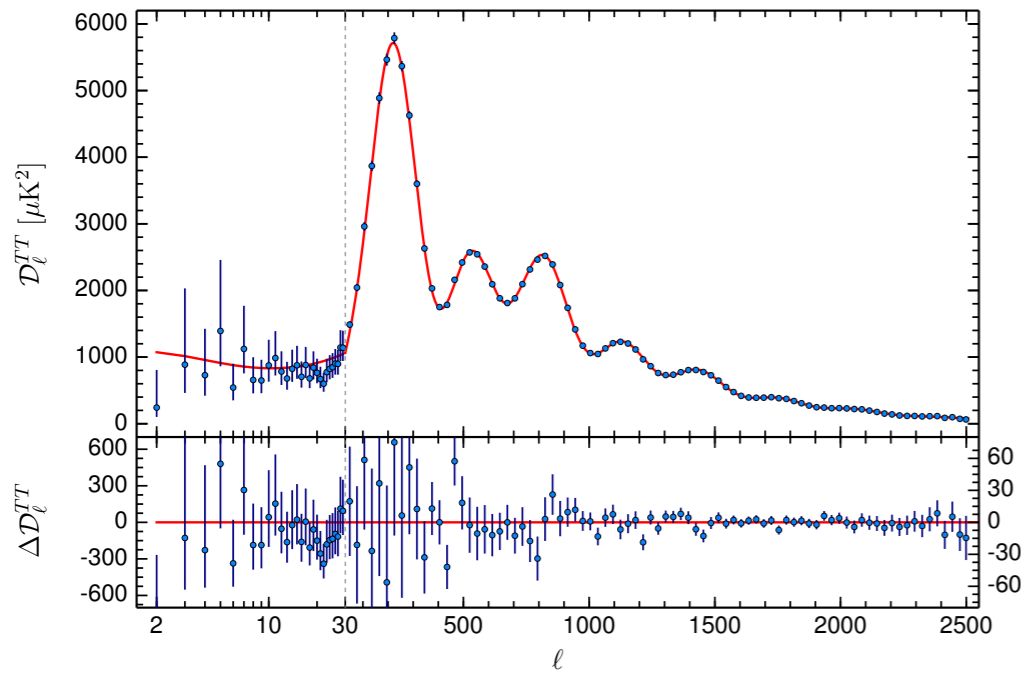
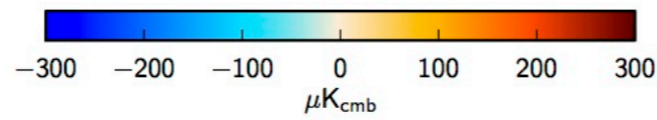
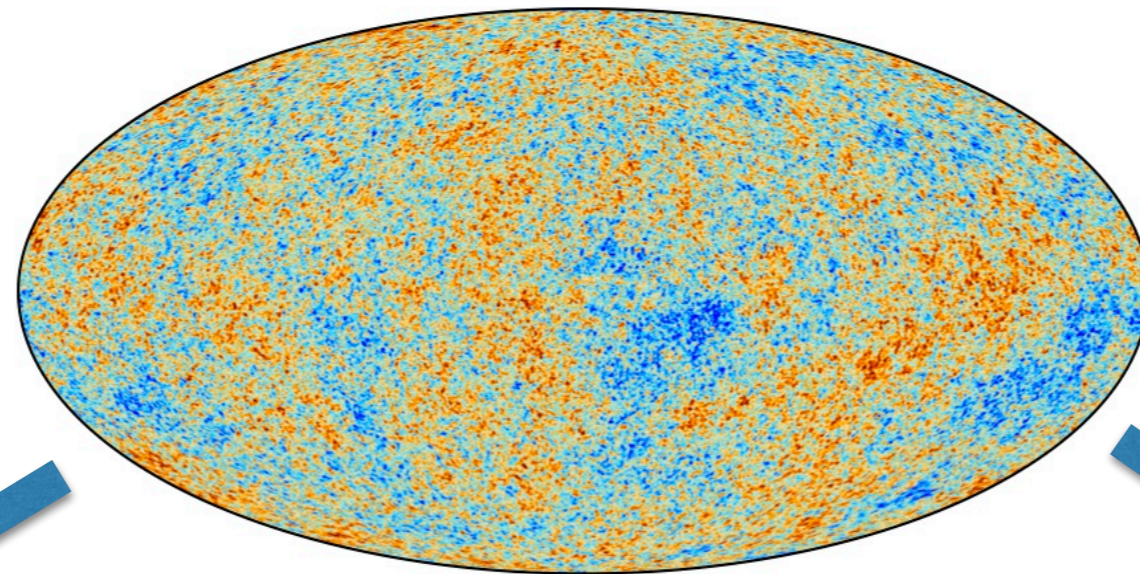


Review of recent dark matter experiments

Yoni Kahn (UIUC)

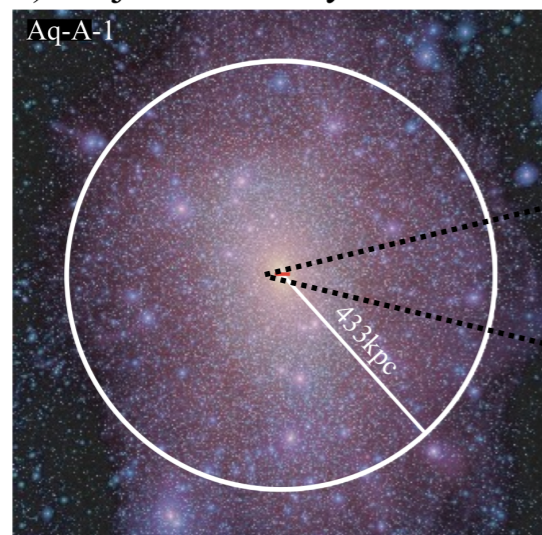
Dark matter (DM) exists!



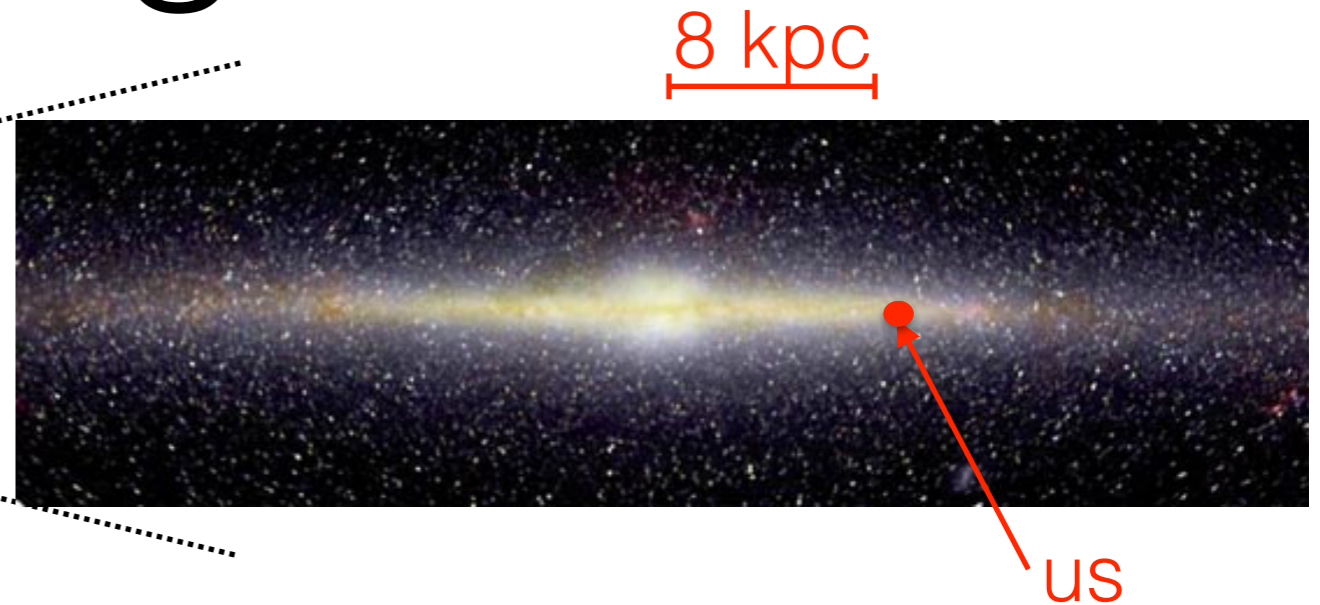
Parameter	[1] <i>Planck</i> TT+lowP
$\Omega_b h^2$	0.02222 ± 0.00023
$\Omega_c h^2$	0.1197 ± 0.0022
$100\theta_{MC}$	1.04085 ± 0.00047
τ	0.078 ± 0.019
$\ln(10^{10} A_s)$	3.089 ± 0.036
n_s	0.9655 ± 0.0062
H_0	67.31 ± 0.96
Ω_m	0.315 ± 0.013
σ_8	0.829 ± 0.014
$10^9 A_s e^{-2\tau}$	1.880 ± 0.014

We have never observed a dark matter particle.

DM in our neighborhood



Springel et al. 2008



Local measurements of stars tell us:

$$\cancel{m_{\text{DM}}} v_{\text{DM}}^2 \sim \cancel{m_{\text{DM}}} \frac{GM(< R)}{R}$$

$$v_{\text{DM}} \sim 10^{-3} c$$

$$(\sim v_{\odot} \sim v_{\text{esc}})$$

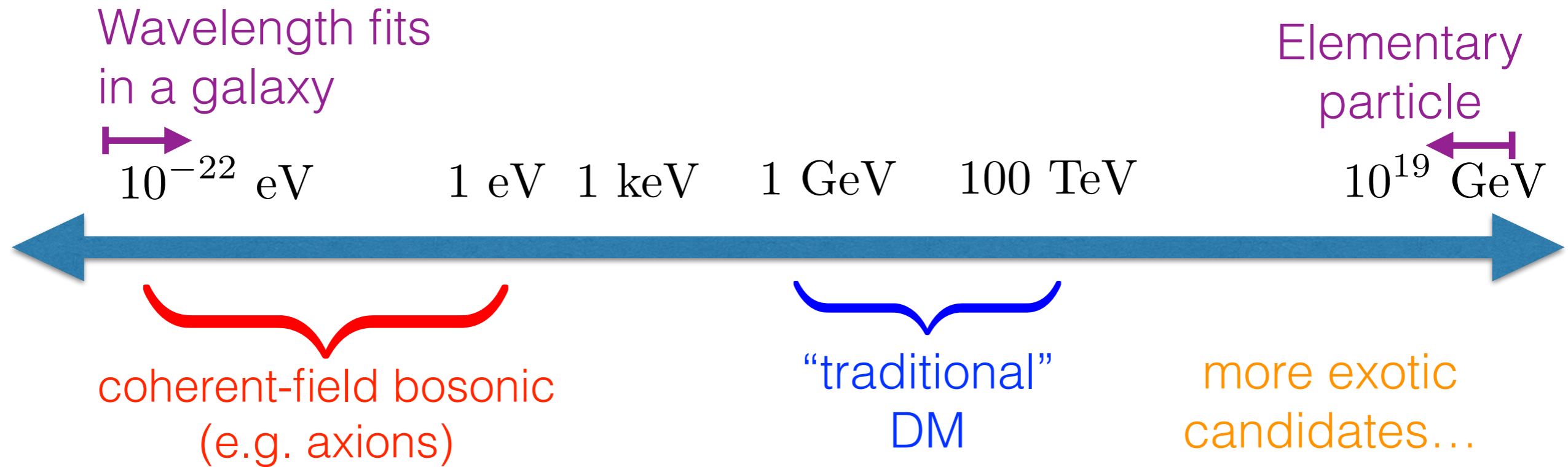
$$\rho_{\text{DM}} \sim 0.3 \text{ GeV}/\text{cm}^3$$

$$\rho_{\text{DM}} = m_{\text{DM}} \times n_{\text{DM}}$$

Few heavy particles,
or lots of light particles...

what is DM mass?

50 orders of magnitude!



50 orders of magnitude

Wavelength fits
in a galaxy

Elementary
particle



10^{-22} eV

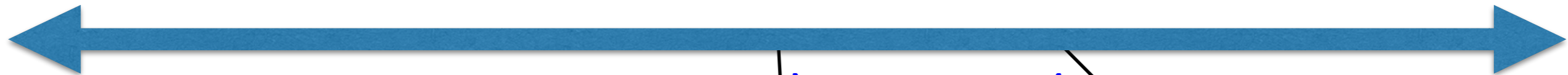
1 eV

1 keV

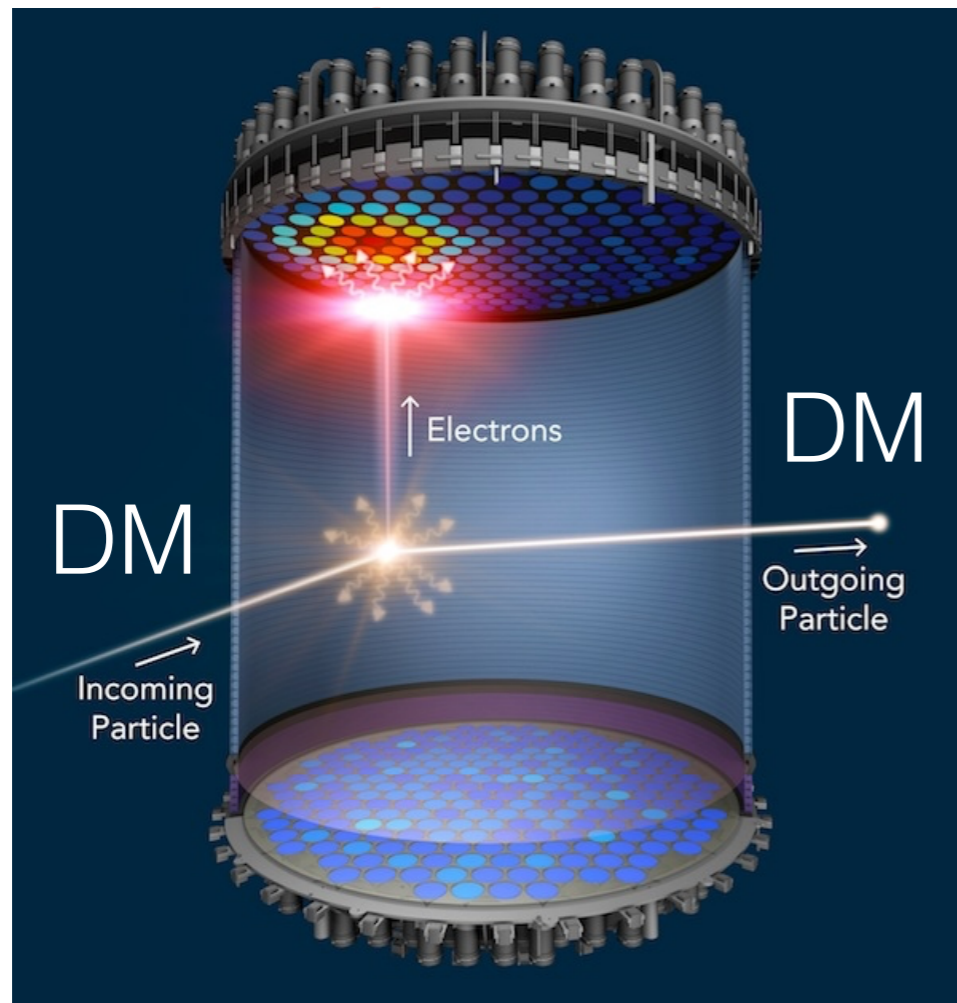
1 GeV

100 TeV

10^{19} GeV



“traditional”
DM



$$n_{\text{DM}} \lambda_{\text{dB}}^3 \ll 1$$
$$\text{KE}_{\text{DM}} = \frac{1}{2} m_{\text{DM}} v_{\text{DM}}^2 > 1 \text{ keV}$$

Rare collisions

50 orders of magnitude

Wavelength fits
in a galaxy



10^{-22} eV

1 eV

1 keV

1 GeV

100 TeV

Elementary
particle

10^{19} GeV



coherent-field bosonic
(e.g. axions)

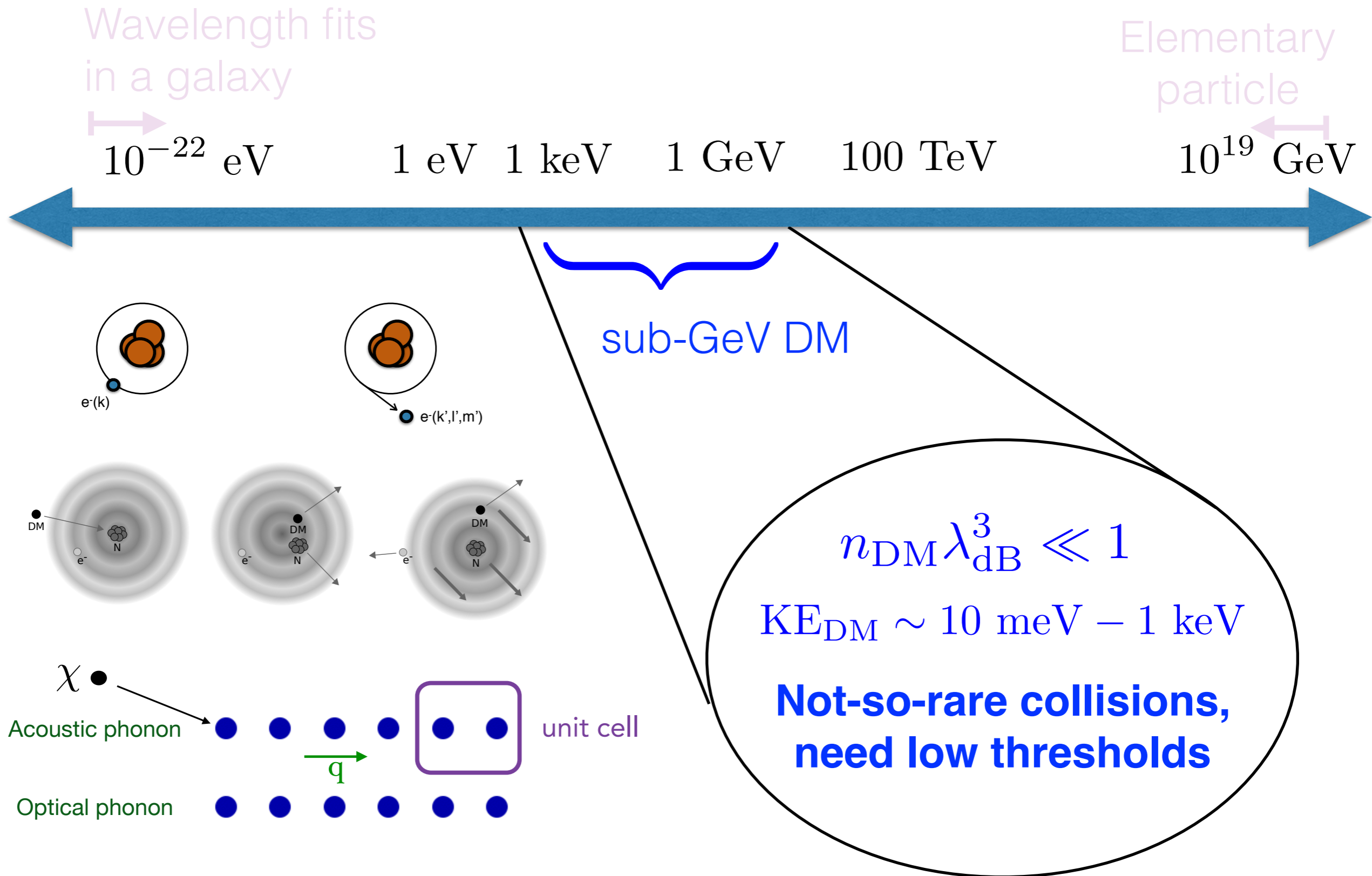
“traditional”
DM

$$n_{\text{DM}} \lambda_{\text{dB}}^3 \gg 1$$

$$\text{KE}_{\text{DM}} = \frac{1}{2} m_{\text{DM}} v_{\text{DM}}^2 \ll 1 \mu\text{eV}$$

Behaves as classical field

50 orders of magnitude



Axion DM experiments

Axion DM Theory

$$a(\mathbf{x}, t) = \frac{\sqrt{2\rho_{\text{DM}}}}{m_a} \cos(m_a t + \mathcal{O}(v_{\text{DM}})\mathbf{x})$$

amplitude set
by local DM density

oscillates at frequency
set by DM mass

e.g. $m_a = 10^{-9}$ eV
 $\lambda_{\text{Comp}} \sim \text{km}$
 $\tau_{\text{Comp}} \sim \mu\text{s}$

Local DM velocity \rightarrow Spatial coherence \rightarrow Temporal coherence

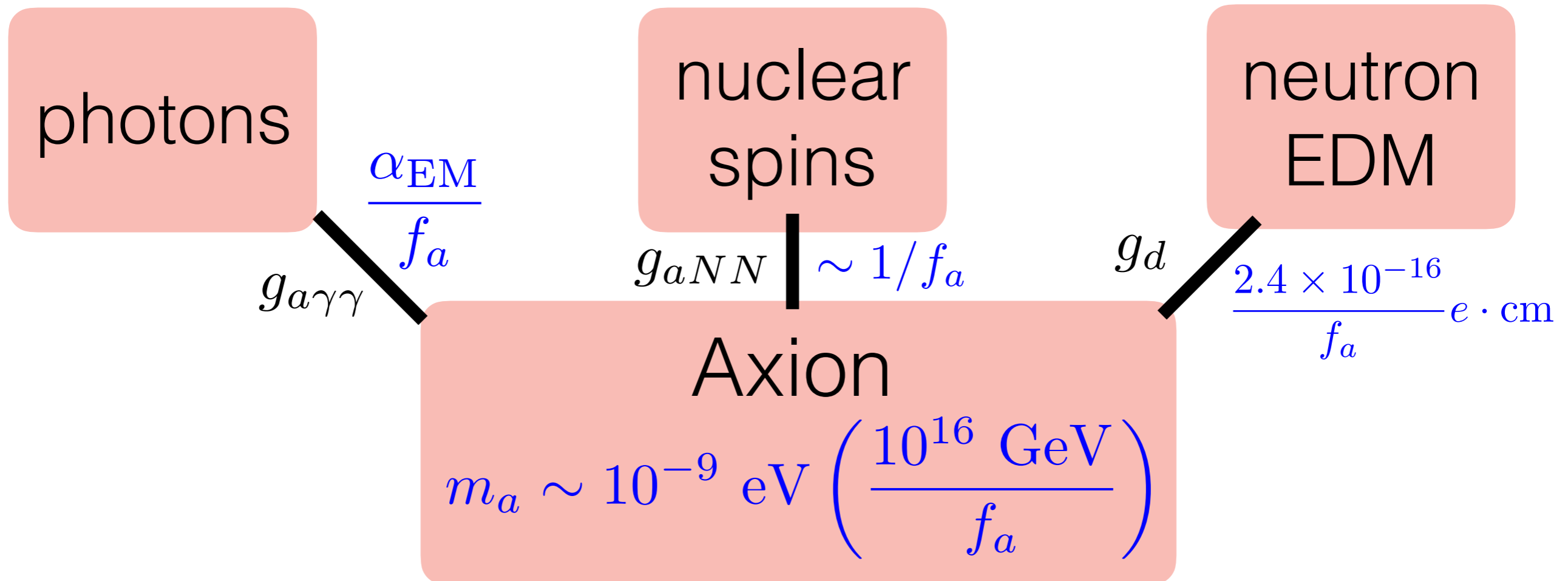
$$\Delta v_{\text{DM}} \sim v_{\text{DM}} \sim 10^{-3}$$

$$\lambda_{\text{dB}} = \frac{\lambda_{\text{Comp}}}{v_{\text{DM}}}$$

$$\tau_{\text{coh}} = \frac{\tau_{\text{Comp}}}{v_{\text{DM}}^2}$$

Classical physics is fine: $m_a = 10^{-9}$ eV $\implies N_a \sim 10^{18}/\text{cm}^3$

How do axions interact?



