## Shear viscosity to entropy ratio in A+A collisions at NICA energies



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## Motivation



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- A.Muronga. PRC 69, 044901 (2004)
- L.Csernai, J.Kapusta, L.McLerran. PRL 97, 152303 (2006)
- P.Romatschke, U.Romatschke. PRL 99, 172301 (2007)
- S.Plumari et al. PRC 86, 054902 (2012)
- ALICE collaboration, CERNCOURIER (14.10.2016)
- J.Rose et al. PRC 97, 055204 (2018)
taken from
R.Rapp, H.Hees. arXiv:0803.0901[hep-ph]

Green-Kubo: shear viscosity $\eta$ may be defined as:

$$
\eta\left(t_{0}\right)=\frac{1}{\hbar} \frac{V}{T} \int_{t_{0}}^{\infty} \mathrm{d} t\left\langle\pi(t) \pi\left(t_{0}\right)\right\rangle_{t}=\frac{\tau}{\hbar} \frac{V}{T}\left\langle\pi\left(t_{0}\right) \pi\left(t_{0}\right)\right\rangle
$$

where

$$
\begin{aligned}
\left\langle\pi(t) \pi\left(t_{0}\right)\right\rangle_{t} & =\frac{1}{3} \sum_{\substack{i, j=1 \\
i \neq j}}^{3} \lim _{\max \rightarrow \infty} \frac{1}{t_{\max }-t_{0}} \int_{t_{0}}^{t_{\max }} \mathrm{d} t^{\prime} \pi^{i j}\left(t+t^{\prime}\right) \pi^{i j}\left(t^{\prime}\right) \\
& =\left\langle\pi\left(t_{0}\right) \pi\left(t_{0}\right)\right\rangle \exp \left(-\frac{t-t_{0}}{\tau}\right)
\end{aligned}
$$

with

$$
\pi^{i j}(t)=\frac{1}{V} \sum_{\text {particles }} \frac{p^{i}(t) p^{j}(t)}{E(t)}
$$

$t_{0}$ : initial cut-off time to start with

## Model setup: cell calculations

- UrQMD calculations, central $\mathrm{Au}+\mathrm{Au}$ collisions at energies $E \in[10,20,30,40] \mathrm{AGeV}$ of the projectile, 51200 events per each
- central cell $5 \times 5 \times 5 \mathrm{fm}^{3} \Rightarrow\left\{\varepsilon, \rho_{\mathrm{B}}, \rho_{\mathrm{S}}\right\}$ at times $t_{\text {cell }}=1 \div 20 \mathrm{fm} / \mathrm{c}$
- statistical model $(\mathrm{SM}):\left\{\varepsilon, \rho_{\mathrm{B}}, \rho_{\mathrm{S}}\right\} \Rightarrow\left\{T, s, \mu_{\mathrm{B}}, \mu_{\mathrm{S}}\right\}$


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## Equilibration in the Central Cell


$\mathbf{t}^{\text {cross }}=2 \mathbf{R} /\left(\gamma_{\mathrm{cm}} \beta_{\mathrm{cm}}\right)$


$$
\mathbf{t}^{\text {eq }} \geq \mathbf{t}^{\text {cross }}+\Delta z /\left(2 \beta_{c \mathrm{~cm}}\right)
$$

## Kinetic equilibrium:

Isotropy of velocity distributions
Isotropy of pressure
L.Bravina et al., PLB 434 (1998) 379; JPG 25 (1999) 351

$$
\frac{d N_{i}}{4 \pi p E d E}=\frac{V g_{i}}{(2 \pi \hbar)^{3}} \exp \left(\frac{\mu_{i}}{T}\right) \exp \left(-\frac{E_{i}}{T}\right)
$$

## Chemical equlibrium:

## Thermal equilibrium: Energy spectra of particles are described by Boltzmann distribution

Particle yields are reproduced by $\mathbf{S M}$ with the same values of $\left(T, \mu_{B}, \mu_{S}\right)$ :

$$
N_{i}=\frac{V g_{i}}{2 \pi^{2} \hbar^{3}} \int_{0}^{\infty} p^{2} d p \exp \left(\frac{\mu_{i}}{T}\right) \exp \left(-\frac{E_{i}}{T}\right)
$$

## Statistical model of ideal hadron gas

 input values$$
\begin{aligned}
\varepsilon^{\text {mic }} & =\frac{1}{V} \sum_{i} E_{i}^{\mathrm{SM}}\left(T, \mu_{\mathrm{B}}, \mu_{\mathrm{S}}\right), \\
\rho_{\mathrm{B}}^{\text {mic }} & =\frac{1}{V} \sum_{i} B_{i} \cdot N_{i}^{\mathrm{SM}}\left(T, \mu_{\mathrm{B}}, \mu_{\mathrm{S}}\right), \\
\rho_{\mathrm{S}}^{\text {mic }} & =\frac{1}{V} \sum_{i} S_{i} \cdot N_{i}^{\mathrm{SM}}\left(T, \mu_{\mathrm{B}}, \mu_{\mathrm{S}}\right) .
\end{aligned}
$$

