Muon final state only
- Particularly useful for L_μ - L_τ
- Limits on the L_μ - L_τ z' coupling as a function of mass, easy to rescale
- Mass range between 6 to 70 GeV

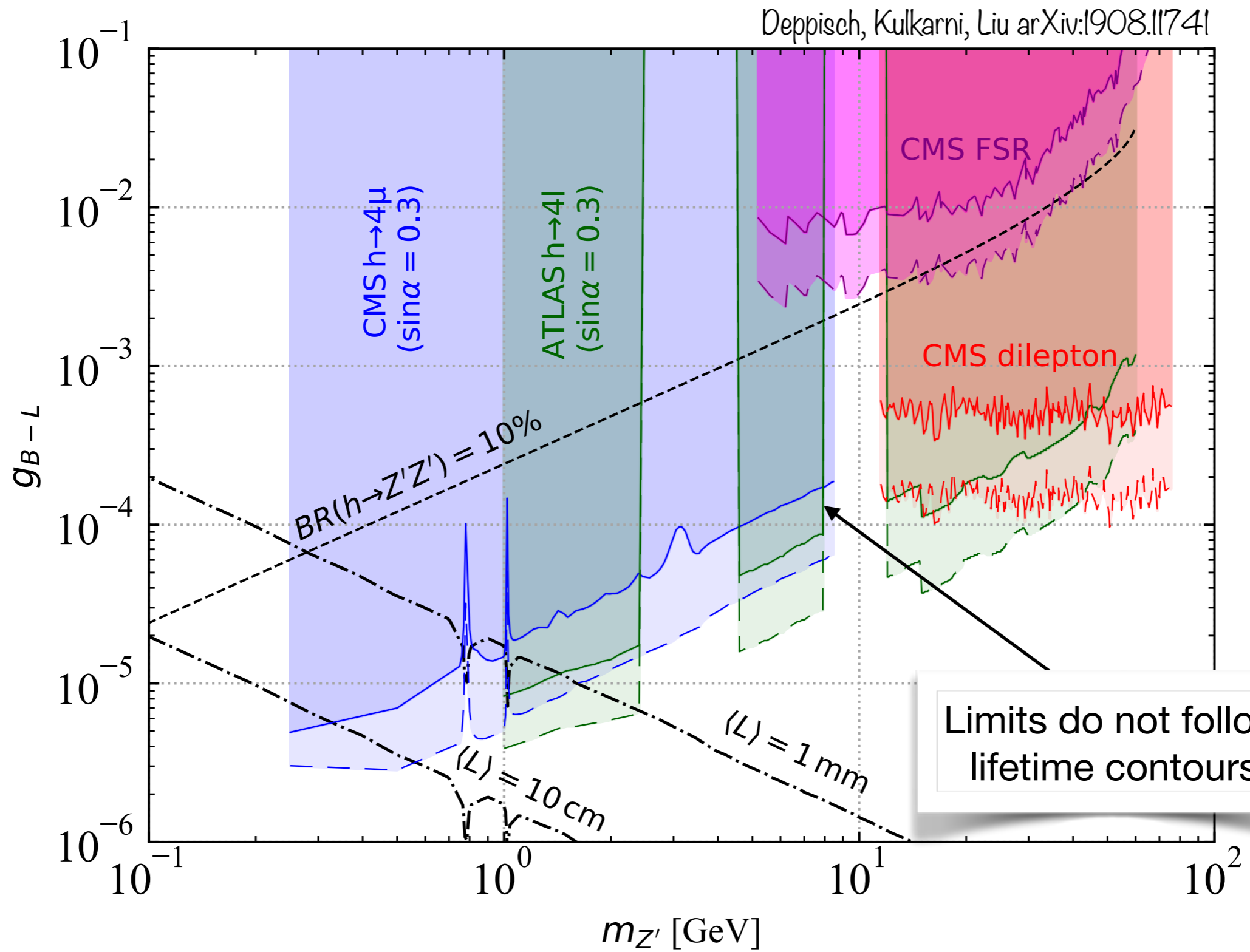


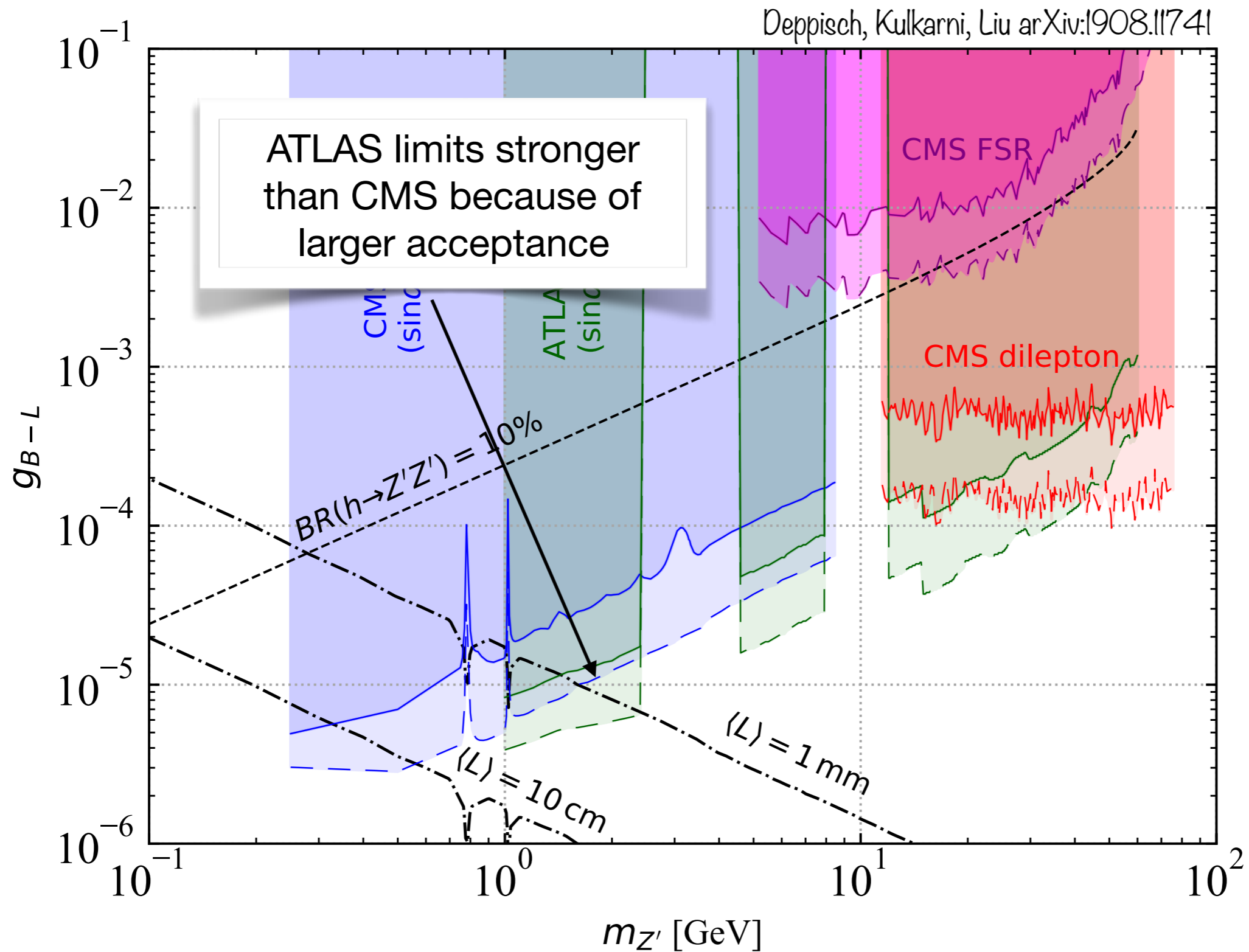
- **CMS-EXO-19-018 (13TeV, 137 fb⁻¹):**

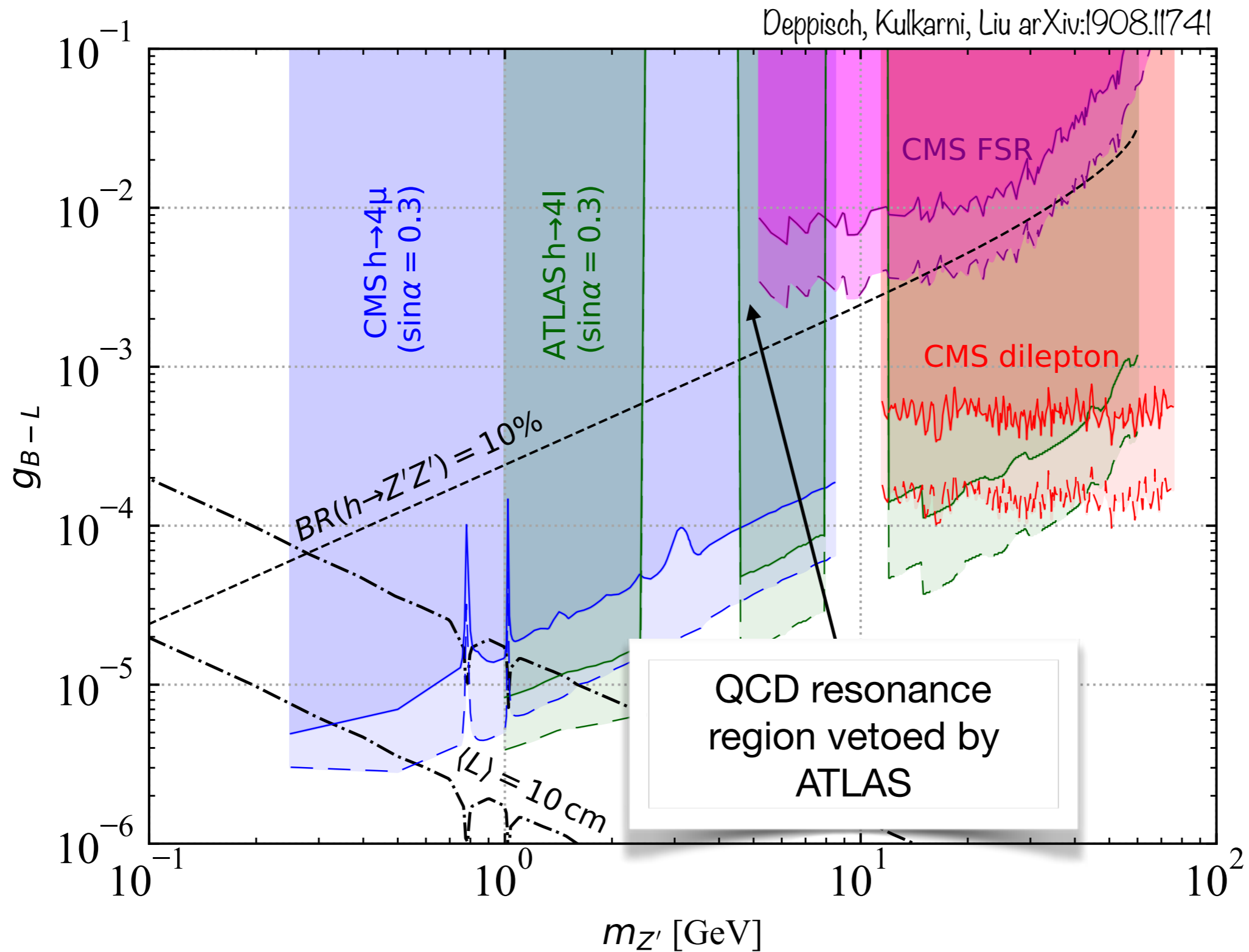
- Latest and greatest CMS muon scouting analysis
- Search for narrow resonance decaying to pair of muons
- Scouting and full reach analysis
- Mass range 10 to 200 GeV

