



NCN

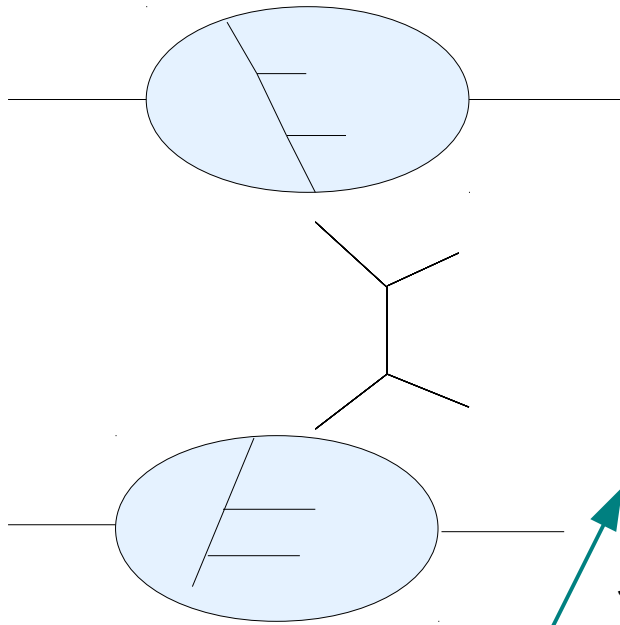
Nuclear unintegrated pdfs

Krzysztof Kutak

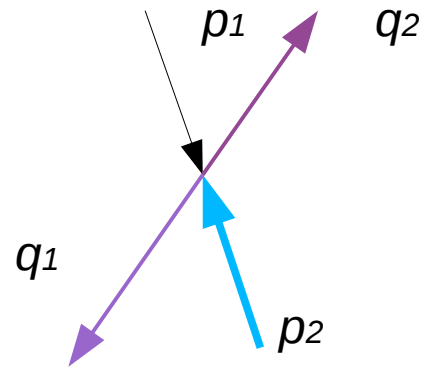


*Based on
E. Blanco, A. van Hameren, H. Jung, A. Kusina
Phys. Rev. D 100, 054023 (2019)*

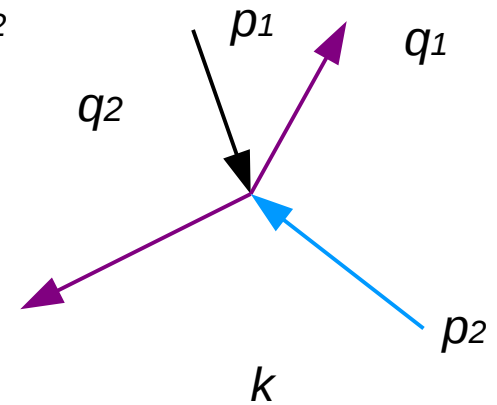
QCD at high energies – k_t factorization



Strongly decreasing
Longitudinal momentum
fractions of off-shell partons



$$p_1 + p_2 = q_1 + q_2$$



$$p_1 + p_2 = q_1 + q_2 + k$$

$$\frac{d\sigma}{dPS} \propto \mathcal{F}_{a^*}(x_1, k_{\perp 1}) \otimes \hat{\sigma}_{ab \rightarrow cd}(x_1, x_2) \otimes \mathcal{F}_{b^*}(x_2, k_{\perp 2})$$

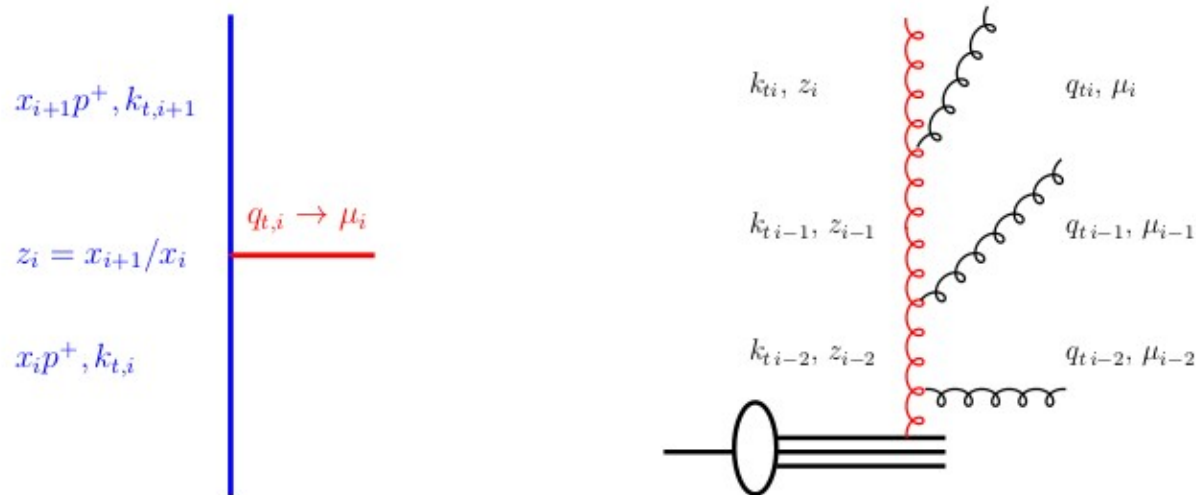
Ciafaloni, Catani, Hautman '93
Collins, Ellis '93

New helicity based methods for ME
Kotko, K.K, van Hameren, '12

Monte Carlo – parton branching method

The idea: construct such parton shower that gives also TMD dependent parton density.
On integrated level the pdf obeys DGLAP equation.

A. Bermudez Martinez, P. Connor, H. Jung, A. Lelek, R. Žlebčík, F. Hautmann, V. Radescu *Phys.Rev. D99* (2019) no.7, 074008



$$\mathcal{A}_a(x, \mathbf{k}, \mu^2) = \Delta_a(\mu^2) \mathcal{A}_a(x, \mathbf{k}, \mu_0^2) + \sum_b \int \frac{d^2 \mathbf{q}'}{\pi \mathbf{q}'^2} \frac{\Delta_a(\mu^2)}{\Delta_a(\mathbf{q}'^2)} \Theta(\mu^2 - \mathbf{q}'^2) \Theta(\mathbf{q}'^2 - \mu_0^2) \\ \times \int_x^{z_M} \frac{dz}{z} P_{ab}^{(R)}(\alpha_s, z) \mathcal{A}_b\left(\frac{x}{z}, \mathbf{k} + (1-z)\mathbf{q}', \mathbf{q}'^2\right)$$

The method provides first consistent and complete set of TMD's which are applicable in Monte Carlo simulations and can be generalized to account for small x effects at least in linear regime. Applies in the regime above saturation scale.