For QCD axion, only one free parameter!

While strong-CP is a great motivation, let's think broadly.

For this talk, "axion" can vary mass/coupling relation

What should we measure?

$$a(\mathbf{x}, t) = \frac{\sqrt{2\rho_{\text{DM}}}}{m_a} \cos(m_a t + \mathcal{O}(v_{\text{DM}})\mathbf{x})$$

In axion DM background, get oscillating observables:

$$\left. \begin{aligned} \nabla \times \mathbf{B}_a &= \frac{\partial \mathbf{E}_a}{\partial t} - g_{a\gamma\gamma} \left(\mathbf{E}_0 \times \nabla a - \mathbf{B}_0 \frac{\partial a}{\partial t} \right) \\ \nabla \cdot \mathbf{E}_a &= -g_{a\gamma\gamma} \mathbf{B}_0 \cdot \nabla a \end{aligned} \right\} \begin{array}{l} \text{Oscillating} \\ \text{response} \\ \text{from static} \\ \text{fields} \end{array}$$

$$H_N \supset g_{aNN} \nabla a \cdot \vec{\sigma}_N \quad \text{Spin-dependent force}$$

$$d_n = g_d a \quad \text{Time-varying EDM}$$

Note: $\nabla a \sim v_{\text{DM}} \sim 10^{-3}$ so some are easier than others

Axion searches with magnetic fields

$$\underbrace{\nabla \times \mathbf{B}_a}_{\text{blue}} = \frac{\partial \mathbf{E}_a}{\partial t} - g_{a\gamma\gamma} \mathbf{B}_0 \frac{\partial a}{\partial t}$$

Cavity regime: $\lambda_{\text{Comp}} \sim R_{\text{exp}}$
e.g. ADMX

$$\nabla \times \mathbf{B}_a = \cancel{\frac{\partial \mathbf{E}_a}{\partial t}} - g_{a\gamma\gamma} \underbrace{\mathbf{B}_0 \frac{\partial a}{\partial t}}_{\mathbf{J}_{\text{eff}}}$$

Quasistatic regime: $\lambda_{\text{Comp}} \gg R_{\text{exp}}$
e.g. ABRACADBRA

$$\cancel{\nabla \times} \mathbf{B}_a = \frac{\partial \mathbf{E}_a}{\partial t} - g_{a\gamma\gamma} \mathbf{B}_0 \frac{\partial a}{\partial t}$$

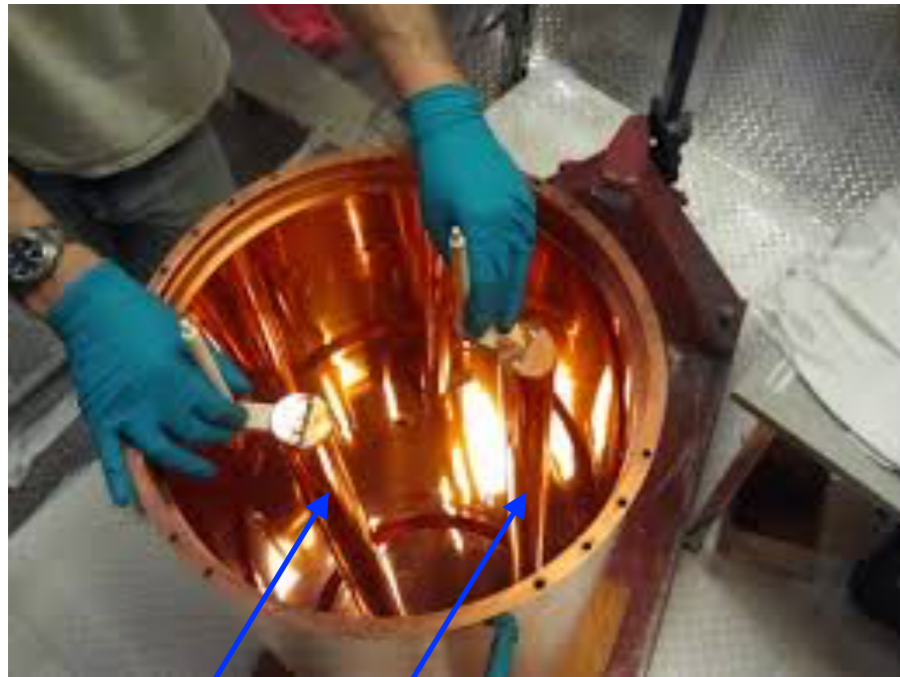
Radiation regime: $\lambda_{\text{Comp}} \ll R_{\text{exp}}$
e.g. MADMAX

Resonant cavity detection

$$\underbrace{\nabla \times \mathbf{B}_a = \frac{\partial \mathbf{E}_a}{\partial t}}_{\text{cavity response}} - g_{a\gamma\gamma} \mathbf{B}_0 \frac{\partial a}{\partial t}$$

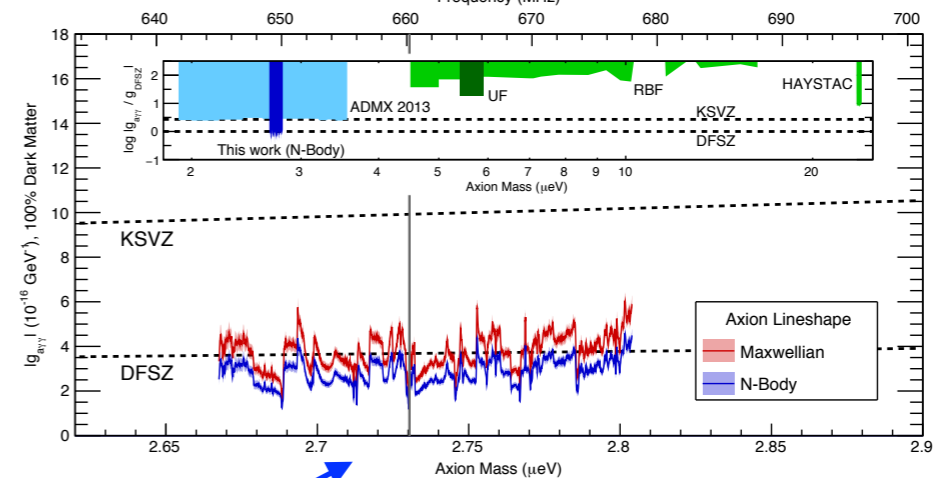
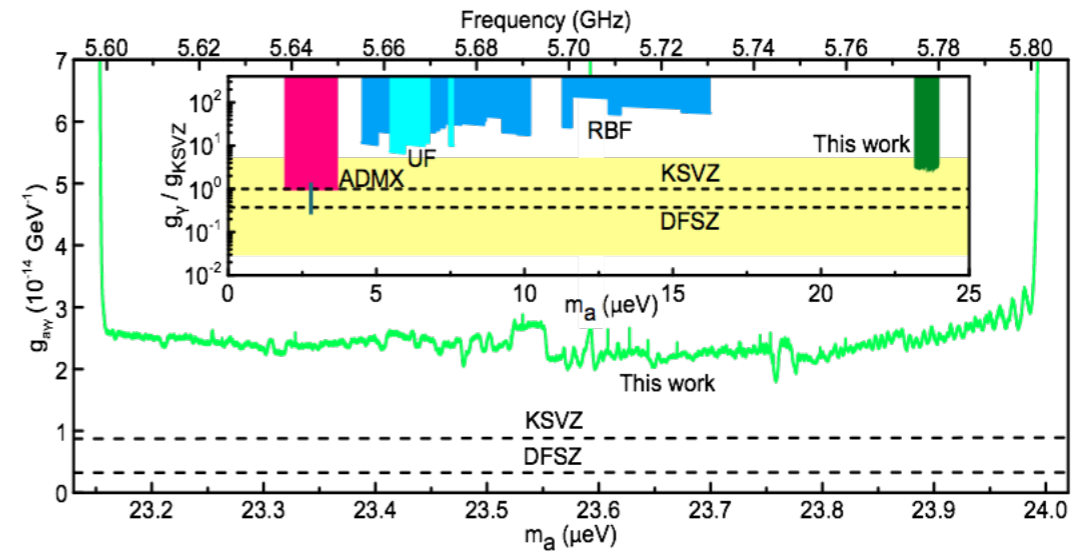
axion source

cavity response



Tune cavity modes to scan axion masses

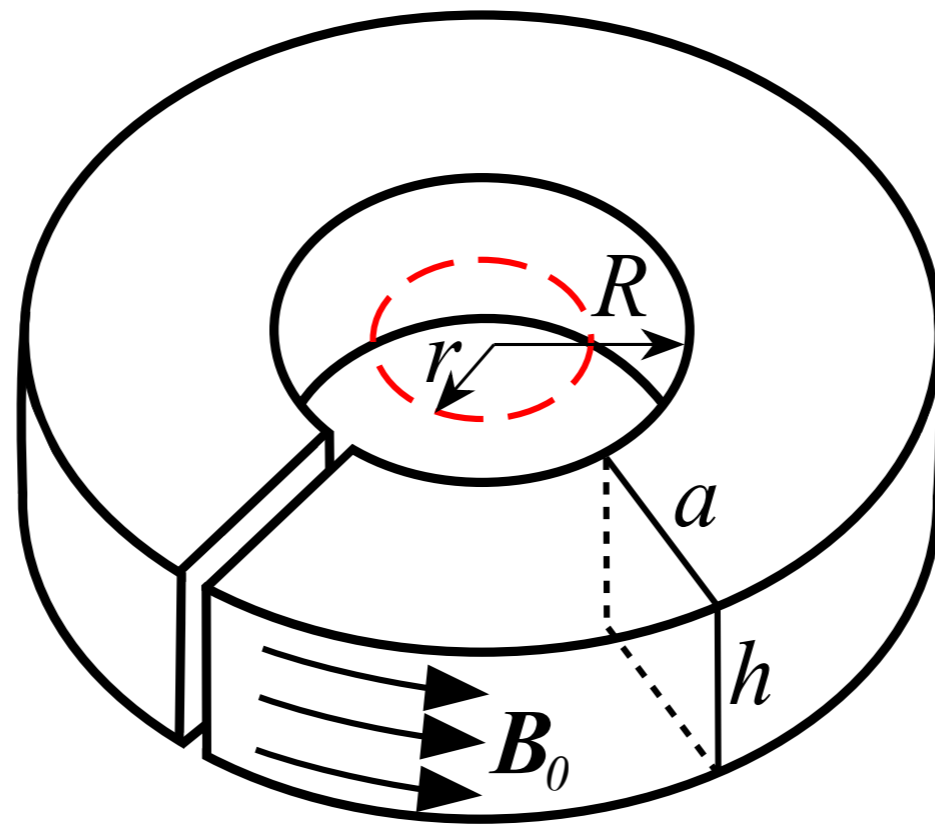
$$P \sim g_{a\gamma\gamma}^2 \frac{\rho_{\text{DM}}}{m_a} B_0^2 V Q$$



Cavity b.c. fix mass range to cavity size; larger masses \rightarrow smaller V

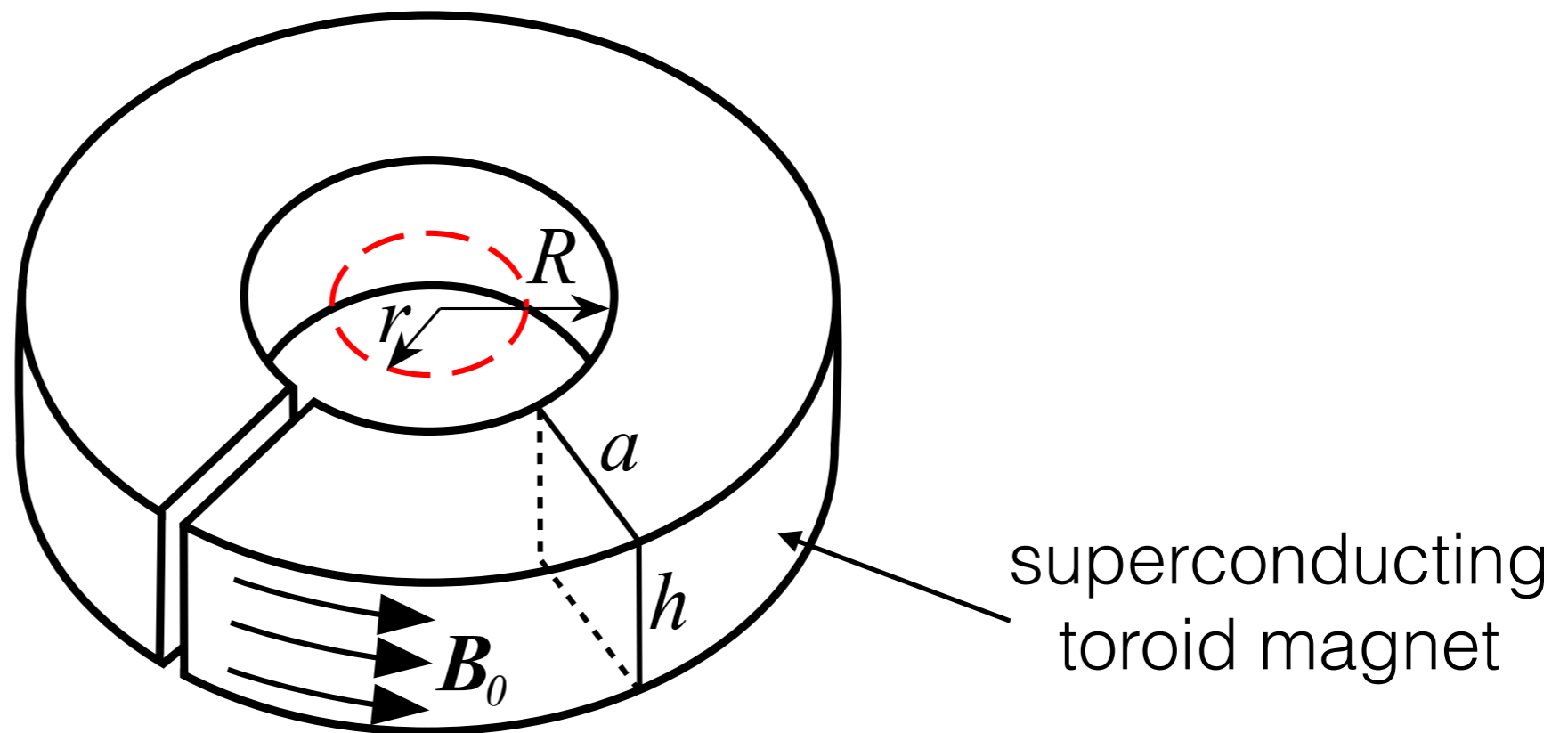
Quasistatic regime: ABRACADABRA

A **B**roadband/**R**esonant **A**pproach to **C**osmic **A**xion **D**etection
with an **A**mplifying **B**-field **R**ing **A**pparatus



Quasistatic regime: ABRACADABRA

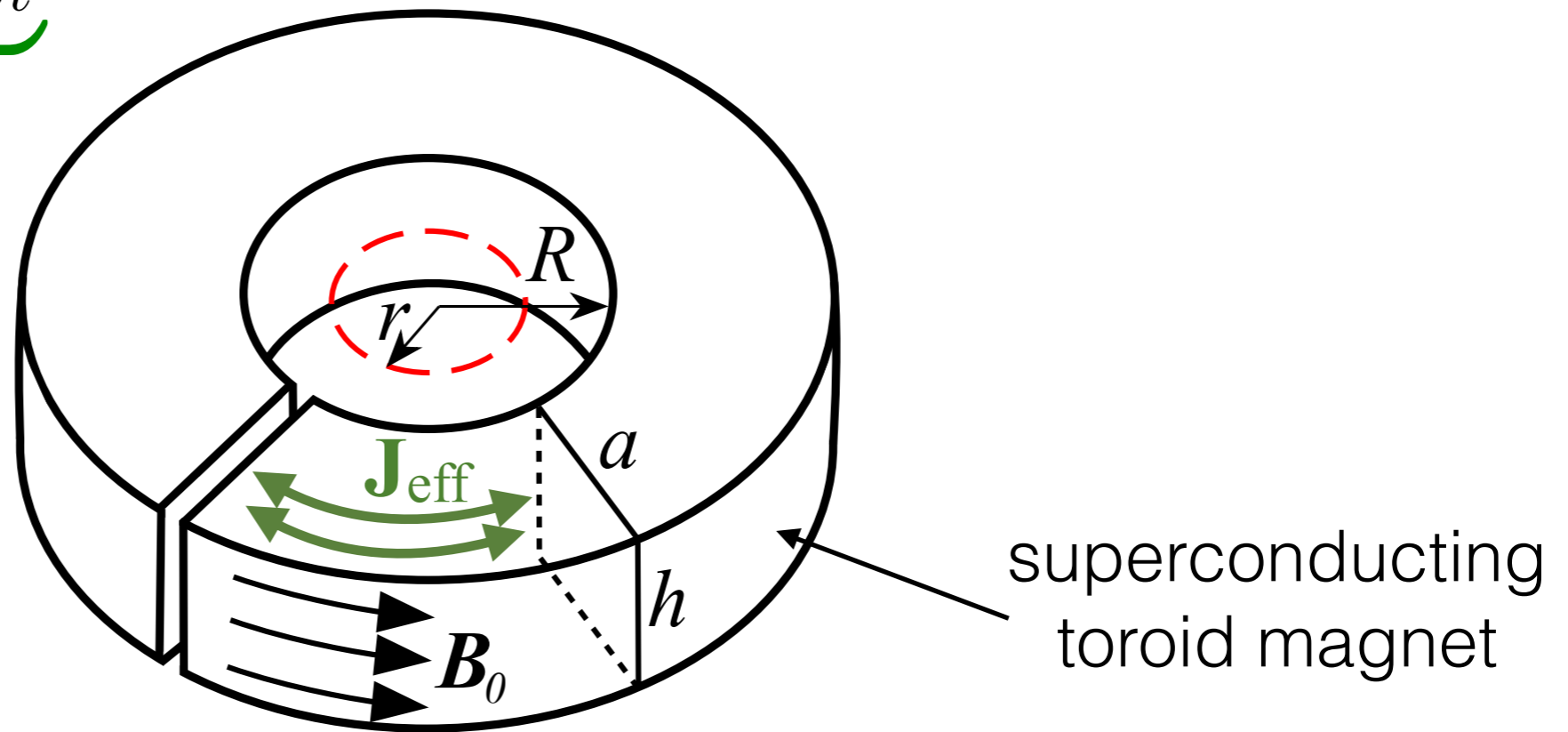
A **B**roadband/**R**esonant **A**pproach to **C**osmic **A**xion **D**etection
with an **A**mplifying **B**-field **R**ing **A**pparatus



Quasistatic regime: ABRACADABRA

A **B**roadband/**R**esonant **A**pproach to **C**osmic **A**xion **D**etection
with an **A**mplifying **B**-field **R**ing **A**pparatus