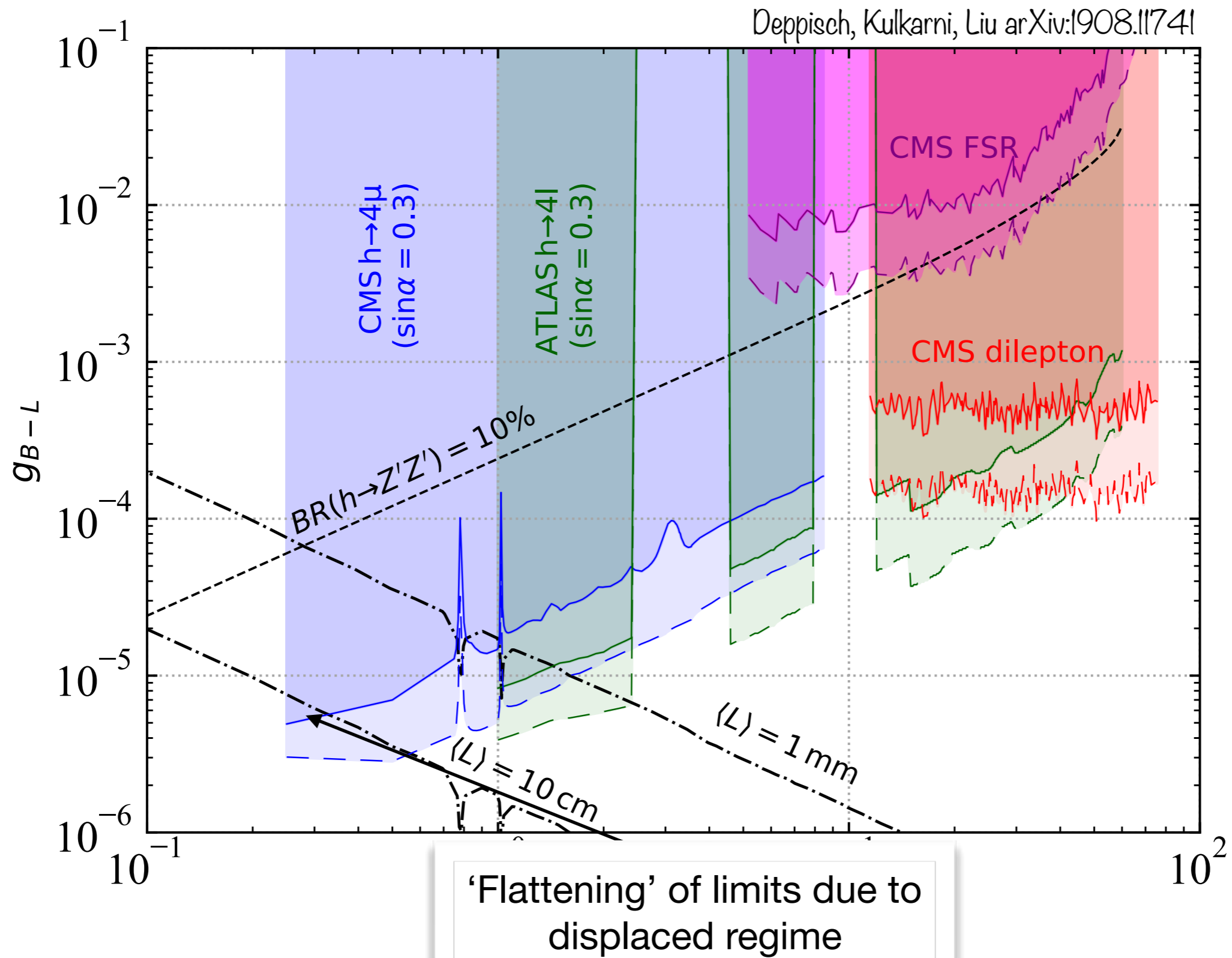


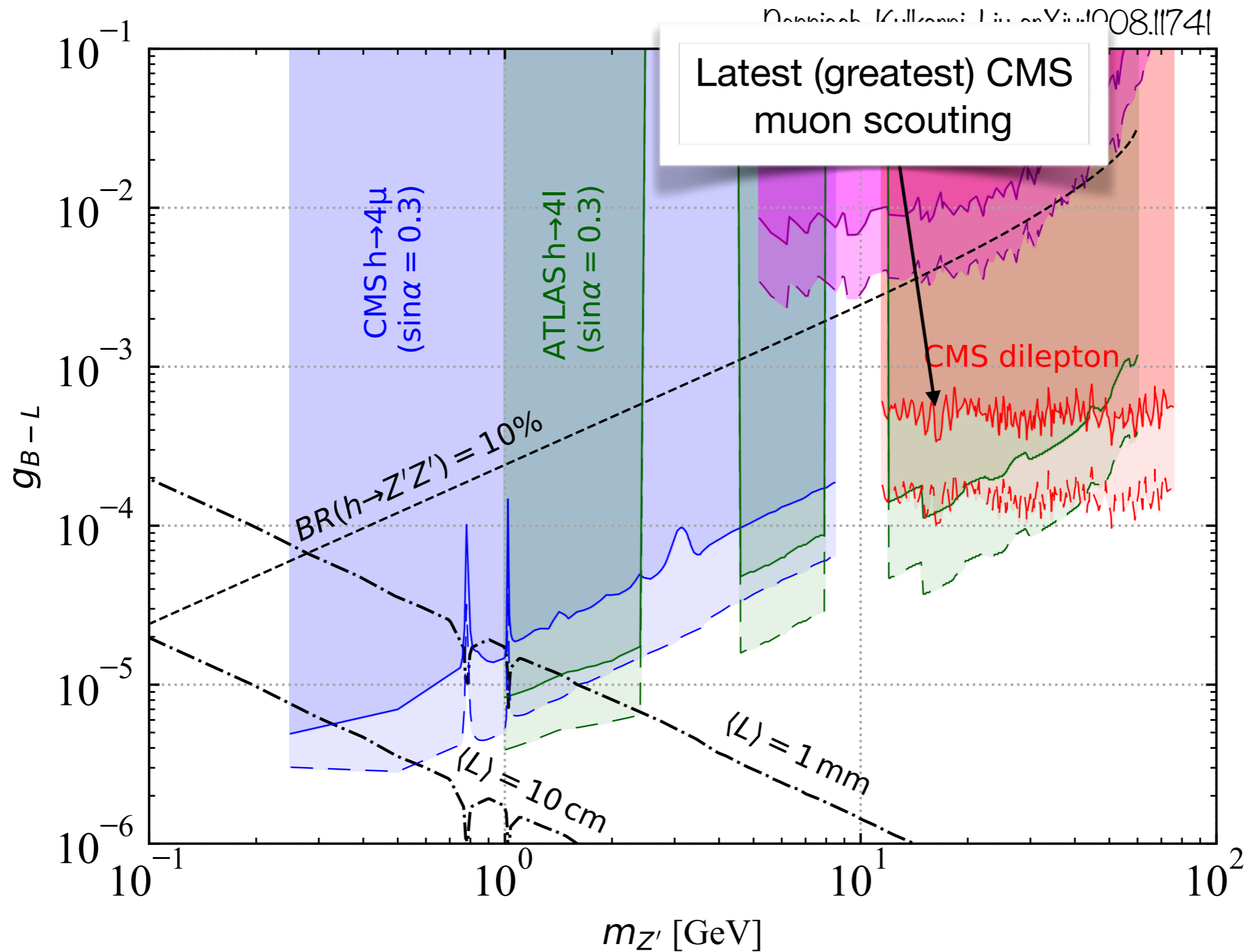


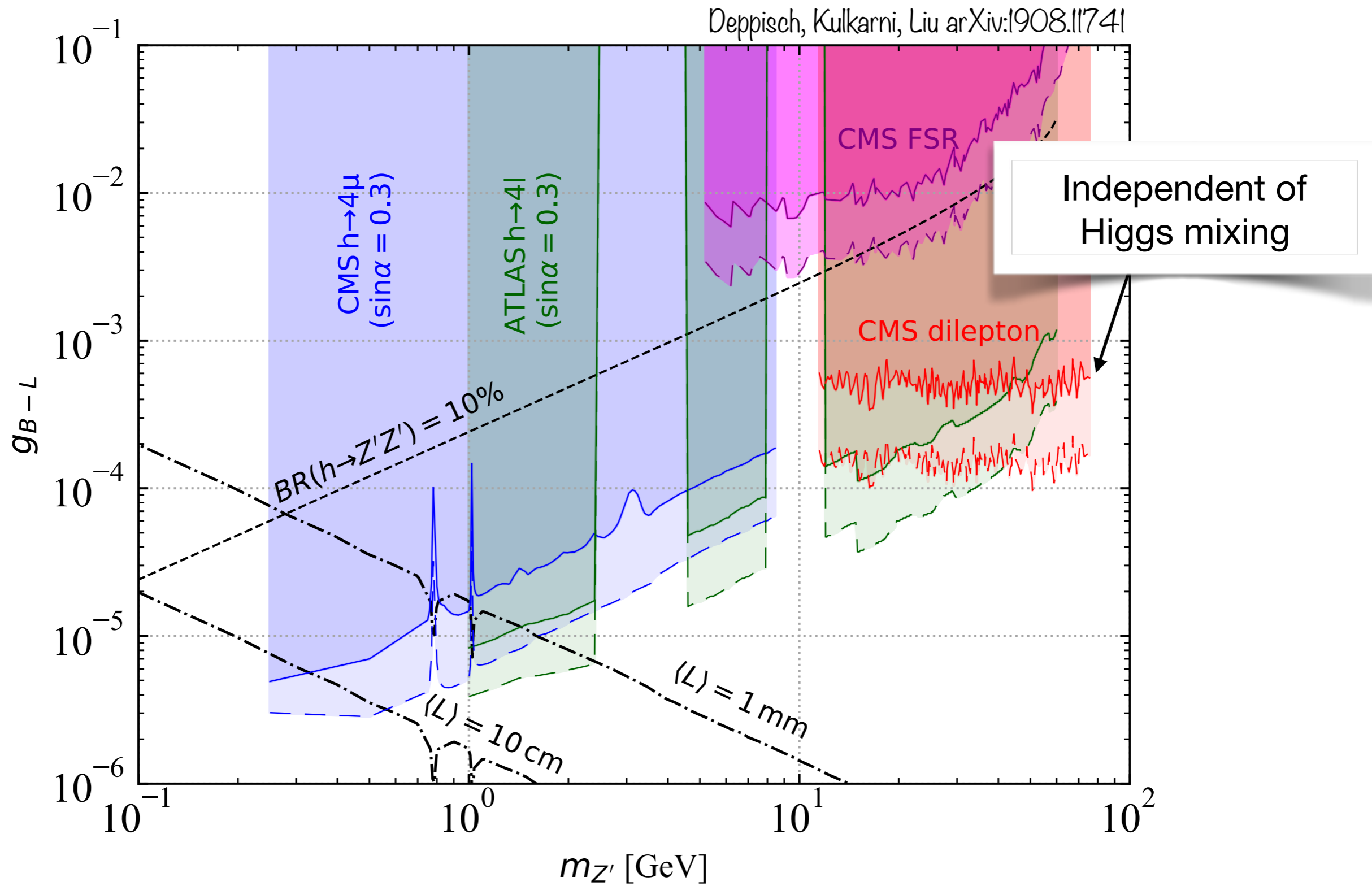


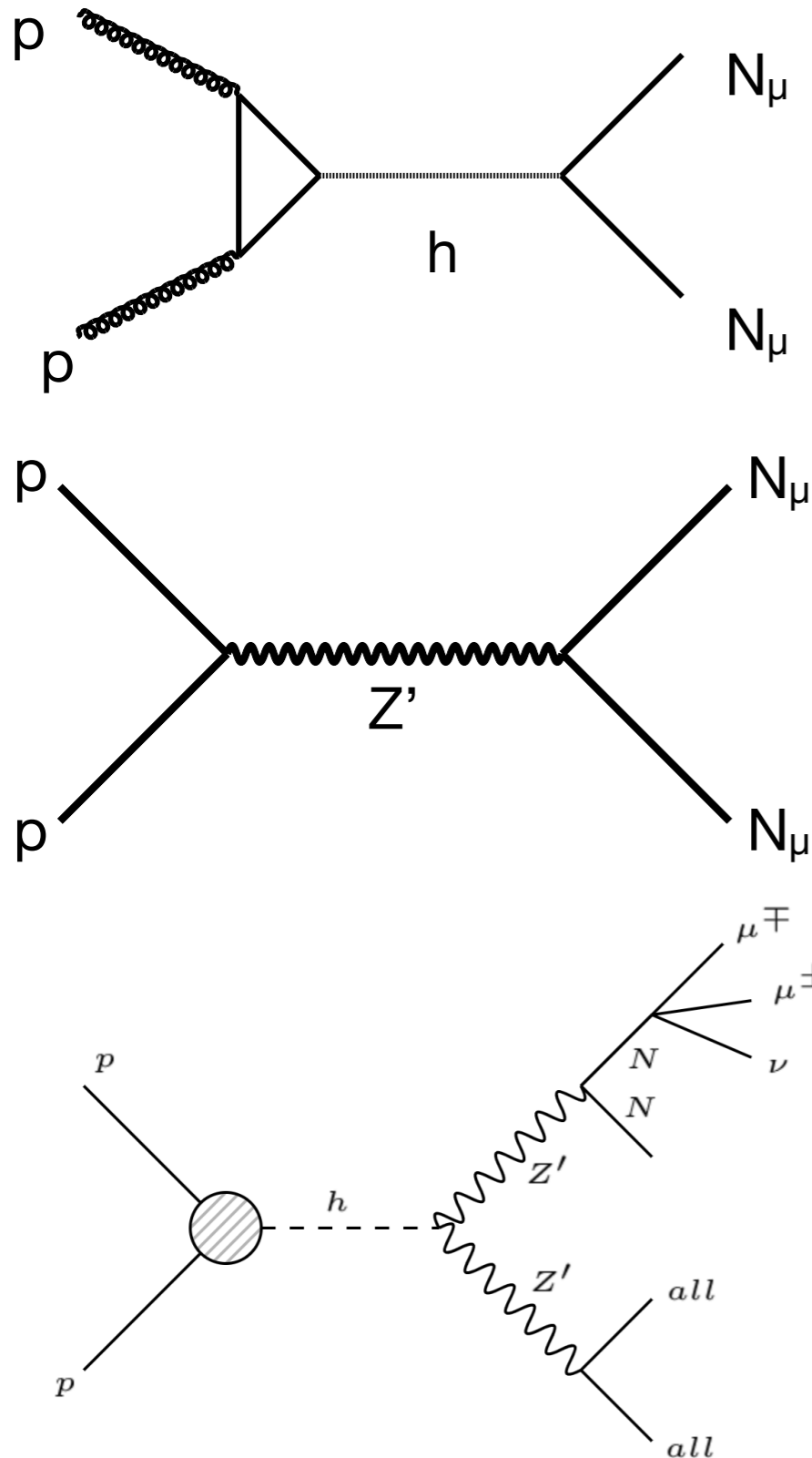






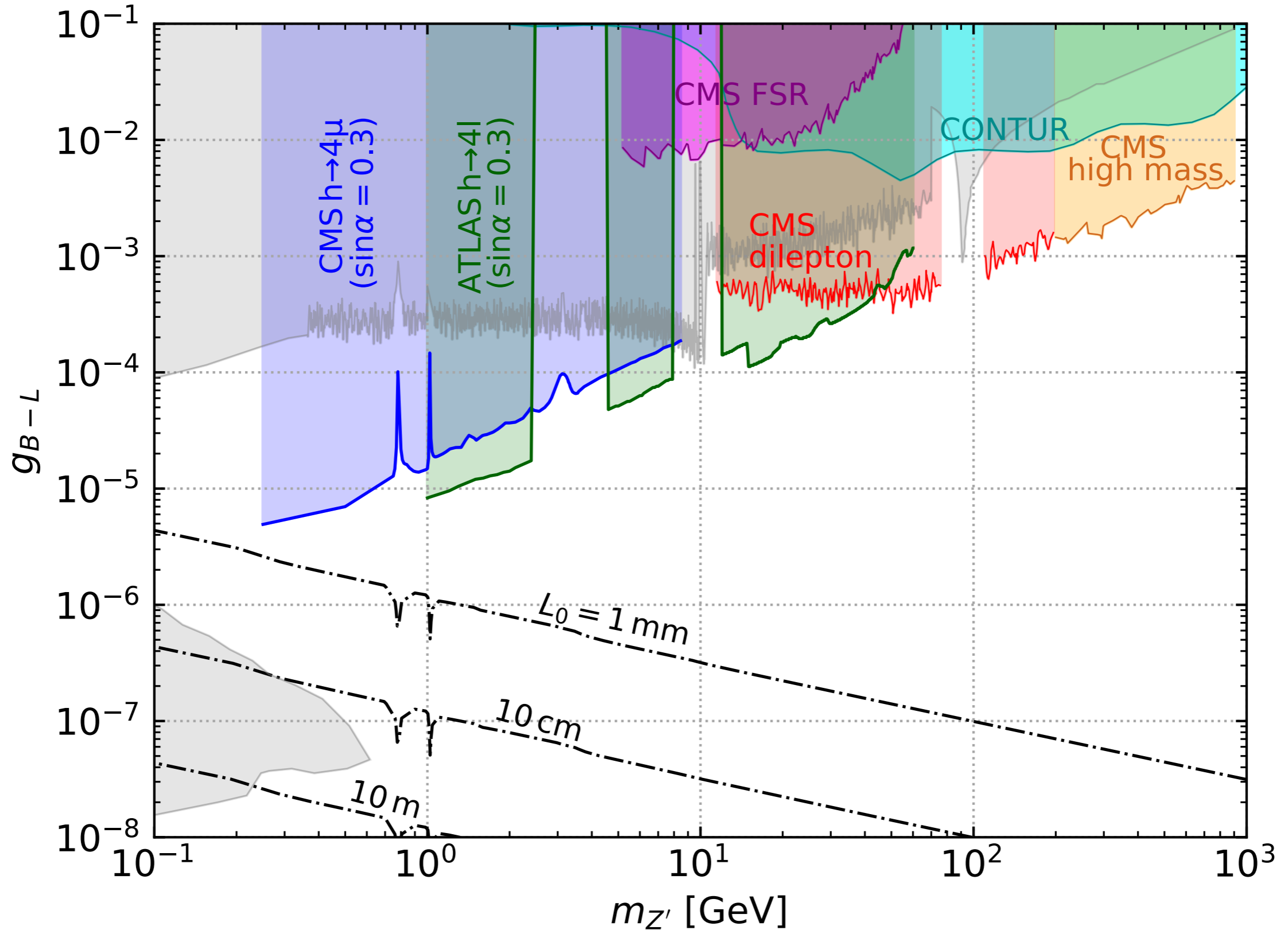