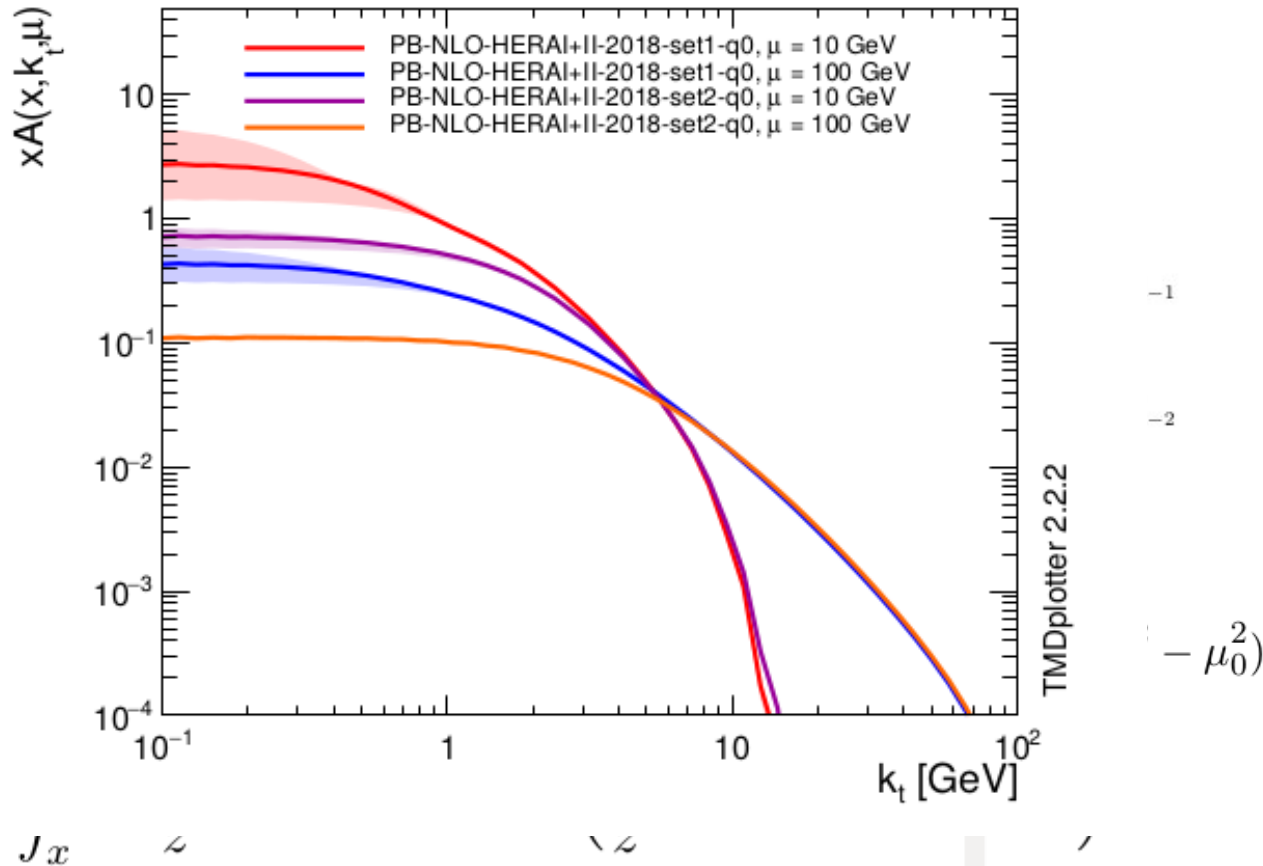
Monte Carlo – parton branching method

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A. Bermudez Martinez,

arXiv:1907.07400 [hep-ph]

gluon, $x = 0.01$



$$\mathcal{A}_a(x, \mathbf{k}, \mu^2) =$$

\times

J_x

\sim

\sim

TMDplotter 2.2.2

-1

-2

$-\mu_0^2$

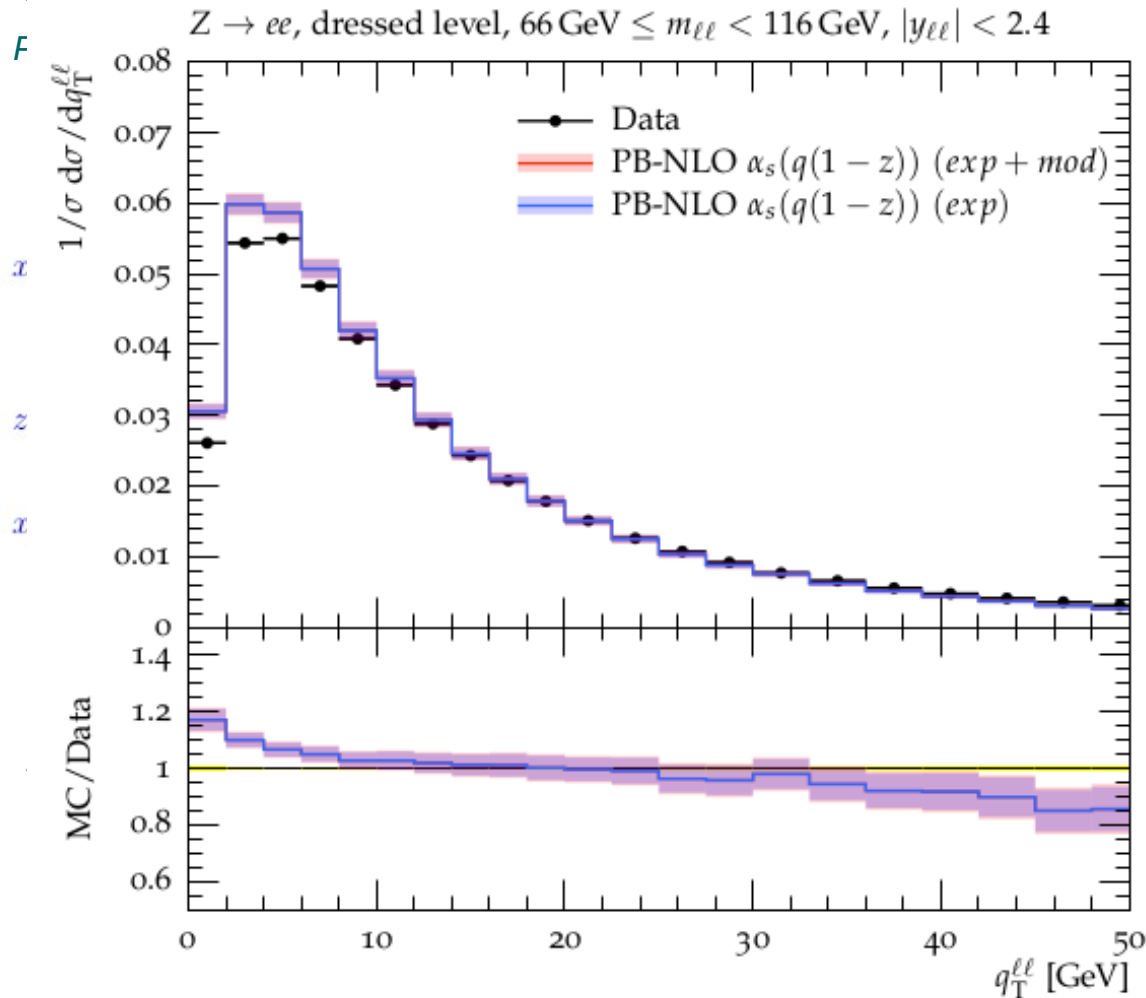
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Monte Carlo – parton branching method

The idea: construct such parton shower that gives also TMD dependent parton density.
On integrated level

A. Bermudez Martinez, F

D99 (2019) no.7, 074008



$$\mathcal{A}_a(x, \mathbf{k}, \mu^2) =$$

$$\cdot \mu_0^2$$

The method provides first consistent and complete set of TMD's which are applicable in Monte Carlo simulations and can be generalized to account for small x effects at least in linear regime. Applies in the regime above saturation scale.