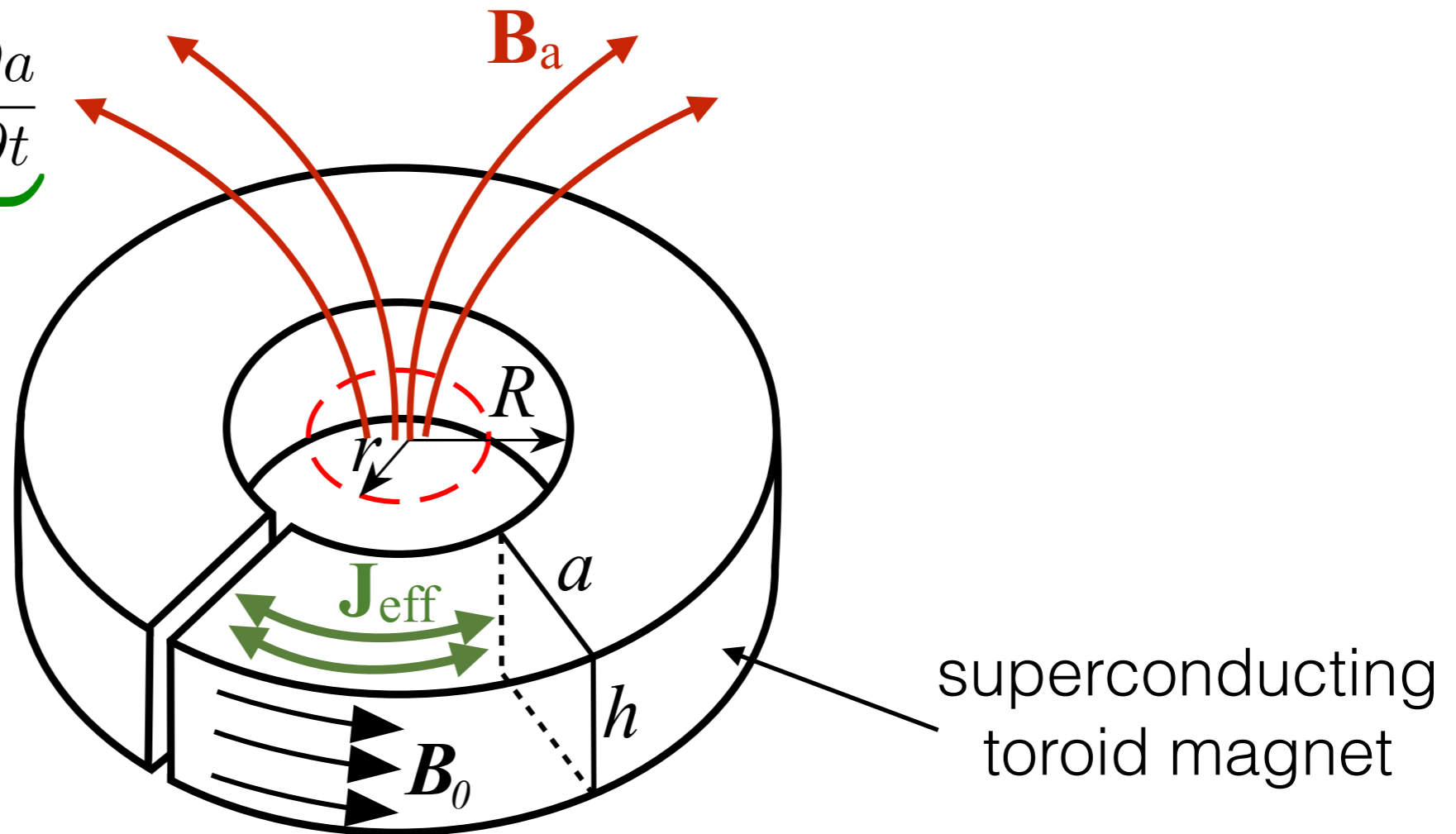
$$\nabla \times \mathbf{B}_a = \cancel{\frac{\partial \mathbf{E}_a}{\partial t}} - \underbrace{g_{a\gamma\gamma} \mathbf{B}_0}_{\mathbf{J}_{\text{eff}}} \frac{\partial a}{\partial t}$$



Quasistatic regime: ABRACADABRA

A **B**roadband/**R**esonant **A**pproach to **C**osmic **A**xion **D**etection
with an **A**mplifying **B**-field **R**ing **A**pparatus

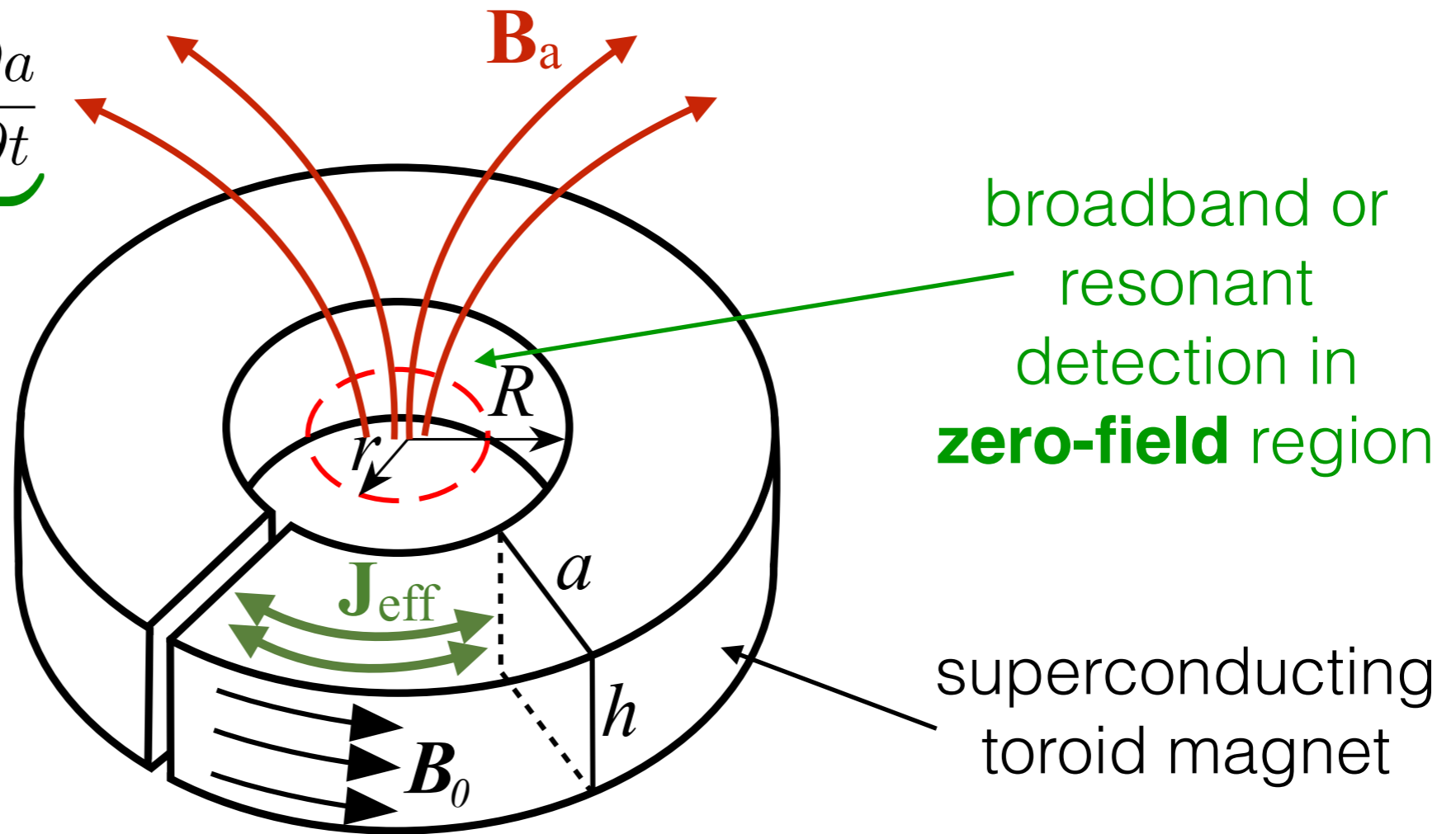
$$\nabla \times \mathbf{B}_a = \cancel{\frac{\partial \mathbf{E}_a}{\partial t}} - \underbrace{g_{a\gamma\gamma} \mathbf{B}_0 \frac{\partial a}{\partial t}}_{\mathbf{J}_{\text{eff}}}$$



Quasistatic regime: ABRACADABRA

A Broadband/**R**esonant **A**pproach to **C**osmic **A**xion **D**etection
with an **A**mplifying **B**-field **R**ing **A**pparatus

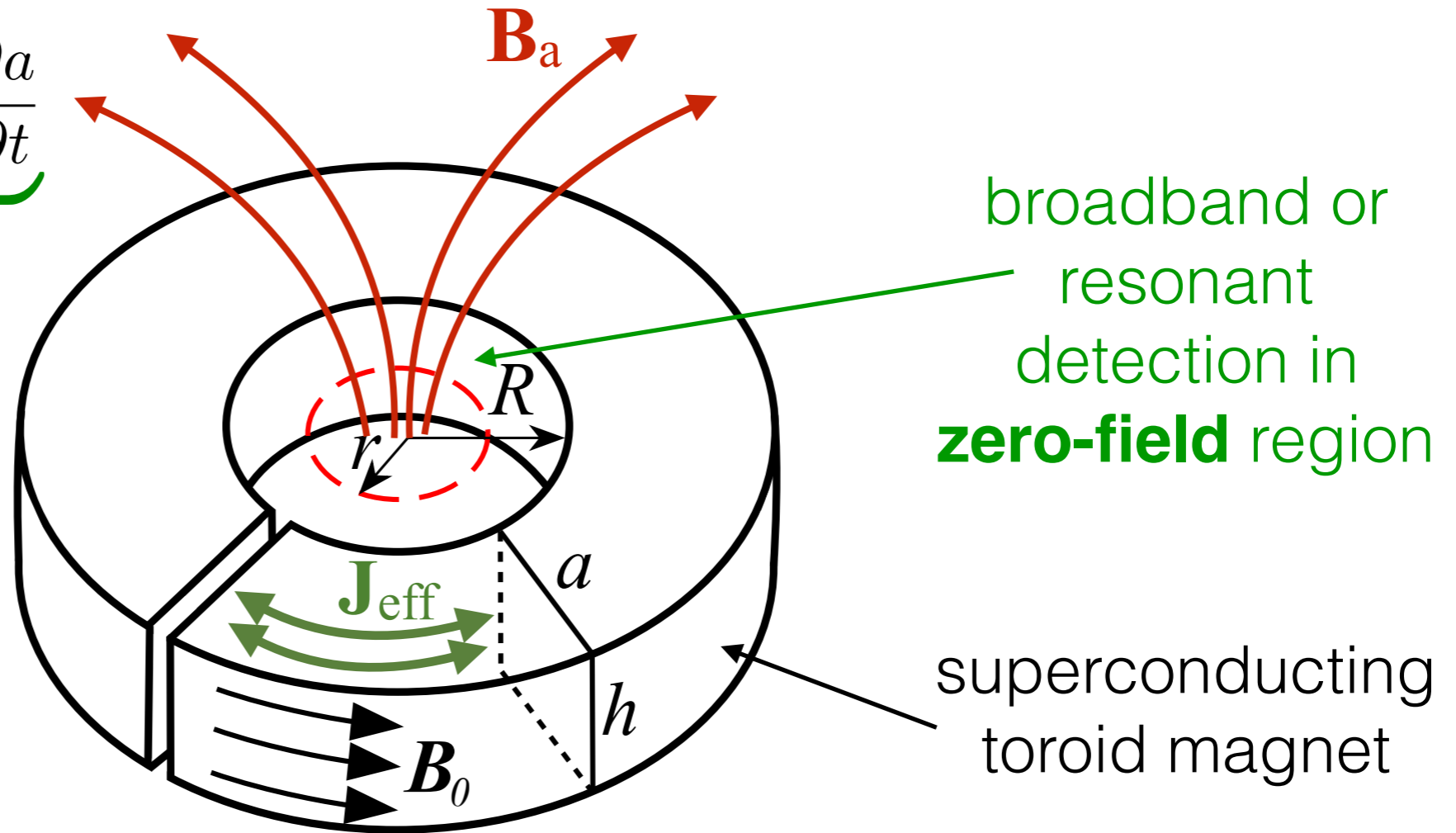
$$\nabla \times \mathbf{B}_a = \cancel{\frac{\partial \mathbf{E}_a}{\partial t}} - \underbrace{g_{a\gamma\gamma} \mathbf{B}_0}_{\mathbf{J}_{\text{eff}}} \frac{\partial a}{\partial t}$$



Quasistatic regime: ABRACADABRA

A **B** broadband/**R**esonant **A**pproach to **C**osmic **A**xion **D**etection
with an **A**mplifying **B**-field **R**ing **A**pparatus

$$\nabla \times \mathbf{B}_a = \cancel{\frac{\partial \mathbf{E}_a}{\partial t}} - \underbrace{g_{a\gamma\gamma} \mathbf{B}_0}_{\mathbf{J}_{\text{eff}}} \frac{\partial a}{\partial t}$$



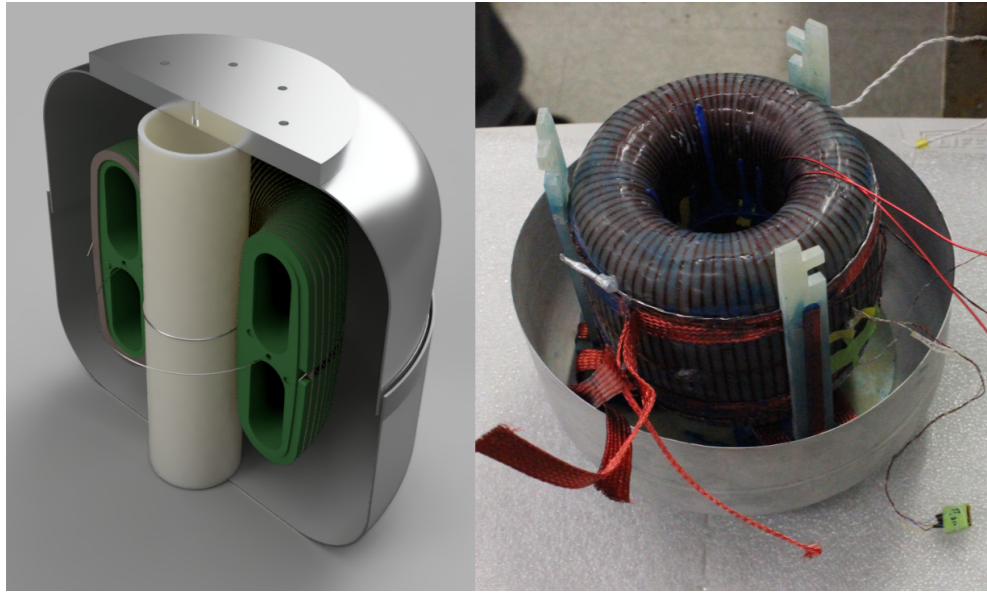
$$\langle \Phi(t)^2 \rangle \sim g_{a\gamma\gamma}^2 \rho_{\text{DM}} B_0^2 V^{5/3}$$

Volume enhancement at low masses



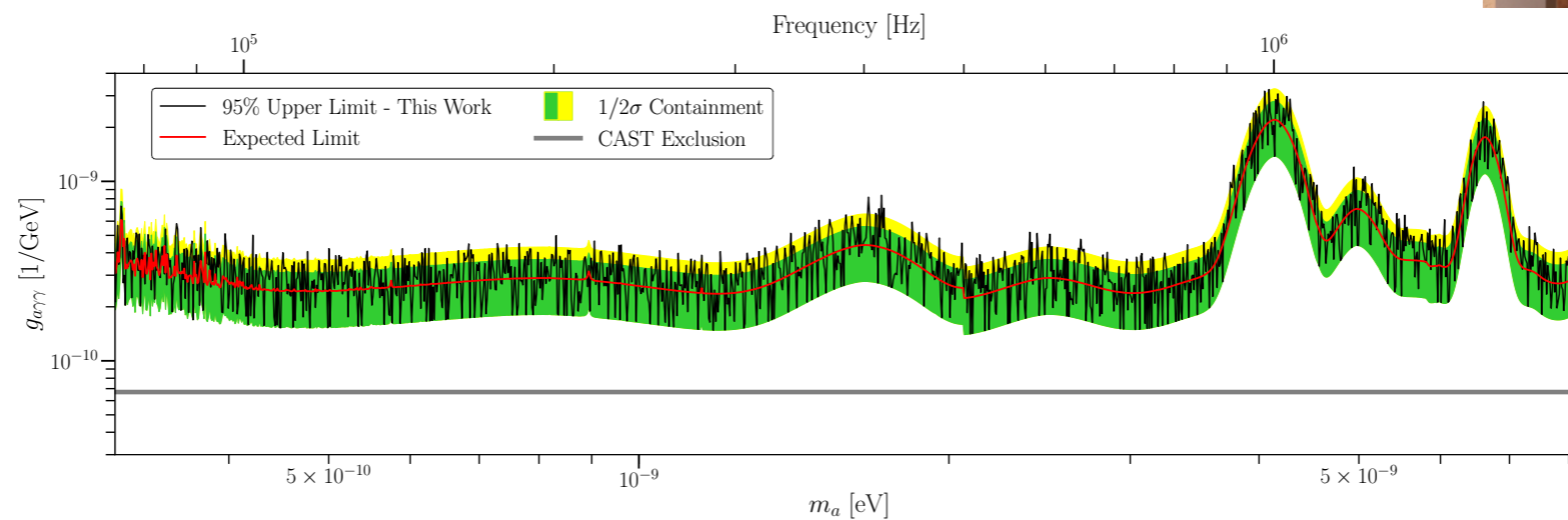
ABRACADABRA

first results



10 cm magnet and pickup

inside a 150 mK
dilution refrigerator

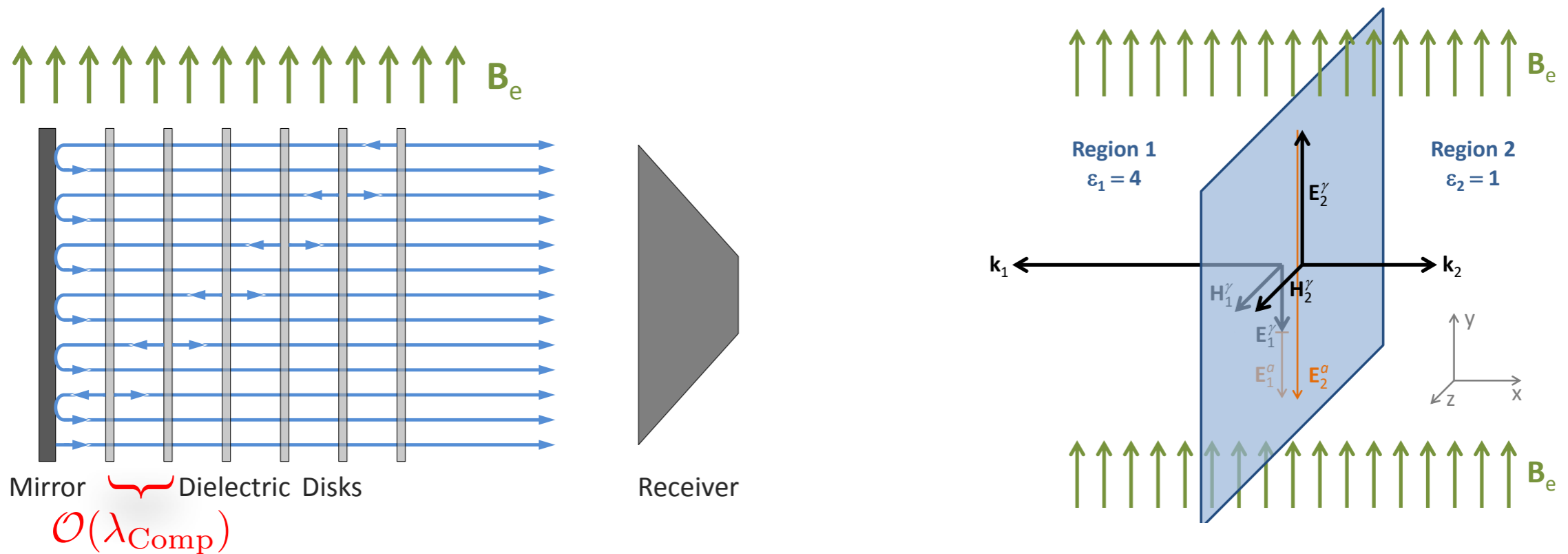


Second run underway, expect better limits soon;
merging w/DM-Radio @ SLAC

Radiation regime: MADMAX

$$\nabla \times \cancel{\mathbf{B}_a} = \frac{\partial \mathbf{E}_a}{\partial t} - g_{a\gamma\gamma} \mathbf{B}_0 \frac{\partial a}{\partial t} \implies \mathbf{D}_a(t) = \epsilon \mathbf{E}_a(t) = -g_{a\gamma\gamma} \mathbf{B}_0 a(t)$$

E+M boundary condition at interfaces forces
radiation to cancel axion-induced **D**



Composite design with **large volume** and **high Q**, design underway

NMR with axion DM

Nuclei immersed in axion DM can have:

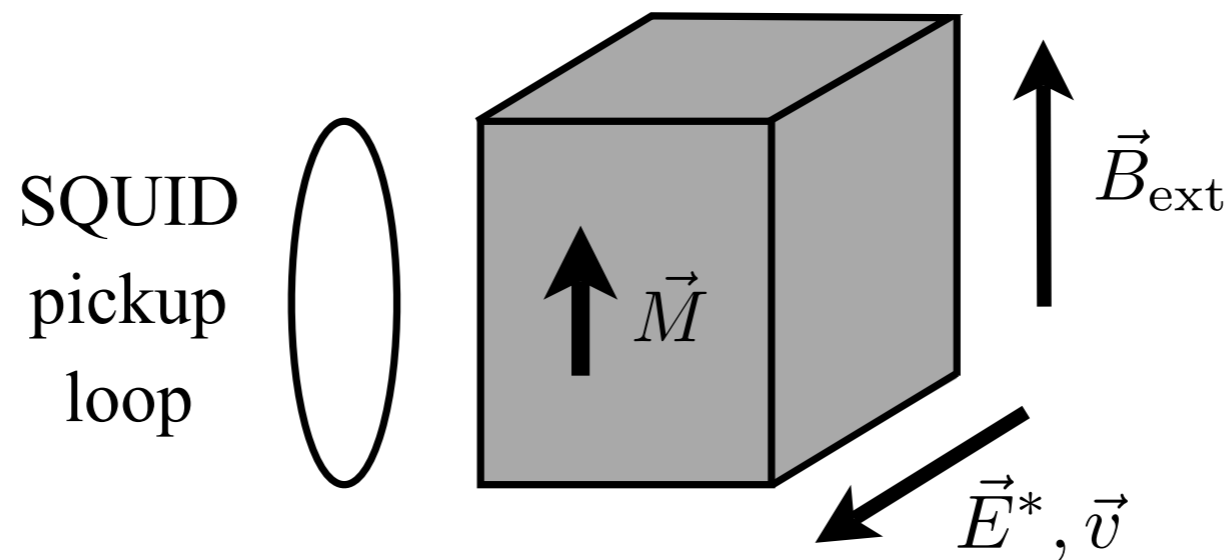
Oscillating EDM

and/or

Spin-dependent force

$$d_n = g_d \frac{\sqrt{2\rho_{DM}}}{m_a} \cos(m_a t)$$

$$H_N \supset g_{aNN} \sqrt{2\rho_{DM}} \cos(m_a t) \vec{v} \cdot \vec{\sigma}_N$$



NMR with axion DM

Nuclei immersed in axion DM can have:

Oscillating EDM

and/or

Spin-dependent force

$$d_n = g_d \frac{\sqrt{2\rho_{DM}}}{m_a} \cos(m_a t)$$

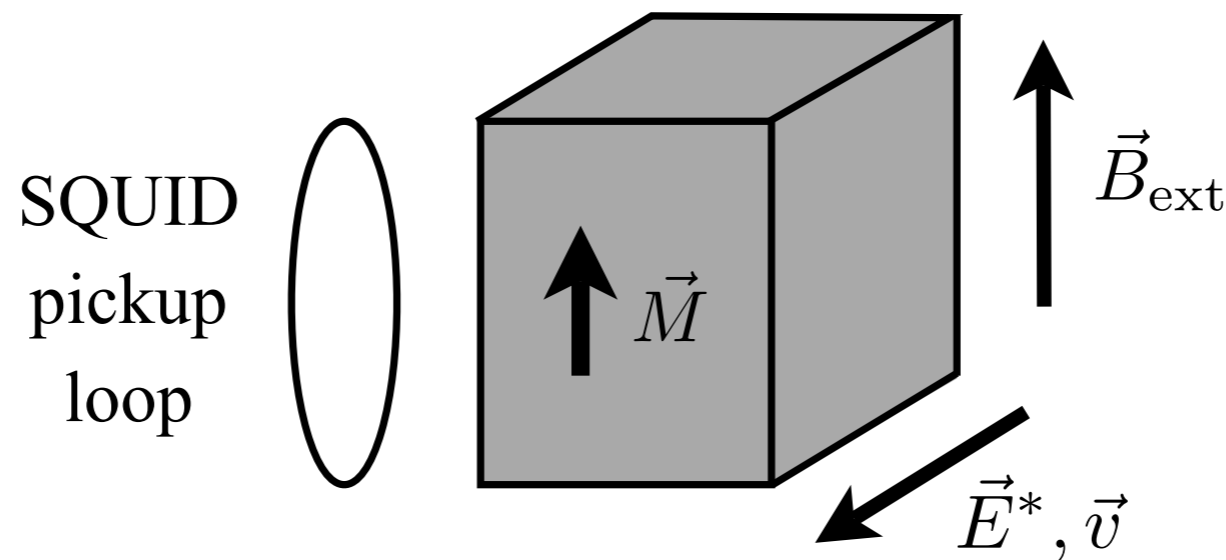
$$H_N \supset g_{aNN} \sqrt{2\rho_{DM}} \cos(m_a t) \vec{v} \cdot \vec{\sigma}_N$$

Polarize some spins, watch them precess around:

External E field

and/or

Axion field velocity



NMR with axion DM

Nuclei immersed in axion DM can have:

Oscillating EDM

and/or

Spin-dependent force

$$d_n = g_d \frac{\sqrt{2\rho_{DM}}}{m_a} \cos(m_a t)$$

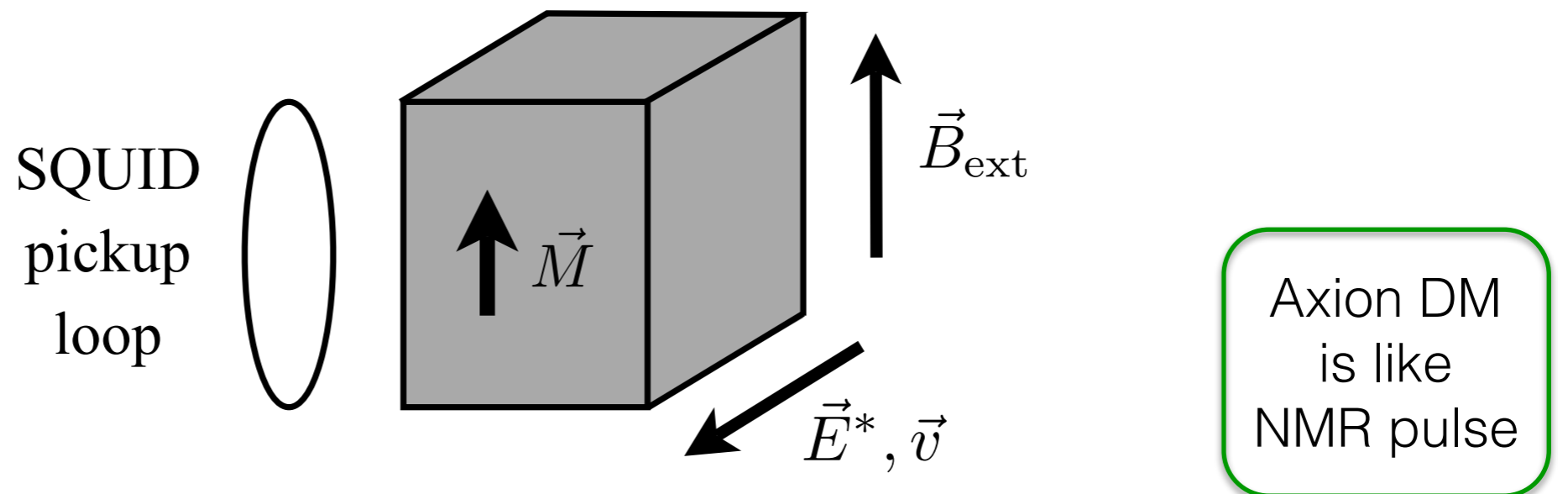
$$H_N \supset g_{aNN} \sqrt{2\rho_{DM}} \cos(m_a t) \vec{v} \cdot \vec{\sigma}_N$$

Polarize some spins, watch them precess around:

External E field

and/or

Axion field velocity

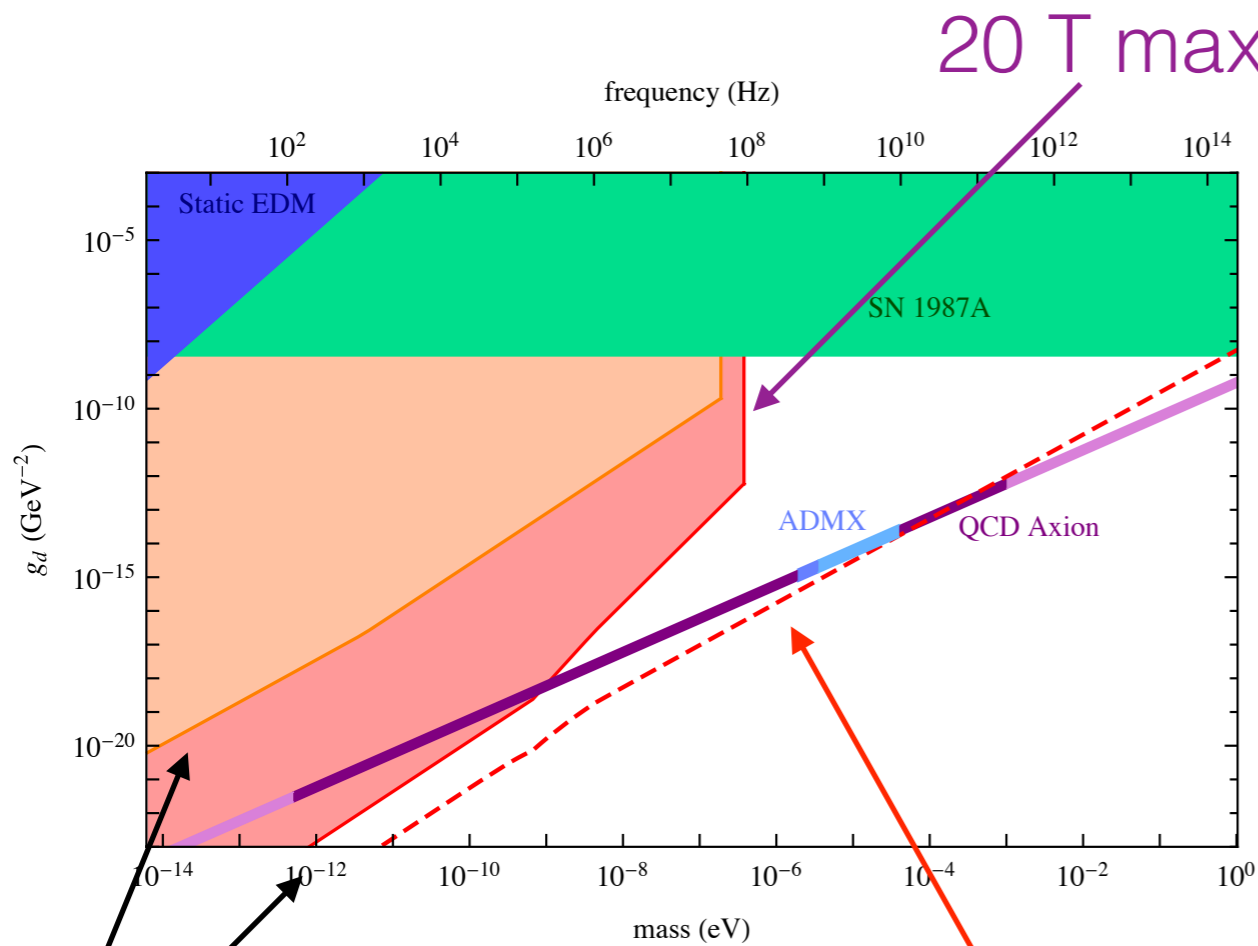


Resonance in transverse magnetization when $2\mu B_{\text{ext}} = m_a$

CASPEr projected Reach

CASPEr-Electric

[Budker et al., Phys. Rev. X 2014]

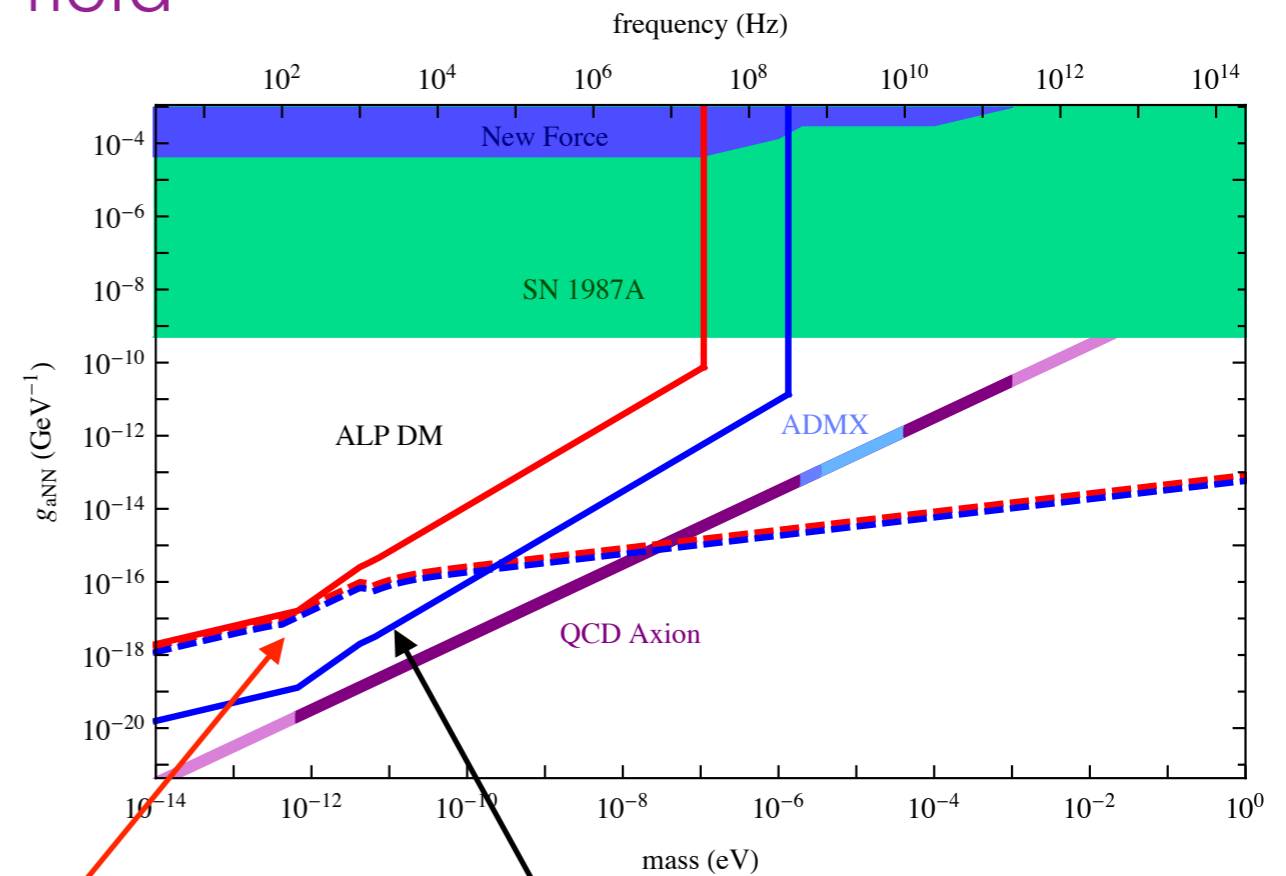


20 T max B-field

magnetization noise

CASPEr-Wind

[Graham and Rajendran, Phys. Rev. D 2013]

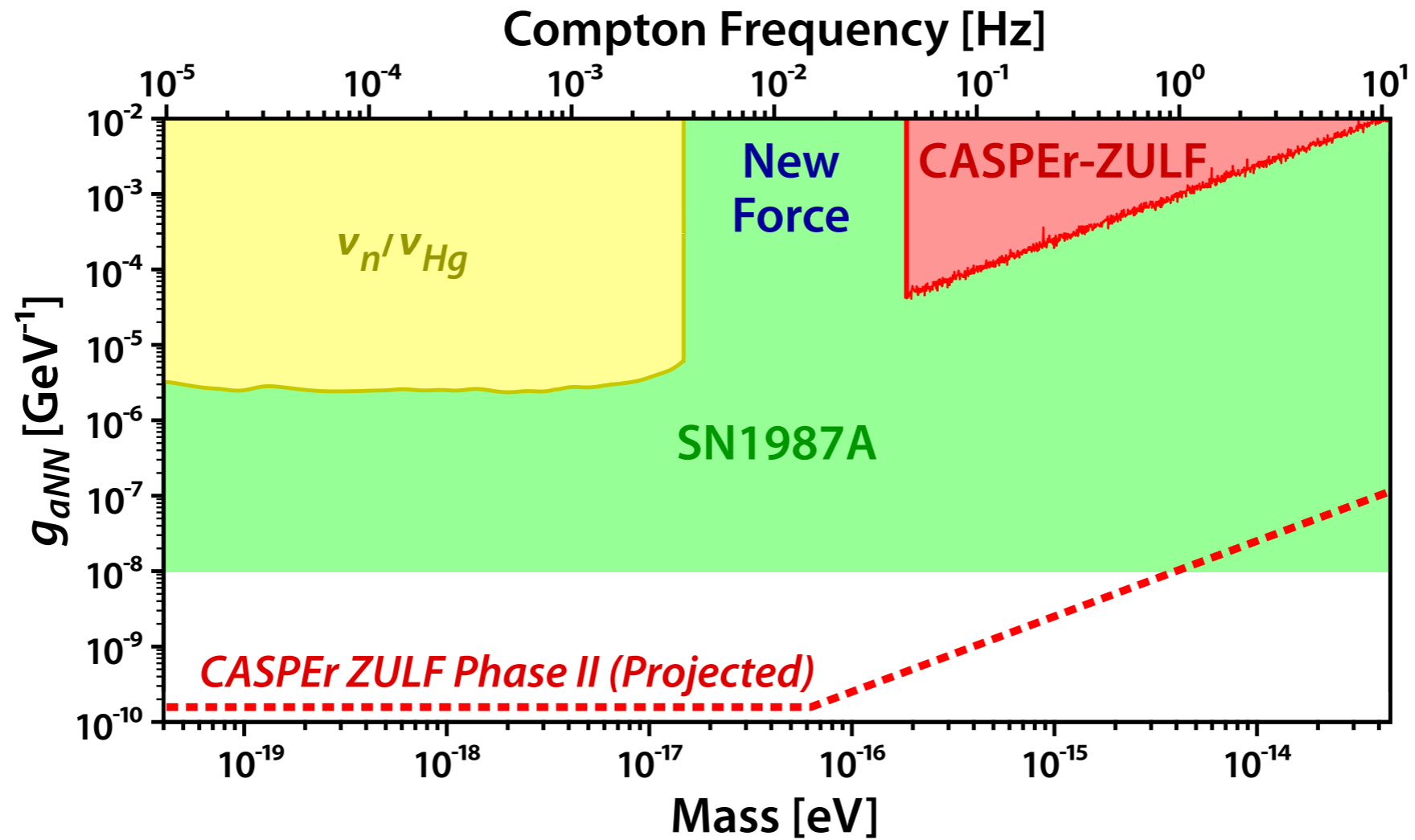


velocity suppression:
can't quite reach QCD axion

non-decoupling signal!

CASPEr ZULF

ALP Wind coupling at very low frequencies



[Garcon et al., Science Advances 2019]

And now
for something
completely different...