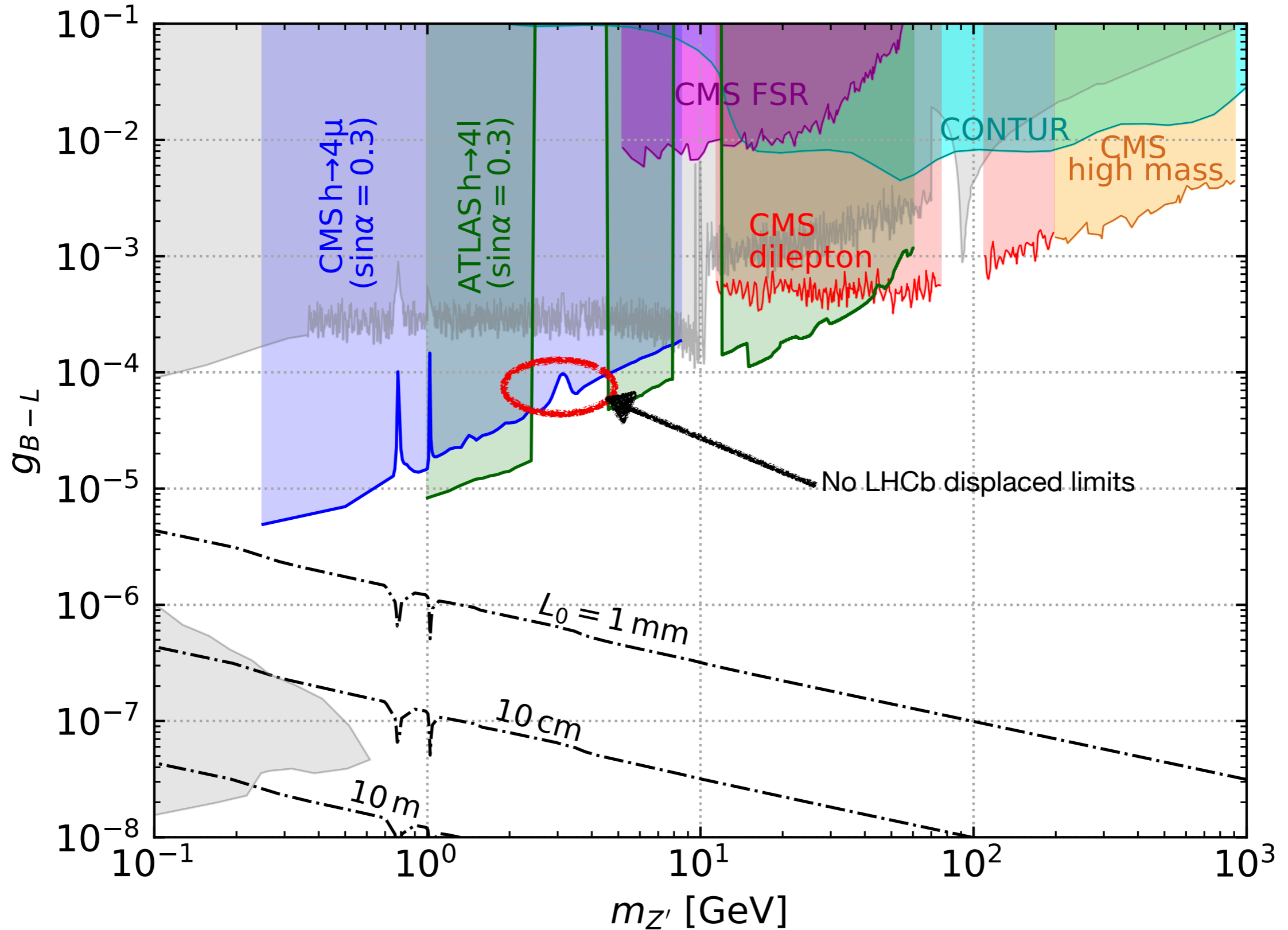


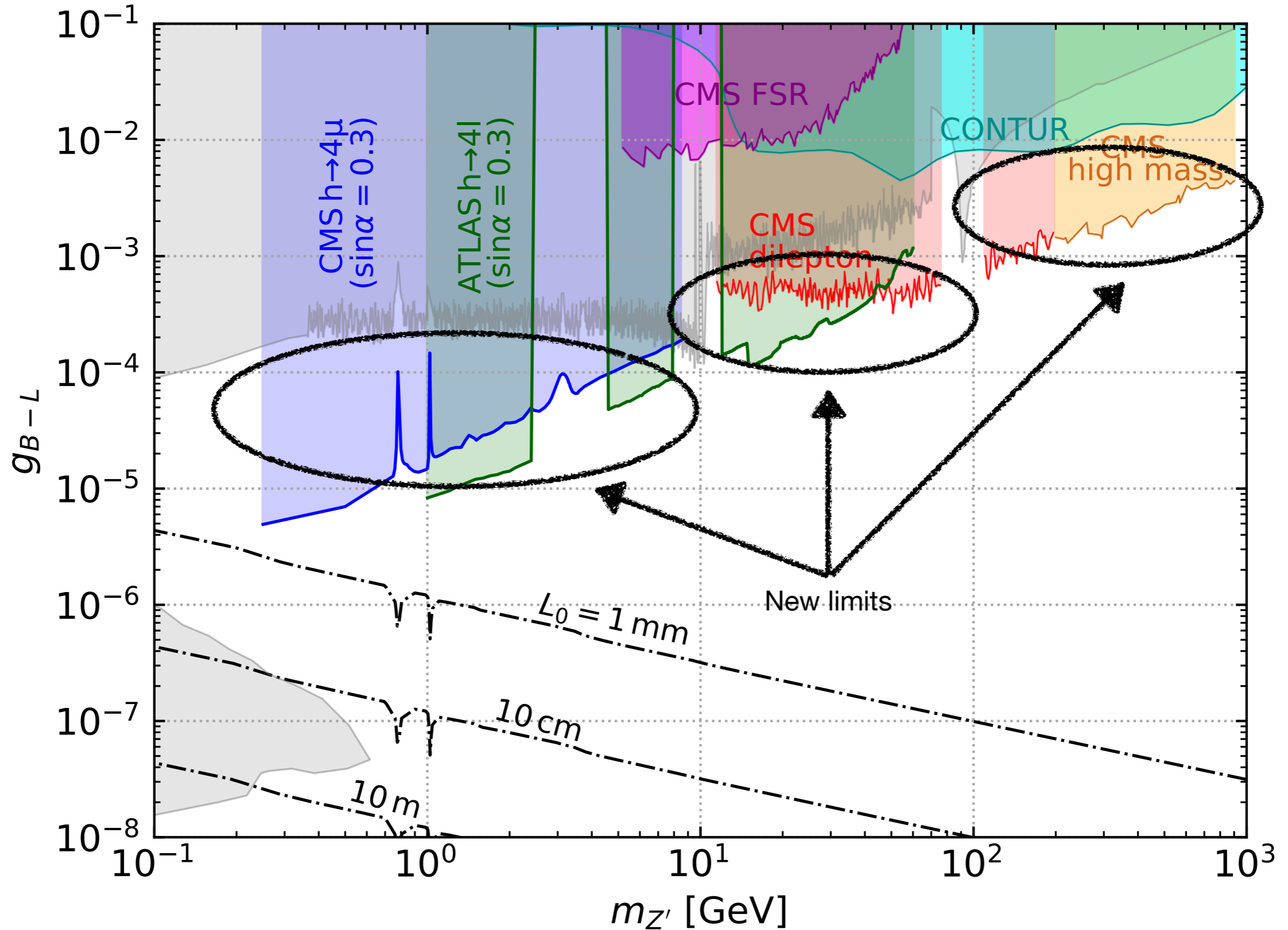


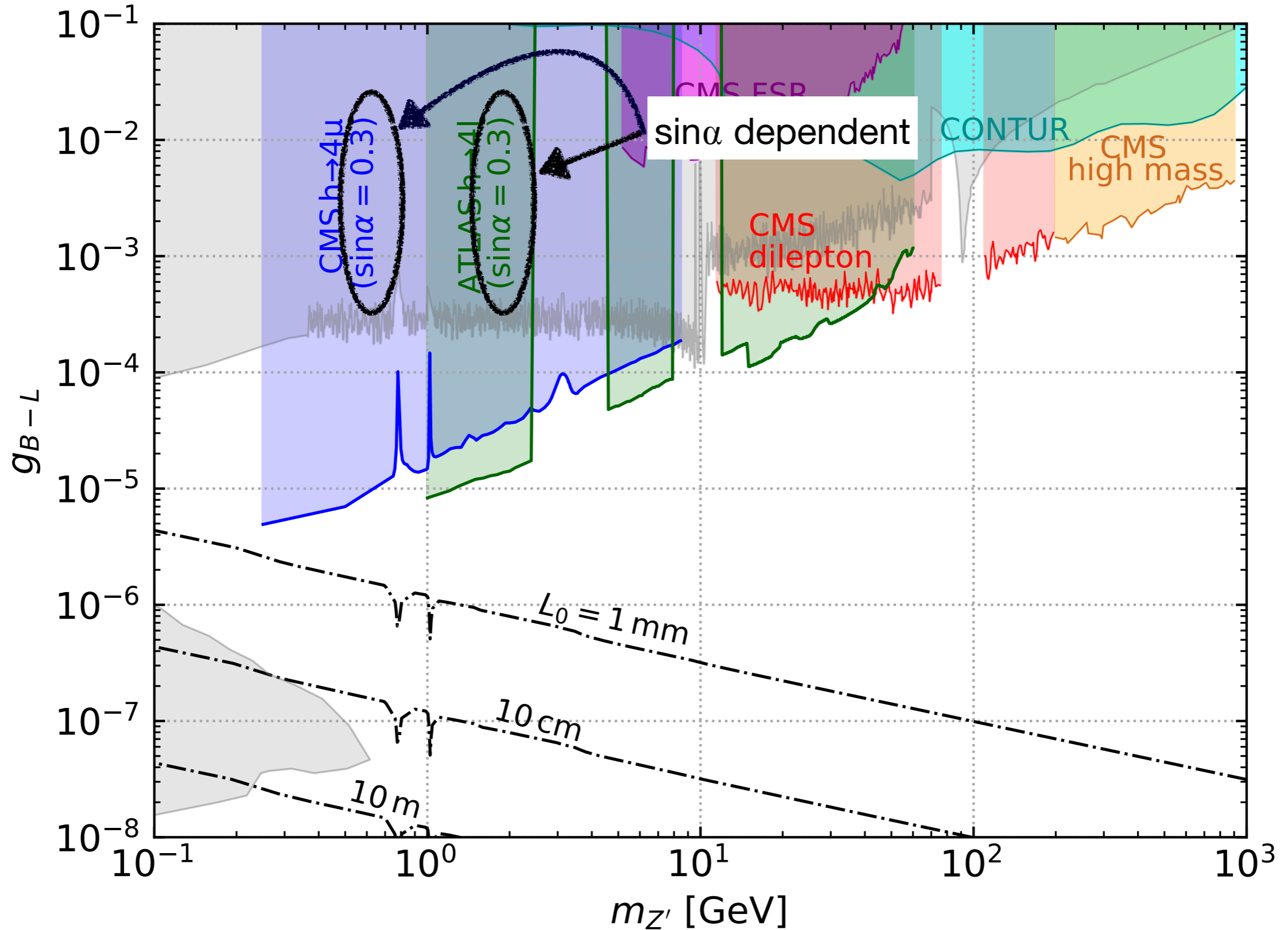
Vertex	Feynman rule
$Z' - l - l$	g_{B-L}
$h - Z' - Z'$	$g_{B-L} \cos\alpha m_{Z'}$
$h - N - N$	$g_{B-L} \cos\alpha m_N/m_{Z'}$
$Z' - N - N$	g_{B-L}

- All vertices $\propto g_{B-L}$
- Improved constraints on $g_{B-L} \rightarrow$ decreased heavy neutrino production rate via h and Z' mediators
- h and Z' mediators remain interesting option in limiting cases

