

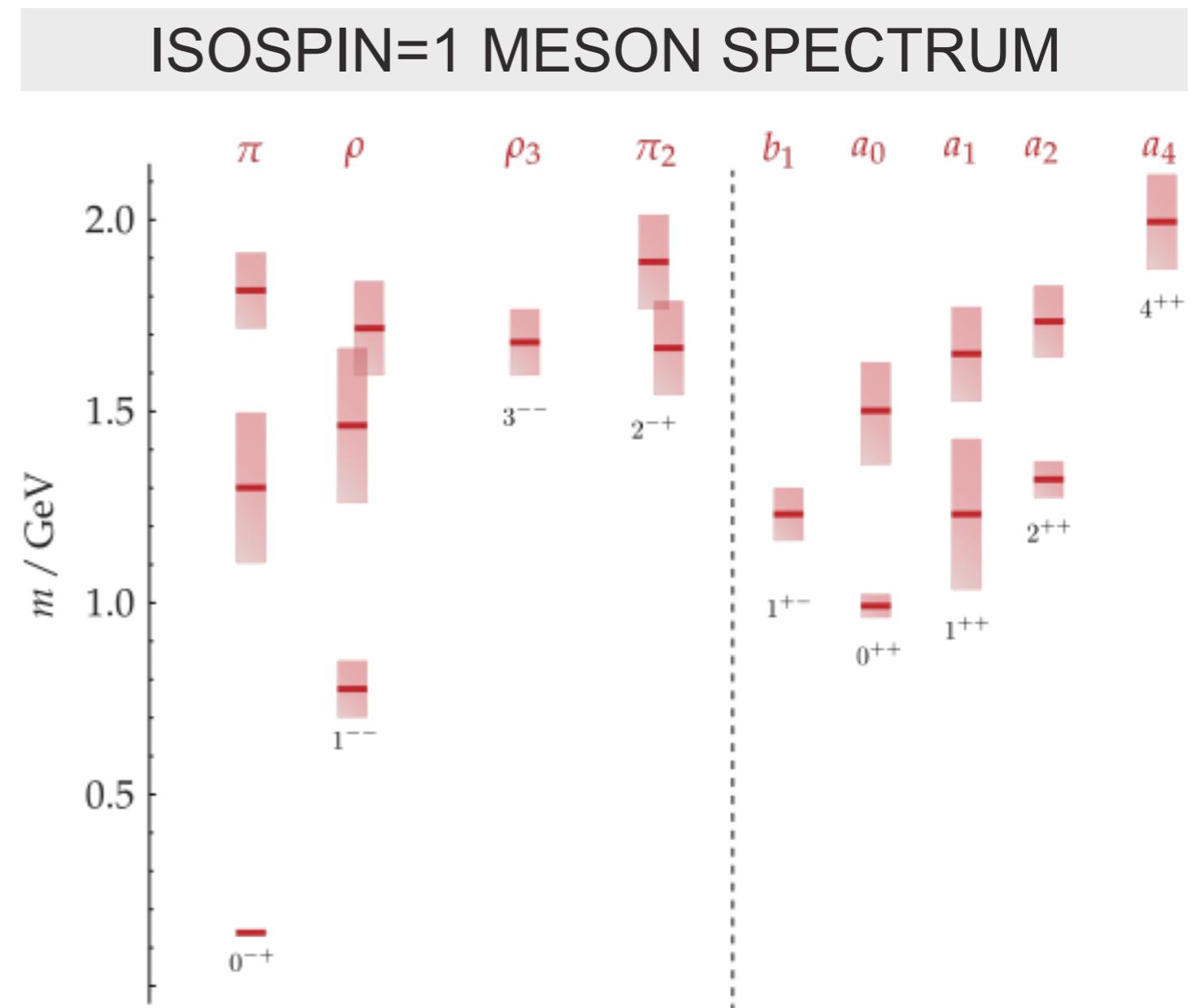
# The power of distillation



Happy 10-th Birthday! Signed, Robert Edwards

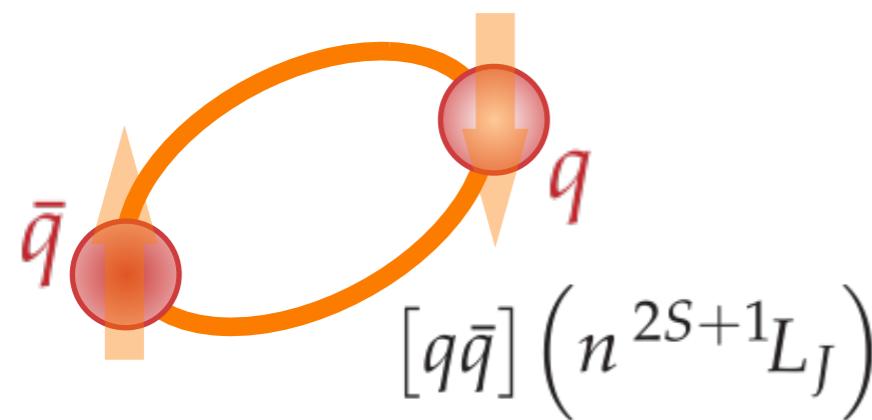
# Experimental meson spectrum

- Mesons classified by their **conserved quantum numbers**
  - Spin, isospin, charge-conjugation **J<sup>P</sup>C**



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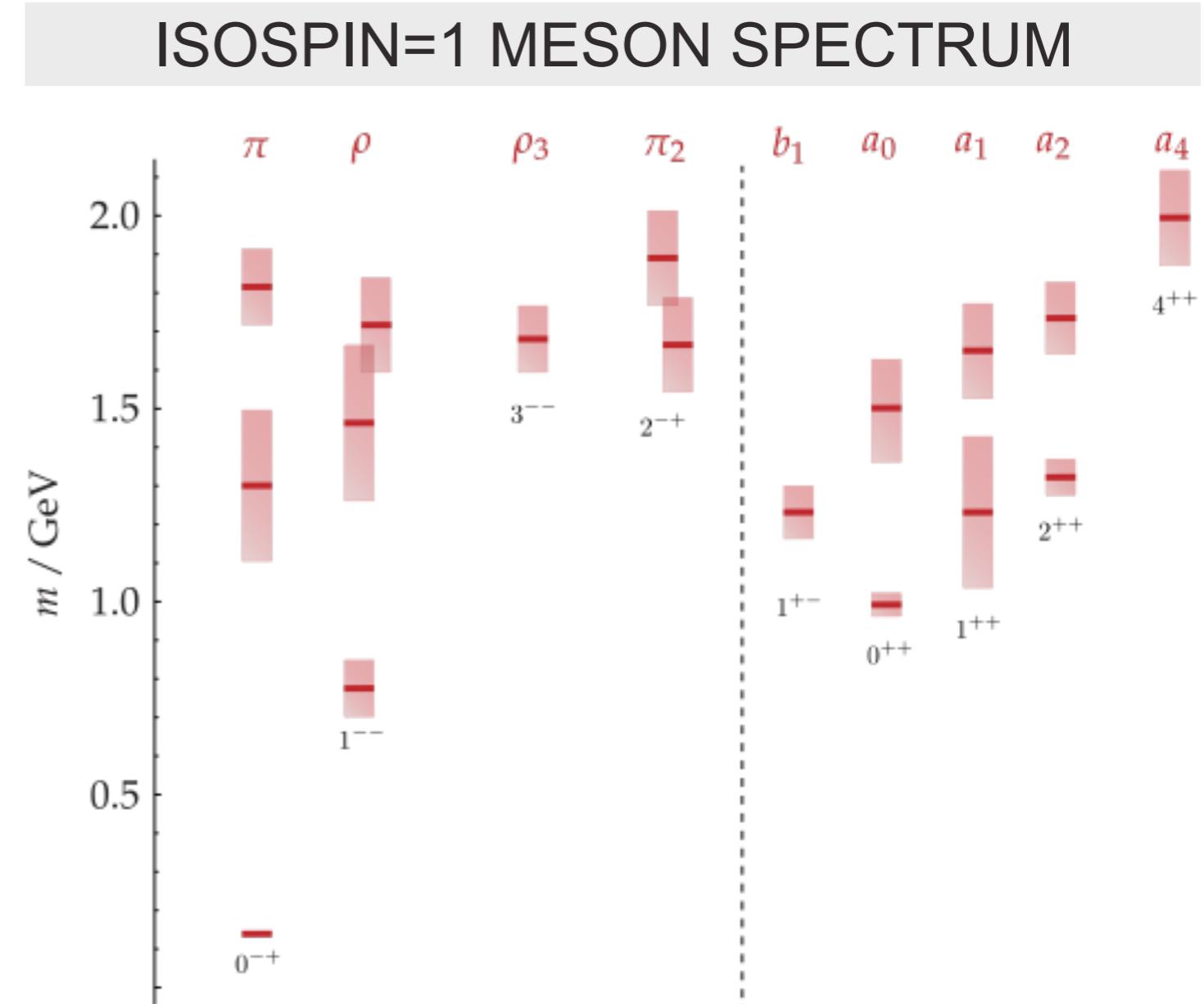


$L = 0 : 0^{-+}, 1^{--}$

$L = 1 : 1^{+-}, (0, 1, 2)^{++}$

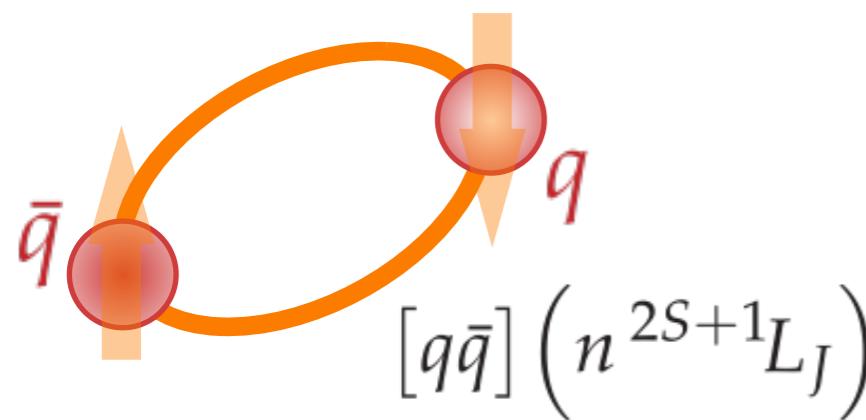
$L = 2 : 2^{-+}, (1, 2, 3)^{--}$

⋮



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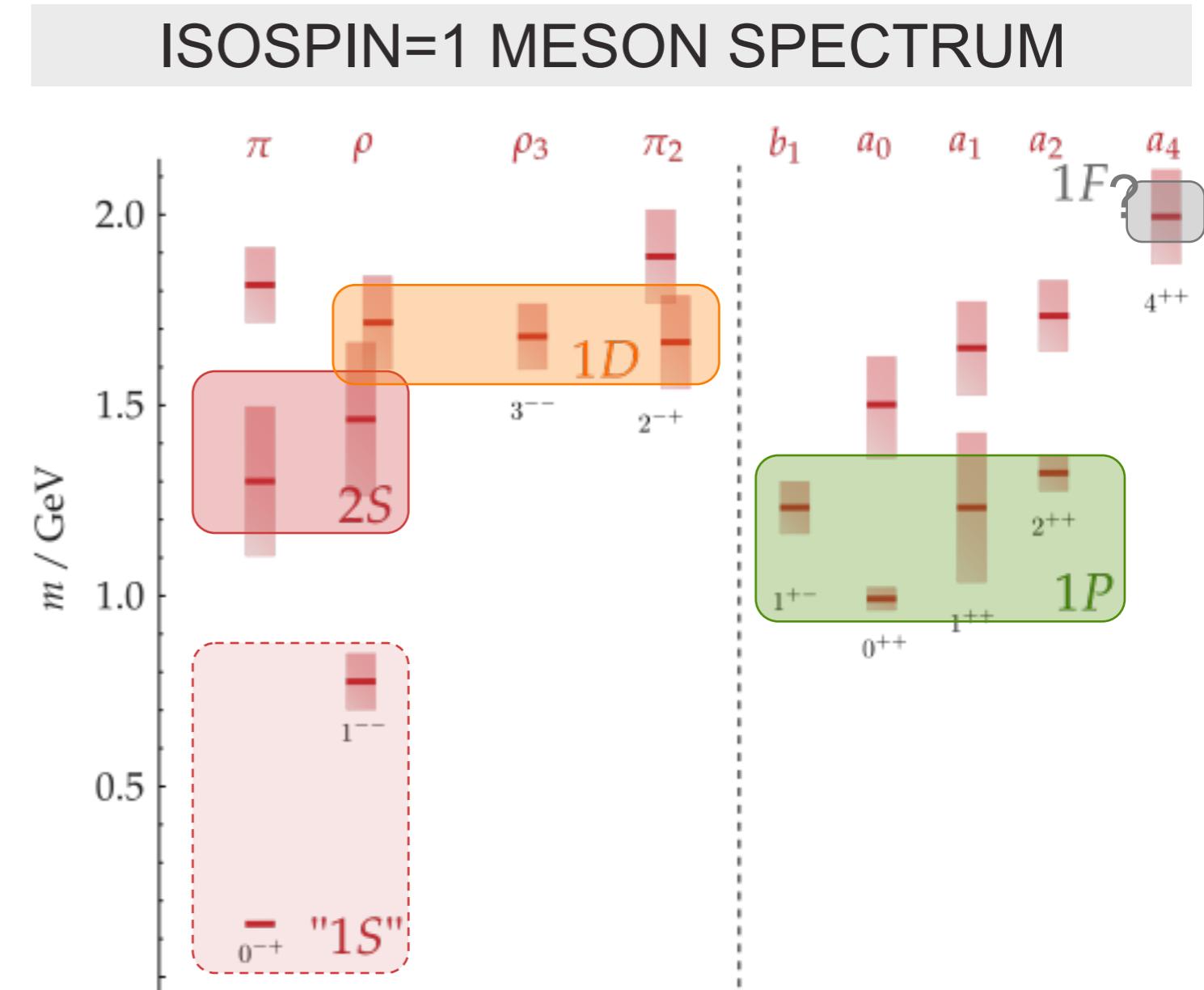


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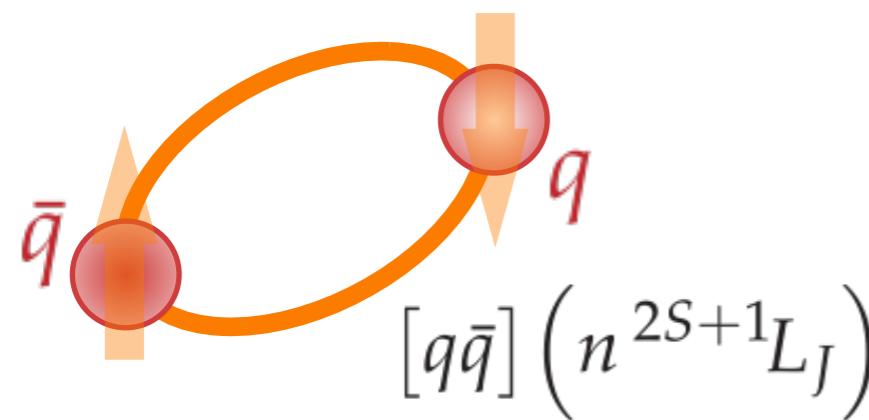
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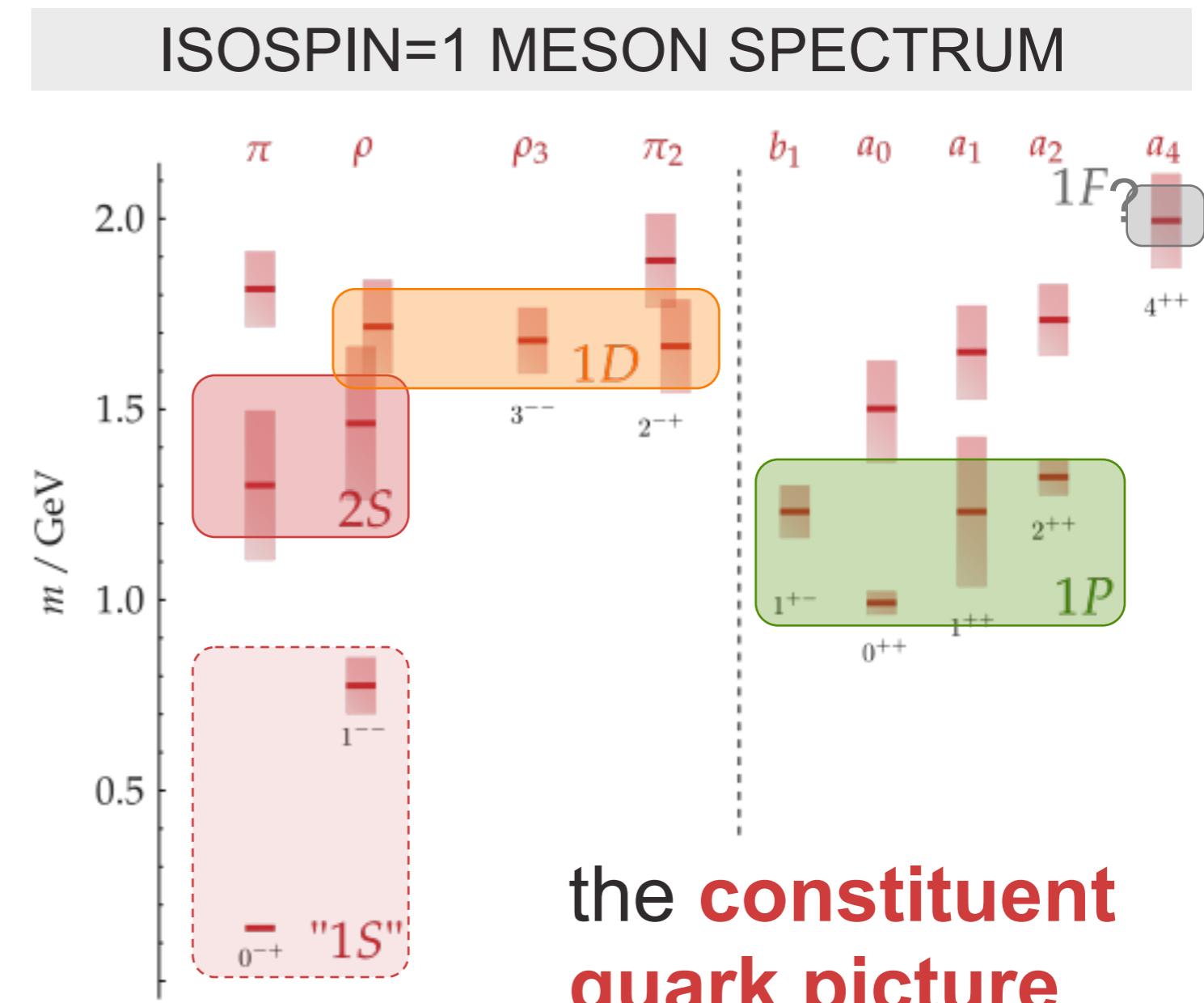
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$L = 1 : 1^{+-}, (0, 1, 2)^{++}$

$L = 2 : 2^{-+}, (1, 2, 3)^{--}$

$\vdots$

**n.b.**  
**absent:**  $0^{--}, 0^{+-}, 1^{-+}, 2^{+-} \dots$

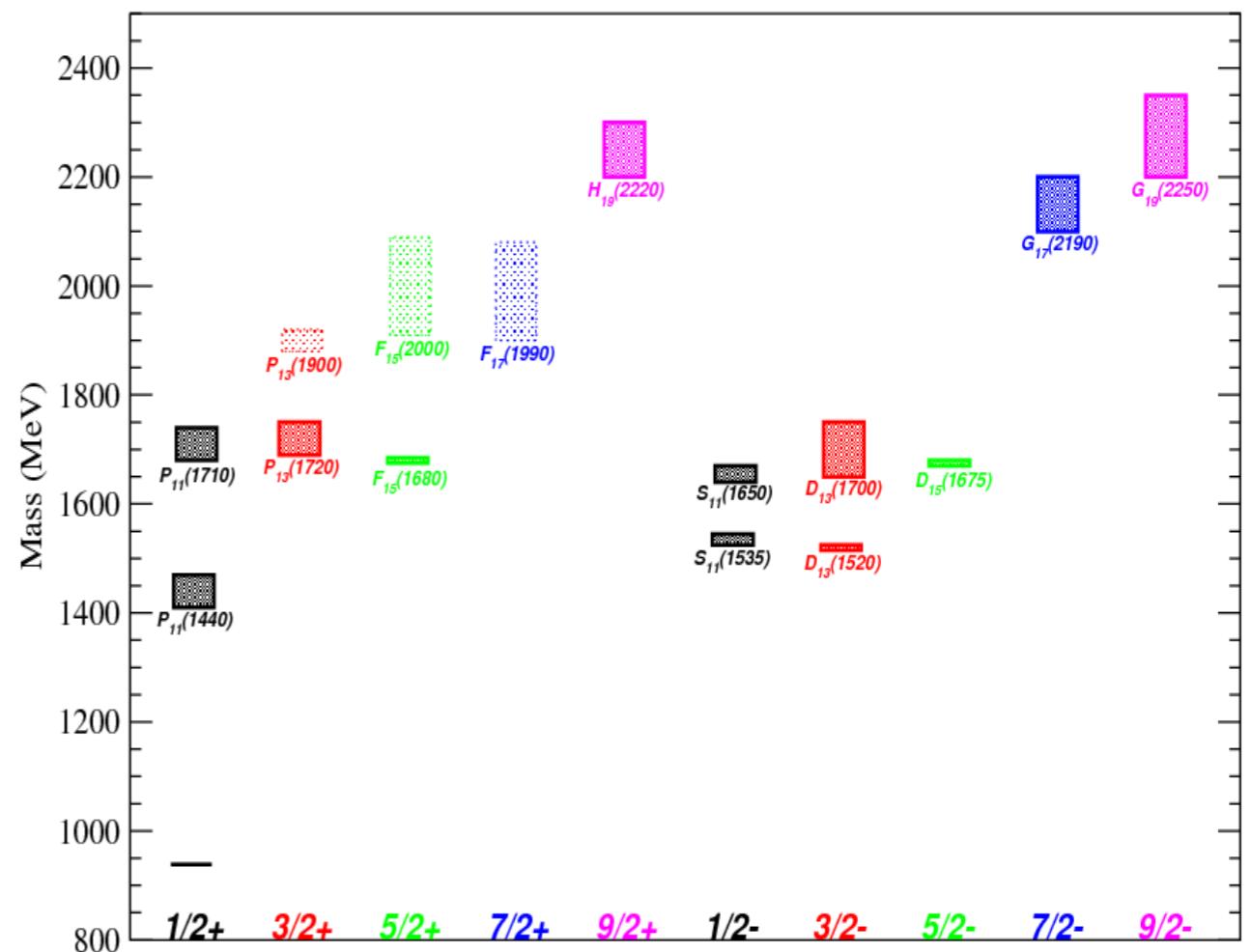


# Experimental baryon spectrum

- Baryons classified by their **conserved quantum numbers**
  - Spin, parity, isospin       $J^P$

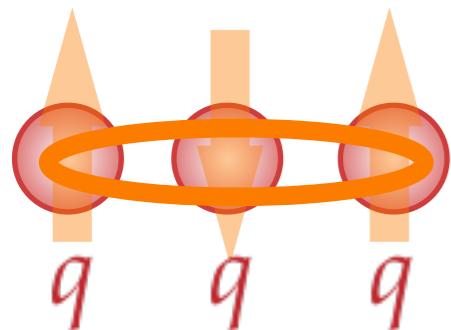
## ISOSPIN=1/2 BARYON SPECTRUM

Nucleon Mass Spectrum (Exp): 4\*, 3\*, 2\*



# Experimental baryon spectrum

- Baryons classified by their **conserved quantum numbers**
  - Spin, parity, isospin       $J^P$



$$[qqq] \left( n^{2S+1} L_\pi \right)$$

Antisymmetric under interchange

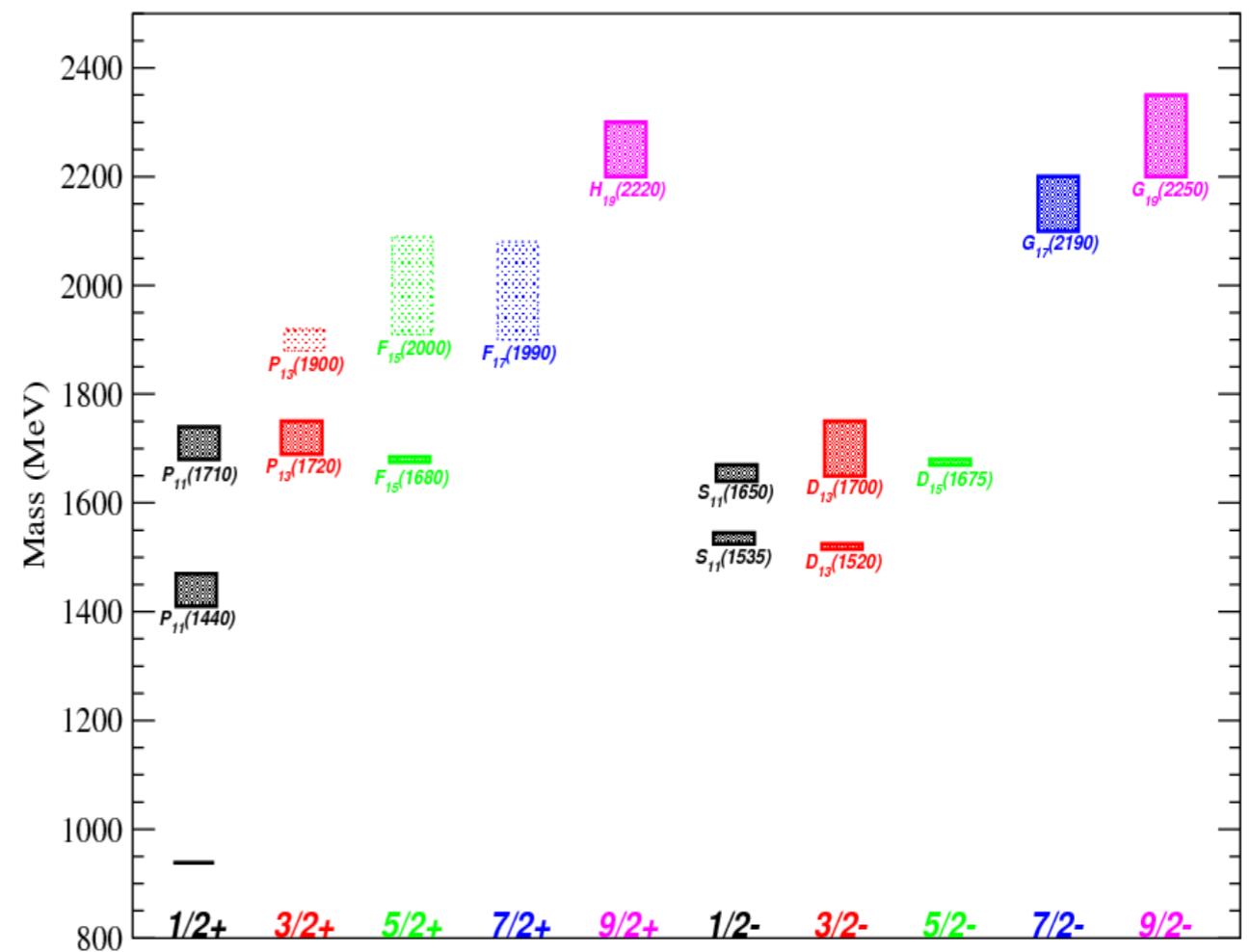
$\pi$  = permutation of quarks in space

$$L = 0_S : \frac{1}{2}^+$$

$$L = 1_M : (\frac{1}{2}, \frac{3}{2}, \frac{1}{2}, \frac{3}{2}, \frac{5}{2})^-$$

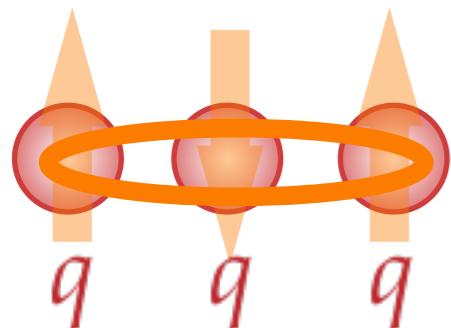
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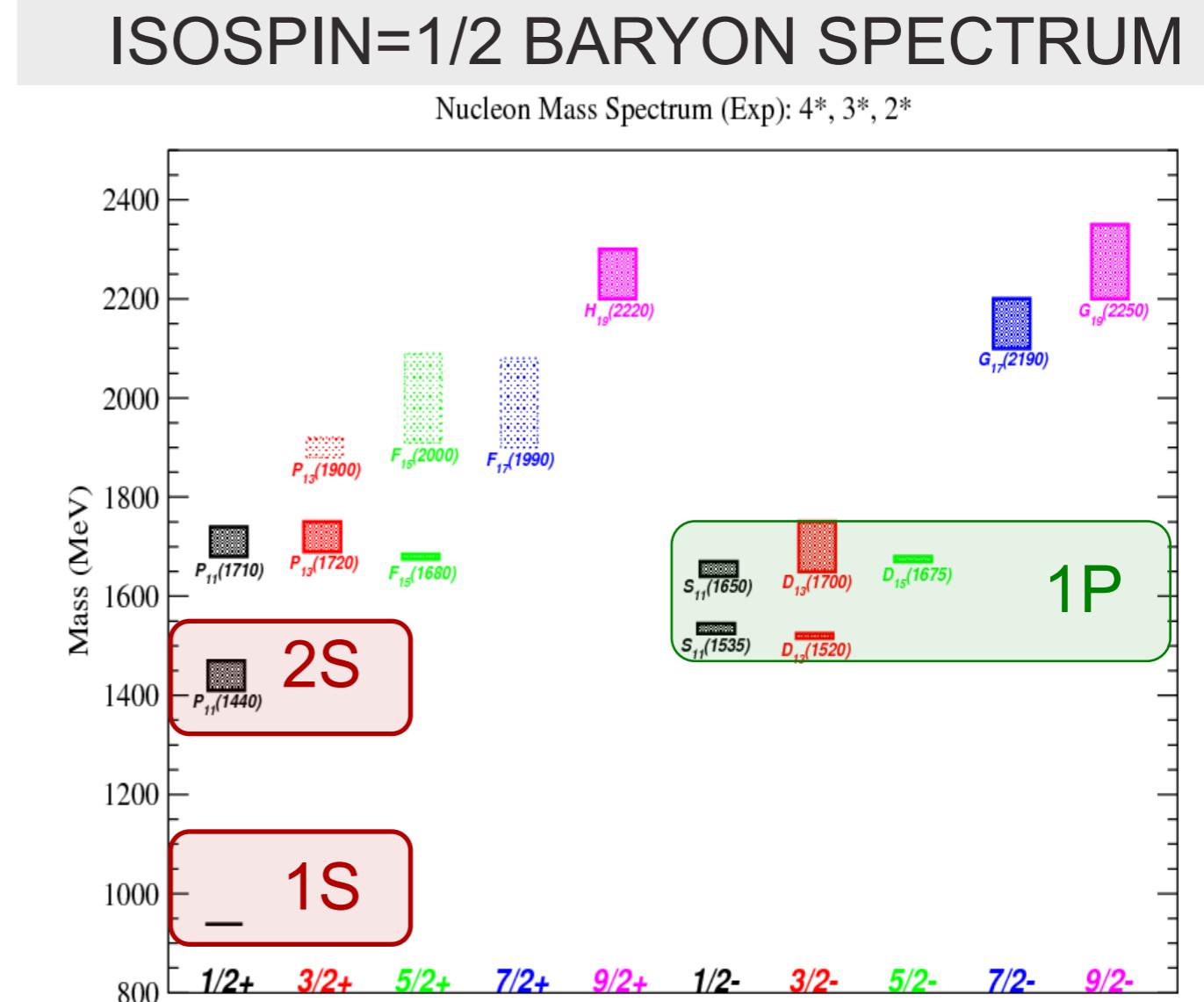
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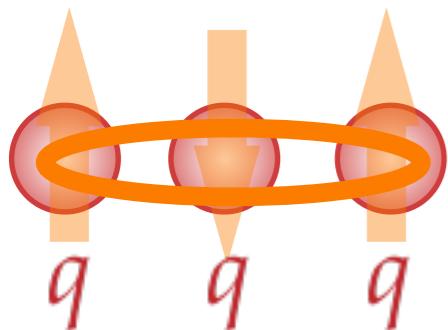
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# Experimental baryon spectrum

- Some states are “missing” ???



$$[qqq] \left( n^{2S+1} L_\pi \right)$$

Antisymmetric under interchange

$\pi$  = permutation of quarks in space

$$L = 0_S : \frac{1}{2}^+$$

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$$L = 2_S : (\frac{3}{2}, \frac{5}{2})^+$$

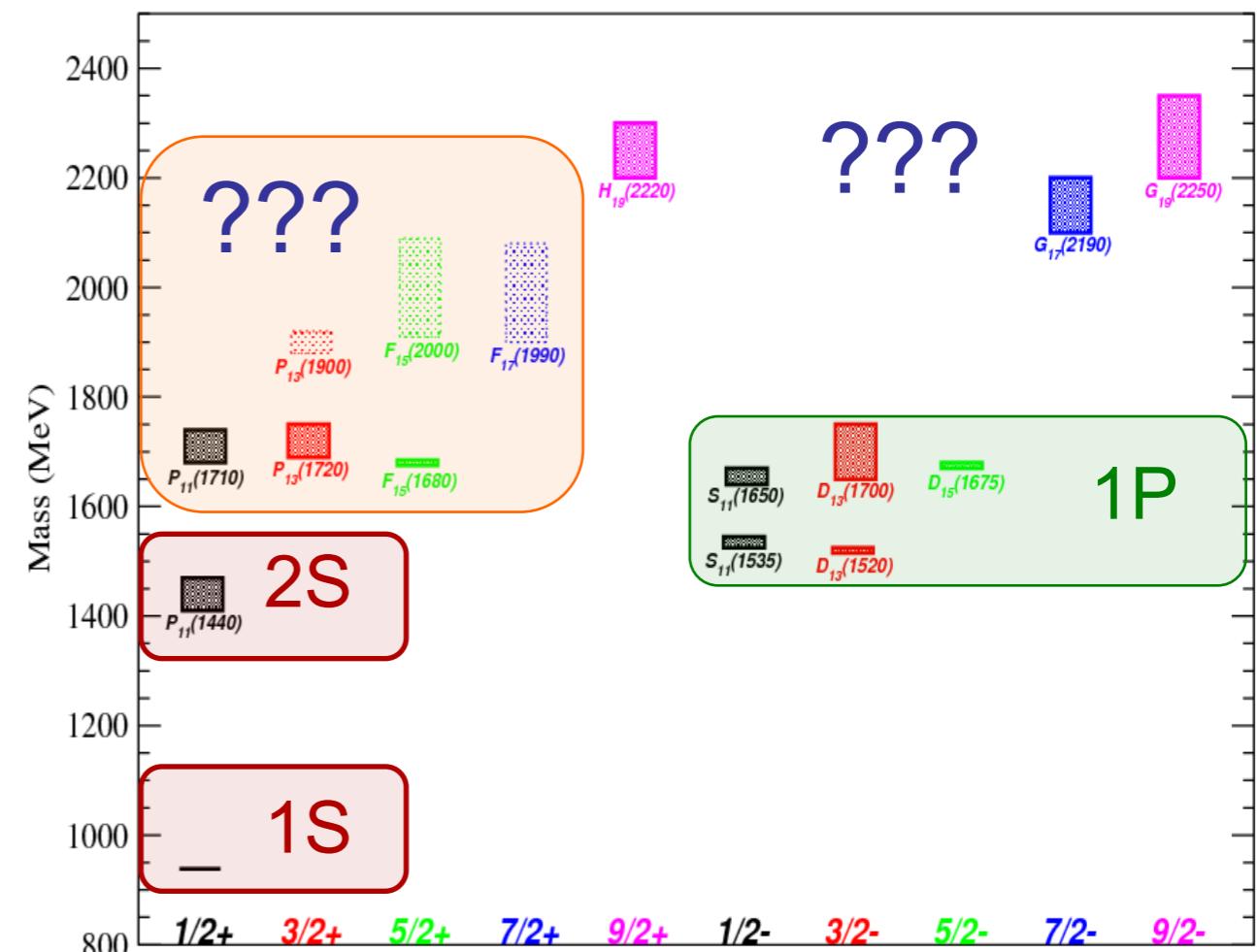
$$L = 2_M : (\frac{3}{2}, \frac{5}{2})^+, (\frac{1}{2}, \frac{3}{2}, \frac{5}{2})^+$$

$$L = 1_A : (\frac{1}{2}, \frac{3}{2})^+$$

$$L = 0_M : (\frac{1}{2}, \frac{3}{2})^+$$

## ISOSPIN=1/2 BARYON SPECTRUM

Nucleon Mass Spectrum (Exp): 4\*, 3\*, 2\*

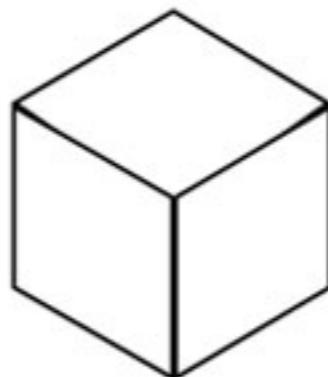


# Finite volume QCD & the hadron spectrum

Compute correlation functions as an average over field configurations

e.g.  $\int \mathcal{D}\psi \mathcal{D}\bar{\psi} \mathcal{D}A_\mu \bar{\psi} \Gamma \psi(t) \bar{\psi} \Gamma \psi(0) e^{-\int d^4x \mathcal{L}_{\text{QCD}}(\psi, \bar{\psi}, A_\mu)}$

‘sum’      ‘field correlation’    ‘probability weight’



*Field integration within a finite, but continuous, hypercube  
Need some kind of ultraviolet regulator....*

Spectrum from two-point correlation functions

$$\begin{aligned} C(t) &= \langle 0 | \mathcal{O}(t) \mathcal{O}^\dagger(0) | 0 \rangle \\ &= \sum_{\mathbf{n}} e^{-E(\mathbf{n})t} \langle 0 | \mathcal{O}(0) | \mathbf{n} \rangle \langle \mathbf{n} | \mathcal{O}^\dagger(0) | 0 \rangle \end{aligned}$$

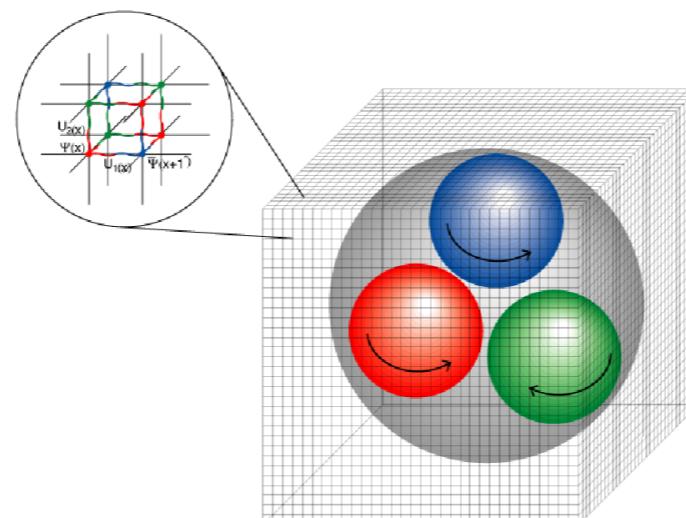
States  $|\mathbf{n}\rangle$  are finite-volume distorted

# Lattice QCD & the hadron spectrum

Compute correlation functions as a Monte Carlo average over field configurations

e.g.  $\int \mathcal{D}\psi \mathcal{D}\bar{\psi} \mathcal{D}A_\mu \bar{\psi} \Gamma \psi(t) \bar{\psi} \Gamma \psi(0) e^{-\int d^4x \mathcal{L}_{\text{QCD}}(\psi, \bar{\psi}, A_\mu)}$

‘sum’      ‘field correlation’    ‘probability weight’



*Discretize the action over sites*

*Serves as an ultraviolet regulator*

Spectrum from two-point correlation functions

$$\begin{aligned} C(t) &= \langle 0 | \mathcal{O}(t) \mathcal{O}^\dagger(0) | 0 \rangle \\ &= \sum_n e^{-E(n)t} \langle 0 | \mathcal{O}(0) | n \rangle \langle n | \mathcal{O}^\dagger(0) | 0 \rangle \end{aligned}$$

States  $|n\rangle$  are finite-volume distorted

# Excited states from correlators

- How to get at excited QCD eigenstates ?