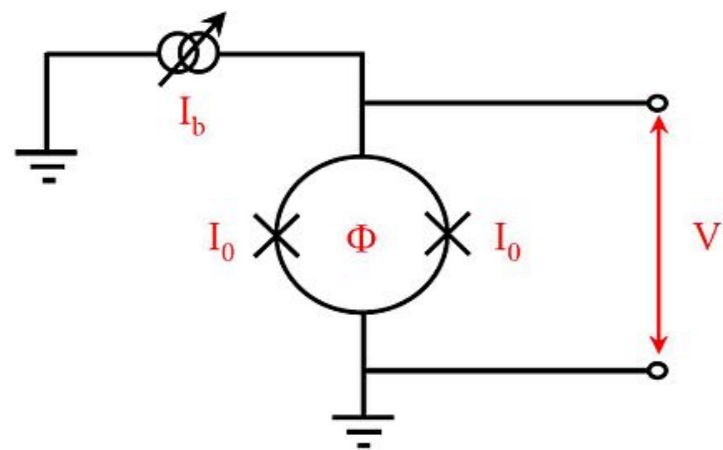
Some musings on direct detection

w/N. Kurinsky, D. Baxter, G. Krnjaic

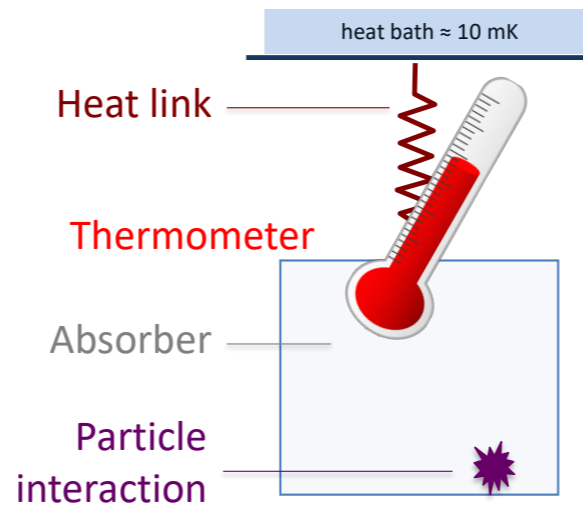
Direct detection principles

Three things detectors can see: charge, light, and heat

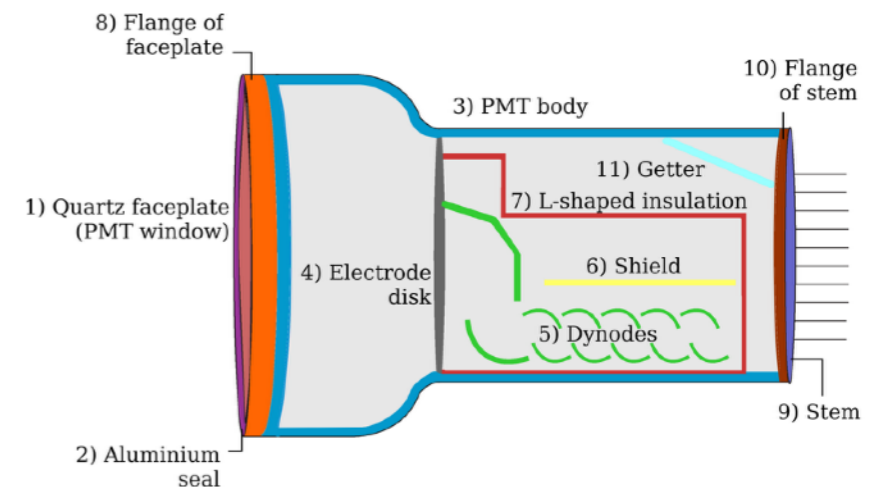
Charge: measure current (e.g. SQUID + HEMT) or voltage (e.g. CCD)



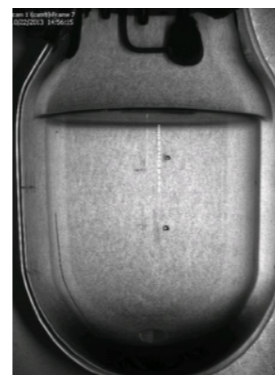
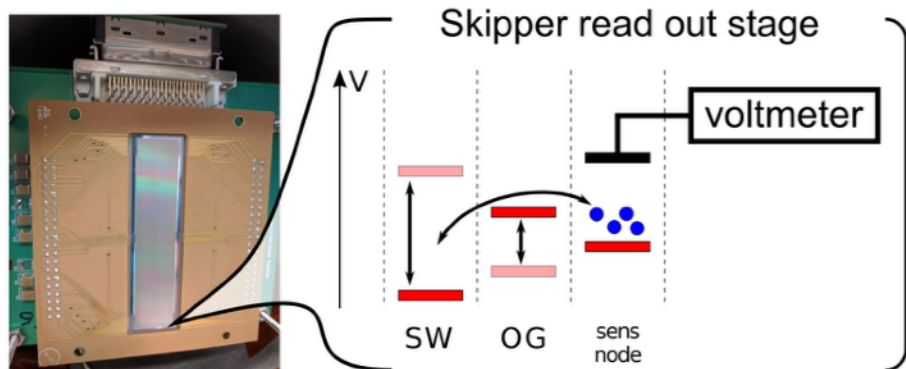
Heat: bolometer (e.g. TES, MKID) or bubble chamber



Light: photosensor (e.g. PMT)

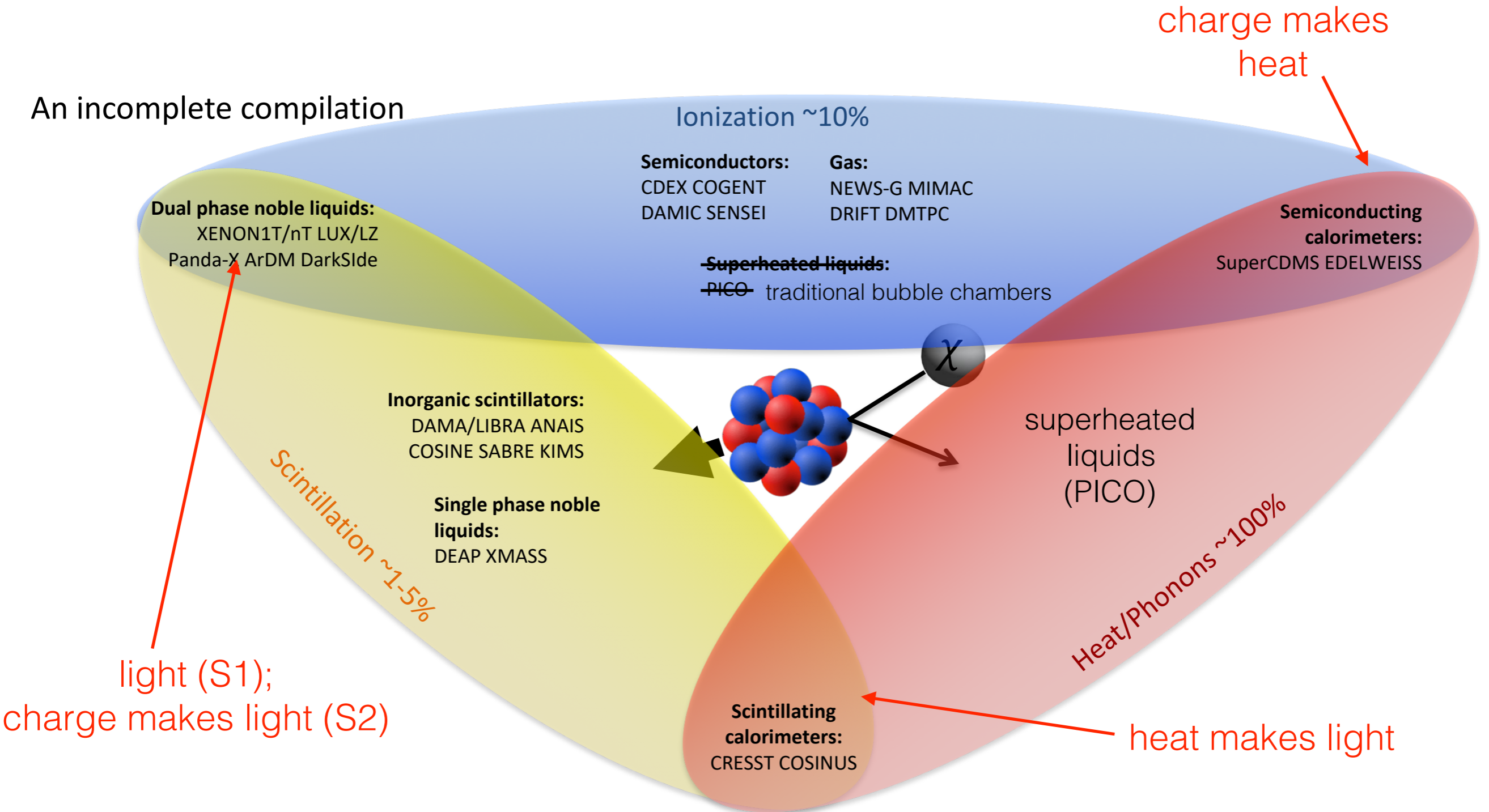


(but really, charge makes light and light makes charge so these are interchangeable: e.g. SiPM is a “photodetector” but triggers on e/h pairs)



The dictionary is complicated!

An incomplete compilation

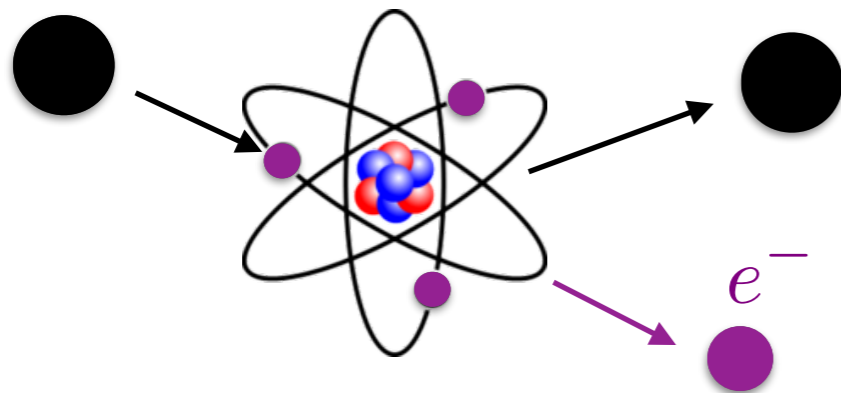


For traditional WIMP DM, signals and mapping between primary event and end-stage signal well-calibrated. For light DM and low thresholds, not so much!

More than just nuclear recoil!

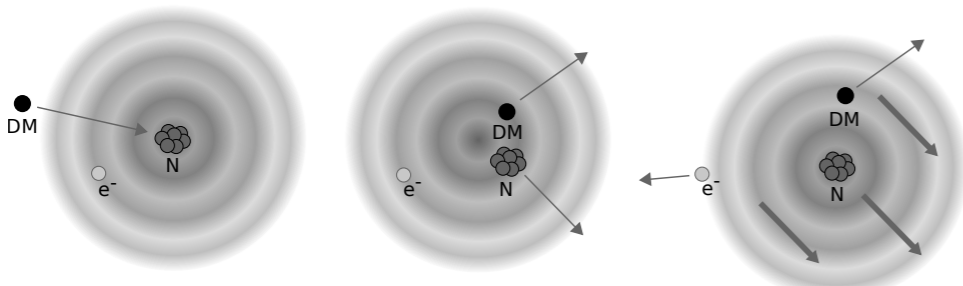
An isolated atom (nucleus + electrons) at rest is a momentum and energy eigenstate. If you whack the atom, both recoil

If you hit the electron directly:

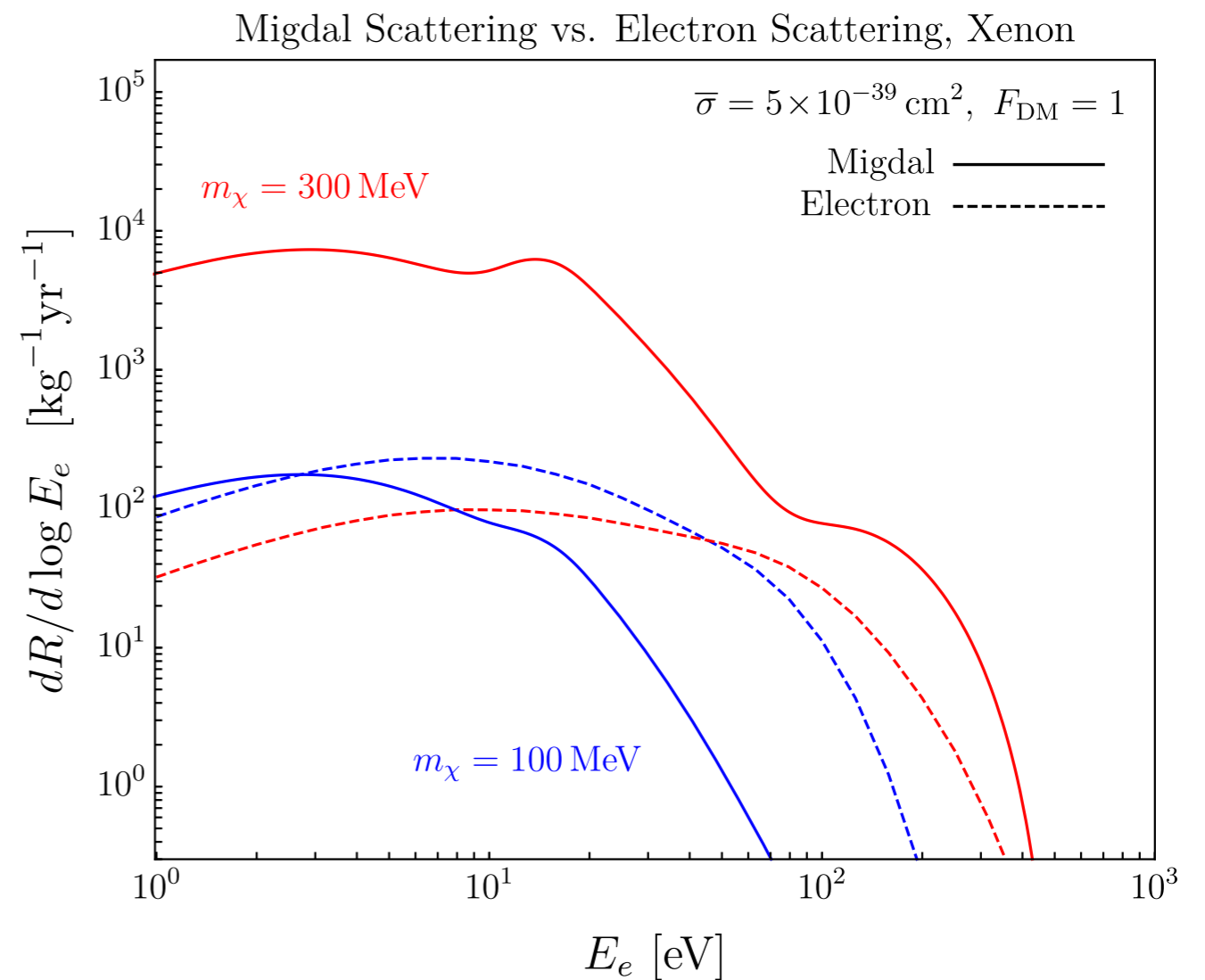


“electron recoil”

If you hit the nucleus:



“Migdal effect”



[Baxter, **YK**, Krnjaic 2019]

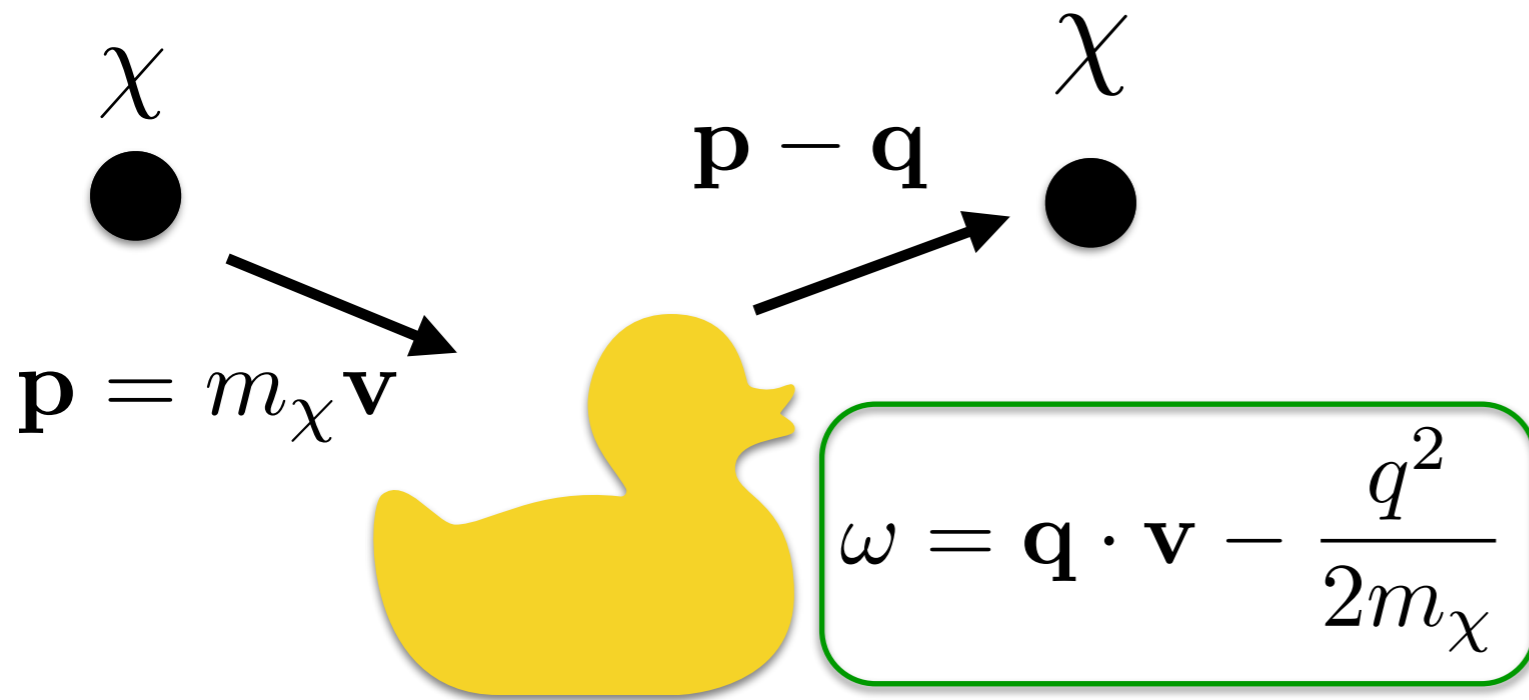
Same inelastic kinematics, vastly different dynamics!

A different perspective...

$$R \sim \int d^3\mathbf{v} f(\mathbf{v}) \int d^3\mathbf{q} F^2(\mathbf{q}) S(\mathbf{q}, \omega_{\mathbf{q}})$$

DM properties

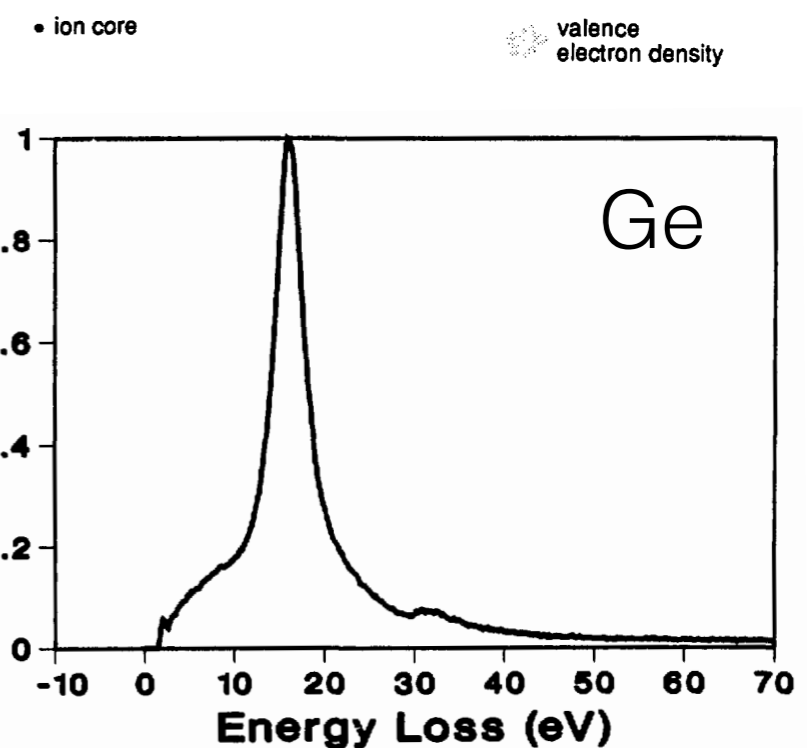
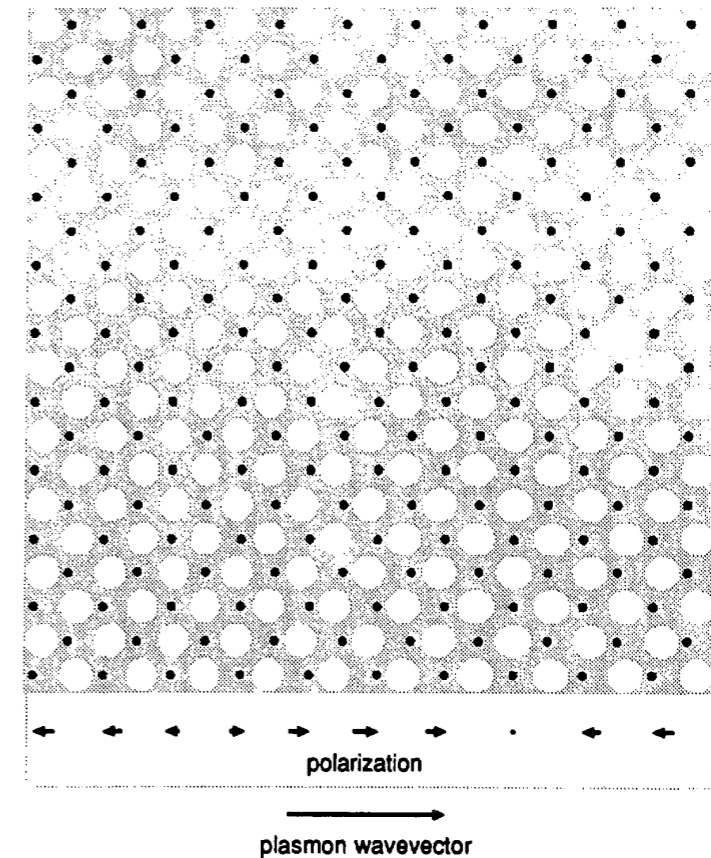
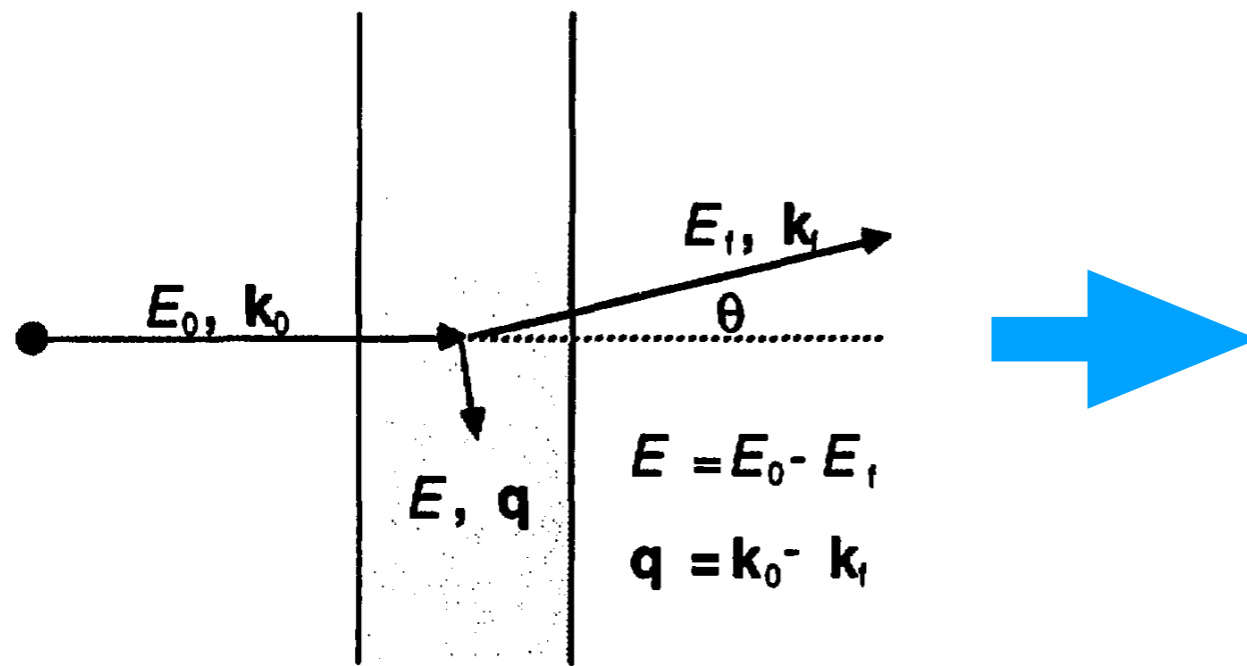
Material properties
(e.g. dielectric function)



(insert your favorite
detector here)

- Something** must respond
at the appropriate (\mathbf{q}, ω) :
- electronic bands
 - phonons
 - magnons
 - “free” nuclei (e.g. defects)
 - atomic orbitals
 - ... and many more
collective effects!