Parton branching method for lead unintegrated pdfs

$$x \mathcal{A}_a^{\text{Pb}}(x, k_t^2, \mu^2) = \int dx' \mathcal{A}_{0,b}^{\text{Pb}}(x', k_{t,0}^2, \mu_0^2) \frac{x}{x'} \mathcal{K}_{ba} \left(\frac{x}{x'}, k_{t,0}^2, k_t^2, \mu_0^2, \mu^2 \right)$$

$$\mathcal{A}_{0,b}^{\text{Pb}}(x, k_{t,0}^2, \mu_0^2) = f_{0,b}^{\text{Pb}}(x, \mu_0^2) \cdot \exp(-|k_{t,0}^2|/\sigma^2)$$

$$\sigma^2 = q_0^2/2$$

$$q_0 = 0.5 \text{ GeV}$$

Momentum conservation → exact kinematics

The splitting function → NLO DGLAP

The shower is consistent with the transversal momentum dependent distribution it uses

Available collinear nuclear PDFs

Multiplicative correction factor

$$f_i^{p/A}(x_N, \mu_0) = R_i(x_N, \mu_0, A) f_i^{\text{free proton}}(x_N, \mu_0)$$

HKN: Hirai, Kumano, Nagai [[PRC 76, 065207 \(2007\)](#)]

DSSZ: de Florian, Sassot, Stratmann, Zurita [[PRD 85, 074028 \(2012\)](#)]

EPS09: Eskola, Paukkunen, Salgado [[JHEP 04 \(2009\) 065](#)]

EPPS16: Eskola, Paakkinen, Paukkunen, Salgado [[EPJC 77 \(2017\) 163](#)]

KT16: Khanpour, Tehrani [[PRD 93, 014026 \(2016\)](#)]

Native nuclear PDFs

$$f_i^{p/A}(x_N, \mu_0) = f_i(x_N, A, \mu_0)$$

$$f_i(x_N, A = 1, \mu_0) \equiv f_i^{\text{free proton}}(x_N, \mu_0)$$

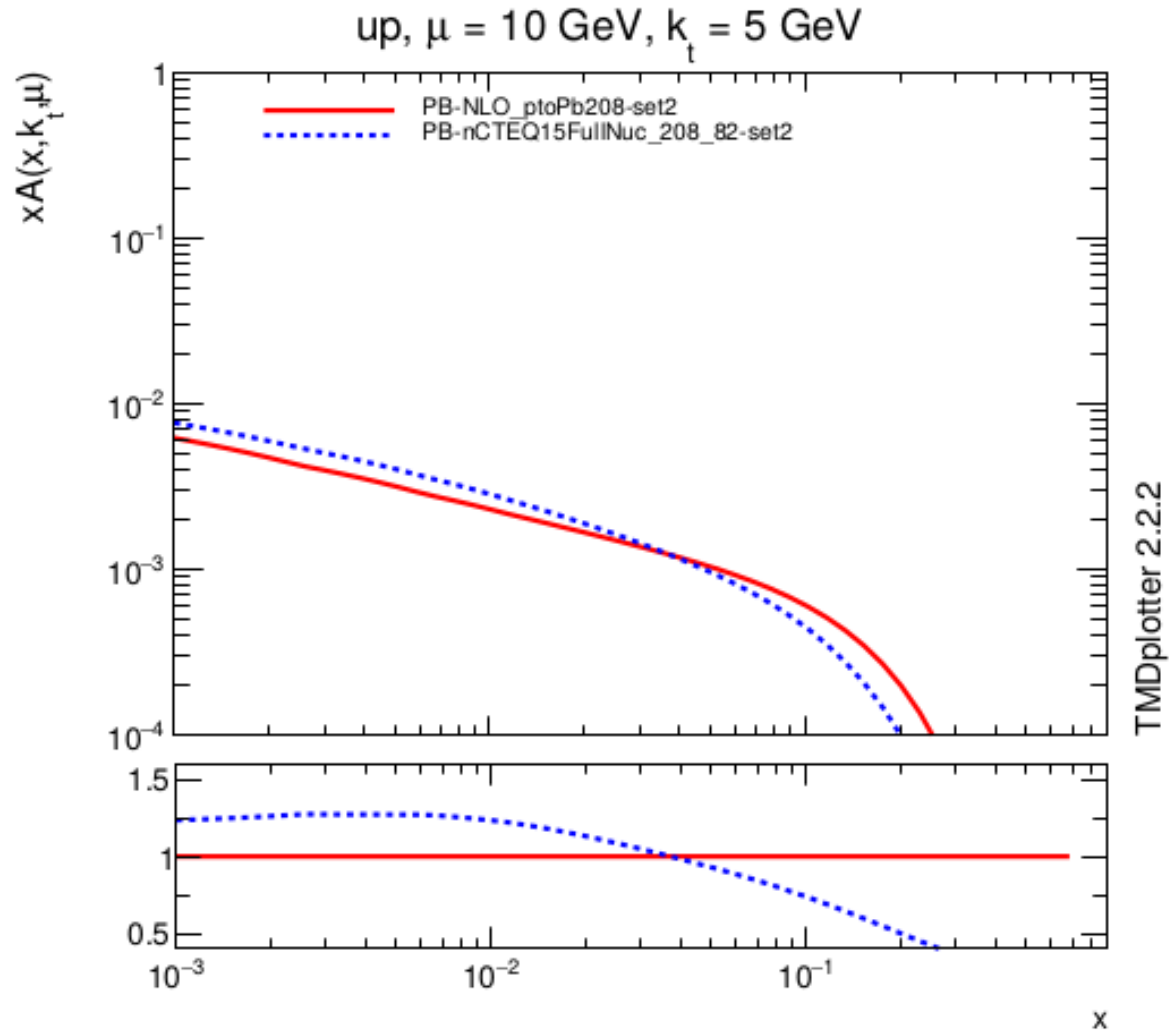
nCTEQ15: Kovarik, Kusina, Jezo, Clark, Keppel, Lyonnet, Morfin, Olness, Owens, Schienbein, Yu [[PRD 93, 085037 \(2016\)](#), [arXiv:1509.00792](#)]

NNPDF Khalek, Ethier, Rojo [arXiv:1904.00018](#)

WHV Marina Walt, Ilkka Helenius, Werner Vogelsang [arXiv:1908.03355](#)