Multiplicity

$$
\begin{aligned}
N_{i}^{\mathrm{SM}} & =\frac{V g_{i}}{2 \pi^{2} \hbar^{3}} \int_{0}^{\infty} p^{2} f\left(p, m_{i}\right) d p, \\
E_{i}^{\mathrm{SM}} & =\frac{V g_{i}}{2 \pi^{2} \hbar^{3}} \int_{0}^{\infty} p^{2} \sqrt{p^{2}+m_{i}^{2}} f\left(p, m_{i}\right) d p \\
P^{\mathrm{SM}} & =\sum_{i} \frac{g_{i}}{2 \pi^{2} \hbar^{3}} \int_{0}^{\infty} p^{2} \frac{p^{2}}{3\left(p^{2}+m_{i}^{2}\right)^{1 / 2}} f\left(p, m_{i}\right) d p \\
s^{\mathrm{SM}} & =-\sum_{i} \frac{g_{i}}{2 \pi^{2} \hbar^{3}} \int_{0}^{\infty} f\left(p, m_{i}\right)\left[\ln f\left(p, m_{i}\right)-1\right] p^{2} d p
\end{aligned}
$$

## Model setup: box calculations

- UrQMD box calculations at $\left\{\varepsilon, \rho_{\mathrm{B}}, \rho_{\mathrm{S}}\right\}$ for every energy and cell time $t_{\text {cell }}$ from cell calculations, 80 points in total, 12800 events per each
$\rho_{\mathrm{B}}$ is included as $N_{p}: N_{n}=1: 1$
$\rho_{\mathrm{S}}$ is included via kaons $\mathrm{K}^{-}$
box size: $10 \times 10 \times 10 \mathrm{fm}^{3}$
box boundaries: transparent
$\pi^{i j}(t)$ data extraction: $t=1 \div 1000 \mathrm{fm} / \mathrm{c}$ in box time, all types of hadrons are taken into account


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## Box with periodic boundary conditions



Initialization: (i) nucleons are uniformly distributed in a configuration space; (ii) Their momenta are uniformly distributed in a sphere with random radius and then rescaled to the desired energy density.

Test for equilibrium: particle yields and energy spectra

## Box: particle abundances


M.Belkacem et al., PRC 58, 1727 (1998)


L.Bravina et al., PRC 62, 064906 (2000)

Saturation of yields after a certain time. Strange hadrons are saturated longer than others (at not very high energy densities)

## Cell + SM





Dependence of $\varepsilon, \rho_{\mathrm{B}}, \rho_{\mathrm{S}}$ (from cell) and of $T, \mu_{\mathrm{B}}, \mu_{\mathrm{S}}$ (from SM ) on $t_{\text {cell }}$

## SM, Boltzmann entropy s




Dynamics of Boltzmann entropy density $s$ and of $s / \rho_{\mathrm{B}}$ in cell

## Results: $\left\langle\pi(t) \pi\left(t_{0}\right)\right\rangle_{t}$ at $E \in[10,20,30,40] \mathrm{AGeV}$



Time dependence of correlators $\left\langle\pi(t) \pi\left(t_{0}\right)\right\rangle_{t}$
$t_{0}=300 \mathrm{fm} / \mathrm{c}$
$t_{\text {cell }} \in\{1 \div 20\} \mathrm{fm} / \mathrm{c}$

## Results: $\left\langle\pi(t) \pi\left(t_{0}\right)\right\rangle_{t}$ at fixed $t_{\text {cell }}$



Time dependence of correlators $\left\langle\pi(t) \pi\left(t_{0}\right)\right\rangle_{t}$
Subplot: the same but at linear scale
$t_{0}=300 \mathrm{fm} / \mathrm{c}$
$t_{\text {cell }}=7 \mathrm{fm} / \mathrm{c}$

## Results: $\tau\left(t_{0}\right)$



Dependence of $\tau$ on $t_{0}$

## Results: $\tau$ from the fit



Dependence of $\tau_{\text {fit }}$ on $t_{0}$

## Results: Comparison of $\tau_{\text {int }}$ and $\tau_{\text {fit }}$



## Results: viscosity $\eta\left(t_{0}\right)$



Dependence of $\eta$ on $t_{0}$

## Results: viscosity $\eta\left(t_{\text {cell }}\right)$



Dynamics of $\eta$ in cell
All curves sit on the top of each other for $t_{\text {cell }} \geq 7 \mathrm{fm} / \mathrm{c}$

## Results: $\eta / s$



Dynamics of $\eta / s$ in cell
$\eta / s$ increases with time for $t_{\text {cell }} \geq 6 \mathrm{fm} / \mathrm{c}$ for all four energies Minimum - for 10 AGeV , corresponding to 4.5 GeV in c.m. frame

## Results: $\eta / s$



## Conclusions

- data from central cell of UrQMD calculations are used as input for SM to calculate temperature $T$ and Boltzmann entropy density $s$, and for UrQMD box calculations in order to estimate shear viscosity $\eta$
- box output data are taken within the range $200 \leq t_{0} \leq 800$ $\mathrm{fm} / \mathrm{c}$ because:
- values at $t_{0}<200 \mathrm{fm} / \mathrm{c}$ are distorted by the initial fluctuation in the box
- values at $t_{0}>800 \mathrm{fm} / \mathrm{c}$ may be disturbed by the analog of Brownian motion
- it is shown that for all four tested energies $\eta$ and $s$ in the cell drop with time
- ratios $\eta / \mathrm{s}$ reach minima about 0.3 at $t \approx 5 \mathrm{fm} / c$ for all energies. Then, the ratios rise to $1.0 \div 1.2$ at $t=20 \mathrm{fm} / \mathrm{c}$
- this increase is accompanied by the simultaneous rise of $\mu_{B}$ and drop of both $T$ and $\mu_{S}$ in the cell