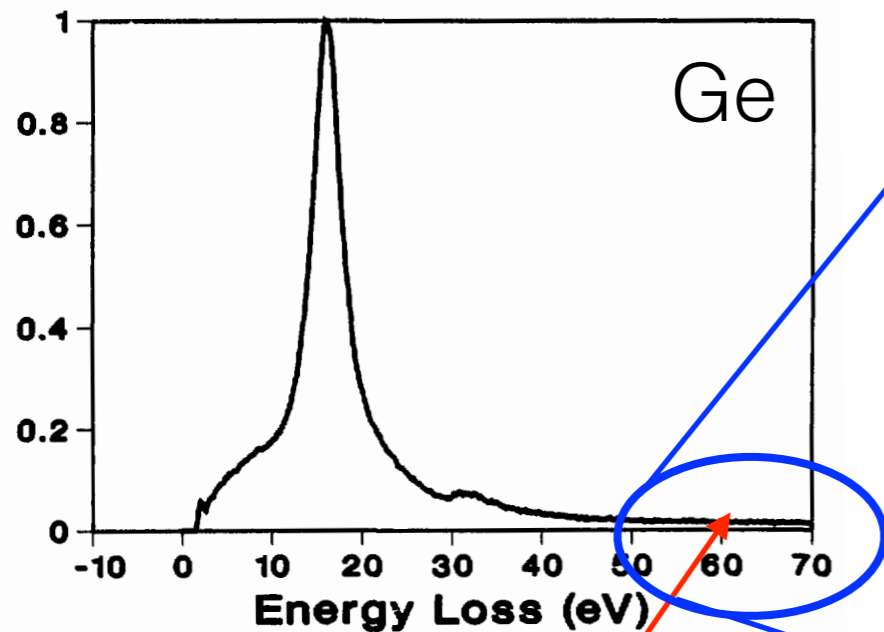
Example: EELS and plasmons



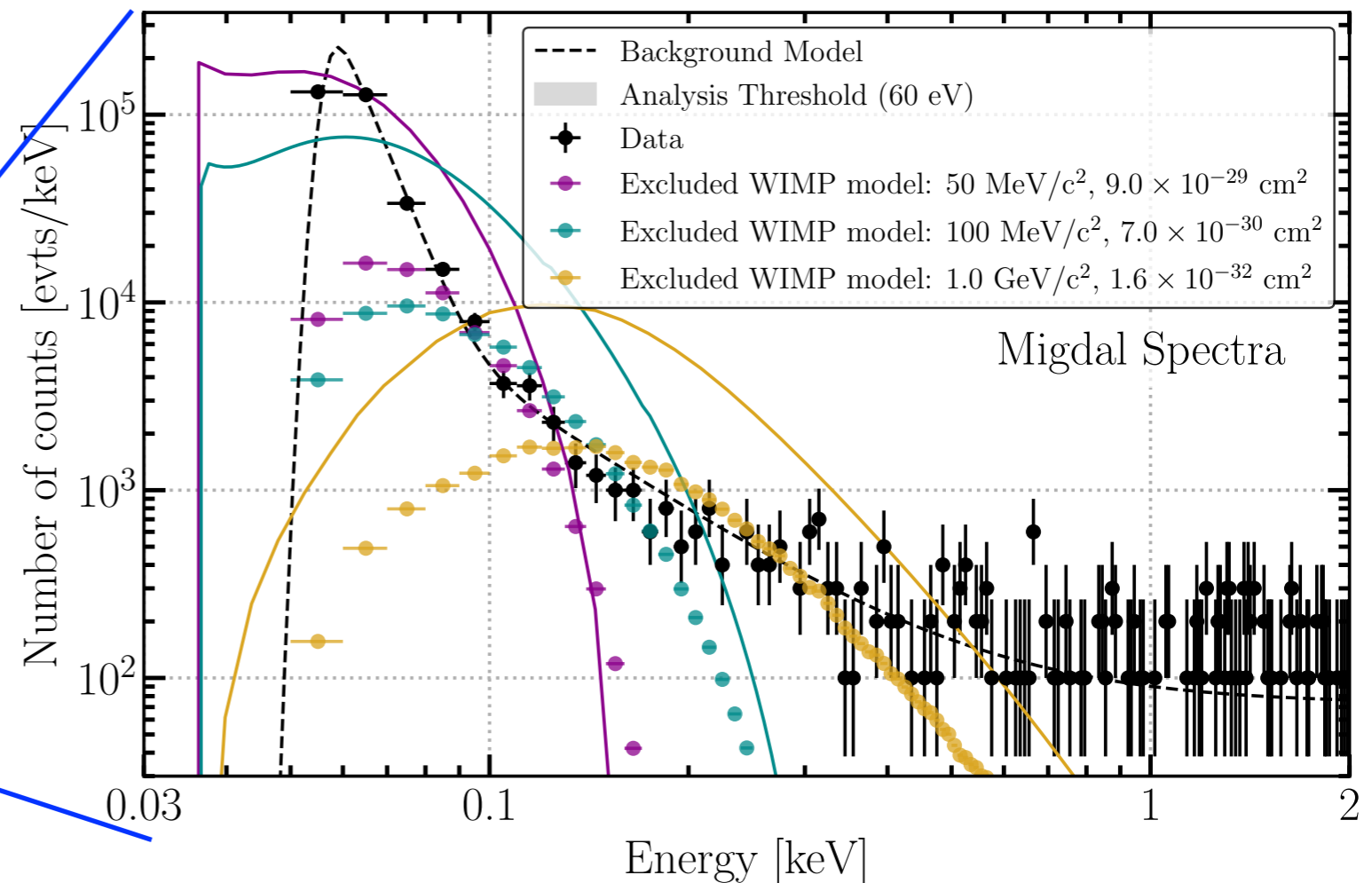
Semi-relativistic electron scattering **not** described by single-particle electron-electron scattering, but by a collective long-range charge wave (plasmon). Electron preferentially deposits ~ 15 eV of energy, **regardless of initial kinetic energy**

[M. Kundmann, Ph.D. thesis 1988]

A plasmon might look like this...

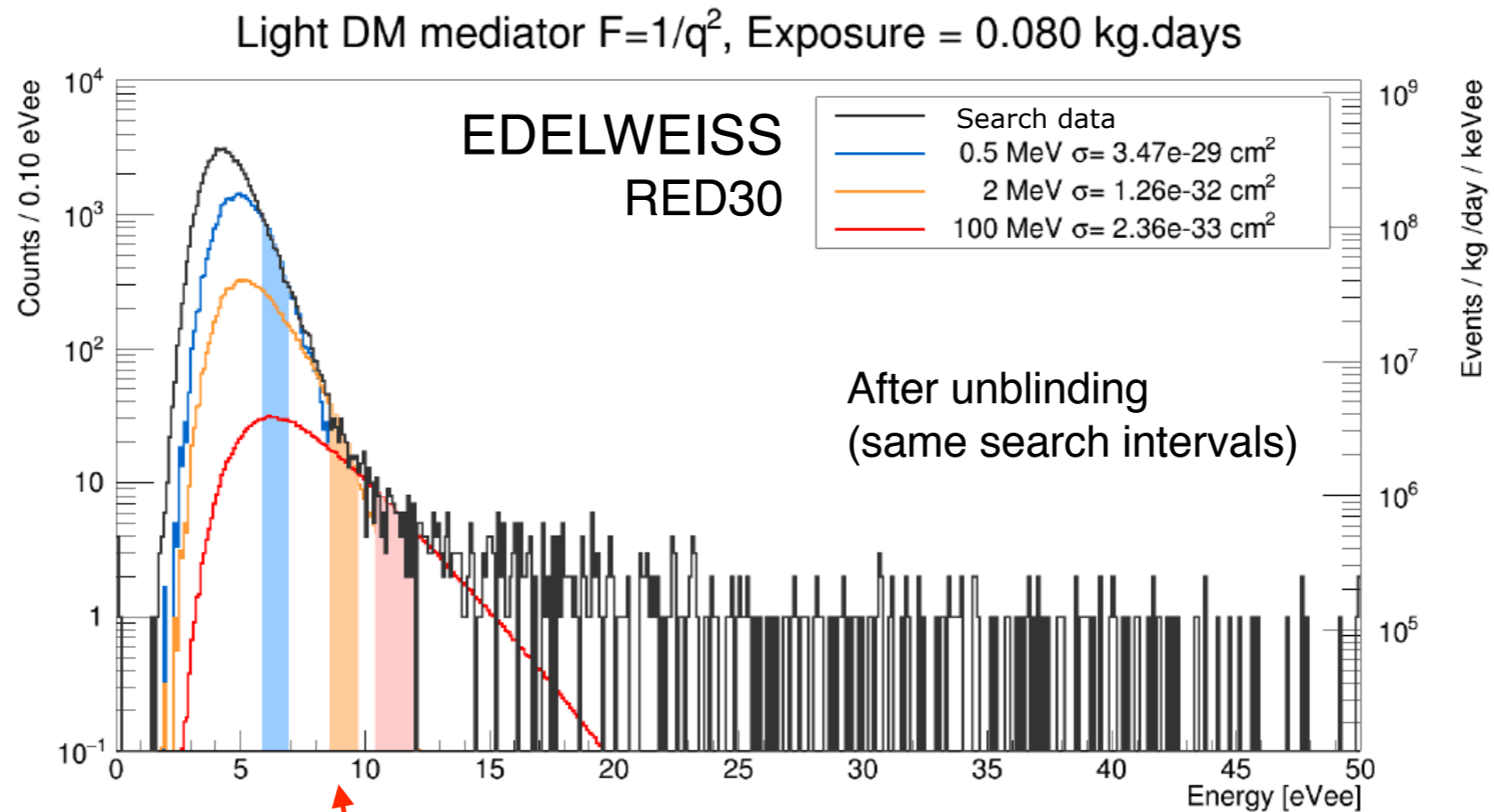


note the huge tail
on a **linear** scale!



EDELWEISS surface (2019)
Germanium detector

Or it might look like this:

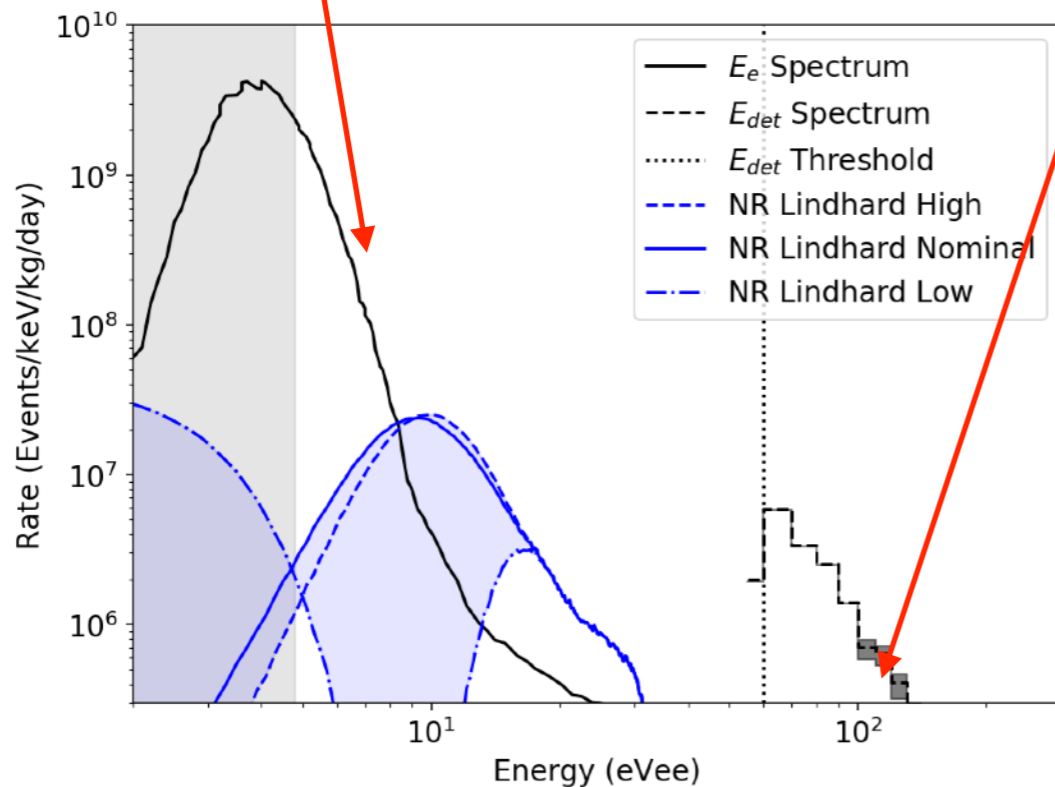


EDELWEISS underground (2020)

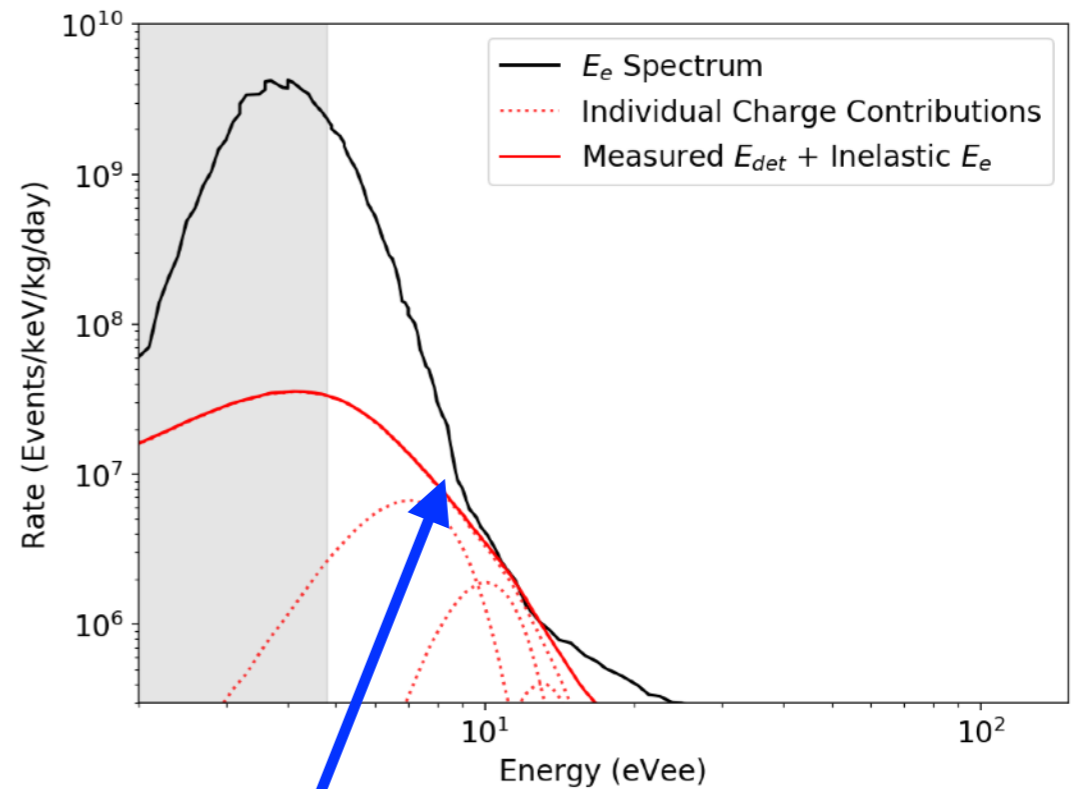
charge makes heat: this is really saying
“on average 1-2 e/h pairs per event”

Compare the two spectra:

electronic energy
only



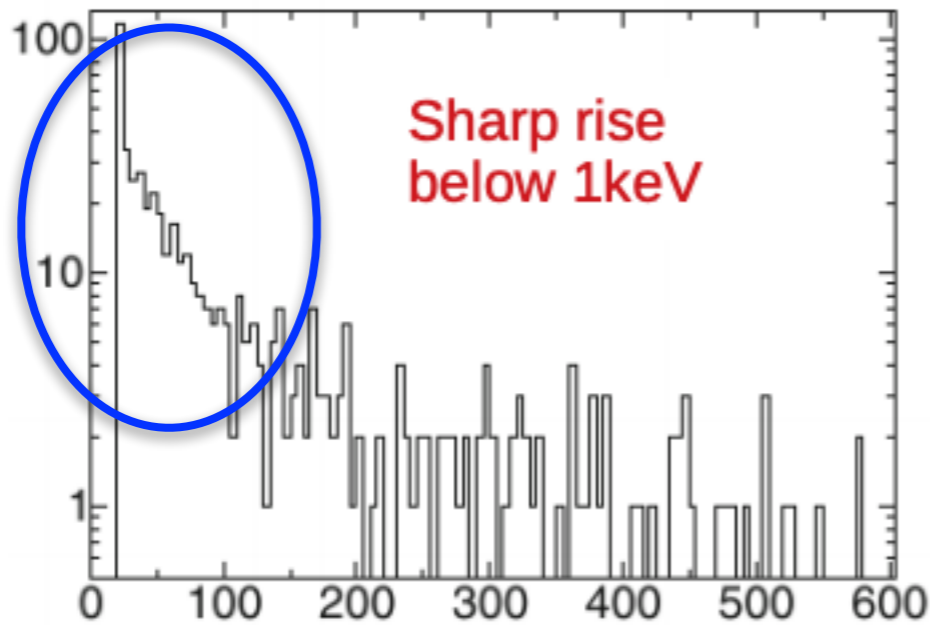
all energy
(heat + charge)



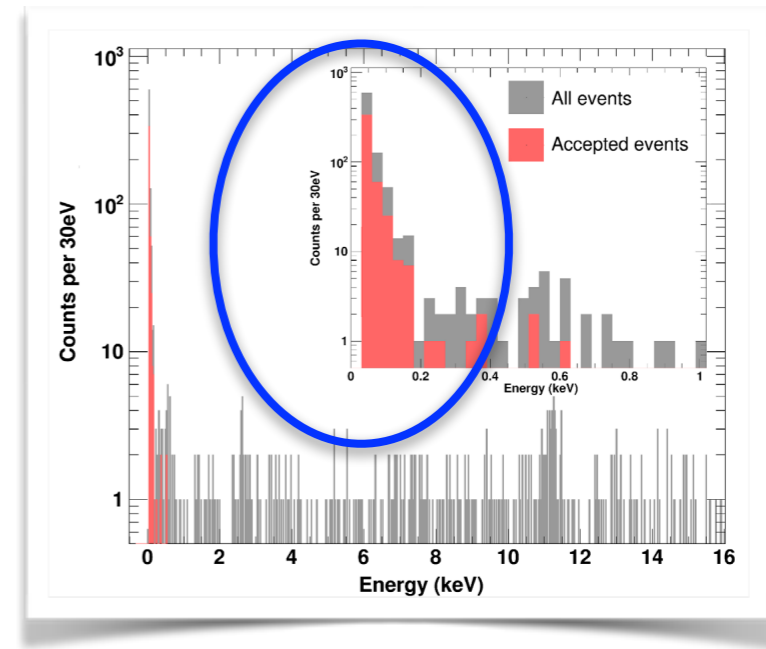
can't be nuclear recoil
or electron recoil...

a model which works: every event makes ~ 2 e/h pairs plus phonons.
Consistent if large plasmon-phonon coupling

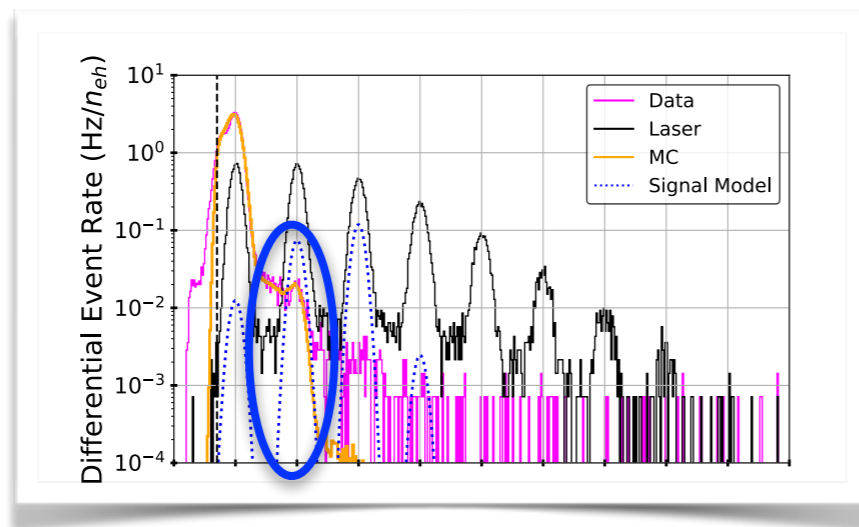
Once you start looking, you see it everywhere...