- optimal operator for state  $|\mathbf{n}\rangle$  :  $\Omega_{\mathbf{n}}^+ \sim \sum_i v_i^{(\mathbf{n})} \mathcal{O}_i^+$   
for a basis of meson operators  $\{\mathcal{O}_i\}$

- can be obtained (in a variational sense) from the matrix of correlators

$$C_{ij}(t) = \langle 0 | \mathcal{O}_i(t) \mathcal{O}_j^\dagger(0) | 0 \rangle$$

- by solving a generalized eigenvalue problem

$$C(t)v^{(\mathbf{n})} = C(t_0)v^{(\mathbf{n})} \lambda_{\mathbf{n}}(t)$$

‘diagonalize the correlation matrix’

eigenvalues

$$\lambda_{\mathbf{n}}(t) \sim e^{-E_{\mathbf{n}}(t-t_0)}$$

- a large basis can be constructed using covariant derivatives :

$$\mathcal{O} \sim \bar{\psi} \Gamma \overleftrightarrow{D} \dots \overleftrightarrow{D} \psi$$

# Operators - quark bilinears

- Scalar/Vector Dirac structures & covariant derivatives

$$\bar{\psi} \gamma_5 \psi \rightarrow J = 0$$

$$\bar{\psi} \gamma_i \psi \rightarrow J = 0$$

$$\bar{\psi} \overleftrightarrow{D}_i \psi \rightarrow J = 1$$

- ⇒ Combine gamma & derivatives

$$\bar{\psi} \Gamma \overleftrightarrow{D} \overleftrightarrow{D} \psi$$

in terms of

$$\langle 1m_1 1m_2 | Jm \rangle \bar{\psi} \overleftrightarrow{D}_{m_1} \overleftrightarrow{D}_{m_2} \psi \rightarrow J = 0, 1, 2$$

$$\langle 1s Ll | Jm \rangle \langle 1l_1 1l_2 | Ll \rangle \bar{\psi} \Gamma_s \overleftrightarrow{D}_{l_1} \overleftrightarrow{D}_{l_2} \psi \rightarrow J = 0, 1, 2, 3$$

- Subduction of continuum to cubic reps.

$$\mathcal{O}_{\Lambda\lambda}^{[J]} = \sum_M S_{\Lambda\mu}^{JM} \mathcal{O}^{JM}$$

# Operators - three quarks

---

- Baryons operators are projectors acting on flavor, Dirac spin, and spatial indices

$$(\text{Flavor}_{\pi_F} \otimes \text{Spin}_{\pi_S} \otimes \text{Space}_{\pi_D}) \{\psi_1 \psi_2 \psi_3\}$$

- Symmetric under quark permutations  $\pi$  & color is antisymmetric
- Now CG-s in permutations  $\pi$  and coupling spin and derivatives/space

e.g., two derivatives

$$\mathcal{O}_i \sim (\text{CGCs})_{i,j,k} \{\overrightarrow{D} \overrightarrow{D}\}_j \{\psi \psi \psi\}_k \quad 1 \times 1 \times \mathcal{S} \rightarrow J = \frac{1}{2}, \frac{3}{2}, \frac{5}{2}, \frac{7}{2}$$

# Distillation

- Define a low rank (spatial) smearing operator

e.g., quark bilinear

$$C_{ij}(t) = \langle \bar{\psi} \square \Gamma_t^i \square \psi \cdot \bar{\psi} \square \Gamma_0^j \square \psi \rangle \quad \square_{xy} = V_x V_y^\dagger$$

- Factorize propagators and operator constructions

$$C_{ij}(t) = \text{Tr} [\Phi^i(t) P(t, 0) \Phi^j(0) P(0, t)]$$

$$\Phi_{\alpha\beta}^i(t) = V(t)^\dagger \Gamma_{\alpha\beta}^i(t) V(t) \quad \text{matrix rep. of operator}$$

$$P_{\alpha\beta}(t, 0) = V(t)^\dagger \mathcal{M}_{\alpha\beta}^{-1}(t, 0) V(0) \quad \text{perambulator}$$

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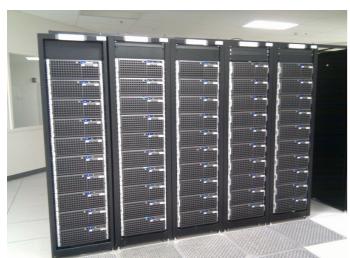
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Multigrid & KNLs+GPUs have been key to construction quark line

arxiv:0905.2160

Trinity 2019

15

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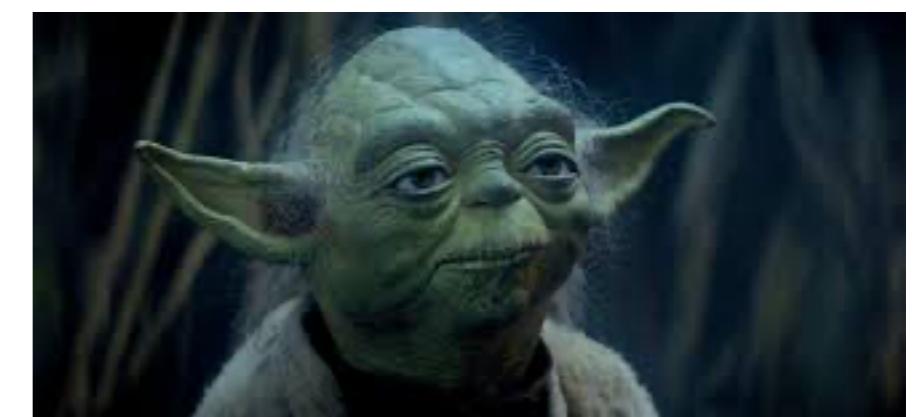
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‘and a powerful ally it is’

arxiv:0905.2160

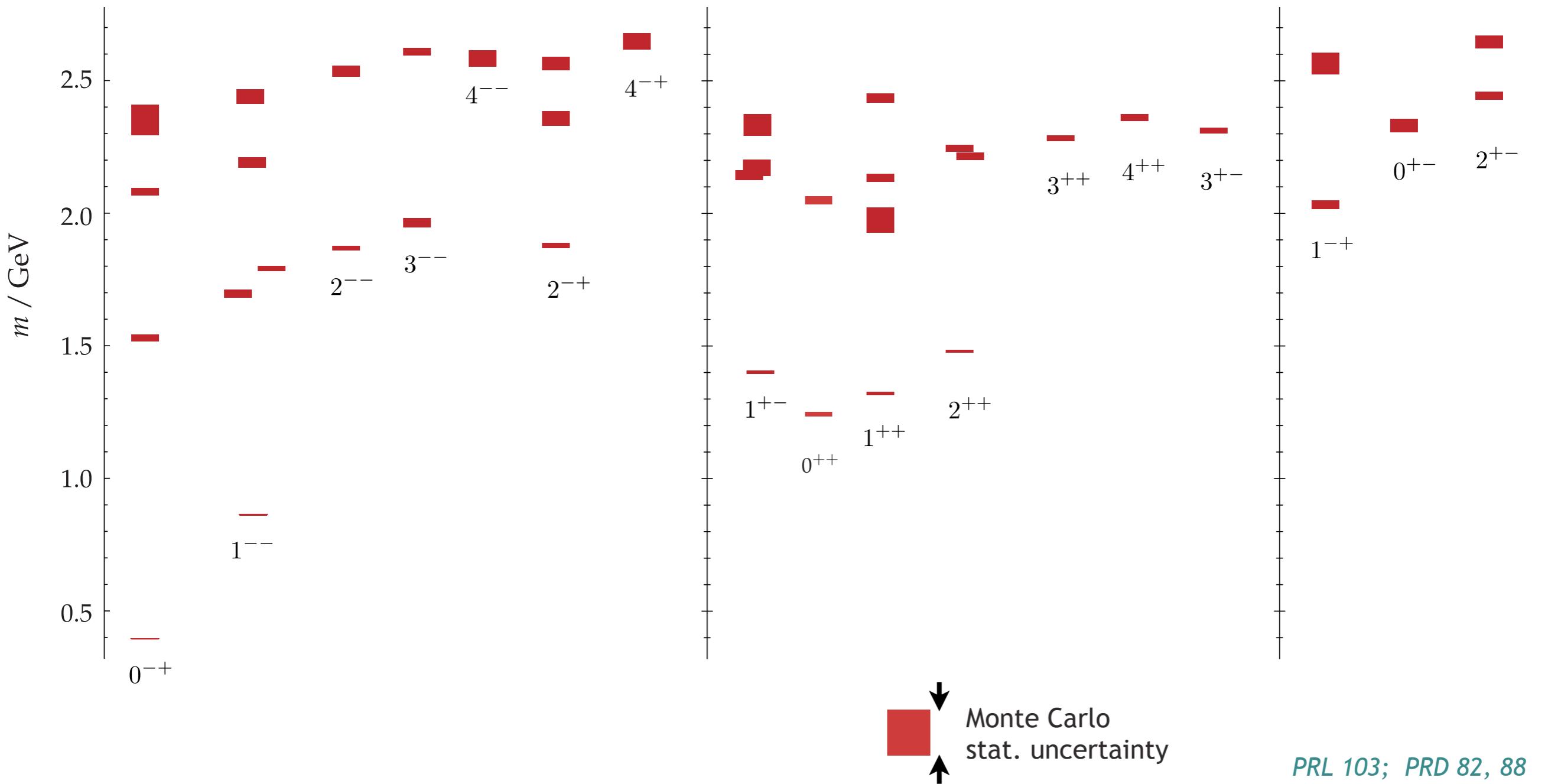
Trinity 2019

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# Glimpse of meson spectrum from lattice QCD

- Appears to be some  $q\bar{q}$ -like near-degeneracy patterns - isovectors

$m_\pi \sim 391$  MeV

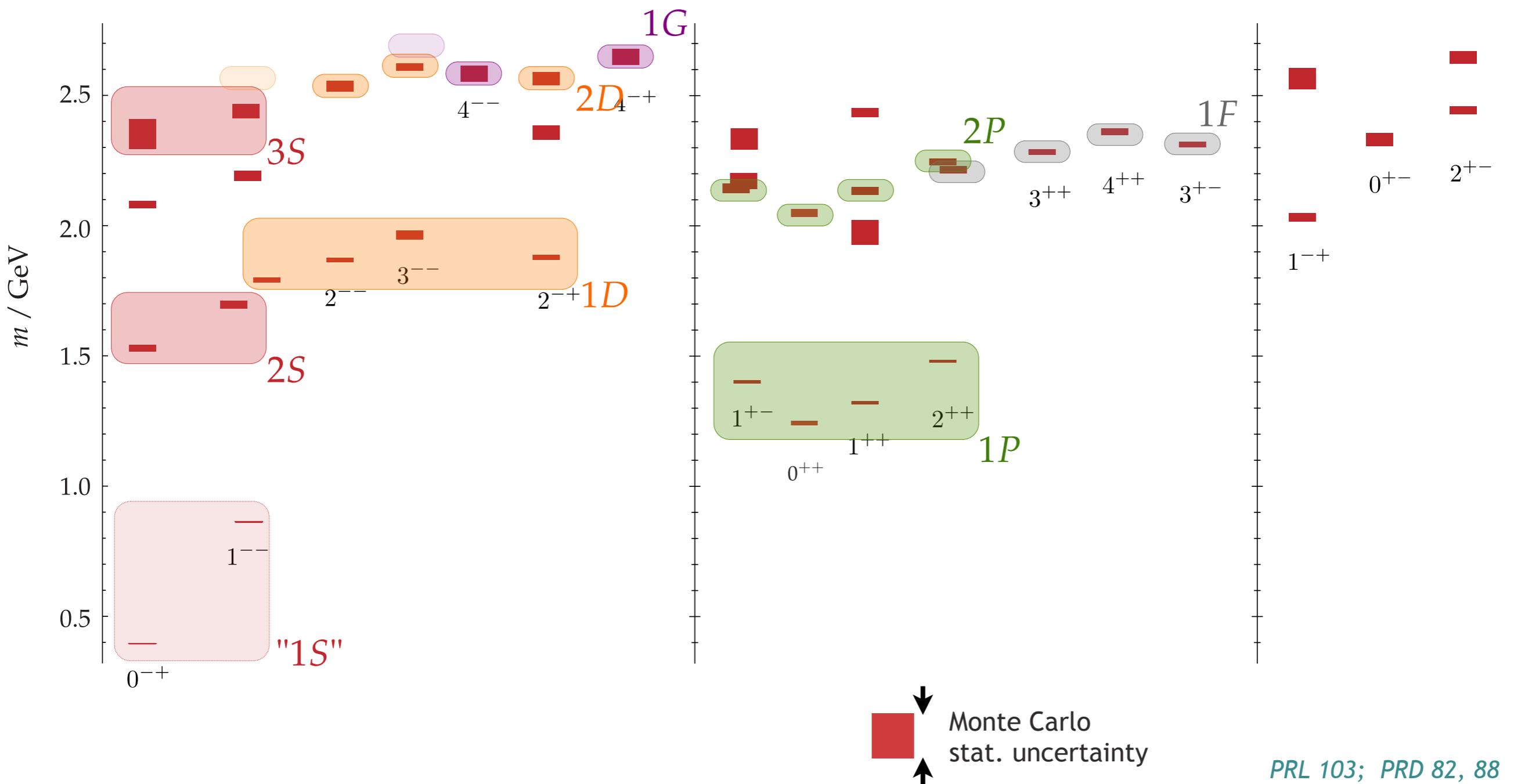


PRL 103; PRD 82, 88

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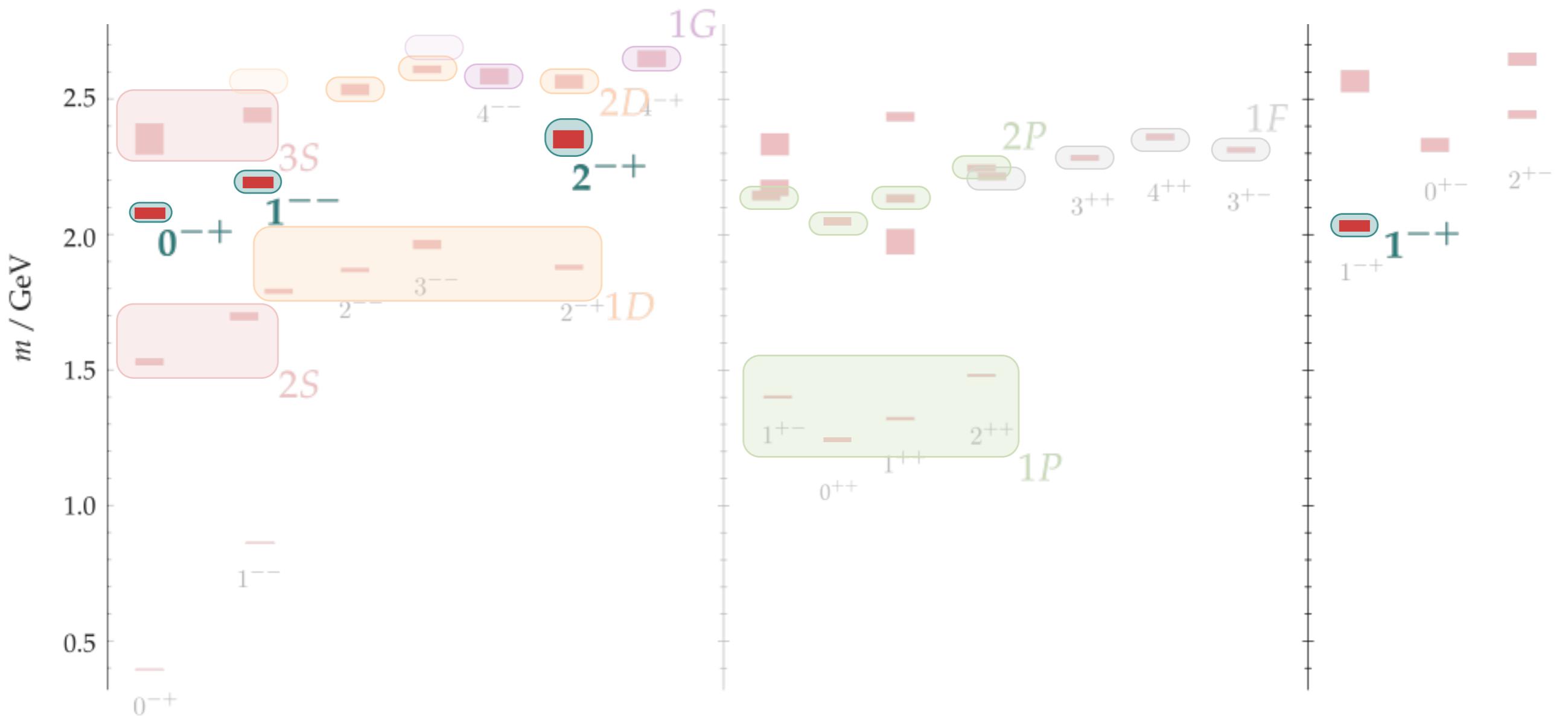
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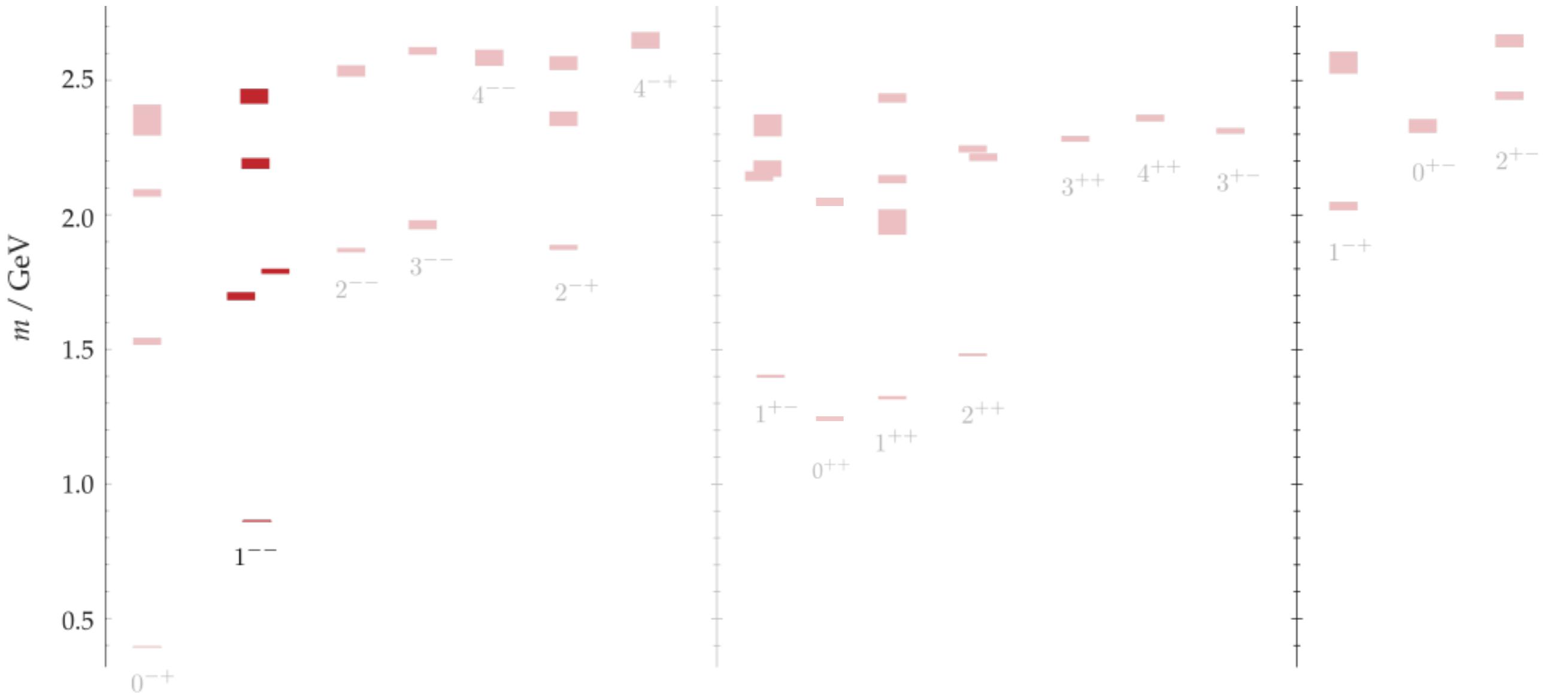
# qq interpretation?

- “Extra” non-exotic states at same energy scale as lightest exotic?

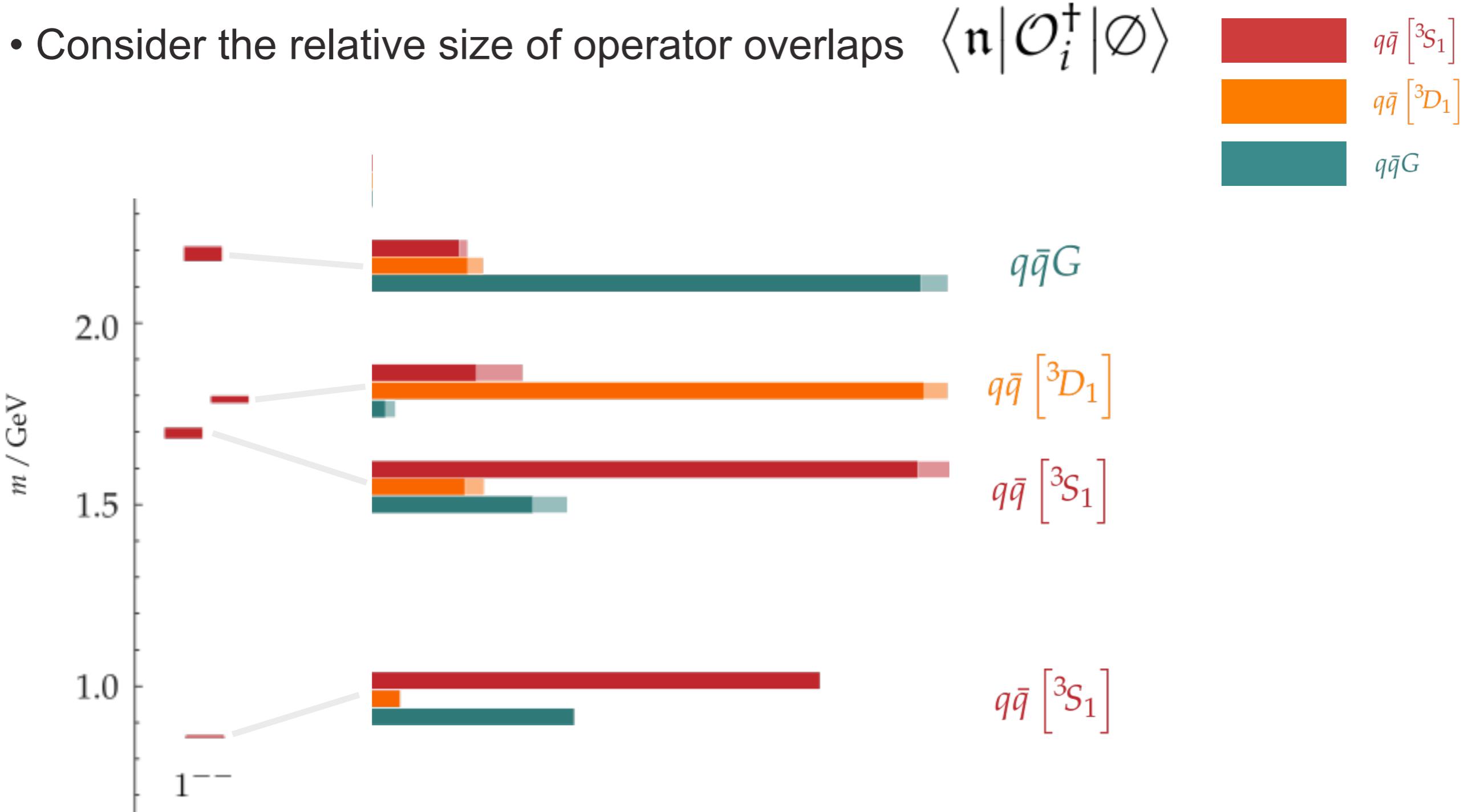


# qq interpretation?

- Consider the relative size of operator overlaps  $\langle \mathbf{n} | \mathcal{O}_i^\dagger | \emptyset \rangle$



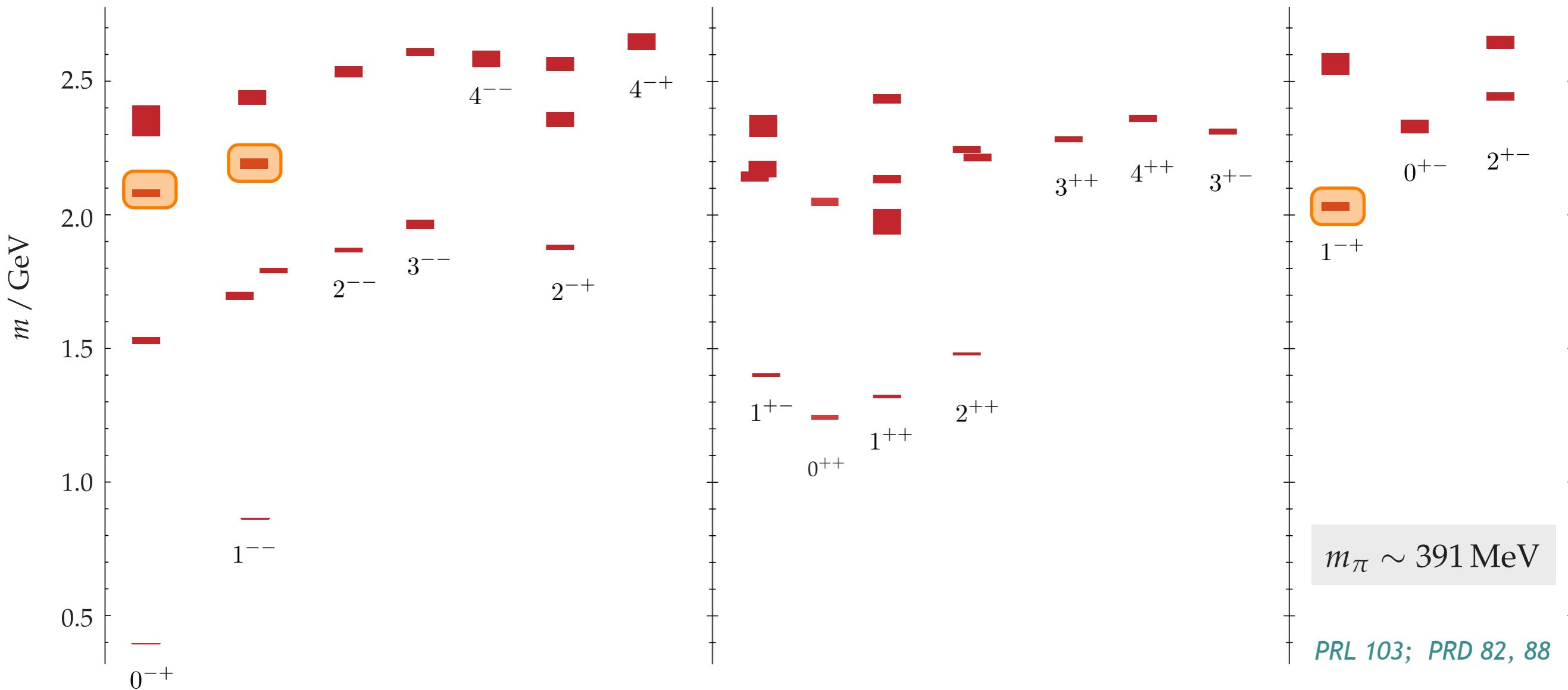
# 1-- operator overlaps



# Glimpse of meson spectrum from lattice QCD

- ‘super’-multiplet of **hybrid mesons** roughly 1.2 GeV above the  $\rho$

$(0, 1, 2)^{-+}, 1^{--}$

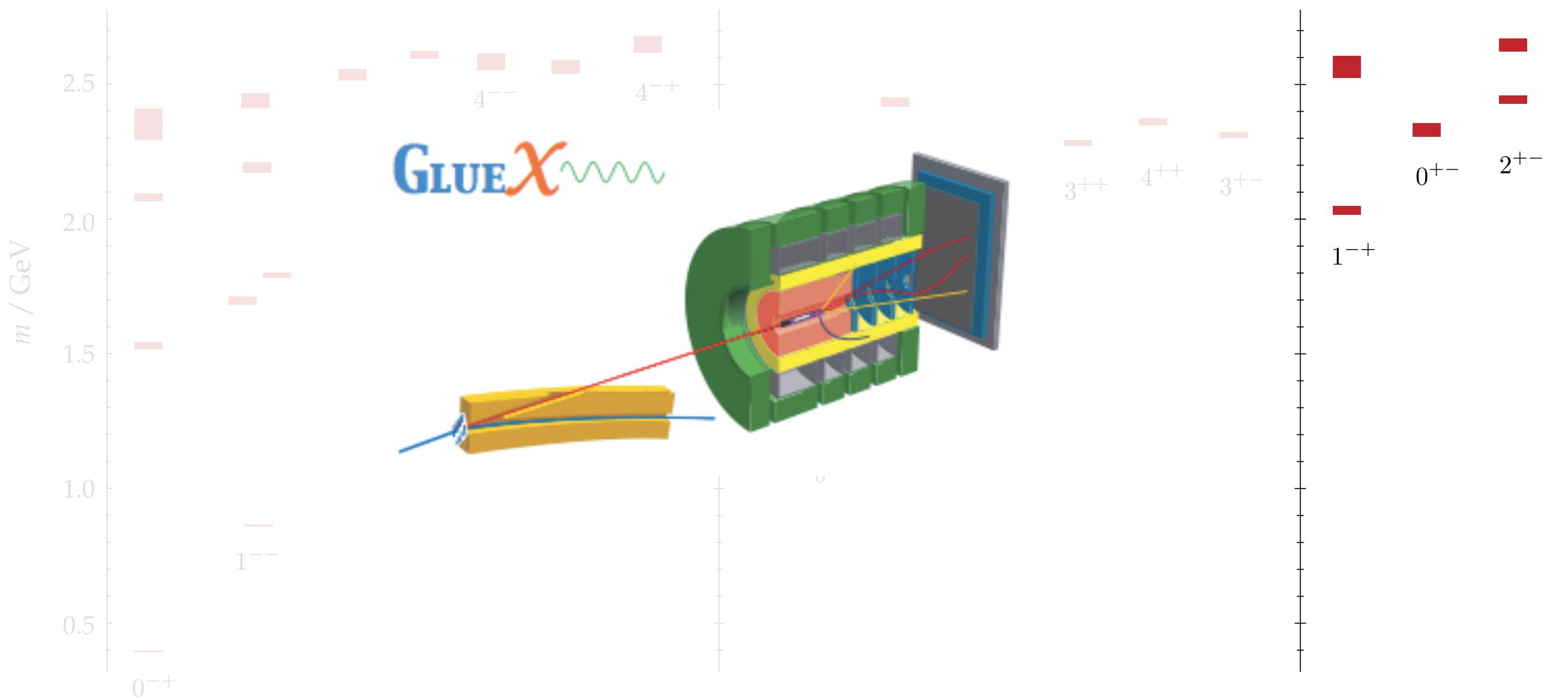


- these states have a dominant overlap onto  $\bar{\psi}\Gamma[D, D]\psi \sim [q\bar{q}]_{8_c} \otimes B_{8_c}$