See also talks by
V. Guzey
M. Zurita

Example of our TMD proton vs. lead

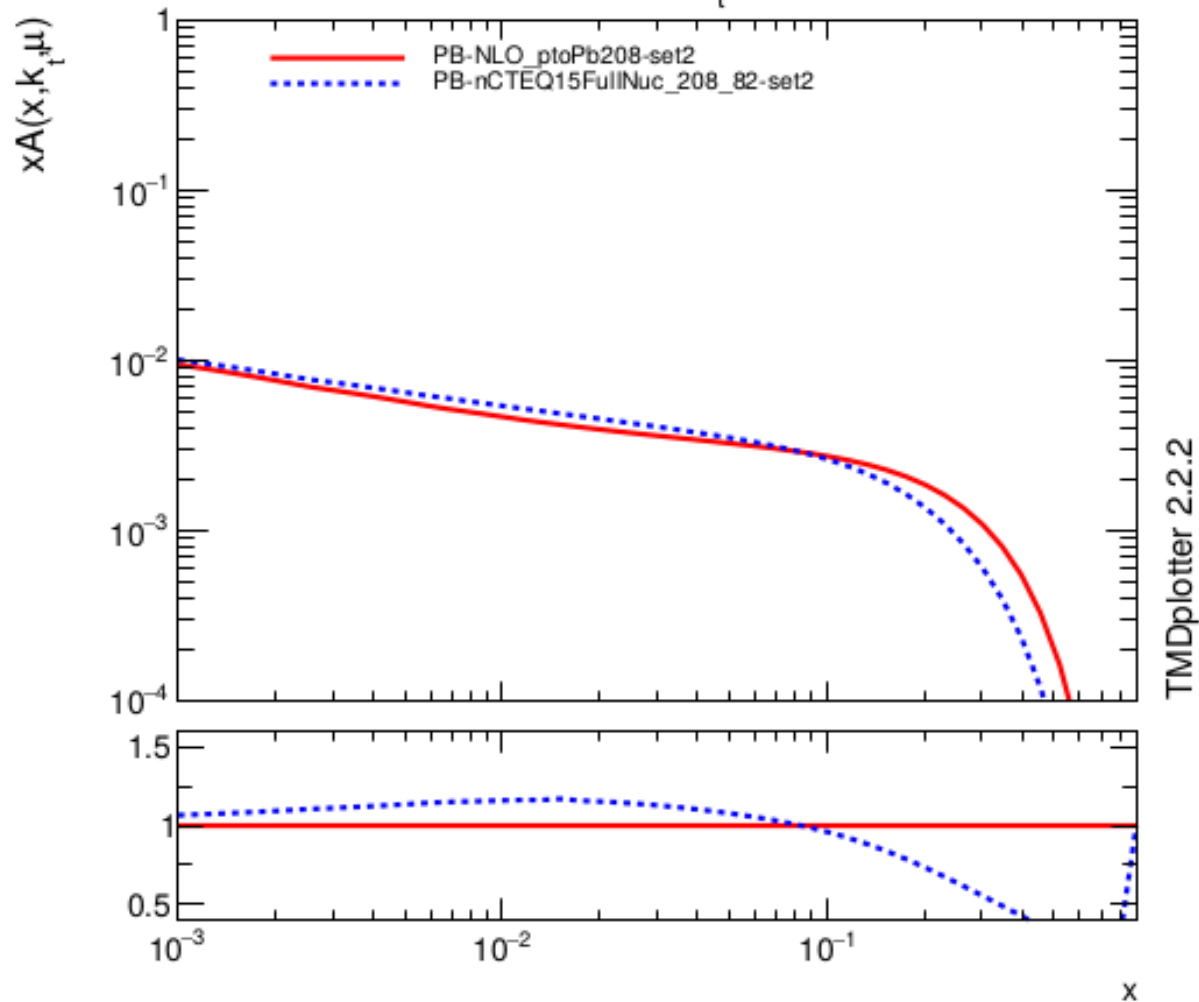


To see how the nuclear effects enter we constructed a test lead TMD from free protons. Neutron distribution is obtained assuming isospin symmetry

$$f^{Pb} = 82/208 f^p + (208-82)/208 f^n$$

Example of our TMD proton vs. lead

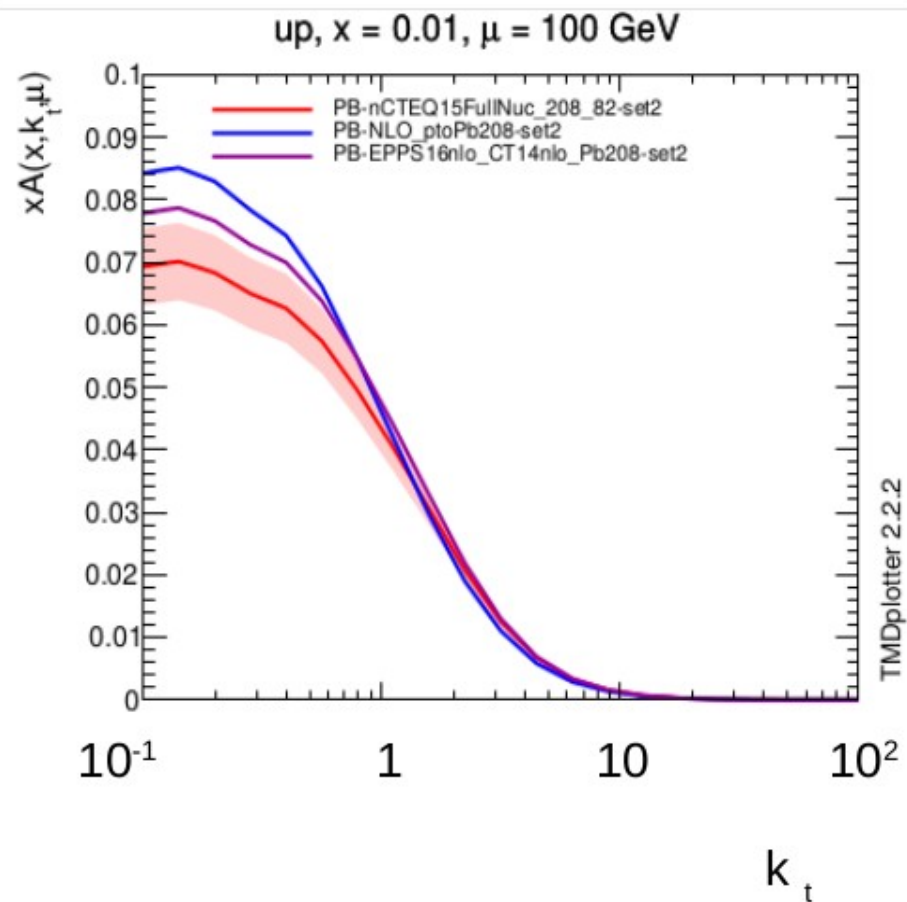
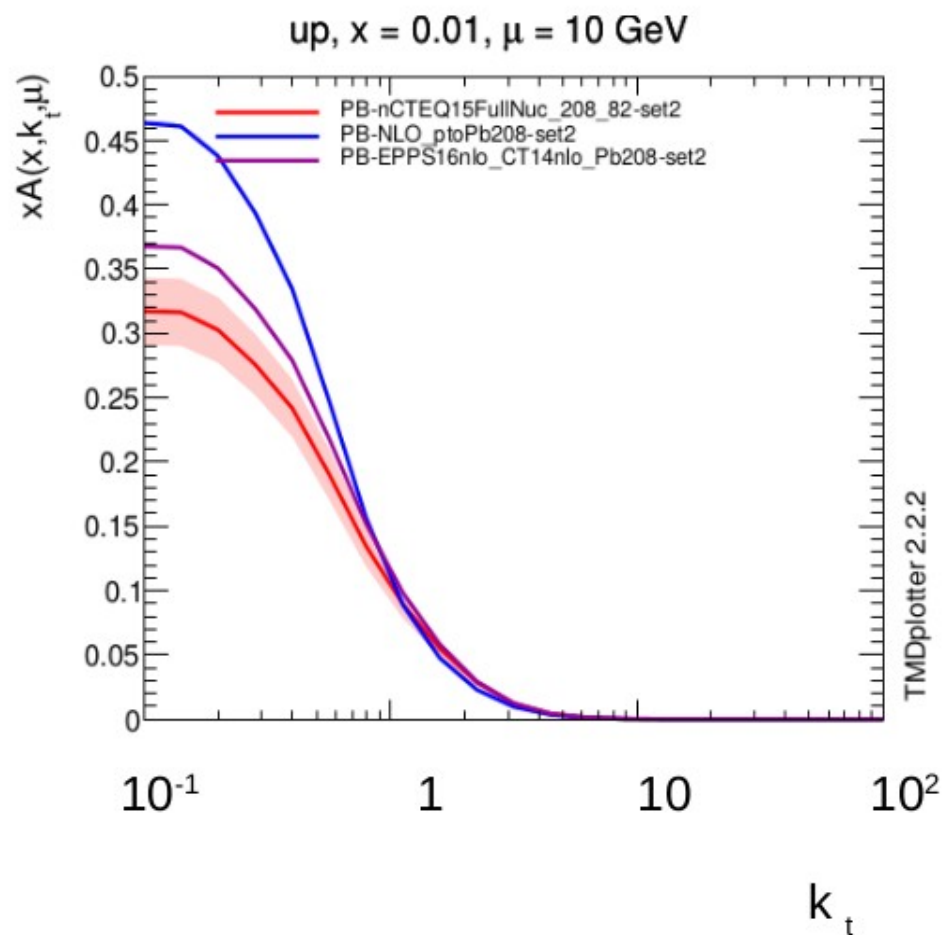
up, $\mu = 100$ GeV, $k_t = 5$ GeV