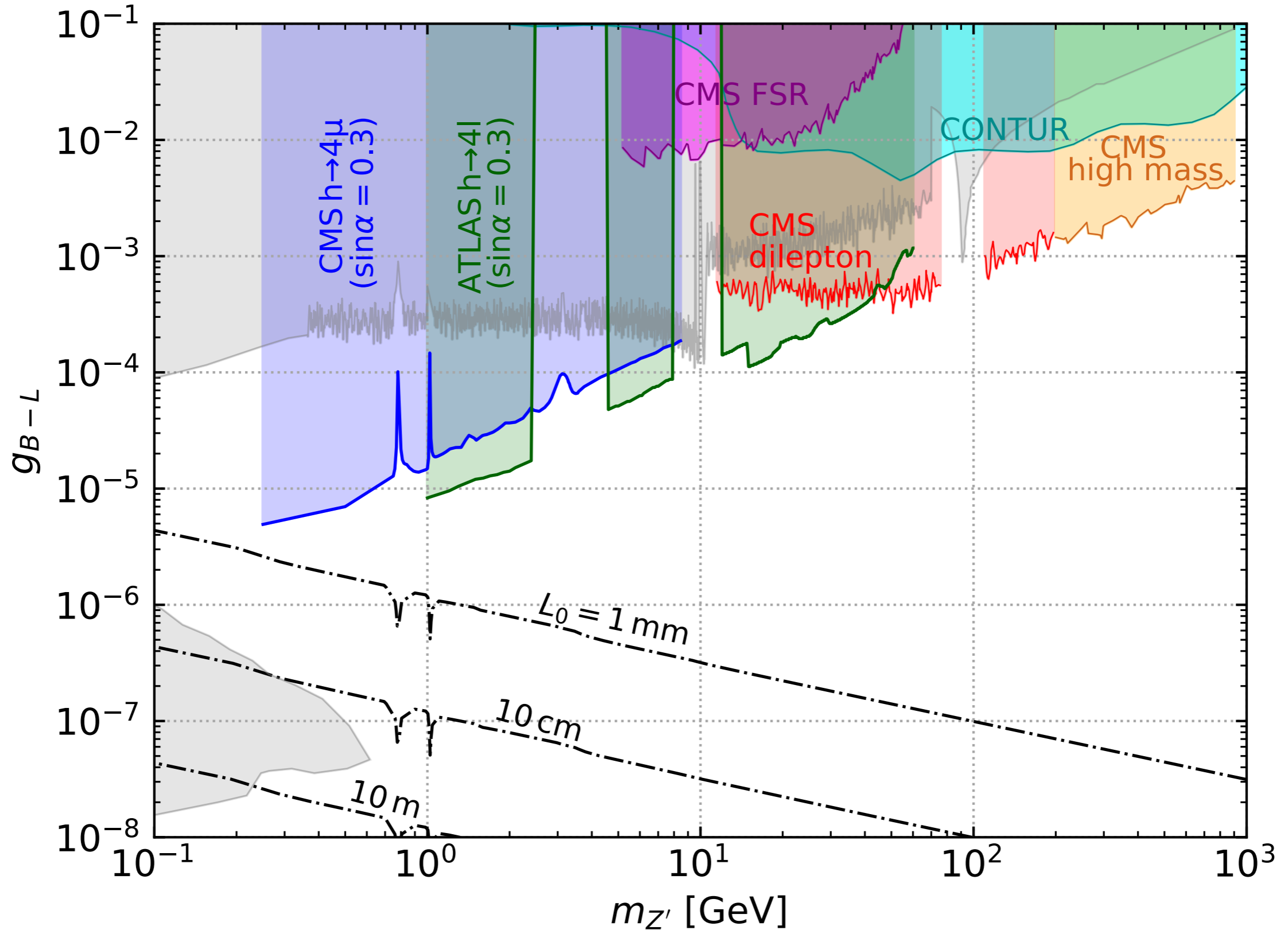








The global picture

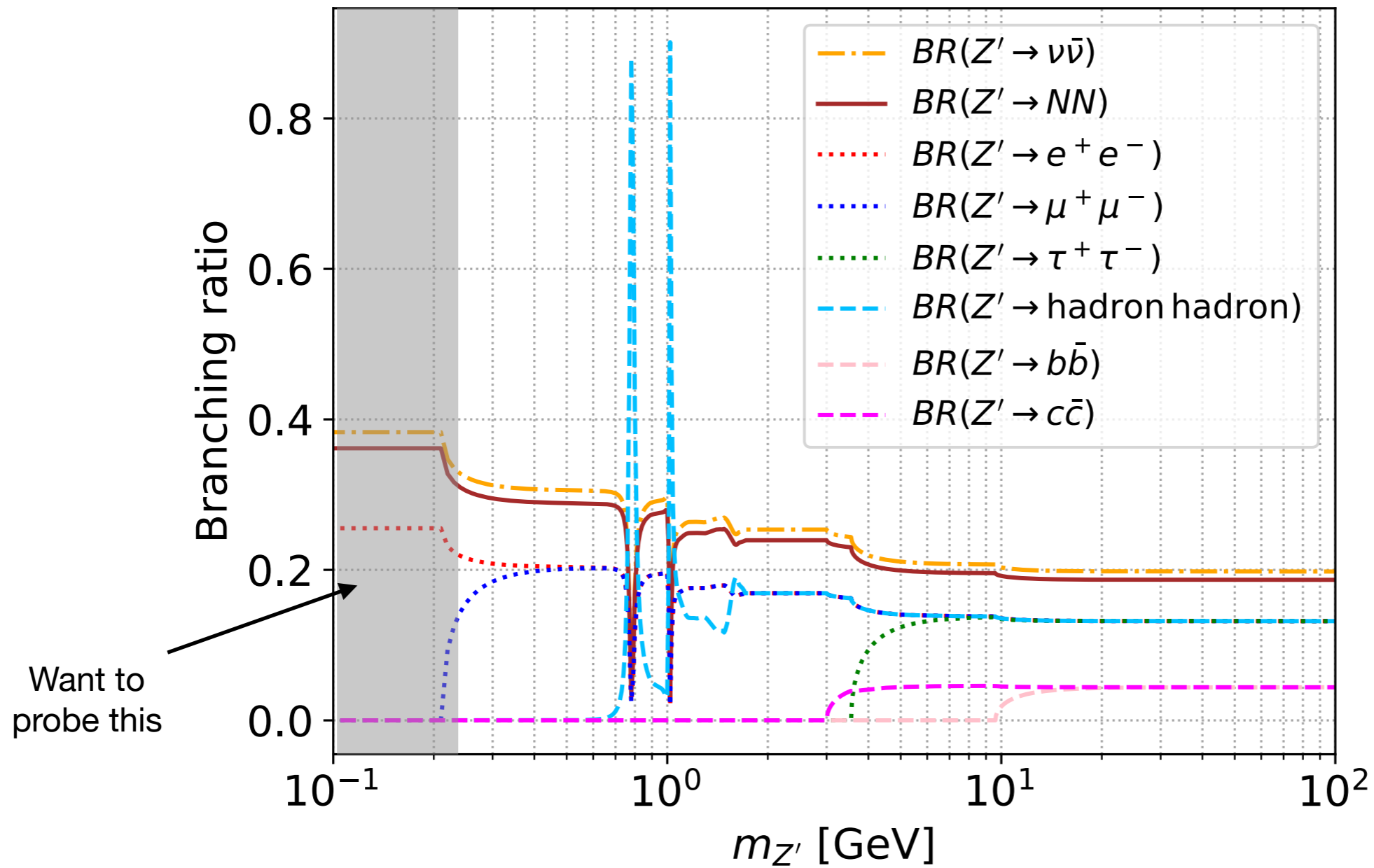


- B-L models one of the simplest extensions of SM physics providing explanations of neutrino masses
- Heavy neutrino production can take place via SM Higgs or B-L Z' decays and probe different regions of B-L parameter space
- Have potential to probe neutrino mixing angles responsible for neutrino mass generations
- Z' production via SM Higgs tightly constrained via ‘lepton-jet’ searches. If Z' production via Higgs is possible, heavy neutrino production is suppressed
 - ATLAS analyses consider both electron and muon lepton-jets
 - Prompt analyses constrain parameter space where Z' is displaced
 - Interpretation for displaced regime not always straightforward
 - Model independent fiducial cross section limits from collaborations are very welcome for reinterpretation exercises
 - Careful reconsideration of information for exotic searches specially trigger and object level efficiencies necessary

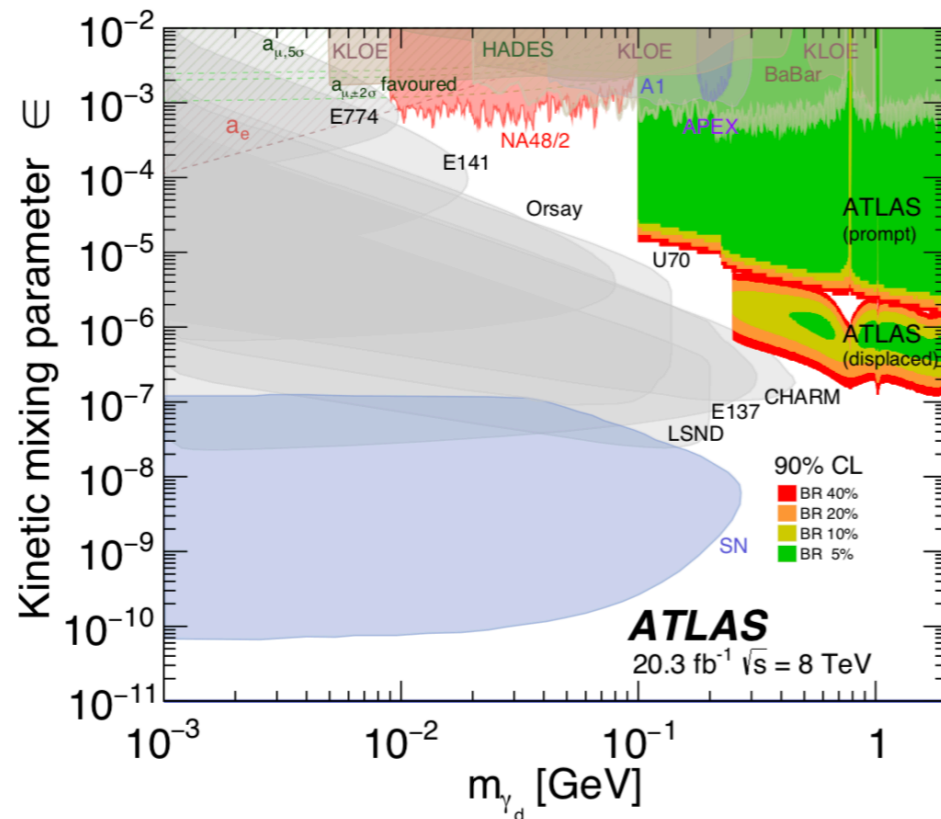
Thank you!