# Glimpse of meson spectrum from lattice QCD

Multiple exotic mesons within range of GlueX

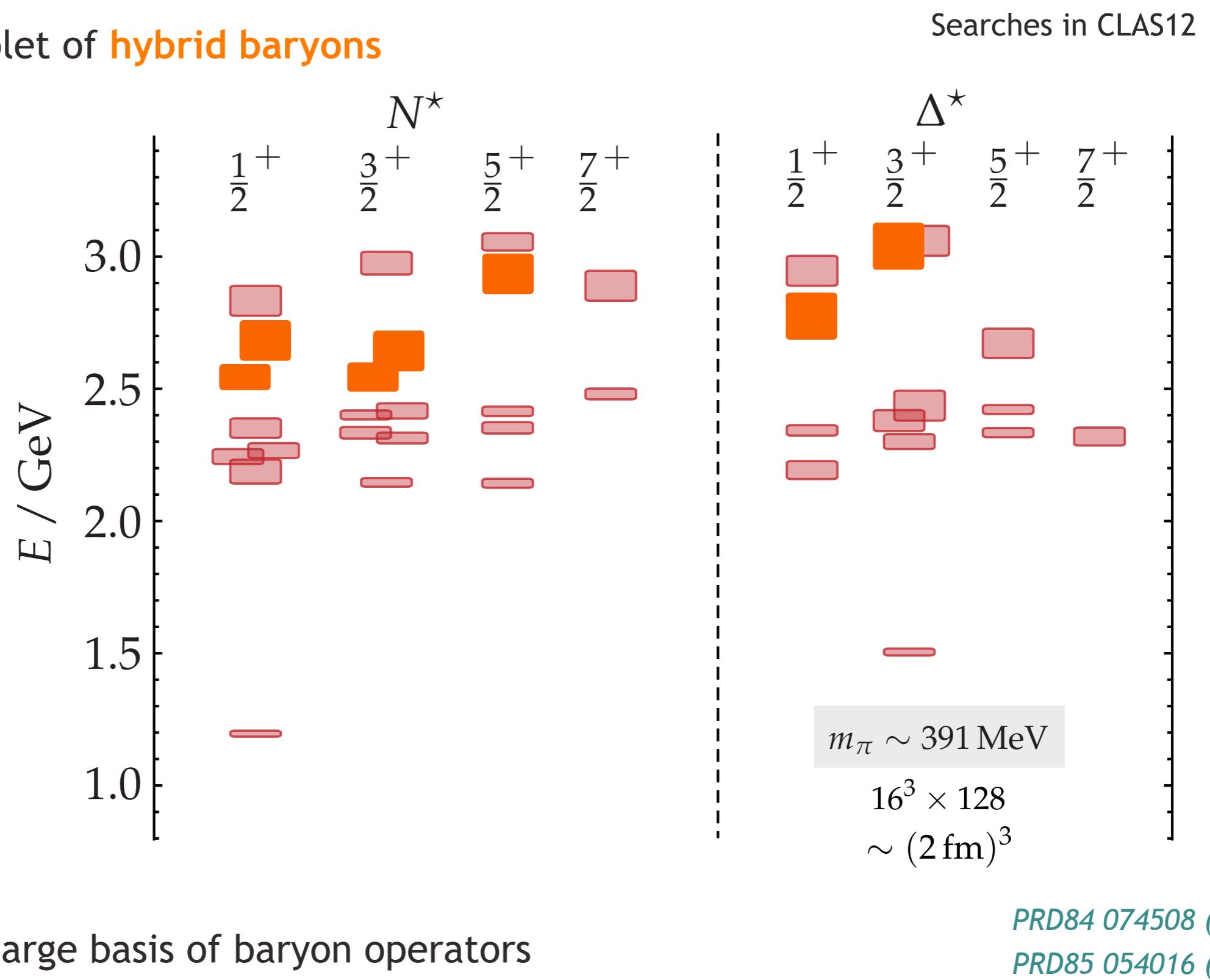
EXOTIC MESONS



PRL 103; PRD 82, 88

# Excited light quark baryons

- A ‘super’-multiplet of **hybrid baryons**

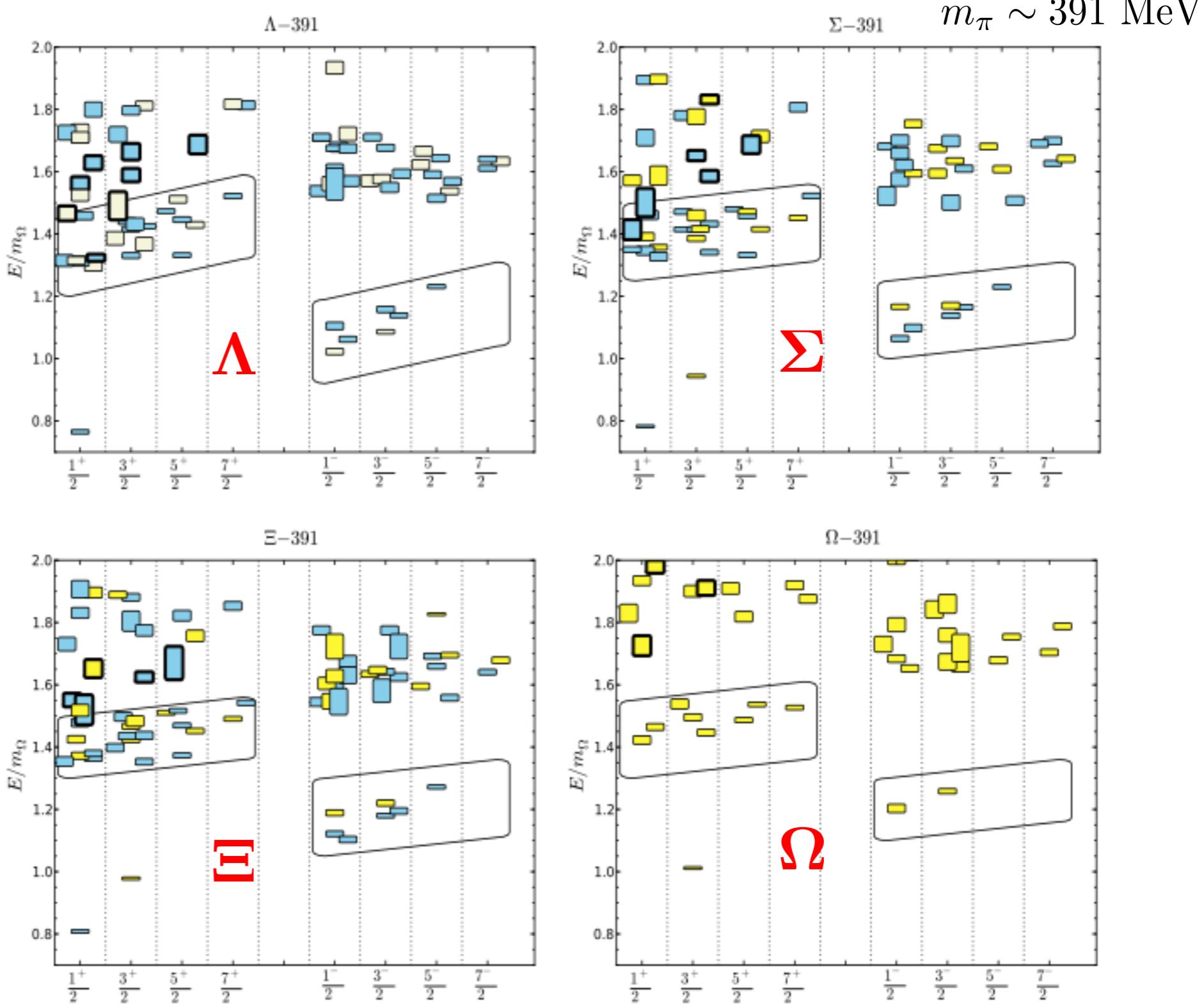


# Excited strange (and charm) quark baryons

Light quarks - SU(3)  
flavor broken

Full non-relativistic  
quark model counting

Some mixing of SU(3)  
flavor irreps

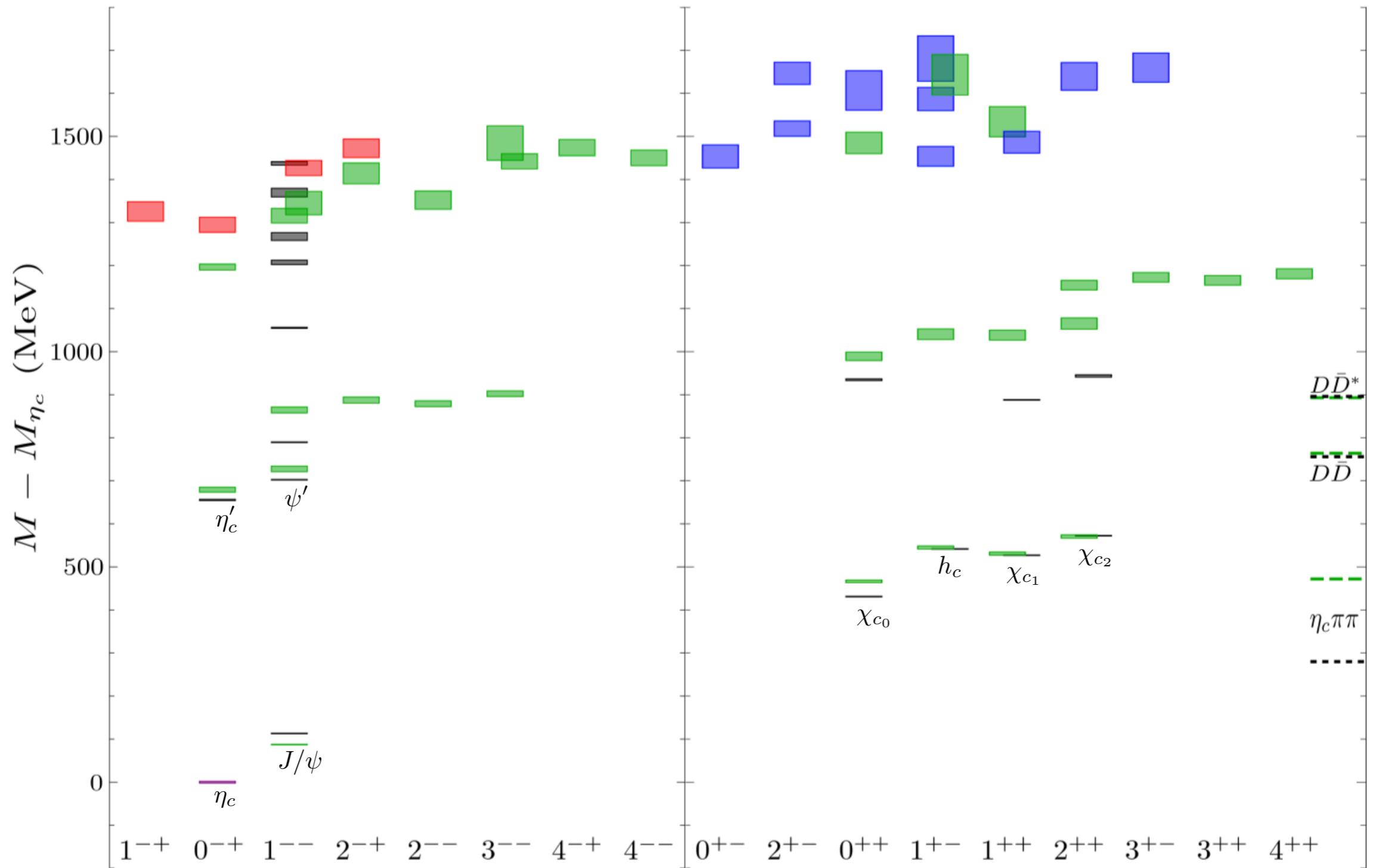


PRD87 054506 (2013)  
PRD90 074504 (2014)  
PRD91 054502 (2015)

# Charmonium spectrum

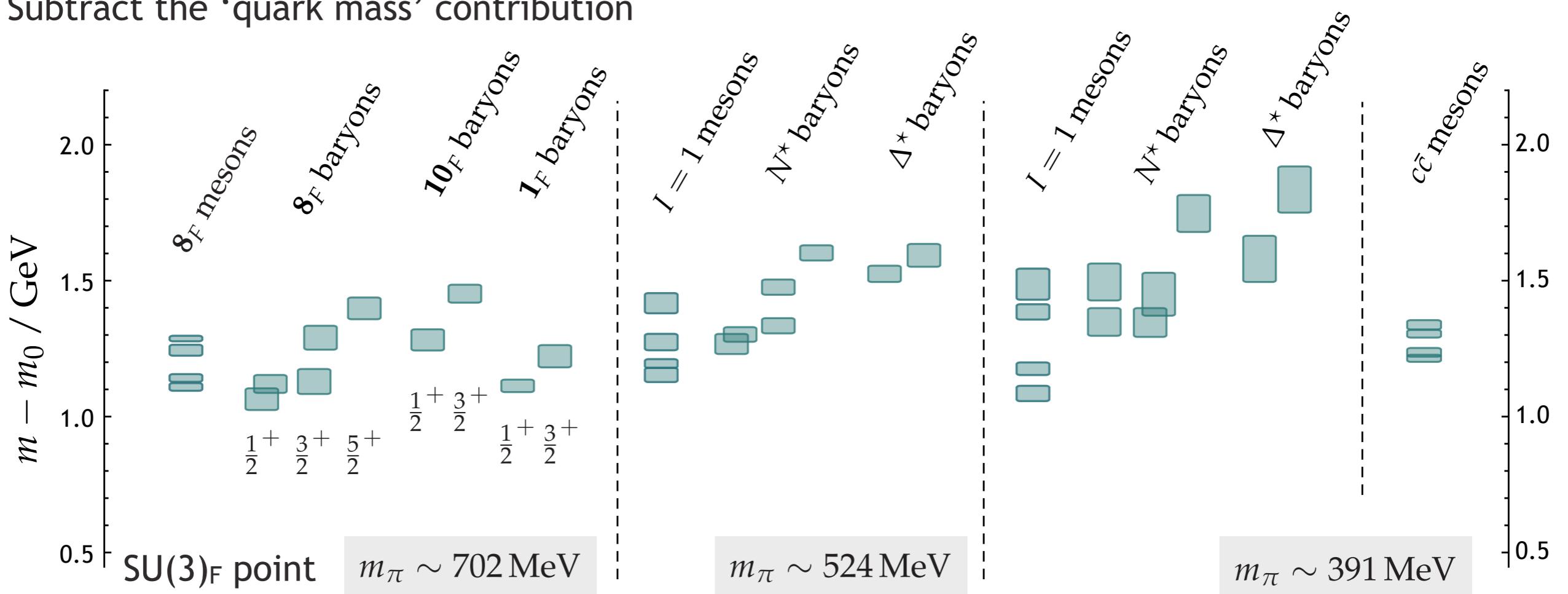
- A ‘super’-multiplet of **hybrid mesons**

Cheung, O’Hara, Tims, Moir, Peardon, Ryan, Thomas (2016)



# Chromo-magnetic excitation

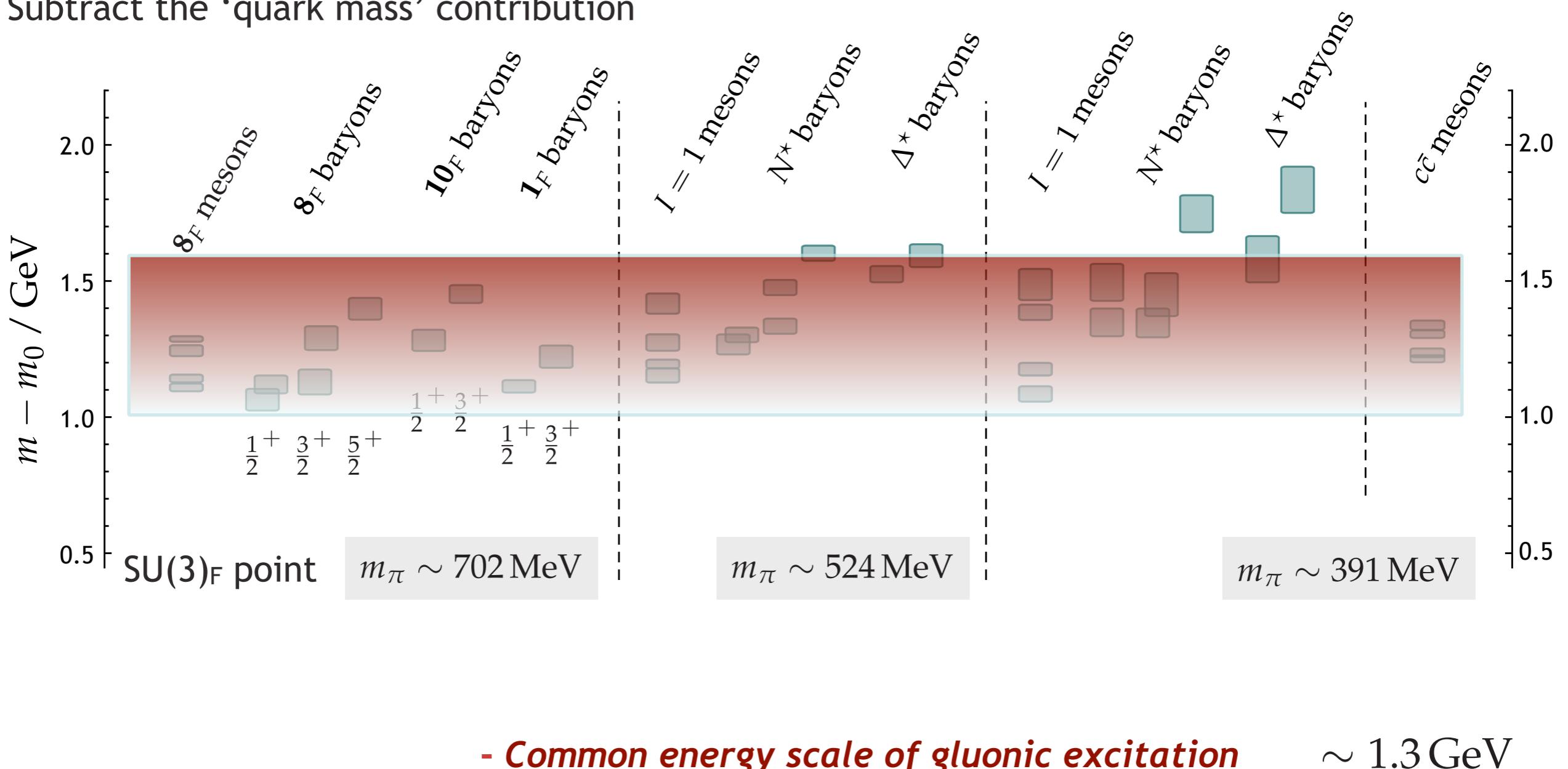
- Subtract the ‘quark mass’ contribution



HADRON SPECTRUM: PRD83 (2011); PRD88 (2013)

# Chromo-magnetic excitation

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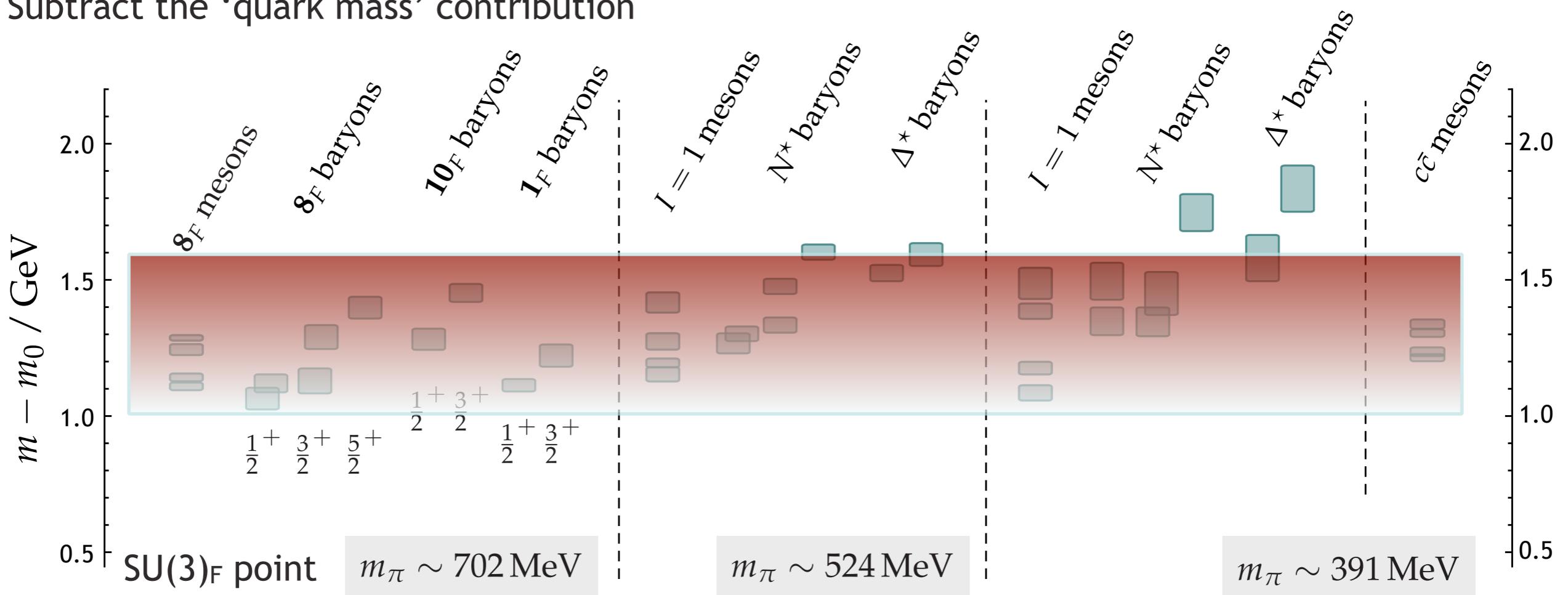
- Common energy scale of gluonic excitation

$\sim 1.3 \text{ GeV}$

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# Chromo-magnetic excitation

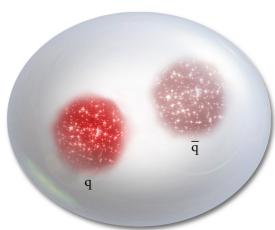
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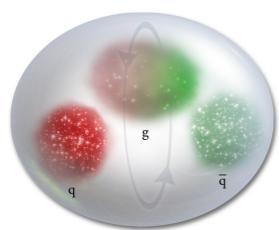
Pattern of states suggest  
gluonic excitations

- *Common energy scale of gluonic excitation*

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Conventional Meson

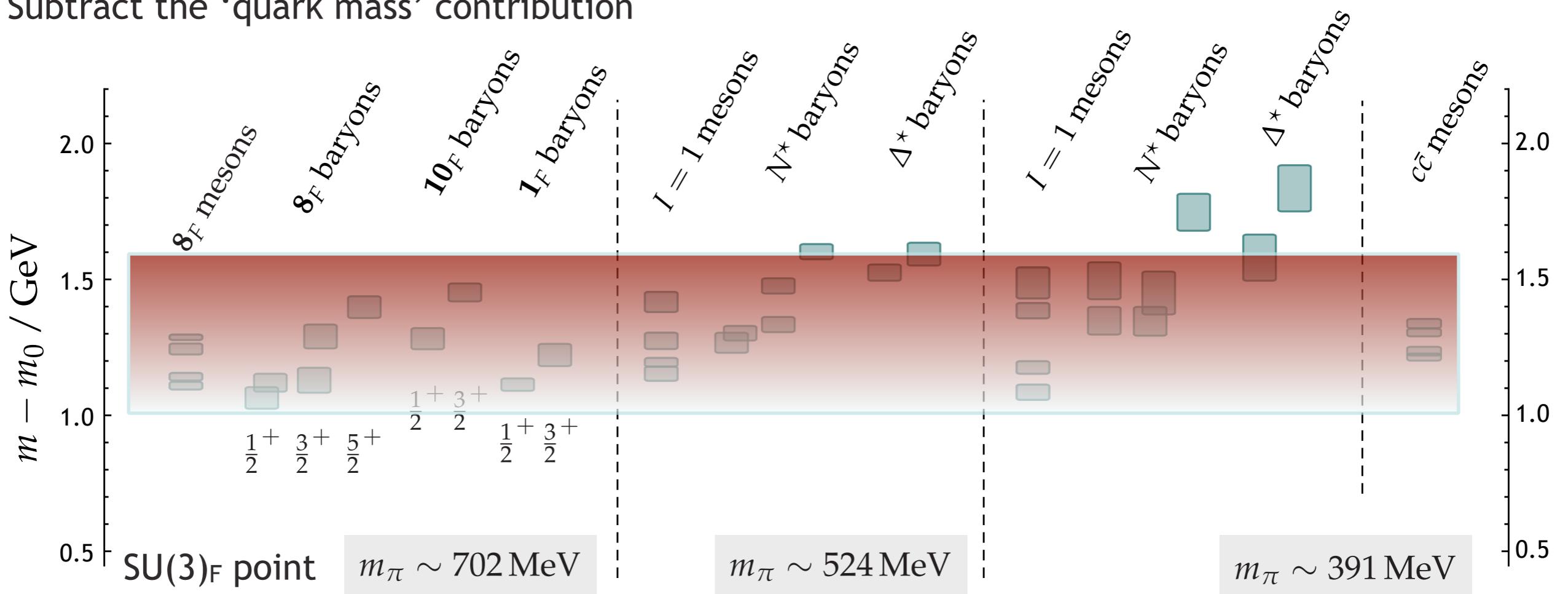


Hybrid Meson

HADRON SPECTRUM: PRD83 (2011); PRD88 (2013)

# Chromo-magnetic excitation

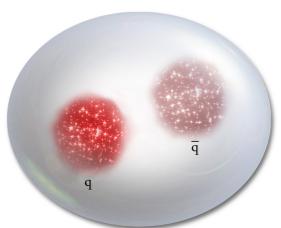
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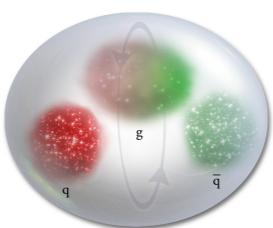
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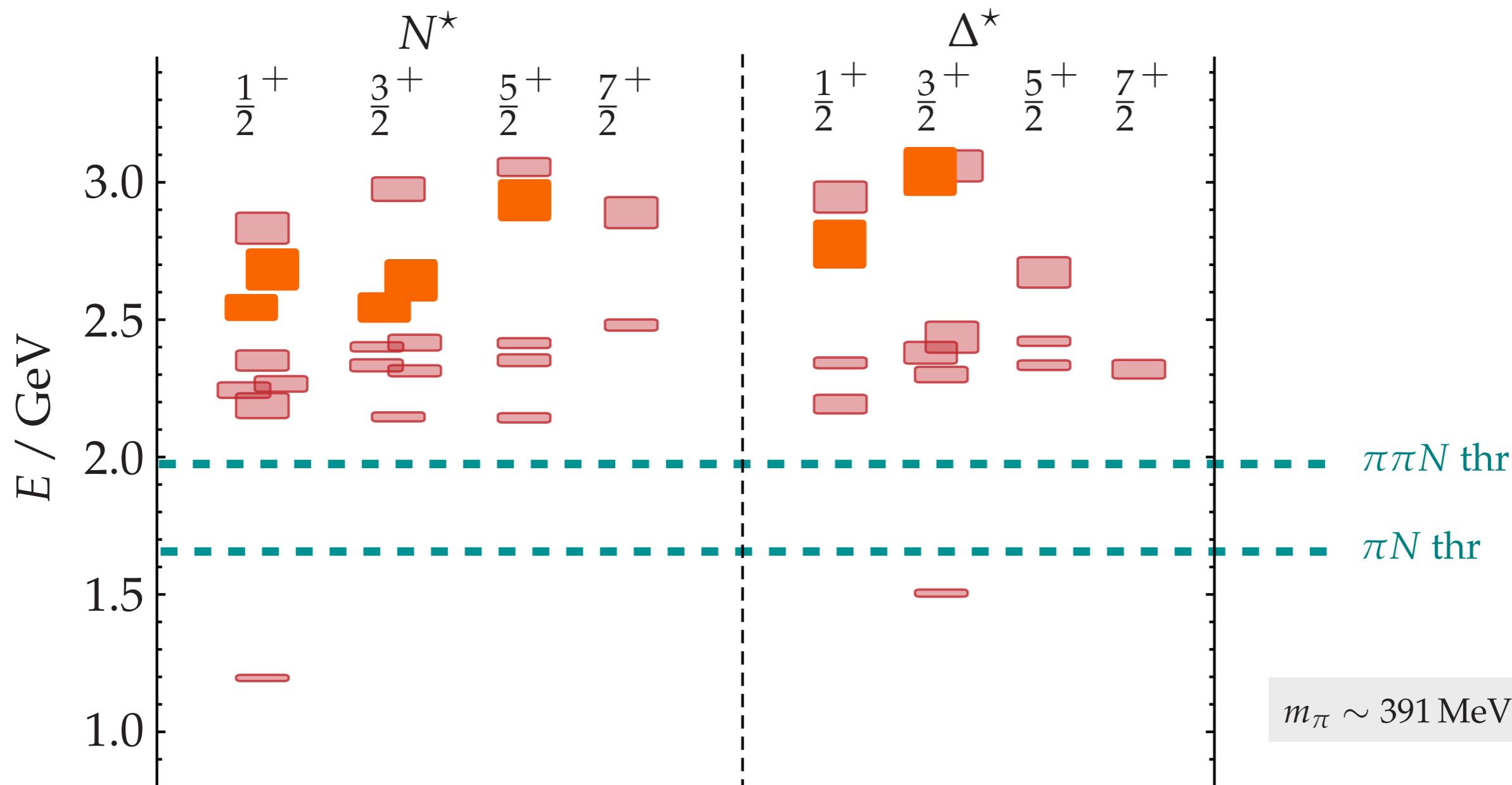
Hybrid Meson

→ Need to know decay modes and rates to compare to expt.

HADRON SPECTRUM: PRD83 (2011); PRD88 (2013)

# Excited states are resonances

- Initial determination of spectrum with only  $qqq$  style operators  
→ missing scattering states

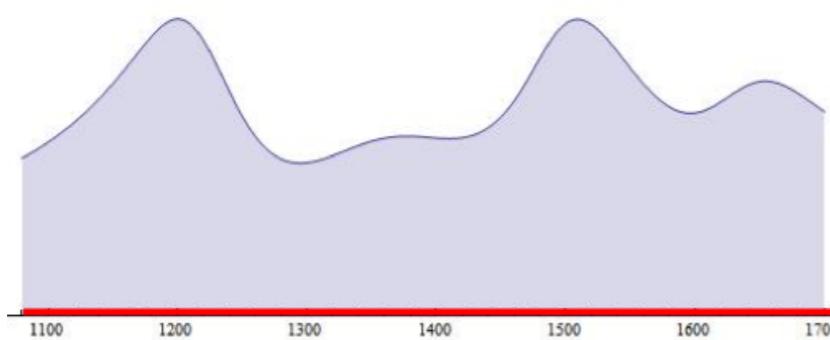


- Some initial results in S11 & P33 have appeared (Graz group)

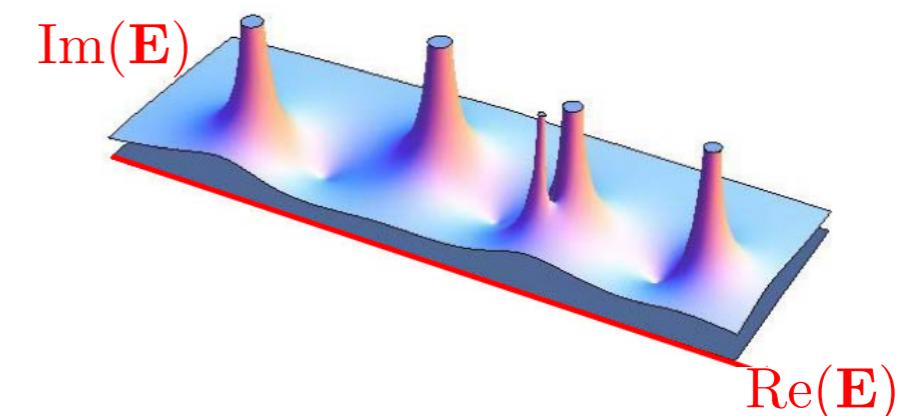
# Resonances

Manifest as behavior of real scattering amplitudes

- E.g.,  $\pi N \pi N$



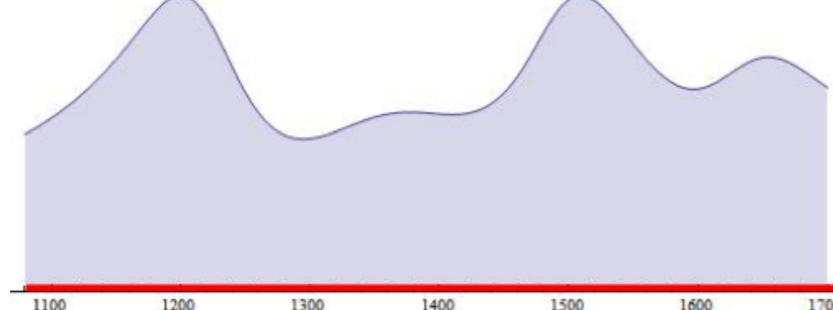
$E$  (MeV)



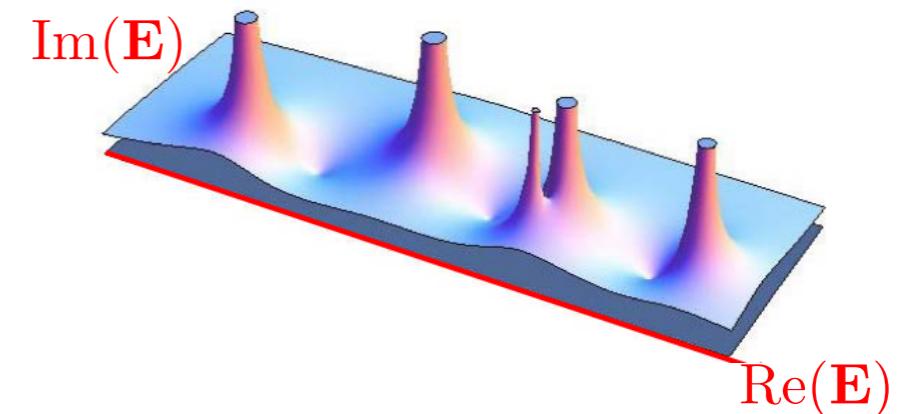
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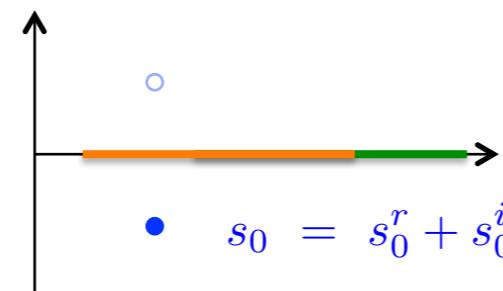


$E$  (MeV)



Formally defined as a pole in a partial-wave scattering amplitude

$$t_l(s) \sim \frac{R}{s_0 - s} + \dots$$



Different channels should have same pole location

Pole structure (location and residue) gives decay information

# QFT in a periodic cube

Lüscher (1986) : application to 3+1dim quantum field theories

*subsequent extensions for moving frames, coupled-channel systems - will come back to this*

**quantization condition:**

solutions,  $E_n$ , of

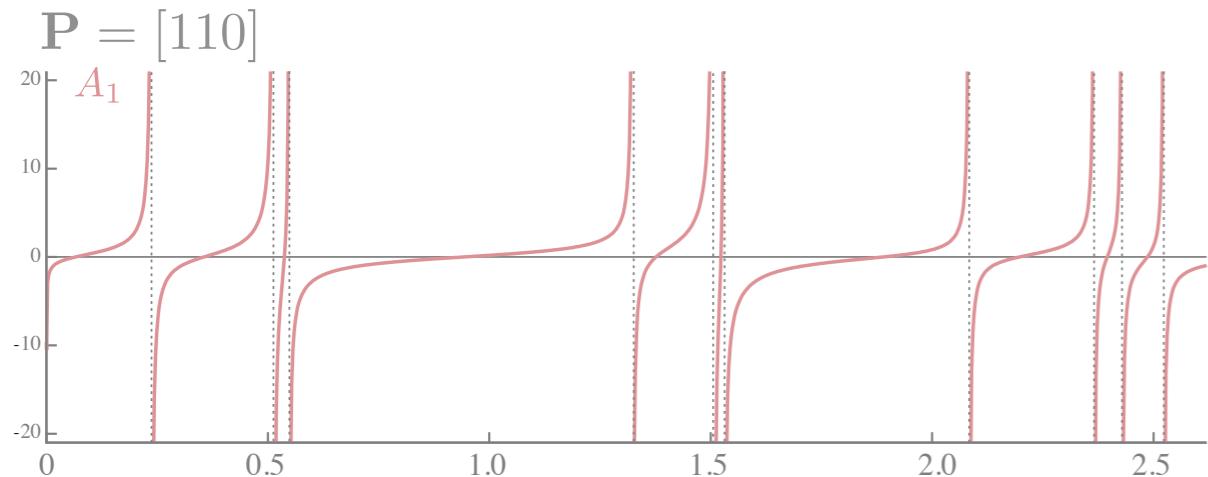
$$\det \left[ \mathbf{1} + i\boldsymbol{\rho} \cdot \mathbf{t} \cdot (\mathbf{1} + i\mathcal{M}) \right] = 0$$

$\boldsymbol{\rho}(E)$  phase-space

$\mathbf{t}(E)$  scattering matrix  $(\mathbf{S} = \mathbf{1} + 2i\sqrt{\boldsymbol{\rho}} \cdot \mathbf{t} \cdot \sqrt{\boldsymbol{\rho}})$

$\mathcal{M}(E, L)$  finite-volume function

technicalities: partial-wave basis not ‘diagonal’ in a cube



# QFT in a periodic cube

Lüscher (1986) : application to 3+1dim quantum field theories

*subsequent extensions for moving frames, coupled-channel systems - will come back to this*

**quantization condition:**

solutions,  $E_n$ , of

$$\det \left[ \mathbf{1} + i\boldsymbol{\rho} \cdot \mathbf{t} \cdot (\mathbf{1} + i\mathcal{M}) \right] = 0$$

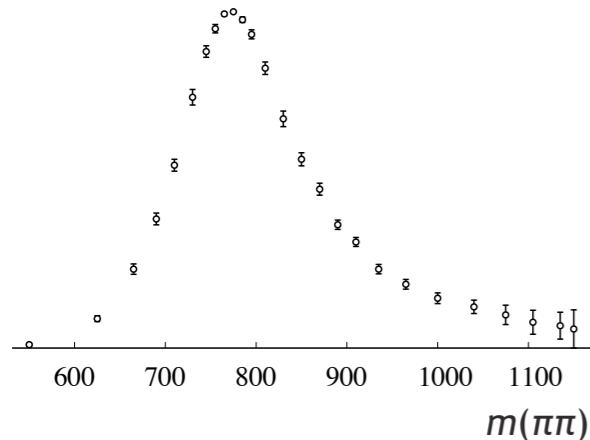
simplest case – elastic scattering of a single partial wave:

$$t^{(\ell)} = \frac{1}{\rho} e^{i\delta_\ell} \sin \delta_\ell \quad \Rightarrow \quad \cot \delta(E) = \mathcal{M}(E, L)$$

$E_n$  value maps to  $\delta(E_n)$

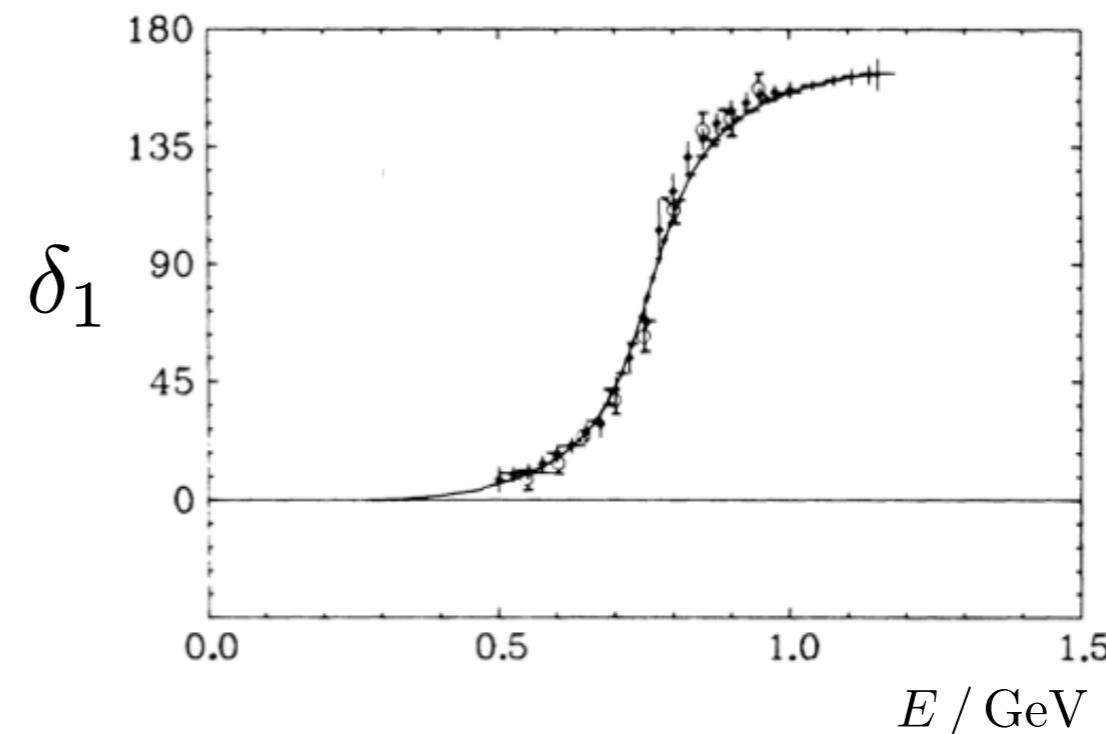
need to compute the spectrum ...