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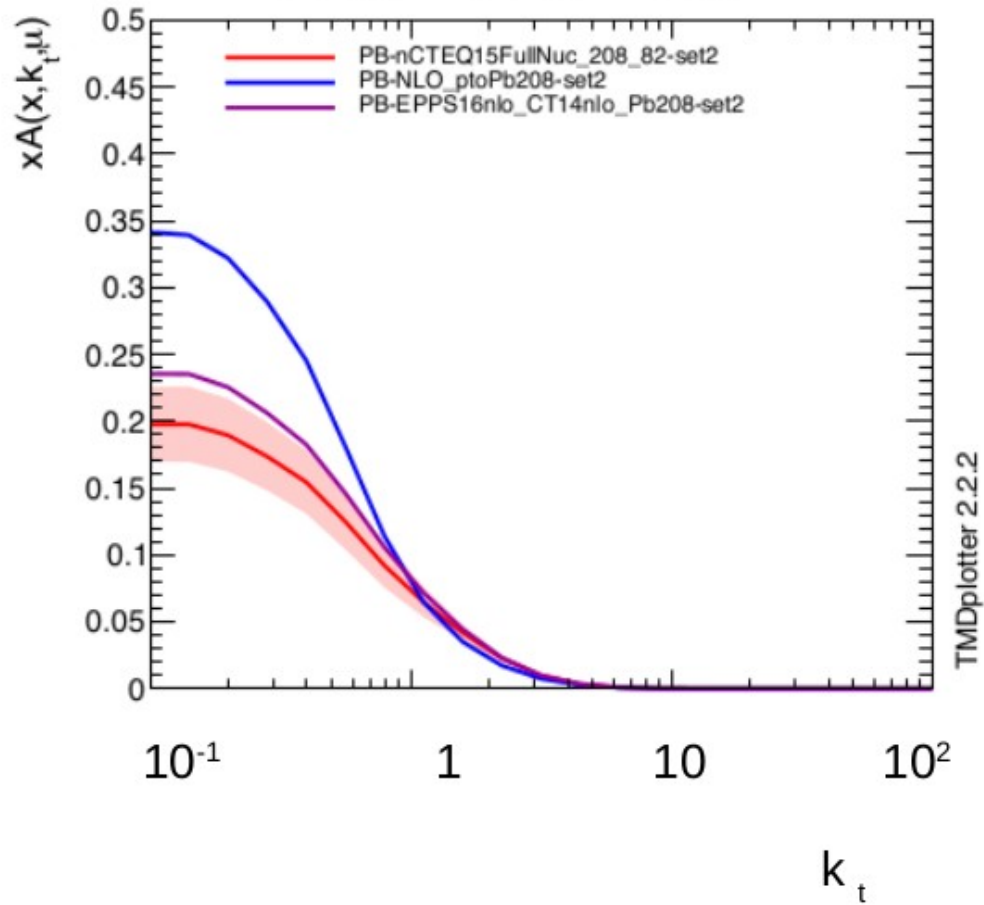
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Transverse momentum dependence