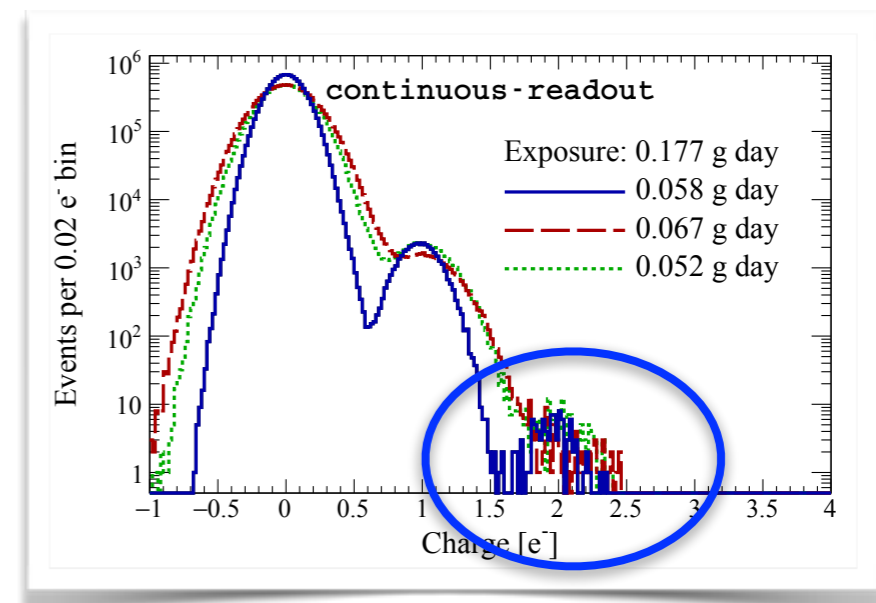
NUCLEUS surface Al_2O_3 (2017)



CRESST-III CaWO_4 (2019)



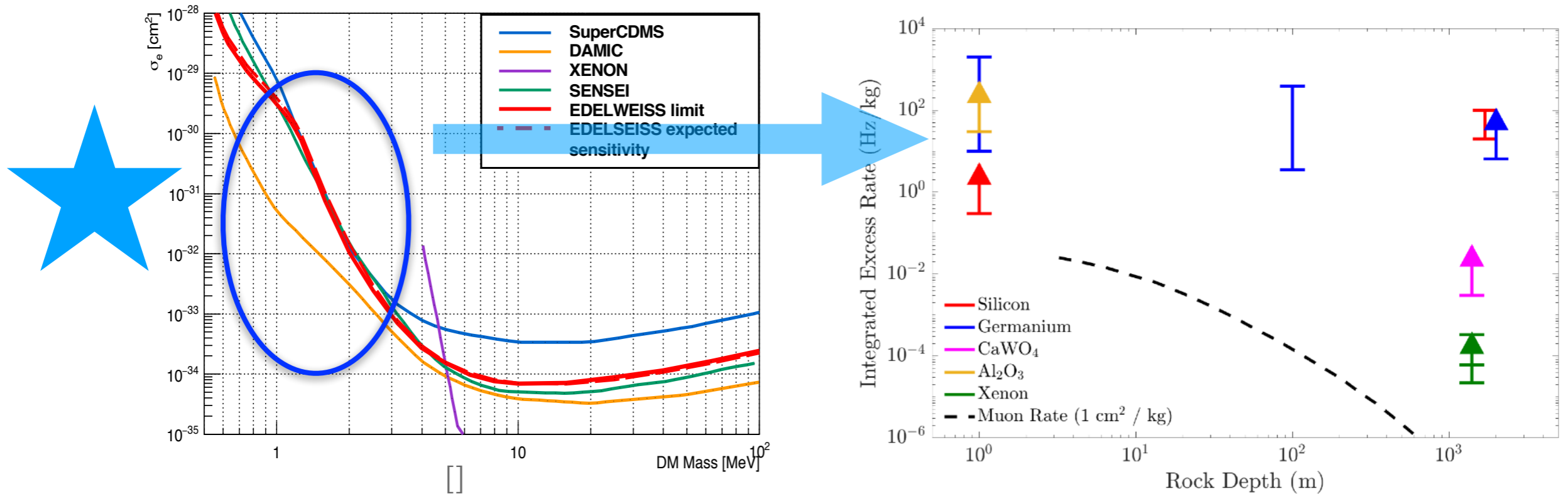
CDMS-HVeV Si (2018)



SENSEI underground Si (2019)

...and at the same rate!

Heavy Dark Sector mediator case (F=1)

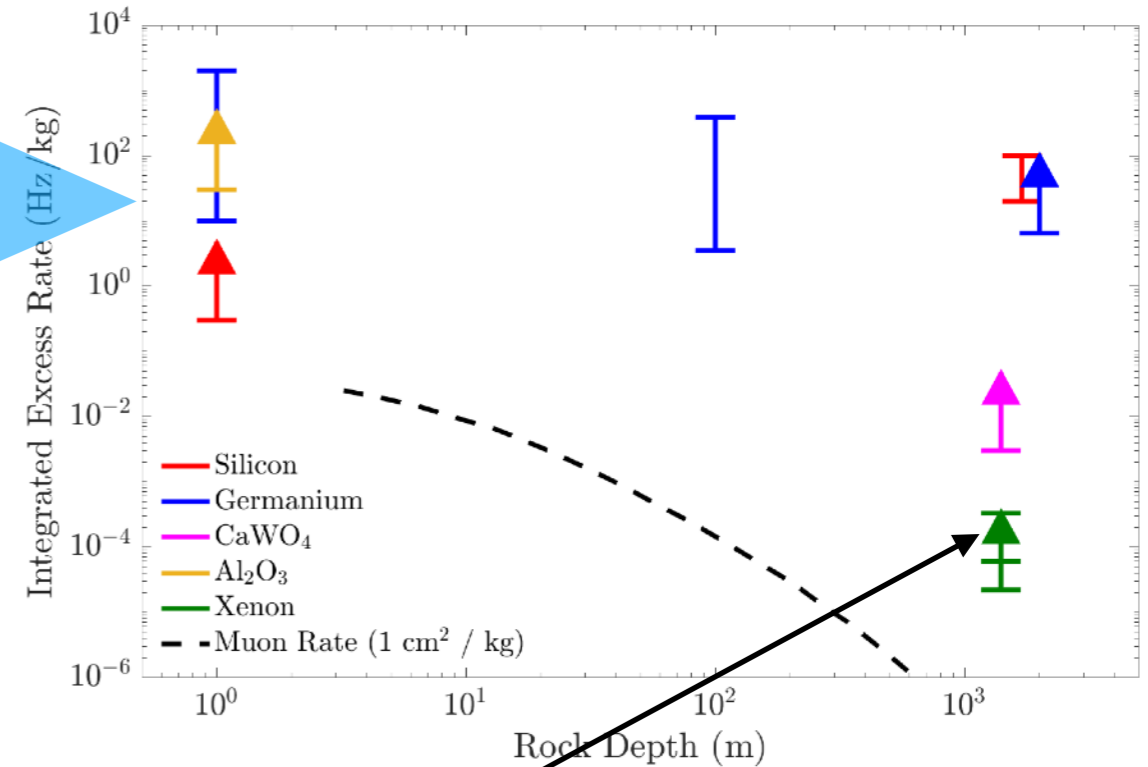
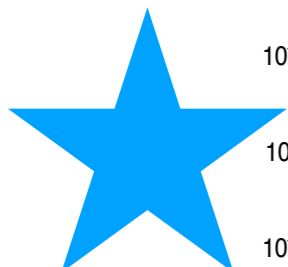
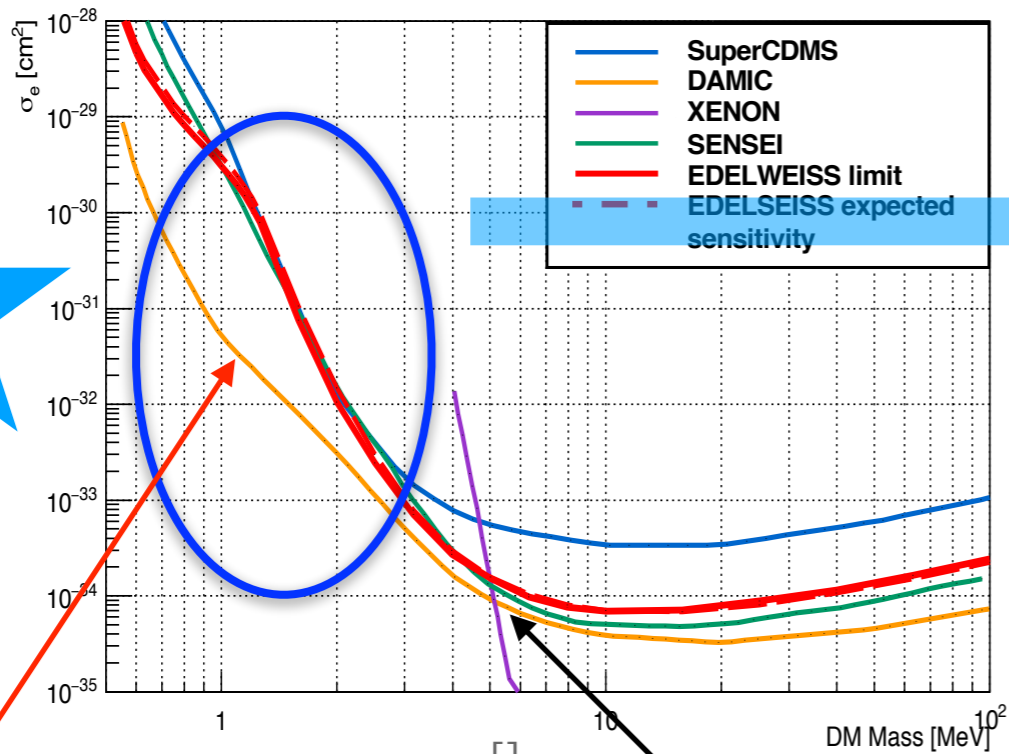


note: this limit plot assumes an electron recoil model.
This **DOES NOT** fit the spectrum!

But, obtaining the same cross section limits means the total 2-electron rate is the same everywhere

...and at the same rate!

Heavy Dark Sector mediator case (F=1)



DAMIC used different ionization model; under same model as CDMS and SENSEI, same limit
 [Ramanathan and Kurinsky, to appear]

no plasmon in Xe, lower rate (bound does not apply)

note: this limit plot assumes an electron recoil model.
 This **DOES NOT** fit the spectrum!

But, obtaining the same cross section limits means the total 2-electron rate is the same everywhere

**Something is making a plasmon at the
same rate everywhere in every
semiconductor detector**

**Something is making a plasmon at the
same rate everywhere in every
semiconductor detector**

Can it be dark matter?

Something is making a plasmon at the same rate everywhere in every semiconductor detector

Can it be dark matter?

To excite plasmon, need small q for fixed $\omega \implies v \gtrsim 10^{-2}$

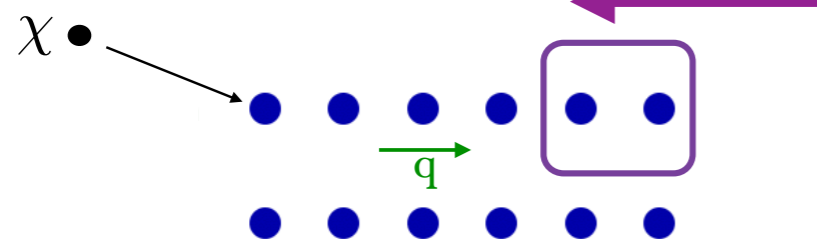
Faster than escape velocity! Either need fast DM, or indirect excitation

Two models:

Heavy and slow

$$m_{\text{DM}} \sim 100 \text{ MeV}$$

$$\text{KE} \sim 100 \text{ eV}$$

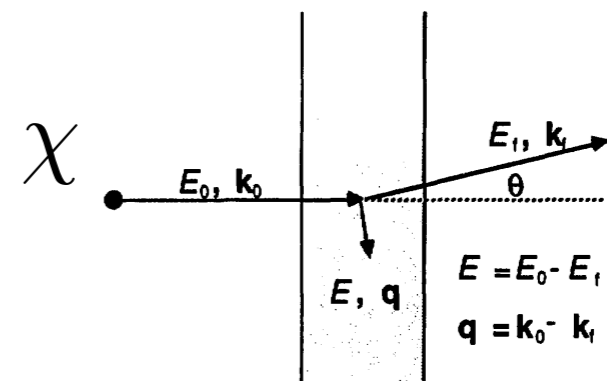


phonons
($q \sim 15 \text{ keV}$)

plasmon
($\omega \sim 15 \text{ eV}, q \sim 0$)

Fast and millicharged

DM is like an electron, with a long-range force. Identical dynamics to EELS, but **no multiple scattering** in a thick detector



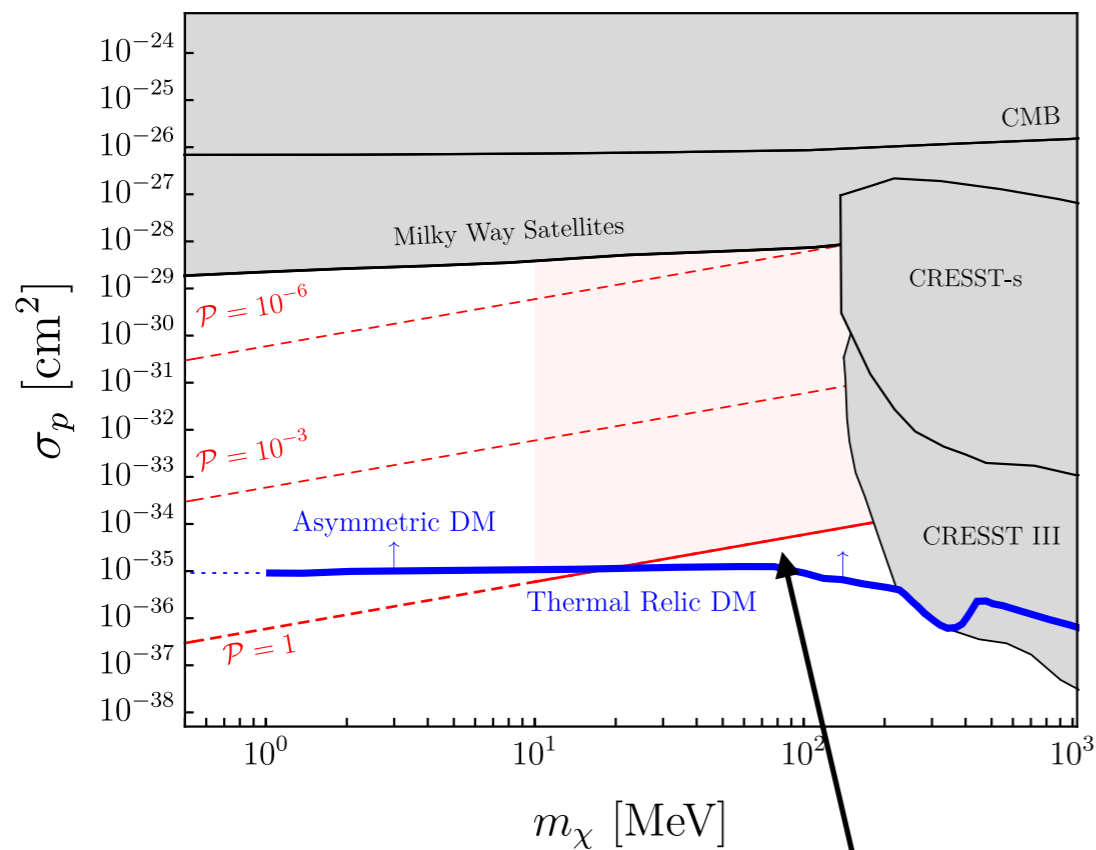
Parameter space

Heavy and slow

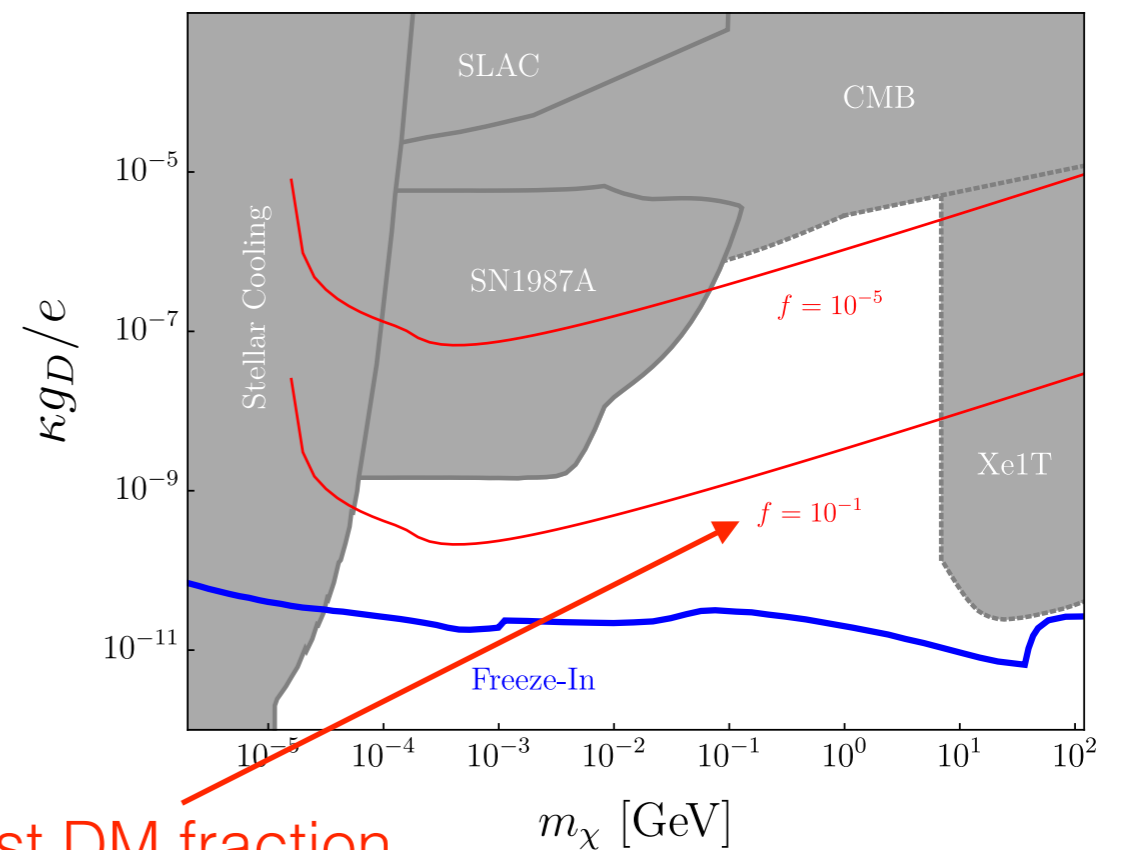
$$R \propto Z^2/A$$

Fast and millicharged

$$R \propto 1/\rho_T$$



Tantalizingly close to the thermal target...

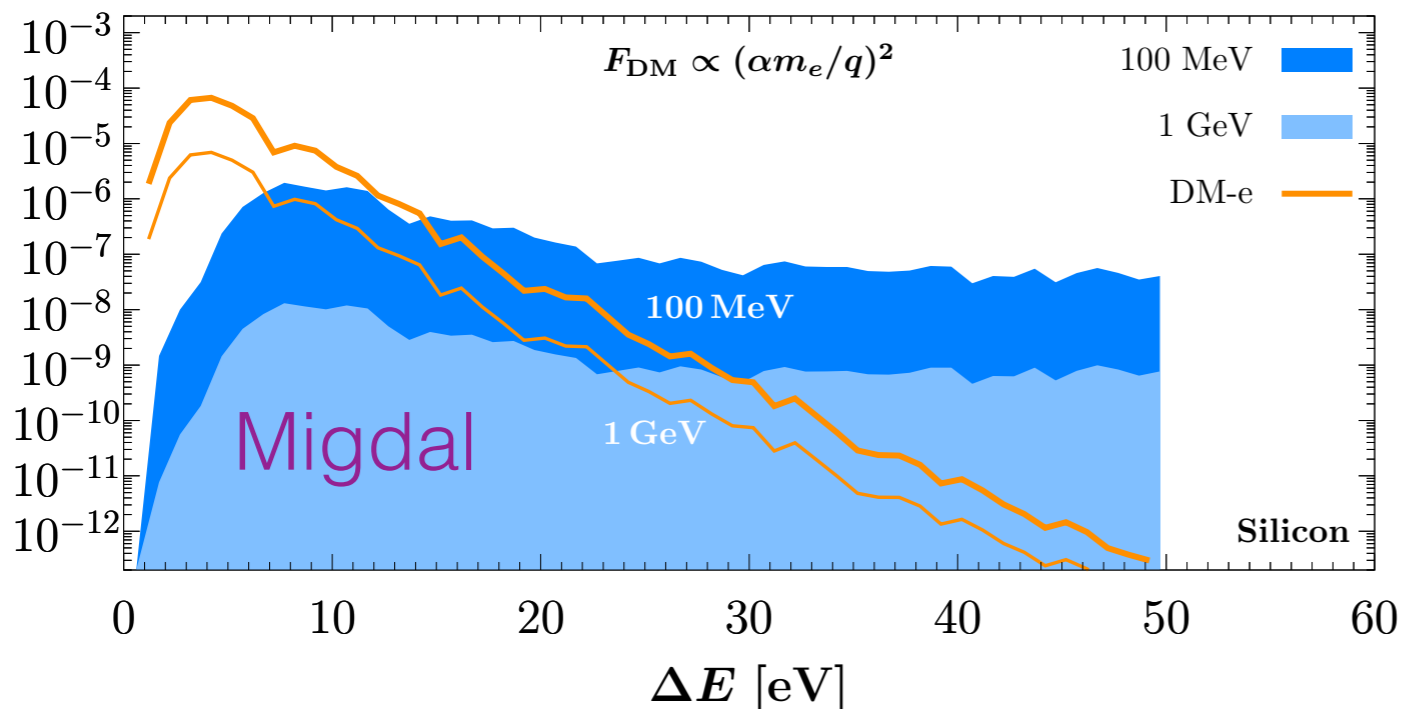


fast DM fraction

Suggests immediate discovery potential with 50 meV thresholds and mg-month exposures

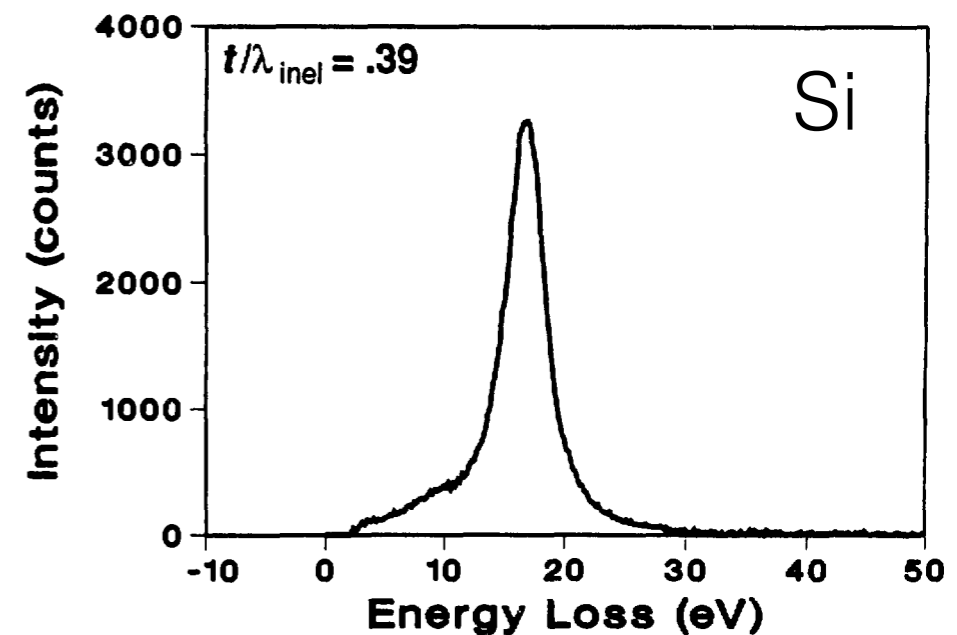
This is a completely new kind of signal

For direct plasmon excitation through a long-range force:



[Essig, Pradler, Sholapurkar, Yu, PRL 2020]

VS.