# An elastic resonance – the $\rho$ in $\pi\pi$



canonical resonance ‘bump’  
described by a rapidly rising phase-shift

scattering phase-shift



PHYSICAL REVIEW D

VOLUME 7, NUMBER 5

1 MARCH 1973

$\pi\pi$  Partial-Wave Analysis from Reactions  $\pi^+p \rightarrow \pi^+\pi^-\Delta^{++}$  and  $\pi^+p \rightarrow K^+K^-\Delta^{++}$  at 7.1 GeV/c†

S. D. Protopopescu,\* M. Alston-Garnjost, A. Barbaro-Galtieri, S. M. Flatté,‡  
J. H. Friedman,§ T. A. Lasinski, G. R. Lynch, M. S. Rabin,|| and F. T. Solmitz  
*Lawrence Berkeley Laboratory, University of California, Berkeley, California 94720*  
(Received 25 September 1972)

# Lattice QCD spectrum

Variational analysis of a matrix of correlation functions

$$C_{ij}(t) = \langle 0 | \mathcal{O}_i(t) \mathcal{O}_j(0) | 0 \rangle$$
$$= \sum_{\mathfrak{n}} e^{-E_{\mathfrak{n}} t} \langle 0 | \mathcal{O}_i | \mathfrak{n} \rangle \langle \mathfrak{n} | \mathcal{O}_j | 0 \rangle$$

operator basis:

‘single-meson’	+	‘meson-meson’
$\bar{\psi} \Gamma \psi$		$\sum_{\hat{\mathbf{p}}_1, \hat{\mathbf{p}}_2} C(\mathbf{p}_1, \mathbf{p}_2; \mathbf{p}) M_1(\mathbf{p}_1) M_2(\mathbf{p}_2)$
( & tetraquark & ... )		maximum momentum guided by non-interacting energies
		$\mathbf{p} = \frac{2\pi}{L} [n_x, n_y, n_z]$
		$\sqrt{m_1^2 + \mathbf{p}_1^2} + \sqrt{m_2^2 + \mathbf{p}_2^2}$

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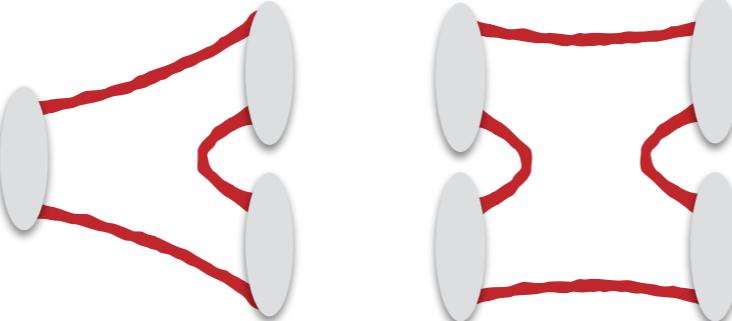
$$\bar{\psi} \Gamma \psi + \sum_{\hat{\mathbf{p}}_1, \hat{\mathbf{p}}_2} C(\mathbf{p}_1, \mathbf{p}_2; \mathbf{p}) M_1(\mathbf{p}_1) M_2(\mathbf{p}_2)$$

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now need to evaluate  
diagrams like



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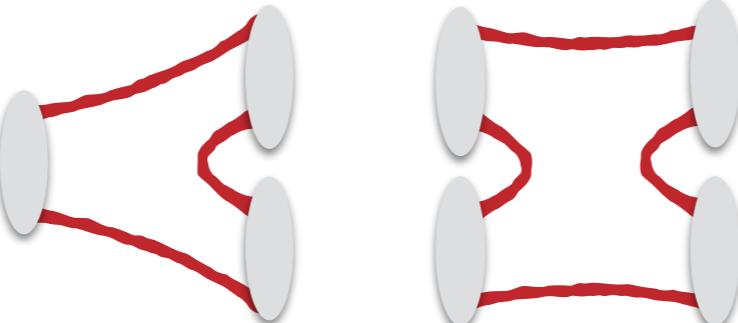
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**Distillation handles  
quark annihilation lines**

**Linear ops from KNLs+GPUs**



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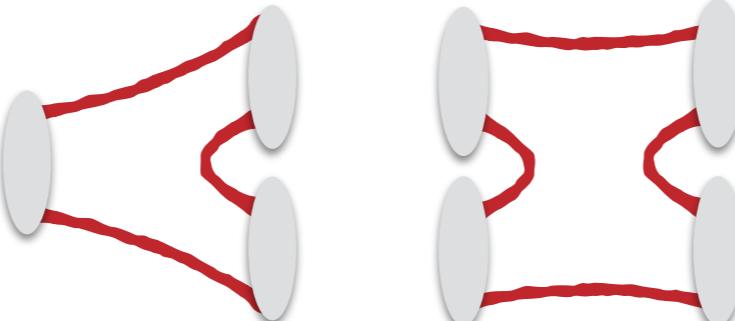
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now need to evaluate diagrams like



Distillation handles  
quark annihilation lines



Linear ops from KNLs+GPUs



‘Don’t underestimate the power...’

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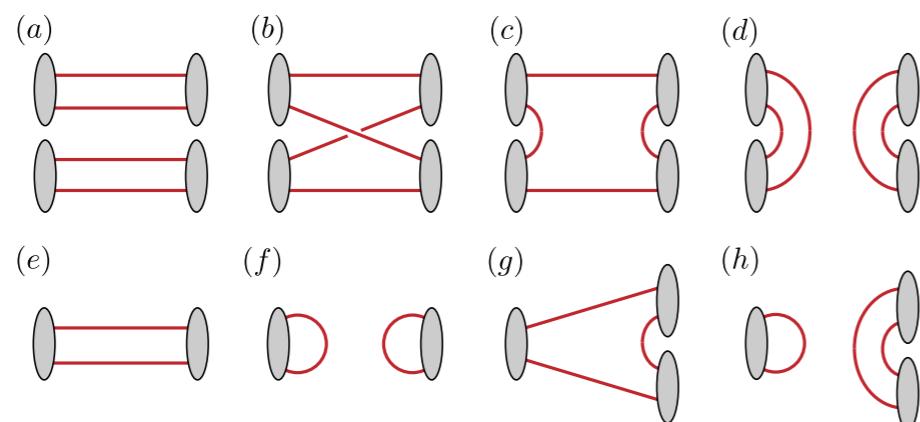
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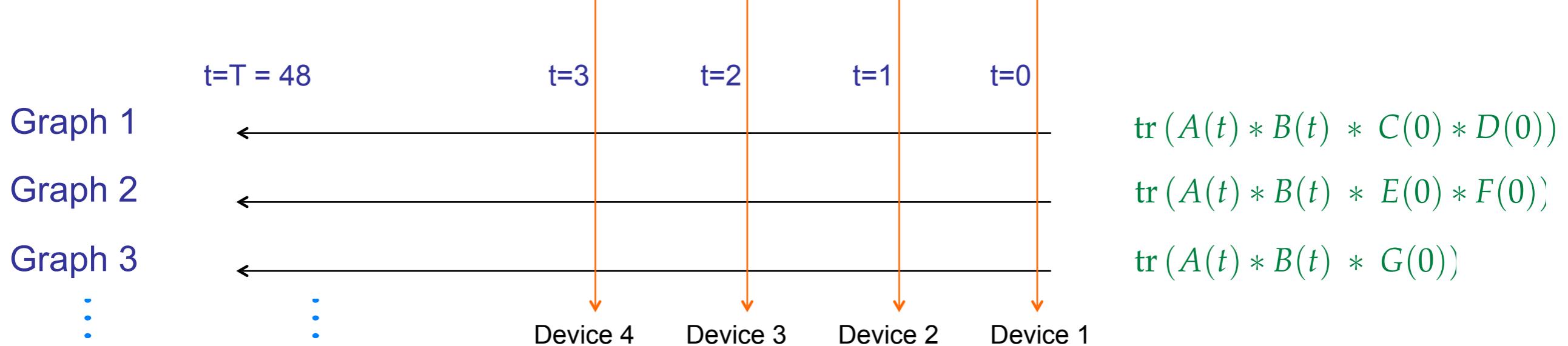
$$\sqrt{m_1^2 + \mathbf{p}_1^2} + \sqrt{m_2^2 + \mathbf{p}_2^2}$$

Can be lots of Wick contractions, and momentum projections

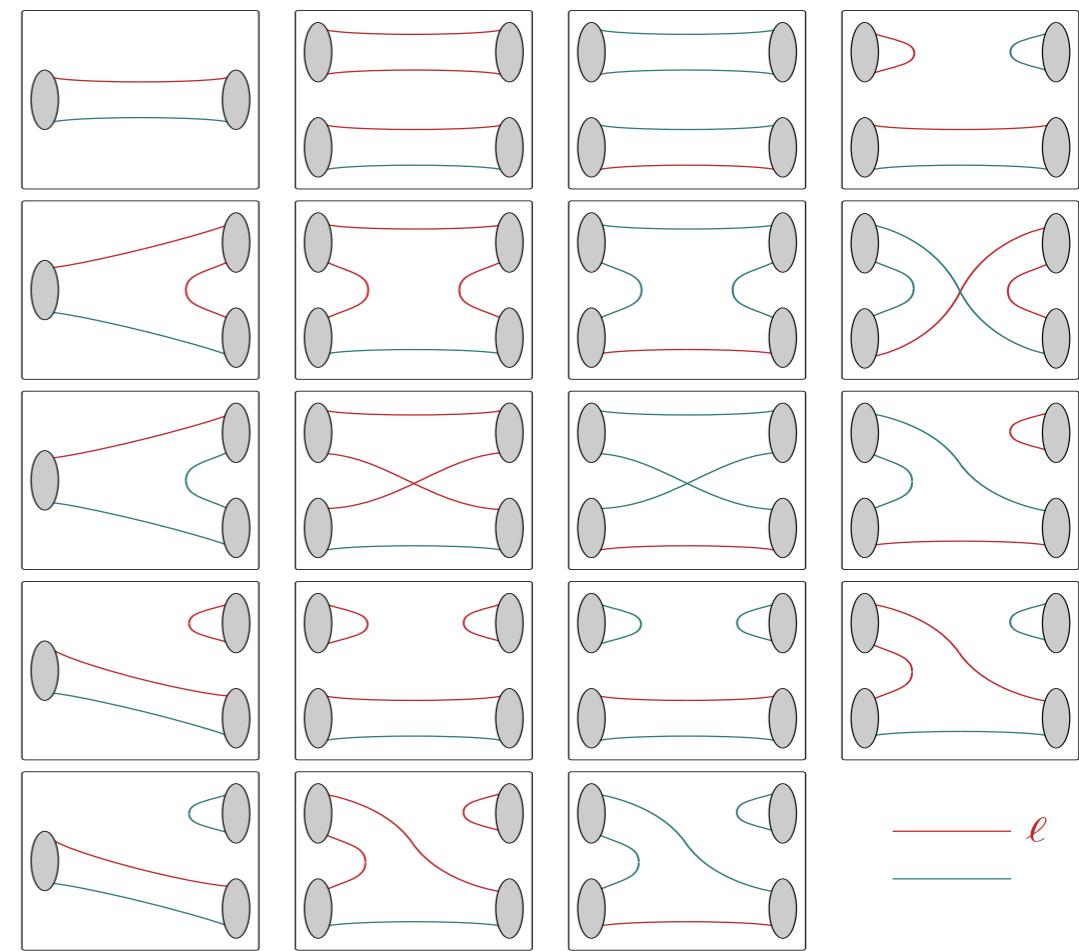
**Worst case: rest-frame  $\rightarrow p=100 \rightarrow 6x, p=110 \rightarrow 12x, p=111 \rightarrow 8x$**



# Optimize order of operations

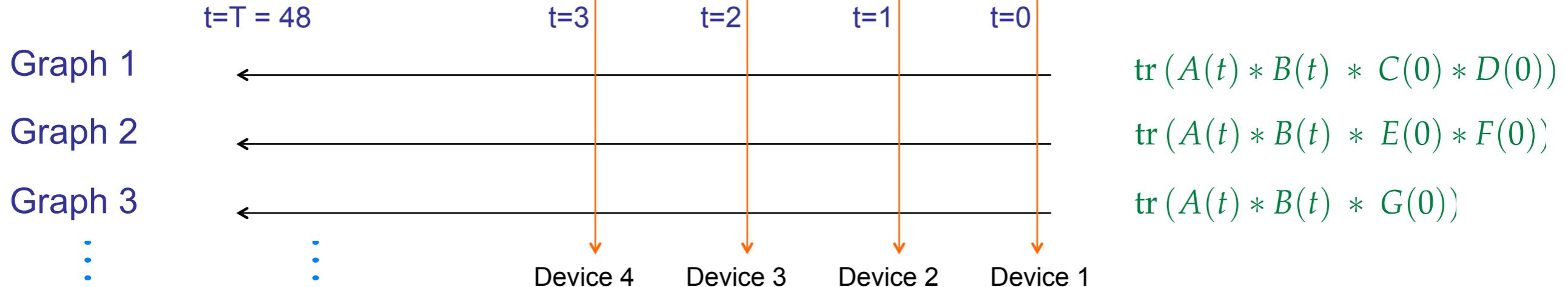


- Traverse graphs along a t-slice
  - 10,000's of graphs
- 3-particles and more...
- Common sub-expression elimination
- For fixed t-slice - 100's vertices



I=1/2 K\*π arXiv:1406.4158

# Optimize order of operations

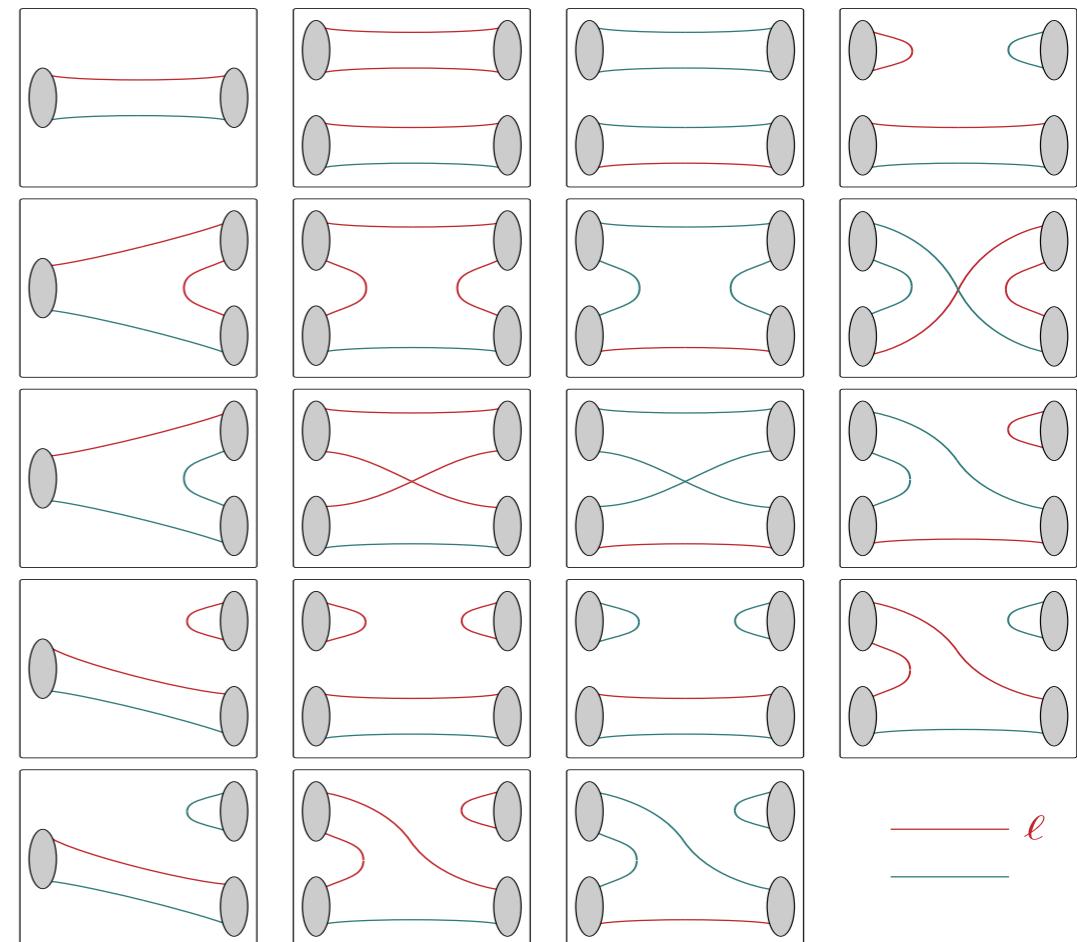


- Traverse graphs along a t-slice
  - 10,000's of graphs
- 3-particles and more...
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Solutions:

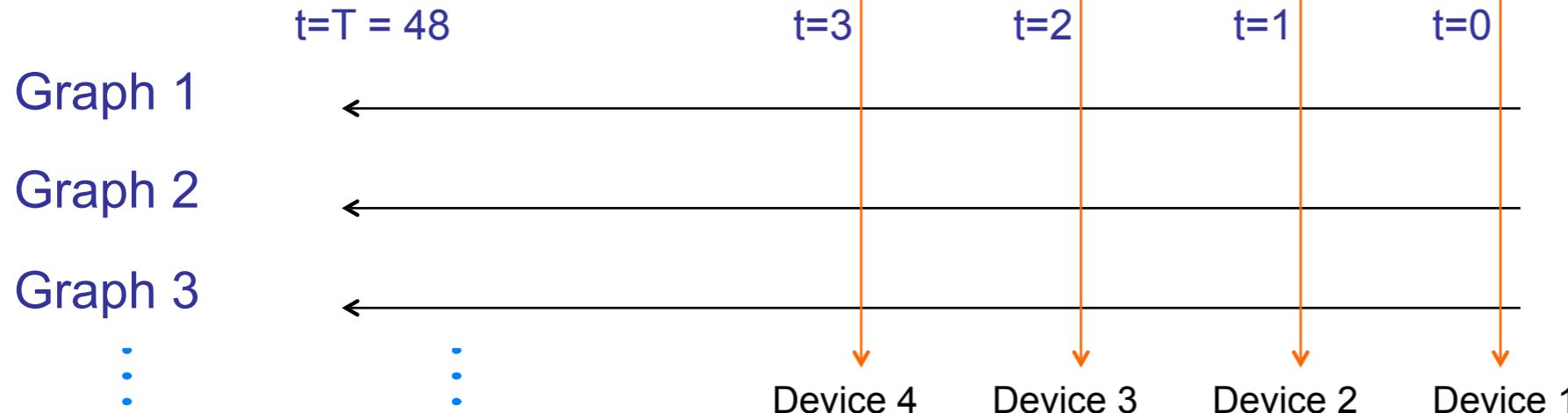
Equivalent graph identification

Dijkstra's method - optimal walk through space of evaluation order

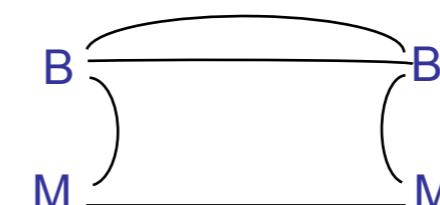


I=1/2 K\*π arXiv:1406.4158

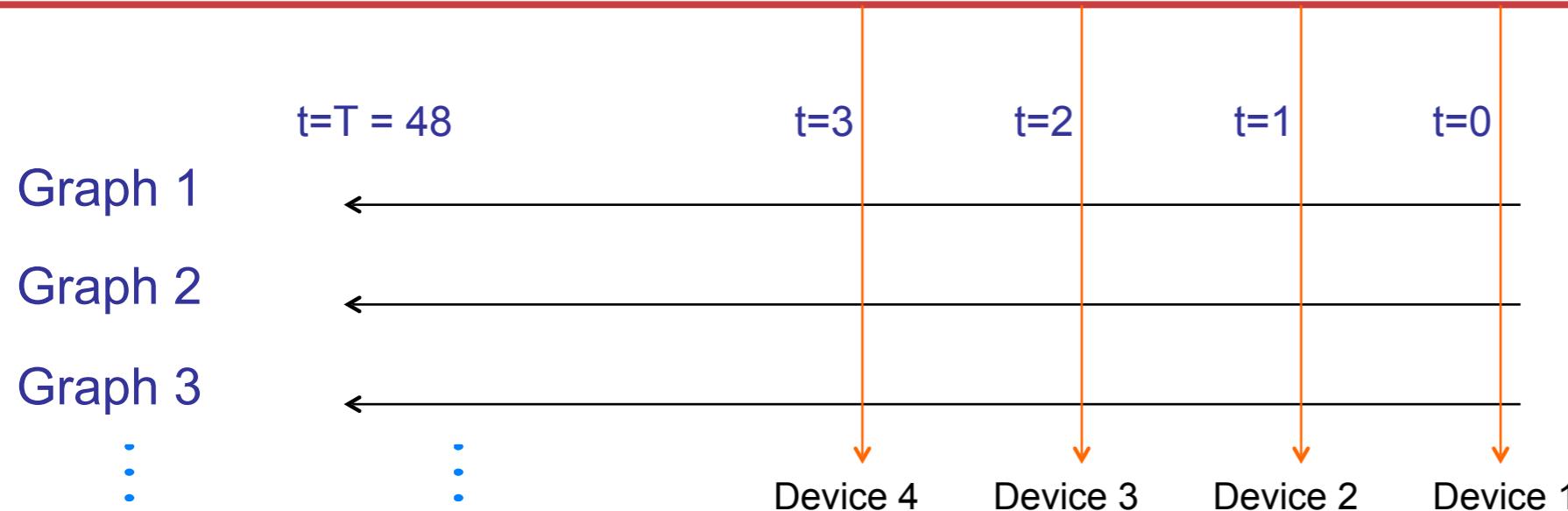
# Workflow is choreographed



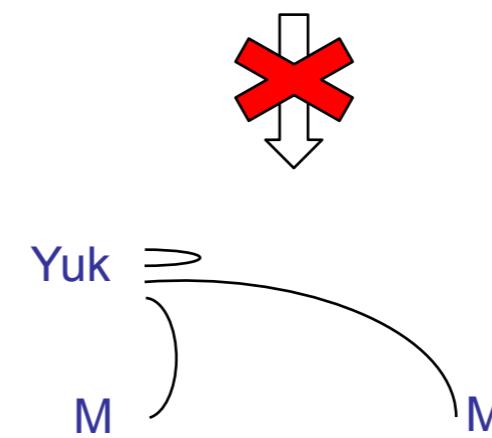
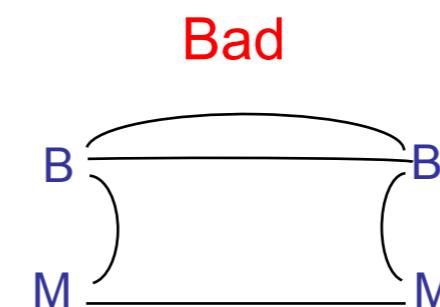
- Topology
  - Vertices can have different ordinalities
- Order of evaluation is important
  - Baryons  $\sim O(N^4)$ , Mesons  $\sim O(N^3)$



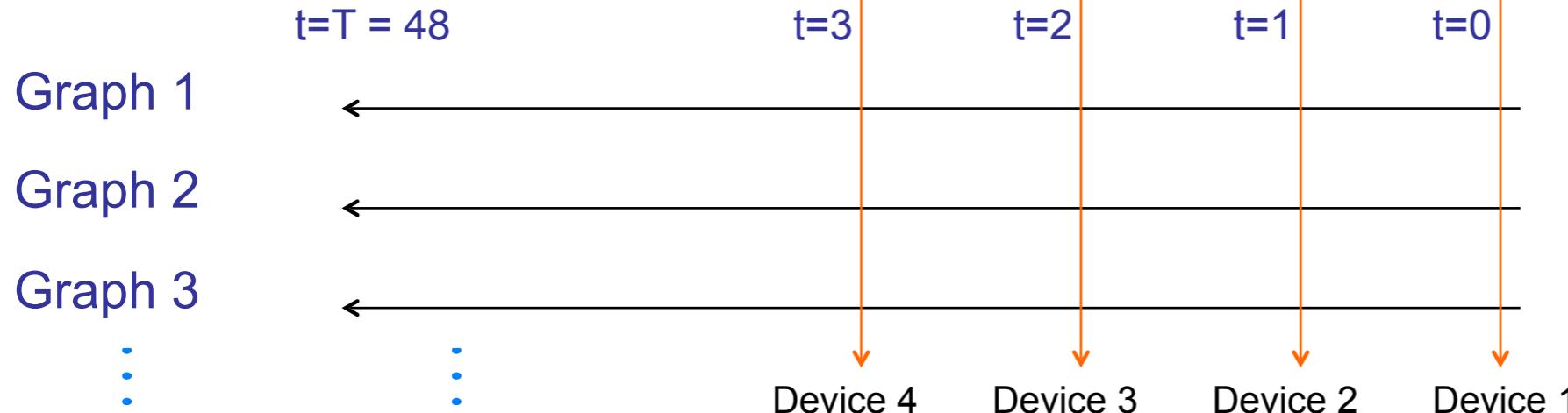
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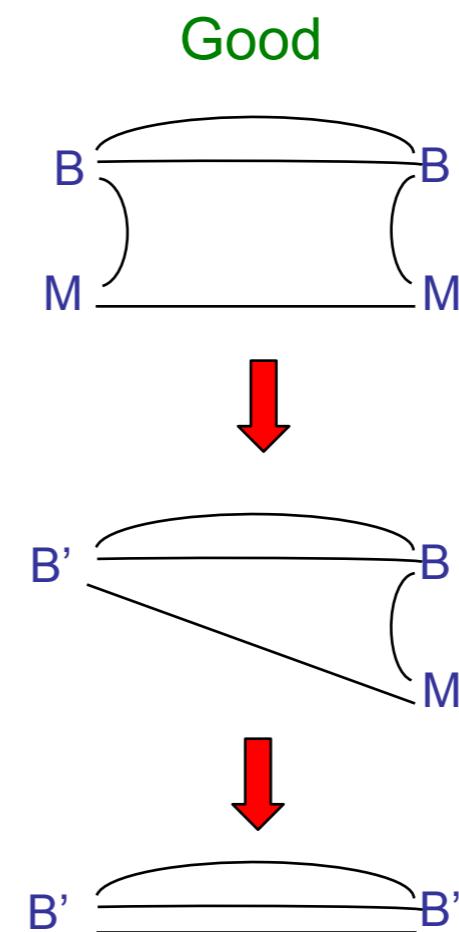
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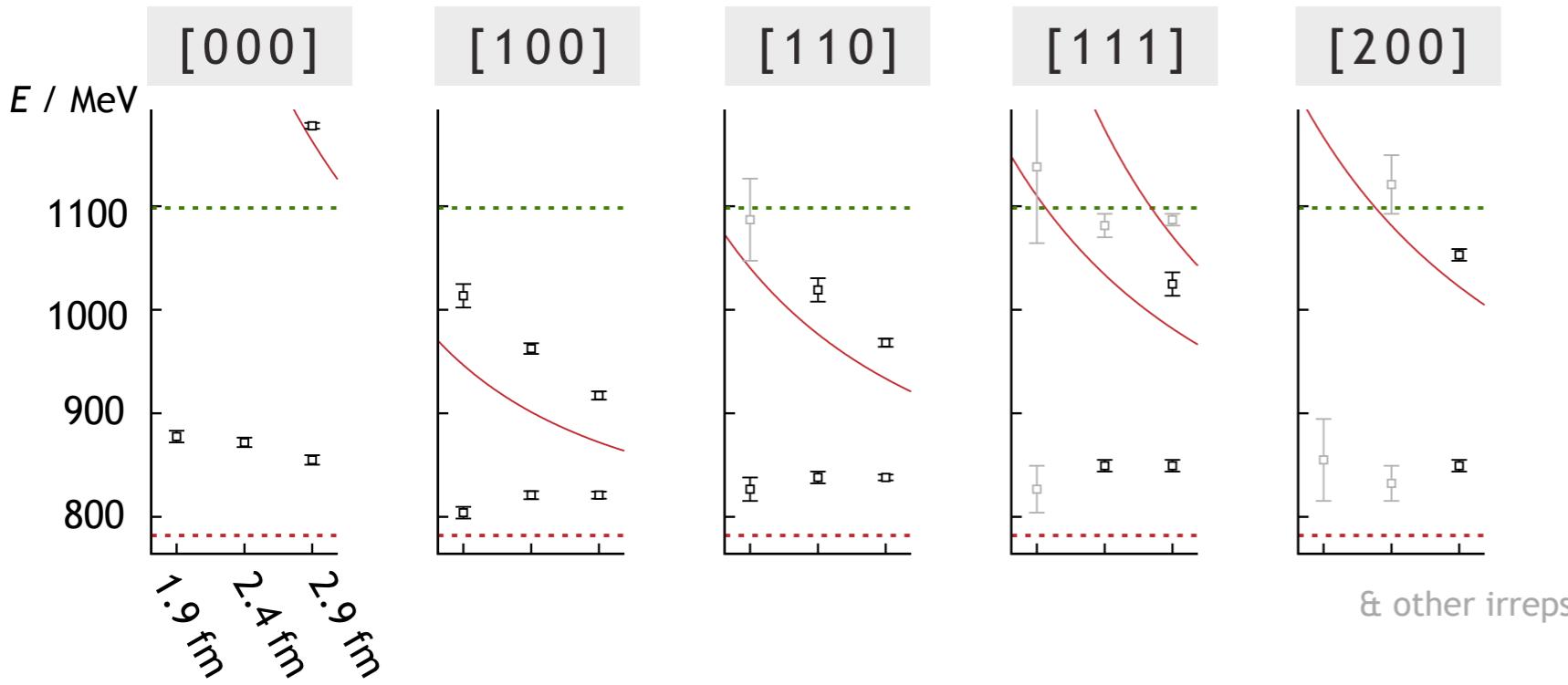
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- Order of evaluation is important
  - Baryons  $\sim O(N^4)$ , Mesons  $\sim O(N^3)$
  - Avoid creating larger ranked objects
- Lots of mat-muls
- Obvious application for accelerators



# An elastic resonance – the $\rho$ in $\pi\pi$ – lattice QCD

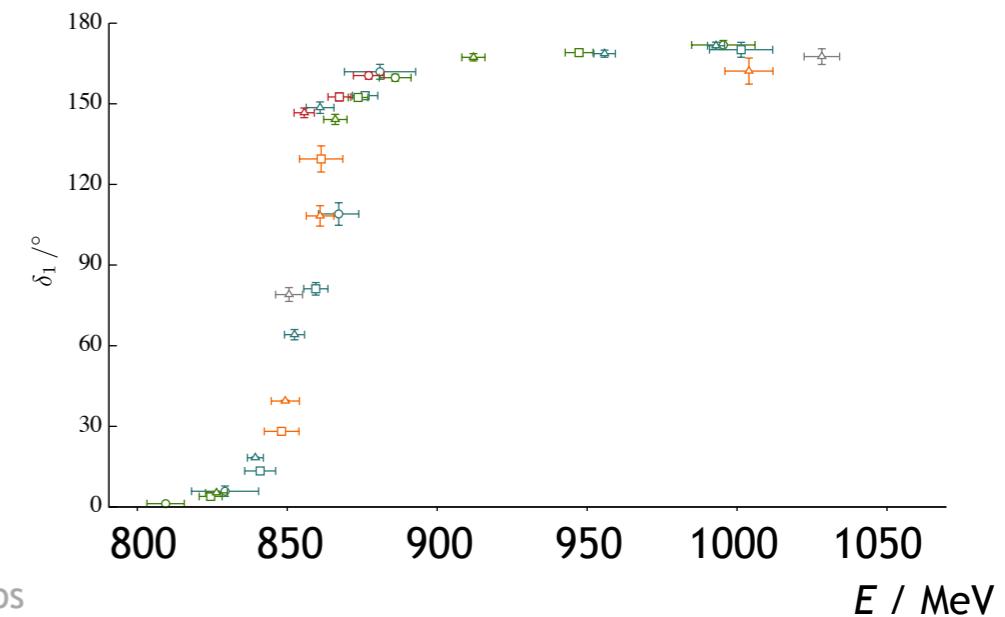
PRD87 034505 (2013)

$m_\pi \sim 391$  MeV



& other irreps

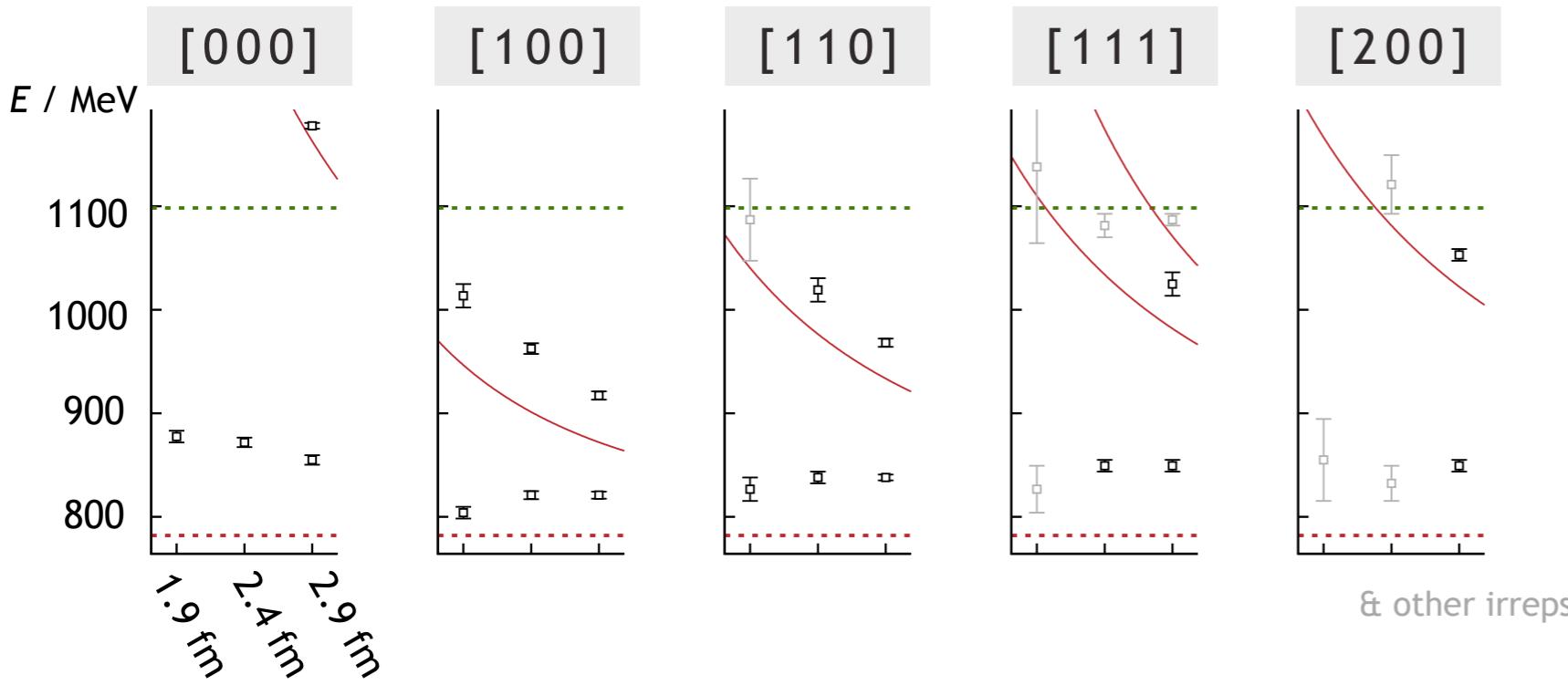
scattering phase-shift



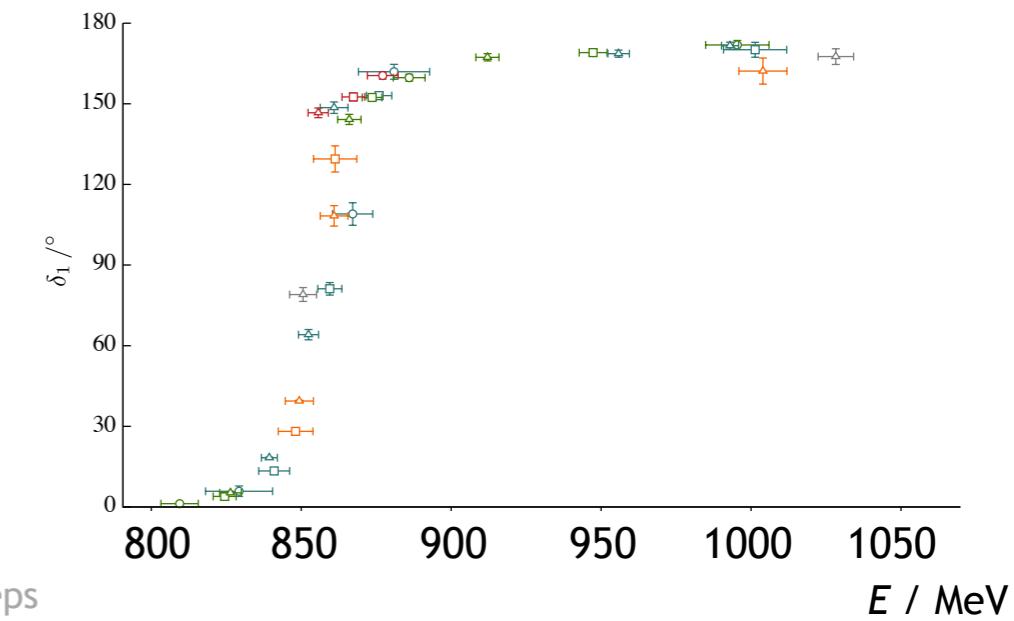
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PRD87 034505 (2013)