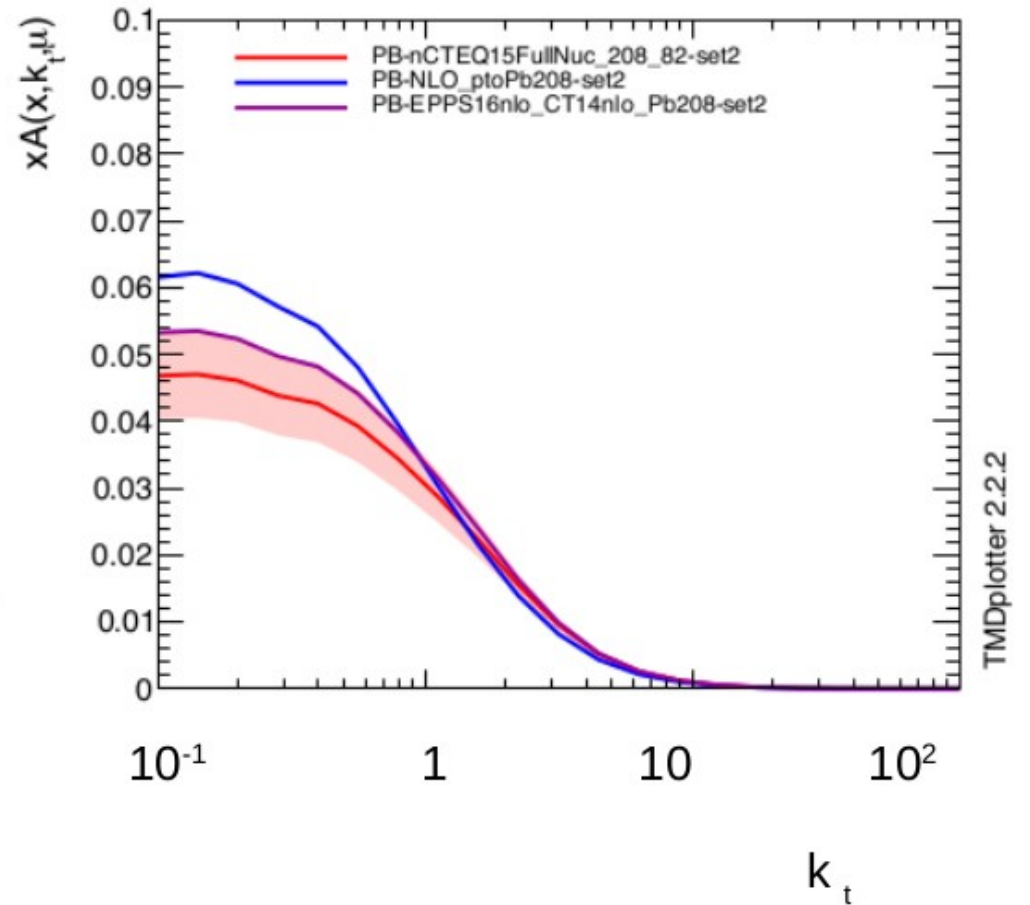


Transverse momentum dependence

anti-up, $x = 0.01$, $\mu = 10$ GeV

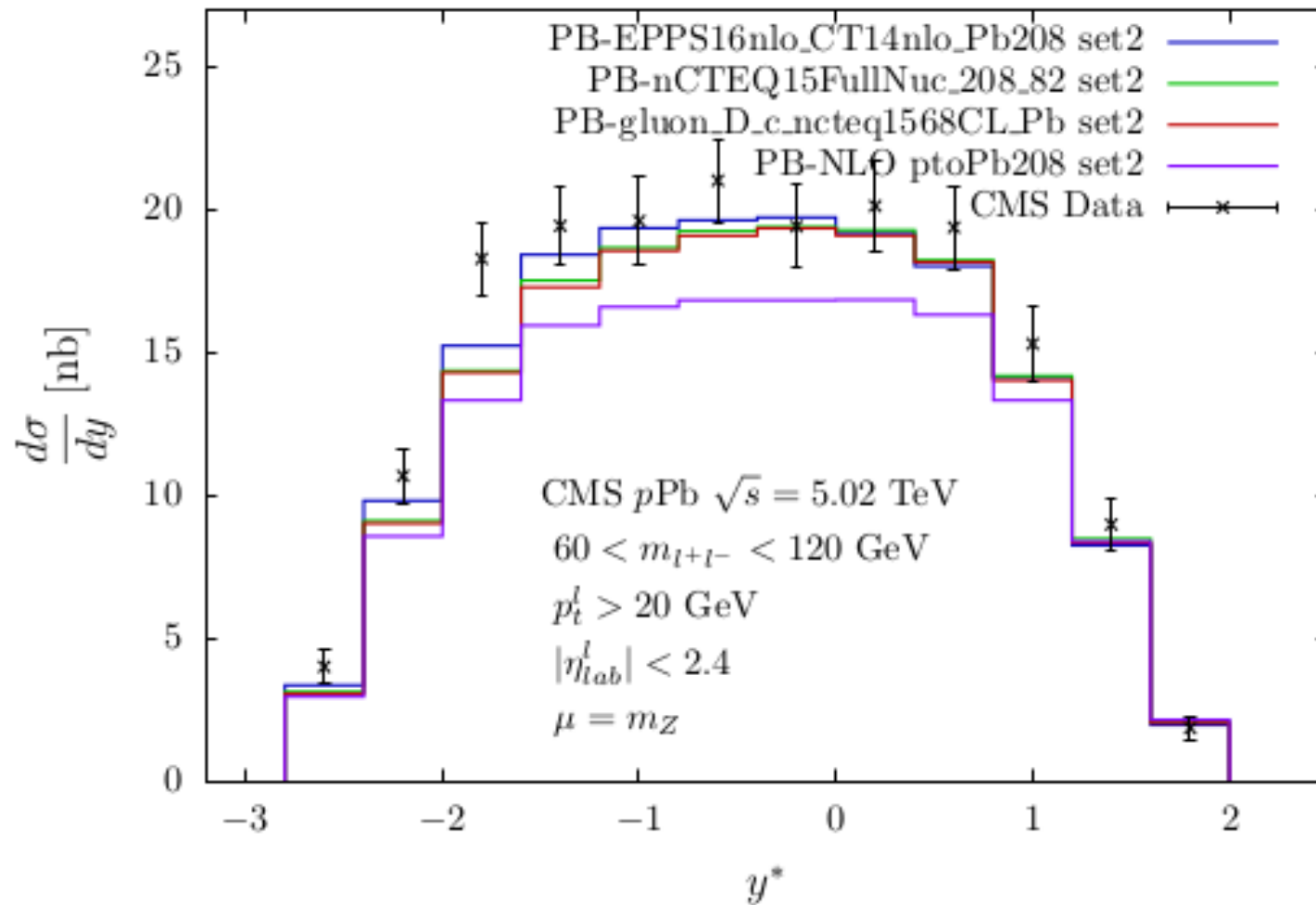


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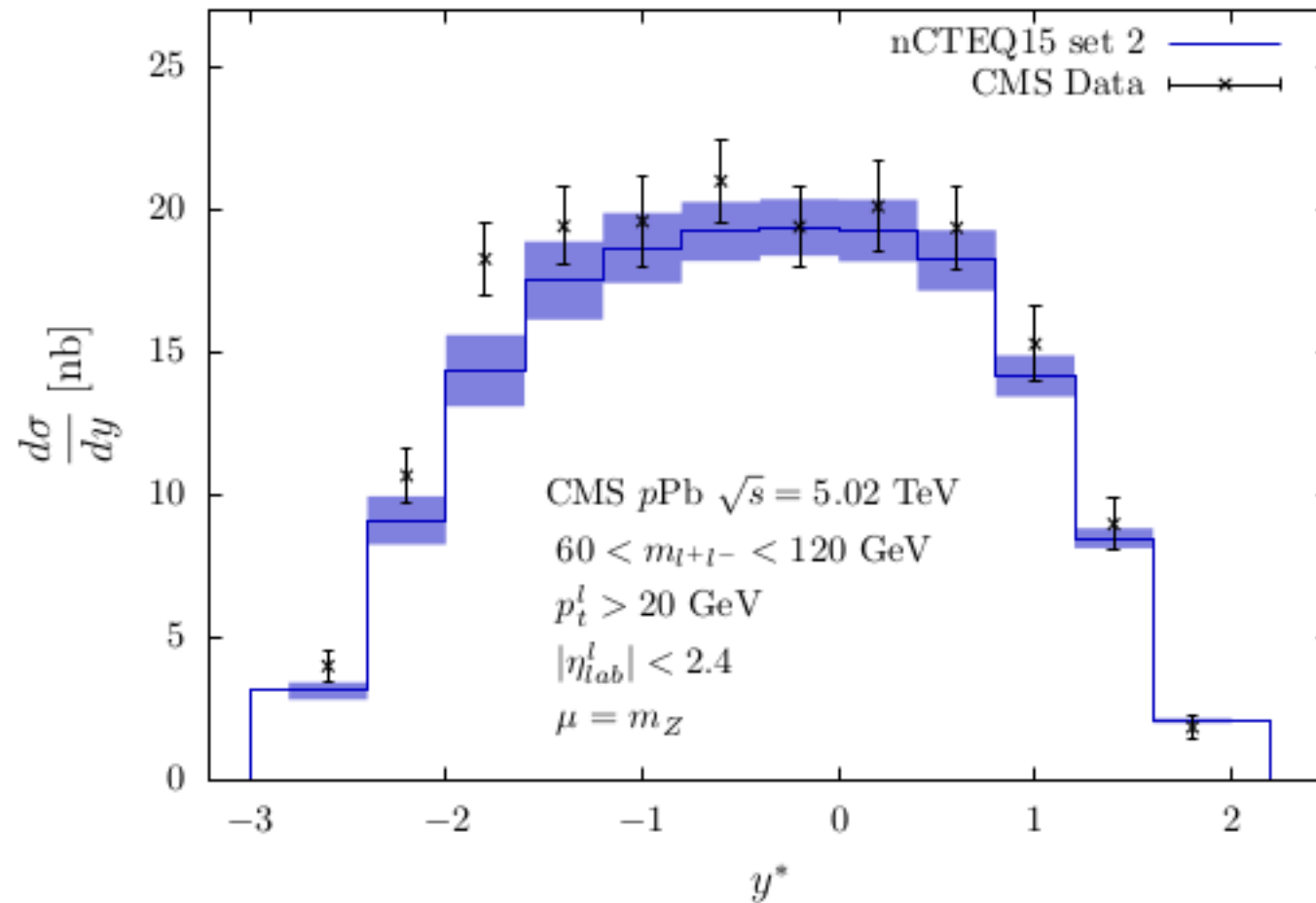
Drell-Yan – rapidity distribution