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CoghillCartooning.com

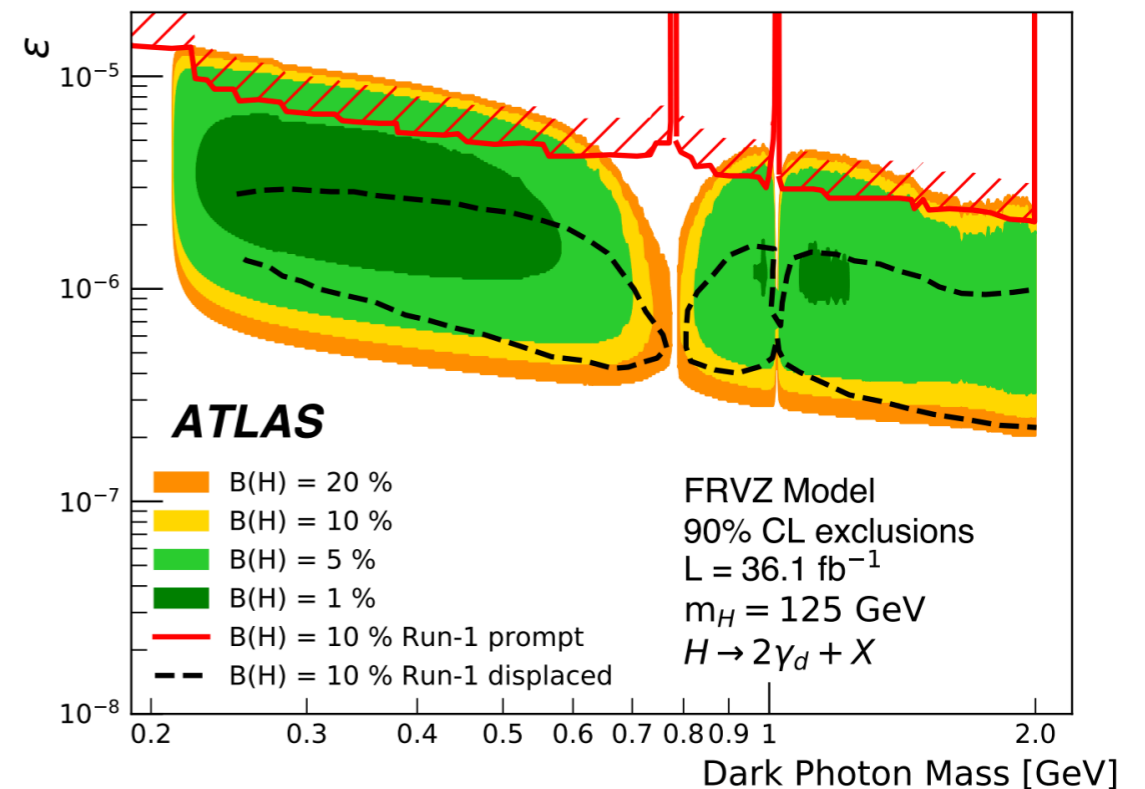


- Possibility of extending Higgs portal analyses for electron final states?
- Mono-jet constraints?

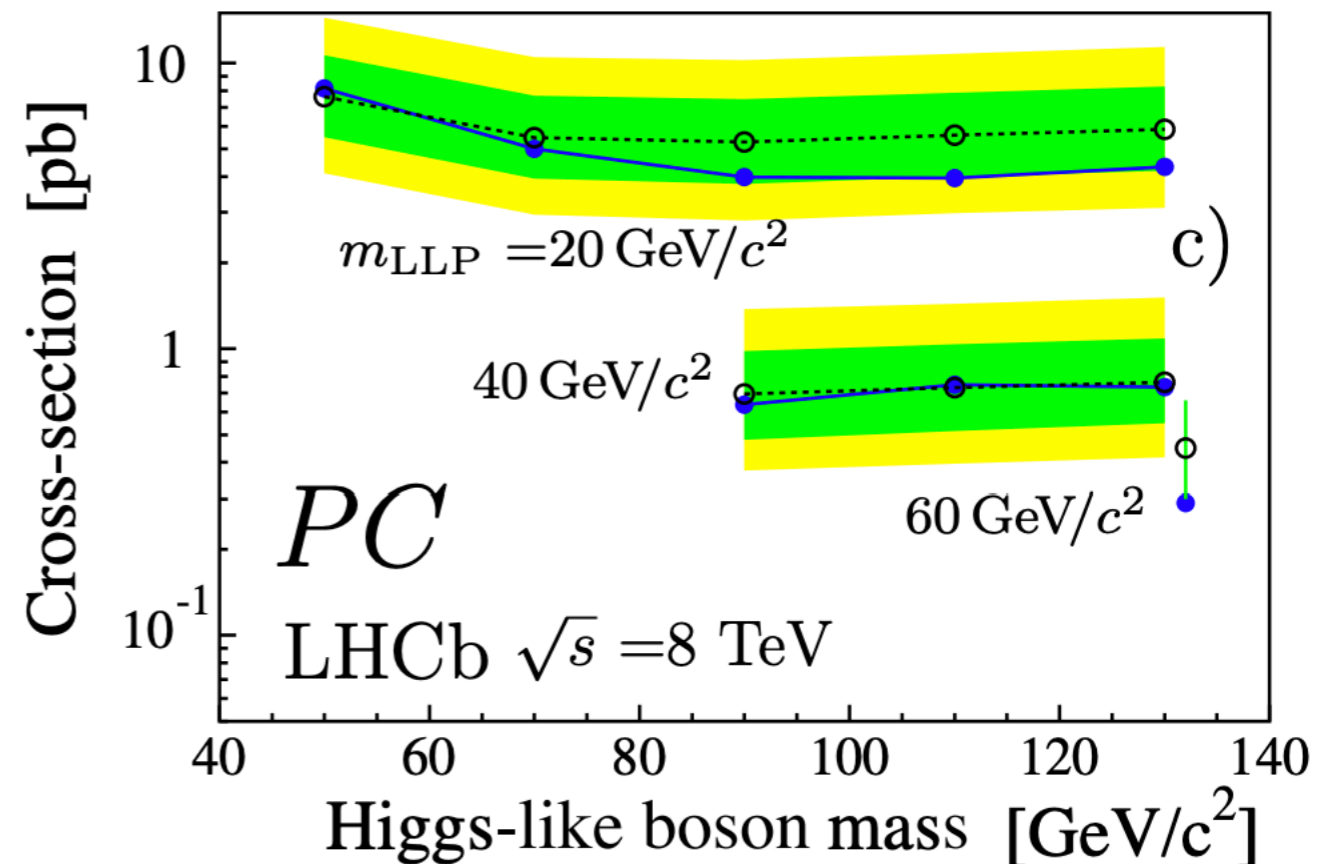


- **ATLAS-EXOT-2014-09 (8 TeV, 20.3 fb⁻¹)**
 - Prompt lepton - jets analysis
 - Limits as a function of FRVZ Z_D mass
 - Both electron and muon final states
 - Mass range from 0.25 to 1.5 GeV
 - Competitive (but not better) limits than CMS at low mass

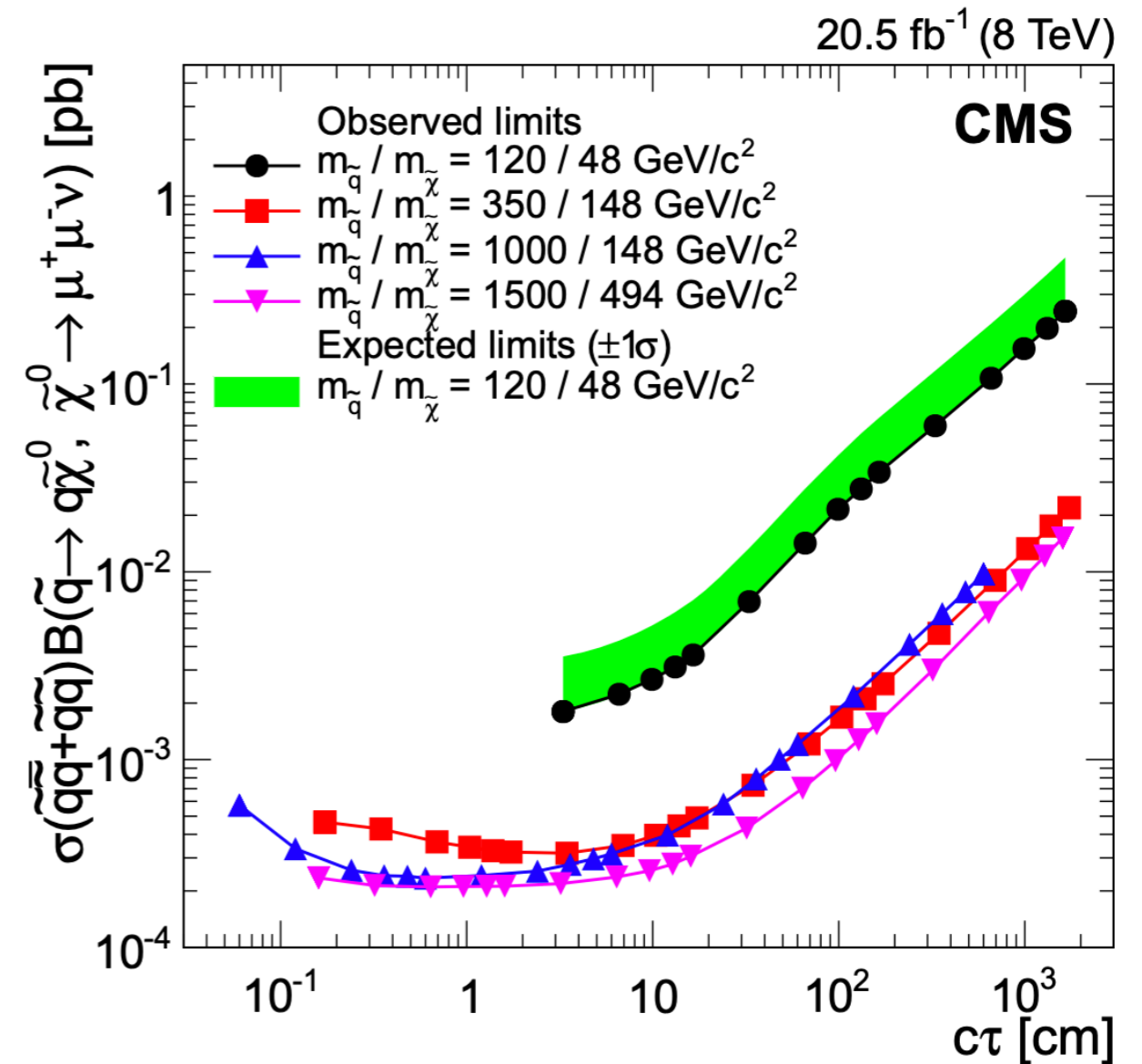
- **ATLAS-EXOT-2017-28 (13 TeV, 36 fb⁻¹)**
 - Displaced lepton - jets analysis
 - Electron and muon LJ
- Prompt analysis as sensitive as CMS analysis
- Displaced analysis 8 TeV not sensitive; 13 TeV potentially sensitive
 - 13 TeV analysis hard to reinterpret