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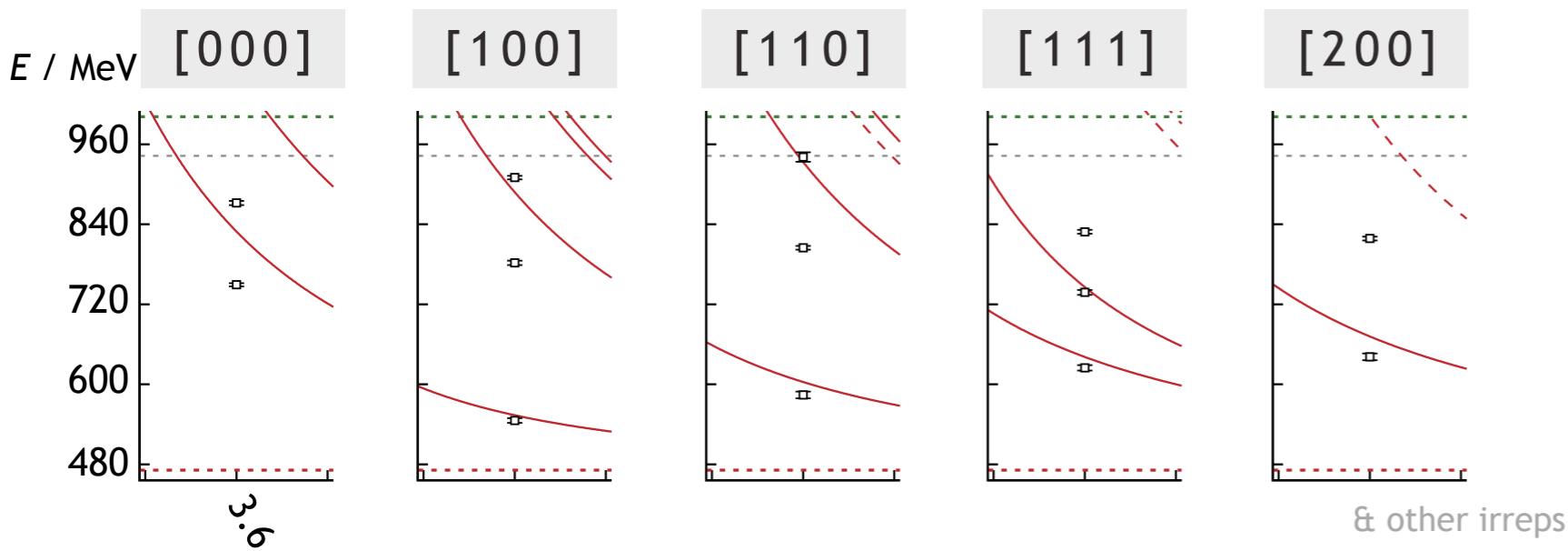


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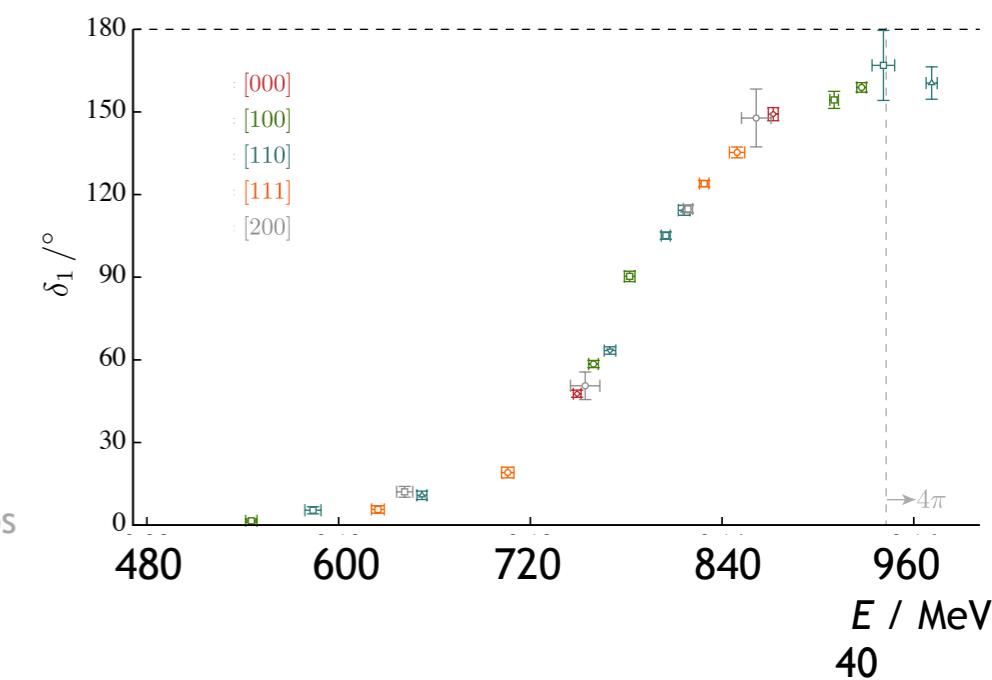


PRD92 094502 (2015)

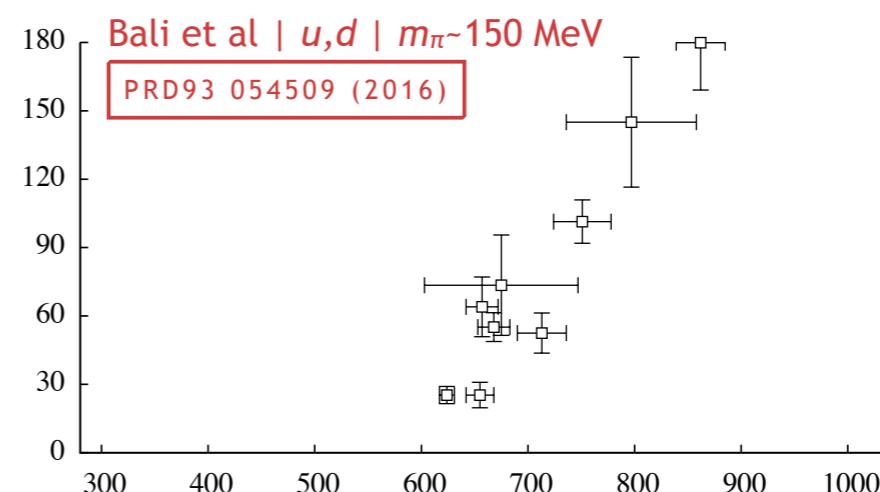
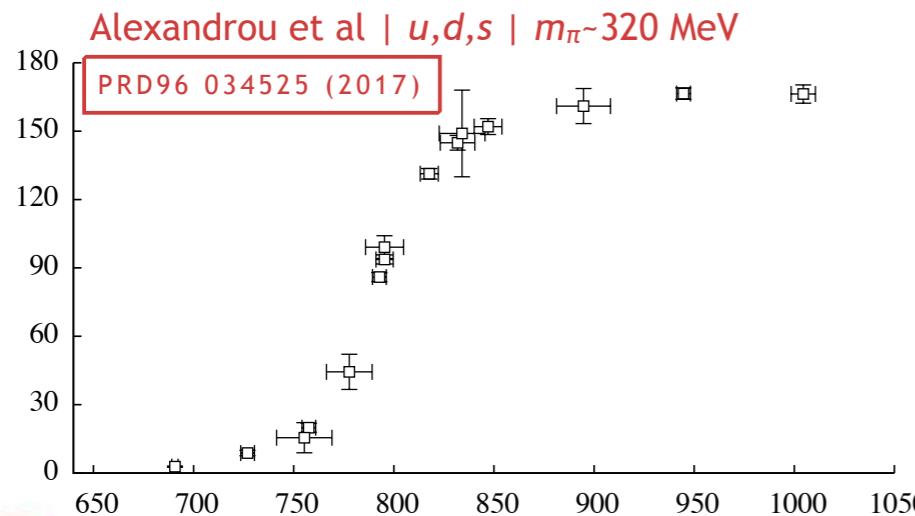
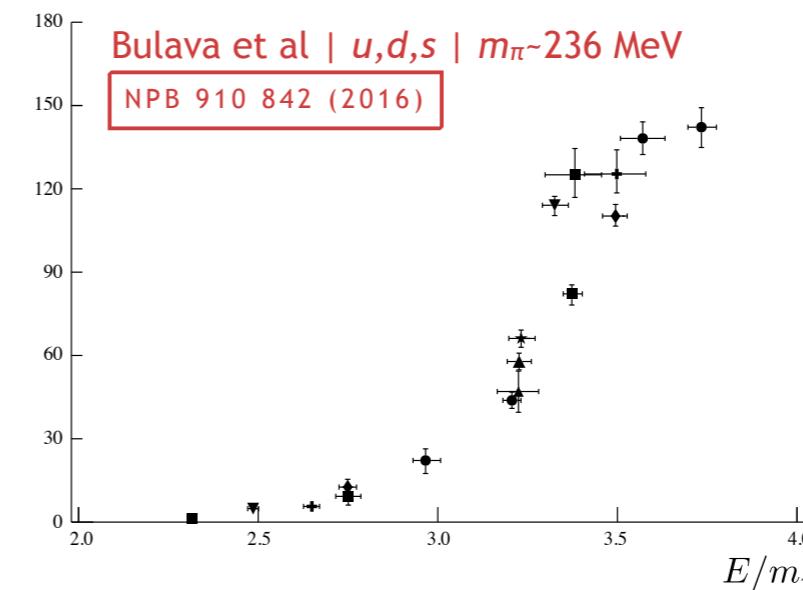
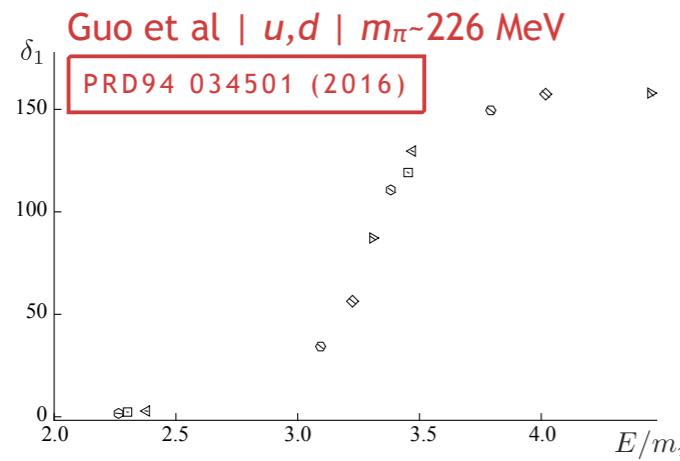
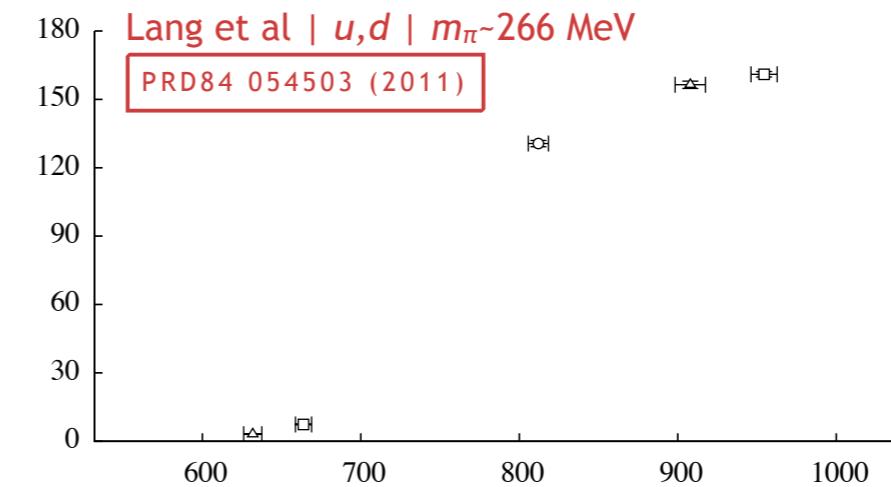
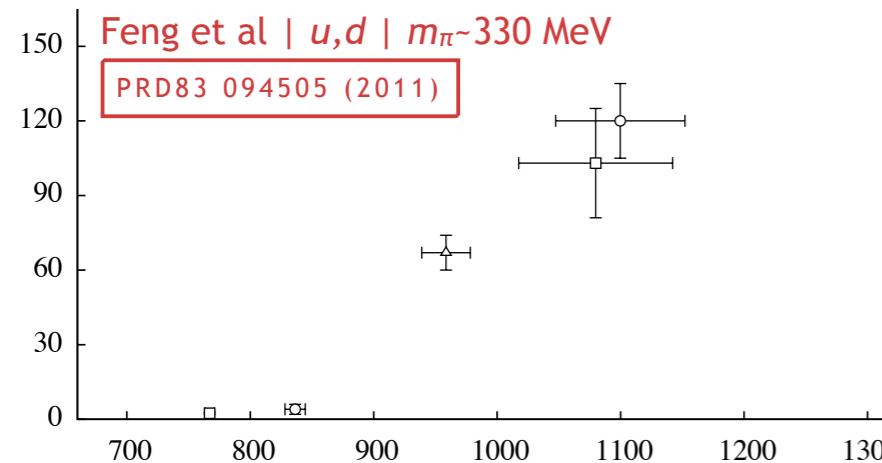
$m_\pi \sim 236$  MeV



scattering phase-shift



# An elastic resonance – the $\rho$ in $\pi\pi$ – lattice QCD



# Coupled-channel resonances

Most resonances decay into more than one final state

e.g. two-channel scattering described by a  $t$ -matrix

$$\mathbf{t}(E) = \begin{pmatrix} t_{11}(E) & t_{12}(E) \\ t_{21}(E) & t_{22}(E) \end{pmatrix}$$

Finite-volume spectrum as a function of scattering becomes more complicated

coupled-channel spectrum  
solutions  $E_n(L)$  of  
 $\det [\mathbf{1} + i\boldsymbol{\rho}(E) \cdot \mathbf{t}(E) \cdot (\mathbf{1} + i\mathcal{M}(E, L))] = 0$

No longer a one-to-one mapping from energy to scattering ...

# Coupled-channels

solutions,  $E_n$ , of

$$\det \left[ 1 + i\boldsymbol{\rho} \cdot \mathbf{t} \cdot (1 + i\mathcal{M}) \right] = 0$$

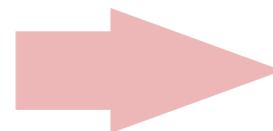
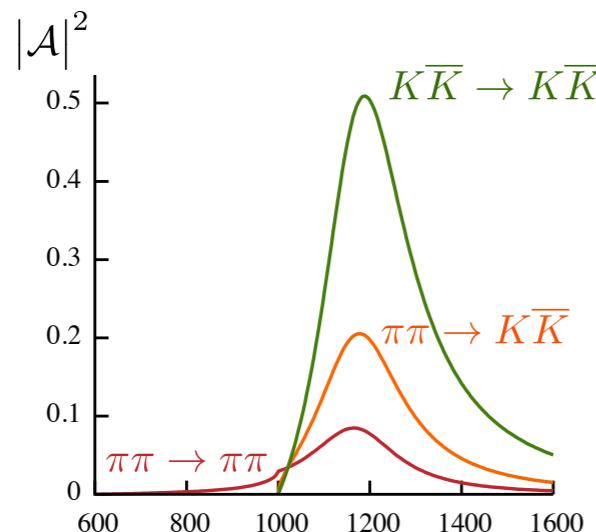
$\boldsymbol{\rho}(E)$  phase-space

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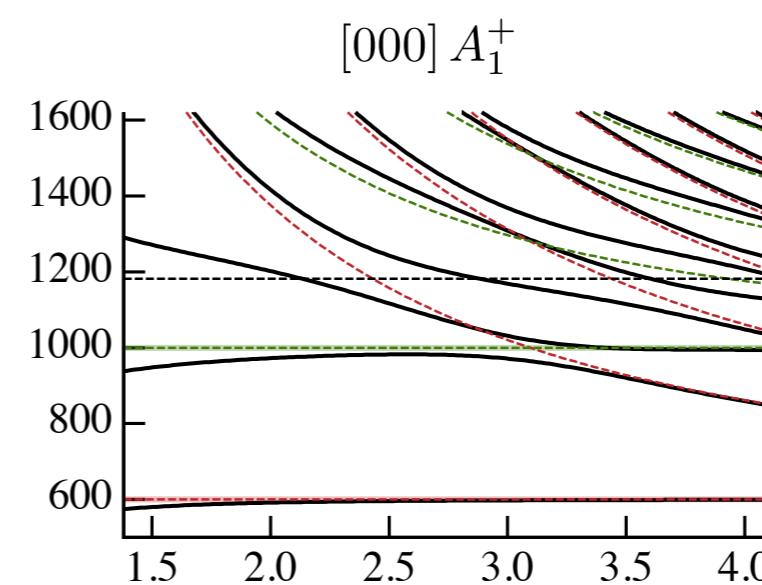
$\mathcal{M}(E, L)$  finite-volume function

$$\mathbf{t} = \begin{pmatrix} \pi\pi \rightarrow \pi\pi & \pi\pi \rightarrow K\bar{K} \\ K\bar{K} \rightarrow \pi\pi & K\bar{K} \rightarrow K\bar{K} \end{pmatrix}$$

for a known  $t(E)$



can predict  $E_n$



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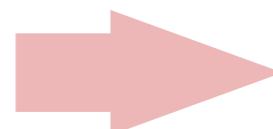
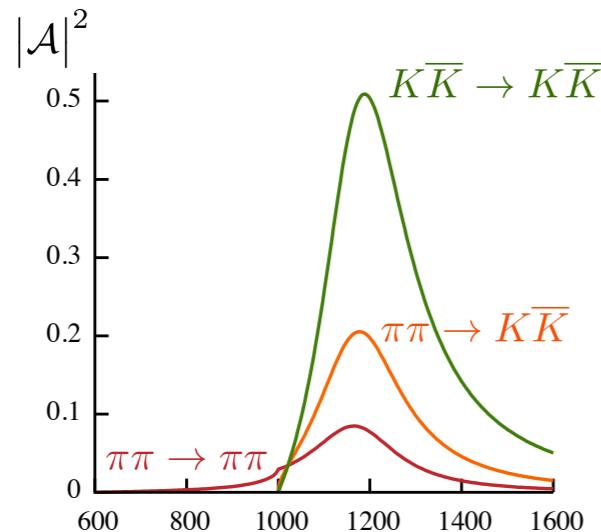
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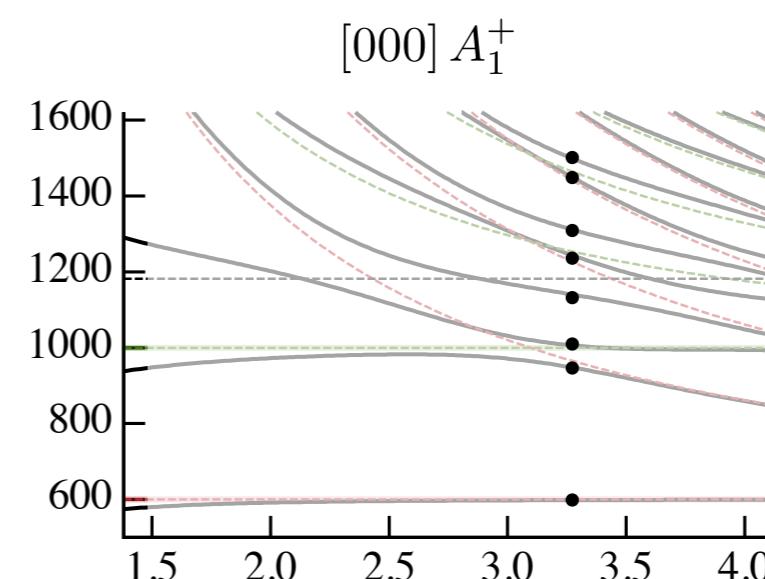
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for a known  $t(E)$



can predict  $E_n$



but how do we perform the inverse mapping ?

# Coupled-channels

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$\mathbf{t}(E)$  scattering matrix

$\mathcal{M}(E, L)$  finite-volume function

parameterize the energy dependence of the  $t$ -matrix

$K$ -matrix is a convenient approach (manifest unitarity)

$$\text{Im } t_{ij}(s) = -\delta_{ij} \rho_i(s) \Theta(s - s_{\text{thr.}}^{(i)})$$

$$\mathbf{t}^{-1}(E) = \mathbf{K}^{-1}(E) - i\boldsymbol{\rho}(E) \quad \text{where } K\text{-matrix is “any” real matrix}$$

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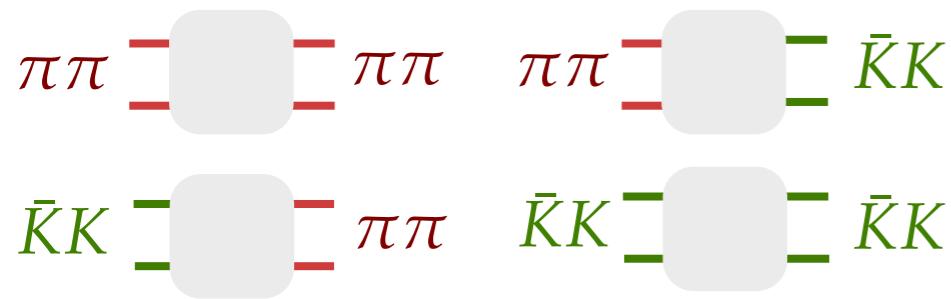
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Want pole mass and residues/couplings of t-matrix

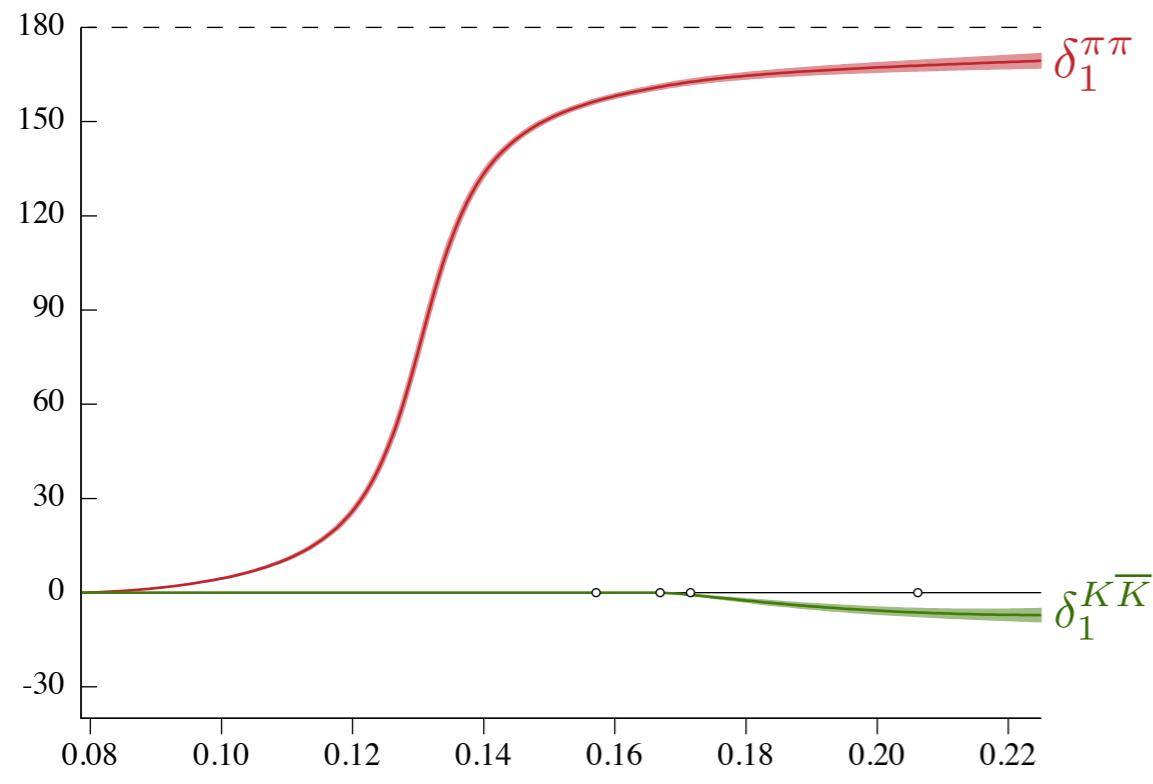
In recent years, progress towards establishing this approach

# $\rho$ resonance as a coupled channel system

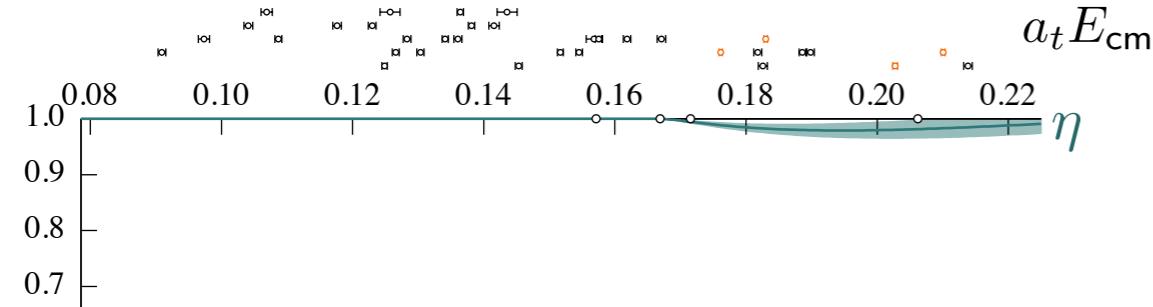


PRD87 034505 (2013)

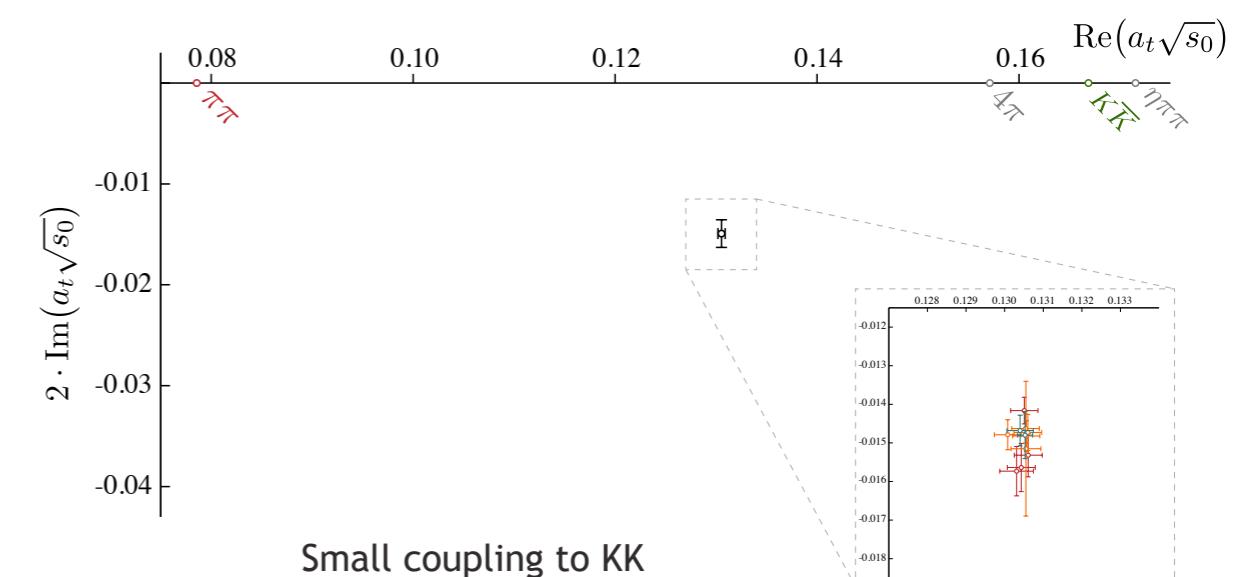
## Phase shifts & inelasticity



$m_\pi \sim 236 \text{ MeV}$

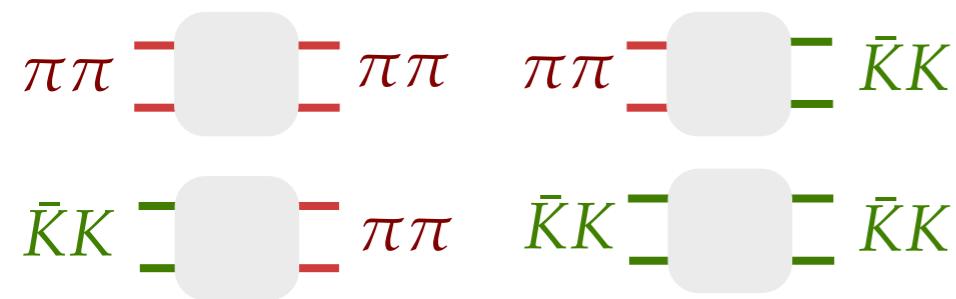


## t-matrix pole location



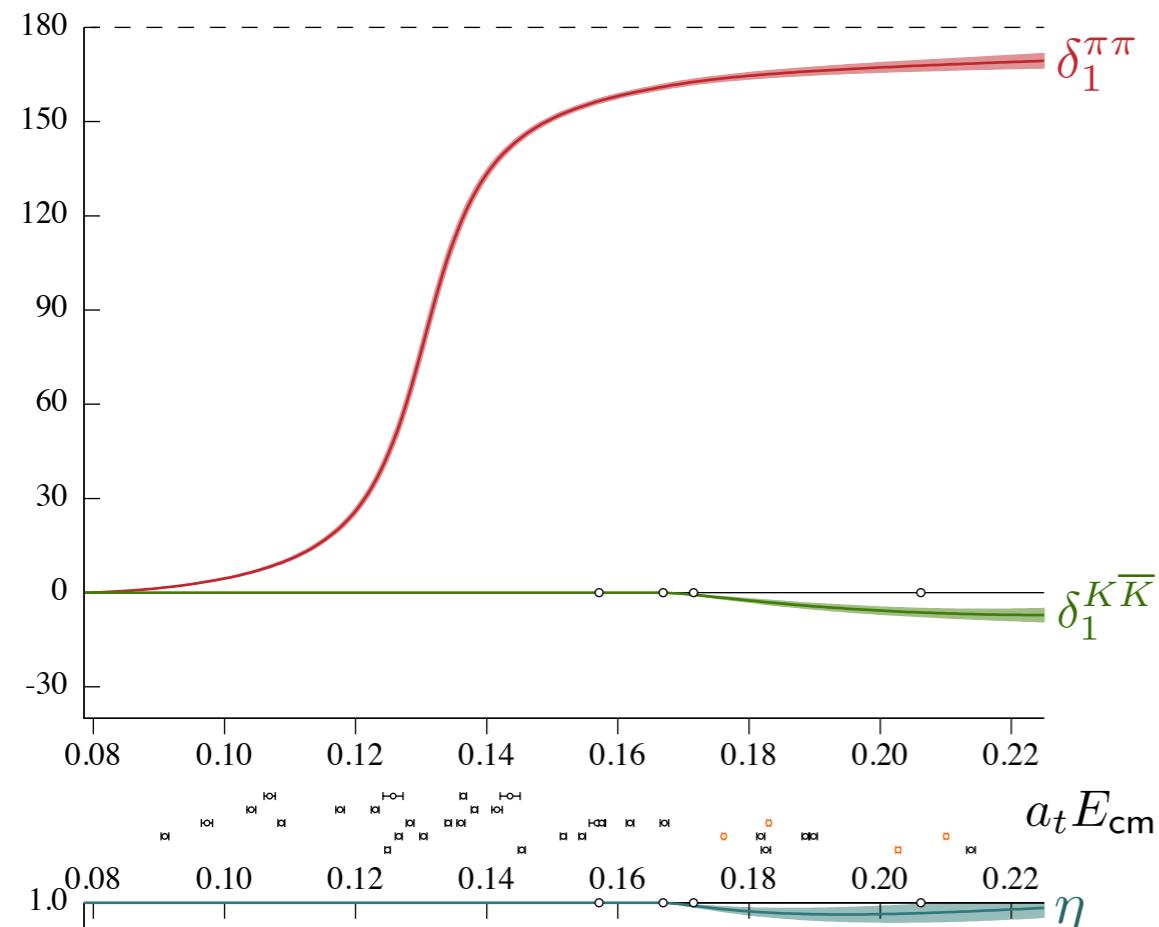
Small coupling to KK

# $\rho$ resonance as a coupled channel system

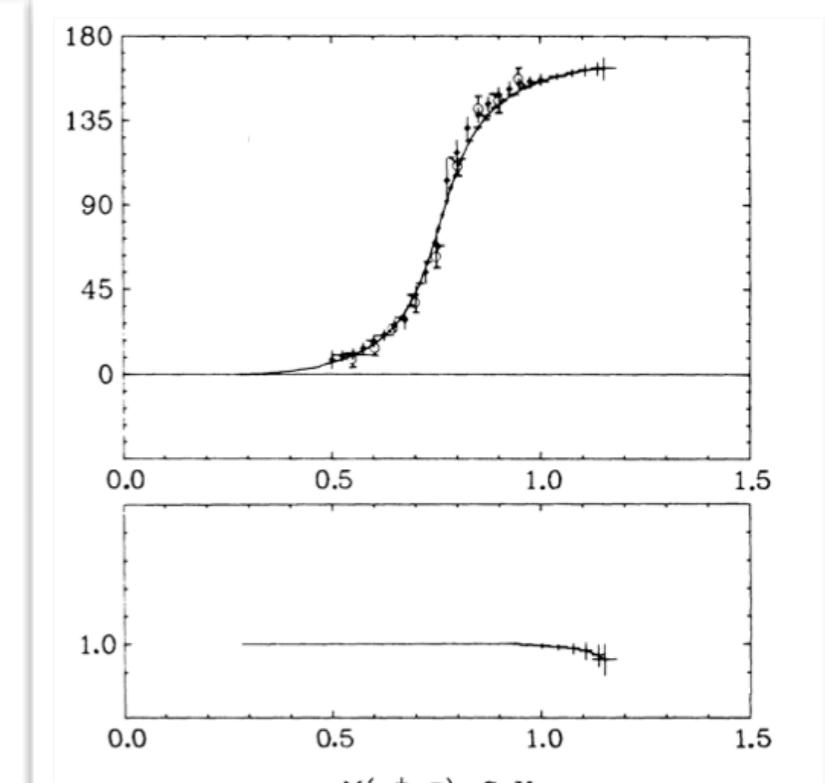


PRD87 034505 (2013)