$$p + Pb \rightarrow Z^* \rightarrow \mu^+ + \mu^-$$



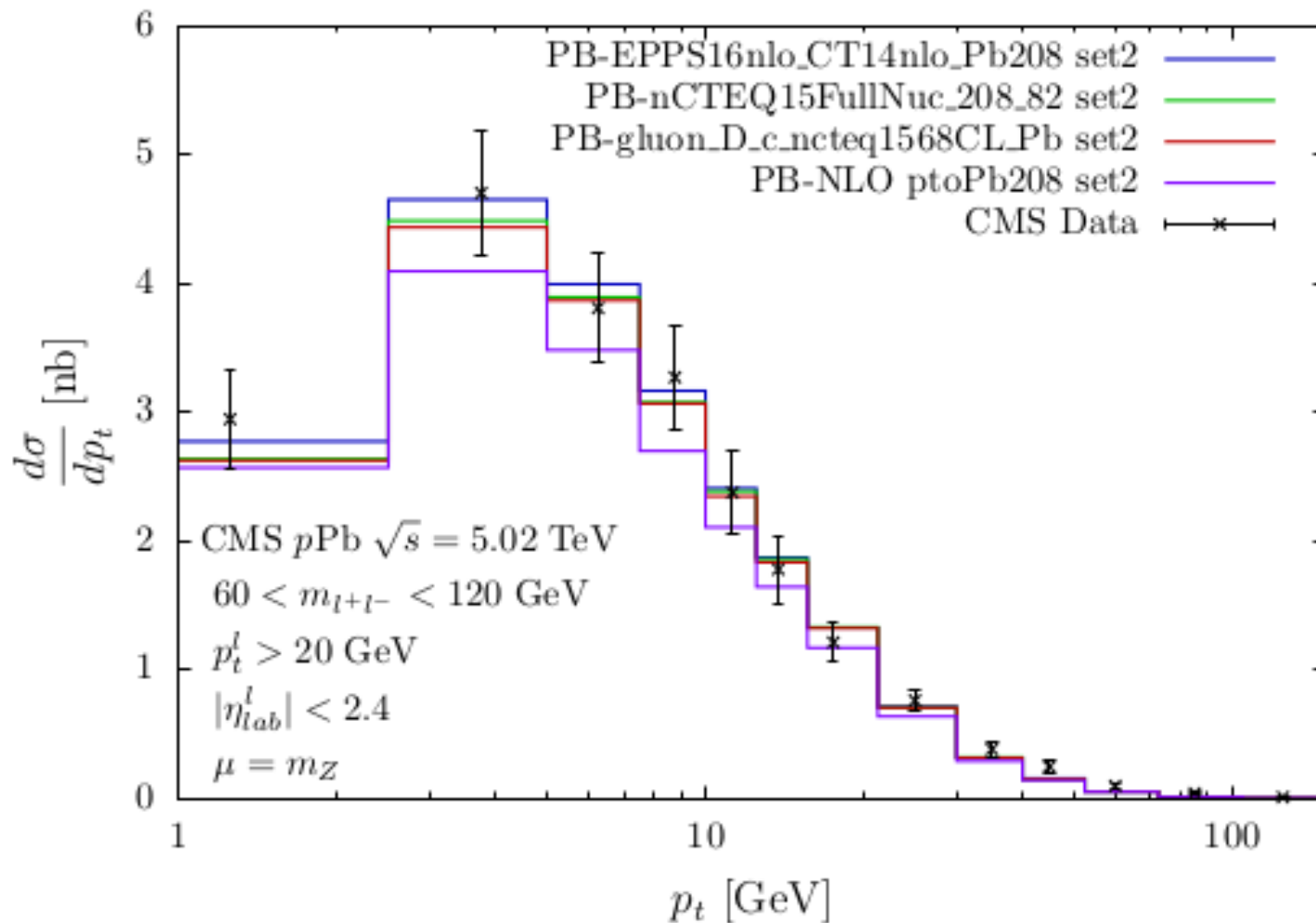
Drell-Yan – rapidity distribution – with PDF uncertainty