- LHCb-PAPER-2016-047: (7+8 TeV, 3 fb⁻¹)
 - ‘Inclusive displaced vertex search’
 - Trigger muons $p_T > 10$ GeV
 - Final state muon and two jets
 - $p_T(\mu) > 12$ GeV, $d_{IP} > 0.25$ mm, $R_{xy} > 0.55$ mm
 - Invariant mass of tracks > 4.5 GeV
 - Interpretation in terms of GUT scale SUSY RPV models



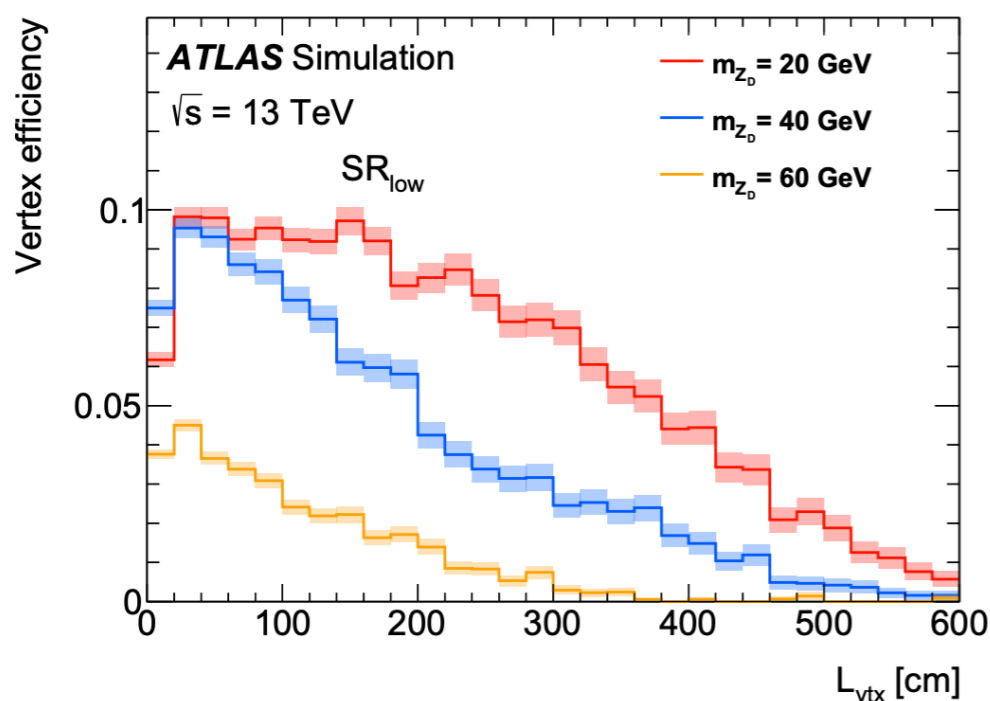
- CMS-EXO-12-037: (8 TeV, 20 fb⁻¹)
 - Inclusive displaced vertex search for pair of electron or muon final states
 - Electron $E_T > 36$ (22) GeV; Muon $p_T > 23$ GeV (reconstructed in muon detectors)
 - Generated $L_{\text{vtx}} < 50$ cm
 - $p_T(\mu) > 12$ GeV, $d_{\text{IP}} > 0.25$ mm, $R_{xy} > 0.55$ mm
 - Interpretation for three body decays



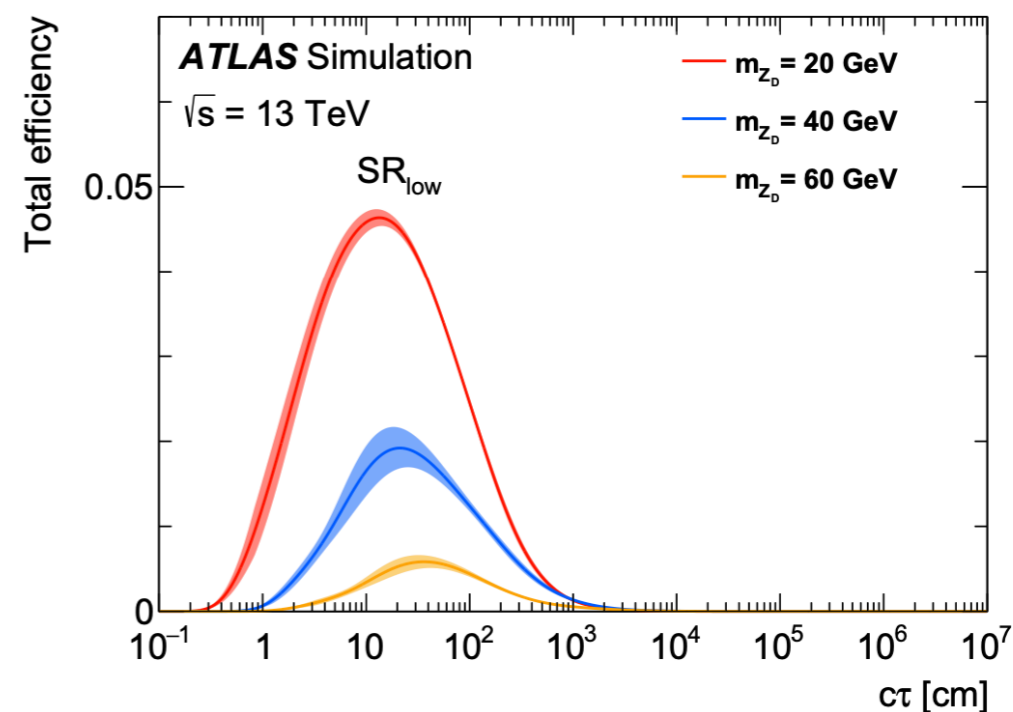
- **ATLAS-EXOT-2017-03 (13 TeV, 32.9 fb⁻¹)**
 - Inclusive search in displaced muon vertex

Signal type	Trigger	Description	Thresholds
High mass	E_T^{miss} single muon	missing transverse momentum single muon restricted to the barrel region	$E_T^{\text{miss}} > 110$ GeV muon $ \eta < 1.05$ and $p_T > 60$ GeV
Low mass	collimated dimuon trimuon	two muons with small angular separation three muons	p_T of muons > 15 and 20 GeV and $\Delta R_{\mu\mu} < 0.5$ $p_T > 6$ GeV for all three muons

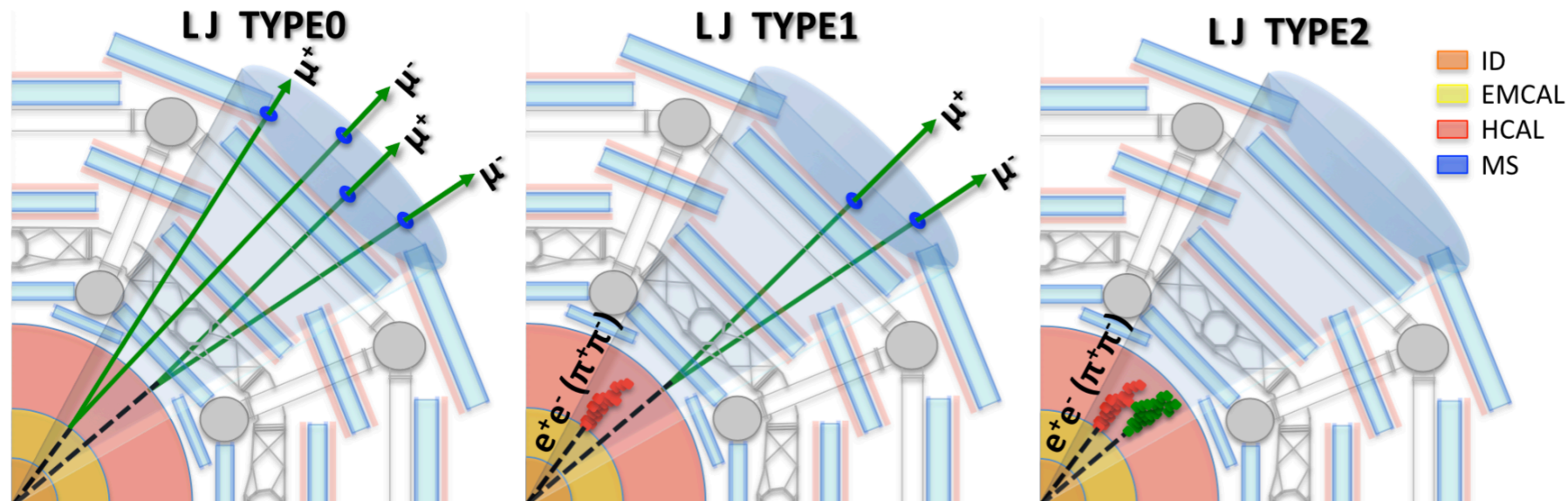
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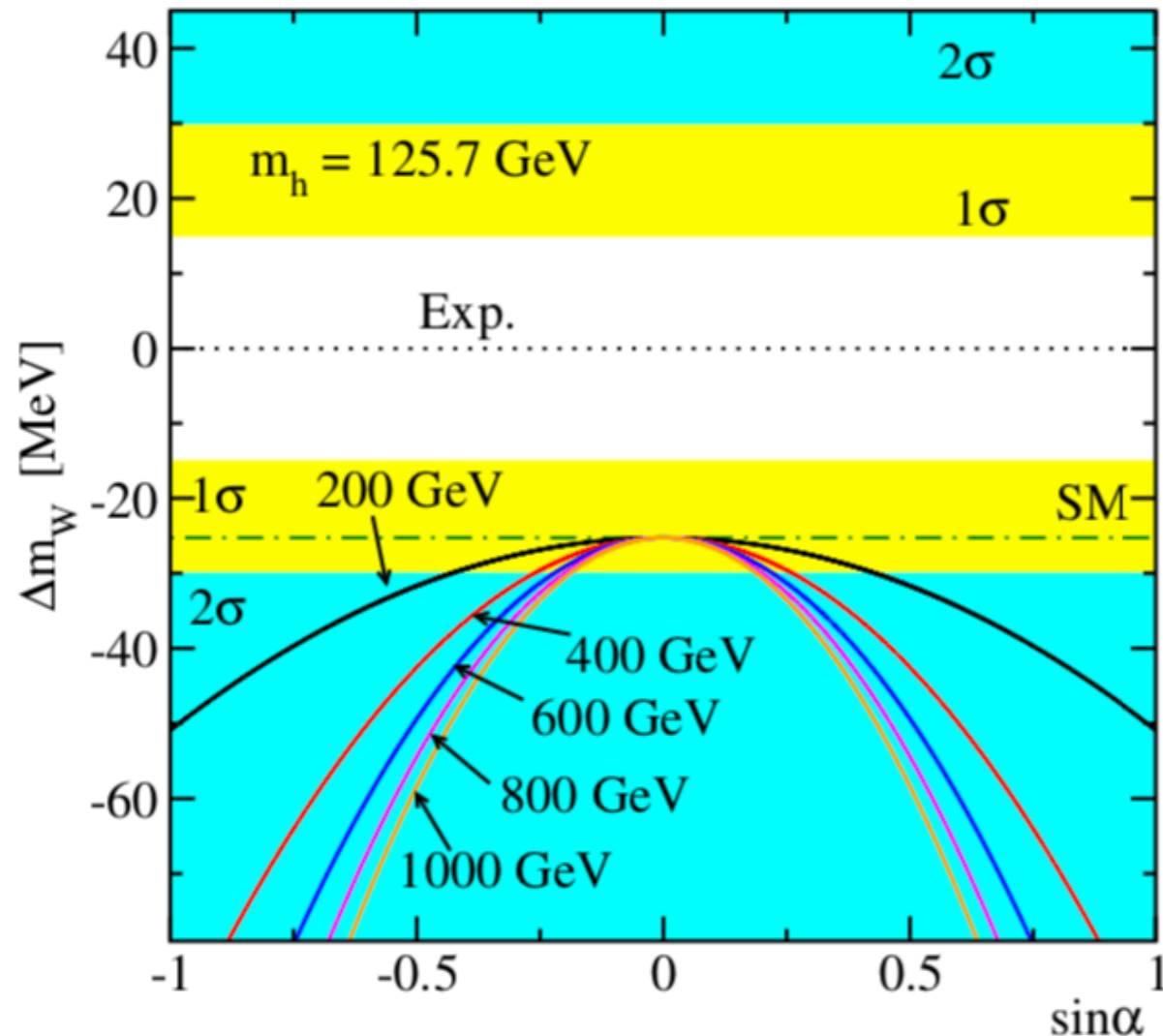