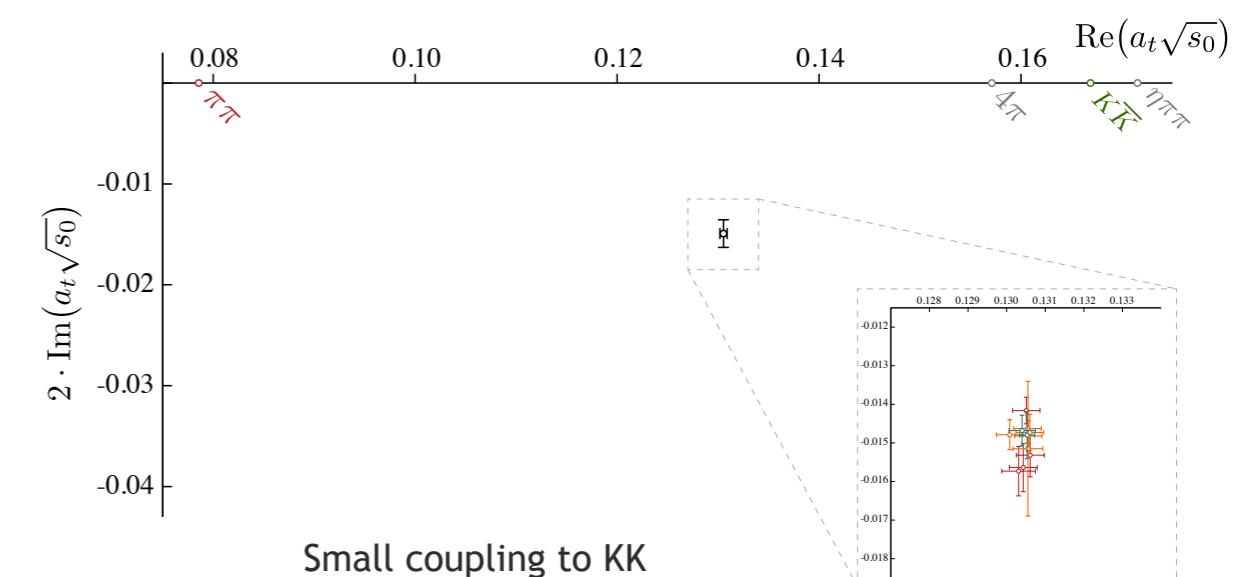
Phase shifts & inelasticity



$m_\pi \sim 236$  MeV



t-matrix pole location



# Side comment: four-particle effects

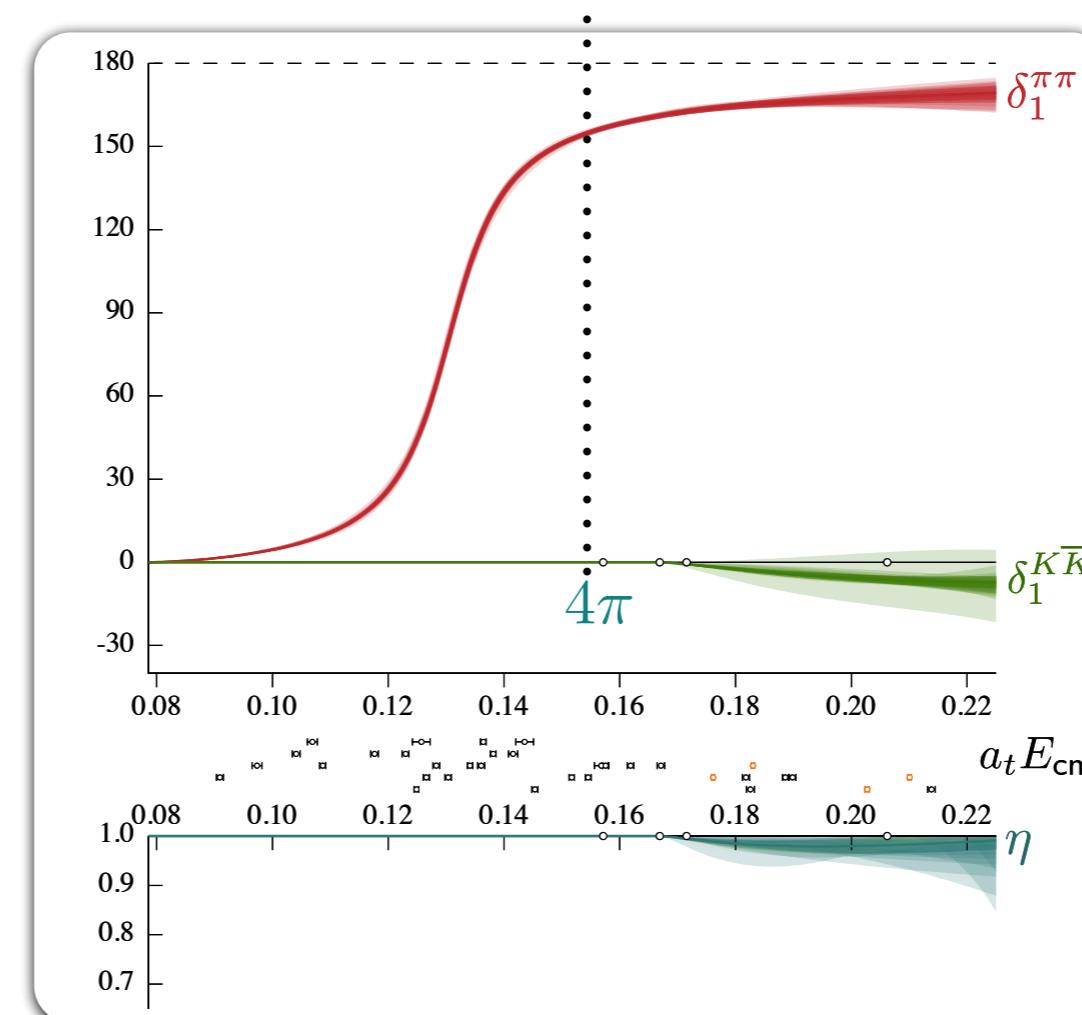
PRD95 074510 (2017)

Don't know the equation that describes  $2\pi - 4\pi$ , but must have the form:

$$\det \left[ \begin{pmatrix} \mathcal{M}_{2\pi} & \\ & \mathcal{M}_{4\pi} \end{pmatrix}^{-1} + \begin{pmatrix} \tilde{K}_{2\pi,2\pi} & \tilde{K}_{2\pi,4\pi} \\ \tilde{K}_{4\pi,2\pi} & \tilde{K}_{4\pi,4\pi} \end{pmatrix} \right] = 0$$

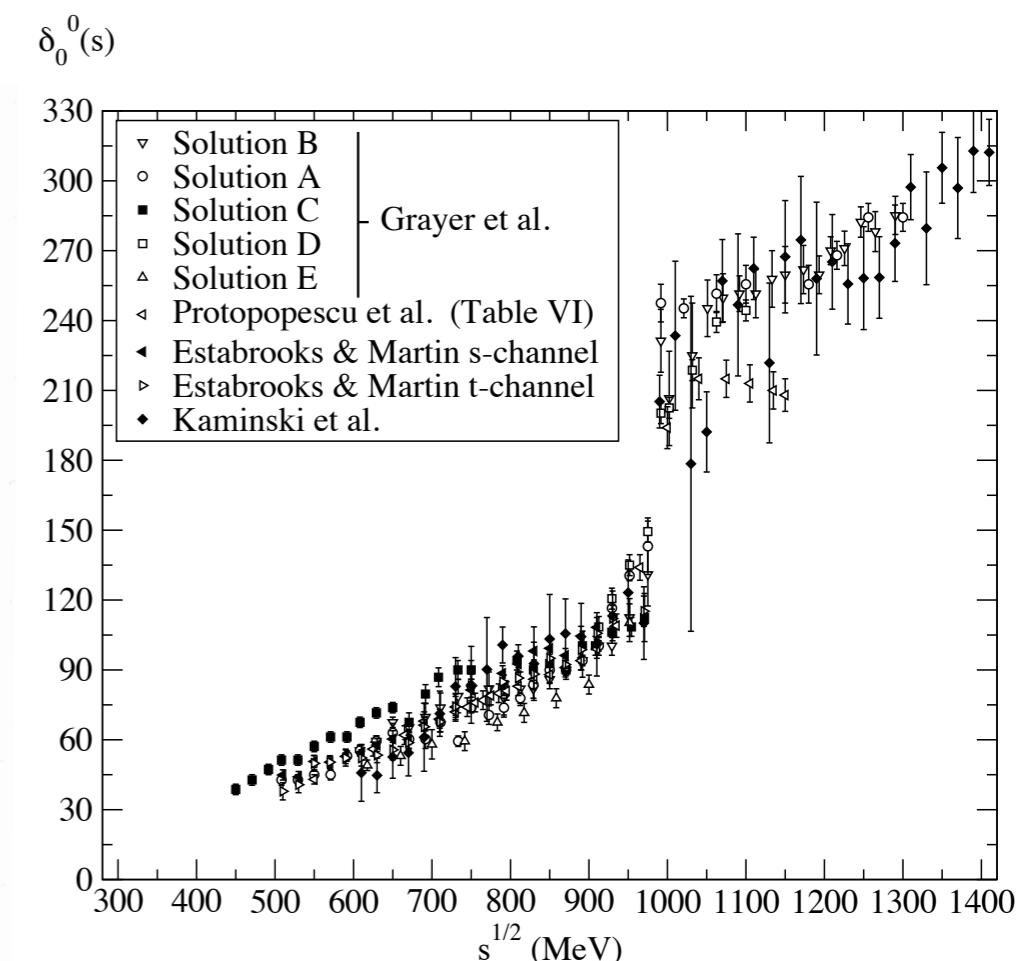
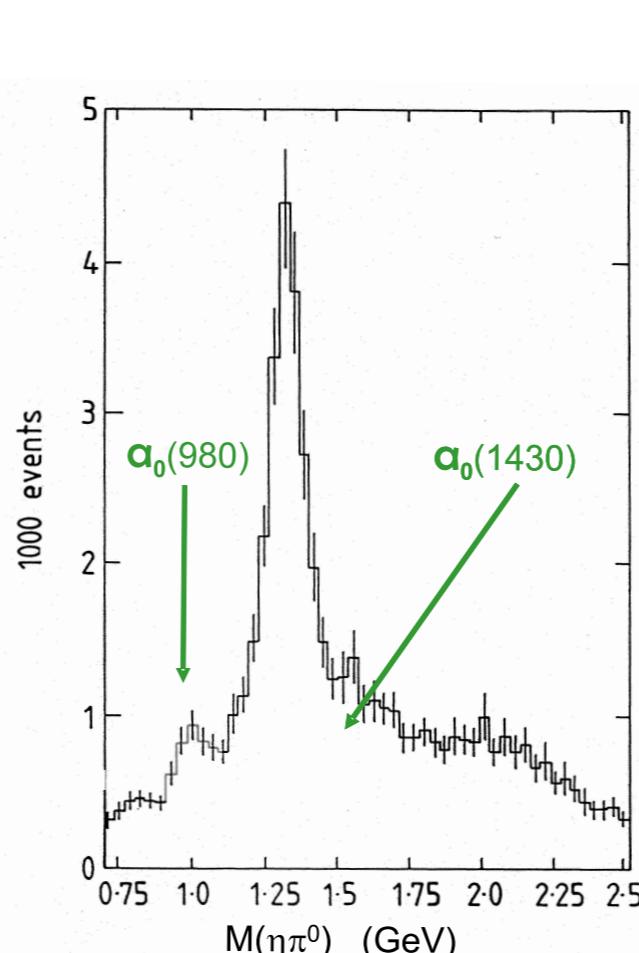
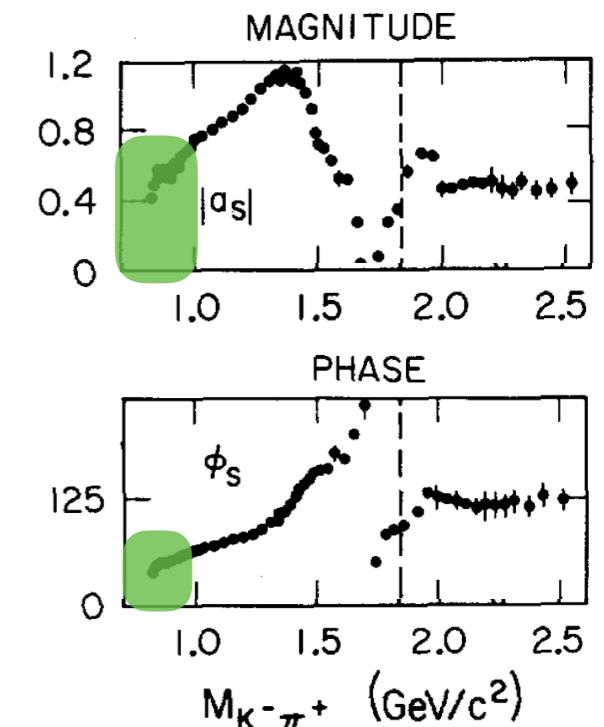
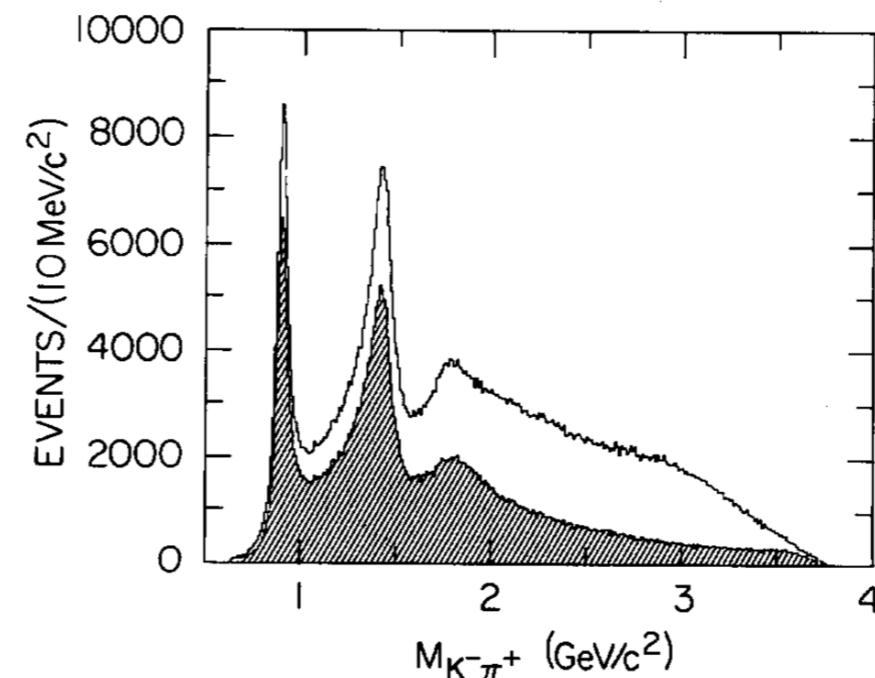
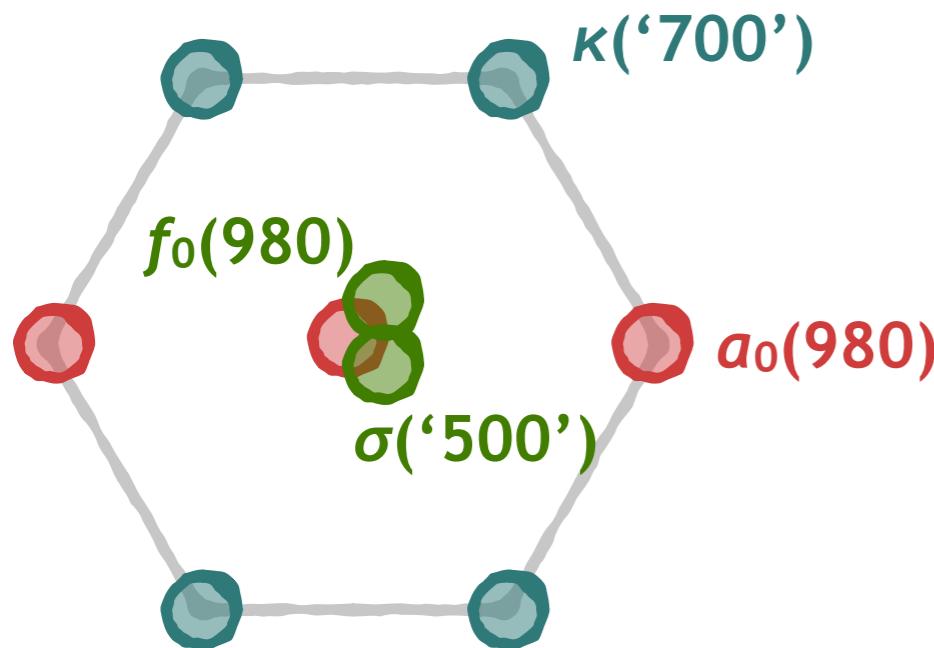
If  $\tilde{K}_{2\pi,4\pi} \sim \mathcal{O}(\epsilon)$ , then factorizes

$$\det \left[ \mathcal{M}_{2\pi}^{-1} + \tilde{K}_{2\pi,2\pi} \right] \times \det \left[ \mathcal{M}_{4\pi}^{-1} + \tilde{K}_{4\pi,4\pi} \right] + \mathcal{O}(\epsilon) = 0$$



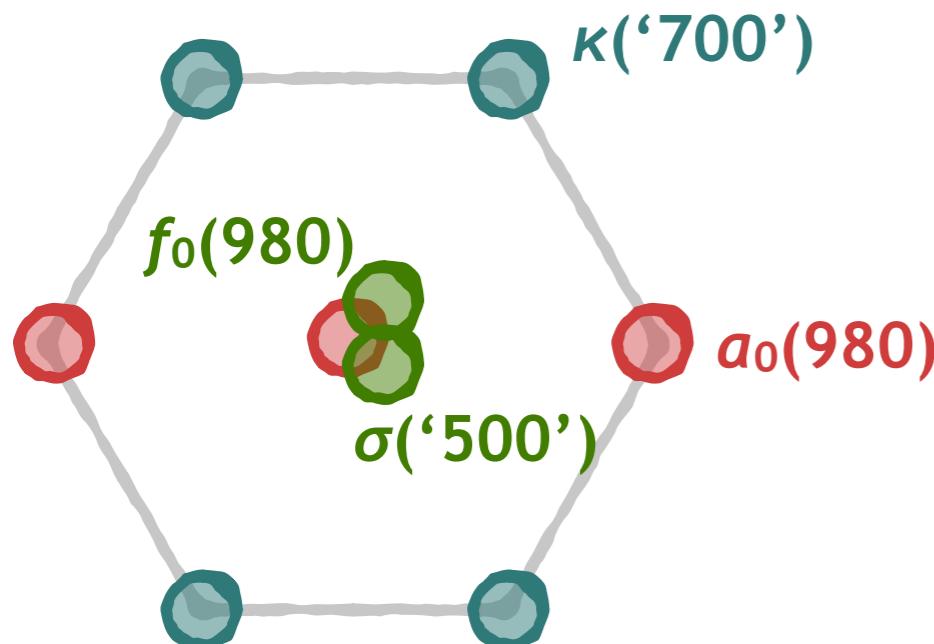
See no clear evidence of  $4\pi$

# Light scalar mesons - empirically

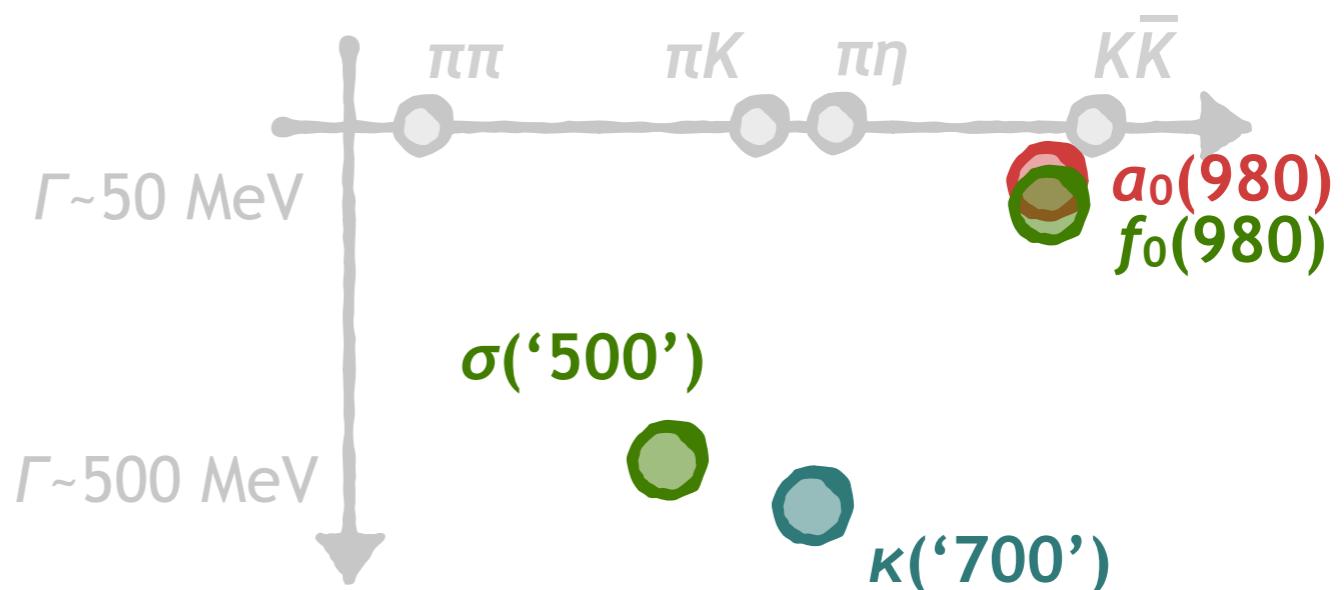


# Light scalar mesons - empirically

Conventional wisdom: an ‘inverted’ mass nonet

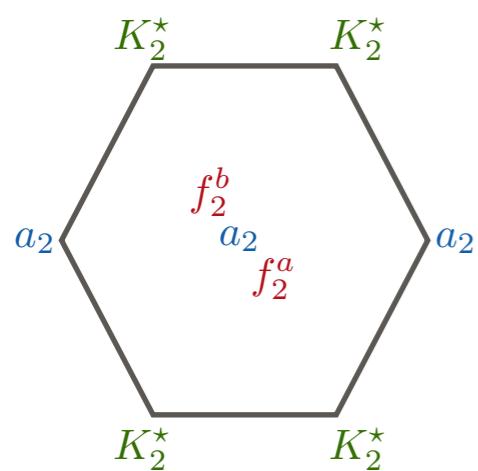


Similar? Vastly different imaginary parts...



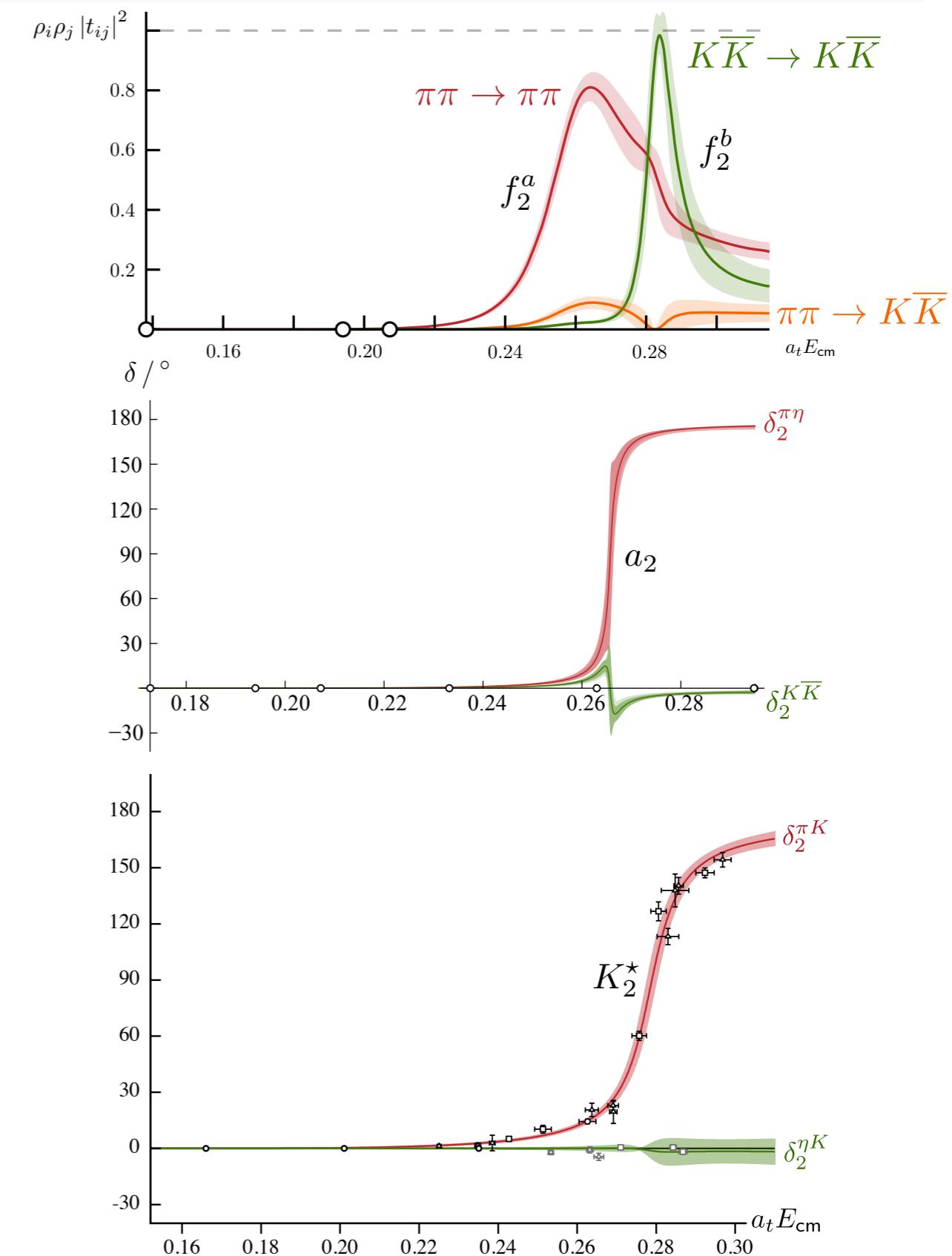
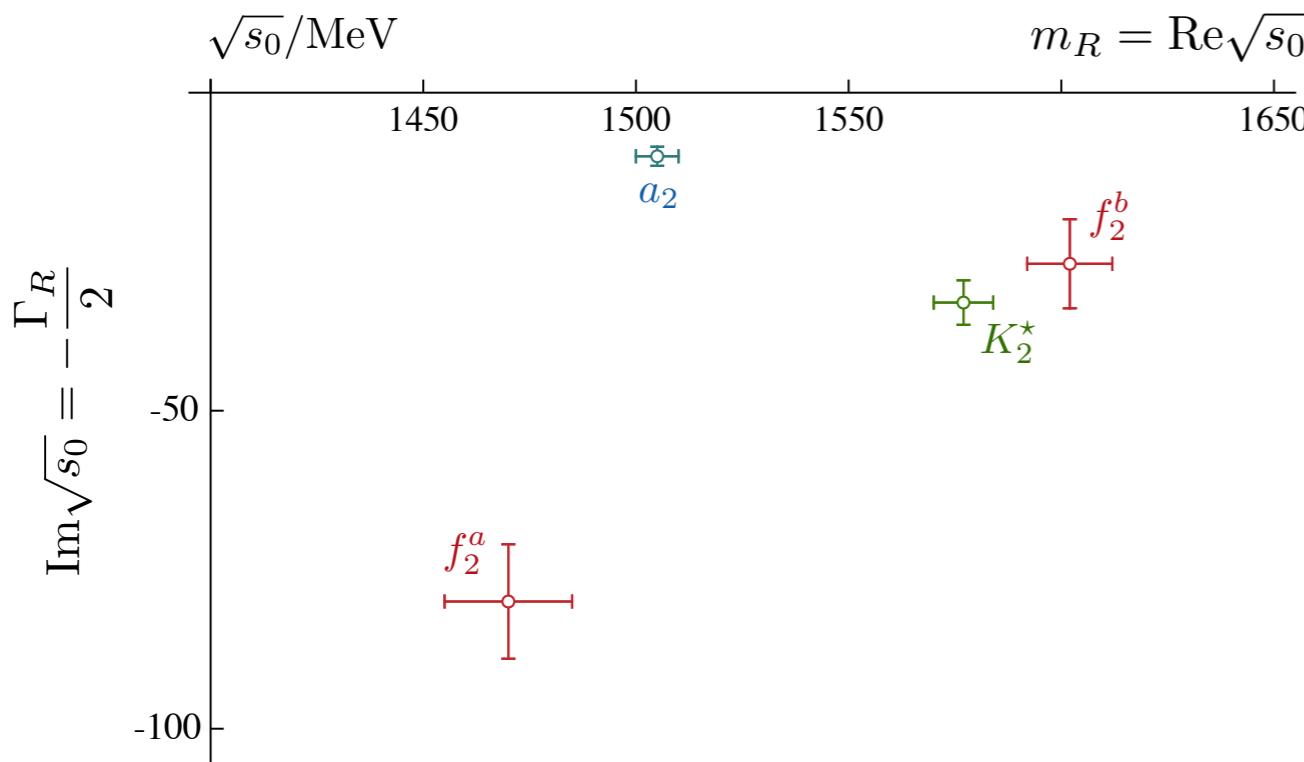
What does QCD have to say?

# Lightest tensors at $m_\pi=391$ MeV

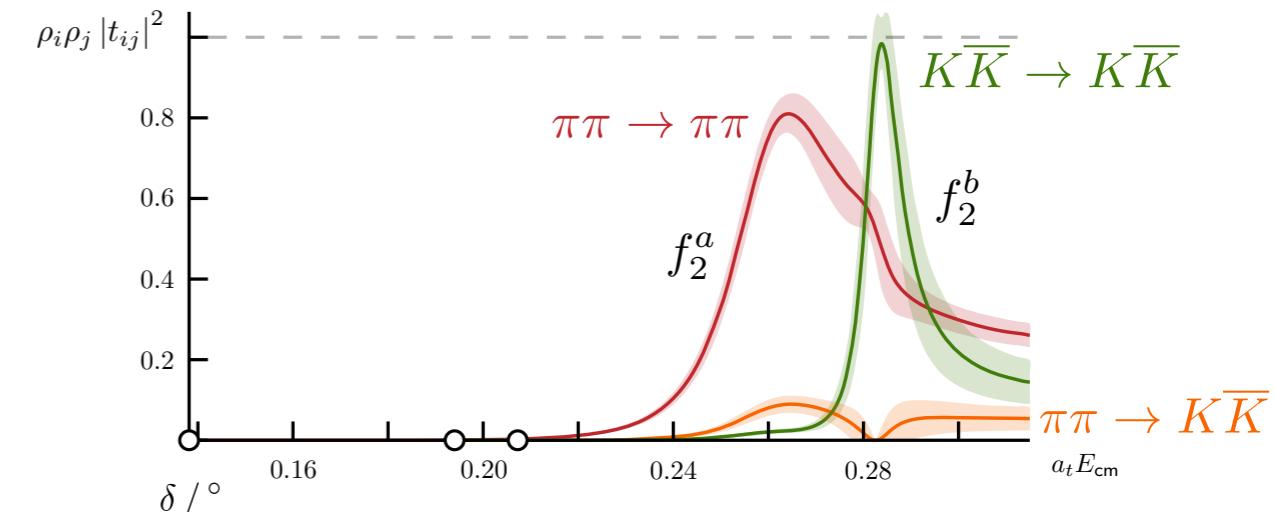
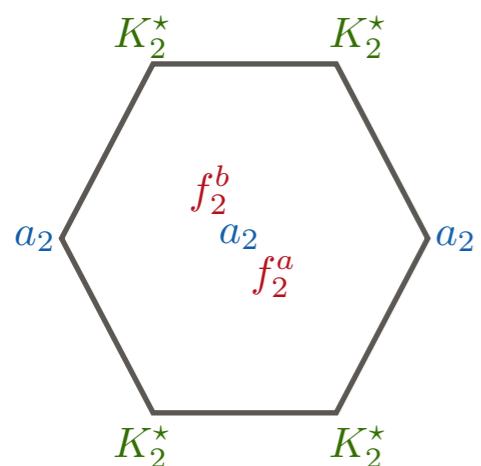


$$t_{ij} \sim \frac{c_i c_j}{s_0 - s}$$

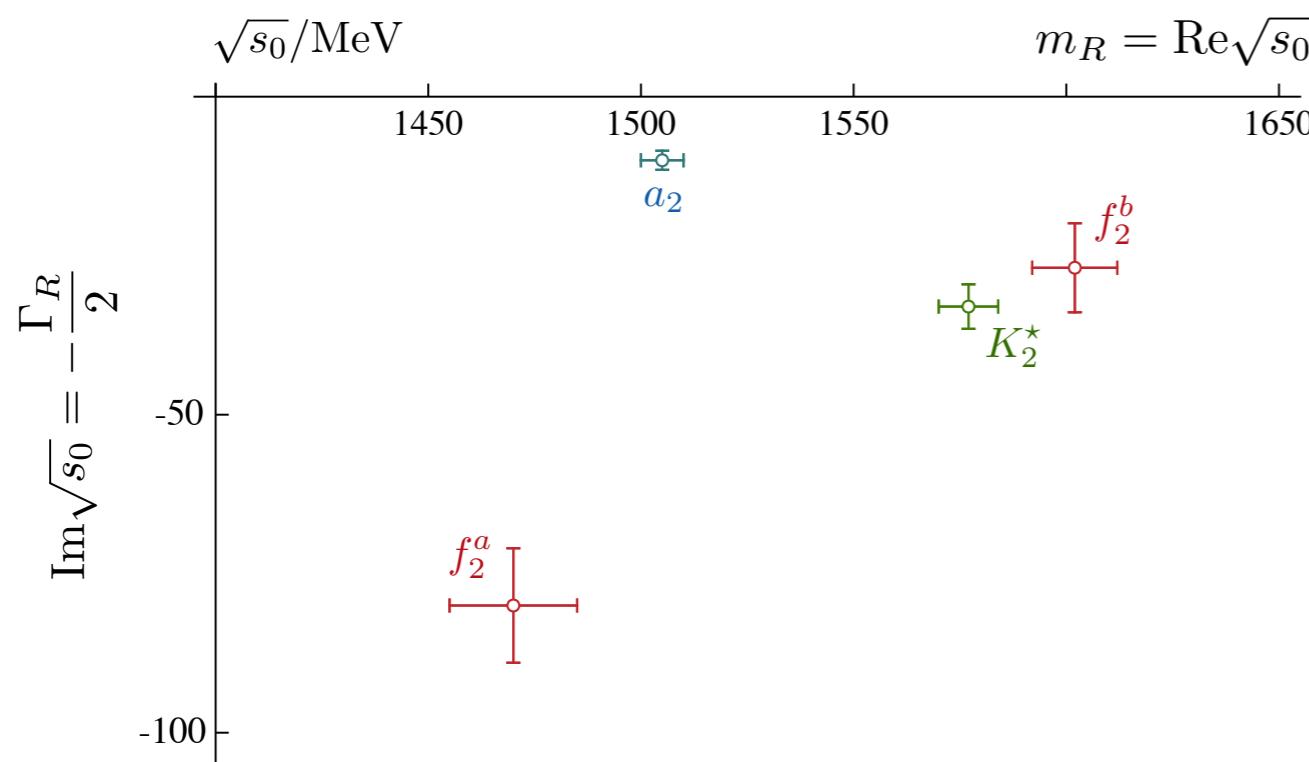
$$\sqrt{s_0} = m_R - i \frac{\Gamma_R}{2}$$