[M. Kundmann, Ph.D. thesis 1988]

Phonon detector: sharp peak at 15 eV with width 3 eV, long tail

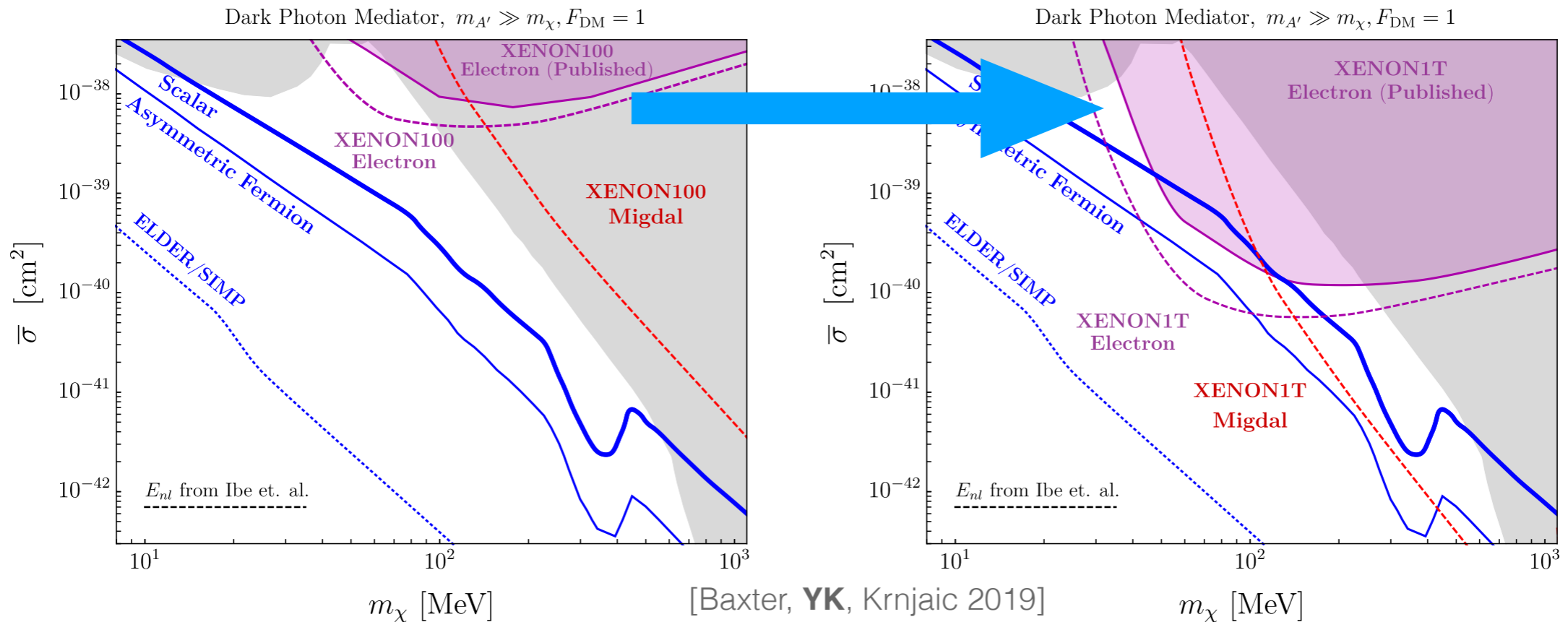
Charge detector: sharp peak in 2-electron bin, tail to 3 electrons

Looks nothing like electron recoil, nuclear recoil, or Migdal!

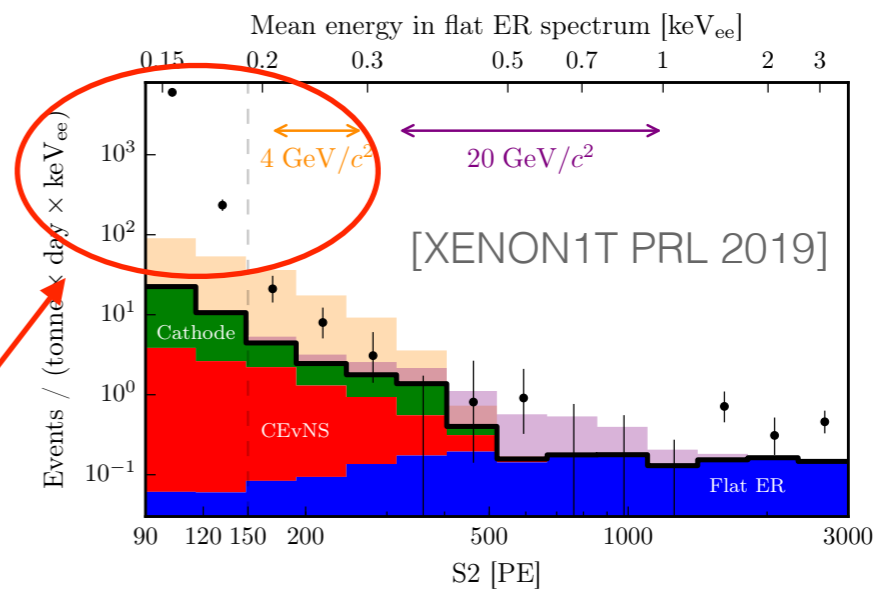
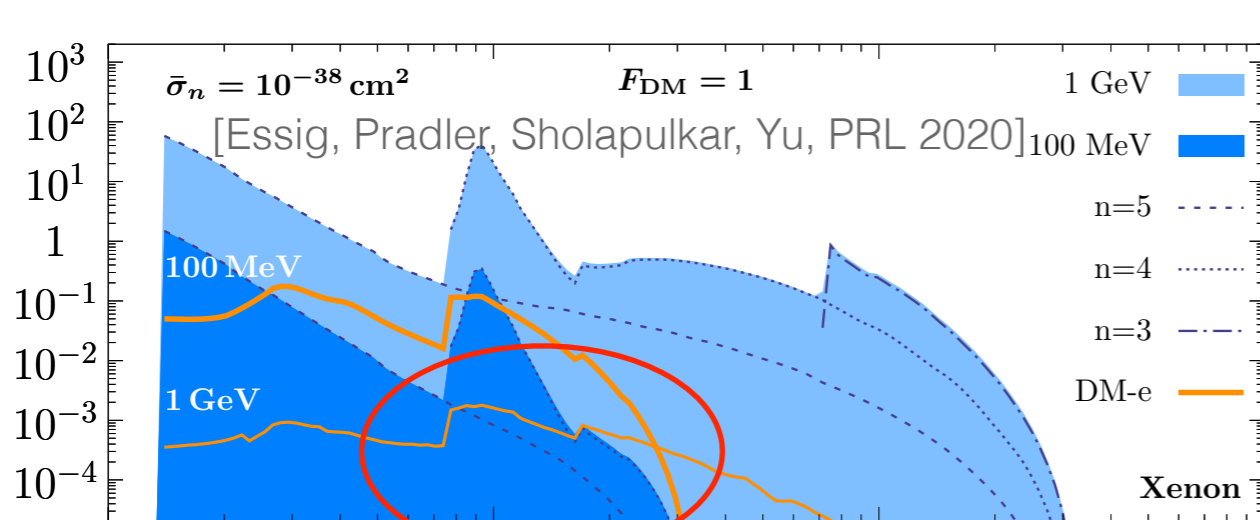
Quintessential many-body effect: only in semiconductors, not in Xe.

Spectrum determined by the material, NOT the DM!!

What about XENON?



XENON1T seems to set a much stronger limit...



But sees an unexplained excess!

If signal is ER + Migdal, completely different signal shape

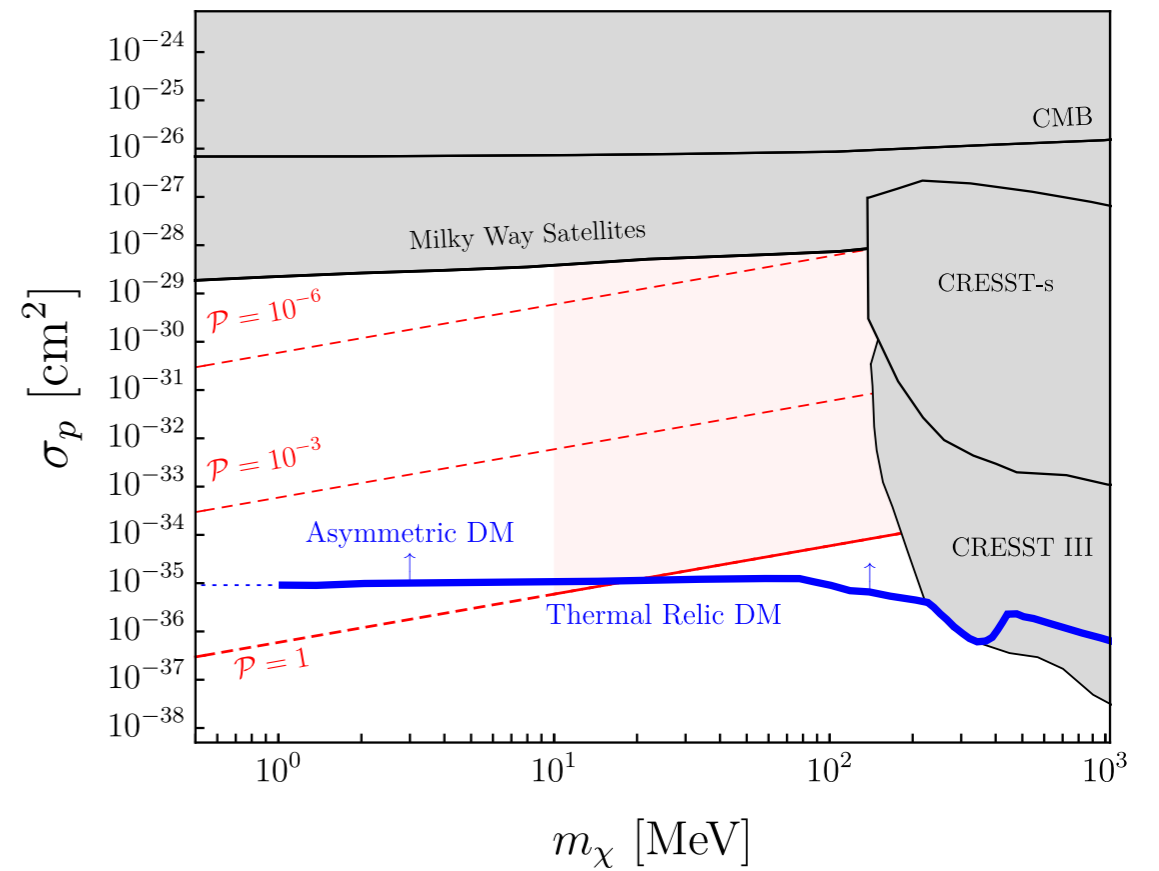
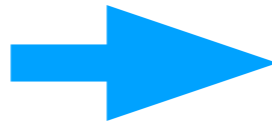
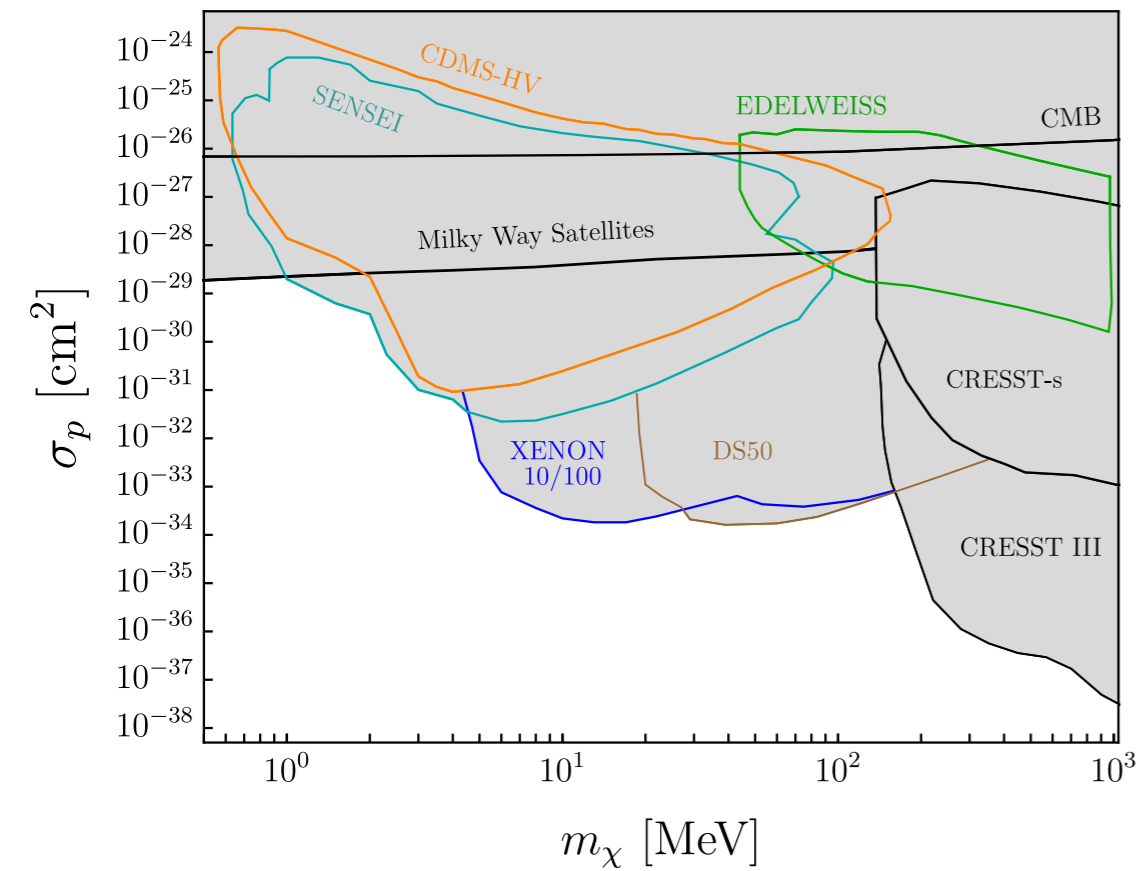
Conclusions

smaller = more speculative

- **There are lots of low-energy excesses**
- **The semiconductor rates are the same everywhere, regardless of overburden, exposure, or detector**
- Every semiconductor event is $\sim 1-2$ e/h pairs plus heat
- Every semiconductor event is a plasmon excitation
- Dark matter is exciting a plasmon in semiconductors and a combination of electron recoil and Migdal in noble liquids
- We've been seeing dark matter for the past 10 years!

Backup slides

Naive nuclear recoil limits

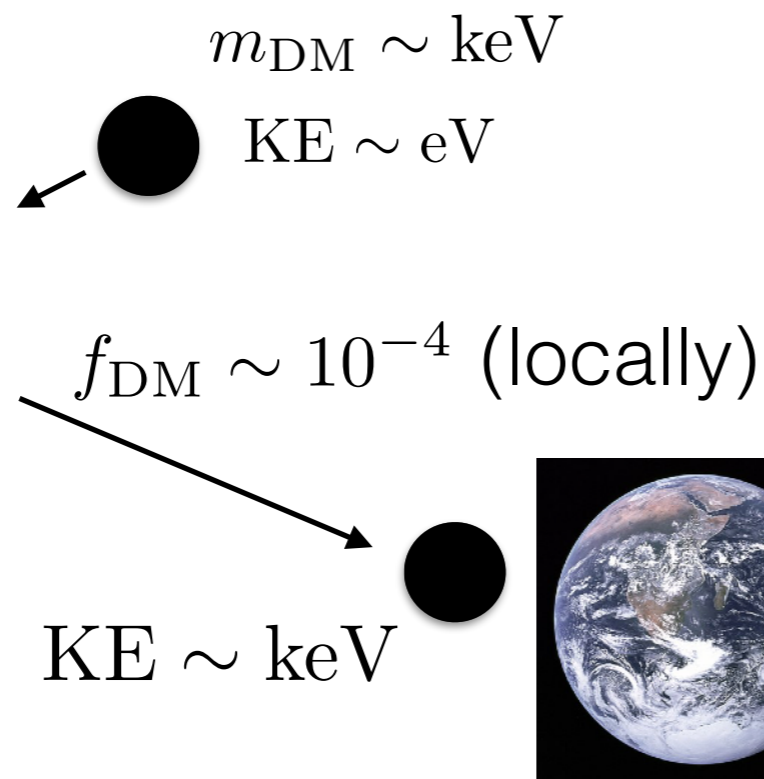
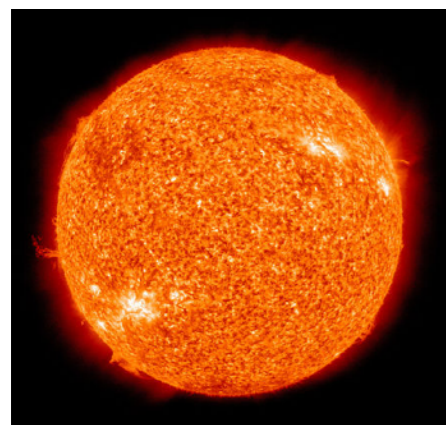


All the excesses

Readout Type	Target	Resolution	Exposure	Threshold	Excess Rate (Hz/kg)	Depth	Reference
Charge (E_e)	Ge	$1.6 e^-$	80 g·d	0.5 eVee ($\sim 1e^-$) ^a	20 ($\geq 2e^-$); 100 ($\geq 1e^-$)	1.7 km	EDELWEISS [12]
	Si	$\sim 0.2 e^-$	0.18 g·d	1.2 eVee ($< 1 e^-$)	6 ($\geq 2e^-$); 400 ($\geq 1e^-$)	100 m	SENSEI [4]
	Si	$0.1 e^-$	0.5 g·d	1.2 eVee ($< 1 e^-$)	10 ($\geq 2e^-$); 2000 ($\geq 1e^-$)	~ 1 m	CDMS HVeV [3]
	Si	$1.6 e^-$	200 g·d	1.2 eVee ($\sim 1e^-$) ^b	~ 6.5	2 km	DAMIC [13]
Energy (E_{det})	Ge	18 eV	200 g·d	60 eV	> 0.3	~ 1 m	EDELWEISS [1]
	CaWO ₄	4.6 eV	5600 g·d	30 eV	$> 3 \times 10^{-3}$	1.4 km	CRESST-III[2]
	Al ₂ O ₃	3.8 eV	0.046 g·d	20 eV	> 30	~ 1 m	ν CLEUS [14]
Photo e^-	Xe	6.7 PE ($\sim 0.25e^-$)	15 kg·d	12.1 eVee (~ 14 PE)	$[0.6, 3.3] \times 10^{-4}$	1.4 km	XENON10 [15]
	Xe	6.2 PE ($\sim 0.31e^-$)	30kg·yr	~ 70 eVee (~ 80 PE)	$> 2.2 \times 10^{-5}$	1.4 km	XENON100 [5]
	Ar		6780 kg·d	50 eVee	$> 6 \times 10^{-4}$	1.4 km	Darkside50 [16]
	C ₆ H ₆	< 1 PE	18.2 kg·d	$\sim 4-6$ eVee	~ 14	~ 1 m	EJ301 + PMT [17]

Models for millicharge acceleration

Solar reflection



Cosmic ray upscattering

Supernova acceleration