$$p + Pb \rightarrow Z^* \rightarrow \mu^+ + \mu^-$$



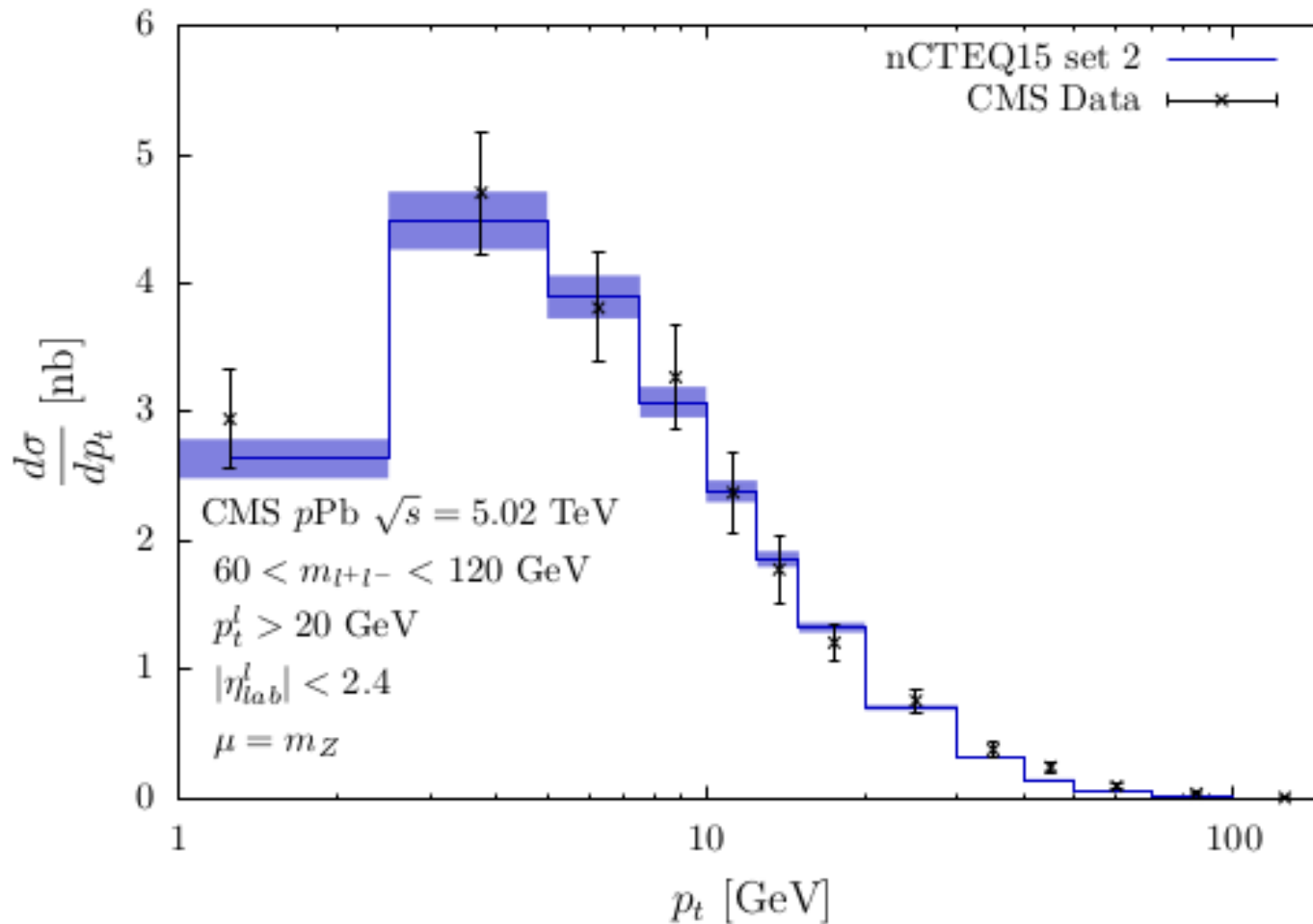
Drell-Yan – distribution in p_t

$$p + Pb \rightarrow Z^* \rightarrow \mu^+ + \mu^-$$



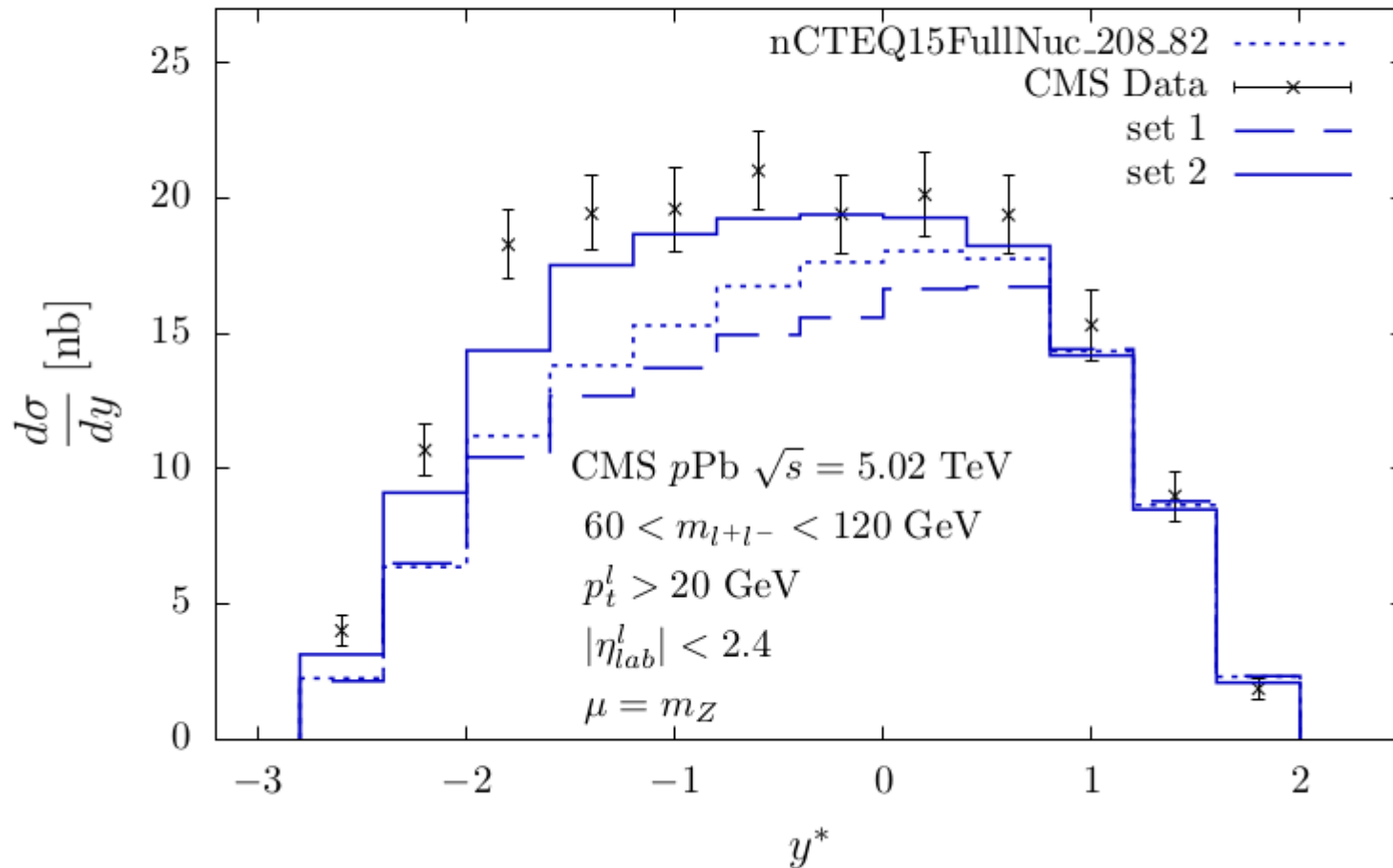
Drell-Yan – distribution in p_t – with PDF uncertainty

$$p + Pb \rightarrow Z^* \rightarrow \mu^+ + \mu^-$$



Comparison to collinear factorization based calculation

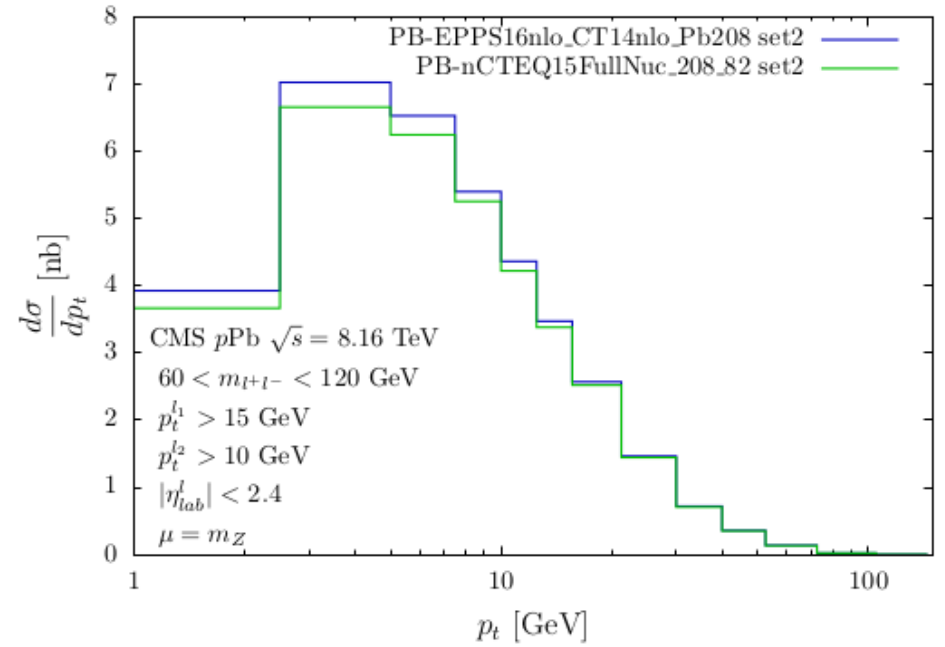
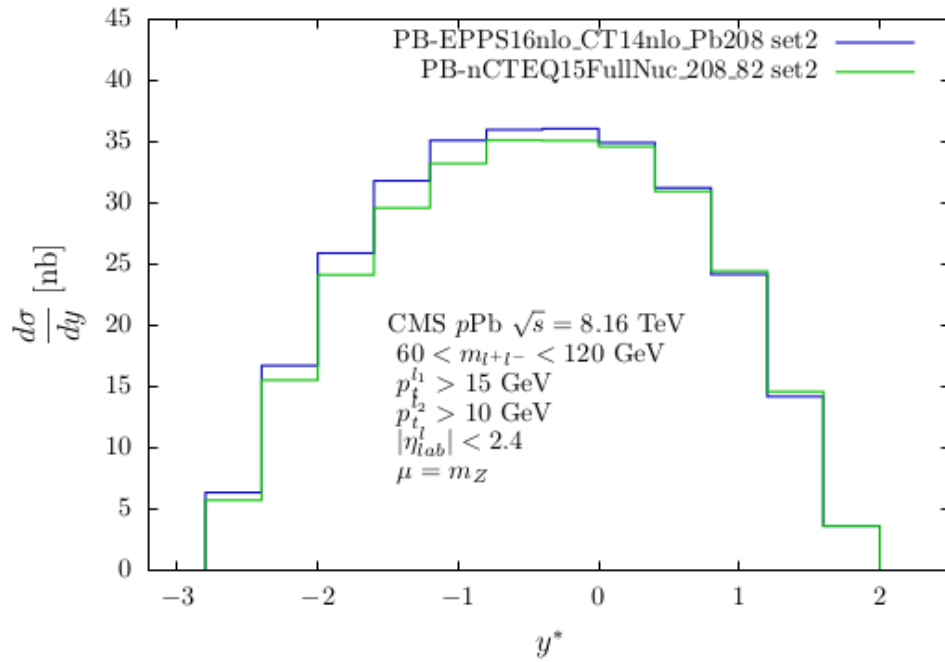
PB-nCTEQ15FullNuc_208_82



Set 1 uses evolution variable as a scale

Set 2 uses $|q_t^2|$ as evolution variable i.e. momentum of emitted gluon

Predictions for 8.16 TeV



Conclusions and outlook

New TMD introduced: i.e. nTMD's using PB method

The set does well describing Drell-Yan data

Application to other processes involving heavy ions

New fit of the initial distribution using off-shell coefficient function