# Lightest tensors at $m_\pi=391$ MeV

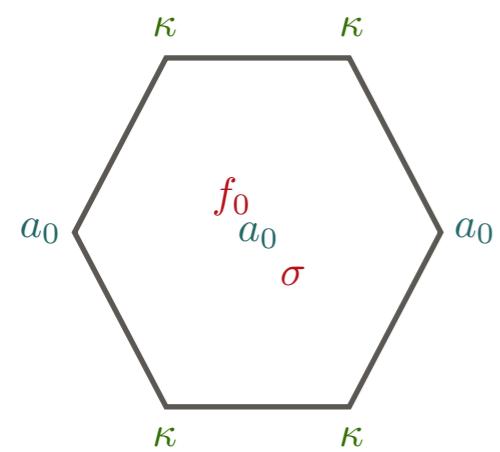


**PDG**  
 $f_2(1270)$   
*flavor tagging  
by decays*  
 $f_2(1525)$   
 $\pi\pi 84\%$   
 $K\bar{K} 89\%$



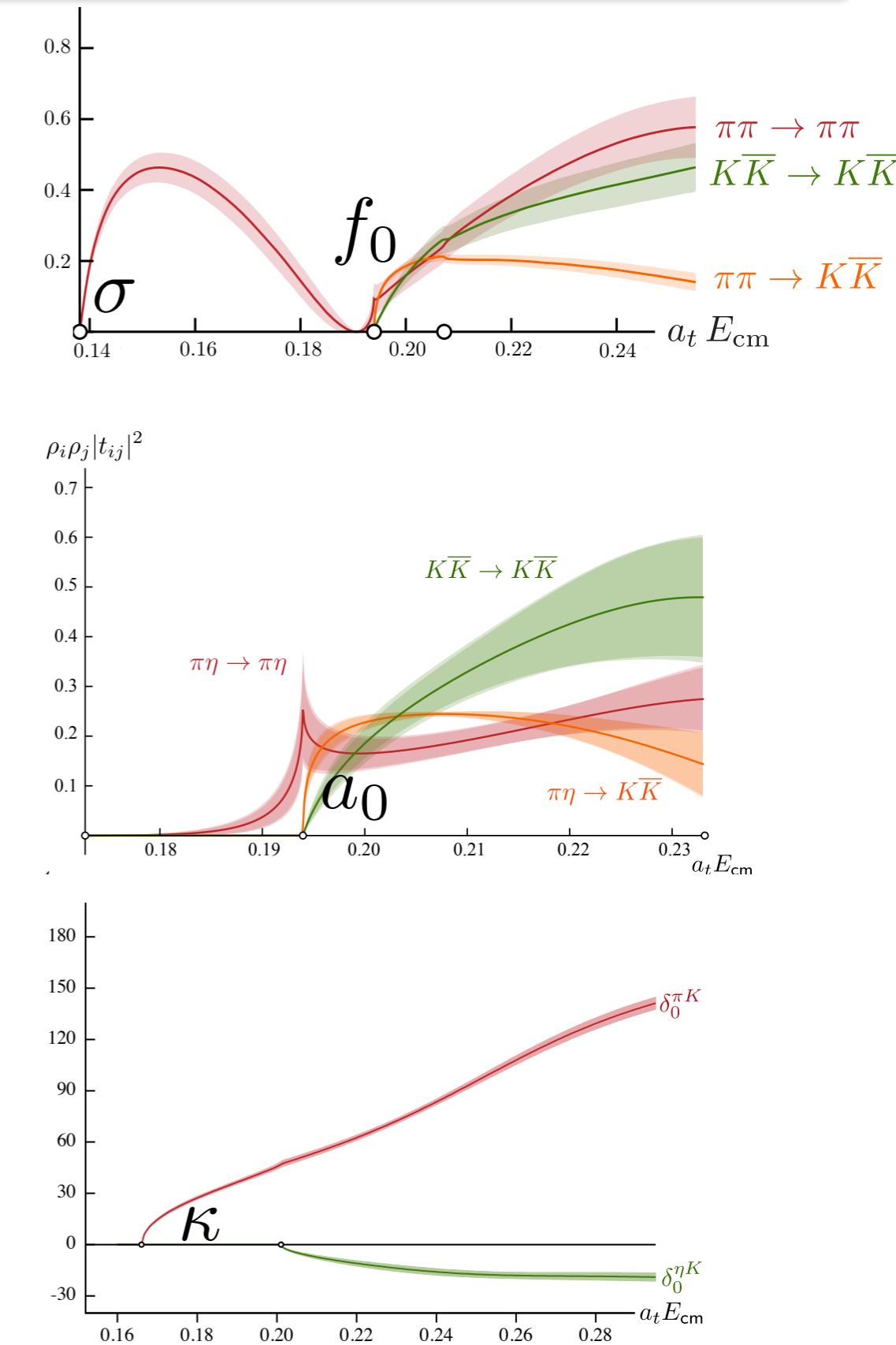
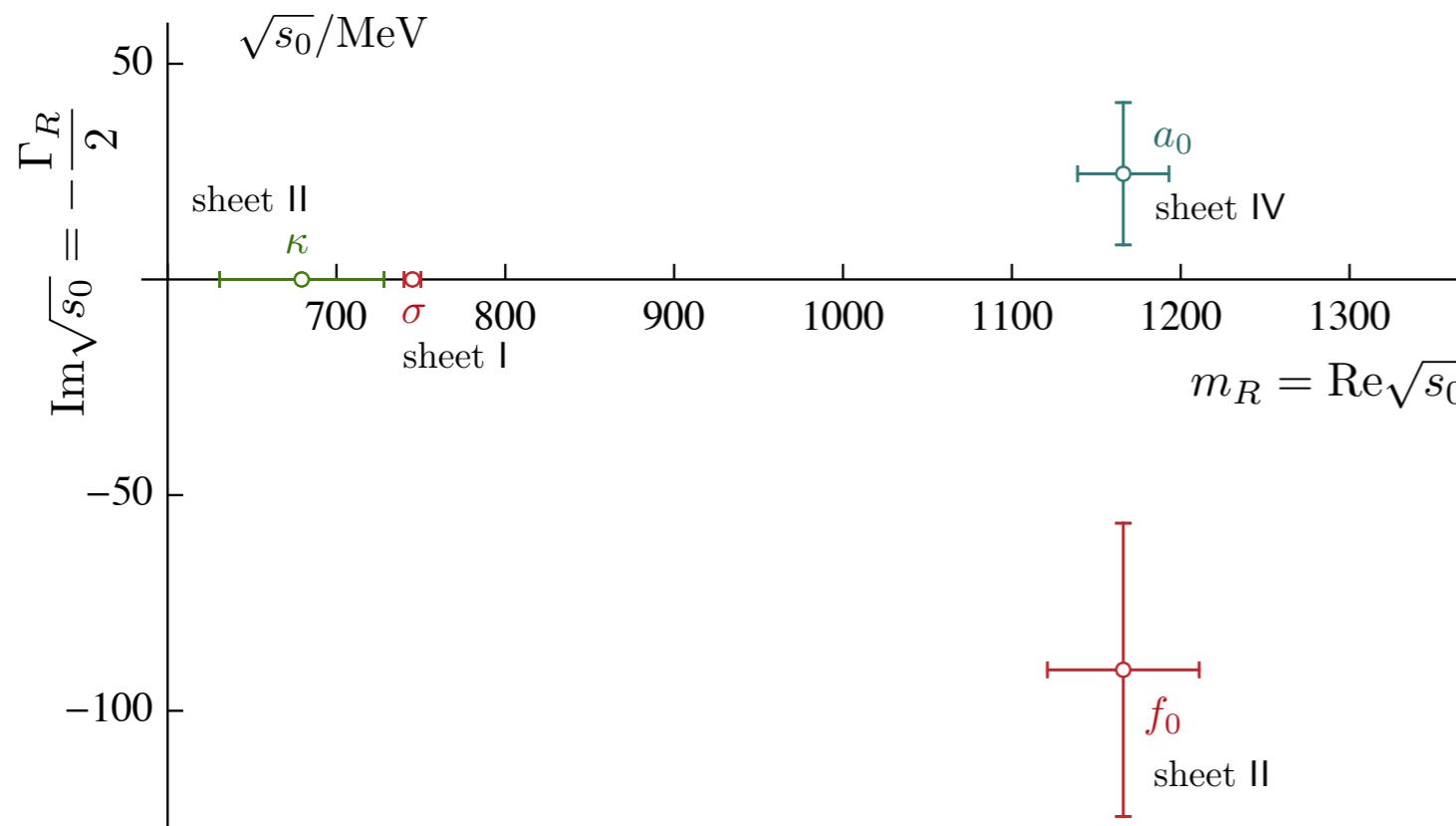
Used as motivation in  
GlueX detector upgrade proposal

# Lightest scalars at $m_\pi=391$ MeV



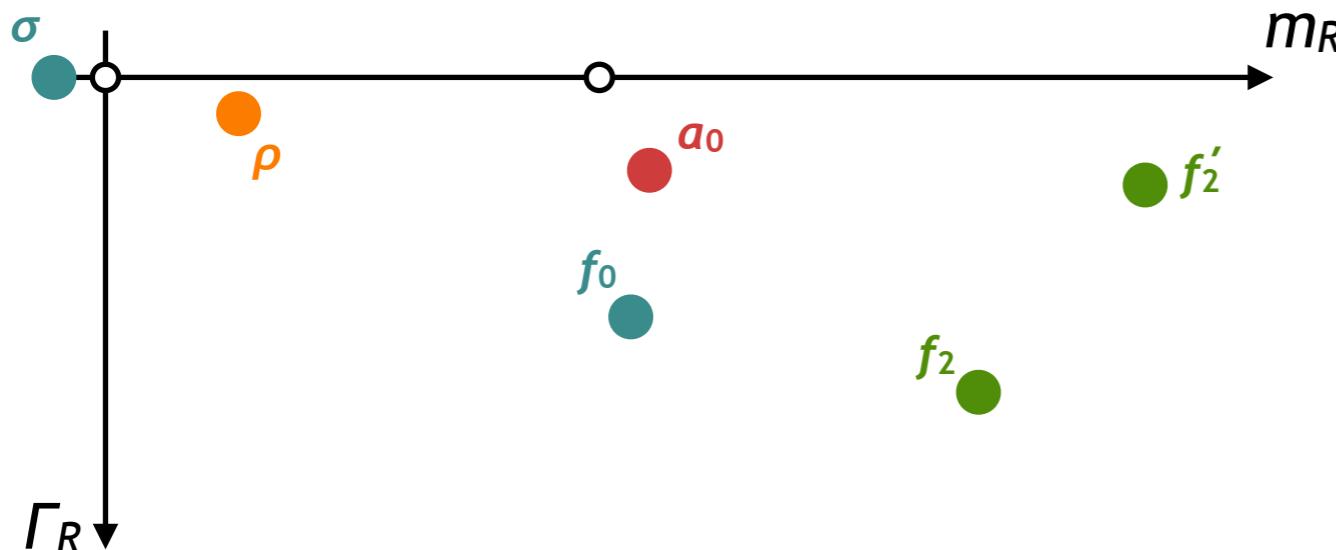
$$t_{ij} \sim \frac{c_i c_j}{s_0 - s}$$

$$\sqrt{s_0} = m_R - i \frac{\Gamma_R}{2}$$

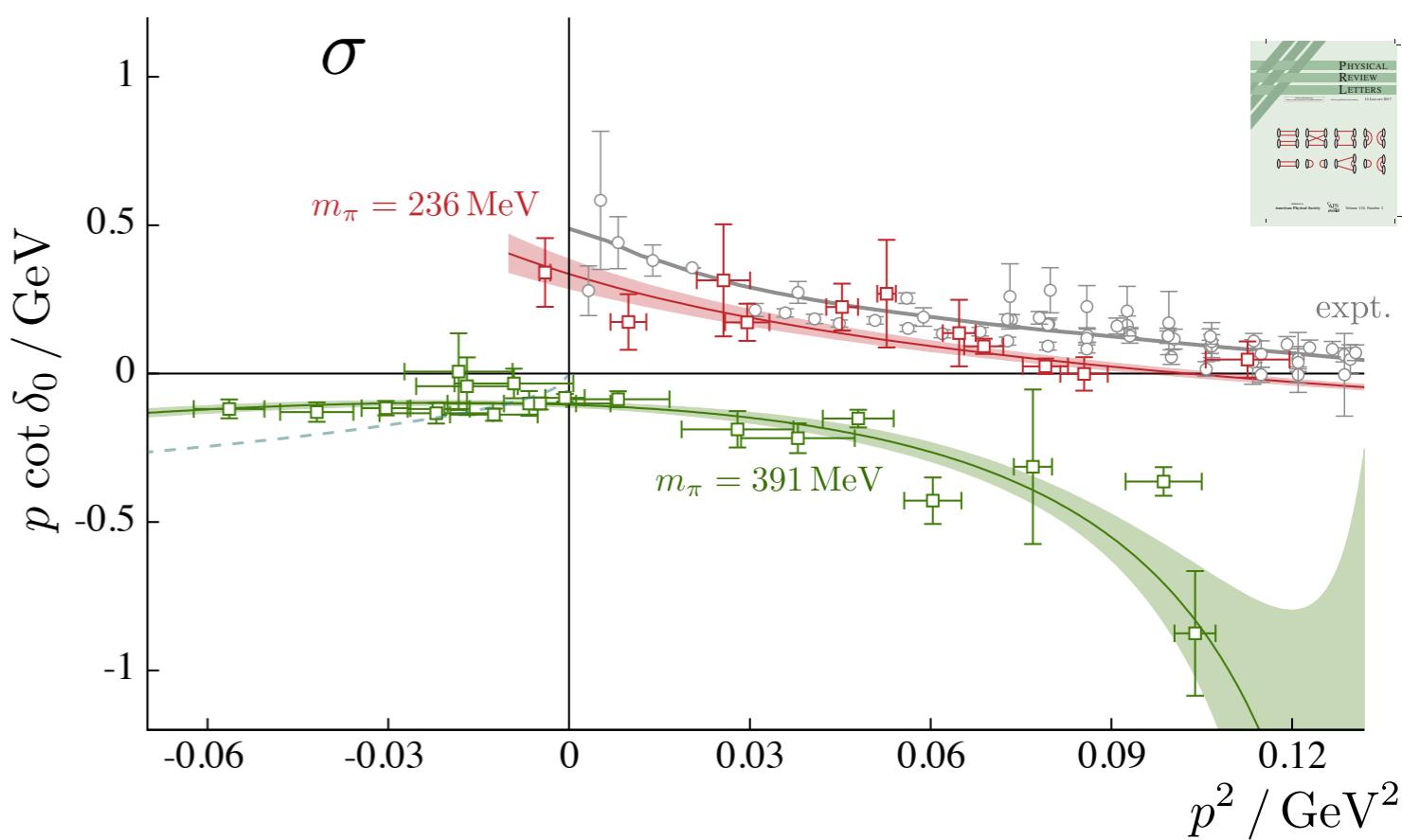
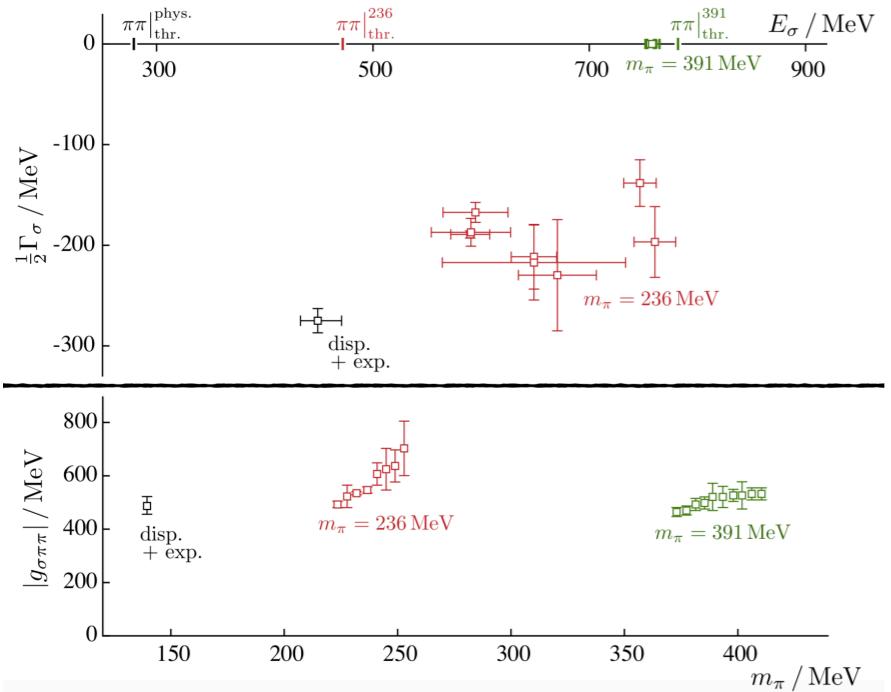
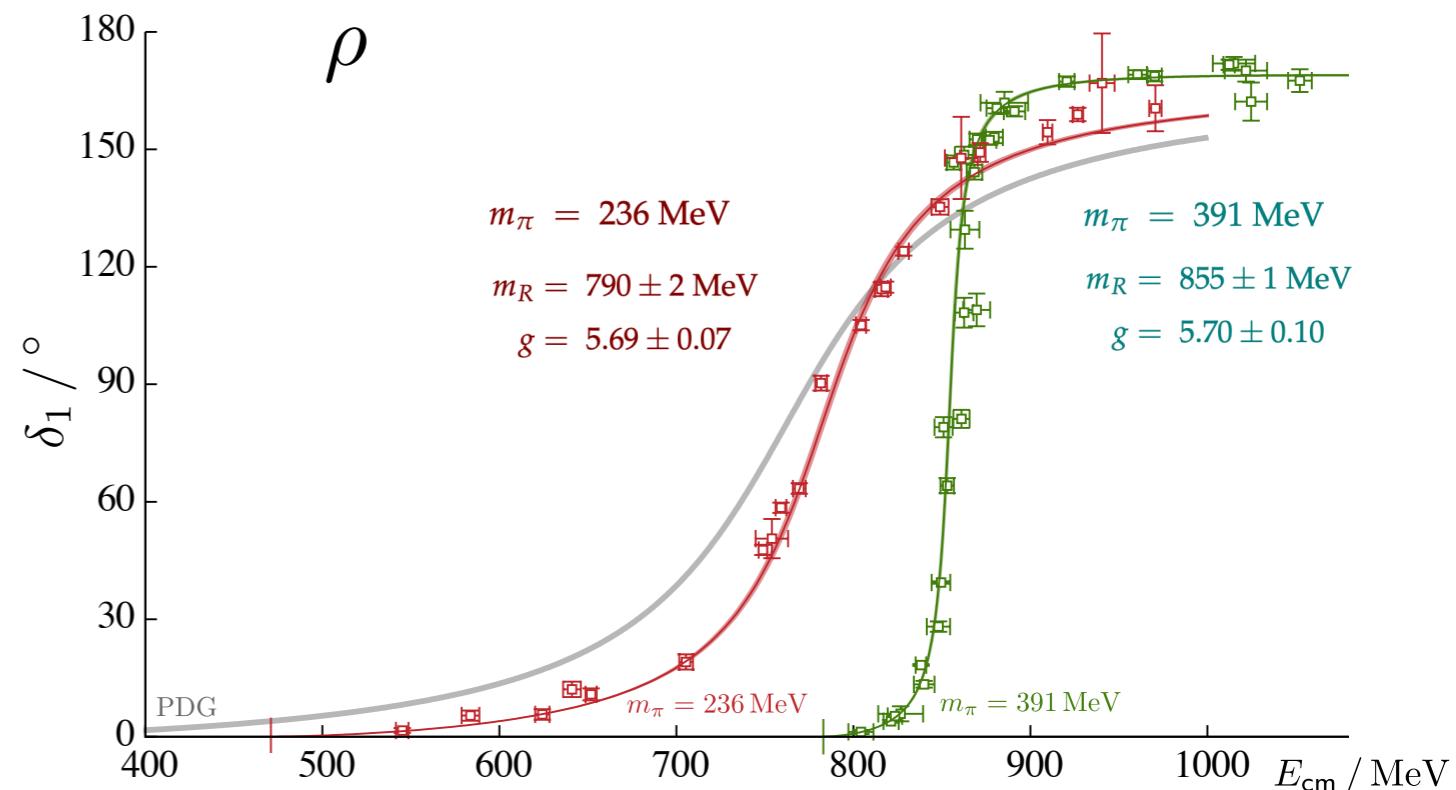
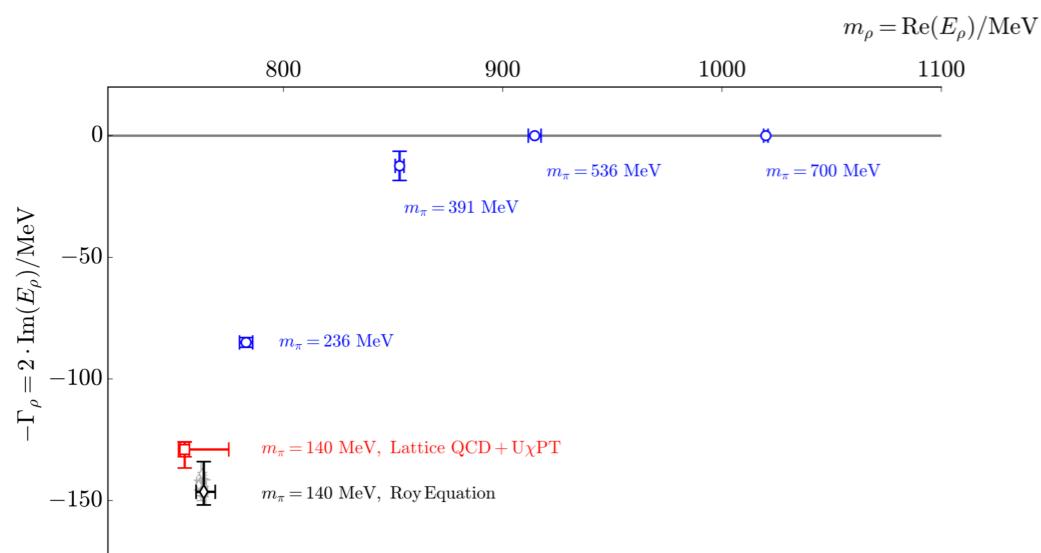


# light meson resonances at $m_\pi \sim 391$ MeV

---



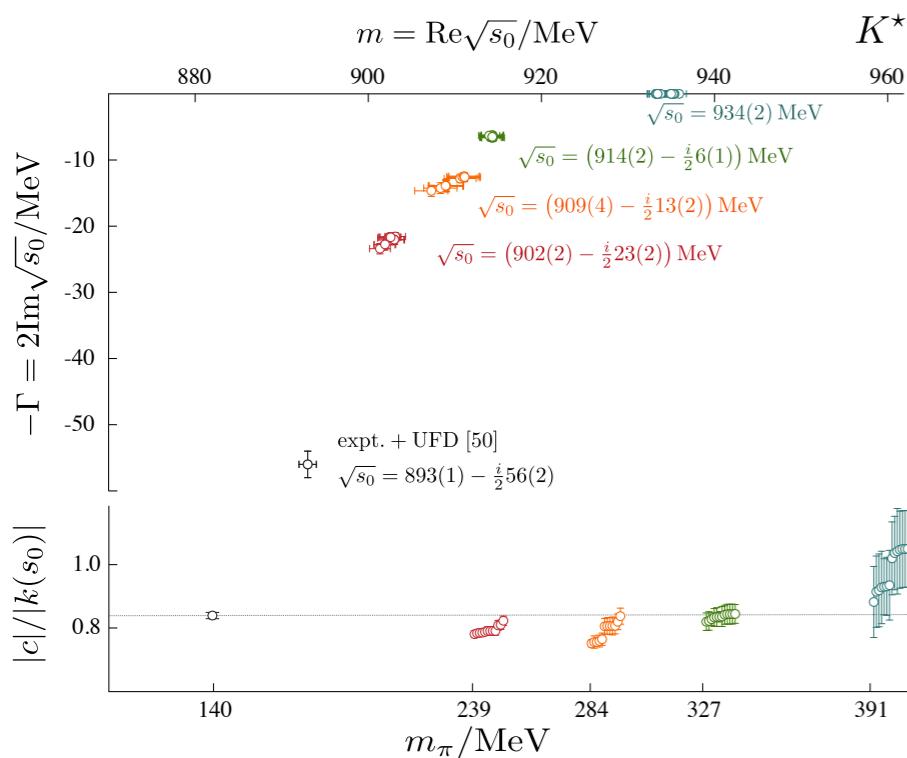
# Quark mass dependence: I=0 & 1



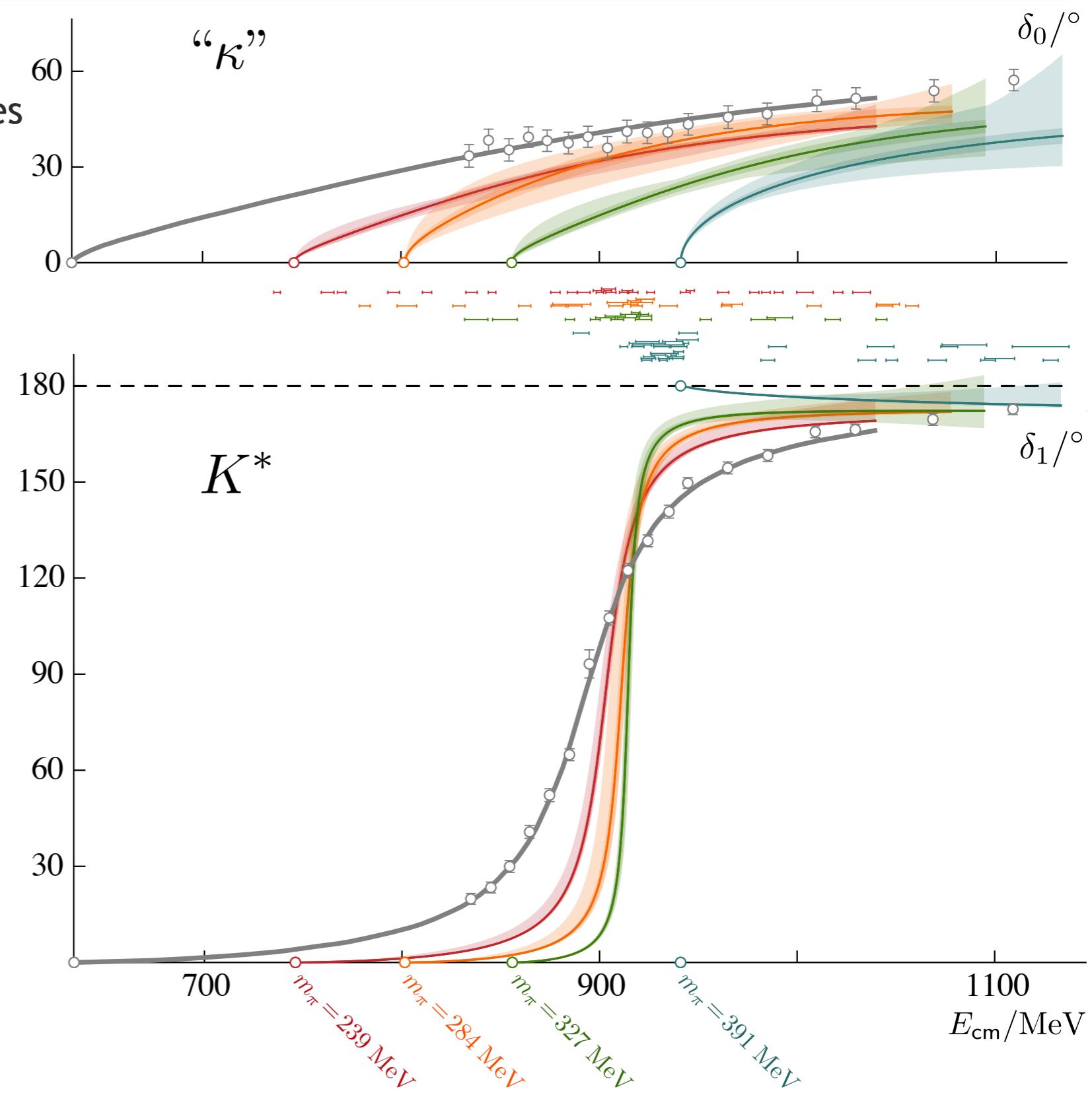
# Quark mass dependence: $|l| = 1/2$

arXiv:1904:03188

A more complicated story...  
need “t” & “u”-channel amplitudes

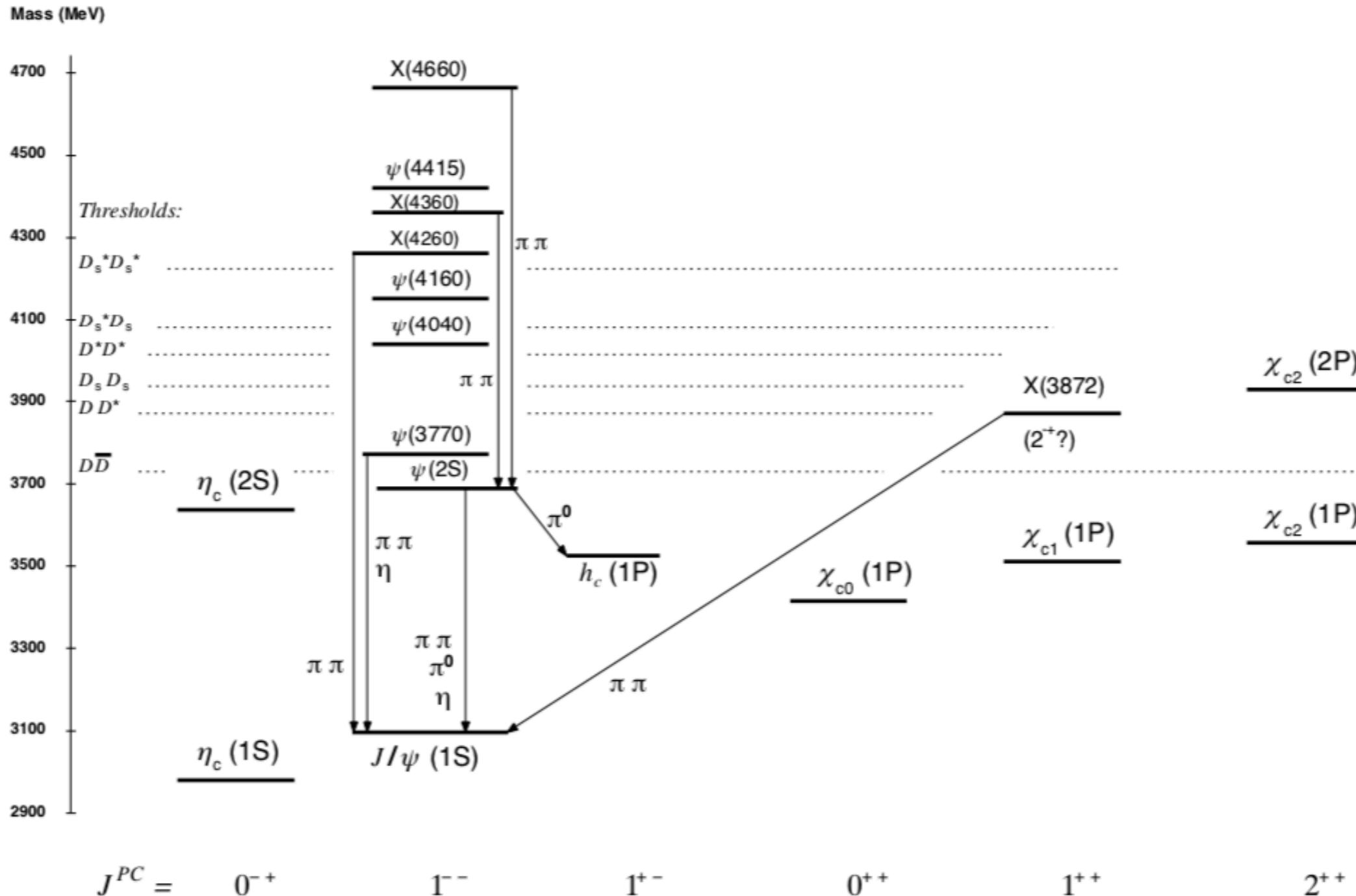


Added new dynamical ensembles



# Charmonium resonances

Several resonances reported near  $D\bar{D}$  thresholds

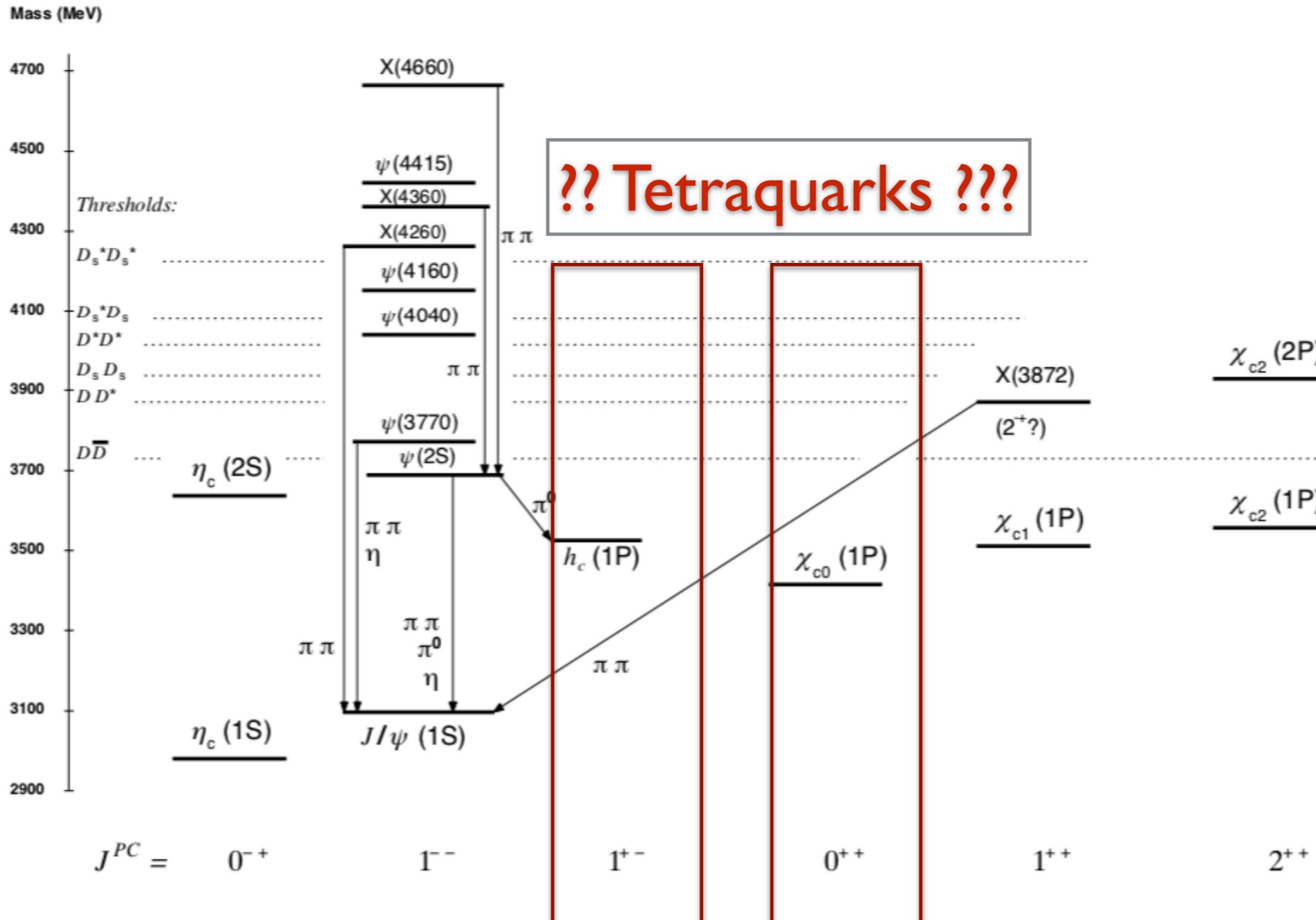


$$J^{PC} = \quad 0^{-+} \quad 1^{--} \quad 1^{+-} \quad 0^{++} \quad 1^{++} \quad 2^{++}$$

# Charmonium resonances

PDG

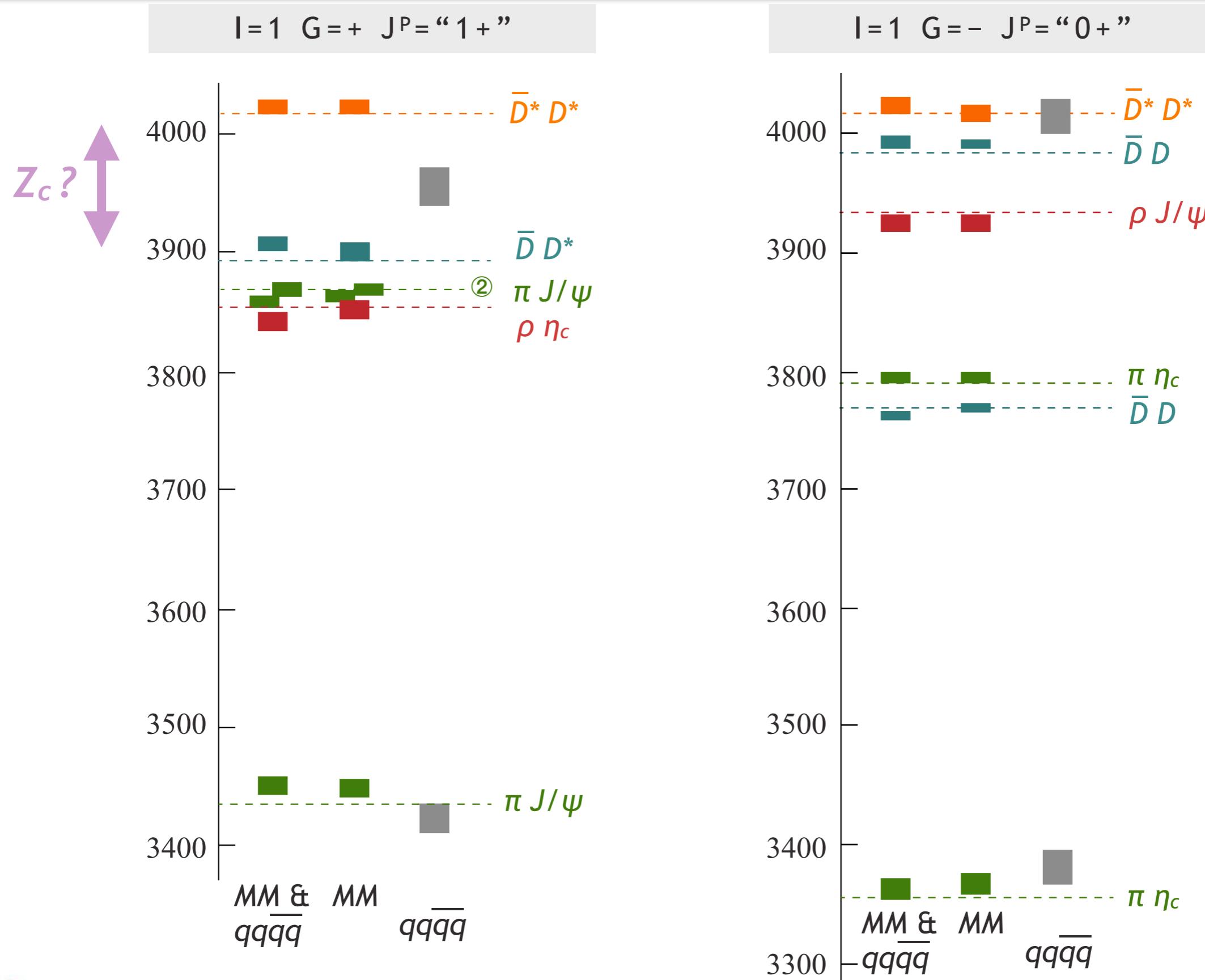
Isospin 1 charmonium?