- **ATLAS-EXOT-2013-22 (sqrt 13 TeV, 20 fb⁻¹)**



- Categorization of lepton jets:
 - Electron-jet if at least one electron candidate with $E_T > 10$ GeV, 2 or more tracks w/ $p_T > 10$ GeV, no muons
 - Muon-jet if at least 2 muons with $p_T > 10$ GeV and no electrons
 - Mixed-jet if at least one electron w/ $E_T > 10$ GeV and at least one muon with $p_T > 10$ GeV
- Triggers:
 - Single e w/ $E_T > 60$ or double e w/ $E_T > 35/25$ GeV
 - Single μ w/ $p_T > 36$ or double μ w/ $p_T > 13/13$ GeV
- **No equivalent CMS electron LJ search yet**

Lopez-Val, Robens arXiv:1406.1043



$$m_W^2 \left(1 - \frac{m_W^2}{m_Z^2} \right) = \frac{\pi \alpha_{EM}}{\sqrt{2} G_F} (1 + \Delta r)$$

$$\Delta m_W = -\frac{1}{2} m_W \frac{\sin^2 \theta_W}{\cos^2 \theta_W - \sin^2 \theta_W} \delta(\Delta r)$$

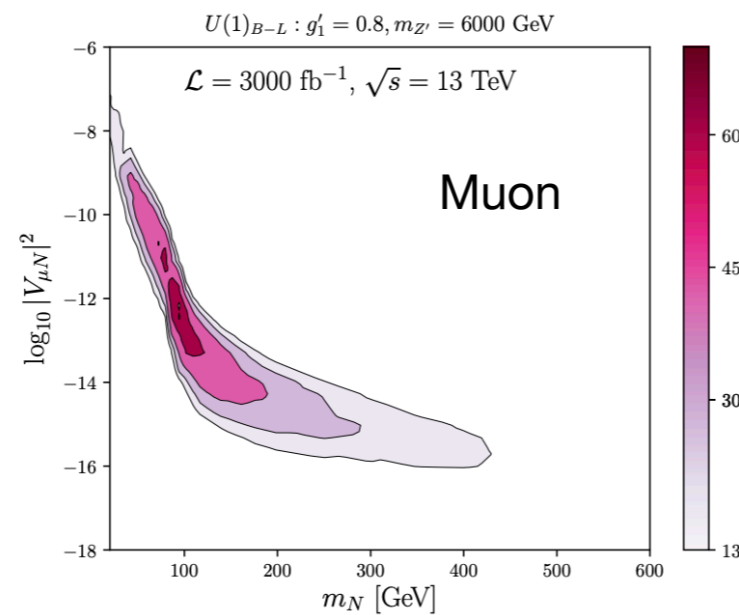
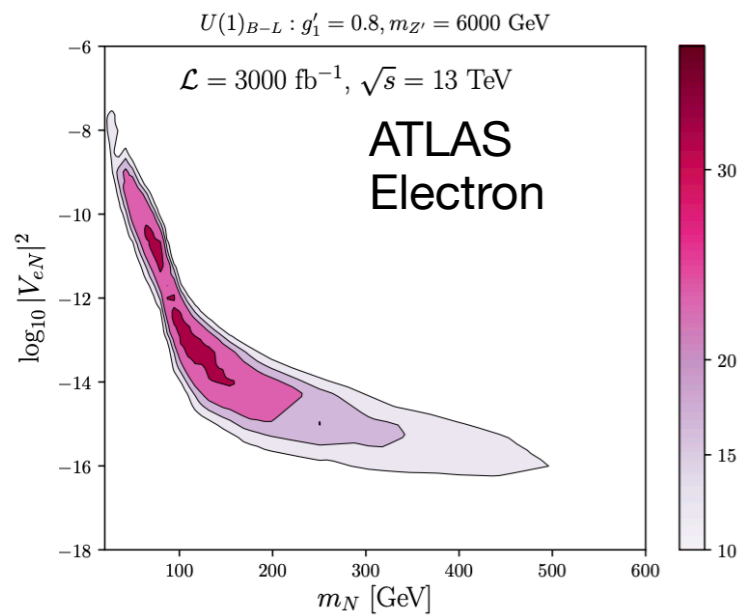
- Constraints can be derived when lighter or heavier Higgs is 125 GeV
- Much stronger constraints when lighter Higgs is 125 GeV and heavier Higgs is heavy
- Driven by discrepancy between observed and predicted value of W mass
- When lighter Higgs is at 125 GeV, higher order EW corrections increase the discrepancy
- When heavy Higgs is at 125 GeV, somewhat better situation however, it is strongly constrained by Higgs signal strengths

ATLAS

Trigger	Muon: $ \eta < 1.07$ and $p_T > 55$ GeV. Electron: $p_T > 120$ GeV
DV region	DV within $4 \text{ mm} < r_{DV} < 300 \text{ mm}$ and $ z_{DV} < 300 \text{ mm}$
DV selection	Made from tracks with $ d_0 > 2 \text{ mm}$ and with $p_T > 1 \text{ GeV}$
	DV track multiplicity $N_{trk} \geq 4$ and invariant mass $m_{DV} \geq 5 \text{ GeV}$

CMS

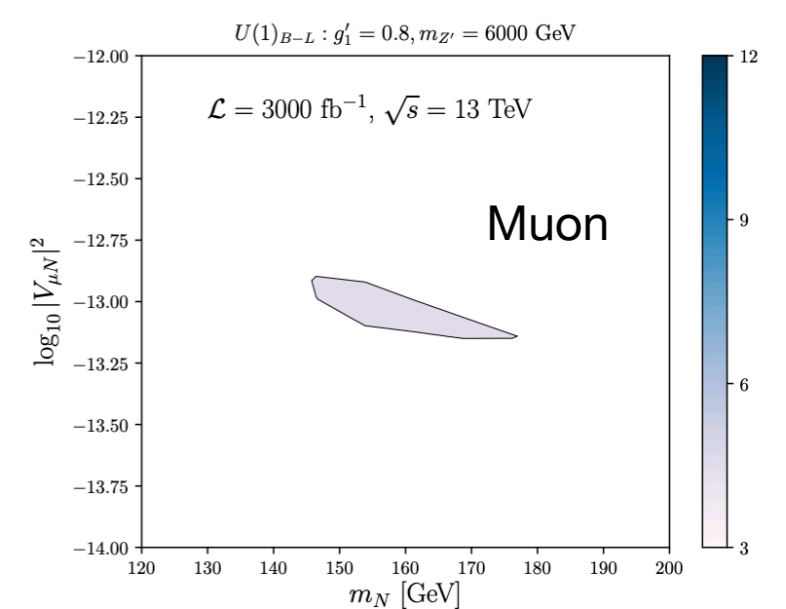
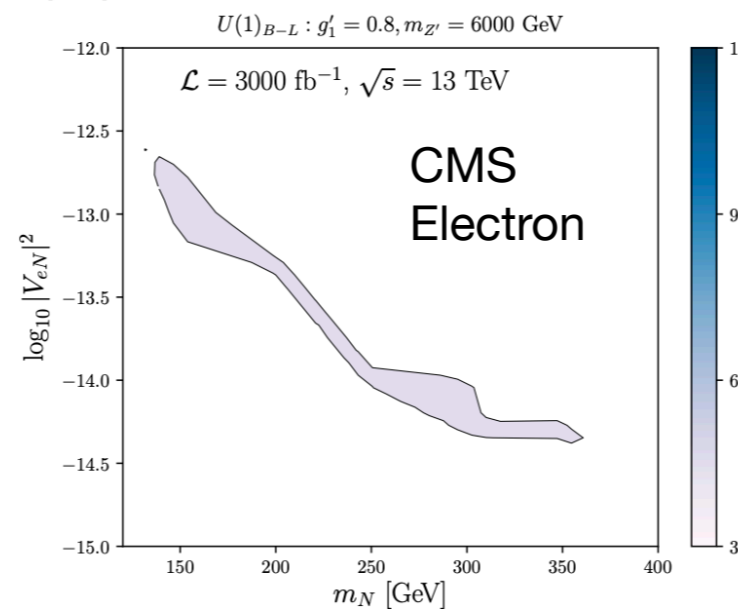
Trigger	$H_T > 1000 \text{ GeV}$
Jet selection	At least 4 jets with $p_T > 20 \text{ GeV}$ and $ \eta < 2.5$
DV region	2 DVs within $0.1 \text{ mm} < r_{DV} < 20 \text{ mm}$ and $d_{VV} > 0.4 \text{ mm}$
DV selection	Made from tracks with $ d_0 \geq 0.1 \text{ mm}$, $p_T > 20 \text{ GeV}$ and $ \eta < 2.5$.
	$\sum p_T \geq 350 \text{ GeV}$, correcting for b quarks.



At least one DV in inner tracker
Efficiencies derived by 'fitting' to the limits

Exactly two DVs

Gen level selection + approximate event level efficiencies



arXiv:1908.09838