# Tetraquarks ?

$m_\pi \sim 391$  MeV

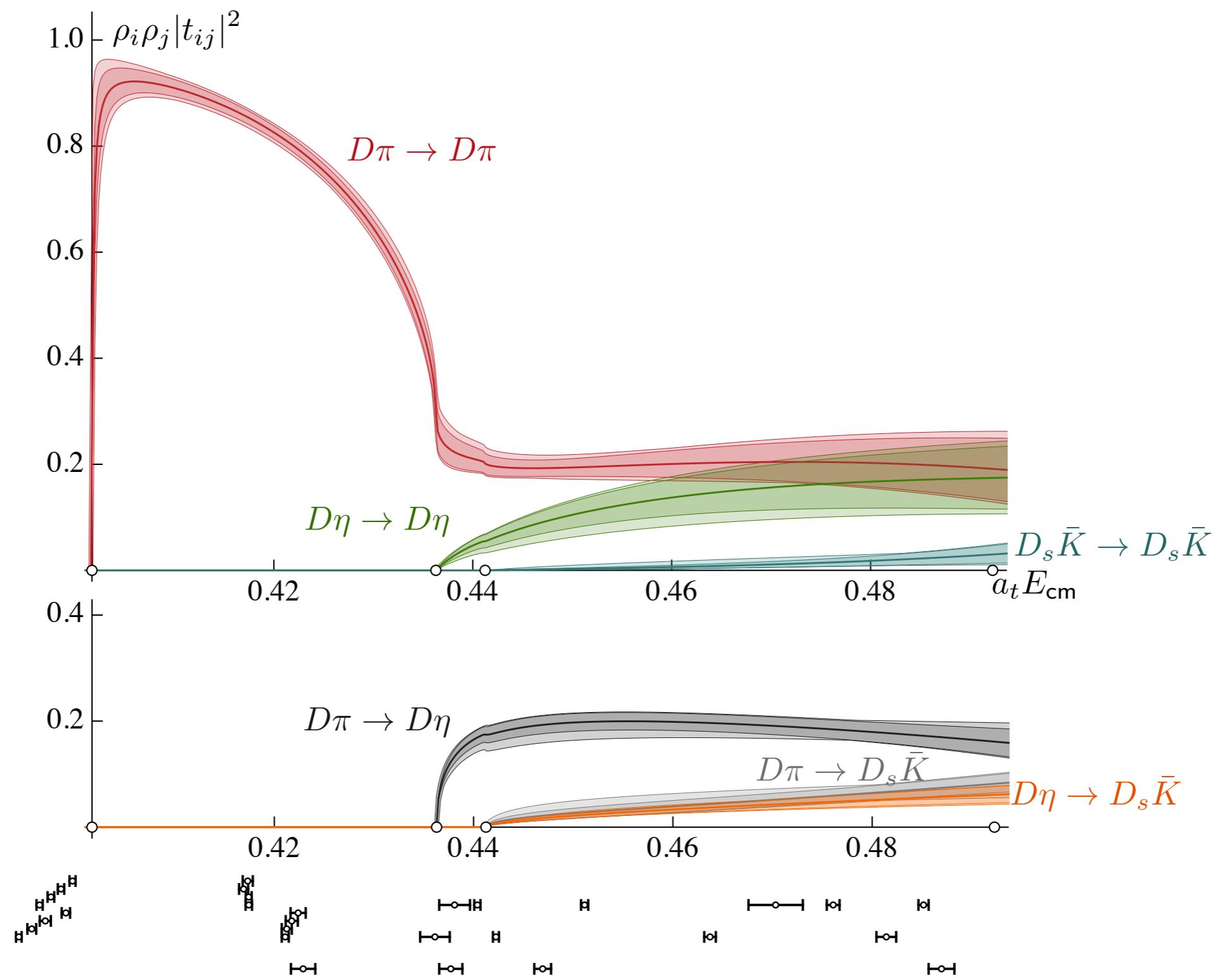
JHEP 1711 033 (2017)



# S-wave D decays

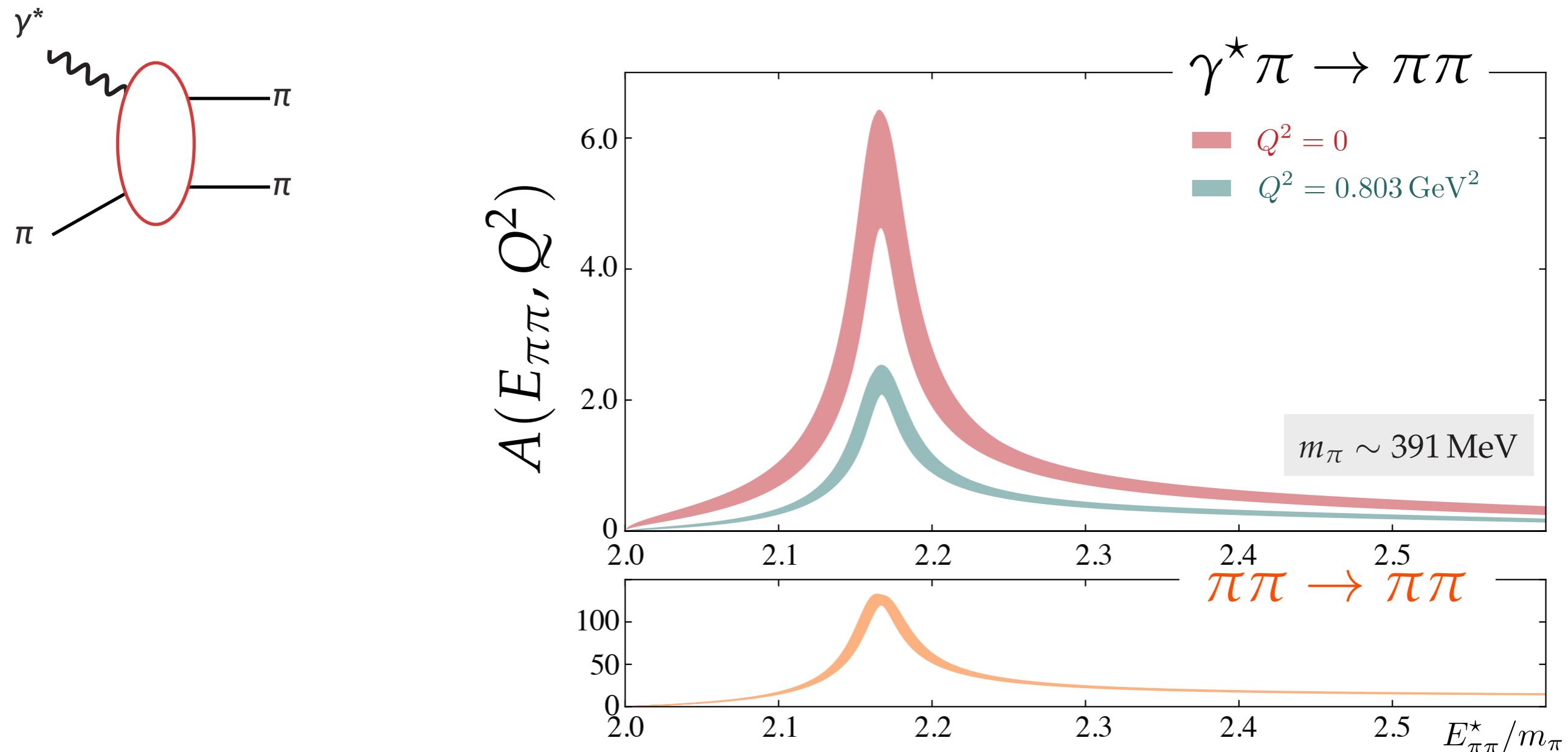
Moir et al, JHEP 1610 (2016) 011

Sharp threshold behavior in D &  $D_s$



# Coupling resonances to currents

- Production mechanisms - e.g., photo-production



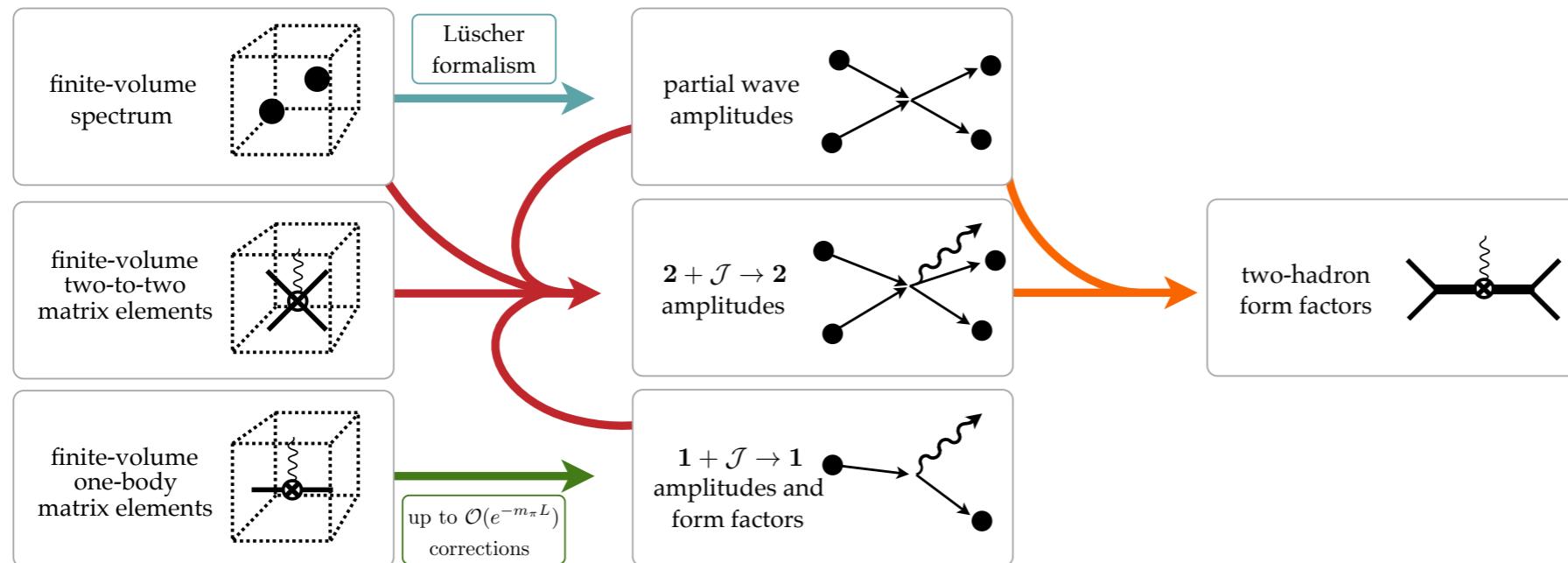
PRL 115 242001 (2015)  
arXiv: 1604.03530

# $2 + \mathcal{J} \rightarrow 2$ transition amplitudes

arXiv:1812.10504

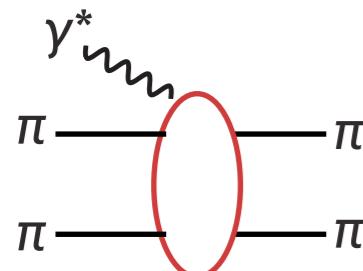
Not directly accessible experimentally

Provides a method to probe internal structure of resonances (e.g., charge radii, form-factors, even PDF-s)



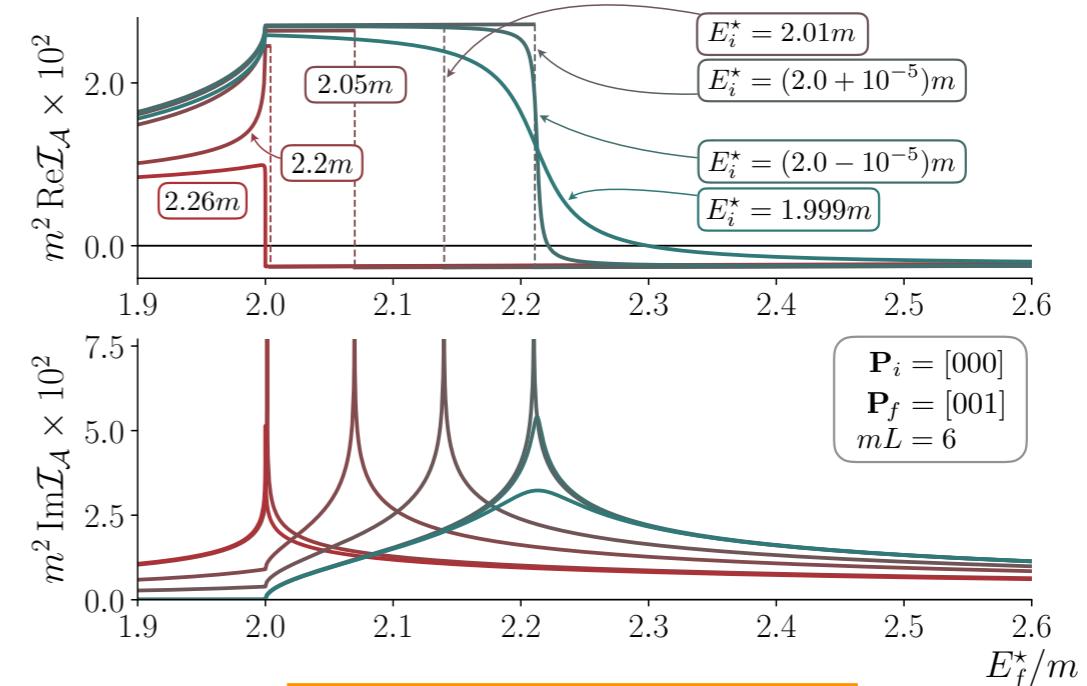
## Establishing formalism

E.g.,  $\gamma^* \pi\pi \rightarrow \pi\pi$



$$\mathcal{A}(E_{\pi\pi}^{\text{in}}, E_{\pi\pi}^{\text{out}}, Q^2)$$

Energy regimes probe diagrams contributing



# Status of scattering formalisms

Two-body coupled channel(s)

- finite-volume formalism ✓
- amplitude formalism (good) except cross-channels (not-good)
  - can provide branching fractions for phenomenologists😊 😐

## ELEMENTARY PARTICLE THEORY

A. D. MARTIN

*Department of Physics,  
University of Durham*

T. D. SPEARMAN

*School of Mathematics,  
Trinity College, Dublin*



1970



冯绍峰 张静初 吴刚 唐嫣 杜淳 张翰  
FENG SHAOFENG ZHANG JINGCHU WU GANG TANG YAN DU CHUN ZHANG HAN

杜志国 邢佳栋 张光北 胡明 程媛媛 胡海锋 李炫臻  
DU ZHIGUO XING JIADONG ZHANG GUANGBEI HU MING CHENG YUANYUAN HU HAIFENG LI XUANZHEN

张番番 导演  
DIRECTOR ZHANG PANPAN

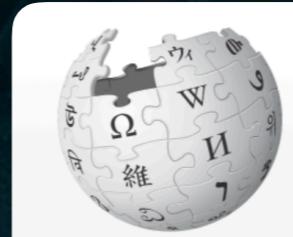
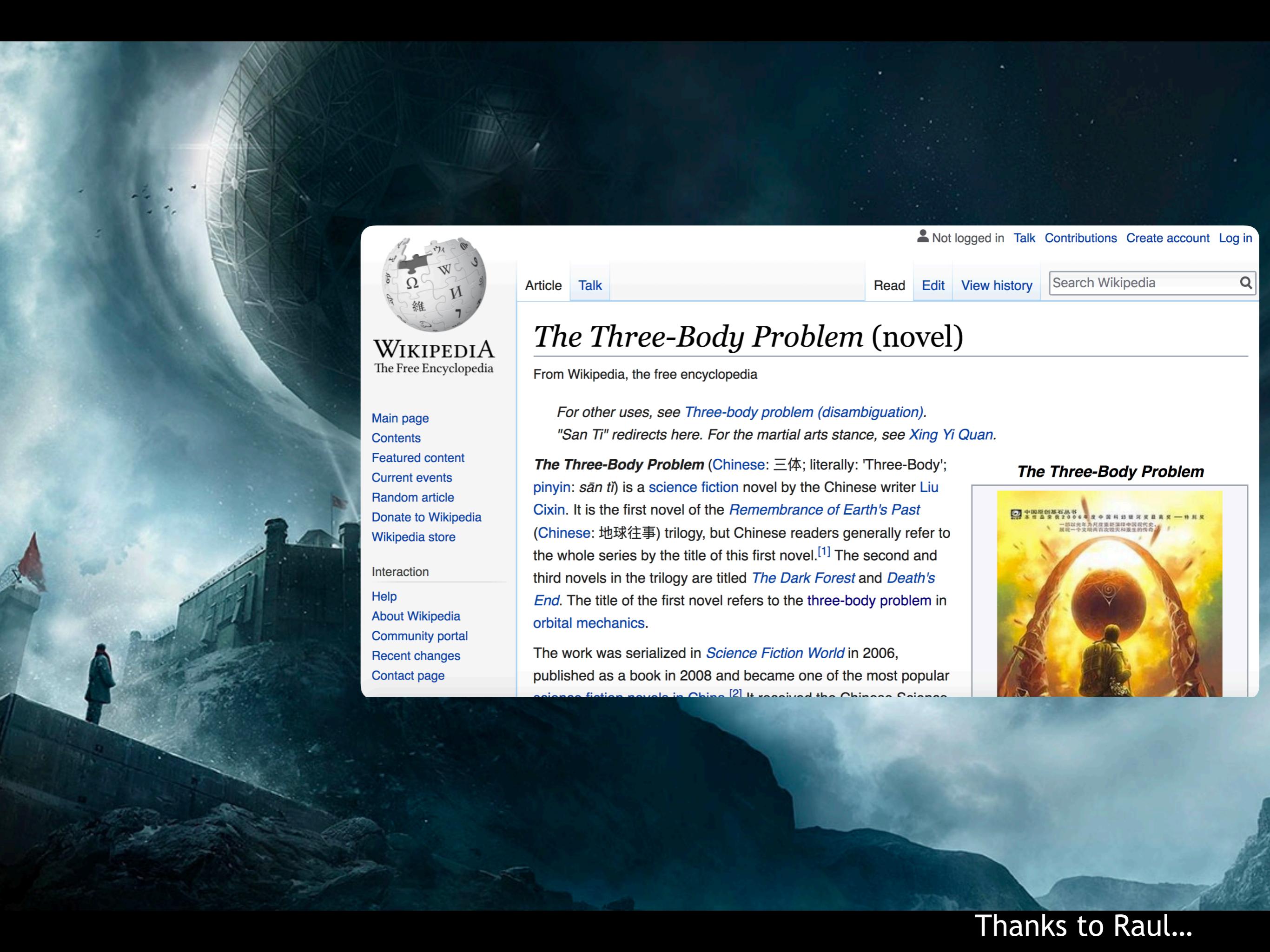
# 三体

刘慈欣 监制  
EXECUTIVE PRODUCER LIU CIXIN

3 BODY IN 3D

*"the three-body problem"*

2019



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# The Three-Body Problem (novel)

From Wikipedia, the free encyclopedia

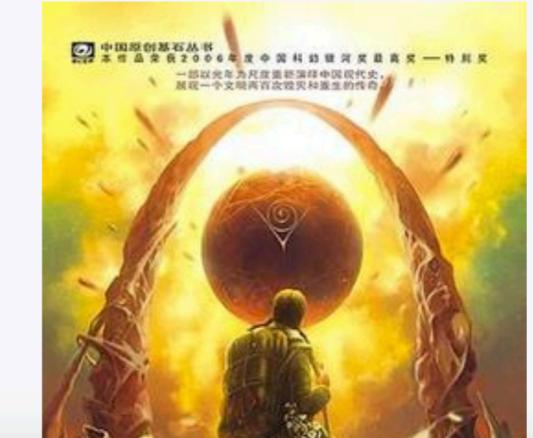
*For other uses, see [Three-body problem \(disambiguation\)](#).*

*"San Ti" redirects here. For the martial arts stance, see [Xing Yi Quan](#).*

**The Three-Body Problem** (Chinese: 三体; literally: 'Three-Body'; pinyin: *sān tǐ*) is a science fiction novel by the Chinese writer Liu Cixin. It is the first novel of the *Remembrance of Earth's Past* (Chinese: 地球往事) trilogy, but Chinese readers generally refer to the whole series by the title of this first novel.<sup>[1]</sup> The second and third novels in the trilogy are titled *The Dark Forest* and *Death's End*. The title of the first novel refers to the three-body problem in orbital mechanics.

The work was serialized in *Science Fiction World* in 2006, published as a book in 2008 and became one of the most popular science fiction novels in China.<sup>[2]</sup> It received the Chinese Science

## The Three-Body Problem



Thanks to Raul...

# Status of scattering formalisms

## Three-body coupled channel

- unitarity in 3-body

[arXiv:1905.11188](#)



- finite-volume formalism for spinless particles

- coupled 2-body & 3-body

[PRD99 014506 \(2018\)](#)



- determinant condition



- integral equations



- 3-body amplitude formalism - not well understood!



- comparisons and cross-checks of existing forms

[arXiv:1905.12007](#)



- need 3-body energies. Codes ready and optimized!

NPLQCD

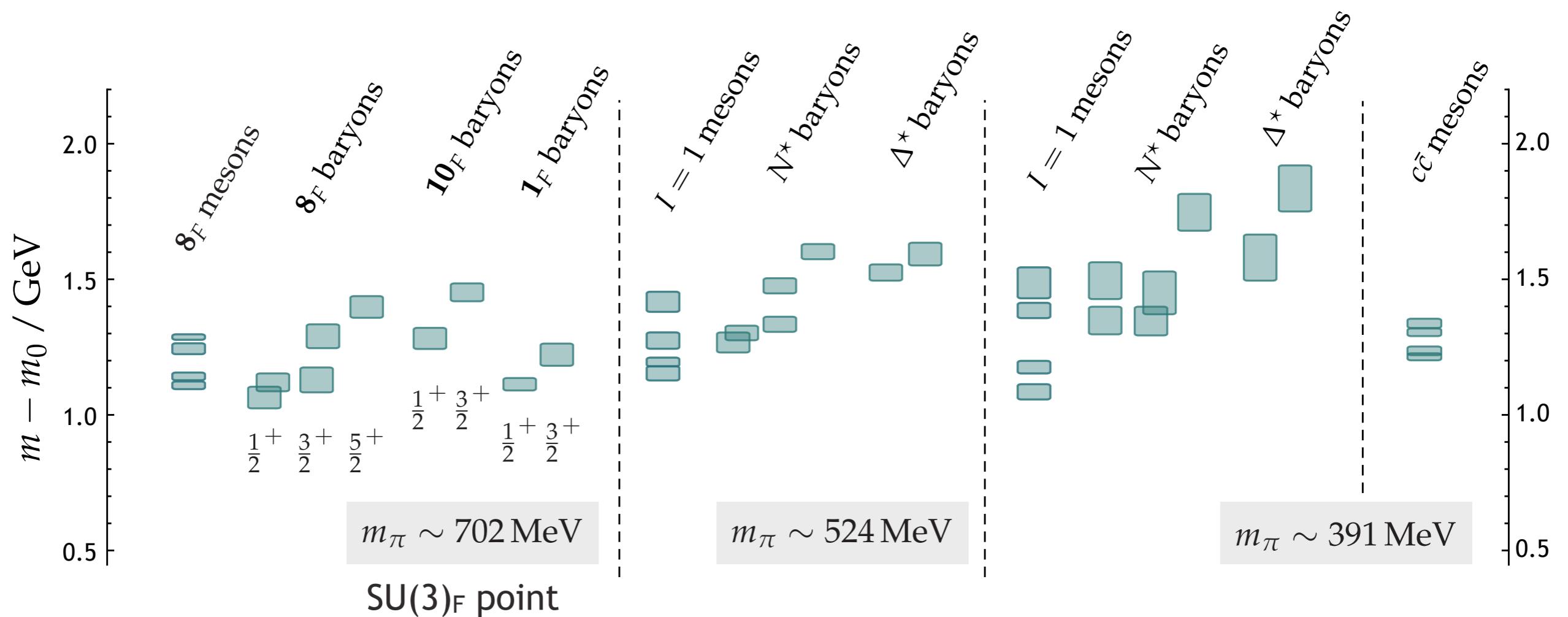
[JHEP 1711 033 \(2017\) - tetraquarks](#)

[arXiv:1904.04136 - b1](#)

[arXiv:1905.04277 -  \$3\pi\$](#)

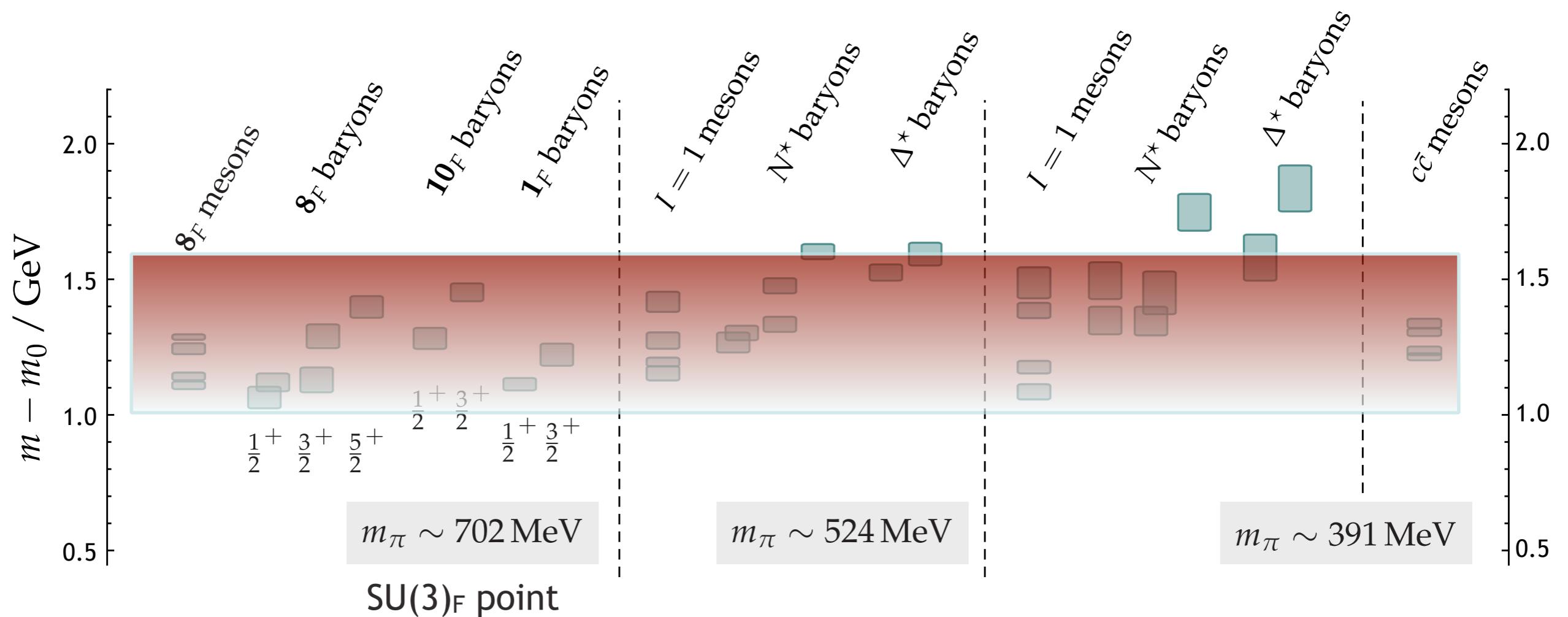
# Hadron Spectroscopy - role of the glue

- Subtract the ‘quark mass’ contribution



# Hadron Spectroscopy - role of the glue

- Subtract the ‘quark mass’ contribution

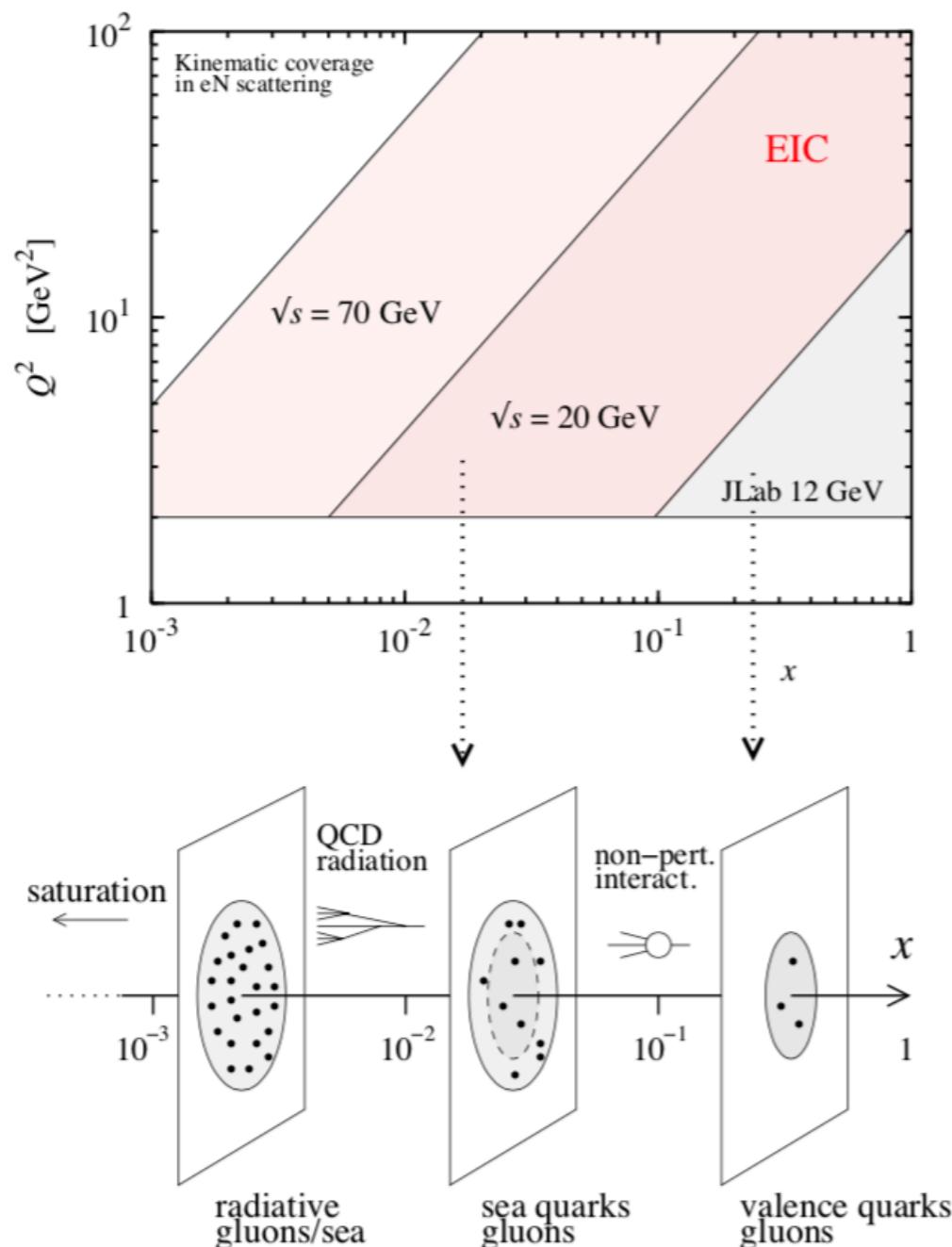


- Common energy scale of gluonic excitation  $\sim 1.3 \text{ GeV}$  Relevant for EIC?
- Strong overlap of gluonic constructions with ground-states

# Electron Ion Collider(s)

Efforts in US (Jefferson Lab & Brookhaven Lab), EicC (IMP, Lanzhou), and eLHC

Particular interest is role of glue in QCD



Partonic structure of gluonic states complementary to such a program

first baby-steps PRD98 014511 (2018)

# Distillation is powerful

---

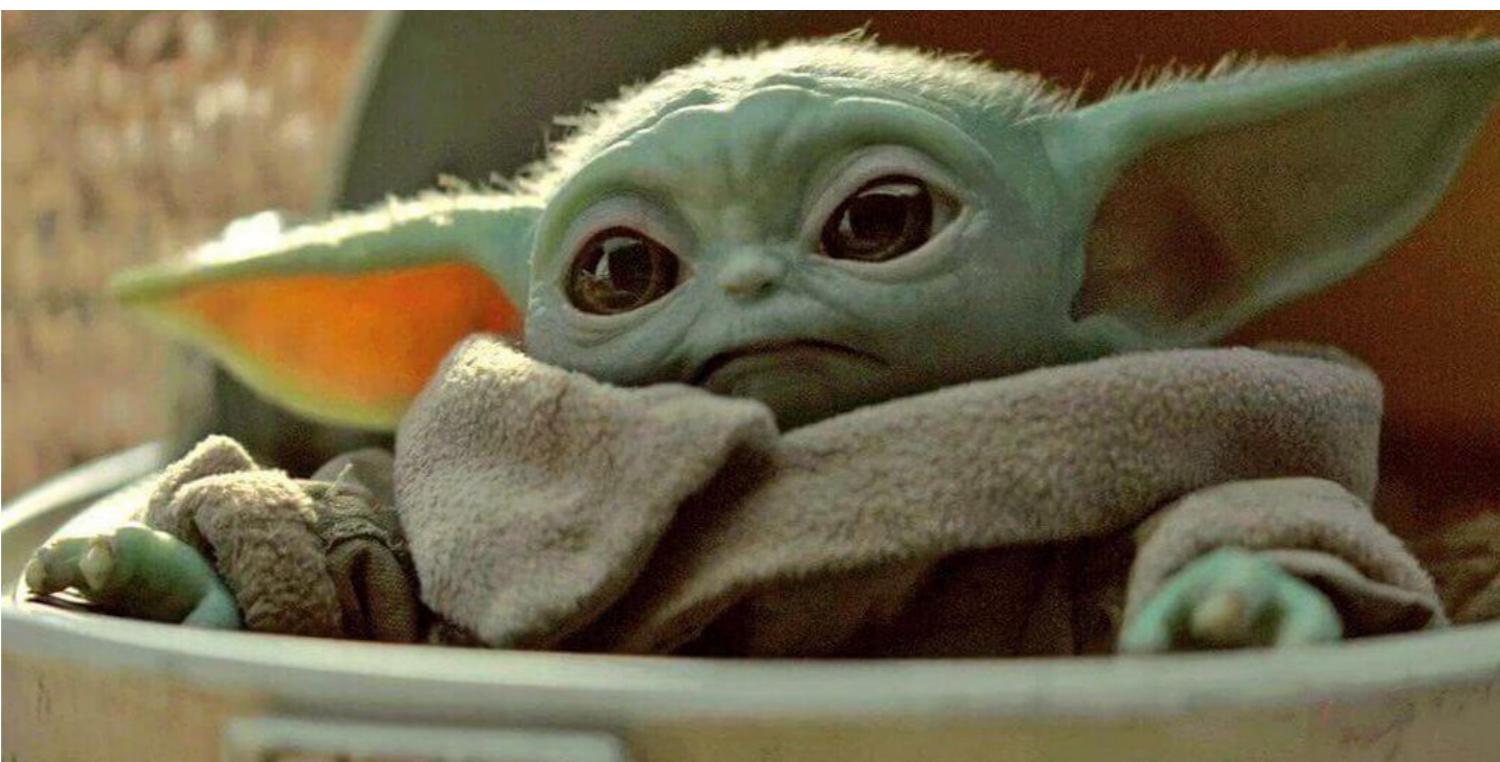
- Early results suggest rich spectrum of mesons & baryons - including hybrids
  - @JLab → new expt. proposals
    - GlueX: kaon PID upgrade, CLAS12: hybrid baryons

# Distillation is powerful

---

- Early results suggest rich spectrum of mesons & baryons - including hybrids
  - @JLab → new expt. proposals
    - GlueX: kaon PID upgrade, CLAS12: hybrid baryons

But this is a program in its infancy



# Distillation is powerful

---

- Early results suggest rich spectrum of mesons & baryons - including hybrids
  - @JLab → new expt. proposals
    - GlueX: kaon PID upgrade, CLAS12: hybrid baryons
- Near term:
  - See projects use range of light/charm-quark masses → poles & couplings
  - Estimates of branching fractions useful for expt. analysis
  - 3-body
- Opportunities & challenges
  - XYZ-s &  $P_c$
  - Hybrid hadrons
  - Quantitative understanding of resonant structures
- Long term:
  - Near light-cone formalism
  - Envision merging spectroscopy and structure projects
    - Understanding role of gluonic structures